

TWO PHOTON ABSORPTION

A Test of Pseudo State Summations to Calculate Polarizability

Spencer Percy and G.W.F. Drake

University of Windsor

Overview

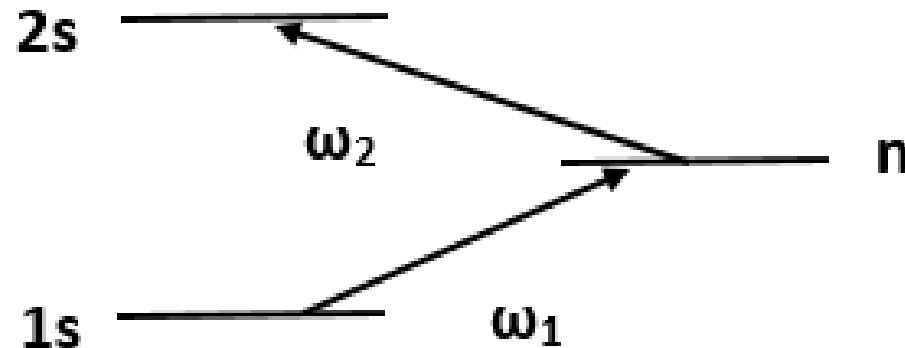
1. **Two Photon Absorption**
2. **Perturbation theory and why we use Pseudostates**
3. **Variational Methods Example: The Polarizability of the $1s$ State of Hydrogen**

Motivation

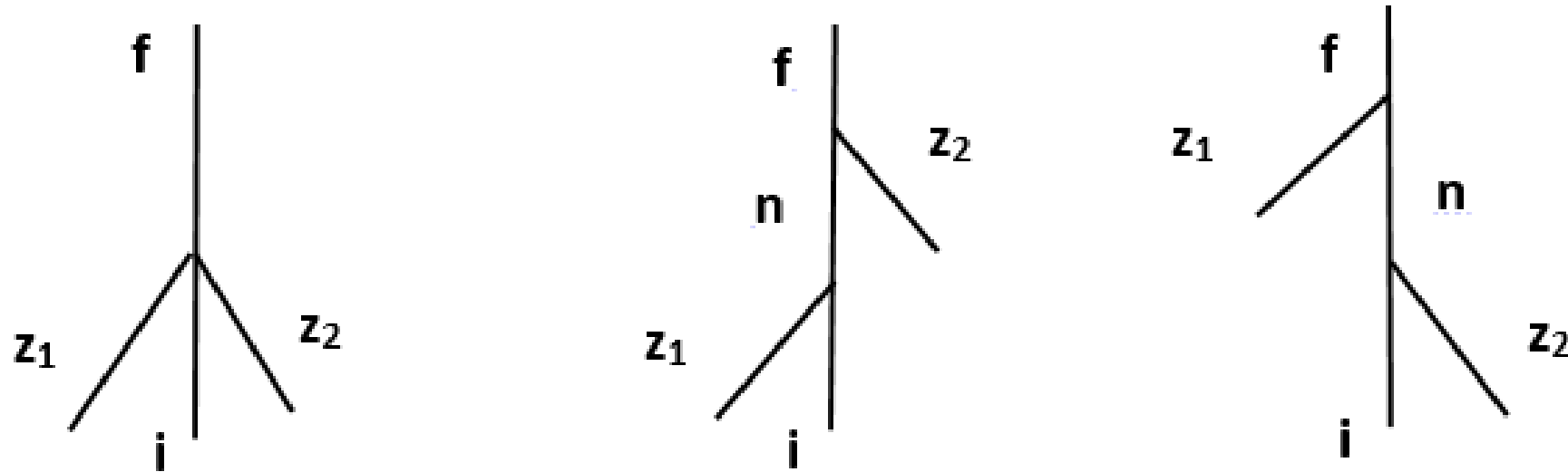
- **2s to 1s transitions with a single electric dipole photon are forbidden by parity. The dominant radiative decay mechanism is two-photon transitions. The 2s state lives for about 1/8 sec, which is a billion times longer than the neighboring 2p state**
- **Using lasers you are able to jump to a higher excited state using a lower frequency laser than the full transition would require**
- **Our group will be checking to see if the process could be a significant correction to astrophysical spectroscopy**

Two Photon Absorption

- Defined as the simultaneous absorption of two photons
- Very similar in operation to scattering processes
- Two different frequency photons can be used as long as the sum of their frequencies is equal to the transition frequency of the given state



Summation Over States



Possible Transitions: Simultaneous absorption and absorbing one photon then the other

$$\sum_n \left(\frac{\langle f|z_2|n\rangle\langle n|z_1|i\rangle}{\omega_1 - \omega_n} + \frac{\langle f|z_1|n\rangle\langle n|z_2|i\rangle}{\omega_2 - \omega_n} \right)$$

Example: Polarizability of Hydrogen

- A very similar equation is found when calculating polarizability
- Defined as the second order perturbation energy due to an external field
- To calculate it you are required to sum over the bound states and then integrate over the continuum

$$\alpha_d = -2 \sum_n \frac{\langle 1s|z|n\rangle\langle n|z|1s\rangle}{E_{1s} - E_n}$$

Example Cont'd

- For an external field $V = eFr \cos \theta$ the first order equation can be solved analytically

$$\psi^{(1)} = -(1/3)(2r + r^2)e^{-r}Y_1^0(\hat{\mathbf{r}})$$

- Which satisfies the first-order perturbation equation

$$(H^{(0)} - E^{(0)})|\psi^{(1)}\rangle + eFz|\psi^{(0)}\rangle = 0$$

- The first order energy $E^{(1)} = eF\langle\psi^{(0)}|z|\psi^{(0)}\rangle$ is zero by parity and thus not included

Example Cont'd

- The second order energy is given by $E^{(2)} = eF\langle\Psi^{(1)}|z|\Psi^{(0)}\rangle$
- Instead of solving $\Psi^{(1)}$ analytically we can multiply the first-order perturbation equation by the inverse operator $(H^{(0)} - E^{(0)})^{-1}$ to obtain the spectral representation of $\Psi^{(1)}$

$$|\Psi^{(1)}\rangle = eF \sum_n \frac{|\Psi_n\rangle\langle\Psi_n|z|\Psi^{(0)}\rangle}{E^{(0)} - E_n}$$

- Where the sum over n represents the summation over the infinite bound states and integration over the continuum

Alternatively

- As an alternative to this lengthy summation we can instead insert a discrete variational set of pseudostates of the form

$$\tilde{X}_p = \sum_i c_i r^i e^{-\beta r} \cos \theta$$

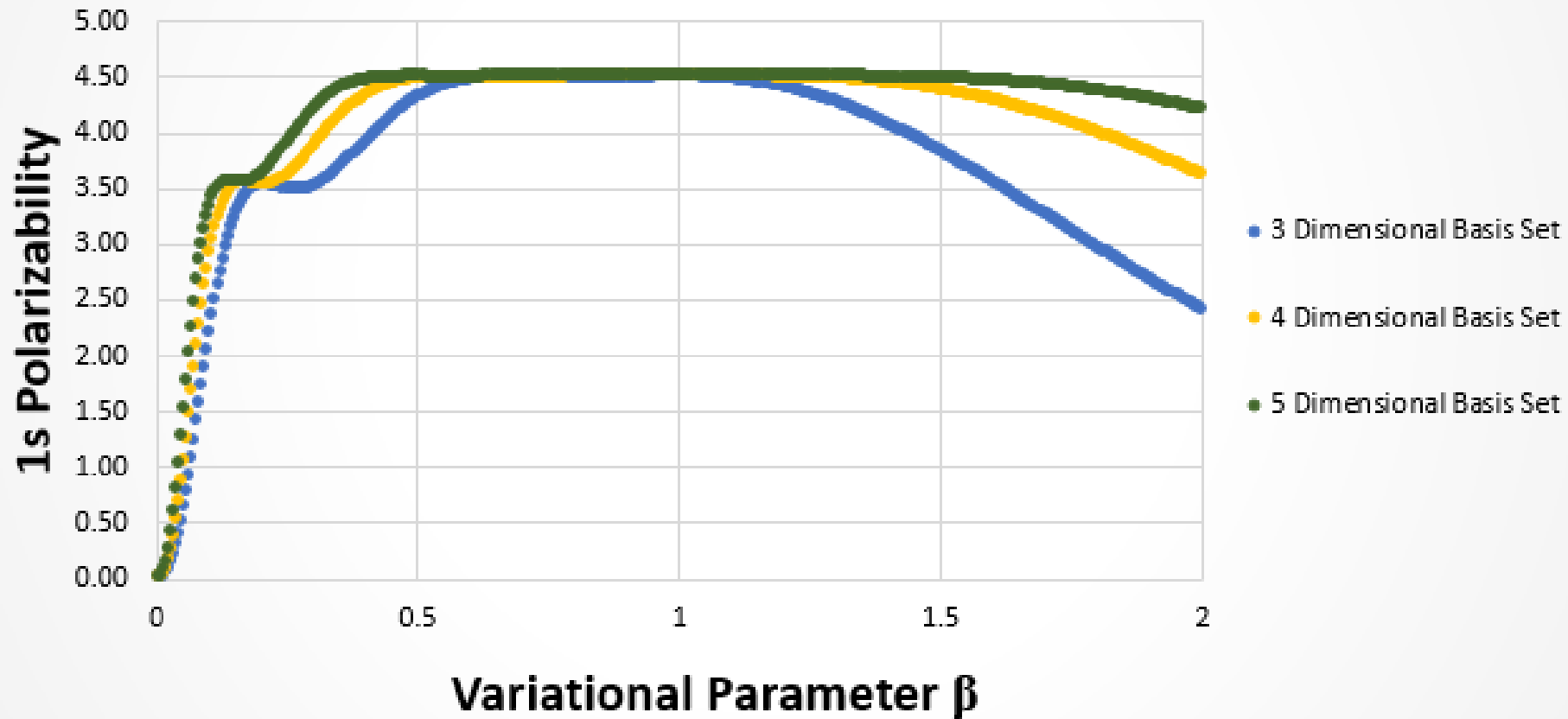
- Where c_i are linear variational parameters found by diagonalizing the Hamiltonian in the generalized eigenvalue problem, β is our non-linear variational parameter and we multiply by $\cos \theta$ because for our example we need P-states

Convenience

- This method will give the exact analytic solution to $\Psi^{(1)}$ for a two term basis set when $\beta=1$
- Using this method with our example we can calculate the solution to the polarizability calculation to be $\alpha_d = 4.5a_0^3$ with a_0 being the Bohr radius
- In addition, one obtains a variational extremum at $\beta = 1$, as demonstrated by varying β and calculating the polarizability. The value obtained from pseudostates for arbitrary β is therefore a lower bound on the exact value. One can expect this also to be true in the two-electron case of helium where the exact analytic solutions are not known

Variational Methods

1s State Polarizability as a function of the Variational parameter β



Thank you

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