



Overview of Axion Searches

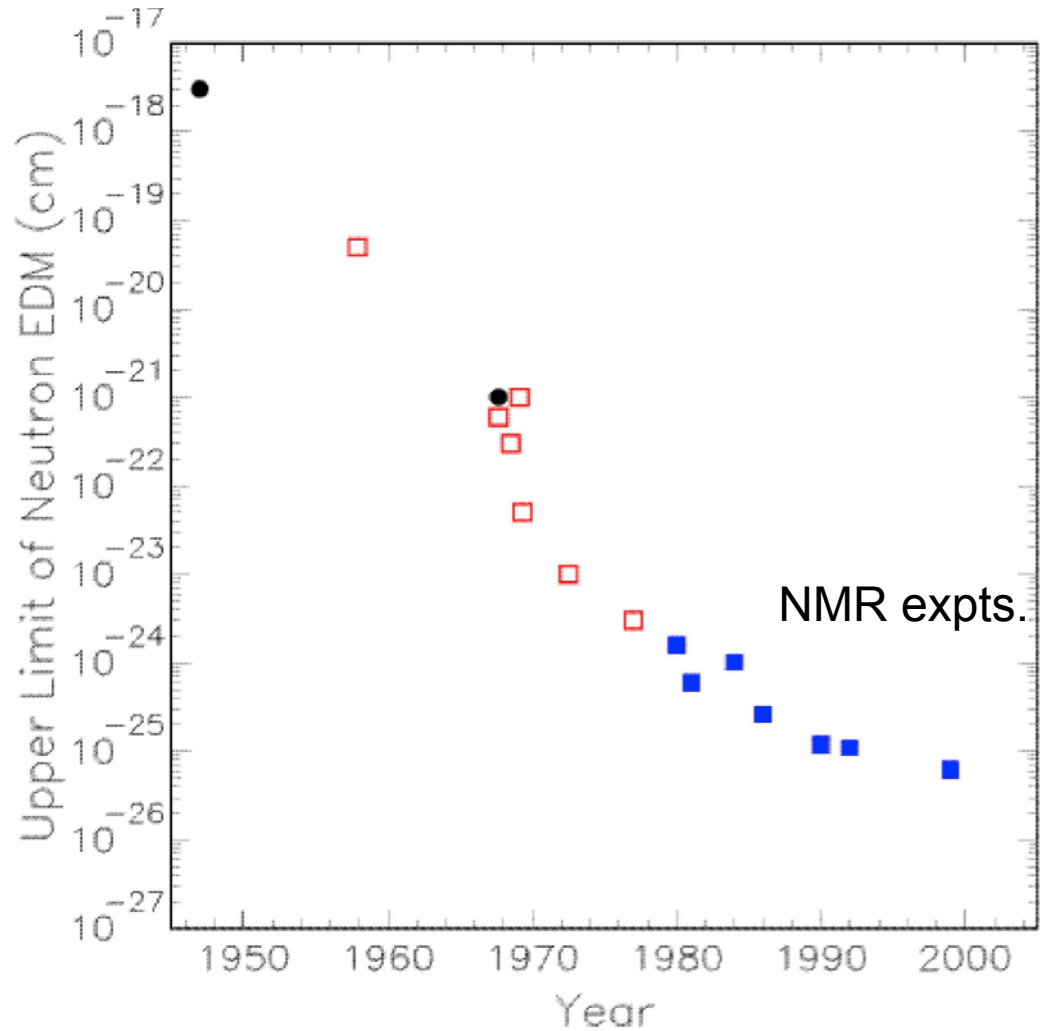
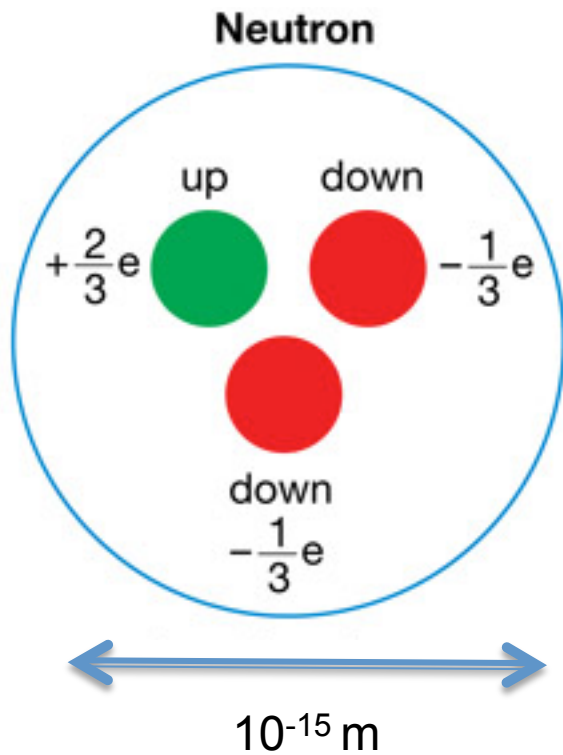
Gianpaolo Carosi
LLNL

TAUP

July 25th, 2017

Why is the neutron electric dipole moment so small?

Naive estimate
gives
 $d_n \approx 10^{-16} \text{ e-cm}$



Peccei-Quinn solution to the strong-CP problem



Roberto Peccei



Helen Quinn



Steven Weinberg



Frank Wilczek

- Peccei & Quinn: Postulate new U(1) symmetry that would be spontaneously broken.
- Weinberg & Wilczek: A new Goldstone boson (the axion)
- Remnant axion VEV nulls QCD CP violation.
- Only free parameter: Symmetry breaking scale (f_a).
- “Invisible Axion”: $f_a \gg$ Weak Scale
- Two general classes of models
 - **KSVZ** [Kim (1979), Shifman, Vainshtein, Sakharov (1980)]: “QCD axion” or “hadronic axion”
 - **DFSZ** [Dine, Fischler, Srednicki (1981), Zhitnitsky (1980)]



“clean up” the Strong-CP problem

Axion couplings

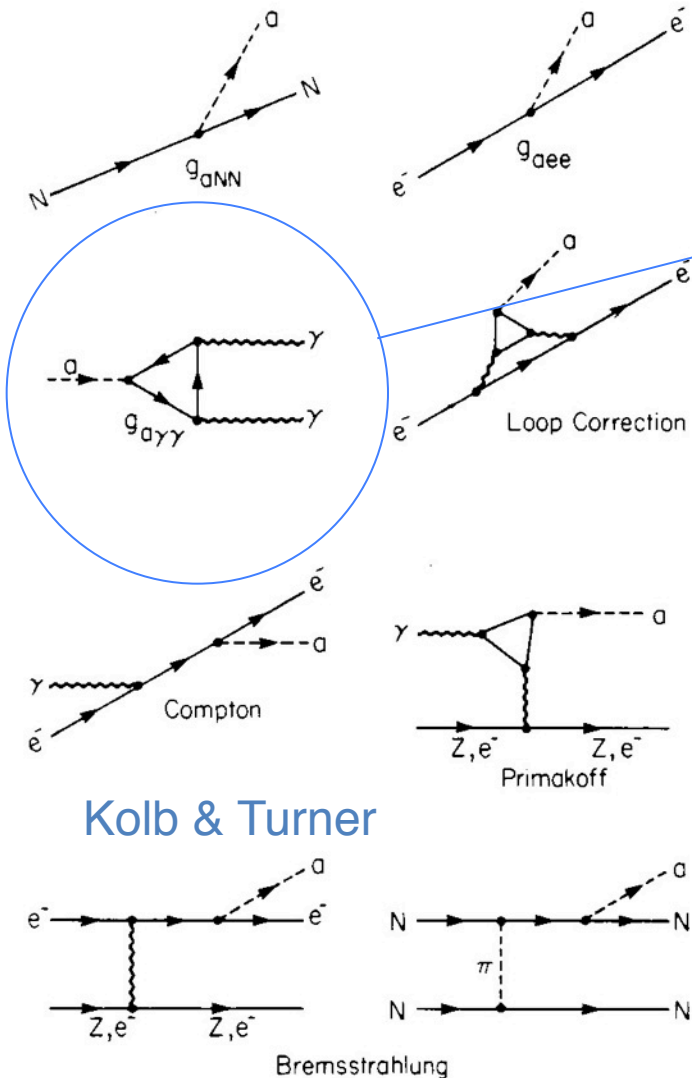
General classes of couplings

- Axion – Nucleon
- Axion – Electron
- Axion – Photon

$g_{a\gamma\gamma}$ is a process with small model uncertainty
 Easily calculable
 Coupling used for haloscopes

Rate depends on “unification group”
 (that is, the particles in the loops),
 ratio of u/d quark masses,
 and mostly f_{PQ}

$$g_{a\gamma\gamma} \sim \frac{\alpha}{f_{PQ}} \left(\frac{E}{N} - 1.95 \right)$$



Kolb & Turner

Variety of experiments



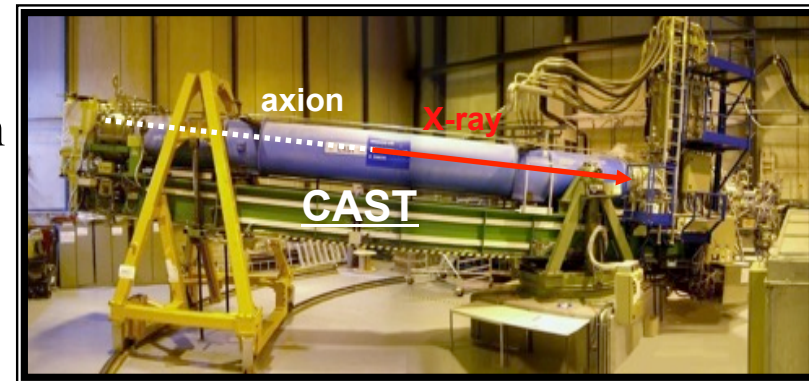
Pierre Sikivie

■ Axions as Dark Matter

- Haloscopes: Microwave cavities in solenoid magnet
- Look for dark matter axions (low mass) converting to photons in B-Field
 - Sikivie PRL 51:1415 (1983)
- **New techniques being explored: NMR, LC-circuit, Axion Wind**

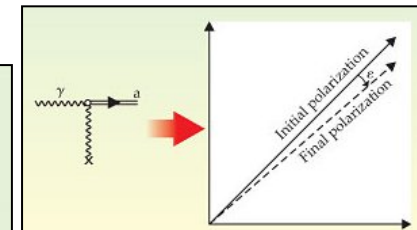
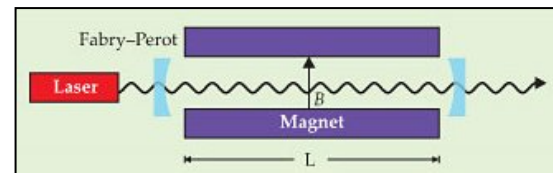
■ Axions from the Sun

- Helioscopes: Axions generated from the sun
 - Sikivie PRL 51:1415 (1983)
 - Van Bibber et al. PhysRevD 39:2089 (1989)
 - **CAST, IAXO**
- Bragg scattering, noble liquids (g_{ae})

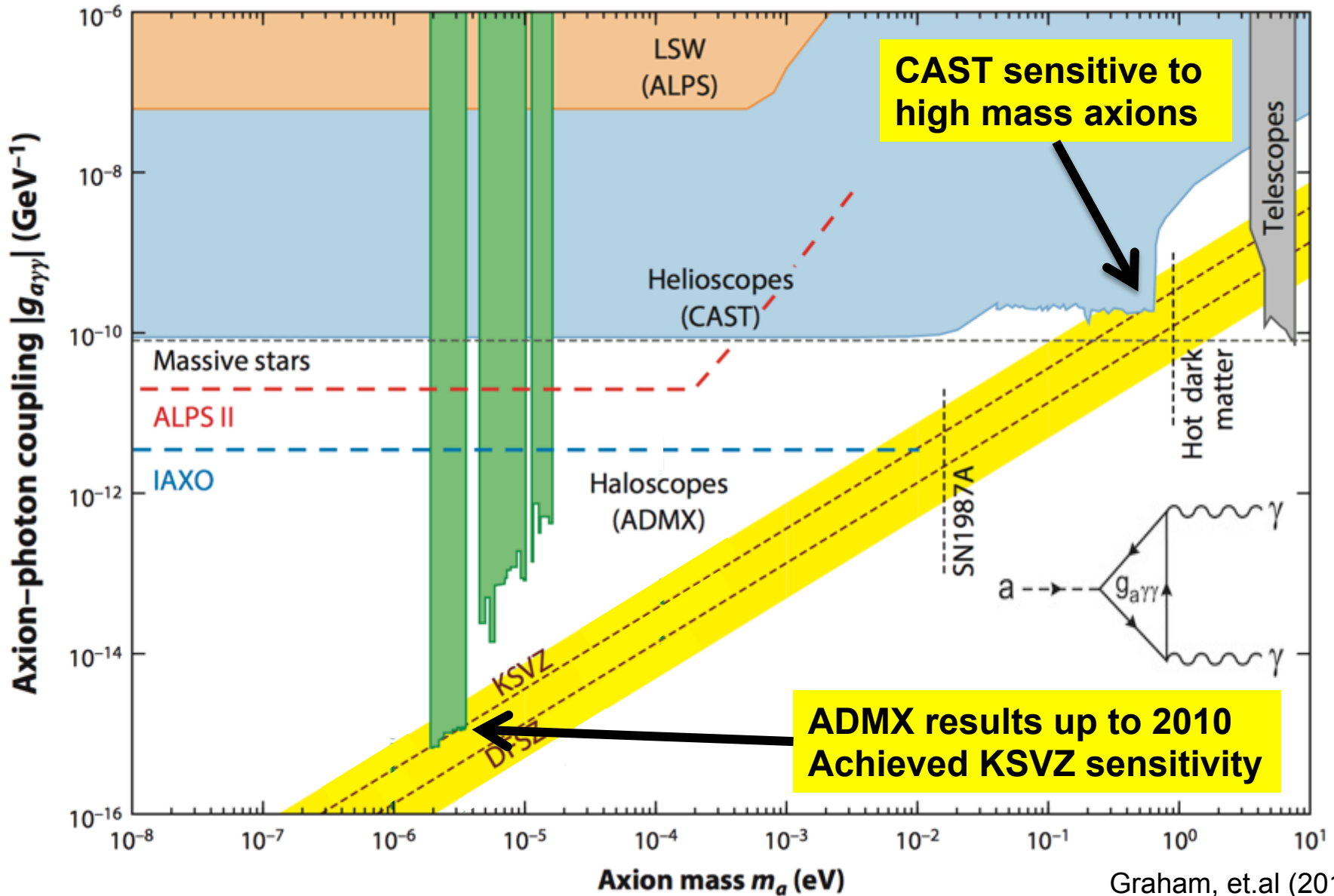


■ Axions in the Lab

- Photon regeneration and polarization changes
 - **PVLAS, ALPS**
- Modifications to short range forces
 - **ARIADNE, Torsion-balance**



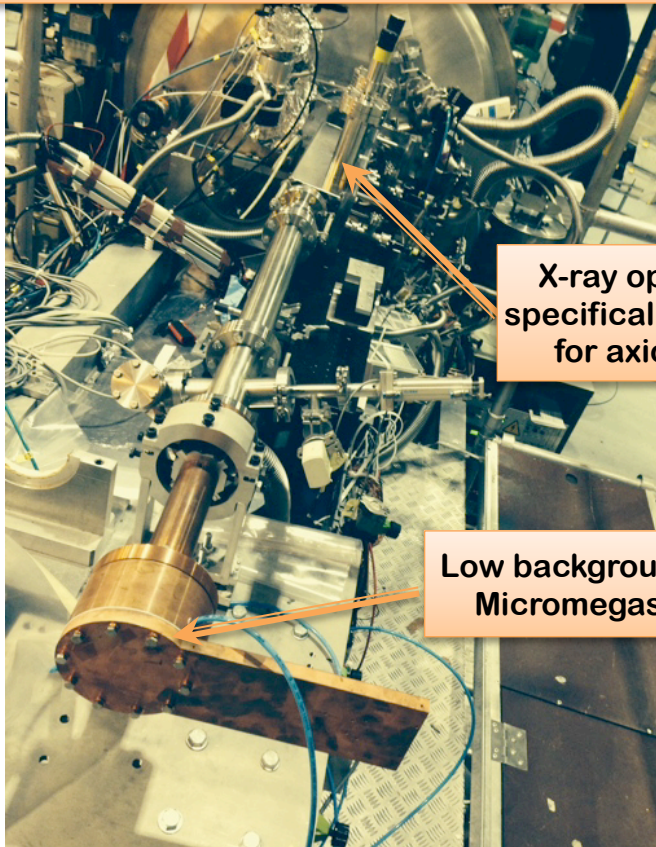
Axion parameter space



Graham, et.al (2016)

New CAST Limit

Enabled by the
IAXO pathfinder system



X-ray optics
specifically built
for axions

Low background
Micromegas

nature
physics

ARTICLES

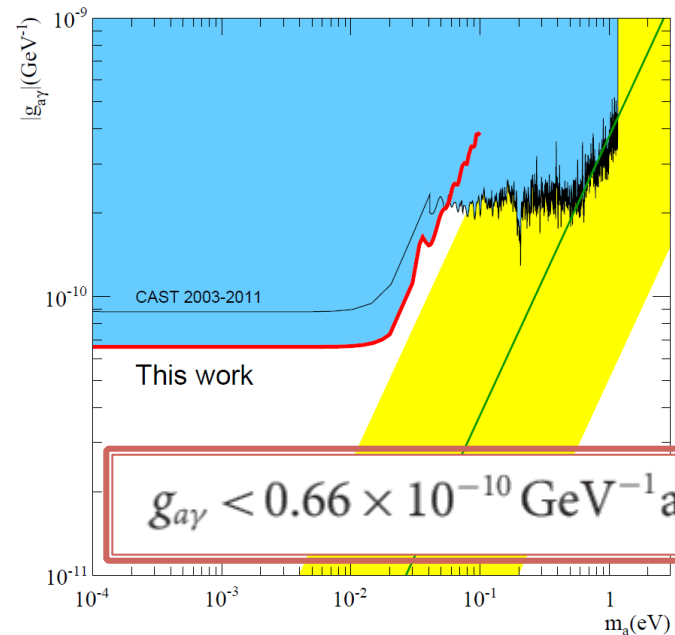
PUBLISHED ONLINE: 1 MAY 2017 | DOI: 10.1038/NPHYS4109

OPEN

New CAST limit on the axion-photon interaction

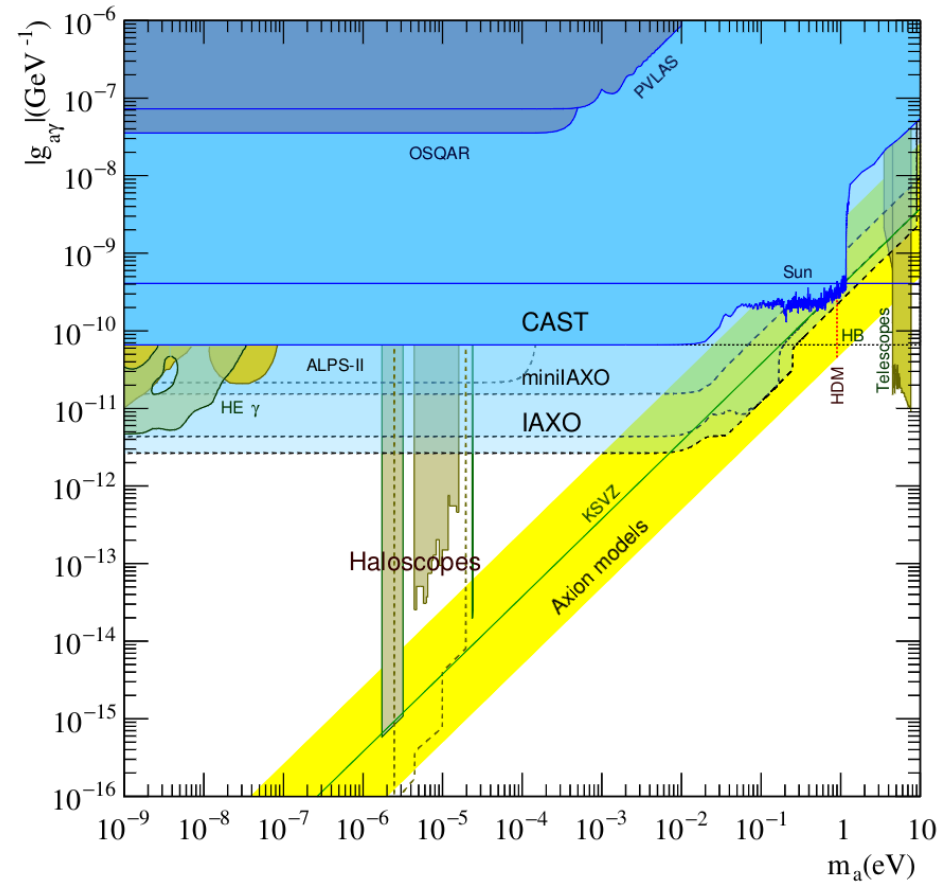
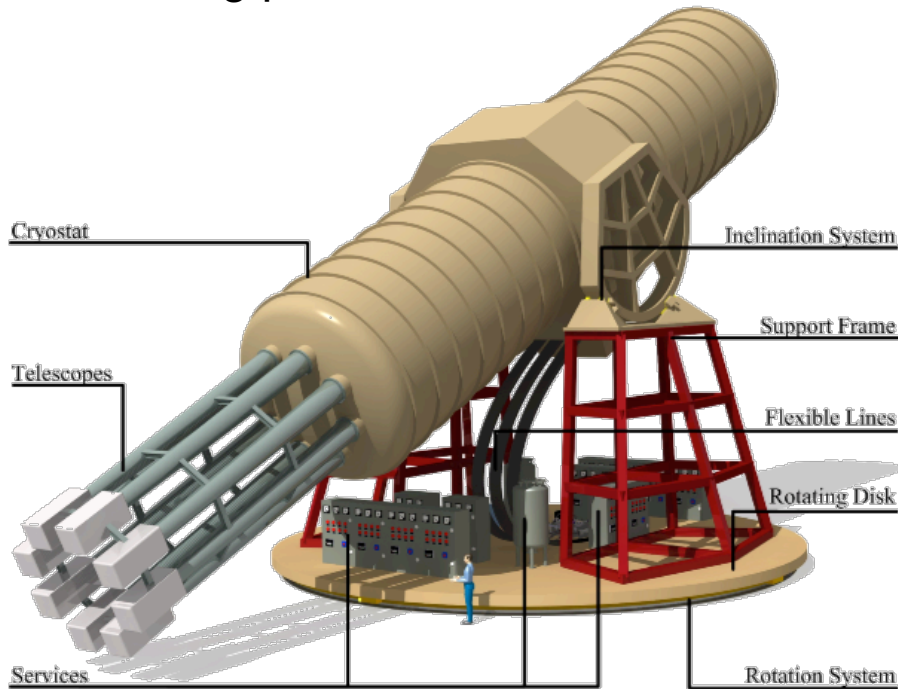
CAST Collaboration[†]

Hypothetical low-mass particles, such as axions, provide a compelling explanation for the dark matter in the universe. Such particles are expected to emerge abundantly from the hot interior of stars. To test this prediction, the CERN Axion Solar Telescope (CAST) uses a 9 T refurbished Large Hadron Collider test magnet directed towards the Sun. In the strong magnetic field, solar axions can be converted to X-ray photons which can be recorded by X-ray detectors. In the 2013-2015 run, thanks to low-background detectors and a new X-ray telescope, the signal-to-noise ratio was increased by about a factor of three. Here, we report the best limit on the axion-photon coupling strength ($0.66 \times 10^{-10} \text{ GeV}^{-1}$ at 95% confidence level) set by CAST, which now reaches similar levels to the most restrictive astrophysical bounds.



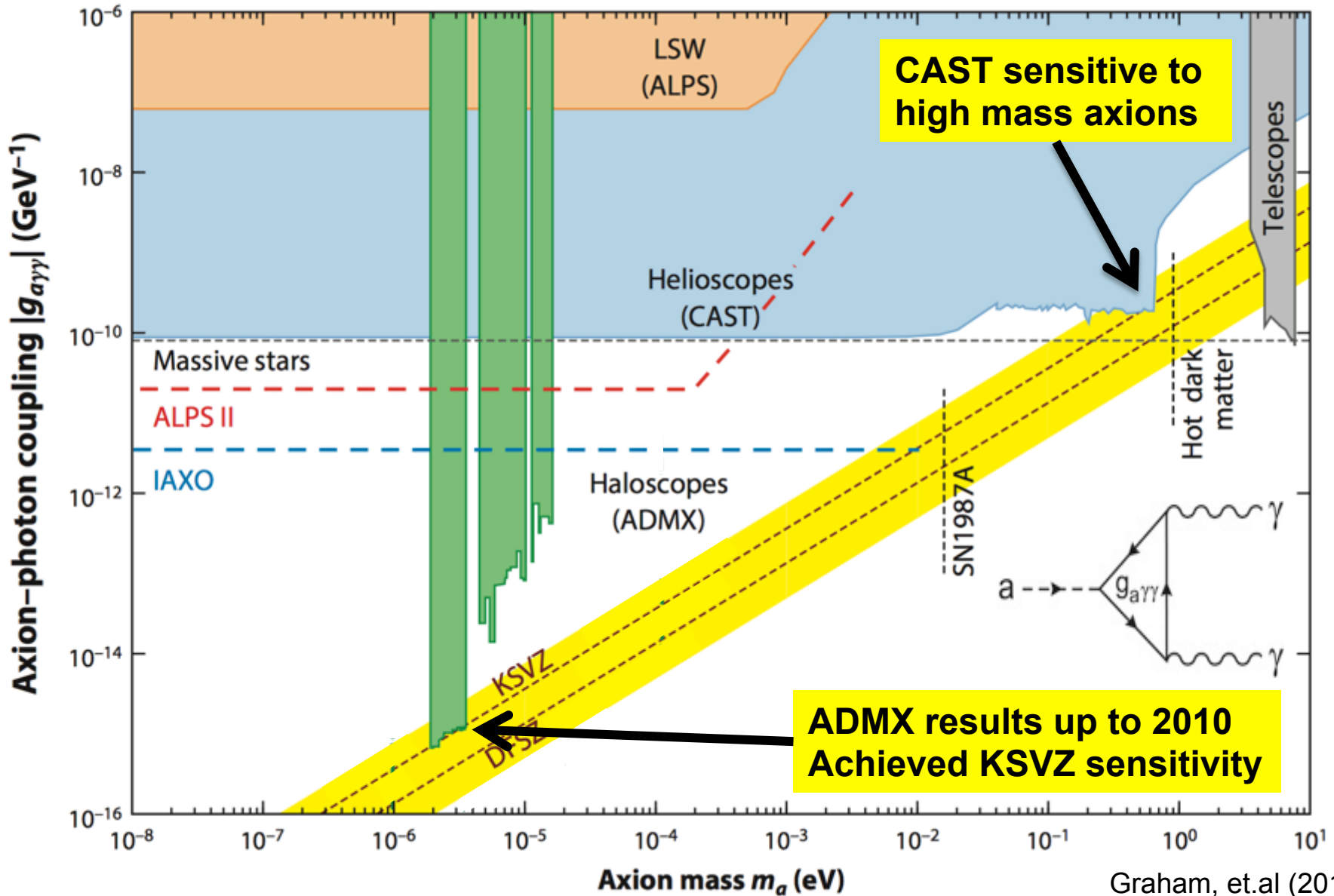
IAXO: the next Helioscope generation

- Large toroidal 8-coil magnet $L = \sim 20$ m
- 8 bores: 600 mm diameter each
- 8 x-ray telescopes + 8 detection system
- Rotating platform with services



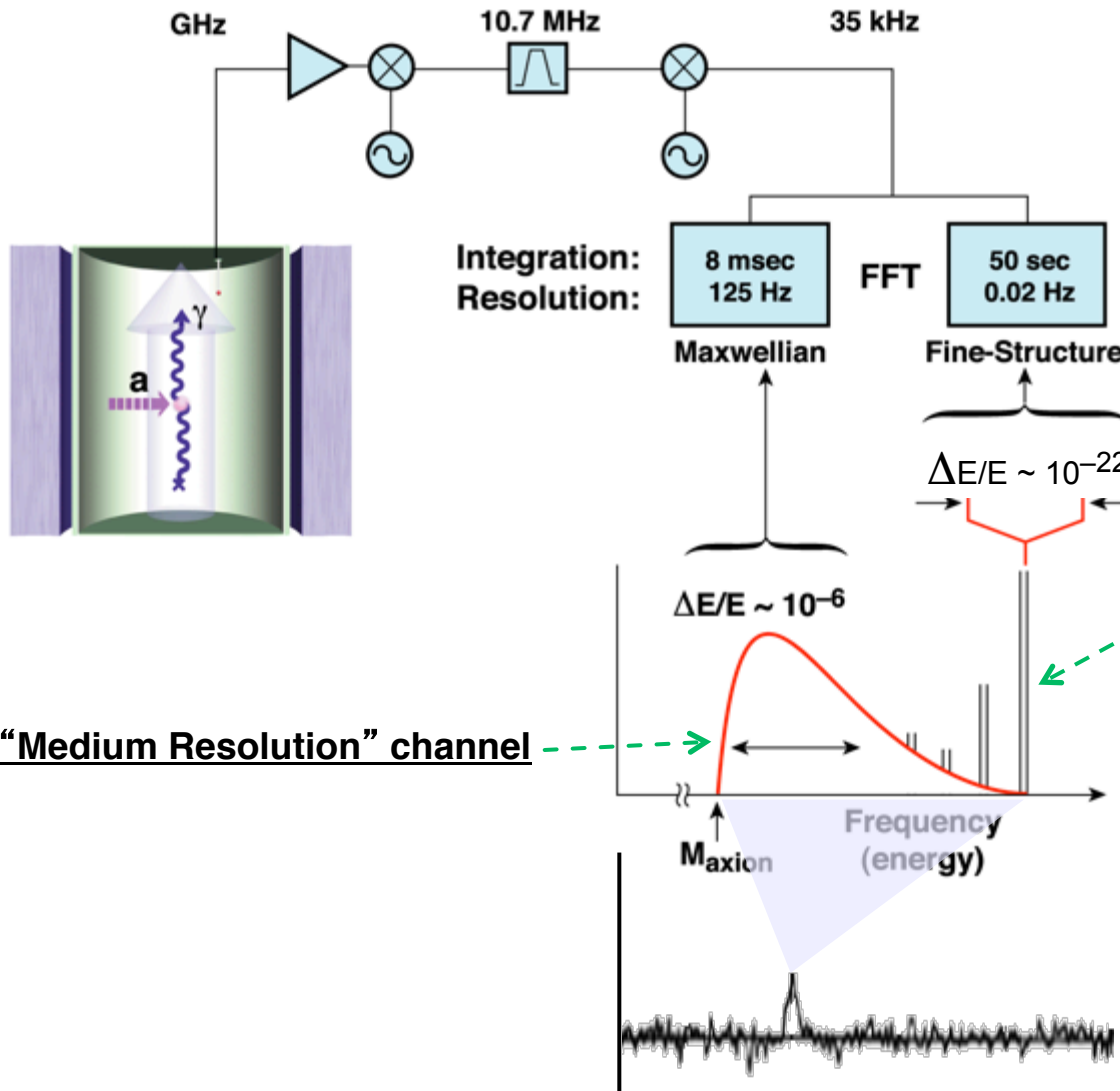
*Please see Beljana Lacic's talk tomorrow in the DM parallel session

Axion parameter space



Graham, et.al (2016)

The ADMX experimental layout



Local Milky Way density:

$$\rho_{halo} \sim 450 \text{ MeV/cm}^3$$

Thus for $m_a \sim 10 \mu\text{eV}$:

$$\rho_{halo} \sim 10^{14} \text{ cm}^{-3}$$

“High Resolution” channel

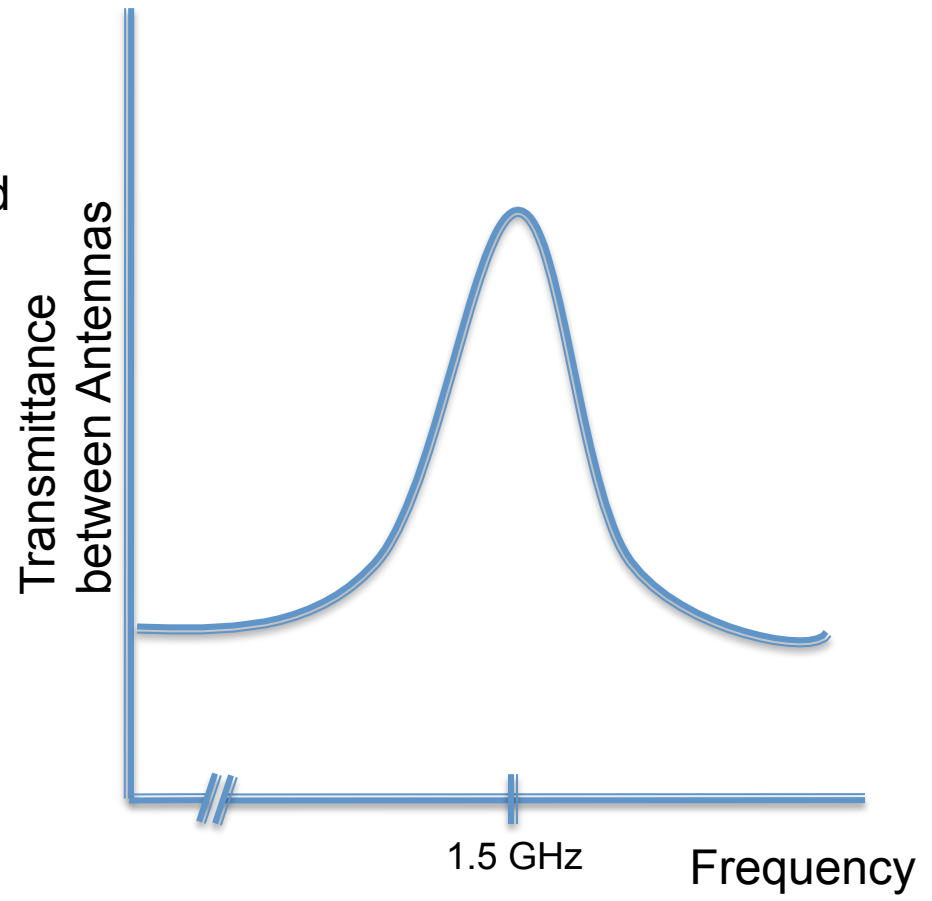
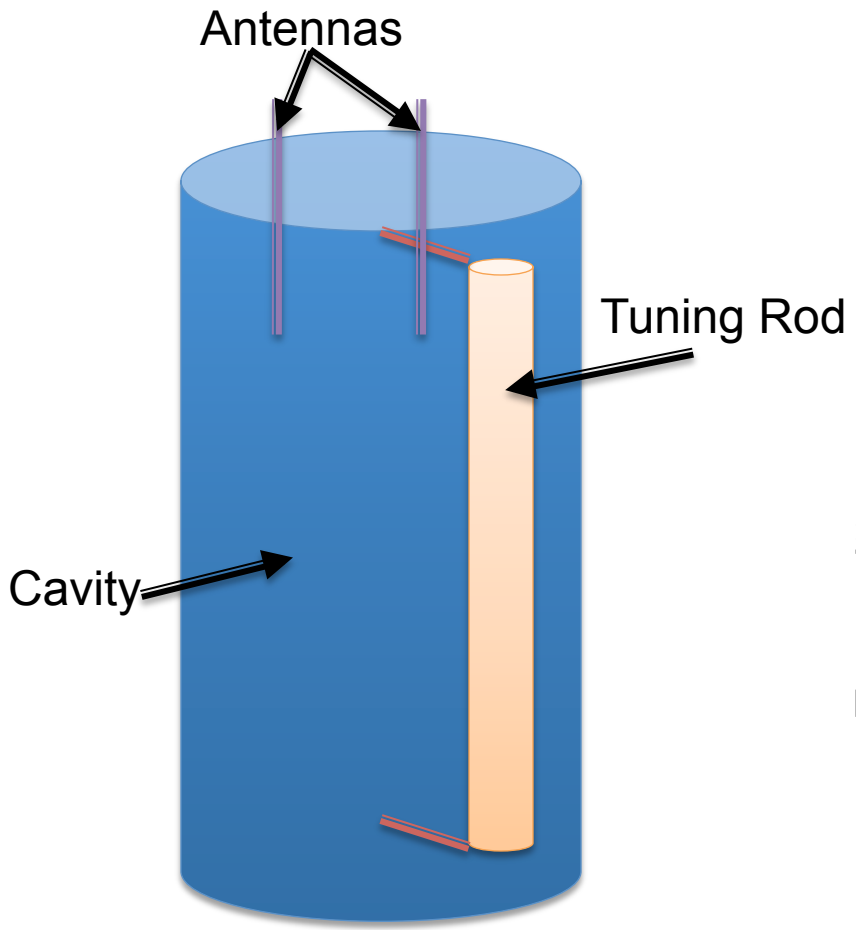
$$\beta_{\text{virial}} \sim 10^{-3} :$$

$$\lambda_{\text{De Broglie}} \sim 100 \text{ m}$$

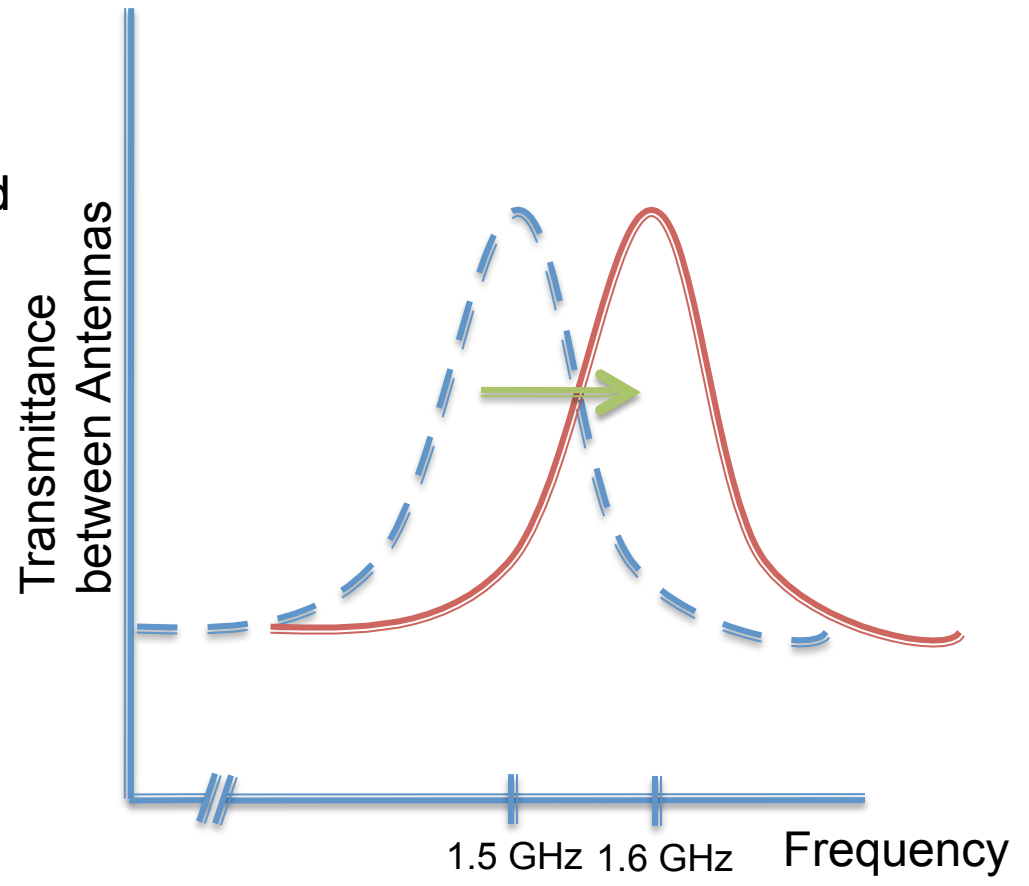
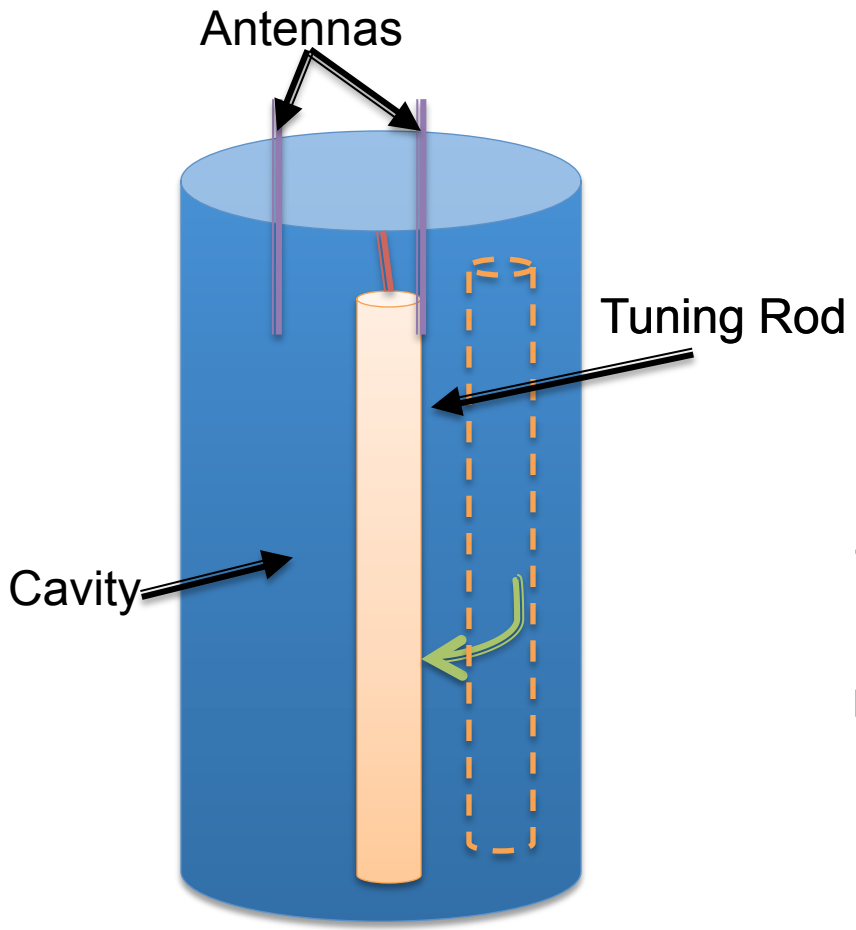
$$\Delta \beta_{\text{flow}} \sim 10^{-11} :$$

$$\lambda_{\text{Coherence}} \sim 1000 \text{ km}$$

Microwave Cavity needs tunable resonance



Microwave Cavity needs tunable resonance



ADMX experimental layout

Field Cancellation Coil

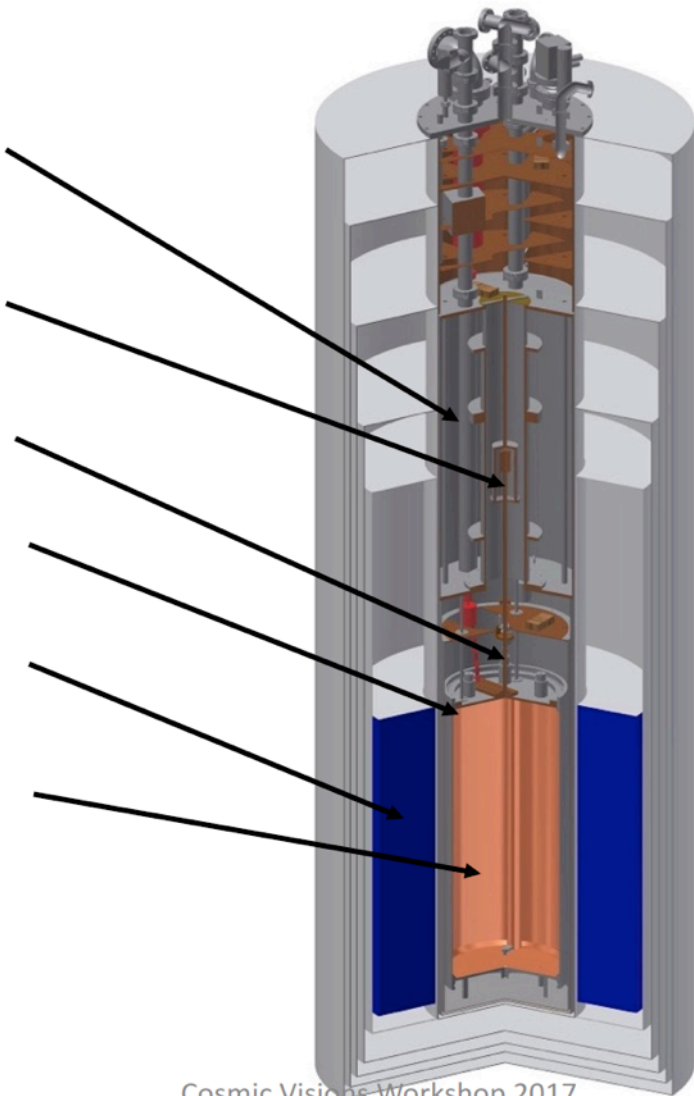
SQUID Amplifier Package

Dilution Refrigerator

Antennas

8 Tesla Solenoid Magnet

Microwave Cavity

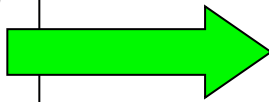


Cosmic Visions Workshop 2017



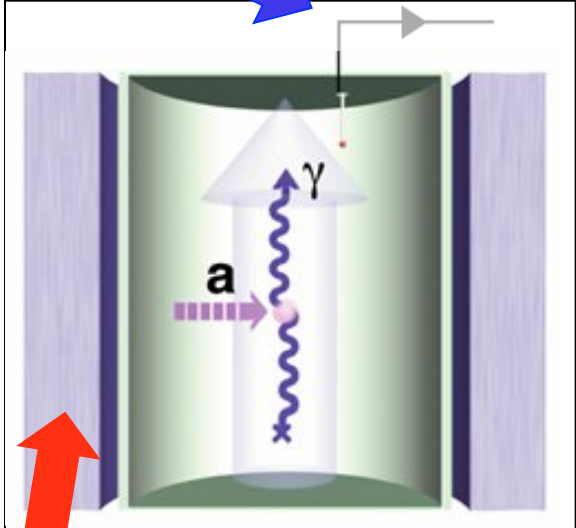
The Radiometer equation dictates strategy

$$\frac{s}{n} = \frac{P_{sig}}{kT_S} \cdot \sqrt{\frac{t}{\Delta\nu}}$$



Integration time limited to ~ 100 sec

* Dicke, 1946



System noise temp.

$$T_S = T_{phys} + T_N$$

T_{Quant} ~ 48 mK @ 1 GHz

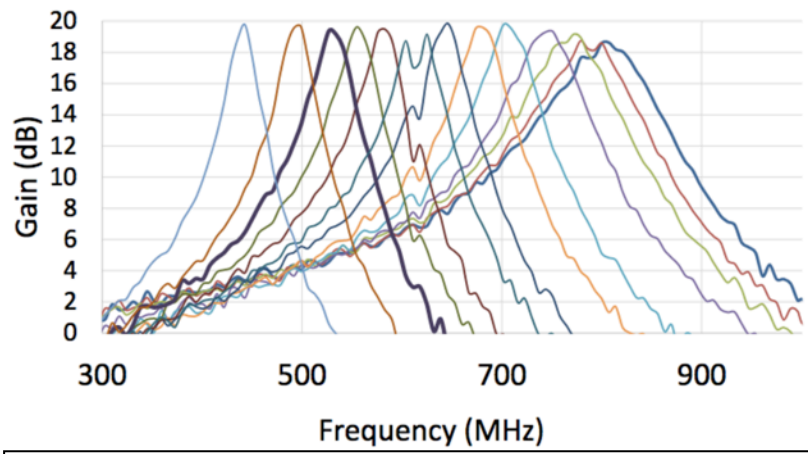
Gen2: Dilution Fridge + Quantum-limited amps

$$P_{sig} \sim (B^2V Q_{cav} C_{010}) (g^2 m_a \rho_a) \sim 10^{-23} \text{ Watts for ADMX}$$

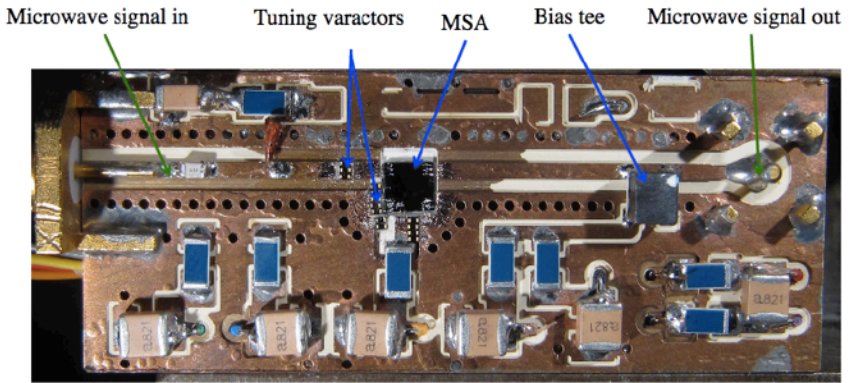
Magnet size, strength B²V ~ \$

Quantum-limited amplifiers

MSA Varactor Tunability

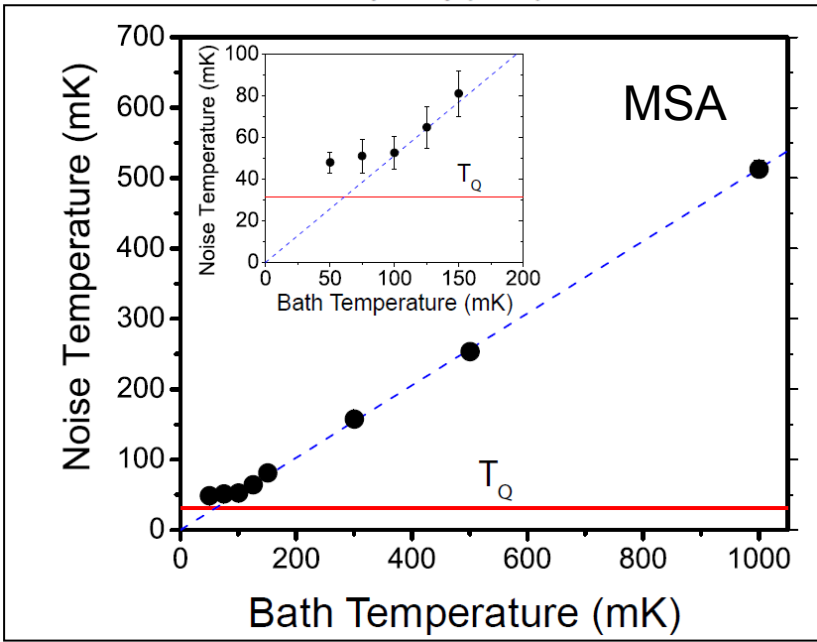


ADMX Tunable MSA



RC filtering for DC lines

Sean O'Kelley,
Clarke Group, UC
Berkeley
< 1 GHz



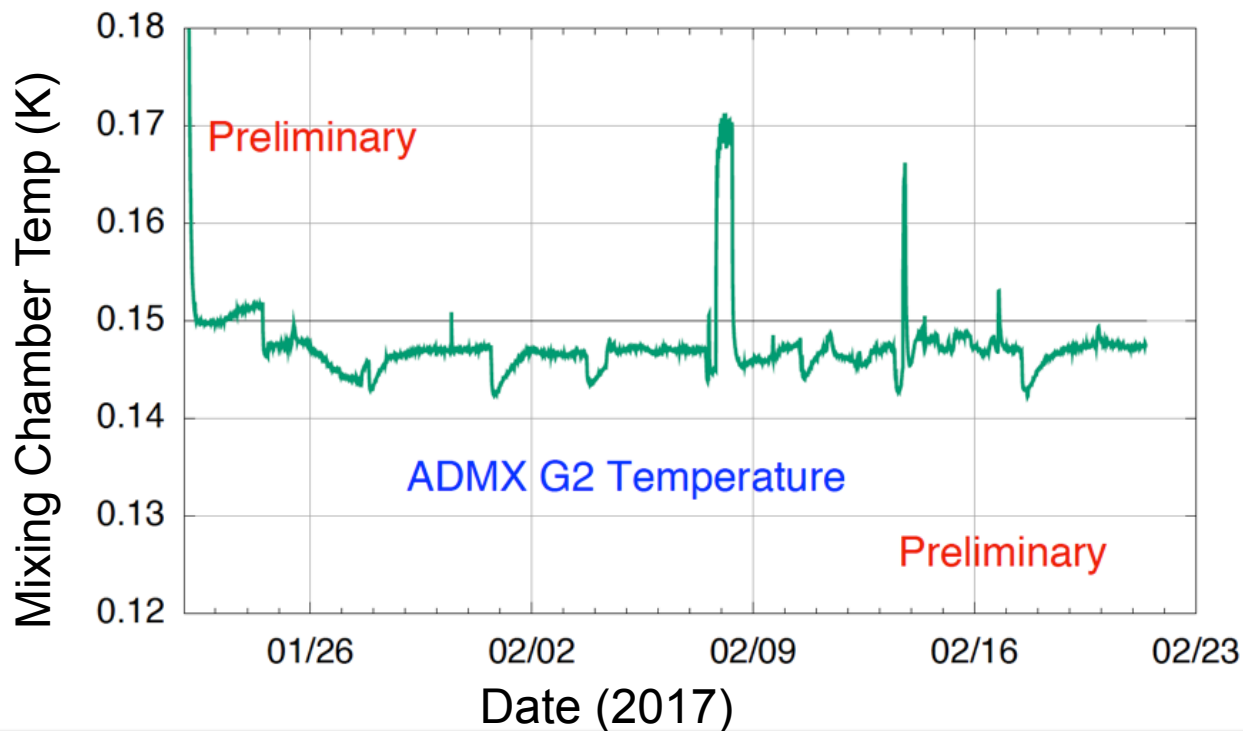
ADMX JPA



Yanjie Qiu, Siddiqi
Group, UC
Berkeley
> 1 GHz

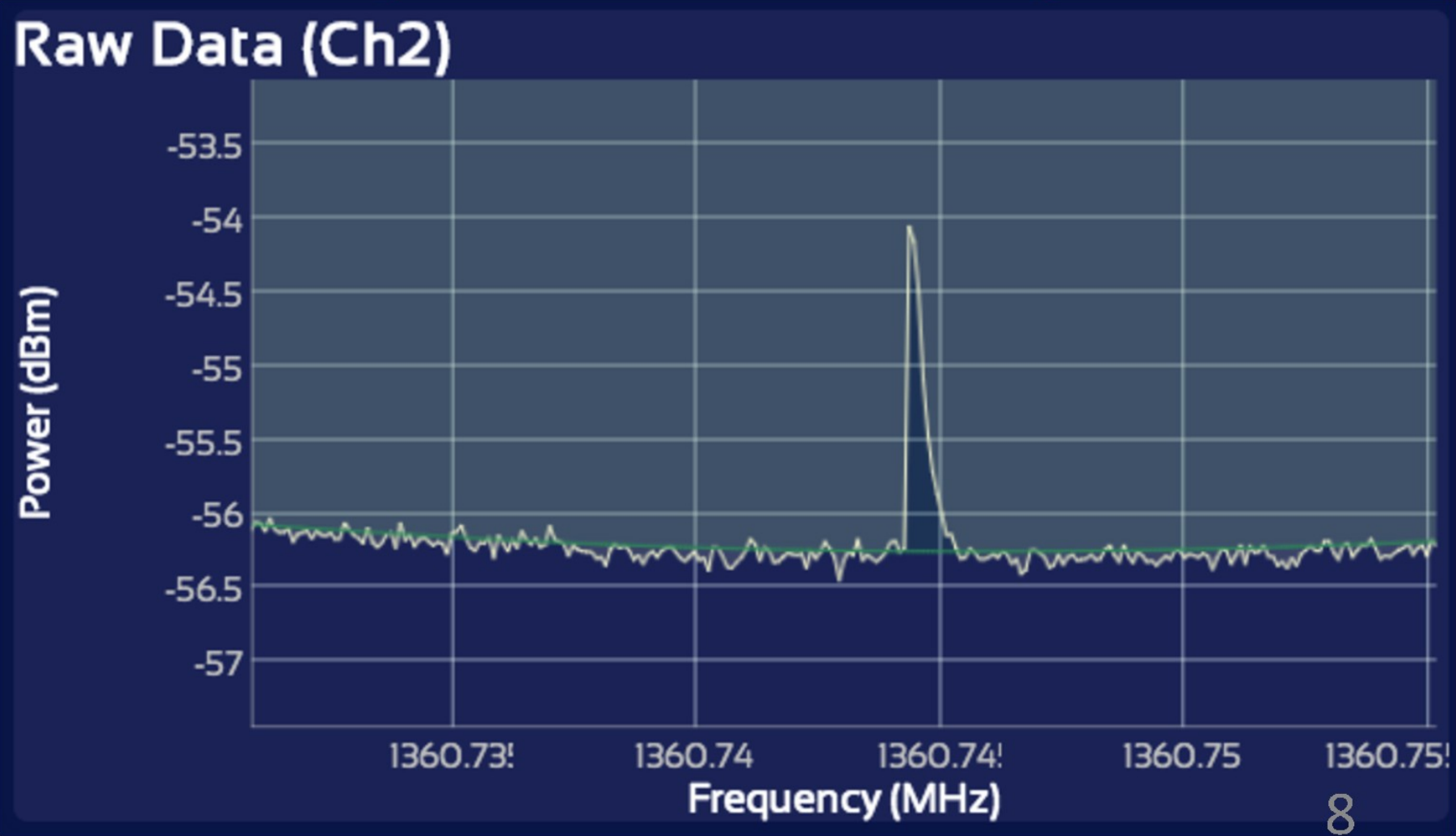
ADMX recently operations

- **Initial commissioning run Aug 9th – Oct 3rd, 2016**
($T_{\text{phys}} \sim 200 \text{ mK}$ and $B \sim 2 \text{ T}$)
- **Oct – Dec, 2016**
 - Upgrades to RF chain and heat sinking
- **First science data run: January - June, 2017**



Raw data and hardware synthetic axion

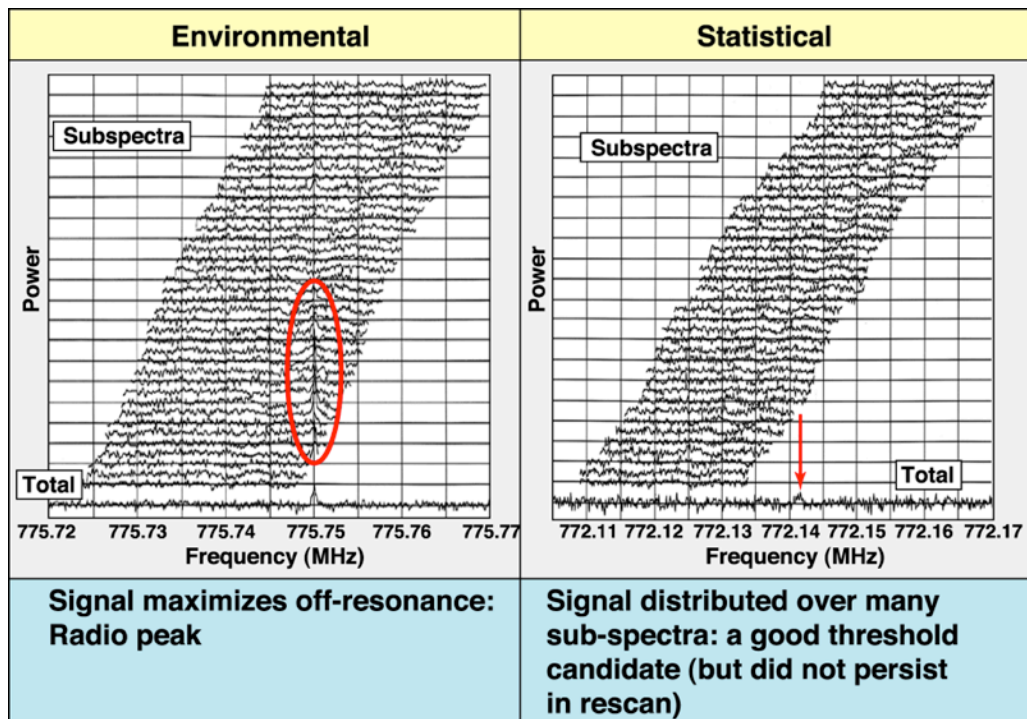
Able to inject custom line-shape through weak port (blinded)



ADMX operations

Live Analysis

1. Cavity frequency scanned until a desired signal-to-noise level is reached.
2. Regions with power above trigger threshold are flagged as potential candidates
 - a. Statistical anomalies, external RF leakage, synthetic injected axions, or AXIONS
3. Rescan candidates; do they persist.
4. If they persist they are transferred to the detection committee
 - a. Several immediate checks...
 - b. Switch to resonant mode that couple to axions (TEM mode).
 - c. Turn B-Field down (power as B^2).



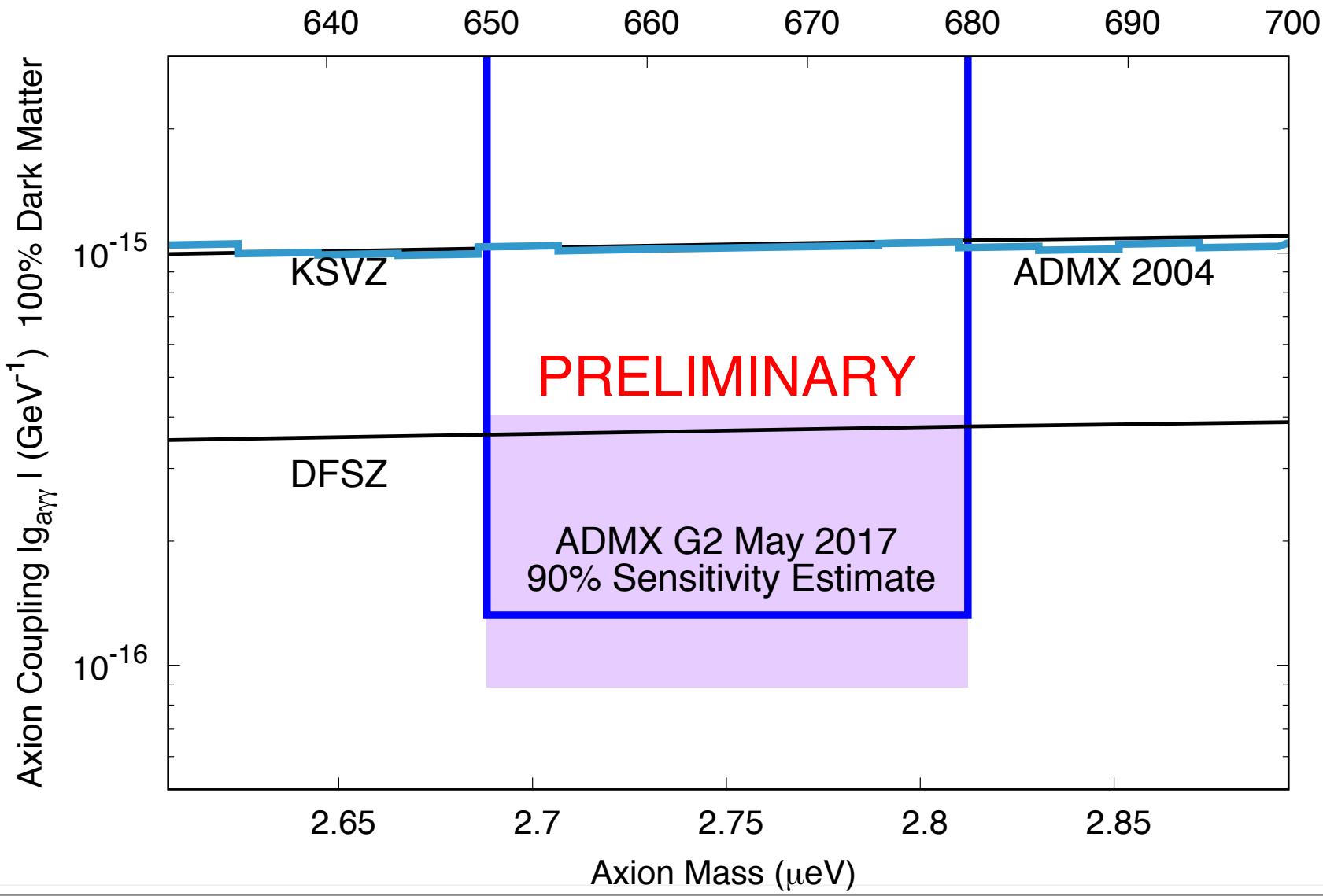
Further Offline Analysis

- Vary the bin size from time-series data.
- High Res analysis for ultra-sharp lines.

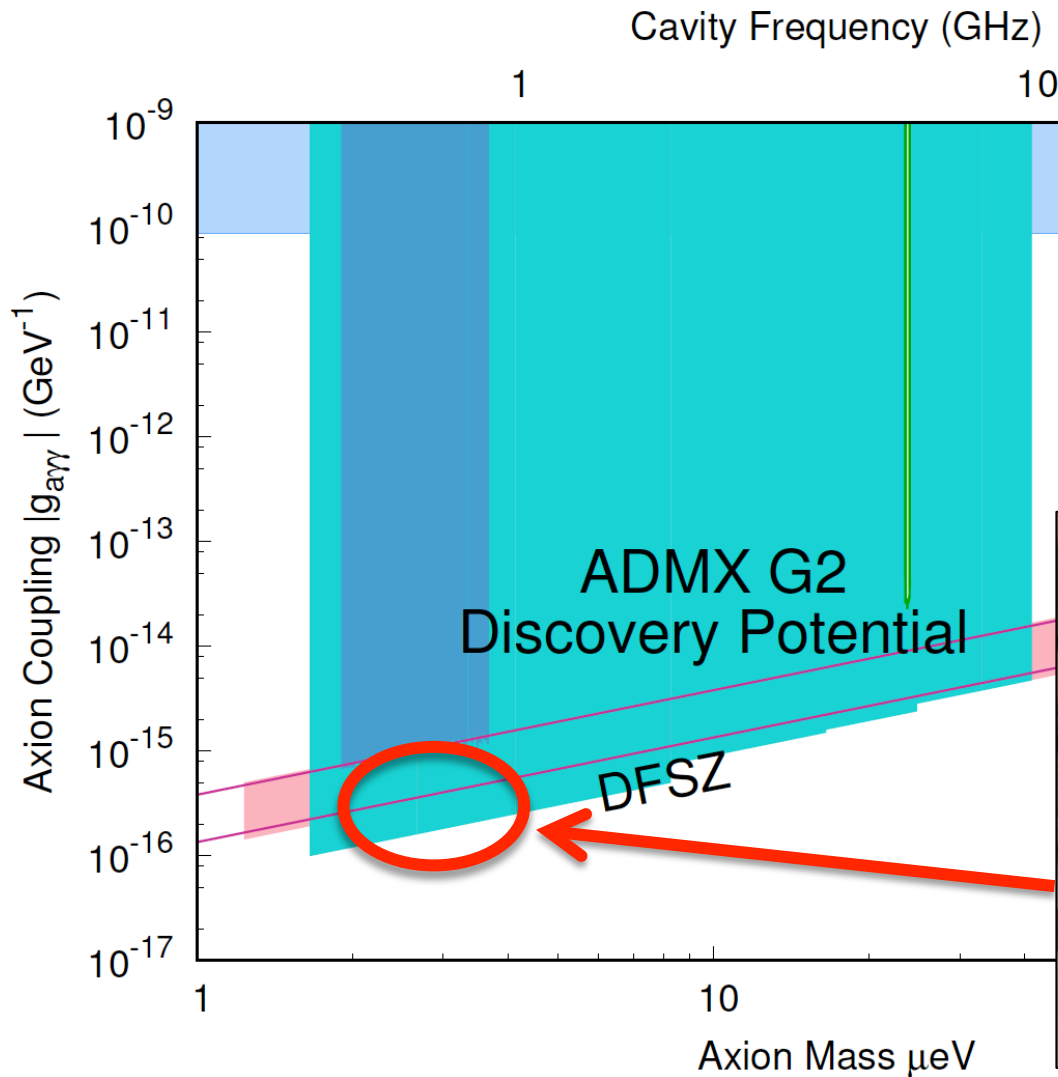
Sensitivity of recent data run

(January – June, 2017)

Cavity Frequency (MHz)



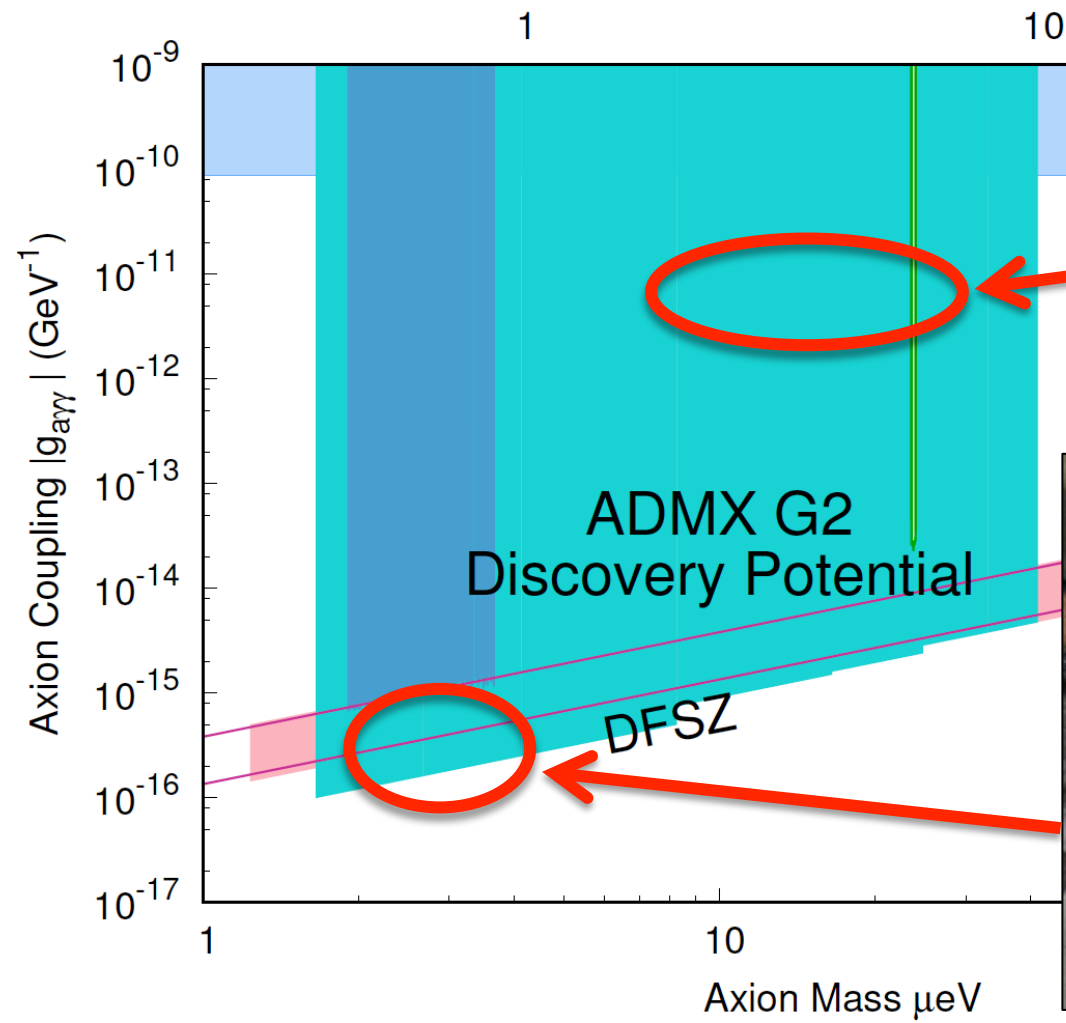
ADMX Main Cavity: Initial run 0.65-1 GHz



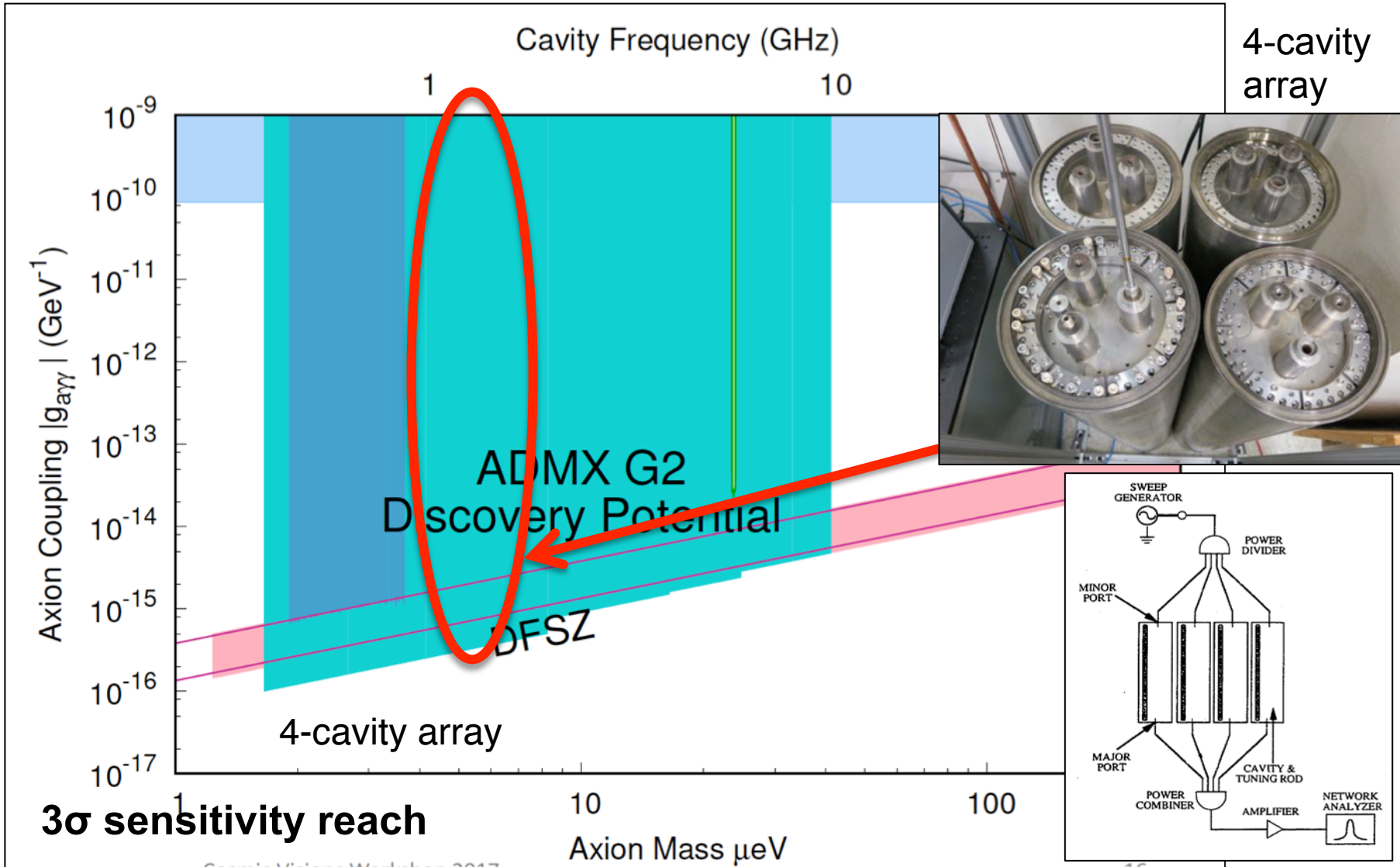
ADMX Sidecar Cavity

4-6 GHz TM_{010} & 6-7 GHz TM_{020}

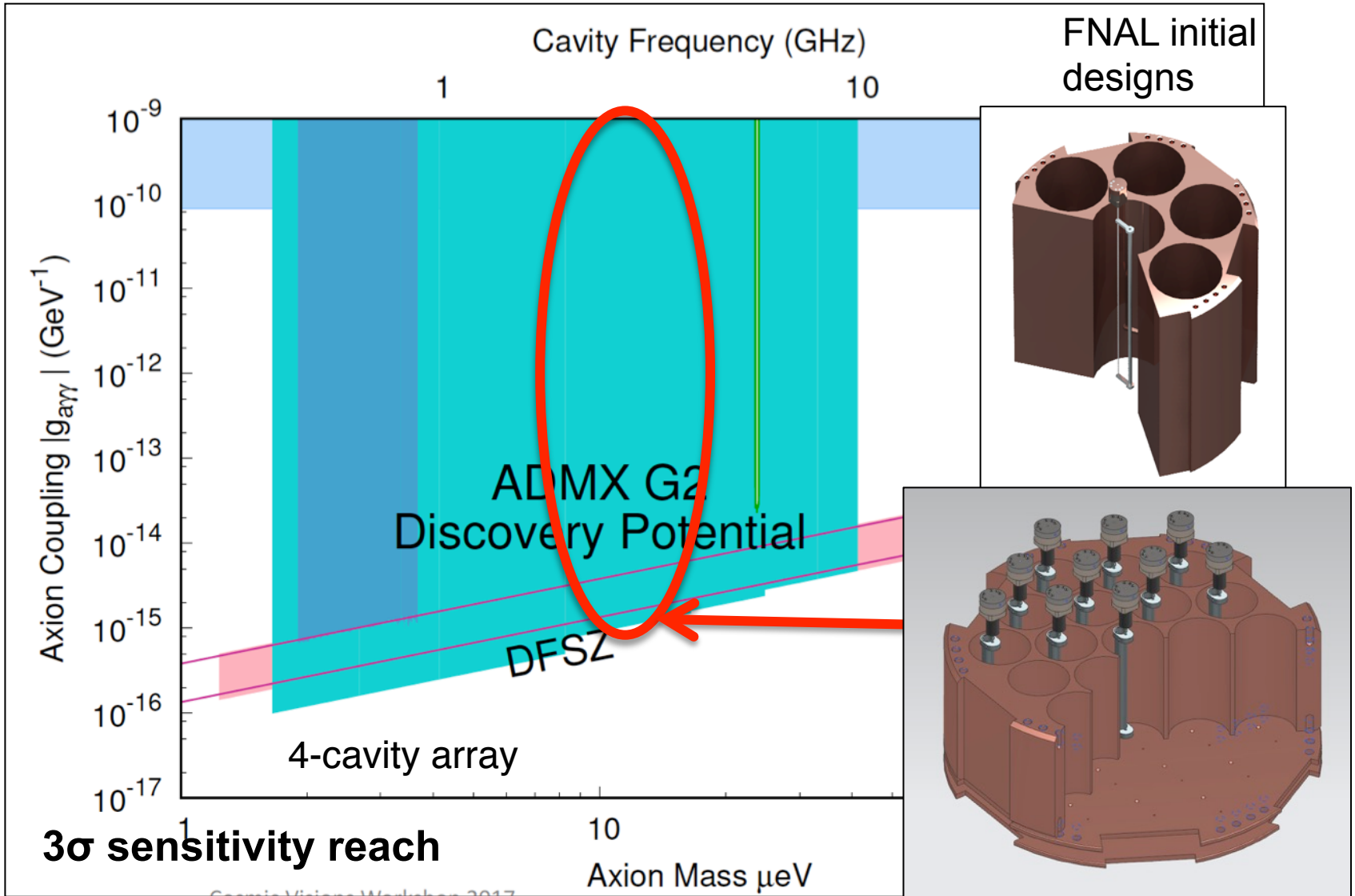
Completely separate system installed above main cavity
Cavity Frequency (GHz)



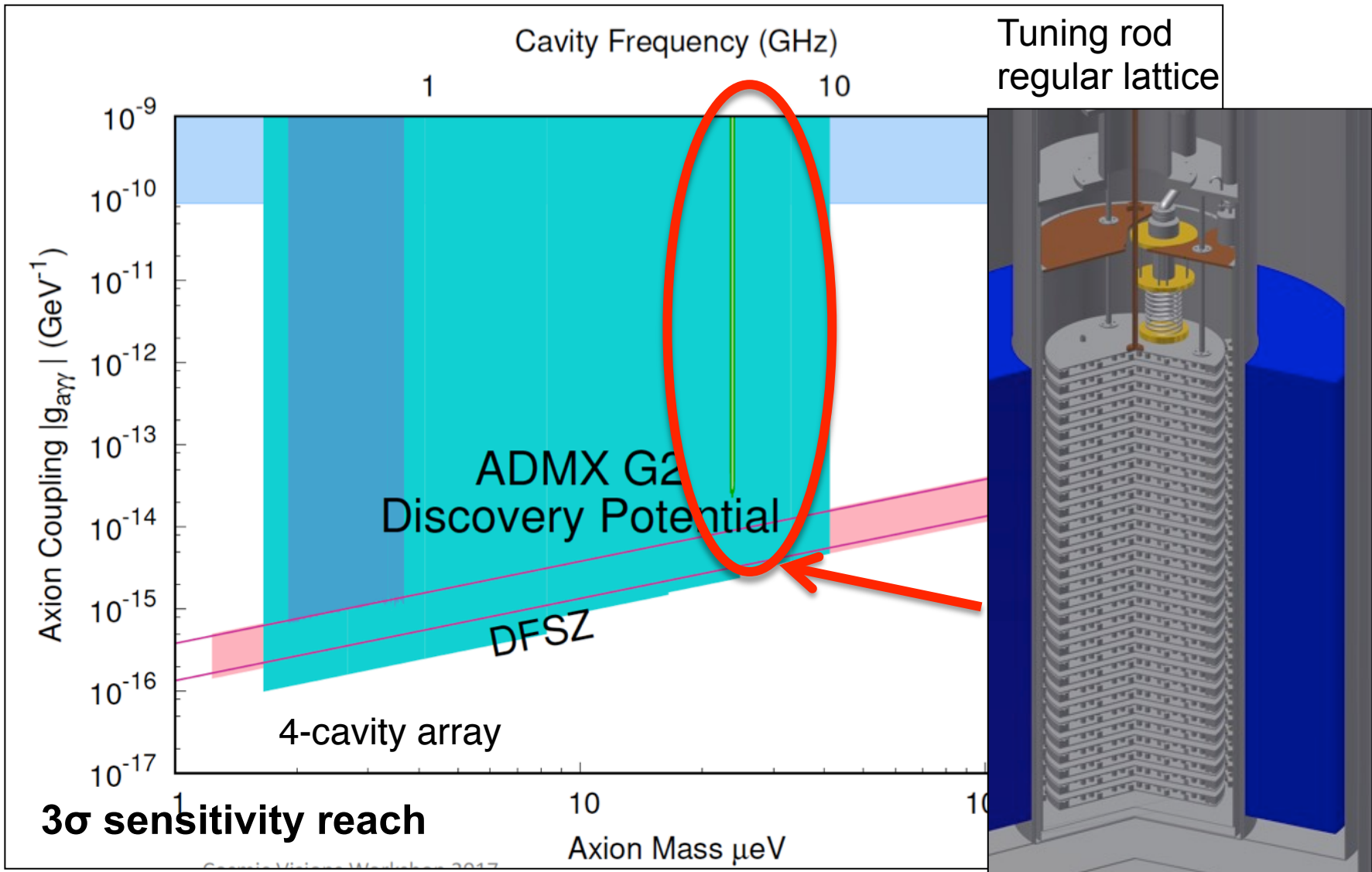
ADMX Science Prospects (1-2 GHz)



ADMX Science Prospects (2-6 GHz)

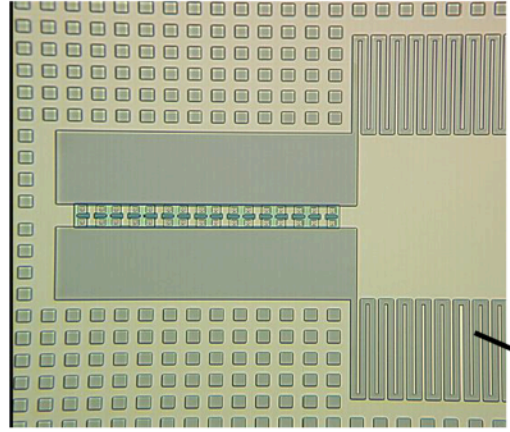


ADMX Science Prospects (6-10 GHz)

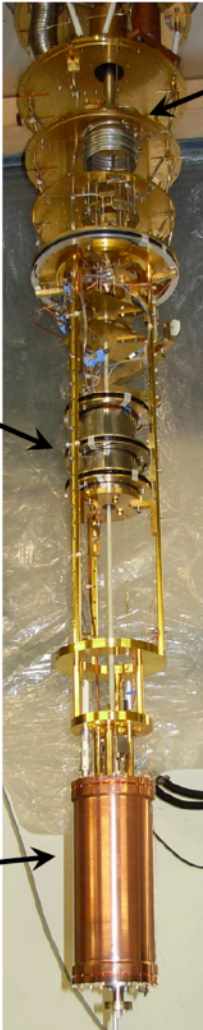


Other groups: HAYSTAC (formally ADMX-High Frequency) 9.4 Tesla magnet and 1.5 liter cavity at Yale U.

Josephson Parametric Amplifier



Microwave Cavity (copper)



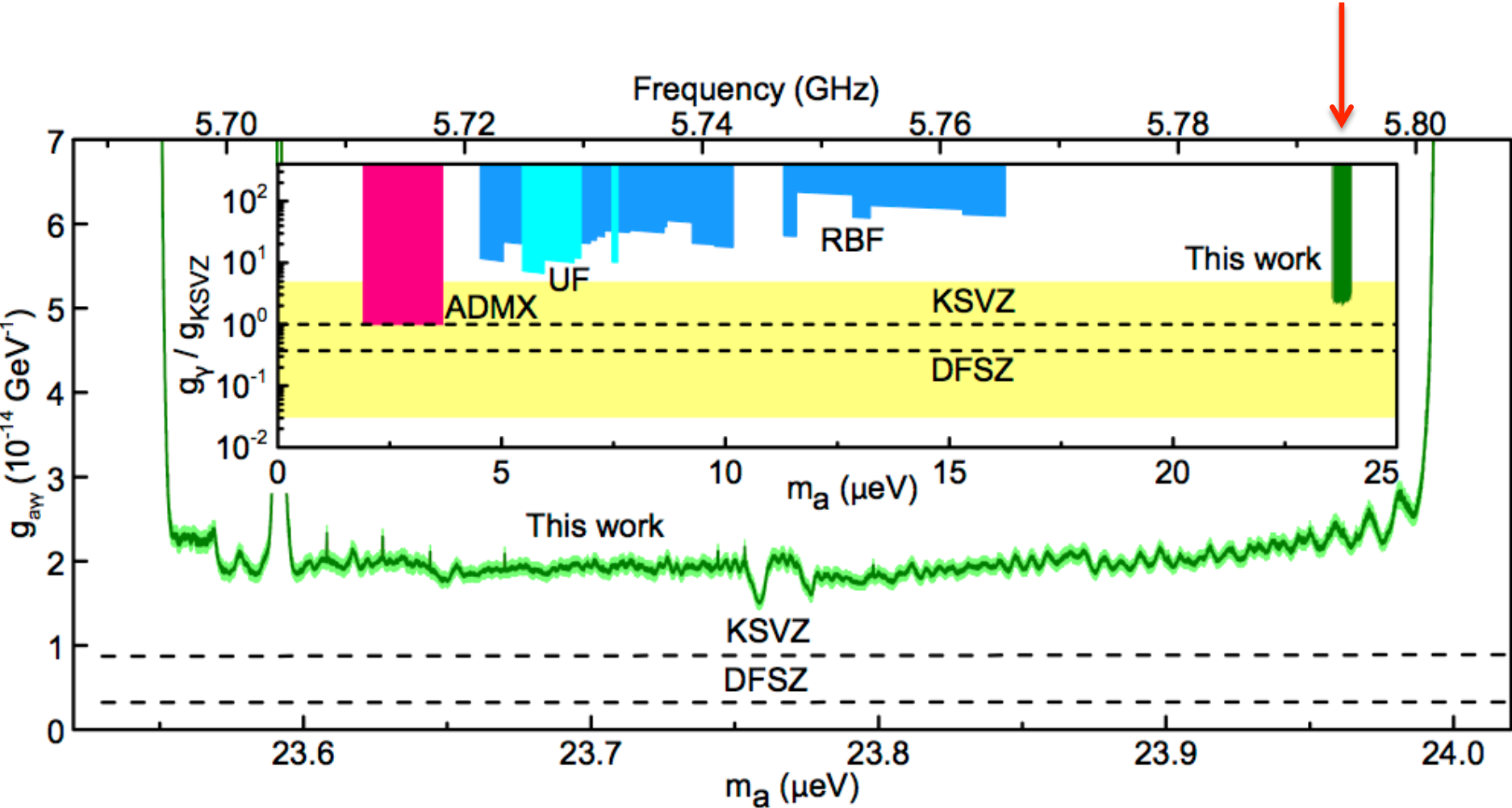
$^3\text{He}/^4\text{He}$ Dilution Refrigerator



9.4 Tesla, 10 Liter Magnet



HAYSTAC recent results

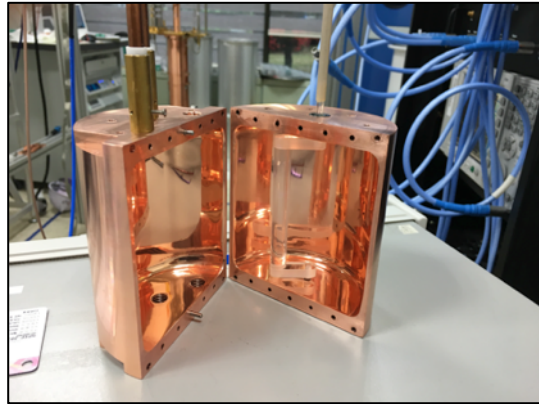


Results: B.M. Brubaker et al., Phys. Rev. Lett. 118 (2017) 061302.
 Design details: S. Al Kenany et al., Nucl. Instrum. Methods A 854 (2017) 11-24.

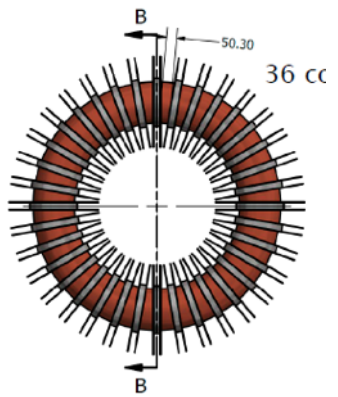
Other Haloscopes coming online

Center for Axion and Precision Physics (CAPP) in South Korea

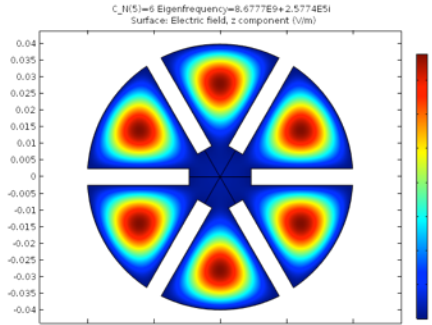
Bring online multiple experimental efforts over the next few years including microwave cavity searches (CULTASK, Toroid & multi-cell cavities)



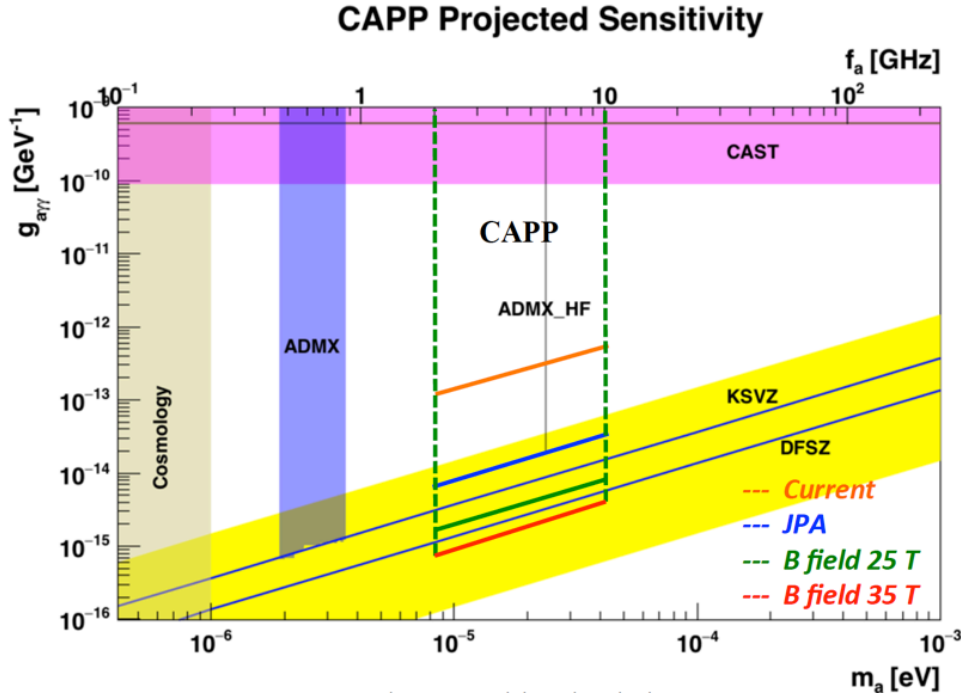
Prototype copper cavity



Toroid cavity



Multi-cell cavities



Other Haloscopes coming online

ORGAN experiment

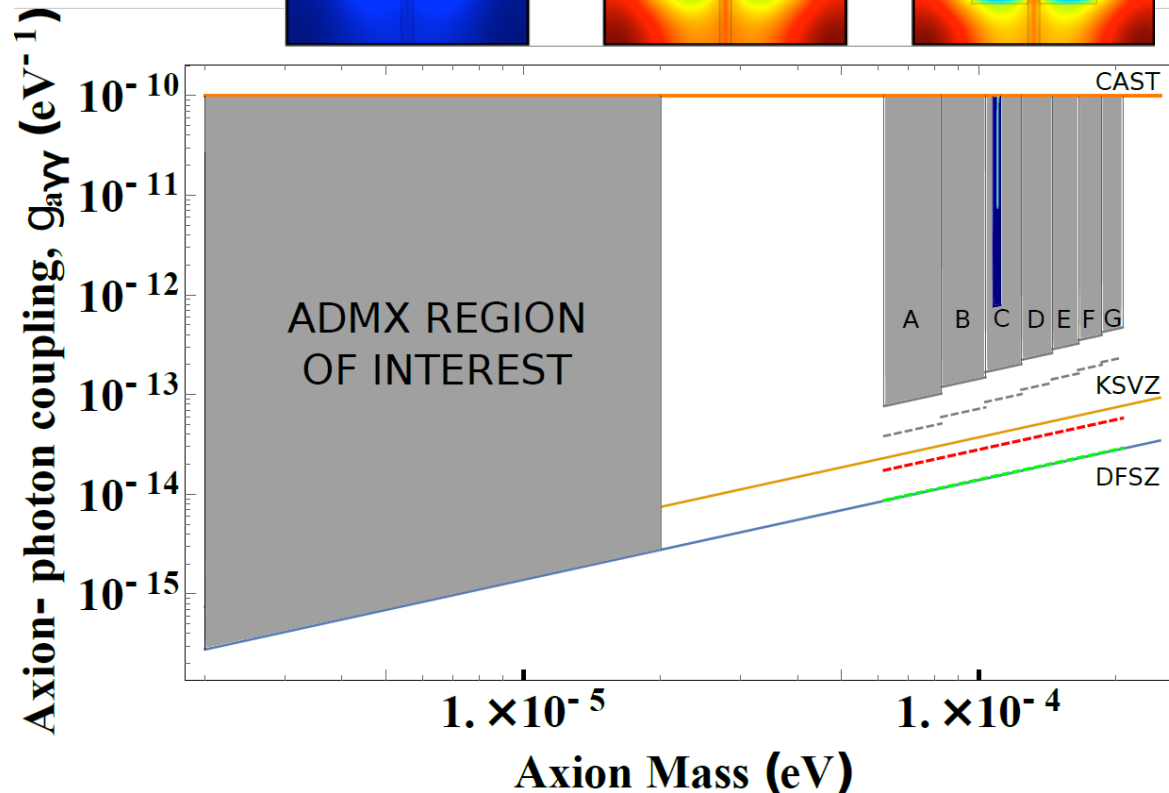
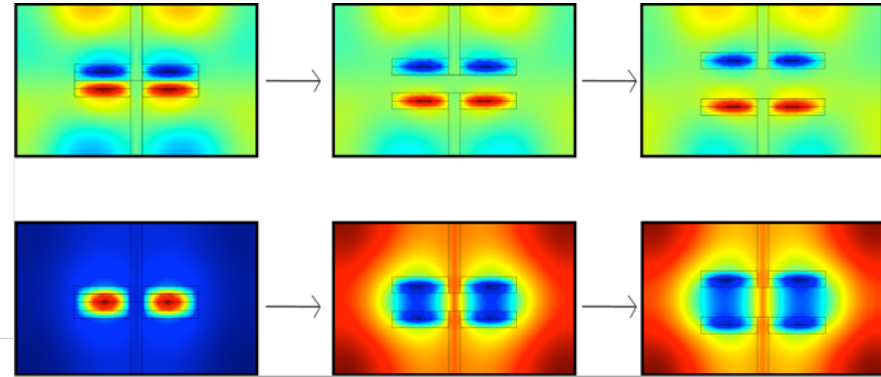
Oscillating Resonant Group AxioN Experiment
(U. of Western Australia)

Exploring new cavity geometries and modes with sapphire disks.

Initial experiments aimed at 26-27 GHz.

Runs A→G are the 2018-2025 runs, with 14 T magnet and SQL Amps

Dashed lines rely on success of squeezed state amplifiers and magnet upgraded to 28 T

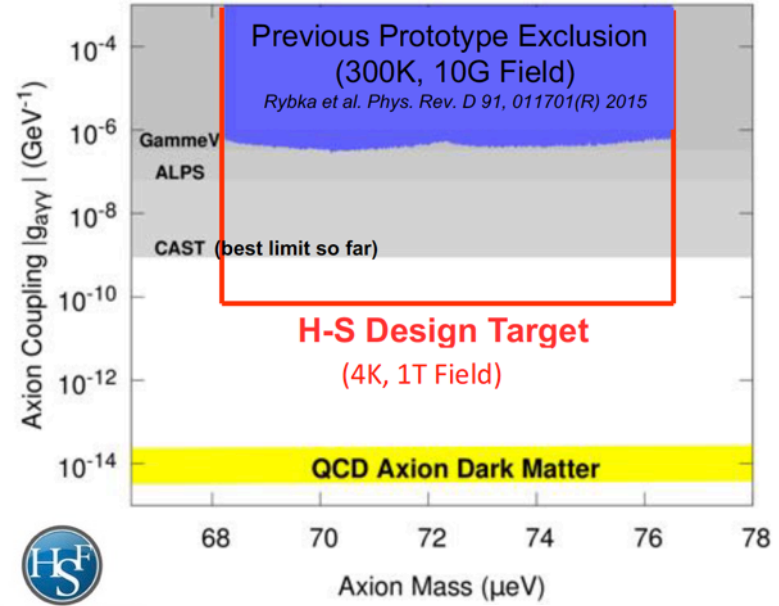
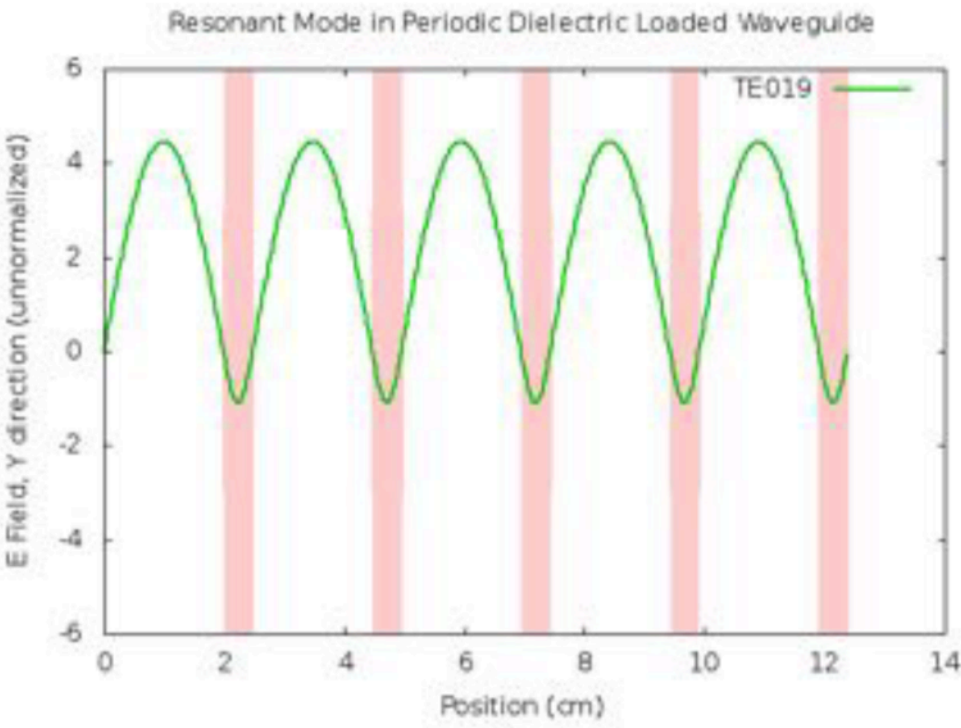
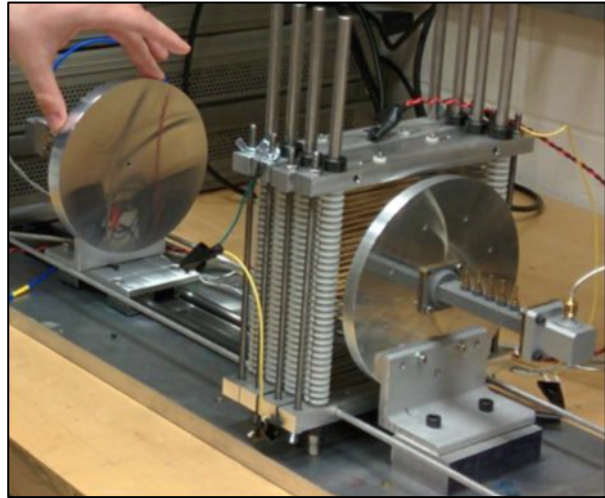


Open resonator design with dipole magnet

Orpheus Project (UW)

Open resonator would usually not couple to axion field (positive and negative E-fields cancel).

Manipulating modes with dielectrics or alternating the magnetic field leads to a net axion coupling.

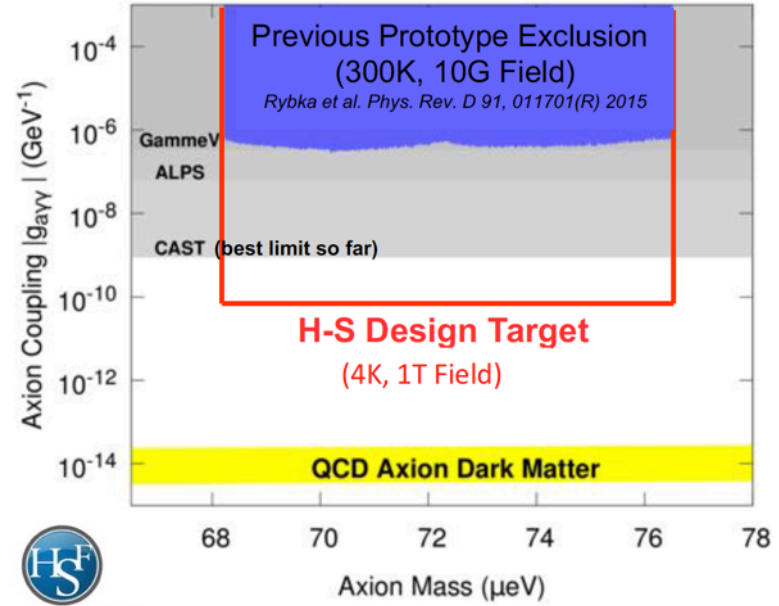
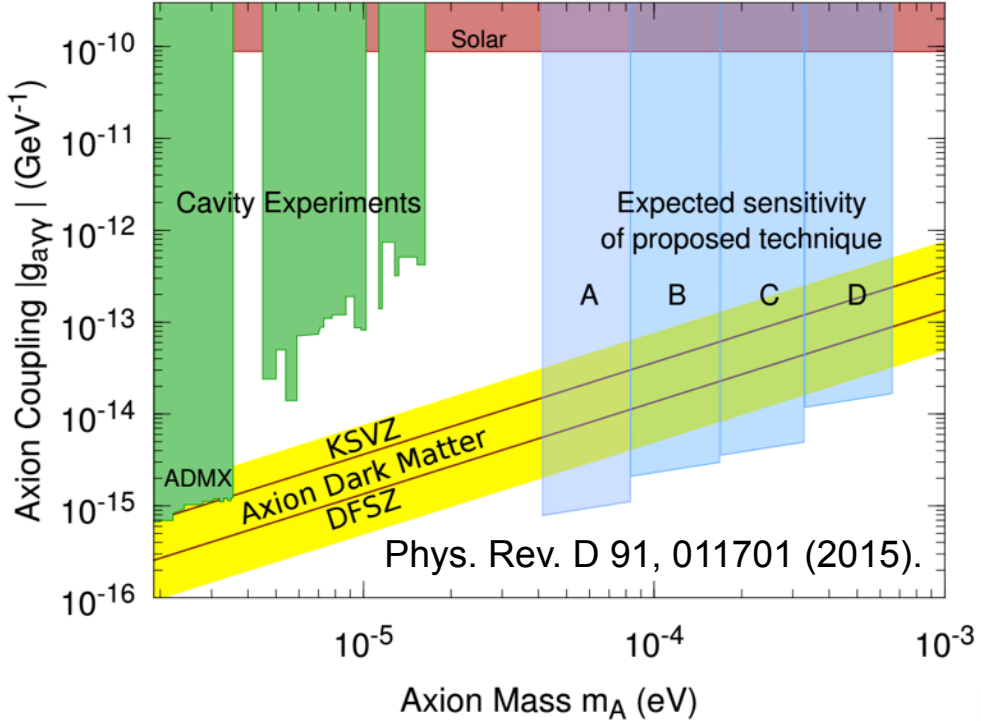
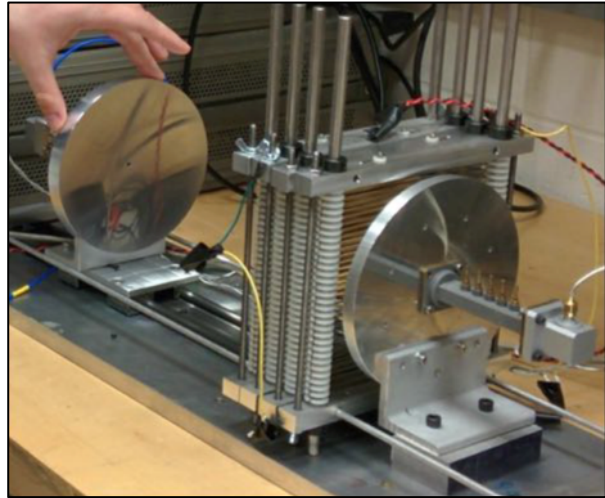


Open resonator design with dipole magnet

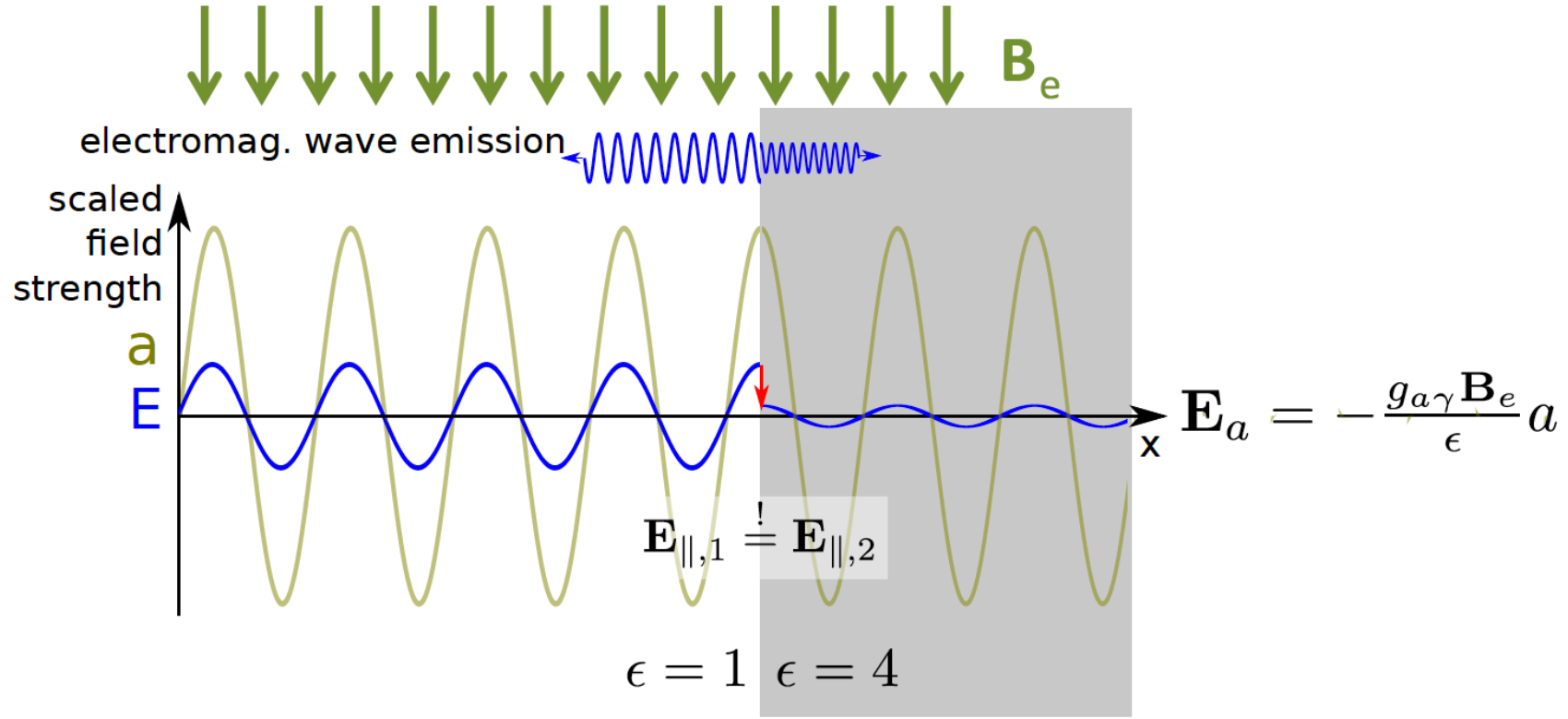
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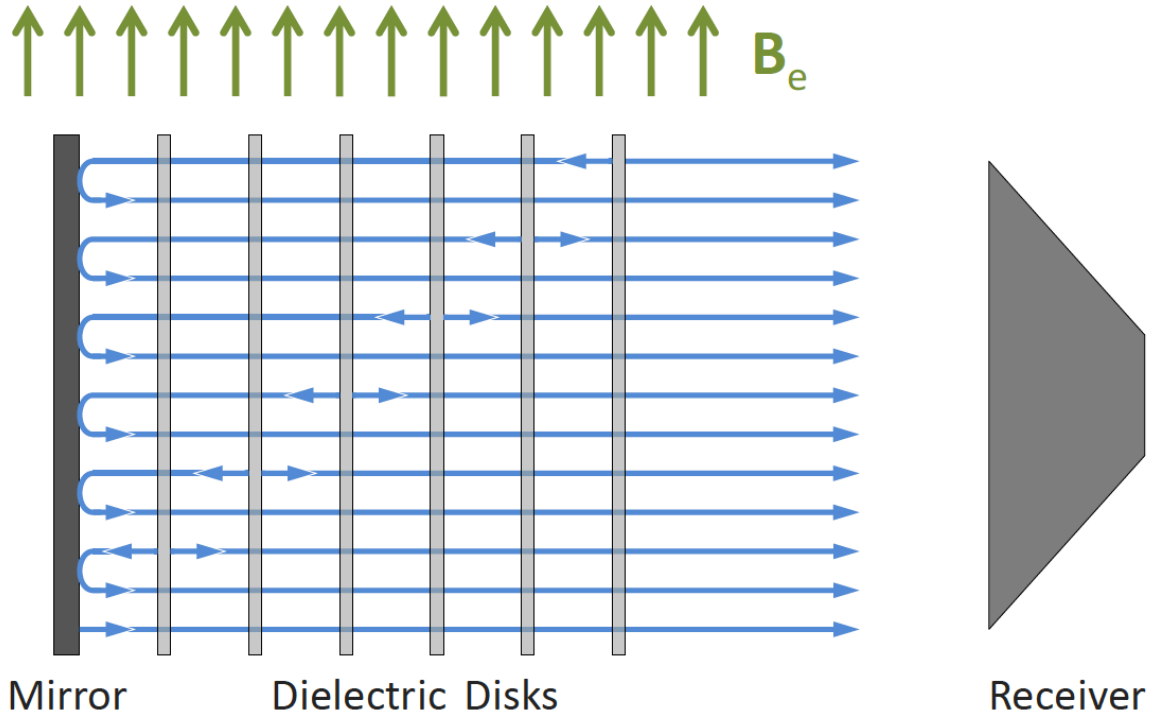
Broadband design with dipole magnet: MADMAX



$$P/A = 2.2 \times 10^{-27} \text{ W m}^{-2} \left(\frac{B_e}{10 \text{ T}} \right) C_{a\gamma}^2 \cdot f(\epsilon_1, \epsilon_2)$$

Broadband design with dipole magnet: MADMAX

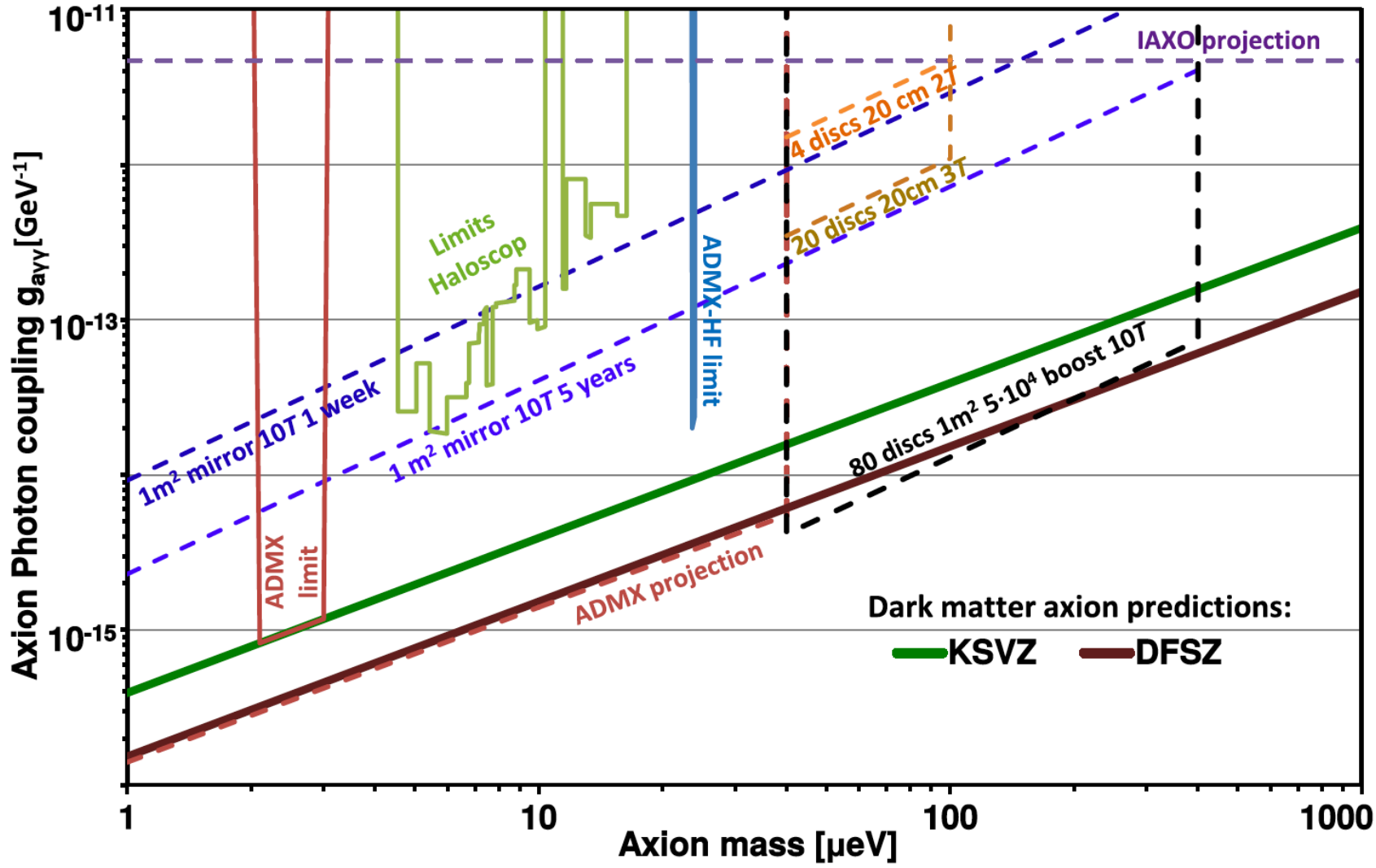
- Multiple dielectrics can constructively interfere



$$P/A = 2.2 \times 10^{-27} \text{ W m}^{-2} \left(\frac{B_e}{10 \text{ T}} \right)^2 C_{a\gamma}^2 \cdot \beta^2$$

β^2 : power emitted by booster / power emitted by single mirror ($\epsilon = \infty$)

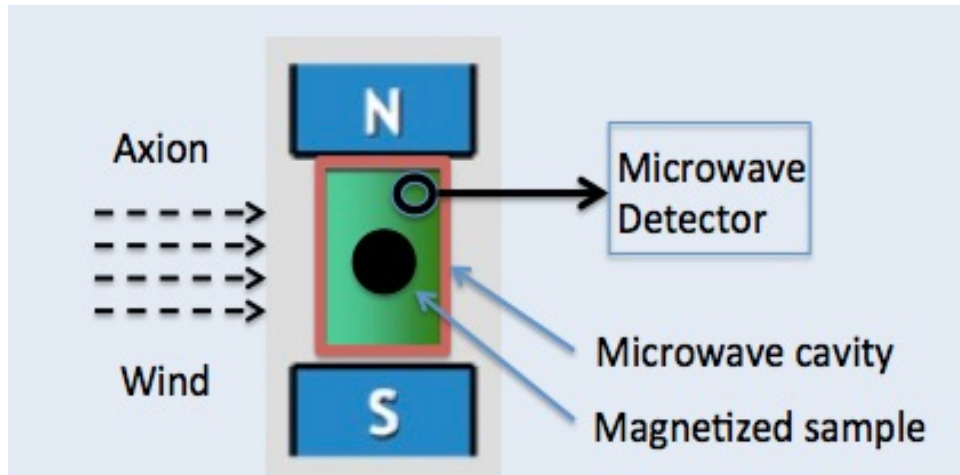
MADMAX sensitivity projections



Bela Majorovits and Stefan Knirck gave more detailed talks yesterday (July 24th)

Alternative strategy for high mass axions

- Look for an axion “wind” which acts as an effective RF magnetic field on electron spin via **electron-axion coupling**
- This axion induced RF excites **magnetic transition in a magnetized sample** (Larmor frequency) and produces a detectable signal
- The **QUAX** (QUest for AXion) experiment



The effective magnetic field associated with the axion wind

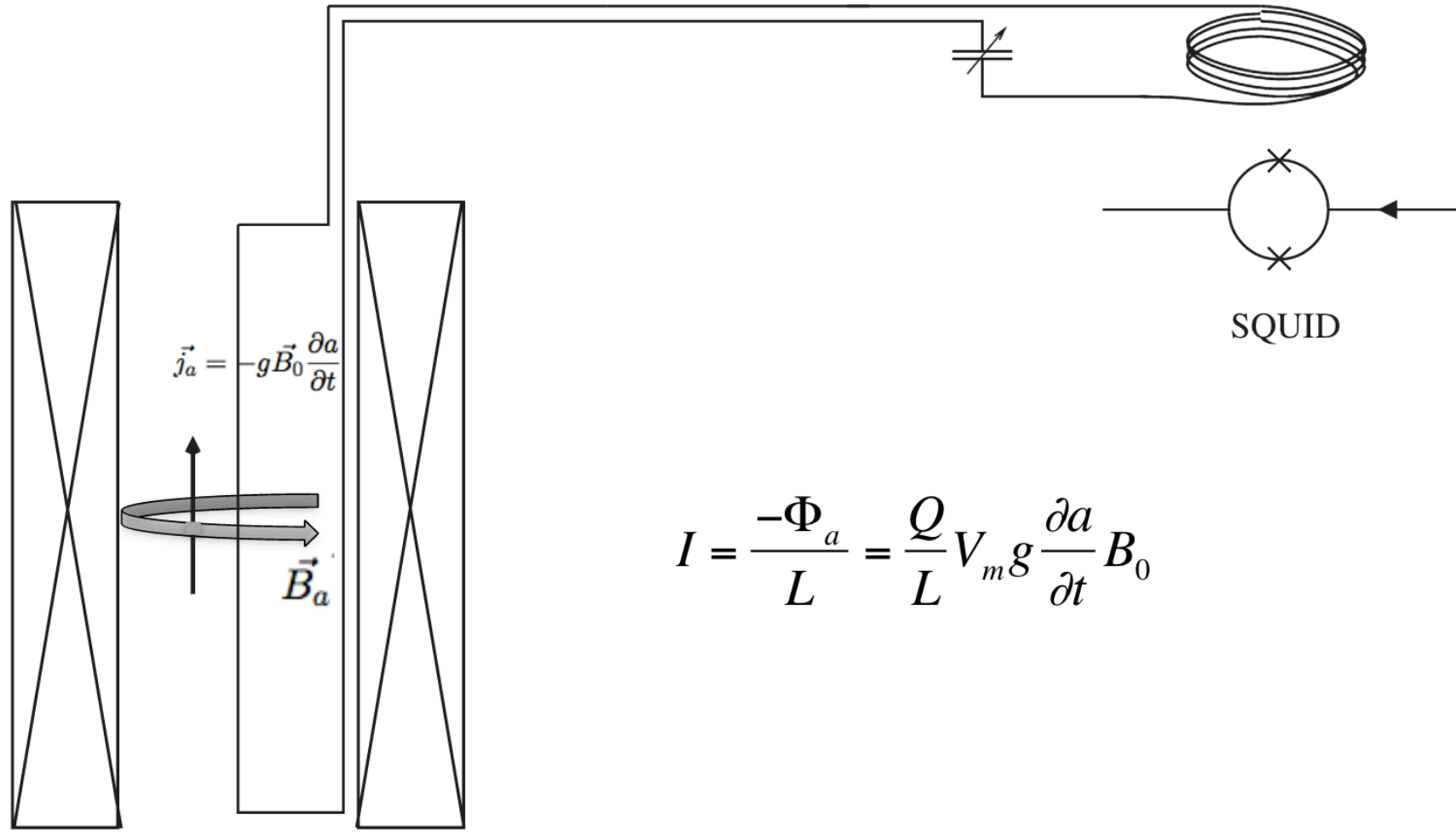
$$B_a = \frac{g_p}{2e} \left(\frac{n_a h}{m_a c} \right)^{1/2} m_a v_E$$

$$B_a = 2.0 \cdot 10^{-22} \left(\frac{m_a}{200 \mu\text{eV}} \right) \text{ T,}$$

$$\frac{\omega_a}{2\pi} = 48 \left(\frac{m_a}{200 \mu\text{eV}} \right) \text{ GHz,}$$

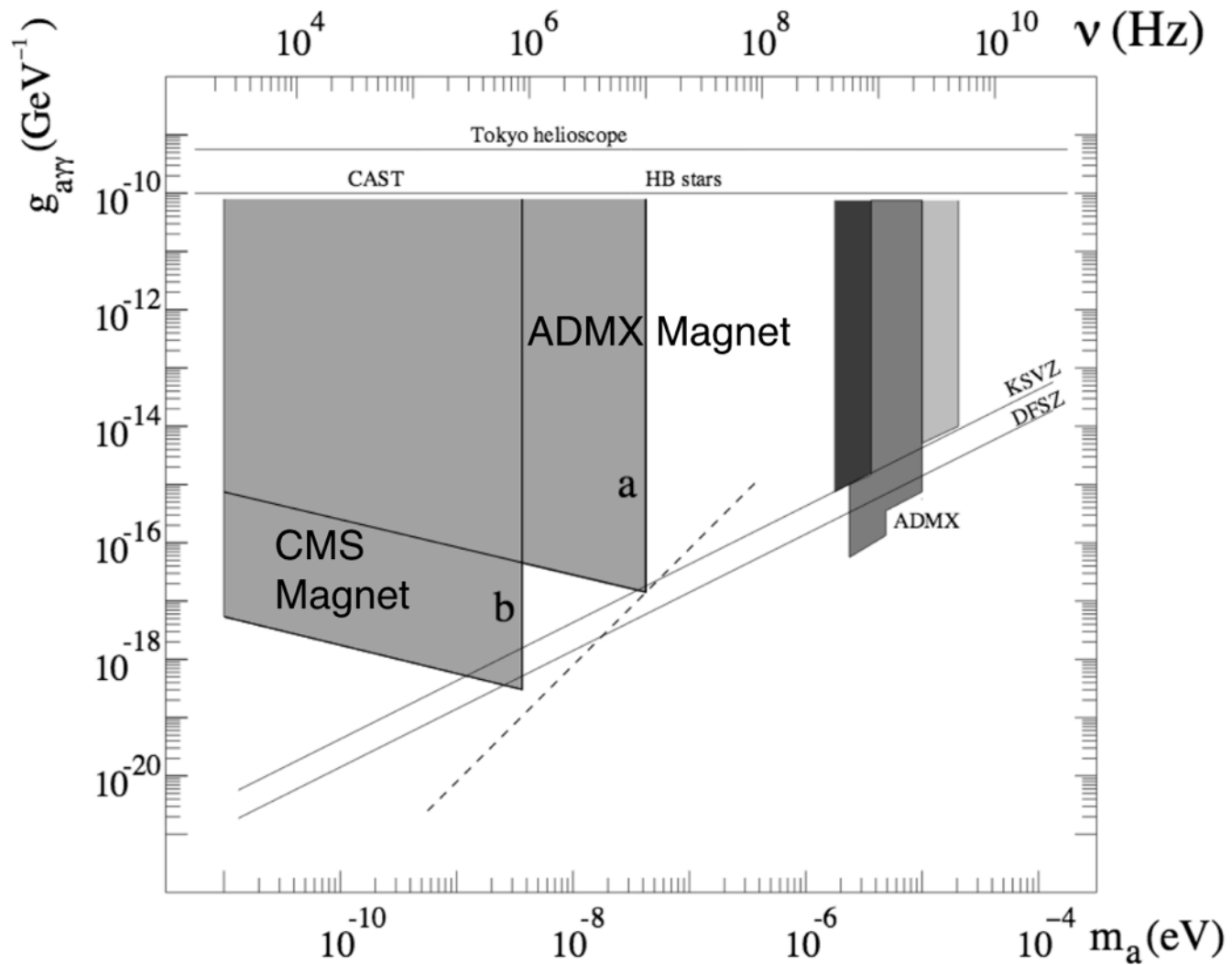
- R. Barbieri et al., *Searching for galactic axions through magnetized media: The QUAX proposal* Phys. Dark Univ. **15**, 135 - 141 (2017)

Going to lower masses (< 1 μeV): LC Circuit



Proposal for Axion Dark Matter Detection Using an LC Circuit
 PRL 112, 131301 (2014)

LC Circuit projections using various magnets

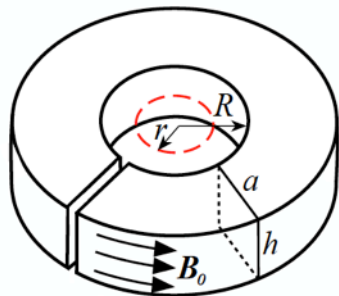


Proposal for Axion Dark Matter Detection Using an LC Circuit
 PRL 112, 131301 (2014)

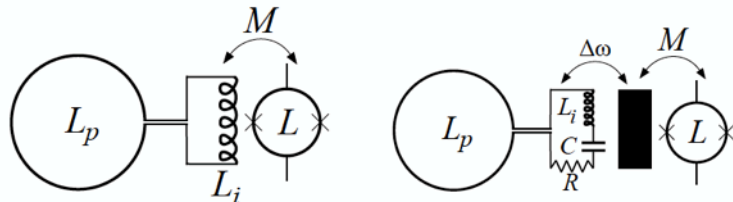
ABRACADABRA Experiment

A Broadband/Resonant Approach to Cosmic Axion Detection with an Amplifying B-field Ring Apparatus

Theory:



Toroidal geometry for zero-field detection

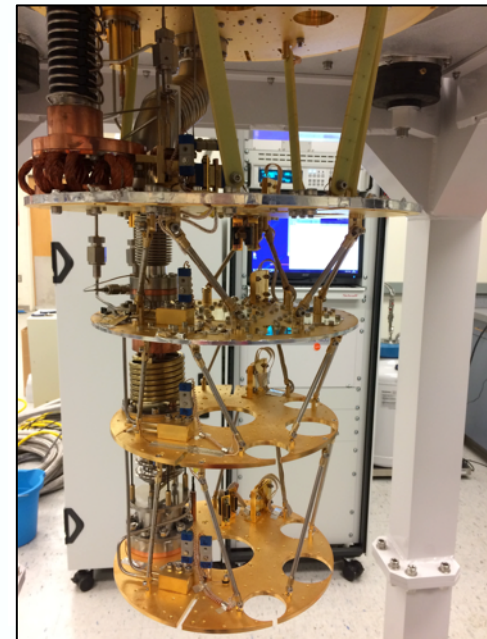
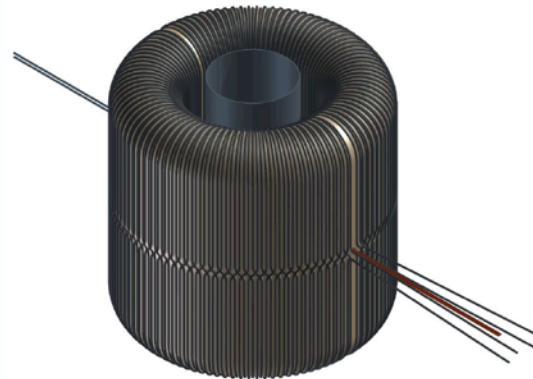


Interchangeable readout:
broadband (low freq.) or
resonant (high freq.)

Experiment:

Prototype specs:

$$R_{in} = 3 \text{ cm}, R_{out} = 6 \text{ cm}, h = 12 \text{ cm}, \\ V = 680 \text{ cm}^3, B_{max} = 1 \text{ T}, G = 0.085$$



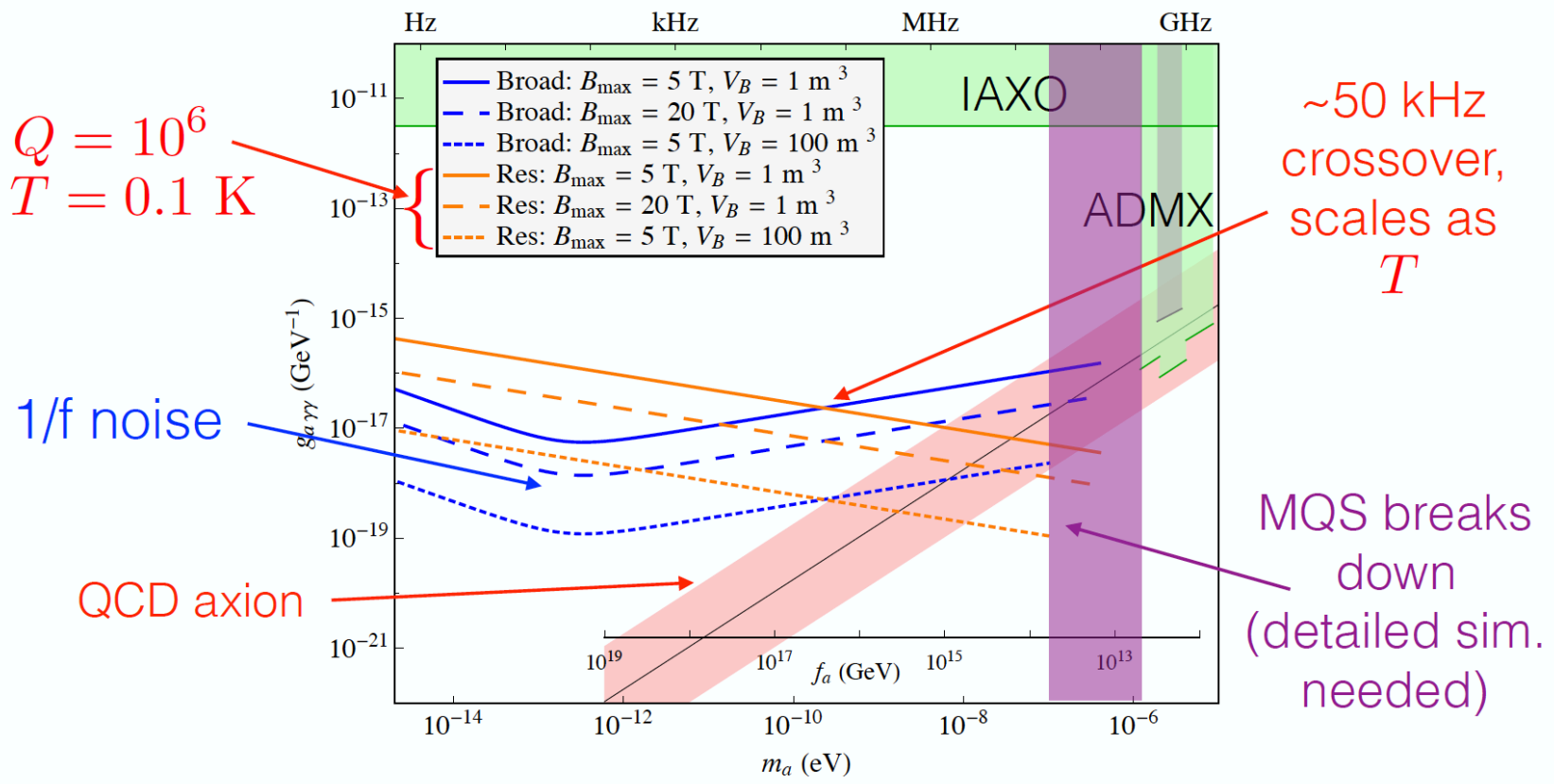
ABRA-10cm @ MIT

YK, Safdi, Thaler, Phys. Rev. Lett. 2016

ABRACADABRA sensitivity projections

1 year **total** measurement time

$$\nu = m_a / 2\pi$$



With same experimental parameters,
 broadband for low frequencies, resonant for high frequencies

YK, Safdi, Thaler, Phys. Rev. Lett. 2016

Very low mass axions (neV)

NMR based experiment: CASPEr

$$\frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu}$$



Coupling to electromagnetic field

$$\frac{a}{f_a} G_{\mu\nu} \tilde{G}^{\mu\nu}$$



Coupling to gluon field

CASPEr Electric

Budker D, et al. Phys. Rev. X4:021030 (2014)

$$\frac{\partial_\mu a}{f_a} \bar{\Psi}_f \gamma^\mu \gamma_5 \Psi_f$$



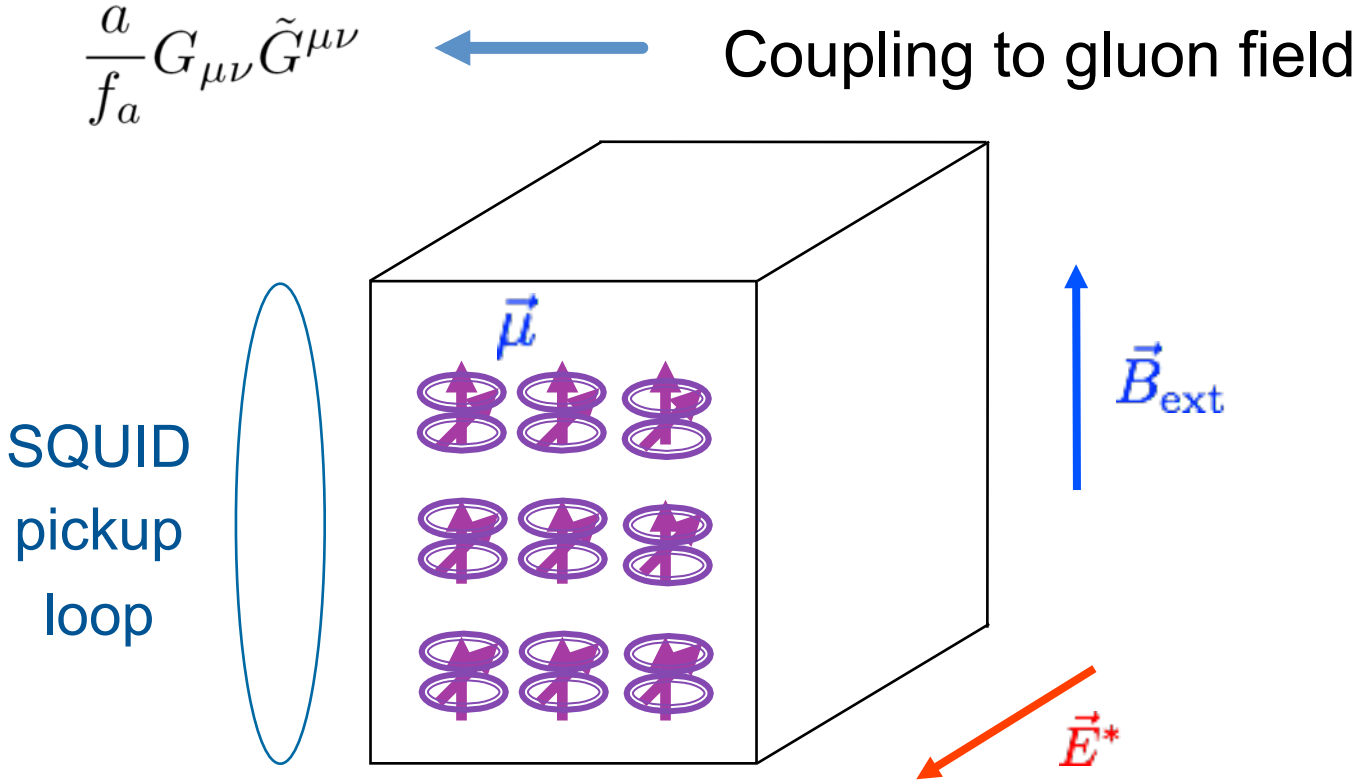
Coupling to fermions

CASPEr Wind

Graham PW, Rajendran S. Phys. Rev. D88:035023 (2013)

CASPER-Electric Experiment

EDM coupling to axion plays role of oscillating transverse B-field



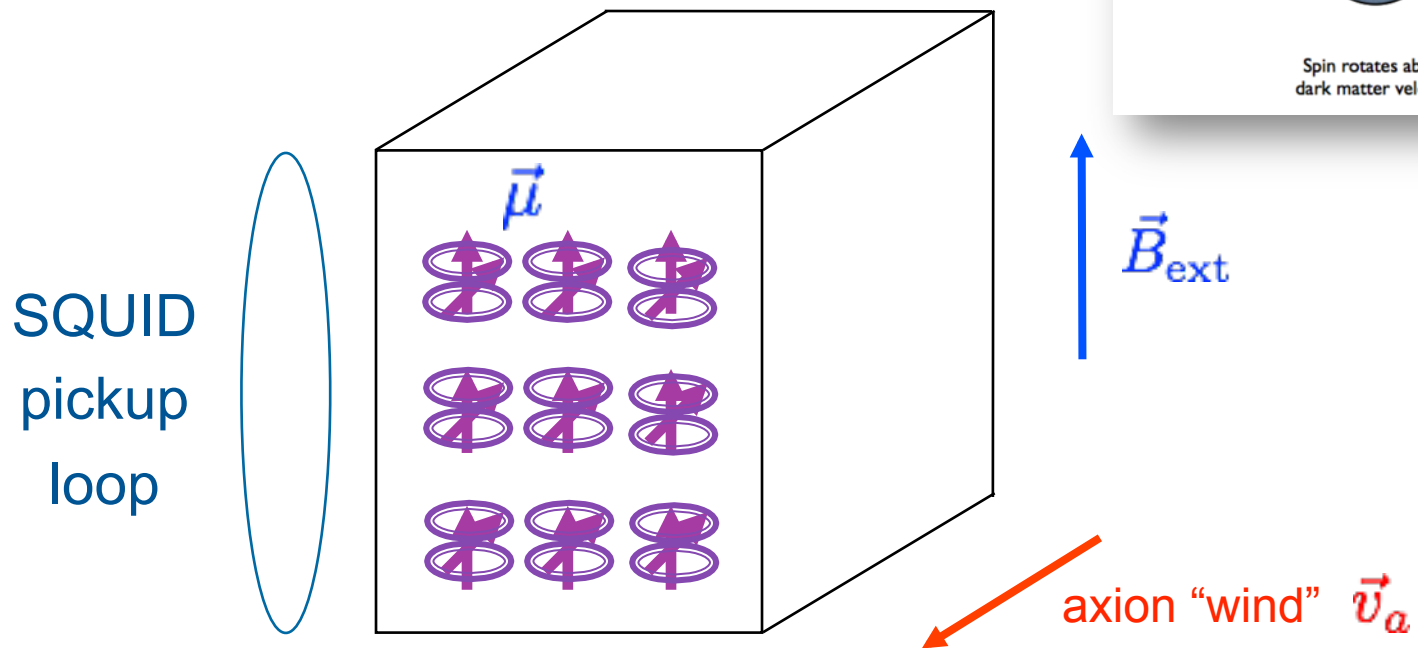
Larmor frequency = axion Compton frequency
→ resonant enhancement.

Axion/ALP-induced spin precession (axion wind)

CASPER-Wind

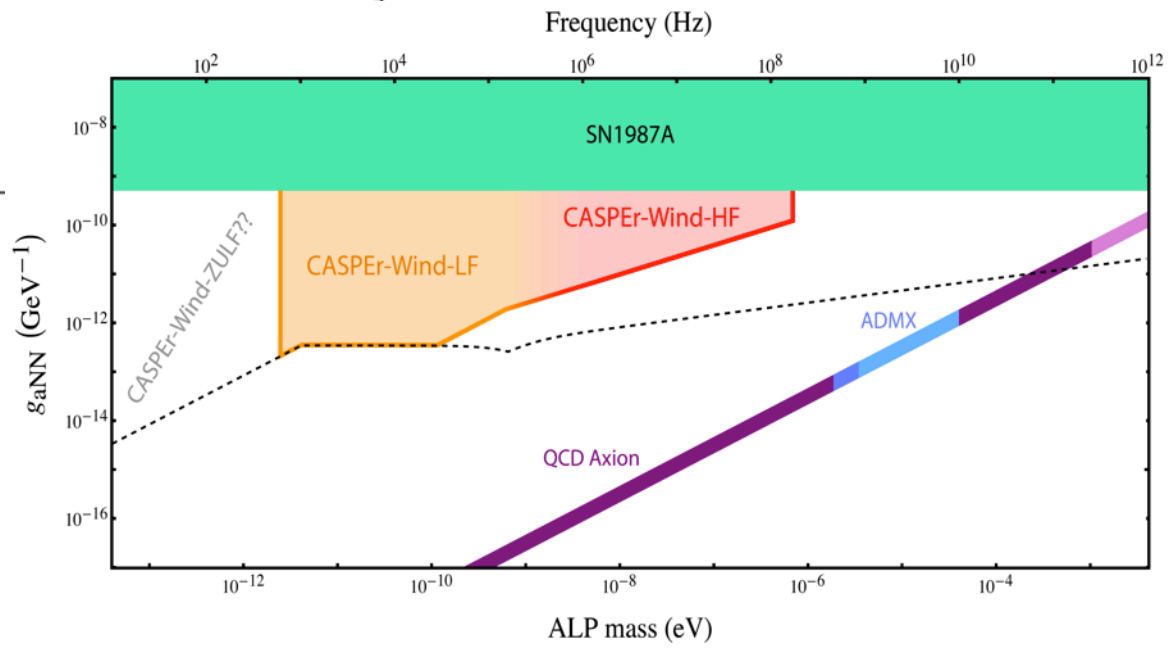
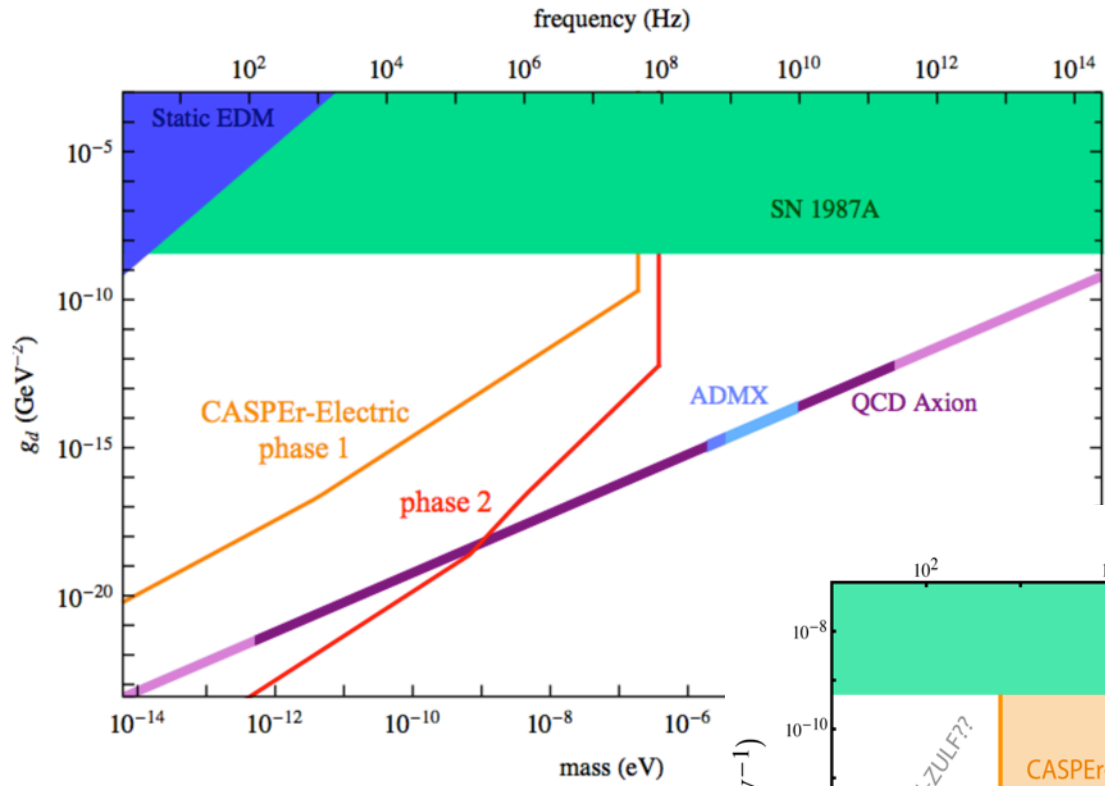
Nonrelativistic limit of the axion-fermion coupling yields a Hamiltonian:

$$H_{\text{wind}} \approx g_{aNN} \nabla a \cdot \boldsymbol{\sigma}_N .$$



Larmor frequency = axion Compton frequency
→ resonant enhancement.

CASPER Anticipated Sensitivities

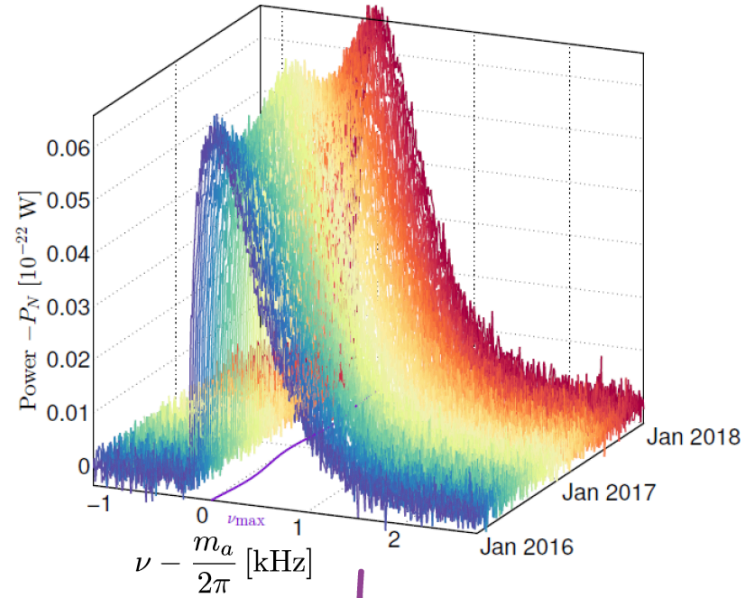


*slides from Derek Kimball

Studying the axion!

Studying the axion is studying the halo!

O'Hare & Green [1701.03118]

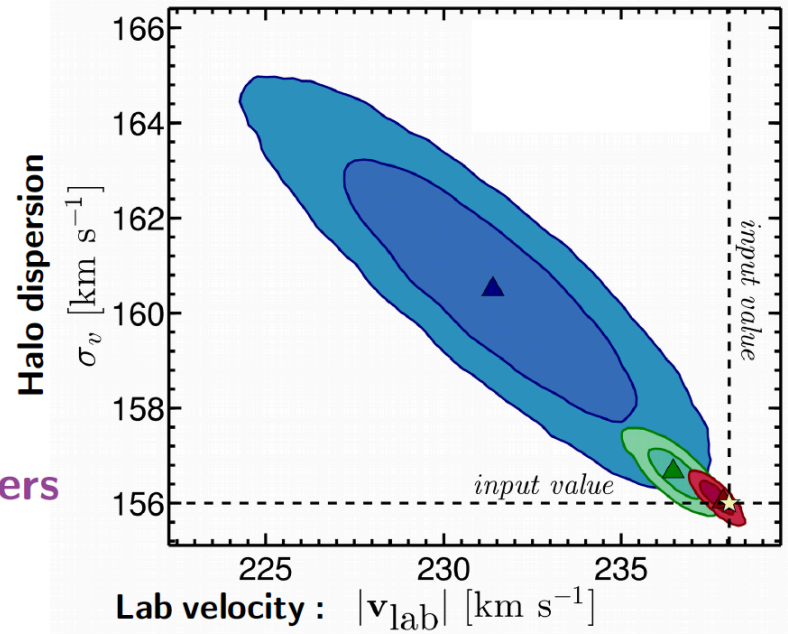


Max likelihood fit:
 → Extract astrophysical parameters

Accuracy < 1 km/s
 Better than astronomical determinations!

Experiment duration:
 (for QCD axion)

- $\tau_{\text{tot}} = 10$ days
- $\tau_{\text{tot}} = 0.5$ yr
- $\tau_{\text{tot}} = 1$ yr



Summary

Axions: solve the Strong-CP problem and are a compelling DM candidate

Haloscopes such as ADMX have sensitivity to DM axions:

ADMX Gen 2: Taking science data since Jan 2017

First experiment to reach sensitivity to DFSZ axions!

Anticipate continuous data taking up to 10 GHz

A number of other haloscopes taking data, coming online

HAYSTAC, CULTASK, ORGAN

New ideas for both high and low mass axion searches already in R&D phases

Orpheus, MADMAX, QUAX, LC-circuit, ABRACADABRA, CASPEr, etc

Information on galaxy dynamics to be learned from studying axion signal.

Discovery of the dark matter axion may be just around the corner!

Supported by DOE Grants DOE grant DE-SC00098000, DOE grant DE-SC0011665, DE-AC52-07NA27344, DE-AC03-76SF00098, the Heising-Simons Foundation, and the Lawrence Livermore National Laboratory, Fermilab and Pacific Northwest National Laboratory LDRD programs. SQUID development was supported by DOE grant DE-AC02-05CH11231.

Backup slides

The CP Problem of Strong Interactions

Characterizes degenerate QCD ground state (Θ vacuum)

Phase of Quark Mass Matrix

Standard QCD Lagrangian contains a CP violating term

$$L_{CP} = -\frac{\alpha_s}{8\pi} \underbrace{(\Theta - \arg \det M_q)}_{0 \leq \bar{\Theta} \leq 2\pi} \text{Tr } \tilde{G}_{\mu\nu} G^{\mu\nu}$$

Induces a neutron electric dipole moment (EDM) much in excess of experimental limits

$$d_n \approx \bar{\Theta} 10^{-16} \text{ e cm} \approx \frac{\bar{\Theta}}{10^2} \mu_n < 3 \times 10^{-26} \text{ e cm}$$

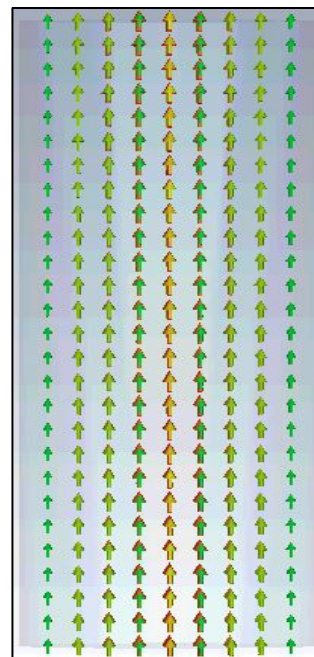
$$\bar{\Theta} \lesssim 10^{-10} \quad \text{Why so small?}$$

Key Microwave Cavity Design Constraints

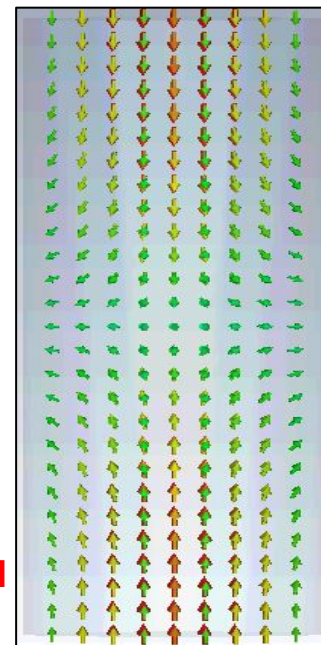
$$\frac{df}{dt} \approx 1.68 \text{ GHz/year} \left(\frac{g_\gamma}{0.36} \right)^4 \left(\frac{f}{1 \text{ GHz}} \right)^2 \left(\frac{\rho_0}{0.45 \text{ GeV/cc}} \right)^2 \cdot \left(\frac{5}{\text{SNR}} \right)^2 \left(\frac{B_0}{8 \text{ T}} \right)^4 \left(\frac{V}{100 \text{ l}} \right)^2 \left(\frac{Q_L}{10^5} \right) \left(\frac{C_{010}}{0.5} \right)^2 \left(\frac{0.2 \text{ K}}{T_{\text{sys}}} \right)^2$$

- Maximize product of $B^2 \cdot V \cdot Q_L \cdot C_{lmn}$ to maximize axion-to-photon conversion power
 - $B^2 V$ set by the magnet bore: $(8\text{T})^2 \cdot (\sim 100 \text{ liters})$
- Loaded Quality factor $Q_L = \text{frequency}/\text{bandwidth}$ ($Q_L \sim 10^5$ for copper cavity $\sim 1 \text{ GHz}$)
- Mode Form Factor C_{lmn}
- Tunability: must be able to shift resonant frequency over an appreciable range (typically 30-50%)

TM₀₁₀ mode
C₀₁₀ ~ 0.69



TM₀₁₁ mode
C₀₁₁ ~ 0.0



↑
B-field

Main Cavity Properties

Volume: 133 liters

Q_{loaded} : 60,000

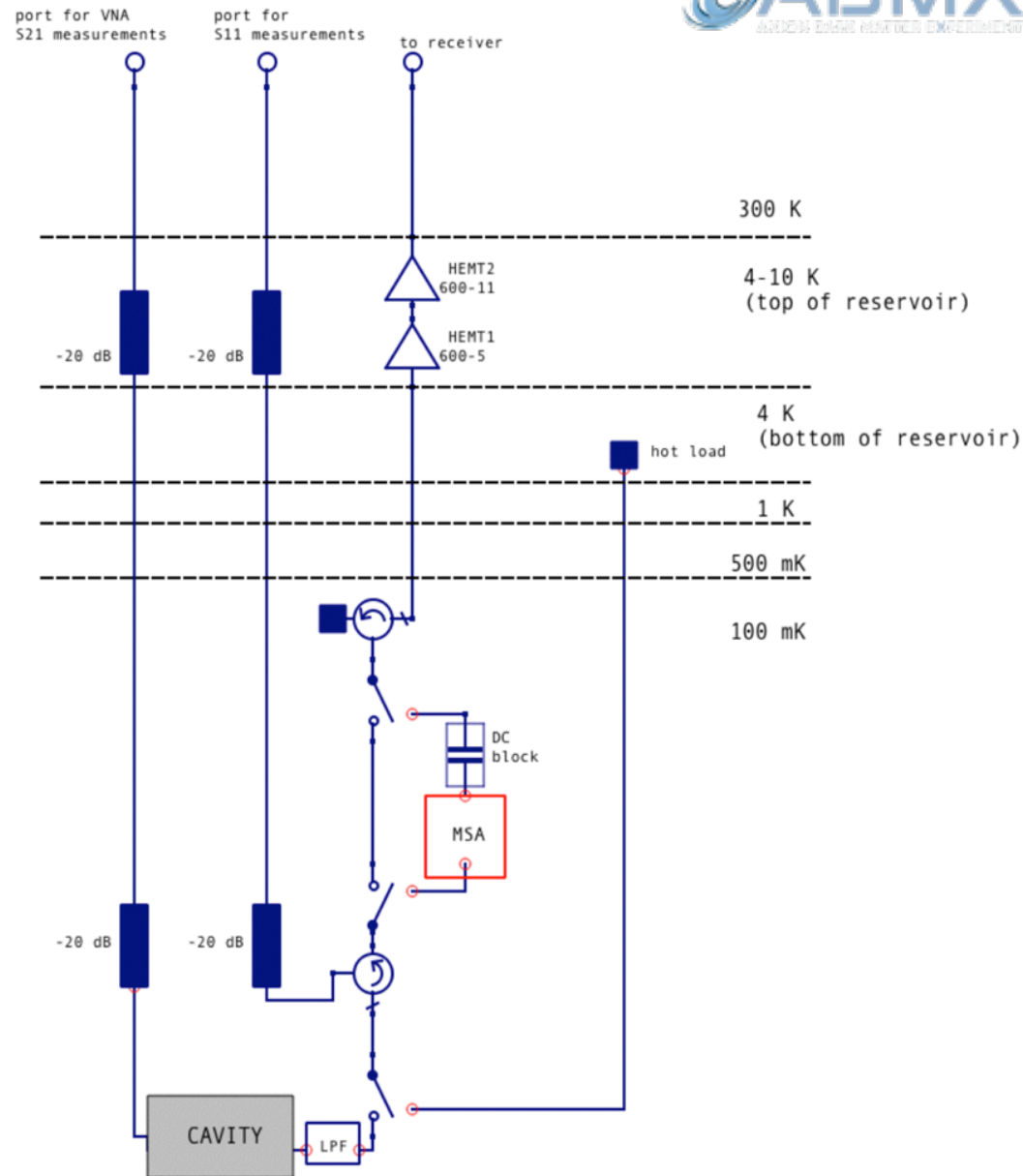
580 – 890 MHz

(with current rods)

Dilution Refrigerator
directly attached to top



Receiver Chain



Injection of swept power & fake axions

Reflection to look at antenna coupling

Hot / Cold load:

Measure system noise temperature

SQUID at $T_{\text{physical}} \sim 300 \text{ mK}$

Cavity at $T_{\text{physical}} \sim 150 \text{ mK}$

Total system noise $\sim 0.5 \text{ K}^*$

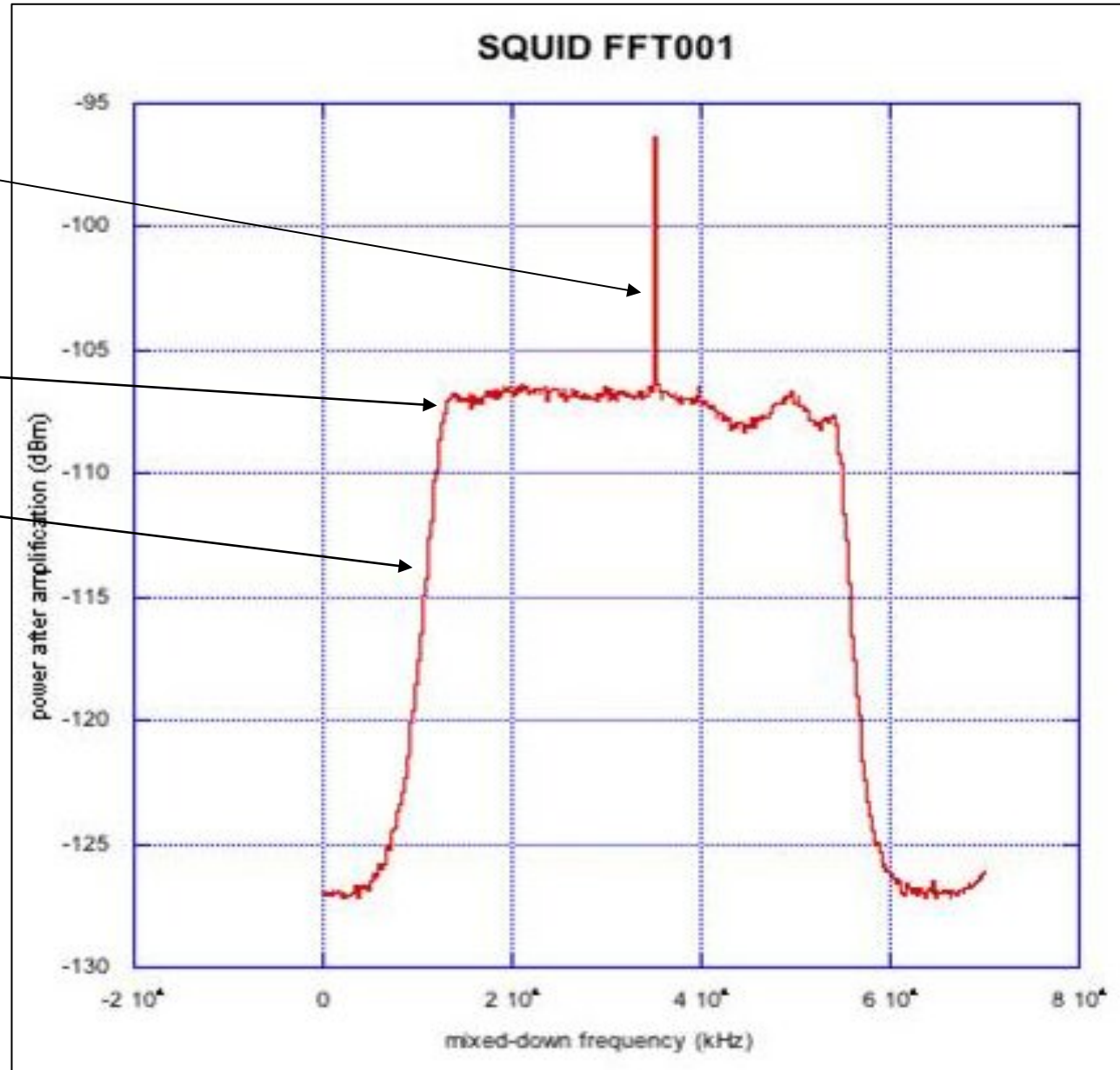
*includes attenuation + post-amplifier contributions.

Example of injected tone (fake axion) (Phase I with SQUID amplifiers)

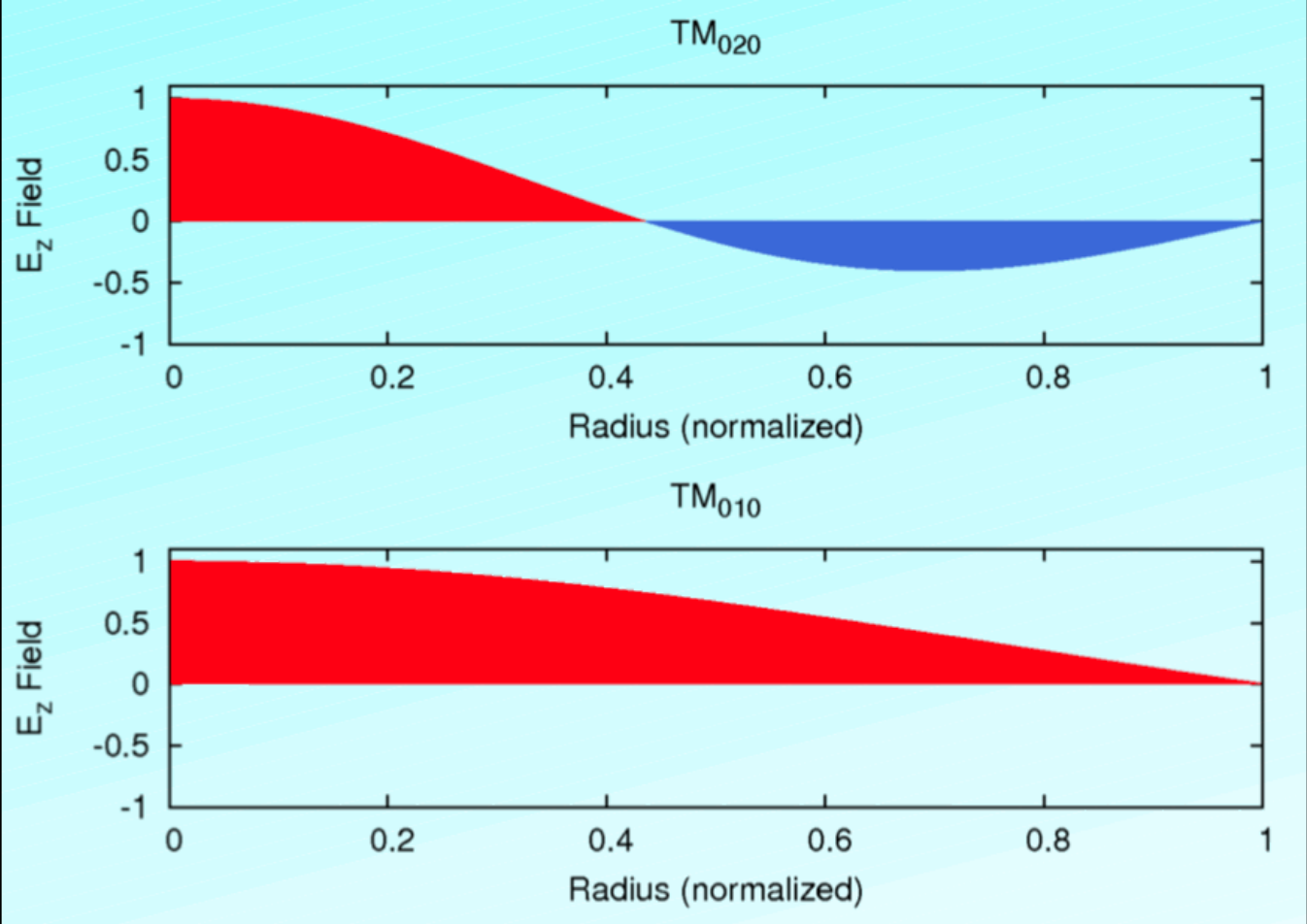
- Injected Power
x100-1000 axion power

- Noise floor

- Bandpass filter



Possible to instrument higher order modes



TM₀₂₀ Mode
Relative Frequency
2.3

Tuning Range
920-2,100 MHz

Relative Power
0.41

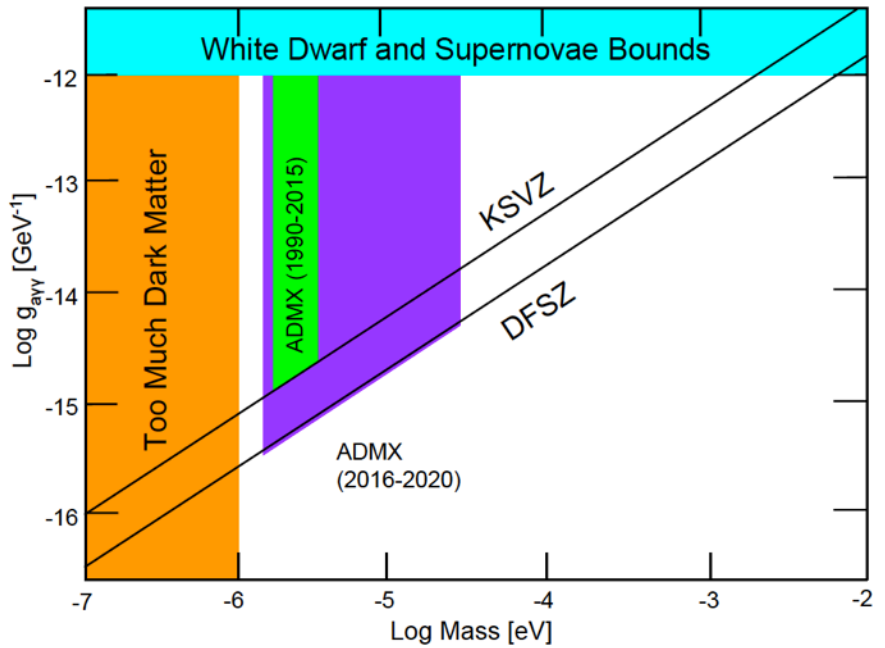
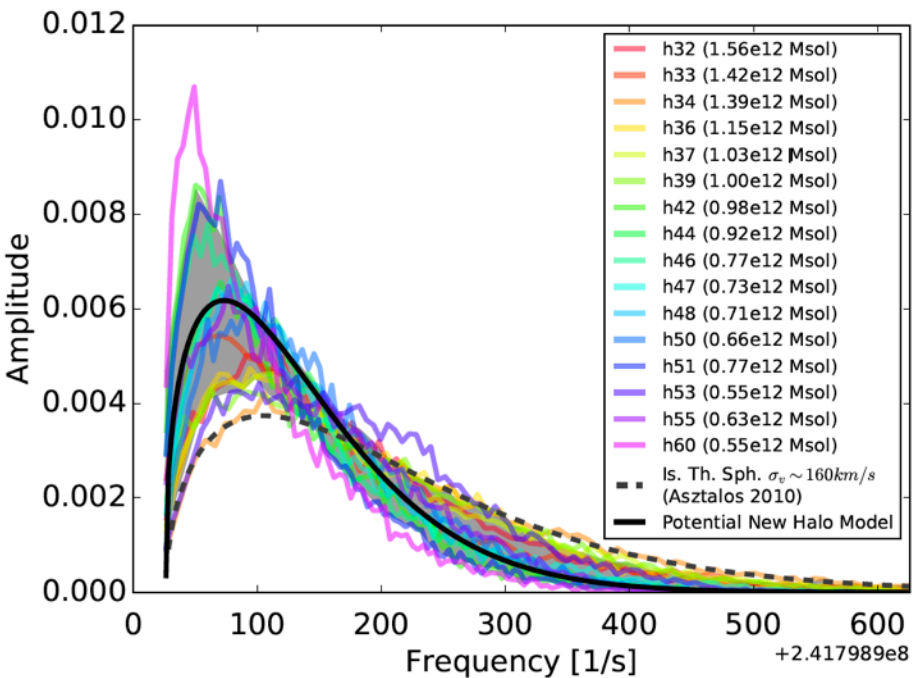
TM₀₁₀ Mode
Relative Frequency
1.0

Tuning Range
400-900 MHz

N-body simulations point to sharper line-shape

If true than haloscope experiments would naturally be even more sensitive

Currently we continue to operate with conservative isothermal halo

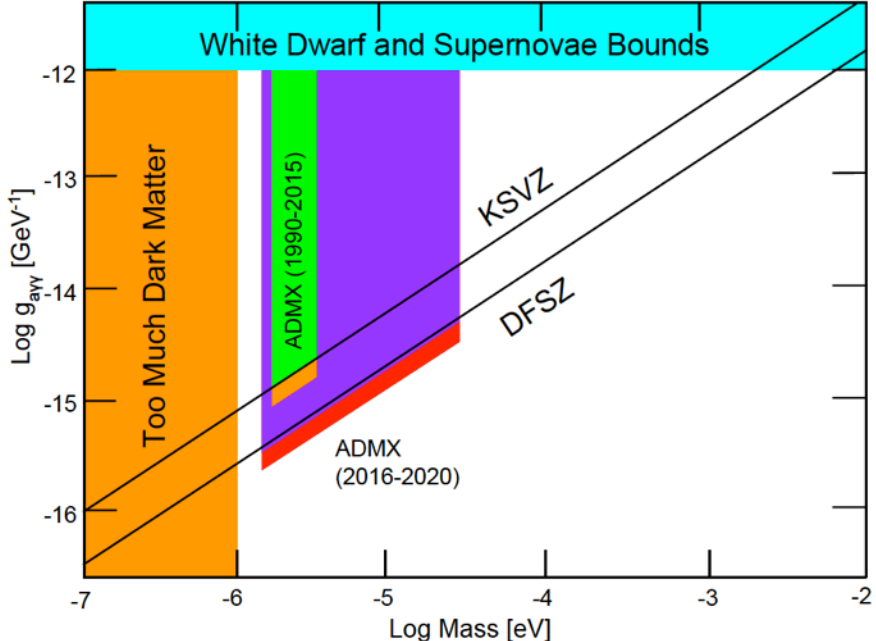
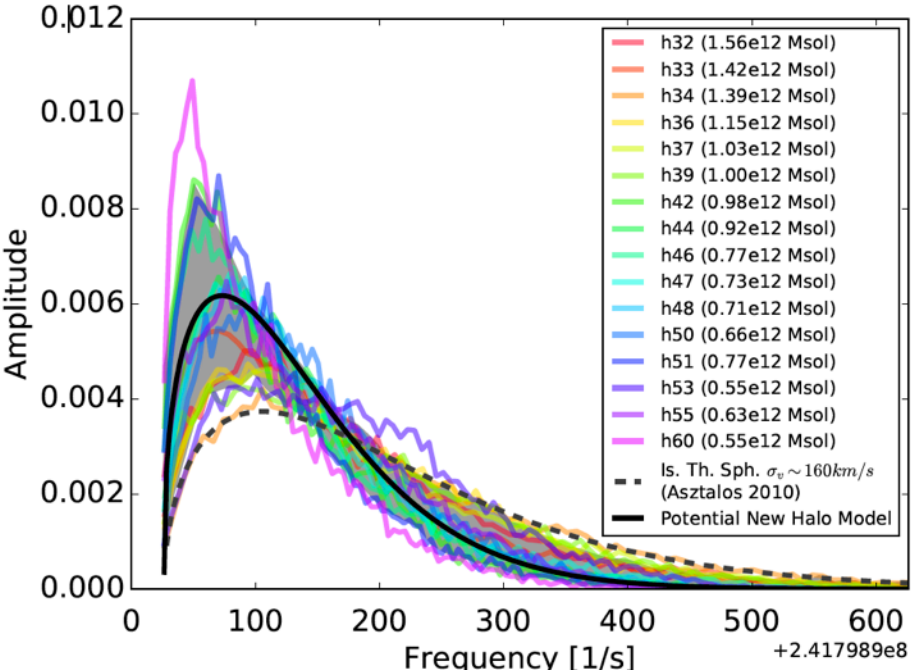


N-body simulations point to sharper line-shape

If true than haloscope experiments would naturally be even more sensitive

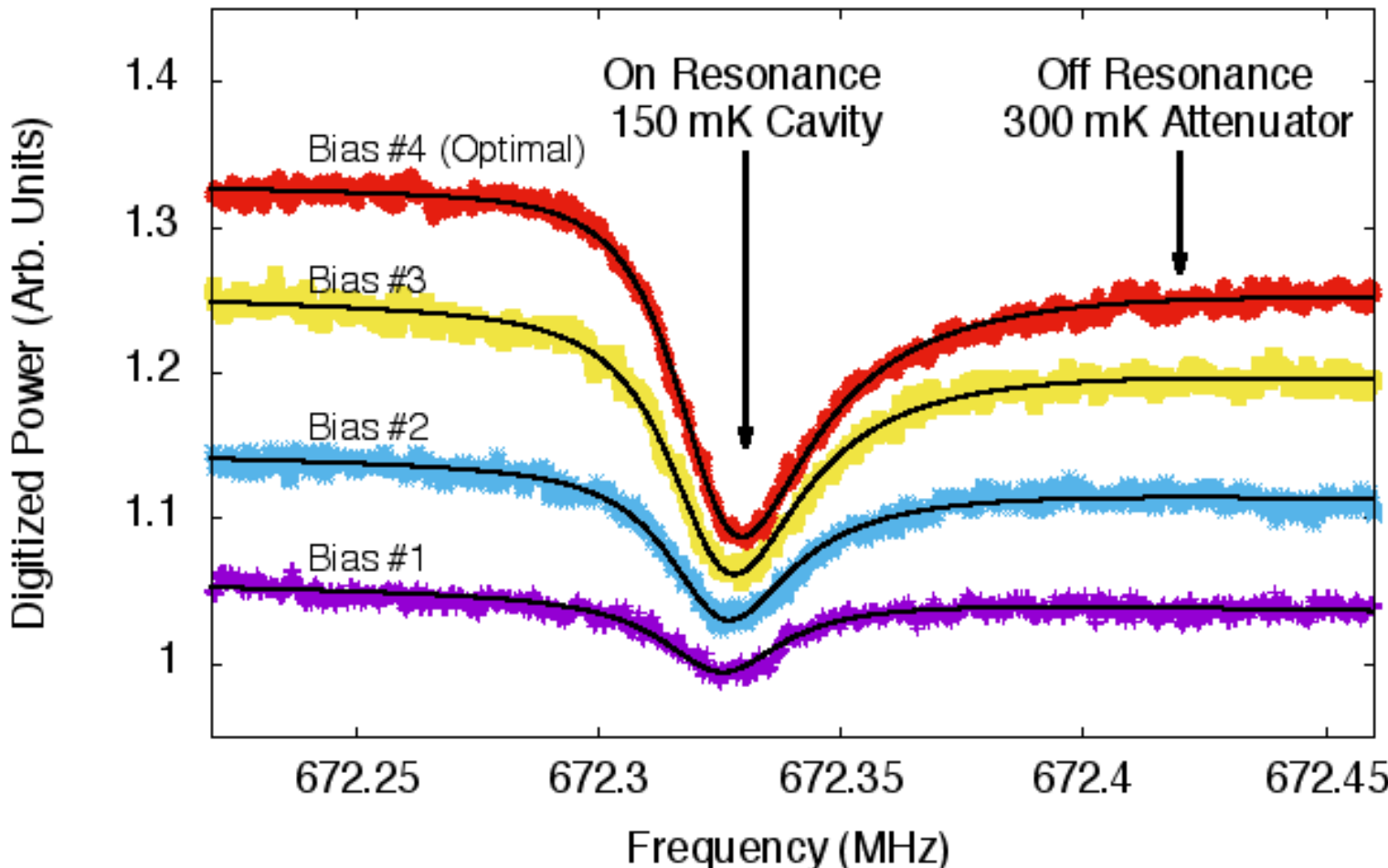
Currently we continue to operate with conservative isothermal halo

If the simulations from the UW “N-body shop” are correct limits can be extended.



Fit receiver chain to determine system noise

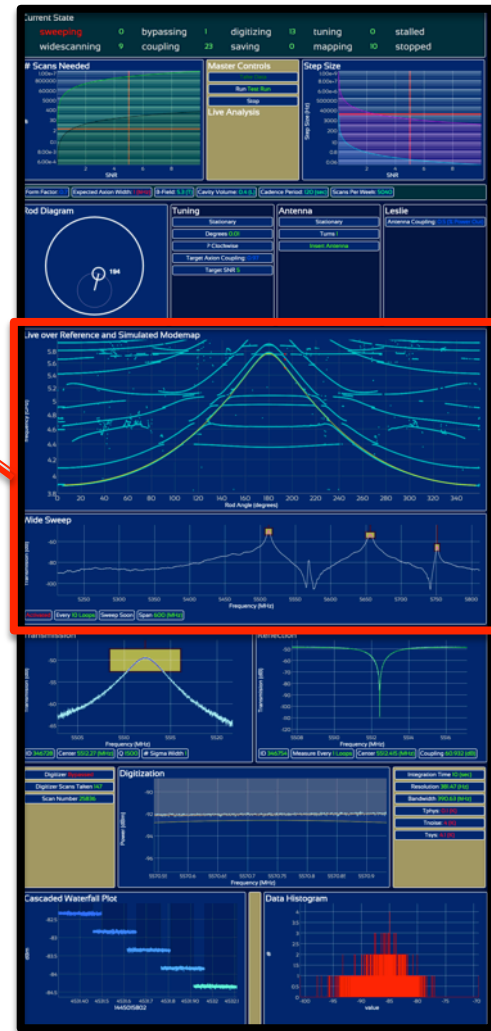
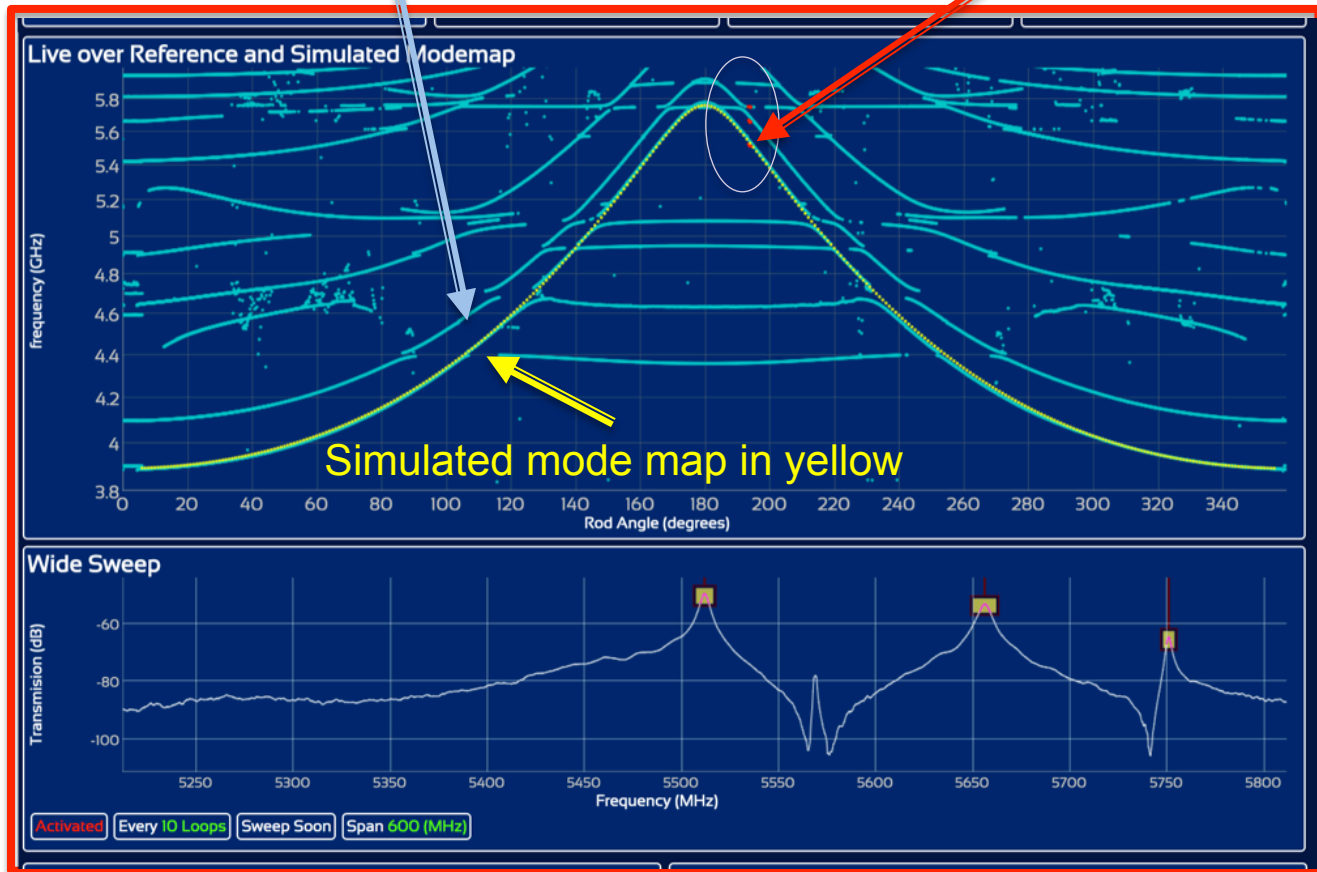
Example Cavity Noise Measurement
Multiple MSA Biases



Sidecar DAQ: Mode Map

Reference mode map in blue

Red Dots are live data taken from wide scan



*Recent Ph.D thesis of Christian Boutan (now at PNNL)

ADMX: Collaboration



Lawrence Livermore
National Laboratory



The
University
Of
Sheffield.



Fermilab



Pacific
Northwest
NATIONAL
LABORATORY



Sponsors

ADMX now DOE Gen 2 project



HEISING - SIMONS
FOUNDATION

R&D support



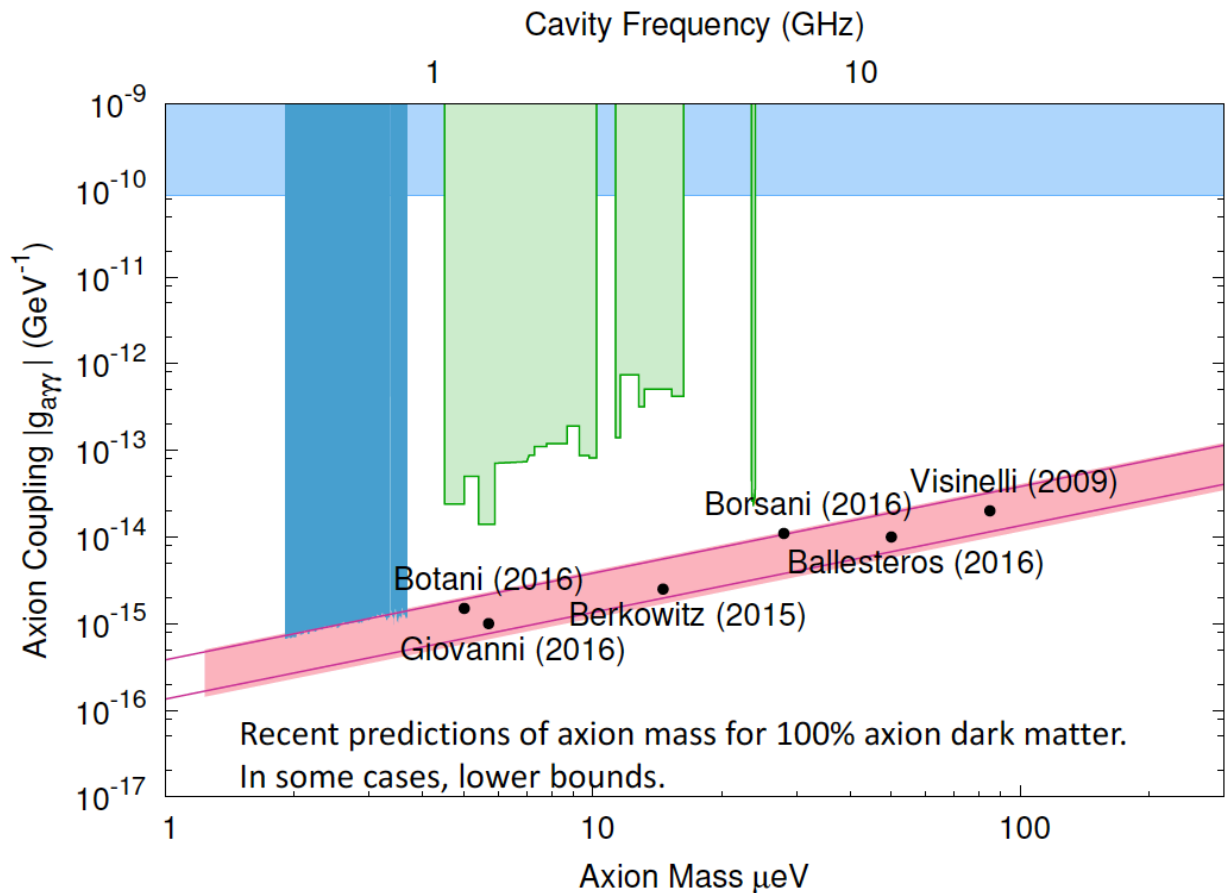
Los Alamos
NATIONAL LABORATORY



Where to look in mass?

For PQ symmetry breaking **after** inflation analytical and lattice results point to 1-100 μeV masses.

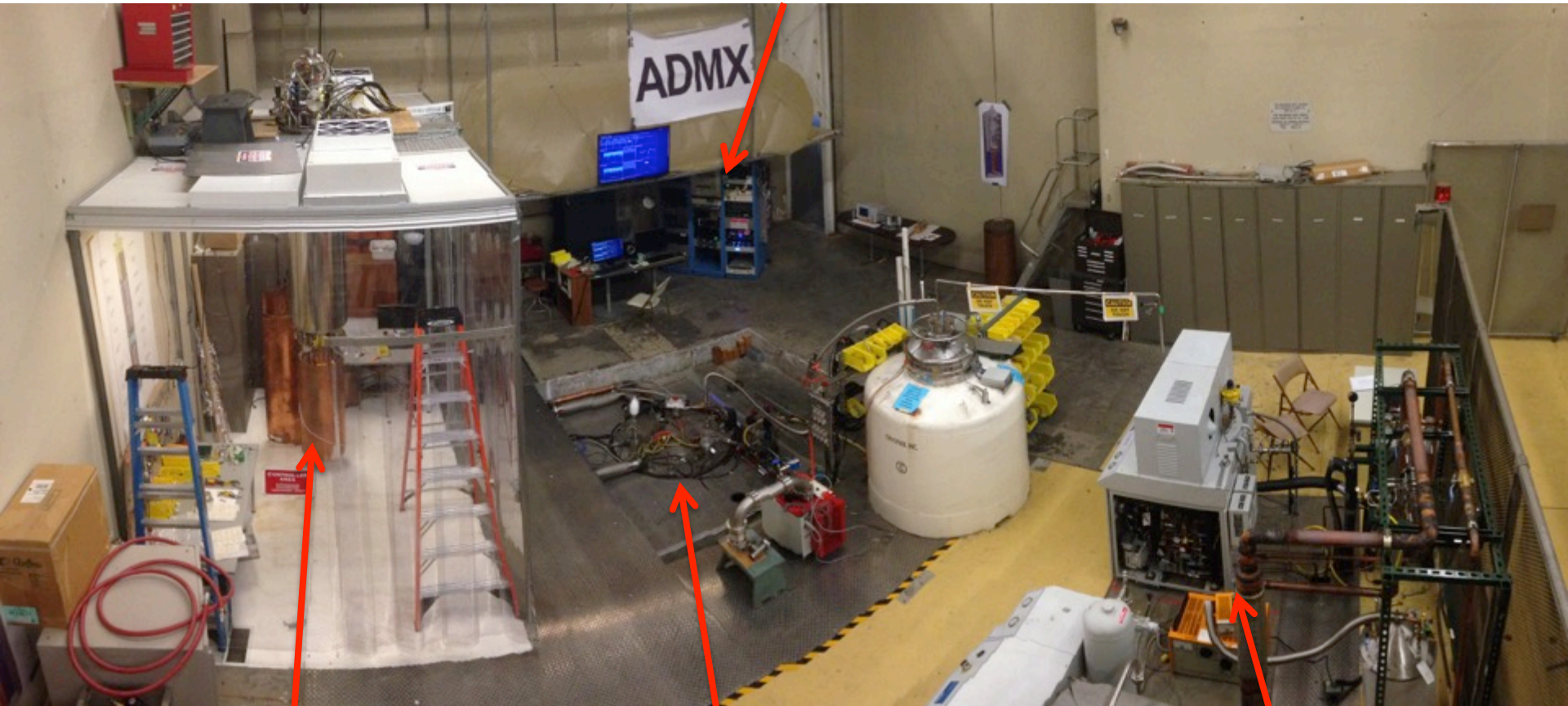
For PQ symmetry breaking **before** inflation less constrained (could be $\ll \mu\text{eV}$ masses)



ADMX site: University of Washington

Center for Experimental Nuclear Physics and Astrophysics (CENPA)

ADMX DAQ & Controls



Cleanroom
(with insert hanging)

ADMX Magnet

Helium liquefier

Scan Rate from Dicke Radiometer equation

- Rate determined from SNR

$$\frac{df}{dt} \approx 750 \text{ MHz/year} \left(\frac{g_\gamma}{0.36} \right)^4 \left(\frac{5}{SNR} \right)^2 \left(\frac{f}{1 \text{ GHz}} \right)^2 \cdot \left(\frac{B_0}{8 T} \right)^4 \left(\frac{V}{100l} \right)^2 \left(\frac{Q_L}{10^5} \right) \left(\frac{C_{010}}{0.5} \right)^2 \left(\frac{0.2 \text{ K}}{T_{sys}} \right)^2$$

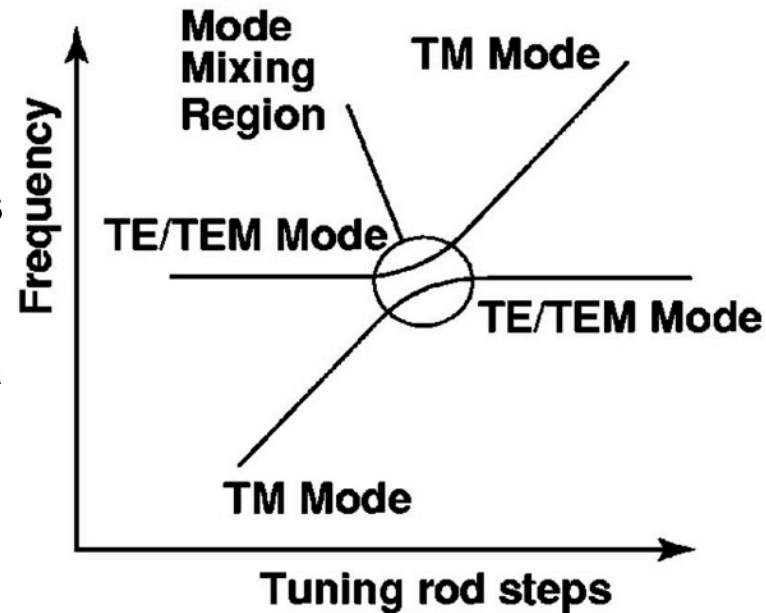
- SNR is the Signal-to-Noise for detection (usually set to 5),
- f is the frequency being searched (where 1 GHz $\sim m_a = 4.1$ ueV)
- T_{sys} is the total system temperature ($T_{sys} = T_{cavity} + T_{amps}$)
- To scan at this sensitivity would take > 100 years with original ADMX

Managing mode-crossings

- Transverse Magnetic (TM) modes move up in frequency as tuning rods are rotated
- Transverse Electric (TE) modes remain relatively static in frequency.
- When both mode's frequencies are degenerate there's a "mode crossing" in which the two modes "mix" and the resonant peak can disappear.
- The longer the cavity, the more TE modes there are in the tuning range.

Keep aspect ratios: *Length / radius* ~ 5 .

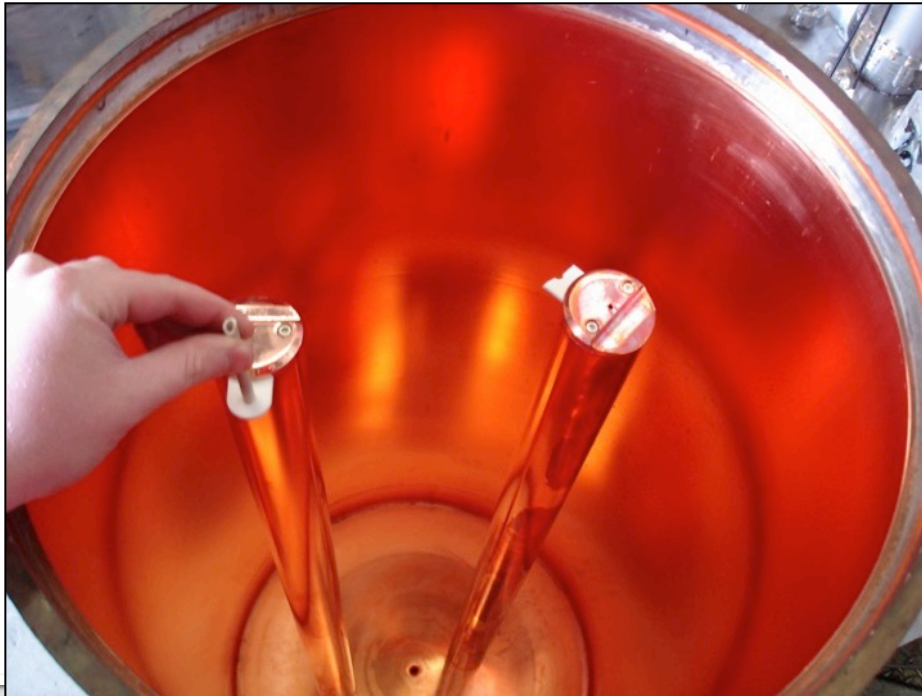
- We step around mode crossings by using multiple tuning rods (metal and dielectric).



Challenge of moving to higher axion mass (frequency)

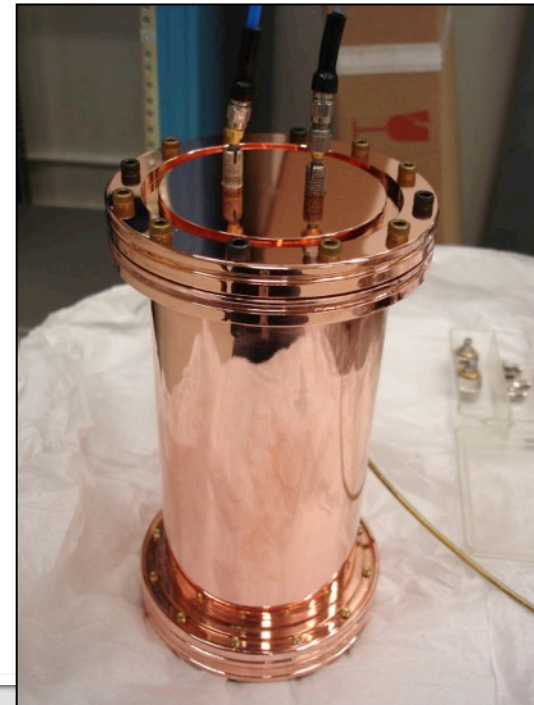
- Scaling single cavity to higher frequencies (f) – Volume $\sim (f)^{-3}$!
- Quality factor also goes down as frequency increases ($Q_L \sim 10^5 \cdot (f)^{-2/3}$)
- Need to move to multi-cavity array's.

Frequency ~ 540 MHz
 $Q_L - 100,000$
Axion Mass $\sim 2 \mu\text{eV}$
Volume – 135 liters



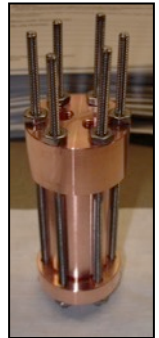
Lawrence Livermore 16" diameter laboratory

Frequency ~ 2.4 GHz
Axion Mass $\sim 9 \mu\text{eV}$
 $Q_L - 60,000$
Volume ~ 2.6 liters



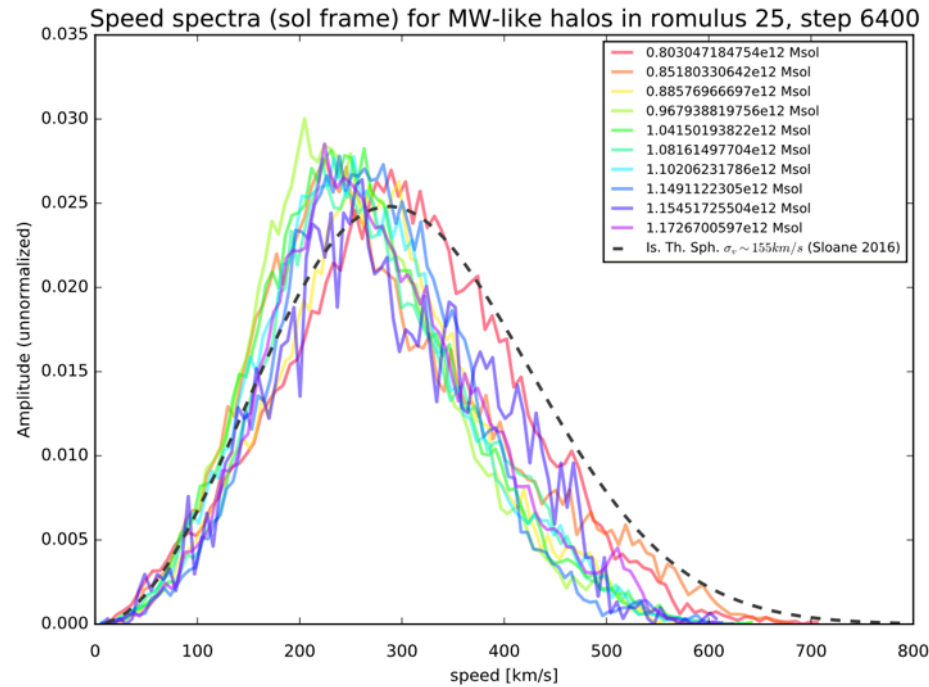
5" diameter Physical Sciences

Frequency ~ 10 GHz
Axion Mass $\sim 36 \mu\text{eV}$
 $Q_L - 25,000$
Volume – 0.025 liters



1" diameter

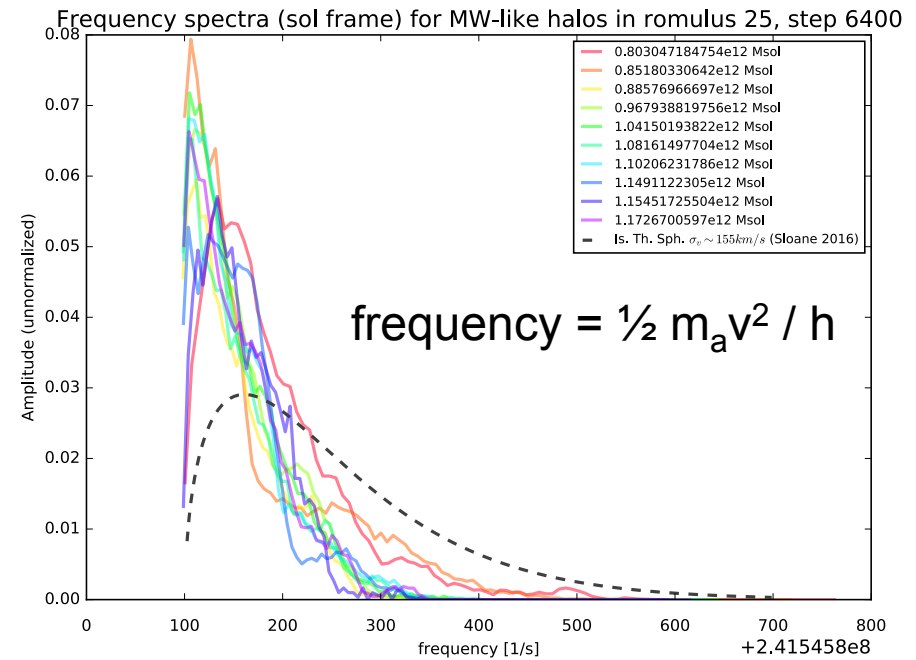
Can obtain axion phase space distribution immediately after discovery + reconfirmation



For example:

Test N-body simulations from T. Quinn predicting long-suspected co-rotating galactic DM \rightarrow Softer spectrum?

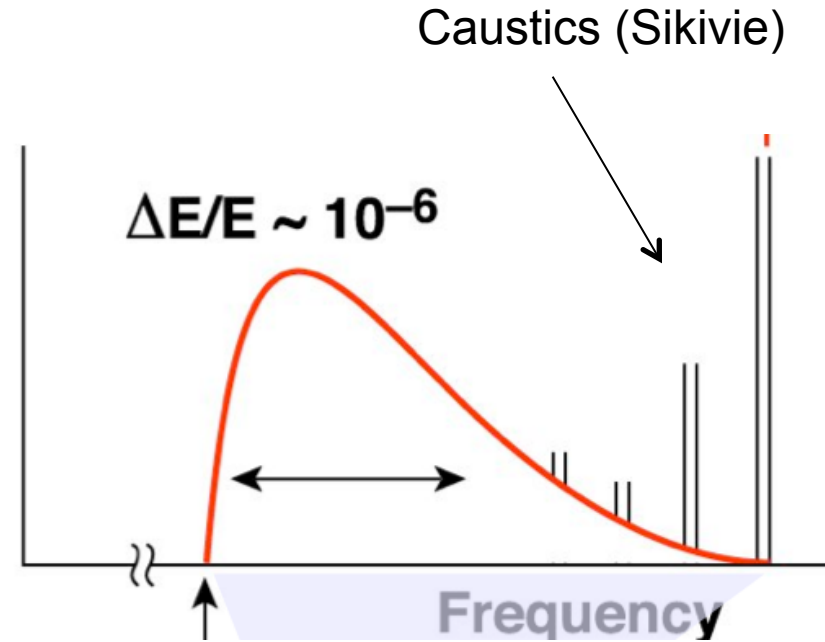
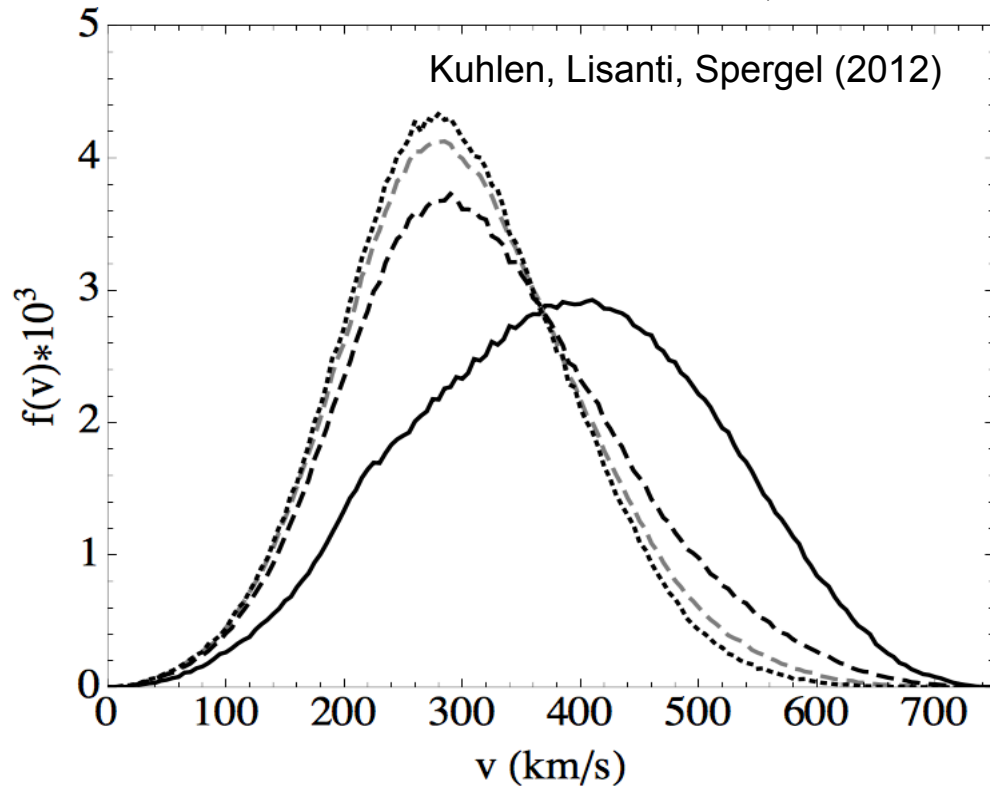
Narrower line = better Signal/Noise?



For 2x frequency resolution, need 2x sample time. Half-power in each bin requires 4x samples. Total 8x integration time is easy!

Test models of debris flow, narrowband caustic structure from recent infalls, etc.

Via Lactea 2 simulation,

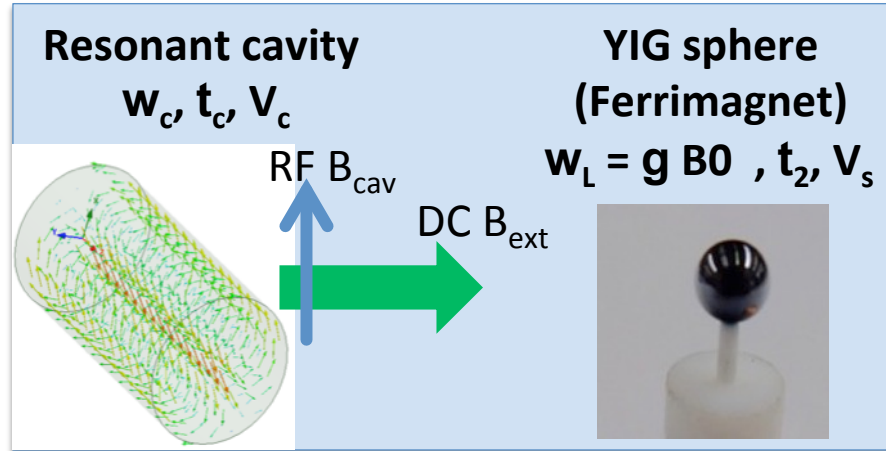


If dark matter is multicomponent, these measurements could also provide valuable input to WIMP detection experiments.

QUAX: Axion induced RF emission

A volume V_s of magnetized material, strong coupled in a microwave resonant cavity, will absorb energy from the axion wind and re-emit as rf power

With magnetizing field
 $B_0 = 1.7 \text{ T} \Rightarrow 48 \text{ GHz}$



$$P_{out} = \frac{P_{in}}{2} = 3.8 \times 10^{-26} \left(\frac{m_a}{200 \mu\text{eV}} \right)^3 \left(\frac{V_s}{100 \text{ cm}^3} \right) \left(\frac{n_S}{2 \cdot 10^{28} / \text{m}^3} \right) \left(\frac{\tau_{min}}{2 \mu\text{s}} \right) W$$

R & D in progress

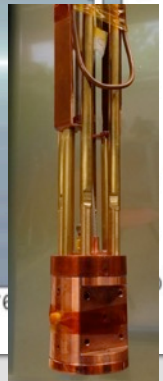
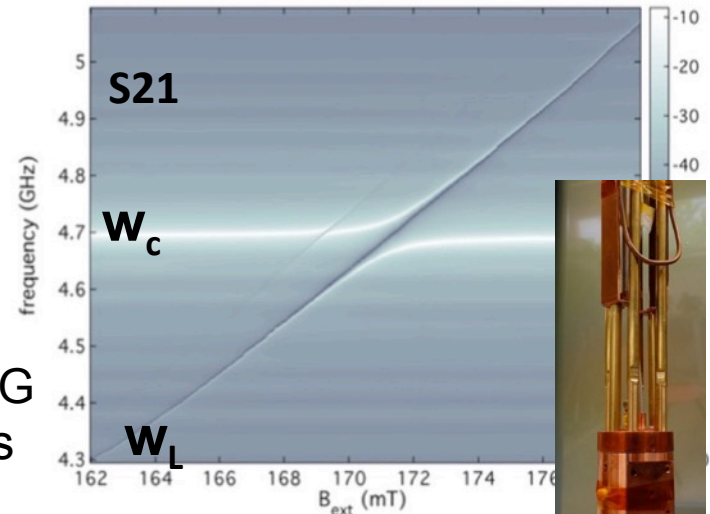


Niobium Cavity

$T = 300\text{K}$
 $f_c = 13.964 \text{ GHz}$
 $Q_0 = 5.0 \cdot 10^3$

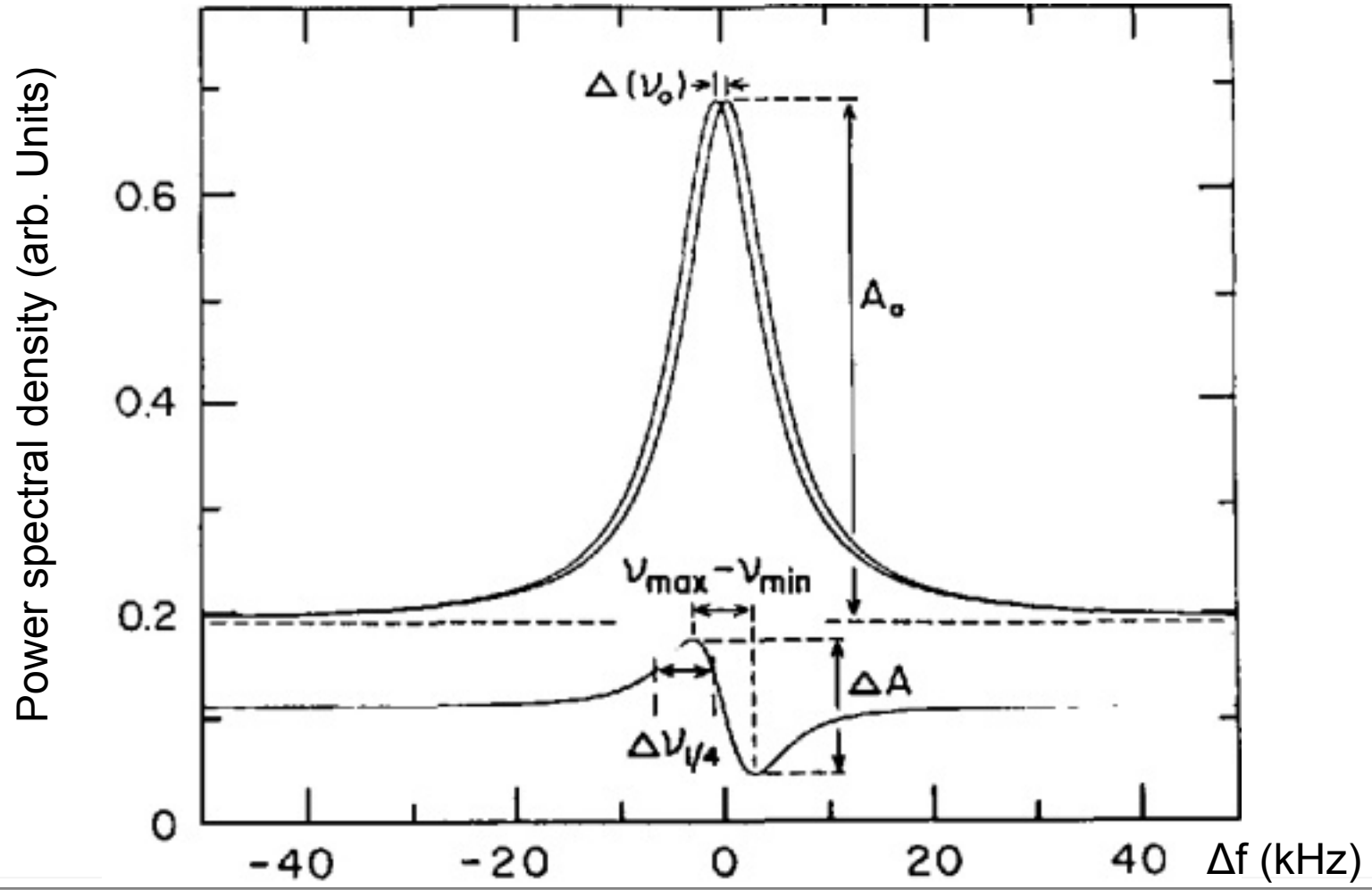
$T = 4.2\text{K}$
 $f_c = 13.960 \text{ GHz}$
 $Q_0 = 5.0 \cdot 10^5$

Cavity – YIG resonances



Annual modulation is also easy

Annual modulation $\delta v/\Delta v = (60 \text{ km/s}) / (300 \text{ km/s}) = 20\%$ of linewidth.
Resolve frequency shift with 5x integration time = few 10's minutes.



CAST experiment @ CERN

