Optomechanical micro-macro entanglement

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SCHRÖDINGER'S CAT IS

An idea





Summary



Background and Motivation

Schrödinger's thought experiment



- *Quantum picture* —> Simultaneous perception of a live and a dead cat. Is it possible?

$$\frac{1}{\sqrt{2}}(|\phi_{NO-decay}\rangle|\psi_{ALIVE-cat}\rangle+|\phi_{YES-decay}\rangle|\psi_{DEAD-cat}\rangle)$$



Possible way out



Elijah Wood and Daniel Radcliffe

Hillary and Bill Clinton

• A pertinent question:" How can we demonstrate macroscopic superposition and micro-macro entanglement?"



Possible way out

!!!!PHOTONS!!!!

- Microscopic state _____ single photon level
- Macroscopic state —> millions or hundreds of millions of photons

So, why not amplify the microscopic quantum state. But?

- Problem —> It is very hard to detect and characterize such a state as the measurements need to have extremely high resolution.
- Solution By local operations, convert the macroscopic state back to microscopic state to demonstrate entanglement.

Scheme

Amplification by displacement



Entanglement at the end proves micro-macro entanglement in the intermediate step.

Micro-macro correlations • Delocalized single-photon state

$$|\psi\rangle = \frac{1}{\sqrt{2}} \left(|0\rangle_A \otimes |1\rangle_B + |1\rangle_A \otimes |0\rangle_B\right)$$

After displacement

$$|\psi_D\rangle = \frac{1}{\sqrt{2}} \left(|0\rangle_A \otimes \hat{D}(\alpha) |1\rangle_B + |1\rangle_A \otimes \hat{D}(\alpha) |0\rangle_B \right)$$

• After Alice's measurement of the X quadrature

$$|\psi_{\rm B}\rangle = \frac{1}{\sqrt{2}} \left(\psi_0(X_{\rm A})\hat{D}(\alpha) |1\rangle_{\rm B} + \psi_1(X_{\rm A})\hat{D}(\alpha) |0\rangle_{\rm B} \right)$$

- Alice's mode microscopic state quadrature measurement.
- Bob's mode macroscopic phase-space displacement phase-space un-displacement quadrature measurement.

(A. I. Lvovsky, R. Ghobadi, A. Chandra, A. S. Prasad and C. Simon, "Observation of micro–macro entanglement of light", Nature Physics 9, 541 (2013)).



An Idea

!!Opto-mechanics!!



Optomechanical storage and retrieval • The basic opto-mechanical Hamiltonian is :

 $H = \hbar \Delta a^{\dagger} a + \hbar \omega_m b^{\dagger} b + \hbar g_0 a^{\dagger} a (b + b^{\dagger})$

- Effective beam splitter Hamiltonian $H_{
 m eff} = g(a^{\dagger}b + ab^{\dagger})$
- The resulting equations of motion are : *a* = −κ*a* − *igb* + √2κ *a_{in} b* = −γ*b* − *iga* + √2γ *b_{in}*
 The input-output relation for the cavity is
 - $a_{out}(t) = -a_{in}(t) + \sqrt{2\kappa} a(t).$

Effects of finite y and mechanical initial state



Entanglement in the final state as a function of the opto-mechanical coupling parameter $y = e^{-G\tau}$ for different values of the initial mechanical phonon number N_{in} .

Entanglement dependence on mechanical noise.



Entanglement in the final state as a function of the mechanical noise parameter $x = \gamma/G$, for different values of the bath mean phonon number *Nth*.

Test of collapse models.

Quantum-gravity induced wavefunction collapse.

• Trampoline resonators Mass = 500 ng, ω_m = 10 kHz, Mechanical quality factor = 10⁶ Temperature = 1 mK Environmentally induced decoherence time scale = 7.6 ms Quantum gravity induced collapse model = 95 µs - 240 µs



Summary

- Proposal to create and observe an opto-mechanical micro-macro entangled state.
- Studied the most important experimental imperfections (phase noise, mechanical decoherence and photon loss), and found the parameter regime where the demonstration of entanglement is possible.
- Found out that the realization is quite possible using state-of-theart opto-mechanical systems and entangled light sources.
- Can test unconventional collapse models, e.g. gravitationally induced collapse.

THANK YOU

Effects of finite y and mechanical initial state



Entanglement in the final state as a function of the opto-mechanical coupling parameter $y = e^{-G\tau}$ for different values of the initial mechanical phonon number N_{in} . In all cases y has to be below a certain threshold value for entanglement to be observable, where the value of the threshold depends on *Nin*. The figure also includes the effect of other imperfections, the relevant parameter values are $x = \gamma/G = 0.01$ and *Nth* = 10 (mechanical noise), $\eta 1 = \eta 2 = \eta c = 0.8$ (losses). The photon number corresponding to the displacement is $ND = |\alpha|^2 = 5000$, and the squeezing parameter is r = 0.55

• :

Entanglement dependence on mechanical noise.



Entanglement in the final state as a function of the mechanical noise parameter $x = \gamma/G$, for different values of the bath mean phonon number *Nth*. The values of the other parameters are *ND* = = 5000, *r* = 0.5, *y* = 0.1, and $\eta_1 = \eta_2 = \eta_c = 0.8$ (losses).

Photon loss



Entanglement in the final state as a function of η_1 for different values of η_2 . Here 1- η_1 and 1 - η_2 are the photon loss before and after the optomechanical system respectively. The values of the other parameters are *ND* = $|\alpha|^2 = 5000$, r = 0.5, y = 0.1, *Nin* = 10, , x = 0.01, *Nth* = 10, and $\eta_c = 0.8$.

Implementation

- ω_m = 3.7 GHz, κ = 500 MHz, and γ = 35 kHz
 T=2 K
 g ≈ 40 MHz, G = g²/κ ≈ 3.2 MHz
 x= γ/G ≈ 0.01
 τ ≈ 100 ns
- $y = e^{-G\tau} \approx 0.1$.