

Detectors for High Radiation and Extreme Environment

PSD13



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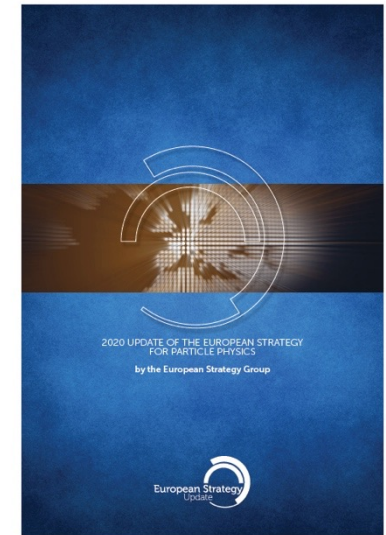
St. Catherine's College
September 3-8, 2023



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OXFORD

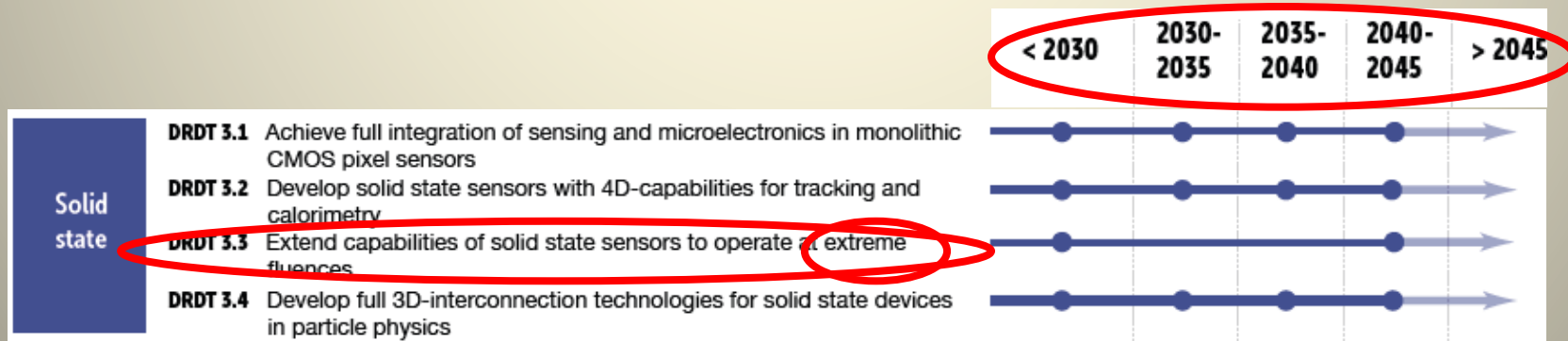
Future (Global) HEP Projects

- The 2020 *European Strategy for Particle Physics* establishes two project initiatives as high-priority
 - “the highest-priority next collider”: “an electron-positron Higgs factory”
 - for the longer term: “a proton-proton collider at the highest achievable energy”, dubbed as the *FCC-hh* project.
- End of 2021 the *ECFA Detector R&D Roadmap* was approved by the CERN Council
 - Long term HEP Detector R&D goals defined
 - Implementation strategy in terms of Detector R&D Collaborations (DRDC) worked out, starting in 2024
- Development cycle towards the use of a new technology in detectors spans over 10 to 20 years.
 - *prospective* detector R&D (“Blue Sky” research) – TRL 1
 - *guided* detector R&D, according to known needs of future projects – TRL 2-5
 - *focussed* detector R&D of approved experiments – TRL 5-7



What's in for PSD's ?

- PSD's are predominantly Solid State... DRD3
 - At least the ones for extreme radiation conditions
- 4 RD Theme's identified in DRD3
 - Monolithic CMOS
 - Precision timing → 4D tracking
 - 3D interconnects
- Operation at **extreme** fluences (DRDT 3.3)
- There could be additional/other extremes in detector operation (e.g. T , a , $p...$) – not addressed in this talk



Extreme ?

- What is extreme ?

extreme (ɪk'stri:m)

adj

1. being of a high or of the highest degree or intensity: *extreme cold*; *extreme difficulty*.
2. exceeding what is usual or reasonable; immoderate: *extreme behaviour*.

- A rather subjective measure

- For LHC 10^{15} n_{eq}/cm^2 was considered extreme

- design was 730/fb @14TeV...

- HL-LHC takes it to $nx10^{16}$ (vertex) or even 10^{17} (FW calo)

- 4000/fb @14TeV

- FCC-hh is *specifying* towards 10^{18} for the tracker (FCC-hh CDR)

- 30/ab @100TeV

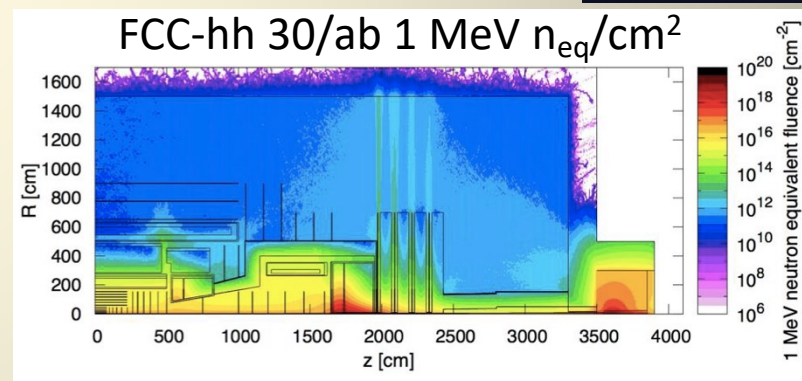
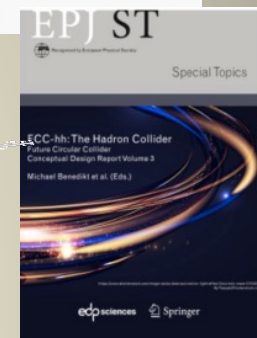
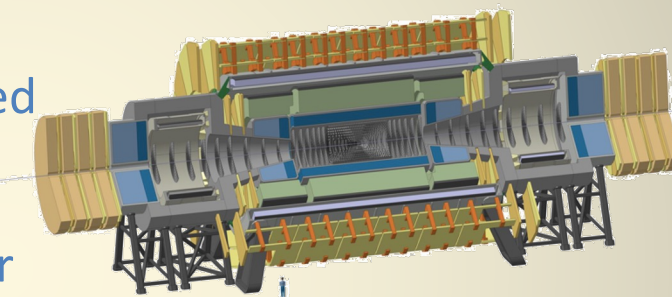
- 300 MGy TID in addition (not addressed)

- Ratio 1:20:600 !

- well, you *need* $\sim 7^2 \approx 50$ in HL/FCC lumi...

- What is the limit of tracking sensors ?

- TRIGA, NPP and ITER are $10^{21} \leftrightarrow 10^{24}$



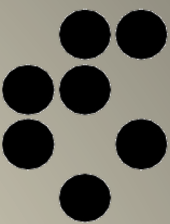
First tracking layer:

- 10 GHz/cm² charged particles
- 10^{18} hadrons/cm² for 30/ab

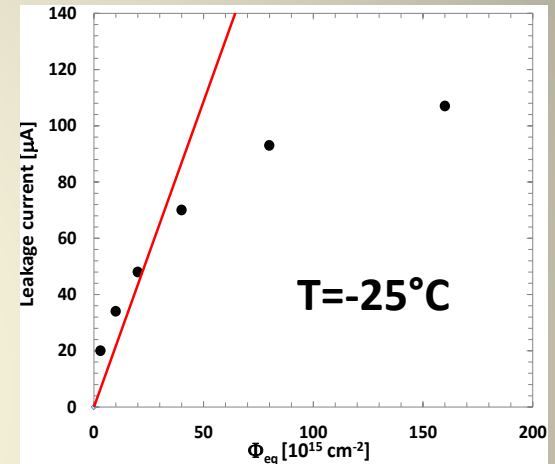
Expectations for $10^{17} n_{eq}/cm^2$

- For a ~yearly replacement of FCC-hh inner tracker !
 - Or a 2-stage operation 5->30/ab
- Linear extrapolation from low fluence data
 - Current: $I_{leak} = 4 \text{ A/cm}^3 @ 20^\circ\text{C}$
 - 2 mA/cm^2 (2W @ 1 kV) for 300 μm thick detector @ -20°C
 - Depletion: $N_{eff} \approx 1.5 \times 10^{15} \text{ cm}^{-3}$
 - $FDV \approx 100 \text{ kV}$
 - Trapping $\tau_{eff} \approx 1/40 \text{ ns} = 25 \text{ ps}$
 - $Q \approx Q_0/d v_{sat} \tau_{eff} \approx 80 \text{ e}/\mu\text{m} \cdot 200 \mu\text{m}/\text{ns} \cdot 1/40 \text{ ns} = 400 \text{ e}$ in very high electric field ($\gg 1 \text{ V}/\mu\text{m}$)
- Looks much like *Mission Impossible* (part n...)

CCE measurements up to $1.6 \times 10^{17} \text{ n}_{\text{eq}}/\text{cm}^2$



- n^+p "spaghetti" strips, 300 μm
- Observed signal not at all compatible with expectations
 - Above 3×10^{15} linear CCE(V_{bias})
 - Power law scaling with fluence, $b \approx -2/3$
 - Leakage current "saturating"

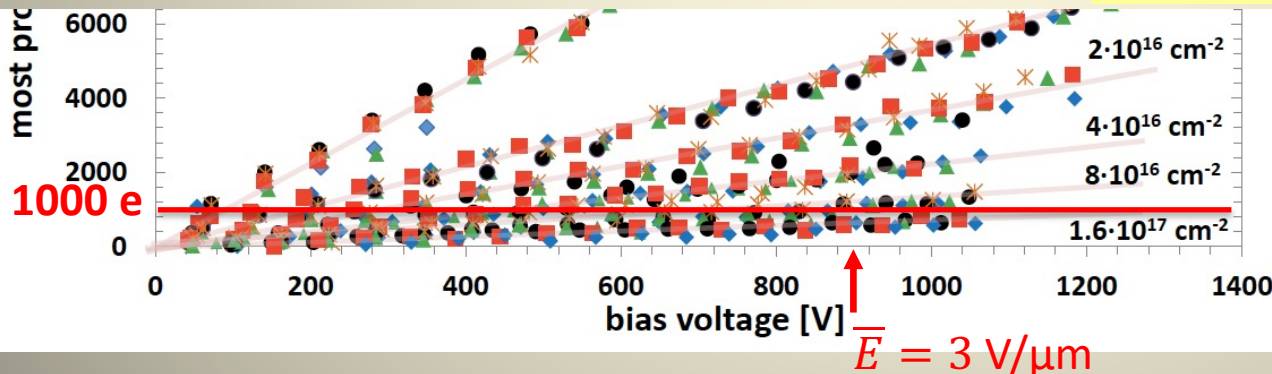


$$Q_{MPV}(V, \Phi) = k \cdot (\Phi / 10^{15} \text{ n}_{\text{eq}}/\text{cm}^2)^b \cdot V$$

$$k = 26.4 \text{ e}_0/\text{V}$$

$$b = -0.683$$

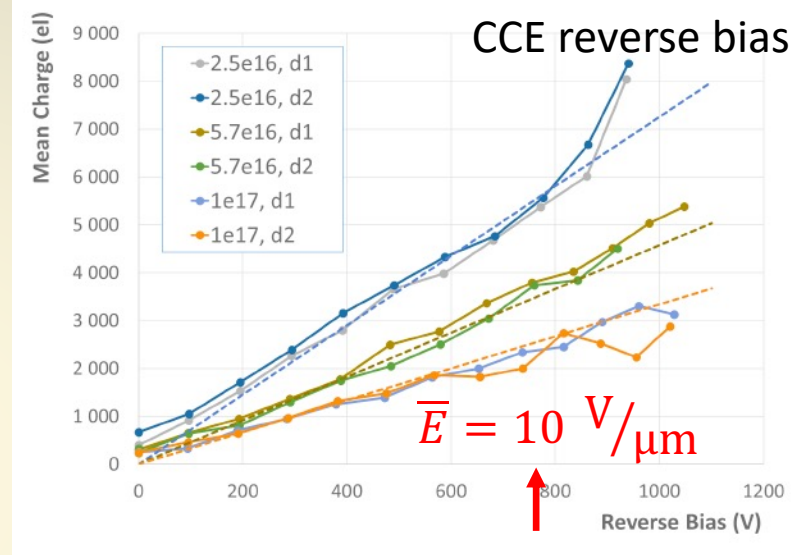
"Magic formula"



From:
**G. Kramberger et al.,
 JINST 8 P08004 (2013).**

More measurements on thin detectors

- 75 μm epi detectors from CNM on low-resistivity substrate
- Irradiated to 0.25, 0.57 and $1.0 \times 10^{17} \text{ n}_{\text{eq}}/\text{cm}^2$
- CCE in reverse and FW
- Annealing 1200 min @ 60°C



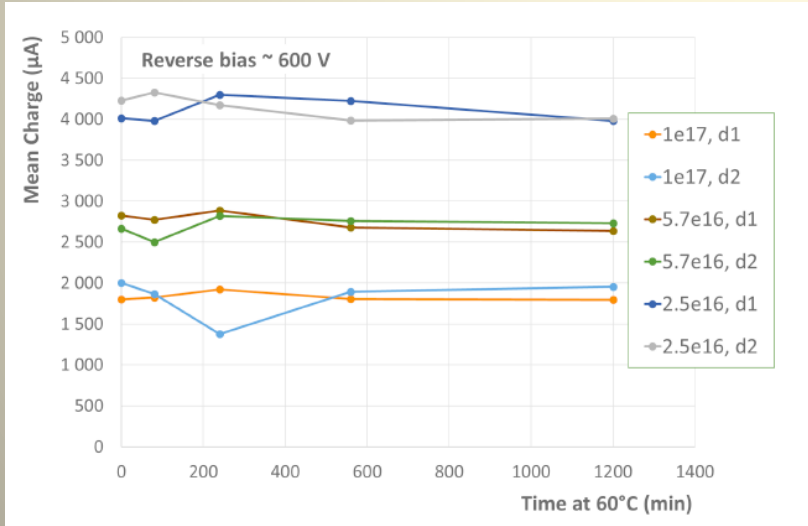
$$Q_{mean} = k \cdot \phi^b \cdot V$$

$$k_{75} = 44 \text{ e}_0/\text{V}$$

$$b_{75} = -0.56$$

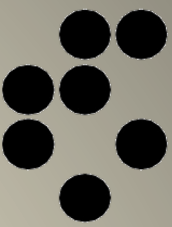
Thinner is better!

From:
**I.Mandić et al.,
 JINST 15 P11018 (2020).**



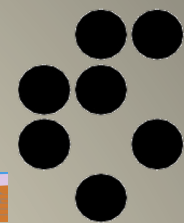


Linear $CCE(V)$??



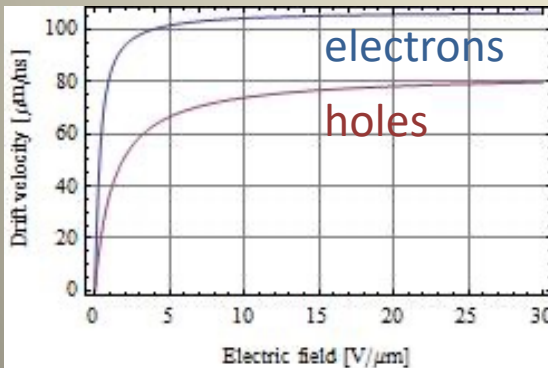
- What could be linear
 - SCR governed $CCE(V)$ after irradiation (\sqrt{V}), highly resistive ENB (\sqrt{V}), without trapping
 - Trapping dominated with non-saturated drift velocity
- What is *not* linear
 - velocity saturation
 - charge multiplication
 - double junction
 - field in ENB
 - ...
- Just a nice coincidence or some physics behind ?
 - look *into* silicon to search for an answer
- Using edge-TCT to probe silicon

Electric Field Measurement

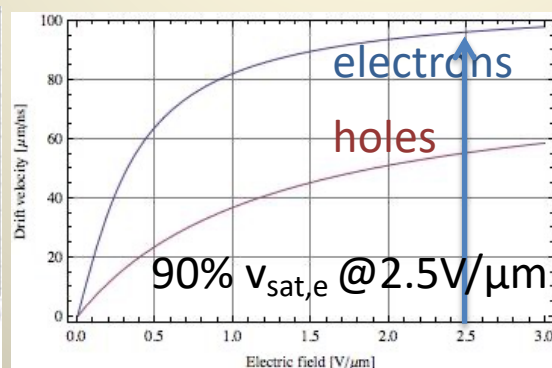


- Initial signal proportional to velocity sum at given detector depth
- Caveats for field extraction
 - Transfer function of electronics smears out signal, snapshot taken at ~ 600 ps
 - Problematic with heavy trapping
 - Electrons with v_{sat} hit electrode in 500 ps
 - Mobility depends on E
 - v saturates for $E \gg 1\text{V}/\mu\text{m}$

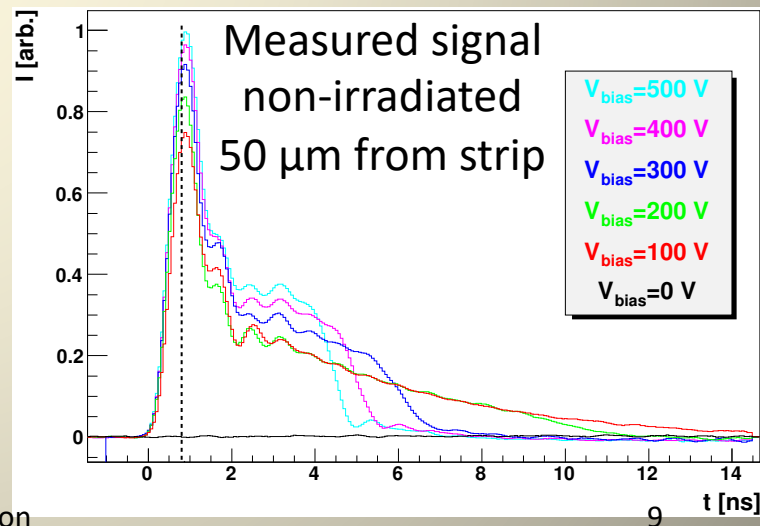
$$\begin{aligned}
 I(t=0) &= q \cdot \vec{v} \cdot \vec{E}_w = \\
 &= N_{e-h} e_0 \cdot (v_e + v_h) / d = \\
 &= N_{e-h} e_0 \cdot (\mu_e + \mu_h) \cdot E(x) / d
 \end{aligned}$$



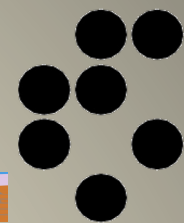
13th PSD, 6/9/2023



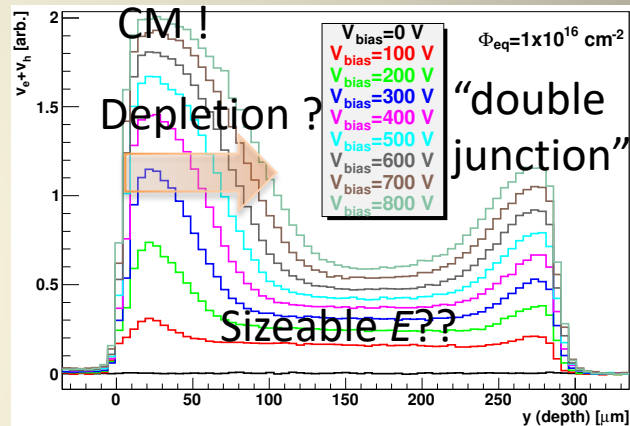
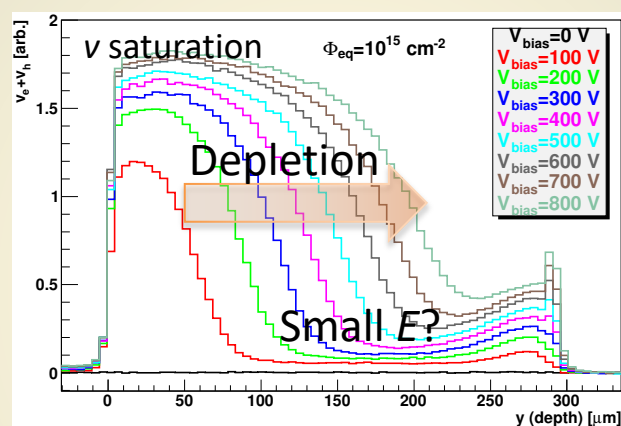
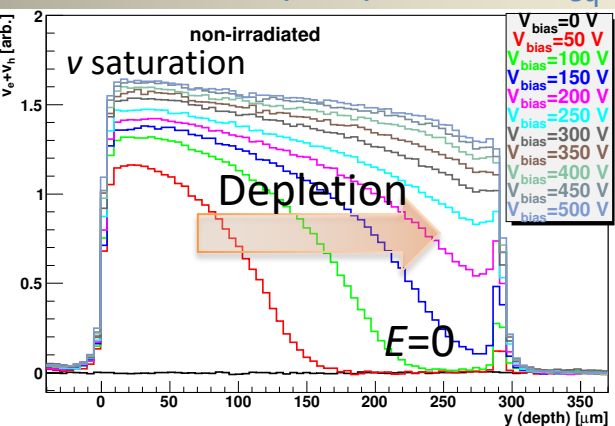
Marko Mikuž: Extreme Radiation



Selected Results from Neutrons



- Hamamatsu ATL07 n⁺ mini-strip, FZ p-type, neutron-irradiated at JSI TRIGA reactor
 - In steps up to 10¹⁶ n_{eq}/cm²

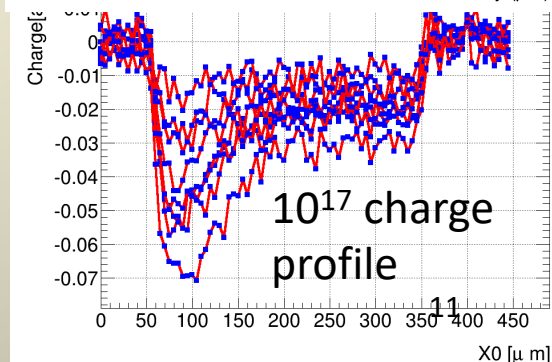
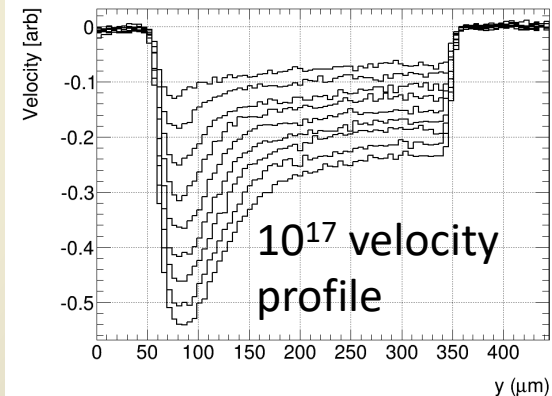
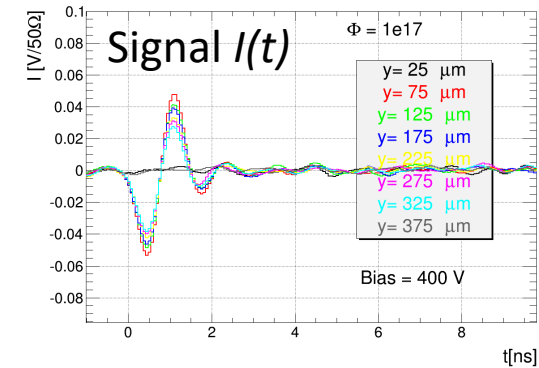


- Very instructive regarding qualitative electric field shape
 - Non-irradiated “by the book” for abrupt junction n⁺p diode
 - SCR and ENB nicely separated, small double junction near backplane
 - Medium fluence ($\Phi=10^{15}$ neutrons): some surprise
 - Smaller space charge than expected in SCR, some field in “ENB”
 - Large fluence ($\Phi=10^{16}$): full of surprises
 - Still lower space charge, sizeable field in “ENB”
 - Charge multiplication (CM) additional trouble for interpretation at large V
- Nice, but let’s try to get *quantitative* !

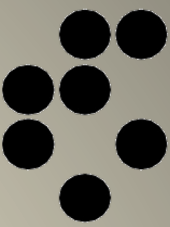
Published in :
**G. Kramberger et al.,
 JINST 9 P10016(2014).**

Extending the Reach

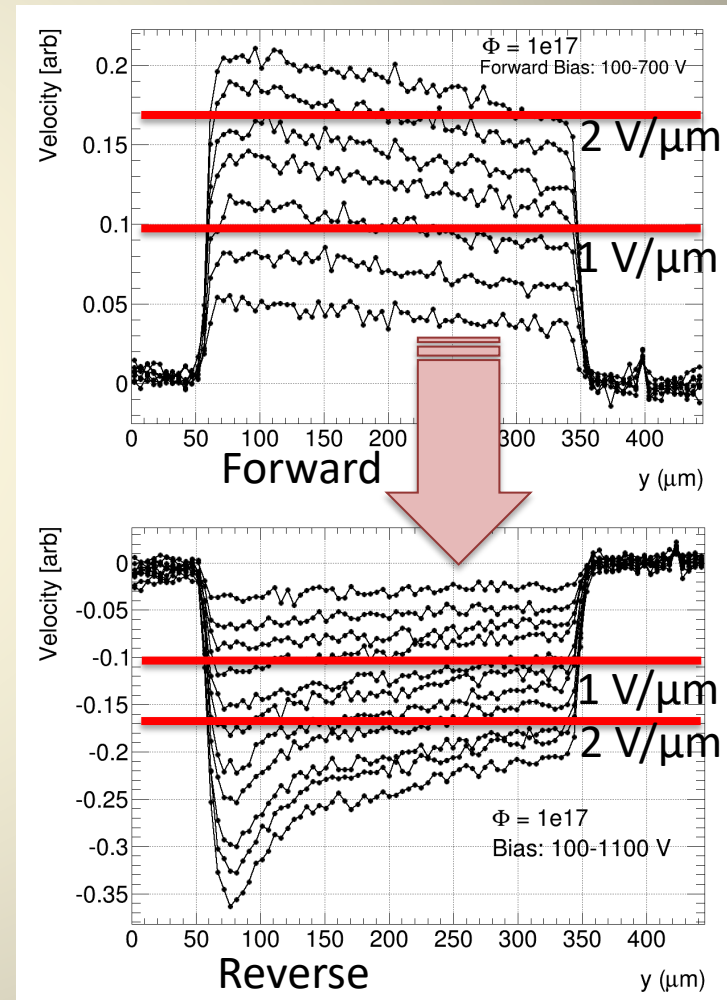
- In 2014 added 5×10^{16} and 10^{17} n_{eq}/cm^2 measurements of the same detector
 - 10^{16} of this fluence fully annealed, the rest 80 min @ 60°C
- Intrinsic feature – signal oscillations
 - period $\sim 5/4$ ns
 - LRC ($C \sim 2\text{pf} \Rightarrow L \sim 20\text{ nH} \sim 1\text{cm}$ of wire)
- Velocity (slope) and charge (integral) yield consistent results
- should be, as $Q \approx Q_0 v_{sum} \tau_{eff} / d$



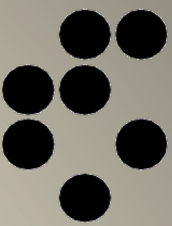
Absolute Field Measurement



- Solution: *concurrent* forward bias v_{sum} measurements
 - Ohmic behaviour with some linear (field) dependence
 - constant (positive) space charge
 - can use $\int E(y) dy = \bar{E}d = V$ to pin down field scale
 - corrections from $v(E)$ non-linearity small
- Use same scale for reverse bias!
- FW measurements up to 700 V
 - know E scale up to 2.33 V/ μm
 - can reveal $v(E)$ dependence



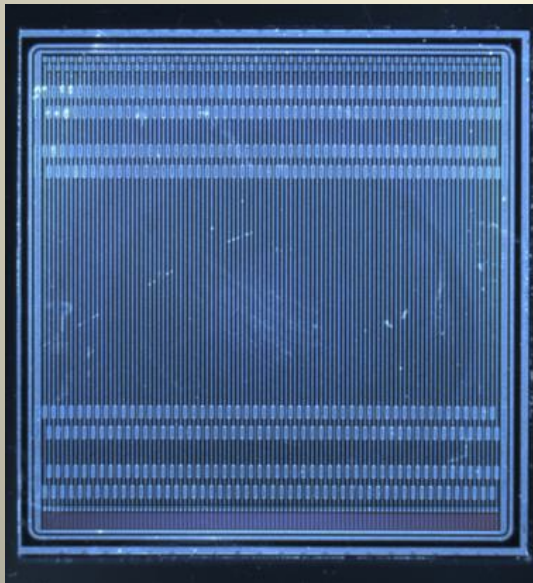
Proton Irradiations



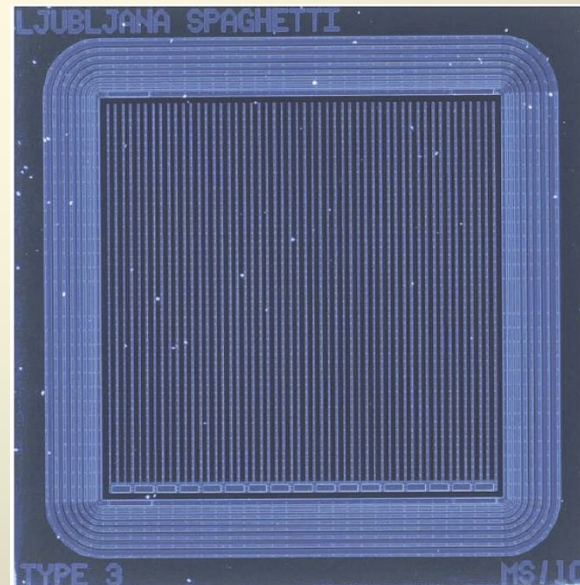
- 5 sample pairs of ATL12 mini-strips irradiated at CERN PS during summer 2015
 - got 0.5, 1.0, 2.9, 11, 28e15 protons/cm², no scanning
 - NIEL hardness factor 0.62
 - thanks to CERN IRRAD team
 - took 41 PS days to reach the highest fluence
- Covers HL-LHC tracker range well
 - does really not look practical for 10¹⁷++
- 2 samples per fluence investigated by E-TCT for all fluences
 - concurrent forward and reverse bias measurements

Additional Irradiations

- $3e17$ n_{eq}/cm^2 , JSI reactor neutrons
 - A12 mini, 7×8 mm², 75 μ m pitch, 300 μ m thick
 - Also to $3e16$, $1e17$
 - Spaghetti: 4×4 mm², n-on-p, strip pitch 80 μ m, 300 μ m thick, strips connected together at side
 - $1.6e17$ received previously, $4.6e17$ total

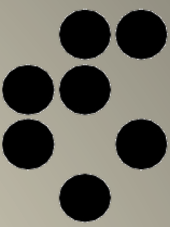


A12

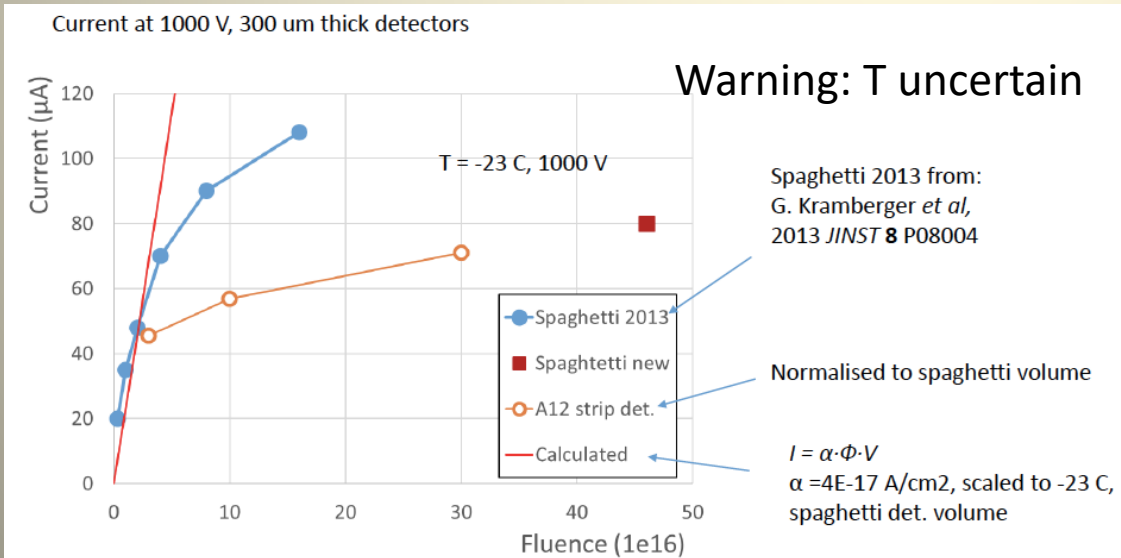
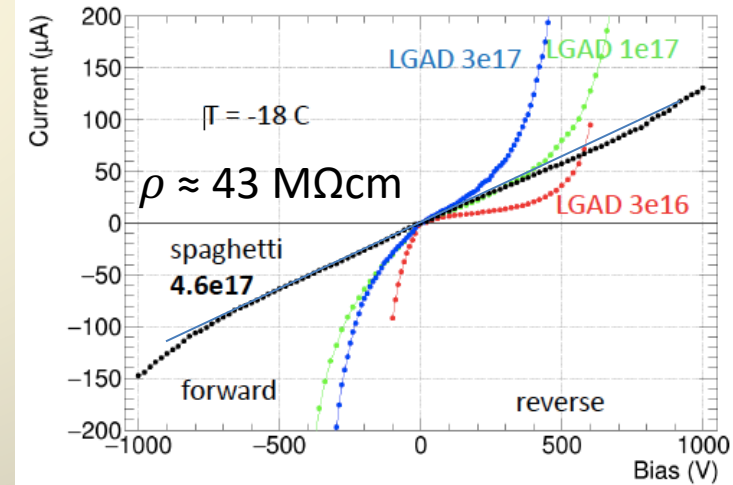
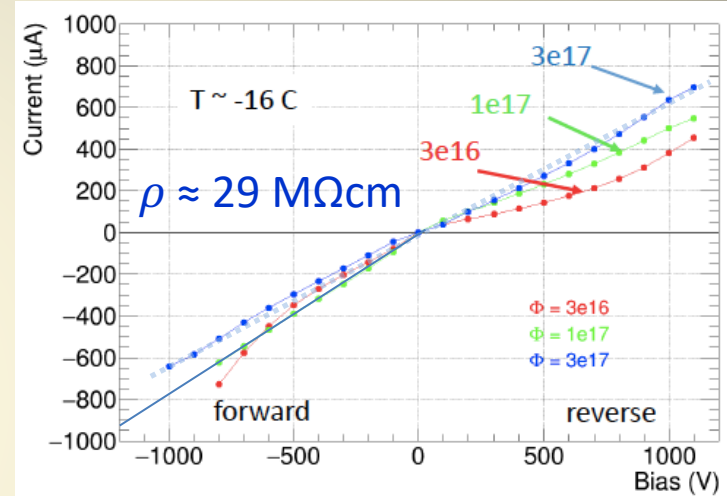


Spaghetti

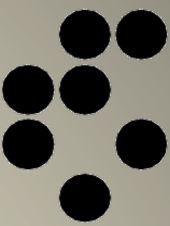
Basic Measurements



- I - V for 3&4.6e17 looks very linear with little difference between reverse/FW bias
 - No breakdown, as observed in LGAD's
- I @1000 V does not scale linearly with fluence !
 - Not governed by generation current ?
- Tried to measure 4.6e17 spaghetti CCE with ^{90}Sr
 - No signal above background observed up to 320 V
 - Magic formula predicts 120e for 4.6e17 @320 V



Mobility Considerations FW bias

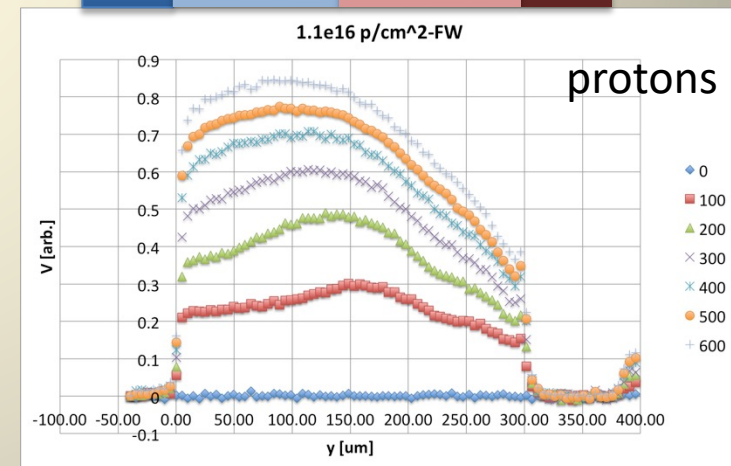
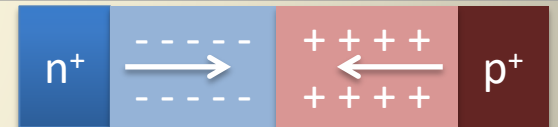
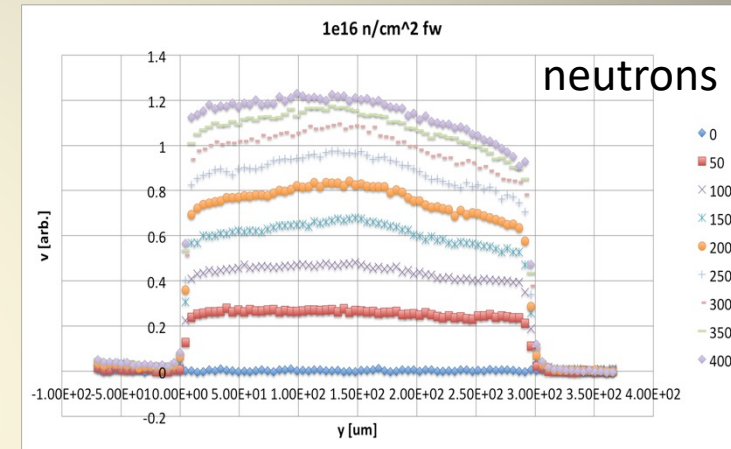


- For forward bias can extract $v(E)$ up to a scale factor
- Observe less saturation than predicted
- Model with

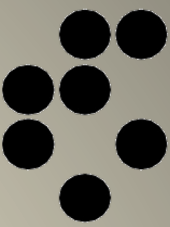
$$v_{sum}(E) = \frac{\mu_{0,e}E}{1 + \frac{\mu_{0,e}E}{v_{e,sat}}} + \frac{\mu_{0,h}E}{1 + \frac{\mu_{0,h}E}{v_{h,sat}}}$$

- keep saturation velocities at nominal values @-20°C ($v_{e,sat} = 107 \mu\text{m/ns}$; $v_{h,sat} = 83 \mu\text{m/ns}$)
- float (common) zero field mobility degradation
- fit $v(E)$ for $\phi_n \geq 5 \times 10^{15}$ and $\phi_p \geq 3 \times 10^{15}$

n.b. FW profiles less uniform for lower fluences & protons; departures from average field still small, corrections $O(\%)$

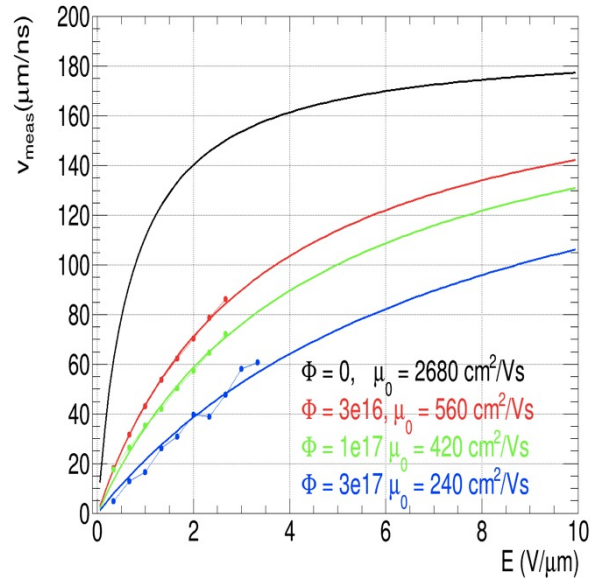
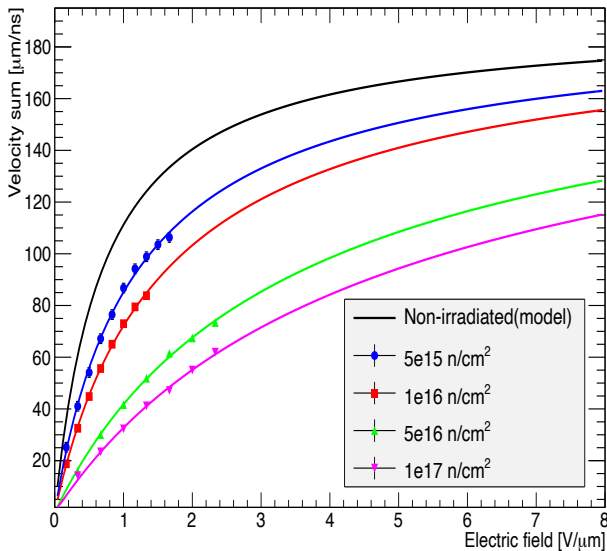


Mobility Fits

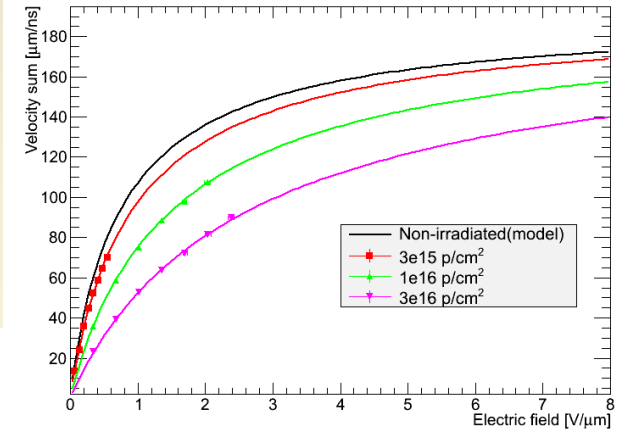


- Model fits data almost perfectly
 - μ_0 degradation the only free parameter, scale fixed by $v_{sum,sat}$
 - At $3e17$ E range too limited ($v(E)$ linear), regard result as upper limit

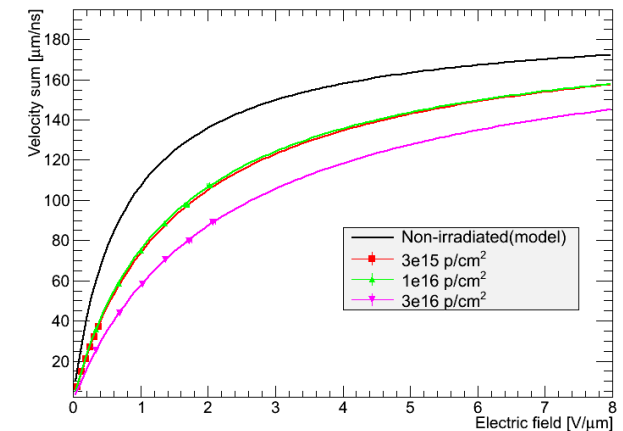
Mobility neutrons



Mobility protons A

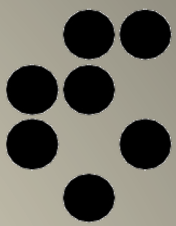


Mobility protons B





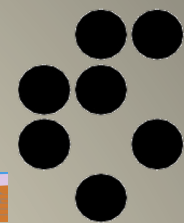
Mobility Results



- Fit to $v_e + v_h$ with common mobility degradation factor
 - factor of **2** at $10^{16} n_{eq}/cm^2$, **6** at $10^{17} n_{eq}/cm^2$, **>10** at $3 \times 10^{17} n_{eq}/cm^2$
 - need **2x/6x/>10** higher E to saturate v !
 - ☠️ correspondingly higher E for charge multiplication !

| Φn | $\mu_{0,sum}$ | Φp | $\mu_{0,sum}$ |
|---------------------------|----------------|---------------------------|----------------|
| [$10^{15} n_{eq}/cm^2$] | [cm^2/Vs] | [$10^{15} n_{eq}/cm^2$] | [cm^2/Vs] |
| non-irr (model) | | 2680 | |
| 5 | 1661 ± 134 | 1.6 | 2063 ± 188 |
| 10 | 1238 ± 131 | 6.1 | 1337 ± 47 |
| 30 | 560 | 15.4 | 817 ± 42 |
| 50 | 555 ± 32 | | |
| 100 | 407 ± 40 | | |
| 100 | 420 | T=-20°C | |
| 300 | <240 | | |

Mobility Analysis



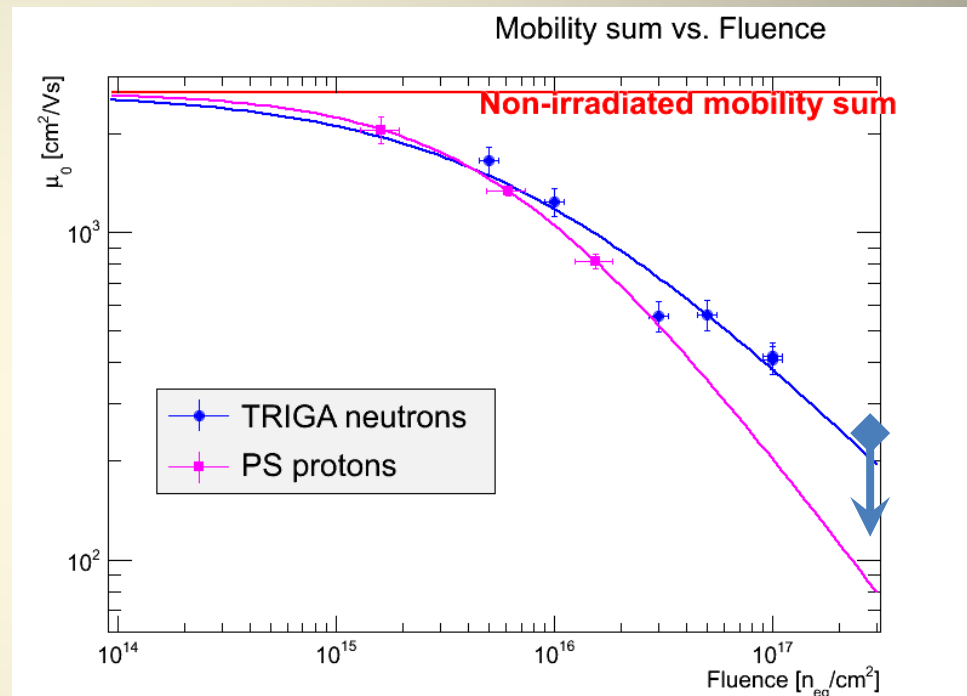
- Mobility governed by hard scattering on acoustic phonons and traps

$$\frac{1}{\tau} = \frac{1}{\tau_{ph}} + \frac{1}{\tau_{trap}}$$

- Fit mobility dependence on fluence with a power law

$$\mu_{0,sum}(\Phi) = \frac{\mu_{0,sum,phonon}}{1 + \left(\frac{\Phi}{\Phi_{1/2}}\right)^a}$$

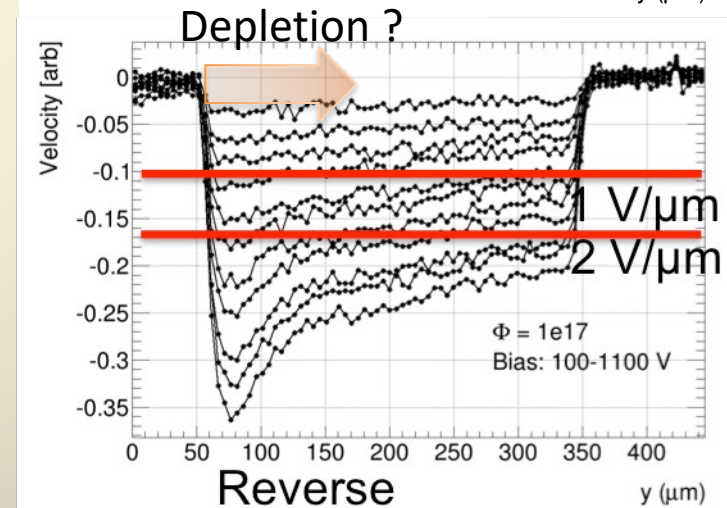
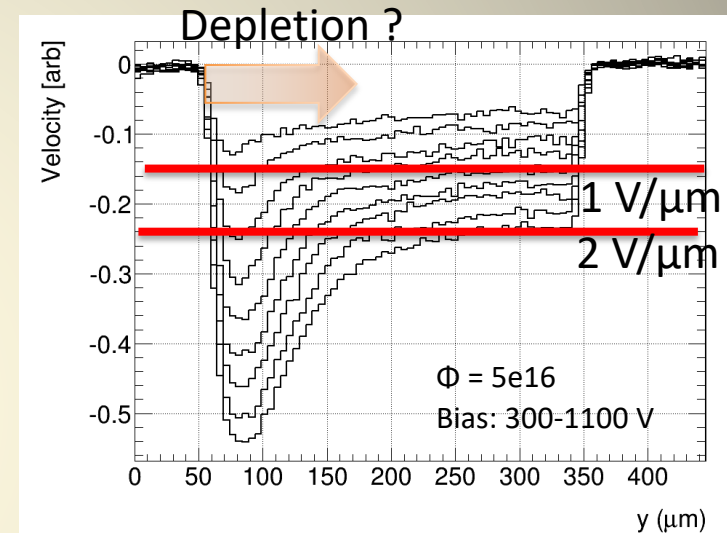
- Fits perfectly, value of a close to linear
 - 10% error assumed for all neutron data
- At same NIEL, mobility decrease worse for protons
 - NIEL violation? Large errors?



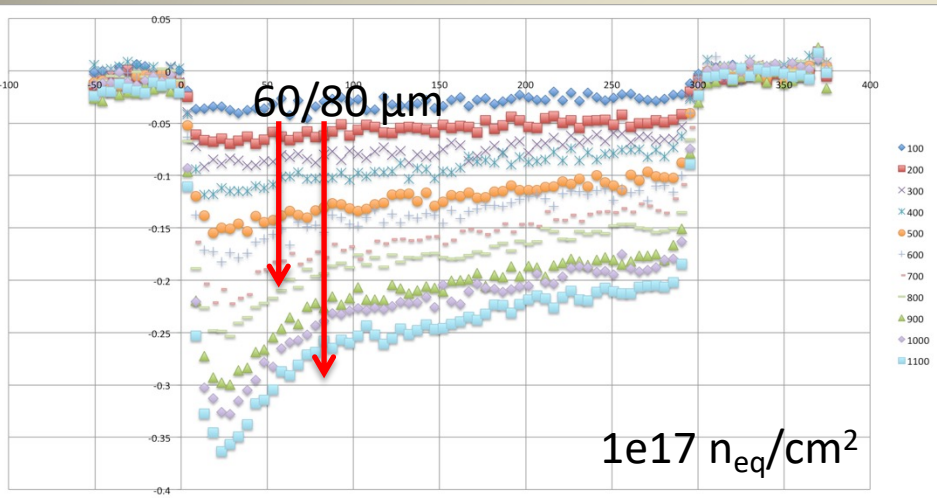
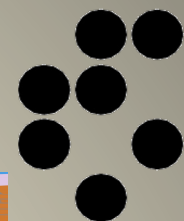
| Irradiation particle | a | σ_a | $\Phi_{1/2} / 10^{15}$ | $\sigma_{\Phi_{1/2}} / 10^{15}$ |
|----------------------|-------|------------|------------------------|---------------------------------|
| Reactor neutrons | -0.68 | 0.08 | 6.9 | 1.7 |
| PS protons | -0.90 | 0.19 | 6.1 | 1.0 |

Reverse Bias Field Profile

- Two distinct regions at high biases
 - Large region from backplane with (small) slope in the field
 - constant (small, negative) space-charge
 - $E = j \cdot \rho$ at junction ? like “ENB” ?
 - indication of thermal (quasi)equilibrium: $np = n_i^2$?
 - thus no current generation ?
 - Small region at junction building up with bias
 - depleted space-charge region ?
 - source of generation current ?



SCR Consistency

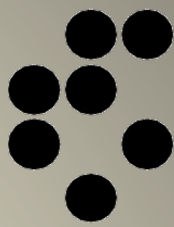


- Hard to estimate SCR extent, especially at lower bias and highest fluence
- A crude estimate
 - $5 \times 10^{16} n_{eq}/cm^2$:
~80 μm @ 600 V; ~120 μm @ 1000 V
 - $10^{17} n_{eq}/cm^2$:
~60 μm @ 600 V; ~80 μm @ 1000 V

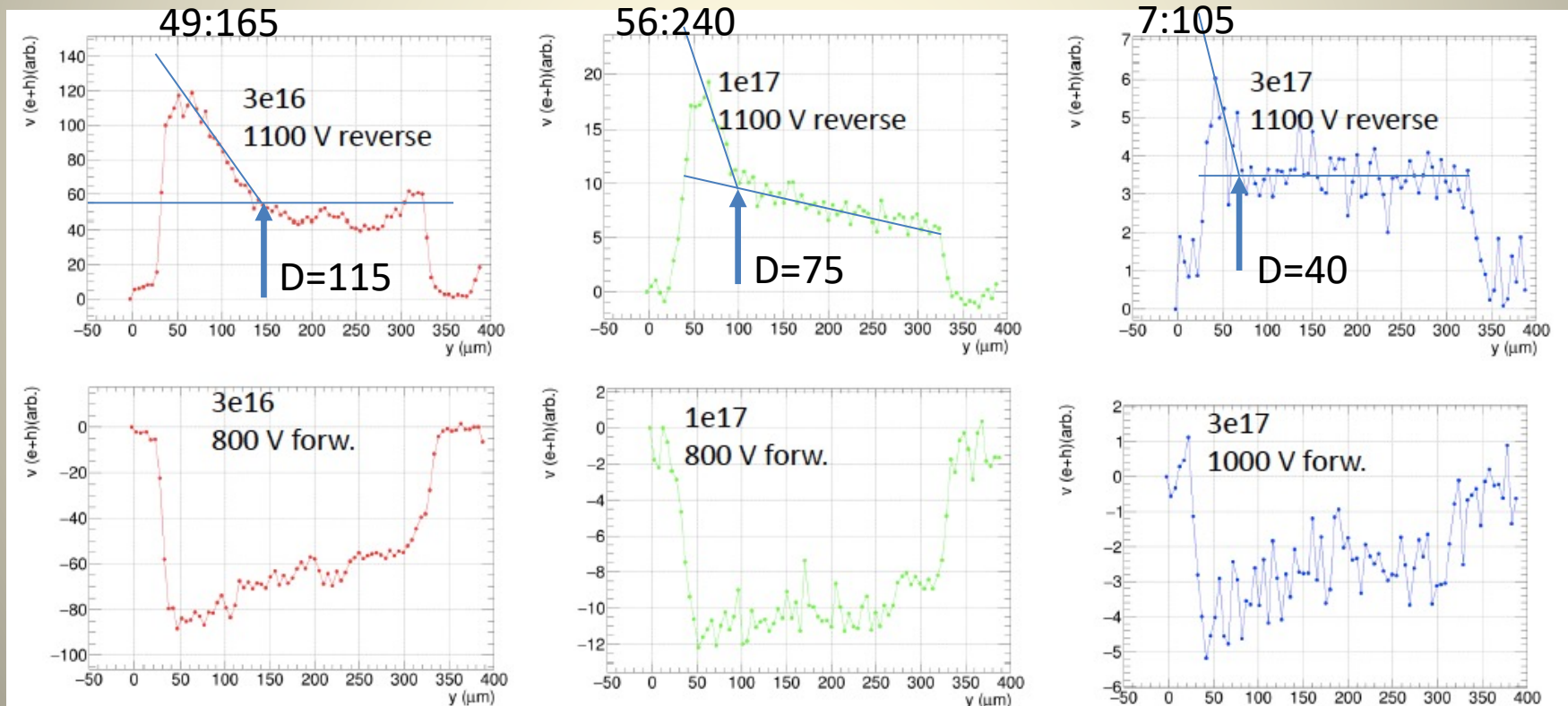
- Predicted/measured currents
 - $5 \times 10^{16} n_{eq}/cm^2$: 300/300 μA @ 600 V; 400/500 μA @ 1000 V
 - $10^{17} n_{eq}/cm^2$: 400/300 μA @ 600 V; 500/600 μA @ 1000 V
 - Not compatible with linear I - V at 3 & 4.6e17 – pure resistor ?
- Reasonable agreement with current generated exclusively in SCR
 - n.b. - current “saturation” observed @1000V in *JINST 8 P08004 (2013)*
- Acceptor introduction rates: $g_c \approx 6/4 \times 10^{-4} cm^{-1}$
 - substantial part (up to 80 %) of voltage drop “spent” in “ENB”
 - matches well data in *JINST 9 P10016(2014)* (up to 10^{16})



ATL12 up to $3e17$

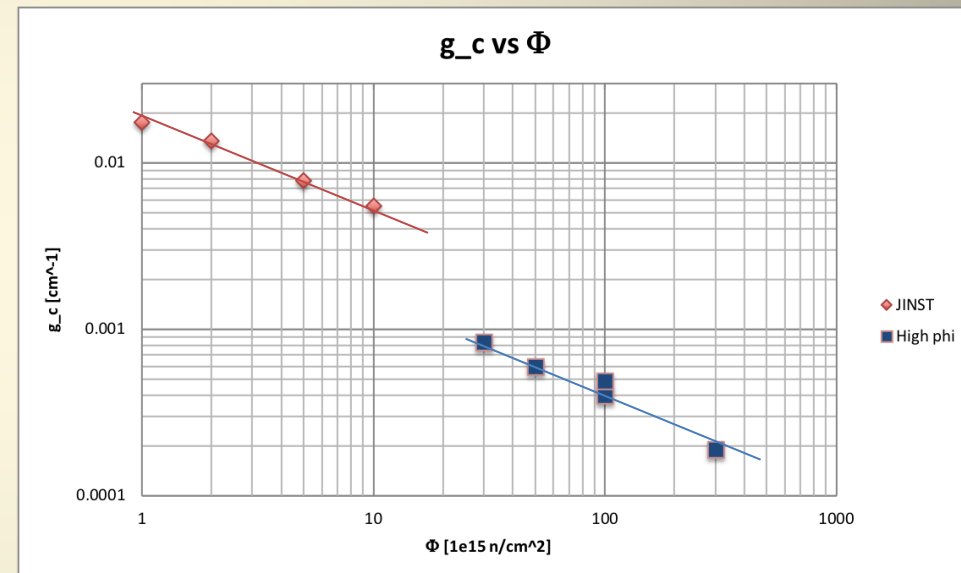


- Estimate of SCR width 115 \rightarrow 75 \rightarrow 40 μm
- V_{drop} in SCR only 23 \rightarrow 19 \rightarrow 6 % of 1100 V

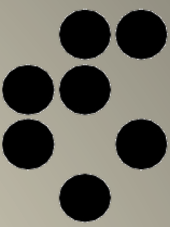


Acceptor introduction in SCR

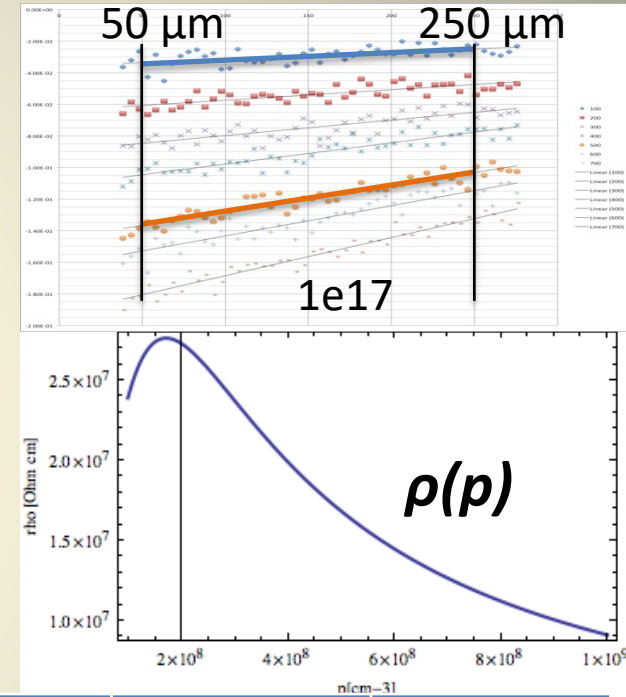
- Stable acceptor introduction rate g_c drops by nearly two orders of magnitude from low fluences to 3×10^{17}
 - Observed up to 10^{16} in *JINST 9 P10016(2014)*
 - Looks like a power law
 - g_c in JINST not taking into account voltage drop out of SCR – higher values of g_c



“ENB” Consistency



- Space charge in “ENB” rising with bias, e.g. for $10^{17} n_{eq}/cm^2$
 - 1.6×10^{11} @ 100 V, $9.2 \times 10^{11} cm^{-3}$ @ 500V
 - c.f. $\sim 4 \times 10^{13} cm^{-3}$ in SCR
 - negative space charge, like in SCR
- Resistivity from $\rho = j/E$ @ 100 V
 - maximum $\rho(p) \approx 2.8 \times 10^7 \Omega cm$ using nominal mobilities @ $p \sim 2 \times 10^8 cm^{-3}$
 - all measured values exceed this limit
 - compatible with measured mobility sum and $p \sim O(10^9) cm^{-3}$
 - Compatible also with ρ from $I-V$ for 3 & 4.6e17

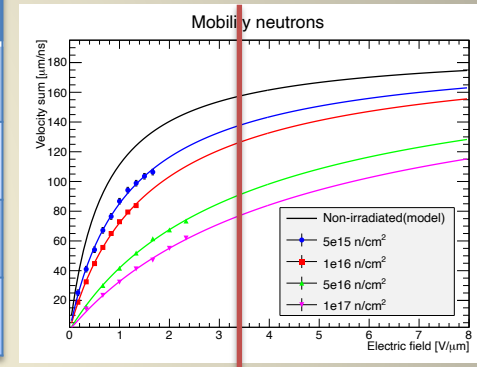


| Φ | ρ | p |
|-----------------|--------------------|------------------|
| $[n_{eq}/cm^2]$ | $[10^7 \Omega cm]$ | $[10^9 cm^{-3}]$ |
| 1e16 | 3.3 | 0.5 |
| 5e16 | 3.0 | 1.5 |
| 1e17 | 2.8 | 2.1 |

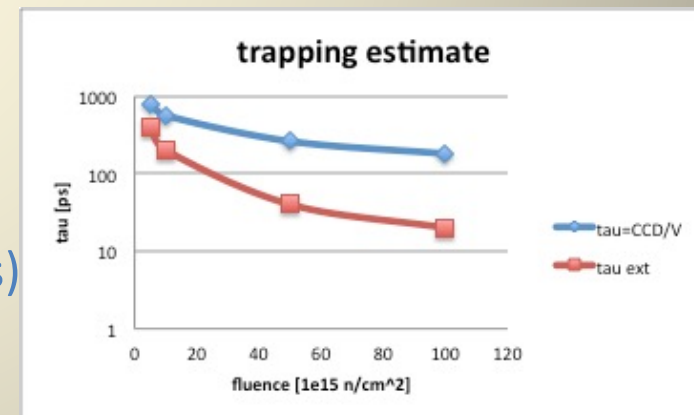
Trapping analysis

- Take v_{sum} at average $E = 3.3 \text{ V}/\mu\text{m}$
- Calculate CCD from “magic formula”

| Φ [1e15] | 5 | 10 | 50 | 100 |
|--------------------------------------|-----|-----|-----|-----|
| $v_{sum}(3.3 \text{ V}/\mu\text{m})$ | 137 | 126 | 90 | 77 |
| $CCD_{1000 \text{ V}} [\mu\text{m}]$ | 110 | 70 | 23 | 14 |
| $\tau \approx CCD/v$ [ps] | 800 | 560 | 260 | 180 |
| τ_{ext} [ps] | 400 | 200 | 40 | 20 |

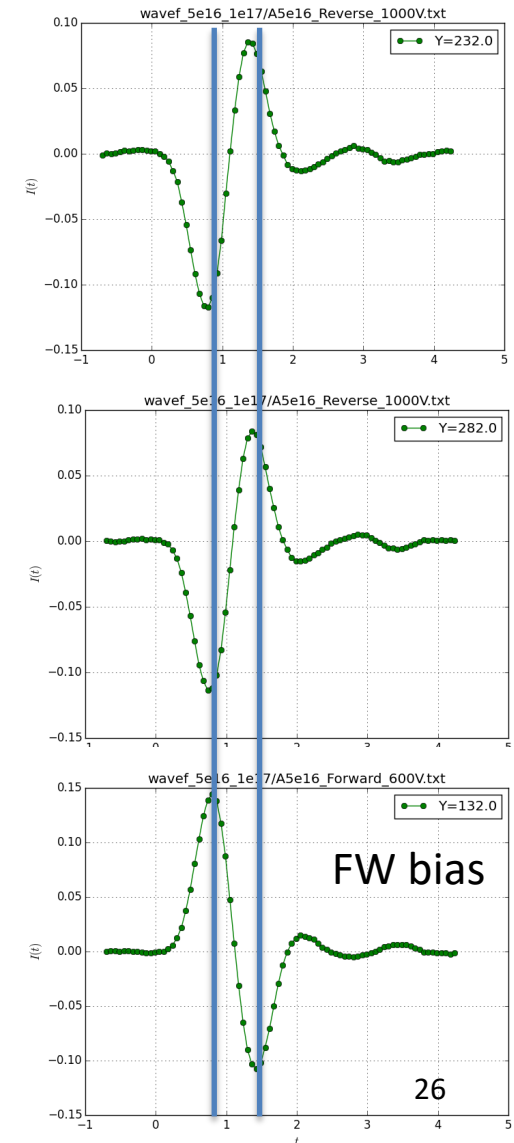
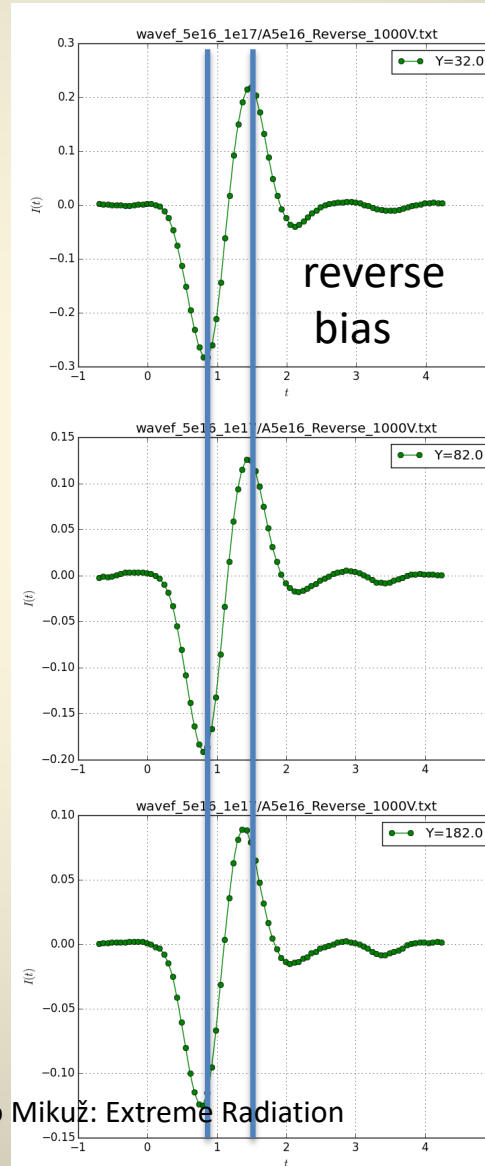


- Implies factor of 6-9 less trapping at highest fluences
 - lowest fluence still x2 from extrapolation
 - weak dependence on fluence as anticipated
 - CM would effectively shorten trapping times
 - not good when large E variations ($v(E)$ saturates)
 - not good when $CCD \approx$ thickness (less signal at same τ)



Trapping – position dependence ?

- Waveforms (WF) plotted every 50 μm in detector depth for reverse bias at 1000 V
- Forward bias in middle of detector added at 600 V
- Very little, if any, WF dependence on position observed
- Trapping not position dependent !?



Trapping revisited

- From *I. Mandić et al., JINST 15 P11018 (2020)*

– FW bias CCE estimated by

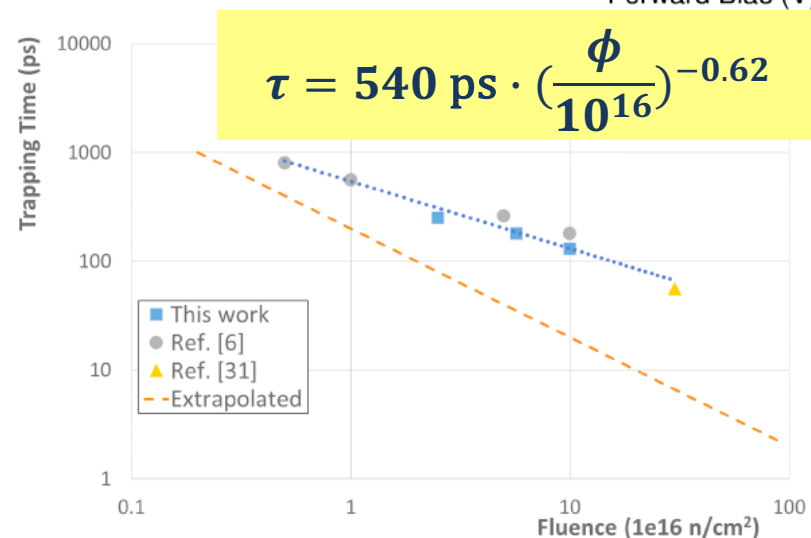
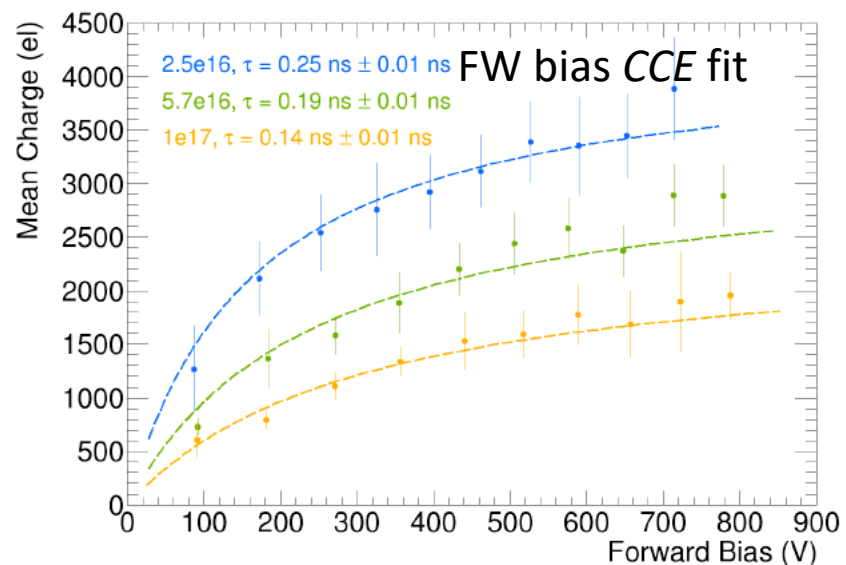
$$Q = \frac{\Delta Q}{\Delta x} \cdot v \cdot \tau$$

- $v(E)$ with fluence dependent μ
- constant $E=V/D$ (FW)

😊 Order of magnitude smaller than extrapolated !

😊 Agrees with estimates from reverse bias CCE

- Trapping independent of bias, seen in wave-forms



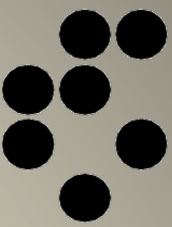
Summary

- Measurements performed on Si detectors irradiated to extreme fluences
 - Neutrons from 10^{15} to 4.6×10^{17} $n_{\text{eq}}/\text{cm}^2$, PS protons from 5×10^{14} to 3×10^{16} p/cm^2
 - Velocity vs. electric field impact observed and interpreted as reduction of zero field mobility
 - Zero field mobility follows power law with $|a| \leq 1$, $\Phi_{1/2} \approx 10^{16}$ n/cm^2
 - Protons degrade mobility more than neutrons
 - Induces resistivity increase in-line with measured I - V
 - Exhibits adverse effect on charge multiplication !
 - Simple field profile for very high neutron fluences
 - Diminishing SCR and highly resistive ENB
 - Effective acceptor introduction rates reduced by factor ~ 100 wrt low fluences
 - Current much lower than anticipated. Generated in SCR only ? Ohmic at highest fluences...
 - Trapping estimates for very high neutron fluences
 - from charge collection in FW and reverse bias
 - from waveforms
 - All estimates point to severe non-linearity of trapping with fluence, 10x lower at 10^{17}
 - Trapping appears independent of electric field
- Conclusion: ***Low fluence extrapolations do not work at all !***
... go out and ***measure*** to get anything working at ***extreme*** fluences !!!

Implications for DRDT 3.3

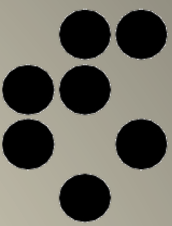
- Basic bulk silicon properties in the fluence range to master are the prerequisite to any inner tracking detector design for *FCC-hh*
- They need to be *measured*
 - Only pioneering consistency checks done so far
- Need resources far beyond current ones
 - Facilities
 - Measurement techniques
 - Peopleat least for the first ~5 of the 20 years
- New DRD3 Collaboration based on the RD50 research line essential for achieving the goal
 - Close to 70 institutes signed up for “WG3 Radiation damage and extreme fluences” !
- EURO-LABS project has 4 neutron irradiations budgeted to $10^{18} n_{eq}/cm^2$
 - Not so obvious how to get high energy protons beyond $10^{17} n_{eq}/cm^2$

Conclusion



NOTHING IS
IMPOSSIBLE,
THE WORD
ITSELF SAYS
“I’M POSSIBLE”!
- AUDREY HEPBURN

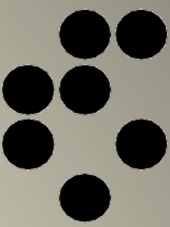
Backup Slides



Proton Irradiations

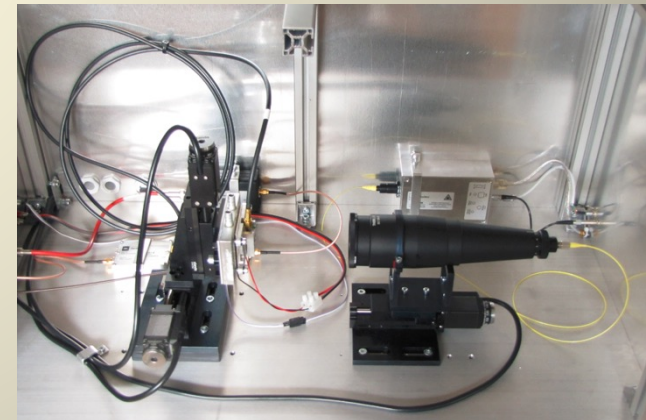
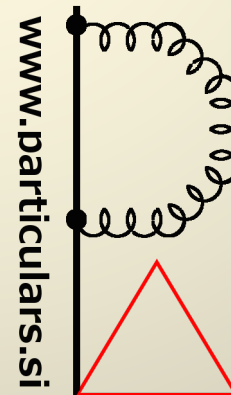
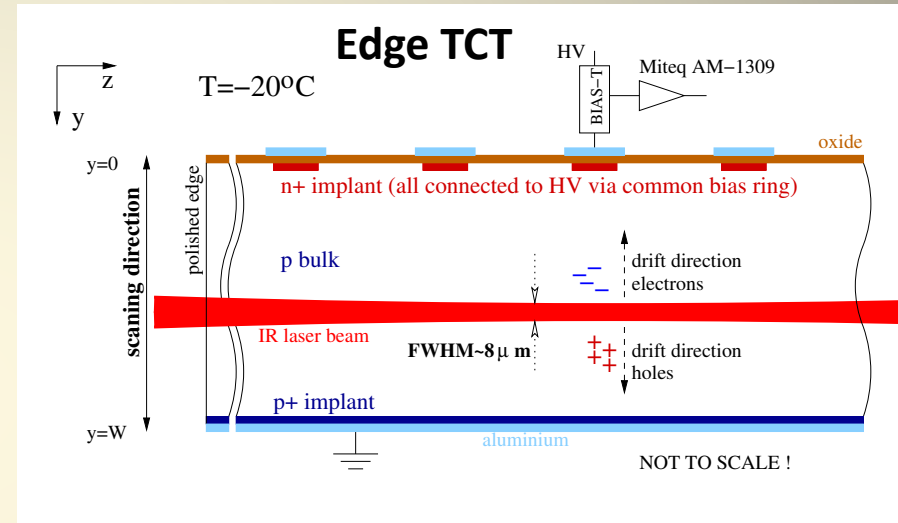
- Several new high energy proton accelerators in construction – spallation sources
 - Energy in GeV range, high currents (mA)
 - 1 μA provides $\sim 10^{18}$ on 1 cm^2 in one day !
- Problem – cooling & radiation safety
 - In 300 μm Si the MIP heating load is $\sim 0.1 \text{ W/cm}^2$
 - Or $\sim 1 \text{ W/g}$, heating rate $\sim 1 \text{ K/s}$
 - Each irradiation site is certified up to a maximum beam current
 - 1 μA needs to be planned, preferably during construction
- Engineering issues that need to be worked on

Edge TCT



Edge-TCT

- Generate charges by edge-on IR laser perpendicular to strips, detector edge polished
- Focus laser under the strip to be measured, move detector to scan
- Measure induced signal with fast amplifier with sub-ns rise-time (Transient Current Technique)
- Laser beam width $8\ \mu\text{m}$ FWHM under the chosen strip, fast (40 ps) and powerful laser
 - Caveat – injecting charge under all strips effectively results in constant weighting (albeit not electric !) field

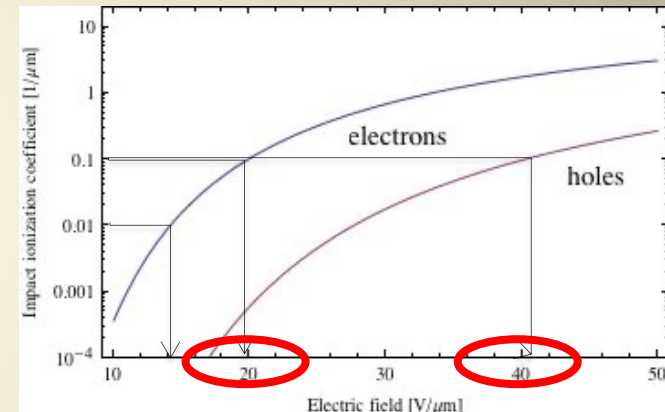


Charge Multiplication

- Multiplication is textbook physics
 - e.g. S.M. Sze, Physics of Semiconductor Devices, Wiley, New York, 1981
 - Ch 1.6.4 High-Field Property
 - Velocity saturation, impact ionization
 - Ch 2.5.3 Avalanche Multiplication
 - Junction break-down
- Measured impact ionization
 - Electrons create 1 pair in 10 μm at $E \sim 20 \text{ V}/\mu\text{m}$ (100 μm at 14 $\text{V}/\mu\text{m}$), holes need $E \sim 40 \text{ V}/\mu\text{m}$
 - Holes need $\sim 1 \text{ mm}$ for pair creation at $E \sim 20 \text{ V}/\mu\text{m}$
 - Neglect hole multiplication in signal creation altogether
 - Need to invoke hole multiplication for junction breakdown
- $\alpha_e \gg \alpha_h$ - Nature gentle to us (in silicon)
 - Large range in E where electrons multiply without inducing breakdown
 - But beware of (too) high electric fields !

$$\alpha_{e,h}(E) = \alpha_{e,h}^{\infty} e^{-b_{e,h}/E}$$

A. G. Chynoweth, Phys. Rev. 109, 1537(1958).



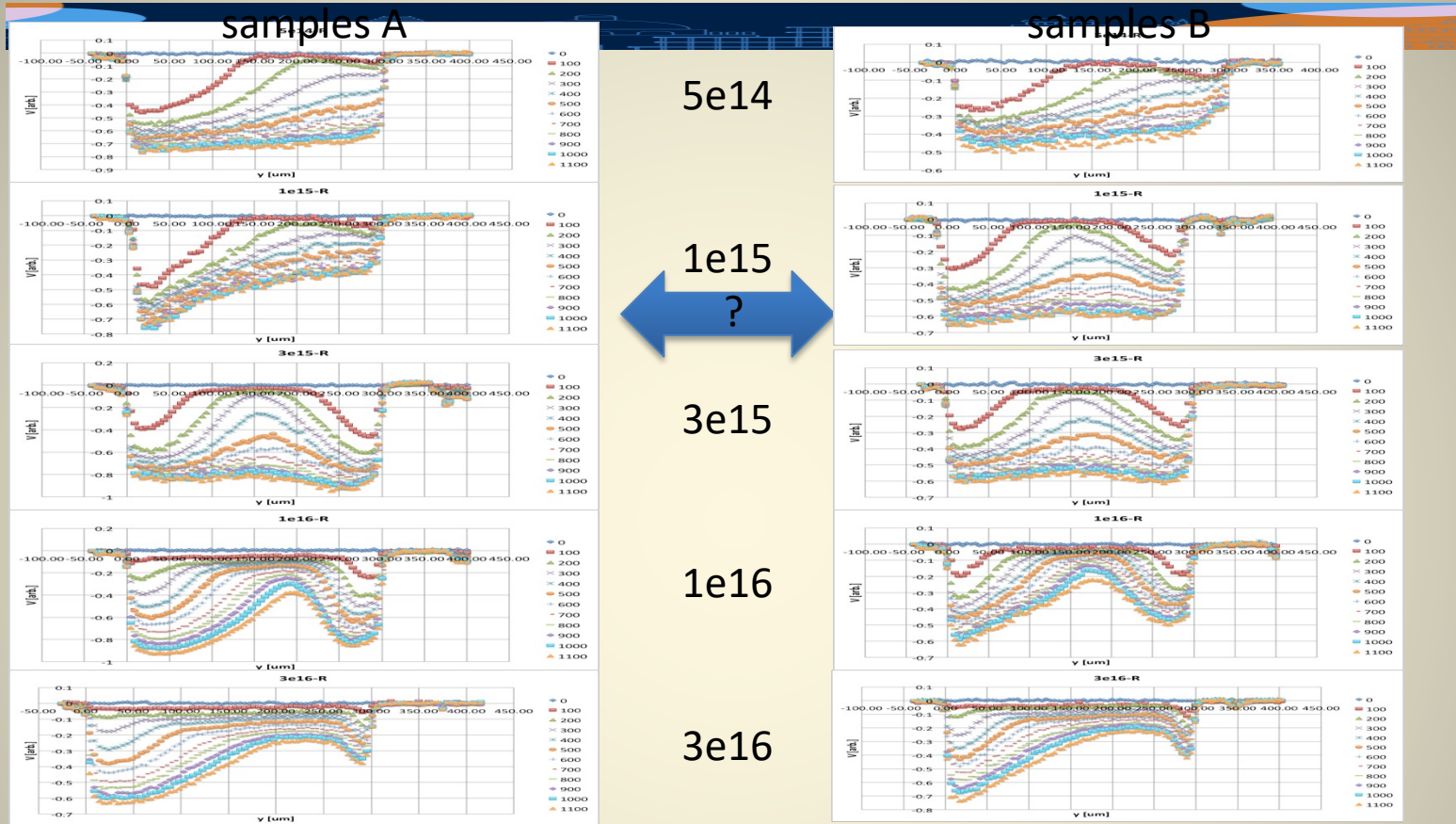
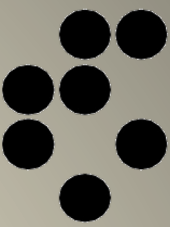
R.VAN OVERSTRAETEN and H.DE MAN, Solid-State Electronics 13(1970),583-608.
 W.MAES, K.DE MEYER, R.VAN OVERSTRAETEN, Solid-State Electronics 33(1990),705-718.

$$\int_0^w dx \alpha_e(x) e^{-\int_0^x (\alpha_e(x') - \alpha_h(x')) dx'} = 1$$

Breakdown condition, can swap α_e with α_h



Reverse velocity profiles



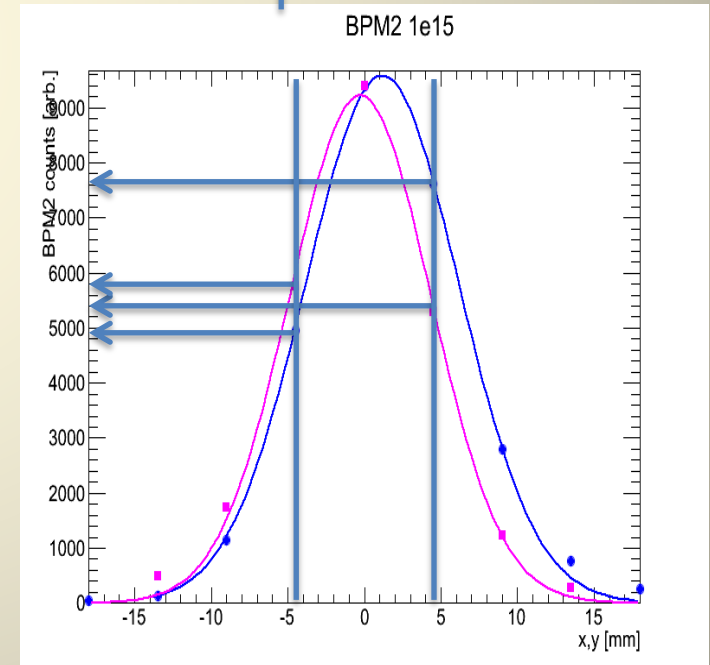
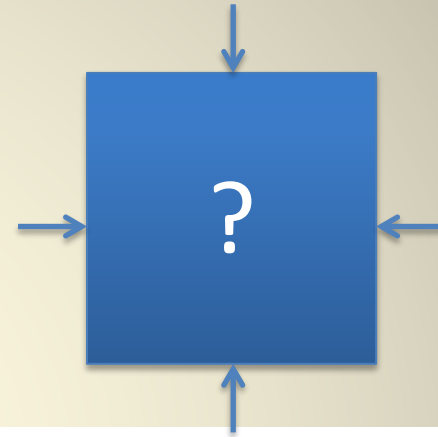
Something's fishy... never repeat experiments ?!



Explained by PS beam profile variation on sample edges

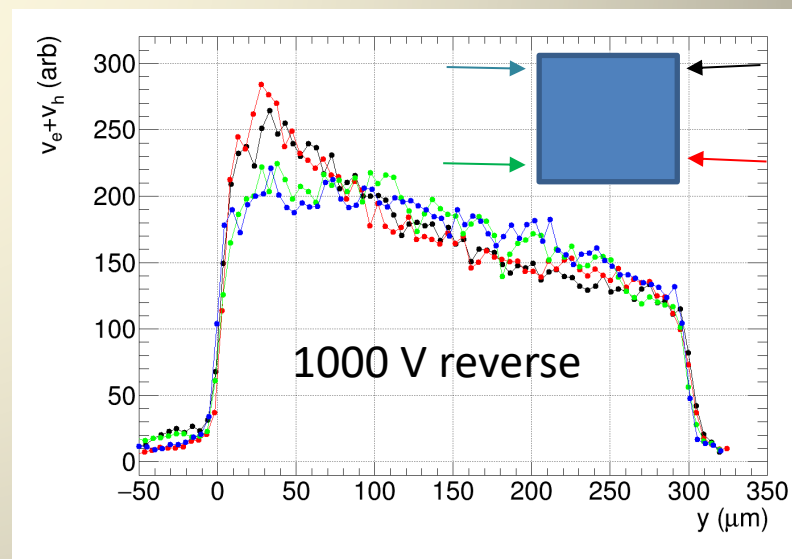
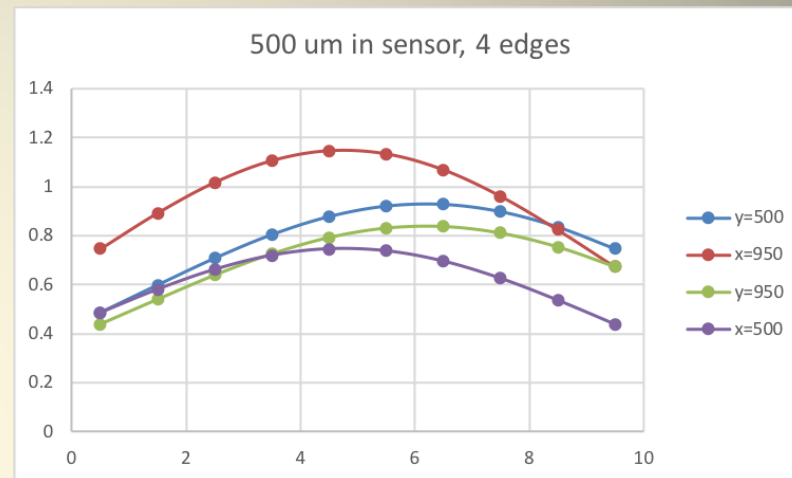
Proton irradiations - details

- Samples irradiated in PS in pairs
 - in series in same sample holder
- Same leakage current in both samples
 - ⇒ same *average* fluence received
- Beam profile asymmetric
 - monitored by BPM2
- Which side did we pick up ?



Protons revisited

- BPM2 results for the 1e15 sample, 0.5 mm in sensor
- 10x10 mm² average to peak: 0.7
 - Values rescaled
- Mid-side to average:
 - 1.17, 0.88, 0.82, 0.74
 - Must be the larger difference
- Correct fluences by -10 %
- Assign 20 % error
- Re-measured one sample from both sides, match with BPM2 data – still in progress
 - Looks like explaining the issue



Mobility Comparison

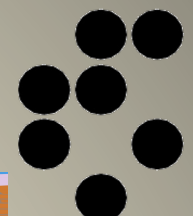


Table 3. Best-fitting parameters for the impurity dependence of electron and hole Ohmic mobilities at room temperature, as given in eqn (6)

| | Electrons | Holes | Units |
|-------------|----------------------|----------------------|---|
| μ_{min} | 92 | 47.7 | $\text{cm}^2 \text{V}^{-1} \text{sec}^{-1}$ |
| μ_{max} | 1360 | 495 | $\text{cm}^2 \text{V}^{-1} \text{sec}^{-1}$ |
| N_{ref} | 1.3×10^{17} | 6.3×10^{16} | cm^{-3} |
| α | 0.91 | 0.76 | — |

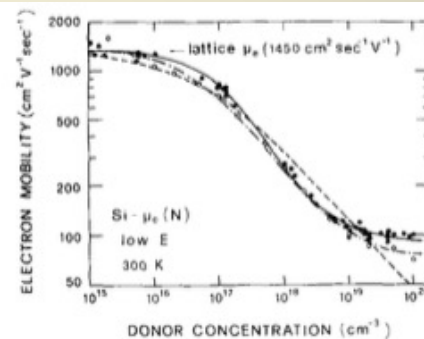


Fig. 5. Electron mobility, μ_e , in silicon at 300 K as a function of impurity concentration. Open and closed circles are the experimental results reported by Irvin[55] and of Mousty *et al.*[56], respectively. The continuous line is the phenomenological best fit (eqn (6)) of Baccarani and Ostoia[53] the broken line the best fit (eqn (7)) of Hilsun[54] the dot-dashed line (eqn (8)) of Scharfetter and Gummel[57] (see Tables 3 and 4).

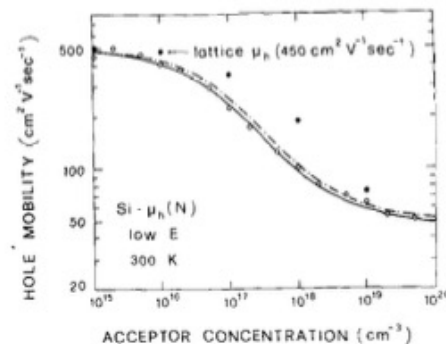


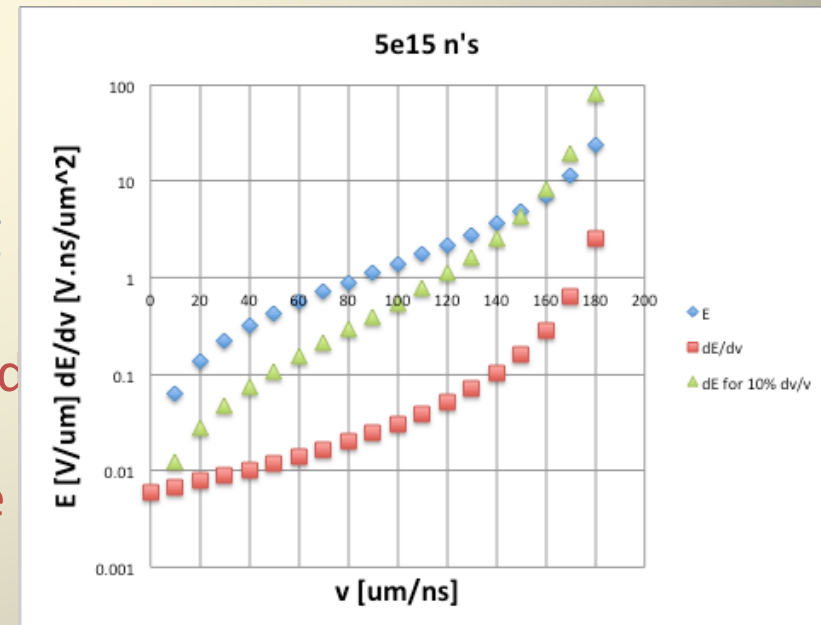
Fig. 6. Hole mobility, μ_h , in silicon at 300 K as a function of impurity concentration. Open circles are experimental results reported by Irvin[55]. Continuous and dot-dashed lines represent the best fitting curves of Caughey and Thomas[58] (eqn (6)) and of Scharfetter and Gummel[57] (eqn (8)), respectively (see Tables 3 and 4).

A REVIEW OF SOME CHARGE TRANSPORT PROPERTIES OF SILICON†
 C. JACOBONI, C. CANALI, G. OTTAVIANI AND A. ALBERICI QUARANTA
 Istituto di Fisica dell'Università di Modena, 41100 Modena, Italy
 (Received 18 March 1976; in revised form 12 July 1976)

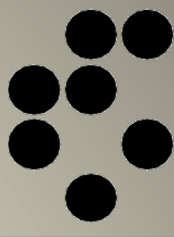
- Dependence on *shallow* dopant concentration
 - Measured in the roaring 60's
- Characteristic trap concentration $N \sim 10^{17} \text{ cm}^{-3}$
 - looks out of reach for typical $g=0(10^{-2})$
- But g refers to $N_{eff} = |N_a - N_d|$
- While N is more like $N_a + N_d$
 - *x-sections for deep and shallow?*
- Power law looks compatible: $\alpha \leq 1$

Velocity and Field Profiles

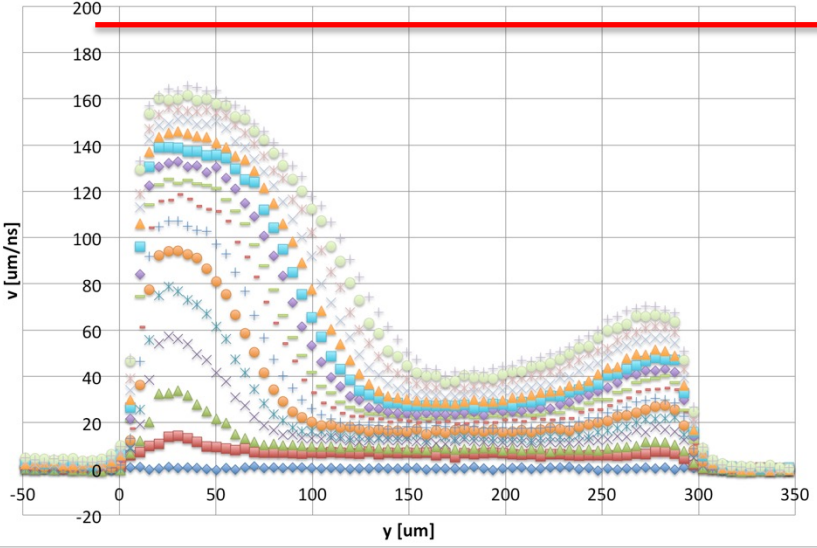
- Knowing $v(E)$ can set scale to velocity profiles
 - assumption: same scale on FW and reverse bias
 - protons: for 5×10^{14} and 10^{15} use same scale, fixed by average field for 5×10^{14} at 1100 V (no good FW data)
- Invert $E(v)$ to get electric field profiles
 - big errors when approaching v_{sat} i.e. at high E
 - exaggerated by CM in high field regions
 - $v > v_{sat}$ not physical, but can be faked by CM



Velocity Profiles Neutrons



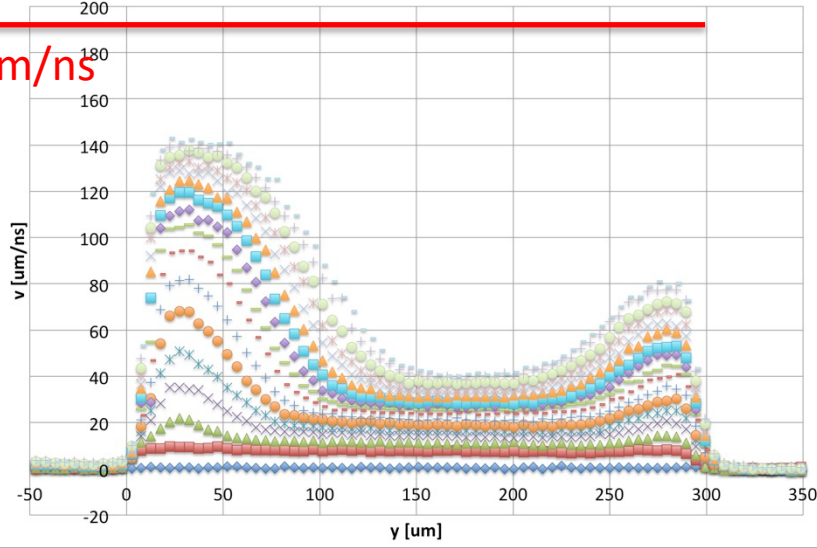
Velocity profile 5e15



$v = 190 \mu\text{m/ns}$

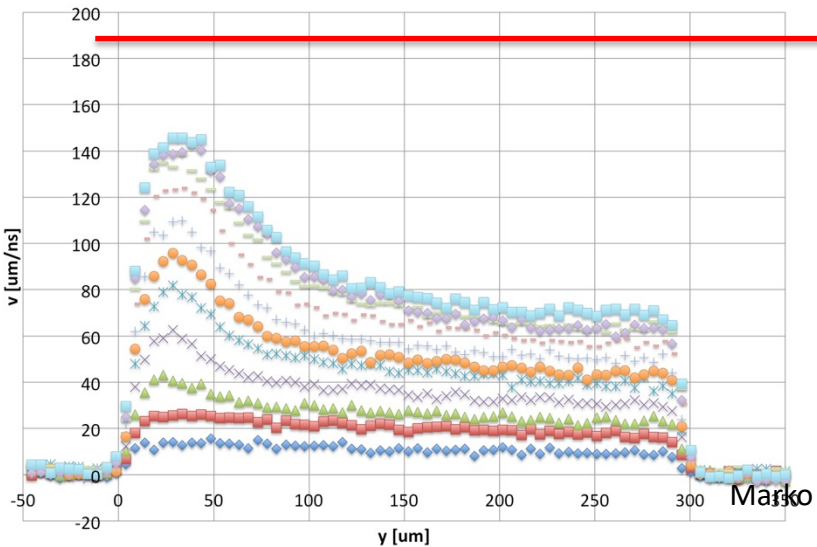
- 0
- 100
- 150
- 200
- 250
- 300
- 350
- 400
- 450
- 500
- 550
- 600
- 650
- 700
- 750

Velocity profile 1e16



- 0
- 50
- 100
- 150
- 200
- 250
- 300
- 350
- 400
- 450
- 500
- 550
- 600
- 650
- 700
- 750
- 800

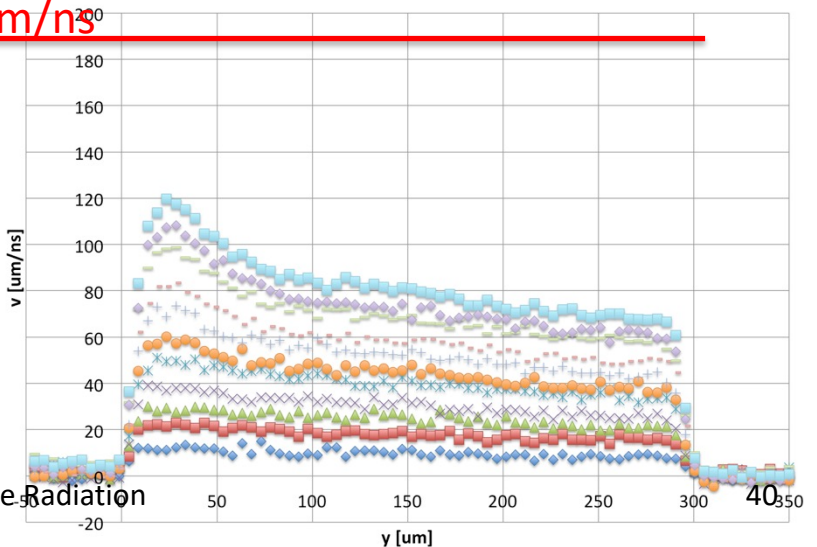
Velocity profile 5e16



$v = 190 \mu\text{m/ns}$

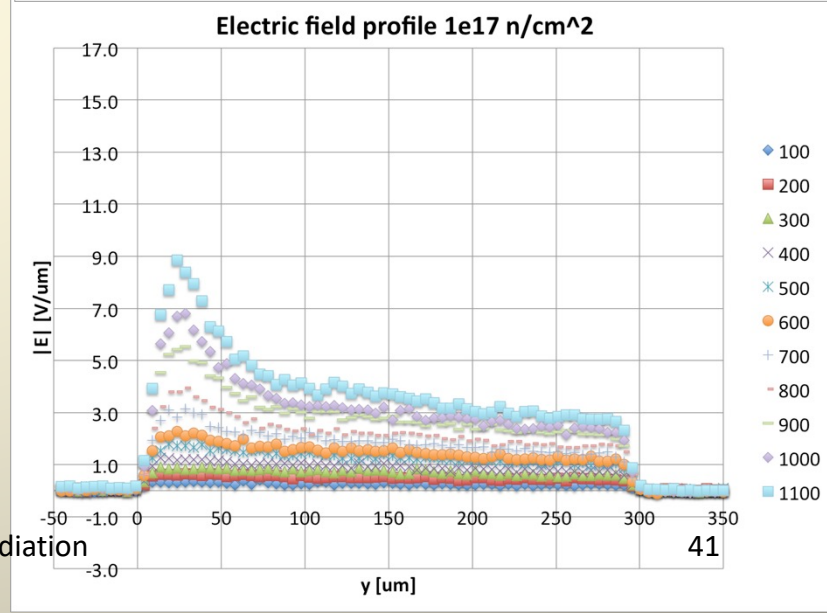
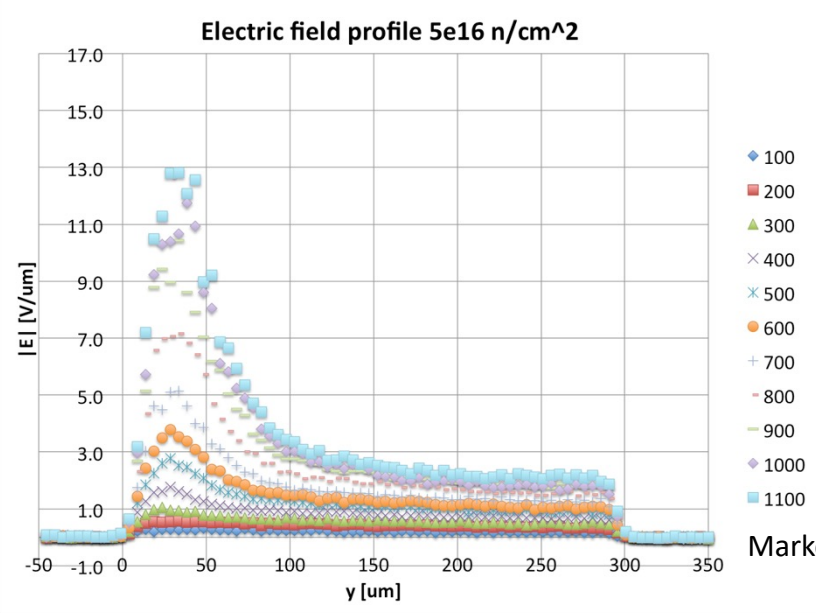
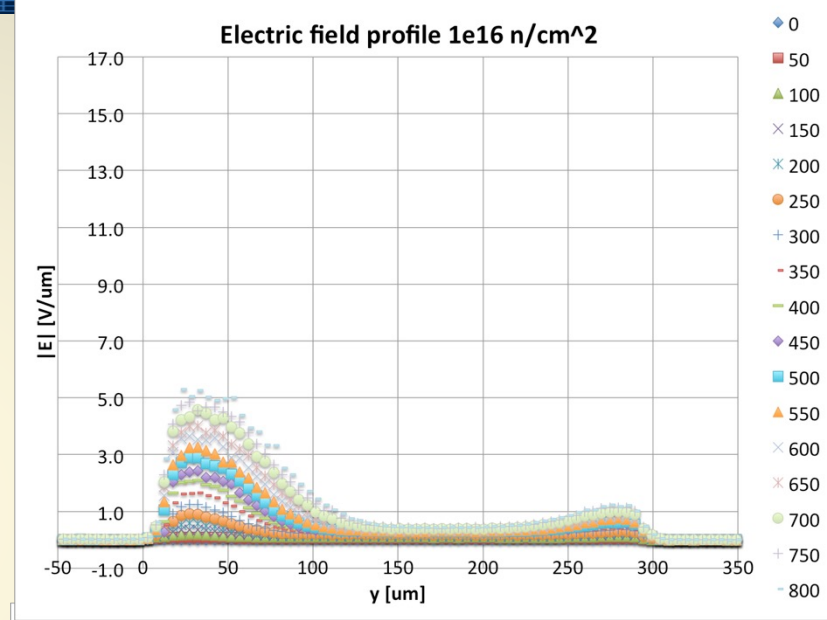
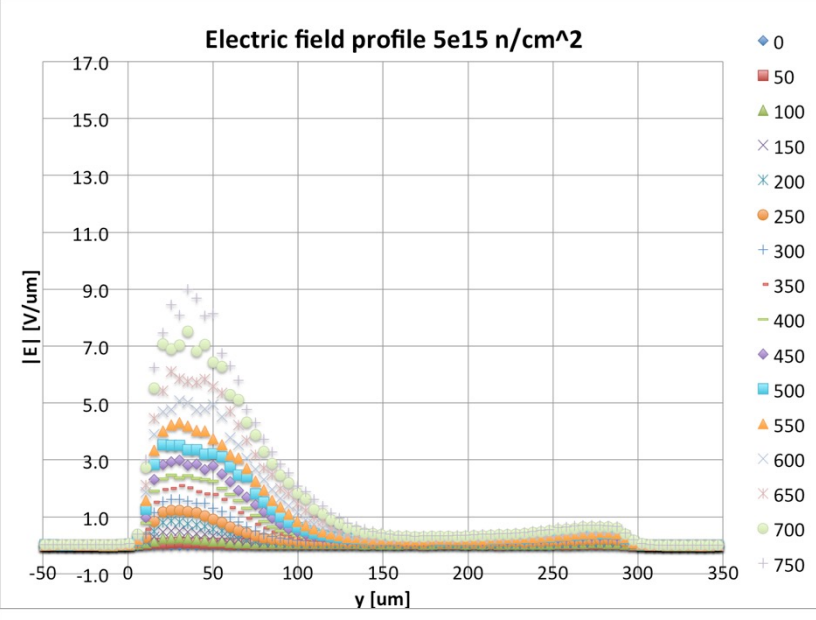
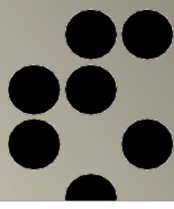
- 100
- 200
- 300
- 400
- 500
- 600
- 700
- 800
- 900
- 1000
- 1100

Velocity profile 1e17

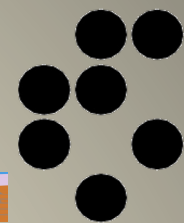


- 100
- 200
- 300
- 400
- 500
- 600
- 700
- 800
- 900
- 1000
- 1100

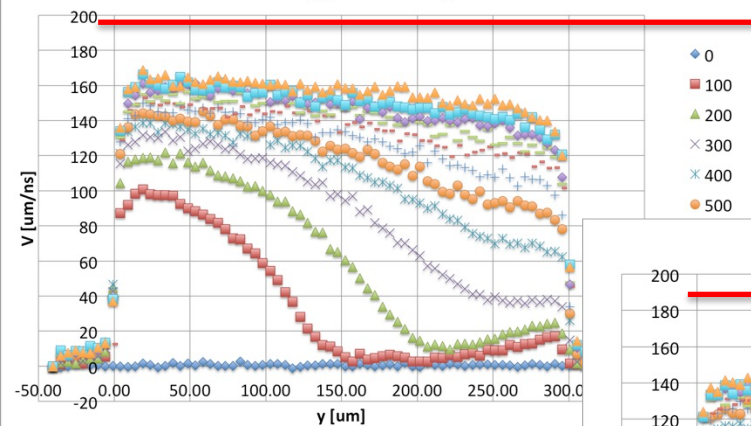
Field Profiles Neutrons



Velocity Profiles Protons

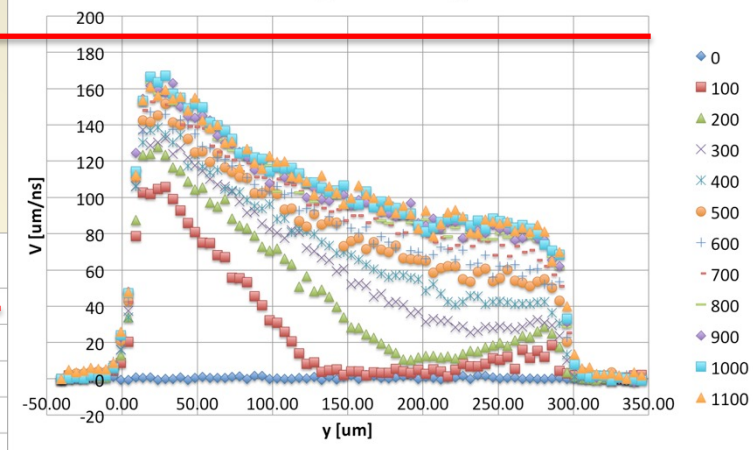


Velocity profile 5e14 p/cm²

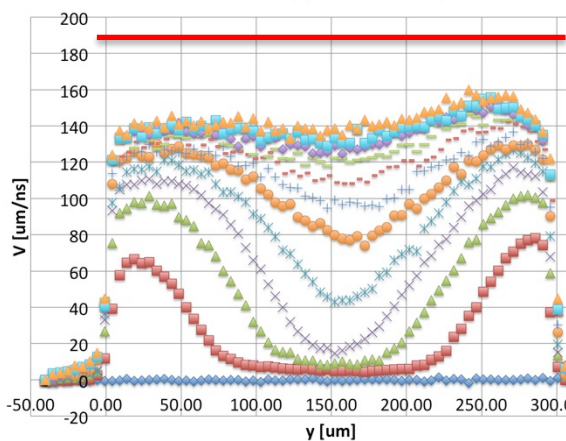


$v = 190 \mu\text{m/ns}$

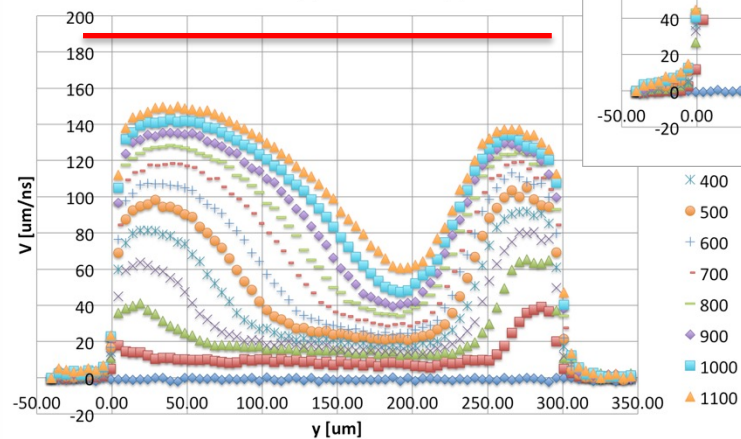
Velocity profile 1e15 p/cm²



Velocity profile 2.9e15 p/cm²

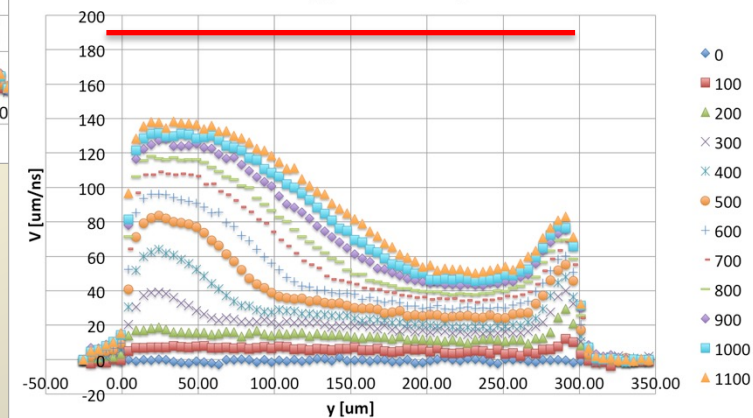


Velocity profile 1.1e16 p/cm²

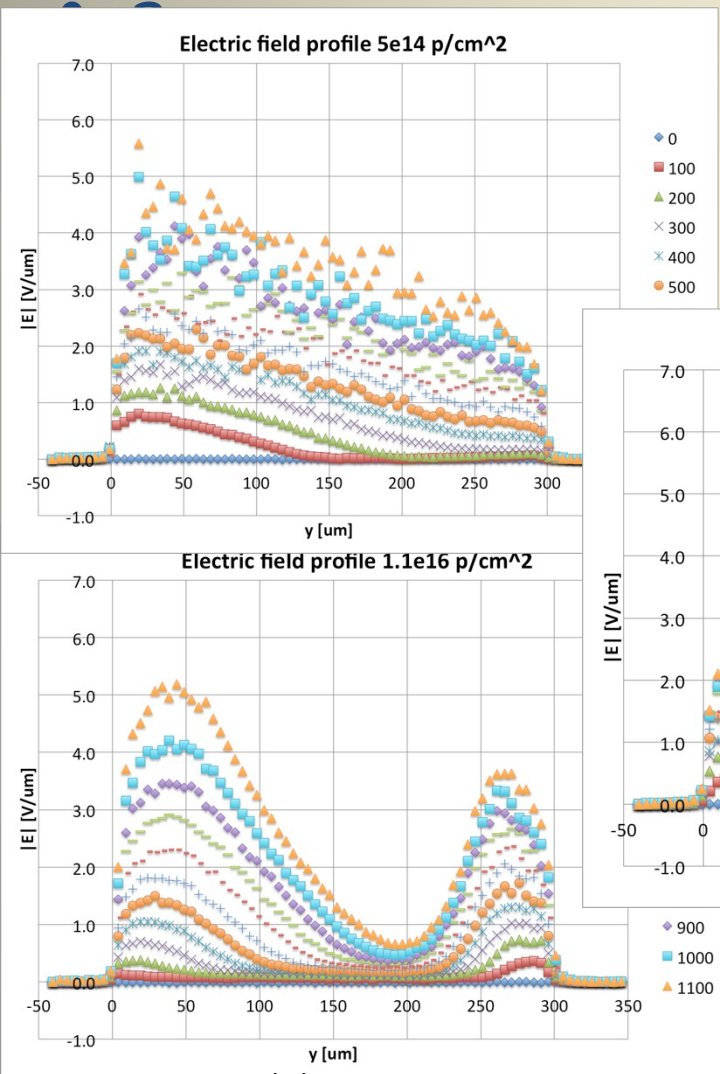
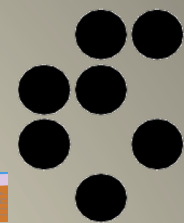


Same scale as
for neutrons

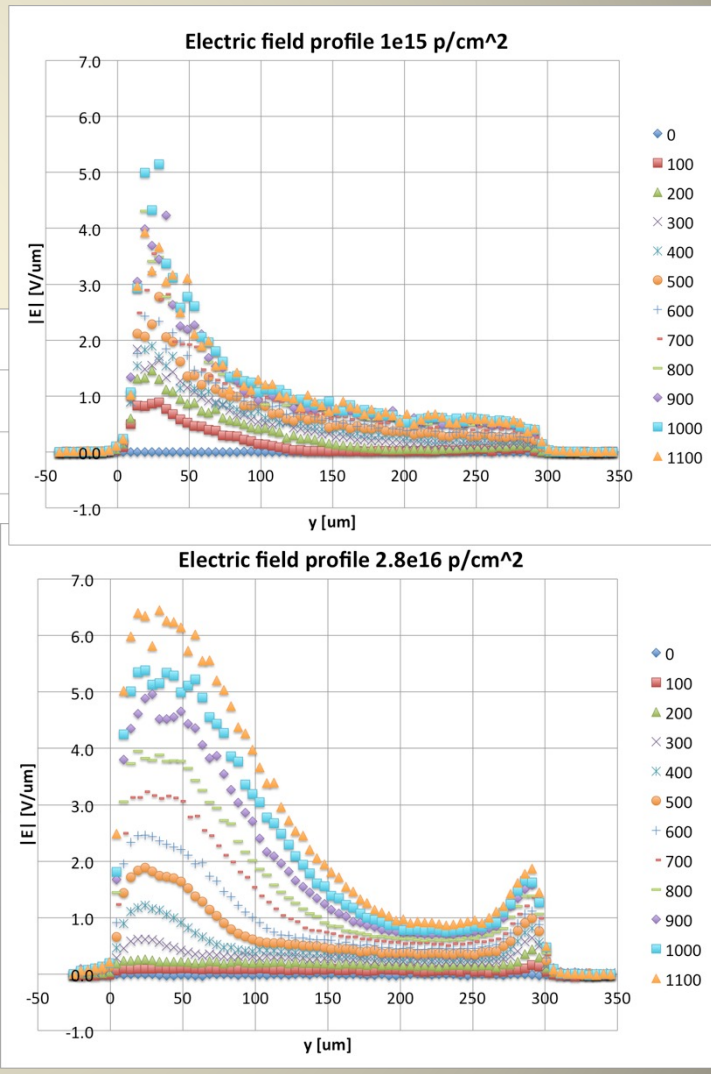
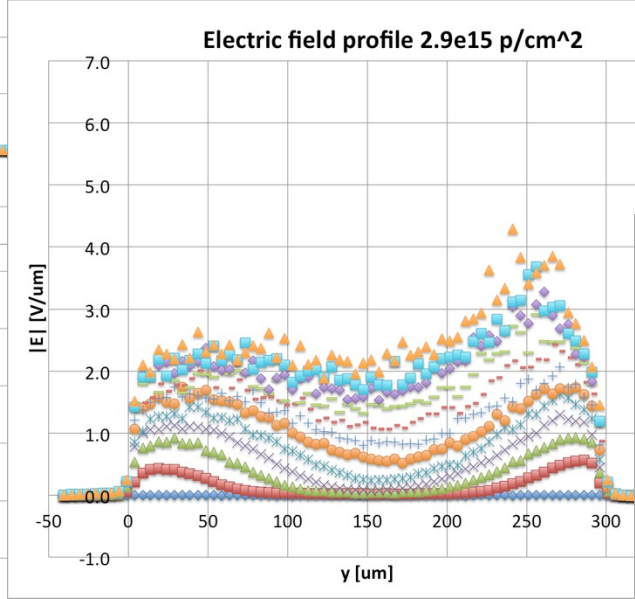
Velocity profile 2.8e16 p/cm²



Field Profiles Protons

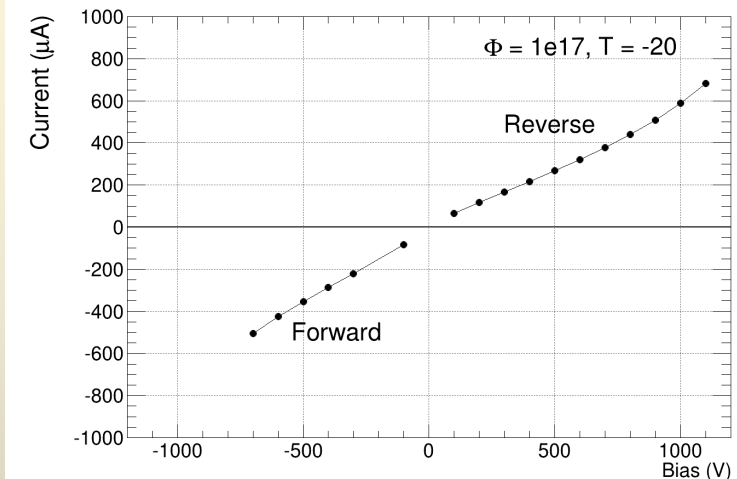
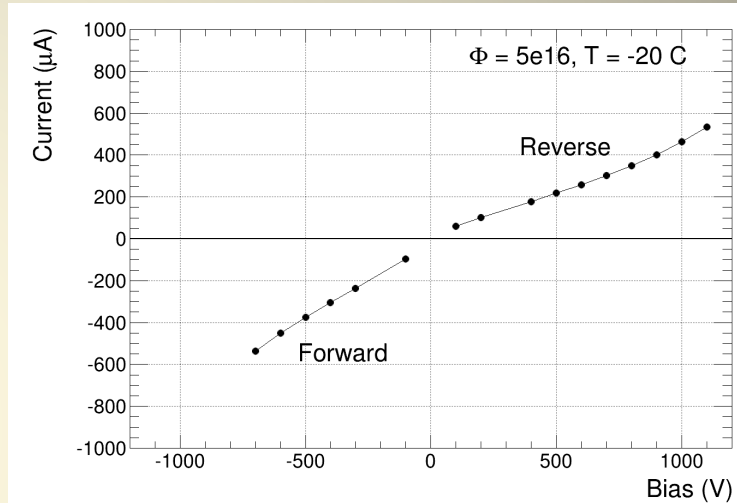


Smaller peak fields than for neutrons
Scale 0-7 V/μm

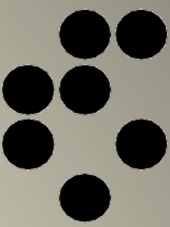


Current Characteristics

- Smooth behaviour in both directions
 - Highly resistive Si limits FW injection
- Reverse current smaller than predicted by an order of magnitude
- Both currents rising \sim linear with bias
 - Slopes FW/reverse more compatible at higher fluences
- Consistent with recent measurements at highest fluences



Trapping Considerations



- Extrapolation from low fluence data with $\beta_{e,h}(-20^\circ\text{C})=4.4, 5.8 \times 10^{-16} \text{ cm}^2/\text{ns}$; $1/\tau = \beta \Phi$

| Φ [1e15] | 5 | 10 | 50 | 100 |
|---|------|------|------|------|
| τ [ps] | 400 | 200 | 40 | 20 |
| $mfp@v_{sat}$ [μm] | 95 | 48 | 9.5 | 4.8 |
| MPV [e_0] | 7600 | 3800 | 760 | 380 |
| MPV@1000 V | 8900 | 5500 | 1800 | 1150 |
| CCD _{1000 V} [μm] | 110 | 70 | 23 | 14 |

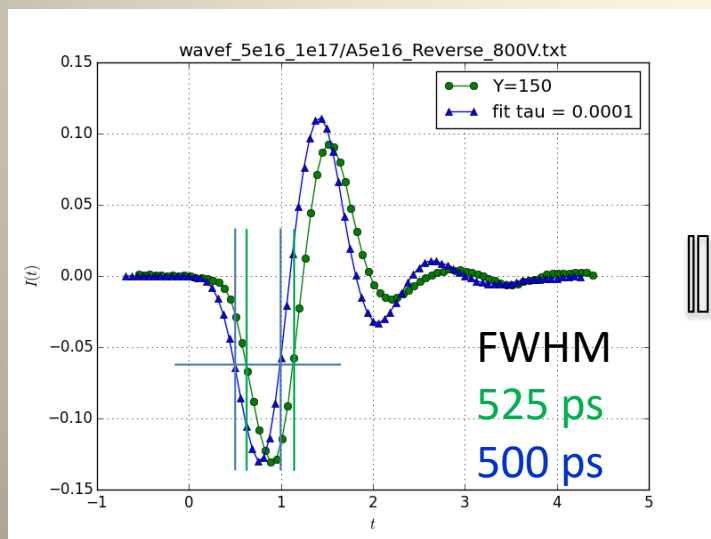
From "magic formula"
JINST 9 P10016(2014)



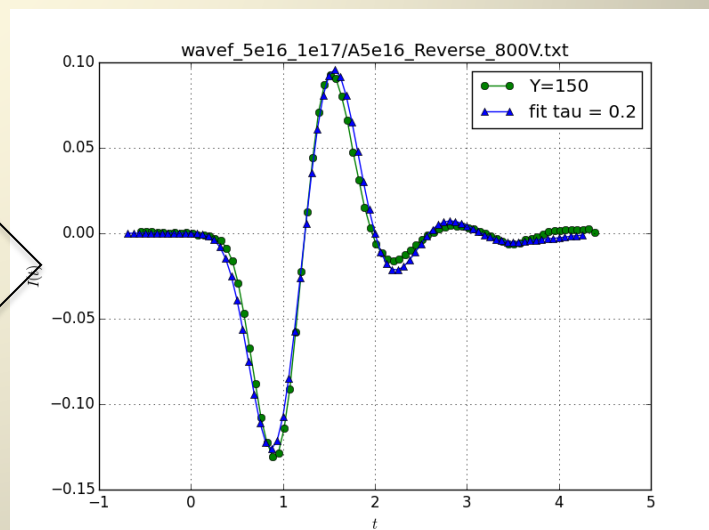
- Measured data exceeds (by far) linear extrapolation of trapping
 - n.b.1: $E \sim 3 \text{ V}/\mu\text{m}$ by far not enough to saturate velocity
 - n.b.2: little sign of CM at highest fluences

Exploiting TCT Waveforms

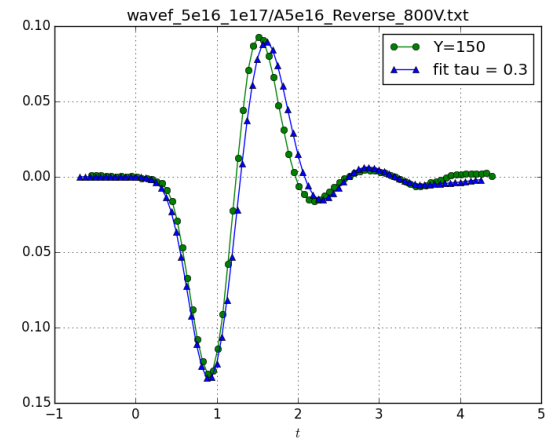
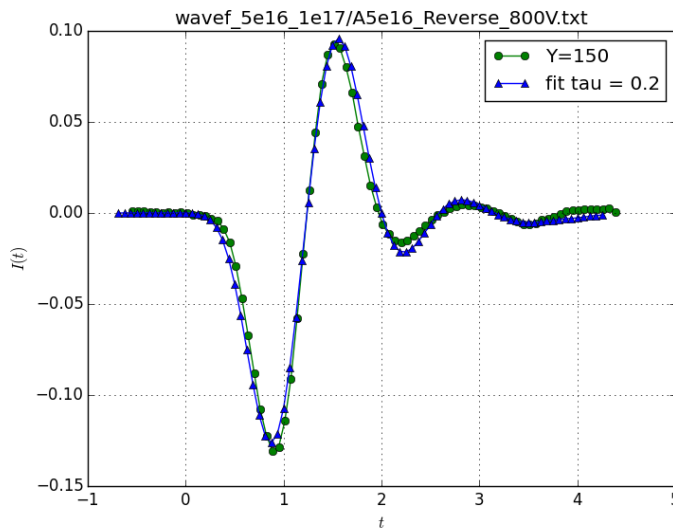
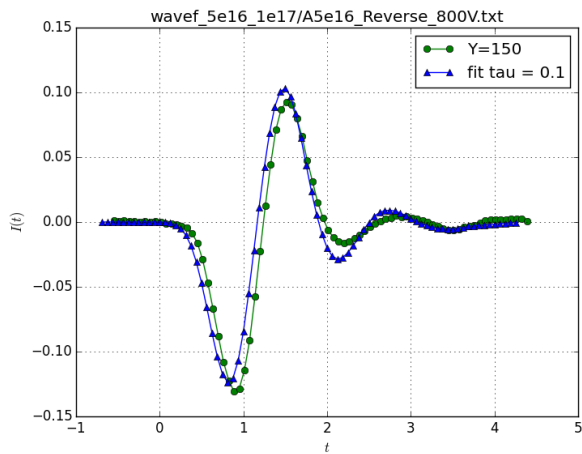
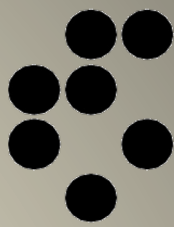
- Waveforms at $y=100\ \mu\text{m}$, 800 V, 5×10^{16} and 10^{17}
 - $E \approx 3\ \text{V}/\mu\text{m}$, CCD/2 implies signal within $\sim 10\ \mu\text{m}$ or $< 0.2\ \text{ns}$
 - the rest you see is the transfer function of the system
- Still distinct signals from the two fluences
 - treat 10^{17} waveform as transfer function of the system
 - convolute with $e^{-t/\tau}$ to match 5×10^{16} response
 - $\tau = 0.2\ \text{ns}$ provides a good match
- In fact, measure $\sim \Delta\tau$, as “transfer” already convoluted with $e^{-t/\tau(1e17)}$!



$\tau = 0.2\ \text{ns}$



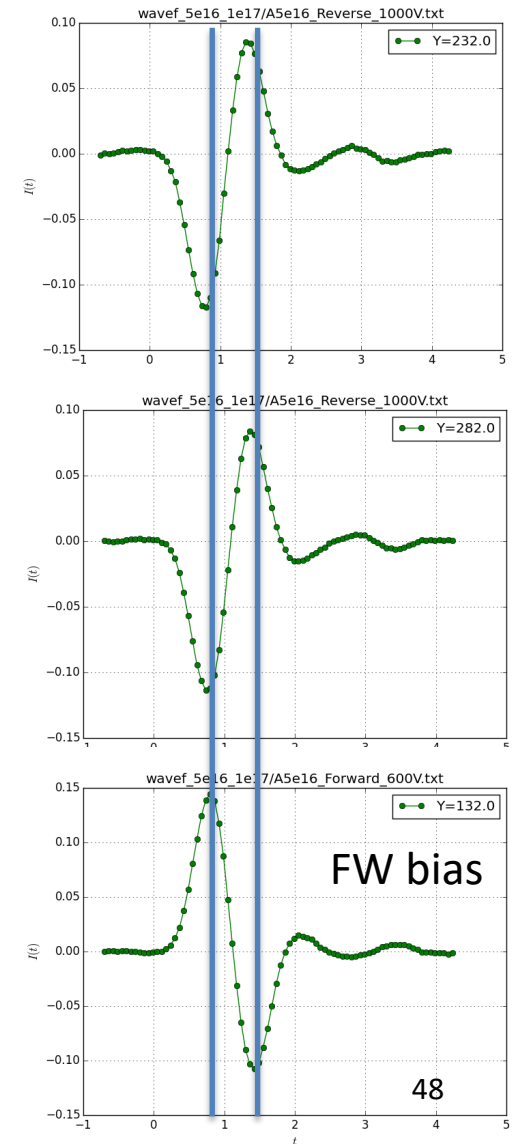
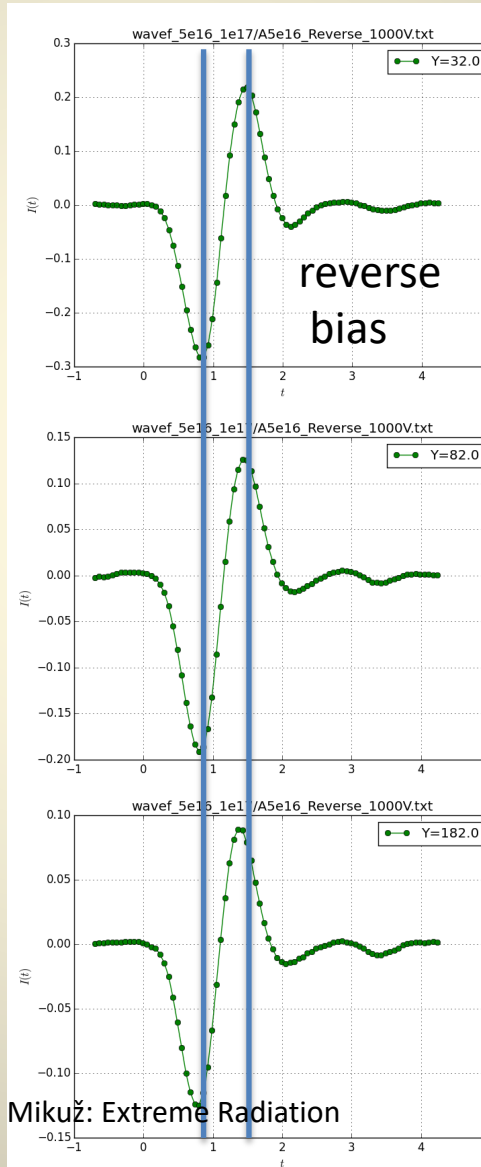
Waveforms: How sensitive ?



- $\Delta\tau = 0.2$ ns certainly best fit, 0.1 too narrow, 0.3 too broad
- precision ~ 50 ps

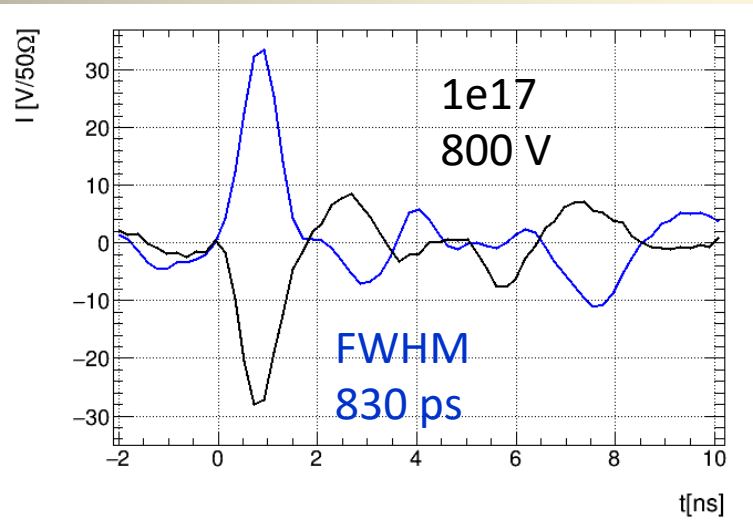
Trapping – position dependence ?

- Waveforms plotted every 50 μm in detector depth for reverse bias at 1000 V
- Forward bias in middle of detector added at 600 V
- Very little, if any, wf dependence on position observed
- Trapping not position dependent !?



Trapping @3e17

- Moved to another setup – different waveforms
 - Widths of reverse and FW similar
 - With decreases 1- \rightarrow 3e17
 - Irregular waveforms with small signal @3e17
 - Hard to state something more quantitative



Blue
Reverse

Black
FW

