

Diffusion Models and Thermodynamic Singularity

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Contents

- Free energy in data generation
- Thermodynamics of diffusion models and its singularity
- Detecting memorization through thermodynamic singularity

Motivation (Free Energy)

$$P(X = H) = 0.9, \quad P(X = T) = 0.1$$

$$P(\underbrace{\text{HHHHHHHHHH}}_{10}) = 0.9^{10}$$

$$P(\underbrace{\text{HHHHHHHTHH}}_{10}) = 0.9^9 \cdot 0.1$$



Among the total N of samples,

Let $P(X = a) > P(X = b)$

$$X_i \in \{a, b\}$$

Number of a samples:

$$N_a = P(a)N$$

Number of b samples:

$$N_b = N - N_a = P(b)N$$

$$\prod_{i=1}^N P(X_i = a) > \prod_{i=1}^{N_a} P(X_i = a) \cdot \prod_{j=N_a+1}^N P(X_j = b)$$

Why do we always see N_1 number of a s and N_2 number of b s?

- Consider $N = 2$ $X_1, X_2 \in \{a, b\}$

- Four assignments cover all possibilities:

$$X_1 = a, X_2 = a \quad P(X_1, X_2) = P(a)^2 \quad \Longleftarrow P(a)^2$$

$$X_1 = a, X_2 = b \quad P(X_1, X_2) = P(a)P(b)$$

$$X_1 = b, X_2 = a \quad P(X_1, X_2) = P(a)P(b) \quad \Longleftarrow 2P(a)P(b)$$

$$X_1 = b, X_2 = b \quad P(X_1, X_2) = P(b)^2 \quad \Longleftarrow P(b)^2$$

Why do we always see N_1 number of a s and N_2 number of b s?

- With N trials, $X_1, X_2, \dots, X_N \in \{a, b\}$
- Probability of having N_a number of a s and N_b number of b s is

$$P(N_a, N_b) = \binom{N}{N_a, N_b} P(a)^{N_a} P(b)^{N_b}$$
$$\binom{N}{N_a, N_b} = \binom{N}{N_a} = \frac{N!}{N_a! N_b!}$$

Why do we always see N_1 number of a s and N_2 number of b s?

- With N trials, $X_1, X_2, \dots, X_N \in \{a, b\}$
- Probability of having N_a number of a s and N_b number of b s is

$$P(N_a, N_b) = \binom{N}{N_a, N_b} P(a)^{N_a} P(b)^{N_b}$$

$$\binom{N}{N_1, N_2} \gg \binom{N}{N, 0}$$

$$P(a)^{N_1} P(b)^{N_2} < P(a)^N$$

Now let's derive the equation like this!

$$P(N_a, N_b) = \binom{N}{N_a, N_b} P(a)^{N_a} P(b)^{N_b}$$

Stirling's approximation for large n,

$$\log n! \approx n \log n - n$$

$$\begin{aligned} \Rightarrow \log \binom{N}{N_a, N_b} &= \log \frac{N!}{N_a! N_b!} \\ &\approx N \log N - \cancel{N} - \{N_a \log N_a - \cancel{N_a}\} - \{N_b \log N_b - \cancel{N_b}\} \\ &= N \cdot \left\{ -\frac{N_a}{N} \log \frac{N_a}{N} - \frac{N_b}{N} \log \frac{N_b}{N} \right\} \end{aligned}$$

Now let's derive the equation like this!

$$\log \binom{N}{N_a, N_b} \approx N \cdot \left\{ -\frac{N_a}{N} \log \frac{N_a}{N} - \frac{N_b}{N} \log \frac{N_b}{N} \right\}$$

Let $Q_i = \frac{N_i}{N}$

$$\log \binom{N}{N_a, N_b} \approx N \sum_{i \in \{a, b\}} -Q_i \log Q_i$$

$$= N \cdot \text{Ent}[Q_i]$$

Entropy of empirical
distribution

Free Energy

- Let

$$\begin{aligned} P(N_a, N_b) &= \exp \left(- \left\{ - N_a \log P(x = a|\lambda) - N_b \log P(x = b|\lambda) \right. \right. \\ &\quad \left. \left. - N \left(- \frac{N_a}{N} \log \frac{N_a}{N} - \frac{N_b}{N} \log \frac{N_b}{N} \right) \right\} \right) \\ &\equiv \exp \left(- \frac{E - TS}{T} \right) \end{aligned}$$

Free Energy

$$P = \exp \left(-\frac{E - TS}{T} \right)$$

$$E = -\frac{N_a}{N} \log P(x = a|\lambda) - \frac{N_b}{N} \log P(x = b|\lambda)$$

Data fitting


$$S = N \cdot \left(-\frac{N_a}{N} \log \frac{N_a}{N} - \frac{N_b}{N} \log \frac{N_b}{N} \right)$$

$$T = \frac{1}{N}$$

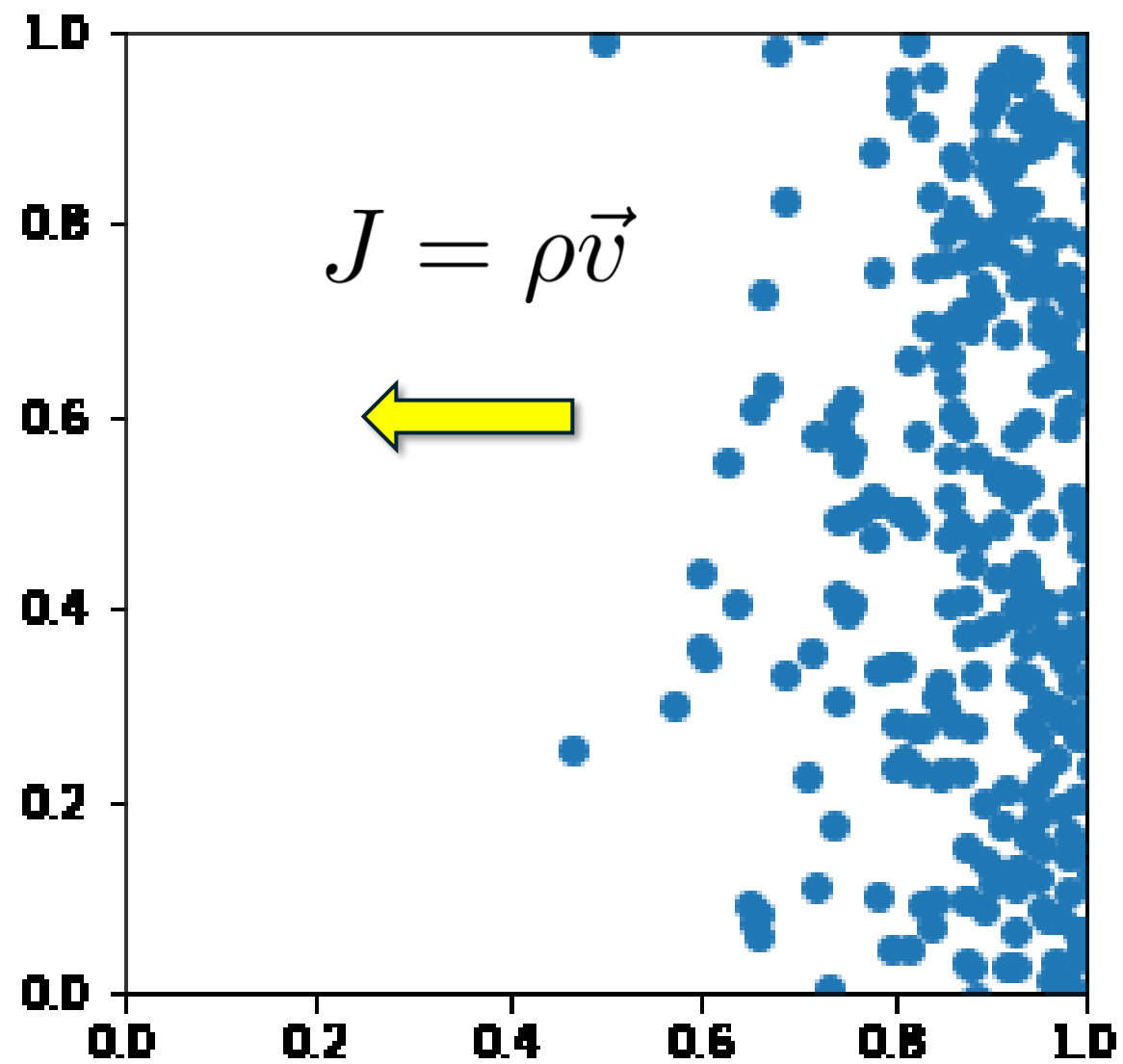
Temperature

$$\frac{1}{T} = \frac{dS}{dE}$$

$$T = \frac{1}{N} \frac{\log P(a) - \log P(b)}{\log Q_a - \log Q_b}$$

At $T = \frac{1}{N}$, the equilibrium occurs at

$$Q_a = P(a), Q_b = P(b)$$

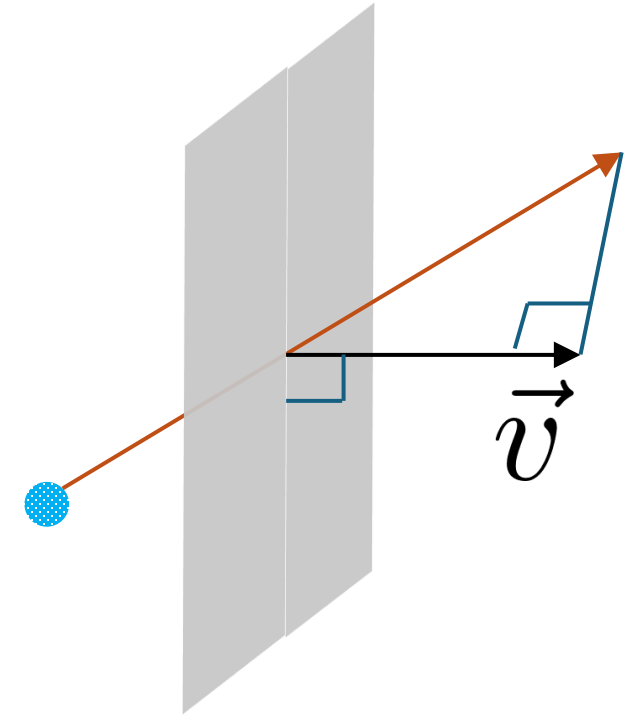


Flux J : Net number of particles passing through the membrane per unit time

$$J \equiv \rho \vec{v}$$

ρ : mass density

$\vec{v}$: net orthogonal velocity



A velocity field $\vec{v}$ associated with J (flux) is of our interest!

Helmholtz Free Energy (Dynamics with Fixed T)

$$F[\rho_t] = \langle U \rangle_{\rho_t} - TS = \int \rho_t(\mathbf{x}) U(\mathbf{x}) d\mathbf{x} + T \int \rho_t(\mathbf{x}) \log \rho_t(\mathbf{x}) d\mathbf{x}$$

$$\frac{dF}{dt} = \int \left. \frac{\partial F}{\partial \rho_t} \right|_{\rho_t} \frac{\partial \rho_t}{\partial t} d\mathbf{x} = - \int \left. \frac{\partial F}{\partial \rho_t} \right|_{\rho_t} \nabla \cdot (\rho_t \vec{v}) d\mathbf{x}$$

$$= \int \rho_t \vec{v}(\mathbf{x})^\top \nabla \left. \frac{\partial F}{\partial \rho_t} \right|_{\rho_t} d\mathbf{x} = \left\langle \vec{v}, \nabla \left. \frac{\partial F}{\partial \rho_t} \right|_{\rho_t} \right\rangle_{\rho_t} \quad \leftarrow \text{Maximally decrease}$$

$$\vec{v} \propto -\nabla \left. \frac{\partial F}{\partial \rho_t} \right|_{\rho_t} \quad \text{with fixed} \quad K = \frac{1}{2} \int \rho_t ||\vec{v}||^2 d\mathbf{x}$$

Flow for Diffusion Models

- Flow that maximally decreases free energy

$$\vec{v}(\mathbf{x}) = -\nabla \frac{\partial F}{\partial \rho_t} = -\nabla U - T \nabla \log \rho_t$$

$$J(\mathbf{x}) = \rho_t(\mathbf{x}) \vec{v}(\mathbf{x}) = -\rho_t(\mathbf{x}) \nabla U - T \nabla \rho_t$$

- Forward process

Fixed potential energy and temperature $U(\mathbf{x}) = \frac{\beta}{4} \mathbf{x}^2$, $T = \frac{\beta}{2}$

$$J_F = -\rho_{\text{eq}} \nabla U - T \nabla \rho_{\text{eq}} = 0 \longrightarrow \rho_{\text{eq}}(\mathbf{x}) \propto \exp(-U(\mathbf{x})/T) = \exp(-\mathbf{x}^2/2)$$

Forward flux

Flow for Diffusion Models

- Flow that maximally decreases free energy

$$\vec{v}(\mathbf{x}) = -\nabla \frac{\partial F}{\partial \rho_t} = -\nabla U - T \nabla \log \rho_t$$

$$J(\mathbf{x}) = \rho_t(\mathbf{x}) \vec{v}(\mathbf{x}) = -\rho_t(\mathbf{x}) \nabla U - T \nabla \rho_t$$

- Reverse process

With learned (time-varying) potential energy,

$$J_R = -J_F$$

Properties of Helmholtz Free Energy Dynamics

[1] Helmholtz free energy does not increase

$$\vec{v}(\mathbf{x}) = -\nabla \frac{\partial F}{\partial \rho_t} = -\nabla U - T \nabla \log \rho_t$$

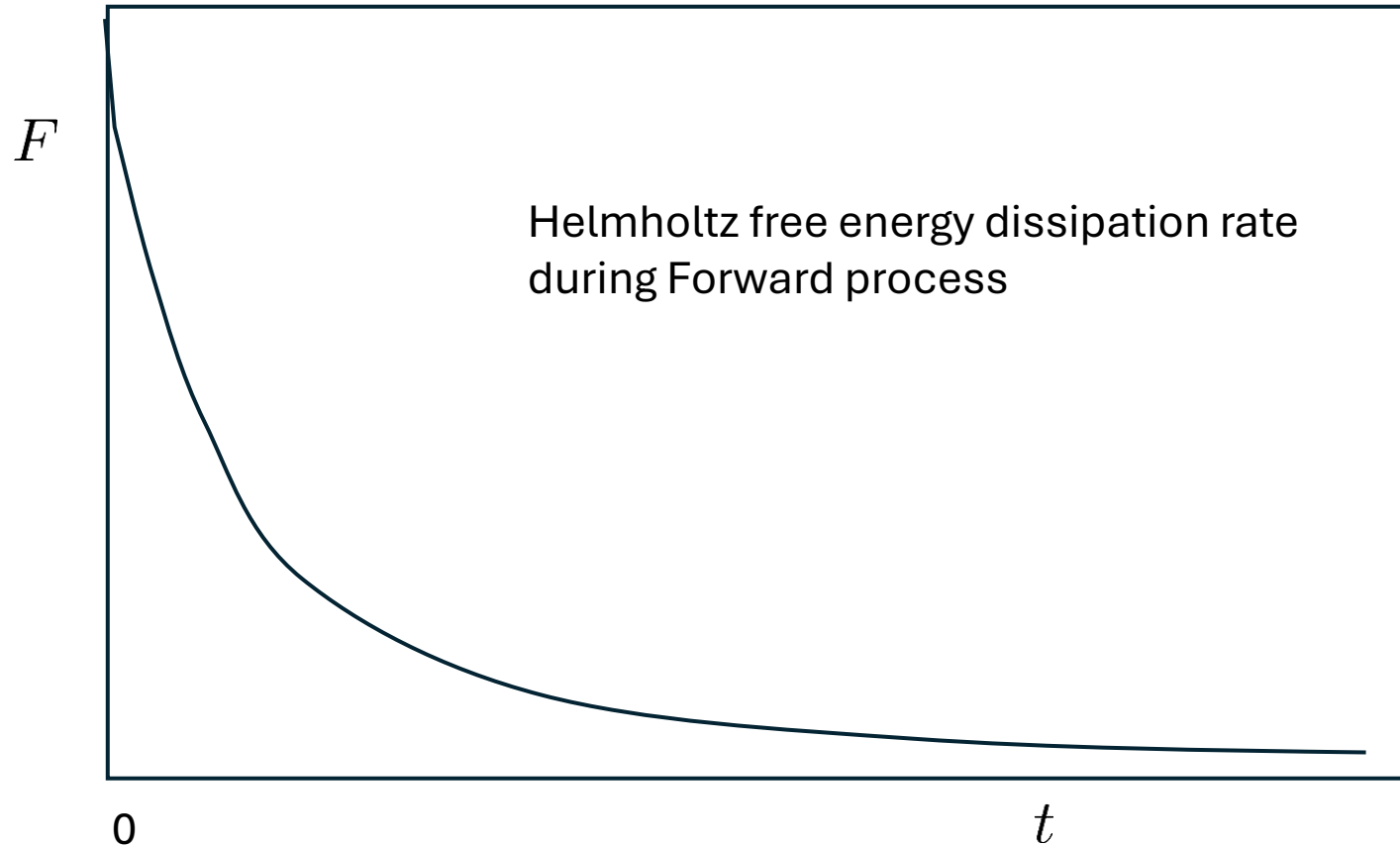
$$\frac{dF}{dt} = -\langle \vec{v}, \vec{v} \rangle_{\rho_t} \leq 0$$

[2] The rate of change decreases over time during the forward process

- The second derivative $\frac{\partial^2 F}{dt^2}$ is positive, as implied by results on the dissipation of Fisher information provided that the potential is “convex”.

Flow to Maximize Free Energy Decrease

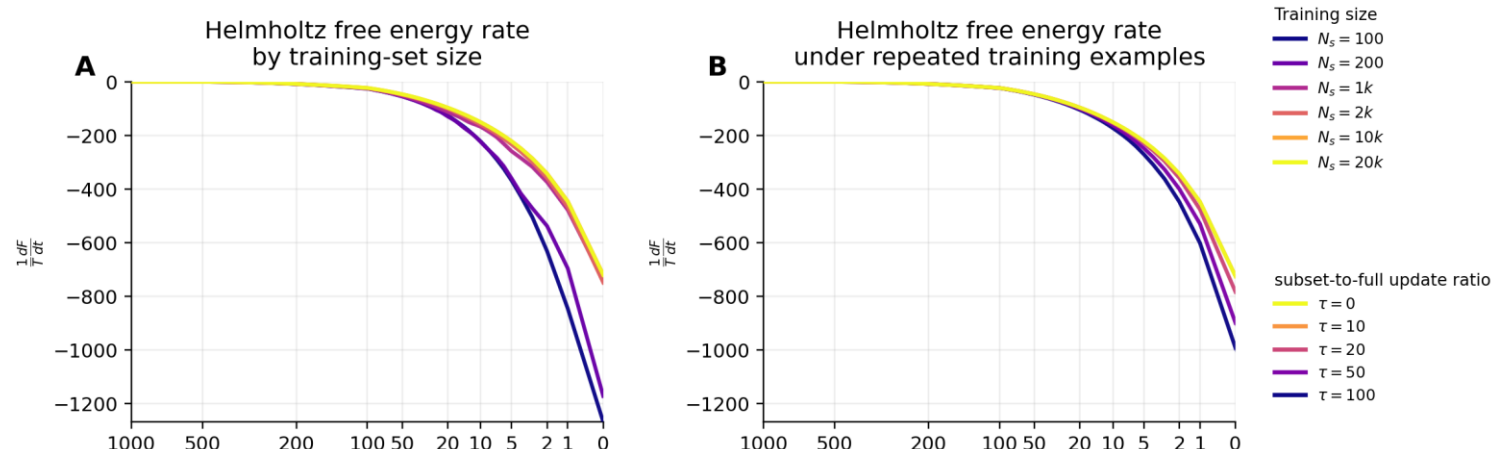
$$W \equiv \frac{1}{T} \left. \widehat{\frac{dF}{dt}} \right|_t = -\frac{\beta}{2} \frac{1}{N} \sum_{i=1}^N \|\hat{\mathbf{v}}_t(\mathbf{x}_{t,i})\|^2$$



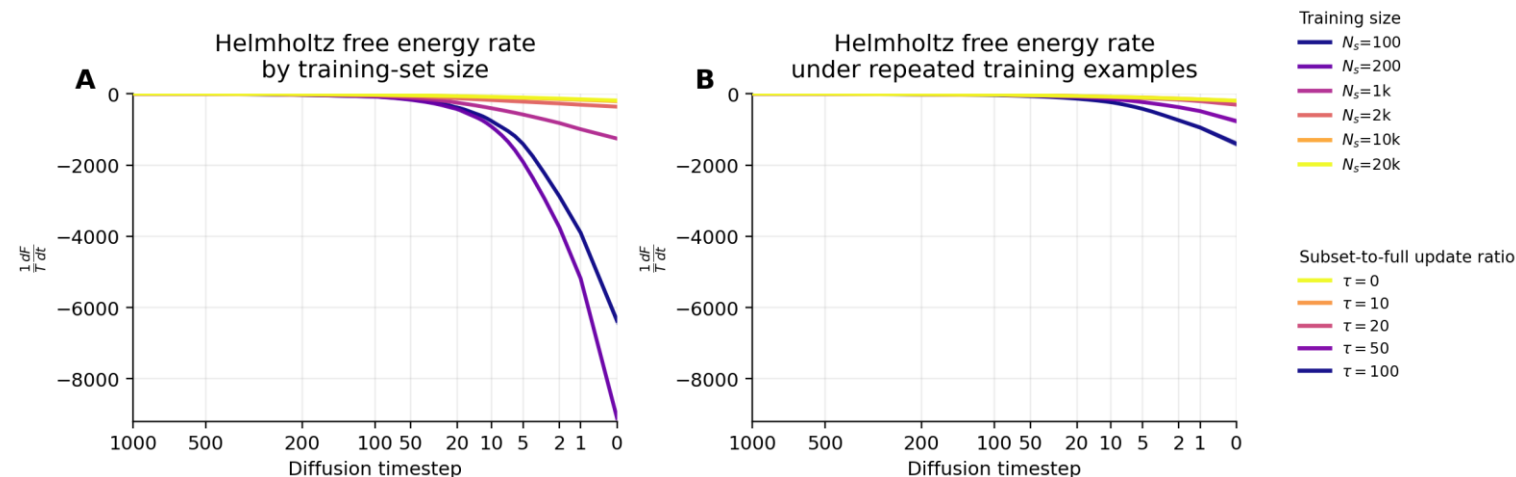
Flow to Maximize Free Energy Decrease

$$W \equiv \frac{1}{T} \frac{d\widehat{F}}{dt} \bigg|_t = -\frac{\beta}{2} \frac{1}{N} \sum_{i=1}^N \|\widehat{\mathbf{v}}_t(\mathbf{x}_{t,i})\|^2$$

- CIFAR-10

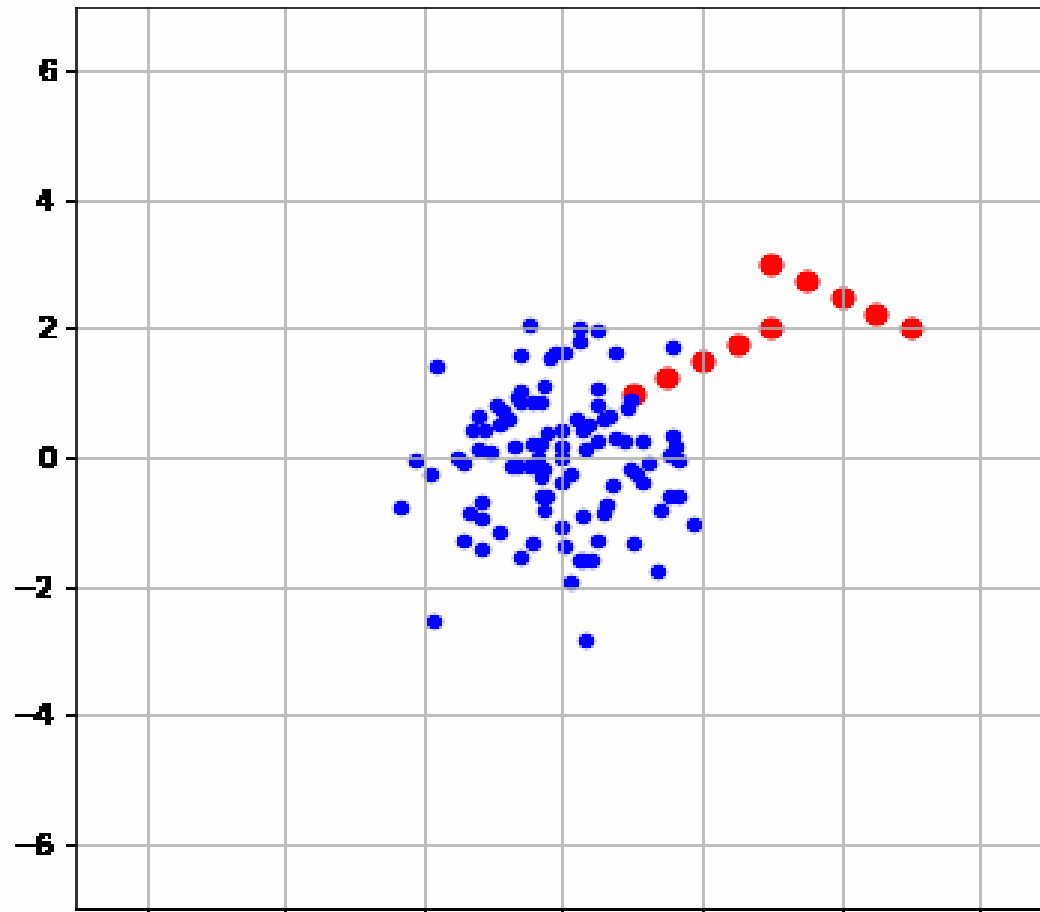


- CelebA-HQ

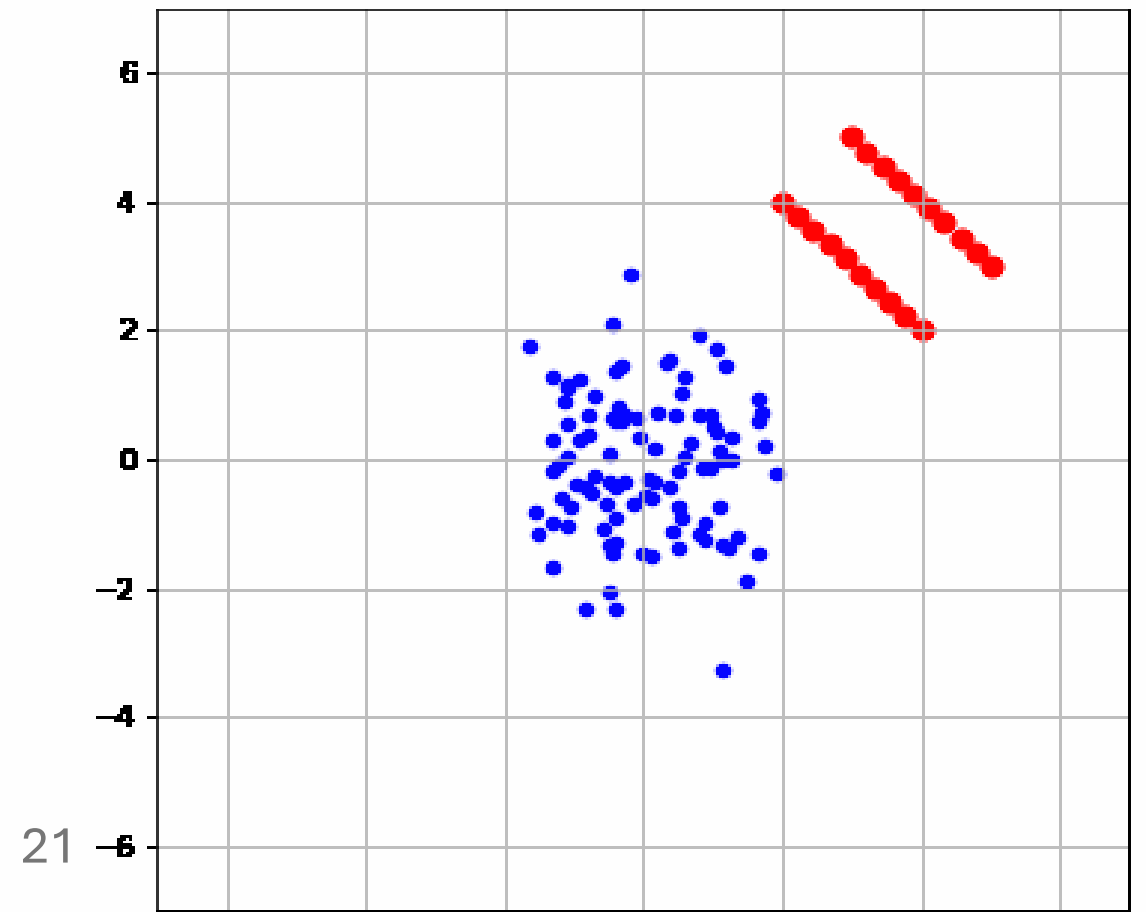




2D Vector Field (Linear beta) at t=500.00



2D Vector Field (Linear beta) at t=500.00



Continuous-time Diffusion of a δ -function

- Diffusion of delta (or mixture of delta) function

$$\rho(\mathbf{x}; t) = k(\mathbf{x}; t) = \mathcal{N}(0, \beta t I)$$

$$\mathbf{x}_{t+\Delta t} = \mathbf{x}_t + \sqrt{\beta \Delta t} \epsilon_t \quad \epsilon_t \sim \mathcal{N}(0, I)$$

Continuous-time Diffusion of a δ -function

- Diffusion of delta function

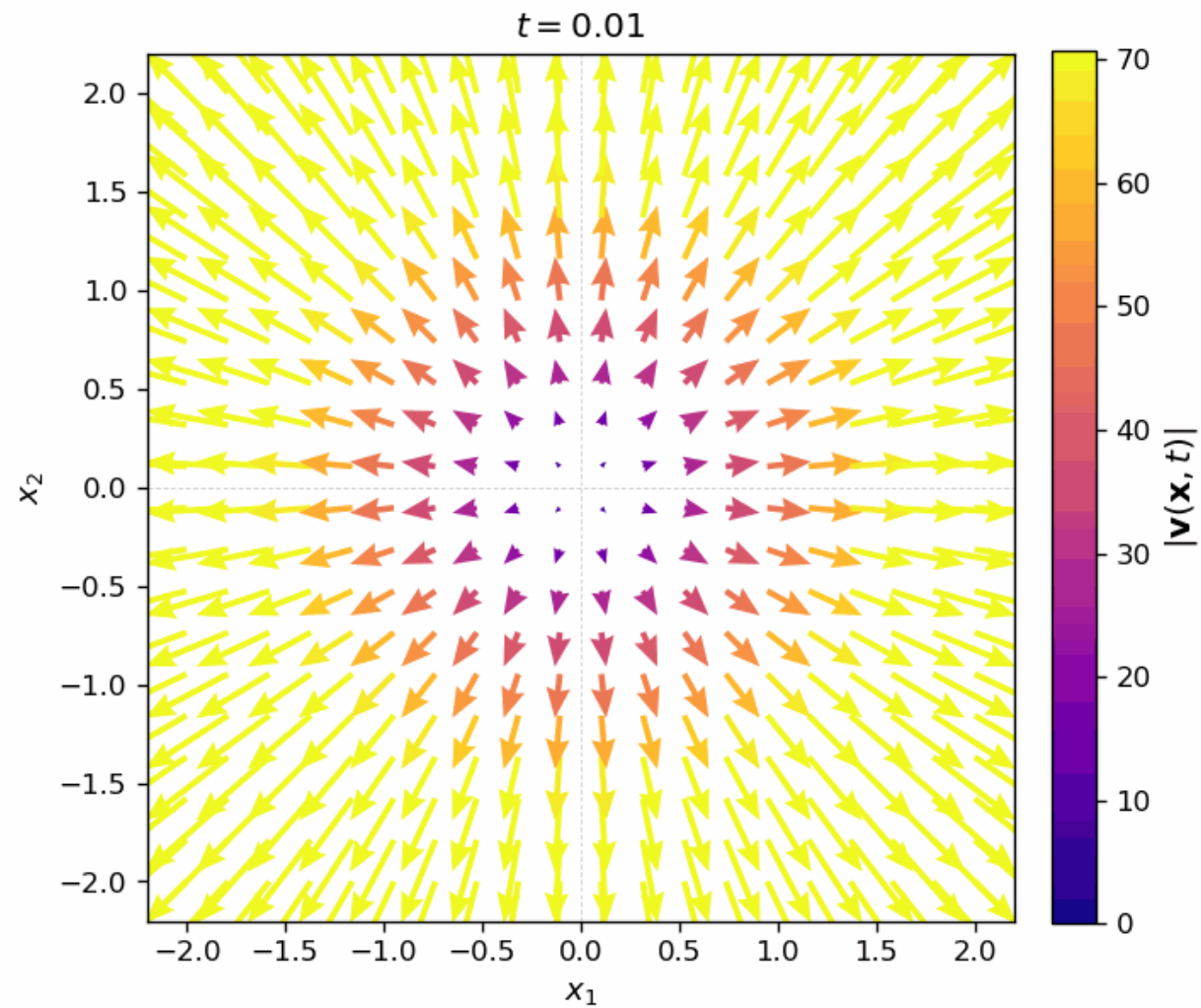
$$\nabla \cdot (\vec{v} k(\mathbf{x}; t)) = -\frac{dk(\mathbf{x}; t)}{dt} \quad \& \quad k(\mathbf{x}; t + \Delta t) = \int k(\mathbf{x}'; t) k(\mathbf{x}' - \mathbf{x}; \Delta t) d\mathbf{x}'$$

$$k(\mathbf{x}; t) \nabla \cdot \vec{v} + \vec{v} \cdot \nabla k(\mathbf{x}; t) = -\frac{dk(\mathbf{x}; t)}{dt}$$

$$k(\mathbf{x}; t) \left[\nabla \cdot \vec{v} - \frac{\mathbf{x}}{\beta t} \cdot \vec{v} \right] = -\frac{1}{2t} \left(\frac{\mathbf{x}^2}{\beta t} - d \right) k(\mathbf{x}; t) \quad \Rightarrow \quad \vec{v}_t = \frac{1}{2t} \mathbf{x}$$

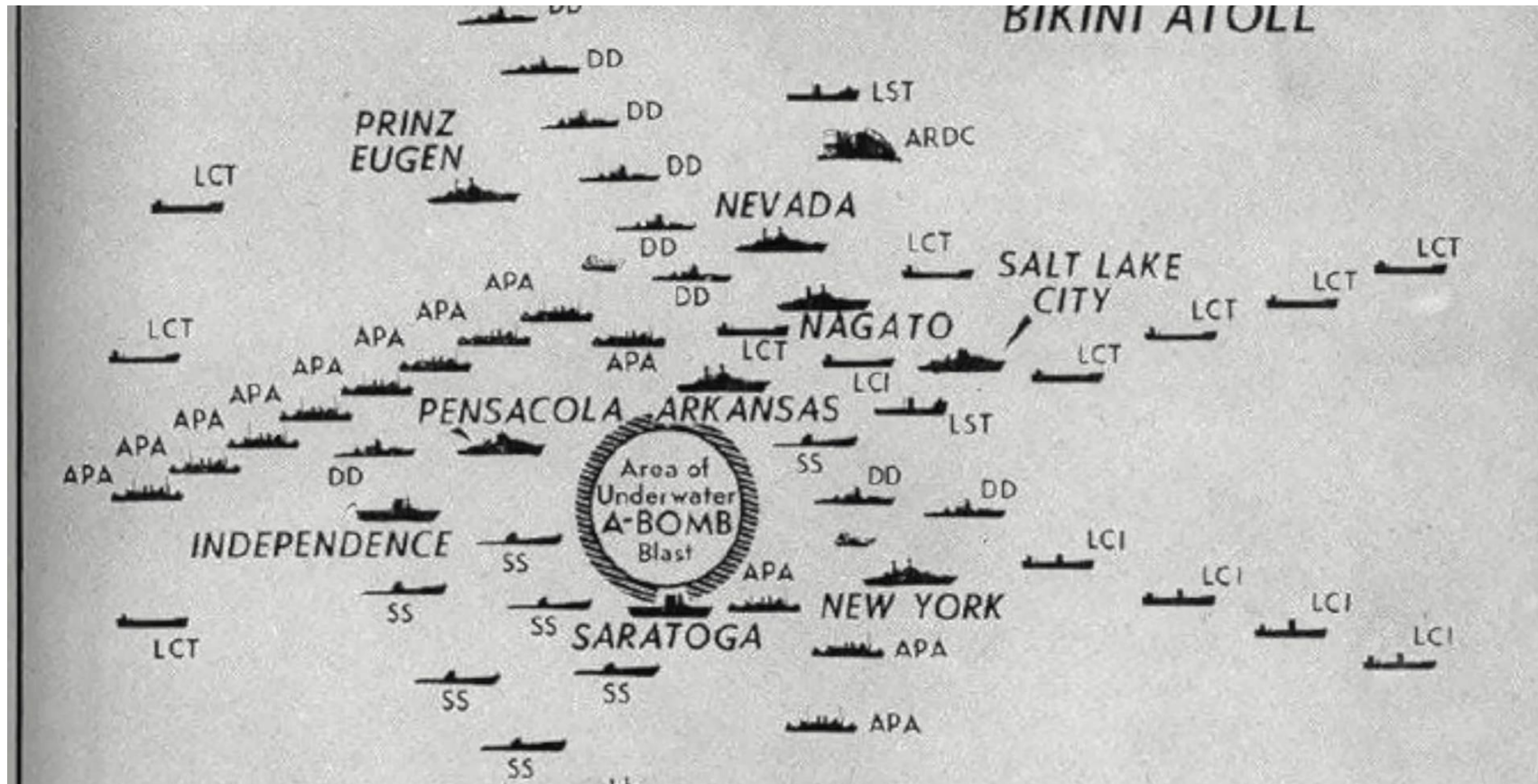
$$\nabla \cdot \vec{v} - \frac{\mathbf{x}}{\beta t} \cdot \vec{v} = -\frac{1}{2} \left(\frac{\mathbf{x}^2}{\beta t} - d \right)$$

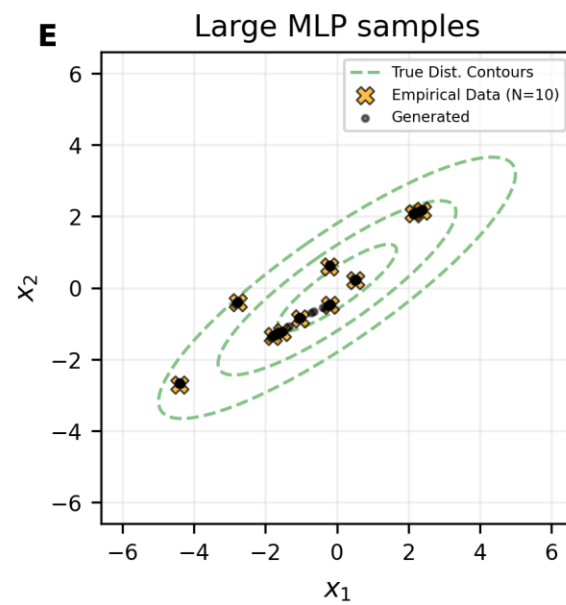
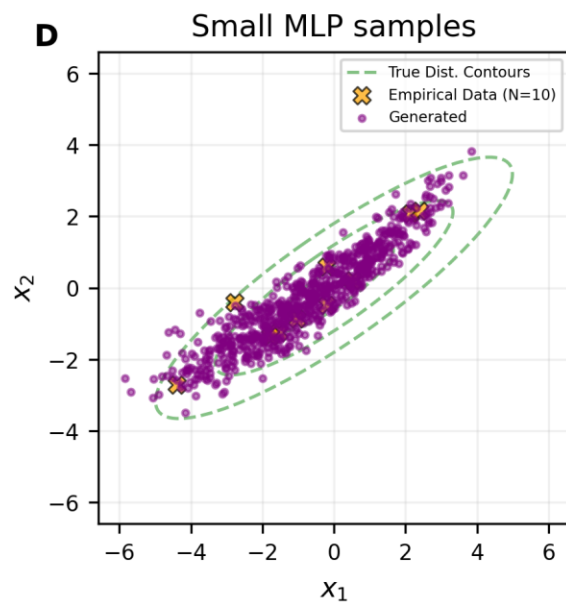
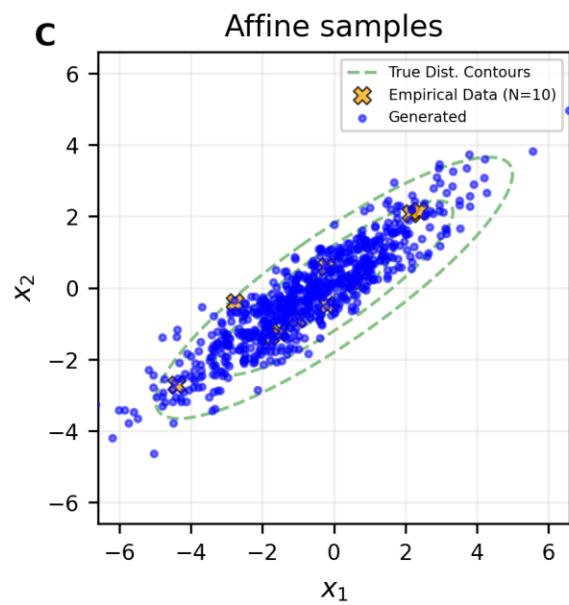
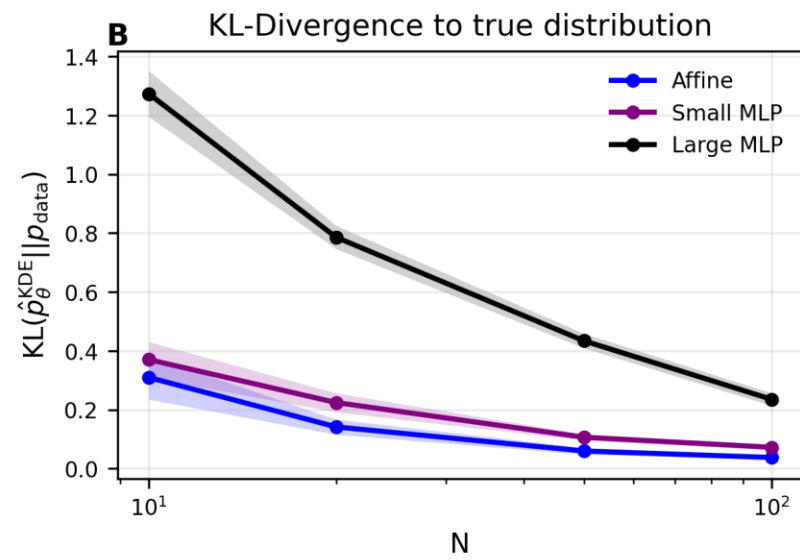
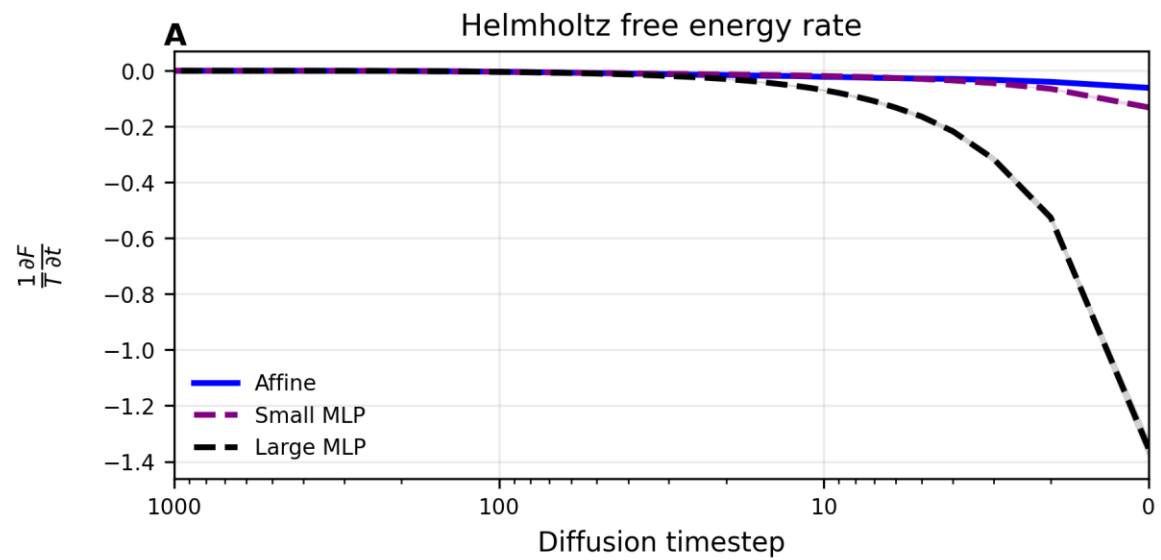
Velocity is β -independent



$$\vec{v}_t = \frac{1}{2t} \mathbf{x}$$

<https://www.youtube.com/watch?v=AL-WRlEr54M>





$N_s=100$

Generated



Nearest
Neighbor

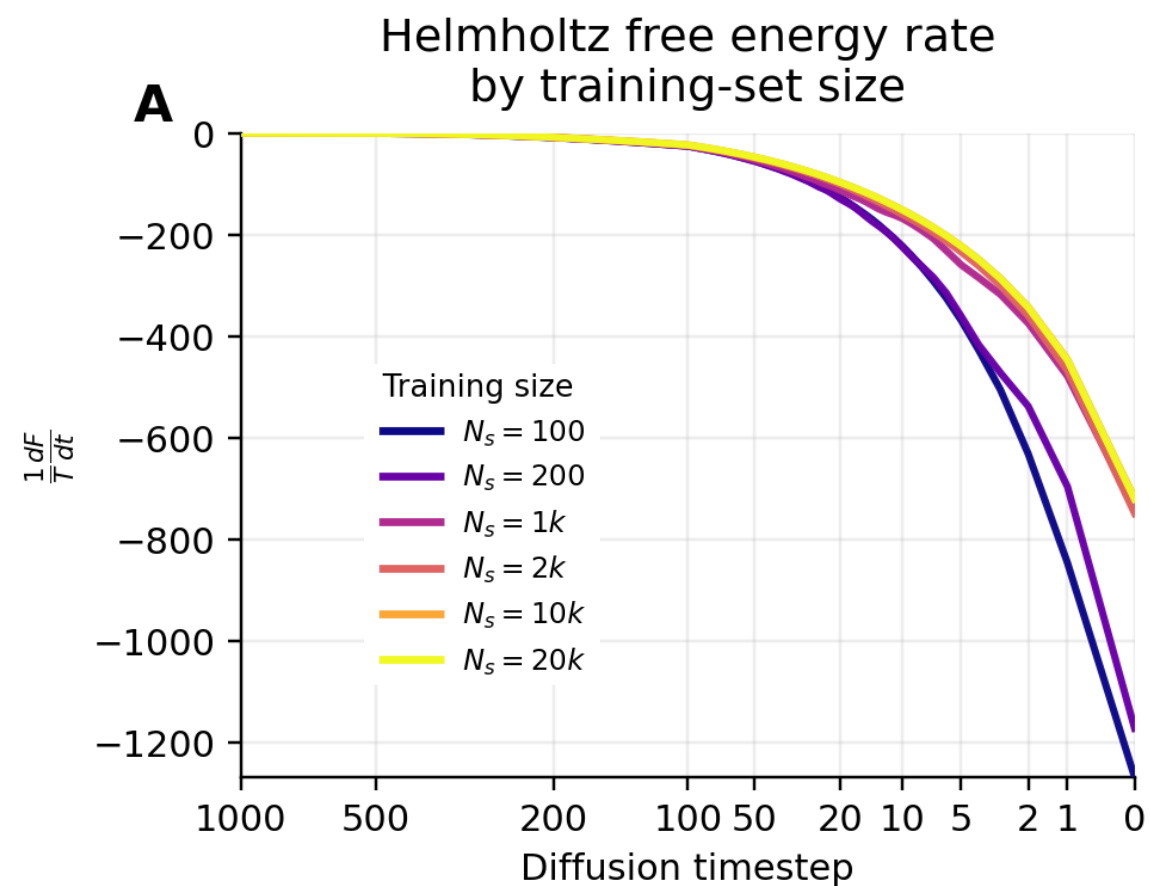


$N_s=20,000$

Generated



Nearest
Neighbor



$\tau=100$

Generated



Nearest
Neighbor

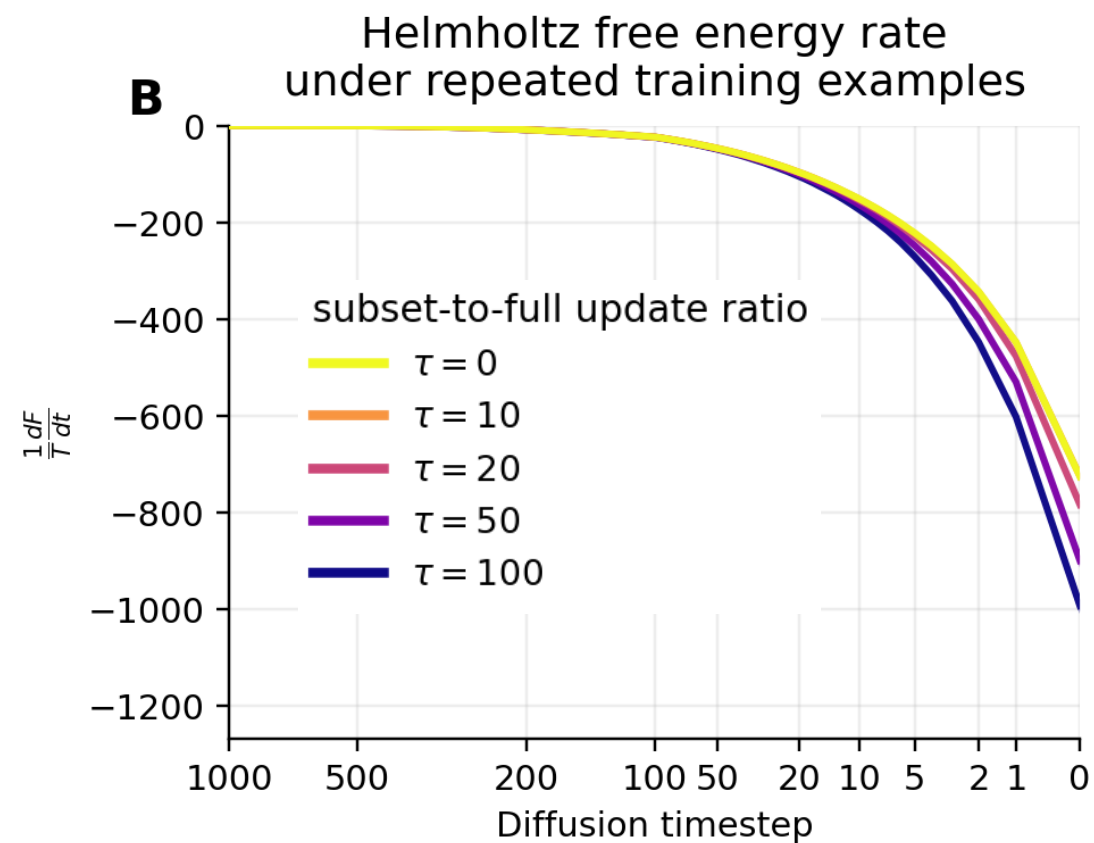


$\tau=10$

Generated



Nearest
Neighbor



$N_s=100$

Generated



Nearest
Neighbor

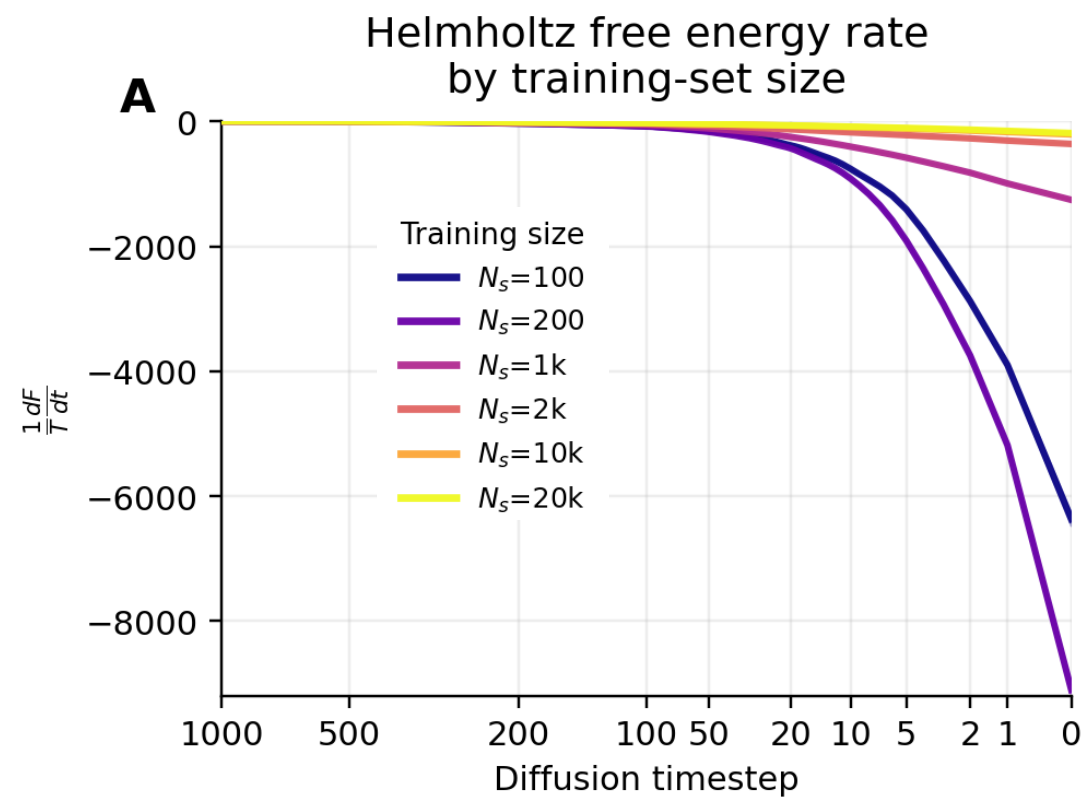


$N_s=20,000$

Generated



Nearest
Neighbor



$\tau=100$

Generated



Nearest
Neighbor

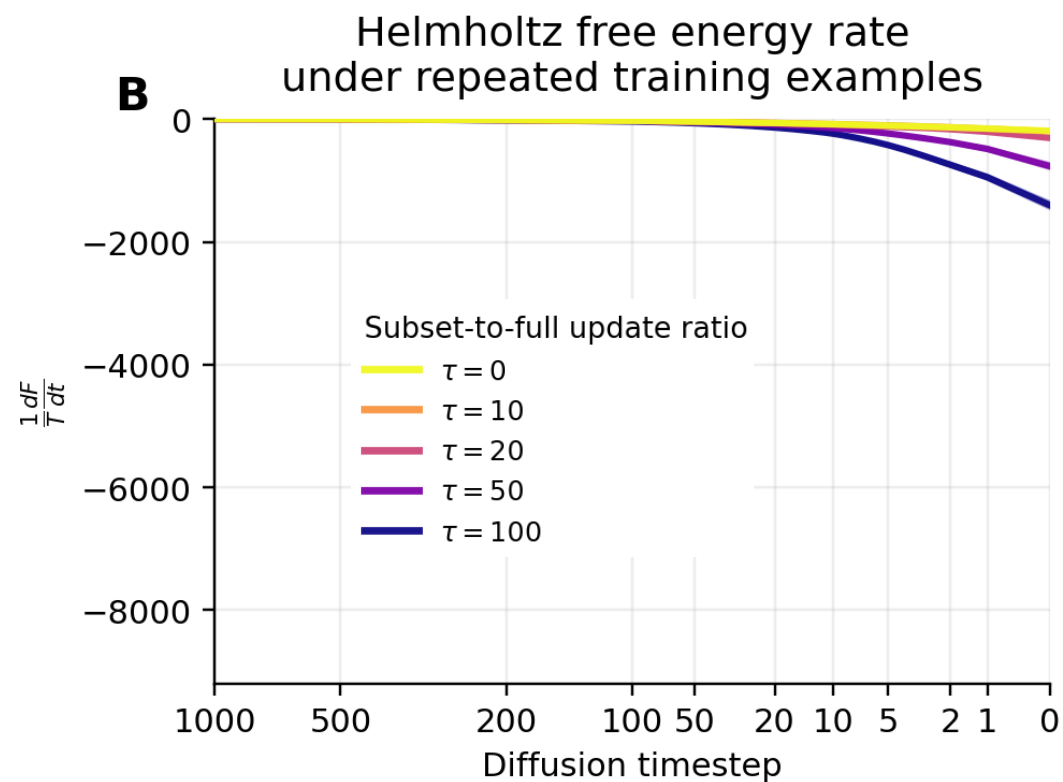


$\tau=10$

Generated



Nearest
Neighbor



Sample-level Memorization Detection

- Stable diffusion v1.4

Method	Time step	AUC	TPR@1%FPR
$\ \mathbf{s}_\theta^\Delta(\mathbf{x}_t)\ $	1,000	0.937 $\pm$ 0.012	0.771 $\pm$ 0.027
$= \ \mathbf{s}_{\theta,\text{cond}} - \mathbf{s}_{\theta,\text{uncond}}\ $ [19]	1	0.972 $\pm$ 0.005	0.909 $\pm$ 0.020
$\ H_\theta^\Delta(\mathbf{x}_T)\mathbf{s}_\theta^\Delta(\mathbf{x}_T)\ ^2$ [7]	1,000	0.962 $\pm$ 0.008	0.856 $\pm$ 0.024
	1	0.970 $\pm$ 0.005	0.874 $\pm$ 0.021
$\ \hat{\mathbf{v}}_{\text{cfg}}\ ^2 - \ \hat{\mathbf{v}}_{\text{uncond}}\ ^2$ (Ours)	1	0.982 $\pm$ 0.003	0.905 $\pm$ 0.019

CelebA-HQ Fine Tunning with Velocity Suppression

Condition	LPIPS ≤ 0.1		LPIPS ≤ 0.3		FID	
Pretrained model	0.00%		0.08%		14.93	
No reference scale	28.53%		33.11%		24.33	
	Full velocity $\hat{\mathbf{v}}_t = x_t + \mathbf{s}_\theta(x_t, t) \mid \lambda = 10^{-3}$			Score-only $\hat{\mathbf{v}}_t = \mathbf{s}_\theta(x_t, t) \mid \lambda = 10^{-5}$		
Condition	LPIPS ≤ 0.1	LPIPS ≤ 0.3	FID	LPIPS ≤ 0.1	LPIPS ≤ 0.3	FID
$\gamma = 1.0$	31.85%	36.79%	26.19	23.14%	31.59%	22.71
$\gamma = 0.8$	11.39%	28.14%	17.99	2.47%	18.01%	53.80
$\gamma = 0.7$	3.50%	20.65%	15.88	1.46%	15.75%	25.74
$\gamma = 0.5$	0.05%	15.31%	14.38	0.29%	2.31%	156.08
$\gamma = 0.3$	0.00%	3.80%	15.87	0.02%	0.23%	295.58

Summary

- Data generation and the shift of distributions are naturally associated with Helmholtz free energy dissipation under diffusion dynamics (using the equation of continuity).
- Thermodynamic singularities in the diffusion of delta functions are characteristic of memorization, and its absence indicates non-memorization.

Thank you!

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