# Effects of Time Ordering in Parametric-Down Conversion

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## Outline

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#### 1 Parametric Down-Conversion and the Magnus expansion

### 2 Entanglement and Time Ordering in SPDC



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### **2** Entanglement and Time Ordering in SPDC











In a 1-D structure  $\hat{H}^{PDC}(t) =$ 



$$\varepsilon \int d\omega_a d\omega_b d\omega_p \Phi e^{i\Delta t} \hat{c}(\omega_p) \hat{a}^{\dagger}(\omega_a) \hat{b}^{\dagger}(\omega_b) + \text{h.c.}$$

- Undepleted pump approximation: operator
   c → scalar α(ω<sub>p</sub>).
- *H*<sup>PDC</sup> is Quadratic Bosonic Hamiltonian (QBH)

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- H<sup>PDC</sup> is a Quadratic Bosonic Hamiltonian.
- Two mode squeezing generator.
- $[H^{\text{PDC}}(t), H^{\text{PDC}}(t')] \neq 0$

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For each k we linearize as follows:

$$k_x(\omega_x) pprox k_x(ar \omega_x) + rac{dk_x}{d\omega_x}\Big|_{\omega_x = ar \omega_x}(\omega_x - ar \omega_x) = k_x(ar \omega_x) + rac{\delta\omega_x}{v_x}$$

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The central frequencies  $\bar{\omega}_x$  are such that:

 $\bar{\omega}_a + \bar{\omega}_b = \bar{\omega}_p$  Energy conservation  $k_a(\bar{\omega}_a) + k_b(\bar{\omega}_b) = k_p(\bar{\omega}_p)$  Momentum conservation

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$$\Delta k = \underbrace{k_{a}(\bar{\omega}_{a}) + k_{b}(\bar{\omega}_{b}) - k_{p}(\bar{\omega}_{p})}_{=0} + \frac{\delta\omega_{a}}{v_{a}} + \frac{\delta\omega_{b}}{v_{b}} - \frac{\delta\omega_{p}}{v_{p}}$$

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Dyson

$$\hat{\mathcal{U}}(t,t_0) = \mathbb{I} - i \int_{t_0}^t dt_1 \hat{H}(t_1) + (-i)^2 \int_{t_0}^t dt_1 \int_{t_0}^{t_1} dt_2 \hat{H}(t_1) \hat{H}(t_2) +$$

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$$\begin{split} \Omega_{1} &= -i \int dt' \hat{H}(t') = -2\pi i \varepsilon \times \\ &\int d\omega_{a} d\omega_{b} J_{1}(\omega_{a}, \omega_{b}) \\ &\times \hat{a}^{\dagger}(\omega_{a}) \hat{b}^{\dagger}(\omega_{b}) + \text{h.c.} \\ J_{1} &= \alpha(\omega_{a} + \omega_{b}) \times \\ &\Phi(\omega_{a}, \omega_{b}, \omega_{a} + \omega_{b}) \end{split}$$

2<sup>nd</sup> Order Magnus Terms

$$\frac{(-i)^2}{2} \int dt_1 \int^{t_1} dt_2 \left[ \hat{H}(t_1), \hat{H}(t_2) \right] \\ \left[ a^{\dagger}(\omega_p) b^{\dagger}(\omega_q) + \text{h.c.}, \\ a^{\dagger}(\omega_r) b^{\dagger}(\omega_s) + \text{h.c.} \right] = \\ \delta(\omega_q - \omega_s) a^{\dagger}(\omega_r) a(\omega_p) \\ + \delta(\omega_p - \omega_r) b^{\dagger}(\omega_s) b(\omega_q) - \text{h.c.}$$

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$$\begin{split} |\psi\rangle &= e^{-2\pi i \int d\omega_1 d\omega_2 \left\{ (J(\omega_1,\omega_2)\hat{a}^{\dagger}(\omega_1)\hat{b}^{\dagger}(\omega_2) + \text{h.c.}) \right\}} \times \\ &e^{-2\pi i \int d\omega_1 d\omega_2 \left\{ G(\omega_1,\omega_2)\hat{a}^{\dagger}(\omega_1)\hat{a}(\omega_2) + H(\omega_1,\omega_2)\hat{b}^{\dagger}(\omega_1)\hat{b}(\omega_2) \right\}} |\text{vac}\rangle \end{split}$$

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$$\underbrace{J(\omega_a, \omega_b)}_{JSA} = \underbrace{\varepsilon J_1(\omega_a, \omega_b)}_{\text{Exact w.o. time order}} + \underbrace{\varepsilon^3(J_3(\omega_a, \omega_b) - iK_3(\omega_a, \omega_b))}_{\text{time order correction}} + \dots$$

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Magnus expansion is a perturbation theory "inside the exponential".

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### A simple Gaussian Model

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 $\Phi = \operatorname{sinc}(\Delta kL/2)$ 



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## First order Magnus term

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$$J_{1}(\omega_{a},\omega_{b}) = \alpha(\omega_{a}+\omega_{b})\Phi(\omega_{a},\omega_{b},\omega_{a}+\omega_{b}) = \frac{\varepsilon\tau}{\sqrt{\pi}}\exp\left(-\mathbf{u}\mathbf{N}\mathbf{u}^{T}\right)$$
$$\mathbf{N} = \begin{pmatrix} \mu_{a}^{2} & \mu^{2} \\ \mu^{2} & \mu_{b}^{2} \end{pmatrix}, \quad \mathbf{u} = (\delta\omega_{a},\delta\omega_{b})$$
$$\eta_{a/b} = s_{p} - s_{a/b}, \quad \mu^{2} = \tau^{2} + \eta_{a}\eta_{b}, \quad \mu^{2}_{a/b} = \tau^{2} + \eta^{2}_{a/b}.$$

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# (First Order) Entanglement properties

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Unentangled photons  $\mu \sim 0$   $(s_a - s_p) \sim -(s_b - s_p) \sim \pm \tau$ 

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Entangled photons det(**N**)  $\sim$  0  $v_p \tau / |v_x - v_p| \ll L/v_x$ 

# (First Order) Entanglement properties





## Magnus correction



time order correction

For an arbitrary  $\Phi$  and  $\alpha$ 

The Joint Spectral amplitude is now complex

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### For the gaussian model

$$r \equiv \frac{\max_{\omega_a,\omega_b} |\varepsilon^3 J_3(\omega_a,\omega_b)|}{\max_{\omega_a,\omega_b} |\varepsilon J_1(\omega_a,\omega_b)|} < \underbrace{\frac{13\pi^2 \varepsilon^2 \tau^2}{6\sqrt{4(\eta_a - \eta_b)^2 \tau^2 + 3(\eta_a \eta_b + \tau^2)^2}}}_{\equiv r_{\max} \sigma}$$

if  $(\det(\mathbf{N}) \sim 0)$ 

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$$s_a-s_p=-(s_b-s_p)=\pm au \ (\mu=0),$$

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$$egin{array}{lll} s_a - s_p &= -(s_b - s_p) = \pm au \; (\mu = 0), \ r_{\mathsf{max}} &= rac{13\epsilon^2}{24} \end{array}$$

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On the Magnus expansion

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- Unitarity and Gaussian preservation come for free!

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- Time ordering effects give the Joint Spectral Amplitude a non-trivial phase structure.
- Time ordering effects are irrelevant for very broad phase matching functions.