

Bianchi IX Dynamics in Dust Time

Based on MA & V. Husain [Phys. Rev. D 96,
044032 (2017), arXiv:1707.07098]

Masooma Ali

University of New Brunswick



Bianchi IX spacetime

$$ds^2 = -dt^2 + e^{-2\Omega(t)} (e^{2\beta(t)})_{ij} \omega^i \omega^j$$

$$\beta_{ij} = \text{diag}\{\beta_+ + \sqrt{3}\beta_-, \beta_+ - \sqrt{3}\beta_-, -2\beta_+\}$$



Bianchi IX spacetime

$$ds^2 = -dt^2 + e^{-2\Omega(t)} (e^{2\beta(t)})_{ij} \omega^i \omega^j$$

$$\beta_{ij} = \text{diag}\{\beta_+ + \sqrt{3}\beta_-, \beta_+ - \sqrt{3}\beta_-, -2\beta_+\}$$

$$a_i = e^{-\Omega} (e^\beta)_{ii}$$



Bianchi IX spacetime

$$ds^2 = -dt^2 + e^{-2\Omega(t)} (e^{2\beta(t)})_{ij} \omega^i \omega^j$$

$$\beta_{ij} = \text{diag}\{\beta_+ + \sqrt{3}\beta_-, \beta_+ - \sqrt{3}\beta_-, -2\beta_+\}$$

$$a_i = e^{-\Omega} (e^\beta)_{ii}$$

- BKL conjecture
- Excellent example of gravitational chaos
- Quantum problem is unsolved



Techniques to study Mixmaster GR dynamics

- Metric variable approach (Covariant Formalism)_[BKL (1970-72)]
- Particle in a box (Hamiltonian formalism) _[Misner (1969), Ryan & Shepley (1975) ..]
- Dynamical systems approach _[Bogoyavlenskii & Noviko, Ellis & MacCallum (1969), Wainwright & Hsu (1989), Ringström (2000),... Hewitt (2016)...]

All these approaches use geometric time



Techniques to study Mixmaster GR dynamics

- Metric variable approach (Covariant Formalism)_[BKL (1970-72)]
- Particle in a box (Hamiltonian formalism) _[Misner (1969), Ryan & Shepley (1975) ..]
- Dynamical systems approach _[Bogoyavlenskii & Noviko, Ellis & MacCallum (1969), Wainwright & Hsu (1989), Ringström (2000),... Hewitt (2016)...]

All these approaches use geometric time

What is the dynamics in matter time?

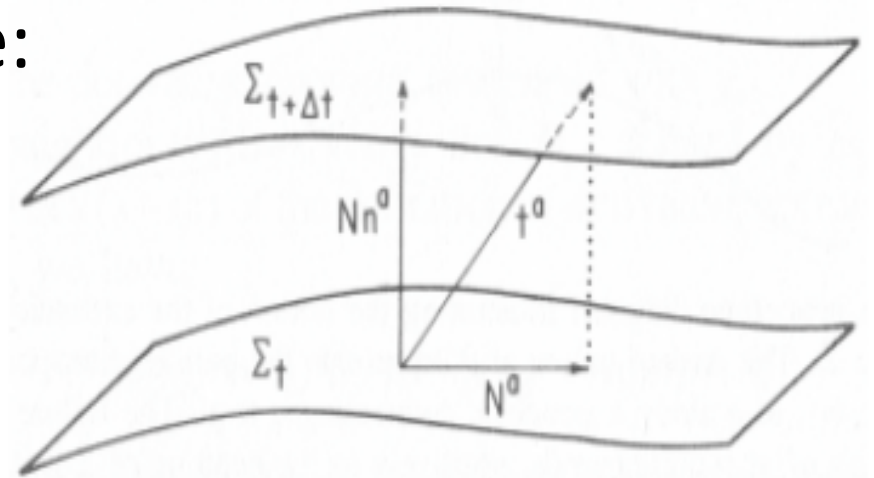


Hamiltonian approach

Dynamical variables are:

$$(q_{ab}, \pi^{ab})$$

$$(\phi, p_\phi)$$

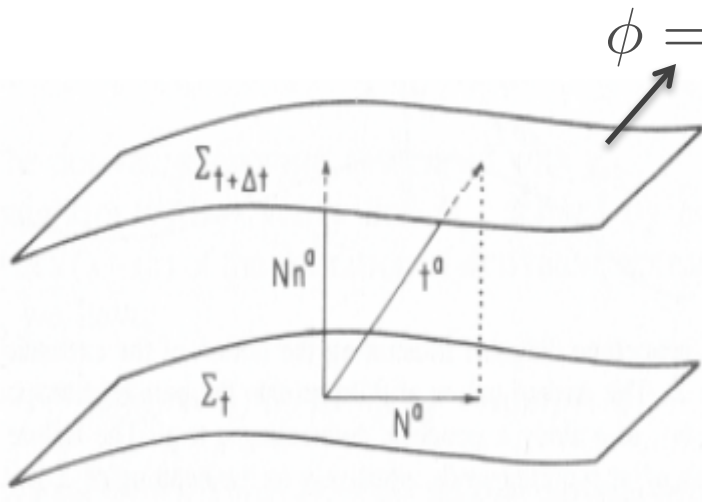


$$S = \int dt d^3x \left[\pi^{ab} \dot{q}_{ab} + p_\phi \dot{\phi} - \underbrace{N\mathcal{H} - N^a \mathcal{C}_a}_{\text{Hamiltonian}} \right]$$

Dust time

Husain & Pawłowski (2012)

Surfaces of constant time = level surfaces of dust



$$\lambda \equiv \phi - t = 0$$
$$\implies N = -1$$

$$H_p = - \int d^3x (\mathcal{H}_G + \mathcal{H}_M + N^a \mathcal{C}_a)$$

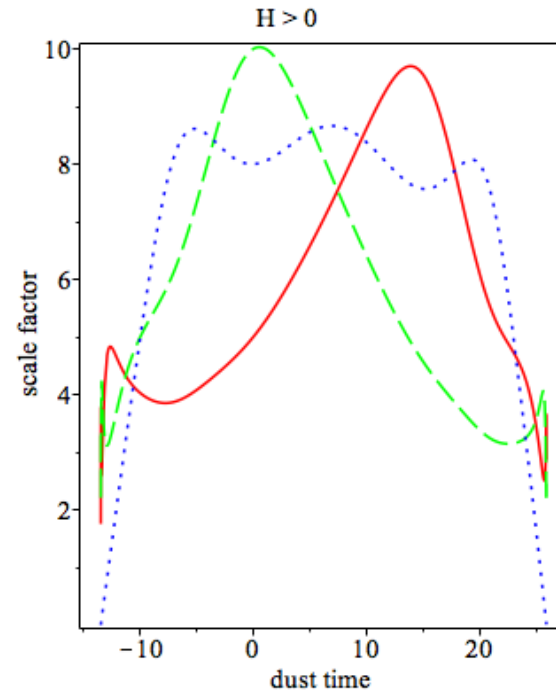
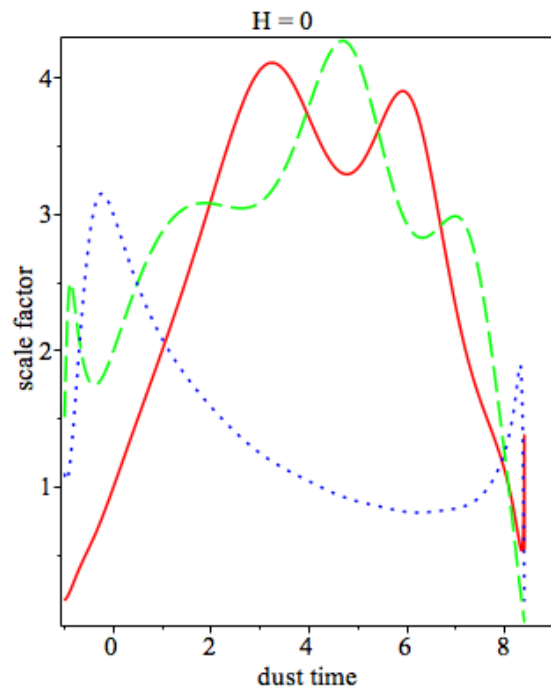
Bianchi IX in dust time

$$H_p = \underbrace{-e^{3\Omega}(p_+^2 + p_-^2 - p_\Omega^2)}_{\text{K(scale factors, momenta)}} + \underbrace{e^{-\Omega}v(\beta_+, \beta_-)}_{\text{V(scale factors)}}$$

K(scale factors, momenta)

V(scale factors)

$$(\Omega, p_\Omega), (\beta_\pm, p_\pm)$$



“Free particle” = (Dust) Kasner

$$H_p = -e^{3\Omega} (p_+^2 + p_-^2 - p_\Omega^2)$$

Kasner Solution:

$$H_p = 0 : a_i = t^{p_i}$$

“Free particle” = (Dust) Kasner

$$H_p = -e^{3\Omega} (p_+^2 + p_-^2 - p_\Omega^2)$$

Kasner Solution:

$$H_p = 0 : a_i = t^{p_i}$$

Dust Kasner Solution:

$$H_p > 0 :$$

$$a_i = (t - \Gamma)^{p_i} (t + \Gamma)^{2/3 - p_i}$$

“Free particle” = (Dust) Kasner

$$H_p = -e^{3\Omega} (p_+^2 + p_-^2 - p_\Omega^2)$$

Kasner Solution:

$$H_p = 0 : a_i = t^{p_i}$$

Dust Kasner Solution:

$$H_p > 0 :$$

$$a_i = (t - \Gamma)^{p_i} (t + \Gamma)^{2/3 - p_i}$$

$$\sum p_i = 1 \quad \sum p_i^2 = 1$$

Characterizing Dust-Kasner

Three relevant integration constants

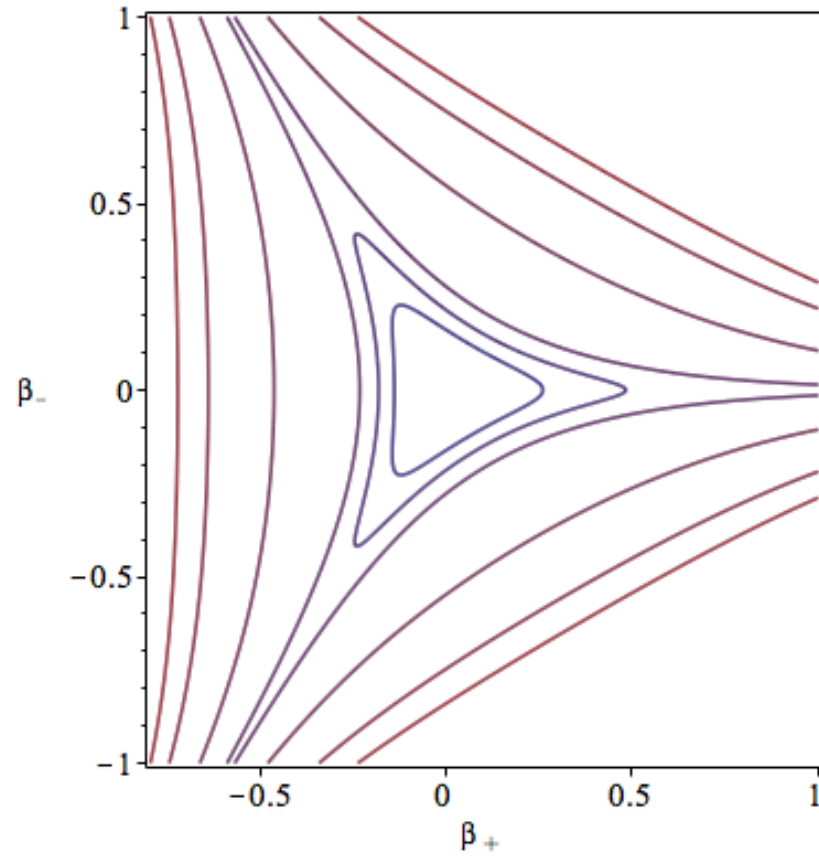
$$(p_{\Omega}^0, v^0, \theta)$$

Define:

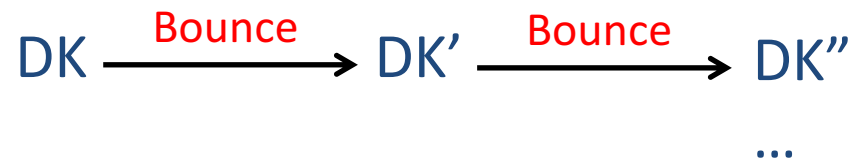
$$\frac{p_{-}^0}{\Gamma} = \sin \theta$$

Particle in a box

$$H_p = -e^{3\Omega}(p_+^2 + p_-^2 - p_\Omega^2) + e^{-\Omega}v(\beta_+, \beta_-)$$

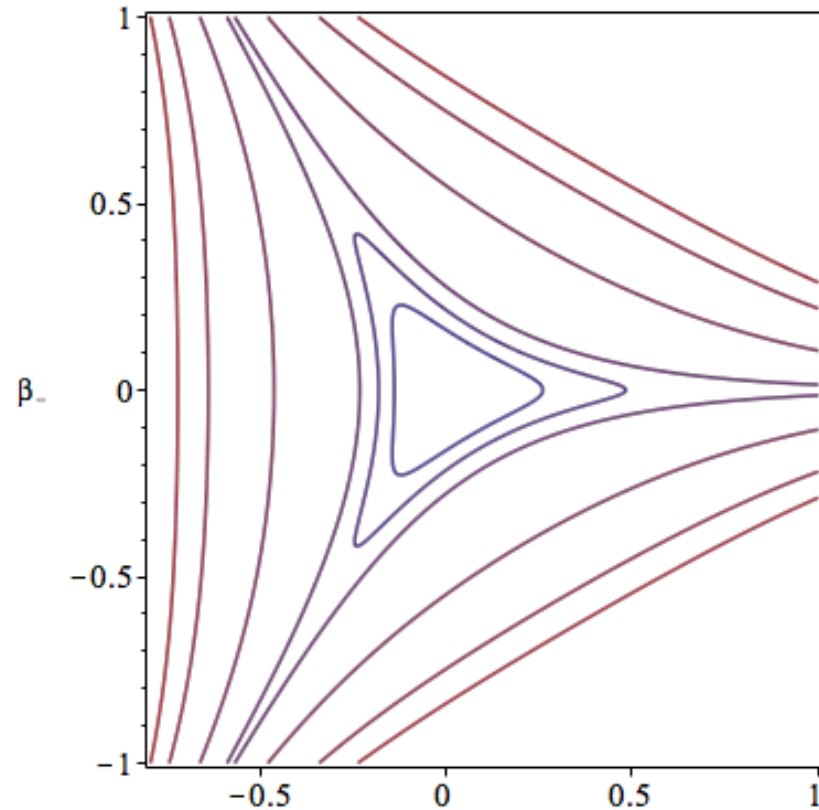


At each wall one term in V is dominant and perturbs the DK Hamiltonian

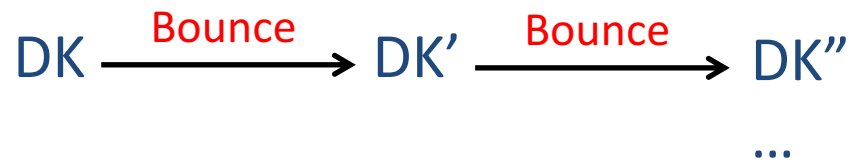


Particle in a box

$$H_p = -e^{3\Omega}(p_+^2 + p_-^2 - p_\Omega^2) + e^{-\Omega}v(\beta_+, \beta_-)$$



At each wall one term in V is dominant and perturbs the DK Hamiltonian



Bianchi IX dynamics can be regarded as a series of bounces between Dust-Kasner solutions

Dynamics near one wall

$$H_p = \text{constant} \quad p_- = \text{constant} \quad p_\Omega^0 - \frac{1}{2}p_+ = \text{constant}$$

Dynamics near one wall

$$H_p = \text{constant} \quad p_- = \text{constant} \quad p_\Omega^0 - \frac{1}{2}p_+ = \text{constant}$$

$$\Gamma^{(i)} \sin \theta_i = \Gamma^{(f)} \sin \theta_f$$

$$\left(\frac{p_\Omega^0}{\Gamma}\right)^{(i)} \sin \theta_f - \left(\frac{p_\Omega^0}{\Gamma}\right)^{(f)} \sin \theta_i = \frac{1}{2} \sin(\theta_i + \theta_f)$$

$$\Delta(p_\Omega^0)^2 - 24H_p\Delta v^0 = \Delta\Gamma^2$$

Dynamics near one wall

$$\left(\frac{p_{\Omega}^0}{\Gamma}\right)^{(i)} \sin \theta_f - \left(\frac{p_{\Omega}^0}{\Gamma}\right)^{(f)} \sin \theta_i = \frac{1}{2} \sin(\theta_i + \theta_f)$$

Vacuum Case: $H_p = 0 \implies \Gamma = p_{\Omega}^0$

$$\sin \theta_i - \sin \theta_f = \frac{1}{2} \sin(\theta_i + \theta_f)$$

Dynamics near one wall

$$\left(\frac{p_{\Omega}^0}{\Gamma}\right)^{(i)} \sin \theta_f - \left(\frac{p_{\Omega}^0}{\Gamma}\right)^{(f)} \sin \theta_i = \frac{1}{2} \sin(\theta_i + \theta_f)$$

Vacuum Case: $H_p = 0 \implies \Gamma = p_{\Omega}^0$

$$\sin \theta_i - \sin \theta_f = \frac{1}{2} \sin(\theta_i + \theta_f)$$

Asymptotic Limit: $v^0 \rightarrow 0 \implies \Gamma \rightarrow p_{\Omega}^0$

Conclusion

- Dust time gives a **new** physical picture.
- Complementary to the other approaches.
- New transition law for dynamics away from the corners
- Bonus: The simple form of the Hamiltonian allows a straightforward construction of the path integral for Bianchi I with dust