



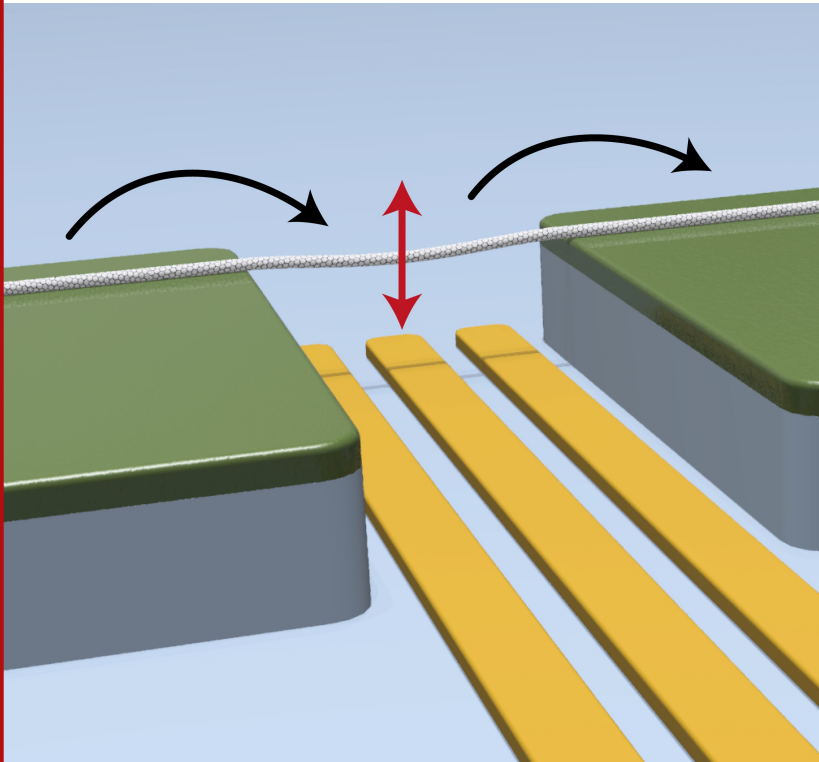
Edward Laird
Quantum electronic sensors

Physics

Lancaster
University



Sensing force and motion

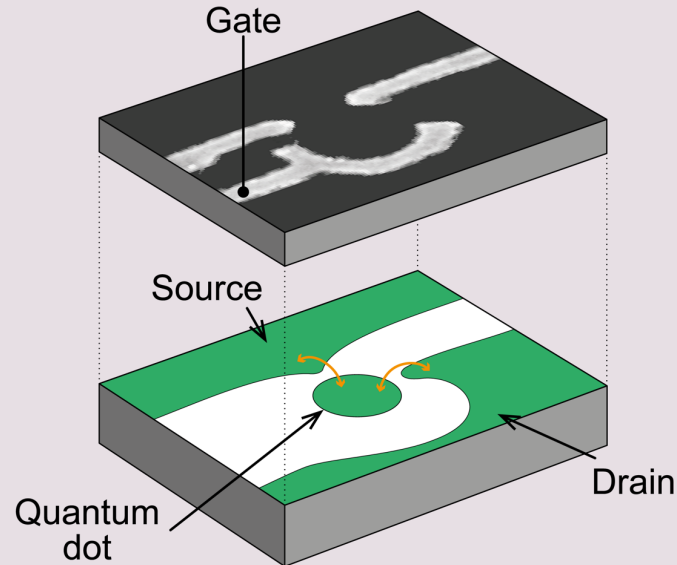


"A coherent nanomechanical oscillator driven by single-electron tunnelling"

Wen et al.

Nature Physics (2020)

Sensing qubits

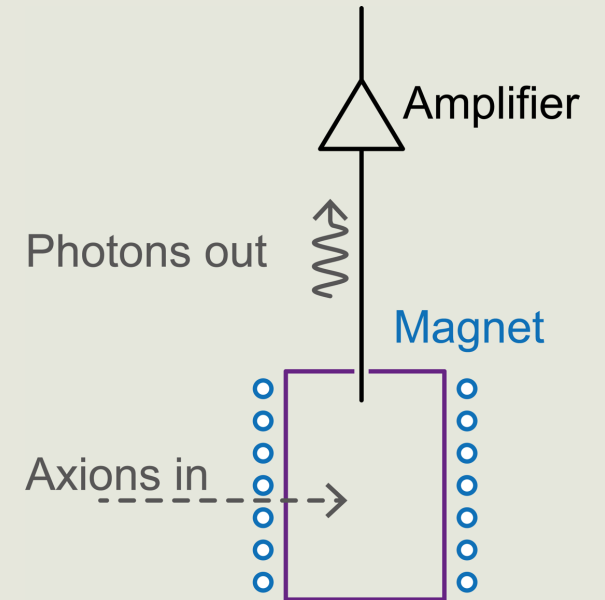


"Sensitive radio-frequency readout of quantum dots using an ultra-low-noise SQUID amplifier"

Schupp et al.

Journal of Applied Physics (2020)

Sensing dark matter



"Searching for wave-like dark matter with QSHS"

Bailey et al.

SciPost Physics Proceedings (2023)

Quantum Sensors for the Hidden Sector: An axion search using quantum electronics

Edward Laird

Physics Department, Lancaster University



Physics

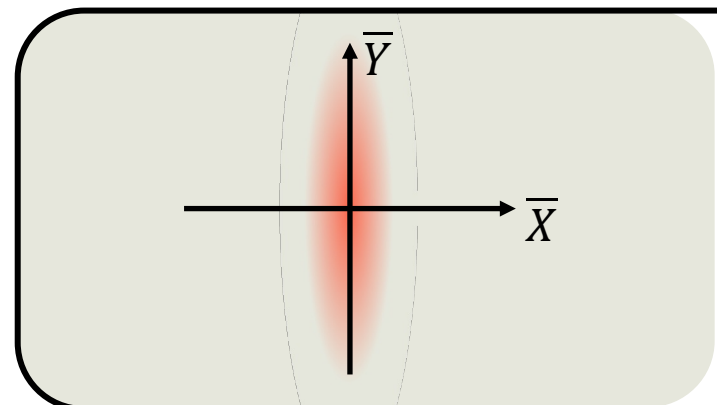
Lancaster
University



Quantum sensors for the hidden sector

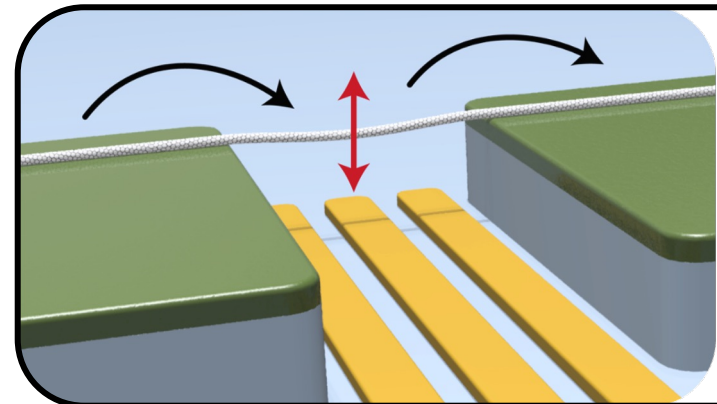


Why should you look for axions?



The QSHS experiment

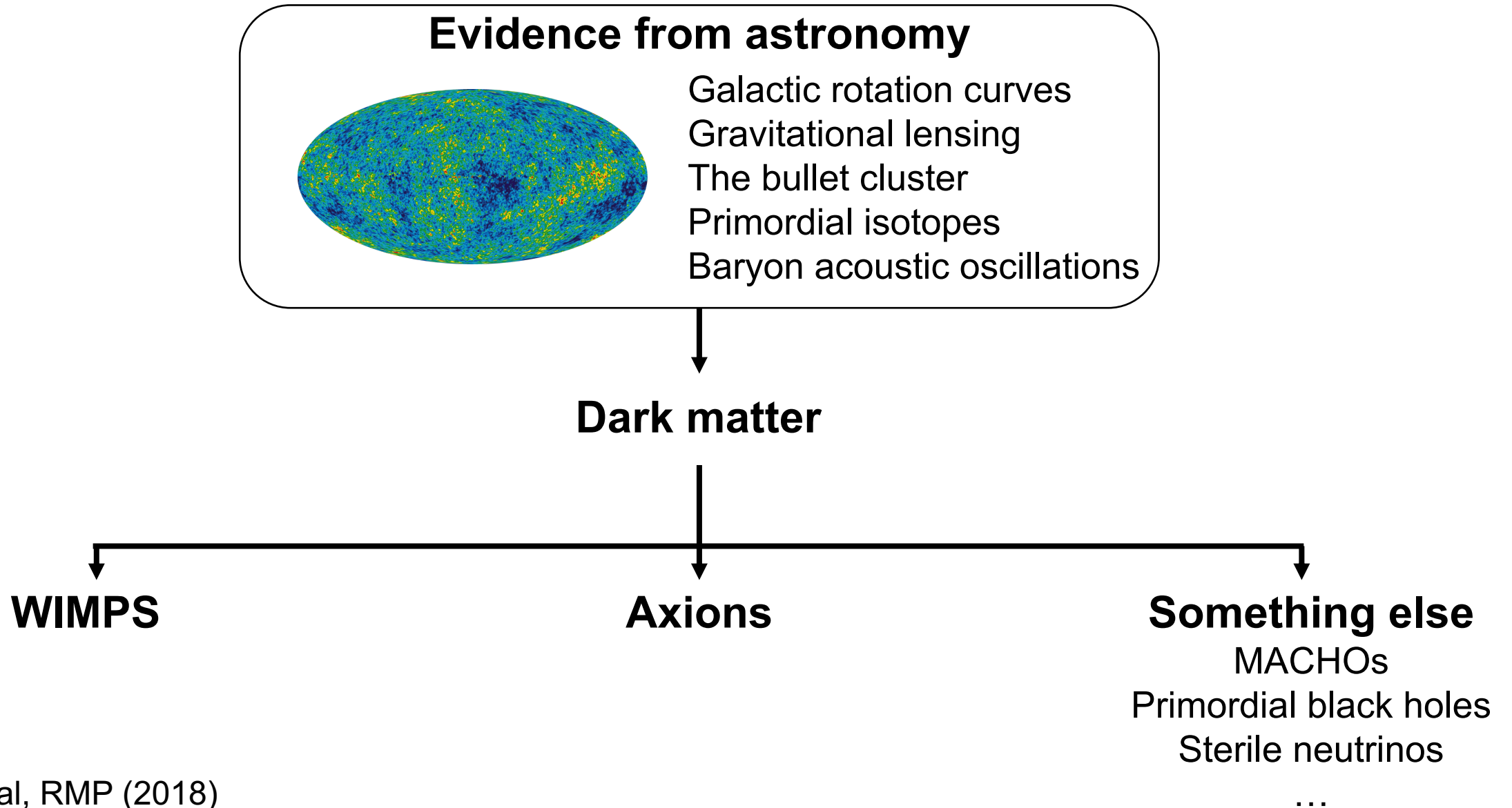
Looking for axions using quantum technology



Vibrating nanotubes in superfluid helium

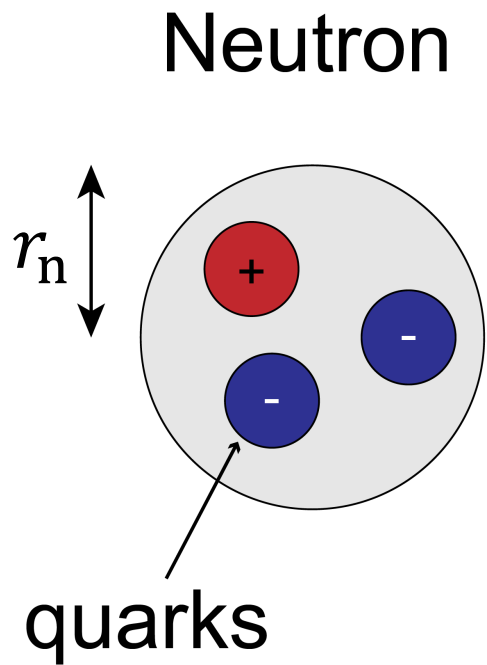
Another device for particle searches?

Is dark matter made from axions?

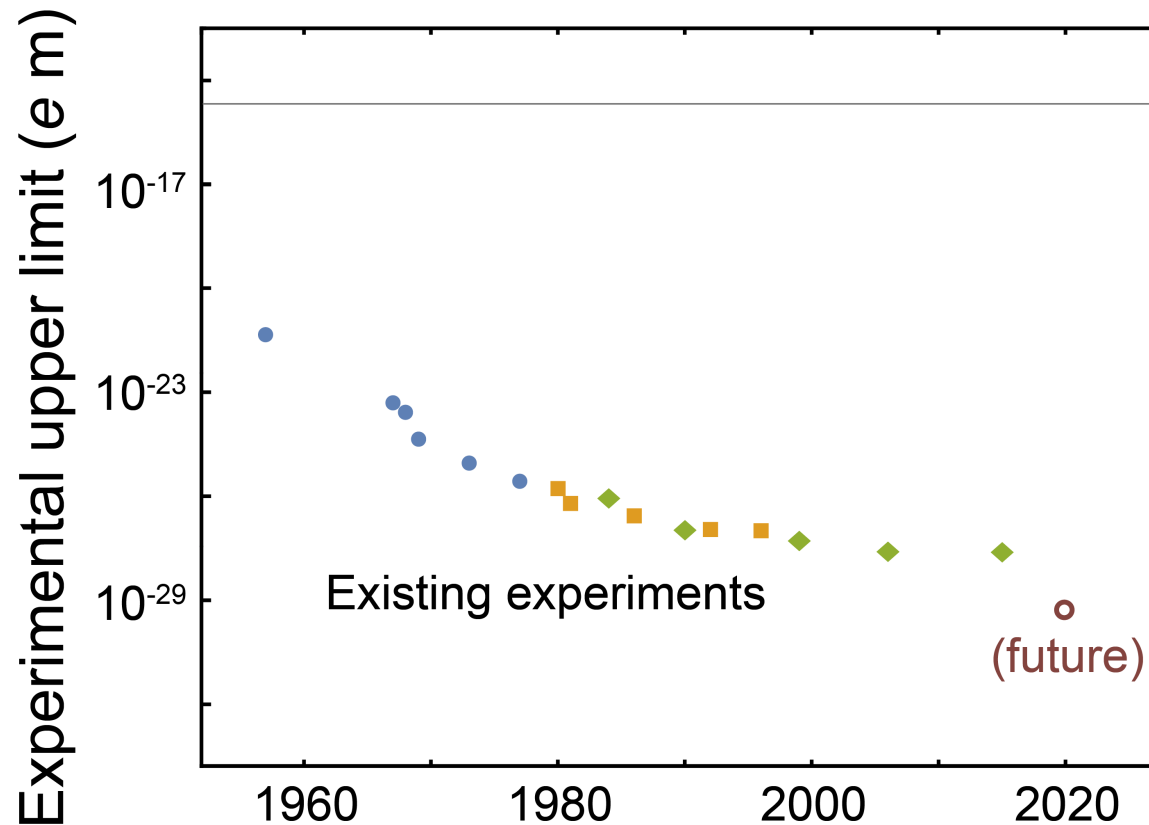


Why should you believe in axions?

The strong CP problem



$$d_{\text{expected}} \sim \frac{e}{3} r_n$$
$$\sim 10^{-15} e m$$



How axions would solve the strong CP problem

Lagrangian of strong nuclear force (simplified):

$$\mathcal{L} = -\frac{1}{4} G_{\mu\nu} G^{\mu\nu} + \theta \frac{\alpha_s}{8\pi} G_{\mu\nu} \tilde{G}^{\mu\nu}$$

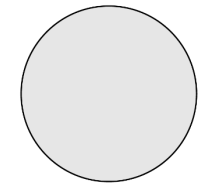
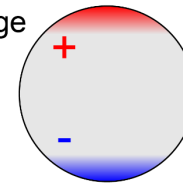
Preserves CP symmetry

Violates CP symmetry

“Dynamic field”

Neutron electric dipole moment

Electric charge



What we expect

What we get

If θ is a number, there is no reason for it to be near zero.

However, if $\theta(\vec{r}, t)$ is a field, then there is a mechanism in QCD theory (the Peccei-Quinn mechanism) that holds it near zero.

The axion: a hypothetical particle

Mechanism
(1977)

Particle
(1978)

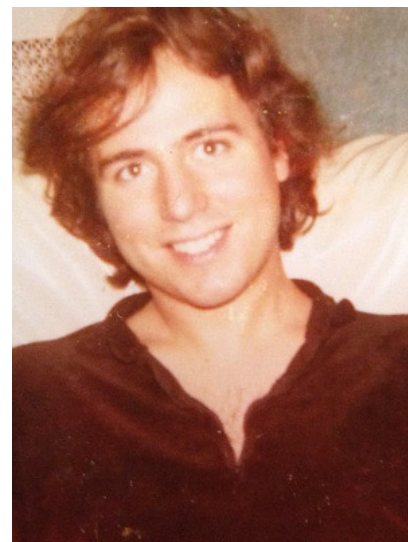


Roberto Peccei

Helen Quinn



Steven Weinberg

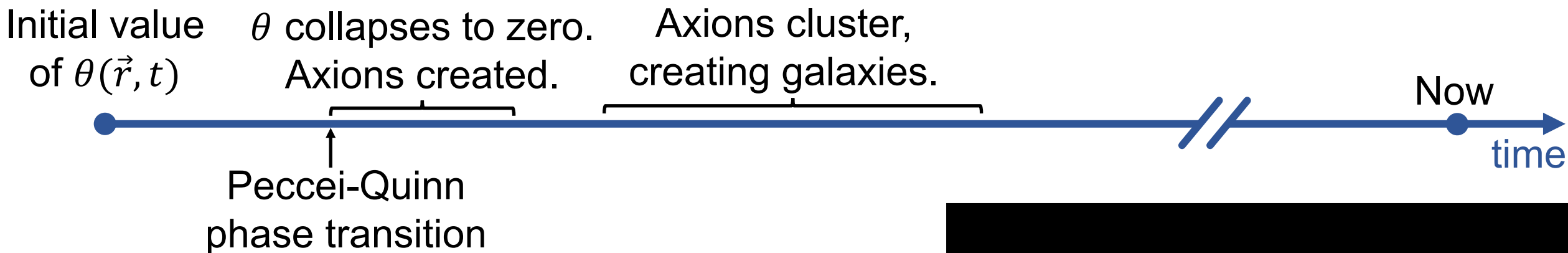


Frank Wilczek

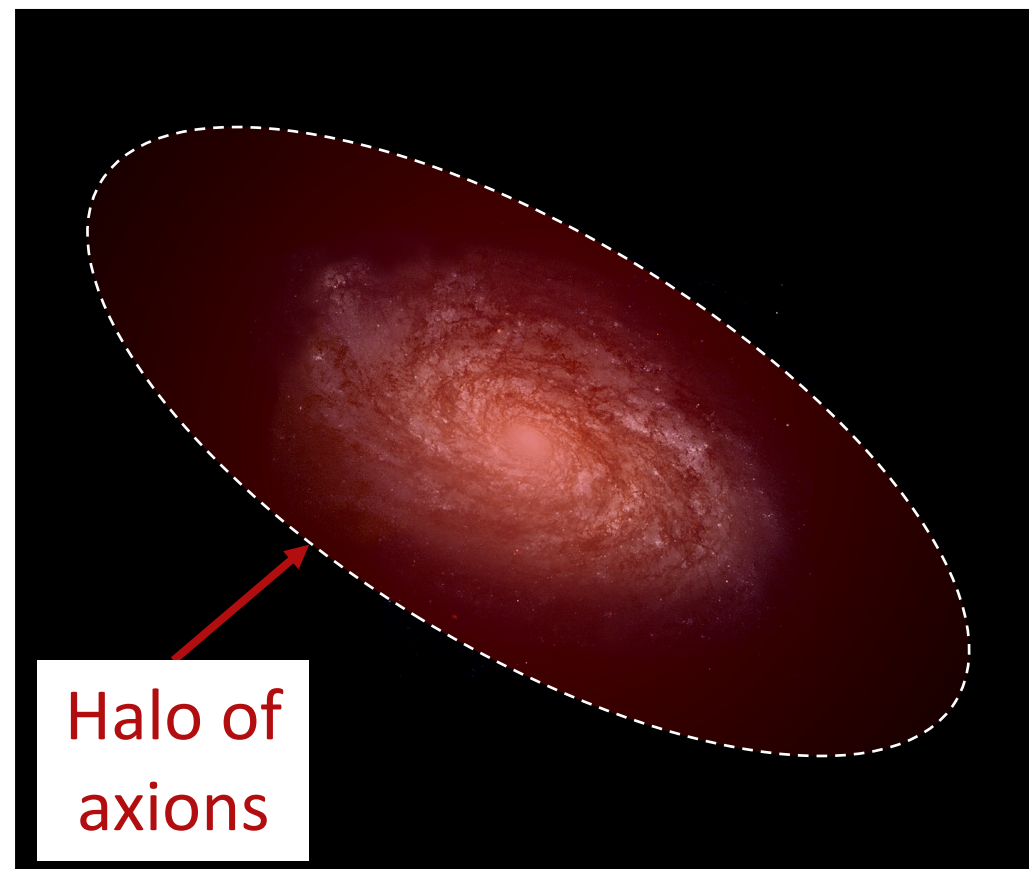


Name

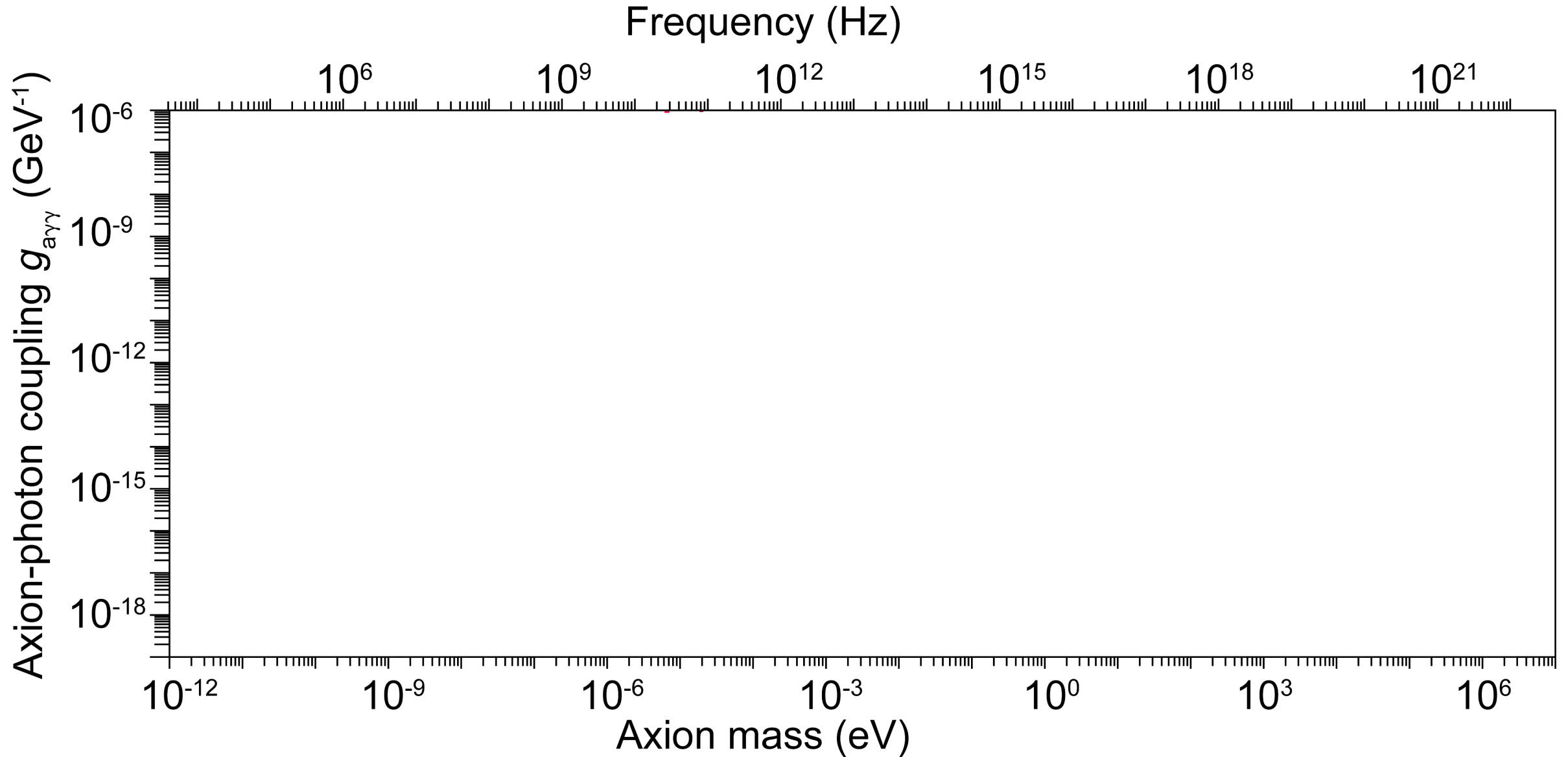
The axion model of dark matter



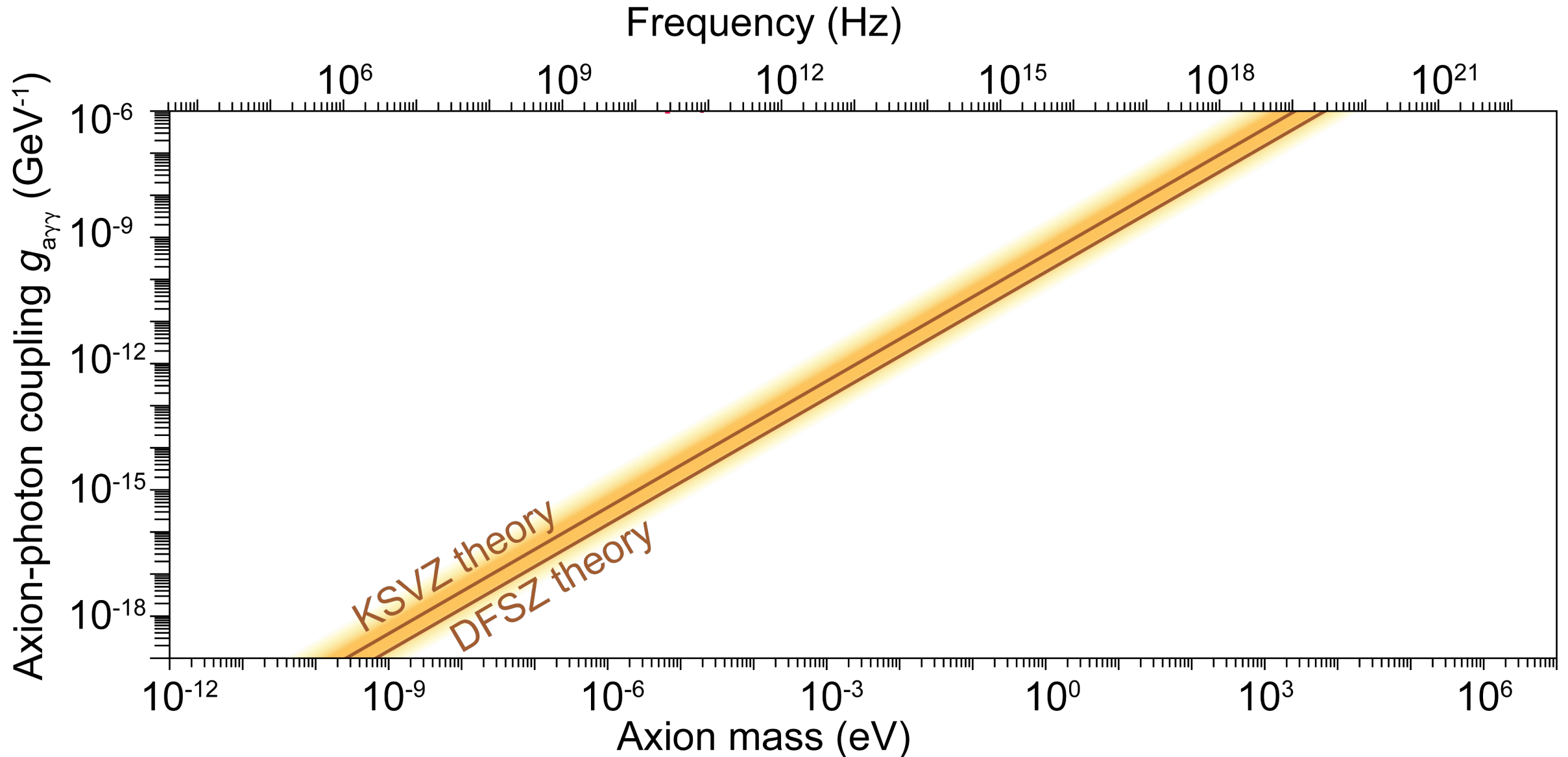
If this model is correct, then about 10^{24} axions pass through you every second (depending on the axion mass).



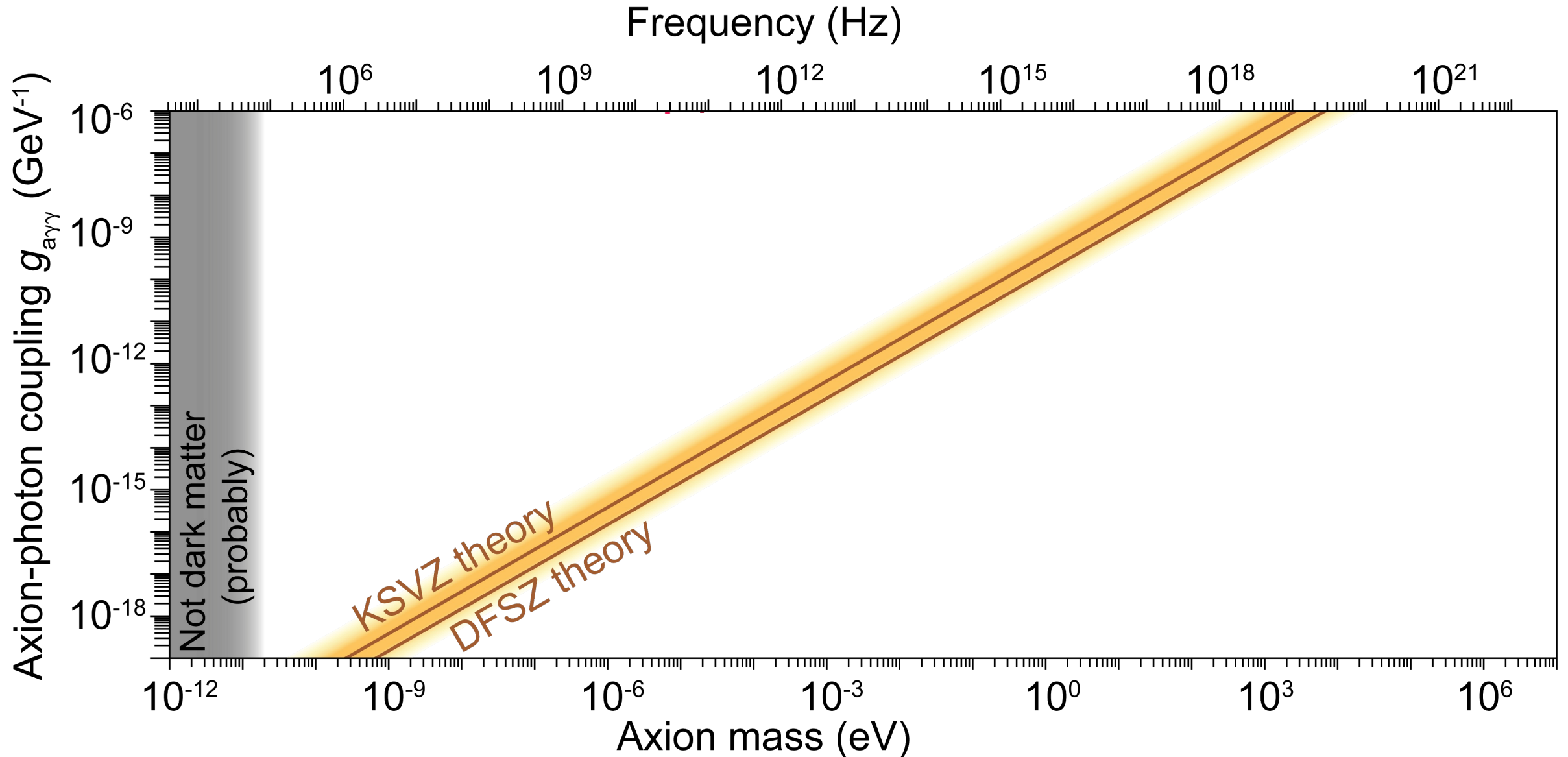
What do we know about the axion's properties?



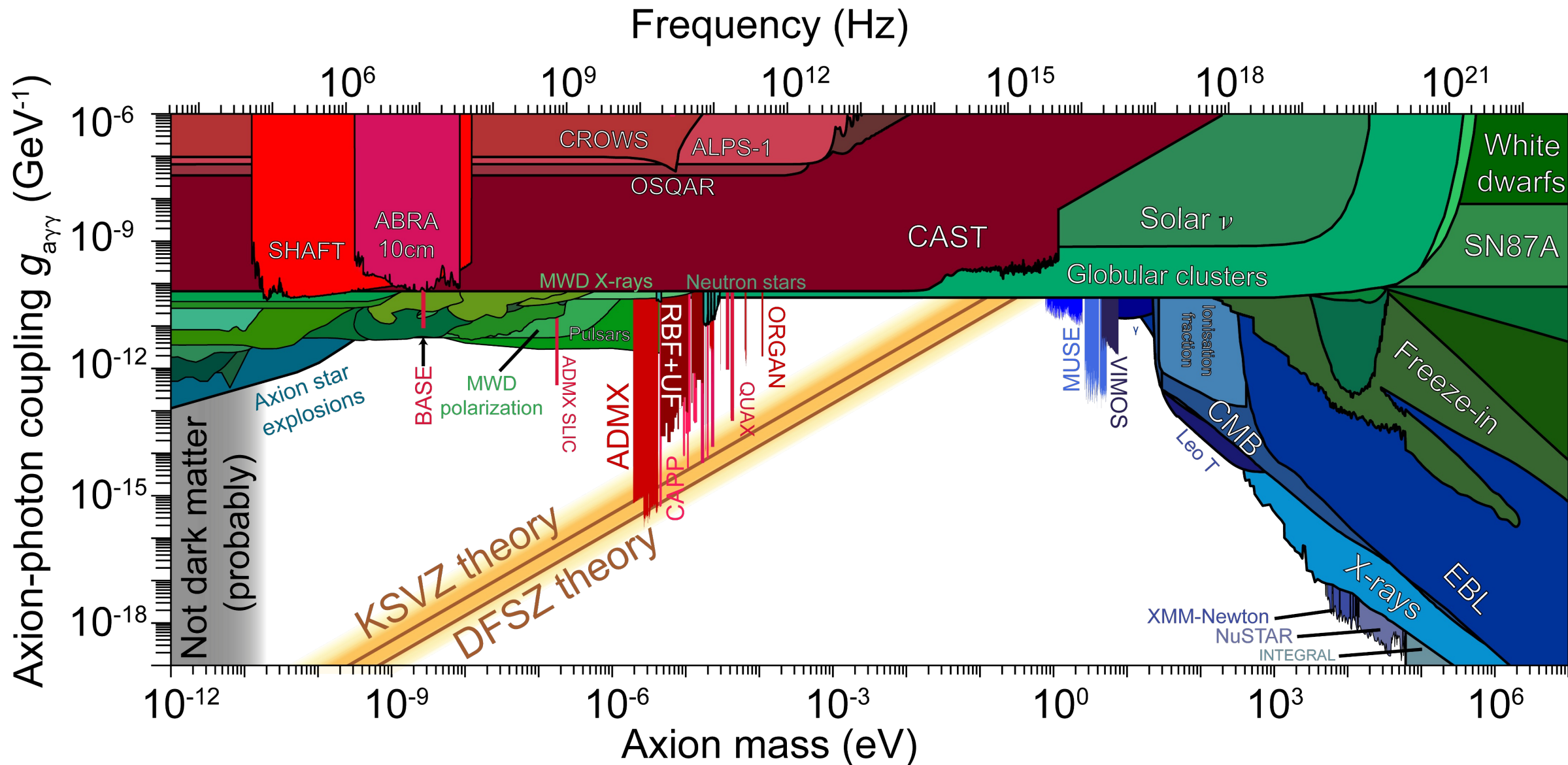
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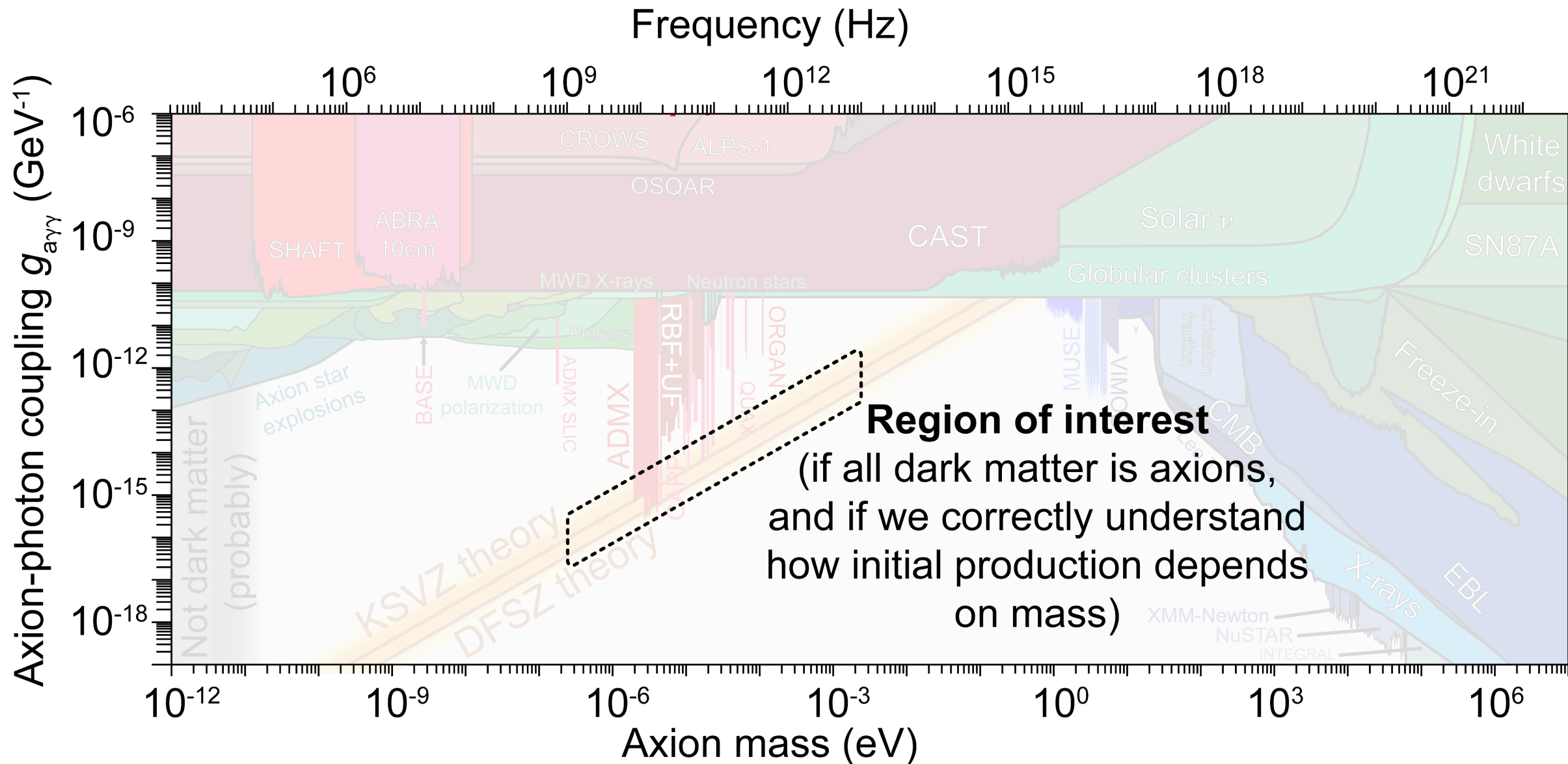
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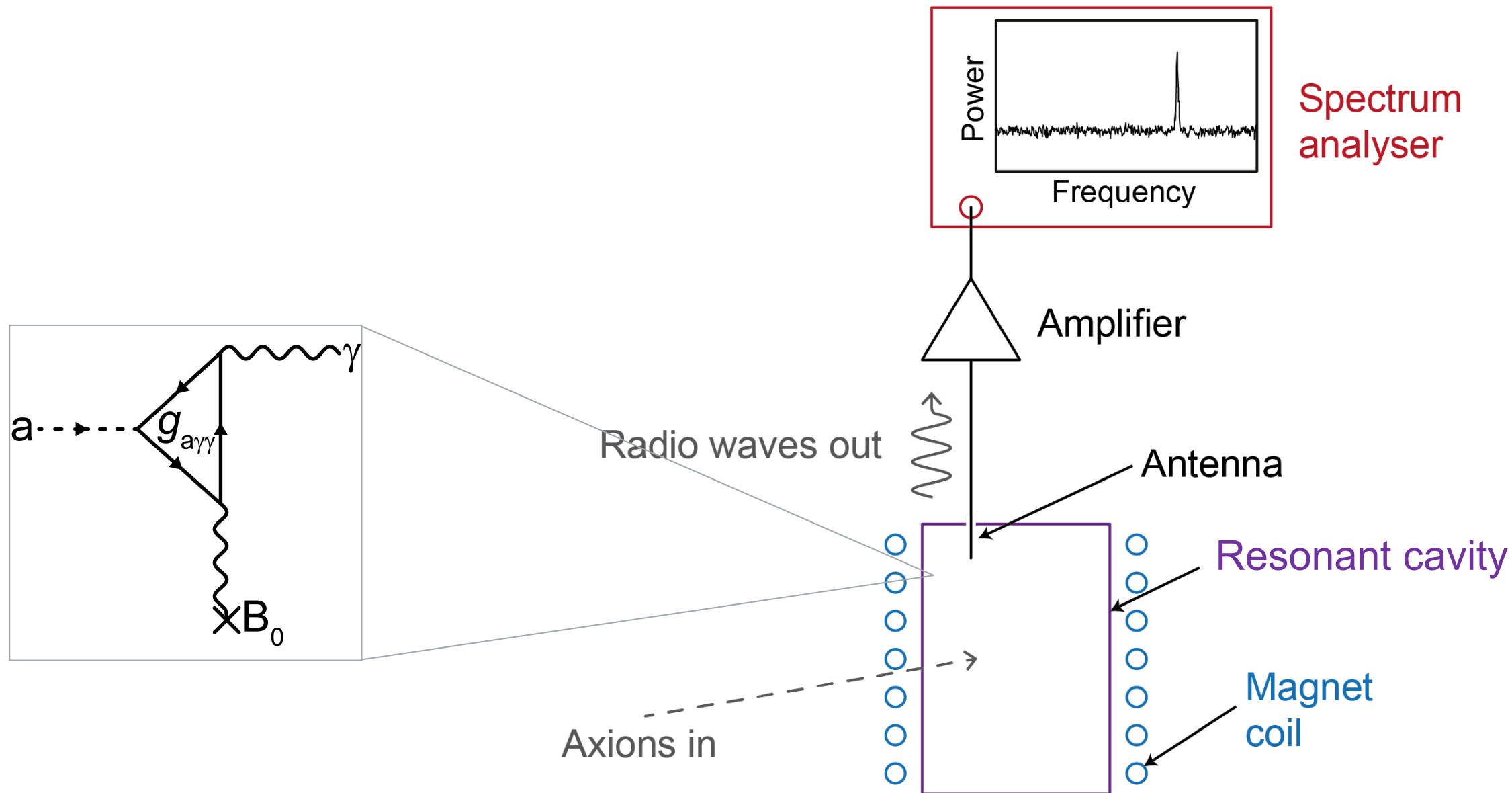
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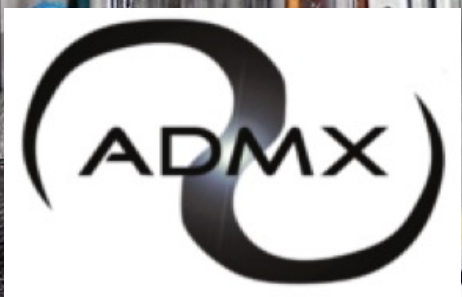
What do we know about the axion's properties?



Detecting axions using a haloscope



Example of a haloscope: the ADMX experiment (our collaborators)



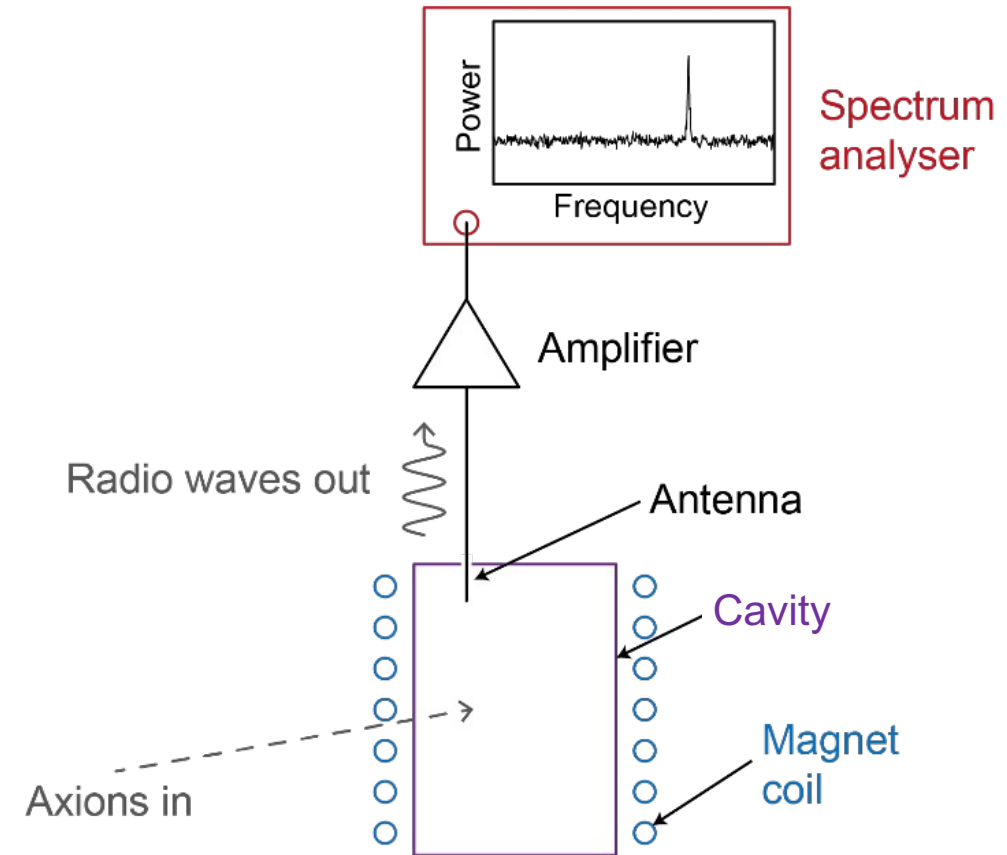
The challenge: measuring a weak radio signal

Expected signal power in a microwave haloscope:

$$P_{\text{axion}} = 2.2 \times 10^{-23} \text{ W} \\ \times \left(\frac{V}{0.136 \text{ m}^3} \right) \left(\frac{B}{7.6 \text{ T}} \right)^2 \left(\frac{C}{0.4} \right) \left(\frac{Q}{30,000} \right) \\ \times \left(\frac{g_\gamma}{0.36} \right)^2 \left(\frac{\rho_a}{0.45 \text{ GeV cm}^{-3}} \right) \left(\frac{f}{740 \text{ MHz}} \right)$$

Signal power in other applications:

- Mobile phone: 10^{-6} W
- Radar: 10^{-12} W



Classical measurement: the Dicke radiometer equation

The signal-to-noise ratio with which an axion peak can be identified is

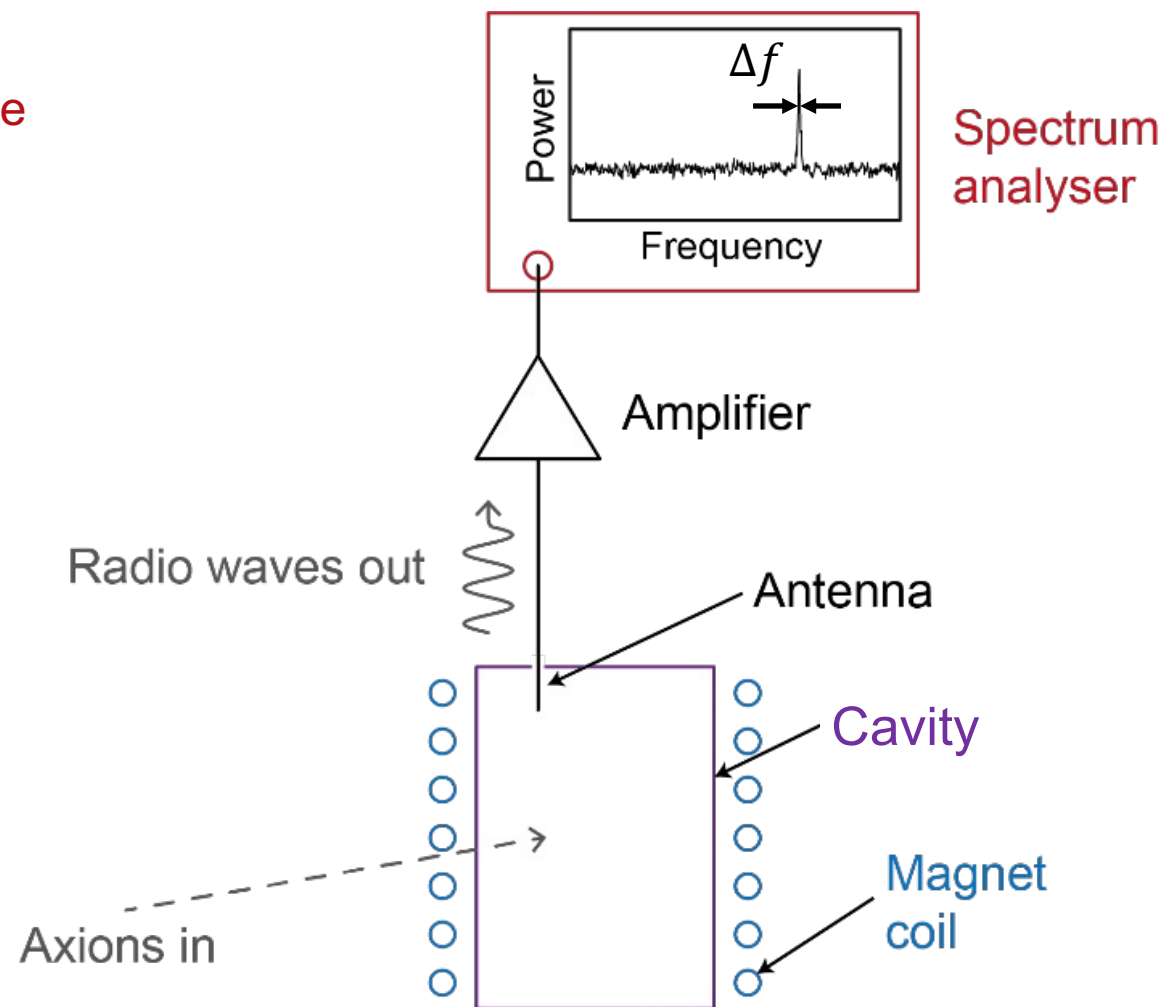
$$\frac{\text{Signal}}{\text{Noise}} = \frac{P_{\text{axion}}}{kT_N} \sqrt{\frac{\tau}{\Delta f}}$$

Labels for the equation:

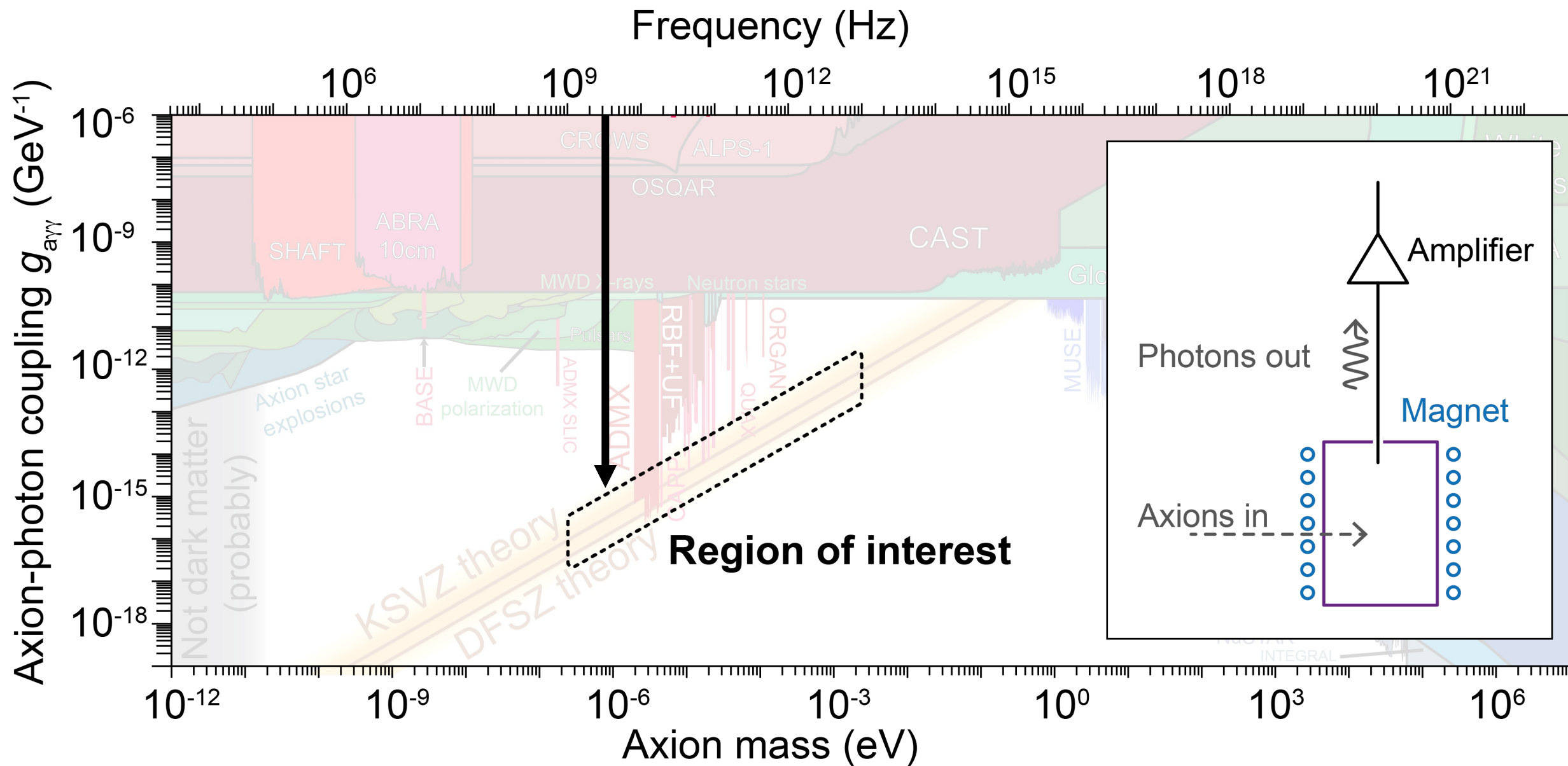
- Signal power: P_{axion}
- Measurement time: τ
- Boltzmann's constant: k
- System noise temperature: T_N
- Signal bandwidth: Δf

To discover an axion, we need to tune to the correct frequency and measure until

$$\frac{\text{Signal}}{\text{Noise}} \gtrsim 5.$$



Searching for axions with a haloscope



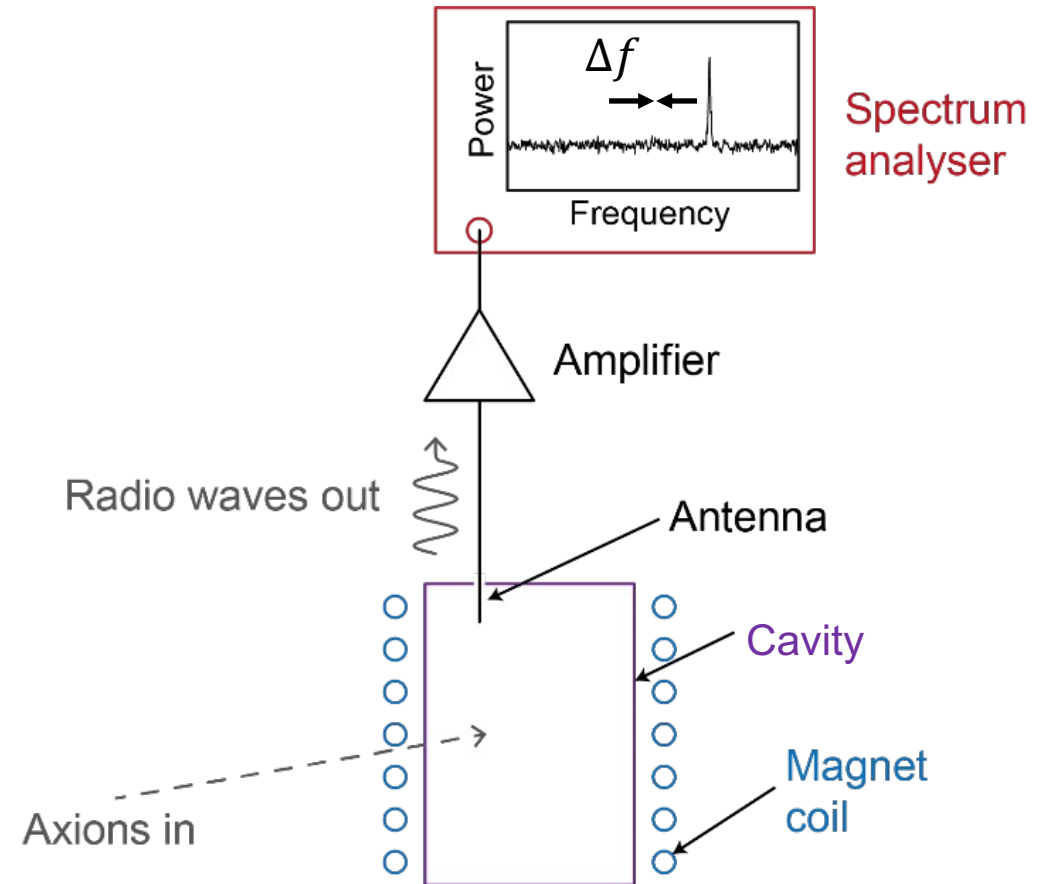
The scan rate and the standard quantum limit

The standard quantum limit for the system noise temperature with a phase-preserving amplifier is:

$$T_N > \frac{hf}{k}$$

Scan rate for an example detector (10L cavity volume, 8T field, 10GHz centre frequency):

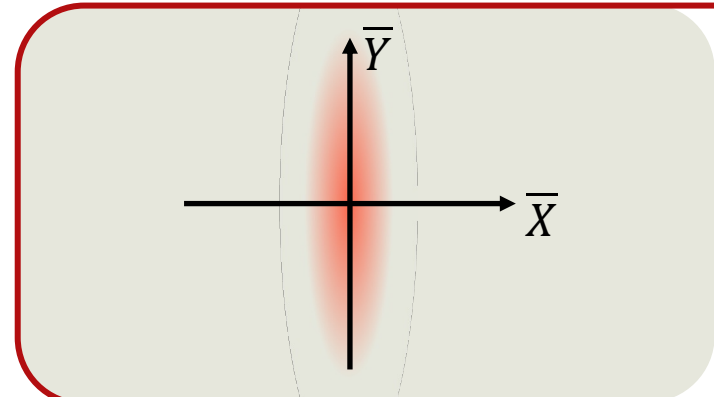
Noise temperature T_N	Scan rate
2 K (A good semiconductor amplifier)	90 yr/GHz
0.48 K (Standard quantum limit)	5 yr/GHz



Quantum sensors for the hidden sector

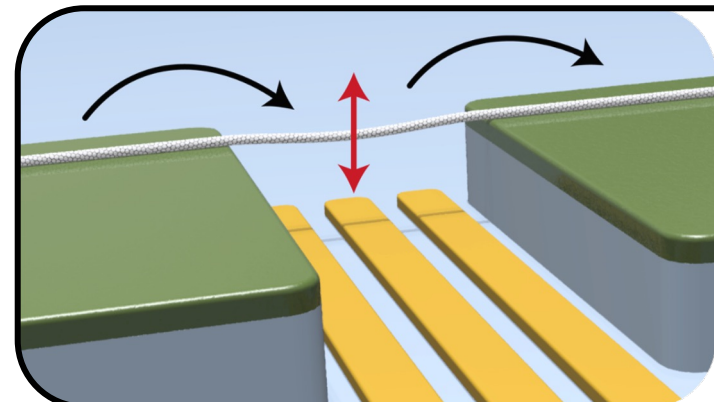


Why should you look for axions?



The QSHS experiment

Looking for axions using quantum technology



Vibrating nanotubes in superfluid helium

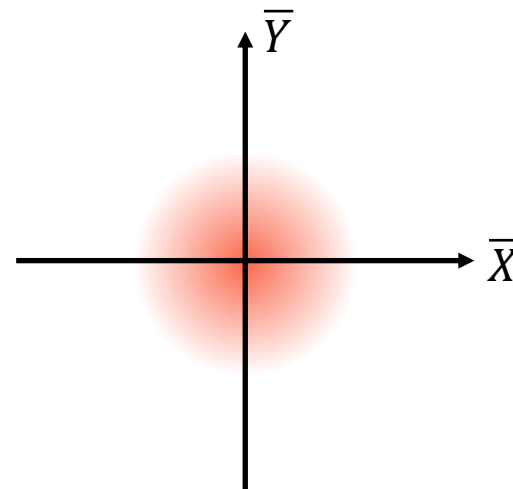
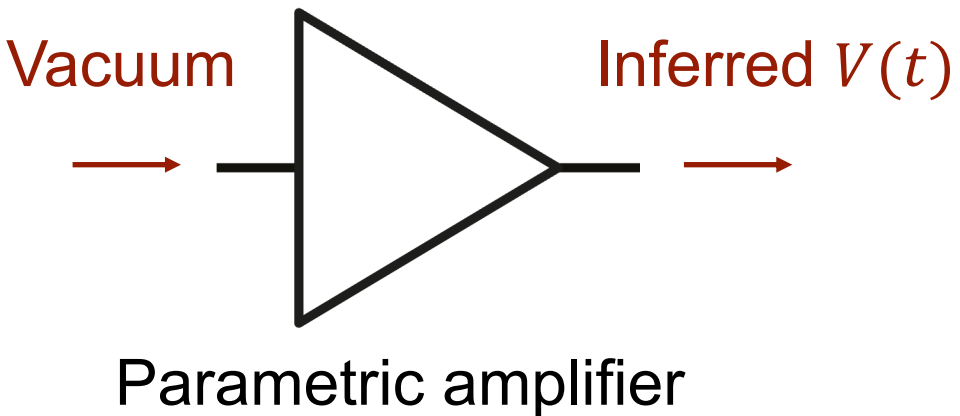
Another device for particle searches?

Squeezing using a parametric amplifier

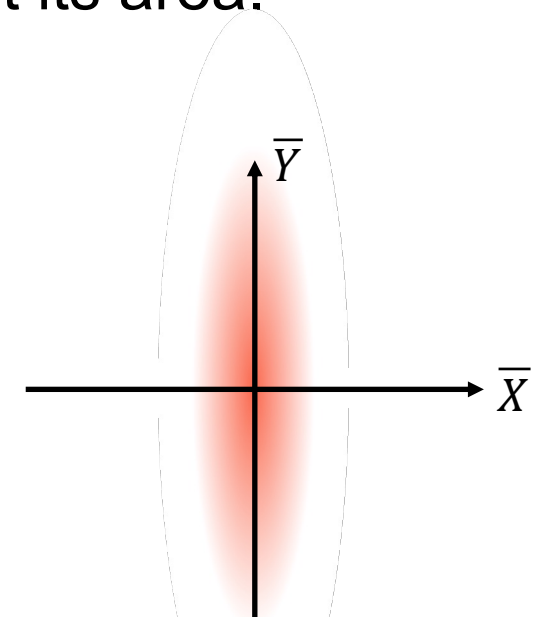
An oscillating voltage near frequency f has the form

$$V(t) = X(t) \cos 2\pi ft + Y(t) \sin 2\pi ft$$

A parametric amplifier can amplify one of these two quadratures while attenuating the other one. This changes the shape of the (\bar{X}, \bar{Y}) histogram, but not its area.

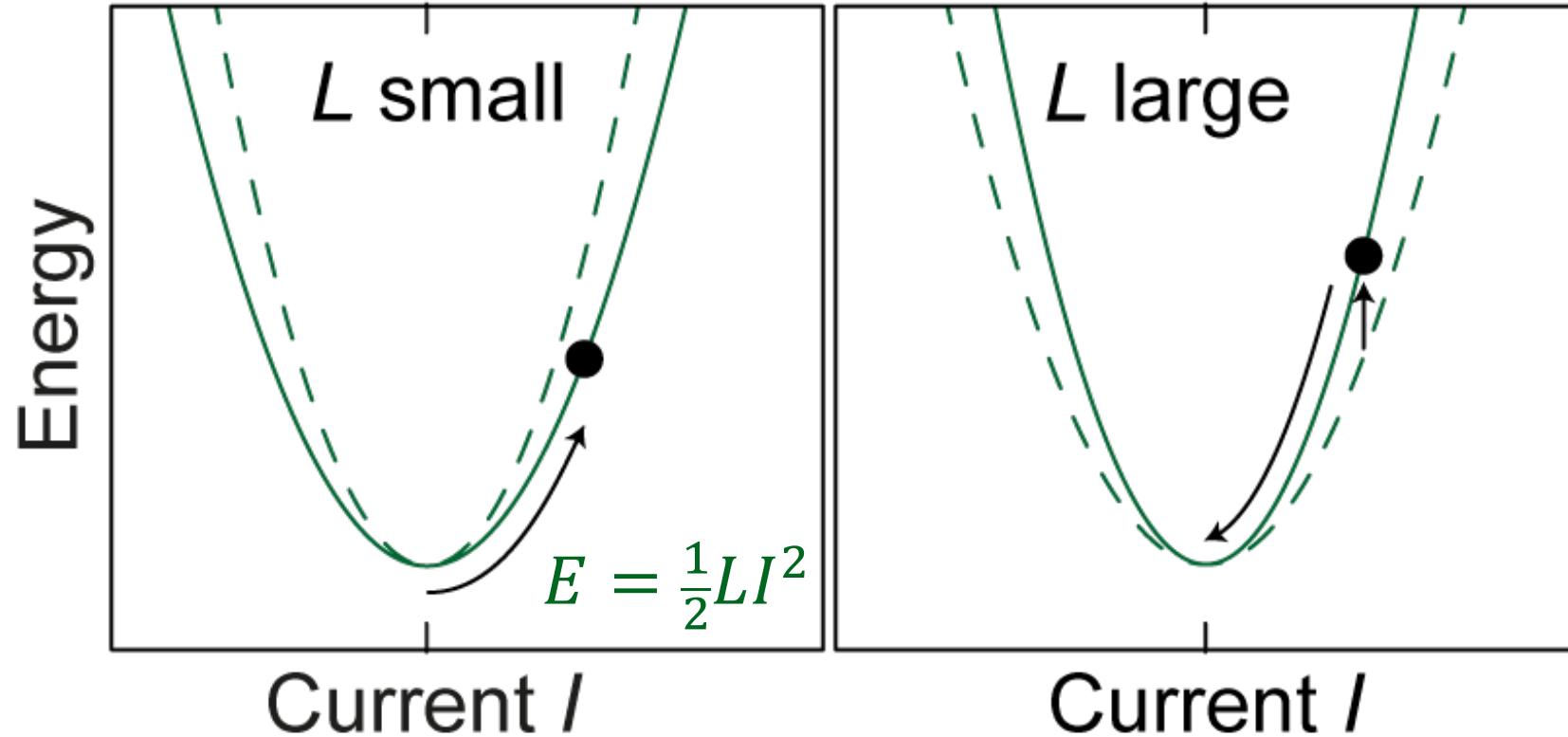
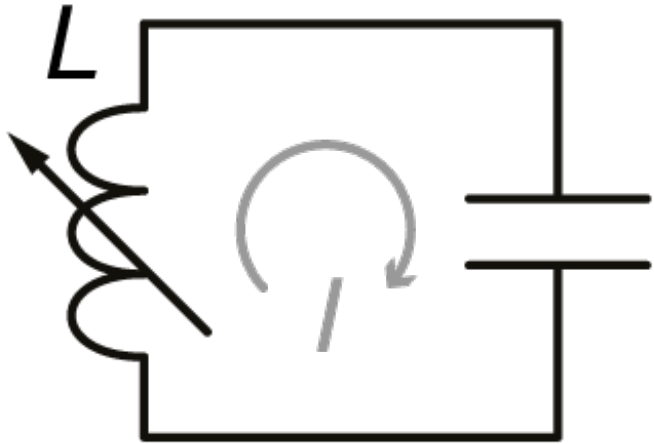


Inferred signal
without squeezing



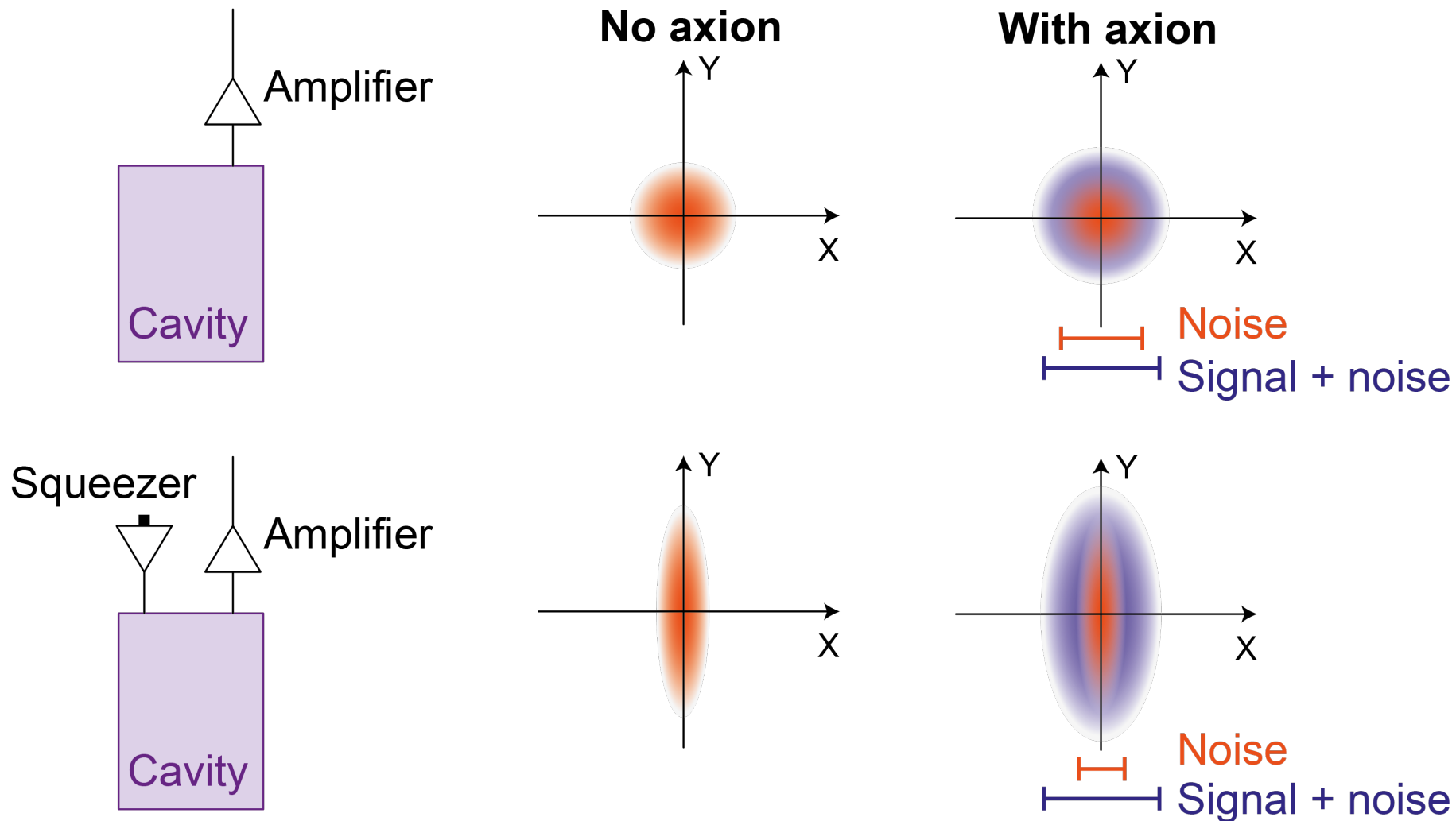
Inferred signal
with squeezing

Principle of parametric amplification



By modulating one parameter of a resonator (such as its inductance) at frequency $2f$, we can pump energy into the signal mode at frequency f .

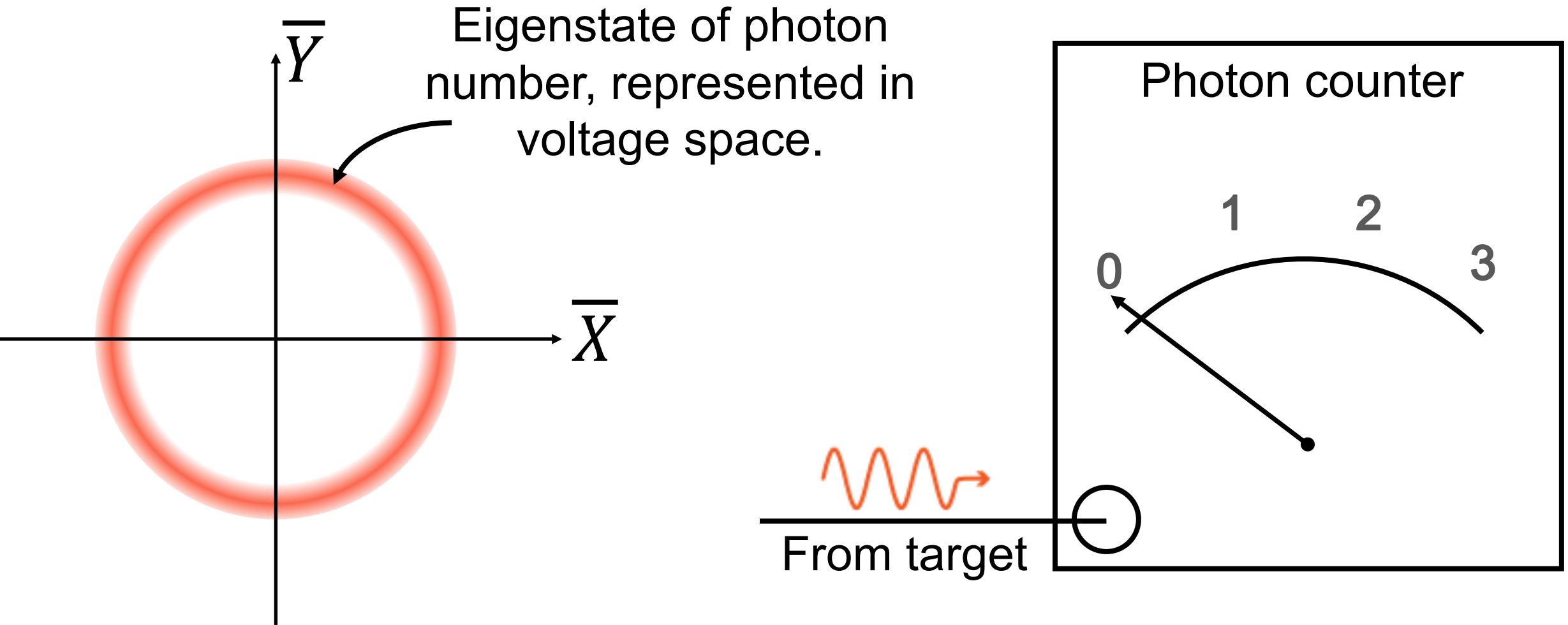
Quantum speed-up method 1: linear squeezing



From the HAYSTAC experiment at Yale

Malnou et al, PRX (2019); Backes et al, Nature (2021)

Quantum speed-up method 2: photon counting



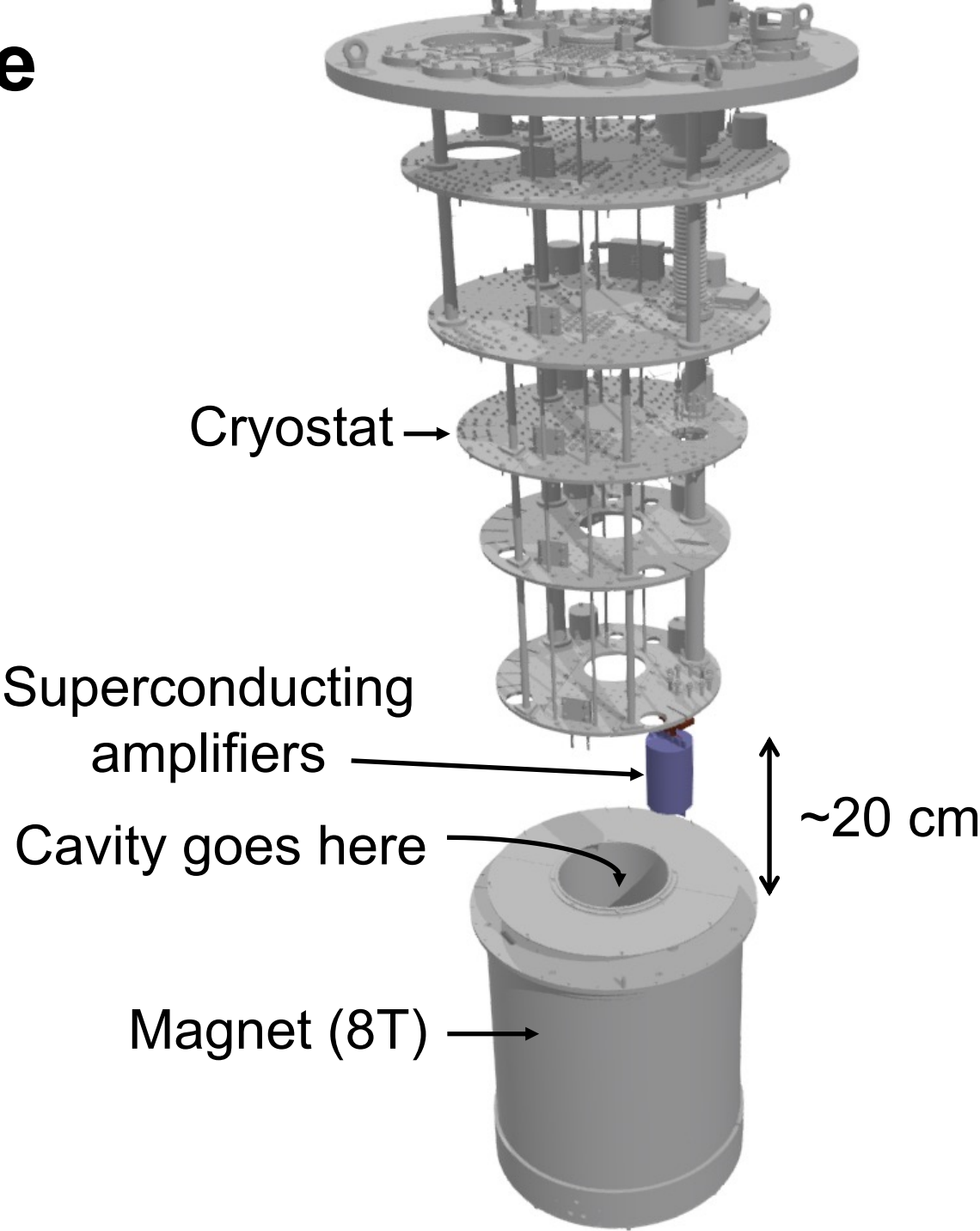
Magnetic field shielding for the QSHS haloscope

Magnetic field shielding is critical to run superconducting amplifiers close to the target cavity.

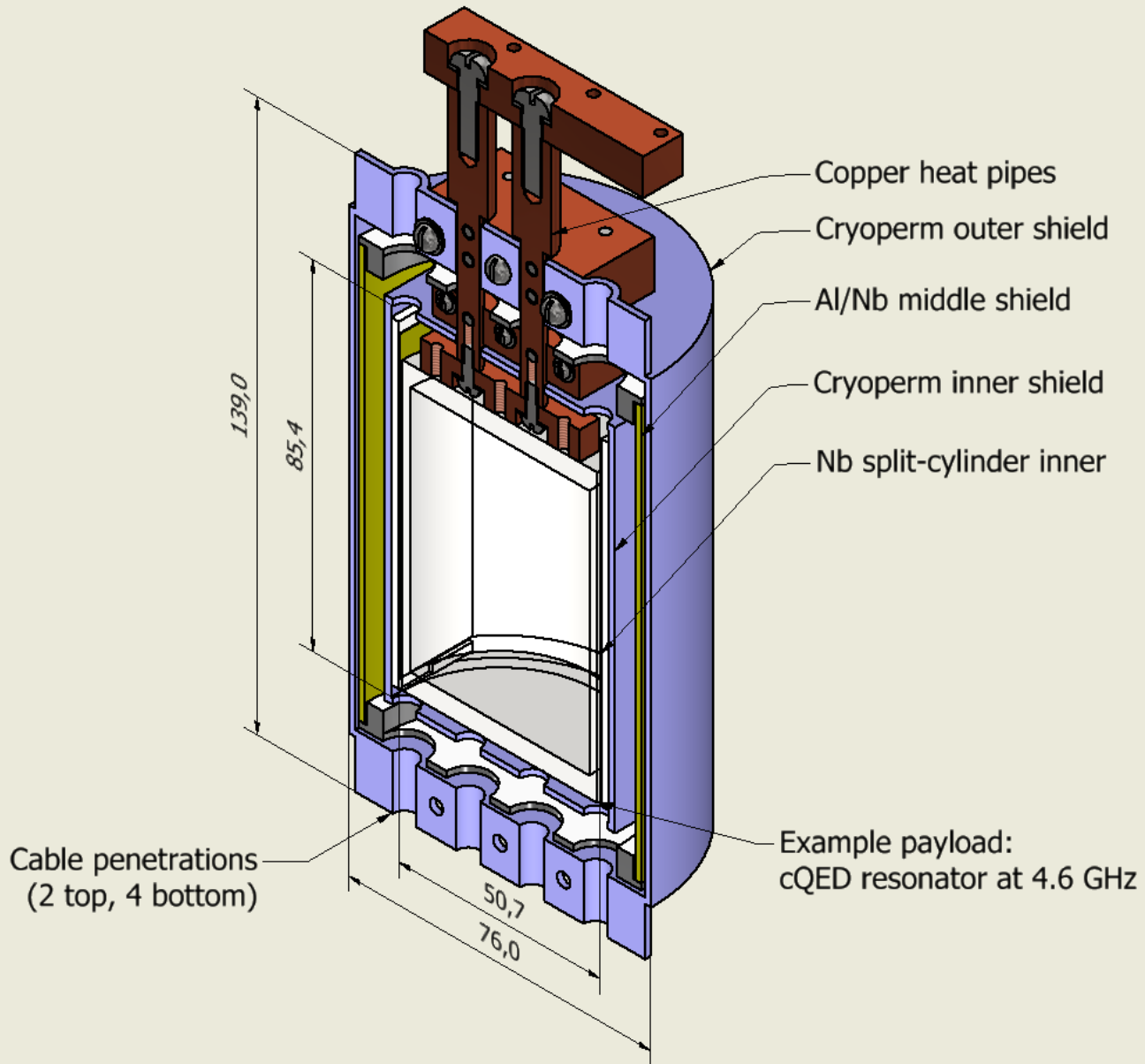
Stray field of main coil:
~10 mT

Operating field of Josephson amplifiers:
~10 nT

(See a thesis from B. Brubaker for information about how this is handled in the HAYSTAC experiment.)



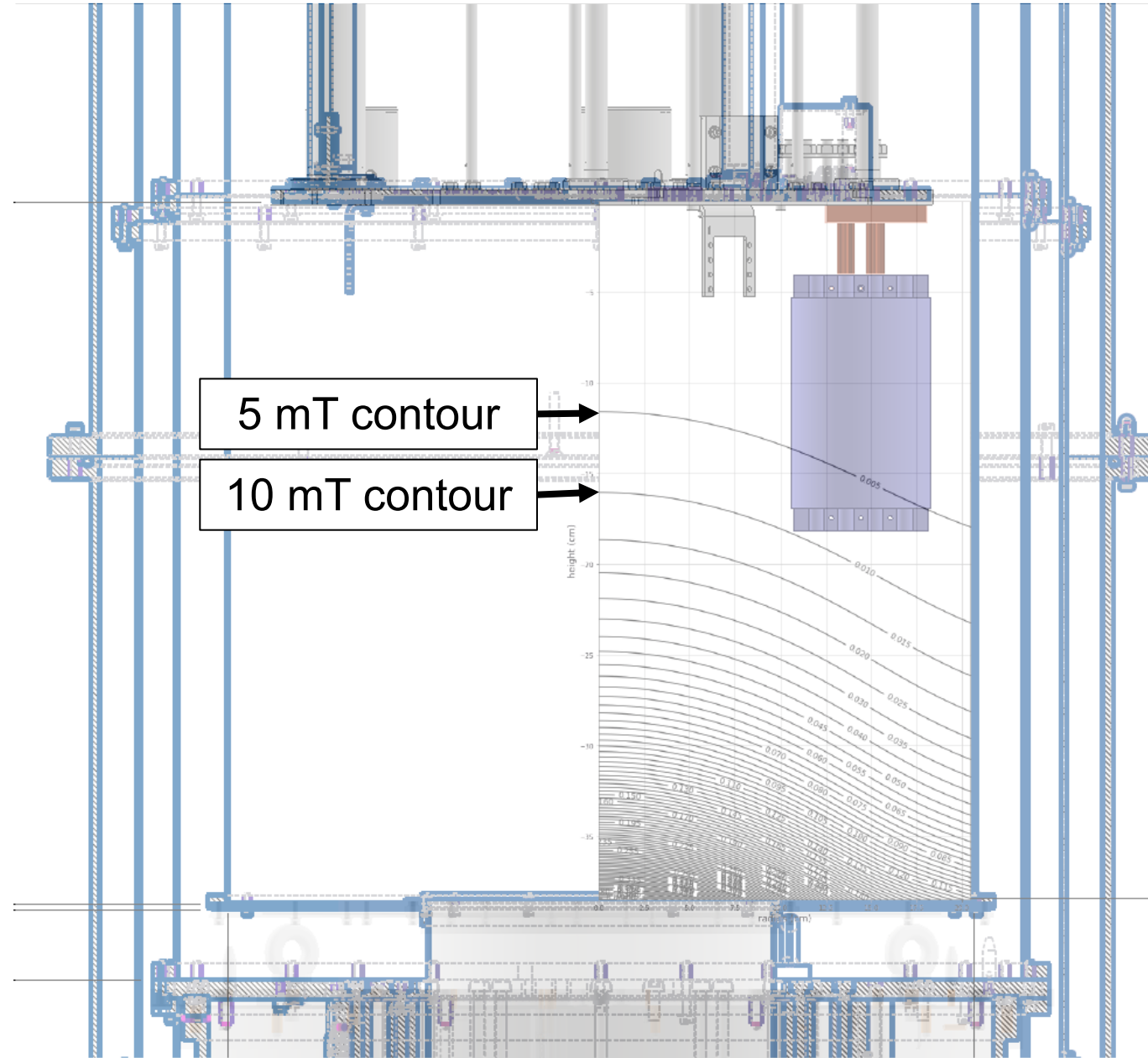
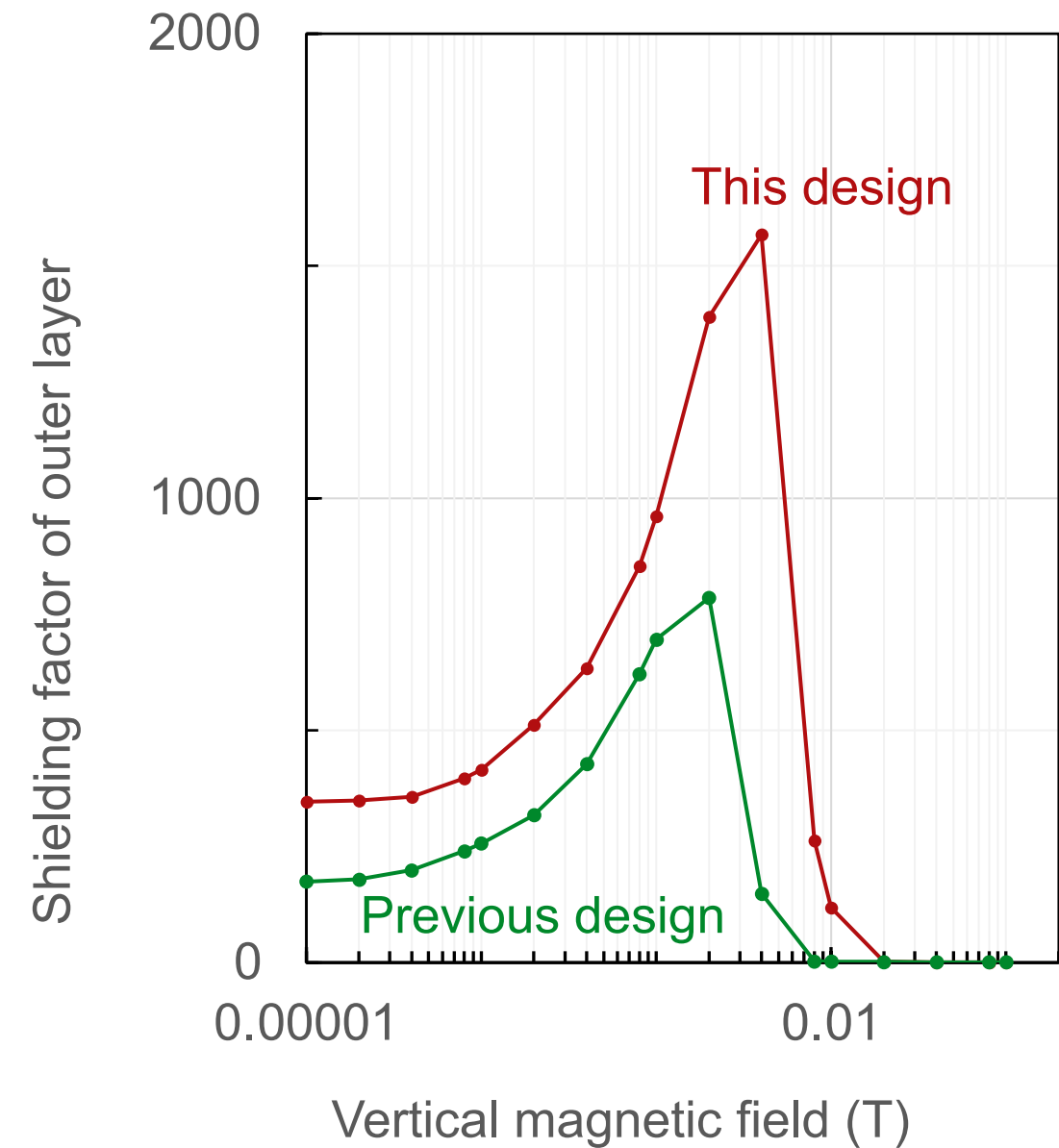
The QSHS field shield



Design criteria:

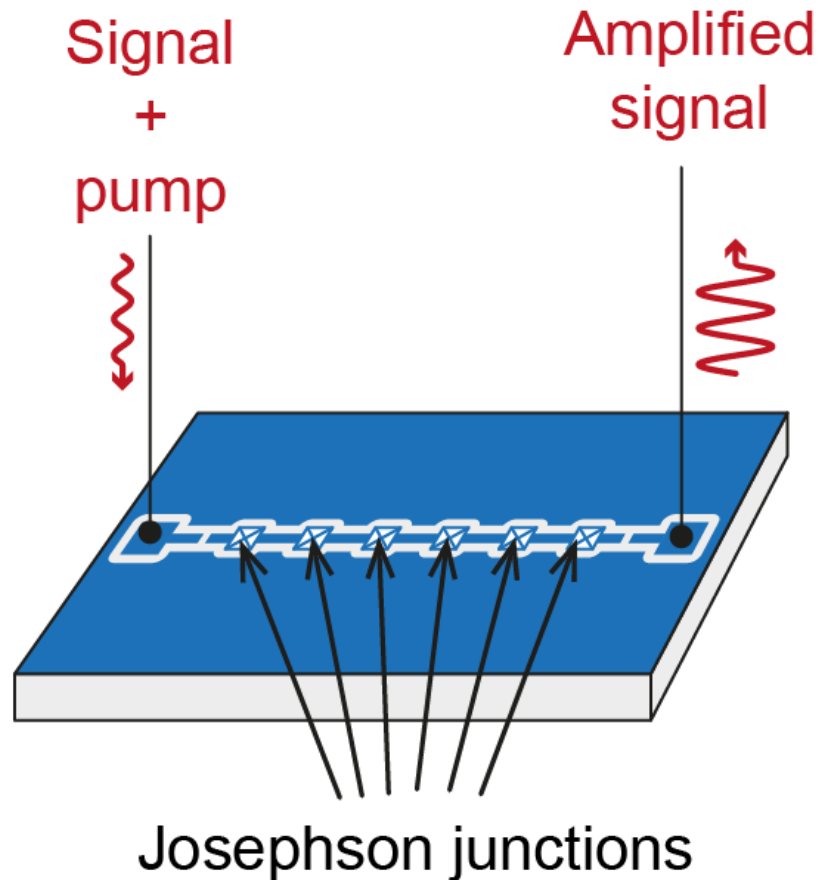
- Field shielding factor of 10^6 in 10 mT
- Easy access for sample exchange, including a clam-shell design.
- Sufficient sample volume (here 0.17 L)
- Small enough to be tested in our vector magnets.

Simulation results

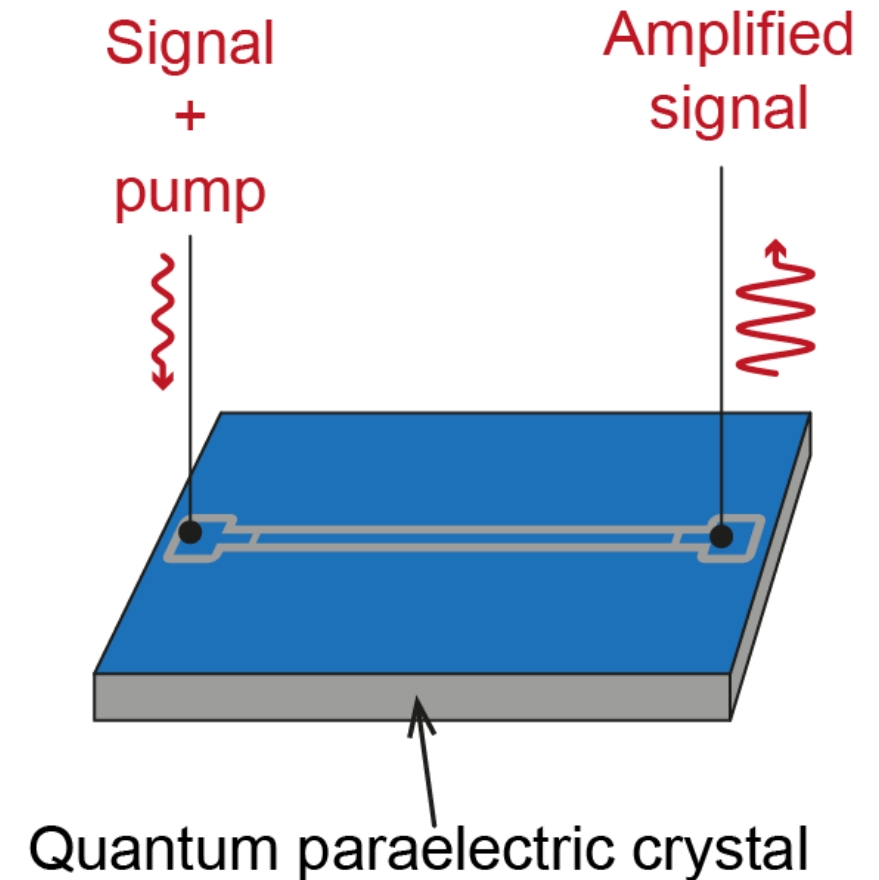


A parametric amplifier that is robust against magnetic fields

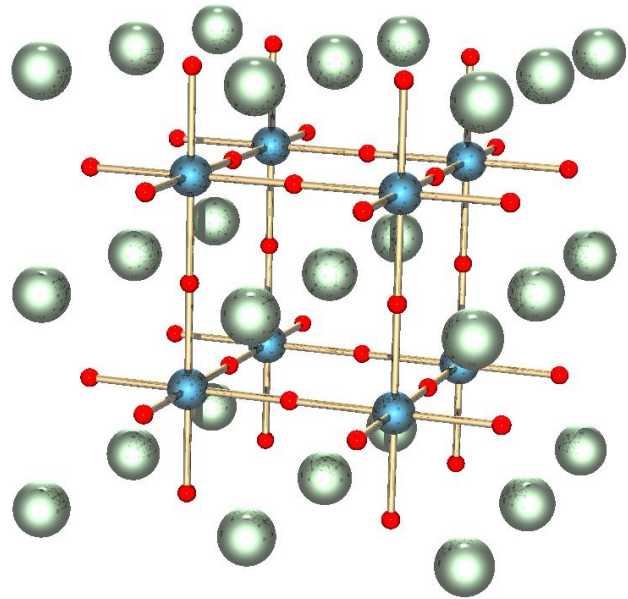
Josephson parametric amplifier



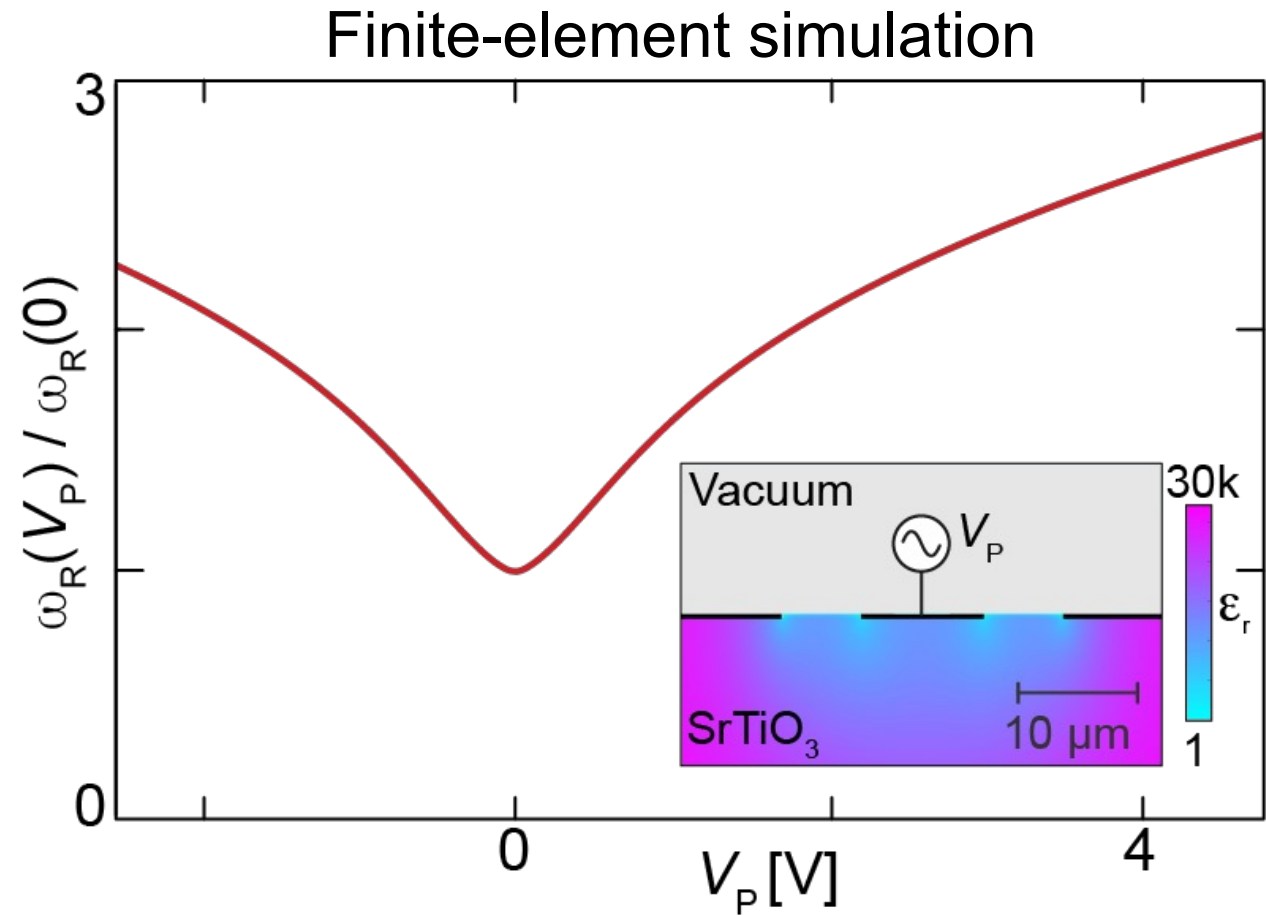
Quantum paraelectric amplifier



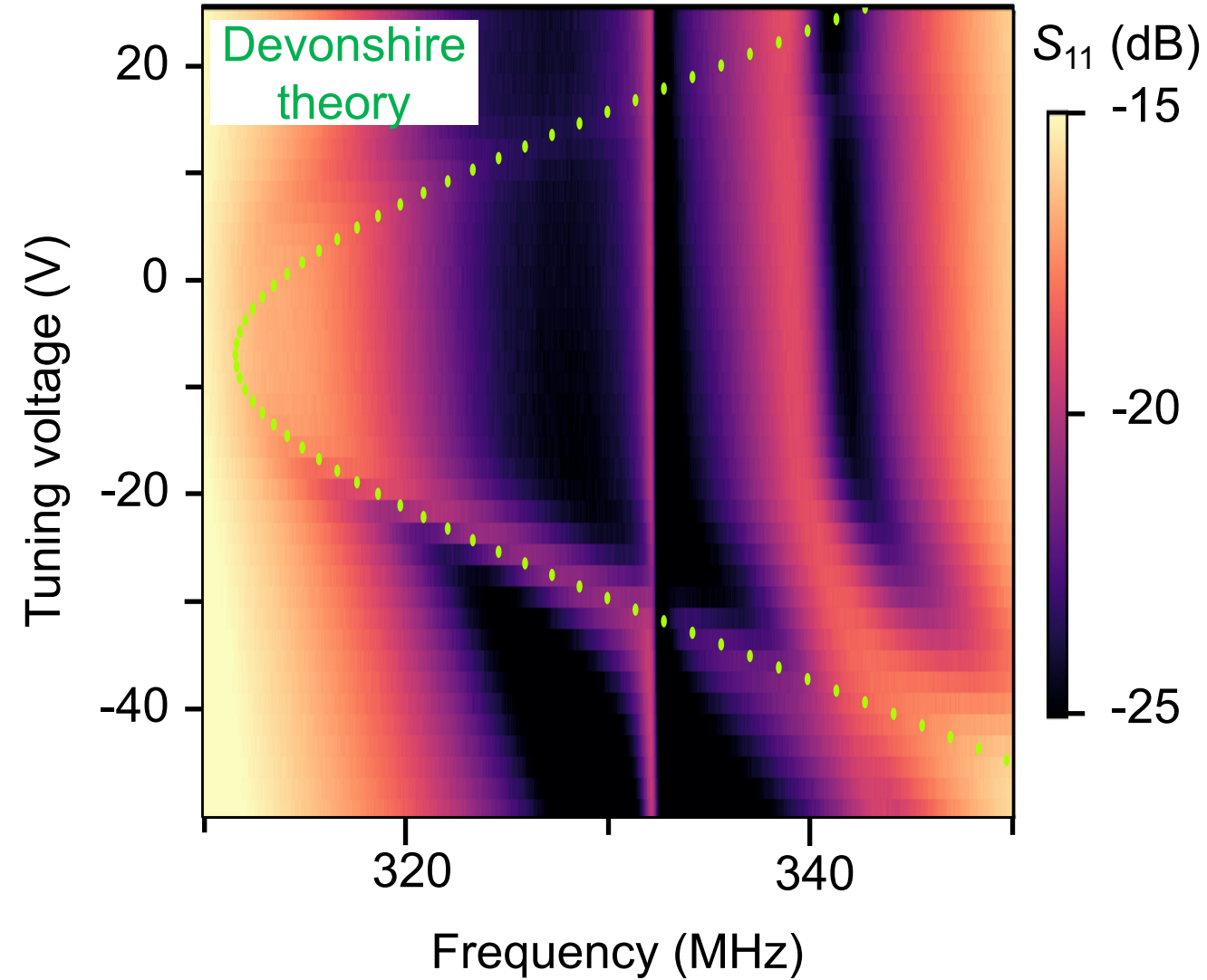
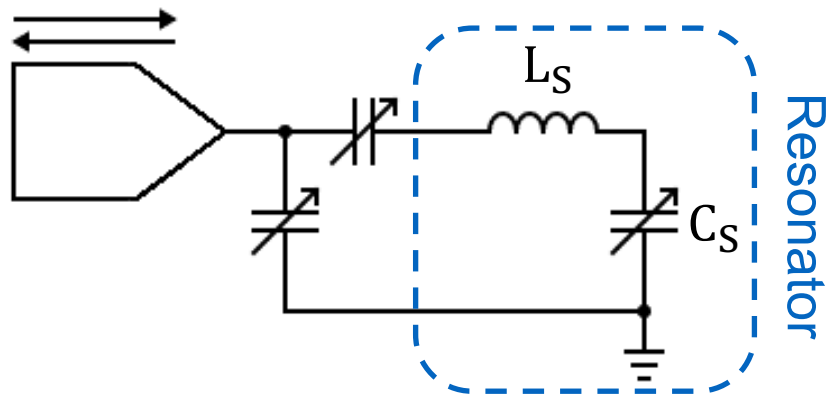
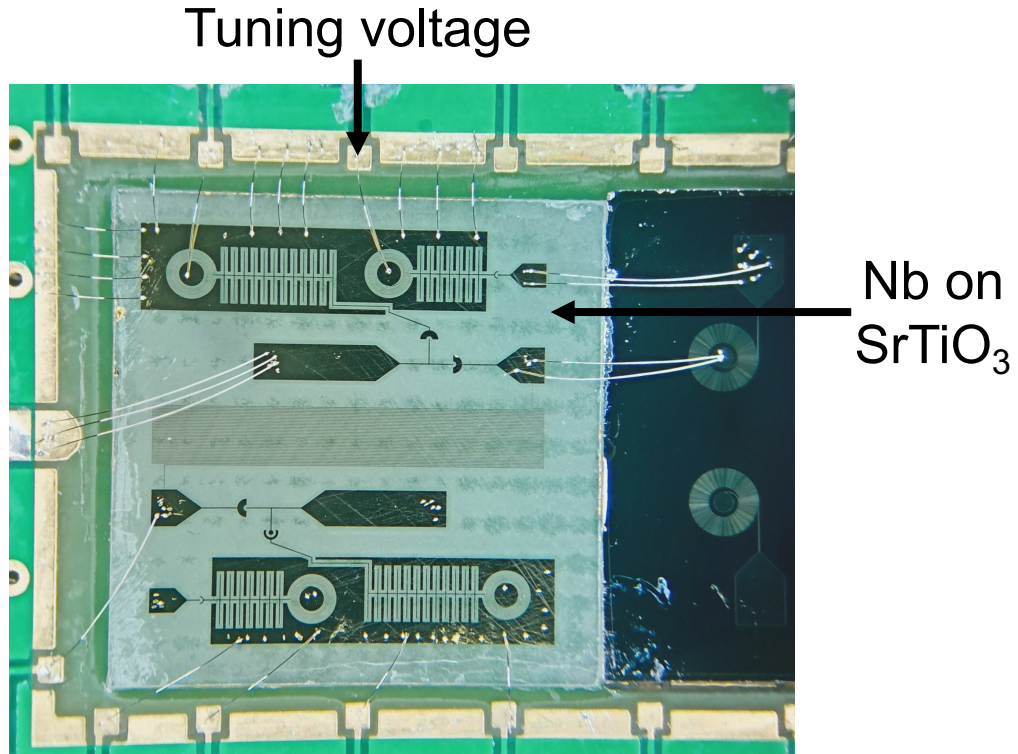
A parametric amplifier that is robust against magnetic fields



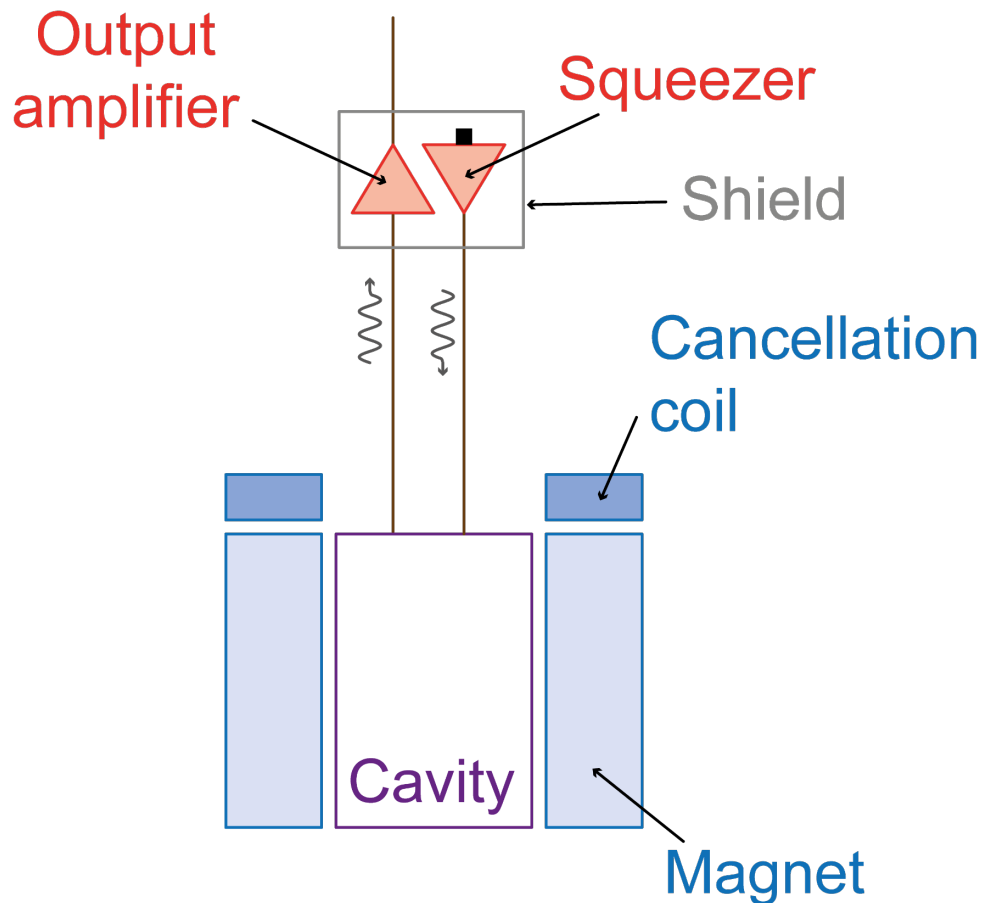
Strontium titanate –
a quantum paraelectric material



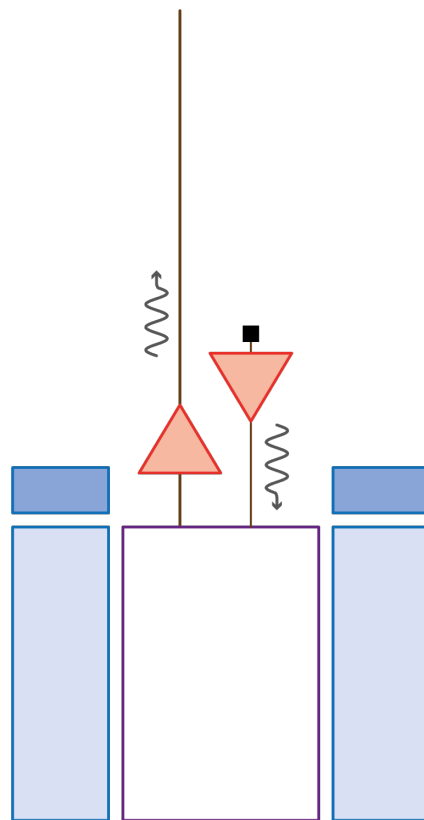
Tuning an RF resonator using paraelectricity



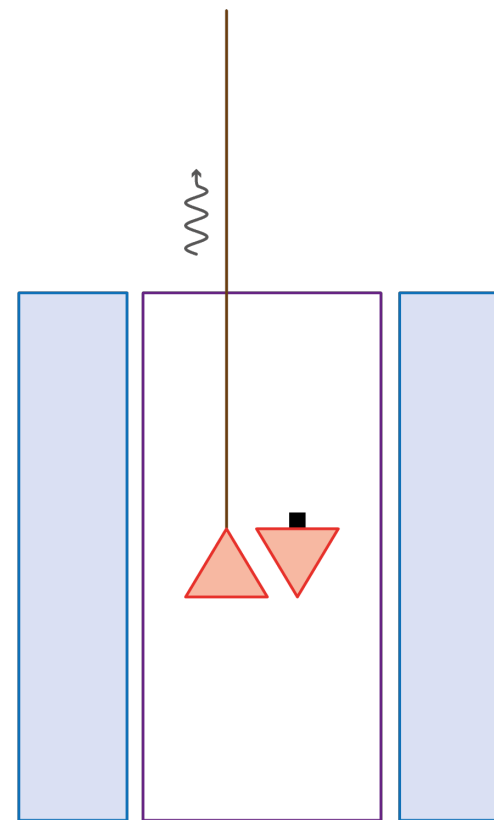
Paraelectric amplifiers for axion searches: the vision



**QSHS-type
haloscope**



**With paraelectric
amplifier**

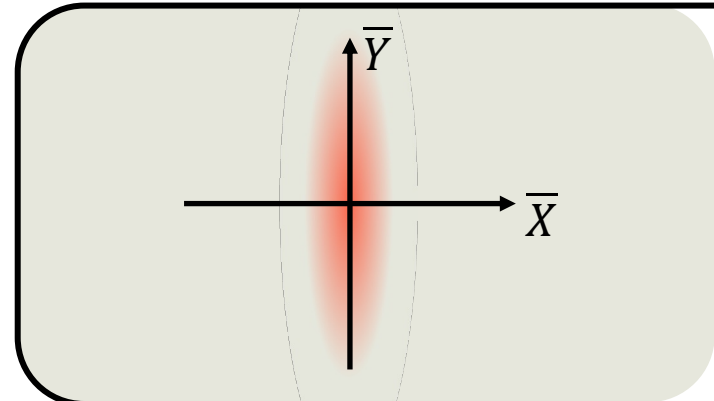


Future haloscope

Quantum sensors for the hidden sector

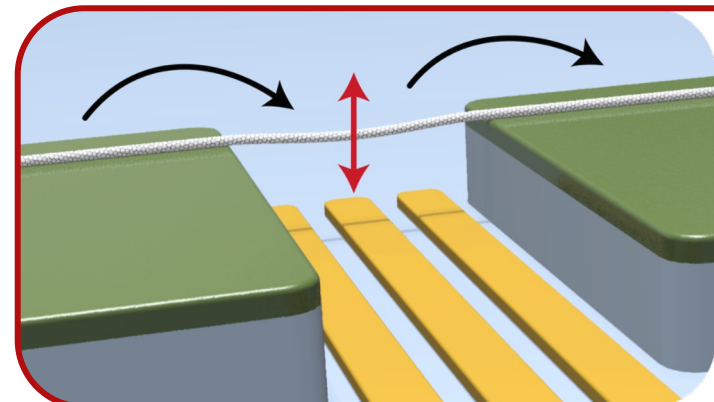


Why should you look for axions?



The QSHS experiment

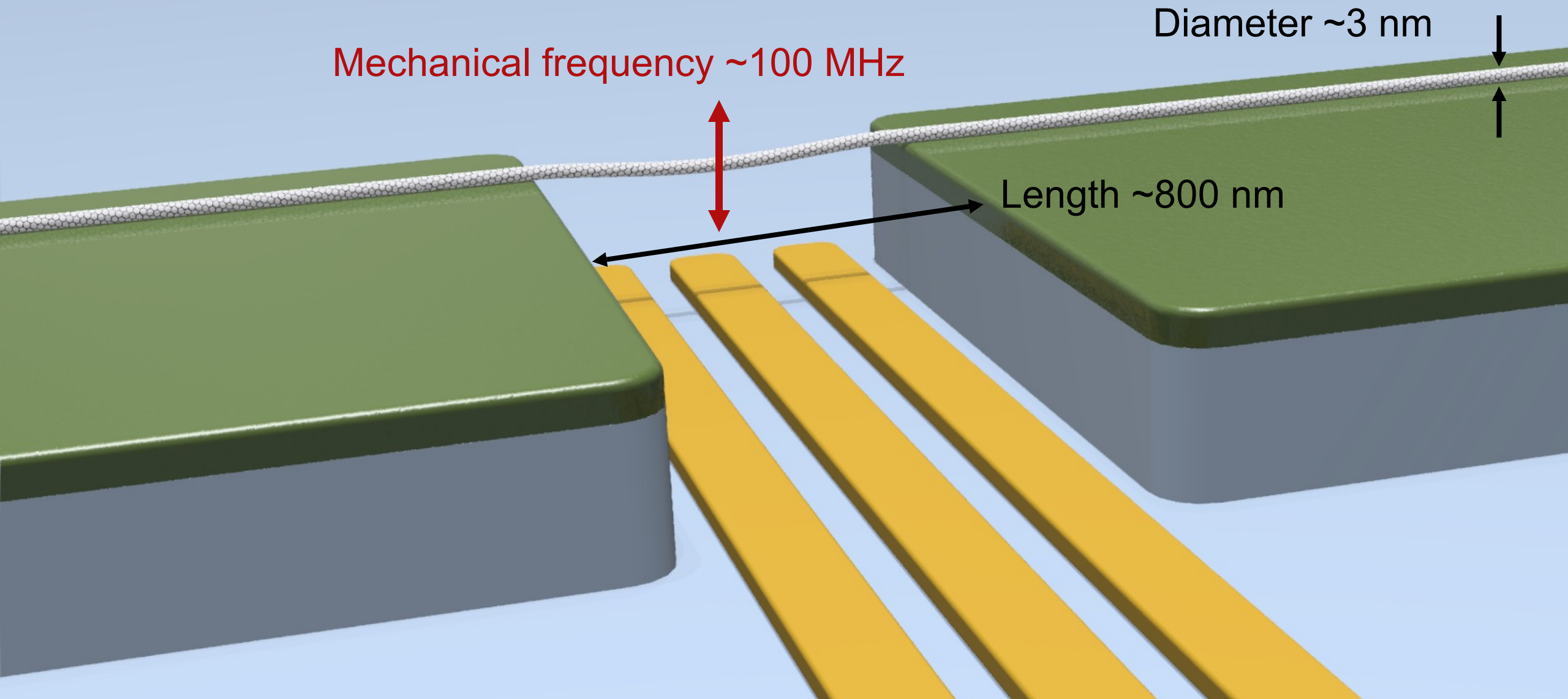
Looking for axions using quantum technology



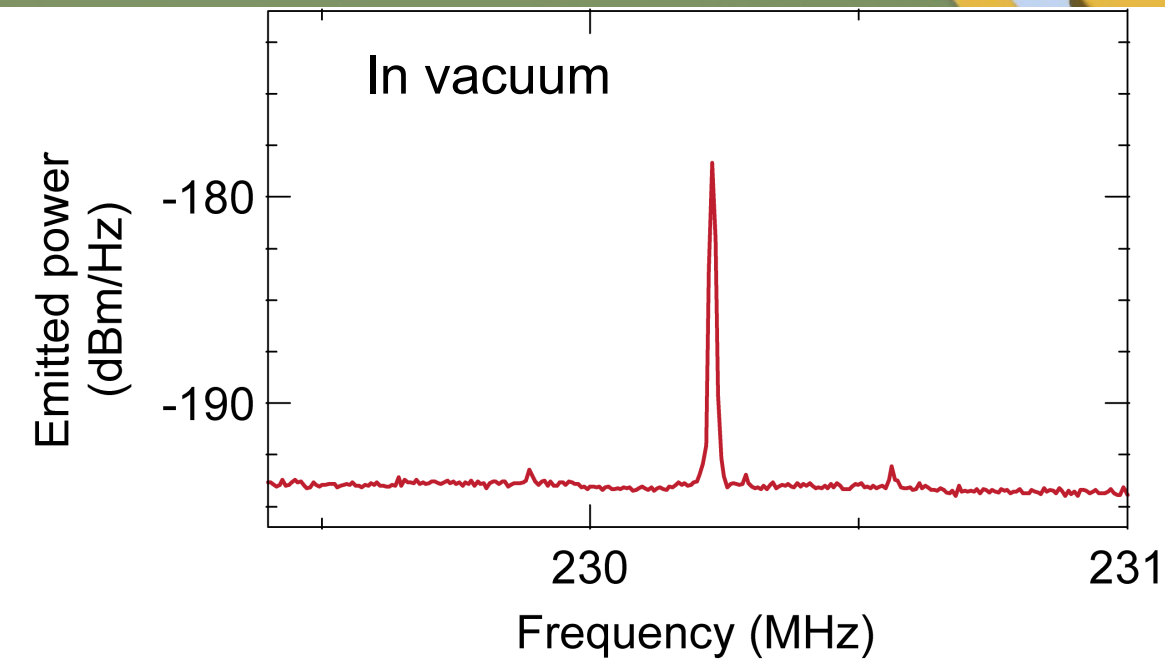
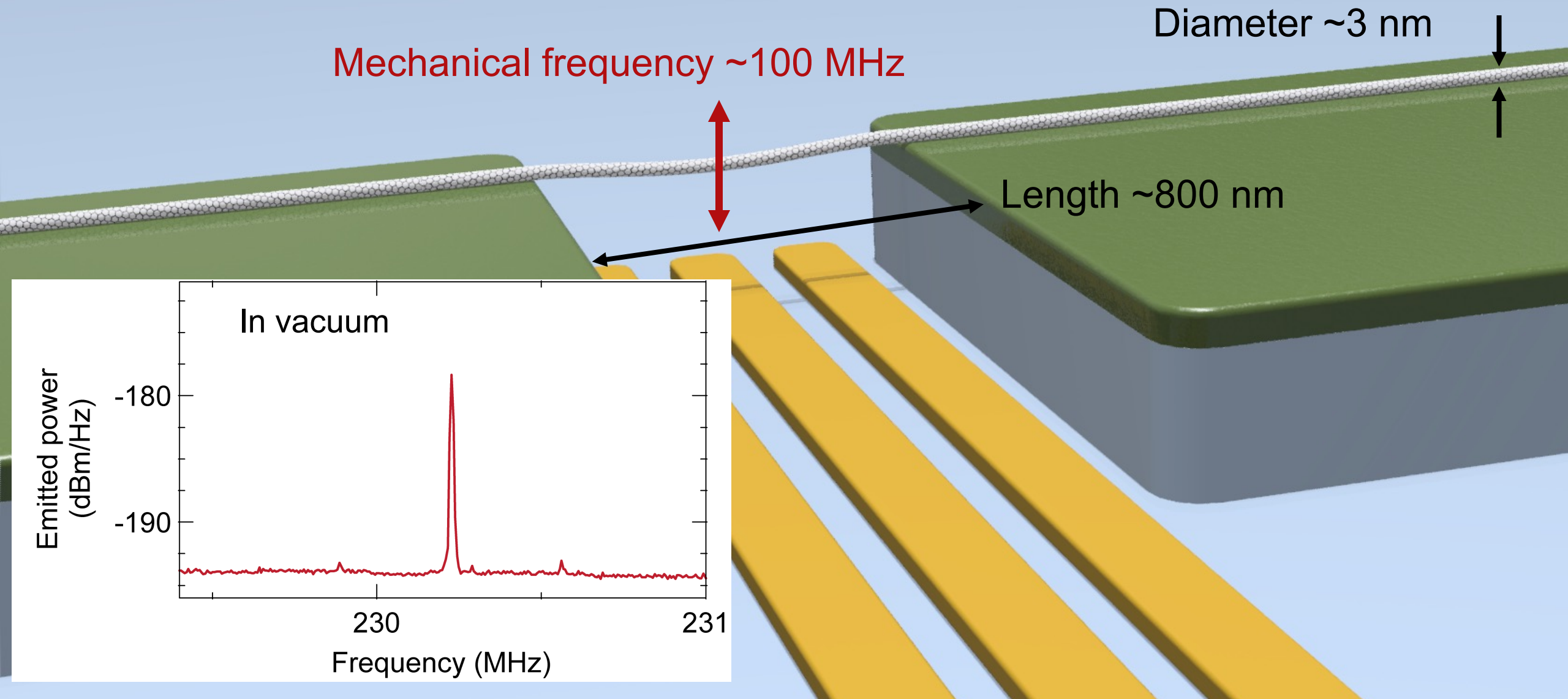
Vibrating nanotubes in superfluid helium

Another device for particle searches?

Carbon nanotube resonators for measuring superfluid helium

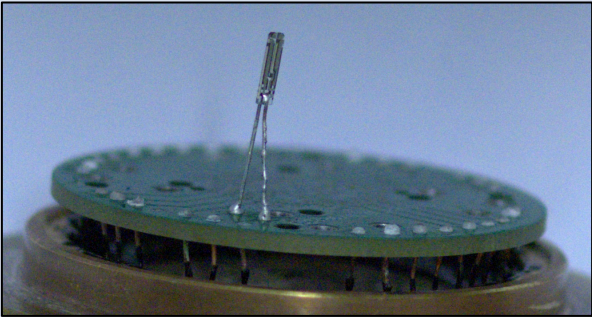


Carbon nanotube resonators for measuring superfluid helium

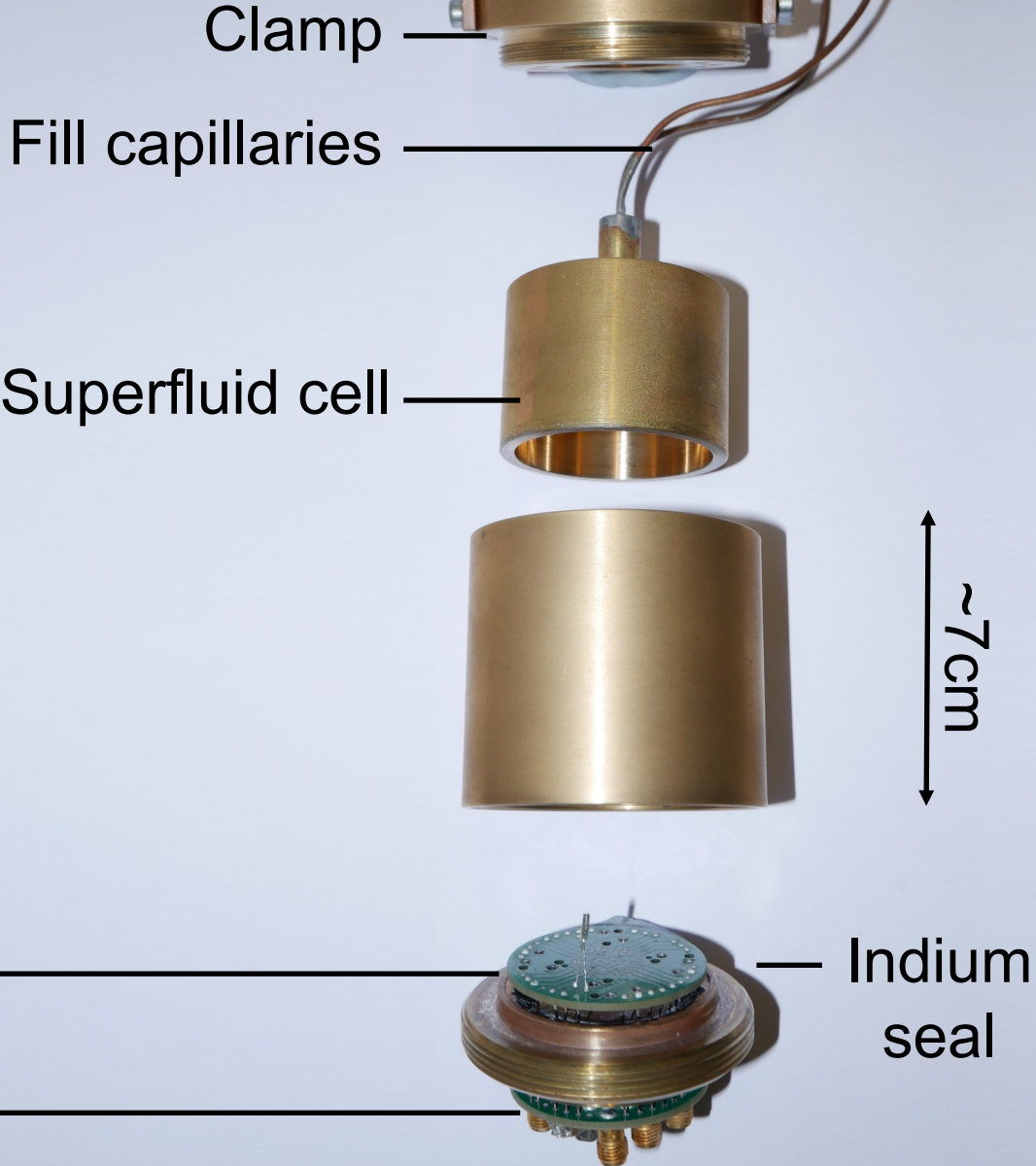
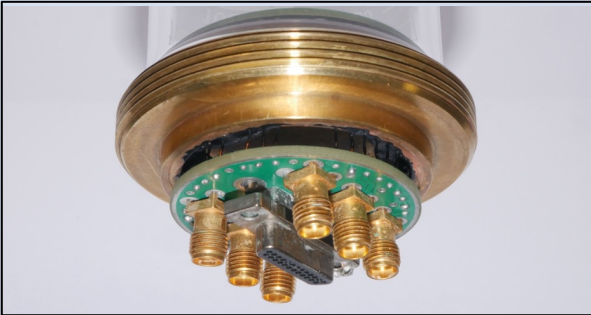


Superfluid plumbing with high-frequency electronics

Sample mounting board
(with quartz
level meter/thermometer)

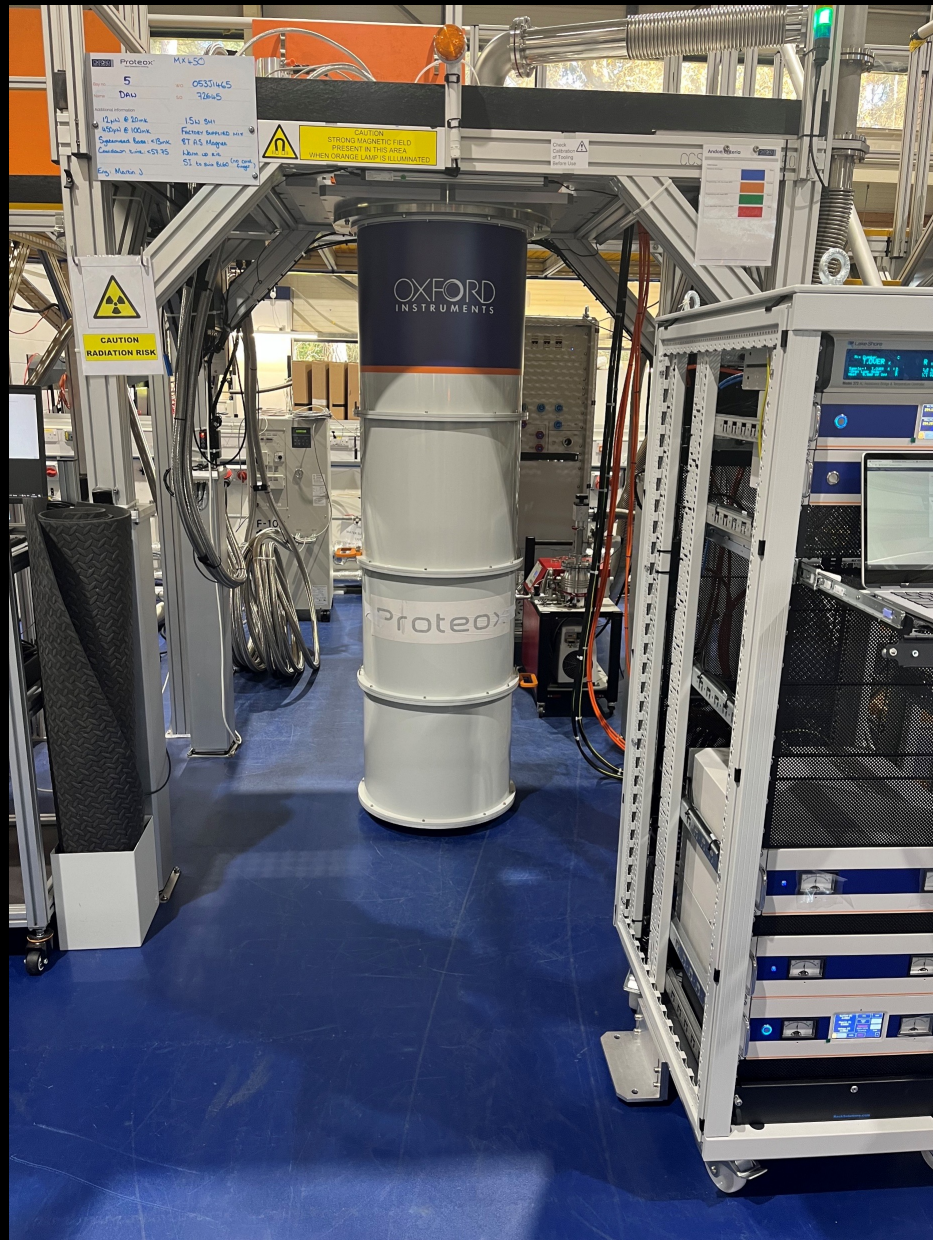


Electrical penetrations

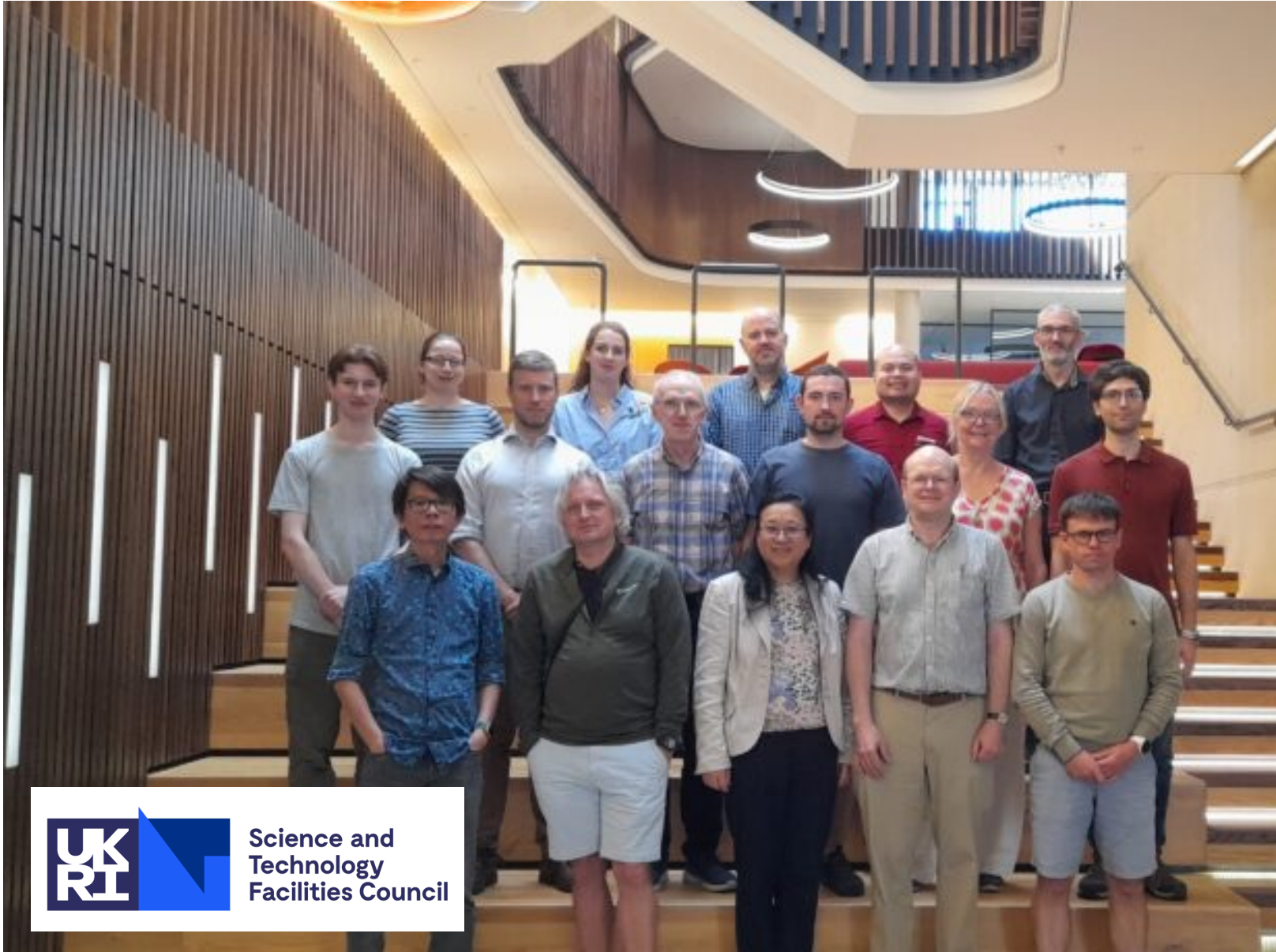




QSHS fridge installation is this month!



The QSHS collaboration



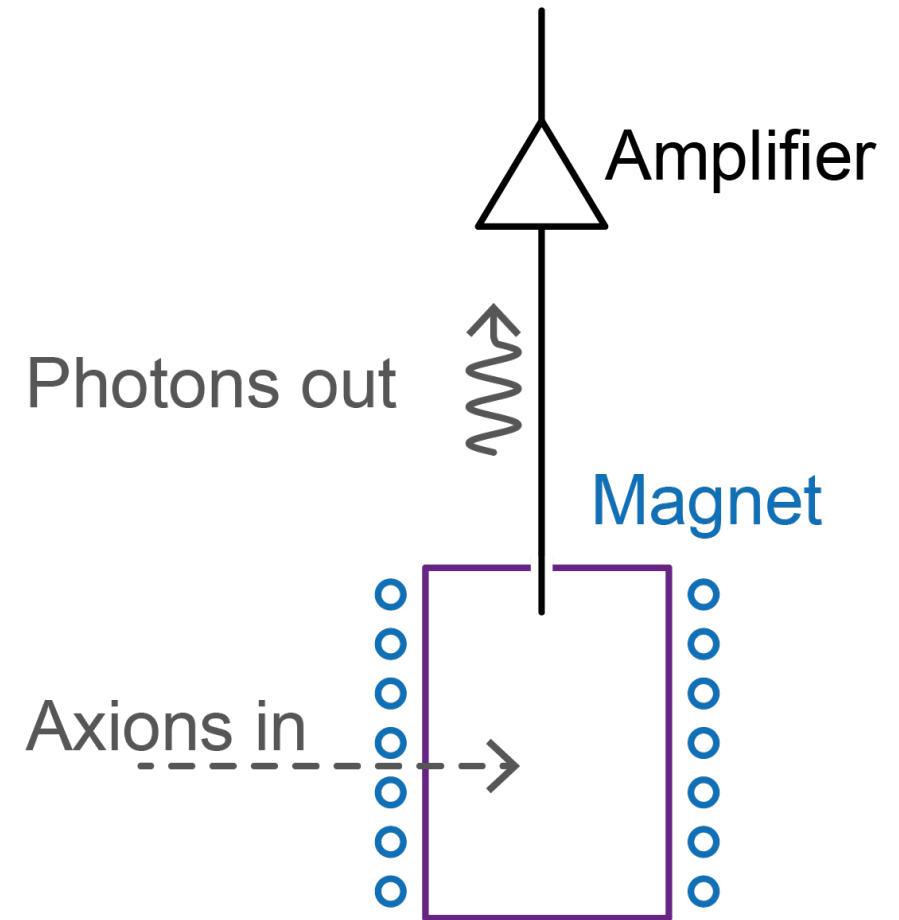
Especially:

- Ed Daw (Sheffield, PI)
- Paul Smith (Sheffield)
- Ian Bailey (Lancaster)
- Yuri Pashkin (Lancaster)
- Ed Romans (UCL)
- Others at qshs.org and also my students
- Deepanjan Das
- Scott Henderson
- Saba Khan

and others at Lancaster including Roch Schaanen, Malcolm Poole, Patrick Steger, and Sergey Kafanov

Quantum Sensors for the Hidden Sector

- Axions are an attractive candidate to be dark matter, with a motivated mass search range of 1-100 μeV ($= 0.25 - 25 \text{ GHz}$).
- Searches are ongoing using haloscope detectors; however, the search rate is limited by the standard quantum limit on conventional amplifiers.
- Quantum amplifiers can circumvent this limit and allow for much faster searches.
- The QSHS consortium is building a haloscope in which novel quantum electronics will be applied to this challenge.





The QSHS collaboration

