

# First Principles Calculations of $p + {}^7\text{Li}$ Radiative Capture (and the X17 anomaly)

**Peter Gysbers**

P. Navrátil, C. Hebborn, G. Hupin,  
K. Kravvaris, S. Quaglioni



CAP Congress: June 6, 2022



# Radiative Capture $A + B \rightarrow C + \gamma$

- Notation:  $B(A, \gamma)C$
- A nuclear reaction that often occurs in astrophysics:
  - Stellar burning:  $d(p, \gamma)^3\text{He}$ ,  $^3\text{He}(\alpha, \gamma)^7\text{Be}$ , ...
  - Big Bang Nucleosynthesis:  $d(p, \gamma)^3\text{He}$ ,  $^4\text{He}(d, \gamma)^6\text{Li}$ , ...
  - Search for new physics:  $^7\text{Li}(p, \gamma)^8\text{Be}$ ,  $^3\text{H}(p, \gamma)^4\text{He}$

More Notation:  
 $d = ^2\text{H}$   
 $\alpha = ^4\text{He}$



adapted from:  
Feng PRD 95, 035017 (2017)

## Calculating Radiative Capture

To calculate the rate of reaction (cross section) we need:

- ▶ initial wavefunction:  $|\Psi_i\rangle$  ( $A + B$ )
- ▶ final wavefunction:  $|\Psi_f\rangle$  ( $C$ )
- ▶ photon interaction (electromagnetic operator):  $\hat{O}_\gamma$

We need to calculate the square of the transition matrix elements:

$$\sigma \sim \sum_{if} |\langle \Psi_f | \hat{O}_\gamma | \Psi_i \rangle|^2$$

**Bound States:**  $|\Psi_f\rangle = \left|J_f^{\pi_f} T_f\right\rangle$

Eigenstate of the nuclear Hamiltonian:

$$H^A |\Psi_k\rangle = E_k |\Psi_k\rangle, \text{ where } H^A = \sum_i^A T_i + \sum_{i < j} V_{ij}^{NN} + \sum_{i < j < f} V_{ijf}^{3N}$$

**Bound States:**  $|\Psi_f\rangle = \left|J_f^{\pi_f} T_f\right\rangle$

Eigenstate of the nuclear Hamiltonian:

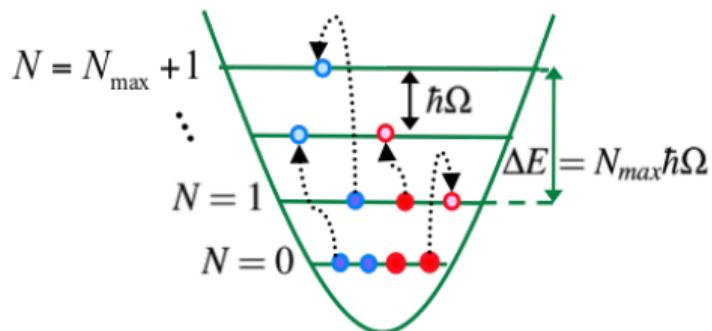
$$H^A |\Psi_k\rangle = E_k |\Psi_k\rangle, \text{ where } H^A = \sum_i^A T_i + \sum_{i < j} V_{ij}^{NN} + \sum_{i < j < f} V_{ijf}^{3N}$$

### The No-Core Shell Model (NCSM)

Expand in anti-symmetrized products of harmonic oscillator single-particle states:

$$|\Psi_k\rangle = \sum_{N=0}^{N_{max}} \sum_j c_{Nj}^k |\Phi_{Nj}\rangle$$

Convergence to exact as  $N_{max} \rightarrow \infty$



**Unbound (Continuum) States:**  $|\Psi_i\rangle = \left\langle [|\Psi_A\rangle |\Psi_B\rangle \psi(\vec{r}_A - \vec{r}_B)]^{(J_i^{\pi_i} T_i)} \right\rangle$

- ▶ The incoming state is made of distinct clusters with relative motion
- ▶ Harmonic oscillator states cannot describe the tail of the wavefunction (long-range physics)
- ▶ A method beyond the NCSM is needed for scattering and reactions

**Unbound (Continuum) States:**  $|\Psi_i\rangle = \left\langle \left[ |\Psi_A\rangle |\Psi_B\rangle \psi(\vec{r}_A - \vec{r}_B) \right]^{(J_i^{\pi_i} T_i)} \right\rangle$

- The incoming state is made of distinct clusters with relative motion
- Harmonic oscillator states cannot describe the tail of the wavefunction (long-range physics)
- A method beyond the NCSM is needed for scattering and reactions

### No-Core Shell Model with Continuum (NCSMC)

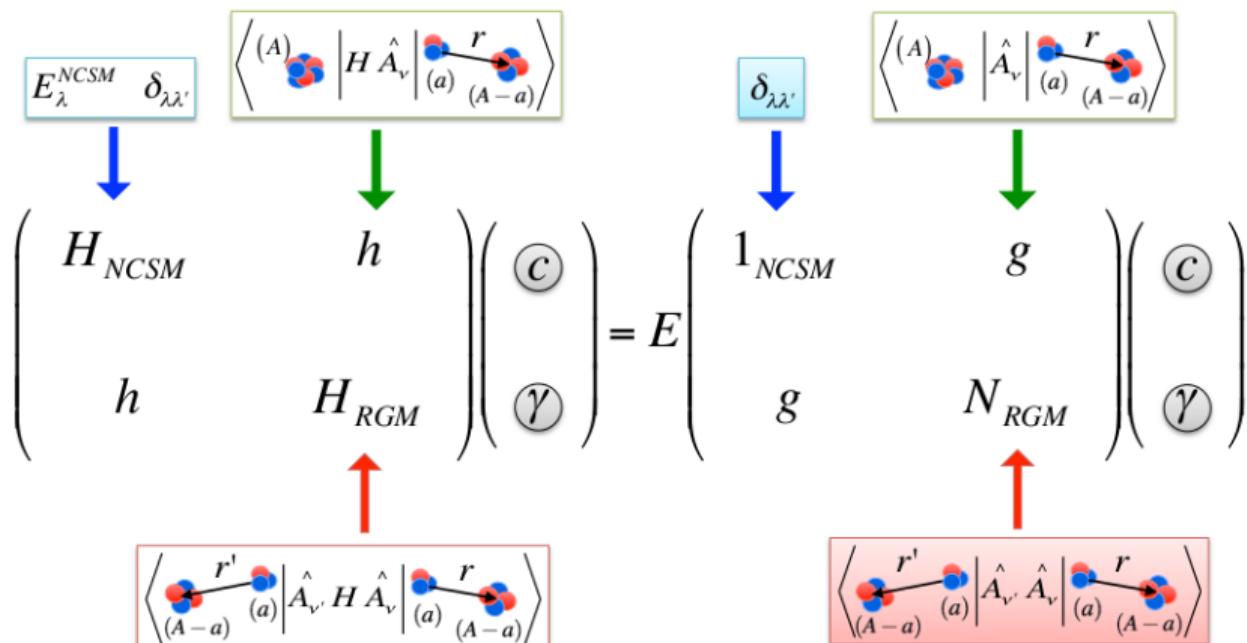
- Solution: extend the NCSM basis!

$$\Psi^{(A)} = \sum_{\lambda} c_{\lambda} \left| {}^{(A)} \text{cluster}, \lambda \right\rangle + \sum_{\nu} \int d\vec{r} \gamma_{\nu}(\vec{r}) \hat{A}_{\nu} \left| {}_{(A-a)}^{\vec{r}}, \nu \right\rangle$$

# NCSMC Equations

$$H \Psi^{(A)} = E \Psi^{(A)}$$

$$\Psi^{(A)} = \sum_{\lambda} c_{\lambda} \left| {}^{(A)} \text{H}_\lambda, \lambda \right\rangle + \sum_v \int d\vec{r} \gamma_v(\vec{r}) \hat{A}_v \left| {}^{(A-a)} \text{H}_{(a)}, v \right\rangle$$



## NCSMC for $p + {}^7\text{Li}$ ( ${}^8\text{Be}$ )

$$\Psi_{\text{NCSMC}}^{(8)} = \sum_{\lambda} c_{\lambda} |{}^8\text{Be}, \lambda\rangle + \sum_{\nu} \int dr \gamma_{\nu}(r) \hat{A}_{\nu} |{}^7\text{Li} + p, \nu\rangle + \sum_{\mu} \int dr \gamma_{\mu}(r) \hat{A}_{\mu} |{}^7\text{Be} + n, \mu\rangle$$

### Process:

- ▶ Solve NCSM for each constituent nucleus:  ${}^8\text{Be}$ ,  ${}^7\text{Li}$  and  ${}^7\text{Be}$ 
  - ▶ 30 eigenstates from  ${}^8\text{Be}$
  - ▶ 5 eigenstates each from  ${}^7\text{Li}$  and  ${}^7\text{Be}$
- ▶ Solve NCSMC for  $c_{\lambda}(E)$ ,  $\gamma_{\nu}(r, E) \rightarrow \Psi(E)$
- ▶ Cross-section depends on transition matrix elements e.g.  $\langle \Psi(E_0) | M1 | \Psi(E) \rangle$

# Results

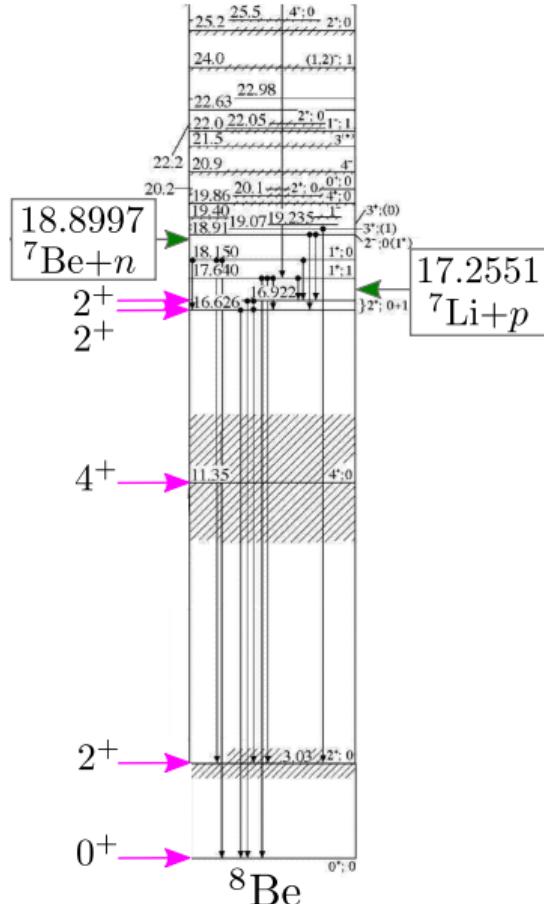
- ${}^8\text{Be}$  Structure
- Scattering:  ${}^7\text{Li}(p, p) {}^7\text{Li}$
- Transfer Reactions:  ${}^7\text{Li}(p, n) {}^7\text{Be}$ ,  ${}^7\text{Be}(n, p) {}^7\text{Li}$
- Radiative Capture:  ${}^7\text{Li}(p, \gamma) {}^8\text{Be}$
- Search for new physics:  ${}^7\text{Li}(p, X) {}^8\text{Be}$

## <sup>8</sup>Be Structure

Calculations of <sup>8</sup>Be “bound” states (w.r.t. <sup>7</sup>Li + *p* threshold) are improved by inclusion of the continuum ( $N_{max} = 9$ )

State	Energy [MeV]		
	NCSM	NCSMC	Experiment
0 <sup>+</sup>	-15.96	-16.13	-17.25
2 <sup>+</sup>	-12.51	-12.72	-14.23
4 <sup>+</sup>	-3.97	-4.31	-5.91
2 <sup>+</sup>	+0.76	-0.10	-0.63
2 <sup>+</sup>	+1.09	+0.31	-0.33

- ▶ Energies likely too high due to neglected  $\alpha + \alpha$  breakup
- ▶ Matches experiment well, except the 3rd 2<sup>+</sup> is still slightly above the <sup>7</sup>Li + *p* threshold

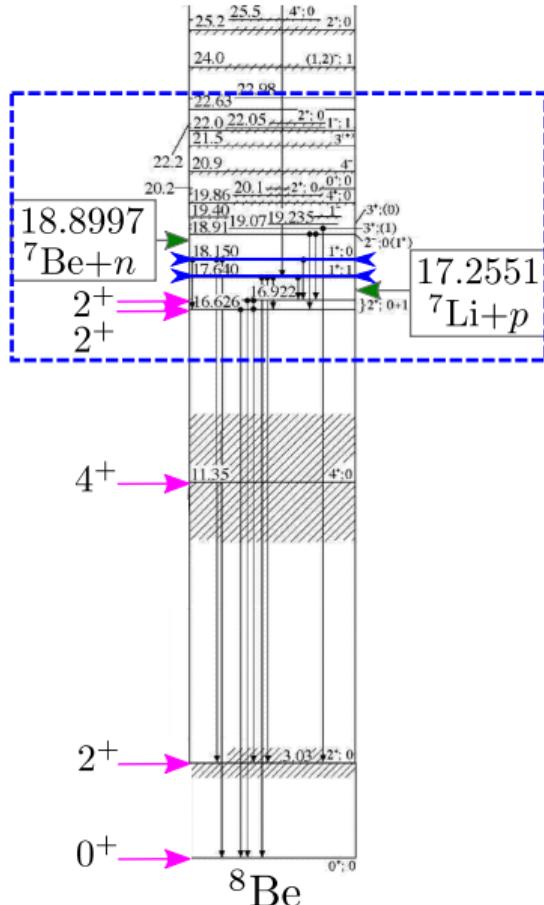


## <sup>8</sup>Be Structure

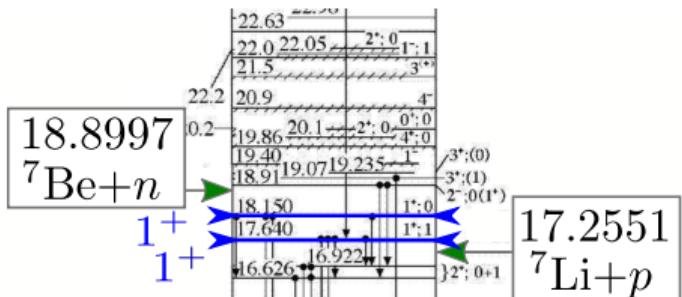
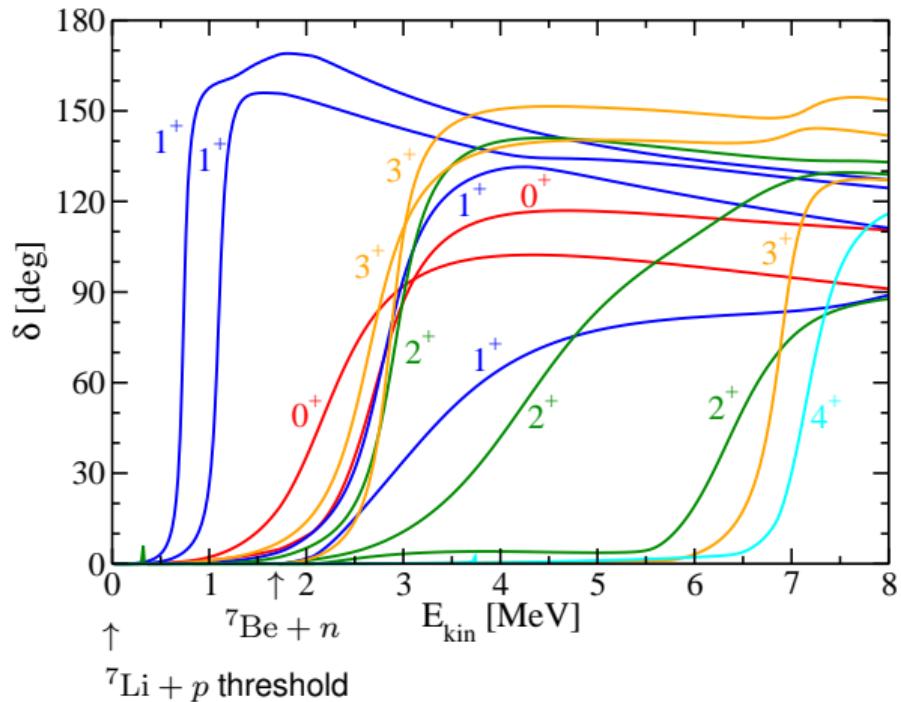
Calculations of <sup>8</sup>Be “bound” states (w.r.t. <sup>7</sup>Li + *p* threshold) are improved by inclusion of the continuum ( $N_{max} = 9$ )

State	Energy [MeV]		
	NCSM	NCSMC	Experiment
0 <sup>+</sup>	-15.96	-16.13	-17.25
2 <sup>+</sup>	-12.51	-12.72	-14.23
4 <sup>+</sup>	-3.97	-4.31	-5.91
2 <sup>+</sup>	+0.76	-0.10	-0.63
2 <sup>+</sup>	+1.09	+0.31	-0.33

- ▶ Energies likely too high due to neglected  $\alpha + \alpha$  breakup
- ▶ Matches experiment well, except the 3rd 2<sup>+</sup> is still slightly above the <sup>7</sup>Li + *p* threshold



## Eigenphase-shift Results (positive parity)

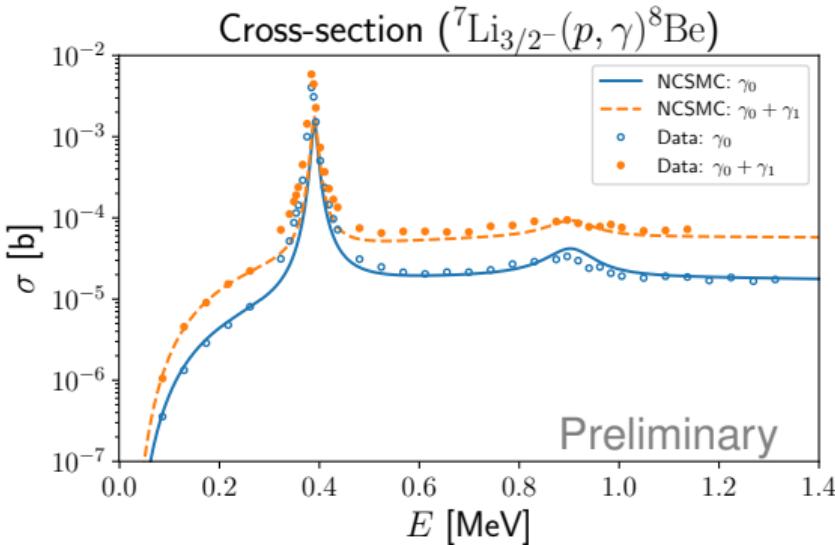


Additional resonances are seen compared to TUNL data

# Radiative Capture

$$^{2S+1}P_J : \left[ \left( |^7\text{Li}\rangle |p\rangle \right)^{(S)} Y_L(\hat{r}) \right]^J \\ P : L = 1$$

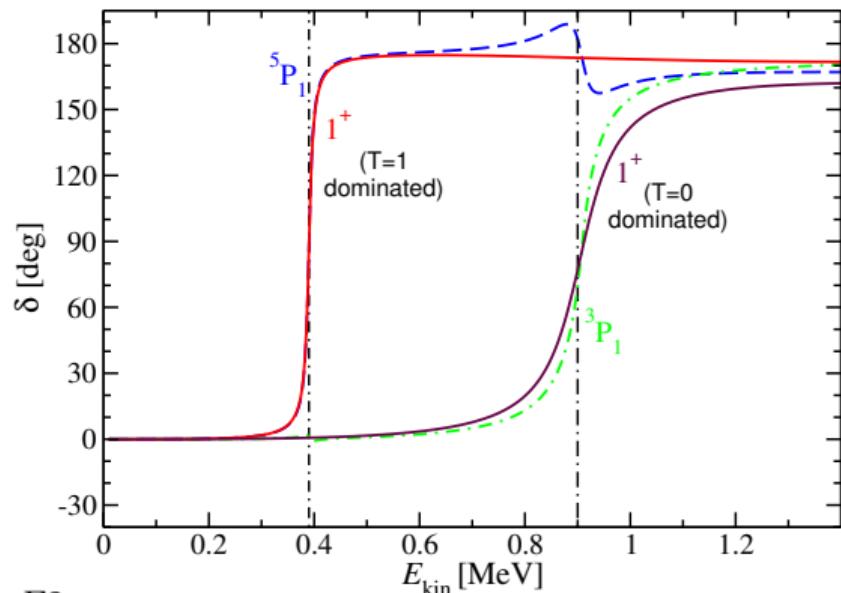
$^7\text{Li} + p$  phase shifts



$\gamma_0$ : decay to ground state ( $0^+$ )  
 $\gamma_1$ : decay to first excited ( $2^+$ )

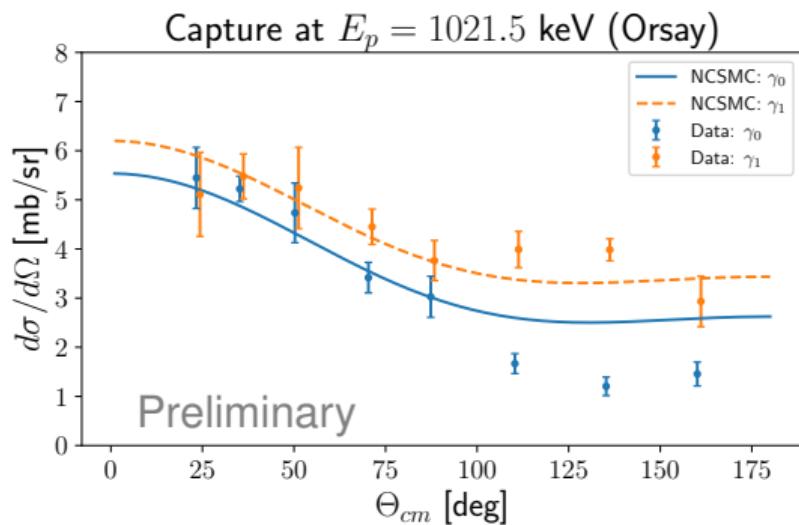
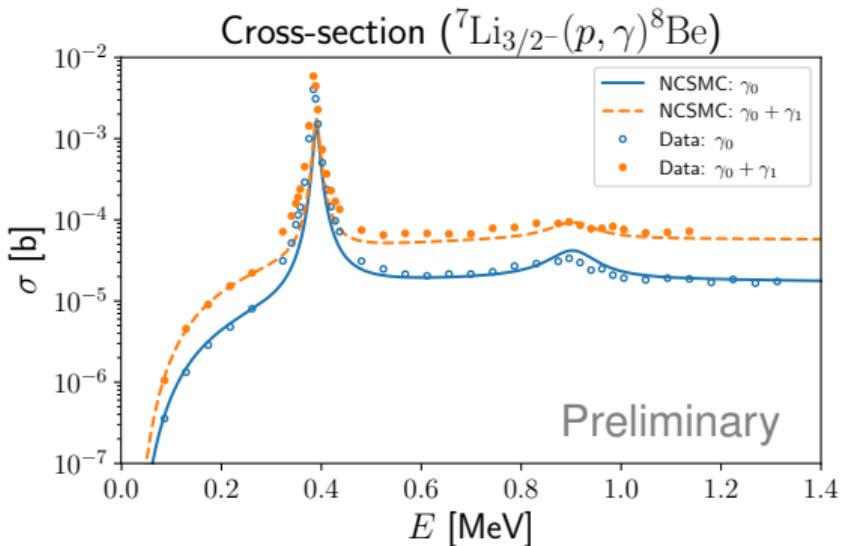
$$\hat{O}_\gamma = E1 + M1 + E2$$

Data: Zahnow et al  
Z.Phys.A 351 229-236 (1995)



Phenomenological adjustment: fit threshold and resonance positions to match experiment

## Radiative Capture (cont.)



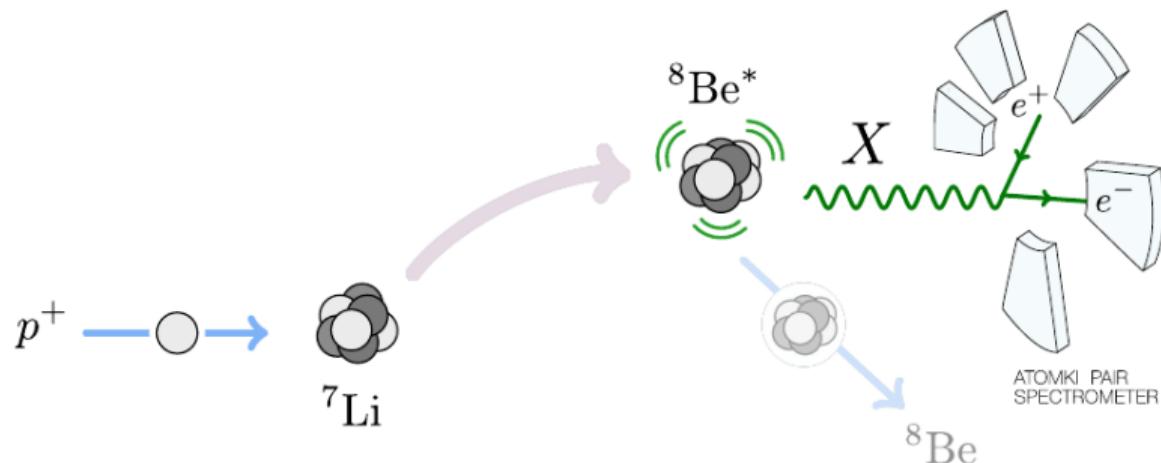
$\gamma_0$ : decay to ground state ( $0^+$ )  
 $\gamma_1$ : decay to first excited ( $2^+$ )

$$\hat{O}_\gamma = E1 + M1 + E2$$

Data: Zahnow et al  
Z.Phys.A 351 229-236 (1995)

## The X17 Anomaly in $p + {}^7\text{Li} \rightarrow {}^8\text{Be} + e^+e^-$

- ${}^7\text{Li}(p, e^+e^-) {}^8\text{Be}$  @ATOMKI (Hungary)
- The decay of  ${}^8\text{Be}$   $1^+$  excited states produces electron-positron pairs

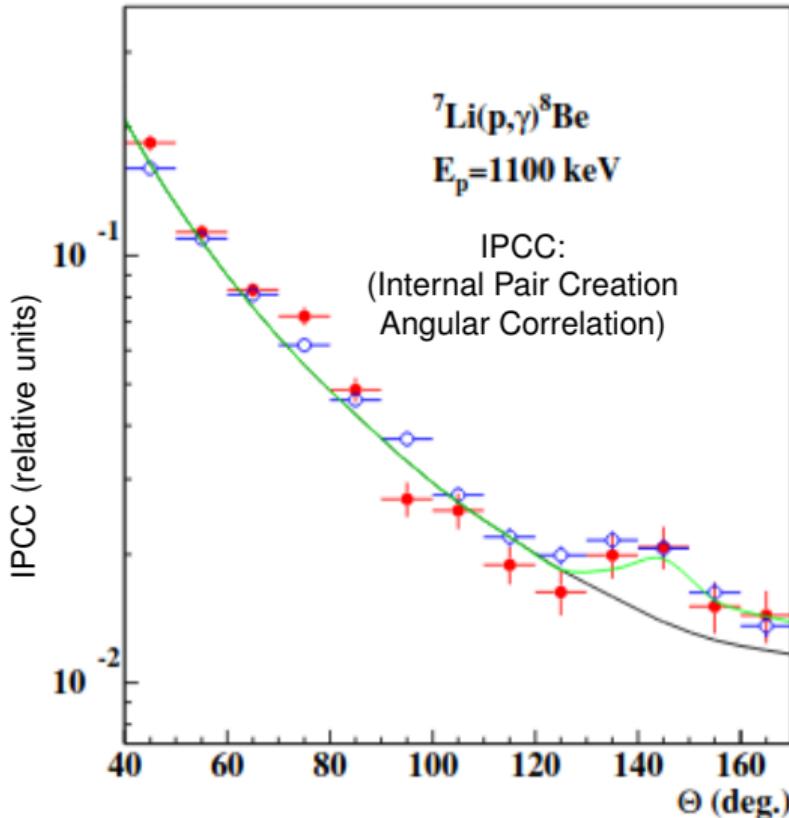


# The X17 Anomaly in $p + {}^7\text{Li} \rightarrow {}^8\text{Be} + e^+e^-$

Firak, Krasznahorkay, et al  
EPJ Web of Conferences 232 04005 (2020)

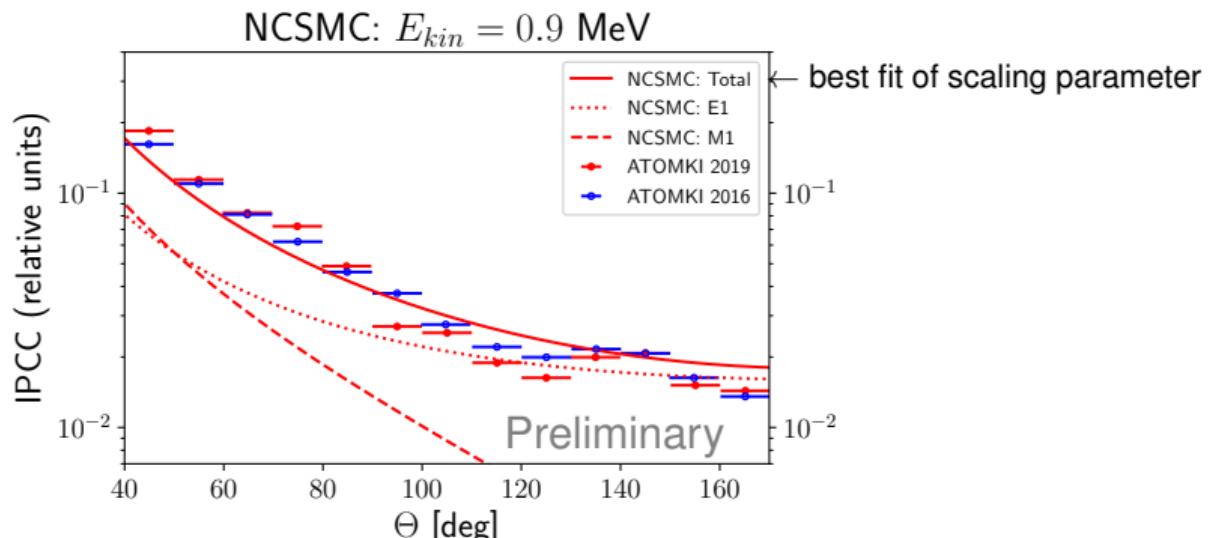
- The angle  $\theta$  between the electron and positron was measured
- Minimum angle from a massive intermediate particle:  $\theta \simeq \sin^{-1}\left(\frac{m_X}{E_X}\right)$
- Bump could be explained by 17 MeV bosons decaying to  $e^+e^-$

Can *ab initio* nuclear physics help interpret the anomaly?

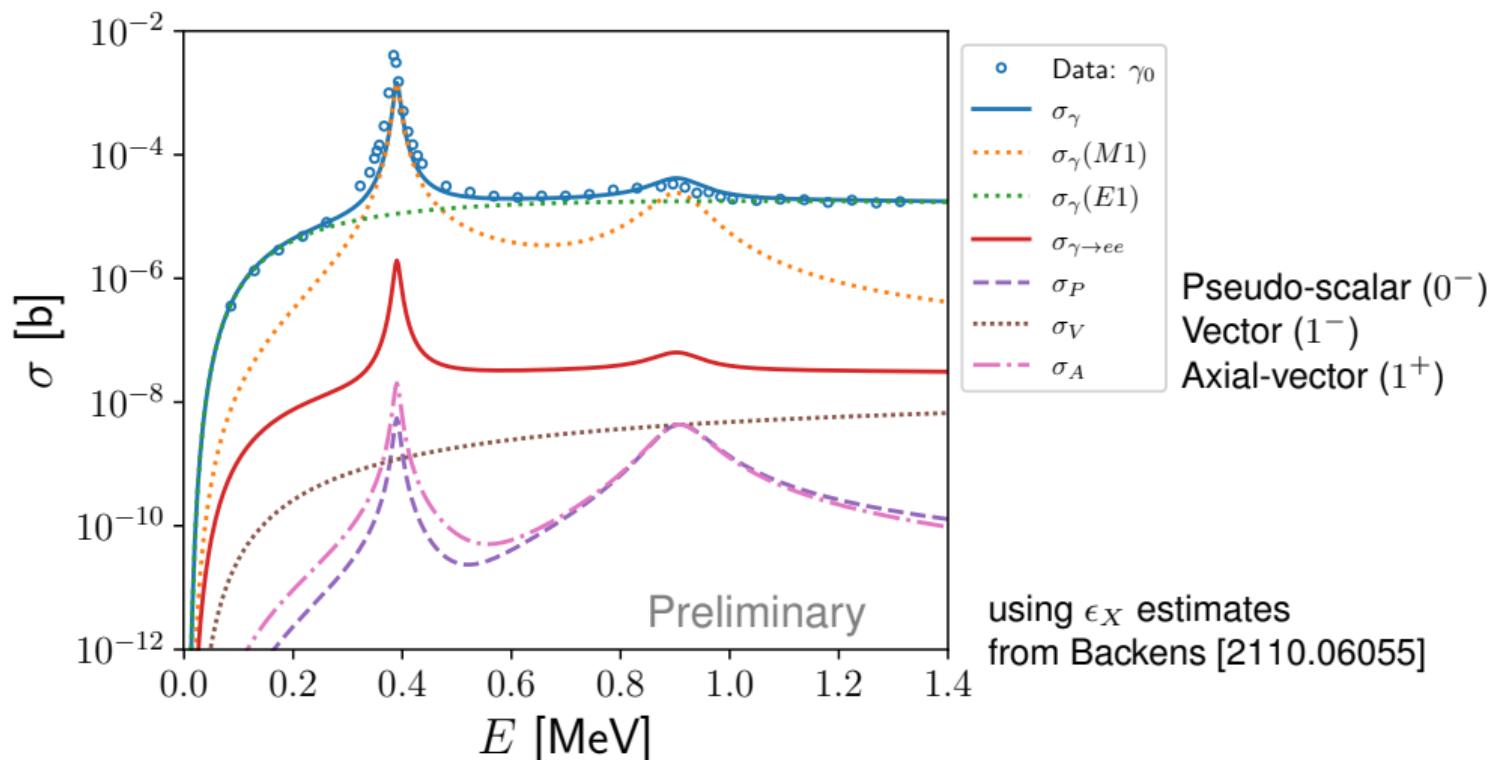


## Pair Production Distribution

- Our calculations (based on 2106.06834 ) under(over)-predict low(high) angles (possible background contamination or missing E1-M1 interference)
- Ongoing and planned experiments at Orsay and Montreal will provide an independent verification of the anomaly
- New ATOMKI data just published (2205.07744), analysis in progress



# Preliminary ${}^7\text{Li}(p, X){}^8\text{Be}$ Cross-sections



## Summary

- ▶ NCSMC successfully describes the spectrum of  ${}^8\text{Be}$  including the  $1^+$  resonances
- ▶ Calculations of  ${}^7\text{Li}(p, \gamma){}^8\text{Be}$  radiative capture match data

## Outlook

- ▶ Compare  ${}^7\text{Li}(p, e^+e^-){}^8\text{Be}$  to data with  $\gamma \rightarrow e^+e^-$  operator and various  $X \rightarrow e^+e^-$  operators (e.g. axions, vector bosons, axial vector bosons)
- ▶ Calculations of  ${}^3\text{H}(p, e^+e^-){}^4\text{He}$  are also relevant to the X17 anomaly
- ▶ Explore charge-exchange reactions relevant for nucleosynthesis:  
 ${}^7\text{Be}(n, p){}^7\text{Li}$ ,  ${}^7\text{Li}(p, n){}^7\text{Be}$

## Backup Slides

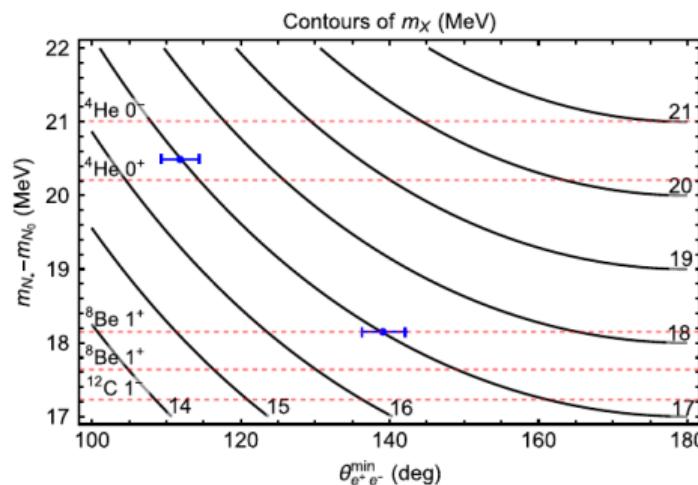
## Constraints on $m_X$

Feng PRD 95, 035017 (2017)

In the frame of the  $X$  boson the electron and positron momenta are anti-parallel.  
Boosted to a minimum separation angle:

$$\theta = 2 \sin^{-1} \left( \frac{m_X}{E_X} \right)$$

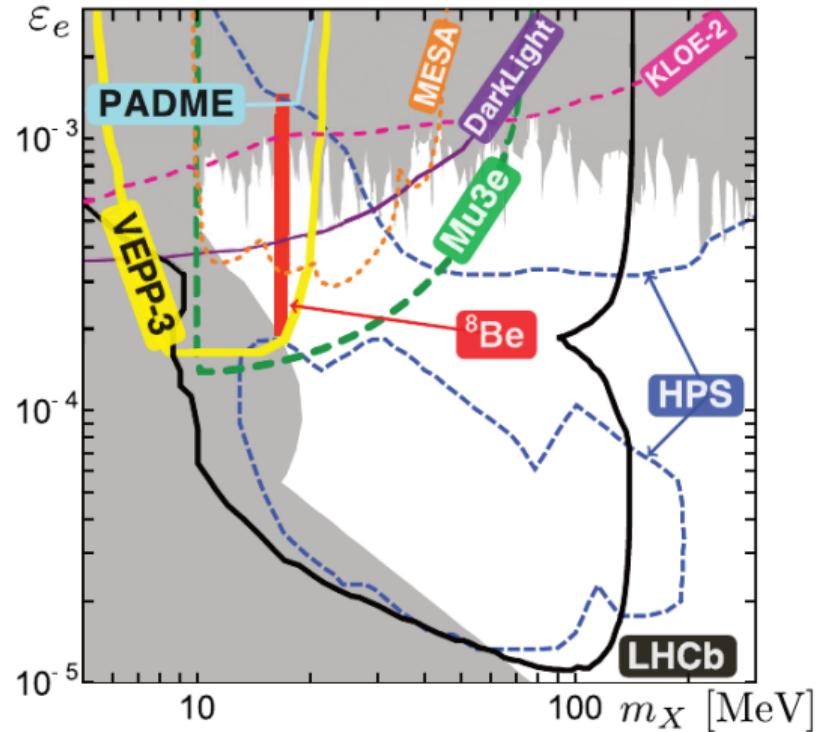
- $^8\text{Be}$  anomaly occurs in the isoscalar transition (decay of  $1^+0$  resonance)
- In-between resonances in  $^4\text{He}$
- Bumps could be explained by 17 MeV bosons decaying to  $e^+e^-$



# Exclusion Plot

Current (gray) and projected sensitivities of future experiments

Feng PRD 95 035017 (2017)



## X17 Candidate Bosons

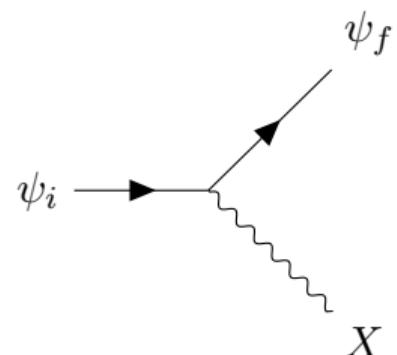
$$(m_X \simeq 17 \text{ MeV}, \Delta E \geq 17.2251 \text{ MeV} [{}^7\text{Li} + p], \\ k_X = \sqrt{\Delta E^2 - m_X^2}, k_\gamma = \Delta E)$$

Operators for  $1^+ \rightarrow 0^+$  decay (in the long-wavelength approximation)

- **Pseudo-scalar** ( $0^-$ ):  $\langle X_P \rangle \sim \epsilon_P \left\langle \hat{S} \right\rangle k_X$
- **Axial-vector** ( $1^+$ ):  $\langle X_A \rangle \sim \epsilon_A \left\langle \hat{S} \right\rangle \sqrt{2 + \frac{m_X^2}{\Delta E^2}}$
- **Vector** ( $1^-$ ):  $\langle X_V \rangle \sim \epsilon_V \left\langle \hat{O}_\gamma \right\rangle \frac{k_X}{k_\gamma}$
- **For comparison:**  $\gamma$  (E1 ( $1^-$ ), M1 ( $1^+$ ), E2 ( $2^+$ ), etc)

$$\langle E1 \rangle \sim \langle rY_1 \rangle k_\gamma$$

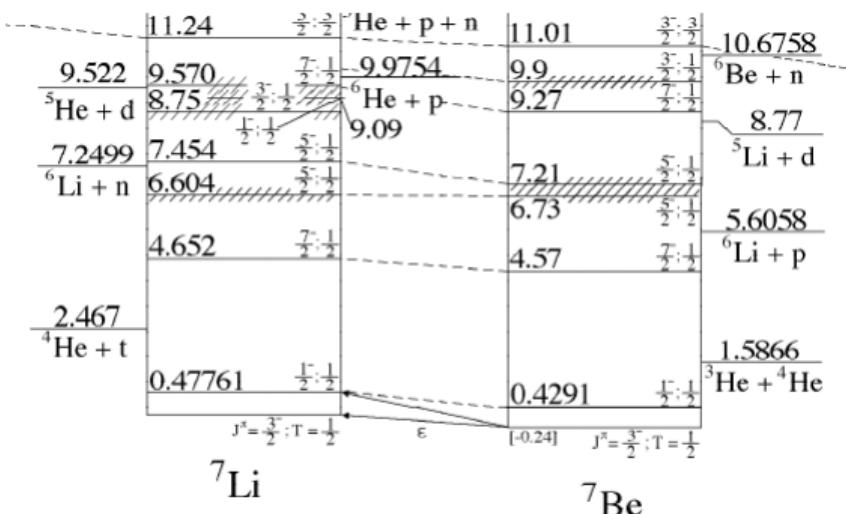
$$\langle M1 \rangle \sim \left( g_l \left\langle \hat{L} \right\rangle + g_s \left\langle \hat{S} \right\rangle \right) k_\gamma$$



# Input States from NCSM

$$\Psi_{\text{NCSMC}}^{(8)} = \sum_{\lambda} c_{\lambda} |{}^8\text{Be}, \lambda\rangle + \sum_{\nu} \int dr \gamma_{\nu}(r) \hat{A}_{\nu} |{}^7\text{Li} + p, \nu\rangle + \sum_{\mu} \int dr \gamma_{\mu}(r) \hat{A}_{\mu} |{}^7\text{Be} + n, \mu\rangle$$

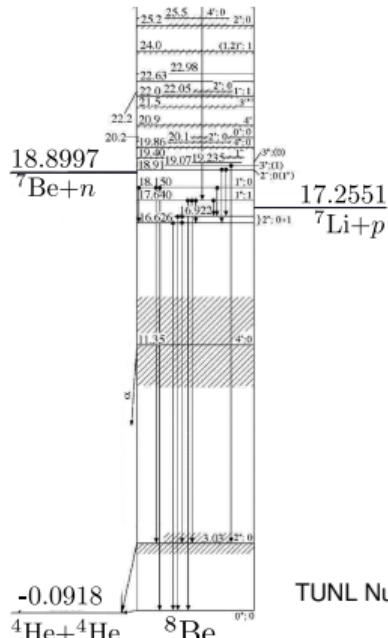
- ▶ 3 NCSM calculations:  ${}^7\text{Li}$ ,  ${}^7\text{Be}$  and  ${}^8\text{Be}$
- ▶  $\{\frac{3}{2}^-, \frac{1}{2}^-, \frac{7}{2}^-, \frac{5}{2}^-, \frac{5}{2}^-\}$   ${}^7\text{Li}$  and  ${}^7\text{Be}$  states in cluster basis
- ▶ 15 positive and 15 negative parity states in  ${}^8\text{Be}$  composite state basis



# Input States from NCSM

$$\Psi_{\text{NCSMC}}^{(8)} = \sum_{\lambda} c_{\lambda} |{}^8\text{Be}, \lambda\rangle + \sum_{\nu} \int dr \gamma_{\nu}(r) \hat{A}_{\nu} |{}^7\text{Li} + p, \nu\rangle + \sum_{\mu} \int dr \gamma_{\mu}(r) \hat{A}_{\mu} |{}^7\text{Be} + n, \mu\rangle$$

- ▶ 3 NCSM calculations:  ${}^7\text{Li}$ ,  ${}^7\text{Be}$  and  ${}^8\text{Be}$
- ▶  $\{\frac{3}{2}^-, \frac{1}{2}^-, \frac{7}{2}^-, \frac{5}{2}^-, \frac{5}{2}^-\}$   ${}^7\text{Li}$  and  ${}^7\text{Be}$  states in cluster basis
- ▶ 15 positive and 15 negative parity states in  ${}^8\text{Be}$  composite state basis



TUNL Nuclear Data Evaluation Project