



Wenzer Qin

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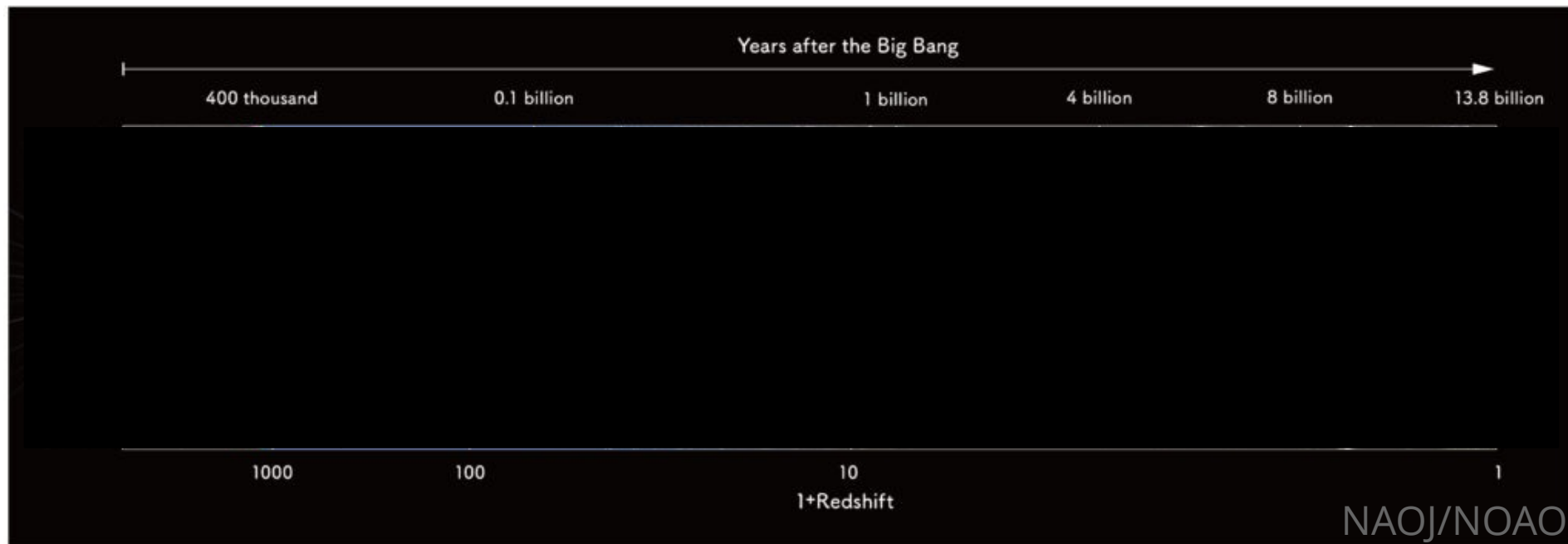
EXTENDING the EFFECTIVE FIELD THEORY of 21CM RADIATION

based on 2205.06270

*TeVPA
AUGUST 8TH , 2022*

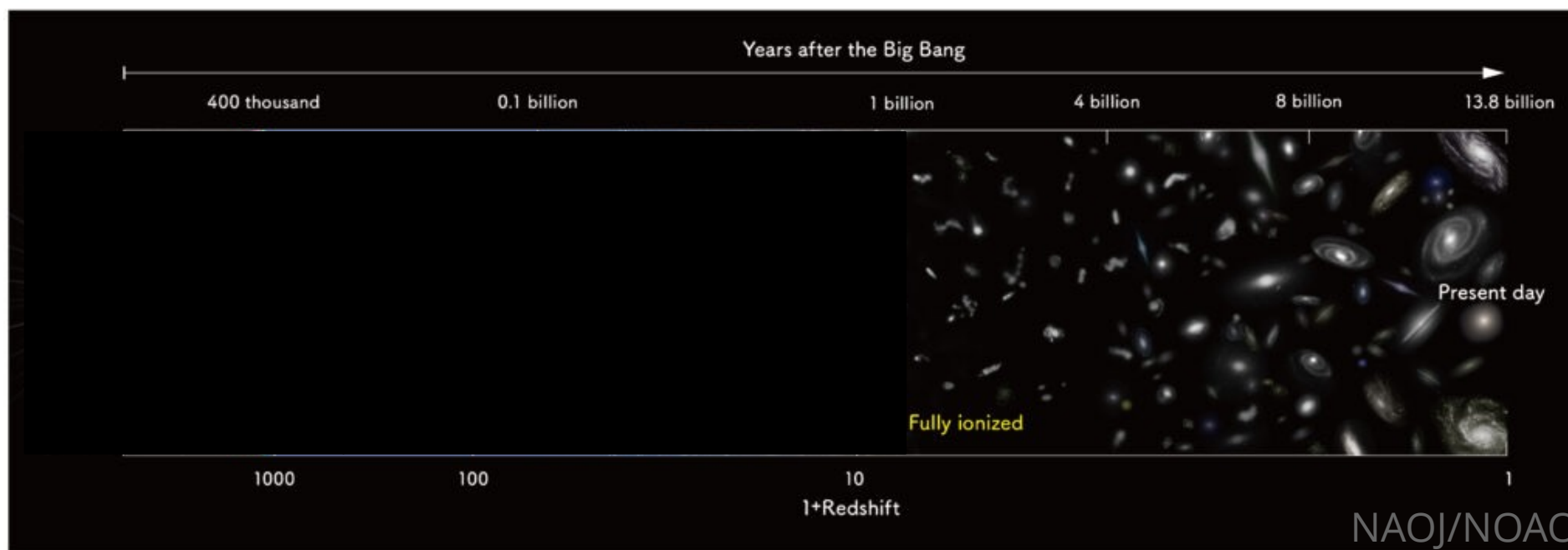
WHAT IS 21CM COSMOLOGY?

- What redshifts have we directly measured?



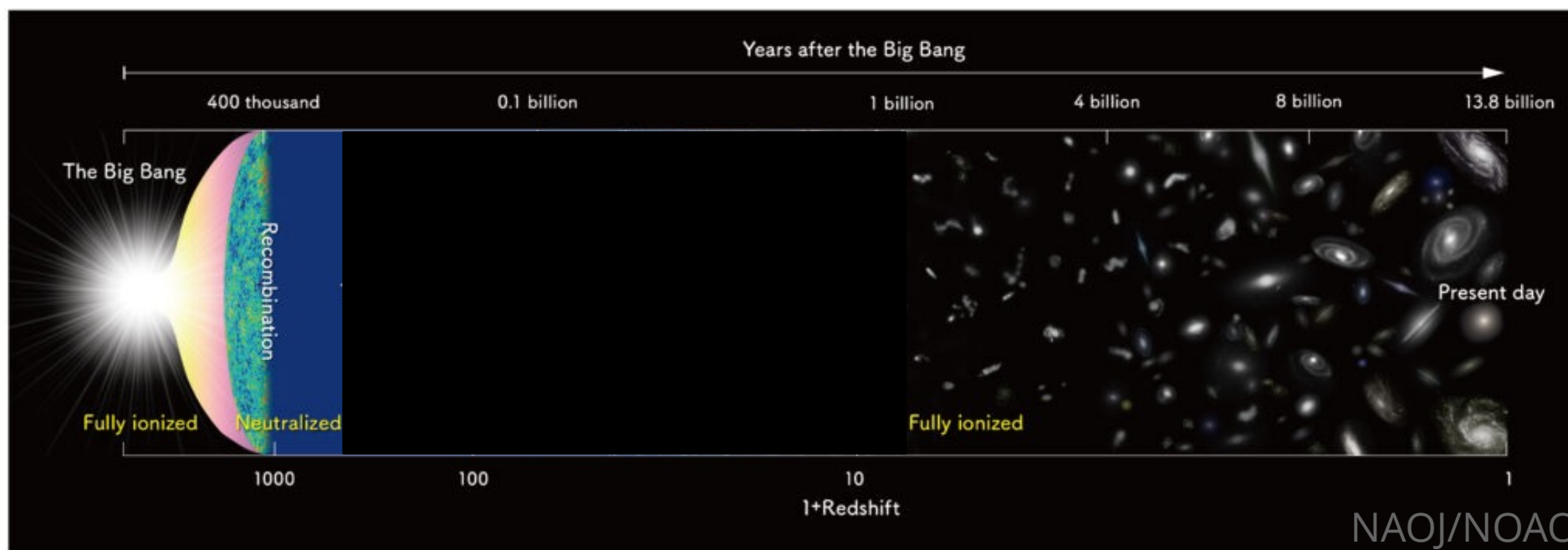
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 - $z \sim \text{few}$: e.g. galaxy surveys



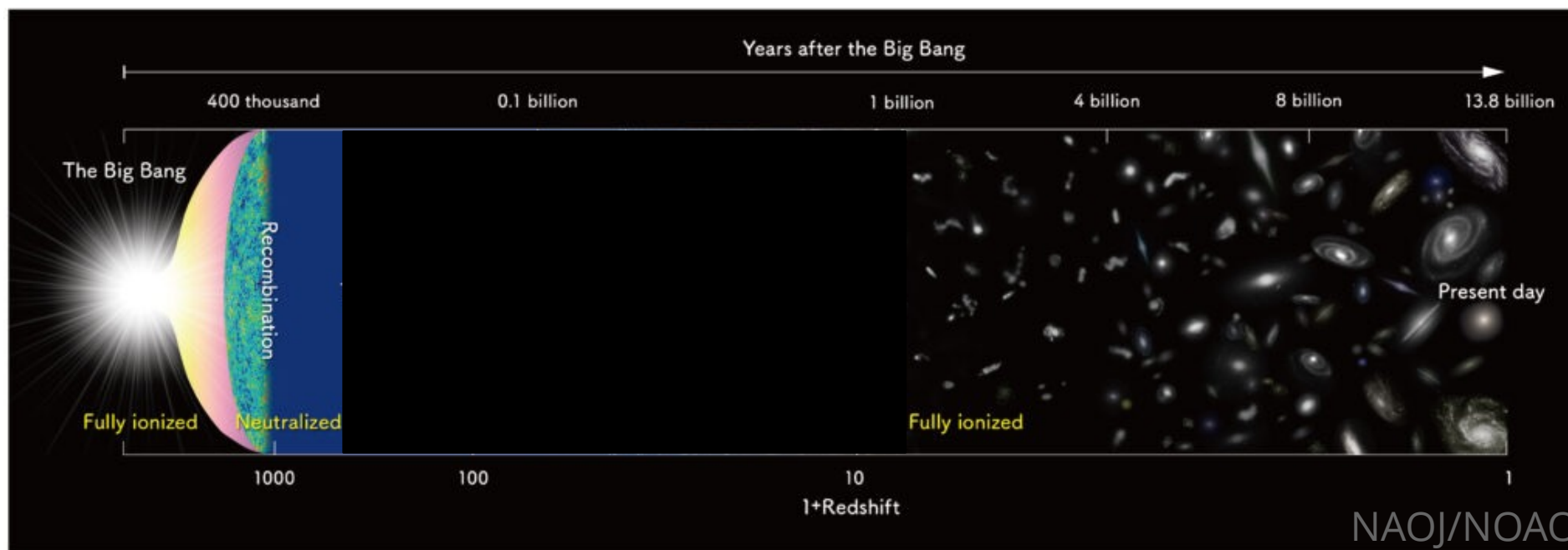
WHAT IS 21CM COSMOLOGY?

- What redshifts have we directly measured?
 - $z \sim \text{few}$: e.g. galaxy surveys
 - $z = 1100$: cosmic microwave background



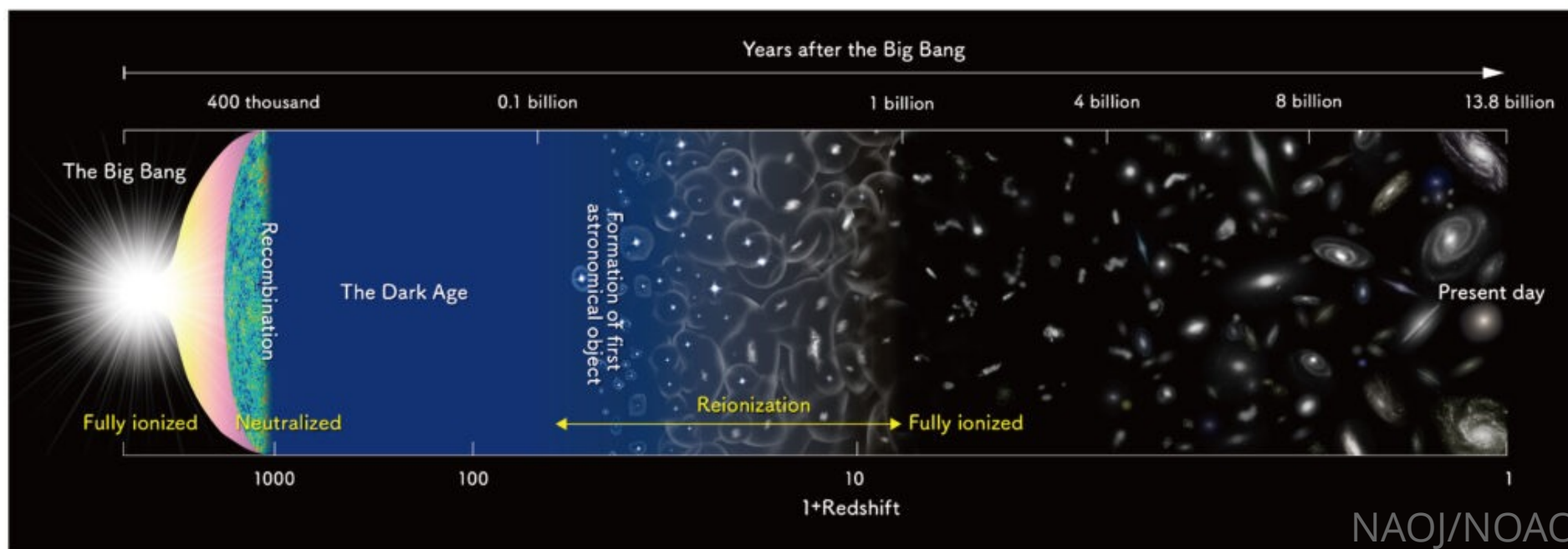
WHAT IS 21CM COSMOLOGY?

- In between, there are few stars/galaxies, only diffuse hydrogen



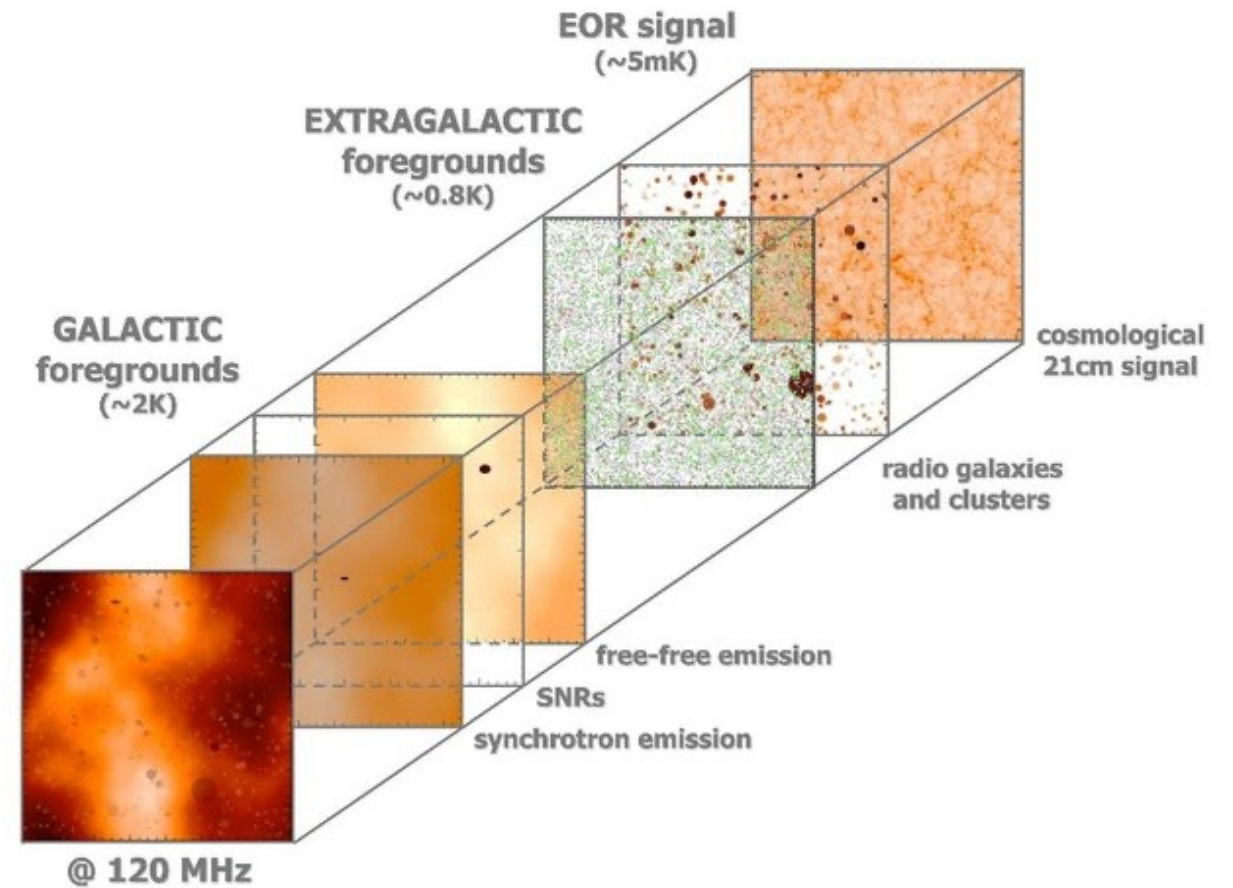
WHAT IS 21CM COSMOLOGY?

- In between, there are few stars/galaxies, only diffuse hydrogen
 - Search for the hyperfine transition of neutral hydrogen → 21cm



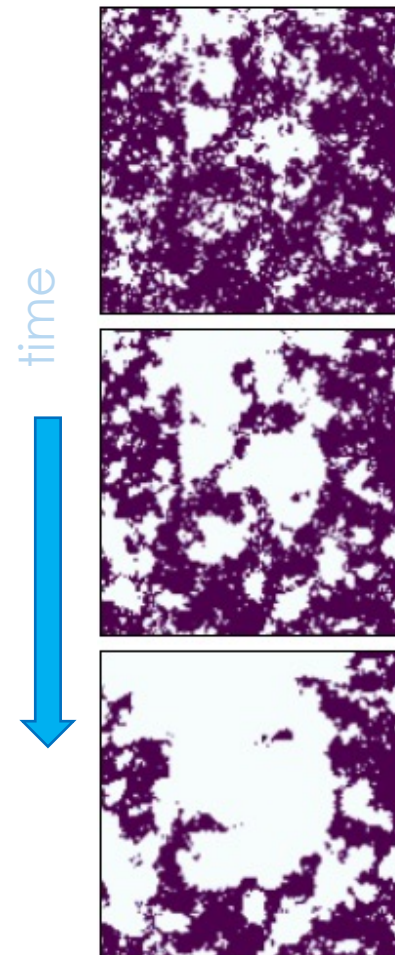
WHAT'S THE CATCH?

- Experimentally: huge foregrounds, e.g. synchrotron radiation

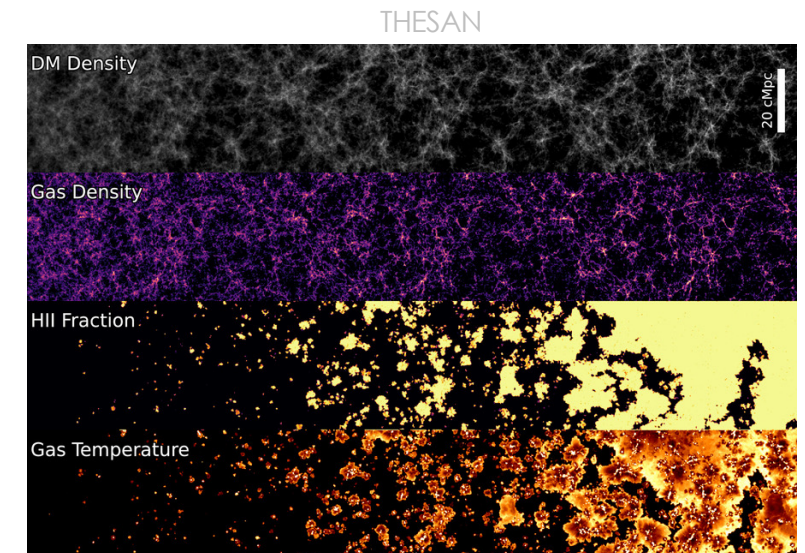
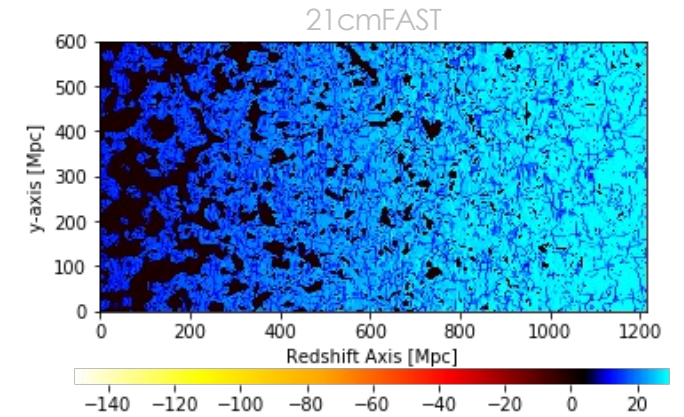


WHAT'S THE CATCH?

- Experimentally: huge foregrounds, e.g. synchrotron radiation
- Theoretically: Prevailing view is that analytic/perturbative methods won't work
 - Reionization is very patchy/nonlinear
 - Instead rely on computationally expensive simulations



McQuinn & D'Aloisio, 2018.



21 CM SIGNAL IS PERTURBATIVE

- McQuinn & D'Aloisio 2018 showed effective field theory (EFT) methods work on observable scales
- Steps in building up our EFT:
 - Relating 21 cm to matter density?
 - Dealing with small nonlinear scales?
 - Redshift space distortions?



FROM DENSITY TO 21CM

- Evolution of matter density/large scale structure is well understood
- To study other quantities, we use ***local bias expansions***
 - E.g. galaxies are biased tracers of matter, form preferentially in overdensities

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- Evolution of matter density/large scale structure is well understood
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 - E.g. galaxies are biased tracers of matter, form preferentially in overdensities
- 21cm brightness temperature is a biased tracer of the matter density field
- Include all operators that respect homogeneity and isotropy

$$\delta_{21} = b_1 \delta - b_{\nabla^2} k^2 \delta + b_2 \delta^2 + \dots$$

SMALL SCALE NONLINEARITIES?

- Small scale modes become nonlinear first → smooth over these

$$\delta_{\text{long}}(\mathbf{x}) = \int d^3x' W_{\Lambda}(\mathbf{x} - \mathbf{x}') \delta(\mathbf{x}')$$

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- Terms with more than one field are still sensitive to small scale modes
 - E.g. using a sharp cutoff...

$$[\delta^2(\mathbf{k})]_{\text{long}} = \int_0^{\Lambda} \frac{dq}{2\pi^2} \delta(\mathbf{q}) \delta(\mathbf{k} - \mathbf{q}) \neq [\delta_{\text{long}}(\mathbf{k})]^2$$

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- Remove the dependence on Λ by *renormalization*



REDSHIFT SPACE DISTORTIONS

- 21 cm signal is radiation
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REDSHIFT SPACE DISTORTIONS

- 21cm signal is radiation
- Frequency is redshifted because sources have peculiar velocities
- Occurs along the line of sight, where interferometers are most sensitive
- Need to include redshift space distortions in perturbative treatments
 - Expand in small k_{\parallel}/H

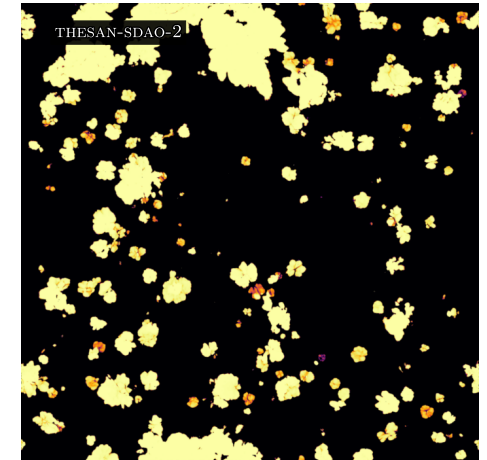
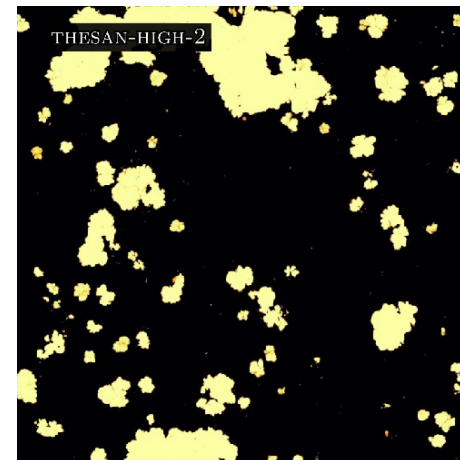
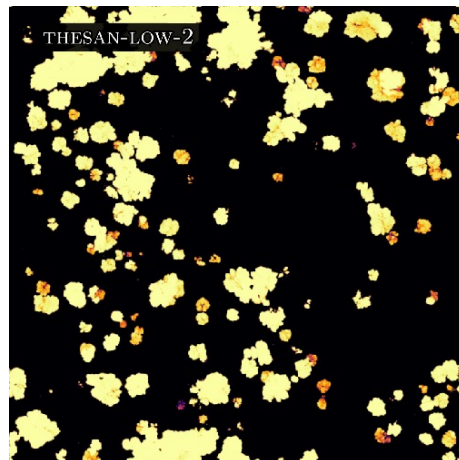
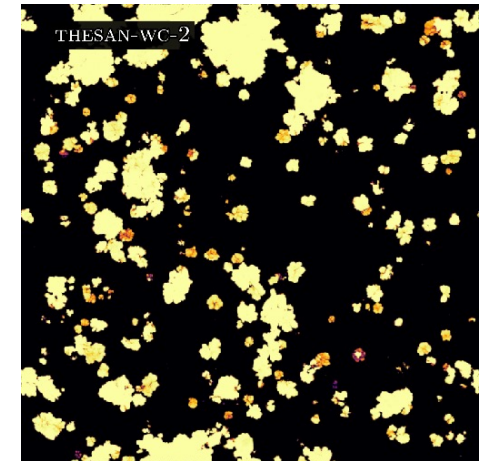
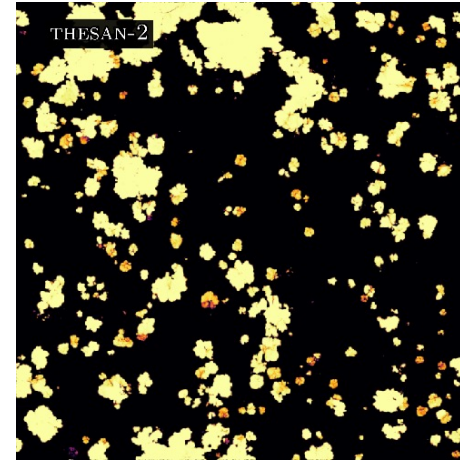
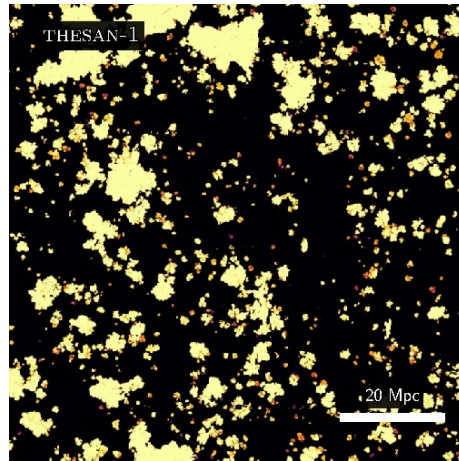
$$\delta_r = \delta(\vec{k}) - i \frac{k_{\parallel}}{\mathcal{H}} v_{\parallel}(\vec{k}) - i \frac{k_{\parallel}}{\mathcal{H}} \delta v_{\parallel}(\vec{k}) - \frac{1}{2} \left(\frac{k_{\parallel}}{\mathcal{H}} \right)^2 v_{\parallel}^2(\vec{k}) - \frac{1}{2} \left(\frac{k_{\parallel}}{\mathcal{H}} \right)^2 \delta v_{\parallel}^2(\vec{k}) + \frac{i}{6} \left(\frac{k_{\parallel}}{\mathcal{H}} \right)^3 v_{\parallel}^3(\vec{k}) + \dots$$



FITTING TO THE
THESAN SIMULATIONS

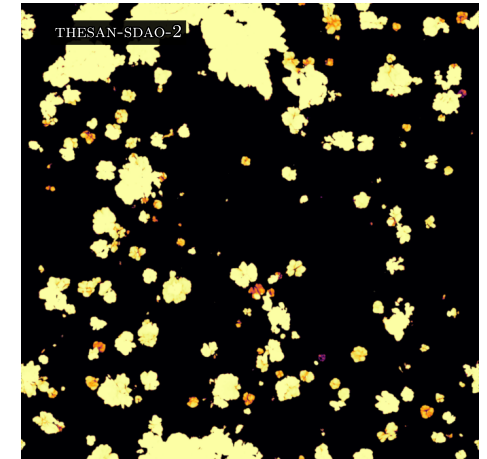
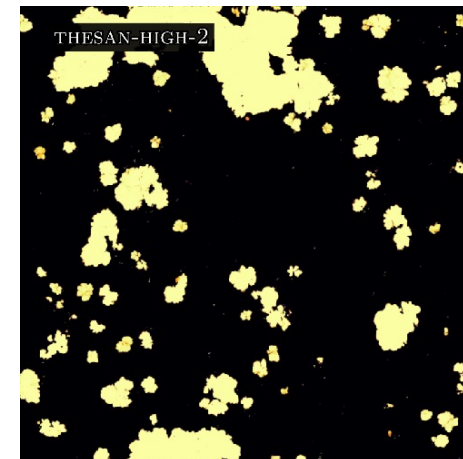
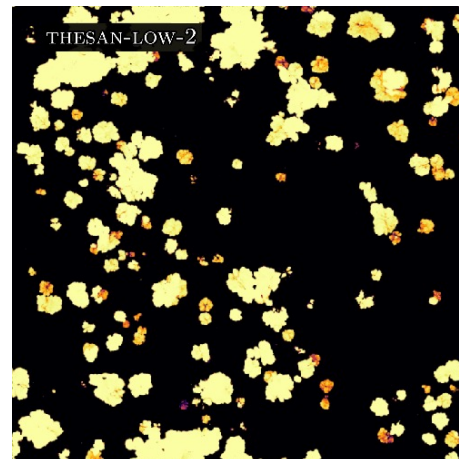
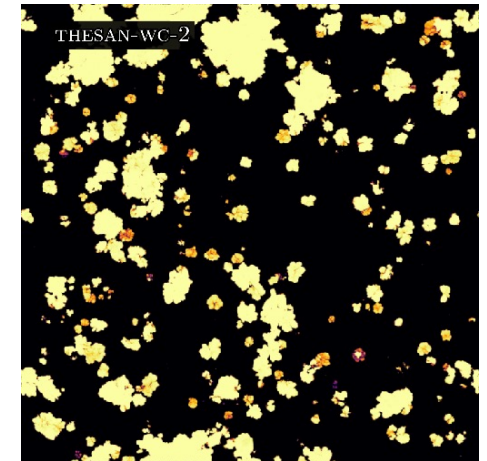
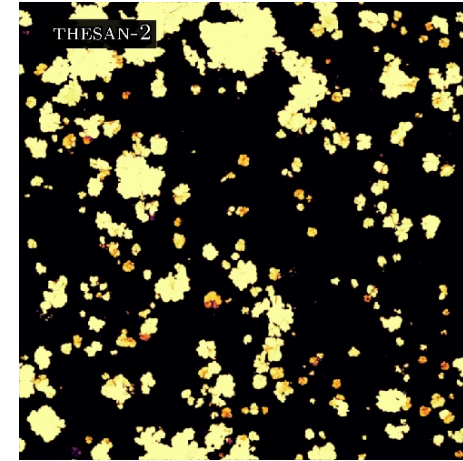
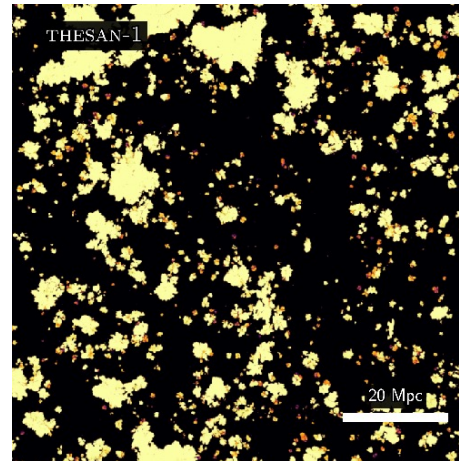
THESAN

- Neutral fraction is ~ 0.7 for each simulation shown



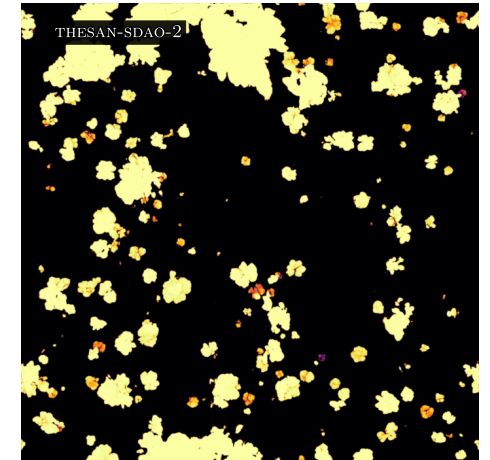
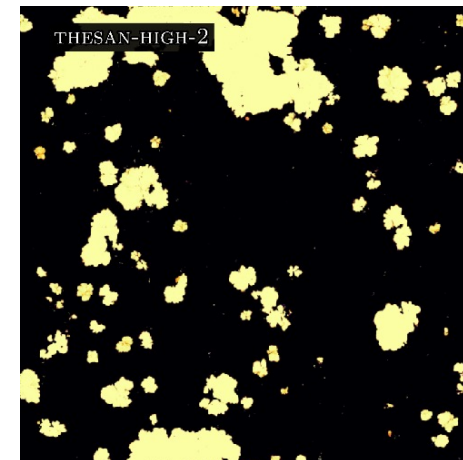
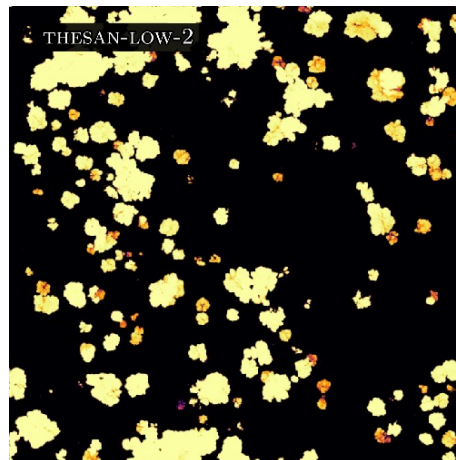
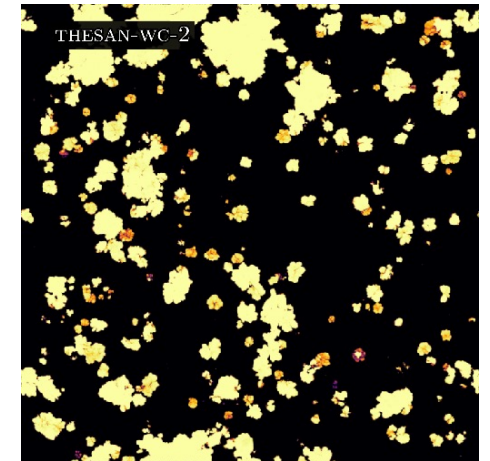
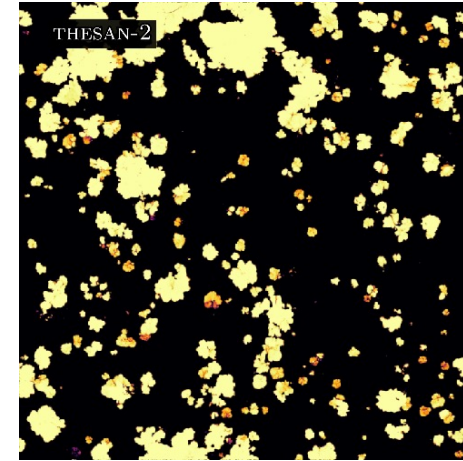
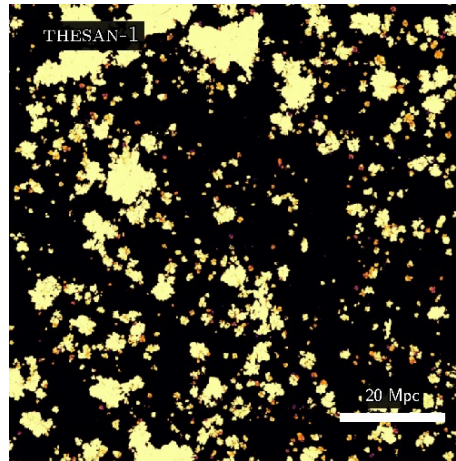
THESAN

- Thesan-1: High resolution
- Thesan-2: Medium resolution
- Thesan-WC-2: Compensates for lower star formation due to less resolution

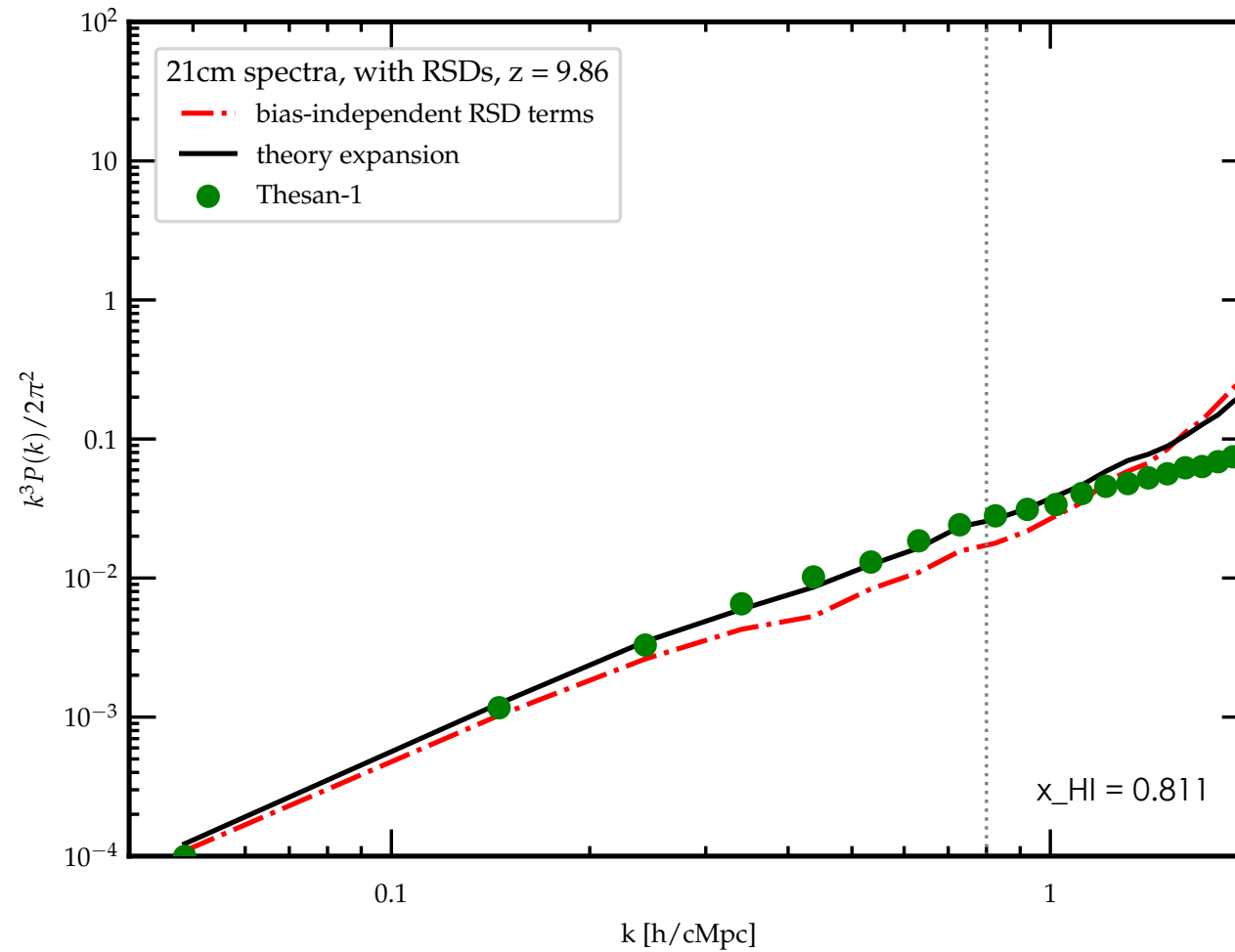


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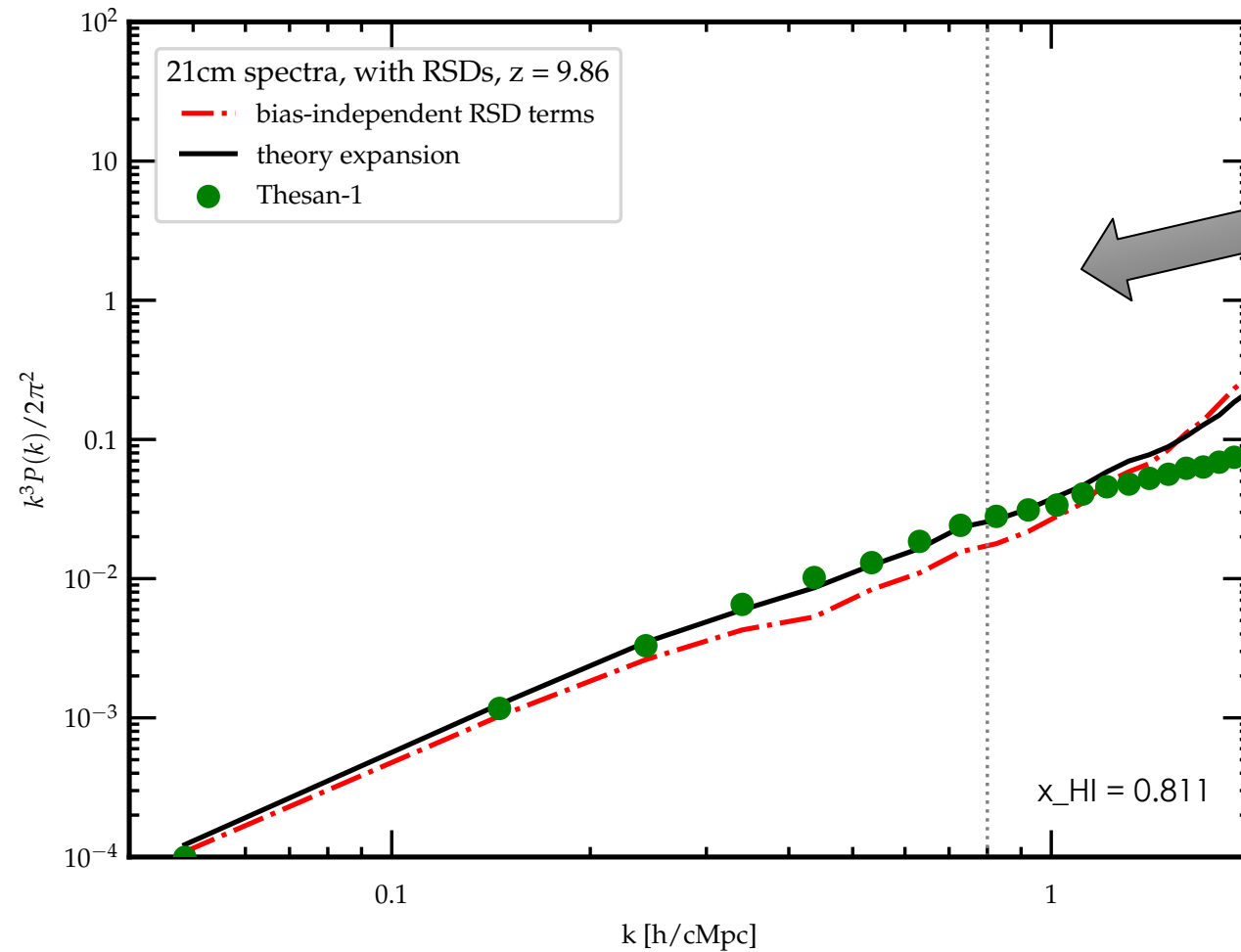
- Thesan-1: High resolution
- Thesan-2: Medium resolution
- Thesan-WC-2: Compensates for lower star formation due to less resolution
- Thesan-Low-2: Small haloes contribute to reionization
- Thesan-High-2: Large haloes contribute to reionization
- Thesan-SDAO-2: Non-standard dark matter model



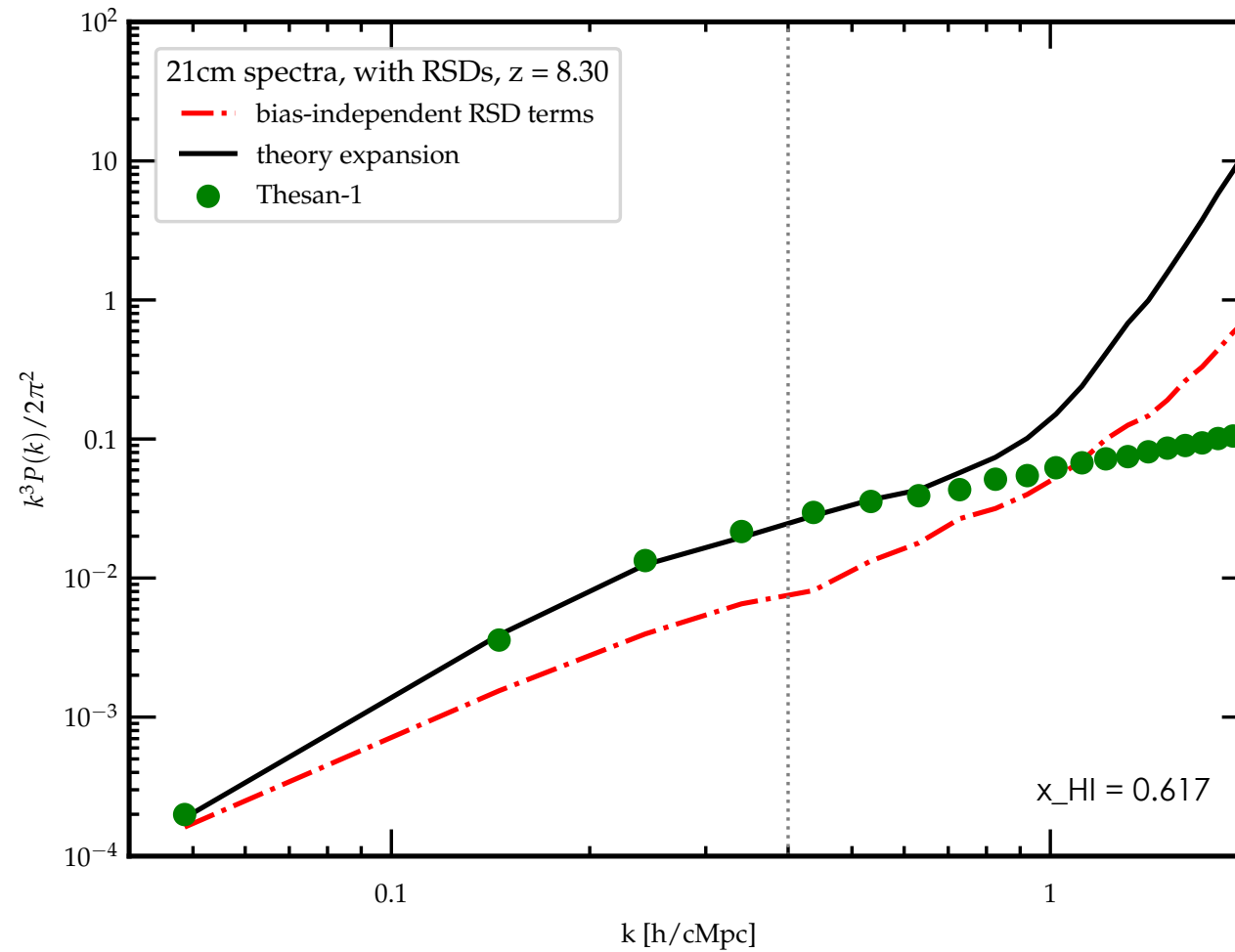
POWER SPECTRUM



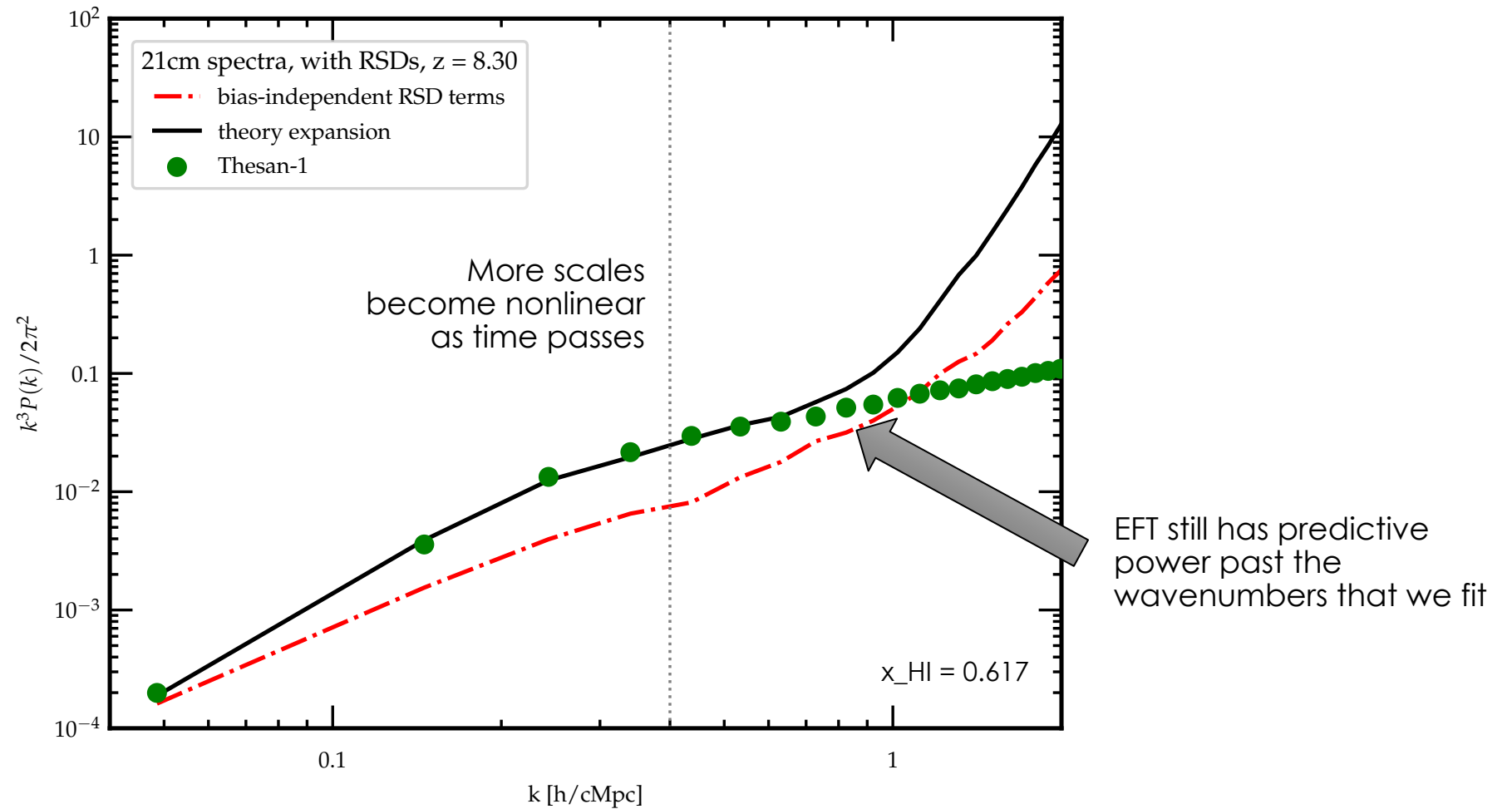
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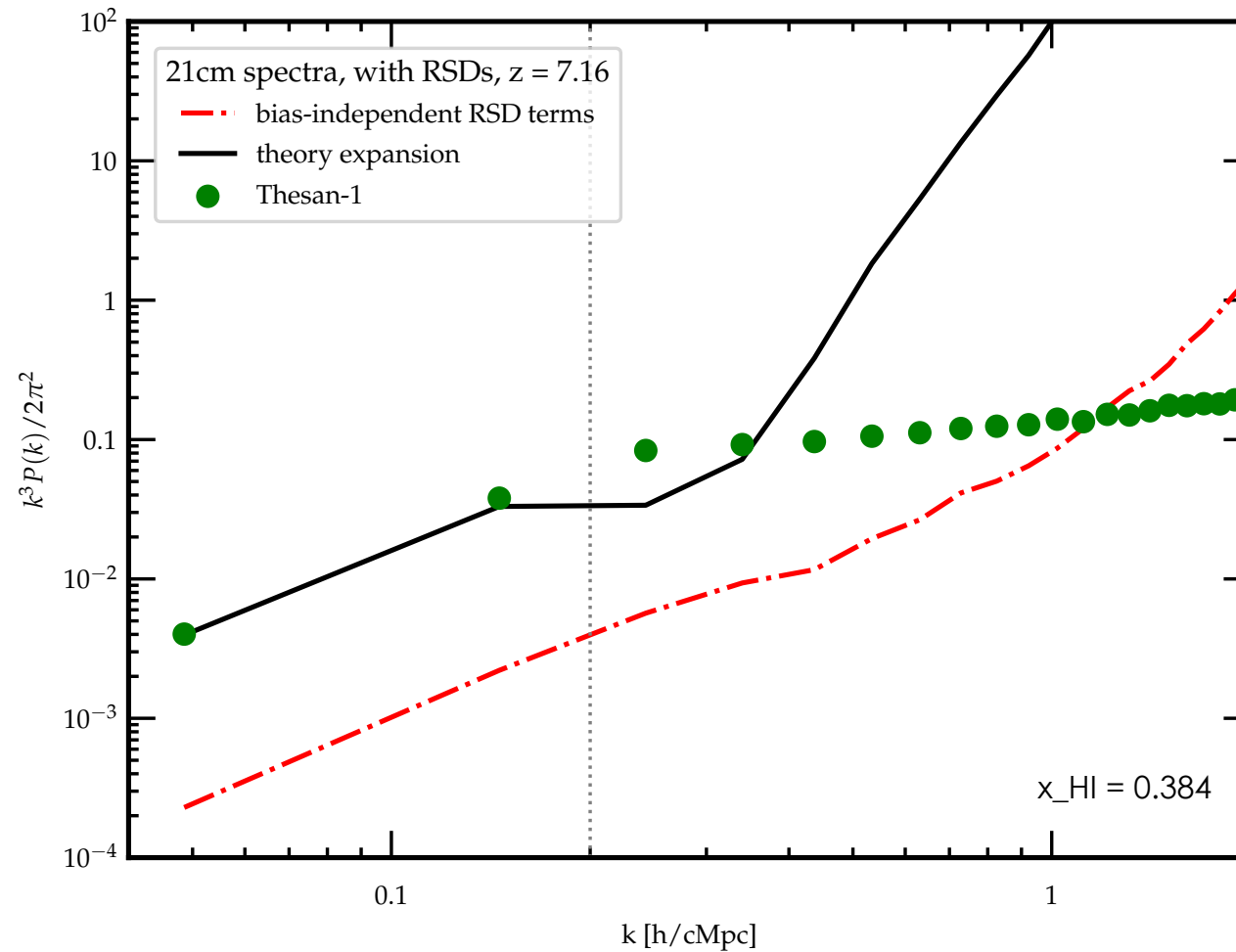
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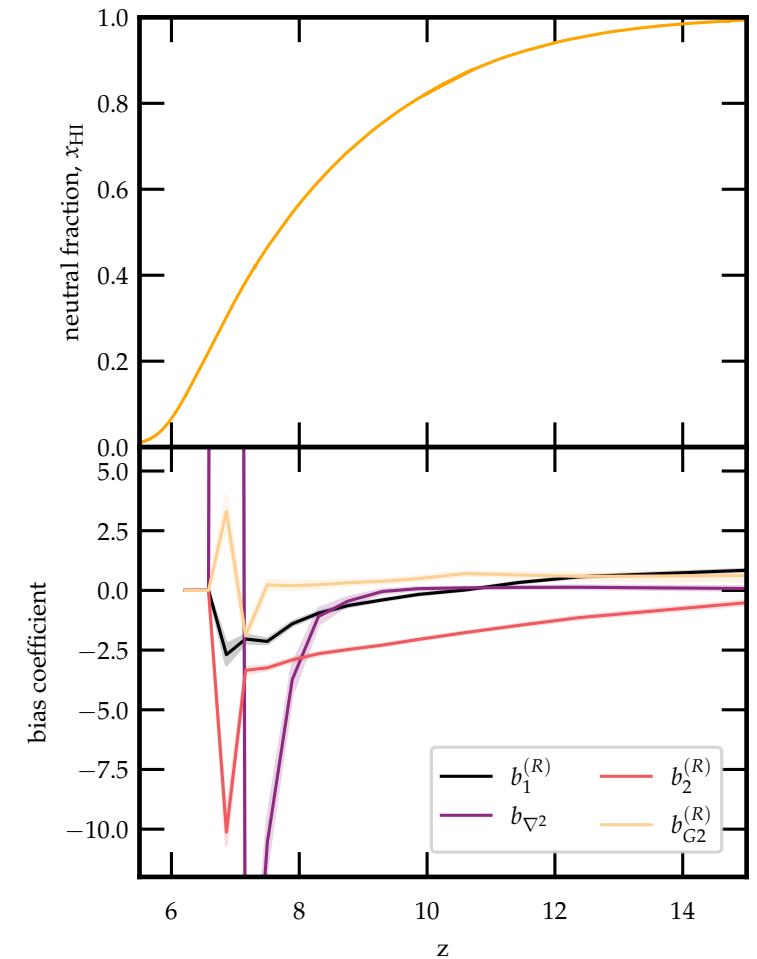
POWER SPECTRUM



At this level of ionization, perturbative theory breaks down

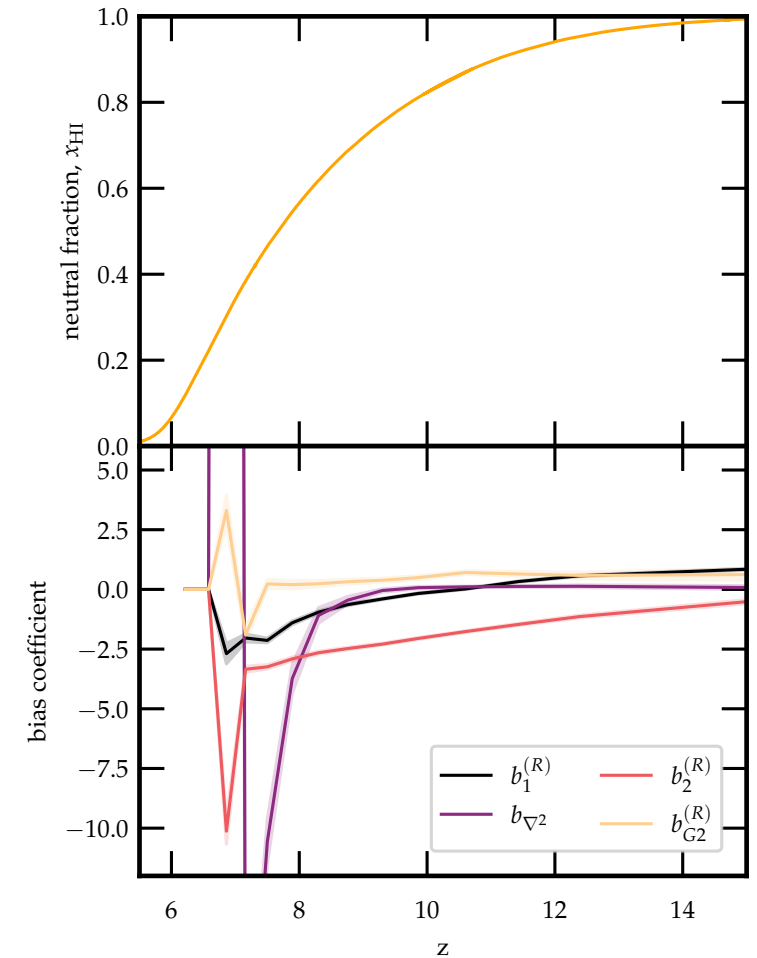
EVOLUTION OF COEFFICIENTS

- How do these coefficients evolve with time?
- Evolution becomes rapid/jagged after a certain time --- theory is breaking down
- At the end of reionization, 21 cm signal vanishes

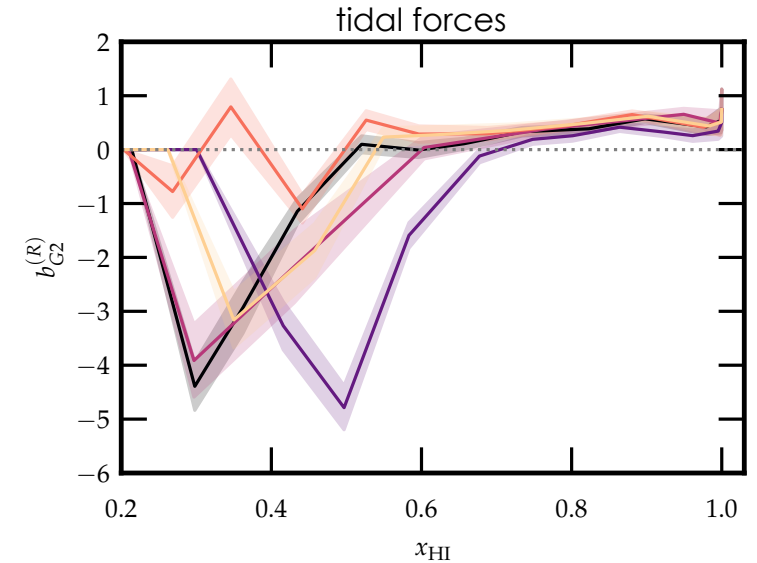
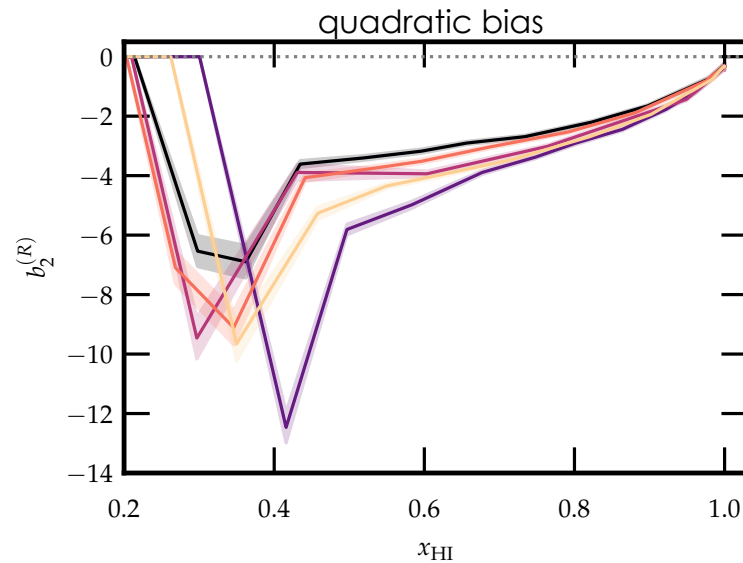
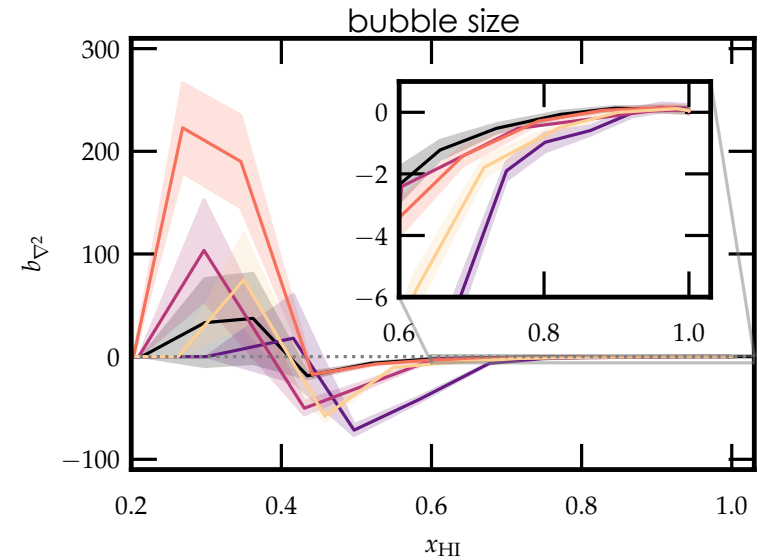
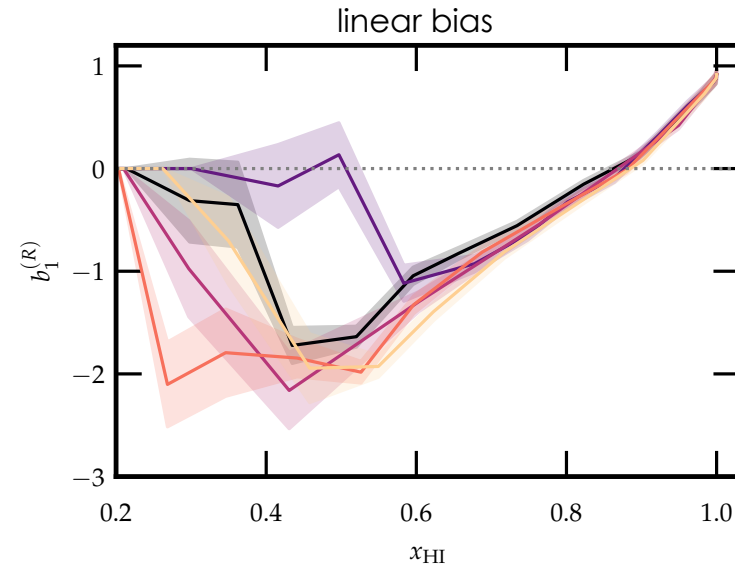
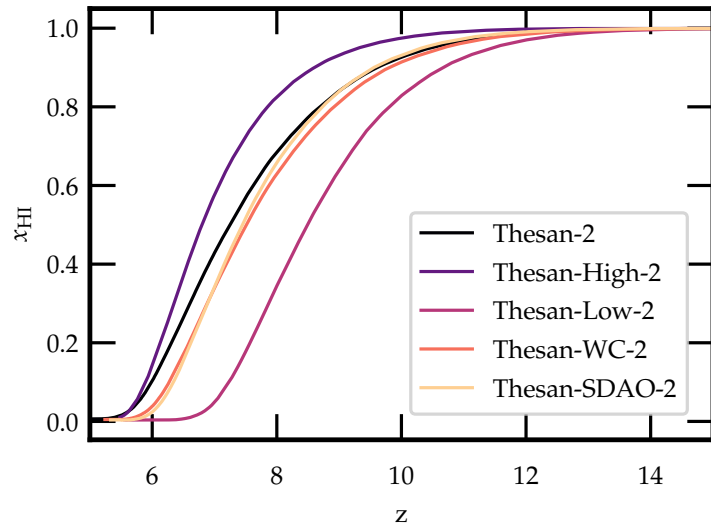


EVOLUTION OF COEFFICIENTS

- Physical interpretations:
 - $b_1^{(R)}$ is linear bias
 - $b_2^{(R)}$ is quadratic bias
 - b_{∇^2} is related to bubble size
 - $b_{G2}^{(R)}$ quantifies anisotropic stresses



SIMULATION DIFFERENCES



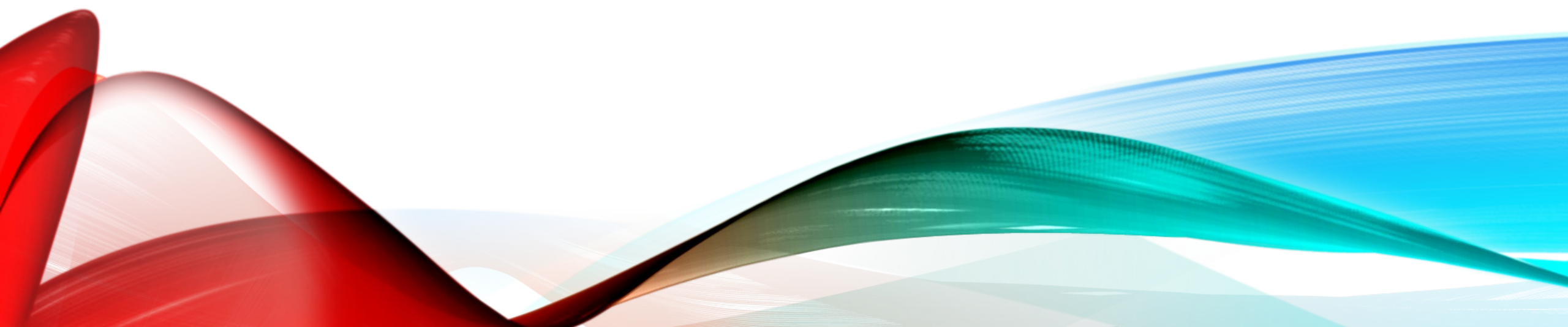


SUMMARY

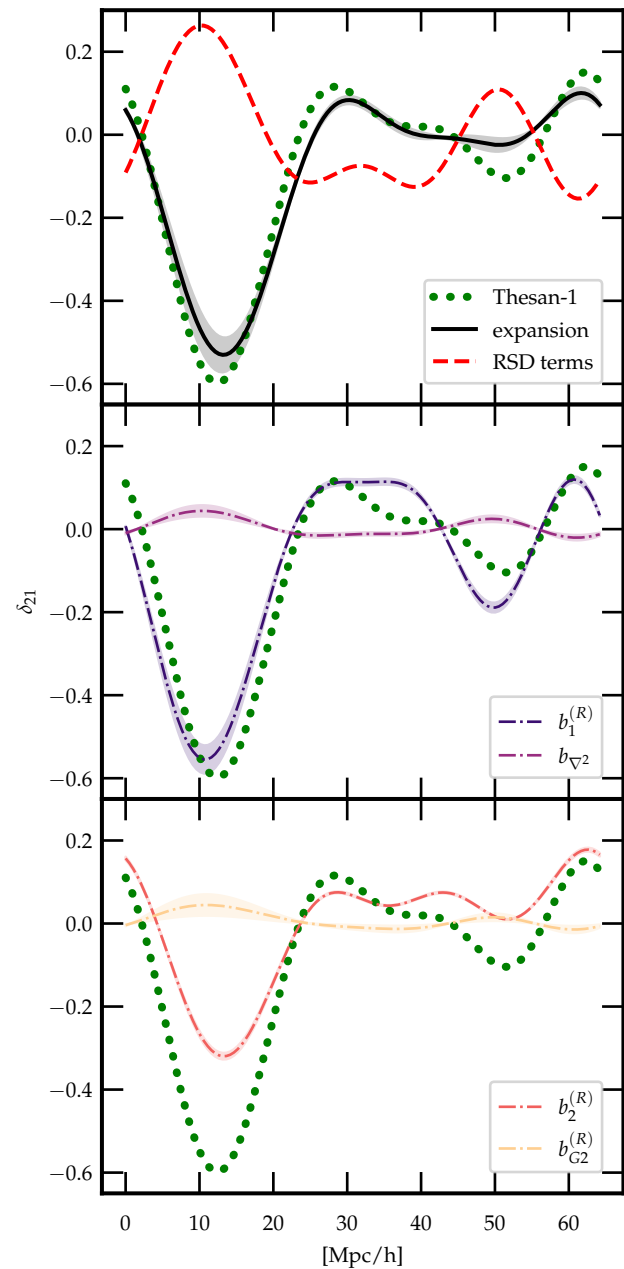
- On observable scales, we can use perturbative methods
- We've extended these EFT methods, e.g. including RSDs
- Theory expansion is a good fit to simulations, at early enough redshifts and large length scales
- Evolution of coefficients reflects different physics

- Future steps:
 - Spin temperature fluctuations
 - Reconstructing modes in the foreground wedge

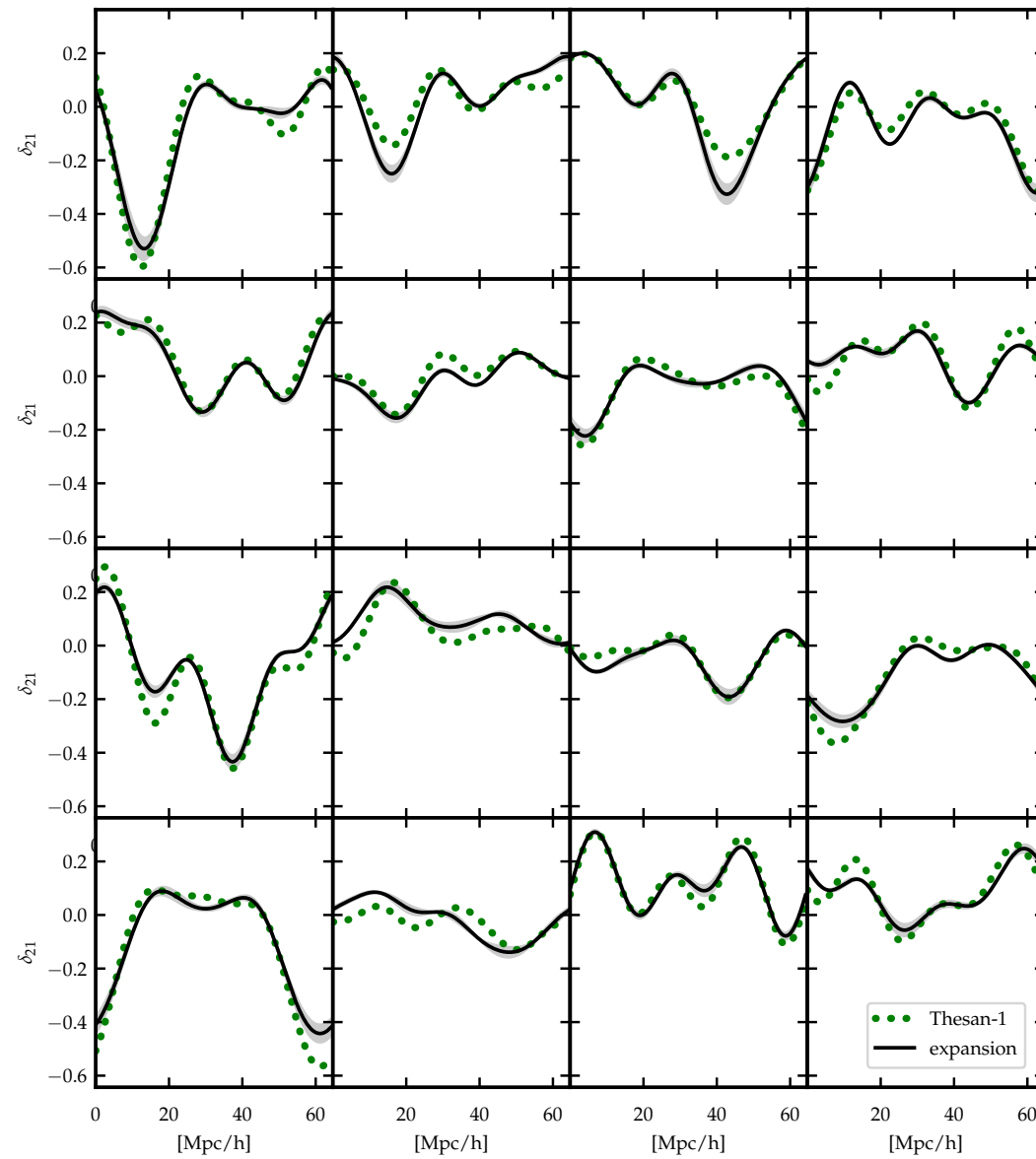
BACKUP



COMPARISON IN CONFIGURATION SPACE



- 1D slices of 21 cm brightness temperature at $z = 8.30$, $x_{\text{HI}} = 0.617$
- Smoothed over $k = 0.4 \text{ h}/\text{cMpc}$

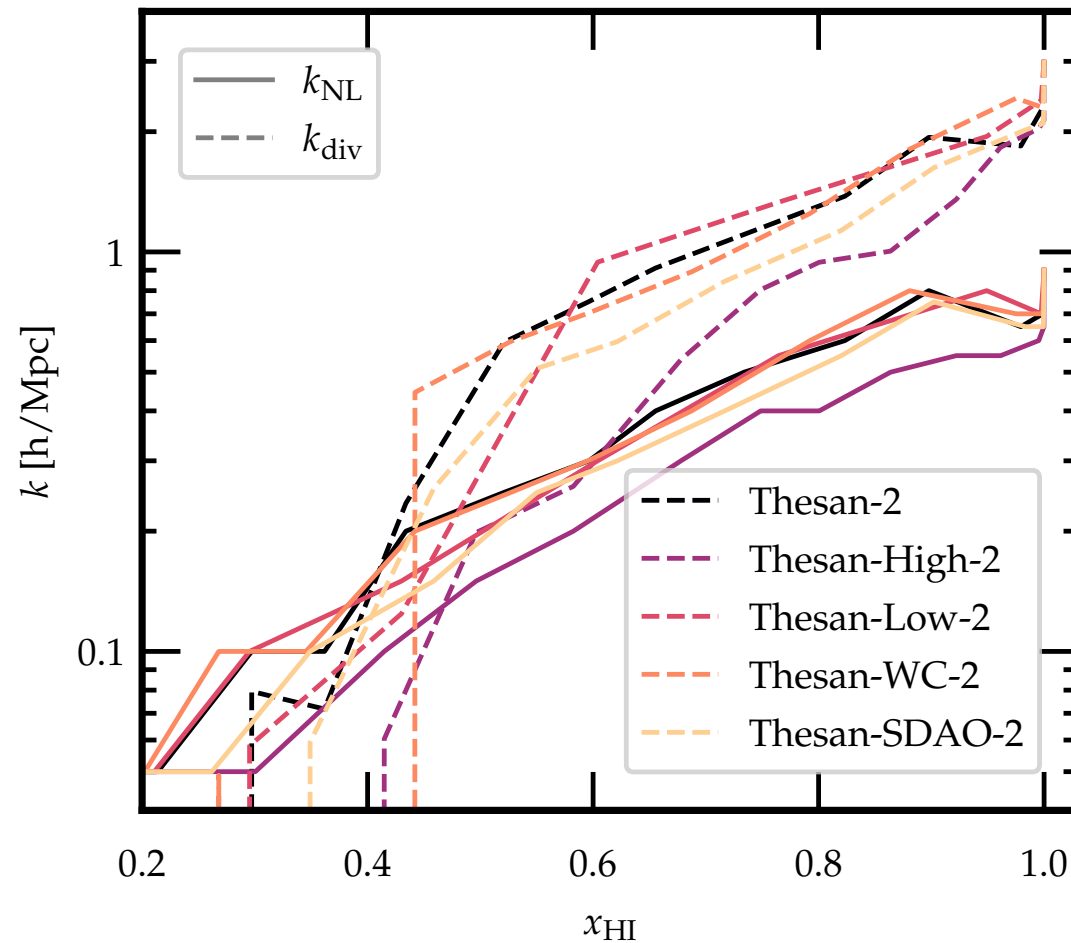


SMALL SCALE NONLINEARITIES?

- Use a systematic method to remove dependence of observables on small scales/cutoff \rightarrow *renormalization*

$$[\delta^2] = \delta^2 - \sigma^2(\Lambda) \left(1 + \frac{68}{21}\delta + \frac{8126}{2205}\delta^2 + \frac{254}{2205}\mathcal{G}_2 \right) + \dots$$

WHEN PERTURBATIVITY BREAKS DOWN



EFT SHAPES AS FUNCTION OF ANGLE

