



Silicon Zoology

Daniel Hynds

Outline

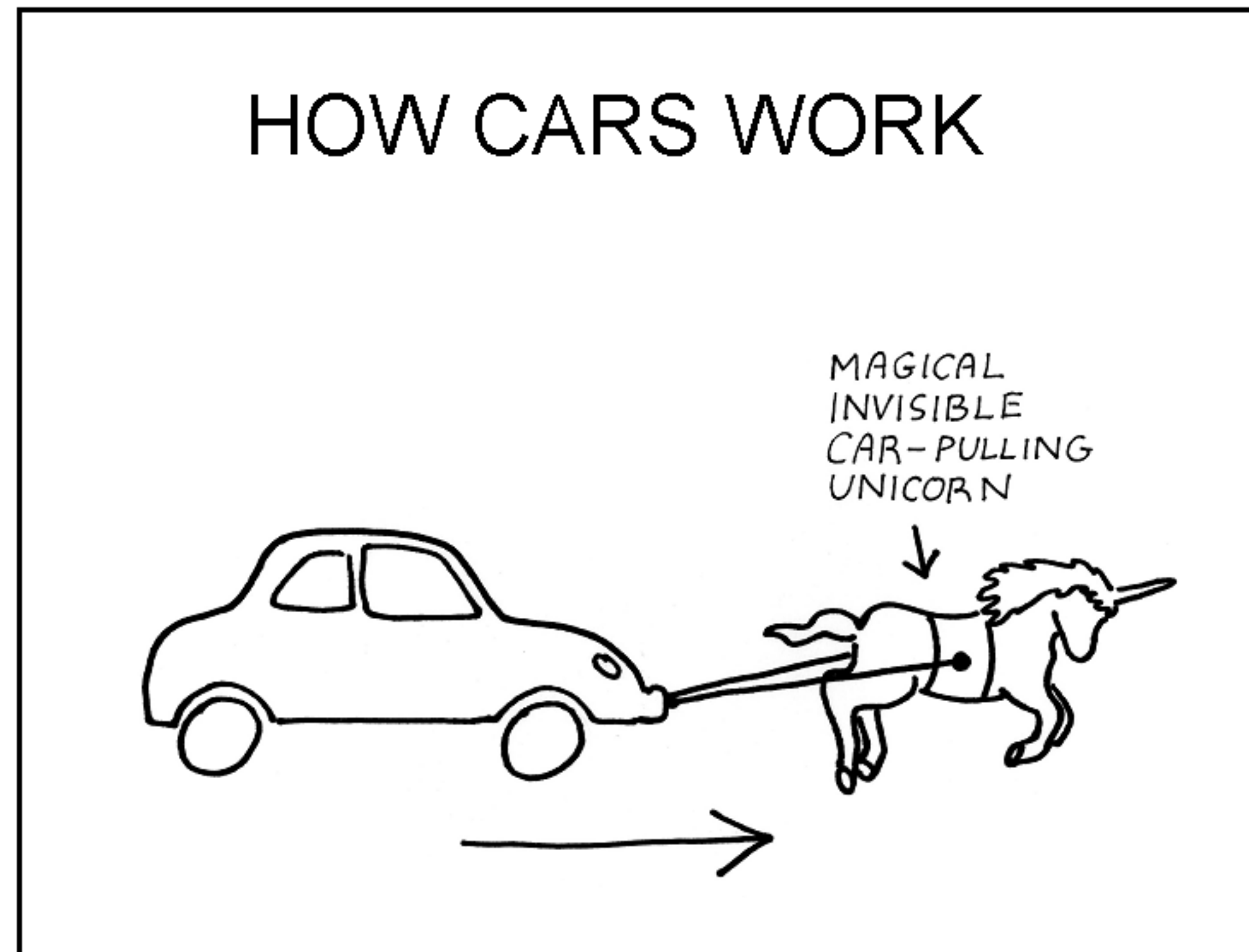
- Electronic structure of solids
- Semiconductor junctions
- Transistor anatomy
- Detector anatomy
- Signal formation
- Charge carrier generation
- Radiation damage
- Semiconductor fabrication
- Silicon pixel detectors

How stuff works

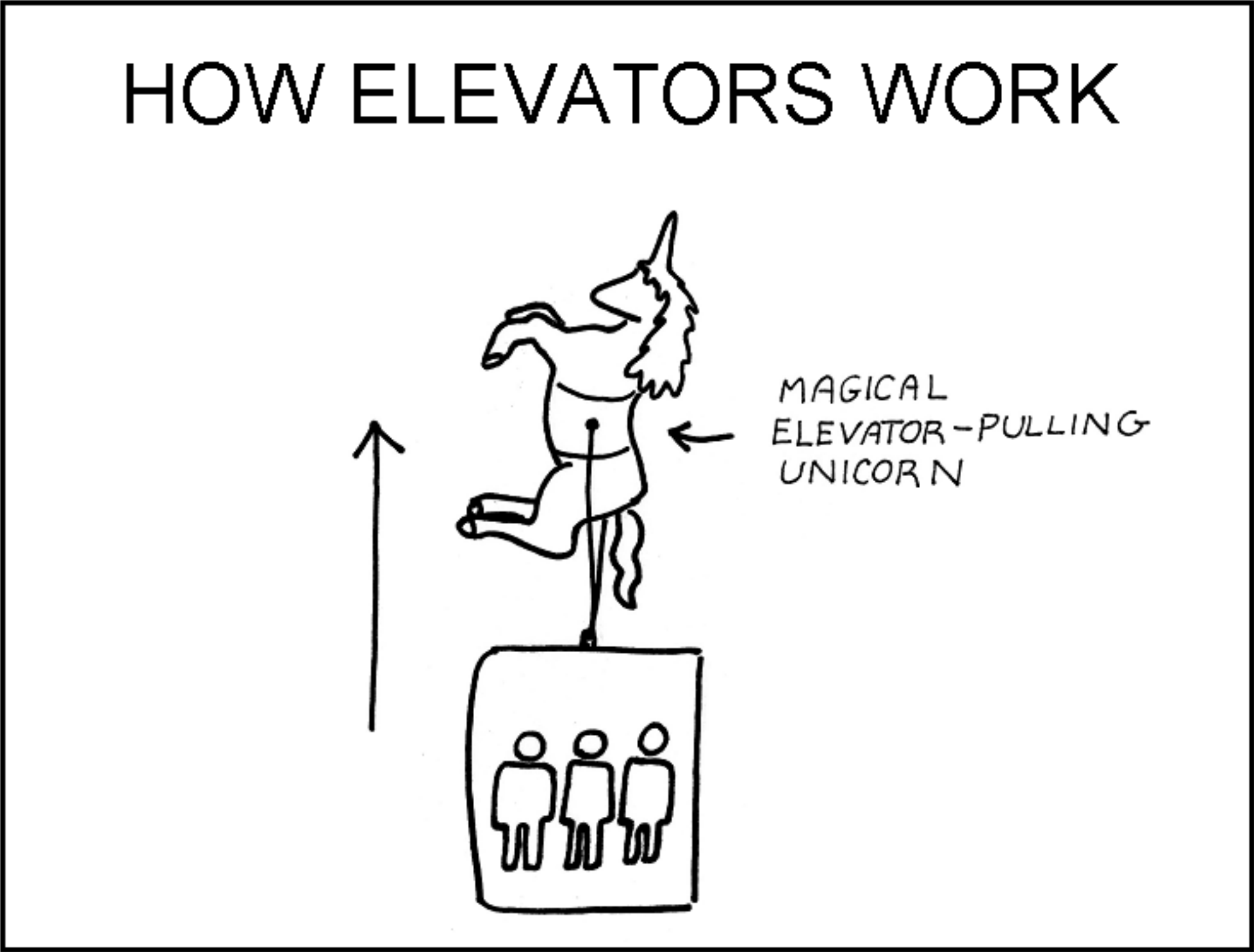
When I was a young boy
and I didn't understand how
something worked, I just
made stuff up in my head.



How stuff works



How stuff works



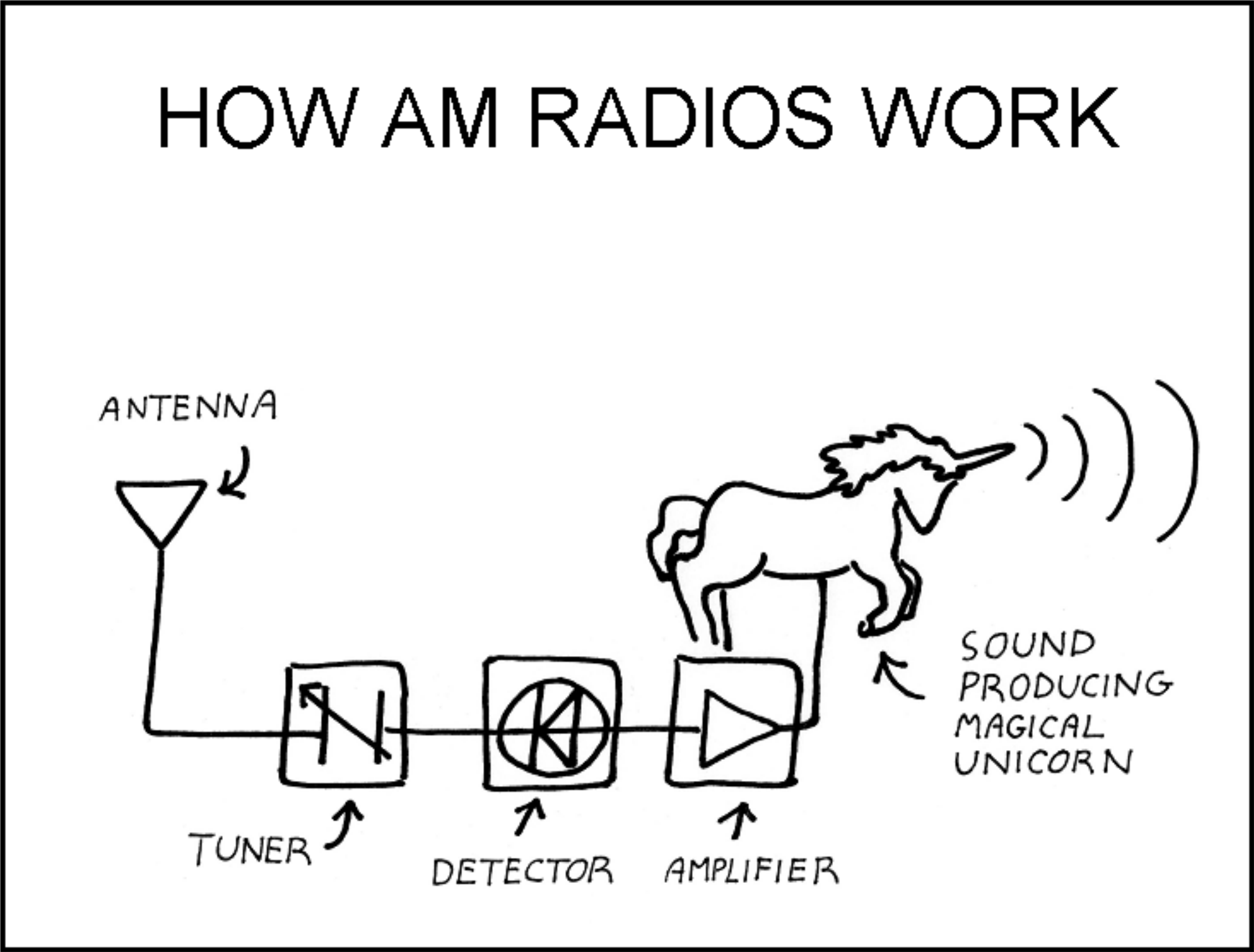
How stuff works

As I got older, I started to learn that science could explain how most things worked.

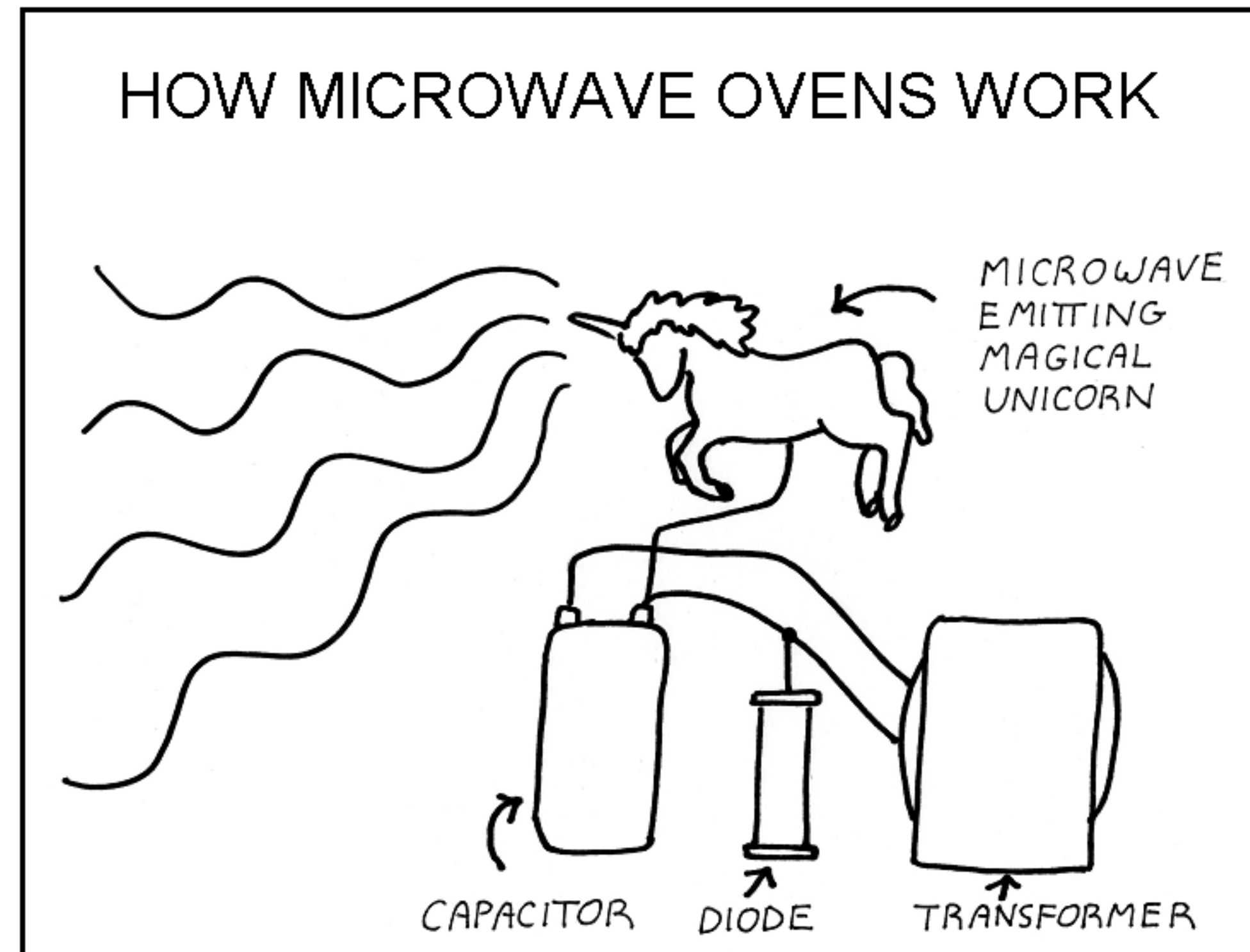
However, when there were holes in my knowledge, my mind still just filled in the gaps.



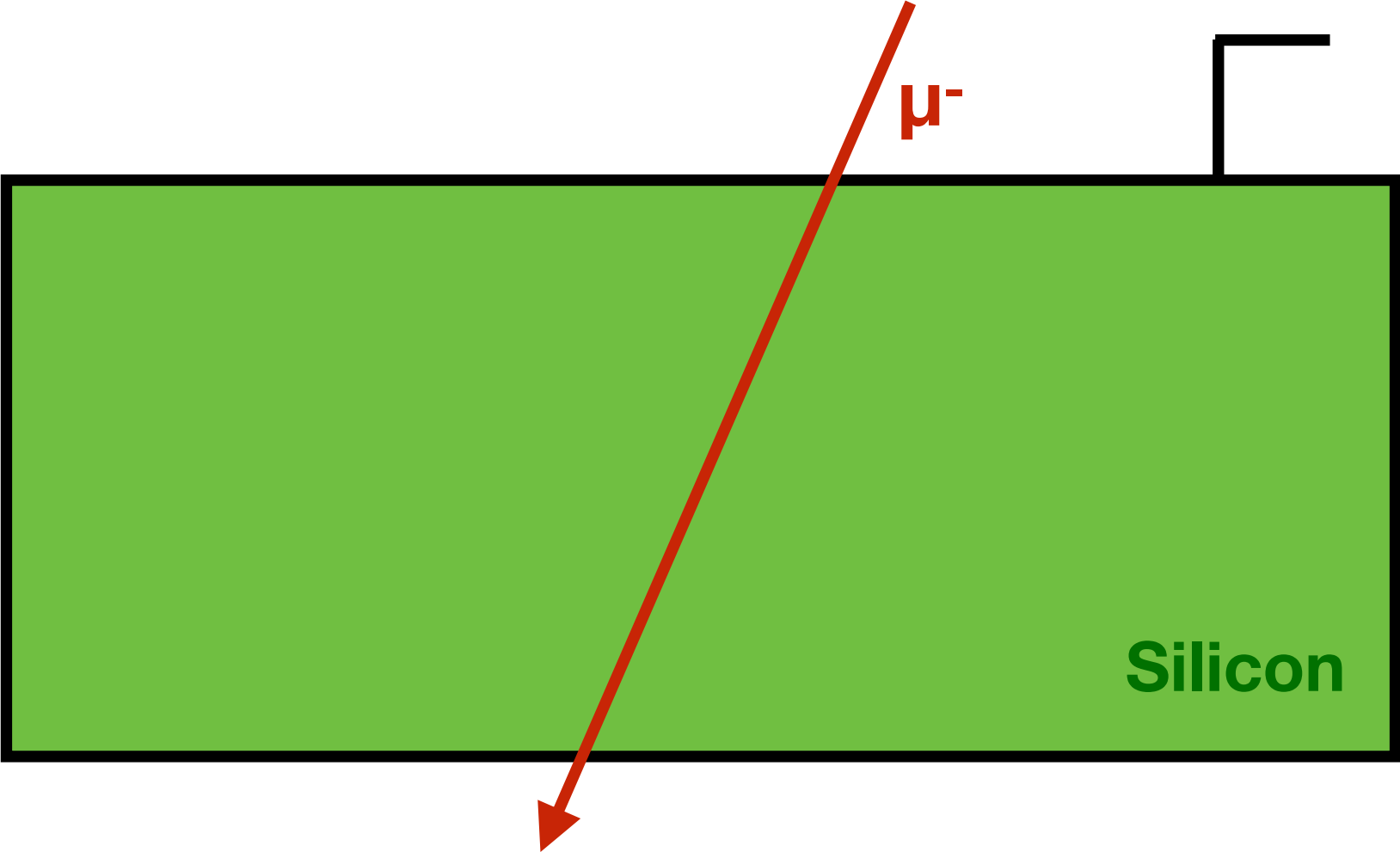
How stuff works



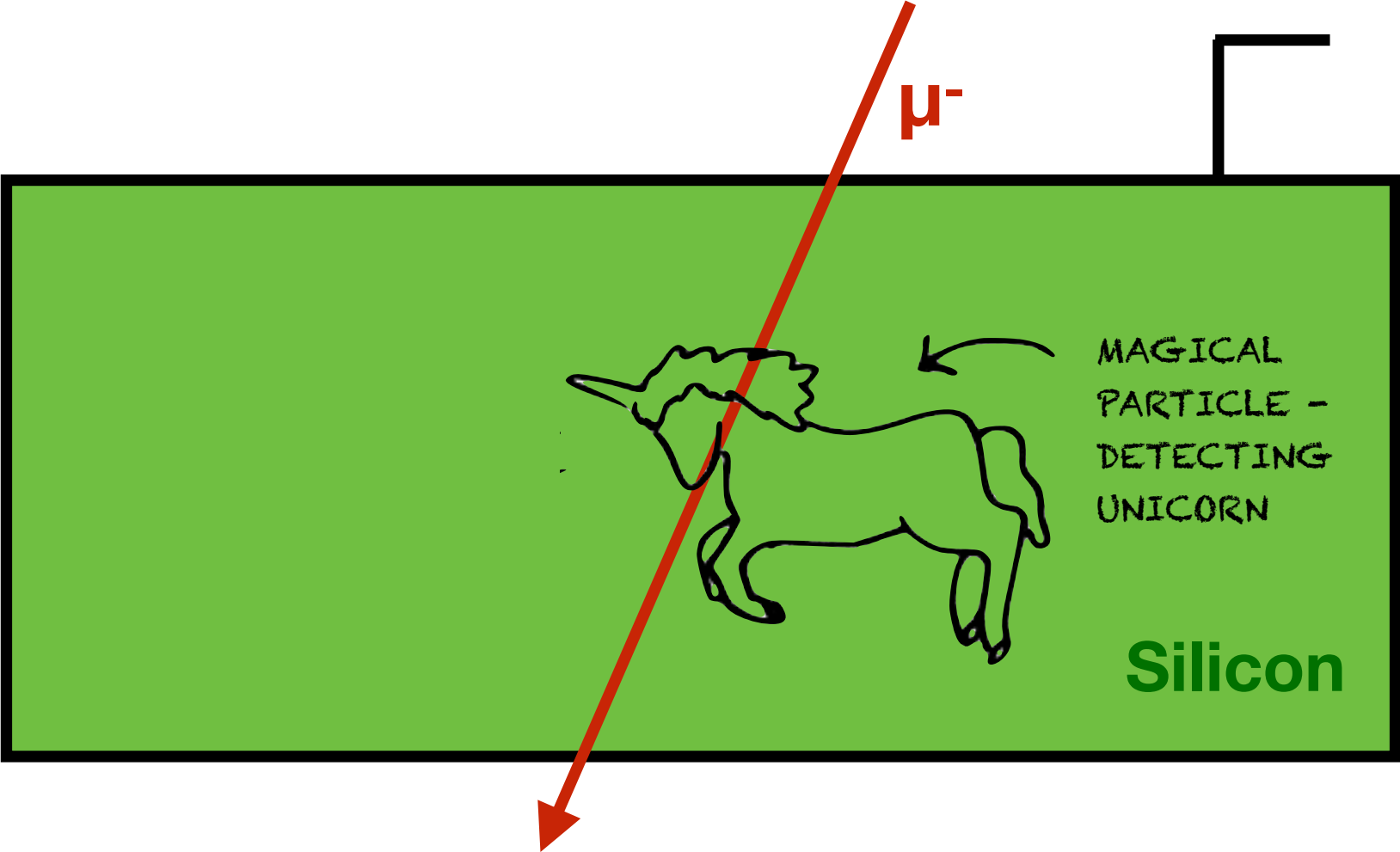
How stuff works



How silicon works...



How silicon works...?



Magic Rocks that Let Tiny Lightning Through Sometimes

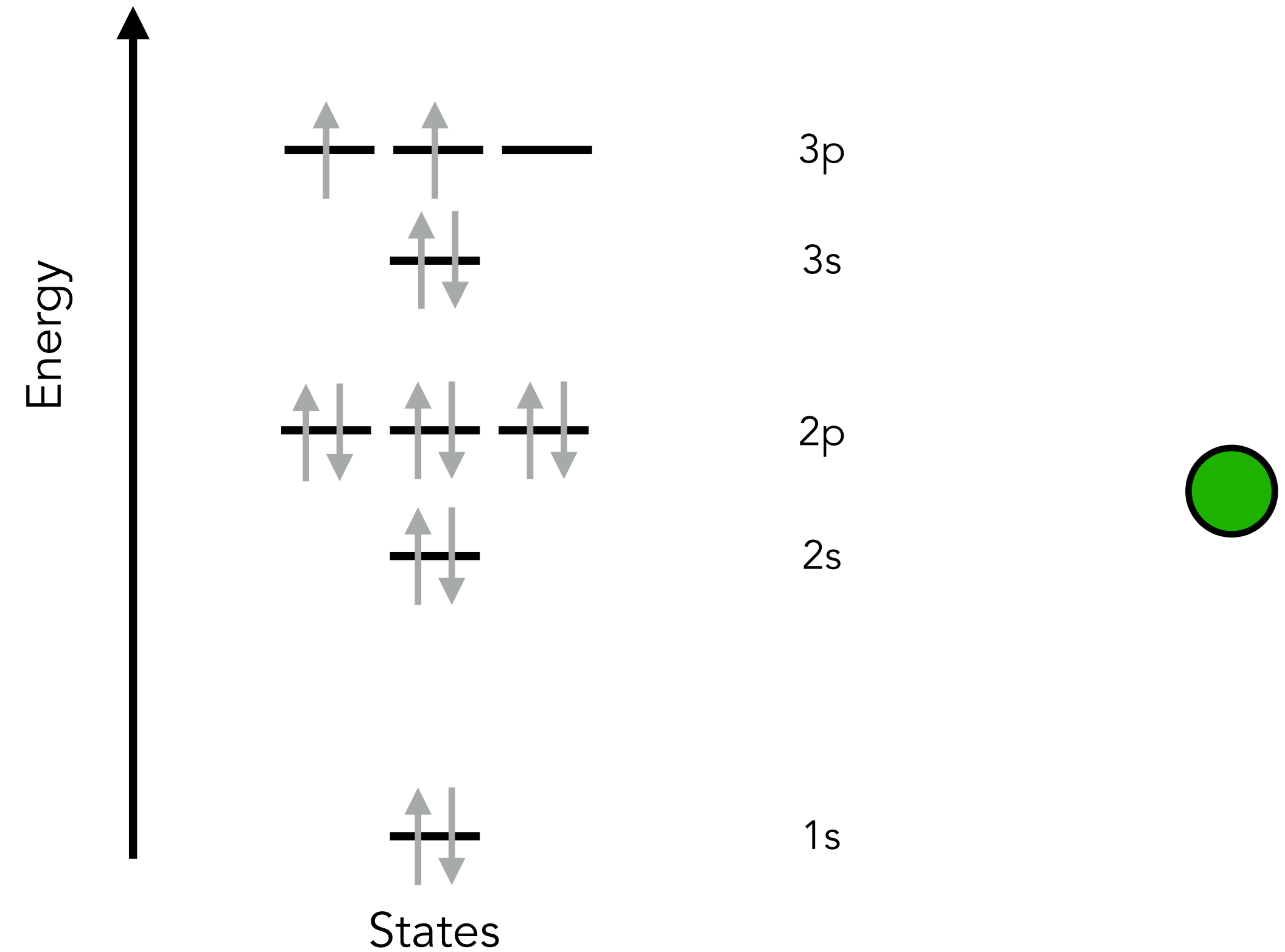
(Semiconductors)

Discrete energy levels

Electrons within an atom sit in discrete energy levels

- Each of these orbitals hosts two electrons - one spin up and one spin down

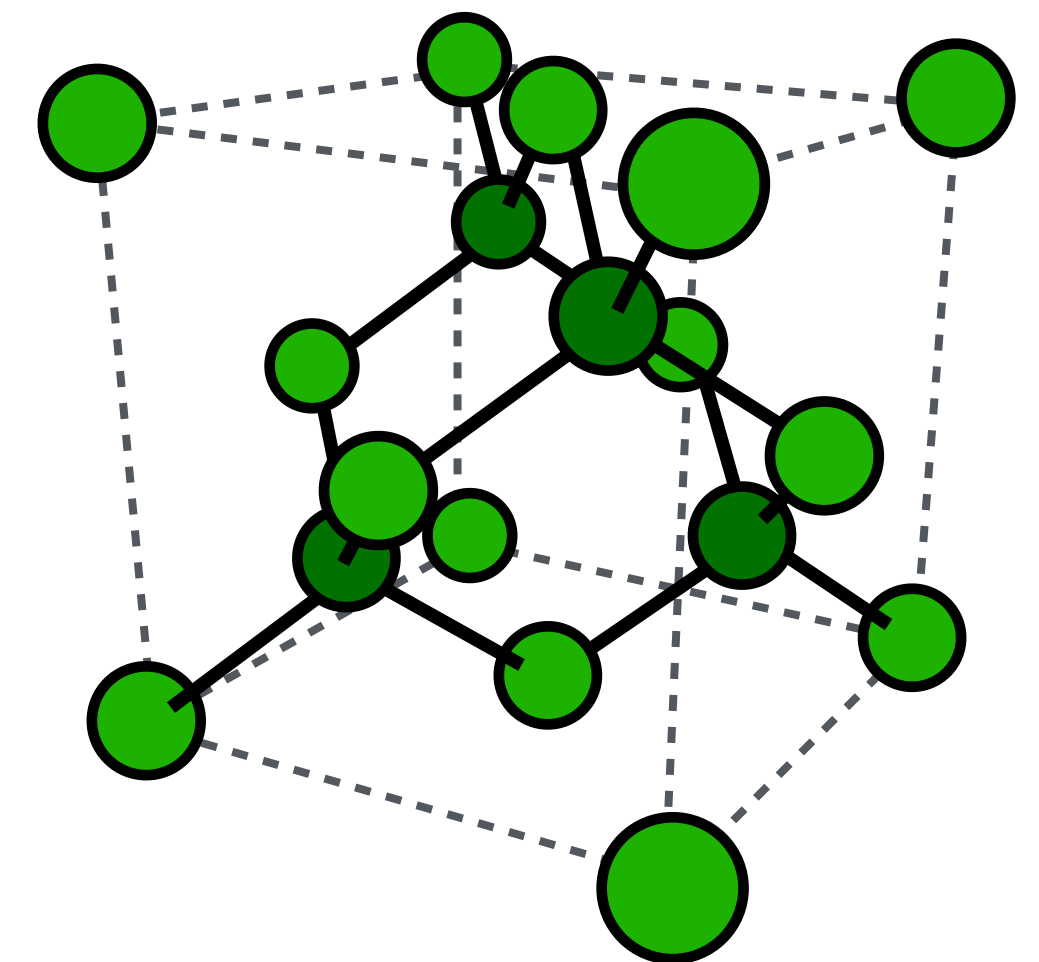
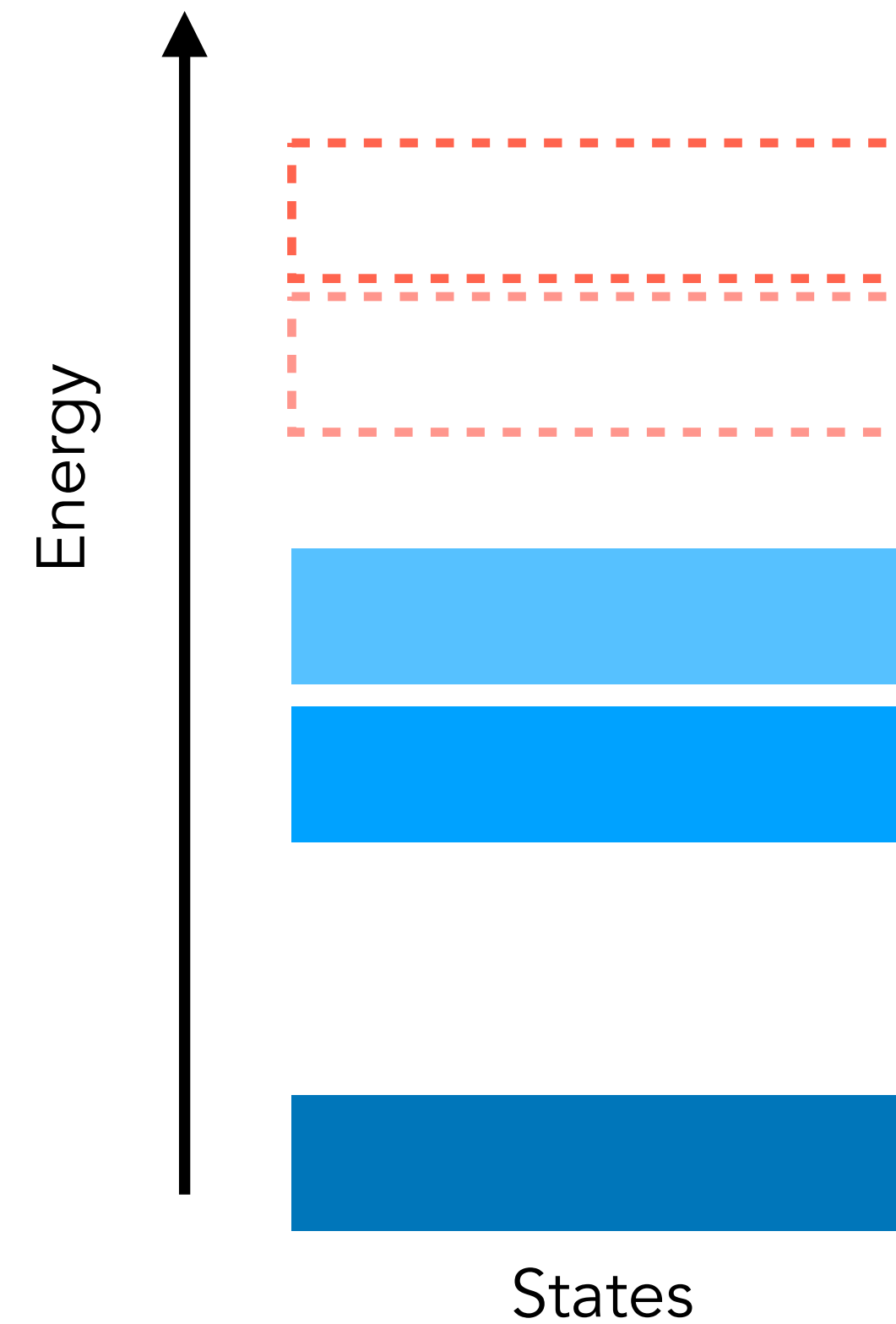
Lower energy states lie closer to the nucleus, while those in the outer shells are the so-called valence electrons that take part in chemical reactions



Band structure

In a periodic lattice, the available energy levels split to produce a continuous energy band

The inner bands will be tightly bound to the nucleus, and electrons there are unable to jump to neighbouring atoms



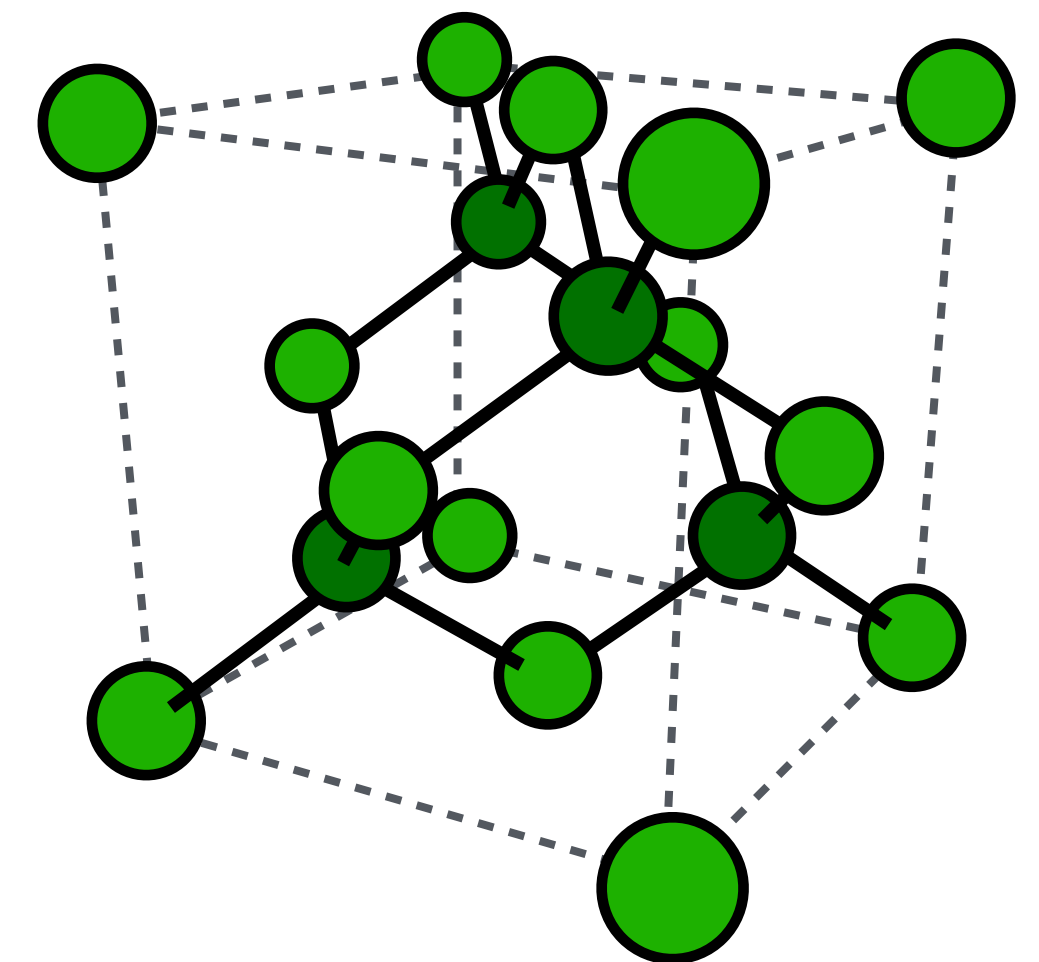
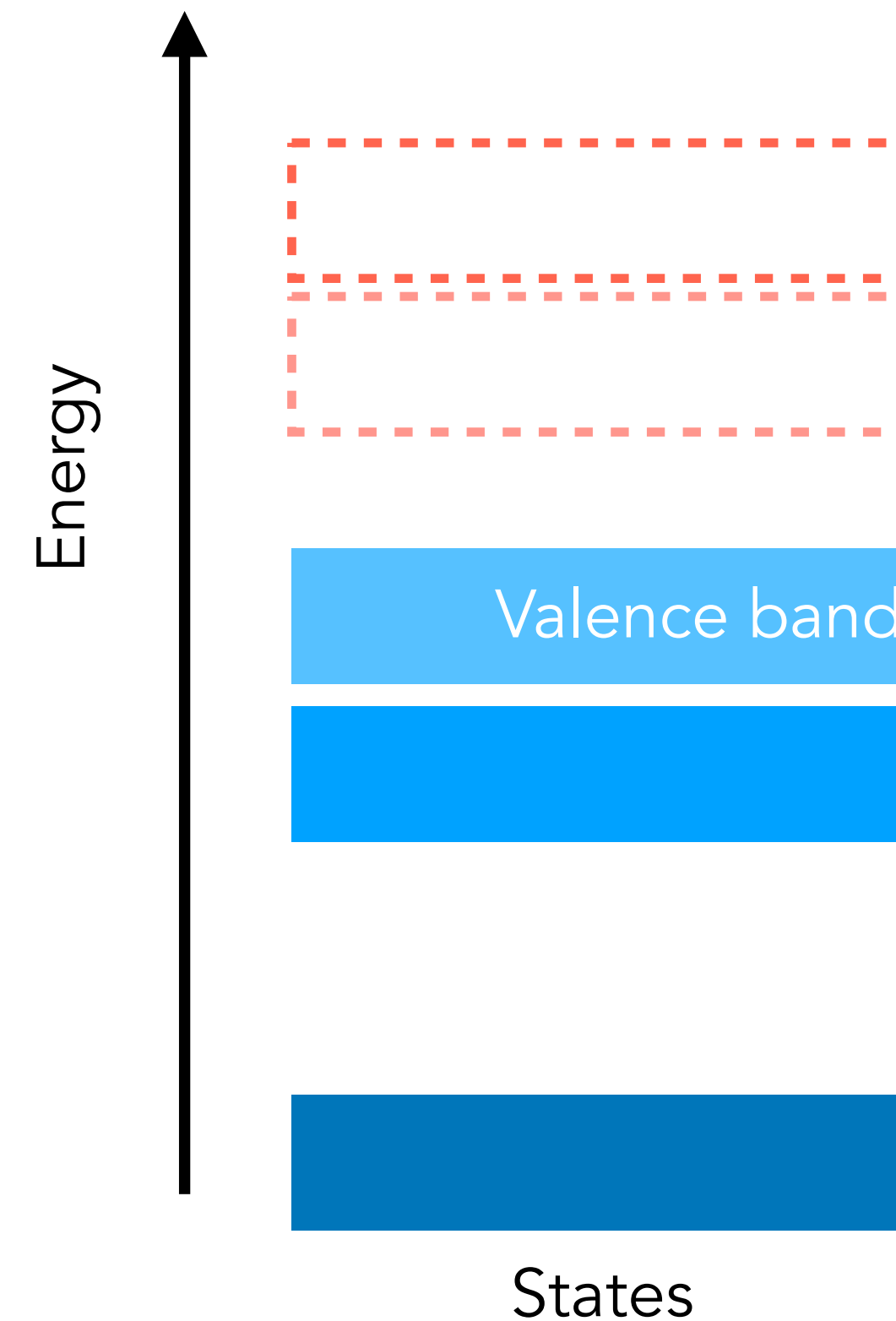
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The most interesting bands are

- The **valence band** - the highest occupied band where electrons are not free to move



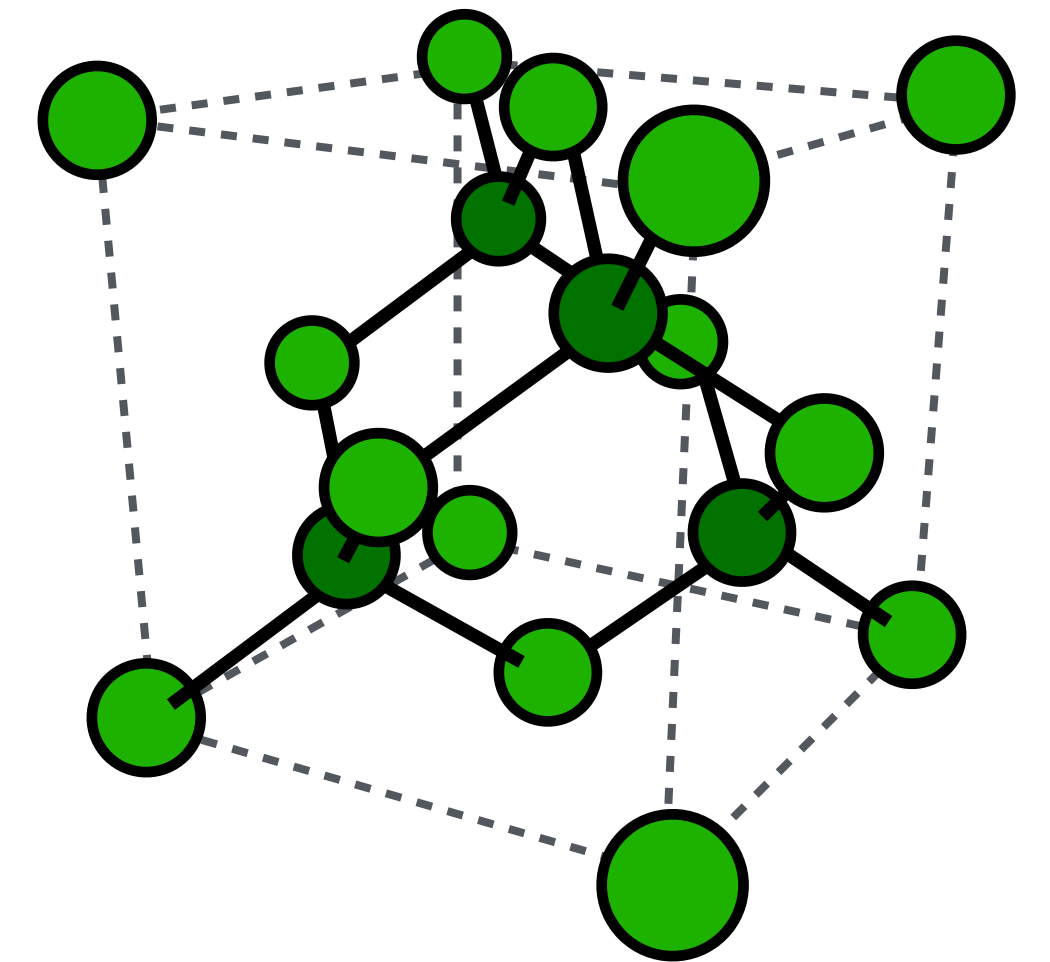
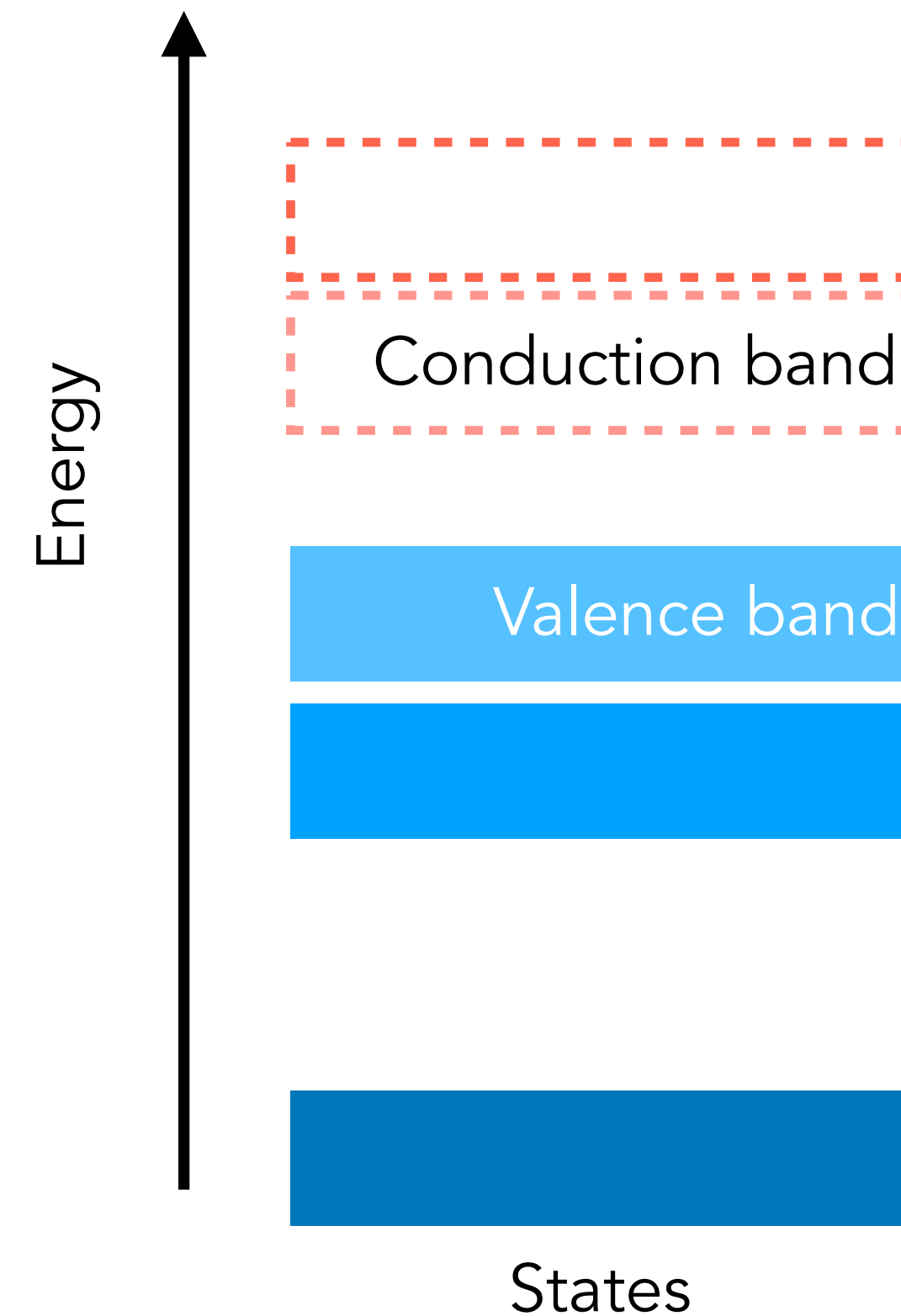
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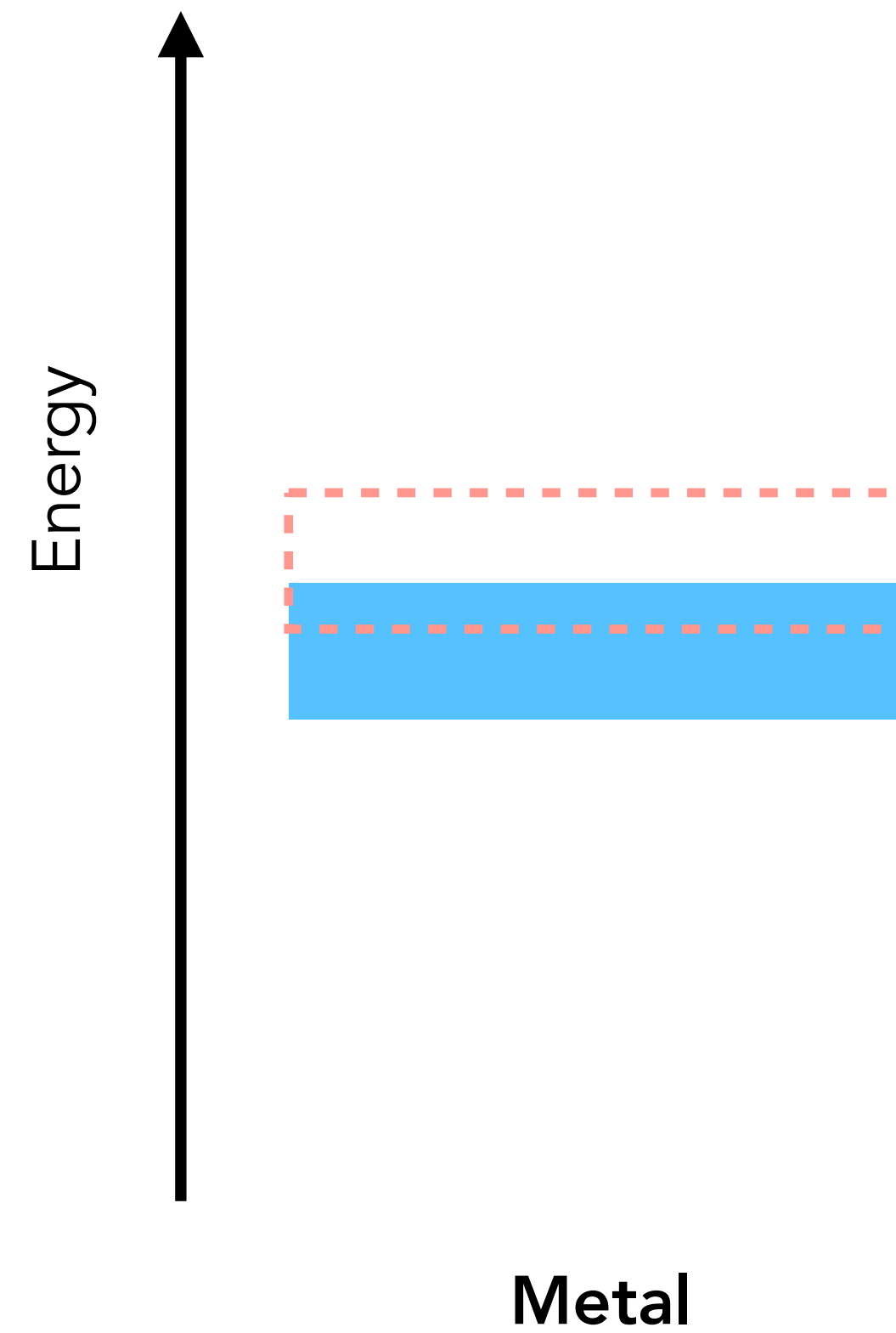
- The **valence band** - the highest occupied band where electrons are not free to move
- The **conduction band** - the lowest energy band where electrons are free to move



Electrical behaviour

The energy gap between the valence and conduction bands will determine the electrical behaviour of the solid

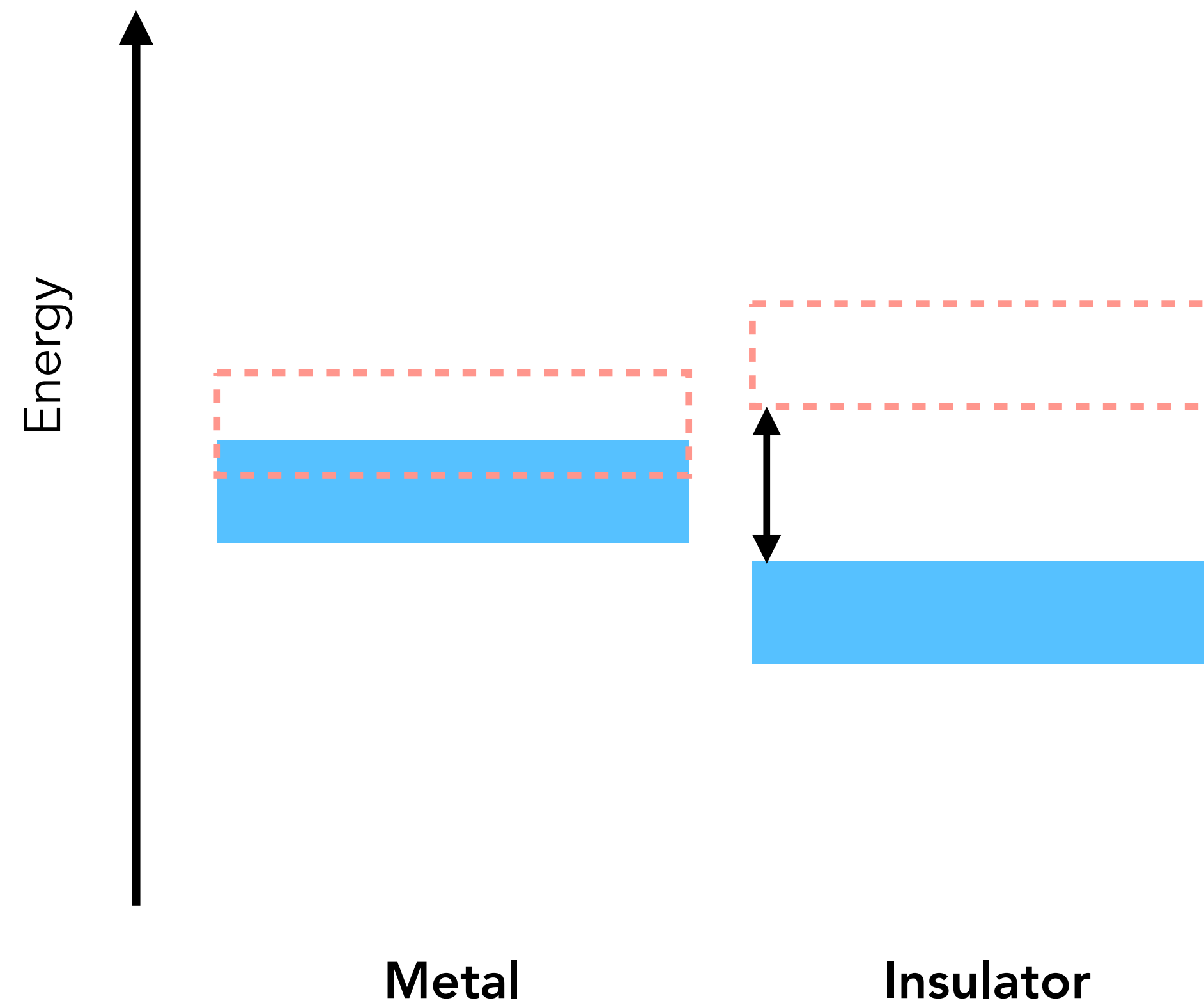
- **Metals** have overlapping states, so already there are charge carriers (electrons) able to move freely throughout the bulk



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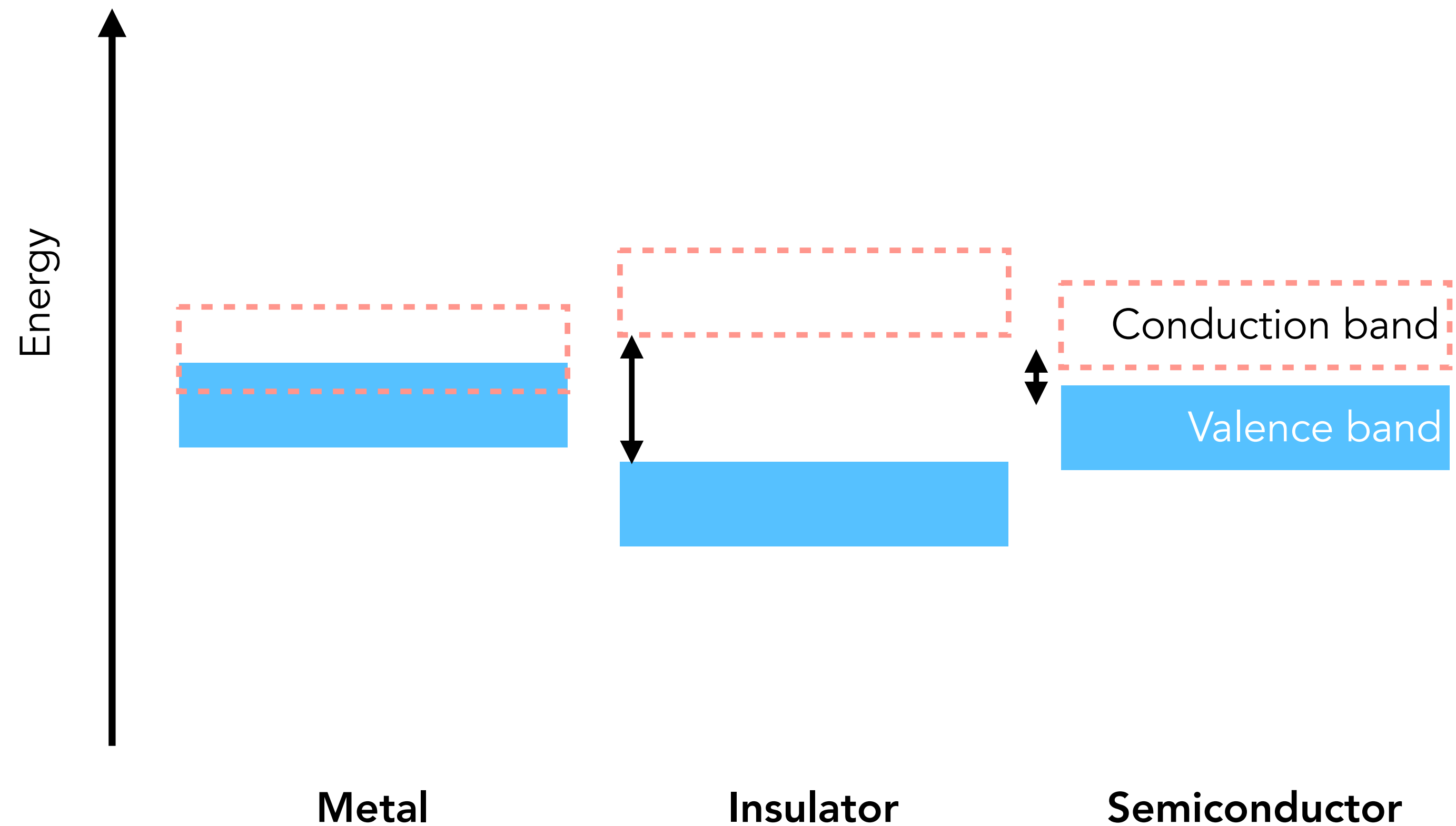
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Electrical behaviour

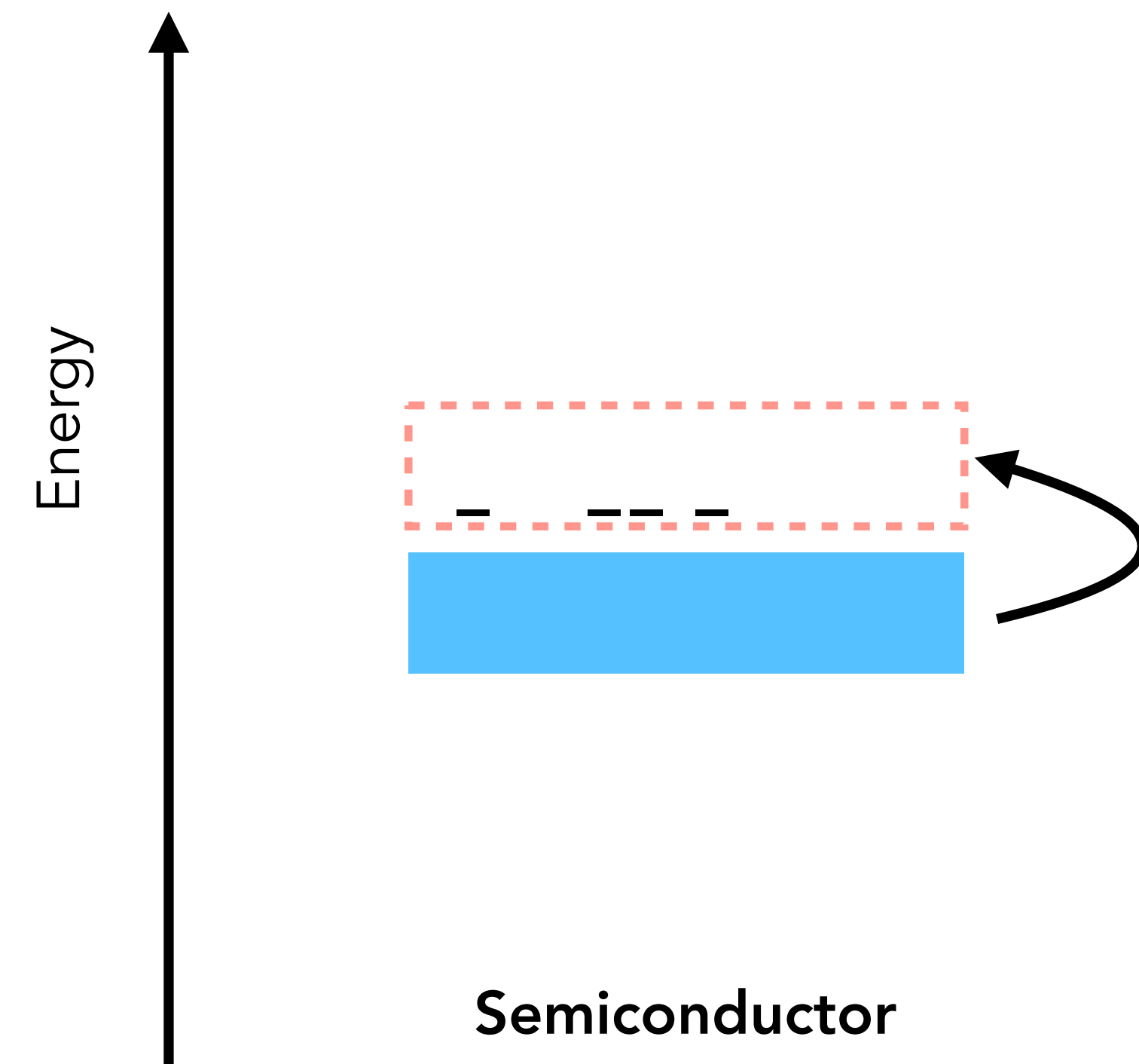
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- **Metals** have overlapping states, so already there are charge carriers (electrons) able to move freely throughout the bulk
- **Insulators** have a large **band gap**, leading to poor charge flow
- **Semiconductors** have a **band gap** of ~few eV, giving some charge carriers at room temperature (eg. Silicon $\sim 10^{10}$ charge carriers cm^{-3} for a density of 5×10^{22} atoms cm^{-3})



Charge carriers in semiconductors

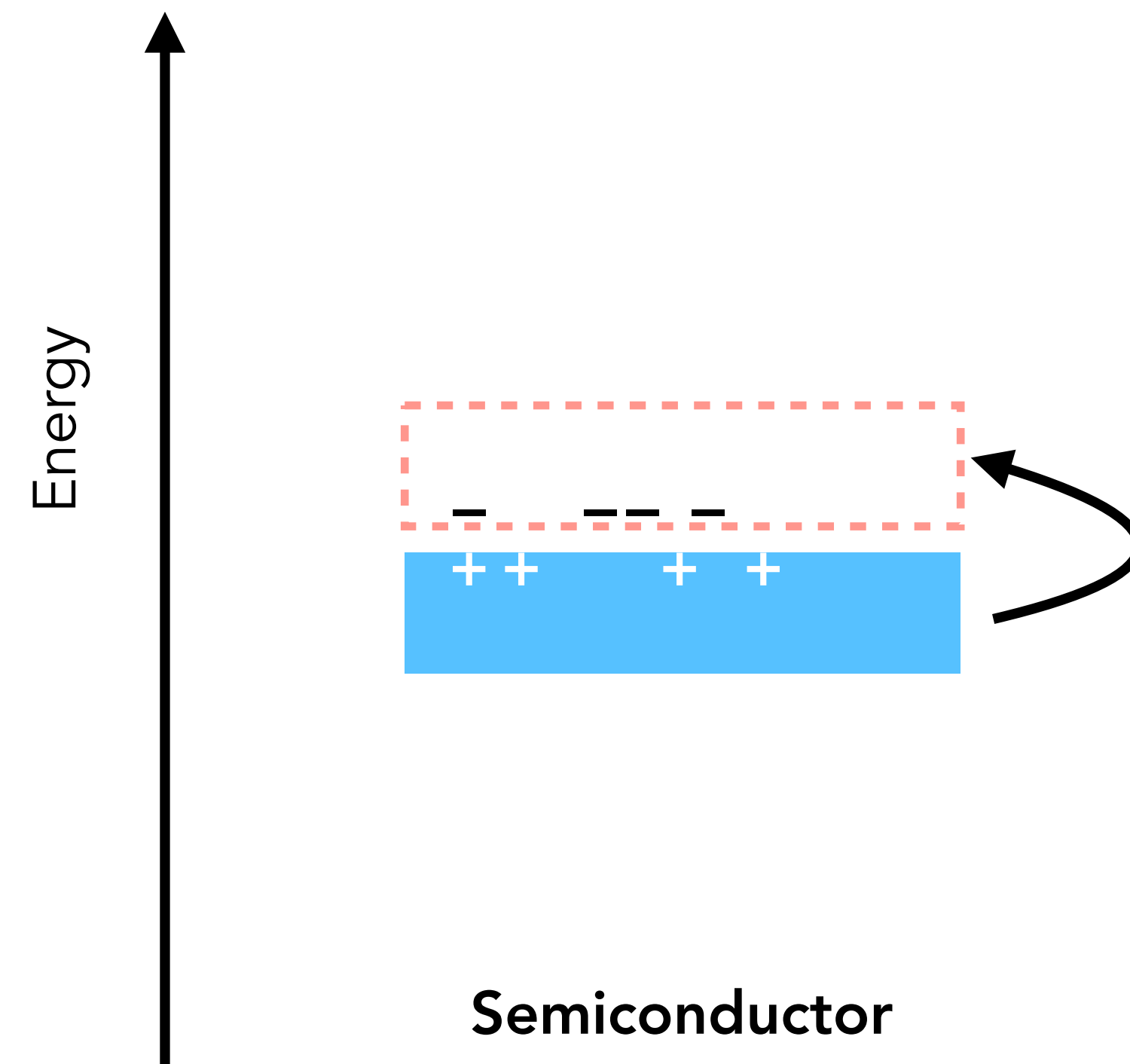
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Charge carriers in semiconductors

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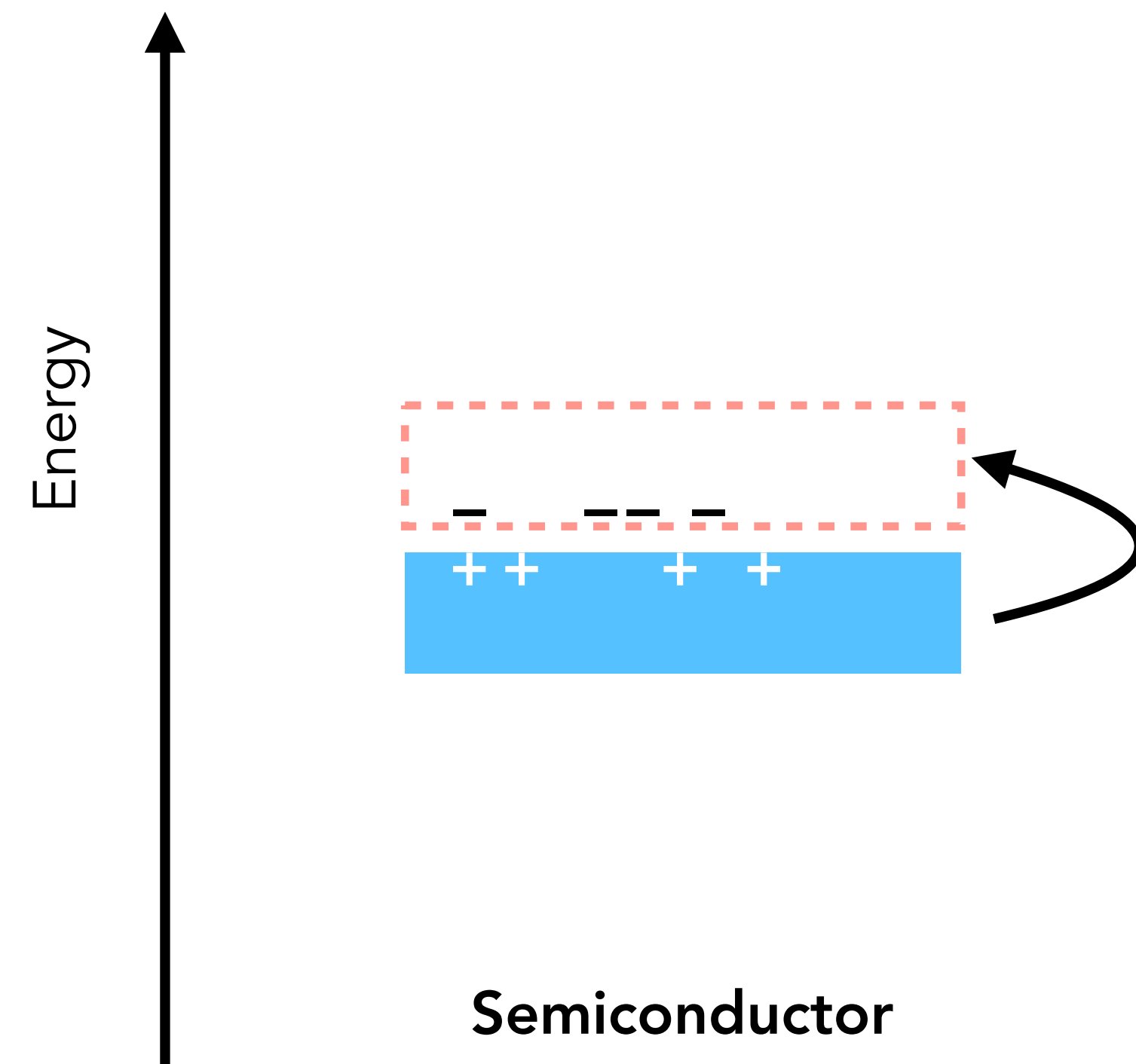
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Other electrons within the valence band can move into these holes, such that the hole appears to travel through the material

- Holes are treated as effective particles which move freely in the valence band and allow a current to flow



Charge carriers in semiconductors

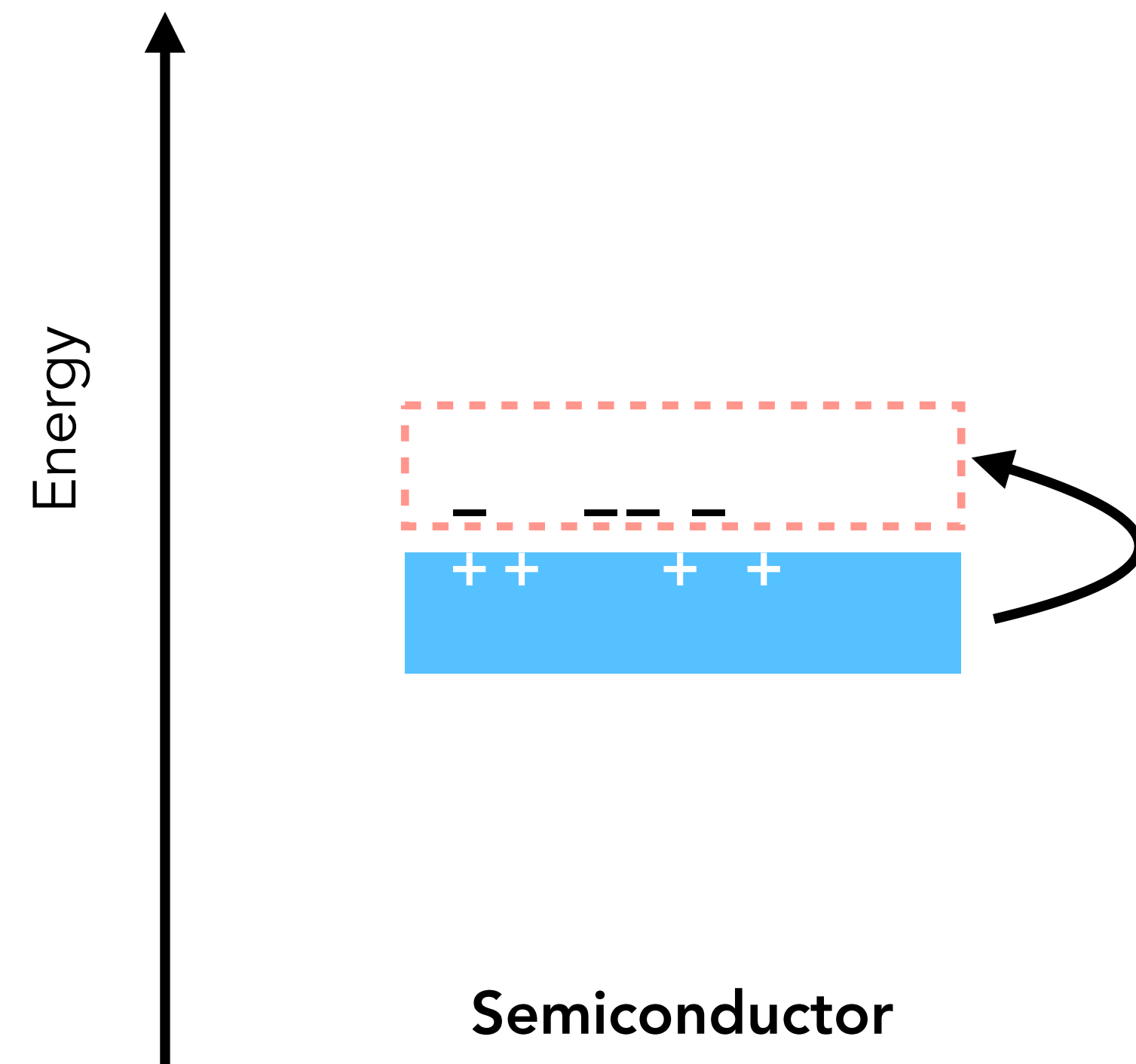
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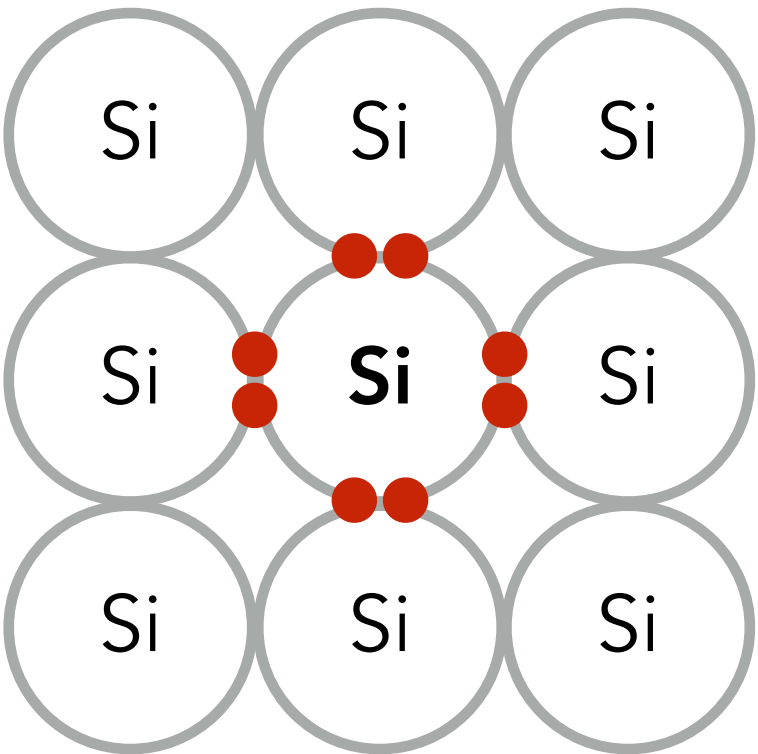
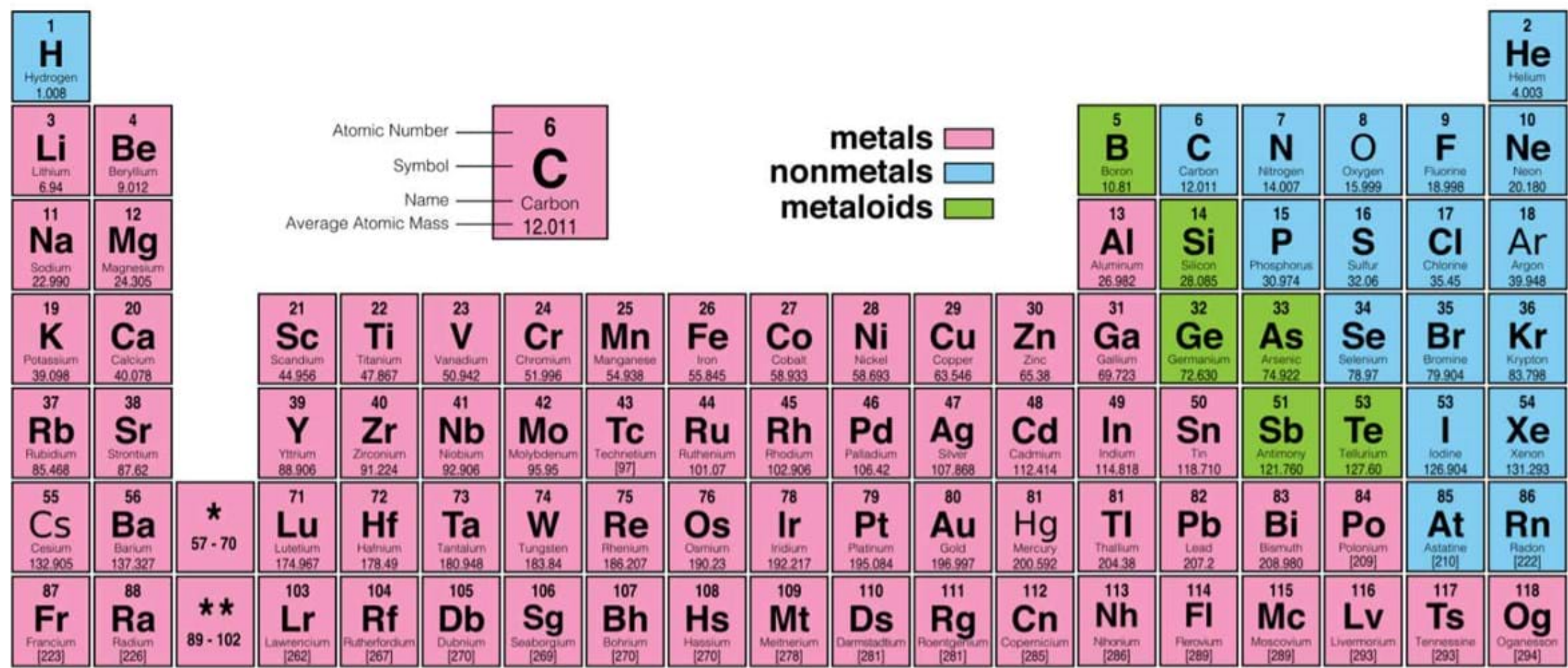
Each movement of an electron from the valence to conduction band produces an **electron-hole pair**



Doping of semiconductors

Silicon is a **group 4** element, and therefore contains 4 valence electrons

- It will therefore form 4 covalent bonds in order to reach a stable configuration



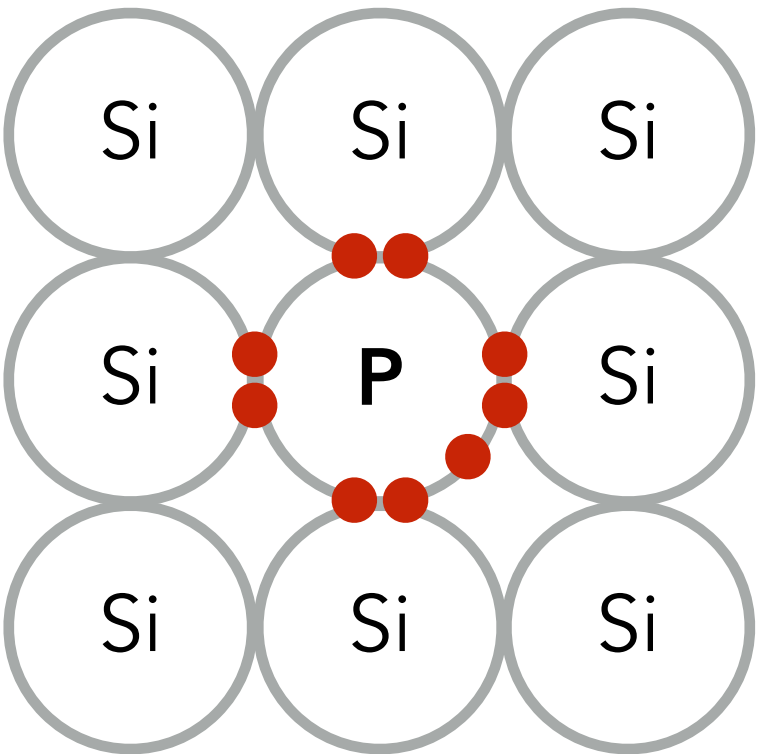
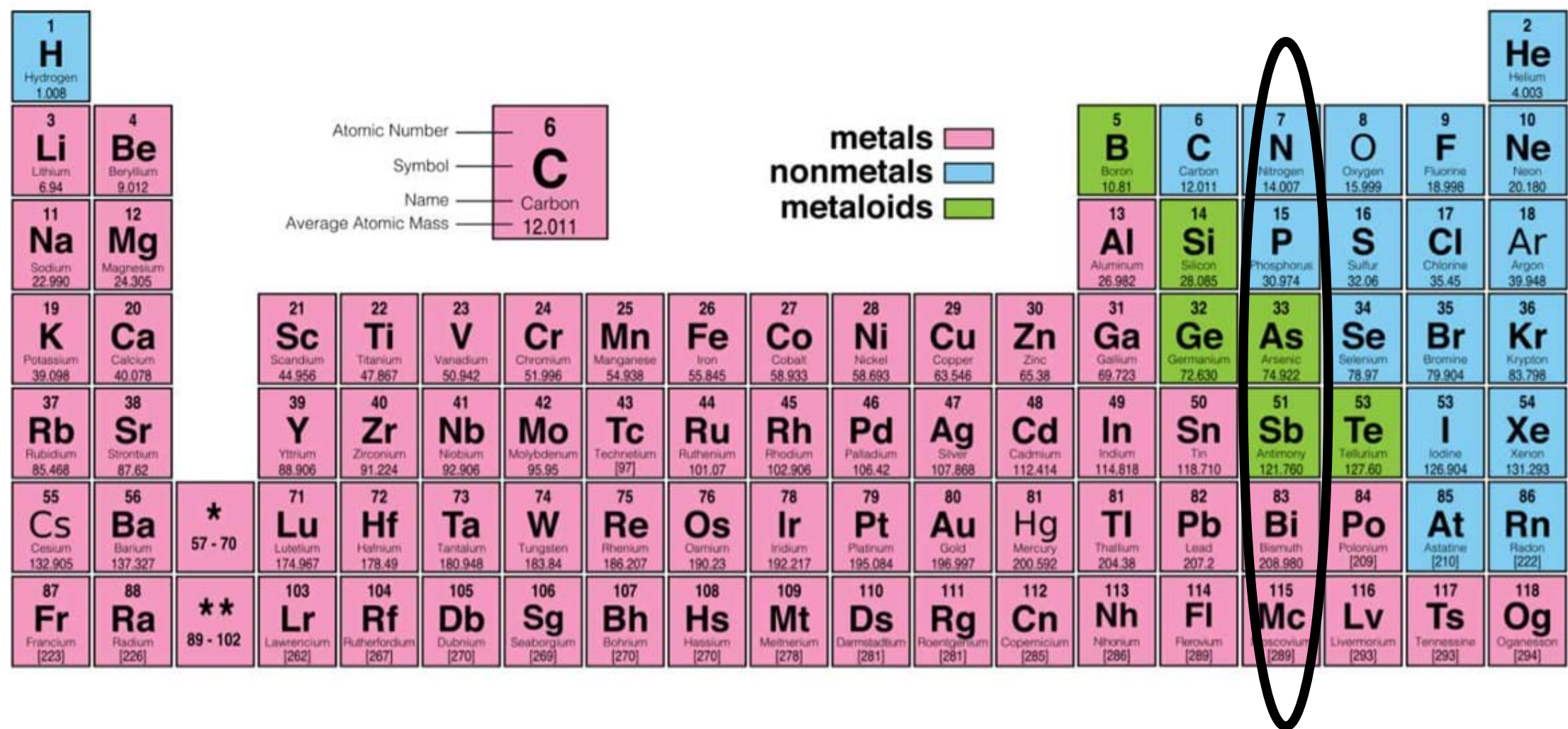
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If we introduce a small amount of a **group 5** element, the crystal lattice will remain intact but there will be some atoms with an excess of electrons - **n-type**

doping



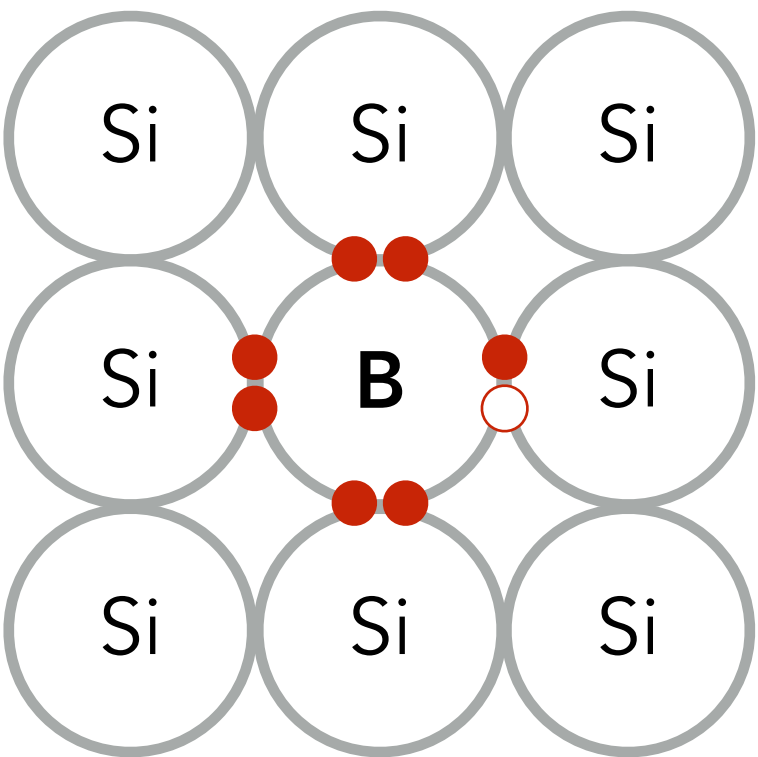
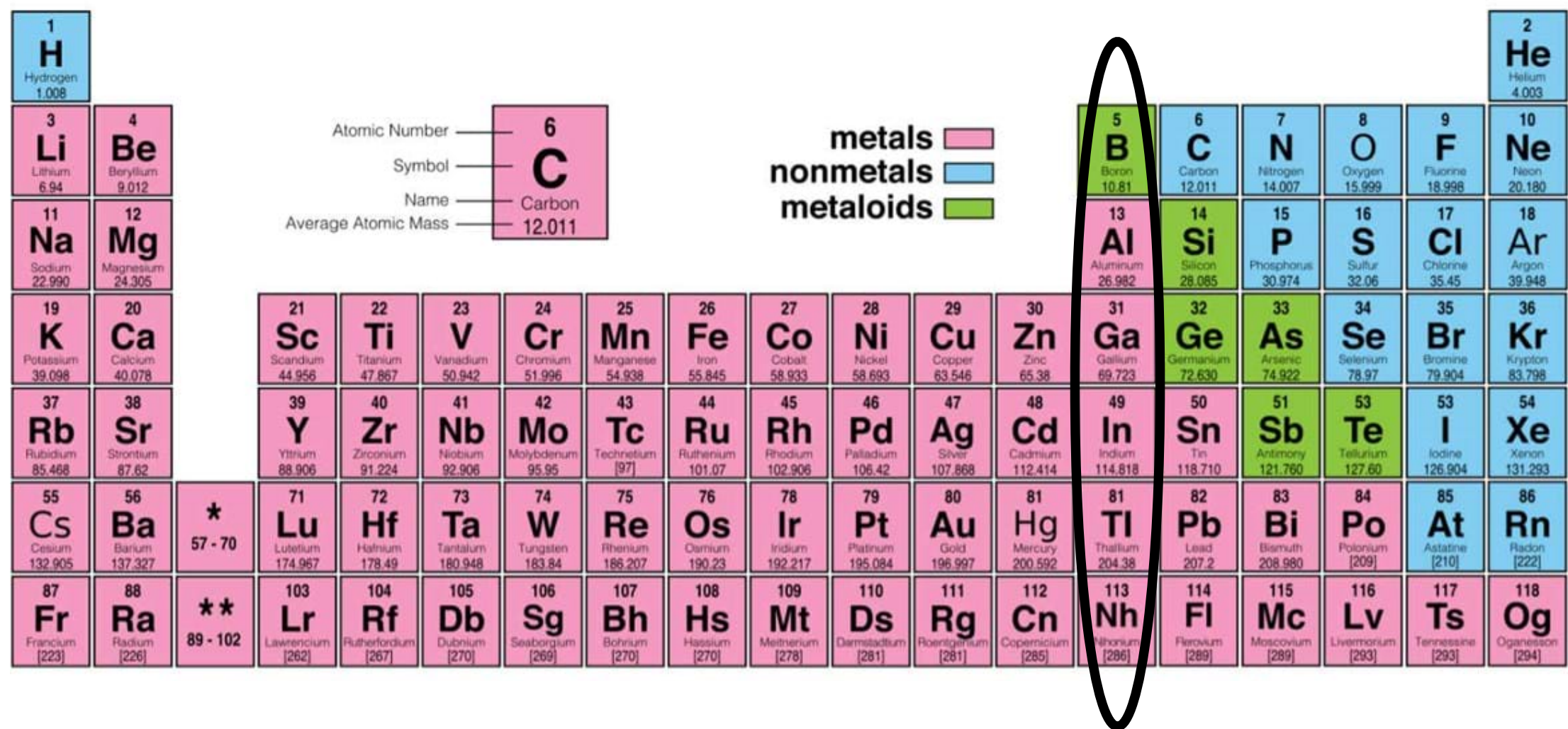
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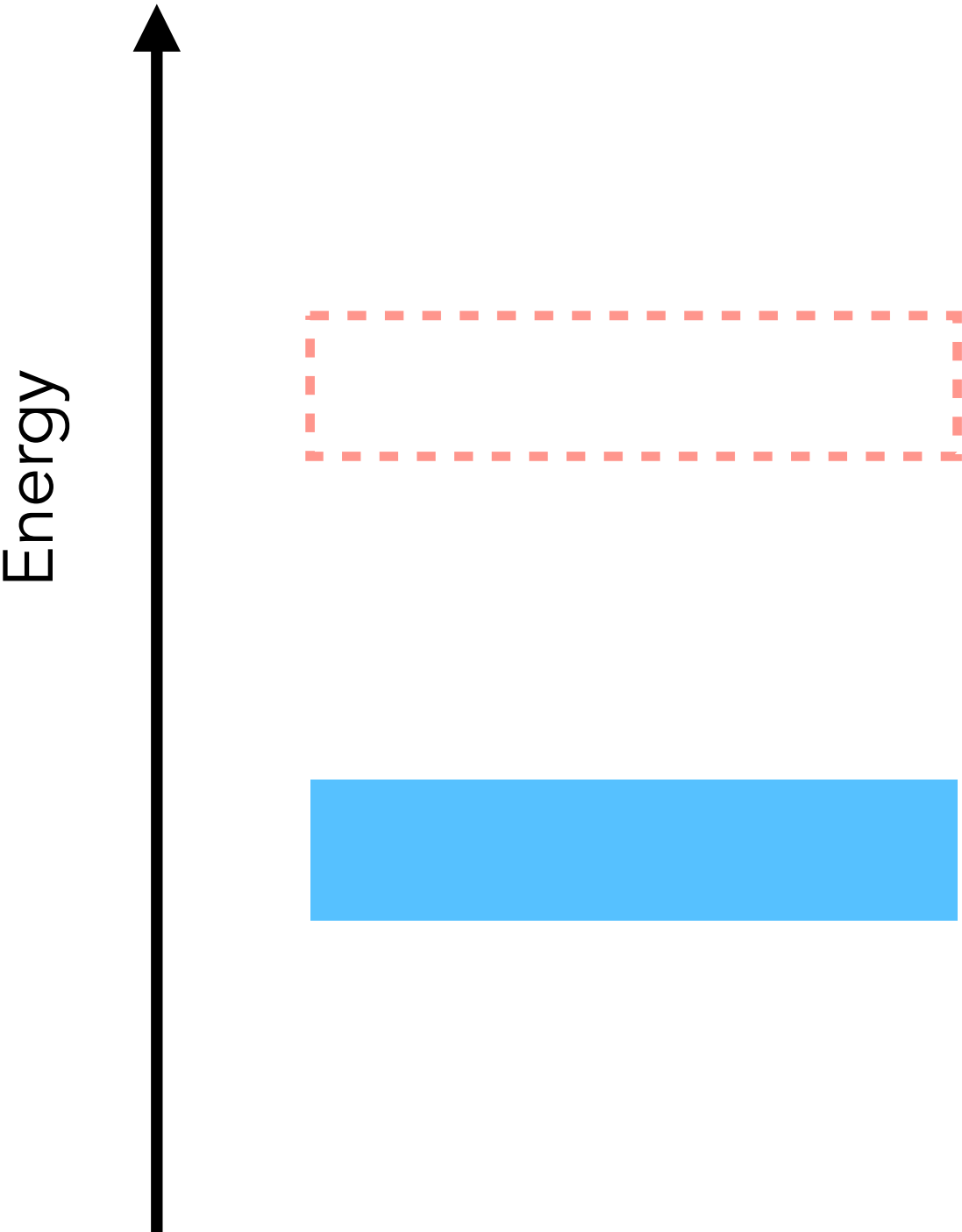
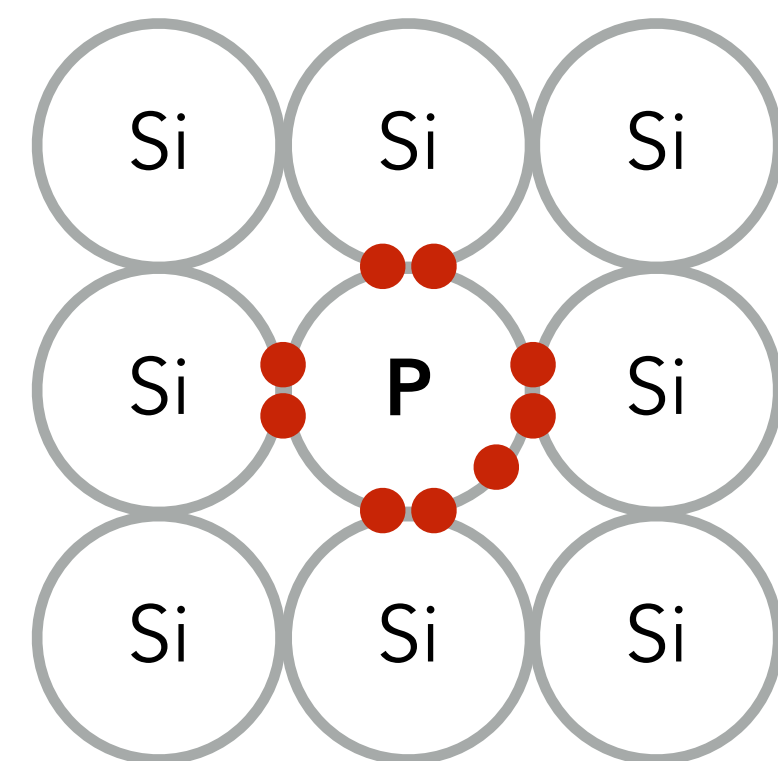
If we introduce a small amount of a **group 3** element, there will be some atoms with an electron deficiency -

p-type doping



Band structure n-type semiconductors

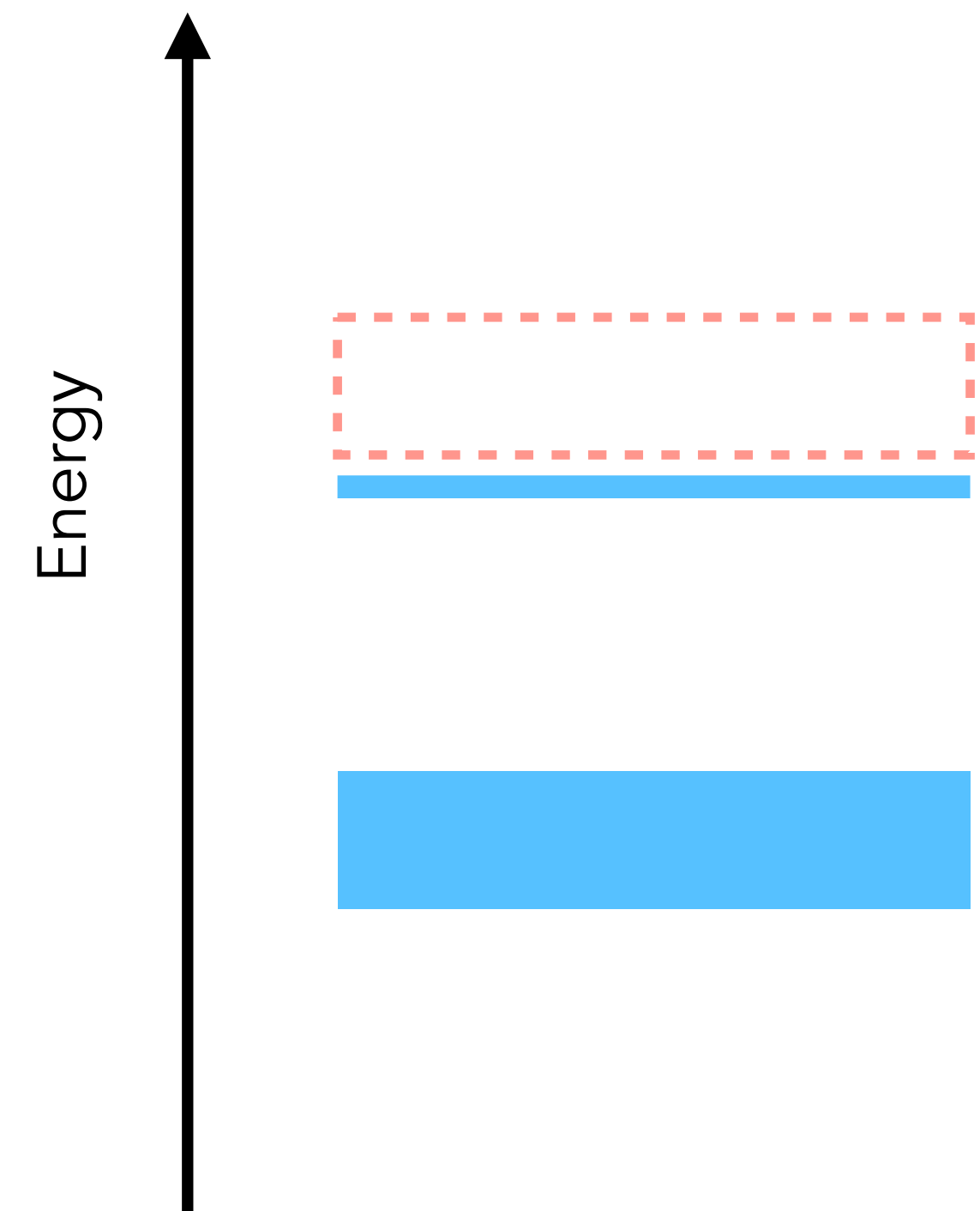
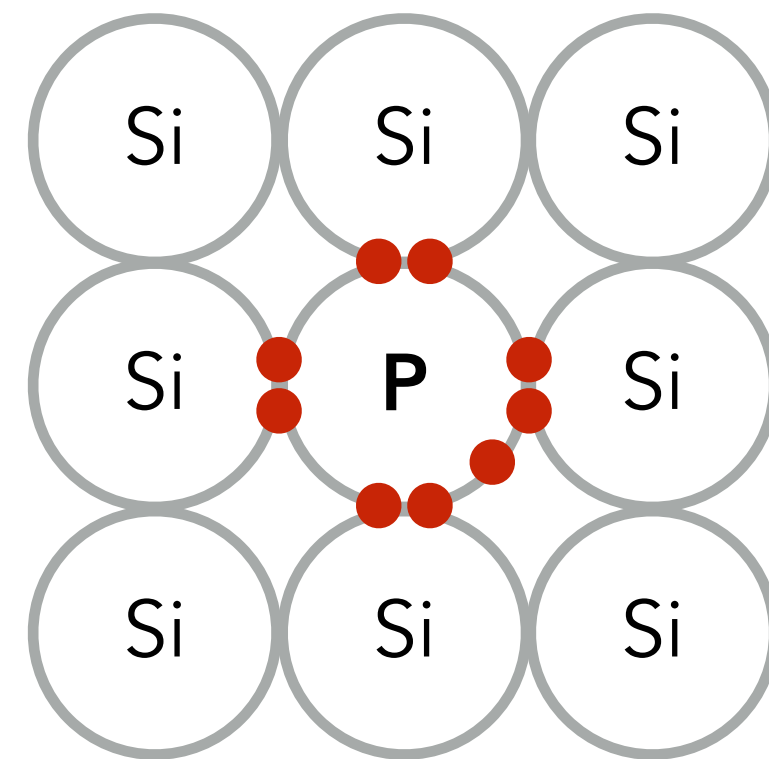
For n-type doping, there is an excess of electrons in the crystal structure



Band structure n-type semiconductors

For n-type doping, there is an excess of electrons in the crystal structure

- These sit within a new band close to the conduction band

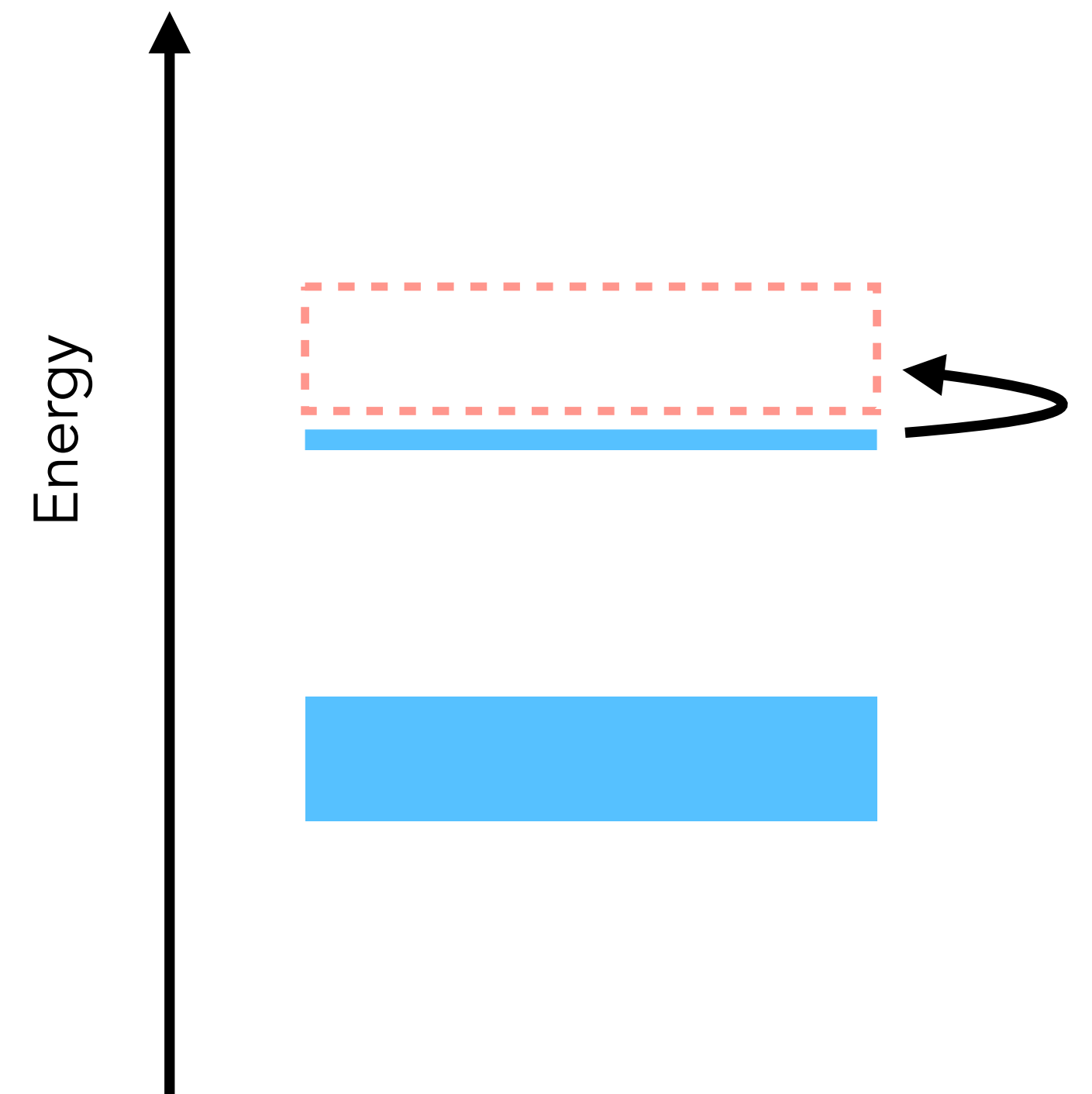


Band structure n-type semiconductors

For n-type doping, there is an excess of electrons in the crystal structure

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The electrons in this band will almost all move to the conduction band with thermal energy alone



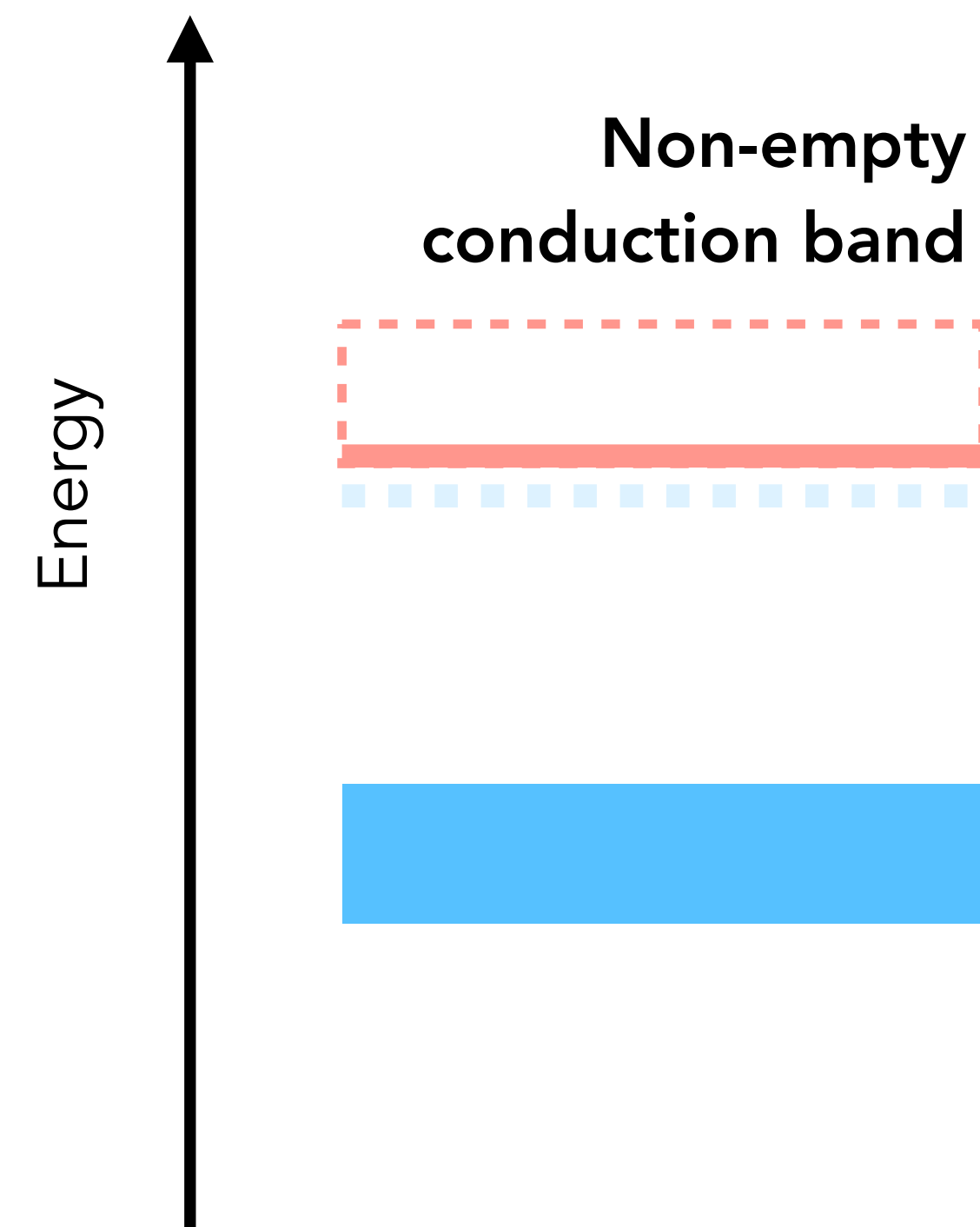
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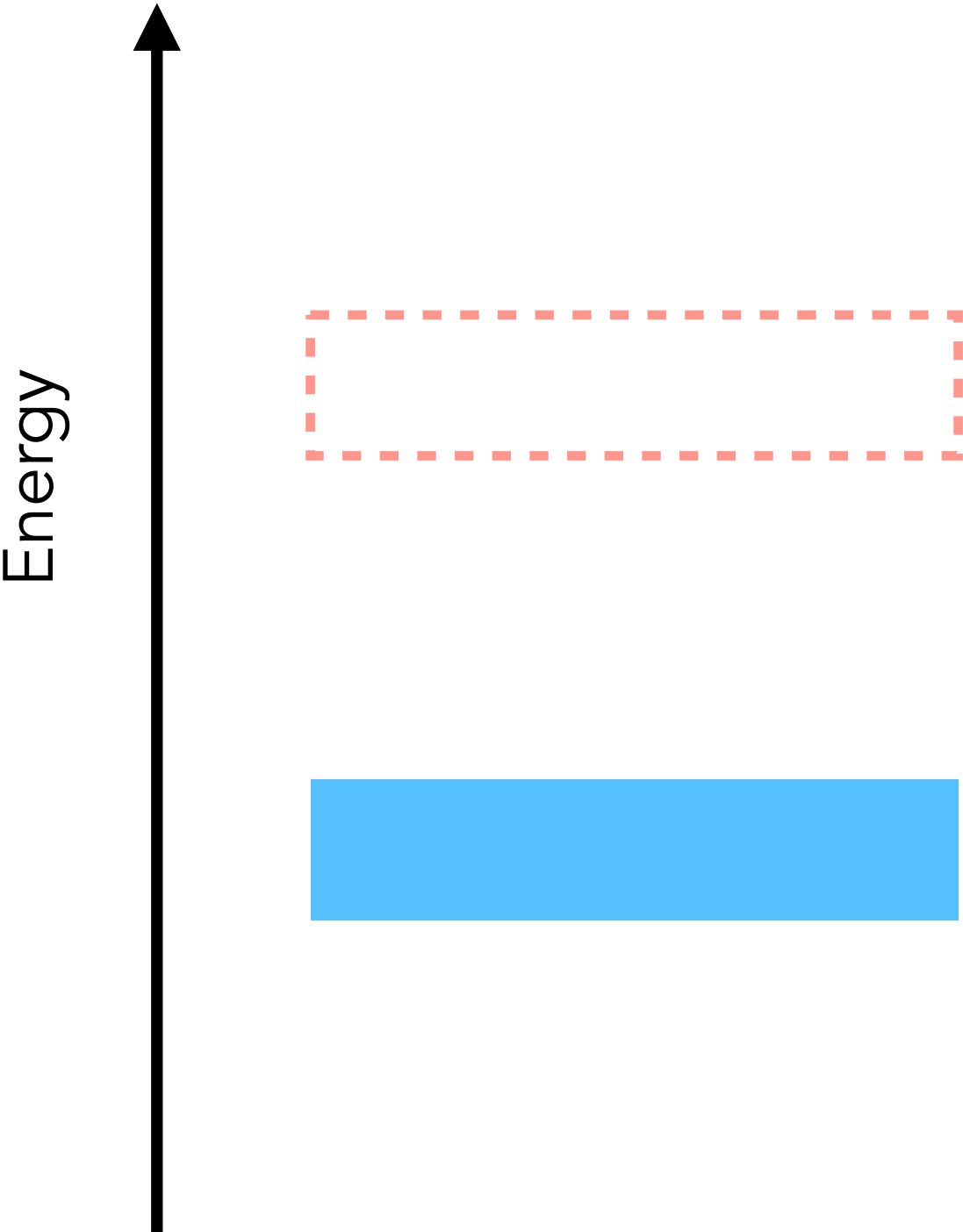
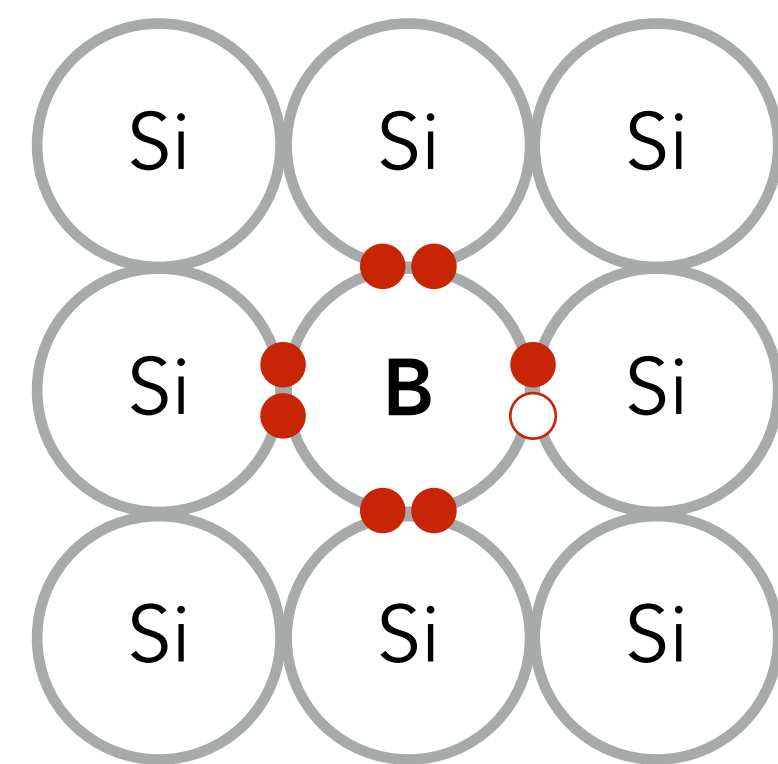
The electrons in this band will almost all move to the conduction band with thermal energy alone

- The majority carriers for n-type materials are electrons in the conduction band
- Significant carrier densities can be achieved introduction tiny fractions of impurities (eg. 1 ppm P = 5×10^{16} electrons cm^{-3} versus 10^{10} cm^{-3} for undoped Si)



Band structure p-type semiconductors

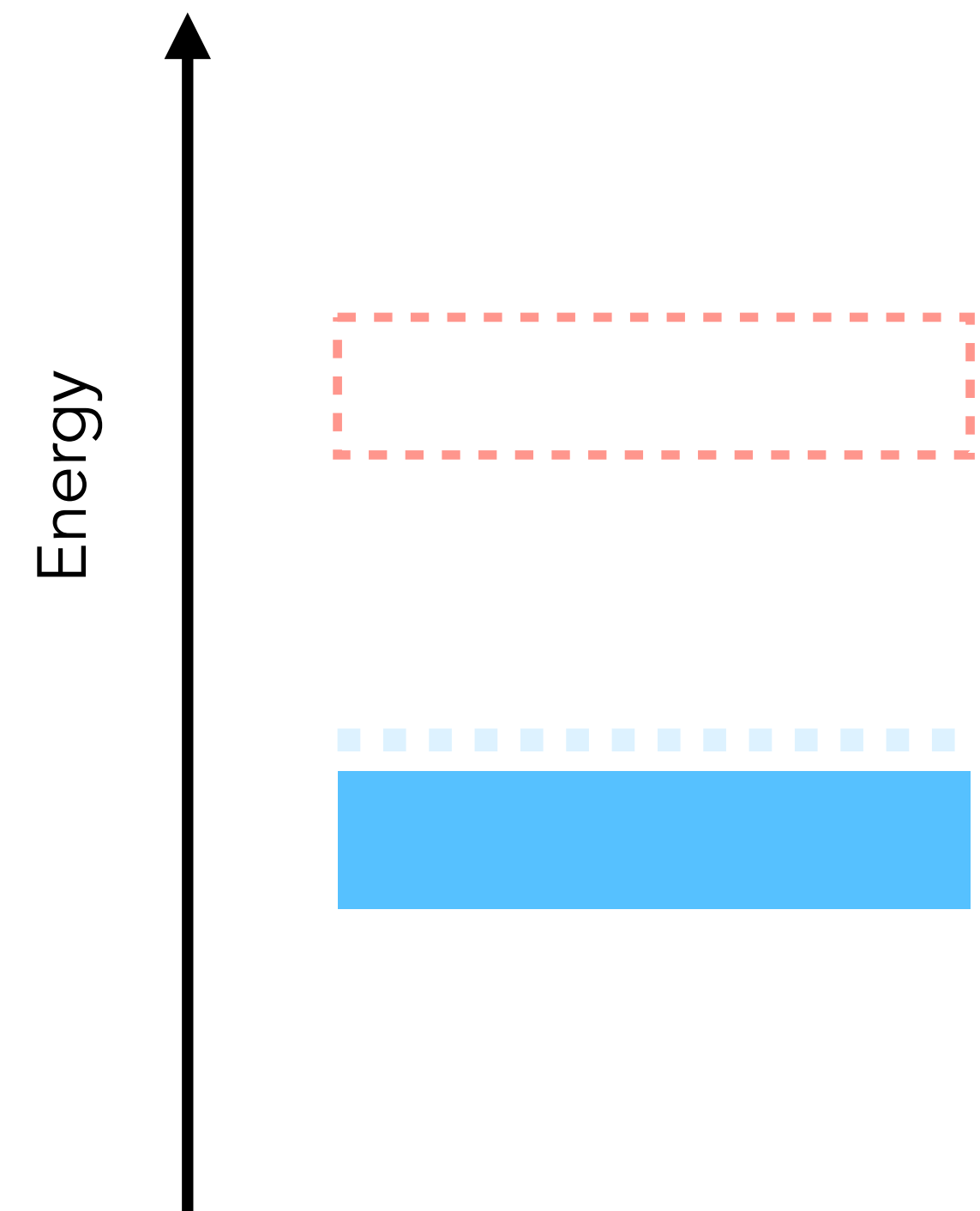
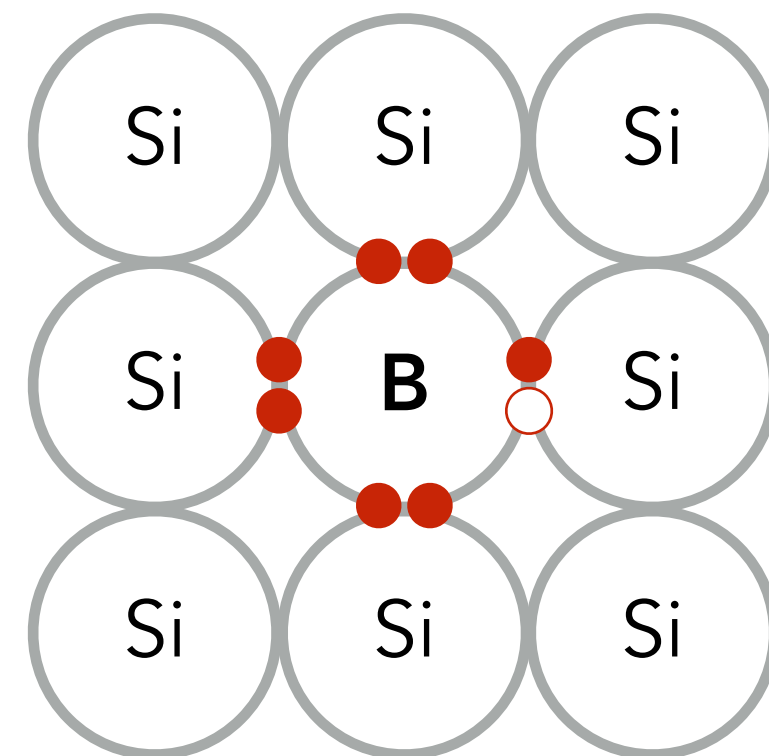
For p-type doping, there is a deficiency of electrons in the crystal structure



Band structure p-type semiconductors

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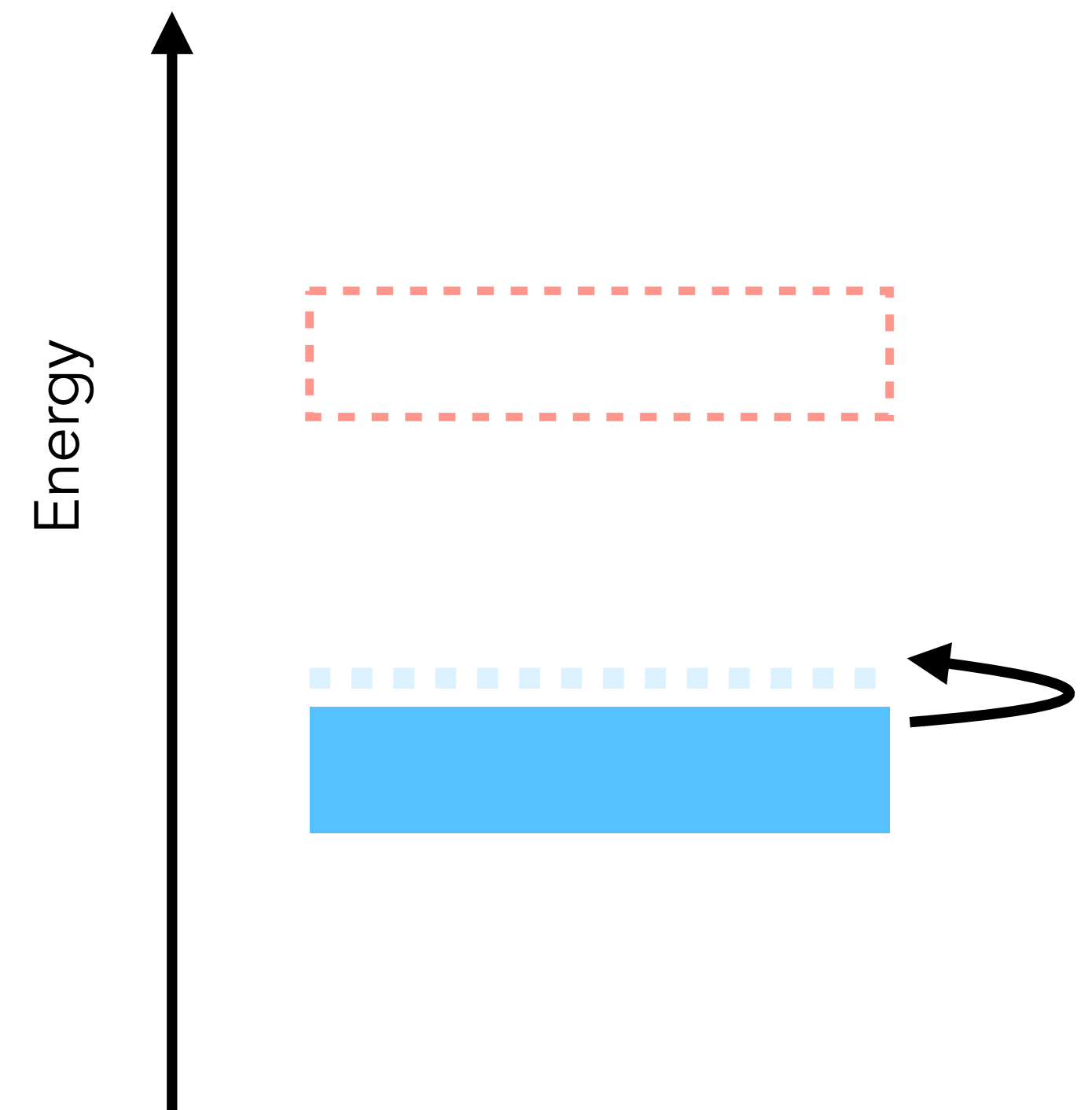


Band structure p-type semiconductors

For p-type doping, there is a deficiency of electrons in the crystal structure

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Electrons in the valence band will move to occupy these holes with thermal energy alone



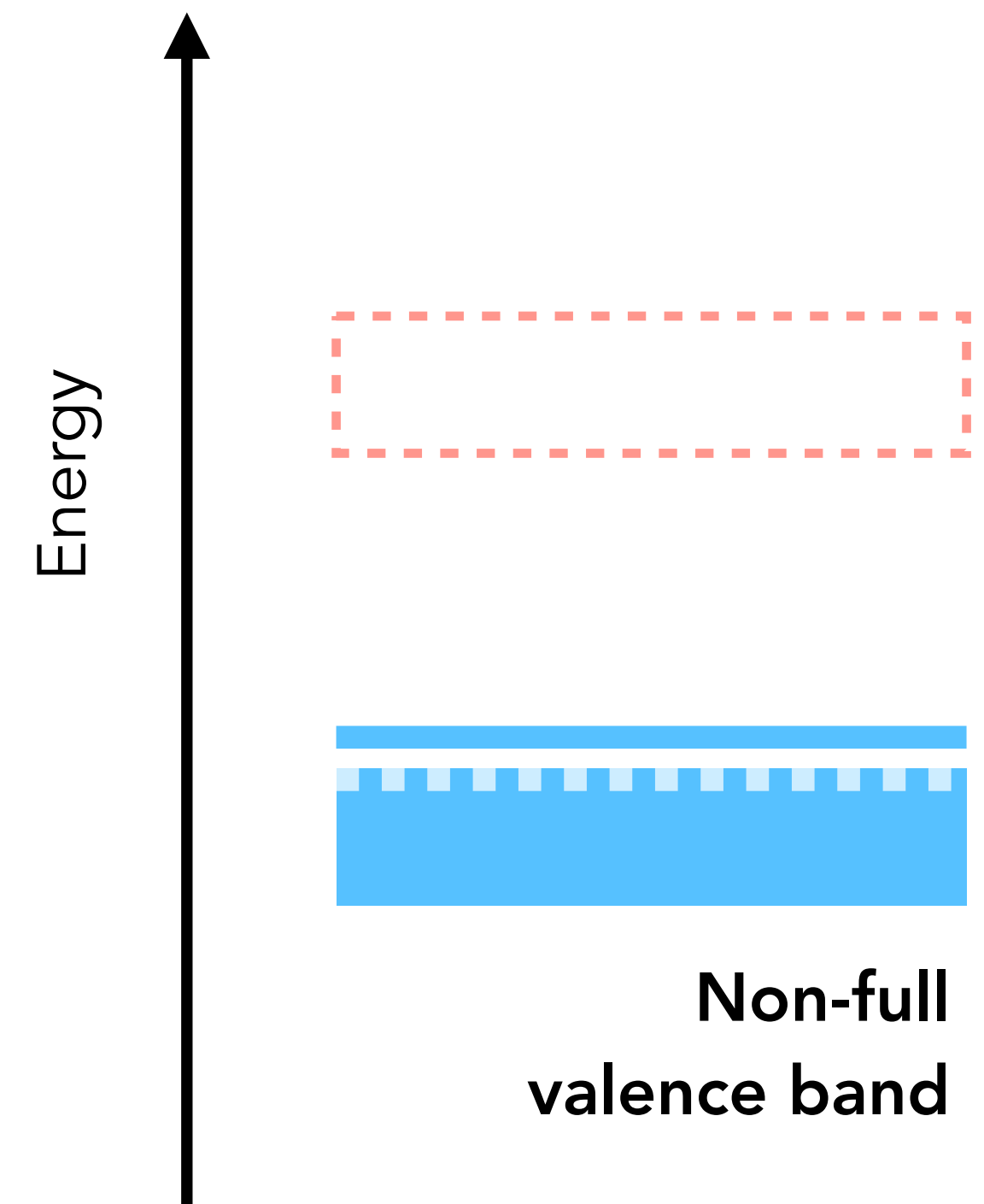
Band structure p-type semiconductors

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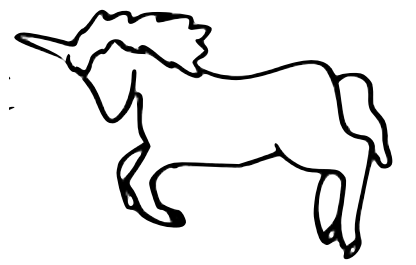
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Electrons in the valence band will move to occupy these holes with thermal energy alone

- The majority carriers for p-type materials are holes in the valence band



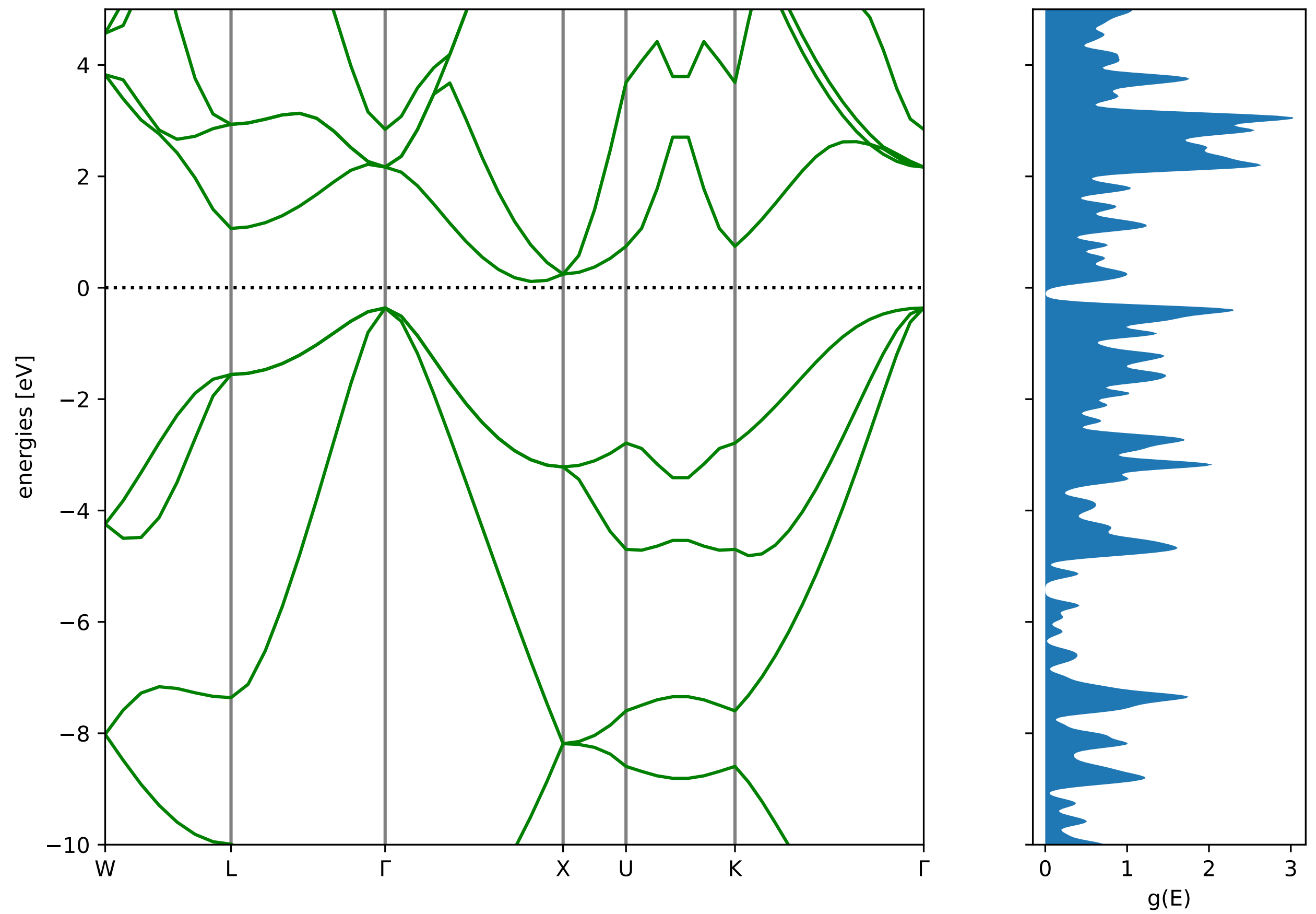
Band theory



All of this is of course very simplified

- Band structure of solids is complicated
- Several models, from Nearly Free Electron (NFE) to tight binding/Linear Combination of Atomic Orbitals (LCAO)

Band structure gives a spread of different possible interactions depending on momentum transfer, phonon interaction, etc



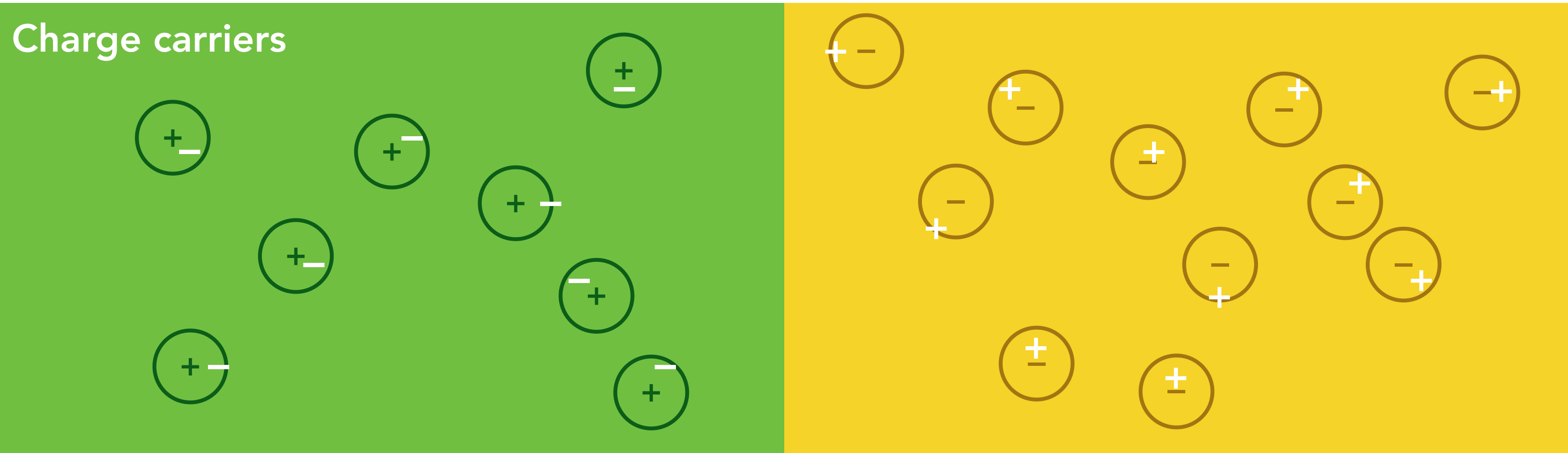
**The Place Where Two Different Kinds of Rock Pieces Meet to
Push Electricity Only One Way**

(P-N junctions)

P-N junctions

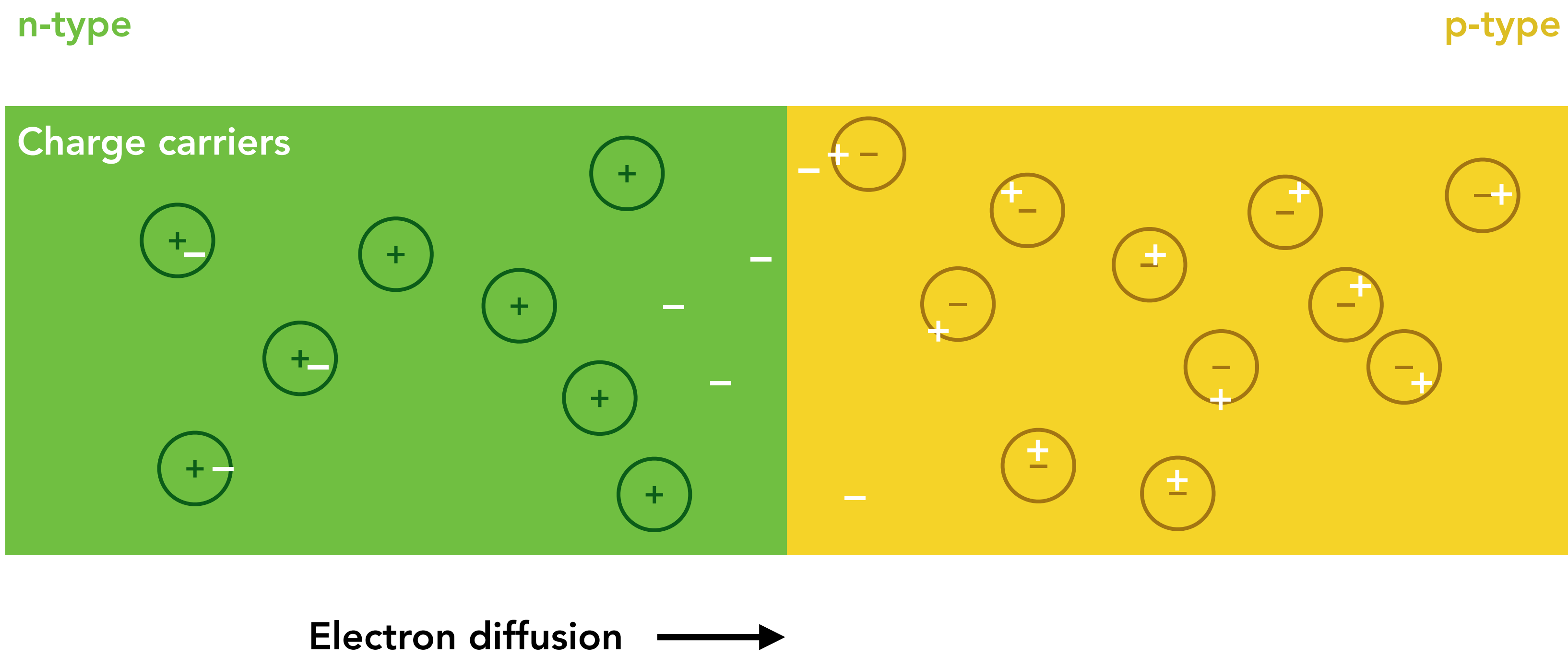
n-type

p-type

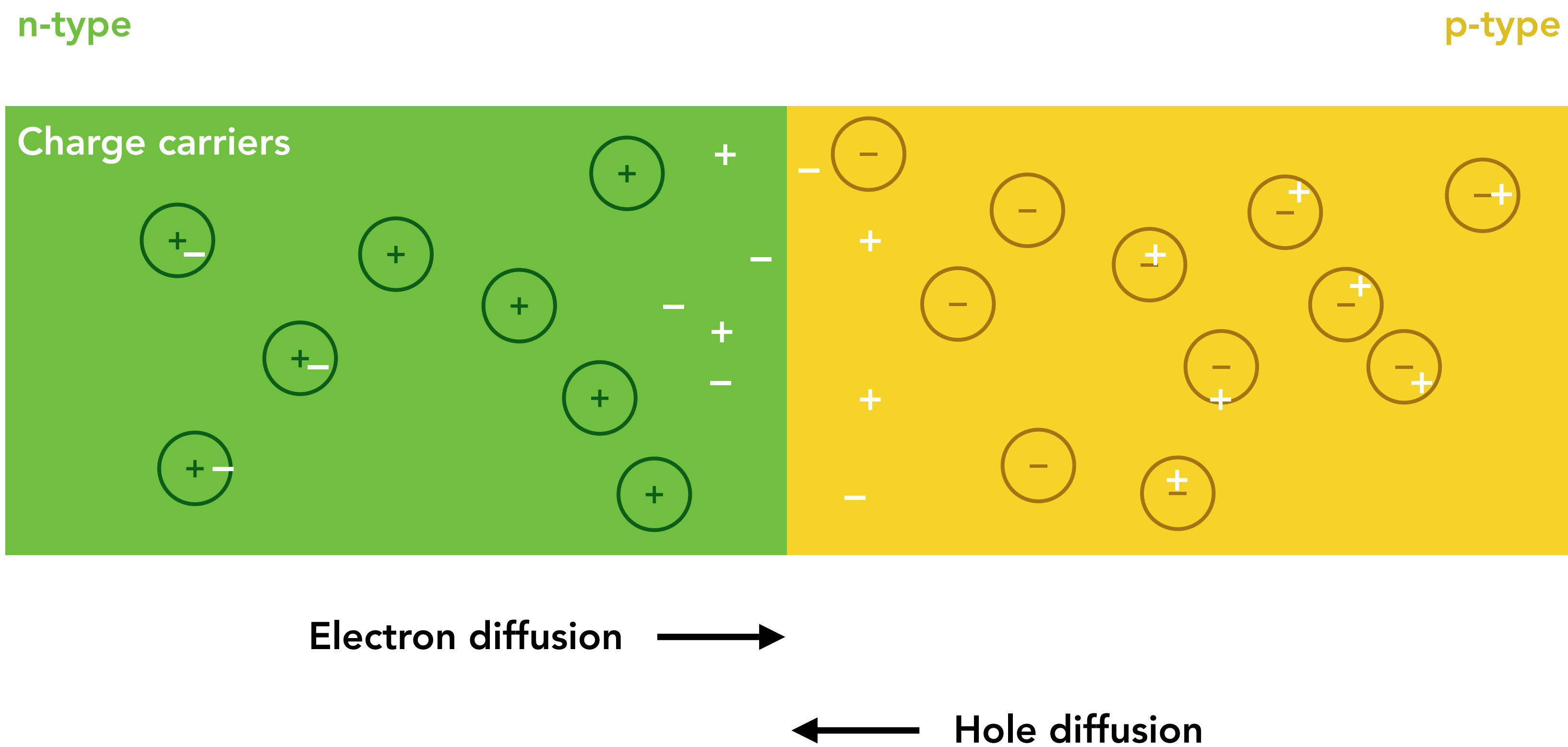


Electrically neutral system

P-N junctions - carrier movement



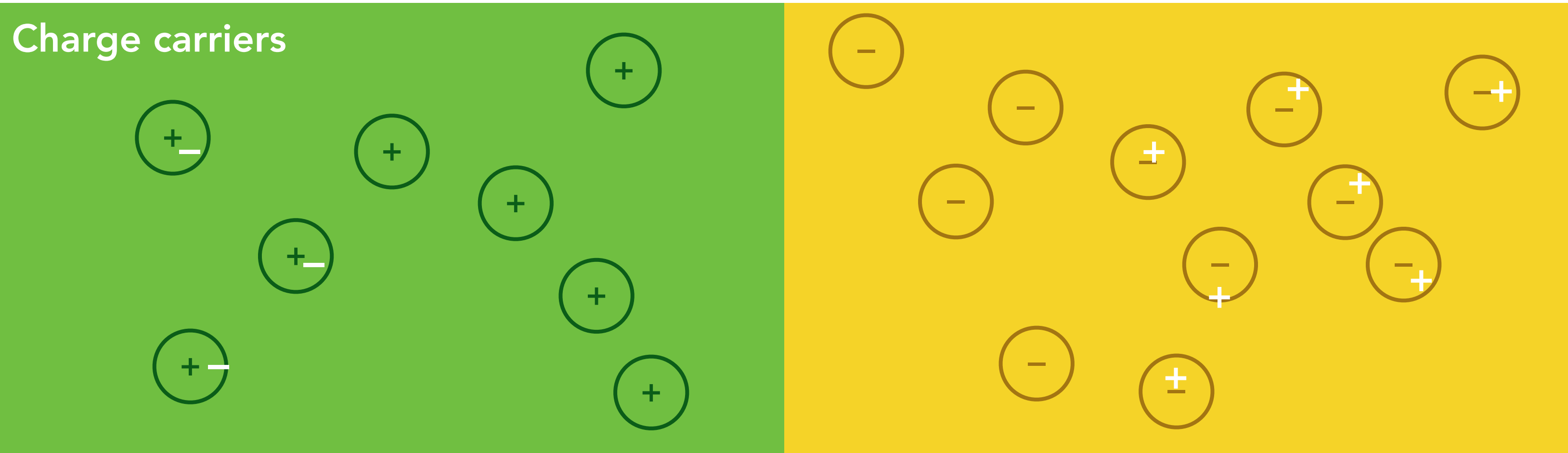
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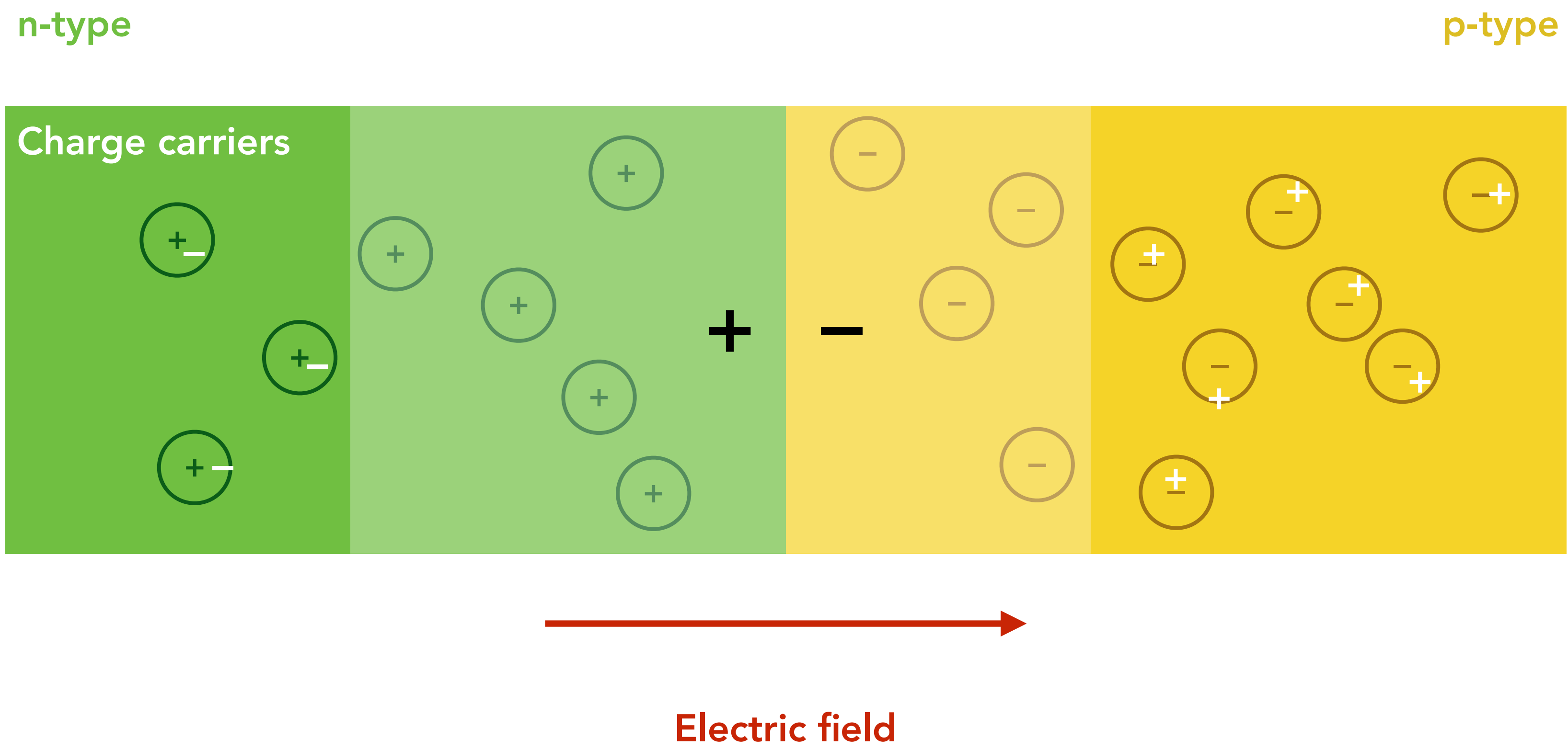
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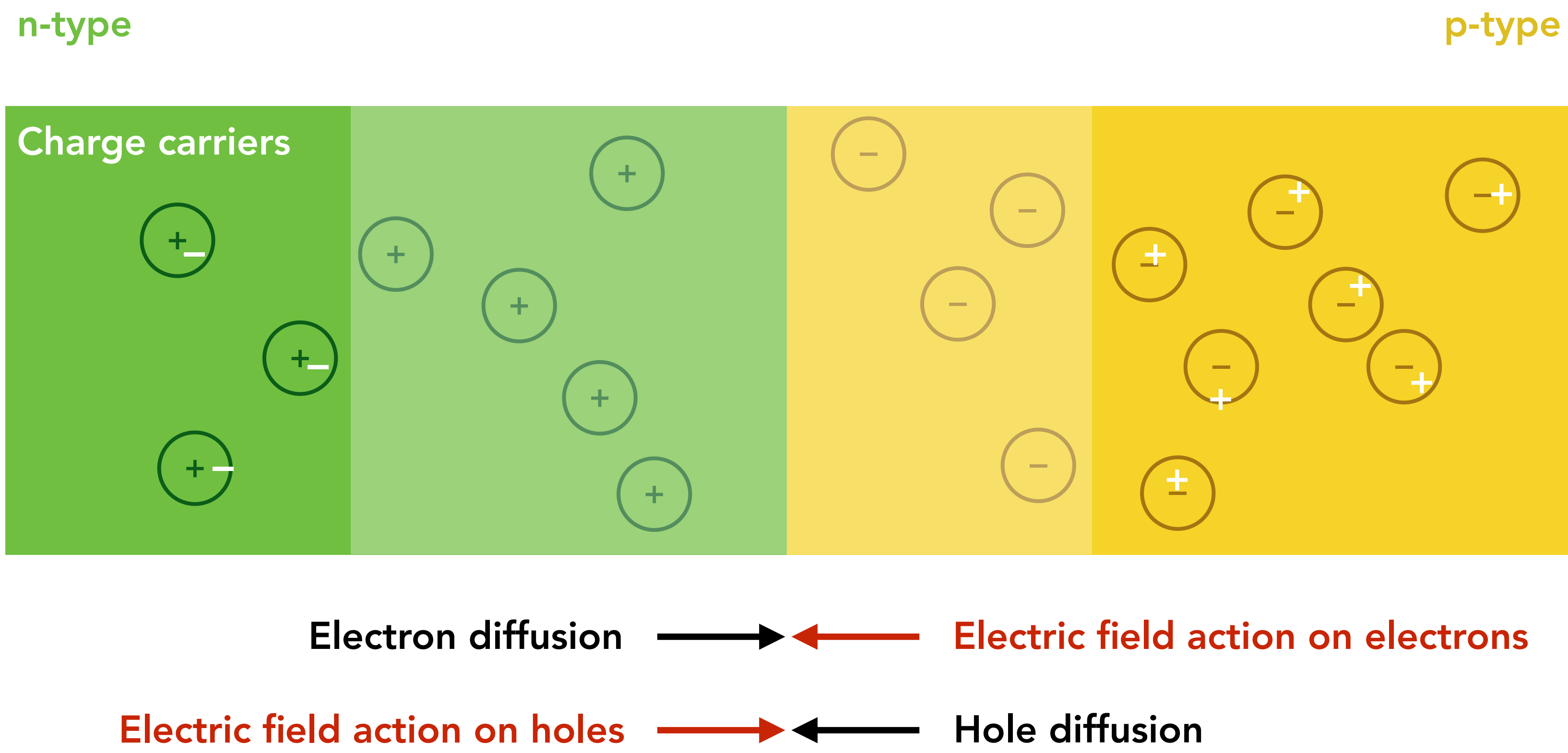


Recombination

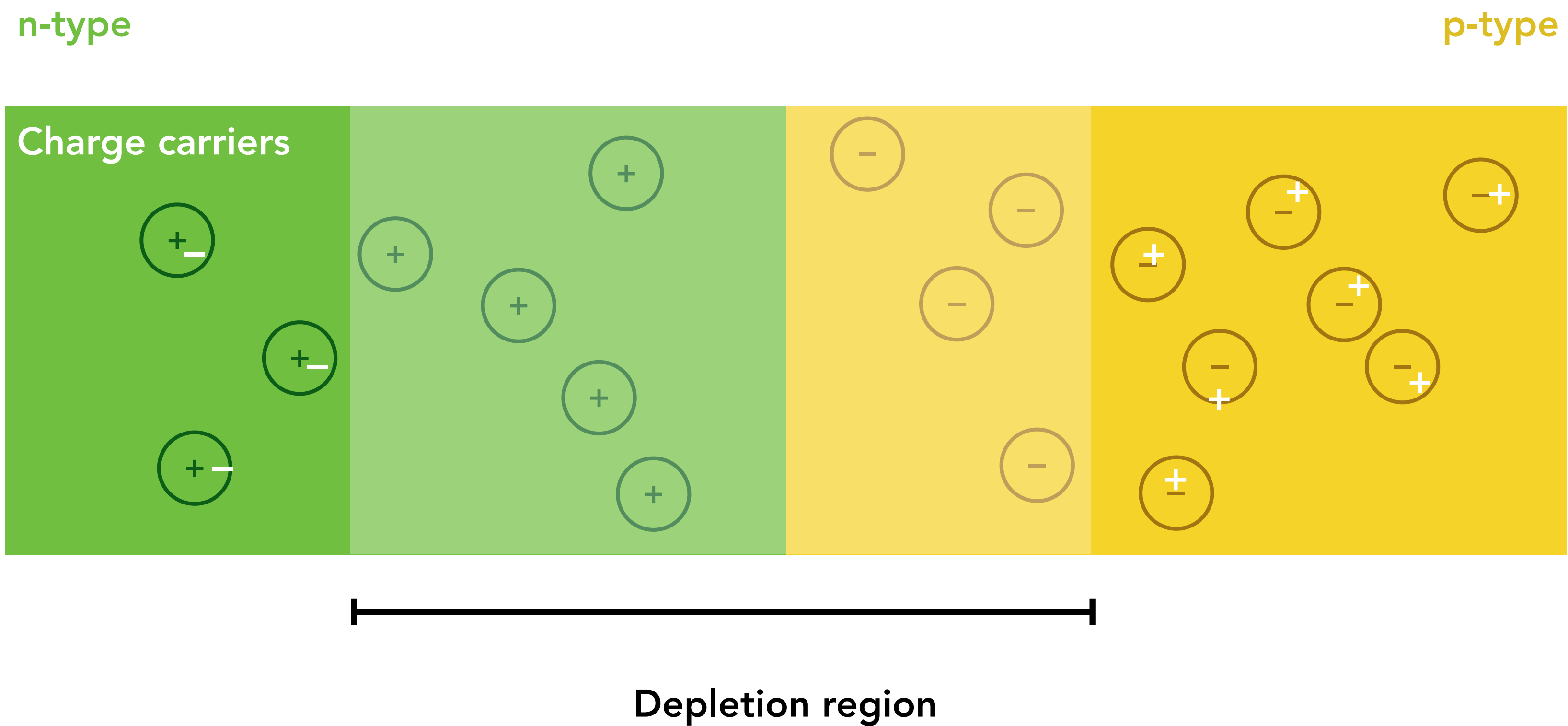
P-N junctions - carrier movement



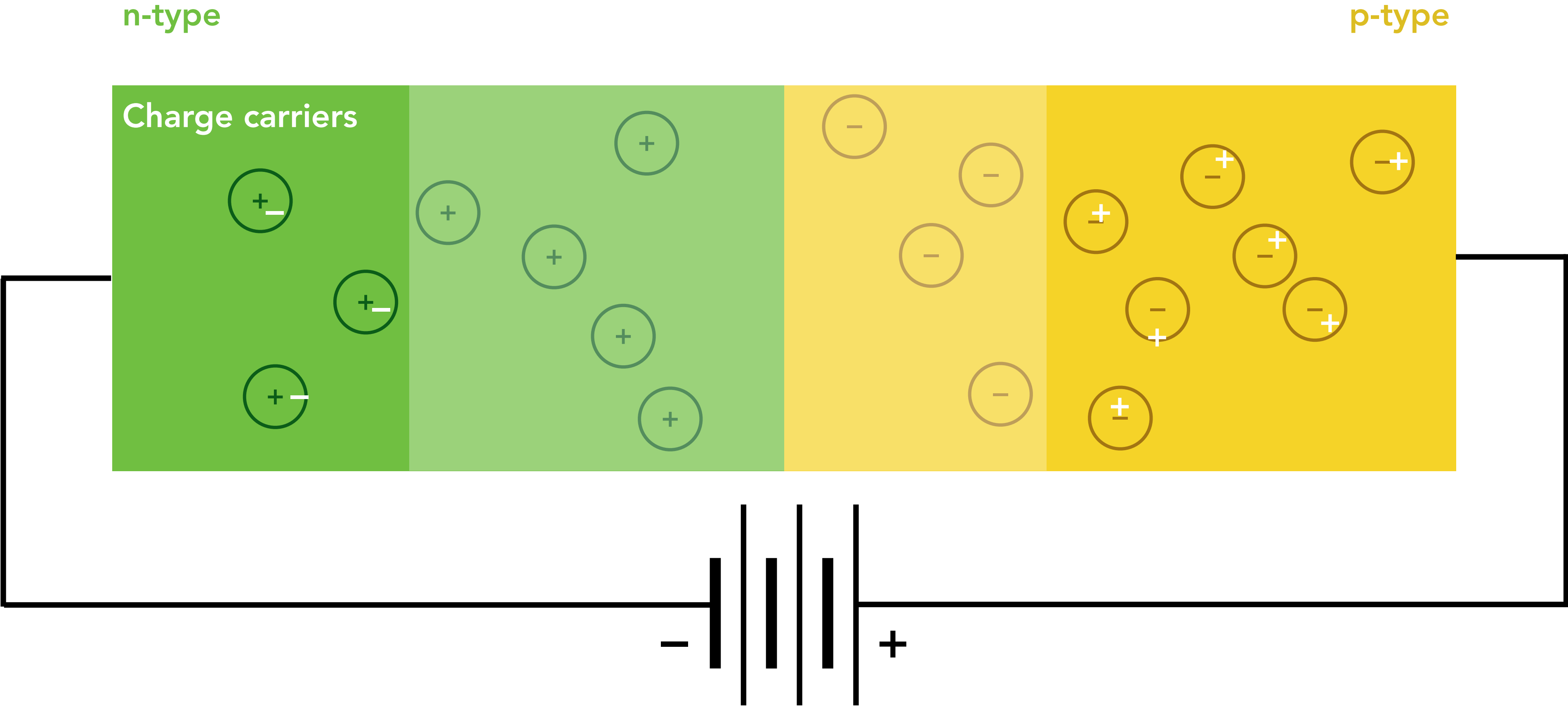
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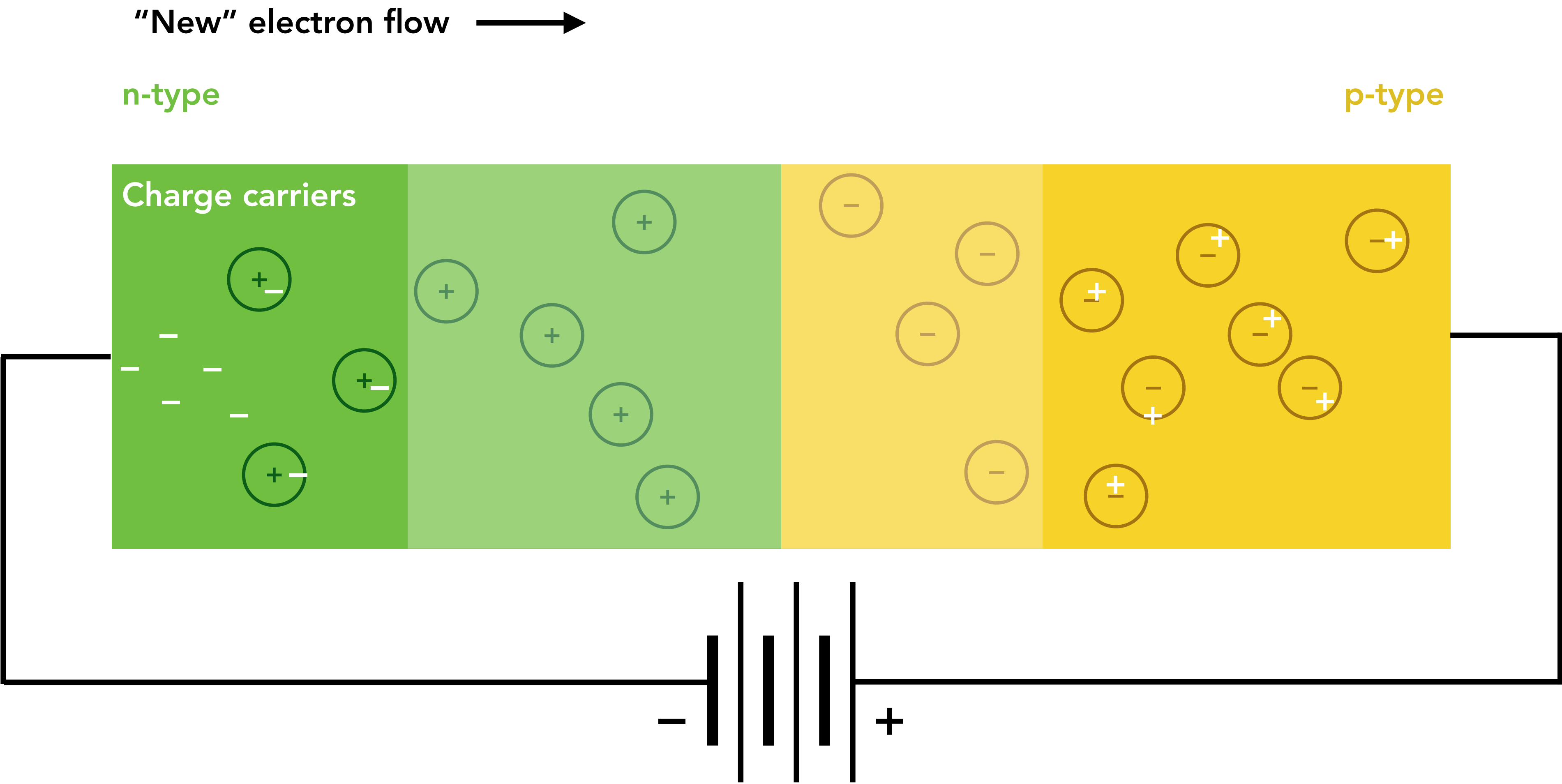
P-N junctions - built-in depletion



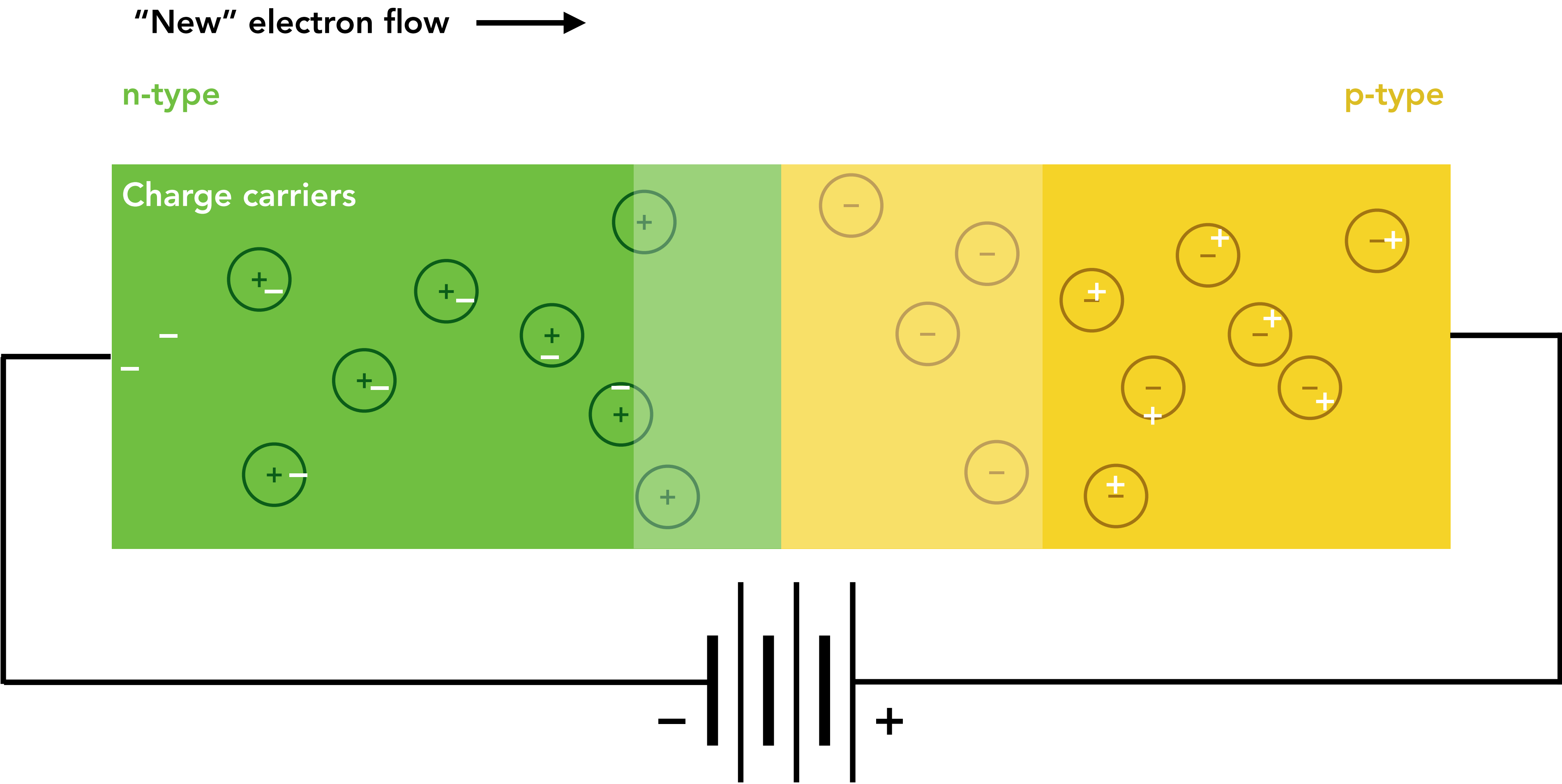
P-N junctions - forward bias



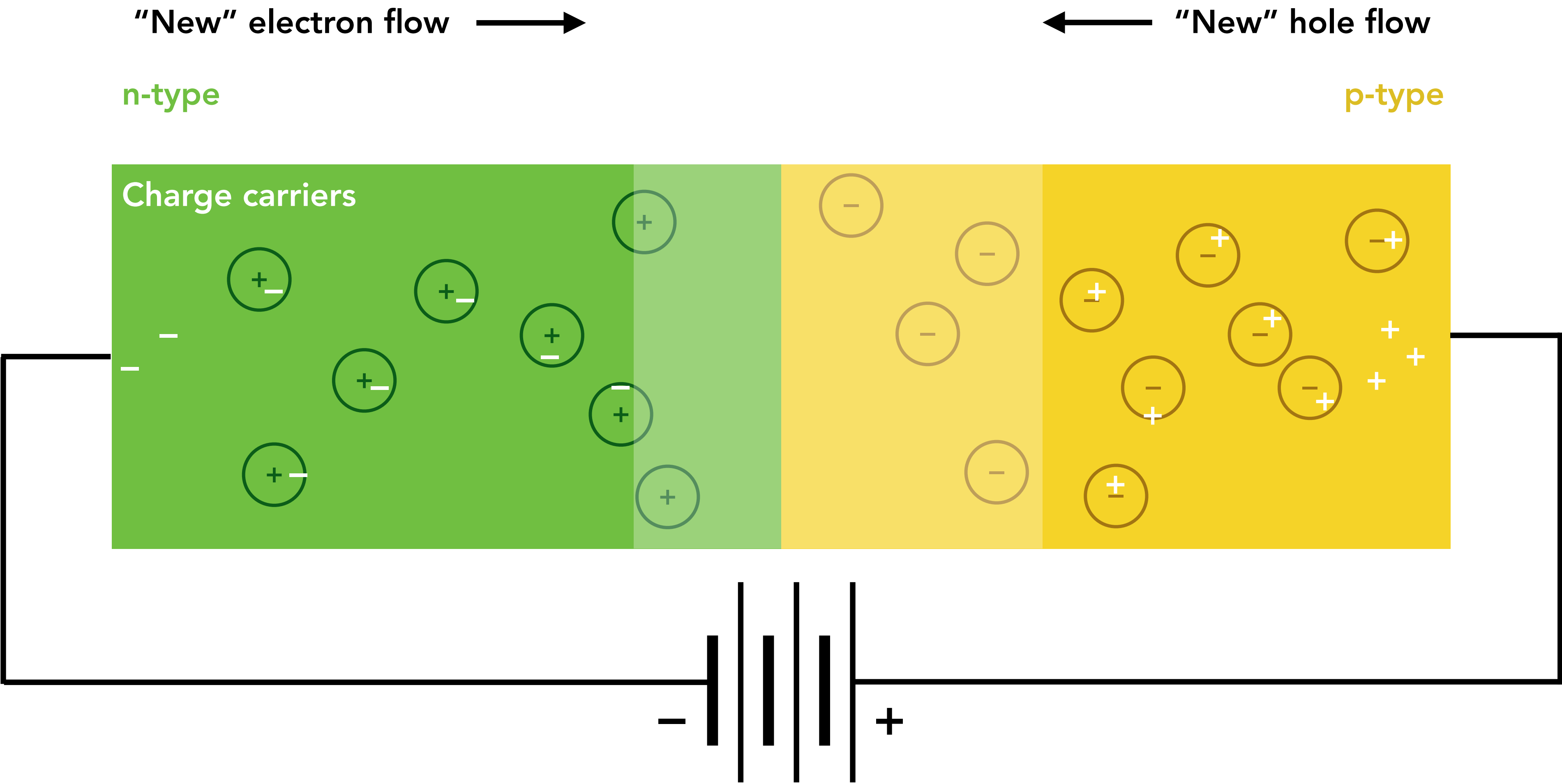
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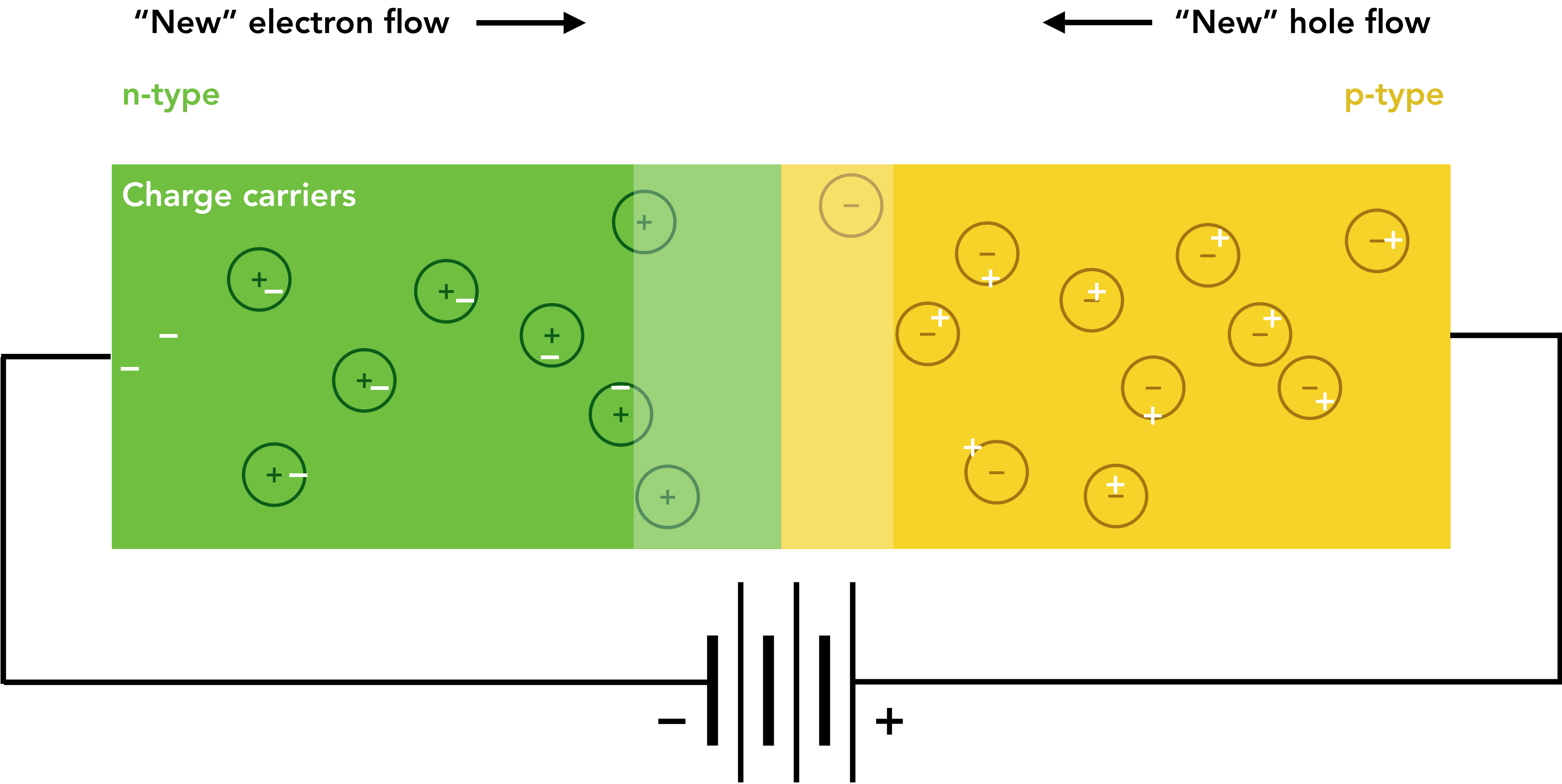
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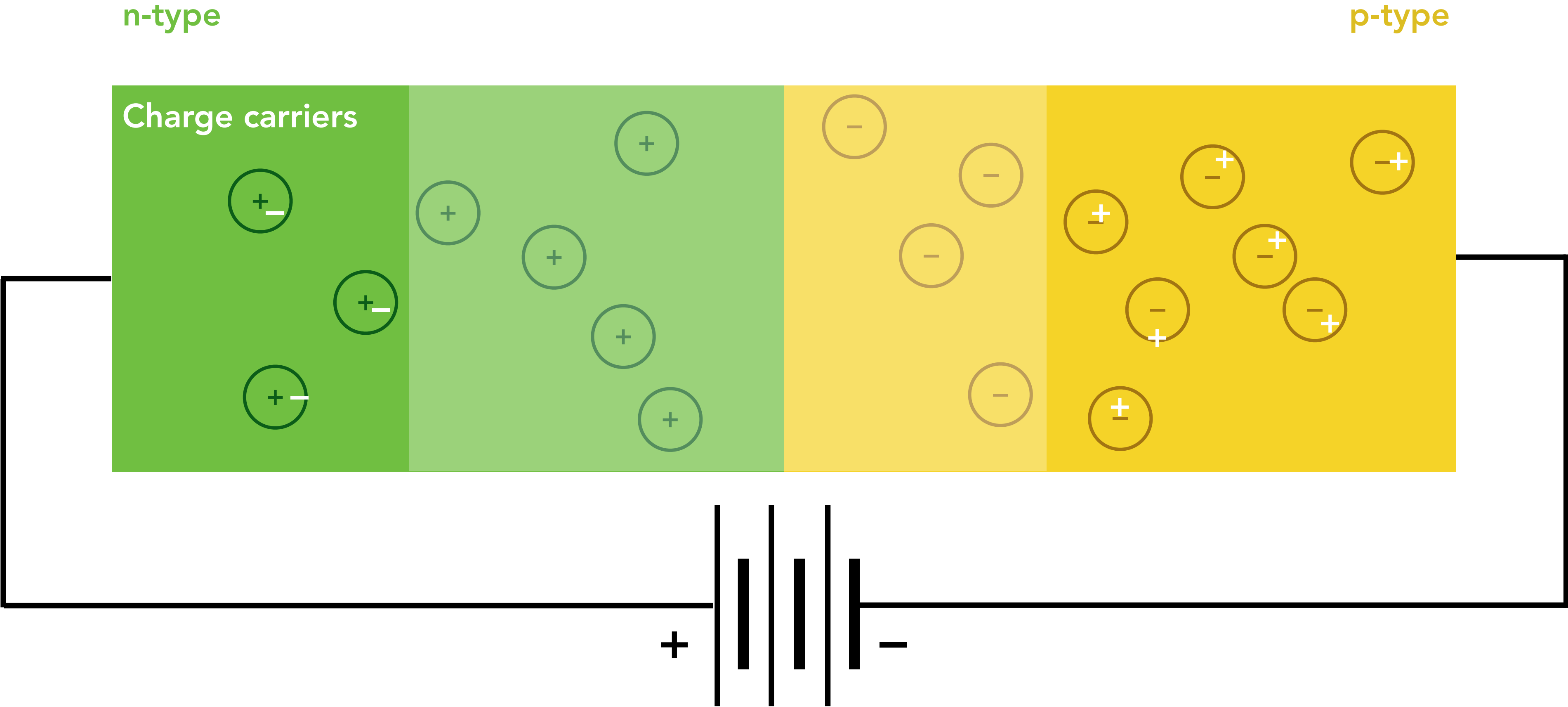
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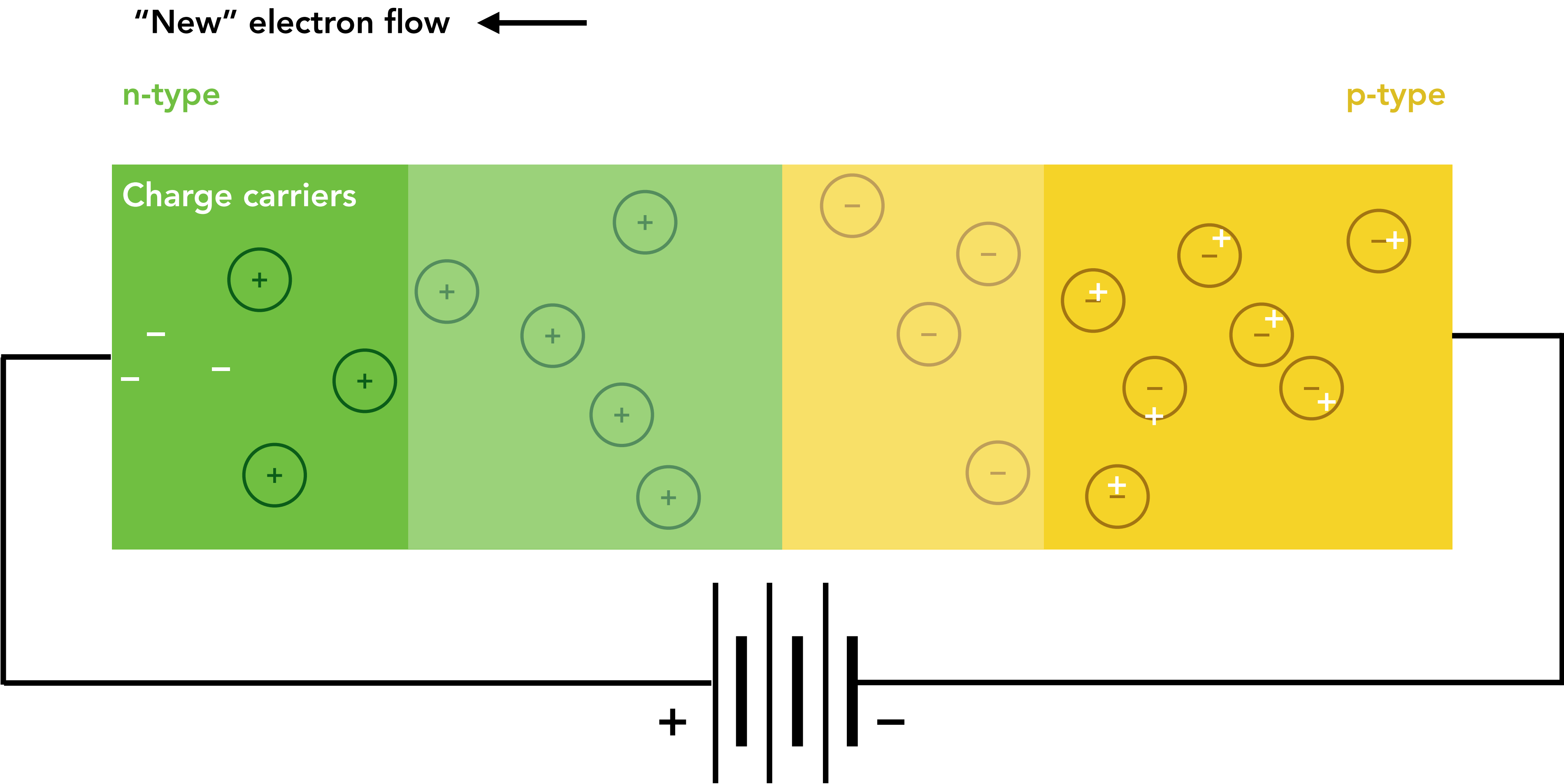
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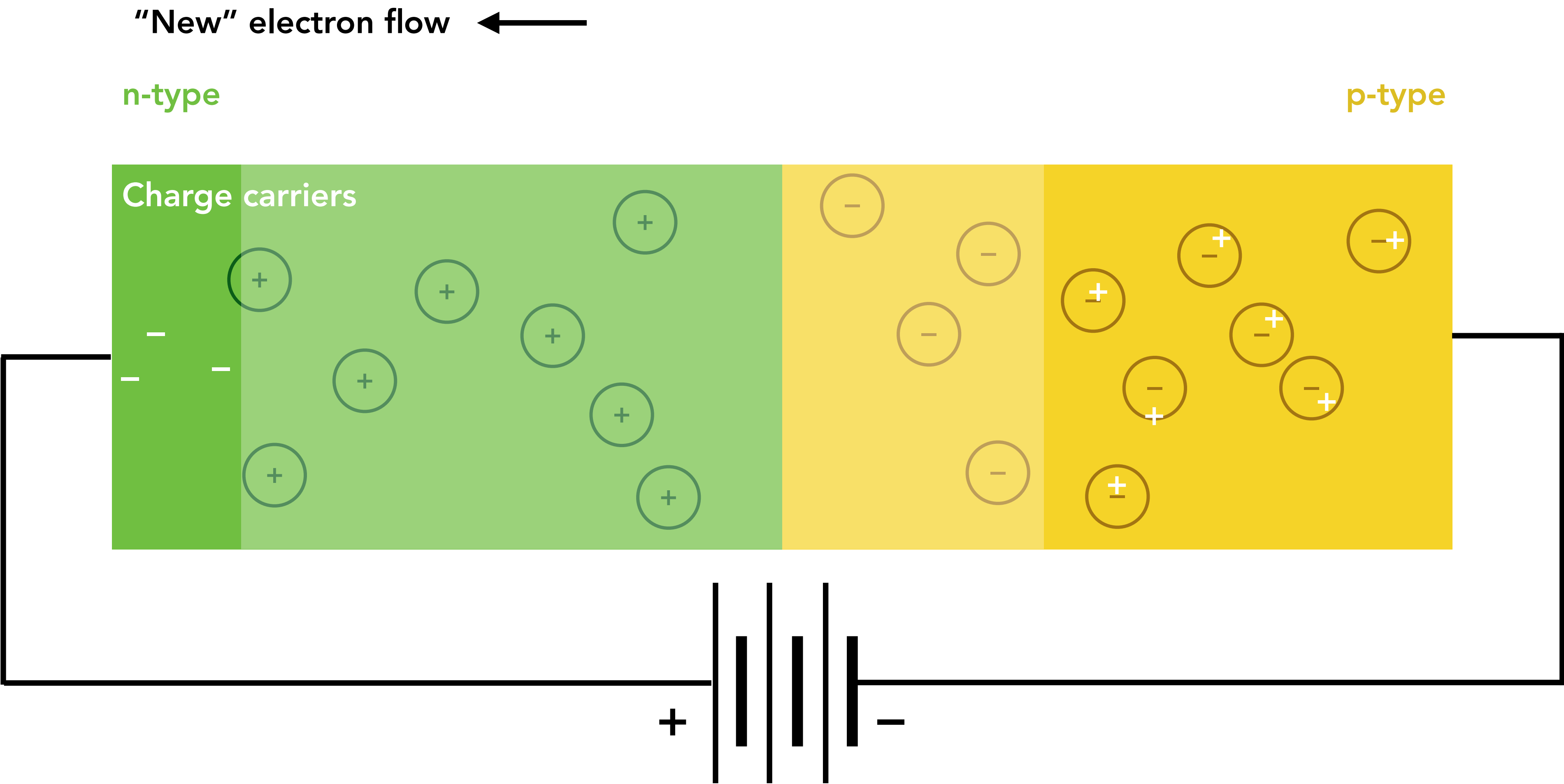
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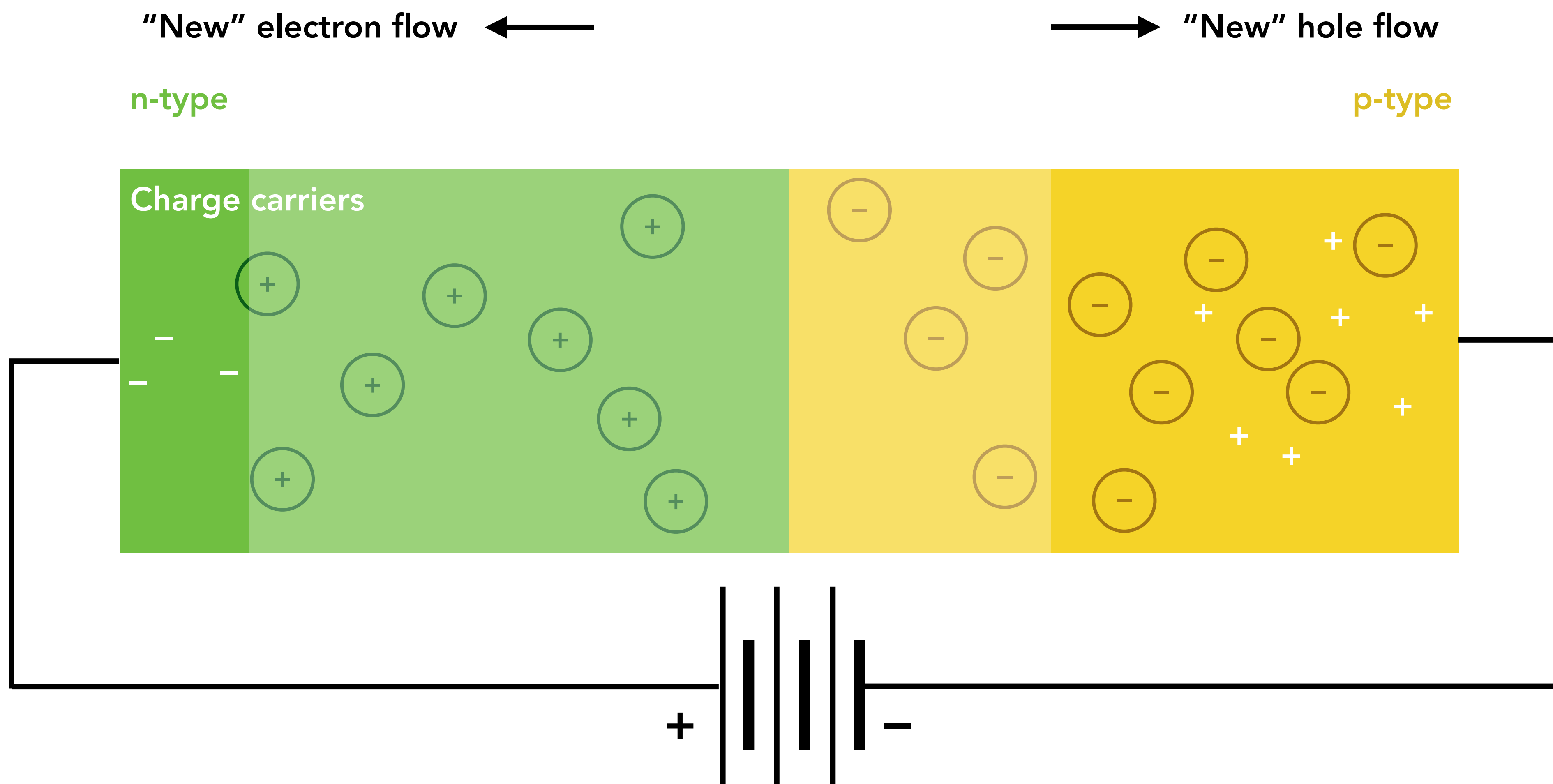
P-N junctions - reverse bias



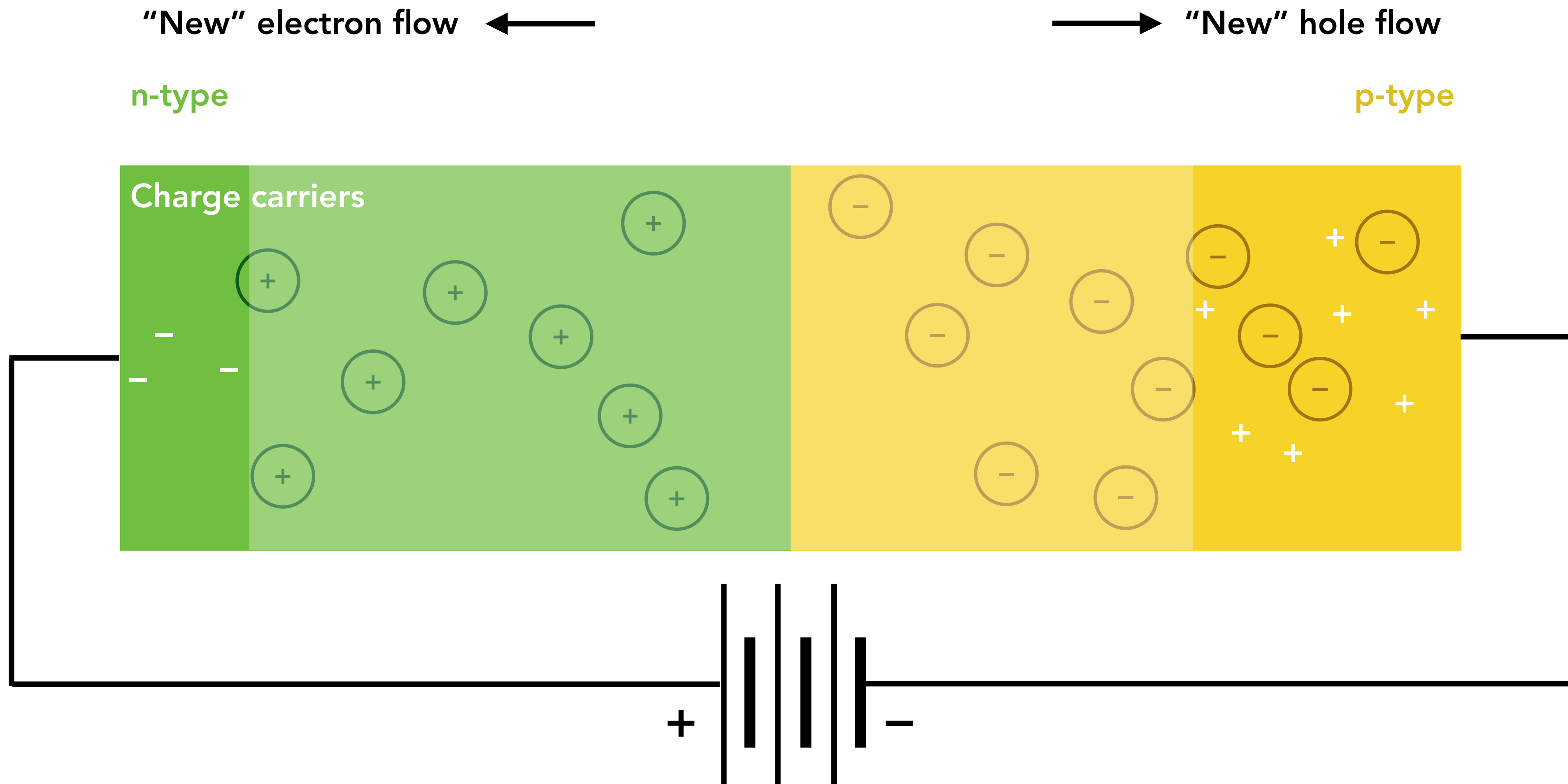
P-N junctions - reverse bias



P-N junctions - reverse bias



P-N junctions - reverse bias



P-N junctions - external field

Bringing together p-doped and n-doped semiconductors will give rise to a natural region free of charge carriers between the two

- This built-in depletion region contains an electric field which counters the diffusion of more carriers across the junction

We can forward bias the junction to reduce the depletion width

- Current will flow through the junction once the depletion width has been reduced to 0

We can reverse bias the junction to increase the depletion width

- If one of the regions has no carriers left then it can be said to be fully depleted



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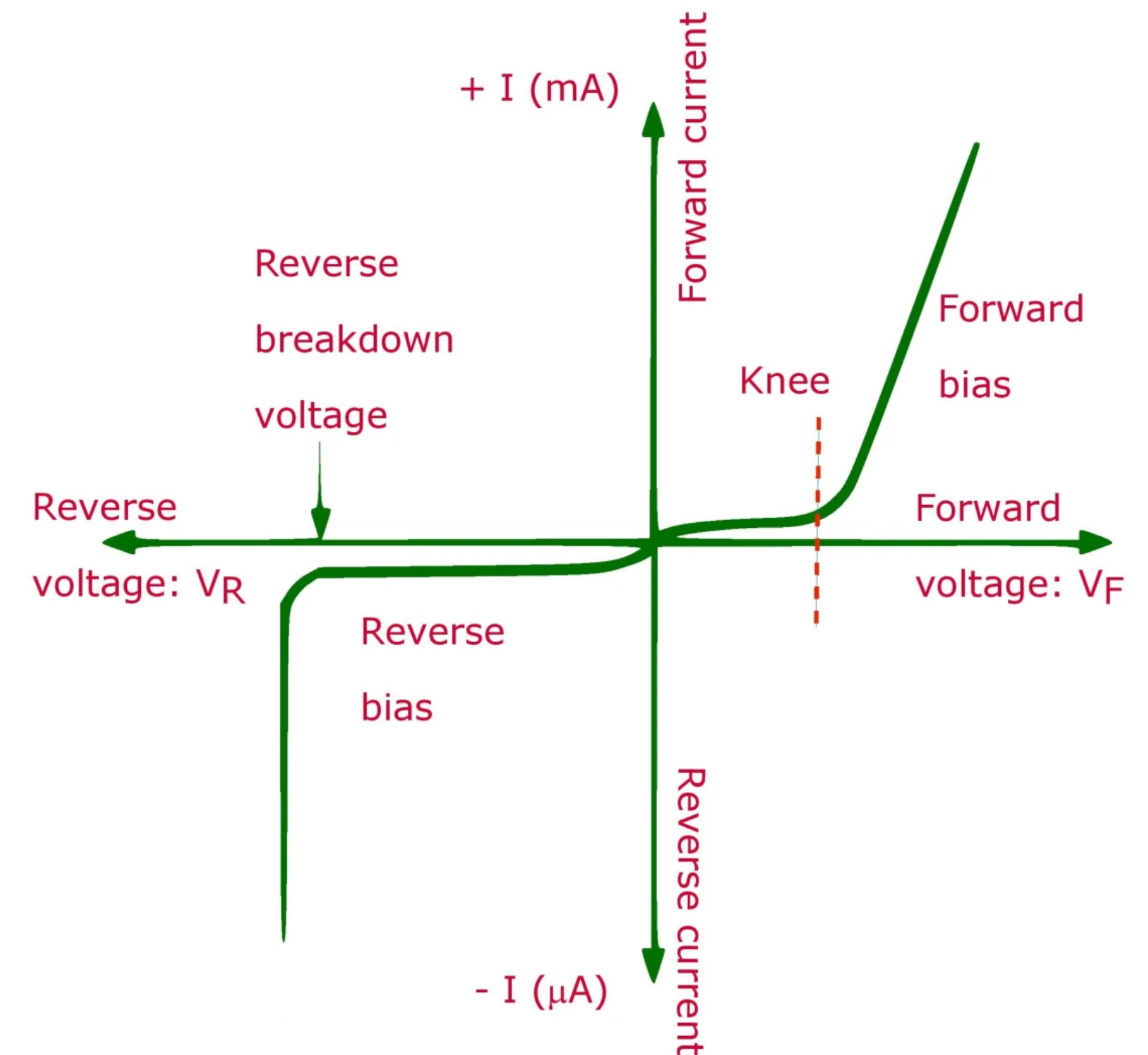
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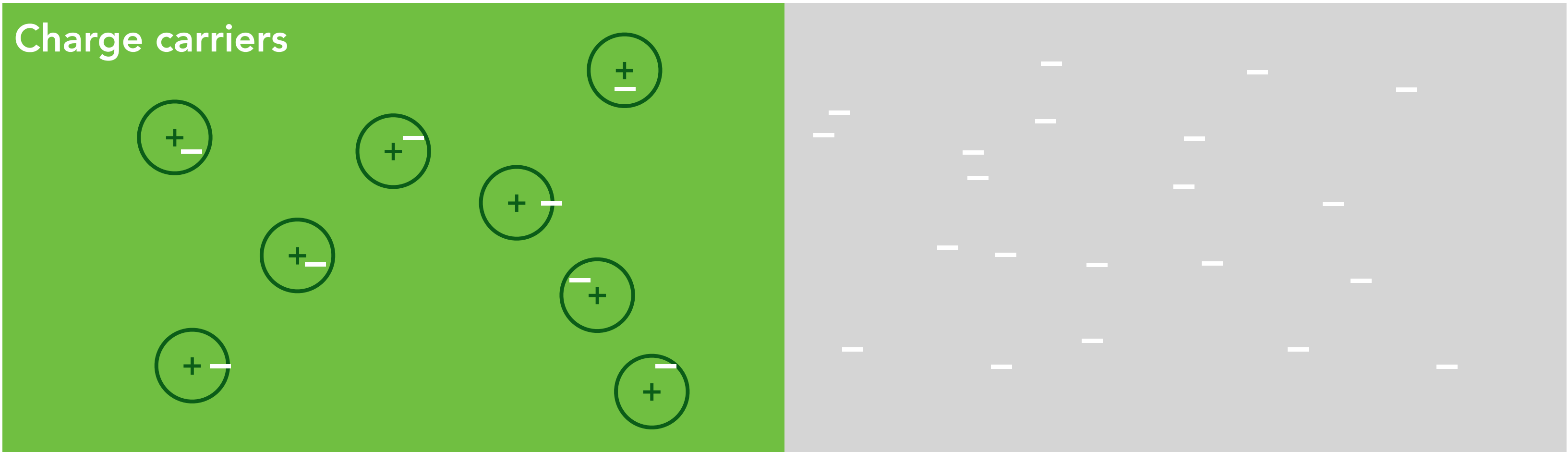
**The Touching Place of Hard Shiny Rocks with Magic Rocks that
Only Sometimes Let Things Through**

(Semiconductor-metal junctions)

S-M junctions

n-type

metal

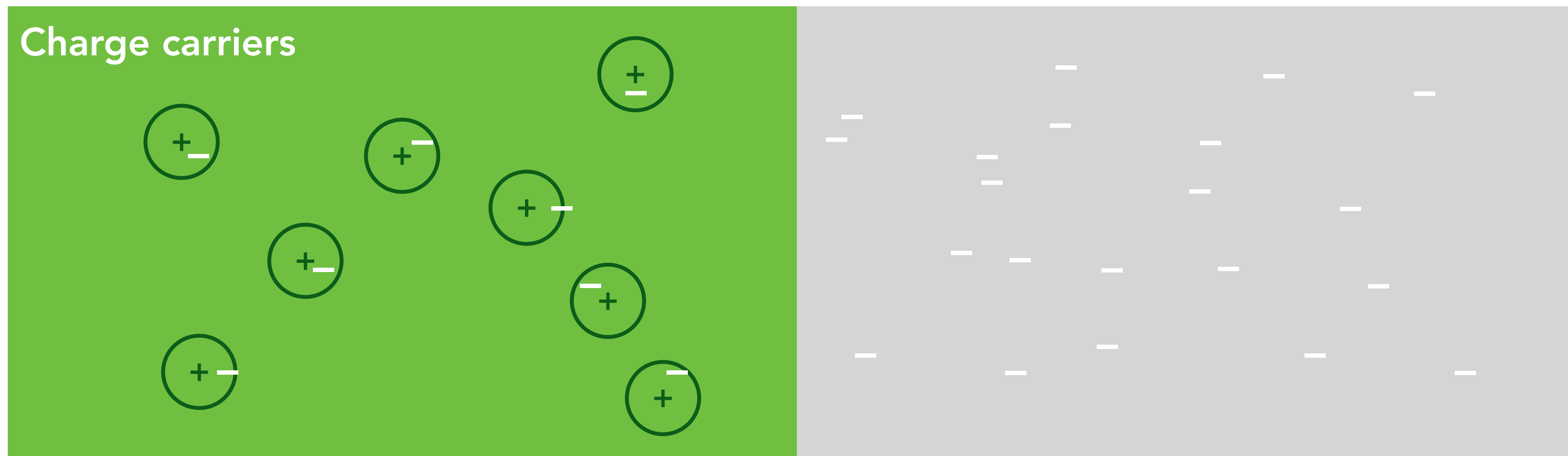


Question: how hard is it to move electrons from (say) the metal to the semiconductor?

S-M junctions

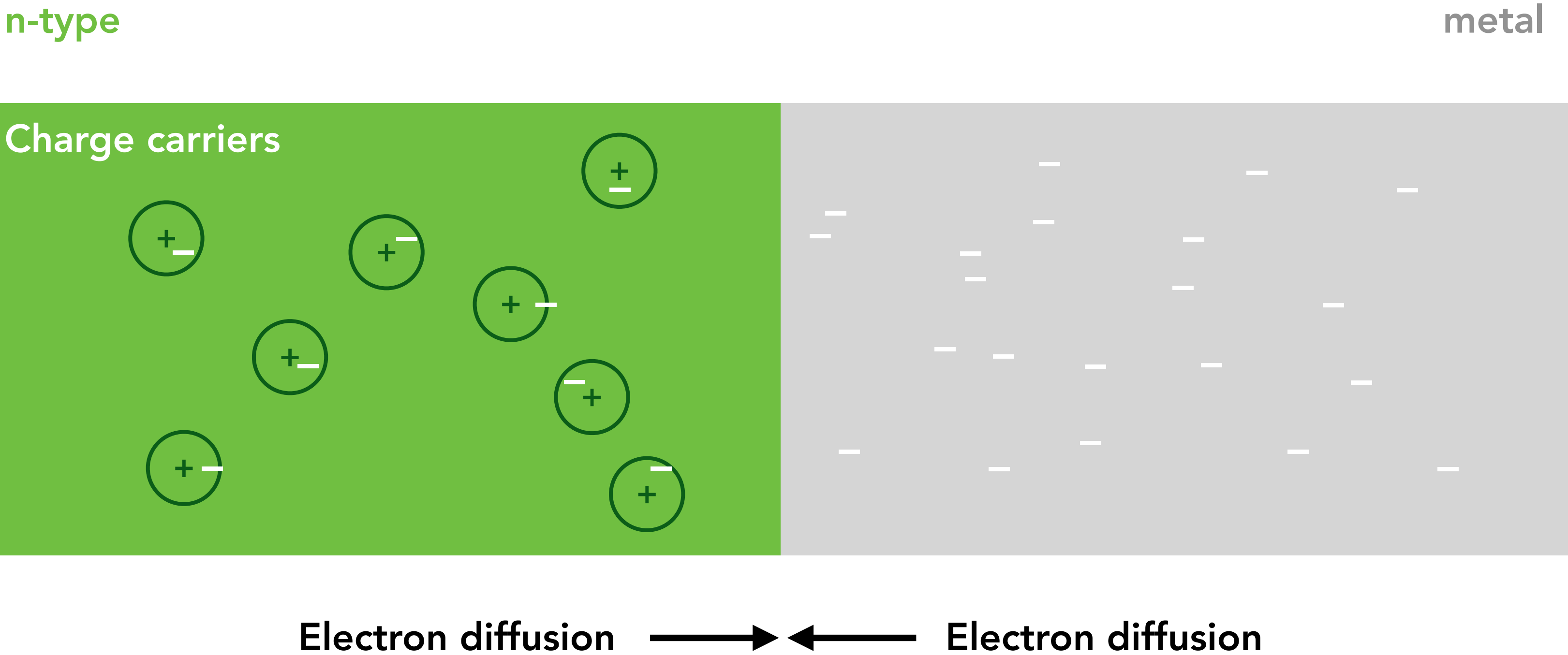
n-type

metal

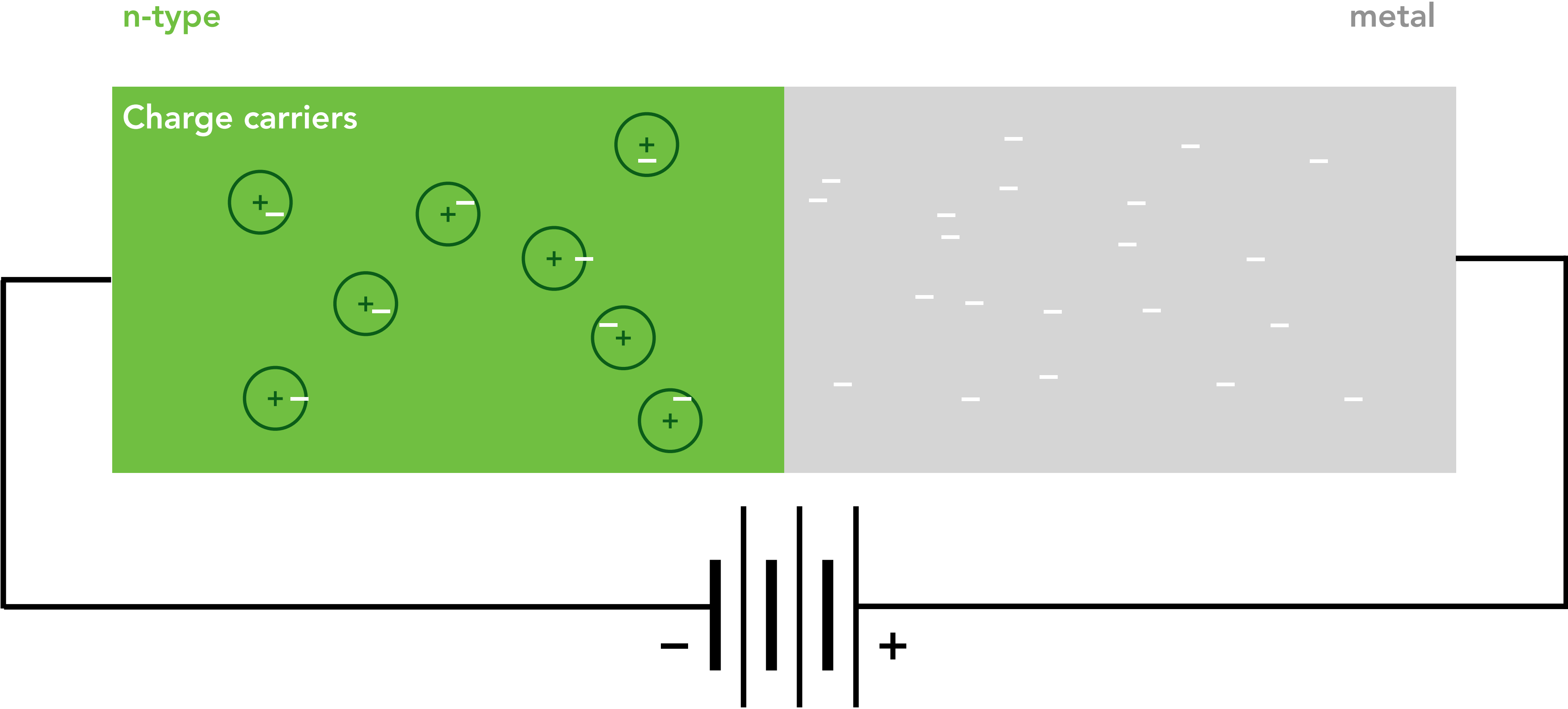


Consider the case where it is **easy** to move from the metal to the semiconductor

S-M junctions - carrier movement



S-M junctions - forward bias

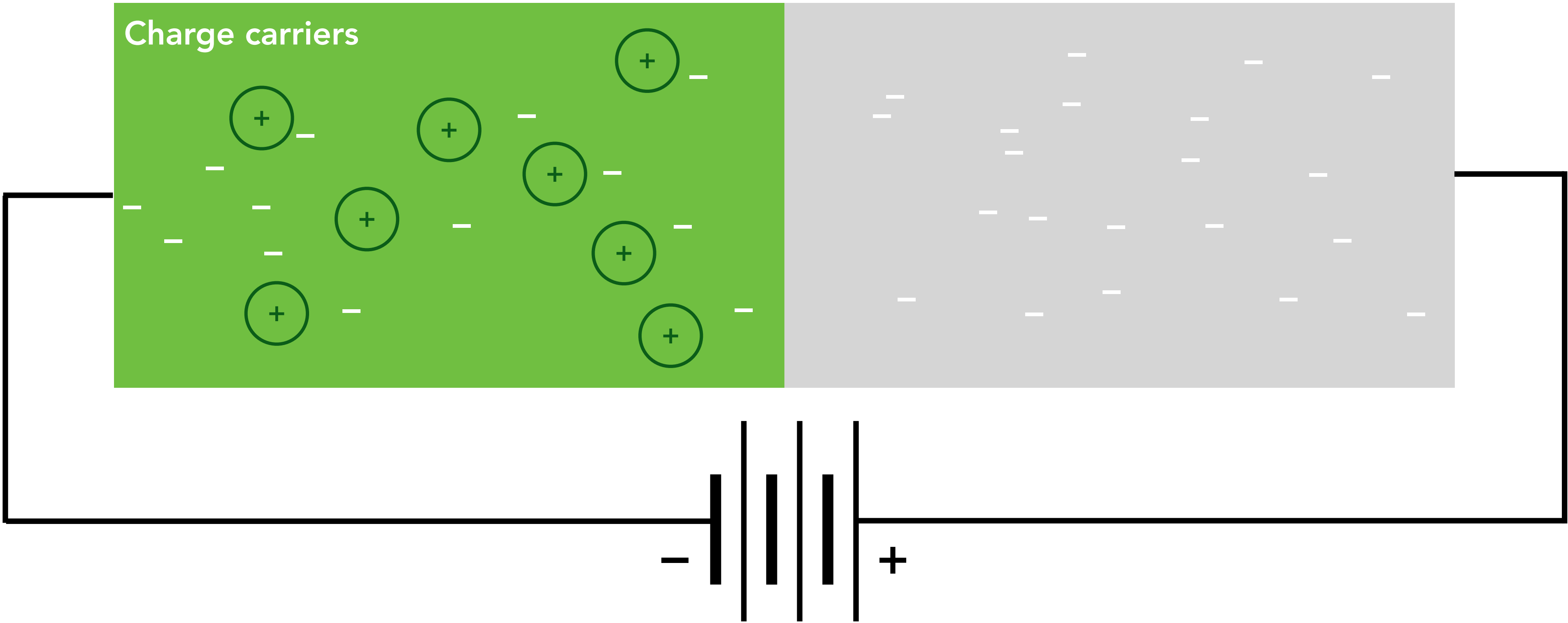


S-M junctions - forward bias

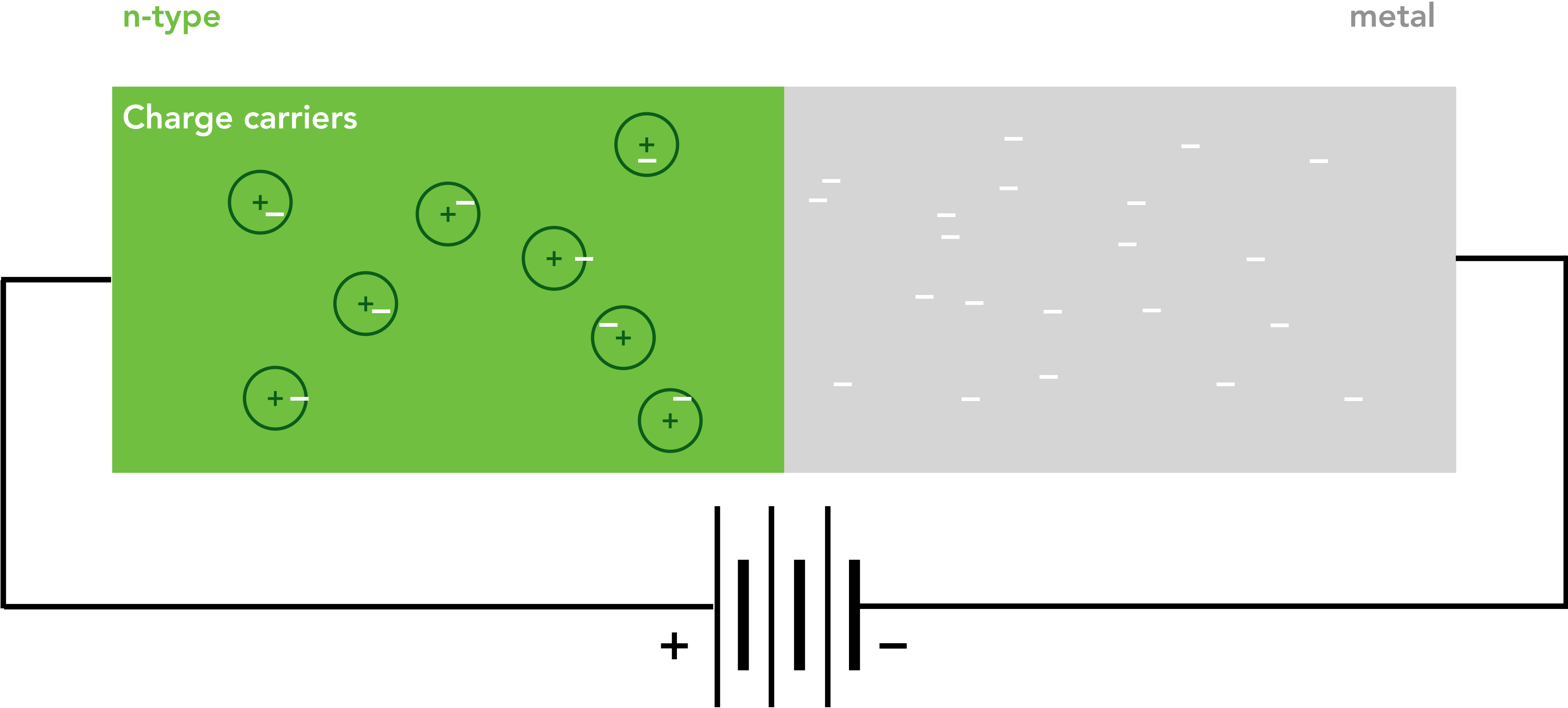
“New” electron flow →

n-type

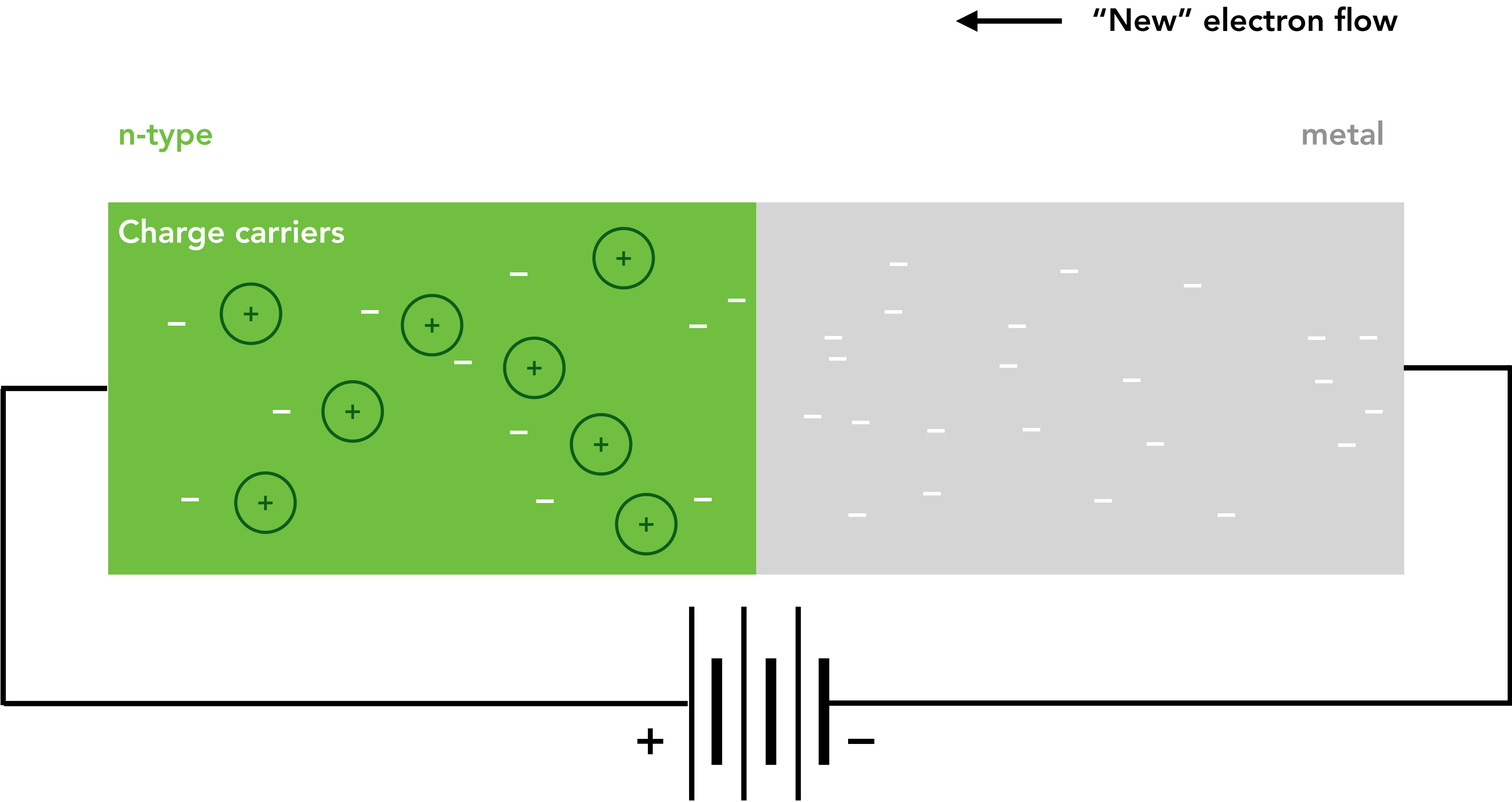
metal



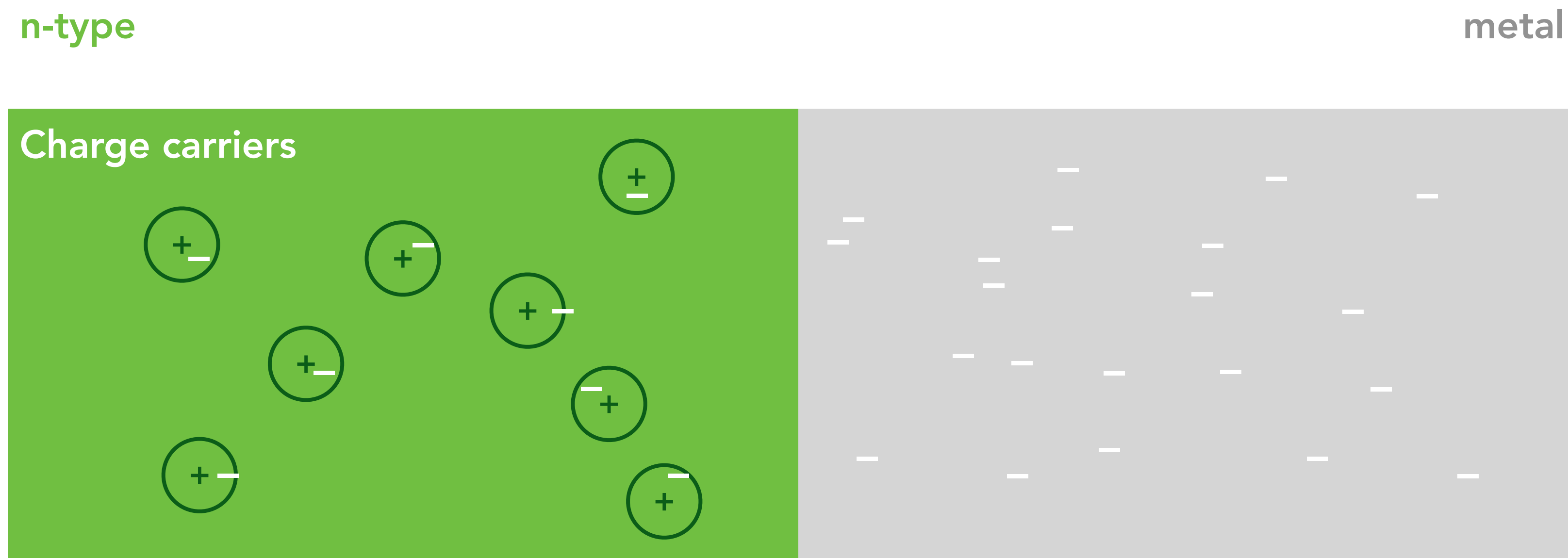
S-M junctions - reverse bias



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S-M junctions



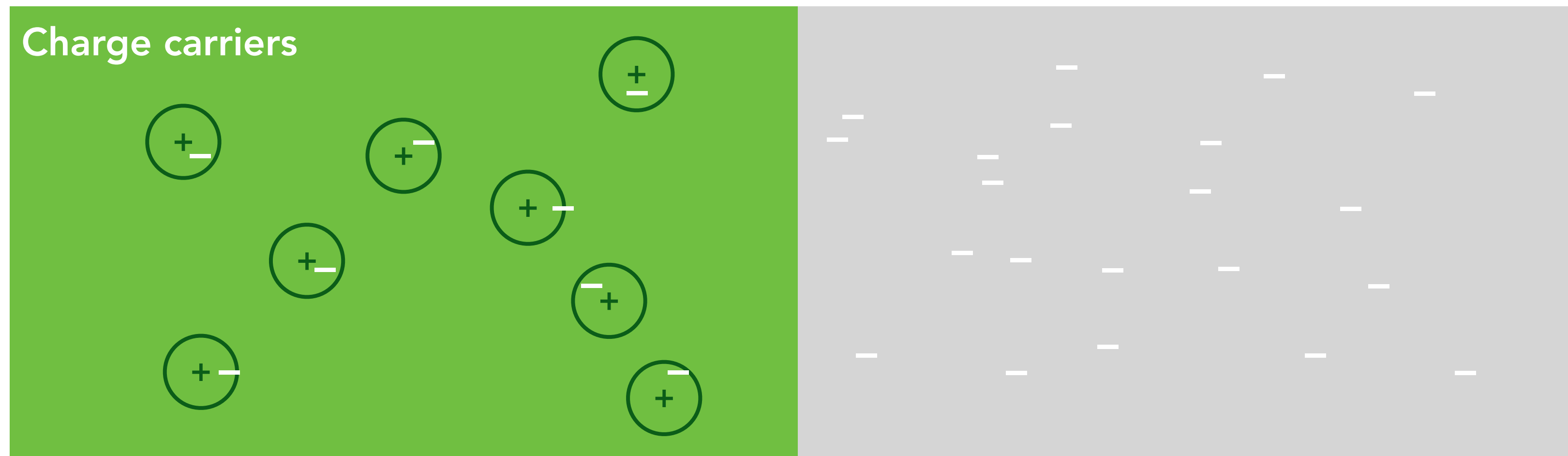
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An **Ohmic contact** is created

S-M junctions

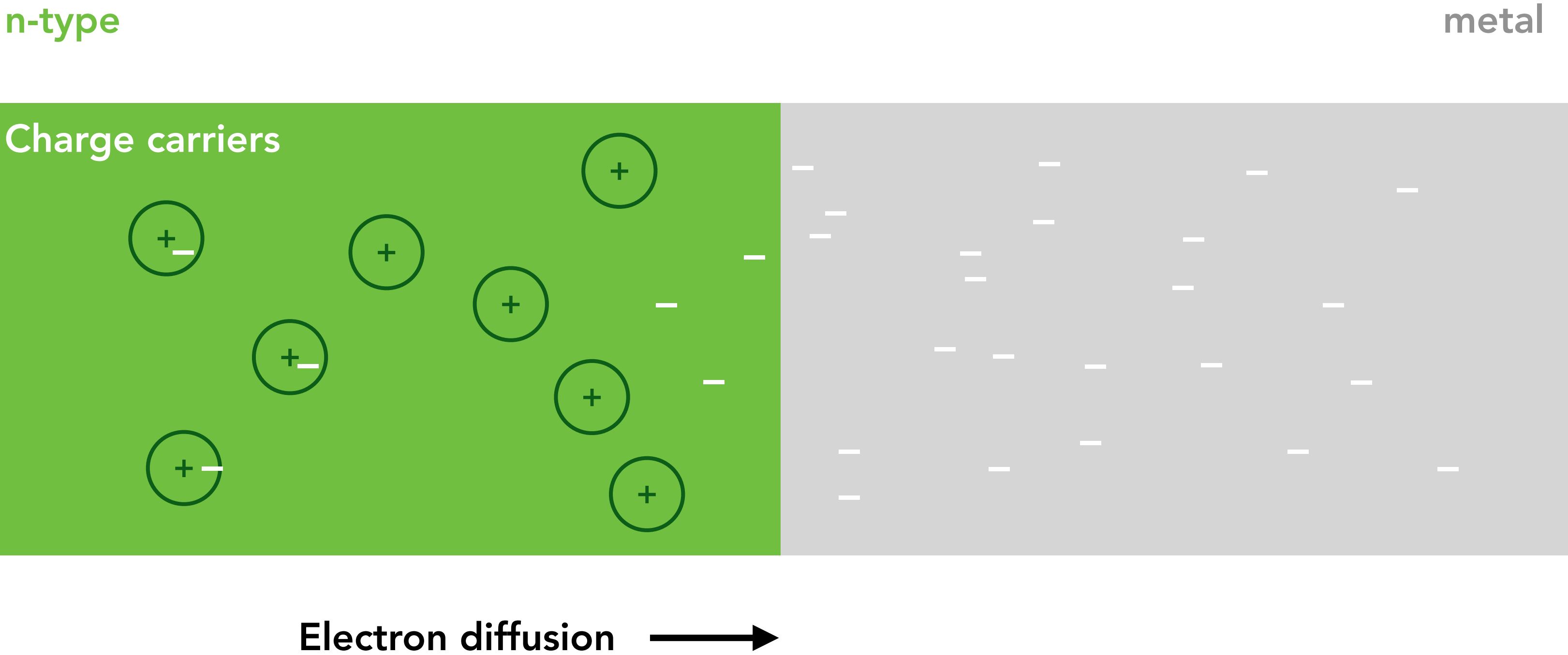
n-type

metal

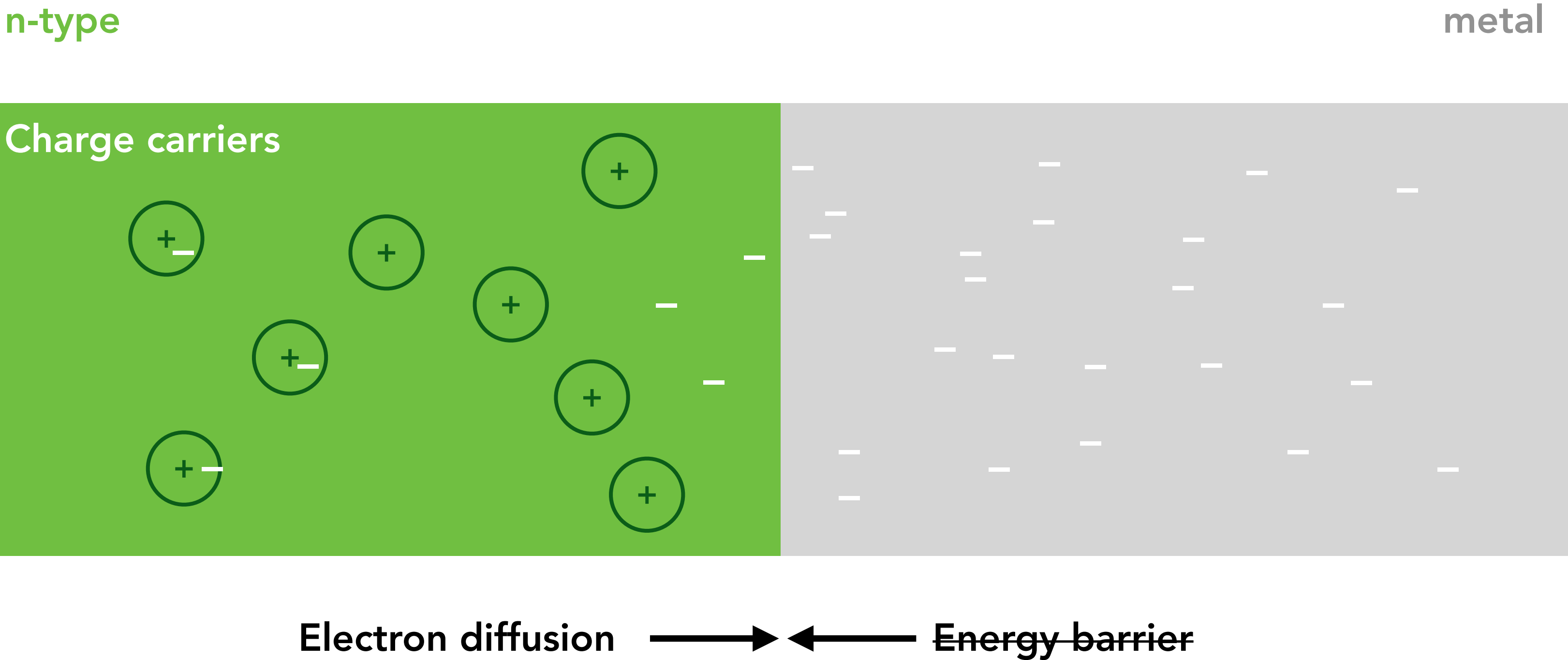


Consider instead the case where it is **difficult** to move from the metal to the semiconductor

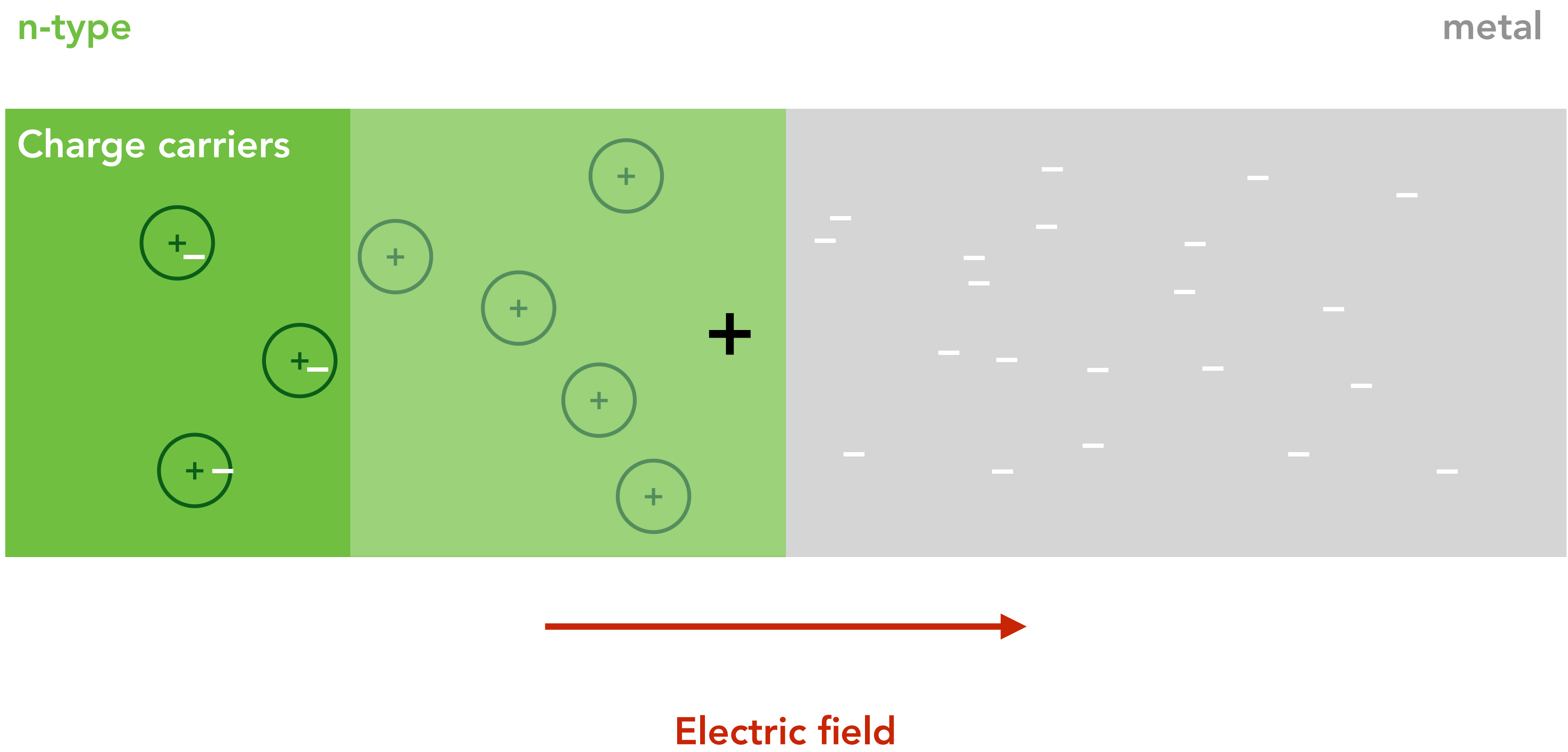
S-M junctions - carrier movement



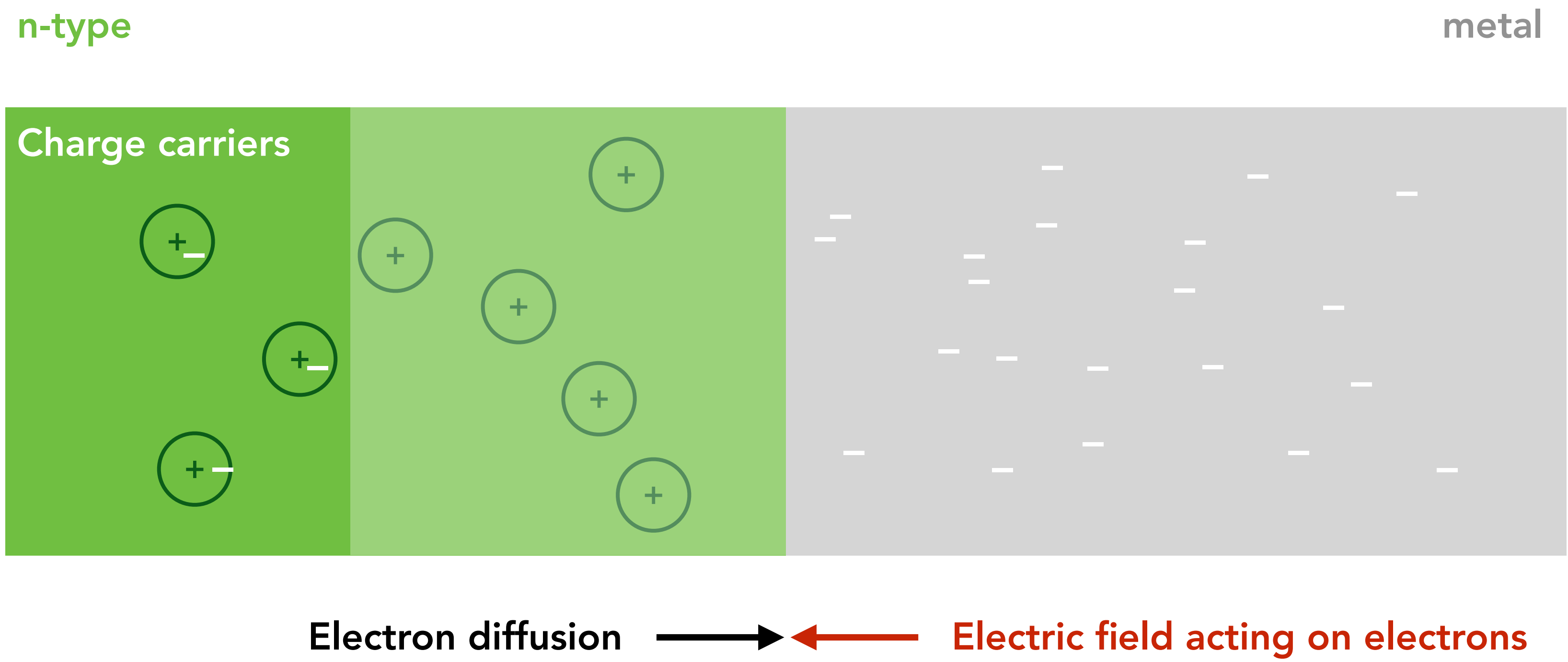
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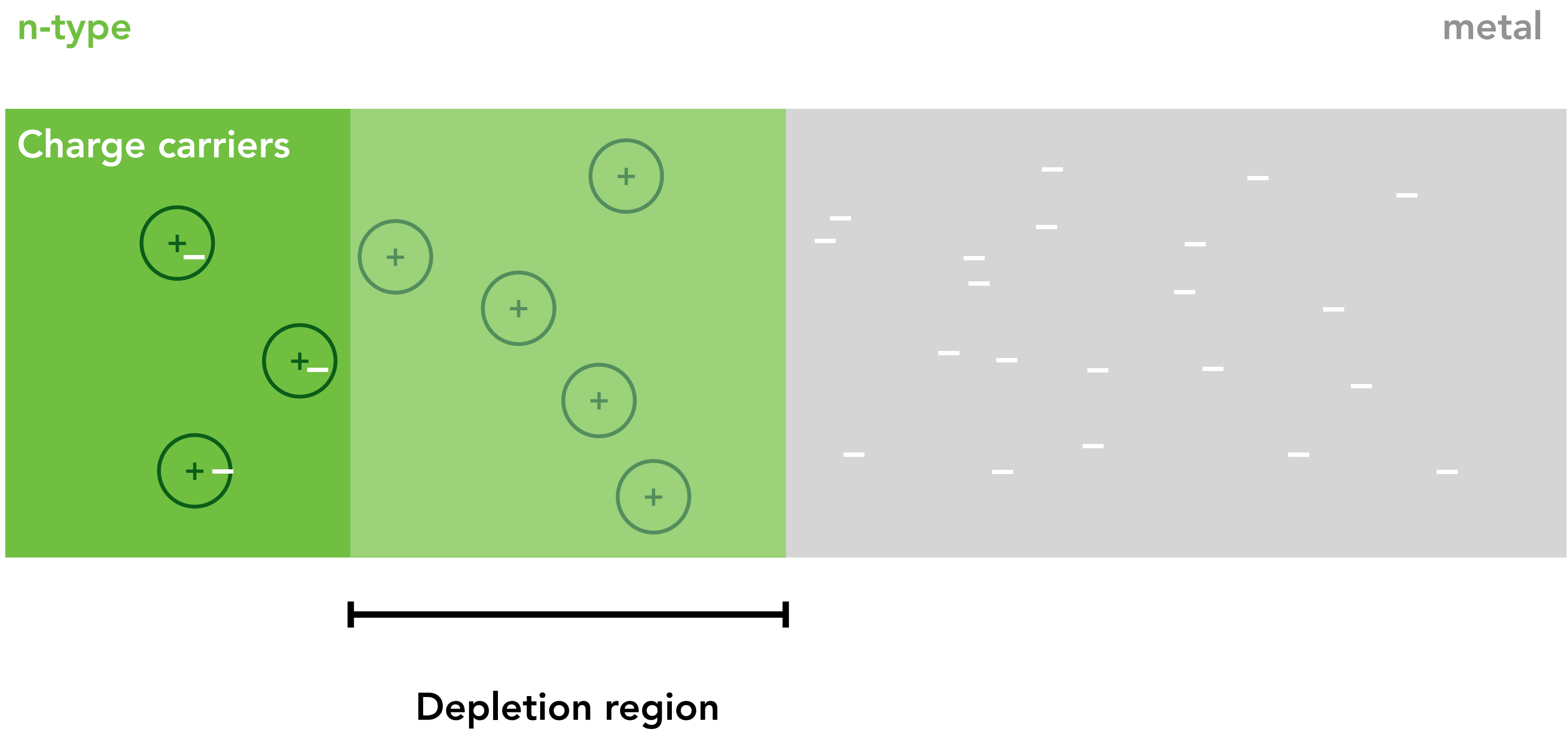
S-M junctions - carrier movement



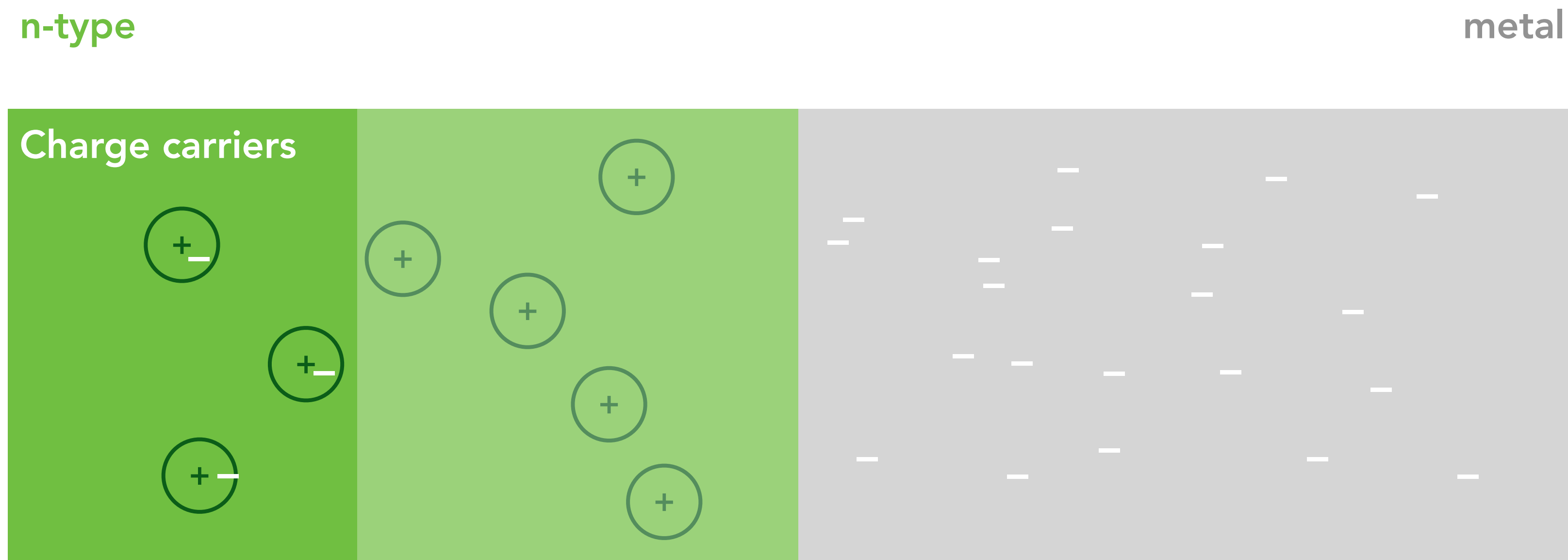
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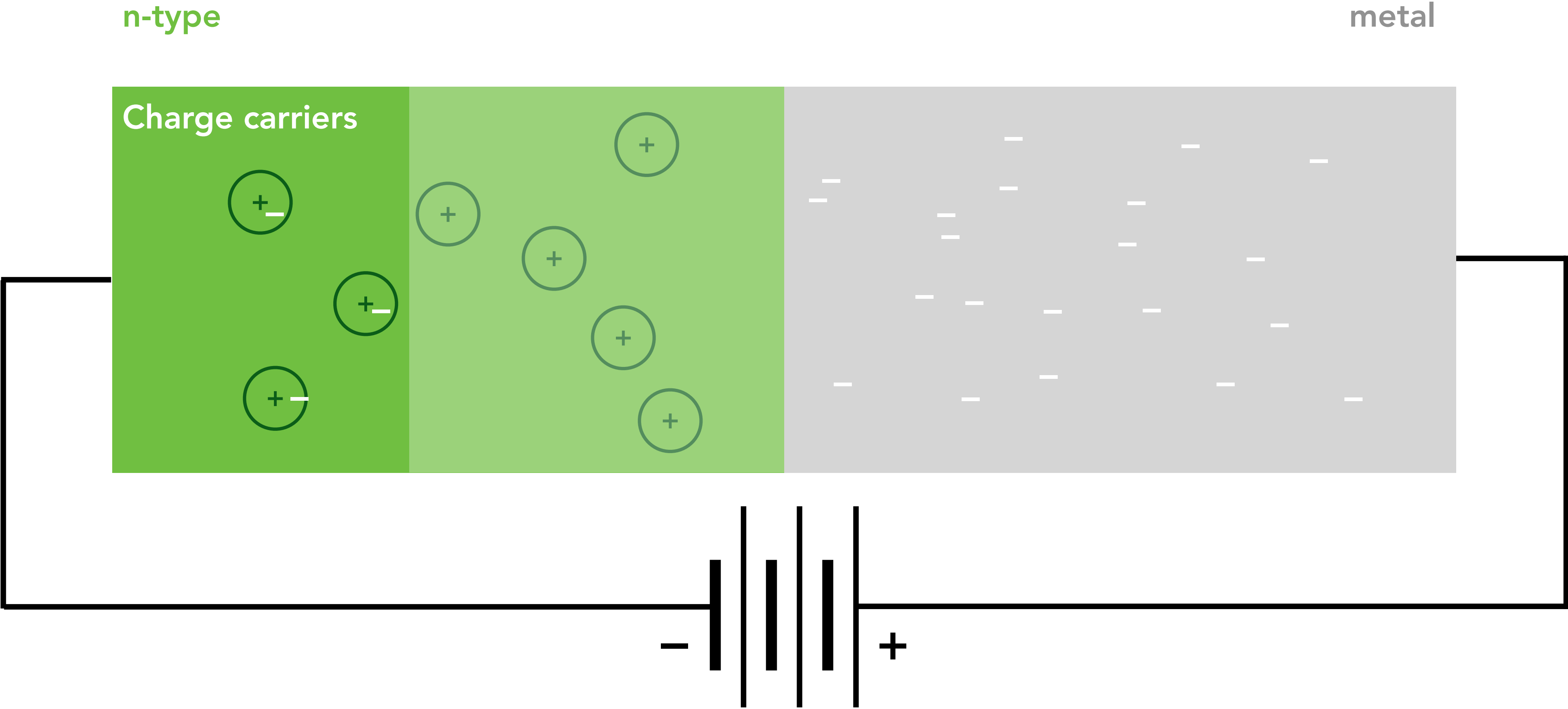
S-M junctions - carrier movement



Consider instead the case where it is **difficult** to move from the metal to the semiconductor

A **Schottky contact** is created

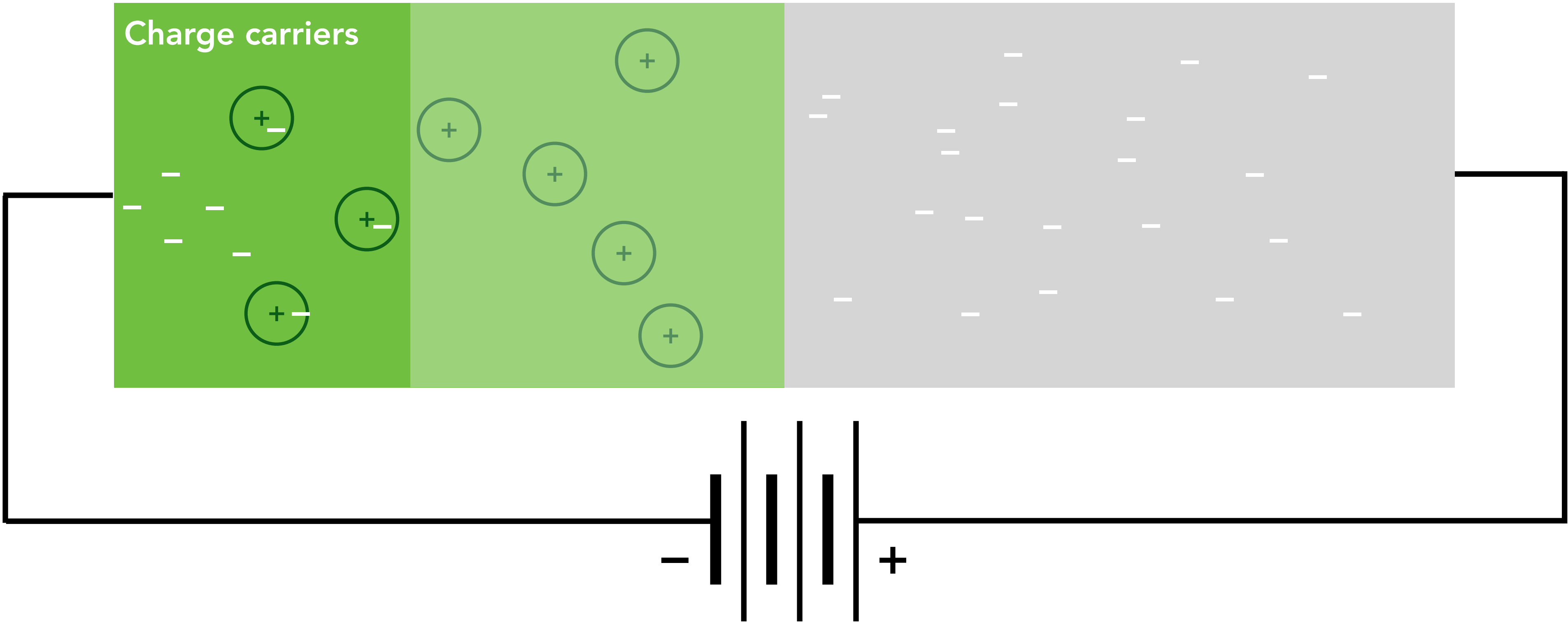
Schottky junctions - forward bias



Schottky junctions - forward bias

“New” electron flow →

n-type metal

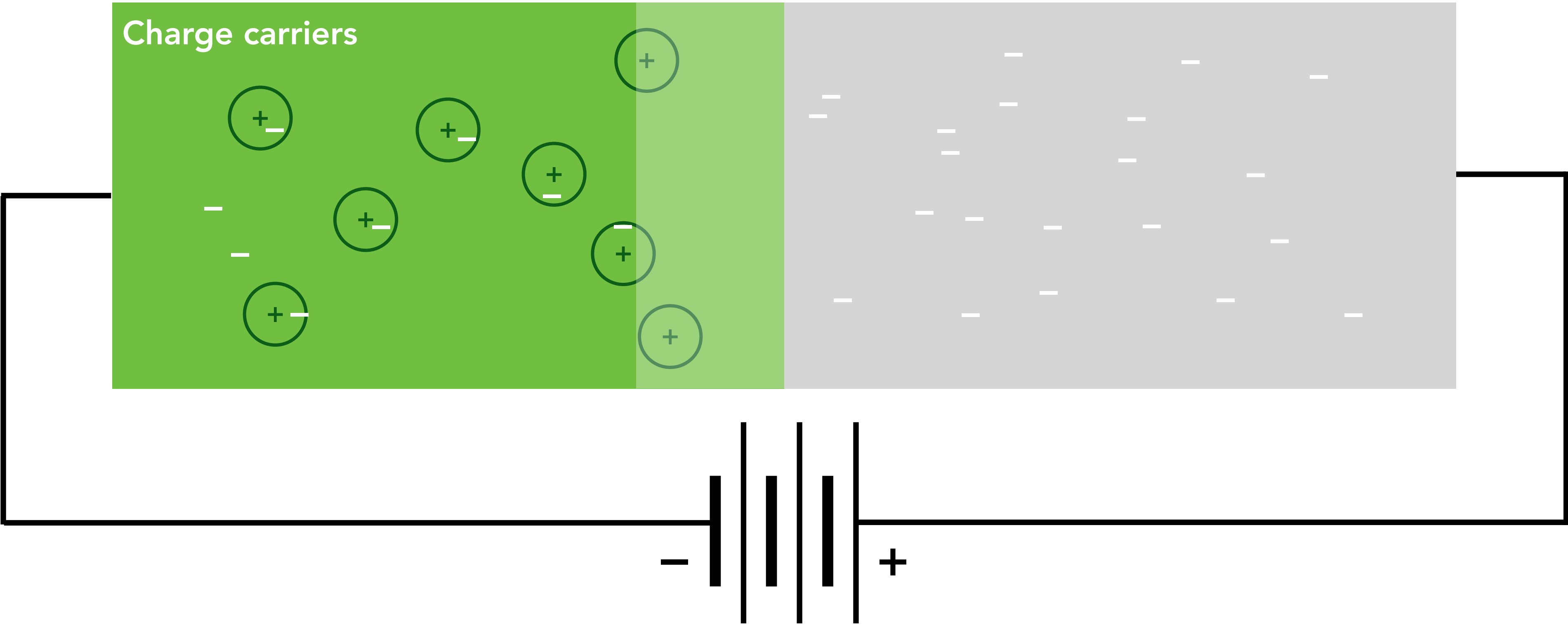


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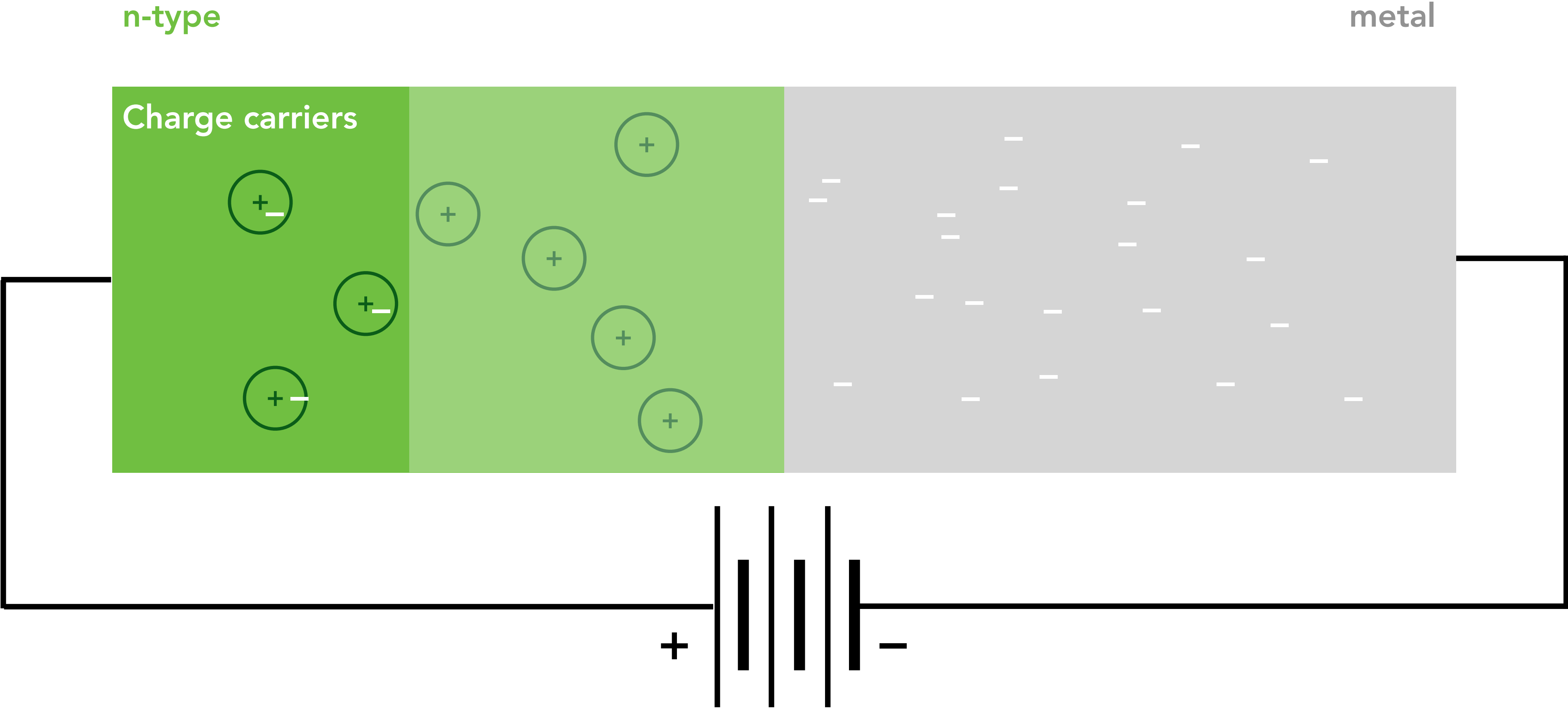
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Schottky junctions - reverse bias

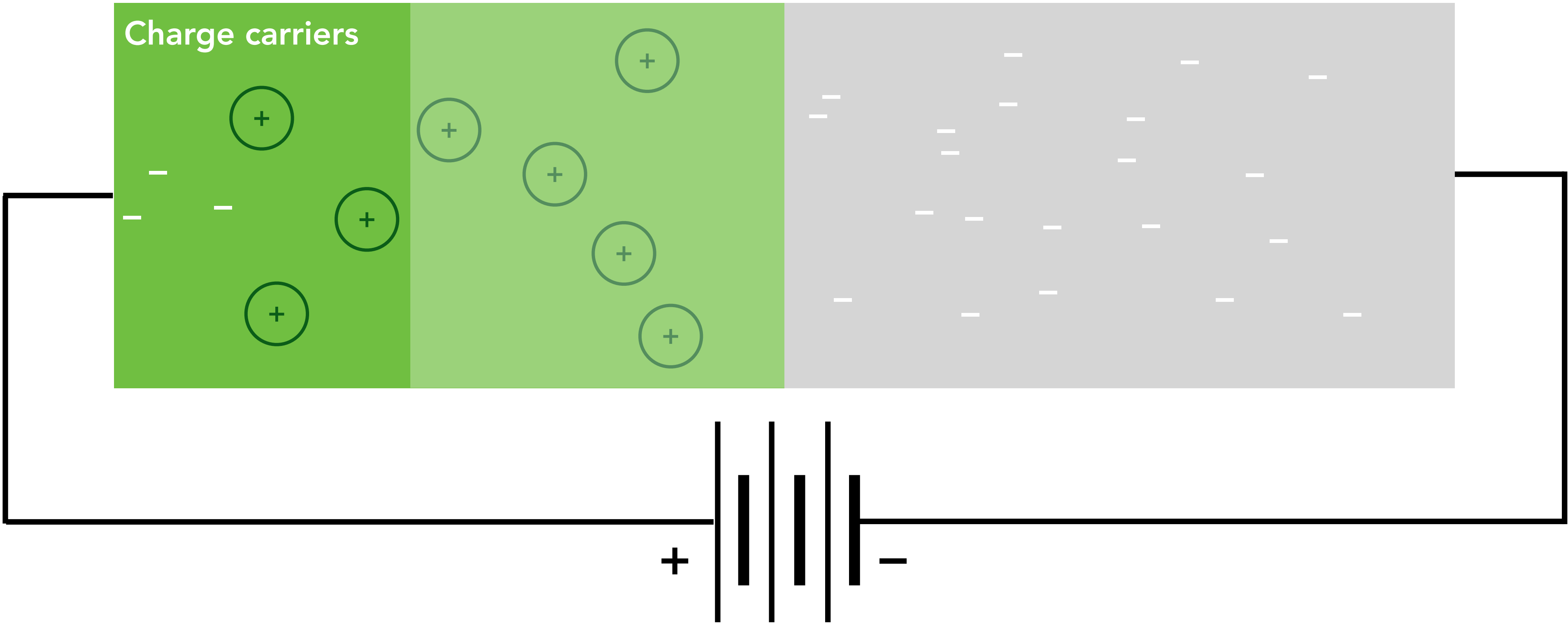


Schottky junctions - reverse bias

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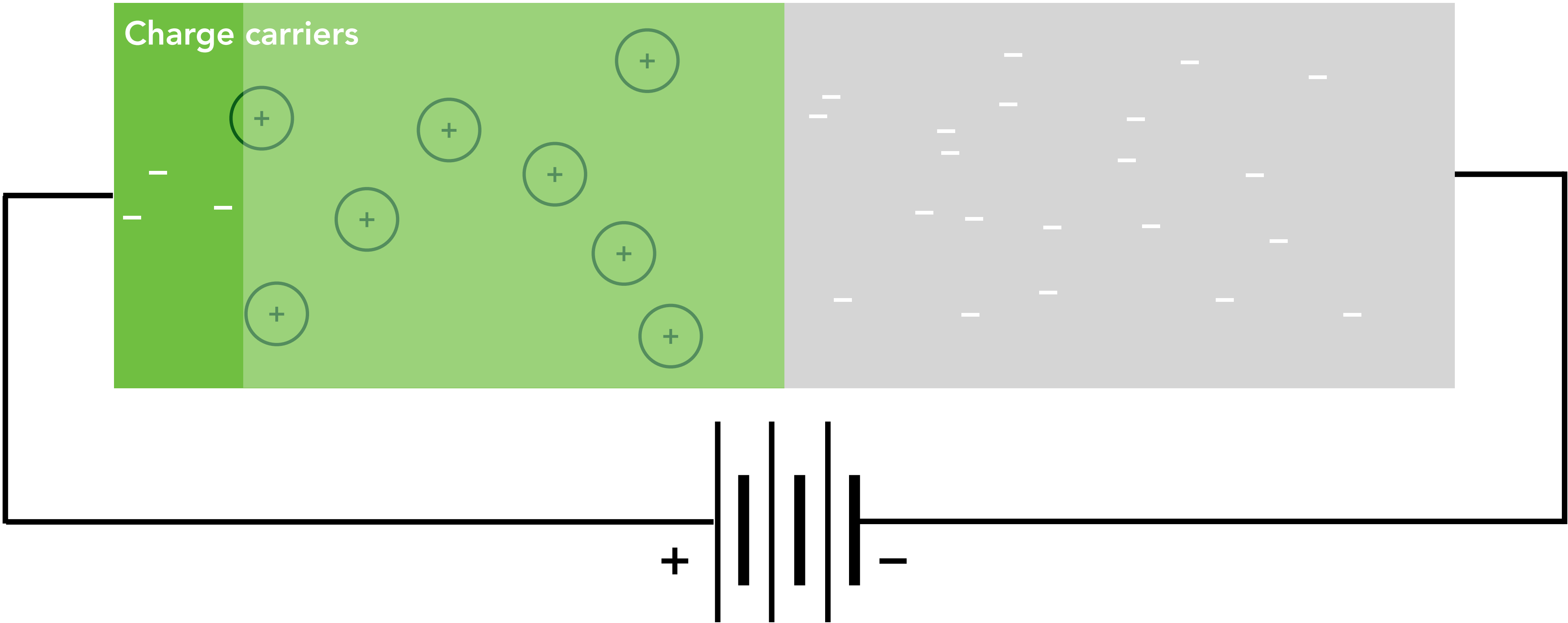


Schottky junctions - reverse bias

"New" electron flow ←

n-type

metal



Schottky junctions - external field

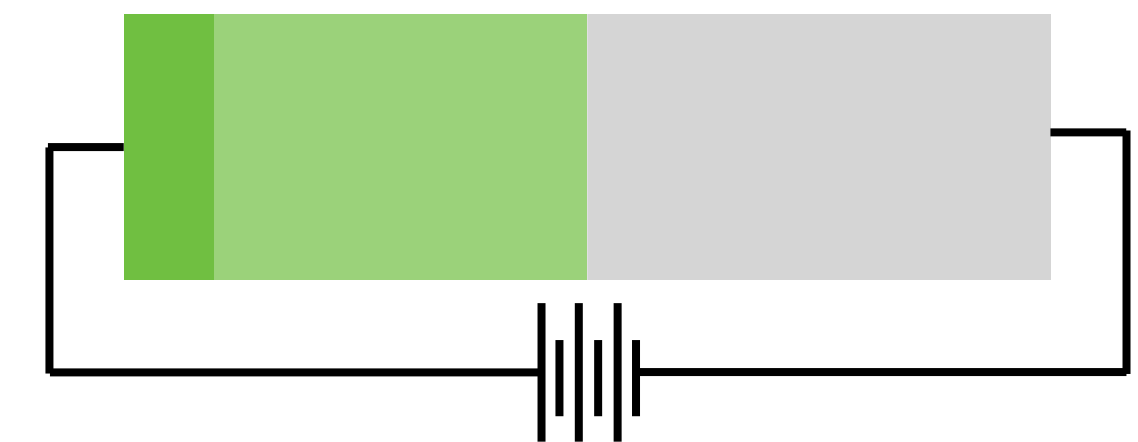
Bringing together a metal and doped semiconductor **can** give rise to a natural region free of charge carriers between the two

- This built-in depletion region contains an electric field which counters the diffusion of more carriers across the junction

We can forward bias the junction to reduce the depletion width

- Current will flow through the junction once the depletion width has been reduced to 0

We can reverse bias the junction as for a pn-junction and no current will flow



Schottky junctions - external field

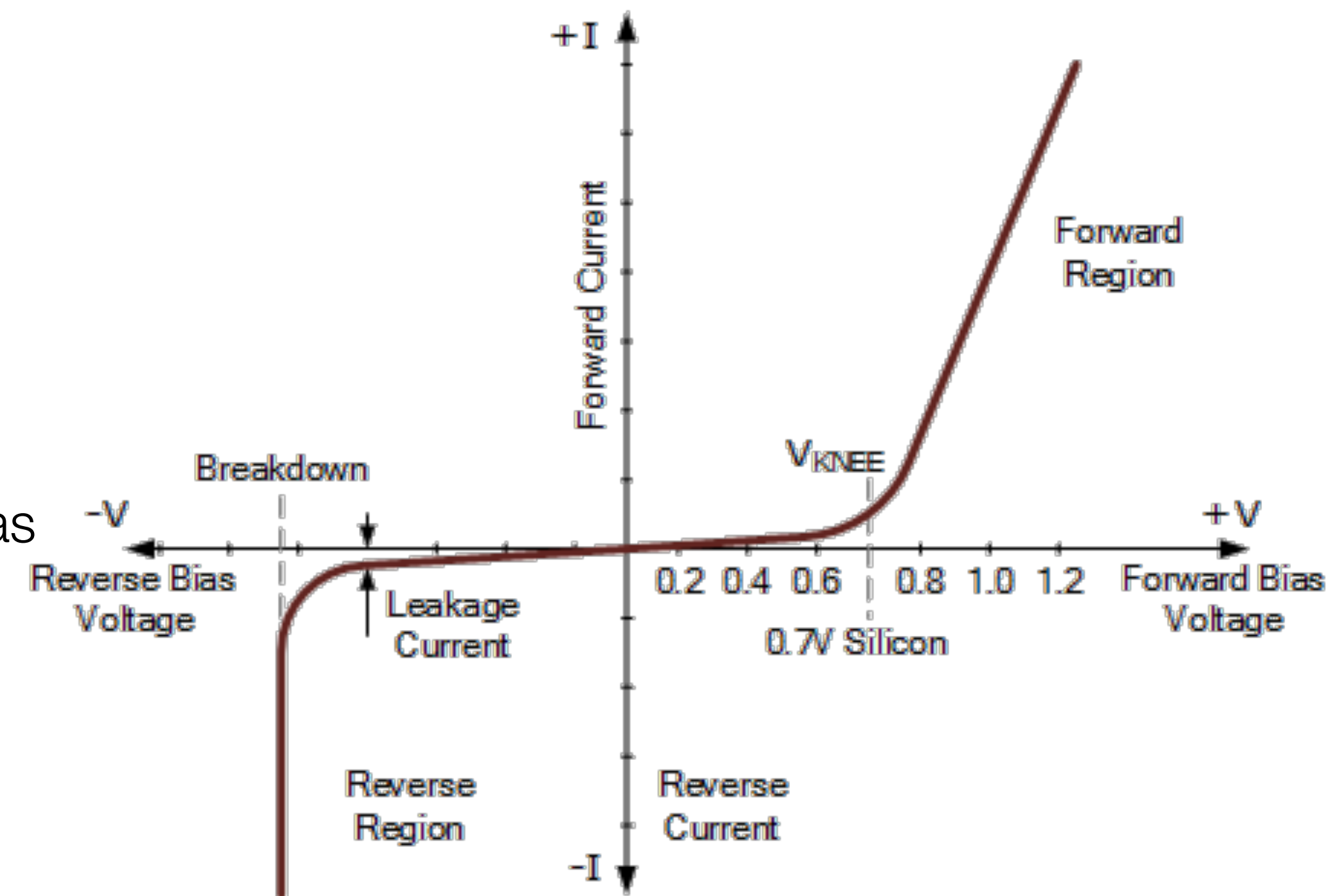
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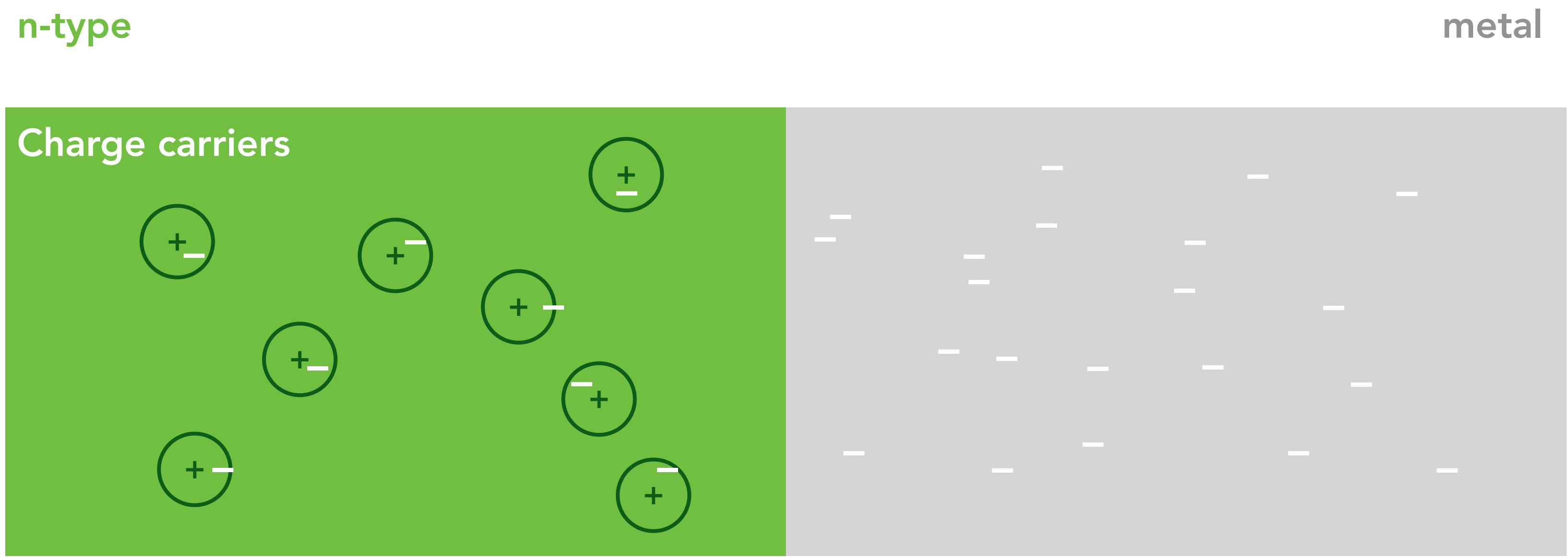
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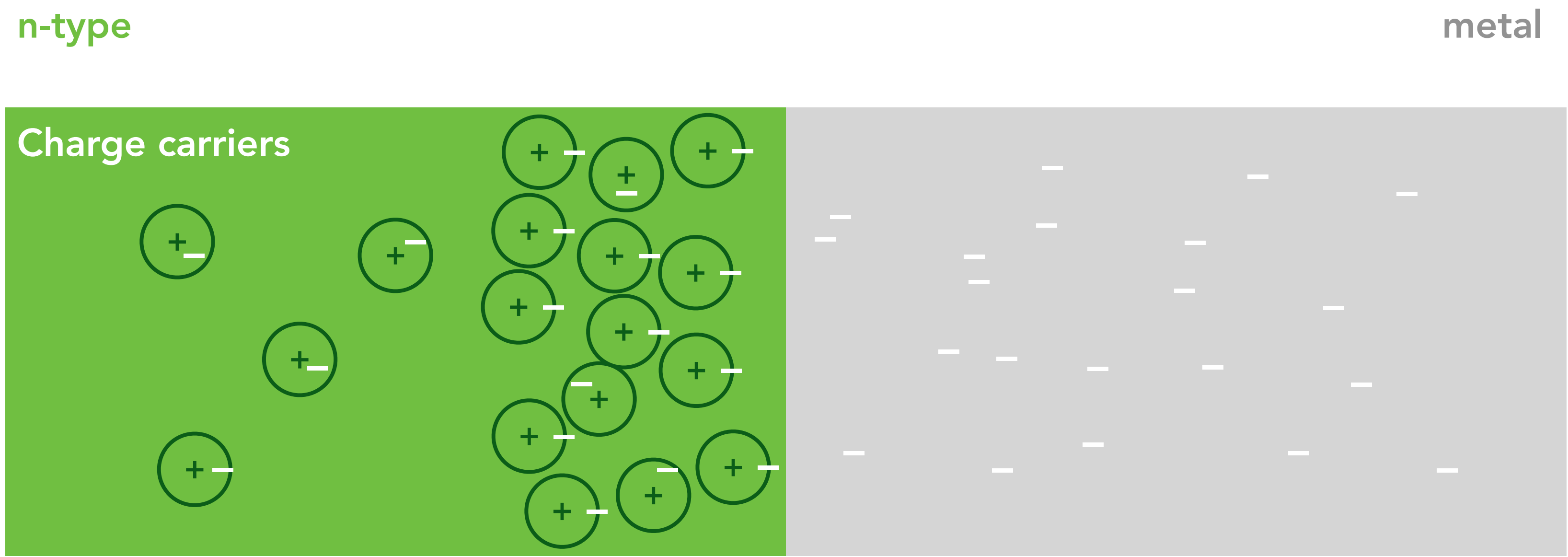


Schottky junctions - NB



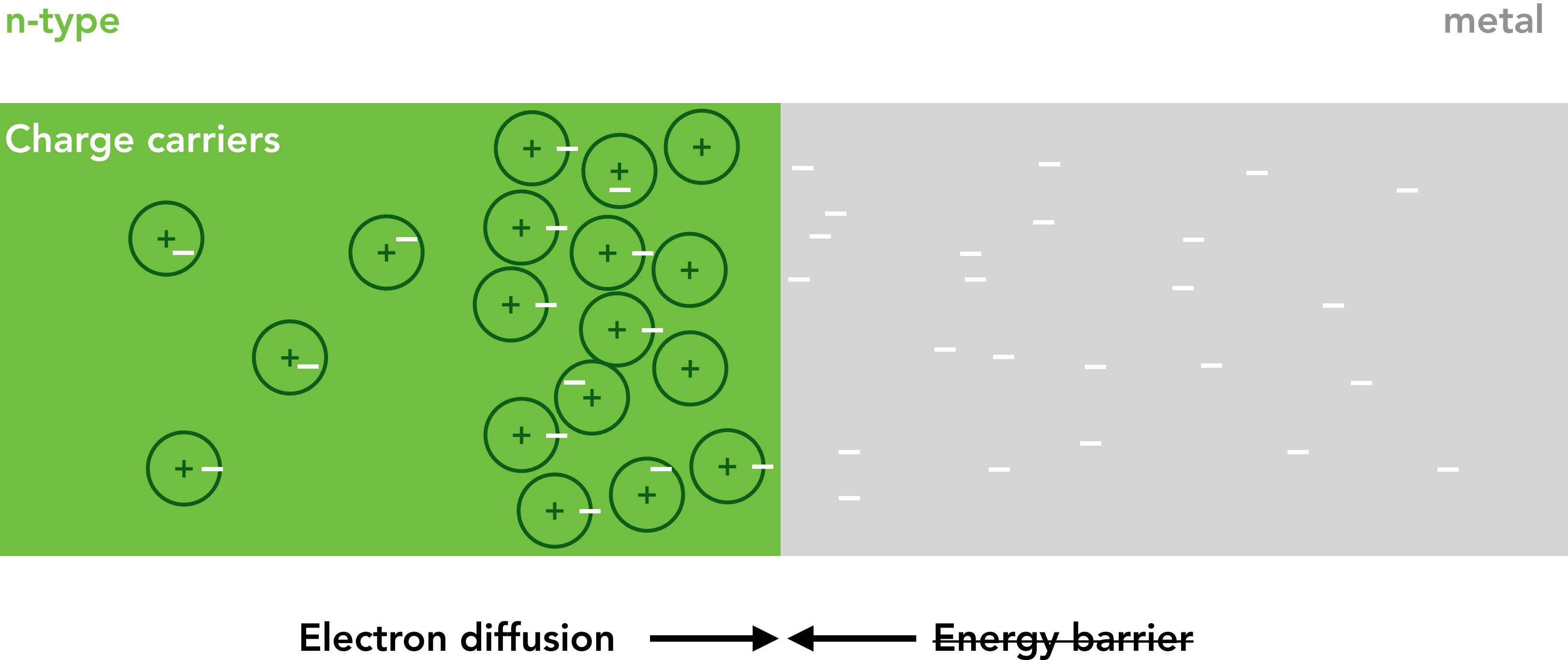
Consider a Schottky diode with a **strongly doped** n-type semiconductor

Schottky junctions - NB

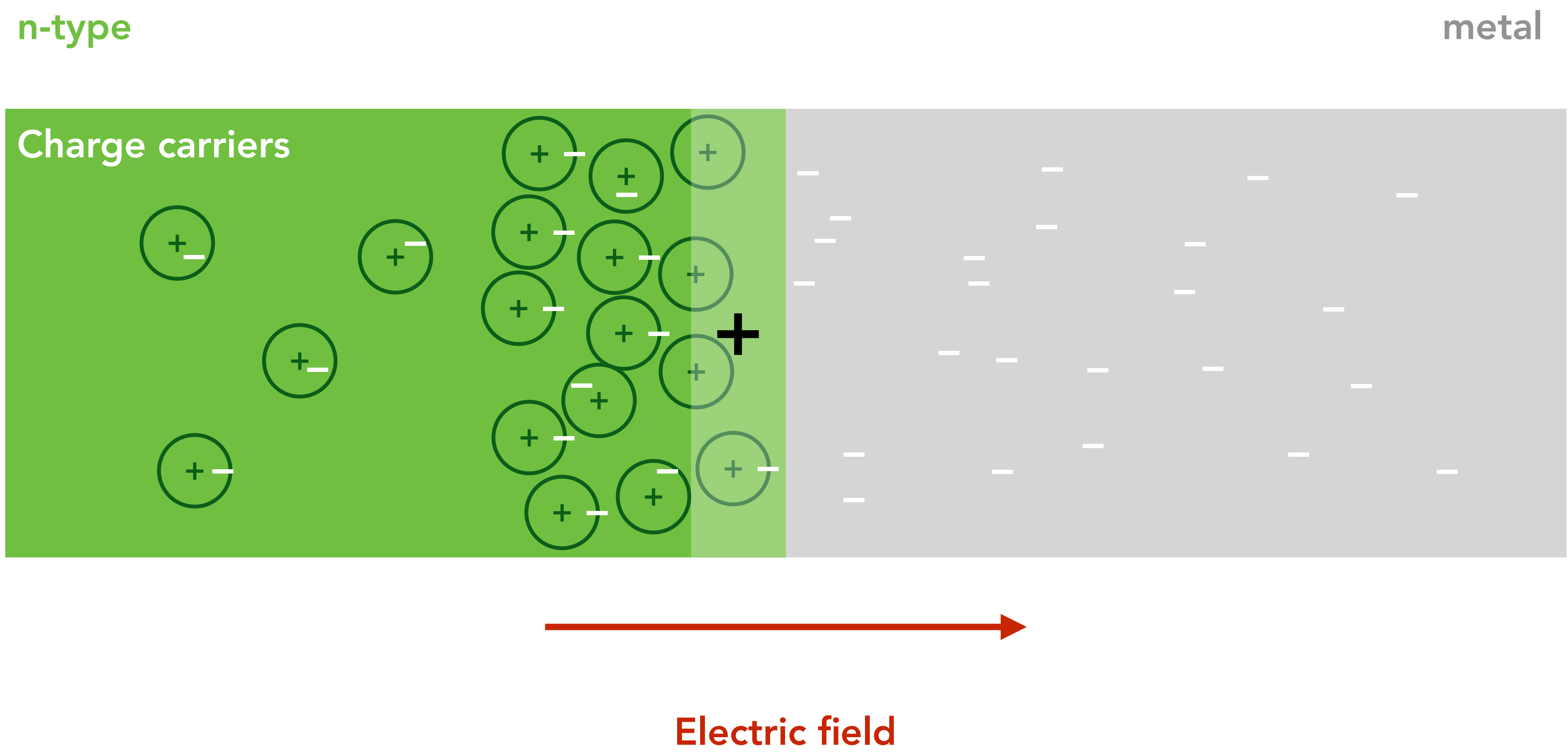


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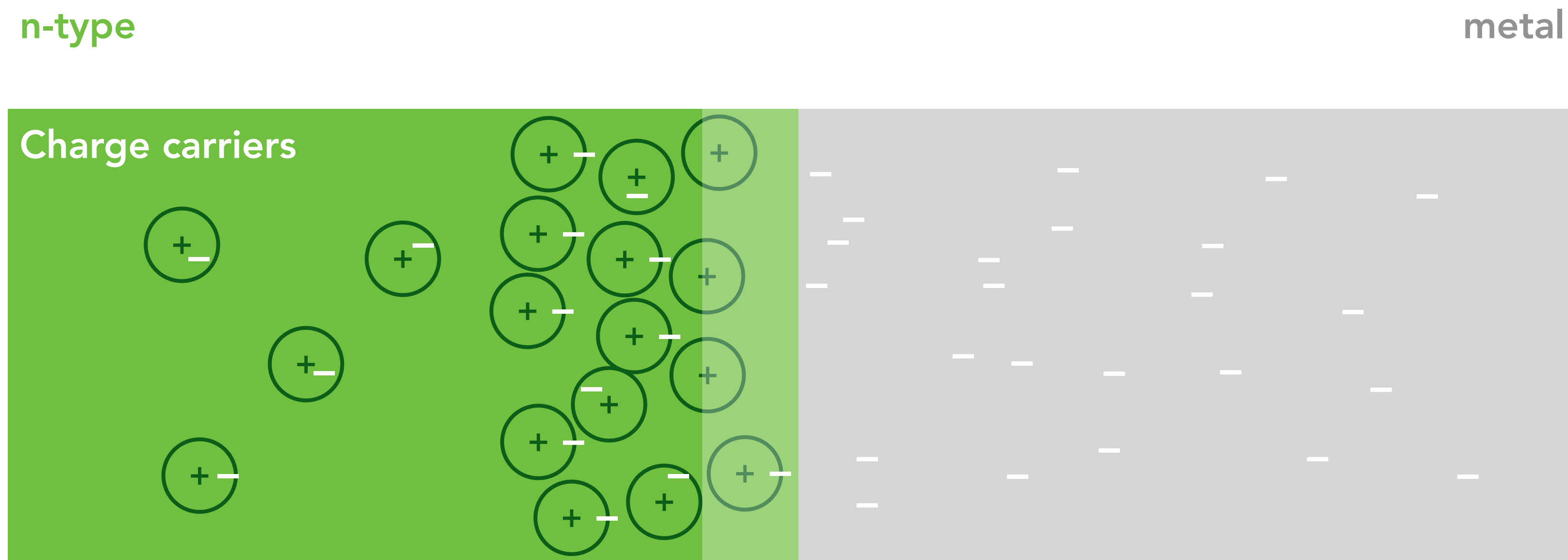
Schottky junctions - NB



Schottky junctions - NB



Schottky junctions - NB



Consider a Schottky diode with a **strongly doped** n-type semiconductor

The depletion width will get smaller and smaller until it looks like an **Ohmic contact**

Electron Light Switches

(Transistors)

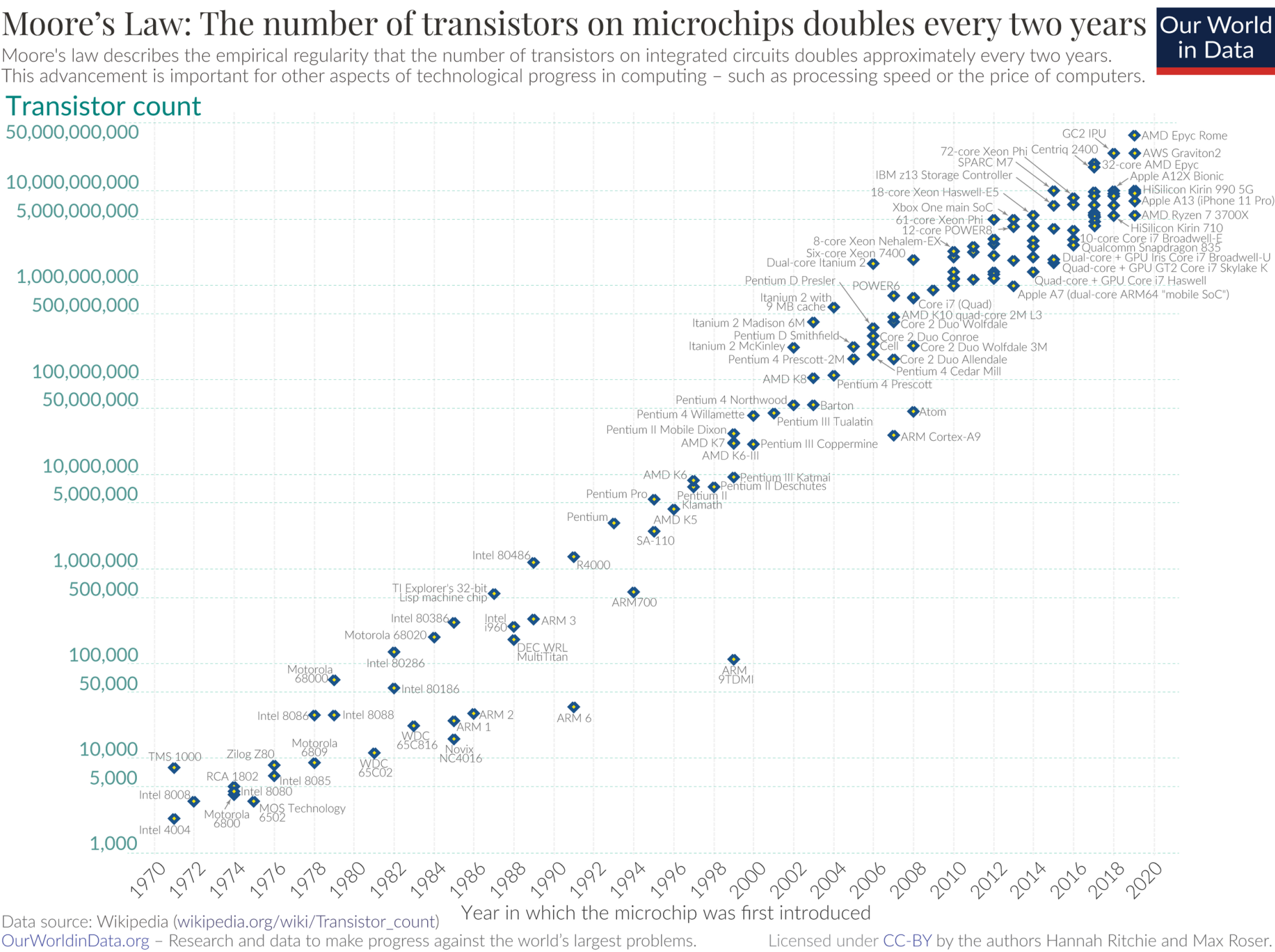
Transistors

Transistors play a crucial role in all modern technology, and are the reason that the semiconductor industry is so vast

- First transistors invented in ~1947
- First MOSFET in 1960

Latest generation of devices contain >100B transistors in an area of 30 x 30 mm²

- Couldn't find more up-to-date numbers, but in 2014 Forbes calculated there had been 3×10^{21} transistors manufactured (cf. 3×10^{13} cells in human body)



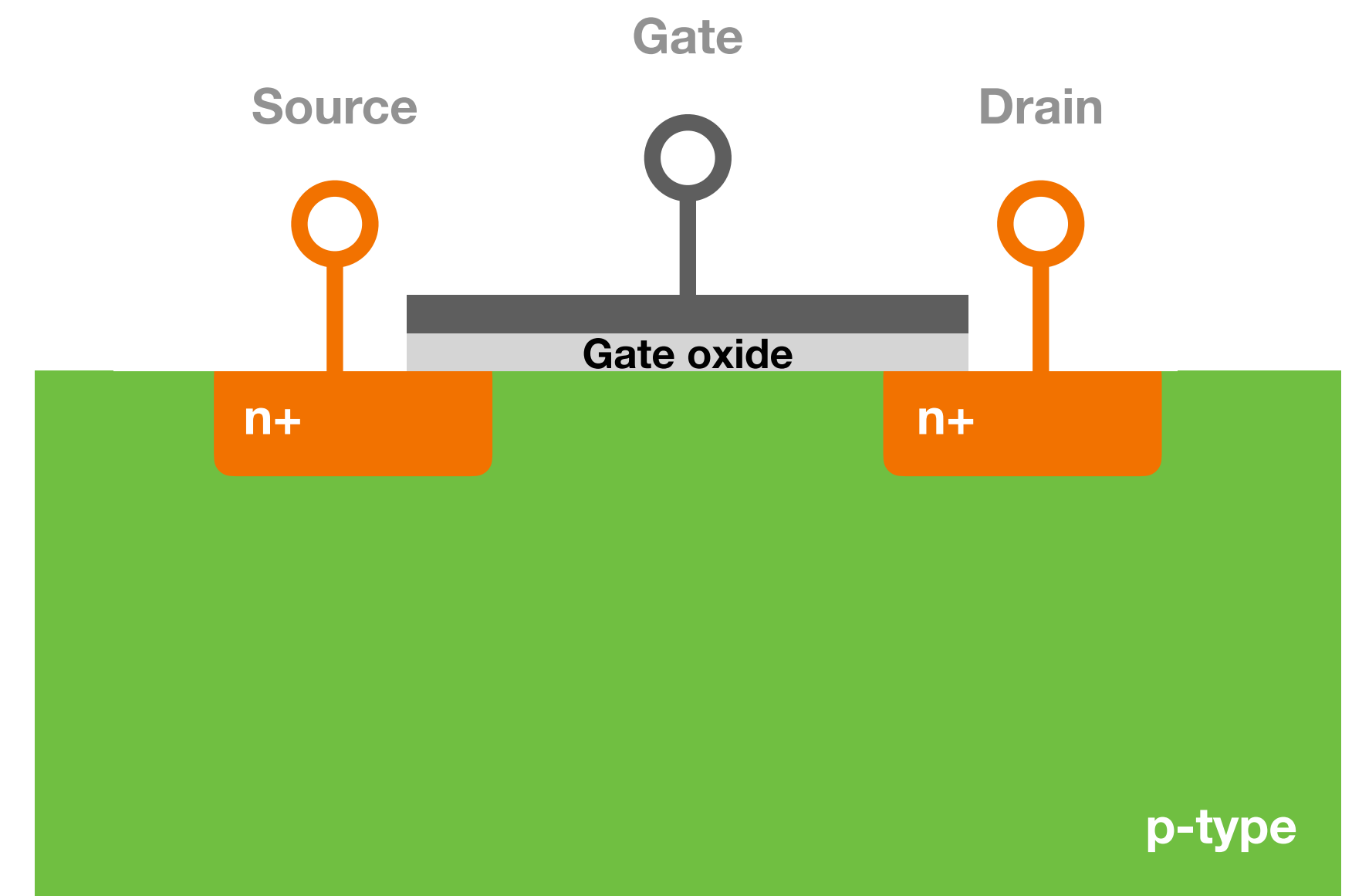
Anatomy of a transistor

Transistors are very much like Bach - variations on a theme...

- There are many kinds of transistors even just in silicon, ranging from the most commonly used MOSFETs, to MISFETs, JFETs, bipolar junction transistors, point-contact transistors, Schottky transistors...

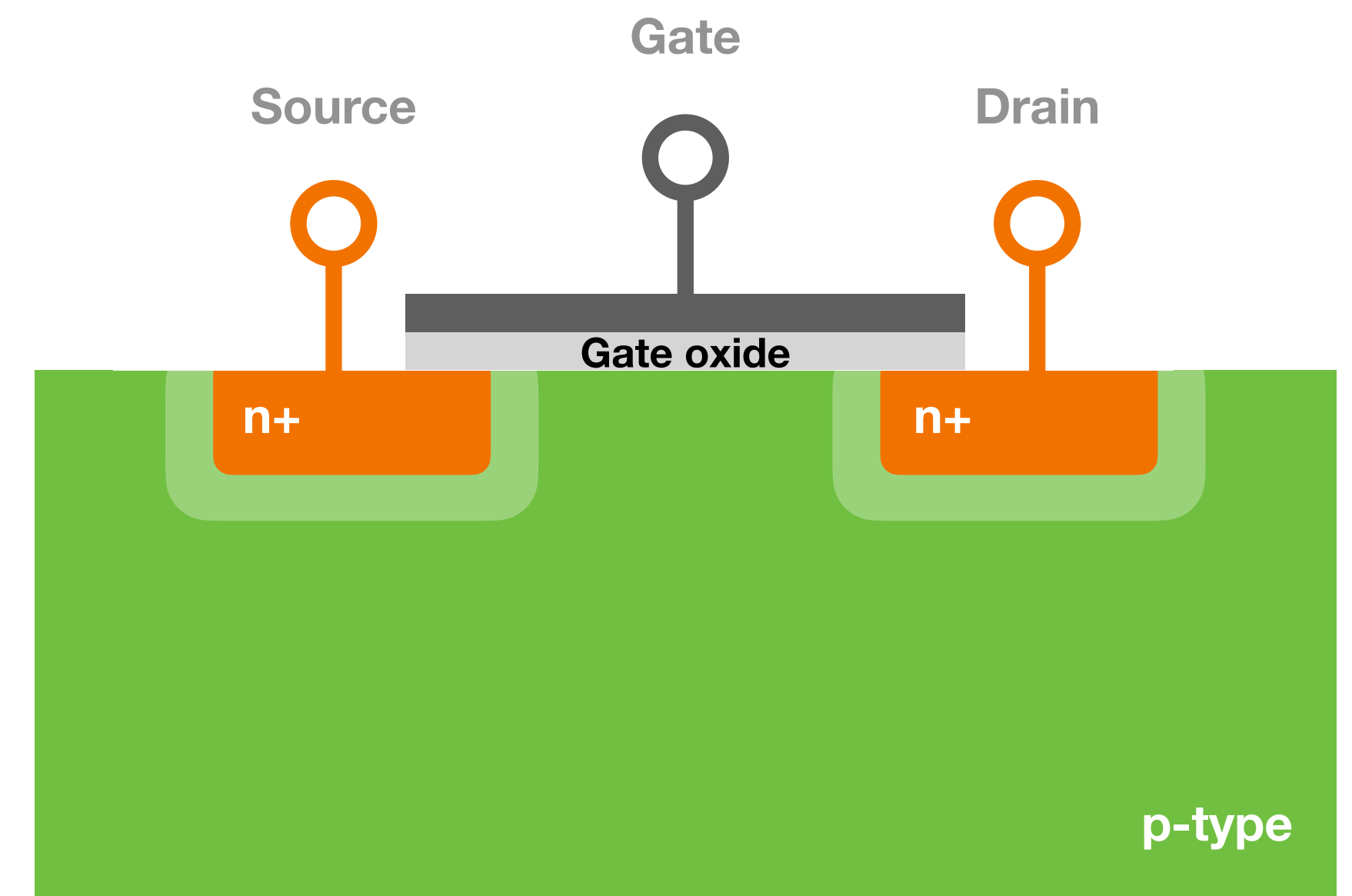
For today we are going to limit ourselves to discussing Metal-Oxide-Semiconductor Field Effect Transistors (MOSFETs)

- Again, there are many variations but we will consider *enhancement-mode MOSFETs*



Operating principles

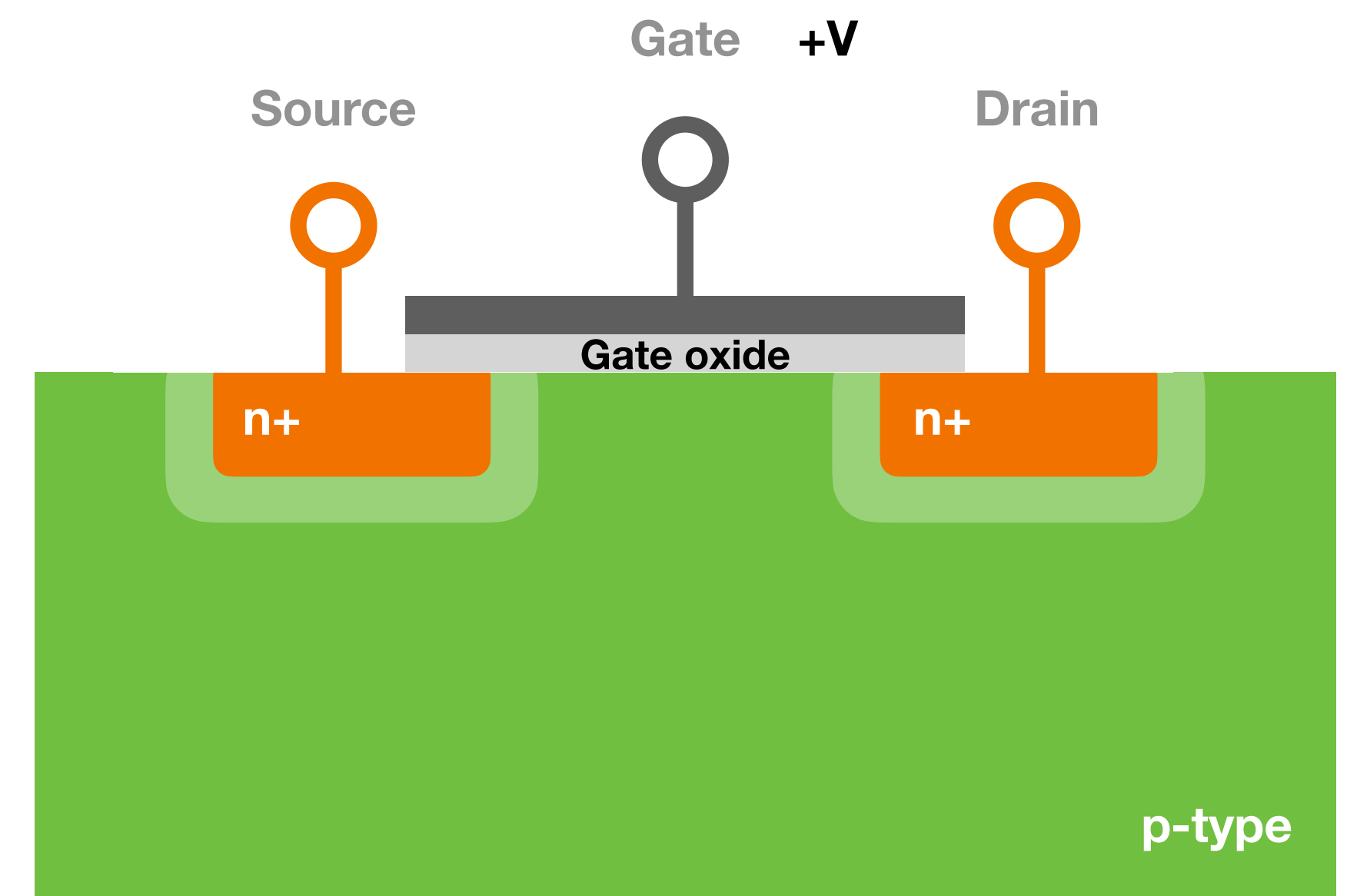
MOSFETs consist of two doped regions - the *drain* and the *source* - on a lightly doped substrate, with insulating *gate oxide* sitting between the doped wells and a conductive *gate*



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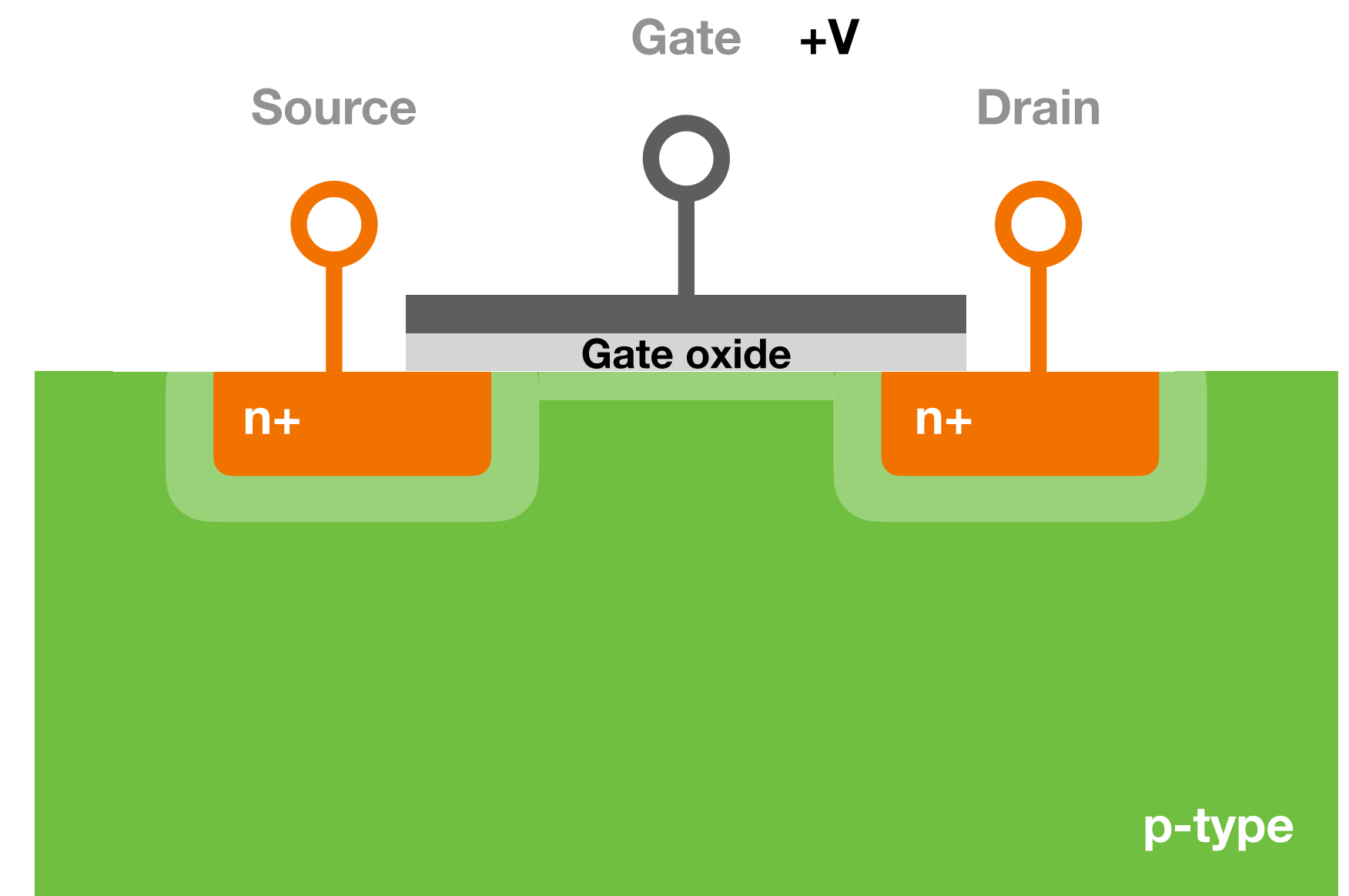
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Operating principles

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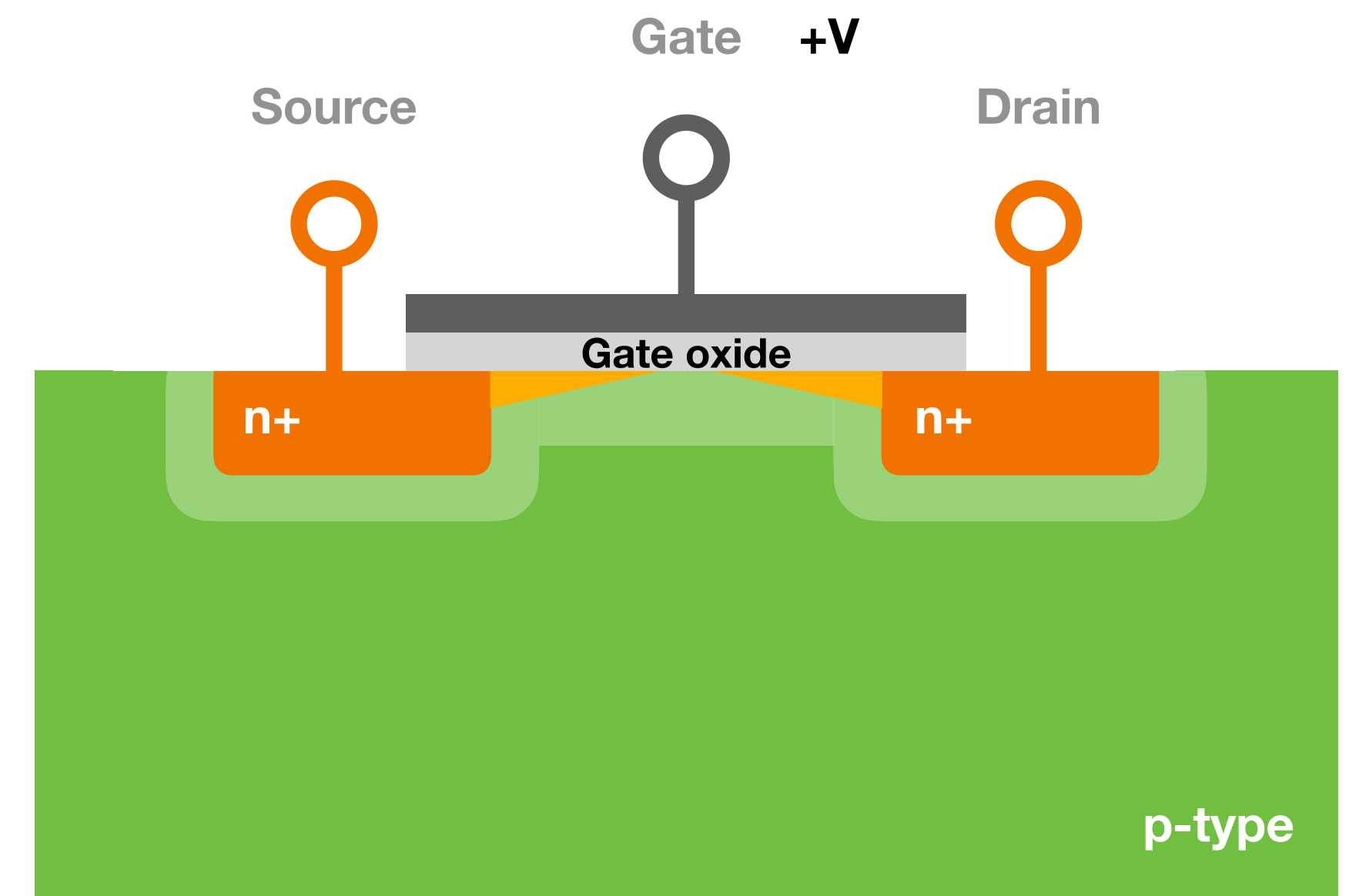
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- This drives holes away from the p-type material below the oxide



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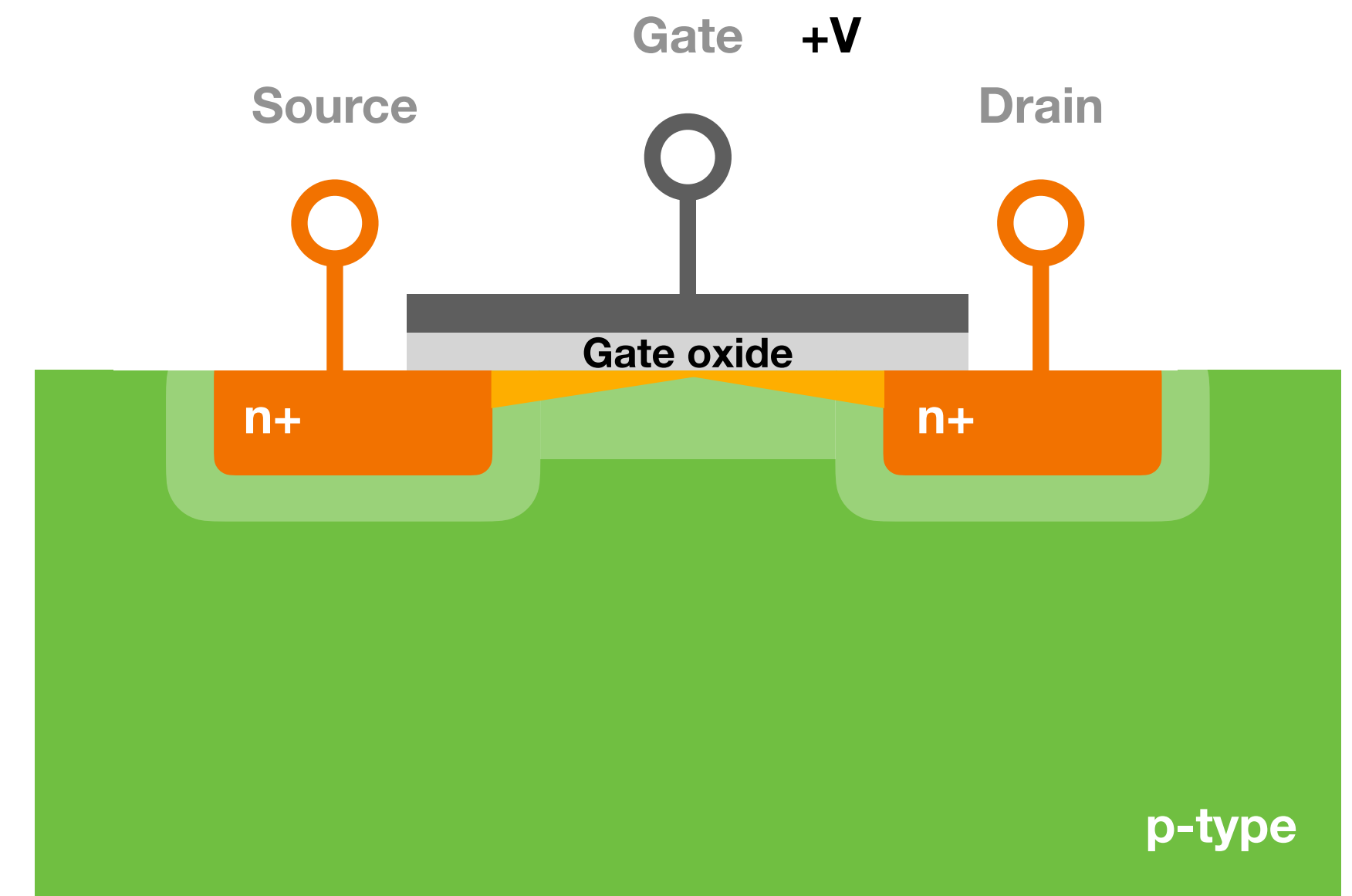
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- With enough voltage it draws electrons from the source and drain



Operating principles

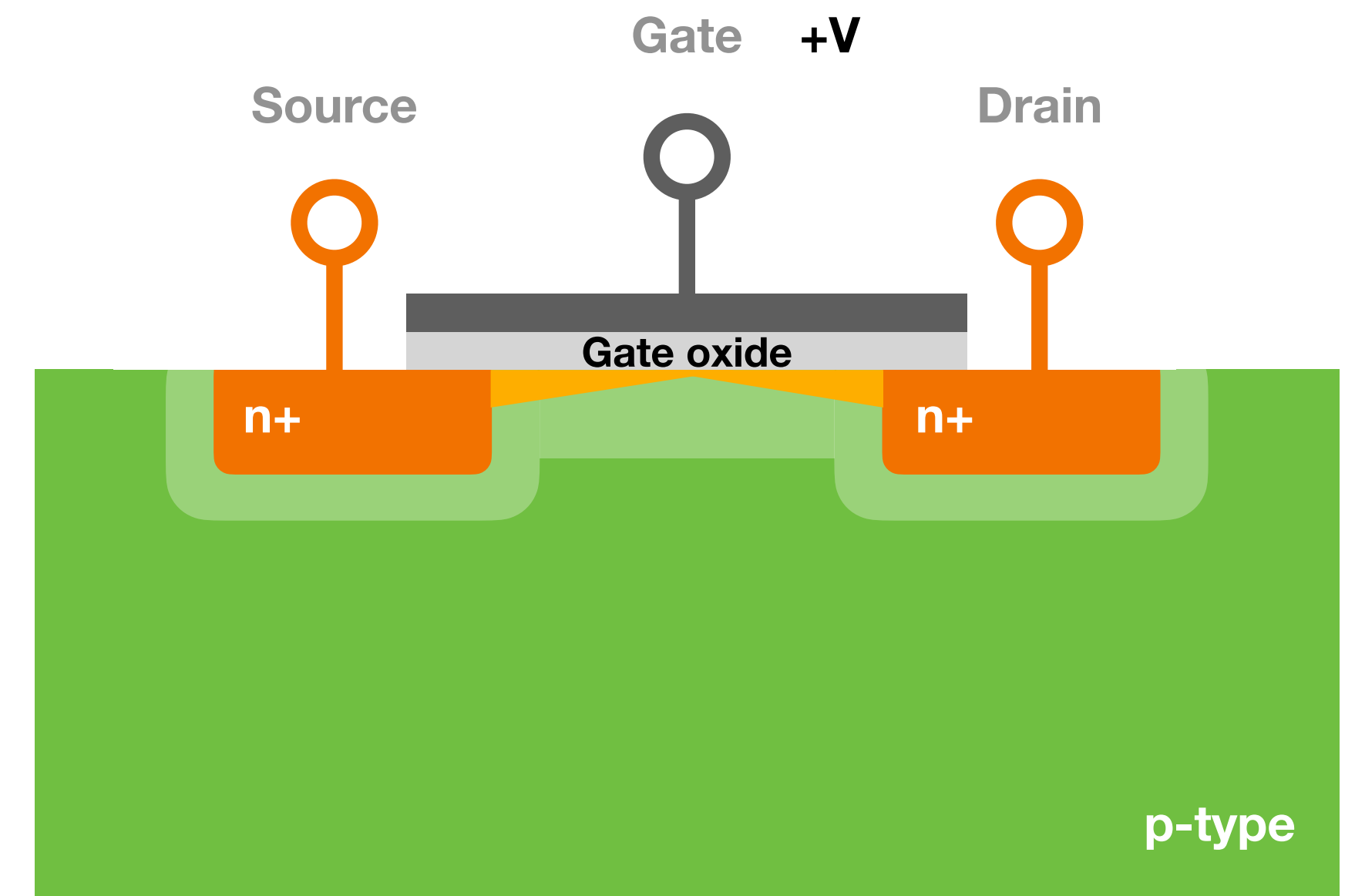
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- A positive voltage is applied to the gate
- This drives holes away from the p-type material below the oxide
- With enough voltage it draws electrons from the source and drain
- Once these two regions join, an *inversion channel* is formed (minority carrier channel) => effectively a switch with a threshold voltage



Operating principles

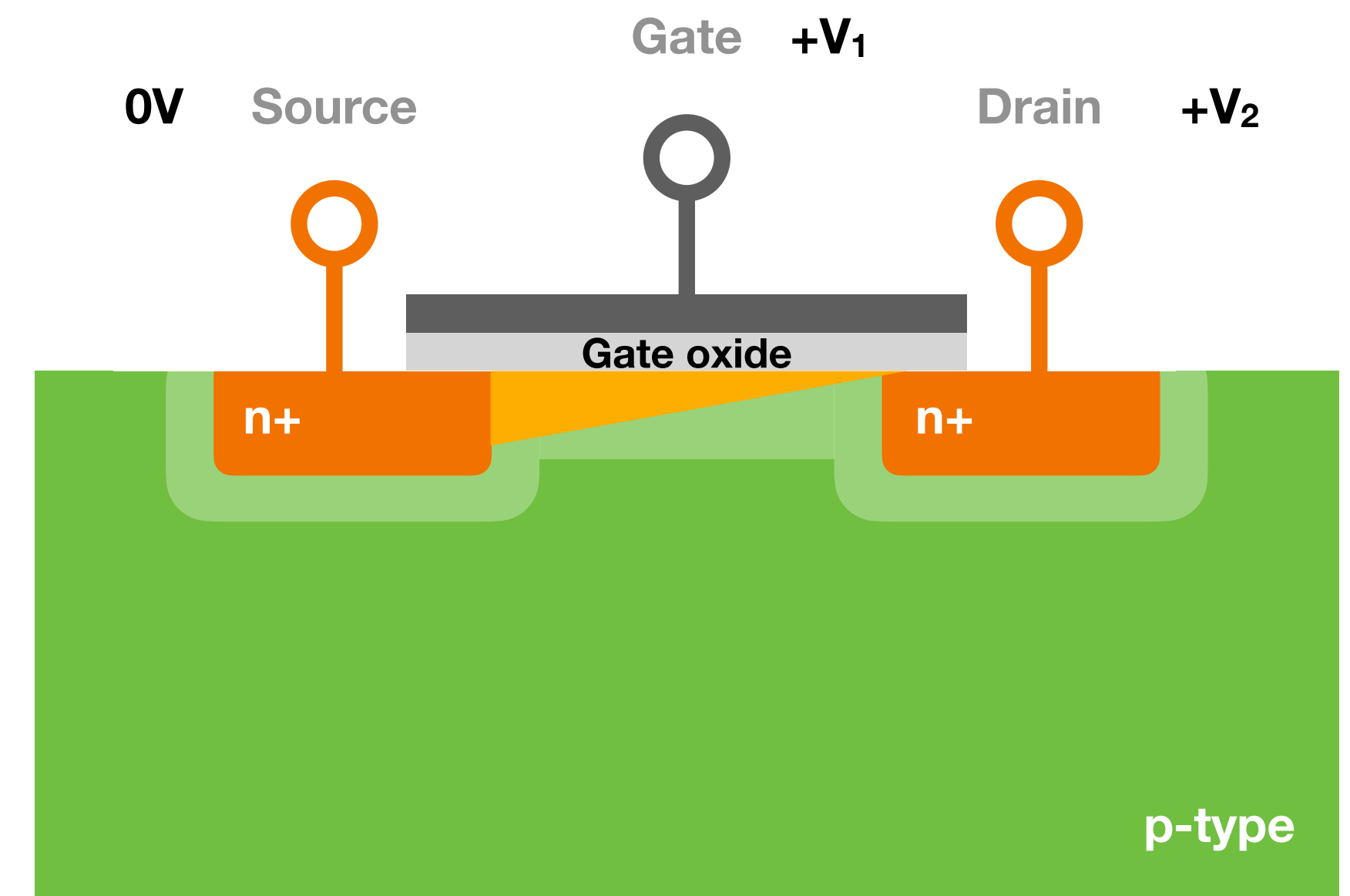
With no potential difference across the source and drain, there is no flow of current (making it pretty useless)



Operating principles

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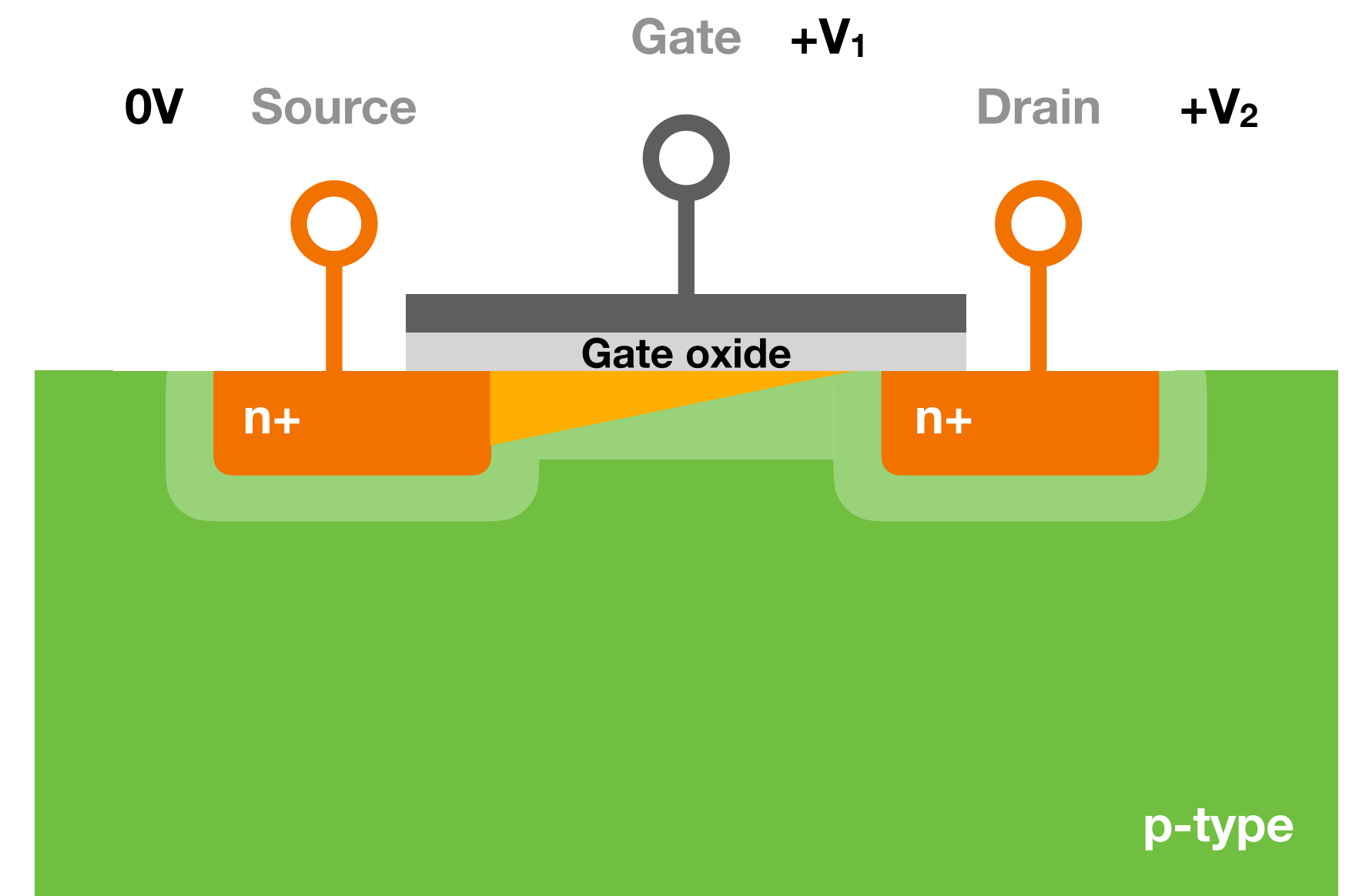
- Typically there will be a potential from source to drain, giving an asymmetric inversion layer



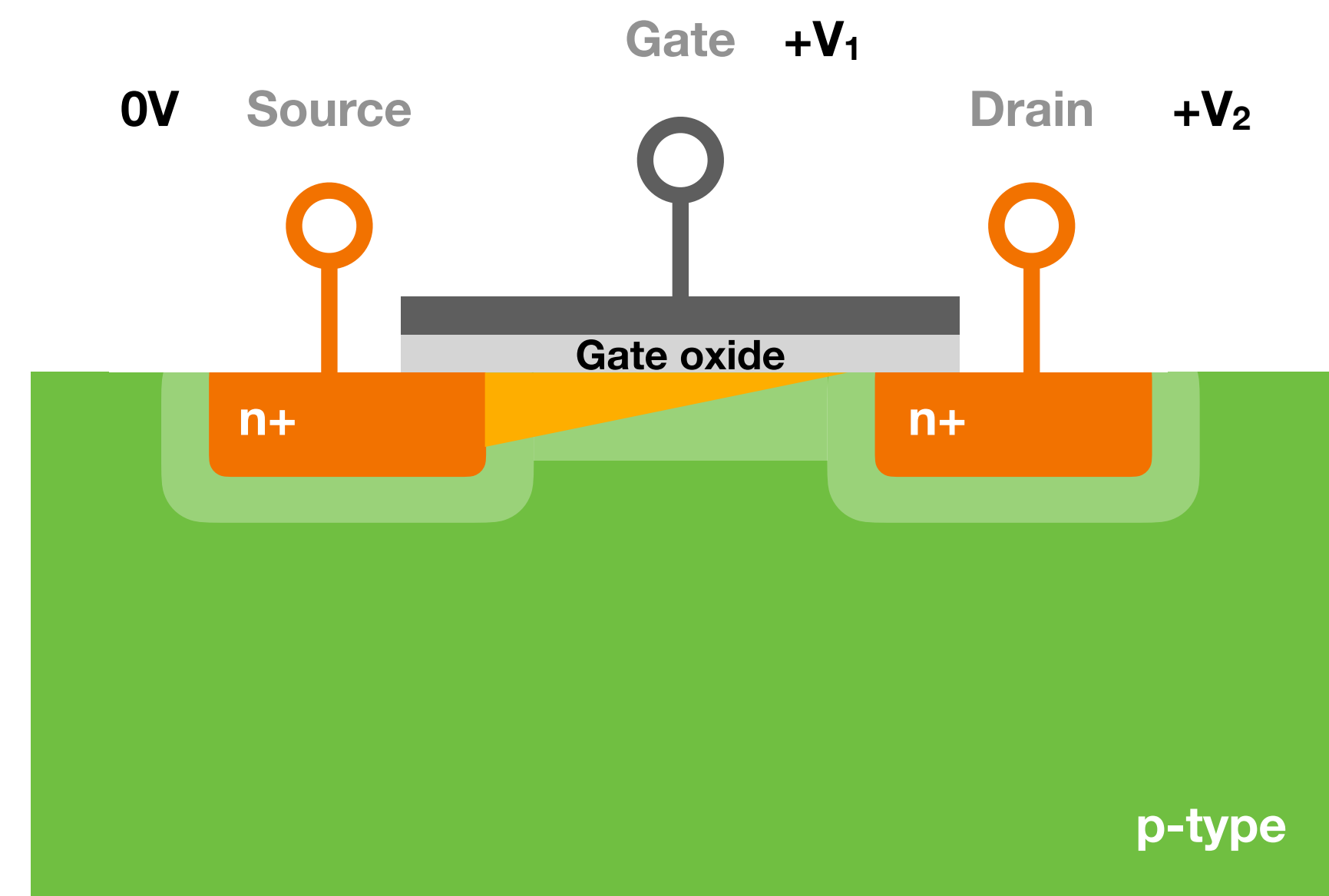
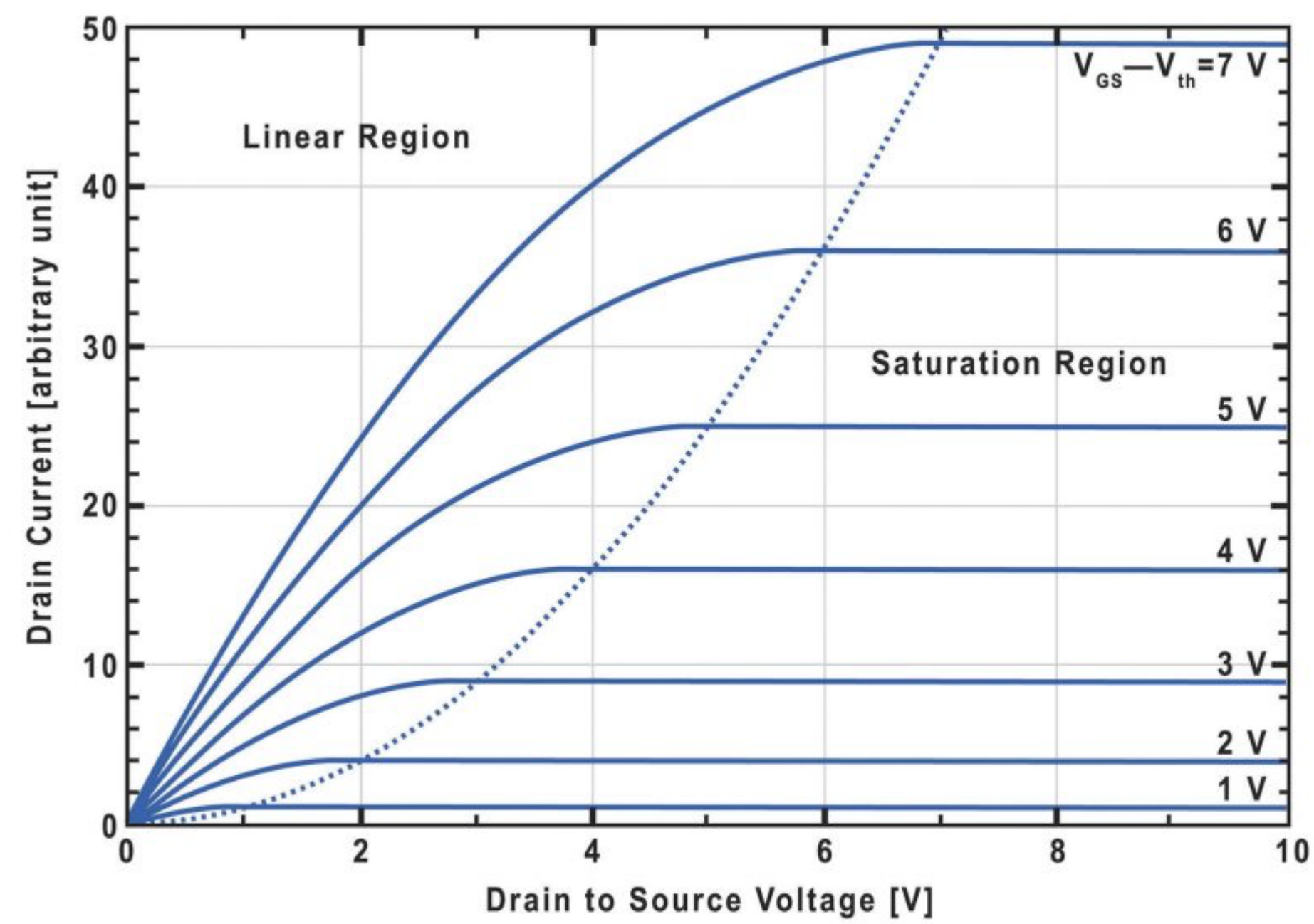
Operating principles

With no potential difference across the source and drain, there is no flow of current (making it pretty useless)

- Typically there will be a potential from source to drain, giving an asymmetric inversion layer
- Increasing the potential difference leads to “pinch off” - the inversion channel is cut off due to the higher potential close to the drain - current still flows but is saturated and no longer increases with voltage



Operating principles

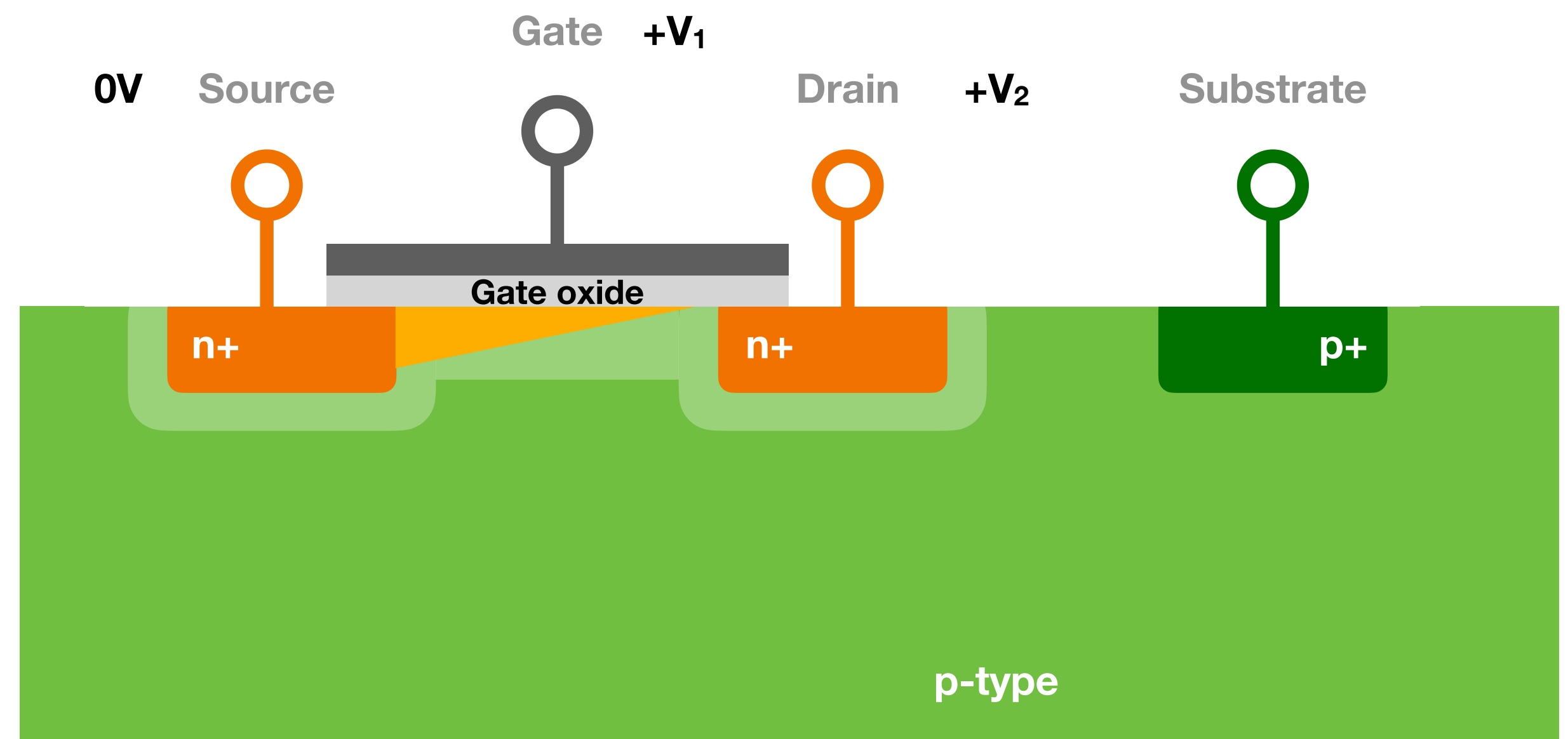


Operating principles

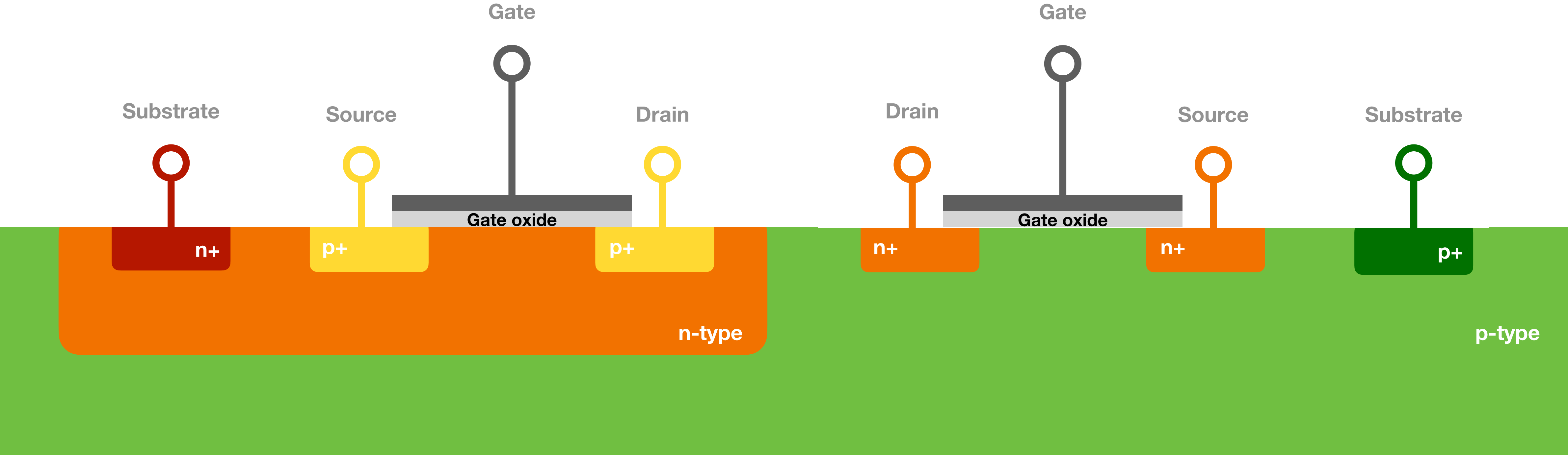
Since we are talking about potential differences it is worth pointing out that the substrate potential can also be controlled

- This is an nMOSFET - analogous devices with **n** and **p** swapped are imaginatively called pMOSFETs

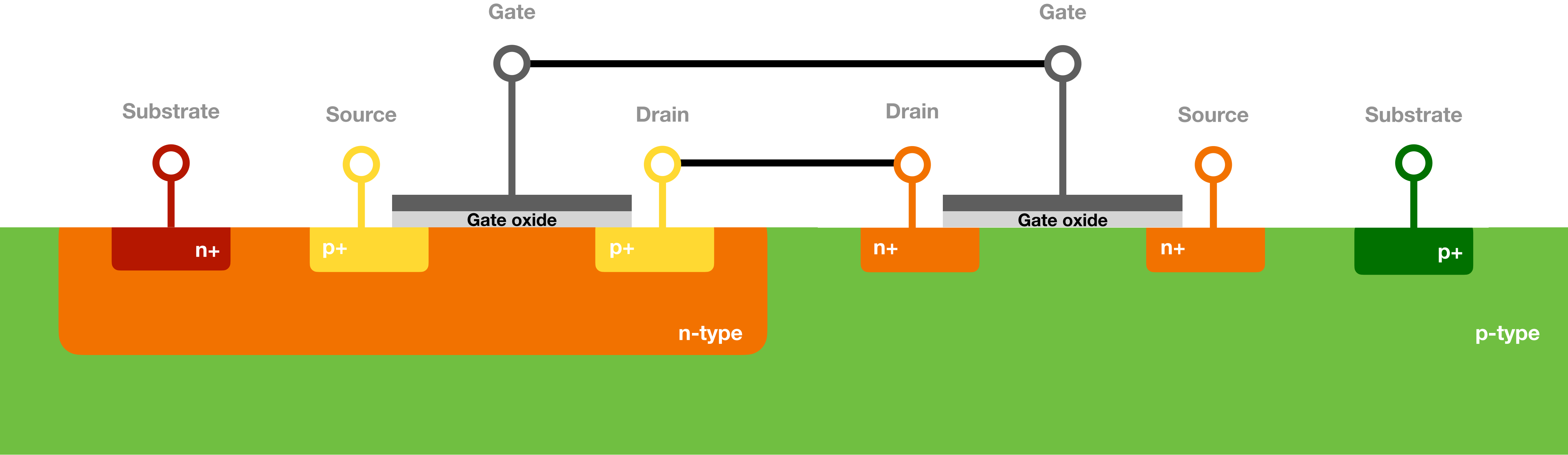
Both behave slightly differently due to the difference in carriers and relative potentials, but the real power comes in combination



CMOS



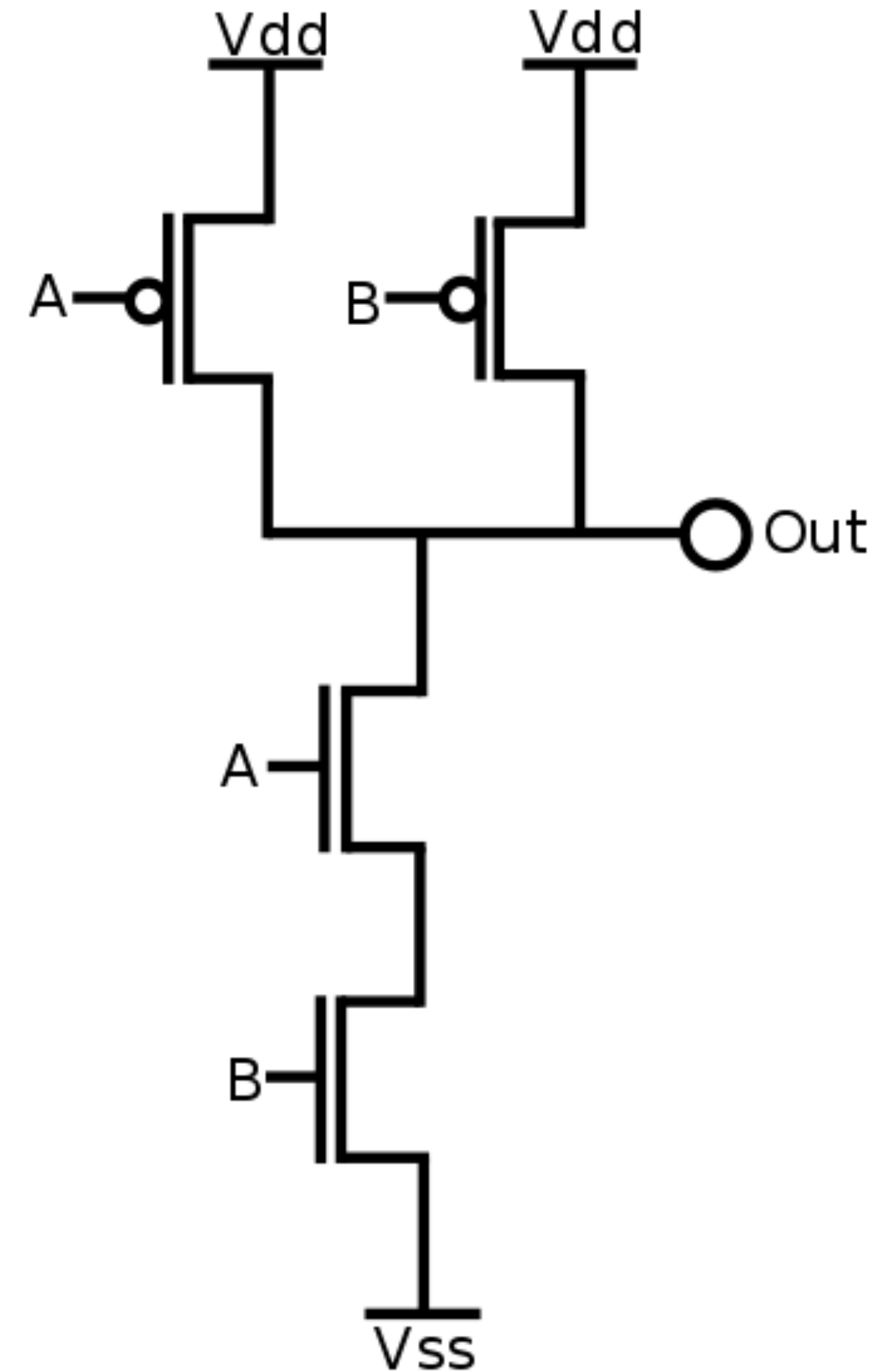
CMOS



CMOS in person

Consider an NAND gate

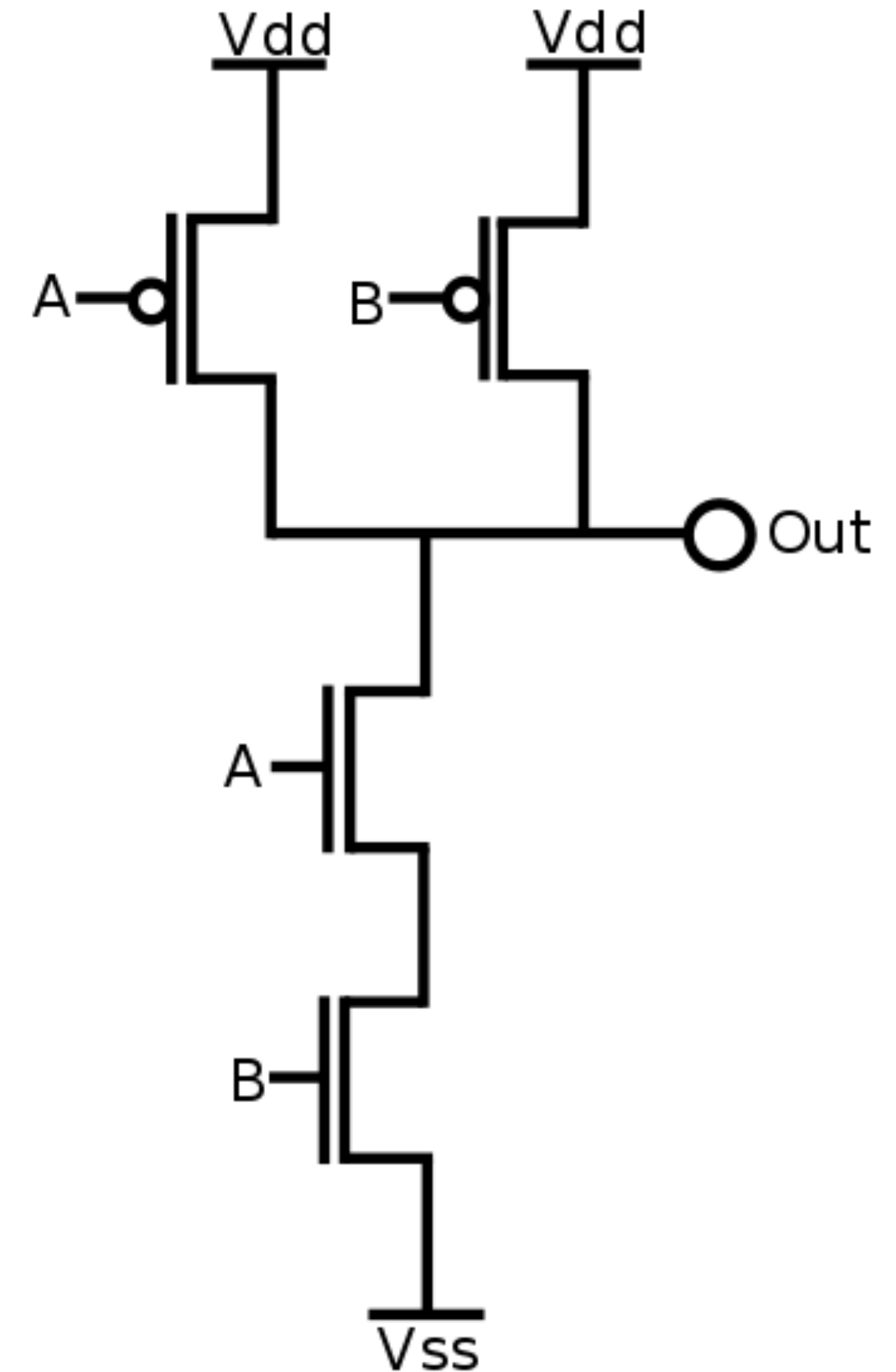
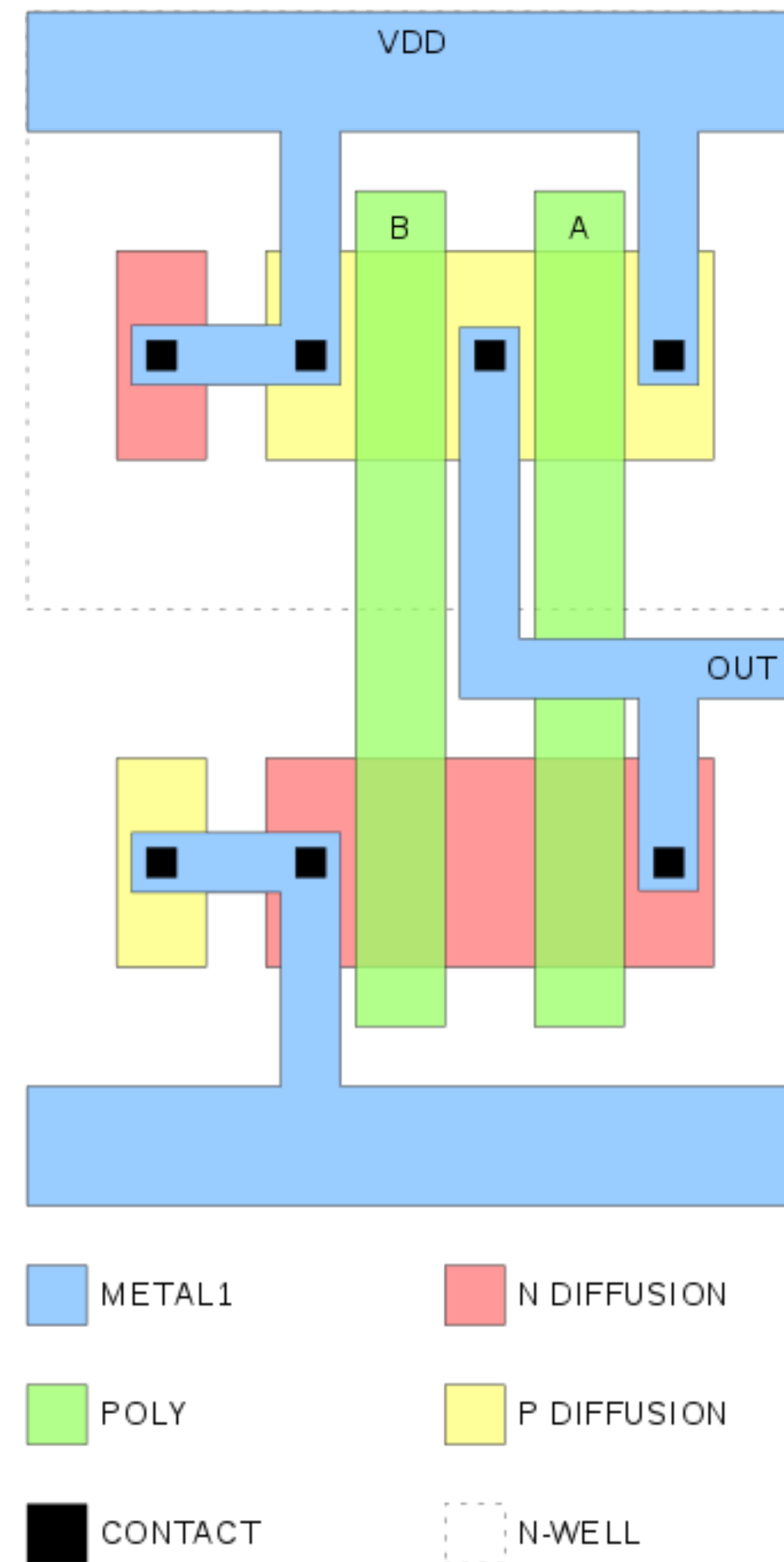
- Effective Out = $\overline{AB}$
- Always **on** unless both inputs are **on**



CMOS in person

Consider an NAND gate

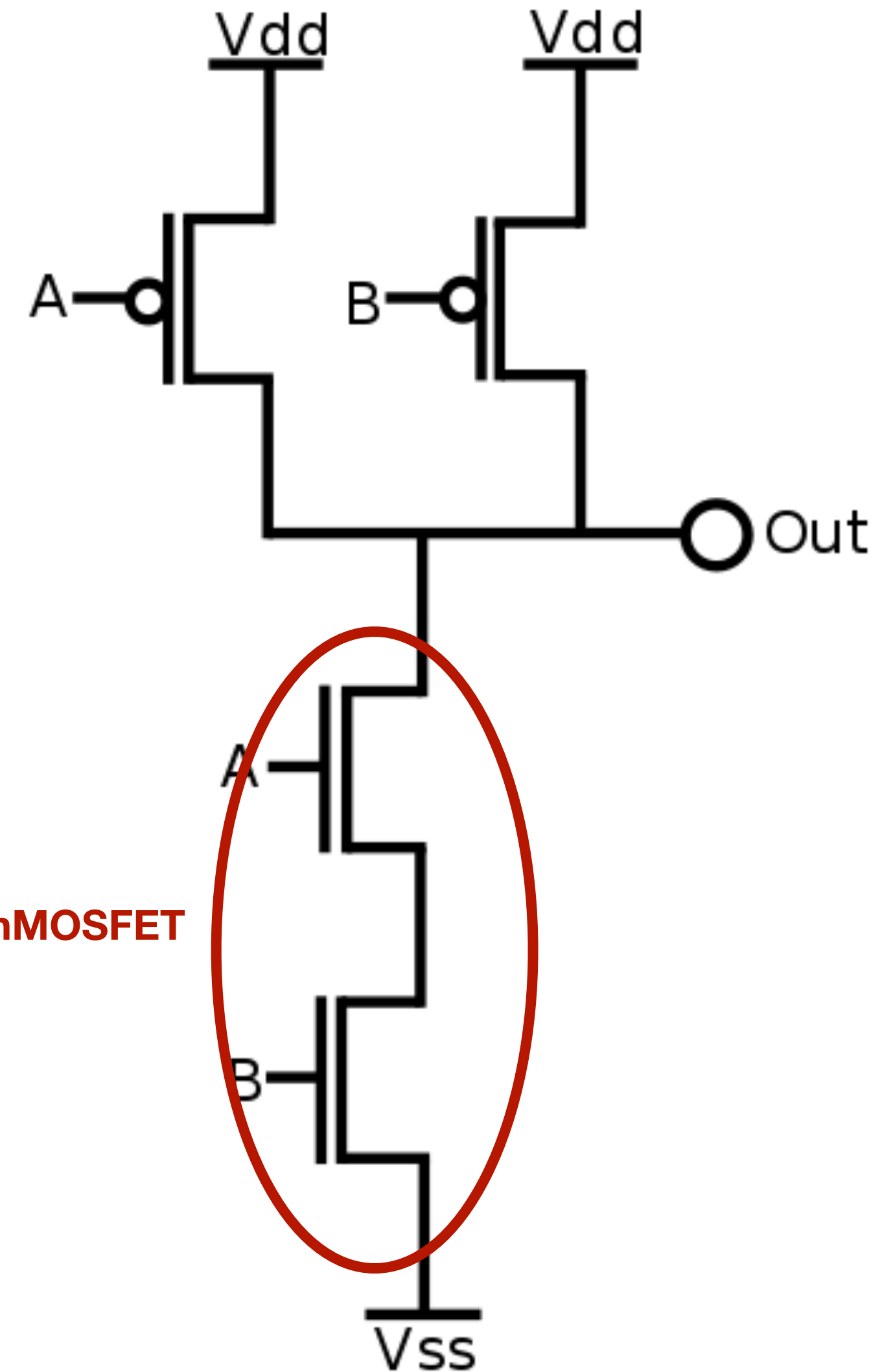
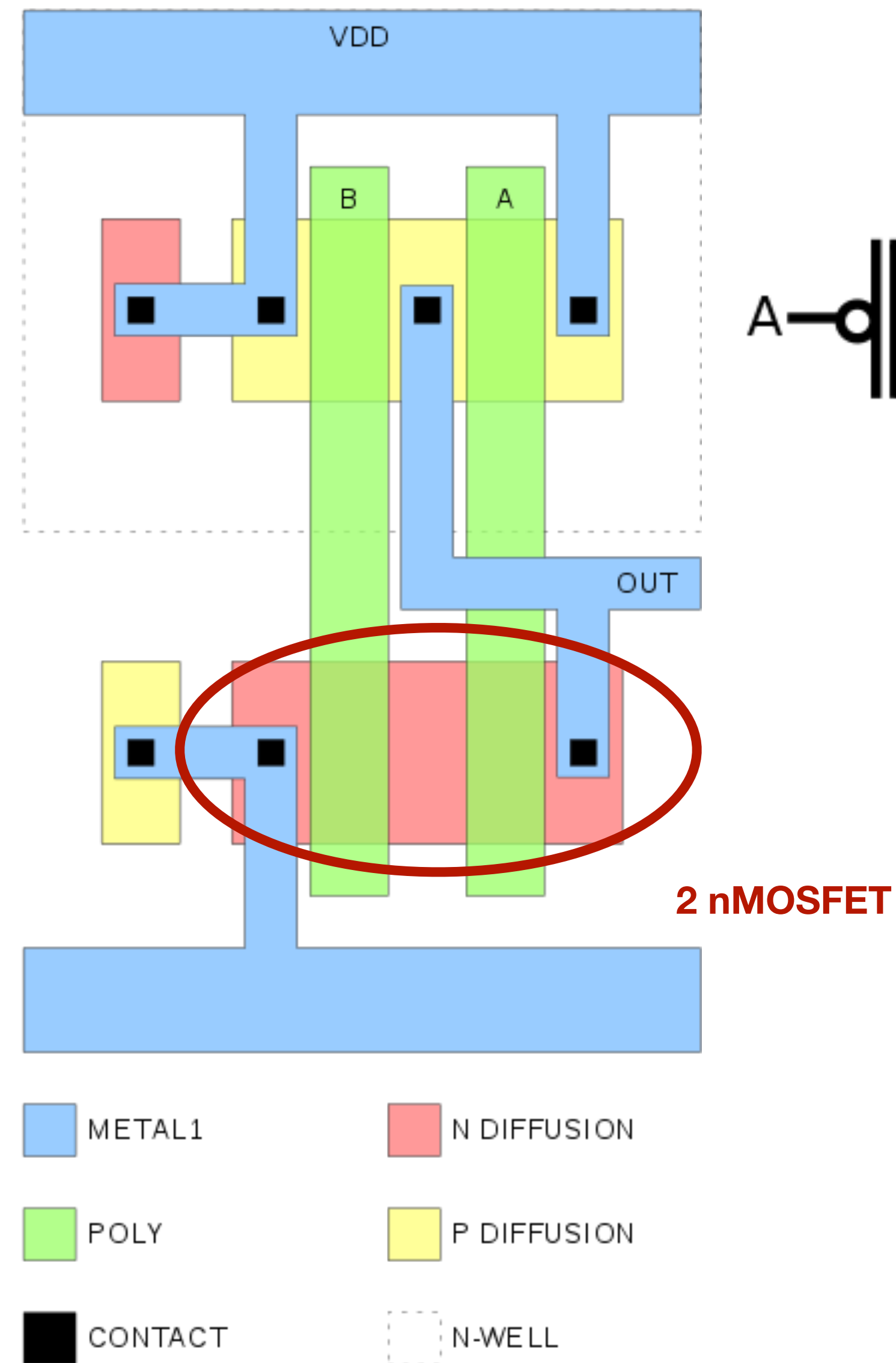
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- Useful to look at the device layout to see what this actually looks like



CMOS in person

Consider an NAND gate

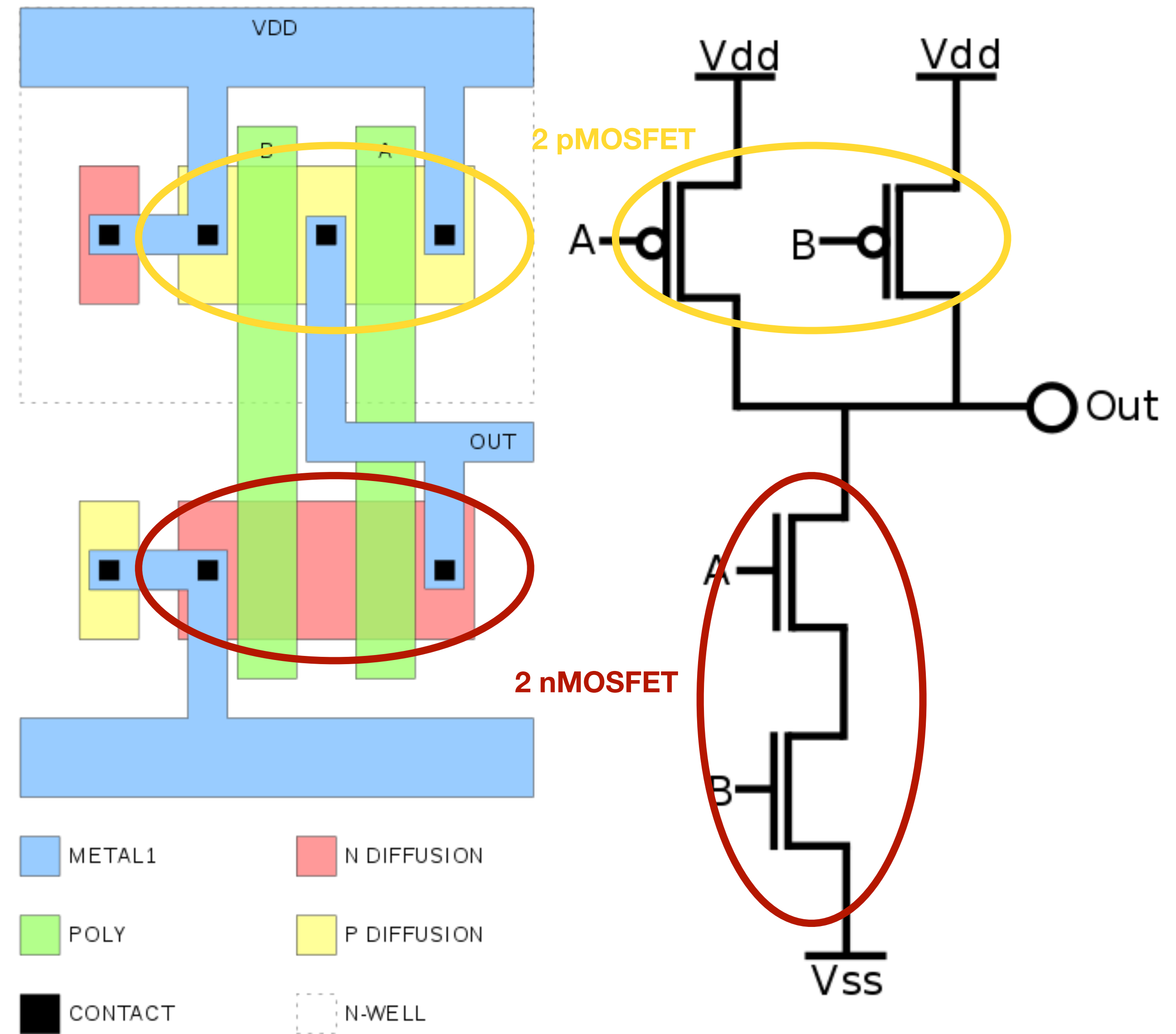
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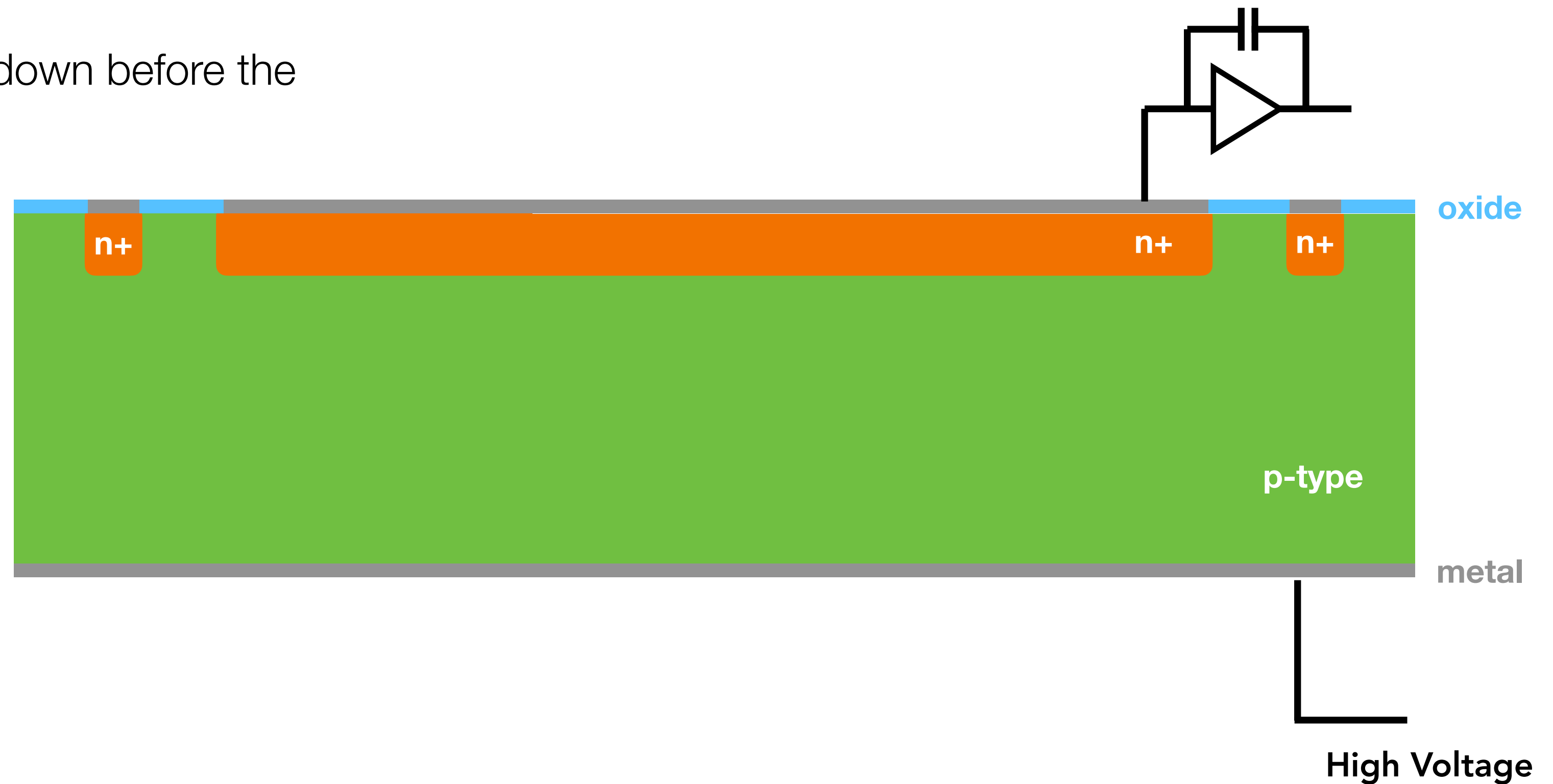
Light-Catching Rock Slices that See the Smallest Pieces of the World

(Semiconductor detectors)

Basic anatomy - 0D

One of the most basic types of detector is simply a PN-diode

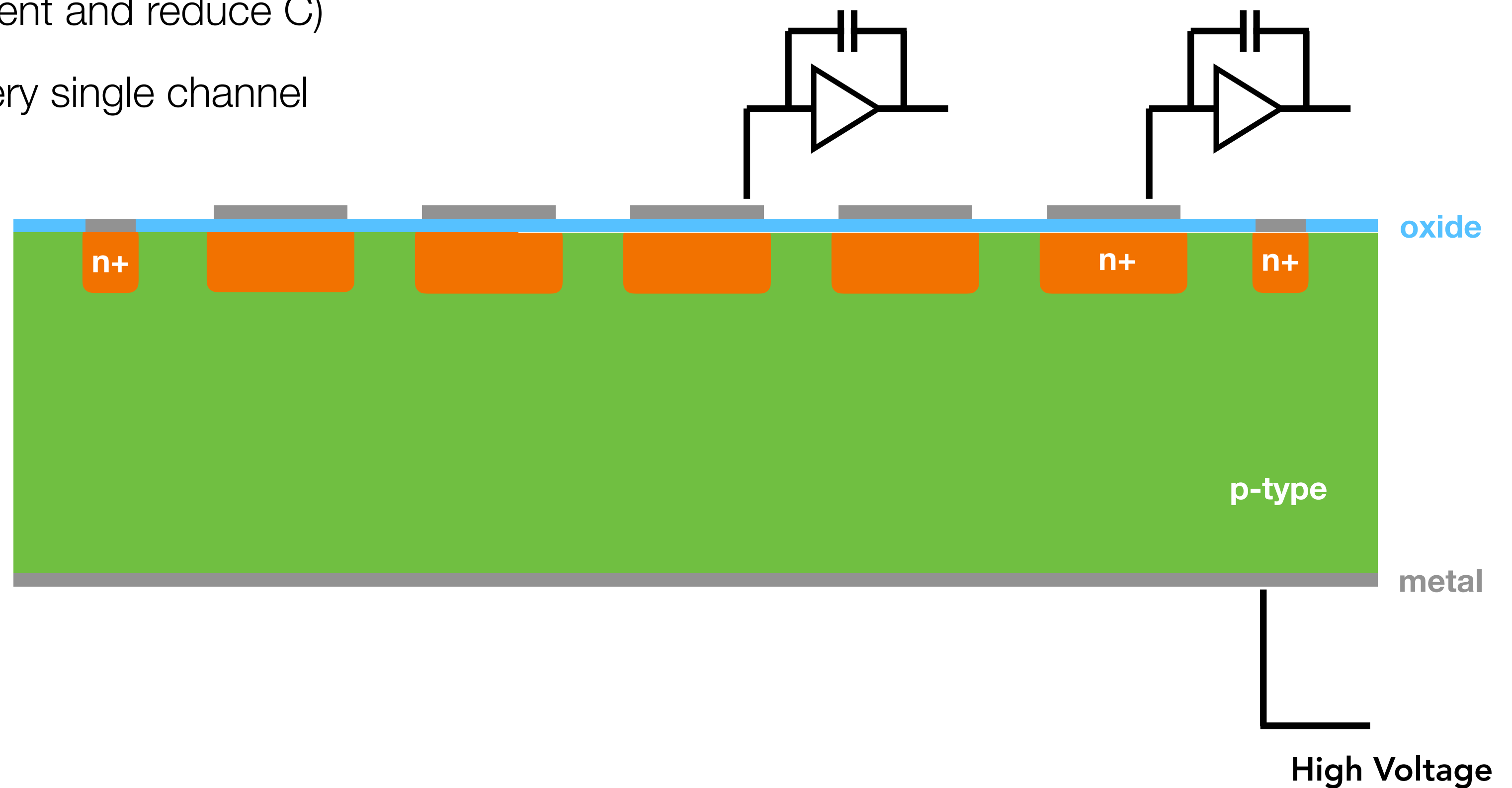
- Large area, highly doped region $\sim 10^{18} \text{ cm}^{-3}$
- High resistivity bulk $\sim 10^{12} \text{ cm}^{-3}$
- *Guard rings* to step the high voltage down before the physical edge
- Back-side HV contact
- Top side contact to read out



Basic anatomy - 1D

Moving on to a single, long *strip detector*

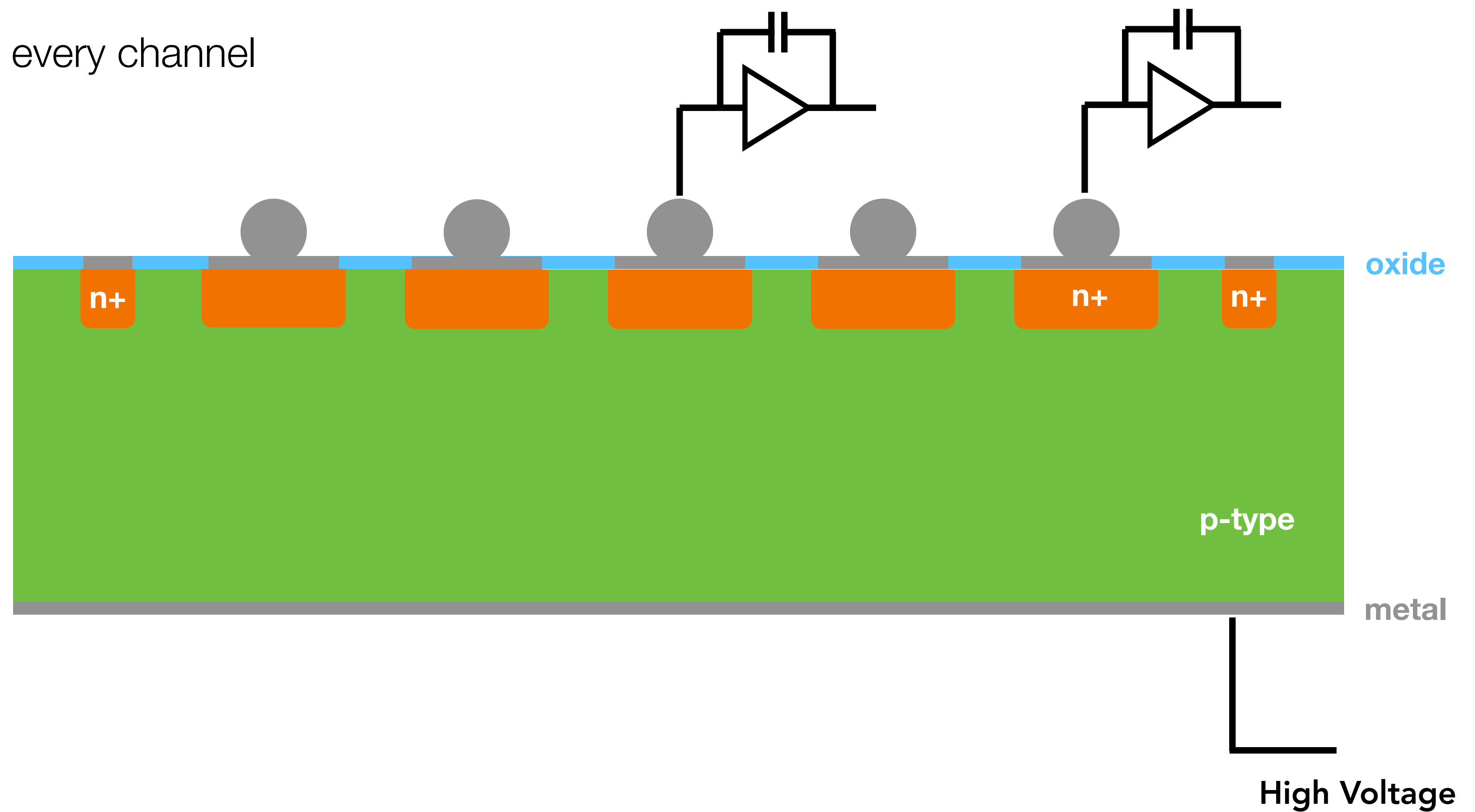
- Strips many cm long, with capacitive readout of the strips rather than DC (to avoid leakage current and reduce C)
- Dedicated electronics to read out every single channel individually on a separate ASIC
- Wire bond each strip “manually”



Basic anatomy - 2D

Finally, pixel detectors

- Now segmented in 2D
- Typically separate readout ASIC, with every channel *bump-bonded* to its own diode



Making Lightning Messages

(Signal formation)

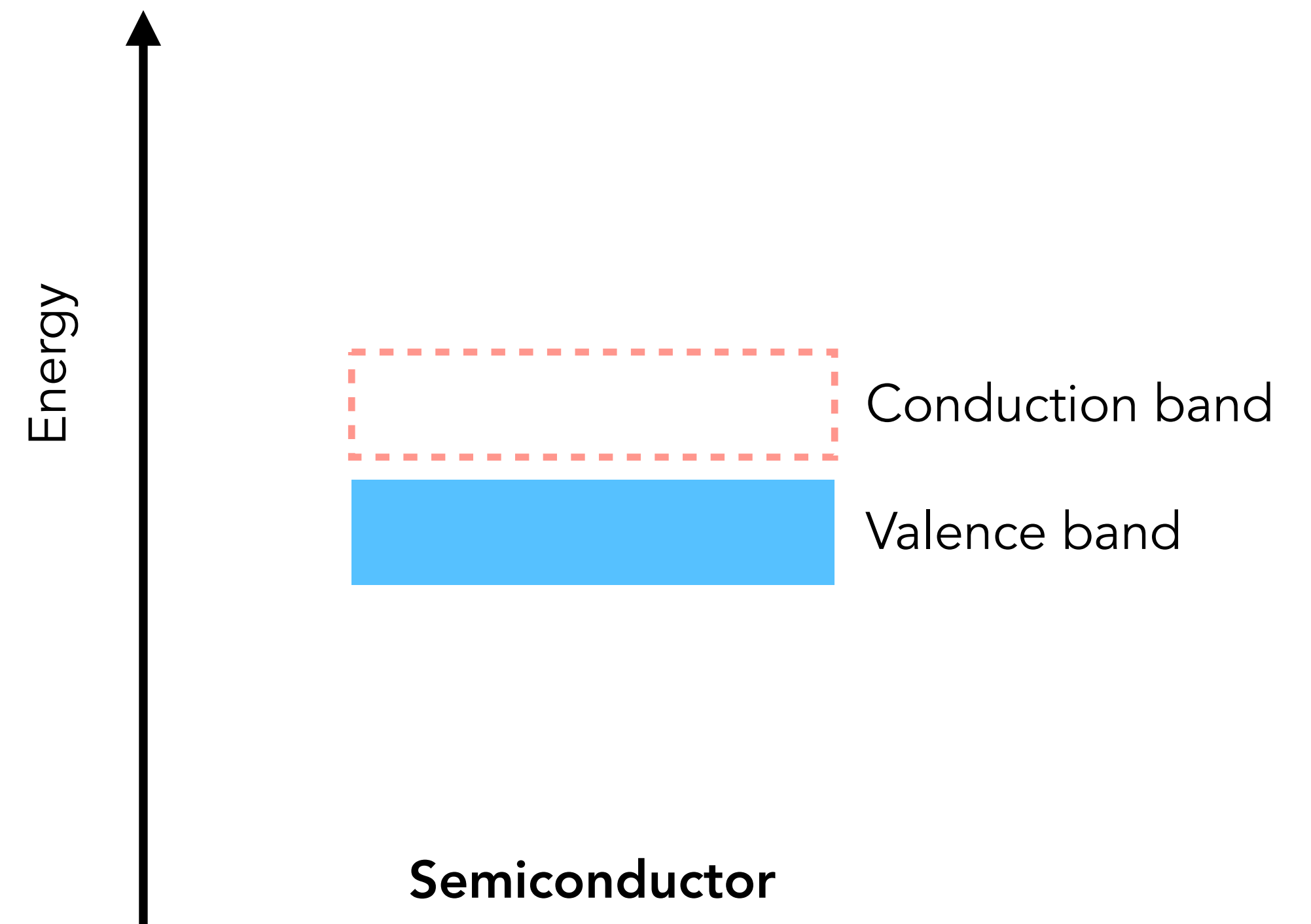
Charge carrier creation

As we depleted our PN-junction, we removed the free charge carriers

- We can continue to do this until the function is *fully depleted*

Under these circumstances the only free charge carriers moving around will be those that have enough thermal energy

- This is a small number for silicon at room temperature



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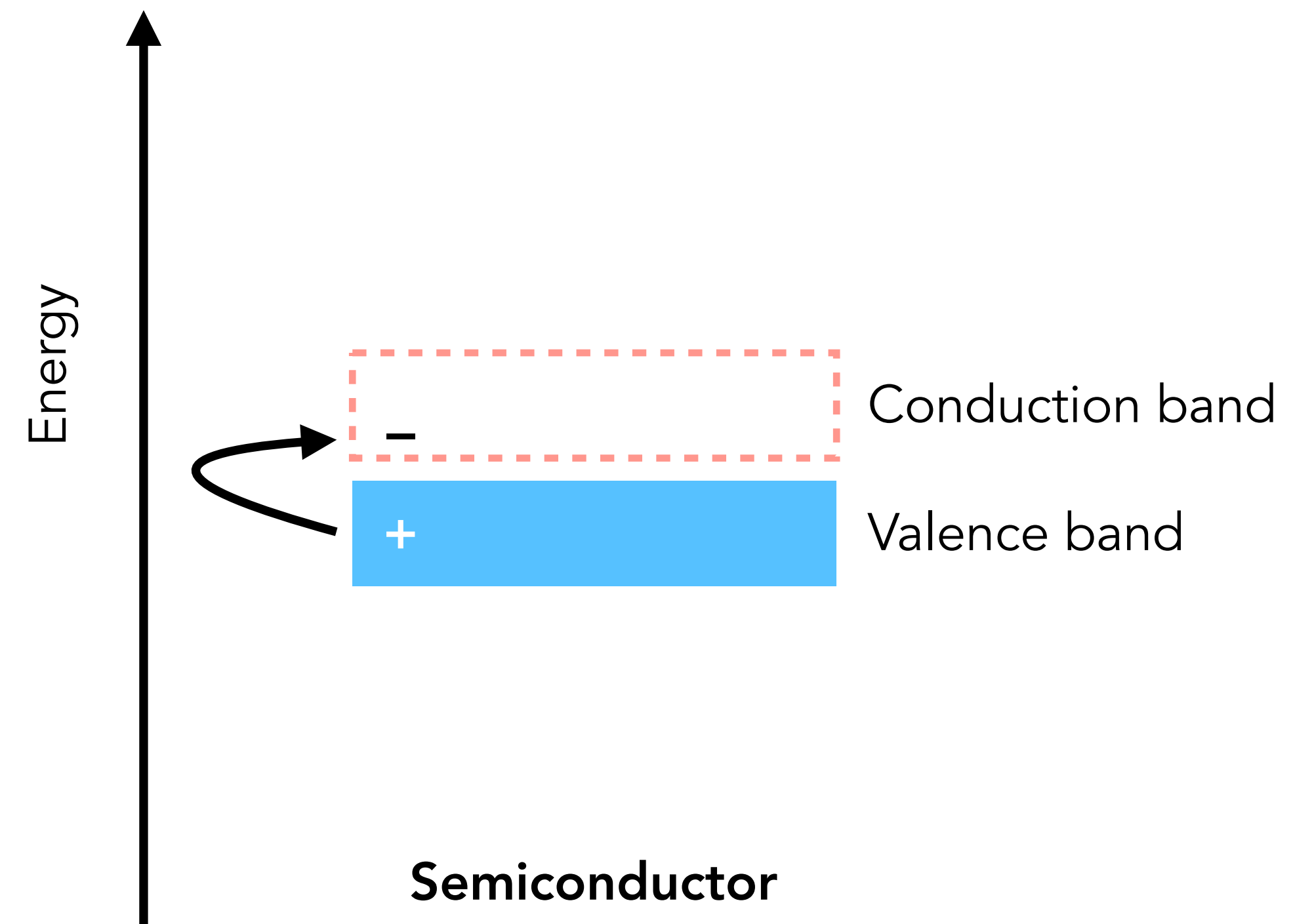
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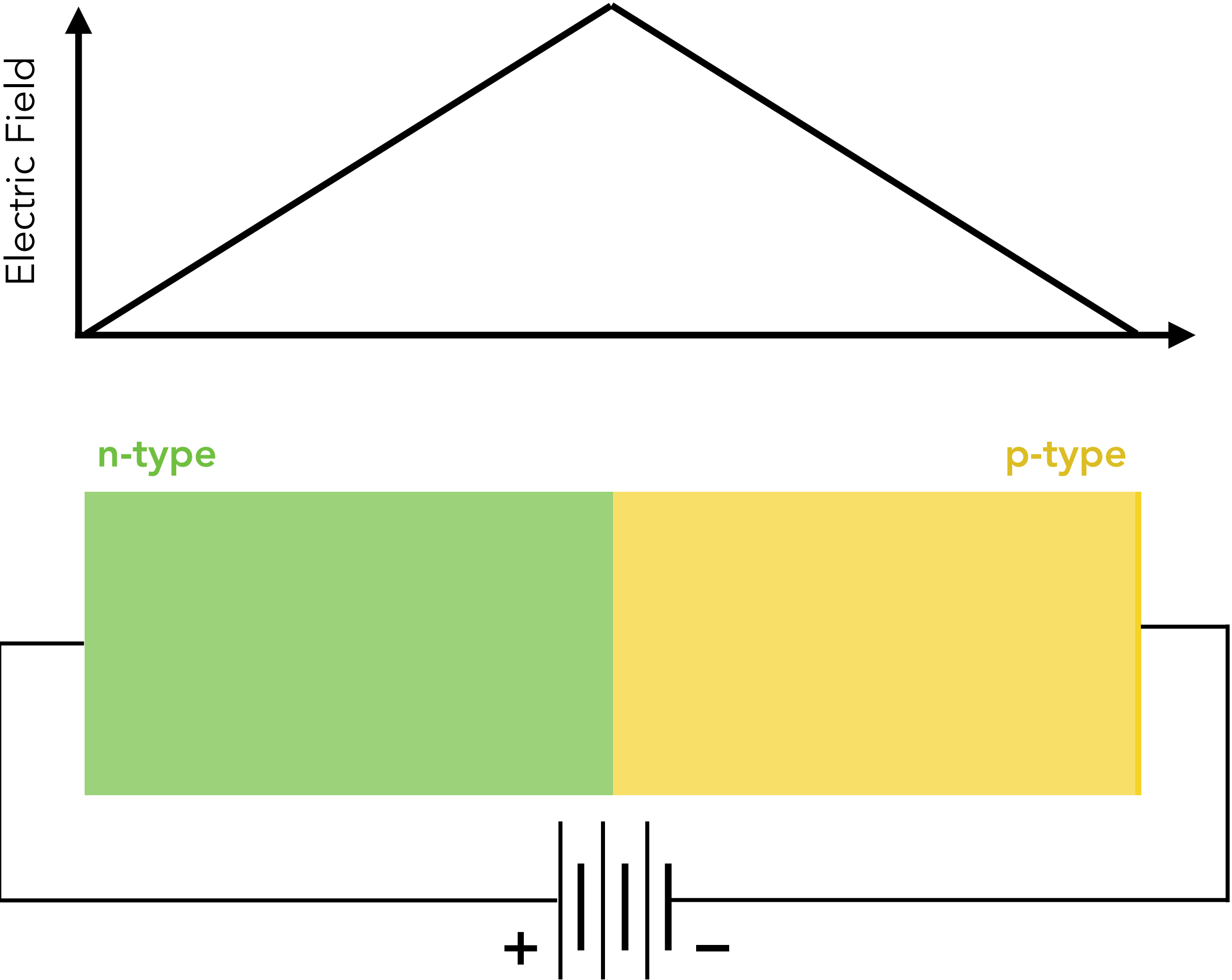
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If a particle interacts with the silicon, it can transfer energy onto individual electrons and promote them to the conduction band

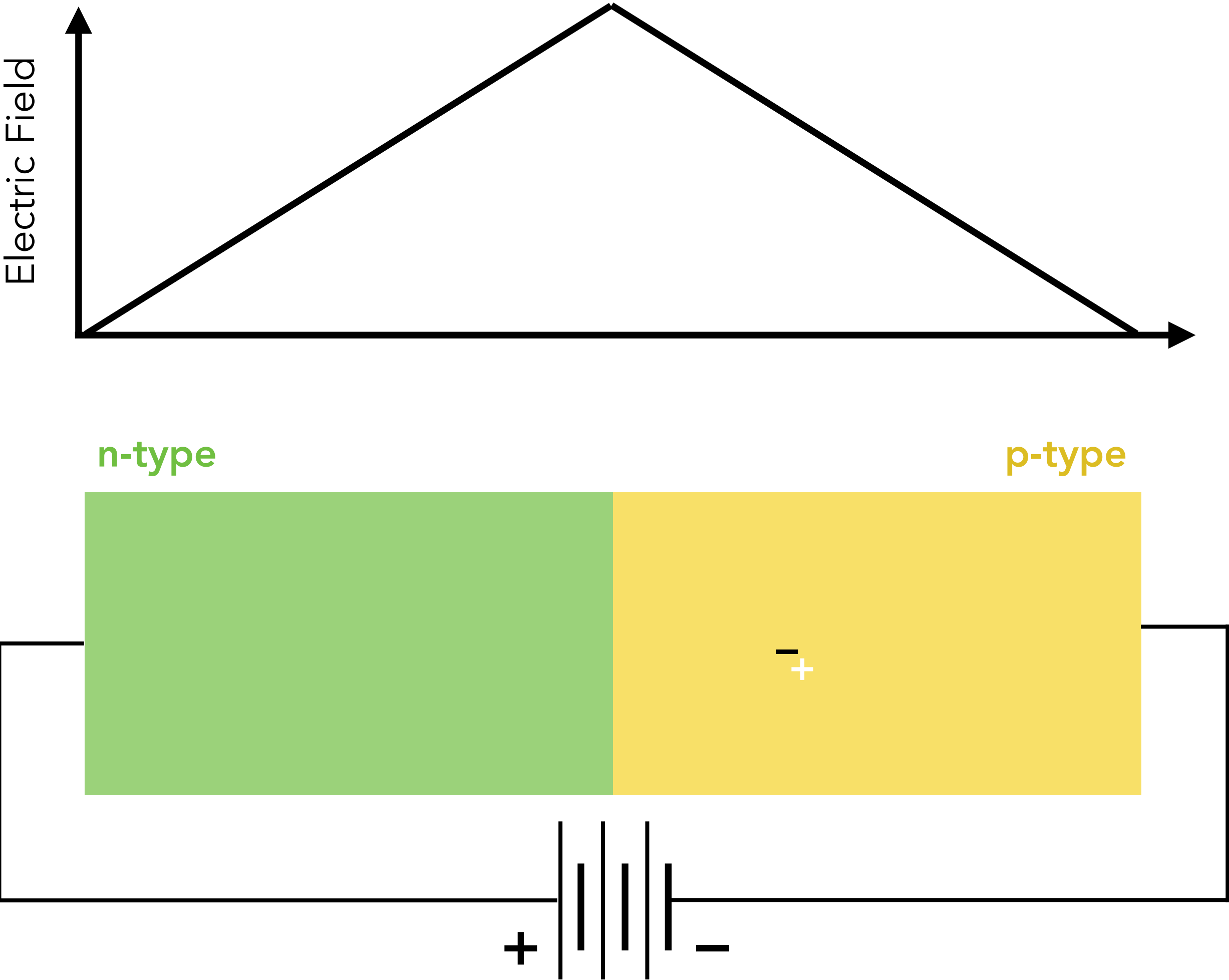
- This is the foundation of solid-state particle detection



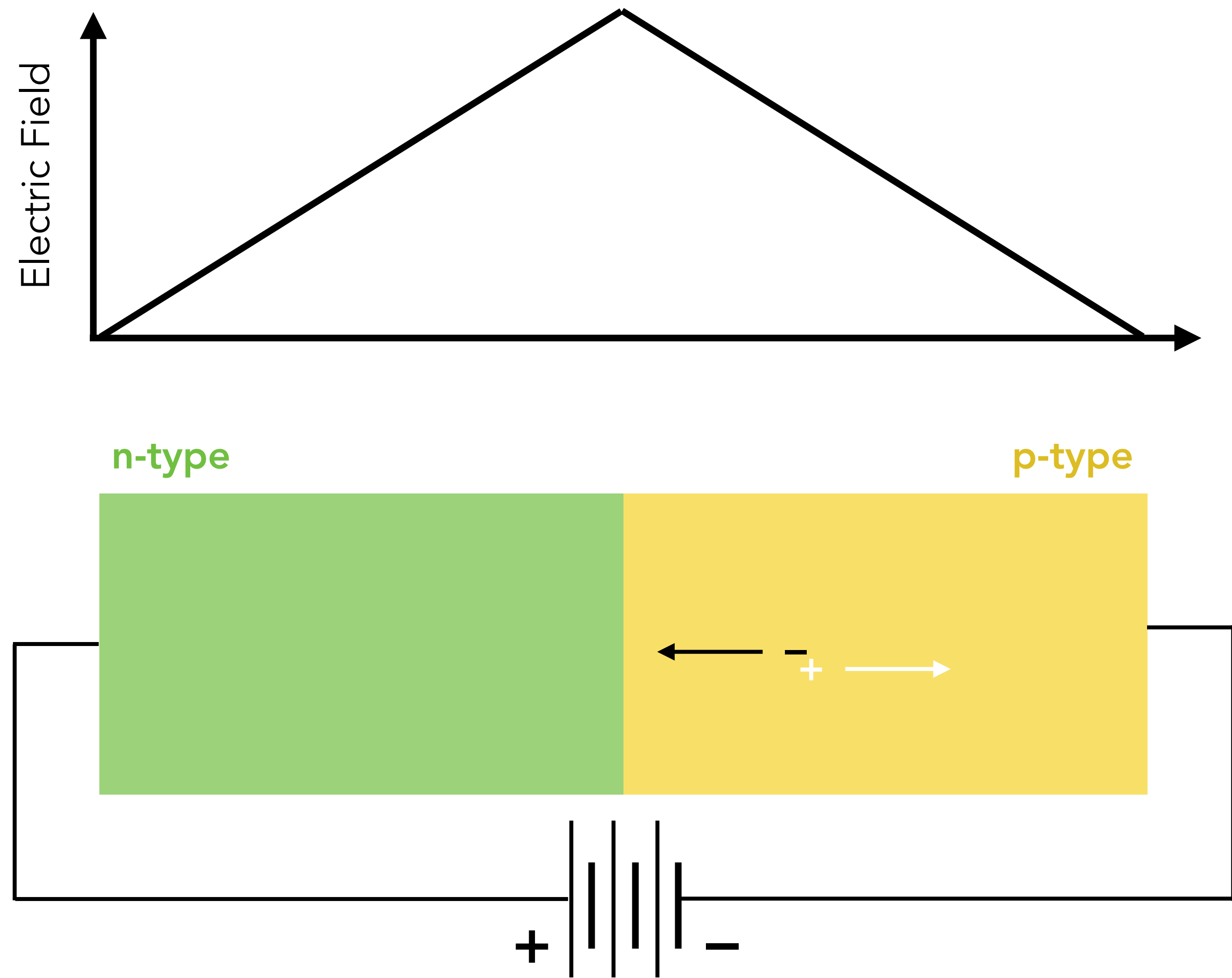
Charge carrier motion



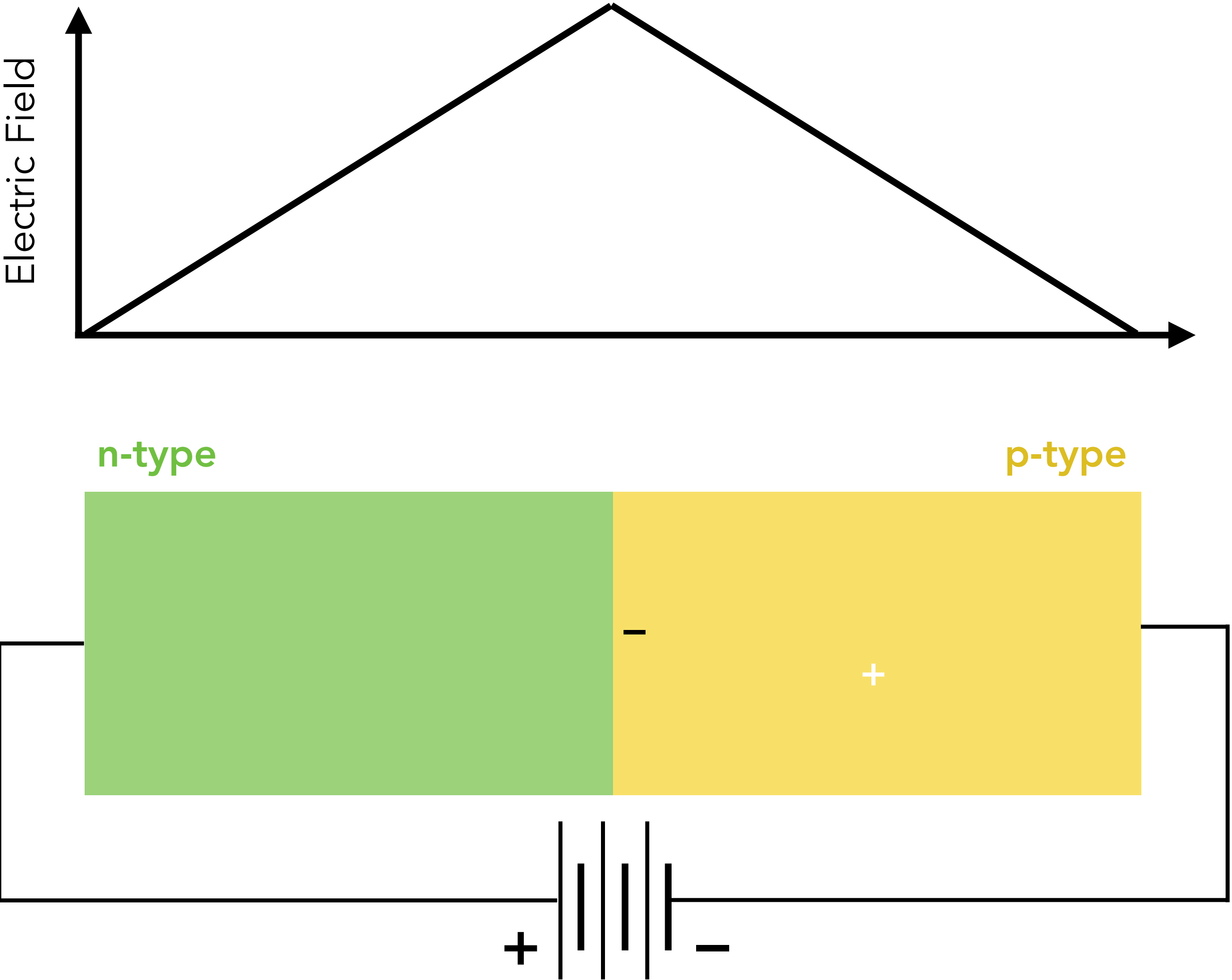
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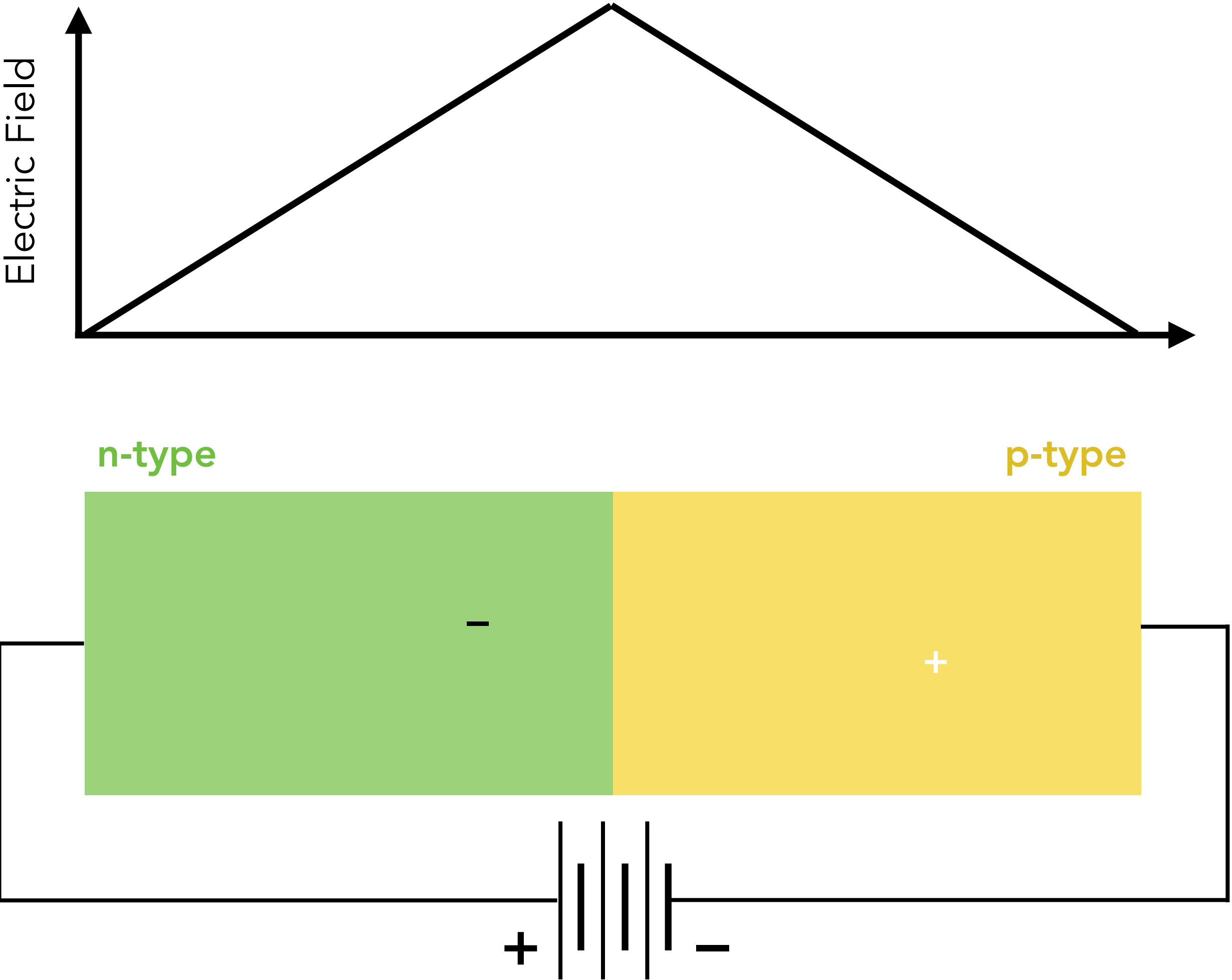
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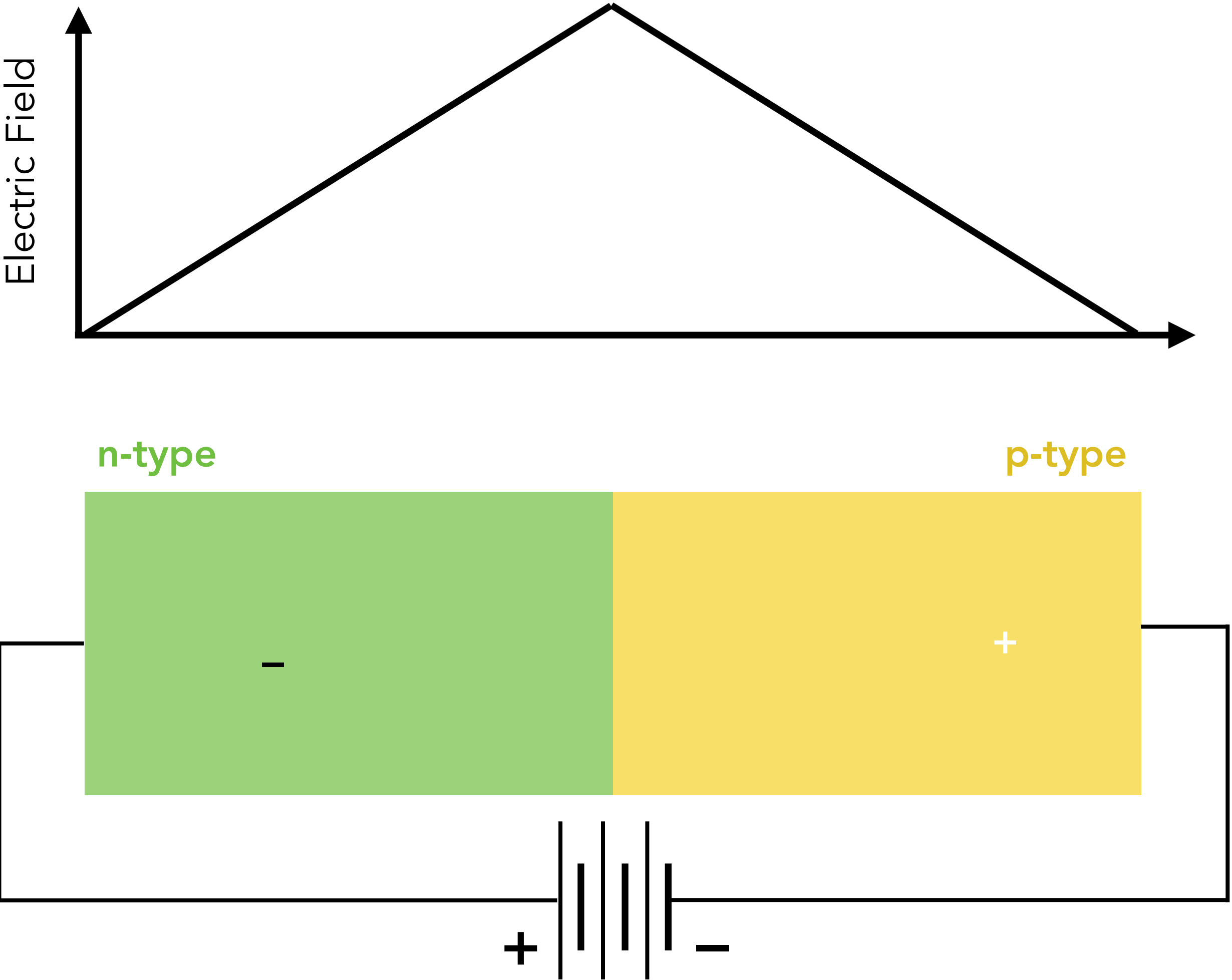
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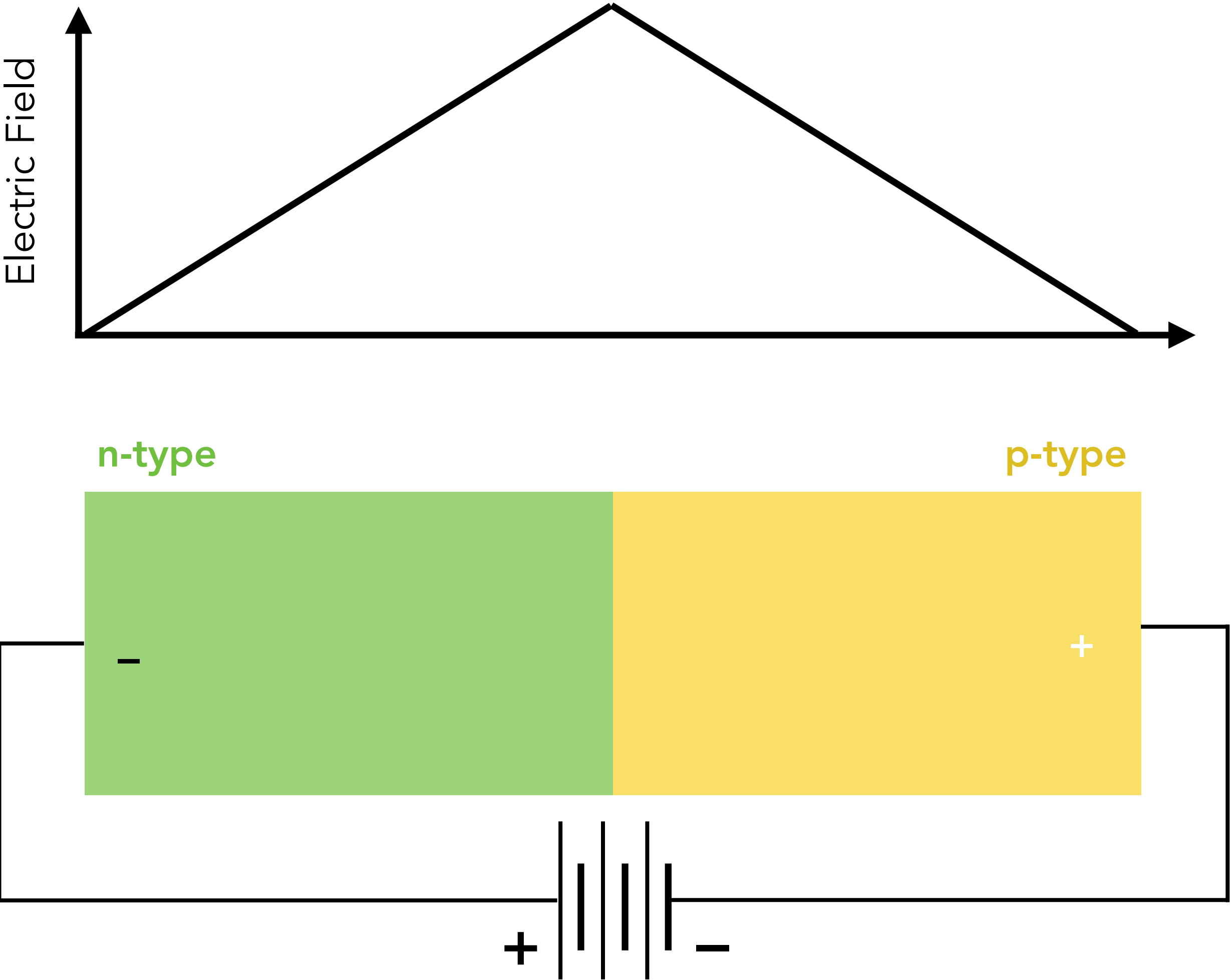
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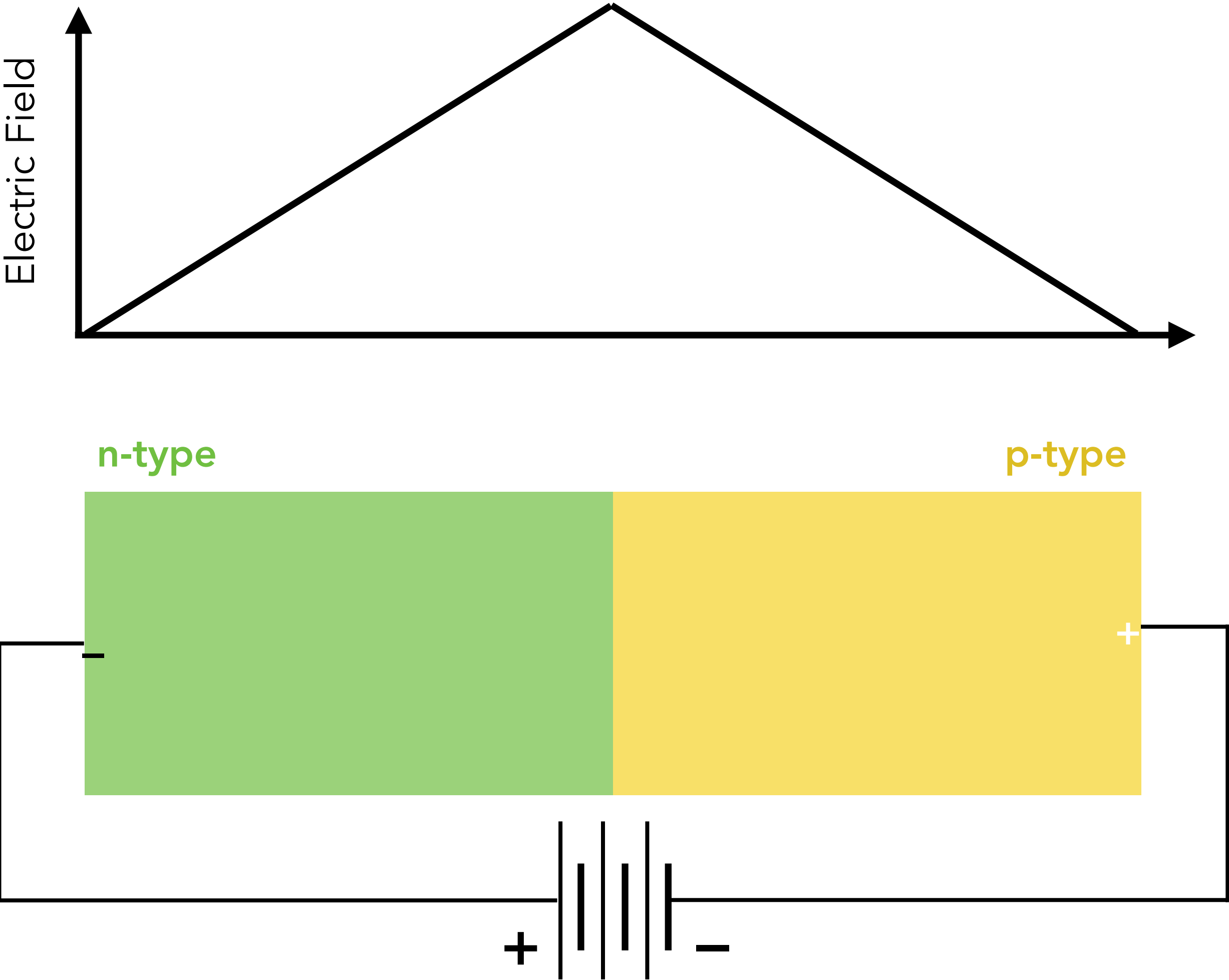
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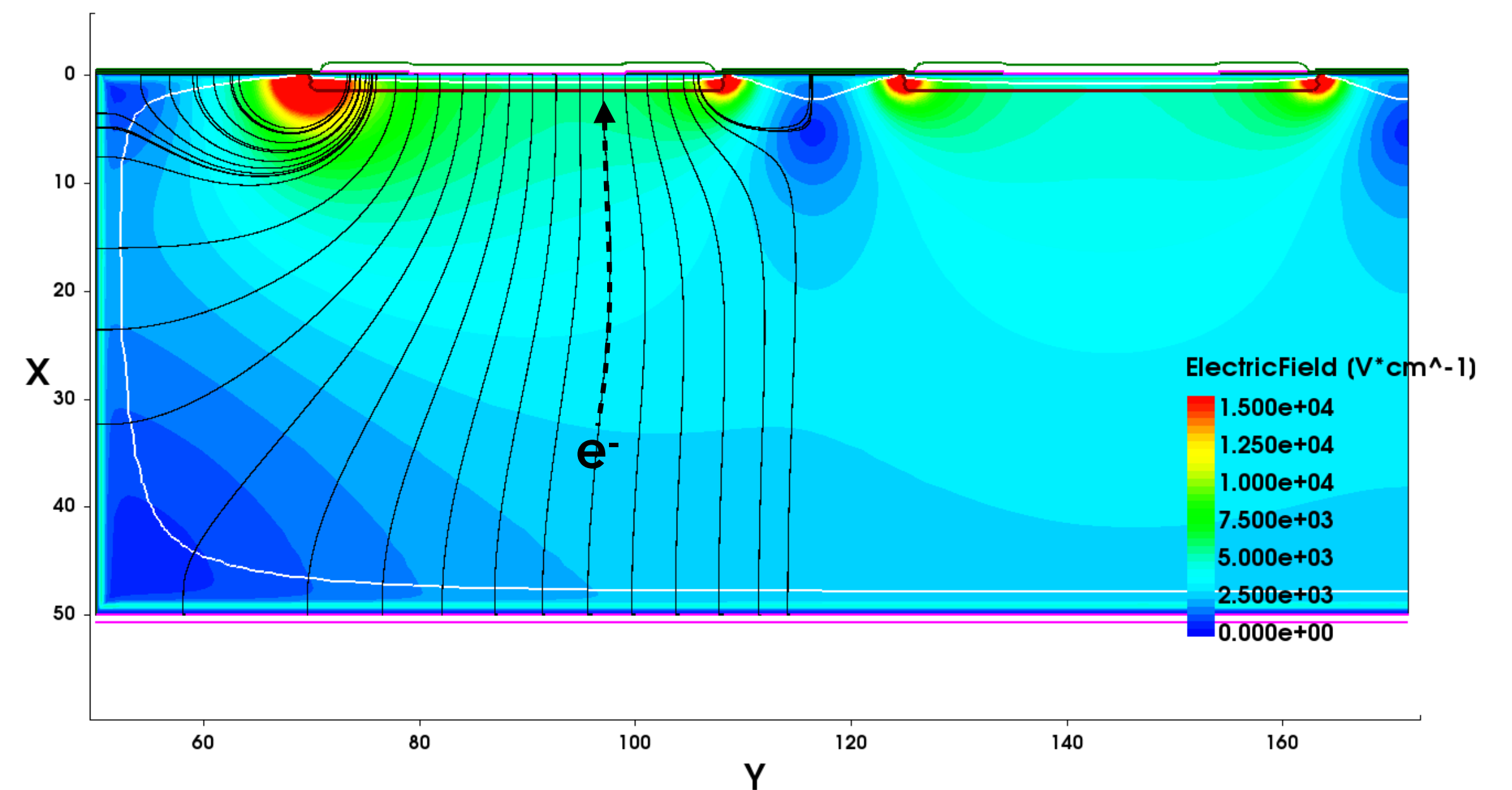
Charge carrier motion



Charge collection

Typically we only collect one species of charge carrier
- often electrons (as they travel faster) but historically p-on-n detectors were popular due to availability of high-resistivity n-type wafers

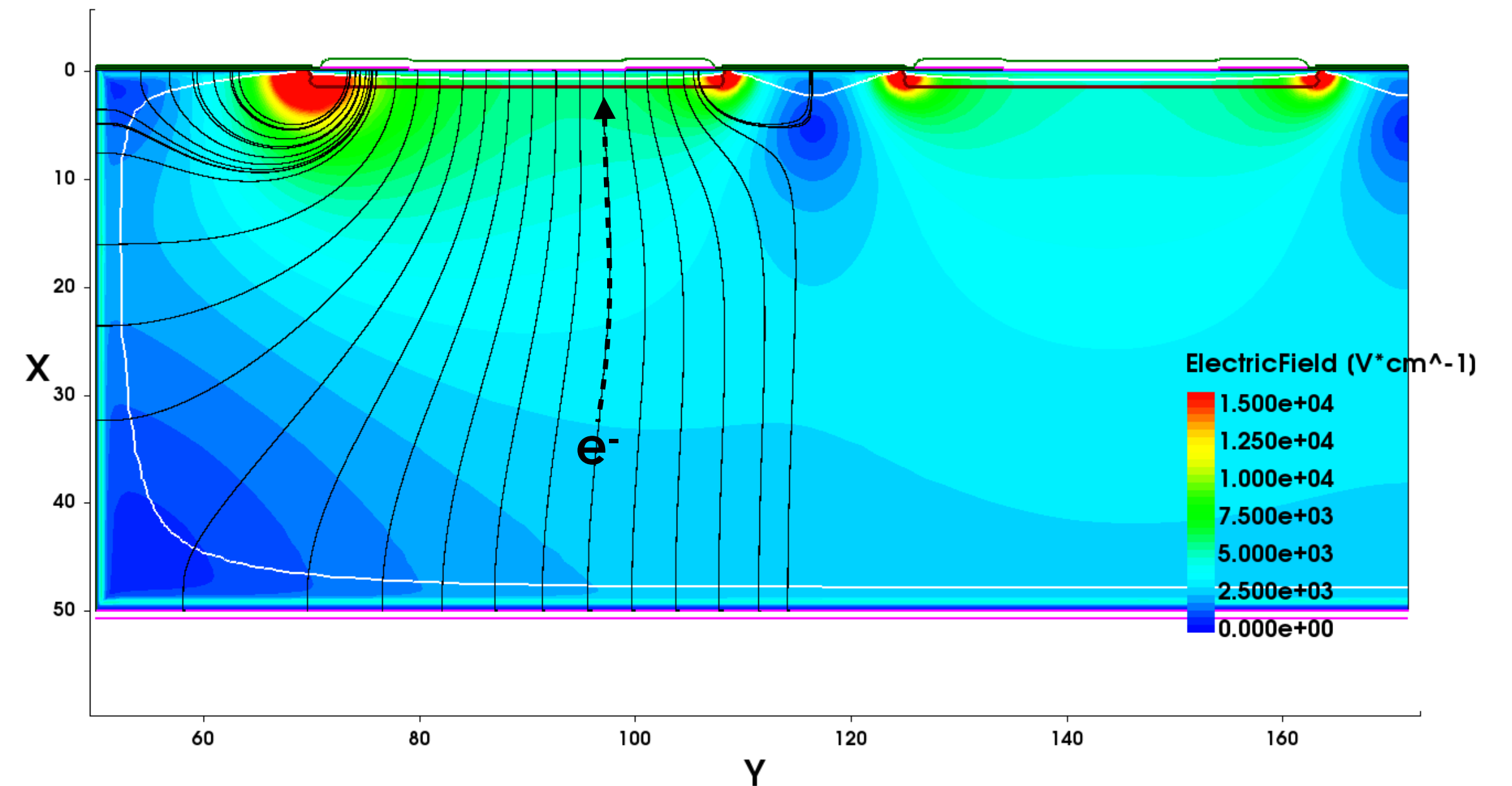
- Built-in depletion widths of only a few microns
- Charge carriers follow the electric field lines within the sensor



Induced current

In a simplistic approach, one can consider the signal at the collection node as being the sum of all of the charges that end up there

- This is true, but in fact a signal is generated as *each charge carrier starts to move*

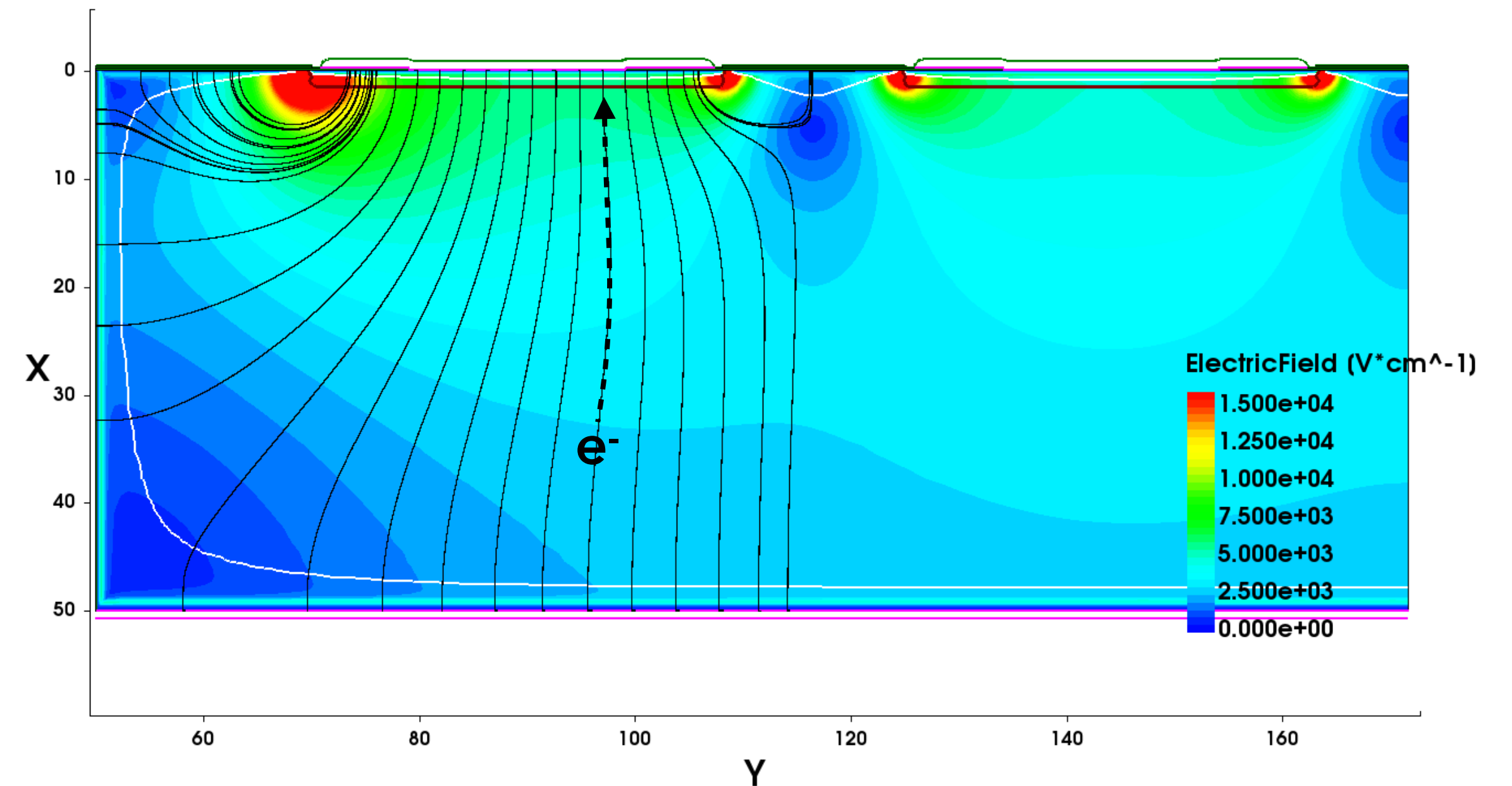


Induced current

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- This is described by Ramo-Shockley theorem

$$I_{ind}(t) = -\frac{q}{V_w} \mathbf{E}_n(\mathbf{x}_q(t); V_w) \cdot \frac{d}{dt} \mathbf{x}_q(t)$$



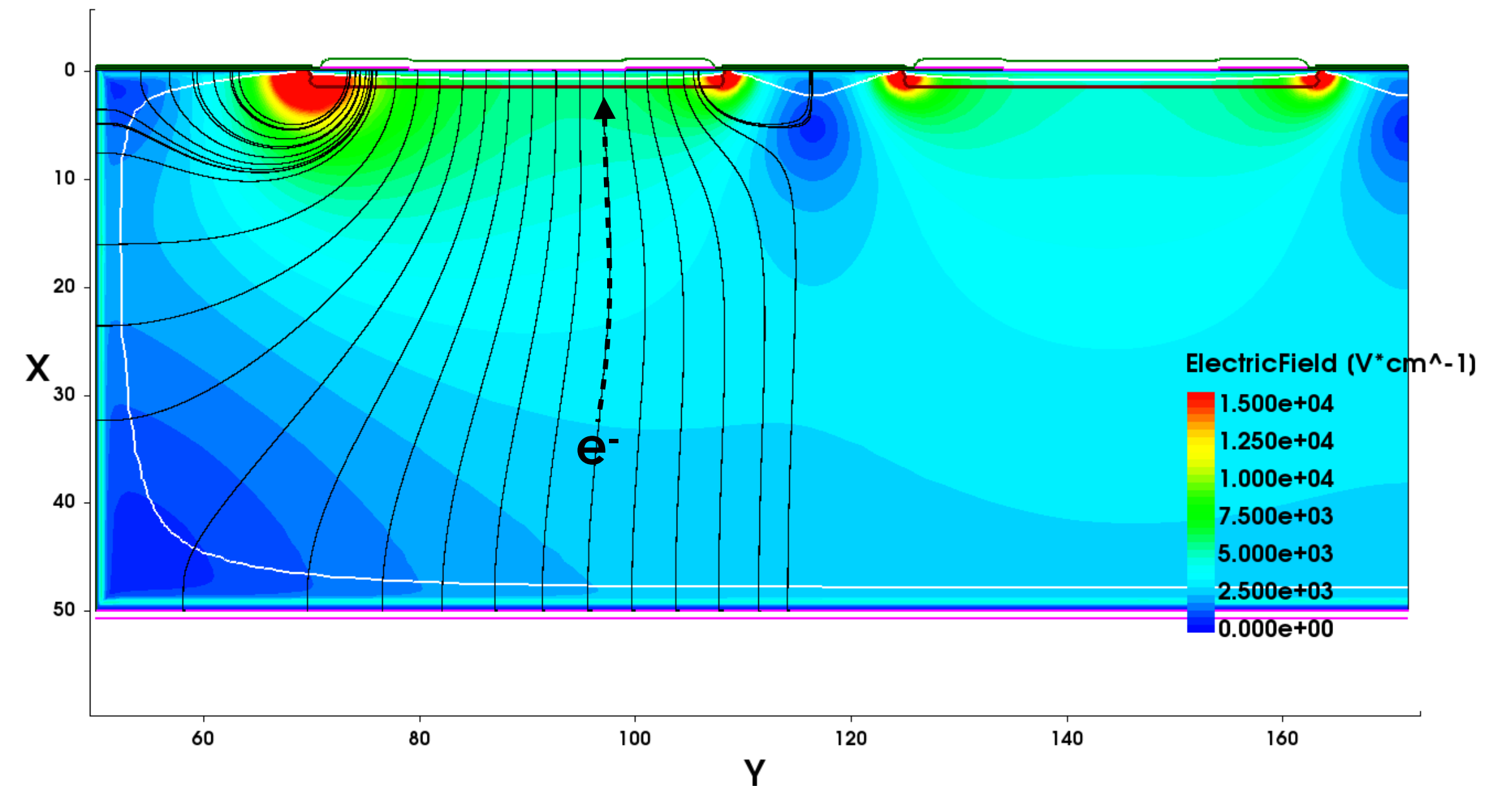
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- This is dictated by the **weighting field**



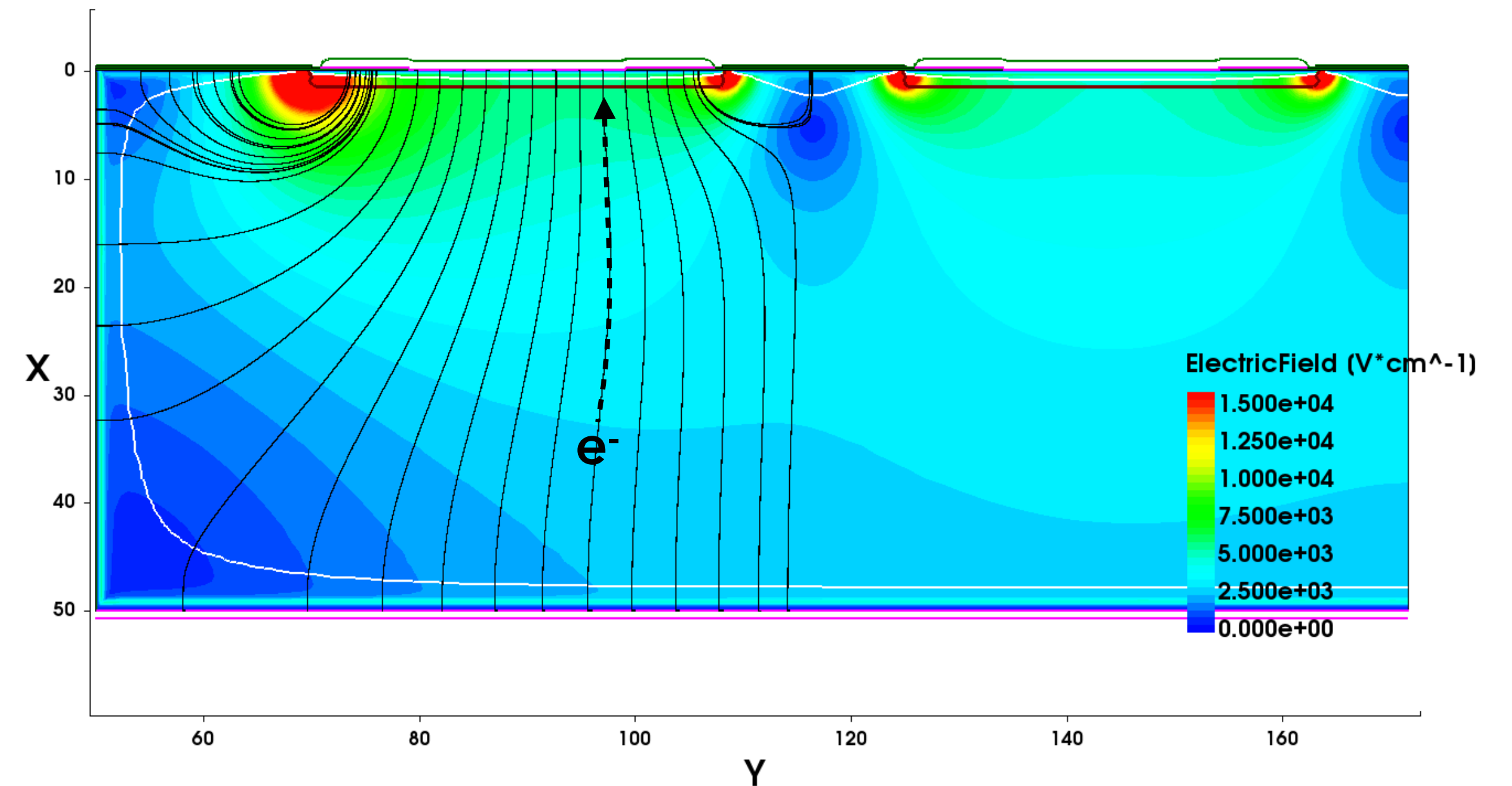
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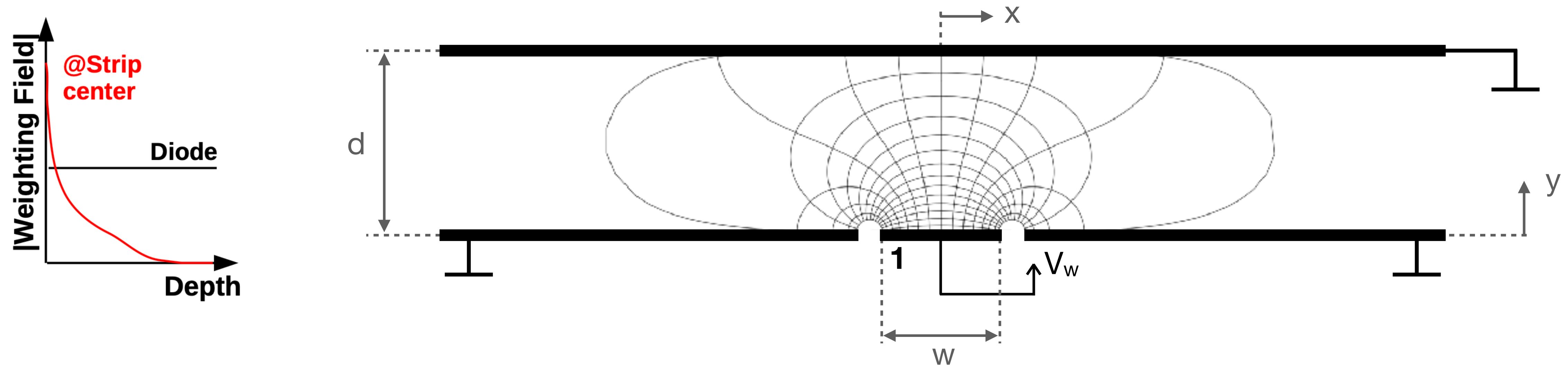
- This is dictated by the **weighting field** and the **velocity**



Weighting fields

The weighting field is effectively the coupling between the moving charges and the ‘antenna’ which is detecting them

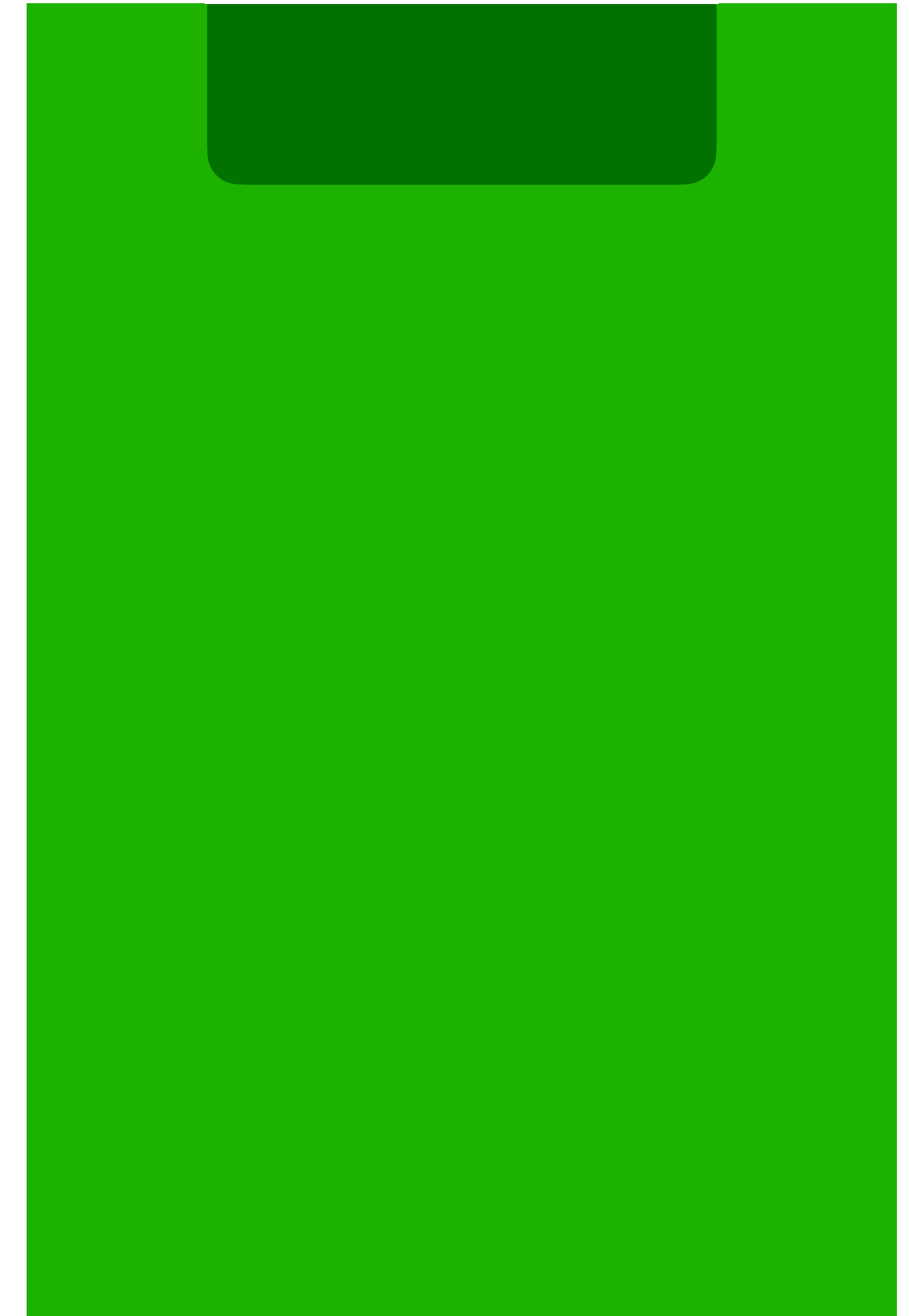
- Calculated by placing a unit potential on the contact of interest and grounding every other contact
- For a large pad detector this is effectively a parallel-plate capacitor: uniform weighting field
- For pixels/strips this is highly non-linear and very sharply peaked close to the pixel implant



Induced current pulses

Transient Current Technique

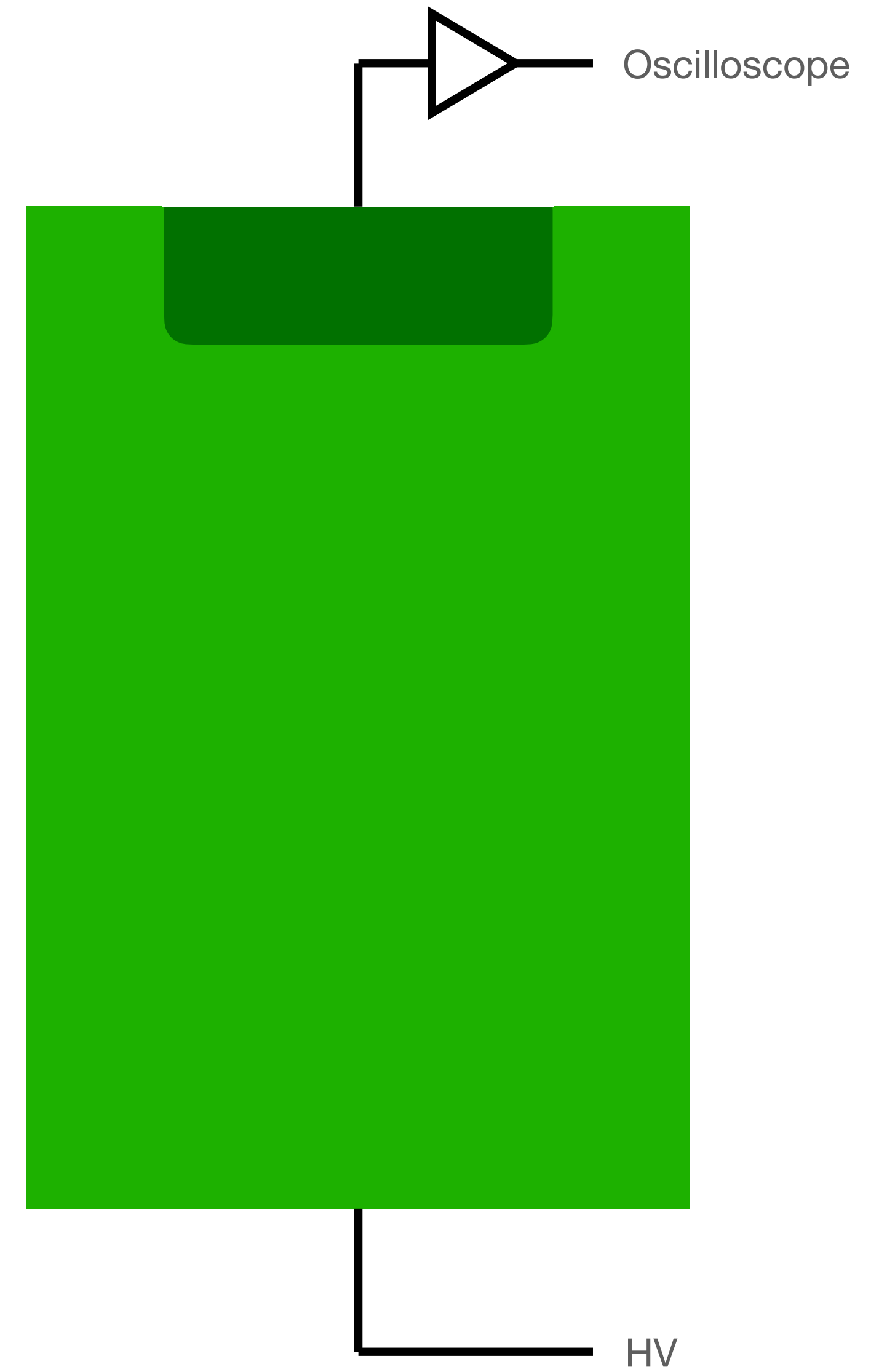
- The whole concept is to look at the induced current on an electrode/cathode while charges move through the device



Induced current pulses

Transient Current Technique

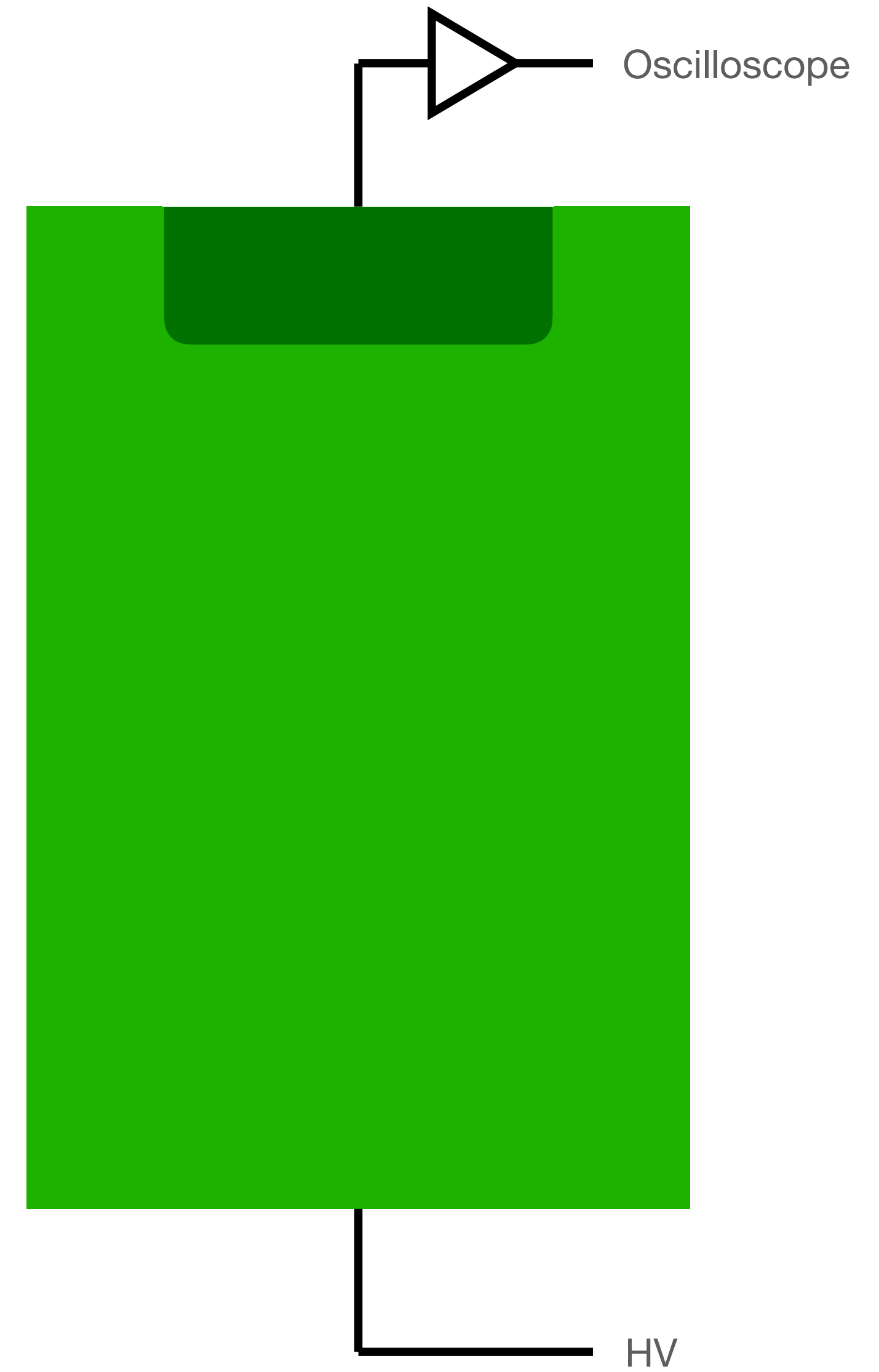
- The whole concept is to look at the induced current on an electrode/cathode while charges move through the device
- Consider charge carriers within the sensor:



Induced current pulses

Transient Current Technique

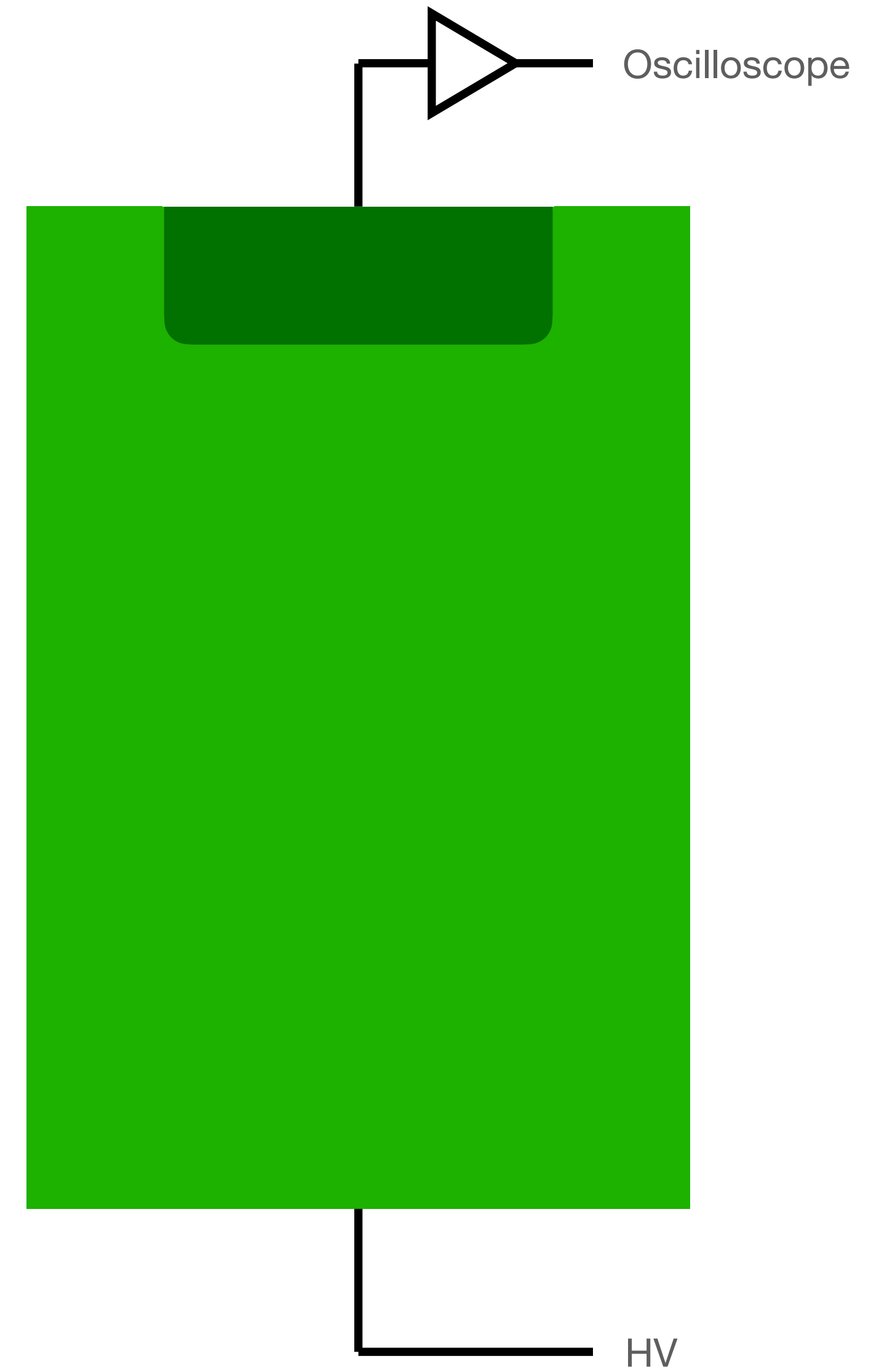
- The whole concept is to look at the induced current on an electrode/cathode while charges move through the device
- Consider charge carriers within the sensor:
 - $v_d = \mu E$



Induced current pulses

Transient Current Technique

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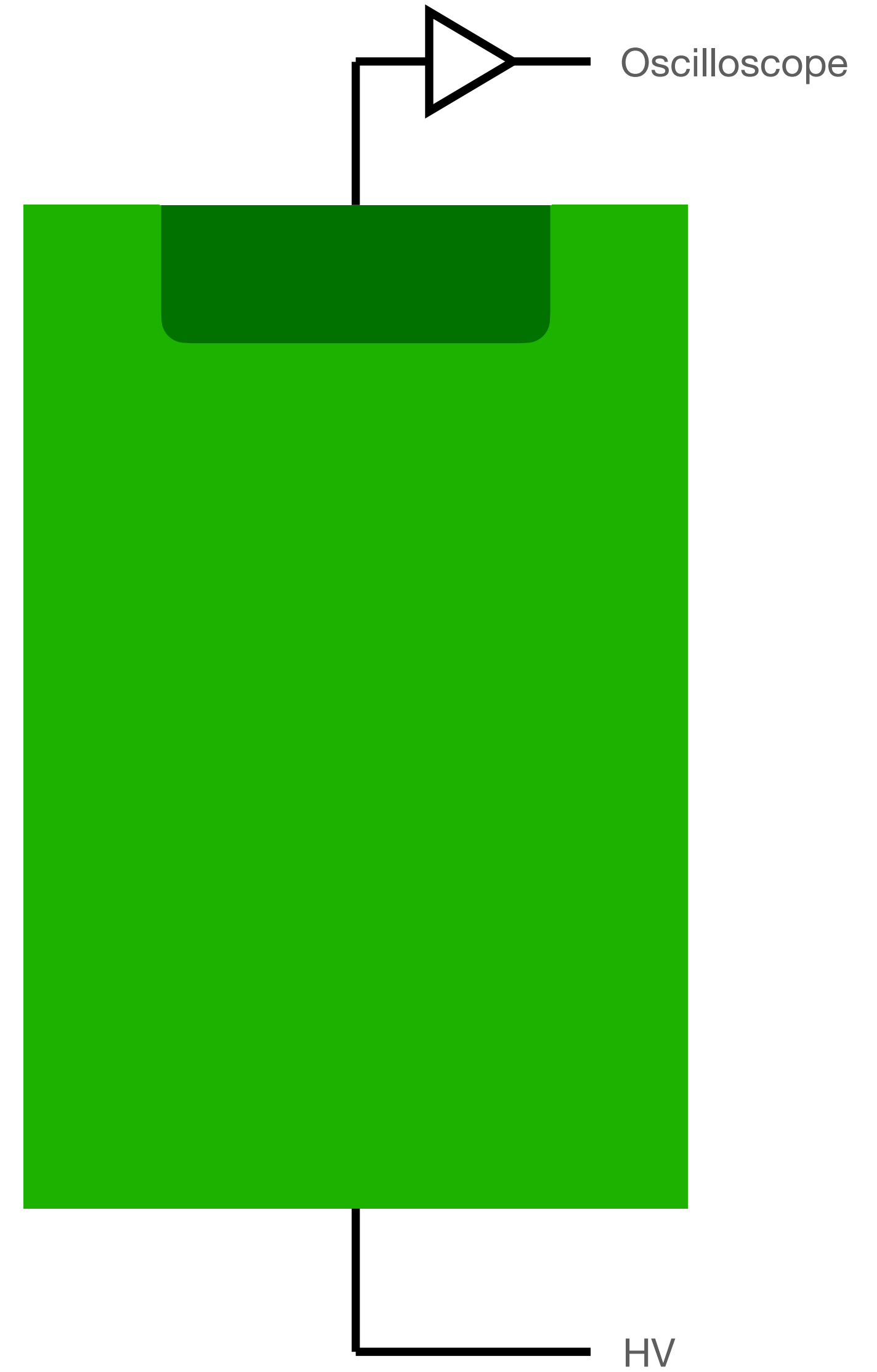
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Induced current pulses

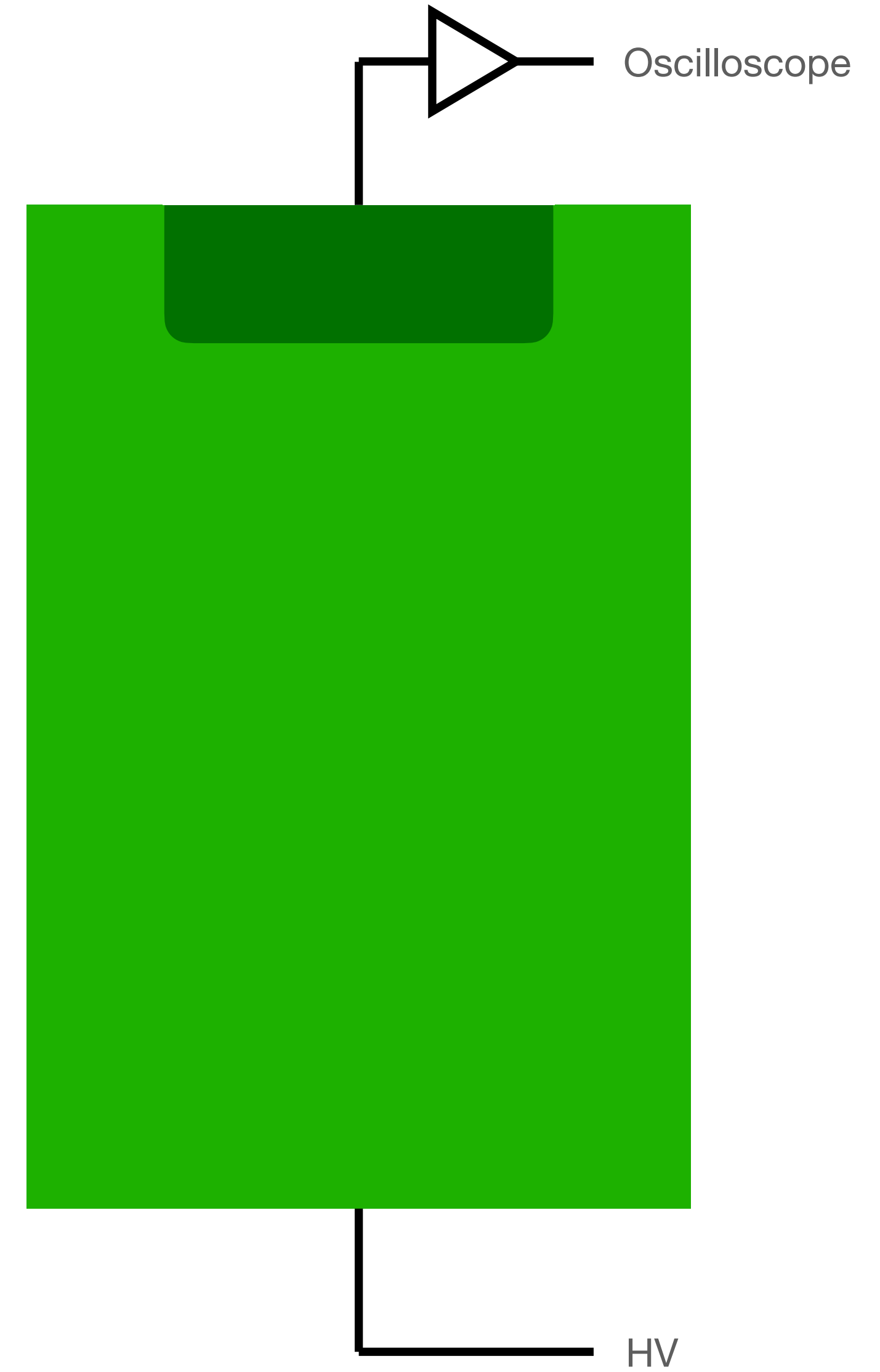
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Weighting field



Induced current pulses

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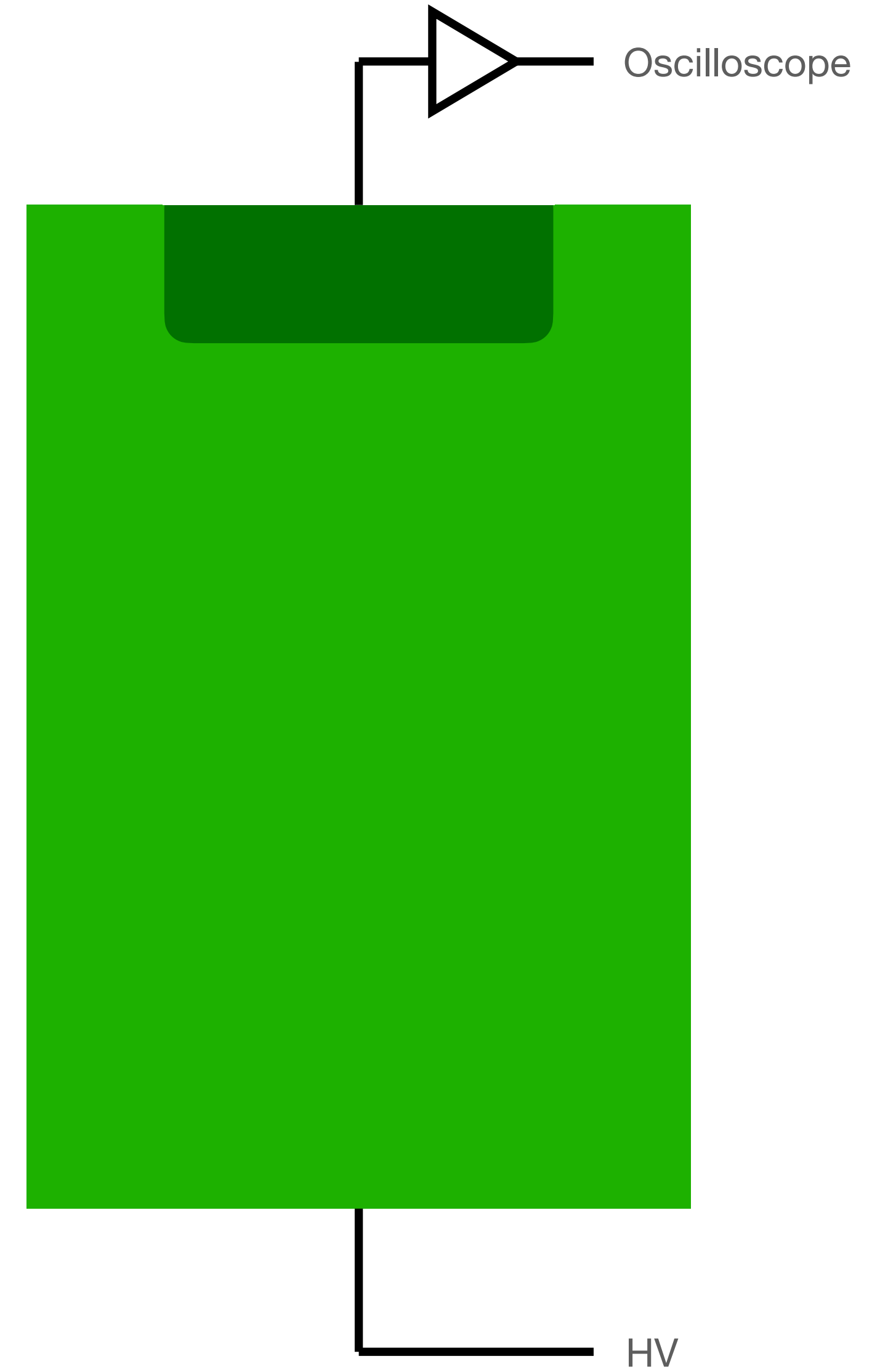
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Weighting field

Electron velocity

Hole velocity

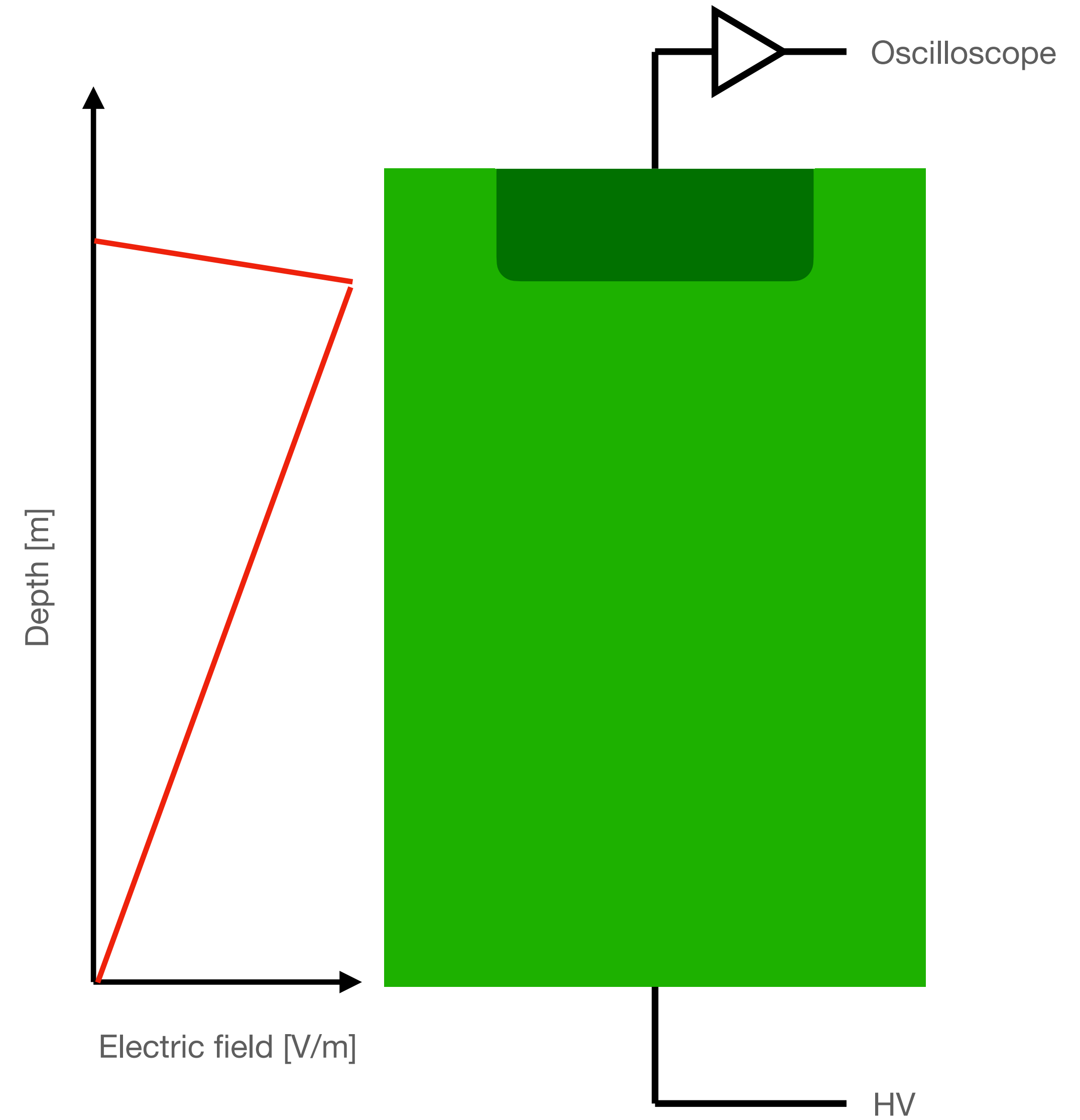
Proportional to Electric field
(below saturation)



Induced current pulses

Transient Current Technique

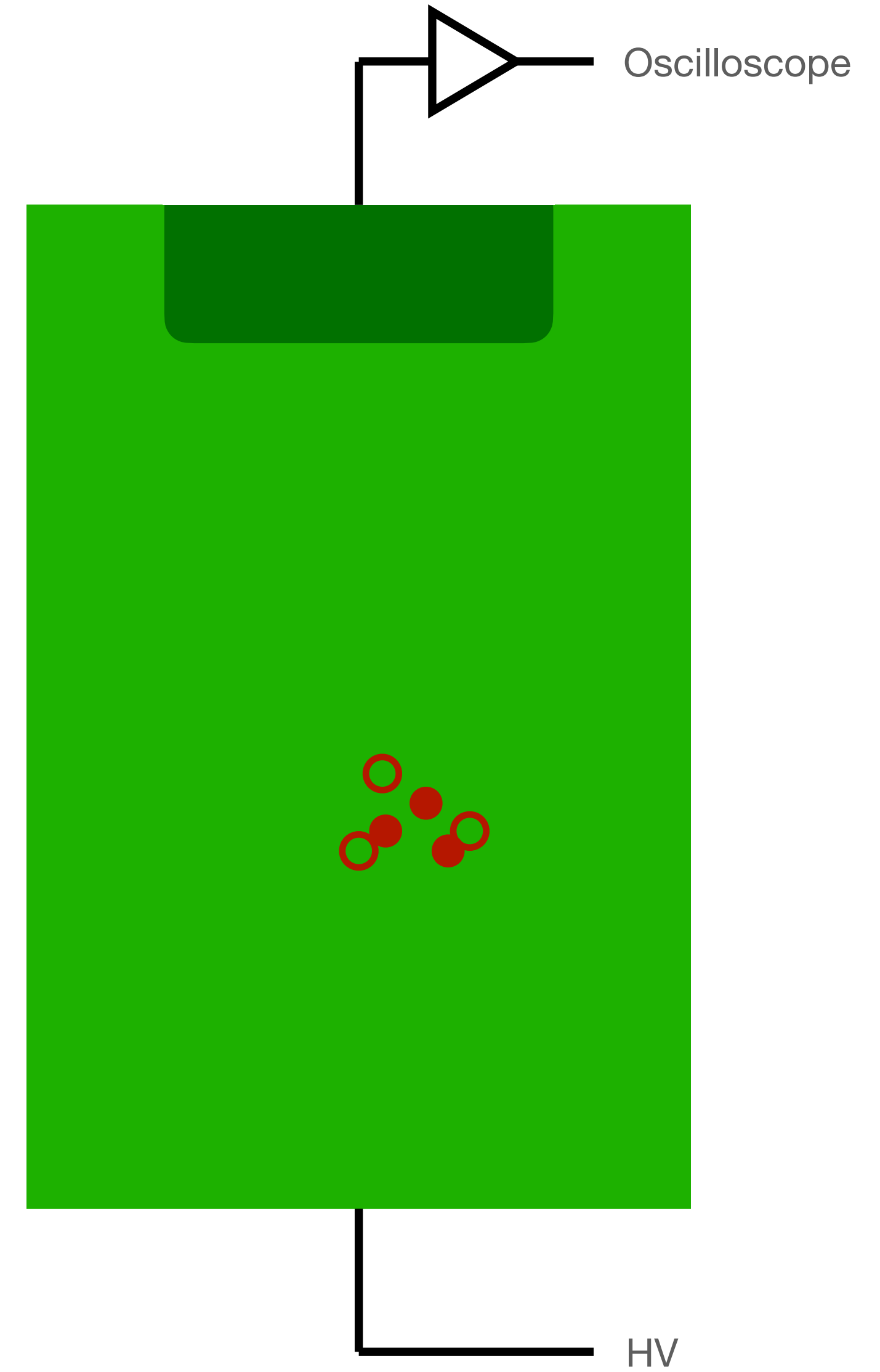
- Consider a typical electric field, with a sensor biased to around depletion



Induced current pulses

Transient Current Technique

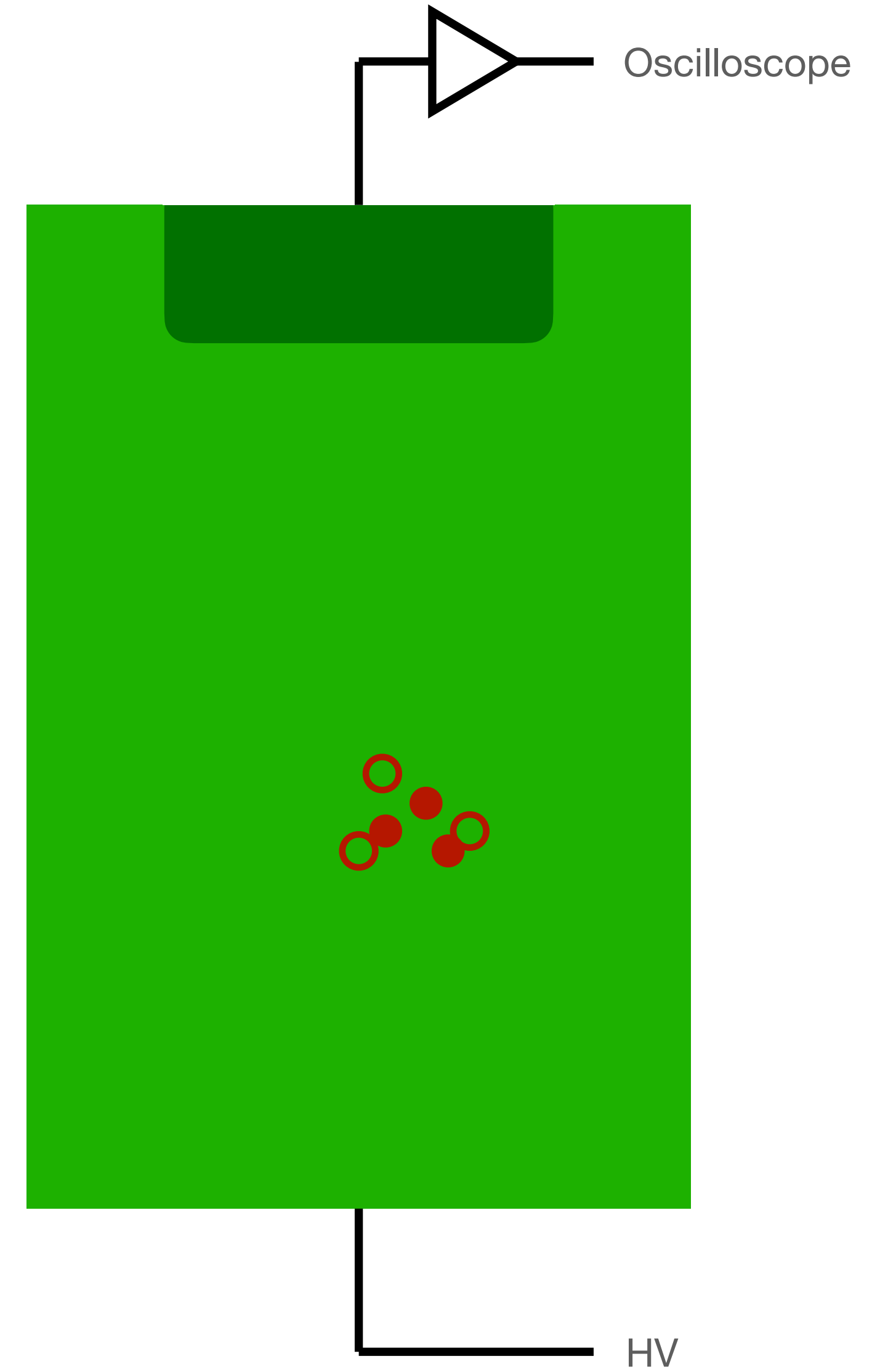
- Look at the induced signal as charge carriers are created at a single spatial point within the sensor



Induced current pulses

Transient Current Technique

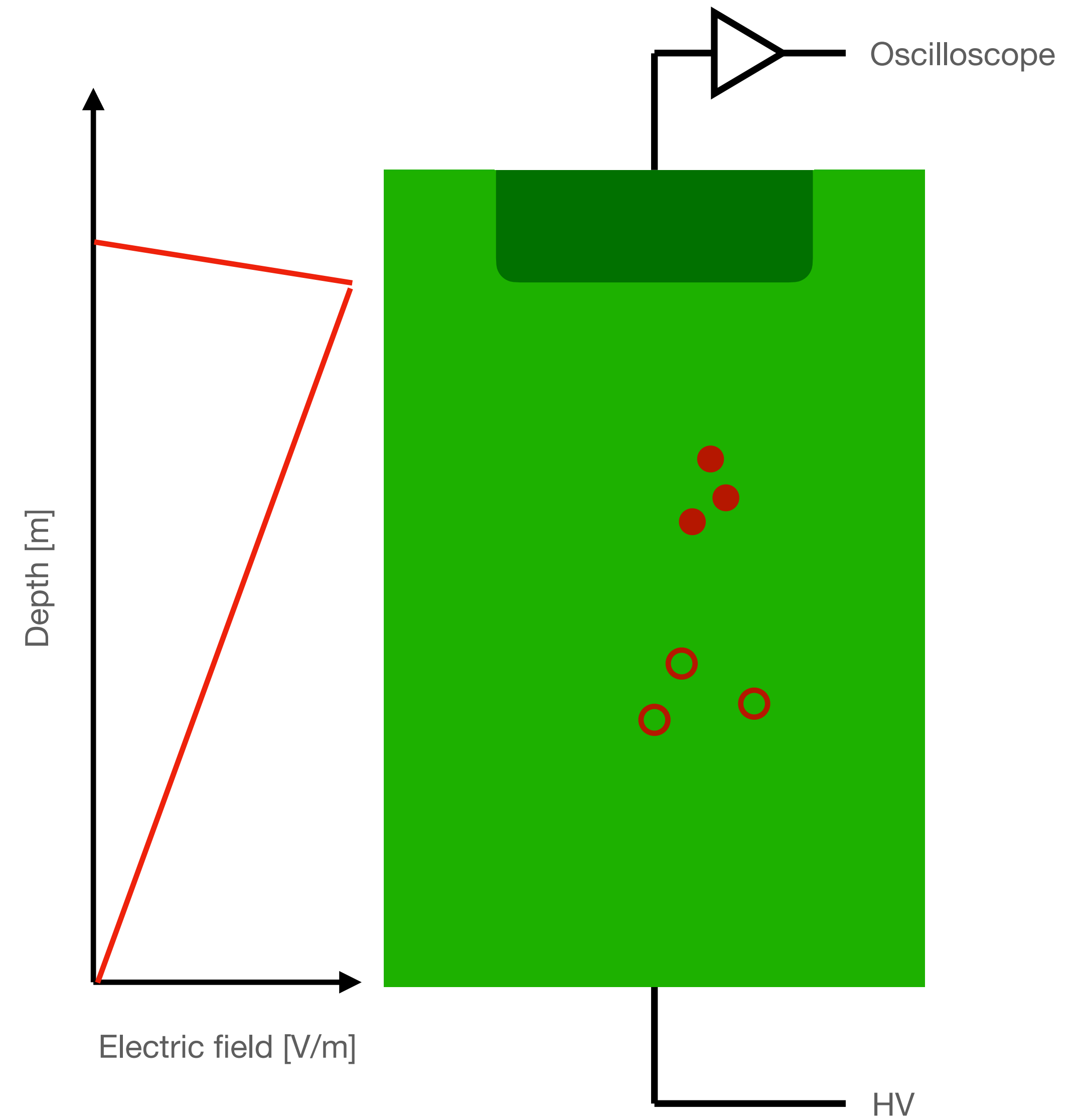
- In most approximations, the acceleration time of the charge carriers is ignored and they are considered to effectively hit their drift velocity instantaneously
- Electron and hole velocities are different, so their current pulses look different (holes move a factor ~ 3 slower)
- Note that the charge carriers are assumed to be independent and do not see each other



Induced current pulses

Transient Current Technique

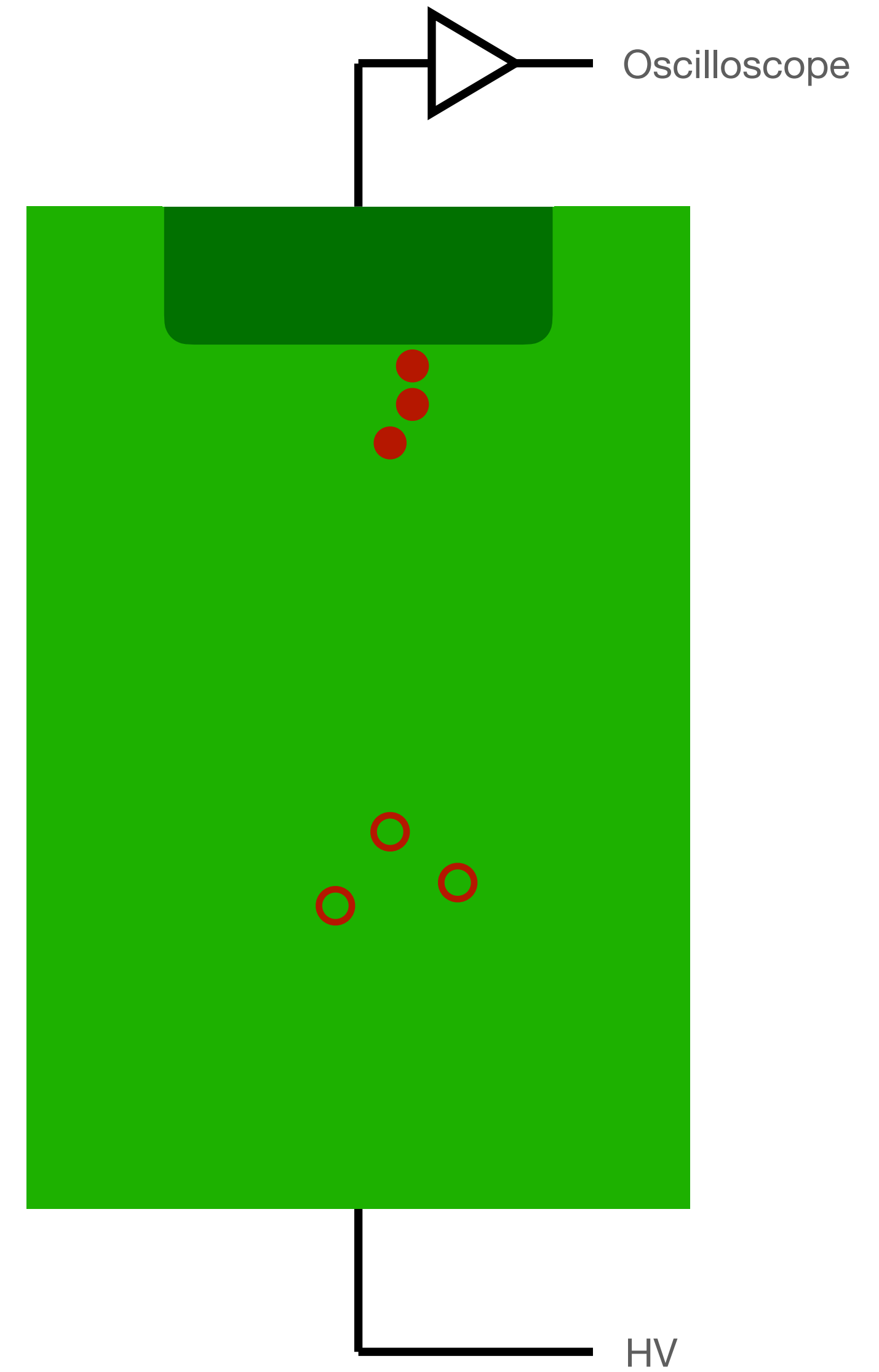
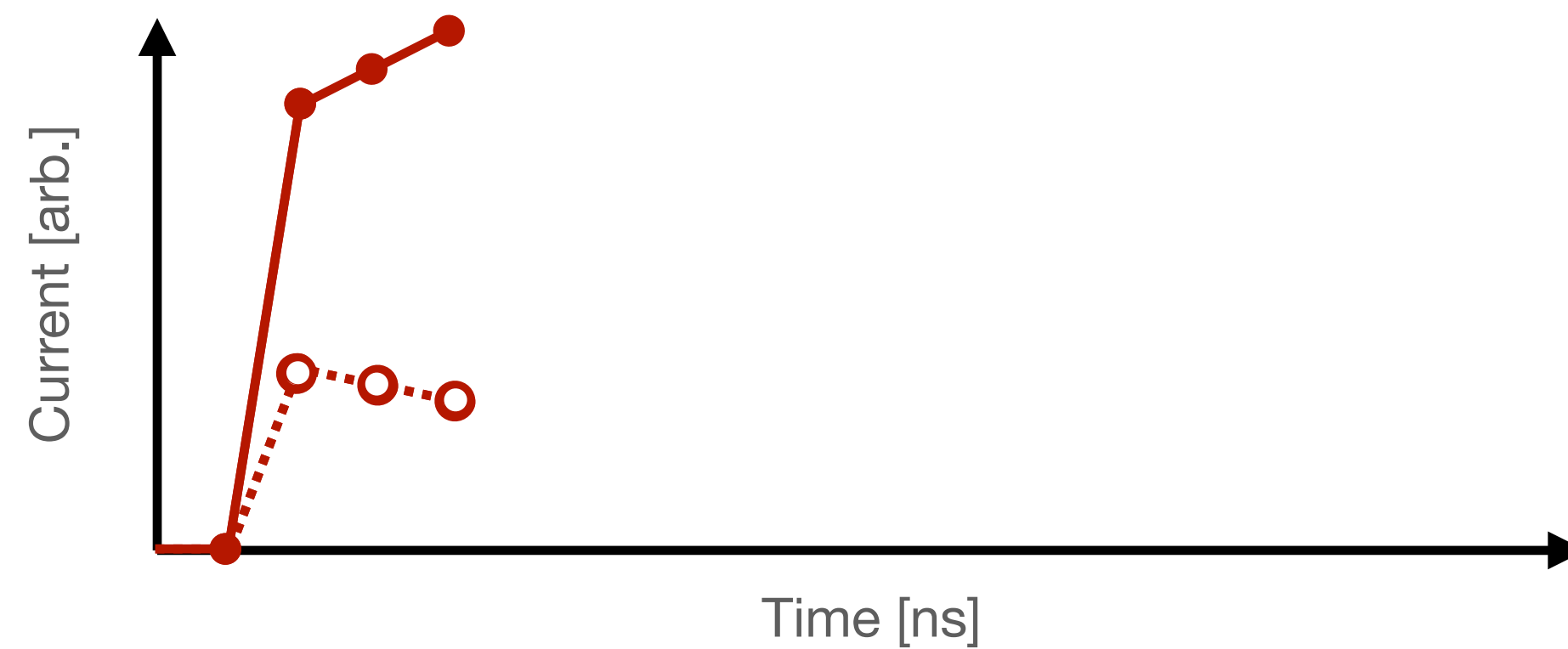
- As the carriers drift, their drift velocities change with the local electric field - in this case the electrons are getting faster and the holes are getting slower



Induced current pulses

Transient Current Technique

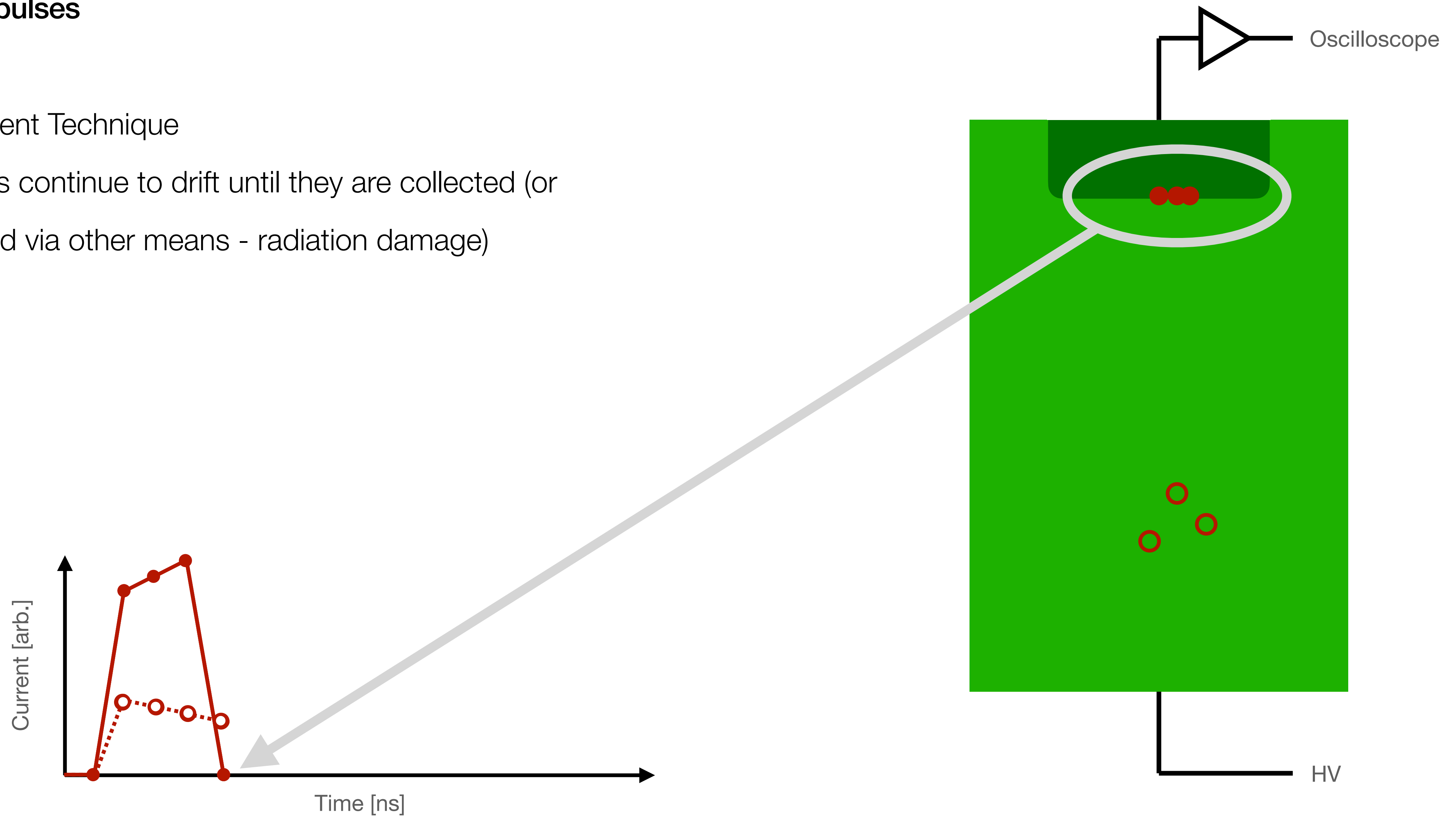
- Carriers continue to drift until they are collected (or stopped via other means - radiation damage)



Induced current pulses

Transient Current Technique

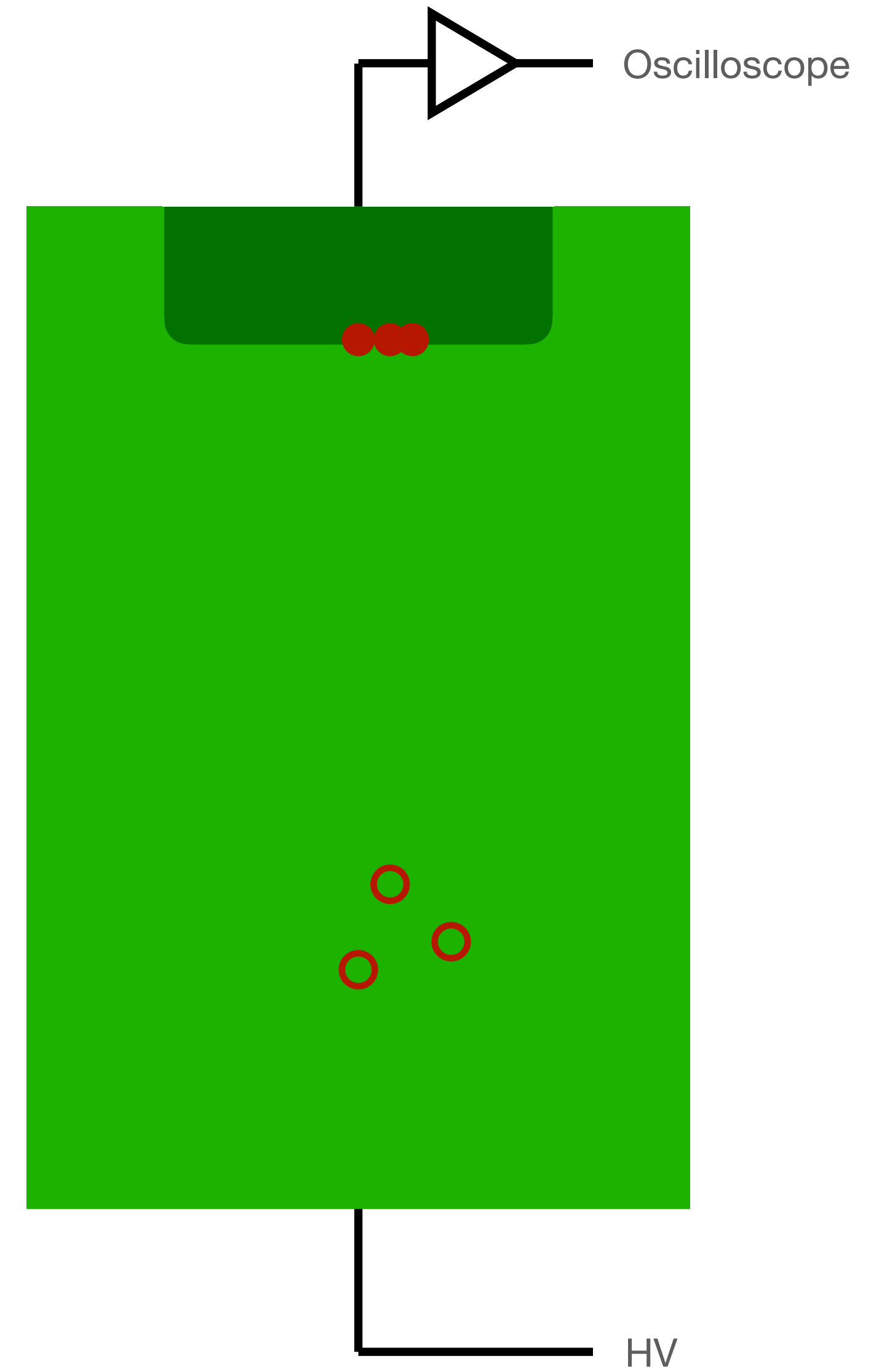
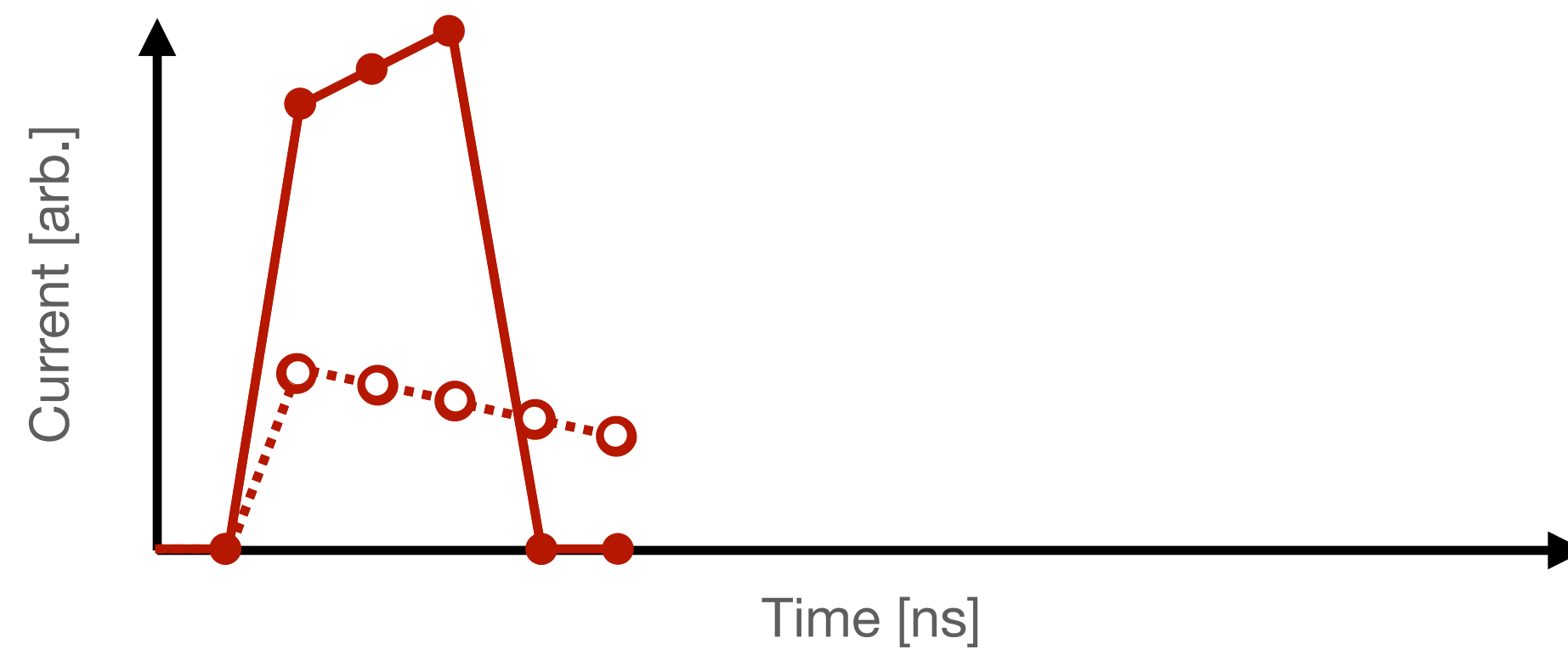
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Induced current pulses

Transient Current Technique

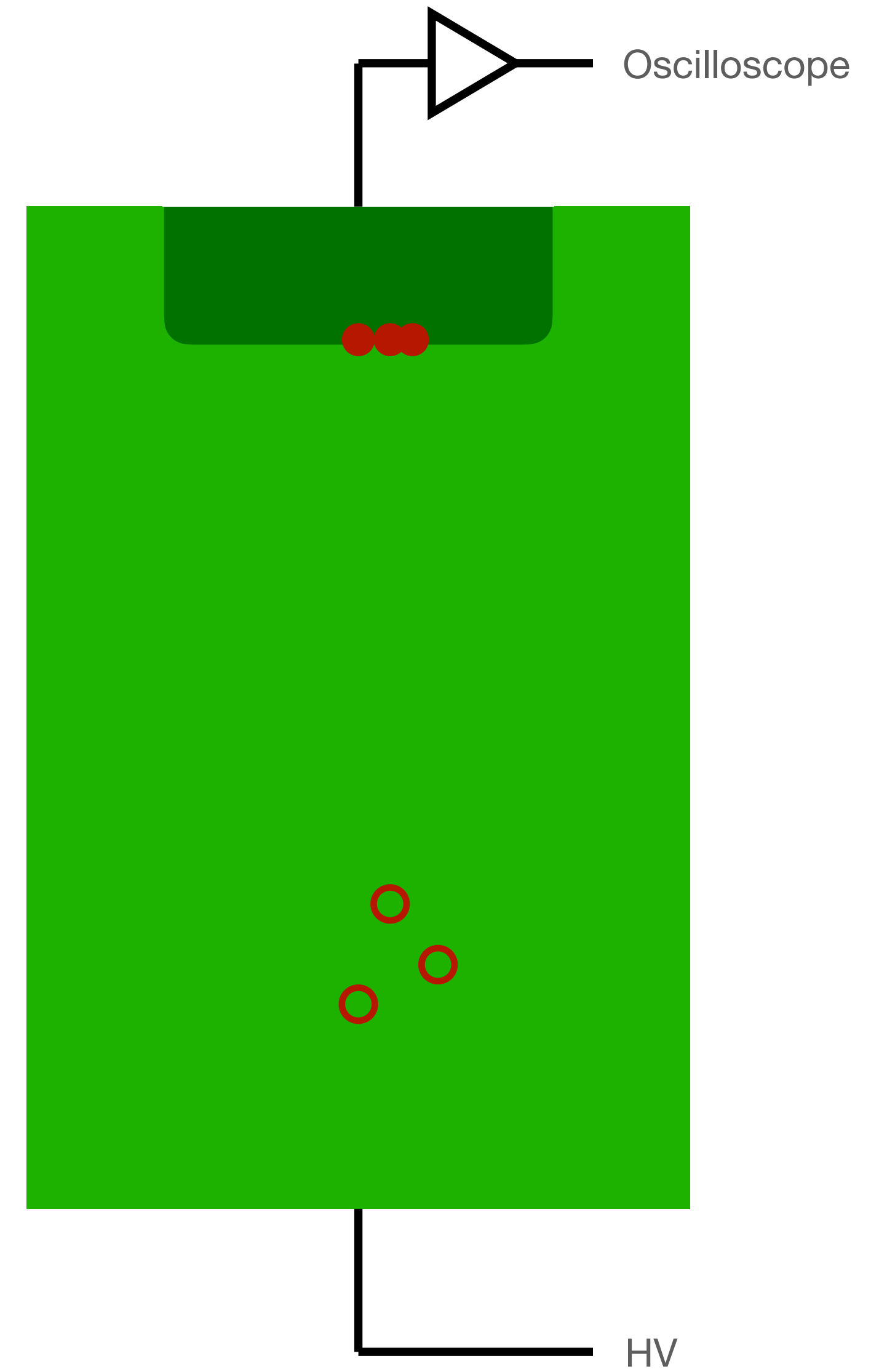
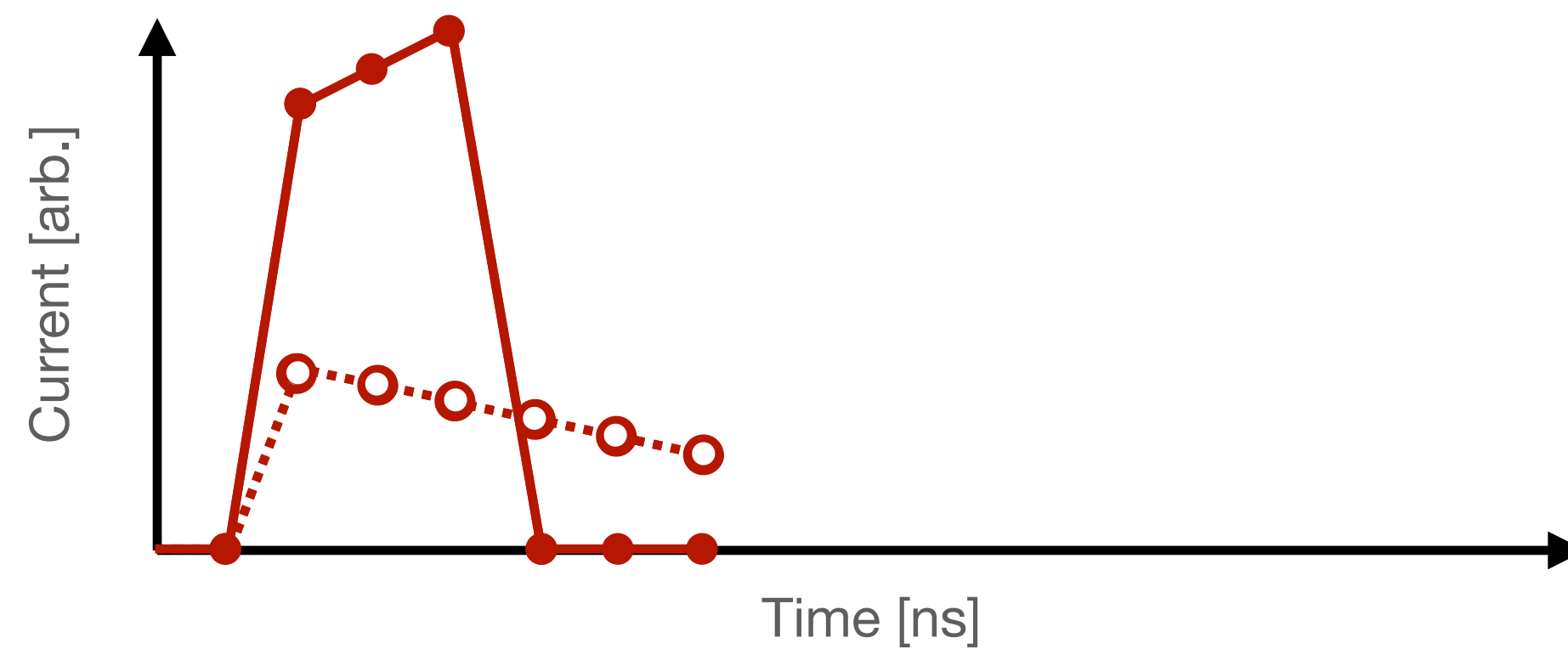
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Transient Current Technique

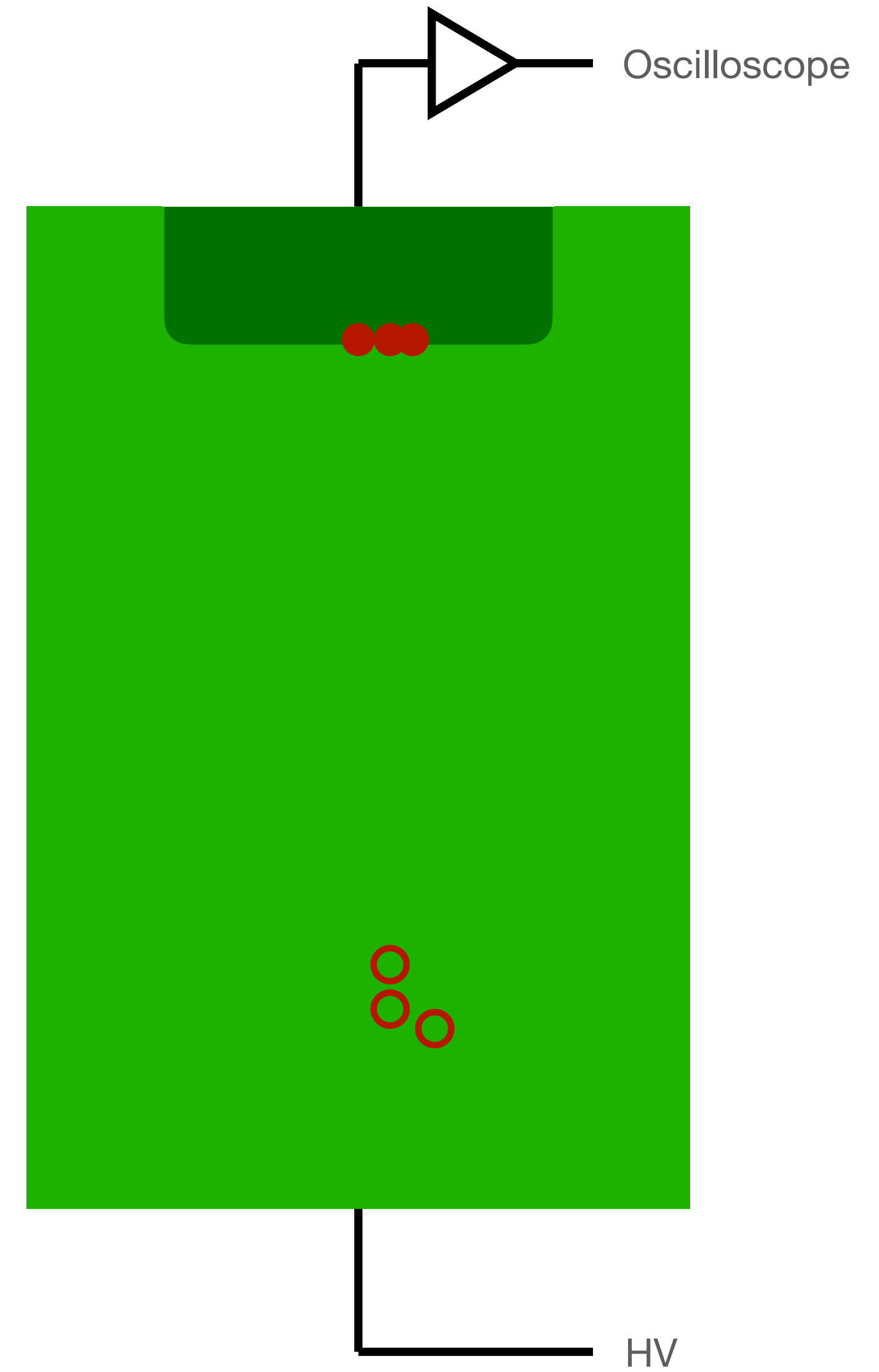
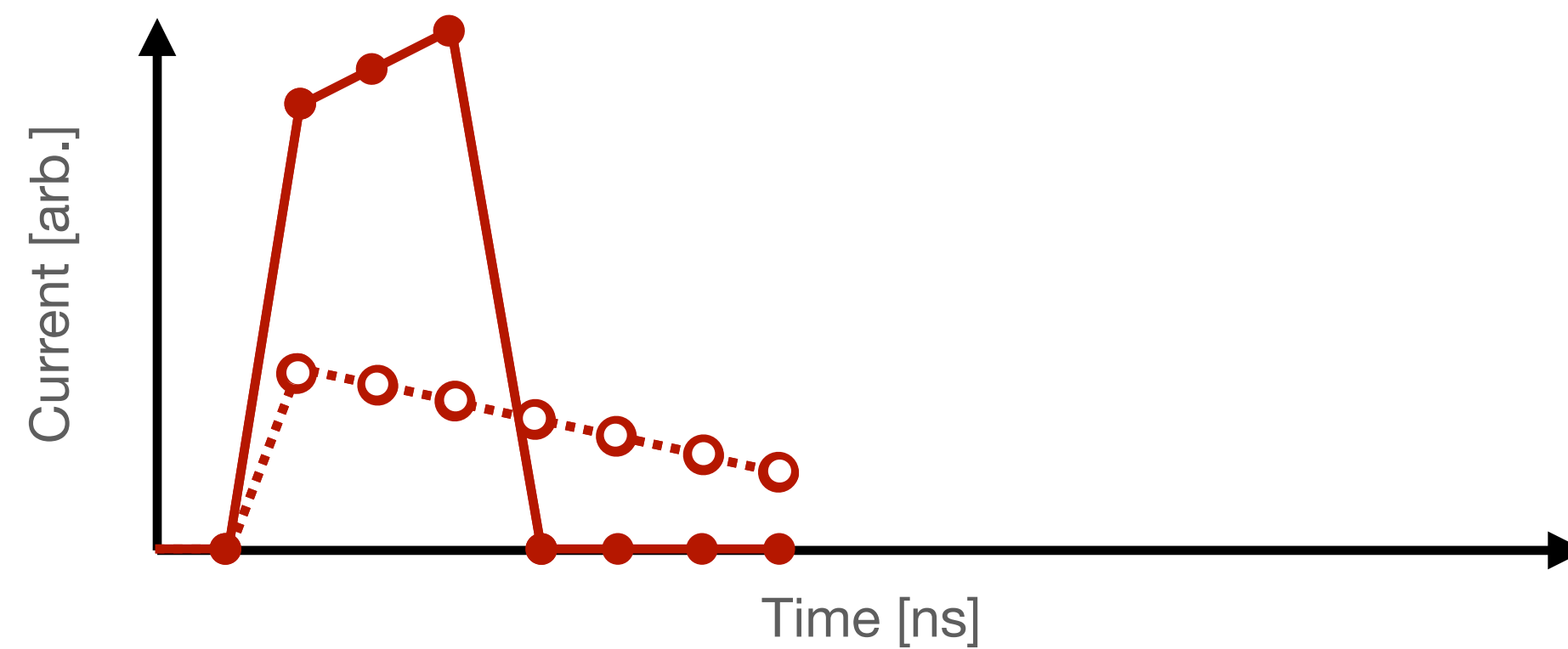
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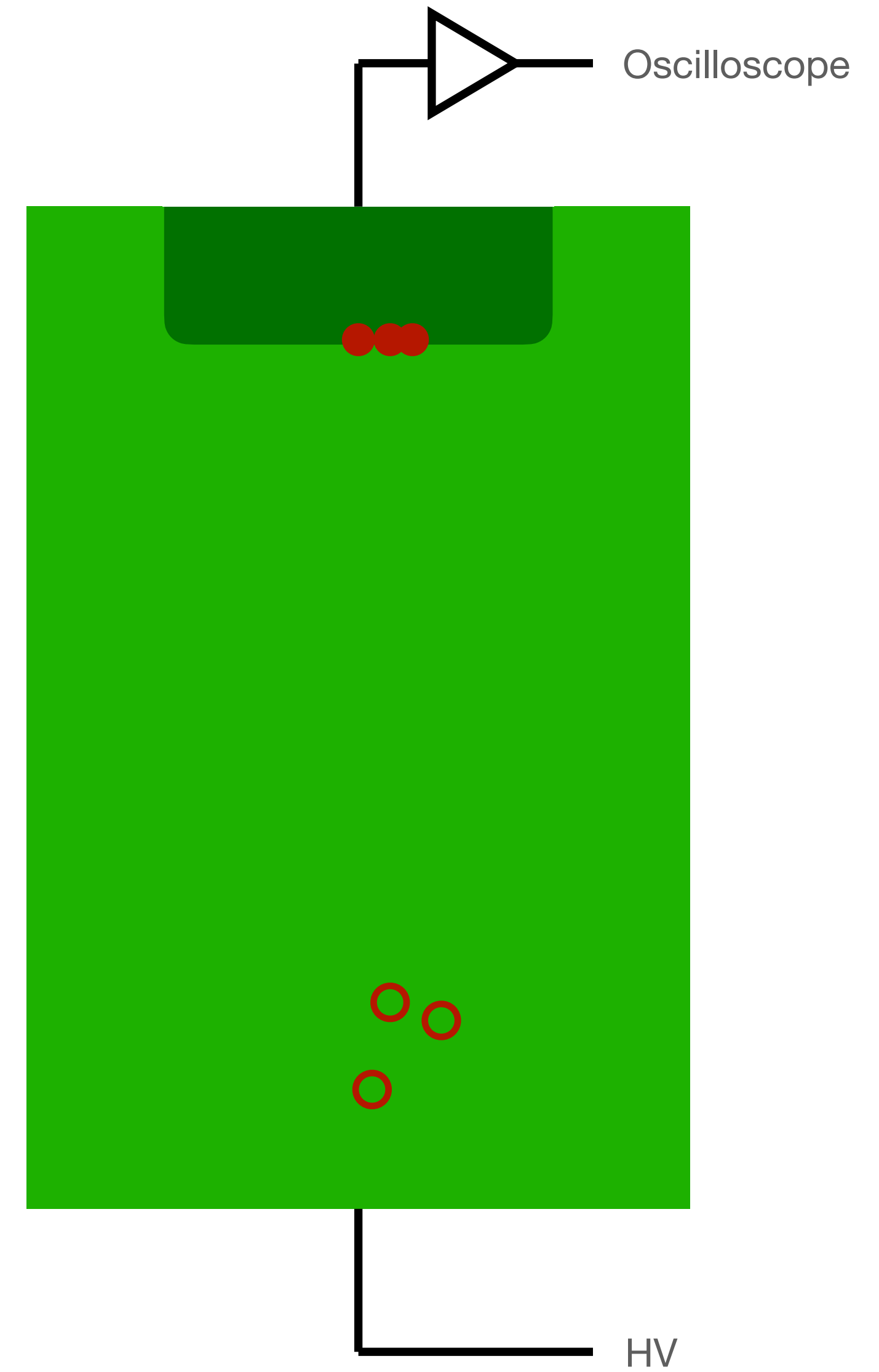
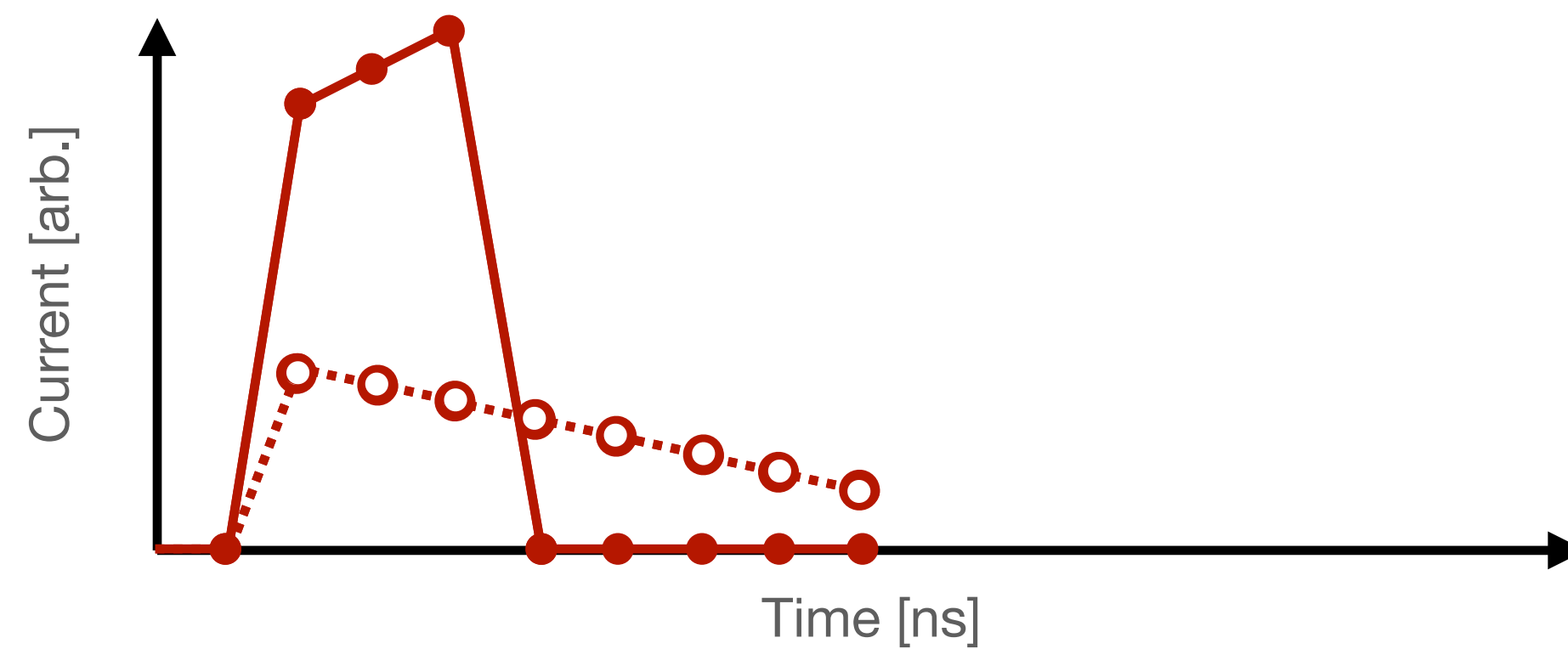
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Induced current pulses

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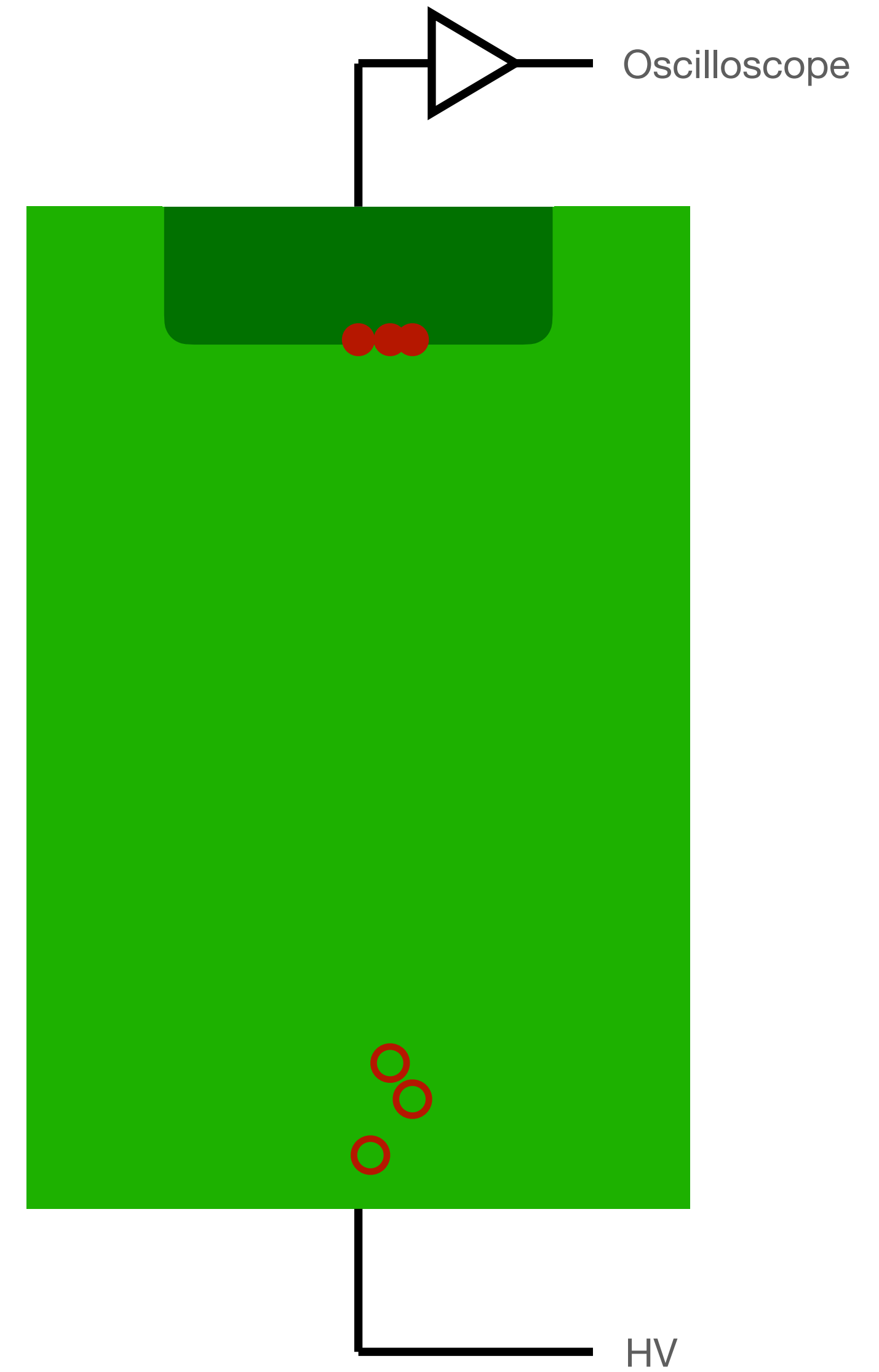
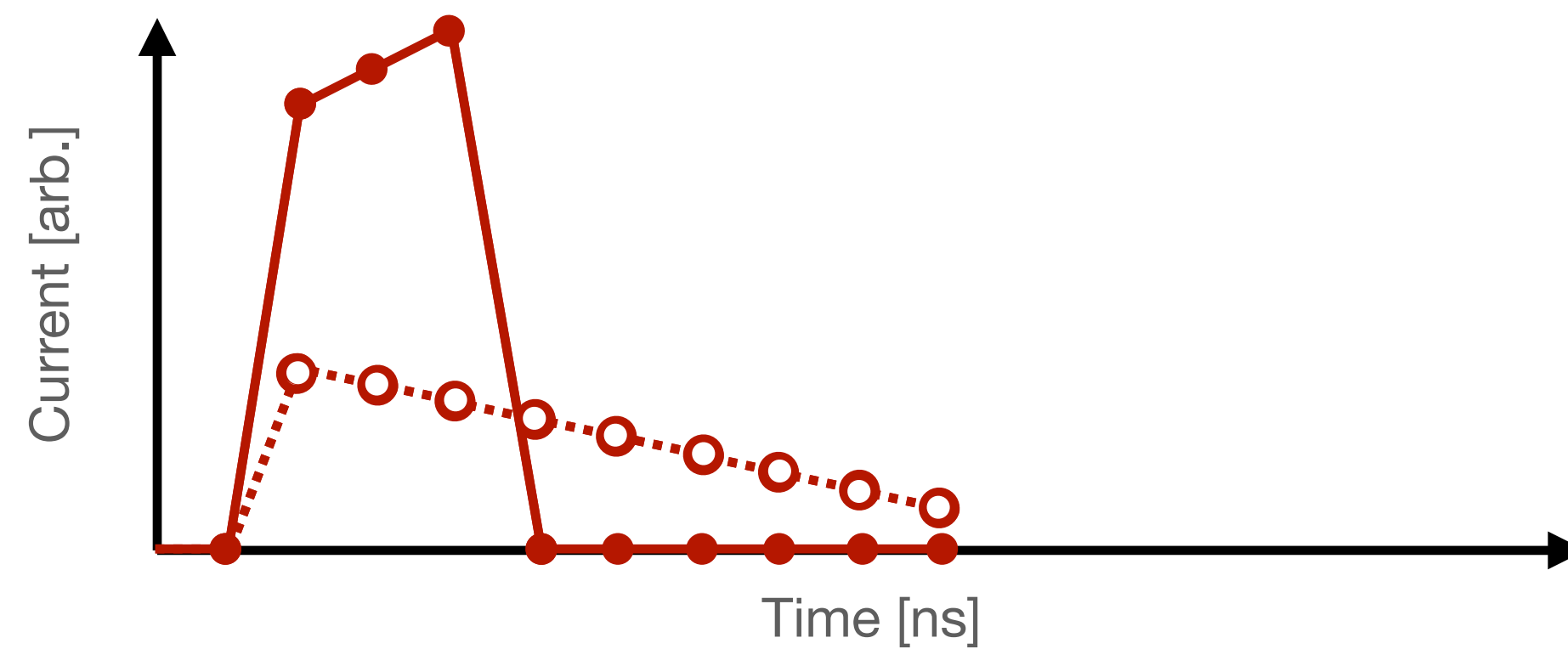
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Induced current pulses

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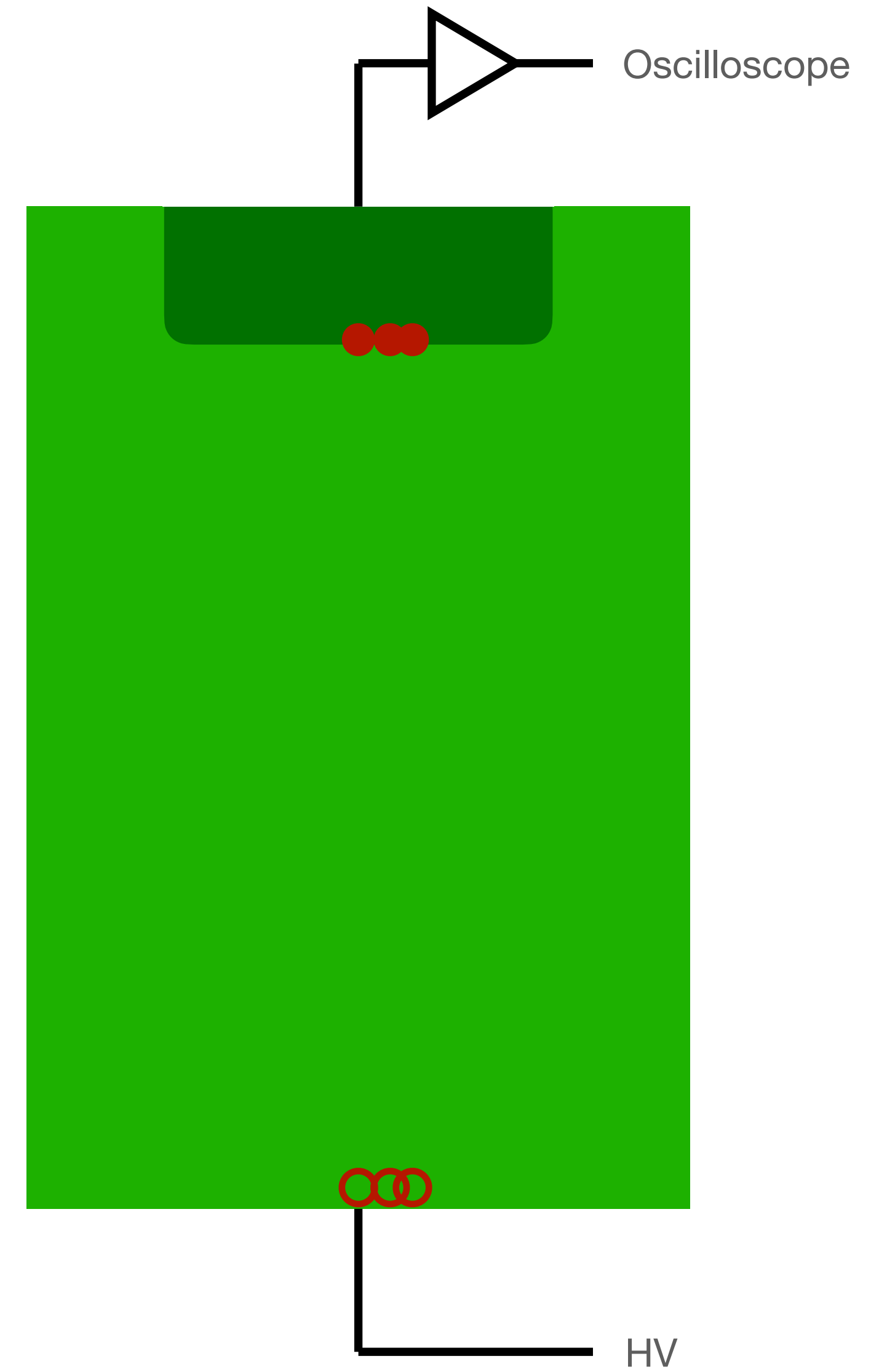
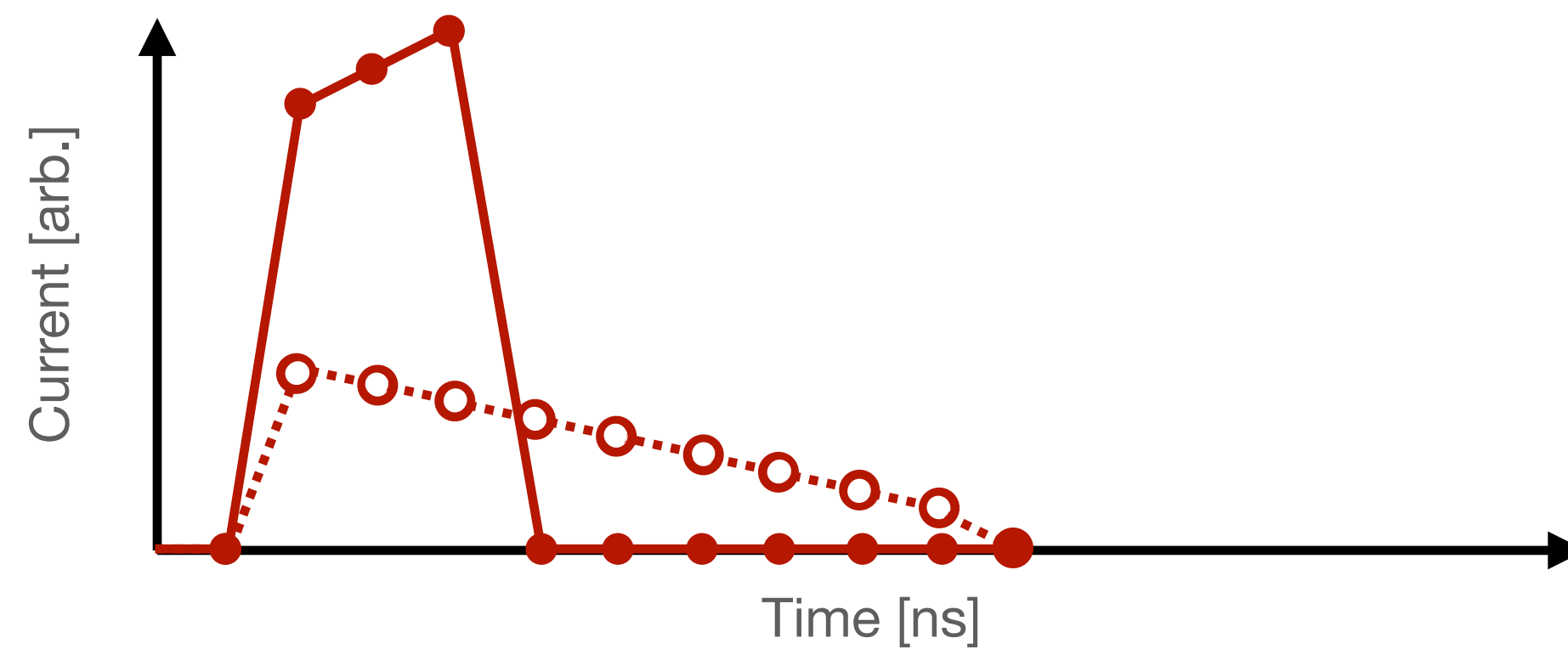
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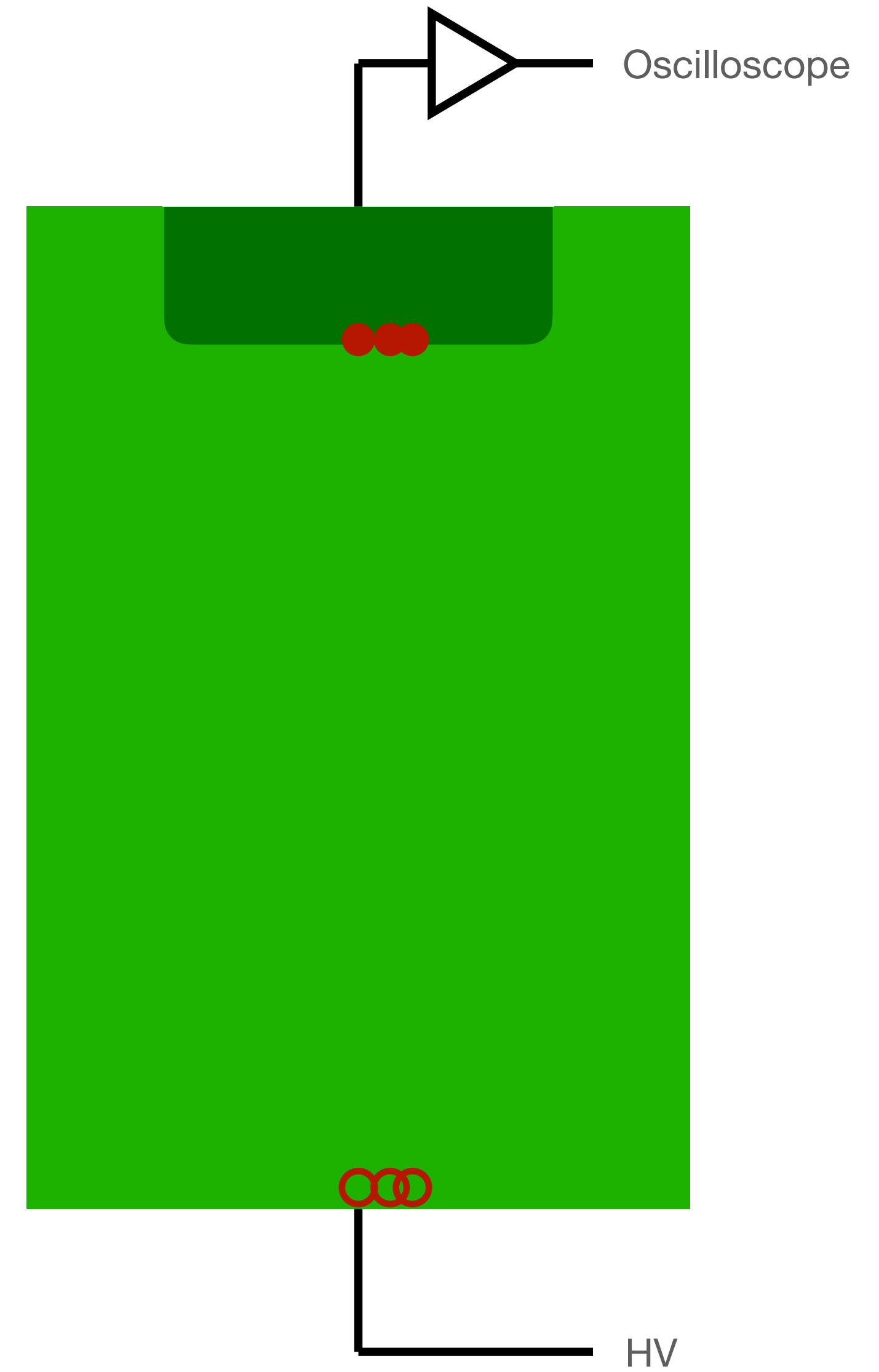
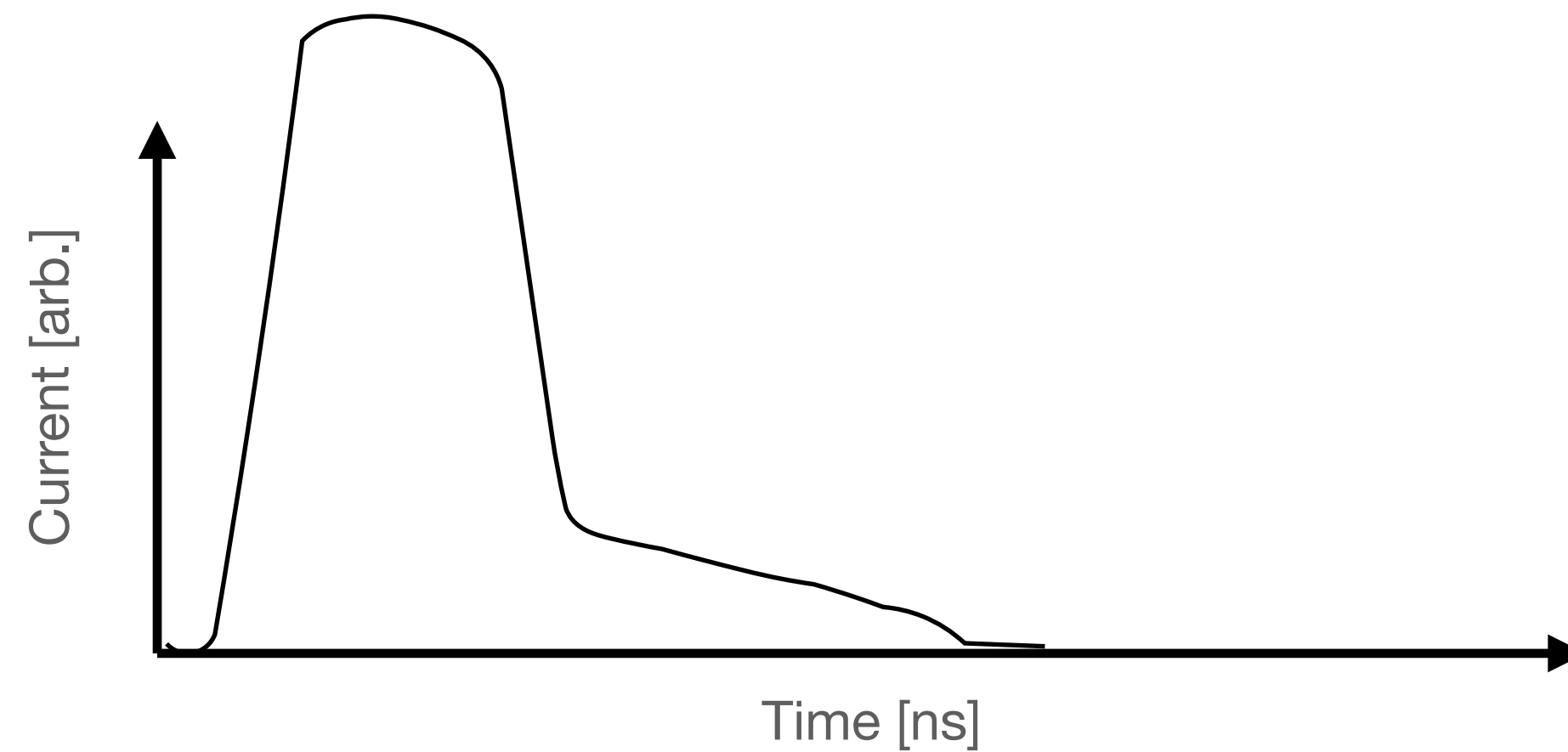
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Induced current pulses

Transient Current Technique

- Finally we have our total current pulse from a single set of generated charge carriers



**Making tiny moving parts inside flat rocks that give us light
and power**

(Charge carrier generation)

Generating charge carriers

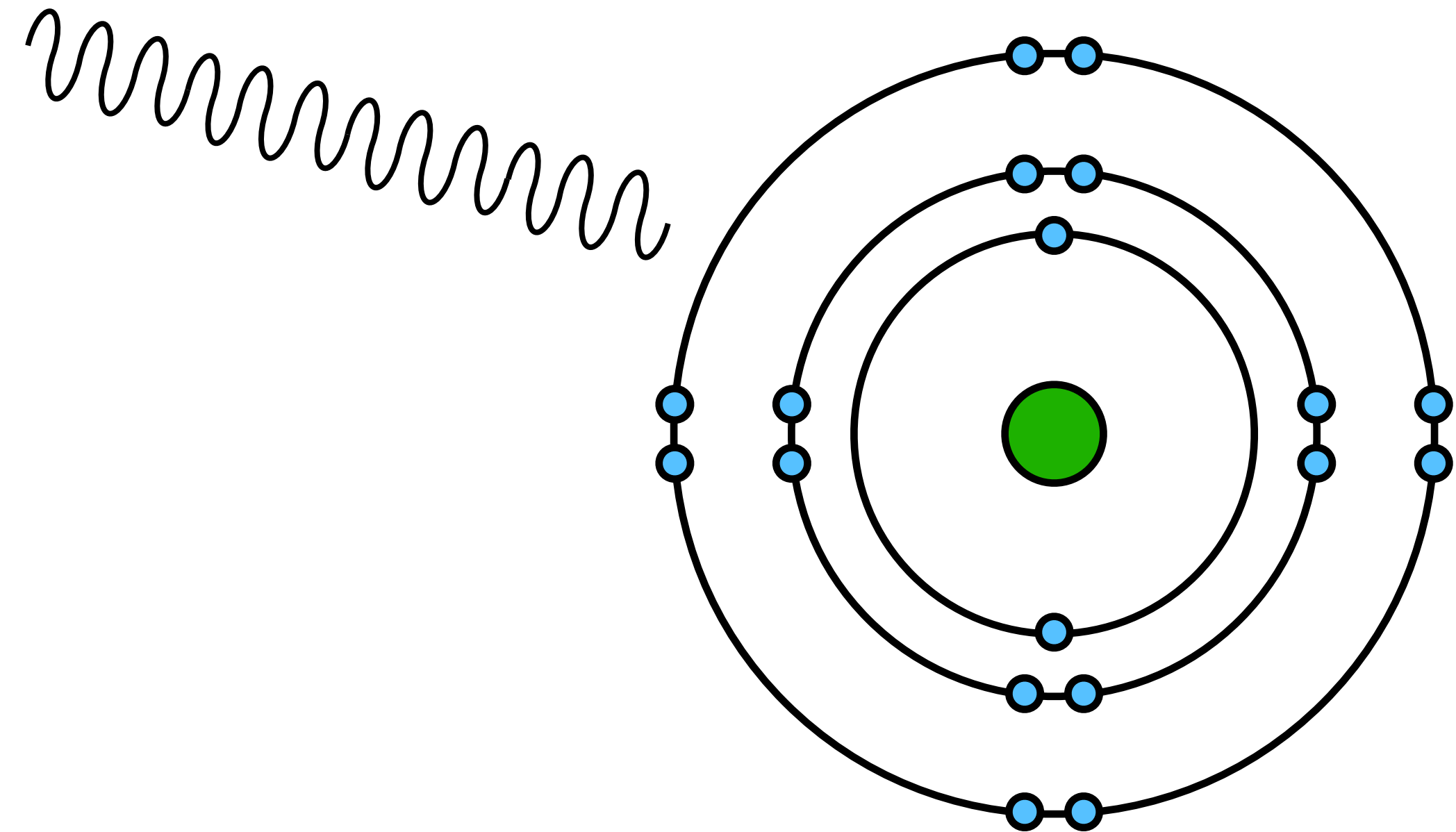
Different particles will interact different with the electrons in our silicon

- Photons have no mass
- Electrons are identical in mass
- Other charged particles are substantially heavier

Photon interactions

Photon interactions depend heavily on the energy of the photon and the density of electrons in the medium, but the following options are available:

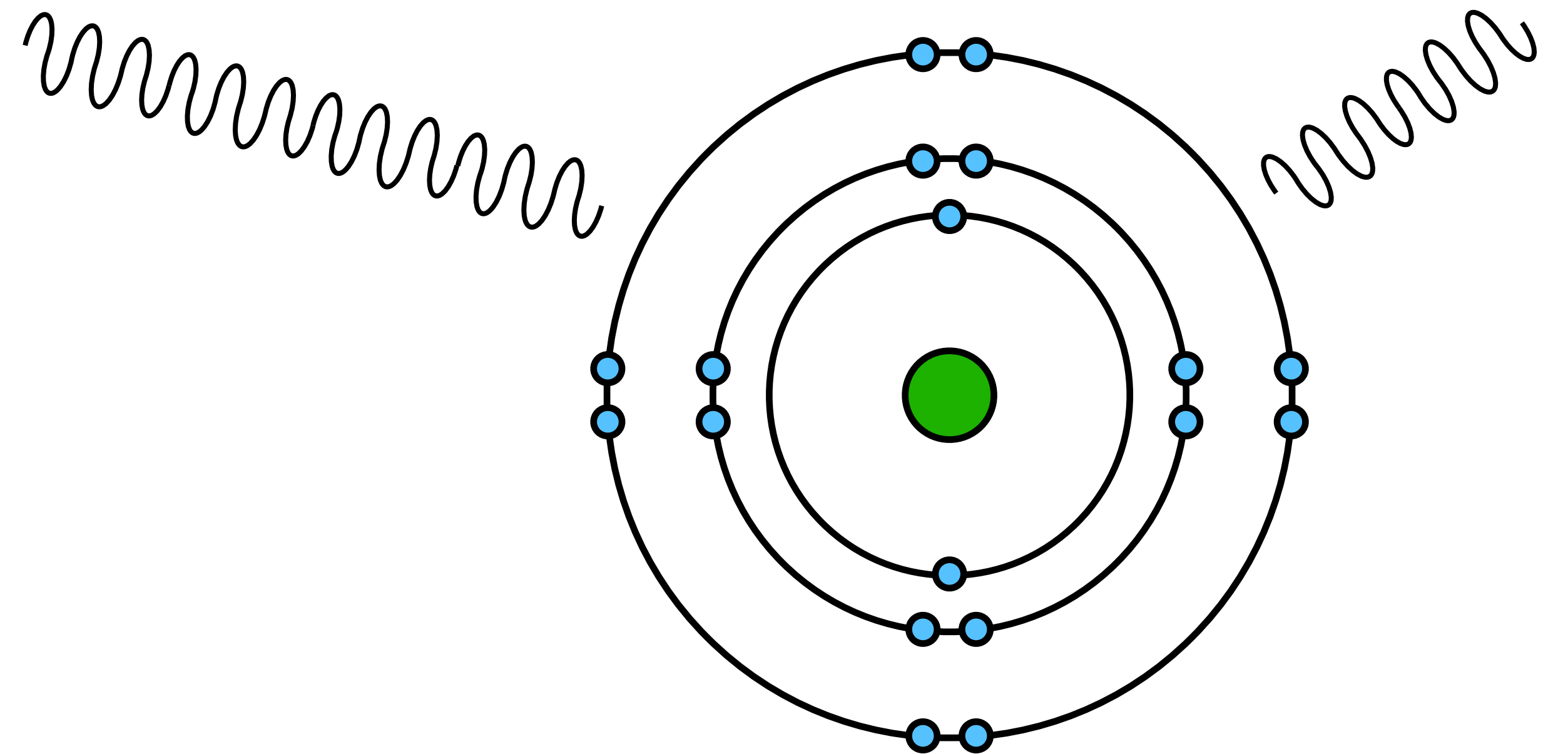
- Elastic scattering (Rayleigh)
- Absorption (photoelectric effect)
- Inelastic scattering (Compton)
- Conversion into matter (pair production)



Photon interactions - Rayleigh scattering

Rayleigh scattering is elastic scattering - so there is no transfer of energy to the medium

- This doesn't help us much to detect anything...

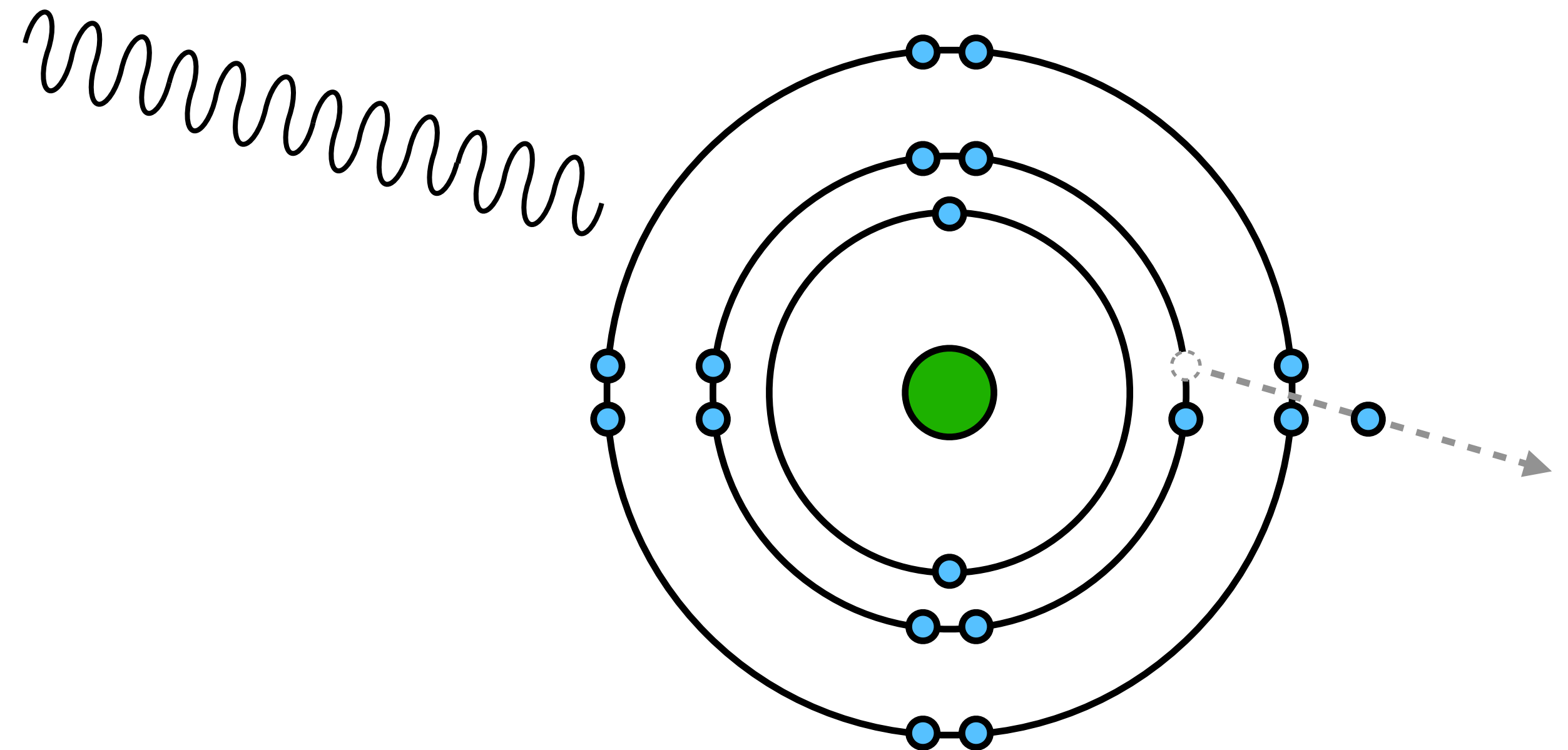


Photon interactions - photoelectric effect

In the photoelectric effect the photon is fully absorbed by an atom, which ejects an electron with energy $E_{\text{photon}} - E_{\text{binding}}$

- The process involves the whole atom to facilitate momentum transfer
- For this reason, it occurs mainly for inner electrons (K shell)
- The cross section depends strongly on atomic number Z , and drops off rapidly with photon energy:

$$\sigma \propto \frac{Z^4}{E^3}$$



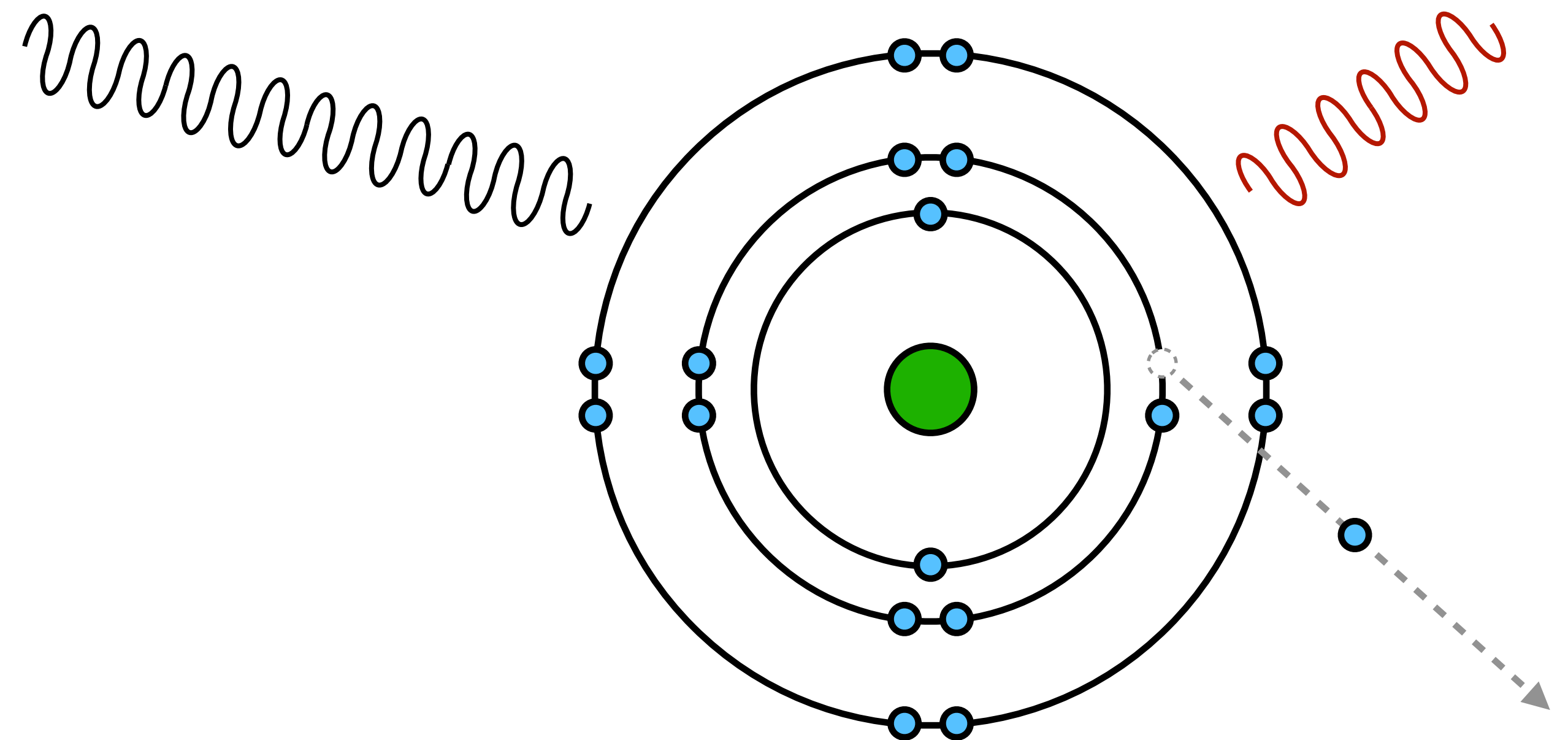
Photon interactions - Compton scattering

Compton scattering involves the *inelastic* scattering of photons from individual electrons

- Some of the photon energy is transferred to the electron

$$\frac{1}{E_f} - \frac{1}{E_i} = \frac{1}{m_e c^2} (1 - \cos\theta)$$

- A lower energy photon will be scattered through angle θ
- The cross section is directly proportional to the electron density

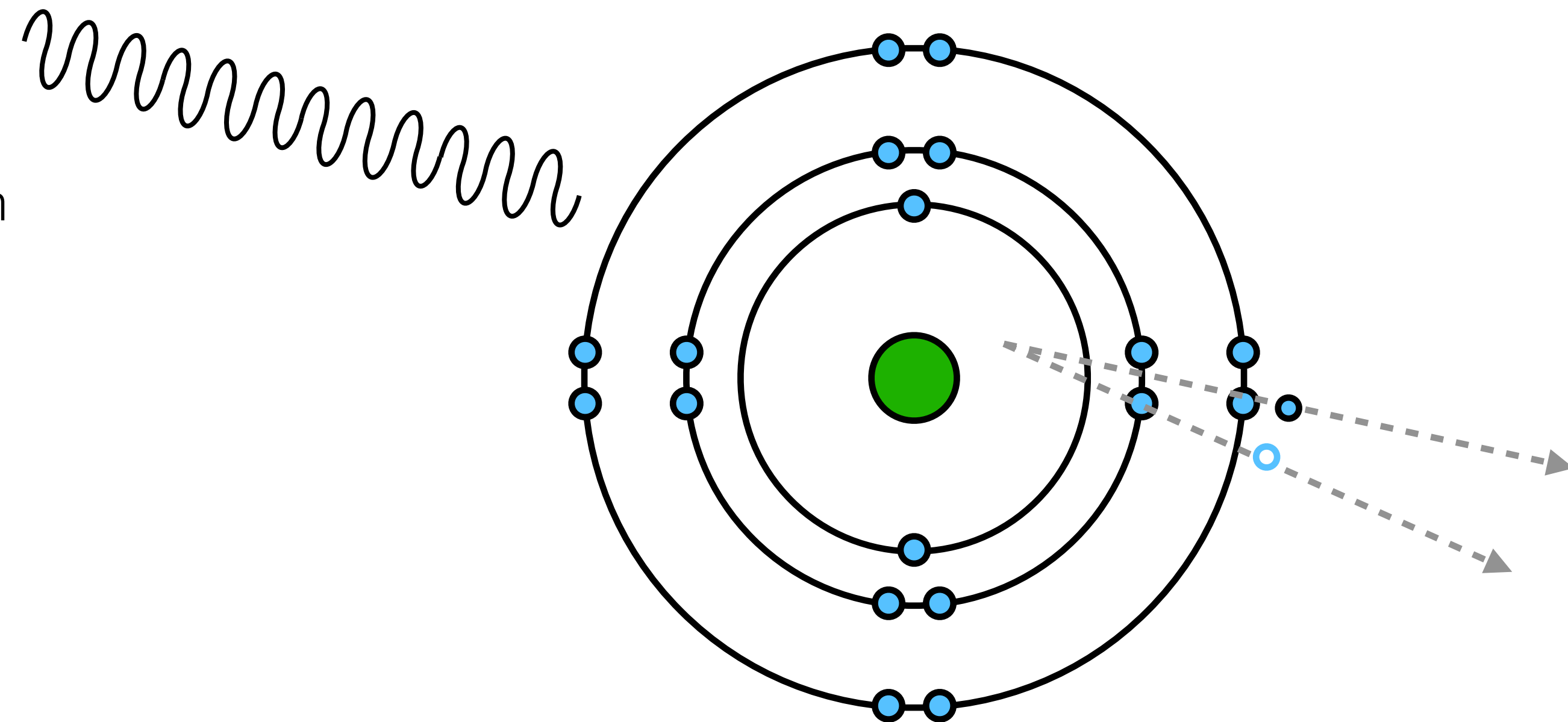


Photon interactions - pair production

Once the photon energy exceeds 1.022 MeV (twice the electron rest mass) then it can interact to produce an electron-positron pair

- As in the photoelectric effect, the nucleus is involved in the interaction to conserve momentum
- The corresponding cross-section varies with the atomic number and grows slowly with photon energy

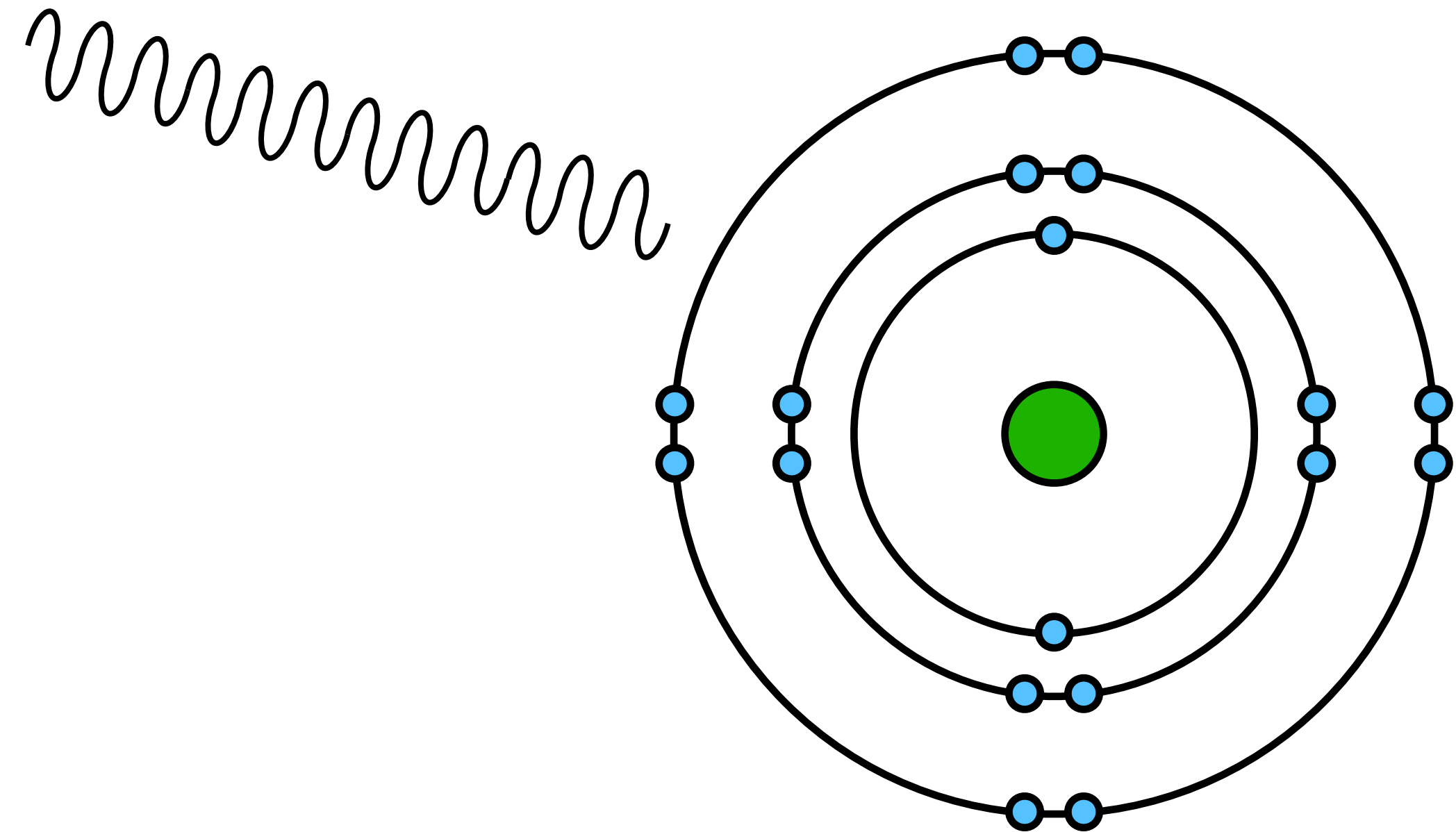
$$\sigma \propto Z^2$$



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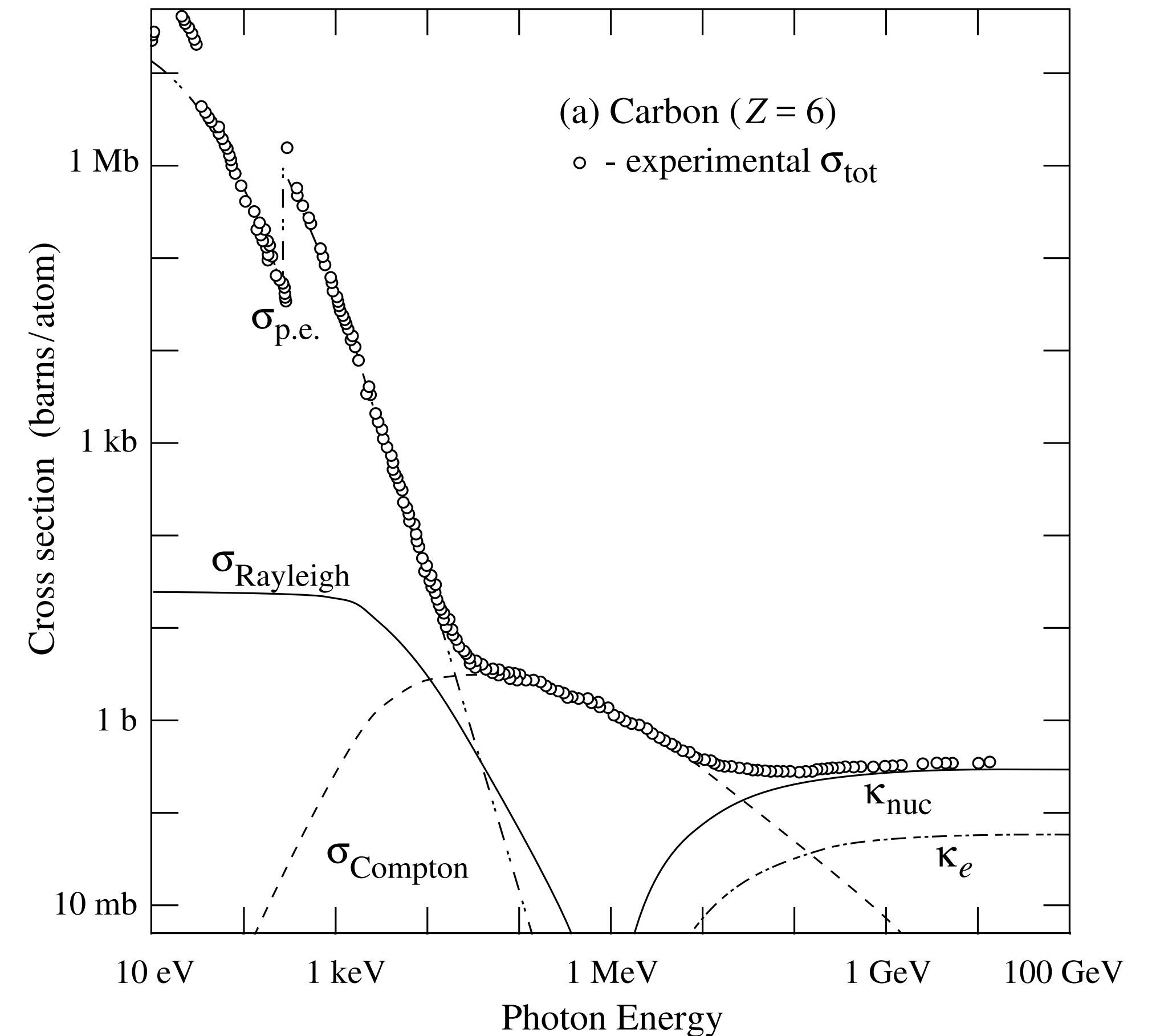
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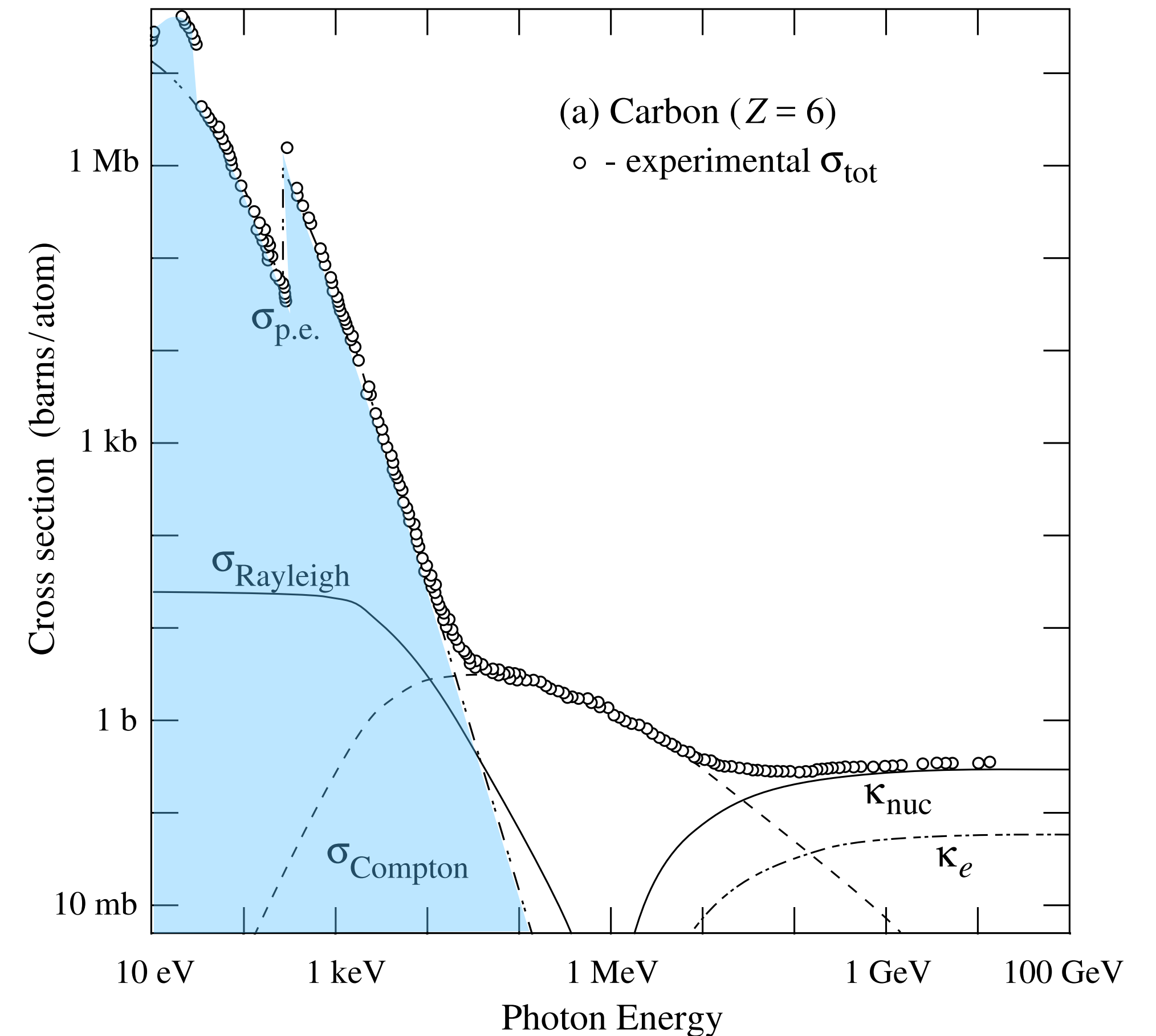
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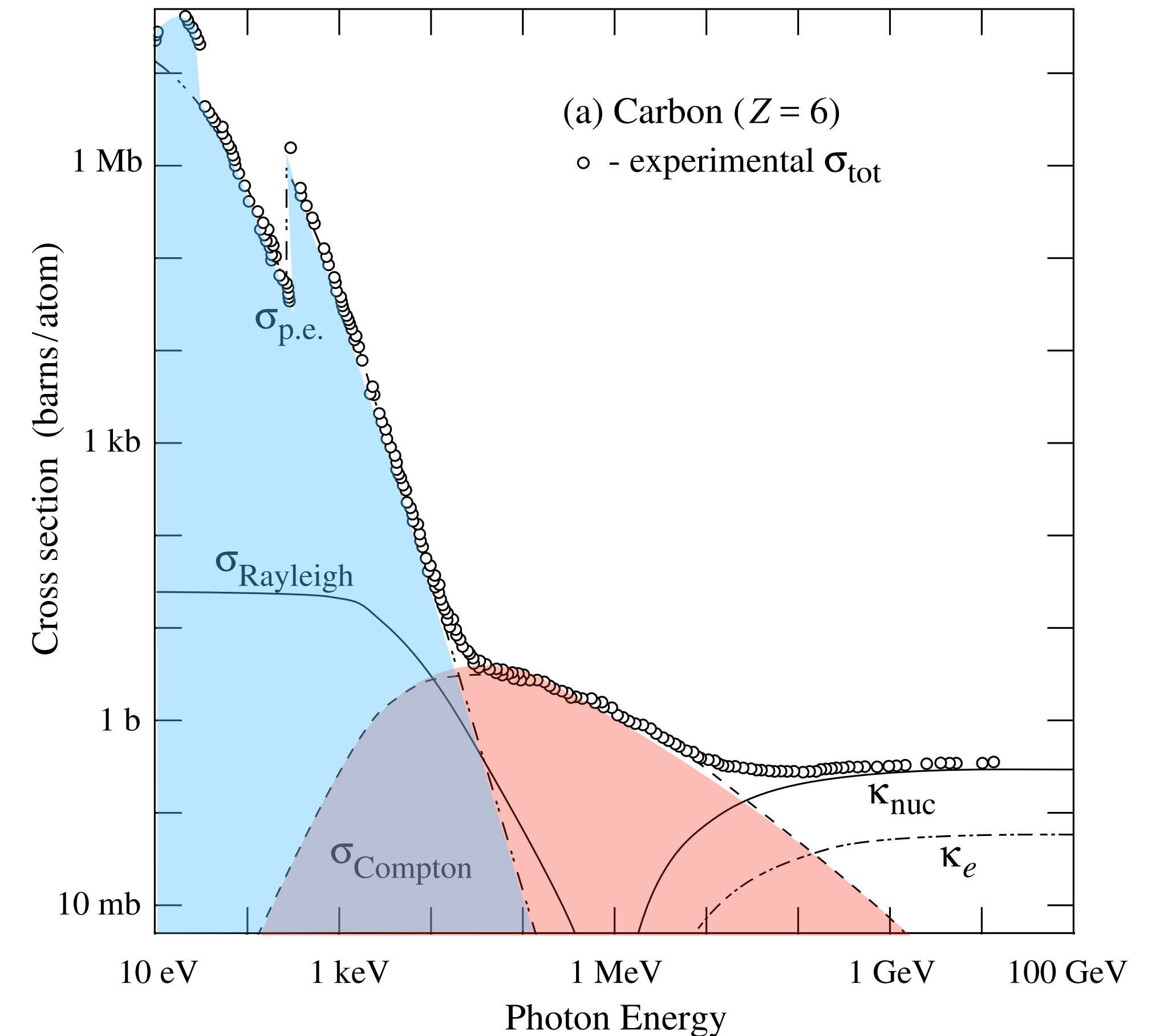
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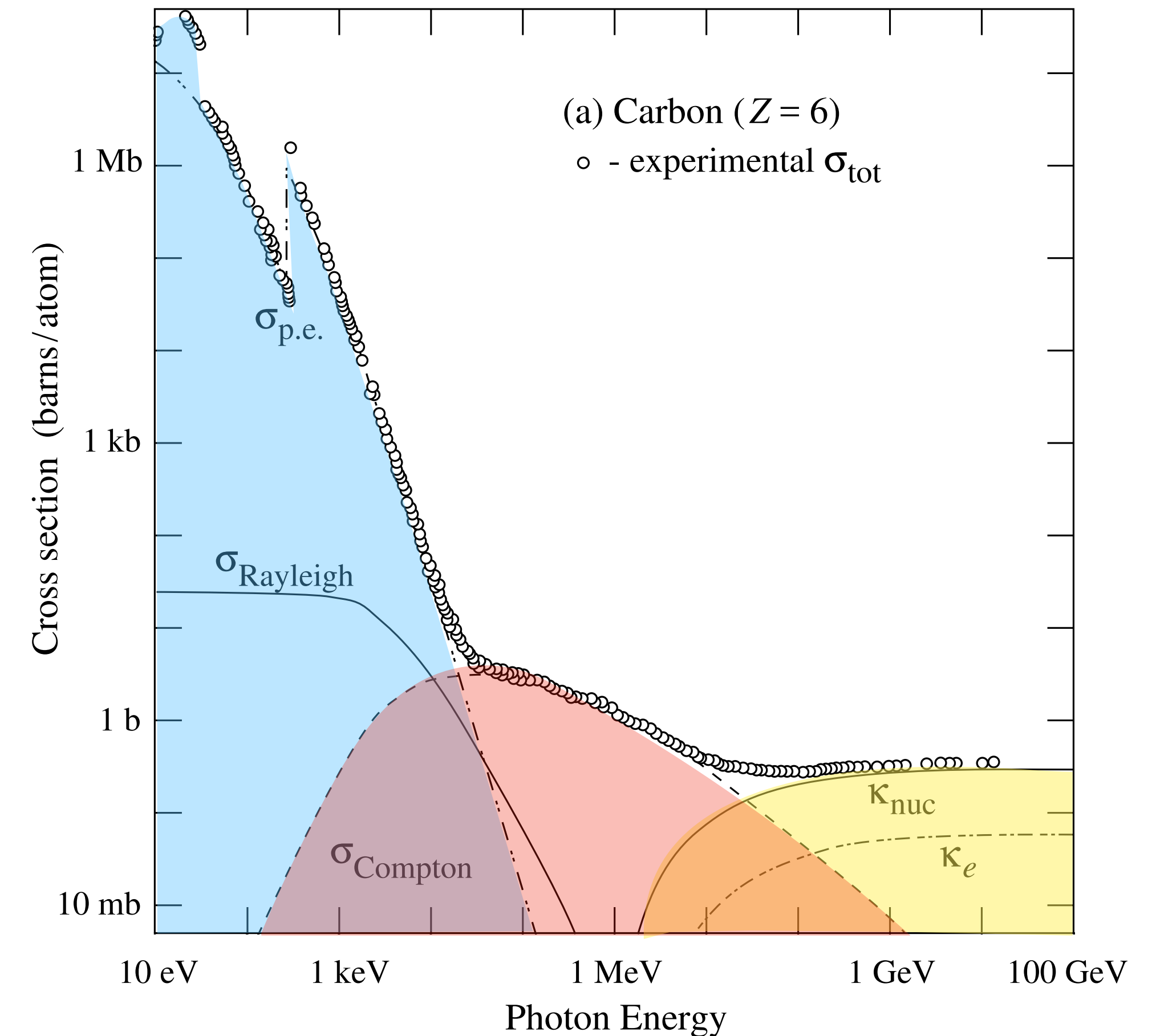
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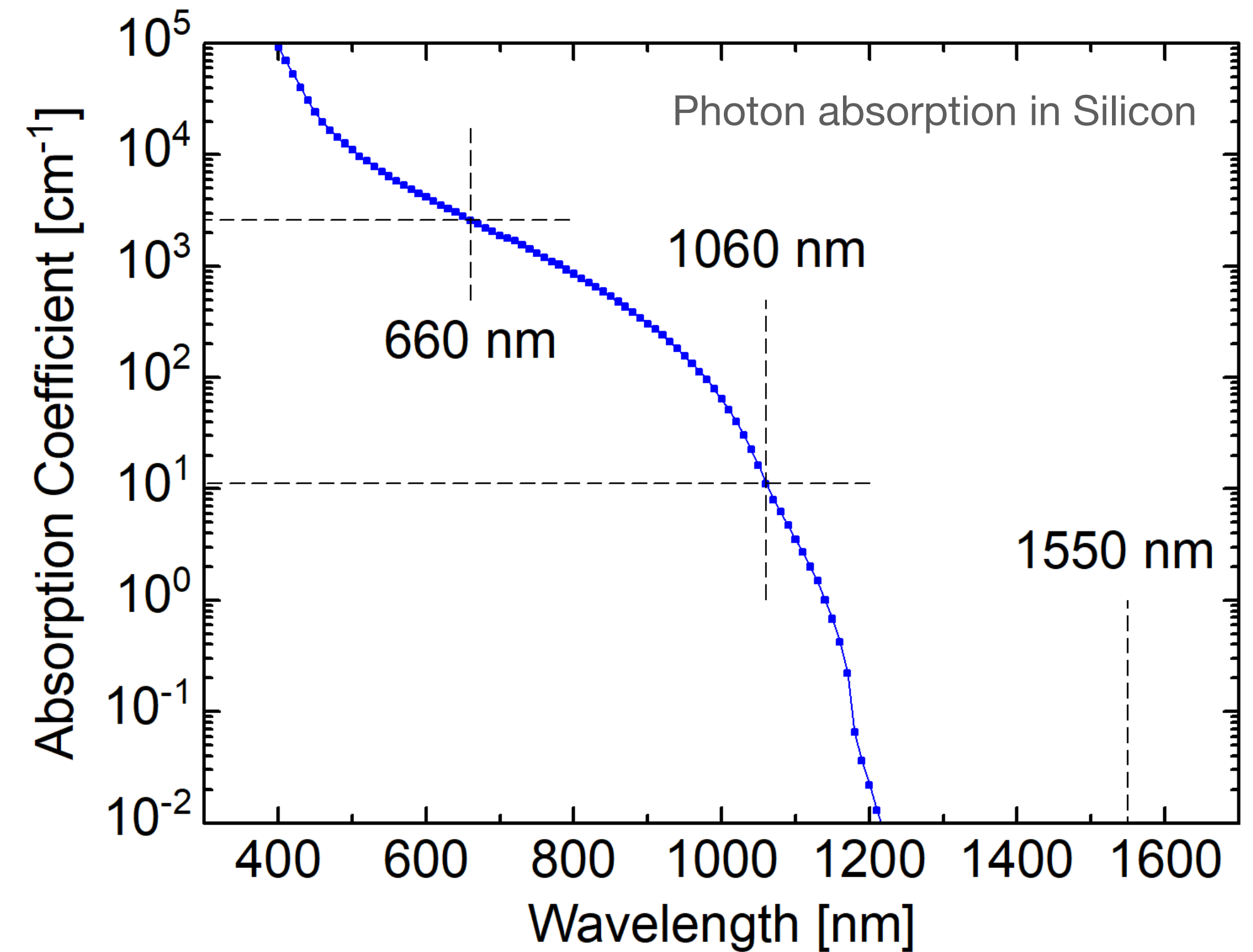
Interaction probability

For the purpose of photon detection, the end-result is that different photon energies interact more or less depending on the detection medium

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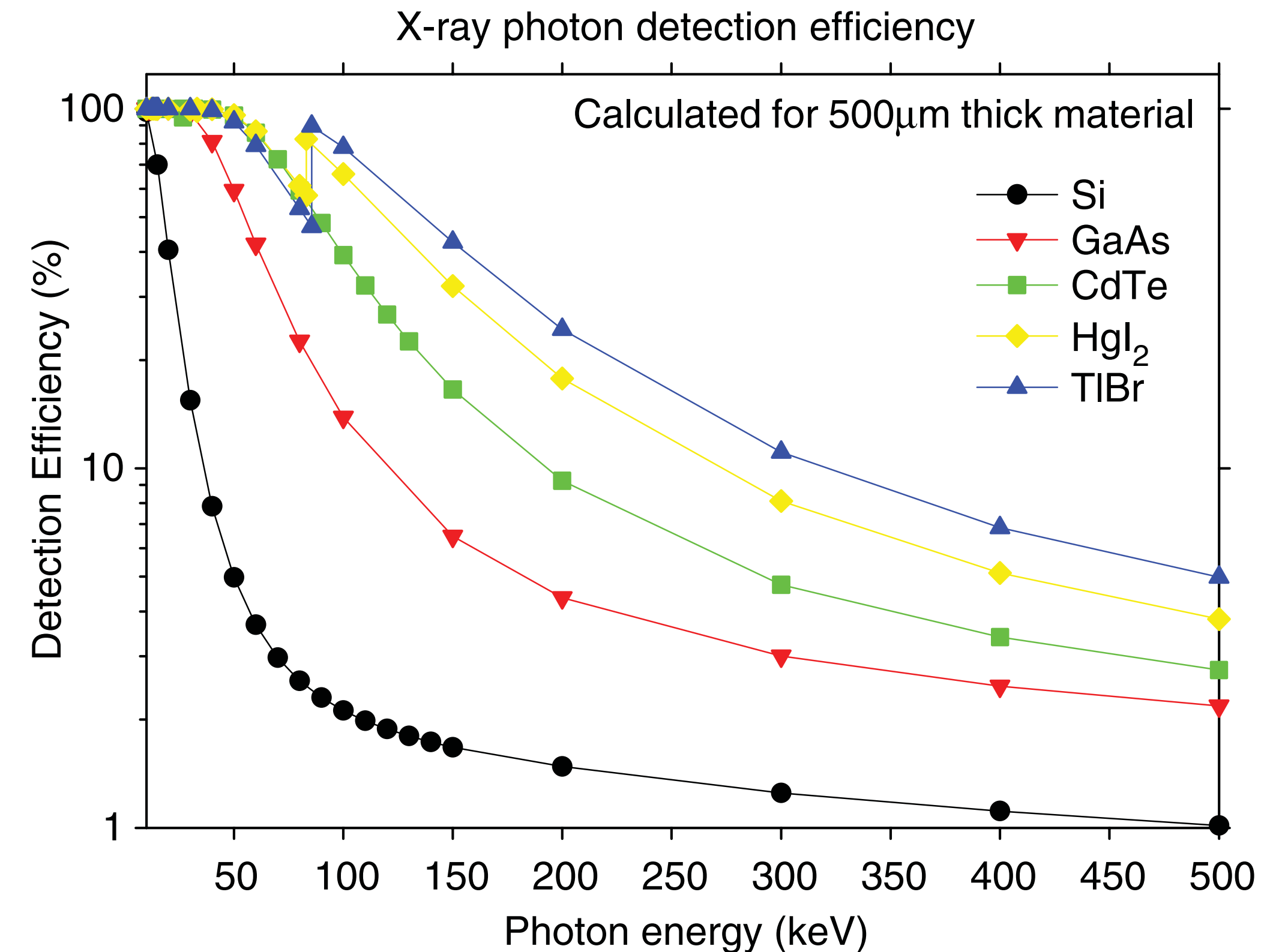
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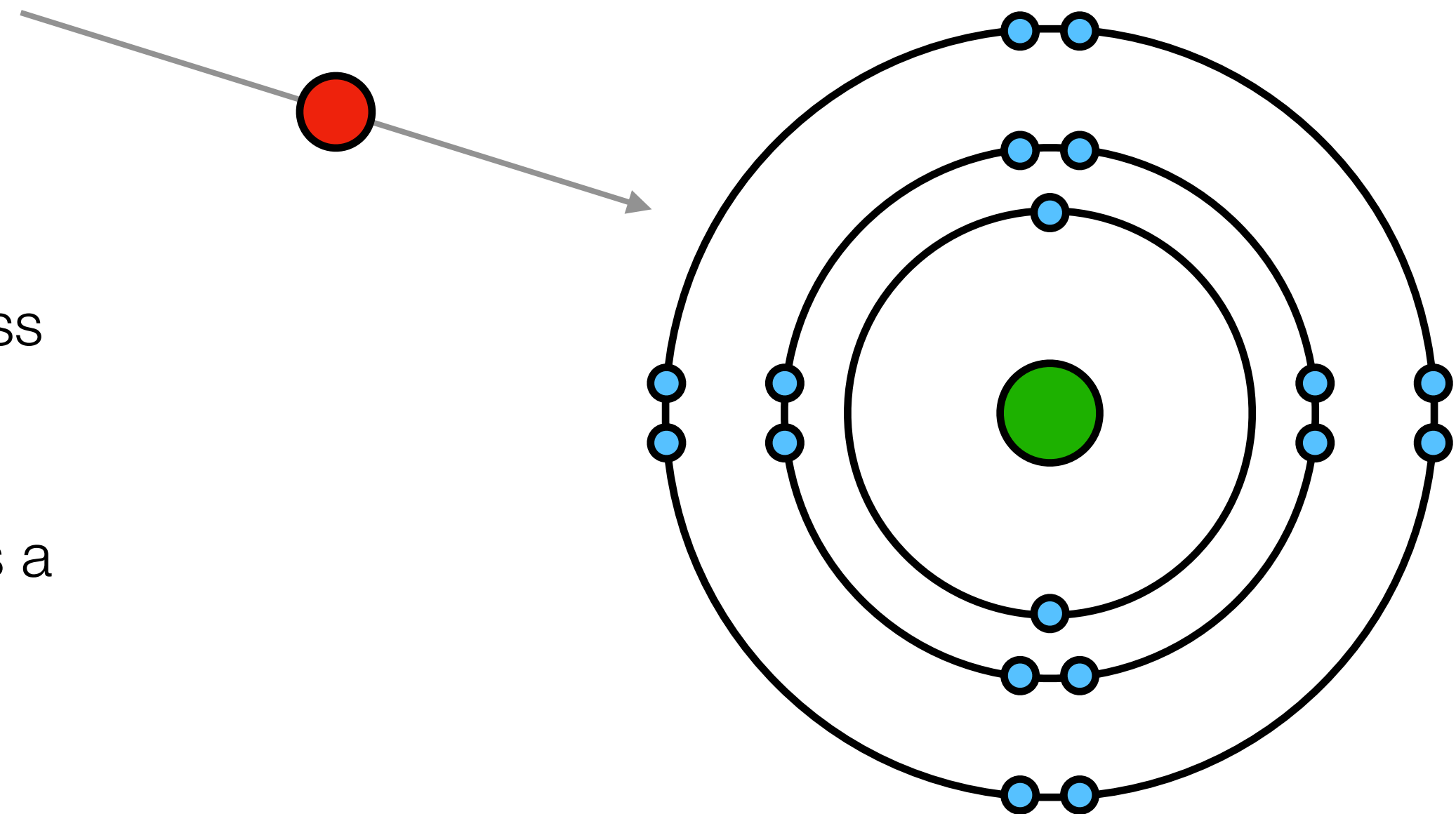
- At low energies materials become transparent
- At higher energies thicker sensors or higher Z materials are often needed in order to ensure that photons interact
- Many x-ray detectors focus on non-silicon sensors such as CdTe, GaAs - but note that silicon benefits from the entire microelectronics industry; everything from single crystal quality to fabrication



Particle interactions

For particles traversing a material, energy is lost through a series of ionisation events, where energy is transferred to individual electrons which are promoted to the conduction band (and may have enough energy to cause further ionisations)

- For electrons this is a bit more tricky since the mass of the two objects is the same
- For all other particles, we can assume that there is a series of interactions, which will distort the particle path and result in energy transfer

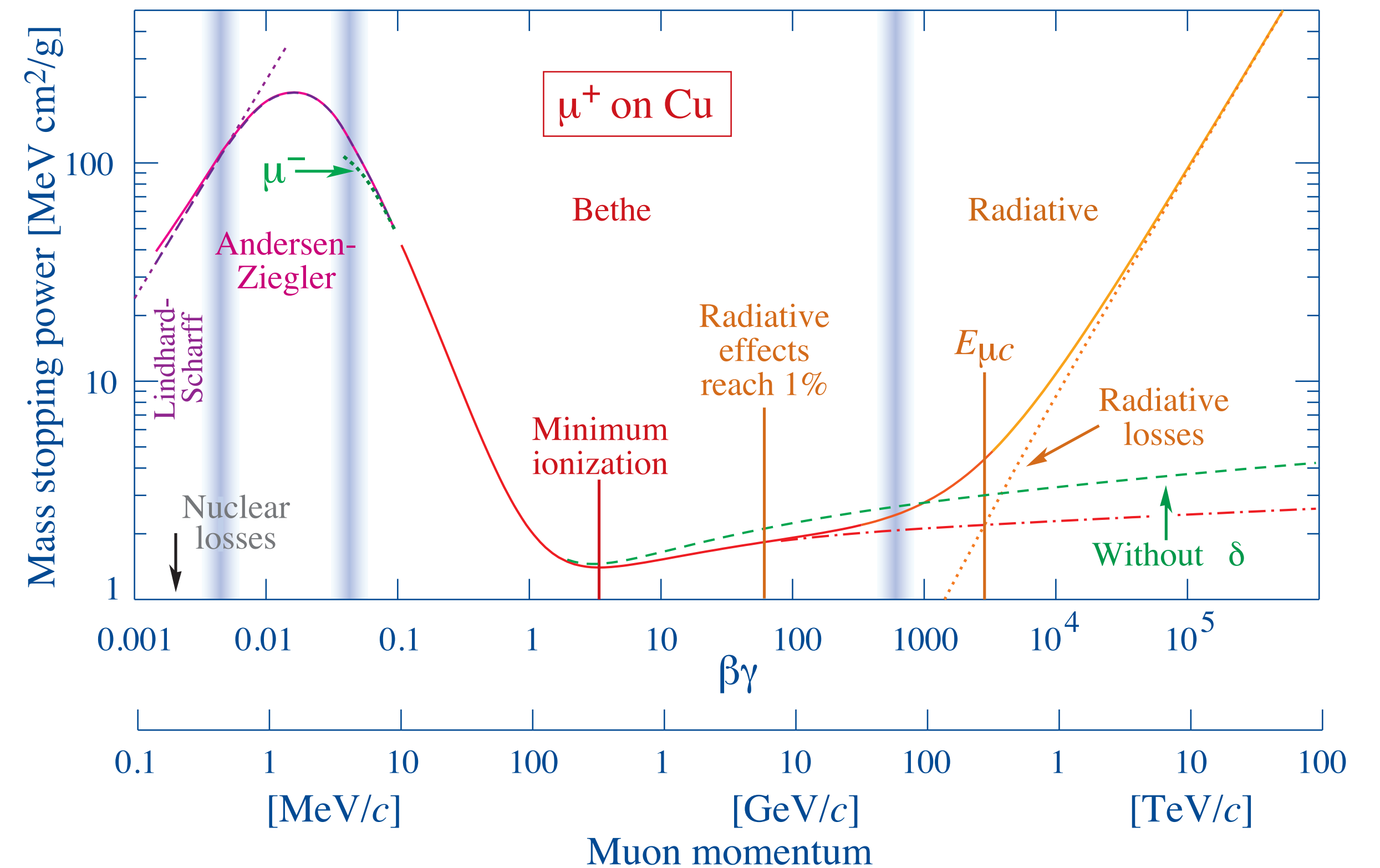


Particle interactions

The energy loss, dE / dx , is described by the Bethe-Bloch equation

$$\frac{dE}{dx} \approx \left(\frac{e^2}{4\pi\epsilon_0} \right)^2 \frac{4\pi z^2}{m_e v^2} N Z \left[\ln \frac{2m_e v^2}{I} - \ln \left(1 - \frac{v^2}{c^2} \right) - \frac{v^2}{c^2} \right]$$

The energy loss is proportional to:



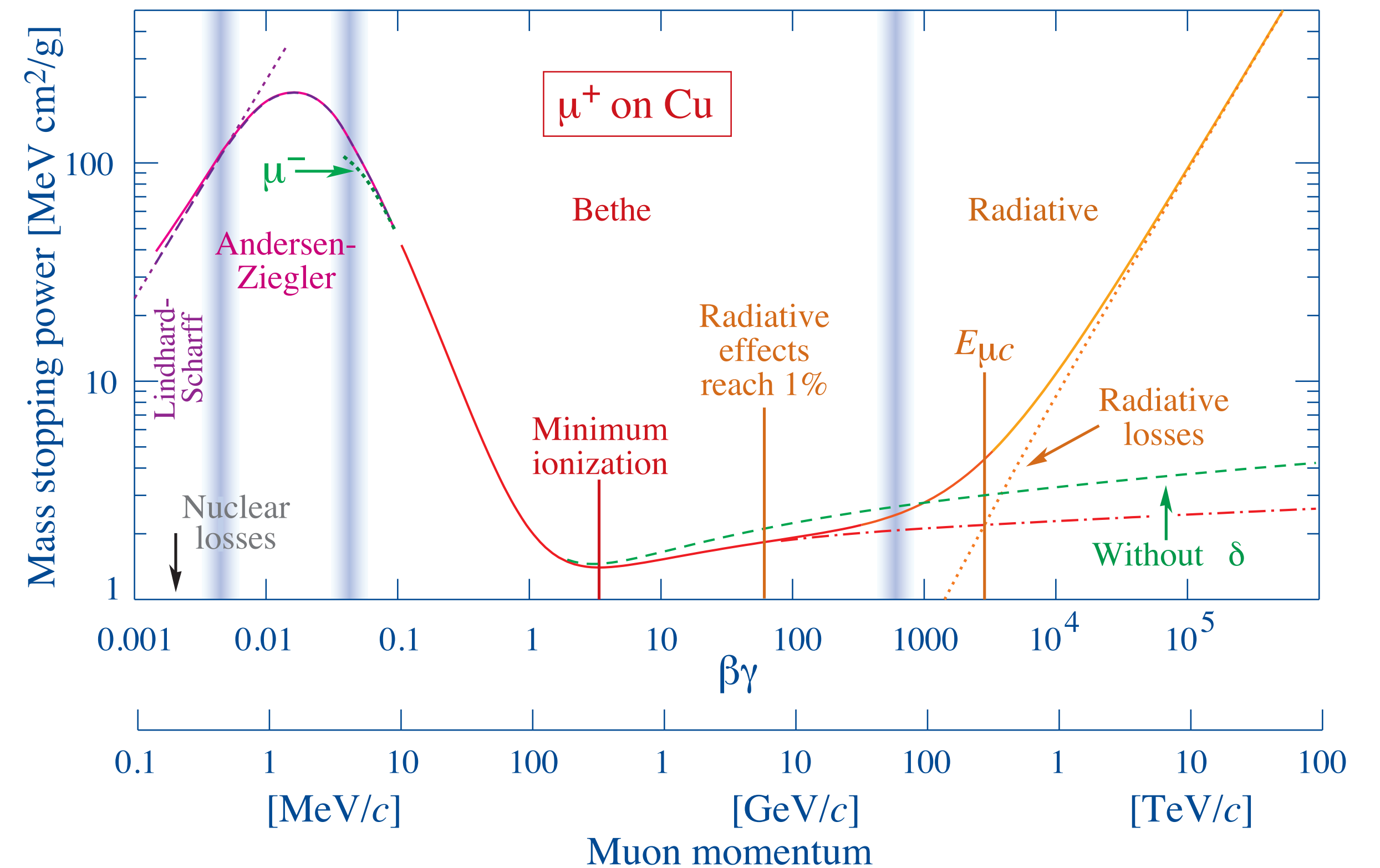
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The energy loss is proportional to:

- The electron density



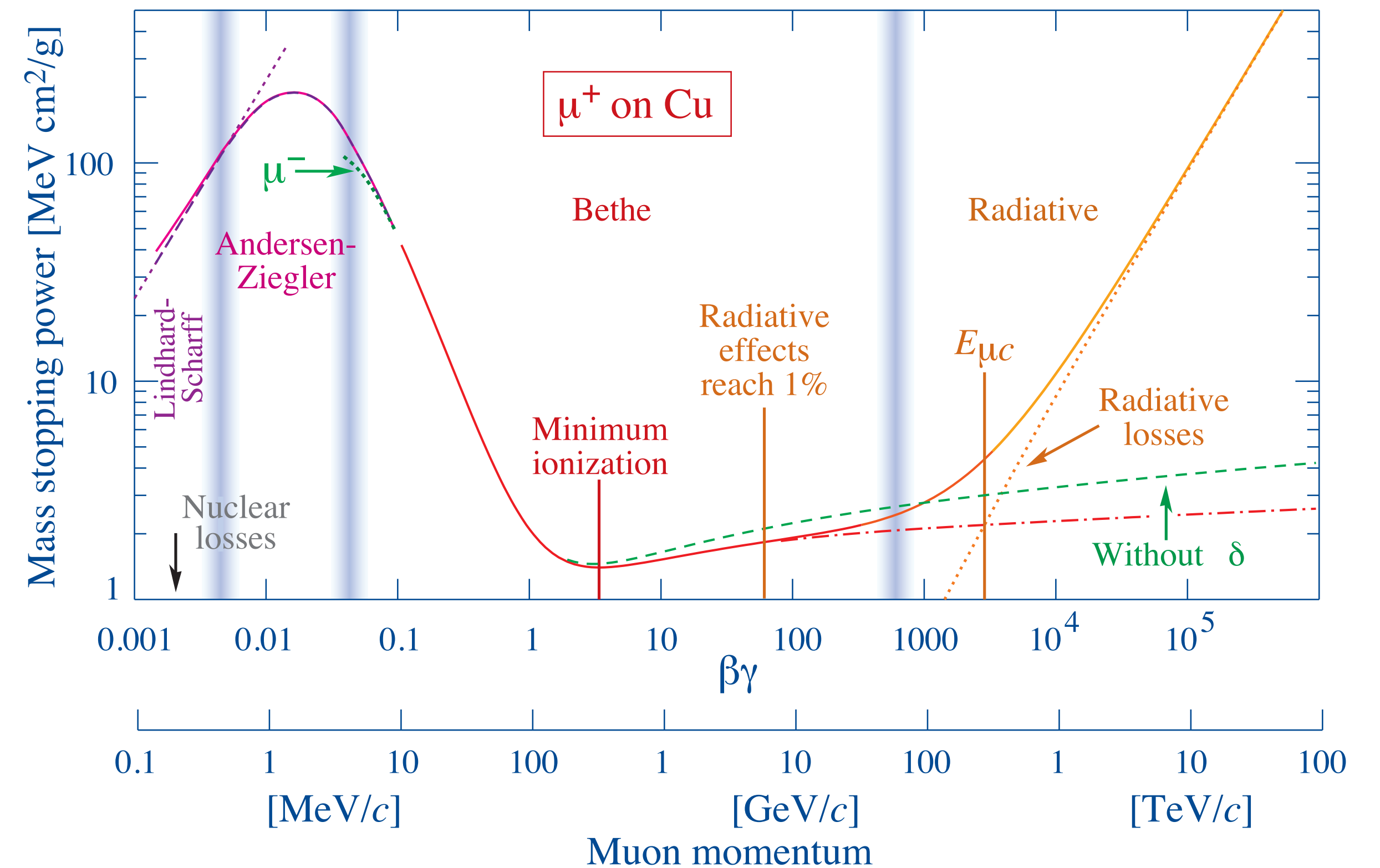
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The energy loss is proportional to:

- The electron density
- The particle mass (at non-relativistic energies)



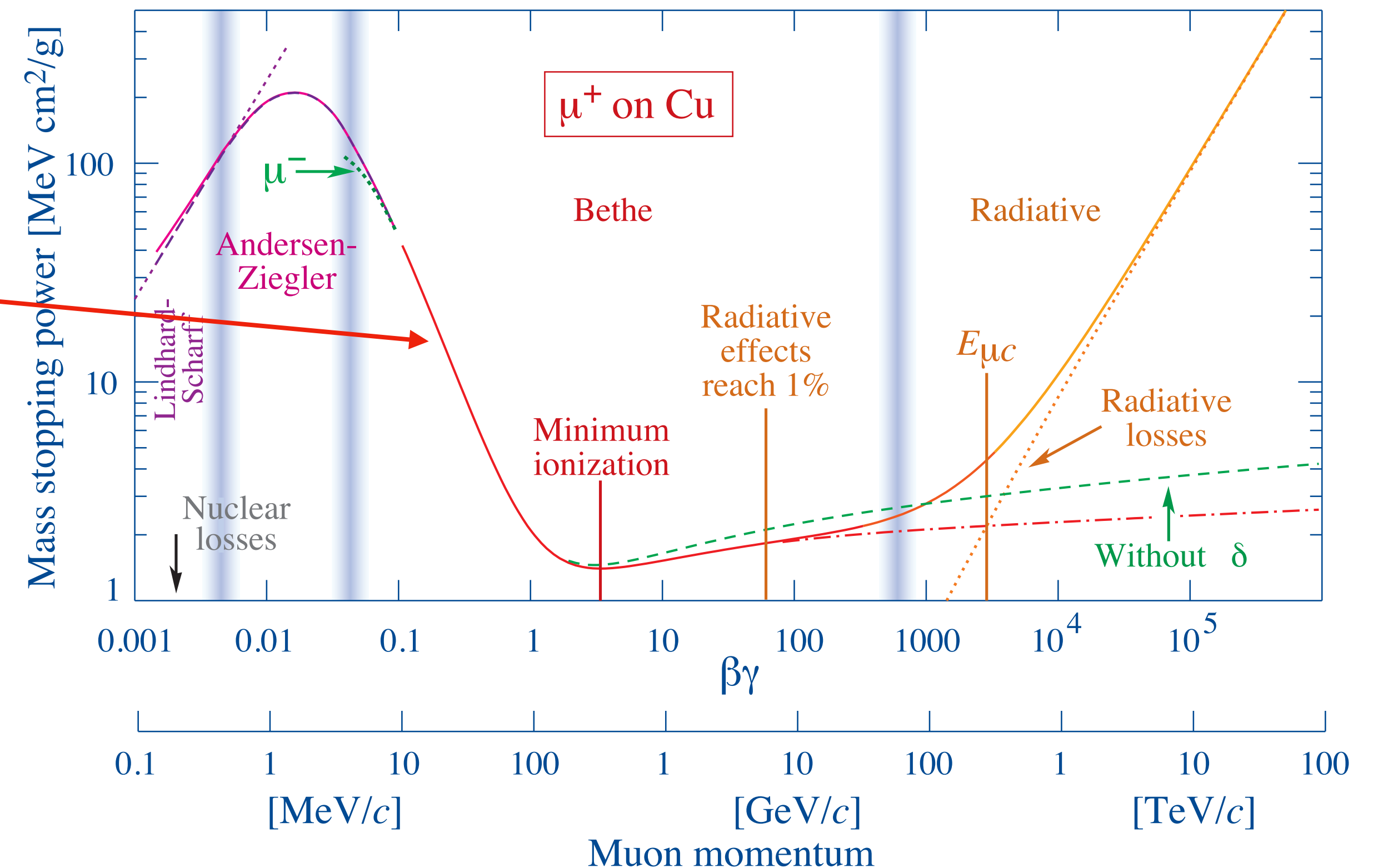
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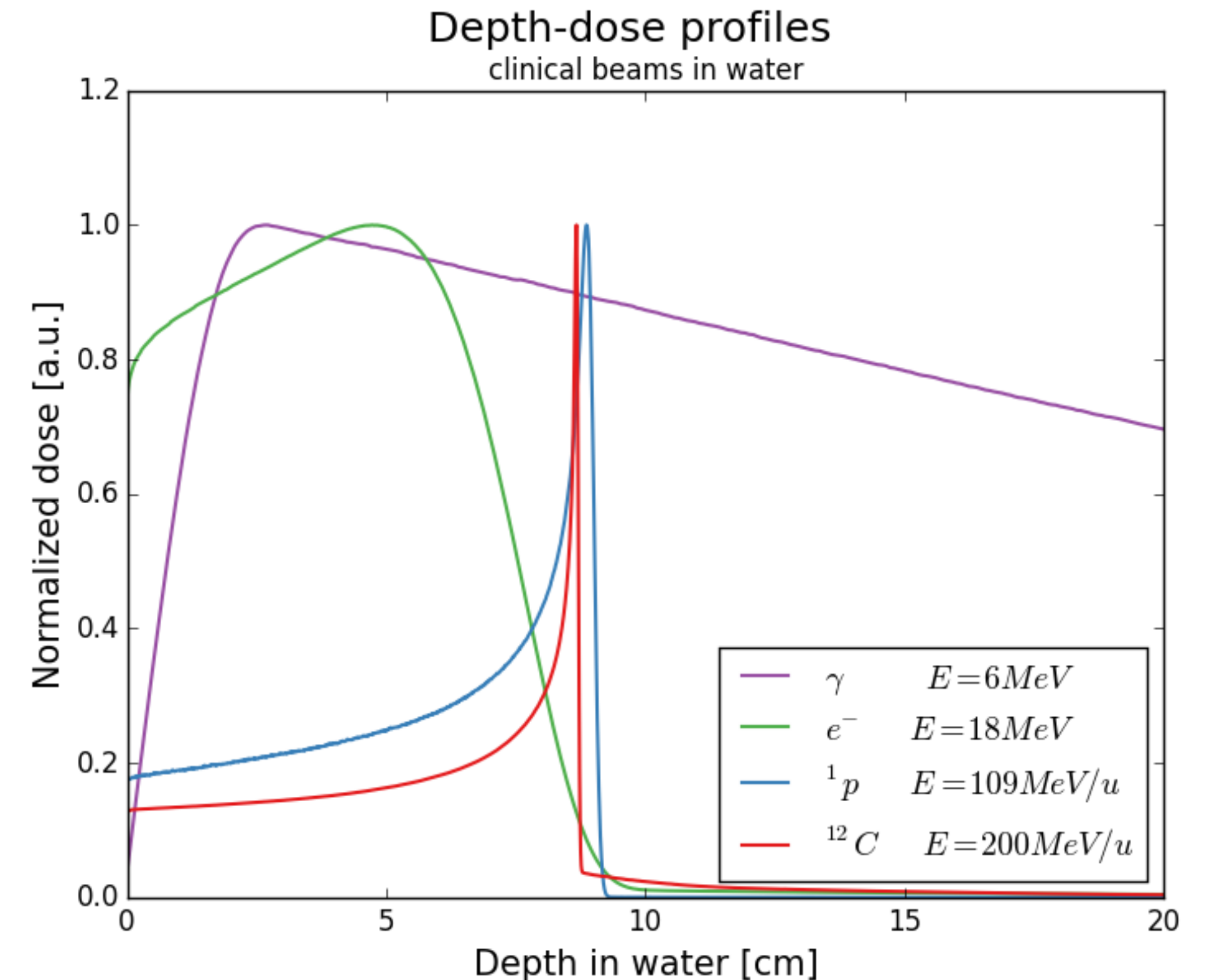
- The electron density
- The particle mass (at non-relativistic energies)
- β^{-2} when relativistic



Bragg peak

In some cases the particle energy is not high enough to pass fully through a given medium

- As the energy gets lower, the probability to interact increases, the higher the energy transfer per unit length
- This leads to a peak in energy as the particle reaches it's end-point - the **Bragg peak**
- This is relevant for detectors, but more important when thinking of targeted radiotherapy: the target zone can have a substantially higher dose than the surrounding healthy tissue



Tiny Fast Flying Things That Break The Inside Parts Of Your Body

(Radiation damage)

Radiation-induced defects

Particle interaction with a detector can lead to either

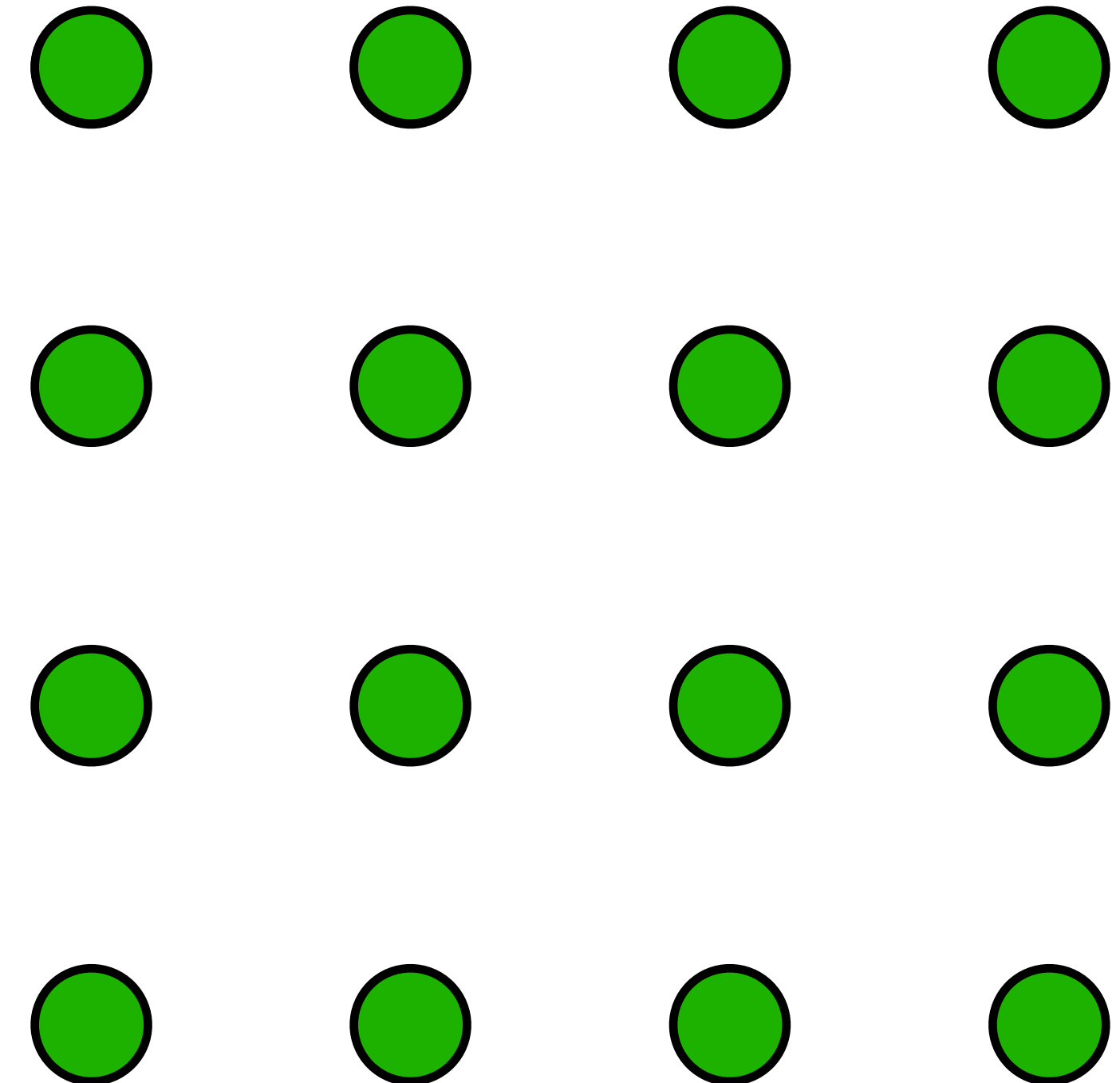
- Non-ionising energy loss
- Ionising energy loss

This could conceivably have detrimental effects

- Non-ionising effects - lattice displacements
- Gradual build-up of charge
- Single large charge deposit effects

Non-ionising energy loss (NIEL)

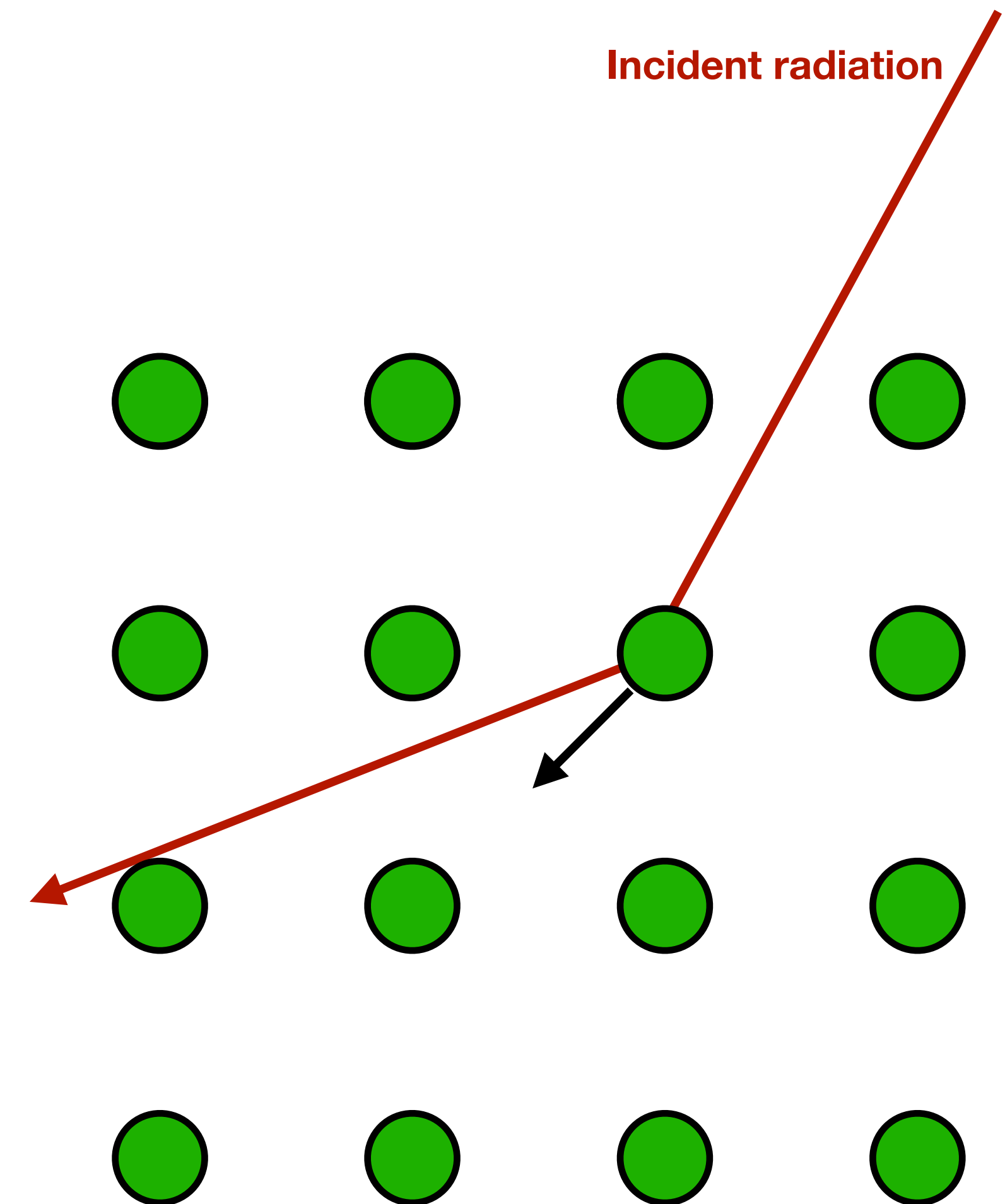
Depending on the incident radiation type, damage can be done to the periodic silicon lattice, representing displacement of atoms from their regular positions



NIEL

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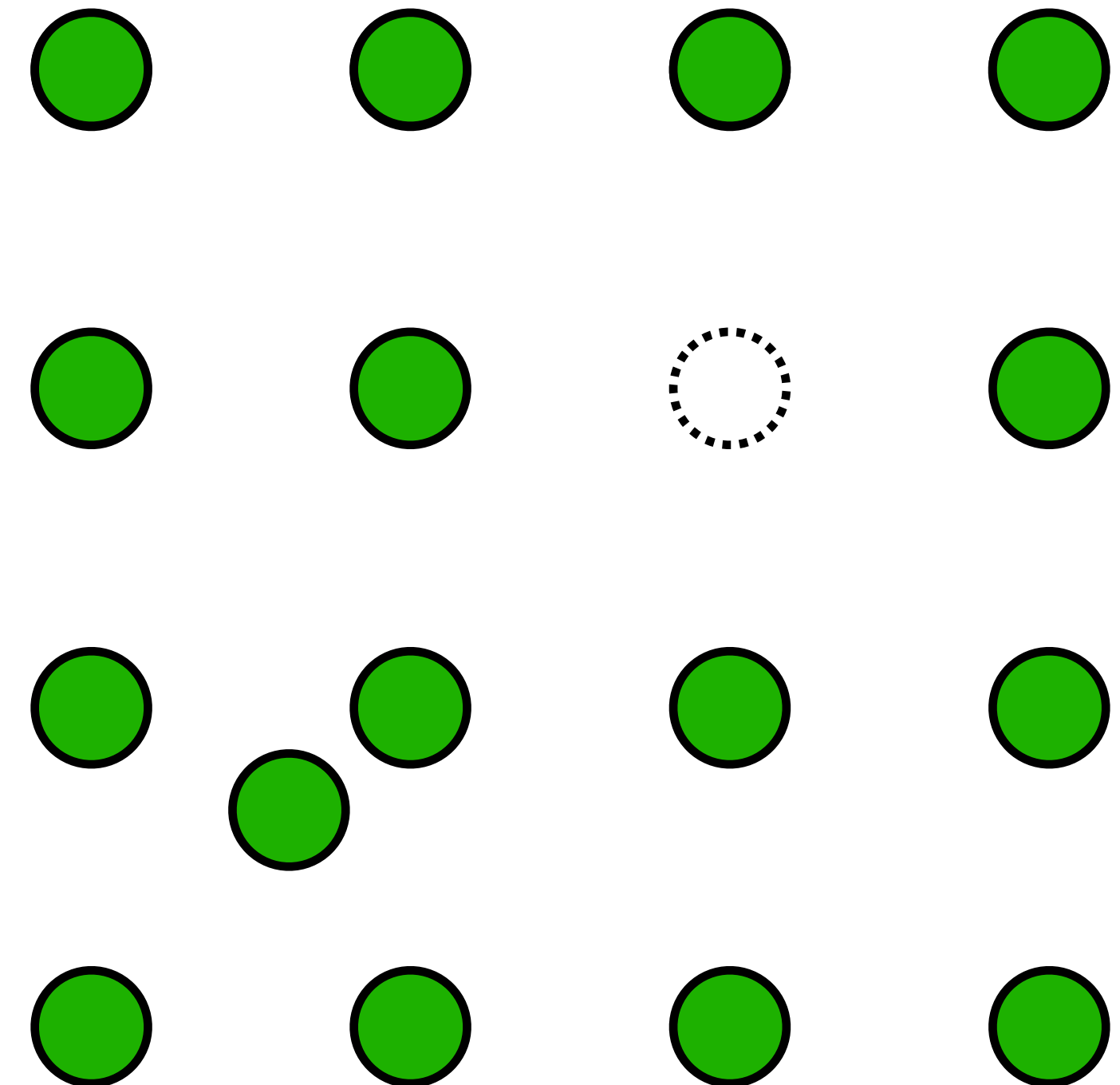
- Nuclear recoil leads to ejection of the silicon atom from the lattice



NIEL

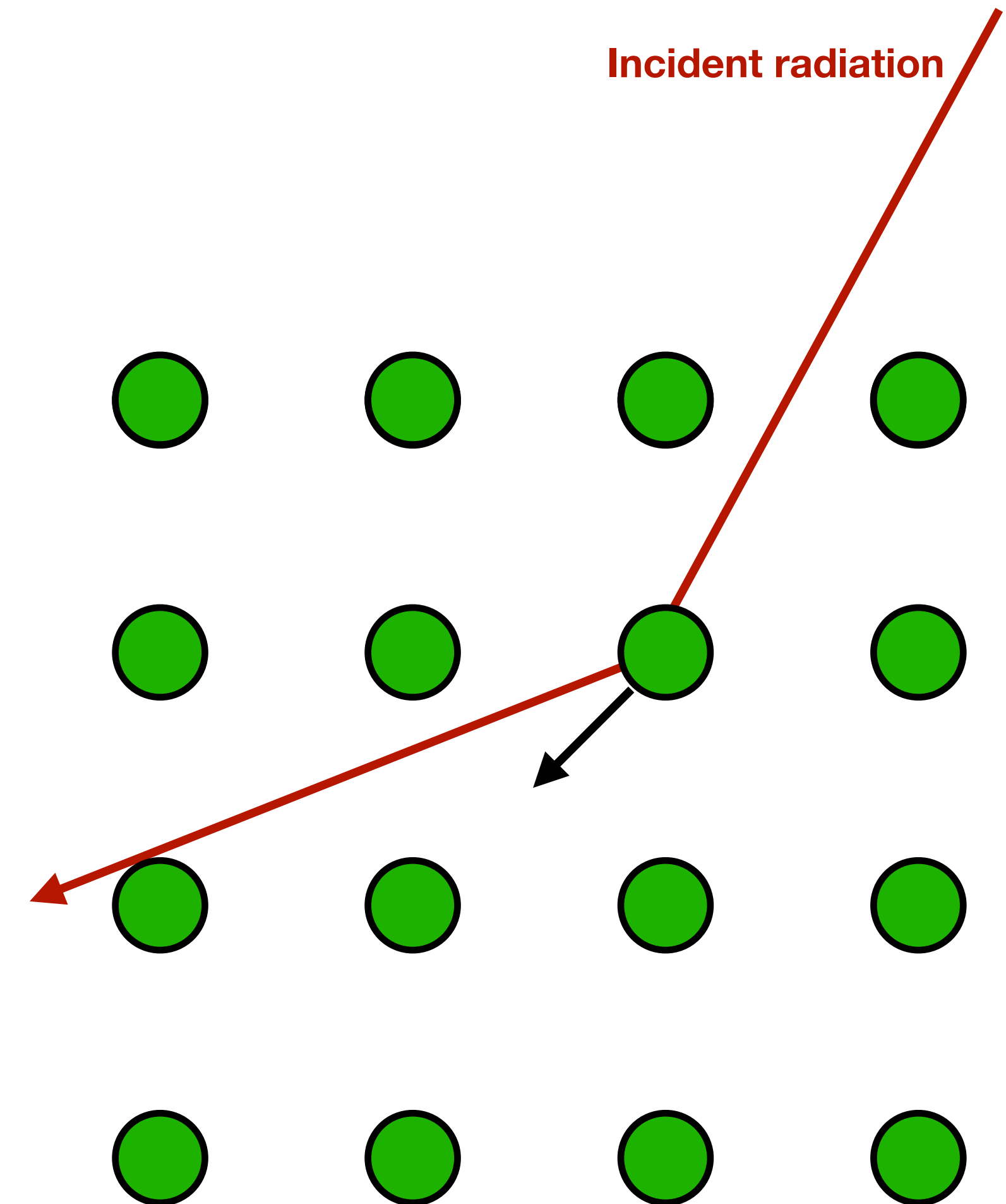
Depending on the incident radiation type, damage can be done to the periodic silicon lattice, representing displacement of atoms from their regular positions

- Nuclear recoil leads to ejection of the silicon atom from the lattice
- We produce **two** defects - an interstitial silicon atom, and a vacancy in the lattice (Frenkel defect)



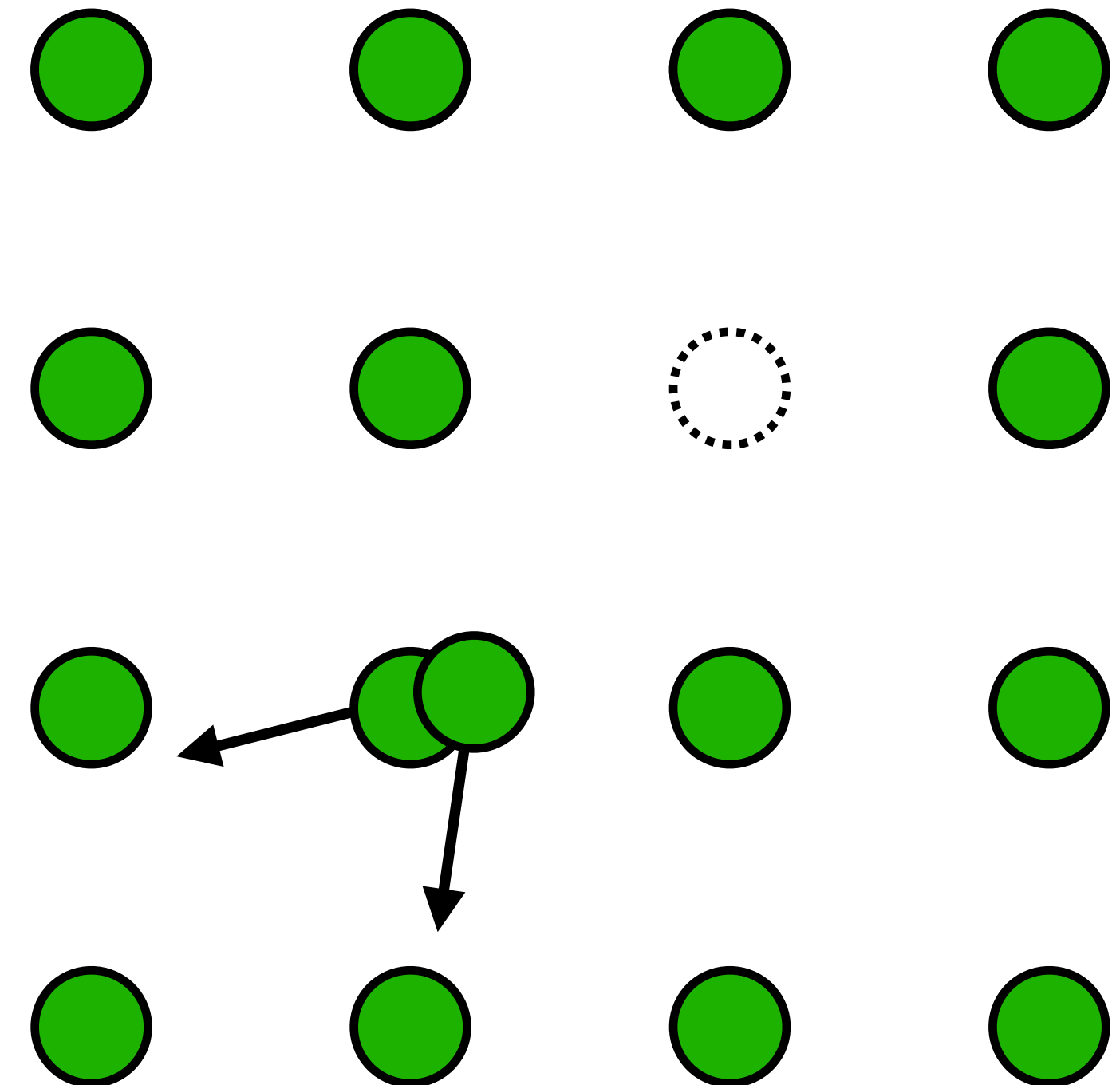
NIEL

Alternatively, our incident radiation could impart enough energy onto the silicon atom that it is able to further displace atoms within the lattice



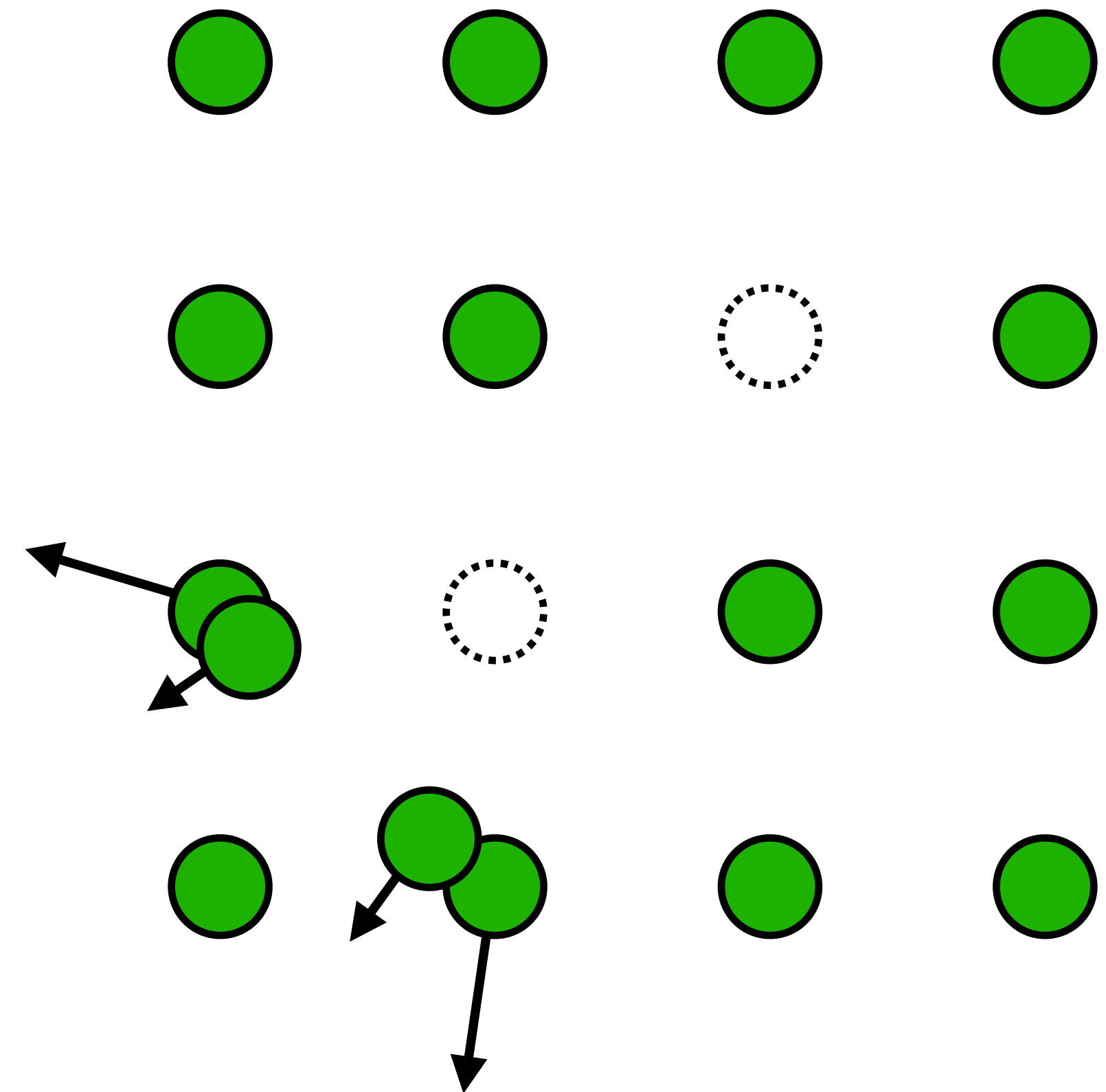
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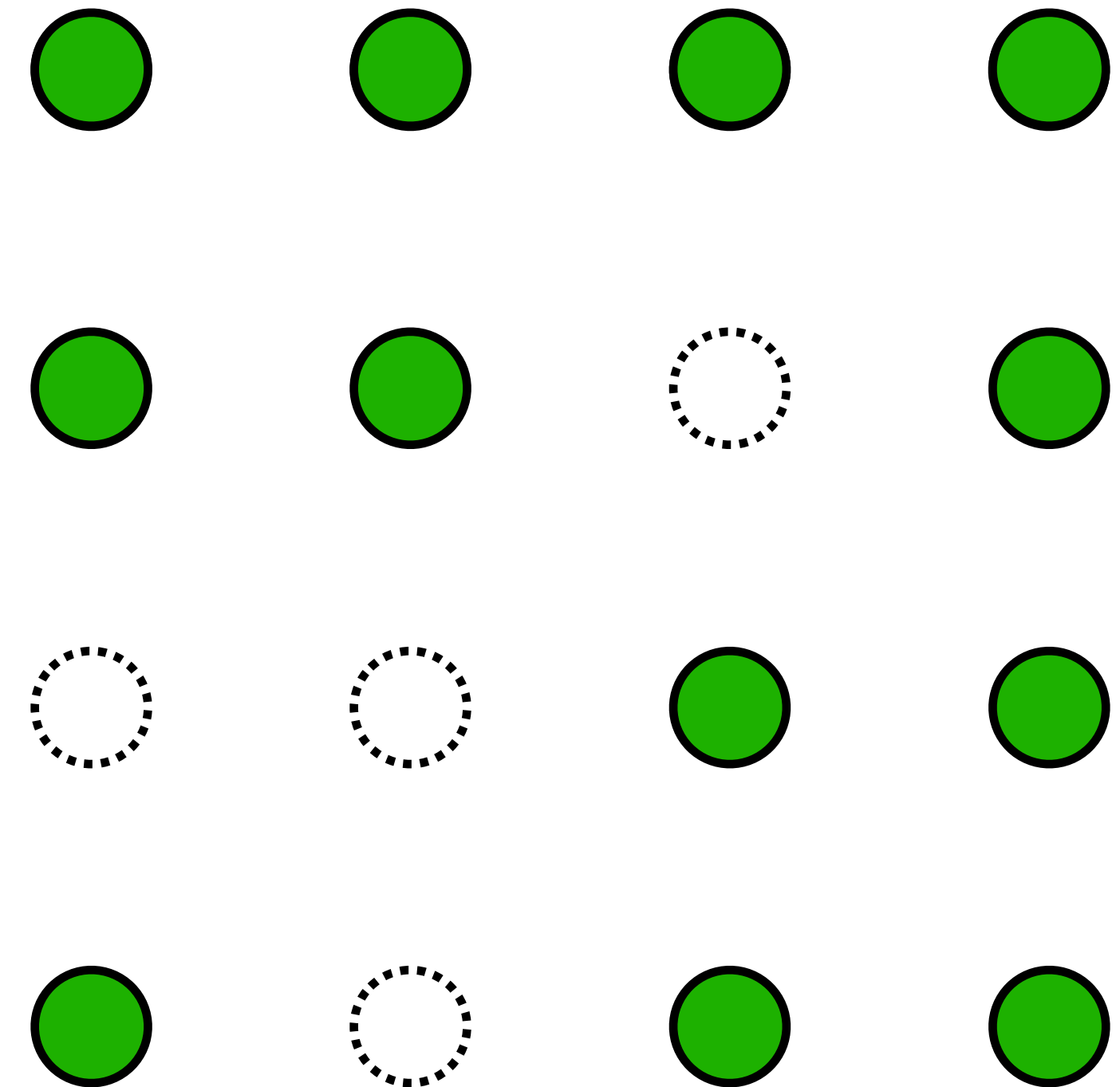
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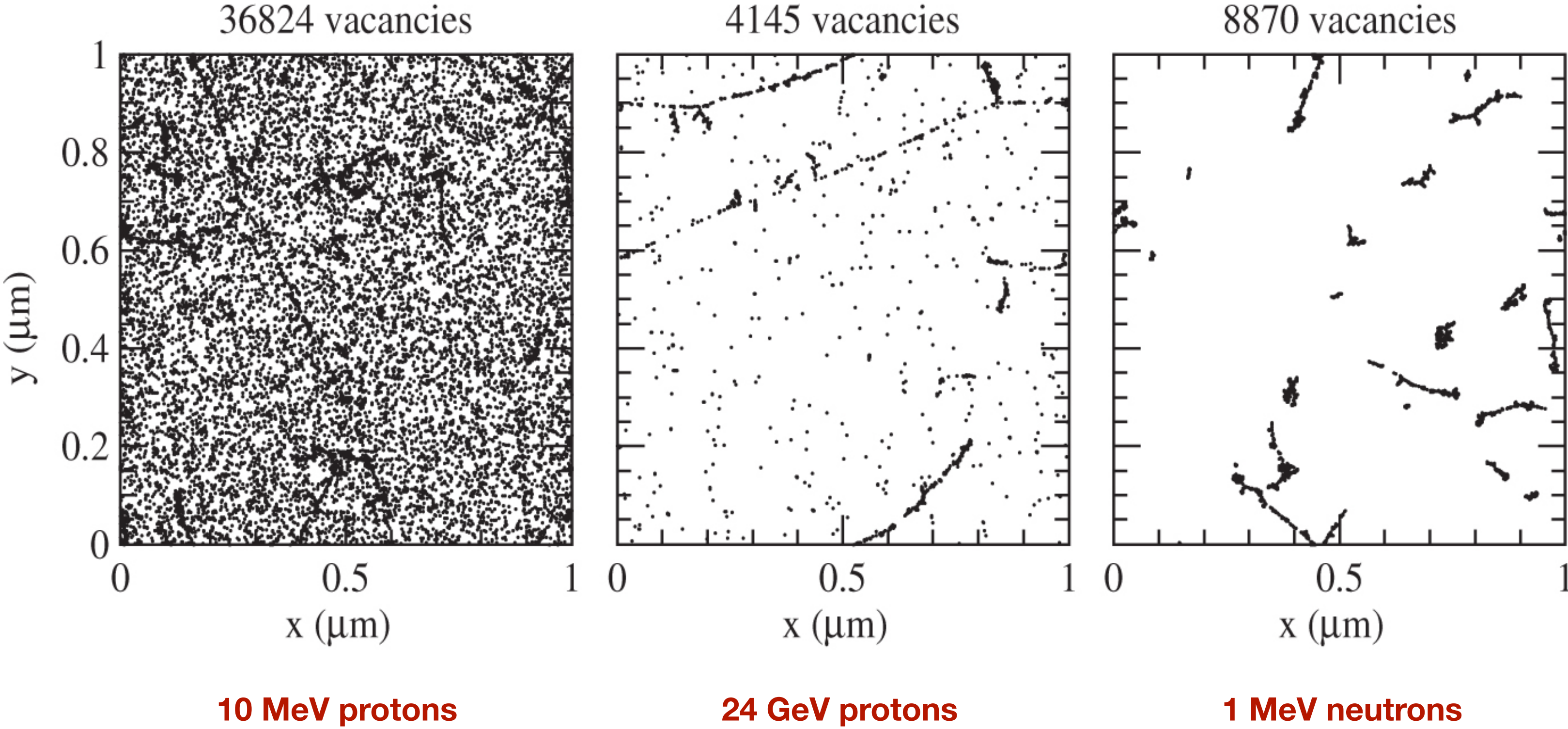
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- We often distinguish between “point” defects and “cluster” defects to highlight these two separate cases



NIEL

All irradiated to 10^{14} cm^{-2}
Depth 1 μm



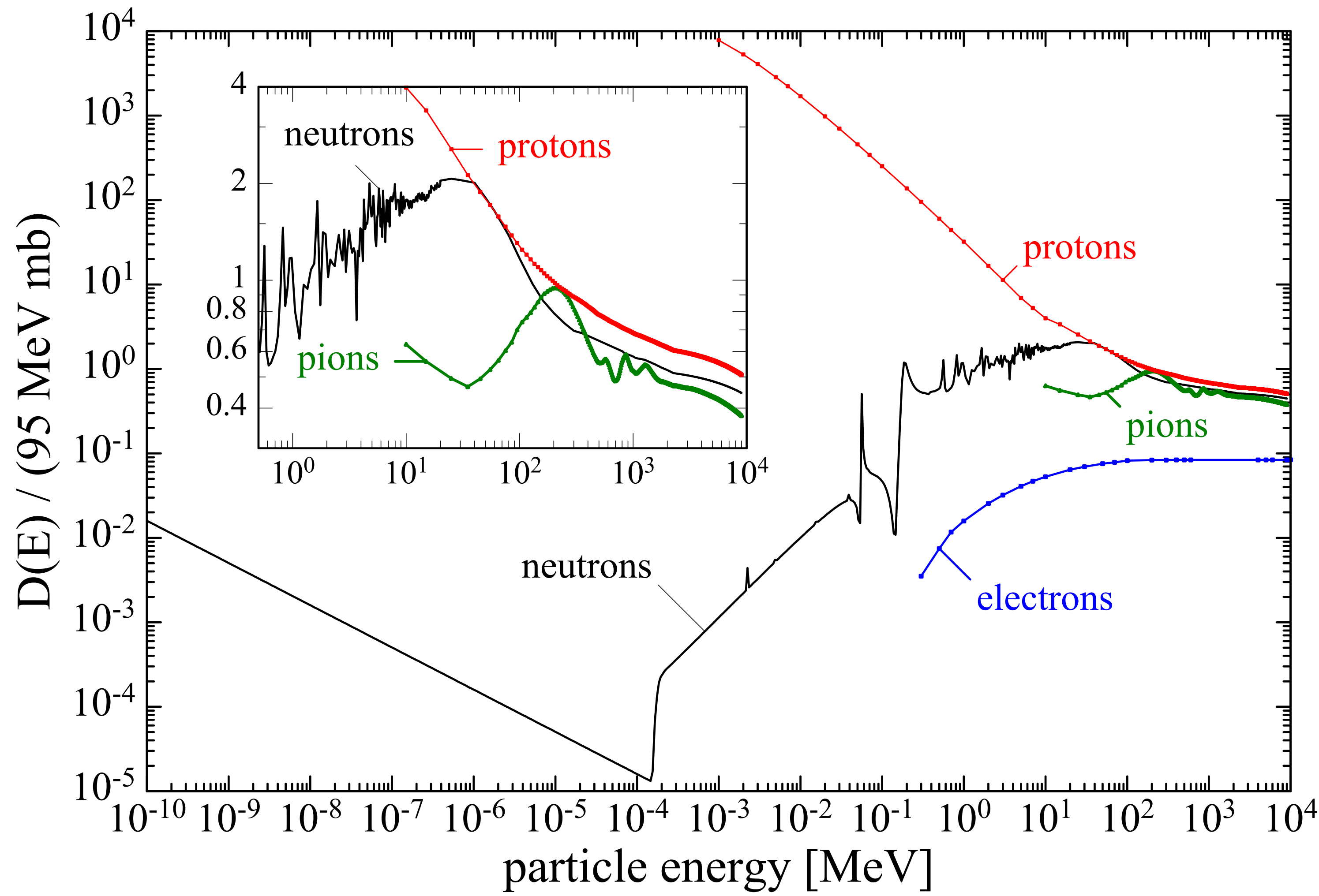
NIEL scaling hypothesis

The NIEL scaling hypothesis is:

Damage effects by energetic particles in bulk material can be described by the total non-ionising energy loss to the lattice

Measurements of NIEL are generally expressed in units of $1 \text{ MeV } n_{\text{eq}} \text{ cm}^{-2}$, and the damage inflicted by any particle of any energy can be expressed using a hardness factor, k

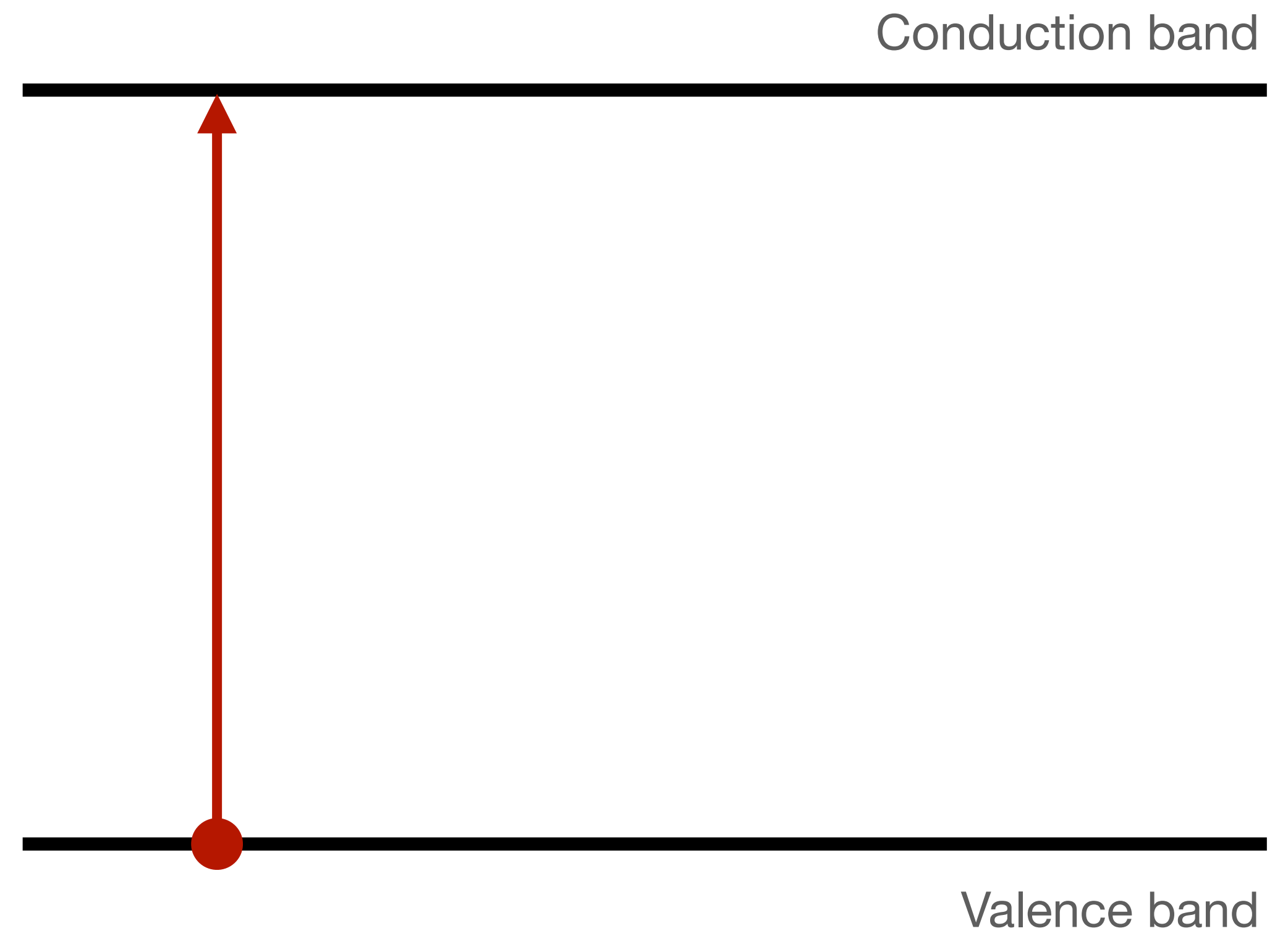
Hardness factors for different particles in silicon



Energy levels in semiconductors

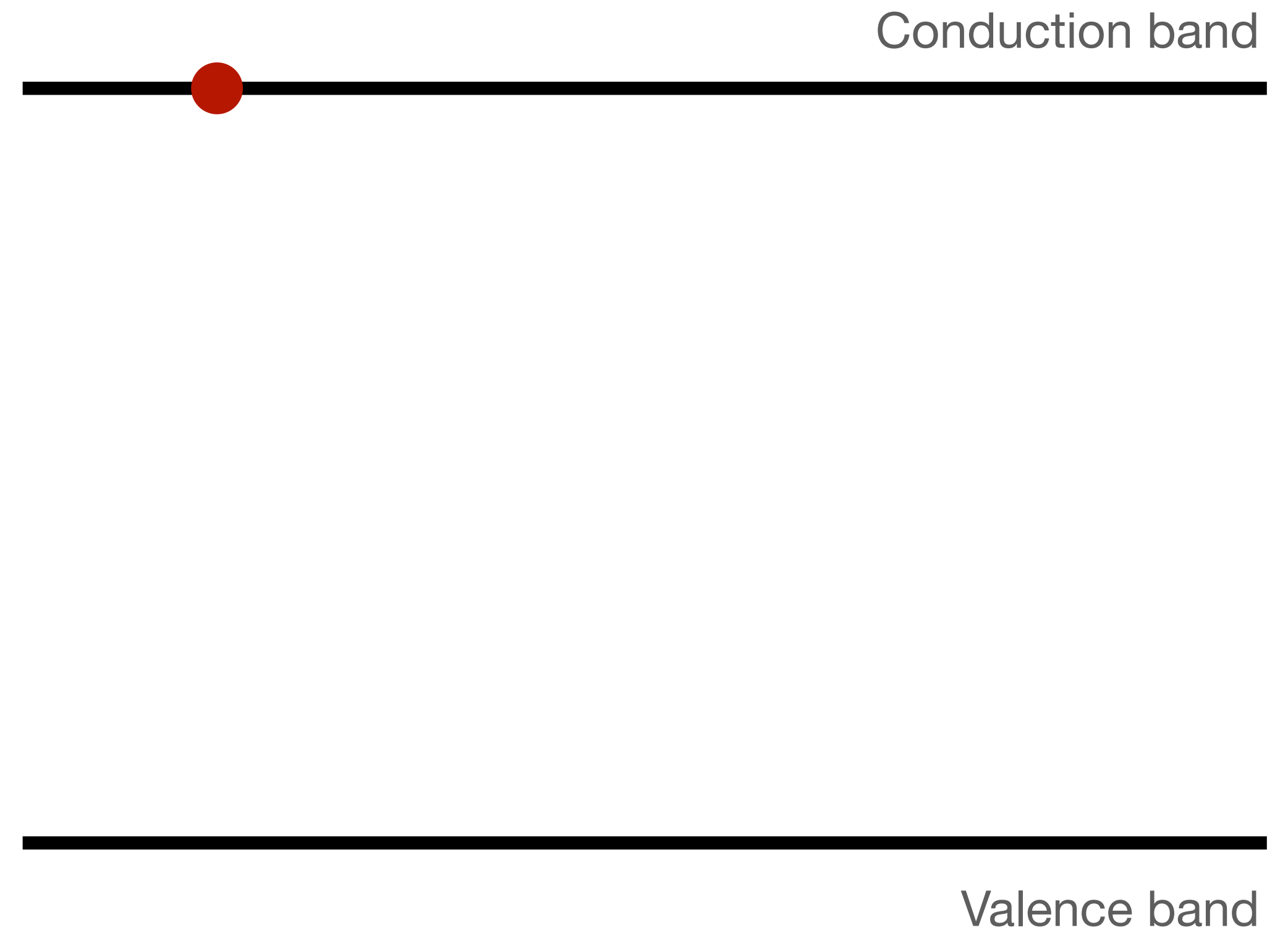
Charge carrier production in semiconductors relates to transitions from the valence to conduction band

- In depleted silicon detectors, this is the mechanism by which interacting particles are observed



Energy levels in semiconductors

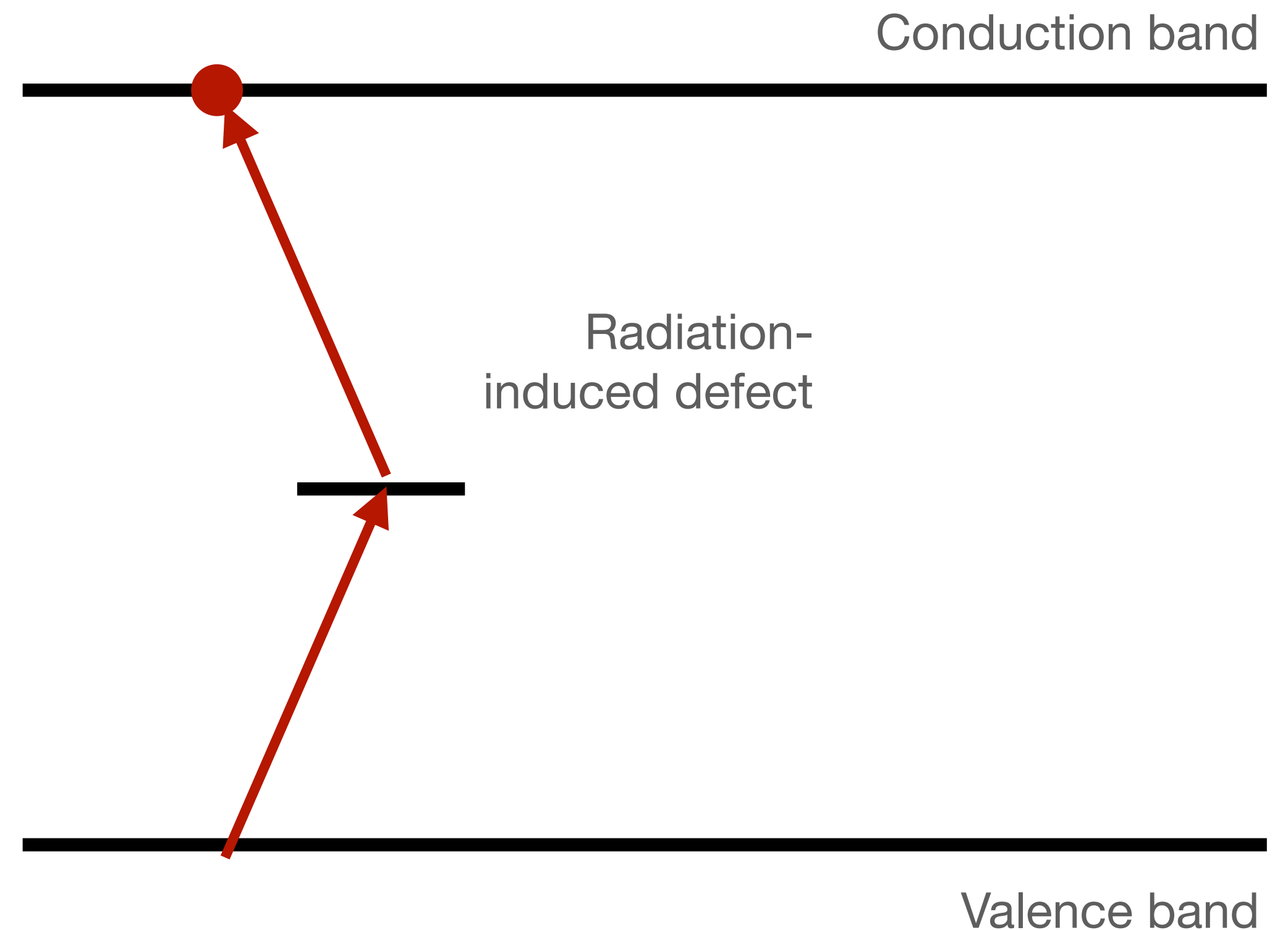
Where this goes wrong after irradiation damage is the introduction of additional energy levels in the band gap



Energy levels in semiconductors

Where this goes wrong after irradiation damage is the introduction of additional energy levels in the band gap

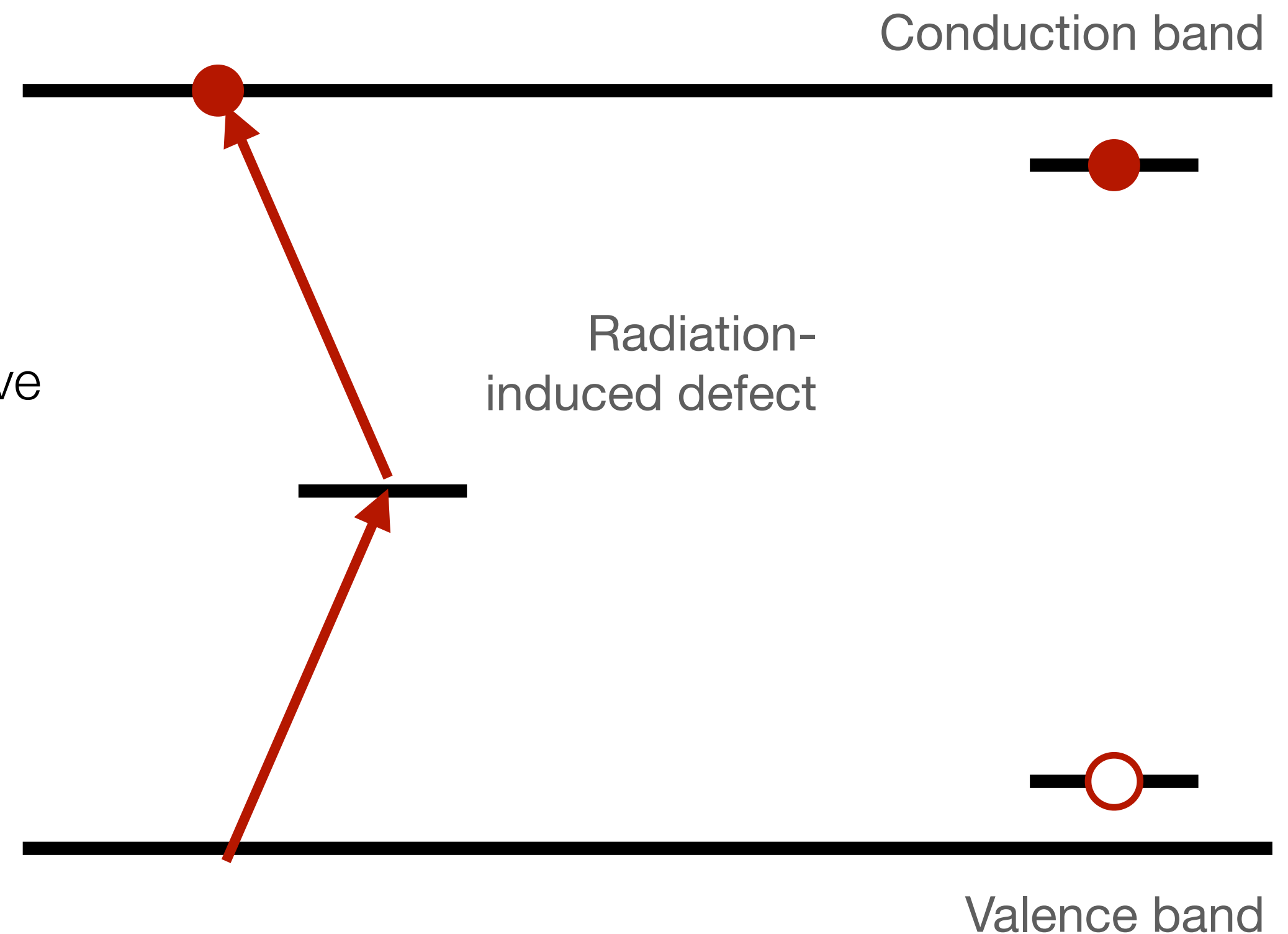
- Mid-level defects provide a shorter path for thermally-generated pairs => increased leakage current



Energy levels in semiconductors

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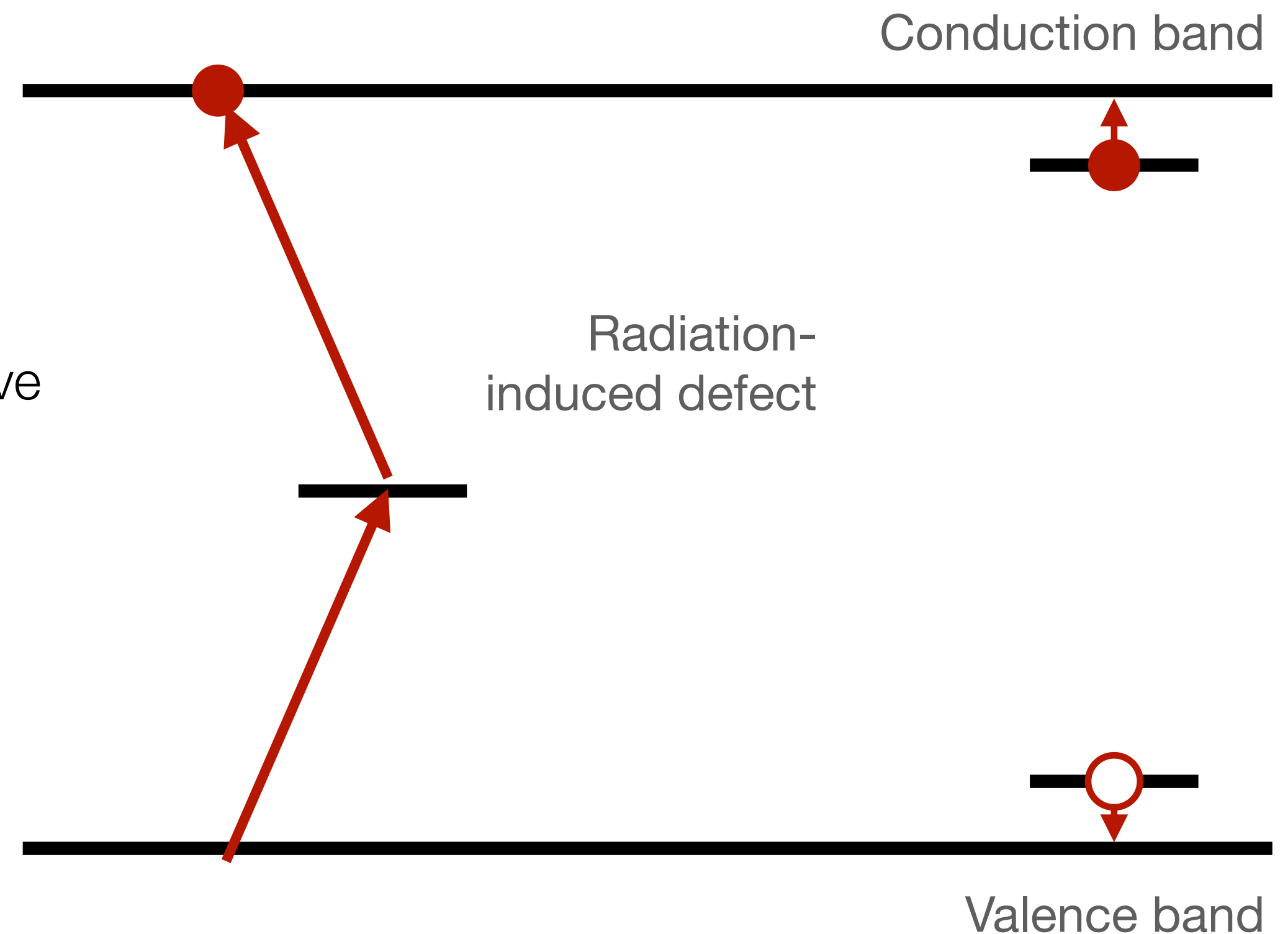
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Energy levels in semiconductors

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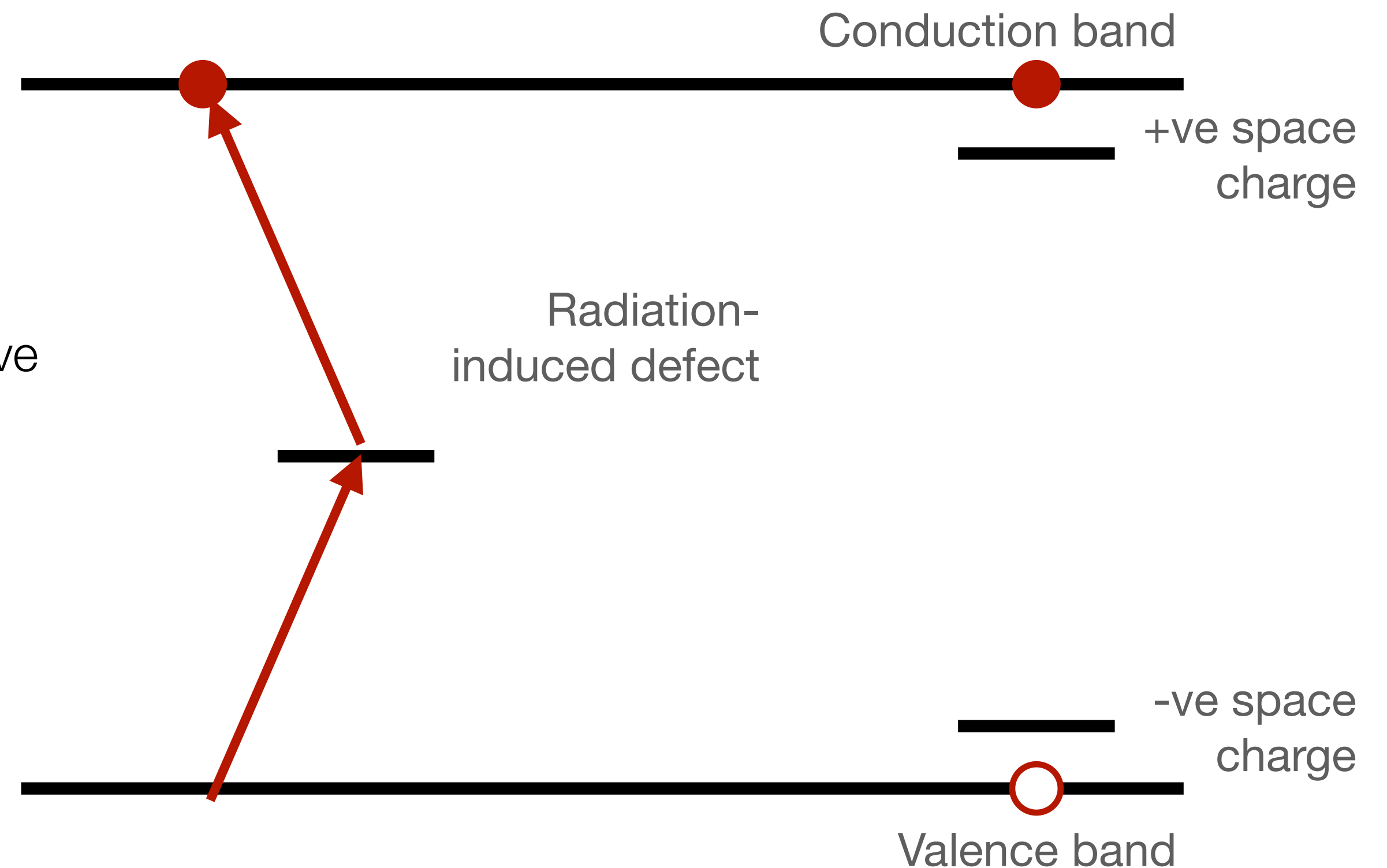
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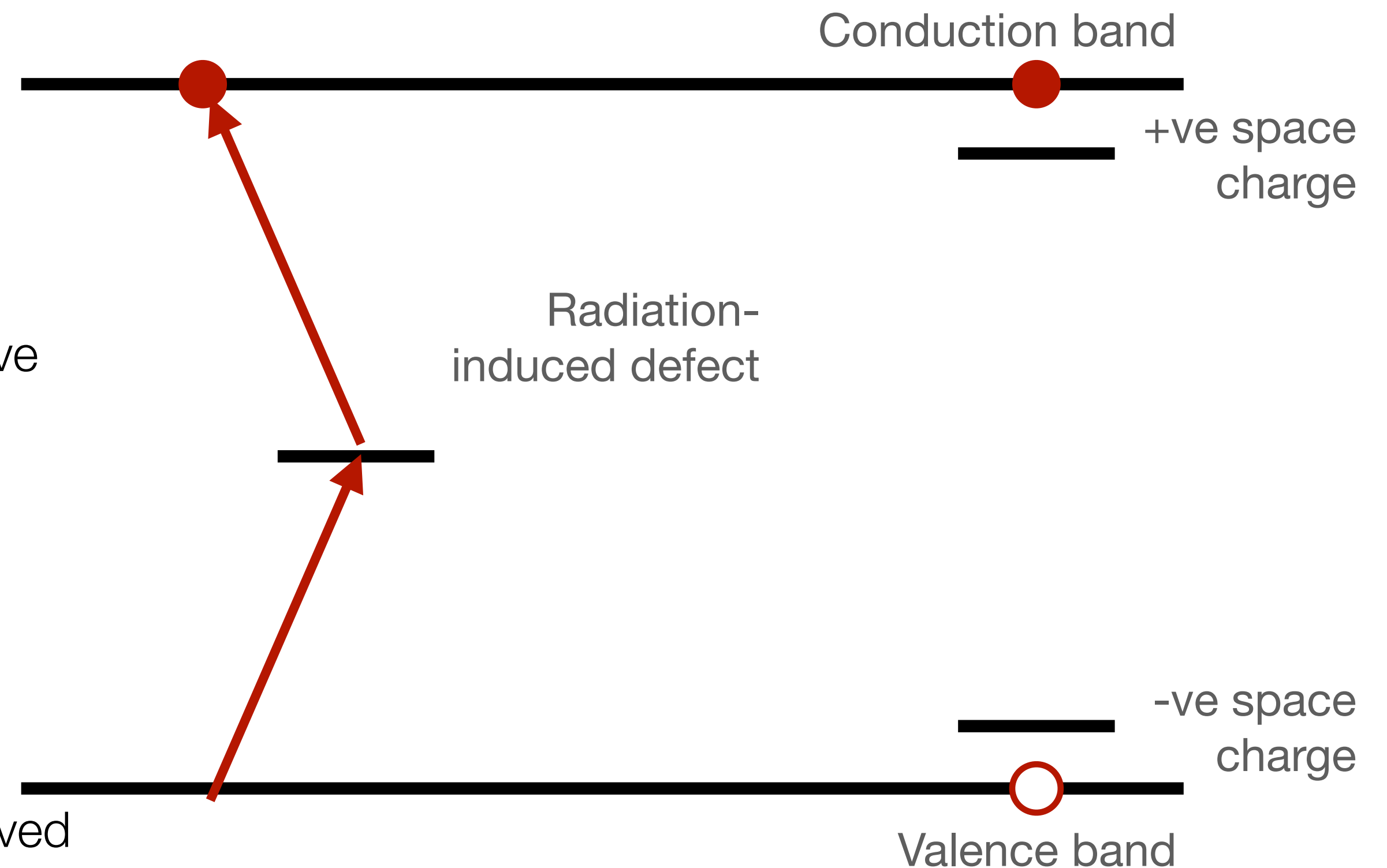
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Where this goes wrong after irradiation damage is the introduction of additional energy levels in the band gap

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In addition, charge carriers can become trapped in defect centres

- Depending on the emission time can reduce the observed signal
- Use carriers with highest mobility!

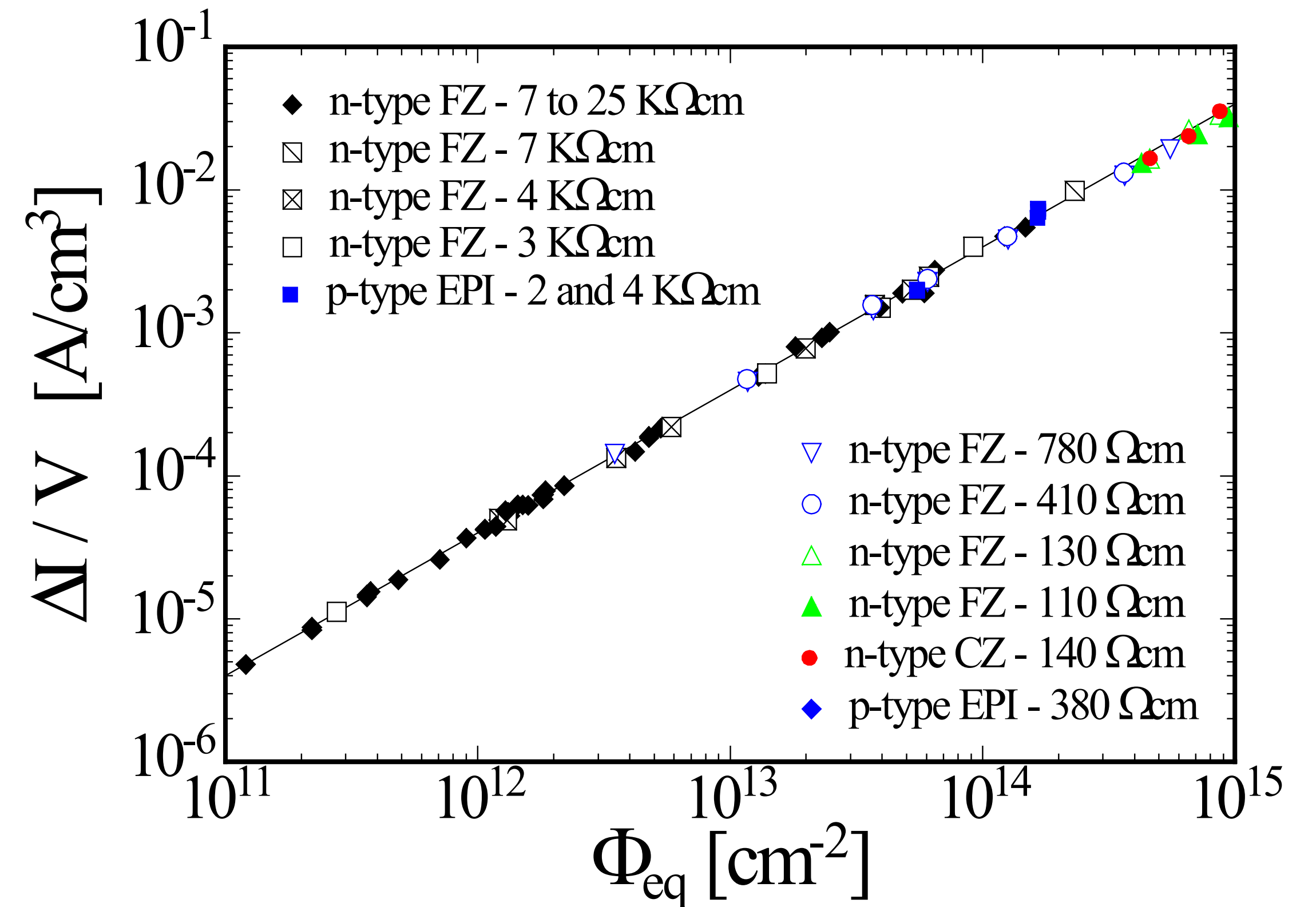
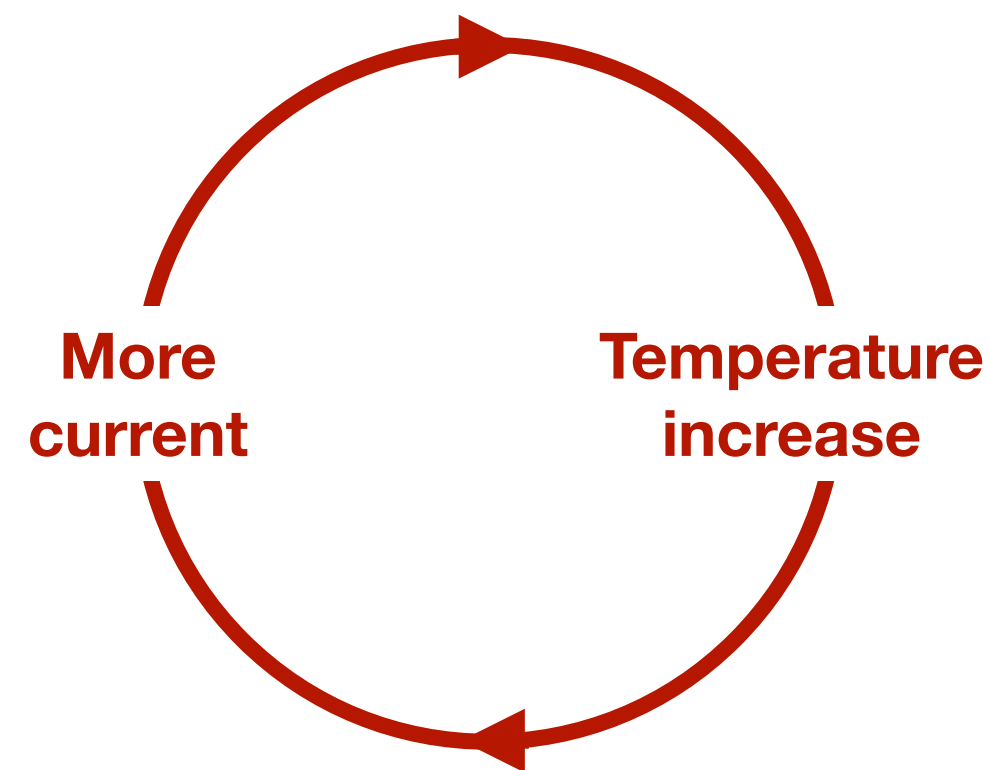


NIEL damage - leakage current

Leakage current increase matches well the NIEL scaling hypothesis over several orders of magnitude

Several significant consequences

- Increased current (noise) in the readout electronics
- Increased power consumption
- Risk of **thermal runaway**



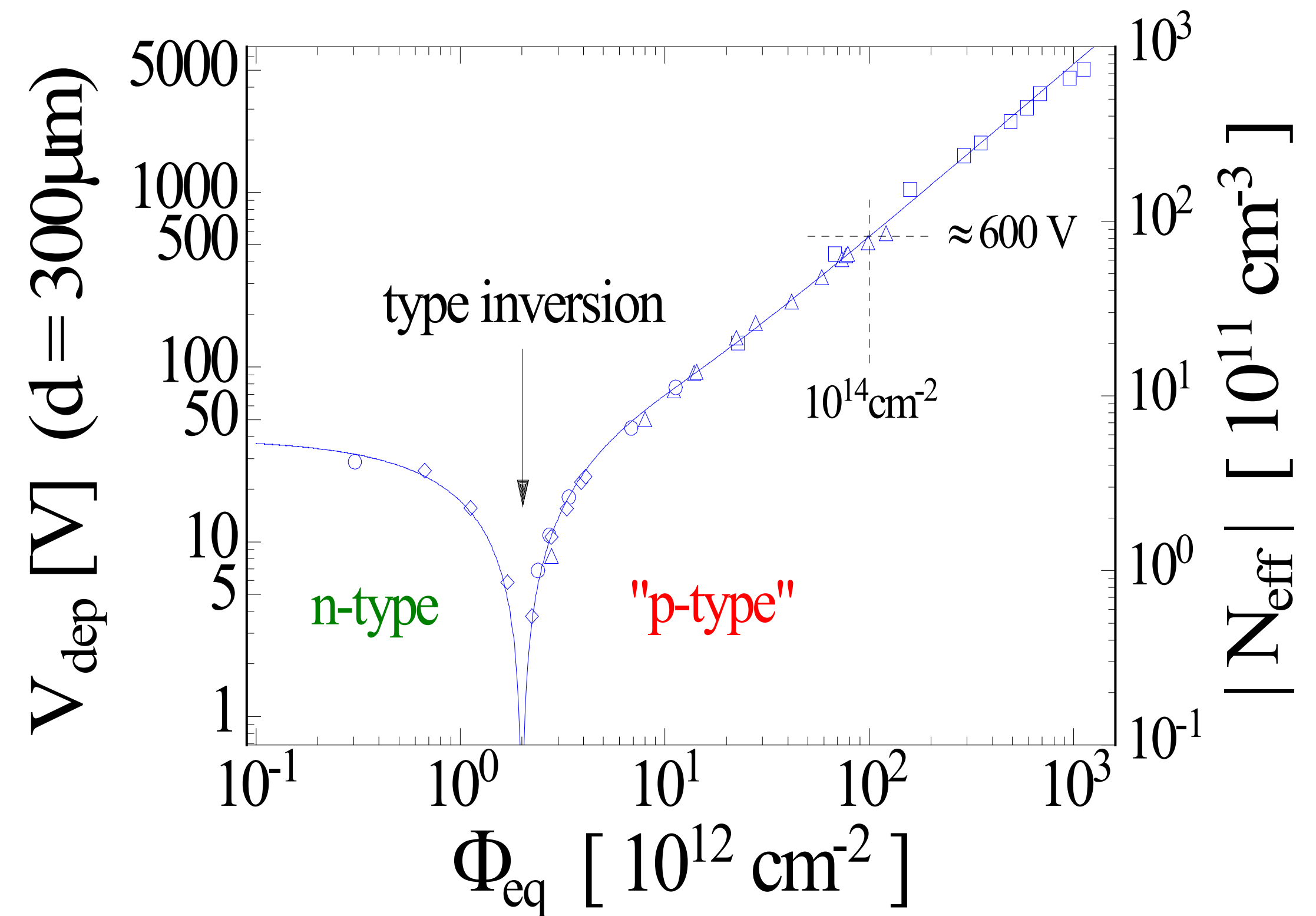
NIEL damage - space charge

The space charge and depletion characteristics are a result of doping - deliberately introduced energy levels that lie in the band gap

- Defects can perform the same role

In silicon, bulk damage produces predominantly *acceptor sites*, effectively turning the material more p-type

- For n-type bulk silicon this means that the device undergoes **type inversion**
- The sensor will get progressively harder to fully deplete after significant doses



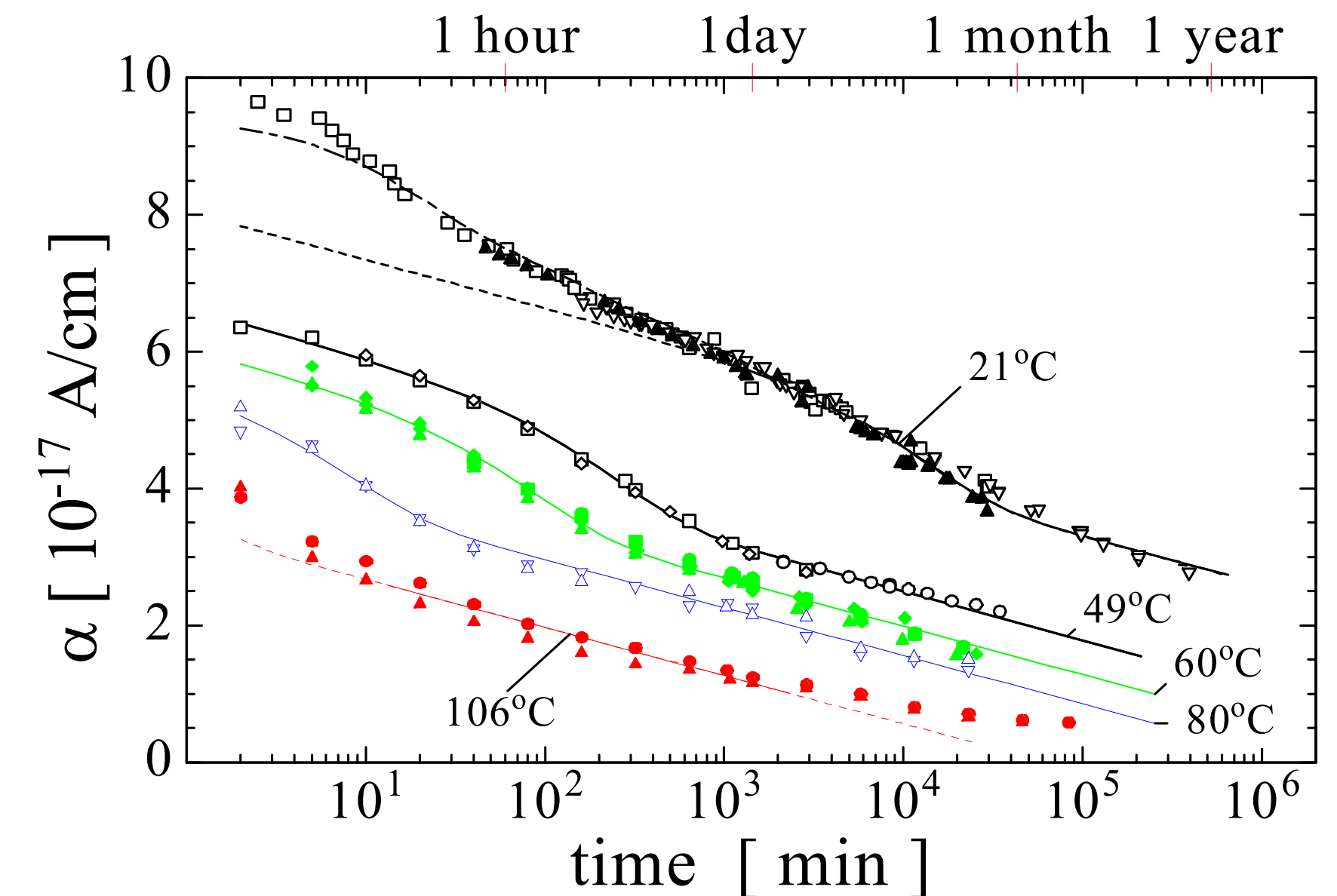
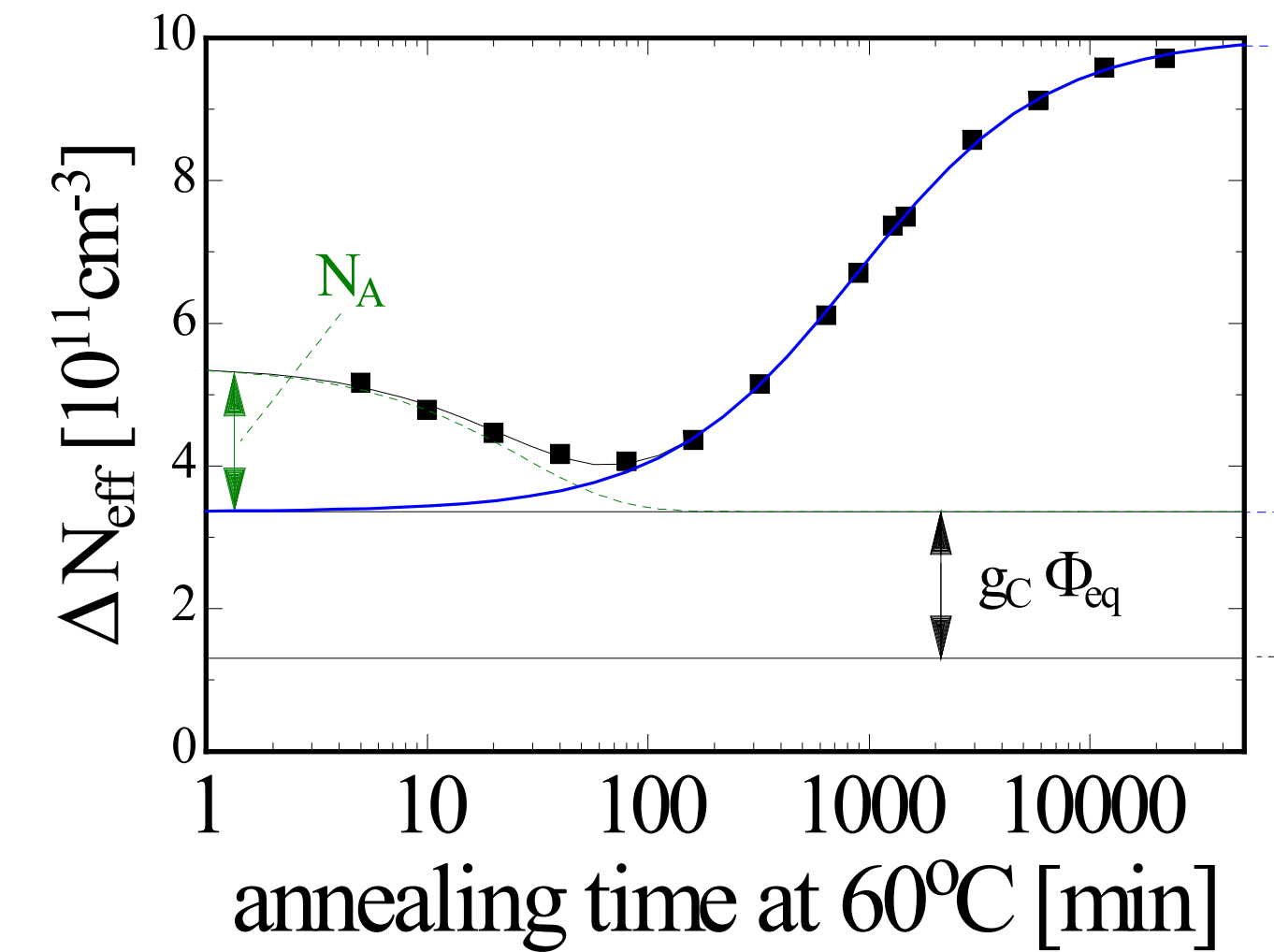
Annealing

Defects in the lattice are not stationary!

- They can move around, recombine, make friends...

Keeping materials at ambient or raised temperatures can be used to repair some of the damage

- However, not all annealing is beneficial - operate detectors cold and anneal for specific set periods



Total ionising dose - TID

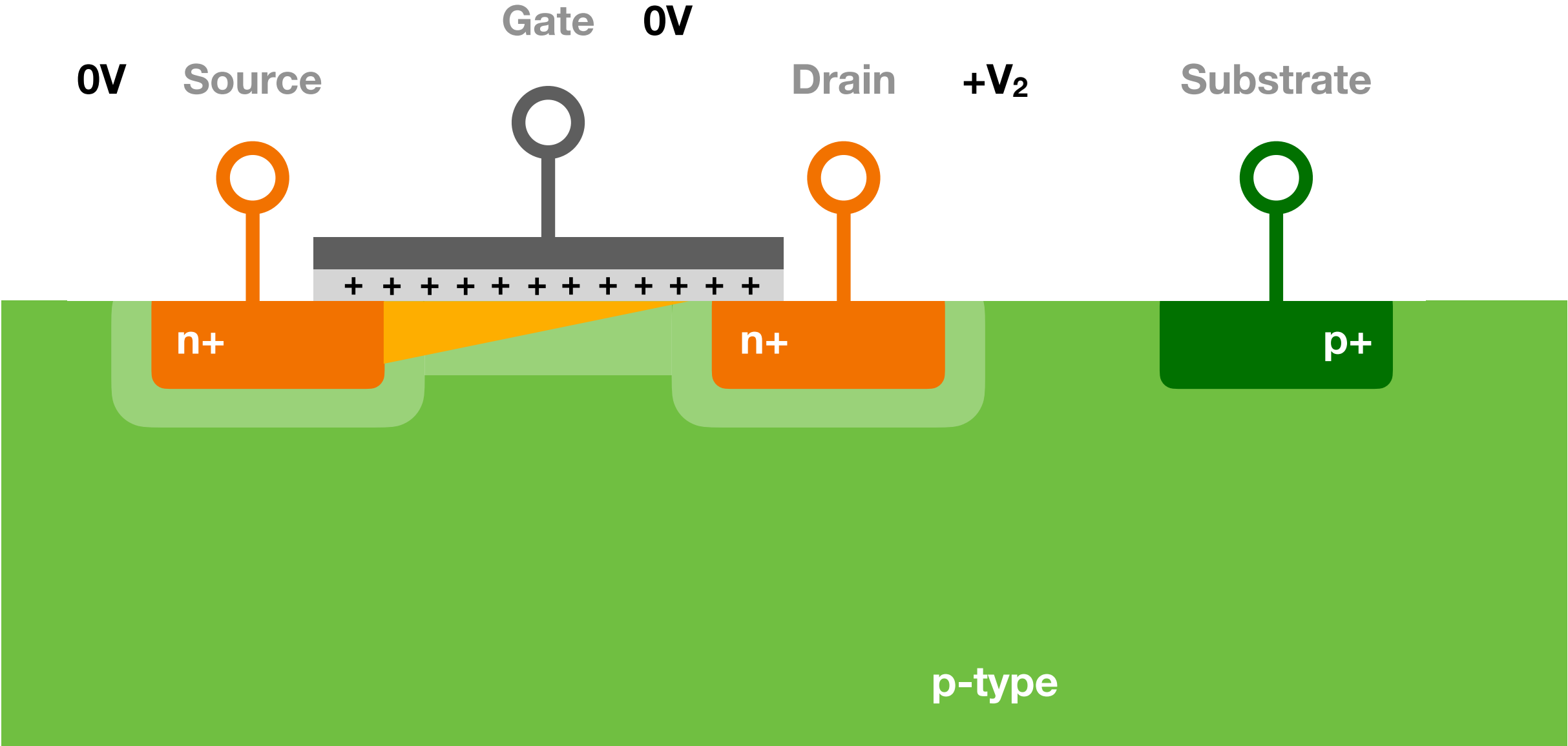
While ionising losses in the bulk constitute signal, and non-ionising losses create crystal defects, a separate issue arises with ionising dose in area susceptible to

- These are typically interfaces and surfaces
- Large effect on electronics - oxide charge build up!
- Total dose measured in *Grays* - joule/kg or *rad* (1 rad = 0.01 Gy)

Oxide charge accumulation leads to several detrimental effects

- Shift in threshold voltage
- Increase in leakage current
- Eventually transistors can't be switched on/off

TID effects



Single event upsets/effects (SEU / SEE)

It is possible for large charge deposits to directly affect transistors or other electronic blocks

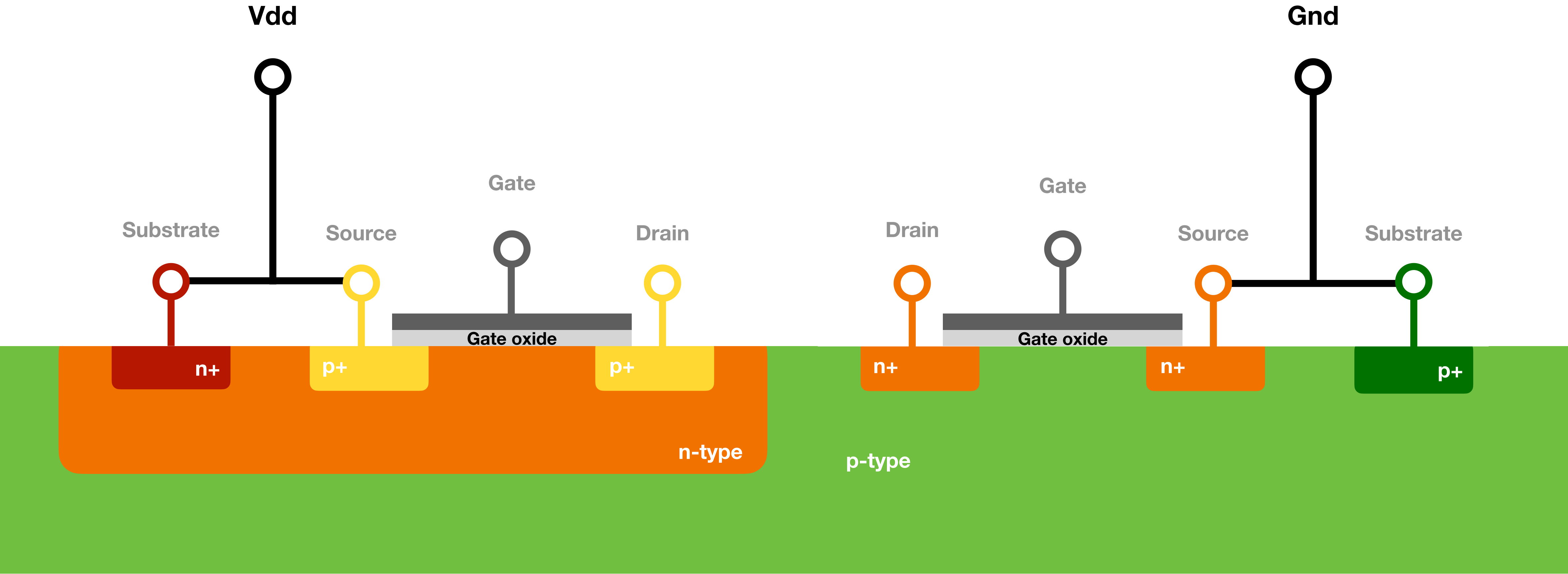
These effects can be as innocent as flipping individual bits in a register

- Triplicate important circuitry and take majority vote
- Live with it, depending on number of bits expected and impact

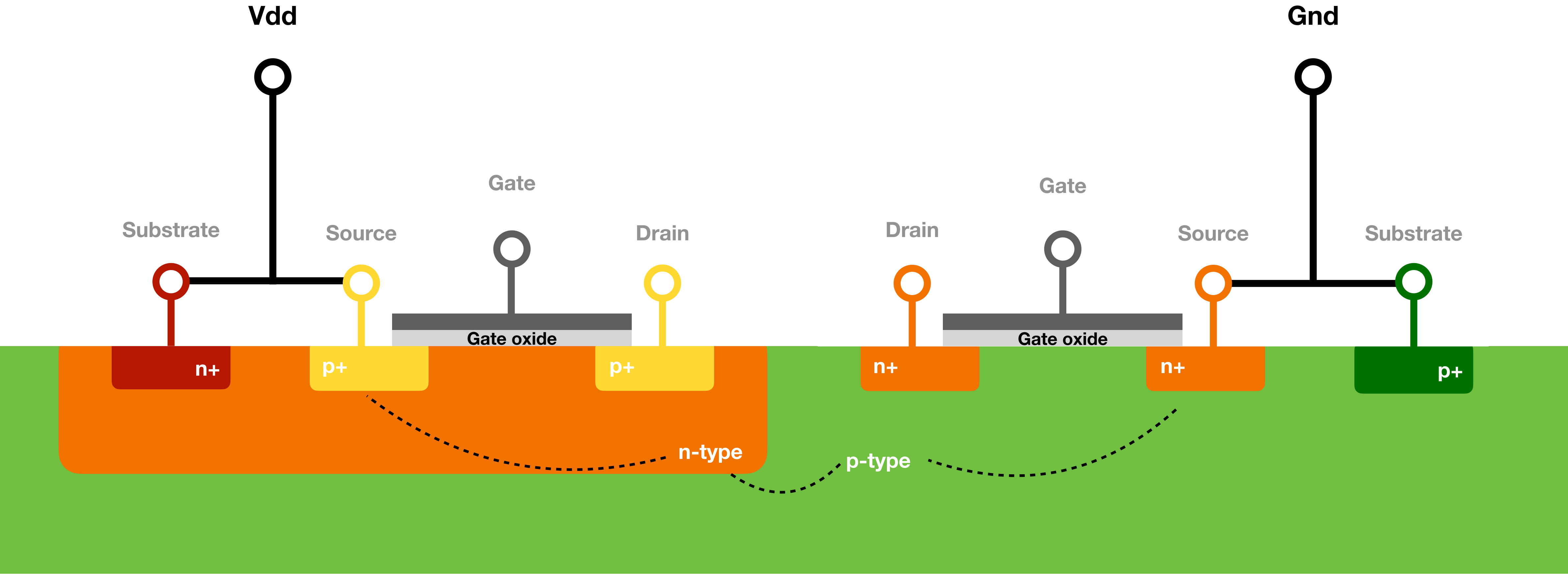
Much more dangerous cases

- Issues like latch-up can physically destroy the electronics!

Latch-up in CMOS



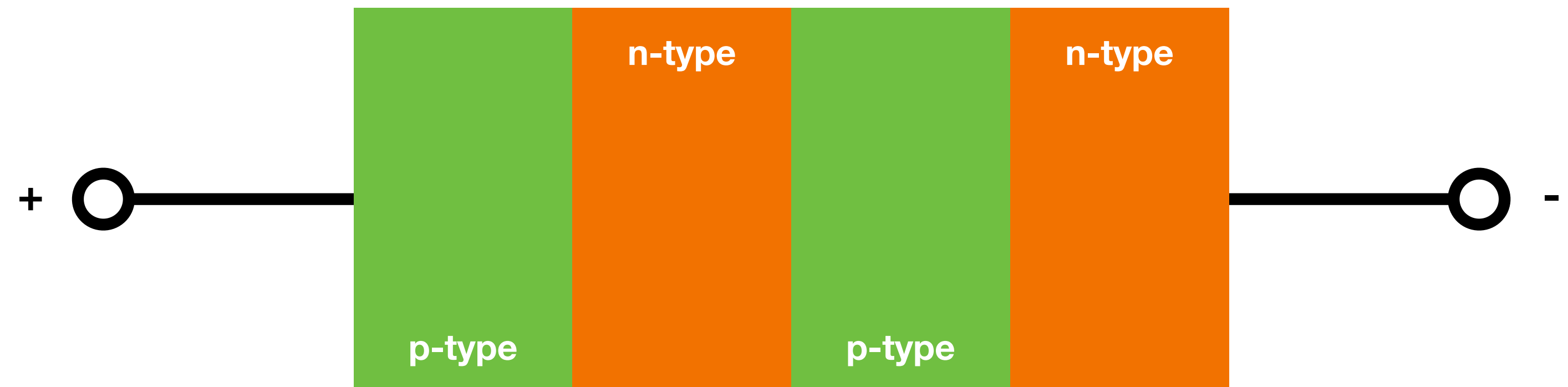
Latch-up in CMOS



Thyristor

The parasitic PNP structure in a CMOS block is equivalent to a device called a *thyristor*

- 3 PN junctions in series



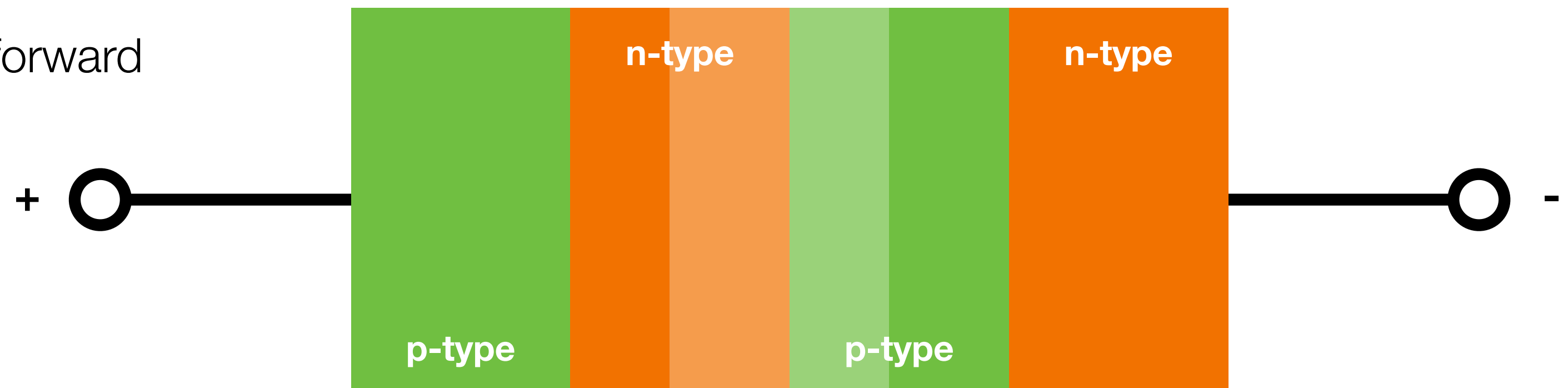
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Under these conditions, two of the junctions are forward biased and one is reverse biased

- No current flows while there is a depleted junction in the centre



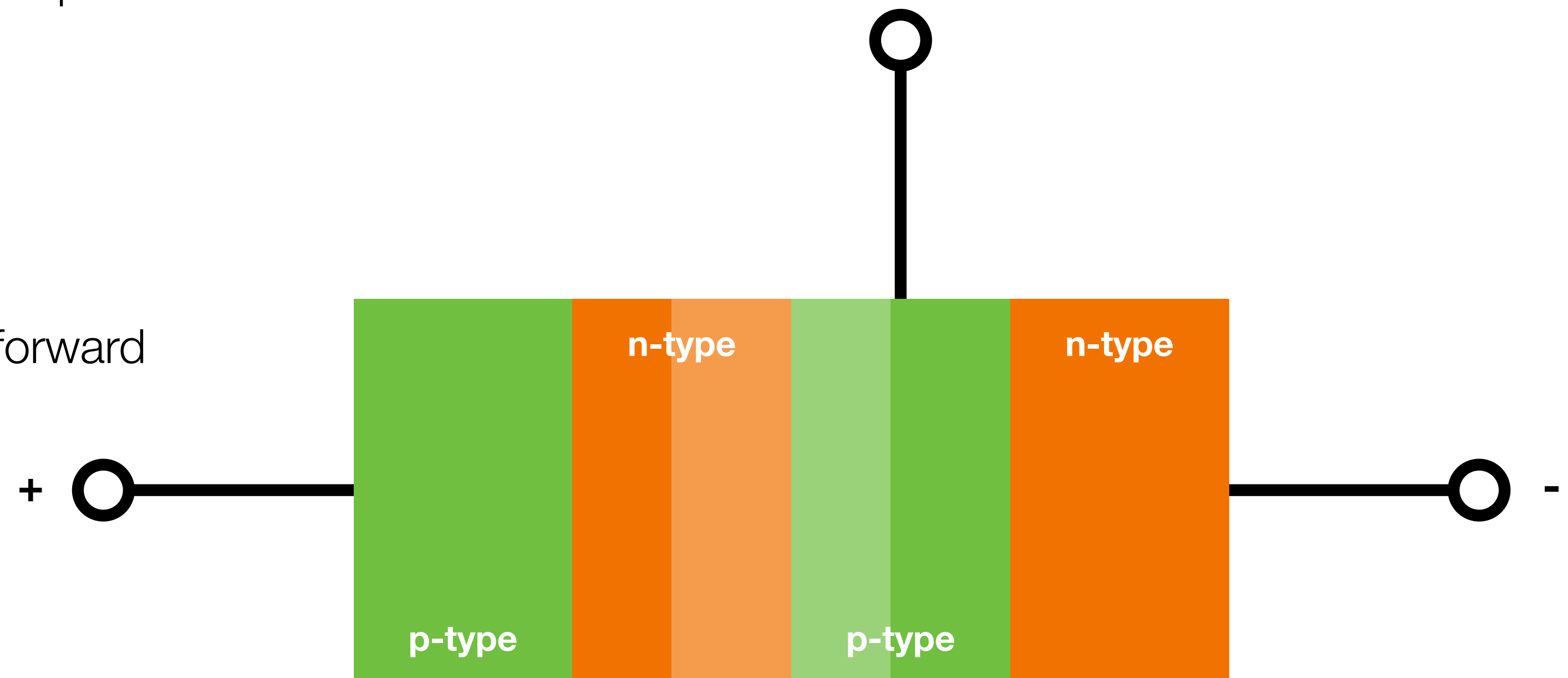
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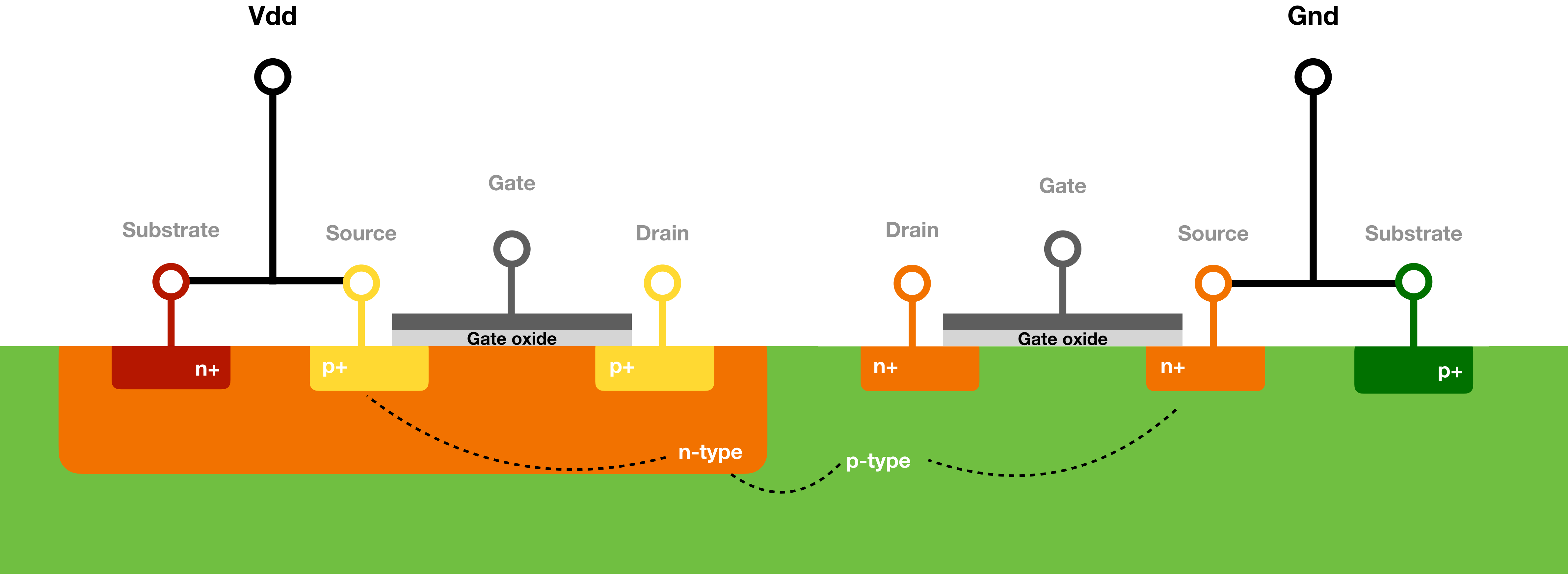
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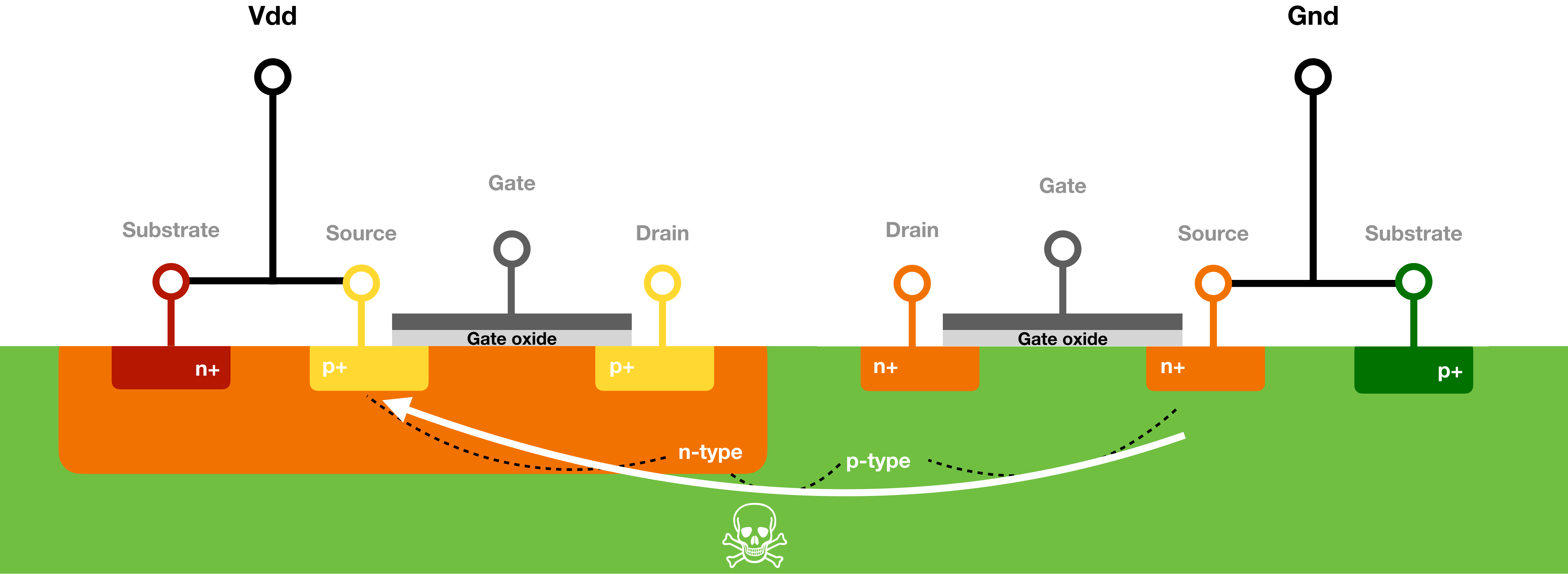
The addition of a gate voltage (or large charge deposit) can reduce the depletion region and lead to a current flow through the whole device

that will not stop while there is current flow

Latch-up in CMOS



Latch-up in CMOS



Building Things from Hot and Cold Rocks to Make Life Work Better

(Silicon fabrication)

Why silicon?

There is a simple reason as to why silicon is the dominant semiconductor used in electronics

- Silicon oxide

The ability to easily grow a native, stable, insulating oxide is what opens up the whole area to exploitation. Next steps are:

- Production of large, single-crystal wafers with low contaminants
- Doping, both n- and p-type
- Electrical contacts (particularly ohmic)

With many/most other semiconductors these are the challenges that need to be overcome before reliable detectors can be manufactured. One then needs to **scale up** (easier said than done!)

Wafer production - the Czochralski method

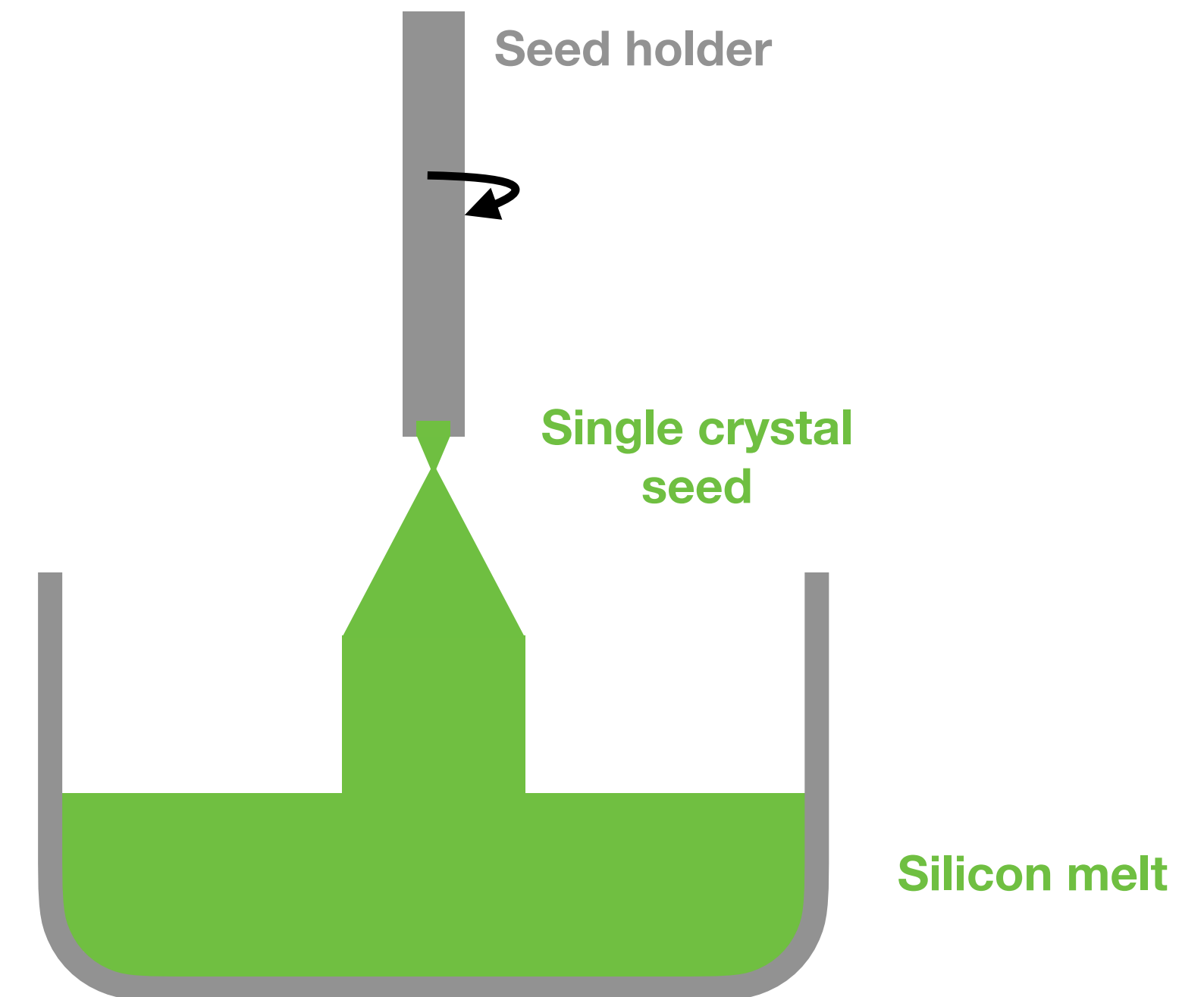
Wafer production still follows the method proposed in 1915 by Czochralski (discovered by accident)

- Very pure silica is placed in a crucible and dissolved to form a melt
- A single crystal seed is dipped into the melt and slowly extracted, while rotating
- The resulting ingot is a single crystal following the same orientation as the seed crystal

This happens at high temperature - 1425°C - and for large diameters

- Commercial processing now done on 12" wafers

Nonetheless, impurities from the crucible (such as O₂) can be present at the level of 10¹⁸ cm⁻³



Wafer production - Float Zone processing

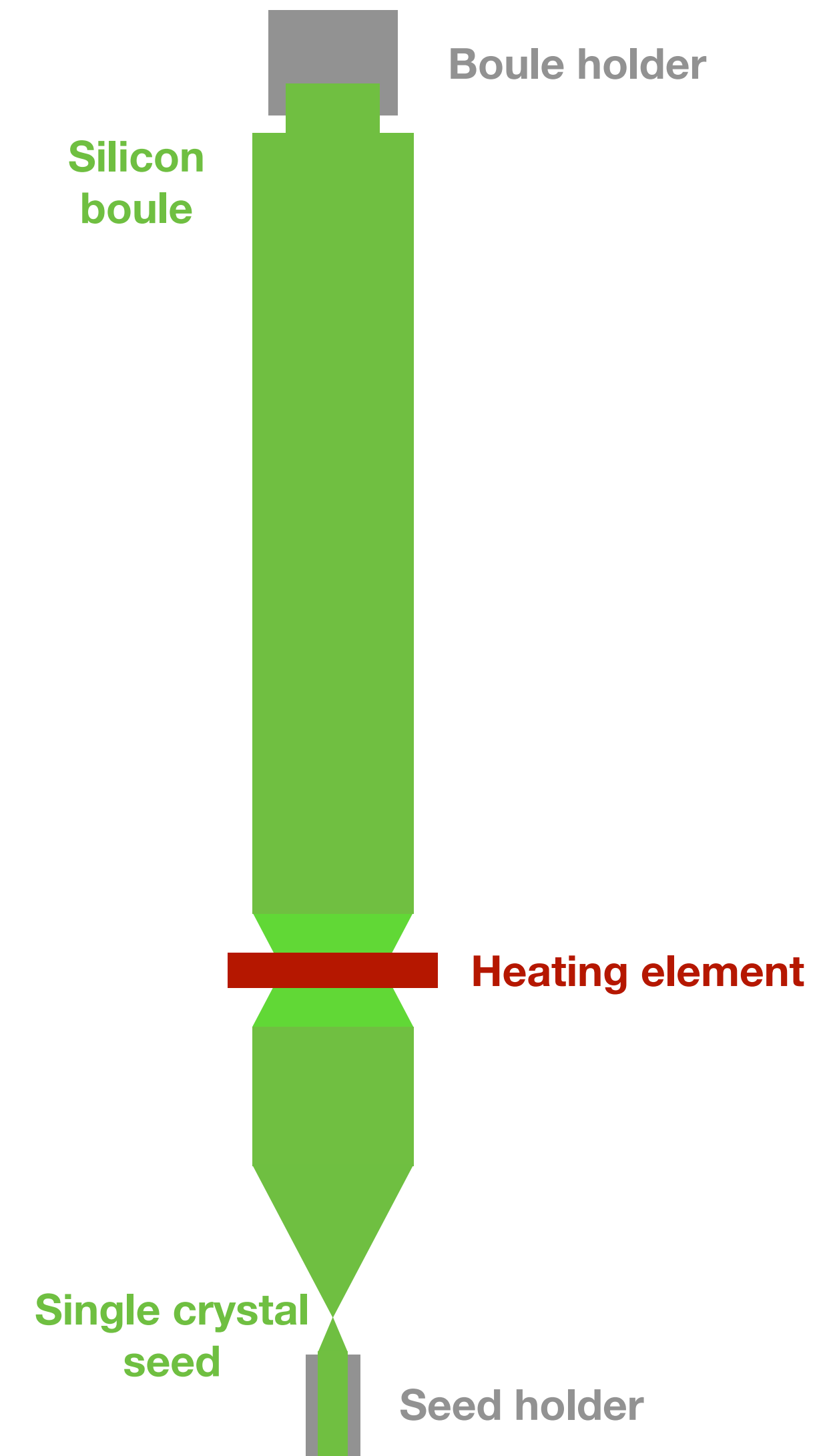
For much purer silicon wafers, float zone processing can be used

A polycrystalline silicon ingot is melted in a narrow band, which is scanned from one end to the other

- Contaminants remain in the melted region rather than crystallising into the now very-pure silicon lattice
- They are filtered all the way to the end of the ingot and then discarded

Produces high-resistivity wafers ideal for detectors

- $n_{\text{eff}} \ll 10^{13} \text{ cm}^{-3}$



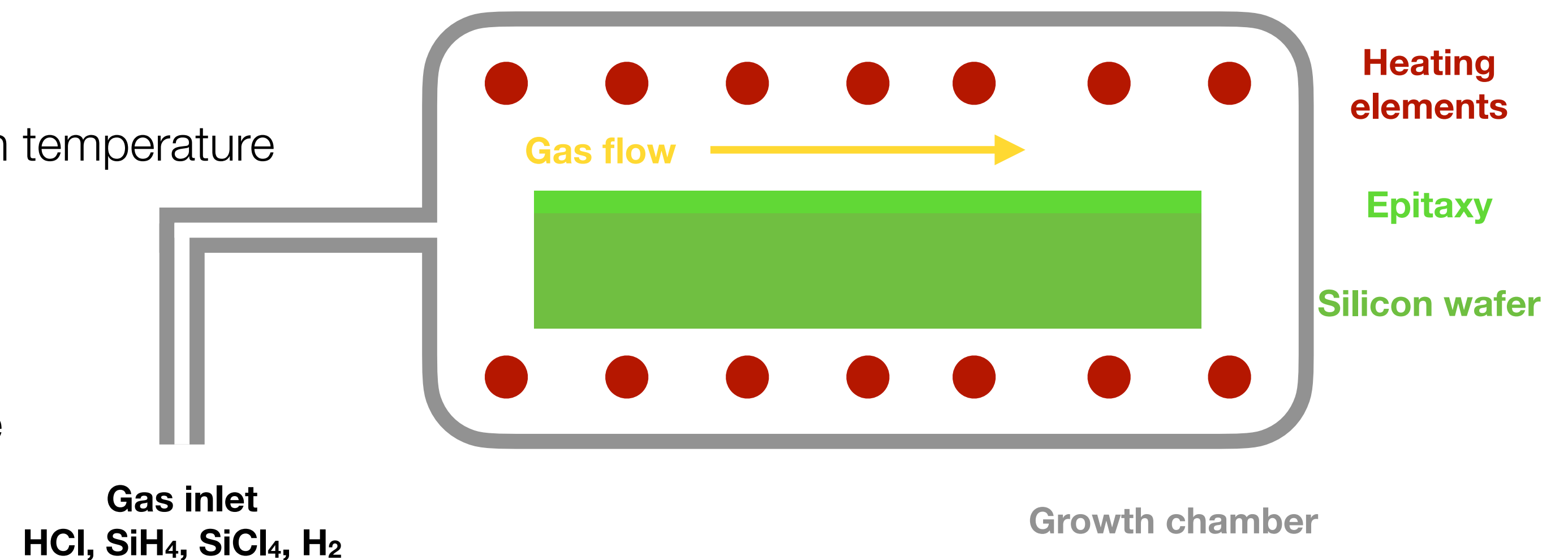
Epitaxial growth

Another method for obtaining a pure silicon crystal with high resistivity is to grow a dedicated pure layer on top of an existing silicon wafer

- This layer is referred to as an *epitaxial layer*

The wafers are placed in a chamber heated to high temperature (800 - 1150°C)

- Wafer surface must be atomically clean
- Gases containing silicon introduced into the chamber
- Growth of 10s of nm per minute

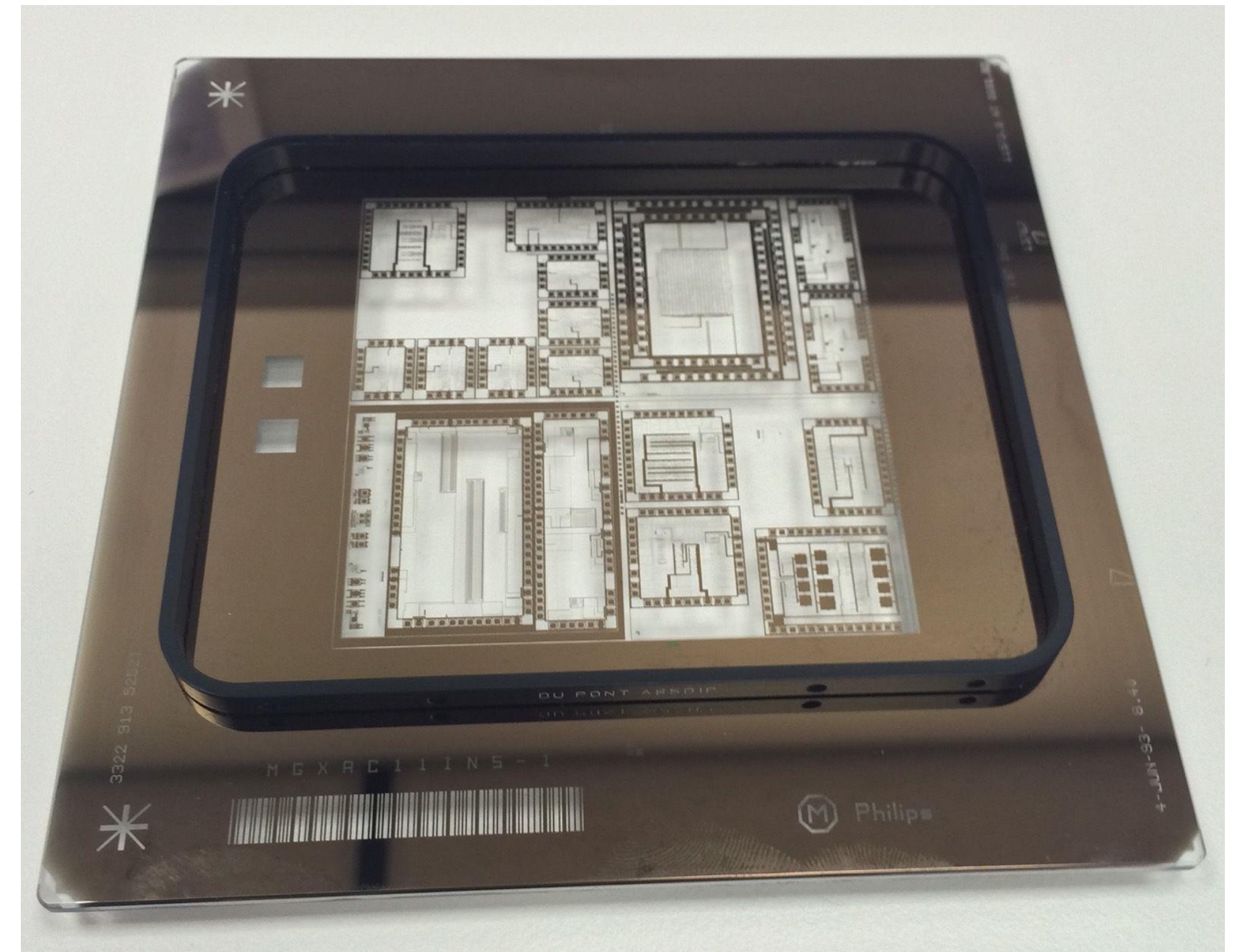


Photolithography

We want to be able to *pattern* our devices => deposit patterned masks

One of the most straightforward ways to do this is via photolithography

- A coating is applied to the full silicon wafer - called a *resist*
- A mask is produced containing the desired pattern
- The pattern is transferred to the resist either by placing the mask in contact and illuminating it, or by projecting light through the mask via a lens
- The resist will either harden or break down, depending on the material chosen

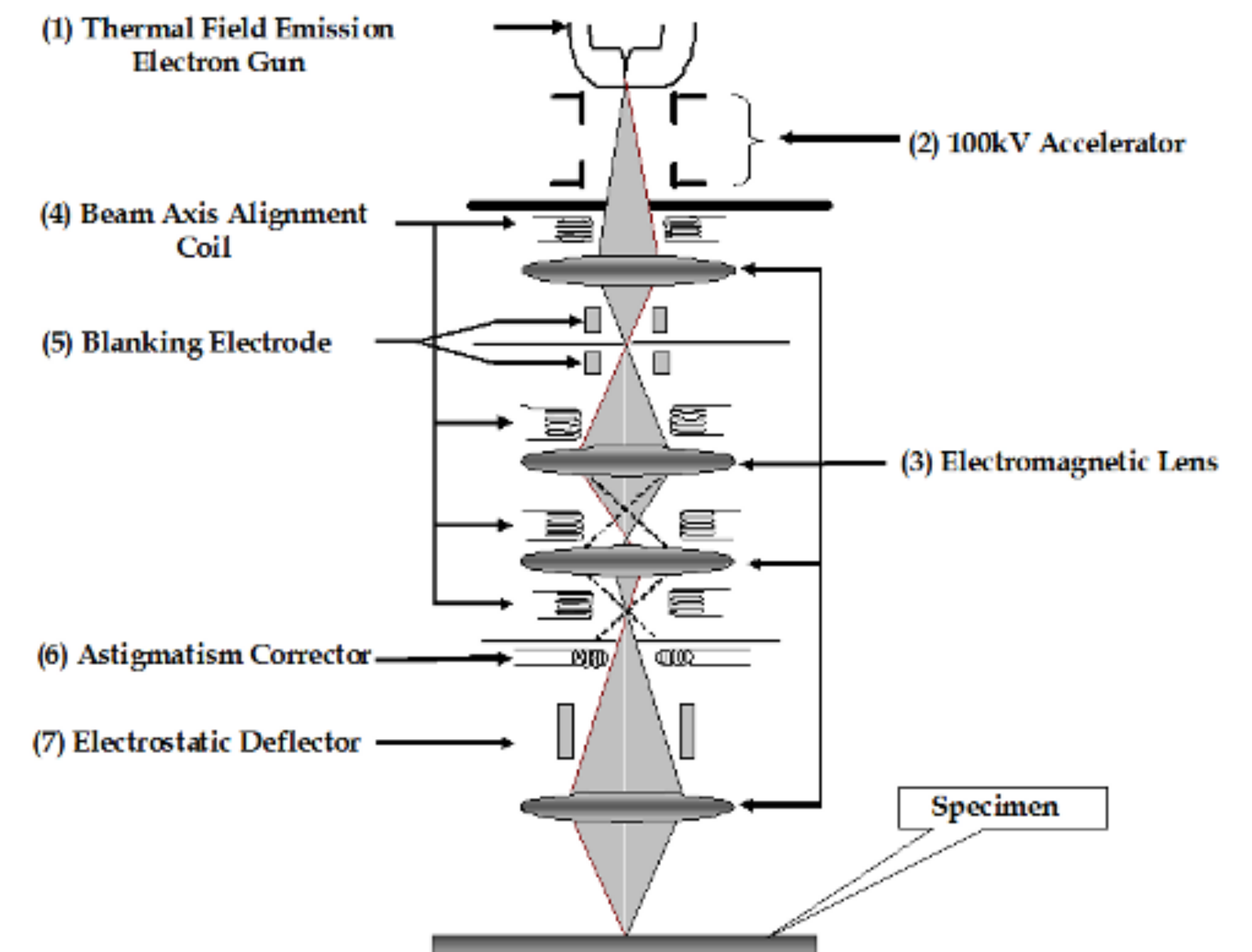


E-beam lithography

Photolithography depends on light and so starts to limit the feature size around the wavelengths used

- For much finer objects use a focussed beam of electrons, scanned across the surface in the desired pattern
- No mask, just direct scanning

Resolutions can be much much better - down to below 10 nm - but at the cost of limited throughput

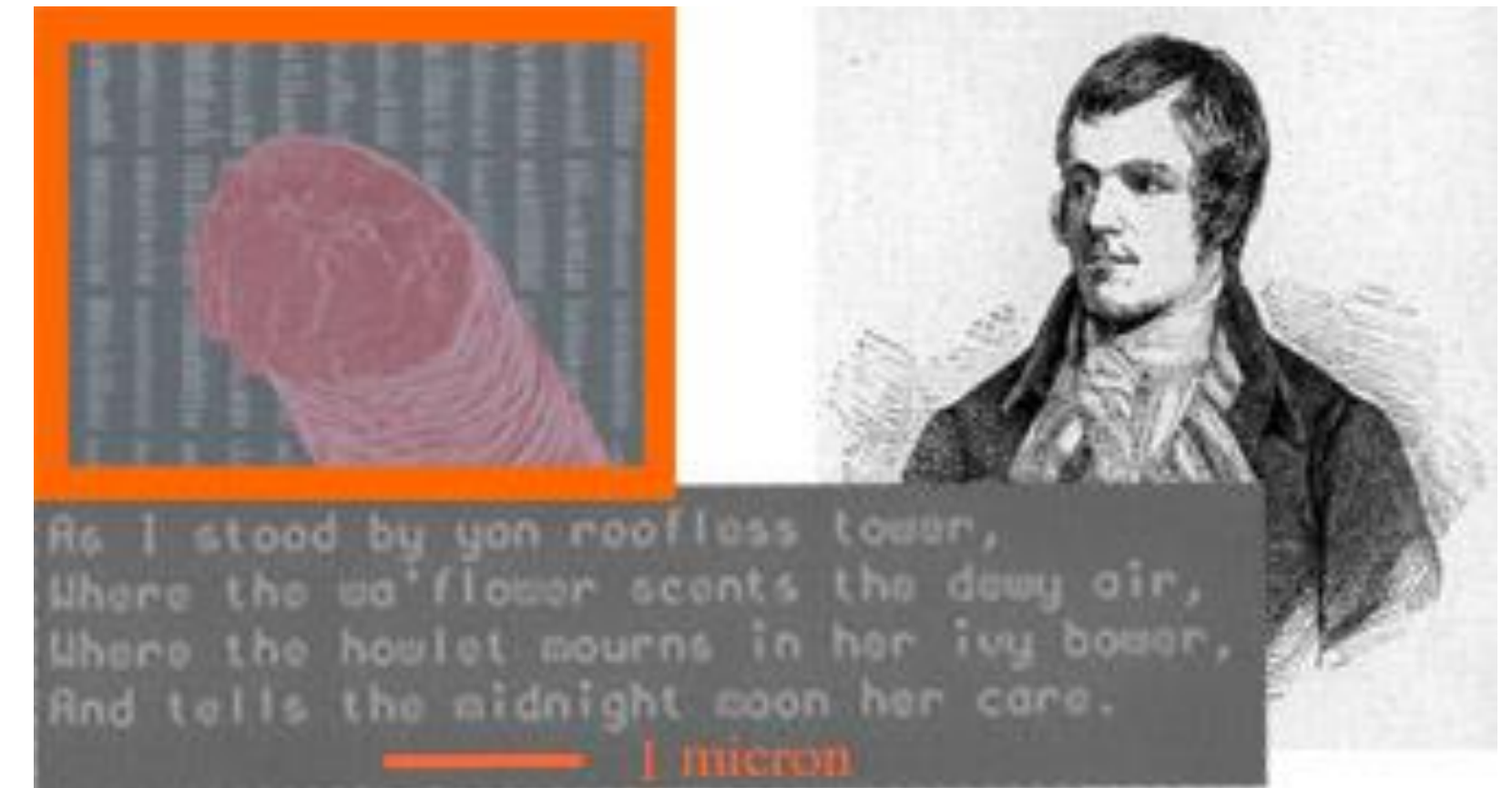


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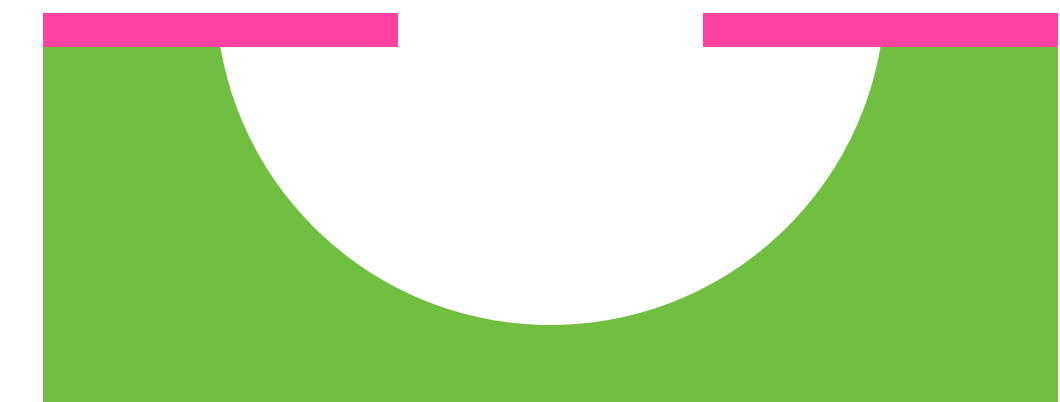
A book of the Complete Works of Robert Burns, Scotland's National Bard, has approximately 480 pages. To publicise the capability of the new Vistec VB6 UHR EWF electron beam lithography tool at the JWNC, we used it to write the Complete Works of Robert Burns on a small piece of silicon. Ten copies would fit on the head of a pin and this is likely to be the world's smallest copy of the works of Burns. The image shows pages of text alongside a human hair plus detailed text from the song "As I stood by yon roofless tower". Each character is approximately 150 nm

Etching

There are two kinds of etching, both with their own advantages/disadvantages: **wet** and **dry**

Wet etching is relatively straightforward

- Immerse your wafer in something hideously corrosive/toxic/carcinogenic until the bits that you don't want have been removed...
- Etching will be isotropic
- Not all materials have suitable chemistry for wet etching (NB. HF used for SiO_2)



Wet etching
Isotropic

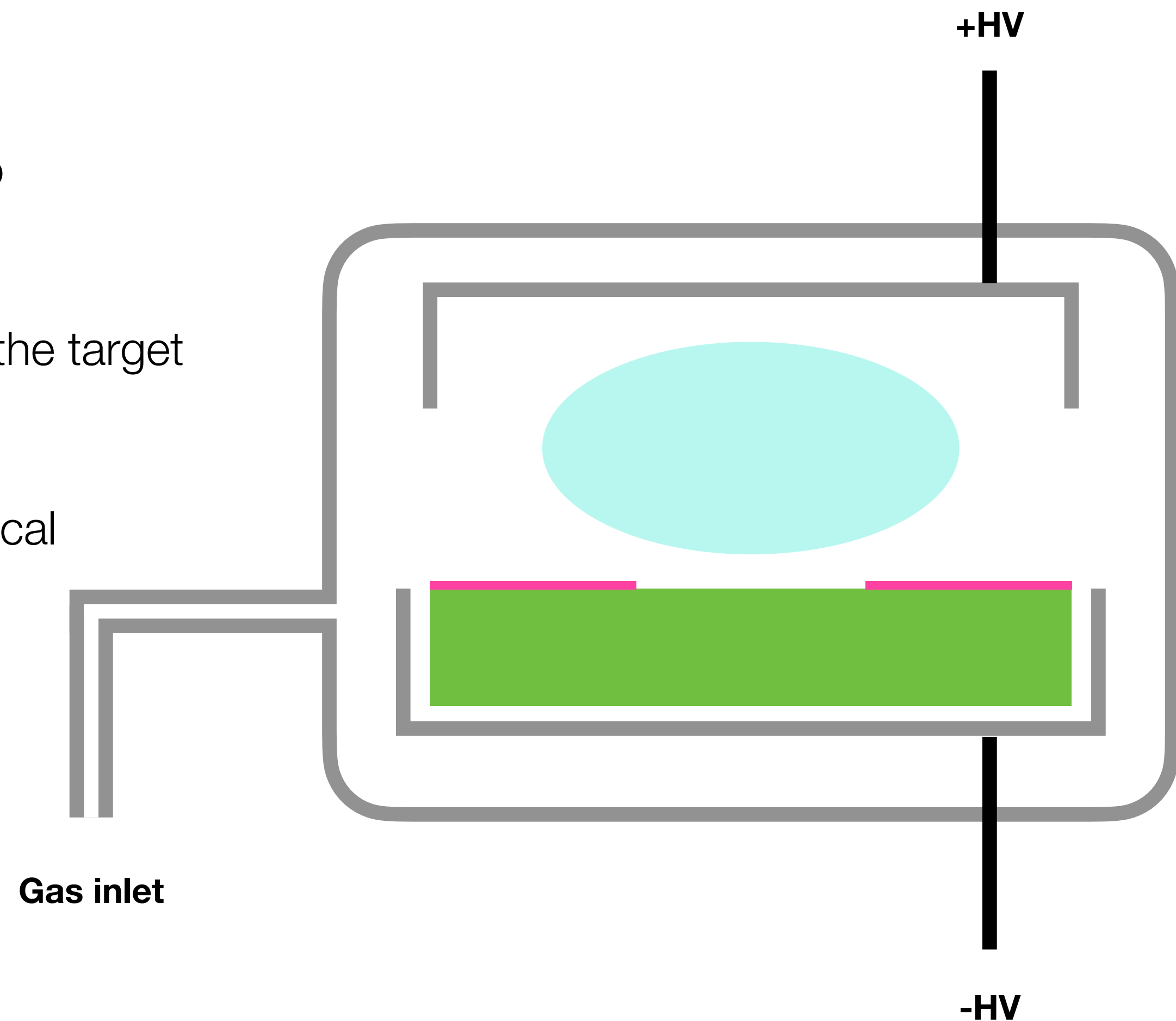


Dry etching
Anisotropic

Etching

Dry etching uses pressurised gas and a high voltage to generate a plasma

- Once ions are created they accelerate towards the target wafer
- A combination of reactant species and mechanical impacts etches anisotropically
- Reactive Ion Etching (RIE) a common technique
- Most materials can be dry etched
- Useful to achieve high aspect ratios



DRIE etching

DRIE etching - the Bosch process - patented in 1994

Start with our silicon wafer, with an area of resist (often silicon oxide)

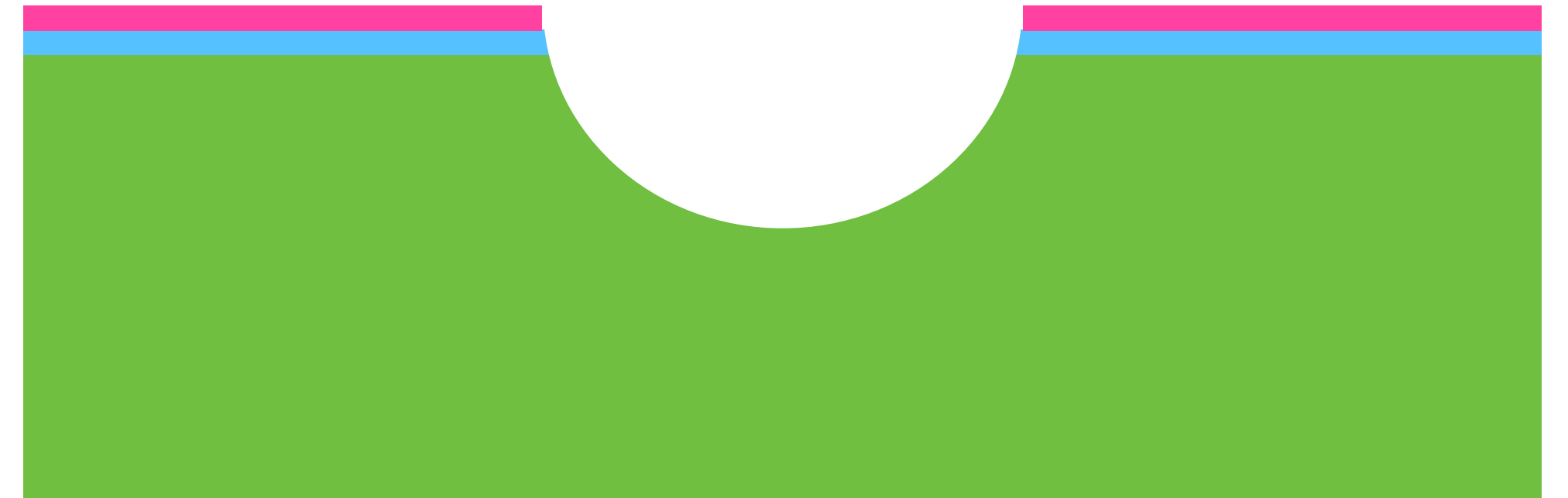


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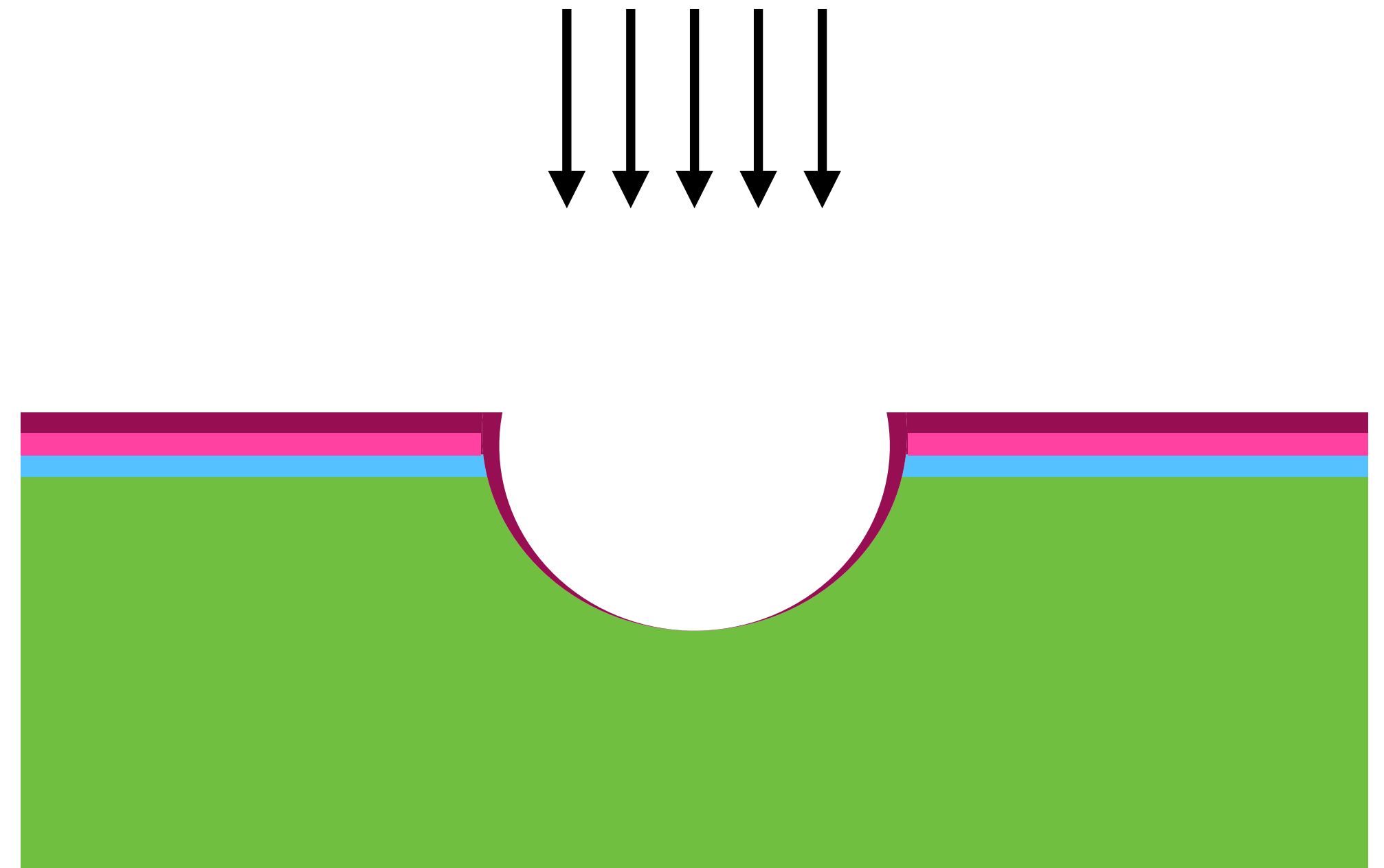


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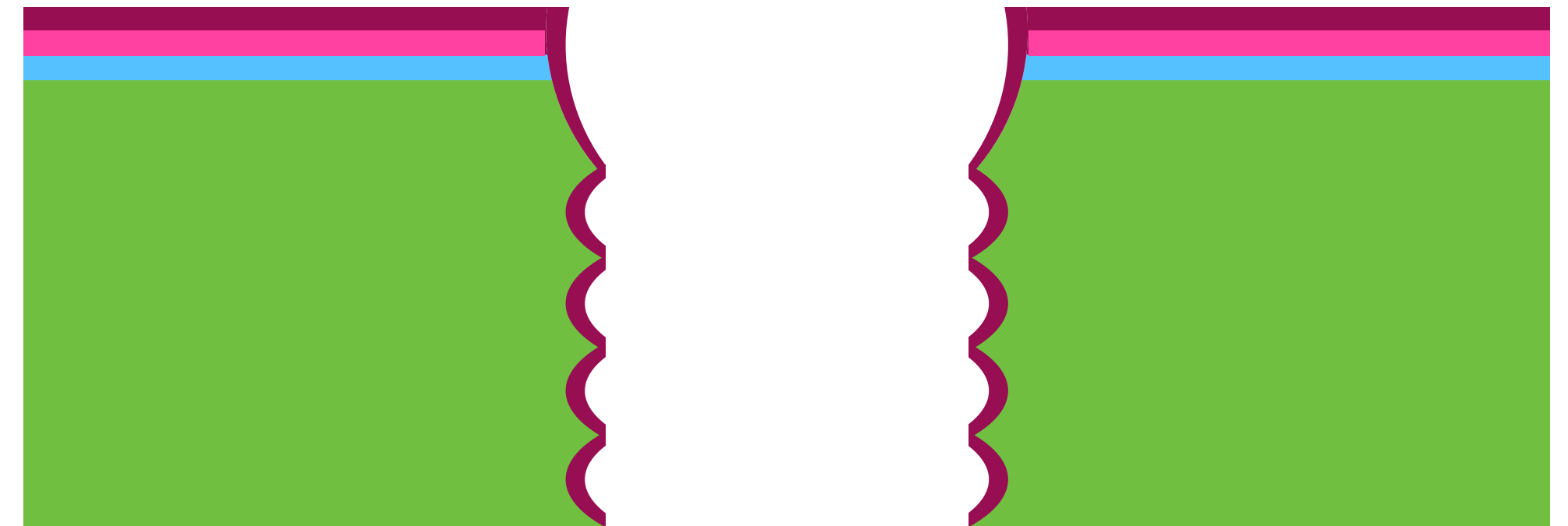


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- Repeat this cycle hundreds or thousands of times!

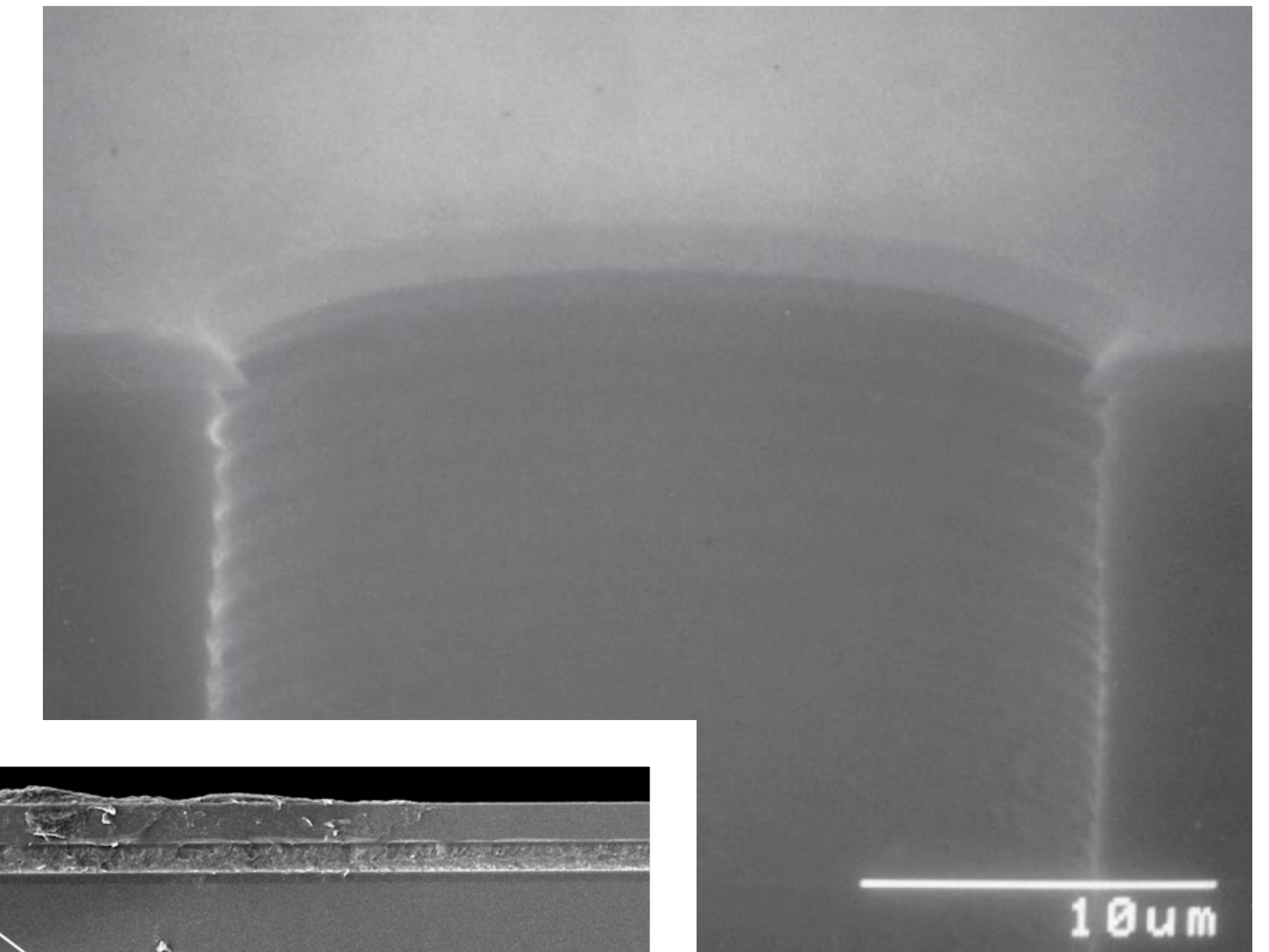
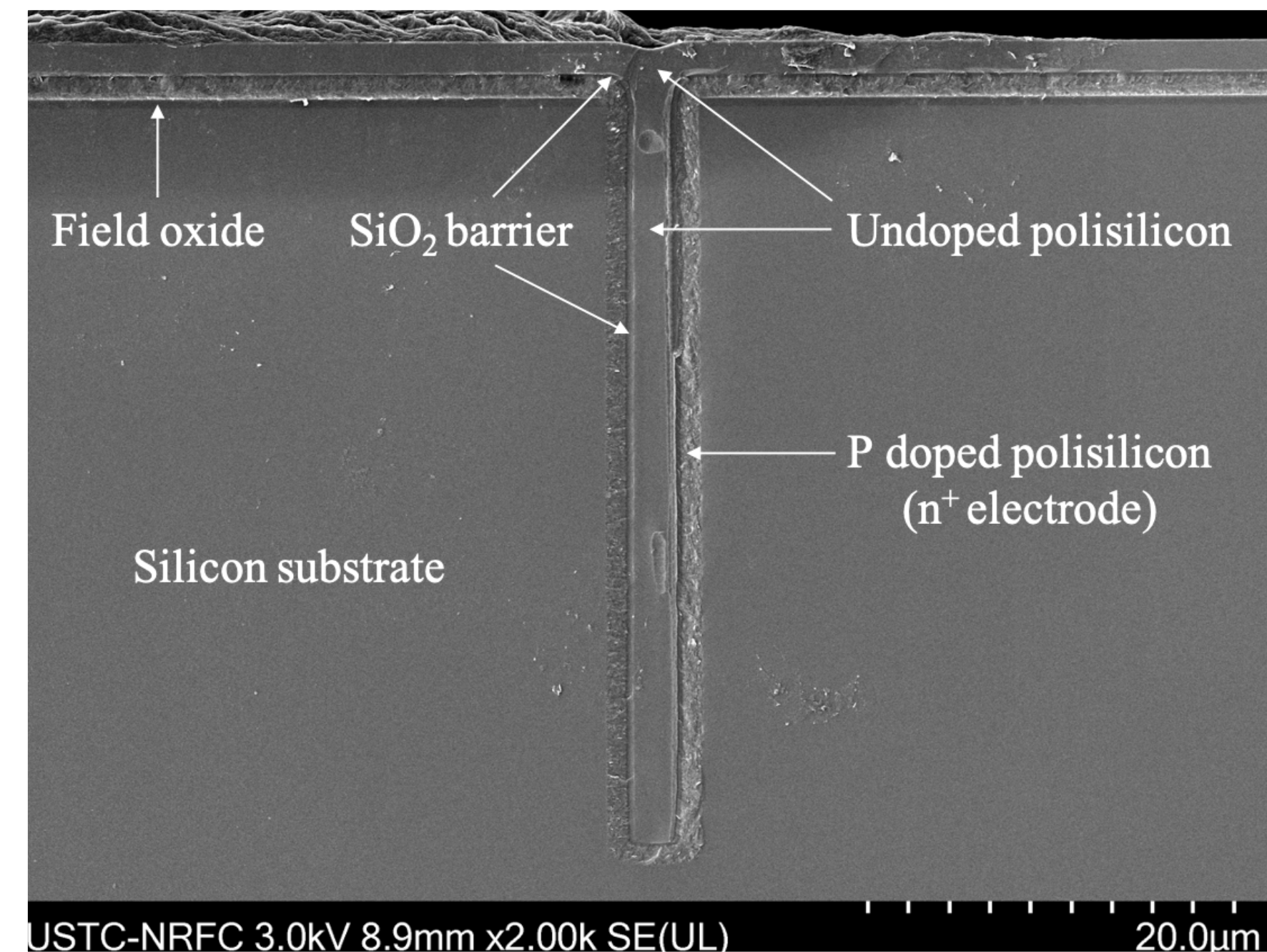


DRIE etching

Widely used in MEMS industry - pressure sensors, triggers, mechanical structures

- Additionally allows for *Through Silicon Vias* (TSVs) - connections from the back-side of a chip to the front-side

Etching rates of $\sim 10 \mu\text{m}$ / minute and aspect ratios of (cutting edge) 70:1 !



Oxidation

Oxidation of the silicon performs multiple roles

- A stable protection layer over the silicon
- An insulating layer for building vertical structures
- *An isolation layer between neighbouring transistors*

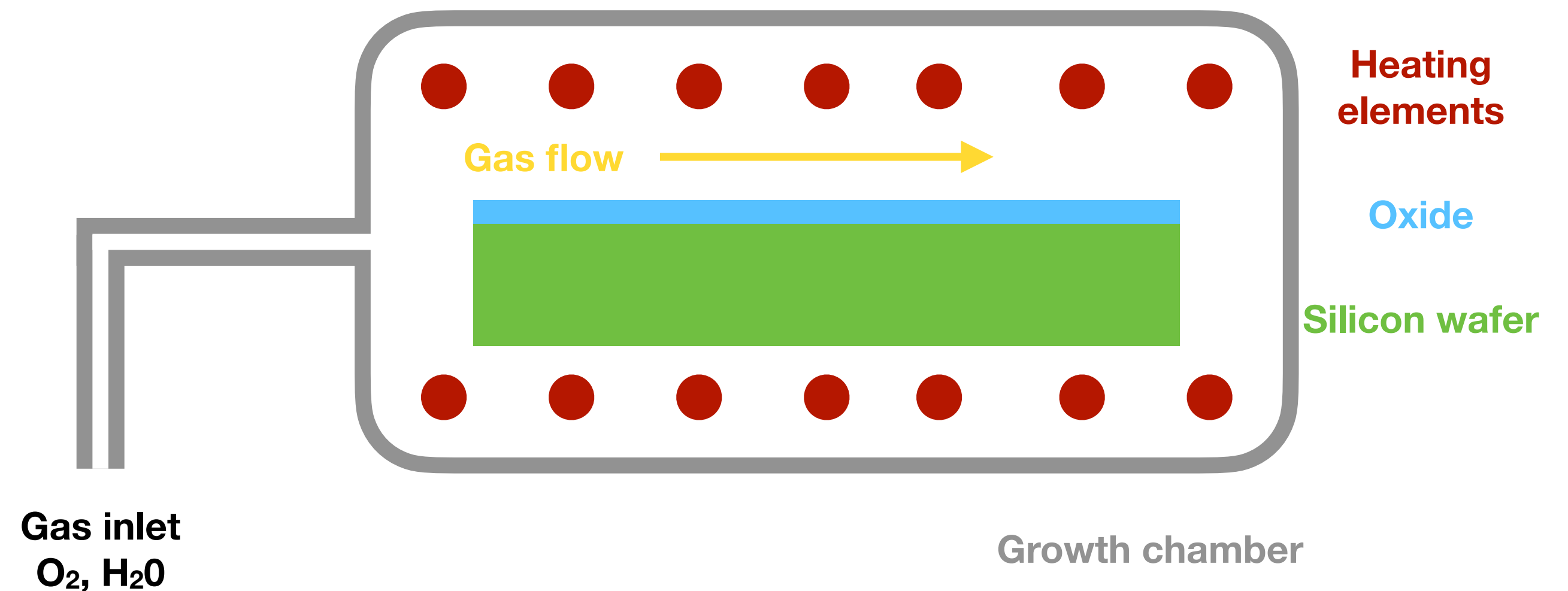
Oxidation - thermal growth

Heating wafers in a chamber and introducing gases is a versatile technique...

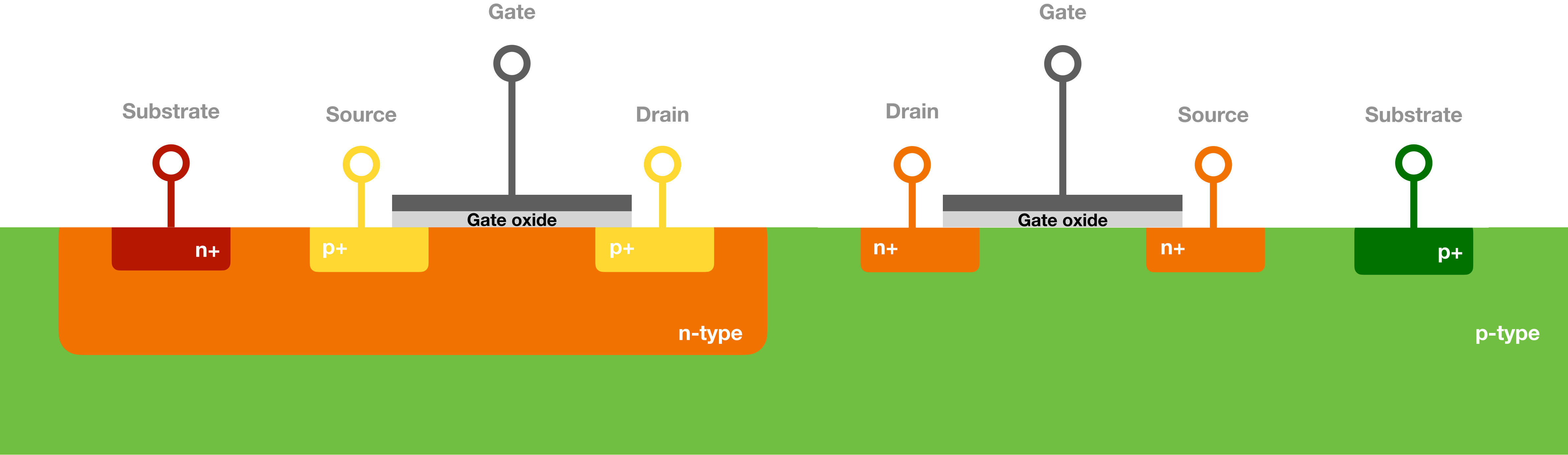
For thermal growth of oxide layers we can do this in two ways:

- **Dry oxidation.** Better quality oxide produced, but slow growth rate ~ 10 nm/hour
- **Wet oxidation.** Lower quality oxide, but 10x growth rate

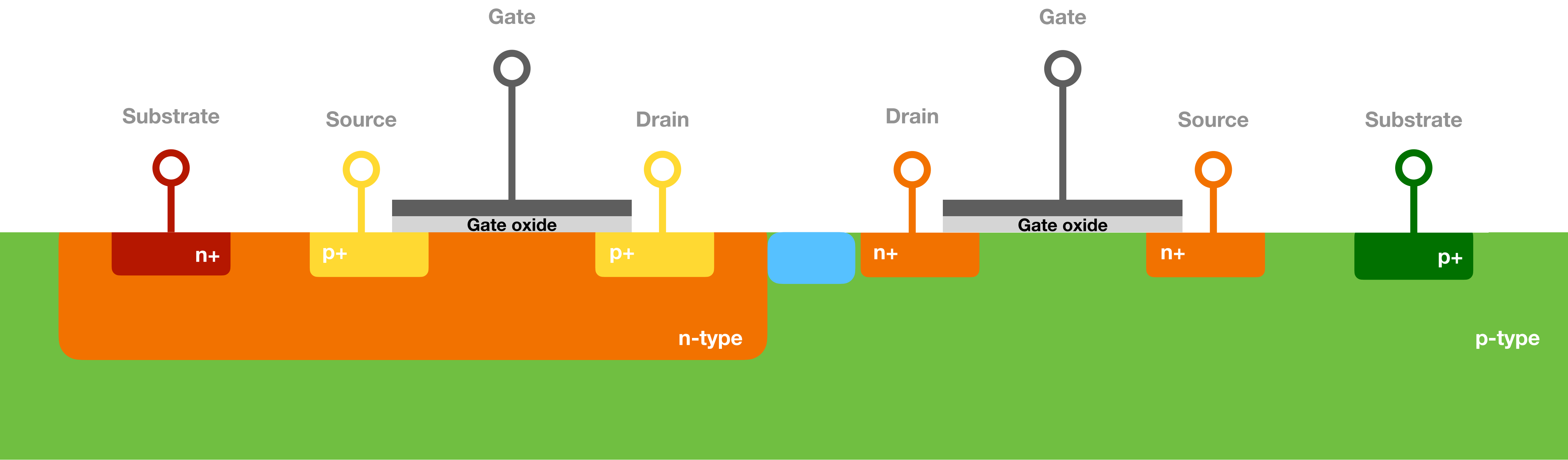
However...



Oxidation for isolation



Oxidation for isolation

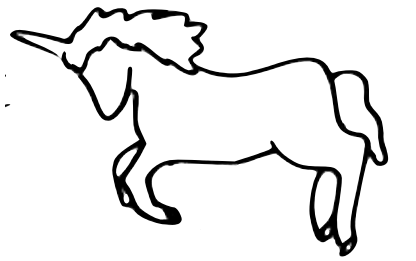
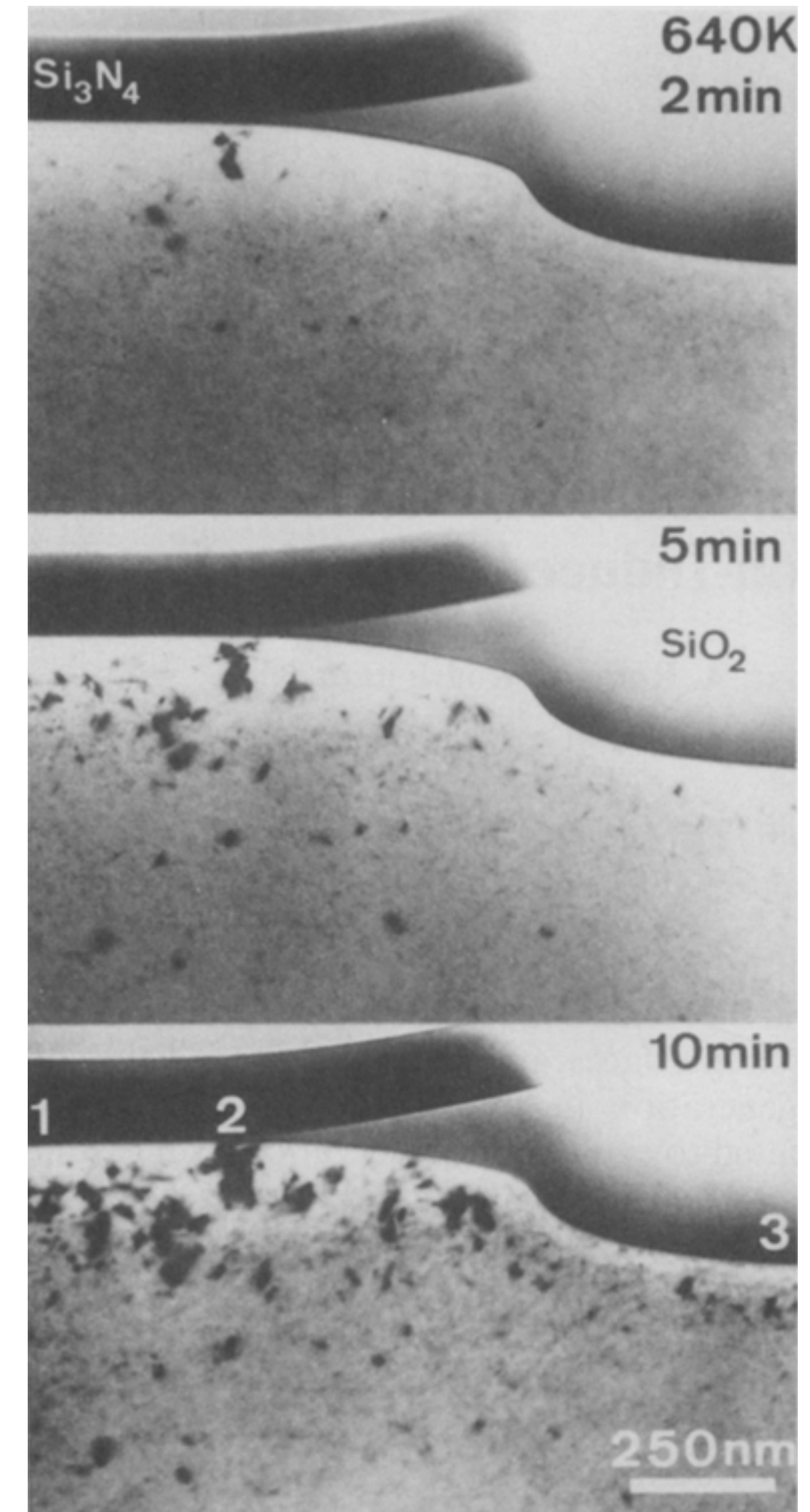


Oxidation - thermal growth

When trying to isolate wells from each other, these shallow trench isolations require filling a large volume with oxide

- The oxide will push up other layers that are on top, and start to wedge deeper along the silicon surface
- So called “bird’s beak”

Things aren’t as simple as we would like them to be...



Oxidation - shallow trench isolation (STI)

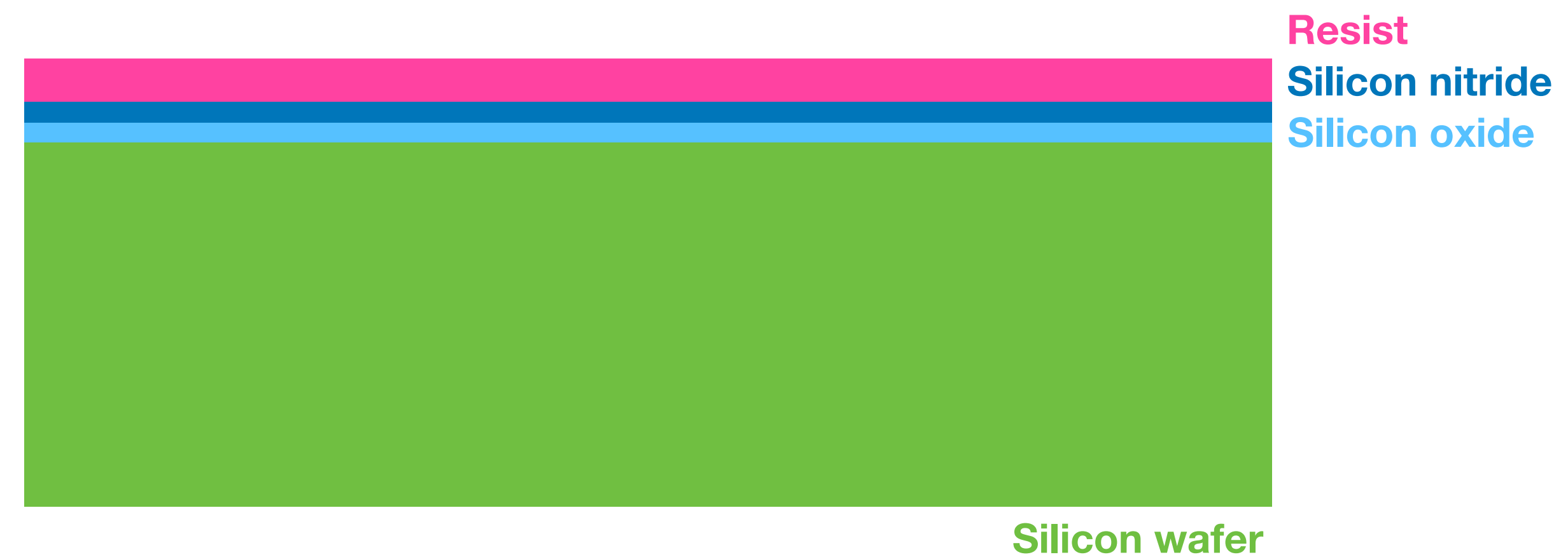
We start with a silicon wafer, with oxide and nitride layers grown on top



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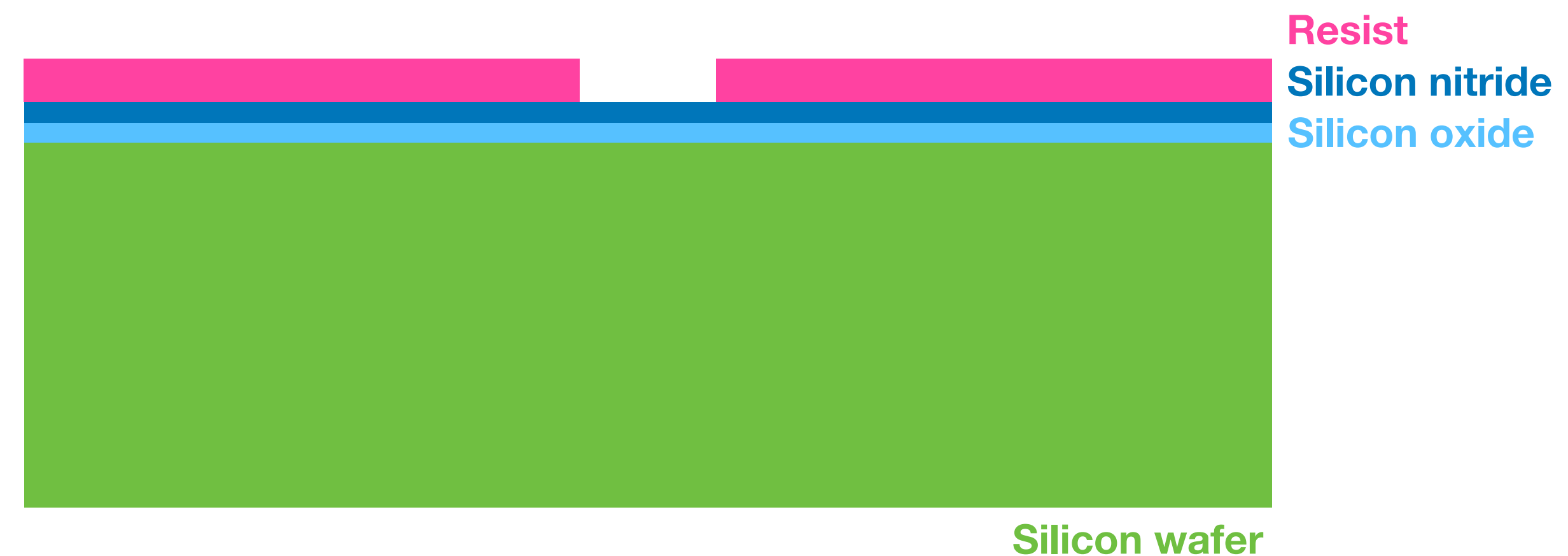
- A resist is grown to provide the etching pattern



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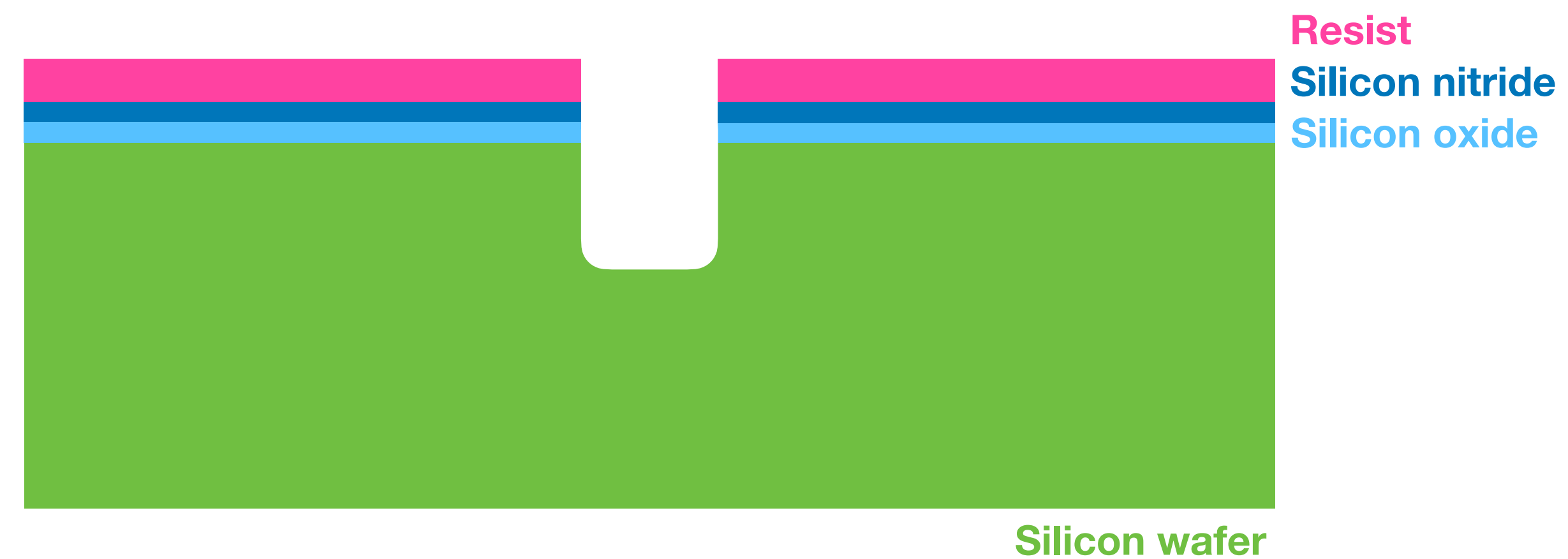
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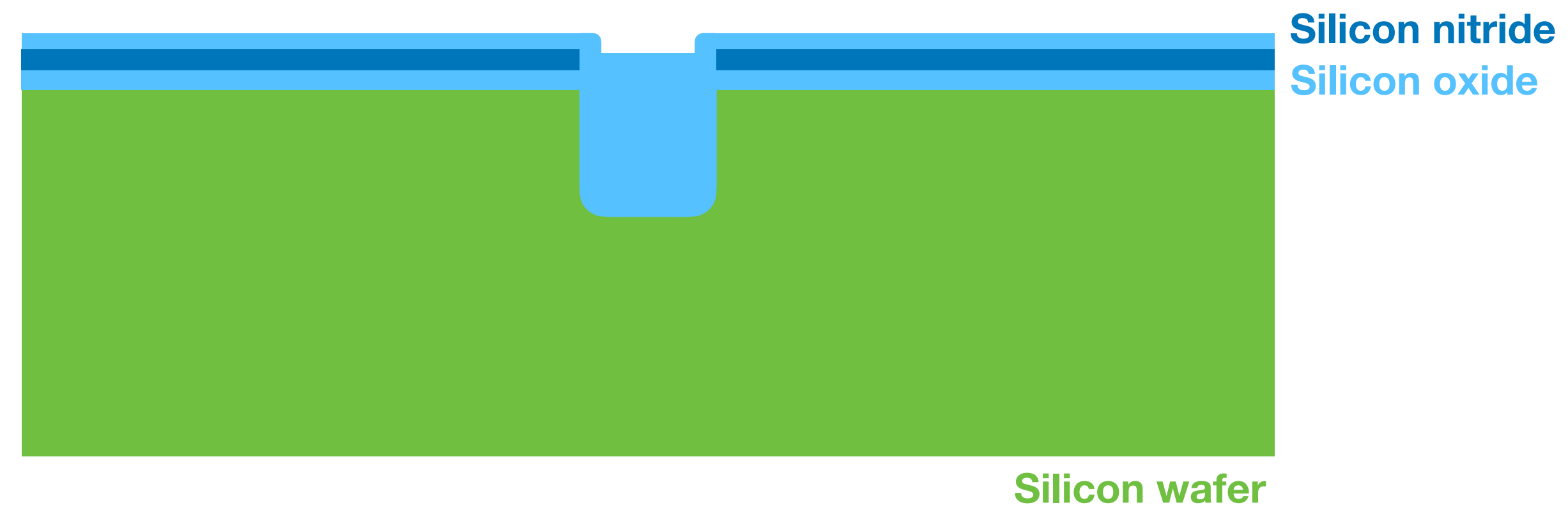
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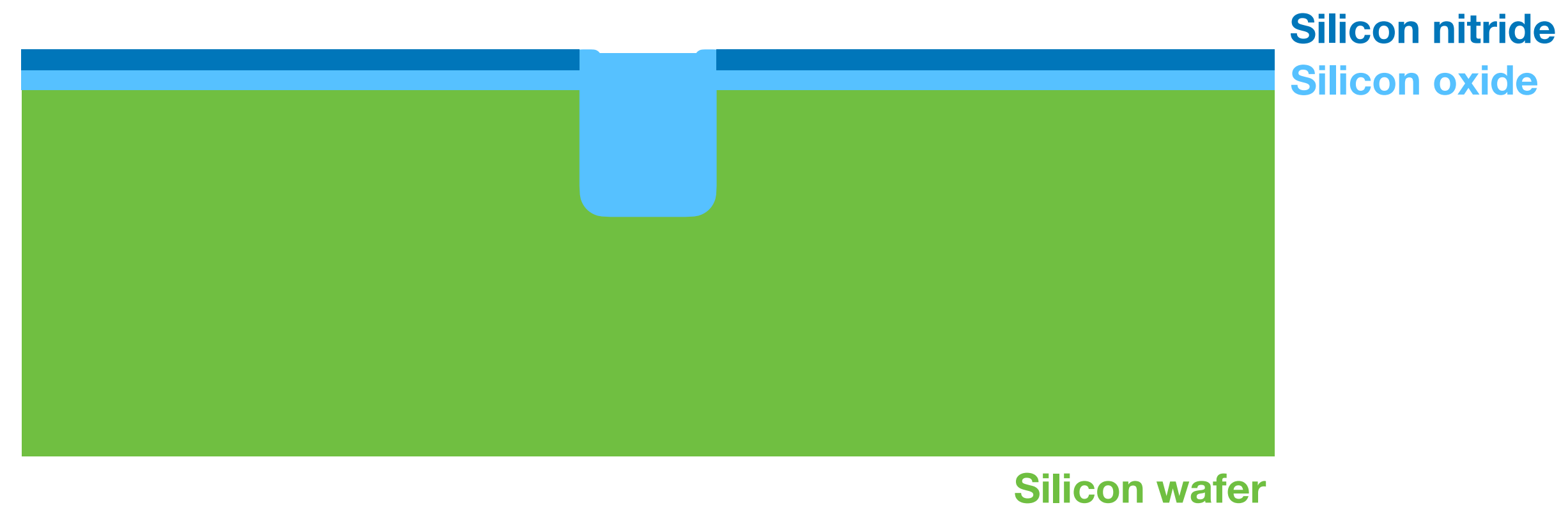
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- **CVD** is used to deposit silicon oxide everywhere



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- A **thermal oxide** is grown
- **CVD** is used to deposit silicon oxide everywhere
- The oxide is polished down
- The nitride is removed



Doping

There are a few methods to dope silicon wafers, from those during growth to gas diffusion in a chamber at high temperature

One particularly useful method is *ion implantation*

- A small particle accelerator is used to create and accelerate dopant ion species to max ~few MeV
- Spectrometer gives high purity beam
- Ions directed onto wafer surface BUT penetration affected by lattice structure

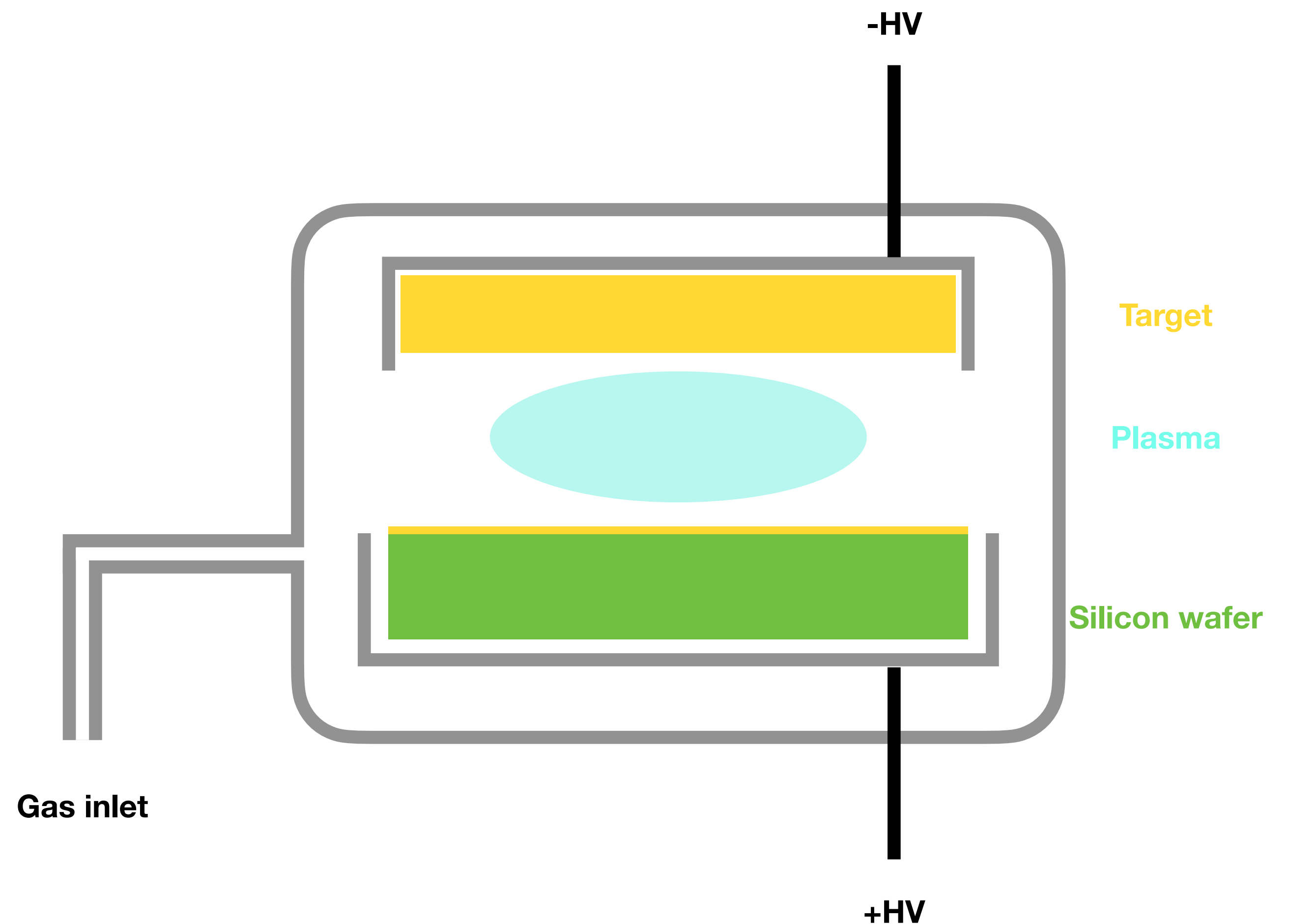


Metallisation

Metal layers can be deposited via CVD, as for other steps, but another common technique is *sputtering*

Sputtering works rather like the dry etch process, but in reverse

- Gas atoms are ionised into a plasma, before being accelerated towards the target metal
- Chunks of the metal are ejected from the impact, spraying all over the chamber



Detector fabrication steps

Just to demonstrate some of the steps involved in making a simple diode:

- Start with a silicon wafer



Detector fabrication steps

Just to demonstrate some of the steps involved in making a simple diode:

- Start with a silicon wafer
- Grow an oxide



Detector fabrication steps

Just to demonstrate some of the steps involved in making a simple diode:

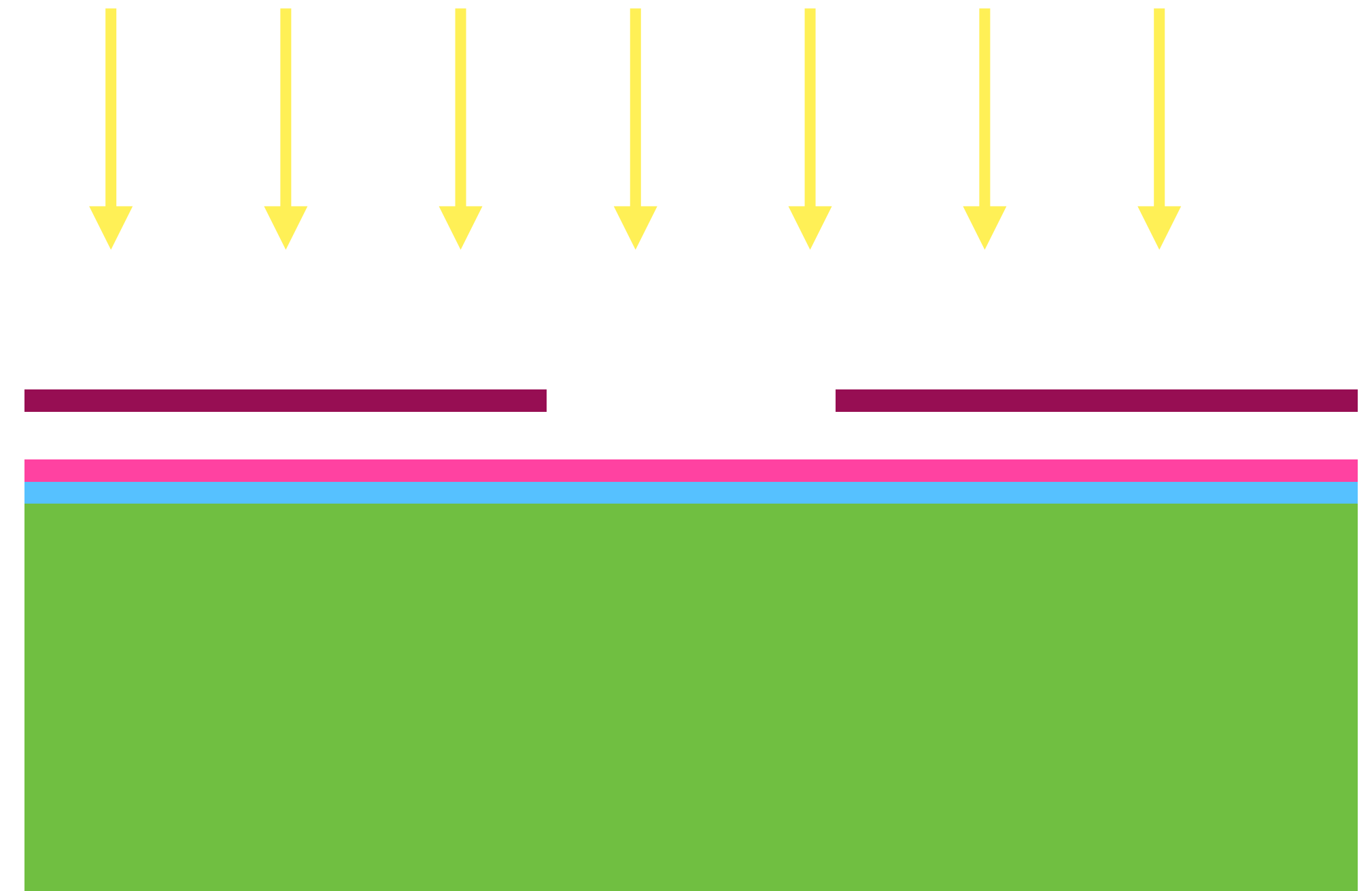
- Start with a silicon wafer
- Grow an oxide
- Spin on a resist



Detector fabrication steps

Just to demonstrate some of the steps involved in making a simple diode:

- Start with a silicon wafer
- Grow an oxide
- Spin on a resist
- Align the lithography mask



Detector fabrication steps

Just to demonstrate some of the steps involved in making a simple diode:

- Start with a silicon wafer
- Grow an oxide
- Spin on a resist
- Align the lithography mask
- Develop the resist



Detector fabrication steps

Just to demonstrate some of the steps involved in making a simple diode:

- Start with a silicon wafer
- Grow an oxide
- Spin on a resist
- Align the lithography mask
- Develop the resist
- Etch the oxide



Detector fabrication steps

Just to demonstrate some of the steps involved in making a simple diode:

- Start with a silicon wafer
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- Develop the resist
- Etch the oxide
- Remove the resist



Detector fabrication steps

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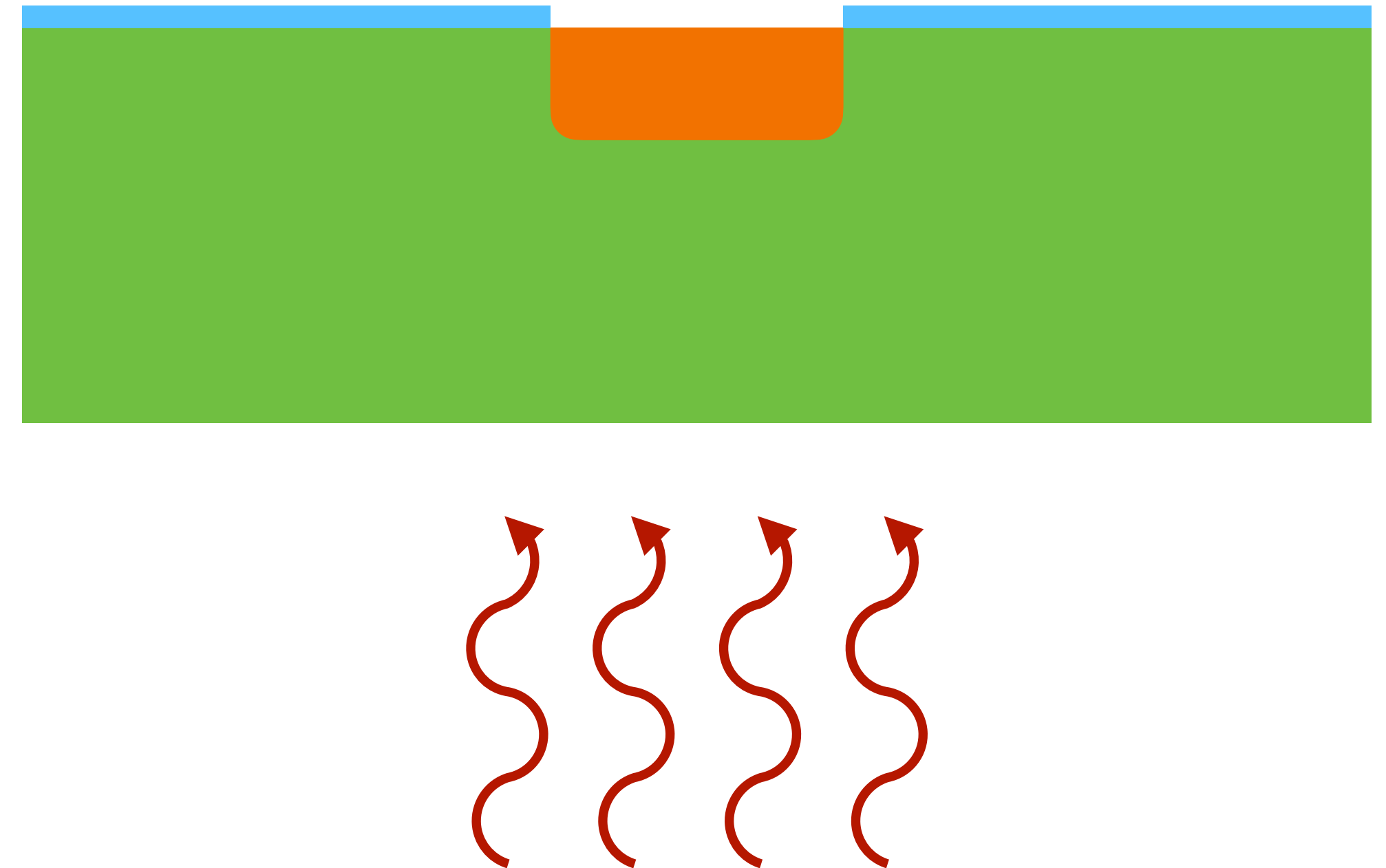
- Start with a silicon wafer
- Grow an oxide
- Spin on a resist
- Align the lithography mask
- Develop the resist
- Etch the oxide
- Remove the resist
- Dope the silicon



Detector fabrication steps

Just to demonstrate some of the steps involved in making a simple diode:

- Start with a silicon wafer
- Grow an oxide
- Spin on a resist
- Align the lithography mask
- Develop the resist
- Etch the oxide
- Remove the resist
- Dope the silicon
- Activate the implant



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 - Activate the implant
- Spin a new resist

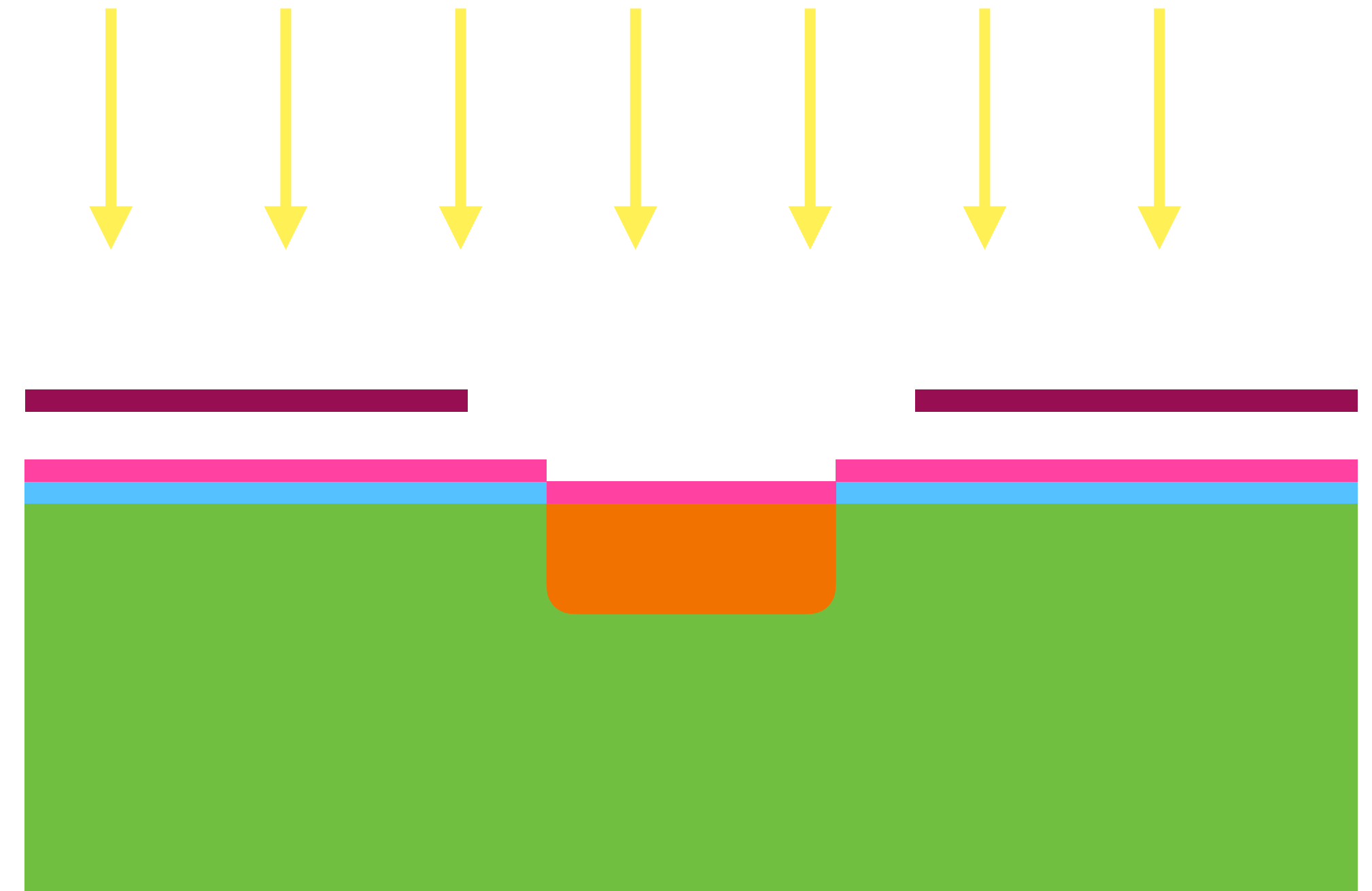


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- Spin a new resist
- Align a new mask



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 - Align a new mask
 - Develop the resist



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- Spin a new resist
 - Align a new mask
 - Develop the resist
 - CVD with metal
 - Remove the resist



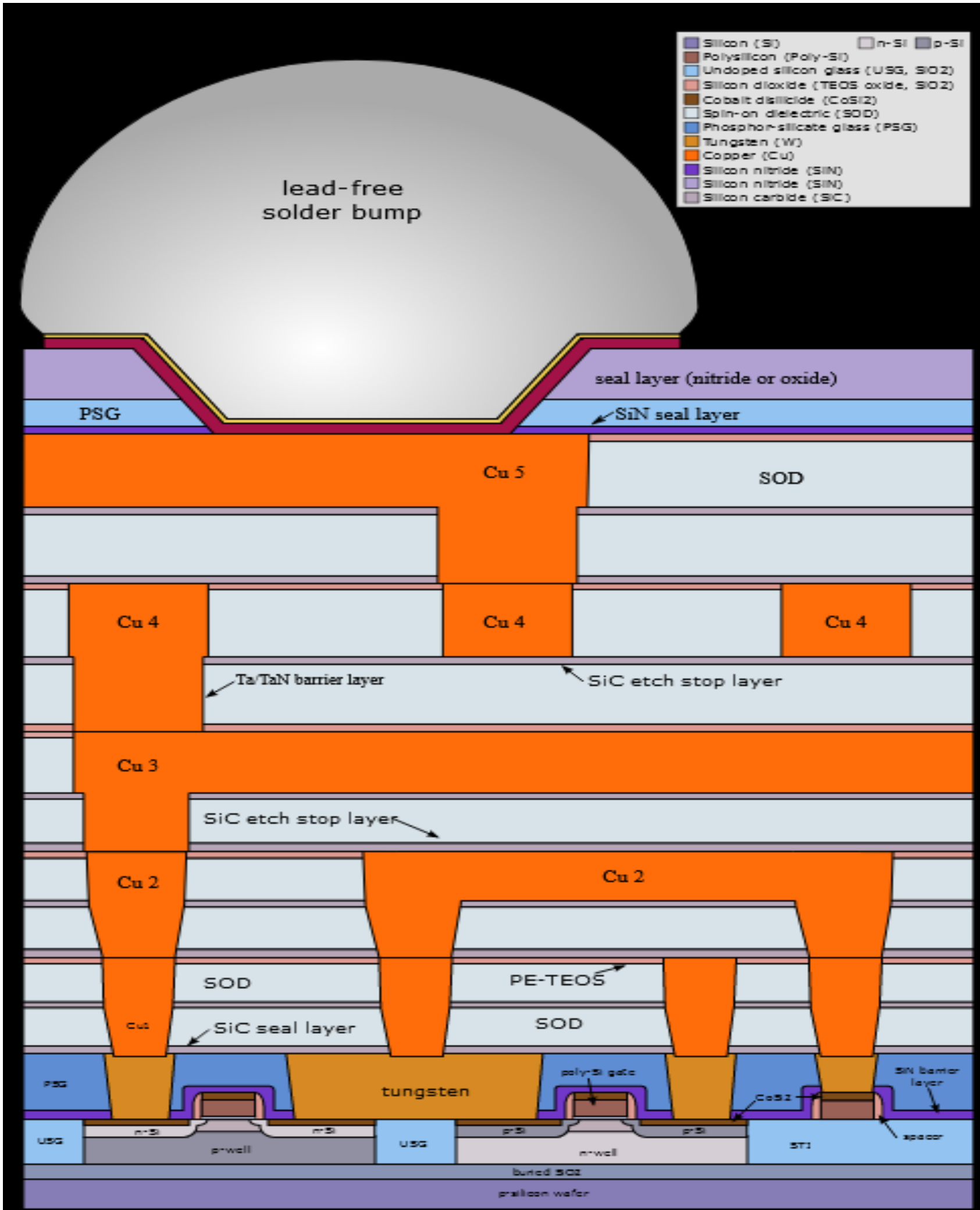
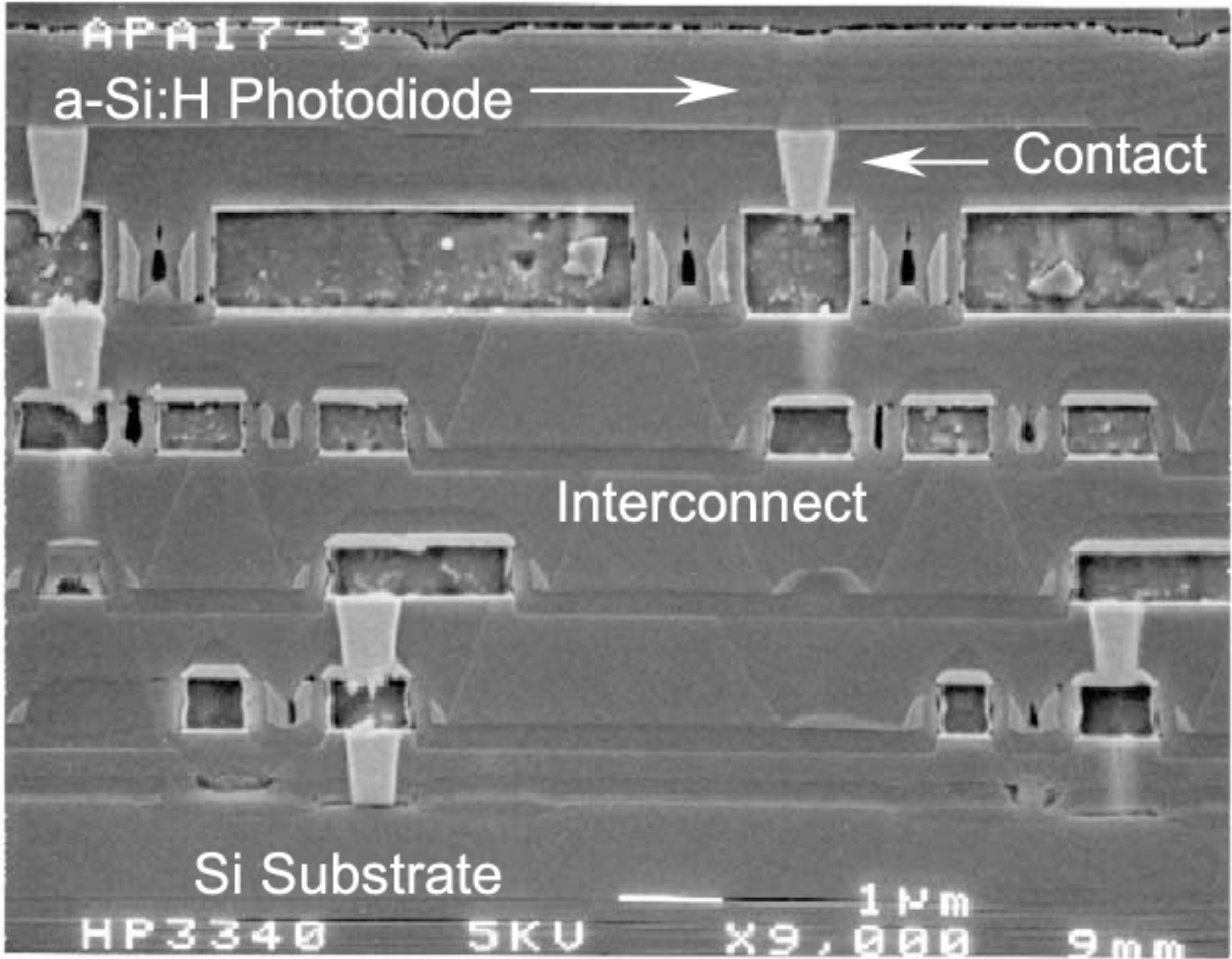
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- Grow an oxide
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- Etch the oxide
- Remove the resist
- Dope the silicon
- Activate the implant
- Spin a new resist
- Align a new mask
- Develop the resist
- CVD with metal
- Remove the resist
- CVD the back side



CMOS cross-section



Tiny Rock Boards that Catch Little Pieces of Light and Fire

(Silicon pixel detectors)

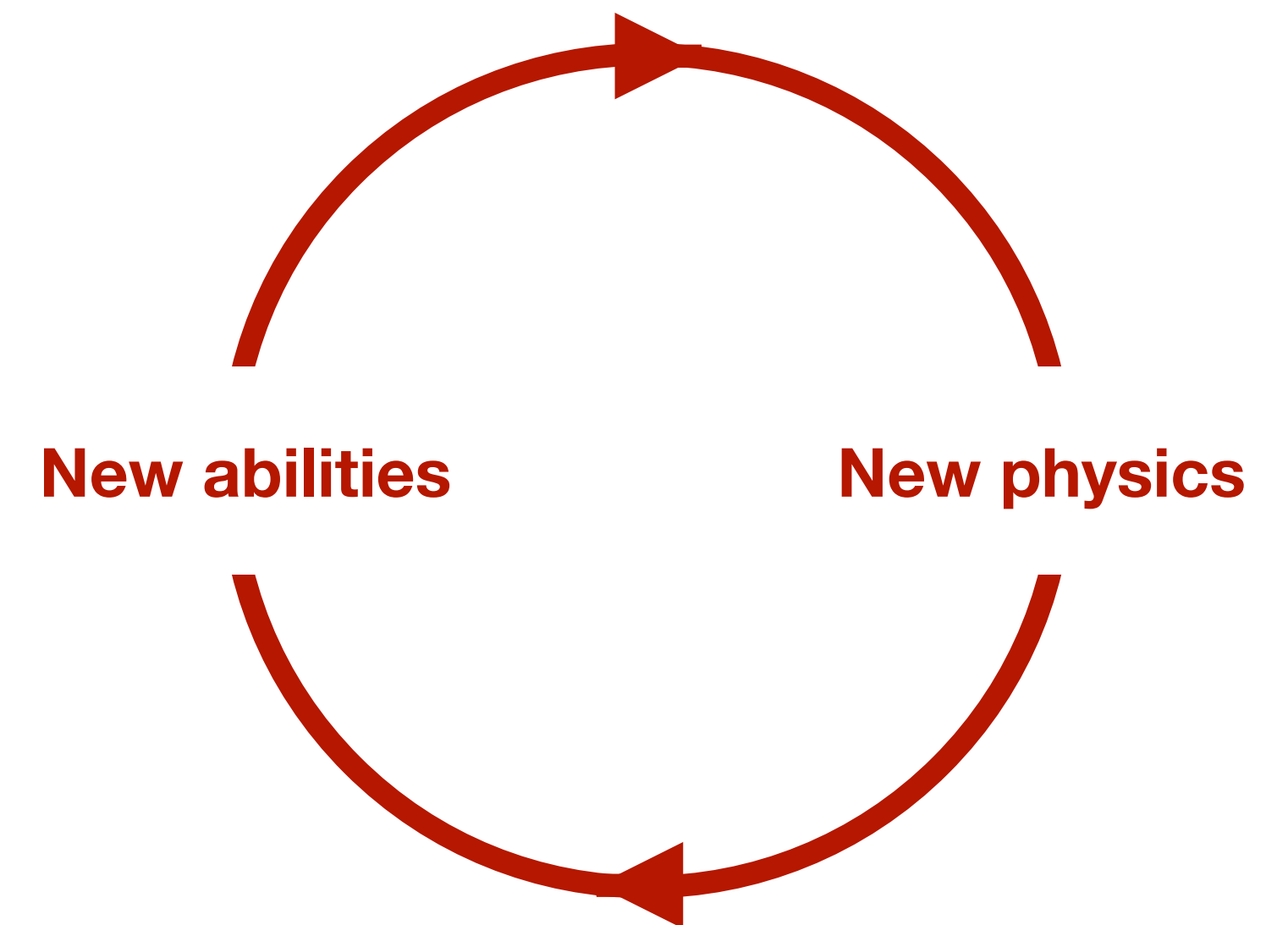
Chickens and eggs

During detector development a common question is: why do we need this? What is the application?

- To some extent a chicken and egg problem. Once we have new measurement abilities we can do more in an experiment, which might push the requirements further

I would argue that this is not true!

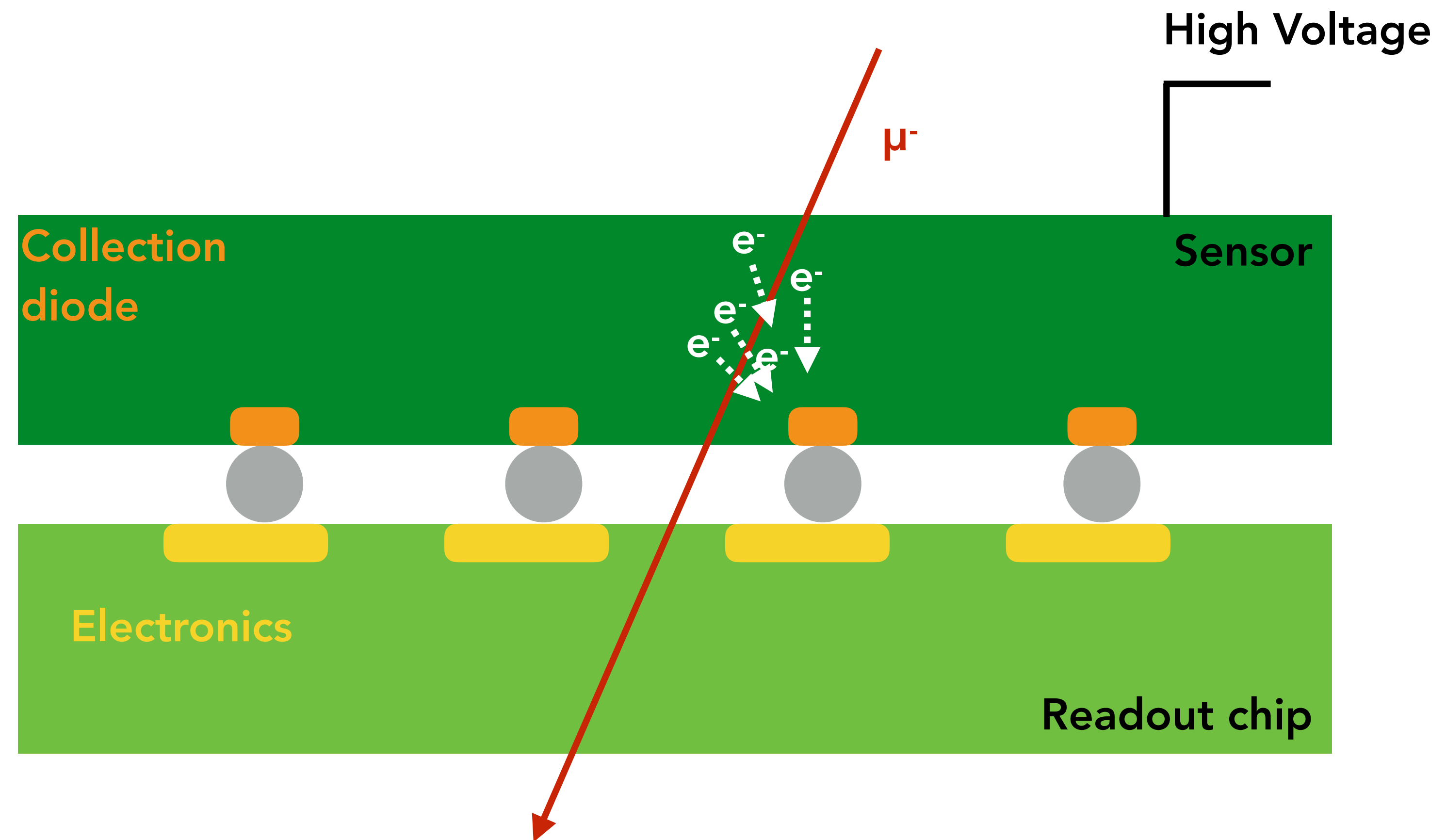
- Huge number of unexpected discoveries in physics directly preceded by improvement in **measurement technology**
- Once we introduce new abilities that we never imagined, we find new ways of exploiting it



Hybrid pixel detector - planar sensor

The technology of choice for most inner (vertex) detectors is still **hybrid pixel detectors**

- Here, the sensor containing our PN-junctions is connected channel by channel (which may be 50 μm in size) to a readout chip
- This decouples what happens on the sensor and in the electronics, and allows us to fully deplete the sensor



Planar sensors: p-on-n

Start with the basic layout of our hybrid pixel detector, revolving around a pn-junction

- **Historically n-type wafers were easier to grow => p-on-n detectors**



Planar sensors: n-on-n

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=> move to electron collection



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Planar sensors: n-on-p

Start with the basic layout of our hybrid pixel detector, revolving around a pn-junction

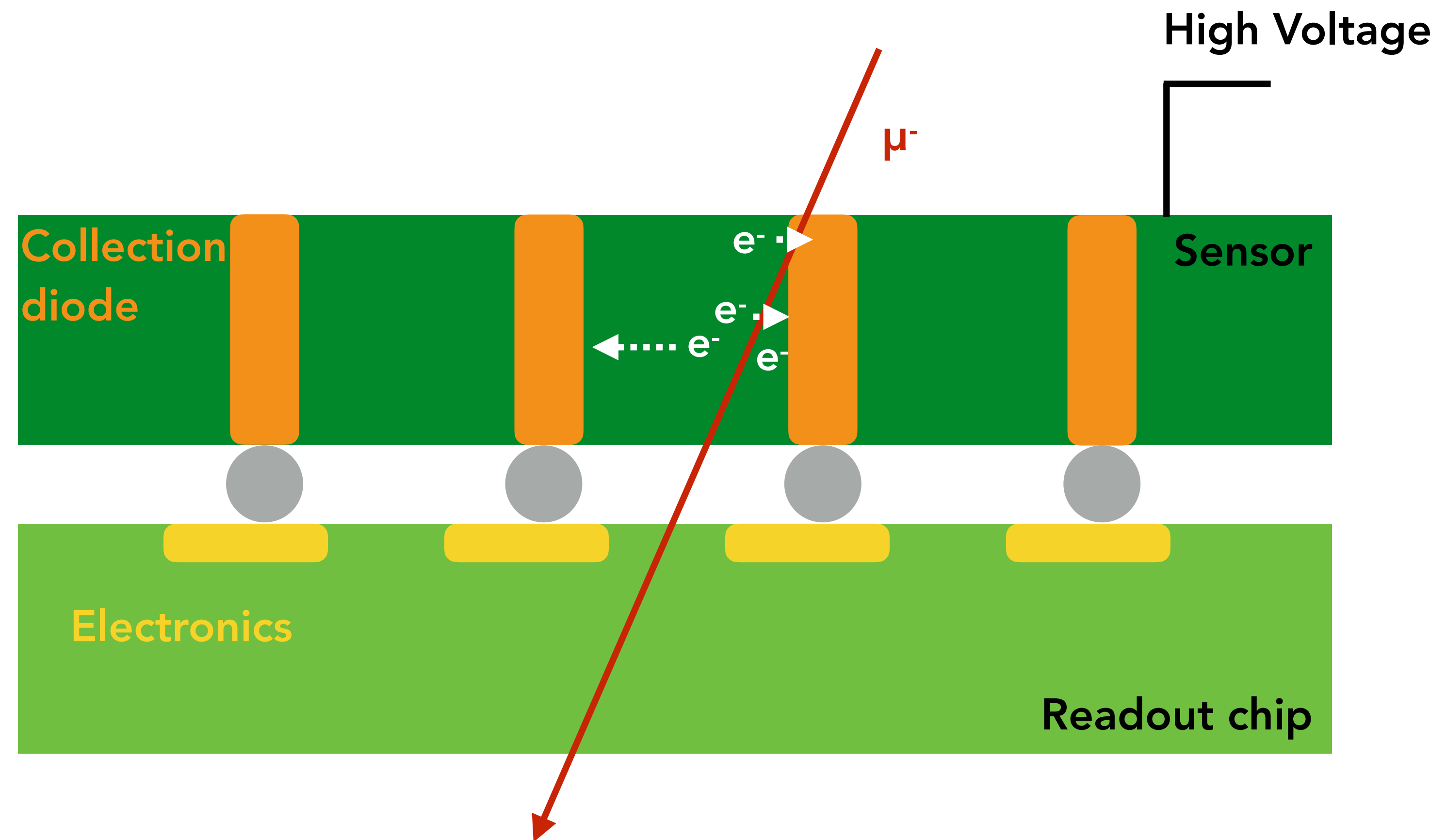
- Historically n-type wafers were easier to grow => p-on-n detectors
- Faster electron mobility in silicon => move to electron collection
- **Improvements in silicon processing + costs of double-sided processing => move to n-on-p**
- In attempts to improve radiation hardness, trend has been to thinner sensors (*but* lower overall signal)



Hybrid pixel detector - 3D sensor

To counter the effects of radiation damage,
reduce the charge collection path

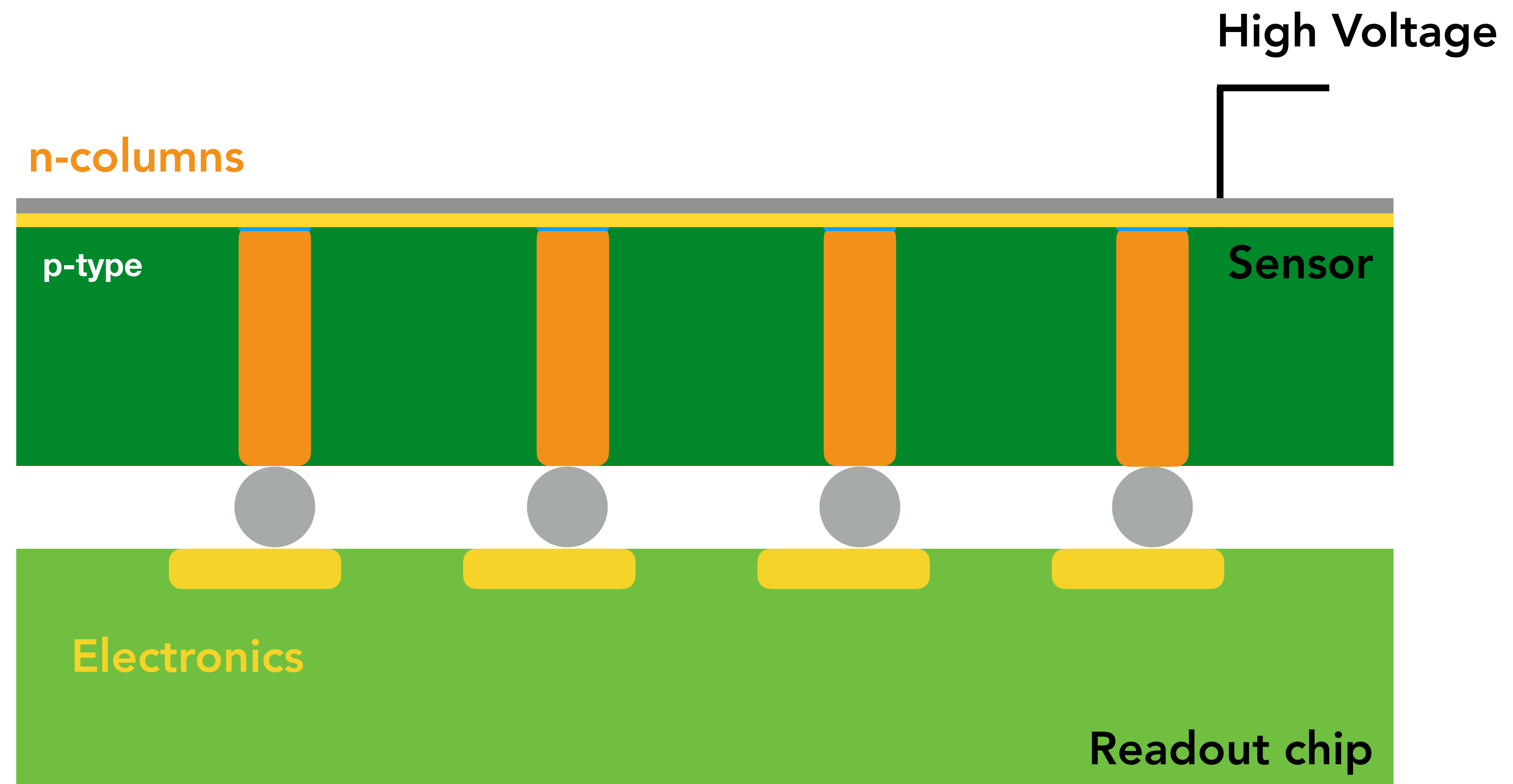
- **3D columns** give short path
irrespective of sensor thickness
- Fast charge collection time and thick
detectors - radiation hard!



Types of 3D sensors

Different 3D sensor layout depending on processing and desired fields

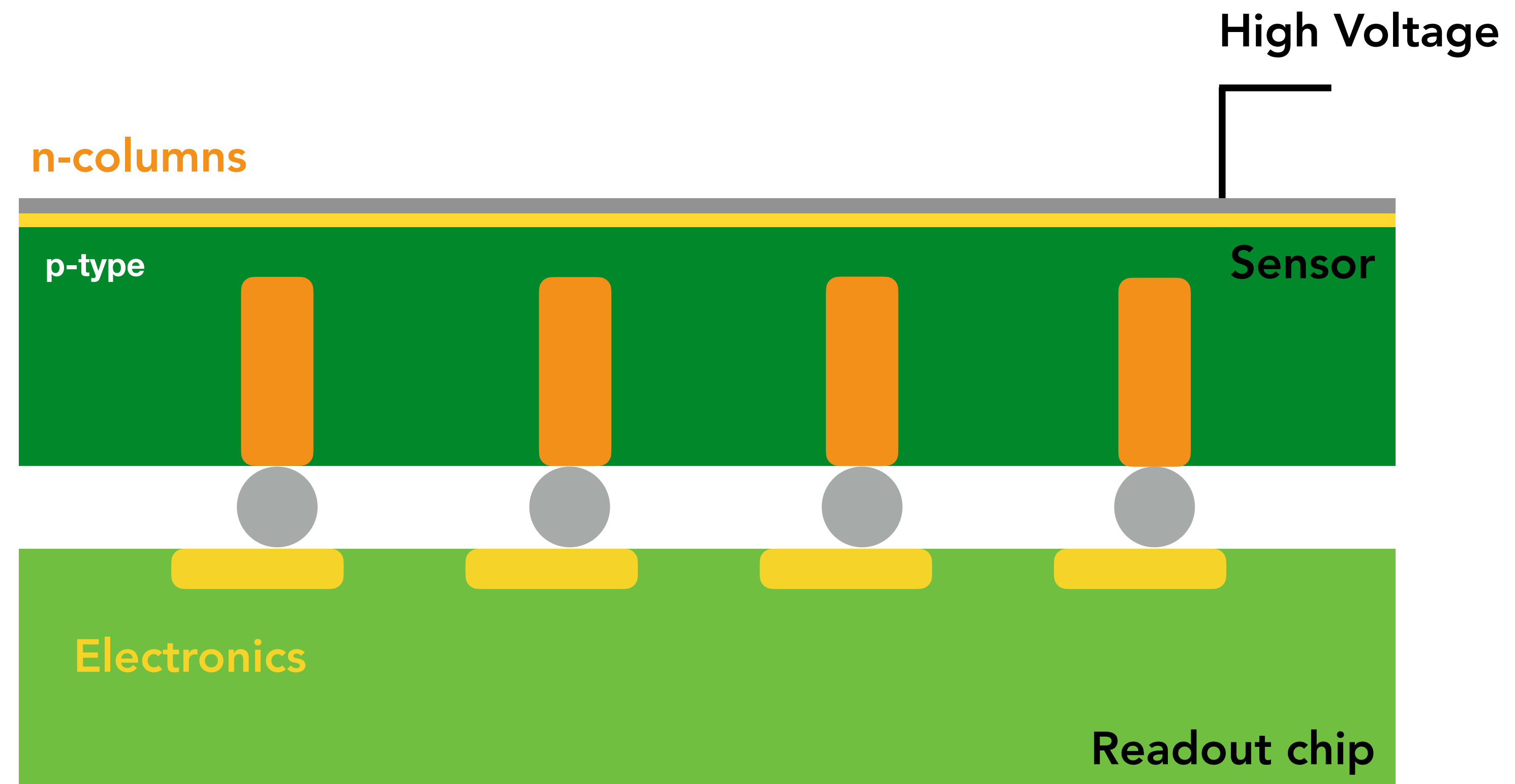
- Single column type, full depth



Types of 3D sensors

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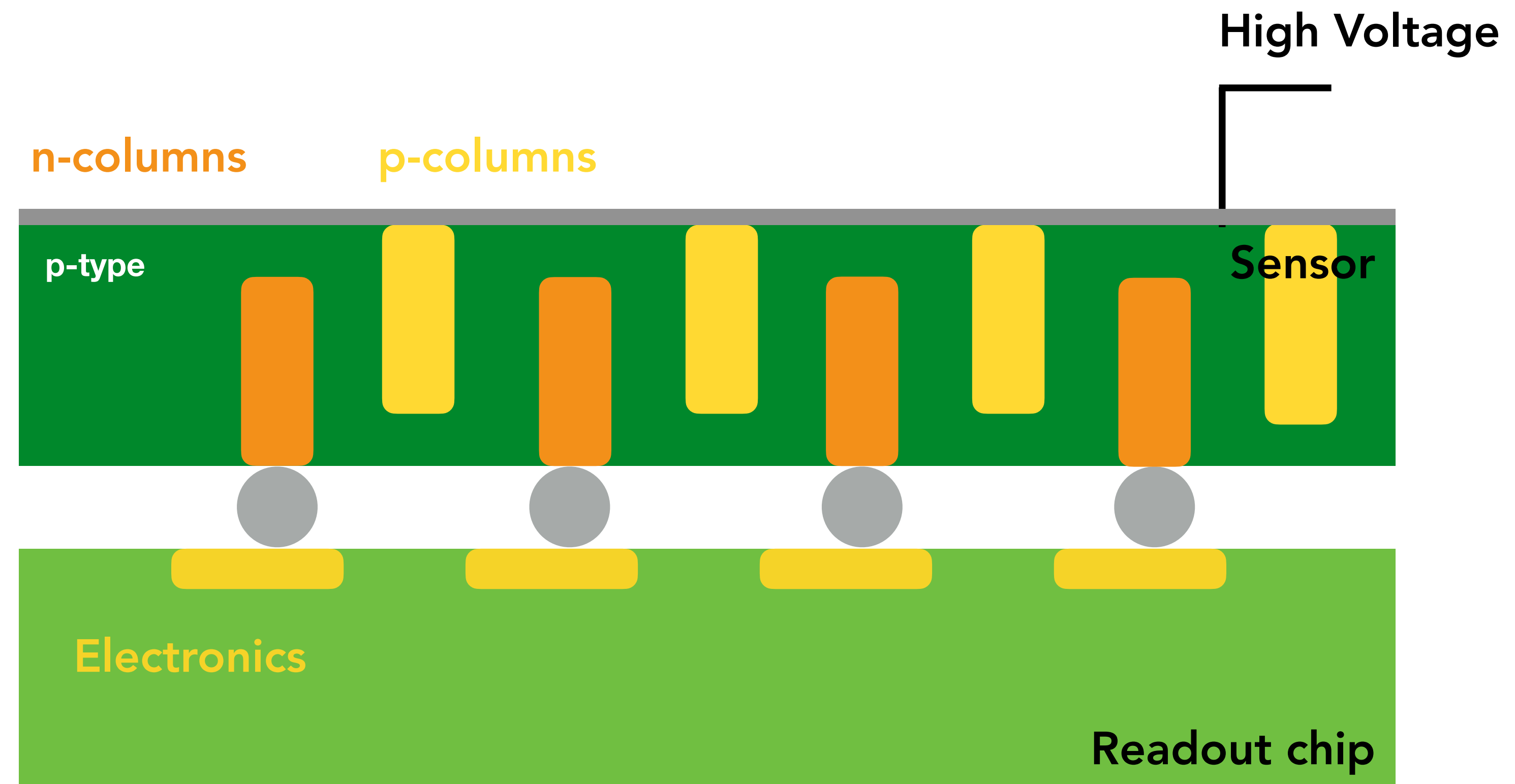
- Single column type, full depth
- Single column type, partial depth



Types of 3D sensors

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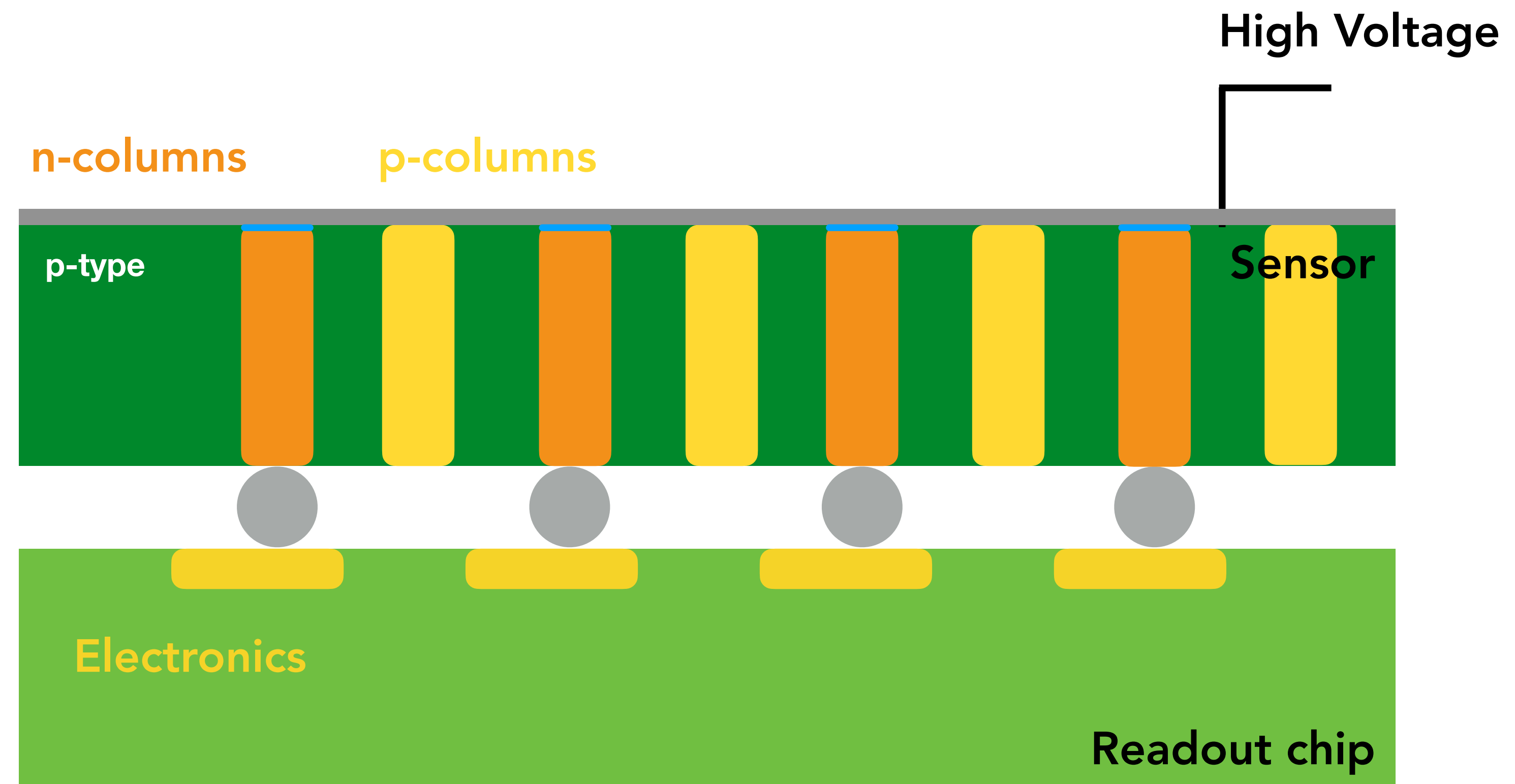
- Single column type, full depth
- Single column type, partial depth
- Dual column types, partial depth



Types of 3D sensors

Different 3D sensor layout depending on processing and desired fields

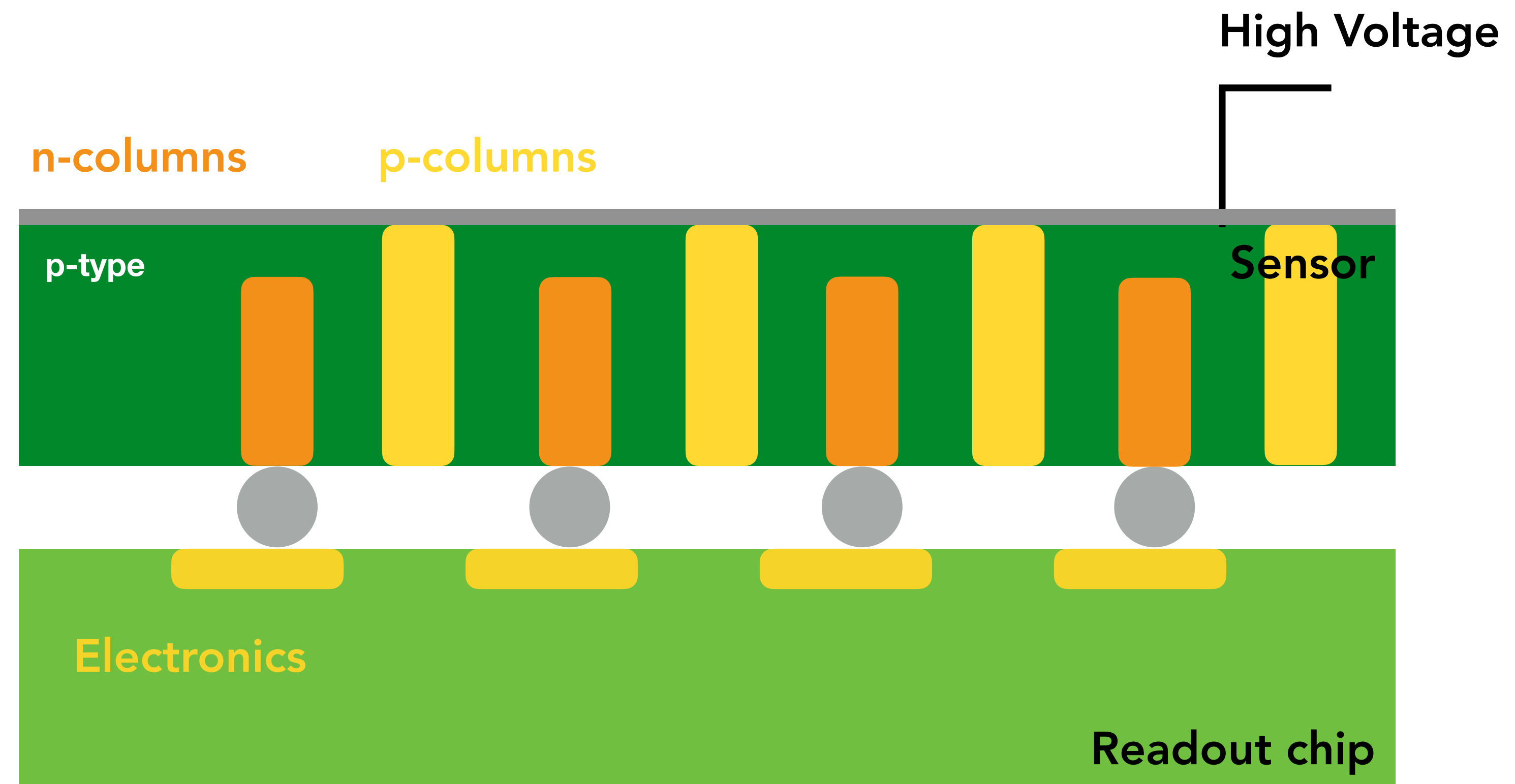
- Single column type, full depth
- Single column type, partial depth
- Dual column types, partial depth
- Dual column types, full depth



Types of 3D sensors

Different 3D sensor layout depending on processing and desired fields

- Single column type, full depth
- Single column type, partial depth
- Dual column types, partial depth
- Dual column types, full depth
- Dual column types, mixed depths

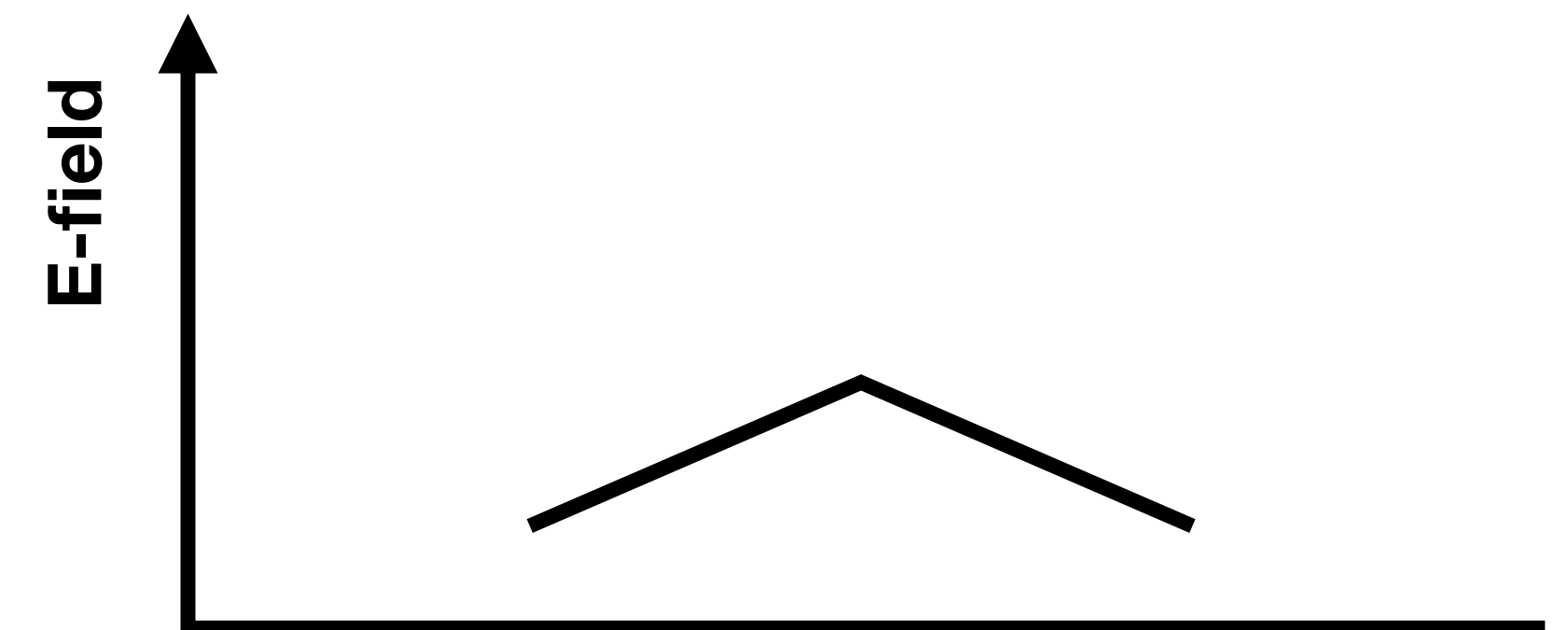


Interlude: Charge multiplication

If we can generate a high enough electric field then charge carriers passing through this region will gain enough energy to cause *further impact ionisation*

- Onset in silicon at around 300 kV/cm
- We want to keep the signal proportional (otherwise see Silicon PhotoMultipliers)

Use doped regions in proximity to each other to localise the high field region

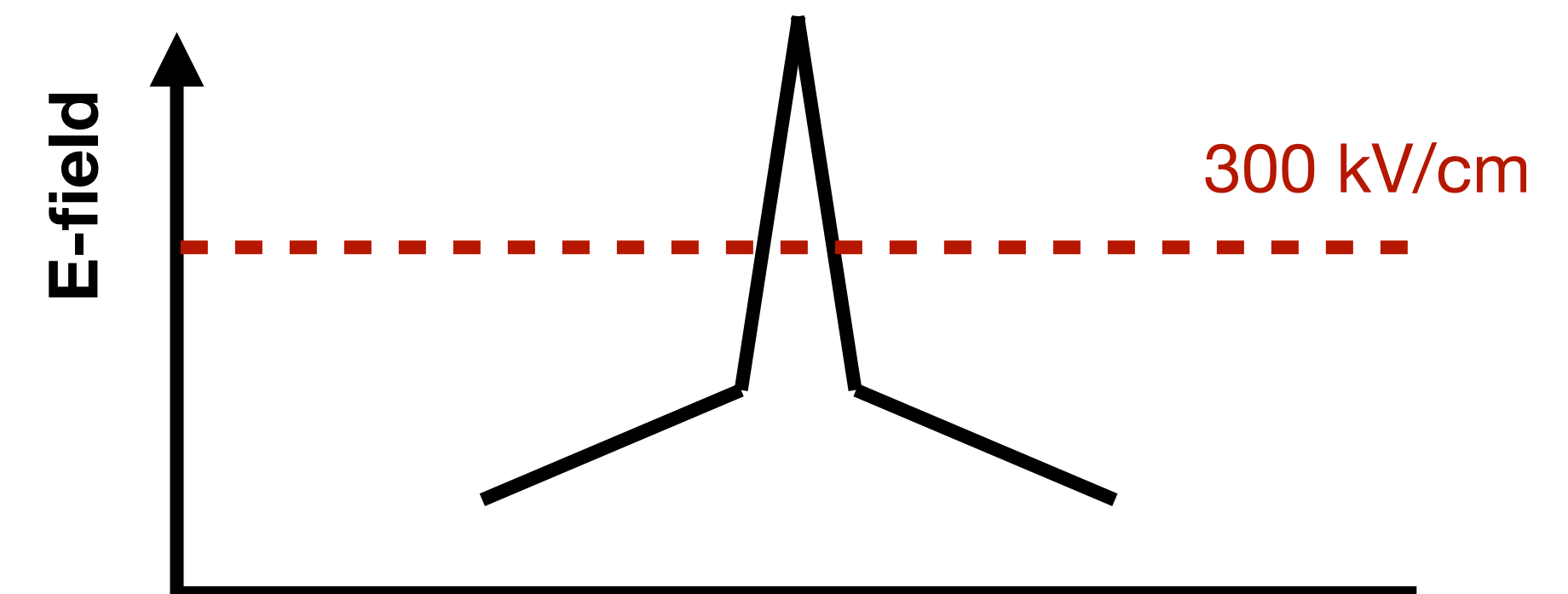
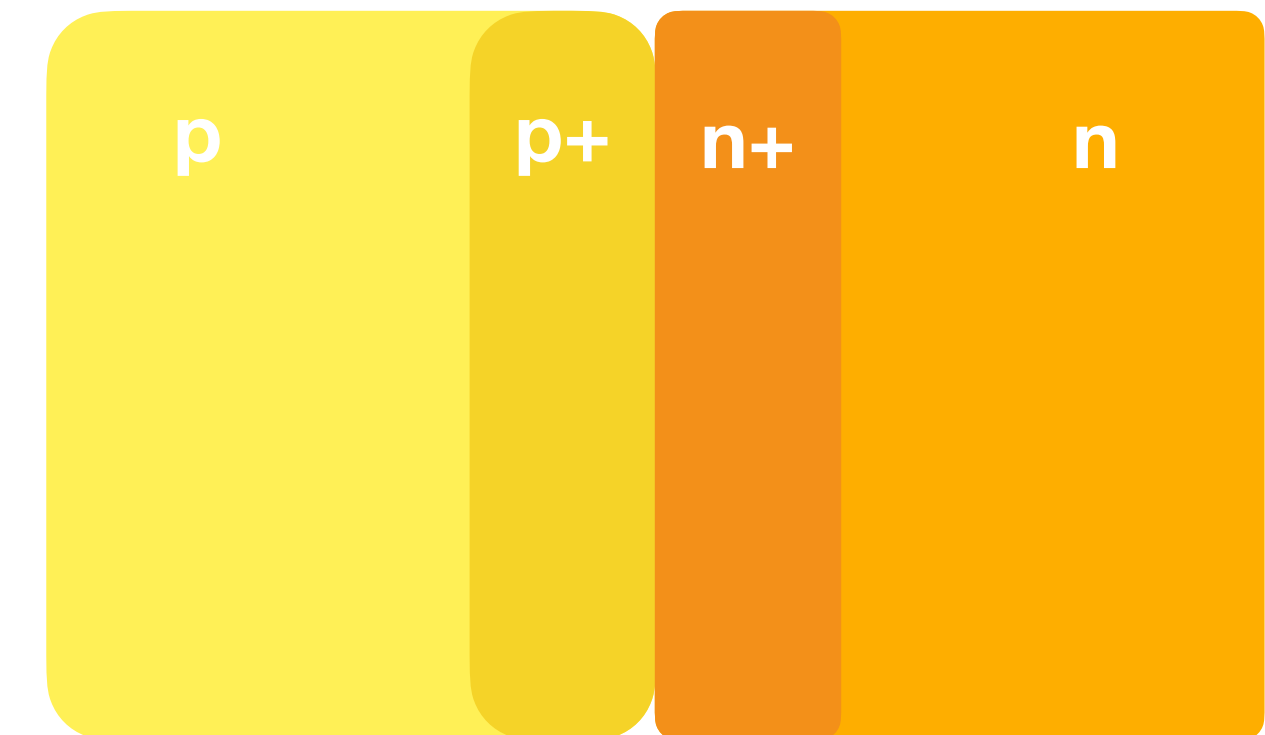


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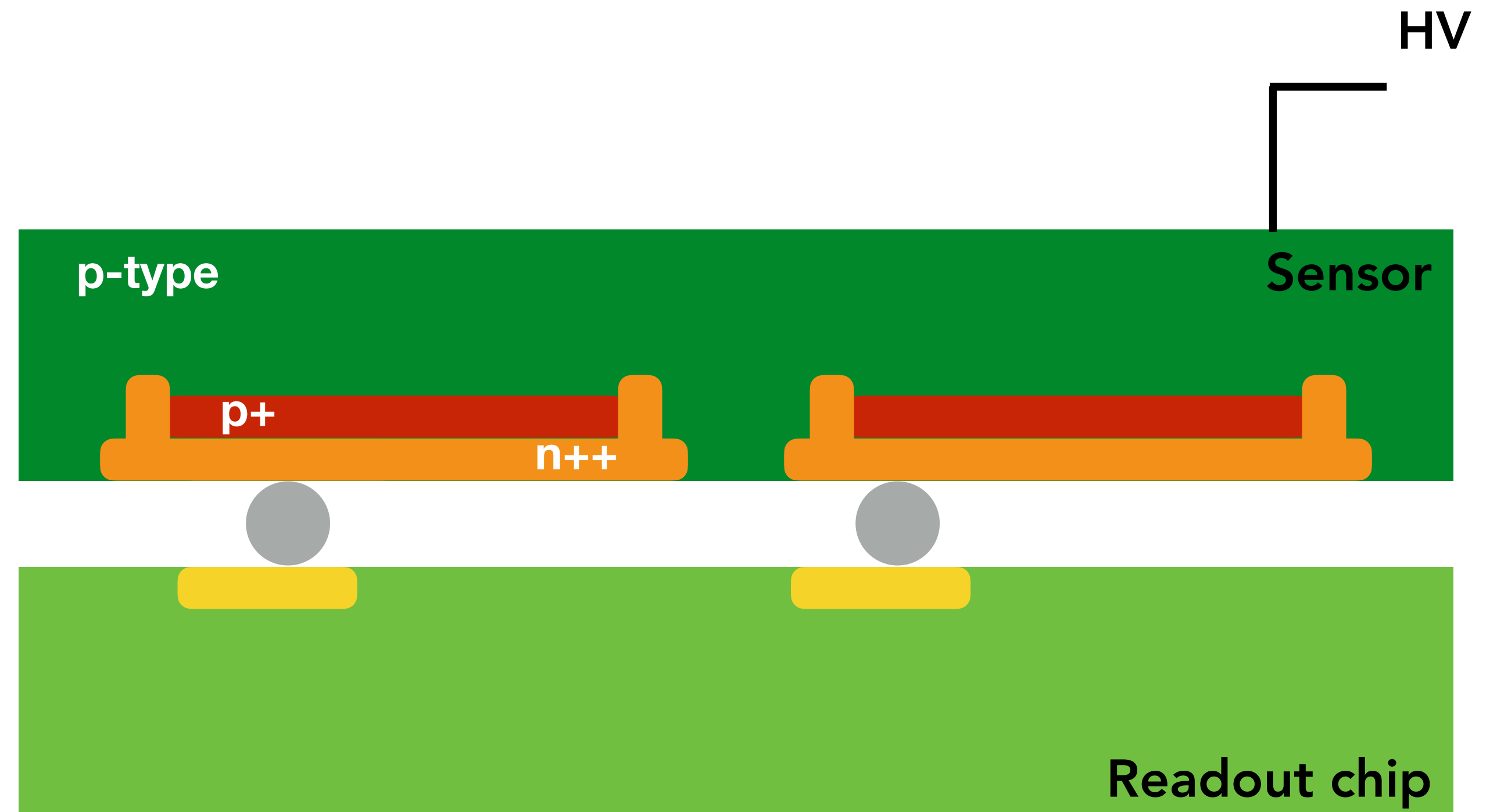


Low Gain Avalanche Diode (LGAD)

The 'classic' LGAD flavour contains a wider n++ collection diode than we have looked at till now (not in itself strange) but **has a deep p-implant that leads to a high field region**

- It also features some structures at the side to step down the electric field at the edge of the gain region - JTE

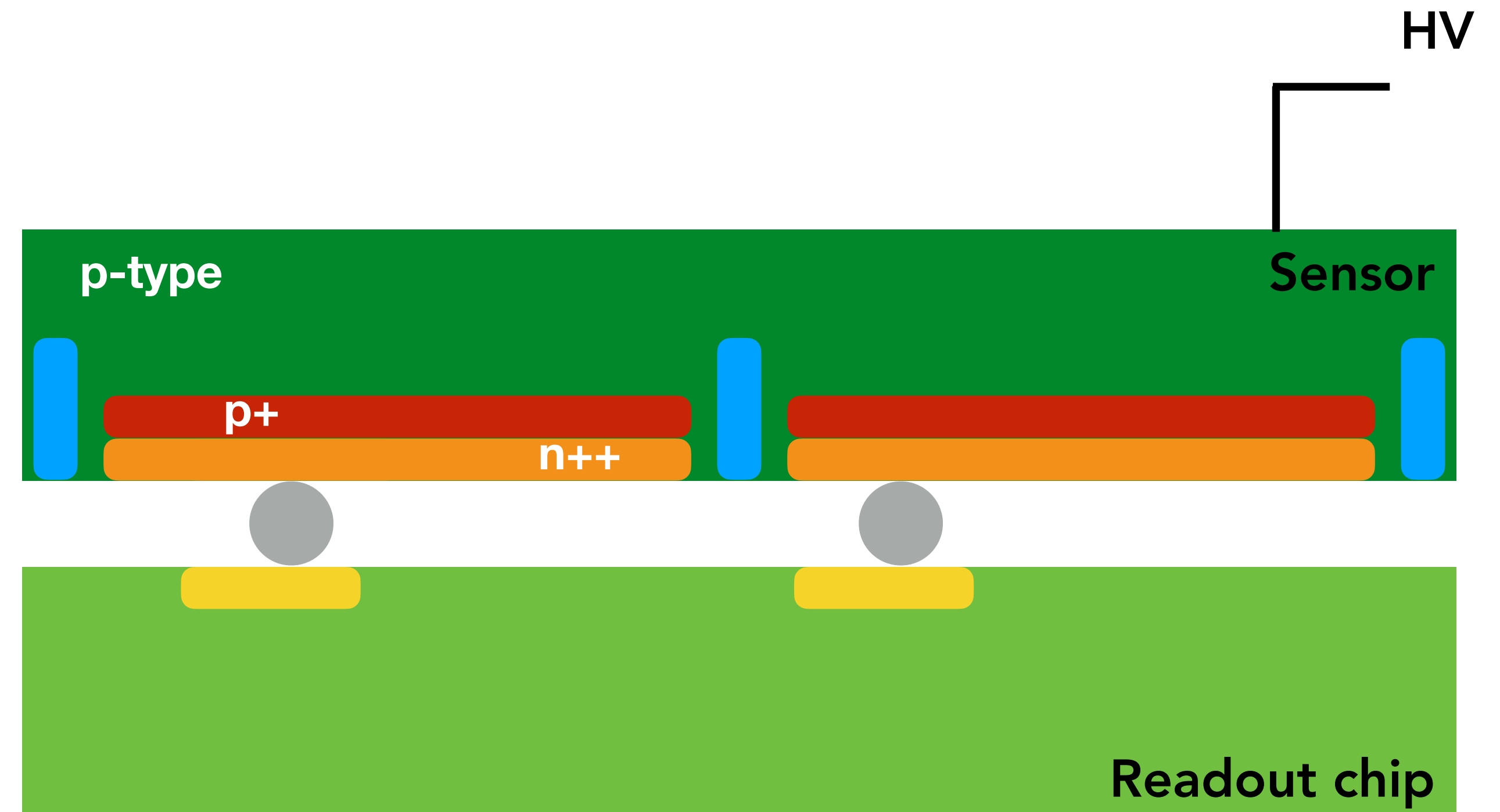
These devices give us gain, but there are clear regions between the pixels that allow charge to reach the collection node without entering the high field zone



Trench Isolated (TI-) LGAD sensors

Many avenues of research focussed on how to minimise the non-gain region between pixels

- Take inspiration from SiPM - etch trenches between the pixels to isolate the electric field on either side

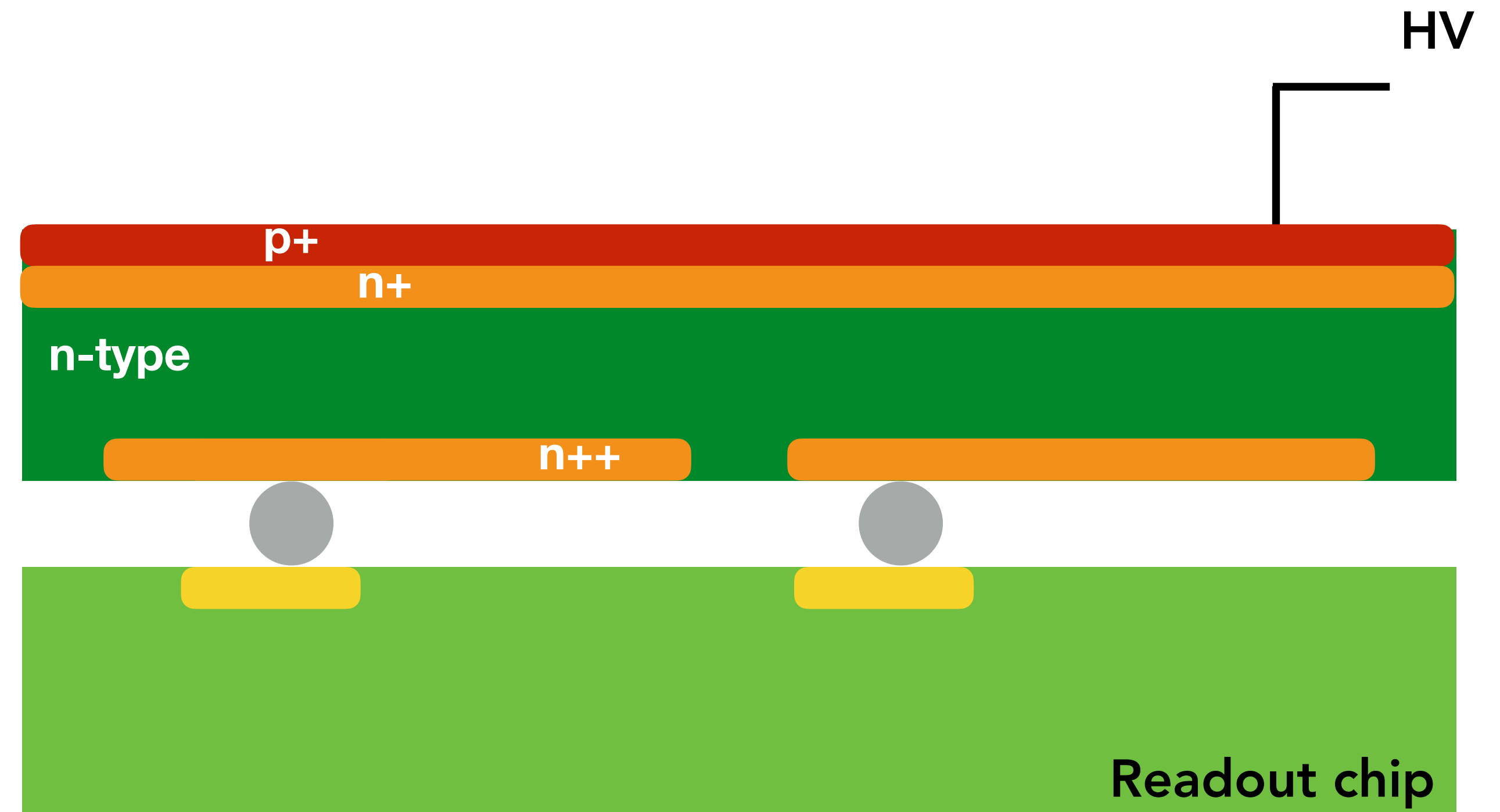


Inverted LGAD sensors (electron collecting)

Look back at the development of planar sensors

- n-on-n sensors used junction on the back side to create the depletion region

If we take our previous LGAD design and move the junction to the back side, then we will collect the electrons immediately, while the holes will travel to the backside and enter the gain region



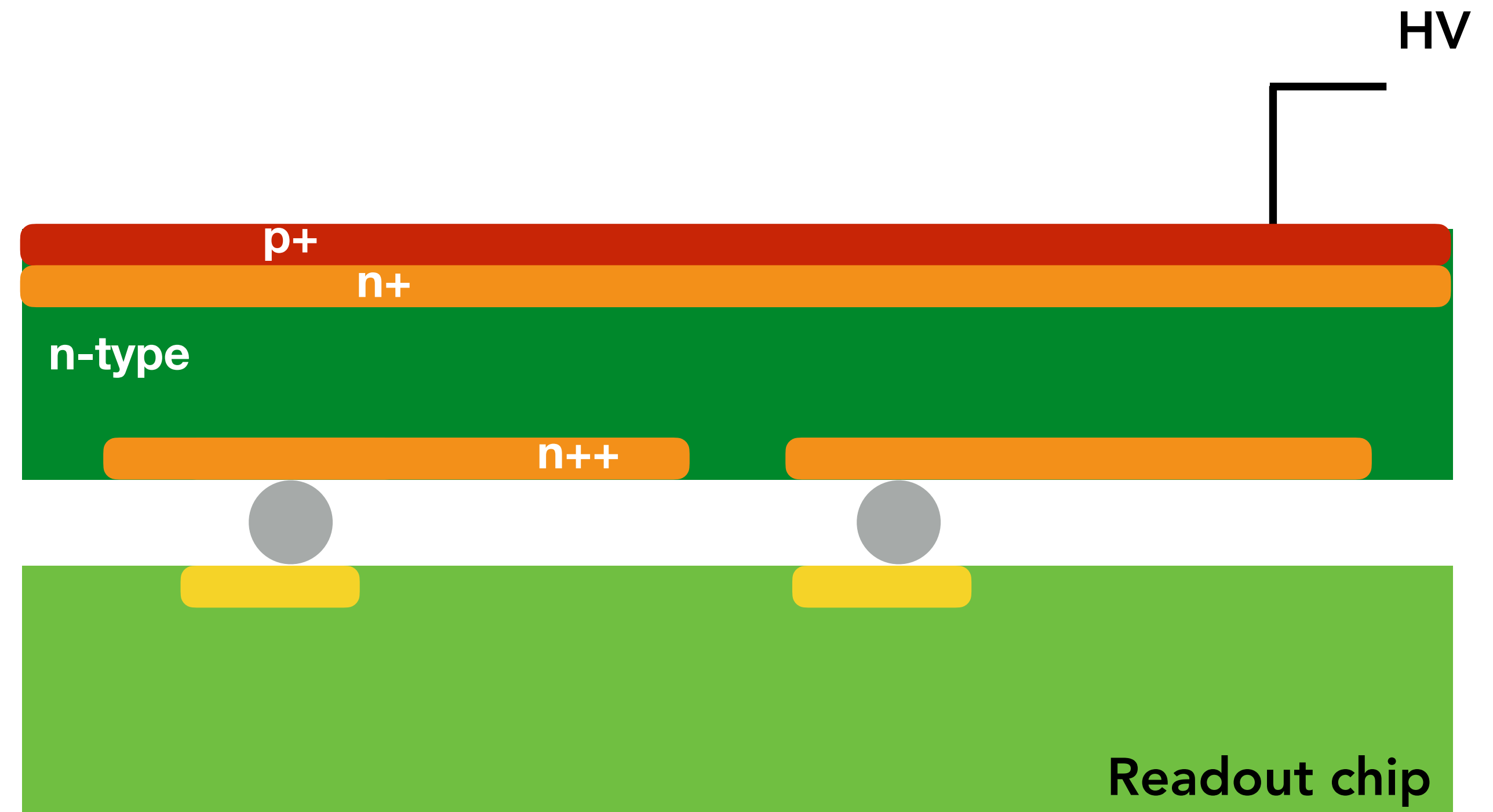
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- BUT... **hole impact ionisation coefficient significantly lower than that of electrons**



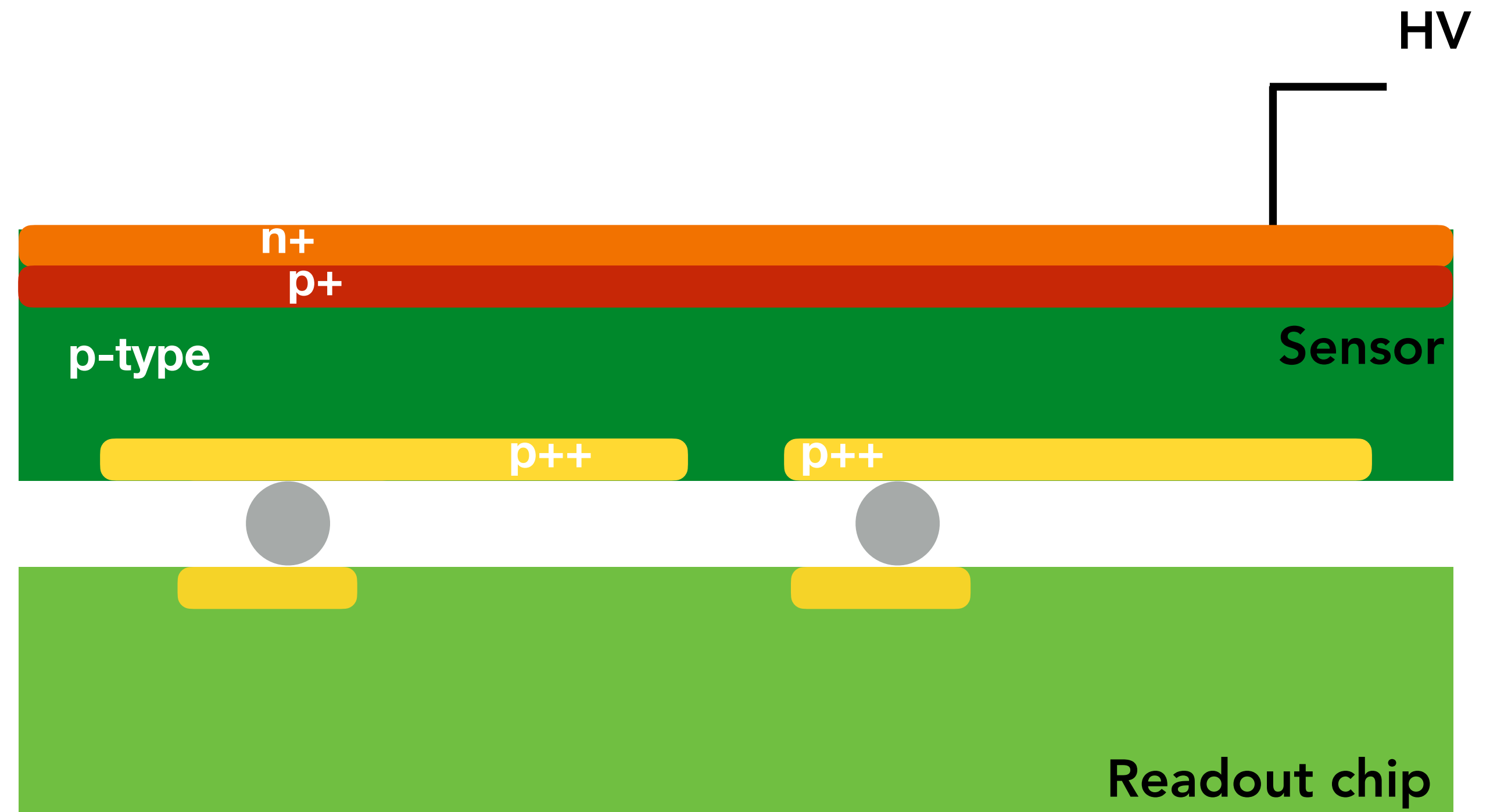
Inverted LGAD sensors (hole collecting)

Ok, so here we go again

- Invert all of the p and n labels - now we collect holes (slower than electrons) but the electrons all travel to the backside and undergo impact ionisation

This works, but now you have to ask how long it takes to drift to the backside...

- Not ideal for fast timing applications



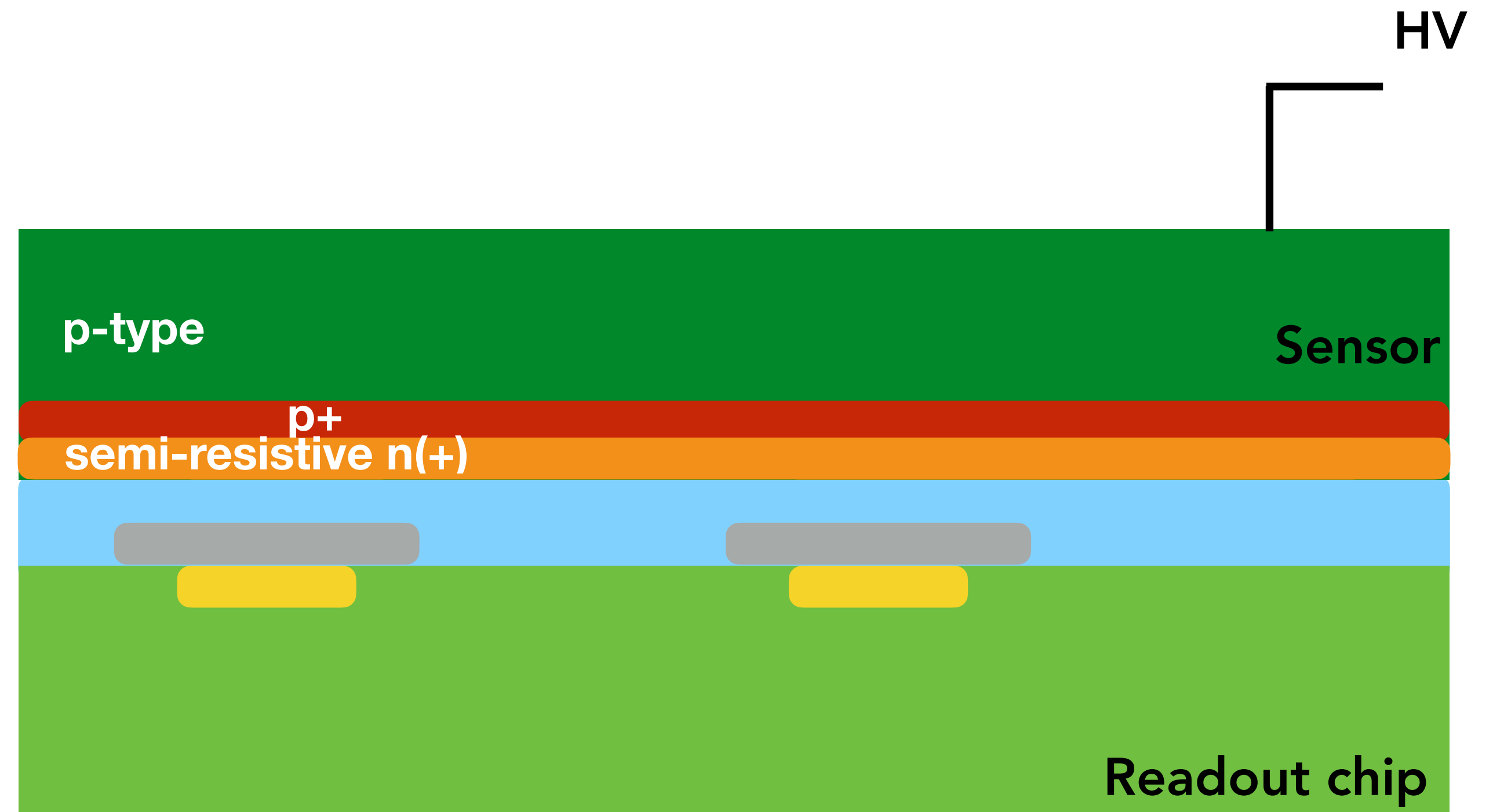
AC-LGADs (Resistive Silicon Detectors)

Now try something similar but quite a bit different

- Don't segment the collection nodes or the gain layer
- Instead, have effectively a massive Pad detector and *capacitively couple your readout electronics*

Uniform weighting field, impact ionisation at the front side, electron collecting!

- Some care to be taken that the n+ layer doesn't build up charge - still need DC connection

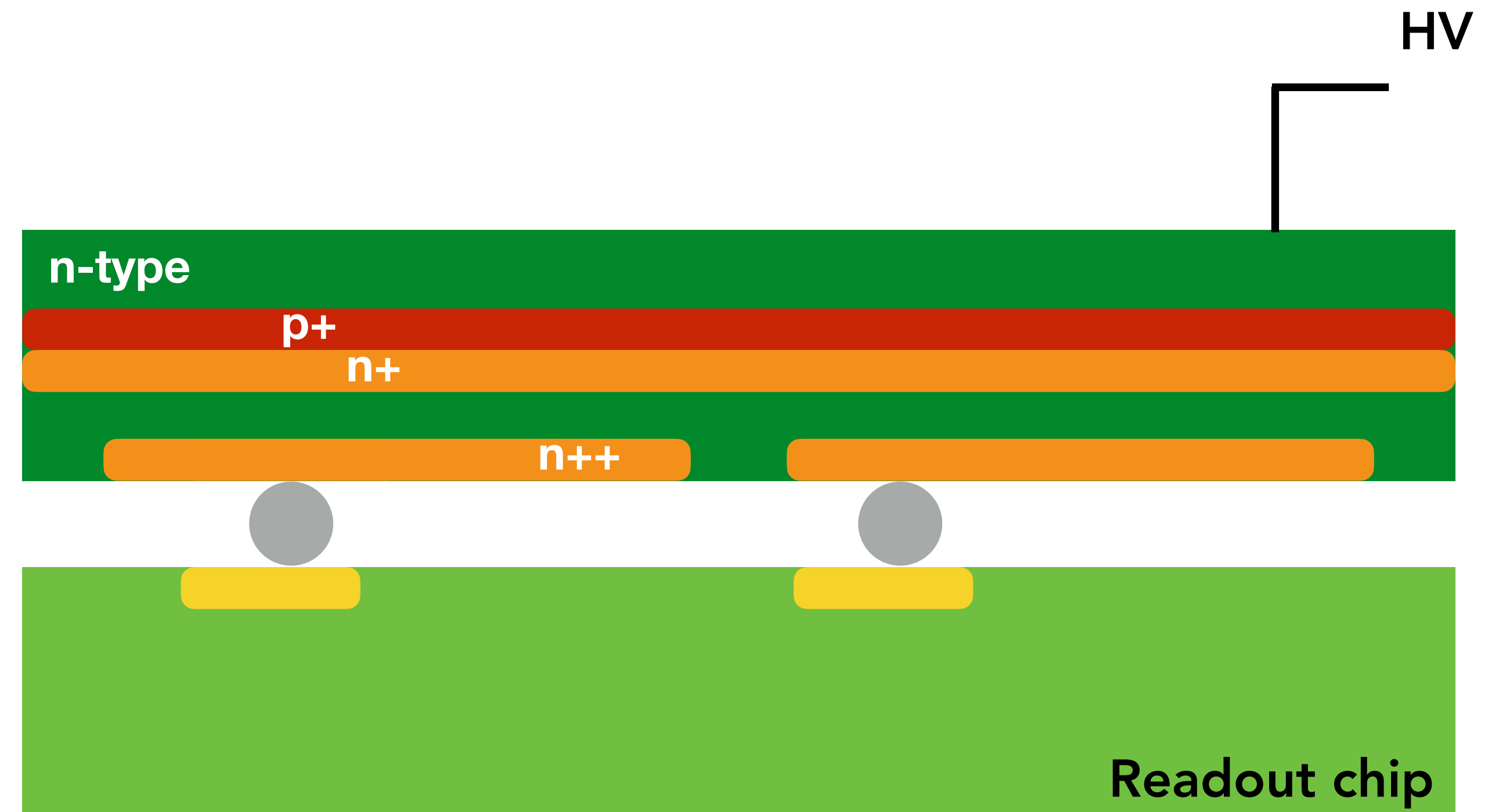


Buried gain layers

If you have the processing ability (doped epi growth/high energy implantation) you could try to create a gain layer *deeper in the silicon*

- This should be deeper than your collection node but doesn't have to be too deep
- If sufficient fraction of generated charges still pass through the gain region could give quite uniform multiplication

Apply the same concept to monolithic devices



Monolithic detector - MAPS

Hybrid pixel detectors give a lot of functionality in-pixel, and allow flexibility for the choice of sensor, but are relatively expensive devices

- If this is all being done in silicon, why not use the same piece? **Monolithic detectors**

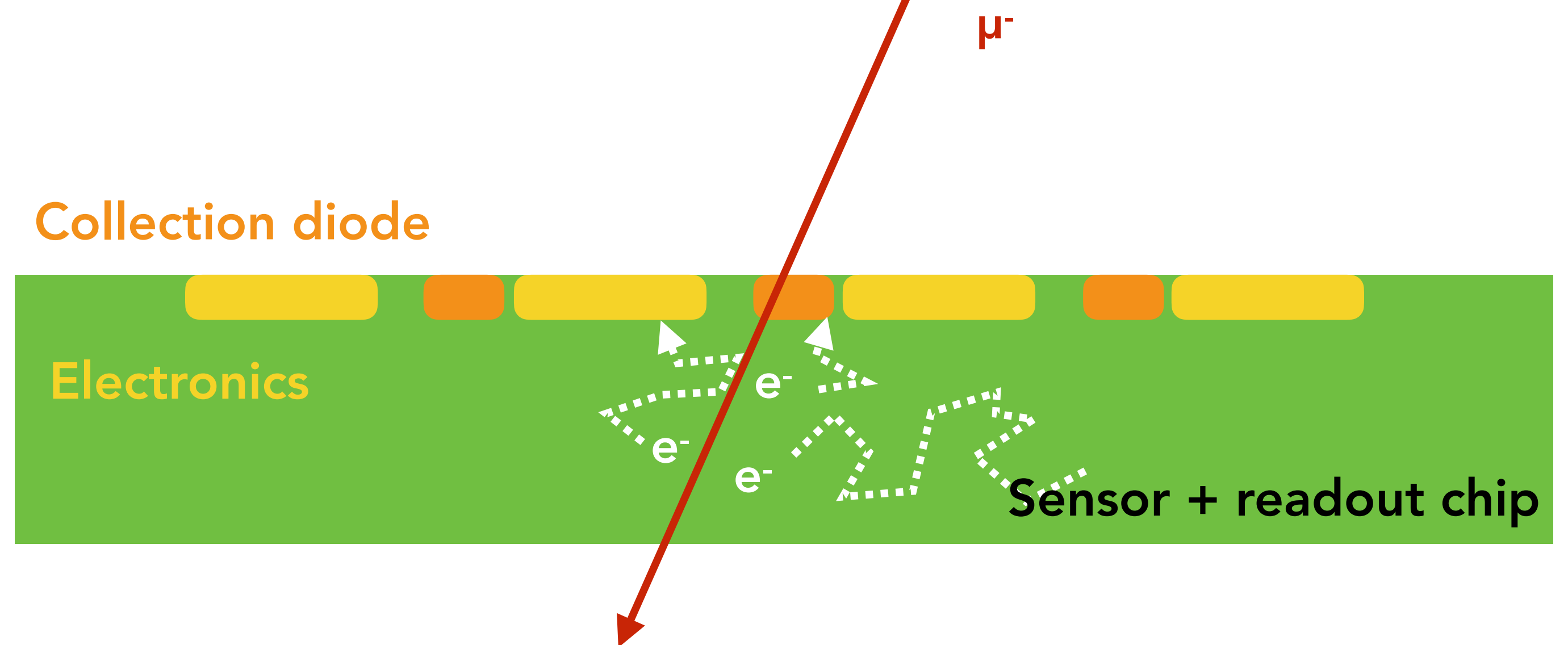
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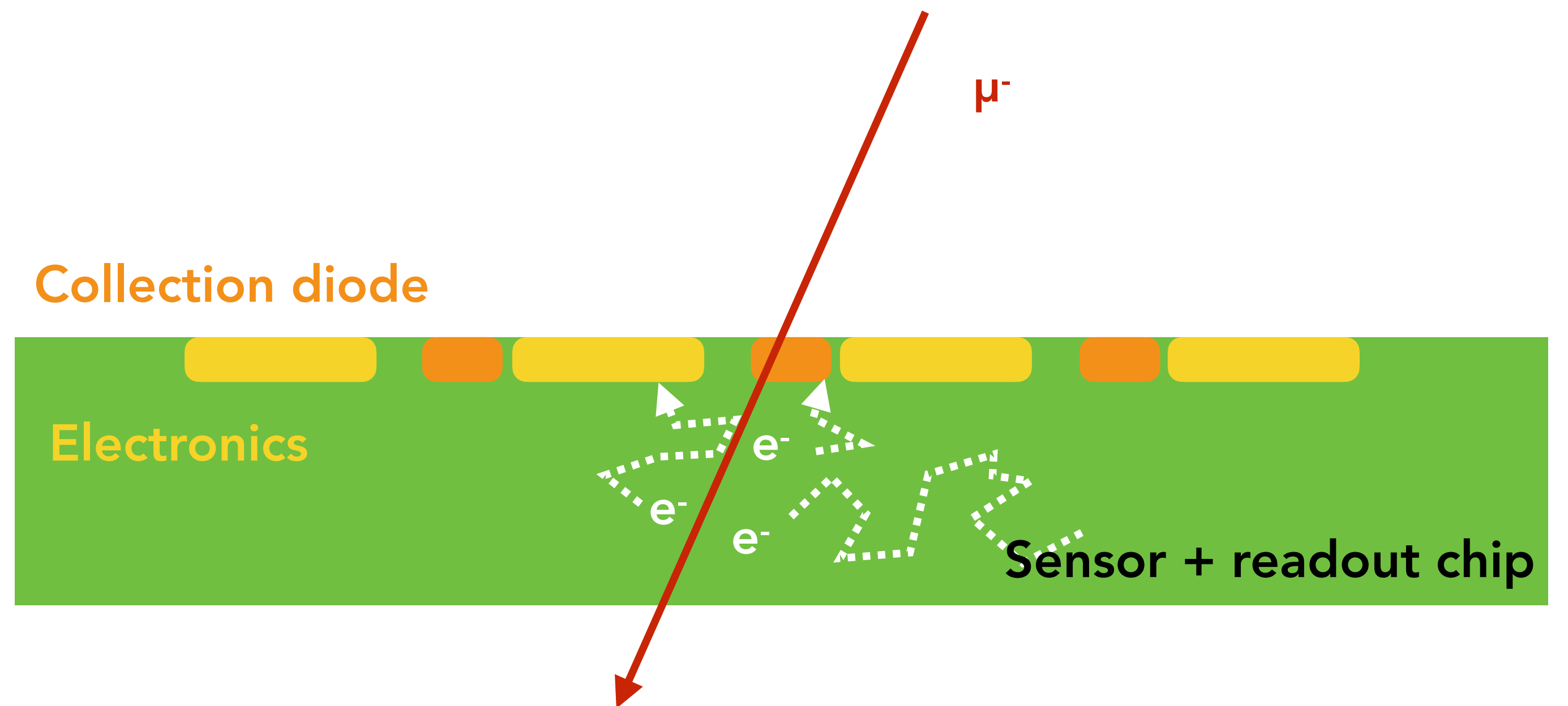
Historically these suffered from slow charge collection, since any voltage applied to the silicon would also upset the electronics

- Small built-in depletion region
- Most charge collected by diffusion
- **Monolithic Active Pixel Sensors (MAPS)**



Monolithic detector - MAPS

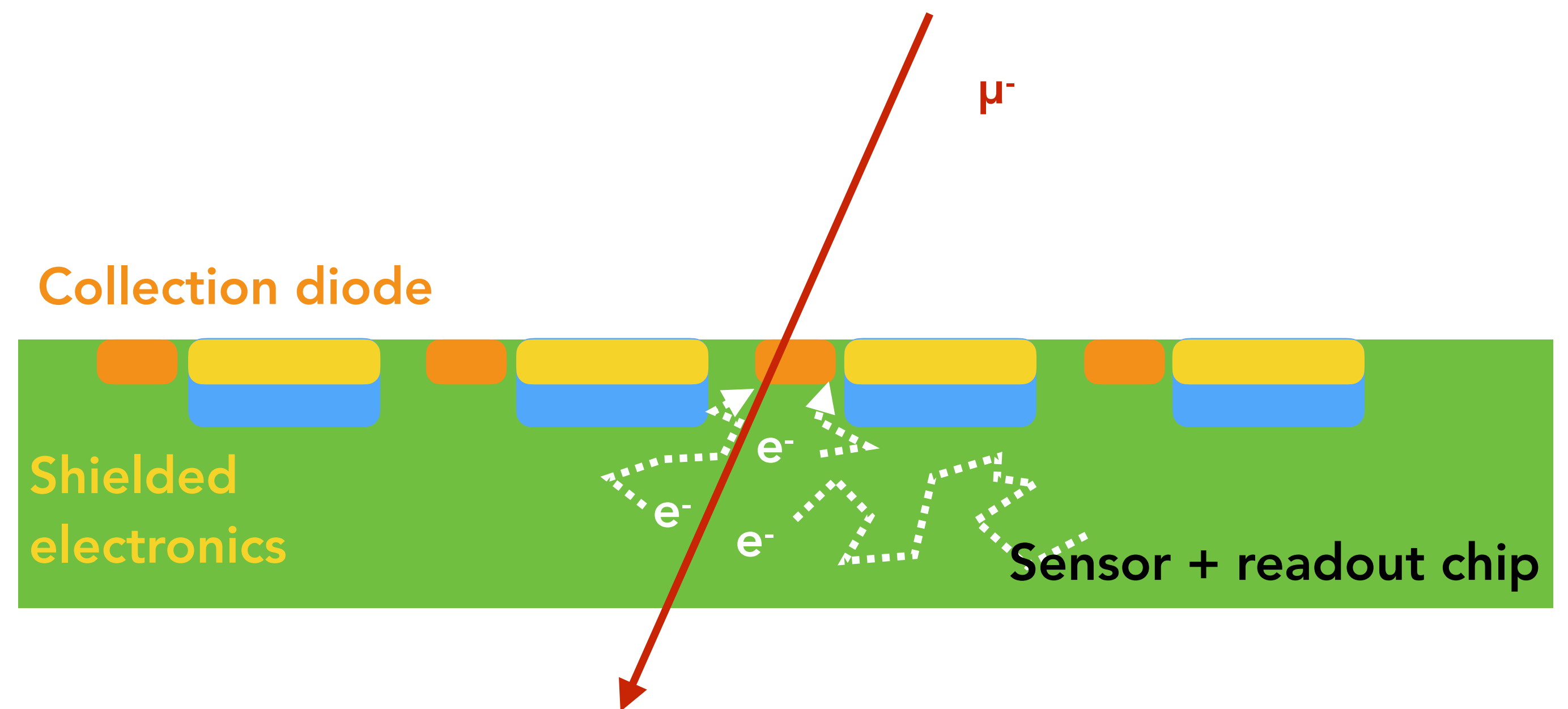
In the last 15 years, monolithic devices have evolved substantially, taking advantage of new processes within the microelectronics industry



Monolithic detector - MAPS

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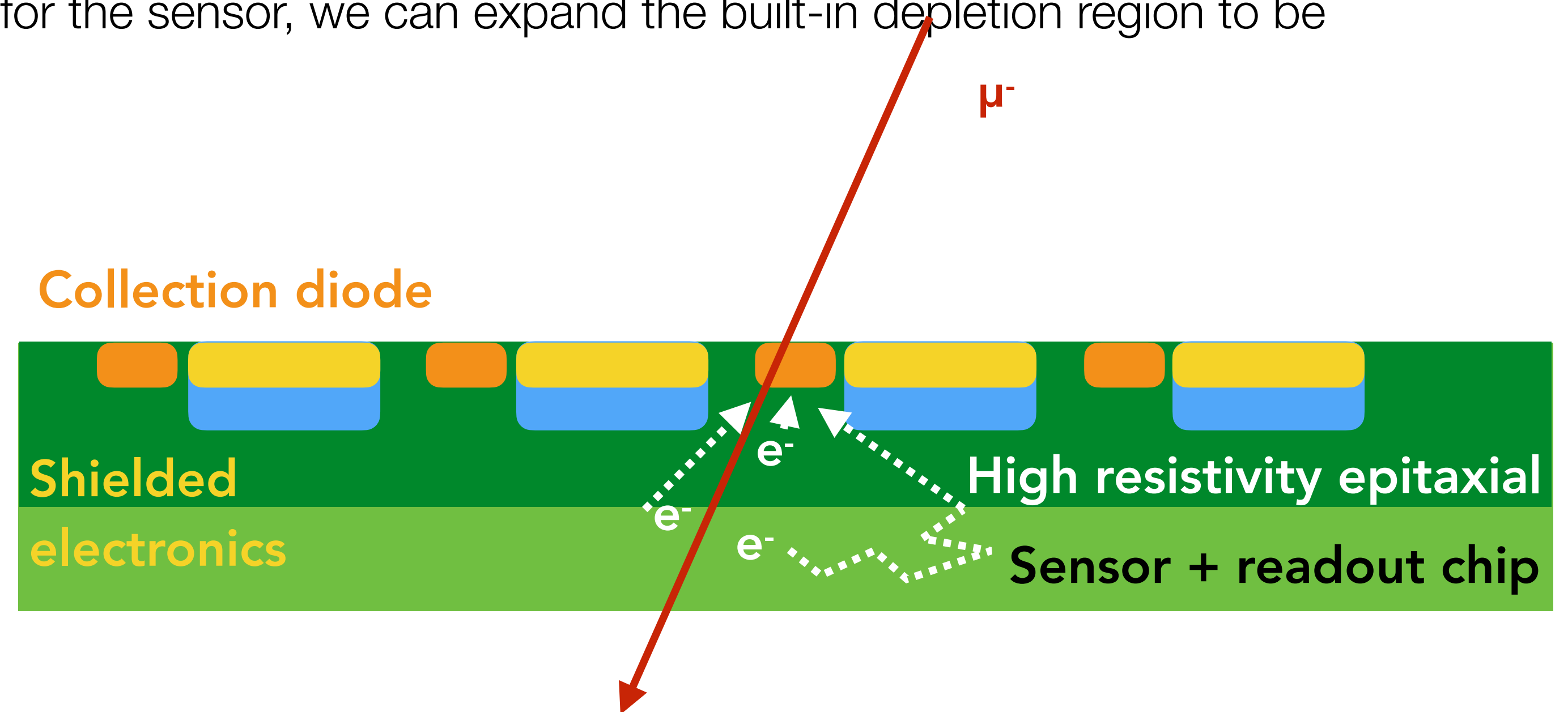
- By shielding our electronics, we can now prevent competition for charge collection between the electronics wells and the collection node



Monolithic detector - HR CMOS

In the last 10 years, monolithic devices have evolved substantially, taking advantage of new processes within the microelectronics industry

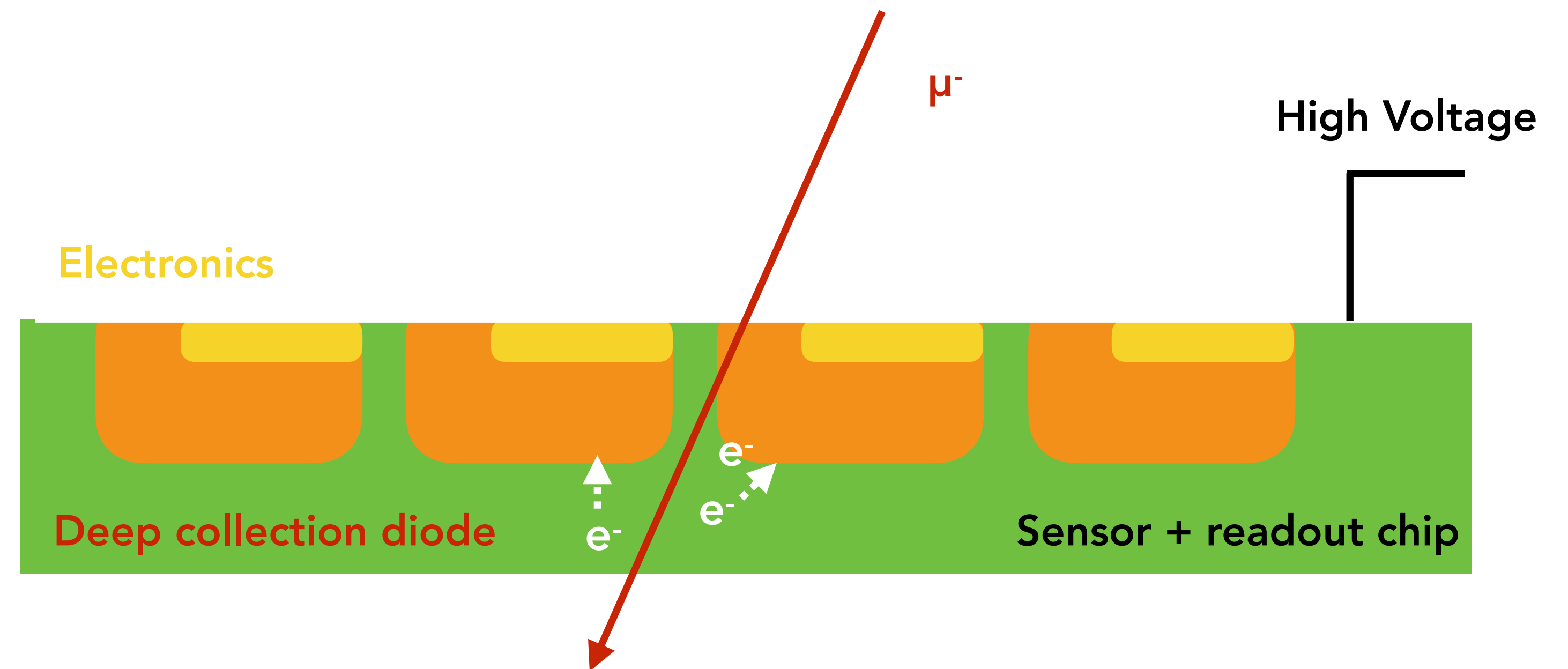
- By shielding our electronics, we can now prevent competition for charge collection between the electronics wells and the collection node
 - By using a high-resistivity epitaxial layer for the sensor, we can expand the built-in depletion region to be substantial in size (10 - 15 μm)
 - These are termed **High Resistivity CMOS** detectors
- 
- The diagram shows a cross-section of a detector. At the top is a thin blue layer labeled 'Passivation'. Below it is a thick yellow layer labeled 'High resistivity epitaxial layer'. At the bottom is a thin grey layer labeled 'Substrate'. A red arrow points upwards from the substrate, passing through the epitaxial layer, and is labeled 'Collection diode' in red. A red dot is located within the epitaxial layer, labeled μ^- in red.



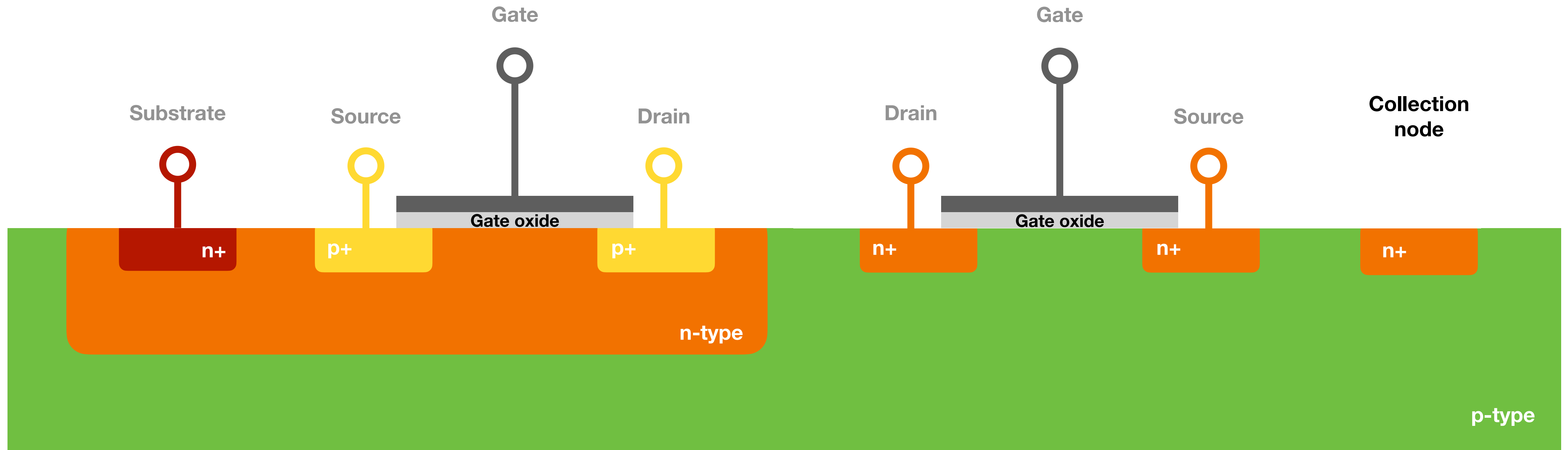
Monolithic detector - HV CMOS

An alternative approach is to tackle the other issue: applying a reverse bias voltage while protecting the electronics

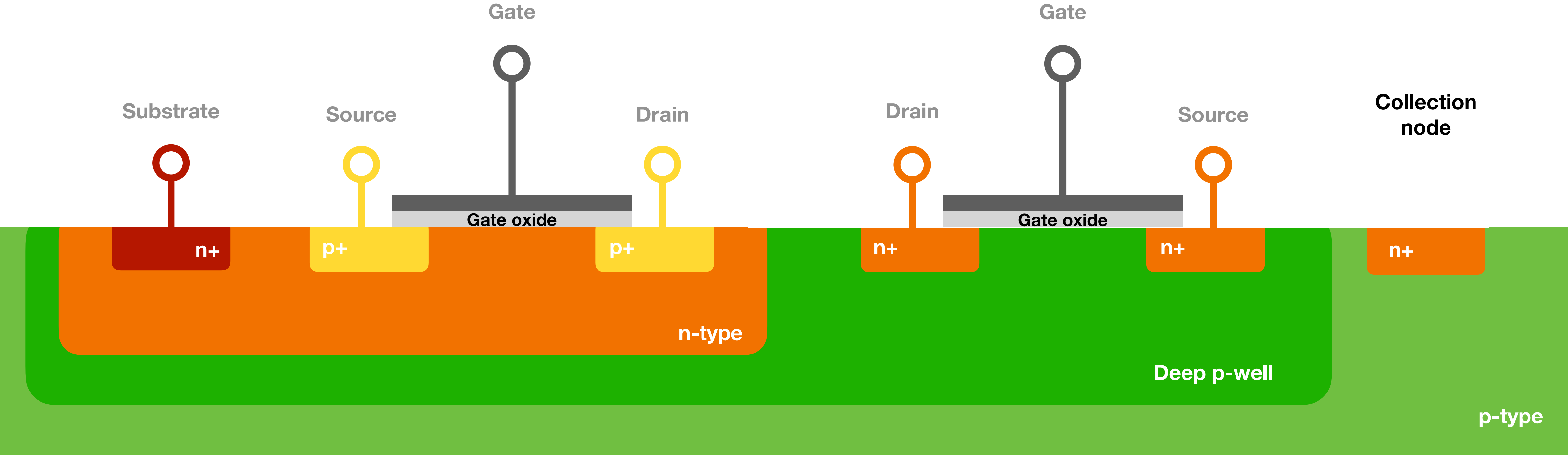
- Many companies have developed deep n-wells to function at high voltages
- We turn our simple collection node into one of these deep n-wells and “hide” all of the electronics inside
- Imaginatively termed **High Voltage CMOS**



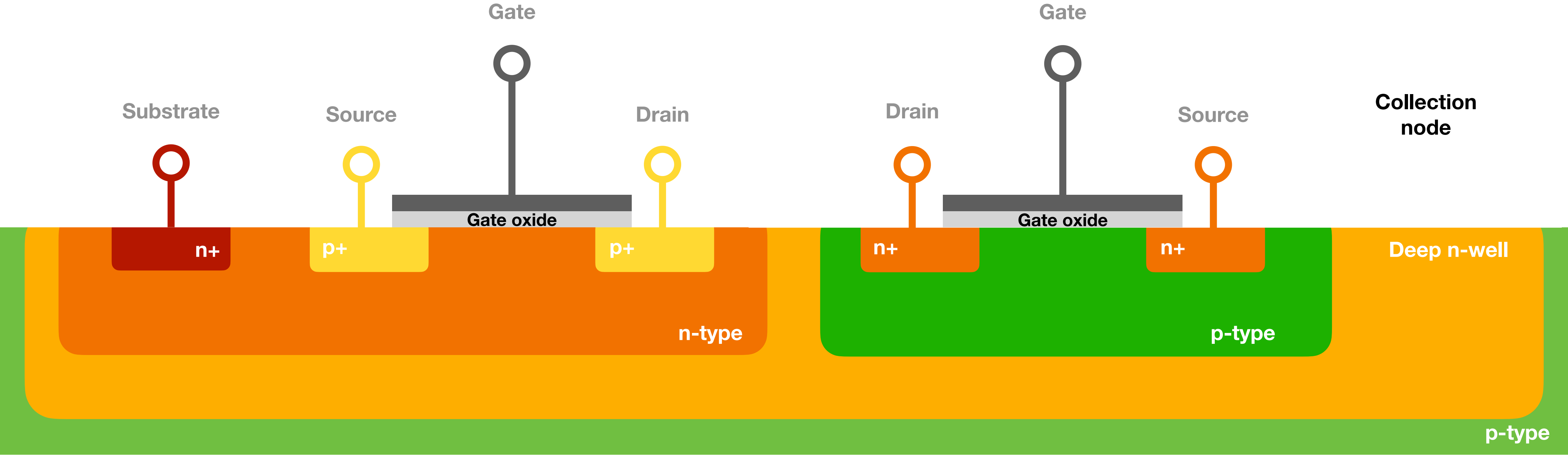
MAPS detectors - classic



HR-CMOS



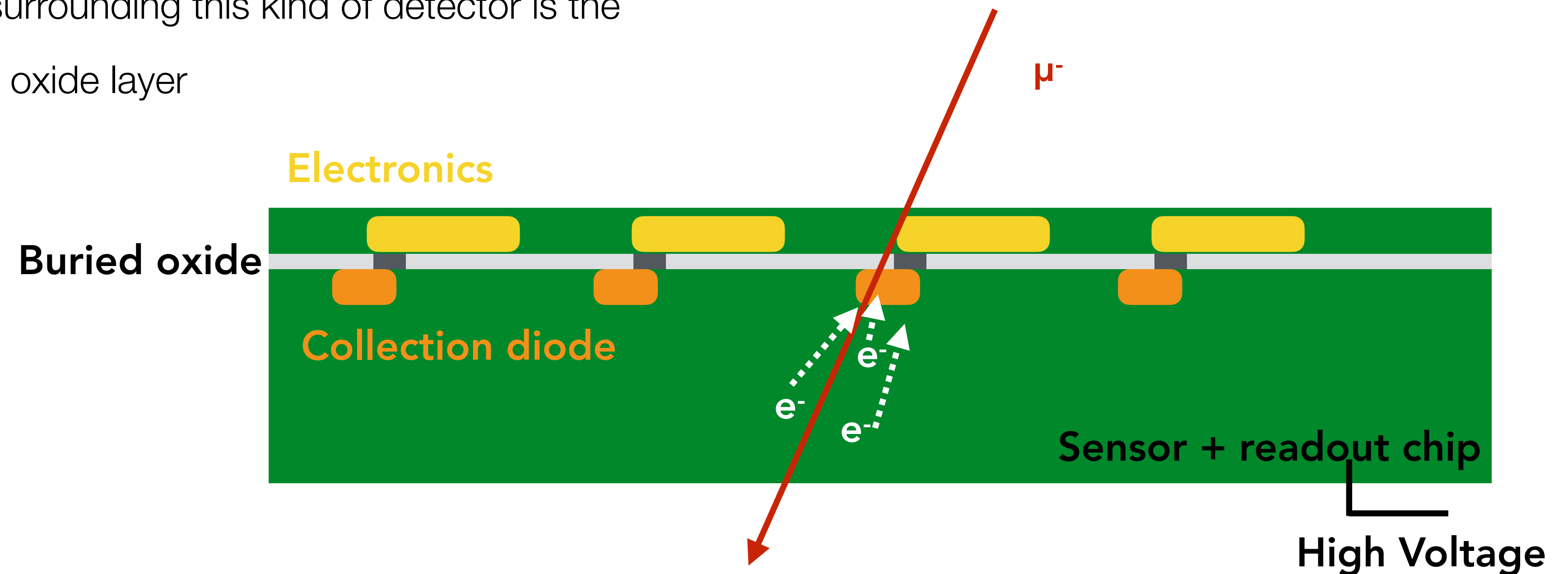
HV-CMOS



Monolithic detector - SOI

Another experimental (but not widely used) approach to solving this problem is to “separate” the silicon containing the collection node from the silicon containing the electronics

- An oxide layer is implanted in order to act as an insulating barrier, with vias connecting the electronics
- Such sensors are called **Silicon on Insulators**
- One of the main issues surrounding this kind of detector is the build-up of charge in the oxide layer



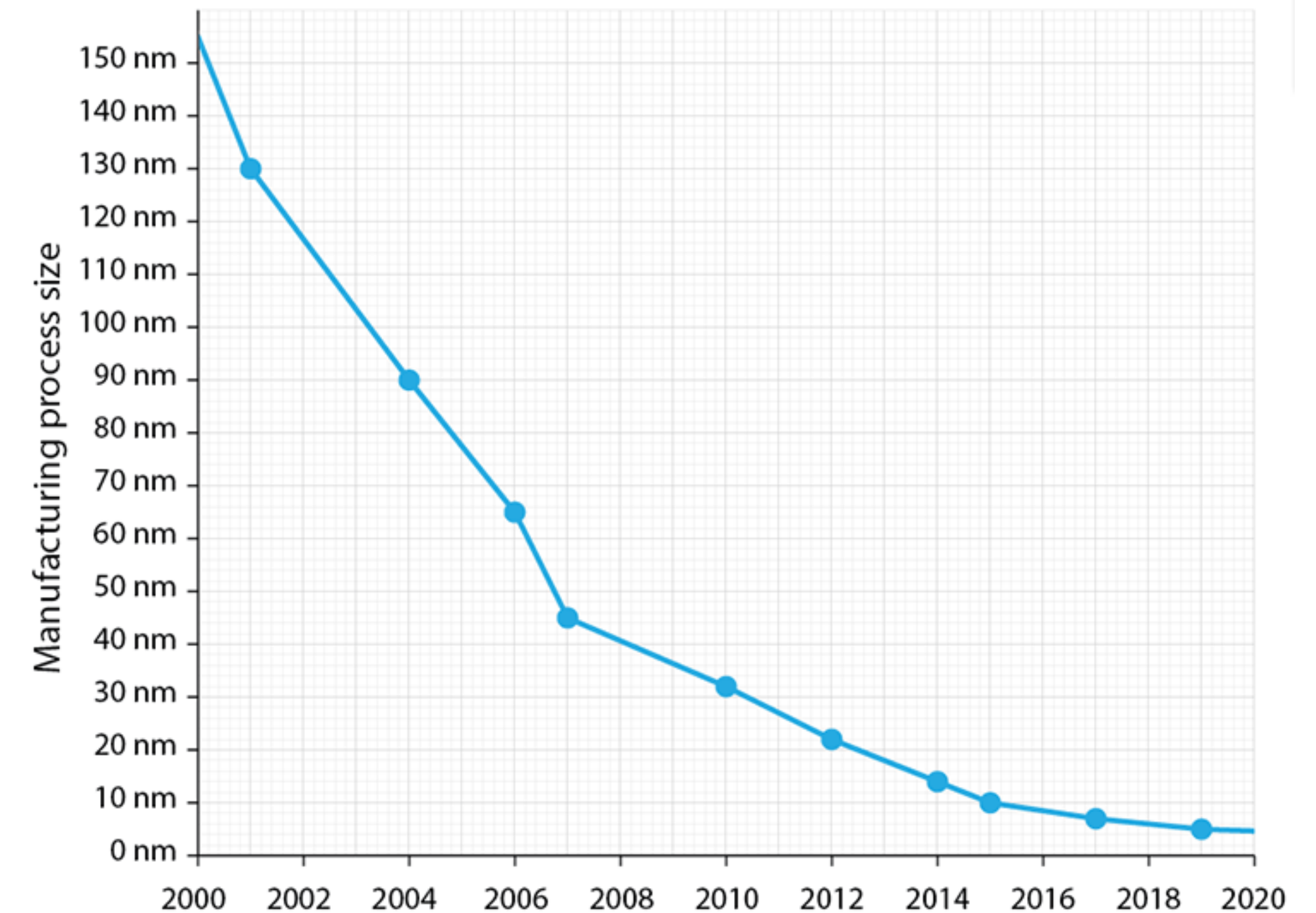
Technology nodes

Microelectronics industry moves much faster (and has much more money) than we do...

- Devices on the market just now with **4 nm** feature size
- In particle physics we are preparing to install the first readout electronics with 65 nm in 2027

Some trends have helped us - smaller devices and thinner oxides give better inherent TID tolerance

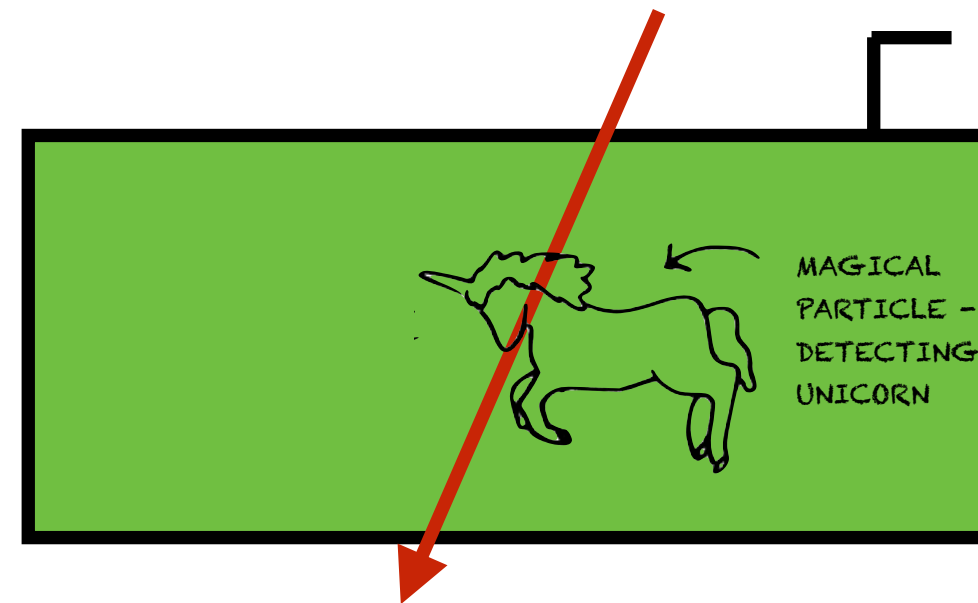
- However, below 28 nm transistor layout changes drastically (FinFETs) - performance TBD



Summary

What kind of detector to choose?

Unfortunately there is no quick answer other than:

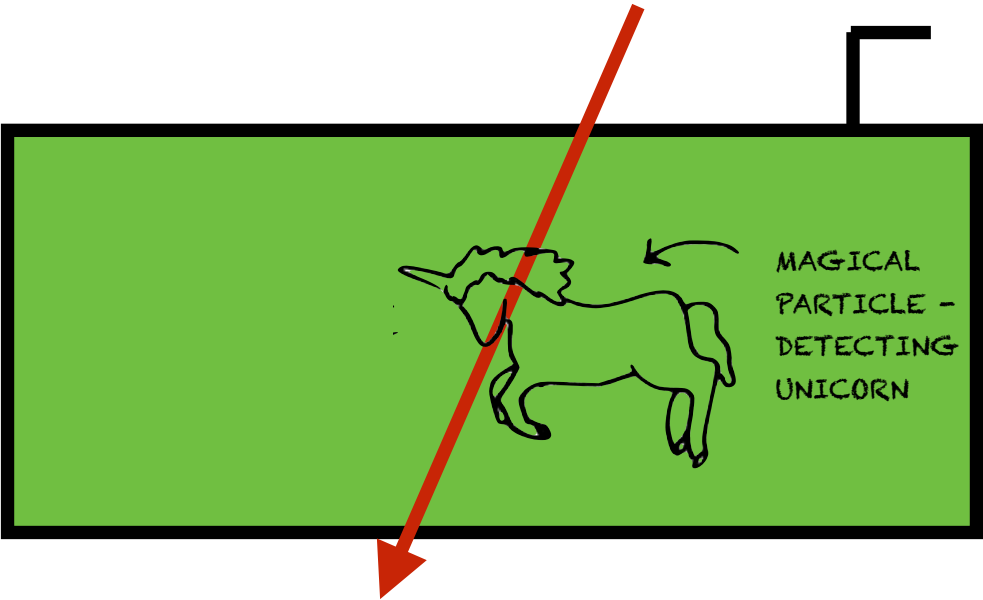


Some broad statements can guide:

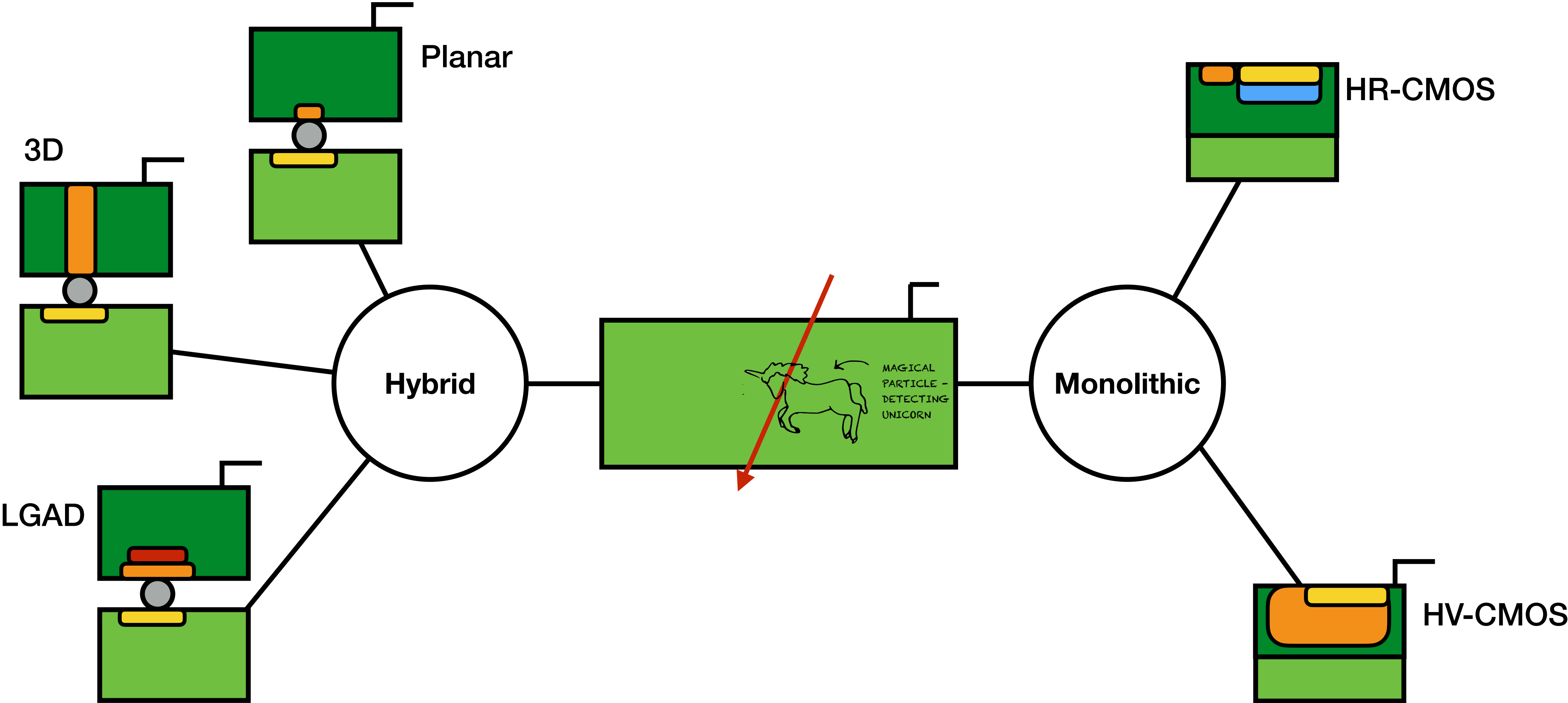
- Hybrid pixel detectors are necessary if you need high-Z sensors (X-rays, gamma rays)
- Monolithic detectors are typically cheaper than hybrid
- Built-in amplification if you have low signal or need very fast timing

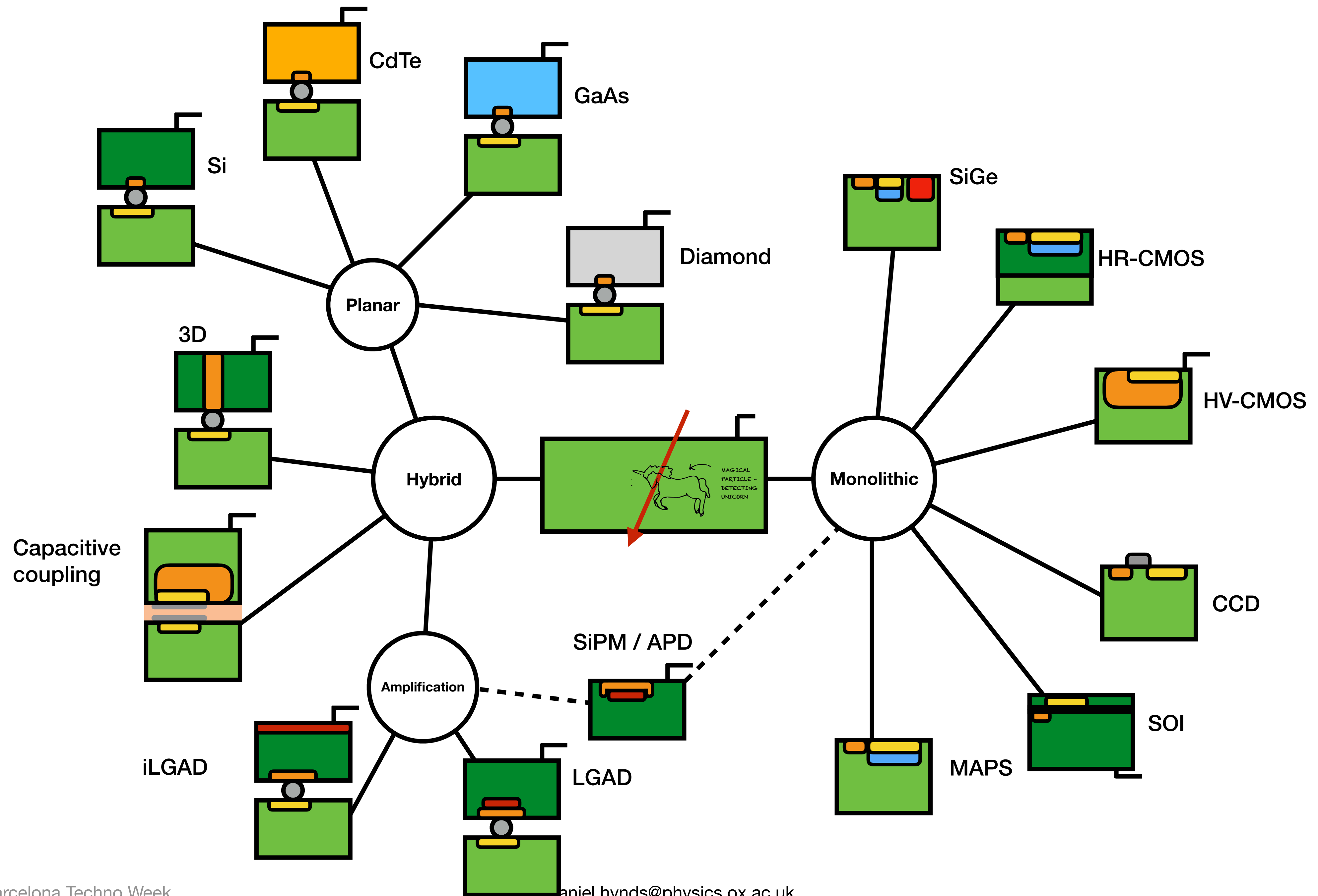
But the details will very much depend on the application

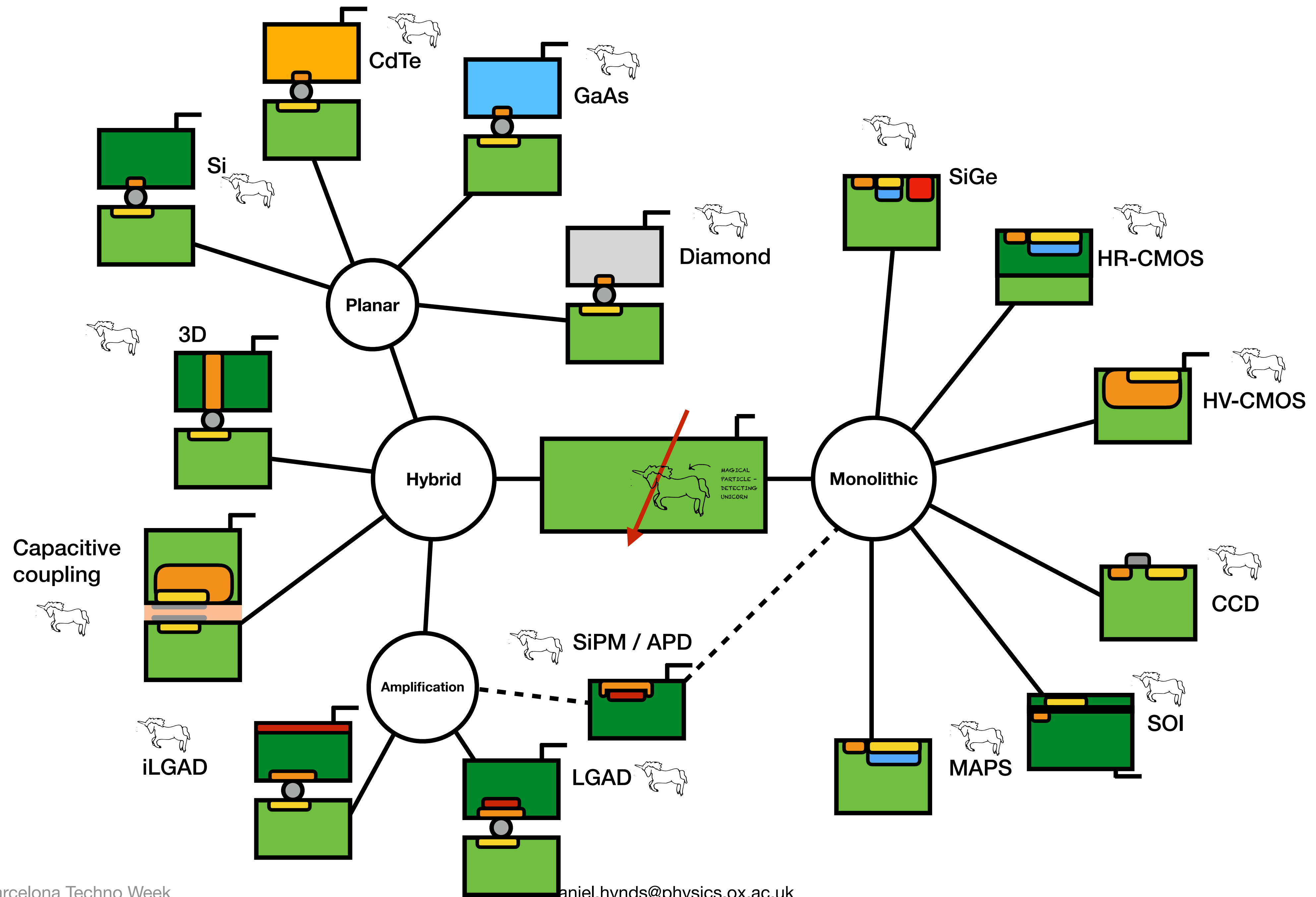
The silicon zoo

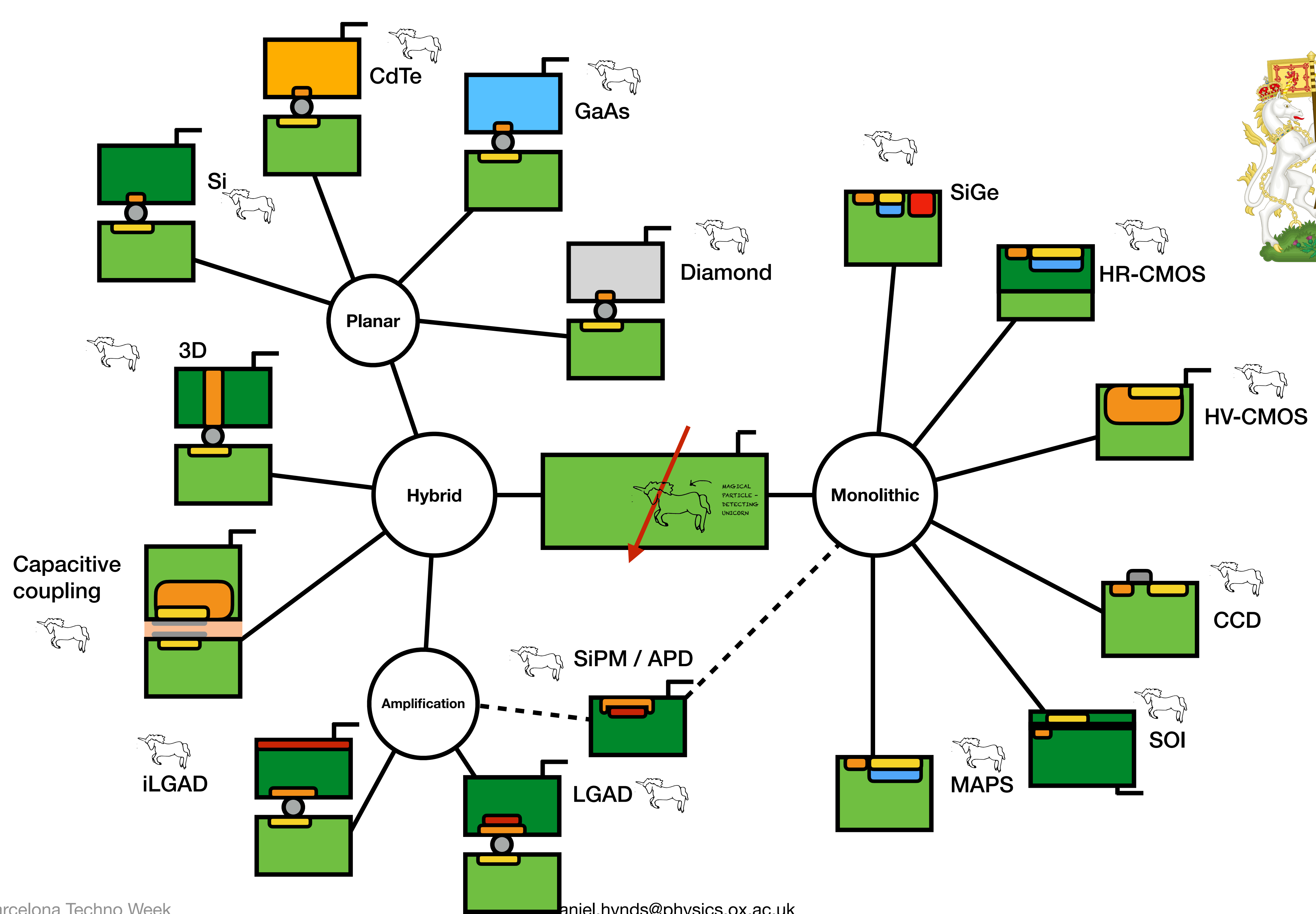


The silicon zoo









To finish

Eventually, I started to realize that science could explain how EVERYTHING worked.

There was no longer any need for my magical unicorns. Science killed my unicorns.

$$\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$\Delta S_{\text{UNIVERSE}} > 0$$

$$\Sigma F_y = T - mg = ma_y$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$



To finish

The physical mechanisms underlying the response of advanced CMOS microelectronics are now relatively well known in broad outline, but they may never be known in full detail. The models for interface trap production, and for the oxide hole trap, both date from the 1980s, although some of the confirming experiments were done in the 1990's. But what the models tell us is that the overall response is extremely complicated. There are multiple processes involved, with complicated time dependences, and different field dependences, and different temperature dependences. How all the pieces fit together will always be a complicated story. Anytime the manufacturing processes change, trench oxides replacing planar isolation structures, or new dielectrics entering production, etc., many of the details of the response will have to be worked out again. Considering how often new manufacturing techniques are introduced, there is little chance the community will ever run out of things to do.

Thank you!