

# Microscopic simulations with modern nuclear forces

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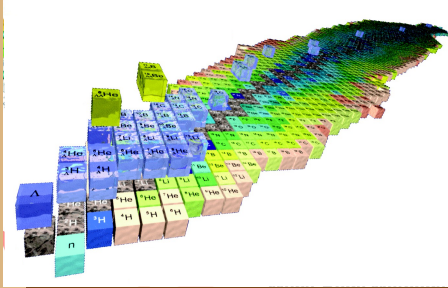
# Getting the TLAs out of the way

**QCD = Quantum Chromodynamics**

**EFT = Effective Field Theory**

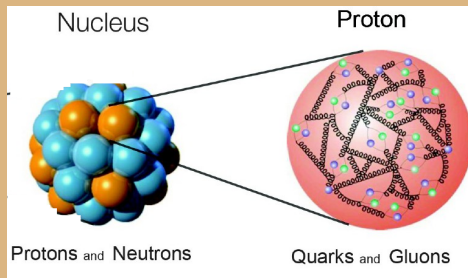
**QMC = Quantum Monte Carlo**

# Outline



## Many-nucleon problem

- Chart of nuclides
- Neutron stars
- Perturbative vs non-perturbative



## Nuclear forces

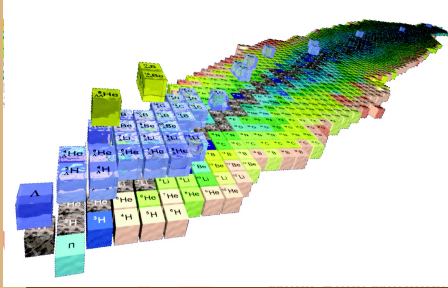
- Chiral Effective Field Theory (EFT)
- Local chiral EFT



## Few- and many-body results

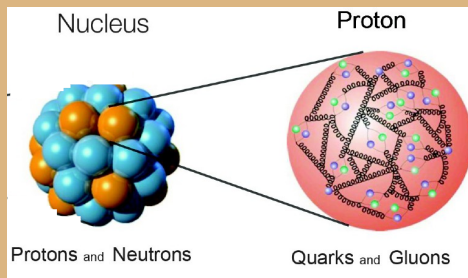
- Pure neutron matter
- Light nuclei

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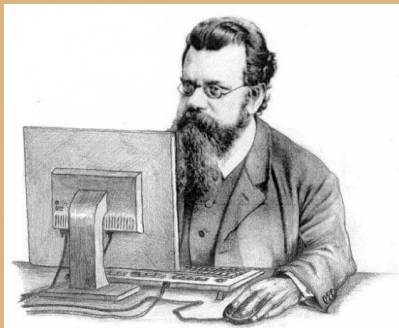
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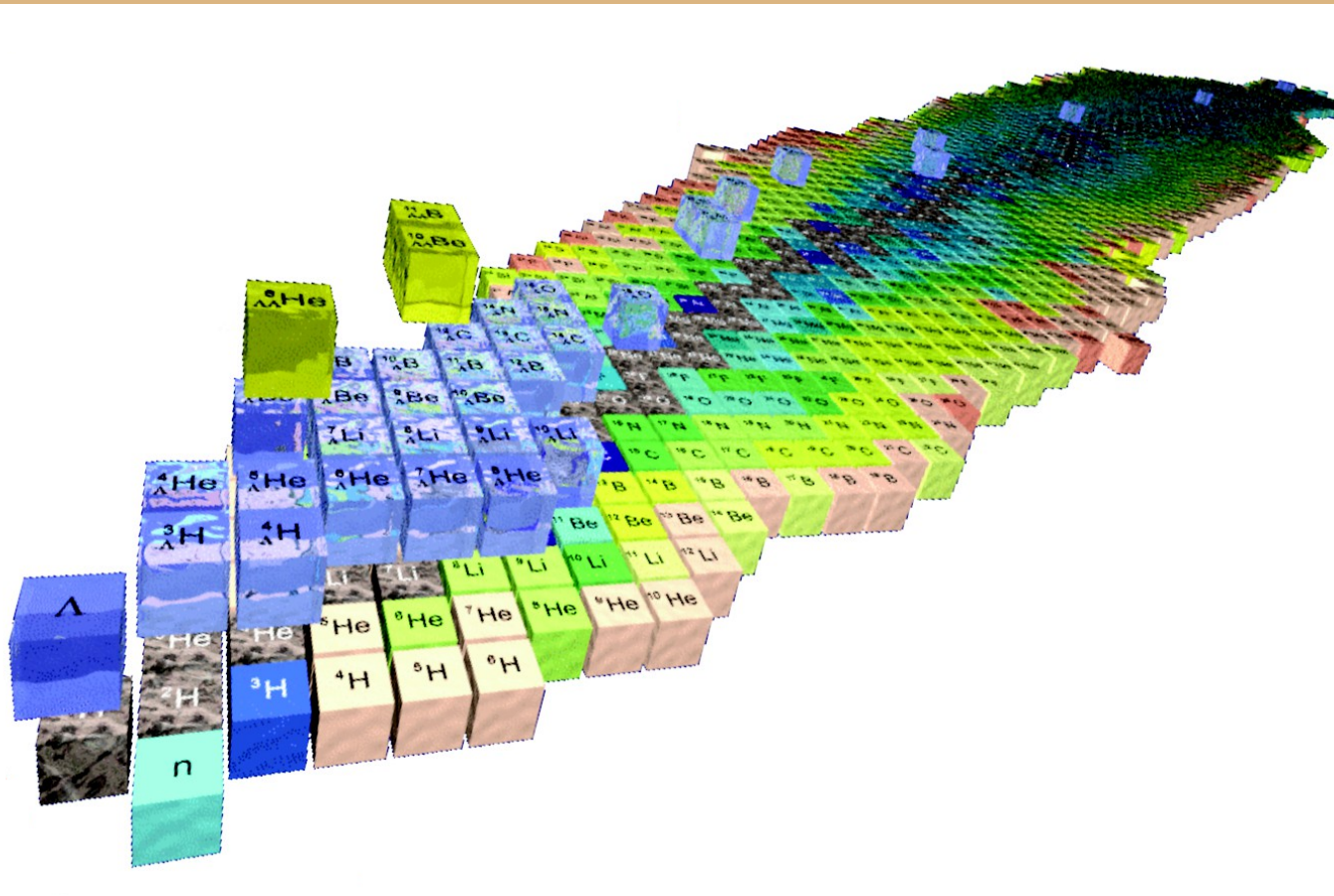


## Few- and many-body results

- Pure neutron matter
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# Goals: terrestrial

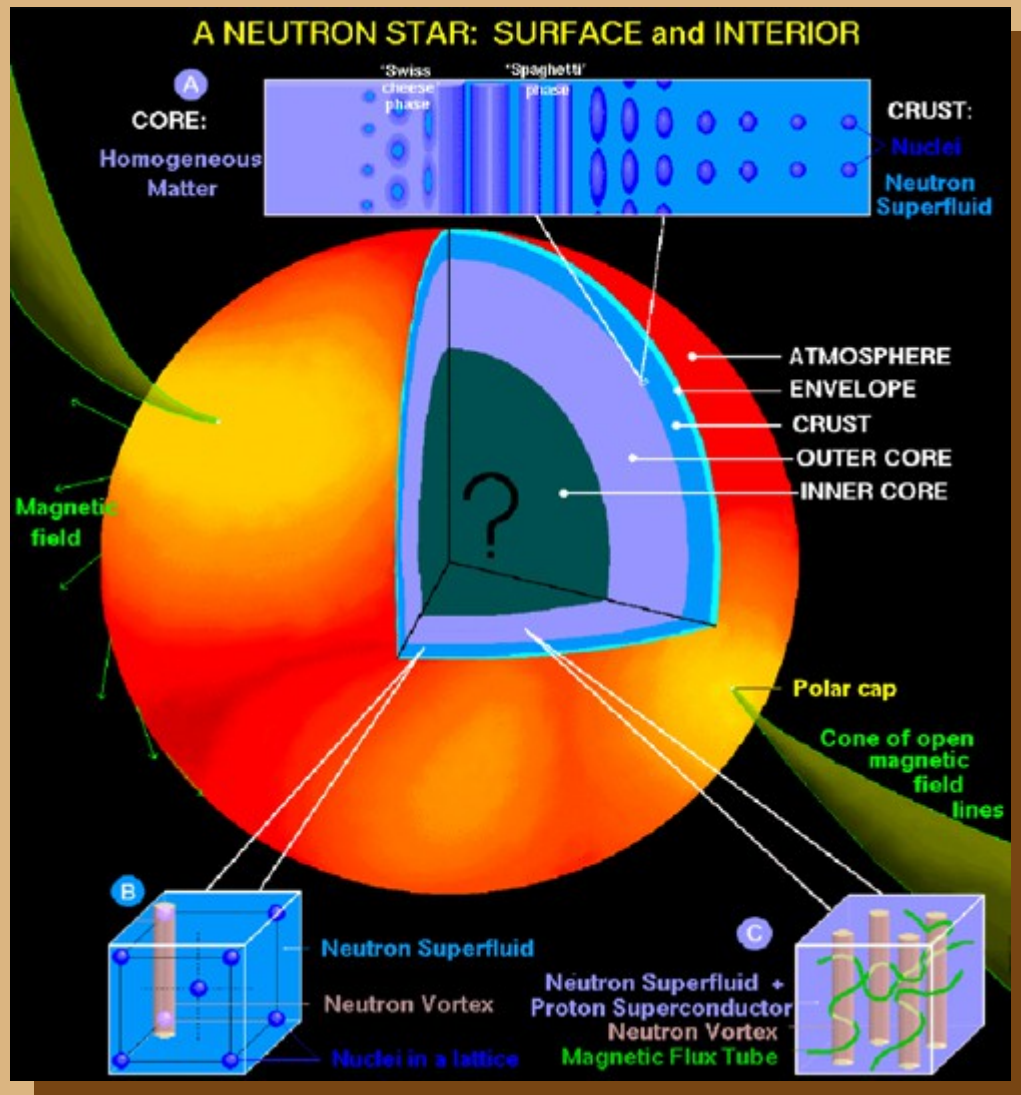
## Chart of nuclides



- Nuclear physics is developed and tested on earth
- Using complicated many-body methods we can try to “build nuclei from scratch”
- We then extrapolate that knowledge to more exotic systems
- Can we describe nuclei *from first principles*?
- What does *first* mean?

# Goals: astrophysical

## Neutron stars as ultra-dense matter laboratories

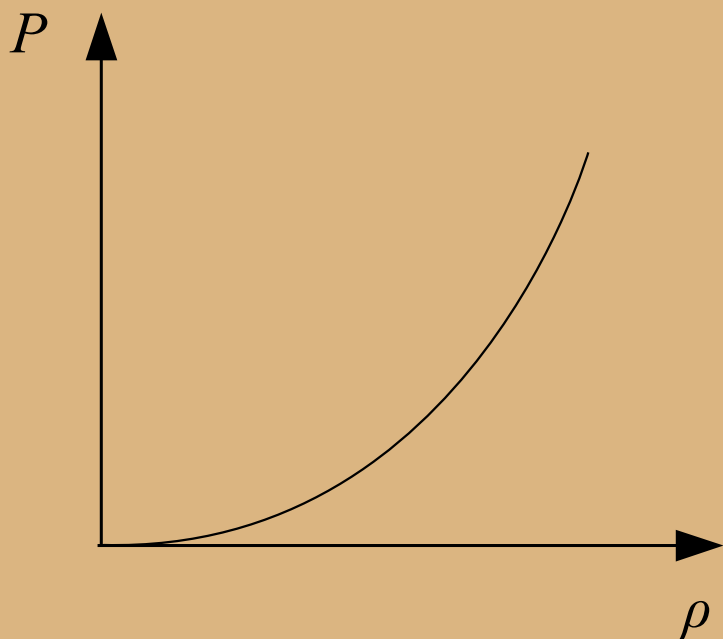


- Ultra-dense: 1.4 solar masses (or more) within a radius of 10 kilometres
- Terrestrial-like (outer layers) down to exotic (core) behaviour
- Observationally probed, i.e., not experimentally accessible
- Can we describe neutron-star matter *from first principles*?
- What does *first* mean?

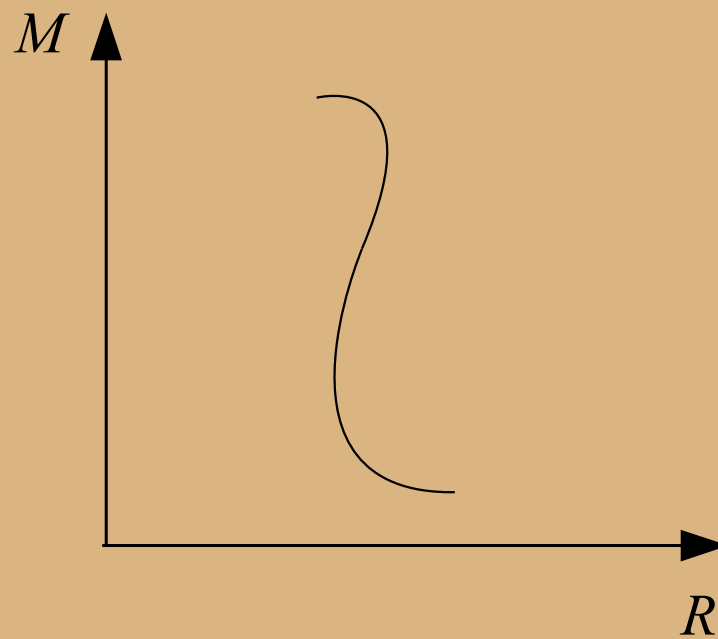
# Neutron stars: micro-macro

## TOV equations (or Hartle-Thorne, etc)

Pressure vs density



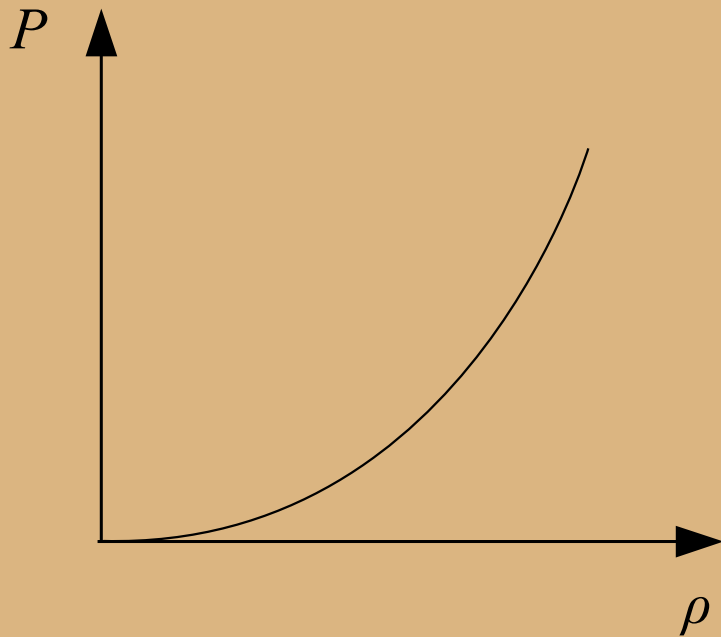
Mass vs radius



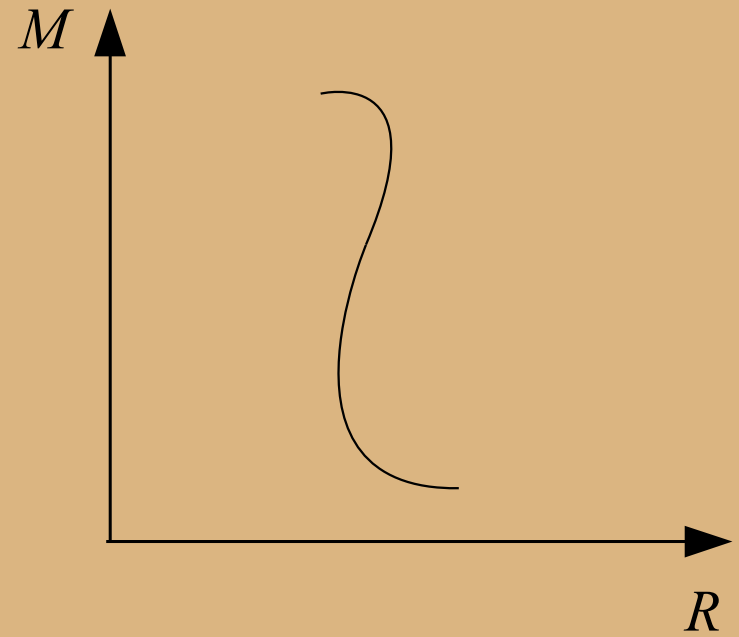
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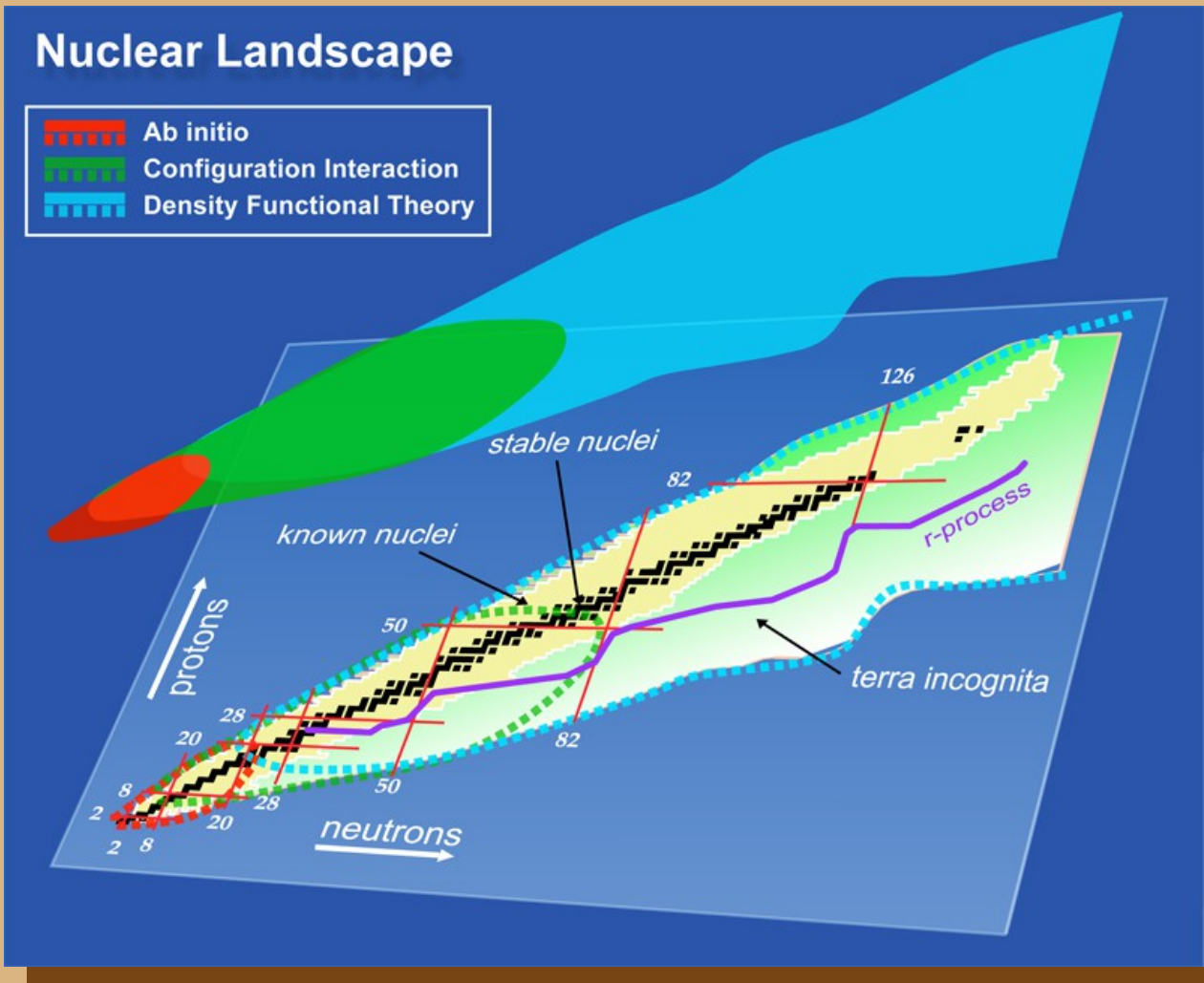
Mass vs radius



**Modern goal: systematic theoretical error bars**



# Many-nucleon problem: methods



- No universal method exists (yet?)
- A lot to be learned if the degrees of freedom are actual particles and there are no free parameters
- Regions of overlap between different methods are crucial
- Is it possible to work at the level of nucleons & pions but still connect to the underlying level?

# Many-nucleon problem

**Lesson from history of physics:**

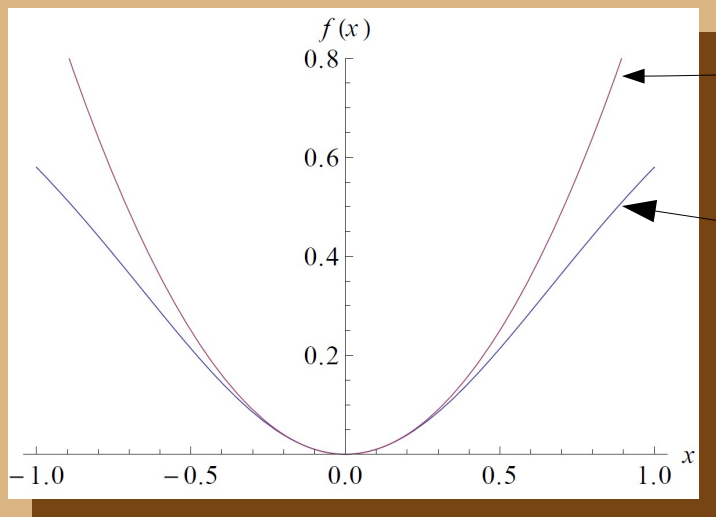
**Find a small expansion parameter  $\alpha$  and use it to organize the different contributions:**

$$\alpha + \alpha^2 + \alpha^3 + \dots$$

**This is called perturbation theory.**

**However, some problems are non-perturbative.**

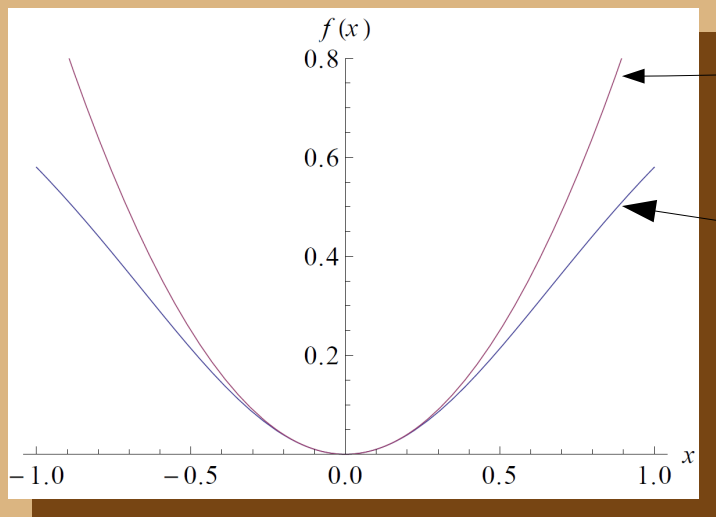
# Perturbative vs non-perturbative



$$1 - \left( \frac{1}{\cosh(x)} \right)^2 = x^2 + \text{something small} \quad ?$$

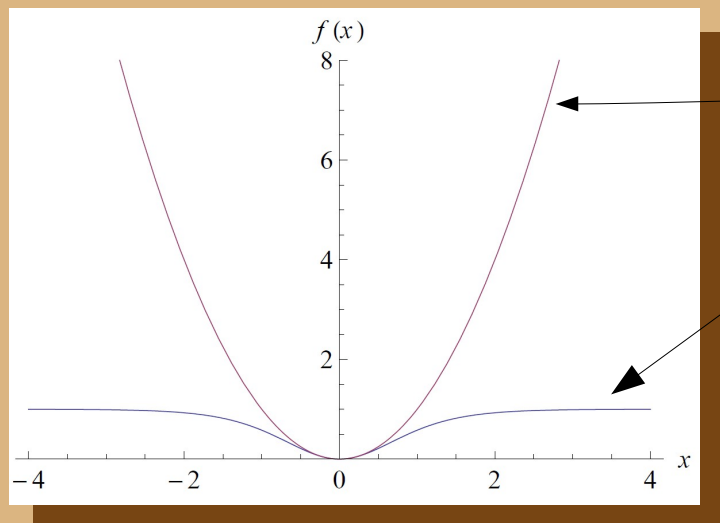
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# Many-body problem: QMC

**Quantum Monte Carlo: stochastically solve the many-body Schrödinger equation in a fully non-perturbative manner**

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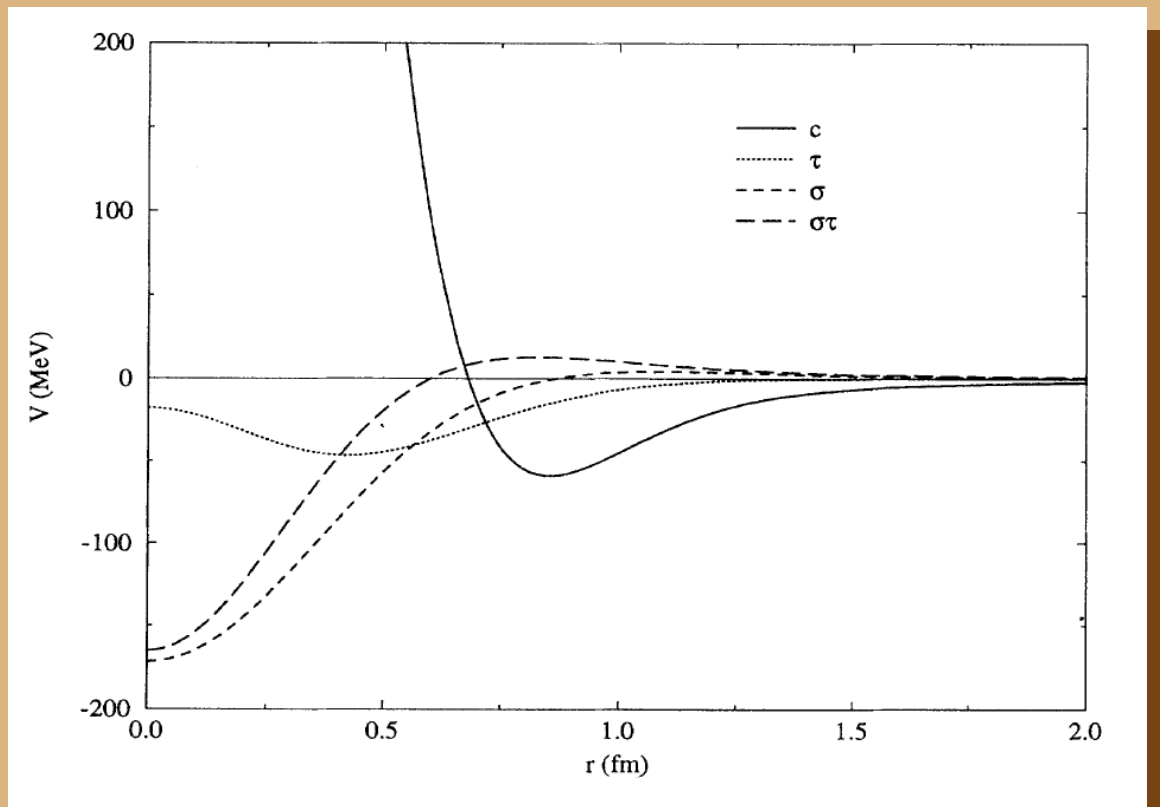
**Quantum Monte Carlo:** stochastically solve the many-body Schrödinger equation in a fully non-perturbative manner

**Rudiments of**

**Diffusion Monte Carlo:** 
$$\Psi(\tau \rightarrow \infty) = \lim_{\tau \rightarrow \infty} e^{-(\mathcal{H} - E_T)\tau} \Psi_V$$
$$\rightarrow \alpha_0 e^{-(E_0 - E_T)\tau} \Psi_0$$

# Many-body problem: QMC

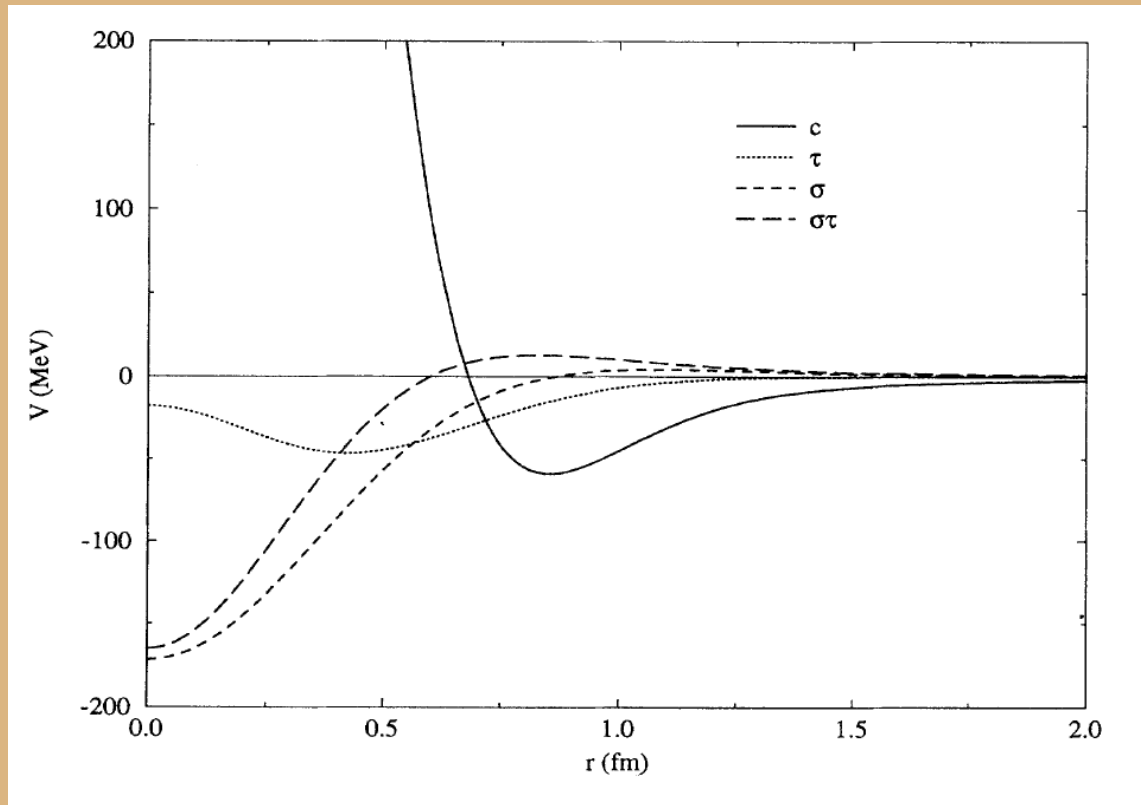
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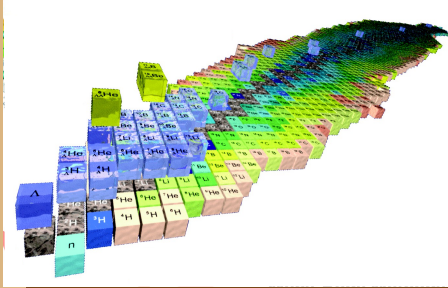


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Such potentials are hard, making them non-perturbative at the many-body level (which isn't a problem for QMC, but is one for almost every other method).

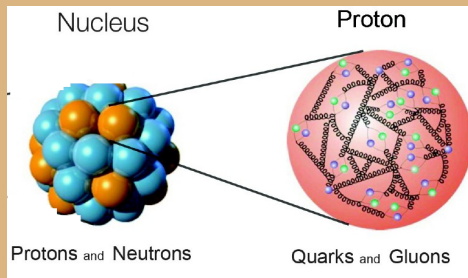


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## Nuclear forces

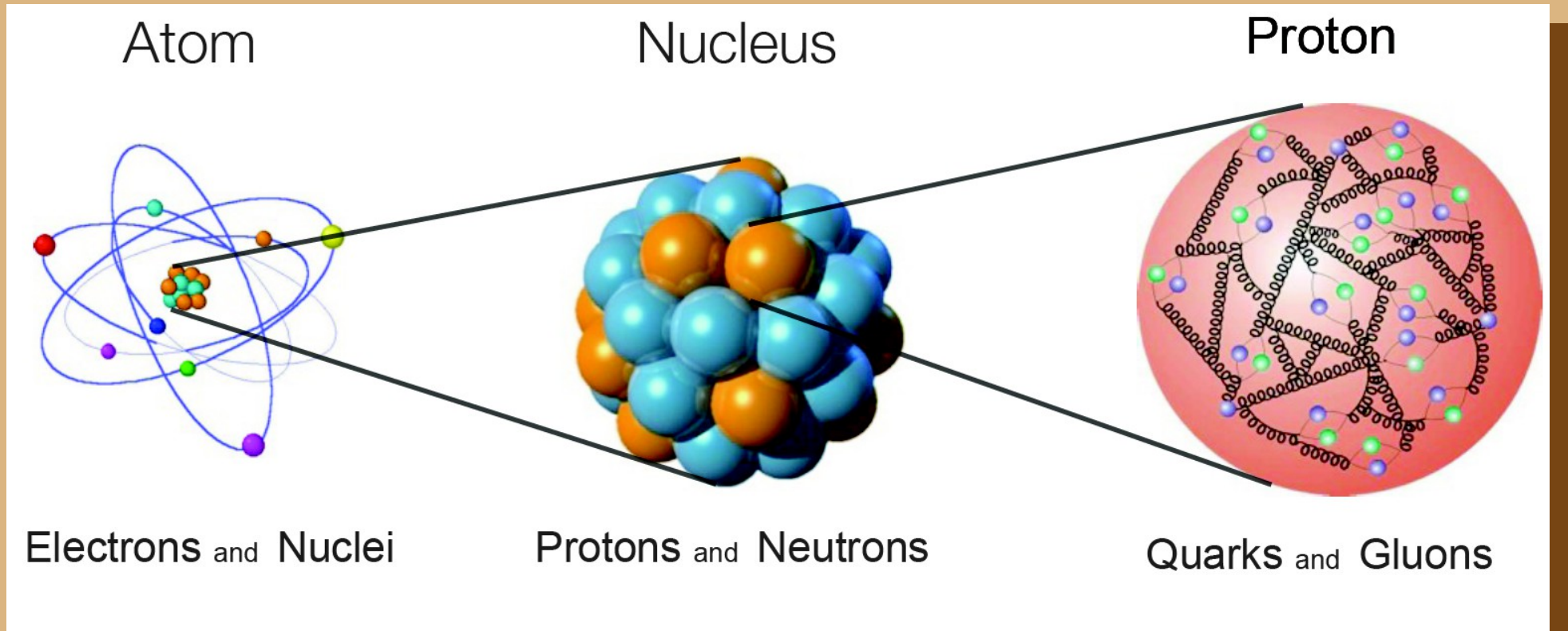
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# Nuclear forces: degrees of freedom



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Quantum Electrodynamics  
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$$\alpha = \frac{1}{137.036}$$

for most practical purposes

(see Gell-Mann & Low paper with 770 citations,  
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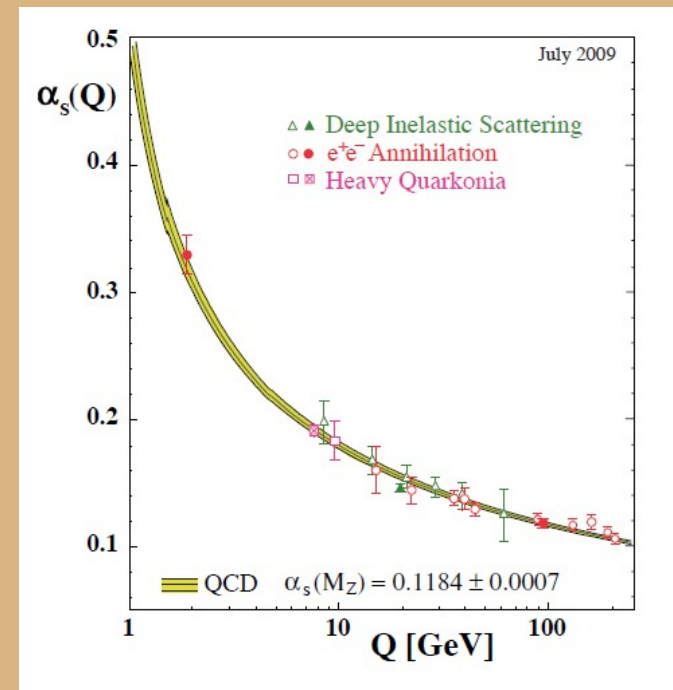
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Quantum Chromodynamics  
(QCD)



# Perturbative and non-perturbative

Faced with an inherent non-perturbativeness of the many-nucleon problem we turned to a non-perturbative method (Quantum Monte Carlo).

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Perhaps a non-perturbative method should also be used for nuclear forces (lattice QCD).

For now, many-nucleon studies are limited to nuclear forces that follow from nucleons and pions (but it would be nice not to ignore the existence of the underlying level).



# Nuclear Hamiltonian: chiral EFT

How to build on QCD in a systematic manner?

Exploit separation of scales:  $a_{1S_0} = (11 \text{ MeV})^{-1}$

$$m_\pi = 140 \text{ MeV}$$

$$\Lambda_\chi \approx m_\rho \approx 800 \text{ MeV}$$

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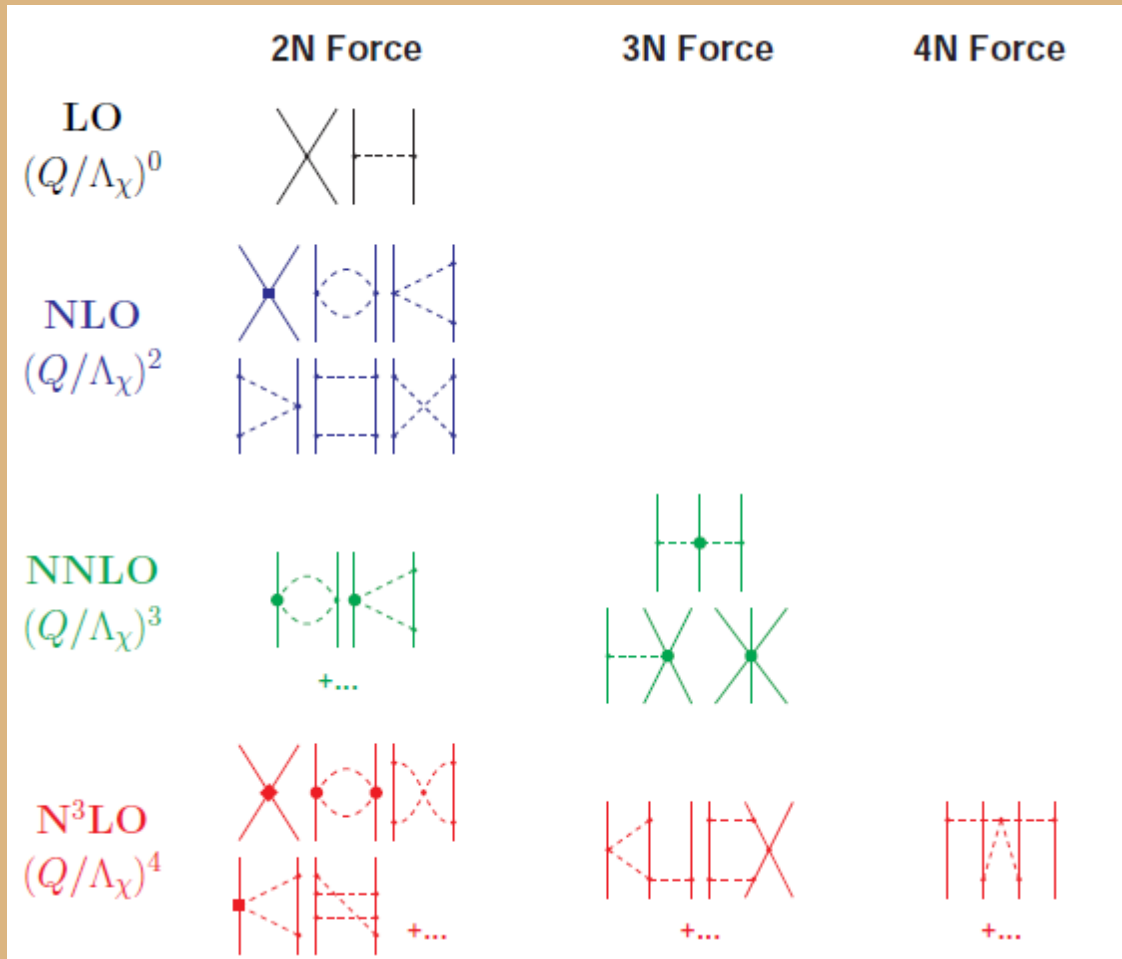
**Chiral Effective Field Theory approach:**

Use nucleons and pions as degrees of freedom

Systematically expand in  $\frac{Q}{\Lambda_\chi}$

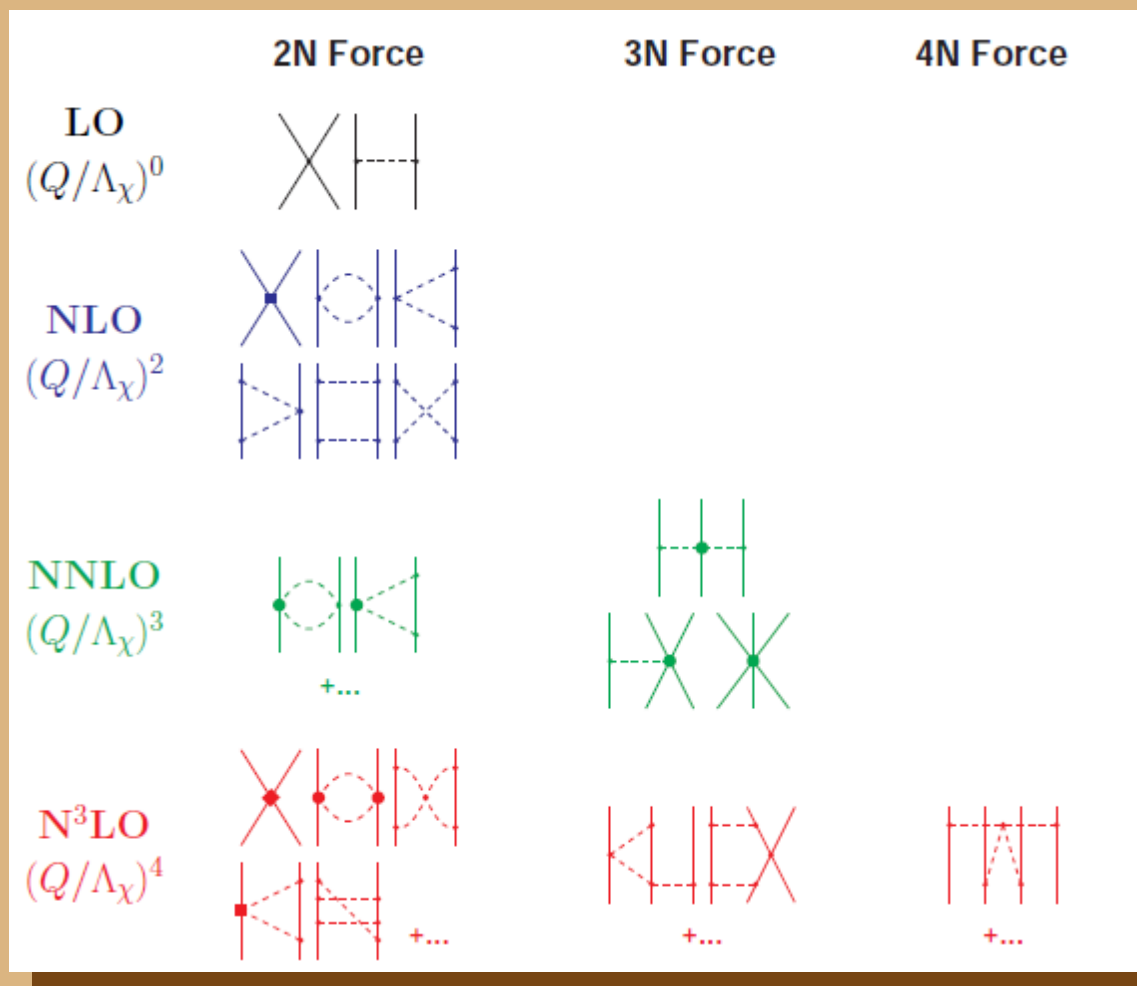
Program introduced by S. Weinberg, now taken over by the nuclear community

# Nuclear Hamiltonian: chiral EFT



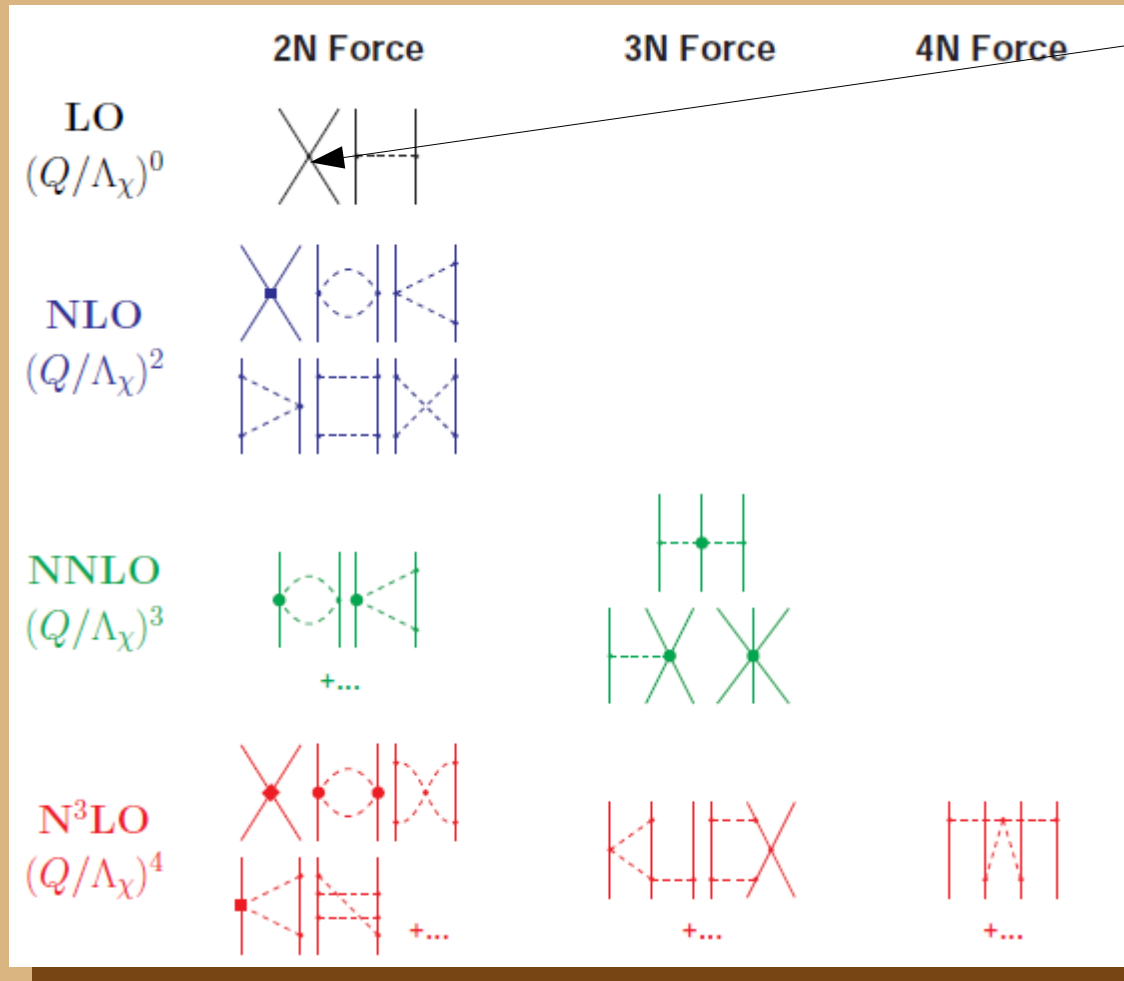
- Attempts to connect with underlying theory (QCD)
- Systematic low-momentum expansion
- Consistent many-body forces
- Low-energy constants from experiment or lattice QCD
- Until now non-local in coordinate space, so unused in continuum QMC
- Power counting's relation to renormalization still an open question

# Nuclear Hamiltonian: chiral EFT

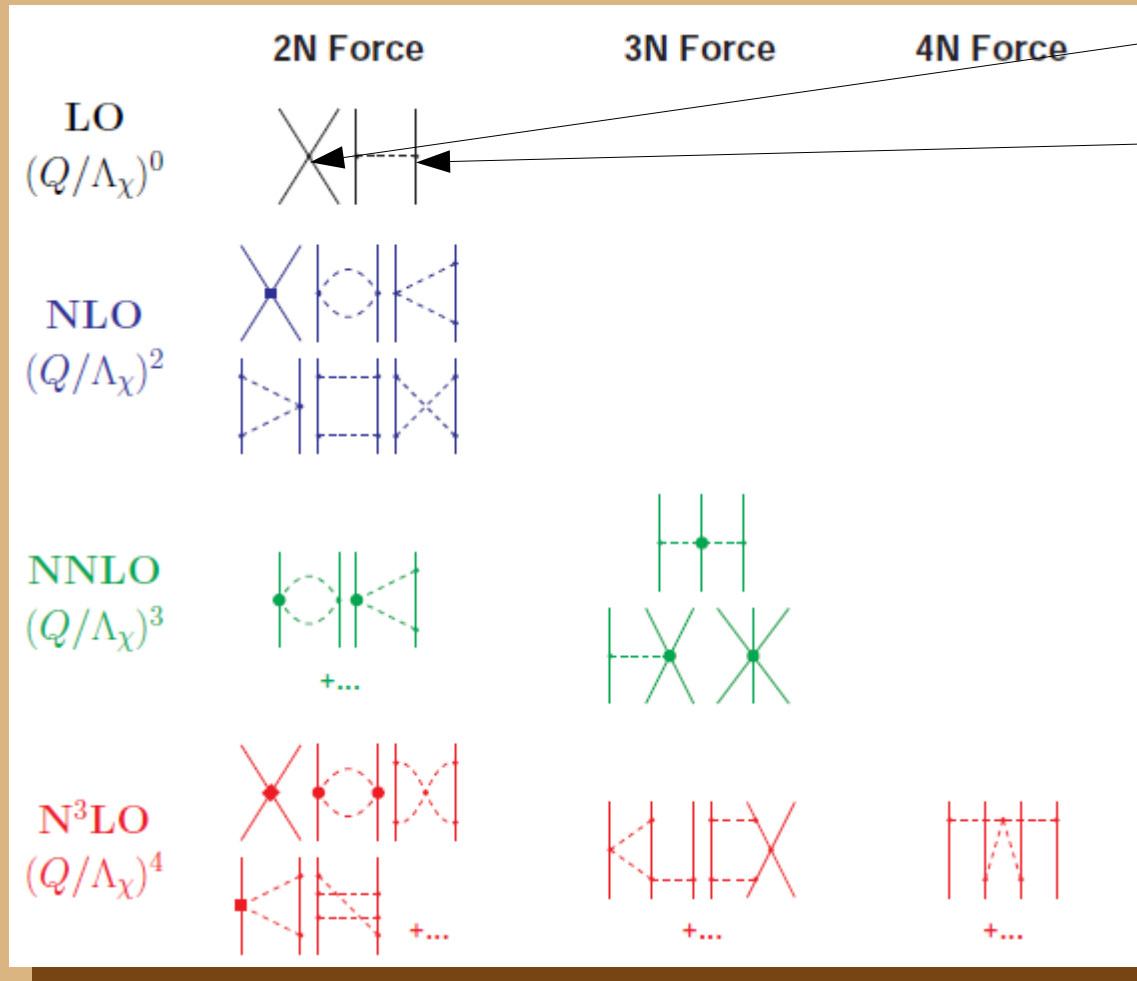


# Nuclear Hamiltonian: chiral EFT

$$V_{\text{ct}}^{(0)} = C_S + C_T \sigma_1 \cdot \sigma_2$$



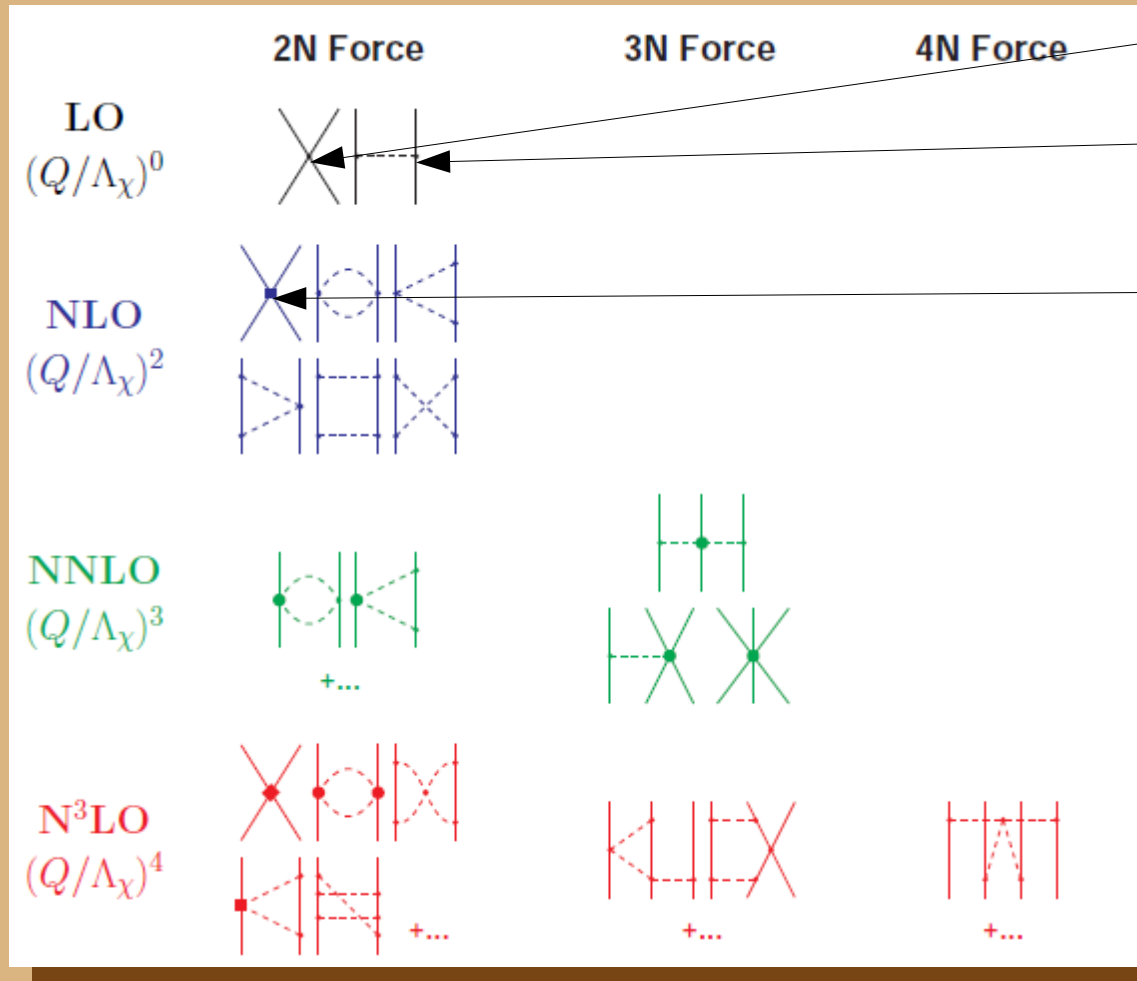
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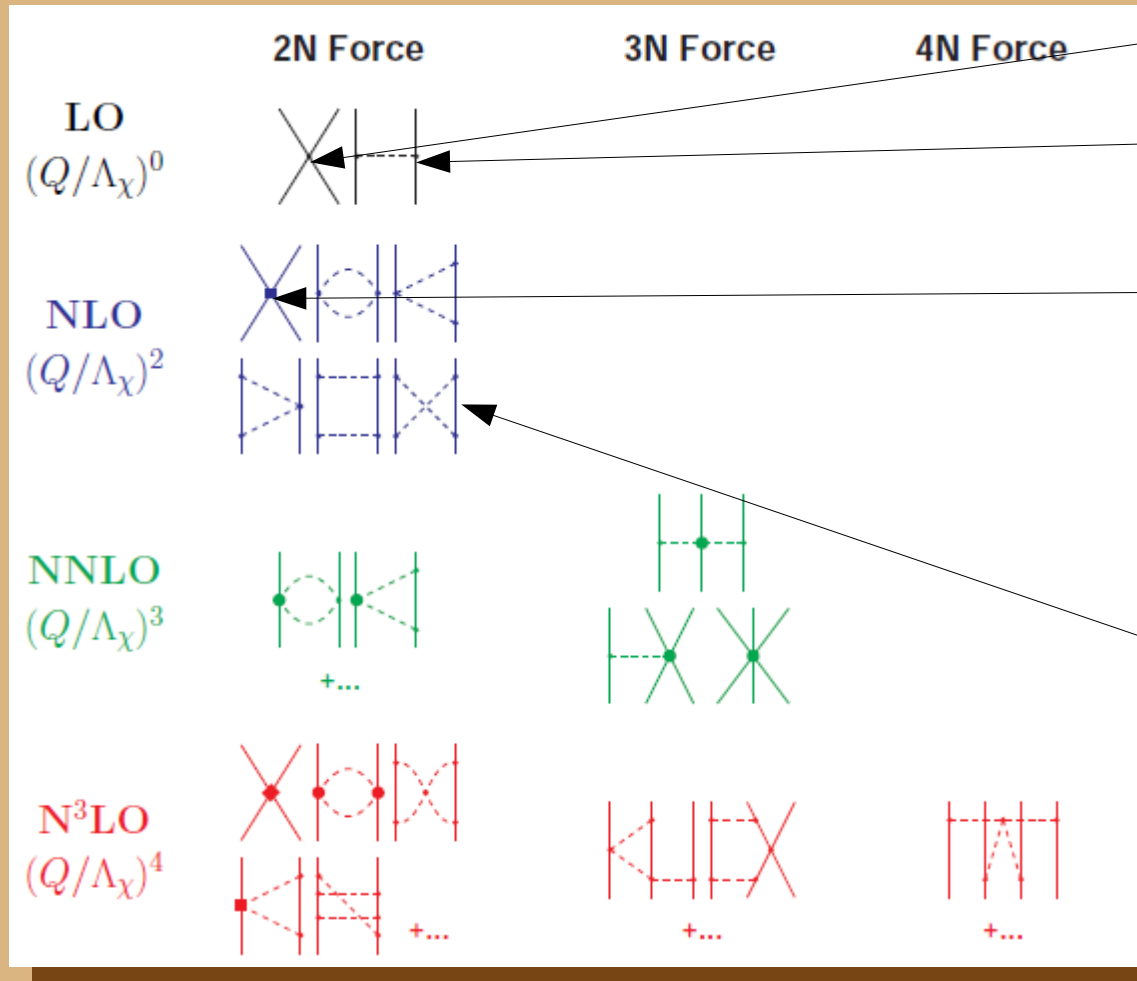
$$+ (C_3 q^2 + C_4 k^2) \sigma_1 \cdot \sigma_2$$

$$+ i \frac{C_5}{2} (\sigma_1 + \sigma_2) \cdot (\mathbf{q} \times \mathbf{k})$$

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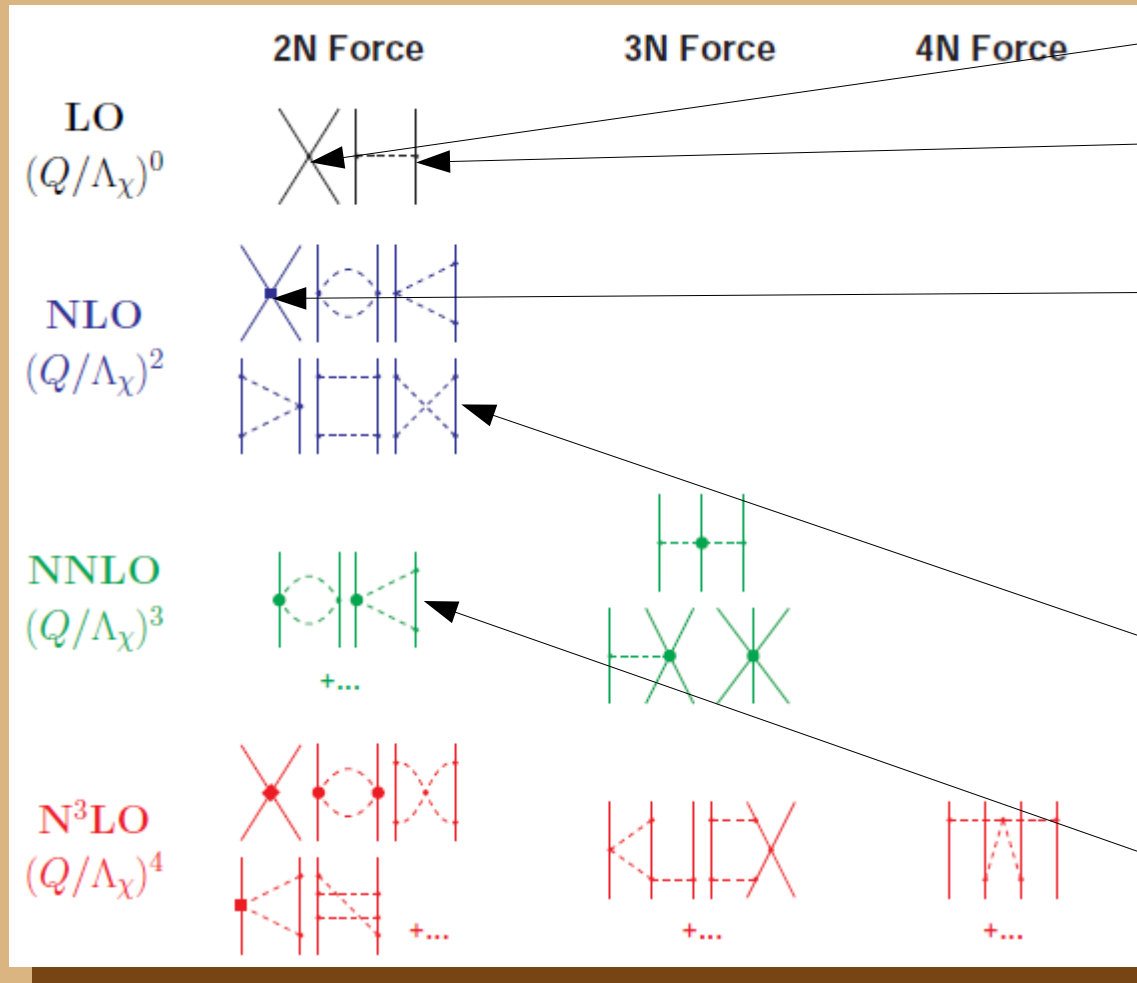
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Long-studied  
two-pion exchange



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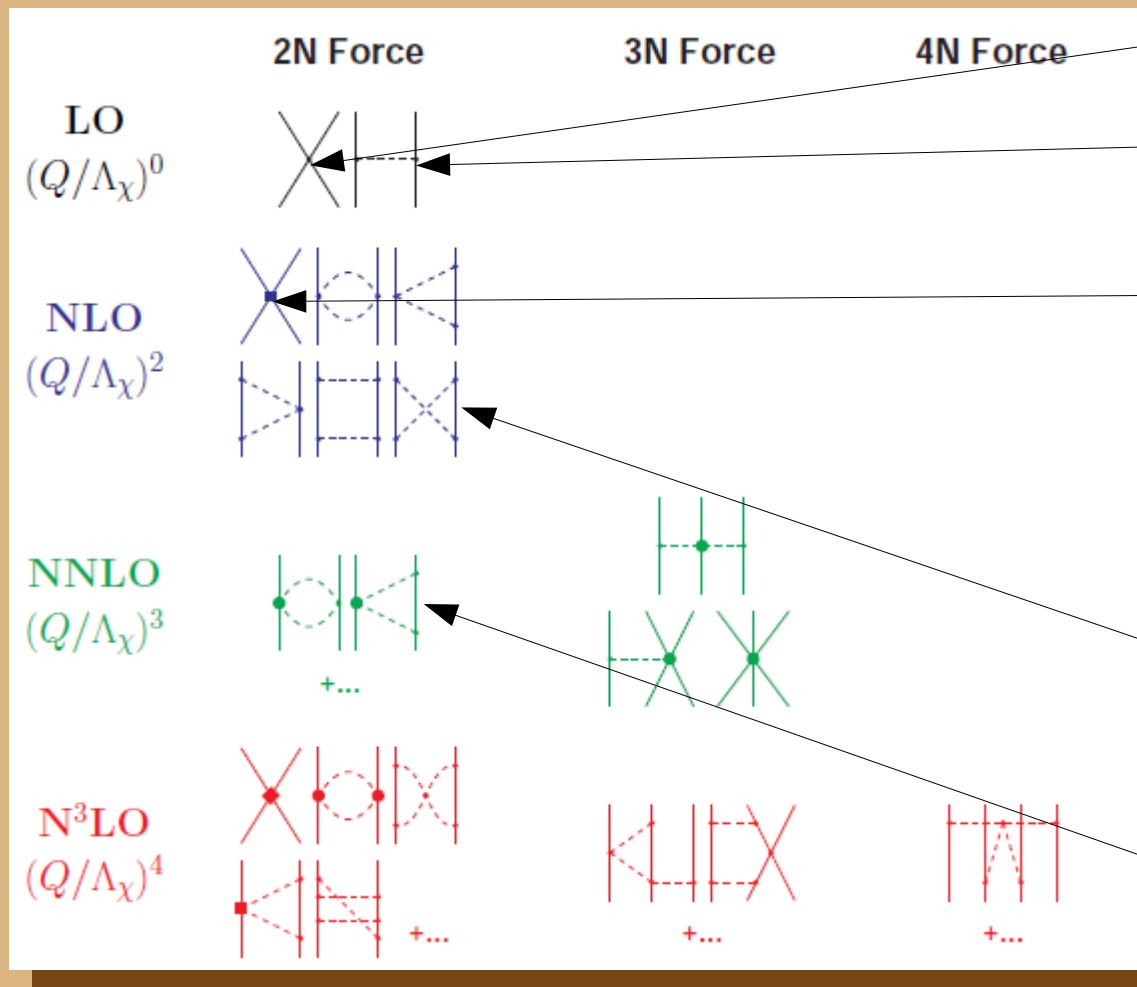
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Contains couplings  
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$$V_{\text{ct}}^{(2)} = C_1 q^2 + C_2 k^2 + (C_3 q^2 + C_4 k^2) \sigma_1 \cdot \sigma_2 + i \frac{C_5}{2} (\sigma_1 + \sigma_2) \cdot (\mathbf{q} \times \mathbf{k}) + C_6 (\sigma_1 \cdot \mathbf{q})(\sigma_2 \cdot \mathbf{q}) + C_7 (\sigma_1 \cdot \mathbf{k})(\sigma_2 \cdot \mathbf{k})$$

Long-studied two-pion exchange

Contains couplings from  $\pi N$  scattering

$\mathbf{q}$  means local

$\mathbf{k}$  means non-local

# Nuclear forces: summary

**Local high-quality phenomenology is hard**

Consubstantial with the successes of nuclear QMC,  
difficult to use in most other many-body methods

**Chiral EFT**

**a) is connected to symmetries of QCD**

**b) has consistent many-body forces, and**

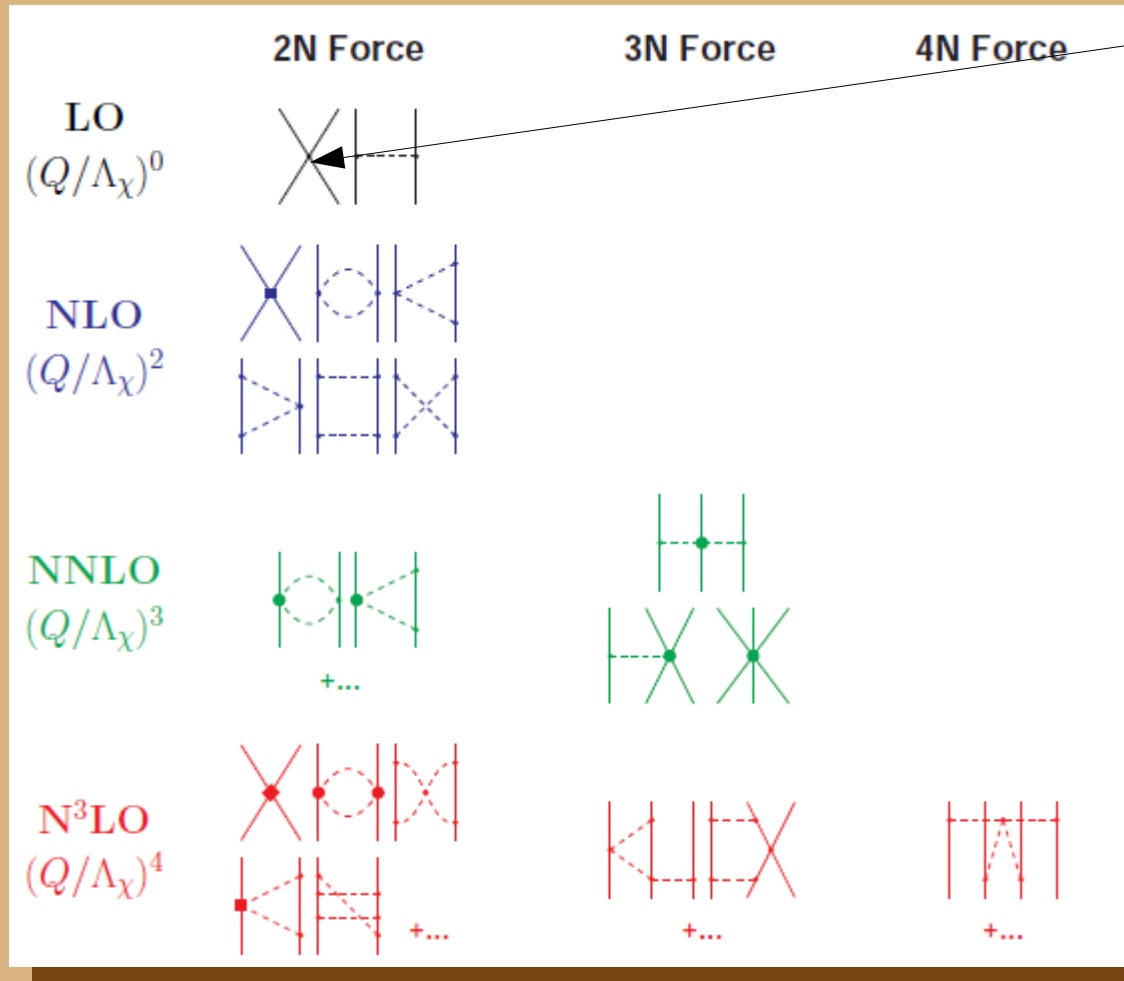
**c) allows us to produce systematic uncertainty bands**

**also happens to be non-local (such are the *sumbebekota*)**

Heavily used in other methods, but not used in nuclear QMC

**Turning to the resolution**

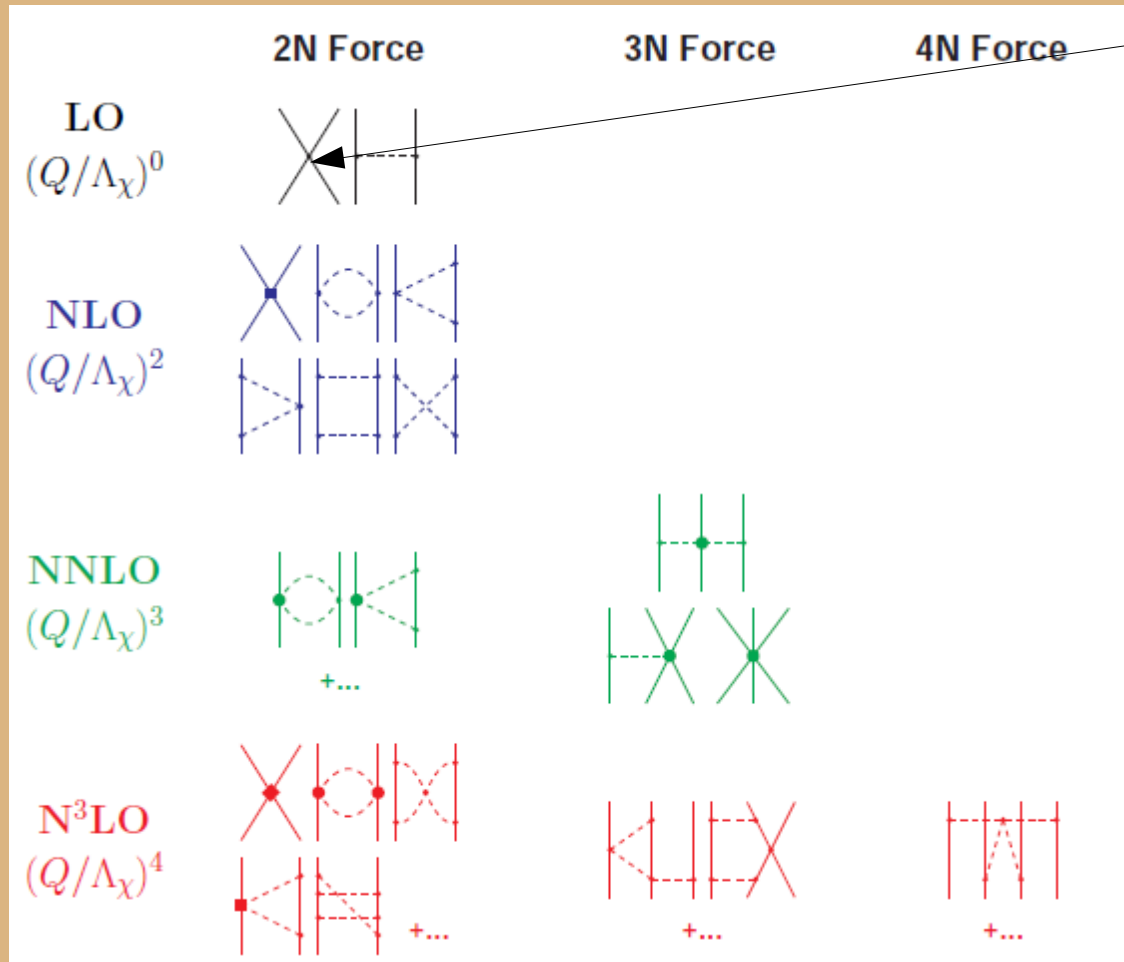
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Merely the standard choice.

Actually 4 terms in full set  
consistent with the symmetries of QCD

$$V_{\text{ct}}^{(0)} = C_1 + C_2 \sigma_1 \cdot \sigma_2 + C_3 \tau_1 \cdot \tau_2 + C_4 \sigma_1 \cdot \sigma_2 \tau_1 \cdot \tau_2$$

Pick 2 and antisymmetrize

# How to go beyond?

## Combine power of Quantum Monte Carlo with consistency of chiral Effective Field Theory

Write down a local energy-independent NN potential

- Use local pion-exchange regulator  $f_{\text{long}}(r) = 1 - e^{-(r/R_0)^4}$   
cf.  $f(p, p') = e^{-(p/\Lambda)^{2n}} e^{-(p'/\Lambda)^{2n}}$

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- Use local pion-exchange regulator  $f_{\text{long}}(r) = 1 - e^{-(r/R_0)^4}$
- Pick 7 different contacts at NLO, just make sure that when antisymmetrized they lead to a set obeying the required symmetry principles

$$\begin{aligned} V_{\text{ct}}^{(2)} = & C_1 q^2 + C_2 q^2 \boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_2 \\ & + (C_3 q^2 + C_4 q^2 \boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_2) \boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2 \\ & + i \frac{C_5}{2} (\boldsymbol{\sigma}_1 + \boldsymbol{\sigma}_2) \cdot \mathbf{q} \times \mathbf{k} \\ & + C_6 (\boldsymbol{\sigma}_1 \cdot \mathbf{q})(\boldsymbol{\sigma}_2 \cdot \mathbf{q}) \\ & + C_7 (\boldsymbol{\sigma}_1 \cdot \mathbf{q})(\boldsymbol{\sigma}_2 \cdot \mathbf{q}) \boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_2 \end{aligned}$$

cf.

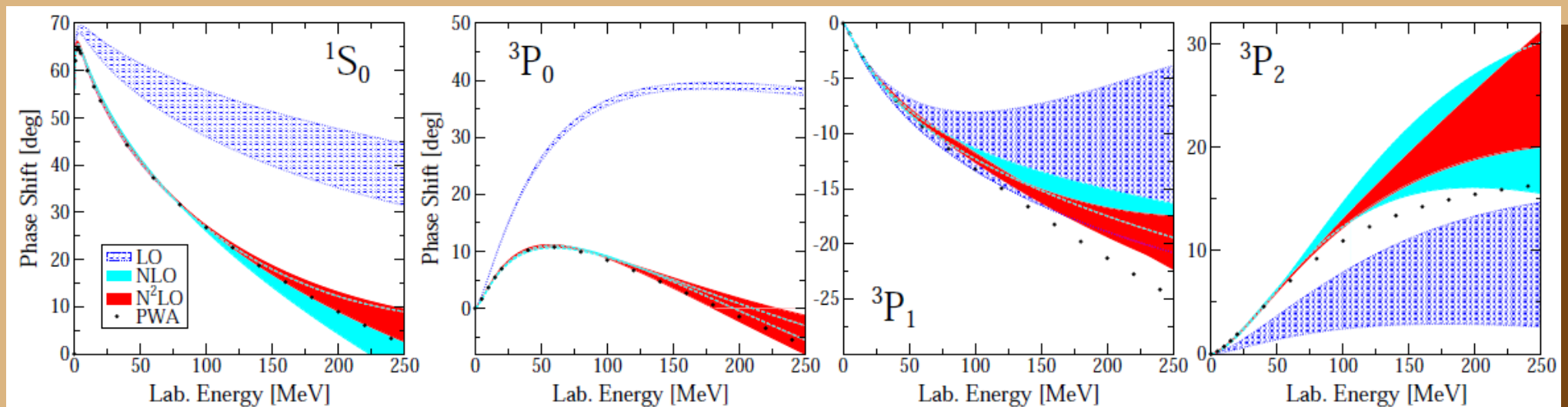
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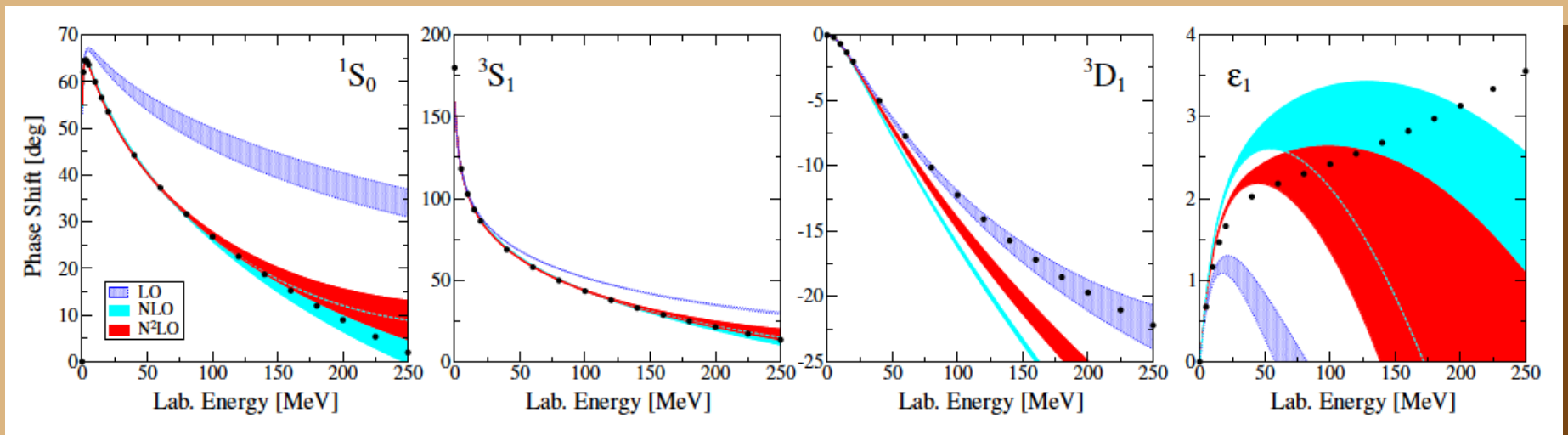
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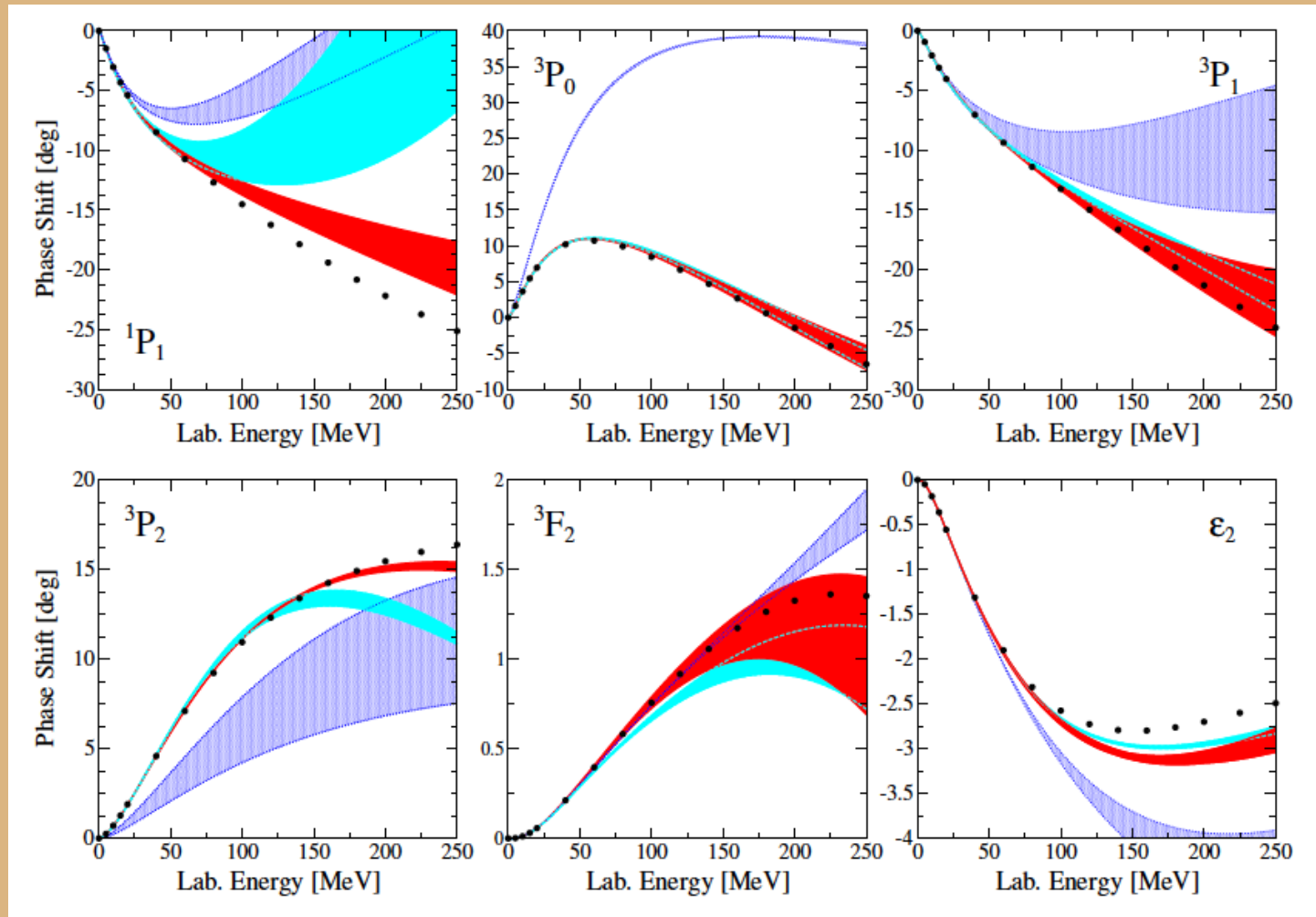
- Write down a local energy-independent NN potential
- Before doing many-body calculations, fit to NN phase shifts



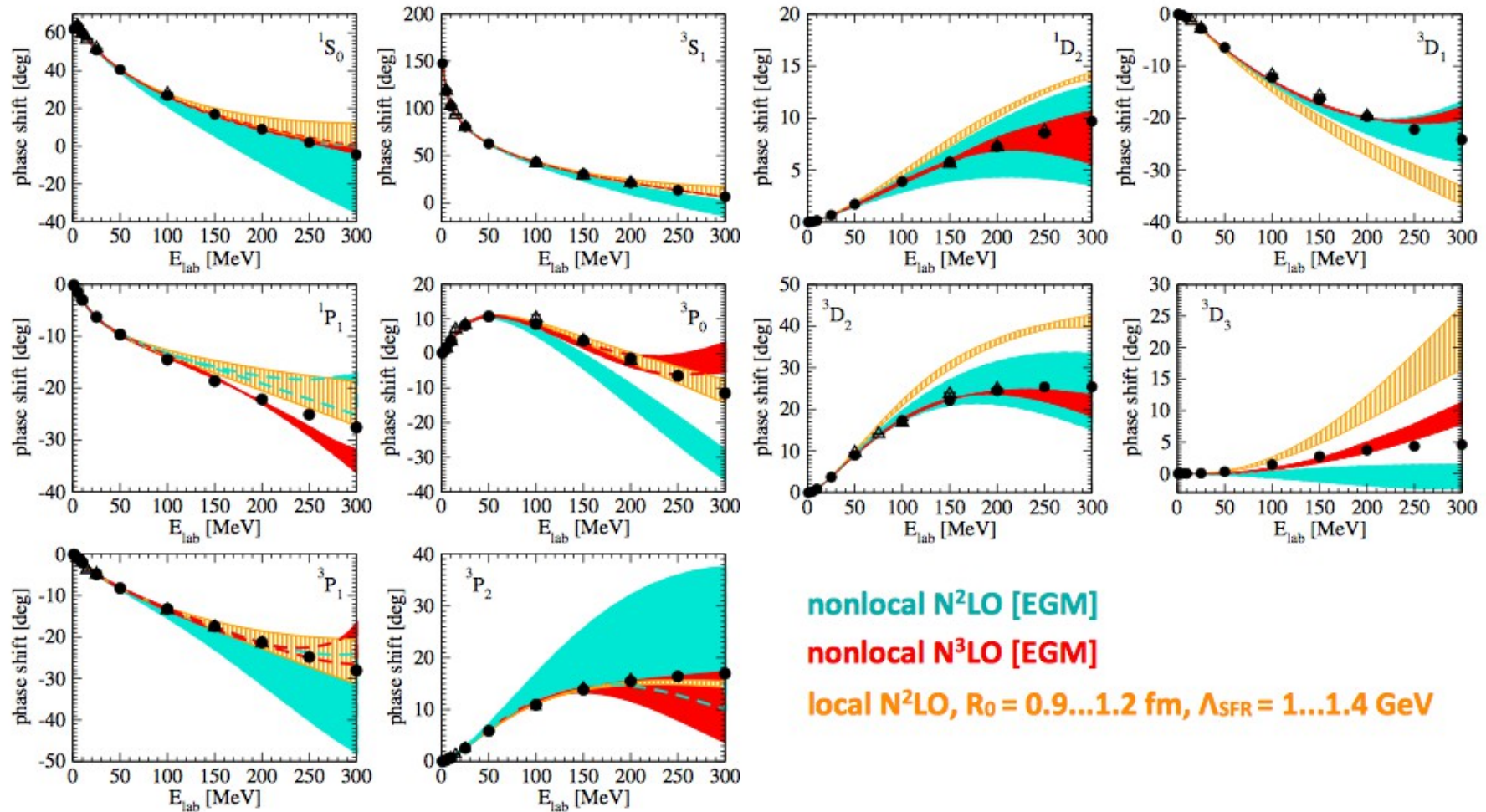
# Updated phase shifts



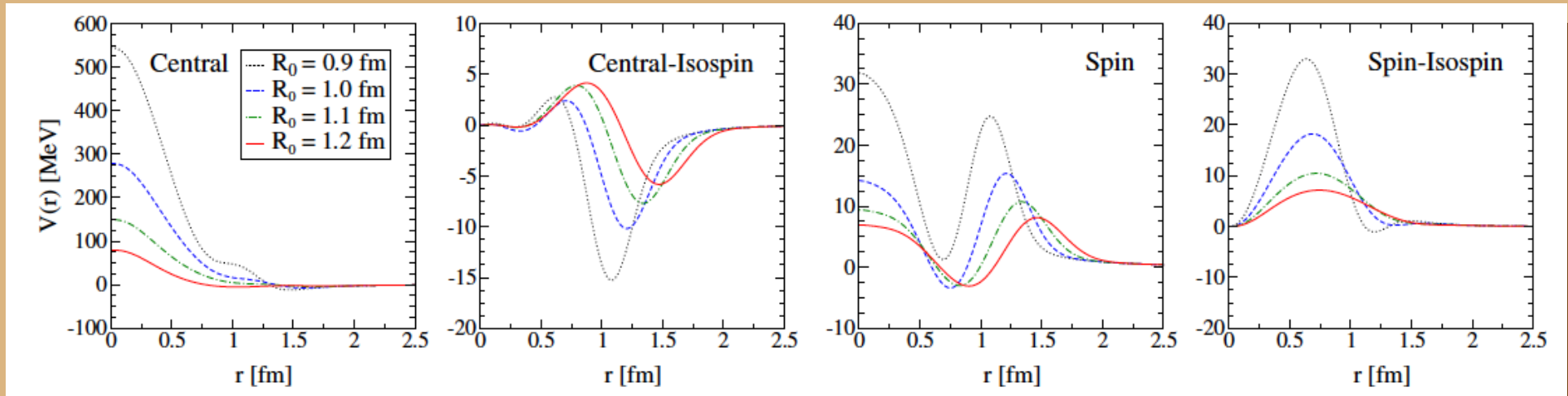
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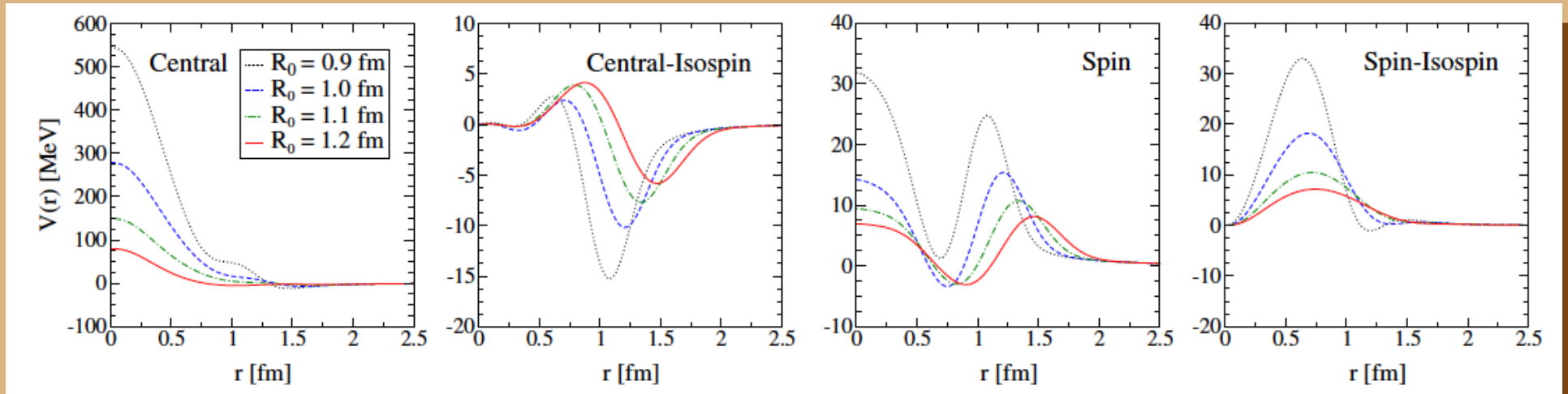
# Compare to non-local EGM



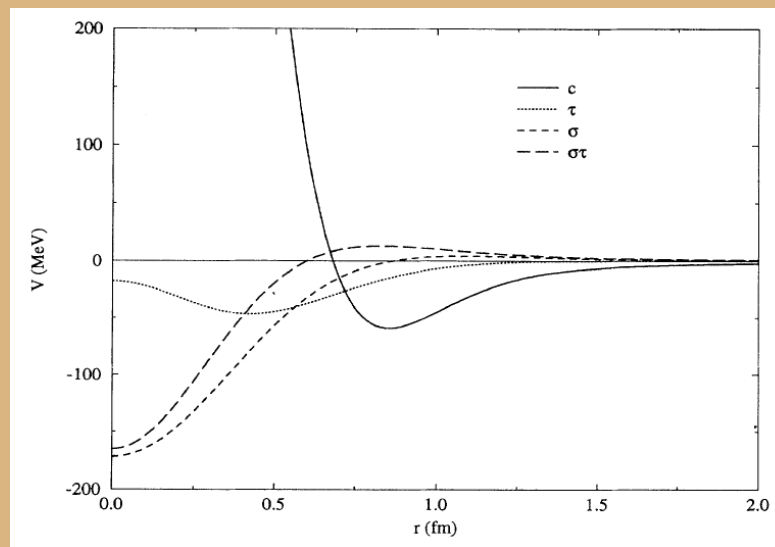
# Since it's local, let's plot it (N<sup>2</sup>LO)



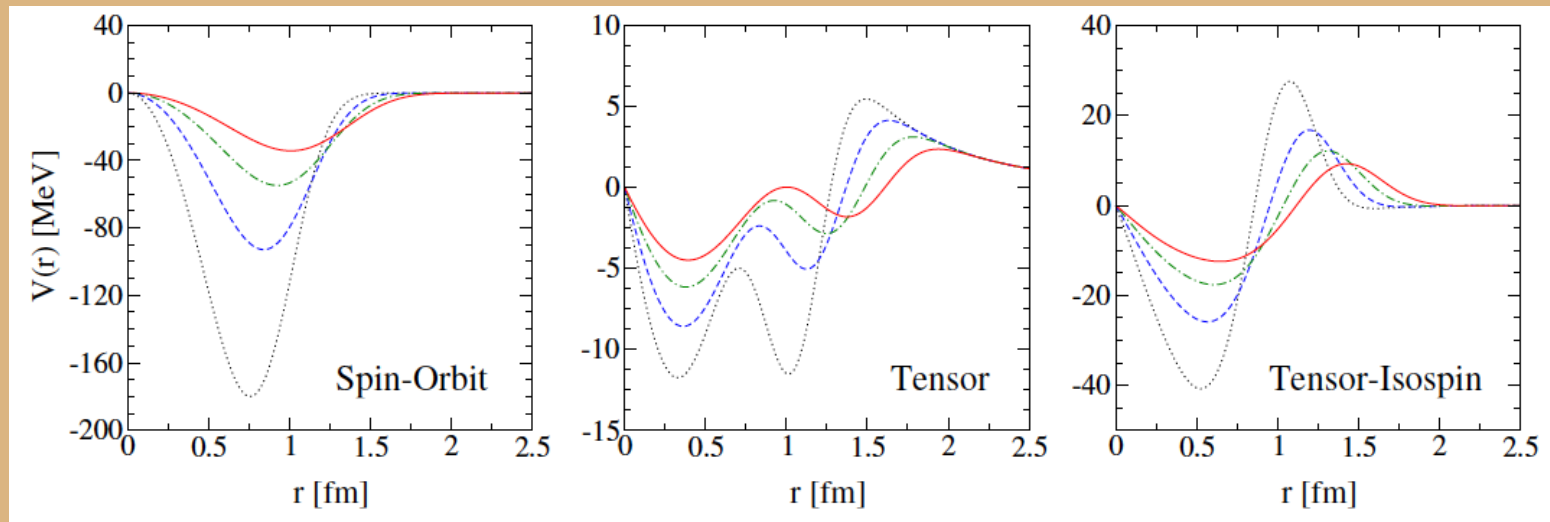
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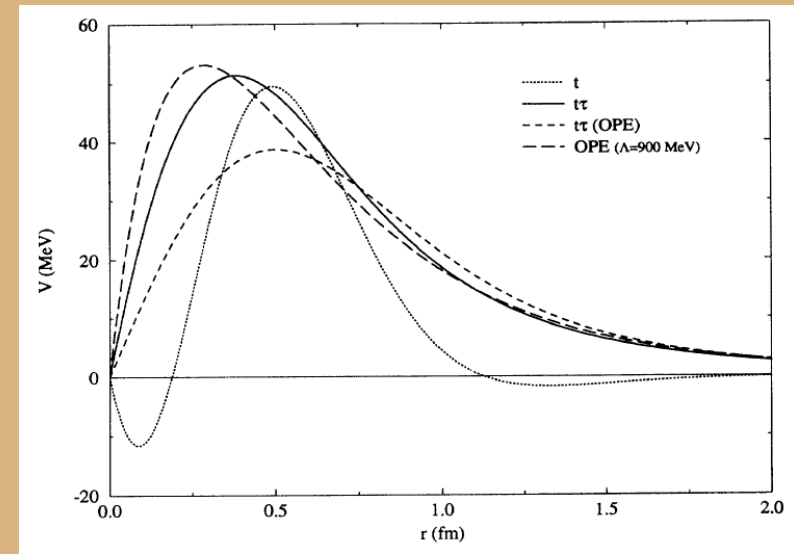
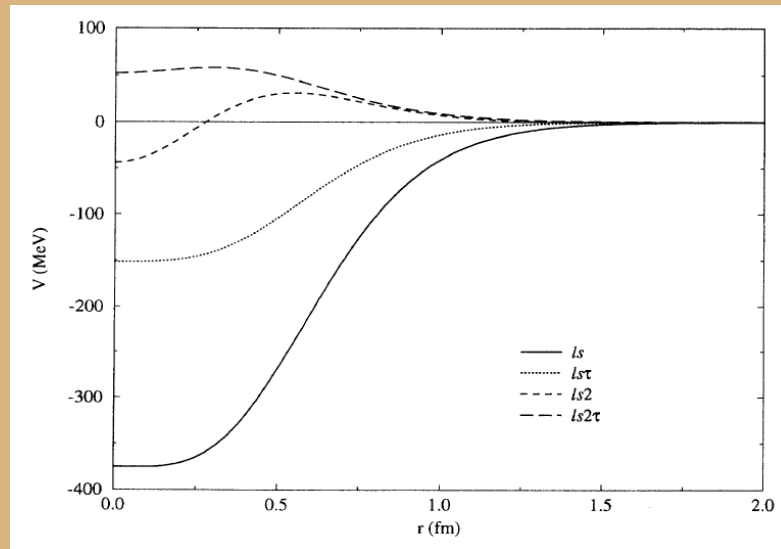
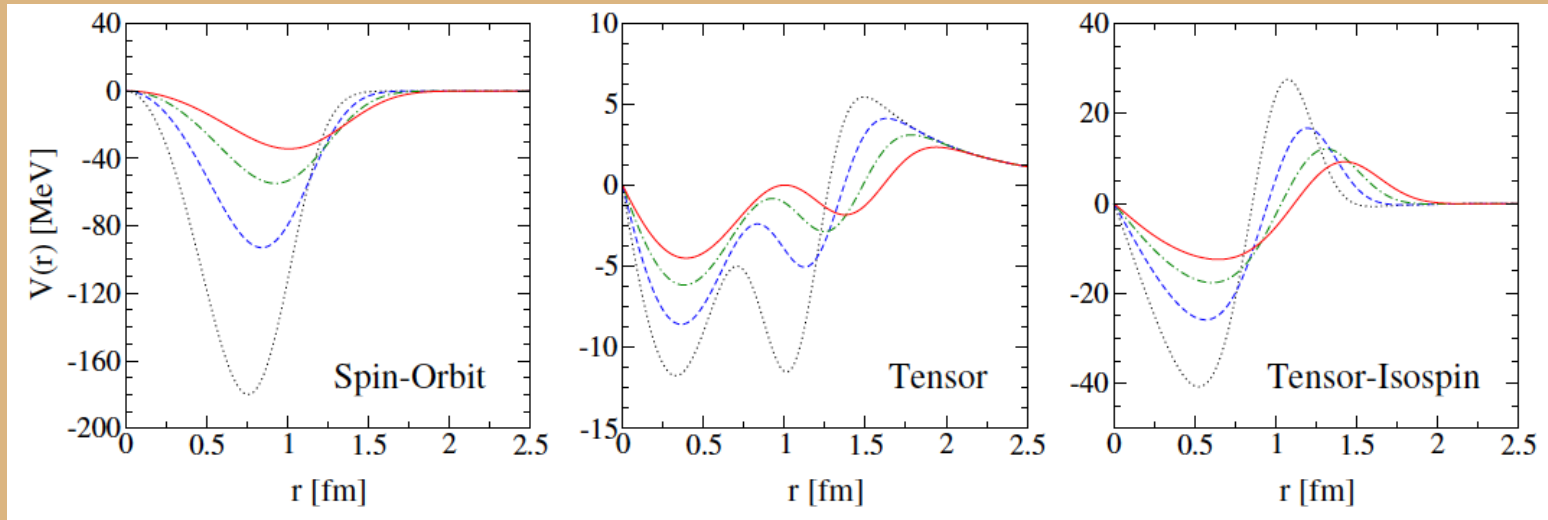
Compare  
with  
AV18



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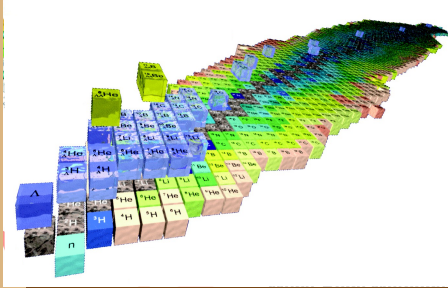
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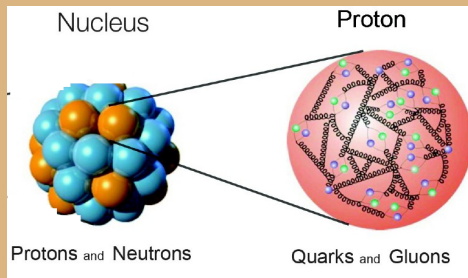


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## Nuclear forces

- Chiral Effective Field Theory (EFT)
- Local chiral EFT

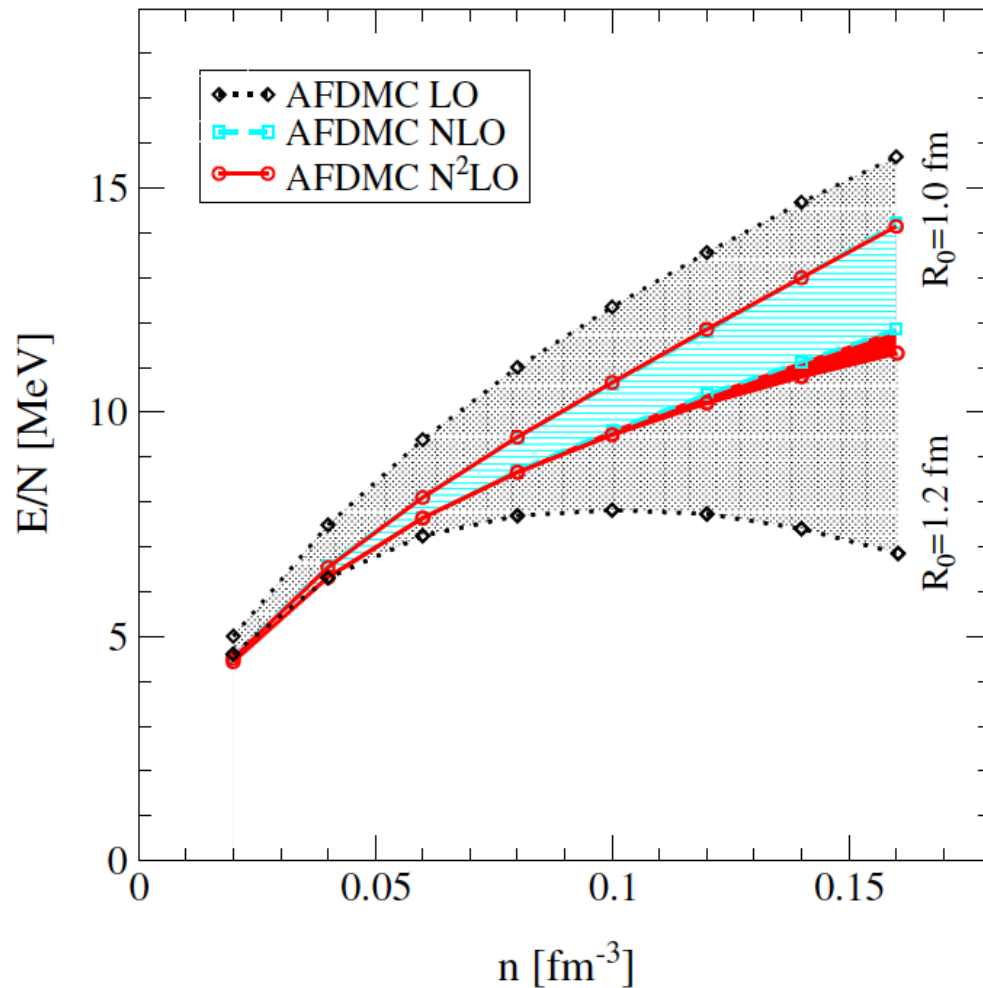


Credit: Bernhard Reischl

## Few- and many-body results

- Pure neutron matter
- Light nuclei

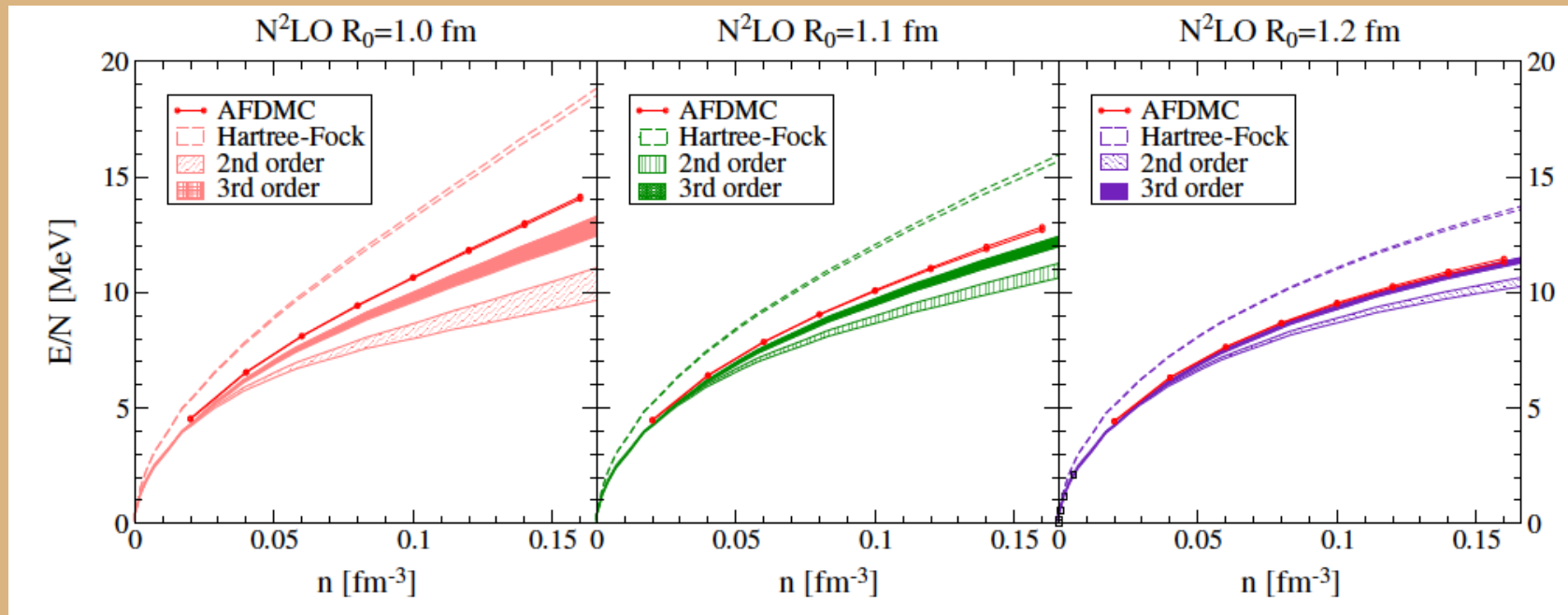
# Chiral EFT in QMC: neutron matter



- Use Auxiliary-Field Diffusion Monte Carlo to handle the full interaction
- First ever non-perturbative systematic error bands
- Band sizes to be expected
- Many-body forces will emerge systematically

**NEUTRONS**

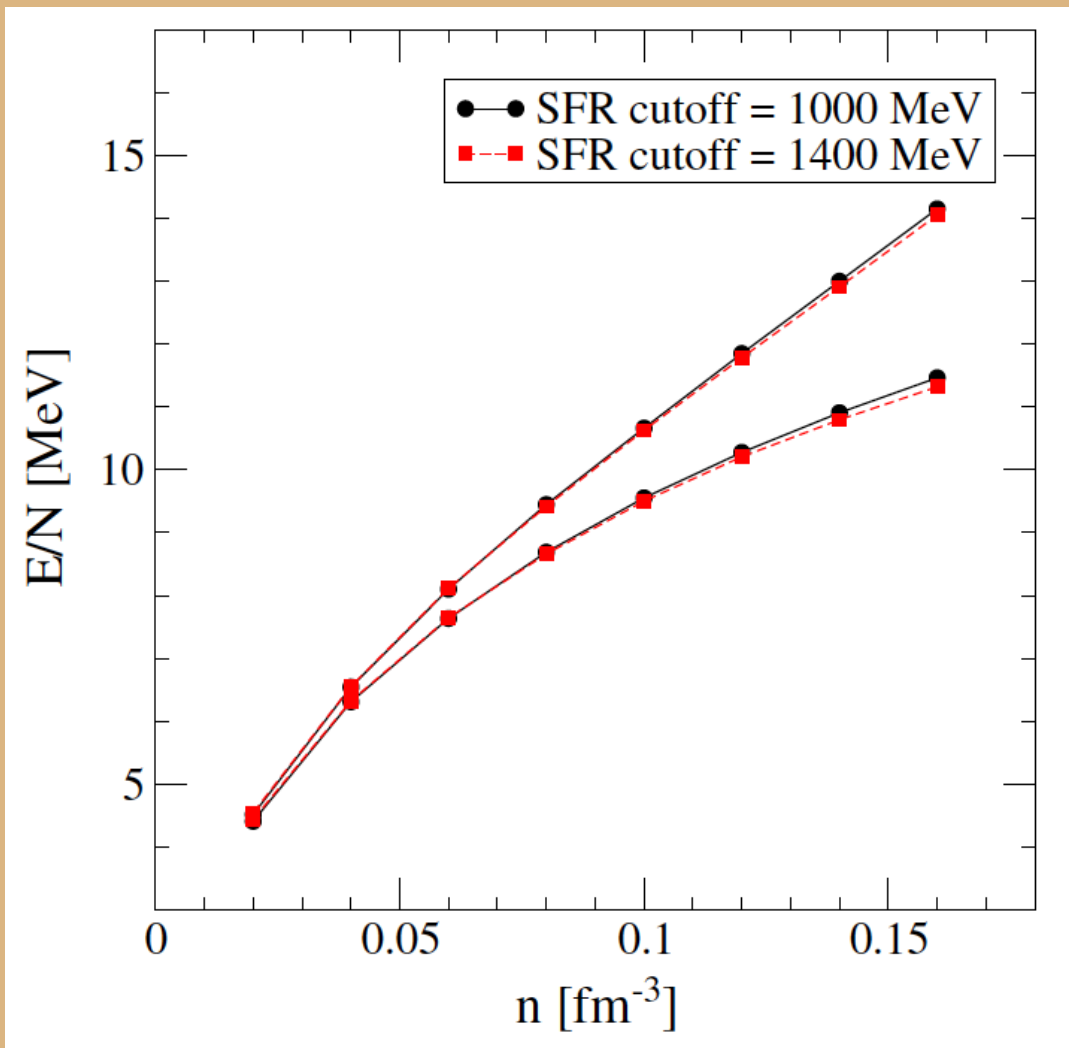
# QMC vs MBPT



- MBPT bands come from diff. single-particle spectra
- Soft potential in excellent agreement with AFDMC

NEUTRONS

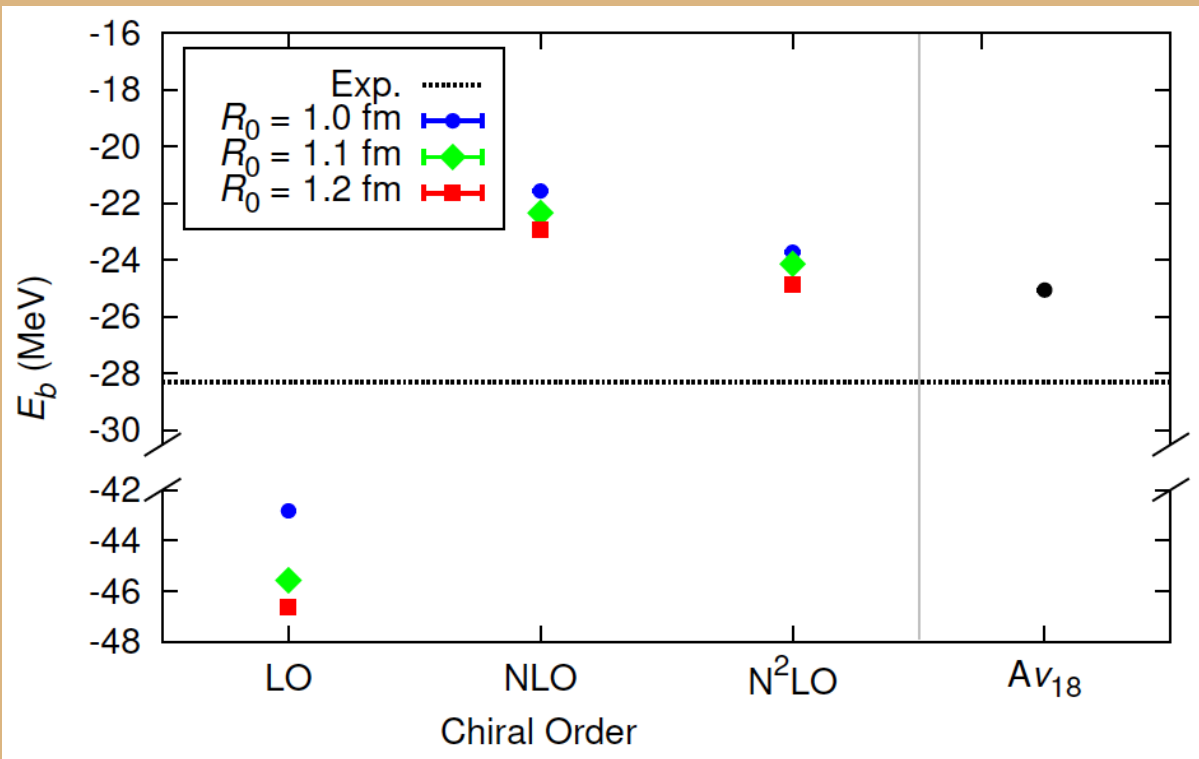
# Varying the SFR cutoff



- SFR cutoff expected not to matter
- Spectral-function regularization cutoff (kept fixed before) now varied
- Clearly shows lack of dependence

**NEUTRONS**

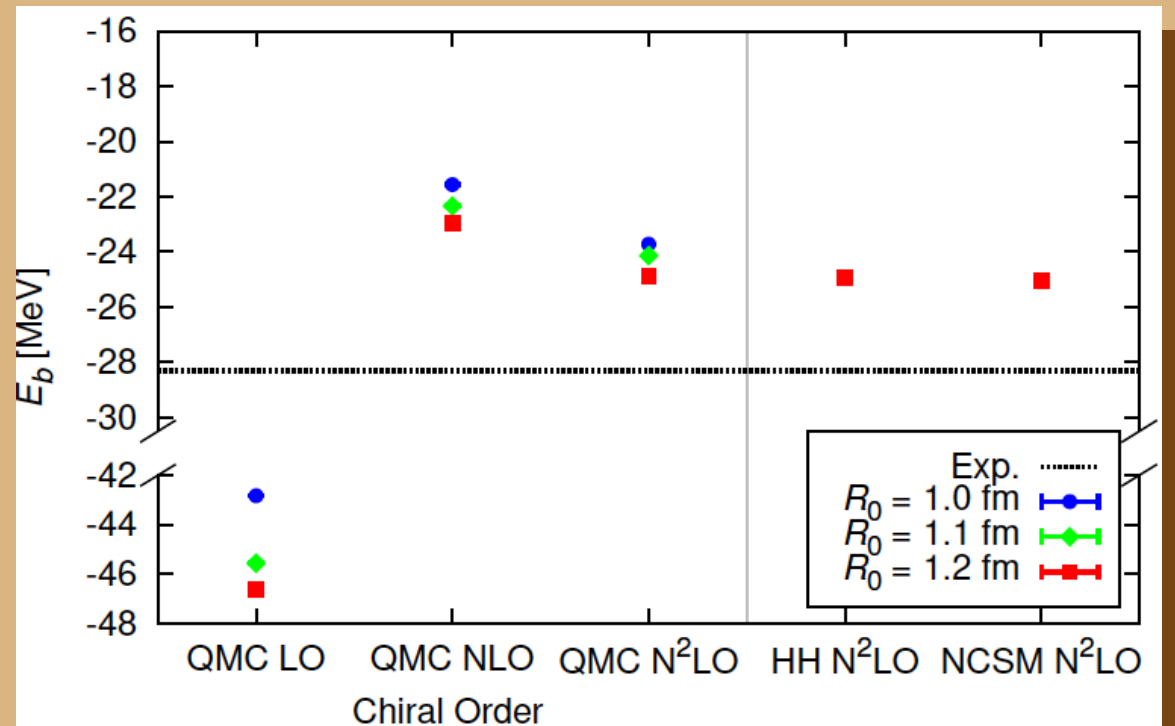
# Nuclear GFMC for light nuclei



- Binding energy of  ${}^4\text{He}$
- Non-perturbative systematic error bands
- All results are strong force + Coulomb, no NNN

NUCLEONS

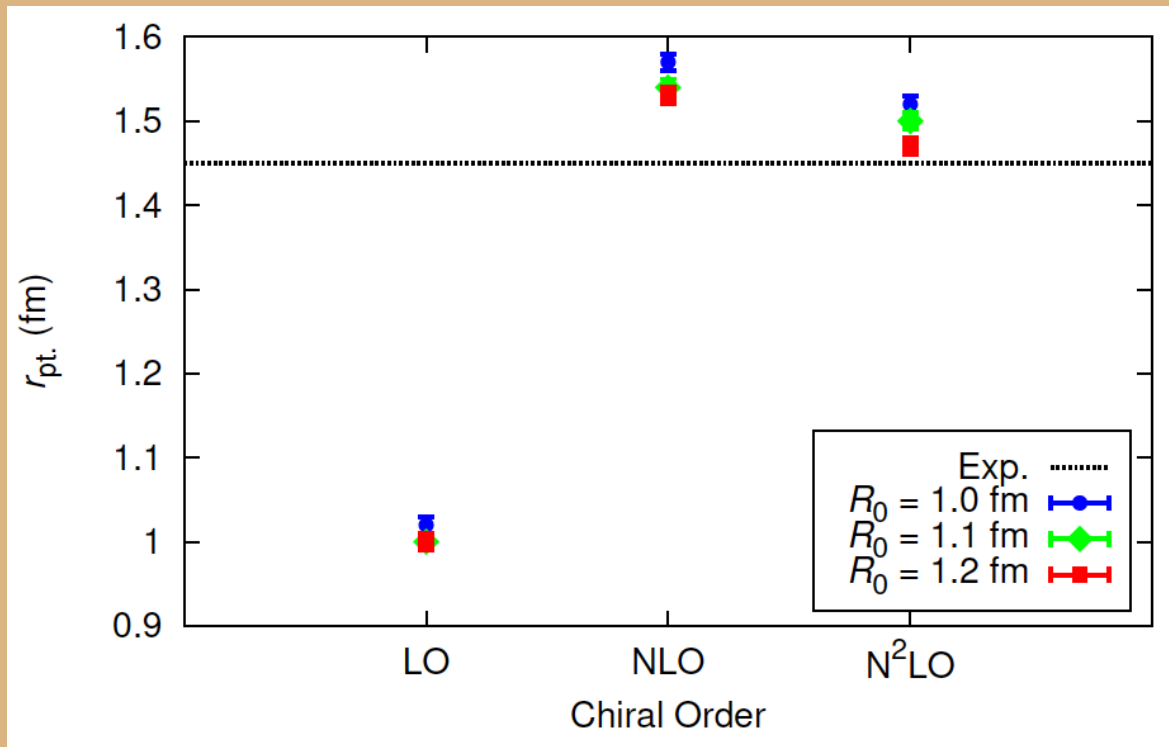
# Local chiral EFT in other methods



- Binding energy of  $^4\text{He}$
- Detailed benchmarking with hyperspherical harmonics (HH) and no-core shell model (NCSM)
- Probing perturbativeness

NUCLEONS

# Nuclear GFMC for light nuclei



- rms point radius of  $^4\text{He}$
- Non-perturbative systematic error bands
- All results are strong force + Coulomb, no NNN

NUCLEONS

# Conclusions

- Chiral EFT can now be used in continuum Quantum Monte Carlo methods
- Non-perturbative systematic error bands can be produced
- Back to TOV as soon as three-neutron forces have been incorporated



# Conclusions

- Chiral EFT can now be used in continuum Quantum Monte Carlo methods
- Non-perturbative systematic error bands can be produced
- Back to TOV as soon as three-neutron forces have been incorporated
- “It is ywrite that every thing Hymself sheweth in the tastyng”

# Acknowledgments

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## Collaborators

- Sonia Bacca (TRIUMF)
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