Electroweak Precision Measurements



Jens Erler





2017 CAP Congress — Testing Fundamental Symmetries —

Kingston, ON May 30, 2017

Advanced Series on Directions in High Energy Physics --- Vol. 14

PRECISION TESTS OF THE STANDARD ELECTROWEAK MODEL

Editor Paul Langacker



World Scientific

• Introduction:

Fundamental symmetries, standard model and beyond



- Introduction: Fundamental symmetries, standard model and beyond
- The weak mixing angle



- Introduction: Fundamental symmetries, standard model and beyond
- The weak mixing angle
- PV experiments



- Introduction: Fundamental symmetries, standard model and beyond
- The weak mixing angle
- PV experiments
- Beyond the Standard Model



- Introduction: Fundamental symmetries, standard model and beyond
- The weak mixing angle
- PV experiments
- Beyond the Standard Model
- Conclusions



Introduction: Fundamental symmetries, standard model and beyond

Quantum Lorentz transformations & translations Weinberg

- Quantum Lorentz transformations & translations Weinberg
- \Rightarrow gauge invariance for m = 0 & h = ±1 (long range force)

- Quantum Lorentz transformations & translations Weinberg
 gauge invariance for m = 0 & h = ±1 (long range force)
- \Rightarrow equivalence principle for m = 0 & h = ±2

- Quantum Lorentz transformations & translations Weinberg
- \Rightarrow gauge invariance for m = 0 & h = ±1 (long range force)
- \Rightarrow equivalence principle for m = 0 & h = ±2
- → local supersymmetry for $m = 0 \& h = \pm 3/2$

- Quantum Lorentz transformations & translations Weinberg
- \Rightarrow gauge invariance for m = 0 & h = ±1 (long range force)
- \Rightarrow equivalence principle for m = 0 & h = ±2
- \Rightarrow local supersymmetry for m = 0 & h = $\pm 3/2$
- \Rightarrow chiral symmetry for m = 0 & h = ±1/2

- Quantum Lorentz transformations & translations Weinberg
- \Rightarrow gauge invariance for m = 0 & h = ±1 (long range force)
- \Rightarrow equivalence principle for m = 0 & h = ±2
- \Rightarrow local supersymmetry for m = 0 & h = $\pm 3/2$
- \Rightarrow chiral symmetry for m = 0 & h = ±1/2
- ➡ CPT symmetry

- Quantum Lorentz transformations & translations Weinberg
- \Rightarrow gauge invariance for m = 0 & h = ±1 (long range force)
- \Rightarrow equivalence principle for m = 0 & h = ±2
- \Rightarrow local supersymmetry for m = 0 & h = $\pm 3/2$
- \Rightarrow chiral symmetry for m = 0 & h = ±1/2
- ➡ CPT symmetry
- no analog for $h = 0 \Rightarrow$ hierarchy problem

• P and CP for QED and pQCD by construction

- P and CP for QED and pQCD by construction
- P violation through chiral fermion representation

- P and CP for QED and pQCD by construction
- P violation through chiral fermion representation
- CP violation due to third chiral fermion generation

- P and CP for QED and pQCD by construction
- P violation through chiral fermion representation
- CP violation due to third chiral fermion generation
- P and T (CP) violation also through $\theta_{QCD} \epsilon^{\mu\nu\rho\sigma} G_{\mu\nu} G_{\rho\sigma}$

- P and CP for QED and pQCD by construction
- P violation through chiral fermion representation
- CP violation due to third chiral fermion generation
- P and T (CP) violation also through $\theta_{QCD} \epsilon^{\mu\nu\rho\sigma} G_{\mu\nu} G_{\rho\sigma}$
- → P a precision tool to study weak interaction (later)

- P and CP for QED and pQCD by construction
- P violation through chiral fermion representation
- CP violation due to third chiral fermion generation
- P and T (CP) violation also through $\theta_{QCD} \epsilon^{\mu\nu\rho\sigma} G_{\mu\nu} G_{\rho\sigma}$
- P a precision tool to study weak interaction (later)
- ➡ CP a possible tool to discover new physics (EDMs)

- P and CP for QED and pQCD by construction
- P violation through chiral fermion representation
- CP violation due to third chiral fermion generation
- P and T (CP) violation also through $\theta_{QCD} \epsilon^{\mu\nu\rho\sigma} G_{\mu\nu} G_{\rho\sigma}$
- P a precision tool to study weak interaction (later)
- \rightarrow CP a possible tool to discover new physics (EDMs)
- CPT a tool to access Planck scale physics (tomorrow)

• Baryon number: p decay, n- \overline{n} oscillations ($\Delta B = 2$)

- Baryon number: p decay, n- \overline{n} oscillations ($\Delta B = 2$)
- Lepton number: $0\nu\beta\beta$ -decay, $\mu^-Ti \rightarrow e^+Ca$, $K^+ \rightarrow \pi^-\mu^+e^+$

- Baryon number: p decay, n- \overline{n} oscillations ($\Delta B = 2$)
- Lepton number: $0\nu\beta\beta$ -decay, $\mu^-Ti \rightarrow e^+Ca$, $K^+ \rightarrow \pi^-\mu^+e^+$
- Charged lepton flavor #: $\mu^- \rightarrow e^-(\gamma), \tau \rightarrow 3\mu, H(Z) \rightarrow \mu e$

- Baryon number: p decay, n- \overline{n} oscillations ($\Delta B = 2$)
- Lepton number: $0\nu\beta\beta$ -decay, $\mu^-Ti \rightarrow e^+Ca$, $K^+ \rightarrow \pi^-\mu^+e^+$
- Charged lepton flavor #: $\mu^- \rightarrow e^-(\gamma), \tau \rightarrow 3\mu, H(Z) \rightarrow \mu e$
- Lepton universality: $B \rightarrow K\mu\mu < B \rightarrow Kee @ 2.6\sigma$ (LHCb)

- Baryon number: p decay, n- \overline{n} oscillations ($\Delta B = 2$)
- Lepton number: $0\nu\beta\beta$ -decay, $\mu^-Ti \rightarrow e^+Ca$, $K^+ \rightarrow \pi^-\mu^+e^+$
- Charged lepton flavor #: $\mu^- \rightarrow e^-(\gamma), \tau \rightarrow 3\mu, H(Z) \rightarrow \mu e$
- Lepton universality: $B \rightarrow K \mu \mu < B \rightarrow K ee @ 2.6\sigma (LHCb)$
- Flavor changing neutral currents: $b \rightarrow s\gamma$, $K^+ \rightarrow \pi^+ \nu \overline{\nu}$

- Baryon number: p decay, n- \overline{n} oscillations ($\Delta B = 2$)
- Lepton number: $0\nu\beta\beta$ -decay, $\mu^-Ti \rightarrow e^+Ca$, $K^+ \rightarrow \pi^-\mu^+e^+$
- Charged lepton flavor #: $\mu^- \rightarrow e^-(\gamma), \tau \rightarrow 3\mu, H(Z) \rightarrow \mu e$
- Lepton universality: $B \rightarrow K \mu \mu < B \rightarrow K ee @ 2.6\sigma (LHCb)$
- Flavor changing neutral currents: $b \rightarrow s\gamma$, $K^+ \rightarrow \pi^+ \nu \overline{\nu}$
- Flavor changing charged currents: $B \rightarrow D \tau v$ (3.9 σ high @ BaBar, Belle, LHCb)

The weak mixing angle $(sin^2\theta_W)$

The weak mixing angle $(\sin^2\theta_W)$

- mixing of $SU(2)_L \times U(1)_Y$
- $W^{\pm} = (W^{\dagger} \mp i W^2) / \sqrt{2}$
- $Z^0 = \cos\theta_W W^3 \sin\theta_W B$

$$A = \sin\theta_{W}W^{3} + \cos\theta_{W}B$$



- $M_W = \frac{1}{2} g v = \cos\theta_W M_Z$
- $\sin^2\theta_W = g'^2/(g^2 + g'^2) = I M_W^2/M_Z^2$ (tree level)

Running weak mixing angle

results and prospects



Running weak mixing angle

results and prospects






 $\begin{array}{c|c} e_{L,R}^{-} & e_{L,R}^{-} \\ & & & \\ & & & \\ & & & \\ f & & & f \end{array}$



atomic parity violation (most precise: Boulder & Paris)



- atomic parity violation (most precise: Boulder & Paris)
- polarized deep inelastic scattering (ēDIS) (SLAC-E122, PVDIS, SoLID)



- atomic parity violation (most precise: Boulder & Paris)
- polarized deep inelastic scattering (ēDIS) (SLAC-E122, PVDIS, SoLID)
- polarized Møller scattering (SLAC-E158 & MOLLER)



- atomic parity violation (most precise: Boulder & Paris)
- polarized deep inelastic scattering (ēDIS) (SLAC-E122, PVDIS, SoLID)
- polarized Møller scattering (SLAC-EI58 & MOLLER)
- polarized elastic ēp (ēC?) scattering (Qweak & Mainz-P2)



- atomic parity violation (most precise: Boulder & Paris)
- polarized deep inelastic scattering (ēDIS) (SLAC-E122, PVDIS, SoLID)
- polarized Møller scattering (SLAC-EI58 & MOLLER)
- polarized elastic ēp (ēC?) scattering (Qweak & Mainz-P2)
- PV in isotope chains (Mainz & KVI Groningen)



- atomic parity violation (most precise: Boulder & Paris)
- polarized deep inelastic scattering (ēDIS) (SLAC-E122, PVDIS, SoLID)
- polarized Møller scattering (SLAC-EI58 & MOLLER)
- polarized elastic ēp (ēC?) scattering (Qweak & Mainz-P2)
- PV in isotope chains (Mainz & KVI Groningen)
- PV in single trapped Ra ions? (KVI Groningen)

The Low-Energy (Fermi) Limit



Effective couplings





Effective couplings

• NC couplings: $g^{ef}_{AV} e \gamma^{\mu}\gamma^{5} e f \gamma_{\mu} f$ $g^{ef}_{VA} e \gamma^{\mu} e f \gamma_{\mu}\gamma^{5} f$

Effective couplings



- NC couplings: $g^{ef}_{AV} e \gamma^{\mu}\gamma^{5} e f \gamma_{\mu} f$
- $|g^{ef}_{AV}| = \frac{1}{2} 2 |Q_f| \sin^2\theta_W$

 $|g^{ef}_{VA}| = \frac{1}{2} - 2 \sin^2\theta_{W}$

 g^{ef} VA $e^{\gamma \mu} e^{f} \gamma_{\mu} \gamma^{5} f$

$\begin{array}{c|c} e_{L,R}^{-} & e_{L,R}^{-} \\ & & & \\ & & & \\ & & & \\ f & & & f \end{array}$

Effective couplings

- NC couplings: $g^{ef}_{AV} e \gamma^{\mu}\gamma^{5} e f \gamma_{\mu} f$ $g^{ef}_{VA} e \gamma^{\mu} e f \gamma_{\mu}\gamma^{5} f$
- $|g^{ef}_{AV}| = \frac{1}{2} 2 |Q_f| \sin^2\theta_W$

 $|g^{ef}_{VA}| = \frac{1}{2} - 2 \sin^2\theta_{W}$

• $f = e \implies |g^{ee}_{AV}| = \frac{1}{2} - 2 \sin^2 \theta_{W} \ll 1$

$\begin{array}{ccc} e_{L,R}^{-} & e_{L,R}^{-} \\ & & & \\ & & & \\ & & & \\ f & & & f \end{array}$

Effective couplings

- NC couplings: $g^{ef}_{AV} e \gamma^{\mu}\gamma^{5} e f \gamma_{\mu} f g^{ef}_{VA} e \gamma^{\mu} e f \gamma_{\mu}\gamma^{5} f$
- $|g^{ef}_{AV}| = \frac{1}{2} 2 |Q_f| \sin^2\theta_W$

 $|g^{ef}_{VA}| = \frac{1}{2} - 2 \sin^2\theta_{W}$

- $f = e \implies |g^{ee}_{AV}| = \frac{1}{2} 2 \sin^2 \theta_{W} \ll 1$
- → in SM: enhanced sensitivity to $sin^2\theta_W$ (compete with ultra-precise Z-pole determinations)

$\begin{array}{c|c} e_{L,R}^{-} & e_{L,R}^{-} \\ & & \\ & & \\ & & \\ f & & \\ & & f \end{array}$

Effective couplings

- NC couplings: $g^{ef}_{AV} e \gamma^{\mu}\gamma^{5} e f \gamma_{\mu} f g^{ef}_{VA} e \gamma^{\mu} e f \gamma_{\mu}\gamma^{5} f$
- $|g^{ef}_{AV}| = \frac{1}{2} 2 |Q_f| \sin^2\theta_W$

 $|g^{ef}_{VA}| = \frac{1}{2} - 2 \sin^2\theta_{W}$

- $f = e \implies |g^{ee}_{AV}| = \frac{1}{2} 2 \sin^2 \theta_{W} \ll 1$
- → in SM: enhanced sensitivity to $sin^2\theta_W$ (compete with ultra-precise Z-pole determinations)
- ➡ Beyond SM: enhanced sensitivity to ∧_{new}



 g^{eq}_{AV} (coherent) Stark induced-Z interference amplitude dominant (spin-independent)



- g^{eq}_{AV} (coherent) Stark induced-Z interference amplitude dominant (spin-independent)
- Q_W(¹³³Cs) ~ 0.6% (incl. theory) c.s. Wood et al. 1997



- g^{eq}_{AV} (coherent) Stark induced-Z interference amplitude dominant (spin-independent)
- Q_W(¹³³Cs) ~ 0.6% (incl. theory) c.s. Wood et al. 1997
- spin-dependent nuclear anapole moment through difference in hyperfine transitions



 $p \rightarrow Q \qquad x_p$

 deuterium target (isoscalar and simple nucleus)



- deuterium target (isoscalar and simple nucleus)
- $A_{LR} \equiv \sigma_L \sigma_R / \sigma_L + \sigma_R \propto Q^2$



- deuterium target (isoscalar and simple nucleus)
- $A_{LR} \equiv \sigma_L \sigma_R / \sigma_L + \sigma_R \propto Q^2$
- large $Q^2 \Longrightarrow A_{LR}(d) \sim 10^{-4}$



- deuterium target (isoscalar and simple nucleus)
- $A_{LR} \equiv \sigma_L \sigma_R / \sigma_L + \sigma_R \propto Q^2$
- large $Q^2 \Longrightarrow A_{LR}(d) \sim 10^{-4}$
- large $y \Longrightarrow g^{eq}_{AV}$ and g^{eq}_{AV}



- deuterium target (isoscalar and simple nucleus)
- $A_{LR} \equiv \sigma_L \sigma_R / \sigma_L + \sigma_R \propto Q^2$
- large $Q^2 \Longrightarrow A_{LR}(d) \sim 10^{-4}$
- large $y \Longrightarrow g^{eq}_{AV}$ and g^{eq}_{AV}
- Q_q weighted (γ -Z interference)



• $A_{LR} \sim 3 \times |0^{-8}|$



- $A_{LR} \sim 3 \times 10^{-8}$
- purely leptonic



- $A_{LR} \sim 3 \times 10^{-8}$
- purely leptonic
- \rightarrow very clean theoretically



- $A_{LR} \sim 3 \times 10^{-8}$
- purely leptonic
- \rightarrow very clean theoretically
- → ultra-high precision



- $A_{LR} \sim 3 \times 10^{-8}$
- purely leptonic
- \rightarrow very clean theoretically
- → ultra-high precision
- need at least one 2-loop electroweak calculation S. Barkanova & A. Aleksejevs (Memorial University of Newfoundland)





• $A_{LR} \propto Q_W(p) + Q^2 B(Q^2, \theta)$



- $A_{LR} \propto Q_W(p) + Q^2 B(Q^2, \theta)$
- **Qweak**: $Q^2 = 0.025 \text{ GeV}^2$



- $A_{LR} \propto Q_W(p) + Q^2 B(Q^2, \theta)$
- **Qweak:** $Q^2 = 0.025 \text{ GeV}^2$
 - extrapolation to $Q^2 = 0$



- $A_{LR} \propto Q_W(p) + Q^2 B(Q^2, \theta)$
- **Qweak:** $Q^2 = 0.025 \text{ GeV}^2$
 - extrapolation to $Q^2 = 0$
 - large γ-Z box Gorchtein, Horowitz, Ramsey-Musolf; Rislow, Carlson; Hall, Blunden, Melnitchouk, Thomas, Young


Polarized elastic scattering

- $A_{LR} \propto Q_W(p) + Q^2 B(Q^2, \theta)$
- Qweak: $Q^2 = 0.025 \text{ GeV}^2$
 - extrapolation to $Q^2 = 0$
 - ¹°0.1' V^d large γ -Z box Gorchtein, 0.2 0.3 0.4 0.5 0.0 0.1 $Q^2 [\text{GeV/c}]^2$ Horowitz, Ramsey-Musolf; **Rislow, Carlson; Hall, Blunden, Melnitchouk, Thomas, Young**

This Experiment

HAPPEX

SAMPLE

→ SM (prediction)

PVA4

G0

6

Ш

θ,

 $B(O^2)$

 $Q_{W}^{p} + Q^{2}$

0.4

0.

0.2

Data Rotated to the Forward-Angle Limit



0.6

Polarized elastic scattering

- $A_{LR} \propto Q_W(p) + Q^2 B(Q^2, \theta)$
- **Qweak:** $Q^2 = 0.025 \text{ GeV}^2$
 - extrapolation to $Q^2 = 0$ ightarrow
 - Ч° 0.1^г large γ -Z box Gorchtein, ightarrow0.2 0.3 0.5 0.0 0.10.4 $Q^2 [\text{GeV/c}]^2$ Horowitz, Ramsey-Musolf; **Rislow, Carlson; Hall, Blunden, Melnitchouk, Thomas, Young**

This Experiment

HAPPEX

SAMPLE

→ SM (prediction)

PVA4 G0

6

Ш

 $Q_{W}^{p} + Q^{2} B(Q^{2}, \theta)$

0.4

0.

0.2

Data Rotated to the Forward-Angle Limit

- $P2: Q^2 = 0.0045 \text{ GeV}^2 (A_{LR} \sim 10^{-8})$
 - γ-Z box correction (error) factor of 8 (5) smaller

0.6

Beyond the Standard Model

New physics discrimination



PV (axial)-electron (vector)-quark couplings







Compositeness Scales

 $[2 g^{eu} - g^{ed}]_{AV}$





Summary

	precision	$sin^2\theta_W$	Λ_{new}
APV Cs-133	0.58 %	0.0019	32.3 TeV
E158	14 %	0.0013	17.0 TeV
Qweak I	19 %	0.0030	17.0 TeV
PVDIS	4.5 %	0.0051	7.6 TeV
Qweak final	4.5 %	0.0008	33 TeV
SoLID	0.6 %	0.00057	22 TeV
MOLLER	2.3 %	0.00026	39 TeV
P2	2.0 %	0.00036	49 TeV
PVES ¹² C	0.3 %	0.0007	49 TeV
APV ²²⁵ Ra	0.5 %	0.0018	34 TeV
APV ²¹³ Ra / ²²⁵ Ra	0.1 %	0.0037	I6TeV

Recent Reviews

Krishna Kumar, Sonny Mantry, William Marciano and Paul Souder Annu. Rev. Nucl. Part. Sci. 63 (2013) 237–67

> Jens Erler and Shufang Su Prog. Part. Nucl. Phys. 71 (2013) 119–149

> > Jens Erler and Ayres Freitas Particle Data Group (2014)

Jens Erler, Charles Horowitz, Sonny Mantry and Paul Souder Annu. Rev. Nucl. Part. Sci. (2014)