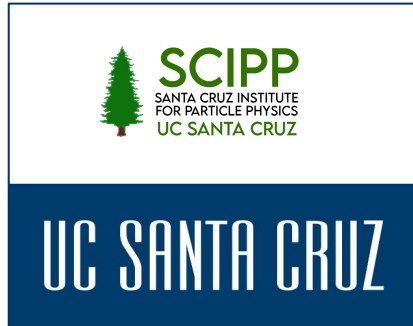


# Fitting Current from InP:Fe Under High Flux Xray Test Beam



Earl Russell Almazan

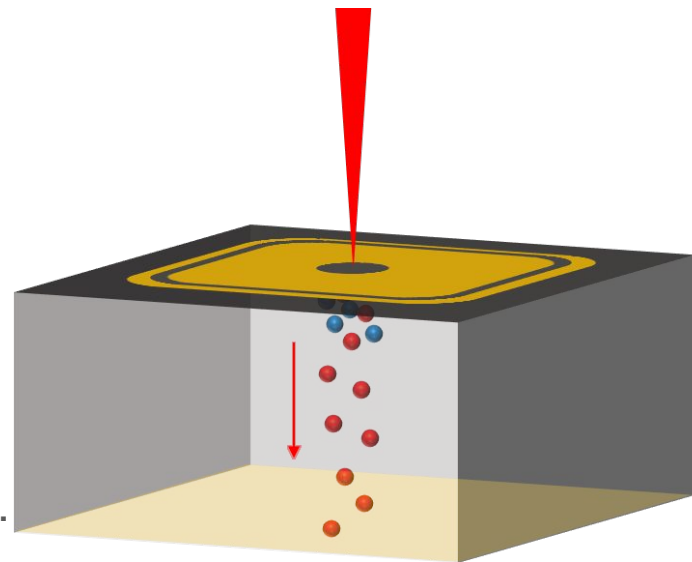
Jason Nielsen, Mike Hance,  
Anirudha Sumant, Anthony Affolder, Dennis Sperlich,  
Ian Dyckes, Jennifer Ott, Jessica Metcalfe, Luise Poley, Manoj  
Jadhav, Sungjoon Kim, Thomas Mccoy, Vitaliy Fadeyev

# Overview

- Xray beam test measurements from DIAMOND Light Source to investigate persistent current phenomenon
- **Persistent Current**
  - minutes-long decay of current after high flux of 15 keV xrays is removed
- In-Progress:
  - **Building a Theoretical Model**
    - What is persistent current? What causes it?
  - **Model Implications**
    - Framework → Predictive Model → Testable Hypotheses and Fits
  - **Data Processing**
    - Fitting model to measured data
    - Extracting parameters
    - Comparing fits to theoretical predictions

# Experimental Setup - DLS

- **Focused xray beam of 15 keV photons**
  - Absorption length  $\sim 100 \mu\text{m}$
  - Carriers generated most near frontside
  - Beam moved on and off device
- **Device biased on backside**
  - Carriers travel through bulk
- **Different Initial Conditions**
  - Temperature, photon flux, applied voltage, etc.
- Records: [Logbook](#) and [Spreadsheet](#)
  - [Plots](#) and [Data](#)
  - Data taken every 2 seconds



Schematic by Jennifer Ott

# Initial Measurements

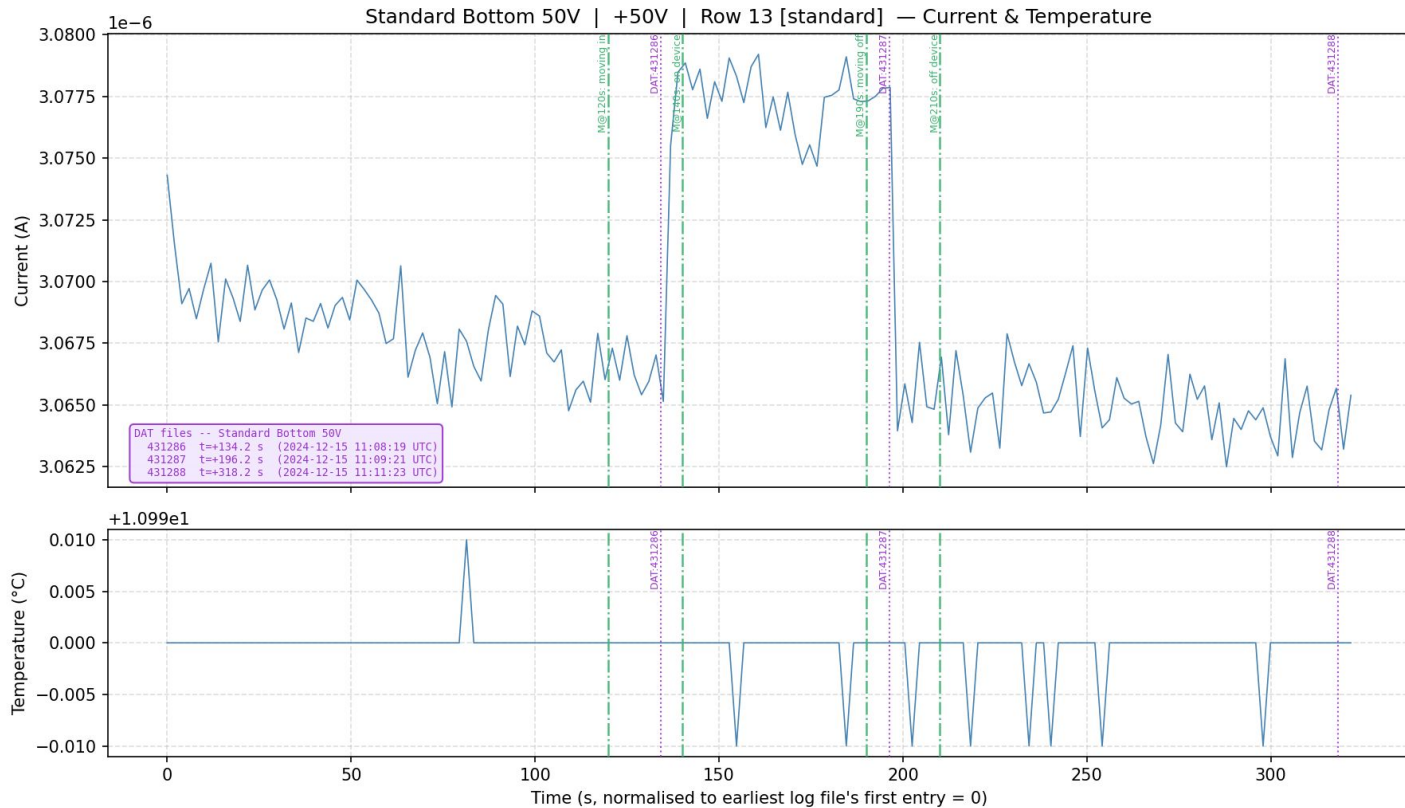
Timescale  
O(5 min)

Baseline  
photocurrent

Stability of  
temperature

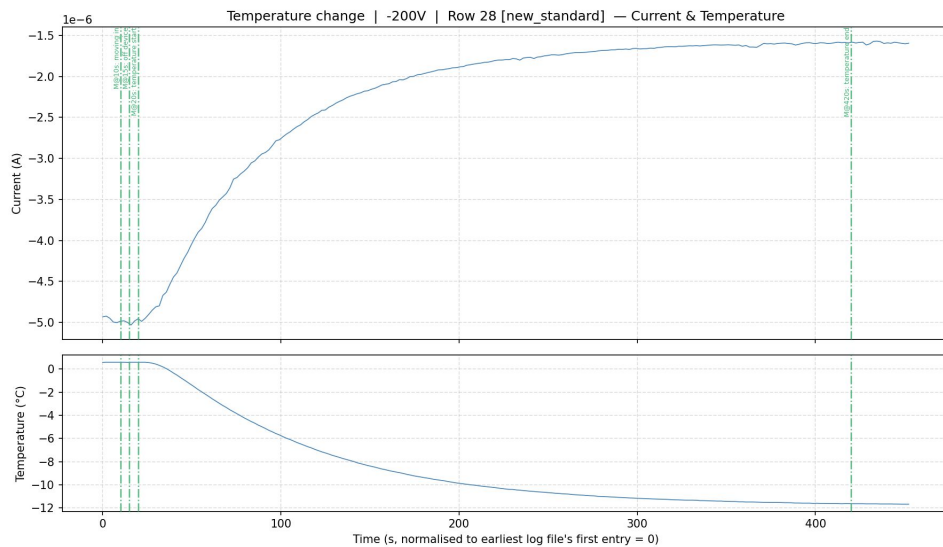
Locate active  
area/device edges

Mostly irrelevant  
for later slides



# Long Scans - Temperature - 1

- Record dark current as a function of temperature
- Get Arrhenius plot and derive activation energy
  - Bandgap or deep Fe traps?



# Long Scans - Temperature - 2

5. Arrhenius (temperature plot, `temperature_arrhenius_fit_parameters.csv`)

$$|I(T)| = A \cdot \exp\left(-\frac{E_a}{k_B T_K}\right) \quad \text{with } T_K = T_C + 273.15$$

Fit performed by linear regression on  $\ln |I|$  vs  $1/T_K$ :

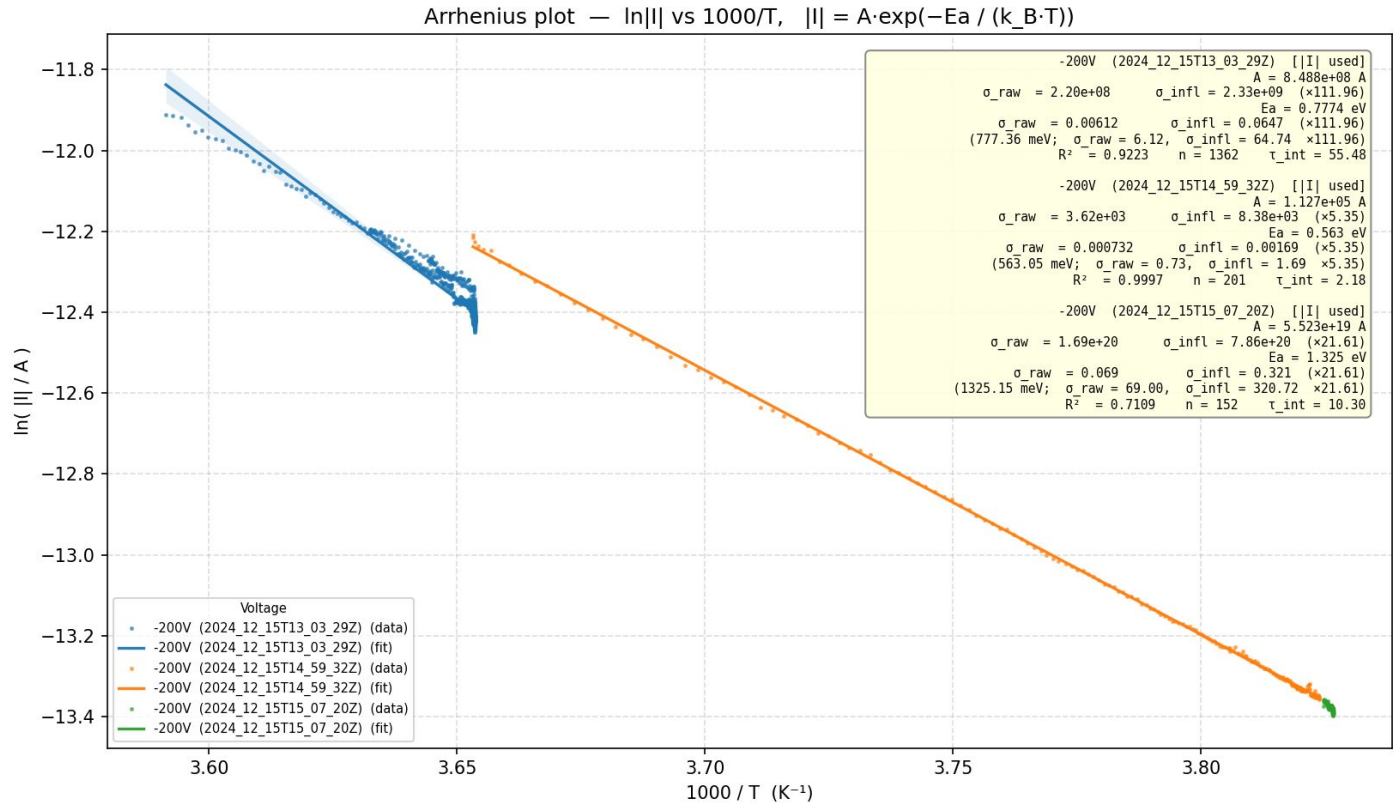
$$\ln |I| = \ln A - \frac{E_a}{k_B} \cdot \frac{1}{T_K}$$

so slope =  $-E_a/k_B$ , intercept =  $\ln A$ .  $E_a$  reported in eV using  $k_B = 8.617333262 \times 10^{-5}$  eV/K.

# Long Scans - Temperature - 3

Remap measured current + temp as functions of time into current vs temp, then fit to Arrhenius

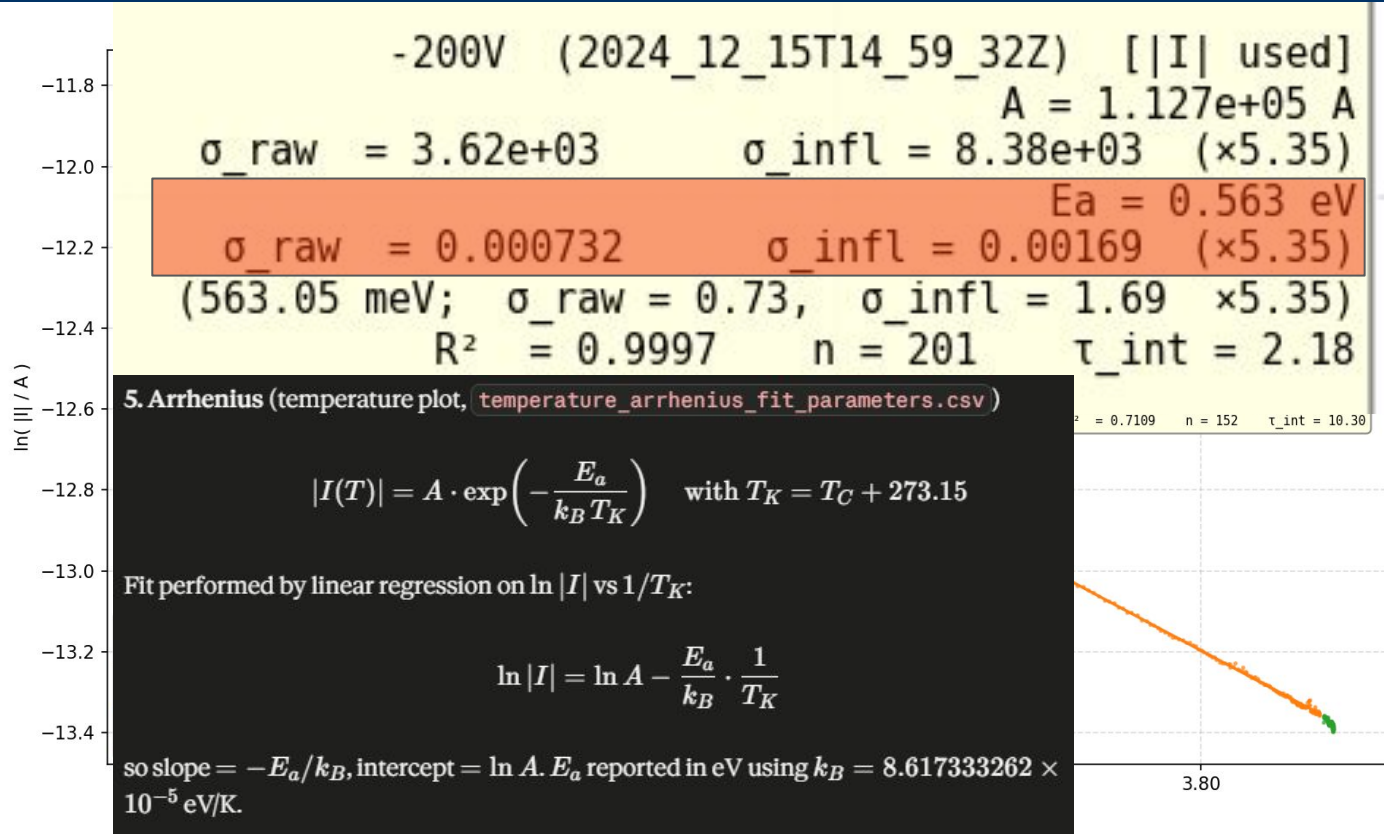
Assume dark current from thermally-driven carrier exchange between energy levels



# Long Scans - Temperature - 4

Activation energy  
 ~trap-to-band gap  
 → Dark current is driven by carrier exchange with the Fe traps, not bandgap

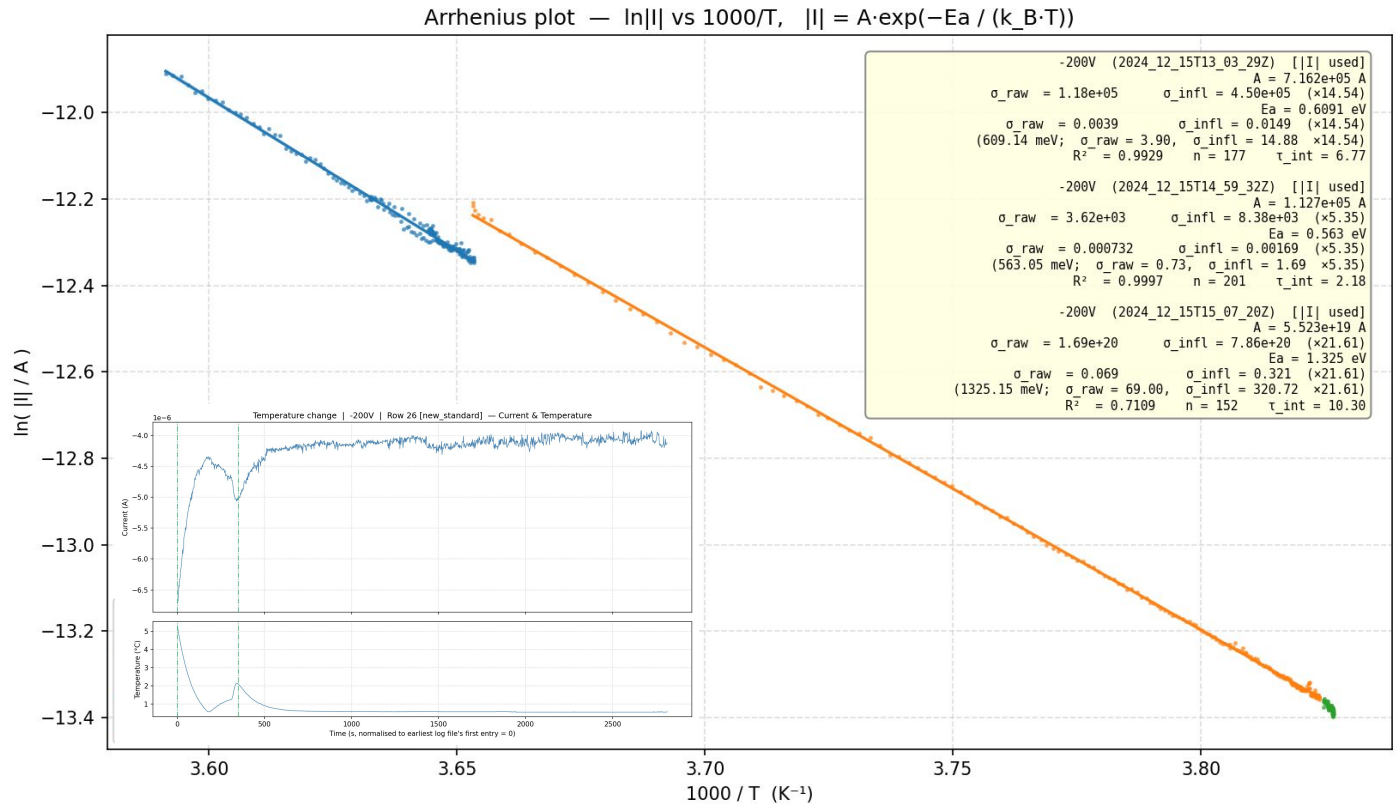
Low uncertainty from scipy fit



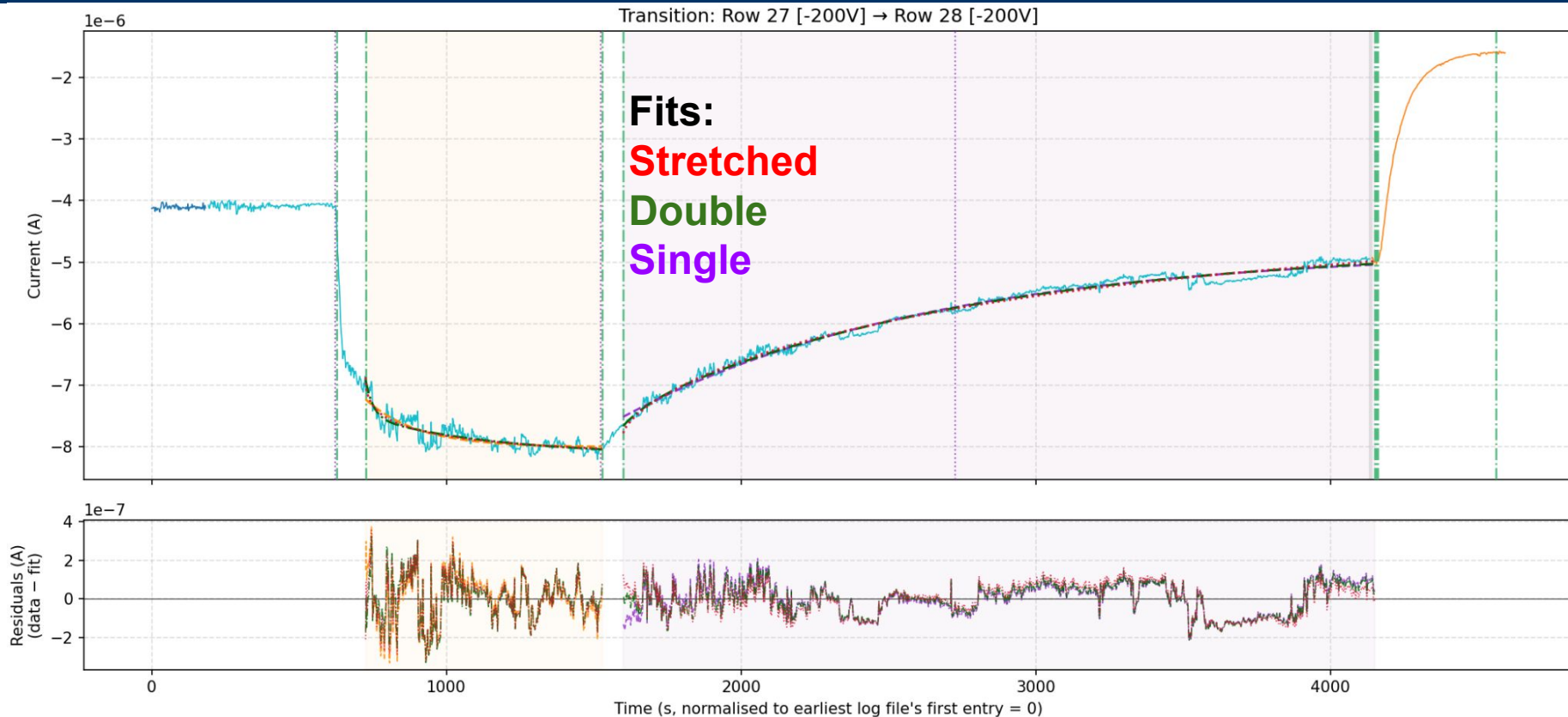
# Long Scans - Temperature - 5

Removing tail end of blue data / large cluster at high current leads to better matching  $E_a$  between runs.

Later tail likely contains other processes not related to thermal exchange of carriers



# Long Scans - Persistent Current - 1



# Long Scans - Persistent Current - 2

1. Single exponential (recovery curve, time domain) — `fit_parameters_single.csv`

$$I(t) = C + A \cdot e^{-kt}$$

with derived quantities  $\tau = 1/k$  and  $Q = A/k$  (charge-equivalent integral).

2. Double exponential (recovery curve, time domain) — `fit_parameters_double.csv`

$$I(t) = C + A_1 \cdot e^{-k_1 t} + A_2 \cdot e^{-k_2 t}$$

with  $\tau_1 = 1/k_1$ ,  $\tau_2 = 1/k_2$ , and  $Q = A_1/k_1 + A_2/k_2$ . Convention in this codebase: subscript 1 is the slow component (smaller  $k$ ), subscript 2 is the fast component.

3. Stretched exponential (recovery curve, time domain) —

`fit_parameters_stretched.csv`

$$I(t) = C + A \cdot e^{-(kt)^\beta}$$

with  $0 < \beta \leq 1$ . The effective time constant reported is  $\tau = 1/k$ , and  $Q = (A/k) \cdot \Gamma(1 + 1/\beta)$  (the proper integral of a stretched exponential).

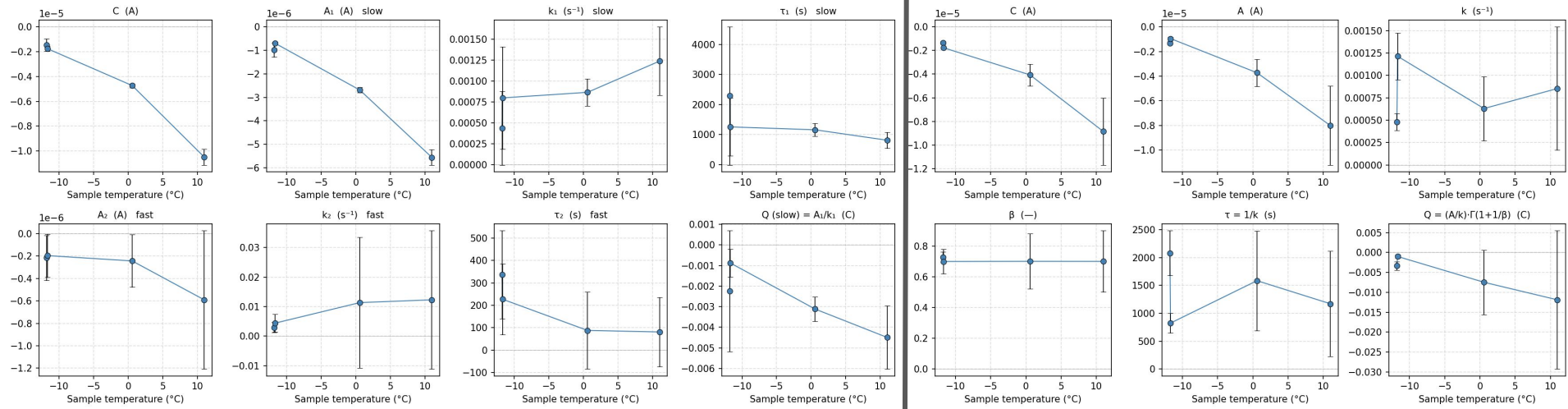
# Long Scans - Temperature Sweep

Constant voltage (-200 V) and beam intensity

Overall trend: characteristic time falls with higher temperature

Double

Stretched



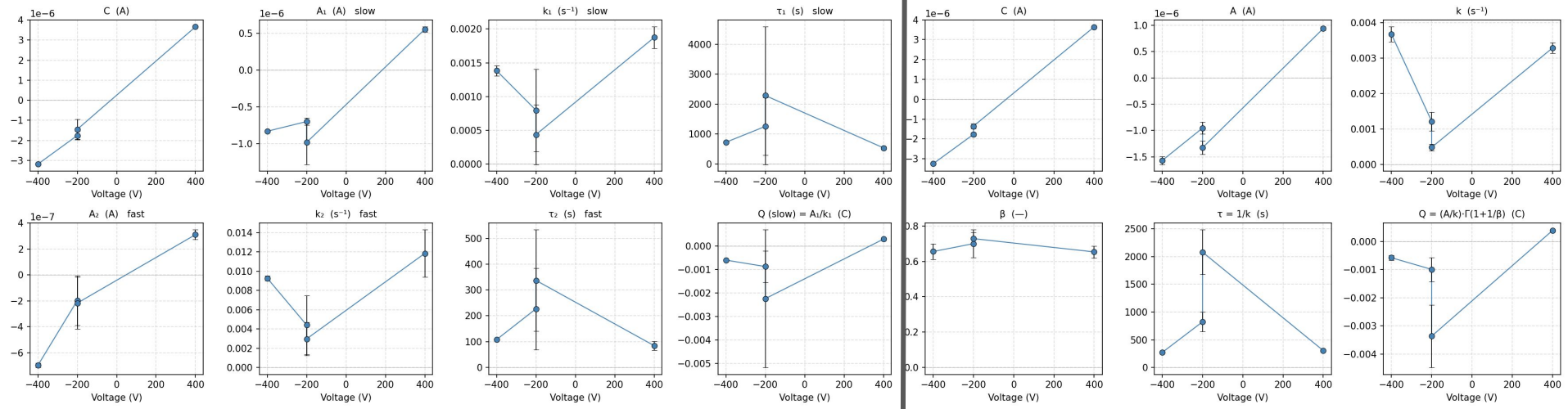
# Long Scans - Voltage Sweep

Constant temperature (-11C) and beam intensity

Overall trend: characteristic time falls with higher voltage magnitude

Double

Stretched



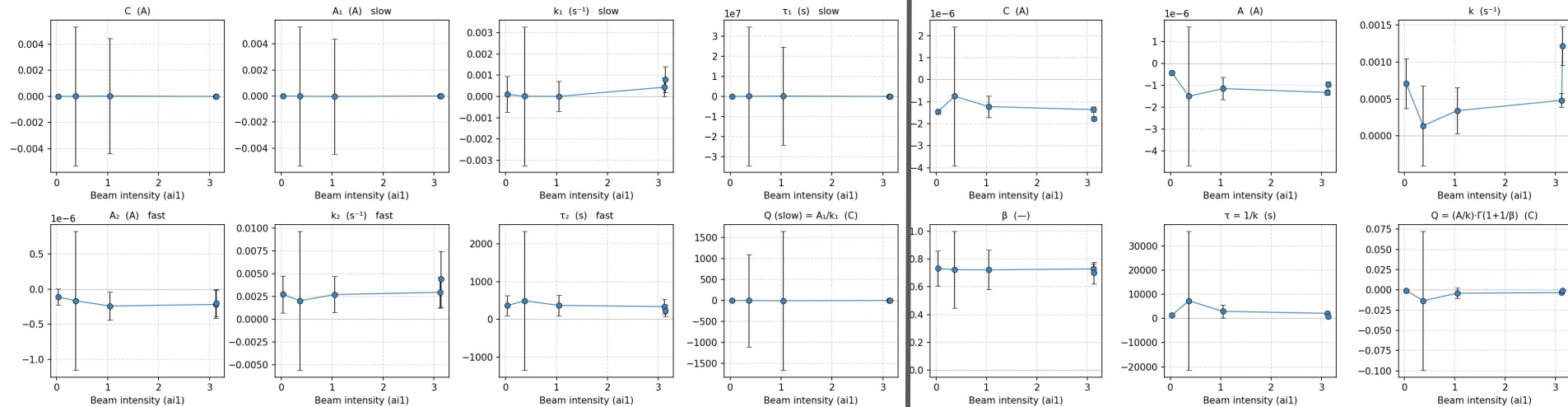
# Long Scans - Intensity Sweep

Constant temperature (-11C) and voltage (-200 V)

Overall trend: Unclear, high uncertainties

Double

Stretched



# Steps Going Forward

- Full set of plots:  
[https://drive.google.com/drive/u/0/folders/1gjqFTu\\_qBRPtUzmwMaO70nDM6\\_Z8W7IO](https://drive.google.com/drive/u/0/folders/1gjqFTu_qBRPtUzmwMaO70nDM6_Z8W7IO)
- Derive conclusions and overview potential theoretical models
  - Need something empirical for deriving theoretical leakage current, especially given Arrhenius plot measurements
- Connect extracted fit parameters to theoretical model

# Backup



# Dark Equilibrium Conditions

- Known Information:

- Trap density:  $1\text{E}16\text{ cm}^{-3}$  (manufacturer)
- Undoped InP carrier density:  $2\text{-}4\text{E}15\text{ cm}^{-3}$  ([literature](#))
- InP bandgap:  $1.34\text{ eV at }300\text{K}$  ([literature](#))  
 $1.41\text{ eV at }77\text{K}$
- Trap State:  $0.6\text{-}0.7\text{ eV below CB}$  (Sungjoon)
- Assumption: Most electrons in conduction band fall into the Fe traps (Sungjoon)
  - Conduction band almost unpopulated
  - InP:Fe is a semi-insulator - confirmed with measurements

# Solid State Calculations

- Where is the Fermi level for this system?
- Using Fermi-Dirac Distribution and assumptions about InP:Fe...

$$f = \frac{1}{1 + e^{\frac{E - E_F}{kT}}}$$

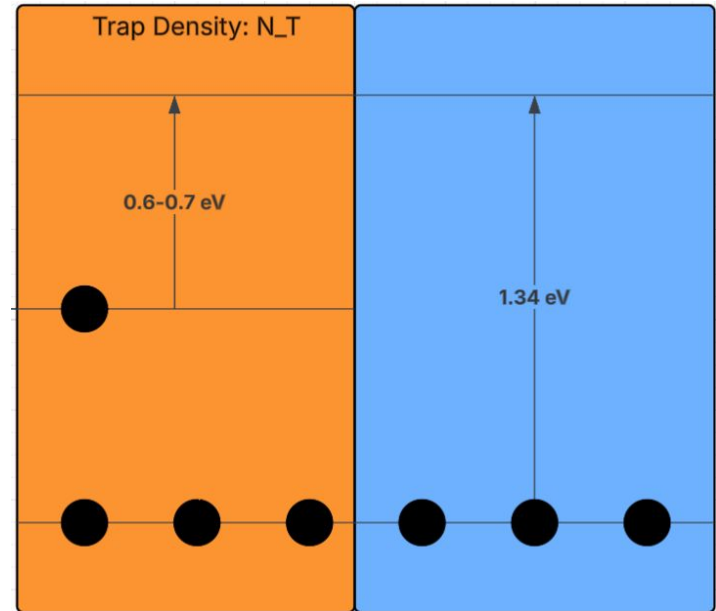
$$n_{total} = N_D^+$$

$$n_{total} = n_{trapped} \text{ at } 300\text{K}$$

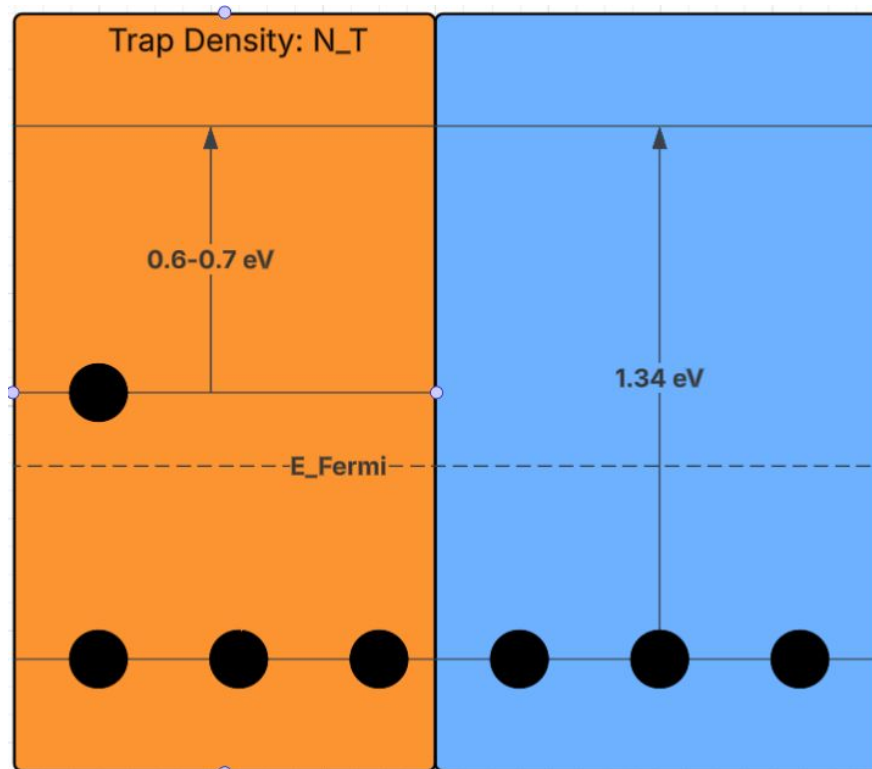
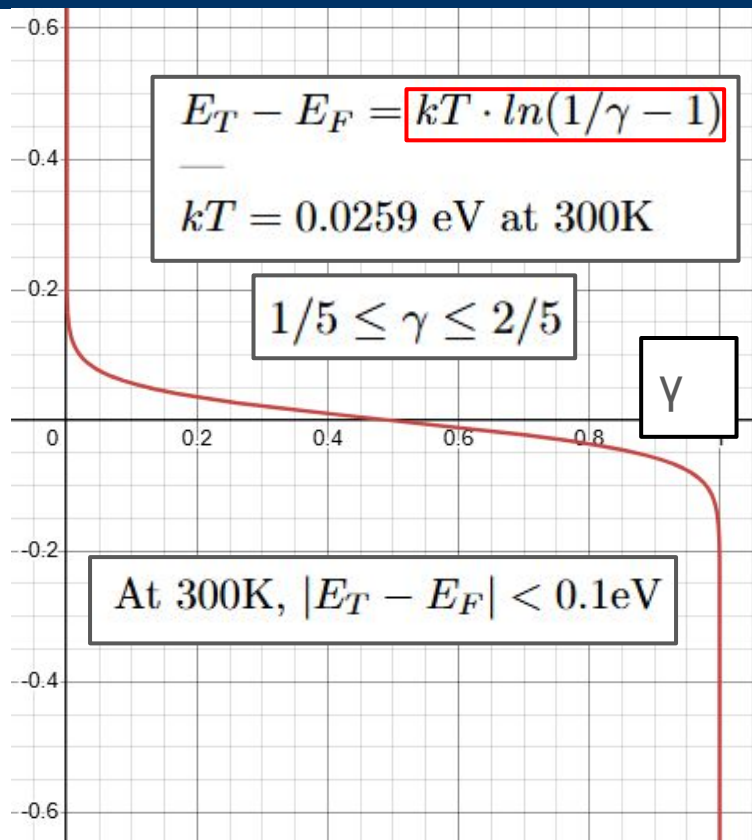
$$E = E_T$$

$$\gamma = N_D^+ / N_T = n_{trapped} / N_T = f$$

$$E_F = E_T - kT \cdot \ln(1/\gamma - 1)$$



# Solid State Calculations



# Solid State Calculations

$$E_T - E_F = kT \cdot \ln(1/\gamma - 1)$$


---


$$kT = 0.0259 \text{ eV at } 300\text{K}$$

$$1/5 \leq \gamma \leq 2/5$$

$$g = \frac{1}{5} \quad \times$$

= 0.2

---


$$E_{30} = \frac{0.0259(30 + 273.15)}{300} \ln\left(\frac{1}{g} - 1\right) \quad \times$$

= 0.0362820267045

---


$$E_{m15} = \frac{0.0259(-15 + 273.15)}{300} \ln\left(\frac{1}{g} - 1\right) \quad \times$$

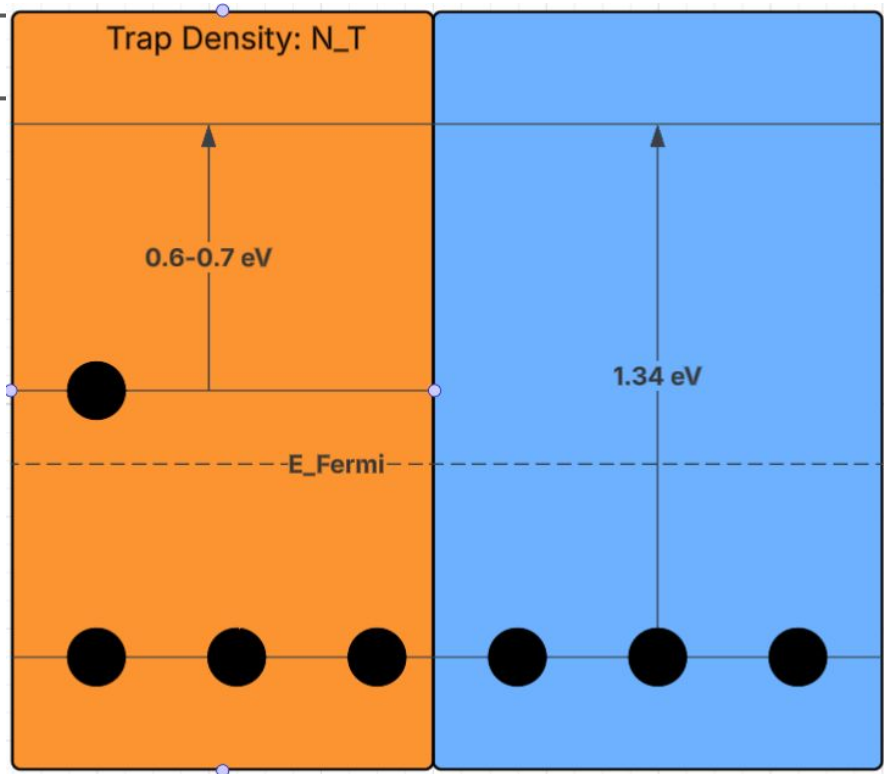
= 0.0308962731116

---


$$E_{30} - E_{m15} \quad \times$$

= 0.00538575359295

30C  
-15C



$\Delta E_F$  from temp change =  $\sim 0.005 \text{ eV}$

# Theoretical Model - Assumptions

- Given TCT measurements and the suspicion that the low amplitude from the hole response is due to fast thermal hole trapping, what occupies traps may be holes
- Assumption: in the illuminated steady state, with the higher carrier density, the traps naturally get a higher density of what they trap most easily: holes
- Upon the removal of illumination, free carriers above the equilibrium quickly evacuate to electrodes within ms from drift.
- What remains for several seconds are a conduction band and valence band mostly at dark equilibrium, but the trap states remain filled and, by thermal processes, release their trapped holes.

# Theoretical Model - Assumptions

- But, what happens to the trapped hole?

Under illumination both processes compete. The relative rates are:

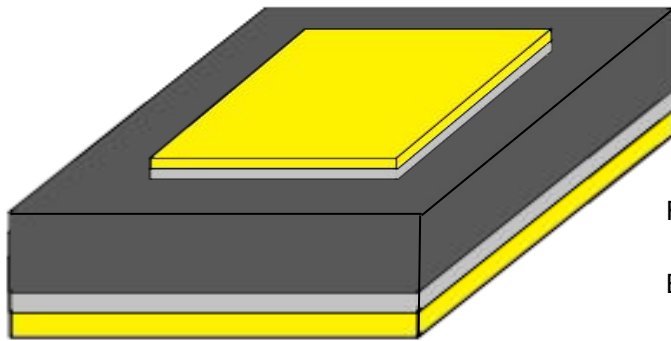
$$\frac{\text{recombination rate}}{\text{re-emission rate}} = \frac{\beta_n n f_T}{e_p f_T} = \frac{\beta_n n}{e_p}$$

- High illumination = higher number of electrons excited to conduction band from impact ionization (main mechanism due to 15 keV photons)
  - Recombination dominates
- No illumination = n population goes to near zero = re-emission more likely
- Trap state filled with more electrons?

# Theoretical Model

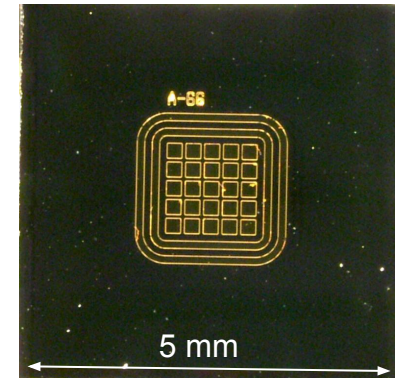
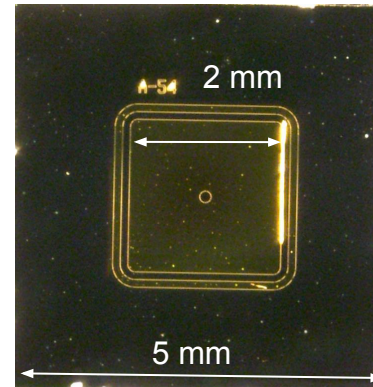
- From TCT measurements, holes do not traverse the bulk fully compared to electrons
  - Assumption 1: Could be explained by fast hole trapping by Fe traps. Assume hole trapping rate  $>$  hole emission
    - Implication: holes release could be responsible for persistent current
- Strategy:
  - Identify approximate steady states with and without illumination
  - Identify mechanisms for transitioning from with to without illumination and their approximate timescales
    - Calculate roughly the timescale for this transition
    - Dependence on this model on measurable/known quantities
      - Bandgap, trap energy, total integrated charge, characteristic time for persistent current, etc.

# Fabricated Device Structure



Gold (100 nm)  
 Chromium (10 nm)  
 InP:Fe (350  $\mu\text{m}$ )  
 Chromium (10nm)  
 Gold (100 nm)

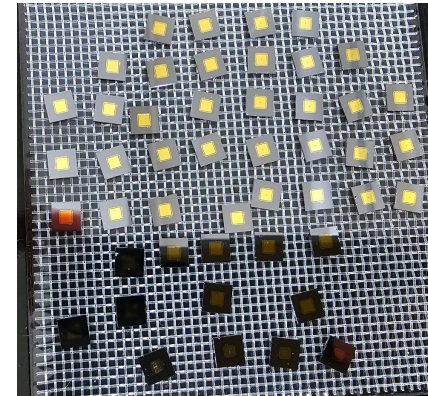
Frontside: e-beam  
 metal deposition  
 Backside: sputter



Argonne  
 NATIONAL LABORATORY

- 2nd iteration of devices
- Check properties with Crystalline InP
  - Commercially Available
  - Test properties with crystalline wafers as a benchmark for thin film devices, whether they be amorphous or crystalline

Device Size	=	5 x 5 mm <sup>2</sup>
Pads	=	2 x 2 mm <sup>2</sup>
Guard Rings	=	100 $\mu\text{m}$
GR Spacing	=	100 $\mu\text{m}$
Hole Diameter	=	150 $\mu\text{m}$
Multipad Pads	=	200 x 200 $\mu\text{m}^2$



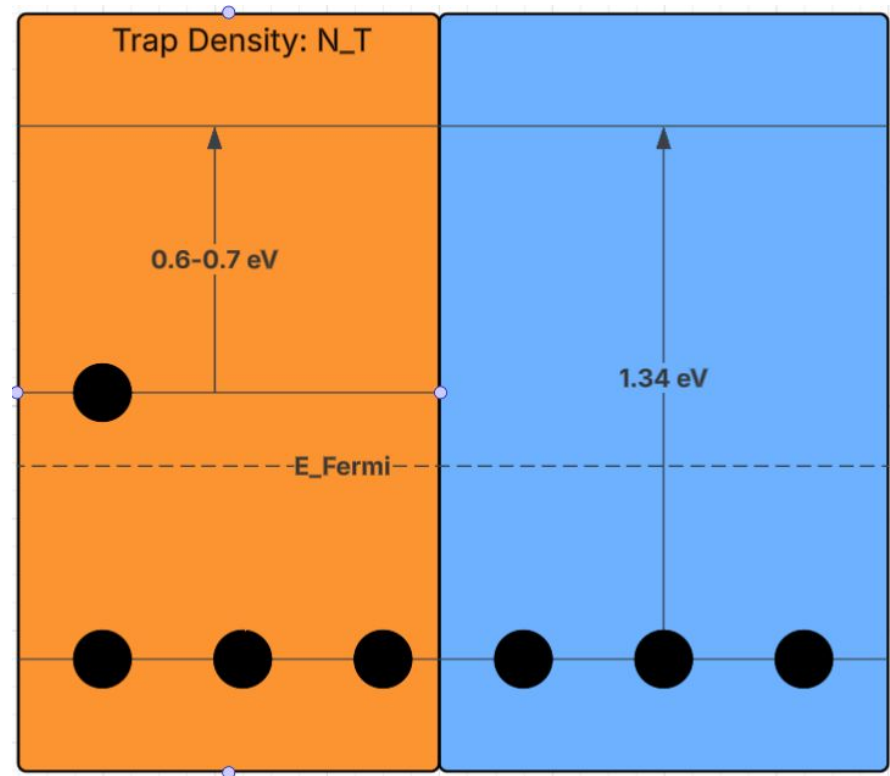
# X-ray Parameters

**Table 5: Beam Line Properties**

Facility	Beam Energy [keV]	Beam Diameter [ $\mu\text{m}$ ]	Operational Temperature [C]	Beam Flux [ph/s]	Sample support
CLS	15	40	2.5	$1.25 \times 10^{11}$	Thick PCB
DLS	15	2	-15	$1.18 \times 10^8$	Thin PCB with enhanced thermal control

# Theoretical Model - Assumptions

- Let's try to intuit the change in carrier/trap-filling distributions after illumination
- While not under illumination and without iron doping, the electron and hole populations are dictated by the intrinsic carrier density  $n_0 p_0 = n_i^2$
- Iron dopants add a trap level, with trap distribution  $N_T$  and fraction of traps filled electrons ( $1-f_T$ ), giving a conservation of carriers equation:



# Theoretical Model - Assumptions - Illumination

- Let's assume majority of carrier generation happens at the bandgap, with this resulting in a split in the fermi-energies for the holes and electrons as their populations rise.
- At illumination steady state,  $n$  is high → recombination dominates as the main way trapped holes

Under illumination both processes compete. The relative rates are:

$$\frac{\text{recombination rate}}{\text{re-emission rate}} = \frac{\beta_n n f_T}{e_p f_T} = \frac{\beta_n n}{e_p}$$

# Theoretical Model - Assumptions - Relaxation

- Now, with no illumination, the electron population goes down → recombination and re-emission are no longer in equilibrium → re-emission dominates as the main method for removing holes in traps

Under illumination both processes compete. The relative rates are:

$$\frac{\text{recombination rate}}{\text{re-emission rate}} = \frac{\beta_n n f_T}{e_p f_T} = \frac{\beta_n n}{e_p}$$

# Theoretical Model - Assumptions

- Law of mass action = under equilibrium,

