



# Solar axion search by annual modulation with XMASS-I detector



K. Ichimura

Kamioka observatory, ICRR, the University of Tokyo

Kavli IPMU

for the XMASS collaboration



# Contents



- Introduction of solar Kaluza-Klein (KK) axion
- XMASS-I detector
- Data Analysis
- Results



# Introduction of KK axion



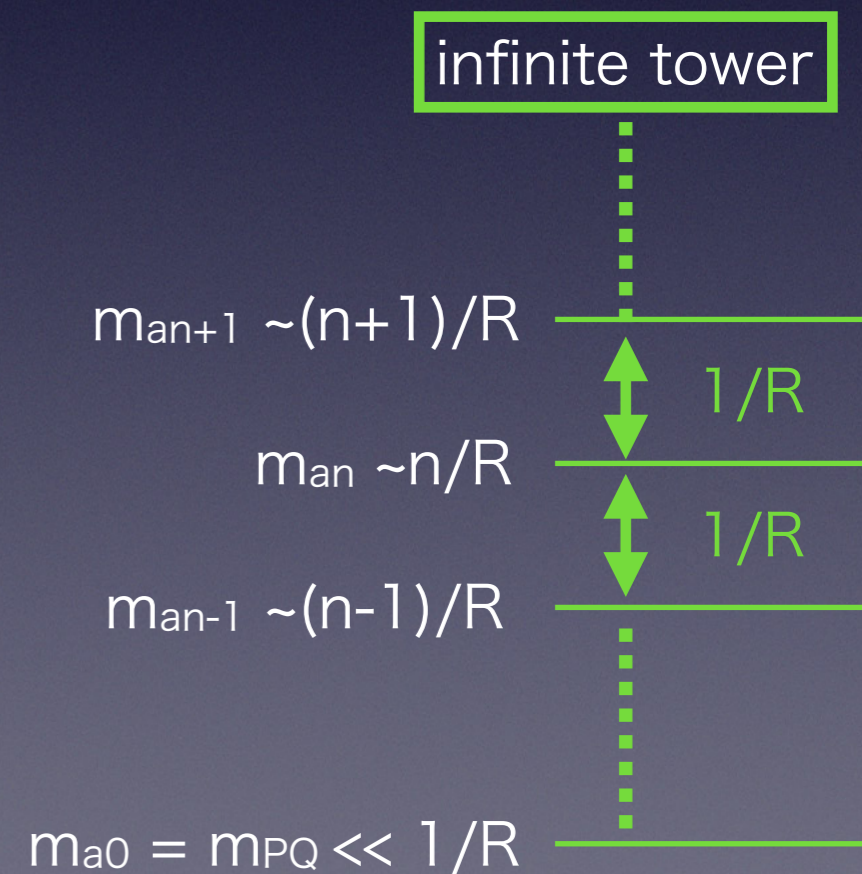
- Axions arise from Peccei-Quinn solution to strong CP problem in QCD.
- On the other hand, large extra dimensions are also proposed to solve the gauge hierarchy problem.
  - Extra dimensions are thought to be “compactified ” in a certain radius R.
  - Motions of a particle in the extra dimensions can be seen as mass state of Kaluza-Klein (KK) excitations, which are separated in  $1/R$

• In theories with  $n$  extra dimensions, axions may be able to propagate and acquire KK excitations.

- axions acquires an infinite tower of KK modes where the lowest KK state is the normal PQ axions.

• Assumed model

- $n= 2$  extra dimensions
- $\delta =2$  extra dimensions that axions can propagate in
- $\rightarrow$  mass spacing of the KK axion  $1/R \sim 1$  eV

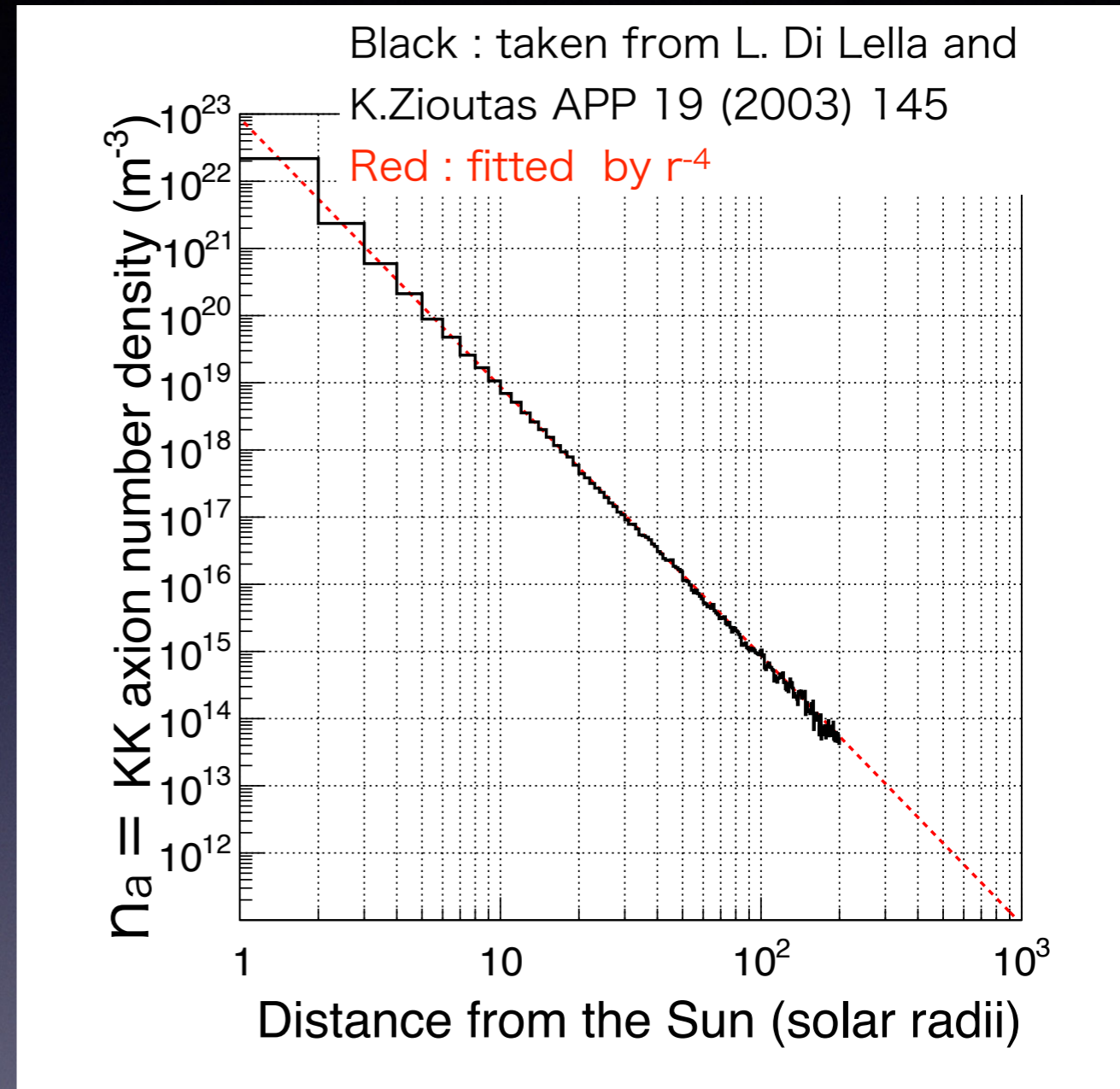




# Introduction of solar KK axion



- KK axion would be produced in the Sun via the Primakoff effect ( $\gamma Z \rightarrow aZ$ ) and a photon coalescence mechanism ( $\gamma \gamma \rightarrow a$ )
- A small fraction is trapped by the gravity of the Sun.
- Such solar/stellar KK axion can explain the solar corona problem and so on by massive axion decay
- In APP 19 (2003) 145, they assume KK axion photon coupling  $\underline{g_{a\gamma\gamma} = 9.2 \times 10^{-14} \text{ GeV}^{-1}}$  by requiring that axion decay is responsible for the X-ray surface brightness of the quiet Sun



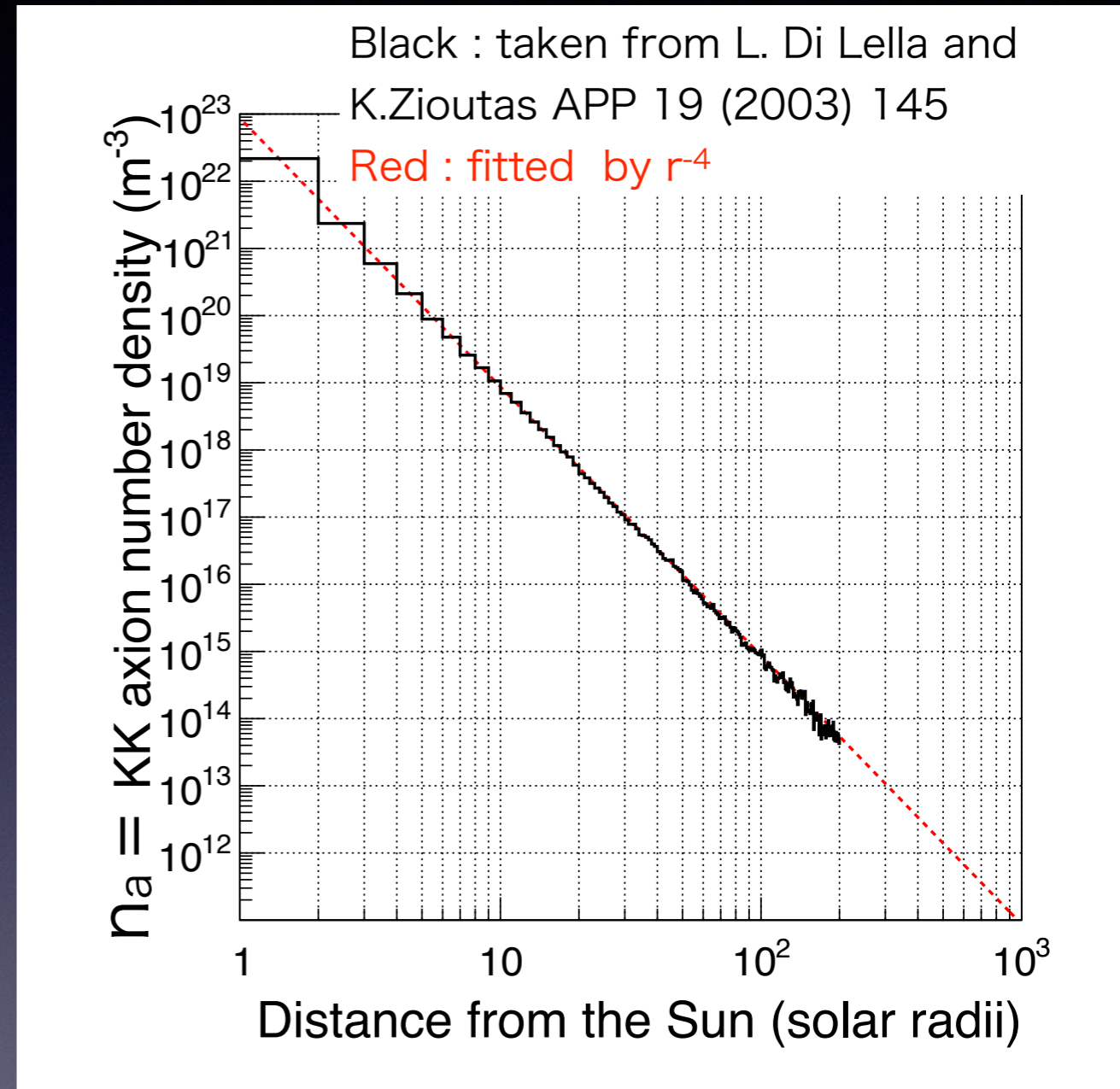
Solar KK axion reference : APP 19 (2003) 145, ApJ 607 (2004) 575 and APP 23 (2005) 287 etc.



# Introduction of solar KK axion



- we can predict the present density of trapped solar axion with the following assumptions
  - KK axion photon coupling  $g_{ar\gamma} = 9.2 \times 10^{-14} \text{ GeV}^{-1}$
  - number of extra dimension  $\delta = 2$
  - Compactification radius  $R = 10^3 \text{ keV}^{-1}$
  - Mass splitting of KK axion : 1 eV



$$n_a = 4.36 \times 10^{13} \text{ m}^{-3} \text{ at perihelion } 211.4 R_{\odot}$$
$$n_a = 3.81 \times 10^{13} \text{ m}^{-3} \text{ at aphelion } 218.6 R_{\odot}$$



# Expected energy spectra



- The expected KK axion decay rate  $R$  :

$$R = (2.5 \times 10^{11} m^{-3} day^{-1}) \left( \frac{g_{a\gamma\gamma}}{GeV^{-1}} \right)^2 \left( \frac{n_a}{m^{-3}} \right)$$

(B.Morgun et al, APP 23 (2005) 287)

- distance between the Sun and the Earth:

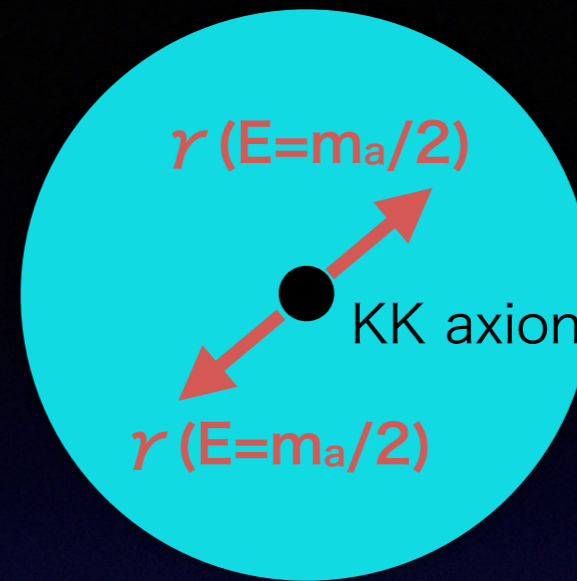
$$r(t) = a \left( 1 - e \cos \frac{2\pi(t - t_0)}{T} \right)$$

- KK axion density  $n_a$  is a function of  $r^{-4}$  :

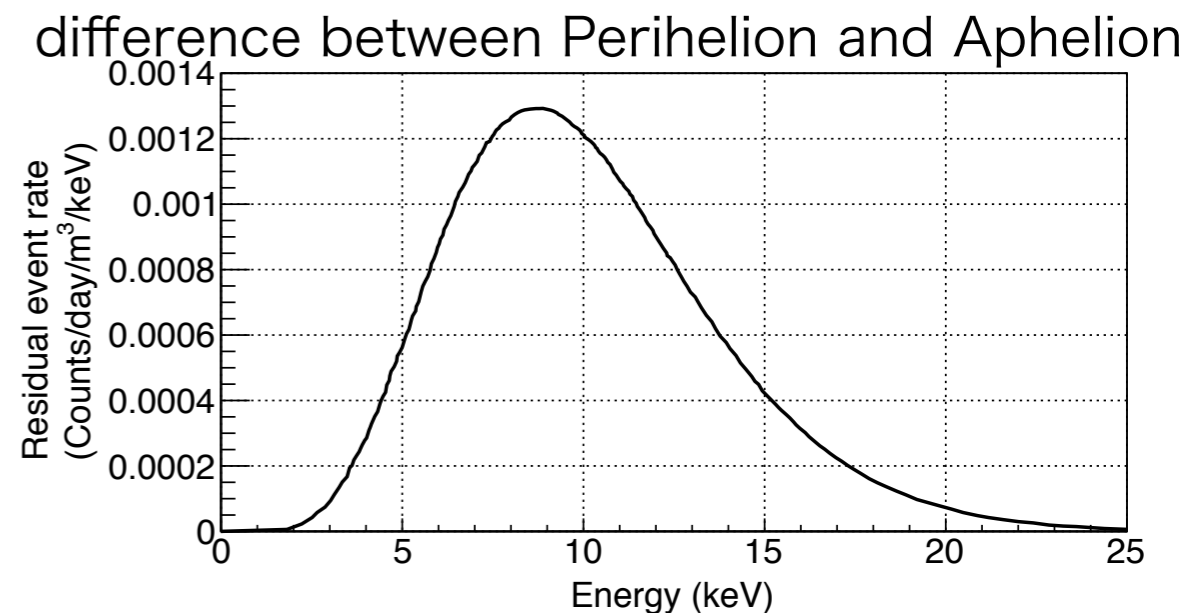
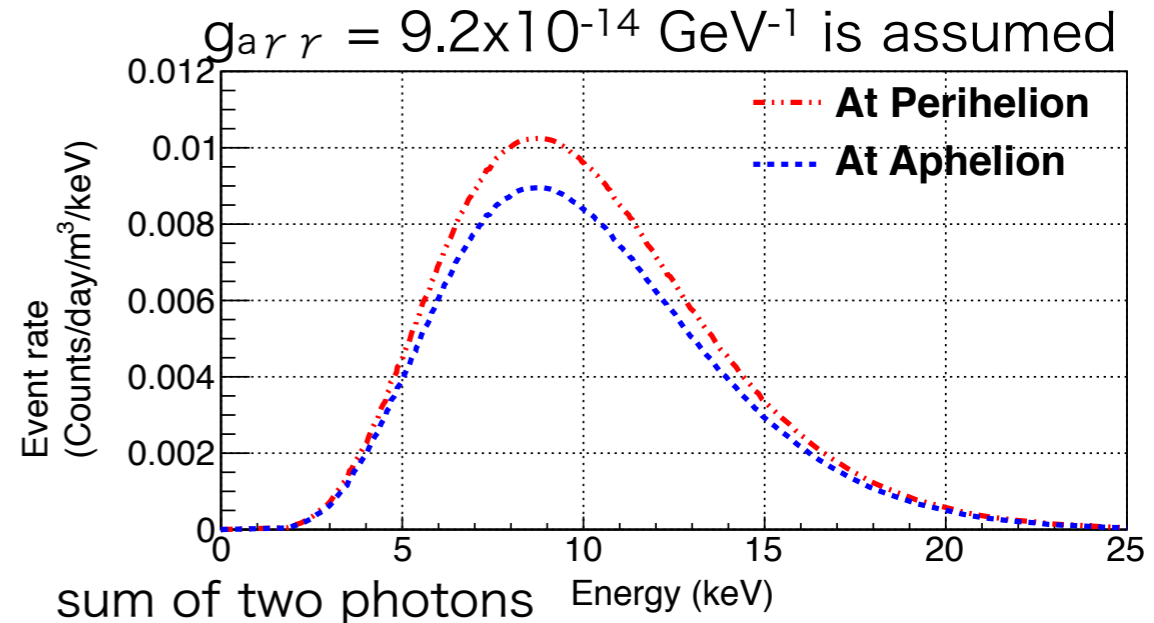
$$n_a(t) = \bar{n}_a \left( 1 - e \cos \frac{2\pi(t - t_0)}{T} \right)^{-4}$$

$$\approx \bar{n}_a \left[ 1 + 4e \left( \cos \frac{2\pi(t - t_0)}{T} \right) + \frac{5}{2} e^2 \cos^2 \frac{2\pi(t - t_0)}{T} \right]$$

- Modulated signal is expected



inside the  
terrestrial det.





# XMASS-I detector

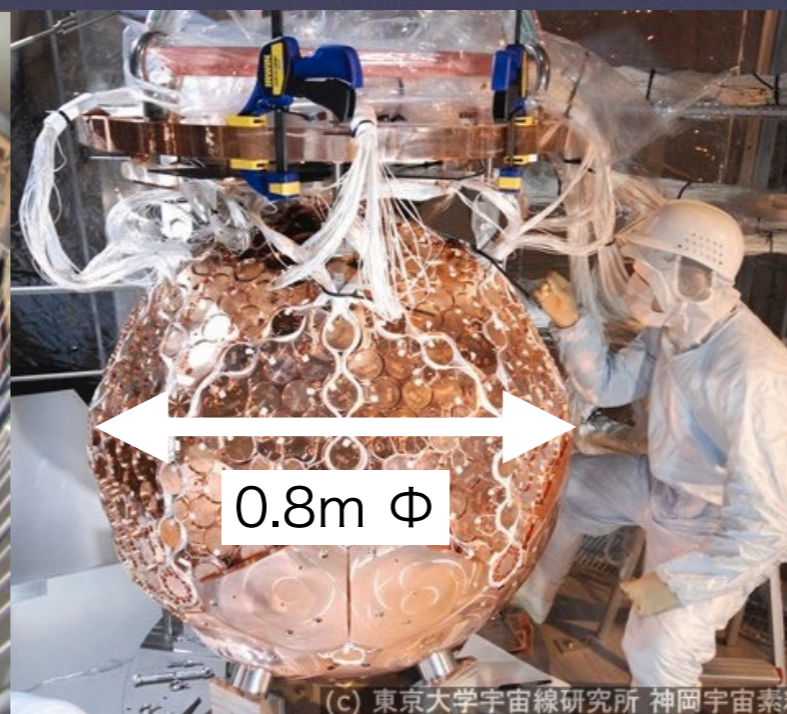


- Located in the Kamioka mine in Japan (~2700 m.w.e.)
- Single phase liquid xenon detector with 832 kg LXe, 0.288m<sup>3</sup> sensitive volume.
- 642 low background 2inch PMTs : 62% photo-cathodes coverage
- High light Yield (~15 p.e. / keV ) and Low energy threshold
  - Energy threshold is enough low to search for solar KK axion.
- High sensitivity for e/ $\gamma$  events as well as nuclear recoil
  - Able to detect Axion Like Particles (ALP), hidden photon, WIMP-Xe inelastic scattering and so on, as well as “Standard” WIMPs
- Stable data taking over years gives the modulation analysis method

$\Phi$ 10m x 10m ultra pure water shield  
70 20-inch PMTs for muon veto



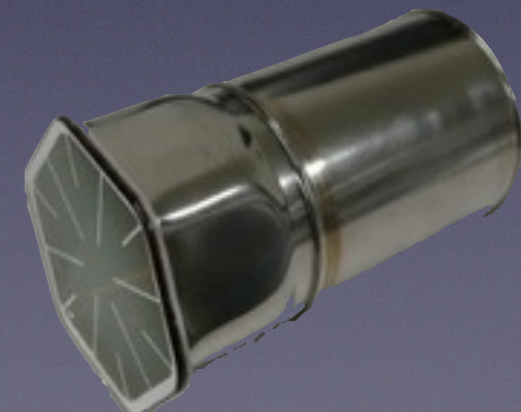
(c) 東京大学宇宙線研究所 神岡宇宙素粒子研究施設



0.8m  $\Phi$

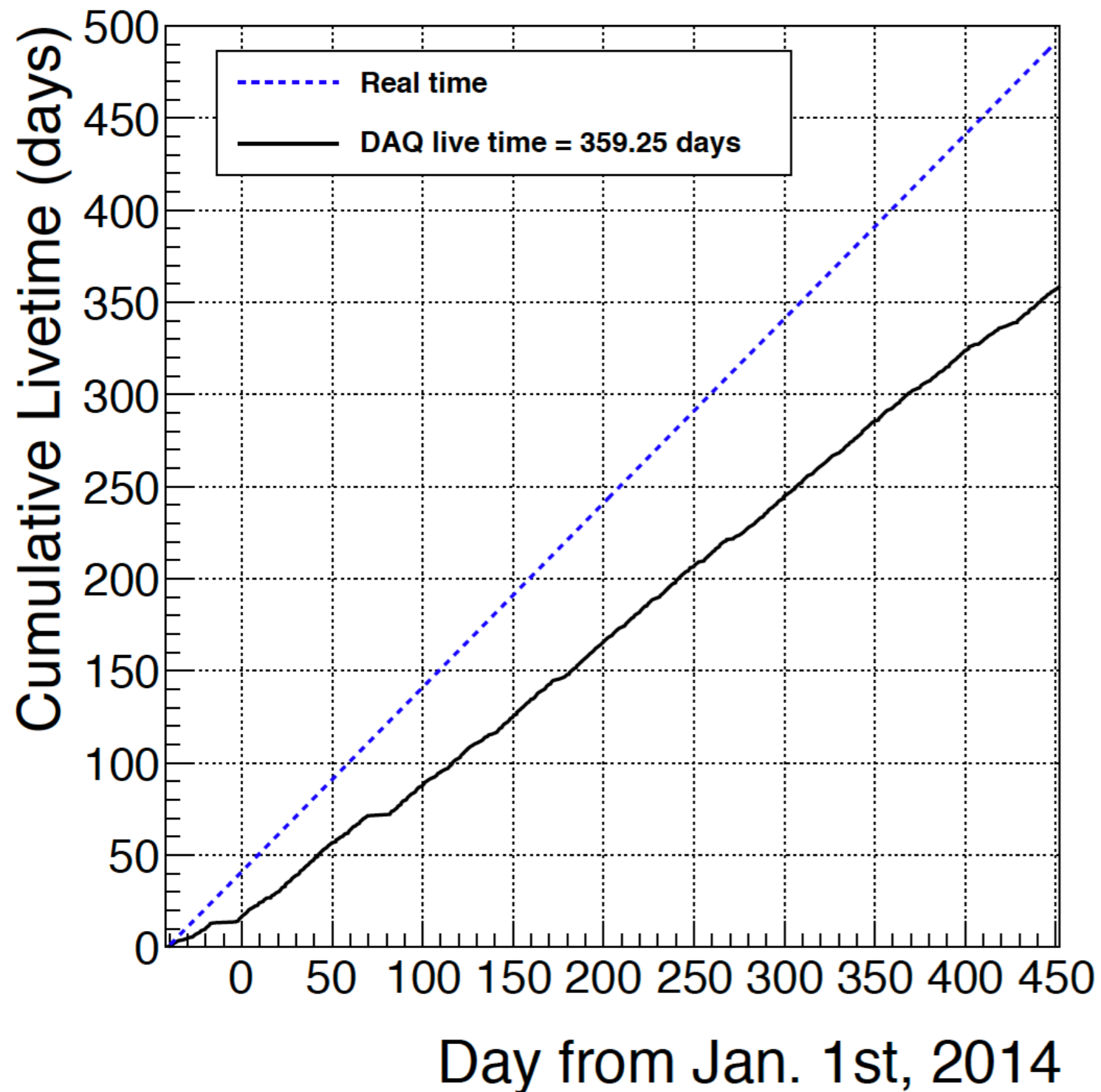
(c) 東京大学宇宙線研究所 神岡宇宙素粒子研究施設

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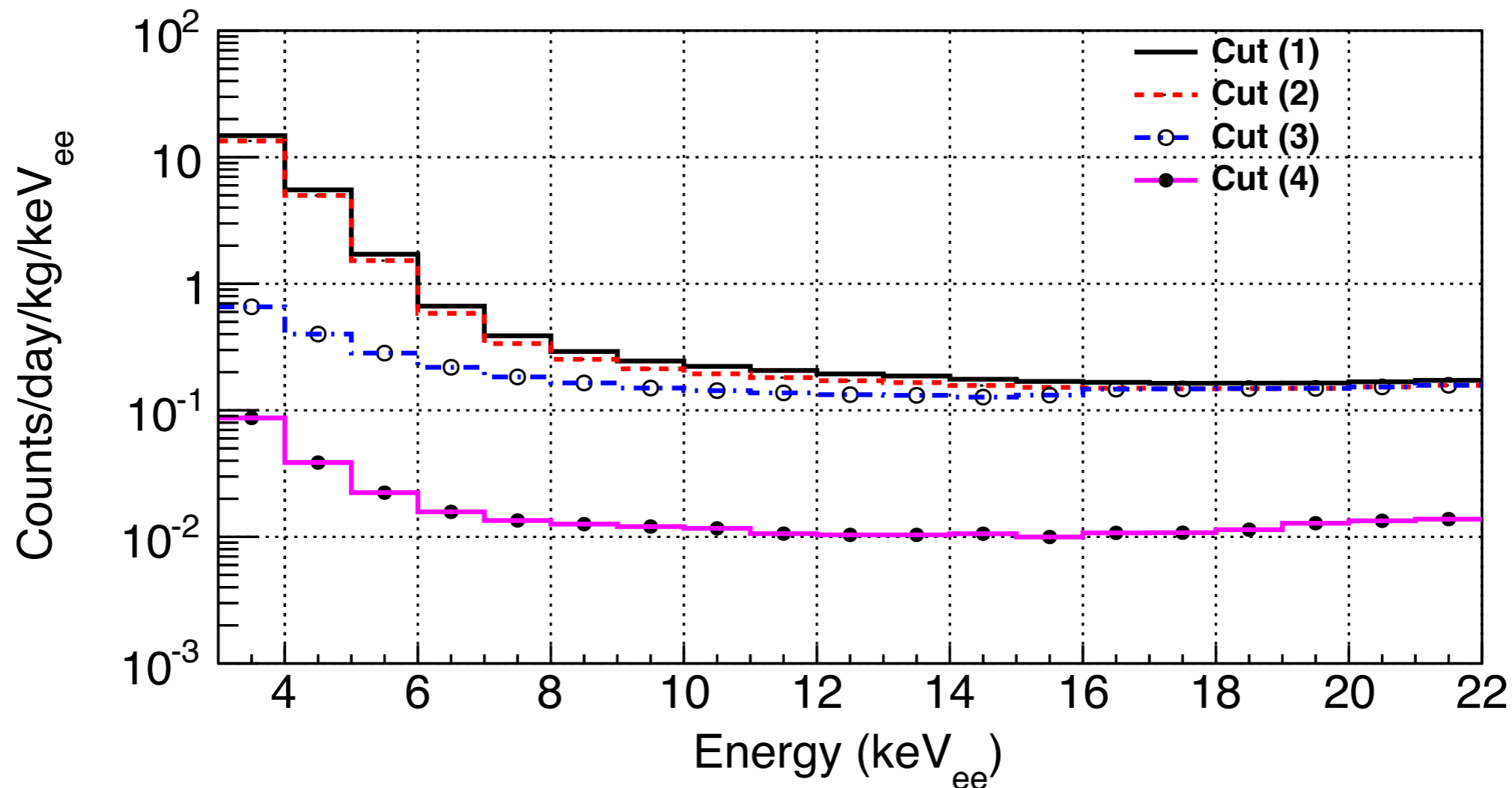
# Data set



- November 2013 - March, 2015
- 504 calendar days with a total live time of 359 days
  - This data set was used for direct DM search by annual modulation. PLB 759 (2016)64
- Data taking is still ongoing.



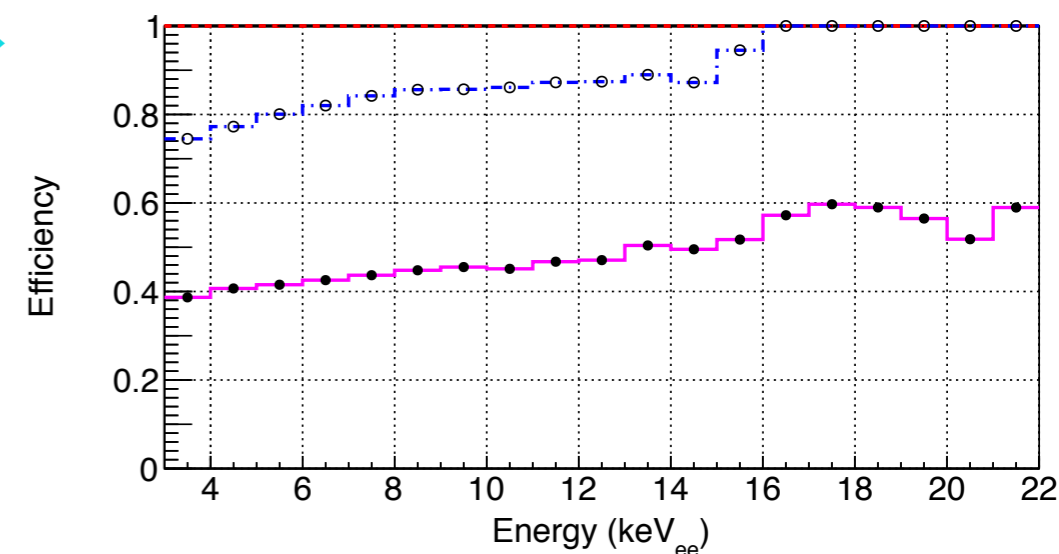
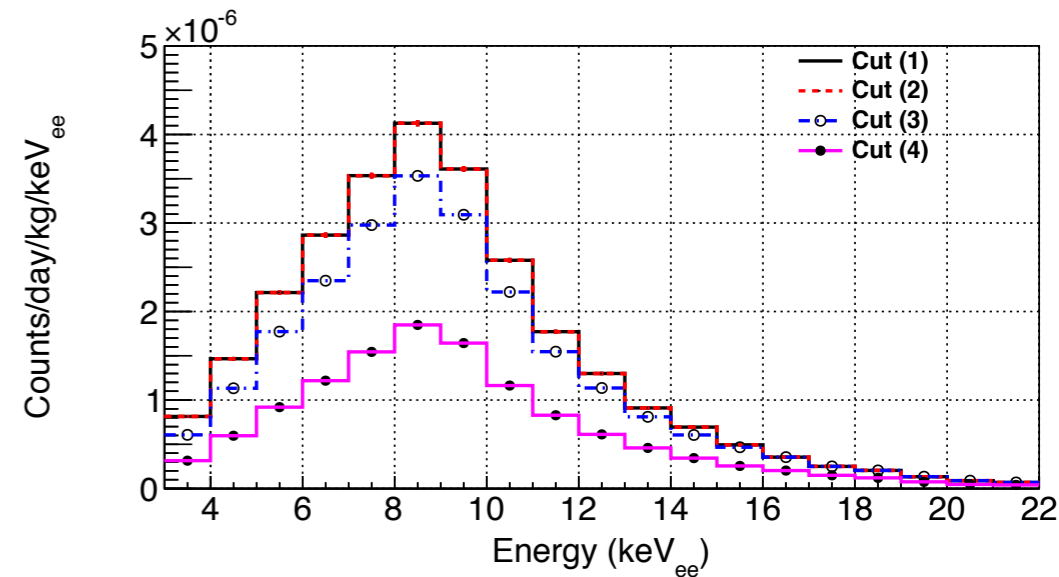
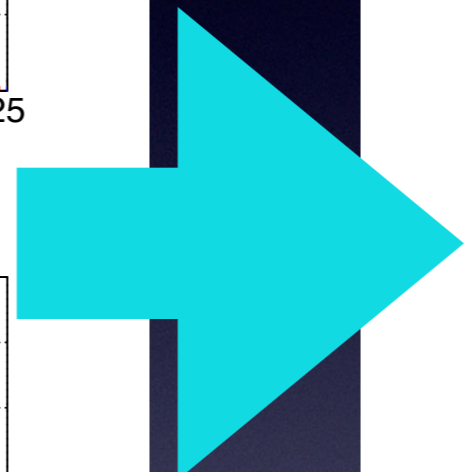
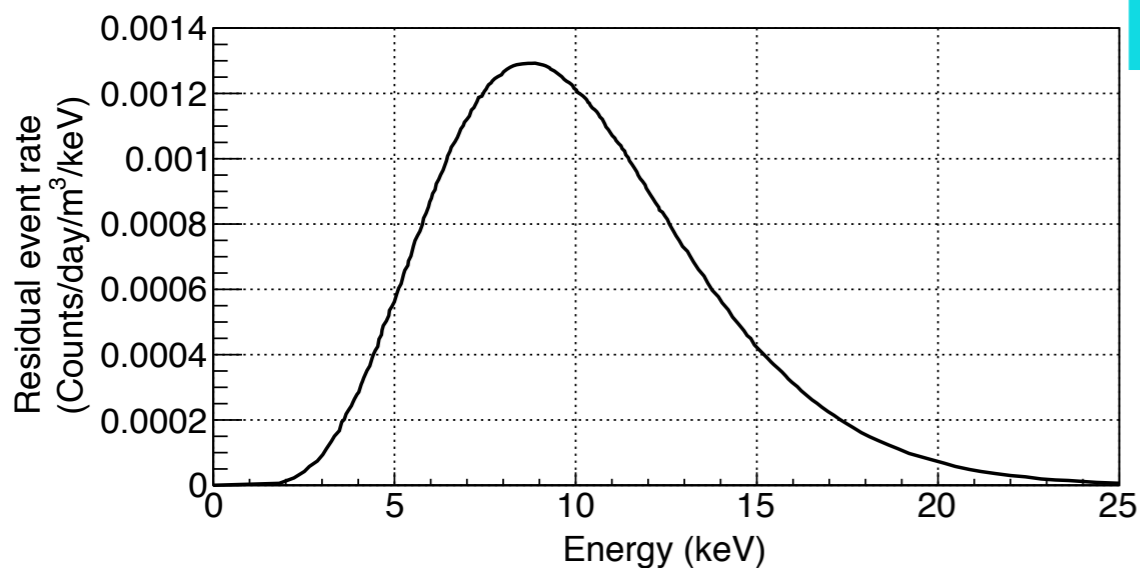
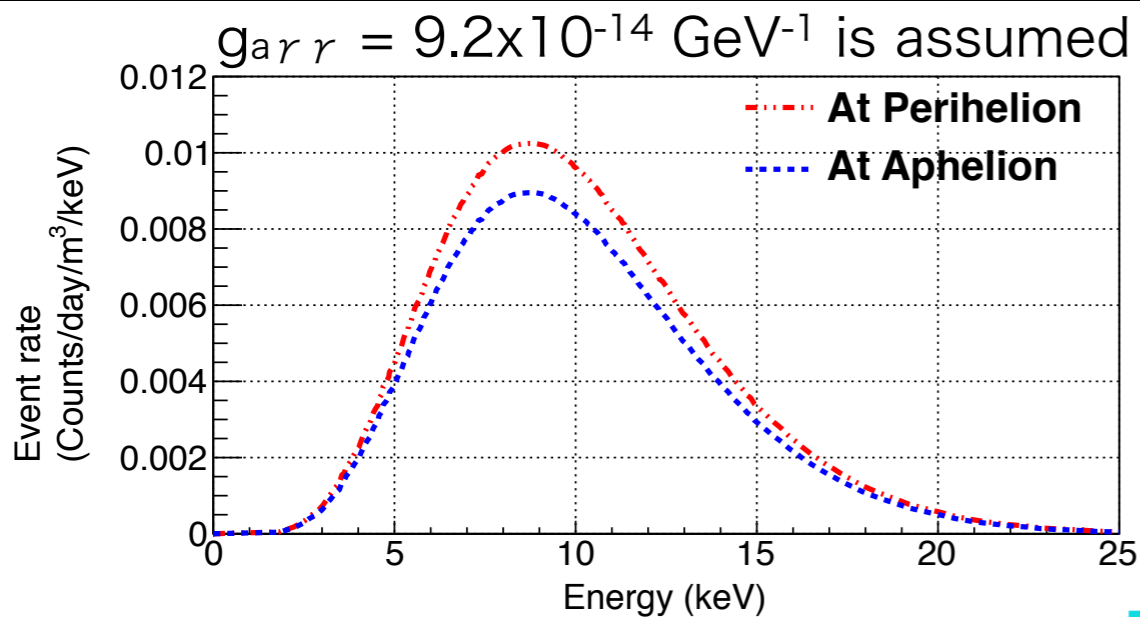
# Event Selection



- Cut1 : triggered only by Inner detector,  $\geq 4$  PMT hits
- Cut2 : Noise events due to afterpulse reduction
  - 10ms veto from previous ID events
  - RMS of hits timing  $< 100$ ns
- Cut3 : Cherenkov events reduction
  - # of hits in the first 20ns is  $< 40\%$  of total hits. which have  $< 200$  observed p.e.
- Cut4 : BG events occurring in front of a PMT window reduction
  - Events which have large MaxPE/TotalPE ratio are removed.



# Signal Simulation



Signal simulation in XMASS-I using Geant4 including detector response after each cut

Assumption :

Distance =  $215.0 R_{\odot}$ , semi major axis

$n_a = 4.07 \times 10^{13} \text{ m}^{-3}$

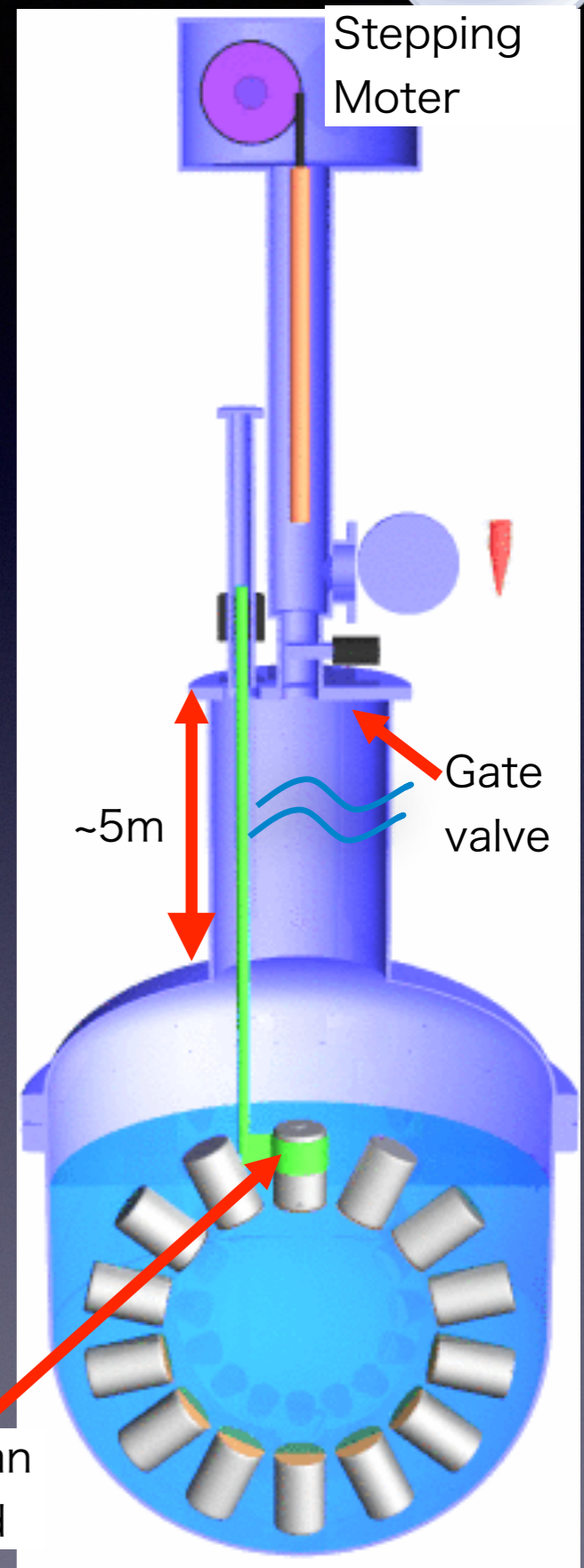
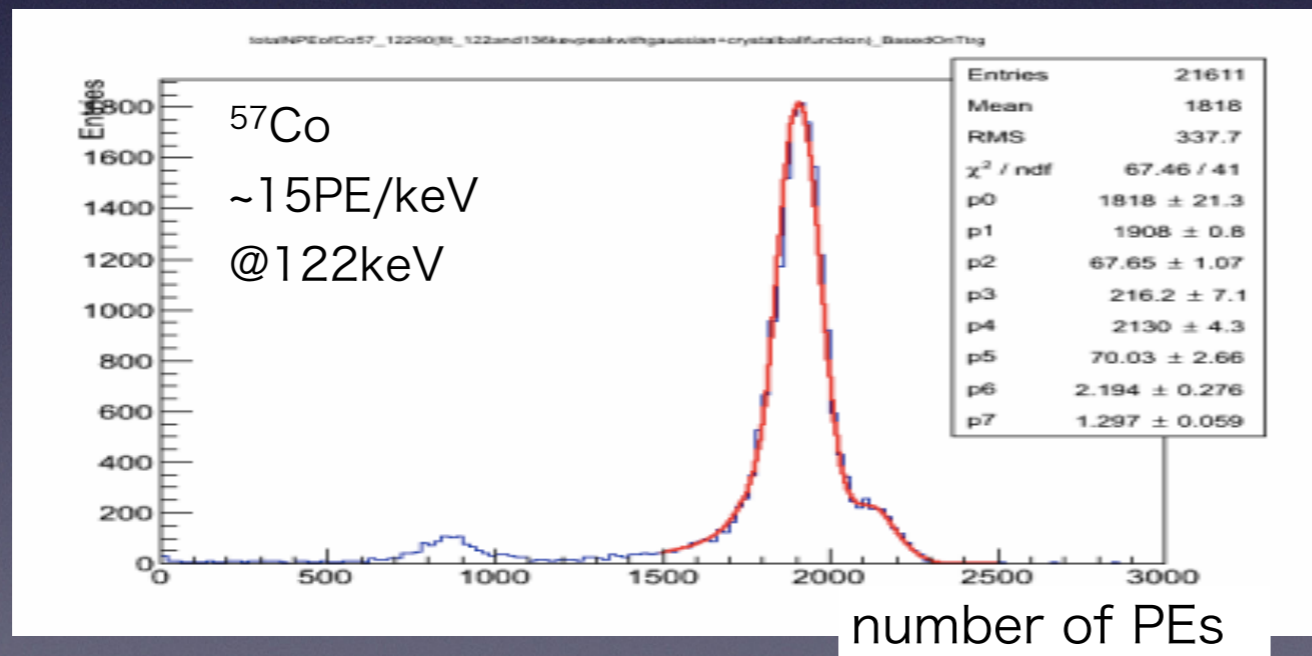
$g_{\gamma\gamma} = 9.2 \times 10^{-14} \text{ GeV}^{-1}$

Input the decay spectra (sum of 2 photons) not including detector response

$n_a = 4.36 \times 10^{13} \text{ m}^{-3}$  at perihelion  $211.4 R_{\odot}$

$n_a = 3.81 \times 10^{13} \text{ m}^{-3}$  at aphelion  $218.6 R_{\odot}$

- Inner Calibration sources :  $^{55}\text{Fe}$ ,  $^{109}\text{Cd}$ ,  $^{241}\text{Am}$ ,  $^{57}\text{Co}$  and  $^{137}\text{Cs}$
- The scintillation light yield response was traced by  $^{57}\text{Co}$  122 keV calibration data taken every (bi-)week, from  $Z=-40\text{cm}$  to  $+40\text{cm}$
- Intrinsic light yield of the liquid xenon scintillator, absorption and scattering length for the scintillation light extracted from the data/MC comparison



Top PMT can be removed

We take  $^{57}\text{Co}$  calibration data every week and we observed p.e. yield changes :

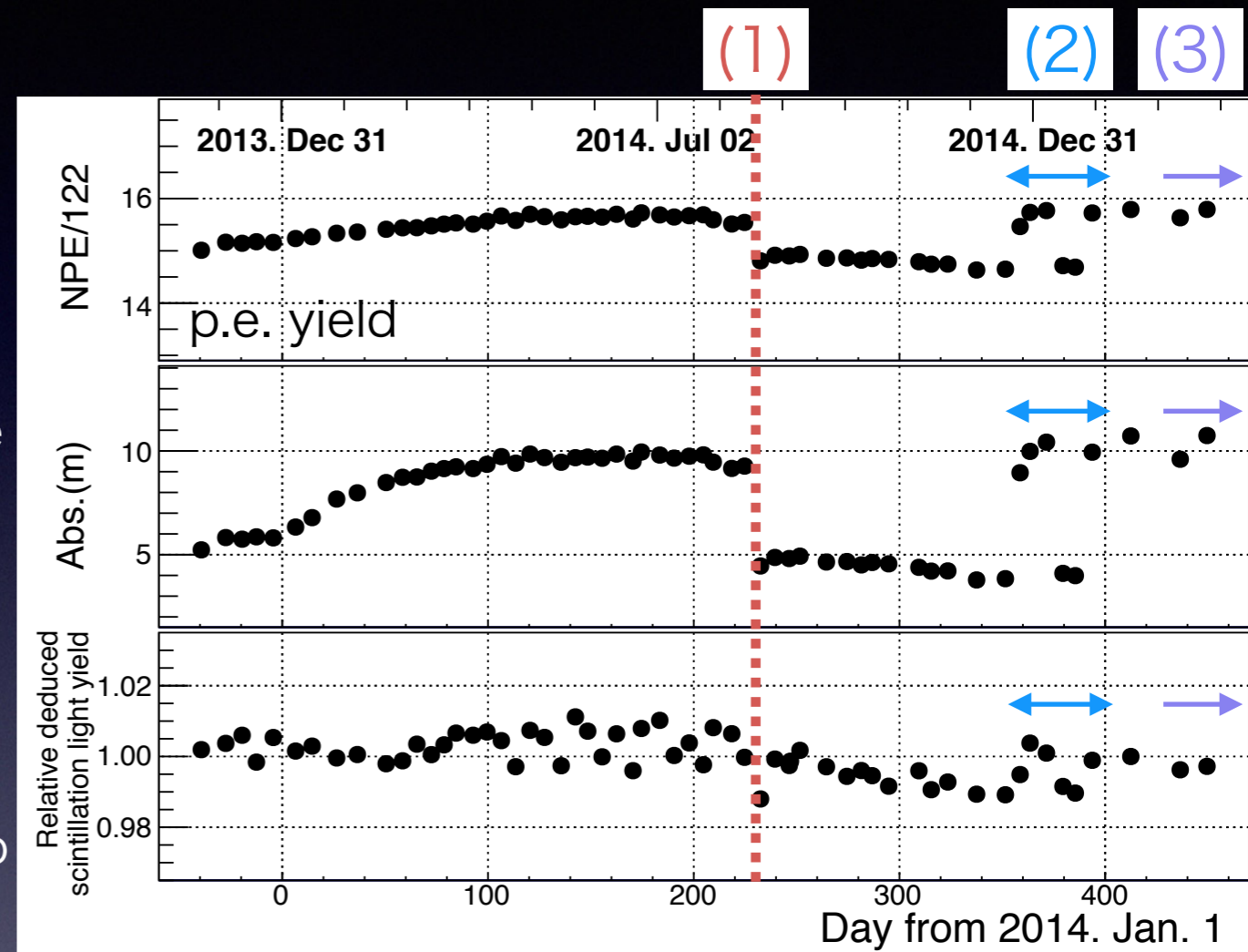
- 1) sudden drop at the power failure
- 2) It recovered after purification work in gas phase
- 3) we continuously circulate the gas purification

We can trace observed p.e. yield change as a changes the absorption length.

Absorption length change : 4m ~ 11m.

Uncertainties due to this instability is taken into account.

Relative intrinsic light yield : stayed within  $\pm 1\%$

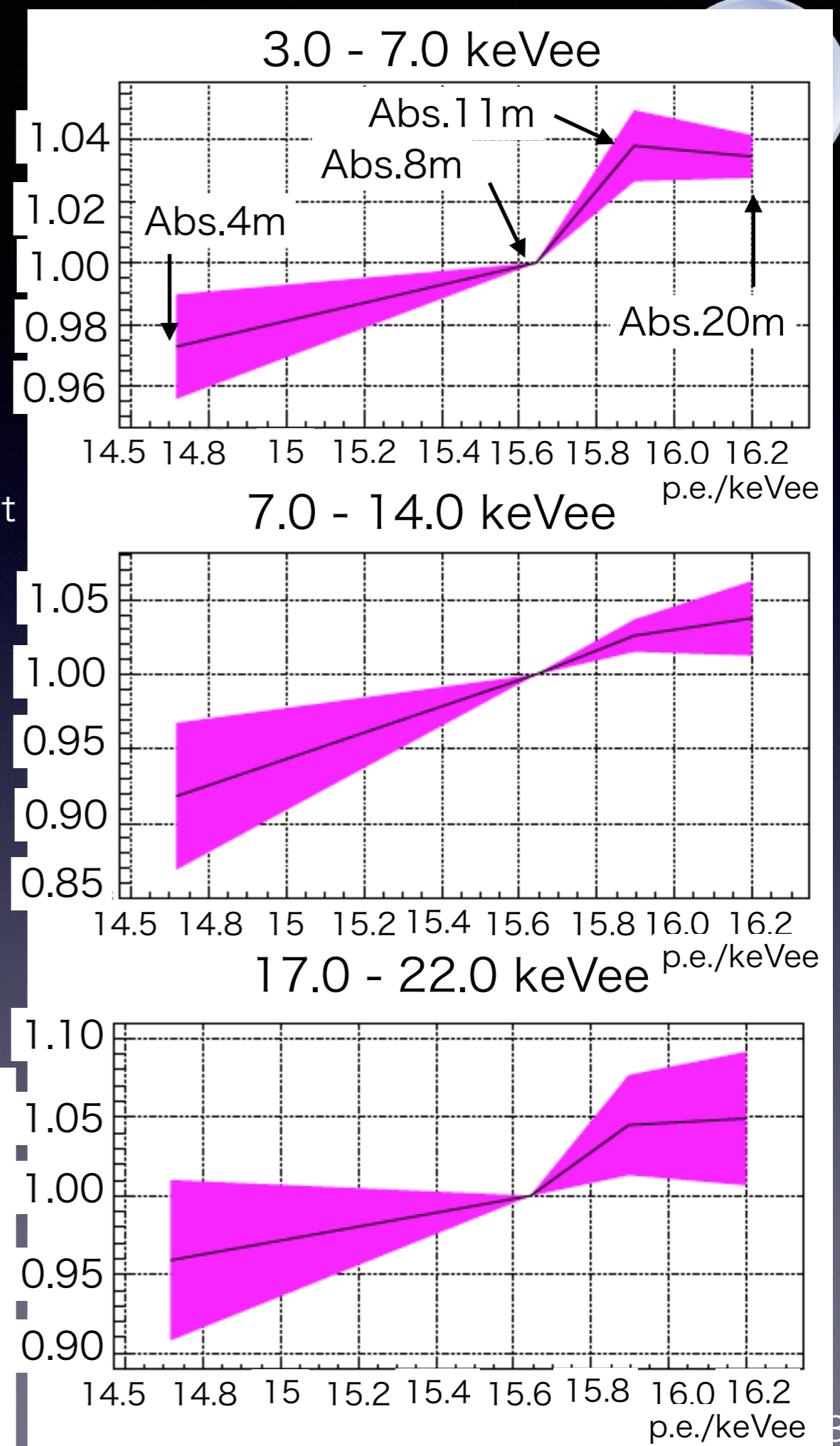
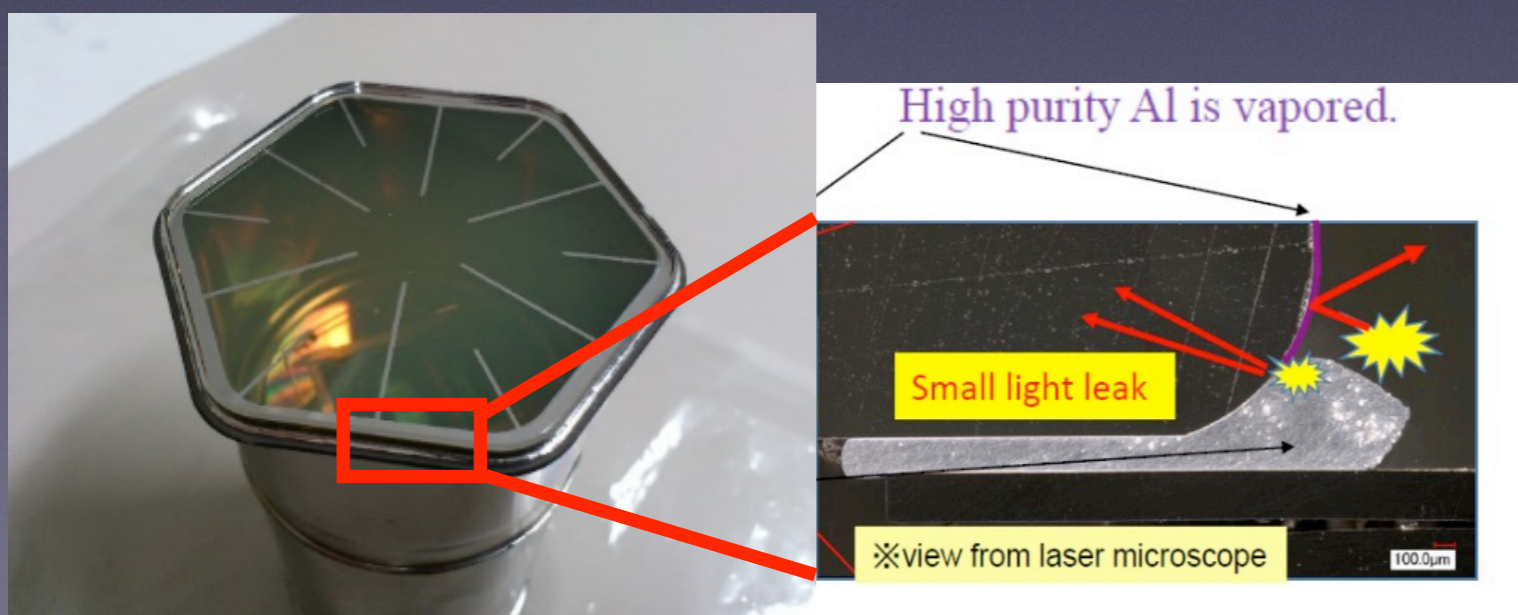




# Relative cut efficiency

- The change of absorption length affects cut efficiency
- The relative change of cut efficiency is evaluated using BGMC, and data is corrected by this obtained relative efficiency
- MC : Dominant BG, summed up corresponds to its amount
  - U-Chain( $^{238}\text{U}$ - $^{230}\text{Th}$ ,  $^{210}\text{Pb}$ ) in the Al seal used for PMT window and body
  - $^{210}\text{Pb}$  in the copper plates on the surface of the ID
- Error band : Al seal shape modeling dependence,  $\pm 5\%$
- Other error, such as non-linearity of scintillation eff are also taken into account.

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# Modulation analysis method



- The data set was divided into 33 time-bin (roughly 15 days each)
- The data in each time-bins were further divided into energy-bin (bin width = 1 keV<sub>ee</sub>)
  - 3-22 keV<sub>ee</sub>, except for 14-17 keV<sub>ee</sub> to avoid systematic effect associated with the end of the range over which the Cherenkov cut is applied
- A least Chi-squares fit all energy/time bins simultaneously to obtain an annual modulation amplitude

$$\chi^2 = \sum_i^{E_{bins}} \sum_j^{t_{bins}} \frac{(R_{i,j}^{data} - R_{i,j}^{ex} - \alpha K_{i,j})^2}{\sigma_{stat;i,j}^2 + \sigma_{sys;i,j}^2} + \alpha^2 + \beta^2$$

K<sub>i,j</sub> : relative eff. systematic error  
α, β : pull term

$$R_{i,j}^{ex} = \int_{t_j - \frac{1}{2}\Delta t_j}^{t_j + \frac{1}{2}\Delta t_j} [C_i + \xi \times (A_i - \beta L_i) (\cos \frac{2\pi(t - t_0)}{T} + \frac{5}{2} e \cos^2 \frac{2\pi(t - t_0)}{T})]$$

**Modulation term**

C<sub>i</sub> : constant term

A<sub>i</sub> : expected amplitude

L<sub>i</sub> : non-linearity of the scintillation eff.

ξ : the ratio of the expected amplitude between the data and the considered model

$$\xi = \frac{g_{a\gamma\gamma}^2}{(9.2 \times 10^{-14} GeV^{-1})^2} \frac{\bar{n}_a}{(4.07 \times 10^{13} m^{-3})}$$

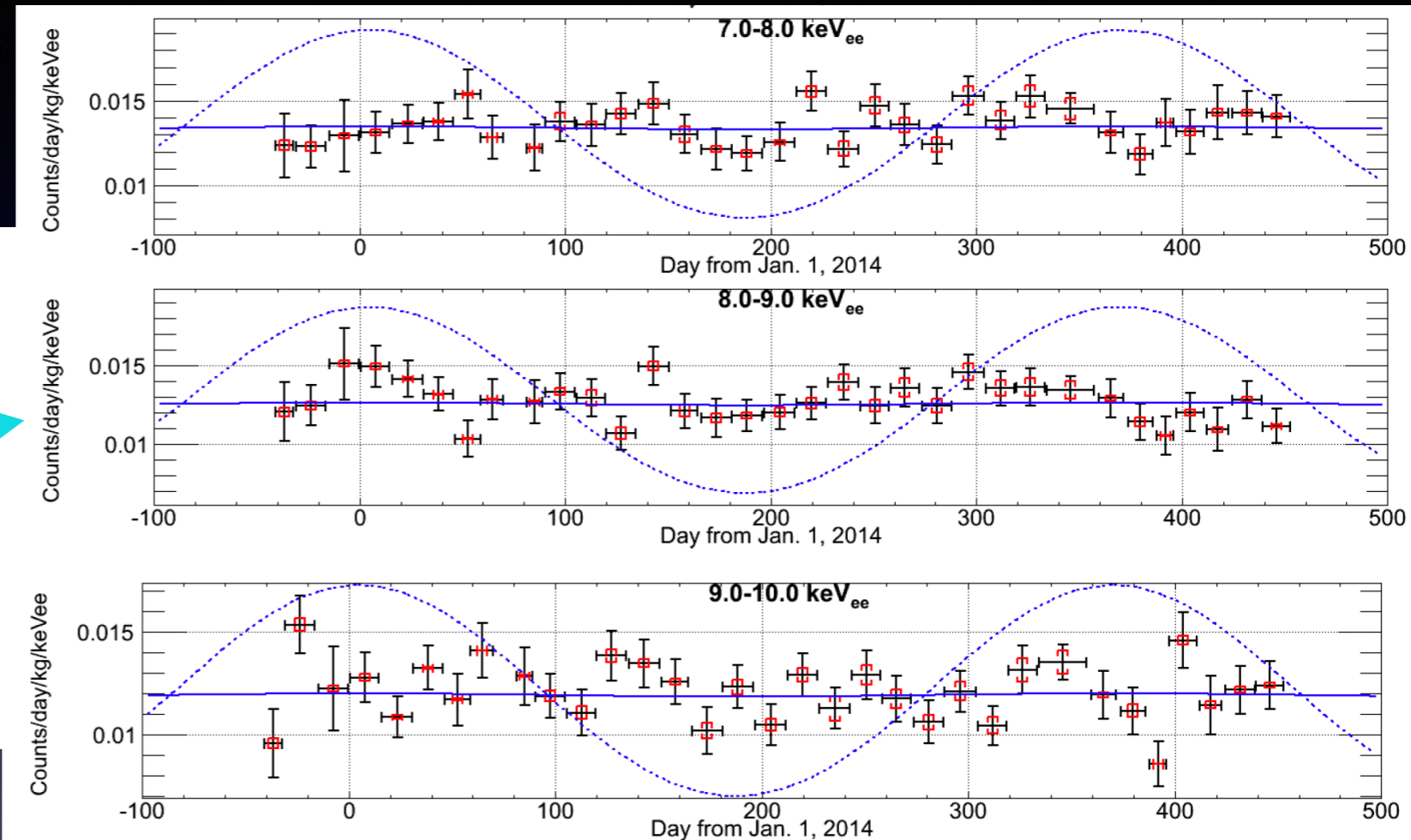
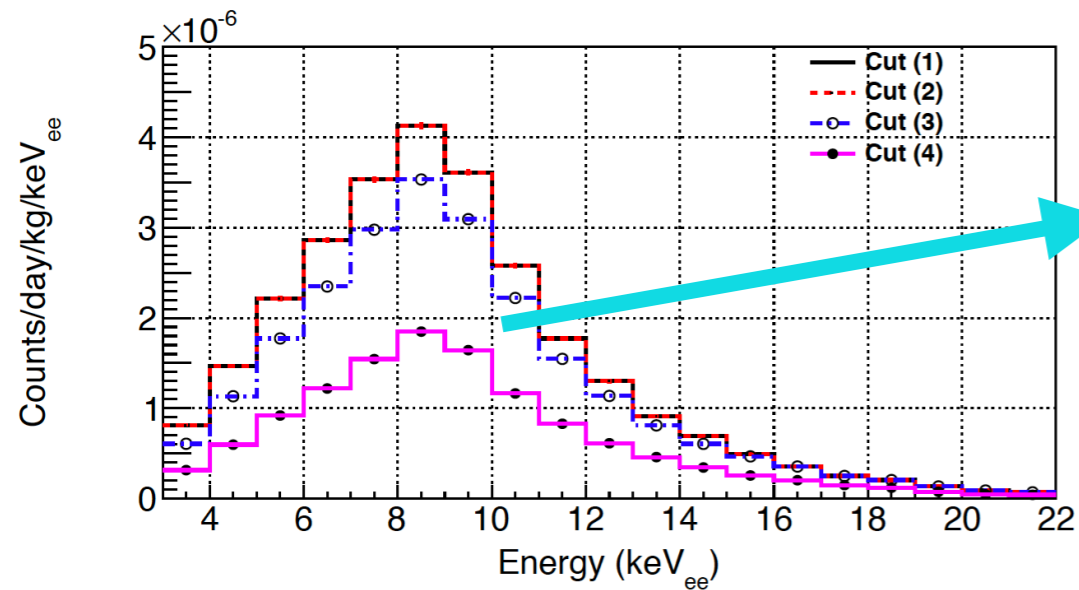
Free parameter : C<sub>i</sub> and ξ



# Time variation of event rate



## Time variation



- No significant excess in amplitude is found
- Best fit  $\xi = 8.2 \times 10^2$  with  $\chi^2/\text{ndf} = 522.4/492$
- 90% CL upper limit  $\xi = 2.7 \times 10^3$
- $g_{ar r} < 4.8 \times 10^{-12} \text{ GeV}^{-1}$  for  $\bar{n}_a = 4.07 \times 10^{13} \text{ m}^{-3}$

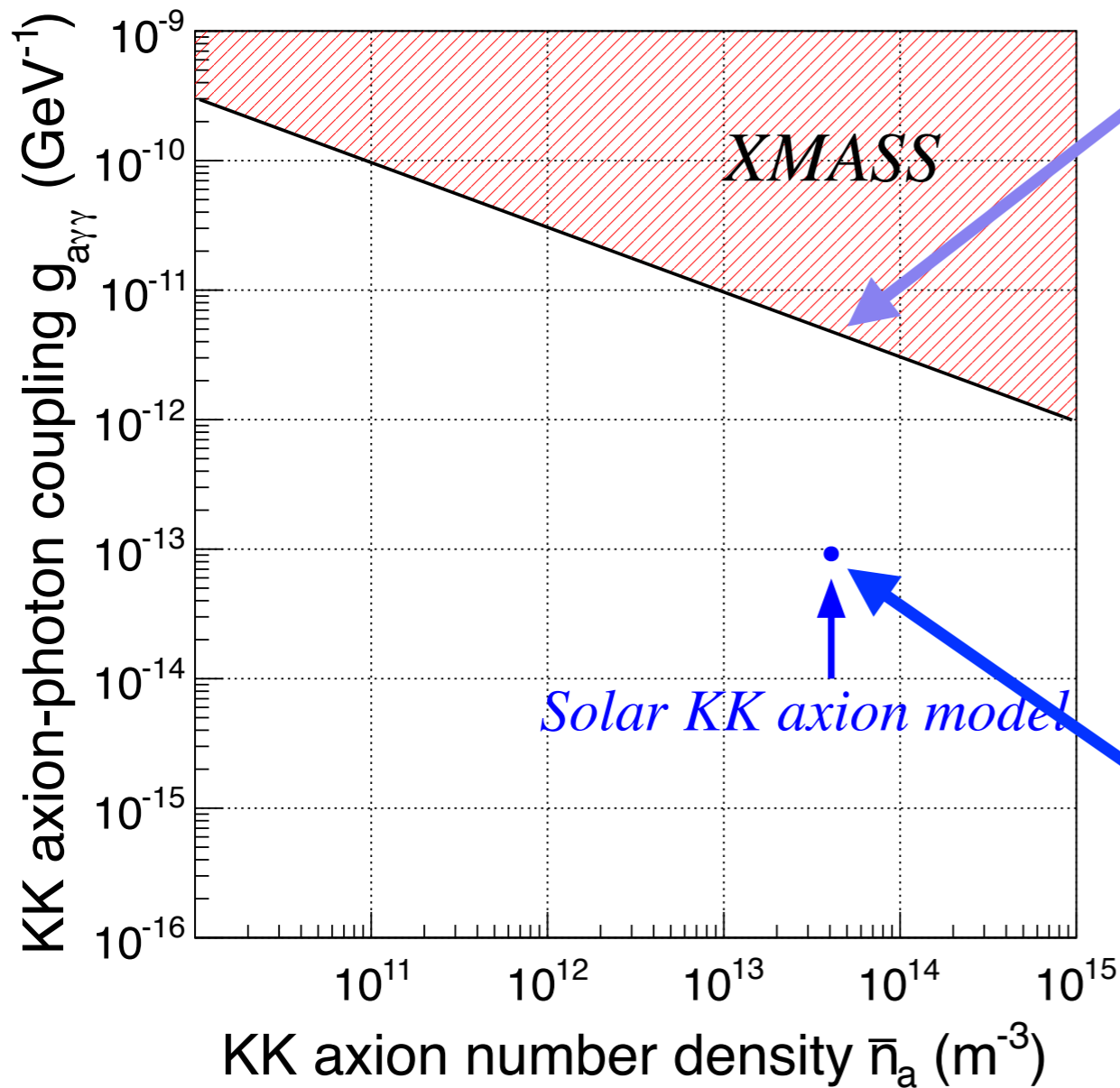
data

syst.err

(solid):best fit result  
(dot) : 90% CL upper  
limit, x20 enhanced



# Result



- $g_{a\gamma\gamma} < 4.8 \times 10^{-12} \text{ GeV}^{-1}$  for  $\bar{n}_a = 4.07 \times 10^{13} \text{ m}^{-3}$
- Rate  $< 234 \text{ m}^{-3} \text{ day}^{-1}$  (90% CL)
- First experimental constraint for KK axions
- Submitted to “Progress of Theoretical and Experimental Physics (PTEP)”

Blue point as a benchmark  
 from L. Di Lella and K.Zioutas APP 19 (2003)  
 145  
 To explain X-ray surface brightness of the  
 quiet Sun

$$R = (2.5 \times 10^{11} \text{ m}^{-3} \text{ day}^{-1}) \left( \frac{g_{a\gamma\gamma}}{\text{GeV}^{-1}} \right)^2 \left( \frac{n_a}{\text{m}^{-3}} \right)$$





# Summary



- Solar KK axion would be produced in the Sun and a small fraction is trapped by the gravity of the Sun. Decays into two photons.
- XMASS searched the decay of solar KK axions by annual modulation (832x359 kg · days)
- No significant excess in amplitude is found, and we set 90% CL upper limit :
- $g_{a\gamma\gamma} < 4.8 \times 10^{-12} \text{ GeV}^{-1}$  for  $\bar{n}_a = 4.07 \times 10^{13} \text{ m}^{-3}$
- **This is the First experimental constraint for KK axions.**
- Submitted to Progress of Theoretical and Experimental Physics (PTEP)



# Backup



# KK axion lifetime

standard electromagnetic coupling

Axion couples to two photons:  $g_{\alpha\gamma\gamma} = \frac{\alpha_{EM}}{\pi} \frac{C_a}{f_{PQ}}$

axion model factor  $\sim 1$

symmetry breaking energy scale

This implies decay to two photons with mean lifetime:  $\tau = \frac{64\pi}{g_{\alpha\gamma\gamma}^2 m_a^3}$

However, astrophysical constraints imply  $\tau$  too long to observe:

$$10^9 \text{ GeV} \leq f_{PQ} \leq 10^{12} \text{ GeV}; 10^{30} \leq t \leq 10^{45} \text{ days}$$

But propagation in **extra dimensions** allows shorter, observable, lifetime

$$m_a = m_{an} \sim \frac{n}{R}$$

# Solar KK axion mass spectrum

## Basis for an experimental search:

**B. Morgan, N. Spooner et al, D. Hoffmann et al., K. Zioutas...**

**B. Morgan et al. *Astrop. Phys* 23 (2005) 287,**

- Leads to differential decay spectrum: 
$$\frac{dR}{dm_a} = \frac{g_{a\gamma\gamma}^2}{64\pi} n_0 m_a^3 f(m_a)$$

$$R = (2.5 \times 10^{11} m^{-3} day^{-1}) \left( \frac{g_{a\gamma\gamma}}{GeV^{-1}} \right)^2 \left( \frac{n_0}{m^{-3}} \right) \longrightarrow \text{Typical rate } \sim 1 m^{-3} day^{-1} \text{ } (\sim keV \text{ events})$$

**Result for trapped axions in orbits around Sun**  $g_{a\gamma\gamma} = 9.2 \times 10^{-14} GeV^{-1}$   $n_0 = 10^{14} m^{-3}$   
(local number density depends on  $g_{a\gamma\gamma}$ )

**Mass spectrum for solar axions trapped in orbits around the sun**

**L. Di Lella and K. Zioutas, *Astrop. Phys.*, 19 (2003) 145**

