University of Sheffield

Particle Astrophysics

SUSAN CARTWRIGHT UNIVERSITY OF SHEFFIELD

2

High energy particle astrophysics: the data

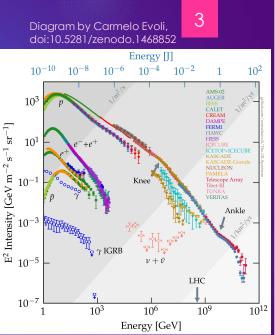
COSMIC RAYS PHOTONS NEUTRINOS EXPECTED RELATIONSHIPS BETWEEN OBSERVABLES

Cosmic rays

Cosmic rays were discovered about 1913 by Victor Hess.

- They consist mostly of protons and heavier ions, and have a power law spectrum with approximate spectral index 2.7.
 - There are two conspicuous slope breaks, the "knee" above 10⁶ GeV and the "ankle" above 10⁹ GeV, and an apparent cut-off around 10¹¹ GeV.

Cosmic rays do not point back to their origin because of deflection by the Galactic magnetic field.



Composition, rigidity and sources

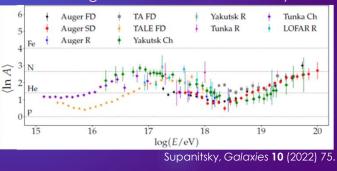
4

For cosmic rays to be accelerated, they must be confined within the accelerating region.

- This can't be done by gravity—it must involve magnetic fields.
- The response of a charged particle to a magnetic field is determined by its rigidity cp/q.

 6
 + Auger FD
 + TA FD
 + Yakutsk R
 + Tu
- Therefore the maximum energy attainable in a given source is higher for heavy ions than for protons.

Increasing mean mass is a signature for a source type cutting off.



5

6

Cosmic ray origins

To identify the origins of cosmic rays we need a neutral messenger.

The possibilities are:

- Photons, if produced by a non-thermal mechanism that requires the presence of high-energy particles
 - \blacktriangleright examples: synchrotron radiation, inverse Compton scattering, π^0 decay.
- High-energy neutrinos
 - produced by π[±] decay (much higher energy than solar or supernova neutrinos).

Photons may only signal the presence of high-energy electrons; neutrinos definitely require high-energy hadrons.

Photons: synchrotron radiation

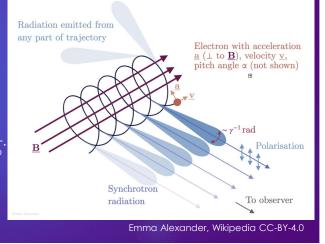
Synchrotron radiation is produced by relativistic particles gyrating in a magnetic field.

 Averaging over the pitch angle α, power emitted is

$$P_{\rm rad} = \frac{4}{3} c \sigma_{\rm T} U_{\rm mag} \beta^2 \gamma^2$$

where $\sigma_{\rm T} = \frac{e^4}{6\pi \epsilon_0^2 c^4 m_e^2}$ and $U_{\rm mag} = \frac{B^2}{2\mu_0}$

• The typical photon energy is $\frac{3}{2}\gamma^2 hv_g \sin \alpha$ where $v_g = eB/(2\pi m)$.



Photons: inverse Compton scattering

A low-energy seed photon backscatters off a high-energy electron.

- Power radiated is $P_{rad} = \frac{4}{3}c\sigma_T U_{rad}\beta^2\gamma^2$ where $U_{rad} = S/c$ is energy stored in radiation field
- ► Typical photon energy is $\frac{4}{2}\gamma^2 hv_0$ where v_0 is the seed photon frequency.

The dependence on the electron energy is the same for both synchrotron and inverse Compton, but $v_0 \gg v_a$.

- Therefore we expect the same shape of spectrum, but at higher energies.
- > The relative normalisation depends on the magnetic field.

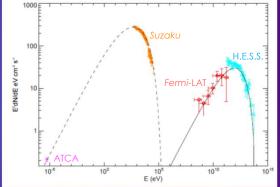
Example: SNR RX J1713.7–3946

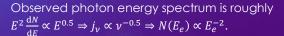
8

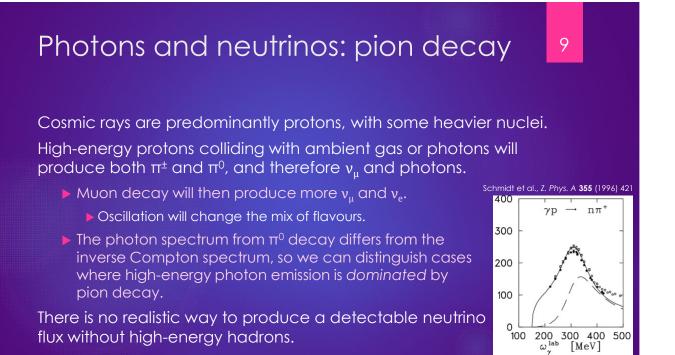
Fit to data using synchrotron radiation (dashed line) plus inverse Compton (solid line).

- Note that for a power-law electron spectrum $N(E_e) \propto E_e^{-\delta}$ we expect a synchrotron radiation spectrum $j_v \propto B^{(\delta+1)/2} v^{-(\delta-1)/2}$.
- Because $P_{rad} \propto \gamma^2$ the electron power law will cut off at high energies owing to rapid energy loss in the high-energy tail, so expect similar cut-off in the photon spectrum, as seen.

Chuyuan Yang and Siming Liu ApJ 773 (2013) 138



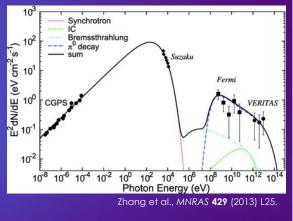




Example: Tycho's supernova

The remnant of the supernova observed by Tycho Brahe in 1572.

- Known to have been Type Ia from reflected spectrum observed in 2008.
- Shape of high-energy spectrum quite different from the synchrotron spectrum, so not inverse Compton.
- **)** Can be well fitted by π^0 decay.



10

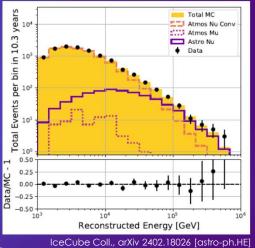
Neutrino detection

High-energy neutrinos are detected by observing neutrino interactions in large water Cherenkov detectors

currently the only one with a large enough active volume is IceCube, but it will be joined by KM3NeT/ARCA.

The principal problem is the background from atmospheric neutrinos

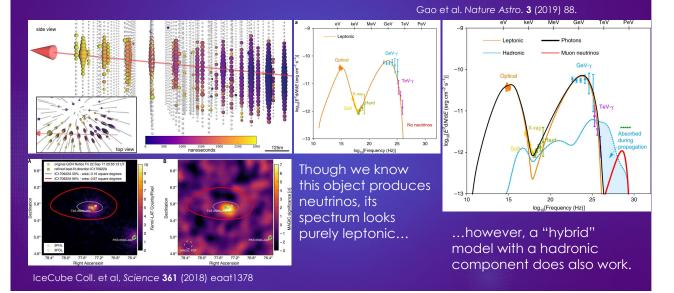
this means that only very high-energy neutrinos can be identified as astrophysical on an event-by-event basis.



11

Example: TXS 0506+056





13

Expected correlations

High-energy neutrinos should be accompanied by high-energy photons, as π^{\pm} production should imply π^{0} production

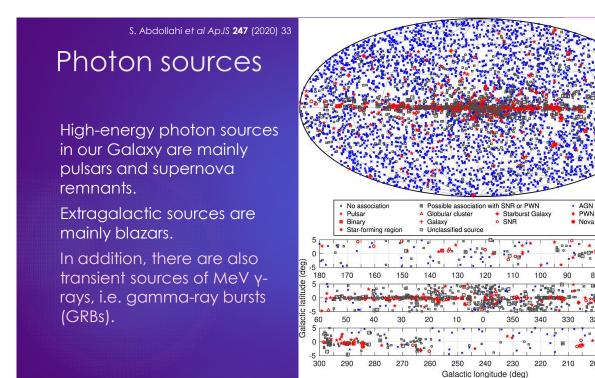
but if the source region (or its surroundings) is very dense, the photons might not escape.

High energy photons from π^0 decays should be accompanied by neutrinos, for the same reason

but the neutrino signal may be too weak to see, or buried in the atmospheric neutrino background

Most acceleration mechanisms should accelerate all charged particles, so high-energy electrons should imply high-energy hadrons

▶ but they might not, if the seed material is mostly e⁺e⁻ pair plasma (which it may be in some environments).



300

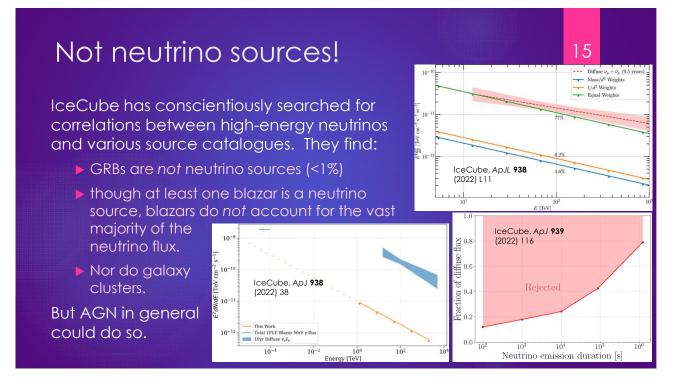
180

320

200

310 6°. .

190



Summary

Cosmic rays, consisting primarily of protons and heavier ions, are observed to have a power-law energy epctrun extending out to $\sim 10^{20}$ eV.

- ▶ This implies the existence of extremely powerful astrophysical accelerators.
- Cosmic rays do not pinpoint these directly because their trajectories are deflected by the Galaxy's magnetic field.

Acceleration of hadrons should be accompanied by γ -ray and neutrino emission from pion decay, and of electrons by photons from synchrotron radiation (radio to X-rays) and inverse Compton (γ -rays).

- These are all seen, but a complete explanation is still lacking
- ▶ in particular, the origins of astrophysical neutrinos are not yet established.

16