

# Strongly-coupled superconductivity from quantum cluster approaches

David Sénéchal

département de physique & Institut quantique  
Université de Sherbrooke

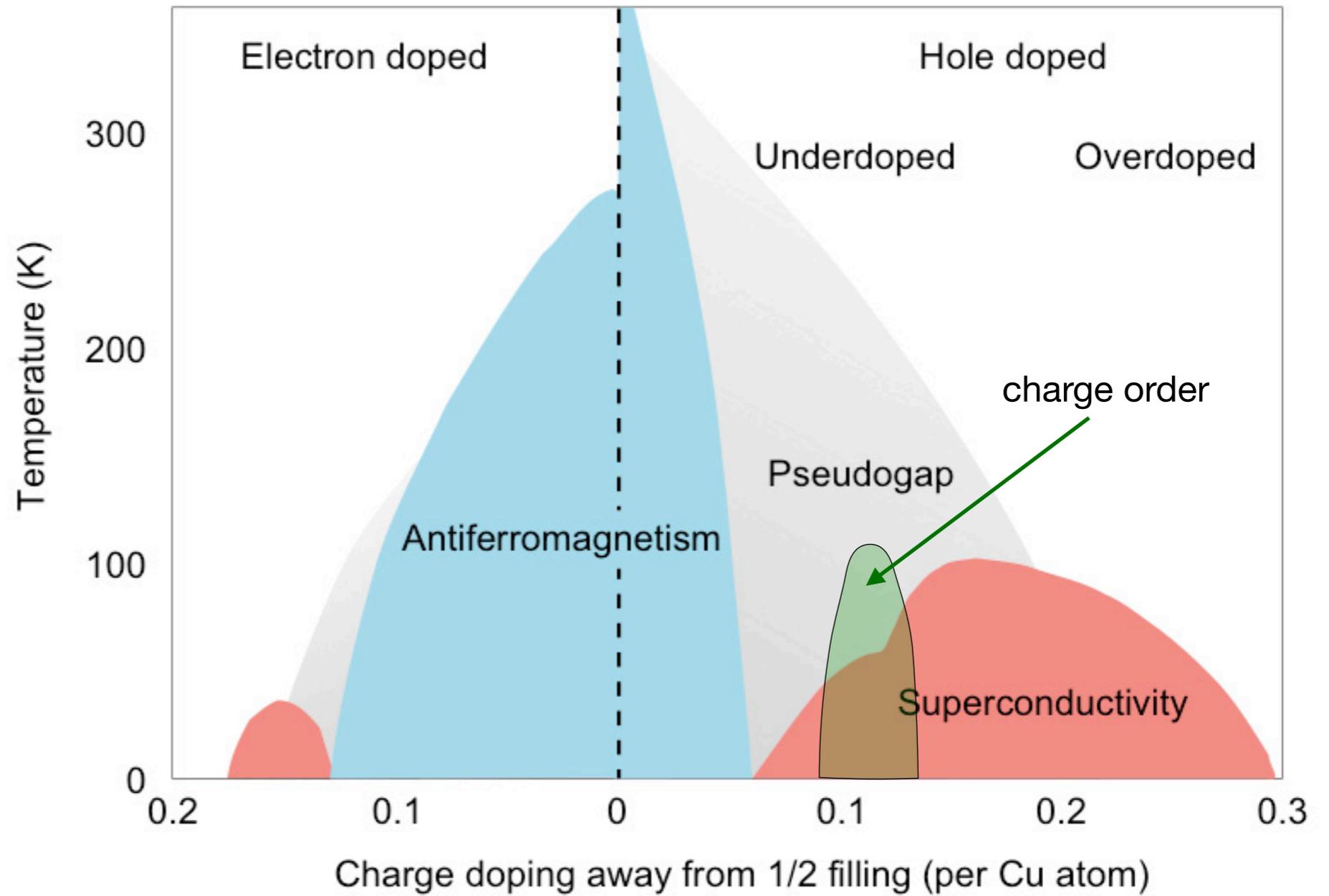
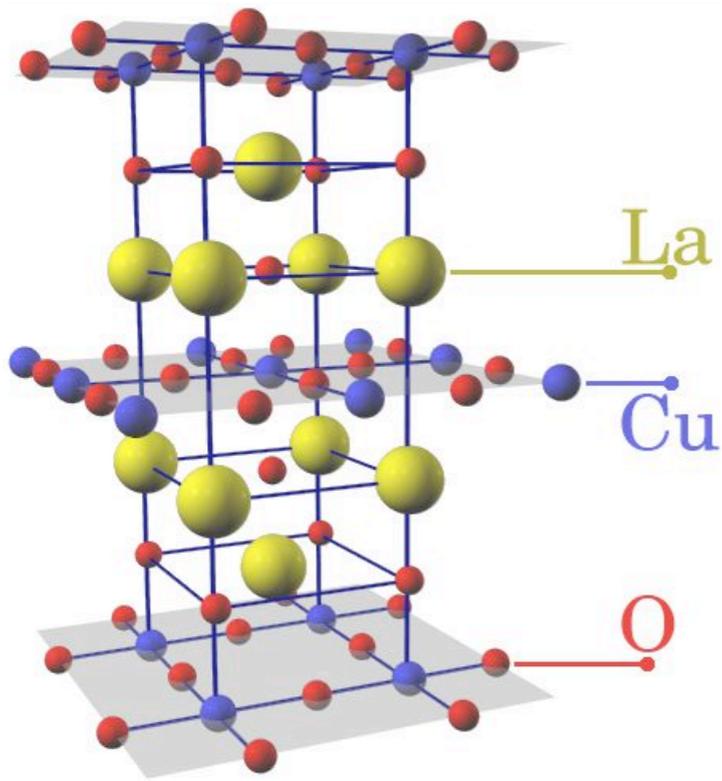
**CAP Congress, June 5 2019**



# Outline

- Hubbard models
- Quantum cluster methods :
  - VCA (variational cluster approximation)
  - CDMFT (cluster dynamical mean field theory)
- Superconductivity in cuprates from CDMFT
  - 1-band Hubbard model
  - 3-band Hubbard model
- Charge order in cuprates
- Loop currents

# High- $T_c$ superconductors



<http://www.mrsec.umn.edu/research/seeds/2011/greven2011.html>

# The Hubbard model

hopping amplitude ←

number of spin ↑ electrons at  $\mathbf{r}$

$$H = \sum_{\mathbf{r}, \mathbf{r}', \sigma} t_{\mathbf{r}, \mathbf{r}'} c_{\mathbf{r}\sigma}^\dagger c_{\mathbf{r}'\sigma} + U \sum_{\mathbf{r}} n_{\mathbf{r}\uparrow} n_{\mathbf{r}\downarrow} - \mu \sum_{\mathbf{r}, \sigma} n_{\mathbf{r}, \sigma}$$

creation operator ←

repulsion

- Simplest model for strong correlations
- AF phase at half-filling + Mott at higher temperature
- Paradigm: *high- $T_c$  superconductors are doped Mott/AF insulators*
- Progress done with quantum cluster methods:
  - Need short-range fluctuations to understand superconductivity
  - Relation with short-range AF fluctuations made evident numerically

# The 3-band Hubbard model

creates a p electron ←

→ creates a d electron

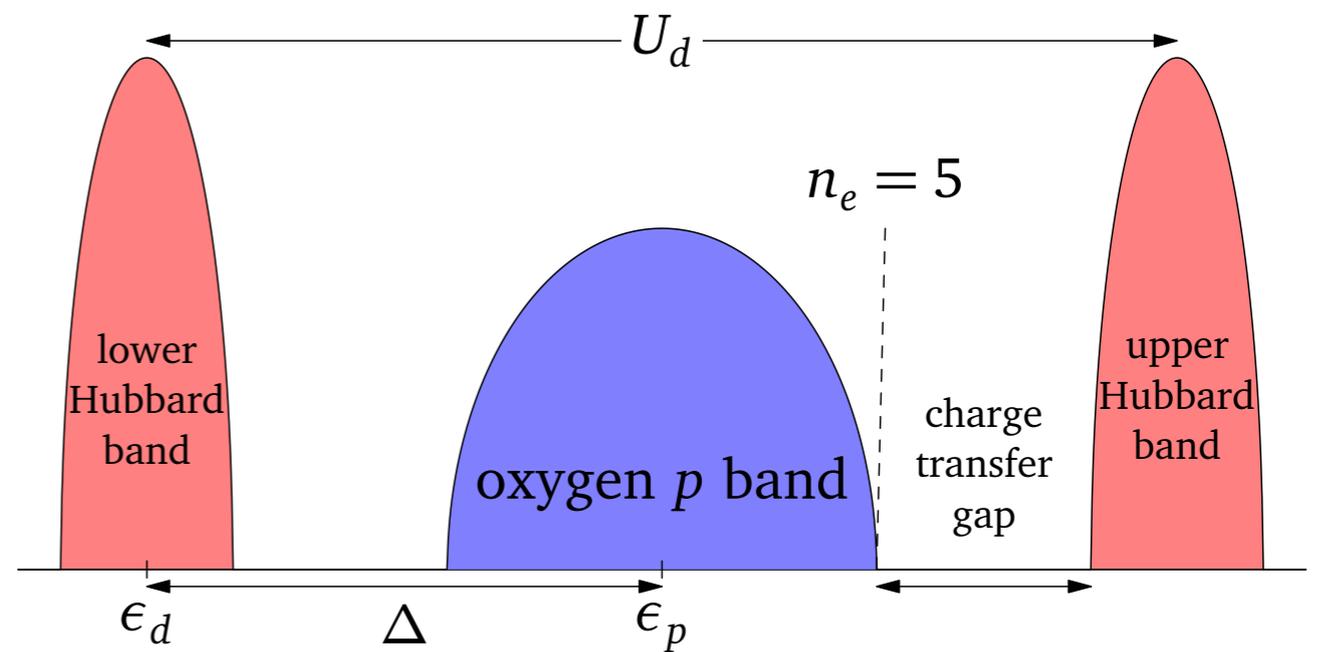
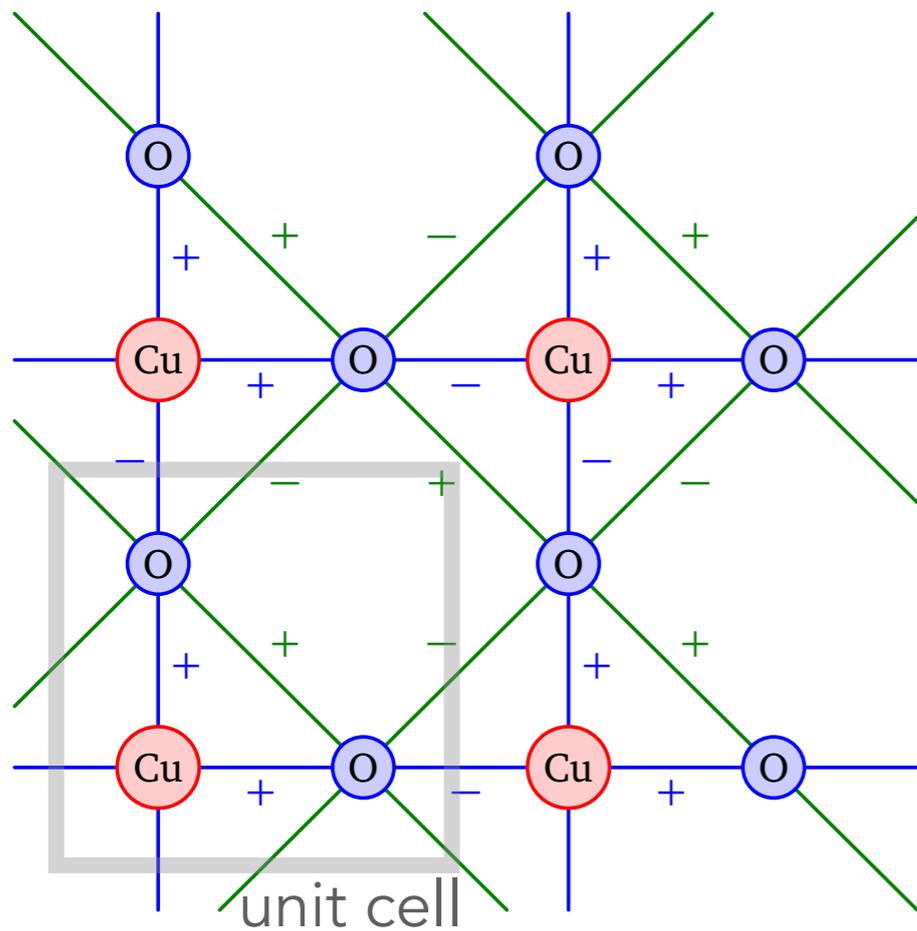
$$H = t_{pd} \sum_{\langle i,j \rangle} \left( p_{i\sigma}^\dagger d_{i\sigma} + d_{i\sigma}^\dagger p_{j\sigma} \right) + t_{pp} \sum_{\langle j,j' \rangle} \left( p_{j\sigma}^\dagger p_{j'\sigma} + \text{H.c.} \right)$$

p orbital level ←

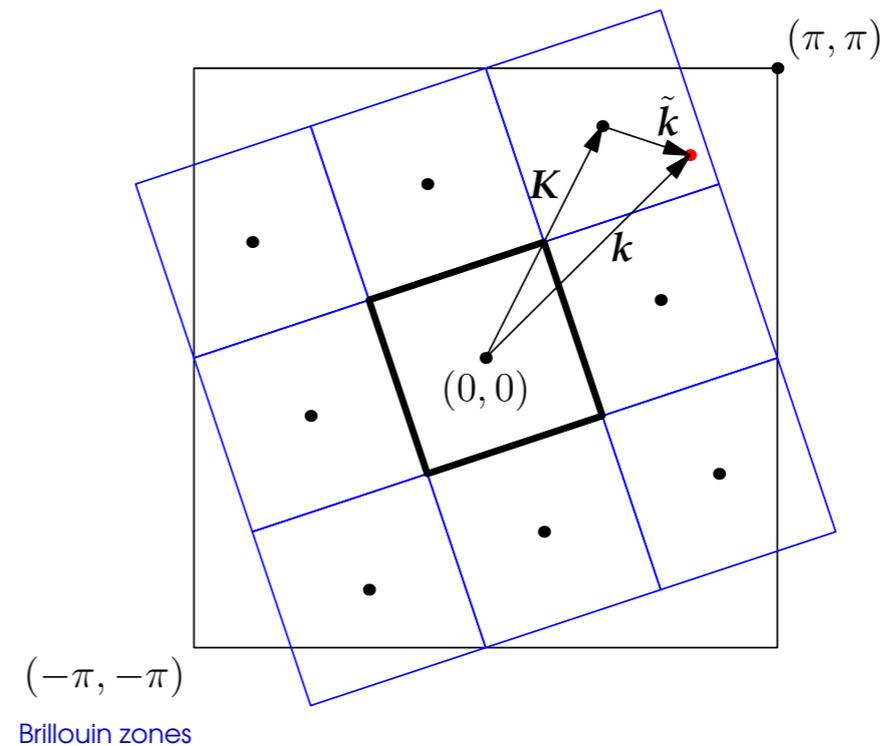
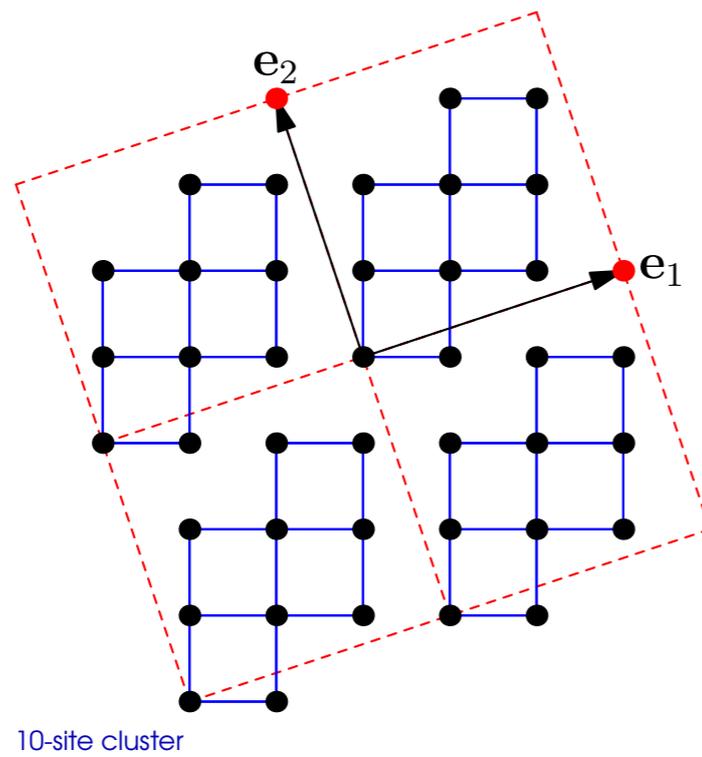
Cu-Cu repulsion ←

$$+ U_d \sum_i n_{i\uparrow}^{(d)} n_{i\downarrow}^{(d)} + (\epsilon_d - \mu) \sum_{i,\sigma} n_{i\sigma}^{(d)} + (\epsilon_p - \mu) \sum_{j,\sigma} n_{j\sigma}^{(p)}$$

→ d orbital level



# Quantum cluster methods



- Based on Green functions
- Lattice tiled by identical units (clusters)
- Clusters used to compute the electron self-energy  $\Sigma$ , otherwise full lattice dispersion relation used
- Different schemes to embed clusters in the lattice model (VCA, CDMFT,...)
- Different **impurity solvers** (ED, CT-HYB, CT-INT, CT-AUX, ...)

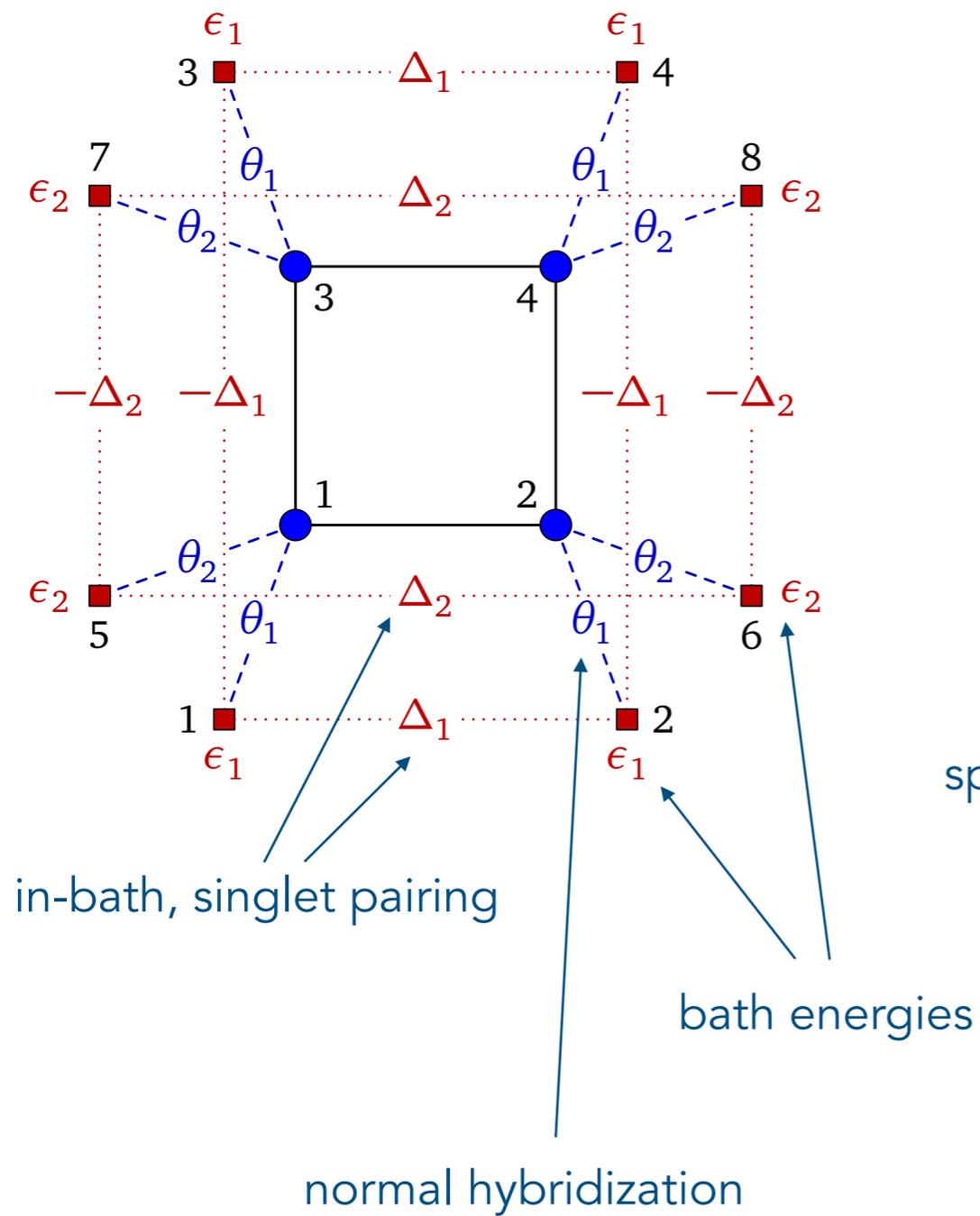
dispersion relation  $\leftarrow$   $\rightarrow$  cluster self-energy

$$\mathbf{G}^{-1}(\tilde{\mathbf{k}}, \omega) = \mathbf{t}(\tilde{\mathbf{k}}) - \Sigma'(\omega)$$

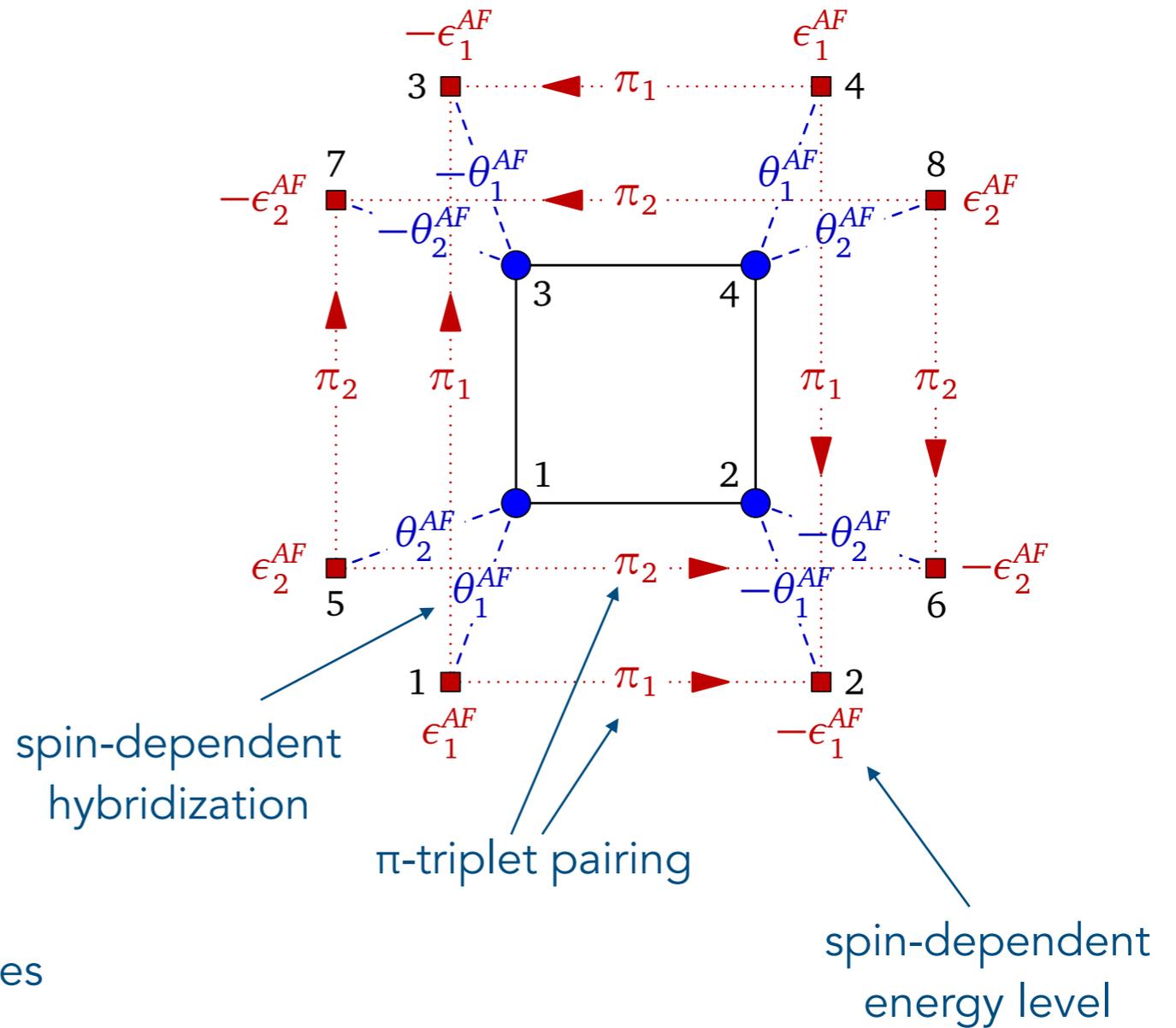
*Exact diagonalization at  $T=0$  in this presentation*

# CDMFT : simple bath parametrization for d-wave SC

pure d-wave SC



d-wave SC + AF

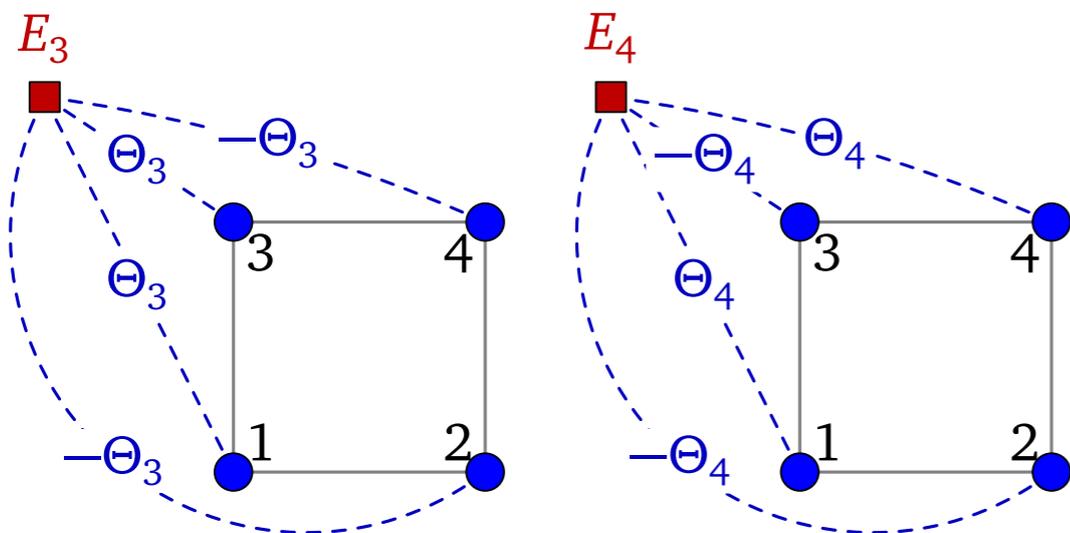
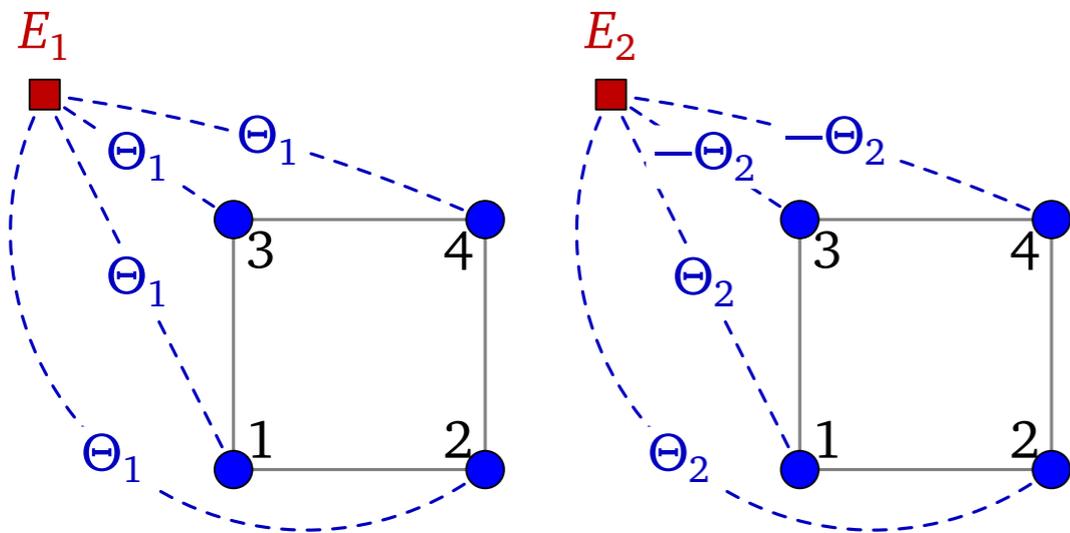


bath is *not* diagonal  
6 or 12 parameters

# Bath parametrization according to cluster symmetry

## pure d-wave SC

$C_{2v}$  : 4 irreducible representations



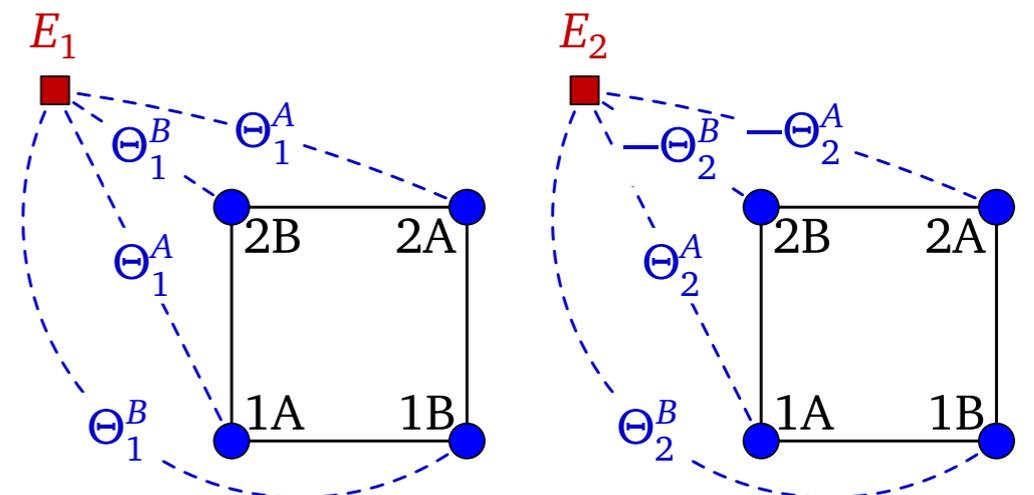
$$\begin{array}{c} \blacksquare \\ E_i \end{array} \text{---} \pm \Theta_i \text{---} \begin{array}{c} \bullet \\ j \end{array} \equiv \pm(\theta_i a_i^\dagger c_j + \Delta_i a_i (i\sigma_y) c_j + \text{H.c.})$$

normal

singlet

## d-wave SC + AF

$C_2$  : 2 irreducible representations



$$\begin{array}{c} \blacksquare \\ E_i \end{array} \text{---} \pm \Theta_i^X \text{---} \begin{array}{c} \bullet \\ jX \end{array} \equiv \pm(\theta_{iX} a_i^\dagger c_{jX} + \Delta_{iX} a_i (i\sigma_y) c_{jX} \\ + \theta_{iX}^s a_i^\dagger \sigma_z c_{jX} + T_{iX} a_i \sigma_x c_{jX} + \text{H.c.})$$

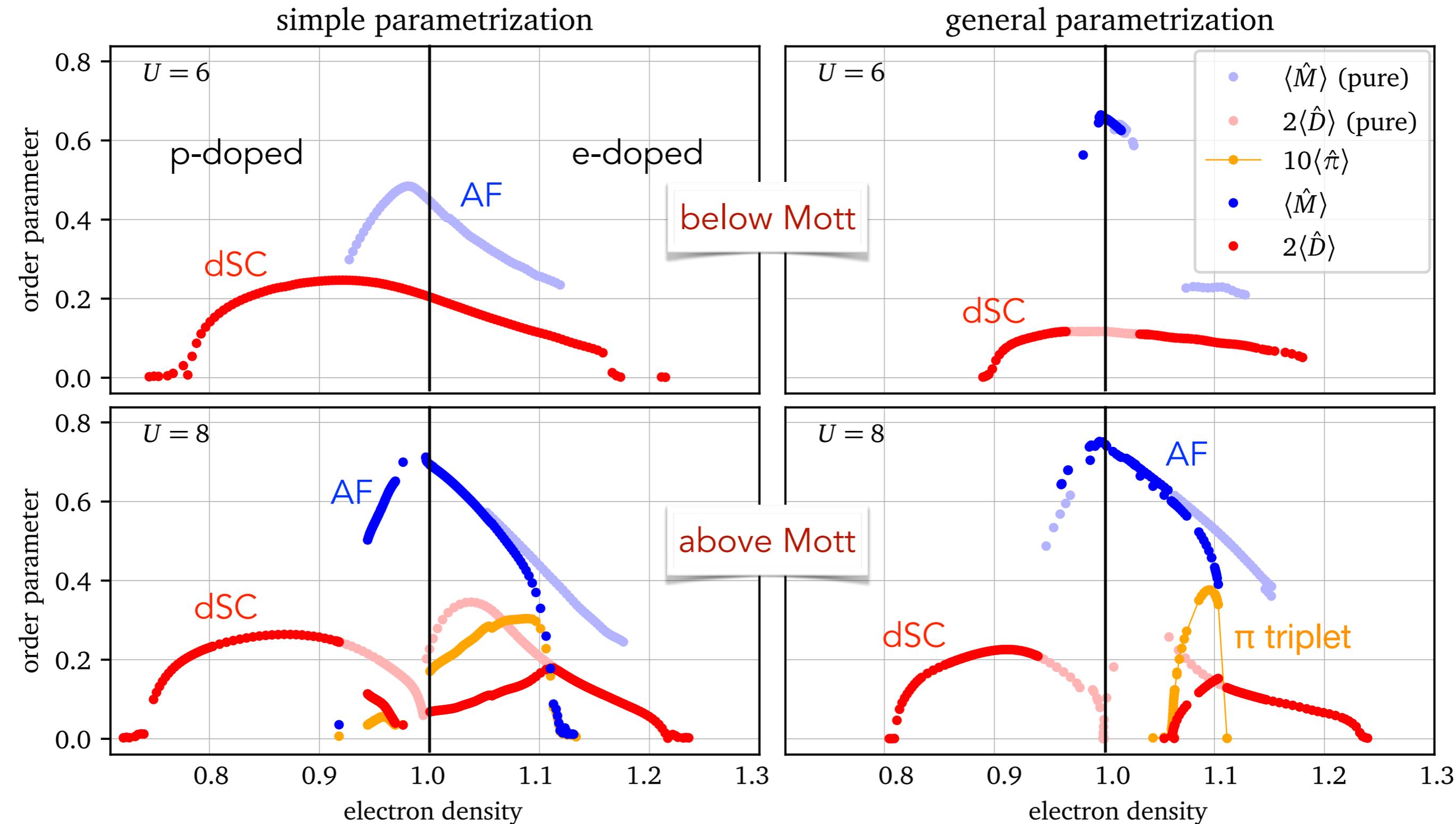
spin dependent

triplet

singlet

bath is diagonal  
24 or 72 parameters

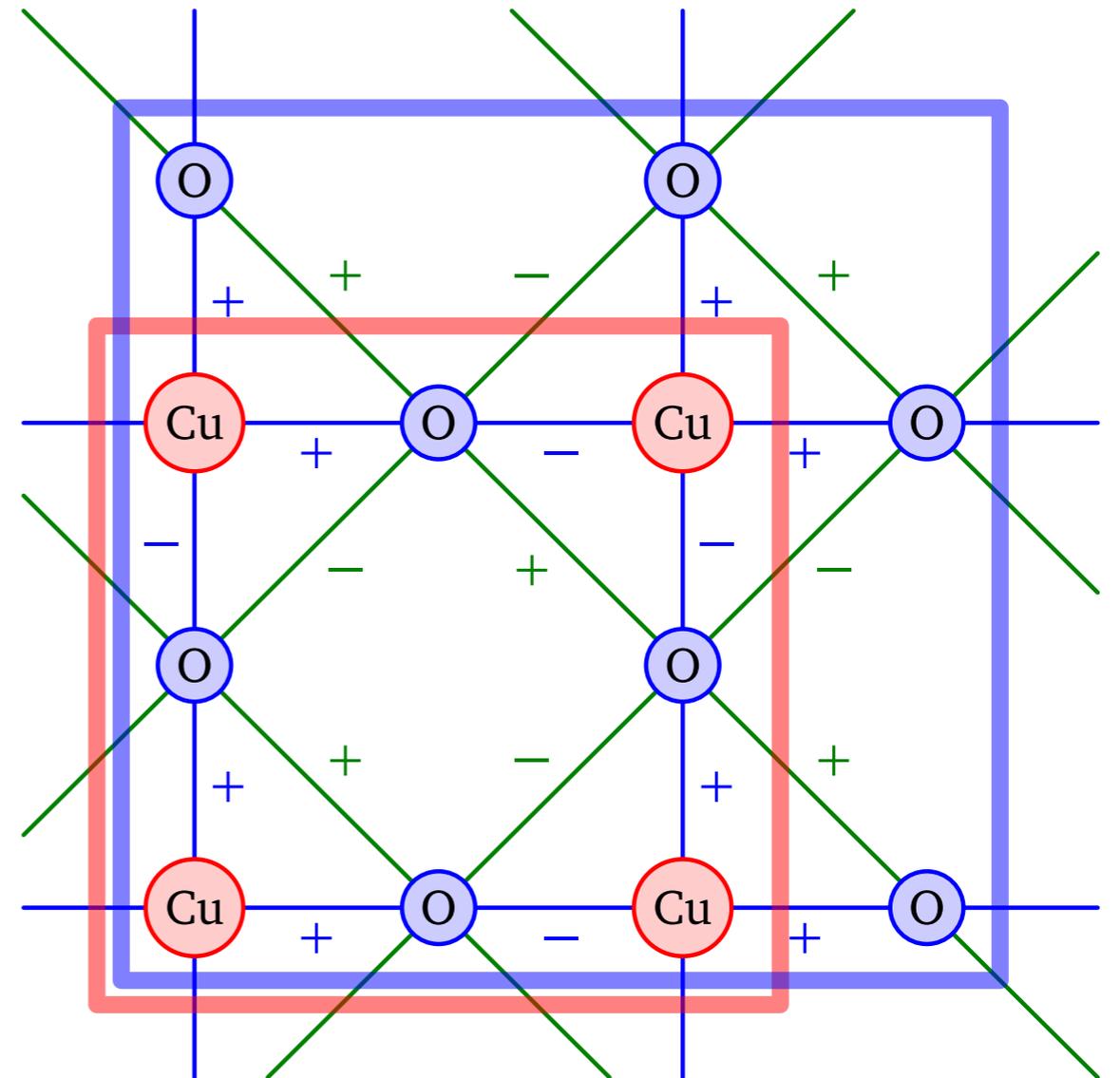
# CDMFT : One-band model for YBCO



A. Foley, S. Verret, AM Tremblay & DS, Phys. Rev. B **99**, 184510 (2019)

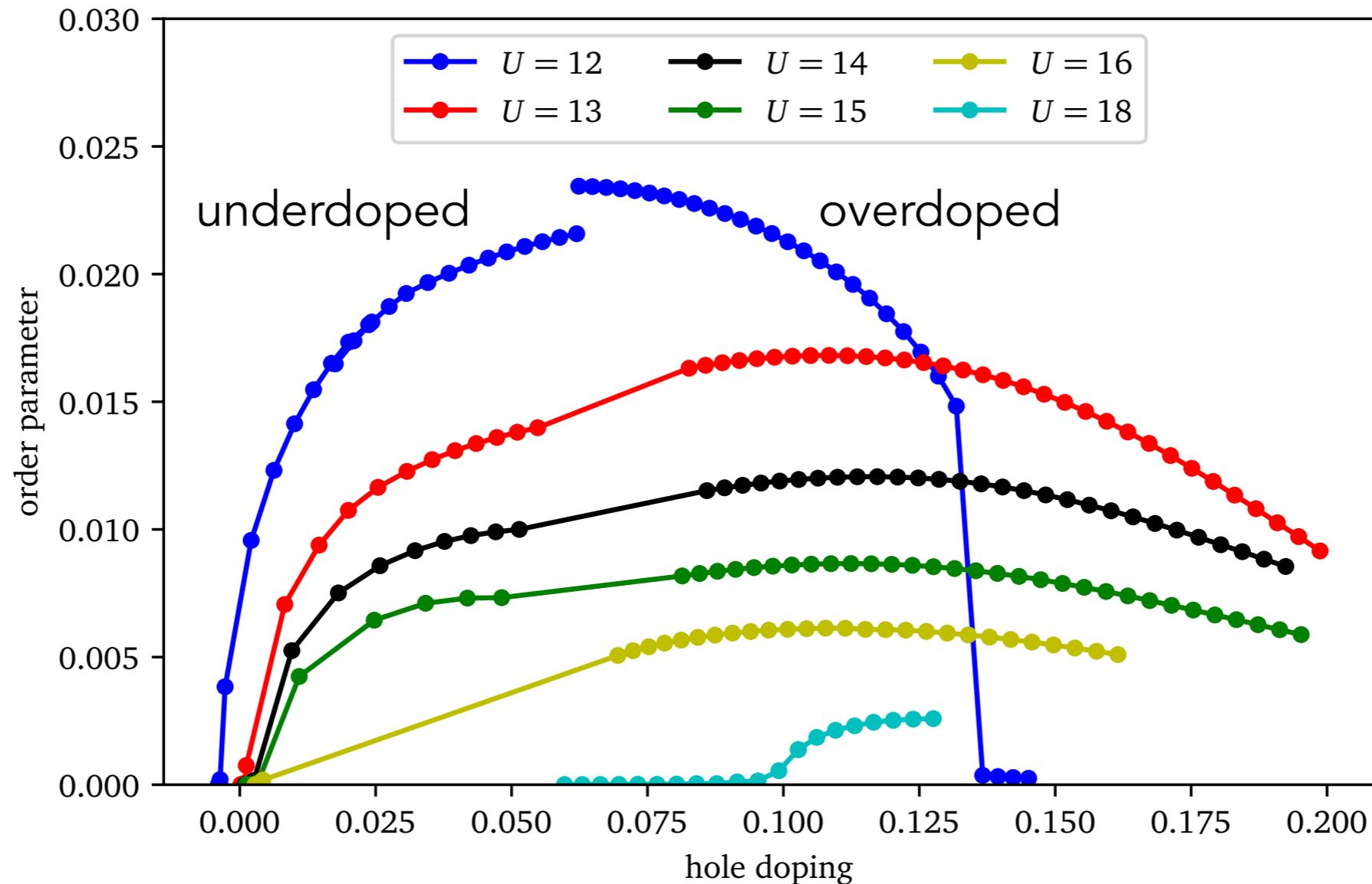
# CDMFT and the 3-band model

- Different, overlapping clusters for Cu and O atoms
- General bath parametrization for Cu cluster
- O cluster is non-interacting: acts like a constant hybridization and has effect only via CDMFT self-consistency relation



parameter set:  $t_{pp} = 1$   $t_{pd} = 1.5$   $\epsilon_p = 7.0$   $\epsilon_d = 0$

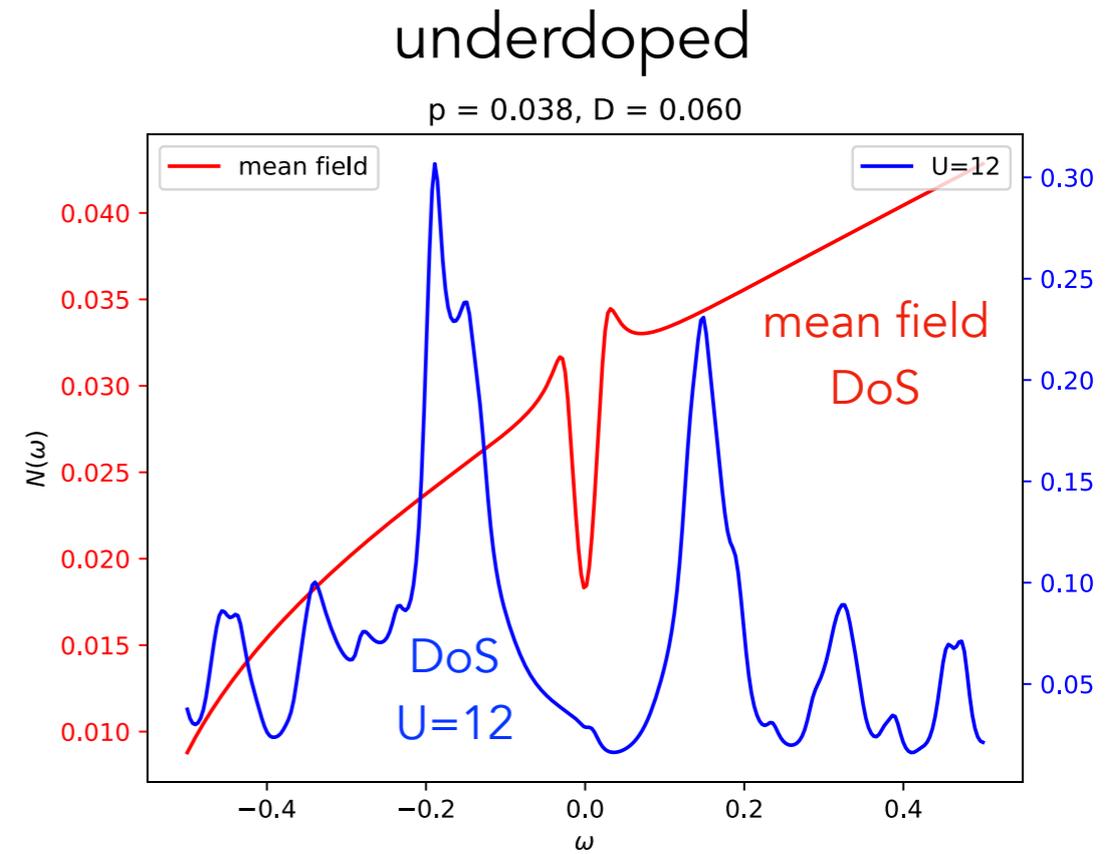
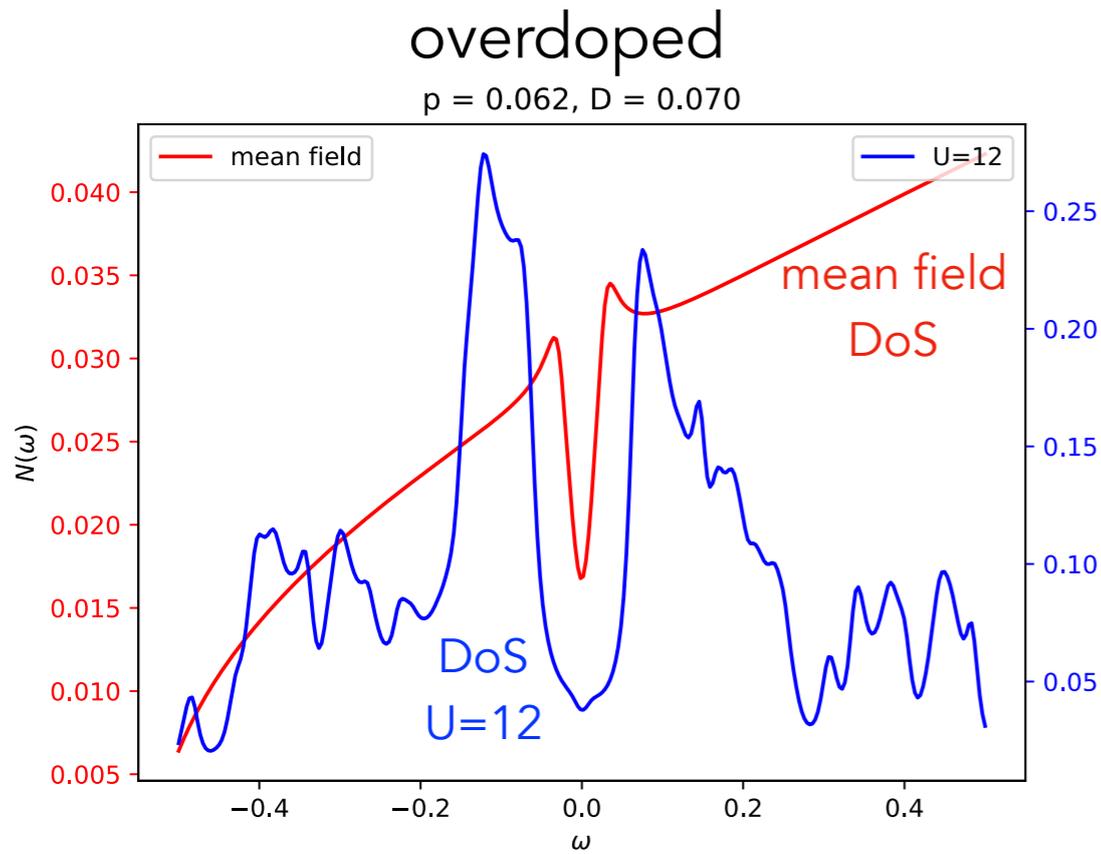
# 3-band model : dSC order parameter (CDMFT)



First-order transition between **underdoped** and **overdoped** SC region of coexistence between two SC solutions (**BEC — BCS**)

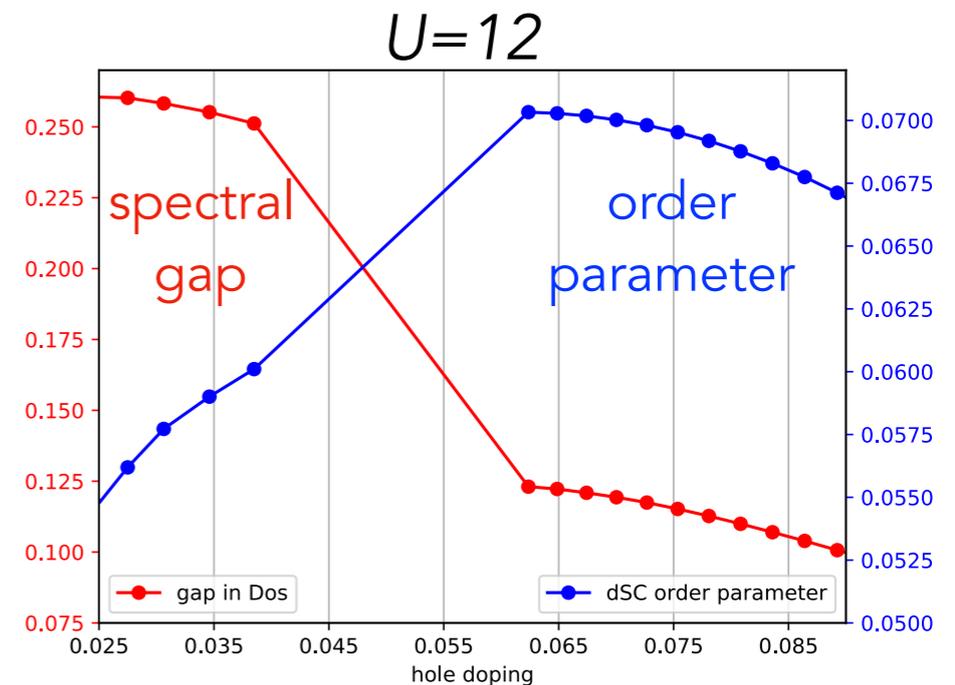
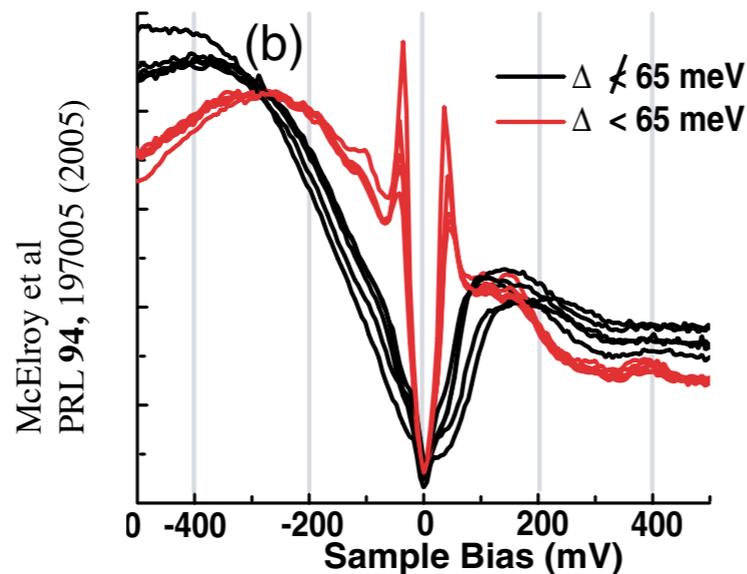
antiferromagnetism **not** probed in this study

# 3-band model : underdoped vs overdoped



**overdoped** : the gap increases with the order parameter

**underdoped** : the gap jumps up (pseudogap) & the order parameter decreases

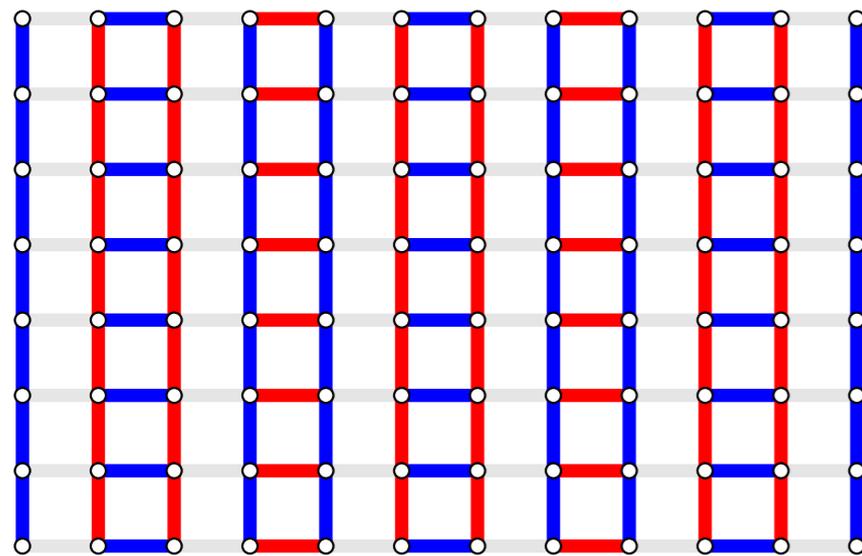


S. Dash & DS, in preparation

# Charge order

for a review : Comin & Damascelli, Annu. Rev. Condens. Matter Phys. 2016. 7:369

- Charge modulation located on the oxygen atoms
- In the 1-band model  $\rightarrow$  bond density wave (BDW)



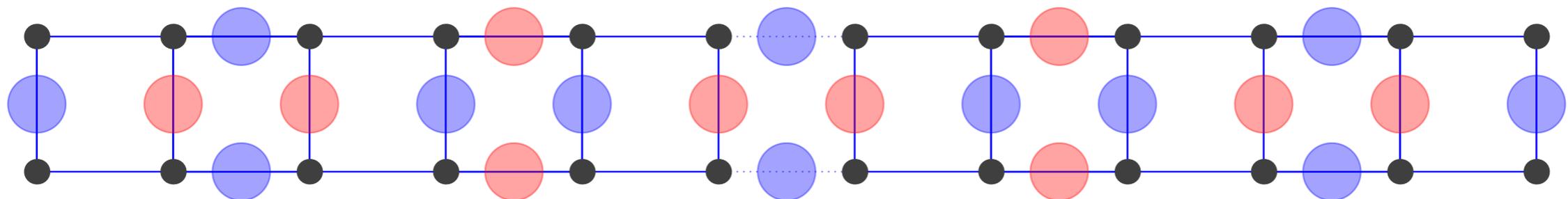
↑  
positive

↑  
negative

VCA clusters

$$\hat{\Psi}_{\text{BDW}} = \sum_{\mathbf{r}, \sigma, \mathbf{a}} t_{\mathbf{q}, \mathbf{a}} c_{\mathbf{r}, \sigma}^{\dagger} c_{\mathbf{r} + \mathbf{a}, \sigma} e^{i(\mathbf{q} \cdot \mathbf{r} + \mathbf{a}/2)} + \text{H.c}$$

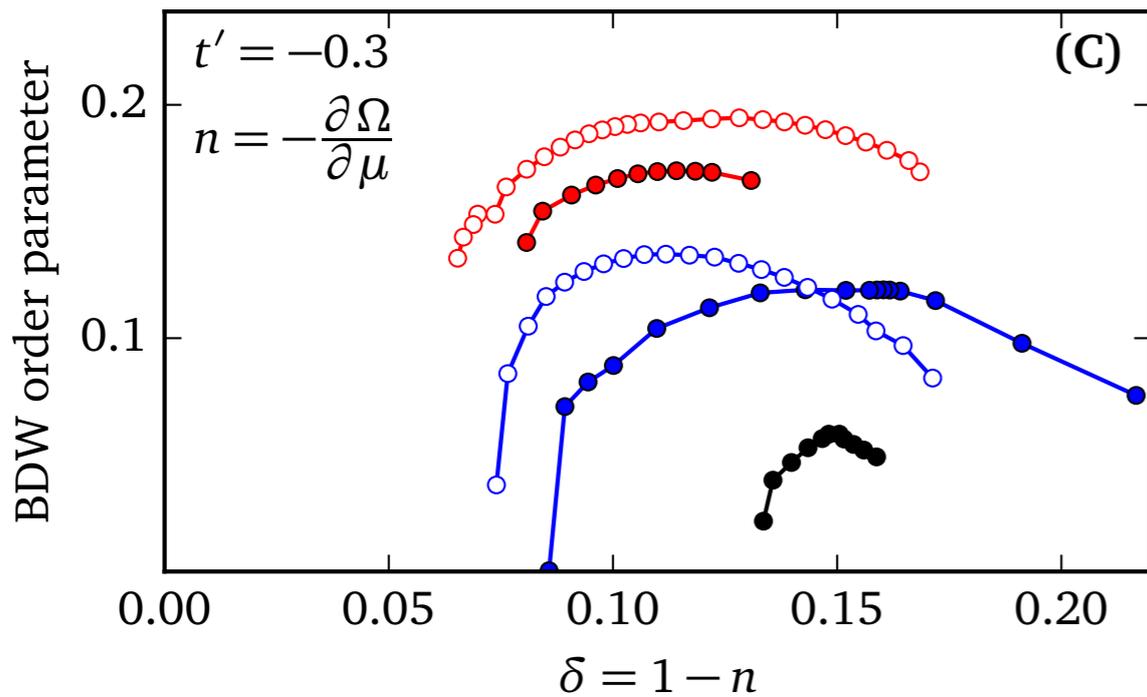
$$\hat{\Psi}_{\text{PDW}} = \sum_{\mathbf{r}, \mathbf{a}} t_{\mathbf{q}, \mathbf{a}} (c_{\mathbf{r}, \uparrow} c_{\mathbf{r} + \mathbf{a}, \downarrow} - c_{\mathbf{r}, \downarrow} c_{\mathbf{r} + \mathbf{a}, \uparrow}) e^{i(\mathbf{q} \cdot \mathbf{r} + \mathbf{a}/2)} + \text{H.c}$$



# Charge order (VCA)

- $U = 5$
- $U = 6$
- $U = 8$
- $U = 6, t'' = 0.2$
- $U = 8, t'' = 0.2$

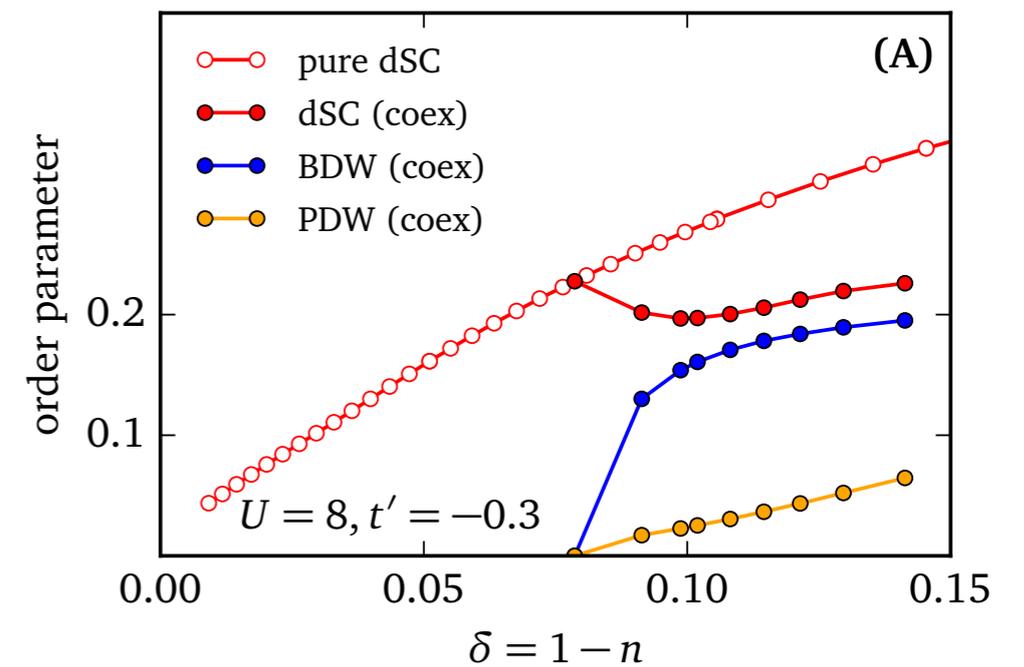
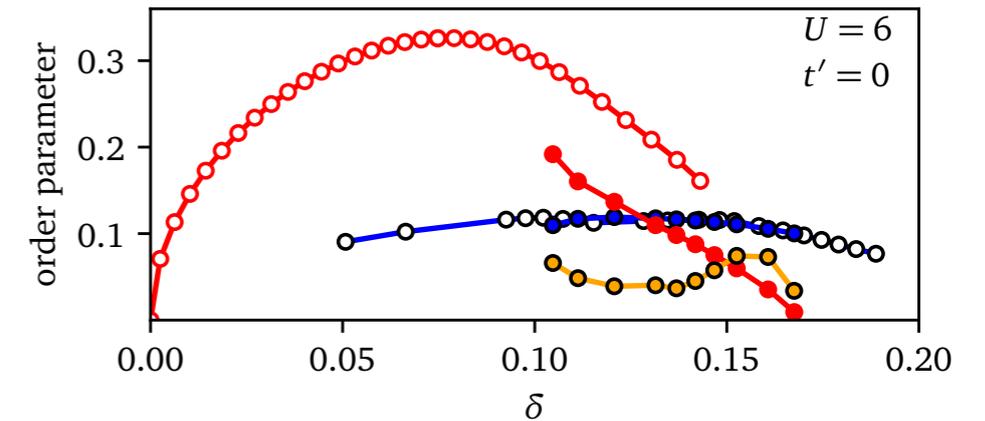
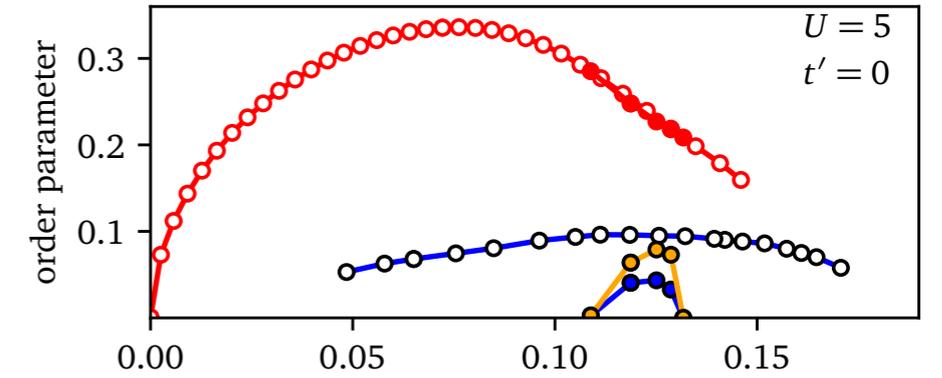
pure bond density wave



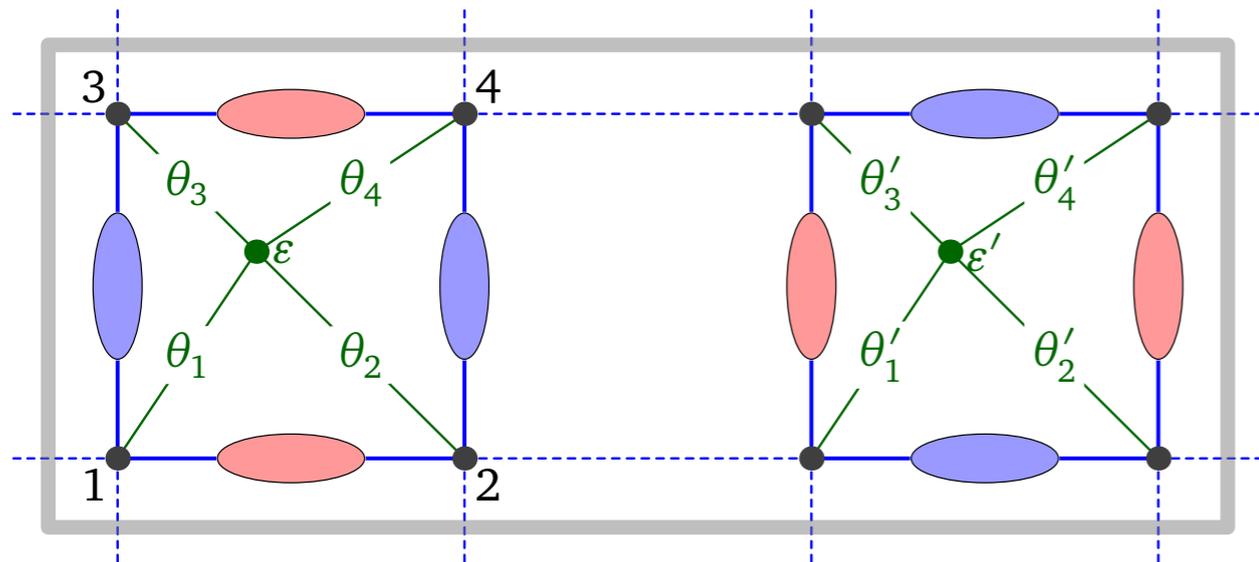
JPL Faye & DS, Phys. Rev. B **95** 115127 (2017)

competition with d-wave superconductivity

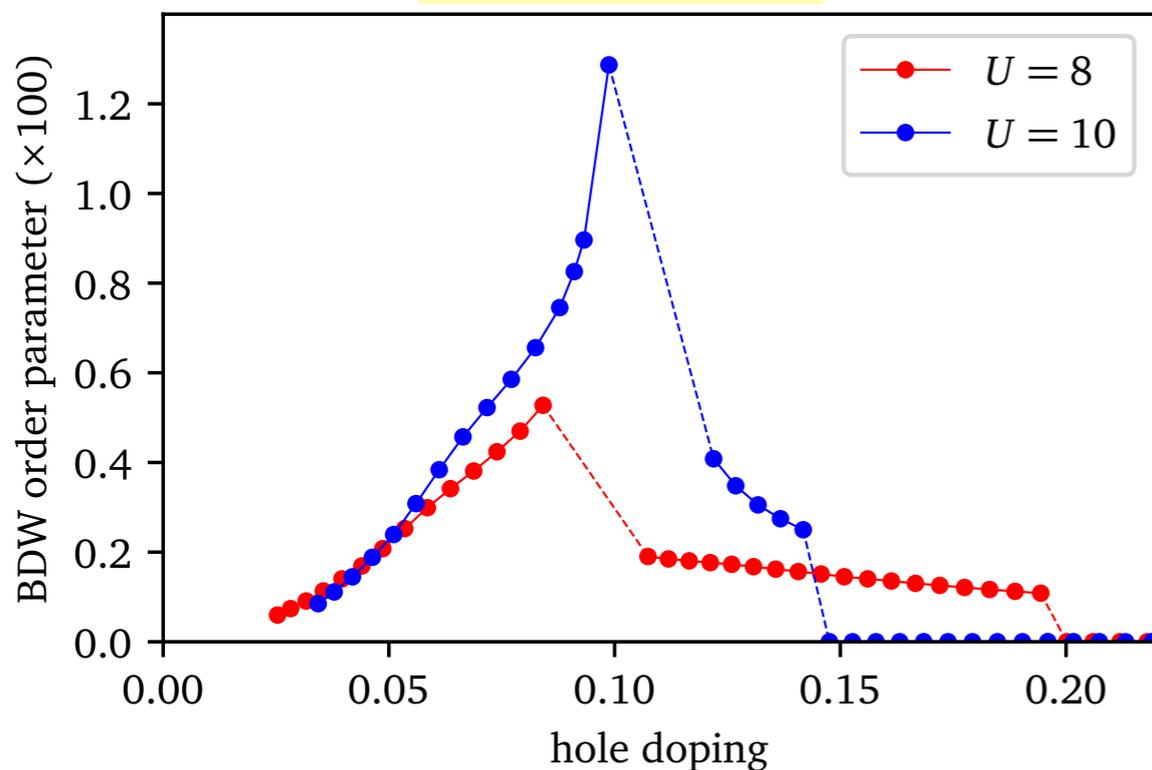
- pure BDW
- BDW (coex)
- pure dSC
- dSC (coex)
- PDW ( $\times 10$ )



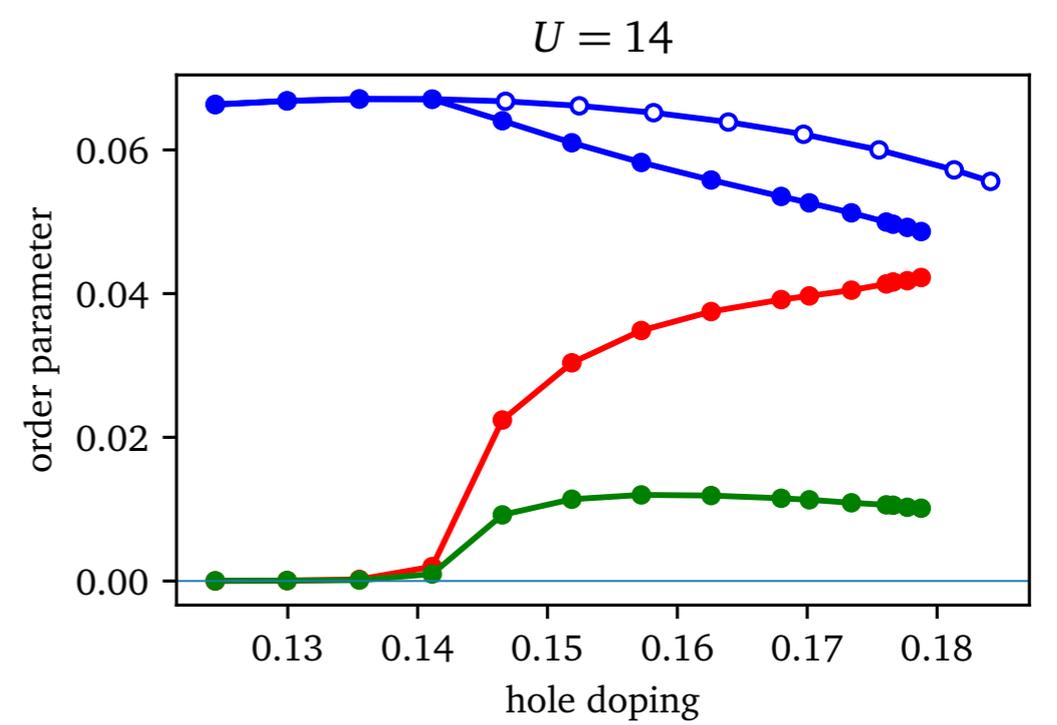
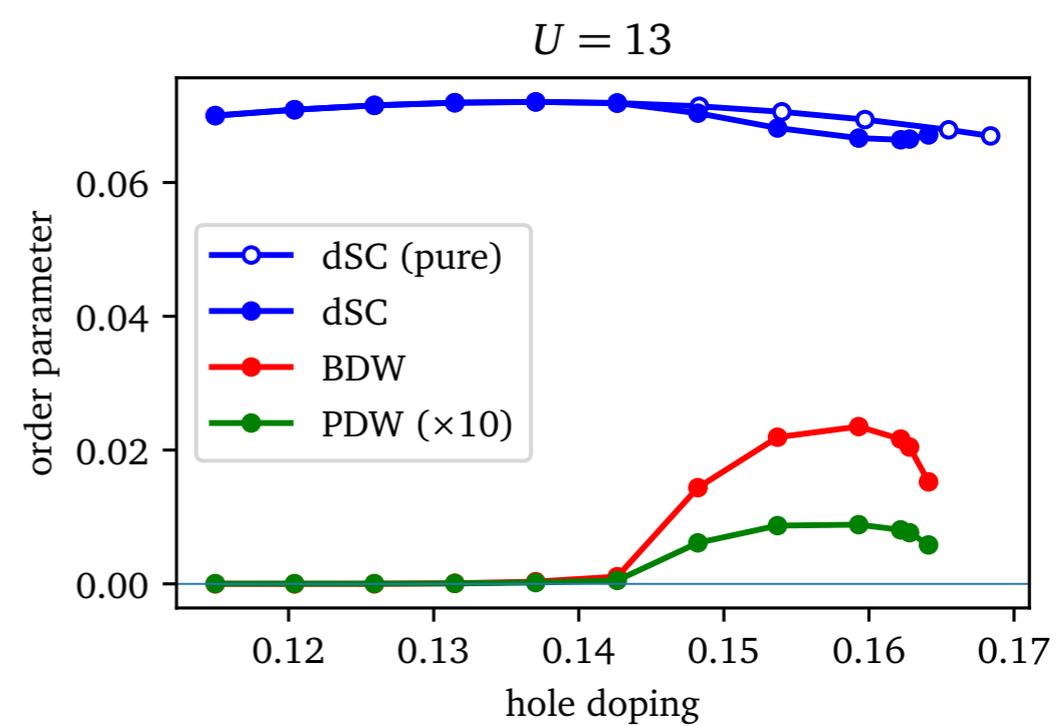
# Charge order (CDMFT)



normal phase



with SC phase



# loop currents

Simon & Varma, PRL **89** 247003 (2002)

current operator along direction  $\mathbf{e}$  :

$$ic_r^\dagger c_{r+\mathbf{e}} - ic_{r+\mathbf{e}}^\dagger c_r$$

$$H \rightarrow H + I\hat{I}$$

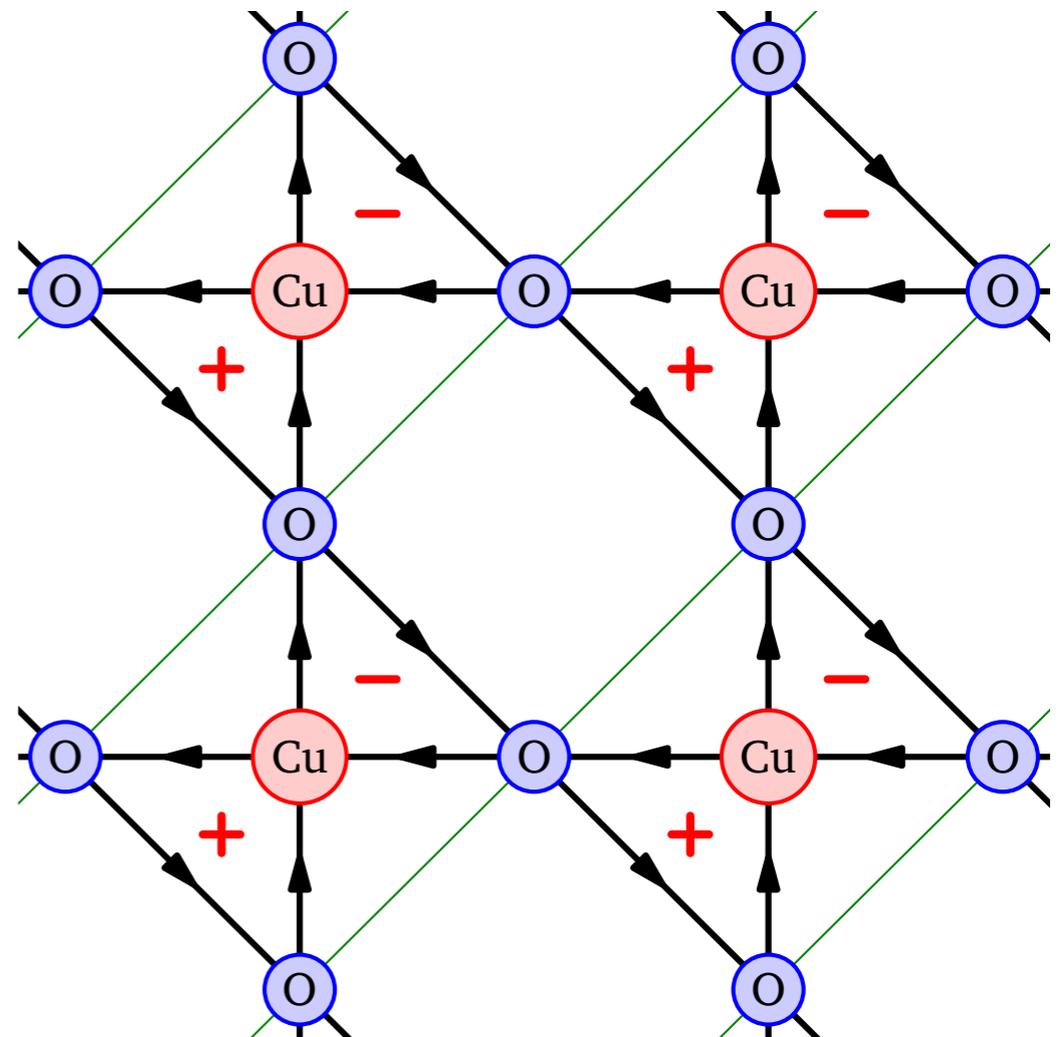
Weiss field  loop current operator

## CDMFT strategy

- Allow complex-valued bath parameters
- Force the system by imposing nonzero  $I$
- solve CDMFT repeatedly while decreasing  $I$
- Measure order parameter  $\langle \hat{I} \rangle$  vs  $I$

Result : order parameter always zero. No loop current order !

But : single set of parameters. more work to be done...



# Conclusions

- Superconductivity in strongly correlated materials can be studied using DMFT-like methods based on small clusters
- There are finite-size effects, and discrete bath effects, but the overall picture is consistent
- Competing orders can also be studied : AFM, charge order
- 1-band & 3-band models capture the physics of cuprates

# Thank you!



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# The Variational Cluster Approximation (VCA)

technical slide

- Add Weiss field(s) to the cluster
- Set their values by optimizing the Potthoff functional

$$\Omega(M) = \Omega'(M) - \int \frac{d\omega}{\pi} \sum_{\tilde{\mathbf{k}}} \ln \det \left[ 1 - \mathbf{V}(\tilde{\mathbf{k}}) \mathbf{G}(i\omega) \right]$$

↖ G.S. energy of cluster
↖ inter cluster hopping

↙ Potthoff functional
↘ cluster Green function

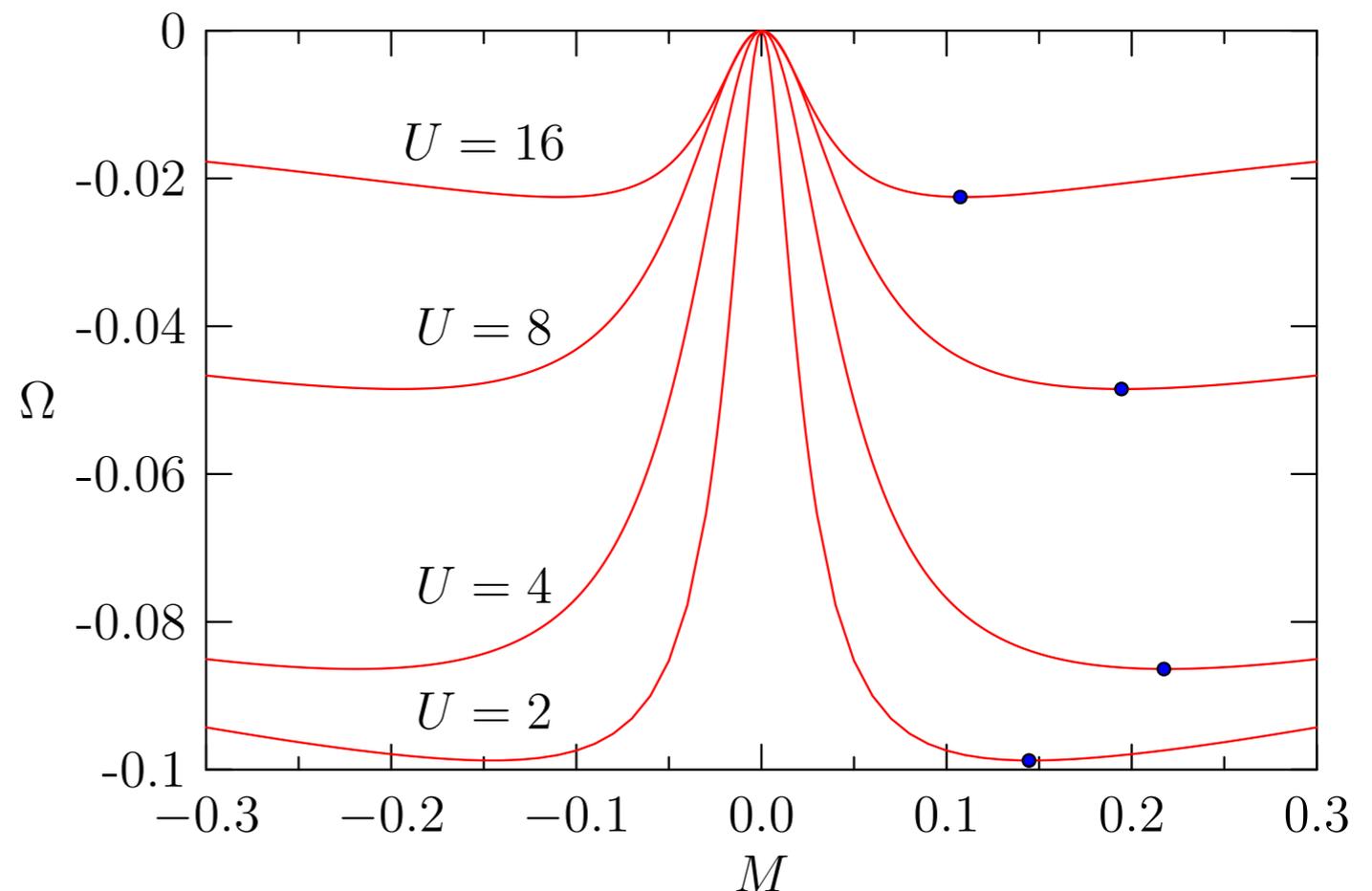
Example : Néel antiferromagnetism

$$H_M = M \hat{M}$$

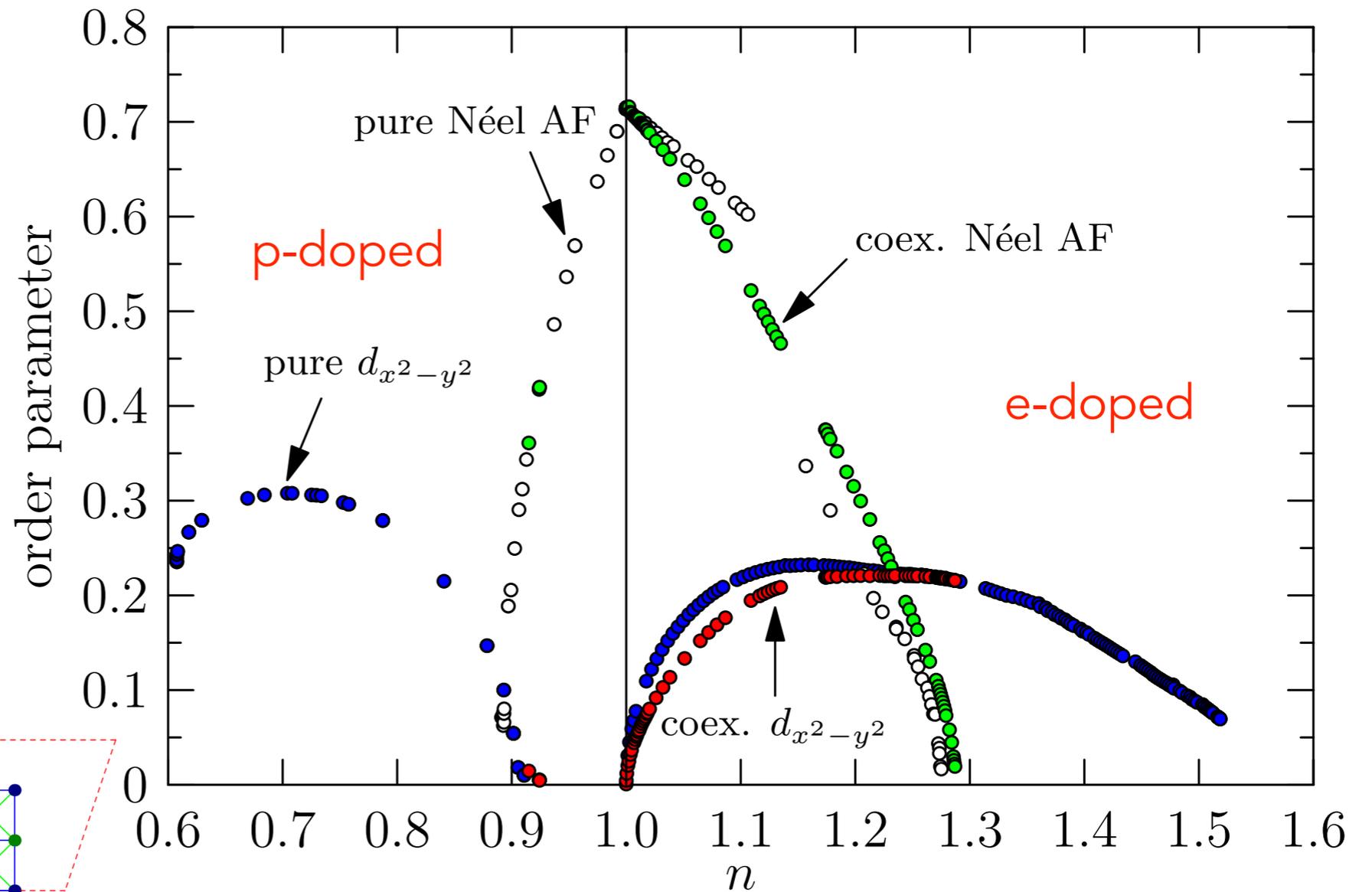
$$\hat{M} = \sum_{\mathbf{r}} e^{i\mathbf{Q}\cdot\mathbf{r}} (n_{\mathbf{r}\uparrow} - n_{\mathbf{r}\downarrow})$$

order parameter :  $\langle \hat{M} \rangle / N_{\text{sites}}$

↖ Weiss field

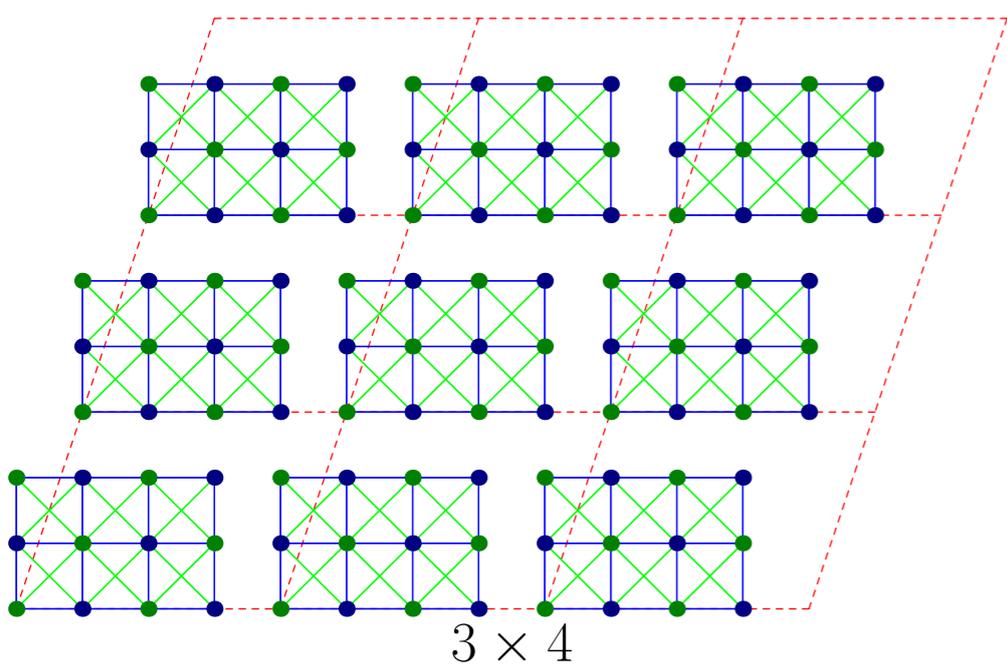


# VCA : phase diagram for YBCO



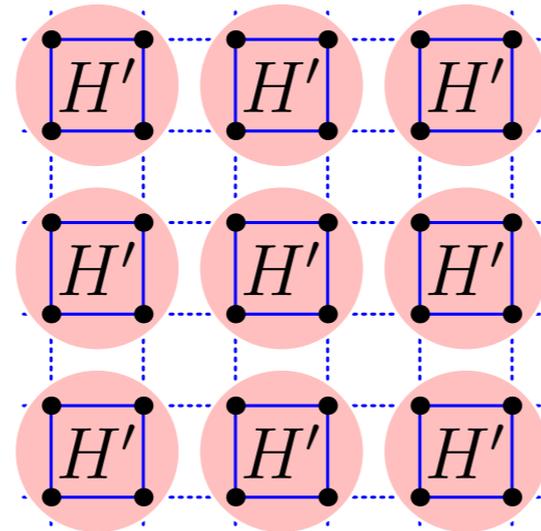
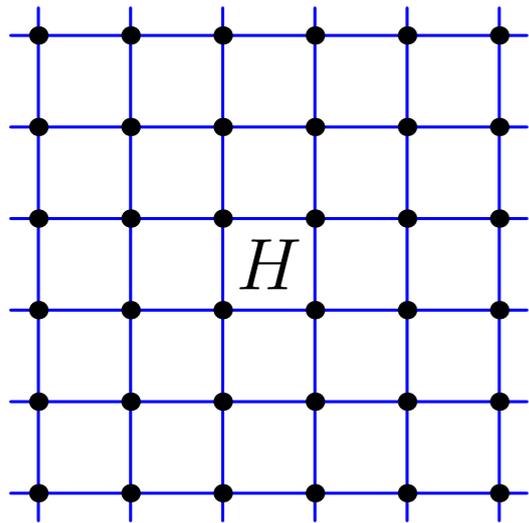
$$t = 1, t' = -0.3, t'' = 0.2$$

12-site cluster



# (Cluster) Dynamical Mean Field Theory

technical slide



- Environment of cluster replaced by uncorrelated « bath »
- Bath parameters determined by self-consistency relation

cluster Hamiltonian

hybridization

bath orbital

$$H' = \sum_{\alpha, \beta} t'_{\alpha\beta} c_{\alpha}^{\dagger} c_{\beta} + U \sum_i n_{i\uparrow} n_{i\downarrow} + \sum_{r, \alpha} \theta_{r\alpha} (c_{\alpha}^{\dagger} a_r + \text{H.c.}) + \sum_r \varepsilon_r a_r^{\dagger} a_r$$

bath energies

Green function :

$$\mathbf{G}'(i\omega_n)^{-1} = i\omega_n - \mathbf{t}' - \mathbf{\Gamma}(i\omega_n) - \mathbf{\Sigma}(i\omega_n)$$

Hybridization function :

$$\Gamma_{\alpha\beta}(i\omega_n) = \sum_r \frac{\theta_{\alpha r} \theta_{\beta r}^*}{i\omega_n - \varepsilon_r}$$

# Adding superconductivity

technical slide

Adding a pairing term  $H_{\text{pair}}$  to the Hamiltonian :

VCA

$$H_{\text{pair}} = \sum_{\alpha, \beta} \Delta_{\alpha\beta} c_{\alpha} c_{\beta} + \text{H.c}$$

Nambu spinor :  $\Psi_{\alpha} = (c_{\alpha}, c_{\alpha}^{\dagger})$

one-body matrix:

$$\begin{pmatrix} \mathbf{t} & \Delta^{\dagger} \\ \Delta & -\mathbf{t}^* \end{pmatrix}$$

CDMFT

cluster ←      → bath

$$H_{\text{pair}} = \sum_{\alpha, r} d_{\alpha r} c_{\alpha} a_r + \sum_{r, r'} \Delta_{r, r'} a_r a_{r'} + \text{H.c}$$

$$\Psi_{\alpha} = (c_{\alpha}, a_r, c_{\alpha}^{\dagger}, a_r^{\dagger})$$

$$\begin{pmatrix} \mathbf{t} & \boldsymbol{\theta} & 0 & \mathbf{d}^{\dagger} \\ \boldsymbol{\theta}^{\dagger} & \boldsymbol{\varepsilon} & -\mathbf{d}^* & \Delta^{\dagger} \\ 0 & -\mathbf{d}^T & -\mathbf{t}^* & -\boldsymbol{\theta}^* \\ \mathbf{d} & \Delta & -\boldsymbol{\theta}^T & -\boldsymbol{\varepsilon}^* \end{pmatrix}$$

# The CDMFT self-consistency loop

