Dark Matter Theory: Status and Updates Kathryn M Zurek





Tuesday, February 16, 16

Universe's Energy Budget

Dynamical selection?



New Dynamics, Definitely BSM

Tuesday, February 16, 16

We have essentially eliminated a SM explanation; need physics BSM



Why particle dark curvature, z_eq matter?



sound speed = baryon to radiation ratio



Why not just ordinary (dark) baryons?

A: BBN and CMB make independent measurements of the baryon fraction. Observations only accounted for with non-interacting matter

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Make baryons non-interacting by binding DM into MaCHOs?

A: looked for those and did not find them; eliminated MACHO range from $\gtrsim 10^{-8} M_{\odot}$,







Make baryons non-interacting by binding DM into MaCHOs?

So A: looked for those and did not find them; from 2005 talk by K. Griest eliminated MACHO range from $\gtrsim 10^{-8} M_{\odot}$ Afshordi, McDonald, Spergel





Why not modify gravity?

 A: Modified gravity theories tend to be sick



 A: Must get the entire range of observations right, not just galactic rotation curves

Why not modify gravity?

 A: Modified gravity theories tend to be sick





X-ray: NASA/CXC/CfA/ <u>M.Markevitch</u> et al.; Lensing Map: NASA/STScI; ESO WFI; Magellan/U.Arizona/ <u>D.Clowe et al.</u> Optical: NASA/STScI; Magellan/U.Arizona/D.Clowe et al

A: Must get the entire range of observations right, not just galactic rotation curves

By contrast, it is easy to explain everything with particle dark matter

From theoretical point of view, theories are compelling, testable.

As the saying goes:



Particle dark matter

No shortage of theories

Supersymmetry

Sector Extra dimensions

Massive neutrino

MeV dark matter

Scalar dark matter

axion



Particle dark matter

No shortage of theories

 Axions and WIMPs (usually, supersymmetric)

Note however: most based on a couple of very popular theories



Dark Matter: Standard Paradigm

Usual picture of dark matter is that it is:
single
stable
(sub-?) weakly interacting
neutral

HIDDEN DARK WORLDS

Our thinking has shifted



$M_p \sim 1 \,\,{\rm GeV}$

Standard Model

From a single, stable weakly interacting particle (WIMP, axion)

> Models: Supersymmetric light DM sectors, Secluded WIMPs, WIMPless DM, Asymmetric DM Production: freeze-in, freeze-out and decay, asymmetric abundance, non-thermal mechanicsms

...to a Hidden Valley with multiple states, new interactions

Models of Dark Matter

The classicSUSY



has all the ingredients and they are present for other reasons DM (sort of) free

IDEA FOCUS: SUPERSYMMETRY

- Provides sharp predictions
- Must be neutral.
- Options sneutrino, bino, wino, higgsino $\tilde{\nu}$ \tilde{B} , \tilde{W}_3 , \tilde{H}
- Sneutrino scatters through Z

Weakly-interacting

Sneutrino, also being neutral, is a good DM candidate.... except for direct detection(!)

 $Q|\text{neutrino}\rangle = |\text{sneutrino}\rangle$ Gauge interaction:



Its couplings are fixed by gauge interactions

Scatters off nucleons through Z boson

Let's compute the rate



Apply to scattering through Z boson

plug in and compare

$$\sigma \approx \frac{g^4 \mu_n^2}{4\pi m_Z^4} \approx 10^{-39} \text{ cm}^2$$

Active $\tilde{\nu}$ DM excluded by direct detection



Can evade constraint by mixing in sterile $\tilde{\nu}$, \tilde{N} . This state does not couple to Z. But is not present in minimal model

What about neutralino?

 2 component fermion χ Majorana fermion
 Ø Possible operators, four Fermi, V-A structure: $\mathcal{O}_{SI} = (\bar{\chi}\gamma_{\mu}\chi)(\bar{q}\gamma^{\mu}q) = 0$ $\mathcal{O}_{SD} = (\bar{\chi}\gamma_{\mu}\gamma_{5}\chi)(\bar{q}\gamma^{\mu}\gamma_{5}q)$ $\mathcal{O}_{\text{vel dep.}} = (\bar{\chi}\gamma_{\mu}\gamma_{5}\chi)(\bar{q}\gamma^{\mu}q)$ SI vanishes identically; others are SD or velocity suppressed

Higgs Scattering

 $\frac{f_{p,n}}{m_{p,n}} = \sum_{q=u,d,s} f_{Tq}^{p,n} \frac{y_q}{m_q} + \frac{2}{27} f_{TG}^{p,n} \sum_{q=c,b} f_{TG}^{p,n} \sum_{q=c$

So neutralino is safe from Z-pole scattering

It scatters predominantly through Higgs boson

 Higgs boson coupling to nucleon comes predominantly through a loop





Shifman, Vainshtein, Zakharov, Phys.Lett. B78 (1978) 443

Higgs scattering crosssection



Are there ways around?

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A bit about neutralino couplings

Supersymmetry fixes what interactions can and cannot occur

Higgs does not interact with a "pure" state



Must have bino-Higgsino or Higgsino-wino mix

Neutral

Mass matrix:



Soft parameters, M_1 and M_2 . Free in SUSY.

In SM, one Higgs works b/c can write field and conjugate $\mathcal{L}_{SM} = \bar{u}y_u Q\phi - \bar{d}y_d Q\phi^* - \bar{e}y_e L\phi^*$

∧ Not so in SUSY:
 $W_{MSSM} = \bar{u}y_uQH_u - \bar{d}y_dQH_d - \bar{e}y_eLH_d$ $\tan \beta = \frac{v_u}{v_d}$ $v_u^2 + v_d^2 = v^2 = (246 \text{ GeV})^2$

WIMP annihilation

processes





Bottom diagrams often dominate if DM is largely wino or largely Higgsino

Escaping direct detection constraints

So even if direct detection constraints are escaped by making neutralino pure

 there may be strong indirect detection constraints

Photons from annihilation in galaxy today constrain pure wino or Higgsino DM



Escaping direct detection constraints

Big cross-section!





Cohen, Lisanti, Pierce, Slatyer

Pure bino DM escapes

- While wino and Higgsino may be constrained by indirect detection, bino escapes
- ${\it @}$ But, even bino has Higgsino component set by μ
- Require $\mu \gg M_1 \sim m_{wk}$ to get rid of Higgsino component

Same parameter enters into Z boson mass

$$m_Z^2 = \frac{|m_{H_d}^2 - m_{H_u}^2|}{\sqrt{1 - \sin^2(2\beta)}} - m_{H_u}^2 - m_{H_d}^2 - 2|\mu|^2$$

Must tune parameters

Loops Matter

 Even if Tree scattering process vanishes, Future
 WIMP DM probes can reach
 1-loop suppressed processes!

 1-loop suppressed wino can be ruled out

I-loop Higgsino harder







Loops Matter

Even 1-loop bino can
 be probed in some
 cases

OM experiments searching for WIMP enter into precision era



Berlin, Robertson, Solon, KZ, 1511.05964



 $\sigma_{\rm SD} \ [{
m cm}^2]$

Waiting for SUSY

ATLAS SUSY Searches* - 95% CL Lower Limits

Status: SUSY 2013

Model

MSUGRA/CMSSM

MSUGRA/CMSSM

MSUGRA/CMSSM

 $\tilde{g}\tilde{g}, \tilde{g} \rightarrow gq(\ell\ell/\ell\nu/\nu\nu)\tilde{\chi}_1^0$

 $\bar{q}\bar{q}, \bar{q} \rightarrow q\tilde{t}_1^0$

28.8→q31

GMSB (7 NLSP)

GMSB (7 NLSP)

Gravitino LSP

 $\tilde{g} \rightarrow t \bar{t} \tilde{\chi}_{j}$

 $\tilde{g} \rightarrow t \bar{t} \tilde{X}$

 $\bar{g} \rightarrow b\bar{t}\bar{t}_1$

 $b_1 b_1, b_1 \rightarrow b \tilde{t}_1$

 $b_1 b_1, b_1 \rightarrow t \tilde{t}_1^d$

 $\tilde{t}_1 \tilde{t}_1$ (light), $\tilde{t}_1 \rightarrow b \tilde{t}_1$

 $\tilde{t}_1 \tilde{t}_1$ (light), $\tilde{t}_1 \rightarrow Wb\tilde{t}_1$

 $\tilde{t}_1 \tilde{t}_1 (medium), \tilde{t}_1 \rightarrow t \tilde{\tilde{t}}_1$

 $\tilde{t}_1 \tilde{t}_1 (\text{medium}), \tilde{t}_1 \rightarrow b \tilde{t}_1$

 $\tilde{t}_1 \tilde{t}_1$ (heavy), $\tilde{t}_1 \rightarrow t \tilde{t}_1^{c}$

 $\tilde{t}_1 \tilde{t}_1$ (heavy), $\tilde{t}_1 \rightarrow t \tilde{\chi}_1^0$

 $\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow c \tilde{t}_1$ ž₁ ž₁(natural GMSB)

 $\tilde{t}_2 \tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$

 $\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell \tilde{\chi}_1$

 $\tilde{\chi}_1^{\pm} \tilde{\chi}_2^0 \rightarrow W \tilde{\chi}_1^0 Z \tilde{\chi}_1^0$

RtR→WRhR

 $\tilde{q}\tilde{q}, \tilde{\chi}_1^0 \rightarrow qq\mu$ (RPV)

Bilinear RPV CMSSM

partial data

full data

 $\tilde{g} \rightarrow q q q$

 $\tilde{g} \rightarrow \tilde{t}_1 t, \tilde{t}_1 \rightarrow bs$

 $\tilde{\chi}_1 \tilde{\chi}_1, \tilde{\chi}_1 \rightarrow \tilde{\ell} \nu(\ell \bar{\nu})$

GGM (bino NLSP)

GGM (wino NLSP)

GGM (higgsino NLSP)

Searches

clusive

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in the

EW

 $\int \mathcal{L} dt = (4.6 - 22.9) \text{ fb}^{-1}$ $\sqrt{s} = 7,8$ TeV e, μ, τ, γ Jets E_{τ}^{miss} [£ dt[fb⁻¹] Mass limit Reference 0 2-6 jets Yes 20.3 1.7 TeV m(q)=m(g) ATLAS-CONF-2013-047 3-6 jets 20.3 any m(q) ATLAS-CONF-2013-062 1 e.µ Yes 1.2 TeV 0 7-10 jets Yes 20.3 1.1 TeV any m(g) 1308.1841 2-6 jets ATLAS-CONF-2013-047 0 Yes 20.3 740 GeV m(R1)=0 GeV 2-6 jets ATLAS-CONF-2013-047 Yes 20.3 1.3 TeV 0 m(x1)=0 GeV 1 e.µ $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^{\pm} \rightarrow qqW^{\pm}\tilde{\chi}_1^{0}$ 3-6 jets ATLAS-CONF-2013-062 Yes 20.3 1.18 TeV $m(\tilde{\chi}_1^2) < 200 \text{ GeV}, m(\tilde{\chi}^*) = 0.5(m(\tilde{\chi}_1^2) + m(\tilde{g}))$ 2 e.µ 0-3 jets 20.3 1.12 TeV m(k1)=0 GeV ATLAS-CONF-2013-089 2 e.µ 2-4 jets Yes 4.7 tan//<15 1208.4688 1.24 TeV 1-2 7 0-2 jets Yes 20.7 tang >18 ATLAS-CONF-2013-026 1.4 TeV 1.07 TeV 2γ Yes m(k1)>50 GeV 4.8 1209.0753 1 e. µ + y Yes 4.8 619 GeV m(k1)>50 GeV ATLAS-CONF-2012-144 GGM (higgsino-bino NLSP) 1bYes 4.8 900 GeV m(x1)>220 GeV 1211.1167 $2e_{,\mu}(Z)$ 0-3 jets Yes 5.8 690 GeV m(H)>200 GeV ATLAS-CONF-2012-152 Yes m(g)>10⁻⁴ eV ATLAS-CONF-2012-147 Ö mono-jet 10.5 645 GeV 0 3 b 20.1 1.2 TeV m(R1)<600 GeV ATLAS-CONF-2013-061 Yes 7-10 jets Yes 20.3 1.1 TeV m(\$1) <350 GeV 1308.1841 0 0-1 e.µ 1.34 TeV ATLAS-CONF-2013-061 3 b Yes 20.1 m(R1)<400 GeV 0-1 e, µ 3 b Yes 20.1 1.3 TeV m(x1)<300 GeV ATLAS-CONF-2013-061 0 2 b Yes 20.1 100-620 GeV m(k1)<90 GeV 1308 2631 2 e, µ (SS) ATLAS-CONF-2013-007 0-3 b Yes 20.7 275-430 GeV $m(\tilde{t}_{1}^{2})=2 m(\tilde{t}_{1}^{2})$ Б. 110-167 GeV 1-2 e. µ 1-2 b Yes 4.7 m(k1)=55 GeV 1208.4305.1209.2102 2 e.µ 0-2 jets Yes 20.3 ATLAS-CONF-2013-048 130-220 GeV ÷. $m(\tilde{t}_{1}^{0}) = m(\tilde{t}_{1}) - m(W) - 50 \text{ GeV}, m(\tilde{t}_{1}) < < m(\tilde{t}_{1}^{0})$ 2 e.µ 2 jets Yes 20.3 Ē, 225-525 GeV m(R1)=0 GeV ATLAS-CONF-2013-065 2 b 0 Yes 20.1 ÷. 150-580 GeV m(ti)<200 GeV, m(ti)-m(ti)=5 GeV 1308.2631 1 e.µ Yes 20.7 200-610 GeV m(t1)=0 GeV ATLAS-CONF-2013-037 1 b Ē. 0 2bYes 20.5 320-660 GeV ATLAS-CONF-2013-024 m(x1)=0 GeV 90-200 GeV ATLAS-CONF-2013-068 0 mono-jet/c-tag Yes 20.3 Ē. m(t
1)-m(t
1)<85 GeV ATLAS-CONF-2013-025 2 e. µ (Z) 1 b Yes 20.7 500 GeV m(x1)>150 GeV 3 e, µ (Z) 1 b Yes 20.7 271-520 GeV m(t
1)=m(t
1)+180 GeV ATLAS-CONF-2013-025 Ē, 2 e.u 20.3 85-315 GeV ATLAS-CONF-2013-049 Ō Yes m(k1)=0 GeV ATLAS-CONF-2013-049 2 e.µ Ö Yes 20.3 125-450 GeV $m(\tilde{\chi}_1^0)=0$ GeV, $m(\tilde{\ell}, \tilde{\tau})=0.5(m(\tilde{\chi}_1^0)+m(\tilde{\chi}_1^0))$ $\tilde{\chi}_{1}^{\uparrow}\tilde{\chi}_{1}^{\uparrow}, \tilde{\chi}_{1}^{\uparrow} \rightarrow \tilde{\tau}\nu(\tau\tilde{\nu})$ $\tilde{\chi}_{1}^{\uparrow}\tilde{\chi}_{2}^{\uparrow} \rightarrow \tilde{\ell}_{L}\nu\tilde{\ell}_{L}\ell(\tilde{\nu}\nu), \ell\tilde{\nu}\tilde{\ell}_{L}\ell(\tilde{\nu}\nu)$ 2τ Yes 20.7 180-330 GeV $m(\tilde{k}_{1}^{0})=0$ GeV, $m(\tilde{\tau}, \tilde{\tau})=0.5(m(\tilde{k}_{1}^{0})+m(\tilde{k}_{1}^{0}))$ ATLAS-CONF-2013-028 0 20.7 $m(\tilde{t}_{1}^{+})=m(\tilde{t}_{2}^{0}), m(\tilde{t}_{1}^{0})=0, m(\tilde{t}, \tilde{v})=0.5(m(\tilde{t}_{1}^{+})+m(\tilde{t}_{1}^{0}))$ ATLAS-CONF-2013-035 3 e.u Yes 600 GeV 3 e, µ Ō 20.7 $m(\tilde{k}_1^n)=m(\tilde{k}_2^0), m(\tilde{k}_1^0)=0$, sleptons decoupled ATLAS-CONF-2013-035 Yes 315 GeV 1 e, µ 2 b Yes 20.3 285 GeV $m(\tilde{t}_1^*)=m(\tilde{t}_2^0), m(\tilde{t}_1^0)=0$, sleptons decoupled ATLAS-CONF-2013-093 Direct $\tilde{x}_1 \tilde{x}_1$ prod., long-lived \tilde{x}_1^* Disapp. trk 1 jet Yes 20.3 270 GeV $m(\tilde{t}_1^*)-m(\tilde{t}_1^0)=160$ MeV, $\tau(\tilde{t}_1^*)=0.2$ ns ATLAS-CONF-2013-069 Stable, stopped g R-hadron 22.9 832 GeV $m(\tilde{k}_1^0)=100 \text{ GeV}, 10 \,\mu\text{s} < r(\tilde{g}) < 1000 \text{ s}$ ATLAS-CONF-2013-057 0 1-5 jets Yes GMSB, stable $\bar{\tau}, \bar{\chi}_{1}^{0} \rightarrow \bar{\tau}(\bar{e}, \bar{\mu}) + \tau(e, \mu) = 1.2 \mu$ 475 GeV 10<tan/l<50 15.9 ATLAS-CONF-2013-058 2γ 230 GeV 0.4<r(21)<2 ns 1304.6310 GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$, long-lived $\tilde{\chi}_1^0$ Yes 4.7 1 µ, displ. vtx -20.3 1.0 TeV 1.5 <cr<156 mm, BR(µ)=1, m(k1)=108 GeV ATLAS-CONF-2013-092 X'_311=0.10, X132=0.05 LFV $pp \rightarrow \bar{\nu}_r + X, \bar{\nu}_r \rightarrow e + \mu$ 2 e, µ 4.6 1.61 TeV 1212.1272 LFV $pp \rightarrow \tilde{v}_r + X, \tilde{v}_r \rightarrow e(\mu) + \tau$ λ'₃₁₁=0.10, λ_{1/2/33}=0.05 1.1 TeV 1 e. µ + 7 4.6 1212.1272 7 jets Yes 4.7 m(q)=m(g), cr15p<1 mm ATLAS-CONF-2012-140 1 e, µ 1.2 TeV $\begin{array}{l} \tilde{\chi}_1^+ \tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow W \tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow e e \tilde{v}_{\mu}, e \mu \tilde{v}_e & 4 \ e, \mu \\ \tilde{\chi}_1^- \tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow W \tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \tau \tau \tilde{v}_e, e \tau \tilde{v}_\tau & 3 \ e, \mu + \tau \end{array}$ 760 GeV Yes 20.7 m(R1)>-300 GeV, A121>0 ATLAS-CONF-2013-036 20.7 350 GeV m(R1)>80 GeV, J111>0 ATLAS-CONF-2013-036 Yes 0 6-7 jets 20.3 916 GeV BR(t)=BR(b)=BR(c)=0% ATLAS-CONF-2013-091 2 e, µ (SS) 20.7 ATLAS-CONF-2013-007 0-3 b Yes 880 GeV Scalar gluon pair, sgluon→qq 0 4 jets 4.6 saluor 100-287 GeV 1210.4826 incl. limit from 1110.2693 Scalar gluon pair, sgluon→tīt 2 e, µ (SS) 1 b Yes 14.3 ATLAS-CONF-2013-051 WIMP interaction (D5, Dirac x) m(y)<80 GeV, limit of<687 GeV for D8 ATLAS-CONF-2012-147 0 mono-jet Yes 10.5 704 GeV √s = 8 TeV √s = 8 TeV 10⁻¹ 1

Mass scale [TeV]

ATLAS Preliminary

$\sqrt{s} = 7 \text{ TeV}$ full data Tuesday, February 16, 16

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LHC is not a DM Machine

• Strong constraints on strong particles



> 1 TeV

• Weak constraints on weak particles



> 200-300 GeV

LHC is a Mediator Machine

• Strong constraints on strong particles



> 1 TeV

• Weak constraints on weak particles



> 200-300 GeV

"Model Independent" Collider Searches for DM









Mono-X is not a discovery mode

Important theory dependence in these plots!

Inappropriate use of higher dimension operators

Failure to take into account direct searches for mediator







Domain of Importance

 LHC is effective on heavy, relatively strongly coupled mediators

Direct detection is unparalleled when mediator is light, small couplings to protons

Direct searches for mediator almost always more constraining than mono-X





Search	Model where it matters	
mono-h	Inelastic DM, 2HDM	
mono- z	Inelastic DM, 2HDM	
nono-jet	squark mediated production, compressed spectrum	
mono-b	sbottom mediated production, compressed spectrum	

stop mediated production, RPV-like

mono-t

q

 $v_{\rm met}$

When Should We Start Looking Elsewhere?

Cannot kill neutralino DM via direct detection, but paradigm does become increasingly tuned

Likewise, LHC will only continue to push constraints on mediating SUSY colored particles up, though relatively weak constraints on DM itself



Dark Matter Model Dynamics

(Looking beyond the vanilla WIMP paradigm)

DM Paradigm: recap

Usual picture of dark matter is that it is:
single
stable
(sub-?) weakly interacting
neutral

Supersymmetry and axions fit the bill.

Hidden Dark Worlds Our thinking has shifted



From a single, stable weakly interacting particle (WIMP, axion)

> Models: Supersymmetric light DM sectors, Secluded WIMPs, WIMPless DM, Asymmetric DM Production: freeze-in, freeze-out and decay, asymmetric abundance, non-thermal mechanisms

 $M_p \sim 1 \text{ GeV}$

Standard Model

...to a hidden world with multiple states, new interactions

HIDDEN VALLEYS

Strassler, KZ 2006

Sub-weak Interactions (DM here.) -HC Standard Model Weak Interactions Inaccessibility

Dark World

Energy

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HIDDEN VALLEYS

Strassler, KZ 2006

Sub-weak Interactions (DM here.)

Torres del Paine



Dark World

Energy

Standard Model

Weak Interactions

THO

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Inaccessibility

Our Thinking Has Shifted: Why?

Simple, attractive, phenomenologically viable models exist

Example: ADM. Start with a single DM particle
 X, and one discovers you need more particles

$$n_X \sim 10^{-10} T^3$$

Our Thinking Has Shifted: Why?

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Example: ADM. Start with a single DM particle X, and one discovers you need more particles



 $n_X \sim 10^{-10} T^3$

Baryon and DM Number Related?

Standard picture: freeze-out of annihilation; baryon and DM number unrelated

Accidental, or dynamically related?

Experimentally, $\Omega_{DM} \approx 5\Omega_b$ Mechanism $n_{DM} \approx n_b$ $m_{DM} \sim 5 \ {\rm GeV}$



Asymmetric DM

"Integrate out" heavy state Higher dimension operators:

 $Xu^c d^c d^c$

Luty, Kaplan, KZ 0901.4117

$m_p \sim 1 \,\,{\rm GeV}$

Standard Model

Inaccessibility

Dark Matter (Hidden Valley)

X

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Astrophysical Implications

DM does not annihilate
It can accumulate in the center of stars
Notable case: neutron stars
Elastically scatter, come to rest in core
High density!

Altering Stellar Inter

 $\sigma_{\rm SI} \propto const.$



Exp. Implications of Dark Sectors

with dark forces
 Direct Detection
 Intensity experiments
 DM self-scattering and halo shapes

Direct Detection

Mediates _large_ scattering cross-sections

 σ_{SI}



$$\simeq \frac{g_n^2 g_{\chi}^2 m_r^2}{\pi m_{A'}^4}$$

~ $10^{-40} \text{ cm}^2 \left(\frac{g_n g_{\chi}}{10^{-4}}\right)^2 \left(\frac{8 \text{ GeV}}{m_{A'}}\right)^4$

0

Simplified model gives rise to many effects

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Connection to Intensity Experiments

Dark sectors may be more efficiently
 produced in low energy intensity experiments

Once above mass scale of mediator, production x-sect scales as $\sigma \sim \frac{g^4}{F^2}$

Low energy, very intense beams generated increased sensitivity

 ${\it \odot}$ Prefer beam energy sitting on mass of mediator $E \sim m_M$

Connection to Intensity Experiments

X



DM Interactions and DM Halos

 Dark matter selfinteractions randomize momenta and isotropize halos

Lead to lower density dark matter halo cores

Dark matter halos (including baryon poor dwarf galaxies) seem to have cores rather than cusps (still controversy as to cause)





Implies Dark Forces!

Very big scattering cross-sections $\sigma/m_X \sim 0.1 \text{ cm}^2/\text{g} \simeq 0.2 \times 10^{-24} \text{ cm}^2/\text{ GeV}$ ($\sigma_{weak} \sim 10^{-39} \text{ cm}^2$) Fits well with new models of DM! $\sigma_T \approx 5 \times 10^{-23} \text{ cm}^2 \left(\frac{\alpha_X}{0.01}\right)^2 \left(\frac{m_X}{10 \text{ GeV}}\right)^2 \left(\frac{10 \text{ MeV}}{m_{\phi}}\right)^4$ Range of dynamics much bigger than
 previously thought

Particle imprints on DM halos

Bound State Scattering



 $\begin{bmatrix} 100 \\ 10 \\ 10 \\ 10 \\ 10^{-2} \\ 10^{-3} \\ 10^{-4} \\ 10^{-5} \\ 10 \\ 10 \\ 10 \\ 10 \\ 10 \\ 100 \\ 100 \\ v (km/s) \\ \end{bmatrix}$

Tulin, Yu, KZ, 1210.0900 Tulin, Yu, KZ, 1302.3898

Quantum Resonances and Strongly Coupled Dynamics in DM Halos

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Terra Incognita





New Phases of Matter

• Nuclear Recoils are fundamentally kinematically limited $E_D \simeq q^2/(2m_{e,N}) \qquad q \sim m_X v$

 30 MeV DM corresponds to 1 eV of energy deposit on nucleons

Electron targets extract more energy

Material Gap

- But target electrons have a gap
- Semi-conductors -- silicon, germanium [CDMS] -- 1 eV gap; Nuclei -- at least 10's of eV
- Does not allow to detect DM lighter than 1 MeV



Need new kinematics...

The second state is the second state of the second state in the second state is the second state of the second state is the second state of the second state is the second state of th

VS

 $E = \frac{1}{2}m_X v^2$

How do we extract all the kinetic energy?

Superconductors have the needed features

Need a nearly gapless material

- Metals are gapless (conduction electrons)
- But also very susceptible to thermal vibrations

Superconductors are perfect: meV gap decouples phonon vibrations from electrons

> Y. Hochberg, Y. Zhao, KZ, 1504.07237 Y. Hochberg, M. Pyle, Y. Zhao, KZ 1512.07630

Target Fermi Velocity

Allows to extract entire
DM kinetic energy $m_X = 1 \text{ keV} \leftrightarrow E_{\text{kin}} = 1 \text{ meV}$

Non-rel (nuclear or electron) target: deposited energy reduced by target mass

$$E_D \simeq \frac{1}{2} \left(\frac{q^2}{m_T} + 2\vec{q} \cdot \vec{v}_{i,T} \right)$$

In metal: $v_F \sim 10^{-2}$

Superconductors and Dark Matter

Ordinary metal undergoes phase transition as temp is cooled

Energetically favorable for electrons in pair up; gap appears



The Idea

 DM scatters with electrons in Cooper pair. If energy deposited is greater than meV, break Cooper pair and create quasi-particles. Detect quasi-particle.

 \bigcirc

Quasiparticles

 \bigcirc

Cooper Pair

Down to the Warm Dark Matter Limit?

- Energy in resulting quasiparticles must be collected
- Concentrate
 quasiparticles onto heat
 sensor
- Heat sensor = TES

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Superconducting Substrate (AI)

SuperConducting Bias Rails (AI)

QuasiParticle Collection Pads (Au)

Insulating layer

TES (W)

Transition Edge Sensor

- Superconducting heat sensor doped at superconducting transition
- Already in use in microwave, xray and DM applications (SPT, ACT, SuperCDMS)



Need energy resolution sufficient to detect meV deposits -- not

there yet

TES	$T_c \; [\mathrm{mK}]$	Volume $[\mu m \times \mu m \times nm]$	Power Noise $[W/\sqrt{Hz}]$	$\sigma_E^{\rm now} \ [{\rm meV}]$	$\sigma_E^{\text{scale}} \text{ [meV]}$
W [18]	125	$25 \times 25 \times 35$	2.72×10^{-18}	120	1.1
Ti [19]	50	$6 \times 0.4 \times 56$	2.97×10^{-20}	47	22
MoCu [20]	110.6	$100\times100\times200$	4.2×10^{-19}	295.4	0.3

Y. Hochberg, M. Pyle, Y. Zhao, KZ 1512.07630

Astrophysically Feasible?



Y. Hochberg, Y. Zhao, KZ, 1504.07237 Y. Hochberg, M. Pyle, Y. Zhao, KZ 1512.07630

Summary

We have some good ideas about the DM sector. A couple of directions have become very well developed: SUSY and axions

New ideas and corresponding search strategies are developing.

Important to keep searches and ideas as broad and inclusive as possible

Summary

Dark Matter has not shown itself yet, but we continue to probe from all sides!





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