

THEORY TARGETS FOR DARK MATTER EXPERIMENTS:

THE CASE OF MAGNETIC DIPOLE DARK MATTER

Based on arXiv:2512.07927 with Asher Berlin and Tanner Trickle

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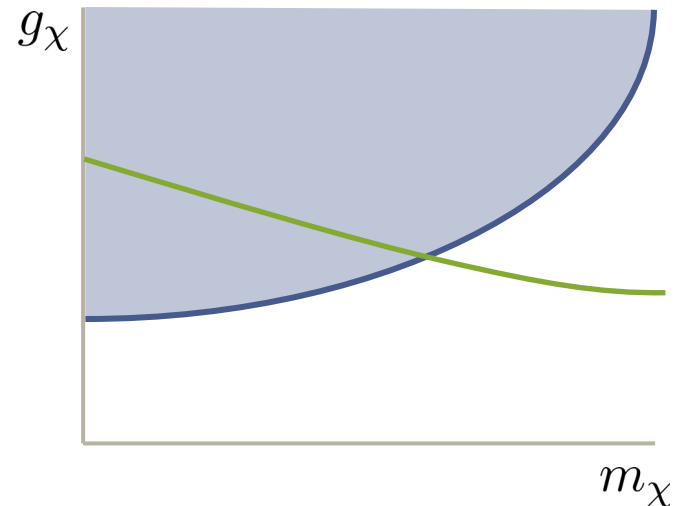
07/17/2026 PARTIQAL 2026

Sensitivity and Theory Target

χ : Dark Matter

m_χ : DM mass

g_χ : Coupling to the SM



Experimental sensitivity

What can the detector reach?

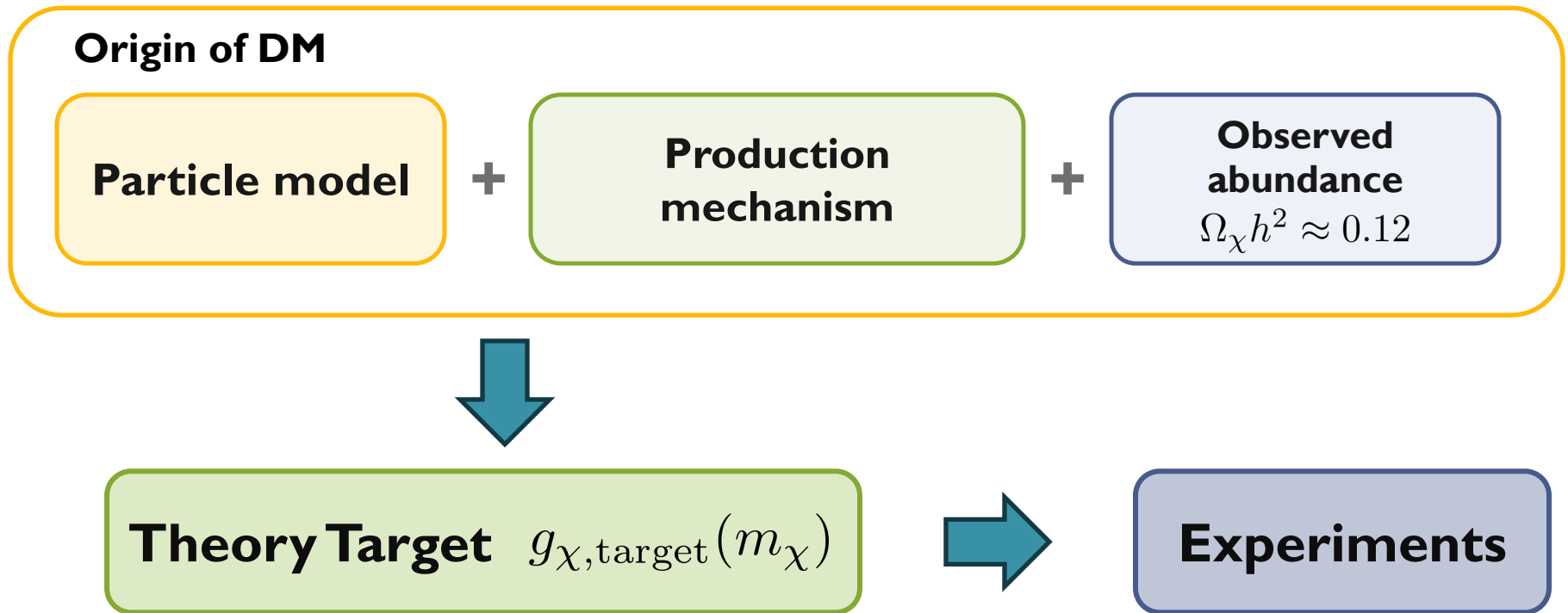
Theory target

What parameter space gives
a physical motivation?

Experimental sensitivity aims to reach theory targets

What Defines a Theory Target?

A theory target is a parameter relation predicted by a specific origin of dark matter



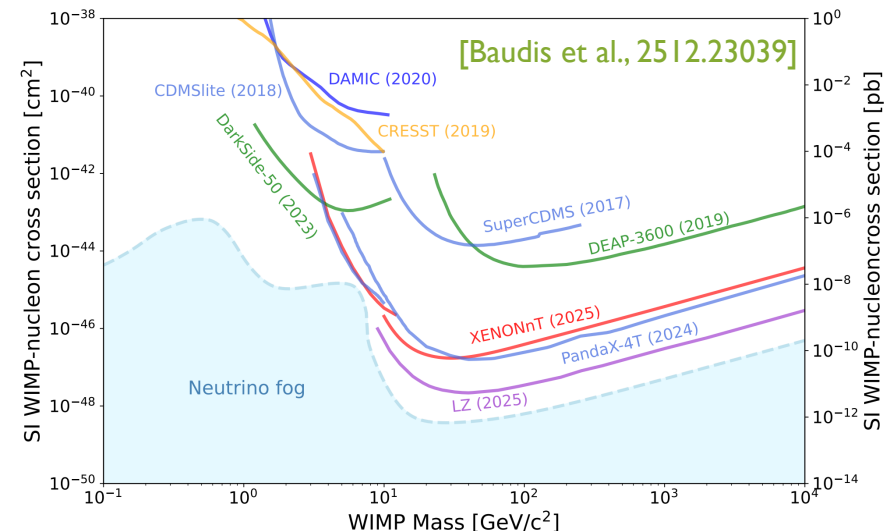
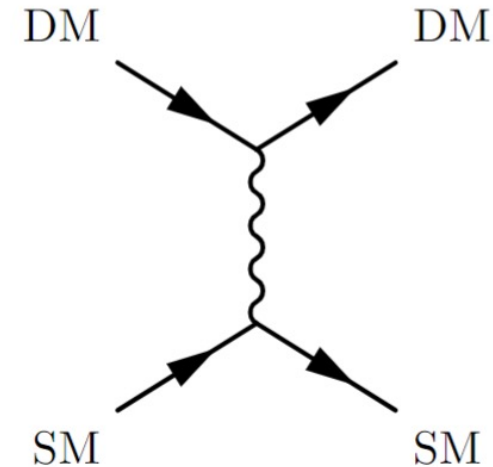
A link between the origin of DM and experiments

What Makes a Good Theory Target?

- **Predictive:** Production and detection are controlled by the same interactions, so the abundance implies a definite target
- **Robust:** Insensitive to initial conditions and the early history of the universe
- **Reachable:** Within the sensitivity of realistic (future) experiments

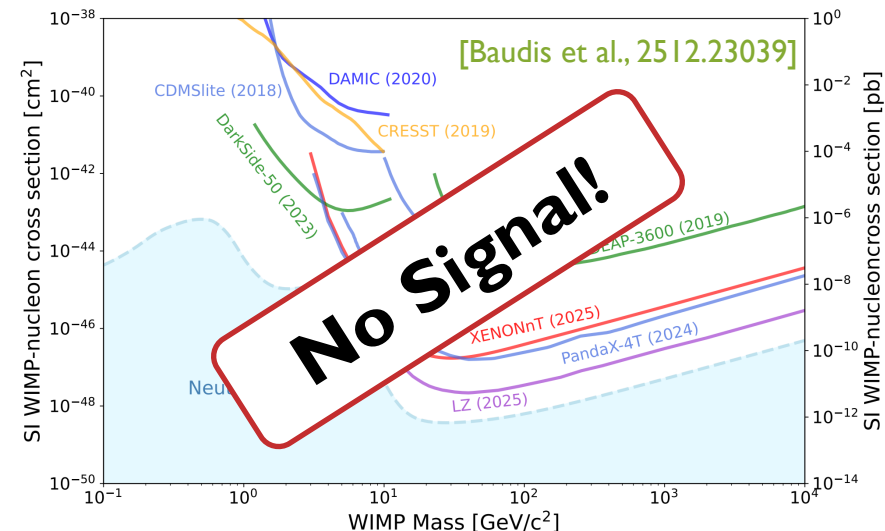
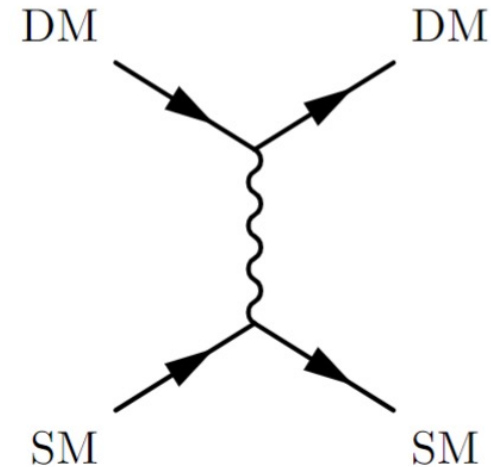
A Good Theory Target: WIMP

- Weak-scale freeze-out naturally yields the observed abundance (WIMP Miracle)
- Predictive ✓
Robust ✓
Reachable ✓
- This target drove decades of nuclear recoil searches



Testing the WIMP Target

- Nearly complete coverage down to the neutrino fog
- No signal, but a decisive result
- A good theory target does its job even without a signal
- Time to look elsewhere: new masses, new interactions

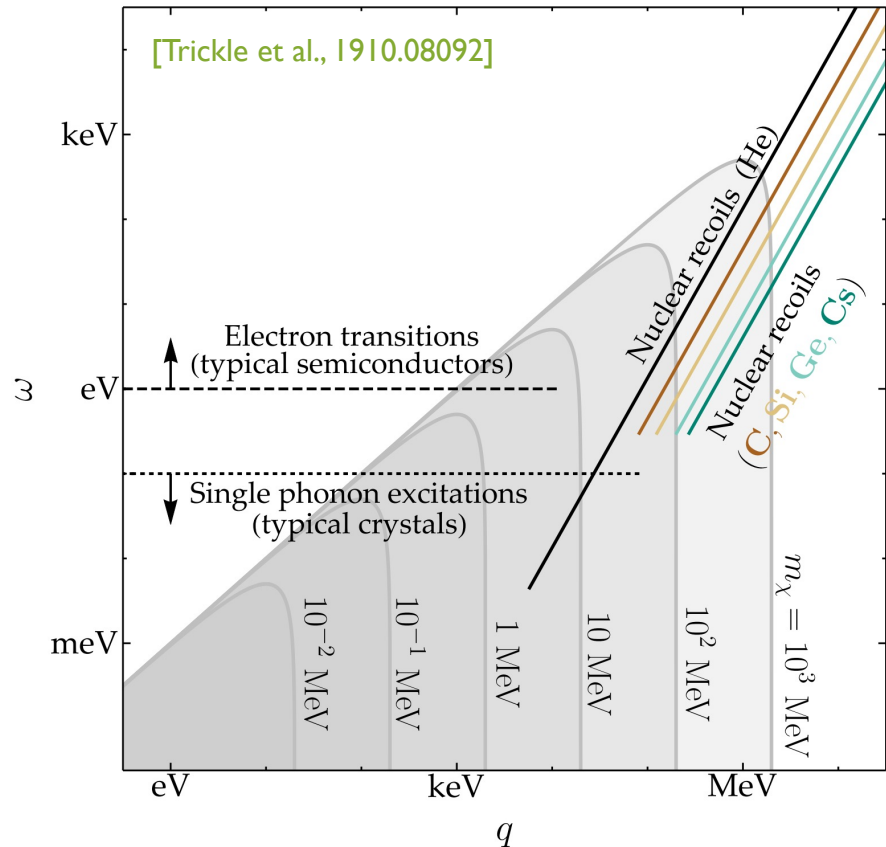
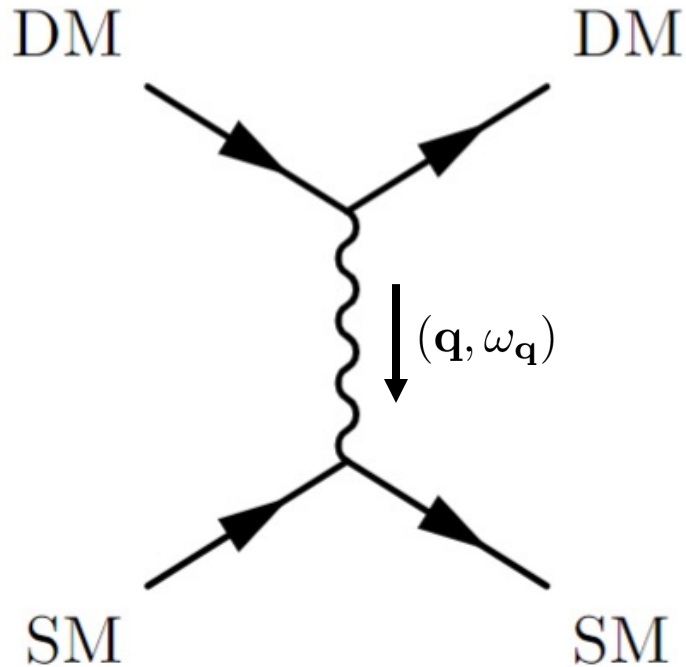


Theory Targets and Experiments Motivate Each Other



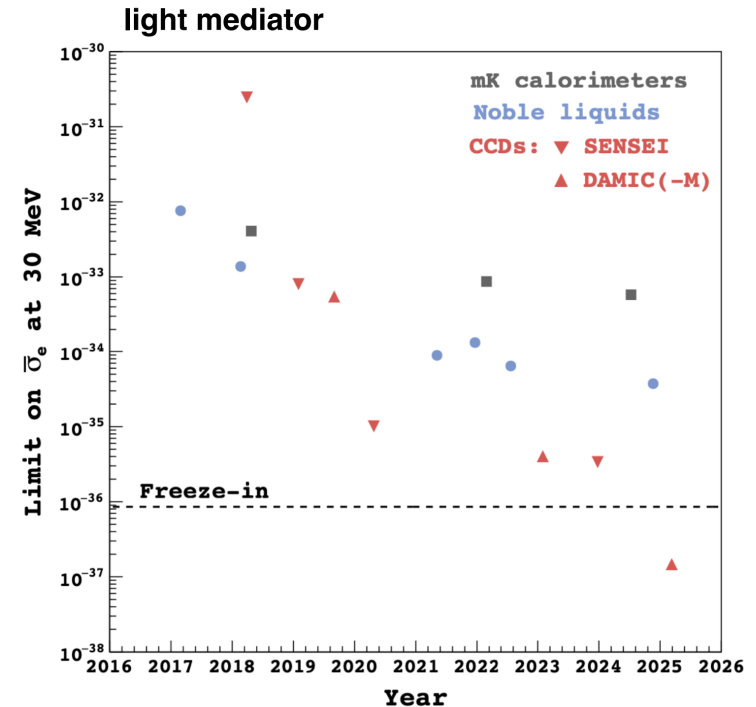
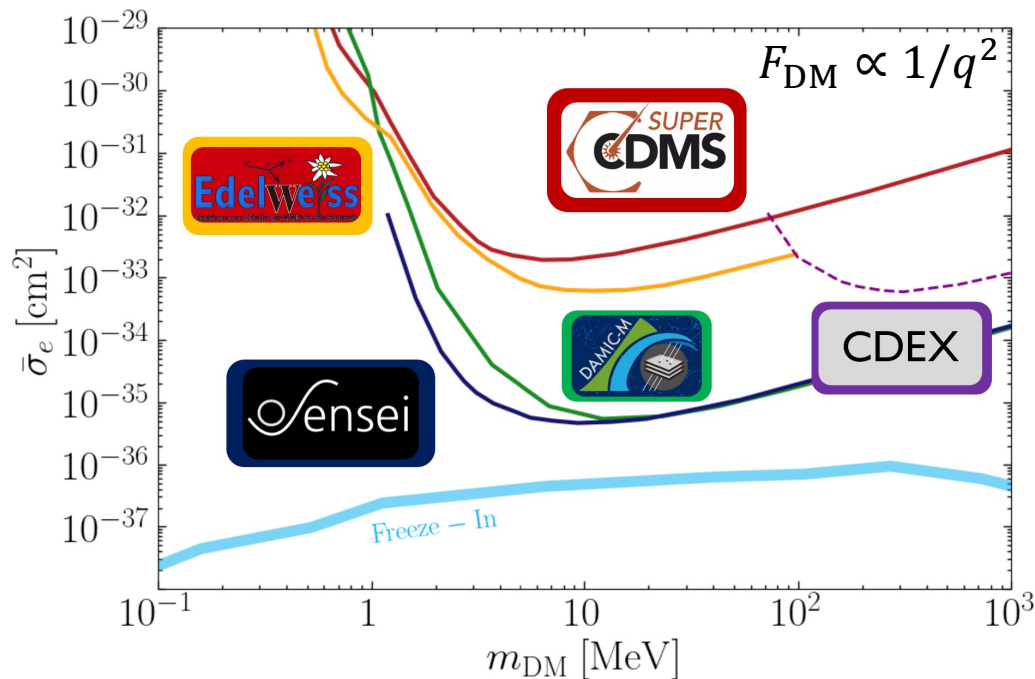
- Theory targets tell experiments where to look
- Experimental results and new technologies tell theory where to think
- The history of DM searches has been built on this interplay

Light DM Direct Detection



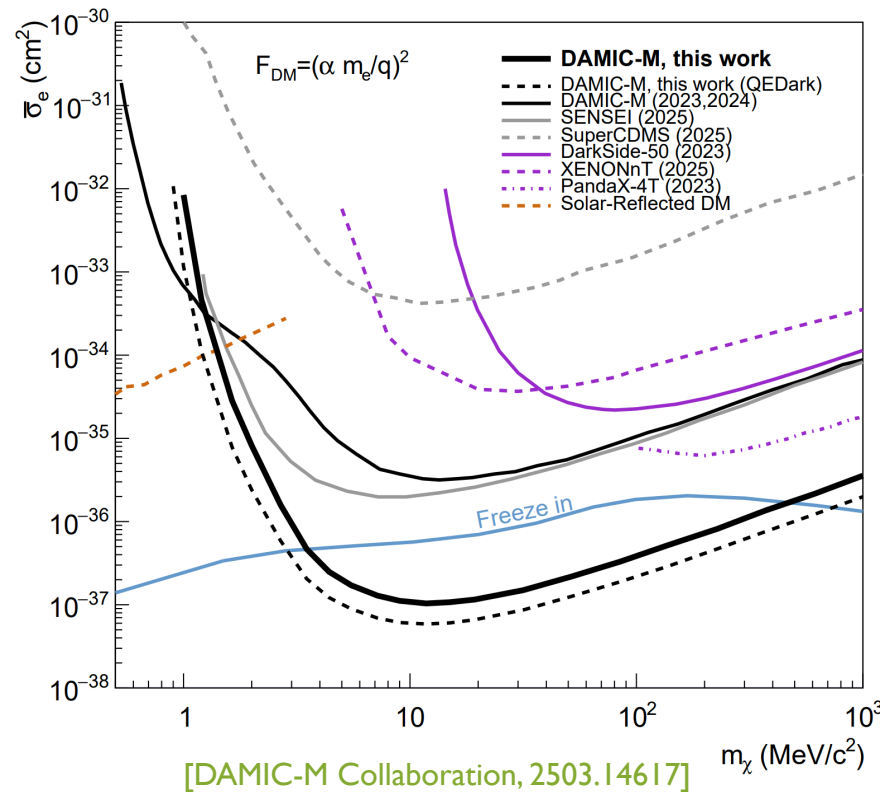
Multiple approaches to light DM detection

Theory Target for Light DM: Freeze-in



- Freeze-out is already ruled out for spin-independent (SI) interactions with a light mediator
- Freeze-in benchmark is being tested
- Predictive ✓ Robust ✓ Reachable ✓

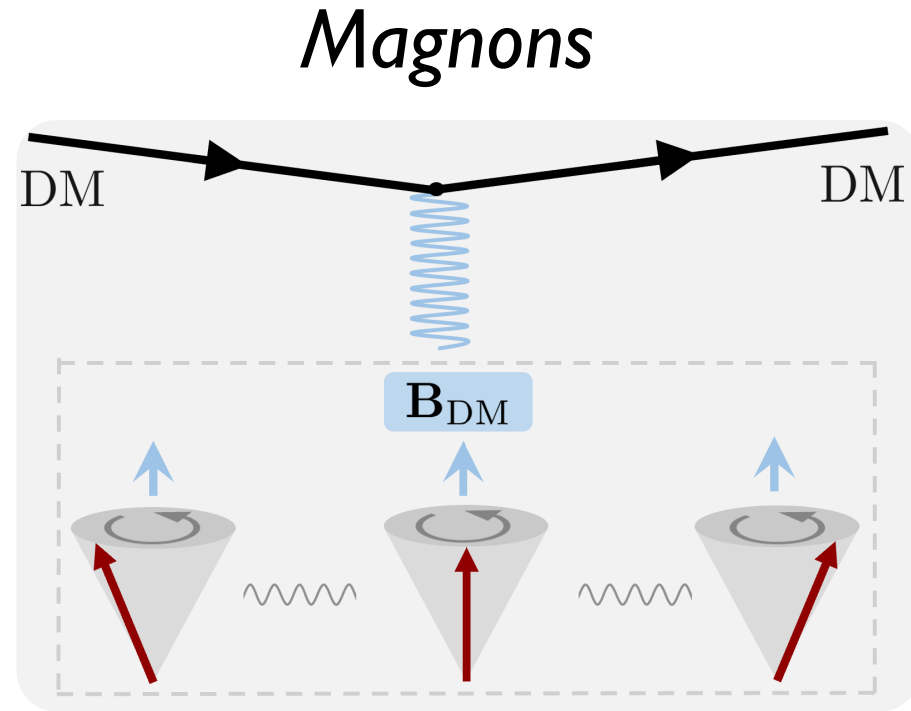
New Results From DAMIC-M



- SI freeze-in parameter space is significantly constrained
- The natural next direction is spin-dependent (SD) interactions

Detection of spin-dependent DM

- **Electron excitation**
paired spins cancel
→ SD suppressed
- **Magnon**
quantized collective spin
precession in ferromagnetically aligned spins
→ coherent SD response
- Single-magnon detection is within reach of current quantum technology

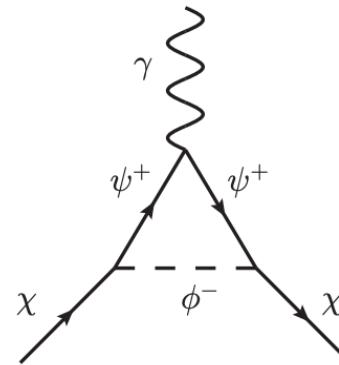
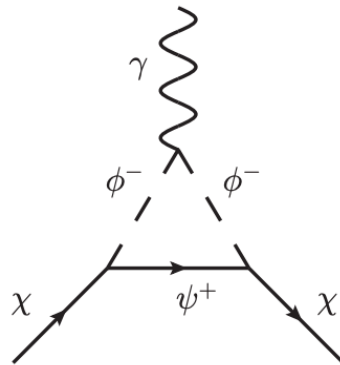


[Trickle, Zhang, Zurek, 1905.13744, 2009.13534]

[Esposito, Pavaskar, 2210.13516]

[Marocco, Wheeler, 2501.18120]

Magnetic Dipole DM



[Graham, Kaplan, Rajendran, Walters, 1203.2531]

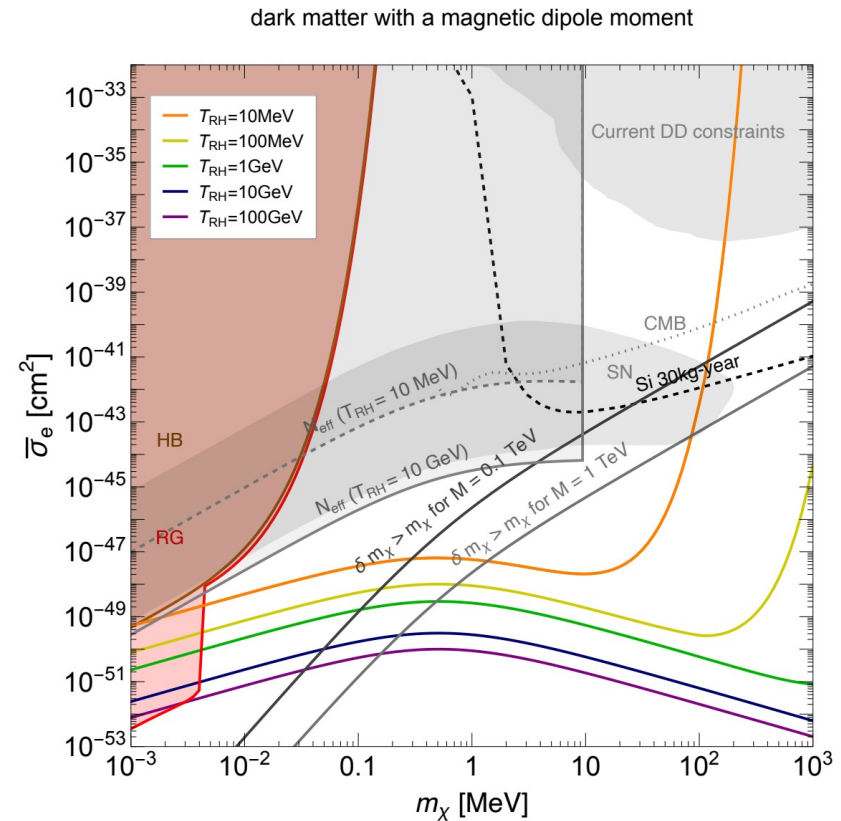
$$\mathcal{L}_{\text{MDM}} = -\frac{1}{2}\mu_\chi \bar{\chi} \sigma_{\mu\nu} \chi F^{\mu\nu}$$

- Naturally generated by heavy charged particles in a loop ($\mu_\chi \sim \frac{ey^2}{32\pi^2 M}$)
- EDM ($\bar{\chi} \sigma_{\mu\nu} \gamma^5 \chi F^{\mu\nu}$) can be obtained from the same diagram with a CP-violating phase
- MDM operator is $\vec{S}_\chi \cdot \vec{B}$ in the non-relativistic limit $\rightarrow$ spin-dependent

Production of Magnetic Dipole DM

$$\mathcal{L}_{\text{MDM}} = -\frac{1}{2}\mu_\chi \bar{\chi} \sigma_{\mu\nu} \chi F^{\mu\nu}$$

- Dimension-5 operator
- UV Freeze-in
→ Depends on T_{RH}
- Predictive ✓
Robust ✗
Reachable Δ

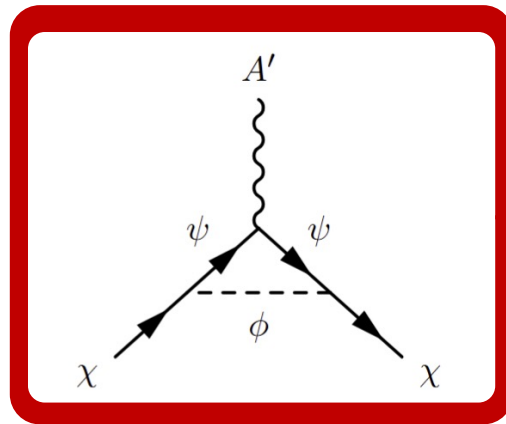


[JHC, Essig, Reinert, 1911.03389]

Magnetic Dipole DM Model in this Work

$$\mathcal{L} \supset A'_\mu \left[g'(\bar{\psi}\gamma^\mu\psi) + g'(\phi^*\partial^\mu\phi) \right] + y'\phi\bar{\psi}\chi$$

Dark sector



$$\mu'_\chi = \frac{g'y'^2}{32\pi^2 M}$$

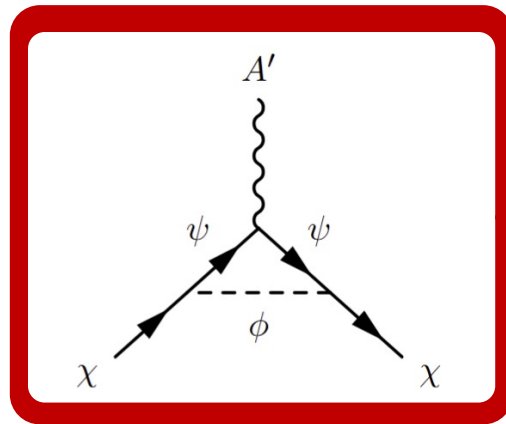
ψ, ϕ Dark Charged Particles

χ Dark Matter

$$\mathcal{L} \supset \frac{\mu'_\chi}{2} F'_{\mu\nu} (\bar{\chi}\sigma^{\mu\nu}\chi)$$

Magnetic Dipole DM Model in this Work

$$\mathcal{L} \supset \underbrace{A'_\mu \left[g'(\bar{\psi}\gamma^\mu\psi) + g'(\phi^*\partial^\mu\phi) \right]}_{\text{Dark sector}} + \underbrace{y'\phi\bar{\psi}\chi + \frac{\varepsilon}{2}F'_{\mu\nu}F^{\mu\nu}}_{\text{Interaction (Kinetic mixing)}}$$



Dark sector

$$\mu'_\chi = \frac{g'y'^2}{32\pi^2 M}$$

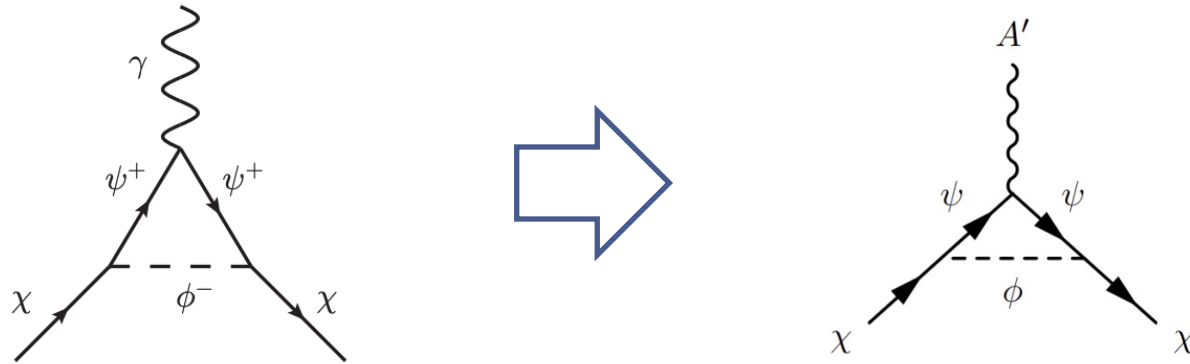
Interaction
(Kinetic mixing)

$$\mu_\chi = \varepsilon\mu'_\chi$$

$$\mathcal{L} \supset \frac{\mu'_\chi}{2} F'_{\mu\nu} (\bar{\chi}\sigma^{\mu\nu}\chi) + \varepsilon e A'_\mu (\bar{e}\gamma^\mu e)$$

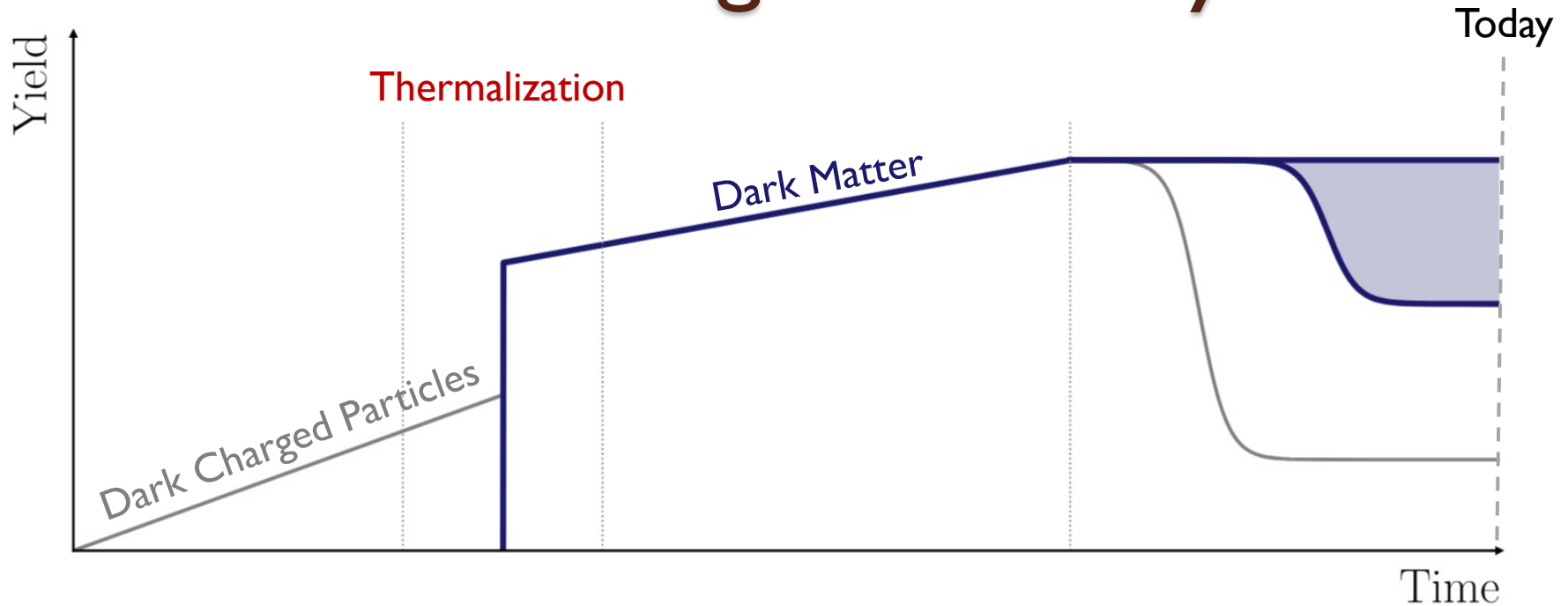
Magnetic Dipole-Coupled Interaction Lagrangian

Magnetic Dipole DM Model in this Work



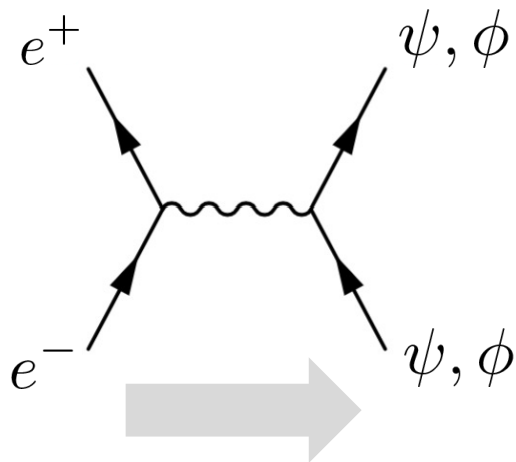
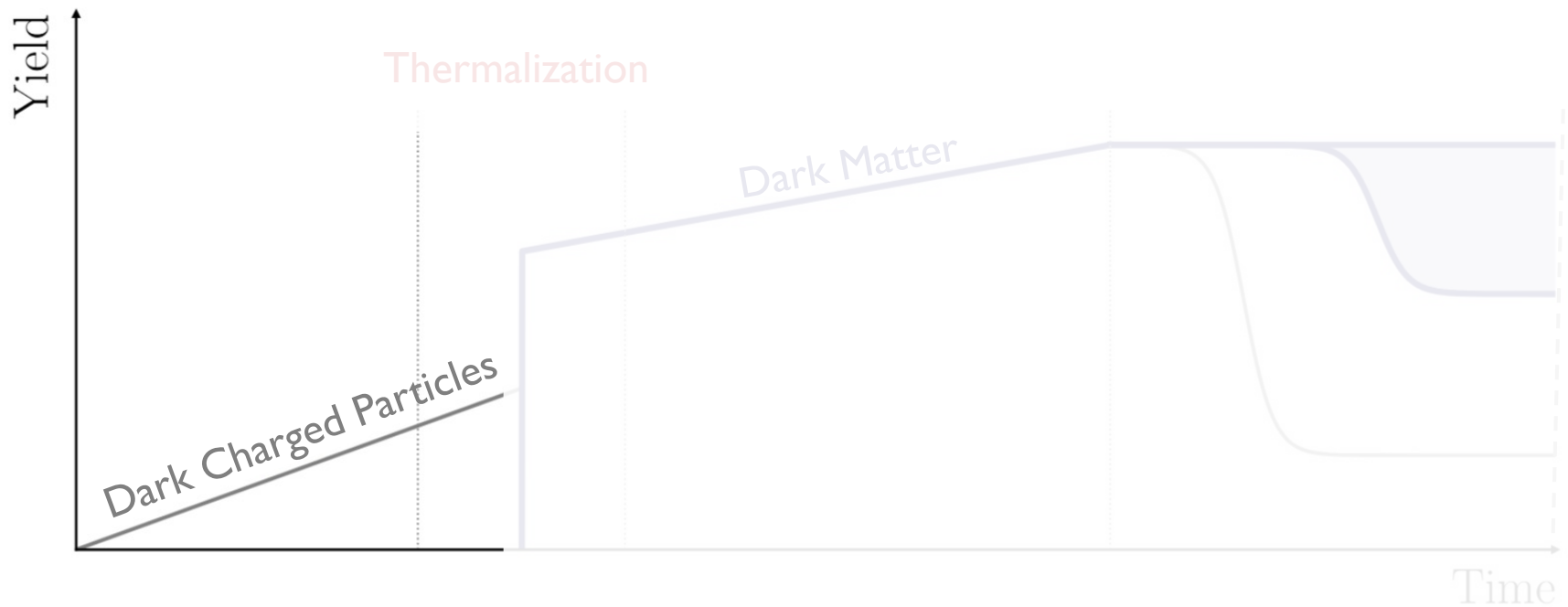
- In the massless A' limit, DCPs (ψ , ϕ) are millicharged particles
- DCPs can be light while χ acquires an MDM at low energies
- Smallness of MDM comes from the small mixing angle ε
- Predictive ✓ Robust ✓ Reachable ✓

Cosmological History



- DCP Freeze-in $\rightarrow$ DS Thermalization $\rightarrow$ DS Evolution $\rightarrow$ DM Freeze-out
- This structure was first explored as “reannihilation” [Chu, Hambye, Tytgat, 1112.0493]
- Further developed in “glaciation” [2111.13709] and “dark sink” [2312.14152, 2408.07744]
- In this work, DM is produced through thermalization, not direct freeze-in

I. DCPs Freeze-in

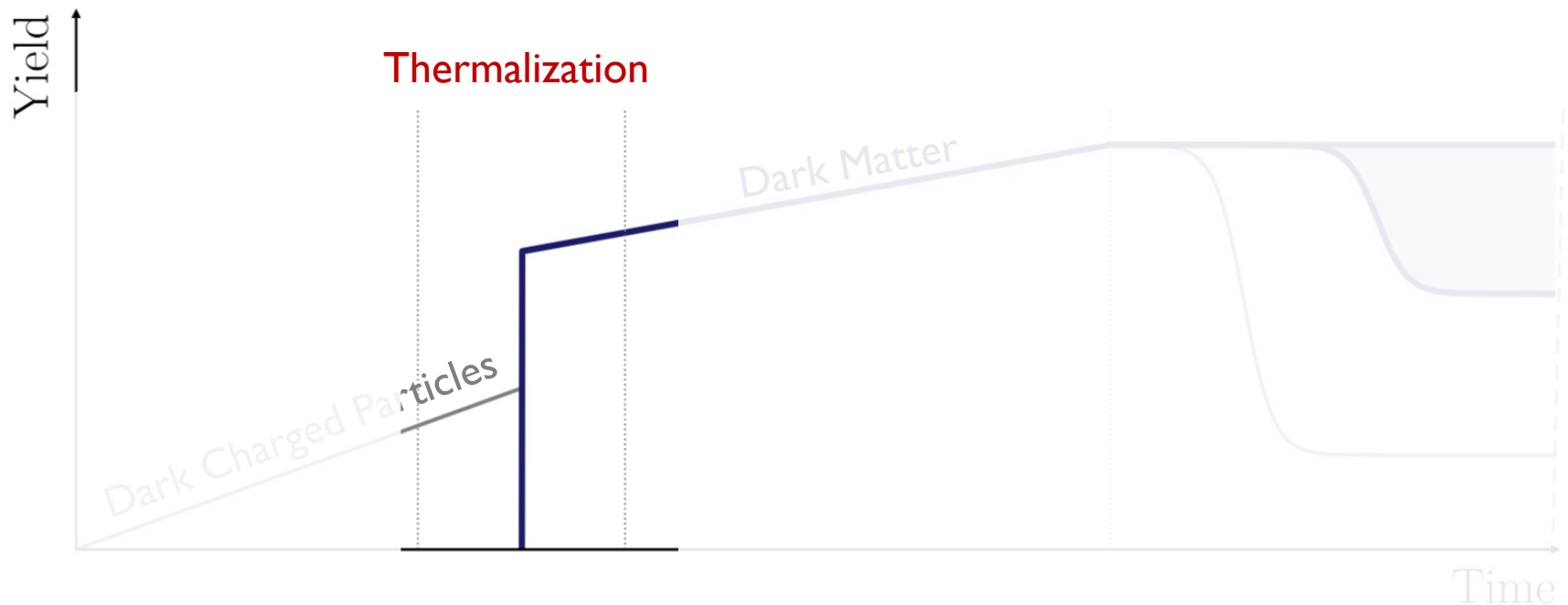


$$Y_X \sim 10^{-4} e^4 \kappa^2 \frac{M_{\text{Pl}}}{T}$$

$$X = \psi, \phi \quad \kappa \equiv \frac{\varepsilon g'}{e}$$

$$m_\psi \approx m_\phi \equiv M > m_\chi$$

2. DS Thermalization



$$\phi \leftrightarrow \psi\chi \quad , \quad \phi\phi \leftrightarrow A'A' \quad , \quad \psi\psi \leftrightarrow A'A' \quad , \quad \phi\phi \leftrightarrow \chi\chi$$

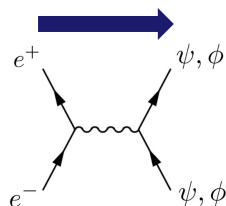
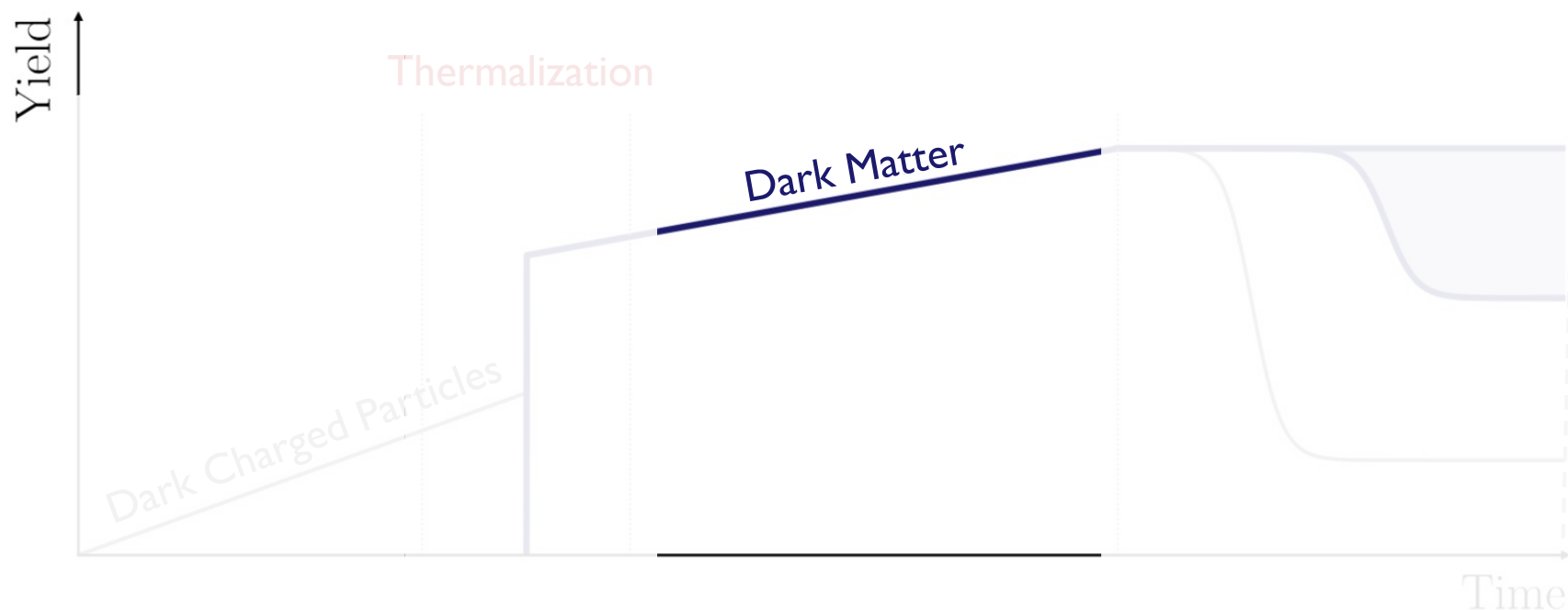
All processes must be faster than Hubble $\Gamma > H$

A few high-energy particles $\rightarrow$ Many low-energy particles

(DS Thermalization from the Sun: Dark Solar Wind)

[JHC, Kaplan, Rajendran, Ramani, Tanin, 2205.11527]

3. DS Evolution



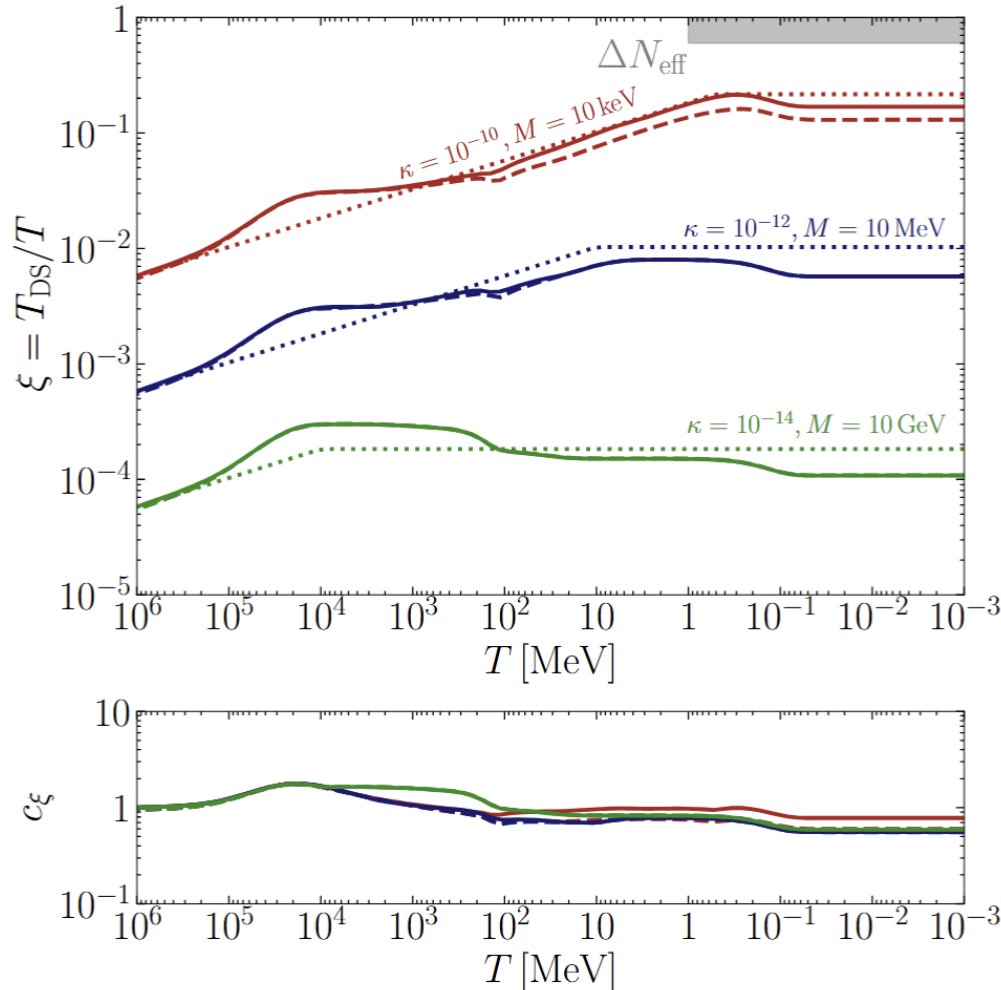
$$\rho_{\text{DS}} \sim T_{\text{DS}}^4$$

$$\frac{d(\rho_{\text{DS}}/s^{4/3})}{dT} \propto -(e^2 \kappa)^2 \frac{M_{\text{Pl}}}{T^2}$$

Energy density
Boltzmann equation

$$\frac{T_{\text{DS}}}{T} \sim \sqrt{e^2 \kappa} \left(\frac{M_{\text{Pl}}}{T} \right)^{1/4}$$

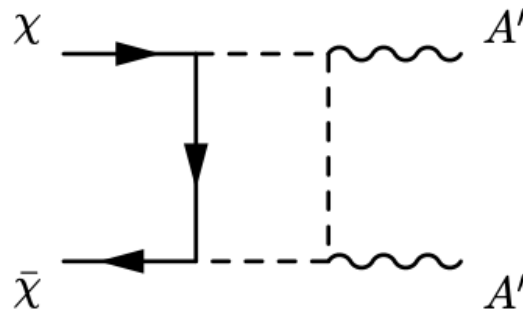
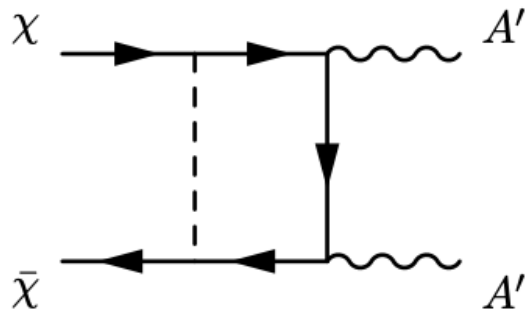
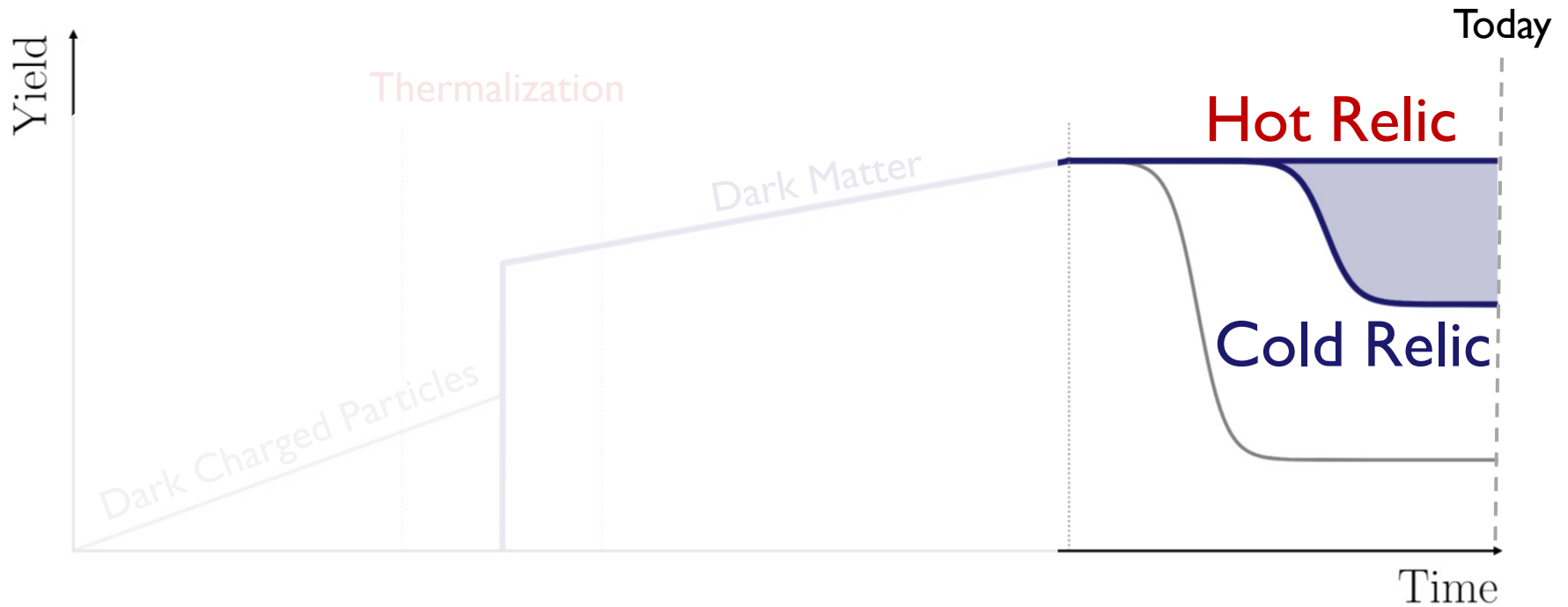
3. DS Evolution



- DS is always colder than SM ($\xi < 1$)
- The the exact shape is determined by various processes and $g_\star$
- Energy injection stops if SM cannot produce DCPs

$$\frac{\rho_{\text{DS}}}{s^{4/3}} = 10^{-4} c_\rho e^4 \kappa^2 \left(\frac{M_{\text{Pl}}}{\max\{T, M, m_e\}} \right)$$

4. DM Freeze-out



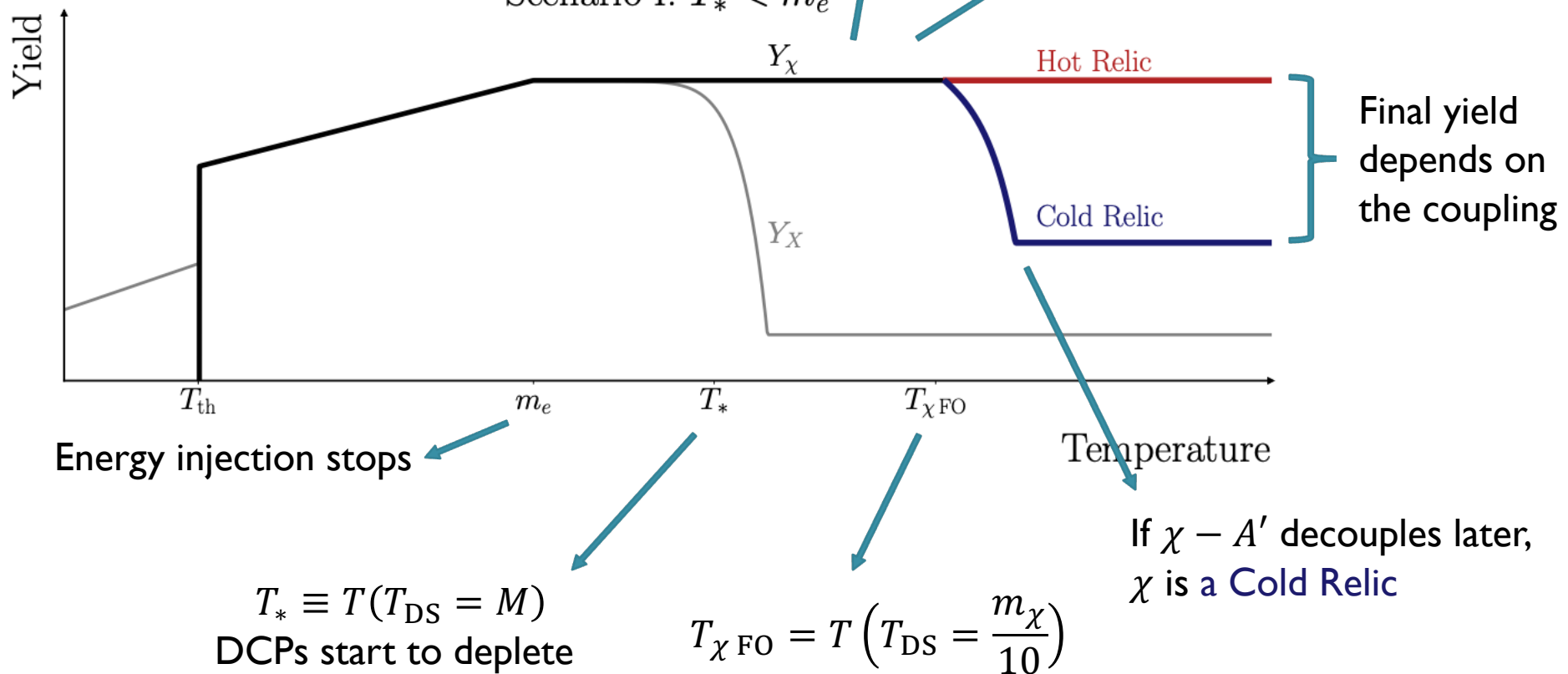
$$\langle \sigma v \rangle_{\chi\bar{\chi} \rightarrow A'A'} \sim \frac{y'^4 g'^4}{768\pi^5} \frac{m_\chi^4}{M^6}$$

4. DM Freeze-out

$$\xi \equiv \frac{T_{\text{DS}}}{T} \sim \sqrt{e^2 \kappa} \left(\frac{M_P}{m_e} \right)^{1/4}$$

Scenario I: $T_* < m_e$

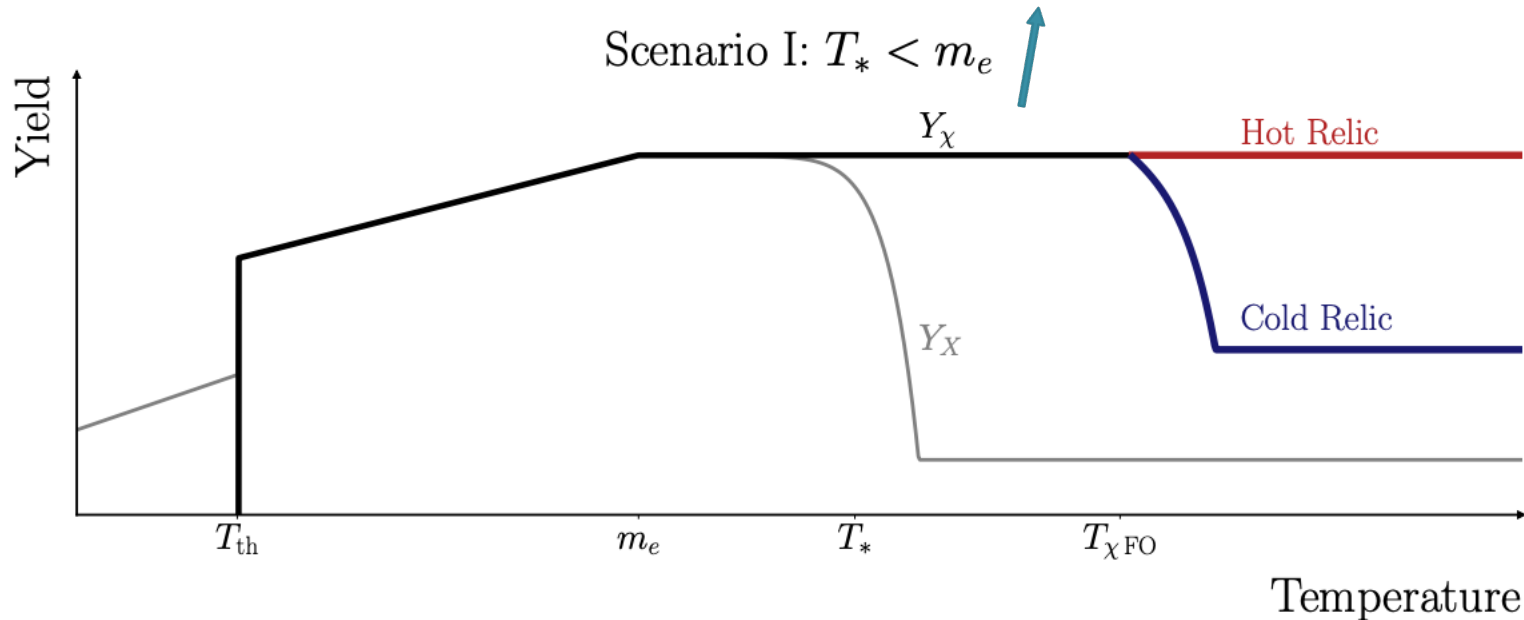
If $\chi - A'$ decouples earlier,
 χ is a **Hot Relic**



4. DM Freeze-out

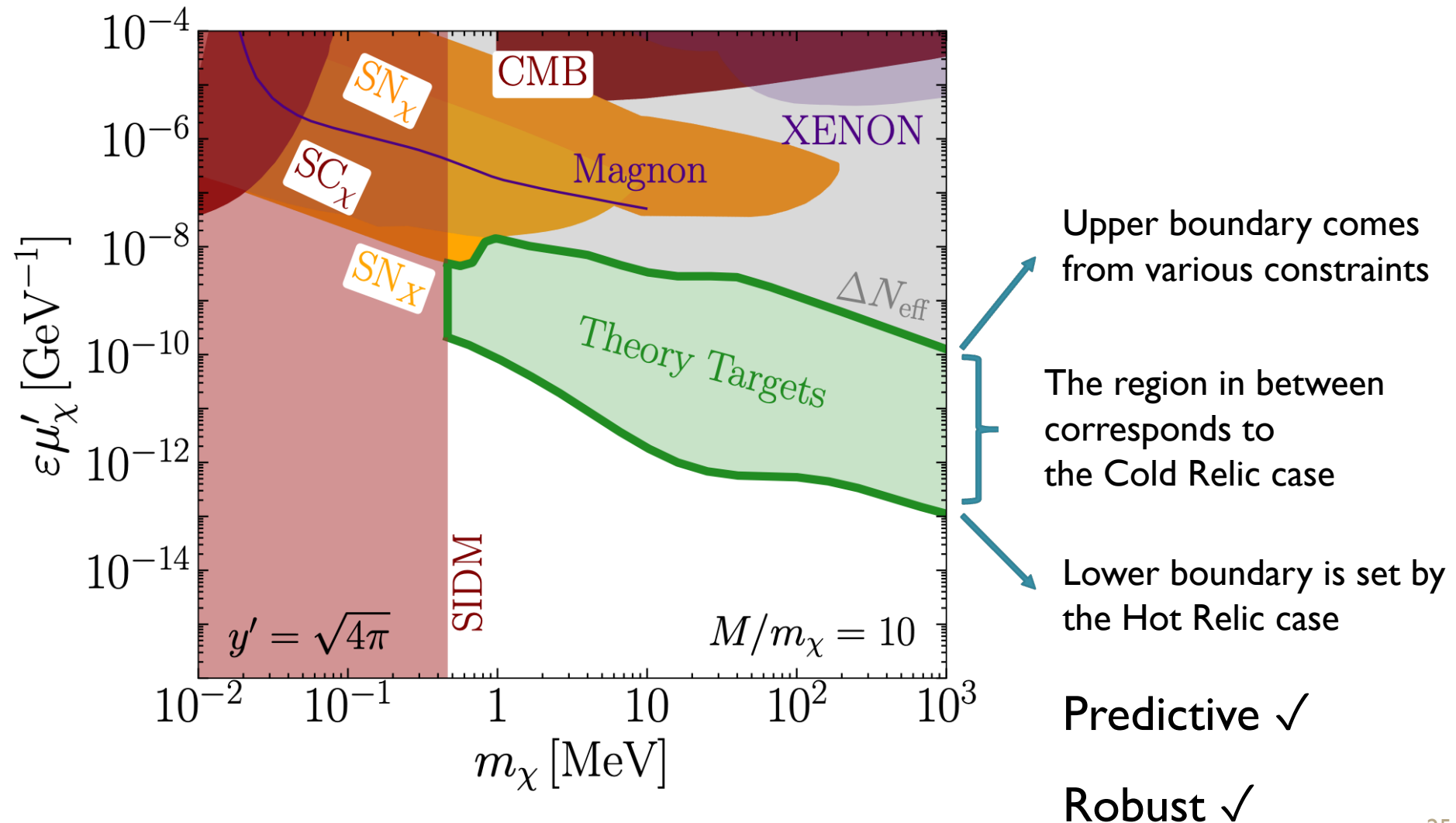
$$\xi \equiv \frac{T_{\text{DS}}}{T} \sim \sqrt{e^2 \kappa} \left(\frac{M_{\text{Pl}}}{m_e} \right)^{1/4}$$

Scenario I: $T_* < m_e$

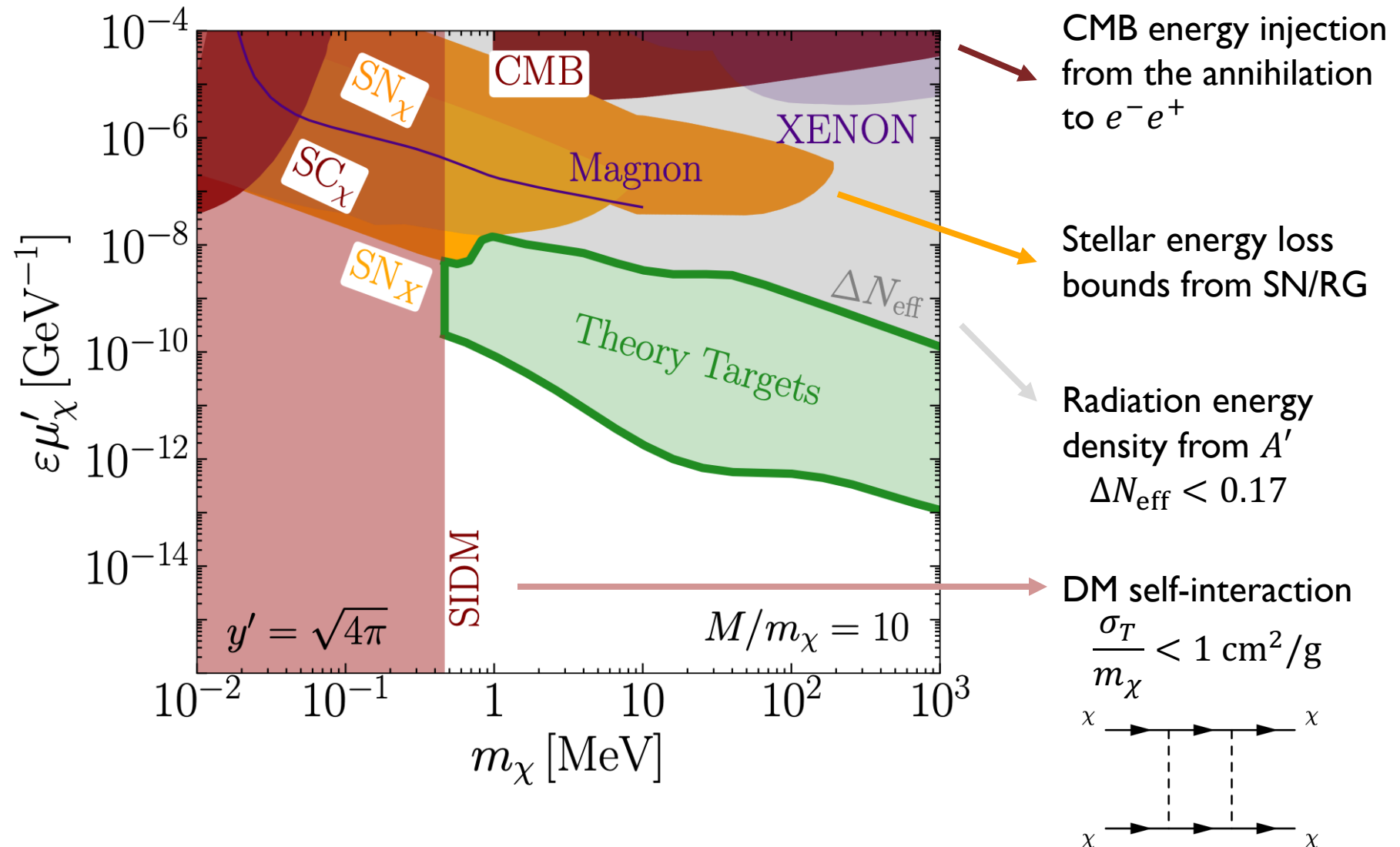


$$Y_X \simeq \begin{cases} \frac{90}{\pi^4 g_{*s}(m_e)} \xi_e^3 \sim 10^{-3} (e^2 \kappa)^{3/2} \left(\frac{M_{\text{Pl}}}{m_e} \right)^{3/4} & \text{(HR)} \\ \left. \frac{H}{s \langle \sigma v \rangle'_{\chi\chi \rightarrow A'A'}} \right|_{T_{XFO}} \sim 10^6 \frac{\sqrt{e^2 \kappa}}{y'^4 g'^4} \left(\frac{M_{\text{Pl}}}{m_e} \right)^{1/4} \left(\frac{m_\chi}{M_{\text{Pl}}} \right) \left(\frac{M}{m_\chi} \right)^6 & \text{(CR)} \end{cases}$$

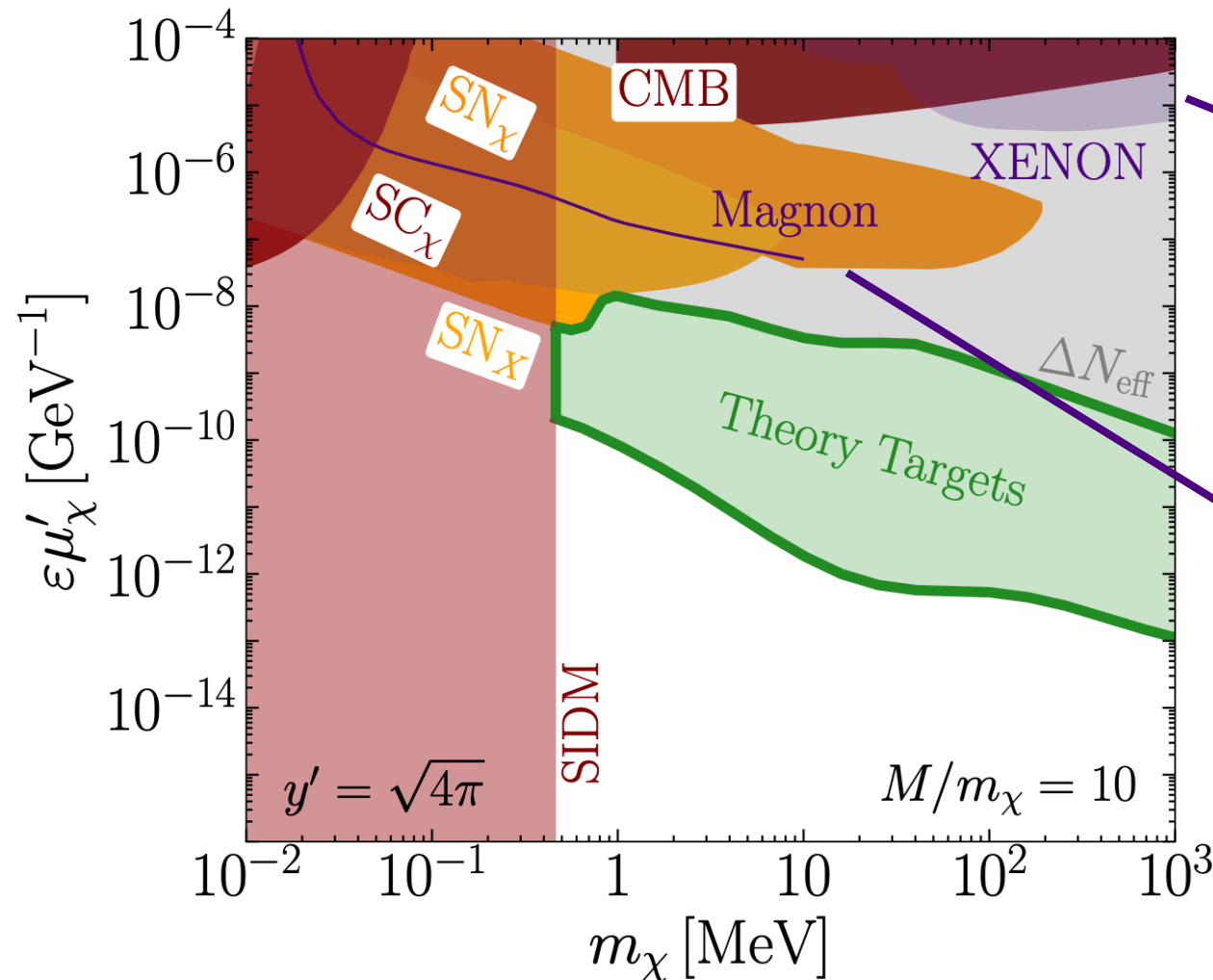
Direct Detection Theory Targets



Various Constraints



Direct Detection Sensitivity

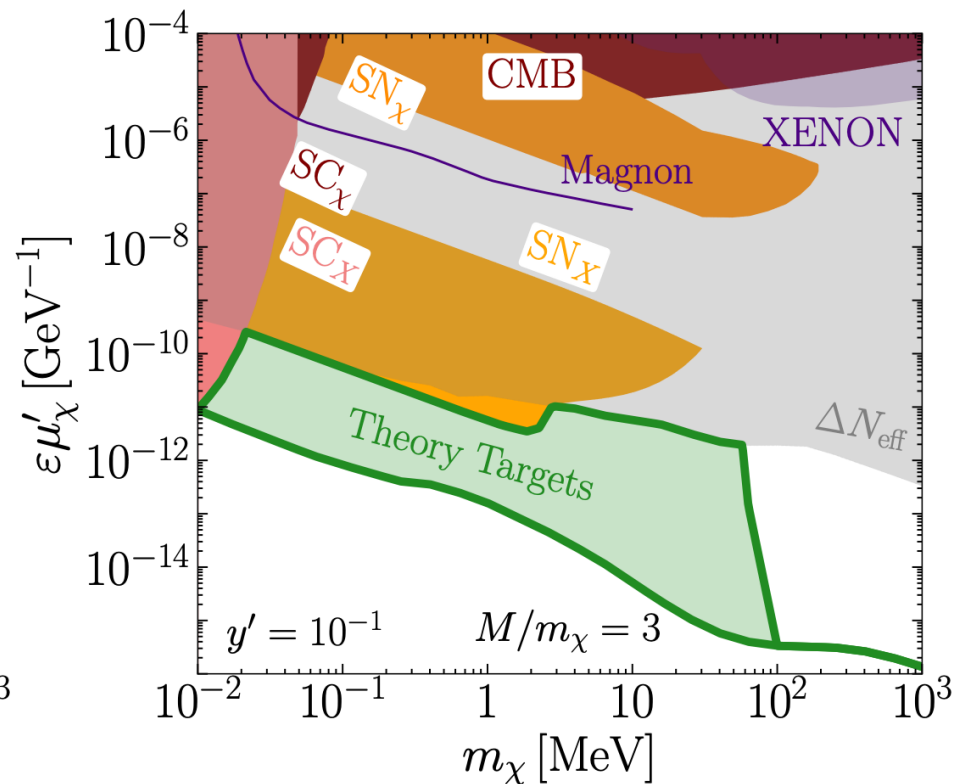
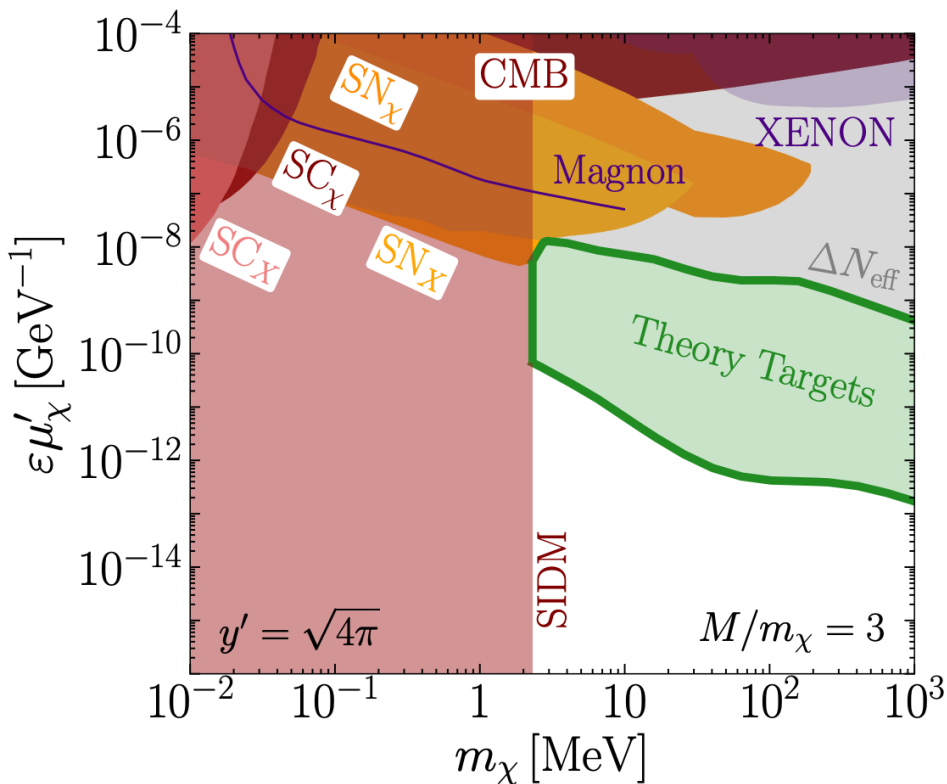


XENON electron-recoil
No coherent enhancement,
but no suppression either

Magnon experiments
Sensitive to SD
Within an order of
magnitude of the theory
target

Reachable ✓

Other parameters



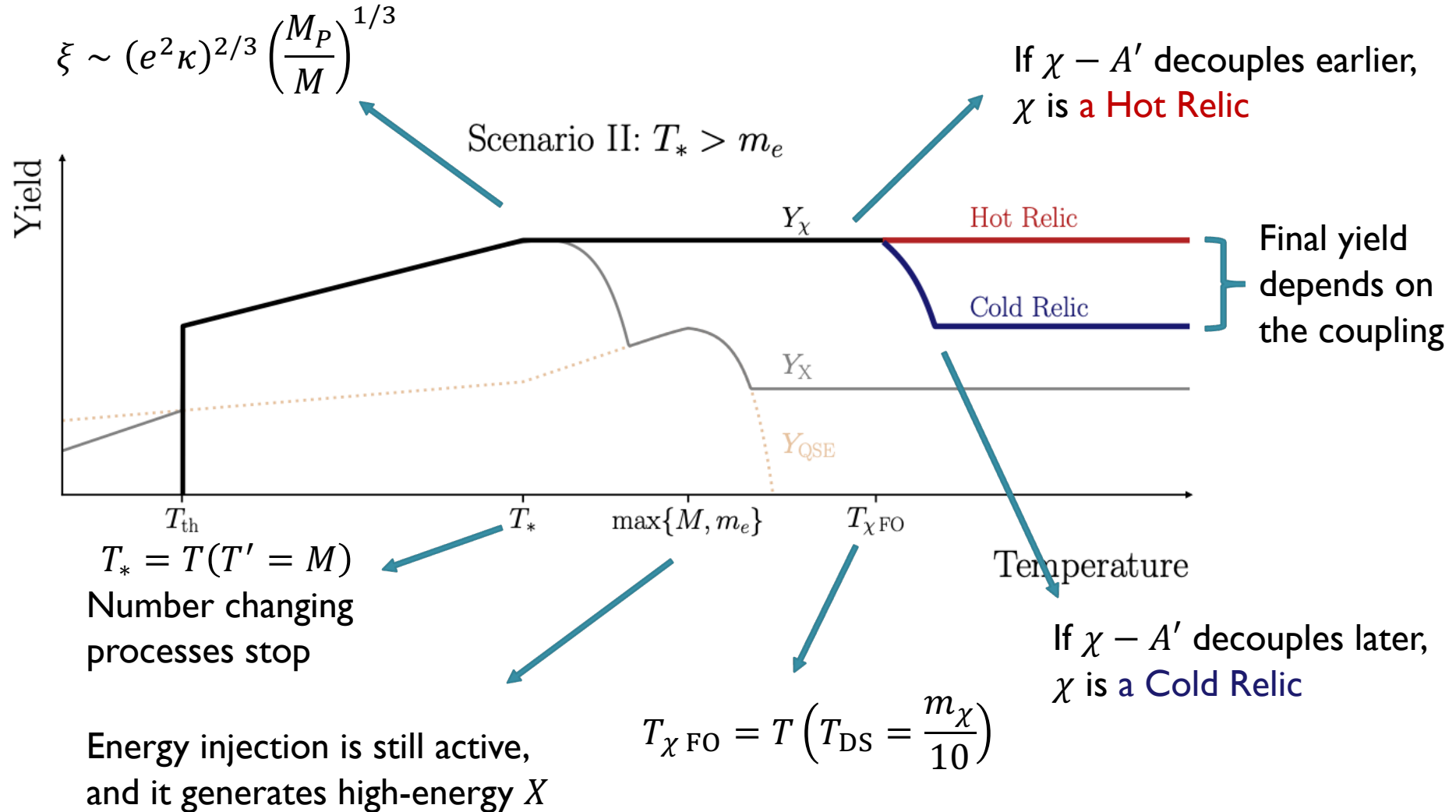
Summary

- Magnon detectors open a new window: coherent sensitivity to spin-dependent DM
- We propose IR freeze-in of dark charged particles, which thermalize the dark sector
- The result is a theory target for magnon experiments:
Predictive ✓
Robust ✓
Reachable ✓

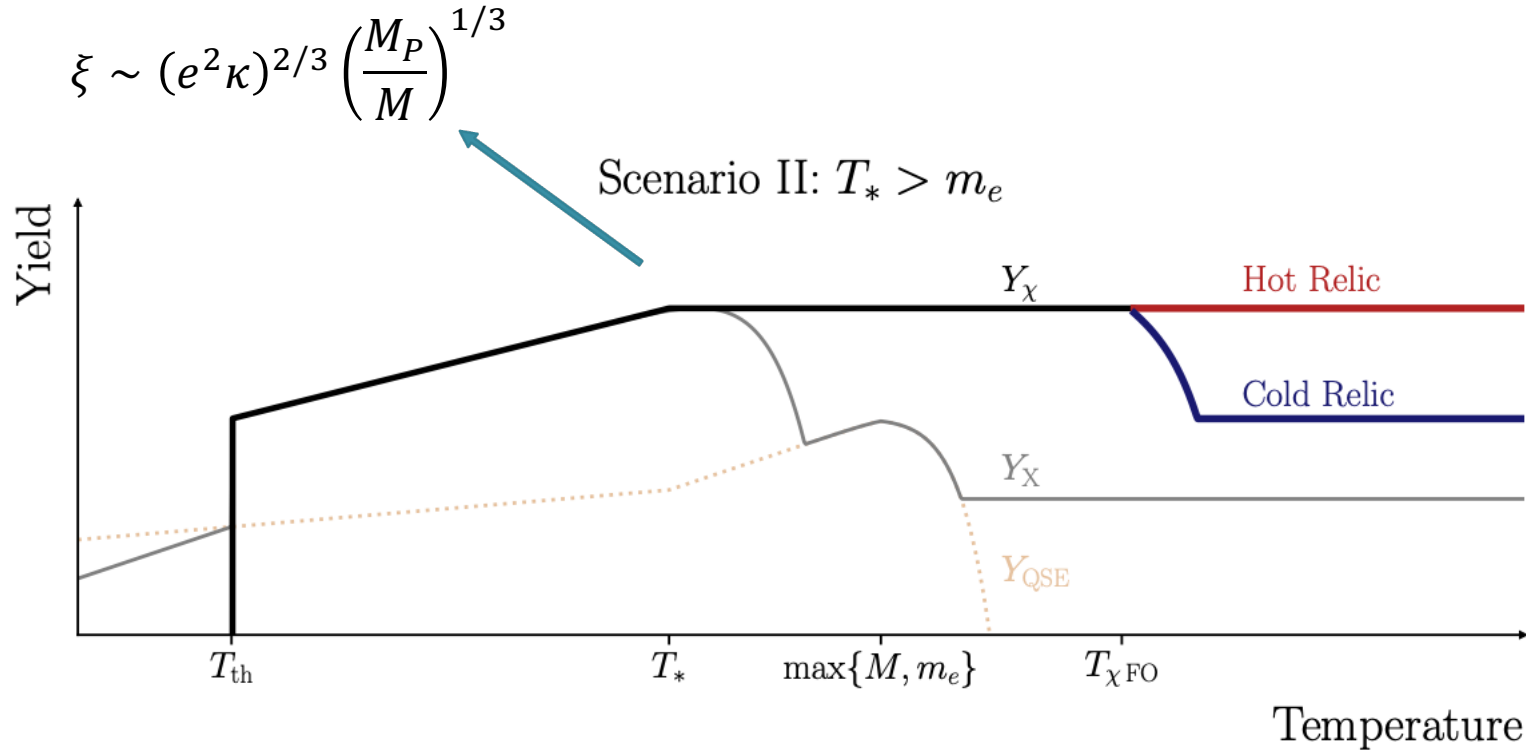
THANK YOU

BACK UP

4. DM Freeze-out: Scenario 2



4. DM Freeze-out: Scenario 2



$$Y_\chi \simeq \begin{cases} \frac{90}{\pi^4 g_{*s}(T_*)} \xi_*^3 \sim 10^{-4} (e^2 \kappa)^2 \left(\frac{M_{Pl}}{M} \right) & \text{(HR)} \\ \left. \frac{H}{s \langle \sigma v \rangle'_{\chi\chi \rightarrow A'A'}} \right|_{\chi FO} \sim \frac{10^5}{y'^4 g'^4} \left(\frac{m_\chi}{M_{Pl}} \right) \left(\frac{M}{m_\chi} \right)^6 \left(\frac{M}{\max\{M, m_e\}} \right) & \text{(CR)} \end{cases}$$