

# $\nu_e$ Cross Section Measurements at the T2K Off-Axis Near Detector

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# Overview

- 1 Introduction
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- 3 Future Work
- 4 Conclusion

## Neutrino Oscillation

- Neutrino oscillation arises from a mixture between the flavour and mass eigenstates of neutrinos, linked by the Maki-Nakagawa-Sakata (MNS) matrix  $U$ :

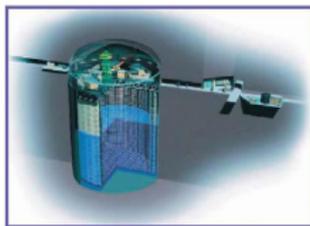
$$|\nu_\alpha\rangle = \sum_{k=1}^3 U_{\alpha k}^* |\nu_k\rangle.$$

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}.$$

- As a neutrino propagates, the quantum mechanical phases of the three different mass states advance at slightly different rates. This results in a changing mixture of mass states as the neutrino travels.

# T2K Experiment

- T2K is a long baseline neutrino oscillation experiment in Japan, that targets the measurement of the mixing angle ( $\theta_{13}$ ), as well as a precision measurement for the mass difference ( $\Delta m_{23}^2$ ) and their mixing angle ( $\theta_{23}$ ).



**Super-Kamiokande**  
(ICRR, Univ. Tokyo)

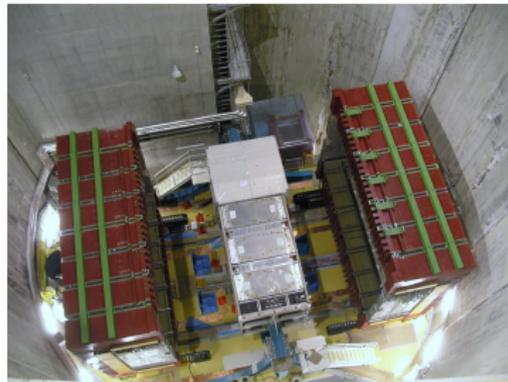
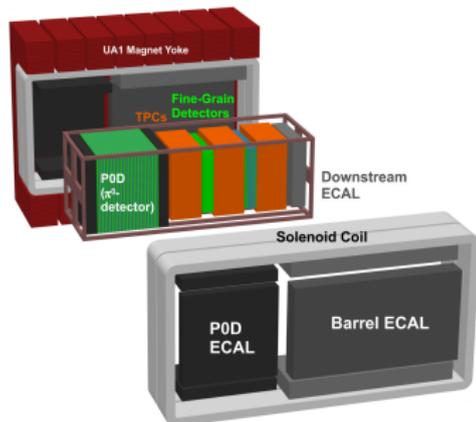


**J-PARC Main Ring**  
(KEK-JAEA, Tokai)



## ND280 Overview

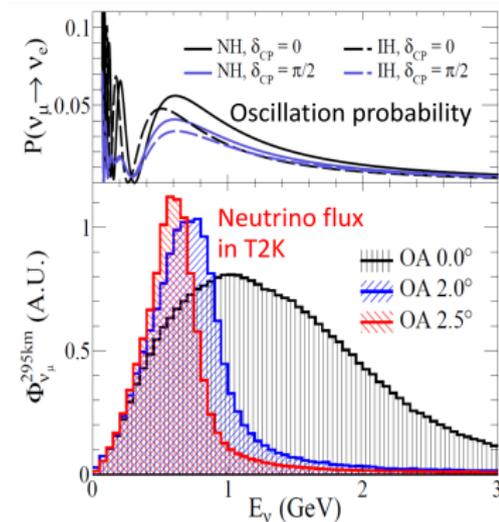
- The experiment uses two detectors; a near detector at 280 m from the neutrino production (in Tokai), and the far detector at 295 km Super-Kamiokande (SK).
- The near detector complex consists of an on-axis beam profile monitor INGRID, and an off-axis detector ND280.



# Neutrino Cross Section

- Neutrino cross section is a very important quantity in itself.
- Cross section measurement plays a vital role in the oscillation analysis, as the oscillation probability “ $P(E_\nu)$ ” is inferred from the number of event (oscillated neutrinos at the far detector) “ $N$ ”, efficiency “ $\epsilon$ ”, cross section “ $\sigma$ ” and the flux “ $\phi$ ”.

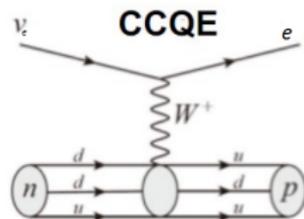
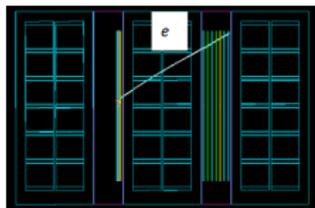
$$N = \int \epsilon \sigma(E_\nu) \phi(E_\nu) P(E_\nu) dE_\nu$$



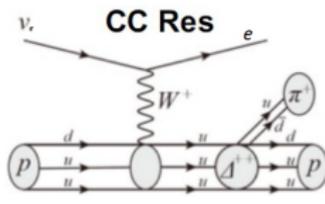
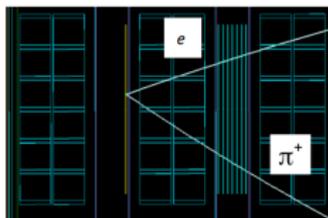
[T.Dorigo]

# $\nu_e$ CC-Inclusive interaction example

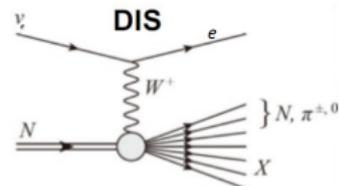
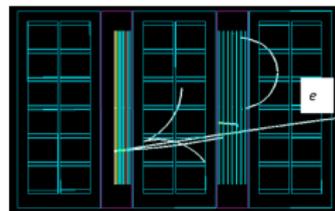
CC0 $\pi$



CC1 $\pi^+$

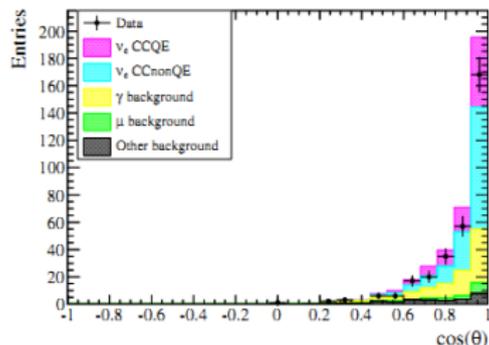
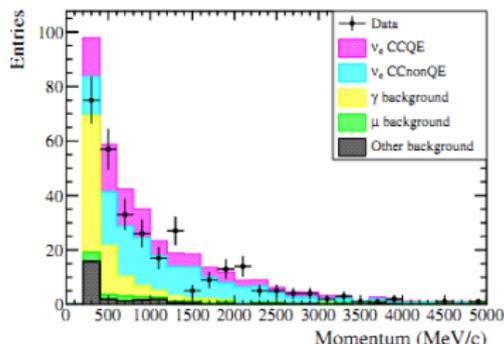


CCother



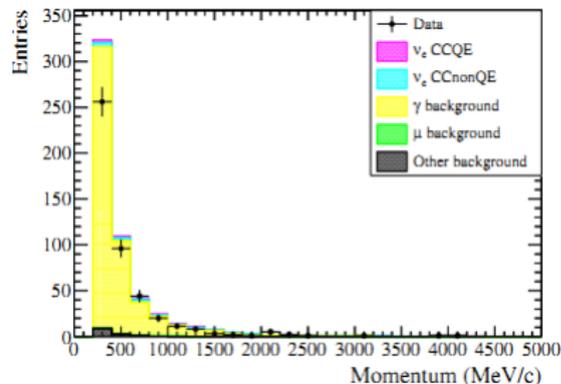
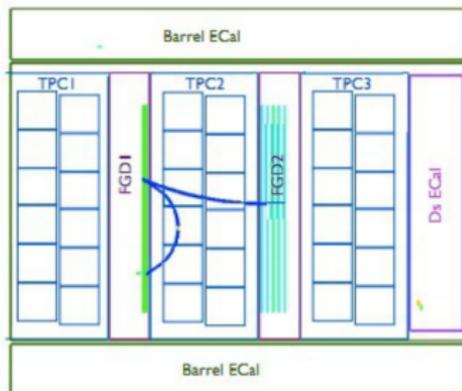
# Differential Cross Section

- select  $\nu_e$  events, e.g: highest momentum negative track, starts at FGD1, TPC and ECAL PIDs
- We measure the Differential Cross Section w.r.t electron momentum "P(e)", electron angle " $\cos(\Theta_e)$ " or the four momentum transfer " $Q_e^2$ ".
- $N_{t_k} = T \phi \int \langle \frac{\partial \sigma}{\partial X} dX \rangle_\phi$   
 with:  
 $N_{t_k}$  : nb of true interaction at bin k  
 T : nb of target nucleons  
 $\phi$  : total flux  
 X : refers to P(e),  $\cos(\Theta_e)$  or  $Q_e^2$
- Total Cross Section  $\langle \sigma_\phi \rangle = \frac{N_{total}}{T \phi}$   
 with  $N_{total}$  : total sum of bins  $t_k$



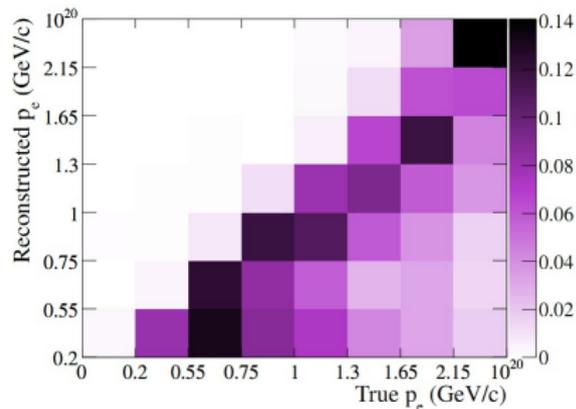
## Gamma Background

- The main background is photon conversion into  $e^+ e^-$  pair  
 $\gamma \rightarrow e^+ e^-$
- Background could be suppressed by rejecting events having tracks with start position not in FGD, or if a particle with mass  $< 100$  MeV was found beside the track.



# Bayesian Unfolding (Correction for detectors inefficiencies and mis-reconstruction.)

- $$N_{t_k}^{m+1} = \frac{1}{\epsilon} \sum_j P_m(t_k|r_j)(N_{r_j}^{meas} - B_{r_j})$$



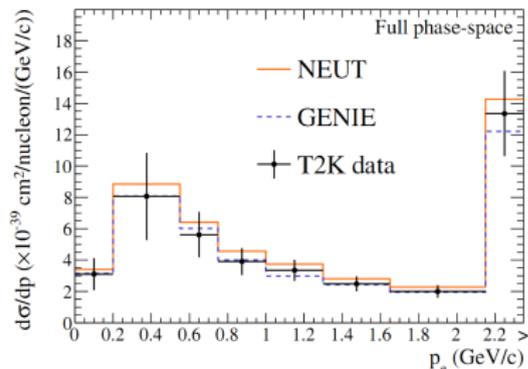
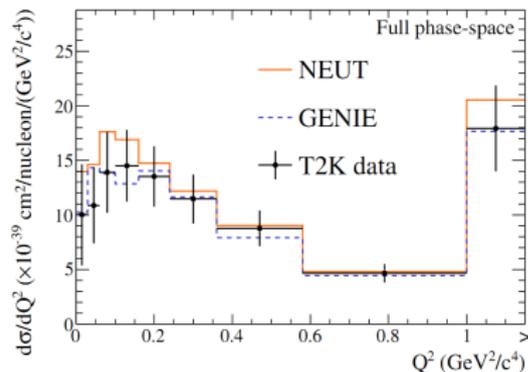
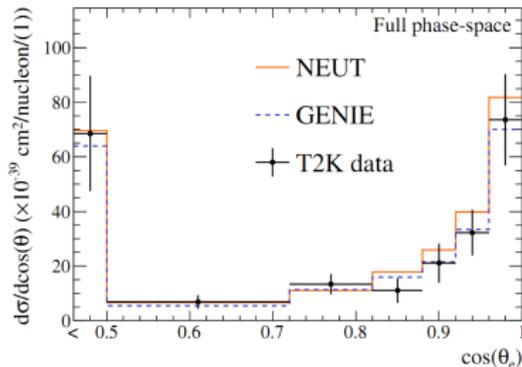
[Smearing Matrix  $P(r_j|t_k)$ ]

with :

- $N_{t_k}$  : Estimator for true signal event in bin k
- $\epsilon_{t_k}$  : Reconstruction efficiency
- $P_m(t_k|r_j)$  : Probability at iteration m to have a true event at bin k given a reconstructed event at bin j
- $B_{r_j}$  : Background

# T2K $\nu_e$ CC-Inclusive XSec results

- $Q_e^2$  Represent the main results as it is the most interesting kinematic variable and has the smallest systematic error.



## T2K $\nu_e$ CC-Inclusive XSec results II

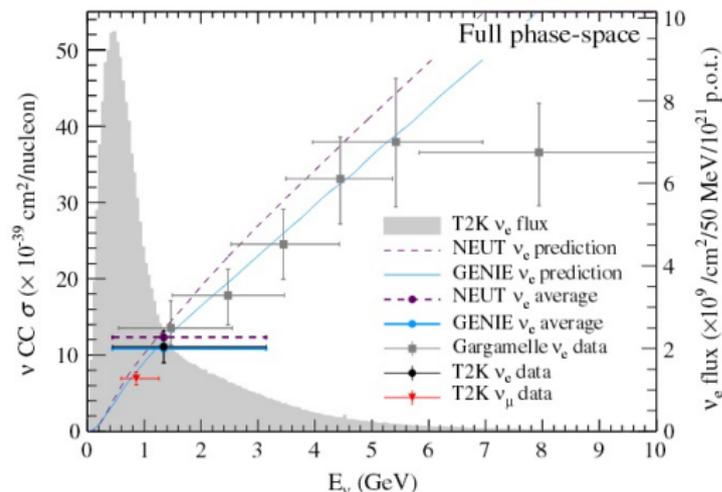
The measured cross section flux average:

$$1.11 \pm 0.20 \times 10^{-38} \text{ cm}^2 / \text{nucleon}$$

is in agreement with:

$$\text{NEUT } (1.23 \times 10^{-38} \text{ cm}^2 / \text{nucleon})$$

$$\text{and GENIE } (1.07 \times 10^{-38} \text{ cm}^2 / \text{nucleon})$$

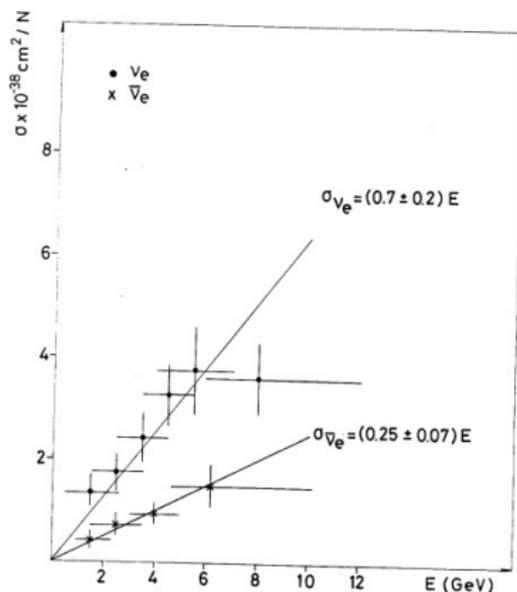


[K.Abe2014]

## $\bar{\nu}_e$ Cross Section Measurement

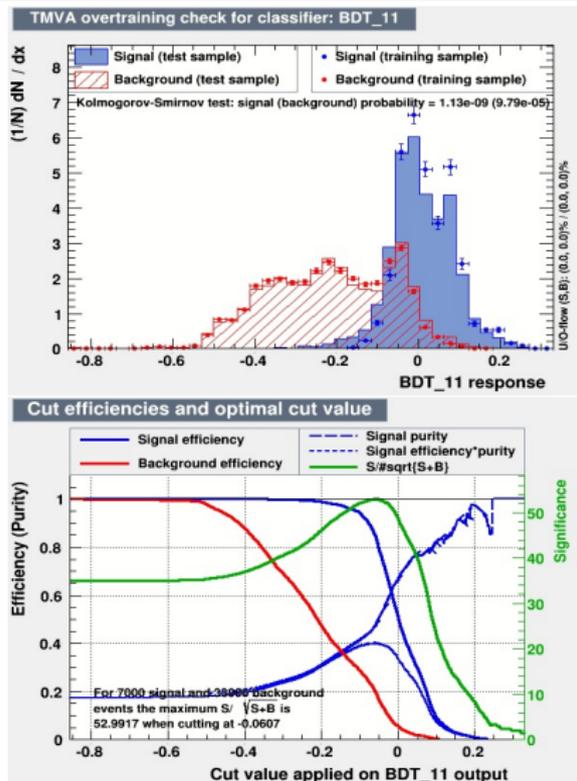
Measurement of the  $\bar{\nu}_e$  Charged Current Inclusive Cross Section using Bayesian Unfolding Technique.

- First time to be measured at T2K.
- Latest published results are from the Gargamelle experiment (1970's).
- Use of machine learning algorithms to classify event and eliminate backgrounds.



[J.Blietschau1978]

# MVA PID



Developed a Multivariate (Boosted Decision Tree) tool for track identification.

## Motivation:

- Look for ways of better identifying gamma and proton background events.
- Use information from FGDs, TPCs, and ECALs.
- Boosted Decision Tree (Root TMVA).
- May help in increasing the selection efficiency and signal purity (not tested on our case yet).
- May help in finding hidden pattern in our data.

## Conclusion

- Neutrino cross section is an important and interesting quantity to be measured.
- T2K  $\nu_e$  CC-Inclusive measured cross section ( $1.11 \pm 0.20 \times 10^{-38} \text{ cm}^2 / \text{nucleon}$ ) is in agreement with that predicated by neutrino event generators like NEUT and GENIE.
- Next step is to measure the  $\bar{\nu}_e$  cross section.
- The use of a machine learning algorithm in particle identification looks promising.

## References



[PDG2010] Review of Particle Physics, Particle Data Group, Journal of Physics G, Nuclear and Particle Physics, Volume 37 Number 7A Article 075021 [July 2010]



[L.Southwell2014] Selecting  $\nu_e$  charged-current events in the near detector at T2K, Luke Southwell



[B.Jamieson2014] T2K Results Blair Jamieson for the T2K Collaboration HSQCD2014



[J.Blietschau1978] Total cross sections for  $\nu_e$  and  $\bar{\nu}_e$  interactions and search for neutrino oscillations and decay, Nuclear Physics B, Elsevier, 1978, 133, pp.205-219 J. Blietschau et al.



[PDG2013] Particle Data Group, Neutrino Cross Section Measurements [2013]



[B.Smith2014] Measurement of the  $\nu_e$  CC Inclusive Cross-section on Carbon Using Data From Runs 1 - 4 , B. Smith [August 30, 2014]



[IntrinsicNueT2K2014] Measurement of the intrinsic electron neutrino component in the T2K neutrino beam with the ND280 detector, the T2K Collaboration [March 11, 2014]



[T.Dorigo] Electron Neutrino In T2K, T.Dorigo [Apr 4, 2013]



Measurement of the Inclusive Electron Neutrino Charged Current Cross Section on Carbon with the T2K Near Detector, K. Abe et al. (T2K Collaboration) Phys. Rev. Lett. 113, 241803 [December 2014]

# T2K Collaboration

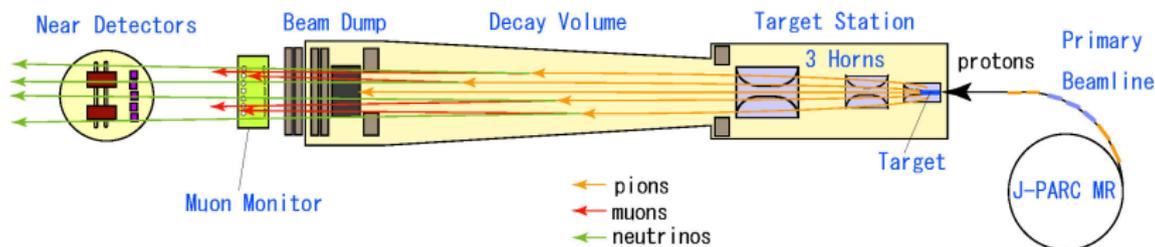


**The T2K collaboration** includes about 500 physicists from 11 countries (Canada, France, Germany, Italy, Japan, Poland, Russia, Spain, Switzerland, UK, USA).

# Thank You

# Backup Slides

## Beam Production



- Produced by (J-PARC), protons are accelerated up to 31 GeV/c, extracted in 5  $\mu$ s long spills with a repetition rate that of 2.6 s.
- The protons strike a 91.4 cm long graphite target, producing hadrons, mainly pions and kaons.
- To produce mainly  $\nu_\mu$ , the charged particles can then be focused by a series of three magnetic horns operating at 250 kA before entering a 96 m long decay volume.
- Most of the surviving charged particles are stopped by the beam dump at the end of the decay volume.
- A muon monitor (MUMON) measures the profile of high energy muons not stopped by the beam dump, monitoring the stability of the beam intensity and the direction of the beam.

# Beam Production II

Different decay modes with their corresponding branching ratios are given below.

$$\pi^- \longrightarrow \mu^- + \bar{\nu}_\mu \quad (\sim 99.98\%). \quad (1)$$

$$K^- \longrightarrow \mu^- + \bar{\nu}_\mu \quad (\sim 63.55\%). \quad (2)$$

$$K^- \longrightarrow \pi^- + \pi^0 \quad (\sim 20.66\%). \quad (3)$$

$$K^- \longrightarrow \pi^- + \pi^- + \pi^+ \quad (\sim 5.59\%). \quad (4)$$

$$K^- \longrightarrow e^- + \bar{\nu}_e + \pi^0 \quad (\sim 5.07\%). \quad (5)$$

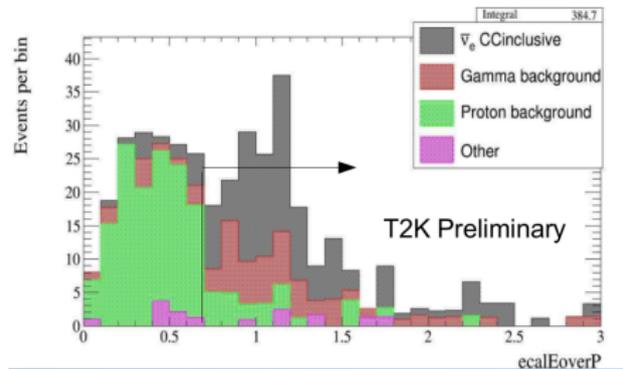
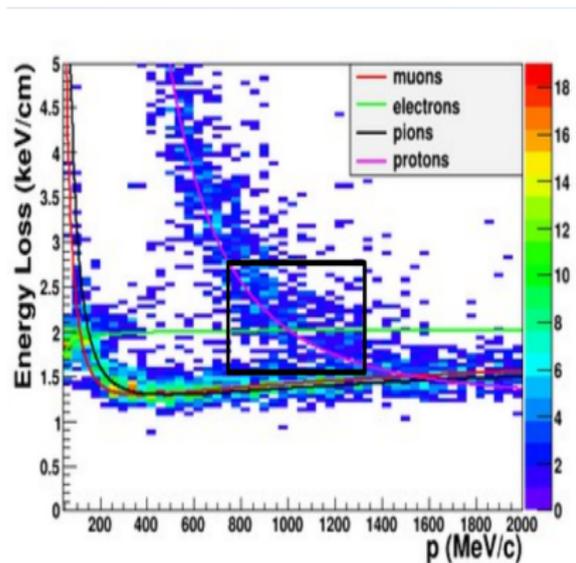
$$K^- \longrightarrow \mu^- + \bar{\nu}_\mu + \pi^0 \quad (\sim 3.35\%). \quad (6)$$

$$K^- \longrightarrow \pi^- + \pi^0 + \pi^0 \quad (\sim 1.76\%). \quad (7)$$

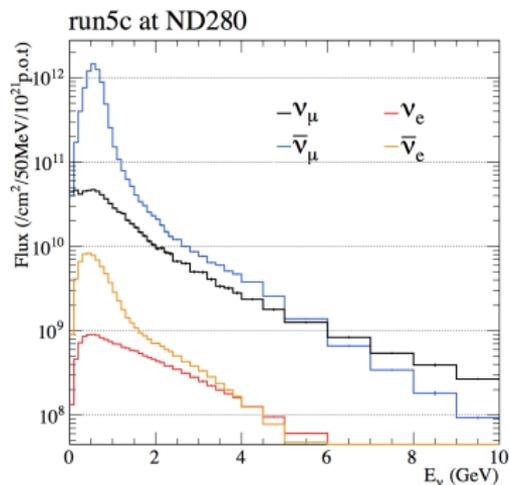
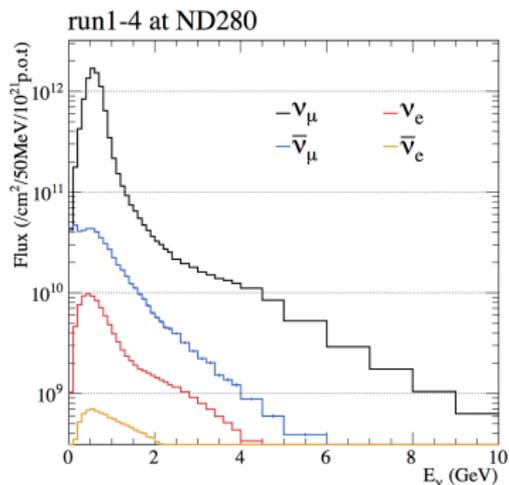
$$\mu^- \longrightarrow e^- + \bar{\nu}_e + \nu_\mu \quad (\sim 100\%). \quad (8)$$

[PDG2010]

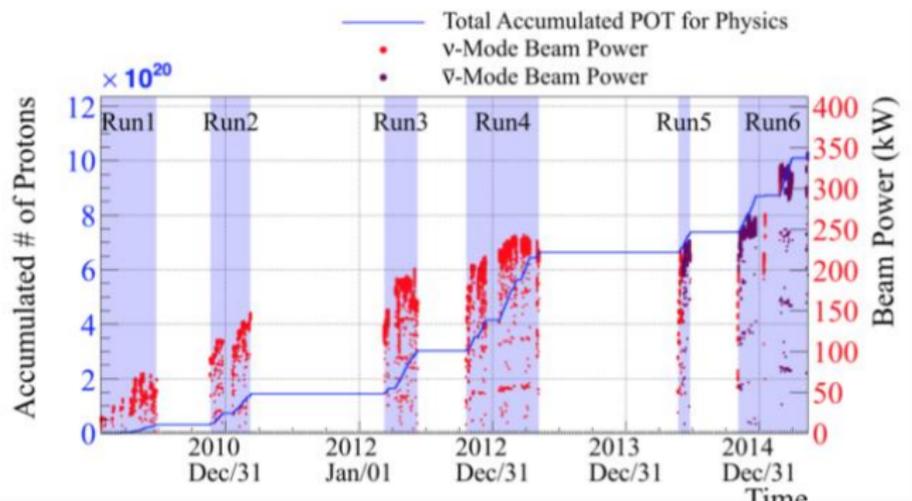
# Proton Background



# Flux Prediction



## $\bar{\nu}_e$ available POT data



- Stable operation at 345kW
- $10.3 \times 10^{20}$  (total) =  $7.0 \times 10^{20} (\nu) + 3.3 \times 10^{20} (\bar{\nu})$  accumulated (Jan 23, 2010 to May 13, 2015)

## Expected $\bar{\nu}_e$ CC inclusive events

- expected to have:  $4.5 \times 10^{20}$  POT by the end of May 2015 run
- expected to have: 6 (Run 5 CC inclusive events Ref. [B.Smith2014])  $\times$   
 $\frac{4.5 \times 10^{20} (\text{expected accumulated POT})}{2.5 \times 10^{19} (\text{Run 5 POT})} \sim 108 \text{ events}$

# Neutrino Interaction Model (slides from Morgan Waskow, Imperial College)

## $\nu$ -nucleus interaction models

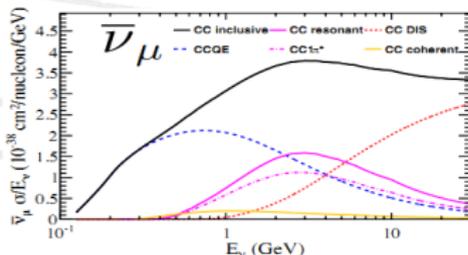
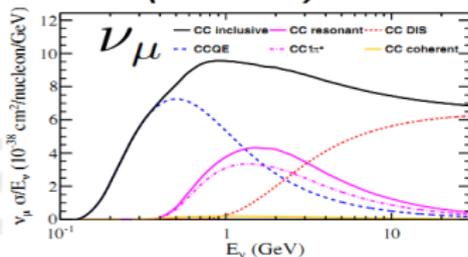
- T2K's primary neutrino generator MC is NEUT
  - Used by SK, SciBooNE, K2K
  - Tuned with fits to external data sets
  - 2012: mainly MiniBooNE CCQE, CC1 $\pi^+$ , CC1 $\pi^0$ , NC1 $\pi^0$ 
    - Fits used to tune model parameters for prior inputs to oscillation analysis
    - Constrained and cross-checked with SciBooNE and K2K data
  - 2014: MiniBooNE and MINERvA  $\nu$  and  $\bar{\nu}$  data sets
    - Fits used to down select default interaction model and tune parameters for prior inputs to oscillation analysis
  - Working on a publication to describe the model and the fit procedure
- Also use GENIE and NuWro for cross-check analyses, systematic errors studies, and deeper inquiries into neutrino interactions

# Neutrino Interaction Model (slides from Morgan Waskow, Imperial College)

## NEUT interaction model

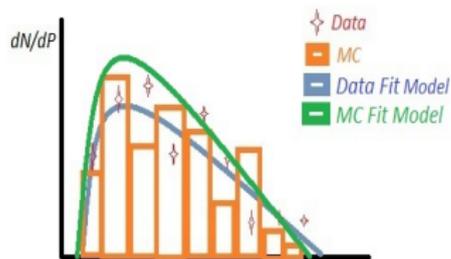
2012 model & parameters (v5.1.4.2)

- CCQE: Llewellyn Smith,  $M_{A}^{QE}=1.2 \text{ GeV}/c^2$
- CC resonant  $\pi$ : Rein-Sehgal,  $M_{A}^{RES}=1.2 \text{ GeV}/c^2$
- 2p2h: not simulated
- Nuclear model: Smith-Moniz RFG
- RPA effects not included
- Coherent pion: Rein-Sehgal with lepton mass effects
- DIS with Bodek-Yang corrections
- Neutrino and antineutrino interactions simulated
- $\nu_\mu$  and  $\bar{\nu}_\mu$  simulated
  - Only differ at low energy

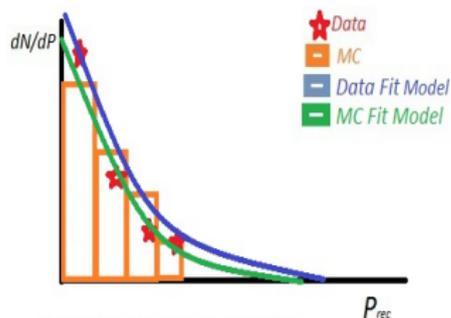




# $\bar{\nu}_e$ CC-Inclusive XSec Measurement strategy



Cartoon for Proton background Xsec contribution



Cartoon for Gamma background Xsec contribution

- Get separate sample of  $\gamma$  and protons.
- Compare momentum event spectrum of background sample to simulation.
- Fit a model  $N_\gamma(P_{rec}, \vec{a}), N_{proton}(P_{rec}, \vec{b})$ .
- From the fit get model uncertainty from difference between data and MC.
- Use a Bayesian unfolding with background subtraction to get cross section as function of  $E_\nu, Q^2$ .

# Bayesian Unfolding Details

- $N_{t_k} = \sum_j S_{r_j t_k} + M_{t_k}$
- $P(r_j | t_k) = \frac{S_{r_j t_k}}{N_{t_k}}$
- $\epsilon_{t_k} = \frac{\sum_j S_{r_j t_k}}{N_{t_k}}$
- $P_m(t_k | r_j) = \frac{P(r_j | t_k) P_{t_k}}{\sum_\alpha P(r_j | t_\alpha) P_{t_\alpha}}$
- $N_{t_k}^{m+1} = \frac{1}{\epsilon} \sum_j P_m(t_k | r_j) (N_{r_j}^{meas} - B_{r_j})$

with :

- $N_{t_k}$  : Estimator for true signal event in bin k
- $S_{r_j t_k}$  : Signal matrix, i.e the number of true simulated signal events in true bin  $t_k$  that were reconstructed in bin  $r_j$
- $M_{t_k}$  : Missed vector, i.e number of simulated signal events in true bin  $t_k$  that were not selected
- $\epsilon_{t_k}$  : Reconstruction efficiency
- $P_m(t_k | r_j)$  : Probability at iteration m to have a true event at bin k given a reconstructed event at bin j
- $B_{r_j}$  : Background

