

The Unruh Effect in a Cavity

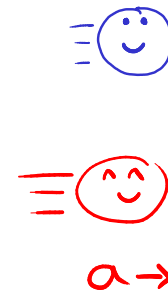
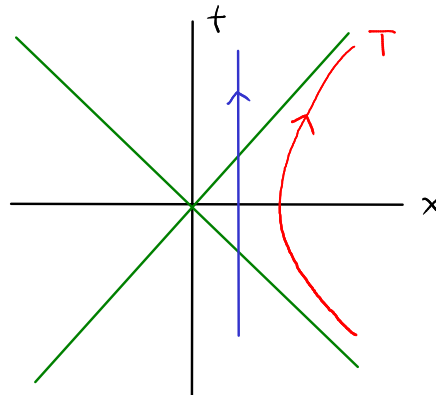
To the Continuum and Beyond

Wilson Brenna

June 18, 2014

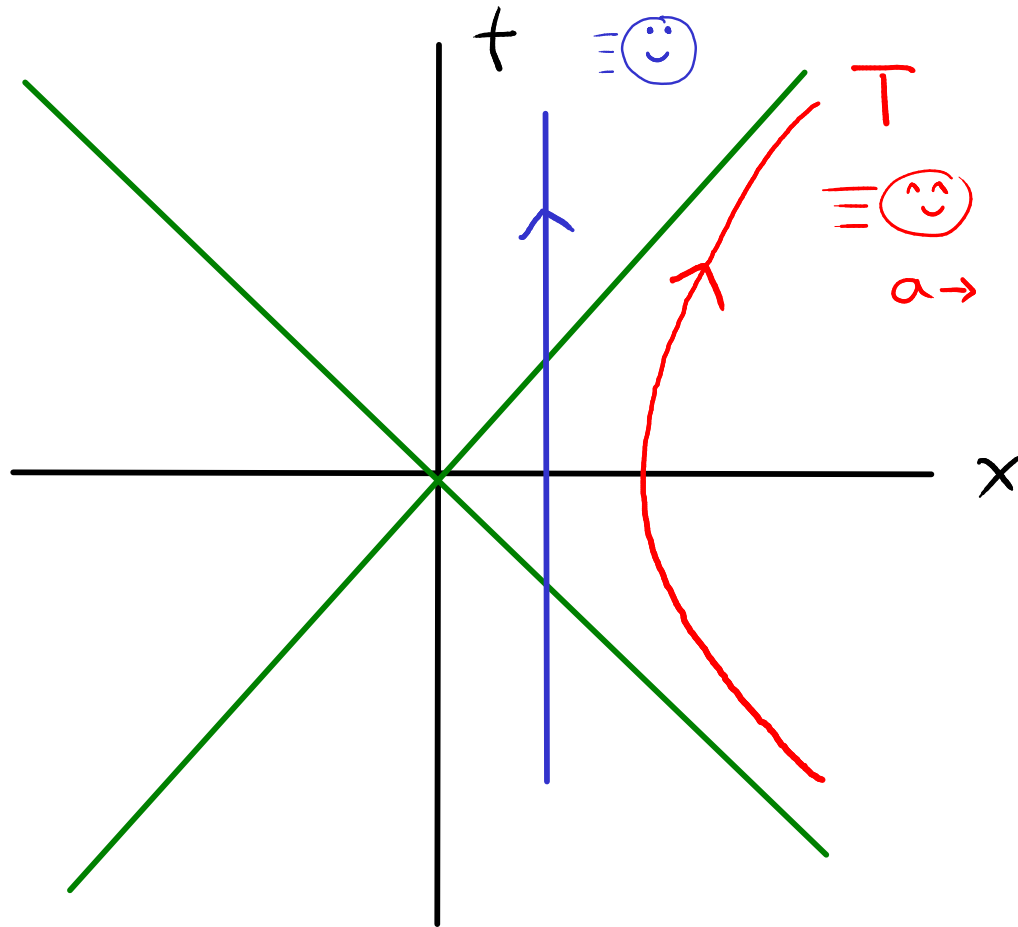
The Unruh Effect

$$T = \frac{a}{2\pi}$$



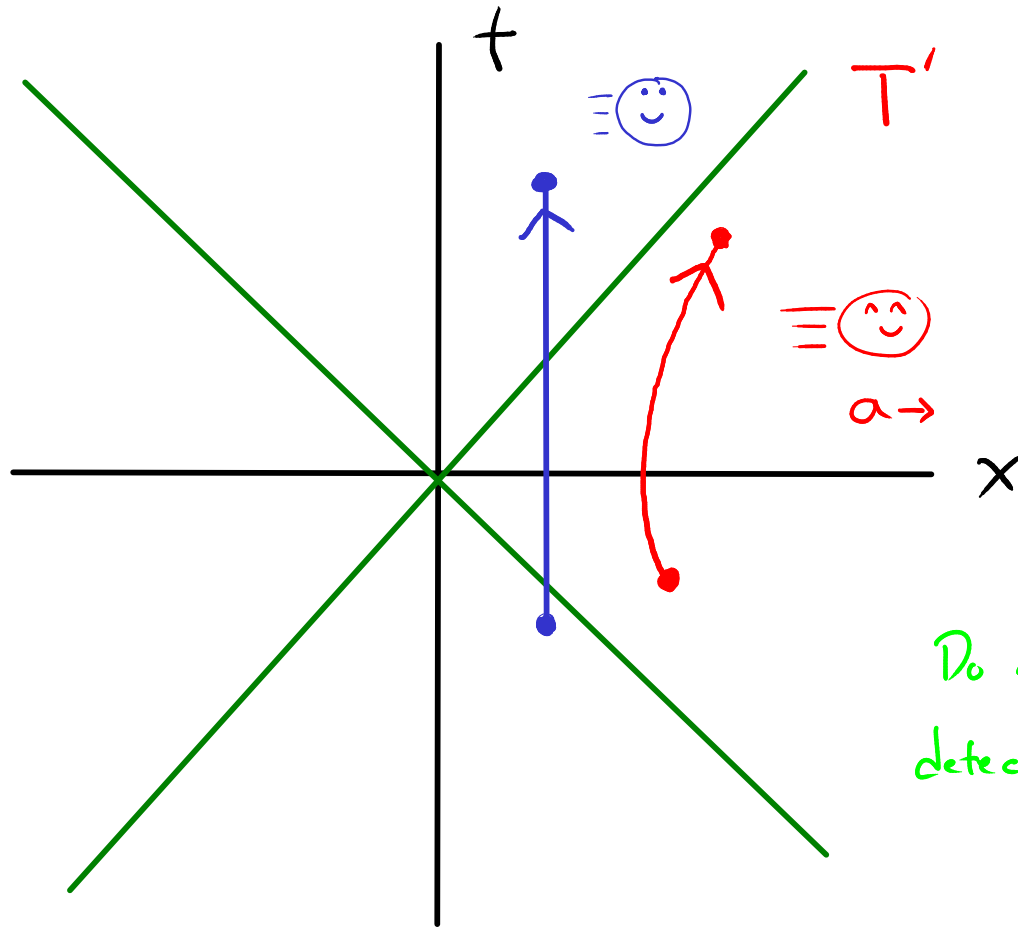
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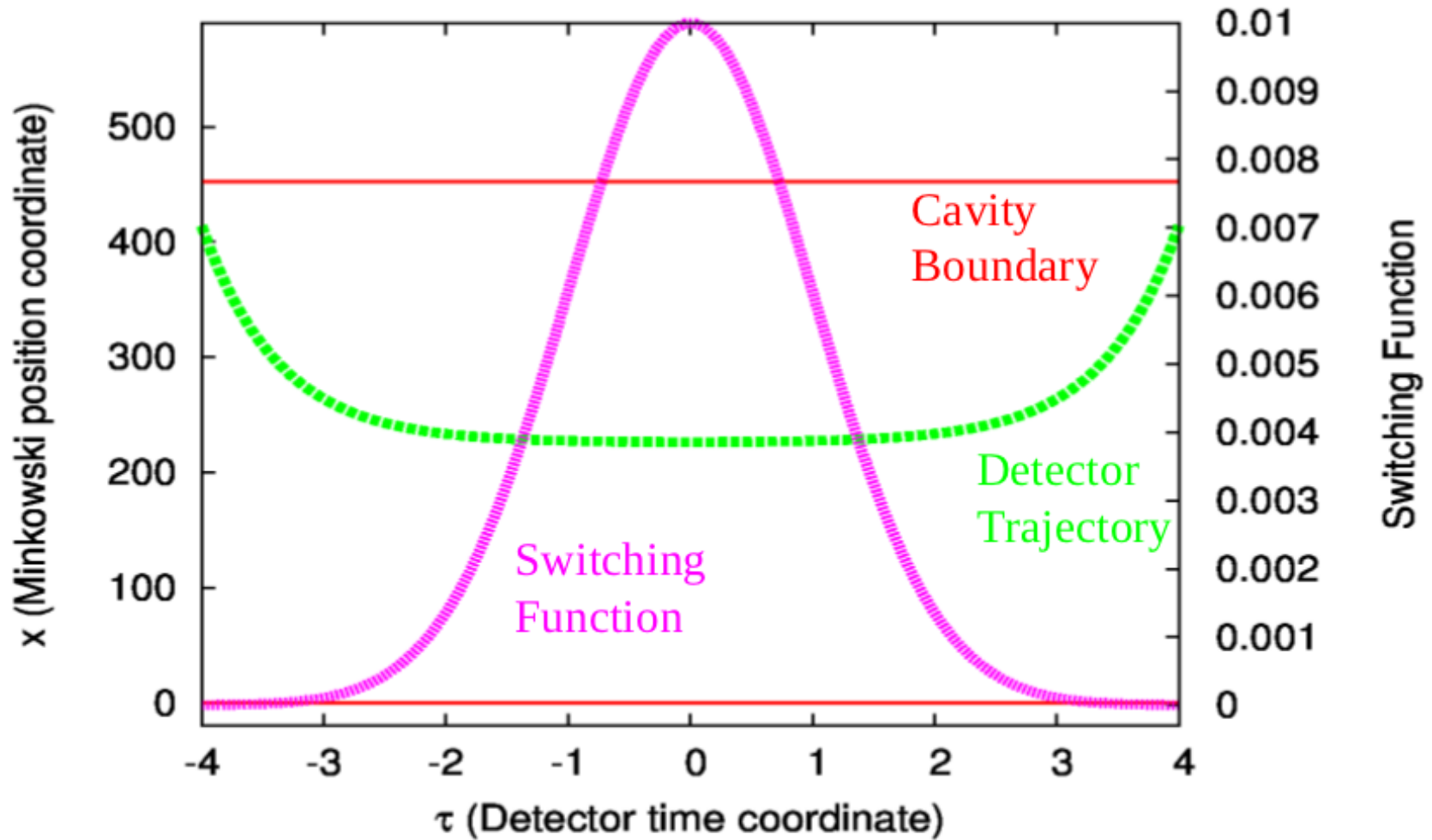
Do accelerated
detectors get HOT?

Outline

1. Review: The Unruh-DeWitt Detector in 2D and Gaussian Quantum Detectors
2. Temperature and Acceleration Plots
 - a. Perturbative
 - b. Non-Perturbative
3. Conclusion

I. Review of UdW and GQI Detectors

Trajectory of the Detector within the Cavity



I. Review of UdW and GQI Detectors

Notation

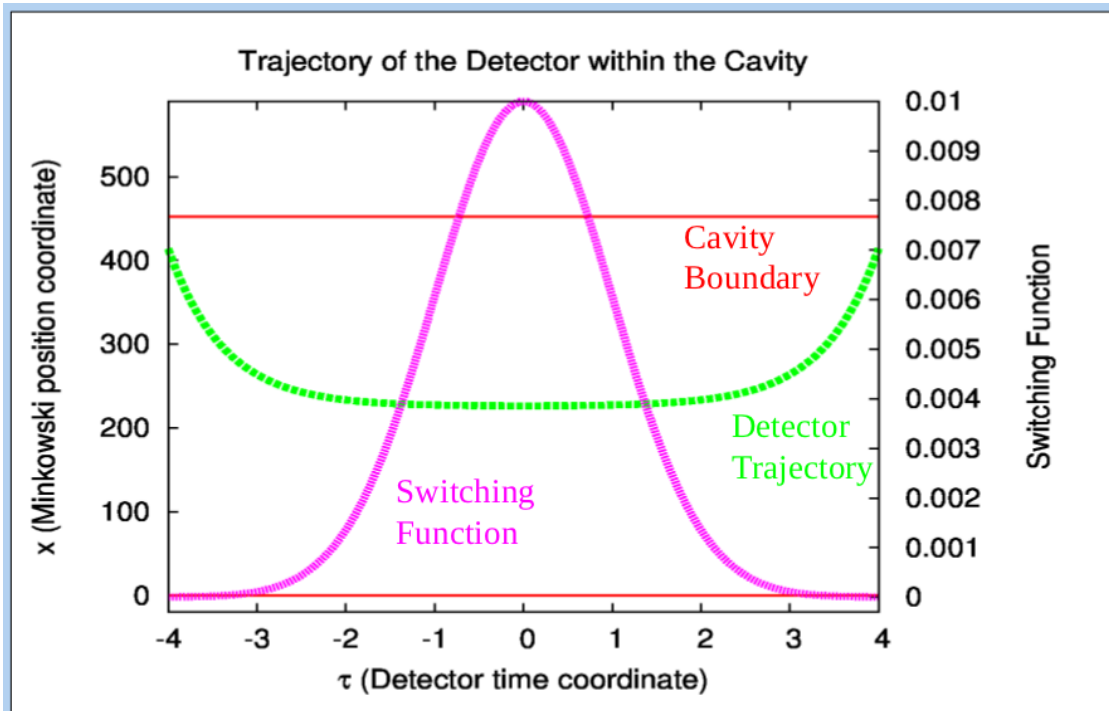
$\lambda(\tau)$: Switching Function = $\lambda_0 e^{-\tau^2/2\sigma^2}$ \uparrow width of switching σ

L : length of cavity

$\Delta\tau$: time of detector's trajectory

N : number of field modes

Ω : energy gap of detector



I. Review of UdW and GQI Detectors

UdW (Minkowski)

"Perturbative"

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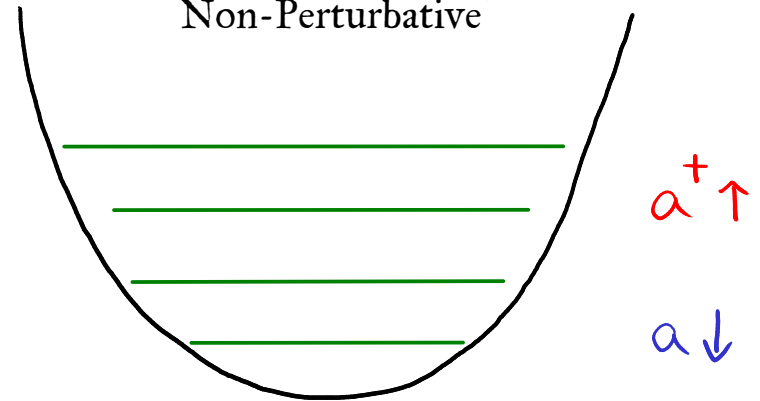
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$\hat{\rho}$: Density Matrix

$$\hat{\rho} = \sum_i p_i |\psi_i\rangle \langle \psi_i|$$

Gaussian

"Non-Perturbative"



σ : Covariance Matrix (State)

$$\sigma = \begin{bmatrix} \langle x^2 \rangle & \langle xp + px \rangle \\ \langle xp + px \rangle & \langle p^2 \rangle \end{bmatrix}$$

I. Review of UdW and GQI Detectors

$$\dot{\mathcal{F}}_{\Delta t}(\Omega) = \int_{-\tau+\tau_0}^{\tau-\tau_0} ds \cdot e^{-i\Omega s} \frac{1}{2\pi} \int \frac{dk}{2\omega} e^{-i\omega t(s)+ik(x(s)-1)}$$

$$\dot{\mathcal{F}}_{\infty}(\Omega) = \int_{-\infty}^{\infty} ds \cdot e^{-i\Omega s - \frac{|s|}{\kappa}} \frac{1}{2\pi} \int \frac{dk}{2\omega} e^{-i\omega\alpha \sinh(s/\alpha) + ik(\alpha \cosh(s/\alpha) - 1)}$$

After some mathematical tricks, this evaluates to

$$\dot{\mathcal{F}}_{\infty} = \frac{1}{\Omega (e^{2\pi\alpha\Omega} - 1)} \rightsquigarrow \text{Boltzmann Dist. at } T = \frac{1}{2\pi\alpha} = \frac{a}{2\pi}$$

We numerically evaluate the above integrals, after adding UV and IR cutoffs:

$$k = \frac{2\pi n}{L} \quad n = [1..N]$$

I. Review of UdW and GQI Detectors

$$\boldsymbol{\sigma}_d = \begin{pmatrix} \langle \hat{q}_d^2 \rangle & \langle \hat{q}_d \hat{p}_d + \hat{p}_d \hat{q}_d \rangle \\ \langle \hat{q}_d \hat{p}_d + \hat{p}_d \hat{q}_d \rangle & \langle \hat{p}_d^2 \rangle \end{pmatrix}$$

$$\hat{H} = \Omega_d \hat{a}_d^\dagger \hat{a}_d + \frac{dt}{d\tau} \sum_n \omega_n \hat{a}_n^\dagger \hat{a}_n + \hat{H}_I(\tau)$$

$$\hat{H}_I = \lambda(\tau) \cdot \hat{m} \cdot \hat{\phi}[x(\tau)]$$

$$\hat{H} = \hat{\boldsymbol{x}}^T \mathbb{F}(\tau) \hat{\boldsymbol{x}}$$

$$\hat{\boldsymbol{x}}^T = (\hat{q}_d, \hat{p}_d, \hat{q}_1, \hat{p}_1, \dots, \hat{q}_N, \hat{p}_N)$$

States are Gaussian, i.e. they can be squeezed, shifted, thermal, etc.

Quadratic Hamiltonians preserve Gaussianity.

I. Review of UdW and GQI Detectors

- Symplectic form Ω_{ij} from $[\hat{\mathbf{a}}_i, \hat{\mathbf{a}}_j] = \Omega_{ij}$
$$\begin{pmatrix} 0 & 1 & & & & & & & \\ -1 & 0 & & & & & & & \\ & & 0 & 1 & & & & & \\ & & -1 & 0 & & & & & \\ & & & & 0 & 1 & & & \\ & & & & -1 & 0 & & & \\ & & & & & & \dots & & \\ & & & & & & & & 0 \end{pmatrix}$$

where $\hat{\mathbf{a}} \equiv (\hat{a}_d, \hat{a}_d^\dagger, \hat{a}_1, \hat{a}_1^\dagger, \hat{a}_2, \hat{a}_2^\dagger, \dots, \hat{a}_N^\dagger)^T$
- Gaussian states and quadratic Hamiltonian⁴

$$\sigma(t) = S(t)\sigma_0 S(t)^T \quad (\text{Unitary Evolution})$$

$$\frac{d}{dt} S(t) = \Omega F^{sym}(t) \cdot S(t)$$

F^{sym} : symmetrization of the Hamiltonian in matrix form, $\mathbb{F}(t)$

Numerically evolve the first order DE from initial conditions!

I. Review of UdW and GQI Detectors

Detector Covariance
Matrix

$$\sigma_d = \begin{pmatrix} (1 + \delta)e^{-r} & 0 \\ 0 & (1 + \delta)e^r \end{pmatrix}$$

Unruh Effect

$$T = \Omega \left[\ln \left(1 + \frac{2}{\delta} \right) \right]^{-1}$$

$$T \sim a$$

• linearity

• slope

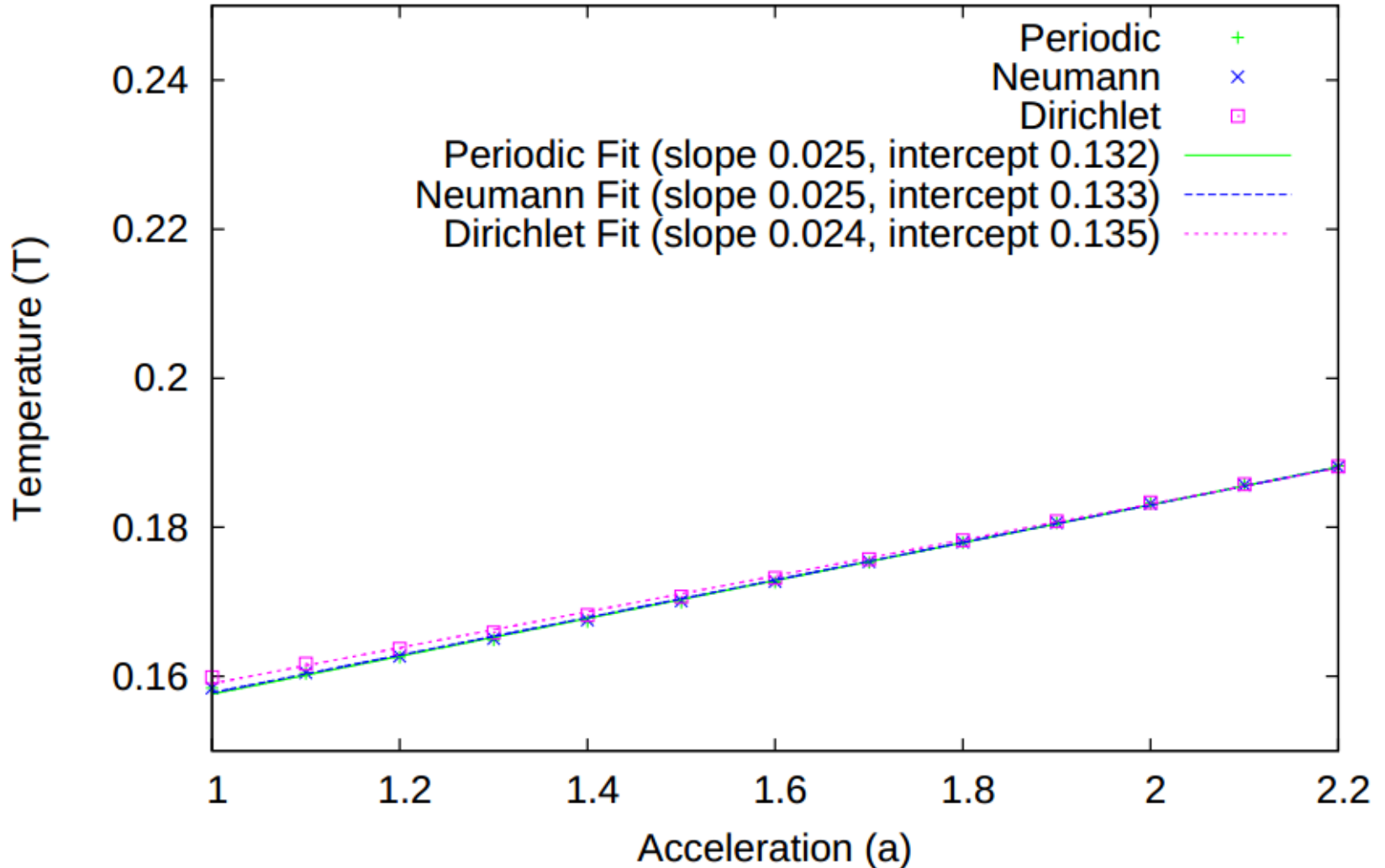
Thermality

$$E \approx \Omega \left(1 + \delta + \frac{1}{2}r^2 \right)$$

$$\therefore r^2 \ll \delta \implies \text{thermality}$$

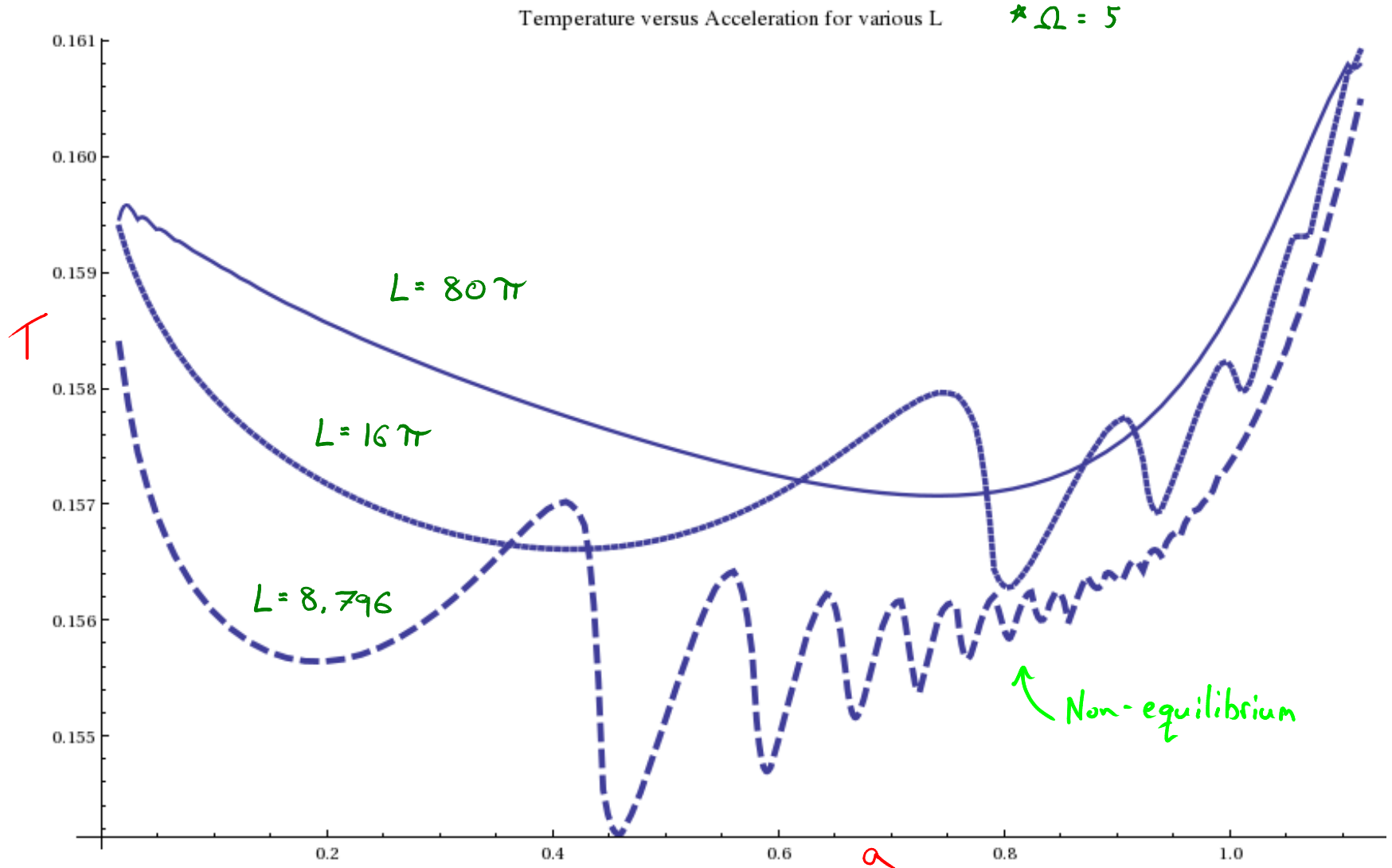
I. Review of UdW and GQI Detectors

Comparing Boundary Conditions



2. Temperature and Acceleration Plots

a. Perturbative (Unruh-DeWitt)

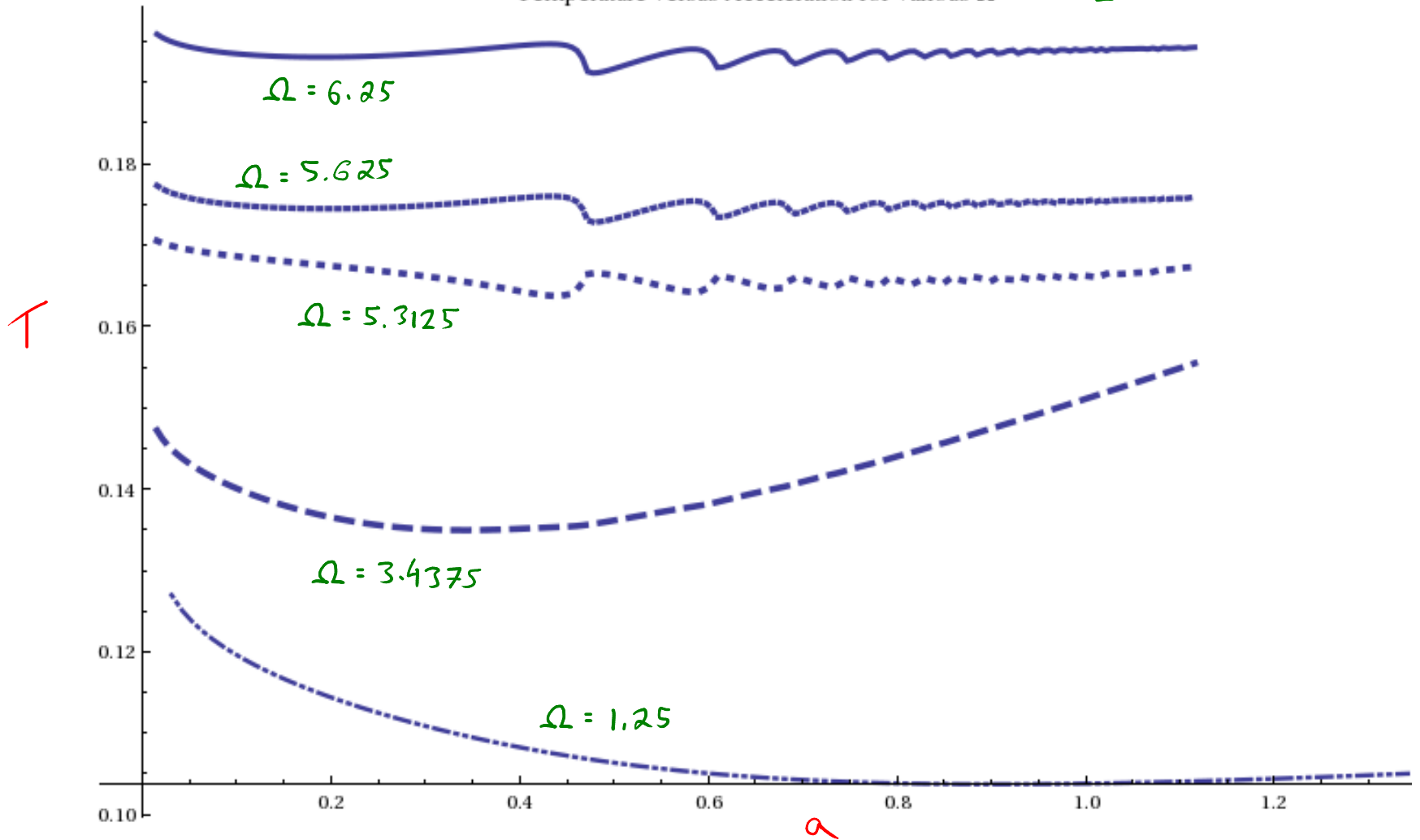


2. Temperature and Acceleration Plots

a. Perturbative

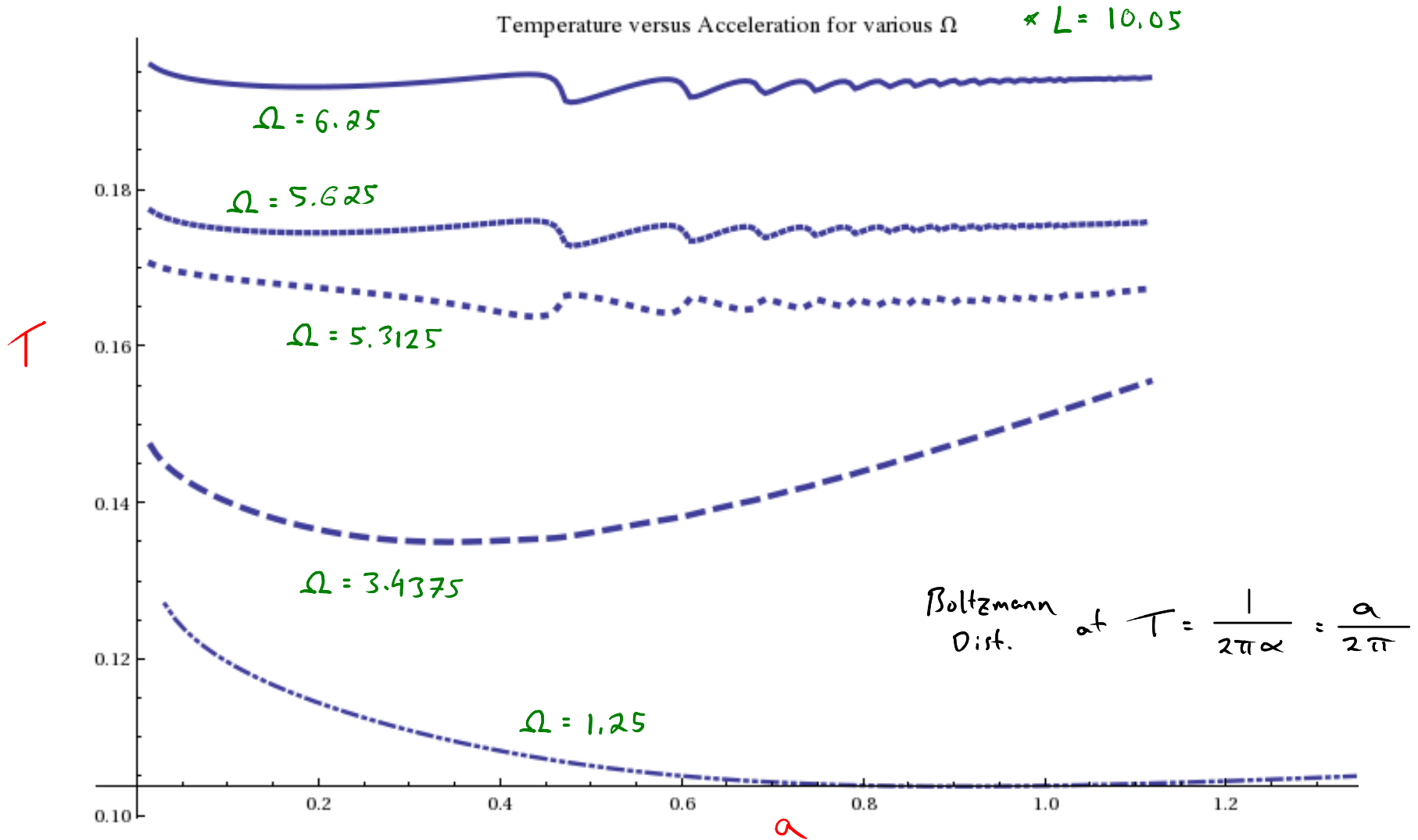
Temperature versus Acceleration for various Ω

* $L = 10.05$



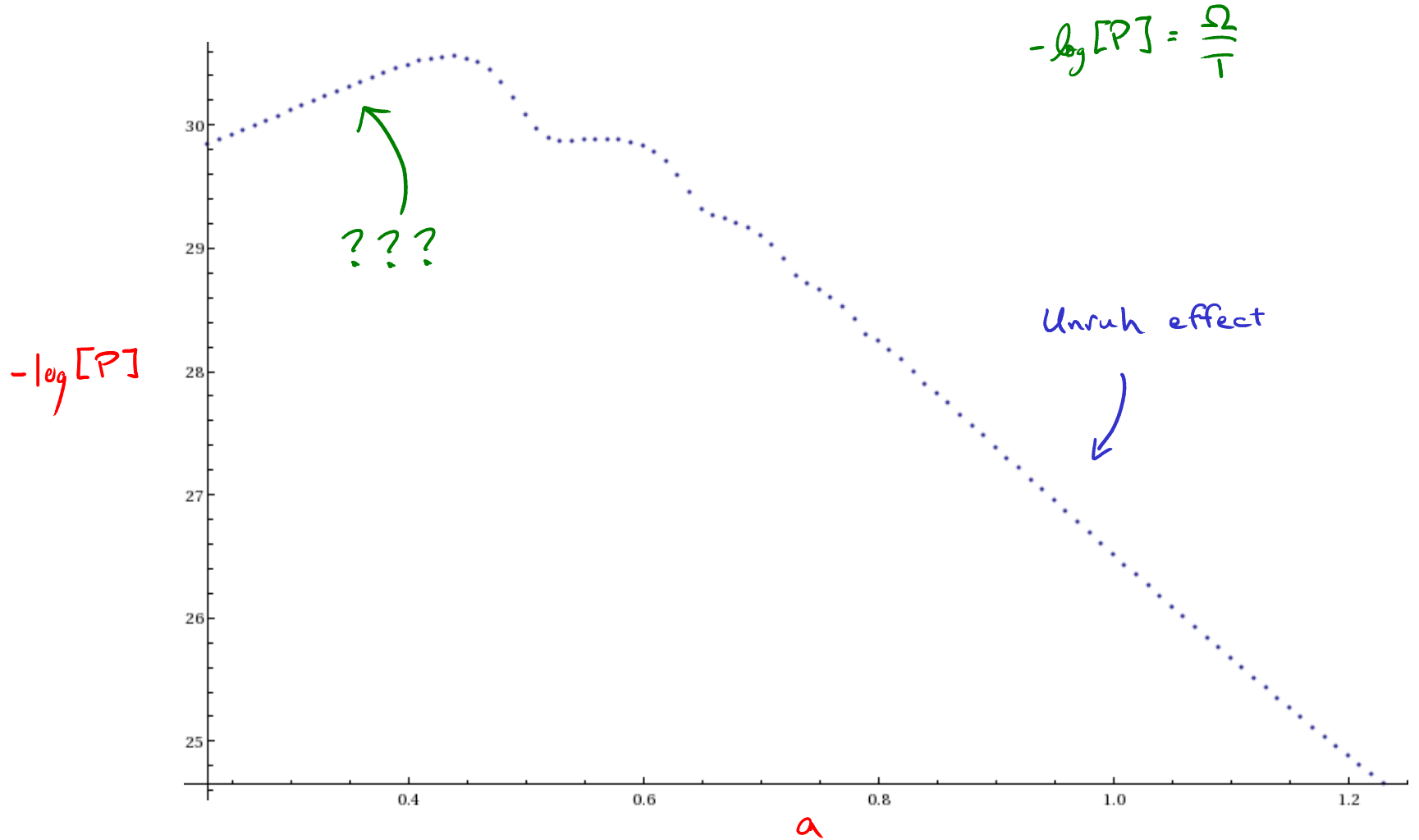
2. Temperature and Acceleration Plots

a. Perturbative



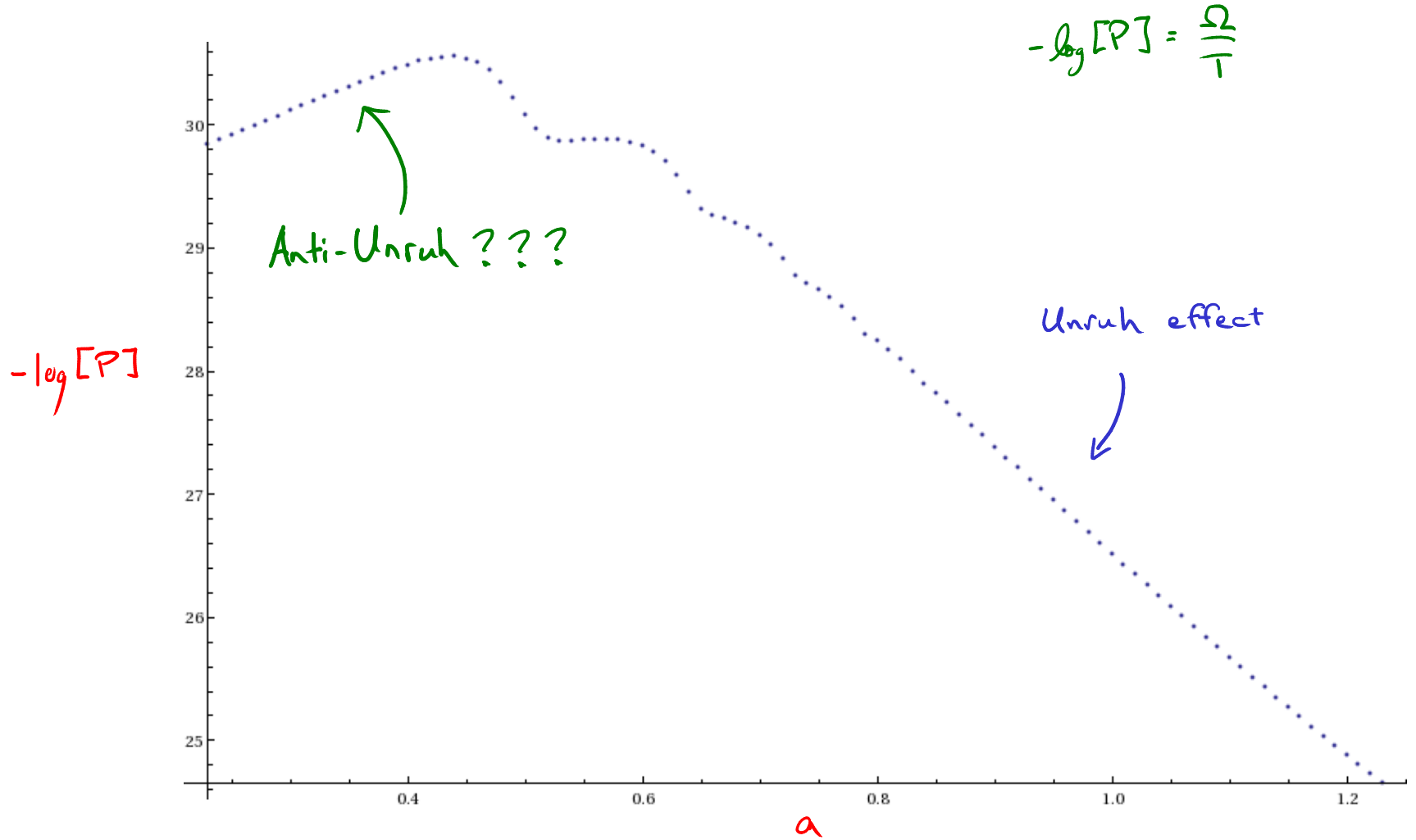
2. Temperature and Acceleration Plots

a. Perturbative



2. Temperature and Acceleration Plots

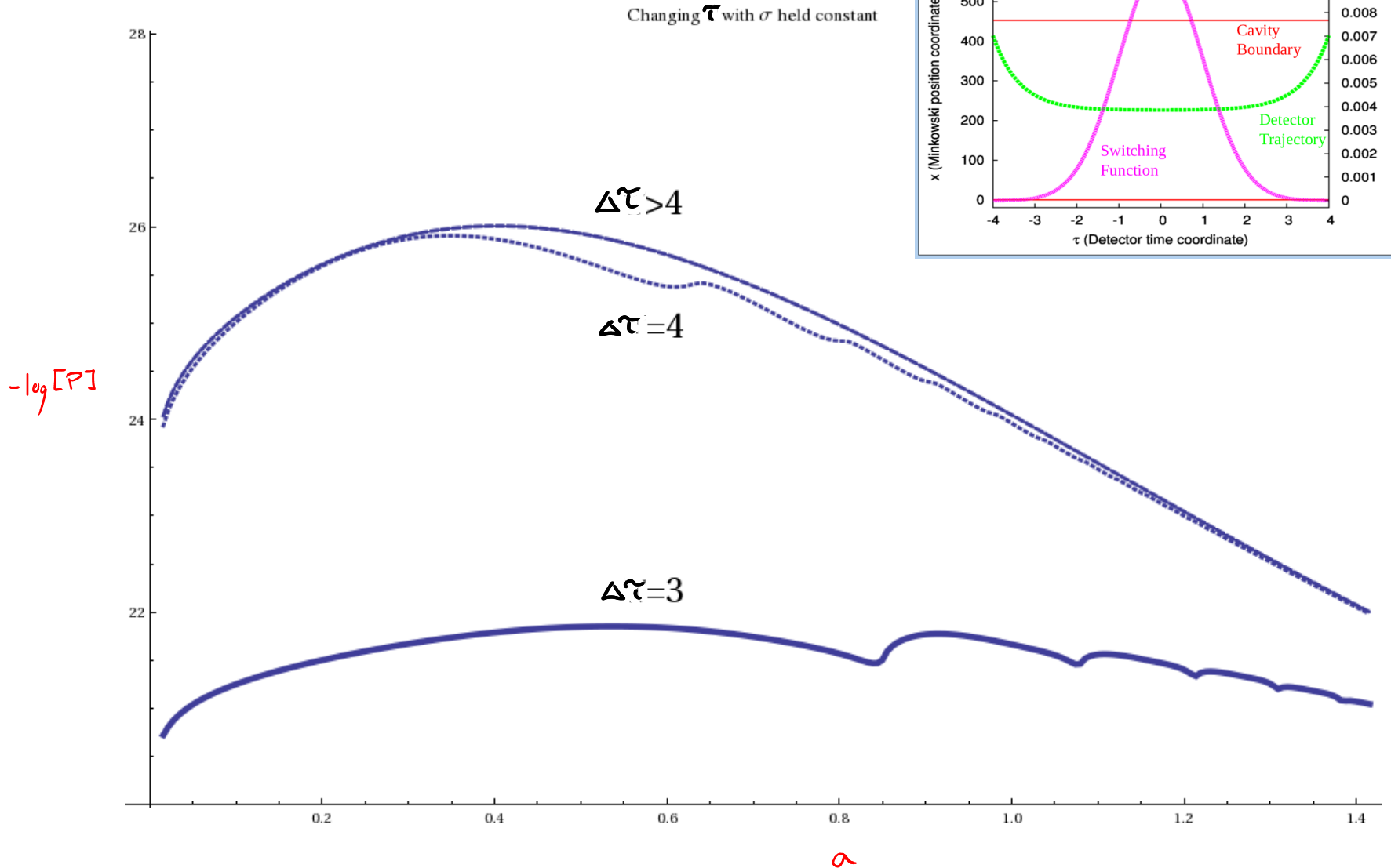
a. Perturbative



2. Temperature and Acceleration Plots

$$L = 4\pi \quad \Omega = 4$$

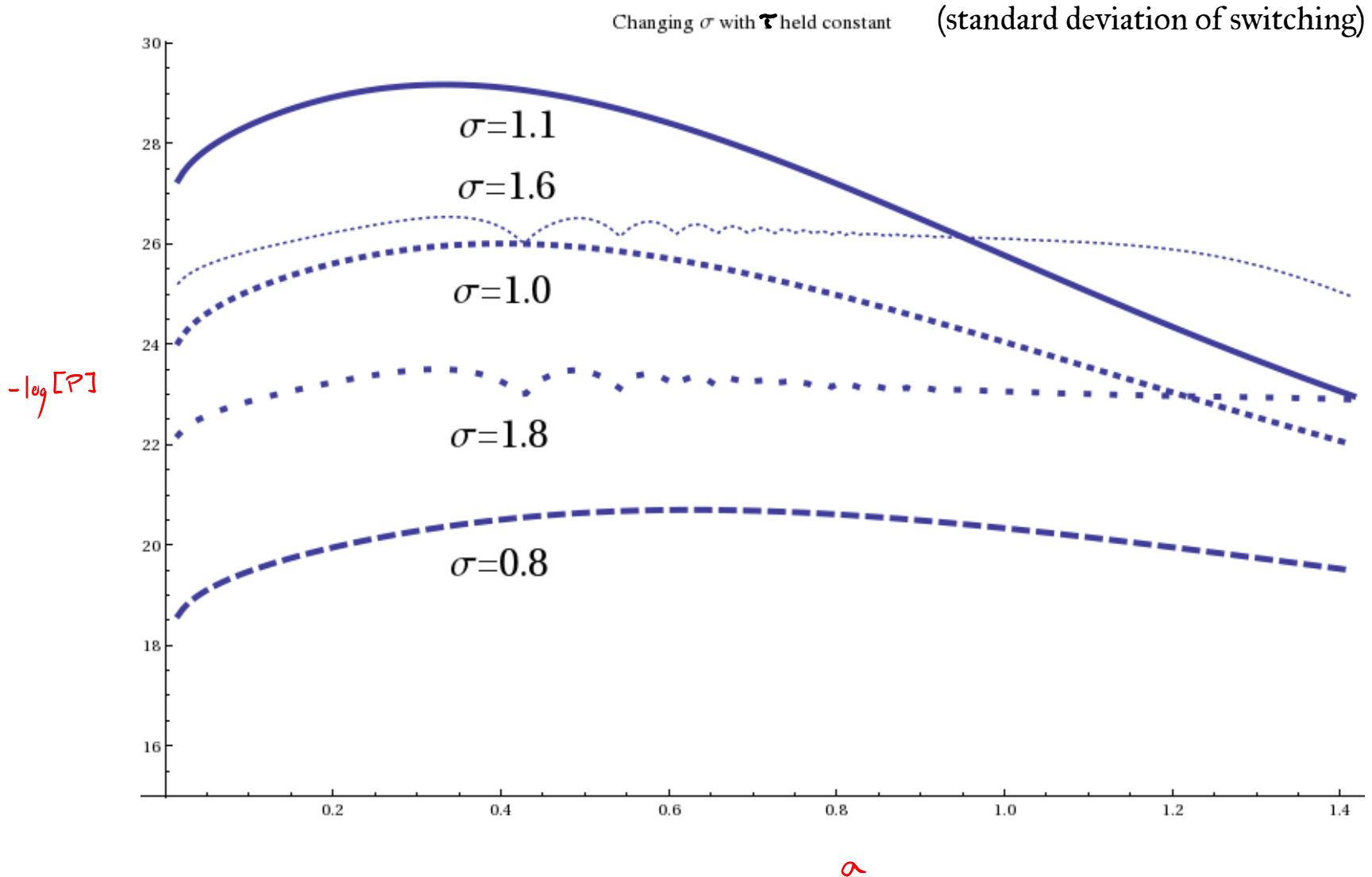
a. Perturbative



2. Temperature and Acceleration Plots

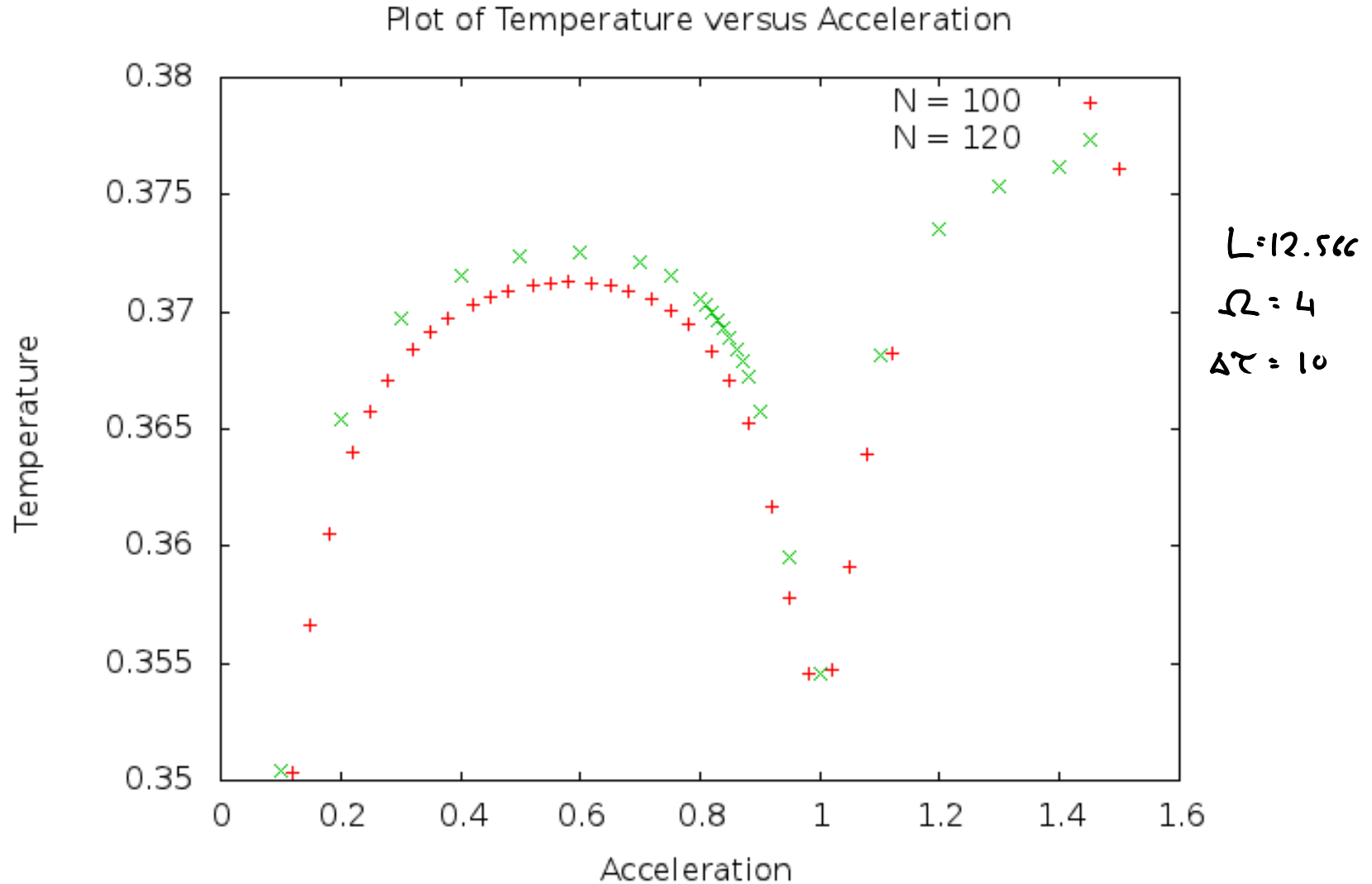
$$L = 4\pi \quad \Omega = 4$$

a. Perturbative



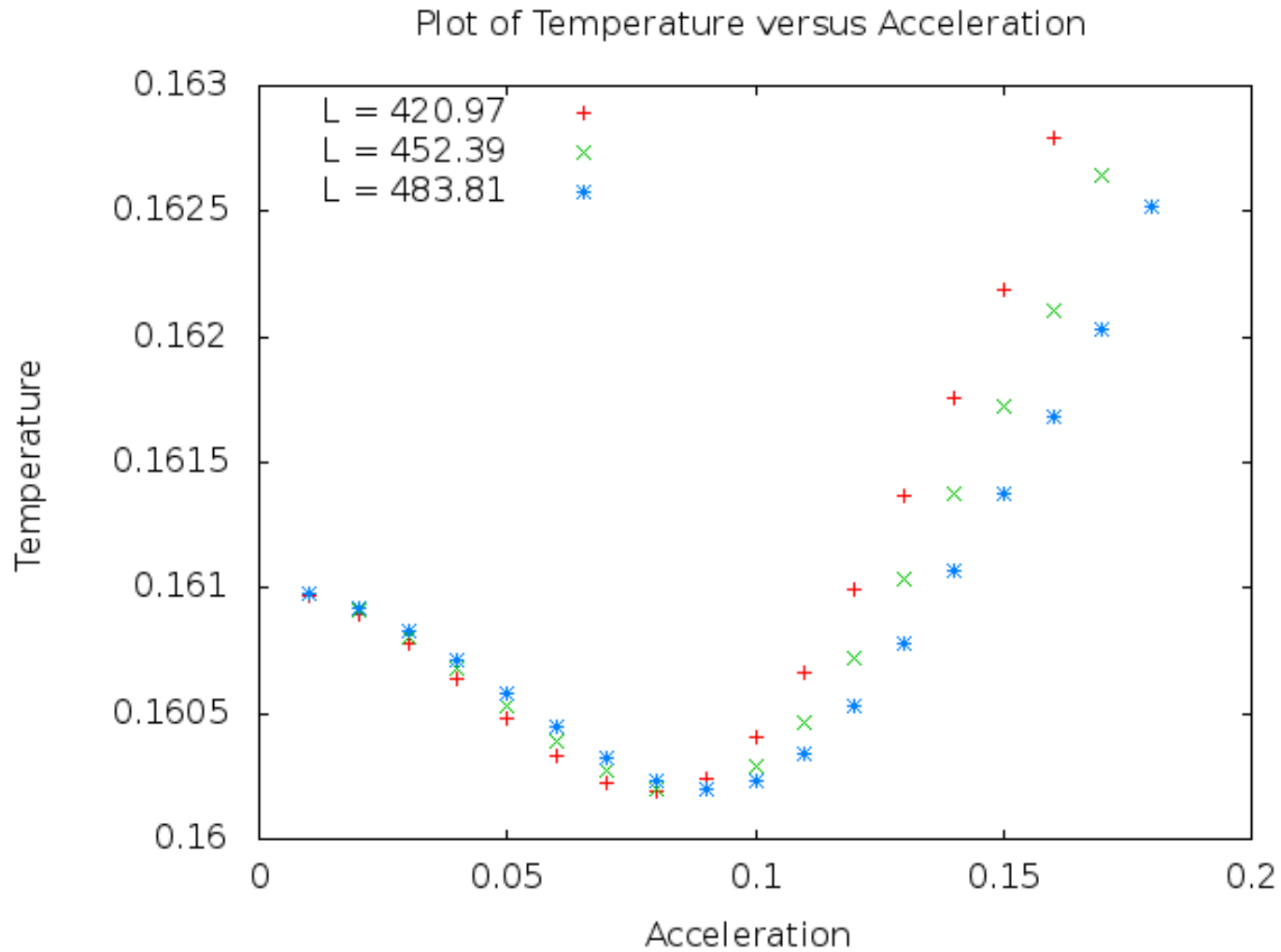
2. Temperature and Acceleration Plots

b. Non-Perturbative (Gaussian Quantum Detector)



2. Temperature and Acceleration Plots

b. Non-Perturbative



Conclusion

Taking the cavity to infinity will remain difficult.
Does a detector get hot when you accelerate it?

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Not always! Detectors can cool down as we accelerate them.

The negative linear temperature vs acceleration behaviour?

It's a non-equilibrium effect!

Appears to be due to a finitely sized cavity/trajectory!

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Taking the cavity to infinity will remain difficult.
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Anti-Unruh...?

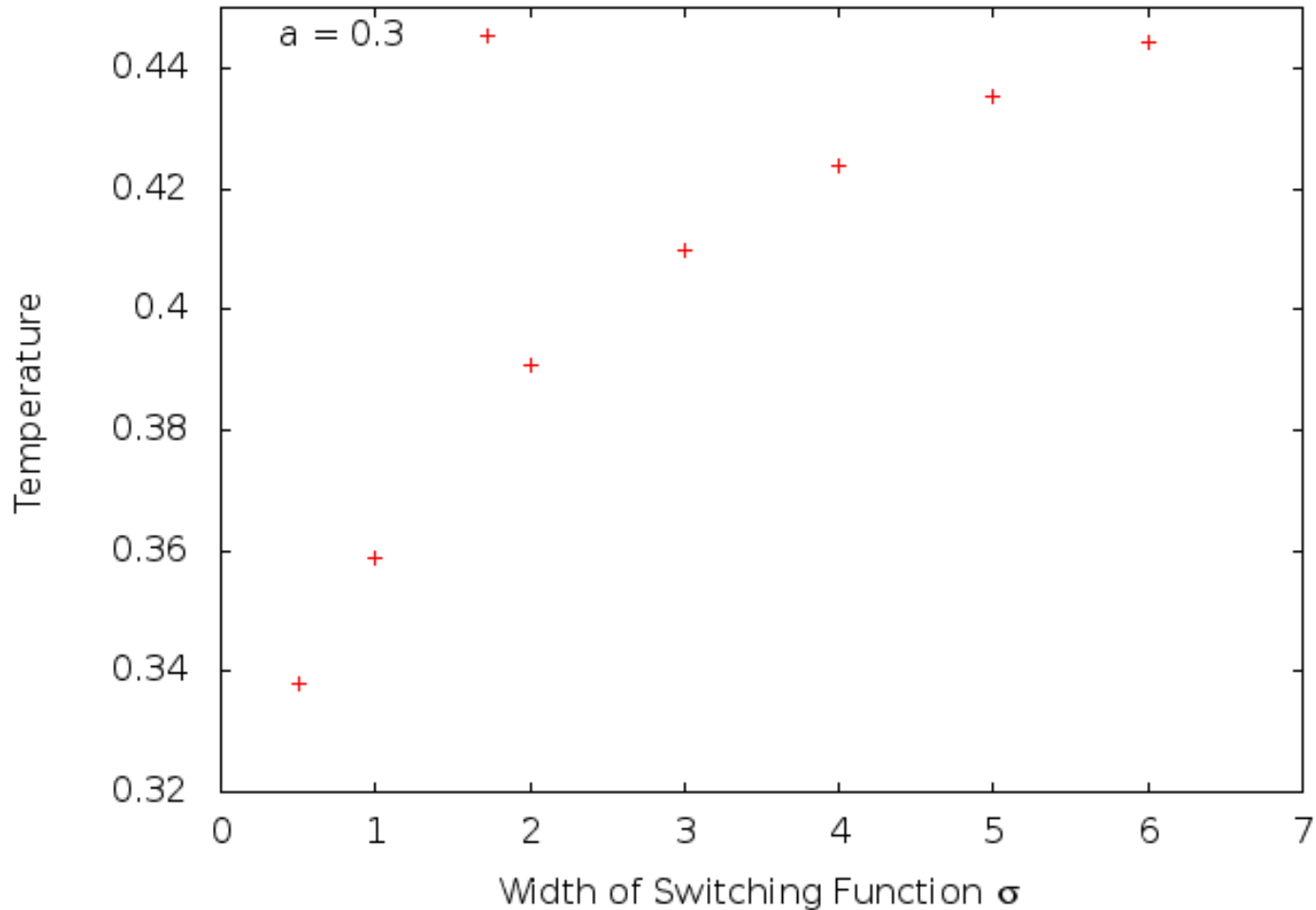
Questions?

2. Temperature and Acceleration Plots

b. Non-Perturbative

ΔT increases too, to
maintain $\Delta T = 12\sigma$

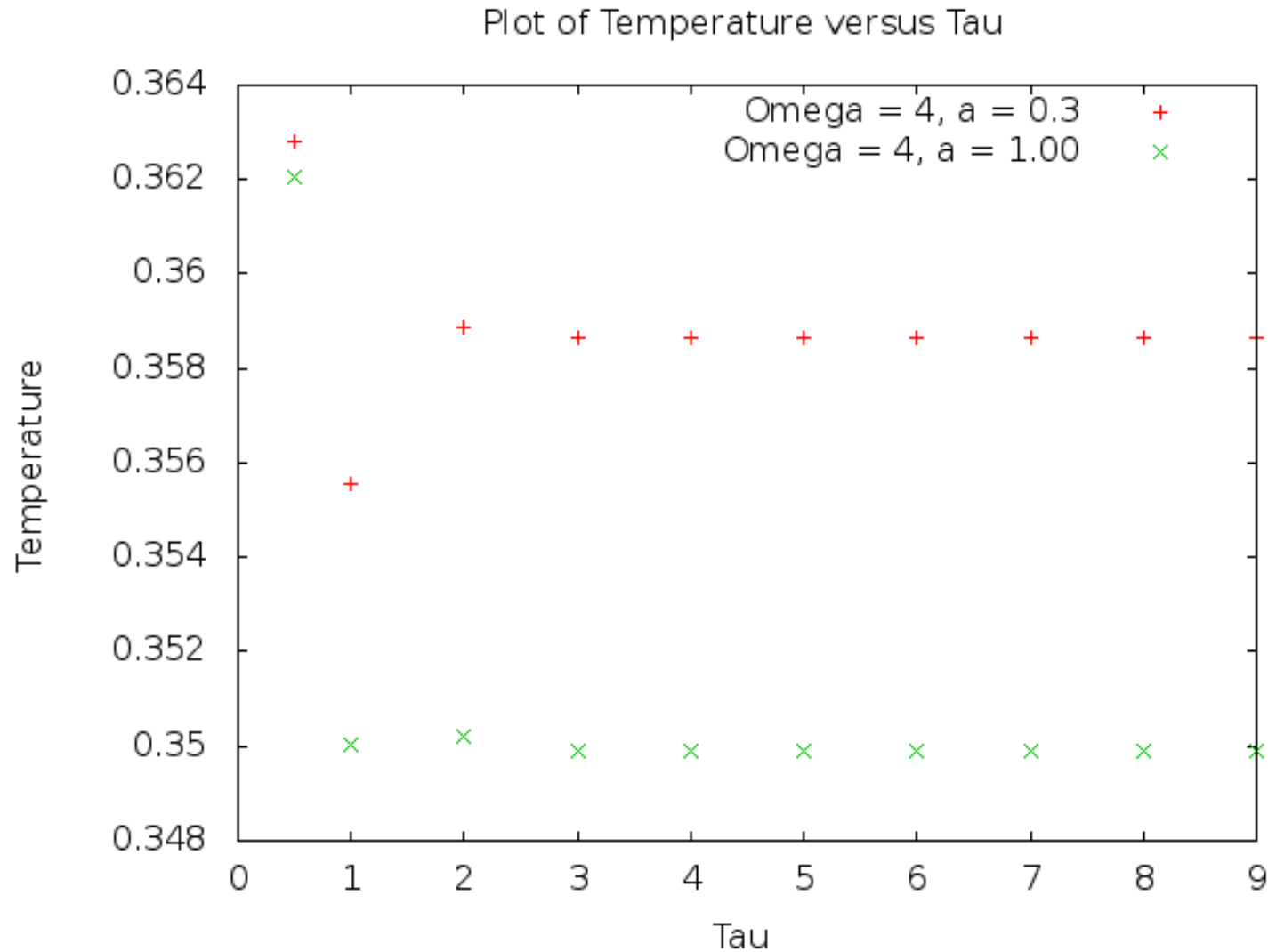
Plot of Temperature versus Width of Switching Function



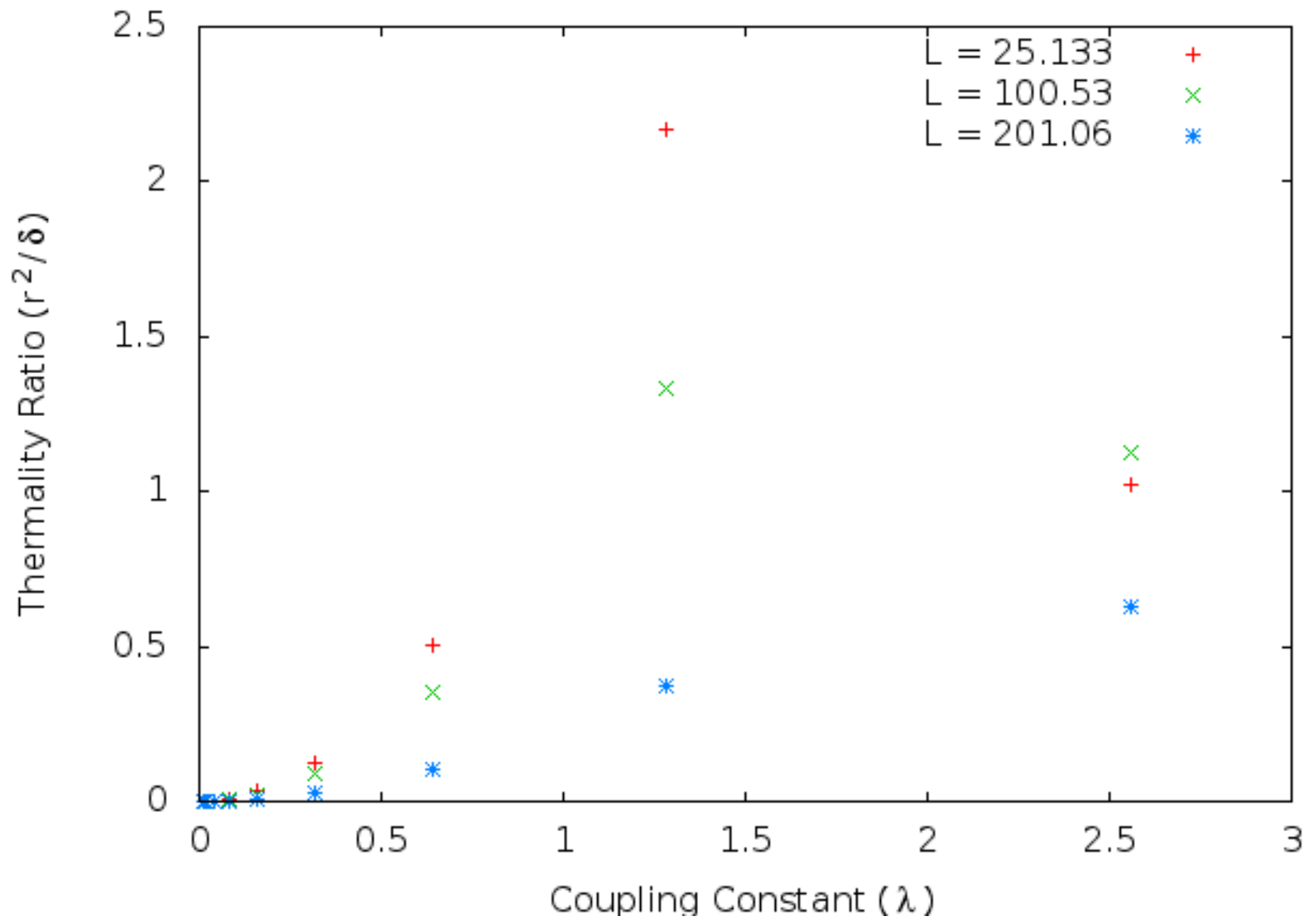
2. Temperature and Acceleration Plots

b. Non-Perturbative

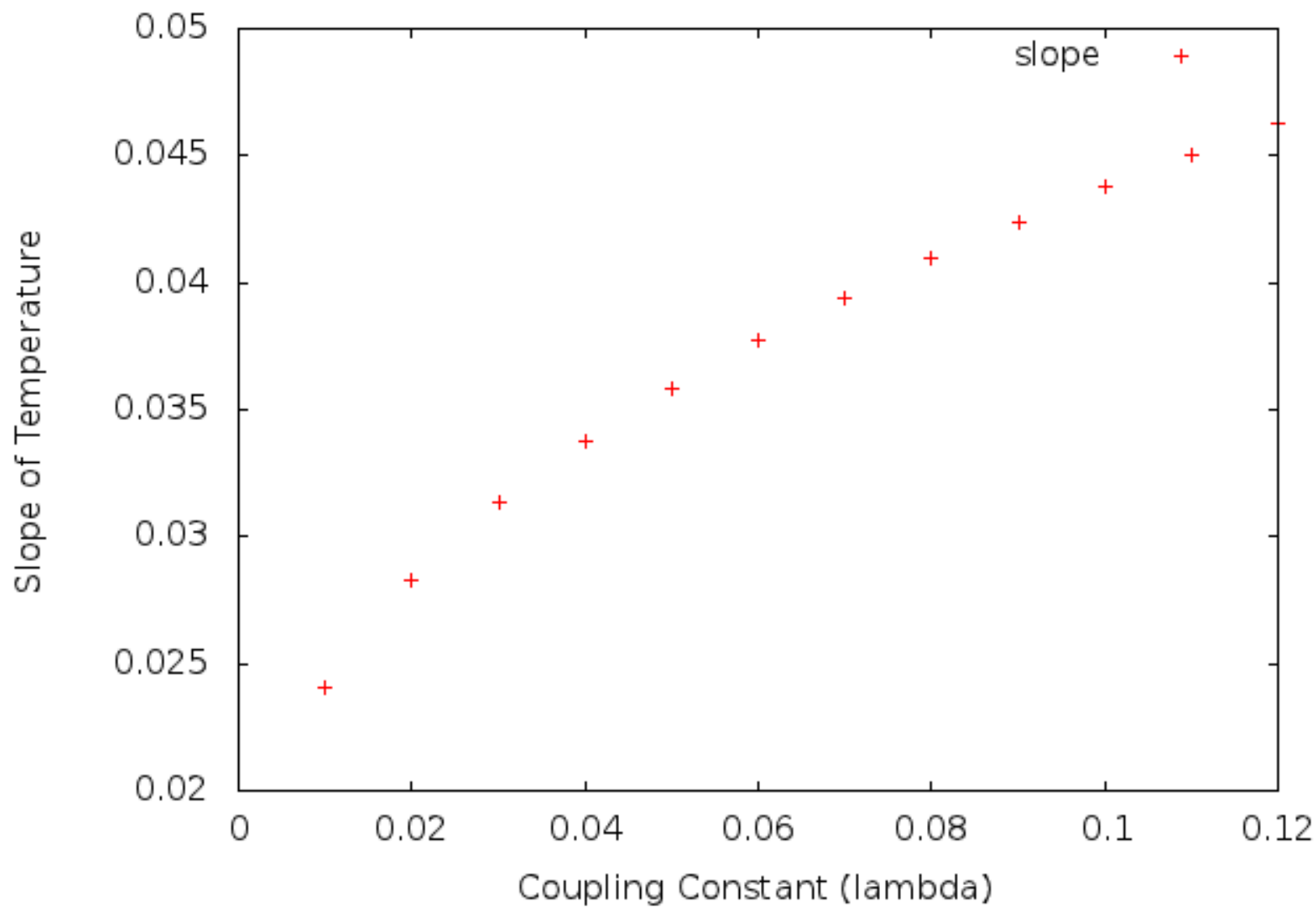
$\Delta T = [-\tau, \tau]$



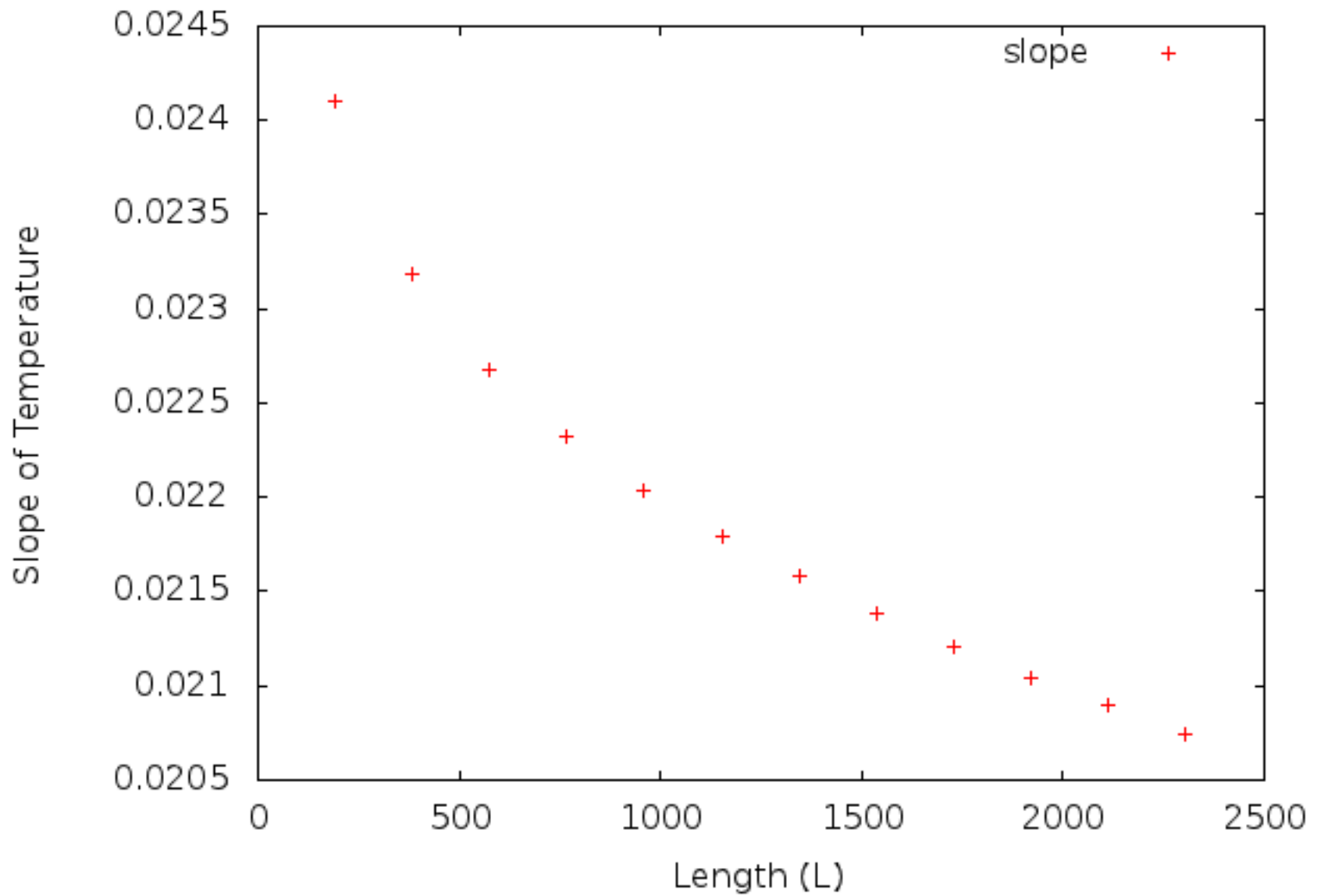
Thermality versus Coupling Constant (λ)



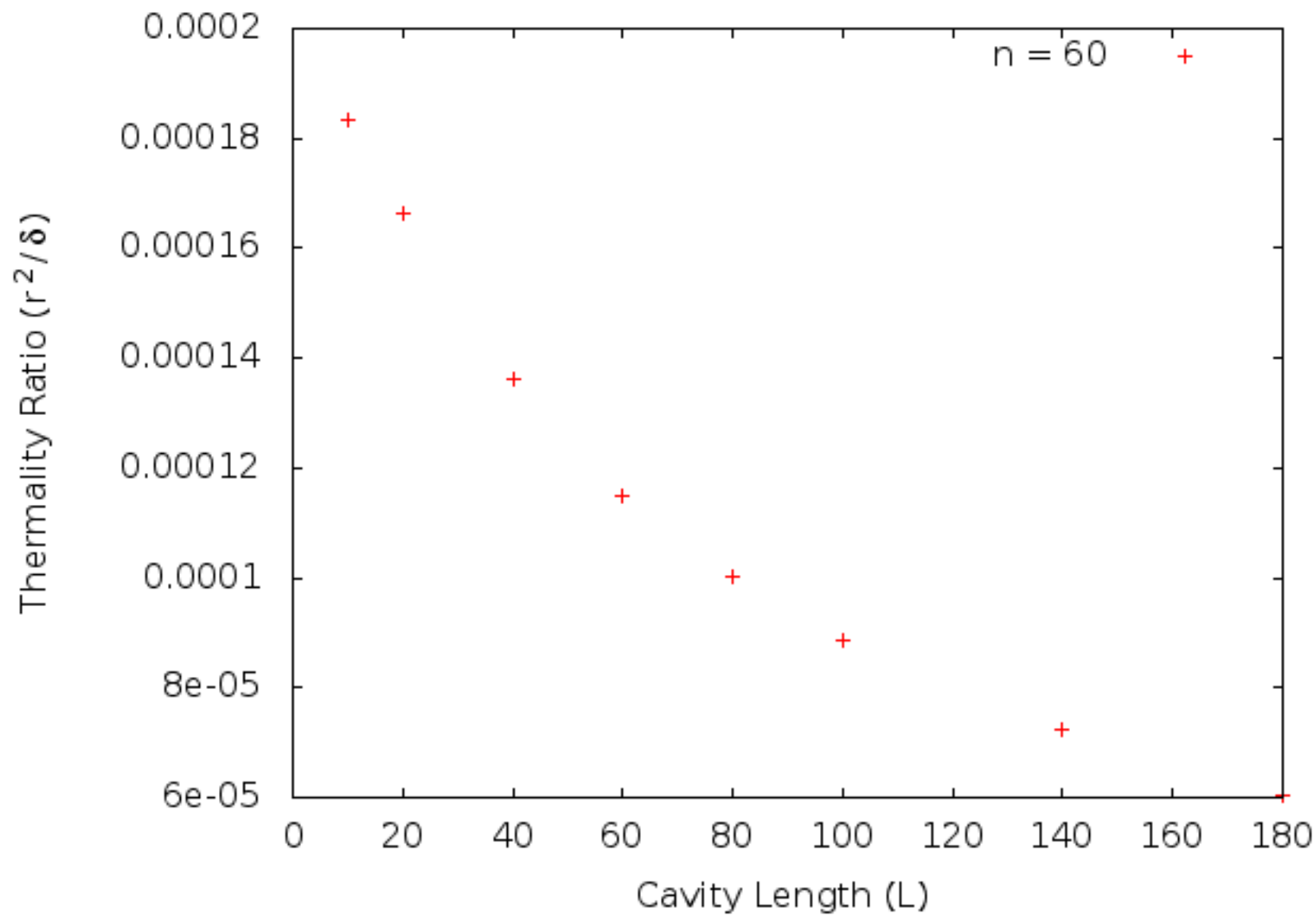
Increasing Coupling



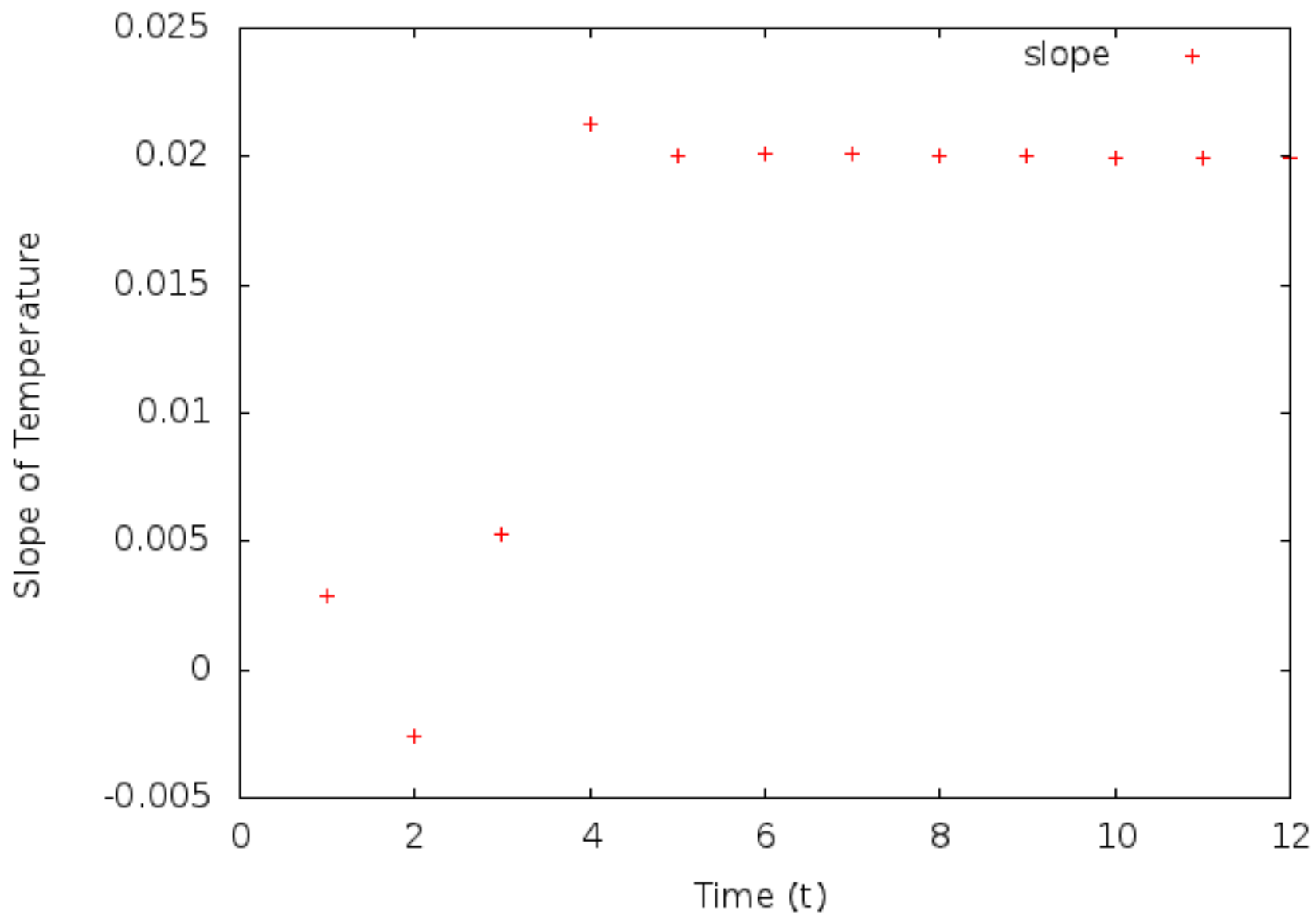
Increasing Length



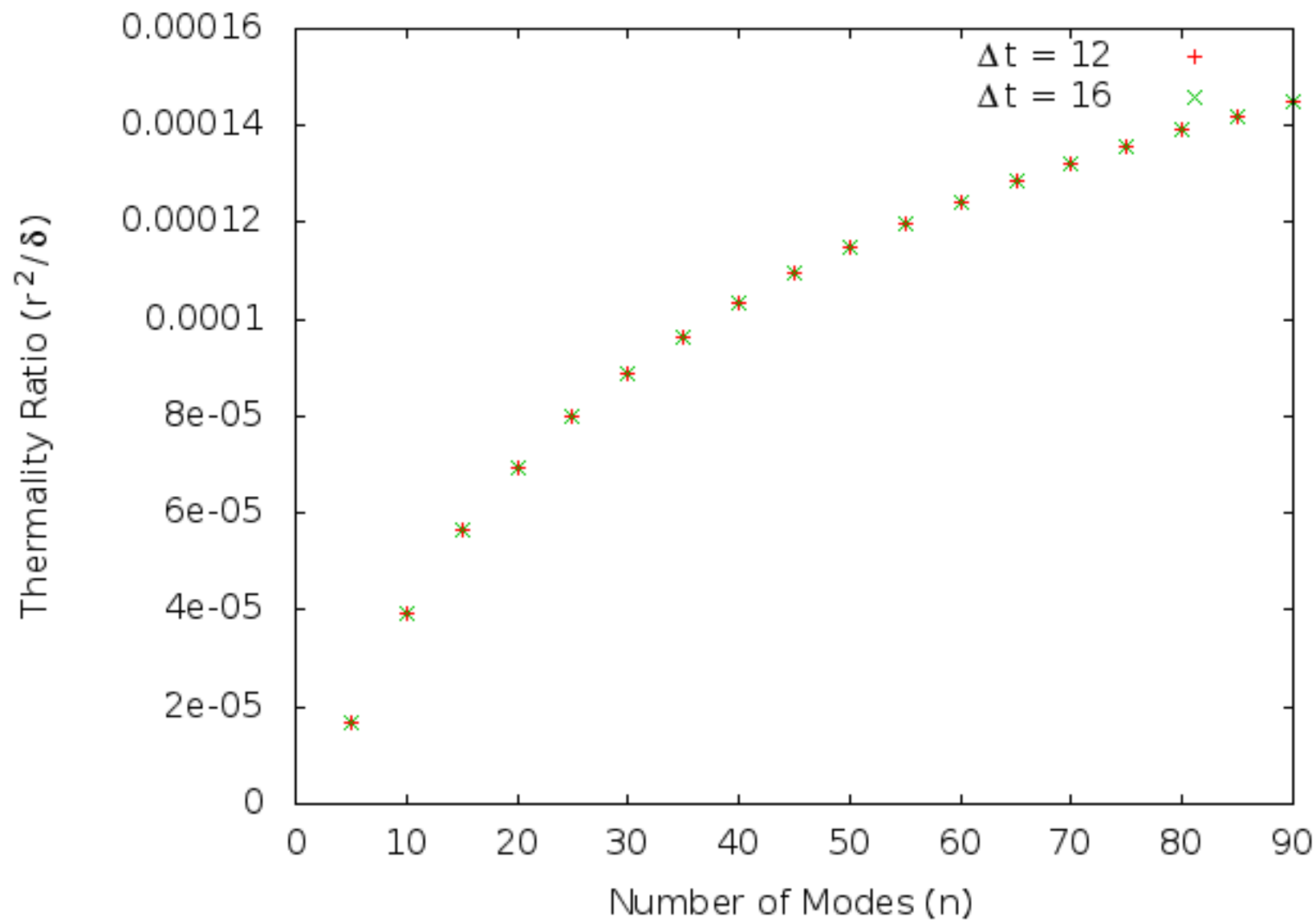
Thermality versus Cavity Length



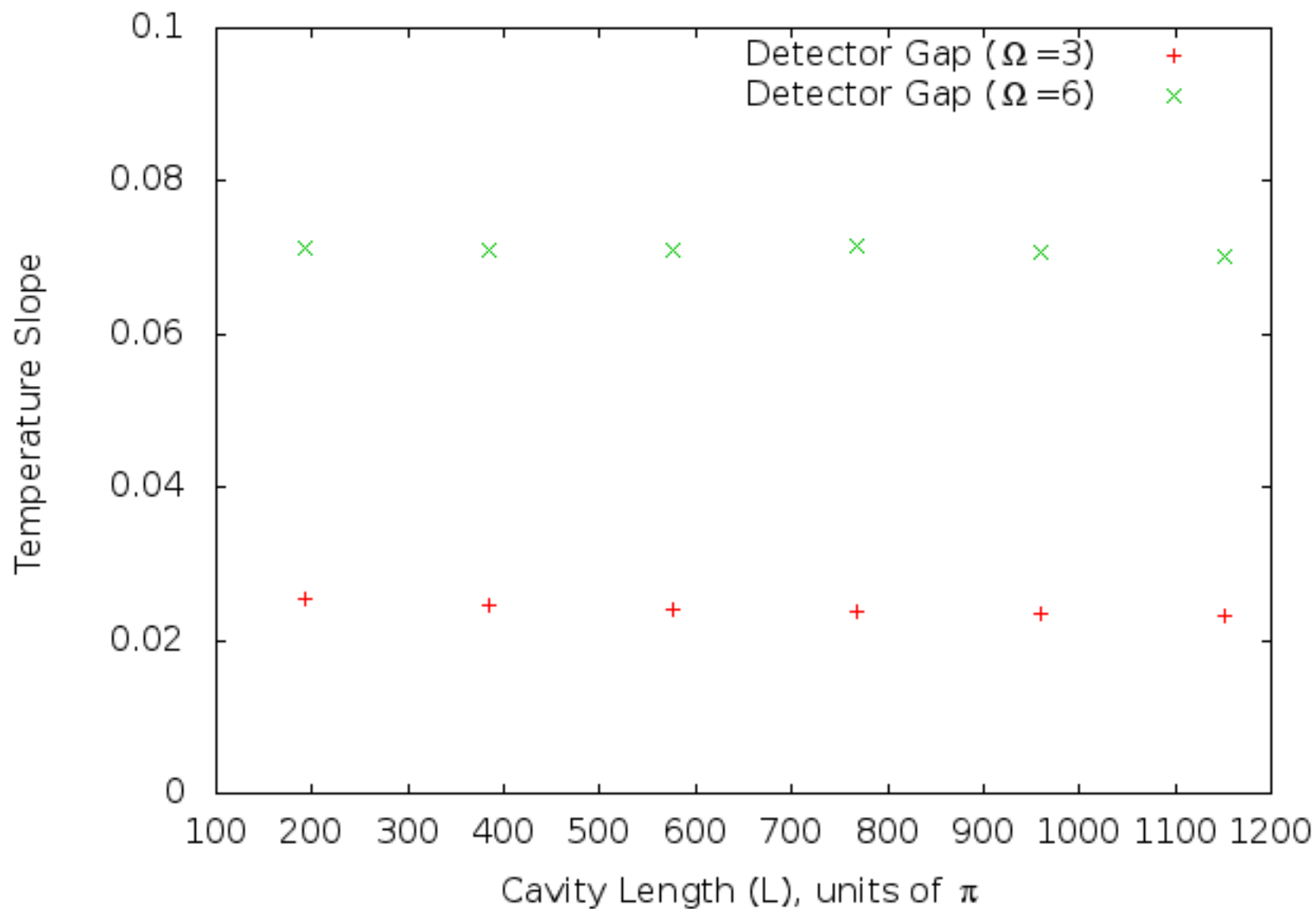
Thermalization (L = 96 pi)



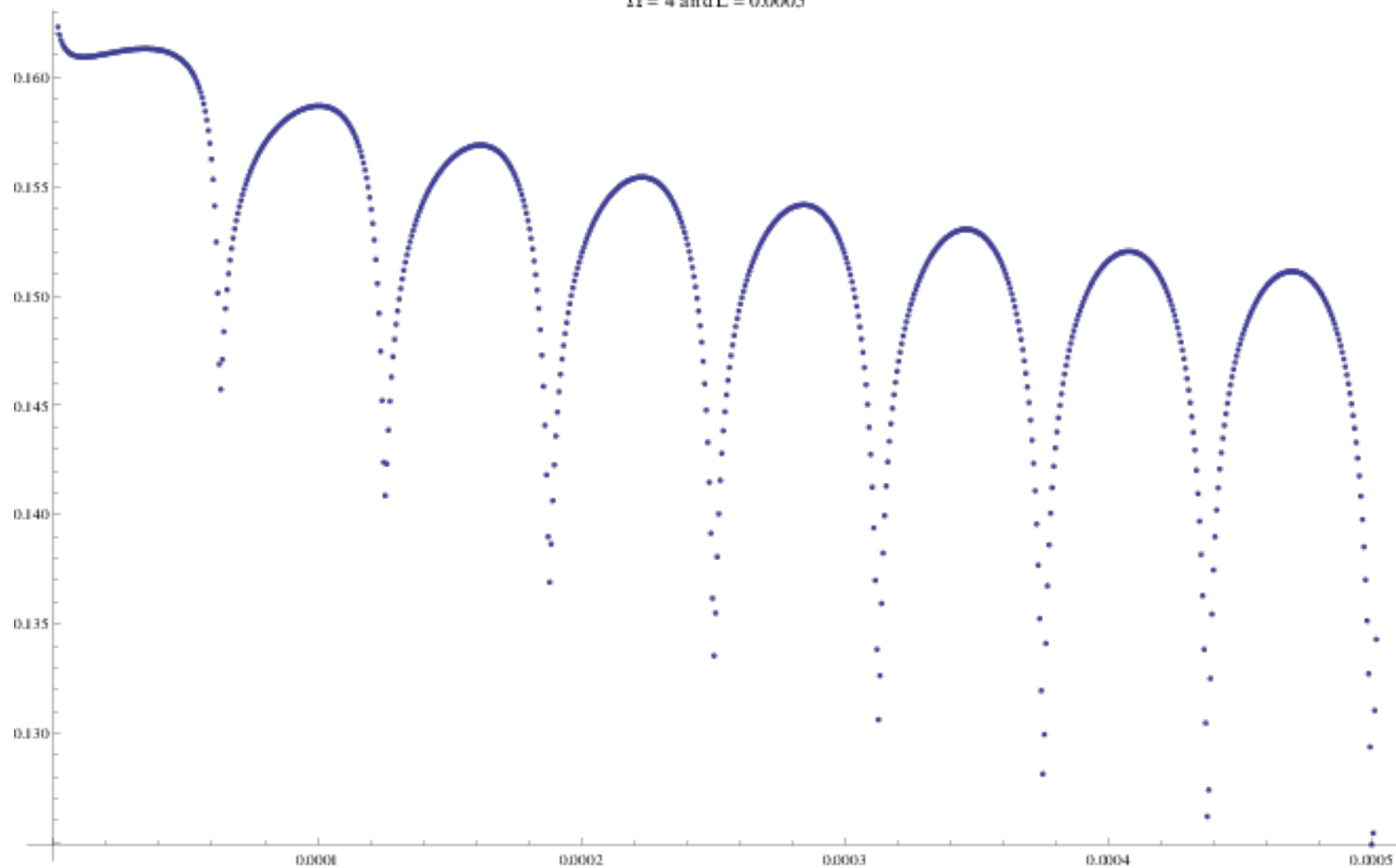
Thermality versus Number of Modes



Temperature Slope vs Cavity Length (L)



$\Omega = 4$ and $L = 0.0005$



Thermality versus Detector Gap (Ω)

