

Exploring Phase Transitions and Gravitational Waves in the Conformal Two-Higgs-Doublet Model (C2HDM)

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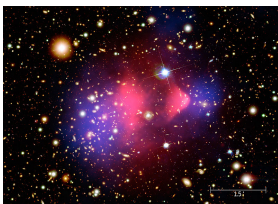


Table Of Contents

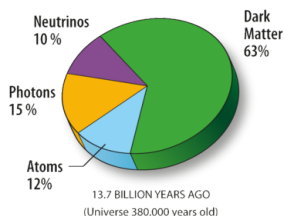
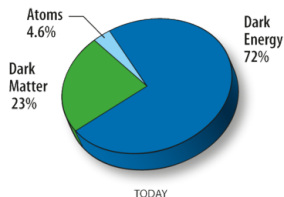
- 1 Introduction
 - Motivation
 - Phase transitions and gravitational waves
- 2 Model and setup
 - Model
 - Effective potential and phase transitions
 - Analysis pipeline
- 3 Preliminary results
 - Preliminary results
 - Comparison to previous work
- 4 Conclusion
- 5 Bibliography

Why go beyond the Standard Model?

- The nature of dark energy
- The nature of dark matter
- The nature and origin of neutrino masses
- The hierarchy problem
- The baryon asymmetry problem
- ⋮



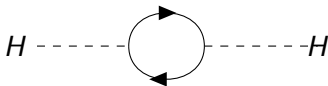
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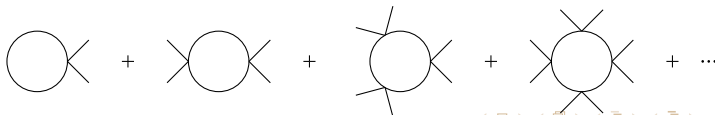
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The hierarchy problem

- Why is the Higgs boson so light?
- How to fix? [P. Saake, 2024]
 - SUSY
 - Composite Higgs
 - Conformal symmetry!
- Zero tree-level mass and radiative generation of the mass of the SM Higgs boson keeps it naturally light → Conformal BSM models
- Radiative mass generation and spontaneous symmetry breaking from the Coleman-Weinberg mechanism



[S. Coleman et al., 1973, E. Gildener et al., 1976]



The baryon asymmetry problem

- Why is there so much more matter than antimatter?
- Sakharov conditions for baryogenesis: [A. Sakharov, 1967]
 - Baryon number violation
 - Breaking of C- and CP-symmetry
 - Processes (strongly) out of thermal equilibrium $\rightarrow$ phase transition (PT)
- The SM cannot account for measured baryon asymmetry: [L. Canetti et al., 2012]
 - Too little CP-violation
 - Electroweak PT (EWPT) not strong enough
- How to fix? Modify the scalar potential! $\rightarrow$ BSM models



Source: Wikipedia

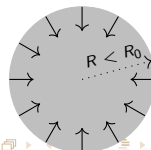
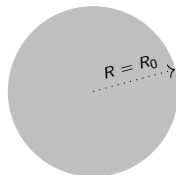
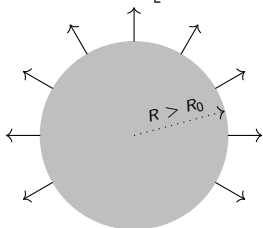
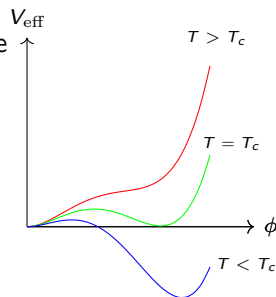
Phase transitions

- Phase: local minimum of thermal effective potential $V_{\text{eff}}(\phi, T)$
- Phase transition: transition from one local minimum to a lower one
- First-order PT: PT with a barrier
- Nucleation rate: [S. Coleman, 1977]

$$\Gamma/V = Ae^{-S_3/T}$$

- Bounce action: [A. Linde, 1983]

$$S_3 = 4\pi \int r^2 dr \left[\frac{1}{2} \left(\frac{d\phi}{dr} \right)^2 + V_{\text{eff}}(\phi, T) \right]$$



Conformal Two-Higgs-doublet model (C2HDM)

- Higgs basis:

$$\langle \mathcal{H}_1 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix}, \quad \langle \mathcal{H}_2 \rangle = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

- Potential in the Higgs basis:

$$V_{\text{C2HDM}}^{\text{Higgs}} = \frac{1}{2} Z_1 (\mathcal{H}_1^\dagger \mathcal{H}_1)^2 + \frac{1}{2} Z_2 (\mathcal{H}_2^\dagger \mathcal{H}_2)^2 + Z_3 (\mathcal{H}_1^\dagger \mathcal{H}_1) (\mathcal{H}_2^\dagger \mathcal{H}_2) + Z_4 (\mathcal{H}_1^\dagger \mathcal{H}_2) (\mathcal{H}_2^\dagger \mathcal{H}_1) \\ + \left[\frac{1}{2} Z_5 (\mathcal{H}_1^\dagger \mathcal{H}_2)^2 + [Z_6 (\mathcal{H}_1^\dagger \Phi_1) + Z_7 (\mathcal{H}_2^\dagger \mathcal{H}_2)] (\mathcal{H}_1^\dagger \mathcal{H}_2) + \text{h.c.} \right]$$

- Parameters: $Z_1, Z_2, Z_3, Z_4 \in \mathbb{R}$, $Z_5, Z_6, Z_7 \in \mathbb{C}$
- Stationary point: $Z_1 = Z_6 = 0$
- CP-violating angle γ : $\tan(2\gamma) = -\text{Im}(Z_5)/\text{Re}(Z_5)$
- Tree-level scalar masses:

$$m_{H_1}^2 = 0, \quad m_{H_2}^2 = \frac{v^2}{2} (Z_3 + Z_4 - |Z_5|), \quad m_{H_3}^2 = \frac{v^2}{2} (Z_3 + Z_4 + |Z_5|), \quad m_{H^\pm}^2 = \frac{v^2}{2} Z_3$$

- Parameter set for scan:

$$Z_2, \text{Im}(Z_5), \text{Re}(Z_7), \text{Im}(Z_7), m_{H_2}, m_{H_3} \quad (\beta, \delta)$$

- NOTE: We eliminate $m_{H^\pm}$ for m_h using the GW-approach.

The Gildener-Weinberg approach

- Generic model with no mass term: [E. Gildener et al., 1976]

$$V = f_{ijkl} \Phi_i \Phi_j \Phi_k \Phi_l$$

- Existence of non-trivial local minimum $\Rightarrow$ flat direction $\mathbf{n}$
- One-loop potential along the flat direction:

$$V_{\text{eff}}(\varphi \mathbf{n}) = A\varphi^4 + B\varphi^4 \ln(\varphi^2/\Lambda_W^2)$$

- The excitation along $\mathbf{n}$ is called the *scalon* $\rightarrow h$
- One-loop scalon mass (approximate):

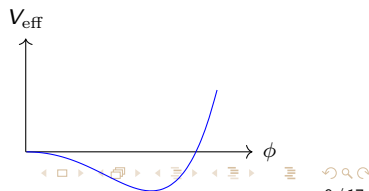
$$m_h^2 = 8Bv^2 = \frac{1}{8\pi^2 v^2} (6m_W^4 + 3m_Z^4 + \sum_H m_H^4 - 4\sum_F m_F^4)$$

- NOTE: True vev may differ:

$$|v^{\text{true}} - v| \leq 5 \text{ GeV}$$

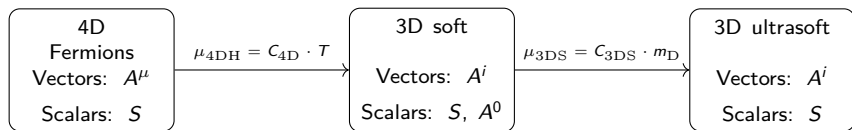
- NOTE: True m_h may differ:

$$|m_h^{\text{true}} - m_h| \leq 2 \text{ GeV}$$



Effective potential and phase transitions

- We use dimensional reduction to construct $V_{\text{eff}}(\phi, T)$.
- Basic idea: Match the 4D theory to a 3D EFT by integrating out the non-zero Matsubara modes: $m^2 + (n\pi T)^2$, $n \neq 0$.
- Advantage over 4D: reduction of theoretical uncertainty
- Dratopi: Constructs the 3D EFT using DRalgo
[A. Ekstedt et al., 2023]
- PhatCat: Phase tracing, bounce action and PT parameters
(currently using CosmoTransitions [C. L. Wainwright, 2012])

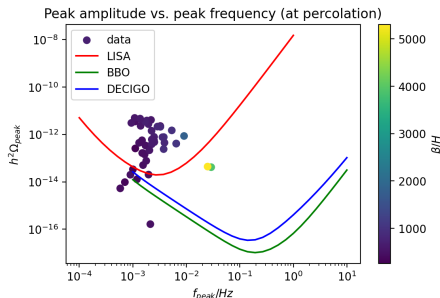
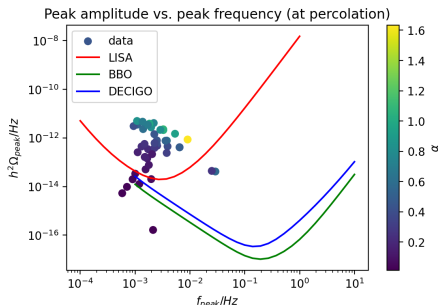


Analysis pipeline

- Scan the parameter space:
 $Z_2, \text{Im}(Z_5), \text{Re}(Z_7), \text{Im}(Z_7), m_{H_2}, m_{H_3} (\beta, \delta)$
- Apply theoretical constraints:
 - Perturbative unitarity
 - Boundedness from below (be careful for conformal models!)
- Apply phenomenological constraints:
 - Oblique parameters S, T, U [Peskin, M. and Takeuchi, T., 1990]
 - Collider constraints: HiggsTools [H. Bahl et al., 2023]
- Construct the effective potential
- Study the thermal history and identify the PTs
- Calculate the PT parameters: $\alpha, \beta/H, T_p$
- Calculate the contribution to the SGW background

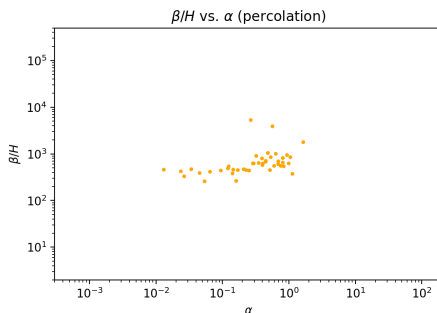
Preliminary results

- GW spectrum characterized by peak amplitude $h^2\Omega_{\text{peak}}$ and peak frequency f_{peak}
- Point above sensitivity curve for detector $\Rightarrow \text{SNR} > 1$
- $\alpha \uparrow \Rightarrow h^2\Omega_{\text{peak}} \uparrow$
- $\beta/H \uparrow \Rightarrow f_{\text{peak}} \uparrow$

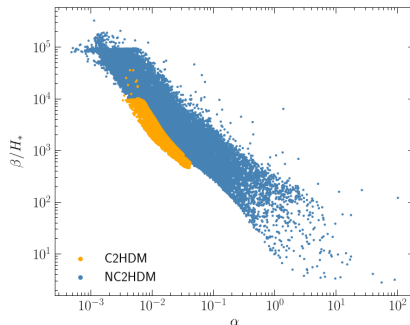


Comparison to previous work

- Comparison to previous work on C2HDM and NC2HDM
[N. Benincasa et al., 2026]
- We find PTs with much larger α !



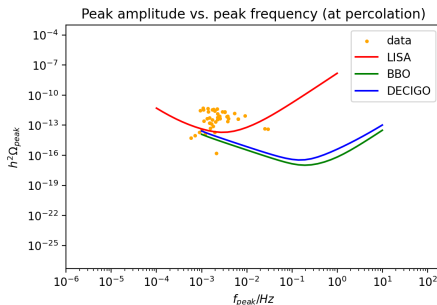
(a) Our results



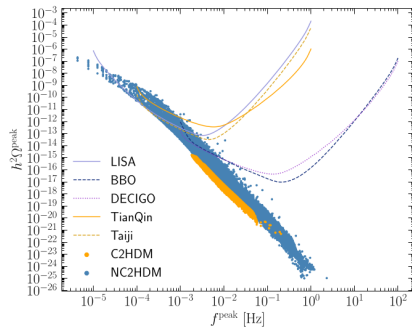
(b) [N. Benincasa et al., 2026]

Comparison to previous work

- Comparison to previous work on C2HDM and NC2HDM
[N. Benincasa et al., 2026]
- We find PTs with much larger $h^2\Omega_{\text{peak}}$!



(a) Our results



(b) [N. Benincasa et al., 2026]

Conclusion

- BSM models with extended scalar sectors are well-motivated e.g. by the baryon asymmetry problem.
- Conformal models are motivated e.g. by the hierarchy problem.
- The C2HDM is a conformal model with additional CP-violation (on top of the SM).
- The results are preliminary but clearly feature first-order phase transitions.
- These phase transitions are stronger than expected from previous work.
- Possible explanations for differences?
 - Construction of effective potential: 4D vs. 3D approach
 - Gildener-Weinberg vs. on-shell scheme
 - CP-violation?
- Further investigations are needed to understand this!

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