

Search for Higgs boson decays
to BSM light bosons in four-lepton events
with the ATLAS detector in 2015-6 and -17

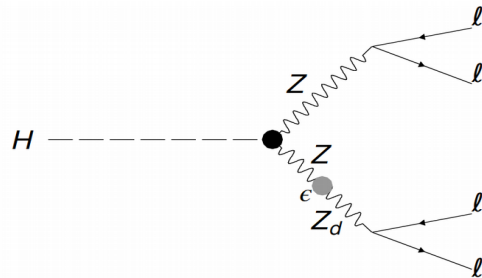
CAP Congress
June 12th, 2018

Justin Chiu



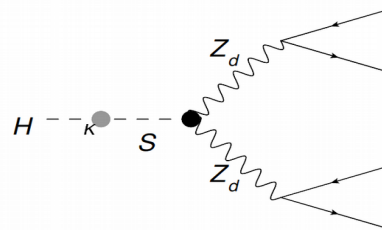
**University
of Victoria**

- This is a search for Higgs decays to four leptons via two intermediate BSM particles of similar mass
- Existing measurements constrain the branching ratio of the Higgs to BSM particles to be less than ~30% (ATLAS & CMS arXiv:1606.02266)
- A small Higgs coupling to a new light state could open up sizable new decay modes
- New particles could couple to the Higgs and provide a “portal” to a hidden (“dark”) sector or extended Higgs sector
- Two benchmark models considered in this analysis (Curtin et. al arXiv:1312.4992 and 1412.0018):
 - Higgs decay to four leptons via one or two intermediate U(1) dark sector particles Z_d
 - Higgs decay to four leptons via two pseudoscalar particles $a = \cos \theta_a S_I + \sin \theta_a A$



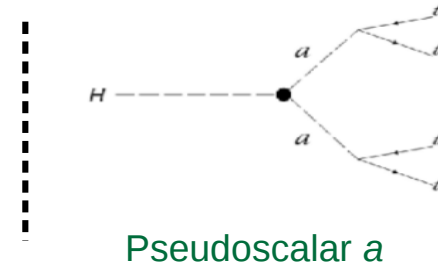
Hypercharge portal

Kinetic mixing ϵ



Higgs portal

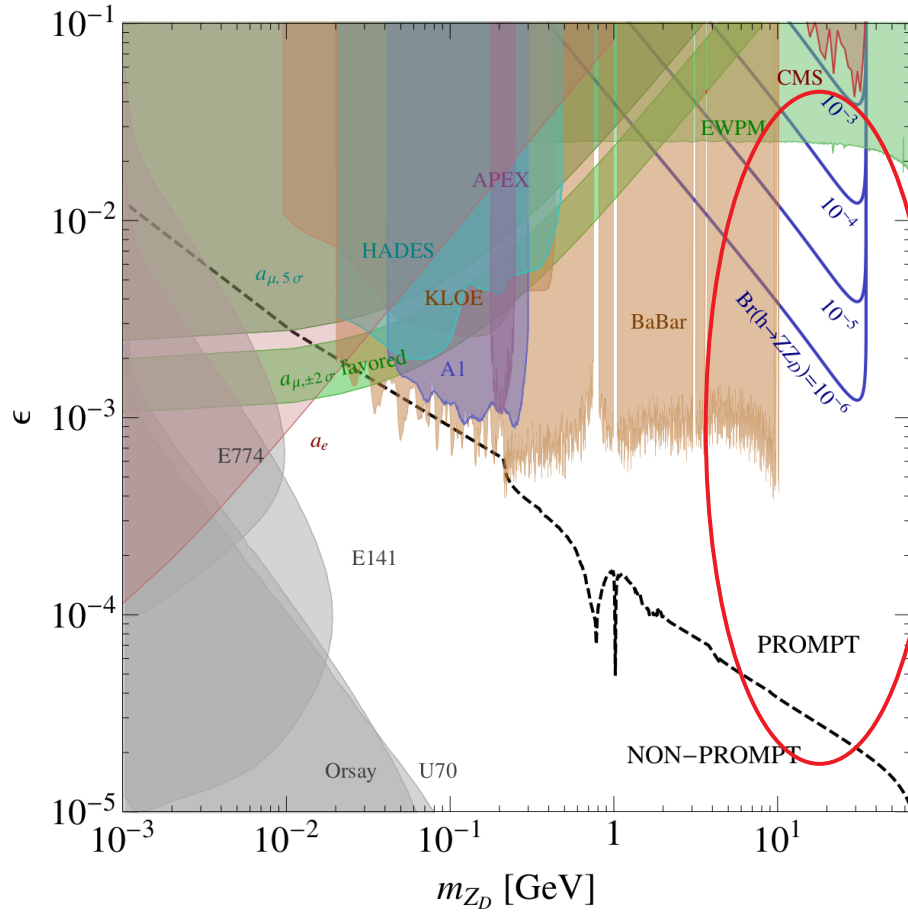
Higgs portal coupling κ
Dark Higgs S



Pseudoscalar a

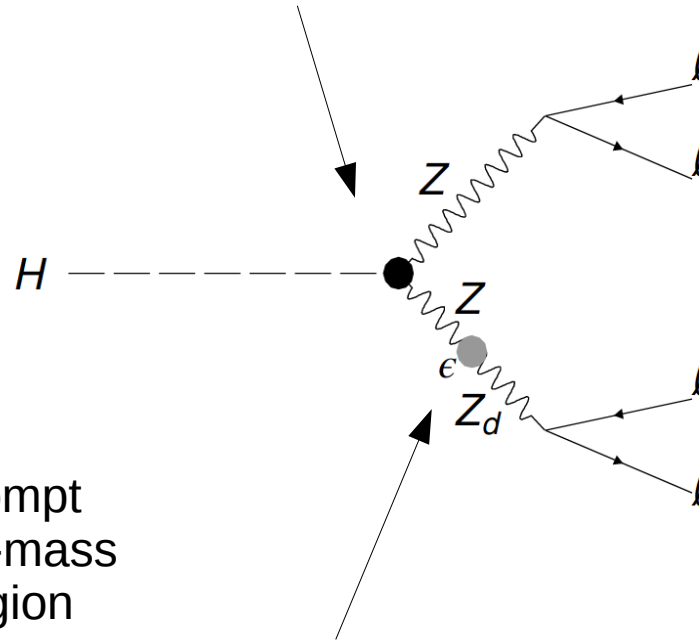
(eigenstate of doublet and singlet in 2HDM+S)

Mixing angle $\theta_a \ll 1$ $S_I = \text{Im}(S)$



Curtin et al. ([arXiv:1312.4992](https://arxiv.org/abs/1312.4992))

Hypercharge scenario

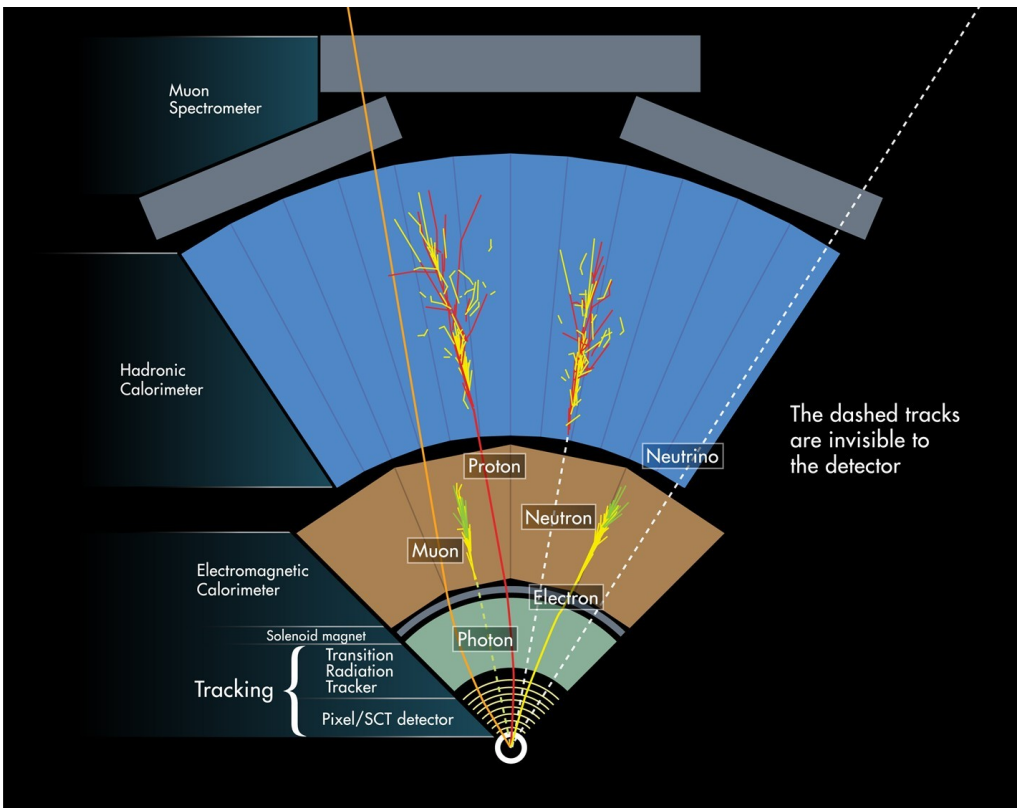


prompt
high-mass
region

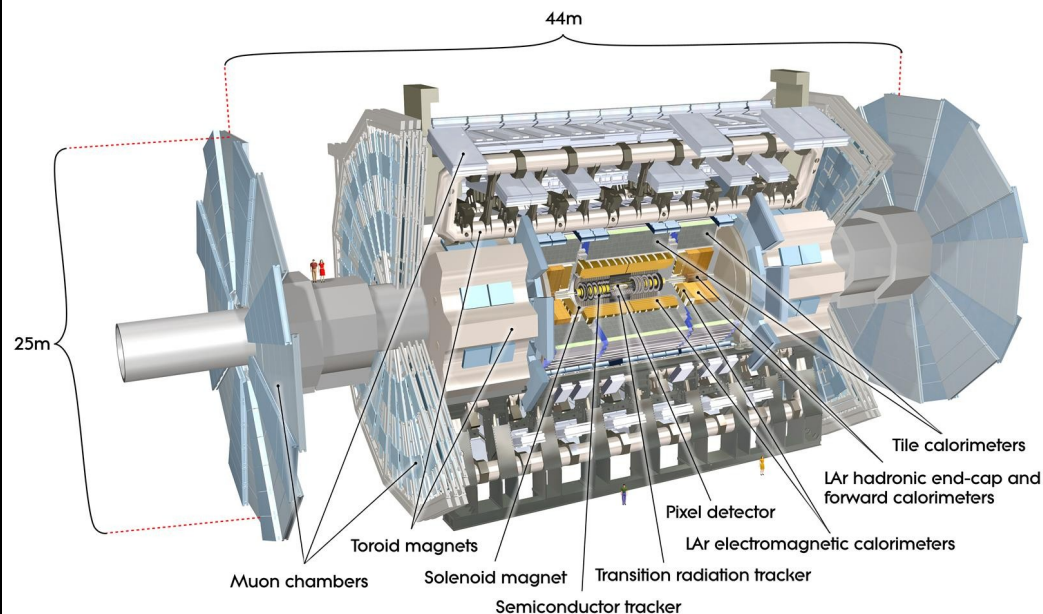
Kinetic mixing parameter ϵ
controls coupling strength between
 Z_d and SM particles

(which in turn determines Z_d lifetime)

Branching ratios determined by gauge coupling



(source: [ATLAS](#))



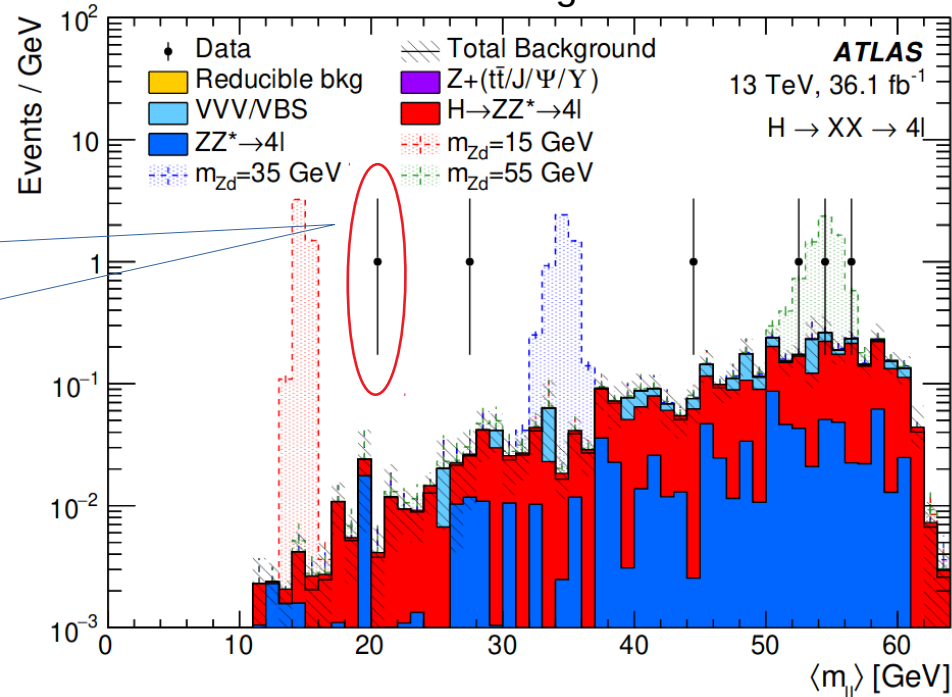
(source: [ATLAS](#))

Look for a light resonance in a channel similar to the “golden” channel of the Higgs ($H \rightarrow ZZ^* \rightarrow 4l$)

- Paper on 2015-6 data covers three analyses ([ATLAS arXiv:1802.0338](#)), where $X = Z_d$ or a :
 - ZX: $H \rightarrow ZX \rightarrow 4l$ ($15 \text{ GeV} < m_X < 55 \text{ GeV}$) ($l = e, \mu$)
 - “Low-mass”: $H \rightarrow XX \rightarrow 4\mu$ ($1 \text{ GeV} < m_X < 15 \text{ GeV}$)
 - “High-mass”: $H \rightarrow XX \rightarrow 4l$ ($15 \text{ GeV} < m_X < 60 \text{ GeV}$) ($l = e, \mu$)

focus today is on high-mass analysis

Published 2015-6 “high-mass” result

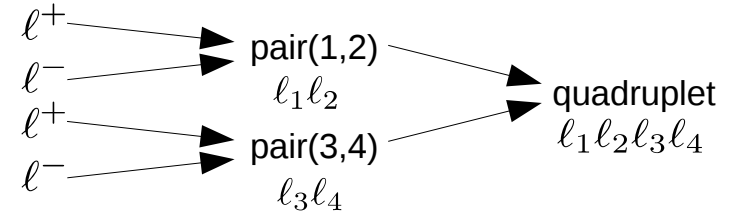


For this bin:
3.2 σ local
1.9 σ global

(N.B. only 1 event!)

Lepton quadruplet formation and selection

- Each lepton must be responsible for firing at least one trigger
- In each event, form quadruplets consisting of two same-flavour opposite sign lepton pairs: “1, 2” & “3, 4”
- Three leading-pT leptons must have $p_T > 20, 15, \text{ and } 10 \text{ GeV}$
- $\Delta R(l, l') > 0.10 \text{ (} 0.20 \text{)}$ for all pairings of same-flavour (different-flavour) leptons; $\Delta R \equiv \sqrt{\Delta(\eta)^2 + \Delta(\phi)^2}$



Quadruplet ranking

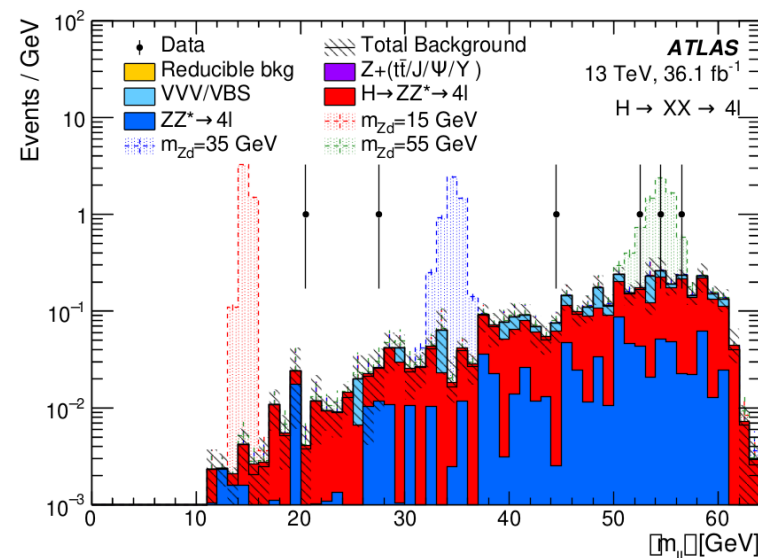
- Select quadruplet with smallest difference in mass between lepton pairs: $\Delta m_{\ell\ell} \equiv |m_{12} - m_{34}|$

Event selection

- Higgs window: $115 \text{ GeV} < m_{4\ell} < 130 \text{ GeV}$
- Compatibility of lepton pair masses: $m_{34}/m_{12} > 0.85$
- Z veto on lepton pairs: $10 \text{ GeV} < m_{12,34} < 64 \text{ GeV}$
- Additional Z veto on alternative pairings (only for 4e and 4 μ quadruplets): $5 \text{ GeV} < m_{14,32} < 75 \text{ GeV}$
- Quarkonia veto: reject event if $(m_{J/\Psi} - 0.25 \text{ GeV}) < m_{12,34,14,32} < (m_{\Psi(2S)} + 0.30 \text{ GeV}) \parallel (m_{\Upsilon(1S)} - 0.70 \text{ GeV}) < m_{12,34,14,32} < (m_{\Upsilon(3S)} + 0.75 \text{ GeV})$

- We take all background predictions from simulation, as in previous iteration ([ATLAS arXiv:1802.03388](#))
- Dominant backgrounds are $H \rightarrow ZZ^*$ and non-resonant ZZ^* :
 - $H \rightarrow ZZ^* \rightarrow 4\ell$: 63%
 - $ZZ^* \rightarrow 4\ell$: 19%
 - Triboson production (VVV): 17%
 - $Z + t\bar{t}, J/\Psi, \text{ or } \Upsilon \rightarrow 4\ell$: $\sim 1\%$
 - Reducible backgrounds ($Z + \text{jets}, t\bar{t}$): $\sim 1\%$
 - working on in-situ estimate for next iteration

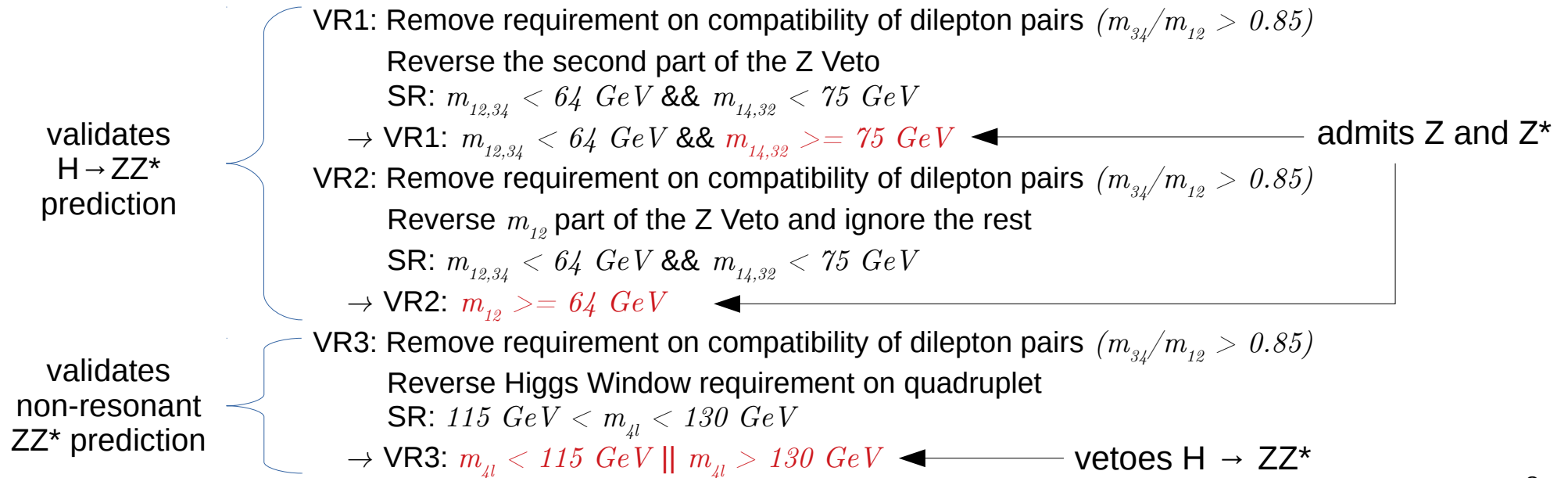
| Process | Yield |
|---|-------------------|
| $ZZ^* \rightarrow 4\ell$ | 0.8 ± 0.1 |
| $H \rightarrow ZZ^* \rightarrow 4\ell$ | 2.6 ± 0.3 |
| VVV/VBS | 0.51 ± 0.18 |
| $Z + (t\bar{t}/J/\Psi) \rightarrow 4\ell$ | 0.004 ± 0.004 |
| Reducible Background | Negligible |
| Total | 3.9 ± 0.3 |
| Data | 6 |



(a) Signal region $\langle m_{\ell\ell} \rangle$ distribution

Previous iteration

- Moving towards a result with 2017 data
- Follow the same approach as the past iteration
- Validate our $H \rightarrow ZZ^*$ and non-resonant ZZ^* background predictions by considering three validation regions:



Validation region 1 (2015-6+7)

VR1: Remove requirement on compatibility of dilepton pairs ($m_{34}/m_{12} > 0.85$)

Reverse the second part of the Z Veto

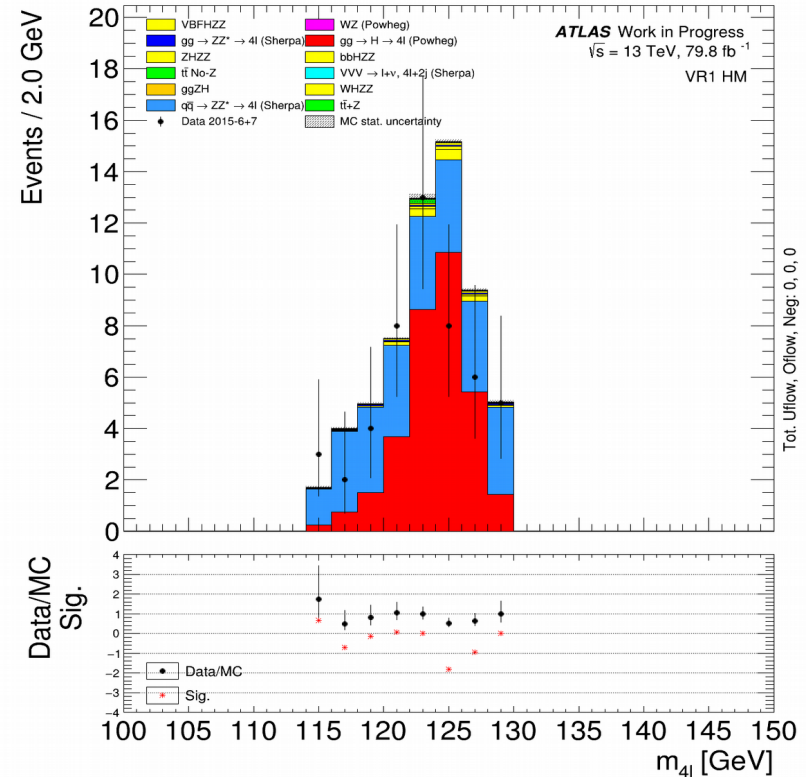
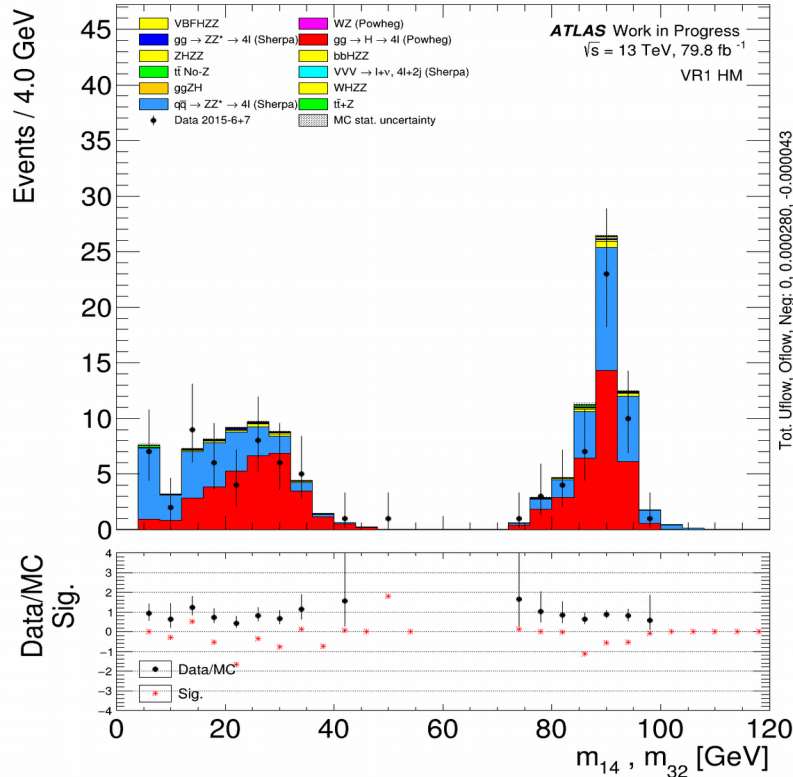
SR: $m_{12,34} < 64 \text{ GeV} \ \&\& \ m_{14,32} < 75 \text{ GeV}$

→ VR1: $m_{12,34} < 64 \text{ GeV} \ \&\& \ m_{14,32} \geq 75 \text{ GeV}$



■ $gg \rightarrow H \rightarrow 4l$ (Powheg)

■ $q\bar{q} \rightarrow ZZ^* \rightarrow 4l$ (Sherpa)



Validation region 2 (2015-6+7)

VR2: Remove requirement on compatibility of dilepton pairs ($m_{34}/m_{12} > 0.85$)

Reverse m_{12} part of the Z Veto and ignore the rest

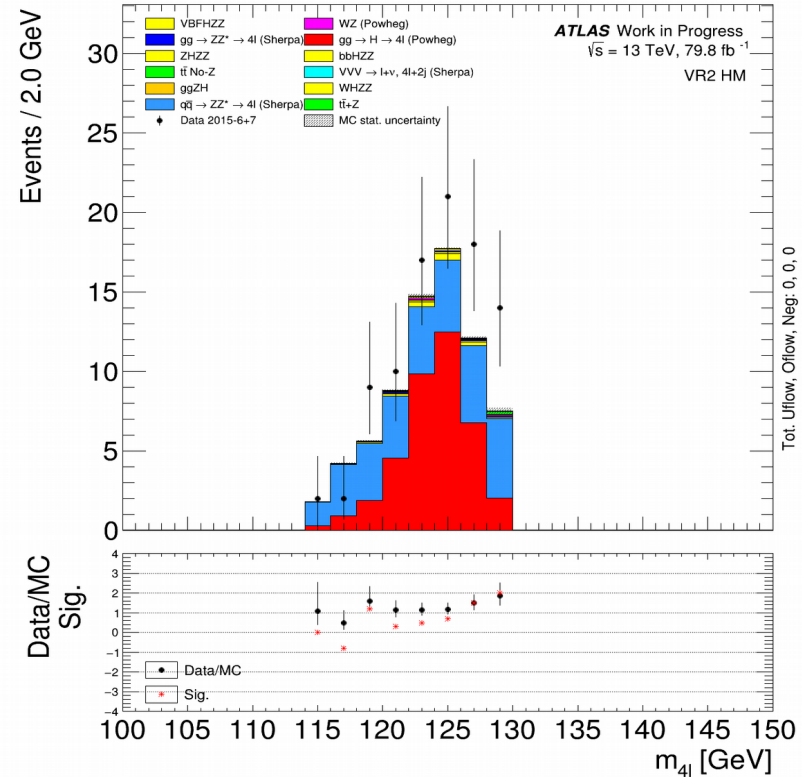
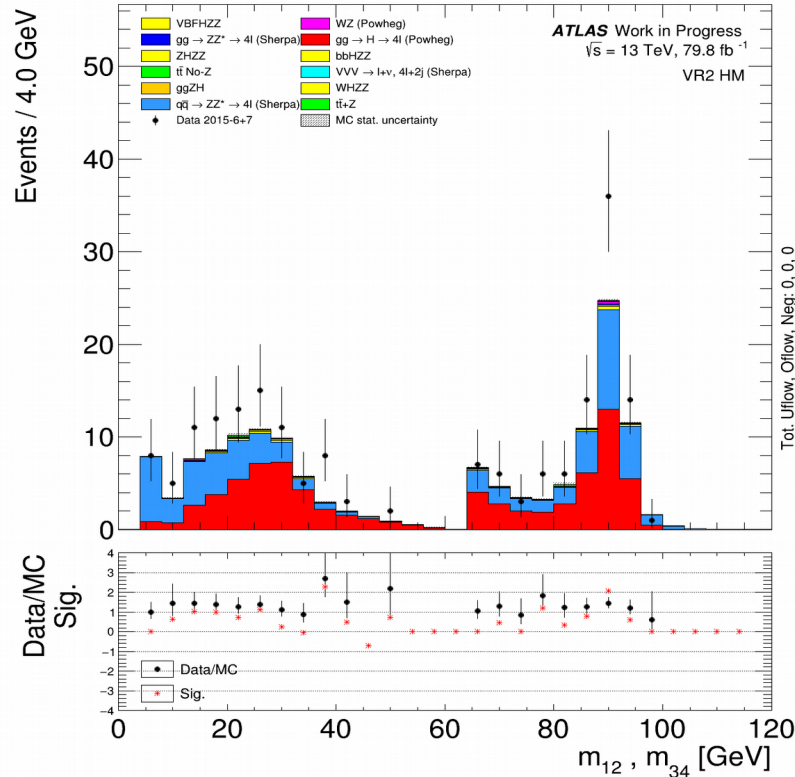
SR: $m_{12,34} < 64 \text{ GeV} \ \&\& \ m_{14,32} < 75 \text{ GeV}$

→ VR2: $m_{12} \geq 64 \text{ GeV}$



■ $gg \rightarrow H \rightarrow 4l$ (Powheg)

■ $q\bar{q} \rightarrow ZZ^* \rightarrow 4l$ (Sherpa)



Validation region 3 (2015-6+7)

VR3: Remove requirement on compatibility of dilepton pairs ($m_{34}/m_{12} > 0.85$)

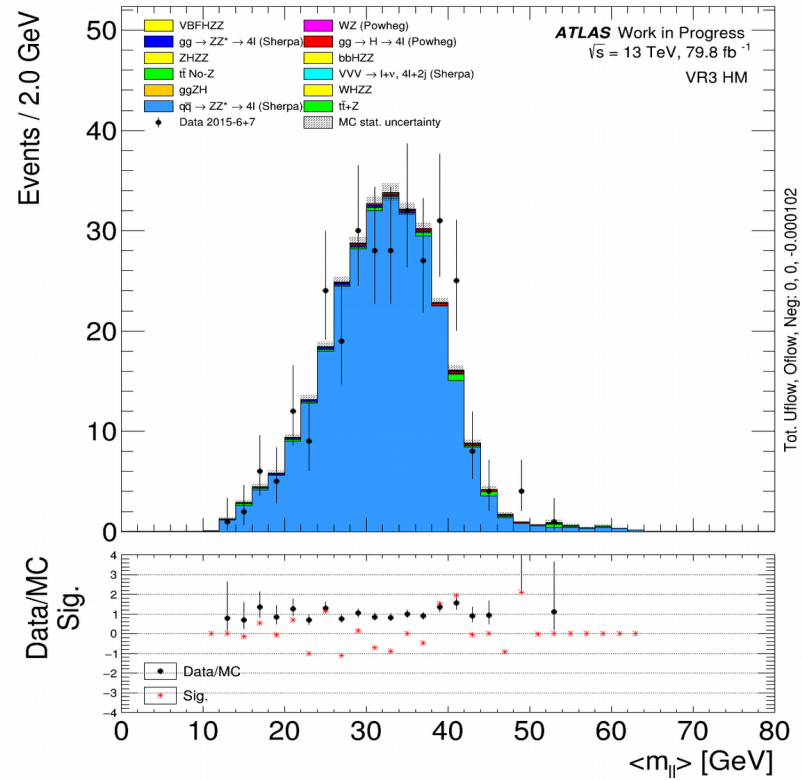
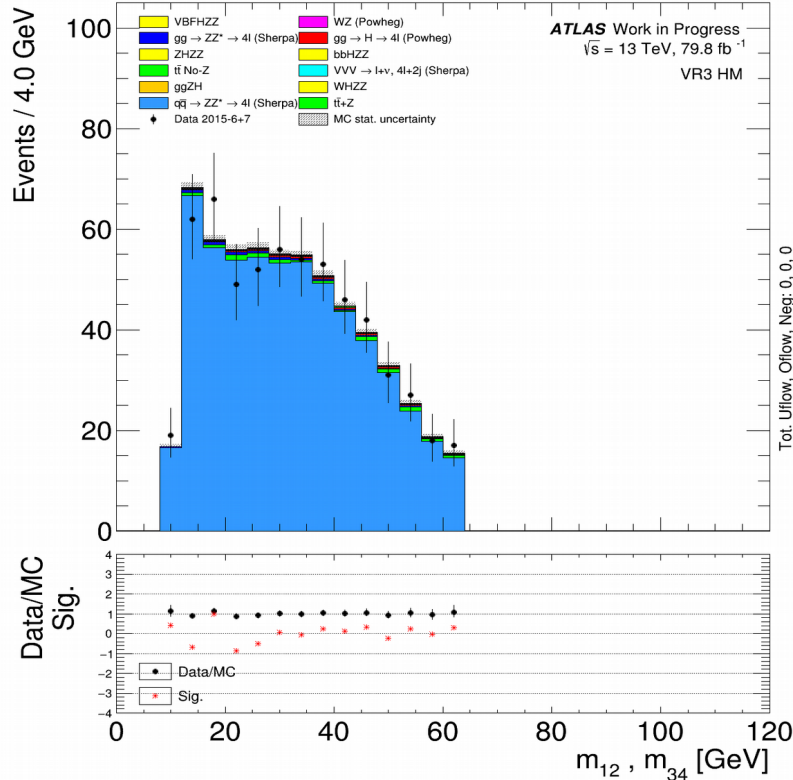
Reverse Higgs Window requirement on quadruplet

SR: $115 \text{ GeV} < m_{4l} < 130 \text{ GeV}$

→ VR3: $m_{4l} < 115 \text{ GeV} \parallel m_{4l} > 130 \text{ GeV}$



$q\bar{q} \rightarrow ZZ^* \rightarrow 4l$ (Sherpa)



- Working towards a result encompassing the full (2015-8) Run-2 dataset
- Other efforts for full Run-2 result not mentioned:
 - In-situ estimate for reducible backgrounds (statistically limited)
 - Extend search to cover wider range in Dark Higgs mass and Z_d
 - Study sensitivity to non-prompt signal
 - Two new channels: $H \rightarrow aa \rightarrow 2\mu 2\tau$ (Run-1) and $H \rightarrow aa \rightarrow 4\tau$ (new)

Thanks!
Questions?

Analysis selection (from paper)

m_{12} is the dilepton pair closest to Z mass and m_{34} is the other di-lepton pair
 m_{14} and m_{32} are the alternative same-flavour opposite sign pairings

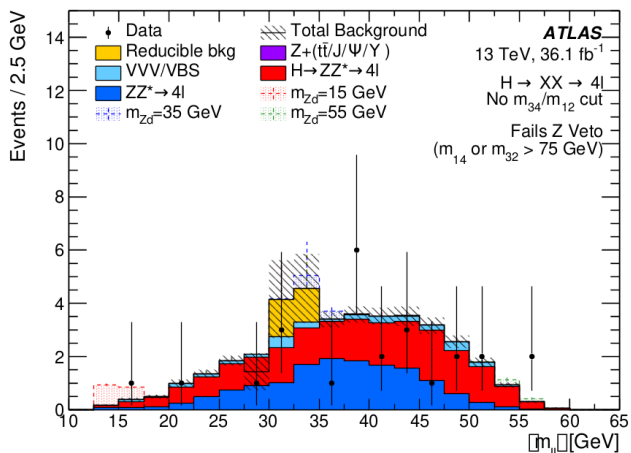
| | $H \rightarrow ZX \rightarrow 4\ell$ (15 GeV < m_X < 55 GeV) | $H \rightarrow XX \rightarrow 4\ell$ (15 GeV < m_X < 60 GeV) | $H \rightarrow XX \rightarrow 4\mu$ (1 GeV < m_X < 15 GeV) |
|----------------------|--|---|---|
| QUADRUPLET SELECTION | <ul style="list-style-type: none"> - Require at least one quadruplet of leptons consisting of two pairs of same-flavour opposite-sign leptons - Three leading-p_T leptons satisfying $p_T > 20$ GeV, 15 GeV, 10 GeV - At least three muons are required to be reconstructed by combining ID and MS tracks in the 4μ channel | | |
| | <ul style="list-style-type: none"> - Select best quadruplet (per channel) to be the one with the (sub)leading dilepton mass (second) closest to the Z mass - 50 GeV < m_{12} < 106 GeV - 12 GeV < m_{34} < 115 GeV - $m_{12,34,14,32} > 5$ GeV | Leptons in the quadruplet are responsible for firing at least one trigger. In the case of multi-lepton triggers, all leptons of the trigger must match to leptons in the quadruplet | |
| | $\Delta R(\ell, \ell') > 0.10$ (0.20) for same-flavour (different-flavour) leptons in the quadruplet | | - |
| QUADRUPLET RANKING | Select first surviving quadruplet from channels, in the order: 4μ , $2e2\mu$, $2\mu2e$, $4e$ | Select quadruplet with smallest $\Delta m_{\ell\ell} = m_{12} - m_{34} $ | |
| EVENT SELECTION | 115 GeV < $m_{4\ell} < 130$ GeV | 120 GeV < $m_{4\ell} < 130$ GeV | |
| | | $m_{34}/m_{12} > 0.85$ | |
| | Reject event if: $(m_{J/\psi} - 0.25 \text{ GeV}) < m_{12,34,14,32} < (m_{\Psi(2S)} + 0.30 \text{ GeV})$, or $(m_{\Upsilon(1S)} - 0.70 \text{ GeV}) < m_{12,34,14,32} < (m_{\Upsilon(3S)} + 0.75 \text{ GeV})$ | | |
| | 10 GeV < $m_{12,34} < 64$ GeV 4e and 4μ channels: 5 GeV < $m_{14,32} < 75$ GeV | 0.88 GeV < $m_{12,34} < 20$ GeV No restriction on alternative pairing | |

(loose leptons, details in backup p30)

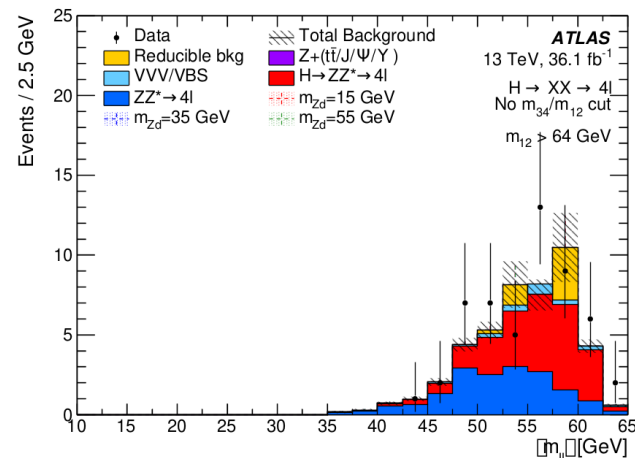


High-mass selection is similar to HZZ4ℓ selection – biggest difference is quadruplet selection

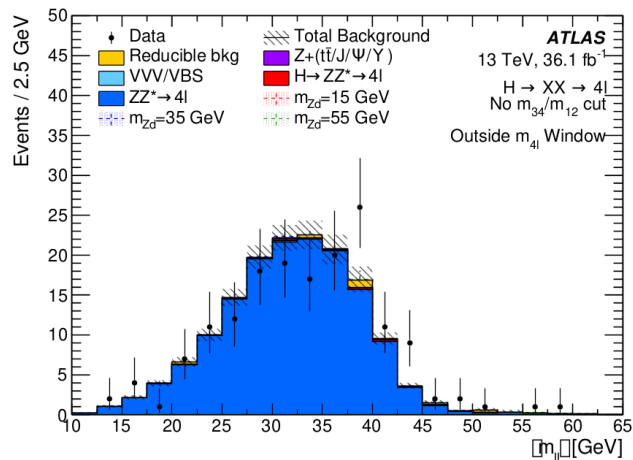
Plots from last paper



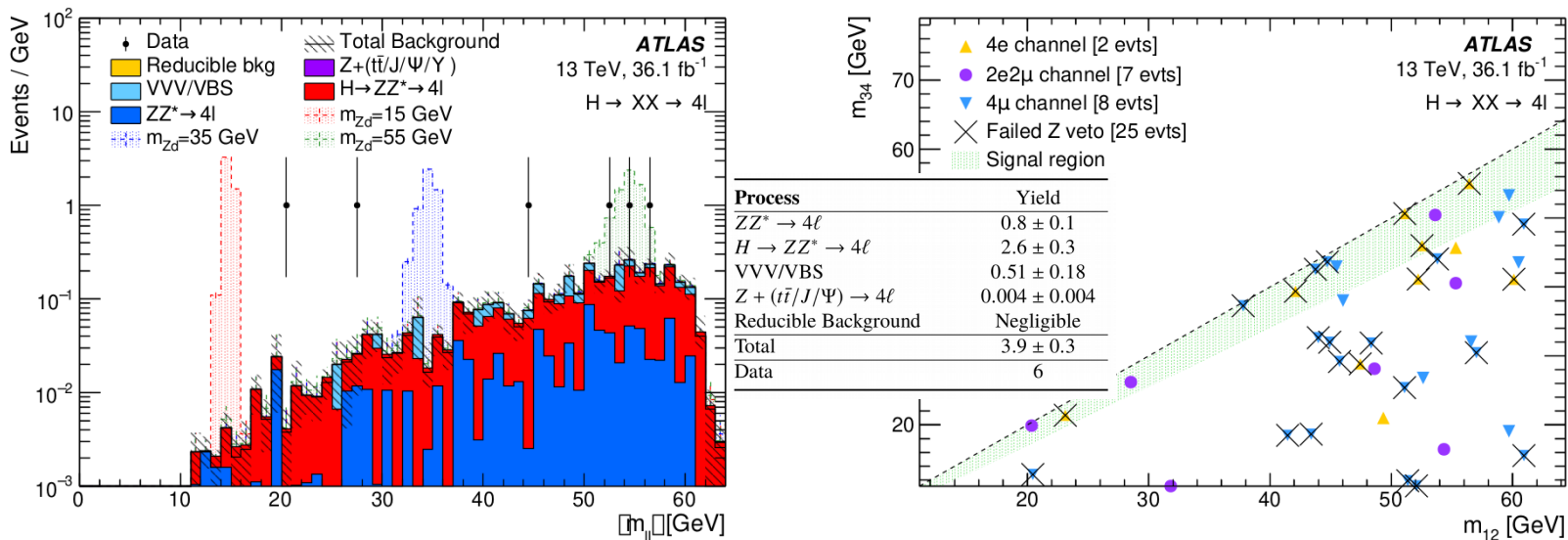
(a) VR1



(b) VR2



(c) VR3



(a) Signal region $\langle m_{\ell\ell} \rangle$ distribution

(b) m_{34} vs m_{12} distribution

Figure 4: Distribution of (a) $\langle m_{\ell\ell} \rangle = \frac{1}{2}(m_{12} + m_{34})$ and (b) m_{34} vs m_{12} , for events selected in the $H \rightarrow XX \rightarrow 4\ell$ ($15 < m_X < 60$ GeV) analysis. The example signal distributions in (a) correspond to the expected yield normalized with $\sigma(pp \rightarrow H \rightarrow Z_d Z_d \rightarrow 4\ell) = \frac{1}{10} \sigma_{\text{SM}}(pp \rightarrow H \rightarrow ZZ^* \rightarrow 4\ell)$. The crossed-through points in (b) fail the Z Veto. The events outside the (shaded green) signal region in figure (b) are events that fail the $m_{34}/m_{12} > 0.85$ requirement. The diagonal dashed line marks where $m_{12} = m_{34}$, and in this range of dilepton masses all events will have $m_{34} < m_{12}$.

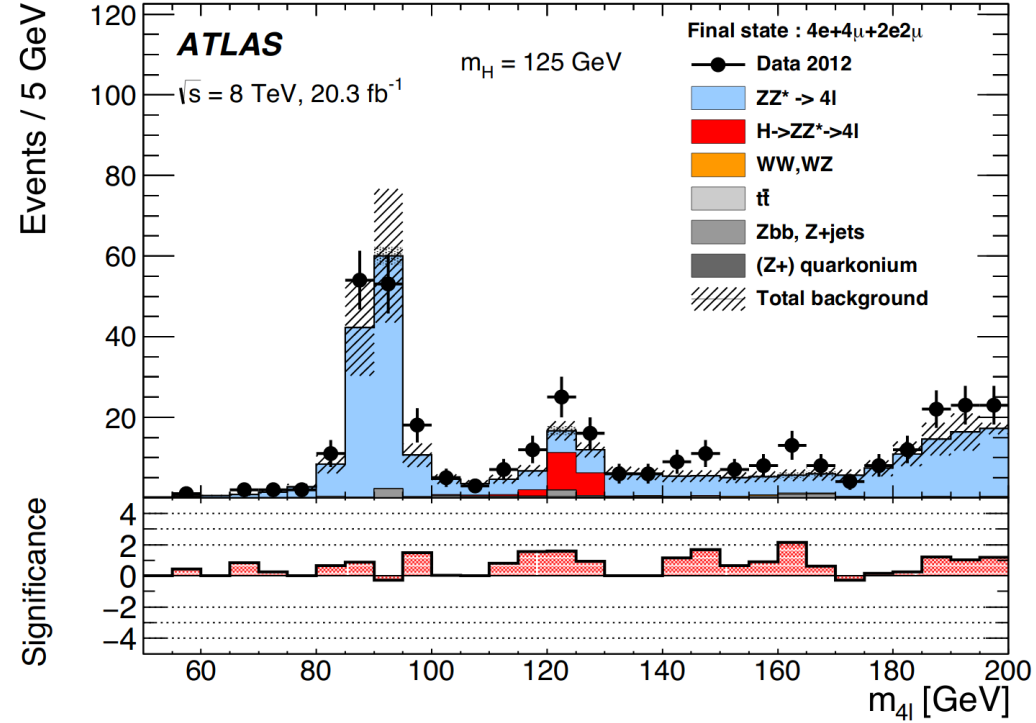
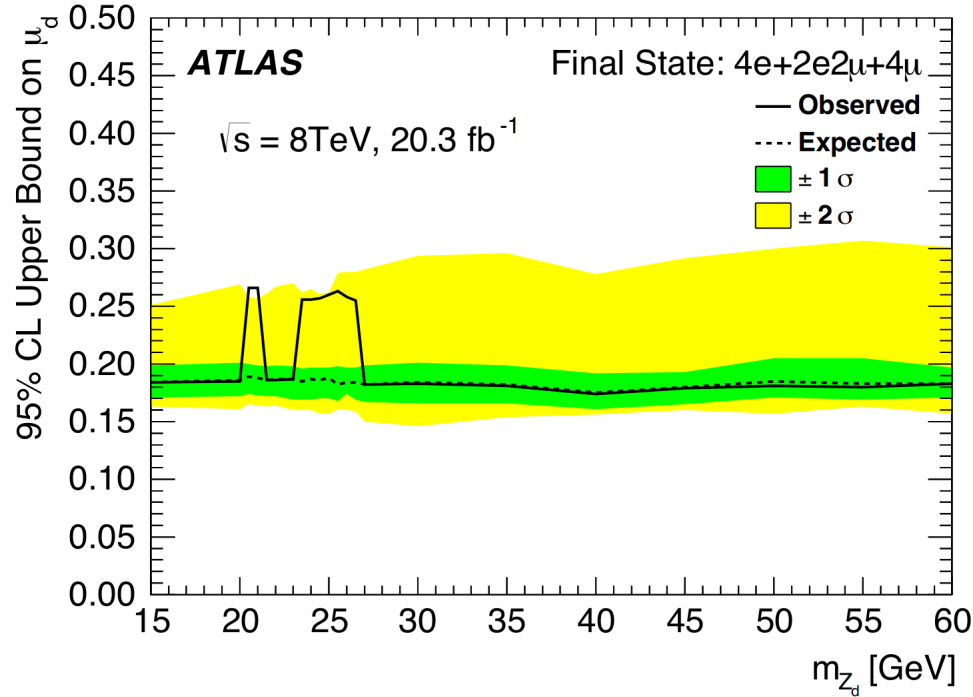


Figure 14: The 95% confidence level upper bound on the signal strength $\mu_d = \frac{\sigma \times \text{BR}(H \rightarrow Z_d Z_d \rightarrow 4\ell)}{[\sigma \times \text{BR}(H \rightarrow ZZ^* \rightarrow 4\ell)]_{\text{SM}}}$ of $H \rightarrow Z_d Z_d \rightarrow 4\ell$ in the combined $4e + 2e2\mu + 4\mu$ final state, for $m_H = 125$ GeV. The $\pm 1\sigma$ and $\pm 2\sigma$ expected exclusion regions are indicated in green and yellow, respectively.

| Couplings | I | II | III (Lepton specific) | IV (Flipped) |
|-----------------|----------------------------|---|-----------------------------|-----------------------------|
| g_{hVV} | $\sin(\beta - \alpha)$ | $\sin(\beta - \alpha)$ | $\sin(\beta - \alpha)$ | $\sin(\beta - \alpha)$ |
| $g_{ht\bar{t}}$ | $\cos \alpha / \sin \beta$ | $\cos \alpha / \sin \beta$ | $\cos \alpha / \sin \beta$ | $\cos \alpha / \sin \beta$ |
| h | $g_{hb\bar{b}}$ | $\cos \alpha / \sin \beta - \sin \alpha / \cos \beta$ | $\cos \alpha / \sin \beta$ | $-\sin \alpha / \cos \beta$ |
| | $g_{h\tau\bar{\tau}}$ | $\cos \alpha / \sin \beta - \sin \alpha / \cos \beta$ | $-\sin \alpha / \cos \beta$ | $\cos \alpha / \sin \beta$ |
| | g_{H^0VV} | $\cos(\beta - \alpha)$ | $\cos(\beta - \alpha)$ | $\cos(\beta - \alpha)$ |
| | $g_{H^0t\bar{t}}$ | $\sin \alpha / \sin \beta$ | $\sin \alpha / \sin \beta$ | $\sin \alpha / \sin \beta$ |
| H^0 | $g_{H^0b\bar{b}}$ | $\sin \alpha / \sin \beta$ | $\cos \alpha / \cos \beta$ | $\cos \alpha / \cos \beta$ |
| | $g_{H^0\tau\bar{\tau}}$ | $\sin \alpha / \sin \beta$ | $\cos \alpha / \cos \beta$ | $\sin \alpha / \sin \beta$ |
| | g_{AVV} | 0 | 0 | 0 |
| A | $g_{At\bar{t}}$ | $\cot \beta$ | $\cot \beta$ | $\cot \beta$ |
| | $g_{Ab\bar{b}}$ | $-\cot \beta$ | $\tan \beta$ | $-\cot \beta$ |
| | $g_{A\tau\bar{\tau}}$ | $-\cot \beta$ | $\tan \beta$ | $-\cot \beta$ |

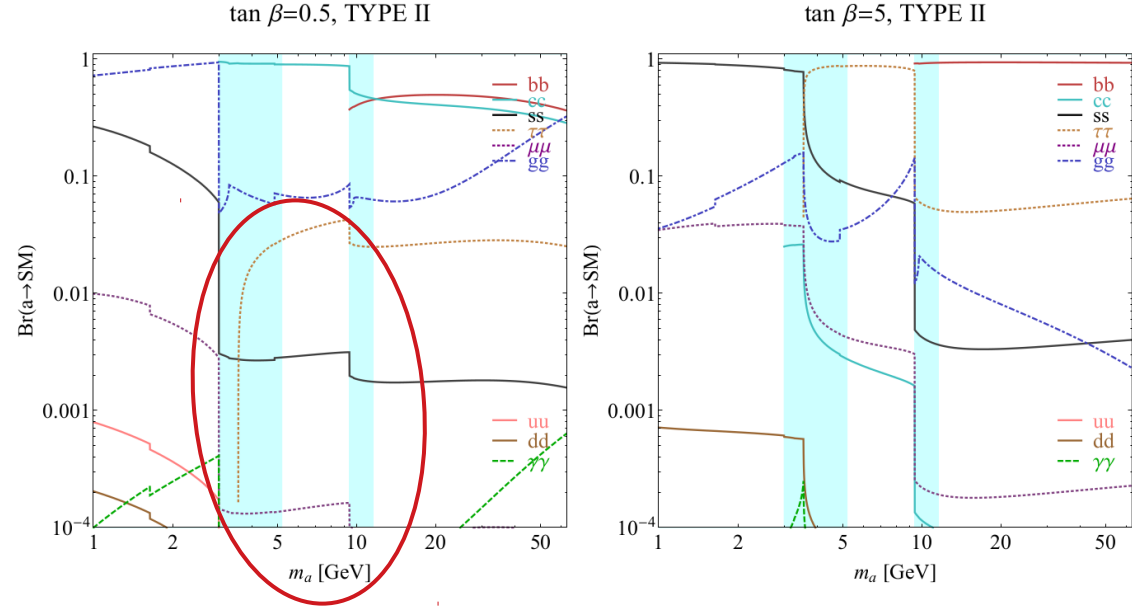


TABLE II: Couplings of the neutral scalar and pseudoscalar mass eigenstates in the four 2HDM with a \mathbb{Z}_2 symmetry, following the notation of [112]. The couplings are normalized to the SM Higgs.

FIG. 7: Branching ratios of a singlet-like pseudoscalar in the 2HDM+S for Type II Yukawa couplings. Decays to quarkonia likely invalidate our simple calculations in the shaded regions.

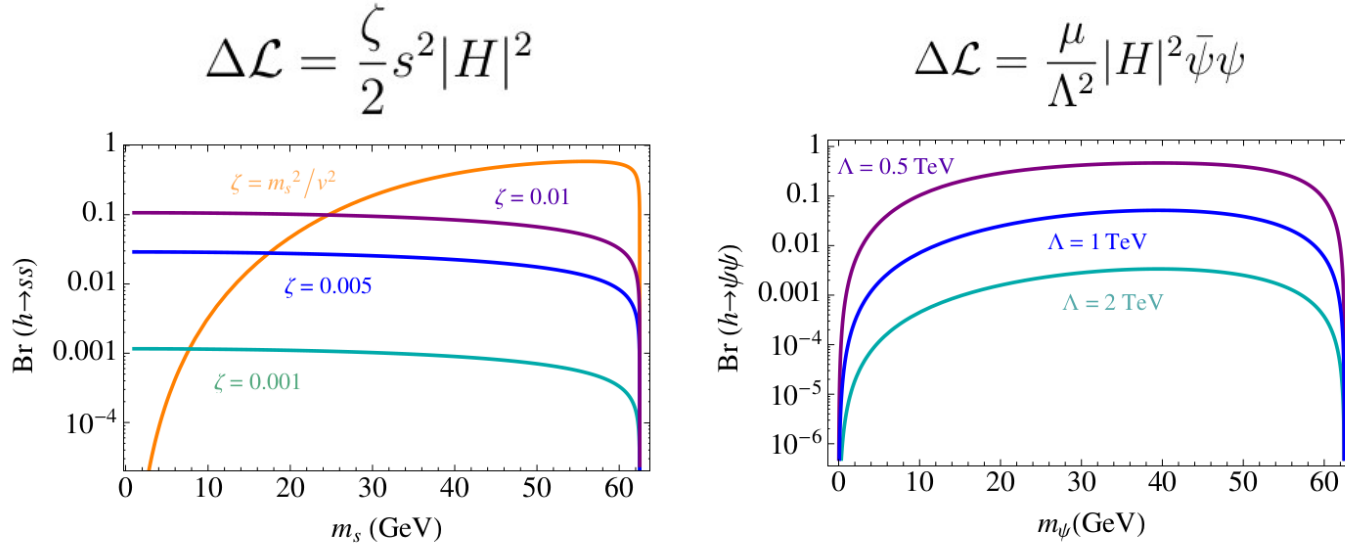
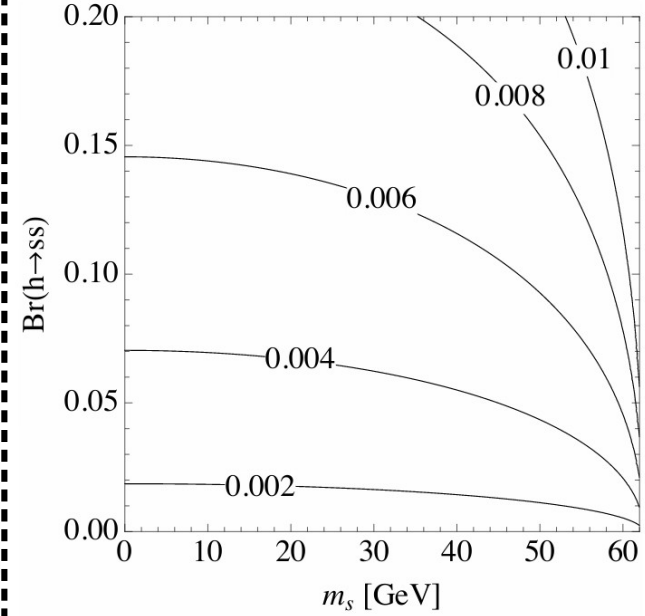
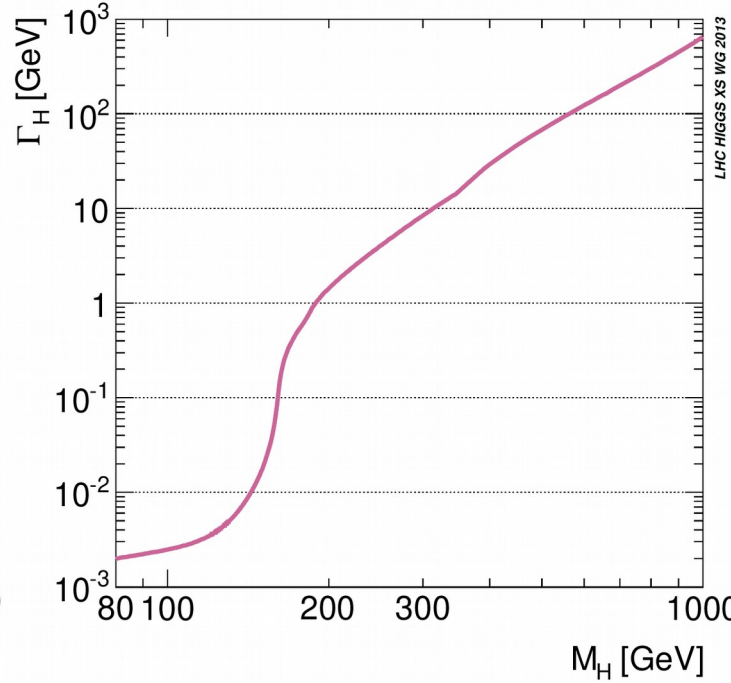
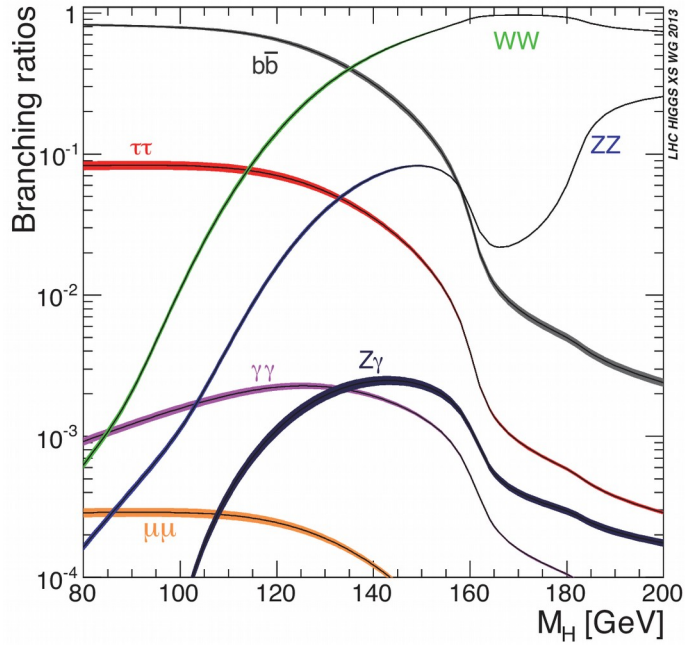


FIG. 1: Sensitivity of a 125 GeV Higgs to light weakly coupled particles. **Left:** Exotic Higgs branching fraction to a singlet scalar s versus the singlet's mass m_s , assuming the interaction Eq. (1) is solely responsible for the $h \rightarrow ss$ decay. If the interaction in Eq. (1) generates the s mass, the result is the orange curve; the other curves are for fixed and independent values of ζ and m_s . **Right:** Exotic Higgs branching fraction to a new fermion ψ interacting with the Higgs as in Eq. (2) to illustrate the sensitivity of exotic Higgs decay searches to high scales, here Λ . We take here $\mu = m_\psi$.



$$\Gamma(h \rightarrow ss) = \frac{1}{8\pi} \frac{\mu_v^2}{m_h} \sqrt{1 - \frac{4m_s^2}{m_h^2}} \approx \left(\frac{\mu_v/v}{0.015} \right)^2 \Gamma(h \rightarrow \text{SM}) \quad 20$$

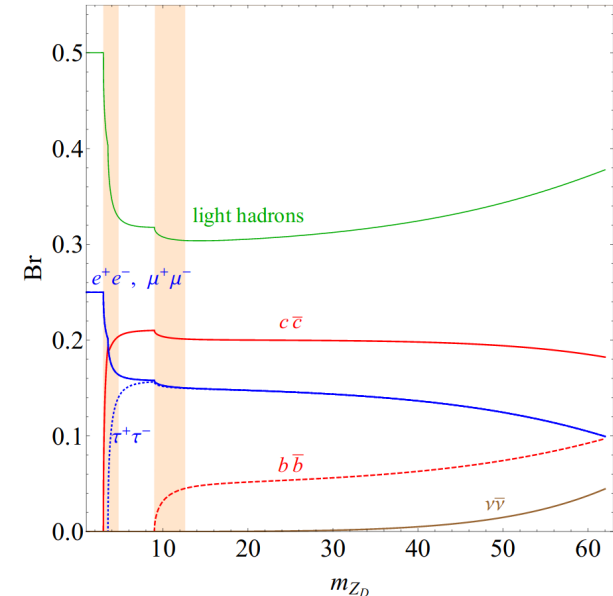
$$\mathcal{L}_{\text{gauge}} = -\frac{1}{4}\hat{B}_{\mu\nu}\hat{B}^{\mu\nu} - \frac{1}{4}\hat{Z}_{D\mu\nu}\hat{Z}_D^{\mu\nu} + \frac{1}{2}\frac{\epsilon}{\cos\theta_W}\hat{B}_{\mu\nu}\hat{Z}_D^{\mu\nu}$$

$$\begin{pmatrix} Z_D \\ B \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ -\frac{\epsilon}{\cos\theta_W} & 1 \end{pmatrix} \begin{pmatrix} \hat{Z}_D \\ \hat{B} \end{pmatrix} \quad \begin{aligned} \tilde{Z} &= Z + \epsilon_Z Z_D \\ \tilde{Z}_D &= Z_D - \epsilon_Z Z, \quad \text{where } \epsilon_Z = \frac{\epsilon \tan\theta_W m_Z^2}{m_Z^2 - m_{Z_D}^2}. \end{aligned}$$

$$\mathcal{L}_{\text{mass}} = \frac{1}{8}w^2 g_D^2 (\hat{Z}_{D\mu})^2 + \frac{1}{8}v^2 (-g\hat{W}_\mu^3 + g'\hat{B}_\mu)^2 \quad \mathcal{L}_{\text{mass}} = \frac{1}{2}m_{Z_D}^2 (Z_{D\mu})^2 + \frac{1}{2}m_Z^2 (Z_\mu - \epsilon \tan\theta_W Z_{D\mu})^2$$

$$\mathcal{L} \supset g_{Z_D f f} Z_D^\mu \bar{f} \gamma_\mu f$$

$$g_{Z_D f f} = -g' \frac{\epsilon}{\cos\theta_W} Y - \epsilon \tan\theta_W \frac{m_Z^2}{m_Z^2 - m_{Z_D}^2} \frac{1}{\sqrt{g'^2 + g^2}} (g^2 T_3 - g'^2 Y)$$



$$V_0 = -\mu^2 |H|^2 + \lambda |H|^4 - \mu_D^2 |S|^2 + \lambda_D |S|^4 + \zeta |S|^2 |H|^2$$

$$\mu^2 = v^2 \lambda + \frac{1}{2} \kappa v_S^2, \quad \mu_S^2 = v_S^2 \lambda_S + \frac{1}{2} \kappa v^2$$

$$\tilde{h} = h - \epsilon_h s$$

$$\tilde{s} = s + \epsilon_h h, \quad \text{where} \quad \epsilon_h = \zeta \frac{\mu \mu_D}{2 \sqrt{\lambda \lambda_D} |\mu^2 - \mu_D^2|}$$

