

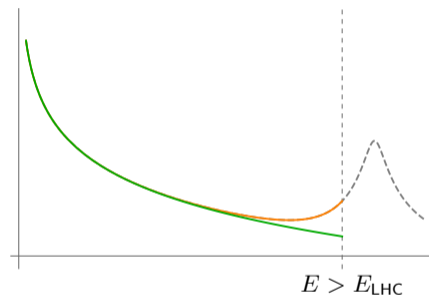
# Search for new physics in top quark production with additional leptons using effective field theory

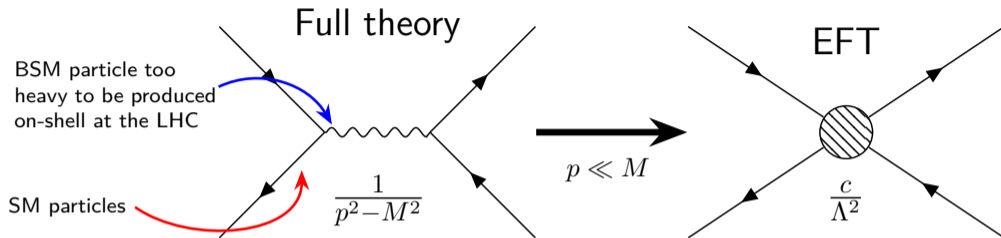
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On behalf of the CMS collaboration

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- Search for new fundamental particles is motivated by the strong evidence for phenomena not described by the SM
- New particles may not be light enough to be produced on-shell at the LHC
- Indirect searches are needed if we want to probe these regimes (or wait until someone builds a larger collider)
- Effective field theory provides a framework for probing these higher energy scales





Since we can't produce heavy particle on-shell at the LHC, it would be hard to find it via a direct search, but EFT can provide discovery potential

The interaction can be described by an EFT operator, with the strength of the interaction determined by a WC  $c$

Important notes:

- EFT validity is only good for energy scales below the cutoff scale  $\Lambda$
- Underlying low energy theory needs to be complete
- Provides systematic and mostly model independent search for BSM effects

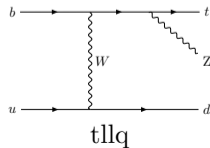
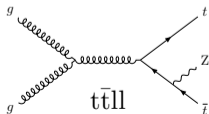
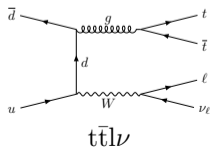
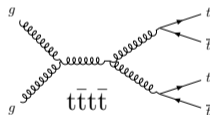
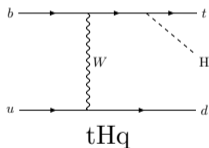
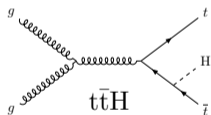
- SM is treated as the lowest order term in an expansion of higher dimensional ( $d > 4$ ) operators, that describe physics above some cutoff scale  $\Lambda$

$$\mathcal{L}_{\text{EFT}} = \mathcal{L}_{\text{SM}} + \sum_{d,i} \frac{c_i^{(d)}}{\Lambda^{d-4}} \mathcal{O}_i^{(d)}$$

- The strength of the interactions introduced by the EFT operators is controlled by a dimensionless parameter called a Wilson coefficient (WC)
- If all WCs are 0, the EFT Lagrangian reduces naturally to the SM
- Only one dimension 5 operator  $\rightarrow$  Violates lepton number and is excluded
- Focus only on dimension 6 operators as the lowest order contributions

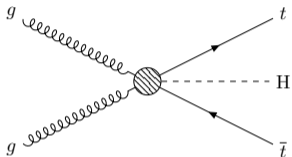
# Analysis Overview

- Aim is to probe the effects of as many EFT operators across all relevant processes simultaneously
- Analysis focuses on operators that couple the top quark to leptons, bosons, and other heavy quarks
- Focus on associated top processes and model how EFT operators affect expected yields
  - Signal processes:  $t\bar{t}H$ ,  $tHq$ ,  $t\bar{t}t\bar{t}$ ,  $t\bar{t}l\nu$ ,  $t\bar{t}ll$ ,  $tllq$
  - Processes are relatively rare and provide for a clean well isolated signal region

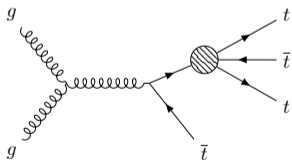


# Analysis Overview

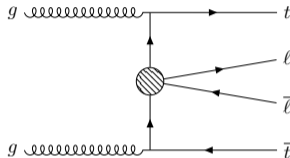
- Focus on 26 operators, which can be grouped together into 4 different categories



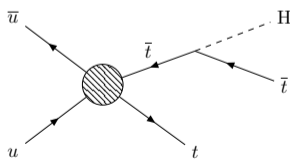
Operators involving **heavy** quarks and **bosons**



Operators involving 4 **heavy** quarks



Operators involving 2 **heavy** quarks and 2 **leptons**



Operators involving 2 **heavy** quarks and 2 **light** quarks

# EFT Parameterization

- Need some way to model EFT contributions
- Matrix element can be written as the sum of SM and new physics components

$$\mathcal{M} = \mathcal{M}_{\text{SM}} + \sum_i \frac{c_i}{\Lambda^2} \mathcal{M}_i \longrightarrow c_i \text{ are the Wilson coefficients}$$

- Since  $\sigma \propto \mathcal{M}^2 \rightarrow$  the cross section will have a **quadratic** dependence on the WCs

$$d\sigma(\vec{c}) \propto \left| \mathcal{M}_{\text{SM}} + \sum_i \frac{c_i}{\Lambda^2} \mathcal{M}_i \right|^2 \propto s_0 + \sum_j s_j \frac{c_j}{\Lambda^2} + \sum_{j,k} s_{jk} \frac{c_j}{\Lambda^2} \frac{c_k}{\Lambda^2}$$

Pure SM
Interference with SM
Pure NP

- Could solve for  $s_0, s_j, s_{jk}$  by generating samples at distinct points in the WC phase space
- Far too computationally intensive. Would need  $\mathcal{O}(100)$  MC samples per signal process!

# EFT Event Weights

- Instead can calculate the dependence on a per event basis
- Parameterization extends to the event level where each event will then have a **weight function**:

The parameterization is similar to the one from before, but is now done per event!

$$w_i(\vec{c}) = s_{0i} + \sum_j s_{ij} \frac{c_j}{\Lambda^2} + \sum_{j,k} s_{ijk} \frac{c_j}{\Lambda^2} \frac{c_k}{\Lambda^2}$$

Pure SM
Interference with SM
Pure NP

- Can sum individual weight functions for events passing a given selection ( $\mathcal{S}$ ):

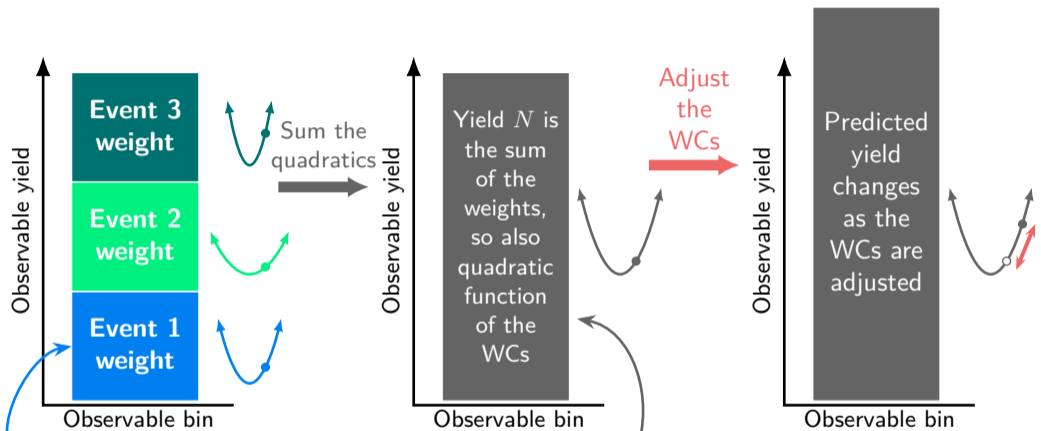
$N$  is a function of the WCs, so can predict the yield for this selection for any arbitrary values of the WCs

$$N(\vec{c}) = \sum_{i \in \mathcal{S}} w_i(\vec{c})$$

$N$  is the sum of a bunch of quadratics, so it is also a quadratic



# Scaling EFT Contributions



$$w_i(\vec{c}) = s_{0i} + \sum_j s_{ij} \frac{c_j}{\Lambda^2} + \sum_{j,k} s_{ijk} \frac{c_j}{\Lambda^2} \frac{c_k}{\Lambda^2}$$

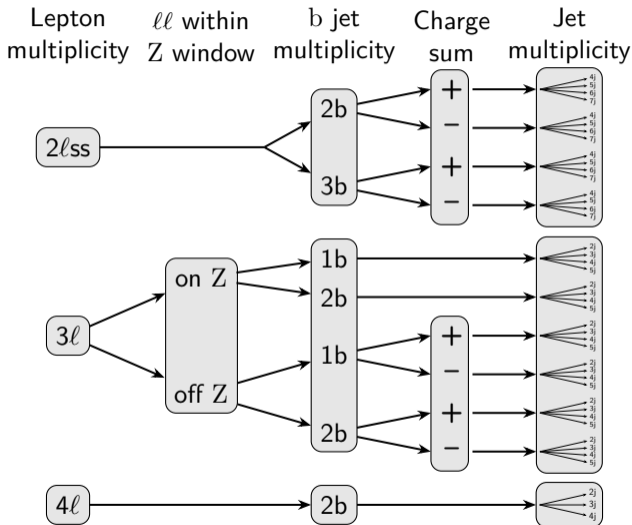
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$$N(\vec{c}) = \sum_{i \in S} w_i(\vec{c})$$

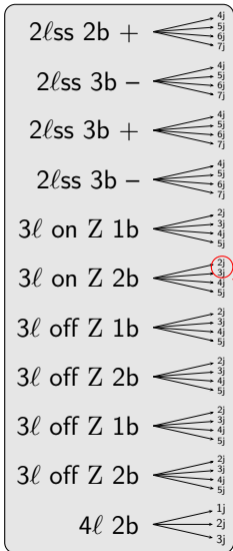
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# Event Selection

- The event selection targets prompt well isolated leptons produced with b jets and additional other jets
- Aims to discriminate between signal processes as much as possible
  - $2lss$ :  $t\bar{t}H$  and  $t\bar{t}W$  (split by charge)
  - $3l$  on Z:  $t\bar{t}ll$  (2b),  $tllq$  (1b)
  - $3l$  off Z: non-resonant  $t\bar{t}ll$  and  $tllq$  (2-quark 2-lepton EFT contributions)
  - $\geq 4l$ :  $t\bar{t}H$  and  $t\bar{t}ll$

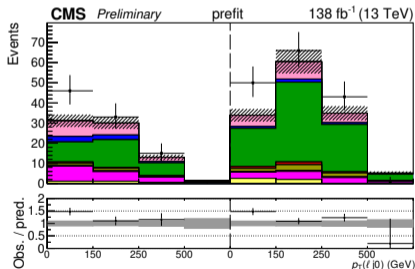


# Event Selection



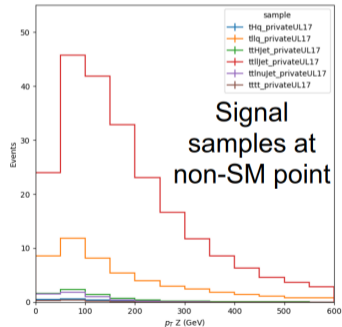
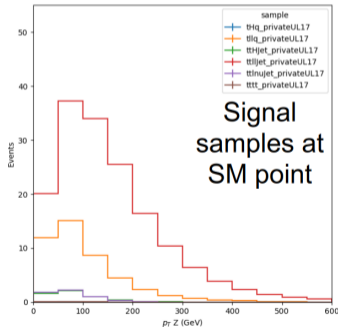
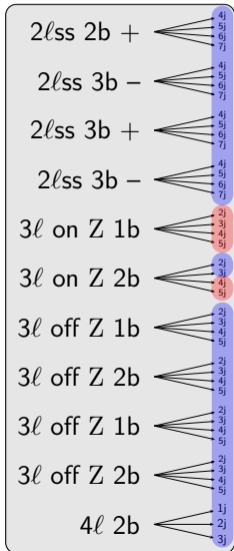
- We use a differential kinematic distribution for each of the 43 jet categories
- Kinematic variable for each jet category is chosen to optimize sensitivity to EFT effects

Differential distribution

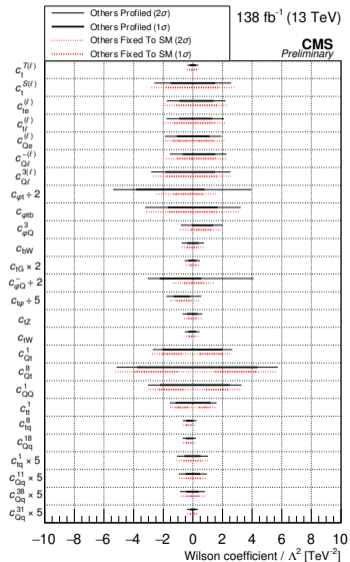


# Event Selection

- Use different variables for different categories to optimize sensitivity to EFT effects
  - $p_T(\ell j_0) \rightarrow p_T$  of the leading lepton plus jet pairs
  - $p_T(Z) \rightarrow p_T$  of the opposite sign lepton pair



- The solid black bars show the limits for each WC where the other WCs are **profiled**
- The dashed red bars show the limits when the other WCs are **frozen to the SM**
- Results are consistent with the SM
- Disjoint  $1\sigma$  intervals for the 4-heavy quark operators is due to the quadratic nature of the EFT parameterization
  - In principle this is true for all WCs
  - Varied signal processes and kinematic distributions help resolve these double minima





- EFT can be a powerful technique for indirect searches of BSM physics at the LHC
- EFT approaches stand to benefit greatly from increased statistics of the current LHC and future HL-LHC runs
- Showcased a search by CMS for new physics in associated top production processes using EFT
  - Data corresponds to  $138 \text{ fb}^{-1}$  of pp collisions collected by CMS
  - Set simultaneous confidence interval limits on 26 WCs associated with top quarks
  - All limits found to be consistent with SM expectations
- There are a number of potential future directions for expanding and improving the analysis:
  - More data!
  - Optimization of the categorization and kinematic variable choices
  - Adding additional signal processes and final states
  - Translate the limits back into constraints on potential BSM models

# Backup

# Modeling EFT Contributions

- The EFT MC samples for the six signal processes ( $t\bar{t}H$ ,  $t\bar{t}l\nu$ ,  $t\bar{t}ll$ ,  $tllq$ ,  $tHq$ ,  $t\bar{t}t\bar{t}$ ) are generated using the **dim6top model** (1802.07237) to estimate relevant EFT effects
  - Uses the Warsaw basis
  - LO calculation, so include an extra jet in the matrix element (when possible) to improve modeling at high jet-multiplicities
  - Include 26 WCs which were found to significantly impact the signal processes

Operator category	WCs
Two heavy quarks	$c_{t\varphi}, c_{\varphi Q}^-, c_{\varphi Q}^3, c_{\varphi t}, c_{\varphi tb}, c_{tW}, c_{tZ}, c_{bW}, c_{tG}$
Two heavy quarks two leptons	$c_{Ql}^{3(\ell)}, c_{Ql}^{-\ell}, c_{Qe}^{(\ell)}, c_{tl}^{(\ell)}, c_{te}^{(\ell)}, c_t^{S(\ell)}, c_t^{T(\ell)}$
Two light quarks two heavy quarks	$c_{Qq}^{31}, c_{Qq}^{38}, c_{Qq}^{11}, c_{Qq}^{18}, c_{tq}^1, c_{tq}^8$
Four heavy quarks	$c_{QQ}^1, c_{Qt}^1, c_{Qt}^8, c_{tt}^1$

- Generated  $\sim 300M$  private MC events in total using the same configurations as central CMS samples



# Event Selection Cuts

Event category	Leptons	$m_{\ell\ell}$	b tags	Lepton charge sum	Jets	Differential variable
2lss 2b	2	No requirement	2	$> 0, < 0$	4,5,6, $\geq 7$	$p_T(\ell j_0)$
2lss 3b	2	No requirement	$\geq 3$	$> 0, < 0$	4,5,6, $\geq 7$	$p_T(\ell j_0)$
3l off-Z 1b	3	$ m_Z - m_{\ell\ell}  > 10 \text{ GeV}$	1	$> 0, < 0$	2,3,4, $\geq 5$	$p_T(\ell j_0)$
3l off-Z 2b	3	$ m_Z - m_{\ell\ell}  > 10 \text{ GeV}$	$\geq 2$	$> 0, < 0$	2,3,4, $\geq 5$	$p_T(\ell j_0)$
3l on-Z 1b	3	$ m_Z - m_{\ell\ell}  \leq 10 \text{ GeV}$	1	No requirement	2,3,4, $\geq 5$	$p_T(Z)$
3l on-Z 2b	3	$ m_Z - m_{\ell\ell}  \leq 10 \text{ GeV}$	$\geq 2$	No requirement	2,3,4, $\geq 5$	$p_T(Z)$ or $p_T(\ell j_0)$
4l	$\geq 4$	No requirement	$\geq 2$	No requirement	2,3, $\geq 4$	$p_T(\ell j_0)$

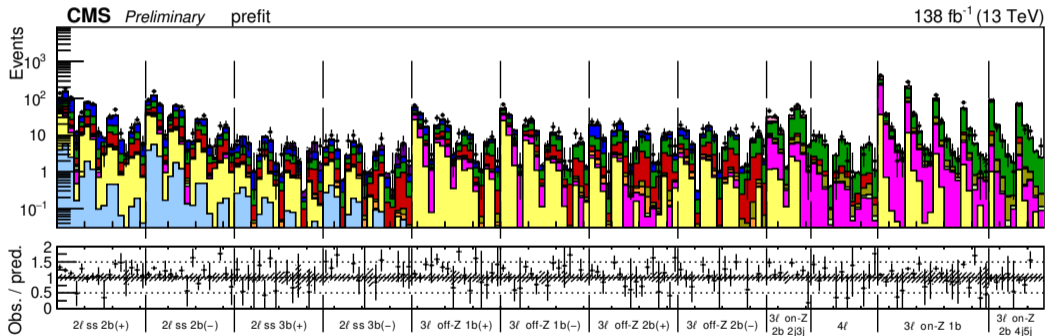
- All jets required to have  $|\eta| < 2.4$  and  $p_T > 30 \text{ GeV}$
- All electrons require  $|\eta| < 2.5$  and all muons require  $|\eta| < 2.4$
- Lepton  $p_T$  cuts ( GeV):
  - 2lss:  $p_T > 25, 15$
  - 3l for e ( $\mu$ ):  $p_T > 25, 15, 15$  (10)
  - 4l for e ( $\mu$ ):  $p_T > 25, 15, 15$  (10), 15 (10)

# Interpretation of Sensitivity

- Sensitivity to most WCs comes from a wide range of bins over all selection categories

Grouping of WCs	WCs	Lead categories
Two heavy two leptons	$c_{Ql}^{3(\ell)}, c_{Ql}^{- (\ell)}, c_{Qe}^{(\ell)}, c_{tl}^{(\ell)}, c_{te}^{(\ell)}, c_t^{S(\ell)}, c_t^{T(\ell)}$	3l off-Z
Four heavy	$c_{QQ}^1, c_{Qt}^1, c_{Qt}^8, c_{tt}^1$	2lss
Two heavy two light “t $\bar{t}$ l $\nu$ -like”	$c_{Qq}^{11}, c_{Qq}^{18}, c_{tq}^1, c_{tq}^8$	2lss
Two heavy two light “tllq-like”	$c_{Qq}^{31}, c_{Qq}^{38}$	3l on-Z
Two heavy with bosons “t $\bar{t}$ ll-like”	$c_{tZ}, c_{\varphi t}, c_{\varphi Q}^-$	3l on-Z and 2lss
Two heavy with bosons “tXq-like”	$c_{\varphi Q}^3, c_{\varphi tb}, c_{bW}$	3l on-Z
Two heavy with bosons with significant impacts on many processes	$c_{tG}, c_{t\varphi}, c_{tW}$	3l and 2lss

# Signal Region Yields



# Signal Region Yields

