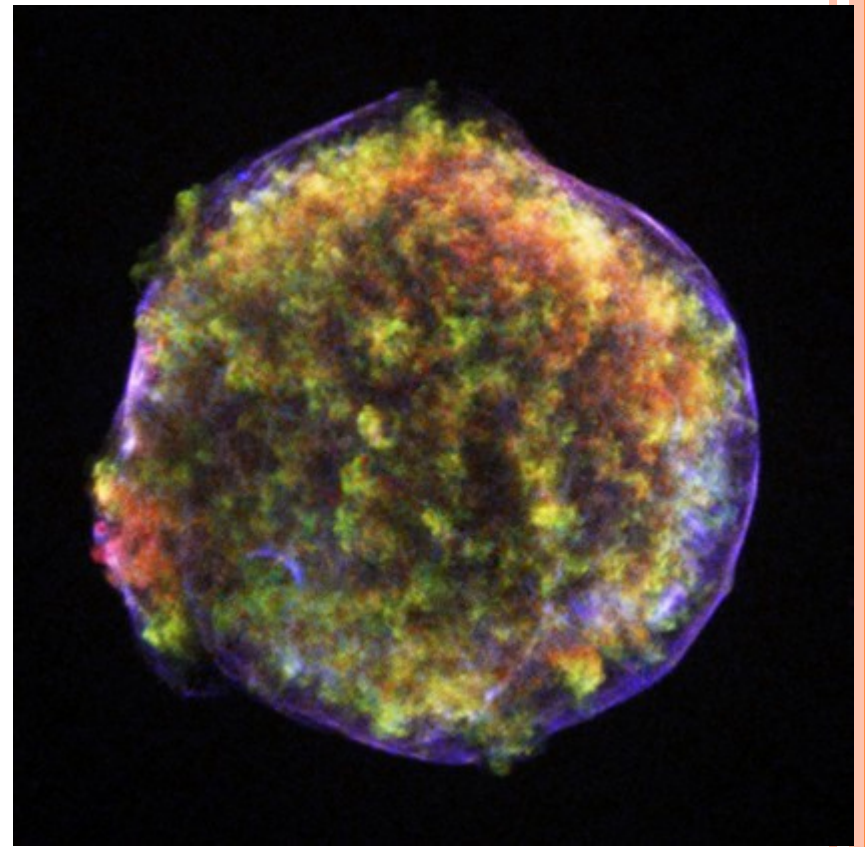


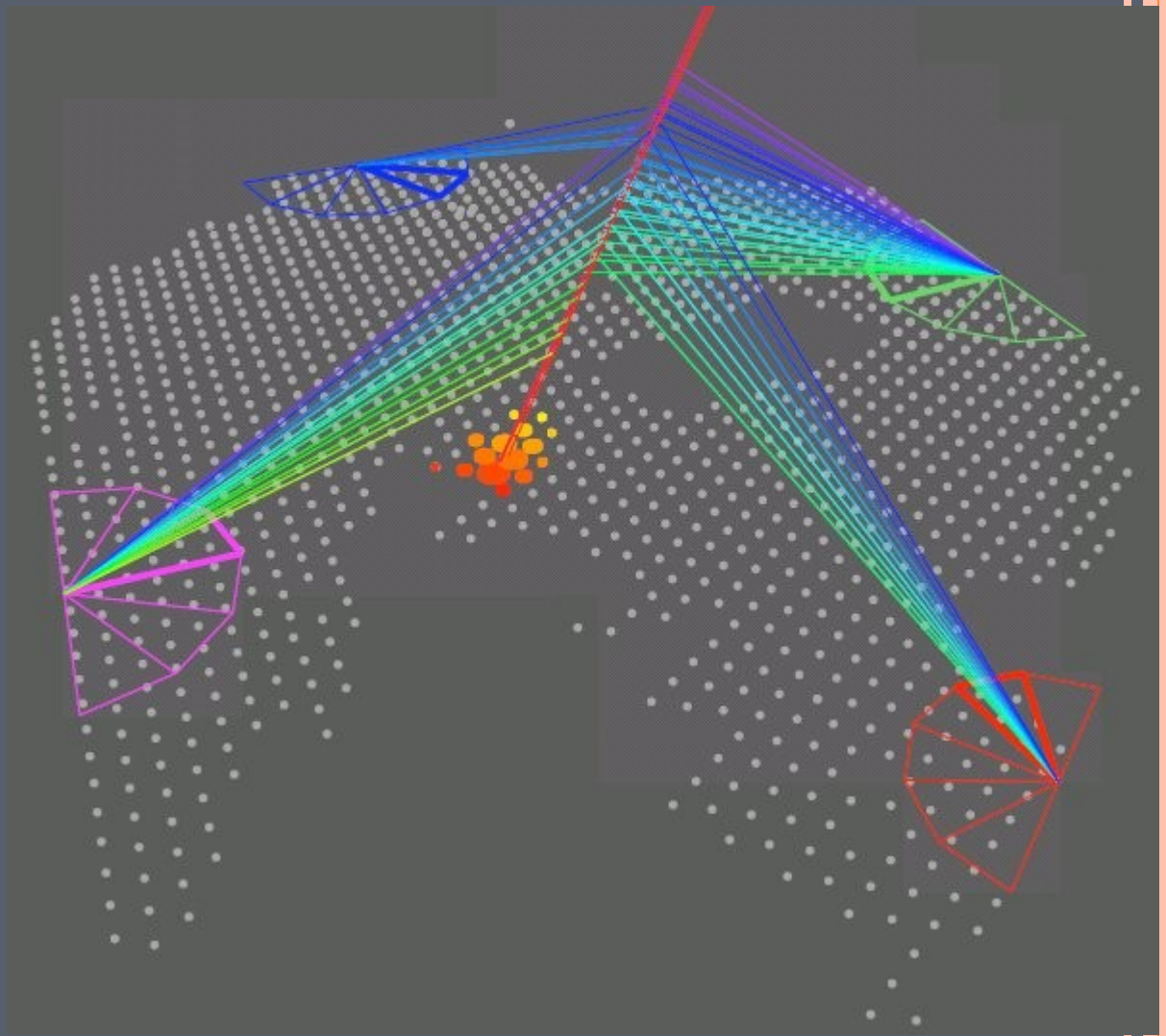


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Sheffield.



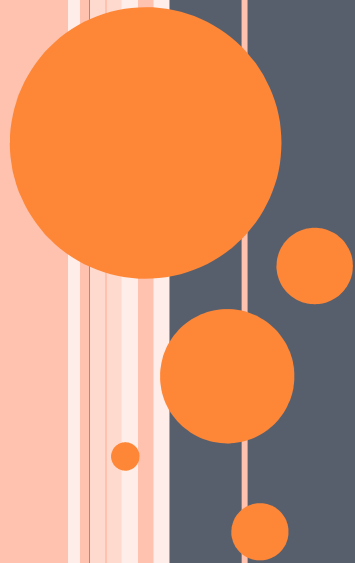
# ASTROPARTICLE PHYSICS LECTURE 1

**Matthew Malek**  
University of Sheffield



# Overview

What is Astroparticle Physics?



# What is Astroparticle Physics?

Various definitions! For this week, we use:

- **The use of particle physics technology to study astrophysical phenomena**

◦ Included:

- neutrino astrophysics
- gamma-ray astronomy
- cosmic rays

coherent field  
with a lot of  
common  
factors

***High Energy  
Astroparticle  
Physics***

- dark matter
- early-universe cosmology

someone  
else's  
problem!

◦ Sometimes also included:

- cosmic microwave background
- gravitational waves
- neutrino masses (especially  $0\nu\beta\beta$ )

not very  
particulate

not very  
astrophysical

# Common Issues

- Low rates
  - fluxes of high-energy particles are small
  - neutrinos and dark matter have weak interactions
- *Need for large detectors*
- No control over “beam”
  - harder to control backgrounds
  - harder to calibrate, e.g., energy resolution
- *Signals can be difficult to establish and/or characterise*
  - *cf. solar and atmospheric neutrino oscillation*



# Related Fields

- Neutrino physics
  - Atmospheric neutrinos are technically “astroparticle physics” but have contributed more to understanding of neutrinos than to astrophysics
  - Somewhat similar situation for solar neutrinos
  - Long-baseline neutrino experiments can do low-energy neutrino astrophysics “for free”
- Nucleon decay
  - many detector technologies useful for both original purpose of Kamiokande (NDE = **N**ucleon **D**ecay **E**xperiment before **N**eutrino **D**etection **E**xperiment!)
  - Planned noble-liquid detectors may be able to do both nucleon decay experiments and dark matter searches



# Topics to be Covered

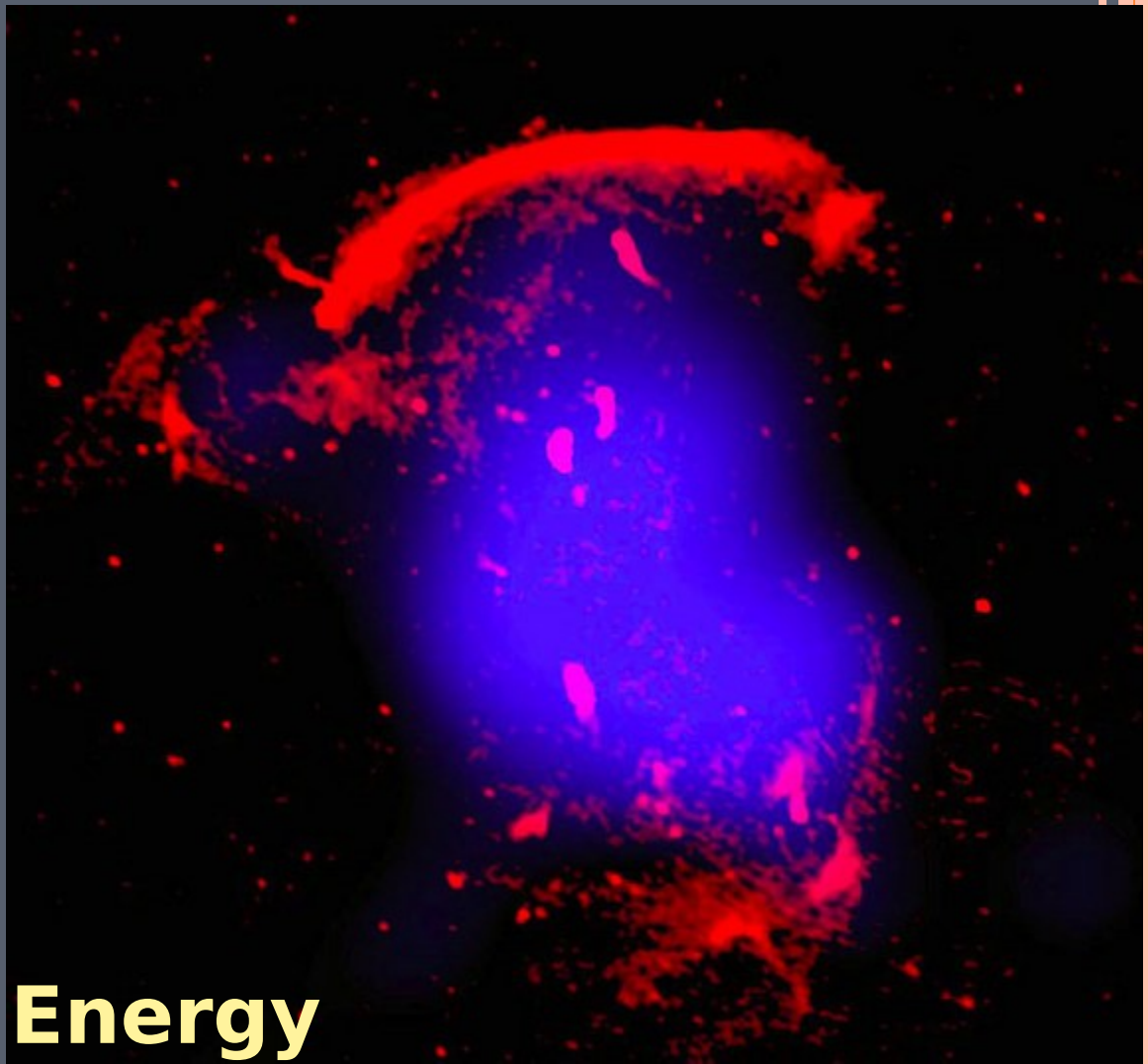
## High energy astroparticle physics

- (cosmic rays, gammas, high-energy neutrinos)
  - sources
  - detection
  - results
  - prospects
- **Dark matter**
  - evidence
  - candidates
  - search techniques

### Not Covering:

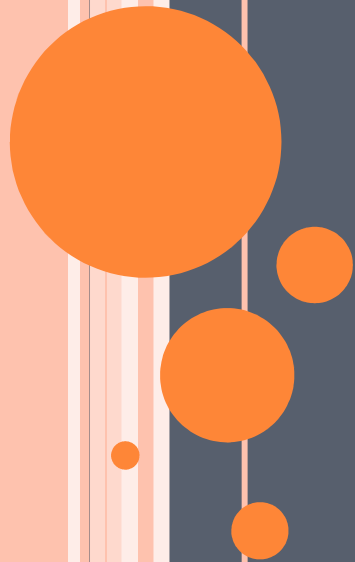
- solar neutrinos (SB)
- neutrino masses (SB)
- supernova neutrinos (no time)





# High Energy Astroparticle Physics

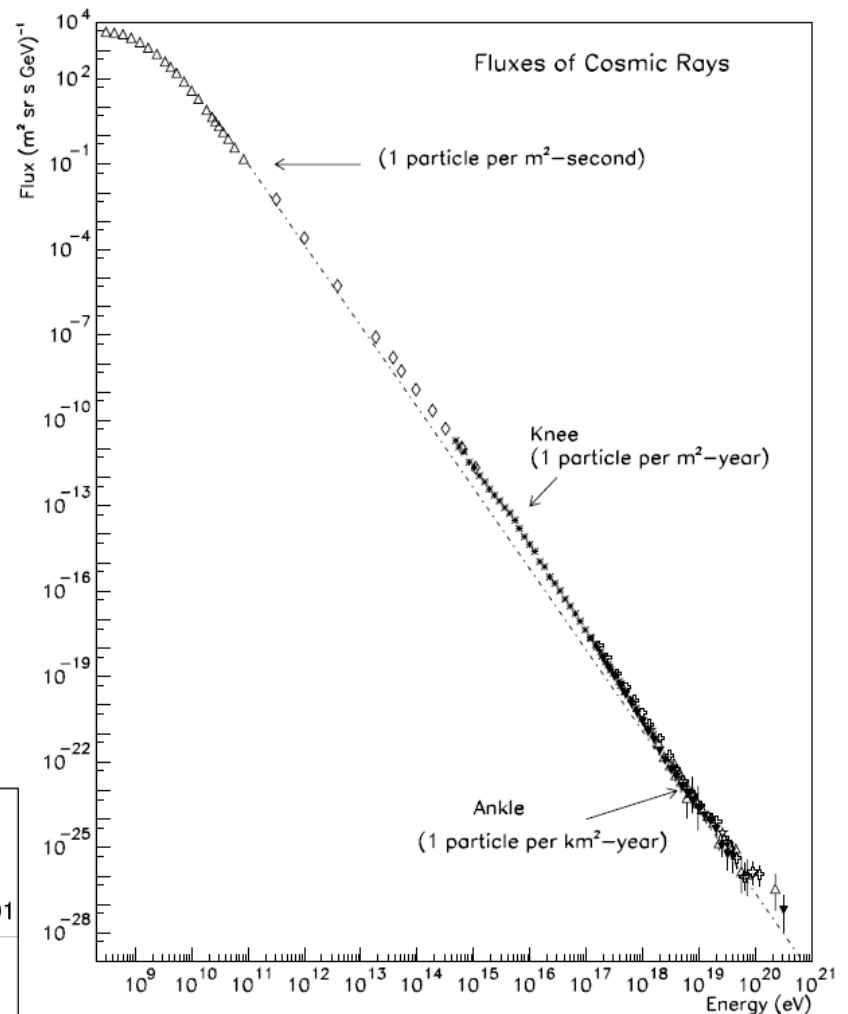
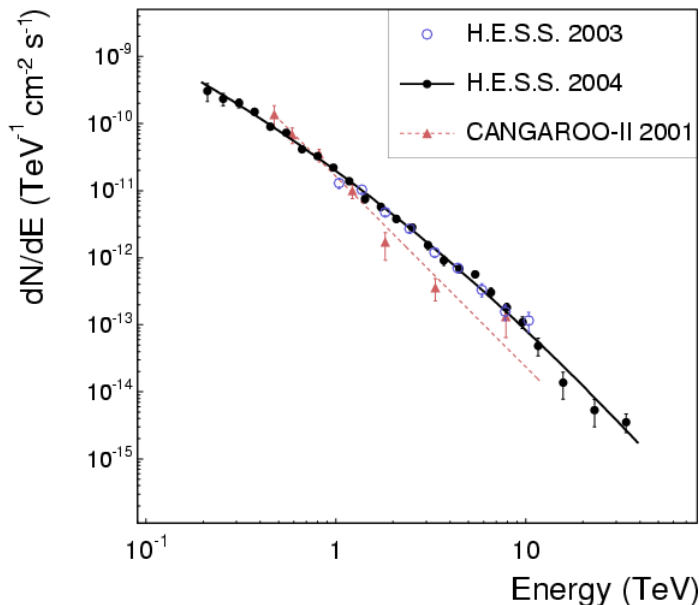
Acceleration Mechanisms  
Sources  
Detection



# Cosmic Accelerators

Cosmic rays and gamma rays are observed up to extremely

- high energies
- something must therefore
- accelerate them



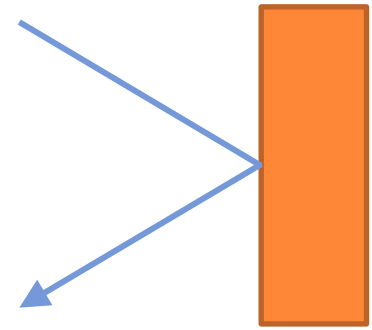
$10^9 \text{ eV} - 10^{21} \text{ eV}$

Note the power-law spectrum





# Acceleration Mechanisms



- Fermi Mechanism

- energetic charged particles can gain energy by scattering off local magnetic turbulence (Fermi 1949)
- Assume particle scatters off much more massive object moving with speed  $\mathbf{u}$ . Then in the c.o.m. frame (i.e., frame of massive object) its energy and momentum before the scatter are

$$E_{\parallel} = \gamma_u (E + up \cos \theta)$$

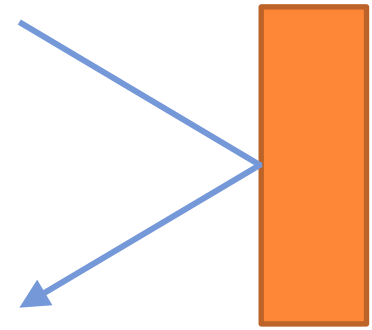
$$p_{\parallel} = \gamma_u (p \cos \theta + uE/c^2)$$

- The particle scatters elastically: its energy is conserved and its x-momentum reversed. In original (lab) frame

$$E_2 = \gamma_u (E_{\parallel} + up_{\parallel}) = \gamma_u^2 E_{\parallel} \left[ 1 + \frac{2uv}{c^2} \cos \theta + \frac{u^2}{c^2} \right]$$



# Acceleration Mechanisms



- Fermi Mechanism

- energetic charged particles can gain energy by scattering off local magnetic turbulence (Fermi 1949)
- We need to average over angle. Head-on collisions are slightly more likely than overtaking collisions, so middle term doesn't just go away. In relativistic limit we find:

$$\left\langle \frac{\Delta E}{E} \right\rangle = \frac{8}{3} \left( \frac{u}{c} \right)^2$$

- Hence this process is known as **second-order Fermi acceleration**.
- The good news
- this produces a power law energy spectrum:  $N(E) \propto E^{-x}$  where
  - $x = 1 + 1/\alpha\tau$ ,  $\alpha$  is the rate of energy increase and  $\tau$  is the residence time of the particle
- The bad news
  - since  $u \ll c$ , it's slow and inefficient

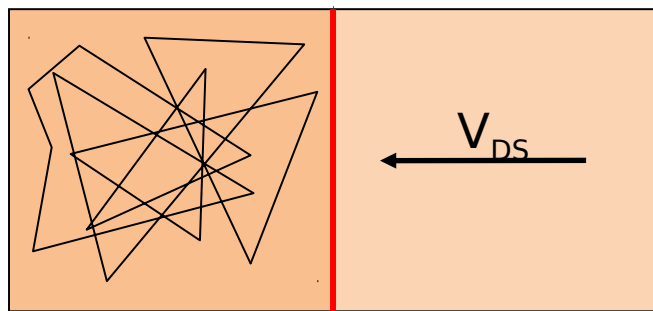
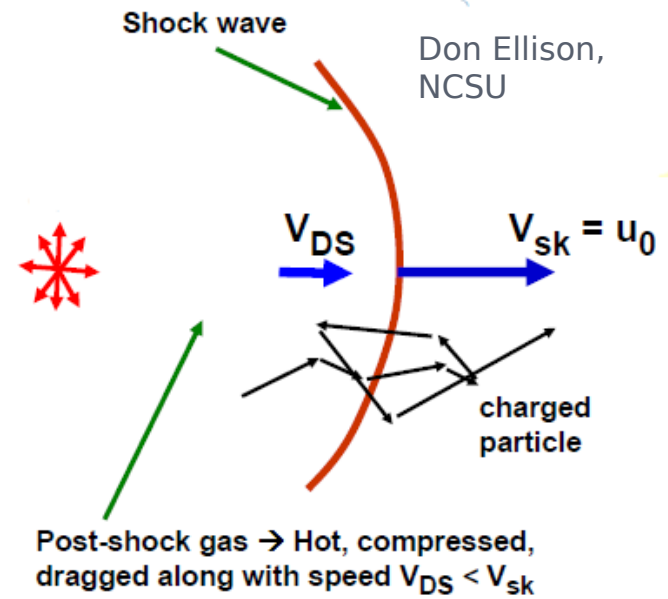


# Acceleration Mechanisms

## First-order Fermi Mechanism

### ◦ (Diffusive Shock Acceleration)

- $O(u/c)$  term gets lost in integral over angles—we could retrieve this if we
- angles—we could retrieve this if we
  - could arrange to have only head-on scatters
  - Consider shock wave as sketched above
  - high-energy particles will scatter so that their distribution is isotropic in the rest frame of the gas

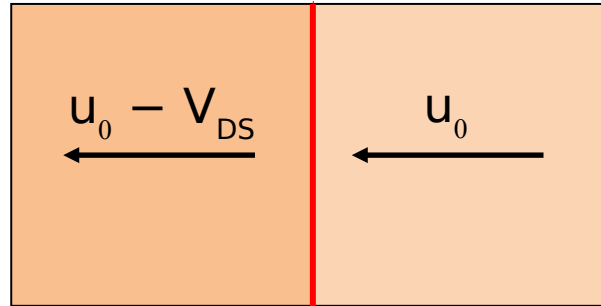


Rest frame of downstream gas

crossing shock **in either direction** produces head-on collision on average



# Acceleration Mechanisms



Rest frame  
of shock

- DSA, continued
- shock compresses gas, so density behind shock is  $\rho^2 > \rho^1$

in rest frame of shock,  $\rho^1 u_0 = \rho^2 u_2$  where  $u_2 = u_0 - V_{DS}$

for strong shock  $\rho^2/\rho^1 = (\gamma + 1)/(\gamma - 1)$  where  $\gamma$  is ratio of specific heats ( $= 5/3$  for hydrogen plasma); therefore expect  $u_2/u_0 \approx 1/4$

And gas approaches shock-crossing particle at speed  $V = 3/4 u_0$

- if high-energy particles move randomly, probability of particle crossing shock at angle  $\theta$  is  $P(\theta) = 2 \sin \theta \cos \theta d\theta$ , and its energy after crossing shock is  $E' \approx E(1 + pV \cos \theta)$  (if  $V \ll c$ )
- therefore average energy gain per crossing is

$$\left\langle \frac{\Delta E}{E} \right\rangle = \frac{V}{c} \int_0^{\pi/2} 2 \cos^2 \theta \sin \theta d\theta = \frac{2V}{3c}$$



# ACCELERATION MECHANISMS

## ○ DSA spectrum

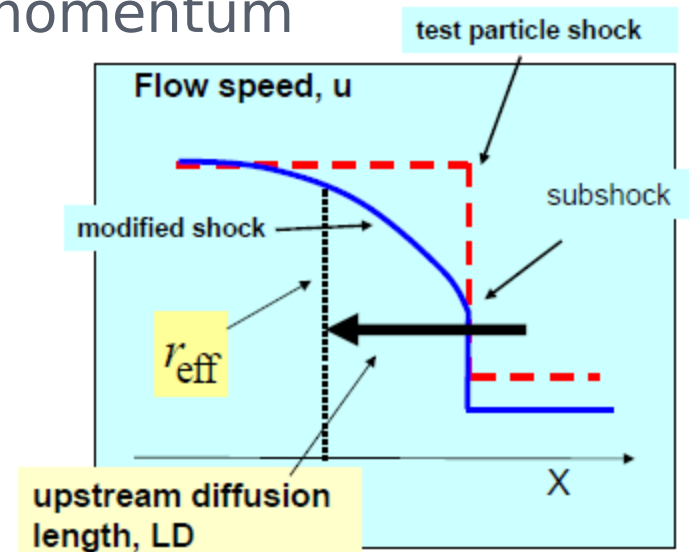
- if average energy of particle after one collision is  $E_1 = fE_0$ , and if  $P$  is probability that particle remains in acceleration region, then after  $k$  collisions there are  $N_k = N_0 P^k$  particles with average energy  $E_k = f^k E_0$ .

- Hence  $\frac{\ln(N/N_0)}{\ln(E/E_0)} = \frac{\ln P}{\ln f}$ , or  $\frac{N}{N_0} = \left(\frac{E}{E_0}\right)^{\ln P / \ln f}$

- This is the number of particles with  $E \geq E_k$  (since some of these particles will go on to further collisions), so differential spectrum is  $N(E) dE \propto E^{(\ln P / \ln f) - 1} dE$
- for DSA this comes to  $N(E) dE \propto E^{-(r+2)/(r-1)} dE$ , where  $r = \rho_2 / \rho_1$ .
  - “universal” power law, independent of details of shock

# Additional Complications

- Above was a “test particle” approach, in which we assume most of the gas is unaffected
  - If acceleration is efficient, high momentum particles will modify the shock
  - Need a consistent treatment which takes proper account of this
    - mathematically challenging
    - but valid across very large range of particle energies
    - Also need to allow for
  - possibility of relativistic shocks

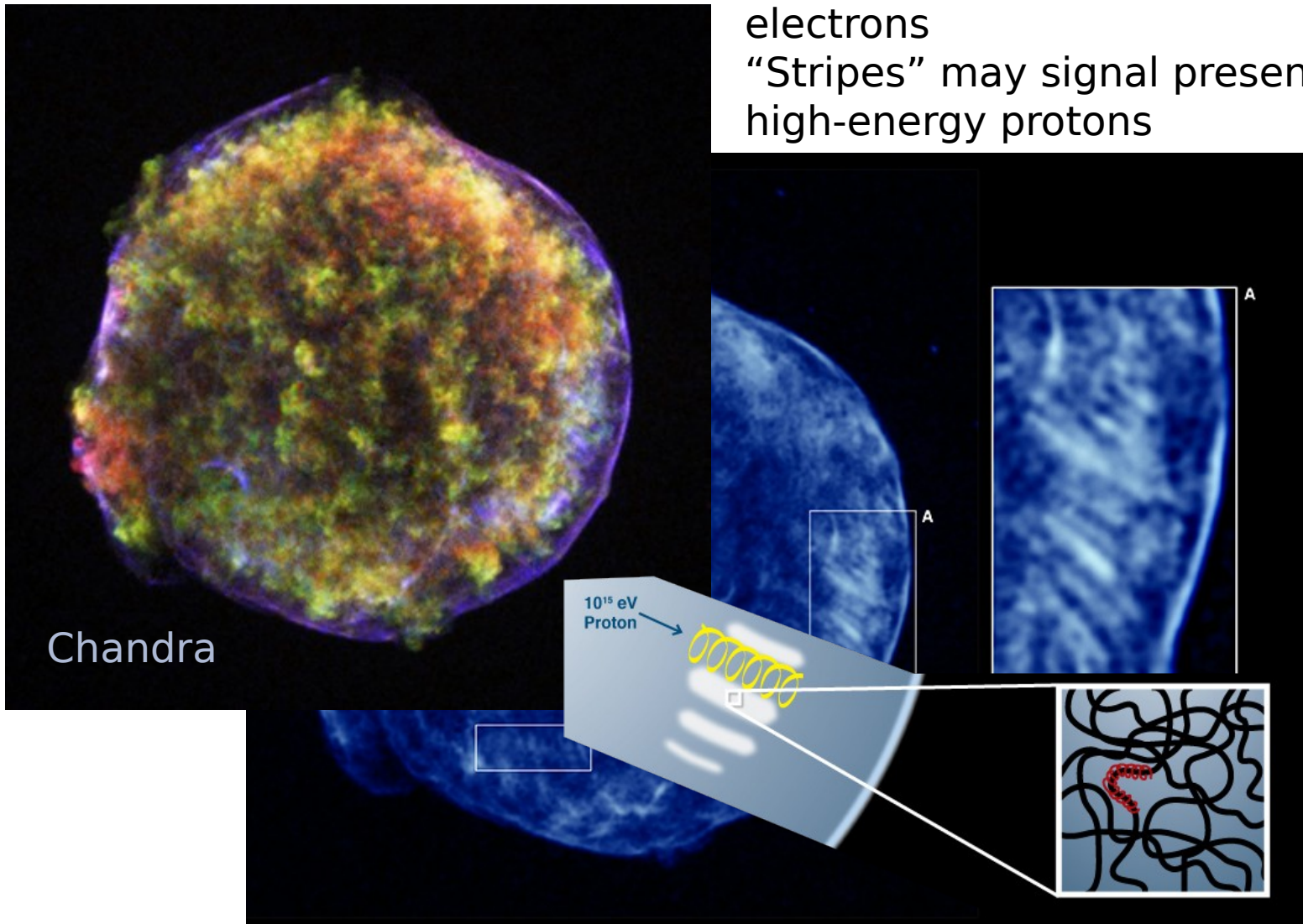


Don Ellison,  
NCSU

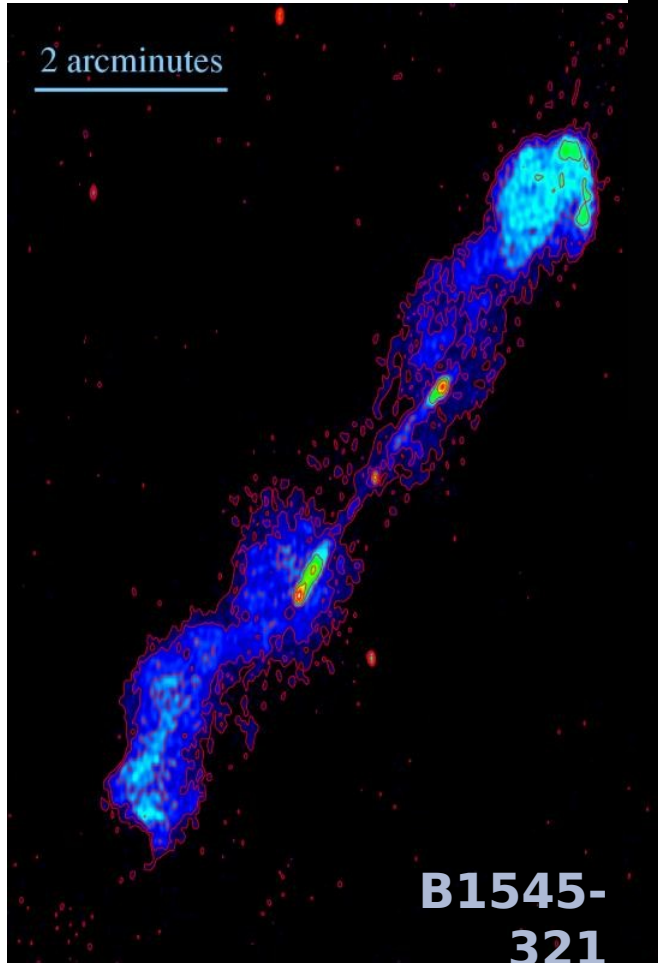


# Tycho's Supernova (SN 1572)

Shock front seen in high-energy electrons  
"Stripes" may signal presence of high-energy protons



# Radio Galaxies

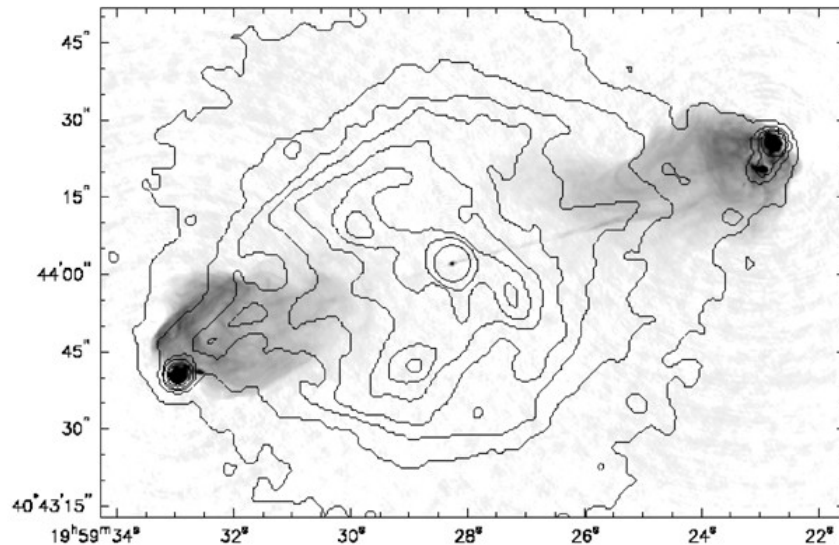


13 cm wavelength ATCA image by L. Saripalli, R. Subrahmanyan and Udaya Shankar

3C 273 jet



Chandra, HST, Spitzer



Cygnus A in X-ray (Chandra) and radio (VLA)





# Acceleration Mechanisms

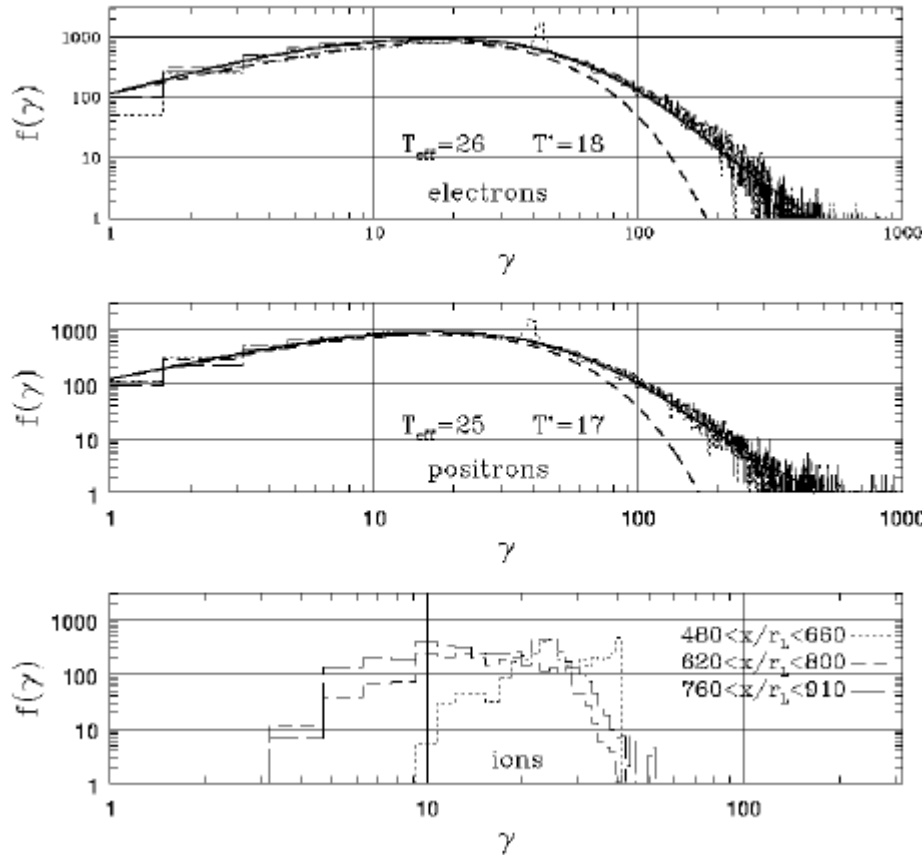
## ◦ Resonant Cyclotron Scattering & Absorption (RSA)

acceleration of  $e^+e^-$  in *relativistic* shock with magnetic field *perpendicular* to particle flow (so DSA doesn't work)

- relevant to pulsar wind nebulae, e.g. Crab
- principle: consider relativistic plasma whose mass is dominated by ions ( $m_i/m_{e\pm} \gg 1$ )
  - ions gyrate *coherently* in magnetic field
  - they therefore radiate ion cyclotron waves (Alfven waves) at shock front
  - positrons and electrons absorb these resonantly and are accelerated to high Lorentz factors with fairly high efficiency (few % of upstream energy density converted to non-thermal  $e^\pm$ )
- mechanism seems to account well for high-energy emission; not so clear that it deals with radio–IR emission
  - two different electron populations?
  - but consistency of spectra suggest otherwise



# RSA Simulations



- Simulation by Amato & Arons (*Apj* **653** (2006) 325)

- Input parameters:

- $N_i/N_{e\pm} = 0.1$
- $m_i/m_{e\pm} = 100$

72% of upstream energy density carried by ions

Result:

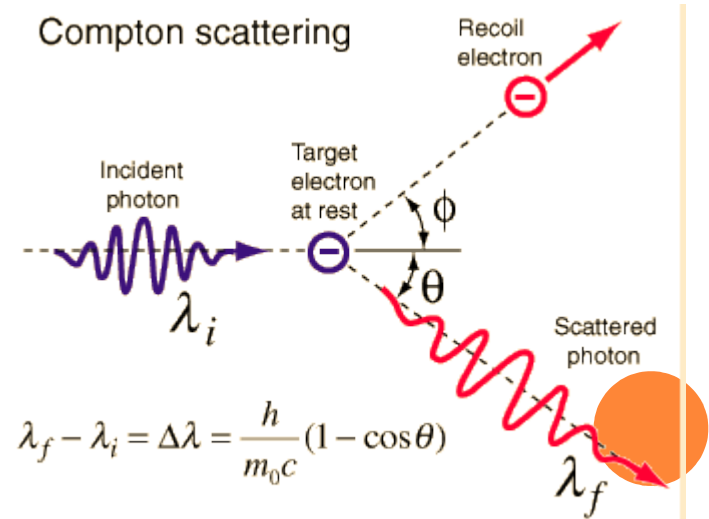
- 5% of upstream energy density winds up in accelerated  $e\pm$

Less extreme ion loading, e.g.,  $m_i/m_{e\pm} = 20$ , preferentially accelerates positrons



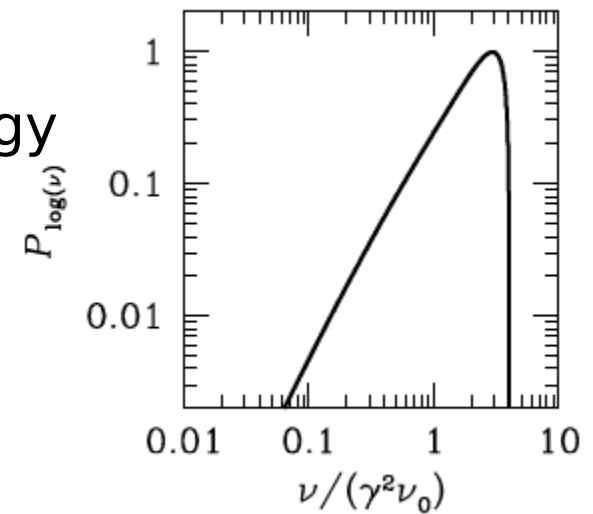
# Photons and Neutrinos

- High-energy photons and neutrinos are **secondary particles** produced by interactions of high-energy primaries.
  - production mechanisms:
    - inverse Compton scattering (photons only)
    - Low-energy photon backscatters off high-energy electron.
    - In electron rest frame we have
      - $\Delta\lambda = h(1 - \cos\theta)/mc^2$ .
  - In lab frame, maximum energy gain
  - occurs in head-on collision:
    - $v \approx 4\gamma^2 v_0$
  - Because of relativistic
  - aberration, spectrum is
    - sharply peaked near maximum

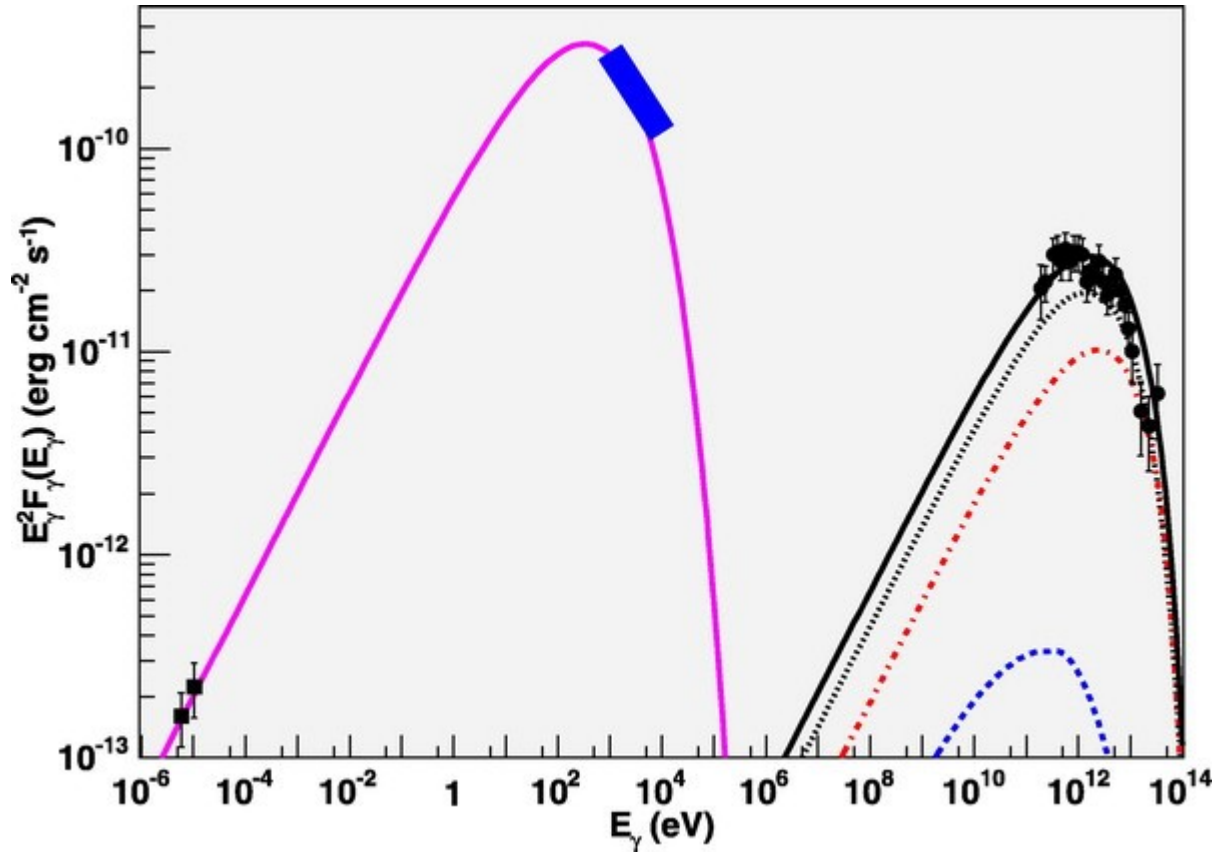


# Photons and Neutrinos

- inverse Compton scattering (continued)
- Plot shows calculated spectrum for
  - monoenergetic photons and electrons.
  - Plenty of potential sources of low-energy photons to be upscattered:
    - synchrotron radiation produced by the same population of fast electrons
    - **(synchrotron-self-Compton, SSC)**
    - cosmic microwave background
    - optical photons from source
  - For real objects, need to integrate over power-law spectrum of electrons and spectrum of photon source



# Spectrum of RXJ 1713.7–3946



Assumed  
distance 1 kpc,  
electron  
luminosity  
 $1.5 \times 10^{30}$  W,  
 $B = 12 \mu\text{G}$

Source  
photons  
include optical,  
IR, CMB

Porter, Moskalenko & Strong, *Apj* **648**  
(2006) L29-L32

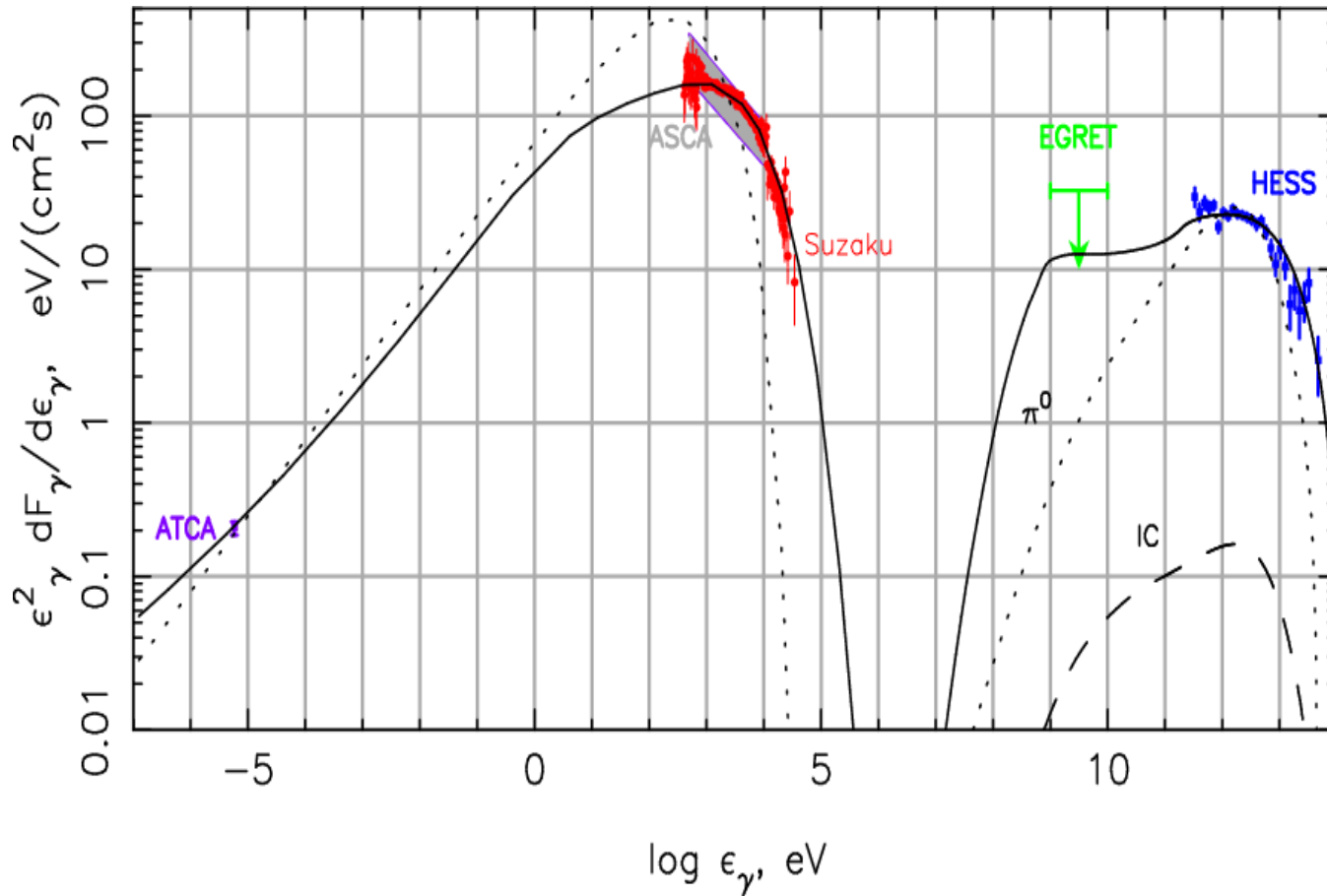


# Photons and Neutrinos

- High-energy photons and neutrinos are **secondary particles** produced by interactions of high-energy primaries.
  - production mechanisms:
    - pion decay (photons and neutrinos)
      - pions produced by high-energy proton colliding with either matter or photons (**pion photoproduction**)
      - neutral pions decay to  $\gamma\gamma$ , charged to  $\mu\nu_\mu$ 
        - mechanism produces both high-energy  $\gamma$ -rays and neutrinos
  - Both mechanisms need population of relativistic charged particles
    - electrons for IC, protons for pion decay
  - Unclear which dominates for observed TeV  $\gamma$ -ray sources



# Spectrum of RXJ 1713.7–3946, take 2



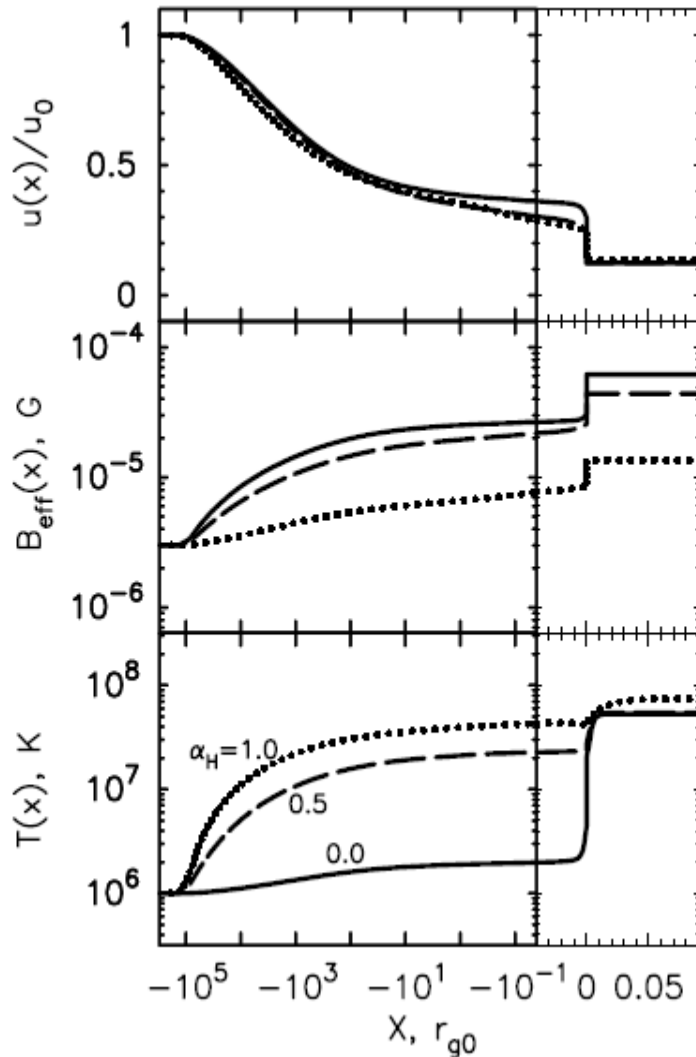
Uses DSA to accelerate protons.

$B = 142 \mu\text{G}$   
downstream of shock.

High B-field enhances synchrotron relative to inverse Compton



# Are high magnetic fields plausible?



- Hadronic model fit to RXJ 1713 needs  $B > 100 \mu\text{G}$ 
  - much larger than ambient Galactic B-fields
  - amplification required to make DSA fits self-consistent
  - fortunately modelling indicates that the interaction of the accelerated CRs with the magnetic field induces turbulence, resulting in amplification
- Direct observational evidence of high B-fields in some SNRs
  - e.g. Cas A,  $B > 500 \mu\text{G}$  from comparing synchrotron & IC/bremsstrahlung contributions
    - (Vink & Laming, *ApJ* **584** (2003) 758)



# Acceleration: Summary

- Observations made in high-energy astroparticle physics require that charged particles be accelerated to very high energies ( $\sim 10^{20}$  eV)
- Likely candidate is diffusive shock acceleration
  - requirement of shocks associated with magnetic fields found in many astrophysical objects, especially supernova remnants and AGN
  - synchrotron radiation from these objects direct evidence for population of fast electrons
  - much less evidence for presence of relativistic hadrons, but there must be some somewhere since we observe them in cosmic rays!
- TeV  $\gamma$ -rays can be produced by fast electrons using inverse Compton scattering, or by fast protons from  $\pi^0$  decay
  - latter will also make TeV neutrinos, not yet observed





# High Energy Astroparticle Physics

Acceleration Mechanisms  
Sources  
Detection



# Gamma-Ray Astronomy

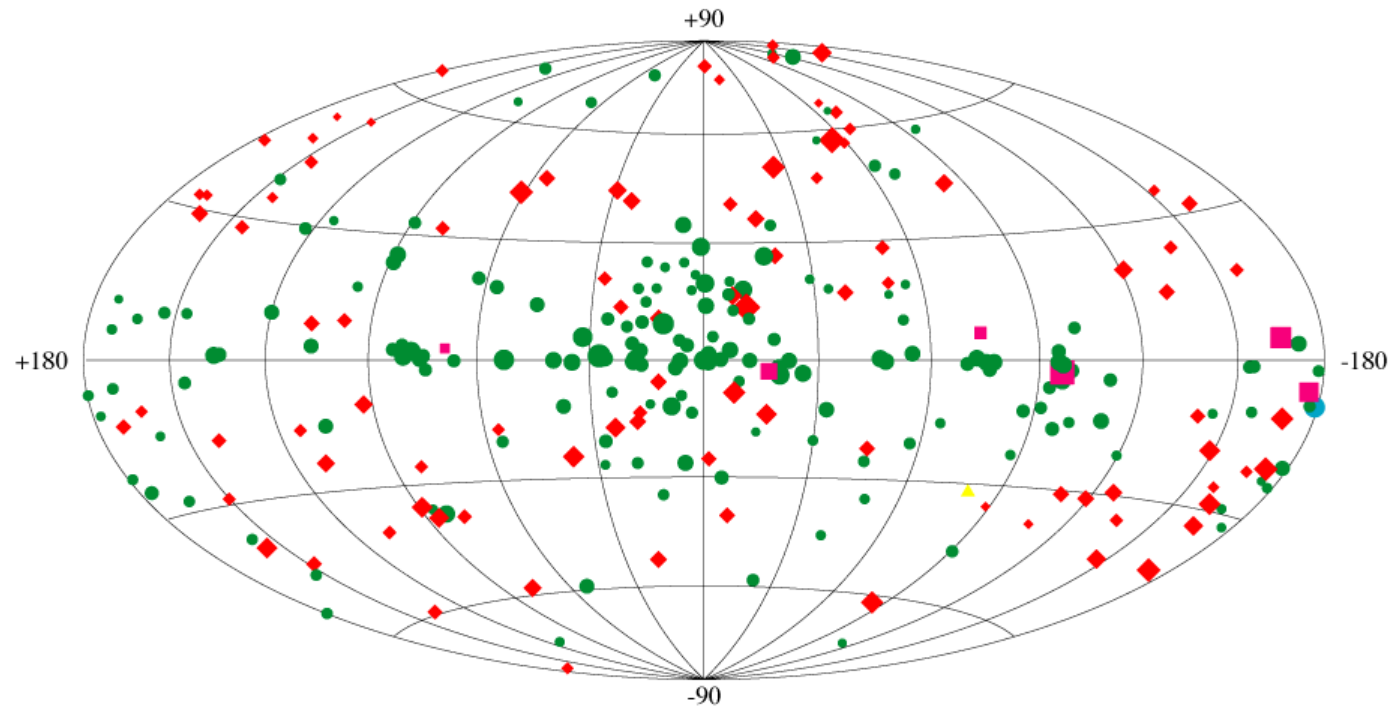
- Well-established branch of high-energy astrophysics
  - most work done at modest energies (few 10s of MeV)
    - some, e.g. EGRET, out to few 10s of GeV
  - this is not usually regarded as astroparticle physics
    - though EGRET catalogue sometimes used as list of candidates for, e.g., neutrino point source searches
- Atmosphere is not transparent to gamma rays
  - low and medium energy  $\gamma$ -ray astronomy is space-based
    - CGRO, SWIFT, GLAST, INTEGRAL, etc.
  - space platforms not suitable for TeV  $\gamma$ -ray astronomy
    - too small!
  - therefore very high energy  $\gamma$ -ray astronomy is a ground-based activity
    - detect shower produced as  $\gamma$ -ray enters atmosphere



# EGRET Point Sources

## Third EGRET Catalog

$E > 100 \text{ MeV}$



- ◆ Active Galactic Nuclei
- Unidentified EGRET Sources
- Pulsars
- ▲ LMC
- Solar FLare

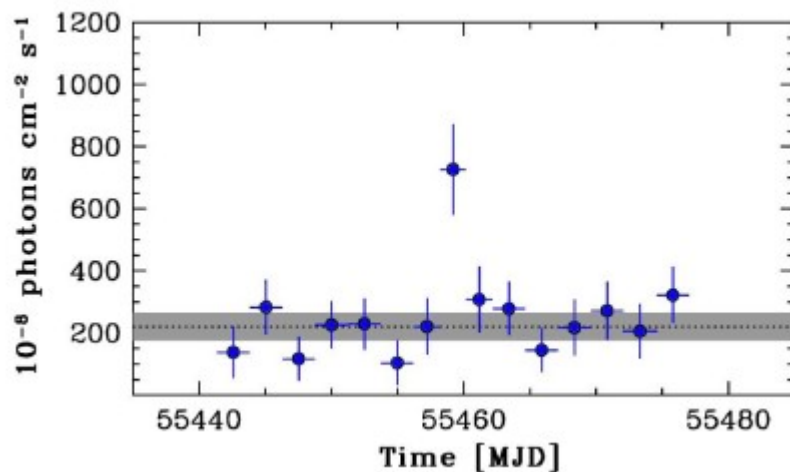
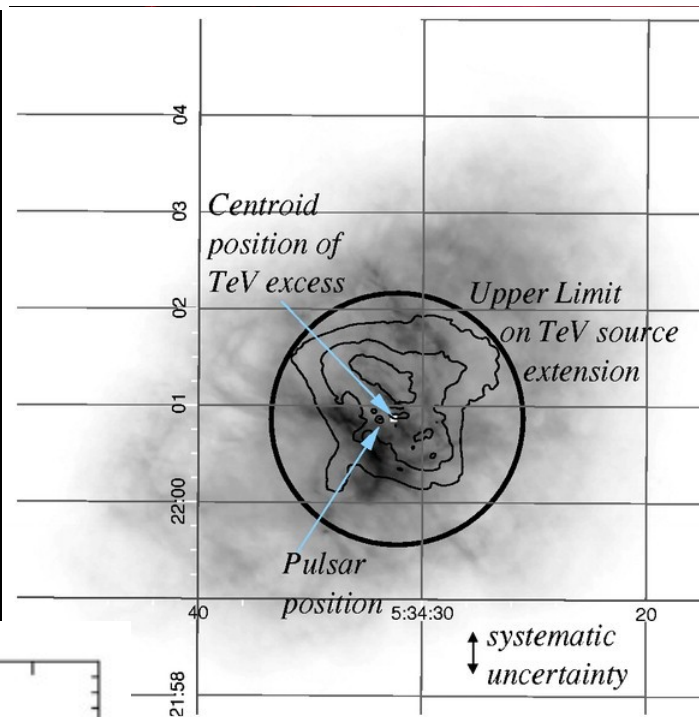


# Gamma-Ray Sources

- From maps, clearly mixed Galactic and extragalactic
  - extragalactic sources of TeV  $\gamma$ s are mostly blazars (a class of AGN where we are looking down the jet)
  - identified Galactic sources are SN-related (supernova remnants and pulsar wind nebulae), plus a few binary compact objects
  - dark/unidentified objects associated with Galactic plane, therefore presumably Galactic
- SNRs and AGN are suitable environments for particle acceleration
  - shocks, magnetic fields, synchrotron emission



# Pulsar Wind Nebula: The Crab

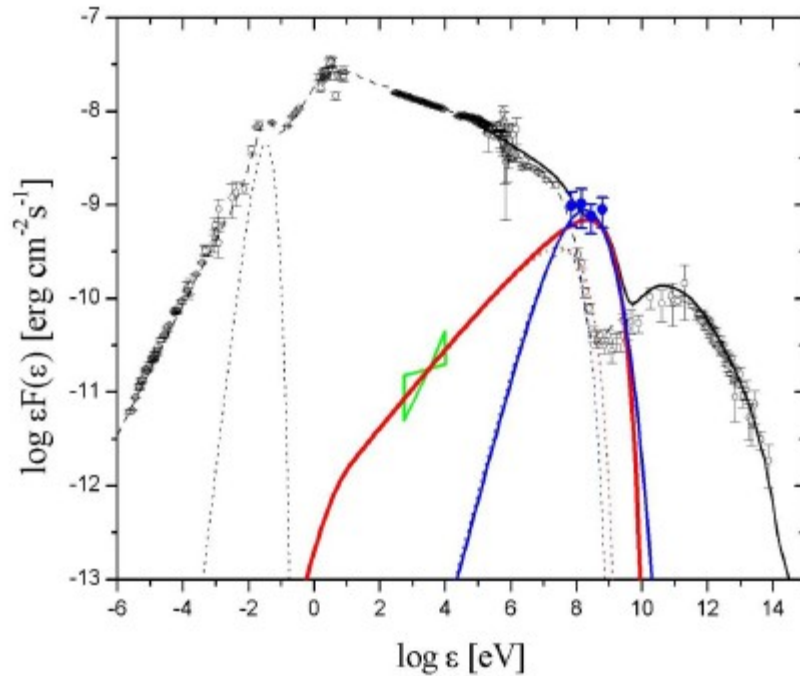


TeV gamma-ray signal as observed by HEGRA (Aharonian et al. 2004)

Medium-energy  $\gamma$ -ray flare observed by AGILE (Tavani et al. 2011)

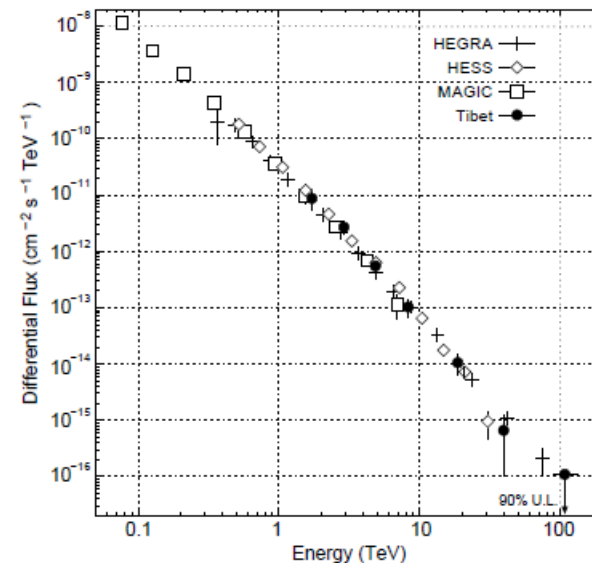


# Pulsar Wind Nebula: The Crab

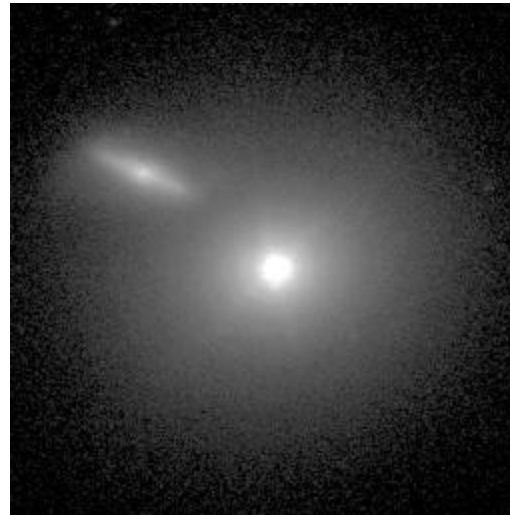
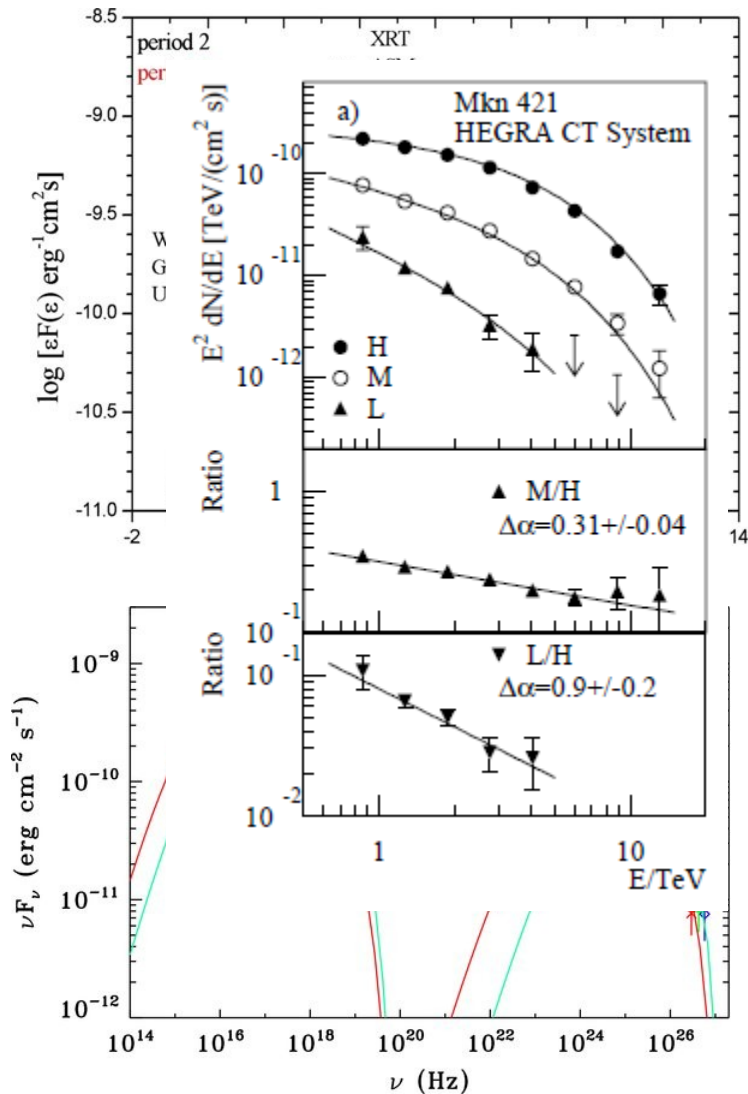


Crab spectral energy distribution showing September 2010 flare

TeV energy spectrum



# Blazar: Mkn 421



Mkn 421 and companion galaxy. Aimo Sillanpaa, Nordic Optical Telescope. (Above: very boring X-ray image by Chandra)

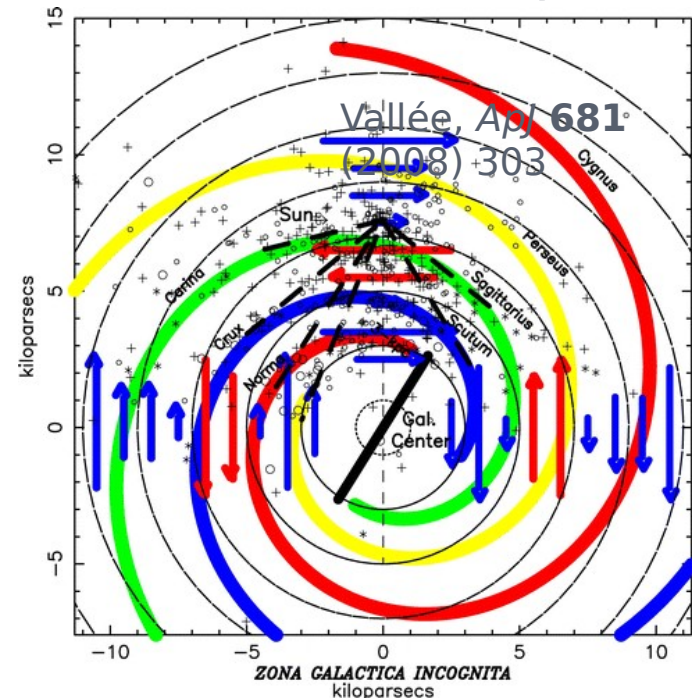
Highly variable (typical of blazars)  
Spectrum varies according to state





# Cosmic Ray Sources

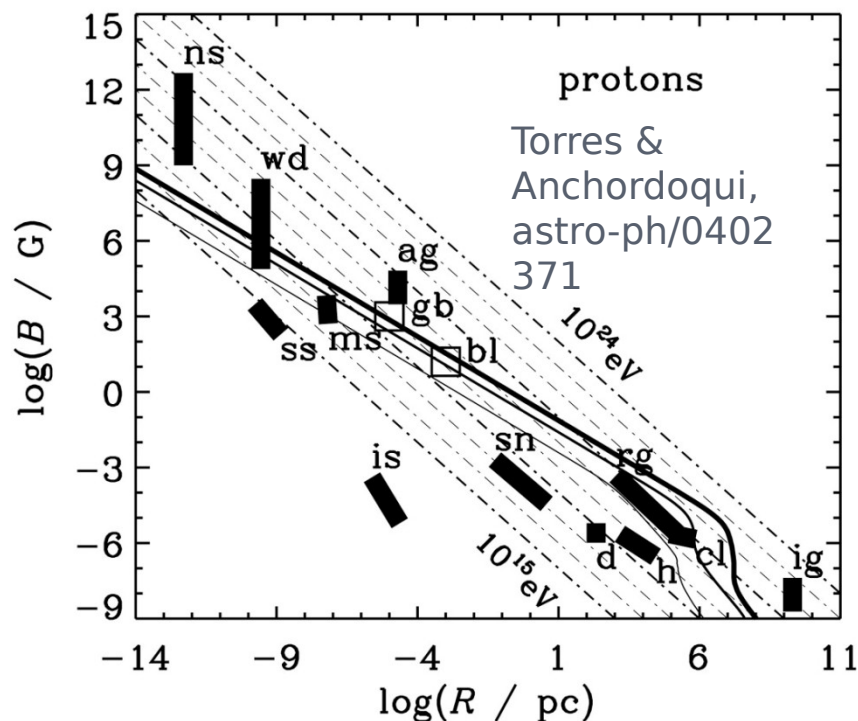
- Observations of cosmic rays now span about 100 years
- However, sources are not definitively established
- Galaxy has a complex magnetic field which effectively
- scrambles direction of
  - charged particles
  - Gamma ray luminosity
  - requires fast particles,
  - but maybe only electrons
  - therefore, observation of
  - $\gamma$ -rays does not definitively
  - establish source as a cosmic
    - ray factory
- Neutrino luminosity *does* require fast hadrons
- but no neutrino point sources yet to be correlated with cosmic rays



# Cosmic Ray Sources

General dimensional analysis suggests

- $E_{\text{max}} [\text{GeV}] \approx 0.03 \eta^{-1} Z R[\text{km}] B[\text{G}]$  (Hillas condition)
- basically requires particles to remain confined in accelerating region
- quite difficult to satisfy for highest-energy CRs
- plot shows
  - neutron stars
  - white dwarfs
  - sunspots
  - magnetic stars
  - active galactic nuclei
  - interstellar space
  - supernova remnants
  - radio galaxy lobes
  - disc and halo of Galaxy
  - galaxy clusters
  - intergalactic medium
  - gamma-ray bursts
  - blazars
  - shock-wave velocities



# Cosmic Ray Sources

- Amount of magnetic deflection decreases with increasing energy
  - highest energy events might remember where they came from...
  - Pierre Auger Observatory observes significant correlation between arrival directions of CRs above 55 EeV and a catalogue of AGN
    - $38 \pm 7\%$  of events within  $3.1^\circ$  of a catalogued
    - nearby AGN, cf. 21% expected for intrinsically isotropic distribution
    - similar results found for SWIFT catalogue—data do however require significant isotropic component (40–80%)



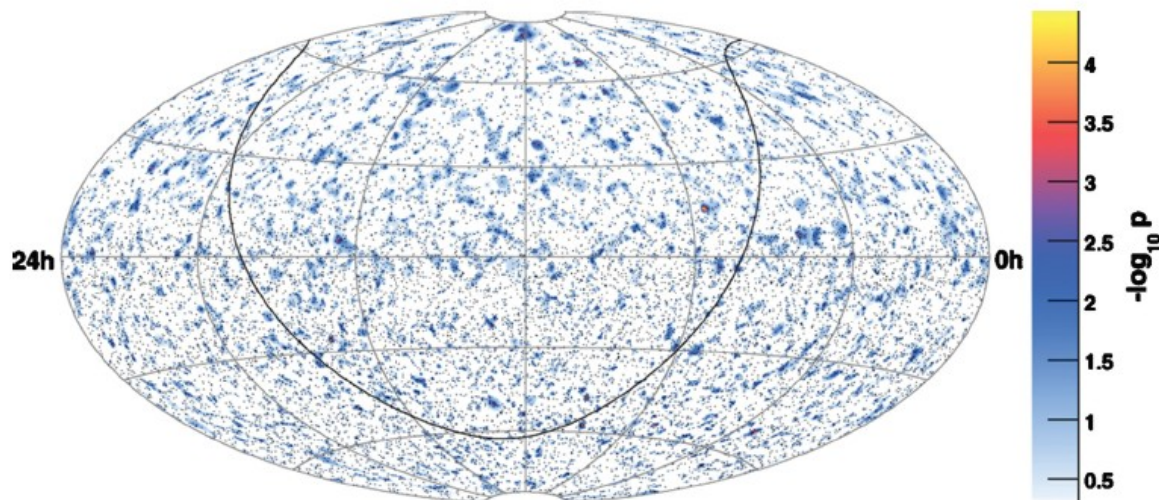
# Cosmic Ray Sources: Summary

- CRs up to about  $10^{15}$  eV or so assumed to come from SNRs
  - but they don't provide good directional information, so this remains to be confirmed
    - neutrino observations, or definitive proof that some SNR  $\gamma$ -rays originate from  $\pi^0$  decay
- Ultra-high energy CRs may come from local AGN
  - statistically significant (but partial) correlation
  - note that intergalactic space is not completely transparent to UHECRs—see later—so *distant* AGN (beyond  $\sim 100$  Mpc) are assumed not to contribute



# Neutrino Sources

- Known sources of low-energy (0.1–100 MeV) neutrinos:
  - Sun
  - SN 1987A
- Known sources of high-energy neutrinos:
  - Starting to develop (since 2018) → see third lecture



IceCube  
search for  
point sources.



# Sources: Summary

- TeV gamma rays are observed from a variety of sources, primarily SNRs within the Galaxy and blazars outside
  - clear evidence of charged particles accelerated to very high energies, but whether electrons or hadrons is unclear
- Cosmic ray sources are difficult to pinpoint because CRs are strongly deflected by the Galactic magnetic field
  - SNRs suspected to be source of CRs at  $<10^{15}$  eV
  - some hints that local AGN may be responsible for highest energy CRs
- Observations of high energy neutrinos would solve the mystery, but need more statistics



