



Quantum computing approaches for simulating parton showers in high energy collisions

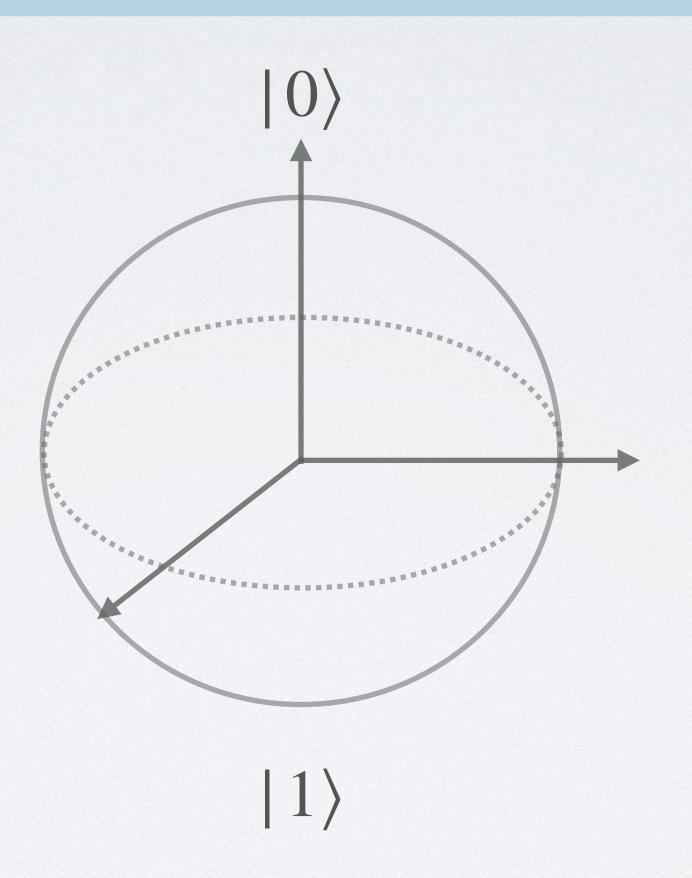
Simon Williams

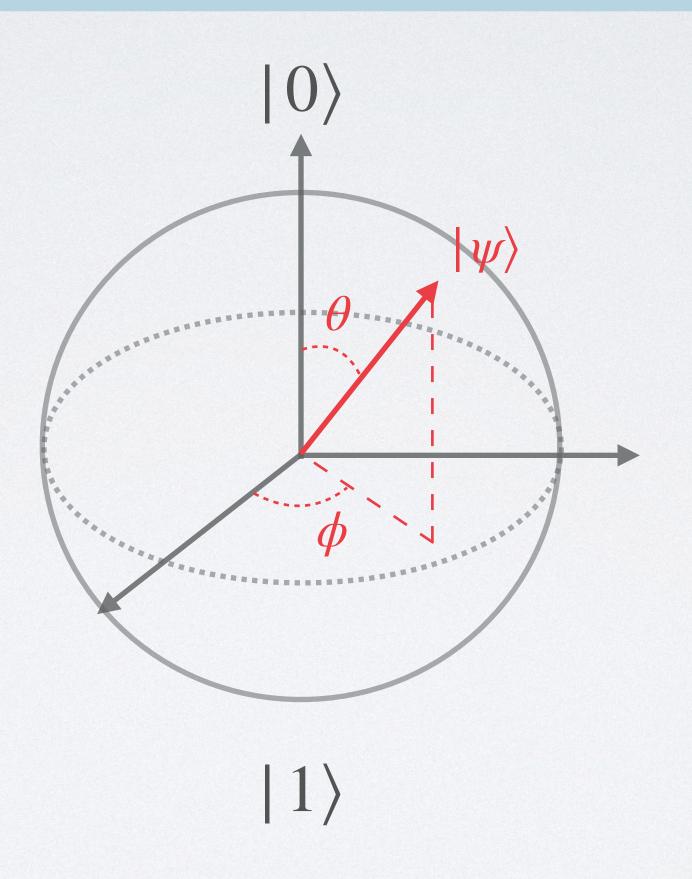
Lake Louise Winter Institute - 25/02/22

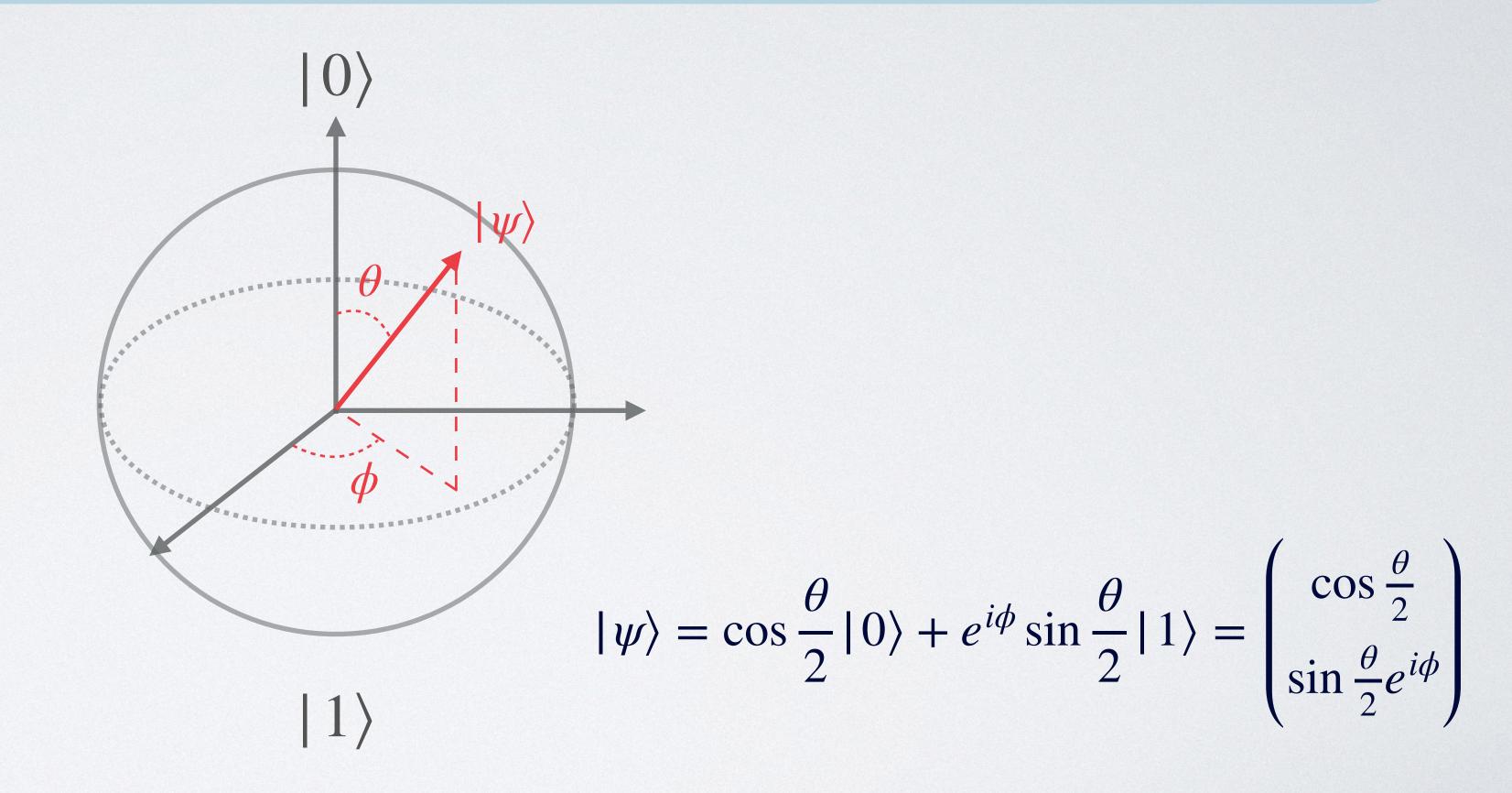
Contents

- The Power of the Qubit!
 - The Quantum Walk Framework
 - Why are we interested in High Energy Physics?
- The Parton Shower
- Quantum Walk approach to the parton shower [1]
- Looking to the Future

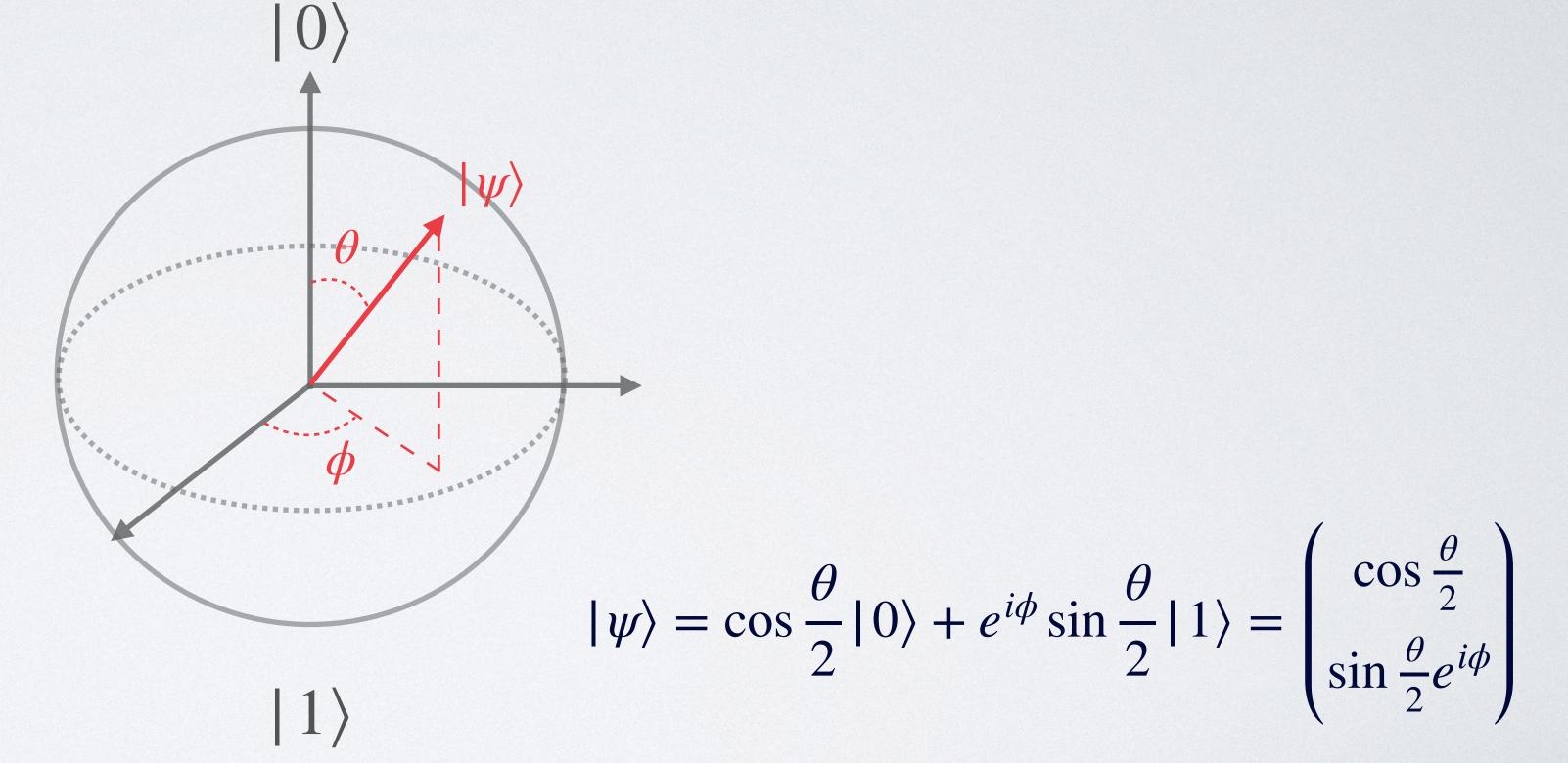
- [1] A quantum walk approach to simulating parton showers, arXiv: 2109.13975
- In collaboration with Sarah Malik (UCL), Michael Spannowsky (IPPP, Durham) and Khadeejah Bepari (IPPP, Durham)







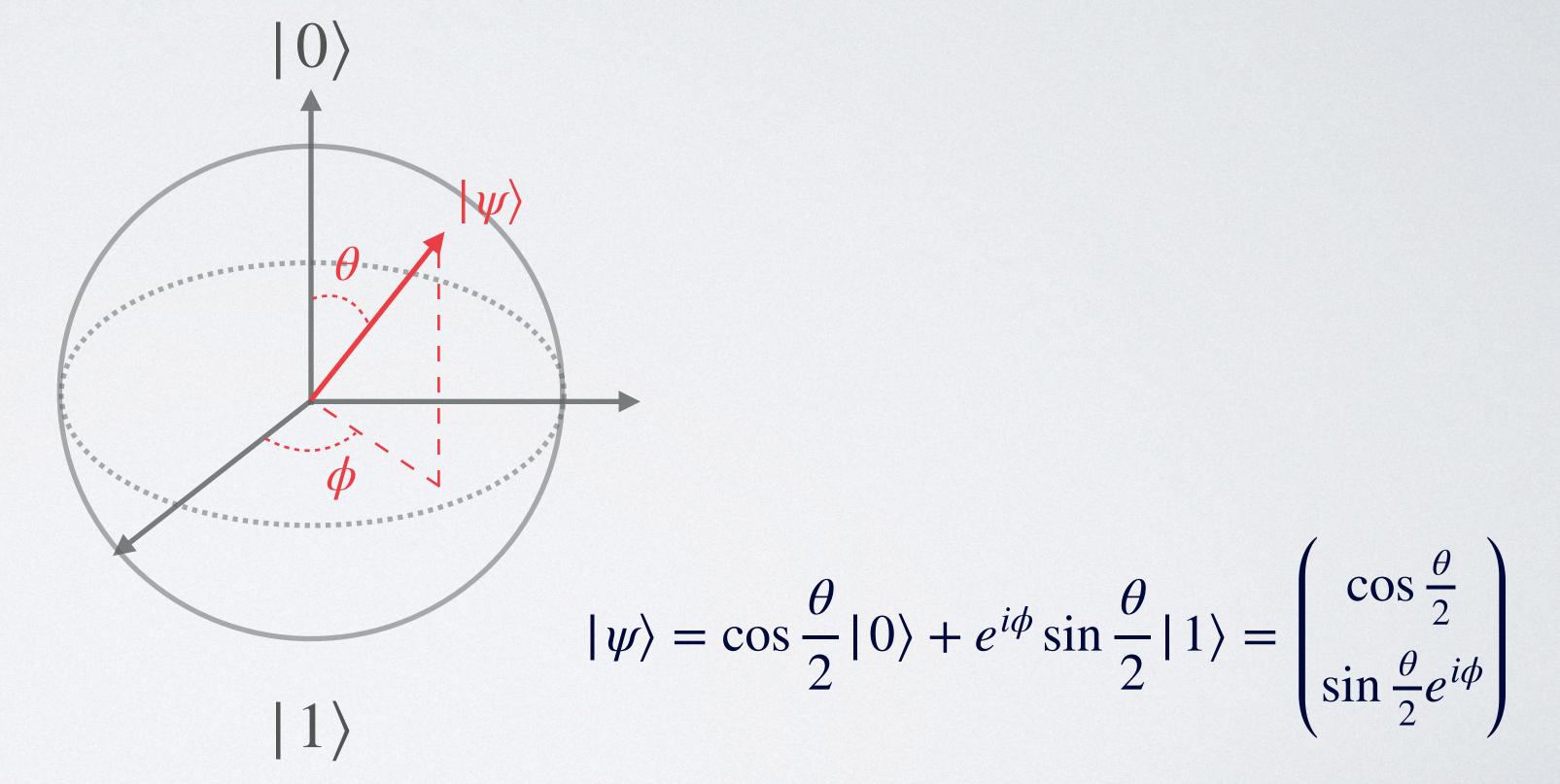
$$U_{3}(\theta, \phi, \lambda) = \begin{pmatrix} \cos(\frac{\theta}{2}) & -e^{i\lambda}\sin(\frac{\theta}{2}) \\ e^{i\phi}\sin(\frac{\theta}{2}) & e^{i(\phi+\lambda)}\cos(\frac{\theta}{2}) \end{pmatrix}$$

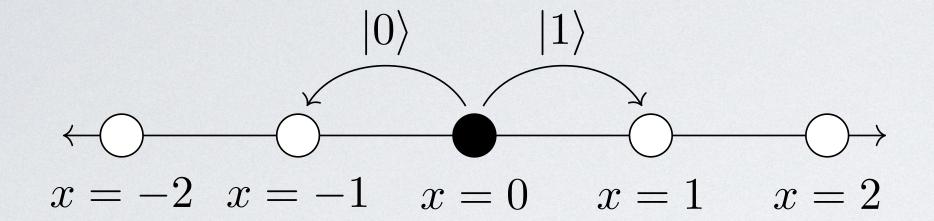


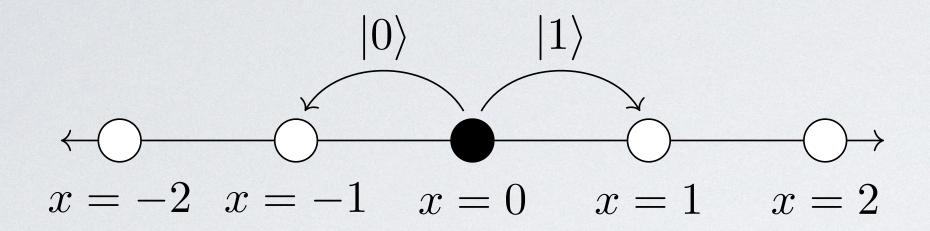
• Qubit: quantum analogue of classical bit, not restricted only to being in either the $|0\rangle$ or $|1\rangle$ state

$$U_{3}(\theta, \phi, \lambda) = \begin{pmatrix} \cos(\frac{\theta}{2}) & -e^{i\lambda}\sin(\frac{\theta}{2}) \\ e^{i\phi}\sin(\frac{\theta}{2}) & e^{i(\phi+\lambda)}\cos(\frac{\theta}{2}) \end{pmatrix}$$

• Extending this to a system of N qubits forms a 2^N -dimensional Hilbert Space



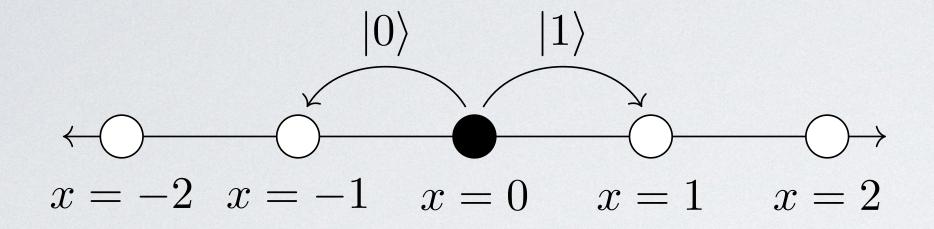




$$\mathcal{H}_{P} = \{ | i \rangle : i \in \mathbb{Z} \}$$

$$\mathcal{H}_{C} = \{ | 0 \rangle, | 1 \rangle \}$$

$$\mathcal{H}_{C} = \{ | 0 \rangle, | 1 \rangle \}$$

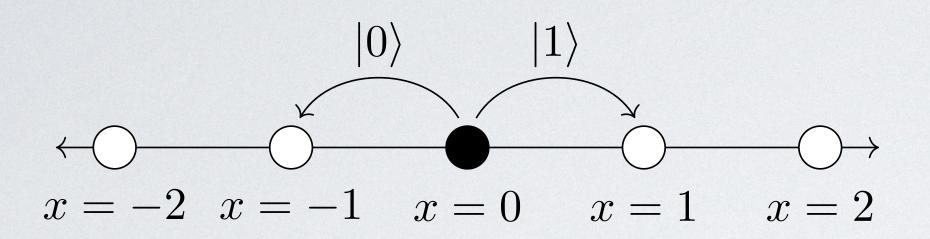


$$\mathcal{H}_{P} = \{ | i \rangle : i \in \mathbb{Z} \}$$

$$\mathcal{H}_{C} = \{ | 0 \rangle, | 1 \rangle \}$$

$$\mathcal{H}_{C} = \mathcal{H}_{C} \otimes \mathcal{H}_{P}$$

$$U = S \cdot (C \otimes I)$$

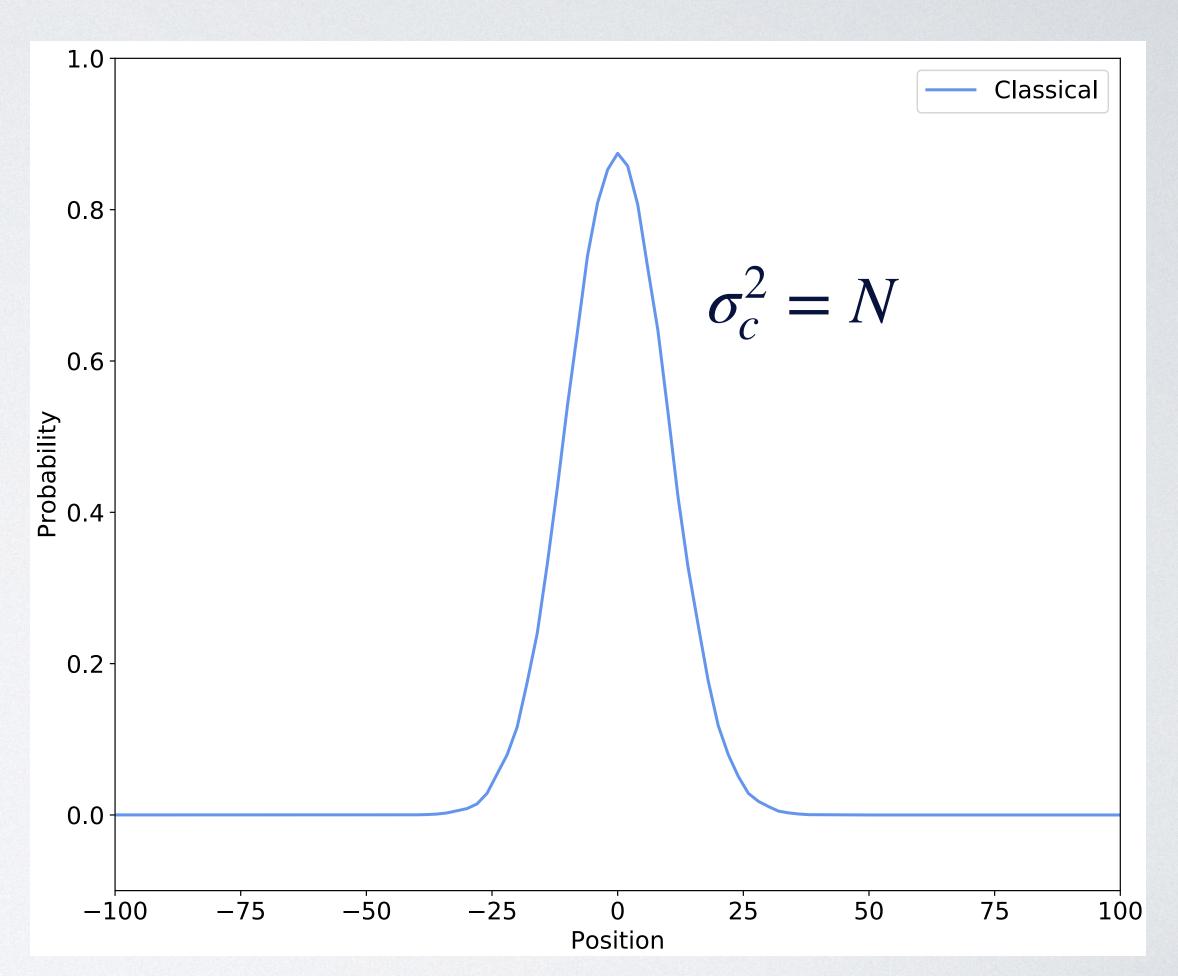


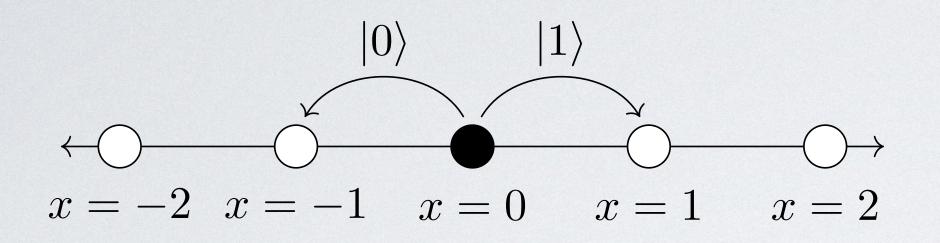
$$\mathcal{H}_{P} = \{ | i \rangle : i \in \mathbb{Z} \}$$

$$\mathcal{H}_{C} = \{ | 0 \rangle, | 1 \rangle \}$$

$$\mathcal{H}_{C} = \{ | 0 \rangle, | 1 \rangle \}$$

$$U = S \cdot (C \otimes I)$$



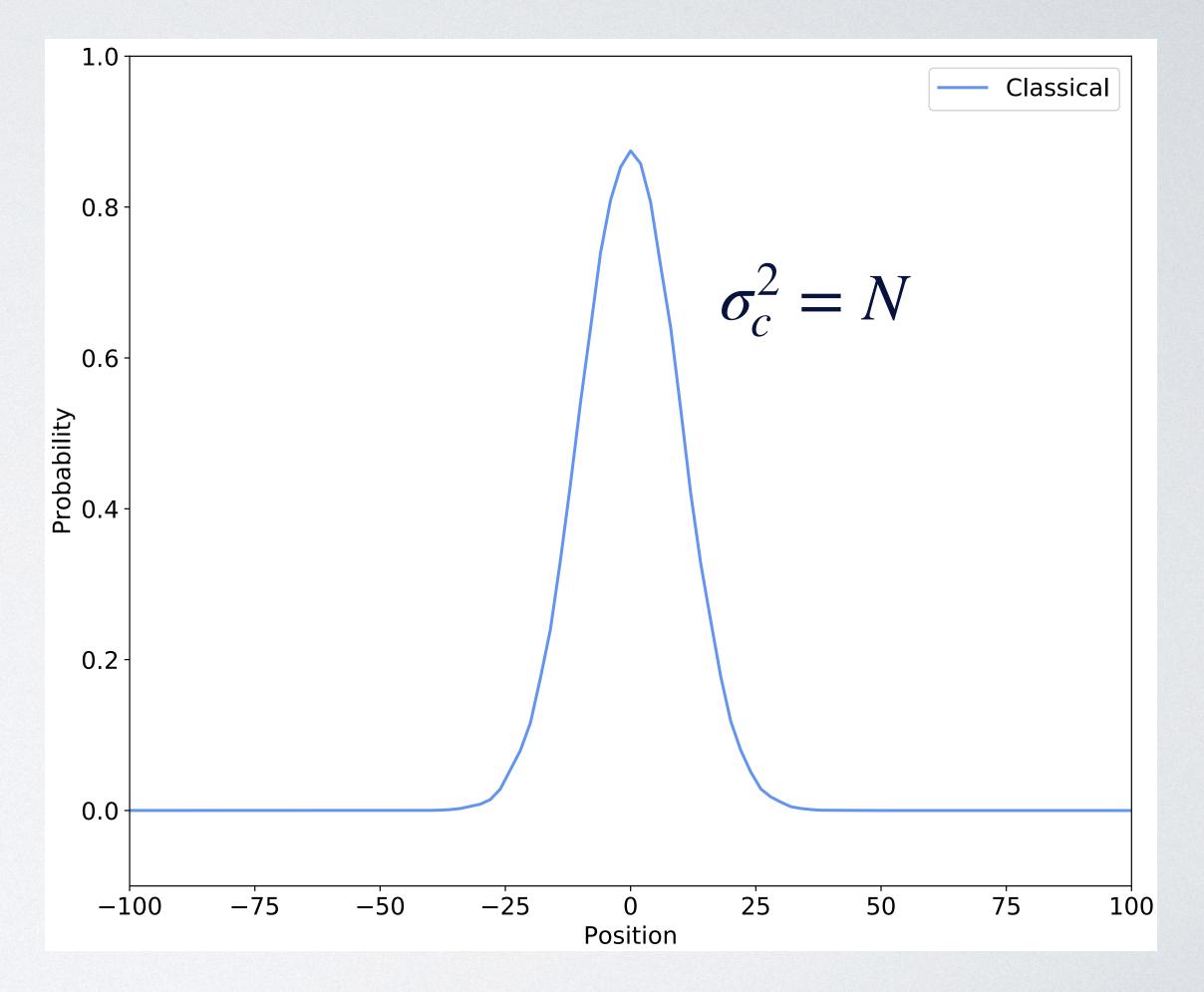


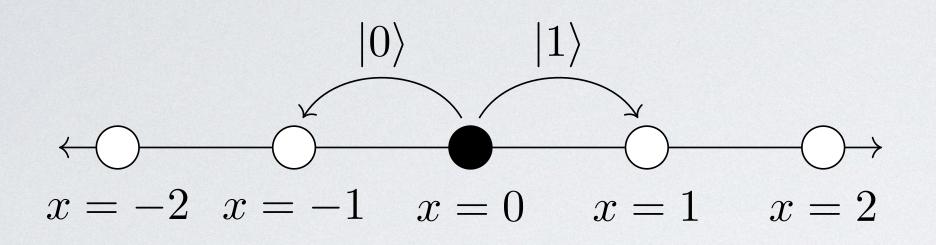
$$\mathcal{H}_{P} = \{ | i \rangle : i \in \mathbb{Z} \}$$

$$\mathcal{H}_{C} = \{ | 0 \rangle, | 1 \rangle \}$$

$$\mathcal{H}_{C} = \{ | 0 \rangle, | 1 \rangle \}$$

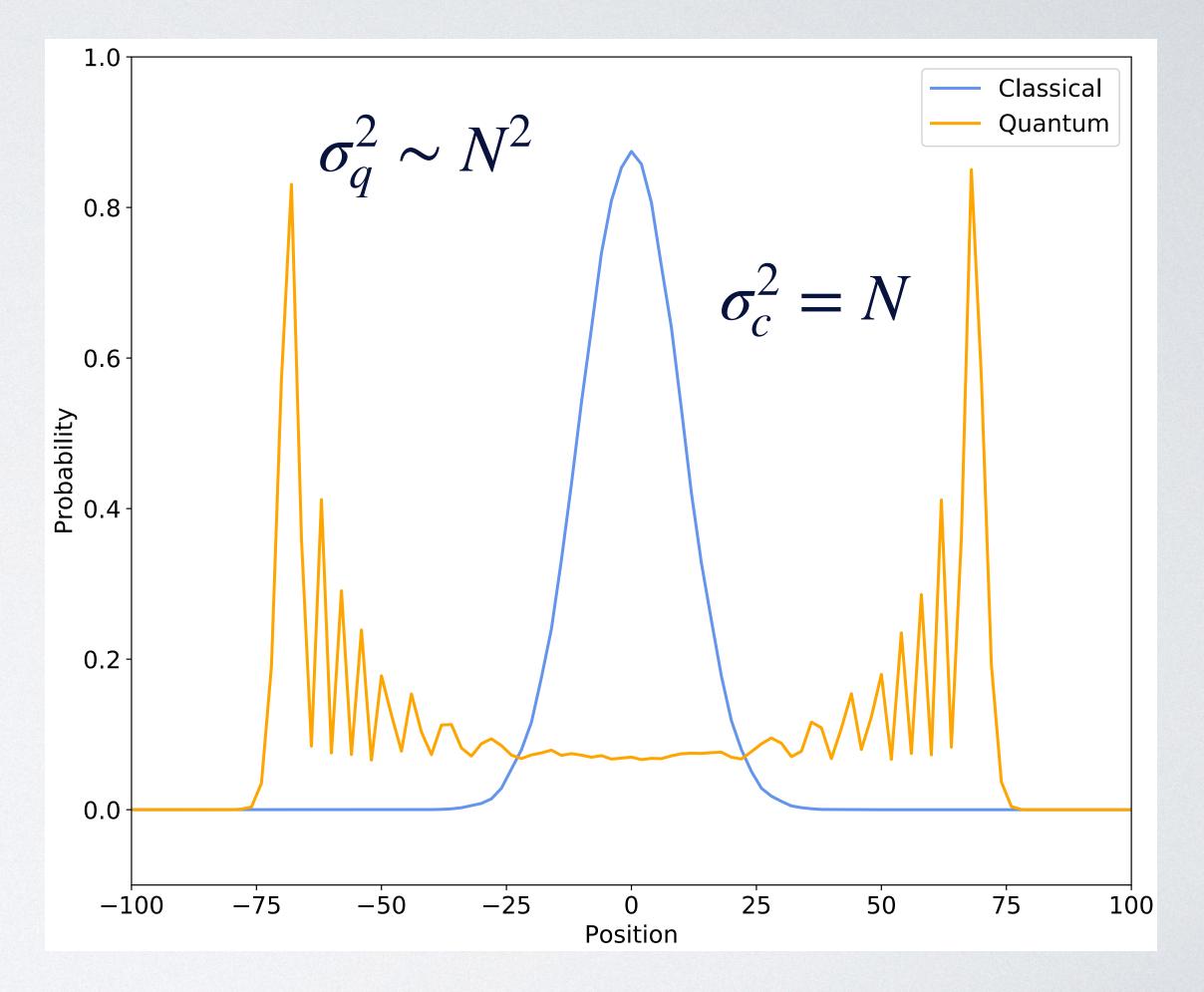
$$U = S \cdot (C \otimes I) \qquad |\varphi\rangle \left\{ \begin{array}{c} |\varphi_0\rangle - \\ \vdots \\ |\varphi_n\rangle - \end{array} \right. \qquad (S \cdot (C \otimes I))^N \qquad \vdots$$

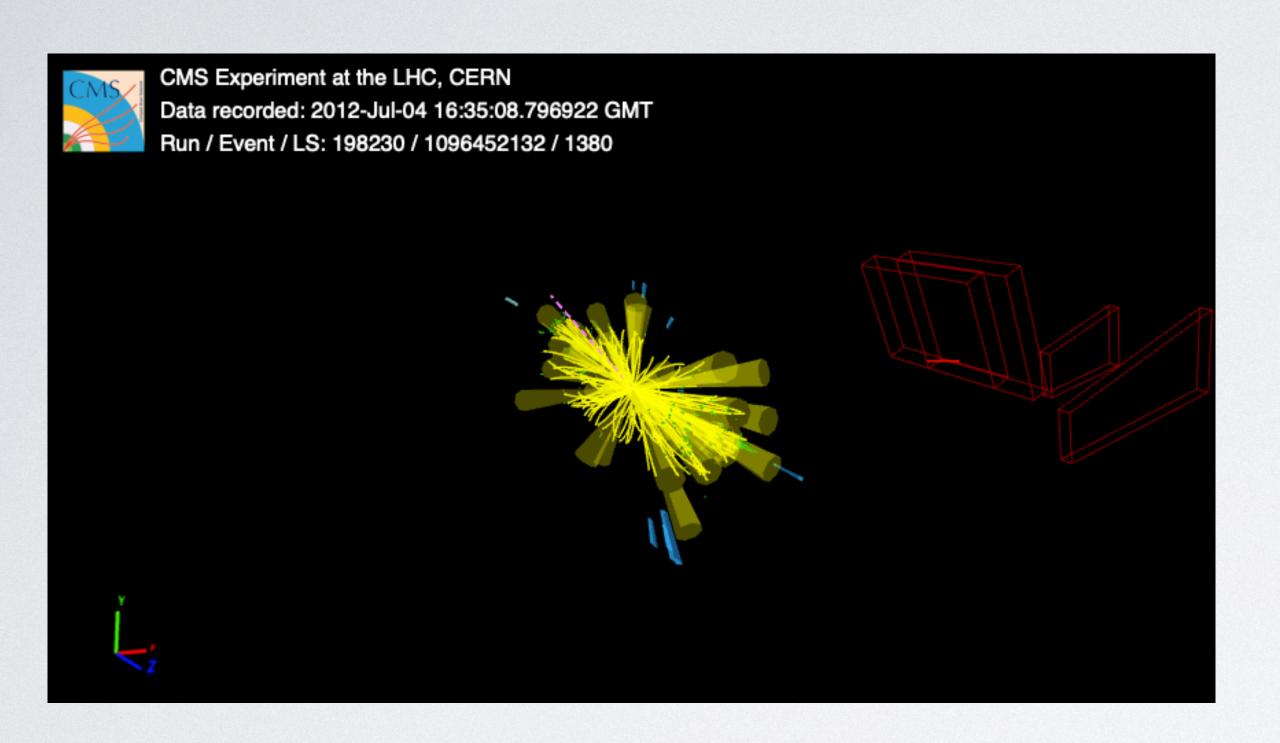


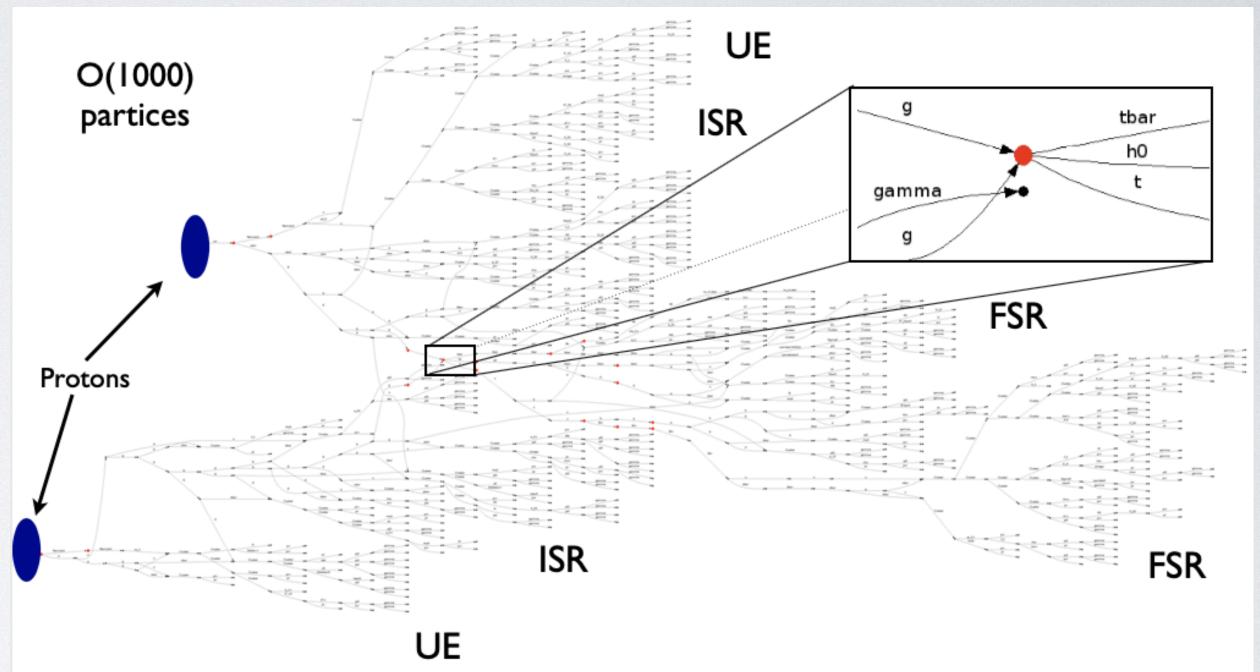


$$\begin{aligned} \mathcal{H}_P &= \{ \mid i \rangle : i \in \mathbb{Z} \} \\ \mathcal{H}_C &= \{ \mid 0 \rangle, \mid 1 \rangle \} \end{aligned} \\ \mathcal{H}_C &= \{ \mid 0 \rangle, \mid 1 \rangle \}$$

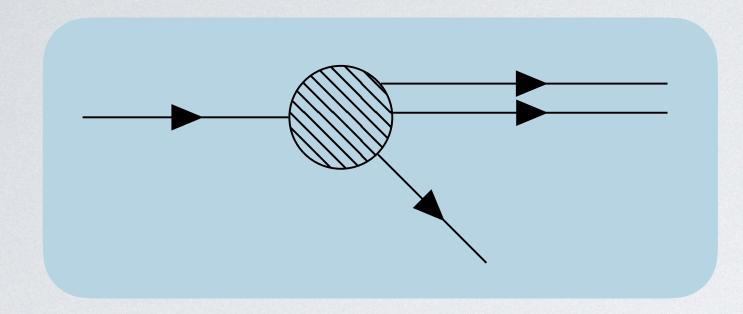
$$U = S \cdot (C \otimes I) \qquad |\varphi\rangle \left\{ \begin{array}{c} |\varphi_0\rangle - \\ \vdots \\ |\varphi_n\rangle - \end{array} \right. \qquad (S \cdot (C \otimes I))^N \qquad \vdots$$

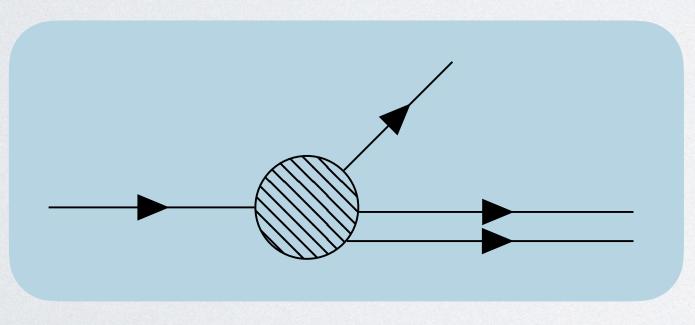






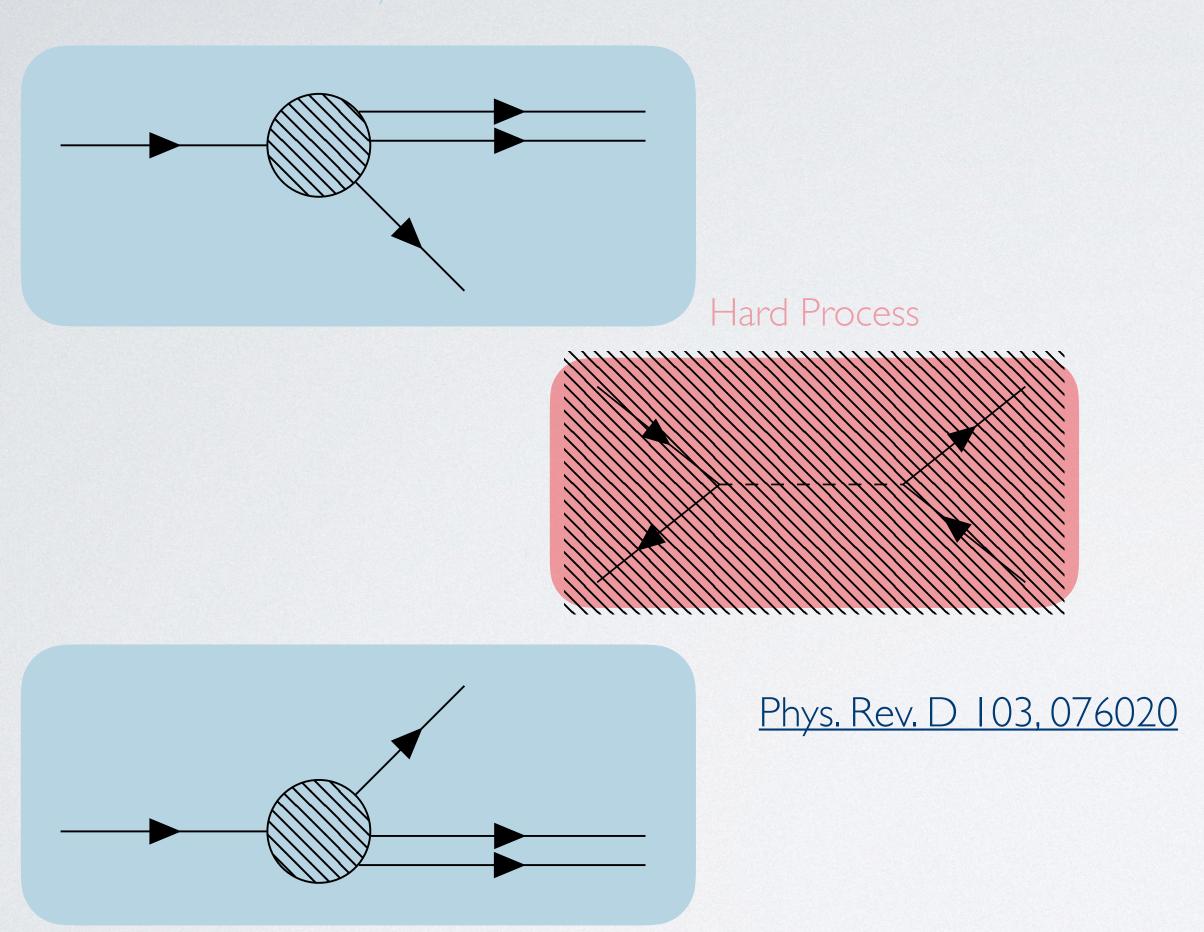
Parton Density Functions





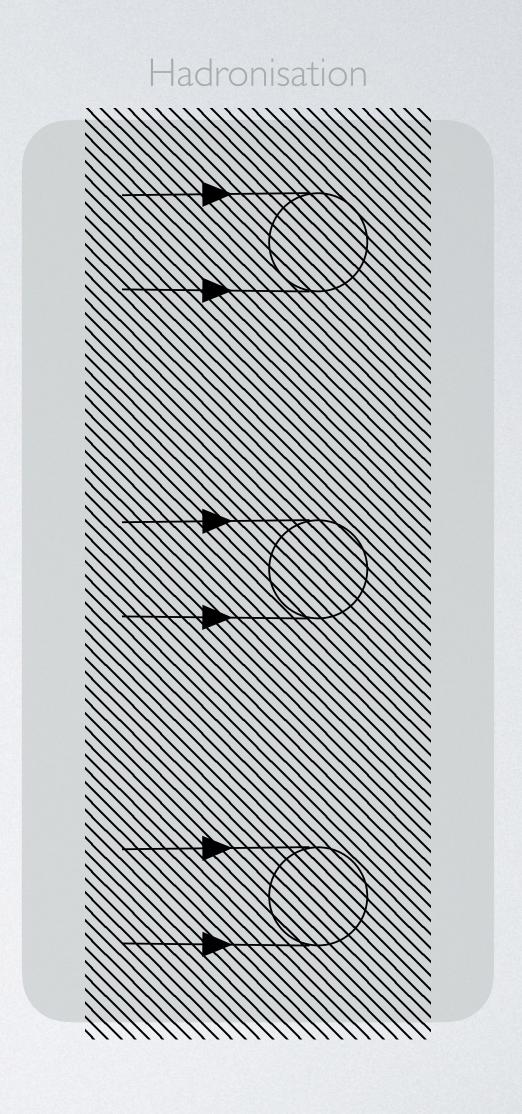
Phys. Rev. D 103, 034027

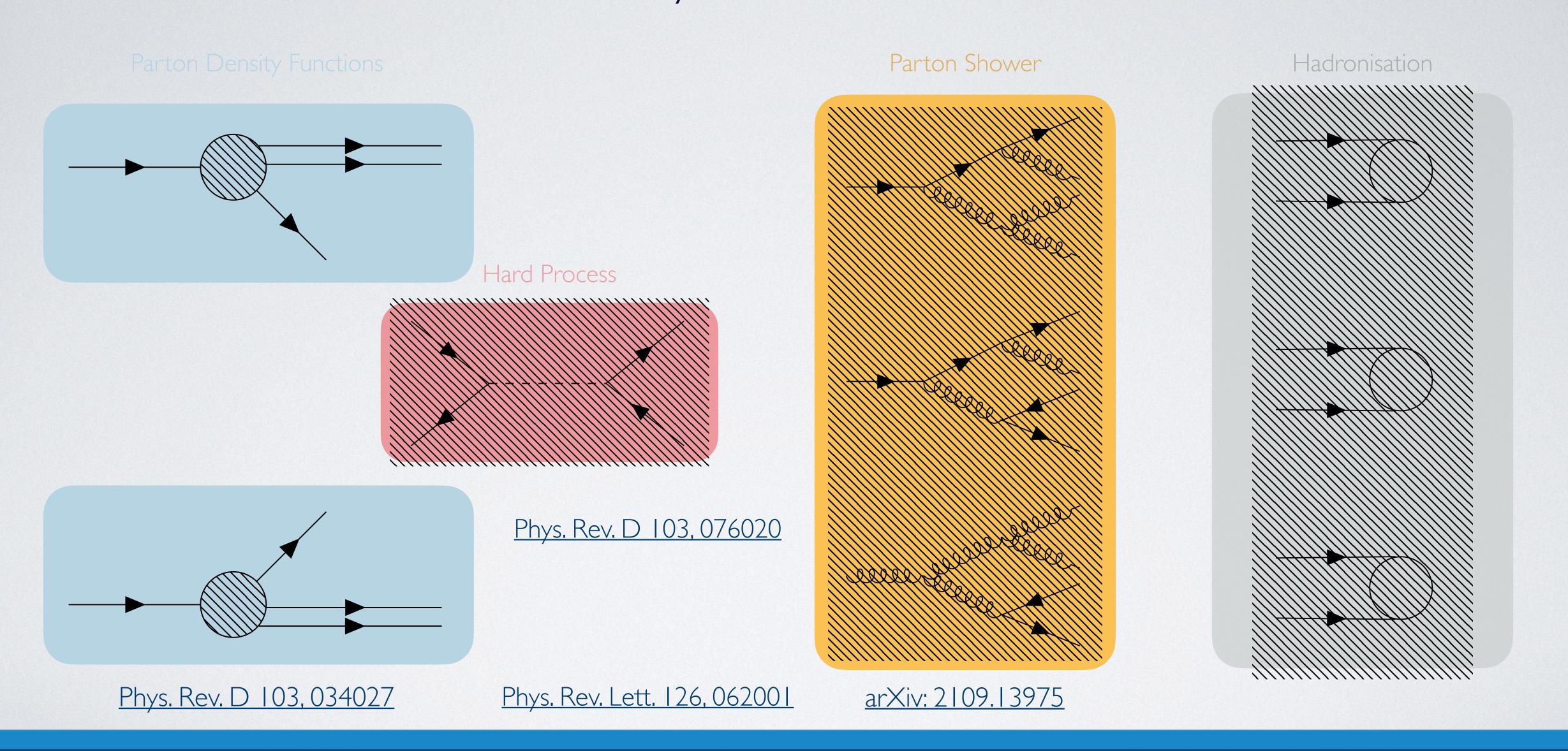
Parton Density Functions

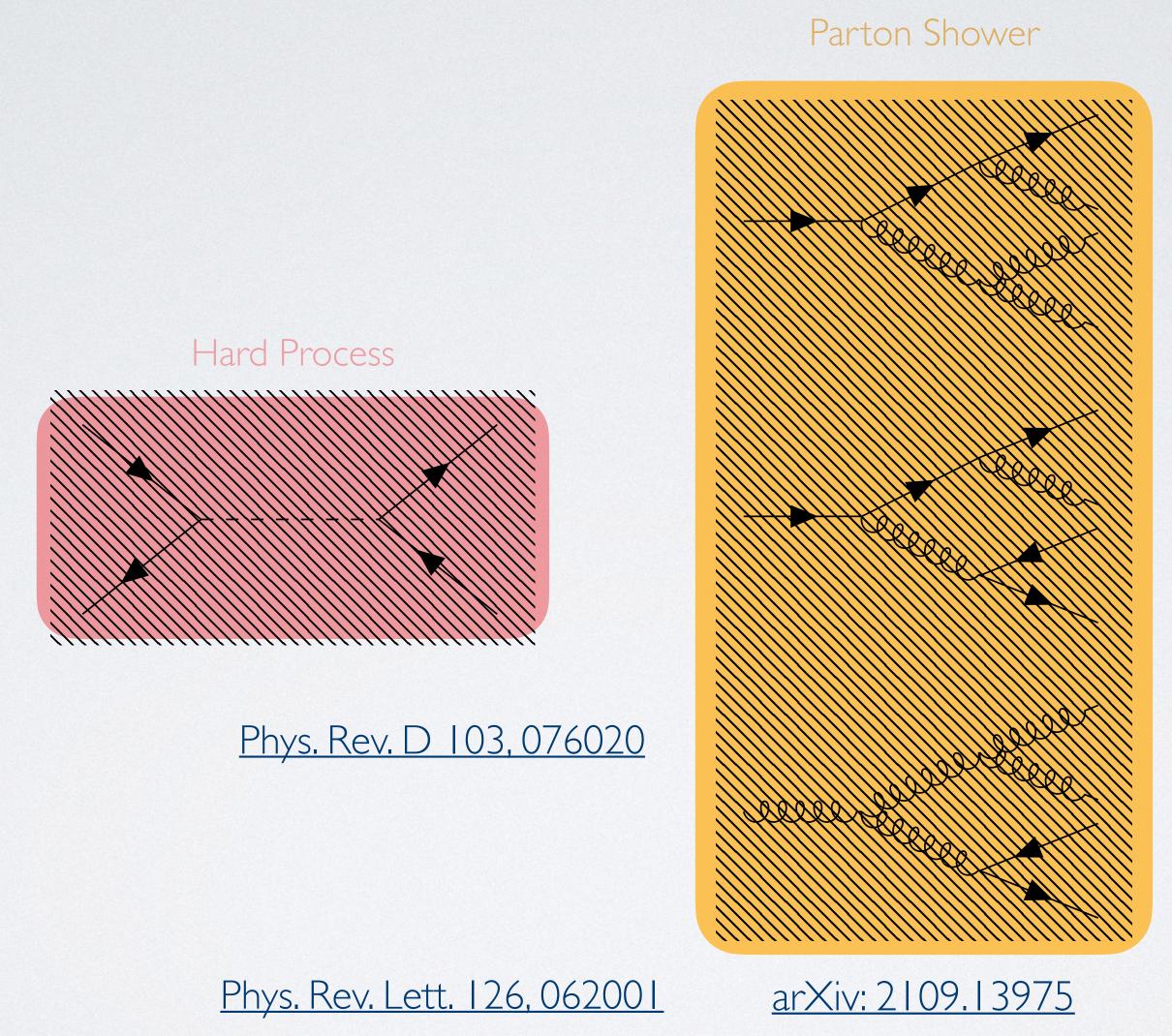


Phys. Rev. D 103, 034027

Hard Process Phys. Rev. D 103, 076020 Phys. Rev. D 103, 034027



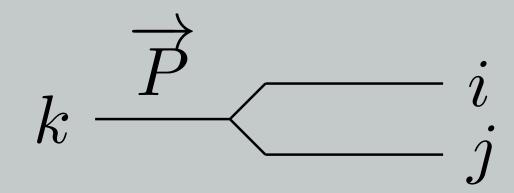




· We present a discrete, collinear toy QCD model comprising one gluon and one quark flavour

- We present a discrete, collinear toy QCD model comprising one gluon and one quark flavour
- To meet current QC qubit restrictions, only collinear splittings have been considered, meaning we do not keep track of individual kinematics

Collinear Condition:



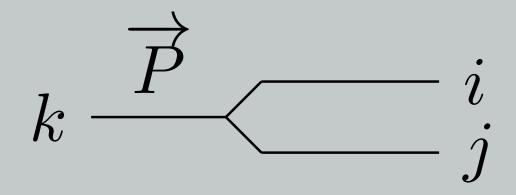
$$p_i = zP$$
,

$$p_i = zP,$$

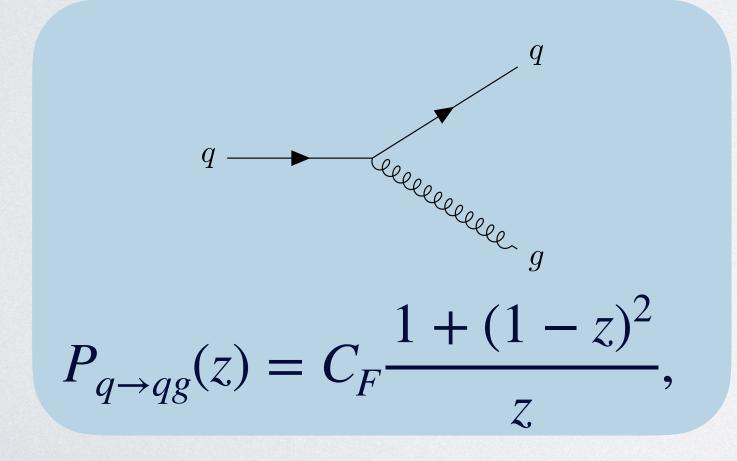
$$p_j = (1 - z)P$$

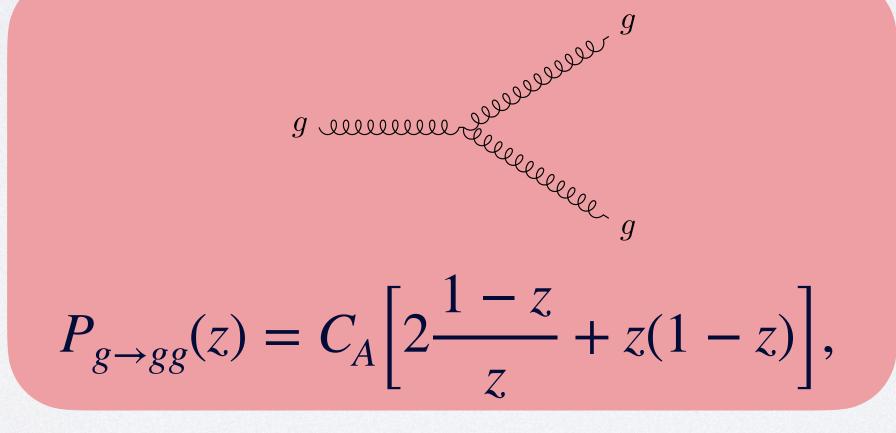
- · We present a discrete, collinear toy QCD model comprising one gluon and one quark flavour
- To meet current QC qubit restrictions, only collinear splittings have been considered, meaning we do not keep track of individual kinematics

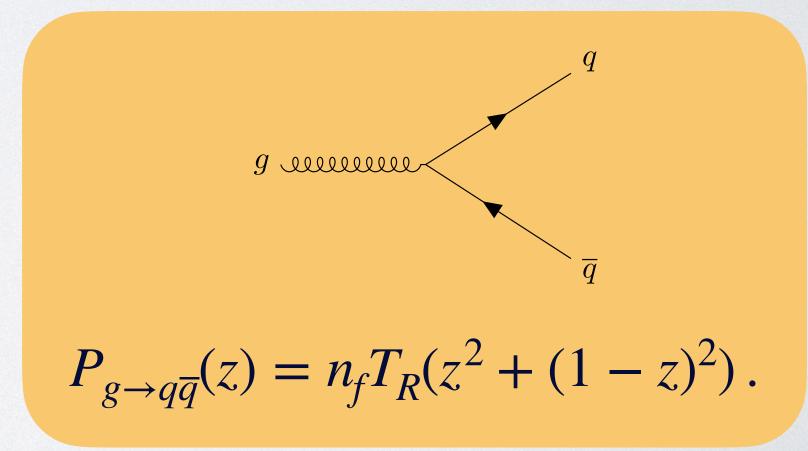
Collinear Condition:

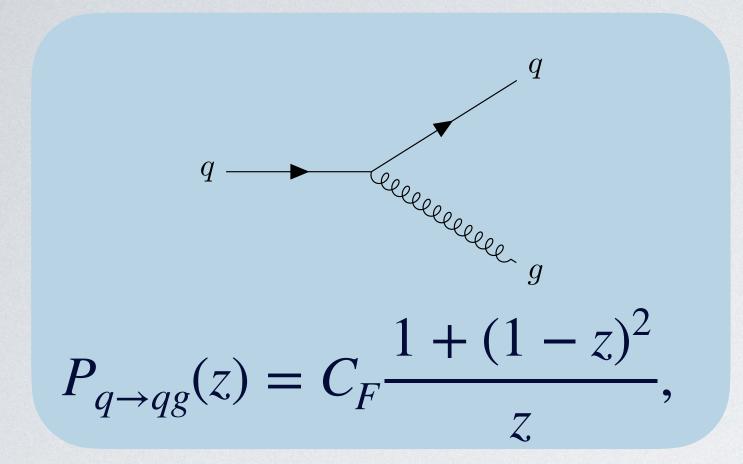


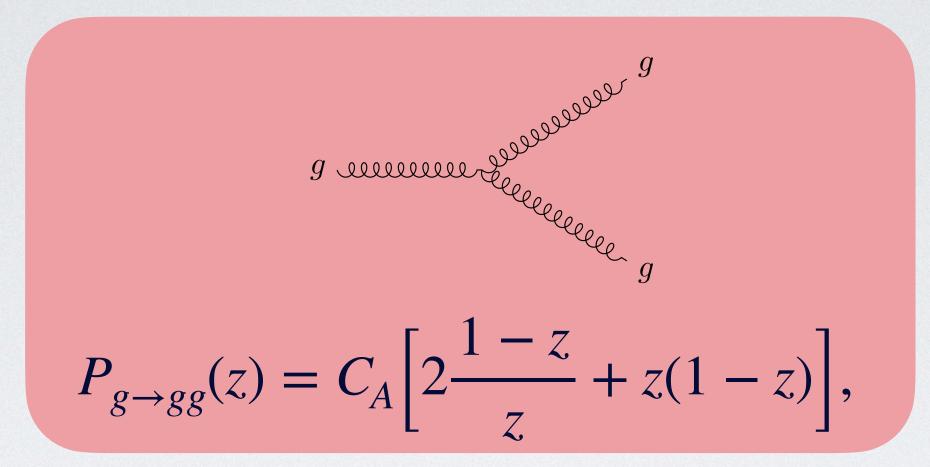
$$p_i = zP,$$

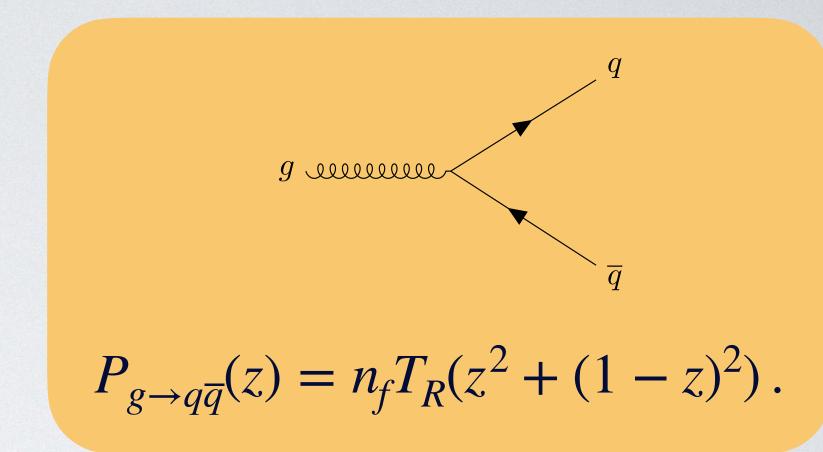


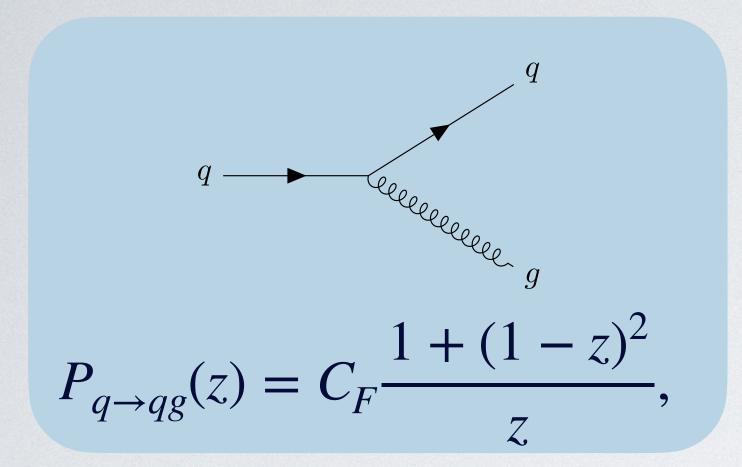


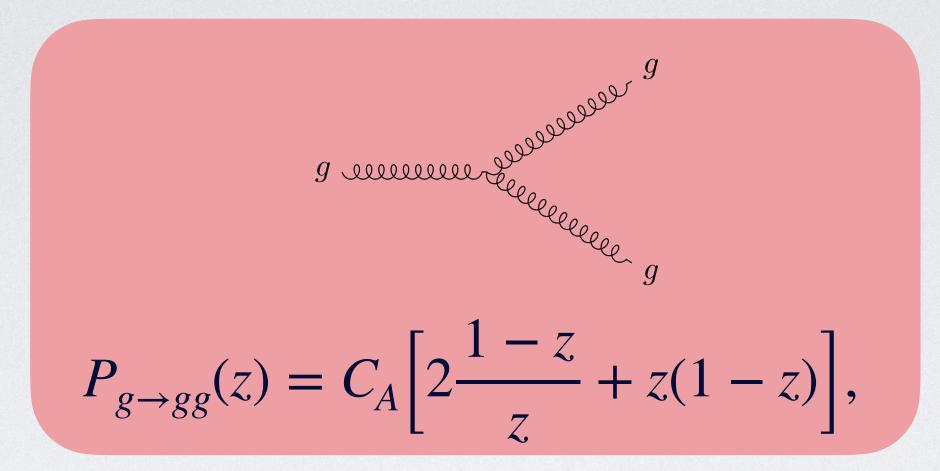


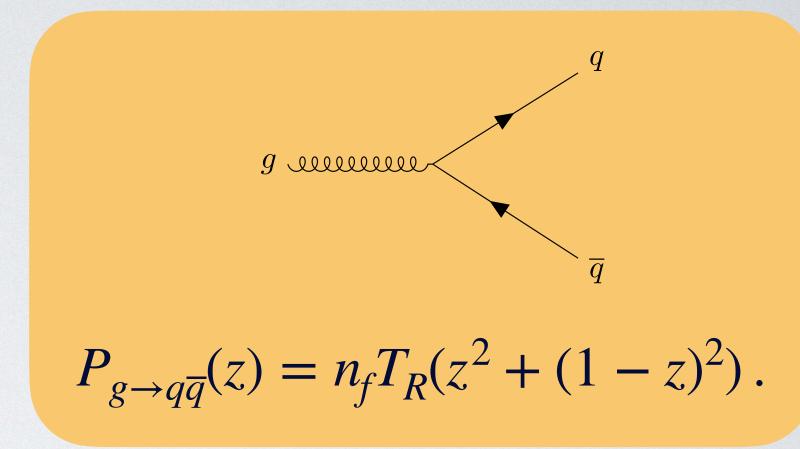






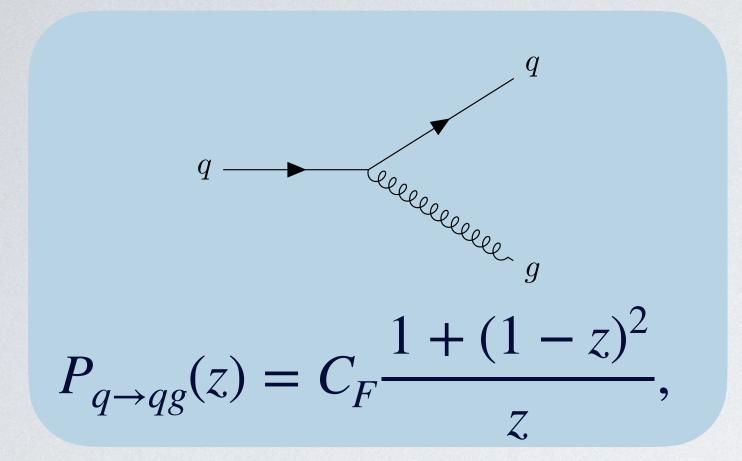


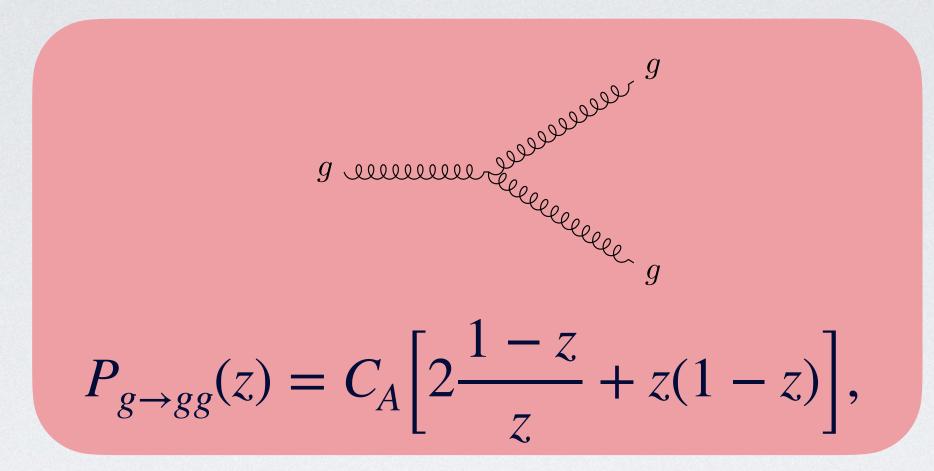


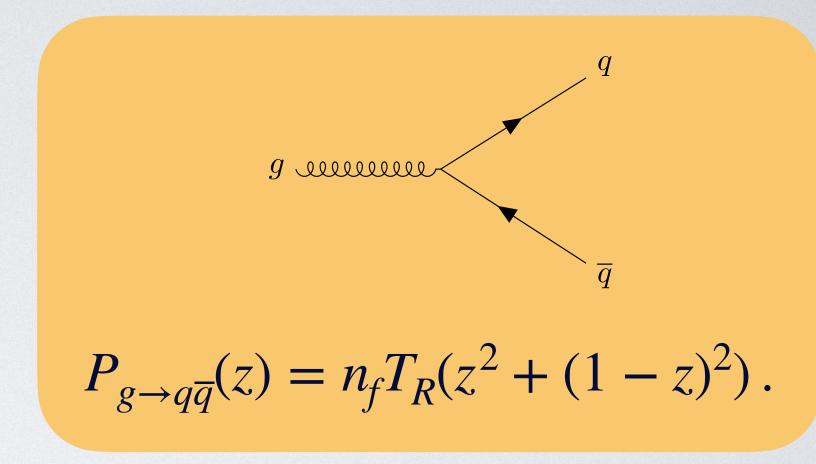


The Sudakov factors have been used to determine whether an emission occurs:

$$\Delta_{i,k}(z_1, z_2) = \exp\left[-\alpha_s \int_{z_1}^{z_2} P_k(z') dz'\right], \qquad \Delta_{tot}(z_1, z_2) = \Delta_g^{n_g}(z_1, z_2) \Delta_q^{n_q}(z_1, z_2) \Delta_{\overline{q}}^{n_{\overline{q}}}(z_1, z_2).$$







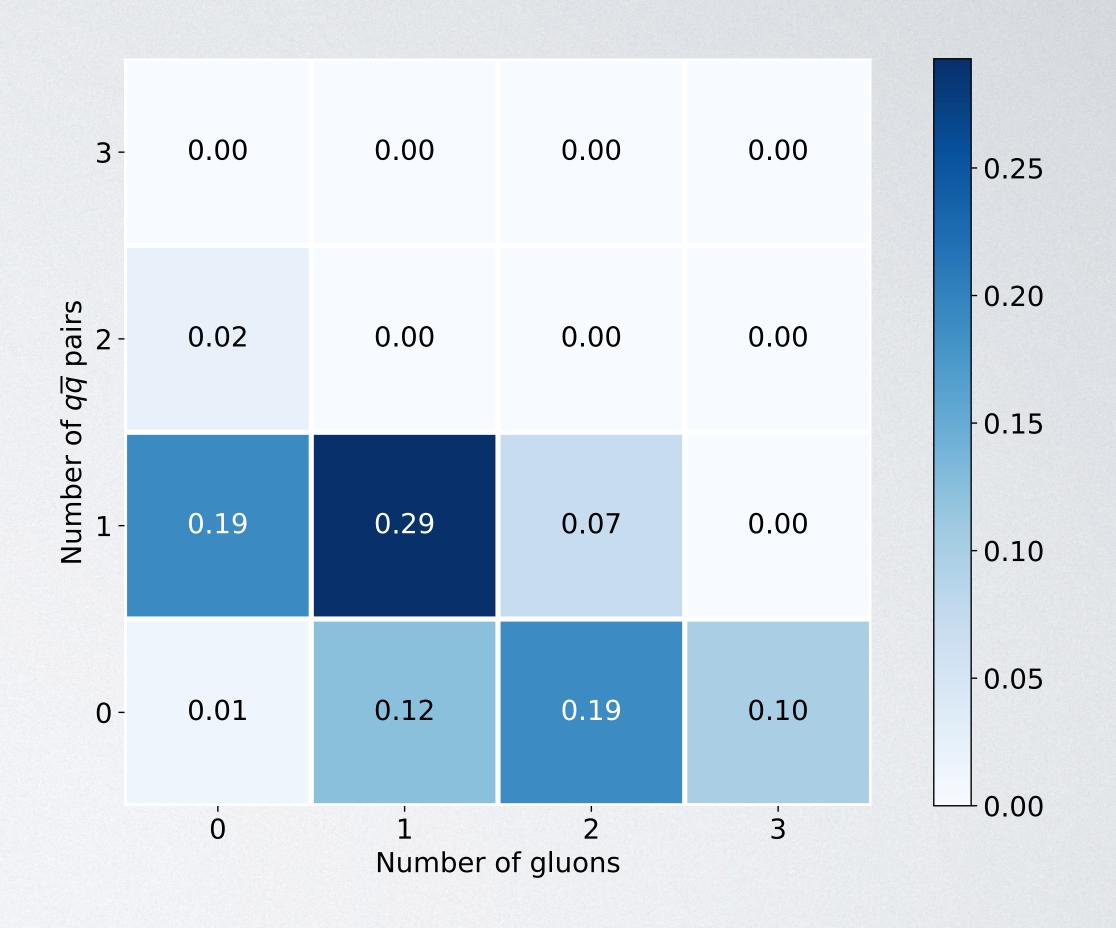
· The Sudakov factors have been used to determine whether an emission occurs:

$$\Delta_{i,k}(z_1, z_2) = \exp\left[-\alpha_s \int_{z_1}^{z_2} P_k(z') dz'\right], \qquad \Delta_{tot}(z_1, z_2) = \Delta_g^{n_g}(z_1, z_2) \Delta_q^{n_q}(z_1, z_2) \Delta_{\overline{q}}^{n_{\overline{q}}}(z_1, z_2).$$

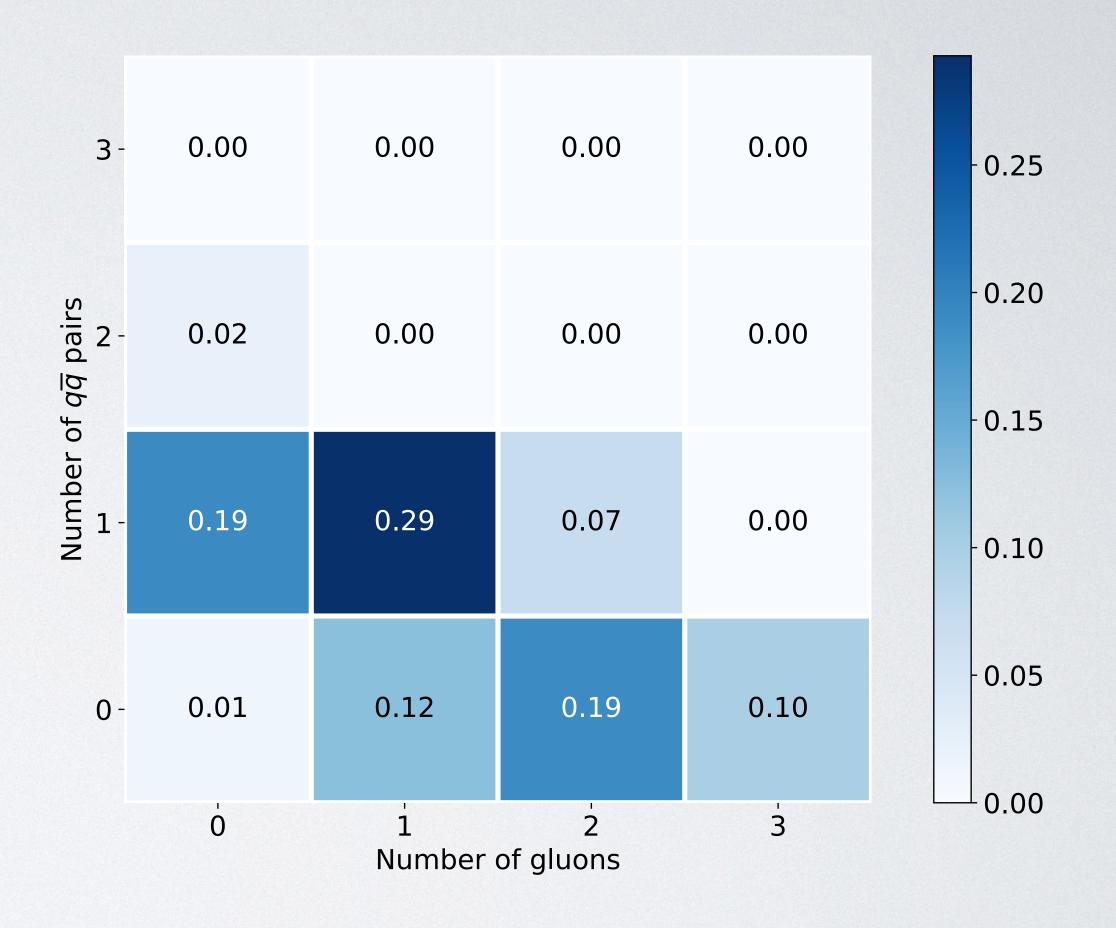
• Combine Sudakov and splitting functions to get splitting probability for $k \to ij$ in a single shower step:

$$Prob_{k\to ij} = (1 - \Delta_k) \times P_{k\to ij}(z)$$

• \mathcal{H}_P : increase dimension of position space to 2D to allow for the simulation of a gluons and quarks

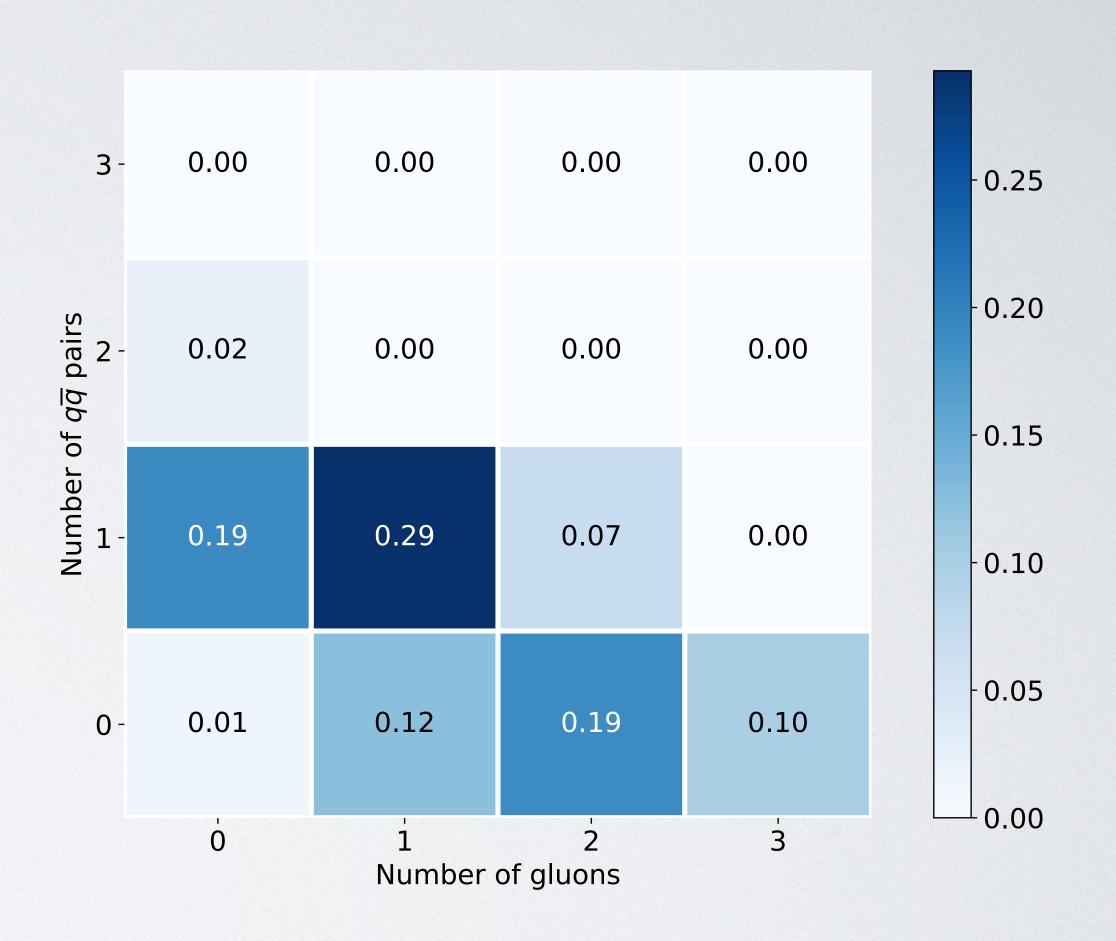


- \mathcal{H}_P : increase dimension of position space to 2D to allow for the simulation of a gluons and quarks
- $\mathcal{H}_{\mathcal{C}}$: increase dimension of coin space to accommodate for the collinear splitting probabilities



- \mathcal{H}_P : increase dimension of position space to 2D to allow for the simulation of a gluons and quarks
- $\mathcal{H}_{\mathcal{C}}$: increase dimension of coin space to accommodate for the collinear splitting probabilities
- C: coin operation is now splitting probability:

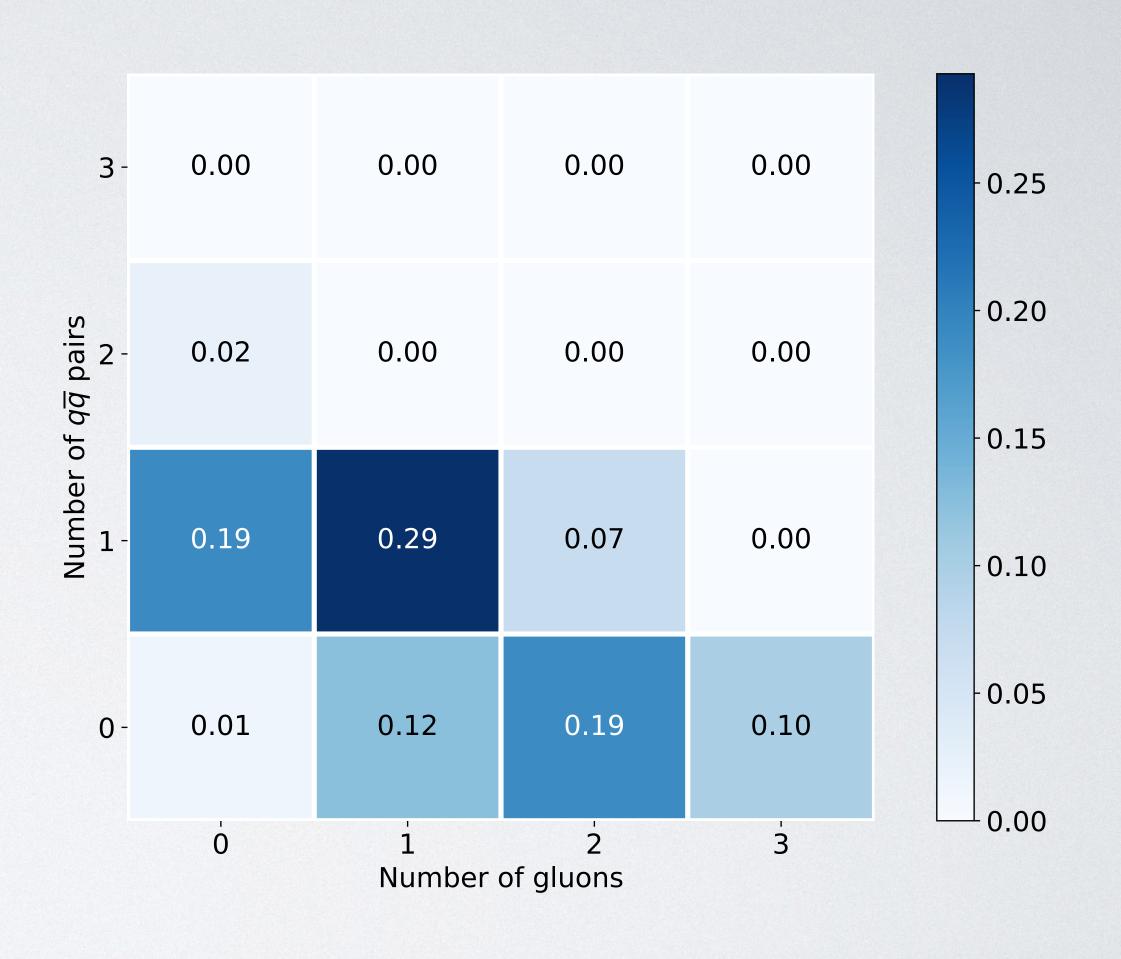
$$P_{ij} = (1 - \Delta_k) \times P_{k \to ij}$$



- \mathcal{H}_P : increase dimension of position space to 2D to allow for the simulation of a gluons and quarks
- $\mathcal{H}_{\mathcal{C}}$: increase dimension of coin space to accommodate for the collinear splitting probabilities
- C: coin operation is now splitting probability:

$$P_{ij} = (1 - \Delta_k) \times P_{k \to ij}$$

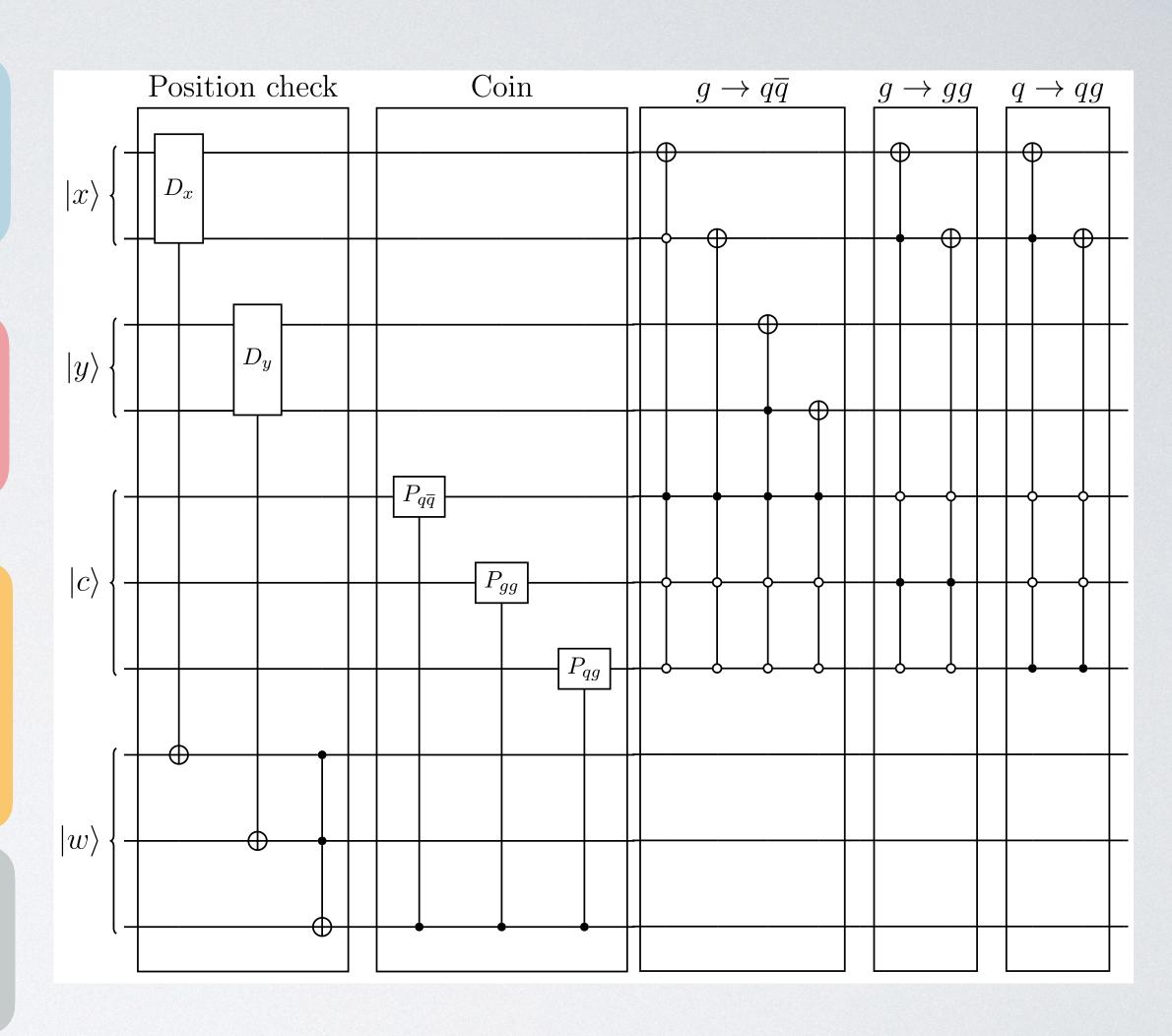
• S: shift operation updates shower content accordingly



- \mathcal{H}_P : increase dimension of position space to 2D to allow for the simulation of a gluons and quarks
- $\mathcal{H}_{\mathcal{C}}$: increase dimension of coin space to accommodate for the collinear splitting probabilities
- C: coin operation is now splitting probability:

$$P_{ij} = (1 - \Delta_k) \times P_{k \to ij}$$

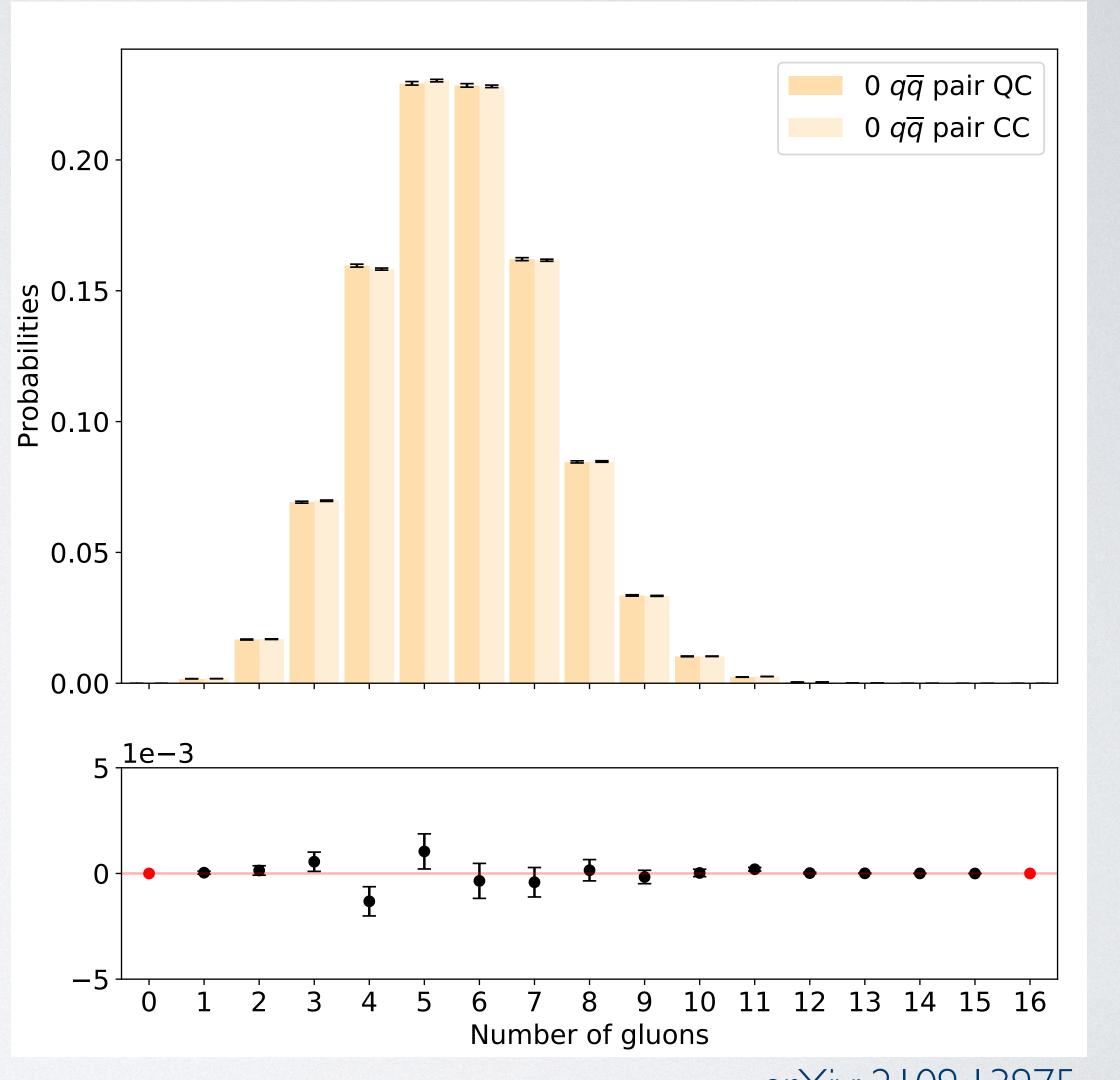
• S: shift operation updates shower content accordingly



| | Previous algorithm | QW |
|----------------|-----------------------|------------------|
| Qubits | 31 | 16 |
| Steps | 2 | 31 |
| Scaling, n_q | $\frac{3N(N+1)}{2}^*$ | $2\log_2(N+1)+6$ |



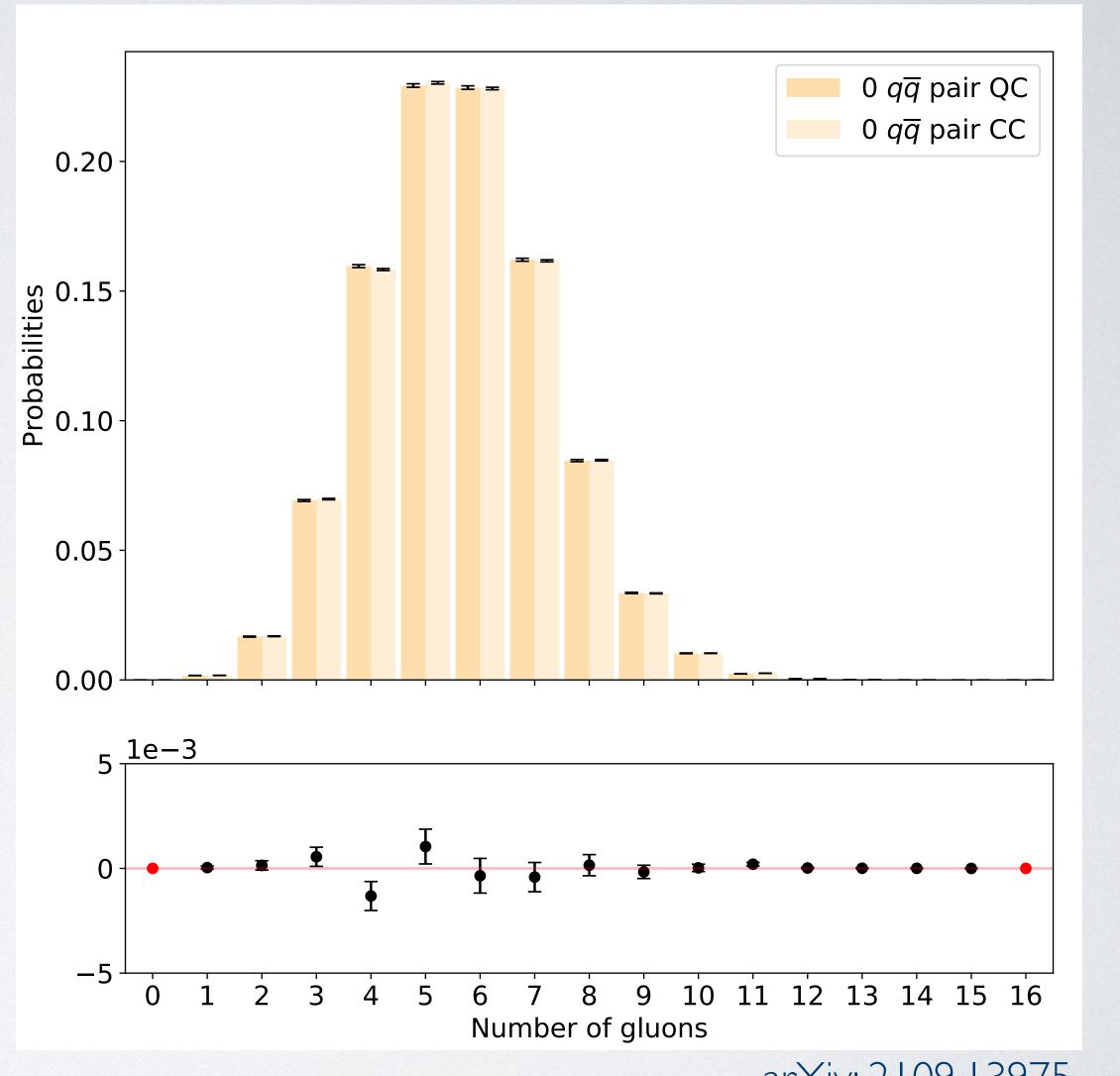
Previous - Phys. Rev. D 103, 076020 (2021)



| | Previous algorithm | QW |
|----------------|-----------------------|------------------|
| Qubits | 31 | 16 |
| Steps | 2 | 31 |
| Scaling, n_q | $\frac{3N(N+1)}{2}^*$ | $2\log_2(N+1)+6$ |



Previous - Phys. Rev. D 103, 076020 (2021)

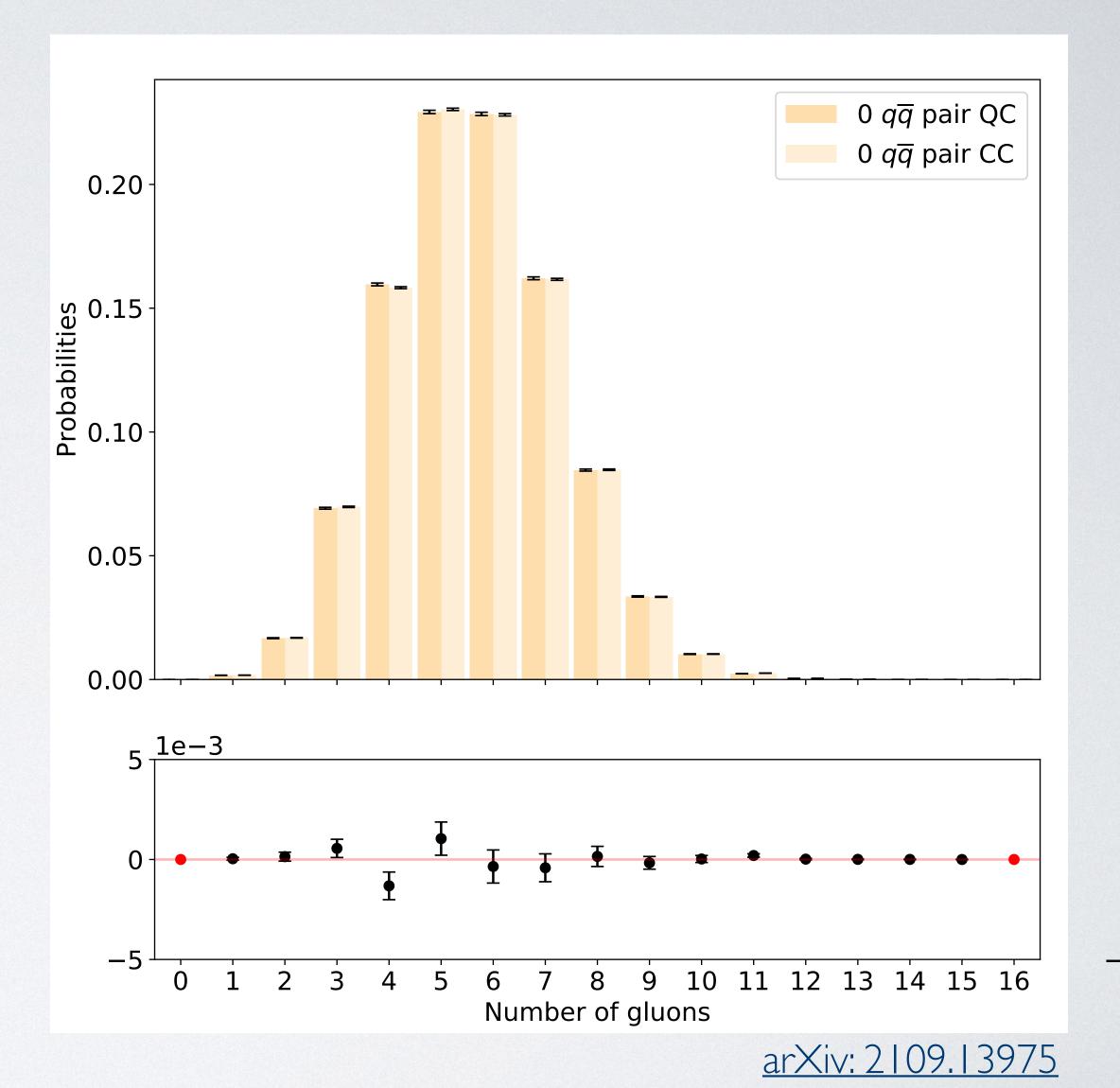


arXiv: 2109.13975

| | Previous algorithm | QW |
|----------------|-----------------------|------------------|
| Qubits | 31 | 16 |
| Steps | 2 | 31 |
| Scaling, n_q | $\frac{3N(N+1)}{2}^*$ | $2\log_2(N+1)+6$ |

*Scaling of a single register, not full circuit!

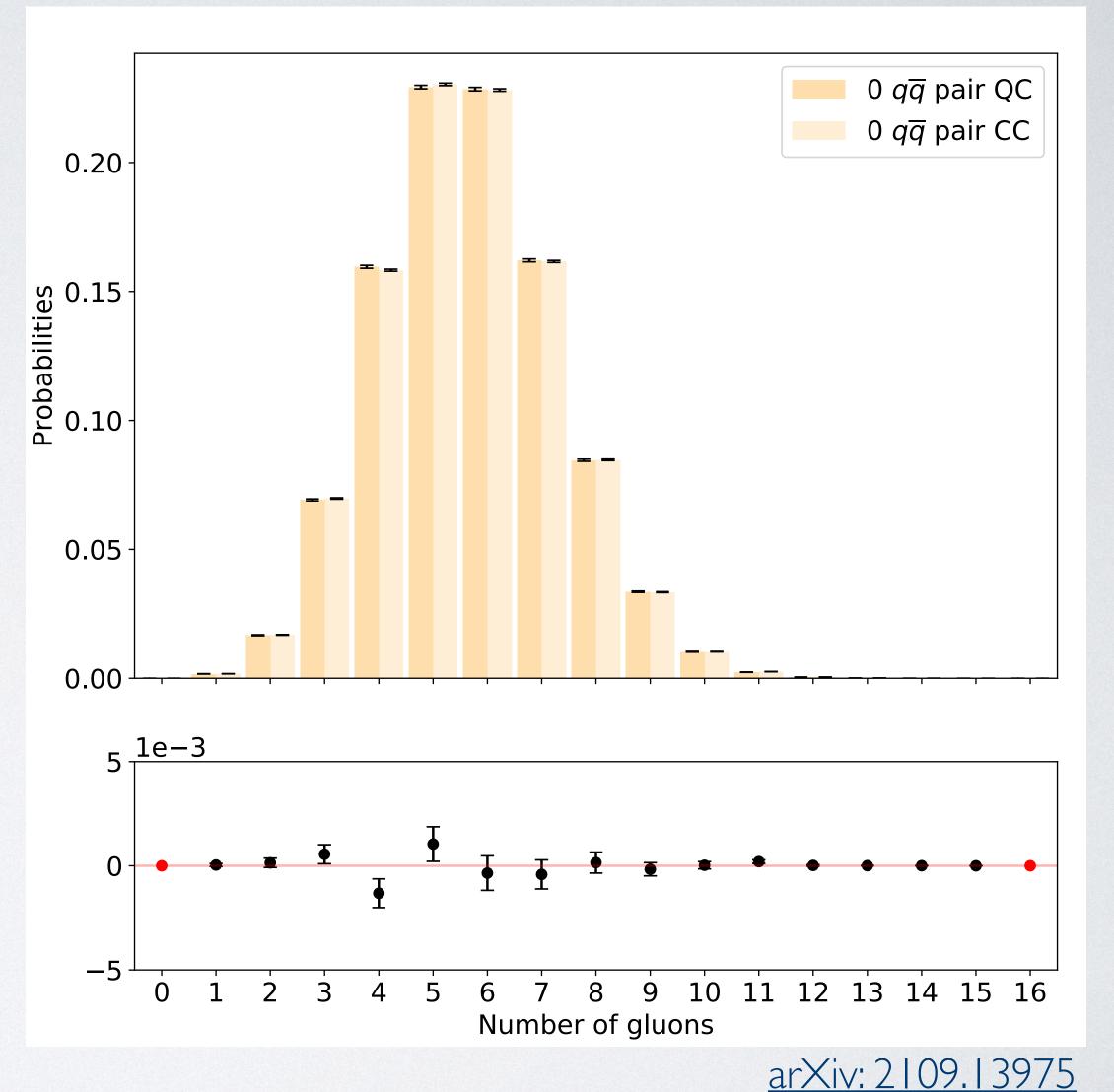
Previous - Phys. Rev. D 103, 076020 (2021)



| | Previous algorithm | QW |
|----------------|-----------------------|------------------|
| Qubits | 31 | 16 |
| Steps | 2 | 31 |
| Scaling, n_q | $\frac{3N(N+1)}{2}^*$ | $2\log_2(N+1)+6$ |

*Scaling of a single register, not full circuit!

Previous - Phys. Rev. D 103, 076020 (2021)



Summary and Looking to the Future

- Present a dedicated quantum algorithm for the simulation of parton showers in high energy collisions:
 - All shower histories calculated in full superposition constructing a final wavefunction containing all possible histories. Measurement projects out a physical quantity.
 - Reframing in the Quantum Walk framework vastly improves the efficiency of the quantum parton shower algorithm and offers a potential quadratic speed up compared to MCMC sampling
- Looking to the future: the introduction of kinematics to the algorithm will be a large step forward in the realism of the algorithm, with the potential of comparison to real data



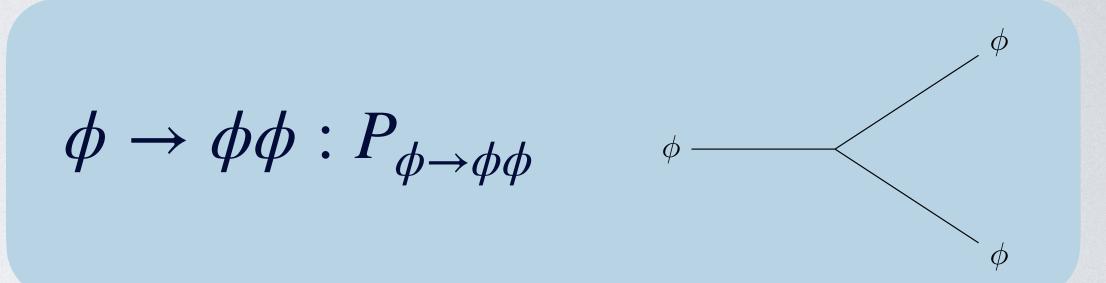


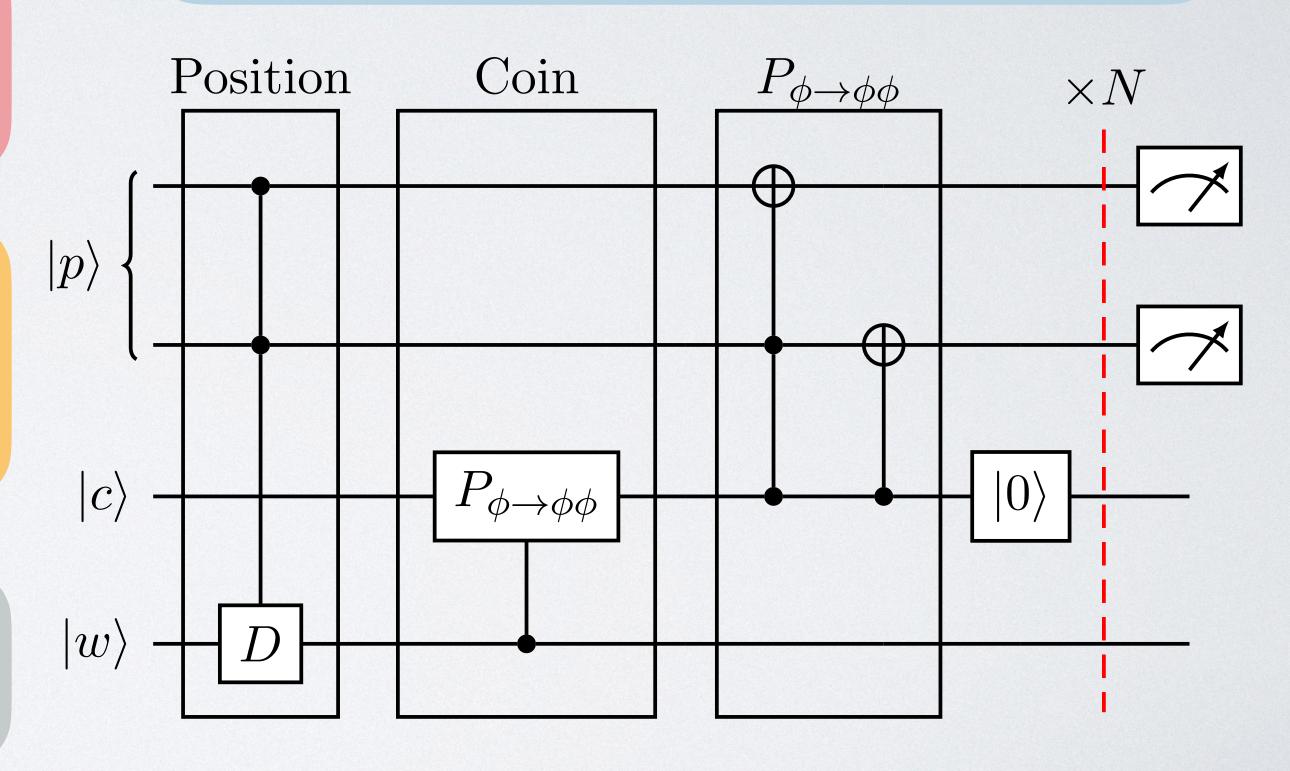


Quantum Walk approach to the parton shower - A Simple Shower

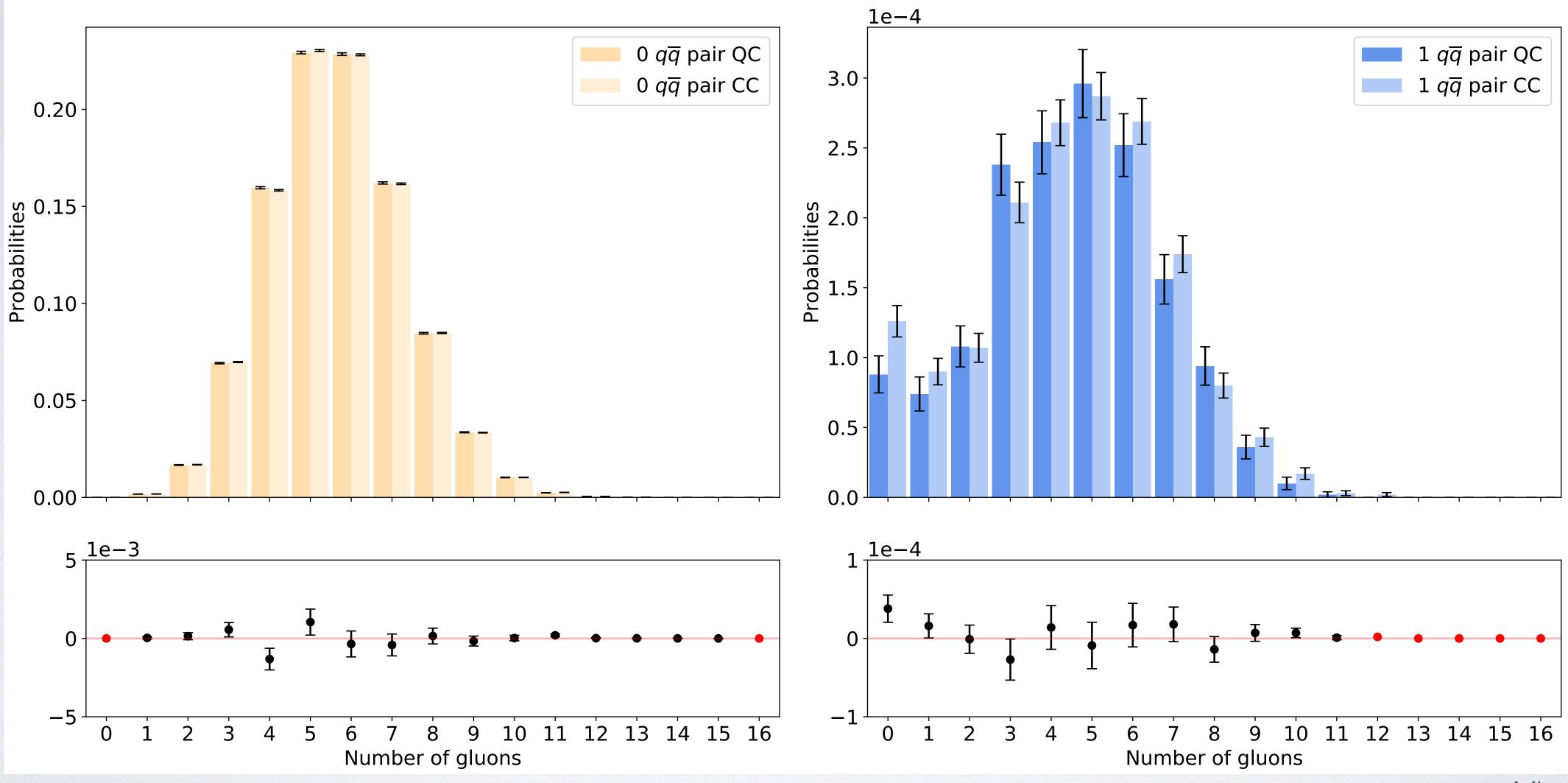
7

- Consider a simple shower with a single particle type ϕ
- \mathcal{H}_c : Here we alter the coin operation to reflect the splitting probability $P_{\phi \to \phi \phi}$
- \mathcal{H}_p : The walker position space now reflects the number of $\pmb{\phi}$ particles present in the shower
- The shift operation only increases the position of the walker, as only $\phi \to \phi \phi$ splittings





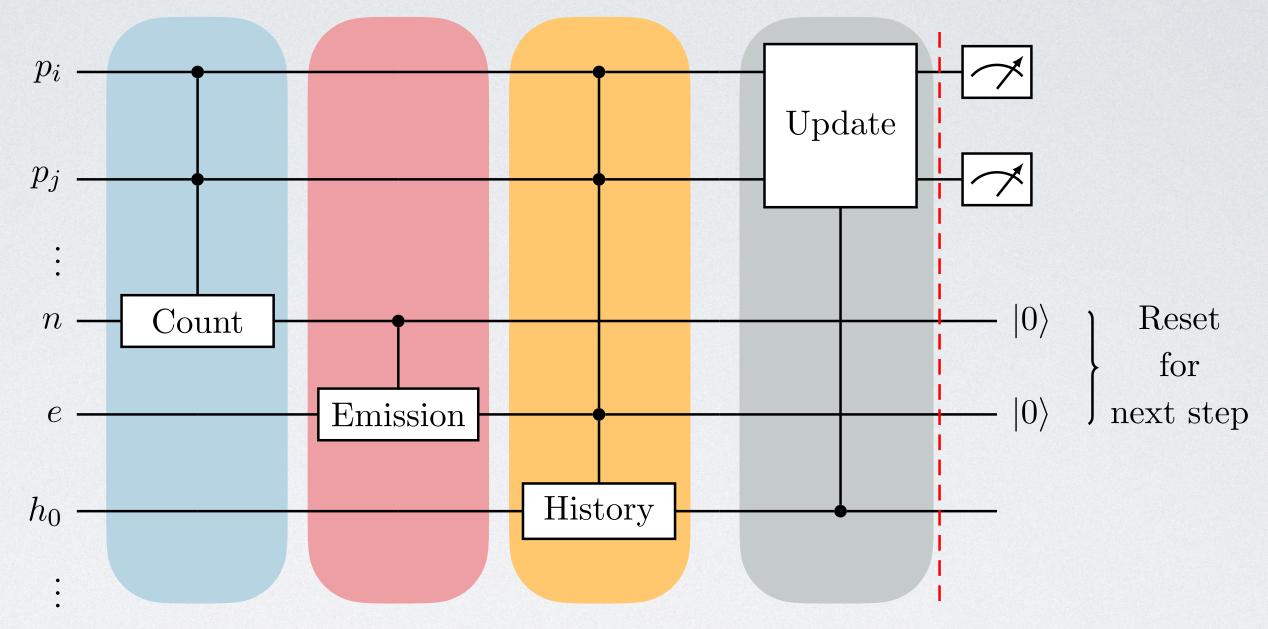
Quantum Walk approach to the parton shower - Results



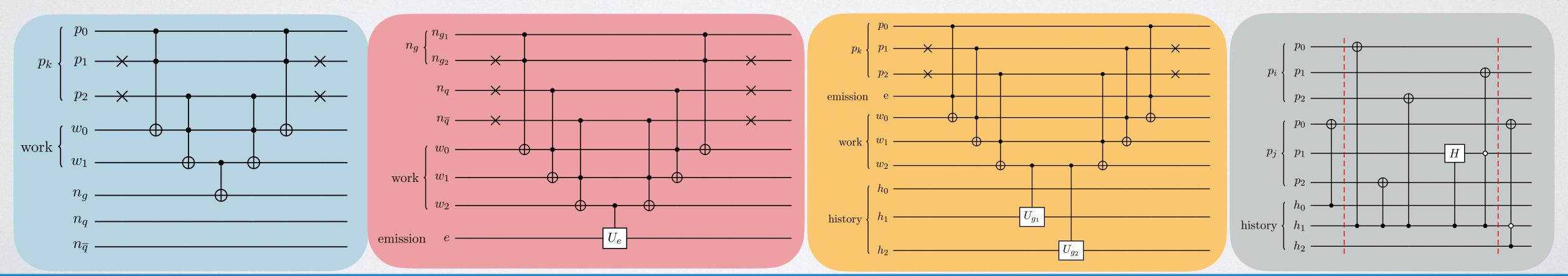
arXiv: 2109.13975

Markov Chain parton shower implementation

Previous algorithm:



Builds on Phys. Rev. Lett. 126, 062001 (2021)



Measurement

- Measurement of an arbitrary qubit system, $|\psi\rangle = \alpha |0\rangle + \beta |1\rangle$, is represented by the projection onto the $|0\rangle$ and $|1\rangle$ state, defining the projection operators $P_0 = |0\rangle\langle 0|$ and $P_1 = |1\rangle\langle 1|$.
- The probability of measuring the $|0\rangle$ state:

$$\operatorname{Prob}(|\hspace{.06cm}0\rangle) = \operatorname{Tr}(P_0|\hspace{.04cm}\psi\rangle\langle\psi\hspace{.04cm}|\hspace{.04cm}) = \langle\psi\hspace{.04cm}|\hspace{.04cm}P_0|\hspace{.04cm}\psi\rangle = |\hspace{.04cm}\alpha\hspace{.04cm}|^2$$

• Qubits are measured in this Projection-Valued Measurement regime and so the final state of the qubit is altered by the measurement. If the qubit is measured in the $|0\rangle$ state, then the final qubit state is:

$$|\psi\rangle \leftarrow \frac{P_0|\psi\rangle}{\sqrt{\langle\psi|P_0|\psi\rangle}} = |0\rangle$$

Looking to the Future of Quantum Computers

- We are on the brink of a 'quantum revolution' - IBM on track to exceed 1000 qubits by 2023
- Quantum Walks have long been conjectured to give a quadratic speed up in the mixing time of Markov Chains
- Quadratic speed up has been proven for several quantum
 MCMC algorithms

