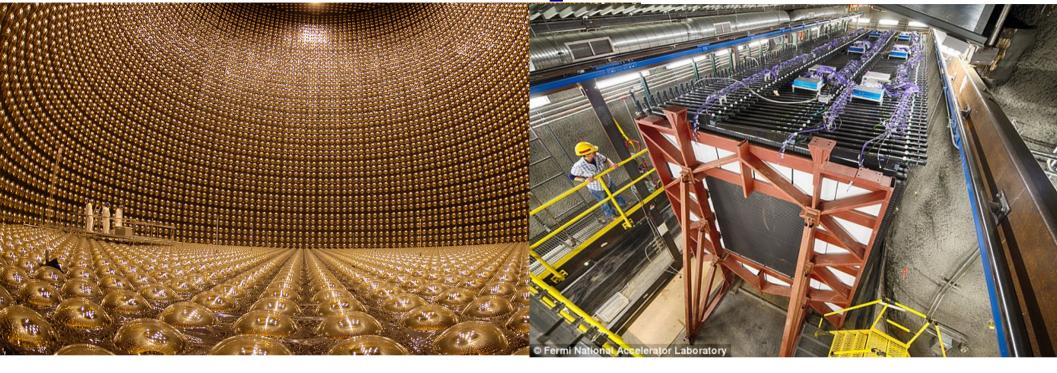
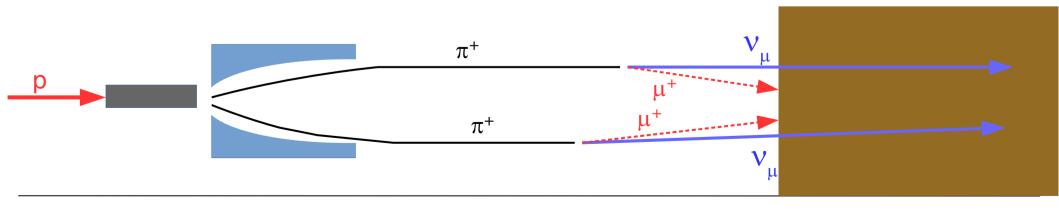
Current Long Baseline Neutrino Experiments





Scott Oser

(UBC/TRIUMF)

TAUP 2017 July 26, 2017

The neutrino 3x3 mixing matrix

Different L/E values pick up different Δm^2 pairs, probing different parts of mixing matrix.

$$\begin{pmatrix}
1 & 0 & 0 \\
0 & c_{23} & s_{23} \\
0 & -s_{23} & c_{23}
\end{pmatrix}$$

$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & e^{i\delta}s_{13} \\ 0 & 1 & 0 \\ -e^{-i\delta}s_{13} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\begin{pmatrix}
c_{12} & s_{12} & 0 \\
-s_{12} & c_{12} & 0 \\
0 & 0 & 1
\end{pmatrix}$$

Atmospheric ν 's:

$$\theta_{23} \approx 46^{\circ}$$

Maximal mixing! (?)

Short baseline reactor ν 's:

$$\theta_{13} \approx 9^{\circ}$$

Small, quark-like mixing

Solar ν 's:

$$\theta_{12} \approx 33^{\circ}$$

Large, non-maximal mixing

Compare to identical parameterization of CKM matrix ...

$$\theta_{23} \approx 2.4^{\circ}$$

$$\theta_{13} \approx 0.2^{\circ}$$

$$\theta_{12} \approx 13^{\circ}$$

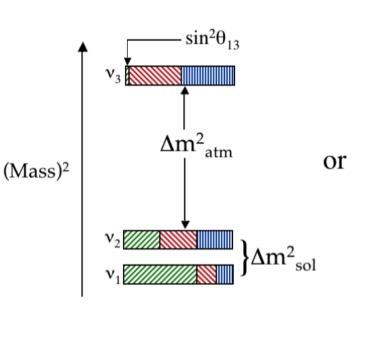
Mass Hierarchy

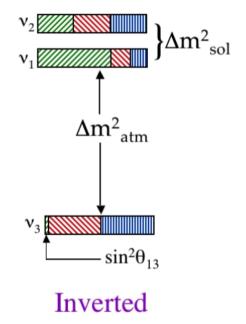
Currently unknown:

- ullet value of $\delta_{\sf CP}$
- sign of the mass hierarchy

$$|\Delta m^2_{32}| = 2.4 \times 10^{-3} \, eV^2$$

$$\Delta m_{21}^2 = 7.5 \times 10^{-5} \text{ eV}^2$$



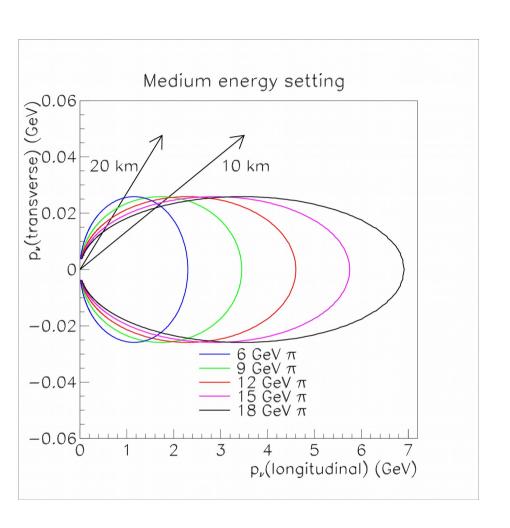


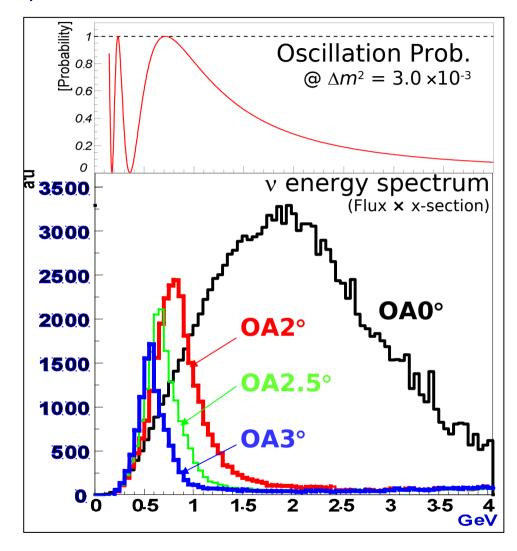
Normal





Off-Axis ν_{μ} Beam





Off-axis beam: more flux near peak oscillation energy, less flux at higher energies where $v_{\rm e}$ backgrounds are produced.

LBL signature #1: ν_{μ} disappearance

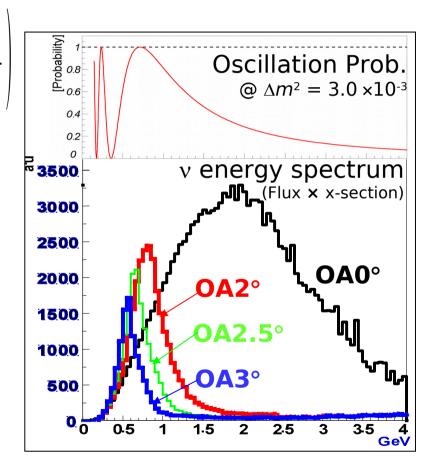
Starting from
$$P_{
u_a o
u_b}(L,E) = \left| \sum_{j,k} U_{aj}^* U_{bj} U_{ak} U_{bk}^* e^{-i rac{\Delta m_{jk}^2 L}{2E}} \right|$$

$$P(\nu_{\mu} \rightarrow \nu_{\mu}) \approx 1 - \cos^{4}\theta_{13} \sin^{2}2\theta_{23} \sin^{2}\left|\frac{\Delta m_{32}^{2} L}{4 E}\right|$$

$$-\sin^{2}2\theta_{13} \sin^{2}\theta_{23} \sin^{2}\left|\frac{\Delta m_{32}^{2} L}{4 E}\right|$$

$$\approx 1 - \sin^{2}2\theta_{23} \sin^{2}\left|\frac{\Delta m_{32}^{2} L}{4 E}\right|$$

Sensitive to Δm_{32}^2 and θ_{23} . Same formula for neutrinos and for antineutrinos, if CPT holds.



LBL signature #2: v_e appearance

$$P(\nu_{\mu} \to \nu_{e}) \sim \frac{\sin^{2} 2\theta_{13}}{-\alpha \sin^{2} 2\theta_{13}} \times \sin^{2} \theta_{23} \times \frac{\sin^{2}[(1-x)\Delta]}{(1-x)^{2}}$$

$$-\alpha \sin 2\theta_{13} \times \sin \delta \sin 2\theta_{12} \sin 2\theta_{23} \times \sin \Delta \frac{\sin[x\Delta]}{x} \frac{\sin[(1-x)\Delta]}{(1-x)}$$

$$+\alpha \sin 2\theta_{13} \times \cos \delta \sin 2\theta_{12} \sin 2\theta_{23} \times \cos \Delta \frac{\sin[x\Delta]}{x} \frac{\sin[(1-x)\Delta]}{(1-x)}$$

$$x \equiv 2\sqrt{2}G_{F}N_{e}\frac{E_{\nu}}{\Delta m_{31}^{2}} \Delta \equiv \frac{\Delta m_{31}^{2}L}{4E_{\nu}}$$

Dominant term corresponds to a ~5% transition probability at the oscillation maximum

LBL signature #2: v_e appearance

$$P(\nu_{\mu} \rightarrow \nu_{e}) \sim \sin^{2}2\theta_{13} \times \sin^{2}\theta_{23} \times \frac{\sin^{2}[(1-x)\Delta]}{(1-x)^{2}} \times \sin\Delta \frac{\sin[x\Delta]}{x} \frac{\sin[(1-x)\Delta]}{(1-x)} \times \sin\Delta \frac{\sin[x\Delta]}{x} \frac{\sin[(1-x)\Delta]}{(1-x)} \times \cos\Delta \frac{\sin[x\Delta]}{x} \frac{\sin[(1-x)\Delta]}{(1-x)} \times \sin\Delta \frac{\sin[x\Delta]}{x} \frac{\sin[(1-x)\Delta]}{(1-x)} \times \cos\Delta \frac{\sin[x\Delta]}{x} \frac{$$

Terms containing δ are sensitive to CP phase. The δ 's flip sign for antineutrinos.

The x parameter (matter effect) also flips sign for antineutrinos. The matter effect is subdominant at T2K due to low beam energy, but larger at NOvA.

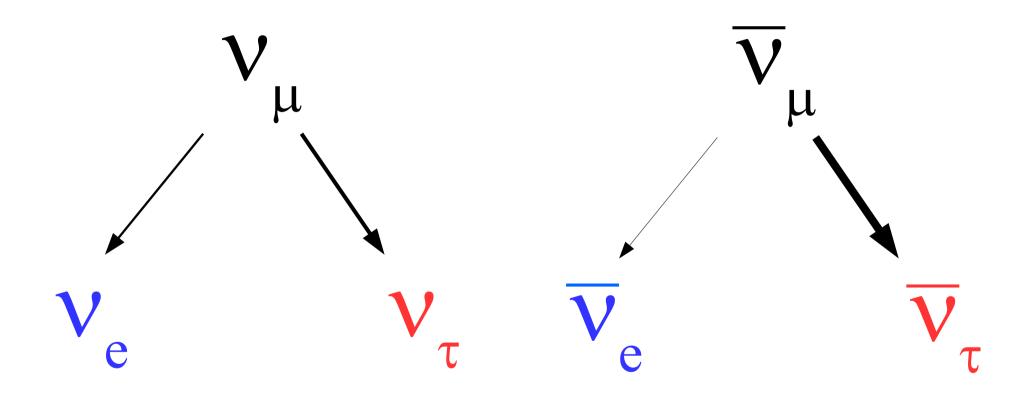
CP, T, and CPT for neutrinos

$$\begin{split} &P\left(\mathbf{v}_{a}\!\rightarrow\!\mathbf{v}_{b}\right) \longleftarrow P\left(\bar{\mathbf{v}}_{a}\!\rightarrow\!\bar{\mathbf{v}}_{b}\right) \\ &P\left(\mathbf{v}_{a}\!\rightarrow\!\mathbf{v}_{b}\right) \longleftarrow P\left(\mathbf{v}_{b}\!\rightarrow\!\mathbf{v}_{a}\right) \\ &P\left(\mathbf{v}_{a}\!\rightarrow\!\mathbf{v}_{b}\right) \longleftarrow P\left(\bar{\mathbf{v}}_{b}\!\rightarrow\!\bar{\mathbf{v}}_{a}\right) \\ &P\left(\mathbf{v}_{a}\!\rightarrow\!\mathbf{v}_{b}\right) \longleftarrow P\left(\bar{\mathbf{v}}_{a}\!\rightarrow\!\bar{\mathbf{v}}_{a}\right) \end{split}$$

If CPT holds in neutrino sector, neutrino survival probability equals antineutrino survival probability.

As a result CP violation is observable only in appearance channels, in which the flavour of the appearing lepton is detected. This is why LBL neutrinos are interesting!

CP and three-flavour mixing



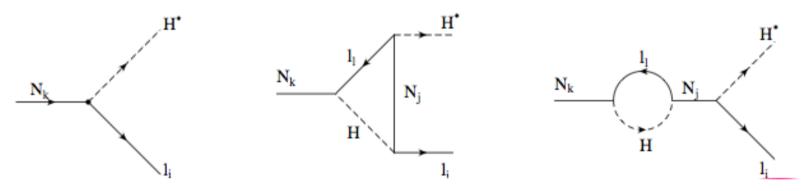
CP violation affects the mix of flavours that results from the oscillation

Leptogenesis

CP violation in quark sector not enough to explain observed matter-antimatter asymmetry in universe.

Neutrino mixing provides another possible source of CPV.

Standard Leptogenesis: decays of RH neutrinos (CPV in decay)
 Quantum interference of tree diagram and one-loop diagram



Usual scenario: decay of heavy Majorana neutrinos Phys.Lett B 174, 45 (1986) Many alternates, eg. leptogenesis with only Dirac v's PRL 89:271601 (2002)

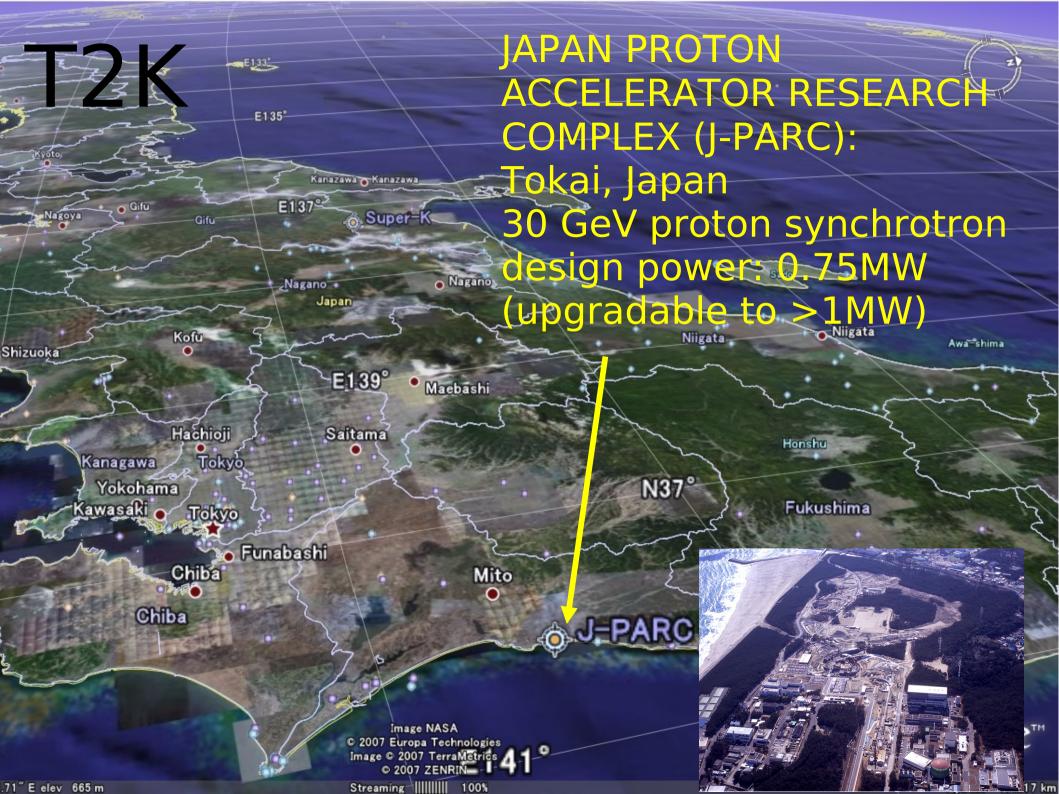
Relation of δ_{CP} to leptogenesis is model-dependent, but observation of leptonic CP violation is an important milestone.

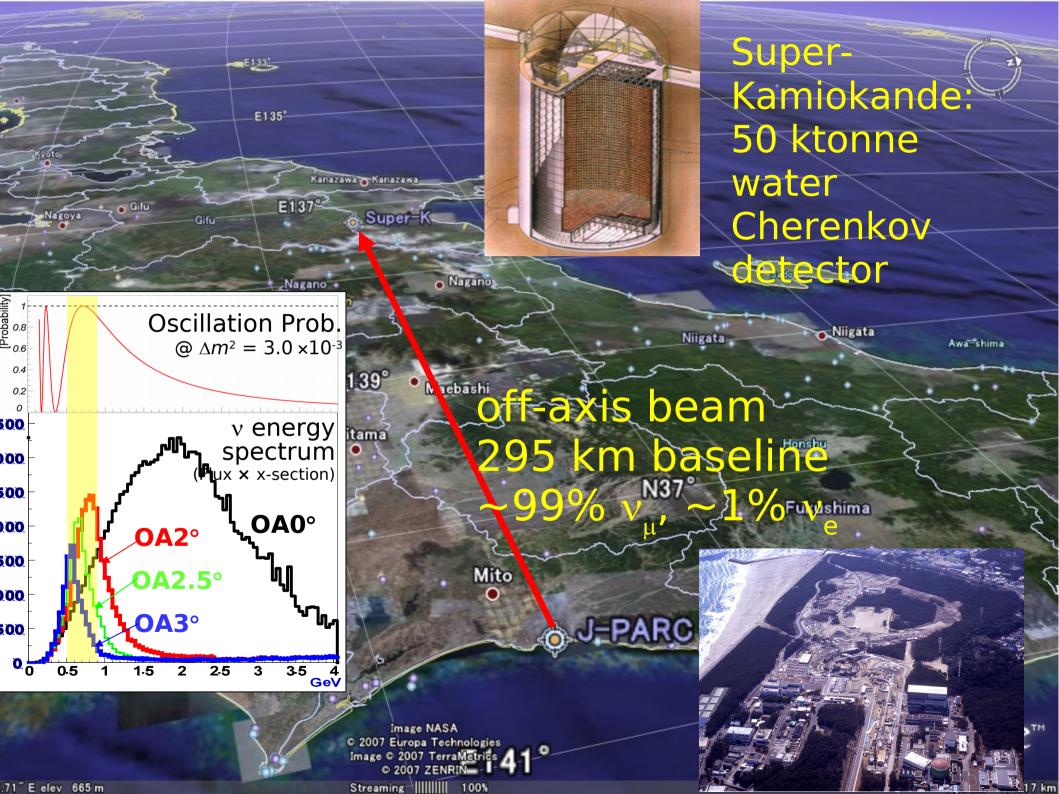
Matter Effects and v_e Appearance

Matter effects modify the oscillation formula. Because the Earth is made of electrons and not heavier leptons, the effective "index of refraction" for v_e is different than that for v_μ . At the oscillation maximum, the v_e appearance probability changes to:

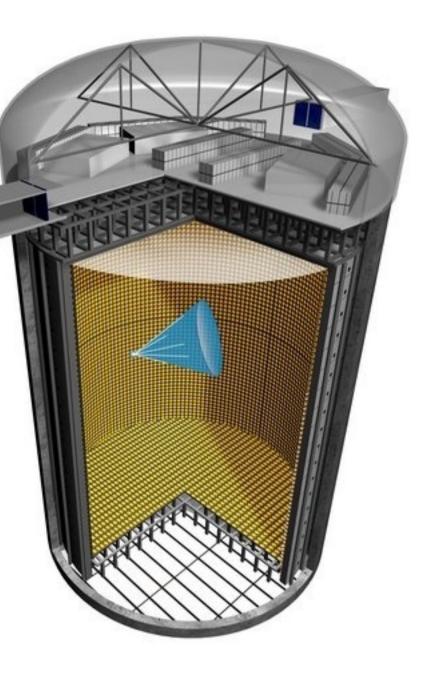
$$P(\mathbf{v}_{\mu} \rightarrow \mathbf{v}_{e}) \approx \left| 1 + 2 \frac{E}{E_{R}} \right| P_{vac}(\mathbf{v}_{\mu} \rightarrow \mathbf{v}_{e})$$
where
$$E_{R} = \frac{\Delta m_{32}^{2}}{2\sqrt{2} G_{F} N_{e}} = \pm 11 GeV$$

The sign of the matter effect is opposite for neutrinos and antineutrinos, and depends on the sign of Δm^2 as well.

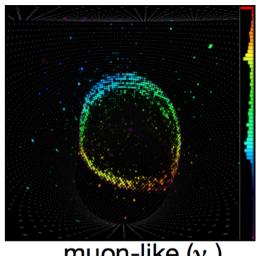




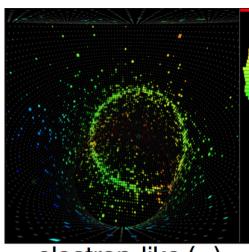




Super-Kamiokande



muon-like (v_{\parallel})

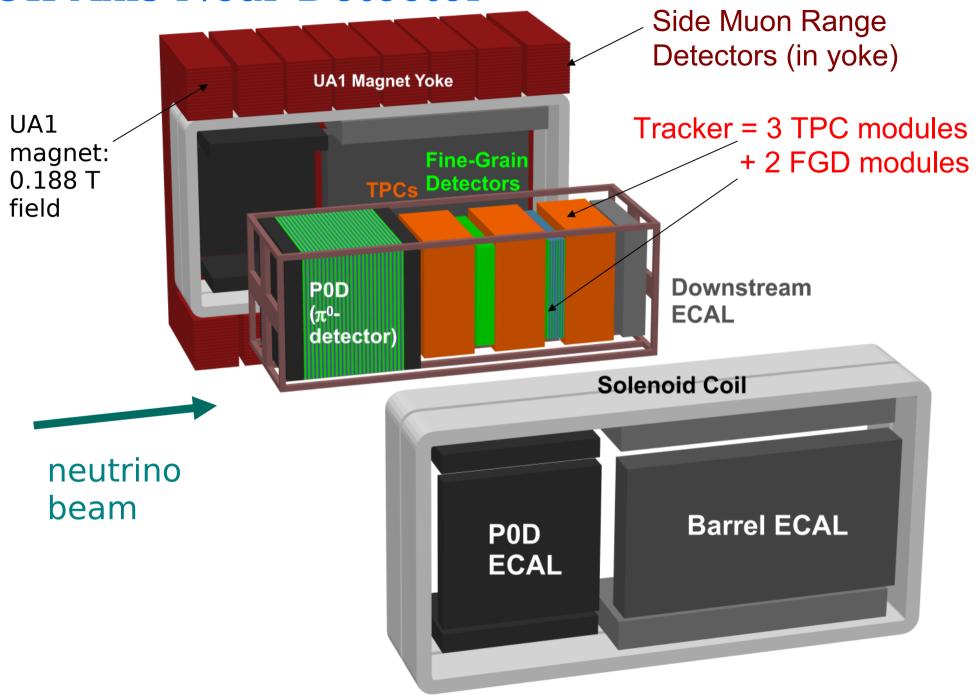


electron-like (v_e)

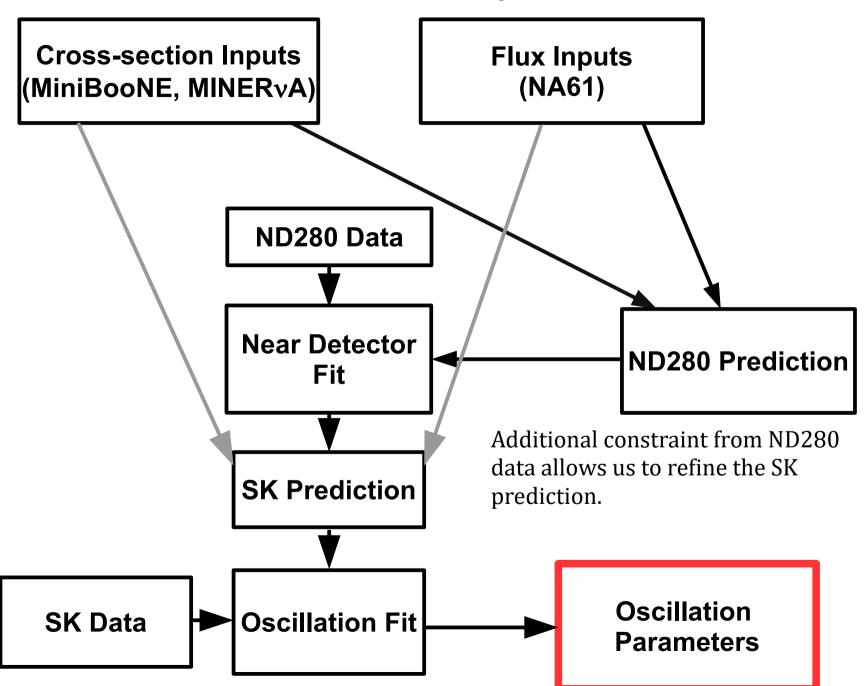
Water Cherenkov detection.

Primary signal channel is CCQE single-ring events. Reconstruct v energy in CCQE hypothesis from lepton kinematics.

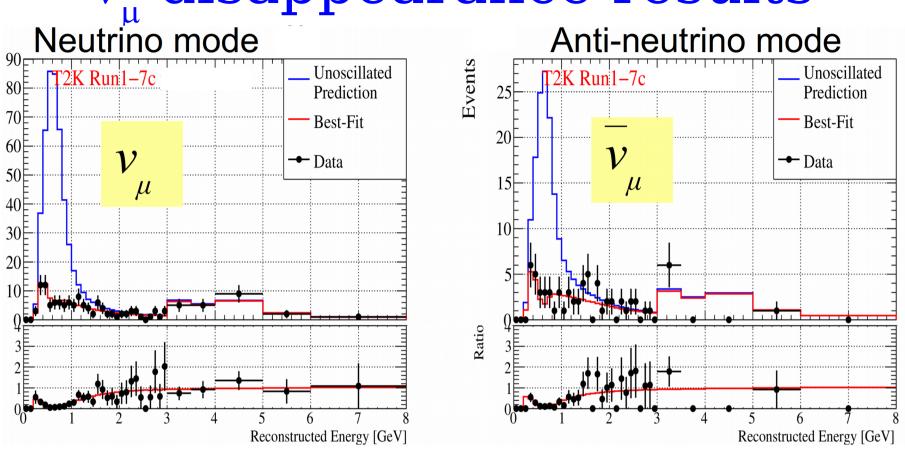
Off Axis Near Detector



T2K Oscillation Analysis Structure



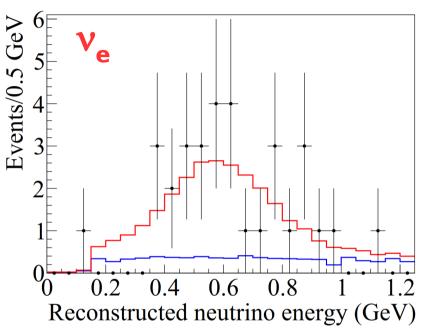
v, disappearance results

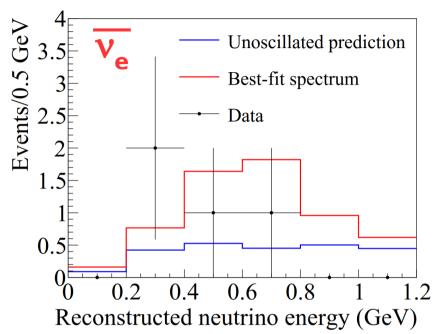


Beam mode	Sample	Exp. Not Osc	Exp. $\delta_{CP} = 0$ (NH)	Observed
neutrino	μ -like	521.8	135.5	135
antineutrino	μ -like	184.8	64.1	66

 7.5×10^{20} protons on target of ν data + 7.5×10^{20} protons on target of anti- ν data

$\nu_{\rm e}$ appearance results



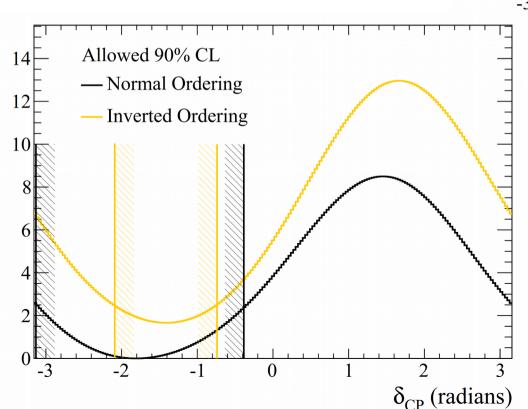


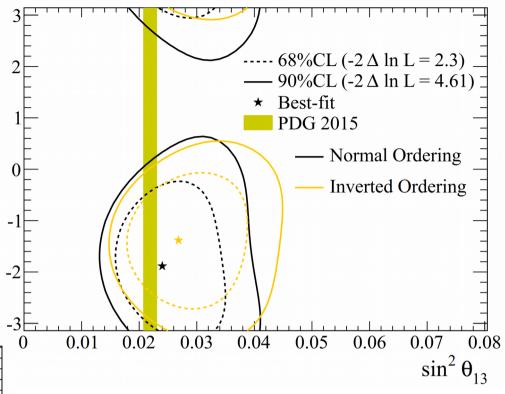
Normal	$\delta_{CP} = -\pi/2$	$\delta_{CP} = 0$	$\delta_{CP} = \pi/2$	$\delta_{CP} = \pi$	Observed
$ u_e$	28.7	24.2	19.6	24.1	32
$\overline{ u}_e$	6.0	6.9	7.7	6.8	4
Inverted	$\delta_{CP} = -\pi/2$	$\delta_{CP} = 0$	$\delta_{CP} = \pi/2$	$\delta_{CP} = \pi$	Observed
$ u_e$	25.4	21.3	17.1	21.3	32
$\overline{ u}_e$	6.5	7.4	8.4	7.4	4

δ_{CP} contours

The best-fit CP phase is close to $-\pi/2$: maximal CP effect. Formally speaking $\delta_{\text{CP}} = 0$ is excluded at 90% CL.

 δ_{CP} (radians)

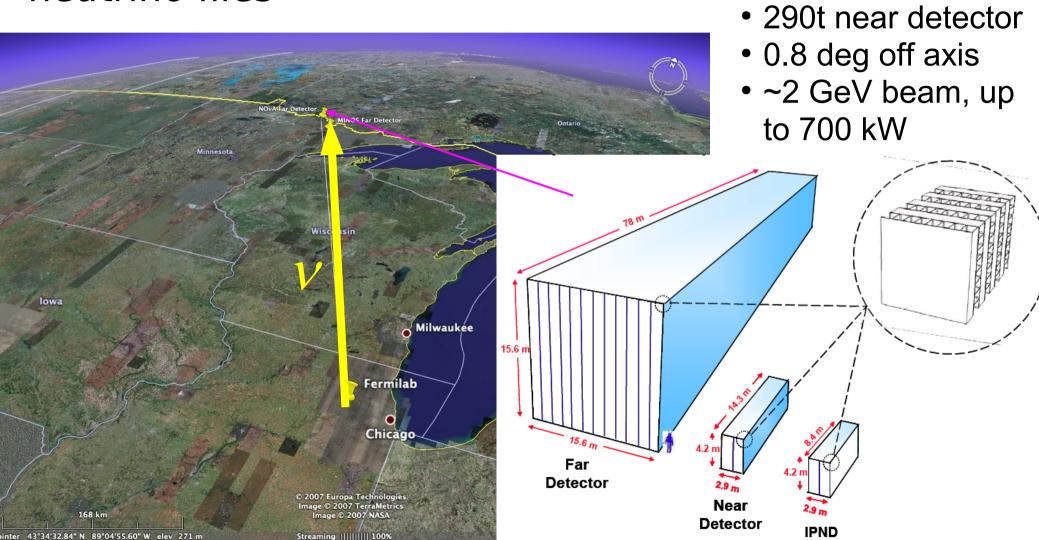




Caveat: the result is primarily driven by very high v_e appearance rate at T2K, beyond expectations of model. In other words, limit is better than expected sensitivity.

NOνA

FNAL to Ash River: 810km as the neutrino flies



• 14 kt liquid

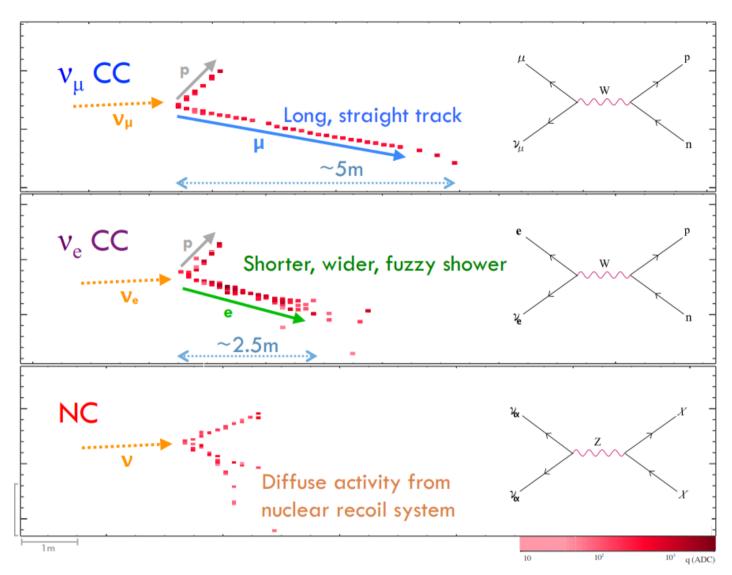
scintillator tracker

NOvA vital characteristics

Flavour ID with NC sensitivity

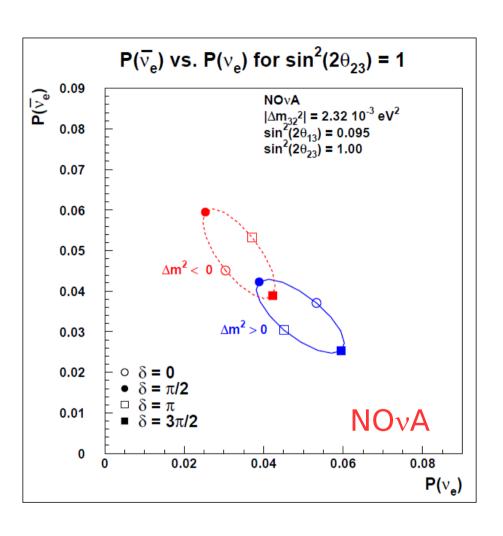
Reconstruction informed by computer vision, machine learning

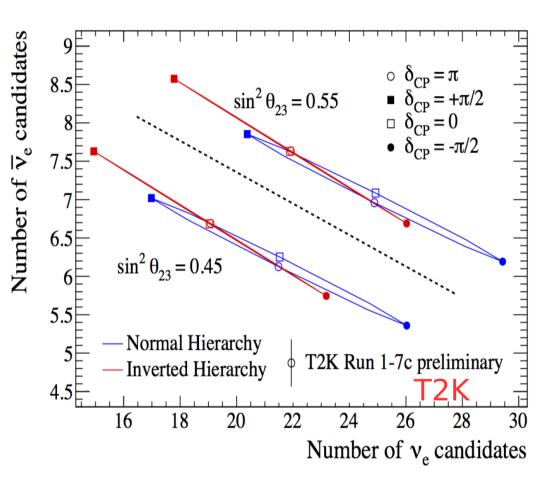
Beam energy of ~2GeV more sensitive than



T2K to hadronic production. Energy estimation is more calorimetric than T2K, where hadrons usually are below detection threshold. (But hadrons are harder to model.)

NOvA matter effect sensitivity

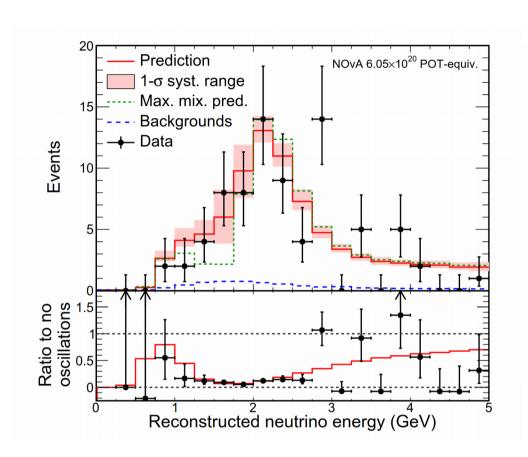


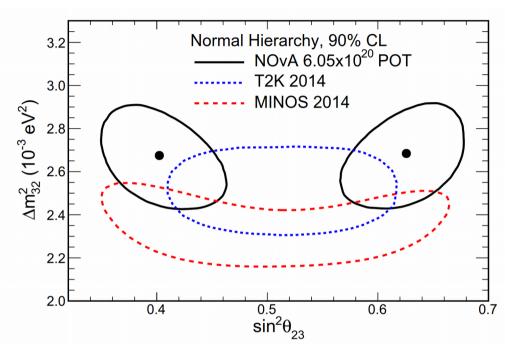


$$P(\nu_{\mu} \rightarrow \nu_{e}) \approx \left| 1 + 2 \frac{E}{E_{R}} \right| P_{vac}(\nu_{\mu} \rightarrow \nu_{e})$$

Higher energy gives NOvA better sensitivity to matter effects than T2K

NOvA Oscillation Results





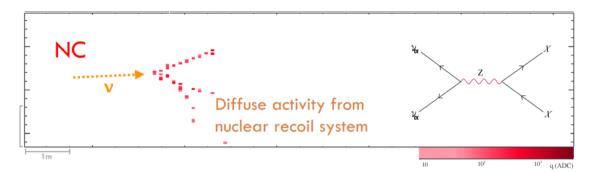
 $(6.05 \times 10^{20} \text{ p.o.t. in neutrino mode})$

33 v_e events on background of 8.2±0.8 Maximal mixing disfavoured at 2.6 σ

Phys. Rev. Lett. 118, 231801 (2017) arXiv: 1703.03328 Phys. Rev. Lett. 118, 151802 (2017), arXiv: 1701.05891

NOvA sterile neutrinos results

Neutral current sensitivity is unique capacity for NOvA. A convolutional neural network is used to distinguish events with no

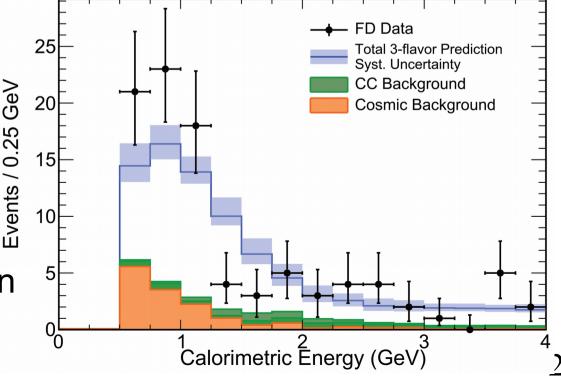


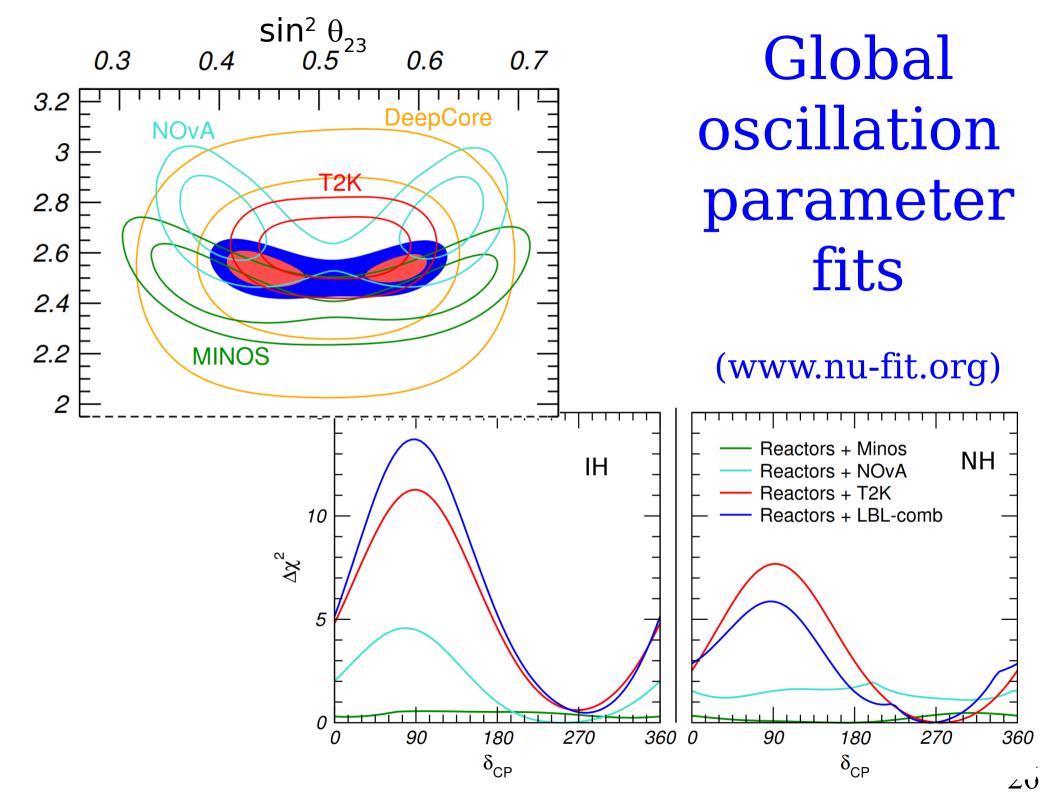
distinguish events with no charged lepton.

If neutrinos mix to a sterile state, then the NC rate may be less

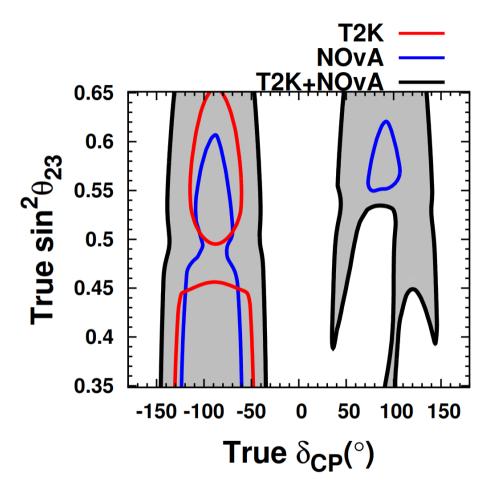
than expected.

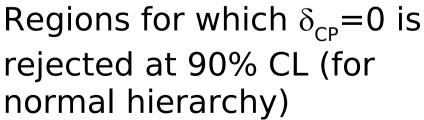
Recent NOvA arXiv submission 1706.04592 saw 95 events where 83.5 ± 9.7(stat) ± 9.4(syst) were predicted assuming mixing only occurs between active neutrino species.

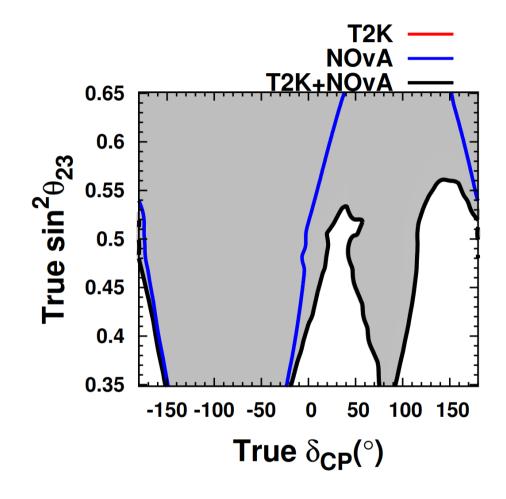




T2K+NOvA combined sensitivity

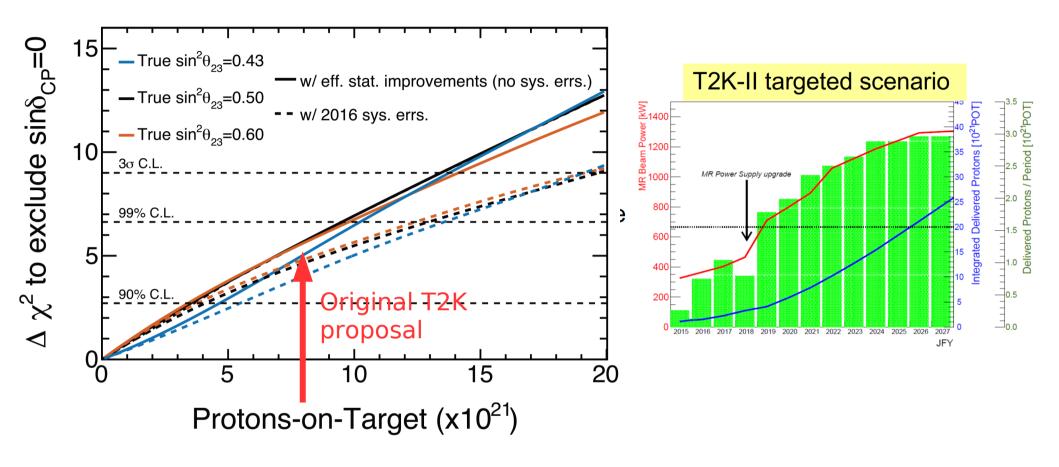






Regions for which the wrong mass hierarchy is rejected at 90% CL (for normal hierarchy)

T2K-II



Proposal to continue running T2K until ~2025, with beam power and near detector upgrades, with an aim to achieve 3σ sensitivity to non-zero δ_{CP} .

arXiv: 1609.04111

Conclusions

Long baseline neutrino experiments with flavour sensitivity are the only window we have on CP violation in the neutrino sector.

T2K and NOvA results are consistent with each other and PMNS paradigm. First limits on CP violation, although still statistically weak, favour maximal CP effect.

LBL experiments dominate Δm_{32}^2 and θ_{23} determination.

See parallel session talks by Mark Scott (T2K—Tuesday, 17:15, Neutrino 4), Nicoletta Mauri (OPERA—Tuesday, 17:00, Neutrino 4) and Kirk Bays (NOvA, Wednesday, 13:00, Neutrino 5) for details.

Backup Slides

CP Violation and v_e Appearance

CP symmetry requires $P(\nu_{\mu} \rightarrow \nu_{e}) = P(\bar{\nu}_{\mu} \rightarrow \bar{\nu}_{e})$

For v_e appearance at Δm_{32}^2 :

$$A_{CP} = \frac{P(\nu_{\mu} \rightarrow \nu_{e}) - P(\bar{\nu_{\mu}} \rightarrow \bar{\nu_{e}})}{P(\nu_{\mu} \rightarrow \nu_{e}) + P(\bar{\nu_{\mu}} \rightarrow \bar{\nu_{e}})} \simeq \frac{\Delta m_{12}^{2} L}{4 E_{\nu}} \frac{\sin 2\theta_{12}}{\sin \theta_{13}} \sin \delta_{CP}$$

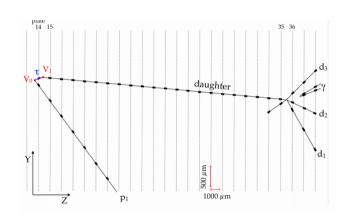
This may be a big asymmetry!

SO WHAT?

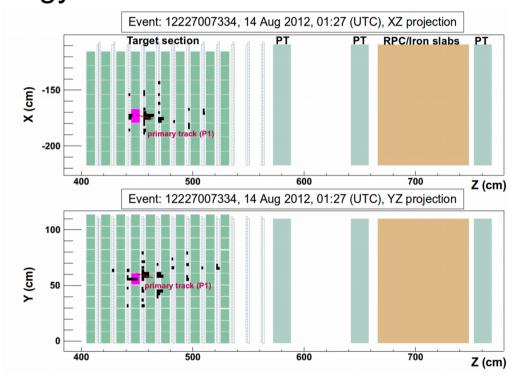
- Our universe is made of matter but not anti-matter.
- A cosmological asymmetry requires CP violation.
- Regular quark CP violation not enough---is this the missing piece?

OPERA

Beam from CERN to Gran Sasso, looking primarily for v_{τ} appearance using emulsion technology.



Five v_{τ} candidates seen on a background of 0.25 ± 0.05, & a claimed significance of 5.1 σ .



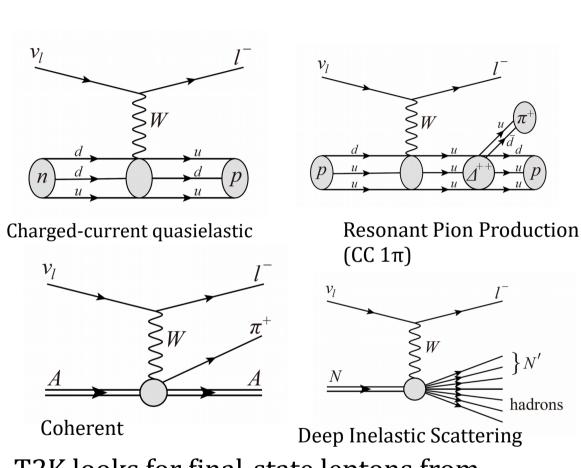
Phys. Rev. Lett. 115, 121802 (2015)

Flavour Oscillation

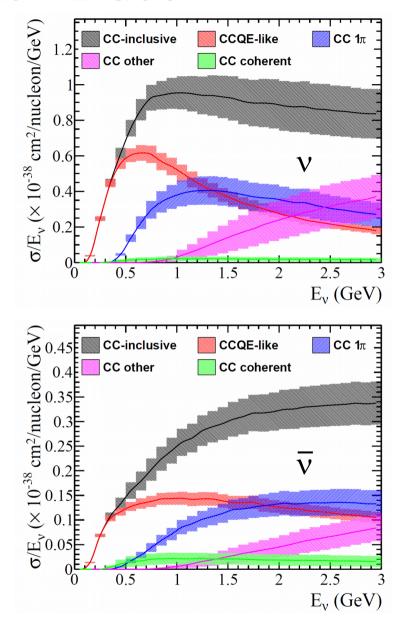
Because a flavour eigenstate produced by a weak interaction is a mix of mass eigenstates which, if $m_1 \neq m_1$, propagate with different kinematics, oscillation can occur.

$$\begin{split} |\mathbf{v}(t=0)\rangle &= |\mathbf{v}_{e}\rangle = \cos\theta \, |\mathbf{v}_{1}\rangle + \sin\theta \, |\mathbf{v}_{2}\rangle \\ |\mathbf{v}(t)\rangle &= \frac{e^{i\sqrt{p^{2}+m_{1}^{2}}t}\cos\theta \, |\mathbf{v}_{1}\rangle}{+e^{i\sqrt{p^{2}+m_{2}^{2}}t}\sin\theta \, |\mathbf{v}_{2}\rangle} \\ Prob \, (\mathbf{v}_{e}\rightarrow\mathbf{v}_{e}) &= 1-\sin^{2}(2\,\theta)\sin^{2}\!\left(\frac{1.27\,\Delta\,m^{2}\,L}{E}\right) \\ \text{Units: [L] = km; [E] = GeV;} \\ \Delta m^{2} &= [\text{eV}^{2}] \end{split}$$

Neutrino interactions at ~1 GeV



T2K looks for final-state leptons from charged-current interactions
Cross-sections not well measured – need to use near detector to normalize.



Neutrino-Nucleon Interactions

The basic neutrino-nucleon interaction model is the dipole form factor model

• For CCQE interactions, this depends on a single physical parameter, the axial mass $M_{_{\rm A}}$.

CCQE interactions are particularly useful as the energy depends only on the outgoing lepton kinematics p_{μ} , θ_{μ}

• This is the main signal for T2K

$$E_{\text{reco}} = \frac{m_p^2 - (m_n - E_b)^2 - m_\mu^2 + 2(m_n - E_b)E_\mu}{2(m_n - E_b - E_\mu + p_\mu \cos\theta_\mu)}$$

Nuclear Effects

Dipole form factor model is sufficient for interactions with nucleons – not nuclear targets. In a nucleus there is binding energy, and Fermi motion of nucleons (no longer at rest). A simplistic nuclear model is used:

Relativistic Fermi Gas

Simple model of nuclear effects for CCQE interactions

 Nucleus is modeled as a Fermi gas of non-interacting neutrons and protons

Uses two nucleus-dependent parameters

- E_B: the nucleon binding energy
- p_F: the Fermi momentum
- Different for each nucleus

Random Phase Approximation

Correction to the RFG model

• Includes first-order nucleon-nucleon correlations not found in the RFG model

Models long range correlations between nucleons at low energies

Not strongly nucleus dependent

Modelling CCQE interactions

Axial form factor for nucleon-neutrino interactions modelled by dipole

parametrization:

$$F_A(Q^2) = g_A \left[1 + \frac{Q^2}{M_A^2} \right]^{-2}$$

Measurements of the axial mass from different experiments are all over the map!

World average: $1.012 \pm 0.031 \pm 0.060 \text{ GeV}$

Deuterium experiments: $0.99 \pm 0.04 \text{ GeV}$

MiniBooNE (carbon): 1.35 ± 0.17 GeV

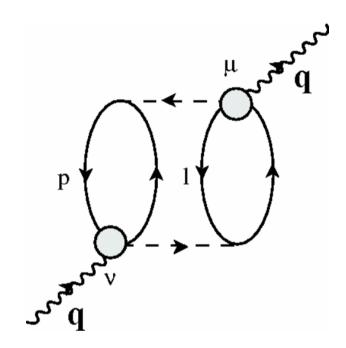
K2K (carbon): 1.20 ± 0.12 GeV

It looks like other effects confound the free nucleon form factor.

Multi-nucleon effects

Other multinucleon correlations are not covered in the RFG + RPA model

- These interactions can produce multiple protons or neutrons in the final state – difficult to identify separately from CCQE
- Irreducible experimental background

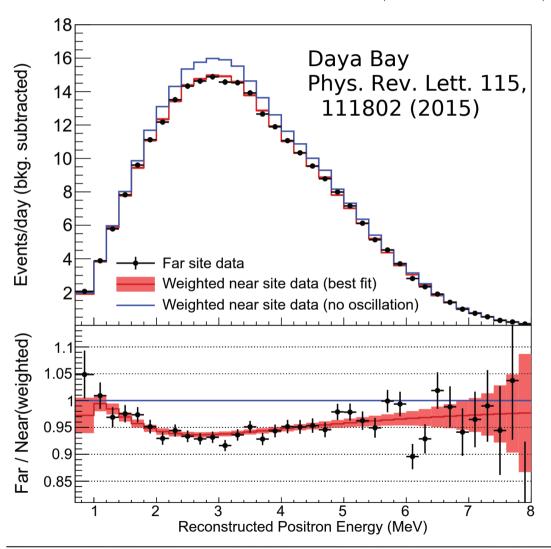


T2K models 2p-2h interactions, where two particle – hole pairs are propagated through the nucleus

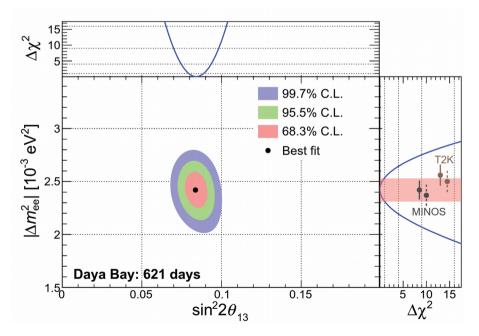
- Can produce two nucleons in the final state
- If not modeled, can have a significant effect on axial mass measurement: originally introduced to solve tensions between the axial mass measured with MiniBooNE and global averages
- T2K uses the Nieves 2p-2h model in the NEUT generator

Reactor neutrinos &
$$\theta_{13}$$

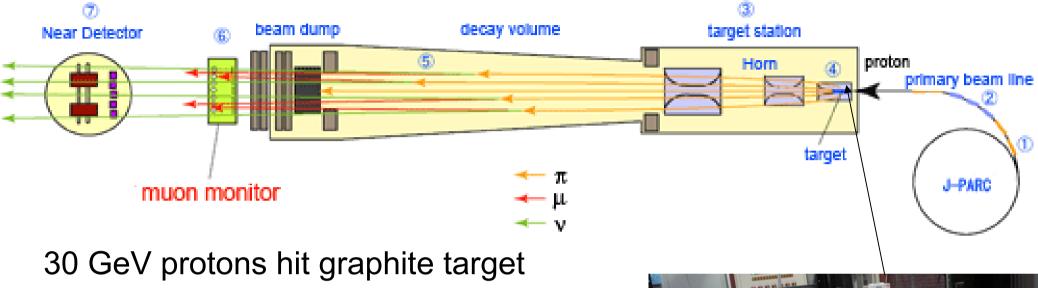
 $P(\bar{v}_e \to \bar{v}_e) = 1 - \sin^2 2\theta_{13} \sin^2 \left(\frac{1.27 \Delta m_{31}^2 L}{E} \right) - \cos^4 \theta_{13} \sin^2 2\theta_{12} \sin^2 \left(\frac{1.27 \Delta m_{21}^2 L}{E} \right)$



Daya Bay, RENO, Double CHOOZ look for disappearance of reactor neutrinos at ~1km baseline.



How To Make A Neutrino Beam



3 magnetic horns focus π^+ , defocus π^- .

 $\pi^{\scriptscriptstyle +} \rightarrow \mu^{\scriptscriptstyle +} + \nu_{_{\mu}}$ in 110m long decay pipe

 μ monitor at far end of beam dump: fluence: $10^8 \,\mu/\text{cm}^2/\text{spill}$ at full power



T2K's 90cm graphite target

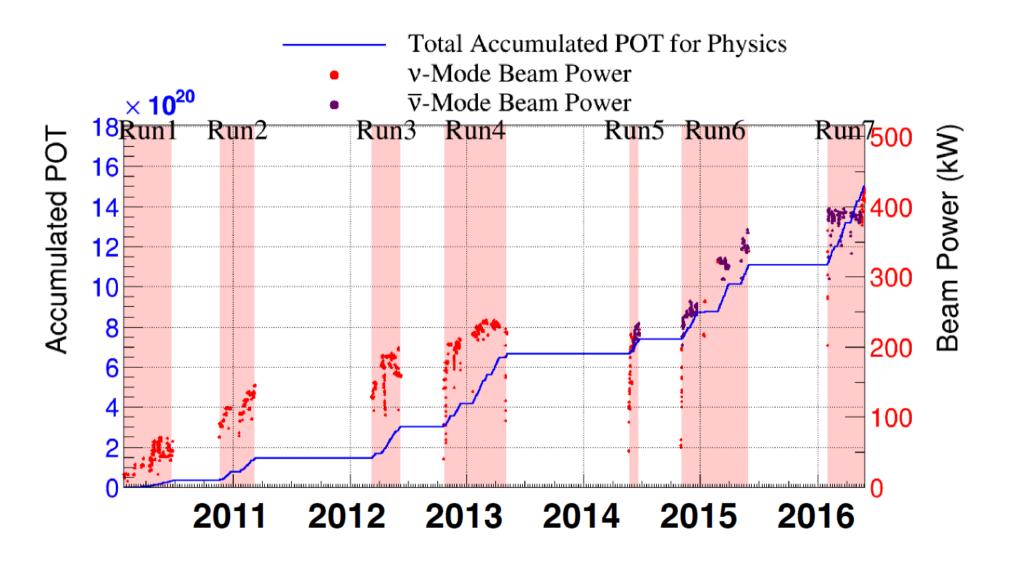




Inside the decay volume

 $\leftarrow The \ 2^{nd} \ focusing \\ horn$

T2K data collection



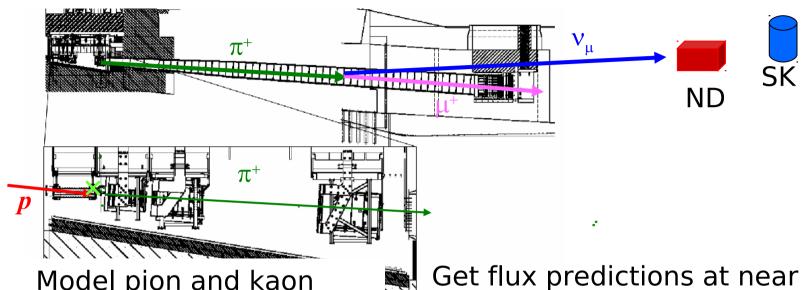
27 May 2016

POT total: 1.510×10²¹

ν-mode POT: 7.57×10²⁰ (50.14%)

v̄-mode POT: 7.53×10²⁰ (49.86%)

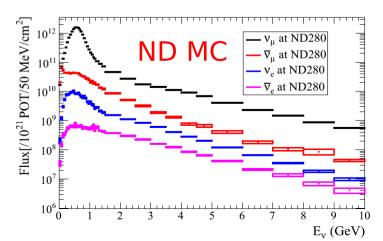
T2K: Flux prediction (Beam MC)

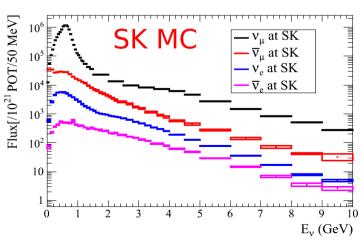


Simulate hadron production on target using FLUKA simulation

Model pion and kaon propagation and decay through horns and beamline

Particle production cross sections tuned to external data from NA61 and others.





detector and SK

Backgrounds to v_e Appearance

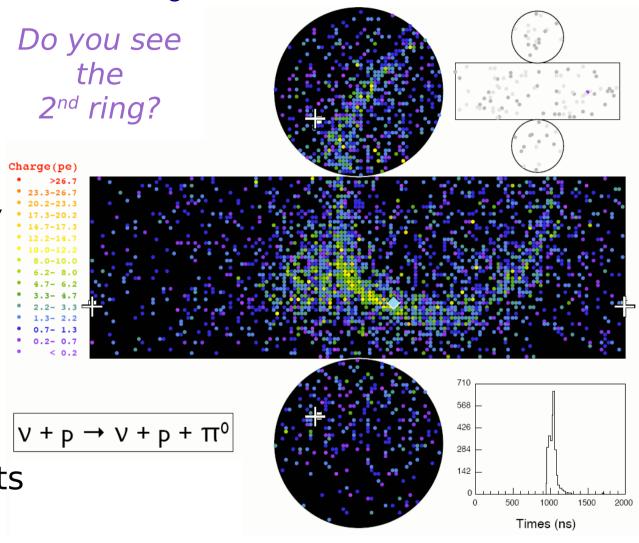
Intrinsic beam v_e :

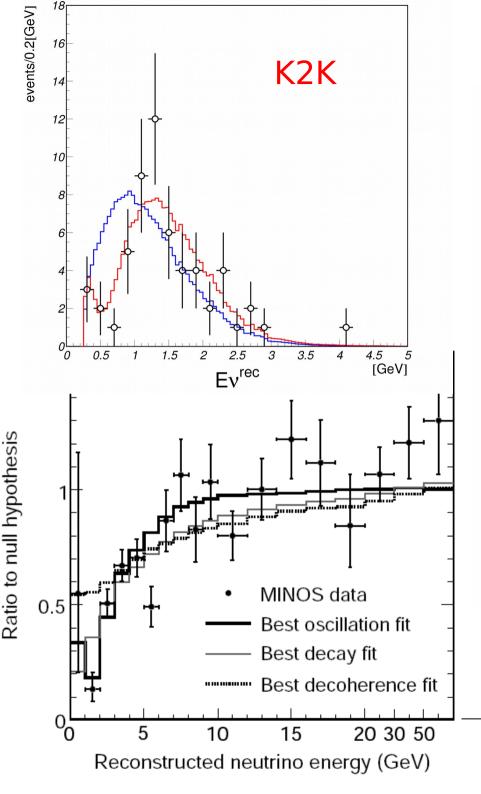
- reduce with E cut
- measure at ND

 π^0 production, with one γ from event not detected at Super-K:

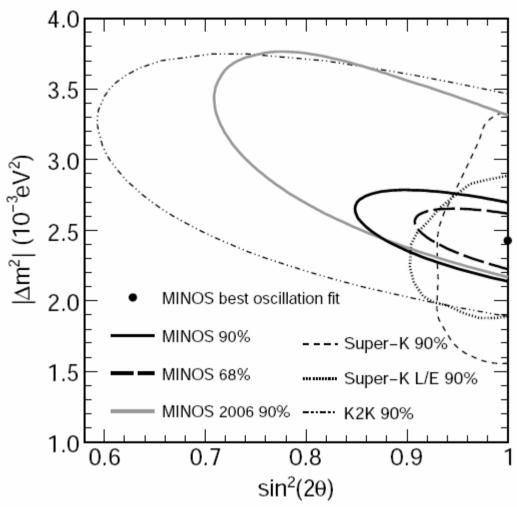
- better ID algorithms
- measure at ND
- measure π^0 in SK

The intrinsic beam events are a more significant background.



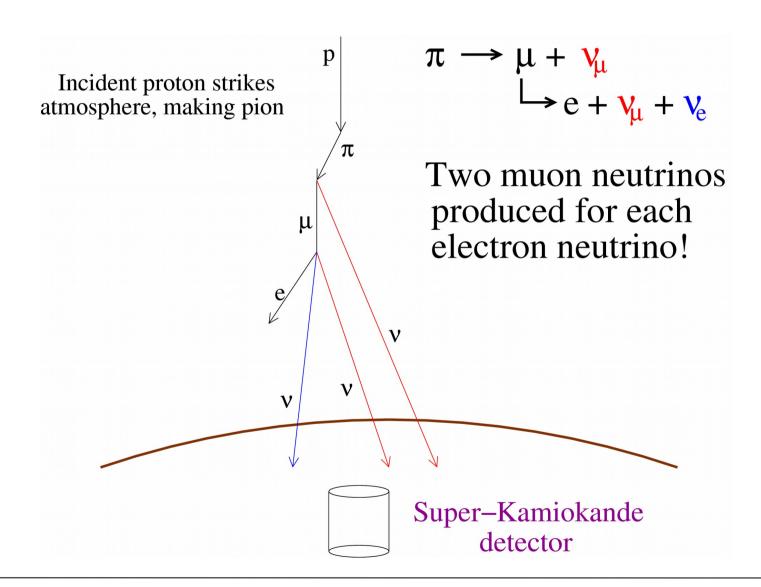


K2K & MINOS

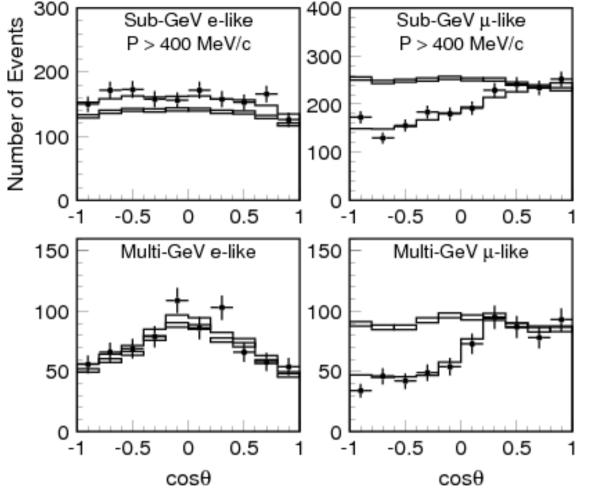


Consistency between atmospheric and long-baseline v oscillation results.

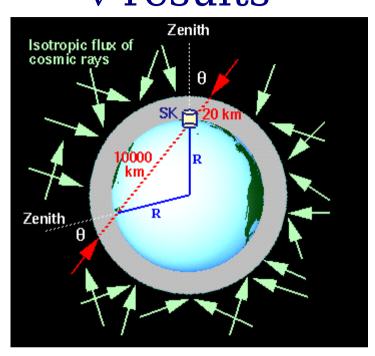
Atmospheric Neutrinos



Super-K atmospheric v results



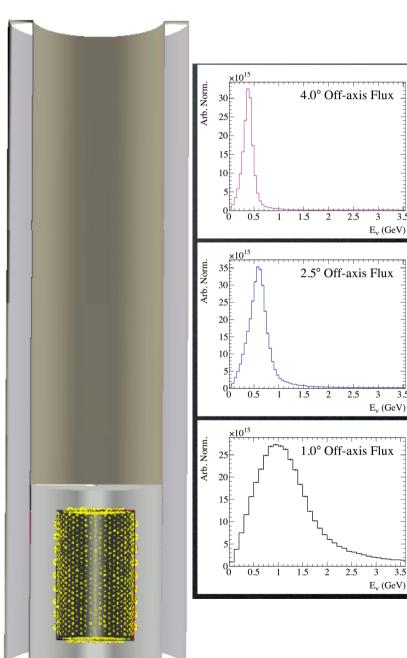
PRL 93:101801, 2004 PRD 71:112005, 2005



No deficit for v_e .

Seems like $\nu_{\mu} \rightarrow \nu_{\tau}$

NuPRISM concept



New concept to exploit the variation in neutrino energy with off-axis angle: tall water Cherenkov near detector spanning range of off-axis angles.

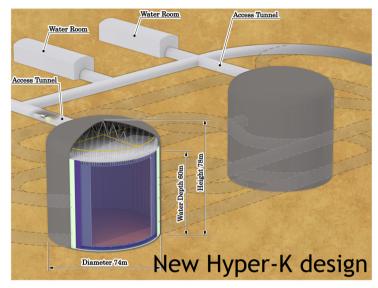
Data taken at different angles can directly predict neutrino interactions with arbitrary neutrino fluxes, including effects from oscillation.

Main proposal to Canadian government under review.

Hyper-K

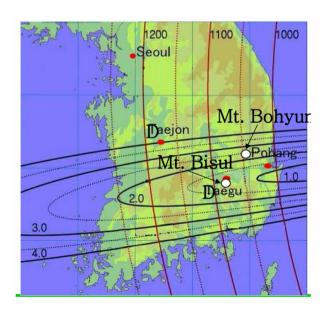
New technical design report: 187kton fiducial volume tank with enhanced photosensor coverage.

Proposal to locate a second detector in Korea at the second oscillation maximum.

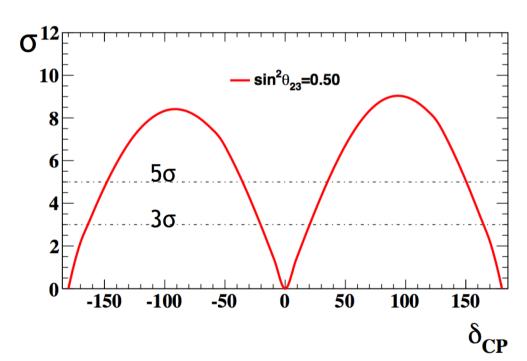


arXiv: 1502.05199

High sensitivity for δ_{CP}

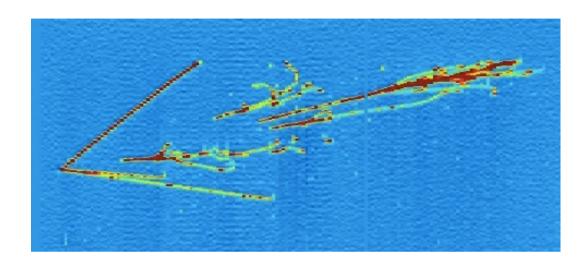


arXiv: 1611.06118



DUNE





Very long baseline experiment aiming at CP violation using liquid argon TPCs as the detector technology.

Higher energy and baseline than Hyper-K.