

Confronting neutrino theories with correlations of oscillation data

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1st Annual MAPS Meeting

Bringing theory and experiments together for precision neutrino physics

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**Swiss National
Science Foundation**

THEORY $\Rightarrow$ **COMBINE** $\Leftarrow$ **EXPERIMENT**

basic concept

$$f_1(x_i) = g(x_k, \dots).$$

$\downarrow$

$$\chi^2(x_1, \dots, x_i, \dots, x_j, \dots, x_n)$$

$\uparrow$

$$f_2(x_j) = (x_m, \dots).$$

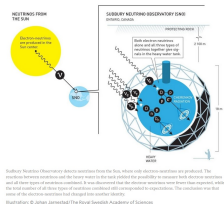
for neutrinos: eg. **mixing scheme relations** + "full range" χ^2 correlated data

a minimal dataset: $\chi^2(\theta_{13}, \theta_{12}, \theta_{23}, \delta_{\text{CP}})$

or better, a dataset with masses:

$$\begin{array}{c}
 \sin^2 \theta_{12} = \dots \\
 \downarrow \\
 \chi^2(\theta_{13}, \theta_{12}, \theta_{23}, \delta_{\text{CP}}, \Delta m_{21}^2, \Delta m_{31}^2) \\
 \uparrow \\
 \cos \delta_{\text{CP}} = \dots
 \end{array}$$

3ν oscillations and masses, $m_0 = m_{\text{lightest}}$



Sudbury Neutrino Observatory detects neutrinos from the Sun, where only electron-neutrinos are produced. The oscillations between electron-neutrinos and the two other types in the sun helped for establishing the neutrino mass and all three types of neutrino oscillations. It was discovered that the electron-neutrinos were fewer than expected, while the total number of all three types of neutrinos combined still corresponded to expectations. The conclusion was that some of the electron-neutrinos had changed into another identity.

Illustration: © Japen, Janssen, The Royal Swedish Academy of Sciences

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 arXiv:1806.11051v1 [hep-ph] 18 Jun 2018. Published for SISSA by SPRINGER.
 www.sciopen.com/journal/100959180611051

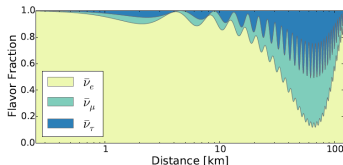


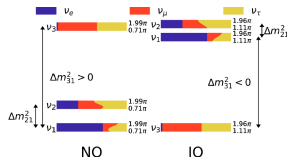
FIG. 1. Illustration of neutrino oscillations. The expected flavor composition of the reactor neutrino flux, for neutrinos of 4 MeV energy used as an example, is plotted as a function of distance to the reactor cores. [arXiv:1503.01059](https://arxiv.org/abs/1503.01059)

Normal mass ordering (NO)

$$\begin{aligned} m_1 &= m_0, \\ m_2 &= \sqrt{m_0^2 + \Delta m_{21}^2}, \\ m_3 &= \sqrt{m_0^2 + \Delta m_{31}^2}, \end{aligned}$$

Inverted mass ordering (IO)

$$\begin{aligned} m_1 &= \sqrt{m_0^2 - \Delta m_{21}^2 - \Delta m_{32}^2}, \\ m_2 &= \sqrt{m_0^2 - \Delta m_{32}^2}, \\ m_3 &= m_0, \end{aligned}$$



Taken from <https://globalfit.astroparticles.es>,
 updated [arXiv:1806.11051](https://arxiv.org/abs/1806.11051), Figure 1

$$\Delta m_{21}^2 > 0$$

$$|\Delta m_{31(32)}^2| > 0$$

NO or IO ?

$$\Sigma m_i < 0.12 \text{ eV at 95\% CL from Planck (CMB)}$$

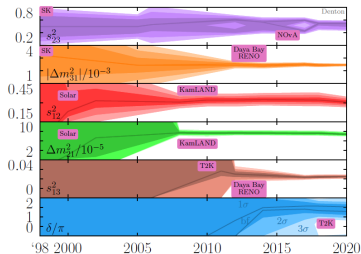
Neutrino flavor and mass eigenstates are related by

$$|\nu_\alpha\rangle = U_{\alpha i} |\nu_i\rangle$$

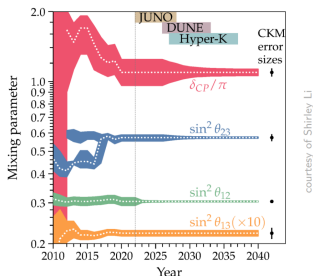
Pontecorvo–Maki–Nakagawa–Sakata parametrization of mixing matrix

$$U_{\text{PMNS}} = \begin{bmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{13}s_{23}e^{i\delta} & c_{12}c_{23} - s_{12}s_{13}s_{23}e^{i\delta} & c_{13}s_{23} \\ s_{12}s_{23} - c_{12}s_{13}c_{23}e^{i\delta} & -c_{12}s_{23} - s_{12}s_{13}c_{23}e^{i\delta} & c_{13}c_{23} \end{bmatrix} \begin{bmatrix} e^{i\alpha_1} & 0 & 0 \\ 0 & e^{i\alpha_2} & 0 \\ 0 & 0 & 1 \end{bmatrix},$$

where $c_{ij} \equiv \cos \theta_{ij}$, $s_{ij} \equiv \sin \theta_{ij}$, $\delta \equiv \delta_{\text{CP}}$.



Peter B. Denton talk, [link here](#).



courtesy of Shirley Li

[arXiv:2204.08668](#), Figure 2.1

Dirac or Majorana ? $(3 + n)\nu$, $n=0$?

NO or IO ? $m_{lightest} \neq 0$?

flavor symmetry ? mixing scheme ?

θ_{23} octant ?

$\delta_{CP} \neq 0$?

$\Sigma m_i < 0.12$ eV $\leftarrow$ this we know from Planck at 95% CL.

$\theta_{13} \neq 0$ $\leftarrow$ this we know from reactor experiments, e.g. Daya Bay.

Experimental data indicate that $\delta_{CP} \neq 0$ and NO are favored.

Closer look at NuFIT data

NuFIT 5.3 (2024)						NuFIT 6.1 (2025)					
without SK atmospheric data						IC23 without SK atmospheric data					
		Normal Ordering (best fit)		Inverted Ordering ($\Delta\chi^2 = 2.3$)				Normal Ordering (best fit)		Inverted Ordering ($\Delta\chi^2 = 1.5$)	
		bfp $\pm 1\sigma$	3 σ range	bfp $\pm 1\sigma$	3 σ range			bfp $\pm 1\sigma$	3 σ range	bfp $\pm 1\sigma$	3 σ range
$\sin^2 \theta_{12}$		$0.307^{+0.012}_{-0.011}$	0.275 to 0.344	$0.307^{+0.012}_{-0.011}$	0.275 to 0.344	$\sin^2 \theta_{12}$		$0.3088^{+0.0067}_{-0.0066}$	0.2893 to 0.3295	$0.3088^{+0.0067}_{-0.0066}$	0.2893 to 0.3295
$\theta_{12}/^\circ$		$33.66^{+0.73}_{-0.70}$	31.60 to 35.94	$33.67^{+0.73}_{-0.71}$	31.61 to 35.94	$\theta_{12}/^\circ$		$33.76^{+0.42}_{-0.41}$	32.54 to 35.03	$33.76^{+0.42}_{-0.41}$	32.54 to 35.03
$\sin^2 \theta_{23}$		$0.572^{+0.018}_{-0.023}$	0.407 to 0.620	$0.578^{+0.016}_{-0.021}$	0.412 to 0.623	$\sin^2 \theta_{23}$		$0.470^{+0.017}_{-0.014}$	0.432 to 0.587	$0.555^{+0.013}_{-0.016}$	0.437 to 0.590
$\theta_{23}/^\circ$		$49.1^{+1.0}_{-1.3}$	39.6 to 51.9	$49.5^{+0.9}_{-1.2}$	39.9 to 52.1	$\theta_{23}/^\circ$		$43.27^{+1.0}_{-0.82}$	41.11 to 50.02	$48.15^{+0.75}_{-0.92}$	41.40 to 50.21
$\sin^2 \theta_{13}$		$0.02203^{+0.00056}_{-0.00058}$	0.02029 to 0.02391	$0.02219^{+0.00059}_{-0.00057}$	0.02047 to 0.02396	$\sin^2 \theta_{13}$		$0.02249^{+0.00057}_{-0.00057}$	0.02070 to 0.02420	$0.02261^{+0.00056}_{-0.00056}$	0.02091 to 0.02433
$\theta_{13}/^\circ$		$8.54^{+0.11}_{-0.11}$	8.19 to 8.89	$8.57^{+0.11}_{-0.11}$	8.23 to 8.90	$\theta_{13}/^\circ$		$8.62^{+0.11}_{-0.11}$	8.27 to 8.95	$8.65^{+0.11}_{-0.11}$	8.31 to 8.97
$\delta_{CP}/^\circ$		197^{+41}_{-25}	108 to 404	286^{+27}_{-32}	192 to 360	$\delta_{CP}/^\circ$		207^{+23}_{-20}	114 to 405	283^{+24}_{-28}	202 to 347
$\frac{\Delta m_{21}^2}{10^{-5} \text{ eV}^2}$		$7.41^{+0.21}_{-0.20}$	6.81 to 8.03	$7.41^{+0.21}_{-0.20}$	6.81 to 8.03	$\frac{\Delta m_{21}^2}{10^{-5} \text{ eV}^2}$		$7.537^{+0.094}_{-0.10}$	7.236 to 7.823	$7.537^{+0.094}_{-0.10}$	7.236 to 7.822
$\frac{\Delta m_{3\ell}^2}{10^{-3} \text{ eV}^2}$		$+2.511^{+0.027}_{-0.027}$	$+2.428$ to $+2.597$	$-2.498^{+0.032}_{-0.024}$	-2.581 to -2.409	$\frac{\Delta m_{3\ell}^2}{10^{-3} \text{ eV}^2}$		$+2.521^{+0.026}_{-0.018}$	$+2.454$ to $+2.592$	$-2.500^{+0.024}_{-0.023}$	-2.569 to -2.430
with SK atmospheric data						IC24 with SK atmospheric data					
		Normal Ordering (best fit)		Inverted Ordering ($\Delta\chi^2 = 9.1$)				Normal Ordering (best fit)		Inverted Ordering ($\Delta\chi^2 = 5.9$)	
		bfp $\pm 1\sigma$	3 σ range	bfp $\pm 1\sigma$	3 σ range			bfp $\pm 1\sigma$	3 σ range	bfp $\pm 1\sigma$	3 σ range
$\sin^2 \theta_{12}$		$0.307^{+0.012}_{-0.011}$	0.275 to 0.344	$0.307^{+0.012}_{-0.011}$	0.275 to 0.344	$\sin^2 \theta_{12}$		$0.3088^{+0.0067}_{-0.0066}$	0.2893 to 0.3295	$0.3088^{+0.0067}_{-0.0066}$	0.2893 to 0.3295
$\theta_{12}/^\circ$		$33.67^{+0.73}_{-0.71}$	31.61 to 35.94	$33.67^{+0.73}_{-0.71}$	31.61 to 35.94	$\theta_{12}/^\circ$		$33.76^{+0.42}_{-0.41}$	32.54 to 35.03	$33.76^{+0.42}_{-0.41}$	32.54 to 35.03
$\sin^2 \theta_{23}$		$0.454^{+0.019}_{-0.016}$	0.411 to 0.606	$0.568^{+0.016}_{-0.021}$	0.412 to 0.611	$\sin^2 \theta_{23}$		$0.470^{+0.017}_{-0.014}$	0.435 to 0.584	$0.550^{+0.013}_{-0.016}$	0.439 to 0.584
$\theta_{23}/^\circ$		$42.3^{+1.1}_{-0.9}$	39.9 to 51.1	$48.9^{+0.9}_{-1.2}$	39.9 to 51.4	$\theta_{23}/^\circ$		$43.29^{+0.96}_{-0.79}$	41.27 to 49.86	$47.90^{+0.73}_{-0.92}$	41.51 to 49.83
$\sin^2 \theta_{13}$		$0.02224^{+0.00056}_{-0.00057}$	0.02047 to 0.02397	$0.02222^{+0.00069}_{-0.00057}$	0.02049 to 0.02420	$\sin^2 \theta_{13}$		$0.02248^{+0.00055}_{-0.00059}$	0.02064 to 0.02418	$0.02262^{+0.00057}_{-0.00056}$	0.02093 to 0.02441
$\theta_{13}/^\circ$		$8.58^{+0.11}_{-0.11}$	8.23 to 8.91	$8.57^{+0.13}_{-0.11}$	8.23 to 8.95	$\theta_{13}/^\circ$		$8.62^{+0.11}_{-0.11}$	8.26 to 8.95	$8.65^{+0.11}_{-0.11}$	8.32 to 8.99
$\delta_{CP}/^\circ$		232^{+39}_{-25}	139 to 350	273^{+24}_{-26}	195 to 342	$\delta_{CP}/^\circ$		212^{+26}_{-36}	125 to 365	274^{+22}_{-25}	203 to 335
$\frac{\Delta m_{21}^2}{10^{-5} \text{ eV}^2}$		$7.41^{+0.21}_{-0.20}$	6.81 to 8.03	$7.41^{+0.21}_{-0.20}$	6.81 to 8.03	$\frac{\Delta m_{21}^2}{10^{-5} \text{ eV}^2}$		$7.537^{+0.094}_{-0.10}$	7.236 to 7.823	$7.537^{+0.094}_{-0.10}$	7.236 to 7.822
$\frac{\Delta m_{3\ell}^2}{10^{-3} \text{ eV}^2}$		$+2.505^{+0.024}_{-0.026}$	$+2.426$ to $+2.586$	$-2.487^{+0.027}_{-0.024}$	-2.566 to -2.407	$\frac{\Delta m_{3\ell}^2}{10^{-3} \text{ eV}^2}$		$+2.511^{+0.021}_{-0.020}$	$+2.450$ to $+2.576$	$-2.483^{+0.020}_{-0.020}$	-2.547 to -2.421

1. NuFIT (1D, 2D + single 3D data set)

<http://www.nu-fit.org>



Home

v6.1: Three-neutrino fit based on data available in November 2025

2. χ^2 profiles from Valencia neutrino global fit (1D, 2D data sets)

<https://globalfit.astroparticles.es>

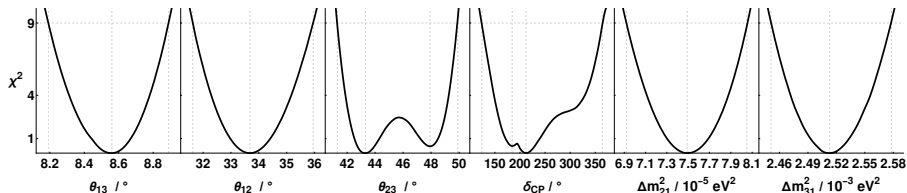
**NEUTRINO
GLOBAL
FIT**

RESULTS FROM THE NEUTRINO
OSCILLATIONS GLOBAL FIT

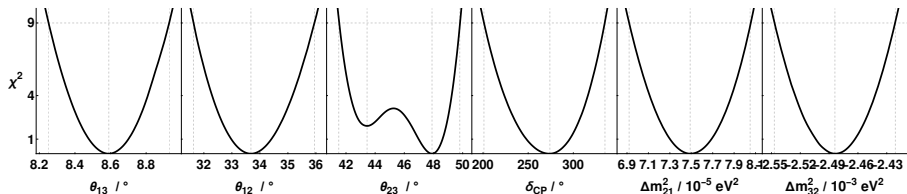
χ^2 profiles

1D profiles of oscillation parameters, 3σ ranges, b-fs, NuFIT v6.0

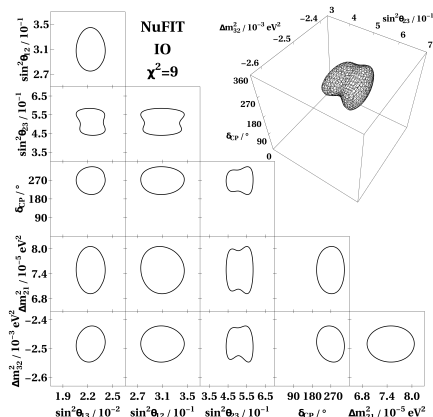
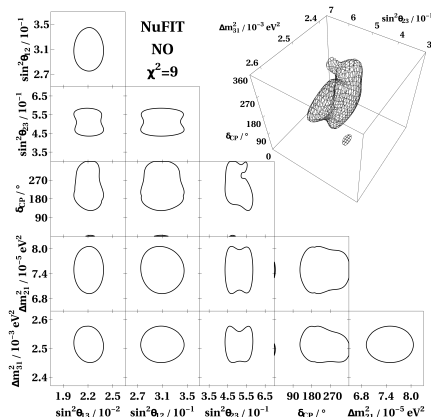
NO



IO



2D and 3D profiles, correlated data, NuFIT v6.0.



Currently, up to 3D datasets are available.

Mathematica output based on NuFIT v6.0 datasets.

```

In[686]:= {{{"", "sin²θ₂₃=0.57", "Δm²₃₁=0.00255 eV²", "δ_CP=290°"},
             {"χ²", deltachisq["NO"] [{"T23"}][5.7], deltachisq["NO"] [{"DMA"}][2.55], deltachisq["NO"] [{"DCP"}][290]}} // TableForm
Out[686]//TableForm=
      sin²θ₂₃=0.57    Δm²₃₁=0.00255 eV²    δ_CP=290°
      χ²    2.71635      3.07873      2.82341 ← treated independently, each χ² ~ 3

In[687]:= {{{"", "sin²θ₂₃=0.57 ∧ Δm²₃₁=0.00255 eV²"}, {"χ²", deltachisq["NO"] [{"T23", "DMA"}][5.7, 2.55]}} // TableForm
Out[687]//TableForm=
      sin²θ₂₃=0.57 ∧ Δm²₃₁=0.00255 eV²
      χ²    4.68546 ←

In[688]:= {{{"", "Δm²₃₁=0.00255 eV² ∧ δ_CP=290°"}, {"χ²", deltachisq["NO"] [{"DMA", "DCP"}][2.55, 290]}} // TableForm
Out[688]//TableForm=
      Δm²₃₁=0.00255 eV² ∧ δ_CP=290°
      χ²    7.56033 ← composition of two, χ² ~ 5 to 13

In[689]:= {{{"", "sin²θ₂₃=0.57 ∧ δ_CP=290°"}, {"χ²", deltachisq["NO"] [{"T23", "DCP"}][5.7, 290]}} // TableForm
Out[689]//TableForm=
      sin²θ₂₃=0.57 ∧ δ_CP=290°
      χ²    13.454 ←

In[690]:= {{{"", "sin²θ₂₃=0.57 ∧ Δm²₃₁=0.00255 eV² ∧ δ_CP=290°"}, {"χ²", deltachisq["NO"] [{"T23", "DMA", "DCP"}][5.7, 2.55, 290]}} // TableForm
Out[690]//TableForm=
      sin²θ₂₃=0.57 ∧ Δm²₃₁=0.00255 eV² ∧ δ_CP=290°
      χ²    17.4936 ← composition of three, χ² ~ 17
    
```

Insight into the χ^2 value is crucial to confront theory with exp.

Extend to full range.

To extend the number of random variables we introduce **conservative estimations**:

$$\chi_{\text{ext}}^2(X_1, \dots, X_i) \equiv \max \left\{ \chi^2(X_1, \dots, X_{k-1}, X_{k+1}, \dots, X_i), \quad k \in \{1 \dots i\} \right\}.$$

and in the context of neutrino parameters e.g.

$$\chi_{\text{ext}}^2(\theta_{13}, \Delta m_{31}^2, \delta_{\text{CP}}) \equiv \max \left\{ \chi_{\text{NuFIT}}^2(\Delta m_{31}^2, \delta_{\text{CP}}), \chi_{\text{NuFIT}}^2(\theta_{13}, \delta_{\text{CP}}), \chi_{\text{NuFIT}}^2(\theta_{13}, \Delta m_{31}^2) \right\}$$

Tested against complete minimized NuFIT data sets, both mass orderings, that

$$\chi^2(X_1, \dots, X_i) \geq \max \left\{ \chi^2(X_1, \dots, X_{k-1}, X_{k+1}, \dots, X_i), \quad k \in \{1 \dots i\} \right\}.$$

Example NuFIT data test:

$$\chi^2(\theta_{23}, \Delta m_{31}^2, \delta_{\text{CP}}) \geq \max \left\{ \chi^2(\Delta m_{31}^2, \delta_{\text{CP}}), \chi^2(\theta_{23}, \delta_{\text{CP}}), \chi^2(\theta_{23}, \Delta m_{31}^2) \right\}$$

$$17.5 \geq \max\{7.6, 13.5, 4.7\} \quad \leftarrow \text{numbers taken from slide 11}$$

due to " $\geq$ " underestimation ($13.5 < 17.5$) $\Rightarrow$ no risk of excluding allowed region (volume).

Reduce.

To reduce the number of random variables, we carry out projections:

$$\chi_{\text{red}}^2(X_1, \dots, X_{k-1}, X_{k+1}, \dots, X_n) \equiv \min \left\{ \chi^2(X_1, \dots, X_k = x_k, \dots, X_n) \text{ for any } x_k \right\}.$$

and in the context of neutrino parameters e.g.

$$\chi_{\text{red}}^2(\theta_{23}, \theta_{13}, \delta_{\text{CP}}) = \min \left\{ \chi^2(\theta_{23}, \theta_{13}, \theta_{12}, \delta_{\text{CP}}) \text{ for any } \theta_{12} \right\}.$$

Tested against complete minimized NuFIT data sets, both mass orderings, that

$$\chi^2(X_1, \dots, X_{k-1}, X_{k+1}, \dots, X_n) = \min \left\{ \chi^2(X_1, \dots, X_k = x_k, \dots, X_n) \text{ for any } x_k \right\}.$$

Example:

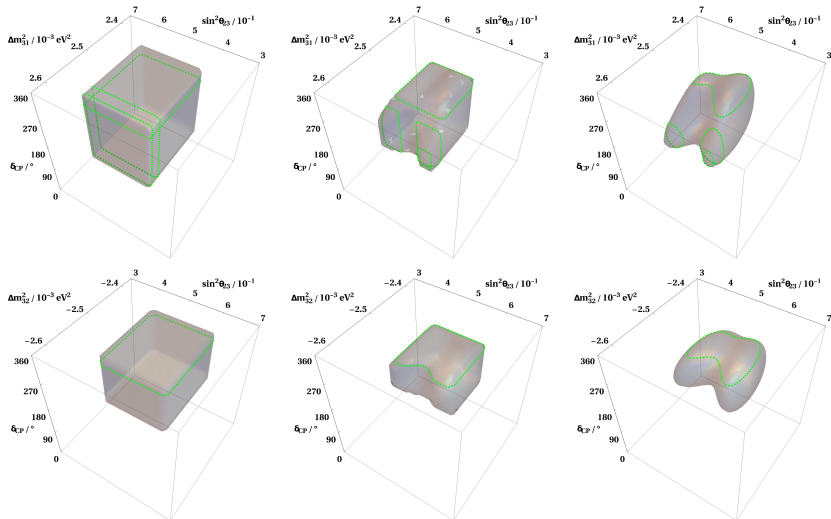
$$\chi_{\text{NuFIT}}^2(\delta_{\text{CP}}) = \min \left\{ \chi_{\text{NuFIT}}^2(\theta_{23}, \delta_{\text{CP}}) \text{ for any } \theta_{23} \right\}.$$

A full (6D) correlated data from global fit (experiment) would further constrain the results.

The extension formula may still be useful due to numerical complexity and memory requirements of 4D-6D data sets.

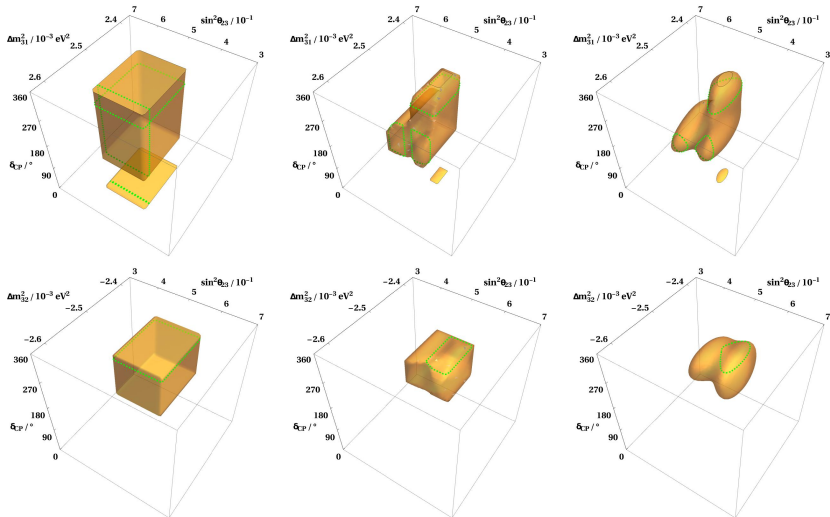
Reconstructed χ^2 volumes from NuFIT (v5.3) 1D and 2D data vs 3D data

Here presented for $\chi^2 \leq 9$, NO (upper panel) and IO (lower panel).
 max(...) extension formula (slide 12) applied to 1D and 2D data to reconstruct the 3D volume
 1D data only (3 σ ranges) 2D correlated data only 3D correlated data



Reconstructed χ^2 volumes from NuFIT (v6.1) 1D and 2D data vs 3D data

Here presented for $\chi^2 \leq 9$, NO (upper panel) and IO (lower panel).
 max(...) extension formula (slide 12) applied to 1D and 2D data to reconstruct the 3D volume
 1D data only (3 σ ranges) 2D correlated data only 3D correlated data



$\chi^2(\theta_{13}, \theta_{12}, \theta_{23}, \delta_{\text{CP}}, \Delta m_{21}^2, \Delta m_{31}^2)$ from global fit (could be without masses)

The NuFIT collaboration performing a global analysis of neutrino oscillation data provides one-dimensional χ^2 projections for $\sin^2 \theta_{ij}$ and δ [49]. We denote them as $\chi_i^2(x_i)$, $i = 1, 2, 3, 4$, where x_i are components of $\vec{x} = (\sin^2 \theta_{12}, \sin^2 \theta_{13}, \sin^2 \theta_{23}, \delta)$. Using these projections, we construct an approximate global χ^2 function as

arXiv:1804.00182

$$\chi^2(\vec{x}) = \sum_{i=1}^4 \chi_i^2(x_i). \quad (4.1)$$

Given the distinct experimental sensitivities, we construct the total χ^2 function as the sum of individual contributions from the solar, reactor, and atmospheric sectors, namely,

$$\chi^2 = \chi_{\text{Solar}}^2 [\theta_{12}(\mathbf{p}), \Delta m_{21}^2(\mathbf{p})] + \chi_{\text{Atmos}}^2 [\theta_{23}(\mathbf{p}), \Delta m_{3\ell}^2(\mathbf{p})] + \chi_{\text{Reactor}}^2 [\theta_{13}(\mathbf{p})], \quad (19)$$

arXiv:2511.15654

$\max(\dots)$ formula at the end always return unchanged NuFIT χ^2 value (except needed interpolation to join the discrete data points - no extrapolation).

$$\chi_{\text{ext}}^2(\theta_{13}, \Delta m_{31}^2, \delta_{\text{CP}})$$

is estimated by

$$\max \begin{cases} \chi_{\text{NuFIT}}^2(\Delta m_{31}^2, \delta_{\text{CP}}) \\ \chi_{\text{NuFIT}}^2(\theta_{13}, \delta_{\text{CP}}) \\ \chi_{\text{NuFIT}}^2(\theta_{13}, \Delta m_{31}^2) \end{cases}$$

The maximum is taken after checking each χ^2 value of global fit available subsets of parameters.

$\theta_{13} \neq 0$, Daya Bay, RENO (2012)

BM, TB, GR, HG disfavored by non-zero θ_{13} .

$$U_{\text{PMNS}} = \begin{bmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{13}s_{23}e^{i\delta} & c_{12}c_{23} - s_{12}s_{13}s_{23}e^{i\delta} & c_{13}s_{23} \\ s_{12}s_{23} - c_{12}s_{13}c_{23}e^{i\delta} & -c_{12}s_{23} - s_{12}s_{13}c_{23}e^{i\delta} & c_{13}c_{23} \end{bmatrix}$$

$$\theta_{13} \neq 0^\circ$$

$\Downarrow$

$$\theta_{23} = 45^\circ$$

$$U_0 = \begin{bmatrix} c_{12} & s_{12} & 0 \\ -\frac{s_{12}}{\sqrt{2}} & \frac{c_{12}}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{s_{12}}{\sqrt{2}} & -\frac{c_{12}}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{bmatrix}$$

$\Downarrow$

$$\theta_{12} = 45^\circ, s_{12} = 1/\sqrt{2}$$

$$\theta_{12} = 35, 26^\circ, s_{12} = 1/\sqrt{3}$$

$$\theta_{12} = 31, 7^\circ$$

$$\theta_{12} = 30^\circ, s_{12} = 1/2$$

Bimaximal Mixing

$$\begin{bmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} & 0 \\ -\frac{1}{2} & \frac{1}{2} & \frac{1}{\sqrt{2}} \\ \frac{1}{2} & -\frac{1}{2} & \frac{1}{\sqrt{2}} \end{bmatrix}$$

Tribimaximal Mixing

$$\begin{bmatrix} \sqrt{\frac{2}{3}} & \frac{1}{\sqrt{3}} & 0 \\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{2}} \\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \end{bmatrix}$$

Golden Ratio Mixing

$$\begin{bmatrix} \frac{1}{\sqrt{2+\varphi}} & \frac{1}{\sqrt{2+\varphi}} & 0 \\ \frac{-1}{\sqrt{4+2\varphi}} & \frac{\varphi}{\sqrt{4+2\varphi}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{4+2\varphi}} & \frac{-\varphi}{\sqrt{4+2\varphi}} & \frac{1}{\sqrt{2}} \end{bmatrix}$$

Hexagonal Mixing

$$\begin{bmatrix} \sqrt{\frac{3}{4}} & \frac{1}{2} & 0 \\ -\frac{1}{2\sqrt{2}} & \frac{\sqrt{3}}{2\sqrt{2}} & -\frac{1}{\sqrt{2}} \\ -\frac{1}{2\sqrt{2}} & \frac{\sqrt{3}}{2\sqrt{2}} & \frac{1}{\sqrt{2}} \end{bmatrix}$$

Golden Ratio Mixing: $\tan \theta_{12} = 1/\varphi$, $\varphi = (1 + \sqrt{5})/2$ being the golden ratio.

Non-zero θ_{13} : Successors of tribimaximal mixing

$$U_{\text{TBM}} = \begin{bmatrix} \frac{2}{\sqrt{6}} & \frac{1}{\sqrt{3}} & 0 \\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{2}} \\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \end{bmatrix}, \quad U_{\text{PMNS}} \simeq \begin{bmatrix} \frac{2}{\sqrt{6}} & \frac{1}{\sqrt{3}} & 0.15 \\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{6}} & -\frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \end{bmatrix},$$

```
In[688]:= Sin[8.6*Pi/180]
Sin[8.6*Pi/180]^2
Out[688]= 0.149535
Out[689]= 0.0223608
```

$$|U_{\text{TM}_1}| = \begin{bmatrix} \frac{2}{\sqrt{6}} & * & * \\ \frac{1}{\sqrt{6}} & * & * \\ \frac{1}{\sqrt{6}} & * & * \end{bmatrix}, \quad |U_{\text{TM}_2}| = \begin{bmatrix} * & \frac{1}{\sqrt{3}} & * \\ * & \frac{1}{\sqrt{3}} & * \\ * & \frac{1}{\sqrt{3}} & * \end{bmatrix},$$

$$U_{\text{TM}_1} = \begin{bmatrix} \frac{2}{\sqrt{6}} & \frac{c_\theta}{\sqrt{3}} & \frac{s_\theta}{\sqrt{3}} e^{-i\gamma} \\ -\frac{1}{\sqrt{6}} & \frac{c_\theta}{\sqrt{3}} - \frac{s_\theta}{\sqrt{2}} e^{i\gamma} & -\frac{s_\theta}{\sqrt{3}} e^{-i\gamma} - \frac{c_\theta}{\sqrt{2}} \\ -\frac{1}{\sqrt{6}} & \frac{c_\theta}{\sqrt{3}} - \frac{s_\theta}{\sqrt{2}} e^{i\gamma} & -\frac{s_\theta}{\sqrt{3}} e^{-i\gamma} + \frac{c_\theta}{\sqrt{2}} \end{bmatrix},$$

$$U_{\text{TM}_2} = \begin{bmatrix} \frac{2c_\theta}{\sqrt{6}} & \frac{1}{\sqrt{3}} & \frac{2s_\theta}{\sqrt{6}} e^{-i\gamma} \\ -\frac{c_\theta}{\sqrt{6}} + \frac{s_\theta}{\sqrt{2}} e^{i\gamma} & \frac{1}{\sqrt{3}} & -\frac{s_\theta}{\sqrt{3}} e^{-i\gamma} - \frac{c_\theta}{\sqrt{2}} \\ -\frac{c_\theta}{\sqrt{6}} + \frac{s_\theta}{\sqrt{2}} e^{i\gamma} & \frac{1}{\sqrt{3}} & -\frac{s_\theta}{\sqrt{3}} e^{-i\gamma} + \frac{c_\theta}{\sqrt{2}} \end{bmatrix}.$$

Mixing-schemes structure: TM_1 and TM_2 , partial μ - τ

Flavor and mass mixing:

$$\begin{bmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{bmatrix} = \begin{bmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{bmatrix} \begin{bmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{bmatrix}.$$

TM_1 and TM_2 mixing-schemes structure

$$|U_{\text{TM}_1}| = \begin{bmatrix} \frac{2}{\sqrt{6}} & * & * \\ \frac{1}{\sqrt{6}} & * & * \\ \frac{1}{\sqrt{6}} & * & * \end{bmatrix}, \quad |U_{\text{TM}_2}| = \begin{bmatrix} * & \frac{1}{\sqrt{3}} & * \\ * & \frac{1}{\sqrt{3}} & * \\ * & \frac{1}{\sqrt{3}} & * \end{bmatrix}.$$

Partial μ - τ ($\mu 1$ - $\tau 1$, $\mu 2$ - $\tau 2$) reflection symmetry mixing structure:

$$|U_{\mu 1}| = |U_{\tau 1}|, \quad |U_{\mu 2}| = |U_{\tau 2}|.$$

TM_1 variant of more general $\mu 1$ - $\tau 1$

TM_2 variant of more general $\mu 2$ - $\tau 2$

TM₁ and TM₂ mixing schemes, partial $\mu - \tau$ symmetry

Comparing the corresponding elements of the first column of U_{PMNS} and U_{TM_1} .

$$|U_{e1}|^2 = c_{12}^2 c_{13}^2 = 2/3 \quad : \quad \theta_{12}^{\text{TM}_1}(\theta_{13}) \quad : \quad s_{12}^2 = \frac{1 - 3s_{13}^2}{3 - 3s_{13}^2},$$

$$|U_{\mu 1}|^2 = |U_{\tau 1}|^2 = 1/6 \quad : \quad \delta_{\text{CP}}^{\text{TM}_1}(\theta_{13}, \theta_{23}) \quad : \quad \cos \delta_{\text{CP}} = \frac{(1 - 5s_{13}^2)(2s_{23}^2 - 1)}{4s_{13}s_{23}\sqrt{2(1 - 3s_{13}^2)(1 - s_{23}^2)}}.$$

Comparing the corresponding elements of the second column of U_{PMNS} and U_{TM_2} .

$$|U_{e2}|^2 = s_{12}^2 c_{13}^2 = 1/3 \quad : \quad \theta_{12}^{\text{TM}_2}(\theta_{13}) \quad : \quad s_{12}^2 = \frac{1}{3 - 3s_{13}^2},$$

$$|U_{\mu 2}|^2 = |U_{\tau 2}|^2 = 1/3 \quad : \quad \delta_{\text{CP}}^{\text{TM}_2}(\theta_{13}, \theta_{23}) \quad : \quad \cos \delta_{\text{CP}} = -\frac{(2 - 4s_{13}^2)(2s_{23}^2 - 1)}{4s_{13}s_{23}\sqrt{(2 - 3s_{13}^2)(1 - s_{23}^2)}}.$$

In partial μ - τ reflection symmetry, the mixing matrix symmetry is given by:

$$|U_{\mu 1}| = |U_{\tau 1}| \quad : \quad \delta_{\text{CP}}^{\mu 1 - \tau 1}(\theta_{13}, \theta_{12}, \theta_{23}) \quad : \quad \cos \delta_{\text{CP}} = \frac{(c_{23}^2 - s_{23}^2)(c_{12}^2 s_{13}^2 - s_{12}^2)}{4c_{12}s_{12}c_{23}s_{23}s_{13}},$$

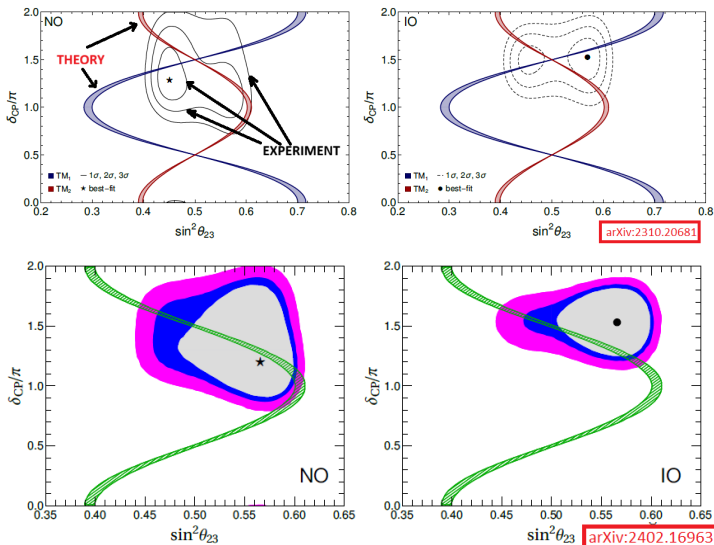
$$|U_{\mu 2}| = |U_{\tau 2}| \quad : \quad \delta_{\text{CP}}^{\mu 2 - \tau 2}(\theta_{13}, \theta_{12}, \theta_{23}) \quad : \quad \cos \delta_{\text{CP}} = \frac{(c_{23}^2 - s_{23}^2)(c_{12}^2 - s_{12}^2 s_{13}^2)}{4c_{12}s_{12}c_{23}s_{23}s_{13}}.$$

Summary of relations arising from mixing-schemes structure

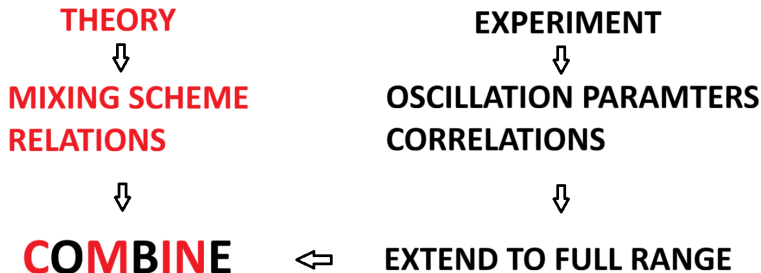
TM ₁	TM ₂
$s_{12}^2 = \frac{1-3s_{13}^2}{3-3s_{13}^2}$	$s_{12}^2 = \frac{1}{3-3s_{13}^2}$
$\cos \delta_{\text{CP}} = \frac{(1-5s_{13}^2)(2s_{23}^2-1)}{4s_{13}s_{23}\sqrt{2(1-3s_{13}^2)(1-s_{23}^2)}}$	$\cos \delta_{\text{CP}} = -\frac{(2-4s_{13}^2)(2s_{23}^2-1)}{4s_{13}s_{23}\sqrt{(2-3s_{13}^2)(1-s_{23}^2)}}$
$\mu 1 - \tau 1$	$\mu 2 - \tau 2$
$\cos \delta_{\text{CP}} = \frac{(c_{23}^2-s_{23}^2)(c_{12}^2s_{13}^2-s_{12}^2)}{4c_{12}s_{12}c_{23}s_{23}s_{13}}$	$\cos \delta_{\text{CP}} = \frac{(c_{23}^2-s_{23}^2)(c_{12}^2-s_{12}^2s_{13}^2)}{4c_{12}s_{12}c_{23}s_{23}s_{13}}.$

How to confront the mixing schemes with experimental data ?

Confronting mixing schemes with oscillation data (examples)



⇒ Next, we will combine theory with experiment which are flashed independently on the plots.



TM₁ mixing relations + extended to full range correlated data, example:

$$\begin{array}{c} \sin^2 \theta_{12} = \frac{1-3 \sin^2 \theta_{13}}{3-3 \sin^2 \theta_{13}} \\ \downarrow \\ \chi^2(\theta_{13}, \theta_{12}, \theta_{23}, \delta_{\text{CP}}, \Delta m_{21}^2, \Delta m_{31}^2) \\ \uparrow \\ \cos \delta_{\text{CP}} = \frac{(1-5 \sin^2 \theta_{13})(2 \sin_{23}^2 - 1)}{4 \sin_{13} \sin \theta_{23} \sqrt{2(1-3 \sin^2 \theta_{13})(1-\sin^2 \theta_{23})}} \end{array}$$

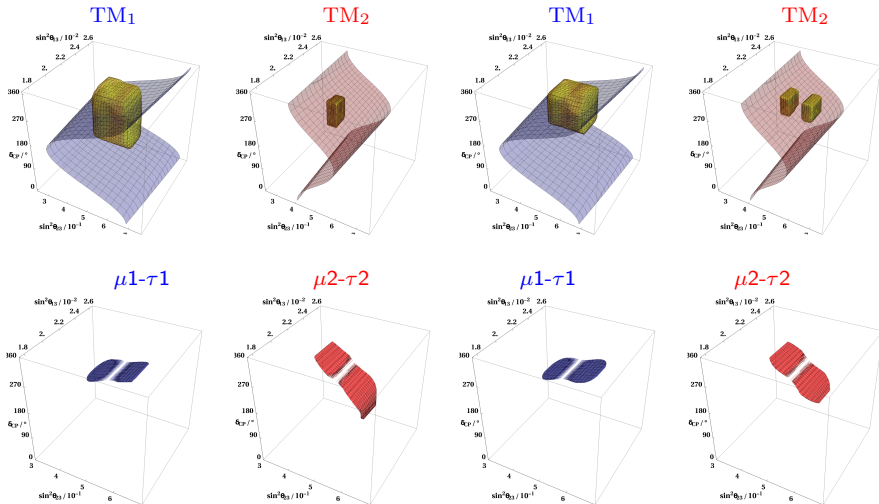
In mixing schemes some parameter values are determined by other parameters values

For any choice of parameters, the equations for these parameters must be solved first.

NO and IO need to be treated separately.

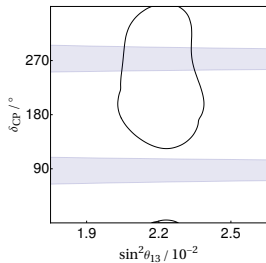
Intersections of correlated data and mixing-scheme relations

$\Delta\chi^2 \leq 9$, NO (left panel) and IO (right panel) (volumes NuFIT v5.3 + $\theta_{12}(\theta_{13})$ relation):

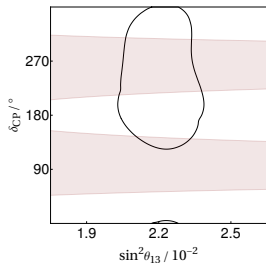
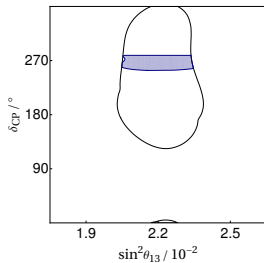


3D projections of 4D intersections for partial μ - τ reflection symmetry.

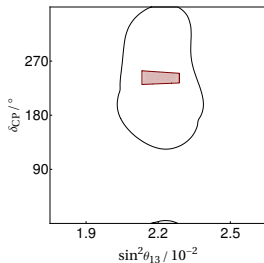
Impact of correlated data on model's predictions (NuFIT v6.0)



THEORY AND EXPERIMENT
 FLASHED
 INDEPENDENTLY
 TM_1
 COMBINED

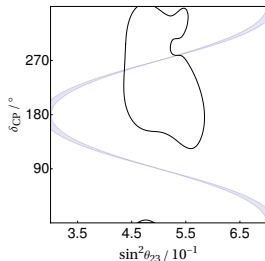


THEORY AND EXPERIMENT
 FLASHED
 INDEPENDENTLY
 TM_2
 COMBINED

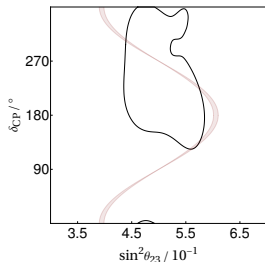
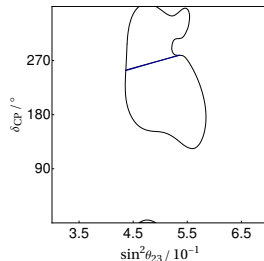


NO, $\chi^2 < 9$

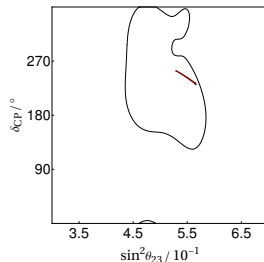
Impact of correlated data on model's predictions (NuFIT v6.0)



THEORY AND EXPERIMENT
 ← **FLASHED INDEPENDENTLY**
TM₁
COMBINED →

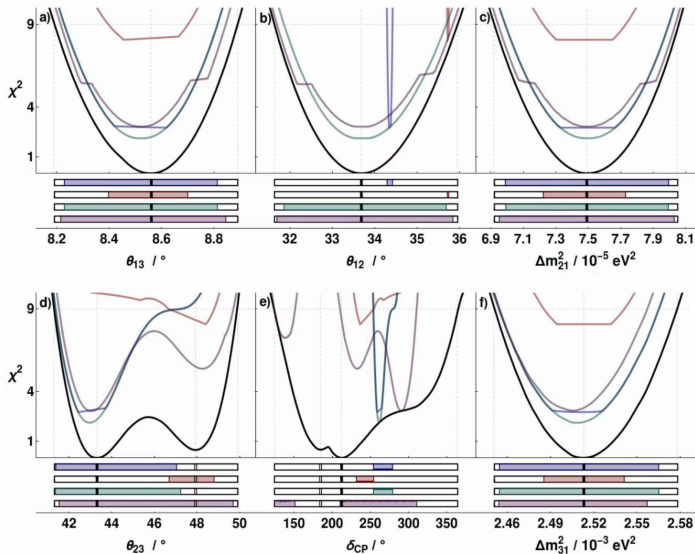


THEORY AND EXPERIMENT
 ← **FLASHED INDEPENDENTLY**
TM₂
COMBINED →



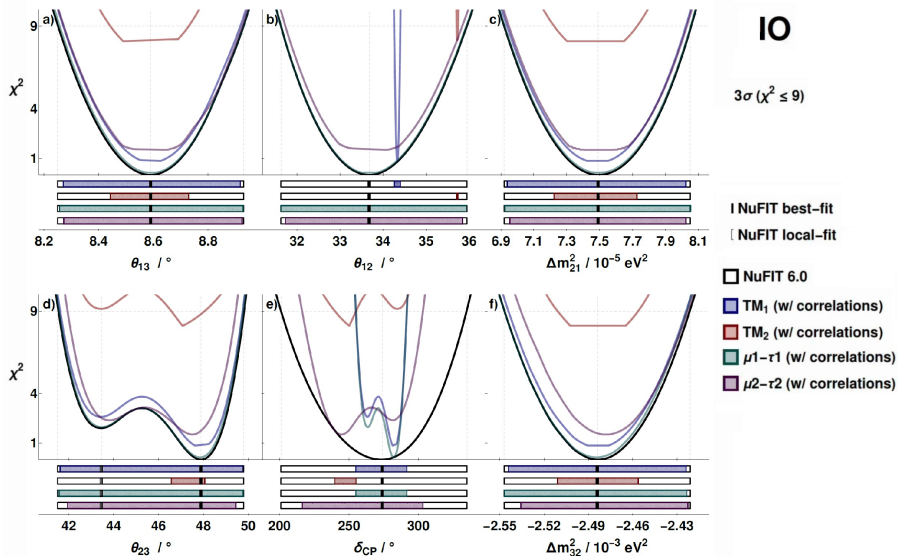
NO, $\chi^2 < 9$

Extracted ranges of model's parameters from correlated data (NuFIT v6.0)



$$\chi^2_{\min}: \text{TM}_1 = 2.8, \text{TM}_2 = 8.1, \mu 1 - \tau 1 = 2.1, \mu 2 - \tau 2 = 2.8$$

Extracted ranges of model's parameters from correlated data (NuFIT v6.0)



$$\chi^2_{min}: \text{TM}_1 = 0.9, \text{TM}_2 = 8.1, \mu 1 - \tau 1 = 0.1, \mu 2 - \tau 2 = 1.5$$

Results for NuFIT 6.0 data sets (w/ atmospheric ν)					
		TM ₁	TM ₂	$\mu 1-\tau 1$	$\mu 2-\tau 2$
χ^2_{min}	NO	2.8	8.1	2.1	2.8
	IO	0.9	8.1	0.1	1.5

Results for NuFIT 6.1 data sets (w/ atmospheric ν)					
		TM ₁	TM ₂	$\mu 1-\tau 1$	$\mu 2-\tau 2$
χ^2_{min}	NO	3.9	21.5	2.4	3.2
	IO	1.9	21.5	0.1	1.6

JUNO results change numbers, not preferences.

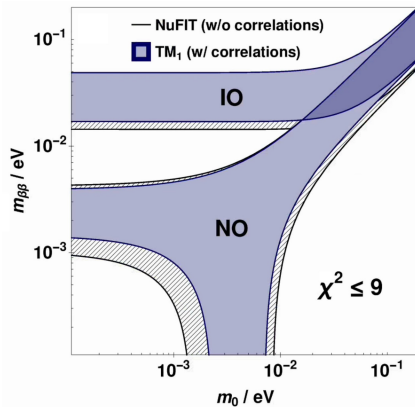
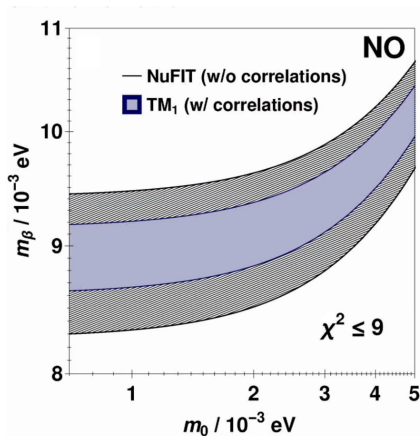
Can we minimize χ^2 ?

We can confront the constraints coming from mixing schemes to global-fit data.

$$\begin{cases} \chi^2(\theta_{13}, \theta_{12}, \theta_{23}, \delta_{\text{CP}}, \Delta m_{21}^2, \Delta m_{31}^2) \leq 9, \\ m_\beta = \sqrt{m_0^2 + \Delta m_{21}^2 c_{13}^2 s_{12}^2 + \Delta m_{31}^2 s_{13}^2}, \\ s_{12}^2 = \frac{1-3s_{13}^2}{3-3s_{13}^2}, \\ \cos \delta_{\text{CP}} = \frac{(1-5s_{13}^2)(2s_{23}^2-1)}{4s_{13}s_{23}\sqrt{2(1-3s_{13}^2)(1-s_{23}^2)}}. \end{cases}$$

$$\left\{ \begin{array}{l} \chi^2(\theta_{13}, \theta_{12}, \theta_{23}, \delta_{\text{CP}}, \Delta m_{21}^2, \Delta m_{31}^2) \leq 9, \\ m_{\beta\beta} = \left| m_0 c_{13}^2 c_{12}^2 e^{i2(\alpha_1 - \delta_{\text{CP}})} + \sqrt{m_0^2 + \Delta m_{21}^2} c_{13}^2 s_{12}^2 e^{i2(\alpha_2 - \delta_{\text{CP}})} \right. \\ \quad \left. + \sqrt{m_0^2 + \Delta m_{31}^2} s_{13}^2 \right|, \\ s_{12}^2 = \frac{1 - 3s_{13}^2}{3 - 3s_{13}^2}, \\ \cos \delta_{\text{CP}} = \frac{(1 - 5s_{13}^2)(2s_{23}^2 - 1)}{4s_{13}s_{23}\sqrt{2(1 - 3s_{13}^2)(1 - s_{23}^2)}}, \\ 0 \leq \alpha_1, \alpha_2 \leq 2\pi. \end{array} \right.$$

Phenomenological predictions: m_β and $m_{\beta\beta}$ (NuFIT v6.0)



Solution of eqs from slides [32](#) and [33](#)

Multidimensional correlated datasets on neutrino oscillation parameters allows for confronting the theoretical models with experimental data

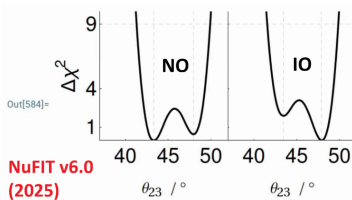
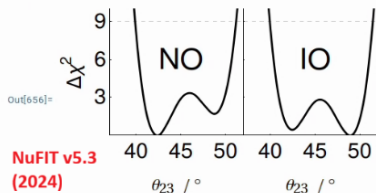
With the presented approach, we

- find χ^2_{min} "floor" for the model
- test TM_1 , TM_2 and partial μ - τ mixing-schemes over the full range of parameters at a given χ^2 level
- some parameters are exceptionally constrained (for NuFIT v6.0)
 - θ_{12} for TM_1 and TM_2 ,
 - θ_{23} for TM_2 ,
 - δ_{CP} for all considered mixing-schemes,
 - Δm^2_{21} and $\Delta m^2_{31(32)}$ for TM_2
- TM_2 is outside the reasonable χ^2_{min} value (for NuFIT v6.1)
- find tighter, model-dependent constraints on the effective neutrino masses
- potentially could be used for χ^2 estimation in $SO(3)$ parametrization

The method is general and can be applied to any model providing mixing-scheme relations in the form analytical expressions involving the neutrino oscillation parameters.

Thank you
for your attention.

3 ν mixing



```

In[658]:= {deltachisq[#][{"T13"}][0] & /@ {"NO", "IO"},
           deltachisq[#][{"DCP"}][0] & /@ {"NO", "IO"}}
Out[658]:= {{1674.27, 1664.84}, {11.0731, 14.4343}}
    
```

```

In[585]:= {deltachisq[#][{"T13"}][0] & /@ {"NO", "IO"},
           deltachisq[#][{"DCP"}][0] & /@ {"NO", "IO"}}
Out[585]:= {{1774.37, 1768.02}, {8.03189, 19.3677}}
    
```

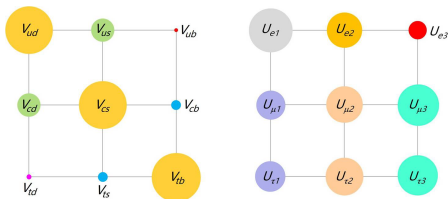
Left: Mathematica output based on NuFIT v5.3 data (2024). **Right:** Mathematica output based on NuFIT v6.0 data (2025).

NO

0.800 → 0.843	0.519 → 0.580	0.143 → 0.155
0.354 → 0.461	0.533 → 0.620	0.635 → 0.769
0.277 → 0.485	0.589 → 0.616	0.620 → 0.759

IO

0.800 → 0.843	0.519 → 0.580	0.143 → 0.156
0.288 → 0.479	0.567 → 0.606	0.635 → 0.772
0.246 → 0.527	0.585 → 0.603	0.616 → 0.759



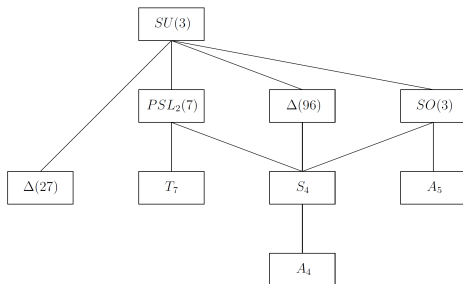
(a) CKM quark flavor mixing

(b) PMNS lepton flavor mixing

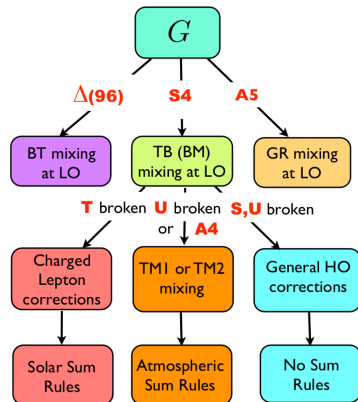
Left: Mixing matrix ($|U_{\text{PMNS}}|$) ranges in 3σ , NuFIT v5.3.

Right: Taken from [arXiv:2210.11922](https://arxiv.org/abs/2210.11922), Figure 2.

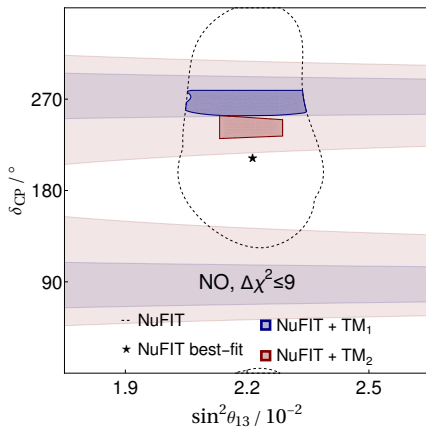
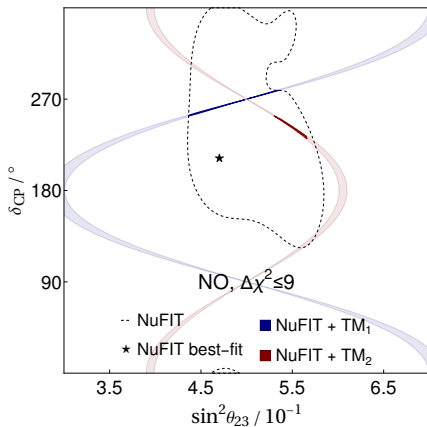
Flavor symmetries and mixing-schemes



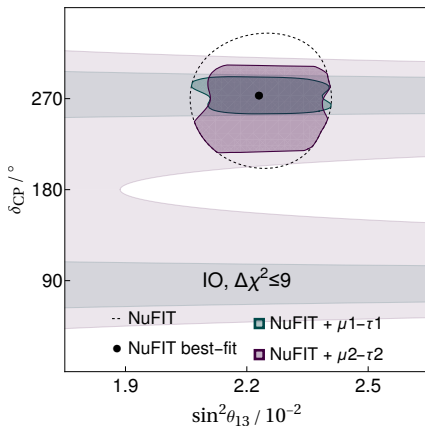
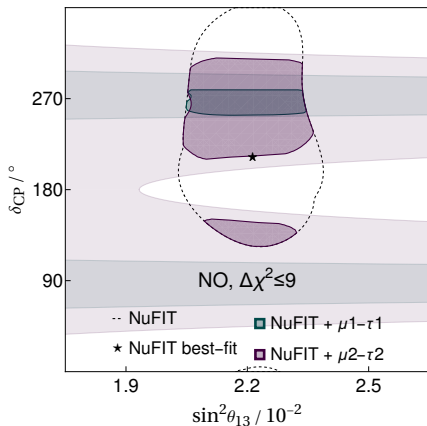
arXiv:1301.1340 S.F. King, C. Luhn



Experiment / Theory / Combined at $\chi^2 < 9$ (NuFIT v6.0)



Experiment / Theory / Combined at $\chi^2 < 9$ (NuFIT v6.0)



Effective neutrino mass

Effective electron neutrino mass:

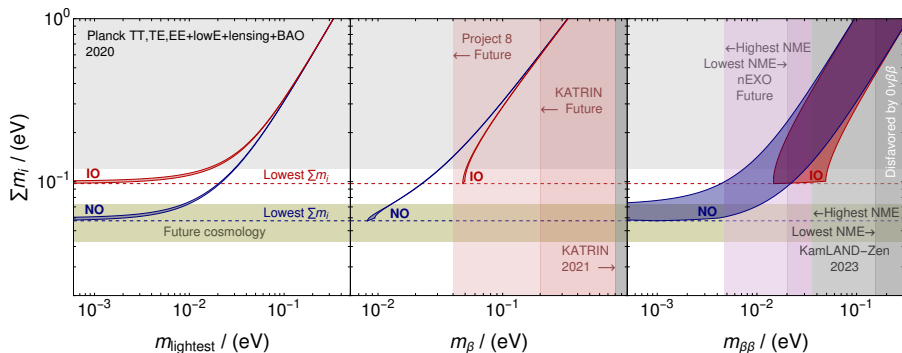
$$m_{\beta}^2 = c_{13}^2 c_{12}^2 m_1^2 + c_{13}^2 s_{12}^2 m_2^2 + s_{13}^2 m_3^2 = \begin{cases} \text{NO:} & m_0^2 + \Delta m_{21}^2 c_{13}^2 s_{12}^2 + \Delta m_{3\ell}^2 s_{13}^2, \\ \text{IO:} & m_0^2 - \Delta m_{21}^2 c_{13}^2 c_{12}^2 - \Delta m_{3\ell}^2 c_{13}^2. \end{cases}$$

Effective Majorana mass:

$$m_{\beta\beta} = \left| \sum_i m_i U_{ei}^2 \right| = \left| m_1 c_{13}^2 c_{12}^2 e^{i2\alpha_1} + m_2 c_{13}^2 s_{12}^2 e^{i2\alpha_2} + m_3 s_{13}^2 e^{-i2\delta_{\text{CP}}} \right|$$
$$= \begin{cases} \text{NO:} & m_0 \left| c_{13}^2 c_{12}^2 e^{i2(\alpha_1 - \delta_{\text{CP}})} + \sqrt{1 + \frac{\Delta m_{21}^2}{m_0^2}} c_{13}^2 s_{12}^2 e^{i2(\alpha_2 - \delta_{\text{CP}})} + \sqrt{1 + \frac{\Delta m_{3\ell}^2}{m_0^2}} s_{13}^2 \right| \\ \text{IO:} & m_0 \left| \sqrt{1 - \frac{\Delta m_{3\ell}^2 + \Delta m_{21}^2}{m_0^2}} c_{13}^2 c_{12}^2 e^{i2(\alpha_1 - \delta_{\text{CP}})} + \sqrt{1 - \frac{\Delta m_{3\ell}^2}{m_0^2}} c_{13}^2 s_{12}^2 e^{i2(\alpha_2 - \delta_{\text{CP}})} + s_{13}^2 \right|. \end{cases}$$

Here $m_0 = m_1$ (m_3) and $\ell = 1$ (2) in $\Delta m_{3\ell}^2$ for NO (IO).

Neutrino mass



Planck: $\sum_i m_i < 0.12$ eV, 95% CL, arXiv:1807.06209

KATRIN: $m_\beta < 0.8$ eV, 90% CL, arXiv:2105.08533

KamLAND-Zen: $m_{\beta\beta} < (0.036 \text{ to}) 0.156$ eV, 90% CL, arXiv:2203.02139