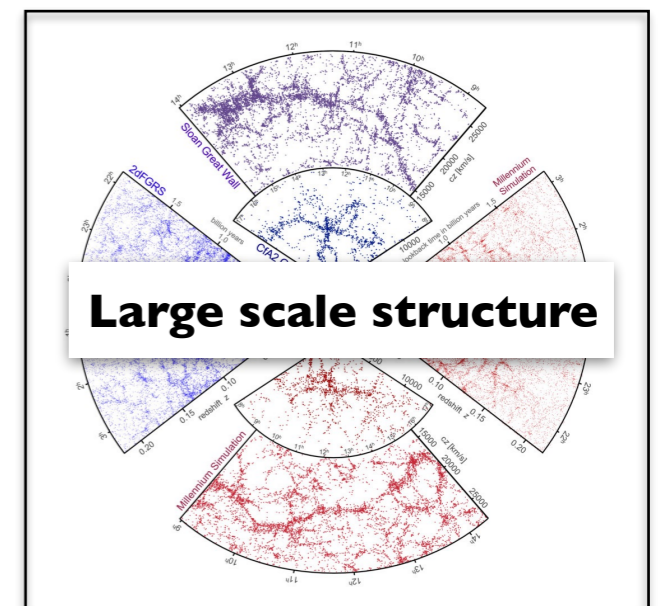
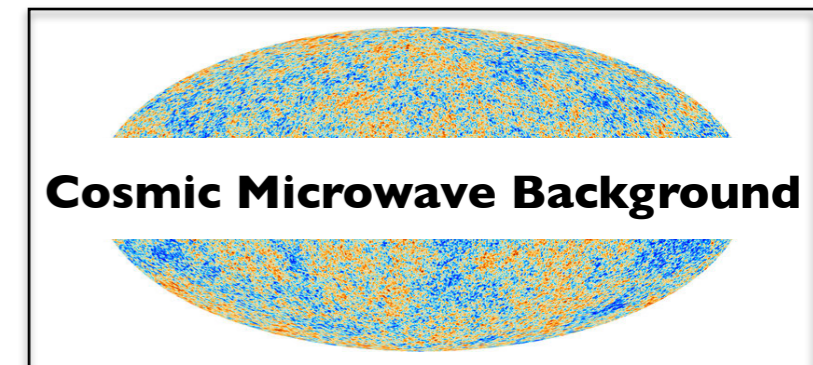
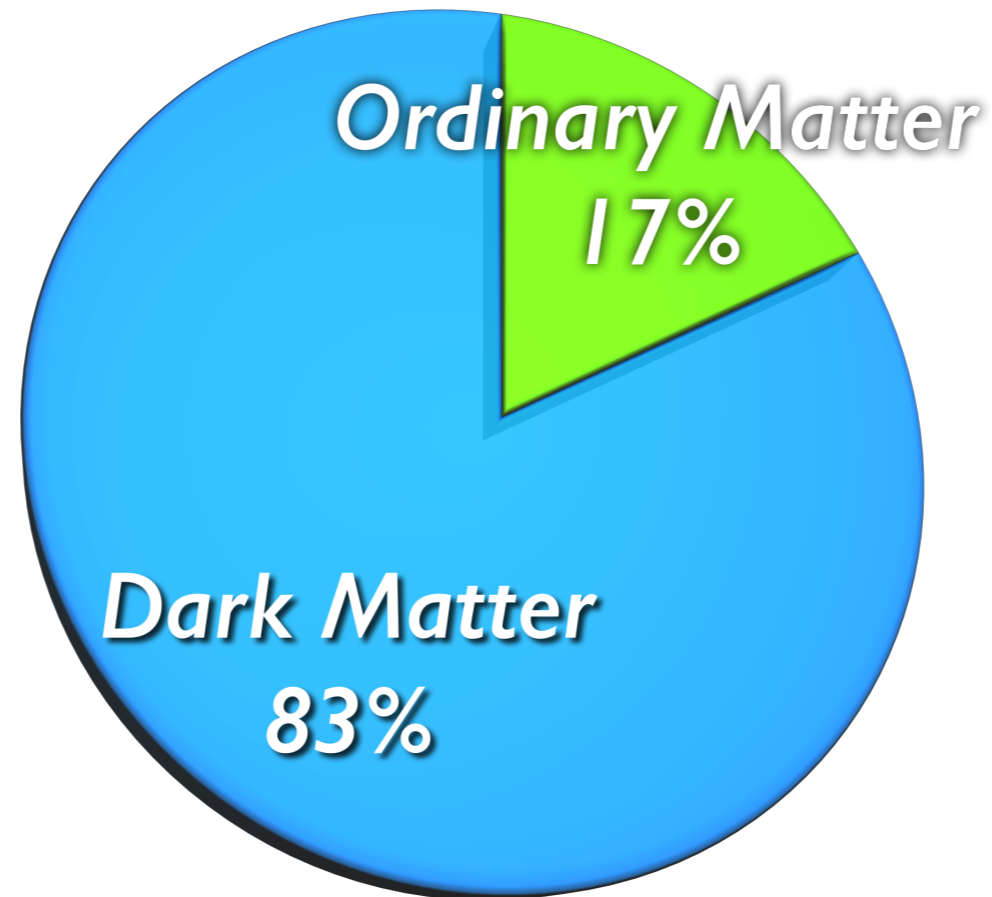
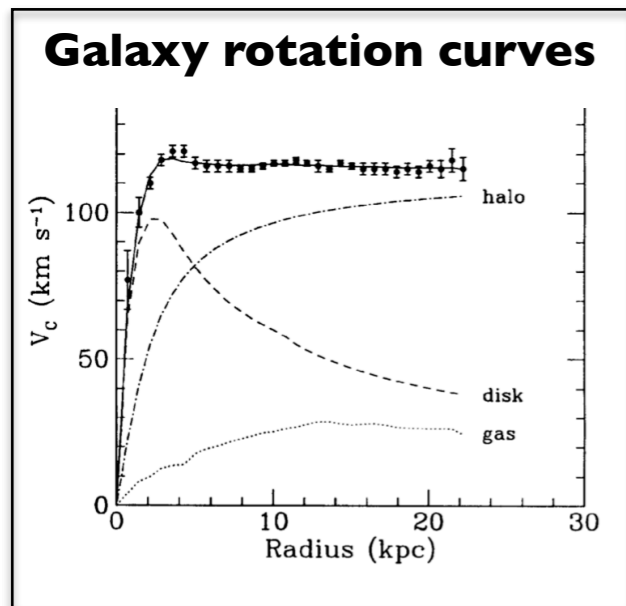


Dark Matter*

Christopher McCabe

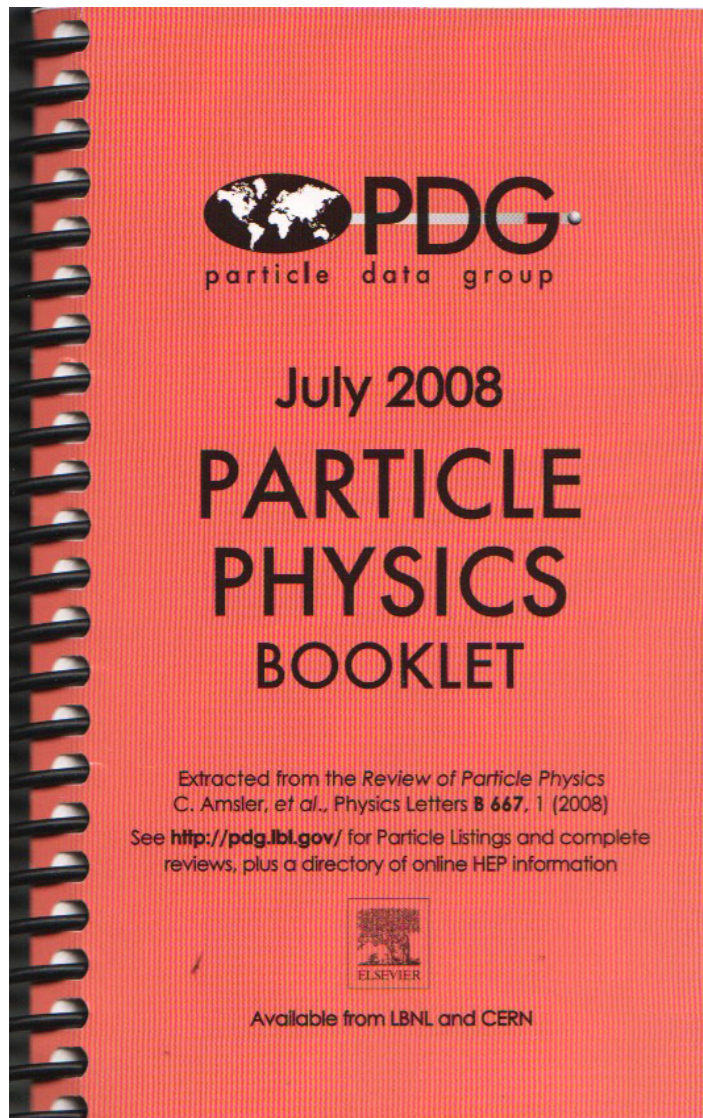
Evidence for dark matter

Matter in the Universe



Evidence from gravitational interactions over many distance scales

What would we like to know?



Dark Matter Particle (X^0)

X^0 mass: $m = ?$

X^0 spin: $J = ?$

X^0 parity: $P = ?$

X^0 lifetime: $\tau = ?$

X^0 scattering cross-section on nucleons: ?

X^0 production cross-section in hadron colliders: ?

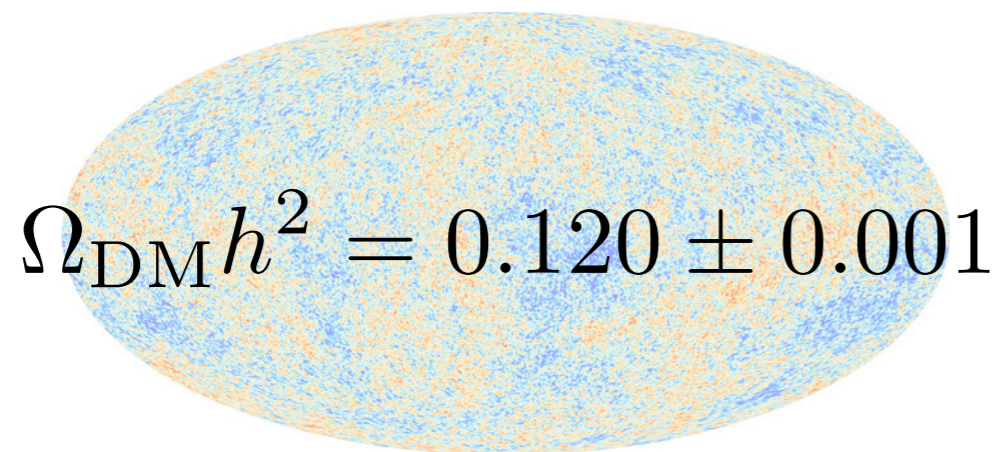
X^0 self-annihilation cross-section: ?

Why should DM interact with the SM?

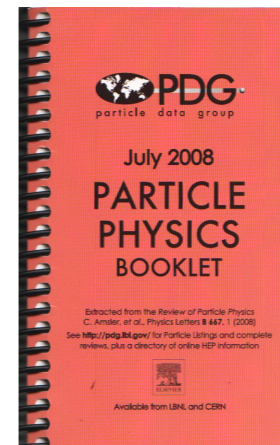
“Up to a point the stories of cosmology and particle physics can be told separately. In the end though, they will come together.”

Steven Weinberg

Cosmology



Particle Physics



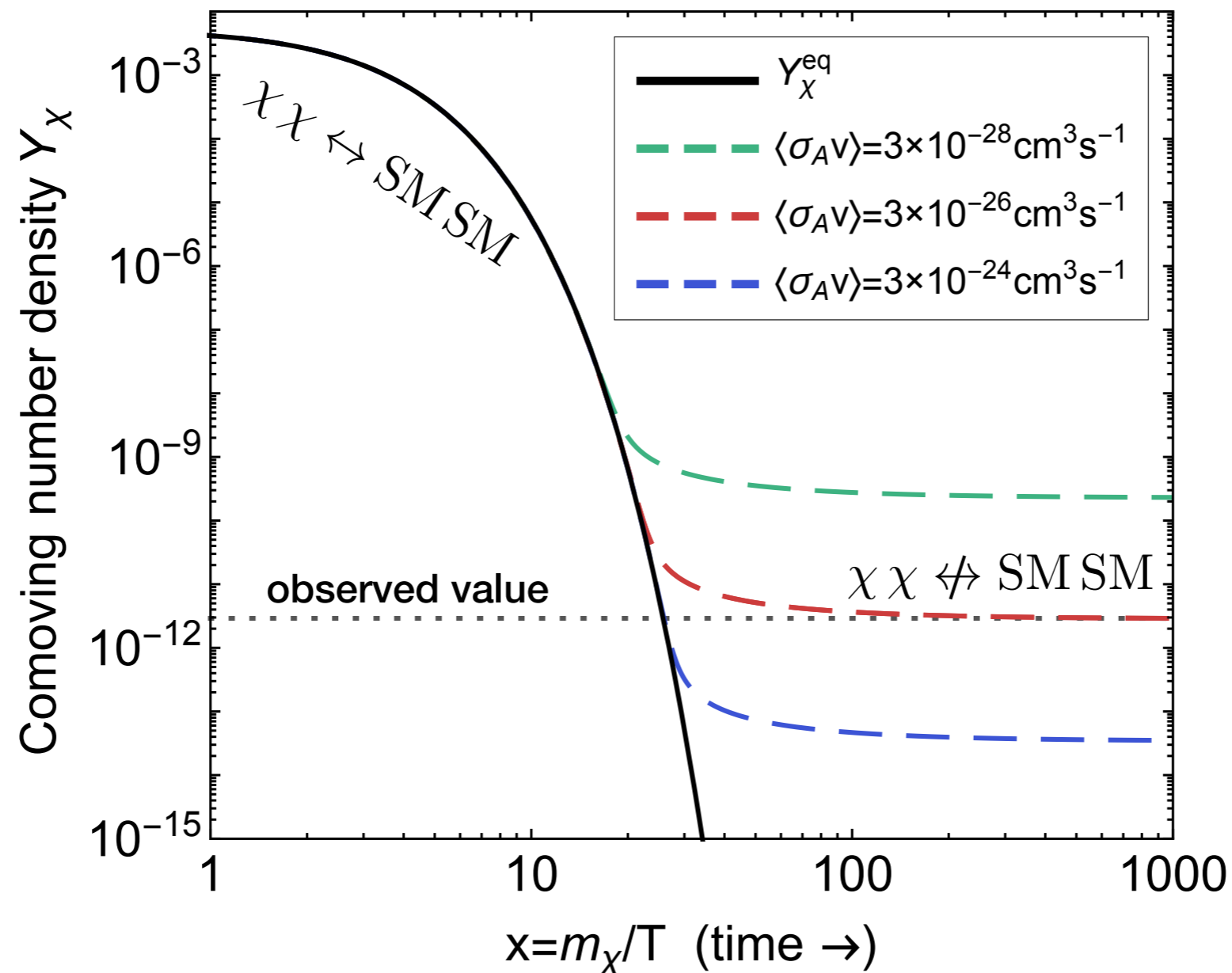
$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{m_q}{\Lambda^3} \bar{\chi} \chi \bar{q} q + \dots$$

Suggests DM - Standard Model interactions are generic
&

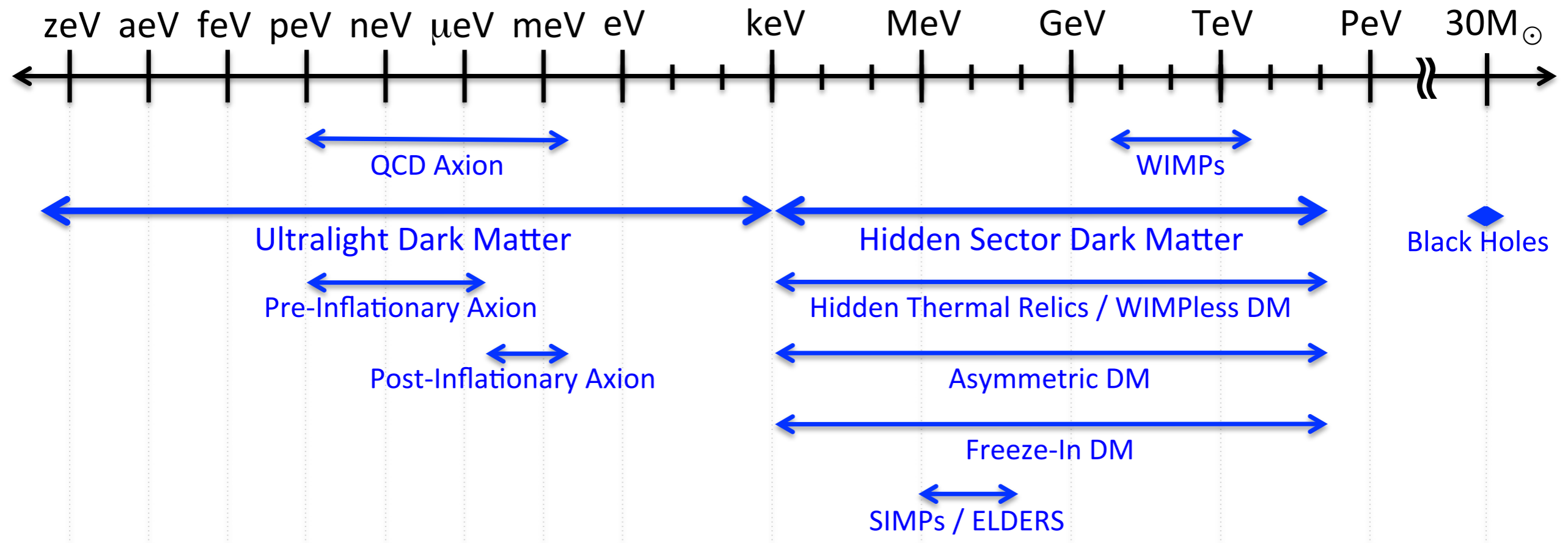
informs and limits the possible interactions

WIMPs: canonical example

Thermal freeze-out mechanism gives observed abundance



Theorists haven't stopped at WIMPs...

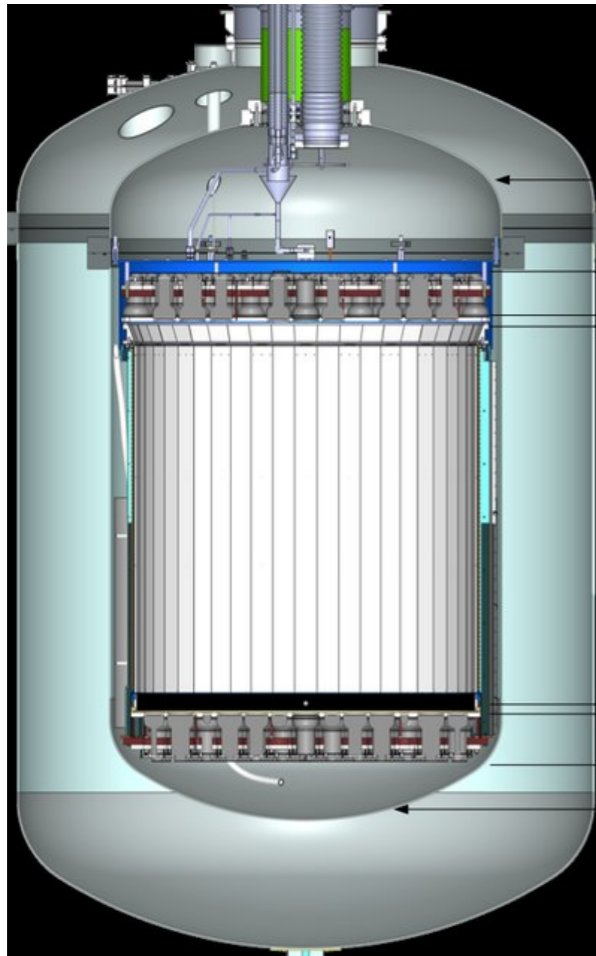


US Cosmic Visions

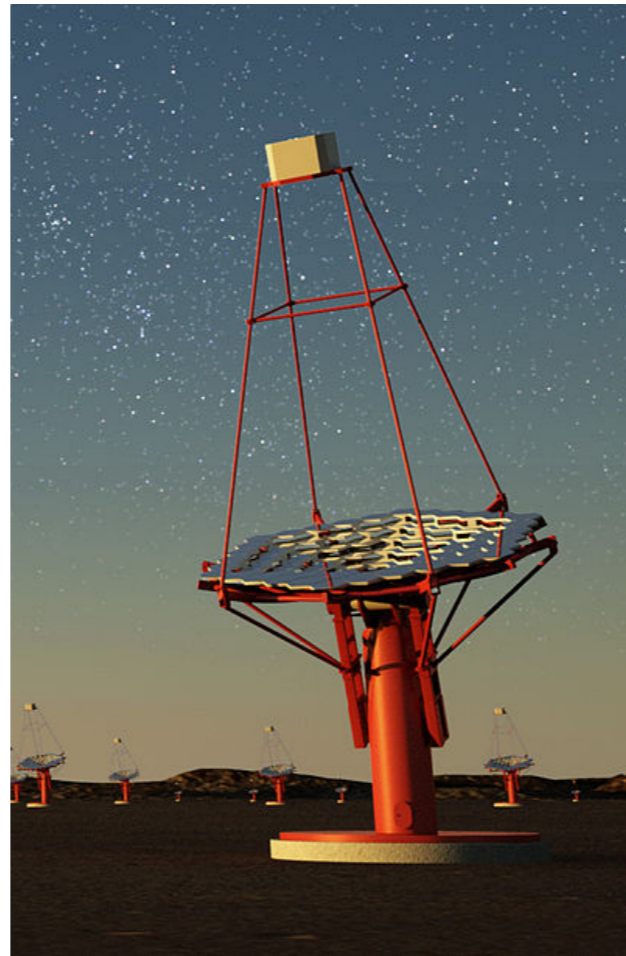
Many candidates outside the WIMP mass range all with SM interactions

Searching for DM - SM interactions

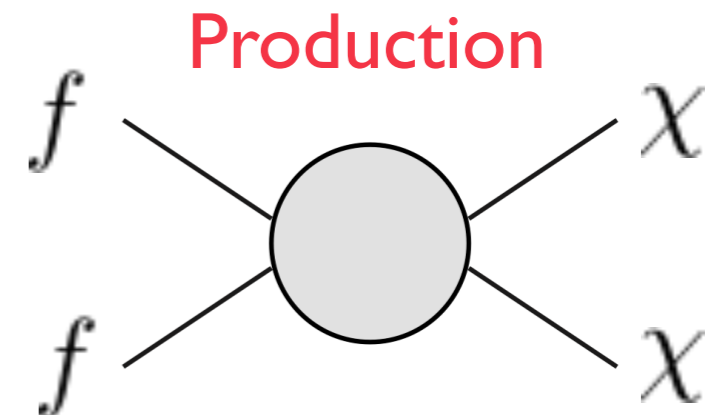
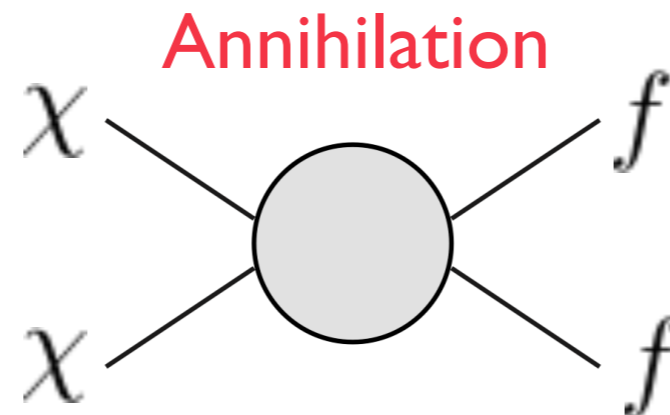
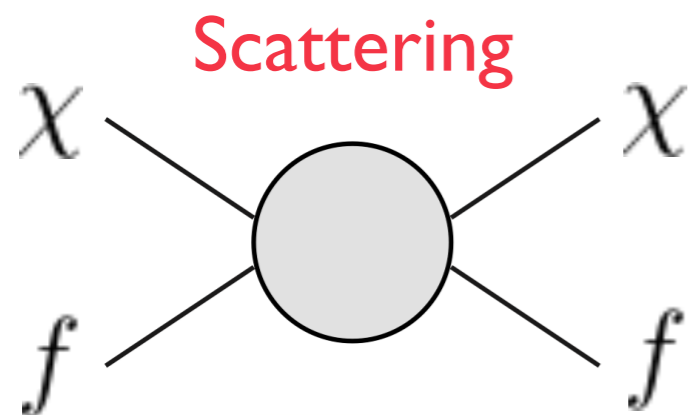
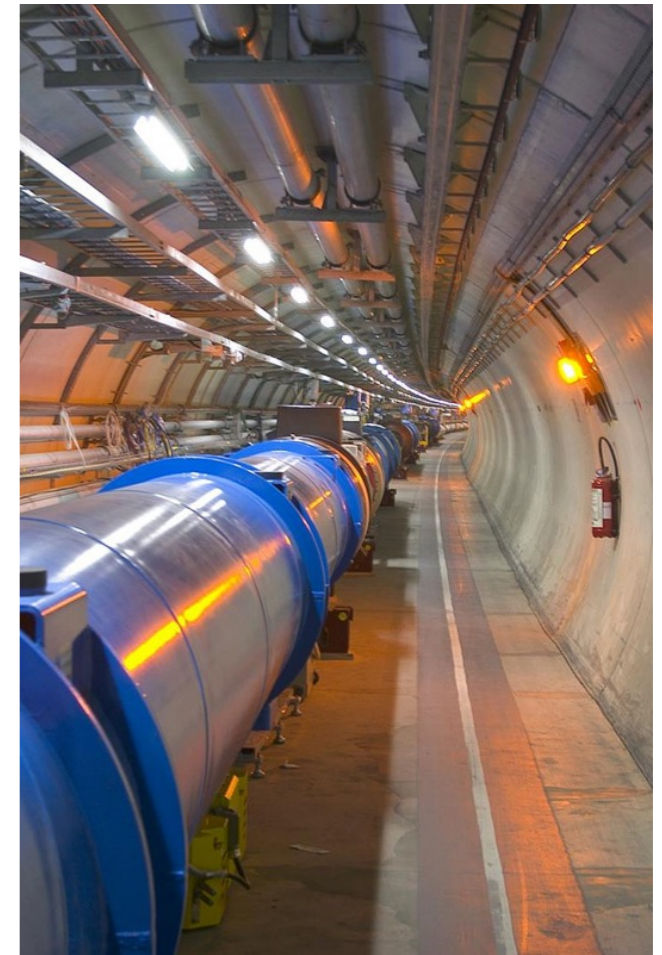
Direct detection



Indirect detection



Collider

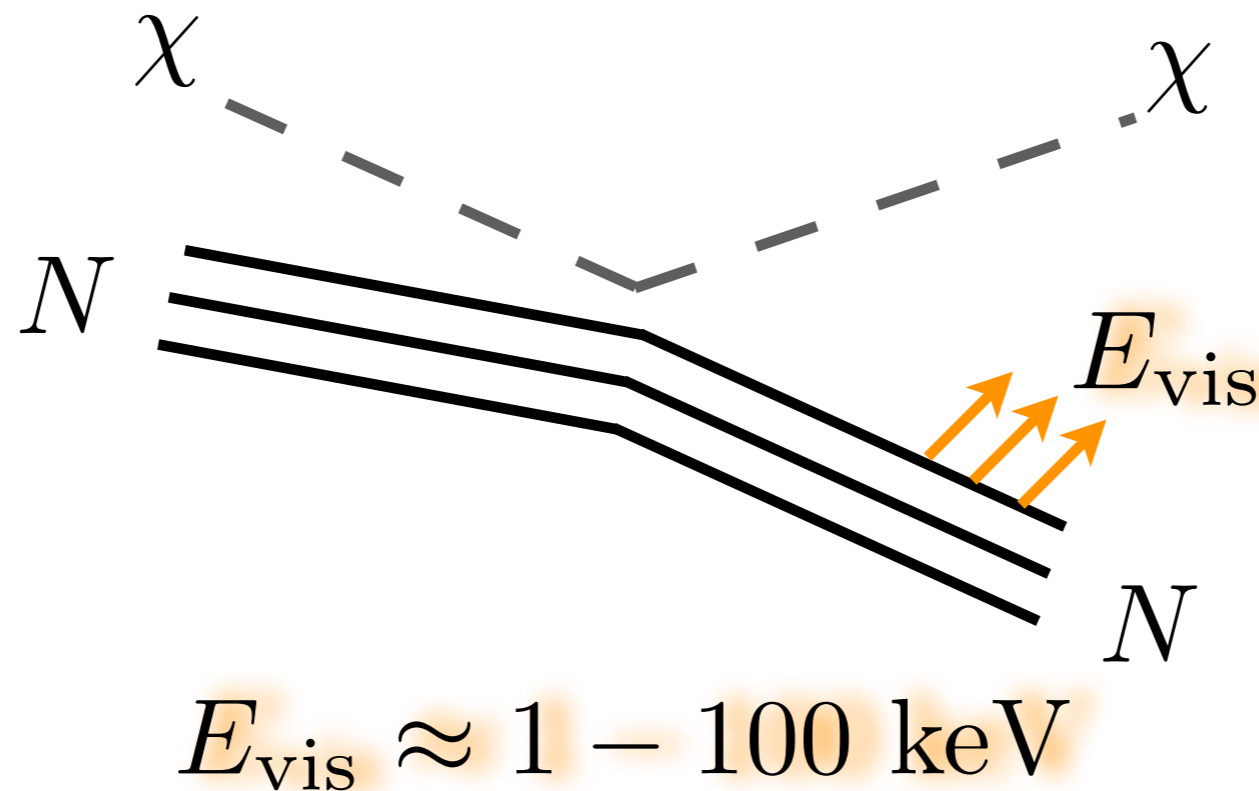


Direct detection

Basics of direct detection

Original aim: detect collisions of dark matter with a nucleus

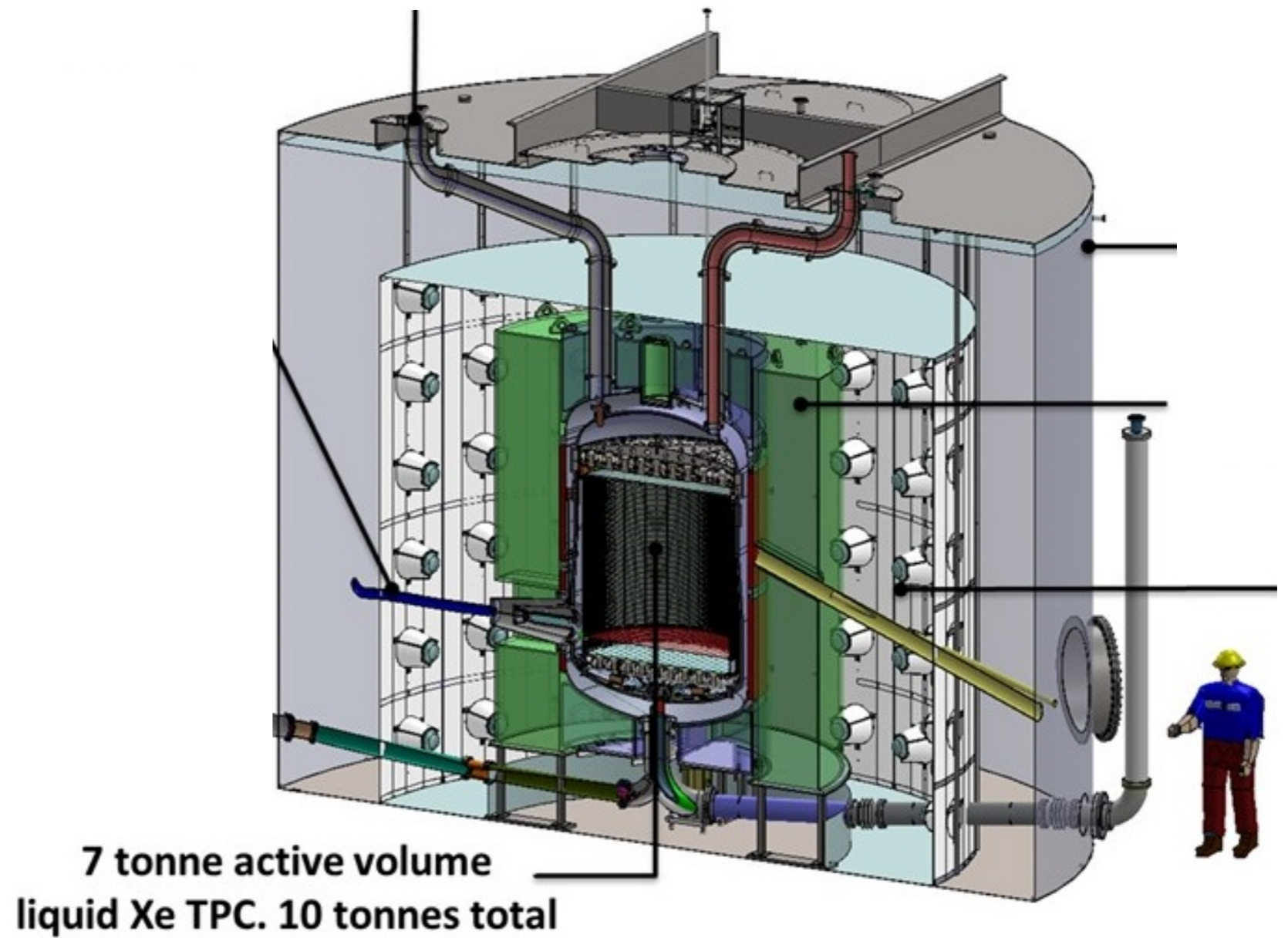
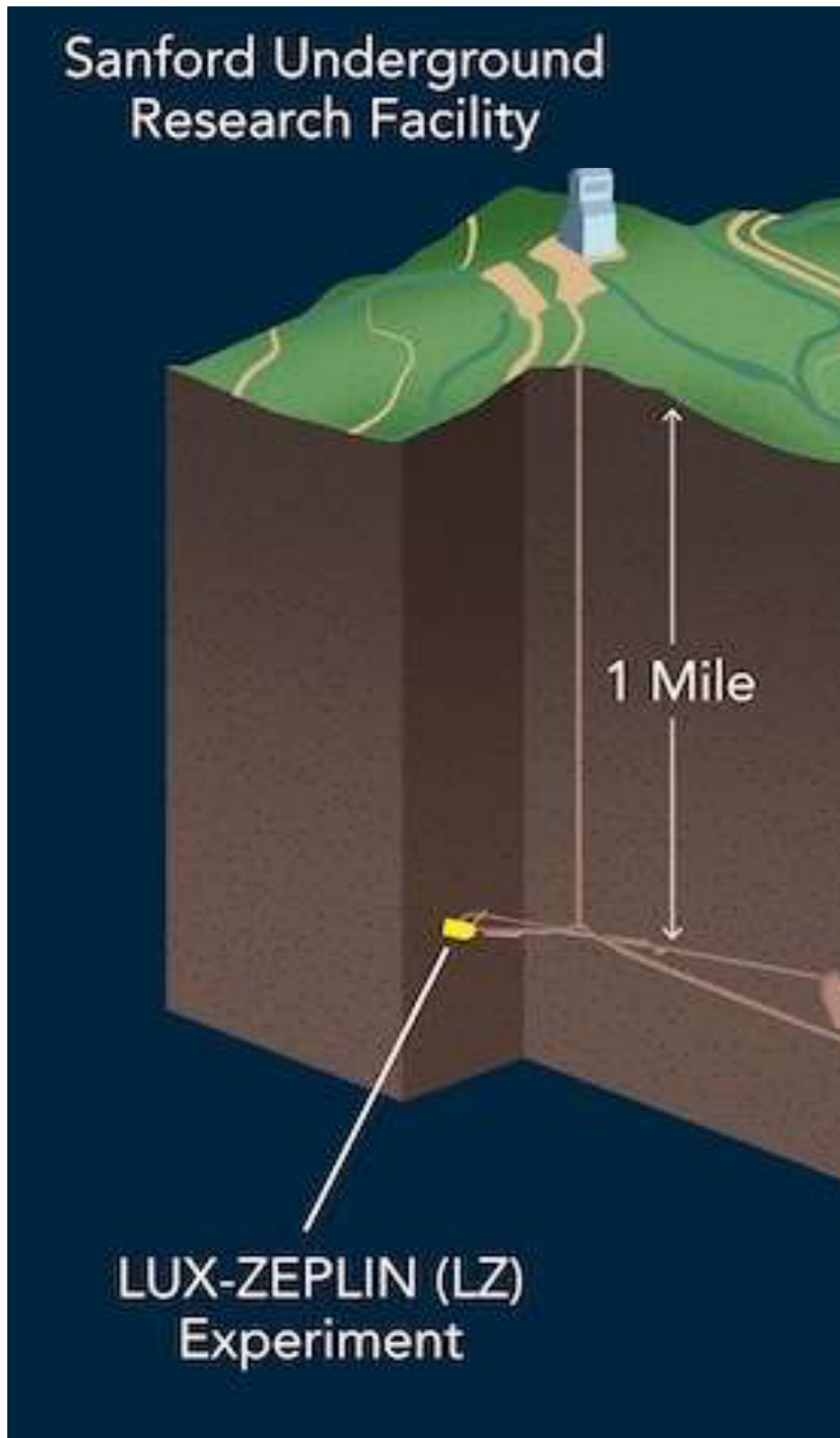
Goodman & Witten (1985)



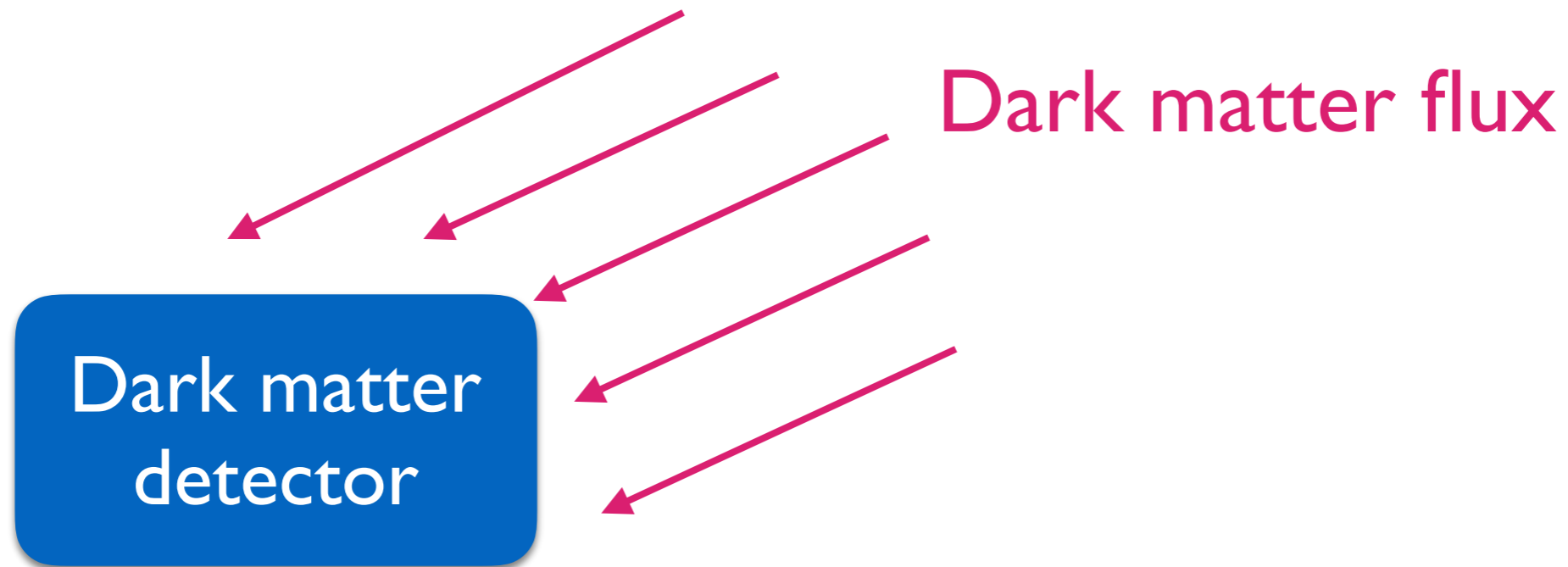
Event rate: few events / year

Ultra-sensitive keV energy detectors

Example: the LZ detector



Basics of direct detection

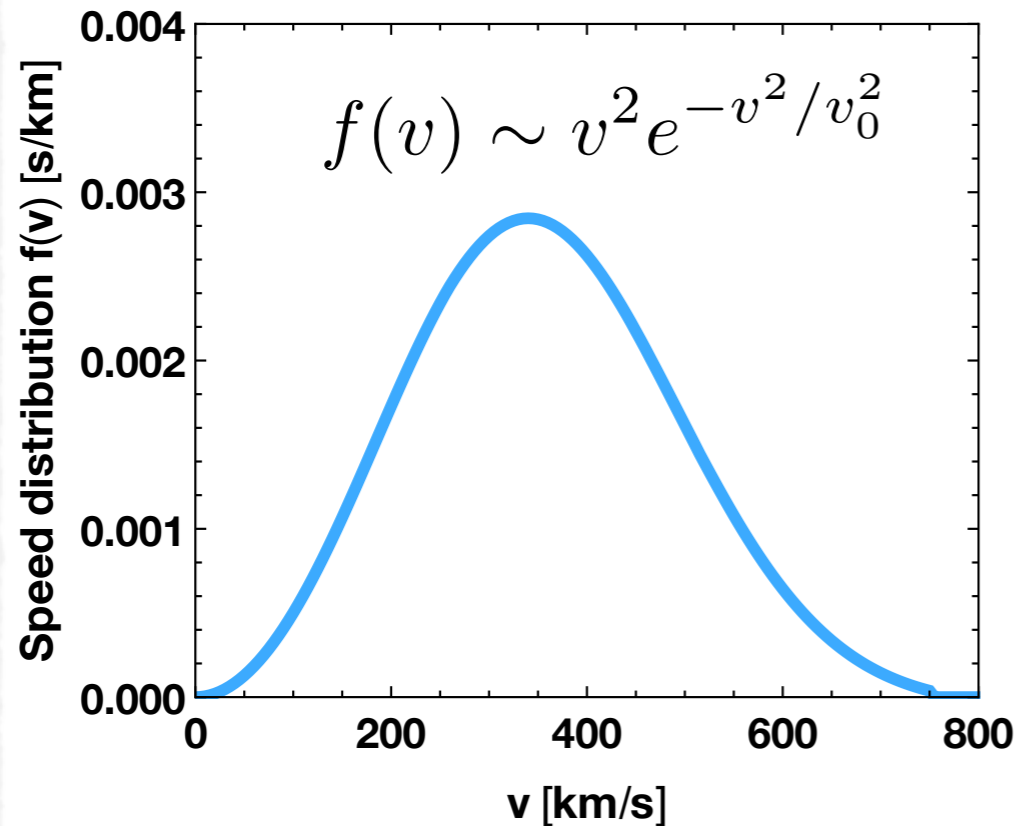


$$\text{Event rate} = \text{DM flux} \times \text{particle physics}$$

Need to model the DM flux to extract the particle physics

SHM: simplest DM flux model

Need to specify velocity distribution and local density to get flux



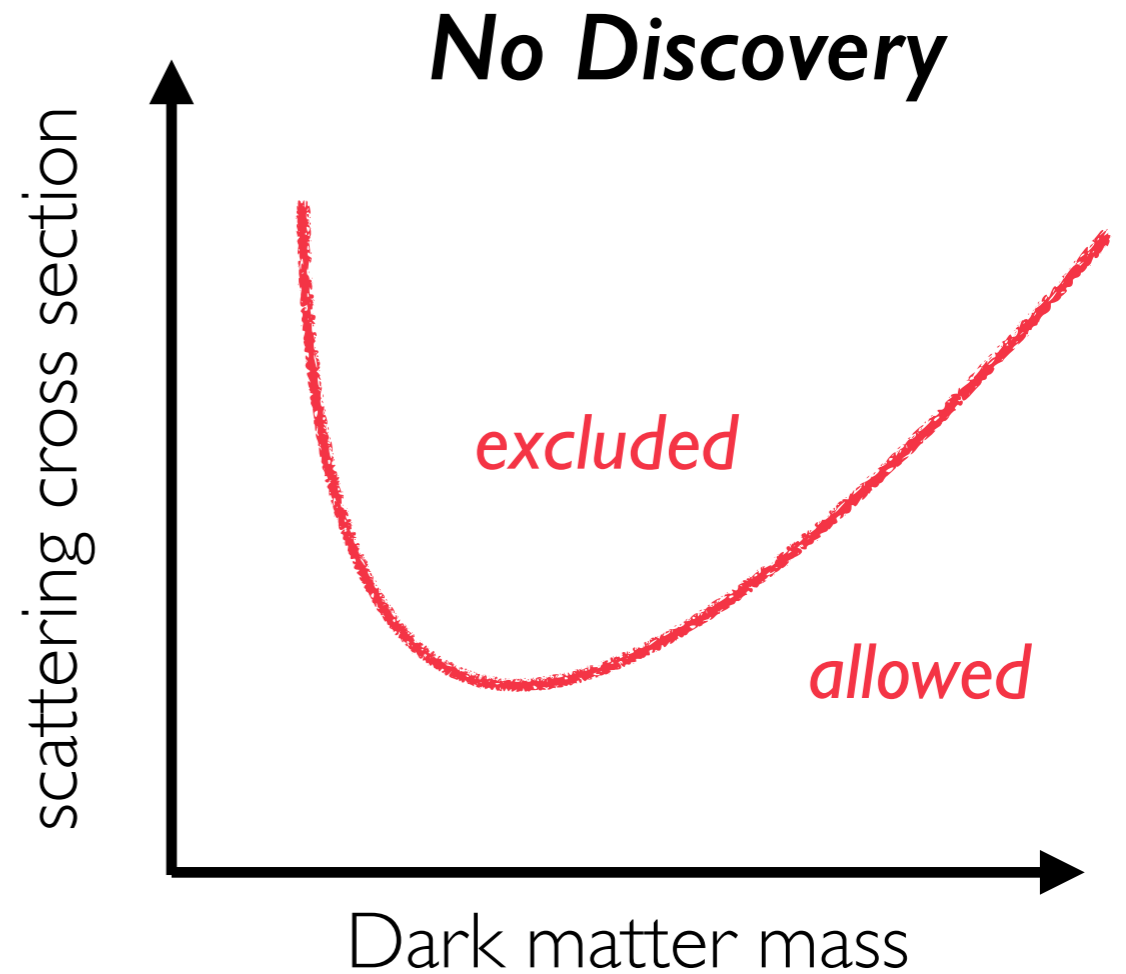
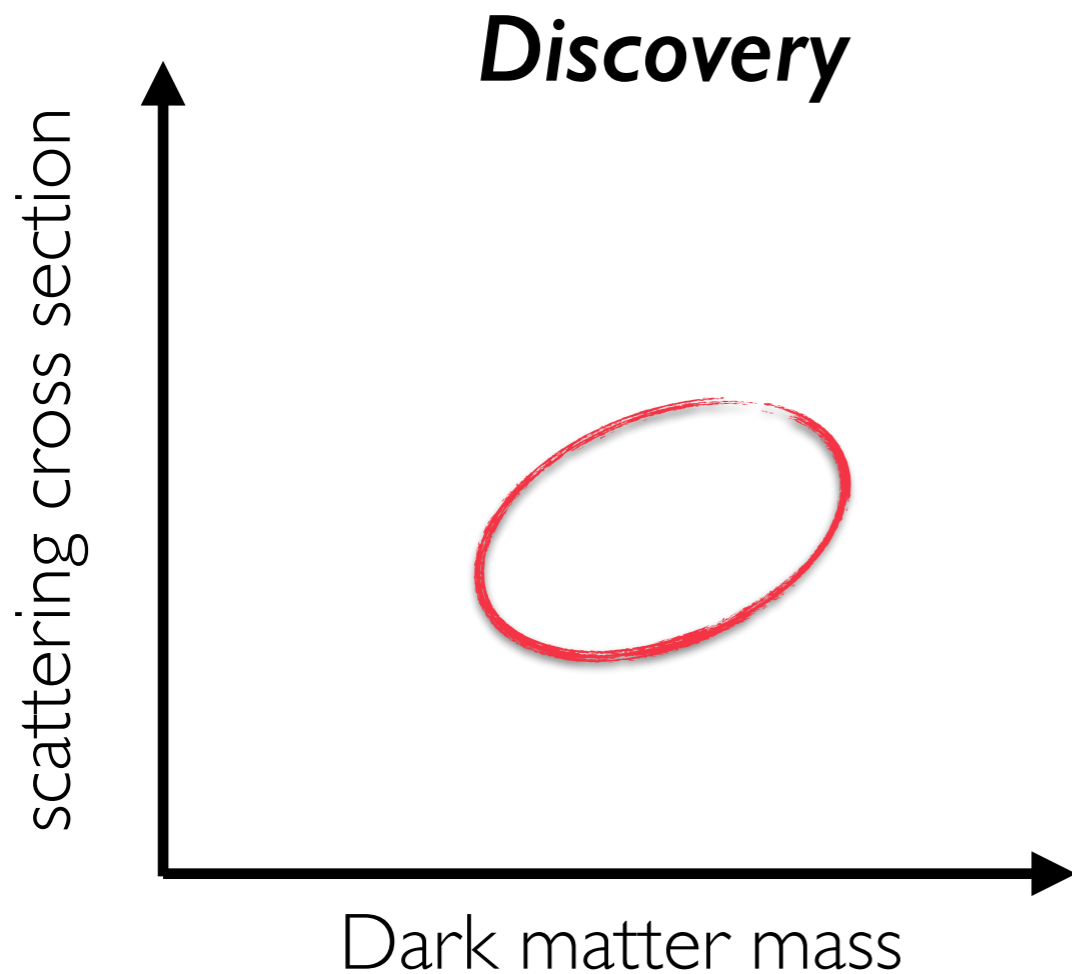
$$\rho_{\text{DM}}(@\text{Sun}) \sim 0.3 \text{ GeV}/\text{cm}^3$$

$$\phi_{\text{thumb}} \sim 10^7 \left(\frac{m_{\text{proton}}}{m_{\text{DM}}} \right) \text{ particles/s}$$

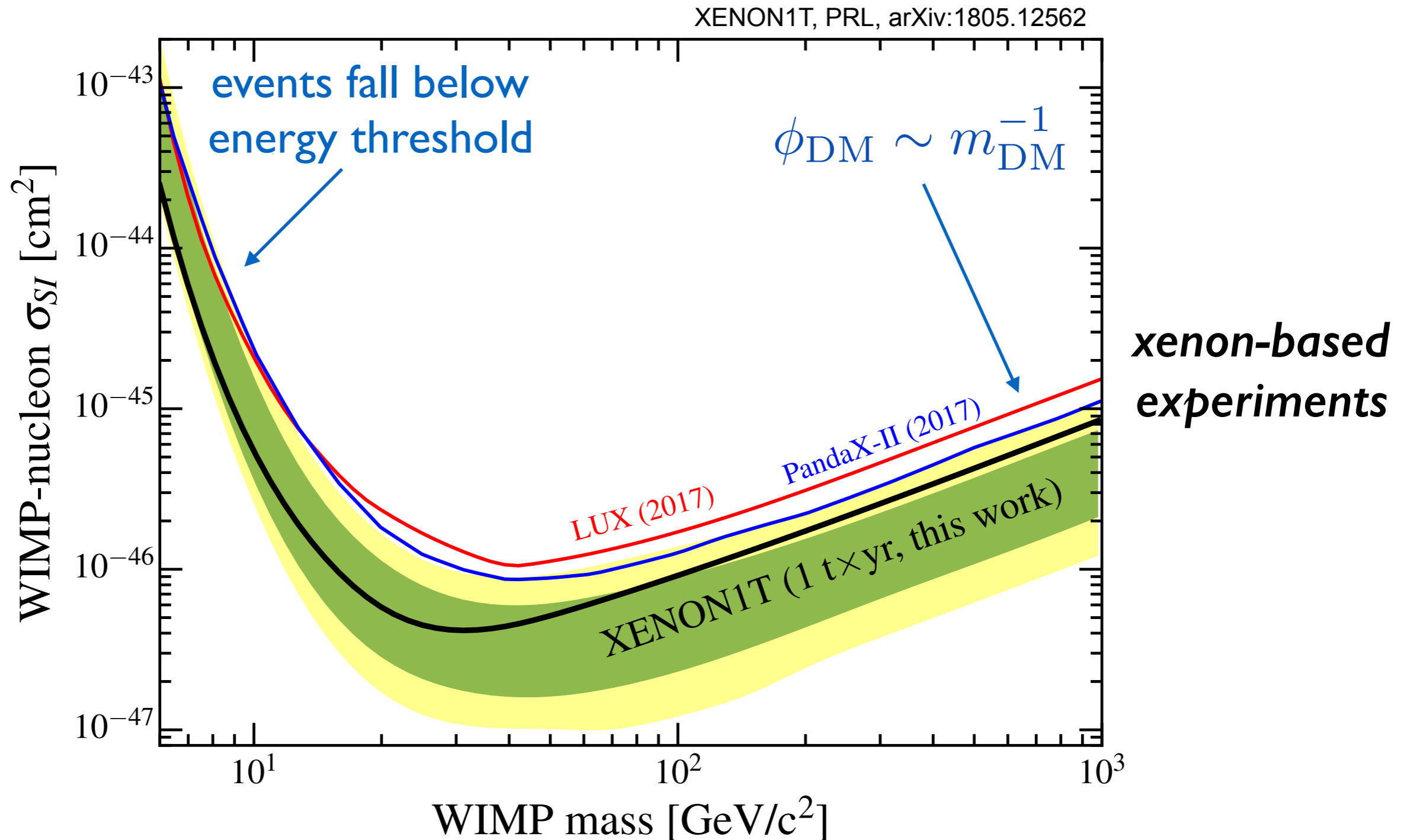
Particle physics signals

Measurement/constraints on

1. *Dark matter mass*
2. *Scattering cross section with nucleons*

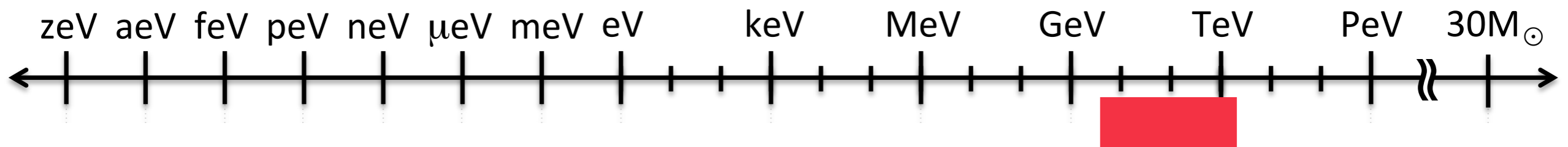
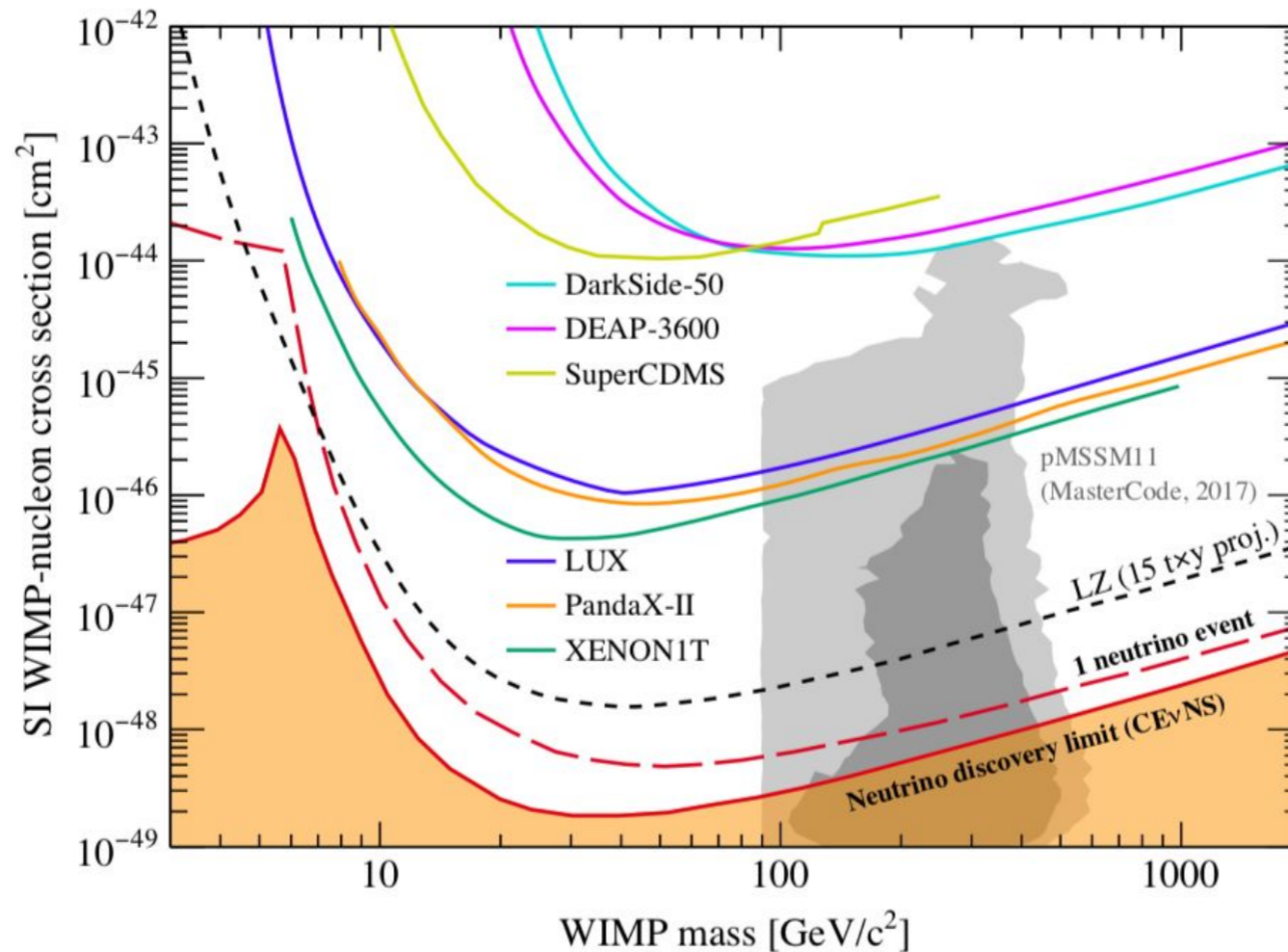


Nuclear recoils: standard WIMP searches



Nuclear recoils: standard WIMP searches

The search for WIMPs continues...

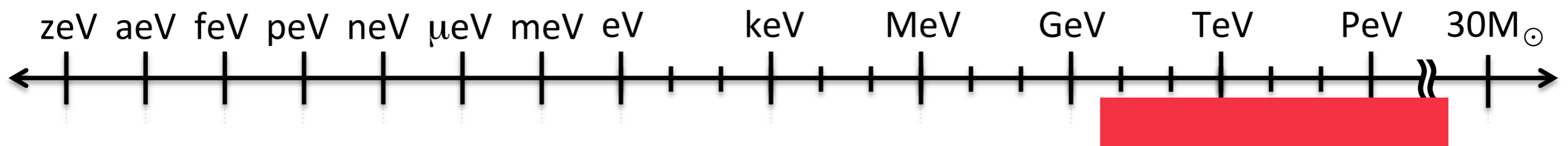
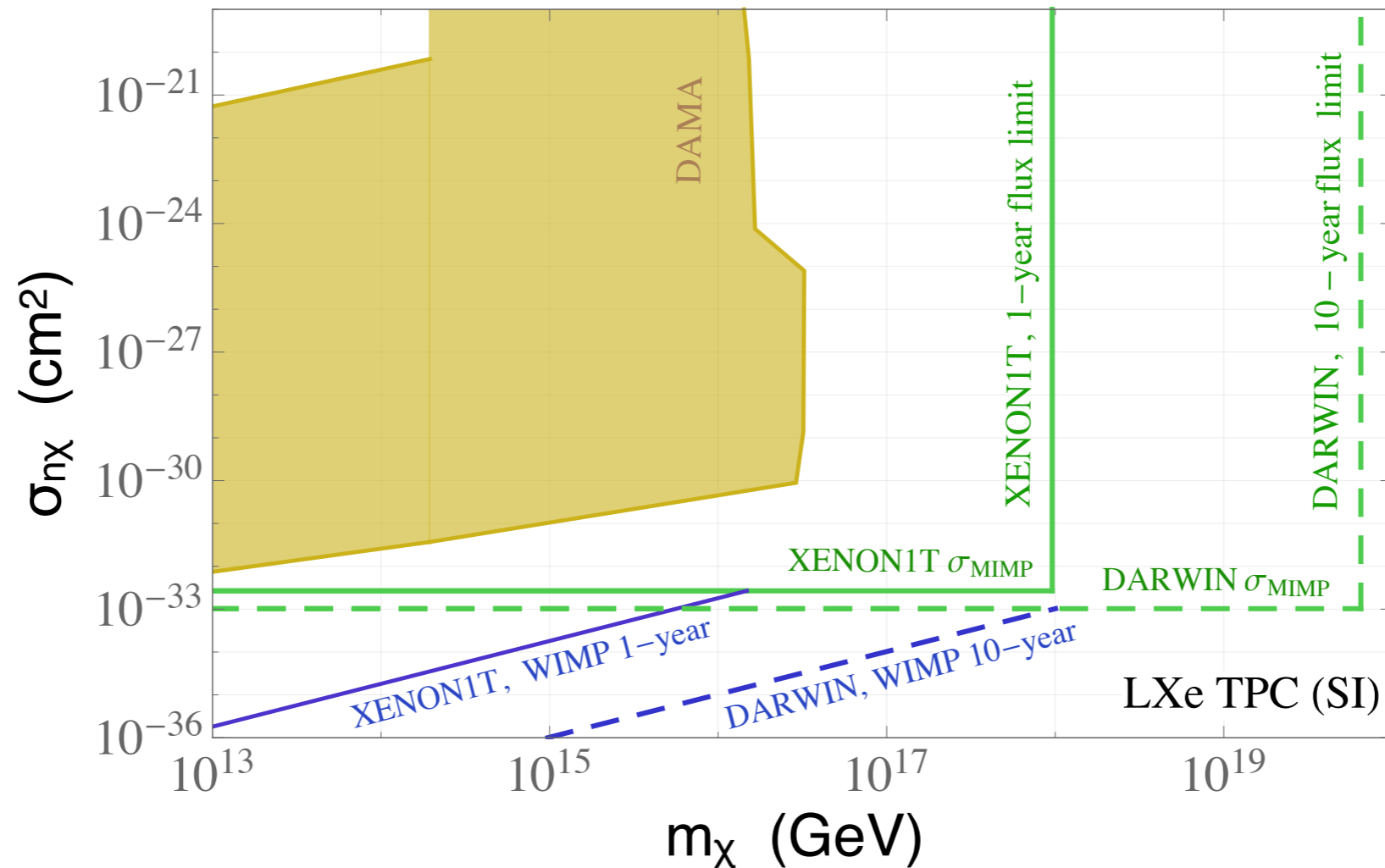


...what happens outside of this normal WIMP mass range?

Nuclear recoils: above the normal WIMP range

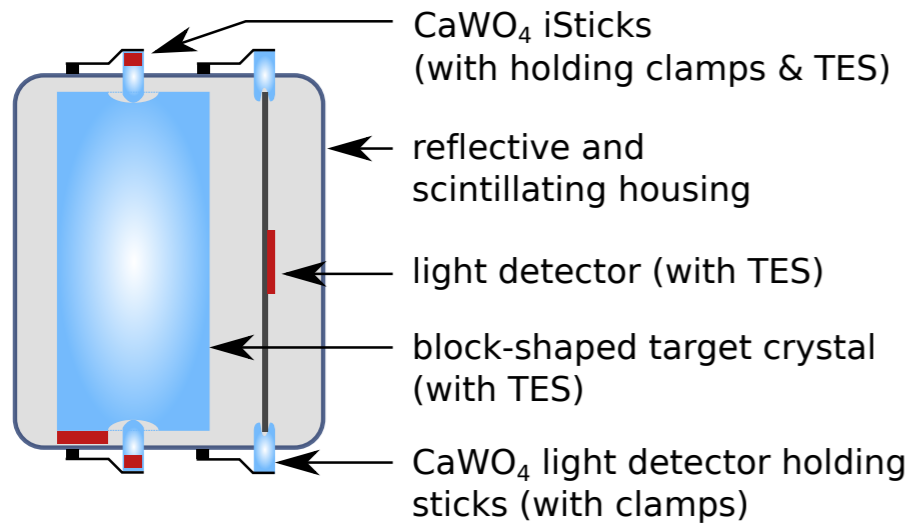
Xenon detectors can probe all the way to the Planck scale masses

Bramante et al, PRD, arXiv:1803.08044



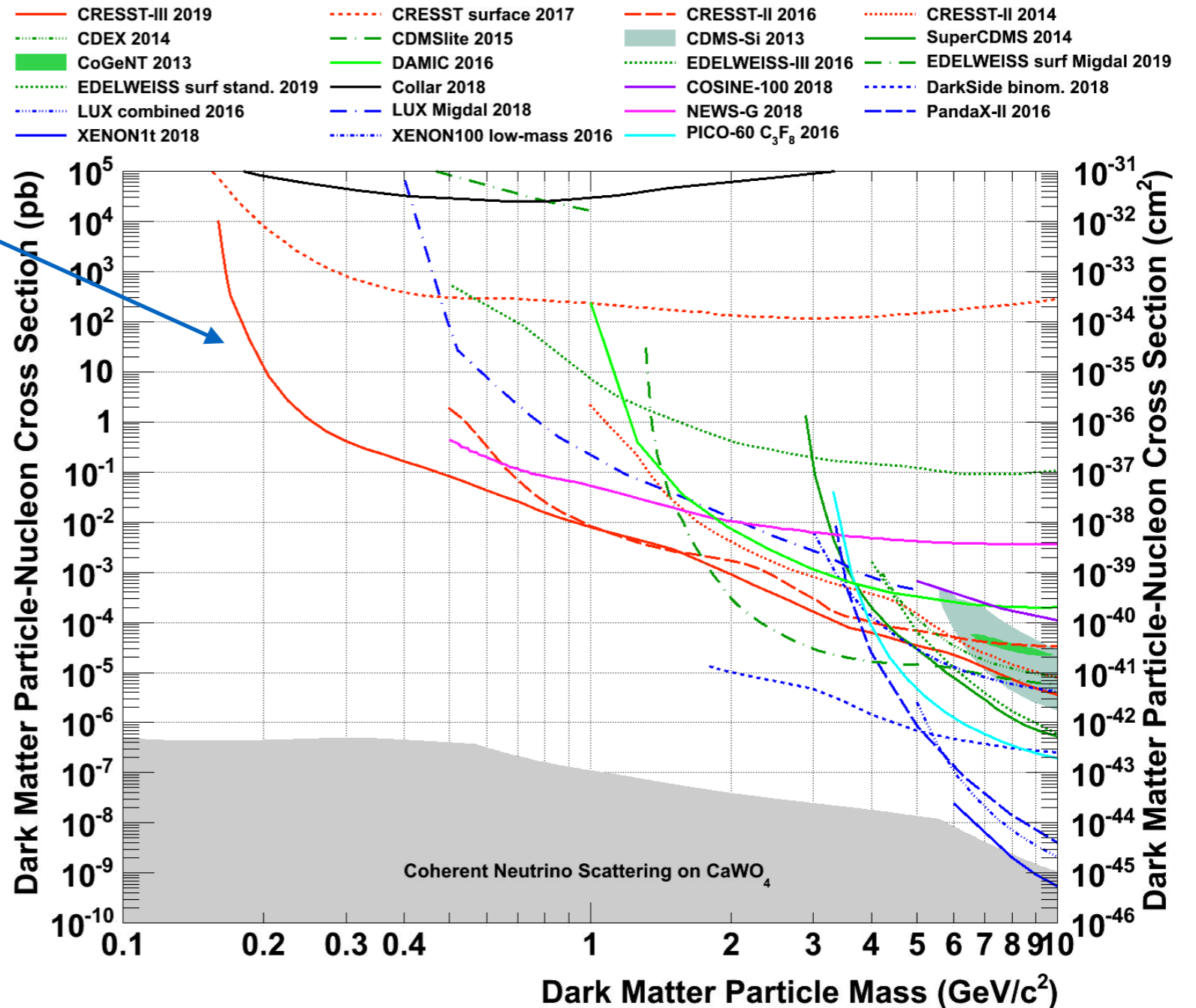
Nuclear recoils: below the normal WIMP range

CRESST-III



Detector mass: 24 grams
Detector threshold: 30 eV

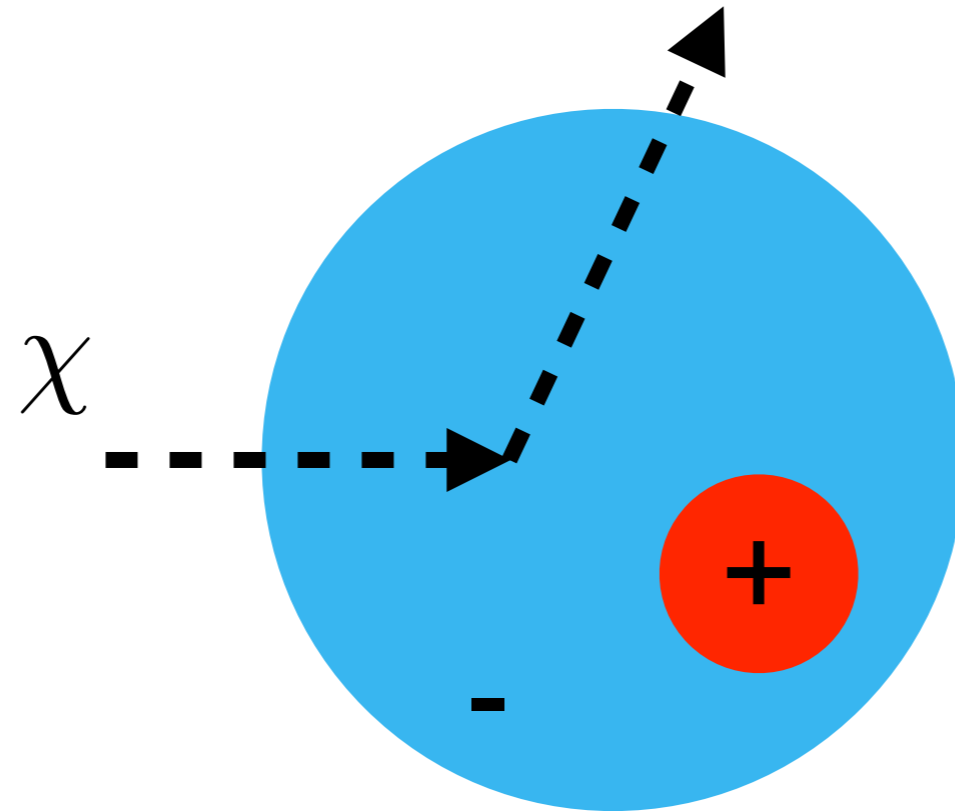
CRESST-III, arXiv:1904.00498



Beyond nuclear recoils: Migdal effect

Emission from the recoiling atom

sub-GeV dark matter in xenon:



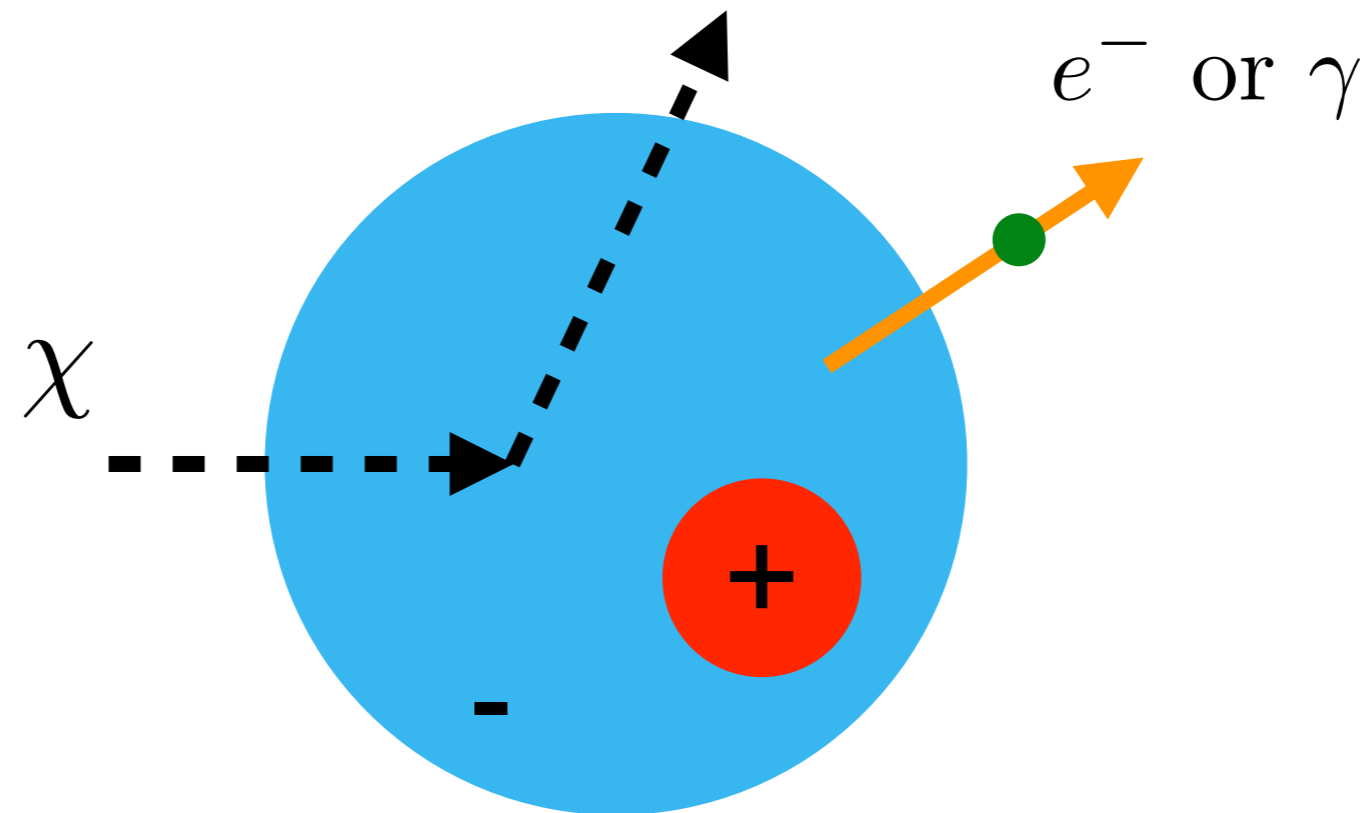
nucleus gets a nudge

$$E_{\text{recoil}} \lesssim 0.1 \text{ keV}$$

nuclear recoil below energy threshold

Emission from the recoiling atom

sub-GeV dark matter in xenon:



nucleus gets a nudge

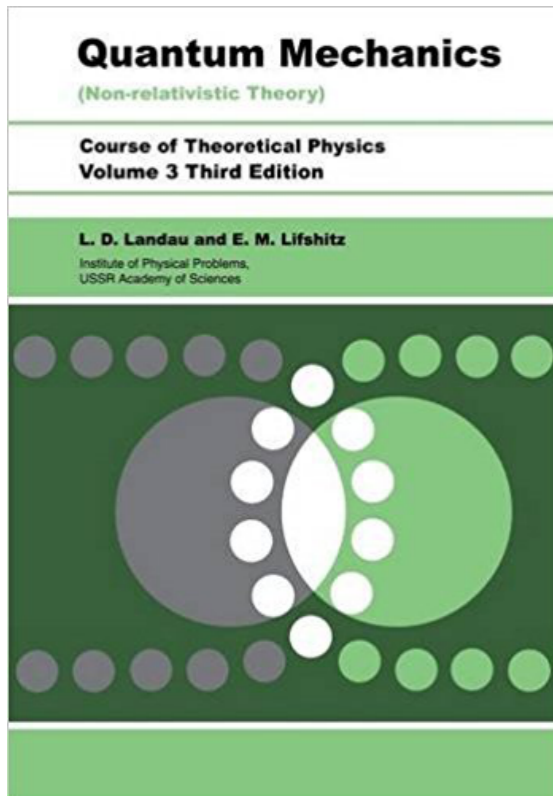
$$E_{\text{recoil}} \lesssim 0.1 \text{ keV}$$

nuclear recoil below energy threshold

...but electrons and photons can be emitted from the atom

Migdal 1939
Kouvaris & Pradler PRL 1607.01789

Electron emission: Migdal effect

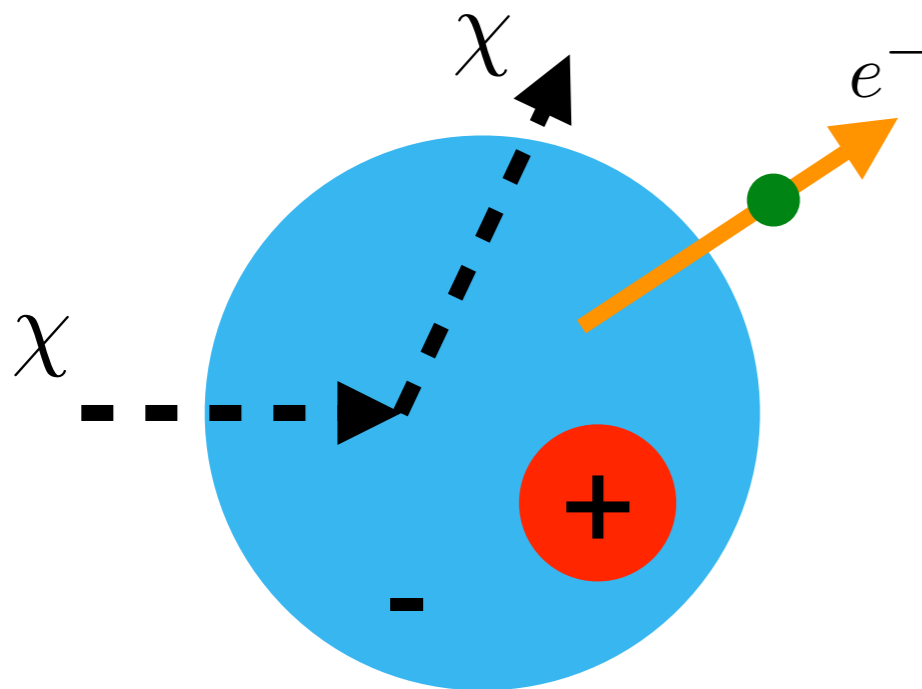


PROBLEM 2. The nucleus of an atom in the normal state receives an impulse which gives it a velocity v ; the duration τ of the impulse is assumed short in comparison both with the electron periods and with a/v , where a is the dimension of the atom. Determine the probability of excitation of the atom under the influence of such a “jolt” (A. B. MIGDAL 1939).

SOLUTION. We use a frame of reference K' moving with the nucleus after the impact. By virtue of the condition $\tau \ll a/v$, the nucleus may be regarded as practically stationary during the impact, so that the co-ordinates of the electrons in K' and in the original frame K immediately after the perturbation are the same. The initial wave function in K' is

$$\psi_0' = \psi_0 \exp(-i\mathbf{q} \cdot \sum_a \mathbf{r}_a), \quad \mathbf{q} = m\mathbf{v}/\hbar,$$

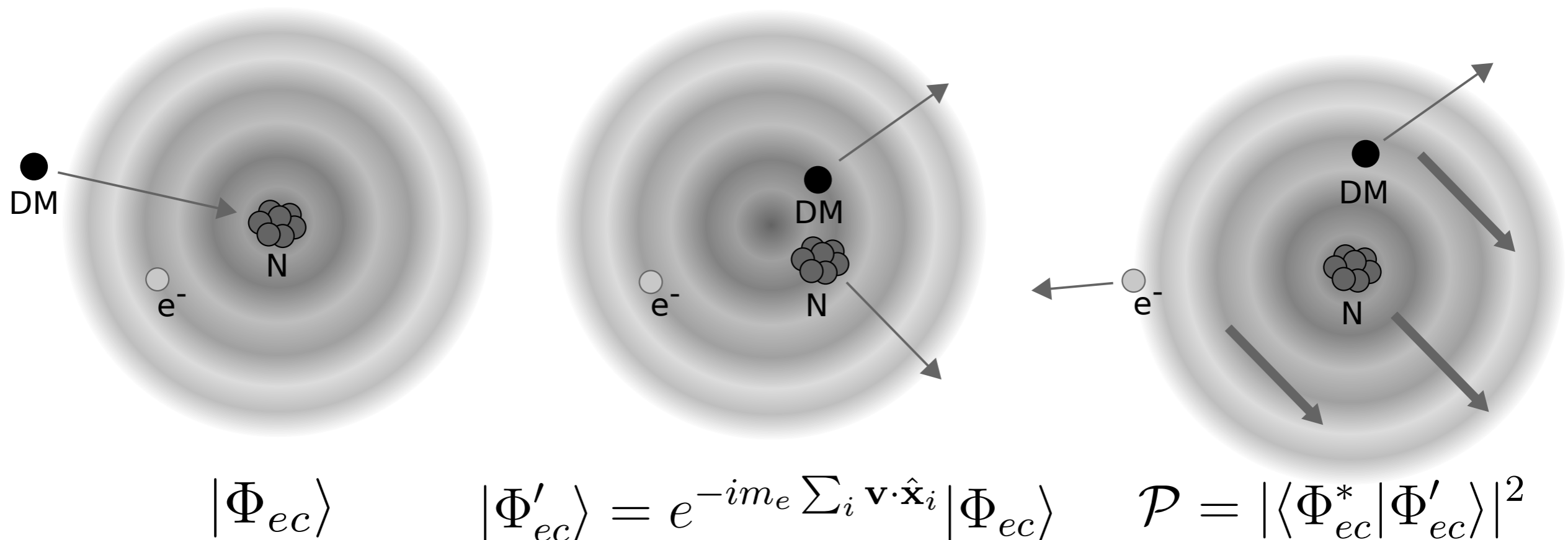
where ψ_0 is the wave function of the normal state with the nucleus at rest, and the summation



Atom emits an electron (Migdal effect)

Ibe, Nakano, Shoji, Suzuki, JHEP, arXiv:1707.07258
Dolan, Kahlhoefer, CM, PRL, arXiv:1711.09906

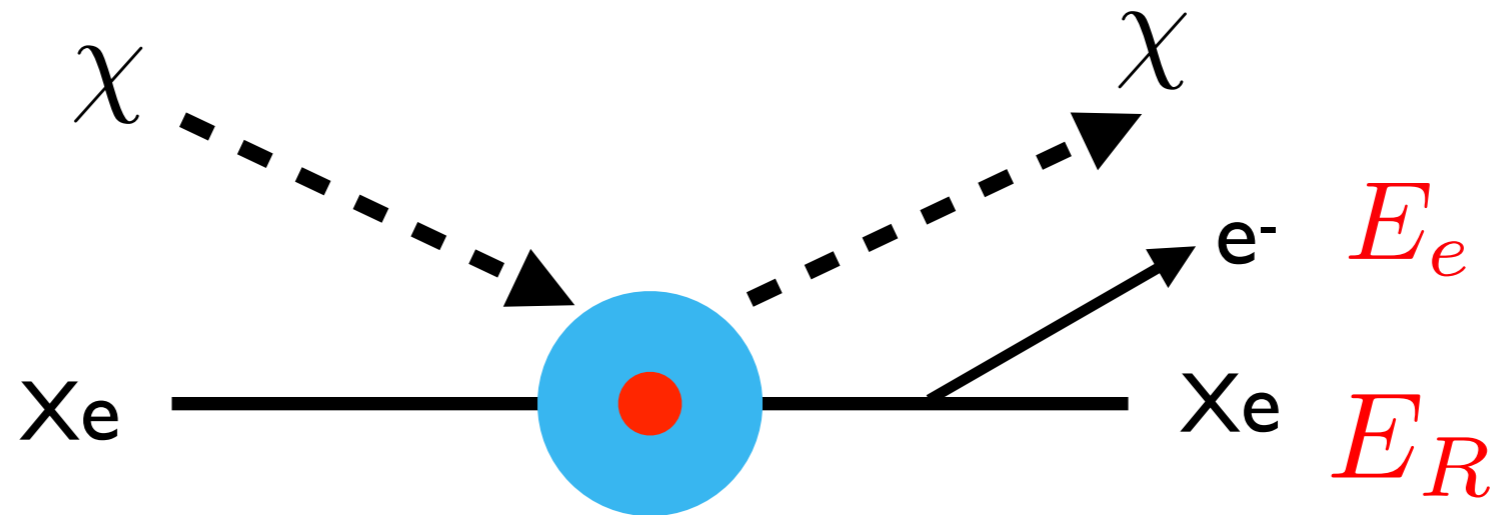
Migdal effect: updated treatment



“...it takes some time for the electrons to catch up, which causes ionisation of the atom.”

Ibe, Nakano, Shoji, Suzuki, JHEP, arXiv:1707.07258
 Dolan, Kahlhoefer, CM, PRL, arXiv:1711.09906

Electron emission: Migdal effect



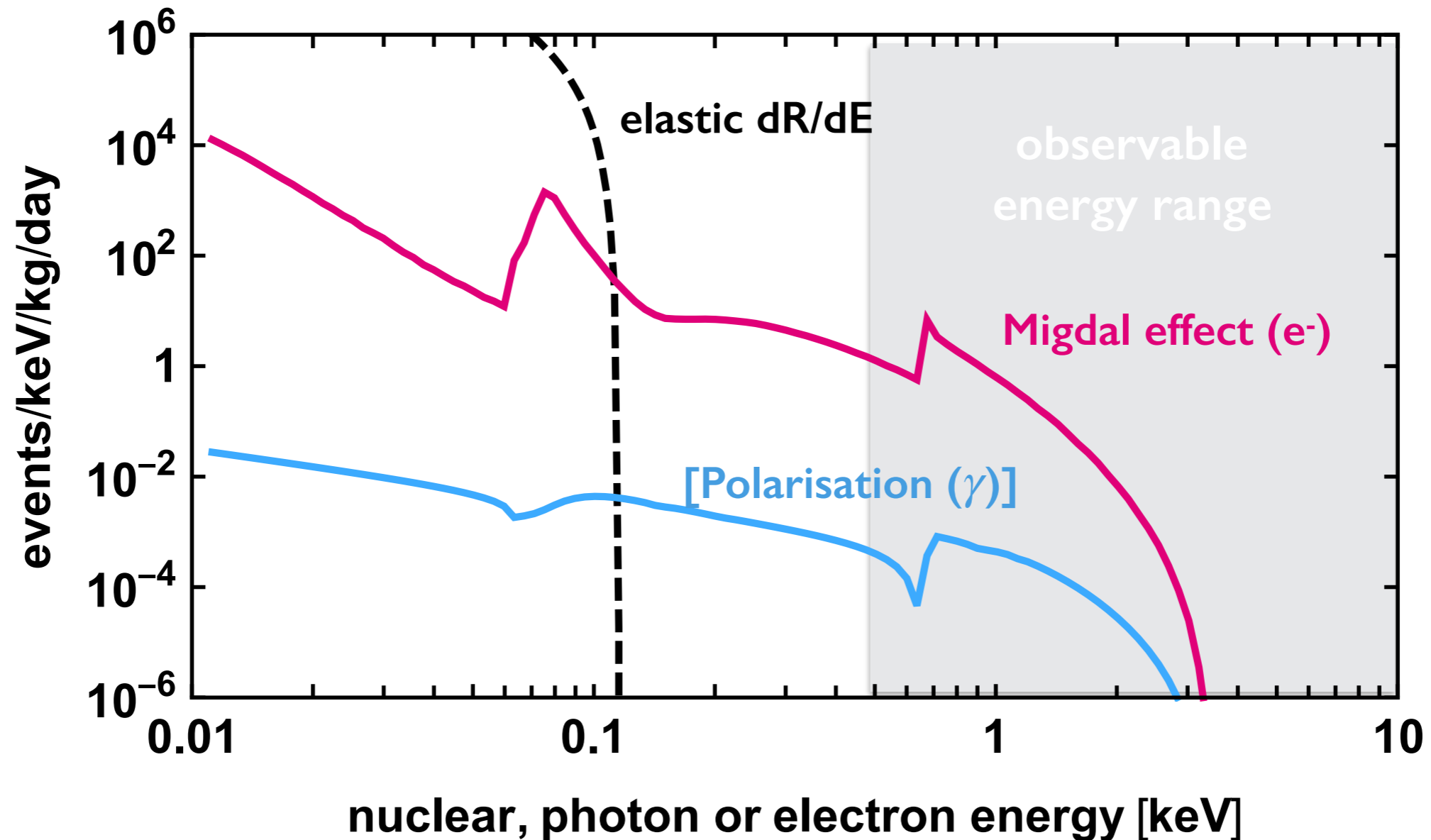
$$\frac{d^2\sigma}{dE_R dE_e} = |Z_{\text{ion}}(E_e, E_R)|^2 \times \frac{d\sigma}{dE_R} \quad \text{Usual 2-}\rightarrow\text{2 cross section}$$

‘atomic form factor’ $\sim \mathcal{P} = |\langle \Phi_{ec}^* | \Phi'_{ec} \rangle|^2 \sim 10^{-3}$

Calculated in Ibe, Nakano, Shoji, Suzuki, JHEP, arXiv:1707.07258 for different elements

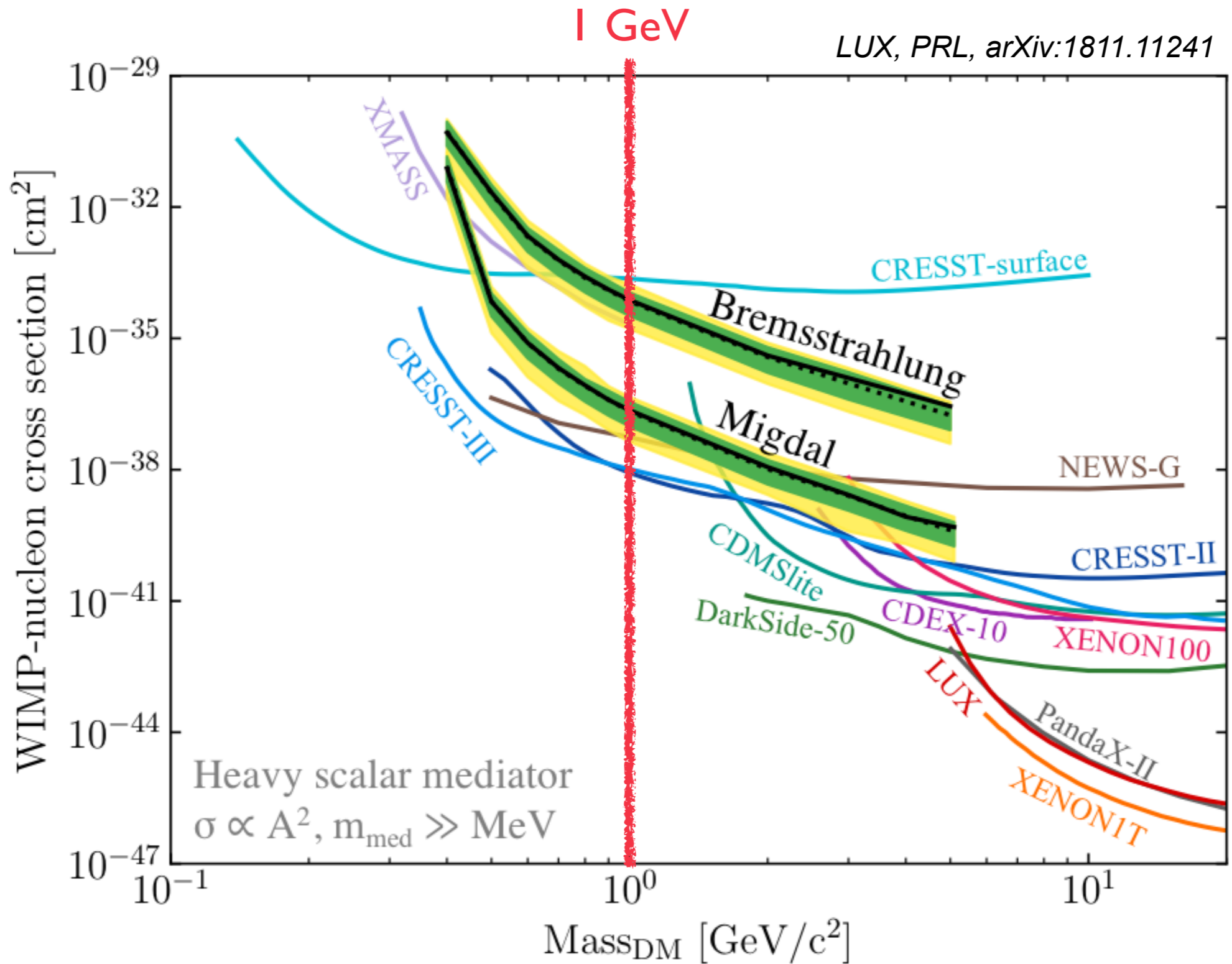
sub-GeV DM signals in xenon

Xenon, $m_{\text{DM}}=1 \text{ GeV}$, $\sigma^0=10^{-35} \text{ cm}^2$



Migdal electrons observable even for $m_{\text{DM}} \lesssim 1 \text{ GeV}$

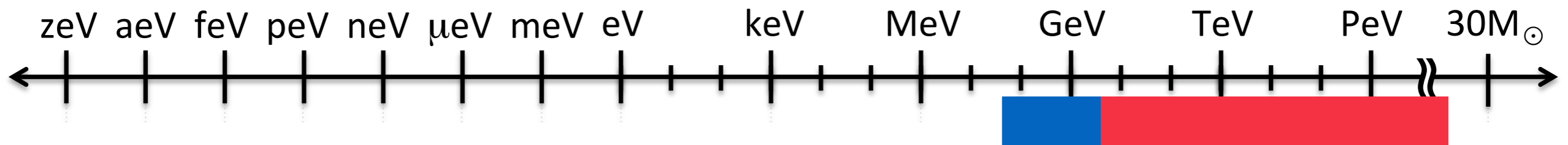
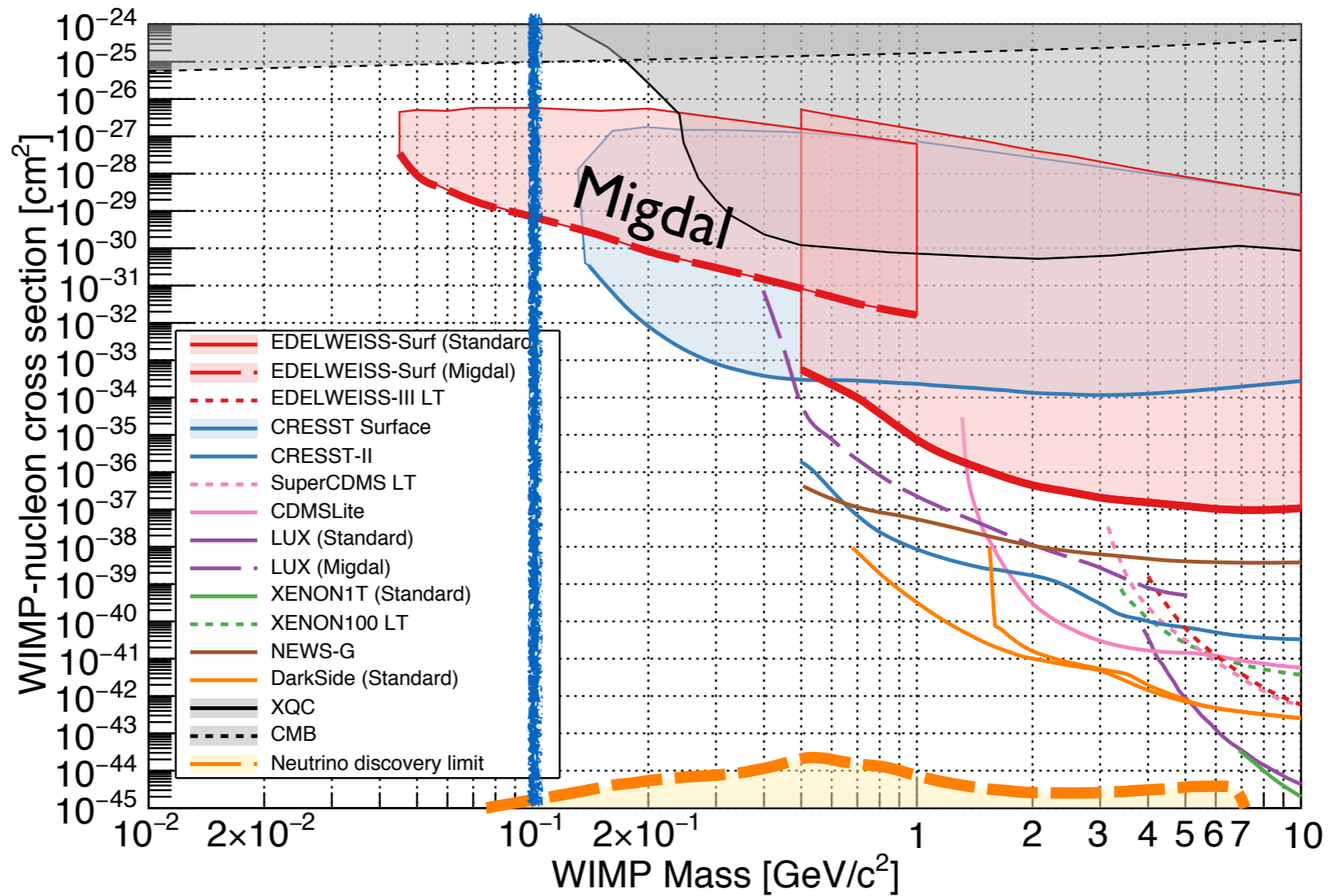
LUX (Xe) sensitive below 1 GeV



Edelweiss (Ge) sensitive below 100 MeV

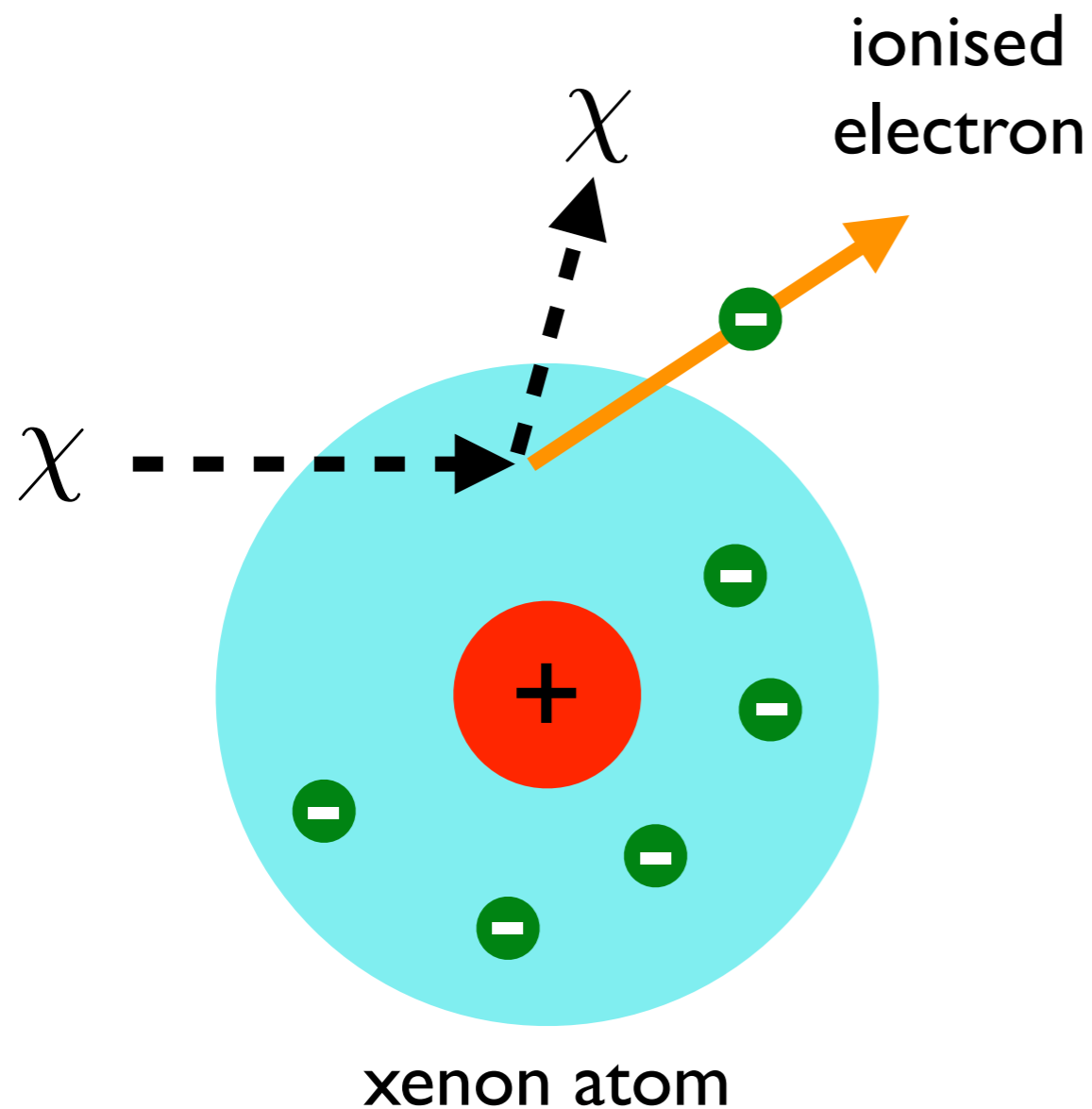
100 MeV

Edelweiss + Kavanagh
arXiv:1901.03588



Beyond nuclear recoils: electron scattering

Electron-ionisation in atoms



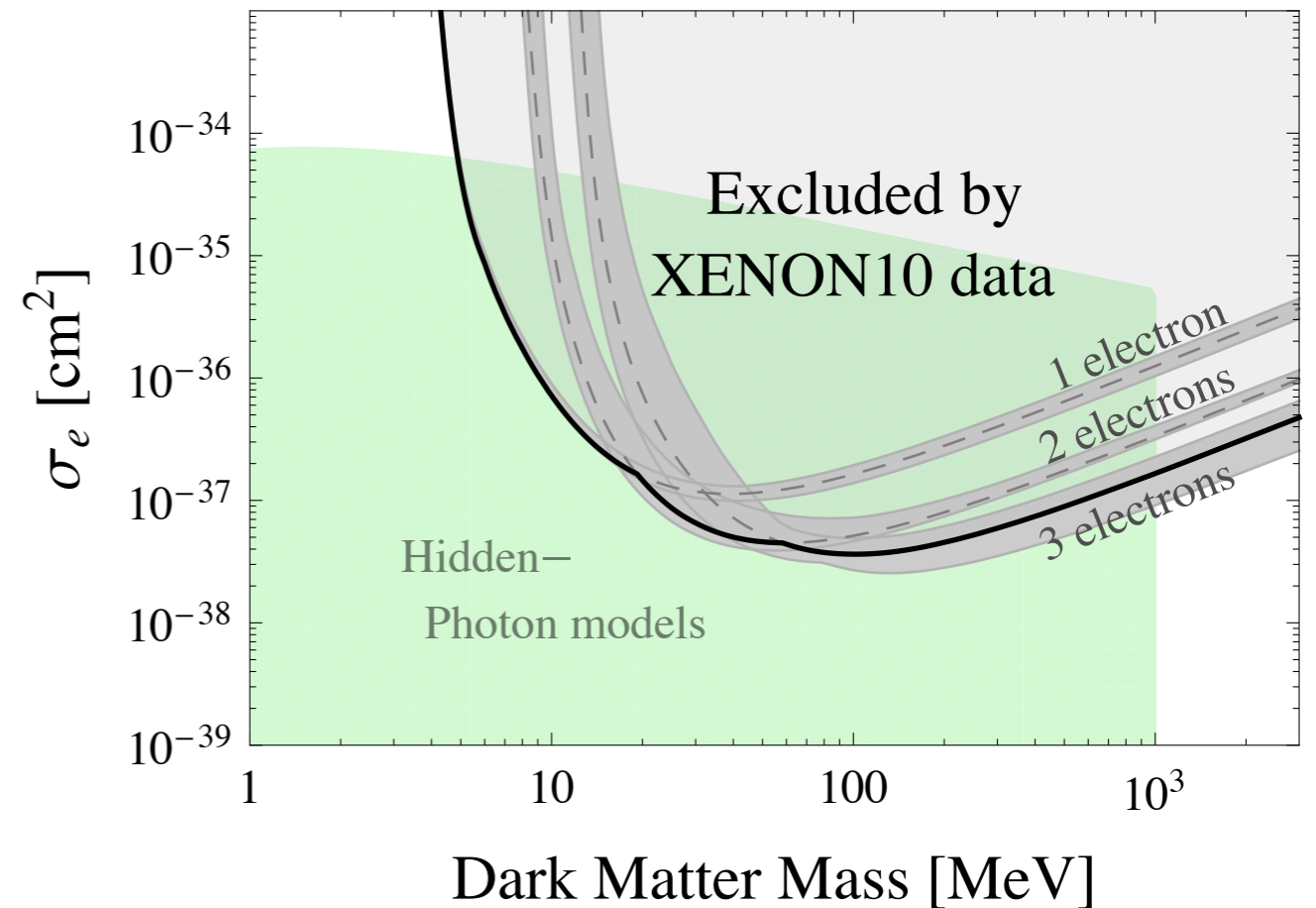
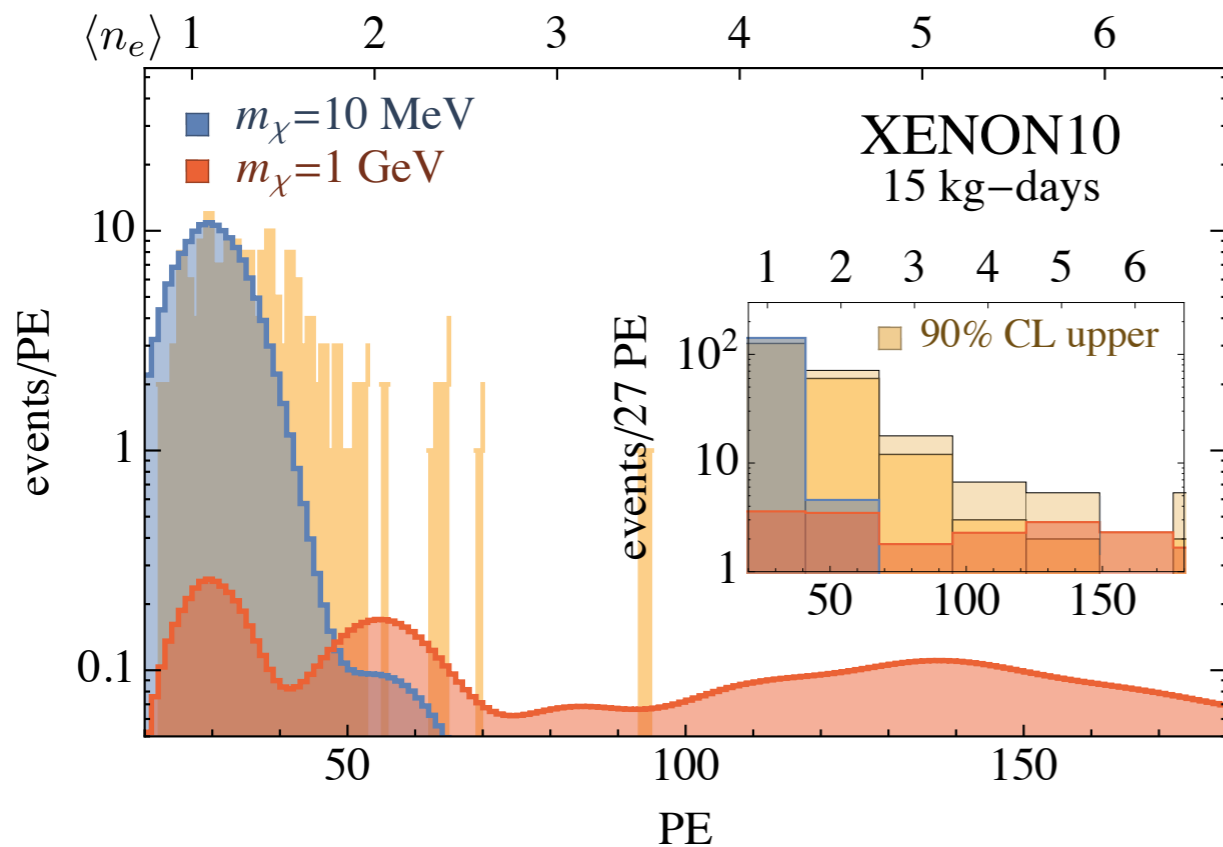
For ionisation, require:

$$\frac{1}{2}m_{\text{DM}}v_{\text{DM}}^2 \gtrsim E_{\text{binding}}(\sim 12 \text{ eV})$$

$$m_{\text{DM}} \gtrsim 5 \text{ MeV}$$

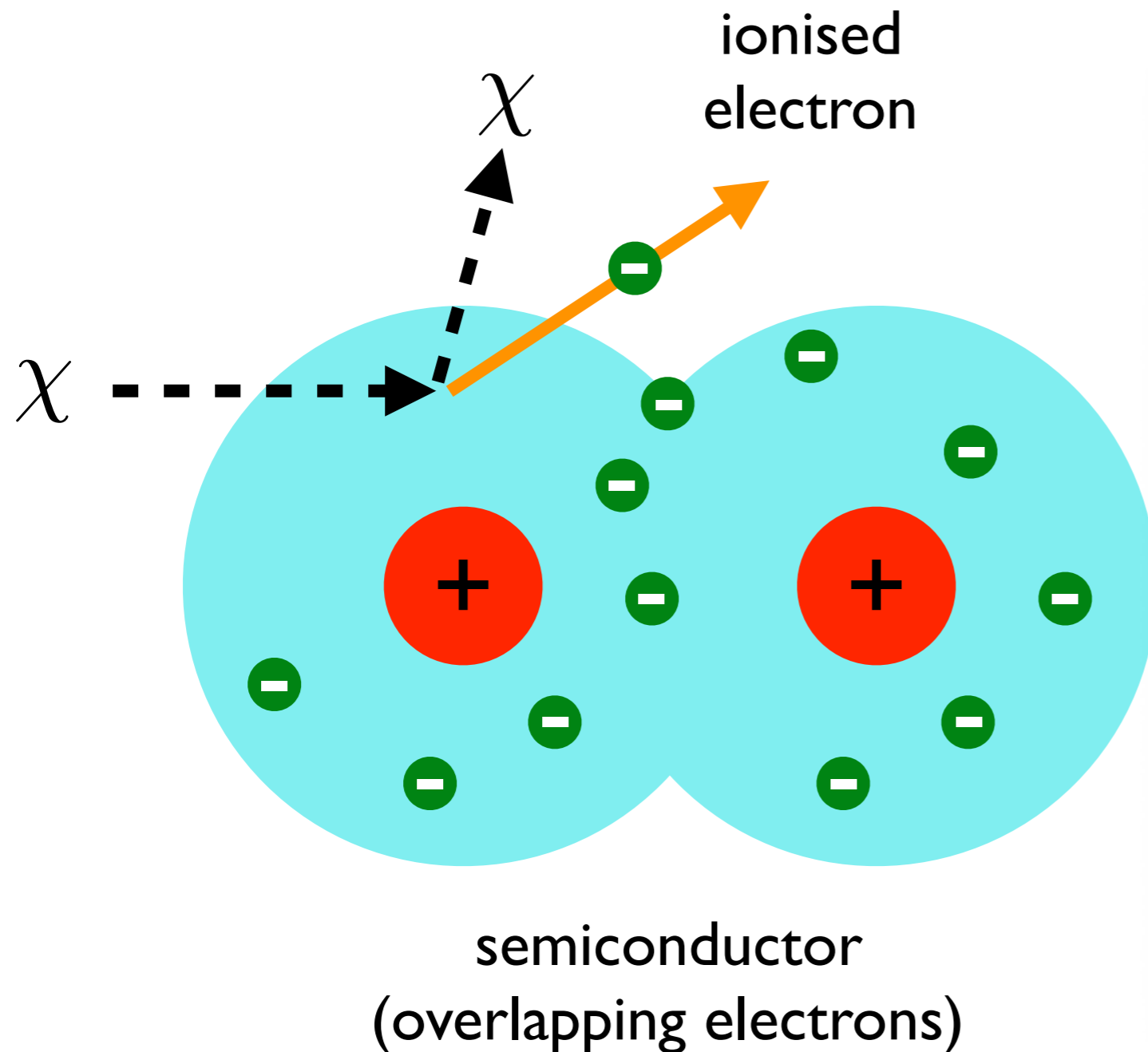
Xenon electron-ionisation constraints

- ZEPLIN, XENON10 & XENON100 can measure single electrons (S2 only)
- Also possible in argon (e.g. DarkSide-50)



Essig et al, PRL, arXiv:1206.2644
 Essig et al, PRD, arXiv:1703.00910

Electron ionisation in semi-conductors



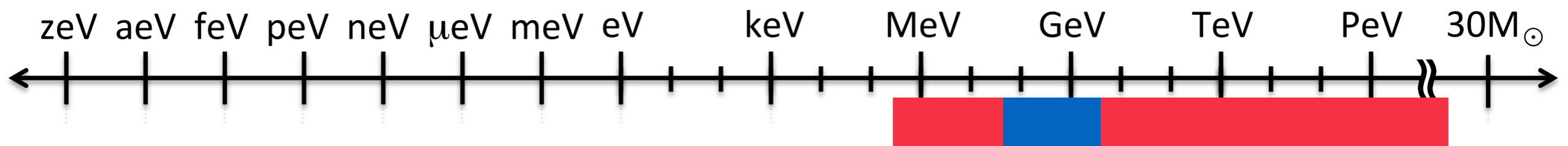
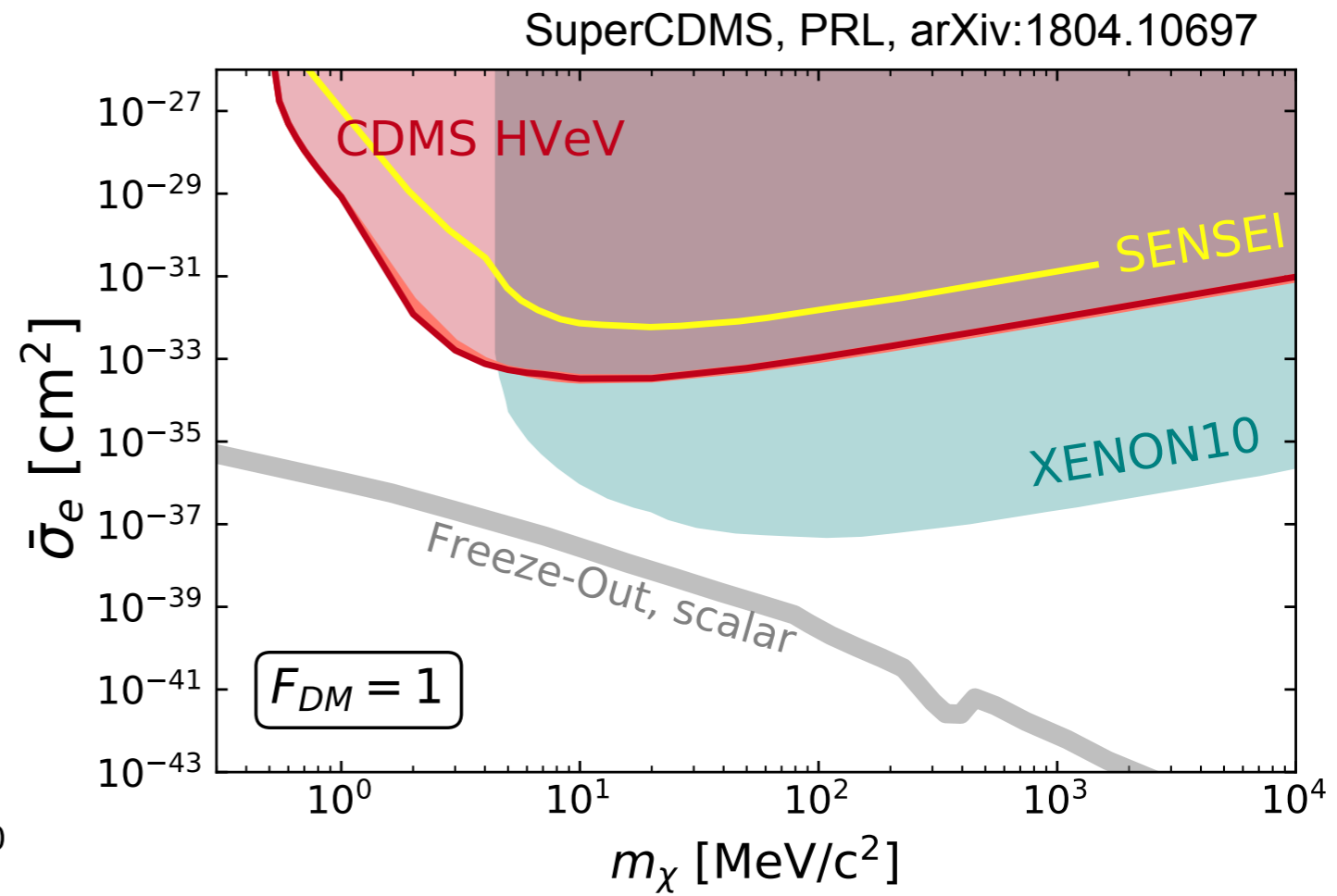
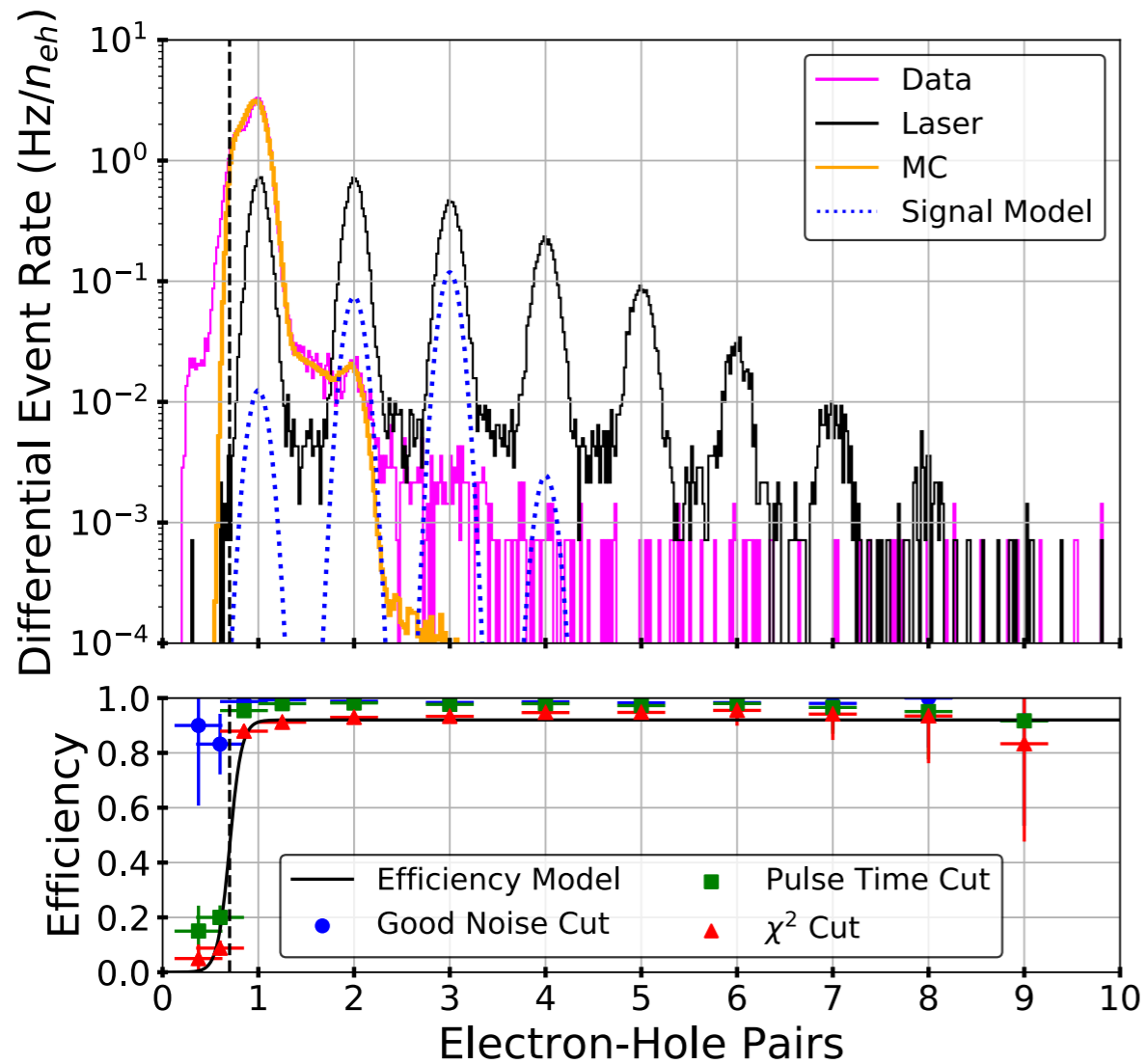
For ionisation, require:

$$\frac{1}{2} m_{\text{DM}} v_{\text{DM}}^2 \gtrsim E_{\text{binding}}$$

$$E_{\text{binding}}^{\text{semi-conduct}} \sim 1 \text{ eV}$$

$$m_{\text{DM}} \gtrsim 0.5 \text{ MeV}$$

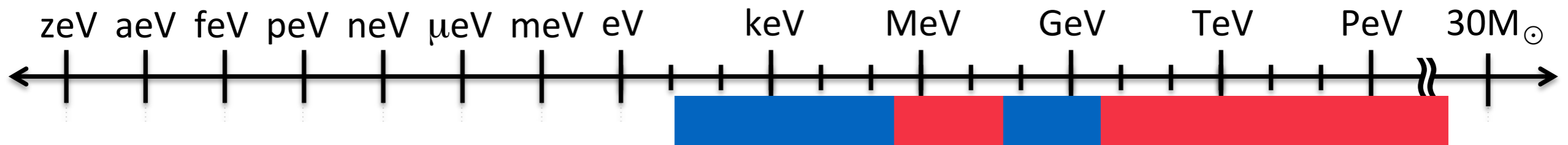
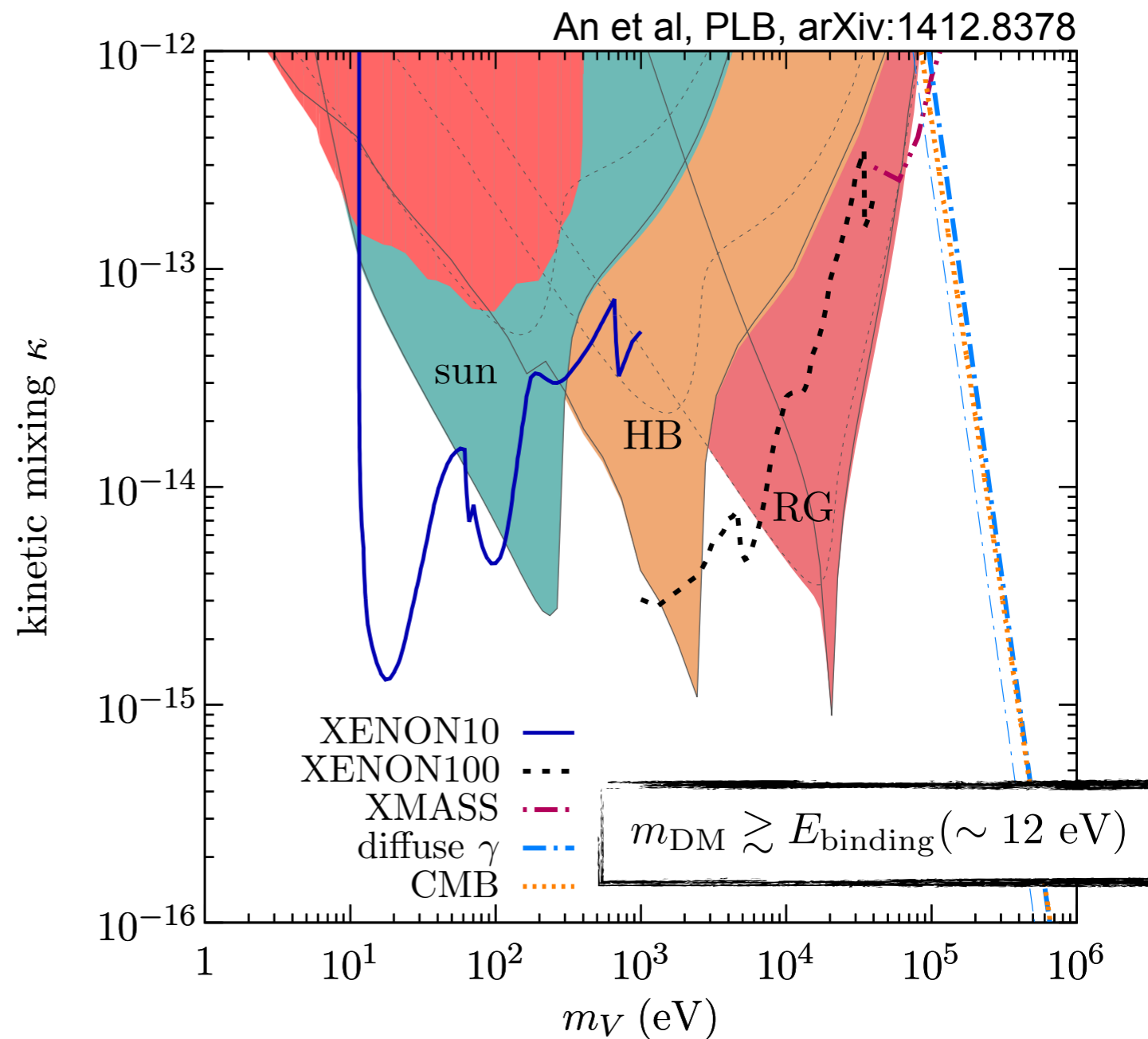
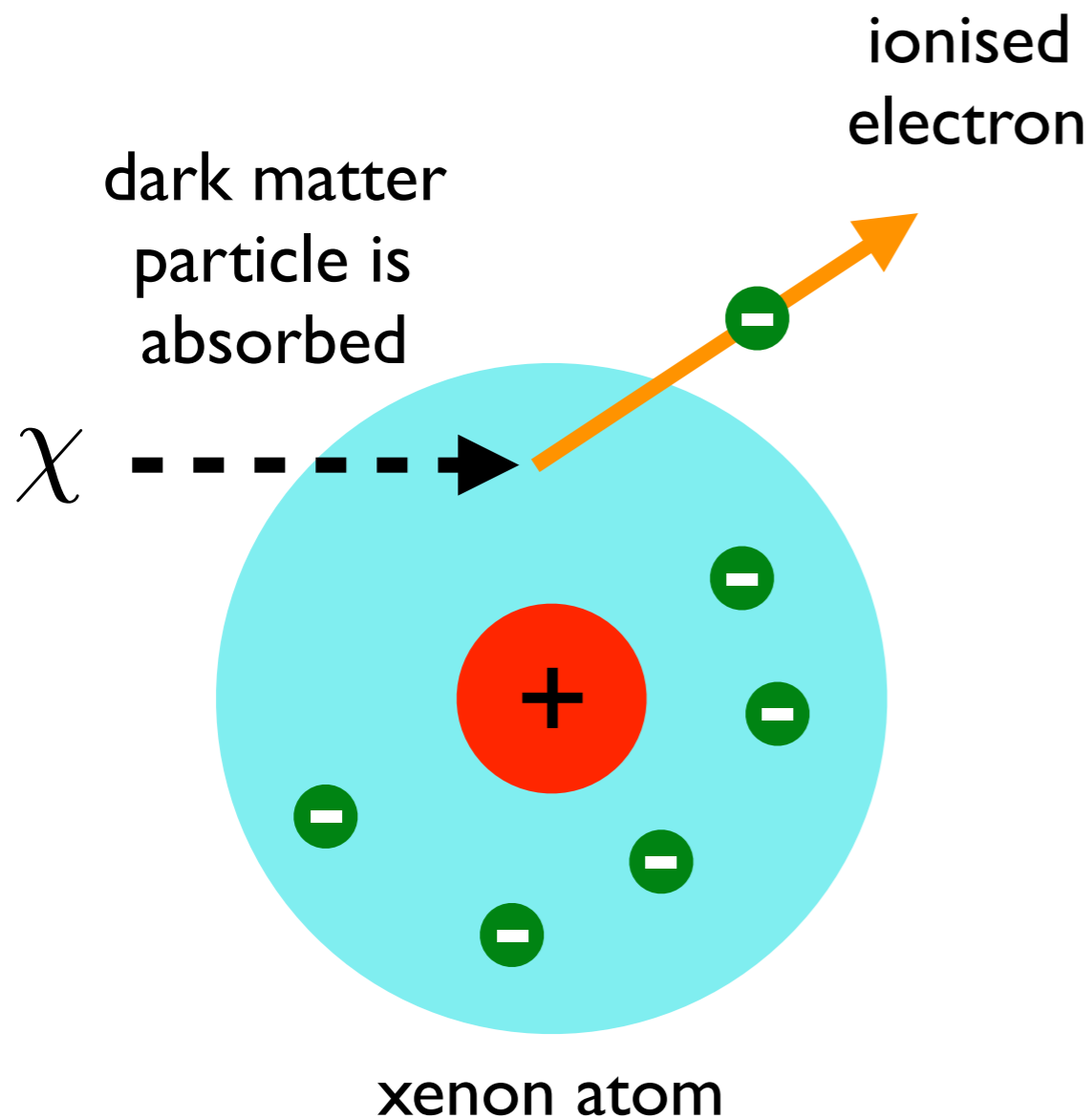
Electron ionisation in semi-conductors



Can we go to even lower masses?

Yes! Dark matter conversion

$$\mathcal{L} \supset V_\mu \bar{\psi}_e \gamma^\mu \psi_e \text{ or } A \bar{\psi}_e \gamma^5 \psi_e$$



Yes! Semi-relativistic dark matter

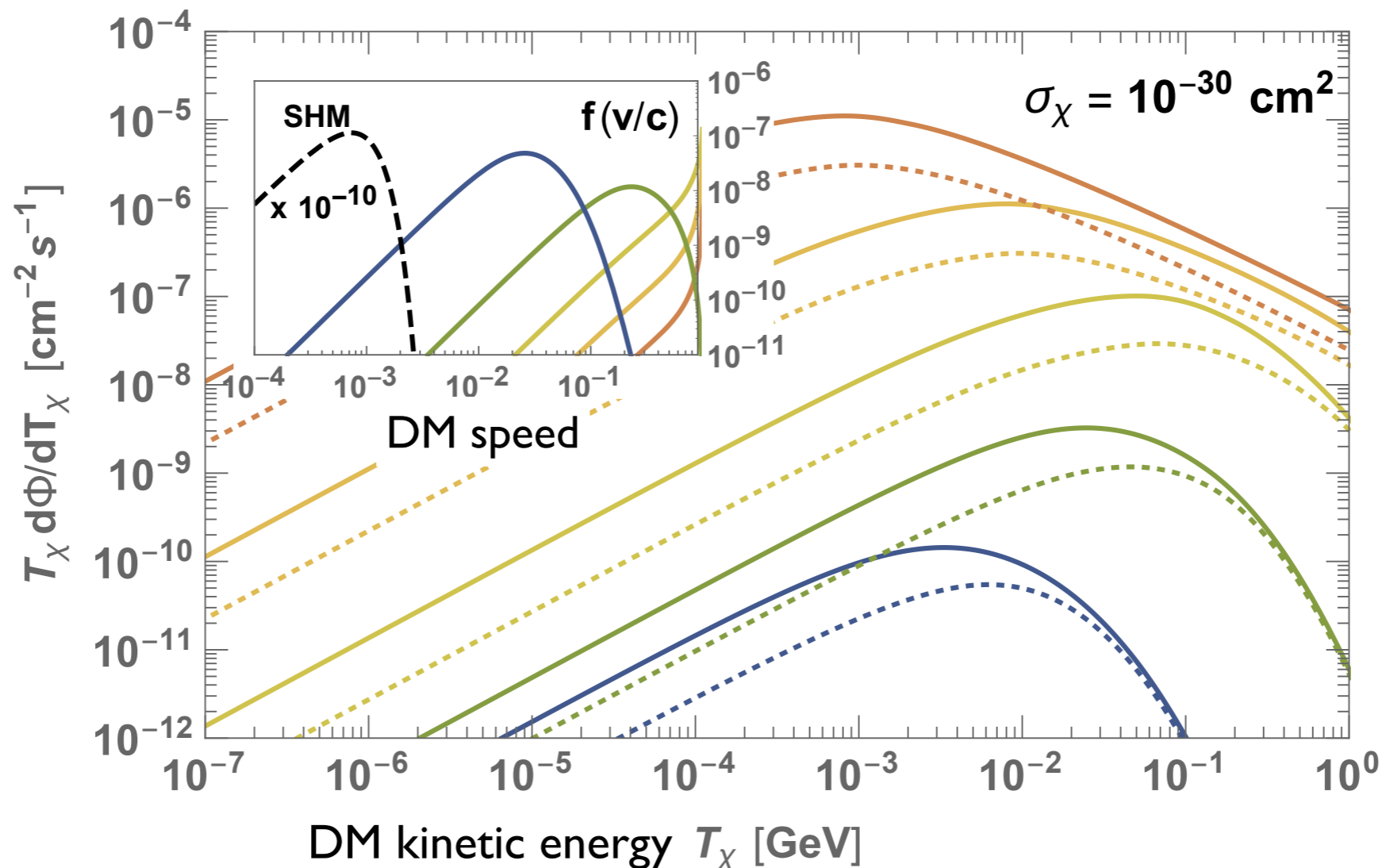
High energy cosmic ray (p or e-)



DM (essentially at rest)



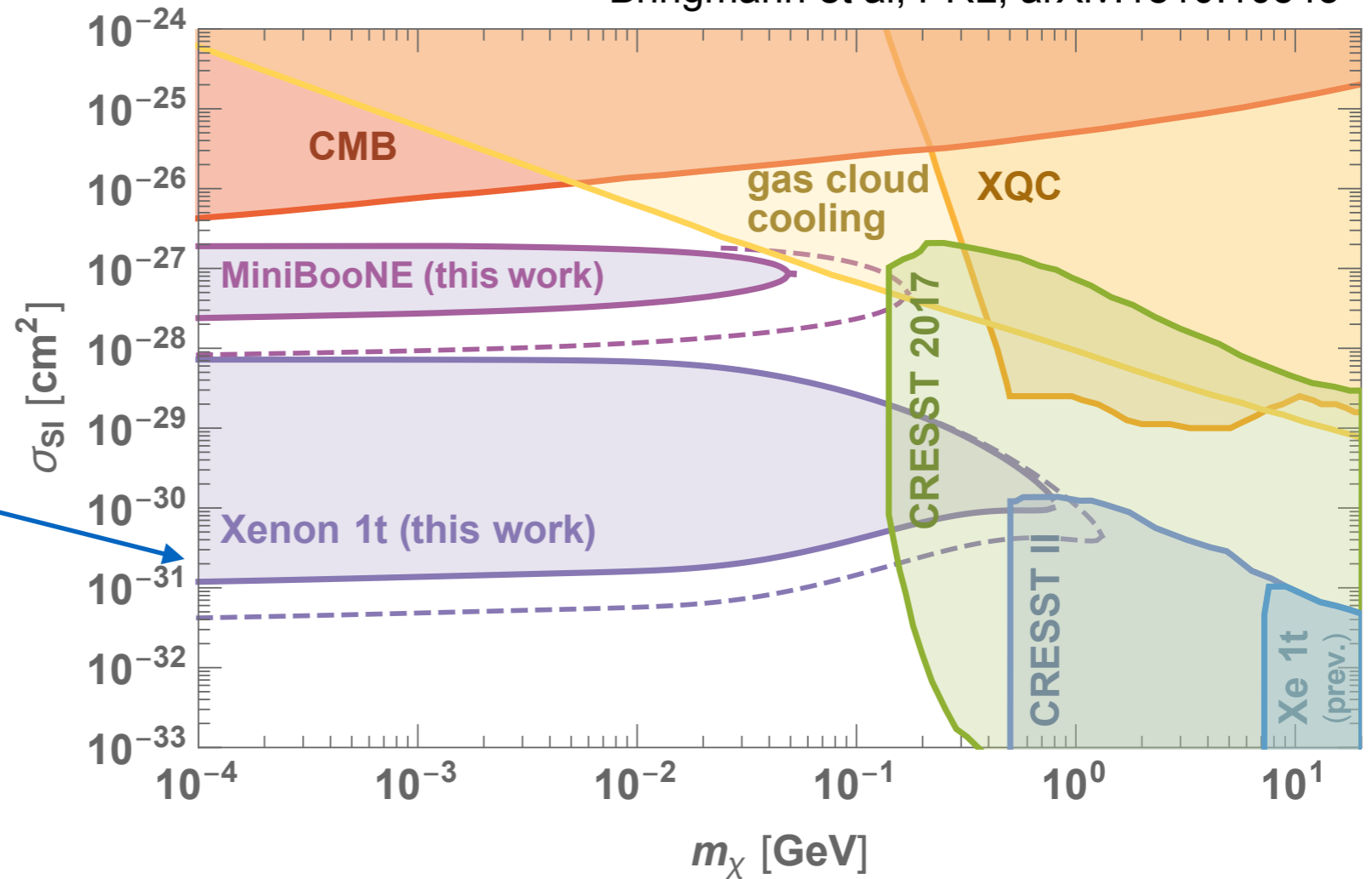
Small irreducible semi-relativistic DM component from cosmic rays scattering



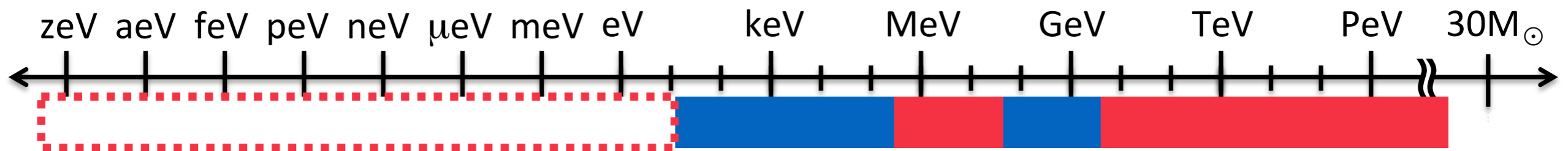
Bringmann et al,
PRL, arXiv:
1810.10543

Yes! Semi-relativistic dark matter

Bringmann et al, PRL, arXiv:1810.10543



In principle, limits extend down to any mass



Take home message

- ‘WIMP’ direct detection experiments are pretty versatile dark matter detectors
- They are sensitive to any DM that gives \sim keV energies
This does not only have to be a WIMP
- Nuclear scattering is important, but also interesting is
 - *electron scattering*
 - *mixed nuclear-electron signals*