



# Constraining Atomic Dark Matter with Cosmological Observables

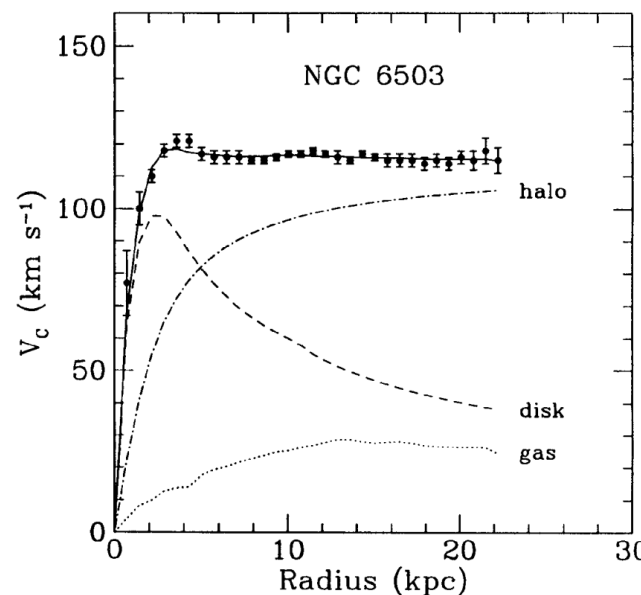
Jared Barron

Based on work in progress with Saurabh Bansal, David Curtin, Yuhsin Tsai

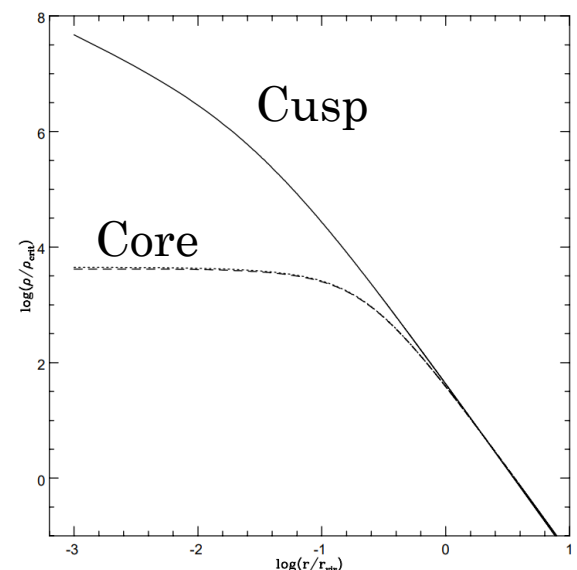
August 11, 2022

# Dark Matter: CDM or something more?

- The cold, collisionless dark matter paradigm explains many observations extremely well.
  - Galaxy rotation curves.
  - The Bullet cluster.
  - The cosmic microwave background spectrum.
  - Large scale structure in the universe.
- But it's not perfect.
  - Core vs cusp problem.
  - Missing satellites.
  - Hubble tension.
- Could a component of dark matter have non-trivial interactions in the dark sector?



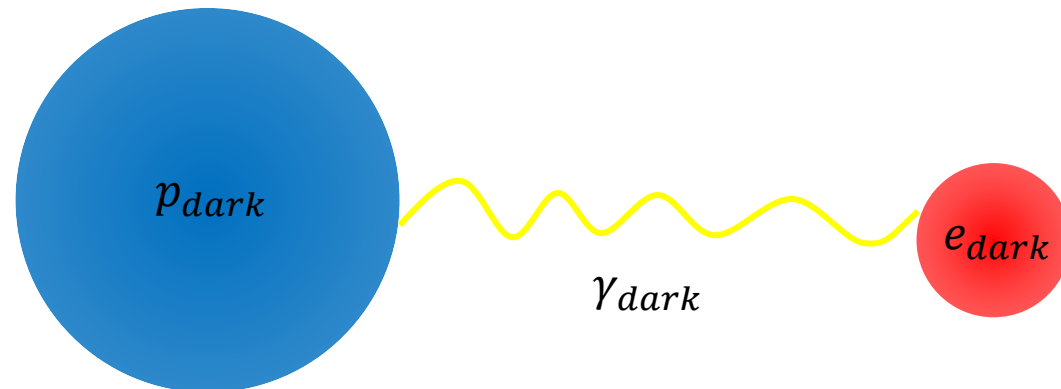
Freese, 2008



Popolo, 2009

# Atomic Dark Matter

- Assume some fraction  $\hat{r} \equiv \frac{\Omega_{ADM}}{\Omega_{DM}}$  of the dark matter is coupled to dark radiation in the early universe.
- This dark matter cannot contribute to growth of density perturbations in the early universe until the dark matter and dark radiation decouple.
- Such a dark sector can also exhibit dissipative behaviour, affecting structure growth on galactic scales.
- Introduce a dark sector with fermions  $p_{dark}$  and  $e_{dark}$  with masses  $m_{p'}$ ,  $m_{e'}$ , neutral under the SM gauge group but oppositely charged under a dark  $U(1)$  gauge force with dark fine structure constant  $\alpha_D$ .
- Assume  $m_{p'} \gg m_{e'}$ .
- Bound state is dark 'hydrogen'.
- No assumption about coupling to SM.



# Why Atomic Dark Matter?

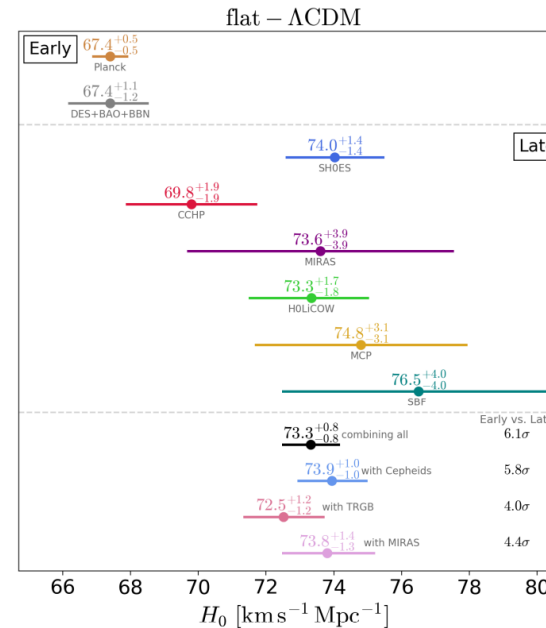
## Theory

- Approximately  $\mathbb{Z}_2$ -symmetric mirror sectors can address the little hierarchy problem.
- e.g. Mirror Twin Higgs.
- The visible sector has atoms – why shouldn't the dark matter?

SM | M2

## Experiment

- Atomic dark matter interactions can address the core-cusp problem.
- Additional dark radiation can ameliorate the Hubble tension.



Verde,  
Treu,  
Riess,  
2019



# Cosmology of Atomic Dark Matter

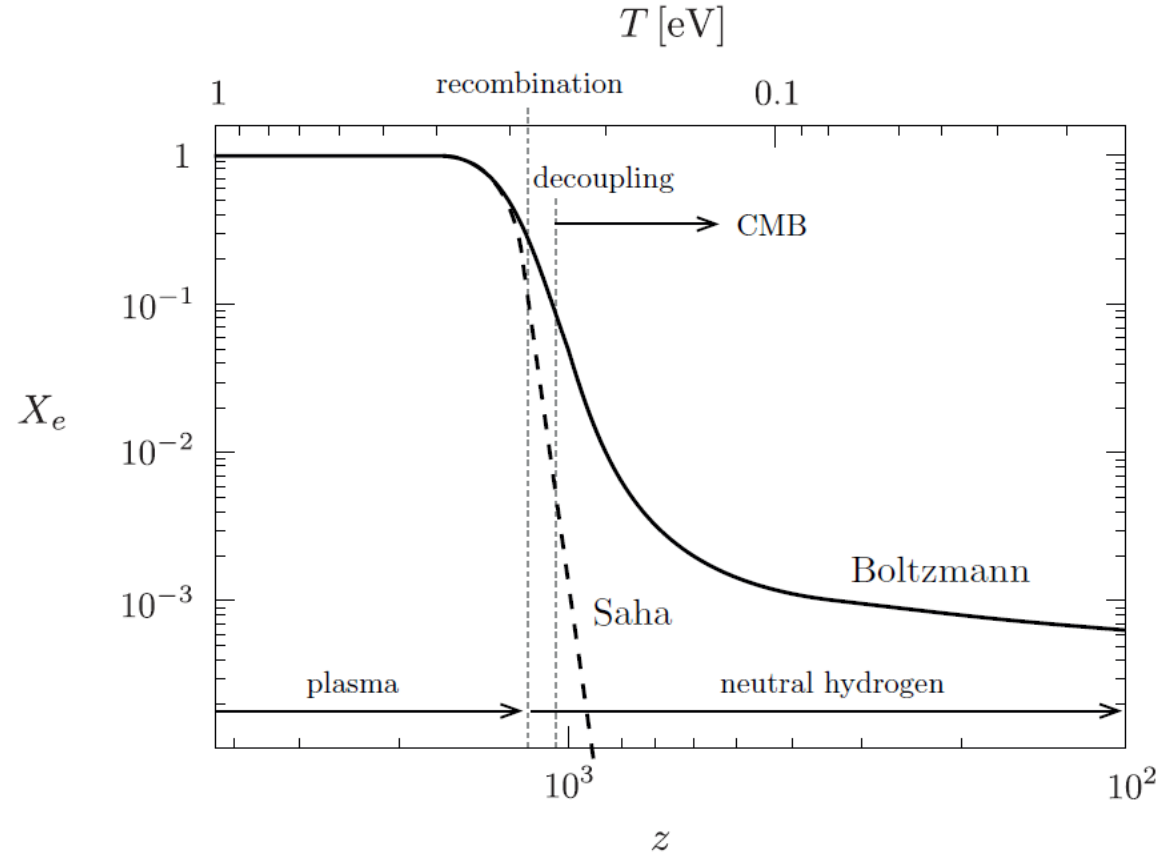
- Unlike CDM, atomic dark matter can undergo **dark recombination** and **dark acoustic oscillations**.
- To comply with bounds on  $\Delta N_{\text{eff}} < 0.3$ , the dark sector must be cold.

$$\xi \equiv \left( \frac{T_D}{T_{SM}} \right) = \left( \frac{7}{8} \frac{4}{11} \Delta N_{\text{eff}} \right)^{\frac{1}{4}} < 0.5$$

- Even accounting for the altered expansion history due to extra relativistic degrees of freedom, ADM produces unique signatures in the matter power spectrum and CMB.
- Cosmology of ADM first systematically reviewed by Cyr-Racine and Sigurdson in *Phys.Rev.D* 87 (2013) 10, and constraints placed in *Phys.Rev.D* 89 (2014) 6.
- We will update constraints with higher precision cosmological observations, and cover the full three-dimensional ADM parameter space.

# (Dark) recombination and decoupling

- Dark recombination occurs after  $T_D$  falls below  $B_D = \frac{\alpha_D^2 m'_e}{2}$ .
- As free dark electron density drops, dark photons stop efficiently exchanging energy and momentum with the dark electrons.
- Dominant energy exchange mechanism is Thomson scattering, with cross-section  $\sigma'_T = \frac{8\pi}{3} \left(\frac{\alpha_D}{m'_e}\right)^2$ .
- When scattering rate  $\Gamma_T \approx n'_e \sigma'_T < H$ , the dark photons decouple from the dark electrons and baryons.



SM evolution of electron ionization fraction during recombination.

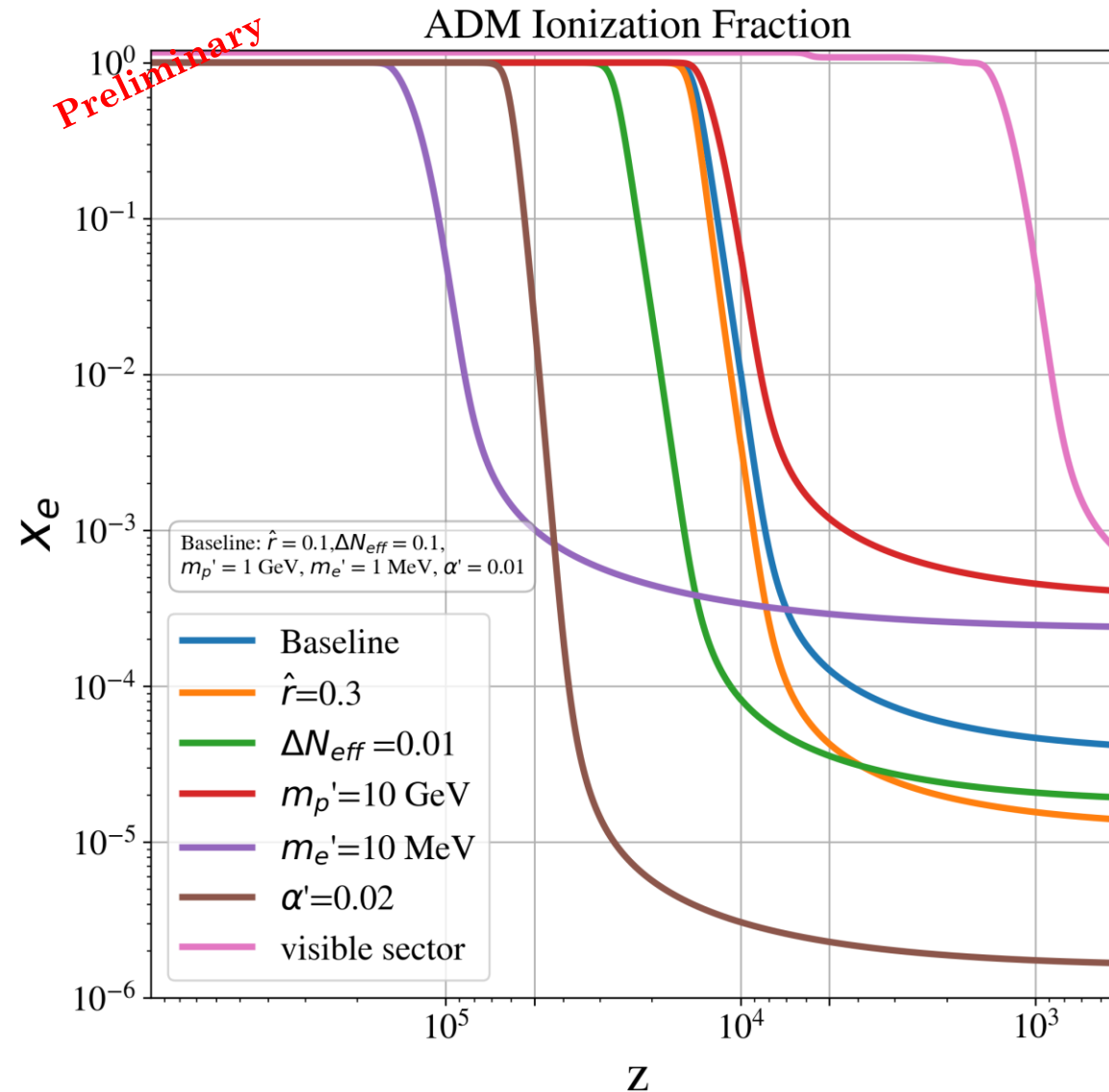


# Implementation of ADM in CLASS

- The Cosmological Linear Anisotropy Solving System (CLASS) is a numerical code that produces CMB and LSS spectra given cosmological parameter inputs.
- Modified to include the dark radiation and dark hydrogen, and solve for the dark recombination.
  - Built on top of modifications by Bansal et al. for MTH. (hep-ph:2110.04317)
- After dark recombination, the dark sector perturbation evolutions are handled by the Effective Theory of Structure (ETHOS) framework in CLASS.

# Dark Recombination

- Higher  $\hat{r} \rightarrow$  lower  $x_e'$ .
  - Why: Higher  $e_{dark}$  abundance.
- Lower  $\Delta N_{eff} \rightarrow$  earlier recombination.
  - Why: Higher  $T_{SM}$  for same  $T_D$ .
- Higher  $m_p' \rightarrow$  higher  $x_e'$ .
  - Why: Lower  $e_{dark}$  abundance.
- Higher  $m_e' \rightarrow$  earlier recombination and higher  $x_e'$ .
  - Why: Higher  $B_D$  and lower  $\sigma_T$ .
- Higher  $\alpha' \rightarrow$  earlier recombination and lower  $x_e'$ .
  - Why: Higher  $B_D$  and higher  $\sigma_T$ .



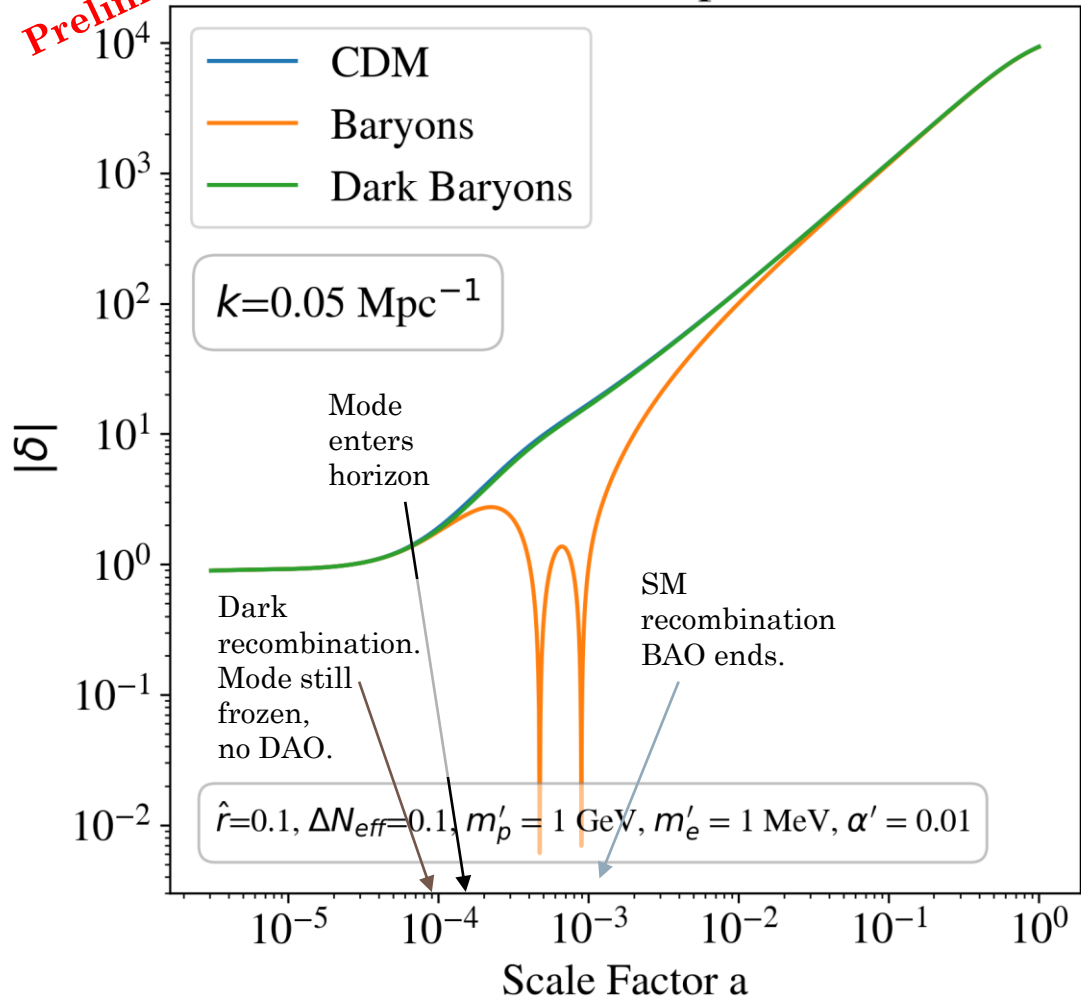


# Dark Acoustic Oscillations



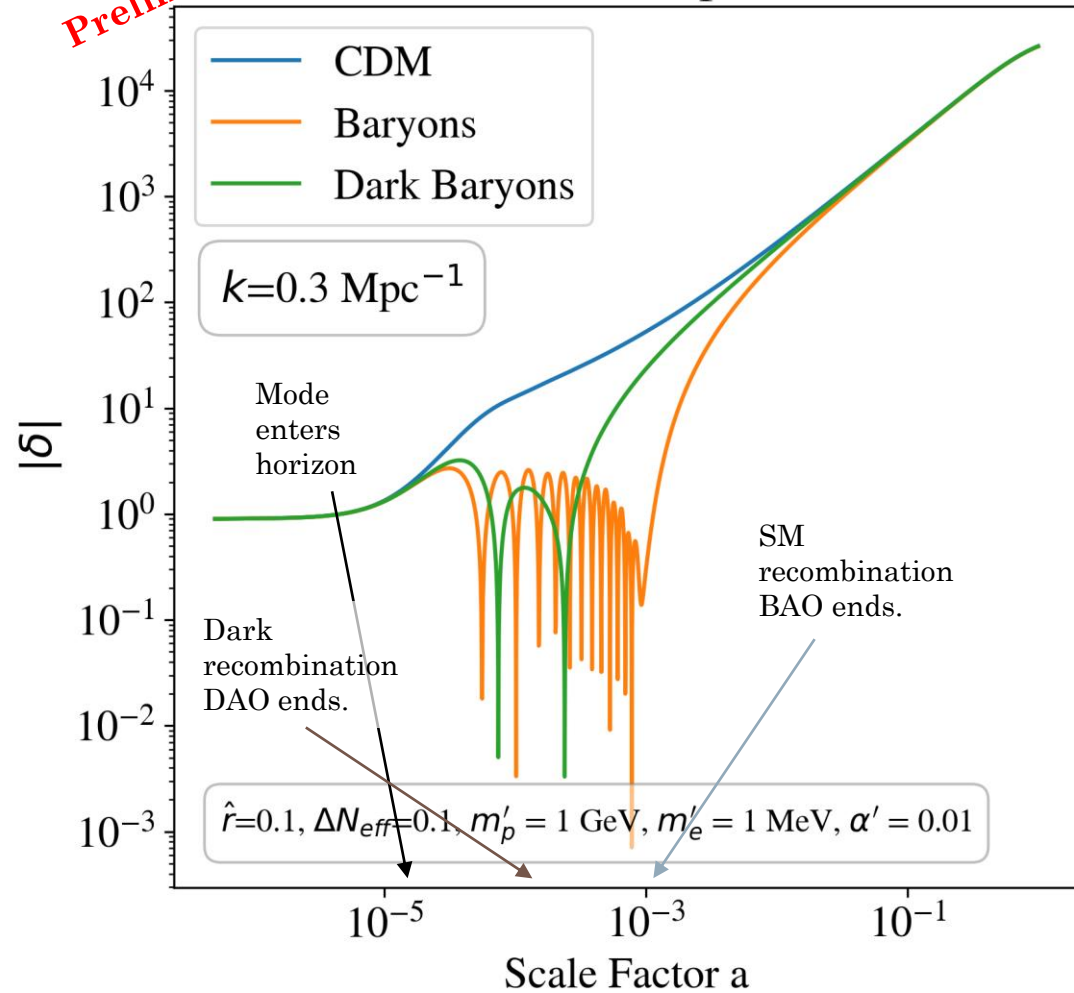
Preliminary

Evolution of matter perturbations



Preliminary

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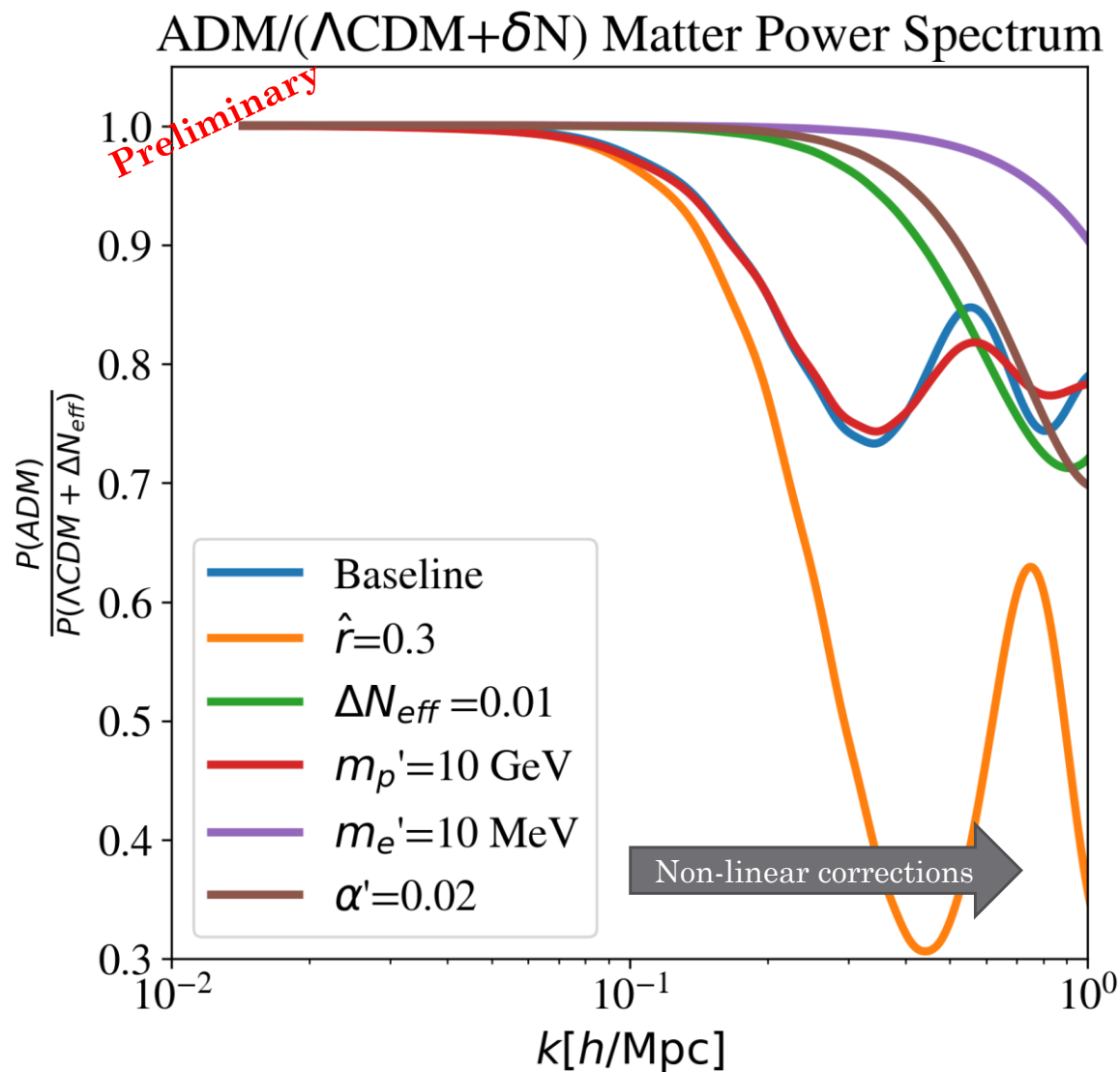
# Modification of Matter Power Spectrum

- Suppression and oscillations for  $k$  that enter horizon before dark decoupling.

- Higher  $\hat{r} \rightarrow$  more suppression.
- Why: Higher fraction of DM undergoes DAO.

- Lower  $\Delta N_{eff}$ , higher  $m_{e'}$ , higher  $\alpha' \rightarrow$  suppression turns on at higher  $k$ .
- Why: Low- $k$  modes enter after recombination, not suppressed.

- Cannot trust linear power spectrum above  $k = 0.1 h \text{ Mpc}^{-1}$ .



Baseline:  $\hat{r} = 0.1, \Delta N_{eff} = 0.1, m_p' = 1 \text{ GeV}, m_e' = 1 \text{ MeV}, \alpha' = 0.01$ .

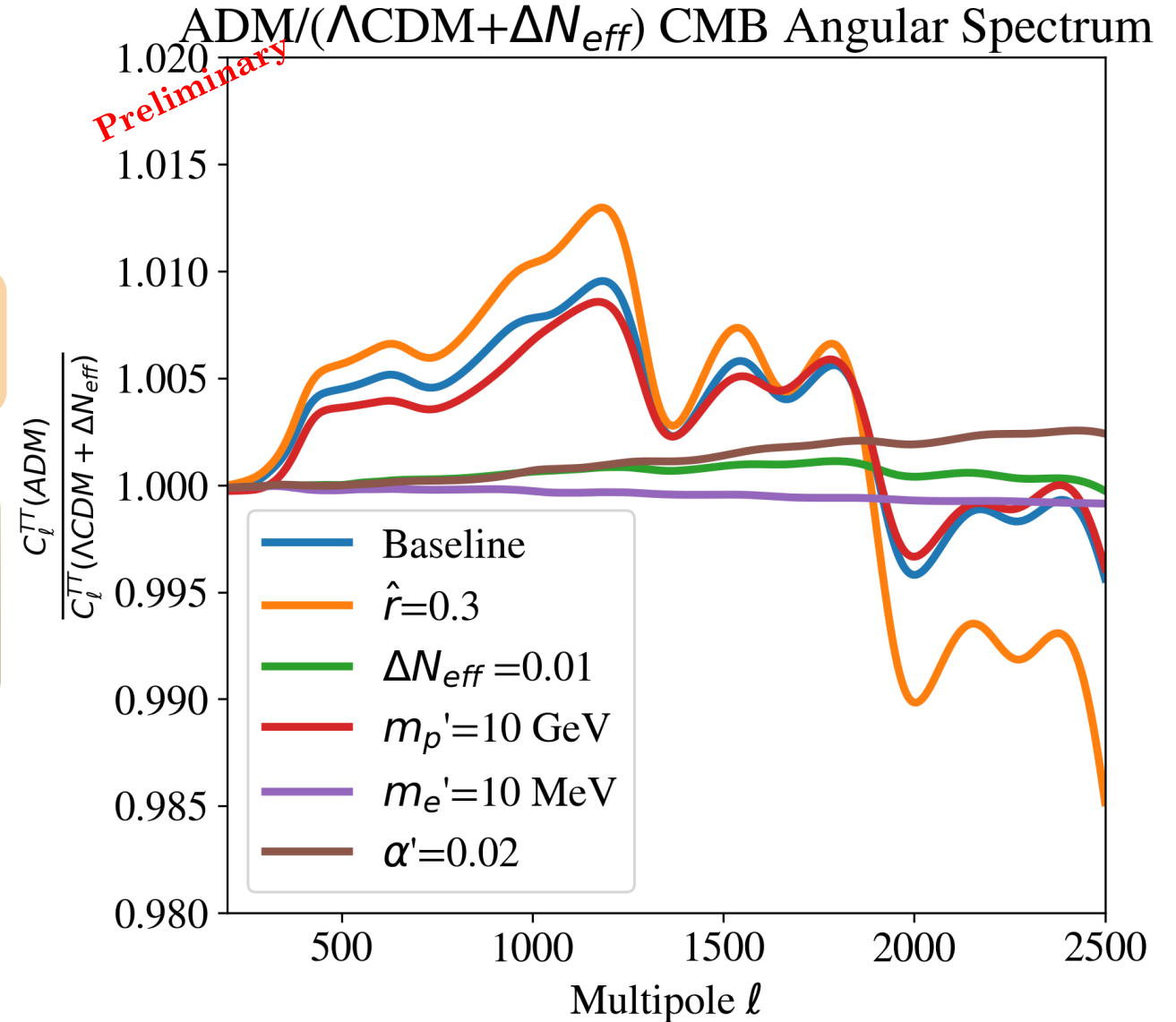
# Modification of CMB

- Complex pattern of modifications to spectrum.

- Higher  $\hat{r} \rightarrow$  Larger deviations.
- Why: Higher fraction of DM undergoes DAO.

- Lower  $\Delta N_{\text{eff}}$ , higher  $m_{e'}$ , higher  $\alpha' \rightarrow$  smaller deviations.
- Why: Dark photons began free-streaming earlier. Gravitational potential altered less.

- Polarization spectrum and lensing also altered.



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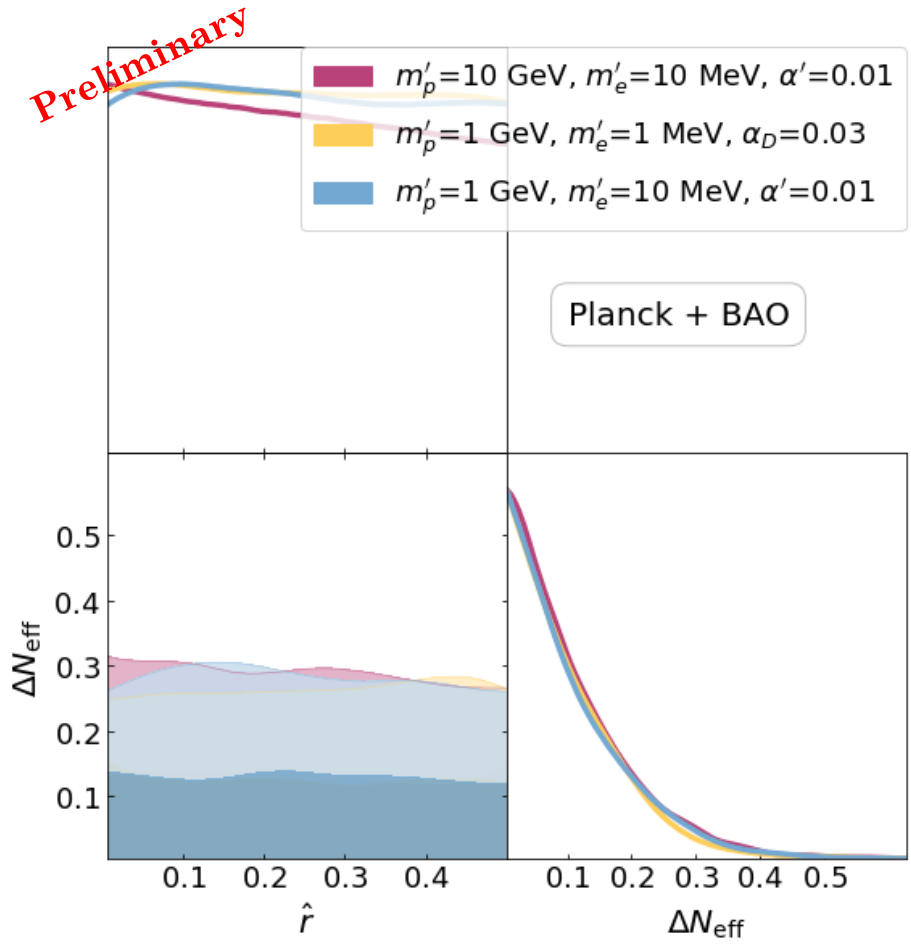
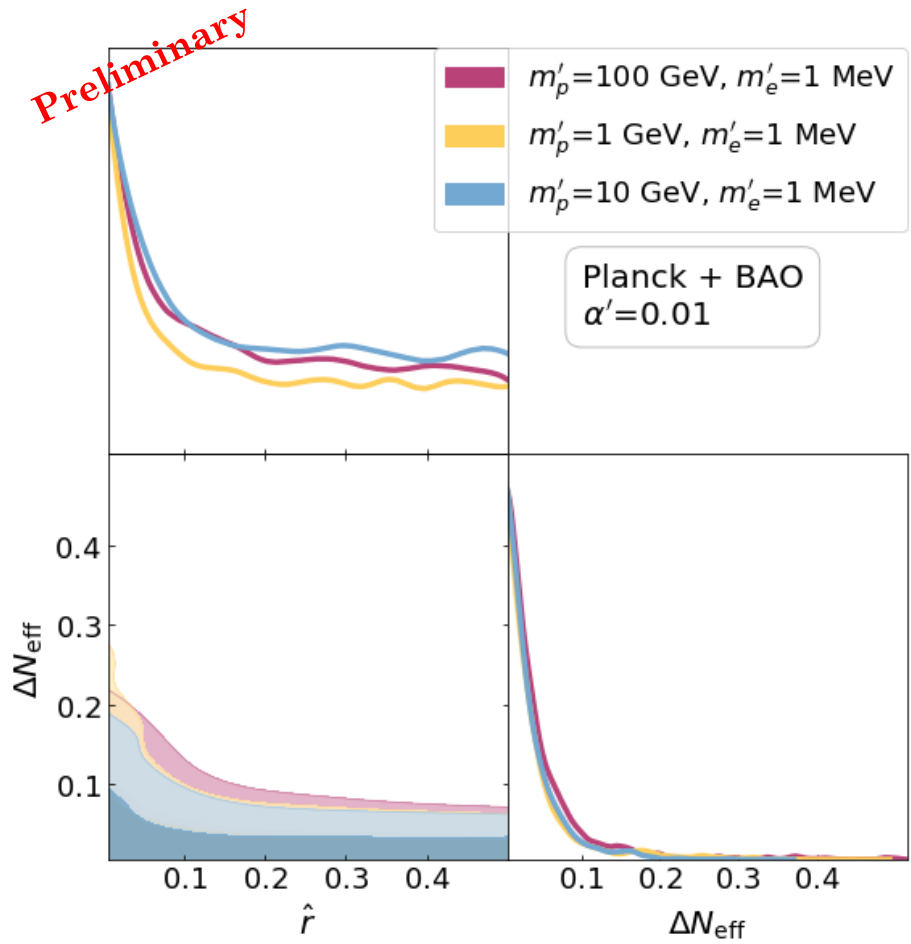




# Constraining the $\Lambda$ CDM parameter space

- Using Planck 2018 high- and low- $l$ , lensing, and BOSS DR12 BAO datasets.
  - Plan to add: KV450 measurements of LSS, SH0ES measurement of  $H_0$ .
- Scan parameter space with Markov Chain Monte Carlo (MCMC). (MontePython)
- Compute 95% C.L. constraints keeping the dark masses and coupling fixed.
- As  $\hat{r} \rightarrow 0$ , recover  $\Lambda$ CDM, so limit on  $\Delta N_{\text{eff}}$  should go to  $\Lambda$ CDM limit.

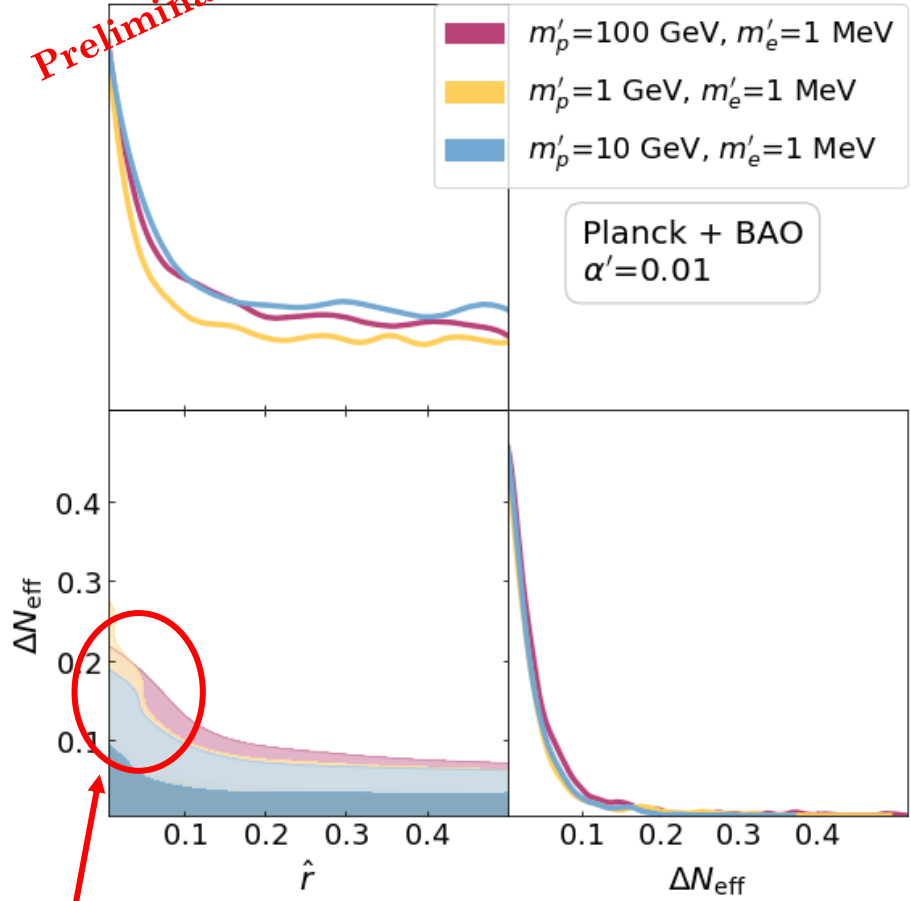
# Constraints on $\Delta N_{\text{eff}}$ and $\hat{r}$



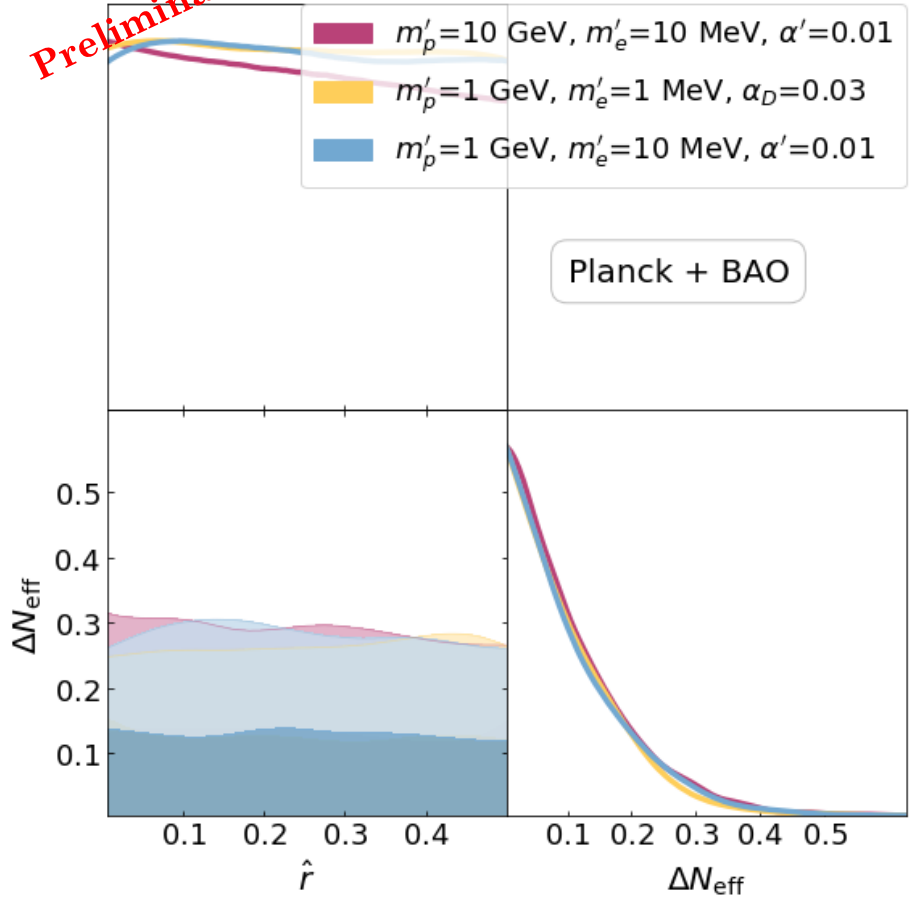
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Preliminary

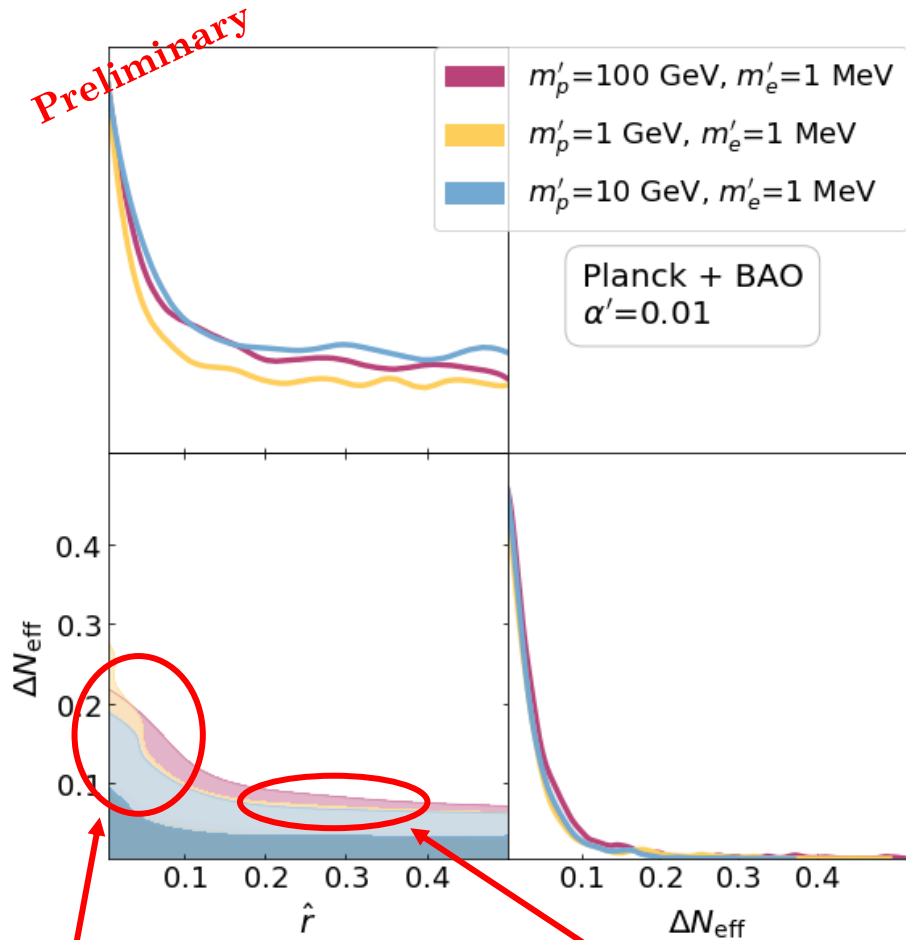


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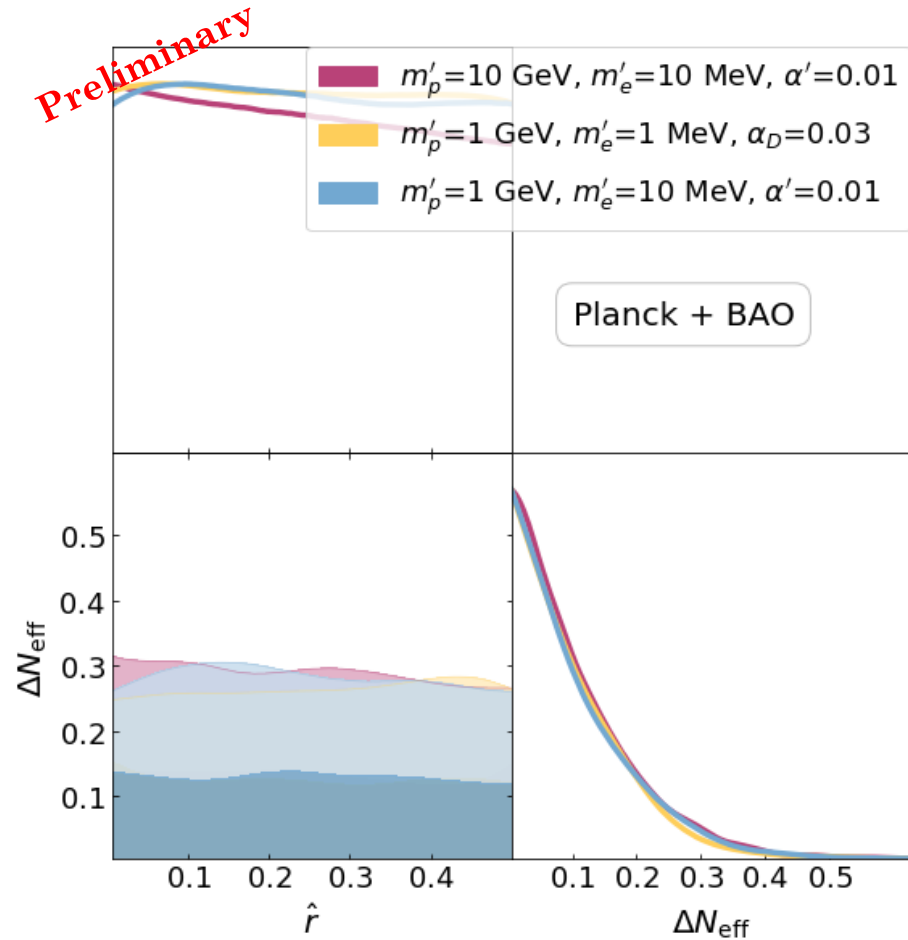
Constraint on  $\Delta N_{\text{eff}}$  weakens for small ADM fraction.

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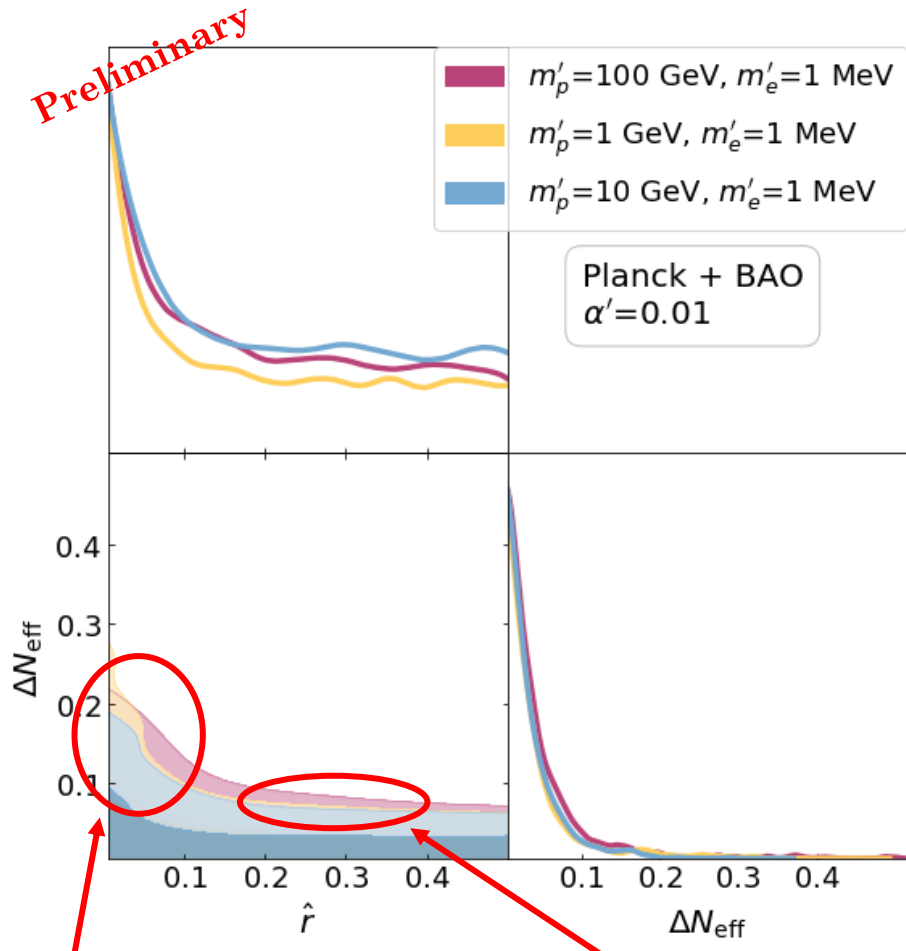


Constraint on  $\Delta N_{\text{eff}}$  weakens for small ADM fraction.

Constraint is mostly insensitive to  $m'_p$ .

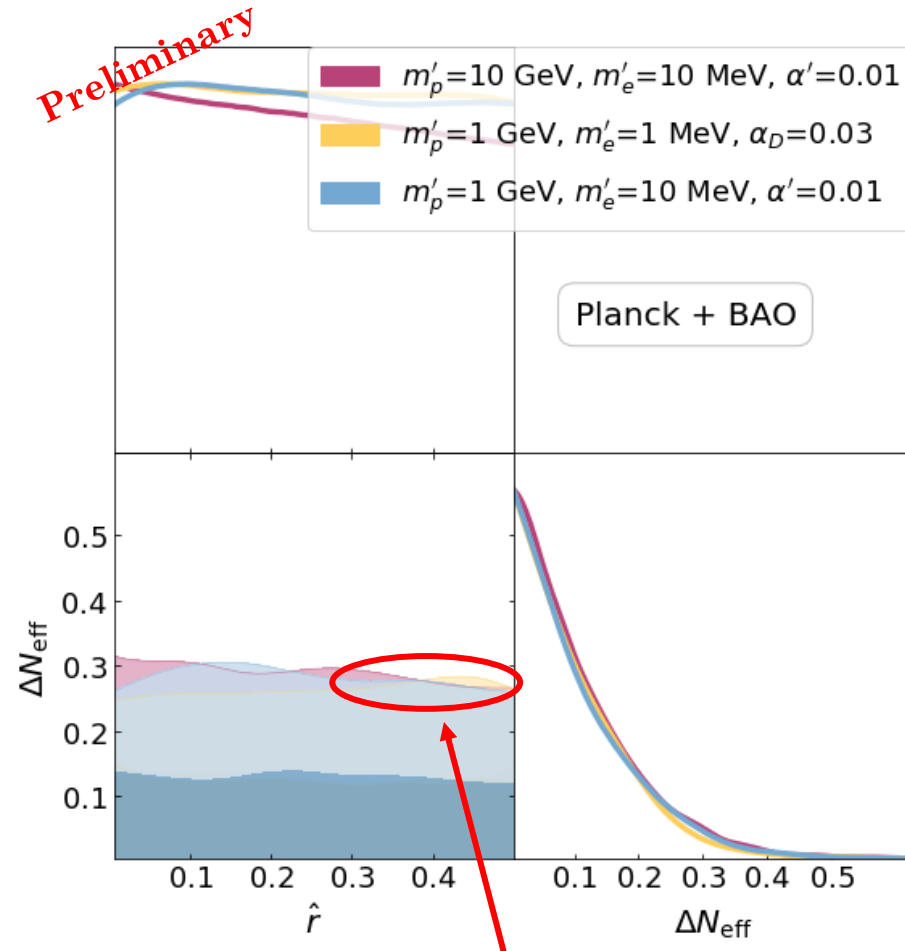


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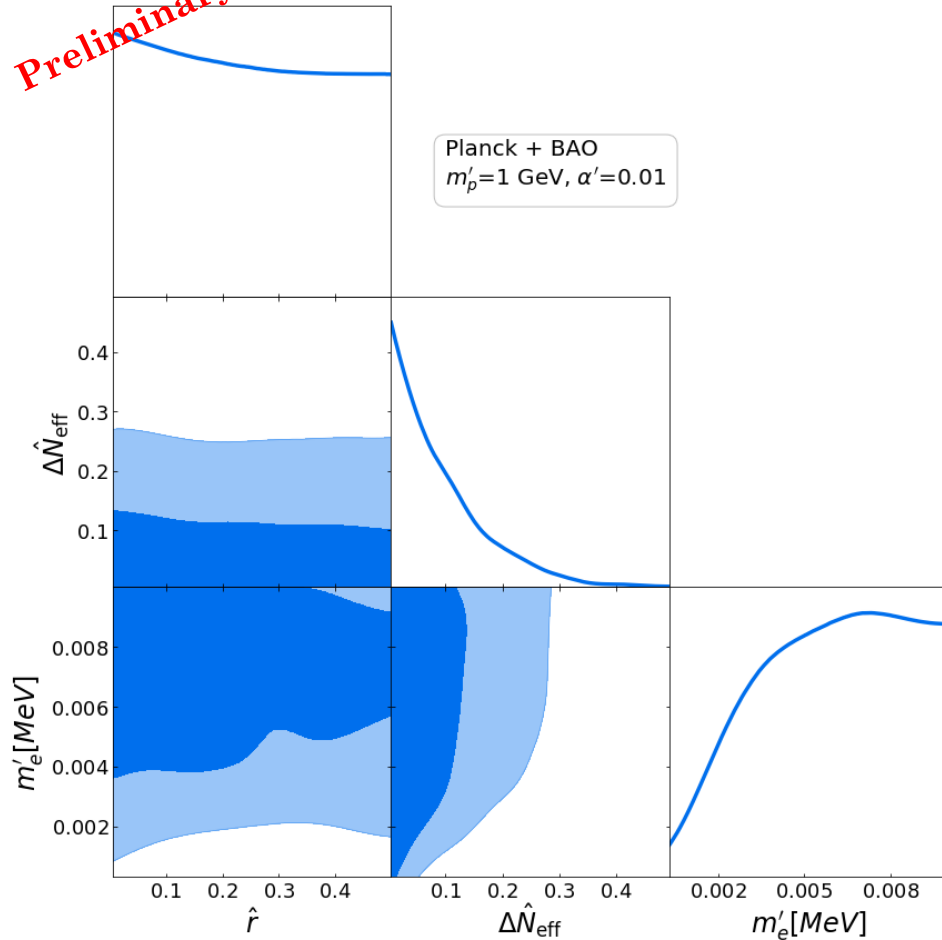


For higher  $m'_e$  or  $\alpha_D$ , large ADM fraction is allowed as long as  $\Delta N_{\text{eff}}$  does not exceed  $\approx 0.3$ .

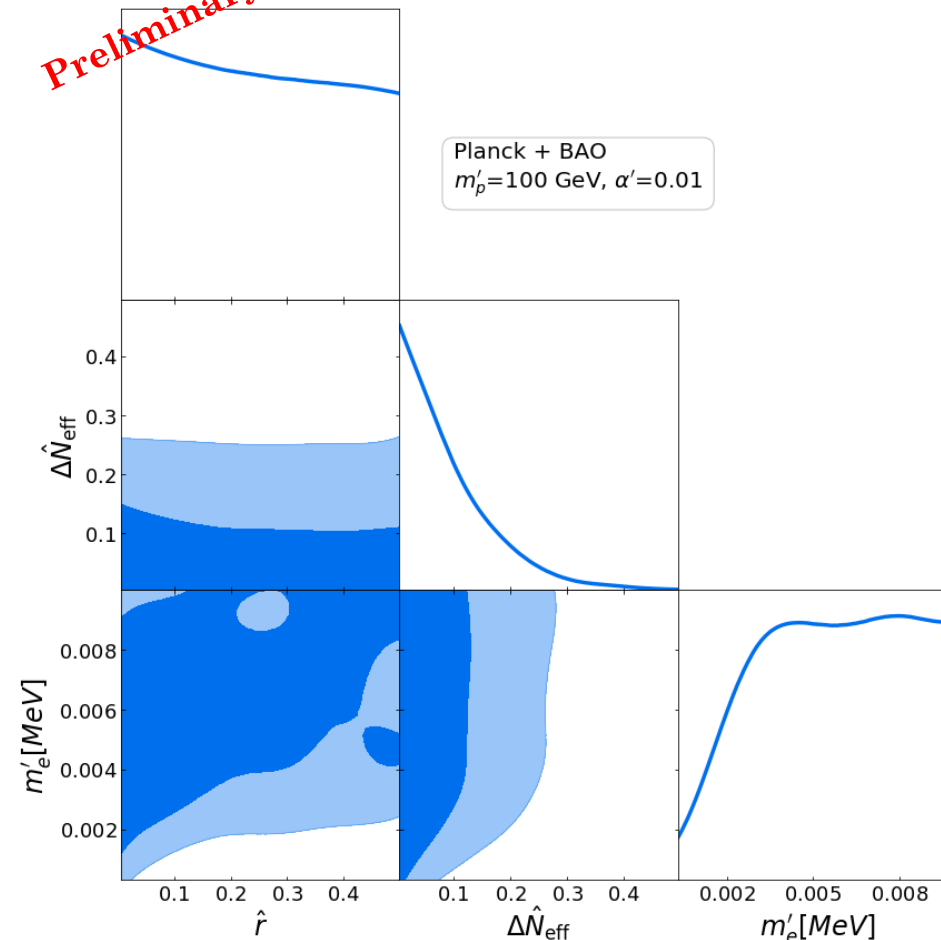


# Constraints on $\Delta N_{\text{eff}}$ , $\hat{r}$ , and $m'_e$

Preliminary



Preliminary



- $\hat{r} > 10\%$  is ruled out at 95% confidence level for  $m'_e < 2 \text{ MeV}$ , with  $m'_p$  from 1 – 100 GeV and  $\alpha' = 0.01$ .
- For higher  $m'_e$ , a large fraction of the DM could be atomic, with  $\Delta N_{\text{eff}}$  up to 0.3.



# Conclusion

- Next steps:
  - Broaden range of parameter scans.
  - Include matter power spectrum and local Hubble constant measurements.
  - Expand ADM scenario to include dark helium.
- Atomic dark matter is a well-motivated, plausible interacting dark sector scenario.
- It can leave distinctive imprints on cosmological observables.
- By modifying CLASS to compute ADM cosmological evolution in generality, we can place new constraints on the atomic dark matter parameter space from precision observations.



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Thank you for your attention!





# References

- Bansal, Kim, Kilda, Low, Tsai: <https://arxiv.org/abs/2110.04317>
- Baumann Cosmology lecture notes
- Blas, Lesgourgues, Tram: <https://arxiv.org/abs/1104.2933>
- Chacko, Craig, Fox, Harnik: <https://arxiv.org/abs/1611.07975>
- Cyr-Racine and Sigurdson: <https://arxiv.org/abs/1209.5752>
- Cyr-Racine and Sigurdson: <https://arxiv.org/abs/1310.3278>
- Gurian, Jeong, Ryan, Shandera: <https://arxiv.org/abs/2110.11964>