

Perturbative and colorful lectures on **Strong Interactions**

lecture 1/3

Laurent Favart

IIHE - Université libre de Bruxelles

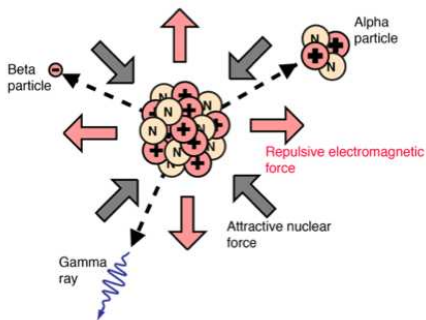
Sinop Gerze Deydaa Particle Physics Summer School

August 18, 2026



- 1 A few words about QCD
 - 1 running of α_S
 - 2 asymptotic freedom and confinement
- 2 quarks and gluons from e^+e^- annihilation
 - 1 quark electric charges
 - 2 quark and gluon spins
 - 3 QCD gauge parameters
- 3 Structure of hadrons
 - 1 elastic scattering
 - 2 deep inelastic scattering
 - 3 proton structure functions
 - 4 DGLAP evolution equation
- 4 hadron - hadron interactions
 - 1 jet production
 - 2 Drell-Yan process

The so rich nucleus

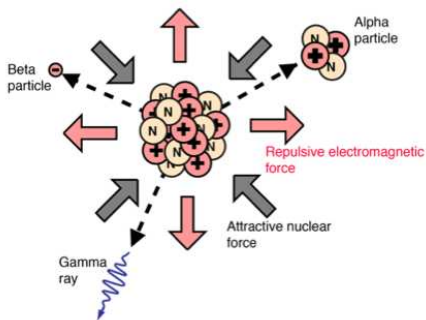


What makes the nucleons hold together in a nucleus ?

$${}^A_ZM = Z \cdot m_p + (A - Z) \cdot m_n - E_L \quad \Rightarrow \quad E_L \simeq 8 \text{ MeV}$$

$\Rightarrow$ named **Strong interaction**

The so rich nucleus



What makes the nucleons hold together in a nucleus ?

$$\frac{A}{Z}M = Z \cdot m_p + (A - Z) \cdot m_n - E_L \quad \Rightarrow \quad E_L \simeq 8 \text{ MeV}$$

$\Rightarrow$ named **Strong interaction**

In a nucleus there is a competition between EM, Strong and Weak interactions

- EM: repulsion between protons
- Strong: bounding between nucleons
- Weak and EM decays

A few preliminary remarks

	proton	neutron
mass	938.280 MeV	939.573 MeV
lifetime	stable	898 ± 16 sec
charge	+1	0
spin	$1/2$	$1/2$
magnetic moment	$2.793\mu_N$	$-1.913\mu_N$

A few preliminary remarks

lifetimes and masses

- $\tau_n \simeq 15$ min - **very long** for wi
 $n \rightarrow p e^- \bar{\nu}_e$

- $\tau_p > 10^{33}$ y - **stable** (if free state)

- $m_n/m_p = 1.0014$!
 $\Rightarrow$ very similar internal bounding fields

Huge bounding energy: GeV !

	proton	neutron
mass	938.280 MeV	939.573 MeV
lifetime	stable	898 ± 16 sec
charge	+1	0
spin	1/2	1/2
magnetic moment	$2.793\mu_N$	$-1.913\mu_N$

A few preliminary remarks

lifetimes and masses

- $\tau_n \simeq 15$ min - **very long** for wi
 $n \rightarrow p e^- \bar{\nu}_e$

- $\tau_p > 10^{33}$ y - **stable** (if free state)

- $m_n/m_p = 1.0014$!
 $\Rightarrow$ very similar internal bounding fields

Huge bounding energy: GeV !

- if $m_n \searrow$: $He/H \nearrow$ - $\nexists$ stars like $\odot$

- if $m_n \nearrow$: all atoms unstable (except H)

	proton	neutron
mass	938.280 MeV	939.573 MeV
lifetime	stable	898 ± 16 sec
charge	+1	0
spin	1/2	1/2
magnetic moment	$2.793\mu_N$	$-1.913\mu_N$

A few preliminary remarks

	proton	neutron
mass	938.280 MeV	939.573 MeV
lifetime	stable	898 ± 16 sec
charge	+1	0
spin	1/2	1/2
magnetic moment	$2.793\mu_N$	$-1.913\mu_N$

magnetic moment

Bohr magneton: $\mu_N = e\hbar/2m_Nc$

we would naively expect $\mu_p = 1$ and $\mu_n = 0$

$\Rightarrow$ first sign of nucleon charged sub-structure !

Proton structure - and QCD

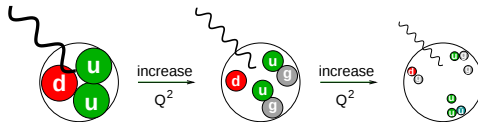
- in baryons, the 3 **valence quarks** are bound by gluon exchanges
- the gluons carry a color and an anti-color (or a linear combination), e.g.

$$r\bar{g}, \quad r\bar{b}, \quad g\bar{r}, \quad g\bar{b}, \quad b\bar{r}, \quad b\bar{g}, \quad \sqrt{1/2}\{r\bar{r} - g\bar{g}\}, \quad \sqrt{1/6}\{r\bar{r} + g\bar{g} - 2b\bar{b}\}$$

but the color singlet state is excluded by hand to avoid free gluons

$$G_0 = \sqrt{1/3}\{r\bar{r} + g\bar{g} + b\bar{b}\}$$

- the proton structure depends on the resolution scale - seeing more or less quantum fluctuations in $q\bar{q}$ (**sea quarks**)



- the QCD coupling constant, α_S , also depends on the resolution scale

- Part 1 -

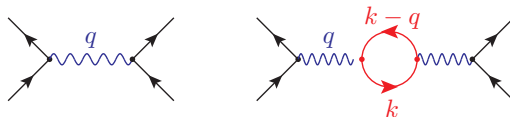
α_S running, asymptotic freedom and
confinement

Screening effect in QED

- in QFT the vacuum fluctuates in virtual particle-antiparticle pairs
- in presence of an electric charge, these pairs get polarised
- this leads to an effective charge which depends on the distance to the probe

At LO :

$$(-i)e_0\gamma^\mu \cdot (-i)\frac{g_{\mu\nu}}{q^2} \cdot (-i)e_0\gamma^\nu = ie_0^2\gamma^\mu\frac{g_{\mu\nu}}{q^2}\gamma^\nu$$



with one loop :

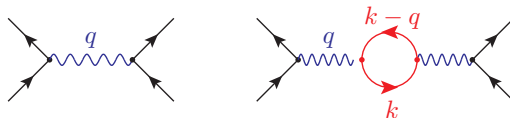
$$(-i)e_0\gamma^\mu (-i)\frac{g_{\mu\rho}}{q^2} (-1) \int \frac{d^4 k}{(2\pi)^4} \text{Tr} \left[(-i)e_0\gamma^\rho \frac{i(\not{k} + m)}{k^2 - m^2} (-i)e_0\gamma^\lambda \frac{i(\not{k} - \not{q} + m)}{(k - q)^2 - m^2} \right] (-i)\frac{g_{\lambda\nu}}{q^2} (-i)e_0\gamma^\nu$$

Screening effect in QED

- in QFT the vacuum fluctuates in virtual particle-antiparticle pairs
- in presence of an electric charge, these pairs get polarised
- this leads to an effective charge which depends on the distance to the probe

At LO :

$$(-i)e_0\gamma^\mu \cdot (-i)\frac{g_{\mu\nu}}{q^2} \cdot (-i)e_0\gamma^\nu = ie_0^2\gamma^\mu\frac{g_{\mu\nu}}{q^2}\gamma^\nu$$



with one loop :

$$(-i)e_0\gamma^\mu (-i)\frac{g_{\mu\rho}}{q^2} \underbrace{(-1)}_{\text{fermion loop}} \int \frac{d^4 k}{(2\pi)^4} \text{Tr} \left[(-i)e_0\gamma^\rho \frac{i(\not{k} + m)}{k^2 - m^2} (-i)e_0\gamma^\lambda \frac{i(\not{k} - \not{q} + m)}{(k - q)^2 - m^2} \right] (-i)\frac{g_{\lambda\nu}}{q^2} (-i)e_0\gamma^\nu$$

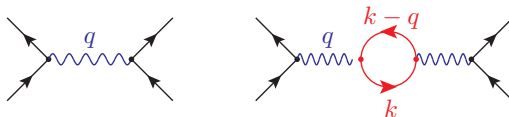
fermion loop

Screening effect in QED

- in QFT the vacuum fluctuates in virtual particle-antiparticle pairs
- in presence of an electric charge, these pairs get polarised
- this leads to an effective charge which depends on the distance to the probe

At LO :

$$(-i)e_0\gamma^\mu \cdot (-i)\frac{g_{\mu\nu}}{q^2} \cdot (-i)e_0\gamma^\nu = ie_0^2\gamma^\mu\frac{g_{\mu\nu}}{q^2}\gamma^\nu$$



with one loop :

$$(-i)e_0\gamma^\mu (-i)\frac{g_{\mu\rho}}{q^2} (-1) \int \frac{d^4 k}{(2\pi)^4} \text{Tr} \left[(-i)e_0\gamma^\rho \frac{i(\not{k} + m)}{k^2 - m^2} (-i)e_0\gamma^\lambda \frac{i(\not{k} - \not{q} + m)}{(k - q)^2 - m^2} \right] (-i)\frac{g_{\lambda\nu}}{q^2} (-i)e_0\gamma^\nu$$

trace because we sum over fermions spin states

does the trace converge ?

$$\int \frac{d^4 k}{(k^2 - m^2)^2} = \int \frac{k^3}{(k^2 - m^2)^2} dk d\Omega$$

divergent for $k \rightarrow \infty$

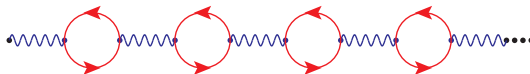
$$\int^{\infty} \frac{dk}{k} \rightarrow \int^{\mu_R} \frac{dk}{k} \quad \rightarrow \text{we introduce a cutoff}$$

the integral give a log, the **effective** charge is given by:

$$e_{eff}^2(1loop) = ie_0^2 \gamma^\mu \frac{g_{\mu\nu}}{q^2} \gamma^\nu \left[1 + \frac{e_0^2}{12\pi^2} \ln \left(\frac{m^2 - q^2}{\mu_R^2} \right) + \frac{e_0^2}{12\pi^2} F(q^2) \right]$$

where $F(q^2)$ is a finite function vanishing for $q^2 \rightarrow \infty$.

Adding more loops

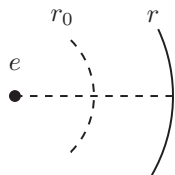


$$e_{eff}^2 = e_0^2 \left[1 + \frac{e_0^2}{3\pi} \ln \left(\frac{m^2 - q^2}{\mu_{UV}^2} \right) + \left(\frac{e_0^2}{3\pi} \ln \left(\frac{m^2 - q^2}{\mu_{UV}^2} \right) \right)^2 + \dots \right]$$

behaves arithmetic sequence, whose sum is (using $Q^2 = -q^2$):

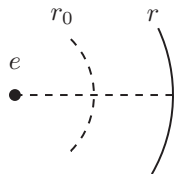
$$e_{eff}^2(Q^2) = \frac{e_0^2}{1 - \frac{e_0^2}{3\pi} \ln \frac{Q^2}{\mu_R^2}}$$

Interpretation : using the distance $r^2 \sim 1/Q^2$ (sign in front of $\ln$ changes)



$$e^2(r) = \frac{e^2(r_0)}{1 + \frac{2e^2(r_0)}{3\pi} \ln \frac{r}{r_0}}$$

Interpretation : using the distance $r^2 \sim 1/Q^2$ (sign in front of $\ln$ changes)



$$e^2(r) = \frac{e^2(r_0)}{1 + \frac{2e^2(r_0)}{3\pi} \ln \frac{r}{r_0}}$$

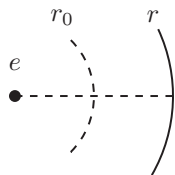
if $r \gg r_0$



$$\simeq \frac{e^2(r_0)}{e^2(r_0) \dots} \rightarrow \alpha = 1/137$$

independ. of r_0

Interpretation : using the distance $r^2 \sim 1/Q^2$ (sign in front of $\ln$ changes)



$$e^2(r) = \frac{e^2(r_0)}{1 + \frac{2e^2(r_0)}{3\pi} \ln \frac{r}{r_0}}$$

if $r \gg r_0$



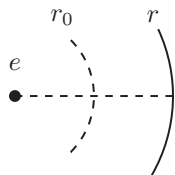
$$\simeq \frac{e^2(r_0)}{e^2(r_0) \dots} \rightarrow \alpha = 1/137 \text{ independ. of } r_0$$



if $r_0 \rightarrow 0$

$$\simeq \frac{e^2(r_0)}{\infty} \text{ only possible if } e^2(r_0 \rightarrow 0) \rightarrow \infty$$

Interpretation : using the distance $r^2 \sim 1/Q^2$ (sign in front of $\ln$ changes)



$$e^2(r) = \frac{e^2(r_0)}{1 + \frac{2e^2(r_0)}{3\pi} \ln \frac{r}{r_0}}$$

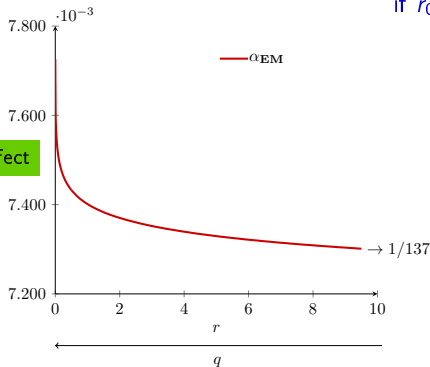
if $r \gg r_0$

$$\simeq \frac{e^2(r_0)}{e^2(r_0) \dots} \rightarrow \alpha = 1/137 \text{ independ. of } r_0$$

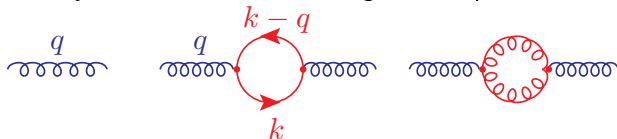
if $r_0 \rightarrow 0$

$$\simeq \frac{e^2(r_0)}{\infty} \text{ only possible if } e^2(r_0 \rightarrow 0) \rightarrow \infty$$

screening effect



in QCD: situation partially similar, on additional gluon loop

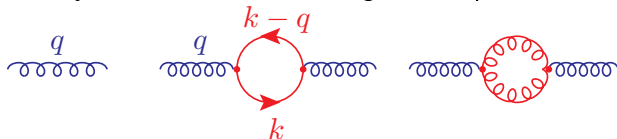


- the **fermion loop** (N_f) is similar to the QED case
- the additional **gluon loop** comes with a **(+1)**
- sum over N_c colors

$$\alpha_s(Q^2) = \frac{\alpha_s(\mu_R^2)}{1 - \frac{2N_f - 11N_c}{6\pi} \alpha_s(\mu_R^2) \ln \frac{Q^2}{\mu_R^2}}$$

- competition between N_f and N_c
- as $b = 2N_f - 11N_c < 0 \Rightarrow$ anti-screening effect

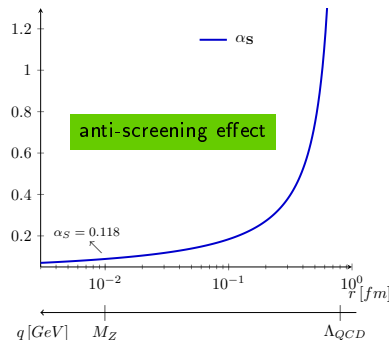
in QCD: situation partially similar, on additional gluon loop



- the **fermion loop** (N_f) is similar to the QED case
- the additional **gluon loop** comes with a **(+1)**
- sum over N_c colors

$$\alpha_s(Q^2) = \frac{\alpha_s(\mu_R^2)}{1 - \frac{2N_f - 11N_c}{6\pi} \alpha_s(\mu_R^2) \ln \frac{Q^2}{\mu_R^2}}$$

- competition between N_f and N_c
- as $b = 2N_f - 11N_c < 0 \Rightarrow$ anti-screening effect



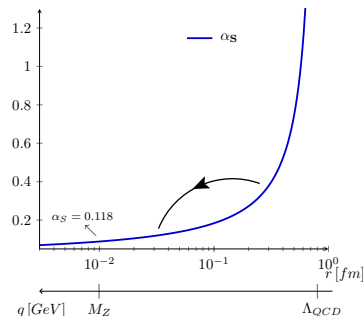
- QFT gives us the **evolution** but not the values
- in one direction or the other

$$\alpha_s(Q^2) = \frac{\alpha_s(\mu_R^2)}{1 - \frac{2N_f - 11N_c}{6\pi} \alpha_s(\mu_R^2) \ln \frac{Q^2}{\mu_R^2}}$$

- Λ_{QCD} is defined as the divergent scale value:

$$\alpha_s(\Lambda_{QCD}^2) = \frac{\alpha_s(\mu_R^2)}{0}$$

- $\Lambda_{QCD} \simeq m_\pi \simeq 100 \text{ MeV}$



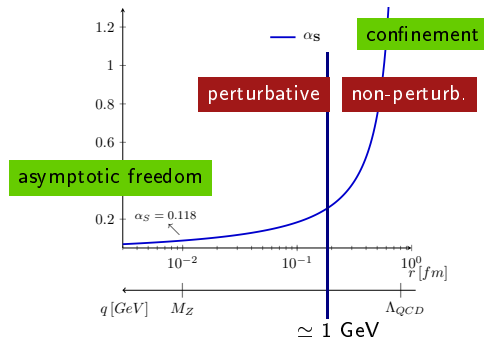
- QFT gives us the **evolution** but not the values
- in one direction or the other

$$\alpha_s(Q^2) = \frac{\alpha_s(\mu_R^2)}{1 - \frac{2N_f - 11N_c}{6\pi} \alpha_s(\mu_R^2) \ln \frac{Q^2}{\mu_R^2}}$$

- Λ_{QCD} is defined as the divergent scale value:

$$\alpha_s(\Lambda_{QCD}^2) = \frac{\alpha_s(\mu_R^2)}{0}$$

- $\Lambda_{QCD} \simeq m_\pi \simeq 100 \text{ MeV}$
- to apply perturbative methods: $\alpha_S \ll 1$
i.e. $Q \gg \Lambda_{QCD}$
- if non-pert. (no Fey. diag, etc) but $\mathcal{L}_{QCD}$ is valid

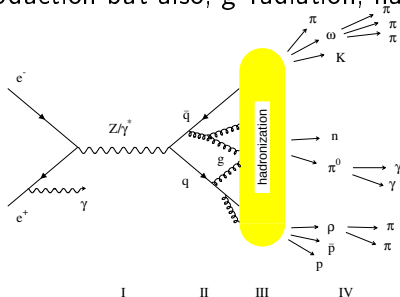


- Part 2 -

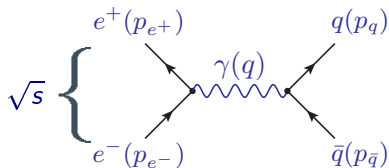
$e^+ e^-$ annihilation and QCD

quark production from e^-e^+ annihilation

- the **cleanest way** to produce quarks is from e^-e^+ annihilations:
 - **no color** charge in the initial state
 - pure **electroweak** process (at LO)
 - initial state momentum fully transferred to the $q\bar{q}$
- allows us to study q production but also, g radiation, hadronisation,...



- at LO, for γ exchange only:

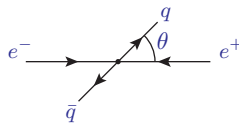


- negl. mass terms :

$$|\overline{\mathcal{M}}|^2 = 8(4\pi)^2 \frac{\alpha^2}{q^4} N_c Q_i^2 \{ (p_{e^+} \cdot p_q)(p_{e^-} \cdot p_{\bar{q}}) + (p_{e^+} \cdot p_{\bar{q}})(p_{e^-} \cdot p_q) \}$$

- only 1 degree of freedom
- in the c.m.s.:

$$\frac{d\sigma}{d\cos\theta} = \frac{\pi\alpha^2}{2s} N_c Q_i^2 (1 + \cos^2 \theta)$$



access to quark electric charges

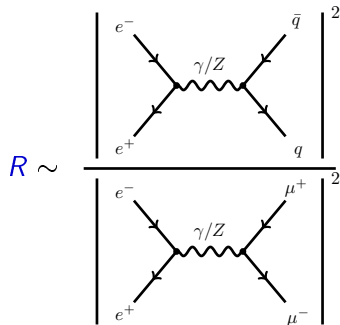
- as all hadrons come from quarks and that quarks always give hadrons :

$$\sigma(e^+e^- \rightarrow X(\text{any hadron})) = \sum_{i=1}^{N_f} \sigma(e^+e^- \rightarrow q\bar{q}) = \frac{4\pi\alpha^2}{3s} N_c \sum_{i=1}^{N_f} Q_i^2$$

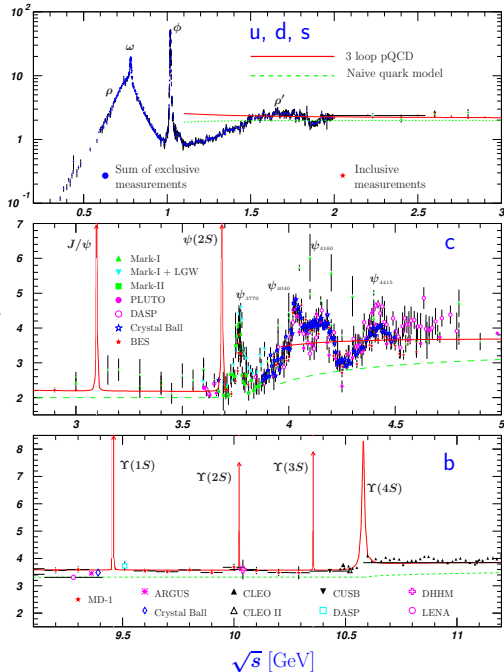
- we can study the observable:

$$R = \frac{\sigma(e^+e^- \rightarrow X)}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)} = \frac{\frac{4\pi\alpha^2}{3s} N_c \sum_{i=1}^{N_f} Q_i^2}{\frac{4\pi\alpha^2}{3s}} = N_c \sum_{i=1}^{N_f} Q_i^2$$

- is this very simple model realistic ?



R

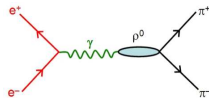


$$\begin{aligned}
 R &= \frac{\sigma(e^+e^- \rightarrow X)}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)} = N_c \sum_{i=1}^{N_f} Q_i^2 \\
 &= 3 \left[\left(\frac{2}{3}\right)^2 + \left(\frac{-1}{3}\right)^2 + \left(\frac{-1}{3}\right)^2 + \dots \right] \\
 &= 2 \quad \text{for } 2m_s \ll \sqrt{s} \ll 2m_c \\
 &= 10/3 \quad \text{for } 2m_c \ll \sqrt{s} \ll 2m_b \\
 &= 11/3 \quad \text{for } 2m_b \ll \sqrt{s} \ll m_Z
 \end{aligned}$$

- amazing success if away from resonances

quark charges are as expected

- if close to a resonance, the important npQCD corrections only in the numerator

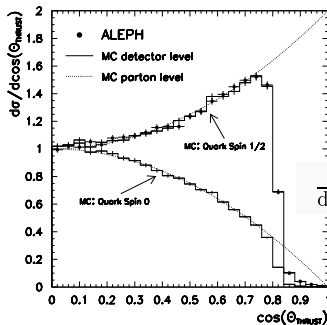
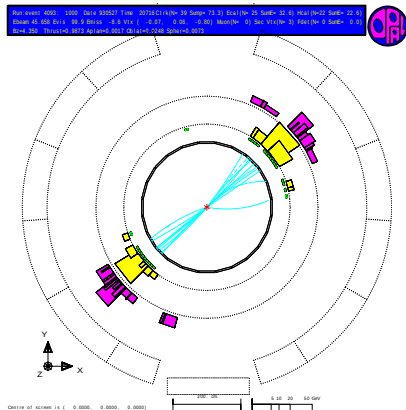
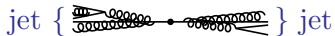


jets - during the had. process, hadrons are forming around the initial quark direction

- two particles below to different jets only if

$$p_i \cdot p_j \gg \Lambda_{QCD}^2$$

- angular measurement:

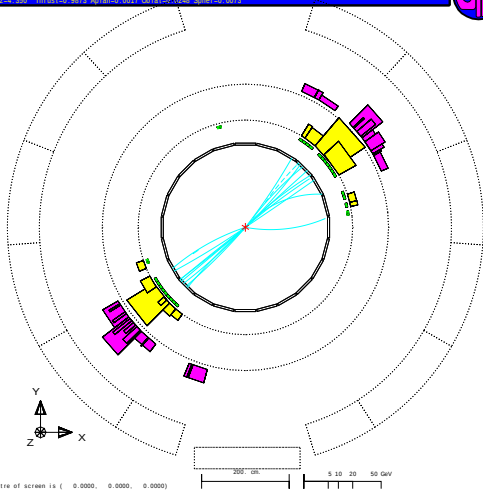


$$\frac{d\sigma}{d\cos\theta} \sim (1 + \cos^2 \theta)$$

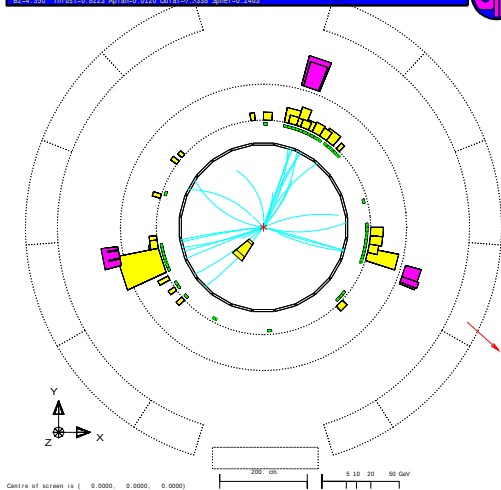
the quark are spin $1/2$!

- excellent agreement between jet and quark directions

Run event 4093 Date 930527 Time 20716 Clrk(N= 39 SunE= 73.3) Ecal(N= 25 SunE= 22.6) Hcal(N=22 SunE= 22.6)
 Ebeam 45.658 Evis 99.9 Emiss -8.6 Vtx (-0.07 0.06 -0.80) Muon(N= 0) Sec Vtx(N= 3) Fdet(N= 0 SunE= 0.0)
 Bz=4.350 Thrust=0.9873 Aplan=0.0017 Obstat=0.7348 Spher=0.6073



Run event 2542 Date 911014 Time 35925 Clrk(N= 28 SunE= 42.1) Ecal(N= 42 SunE= 59.8) Hcal(N= 8 SunE= 12.7)
 Ebeam 45.609 Evis 86.2 Emiss 5.0 Vtx (-0.05 0.12 -0.90) Muon(N= 1) Sec Vtx(N= 0) Fdet(N= 2 SunE= 0.0)
 Bz=4.350 Thrust=0.8223 Aplan=0.0120 Obstat=0.7339 Spher=0.2463



That's all for today