



Loop-mediated

Dark Matter – neutrino interactions

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The Standard Model is incomplete

The Standard Model is incomplete



electron
neutrino



muon
neutrino



tau
neutrino

Neutrino mass
mechanism

The Standard Model is incomplete



electron neutrino



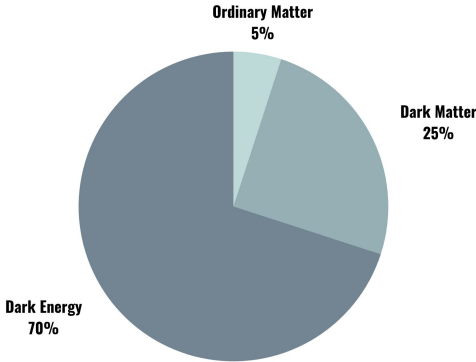
muon neutrino



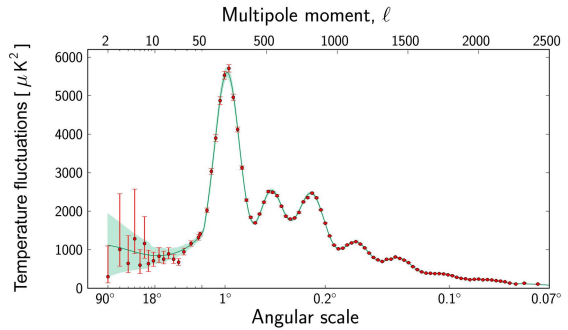
tau neutrino



Neutrino mass mechanism

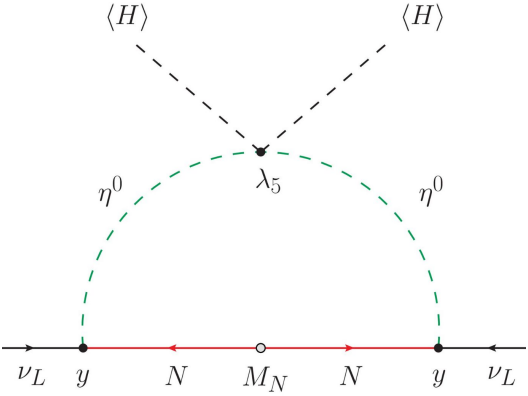


Astrophysical and cosmological evidence of dark matter



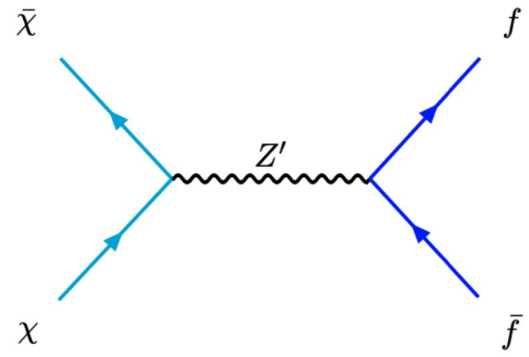
Theoretical Framework

Dark matter interactions with the standard model through the neutrino sector



Scotogenic model

Ma, E. (2006).
+many more extensions



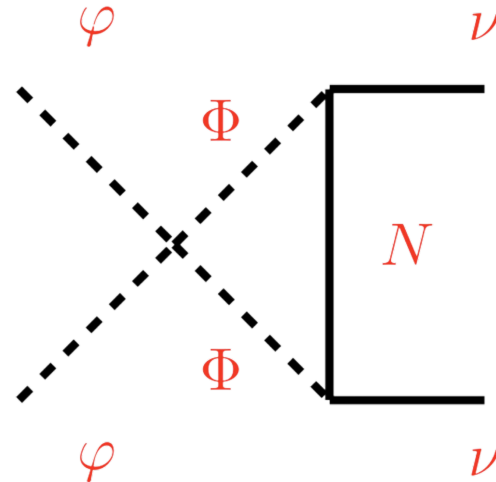
LR/U(1)' models

L, B-L symmetries

Senjanovic, G., & Mohapatra, R. N. (1975). PRD, 12(5), 1502.
 Mohapatra, R. N., & Pati, J. C. (1975). PRD, 11(9), 2558.
 G.B. Gelmini and M. Roncadelli, Phys. Lett. 99B (1981) 411
 +many more

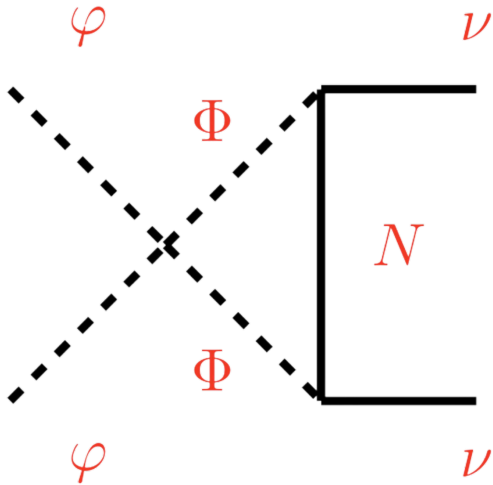
Research Question

Can dark matter interact with neutrinos via a 1-loop diagram, reproduce the observed relic abundance and still be consistent with current constraints?



Chao, W. (2020). arXiv: 2009.12002

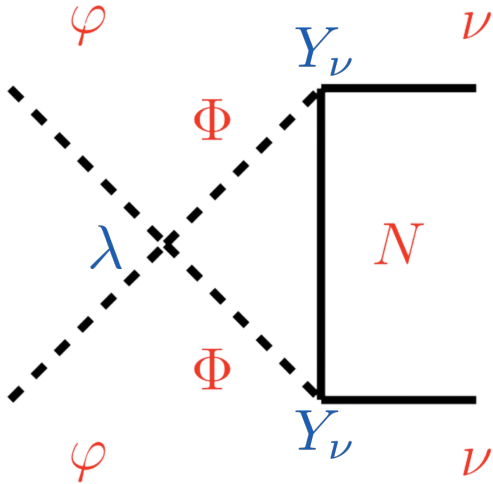
Dark matter annihilation into neutrinos



Real scalar DM	φ
Complex scalar mediator	Φ
Heavy neutrino mediator	N

Chao, W. (2020). arXiv: 2009.12002

Dark matter annihilation into neutrinos



Chao, W. (2020). arXiv: 2009.12002

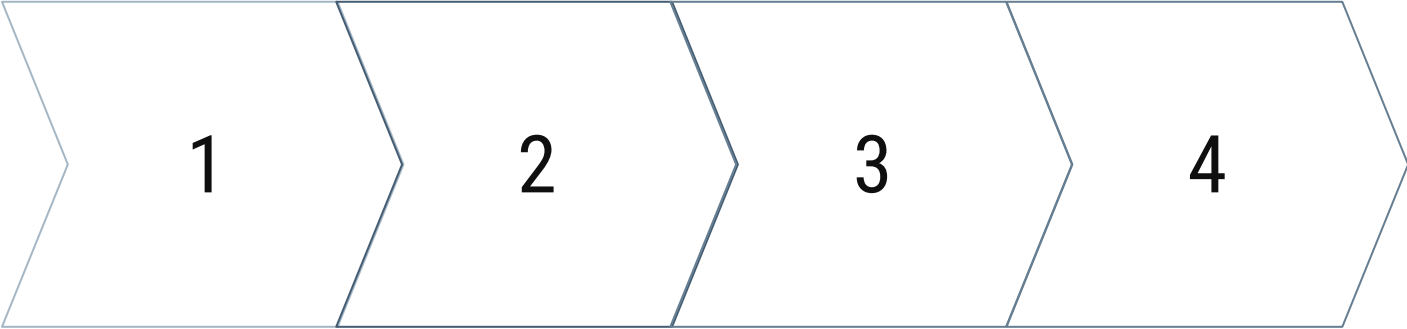
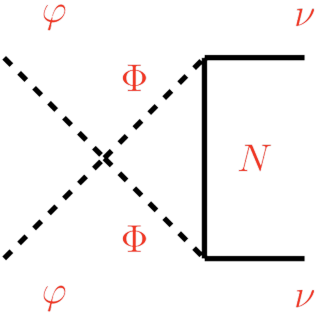
Real scalar DM φ

Complex scalar mediator Φ

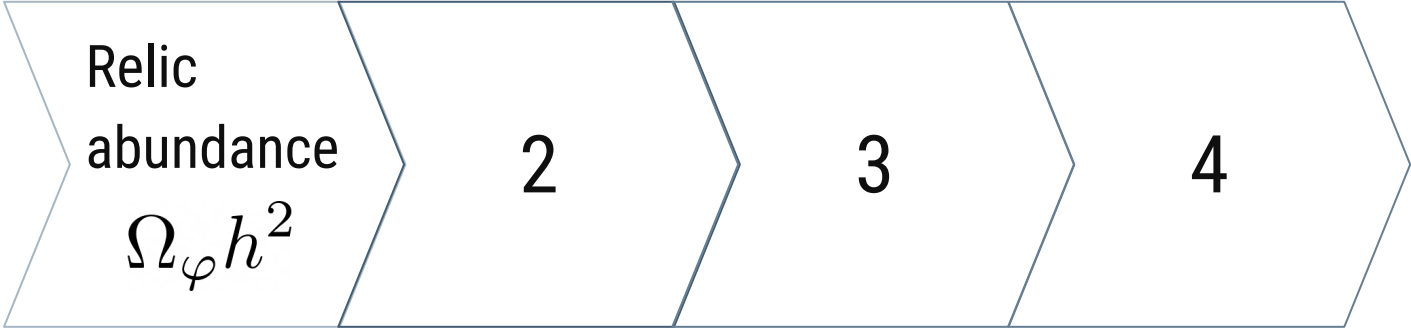
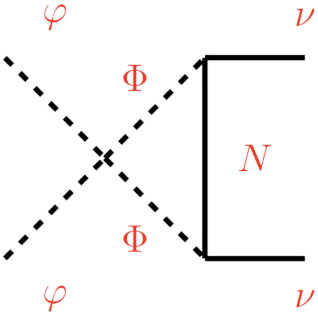
Heavy neutrino mediator N

Couplings λ, Y_ν

The Dark Matter Escape Room

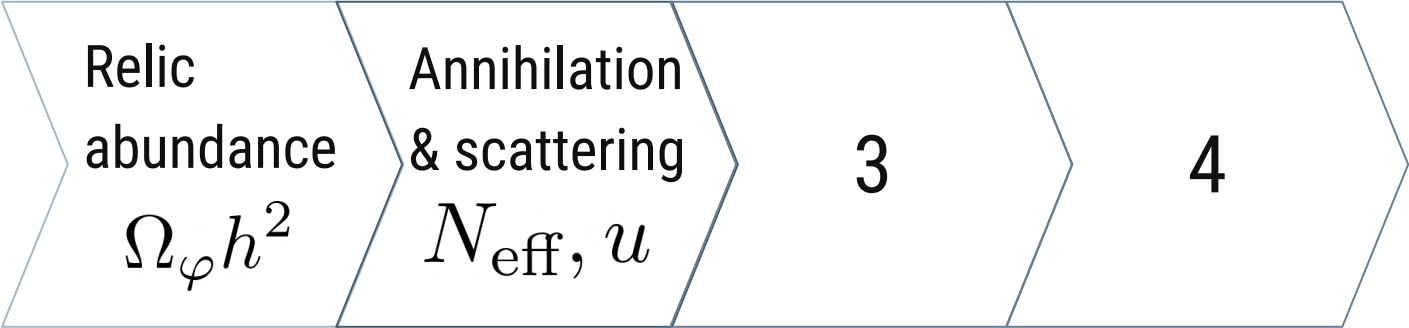
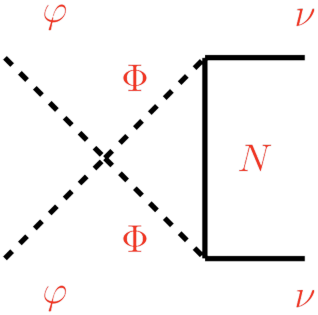


The Dark Matter Escape Room



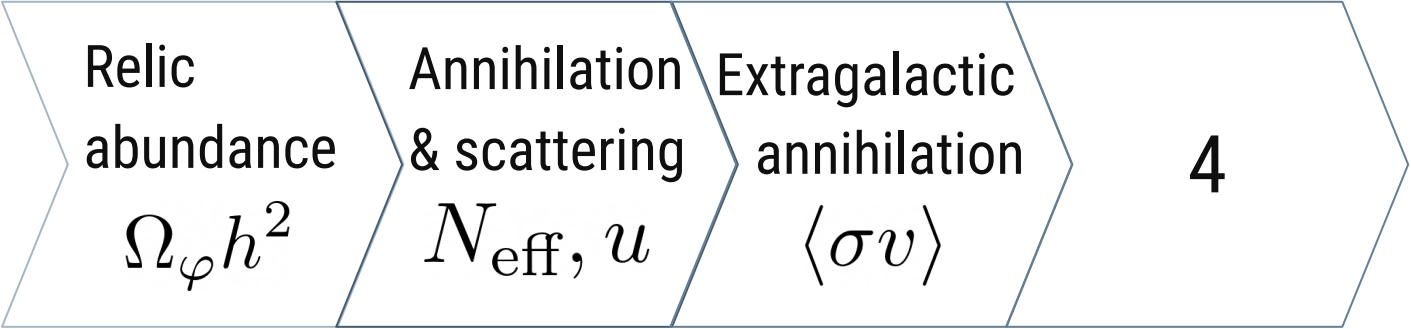
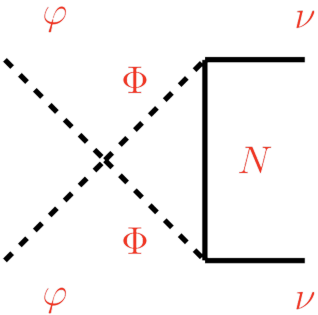
Early universe
DM-neutrino
interactions

The Dark Matter Escape Room



Early universe
DM-neutrino
interactions

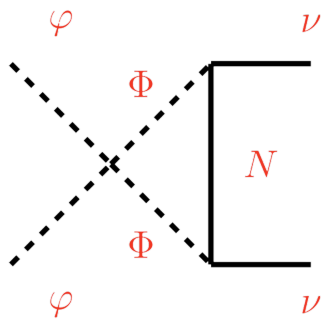
The Dark Matter Escape Room



Early universe
DM-neutrino
interactions

Supernova
Relic Neutrinos
(DSNB)

The Dark Matter Escape Room



Relic
abundance

$$\Omega_\phi h^2$$

Annihilation
& scattering

$$N_{\text{eff}}, u$$

Extragalactic
annihilation

$$\langle \sigma v \rangle$$

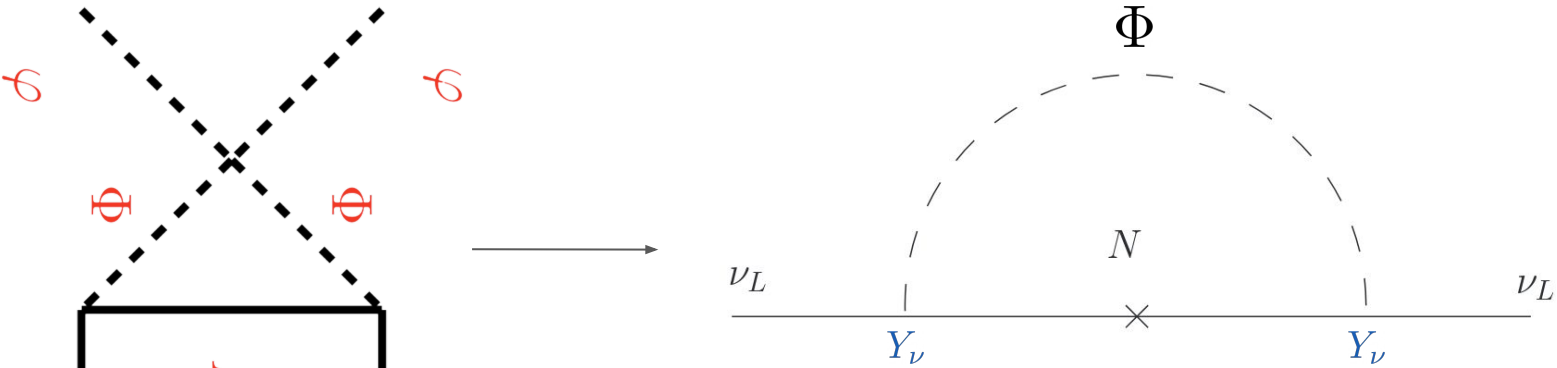
Lightest ν
mass

$$m_{\nu_0}$$

Early universe
DM-neutrino
interactions

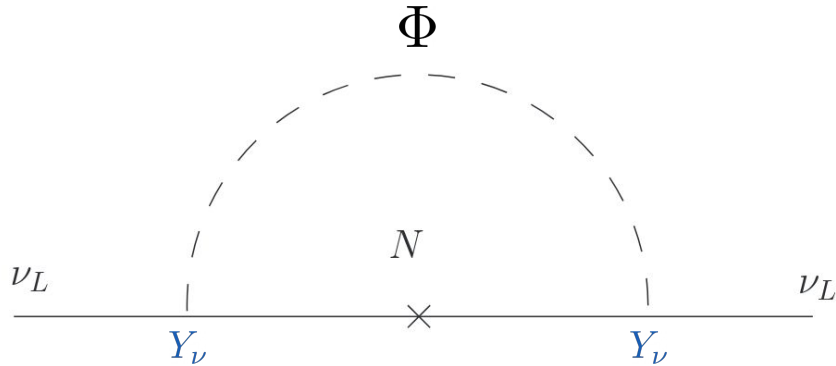
Supernova
Relic Neutrinos
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Neutrino mass mechanism



Böehm, C. et al. (2007). arXiv: 0612228

Neutrino mass mechanism

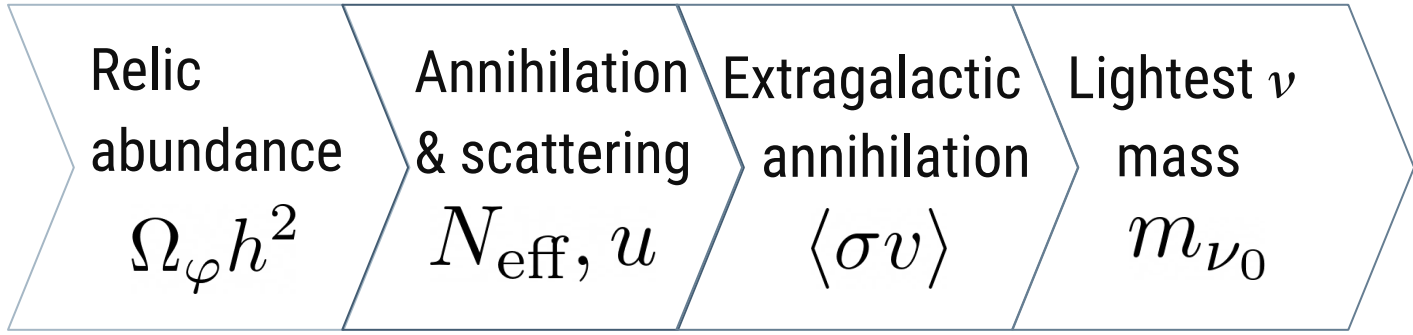
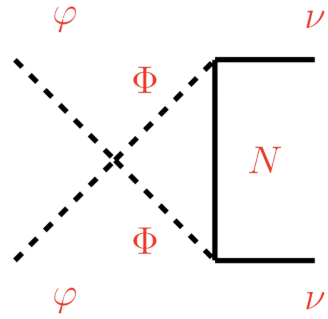


Böehm, C. et al. (2007). arXiv: 0612228

Neutrino mass m_ν
 Complex scalar
 mediator $\Phi = \frac{\phi_1 + i\phi_2}{\sqrt{2}}$
 Heavy neutrino
 mediator N

Couplings Y_ν

The Dark Matter Escape Room

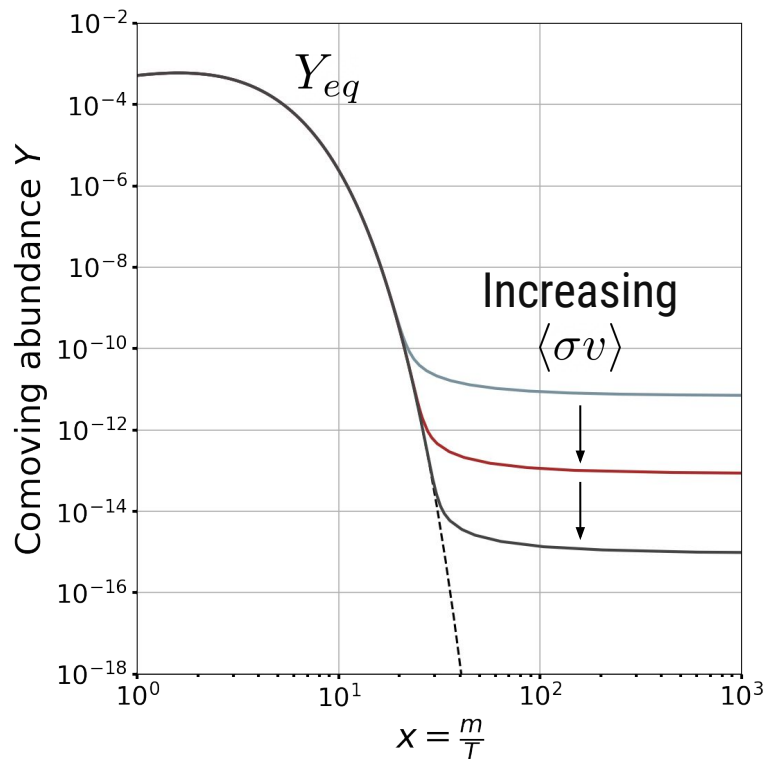


Early universe
DM-neutrino
interactions

Supernova
Relic Neutrinos
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MCMC

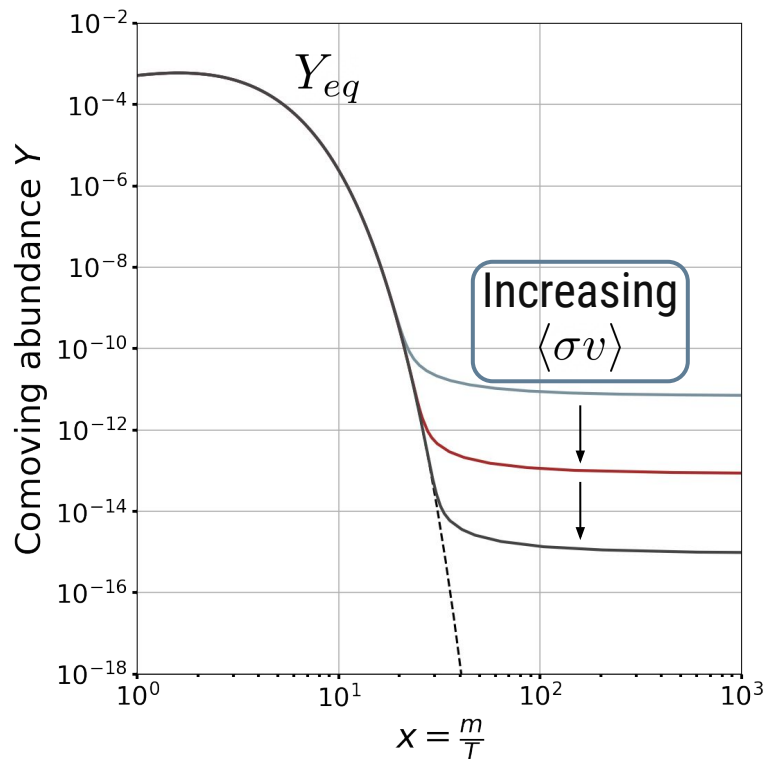
First Test: Dark Matter Relic Abundance



Boltzmann equation

$$\frac{dY}{dx} = \frac{s\langle\sigma v\rangle}{Hx} \left[1 + \frac{1}{3} \frac{d(\ln g_s)}{d(\ln T)} \right] (Y_{eq}^2 - Y^2)$$

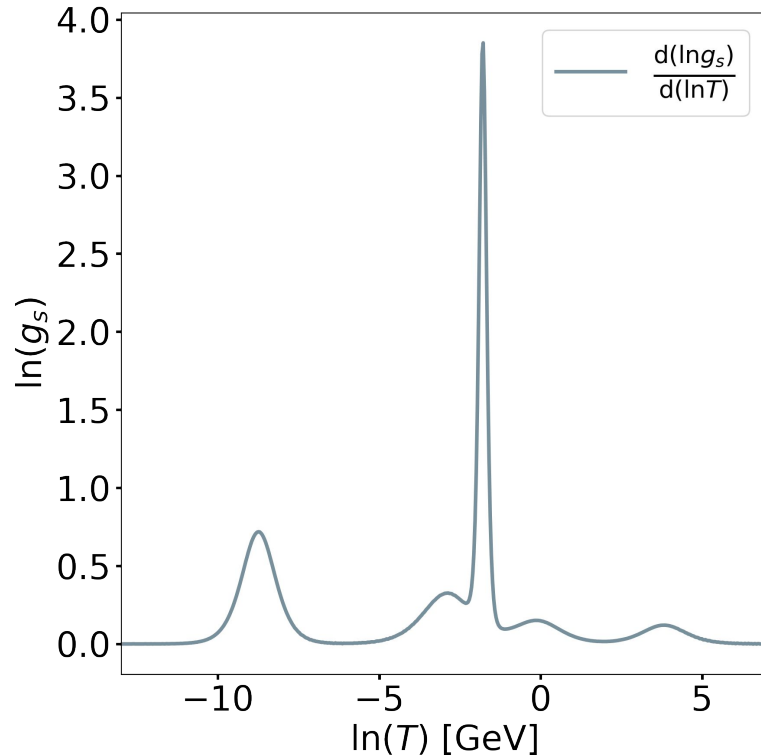
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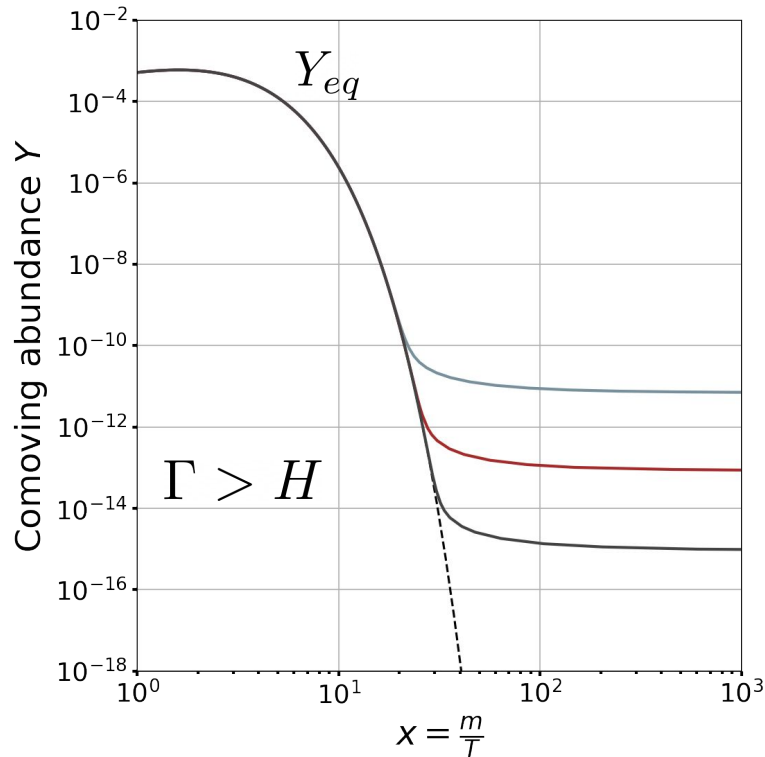
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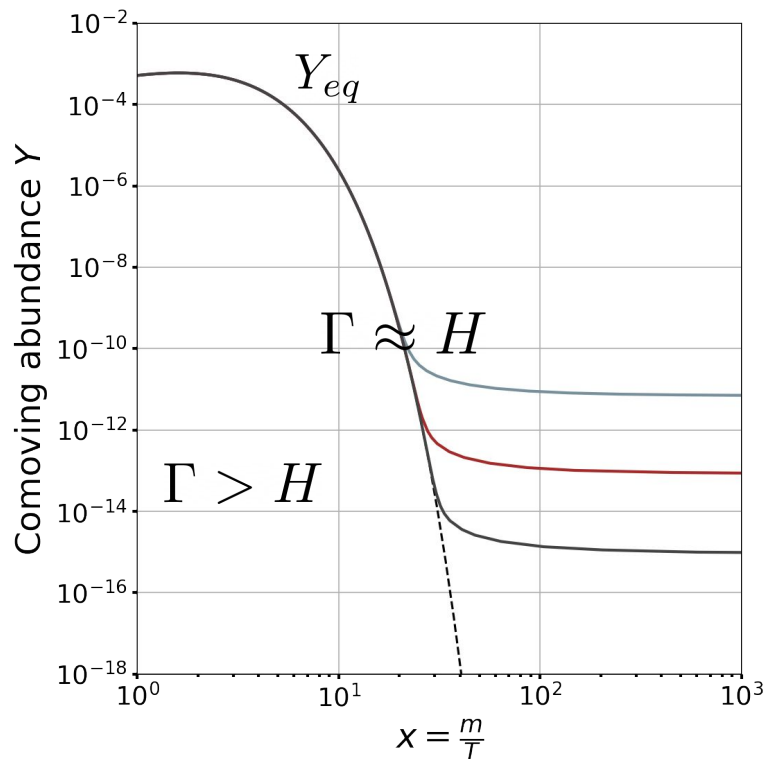


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$\Gamma > H$ Follows equilibrium abundance

First Test: Dark Matter Relic Abundance



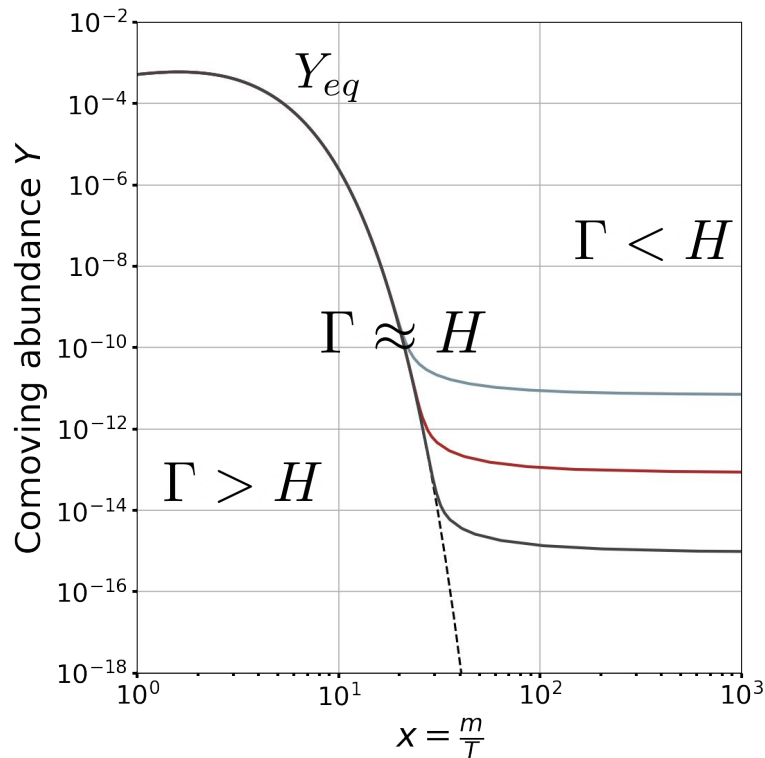
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$\Gamma \approx H$ Freeze-out

First Test: Dark Matter Relic Abundance



Boltzmann equation

$$\frac{dY}{dx} = \frac{s\langle\sigma v\rangle}{Hx} \left[1 + \frac{1}{3} \frac{d(\ln g_s)}{d(\ln T)} \right] (Y_{eq}^2 - Y^2)$$

$\Gamma > H$ Follows equilibrium abundance

$\Gamma \approx H$ Freeze-out

$\Gamma < H$ Away from equilibrium, sets
relic abundance

First Test: Dark Matter Relic Abundance

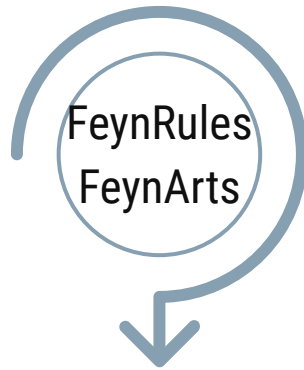
Interaction Lagrangian

$$-\mathcal{L}_{\text{int}} = \frac{1}{2}m_{\varphi}^2\varphi^2 + \frac{1}{2}\lambda\Phi^2\varphi^2 + Y_{\nu}\overline{\nu}_L\Phi N_R + \text{h.c.}$$

First Test: Dark Matter Relic Abundance

Interaction Lagrangian

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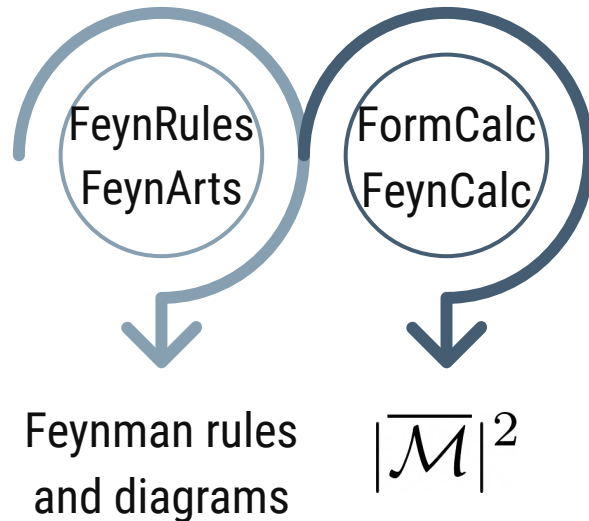


Feynman rules
and diagrams

First Test: Dark Matter Relic Abundance

Interaction Lagrangian

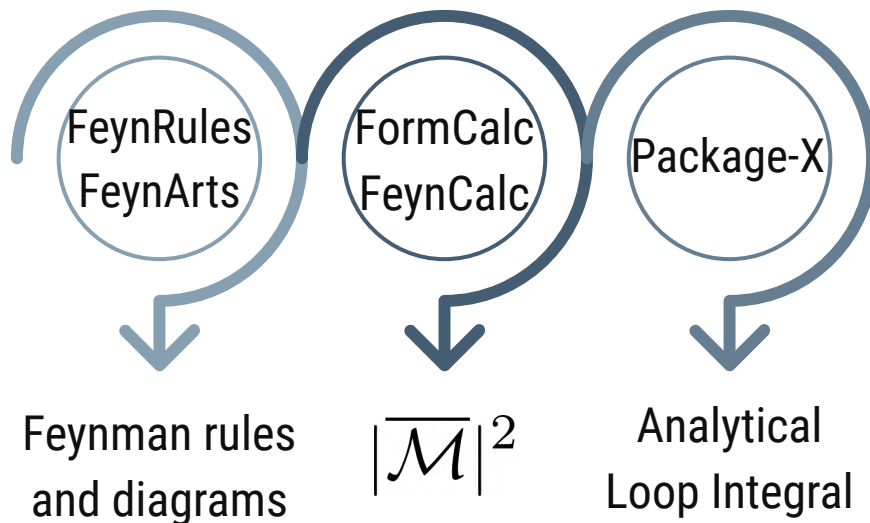
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First Test: Dark Matter Relic Abundance

Interaction Lagrangian

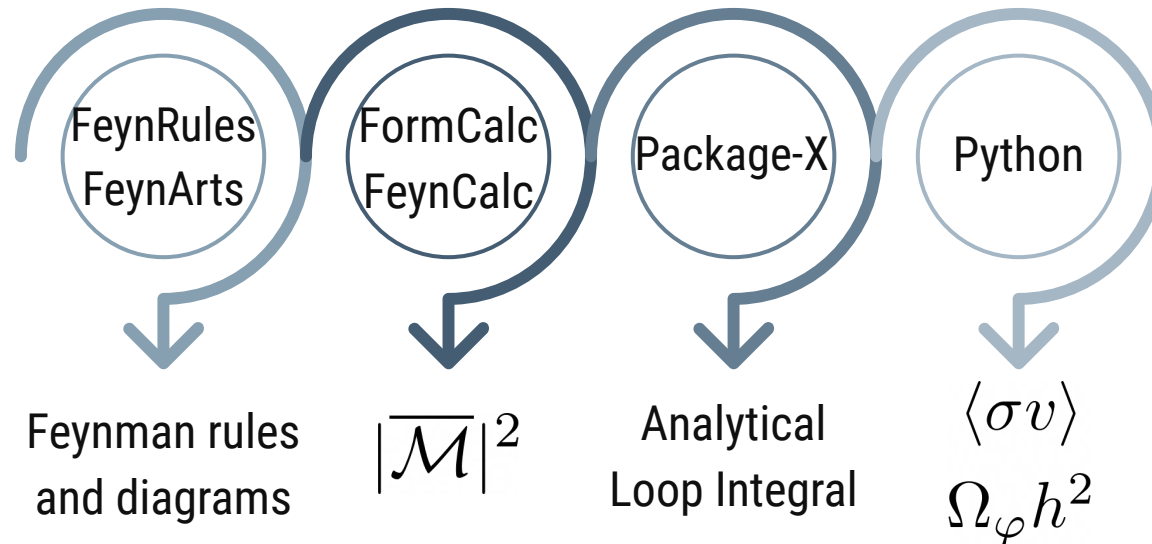
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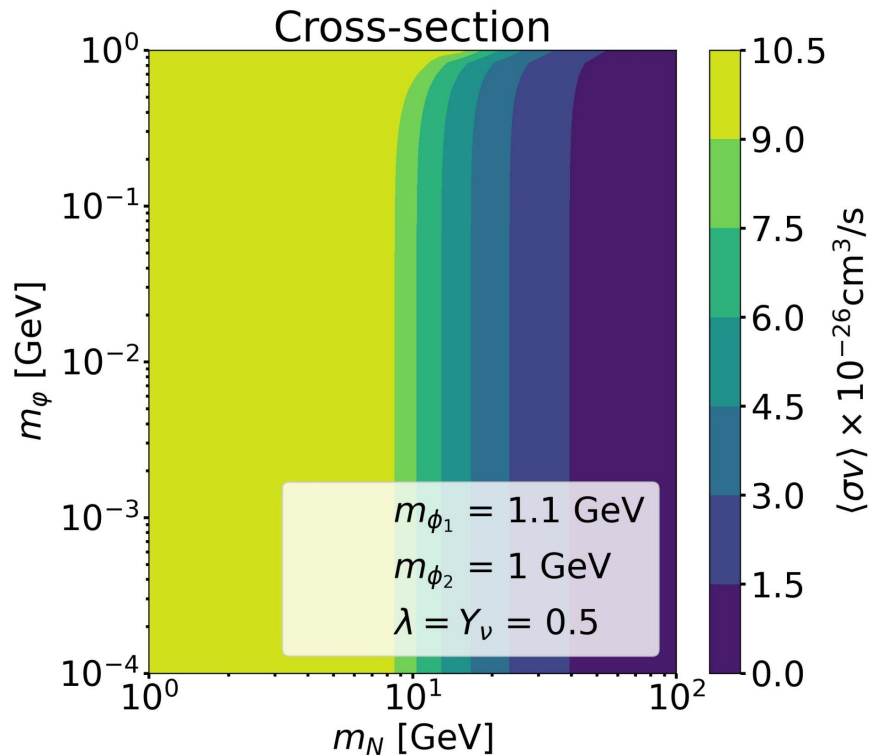
First Test: Dark Matter Relic Abundance

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First Test: Dark Matter Relic Abundance



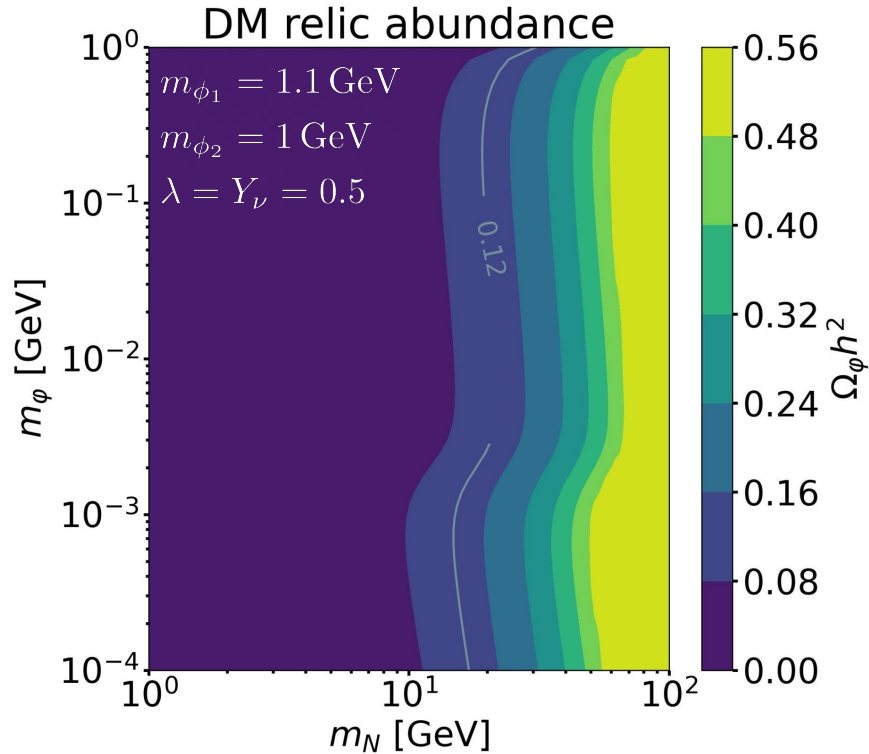
Stability constraint

$$m_N > m_\phi > m_\varphi$$

Nearly independent of

$$m_\varphi \leq 100 \text{ MeV}$$

First Test: Dark Matter Relic Abundance



Constraint from Planck 2018

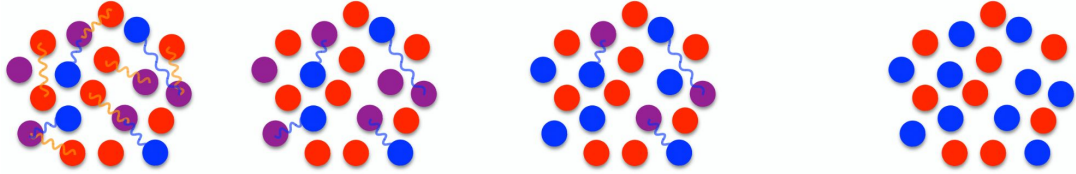
$$\Omega_c h^2 = 0.1200 \pm 0.0012$$

(68%, Planck TT,TE,EE + lowE + lensing)

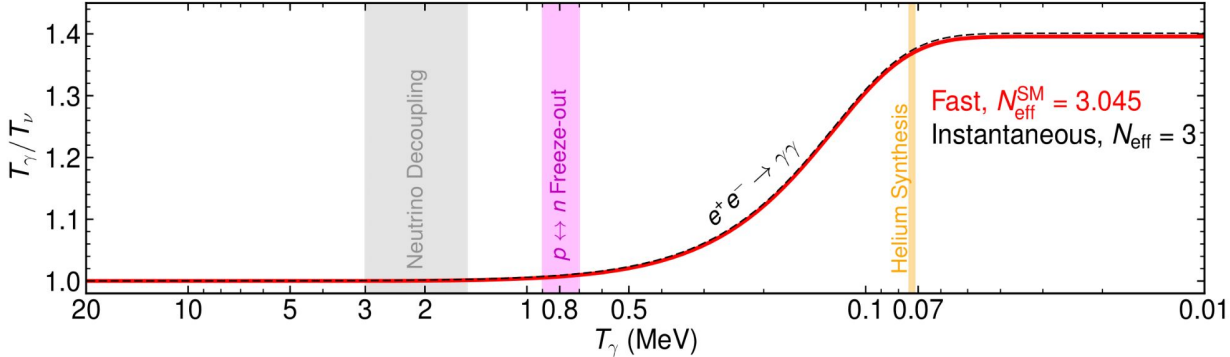
Second Test: Primordial abundances and N_{eff}

Big Bang Nucleosynthesis (BBN)

ν e^\pm γ W/Z



Evolution in the Standard Model



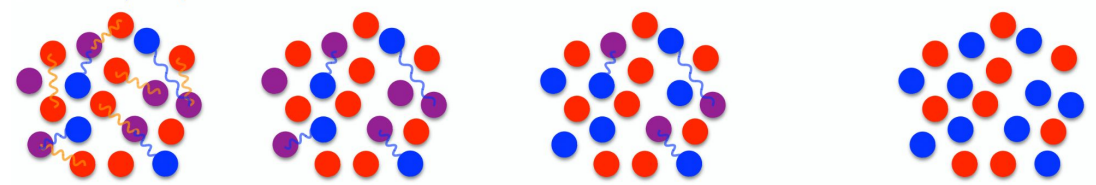
$$T_\gamma/T_\nu = (11/4)^{1/3}$$

$$\rho_{\text{rad}} = \rho_\gamma \left[1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{\text{eff}} \right]$$

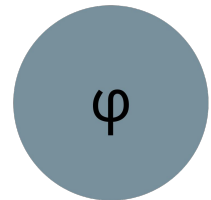
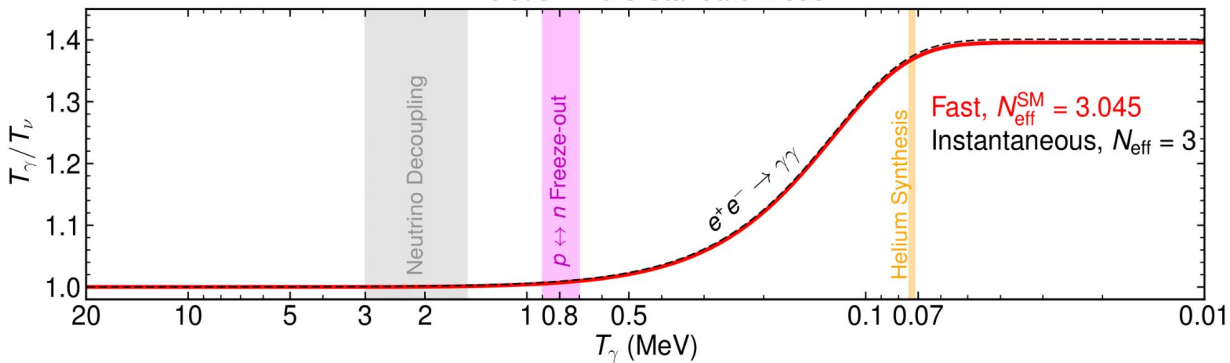
Second Test: Primordial abundances and N_{eff}

BBN and light dark matter

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Evolution in the Standard Model

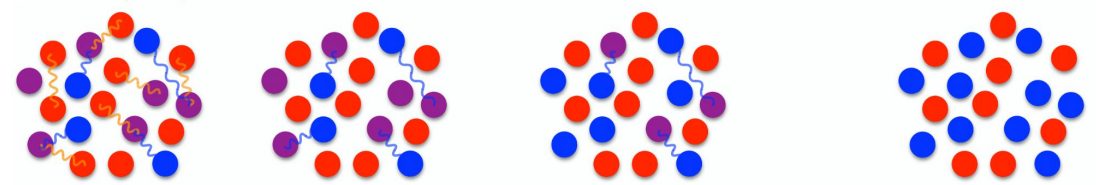


$$m_\phi \leq 20 \text{ MeV}$$

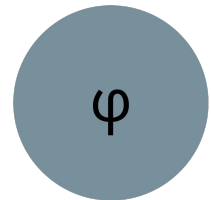
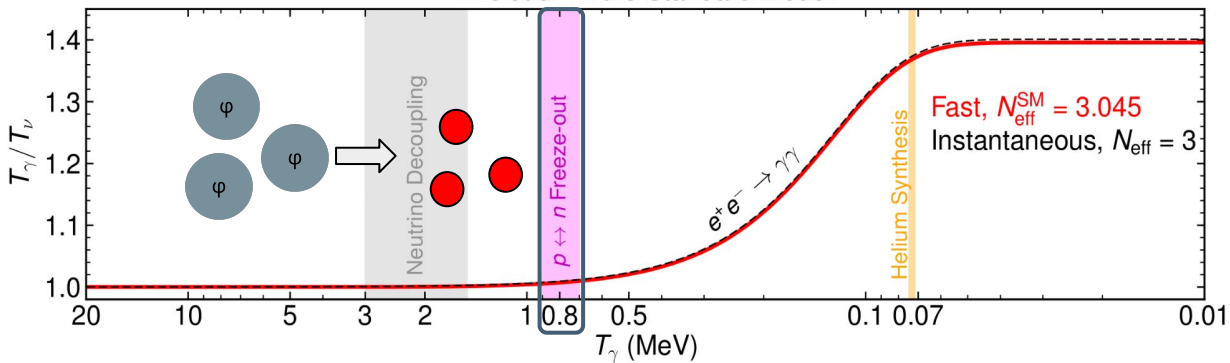
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Evolution in the Standard Model



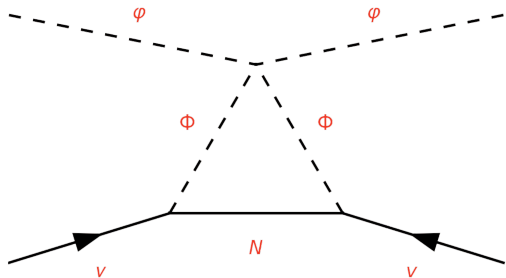
$$m_\phi \leq 20 \text{ MeV}$$

$$Y_P \quad Y_D \quad N_{\text{eff}}$$

Second Test: DM-neutrino scattering

Thermally averaged
scattering cross-section

$$\langle \sigma \rangle_{\text{DM}-\nu} = \frac{\int d^3 \mathbf{p}_\nu f_\nu \sigma}{\int d^3 \mathbf{p}_\nu f_\nu}$$

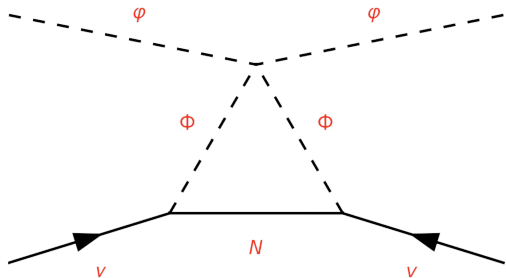


Fermi-Dirac distribution $f_\nu(k) = \frac{1}{e^{-k/T_\nu} + 1}$

Second Test: DM-neutrino scattering

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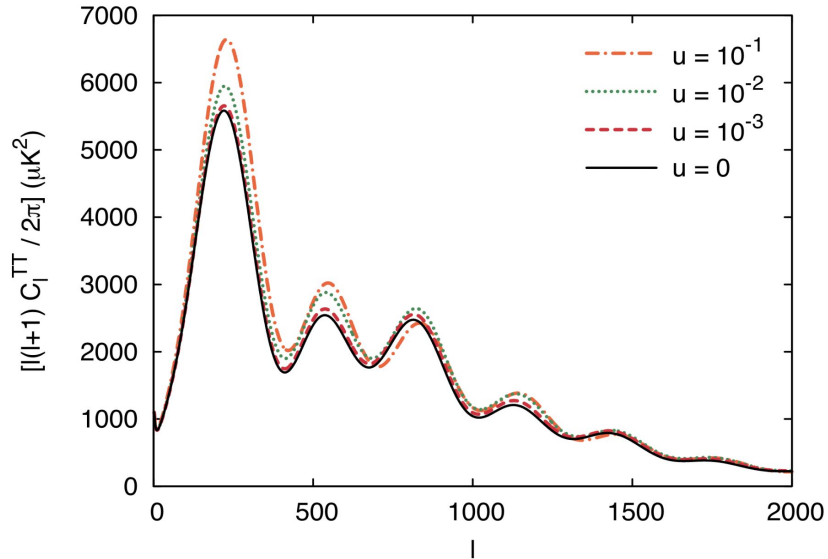
Fermi-Dirac distribution $f_\nu(k) = \frac{1}{e^{-k/T_\nu} + 1}$

Strength parameter $u \equiv \left[\frac{\langle \sigma \rangle_{\text{DM}-\nu}}{\sigma_{\text{Th}}} \right] \left[\frac{100 \text{ GeV}}{m_{\text{DM}}} \right]$

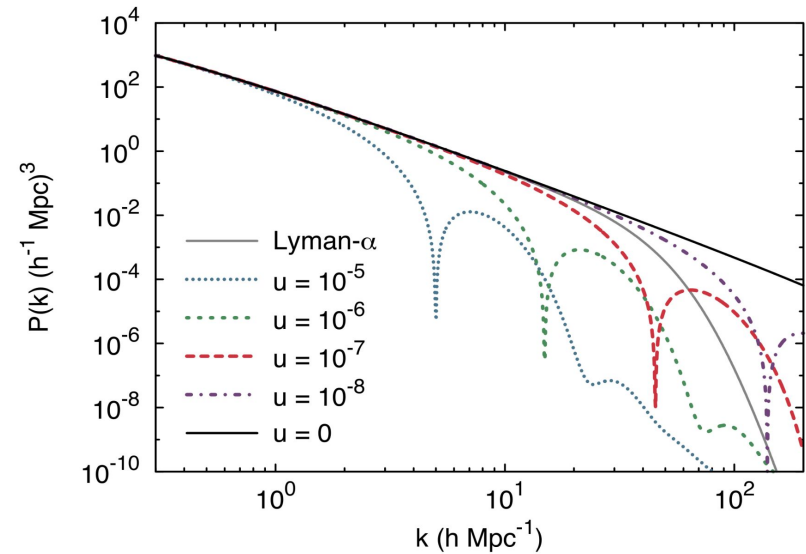
Wilkinson, R.J. et al. (2014). arXiv: 1401.7597

Second Test: DM-neutrino scattering

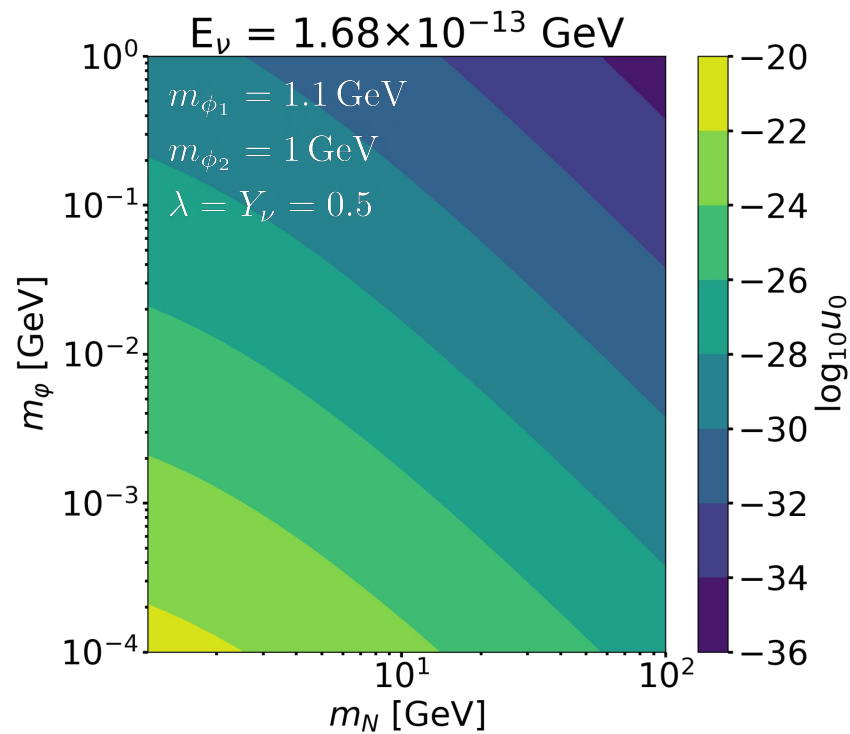
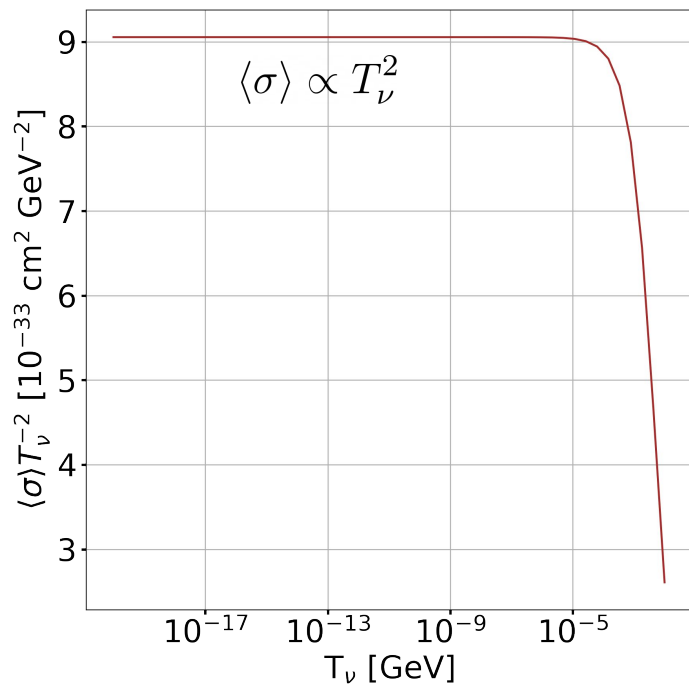
CMB Power Spectrum



Matter Power Spectrum



Second Test: DM-neutrino scattering



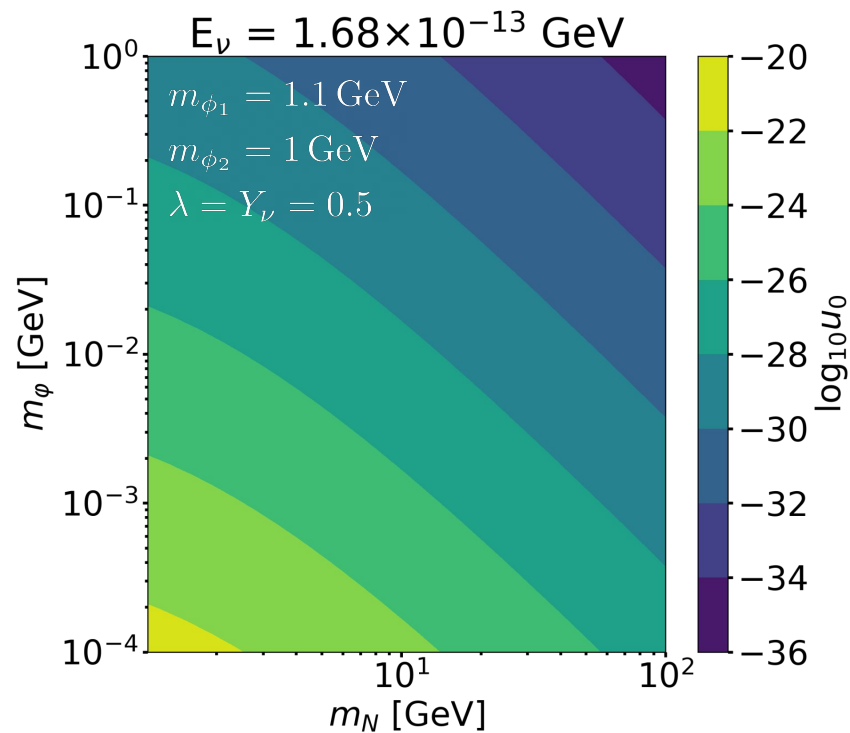
Second Test: DM-neutrino scattering

CMB & Matter Power Spectrum

$$\sigma_{\text{DM}-\nu,0} \lesssim 10^{-45} (m_{\text{DM}}/\text{GeV}) \text{ cm}^2$$

$$\sigma_{\text{DM}-\nu} \propto T^2, 10^{+13} u_0 < 2.56$$

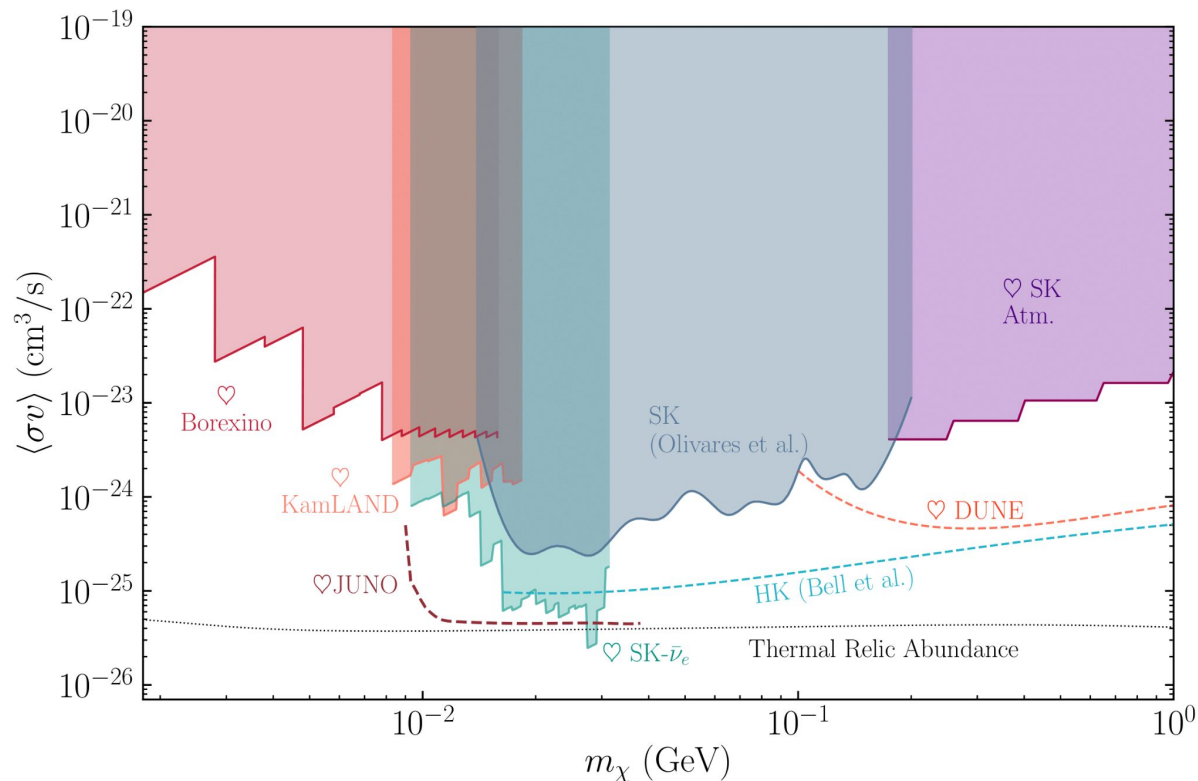
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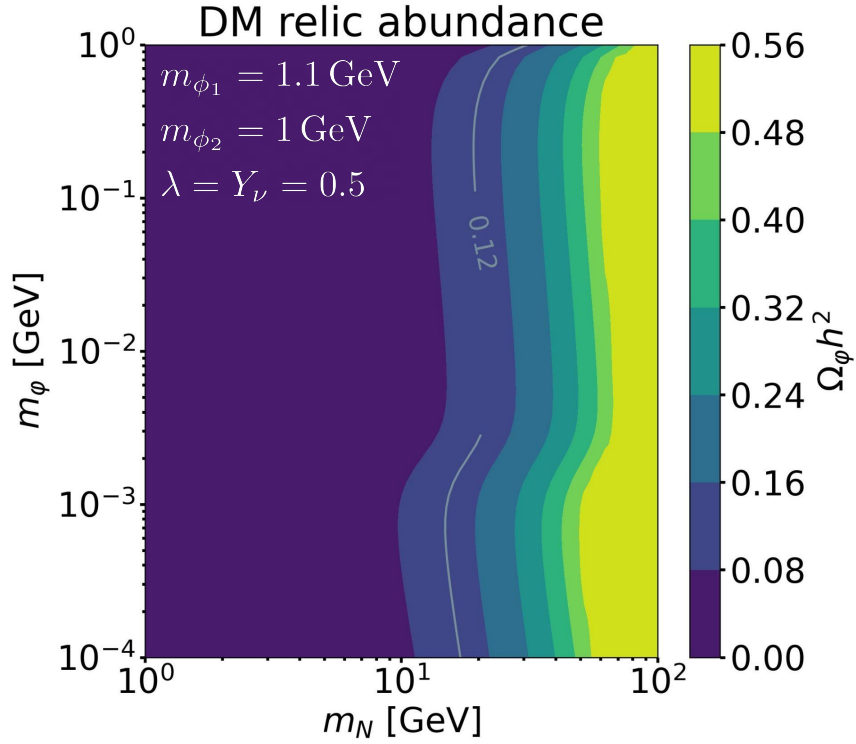
Third Test: Extragalactic DM Annihilation

Strongest constraint comes from the relic supernova electron antineutrino flux detected by Super-Kamiokande

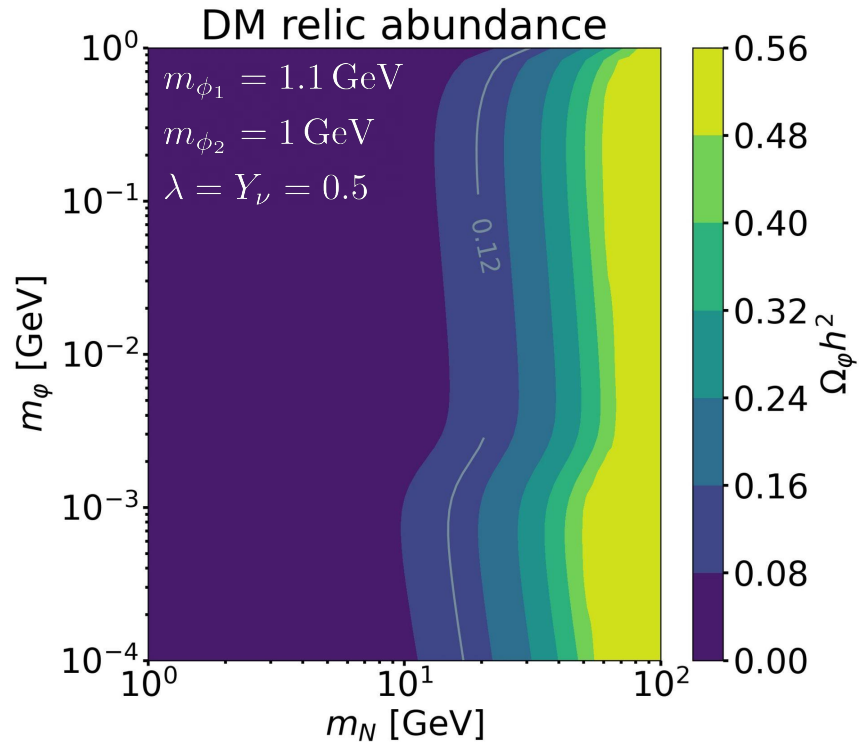
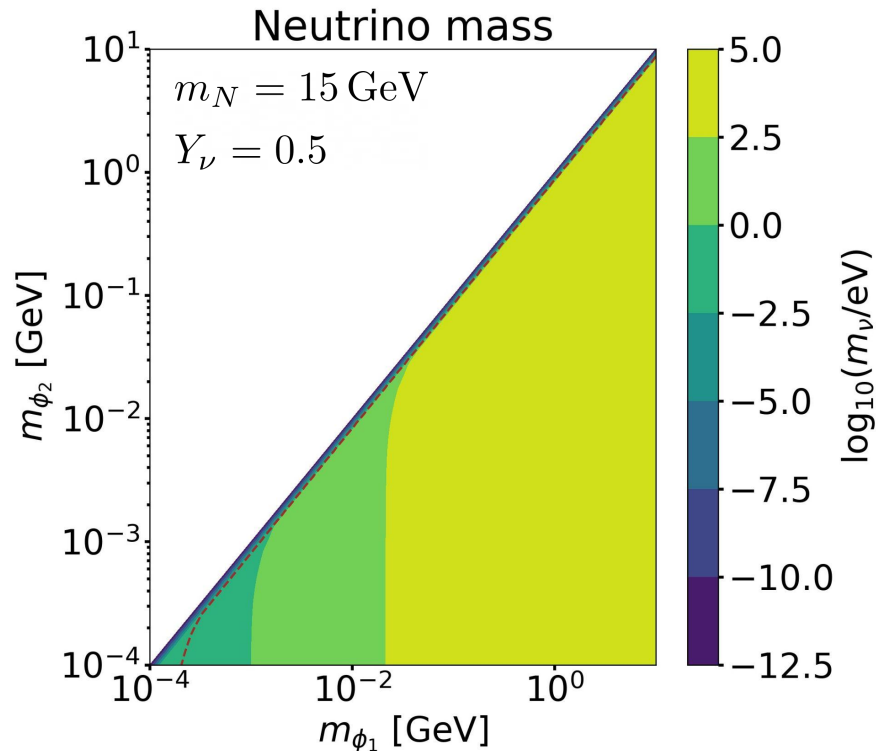
$$27 < m_\chi < 30 \text{ MeV}$$



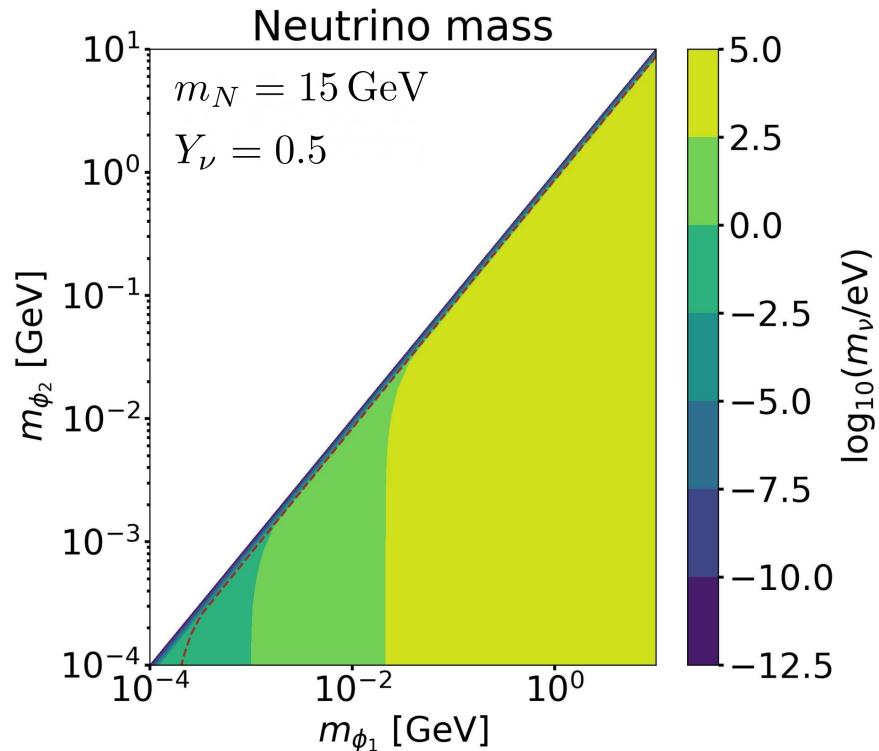
Fourth test: Adding neutrino masses



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Fourth test: Adding neutrino masses

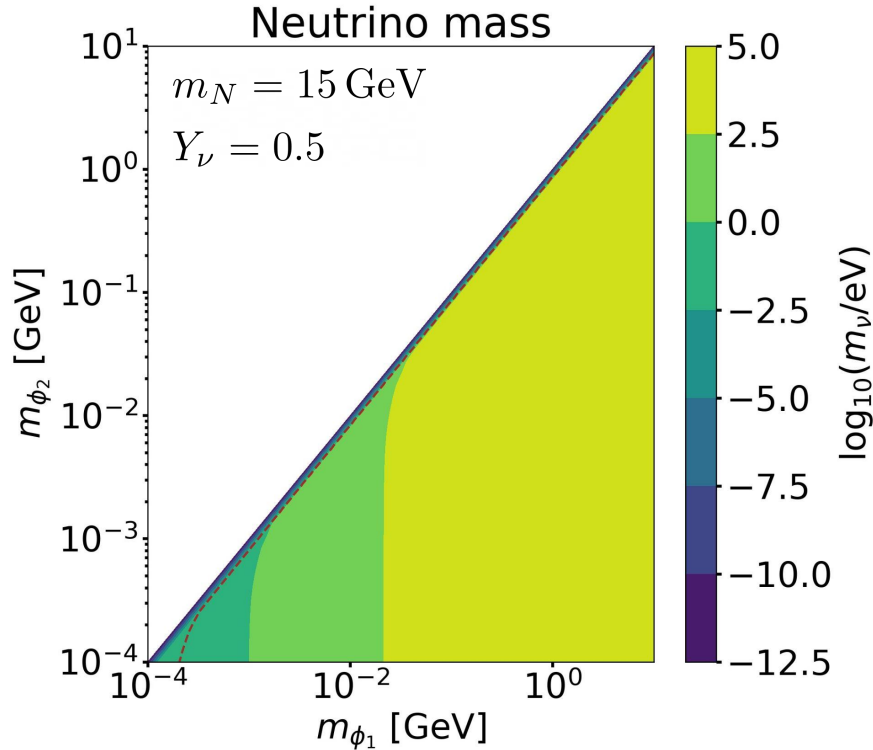


Mass of the lightest neutrino

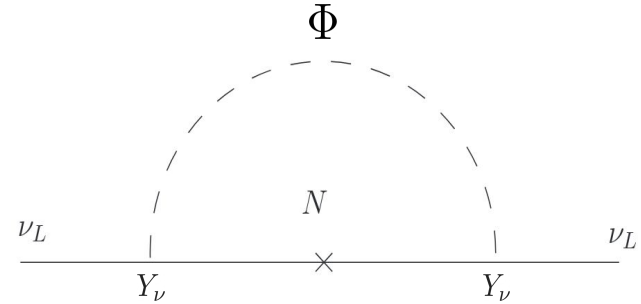
$$m_{\nu_0} < 0.037 \text{ eV}$$

95%, Normal Ordering. The GAMBIT Cosmology Workgroup. (2021). arXiv: 2009.03287

Fourth test: Adding neutrino masses

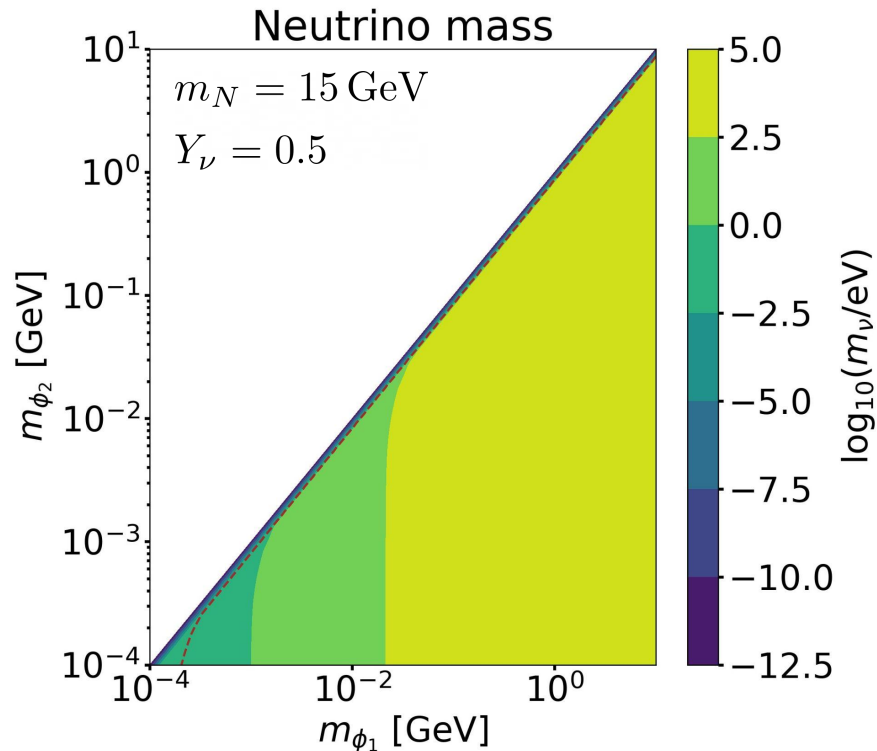


$$m_\nu = \frac{Y_\nu^2}{32\pi^2} m_N \left[\frac{m_{\phi_1}^2}{(m_N^2 - m_{\phi_1}^2)} \ln \left(\frac{m_N^2}{m_{\phi_1}^2} \right) - \frac{m_{\phi_2}^2}{(m_N^2 - m_{\phi_2}^2)} \ln \left(\frac{m_N^2}{m_{\phi_2}^2} \right) \right]$$



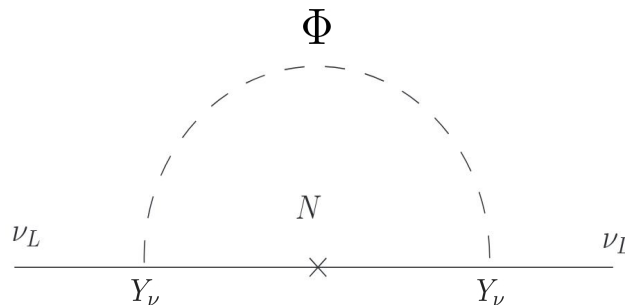
Böhm, C. et al. (2007). arXiv: 0612228

Fourth test: Adding neutrino masses



$$m_\nu = \frac{Y_\nu^2}{32\pi^2} m_N \left[\frac{m_{\phi_1}^2}{(m_N^2 - m_{\phi_1}^2)} \ln \left(\frac{m_N^2}{m_{\phi_1}^2} \right) - \frac{m_{\phi_2}^2}{(m_N^2 - m_{\phi_2}^2)} \ln \left(\frac{m_N^2}{m_{\phi_2}^2} \right) \right]$$

$$m_\nu \approx m_{\phi_1}^2 - m_{\phi_2}^2$$



Böehm, C. et al. (2007). arXiv: 0612228

Final Test: Combining All Constraints

Supernova relic neutrino flux

$$\text{SK} - \bar{\nu}_e, \langle \sigma v \rangle$$

$$27 < m_\varphi < 30 \text{ MeV}$$

Argüelles, C.A. et al. (2021). arXiv: 1912.09486

BBN primordial abundances

$$N_{\text{eff}}, m_\varphi \leq 20 \text{ MeV}$$

Obtained using AlterBBN

CMB & Matter Power Spectrum

$$\sigma_{\text{DM}-\nu,0} \lesssim 10^{-45} (m_{\text{DM}}/\text{GeV}) \text{ cm}^2$$

$$\sigma_{\text{DM}-\nu} \propto T^2, 10^{+13} u_0 < 2.56$$

Wilkinson, R.J. et al. (2014). arXiv: 1401.7597

Relic abundance

$$\Omega_c h^2 = 0.1200 \pm 0.0012$$

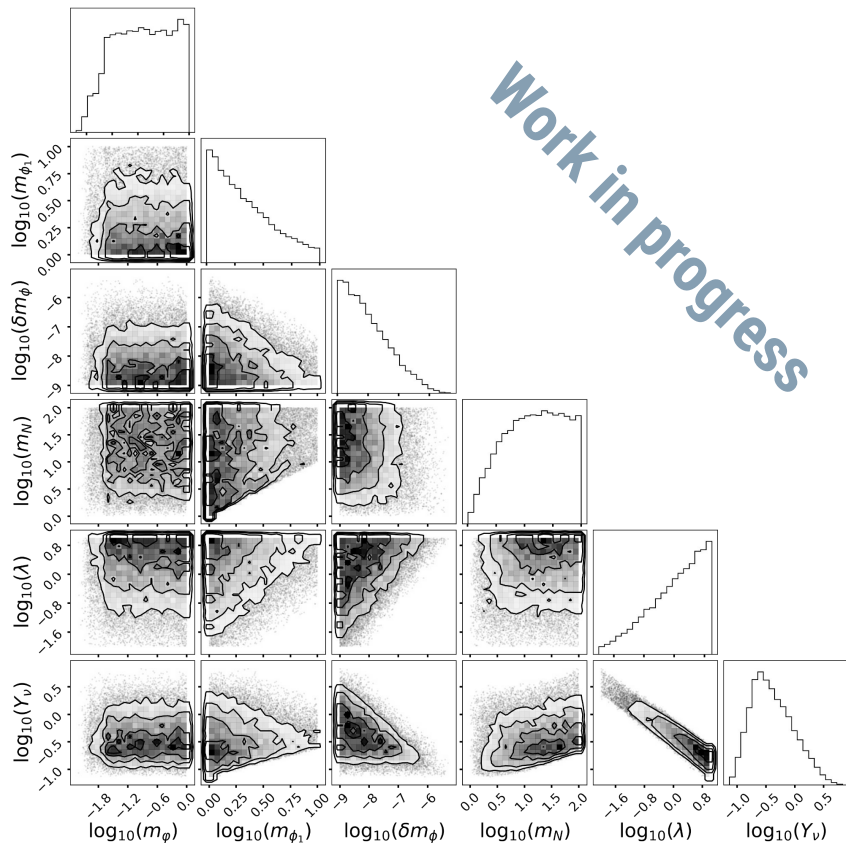
(68%, Planck TT,TE,EE + lowE + lensing)

Mass of the lightest neutrino

$$m_{\nu_0} < 0.037 \text{ eV}$$

95%, Normal ordering. The GAMBIT Cosmology Workgroup. (2021). arXiv: 2009.03287

Final Test: Combining All Constraints



Preliminary MCMC

emcee run in the Frontenac cluster

Prior $m_N > m_\phi > m_\varphi$

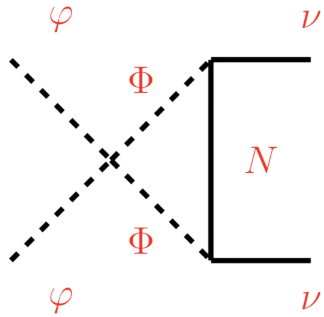
$$\delta m_\phi = 1 - \frac{m_{\phi_2}}{m_{\phi_1}}$$

Strongest constraints

$$N_{\text{eff}} + m_{\nu_0} + \Omega_\varphi h^2$$

More tests to be done!

Did this DM model escape the room?



Relic
abundance

$$\Omega_{\phi} h^2$$

Annihilation
& scattering

$$N_{\text{eff}}, u$$

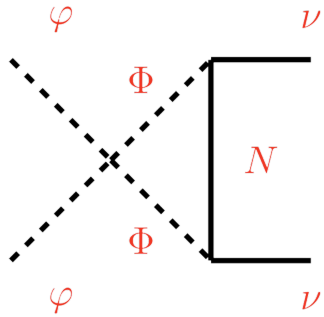
Extragalactic
annihilation

$$\langle \sigma v \rangle$$

Lightest ν
mass

$$m_{\nu_0}$$

Did this DM model escape the room?



Relic
abundance

$$\Omega_{\phi} h^2$$

Annihilation
& scattering

$$N_{\text{eff}}, u$$

Extragalactic
annihilation

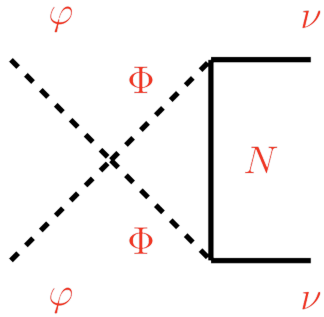
$$\langle \sigma v \rangle$$

Lightest ν
mass

$$m_{\nu_0}$$



Did this DM model escape the room?



Relic
abundance

$$\Omega_\phi h^2$$



Annihilation
& scattering

$$N_{\text{eff}}, u$$



Extragalactic
annihilation

$$\langle \sigma v \rangle$$

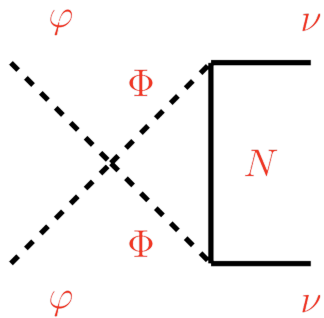


Lightest ν
mass

$$m_{\nu_0}$$



Did this DM model escape the room?



Relic
abundance

$$\Omega_{\phi} h^2$$



Annihilation
& scattering

$$N_{\text{eff}}, u$$



Extragalactic
annihilation

$$\langle \sigma v \rangle$$



Lightest ν
mass

$$m_{\nu_0}$$



Future work:

- MCMC for + masses, constraints & errors - Allow for DM underabundance
- Statistical analysis: profile likelihood ratios & p-values



Thank you!

Questions?

Collaborators: Aaron Vincent (Queen's) and Gopolang Mohlabeng (UCI)

E-mail: karen.maciascardenas@queensu.ca

BACKUP SLIDES

Squared Amplitude

FeynRules

$$-\mathcal{L}_{\text{int}} = \frac{1}{2}m_\varphi^2\varphi^2 + \frac{1}{2}\lambda\Phi^2\varphi^2 + Y_\nu\bar{\nu}_L\Phi N_R + \text{h.c.}$$

Squared amplitude with
FormCalc / FeynCalc

$$|\overline{\mathcal{M}}|^2 = \frac{\lambda^2 m_N^2 s Y_\nu^4 C_0(0, 0, s, m_\Phi^2, m_N^2, m_\Phi^2)^2}{64\pi^4}$$

Analytical form of the
scalar Passarino-Veltman
integral with Package-X

$$C_0(0, 0, s, m_\Phi^2, m_N^2, m_\Phi^2) =$$

$$\frac{\text{DiLog}\left(\frac{2(m_N^2 - m_\Phi^2)}{2m_N^2 - \sqrt{s(s - 4m_\Phi^2)} - 2m_\Phi^2 + s}, s\right)}{s} + \frac{\text{DiLog}\left(\frac{2(m_N^2 - m_\Phi^2)}{2m_N^2 + \sqrt{s(s - 4m_\Phi^2)} - 2m_\Phi^2 + s}, s\right)}{s}$$

$$- \frac{\text{DiLog}\left(\frac{2(m_N^2 - m_\Phi^2 + s)}{2m_N^2 - \sqrt{s(s - 4m_\Phi^2)} - 2m_\Phi^2 + s}, s\right)}{s} - \frac{\text{Li}_2\left(\frac{2(m_N^2 - m_\Phi^2 + s)}{2m_N^2 - 2m_\Phi^2 + s + \sqrt{s(s - 4m_\Phi^2)}}\right)}{s}$$

$$+ \frac{\text{Li}_2\left(\frac{(m_N^2 - m_\Phi^2)(m_N^2 - m_\Phi^2 + s)}{m_N^4 - 2m_\Phi^2 m_N^2 + s m_N^2 + m_\Phi^4}\right)}{s} - \frac{\text{Li}_2\left(\frac{(m_N^2 - m_\Phi^2)^2}{m_N^4 - 2m_\Phi^2 m_N^2 + s m_N^2 + m_\Phi^4}\right)}{s}$$

Cross-section and relic abundance approximations

s-wave cross-section

Wells, J.D. (1994). arXiv: 940219

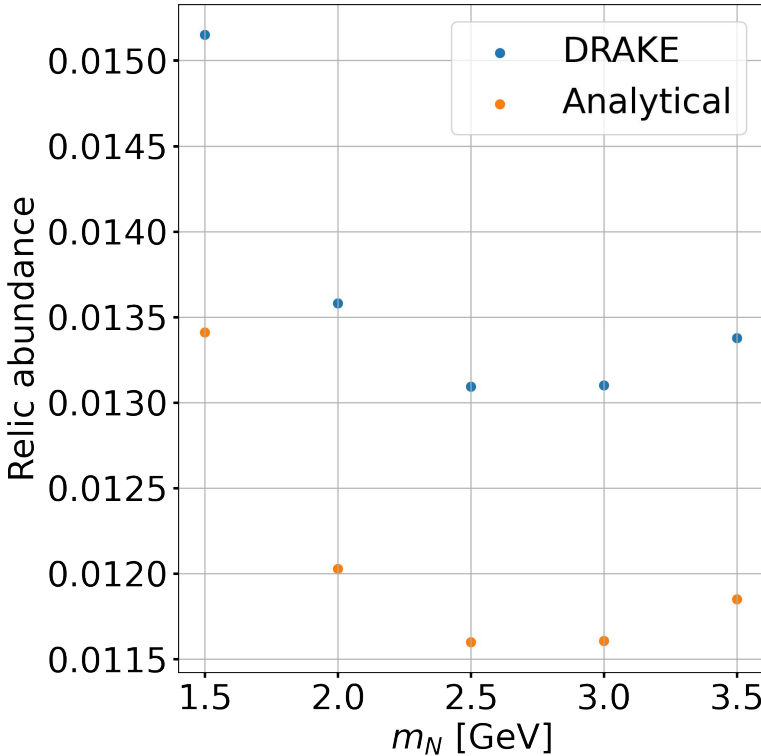
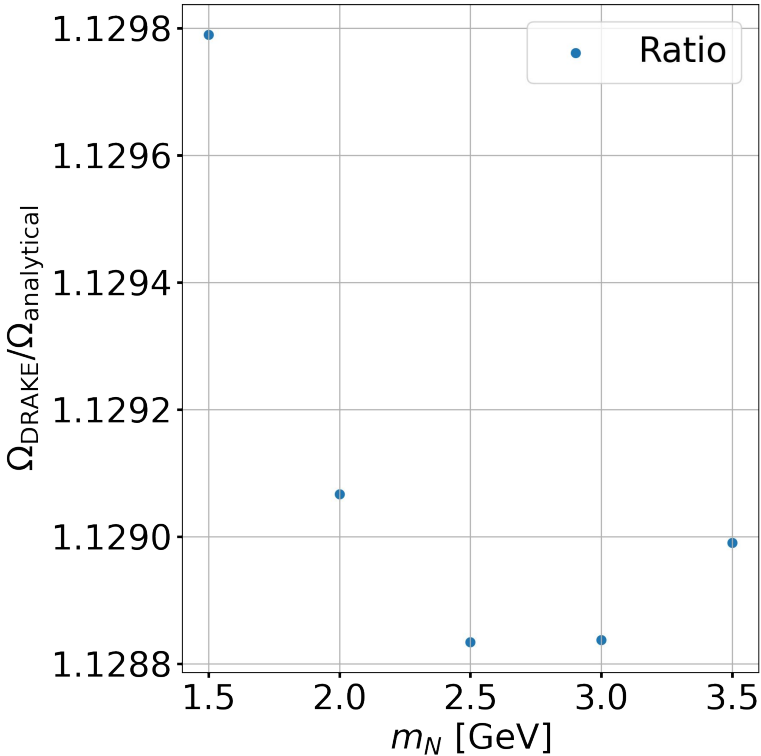
$$\langle\sigma v\rangle \approx \frac{\lambda^2 Y_\nu^4 m_N^2}{512\pi^5} |C_0(0, 0, 4m_\varphi^2, m_\Phi^2, m_N^2, m_\Phi^2)|^2$$

DM relic abundance approximation

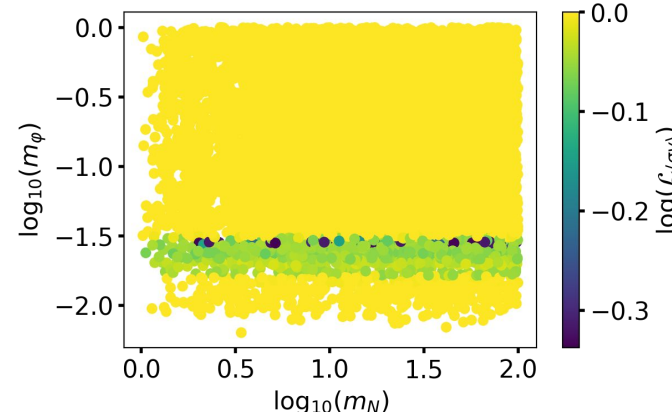
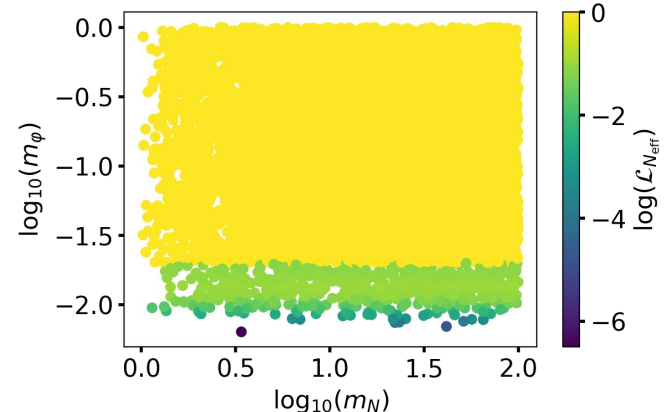
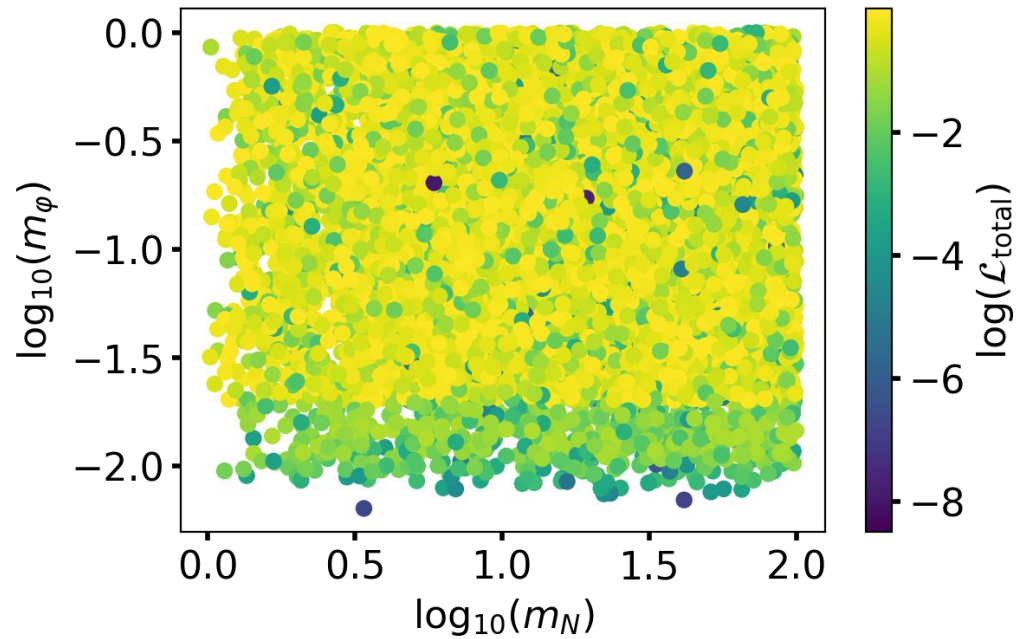
Steigman, G. et al. (2012). arXiv: 1204.3622

$$\Omega_c h^2 = \frac{9.92 \times 10^{-28}}{\langle\sigma v\rangle} \left(\frac{x_*}{g_*^{1/2}}\right) \left(\frac{(\Gamma/H)_*}{1 + \alpha_*(\Gamma/H)_*}\right)$$

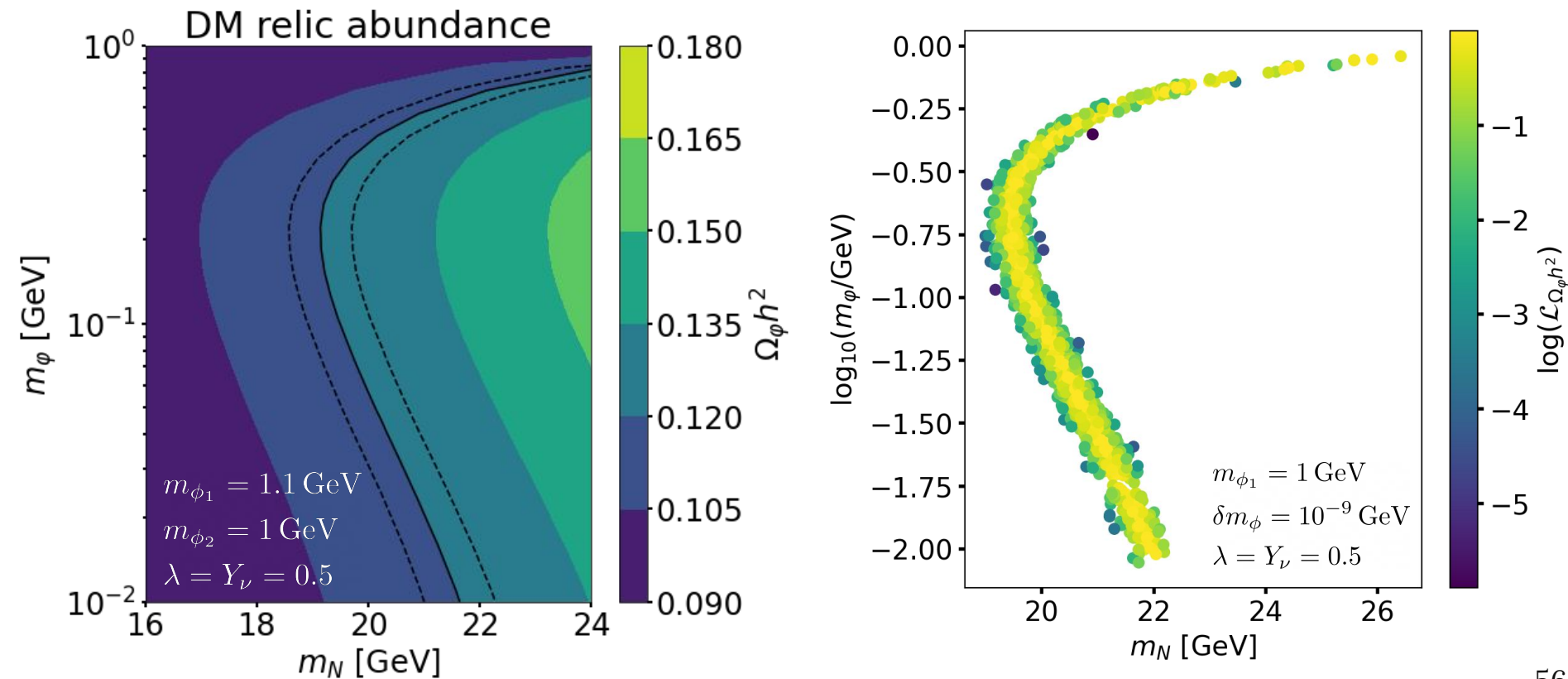
Comparison to the Boltzmann code DRAKE



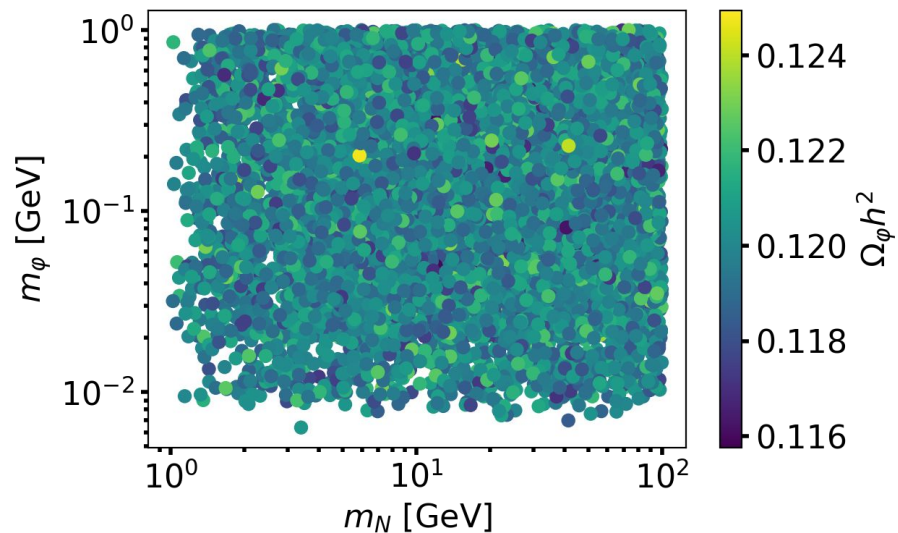
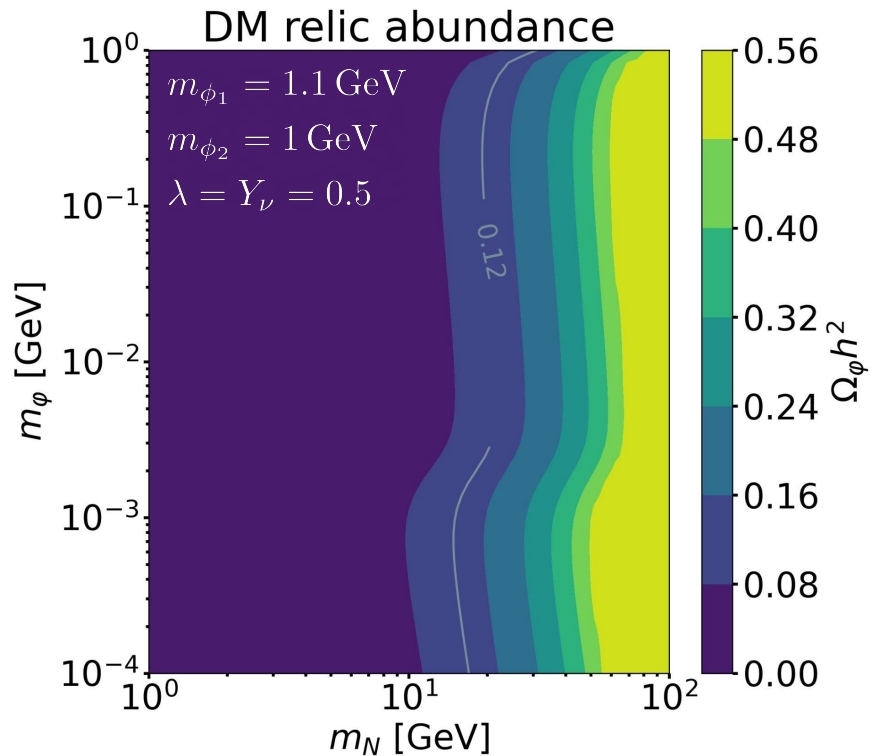
DM mass vs heavy neutrino mass likelihoods



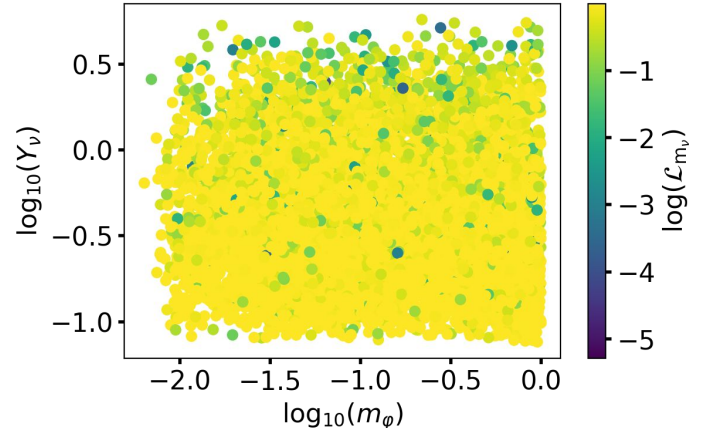
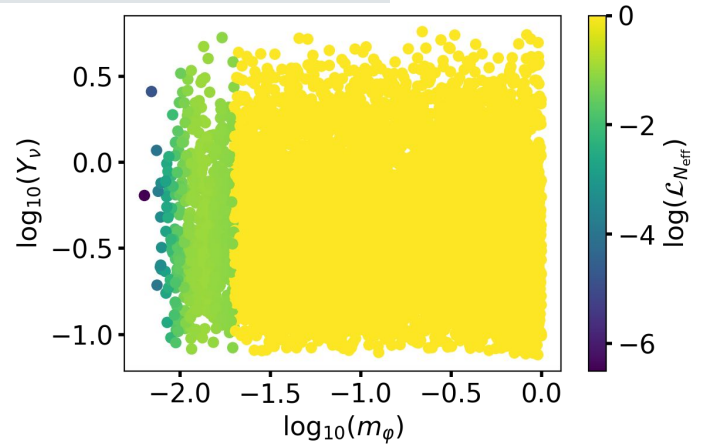
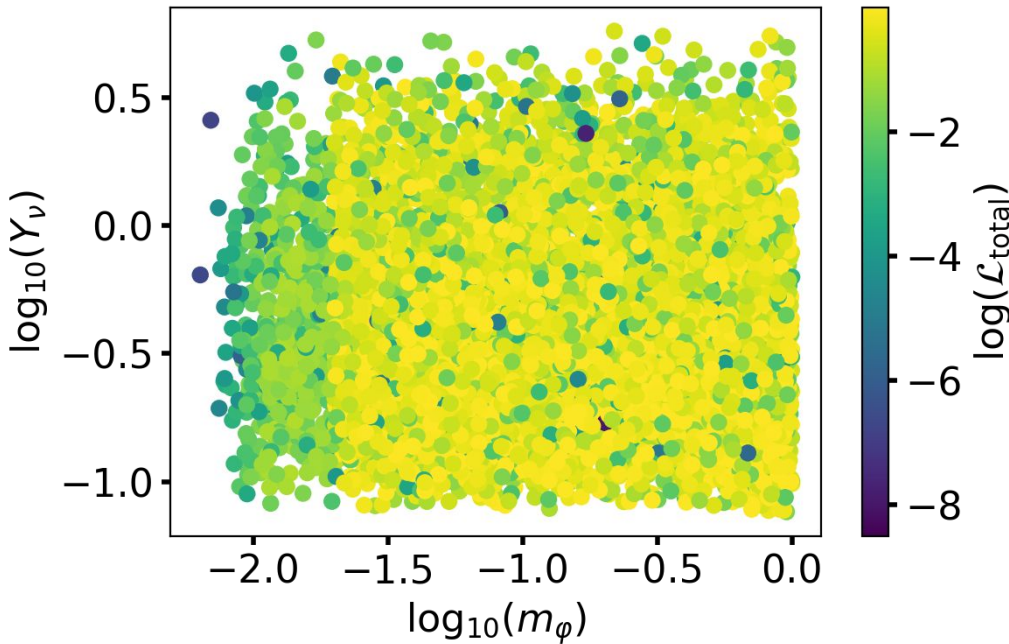
2D Relic Abundance Likelihood



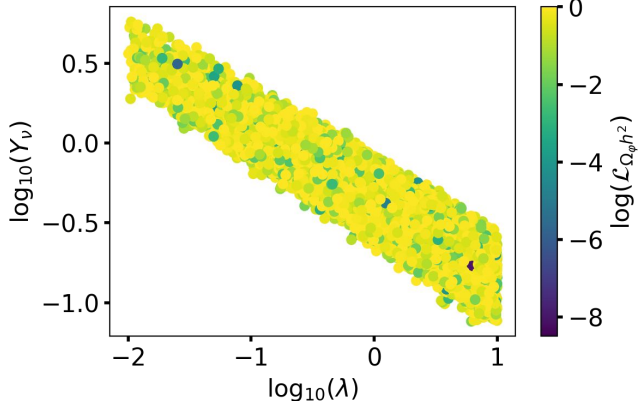
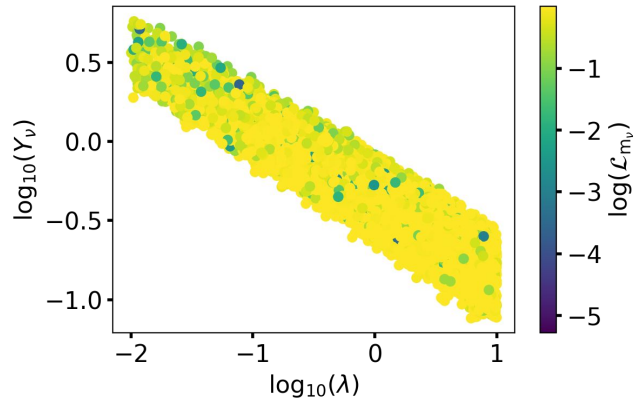
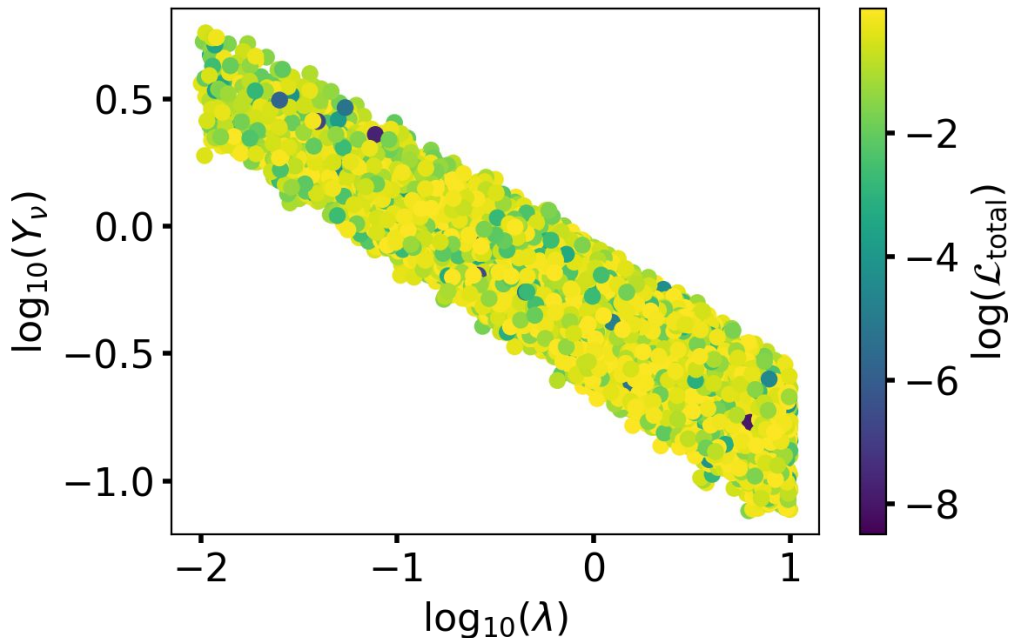
Relic abundance for fixed + varying parameters



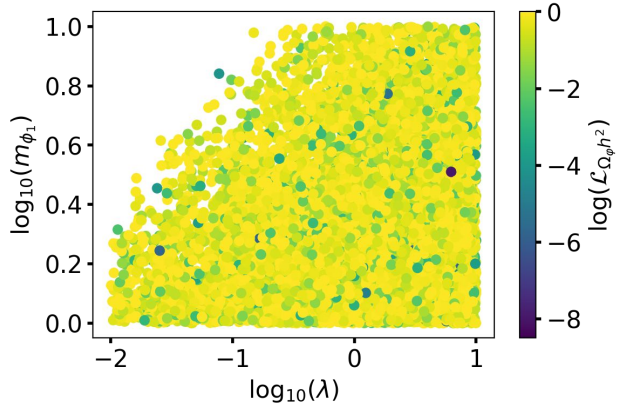
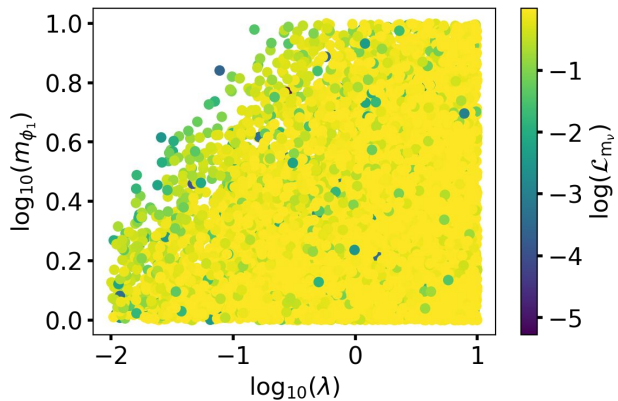
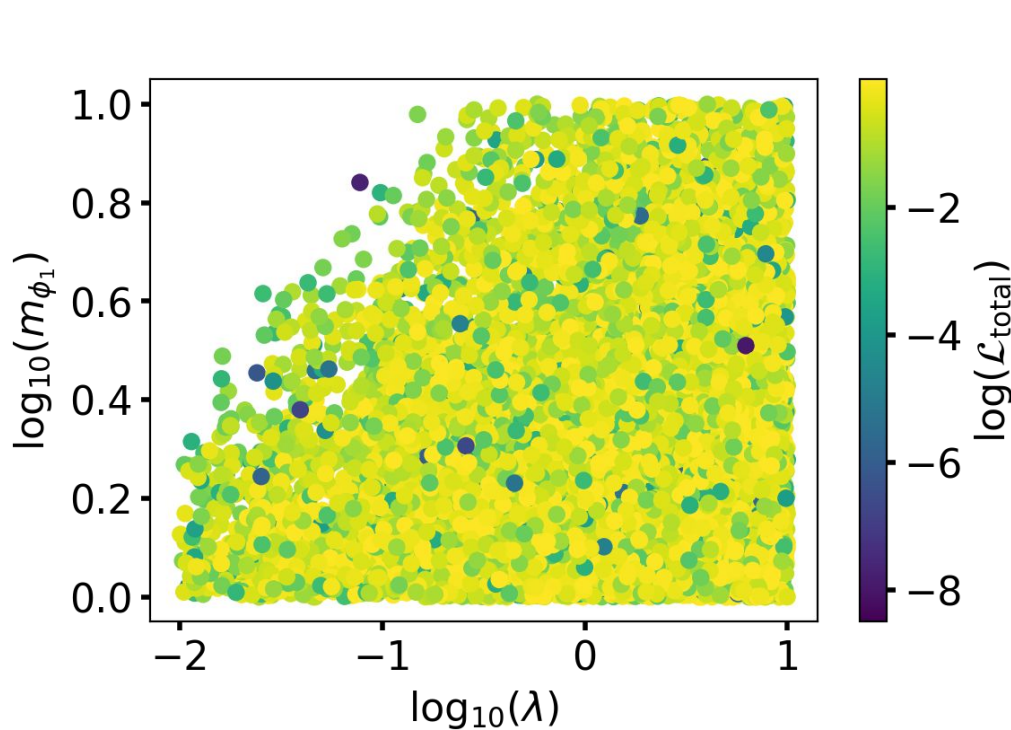
Yukawa coupling vs DM mass likelihoods



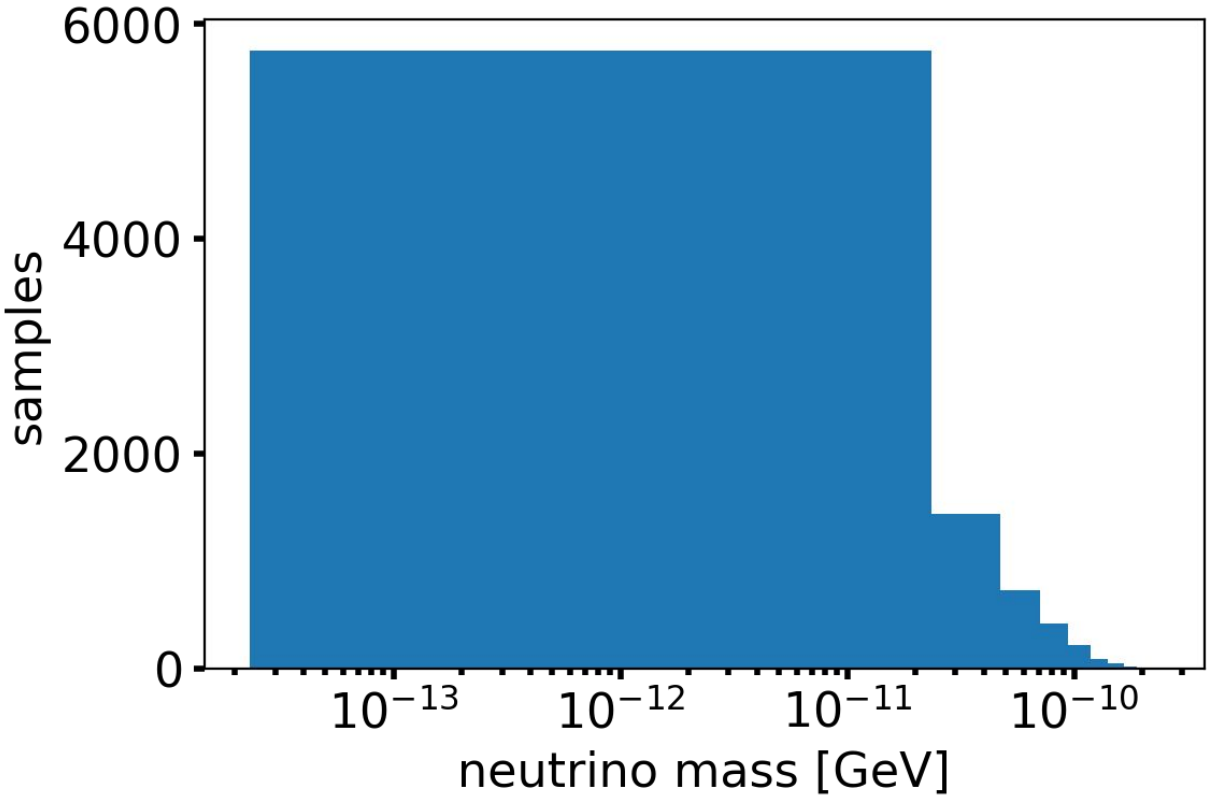
Yukawa coupling vs λ coupling likelihoods



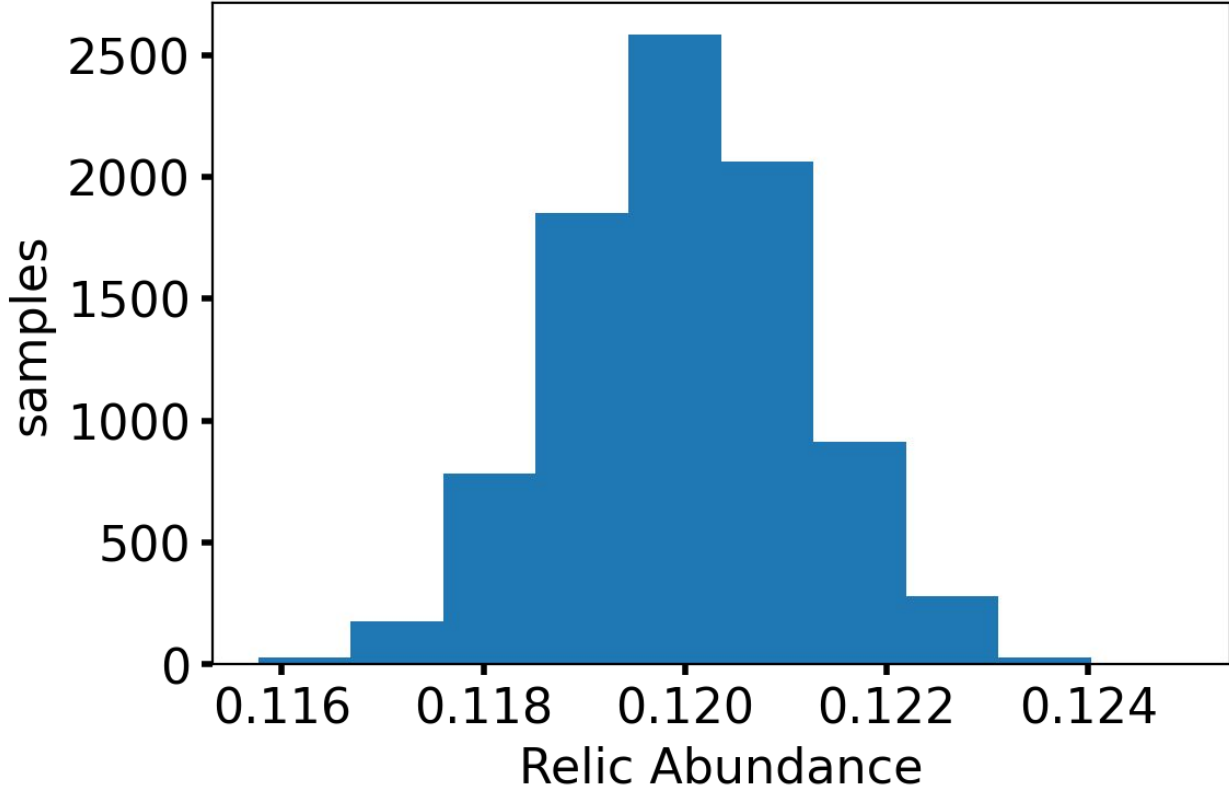
Scalar mediator mass vs λ coupling likelihoods



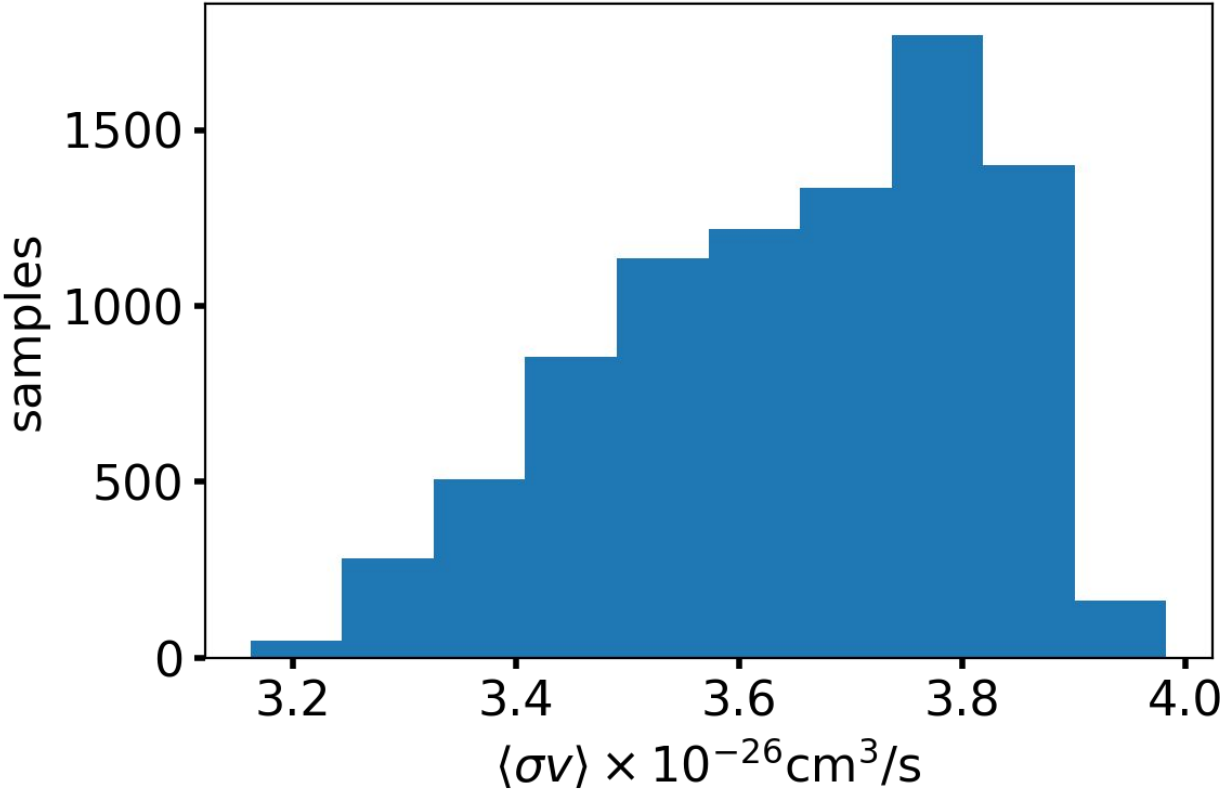
Neutrino mass sampling



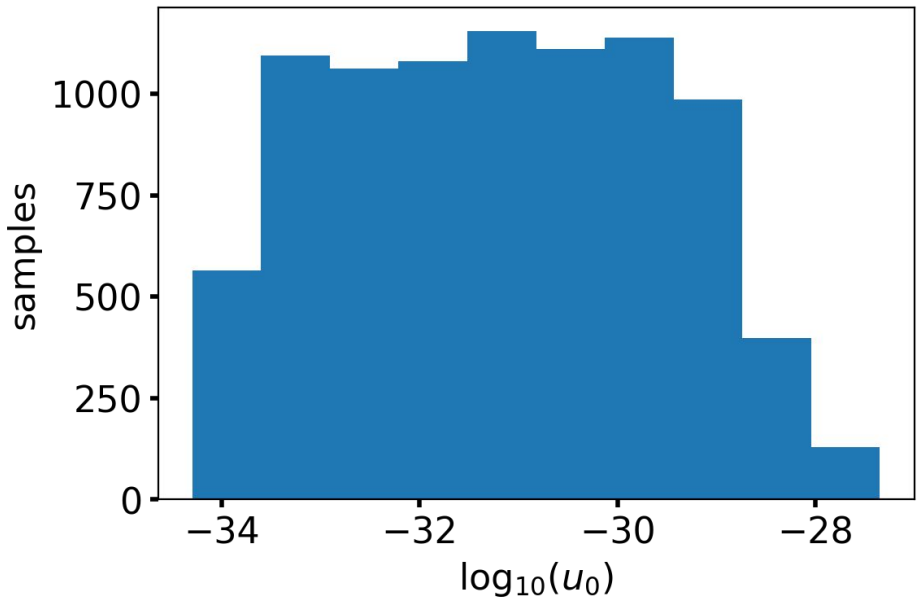
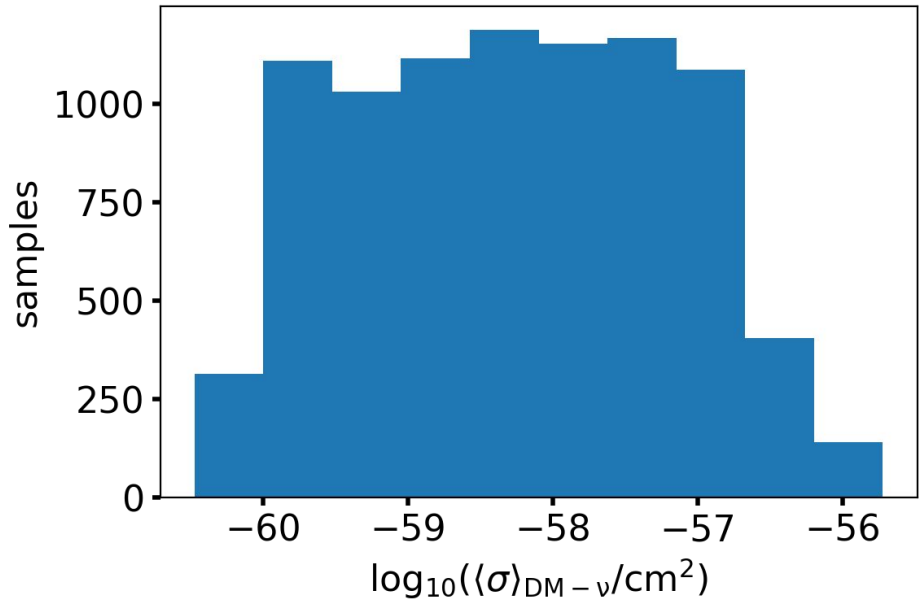
Relic abundance sampling



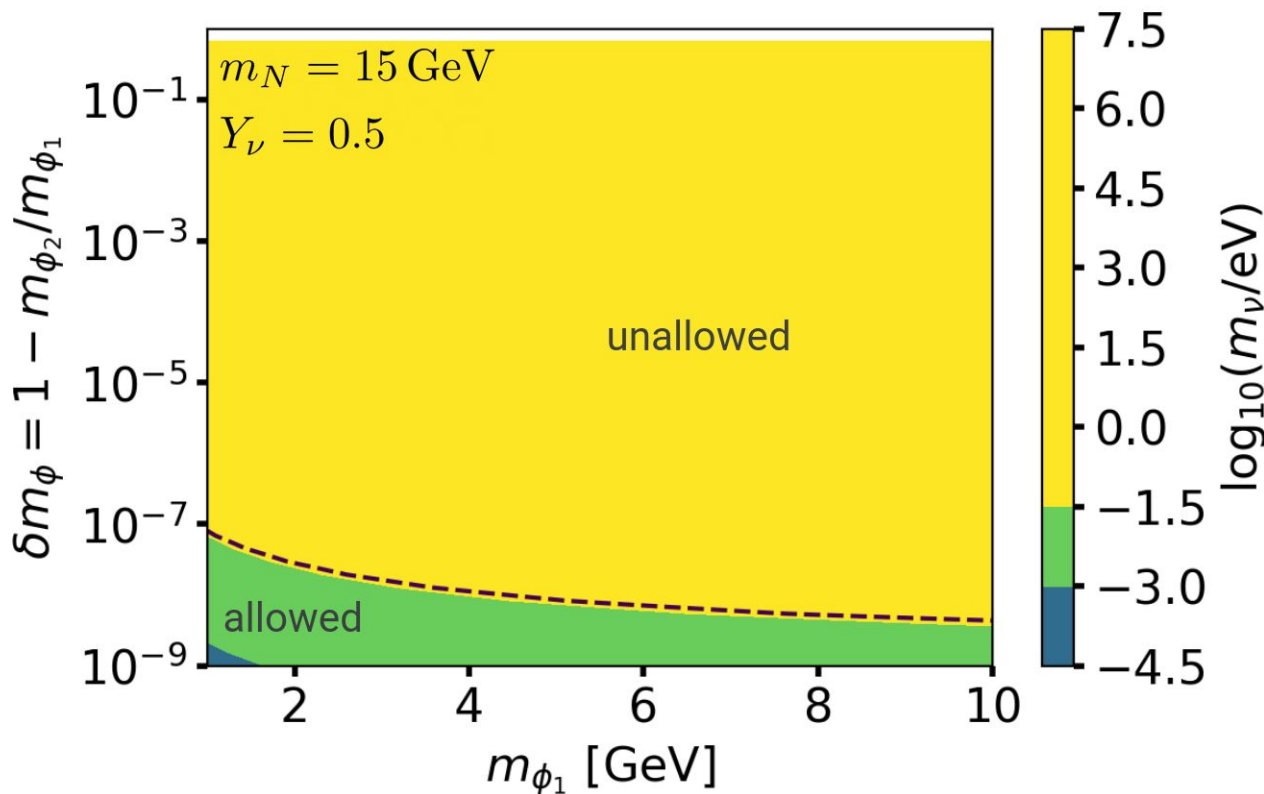
Annihilation cross-section sampling



Scattering cross-section sampling



Neutrino mass parameter space



Structure formation

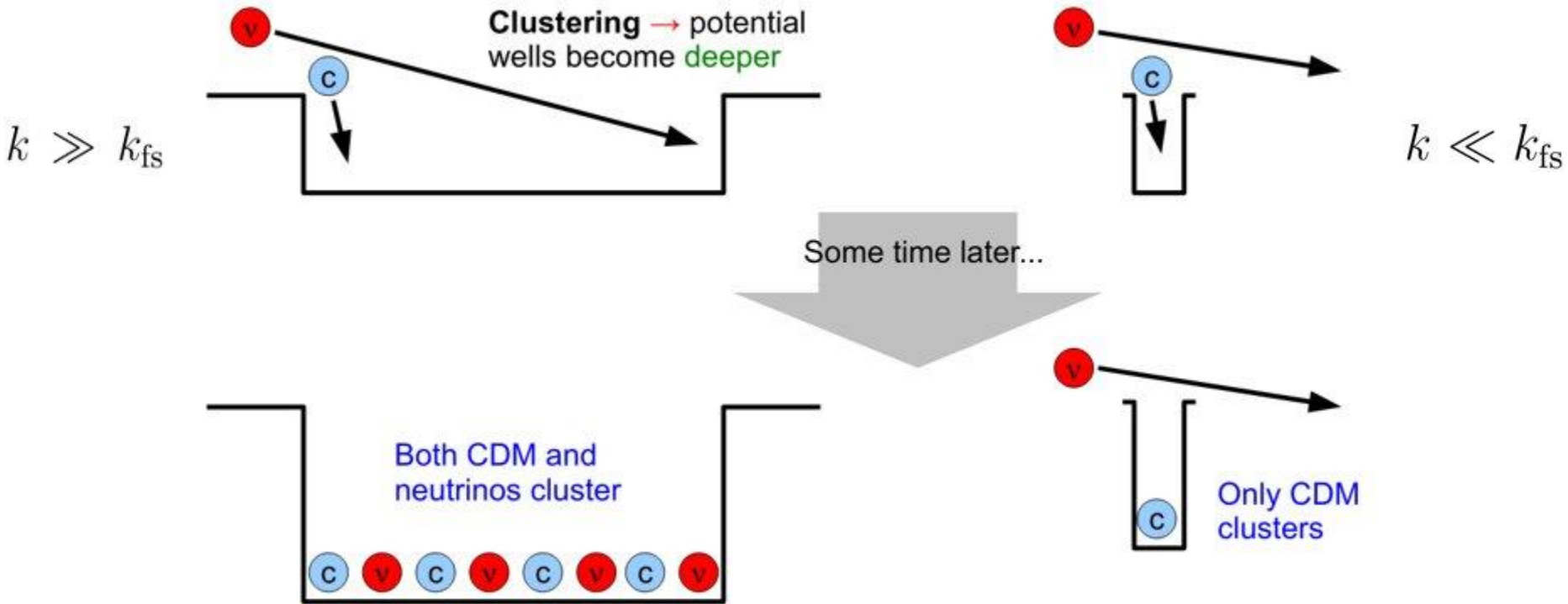


Diagram by Yvonne Y. Y. Wong