

Lattice Gauge Theory Challenges for Quantum Computers

Where the Euclidean paradigm hits exponential walls — and what unitarity does (and does not) buy

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Fifty years of the Euclidean paradigm — and five exponential walls

$$\langle O \rangle = Z^{-1} \int \mathcal{D}U \mathcal{D}\bar{\psi} \mathcal{D}\psi \, O e^{-S} \quad \text{importance sampling works iff } e^{-S} \geq 0 \text{ and the chain is ergodic}$$

1 Sign problem

$\mu \neq 0, \theta \neq 0$: the measure turns complex

2 Real-time

Minkowski weight e^{iS} : no probability at all

3 Signal-to-noise

Variance decays more slowly than the signal

4 Small overlaps

Trial states barely touch the states we want

5 Tunneling times

Ergodicity lost: 1st-order transitions, frozen topology

Five distinct pathologies: 1–2 live in the measure, 3–4 in the estimator, 5 in the Markov chain. All are exponential. None is negotiable.

What unitarity buys by construction

$$|\psi(t)\rangle = e^{-iHt} |\psi(0)\rangle \quad \text{— no measure, no weights, no configurations to sample}$$

SOLVED BY CONSTRUCTION

1 · Sign problem

Finite μ is just a term in H ; θ is just a parameter. Hamiltonian evolution never sees a Boltzmann weight — solved structurally, not by ingenuity.

Classical workarounds (reweighting, complex Langevin, density of states) remain exponentially costly or hard to control; a generic classical solution is NP-hard.

SOLVED BY CONSTRUCTION

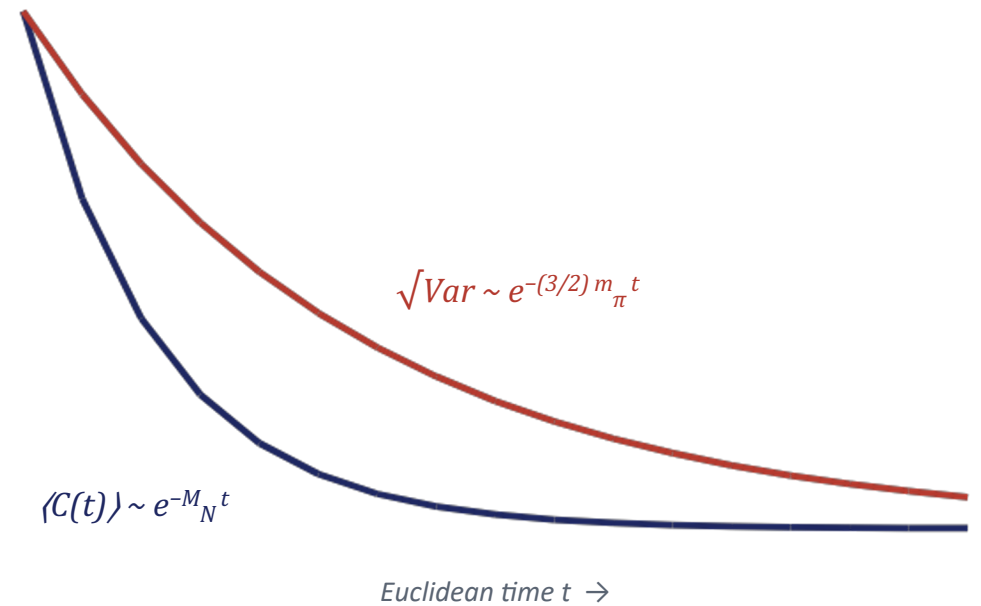
2 · Real-time dynamics

e^{-iHt} is the native hardware operation: Trotterised evolution costs $\text{poly}(V, t, 1/\epsilon)$. No analytic continuation, no ill-posed inversion of $G(\tau) = \int d\omega e^{-\omega\tau} \rho(\omega)$.

Scattering, transport, thermalisation, fragmentation become direct measurements rather than inverse problems.

Signal-to-noise: the Parisi–Lepage wall

- The variance of a correlator is itself a correlator — with its own, lighter, spectrum
- Nucleon: signal falls with M_N , noise with $(3/2) m_\pi$
- Glueballs and disconnected diagrams face constant noise — vacuum quantum numbers
- On a quantum computer, Monte Carlo noise becomes shot noise, $\sigma \sim 1/\sqrt{N_{\text{shots}}}$ — but the variance of the physical estimator can still grow exponentially with t or V

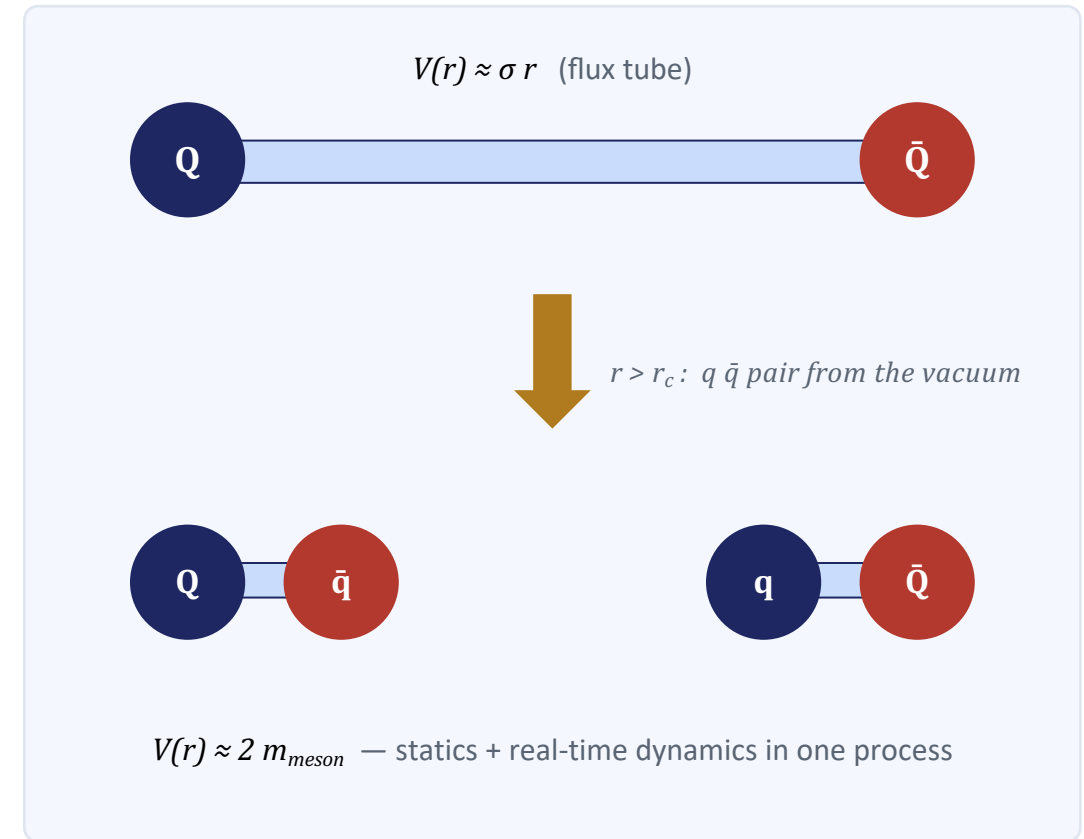


The noise outlives the signal: an exponential race that the signal always loses.

Small overlaps: string breaking as the showcase

- Spectroscopy rests on trial operators: contamination dies only like $e^{-\Delta E t}$
- String breaking is the classic case: a Wilson-line operator has tiny overlap onto the broken two-meson state — it took ~20 years of classical effort to resolve
- On quantum hardware the problem returns as state-preparation cost: repetitions $\propto 1/|\langle \psi_{\text{trial}} | \Omega \rangle|^2$; amplitude amplification helps only quadratically
- First quantum-simulator observations of string breaking appeared in 2024–25 — in simplified models, far from continuum QCD: the benchmark is set, not met

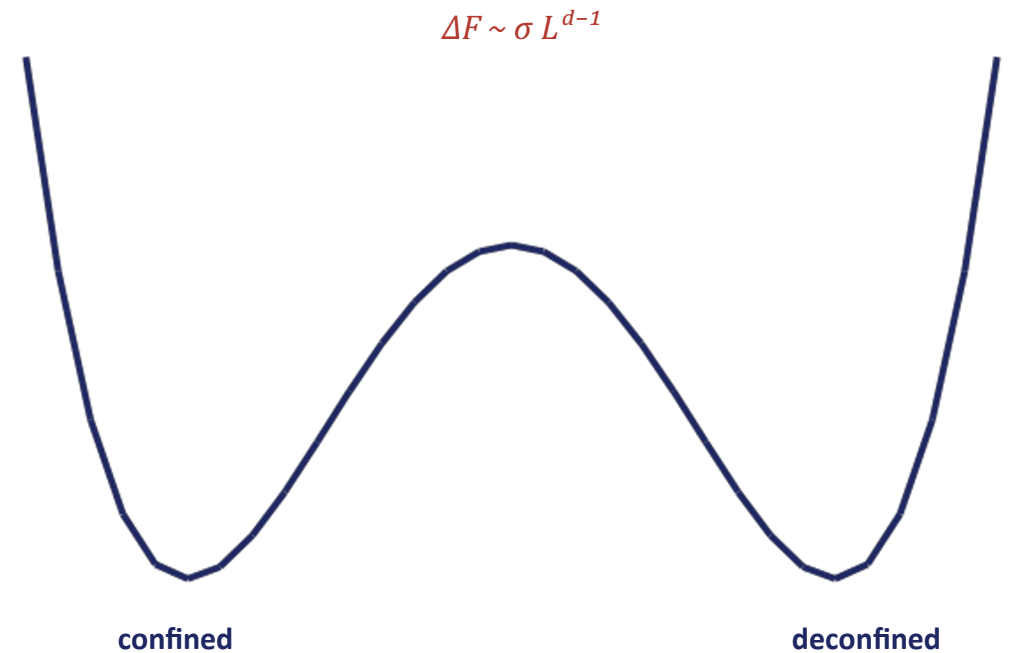
$$N_{\text{shots}} \sim 1 / |\langle \psi_{\text{trial}} | \Omega \rangle|^2$$



Exponential tunneling times & broken ergodicity

- First-order transitions: phase flips suppressed by the interface free energy — supercritical slowing down (SU(N) deconfinement is first order for $N \geq 3$)
- Topological freezing: τ_Q diverges as $a \rightarrow 0$; the Markov chain gets stuck in one sector
- Classical lesson: algorithmic ingenuity can flatten barriers — density-of-states (LLR) methods sample across coexisting phases where importance sampling fails
- On quantum hardware the barrier returns through state preparation: adiabatic evolution needs $T \gtrsim \Delta^{-2}$, and at a first-order transition the gap closes exponentially

$$\tau_{\text{flip}} \sim e^{2\sigma L^{d-1}}, \quad \Delta_{\text{min}} \sim e^{-c L^{d-1}}$$



The classical tunneling barrier returns as an exponentially small level splitting.

Scoreboard: classical vs quantum

Challenge	Classical Monte Carlo	Quantum computer
1 · Sign problem	$\exp(V)$ reweighting; generically NP-hard	SOLVED — no weights in unitary evolution
2 · Real-time dynamics	ill-posed inversion of Euclidean correlators	SOLVED — $e^{(-iHt)}$ is native
3 · Signal-to-noise	$S/N \sim \exp(-(M_N - 3/2 m_\pi)t)$	REPHRASED — shot cost of measurement
4 · Small overlaps	excited-state contamination, operator bases	REPHRASED — state preparation $\propto 1/ \text{overlap} ^2$
5 · Tunneling times	$\exp(\sigma L^{(d-1)})$ flips; frozen topology	REPHRASED — adiabatic gap closes exponentially

2 solved structurally, 3 transformed: the exponentials reappear elsewhere — from the measure and the Markov chain into state preparation and measurement.

“The sign problem dies — the other exponentials just get reframed”

A showcase: string breaking

Statics (overlaps) plus dynamics (real time) in one process. 2024–25 simulator results set the benchmark — now towards non-Abelian theories and controlled errors.

Honest benchmarks

Compare against the best classical method: in 1+1D and much of 2+1D, tensor networks also evade the sign problem. Advantage is a race, not a right.

A verification culture

If the payoff is regimes no classical method reaches, cross-check in classically accessible corners — with error bars that mean what lattice error bars mean.

*Welcome to QuantHEP 2026 — the lattice community is watching, with sympathy and with error bars.
Have a great conference!*