



# The Nuclear Delta Force in Quadrupole Deformed Nuclei

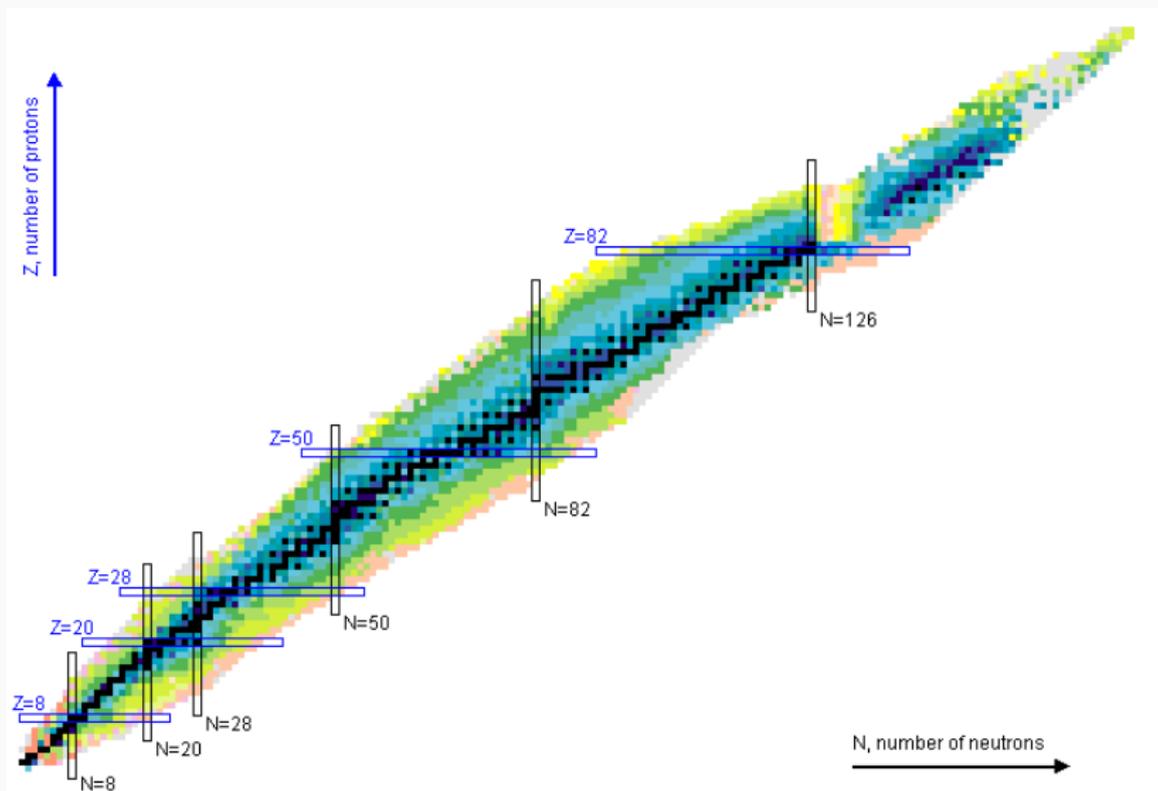
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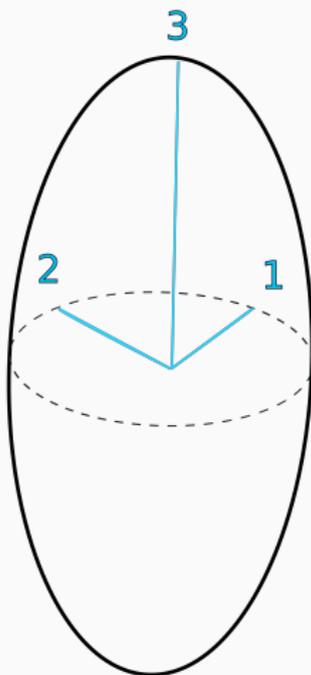
# Chart of Nuclides



# Collective Degrees of Freedom

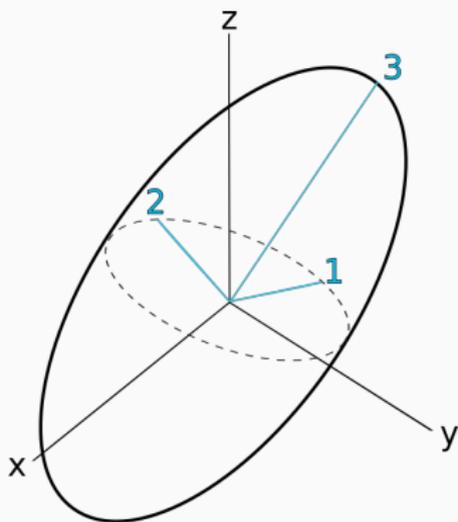
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## Quadrupole Deformation: Axial Symmetry



# Rotor Wavefunction

$$|RMK\rangle = \frac{1}{N} D_{MK}^R(\phi_E, \theta_E, \psi_E)$$



# Applying the Axial Symmetry to the Rotor Wavefunction

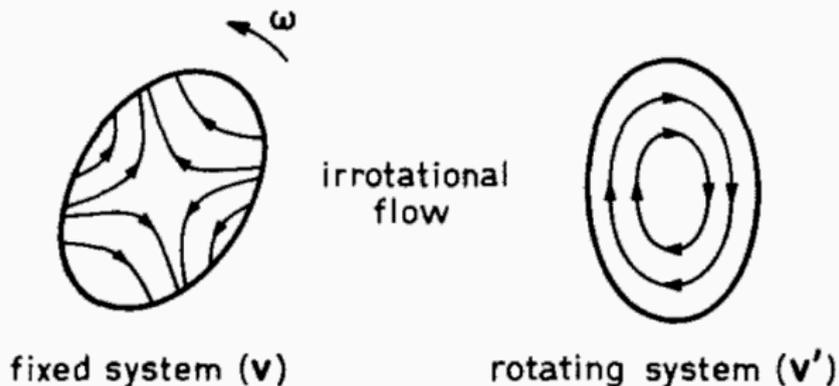
The characteristic symmetry elements is the  $\pi$  rotation about the short axis and the identity element:

$$\hat{S} = \hat{R}(\mathbf{e}_1, \pi) + \hat{\mathbb{I}}.$$

Applying this operator to  $|RMK\rangle$  and normalizing, we obtain the rotor wavefunction:

$$|\Psi_{RMK}\rangle = \sqrt{\frac{2R+1}{16\pi^2(1+\delta_{K,0})}} (D_{MK}^R + (-1)^R D_{M\bar{K}}^R).$$

# Liquid Drop Model



**Left:** The surface flow patterns are visualized from the fixed lab frame for a liquid drop.

**Right:** The surface flow patterns are visualized in the intrinsic rotating frame for a liquid drop<sup>†</sup>.

$$\mathcal{I}_{33} = 0, \mathcal{I}_{11} = \mathcal{I}_{22} = \frac{3}{4}\mathcal{I}_0$$

<sup>†</sup>A. Bohr and B. Mottelson, *Nuclear Structure, Vol. II*, W. A. Benjamin Inc., 1975

## Axial Rotor Hamiltonian

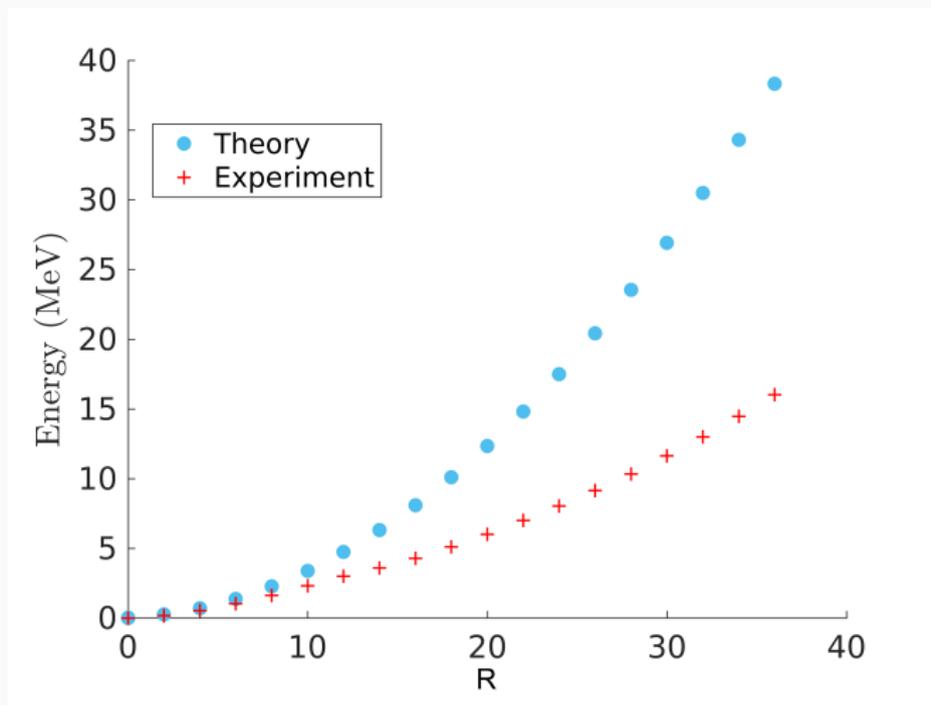
$$H_{Rot} = \frac{1}{2} \left[ \frac{\hat{R}_1^2}{\mathcal{I}_{11}} + \frac{\hat{R}_2^2}{\mathcal{I}_{22}} + \frac{\hat{R}_3^2}{\mathcal{I}_{33}} \right]$$

But,  $\mathcal{I}_{11} = \mathcal{I}_{22} = \frac{3}{4}\mathcal{I}_0$  and  $\mathcal{I}_{33} = 0 \implies R_3 = 0$

$$H_{Rot} = \frac{2}{3} \left[ \frac{\hat{R}^2}{\mathcal{I}_0} \right]$$

$$H_{Rot}|\Psi RM0\rangle = \frac{2\hbar^2}{3\mathcal{I}_0} R(R+1)|\Psi RM0\rangle$$

## Comparison To Data: $^{126}\text{Ce}$

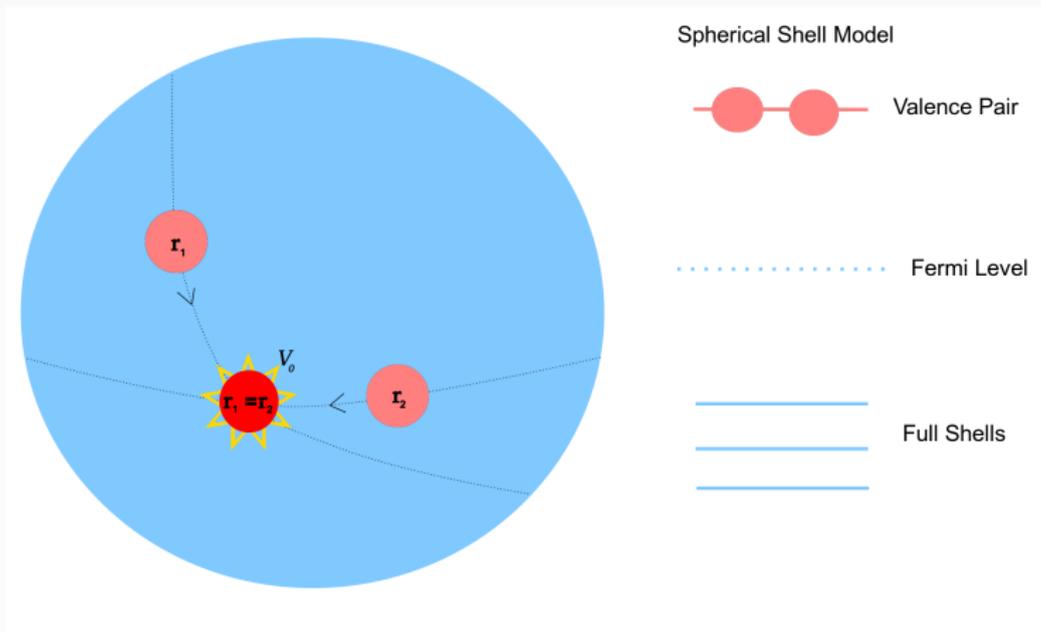


The parameter  $\mathcal{I}_0$  was estimated using the  $R = 2$  state, to be  $23.6 \hbar^2 / \text{MeV}$ .

# Single Particle Degrees of Freedom

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# Nuclear Delta Force



$$H = H_{Spherical} + H_{\delta}$$

$$E = (E_1 + E_2) \langle \Psi J' \Omega' | \Psi J \Omega \rangle + \langle \Psi J' \Omega' | V_{\delta} \delta(\mathbf{r}_2 - \mathbf{r}_1) | \Psi J \Omega \rangle$$

$$|\Psi_{J\Omega}\rangle \triangleq \frac{1 - (-1)^{j_1+j_2-J}}{2} |\mathcal{R}\rangle |J\Omega\rangle$$

where

$$|J\Omega\rangle = \sum_{\Omega_1, \Omega_2} \langle j_1\Omega_1 j_2\Omega_2 | J\Omega \rangle |j_1\Omega_1\rangle |j_2\Omega_2\rangle$$

and

$$|\mathcal{R}(r_1)\rangle |\mathcal{R}(r_2)\rangle \triangleq |\mathcal{R}\rangle$$

$$H_\delta = V_\delta \delta(\mathbf{r}_2 - \mathbf{r}_1) = V_\delta \delta(r_2 - r_1) \delta(\theta_2 - \theta_1) \delta(\phi_2 - \phi_1)$$

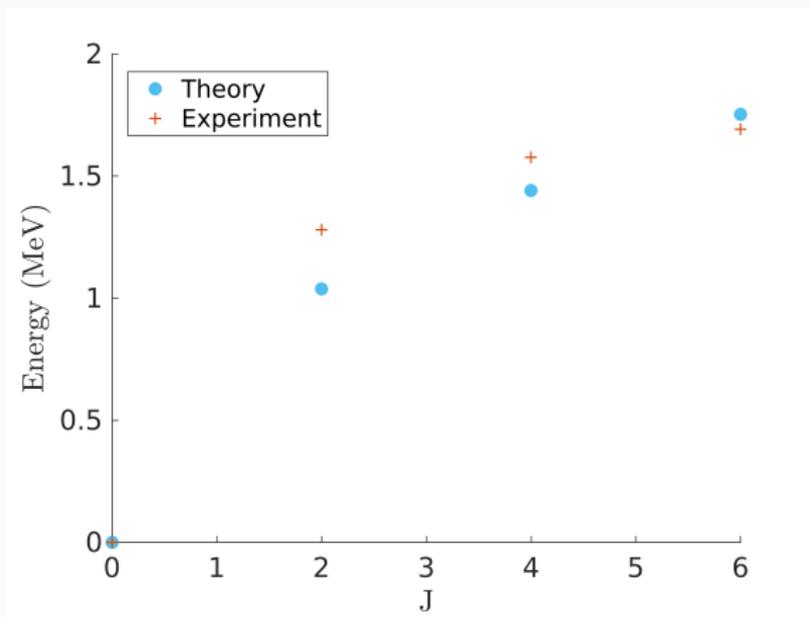
Applying this, we derive

$$\langle \Psi J' \Omega' | H_\delta | \Psi J \Omega \rangle = V_\delta \langle \mathcal{R}' | \delta(r_2 - r_1) | \mathcal{R} \rangle \langle J' Q' | \delta(\theta_2 - \theta_1) \delta(\phi_2 - \phi_1) | J Q \rangle$$

where

$$V_\delta \langle \mathcal{R}' | \delta(r_2 - r_1) | \mathcal{R} \rangle \triangleq \xi$$

## Comparison to Data: $^{134}\text{Te}$



By fitting the model predictions to the data,  $\xi$  was found to be  $-7.0$  MeV for the two  $g_{9/2}$  valence protons of  $^{134}\text{Te}$ .

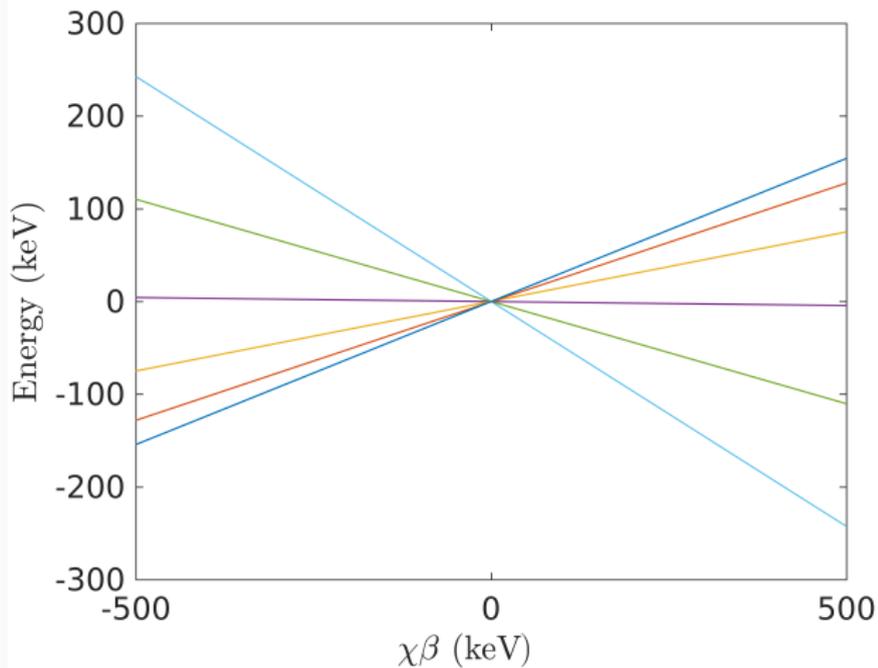
Data extracted using the NNDC On-Line Data Service from the ENSDF database.

## Two Coupled Nucleons in an Axially-Deformed Potential

$$H = H_{spherical} + H_{\beta}$$

$$H_{\beta} = \pm \chi \beta [Y_{20}(\theta_1, \phi_1) + Y_{20}(\theta_2, \phi_2)]$$

# Model



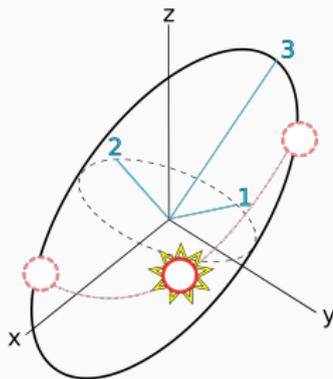
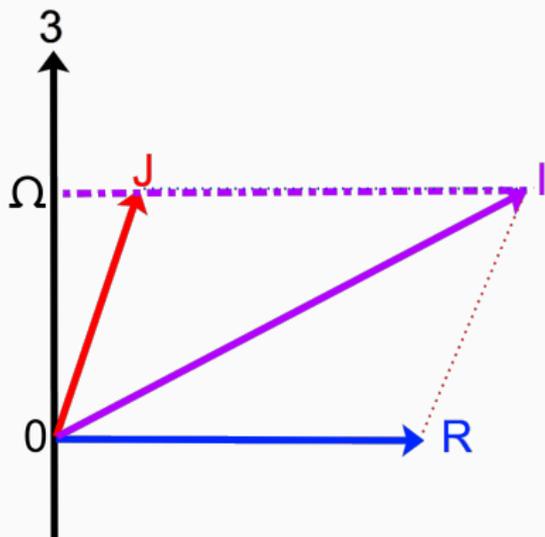
For two  $h_{11/2}$  protons, the resultant energies from  $H_\beta$  are shown for the  $\Omega = 0$  state.

# Pair Rotor Coupling Model with Delta Force Interaction

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# The Basis Wavefunction

$$|IMj_1j_2J\Omega\rangle = \frac{1 - (-1)^{j_1+j_2-J}}{2} \sqrt{\frac{2I+1}{16\pi^2(1+\delta_{\Omega,0})}} \left[ D'_{M\Omega} |J\Omega\rangle + (-1)^{I-J} D'_{M\bar{\Omega}} |J\bar{\Omega}\rangle \right]$$



$$H = H_{Rotor} + H_{\delta} + H_{\beta}$$

$$H_{Rotor} = \frac{2}{3} \left[ \frac{\mathbf{R}^2}{\mathcal{I}_o} \right]$$

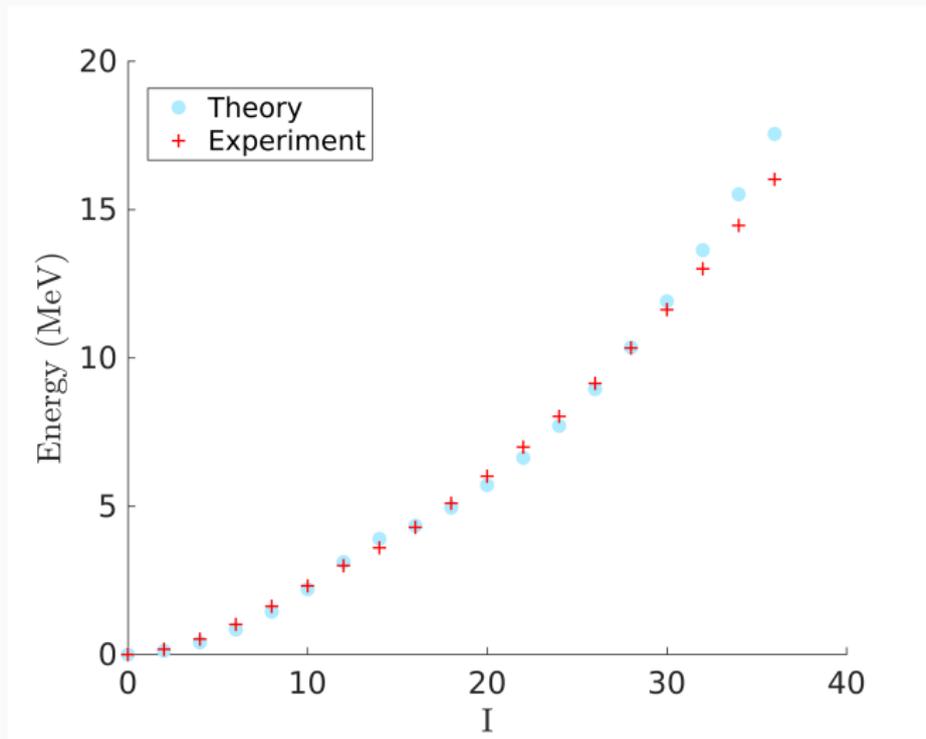
$$H_{Rotor} = \frac{2}{3\mathcal{I}_o} [\mathbf{I} - \mathbf{J}]^2$$

$$H_{Rotor} = \frac{2\hbar^2}{3\mathcal{I}_o} [\hat{I}^2 - \hat{I}_3^2 + \hat{J}^2 - \hat{J}_3^2] + \frac{4\hbar^2}{3\mathcal{I}_o} [\hat{I}_{+1}\hat{J}_{-1} + \hat{I}_{-1}\hat{J}_{+1}]$$

$$H_{Rotor} = H_{Diagonal} + H_{Coriolis}$$

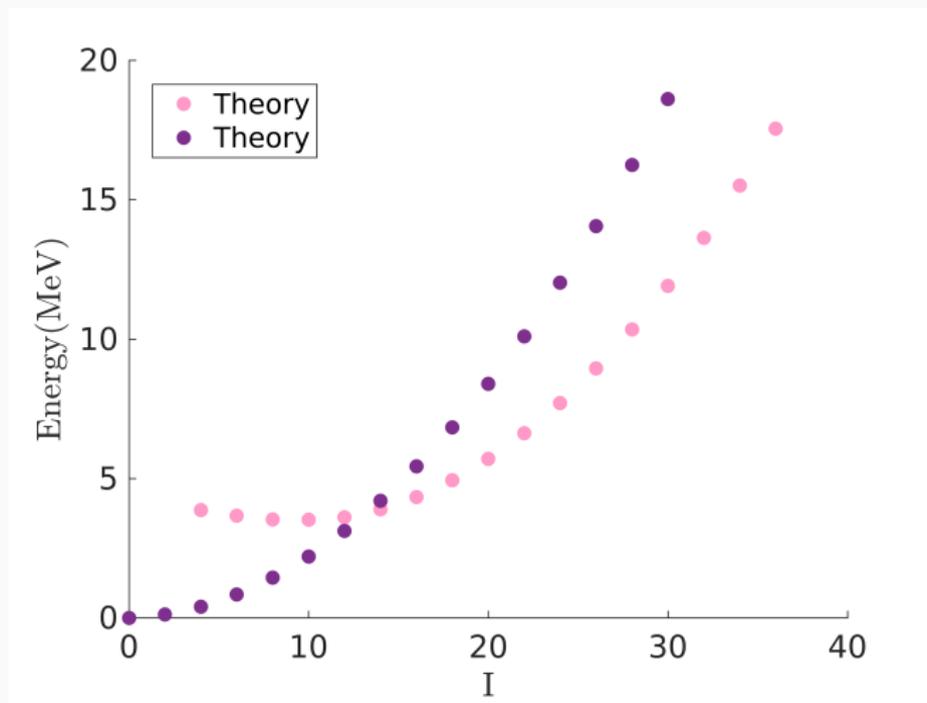
$$H = H_{Diagonal} + H_{Coriolis} + H_{\delta} + H_{\beta}$$

# Comparison to Data: $^{126}\text{Ce}$



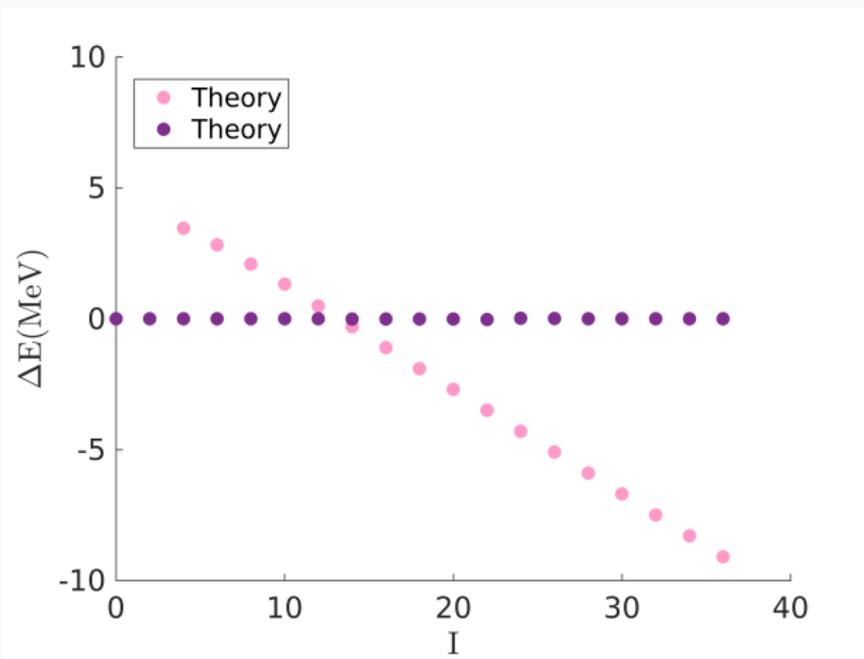
Two  $h_{11/2}$  protons taken as the valence particles for  $^{126}\text{Ce}$ .  
 $\chi\beta = 0.2$  MeV,  $\xi = -8.4$  MeV,  $\mathcal{I}_0 = 33.2 \hbar^2/\text{MeV}$

# Band Crossing

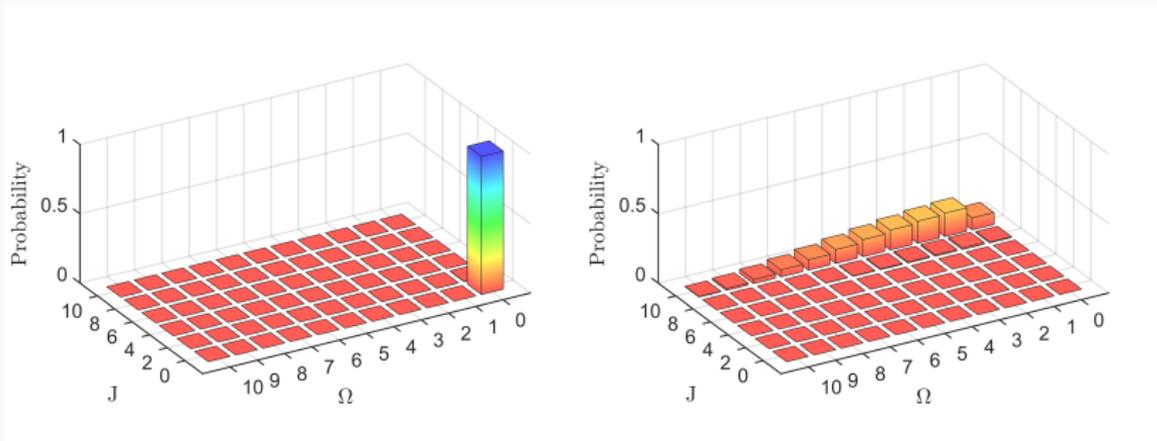


# Single Particle Effects Vs Collective Effects

$$\Delta E = E - \frac{2\hbar^2}{3\mathcal{I}_0} I(I+1)$$



# Probability Distribution



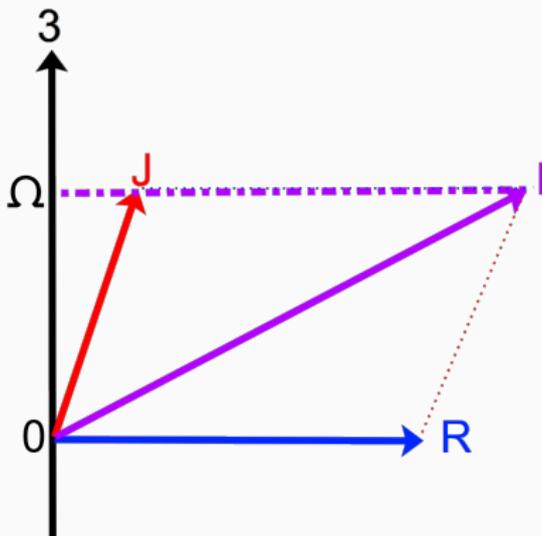
**Left:** The probability distribution over the angular momentum states for  $l = 12$ .

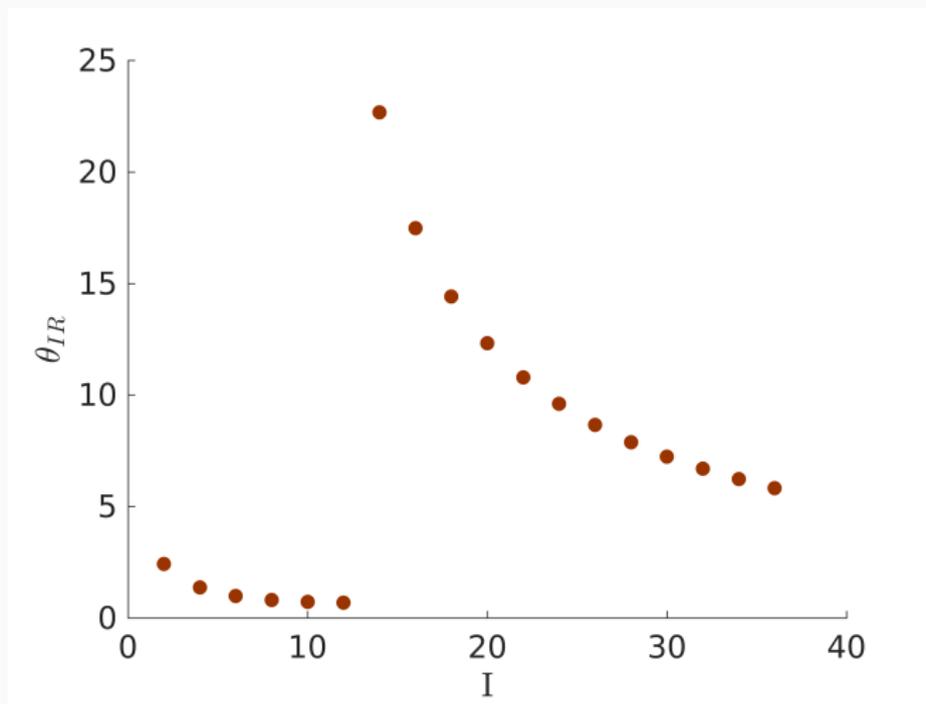
**Right:** The probability distribution over the angular momentum states for  $l = 14$ .

# Angles

What if we want to find the angles between the vectors  $I$  and  $R$ ?

$$\cos(\theta_{IR}) = \frac{\langle \mathbf{I} \cdot \mathbf{R} \rangle}{\sqrt{\langle I^2 \rangle} \sqrt{\langle R^2 \rangle}}$$





Using the lowest energy wavefunction for a given spin, the expectation value of the angle between  $\mathbf{I}$  and  $\mathbf{R}$  is calculated as a function of  $I$ .

$$\begin{aligned}\boldsymbol{\mu} &= g \frac{\mu_N}{\hbar} \mathbf{I} \\ \boldsymbol{\mu} &= \frac{\mu_N}{\hbar} (g_R \mathbf{R} + g_p \mathbf{J})\end{aligned}$$

Equating both sides and taking the dot product with  $\vec{I}$ ,

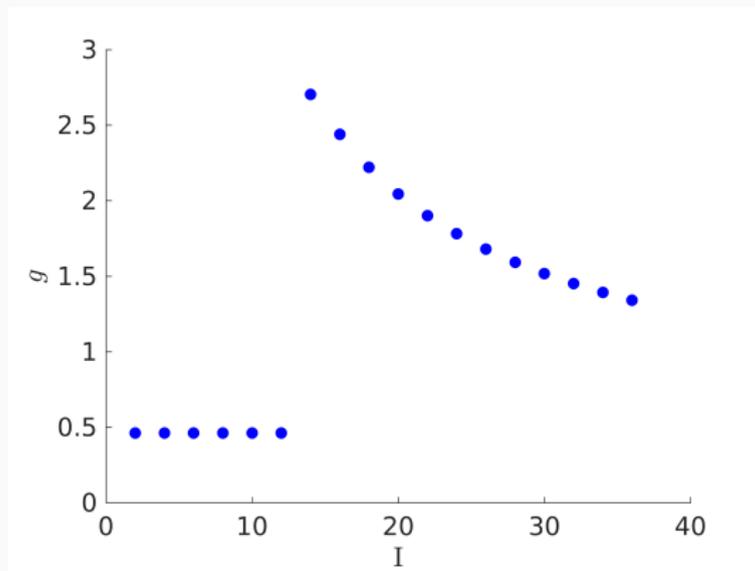
$$g = \frac{g_R \mathbf{I} \cdot \mathbf{R} + g_p \mathbf{I} \cdot (\mathbf{I} - \mathbf{R})}{I^2},$$

where we can estimate

$$\begin{aligned}g_R &= \frac{Z}{A} \\ g_p &= a \times g_{\text{Free Proton}}.\end{aligned}$$

# G-Factor

$$g = \frac{(g_R - g_p)\langle \mathbf{l} \cdot \mathbf{R} \rangle + g_p \langle I^2 \rangle}{\langle I^2 \rangle}$$



The expectation value of  $g$  predicted for a given spin,  $I$ , using the lowest energy wavefunction in the model. Here,  $a = 0.65$ ,  $g_{\text{Free Proton}} = 5.59$ ,  $g_R = \frac{58}{126}$ .

# Summary, Future Work, and Thanks

1. The model shows that the single particle degrees of freedom are of comparable impact with the collective degrees of freedom.
2. Transition rates and  $B(E2)$  to be calculated.
3. New deformations to be explored : Triaxial quadrupole, octupole, hexadecapole, etc.



**Questions?**