

Harvesting Entanglement with Inertially Moving Detectors

Laura Henderson

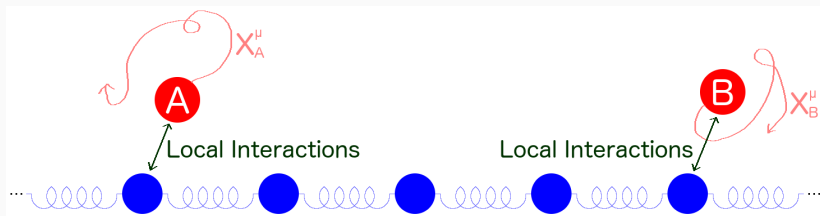
R. B. Mann

E. Martín-Martínez

University of Waterloo

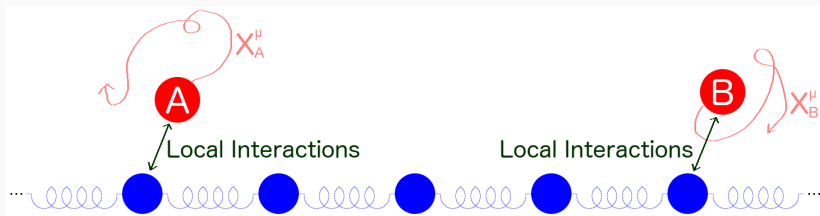
What is “Entanglement Harvesting”?

1. Take two *uncorrelated* particle detectors and allow them to locally interact with a free quantum field.



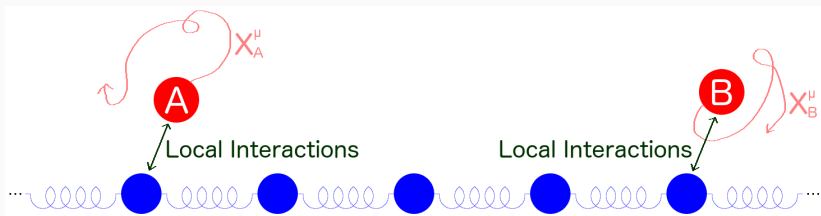
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2. After some time, these two detectors will become *entangled*, even if they remain spacelike separated.
3. This entanglement has been swapped from the field to the detectors.



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¹S. J. Summers and R. F. Werner, Phys. Lett. A 110, 257 (1985).

²A. Valentini, Phys. Lett. A 153, 321 (1991).

³G. Salton, R. B. Mann, and N. C. Menicucci, New J. Phys. 17, 035001 (2015).

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History

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- In 2016 Pozas-Kerstjens and Martín-Martínez calculated the entanglement for a pair of stationary detectors over a variety of spacial profiles and switching functions⁴.

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What happens to entanglement harvesting when two detectors move with a **constant velocity**?

The Field and the Detector

- The field is a massless scalar field, $\hat{\phi}(\mathbf{x}, t)$, in vacuum state, $|0\rangle$.

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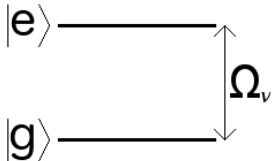
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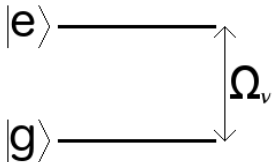
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- They detect a particle when they move from the ground, $|g\rangle$, to the excited state, $|e\rangle$.

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The Interaction Hamiltonian

The detectors (A and B) are coupled to the scalar field by the interaction Hamiltonian⁶

$$\hat{H}_I = \lambda \sum_{\nu \in \{A, B\}} \chi(t) \hat{\mu}_\nu(t) \hat{\phi}(\mathbf{x}_\nu(t)),$$

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- $\hat{\mu}_\nu(t)$ is the monopole moment of each detector.
- $\mathbf{x}_\nu(t)$ is the spacetime position of each detector.

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Evolution of the System

1. The detector-field system is initially in the state

$$\rho_0 = |0\rangle \langle 0| \otimes |g_A\rangle \langle g_A| \otimes |g_B\rangle \langle g_B|.$$

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3. Trace out the field to get the final state of the two detectors

$$\rho_{AB} = \text{Tr}_\phi [\rho] = \begin{pmatrix} 1 - \mathcal{L}_{AA} - \mathcal{L}_{BB} & 0 & 0 & \mathcal{M}^* \\ 0 & \mathcal{L}_{AA} & \mathcal{L}_{AB} & 0 \\ 0 & \mathcal{L}_{BA} & \mathcal{L}_{BB} & 0 \\ \mathcal{M} & 0 & 0 & 0 \end{pmatrix} + \mathcal{O}(\lambda_\nu^4).$$

Measuring Entanglement

The entanglement is measured through the the negativity \mathcal{N} (an entanglement monotone)⁷.

$$\mathcal{N} = \frac{\text{Tr} |\rho^{T_A}| - 1}{2}.$$

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It is, to second order in λ

$$\mathcal{N}^{(2)} = -\frac{1}{2} \left(\mathcal{L}_{AA} + \mathcal{L}_{BB} - \sqrt{(\mathcal{L}_{AA} - \mathcal{L}_{BB})^2 + 4|\mathcal{M}|^2} \right).$$

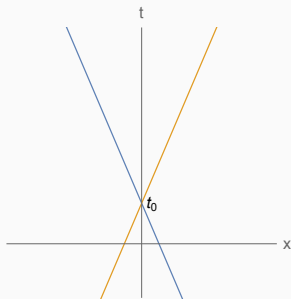
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Uniform Velocity

The detectors have opposite velocities (equal speeds) and pass by each other at $t = t_0$.

$$x_A(t) = v(t - t_0)$$

$$x_B(t) = -v(t - t_0).$$

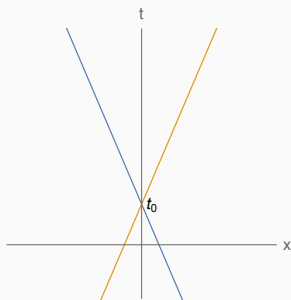


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This is a 1+1 dimensional problem, so it requires an infrared (IR) cutoff, Λ .

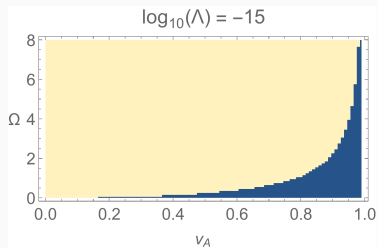
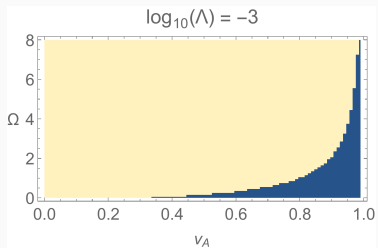
$$\mathcal{L}_{AA} = \frac{\lambda^2 \sigma^2}{2} (1 - v^2) \int_{\Lambda} \frac{dk}{|k|} e^{-\sigma^2 G_-(k)^2}$$
$$\mathcal{L}_{BB} = \frac{\lambda^2 \sigma^2}{2} (1 - v^2) \int_{\Lambda} \frac{dk}{|k|} e^{-\sigma^2 G_+(k)^2}$$

$$\mathcal{M} = \frac{\lambda^2 \sigma^2}{4} (1 - v^2) \int_{\Lambda} \frac{dk}{|k|} \operatorname{erf}(i\sigma|k|)$$
$$\times \left[e^{-2ikvt_0} e^{-\frac{\sigma^2}{2}(G_+(k)^2 + P_-(k)^2)} + e^{2ikvt_0} e^{-\frac{\sigma^2}{2}(G_-(k)^2 + P_+(k)^2)} \right]$$

where

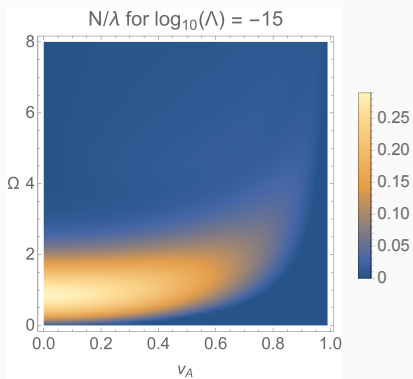
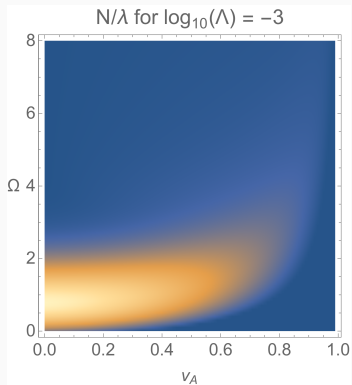
$$G_{\pm}(k) = |k| + (1 - v^2)\Omega \pm vk, \quad P_{\pm}(k) = |k| - (1 - v^2)\Omega \pm vk.$$

Negativity



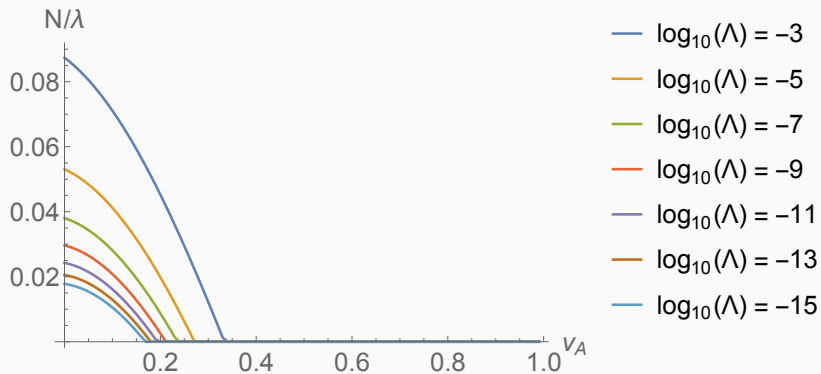
Yellow = Entanglement; Blue = No Entanglement

Negativity

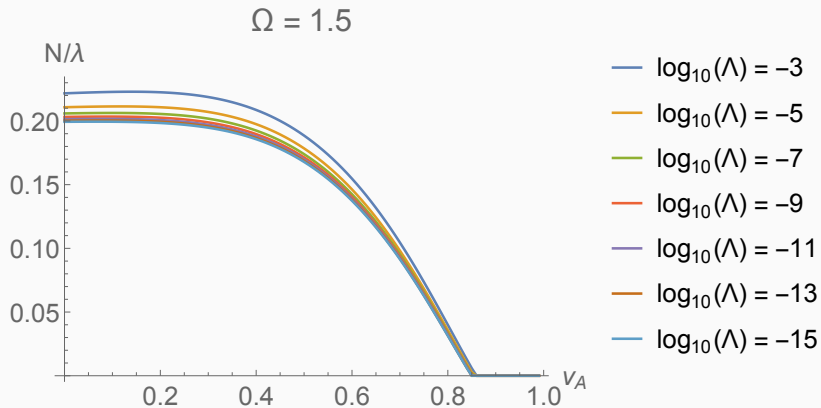


$$\Omega = 0$$

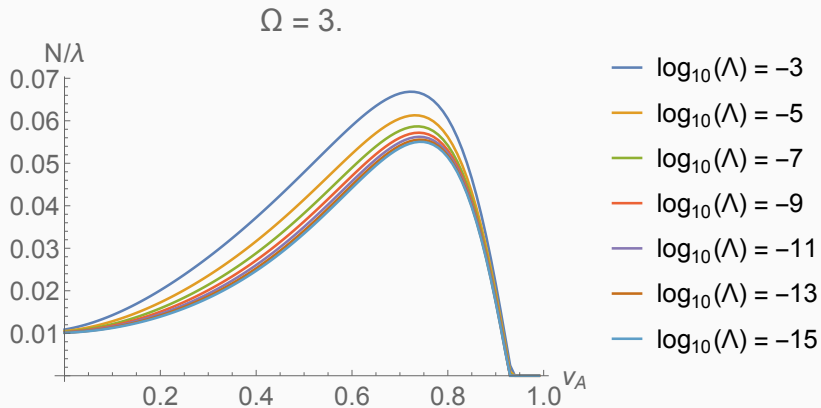
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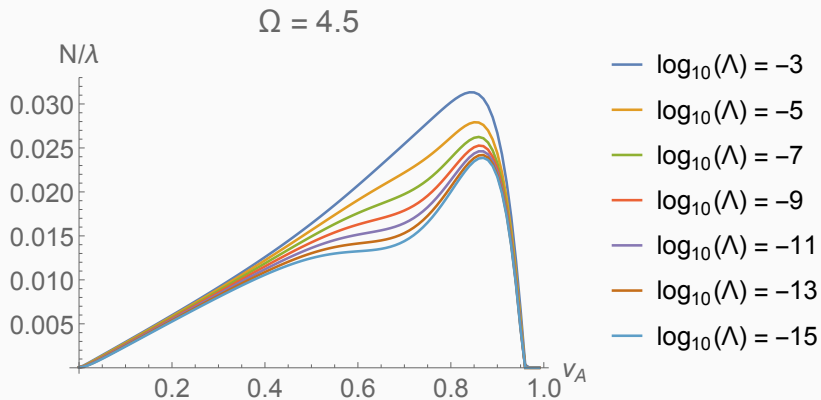
$$\Omega = 1.5$$



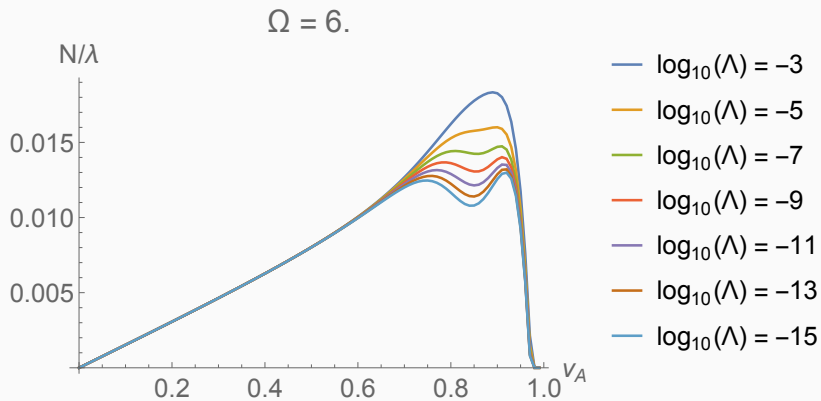
$$\Omega = 3$$



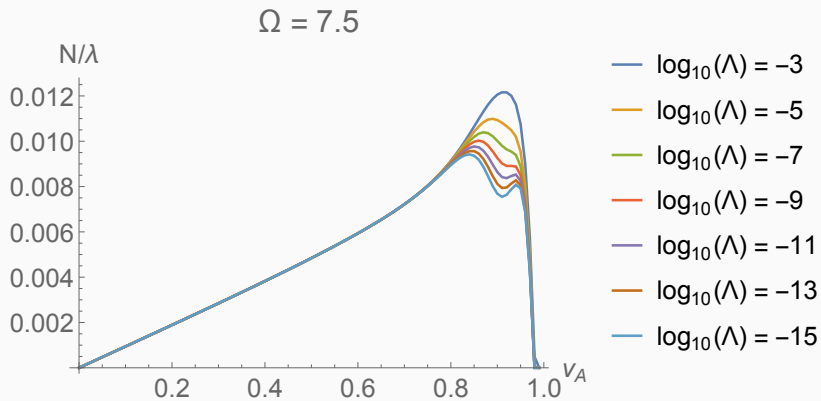
$$\Omega = 4.5$$



$$\Omega = 6$$



$$\Omega = 7.5$$



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- There is no entanglement for very large velocities.
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- As the energy gap increases more local maxima form in the velocity curves.
- A sufficiently small IR cutoff is required to observe this extra structure.