

# Charged Higgs Bosons in different Left-Right Models at the LHC

Poulose Poulose

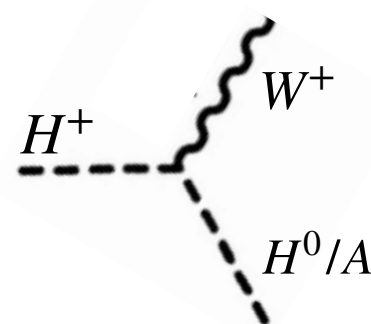
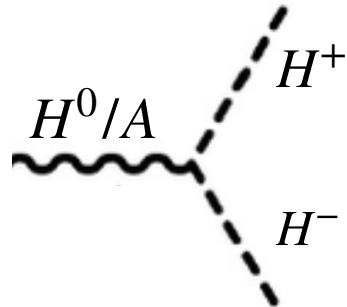
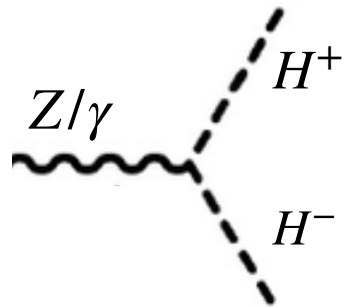
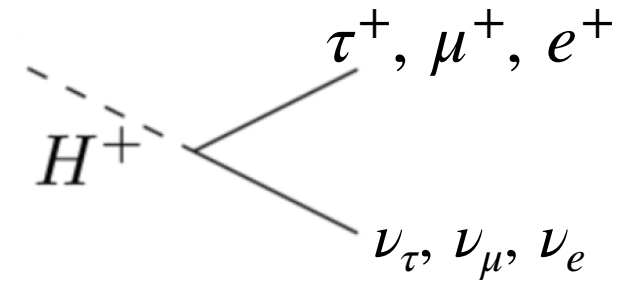
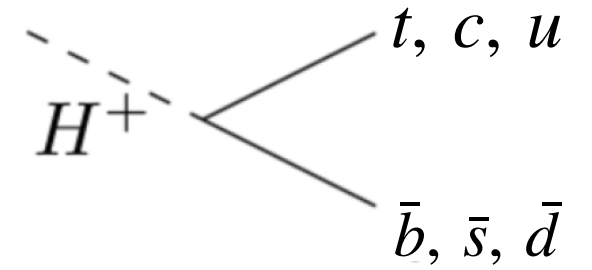
IIT Guwahati, India

# Charged Higgs in BSM scenarios

## Multi-Higgs Extensions:

MSSM, 2HDM, 3HDM, LRSM, Alternate Left-Right Models, Triplet-Higgs extensions, Singlet Charged Higgs extensions.

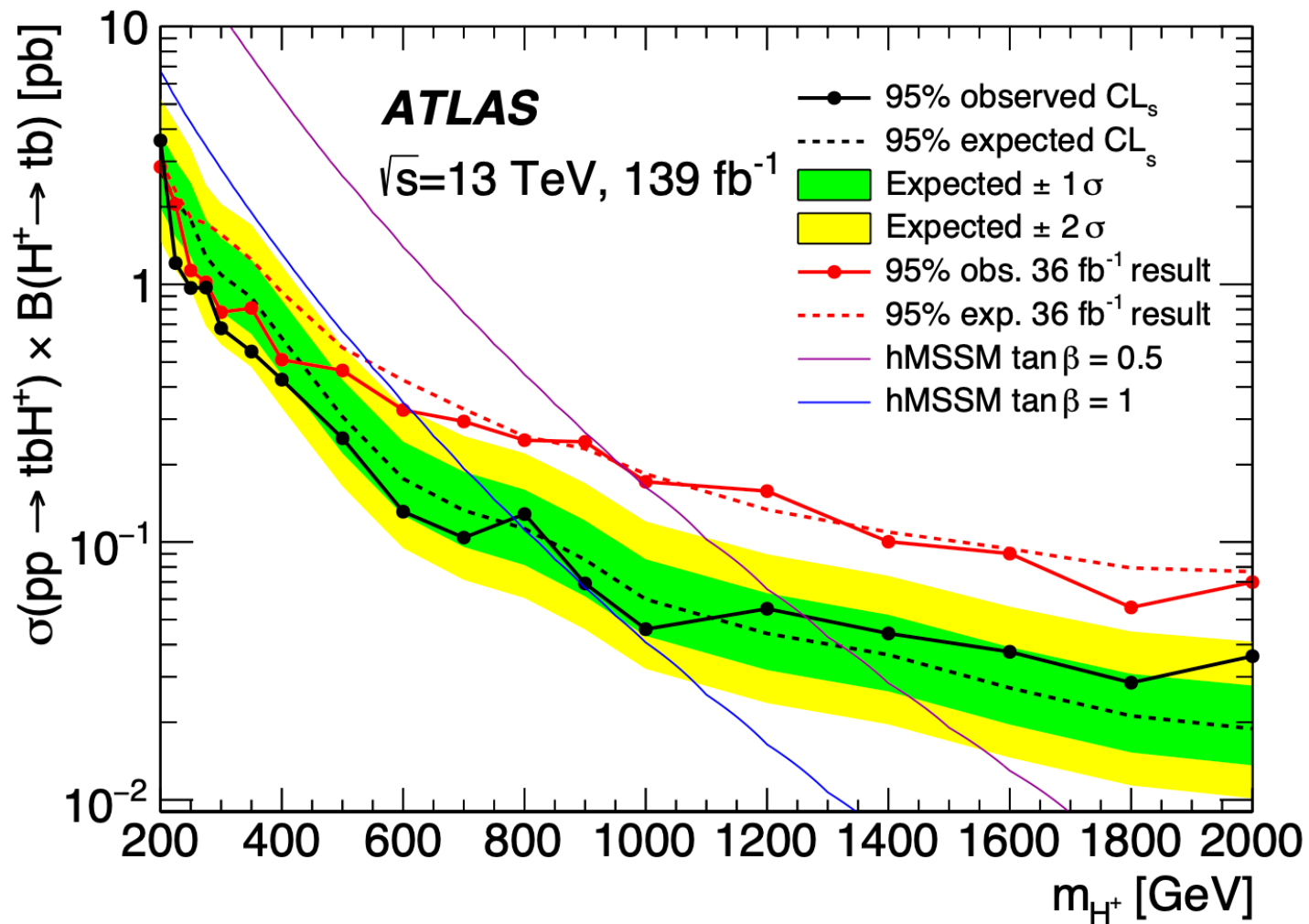
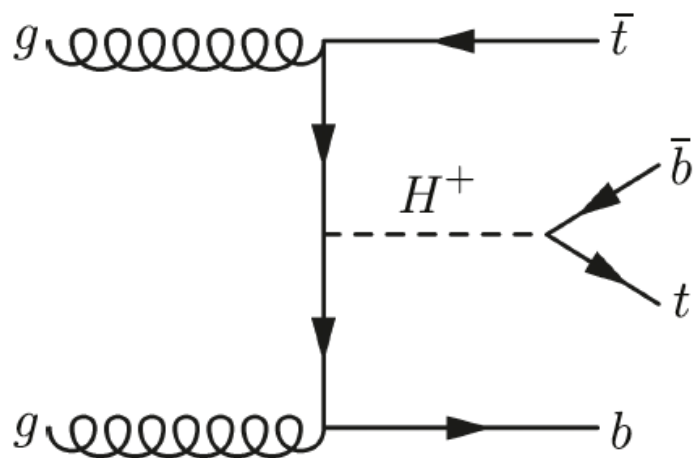
Non-singlet gauge multiplets necessarily have charged current interactions. This has become a strategy to search for such charged Higgs bosons at the colliders.



# Charged Higgs in BSM scenarios

For  $M_{H^+} > m_t$ ,

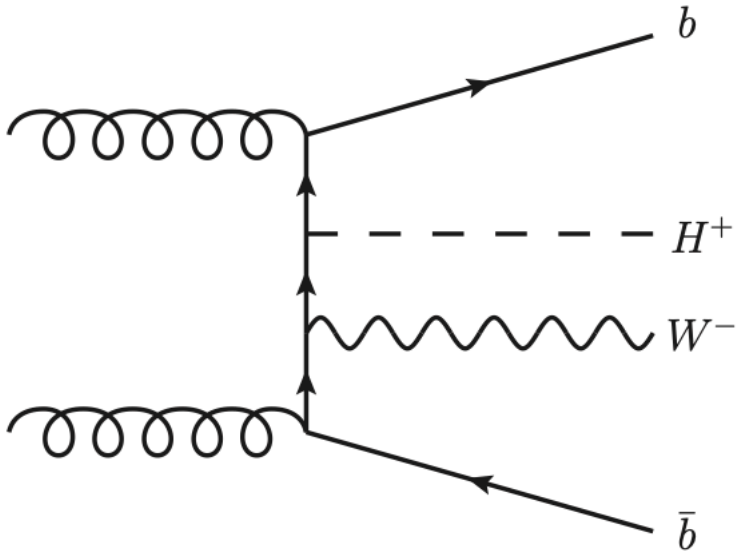
dominant decay channel is  $H^+ \rightarrow t\bar{b}$



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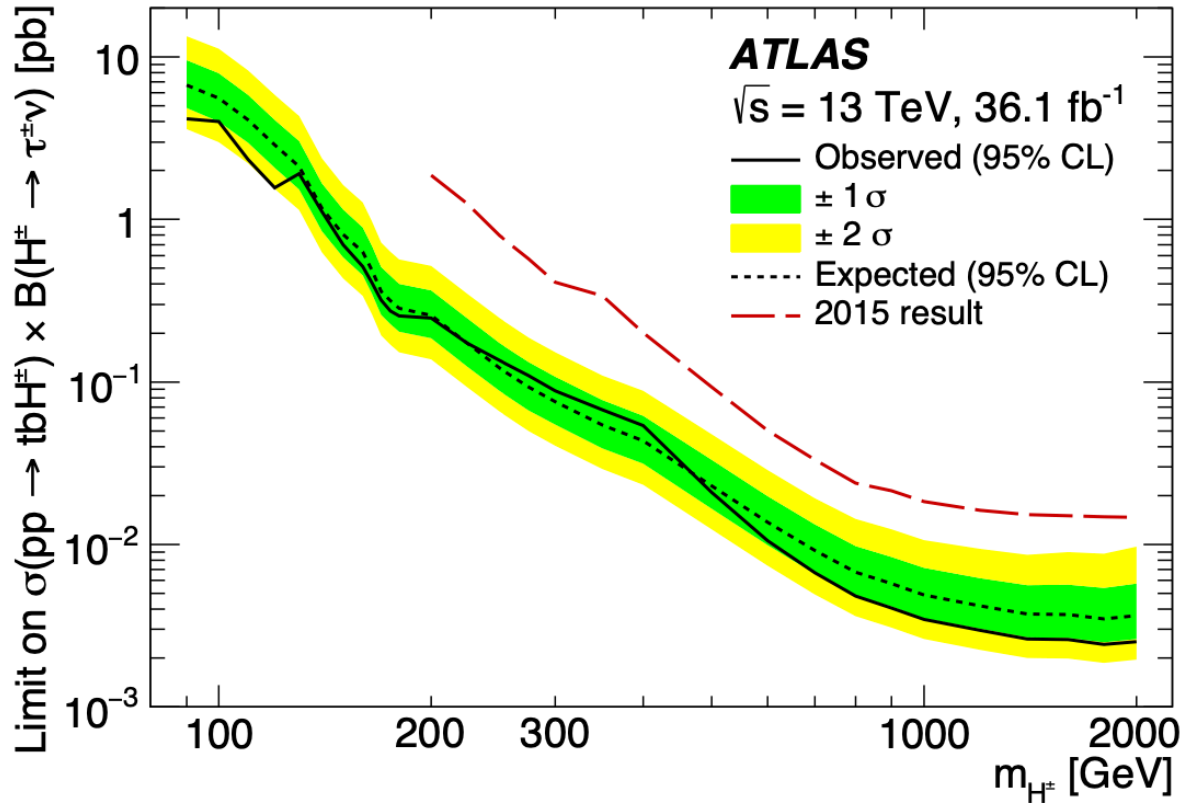
# Charged Higgs in BSM scenarios

Another influential interaction for non-singlet gauge multiplets is the interaction with  $W^\pm$  and neutral Higgs boson.



$$pp \rightarrow H^+ W^- b \bar{b}$$

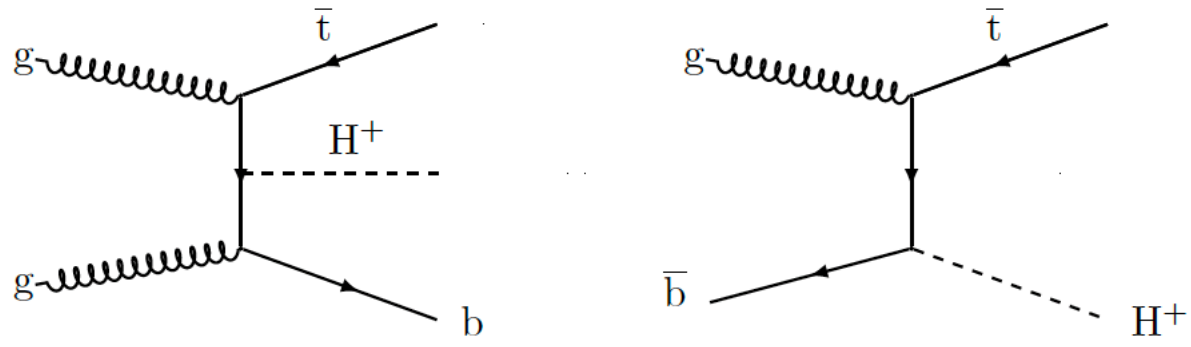
$$H^+ \rightarrow \tau^+ \nu_\tau$$



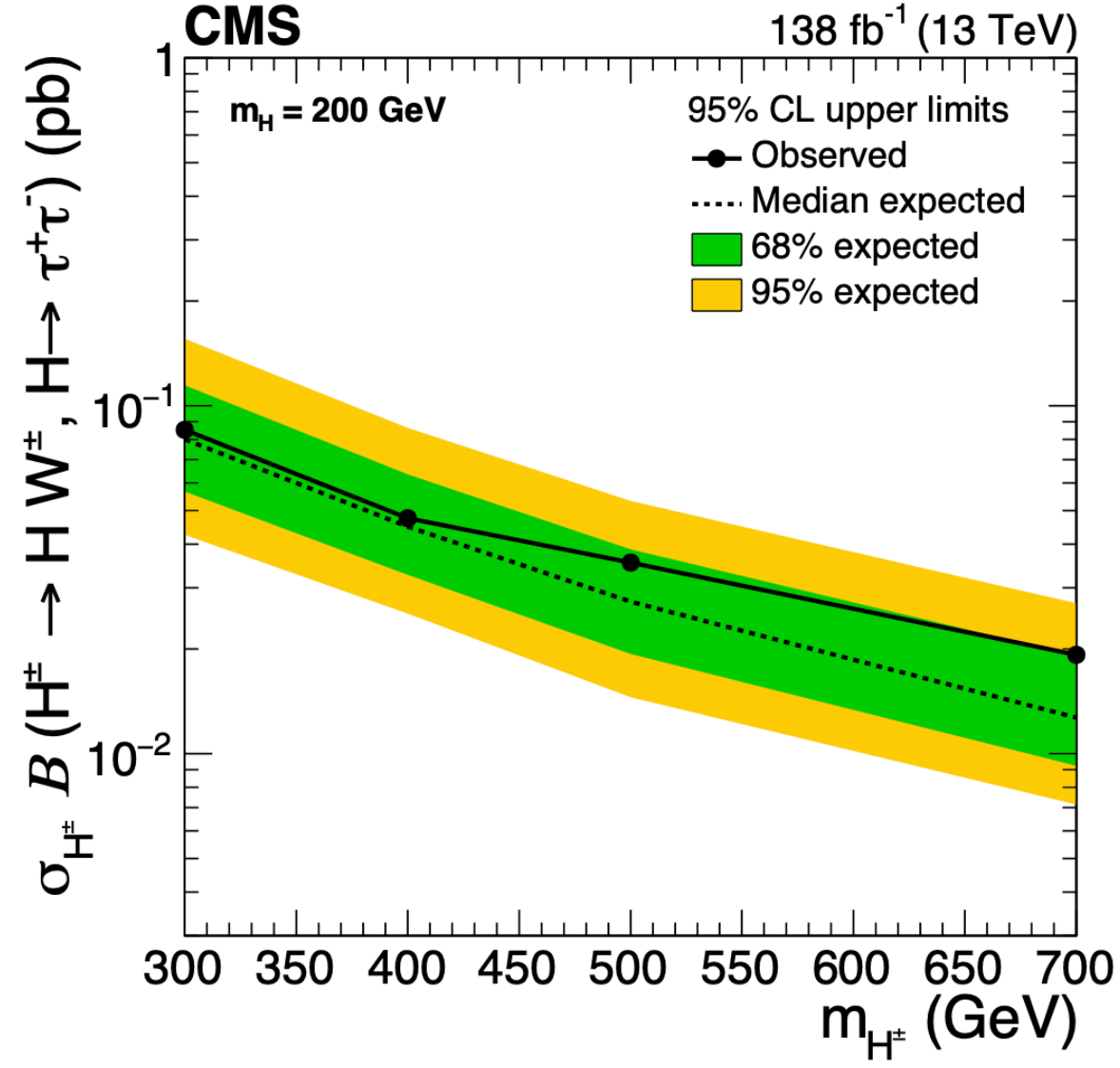
ATLAS: [JHEP 09 \(2018\) 139](#)

# Charged Higgs in BSM scenarios

Another influential interaction for non-singlet gauge multiplets is the interaction with  $W^\pm$  and neutral Higgs boson.



$$H^+ \rightarrow HW^+, \quad H \rightarrow \bar{\tau}\tau$$



JHEP09(2023)032

# Charged Higgs in BSM scenarios

Lighter ones decay to lighter quarks

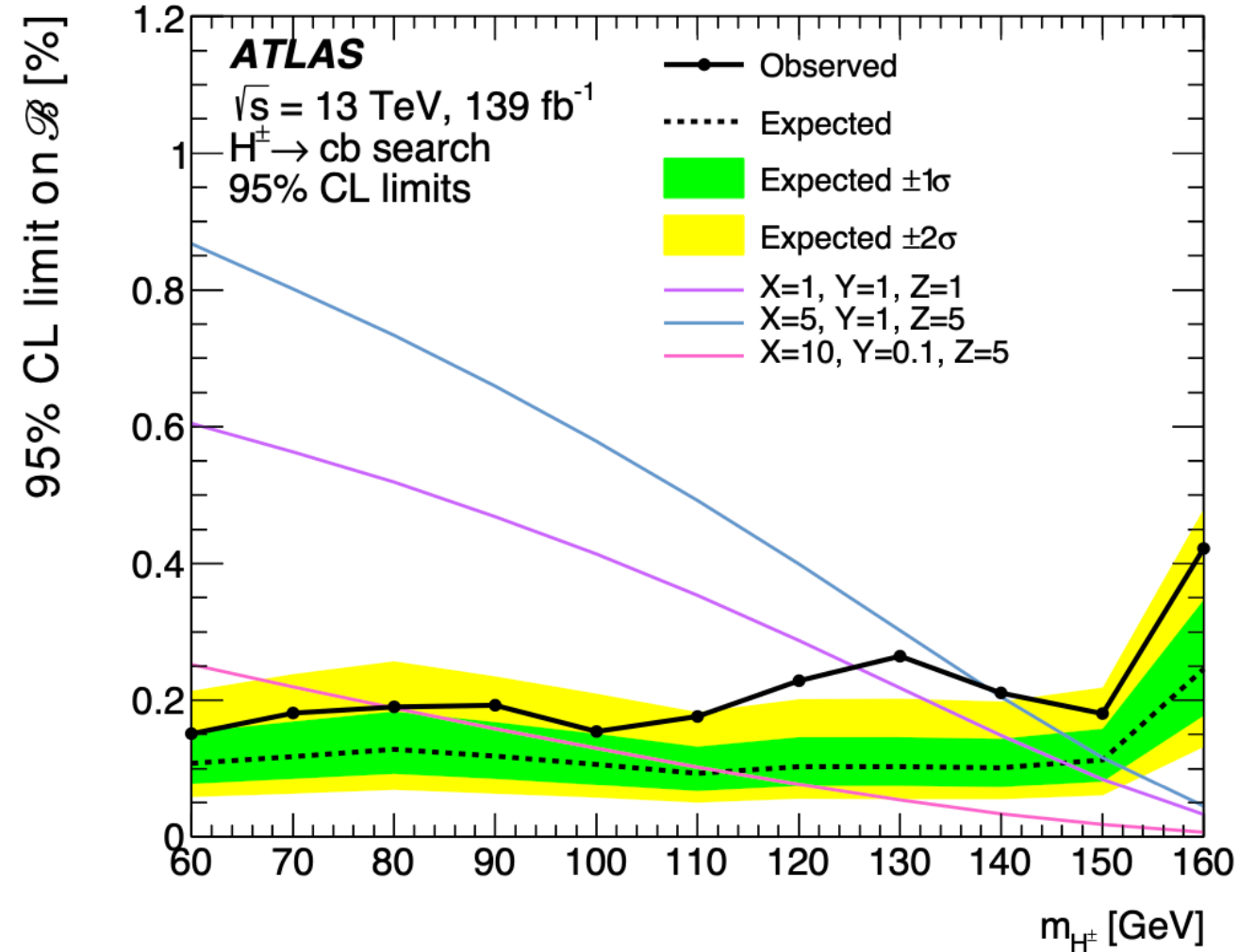
$$pp \rightarrow t\bar{t}$$

$$t \rightarrow H^+ b$$

$$t \rightarrow W^+ b$$

$$H^+ \rightarrow c\bar{b}$$

$$W^+ \rightarrow \ell^+ \nu$$



$$\mathcal{B} = BR(t \rightarrow H^+ b) BR(H^+ \rightarrow c\bar{b})$$

# Left-Right Symmetric Models

SM is chiral, the left-handed and right-handed particles differ in their fundamental interaction. Parity is violated maximally.

Left-Right Symmetric models are attempts to understand the origin of the parity violation. We can start with a Left-Right symmetric model at higher energies:

$$SU(3)_C \times SU(2)_L \times SU(2)_R \times U(1)$$

# Left-Right Symmetric Models

SM is chiral, the left-handed and right-handed particles differ in their fundamental interaction. Parity is violated maximally.

Left-Right Symmetric models are attempts to understand the origin of the parity violation. We can start with a Left-Right symmetric model at higher energies:

$$SU(3)_C \times SU(2)_L \times \underbrace{SU(2)_R \times U(1)}_{U(1)_Y}$$

This is broken to the SM gauge symmetry

$$U(1)_Y$$

$$\implies SU(3)_C \times SU(2)_L \times U(1)_Y$$



The Standard LRSM

Needs additional scalar to break the LR symmetry.

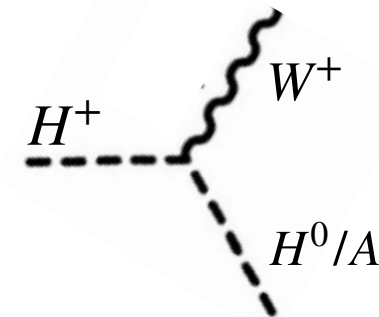
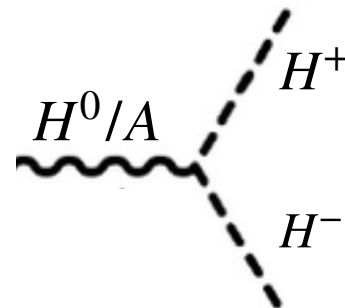
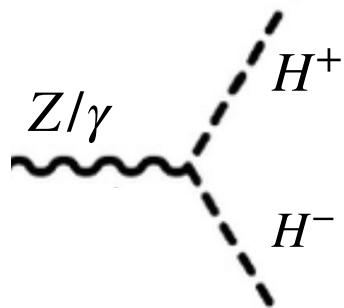
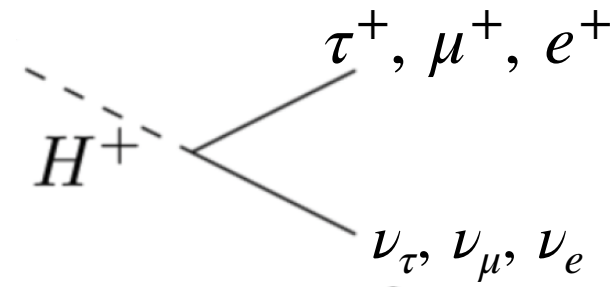
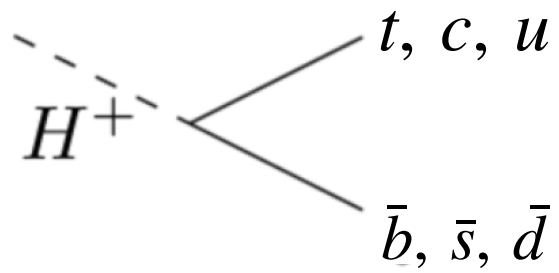
Leads to charged Higgs

	Particles	$SU(3)_C$	$SU(2)_L$	$SU(2)_R$	$U(1)_{B-L}$
Quarks	$Q_L = \begin{pmatrix} u_L \\ d_L \end{pmatrix}$	3	2	1	$\frac{1}{3}$
	$Q_R = \begin{pmatrix} u_R \\ d_R \end{pmatrix}$	3	1	2	$\frac{1}{3}$
Leptons	$L_L = \begin{pmatrix} \nu_L \\ e_L \end{pmatrix}$	1	2	1	-1
	$L_R = \begin{pmatrix} \nu_R \\ e_R \end{pmatrix}$	1	1	2	-1
Scalars	$\Phi = \begin{pmatrix} \phi_1^0 & \phi_1^+ \\ \phi_2^- & \phi_2^0 \end{pmatrix}$	1	2	2*	0
	$\Delta_L = \begin{pmatrix} \frac{\delta_L^+}{\sqrt{2}} & \delta_L^{++} \\ \delta_L^0 & -\frac{\delta_L^+}{\sqrt{2}} \end{pmatrix}$	1	3	1	2
	$\Delta_R = \begin{pmatrix} \frac{\delta_R^+}{\sqrt{2}} & \delta_R^{++} \\ \delta_R^0 & -\frac{\delta_R^+}{\sqrt{2}} \end{pmatrix}$	1	1	3	2

TABLE I: The particle content of LRSM.

# LRSM charged Higgs

The LHC search results are applicable



## Another version of Left-Right Symmetric model

# LRSM

Singlet

# Alternate LR Model

	Particles	$SU(3)_C$	$SU(2)_L$	$SU(2)_R$	$U(1)_{B-L}$
Quarks	$Q_L = \begin{pmatrix} u_L \\ d_L \end{pmatrix}$	3	2	1	$\frac{1}{3}$
	$Q_R = \begin{pmatrix} u_R \\ d_R \end{pmatrix}$	3	1	2	$\frac{1}{3}$
Leptons	$L_L = \begin{pmatrix} \nu_L \\ e_L \end{pmatrix}$	1	2	1	-1
	$L_R = \begin{pmatrix} \nu_R \\ e_R \end{pmatrix}$	1	1	2	-1
Scalars	$\Phi = \begin{pmatrix} \phi_1^0 & \phi_1^+ \\ \phi_2^- & \phi_2^0 \end{pmatrix}$	1	2	2*	0
	$\Delta_L = \begin{pmatrix} \frac{\delta_L^+}{\sqrt{2}} & \delta_L^{++} \\ \delta_L^0 & -\frac{\delta_L^+}{\sqrt{2}} \end{pmatrix}$	1	3	1	2
	$\Delta_R = \begin{pmatrix} \frac{\delta_R^+}{\sqrt{2}} & \delta_R^{++} \\ \delta_R^0 & -\frac{\delta_R^+}{\sqrt{2}} \end{pmatrix}$	1	1	3	2

TABLE I: The particle content of LRSM.

	Particles	$SU(3)_C$	$SU(2)_L$	$SU(2)'_R$	$U(1)_{B-L}$	S
Quarks	$Q_L = \begin{pmatrix} u_L \\ d_L \end{pmatrix}$	3	2	1	$\frac{1}{6}$	0
	$Q_R = \begin{pmatrix} u_R \\ d'_R \end{pmatrix}$	3	1	2	$\frac{1}{6}$	$-\frac{1}{2}$
	$d'_L$	3	1	1	$-\frac{1}{3}$	-1
	$d_R$	3	1	1	$-\frac{1}{3}$	0
Leptons	$L_L = \begin{pmatrix} \nu_L \\ e_L \end{pmatrix}$	1	2	1	$-\frac{1}{2}$	1
	$L_R = \begin{pmatrix} n_R \\ e_R \end{pmatrix}$	1	1	2	$-\frac{1}{2}$	$\frac{3}{2}$
	$n_L$	1	1	1	0	2
	$\nu_R$	1	1	1	0	1
Scalars	$\Phi = \begin{pmatrix} \phi_1^0 & \phi_1^+ \\ \phi_2^- & \phi_2^0 \end{pmatrix}$	1	2	2*	0	$-\frac{1}{2}$
	$\chi_L = \begin{pmatrix} \chi_L^+ \\ \chi_L^0 \end{pmatrix}$	1	2	1	$\frac{1}{2}$	0
	$\chi_R = \begin{pmatrix} \chi_R^+ \\ \chi_R^0 \end{pmatrix}$	1	1	2	$\frac{1}{2}$	$\frac{1}{2}$

TABLE I: The particle content of ALRM.

E. Ma, Phys. Rev. D36, 274 (1987).

Particles		$\begin{pmatrix} u_R \\ d_R \end{pmatrix}$		$SU(2)_R$	$U(1)_{B-L}$
Quarks	$Q_L = \begin{pmatrix} u_L \\ d_L \end{pmatrix}$ $Q_R = \begin{pmatrix} u_R \\ d_R \end{pmatrix}$			1	$\frac{1}{3}$
Leptons	$L_L = \begin{pmatrix} \nu_L \\ e_L \end{pmatrix}$ $L_R = \begin{pmatrix} \nu_R \\ e_R \end{pmatrix}$	1	2	1	-1
Scalars	$\Phi = \begin{pmatrix} \phi_1^0 & \phi_1^+ \\ \phi_2^- & \phi_2^0 \end{pmatrix}$	$\begin{pmatrix} \nu_R \\ e_R \end{pmatrix}$		2*	0
	$\Delta_L = \begin{pmatrix} \frac{\delta_L^+}{\sqrt{2}} & \delta_L^{++} \\ \delta_L^0 & -\frac{\delta_L^+}{\sqrt{2}} \end{pmatrix}$			1	2
	$\Delta_R = \begin{pmatrix} \frac{\delta_R^+}{\sqrt{2}} & \delta_R^{++} \\ \delta_R^0 & -\frac{\delta_R^+}{\sqrt{2}} \end{pmatrix}$			1	1

TABLE I: The particle content of LRSM.

Particles		$\begin{pmatrix} u_R \\ d'_R \end{pmatrix}$		$SU(2)'_R$	$U(1)_{B-L}$	S		
Quarks	$Q_L = \begin{pmatrix} u_L \\ d_L \end{pmatrix}$ $Q_R = \begin{pmatrix} u_R \\ d'_R \end{pmatrix}$ $d'_L$ $d_R$			1	$\frac{1}{6}$	0	2	$\frac{1}{6}$
Leptons	$L_L = \begin{pmatrix} \nu_L \\ e_L \end{pmatrix}$ $L_R = \begin{pmatrix} n_R \\ e_R \end{pmatrix}$ $n_L$ $\nu_R$	$\begin{pmatrix} n_R \\ e_R \end{pmatrix}$		1	2	1	$\begin{matrix} d_R \\ d'_L \\ \nu_R \\ n_L \end{matrix}$	
	Scalars			$\Phi = \begin{pmatrix} \phi_1^0 & \phi_1^+ \\ \phi_2^- & \phi_2^0 \end{pmatrix}$ $\chi_L = \begin{pmatrix} \chi_L^+ \\ \chi_L^0 \end{pmatrix}$ $\chi_R = \begin{pmatrix} \chi_R^+ \\ \chi_R^0 \end{pmatrix}$	1	2		1
					1	1		2

TABLE I: The particle content of ALRM.

## Some interesting features

$\nu_R$  is a Majorana neutrino, as against being part of the couple in the standard LRSM

=> Possibility of Leptogenesis

$\nu_R$  along with presence of light ( $\sim 100$  GeV)  $H^+$  can impact the  $0\nu\beta\beta$  processes

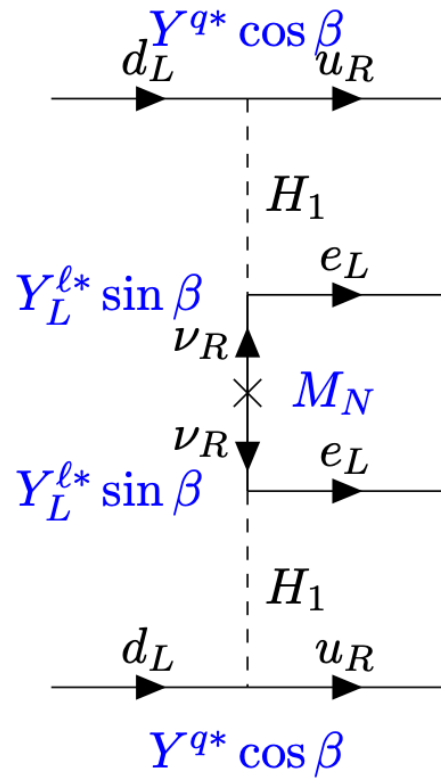
*Phys.Rev.D* 102 (2020) 7, 075020 arXiv: 2008.12270

$n / H_1^0$  : The lightest is a dark matter candidate.

JHEP 12, 032, arXiv:2211.04286

$\nu_R$  along with presence of light ( $\sim 100$  GeV)  $H^+$  can impact the  $0\nu\beta\beta$  processes

*Phys.Rev.D* 102 (2020) 7, 075020 arXiv: 2008.12270



$$\begin{aligned}
 -\mathcal{L}_Y = & \bar{Q}_L Y^q \tilde{\Phi} Q_R + \bar{Q}_L Y_L^q \chi_L d_R + \bar{Q}_R Y_R^q \chi_R d'_L + \bar{L}_L Y^l \Phi L_R \\
 & + \bar{L}_L Y_L^l \tilde{\chi}_L \nu_R + \bar{L}_R Y_R^l \tilde{\chi}_R n_L + \text{h.c.}
 \end{aligned}$$

For details of the model:

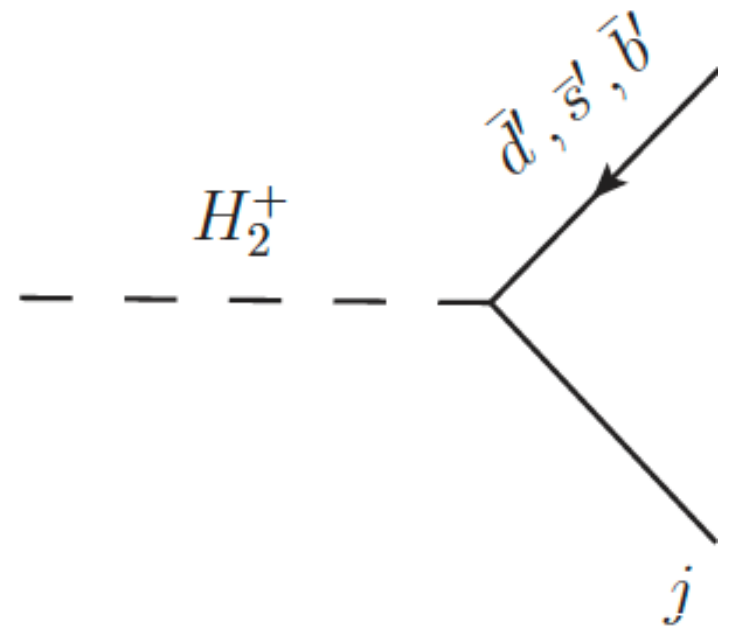
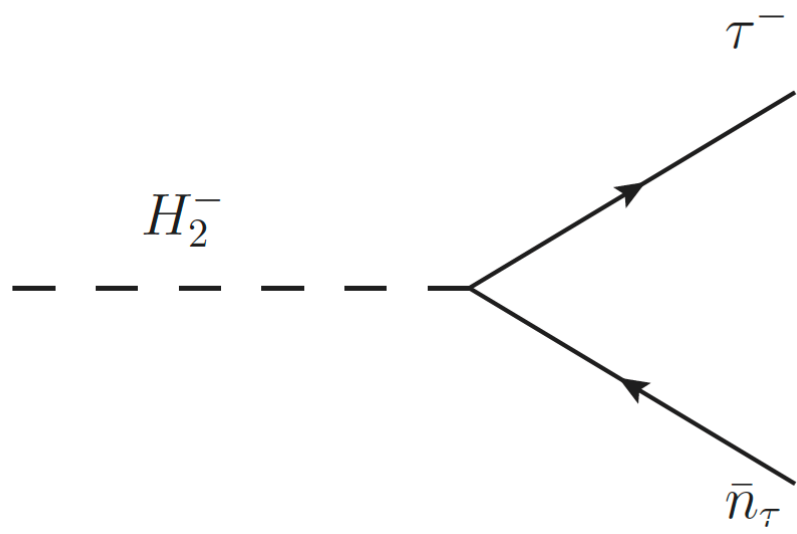
M. Frank, C. Majumdar, P. Poullose, S. Senapati, and U. A. Yajnik

JHEP 03, 065, arXiv:2111.08582 [hep-ph]

JHEP 12, 032, arXiv:2211.04286 [hep-ph]

*Phys.Rev.D* 102 (2020) 7, 075020 arXiv: 2008.12270

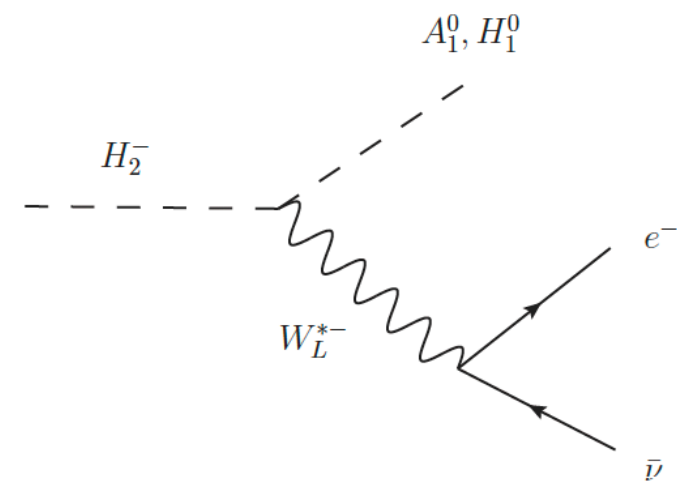
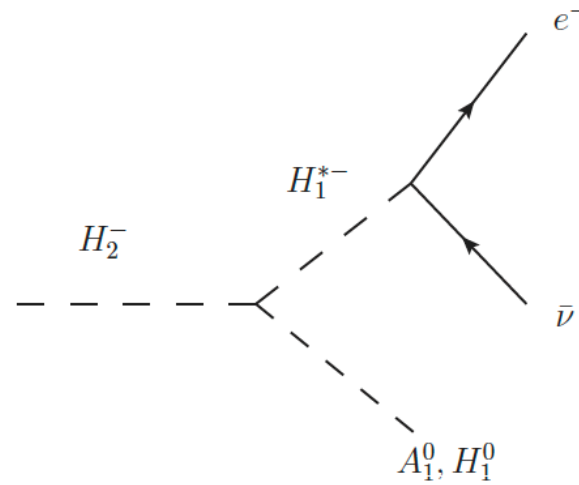
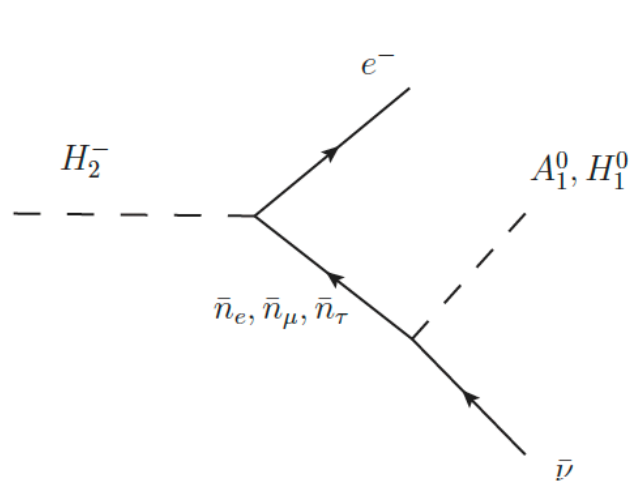




Non-standard decays, which are not searched for at the LHC

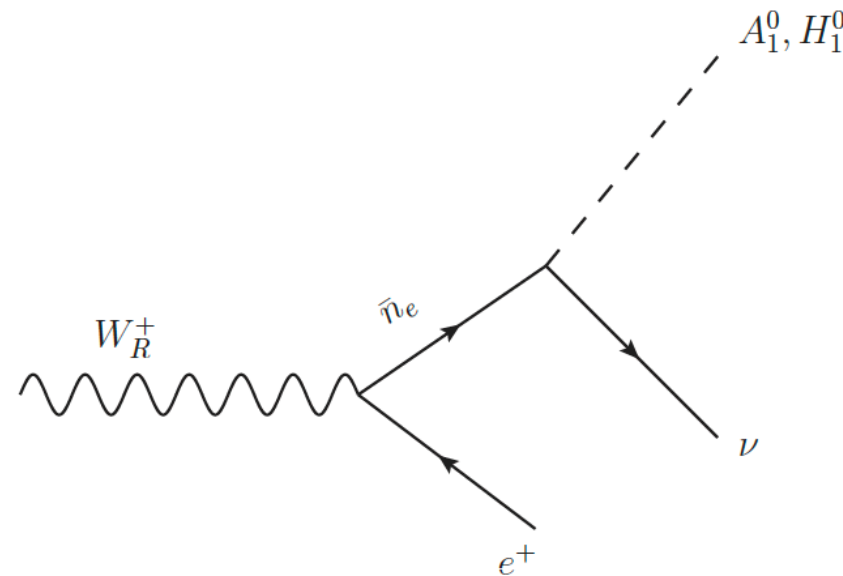
# Our study: Search for such an $H^+$ at the LHC and future colliders

$$pp \rightarrow H^+H^-, W_R W_R, W_R H^\pm$$



$$M_{H^+} < m_{n_\ell} < M_{W_R}$$

Completely different topology compared to the Standard LRSM (or 2HDM / MSSM) case



Work in progress With:  
M. Frank, C. Majumdar,  
S. Senapati, U. Yagnik

An example:

$$pp \rightarrow H^+H^-, W_R W_R, W_R H^\pm$$

$$M_{W_R} = 1146 \text{ GeV}$$

$$M_{H^+} = 187 \text{ GeV}$$

$$m_{d'} = 1.3 \text{ TeV}$$

$$m_n = 130 \text{ GeV}$$

$$BR(H^+ \rightarrow \tau^+ n_\tau) = 59.6 \%$$

$$BR(n \rightarrow H_1^0 \nu_e) = 25.35 \%$$

$$\sigma \times BR = 108.7 \text{ fb}$$

$$@ \sqrt{s} = 13 \text{ TeV}$$

Detailed detector level analyses is being done.

## Summary

Search strategy of  $H^+$  at the LHC presume their standard charged current interaction.

Considering a well motivated alternate Left-Right Model,  $H^+$  can have non-standard charged current interactions involving exotic fermions (quarks and neutral leptons).

Collider direct searches of  $H^+$  should include such possibilities.

Indirect constraints coming from flavour sector can also be different compared to the standard LRSM.