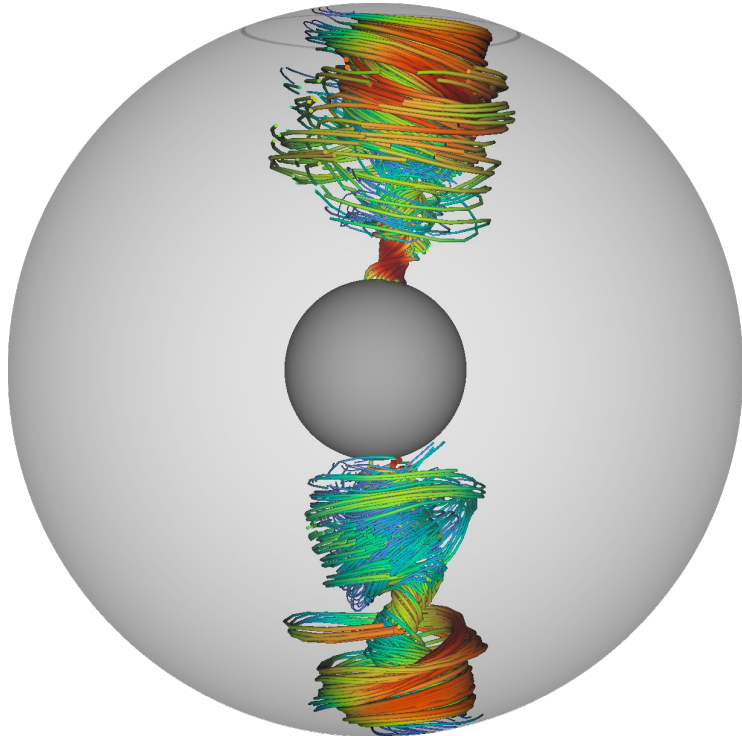


# Magnetar formation from Tayler-Spruit dynamo in protoneutron stars



Jérôme Guilet  
(CEA Saclay)

Collaborators:

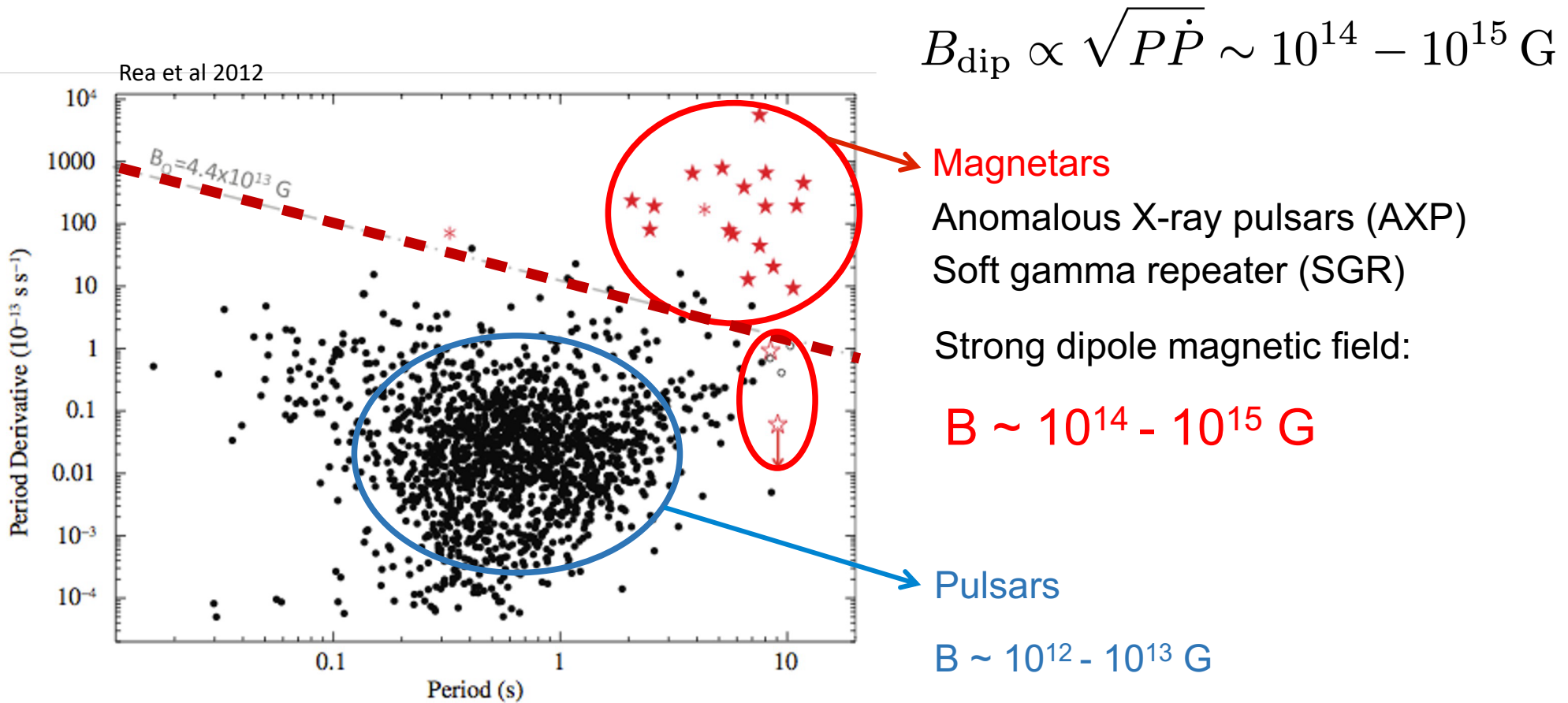
Paul Barrère

Andrei Igoshev

Raphaël Raynaud

Alexis Reboul-Salze

# Magnetars: the most intense known magnetic fields



# Galactic magnetars: observational constraints

30 magnetars in Galaxy and magellanic clouds

Magnetic dipole  $B_{\text{dip}} \sim 10^{14} - 10^{15} \text{ G}$

Indications of **stronger multipolar/toroidal field**

Young: typically  $10^4 - 10^5$  years

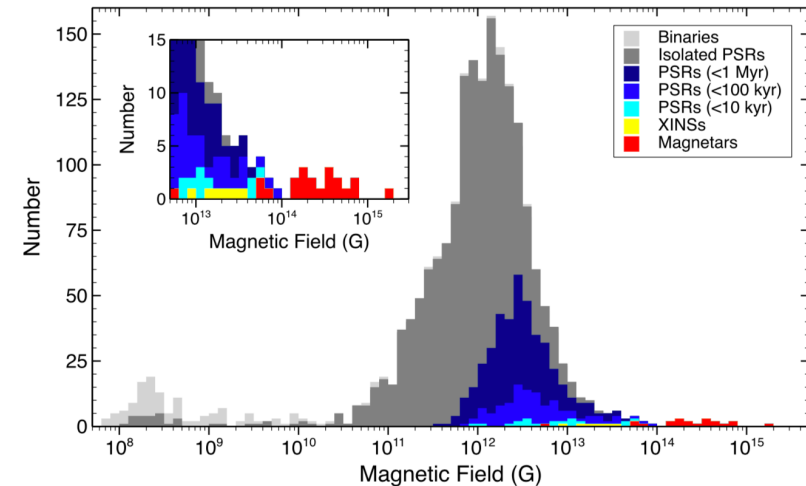
Birth rate :  $\sim 10\text{-}50$  % of core-collapse SNe/neutron star birth

Slow rotation  $P \sim 1 - 10 \text{ s}$

-> But rotation at birth unknown

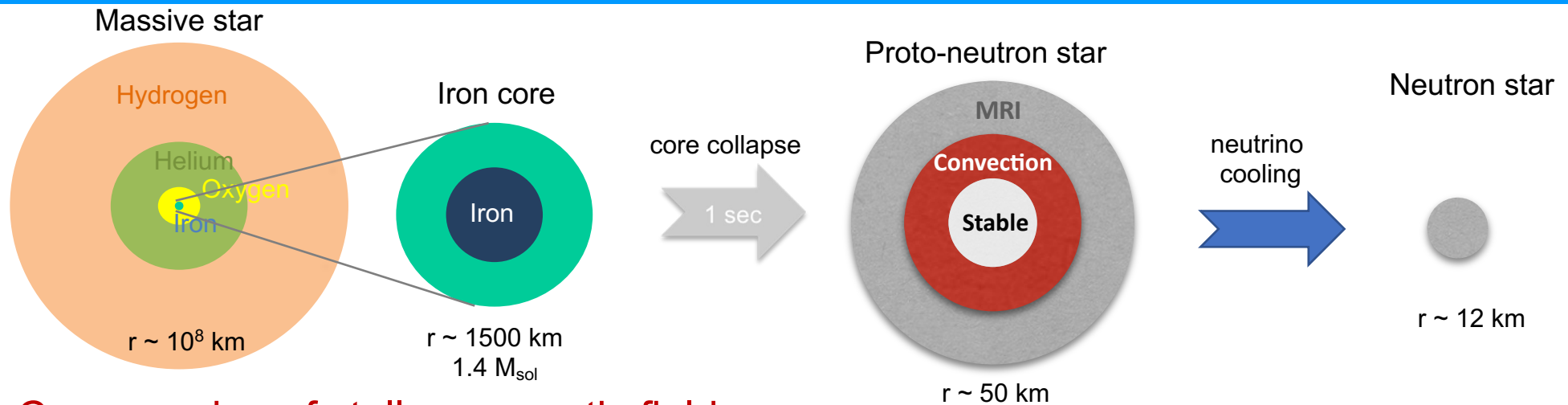
$\sim 10$  magnetars associated to a supernova remnants

=> Normal explosion kinetic energy =>  $P > 5 \text{ ms}$



Olausen & Kaspi 2014

# Different scenarios for magnetar formation



## Compression of stellar magnetic field :

Amplification by a few  $\sim 10^4$  during core collapse

Very magnetised stars on surface ( $B > 1 \text{ kG}$ ) : also need a  $10^{10}$ - $10^{11} \text{ G}$  in the iron core

## Protoneutron star dynamos

### Magnetorotational instability

Similar to accretion disks

Reboul-Salze+ 22, Guilet+ 22

### Convective dynamo

Similar to planetary  
& stellar dynamos

Raynaud+20, 22, Verma+26,  
Masada+22

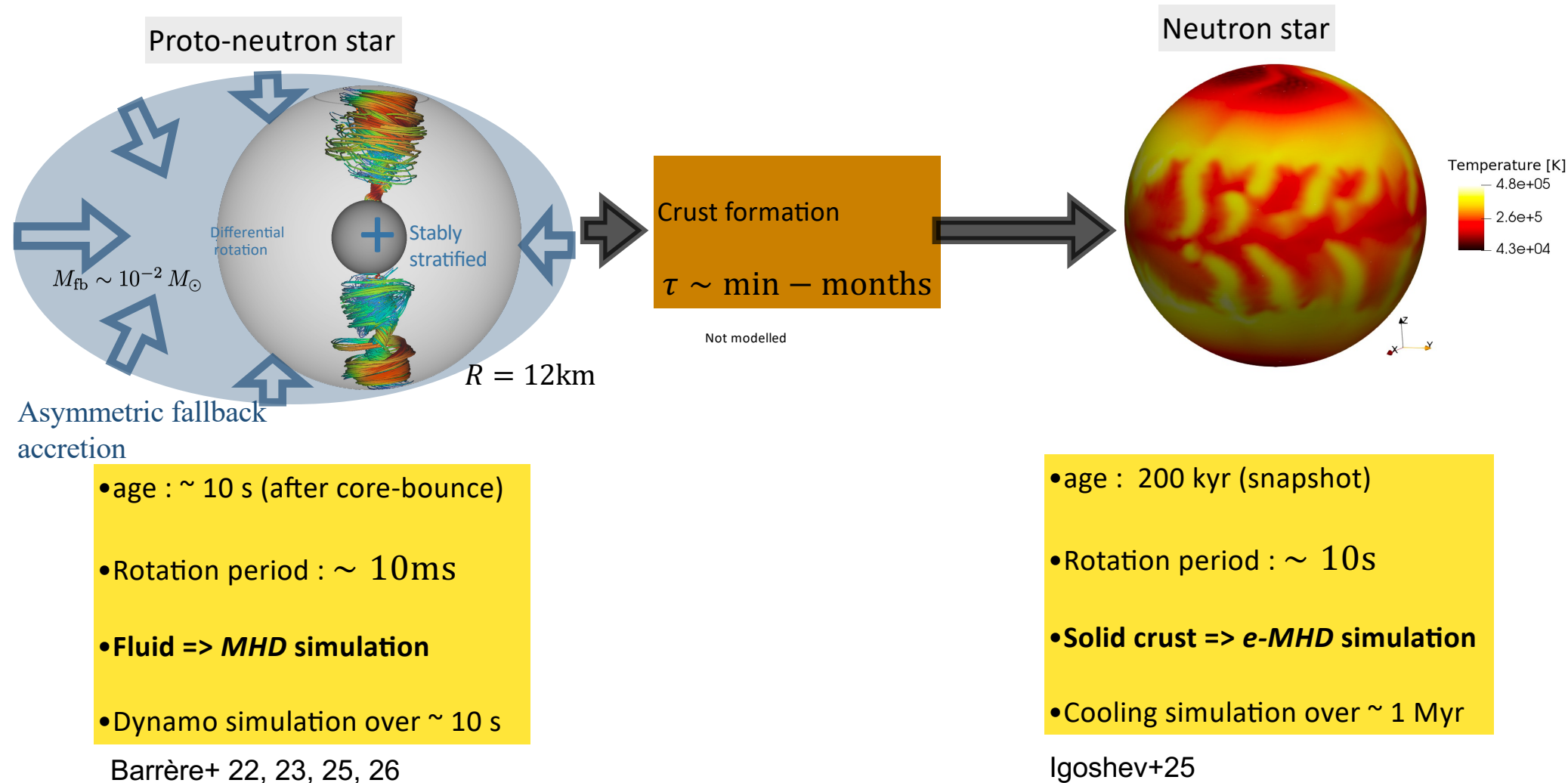
### Taylor-Spruit dynamo

Similar to stellar radiative  
zones

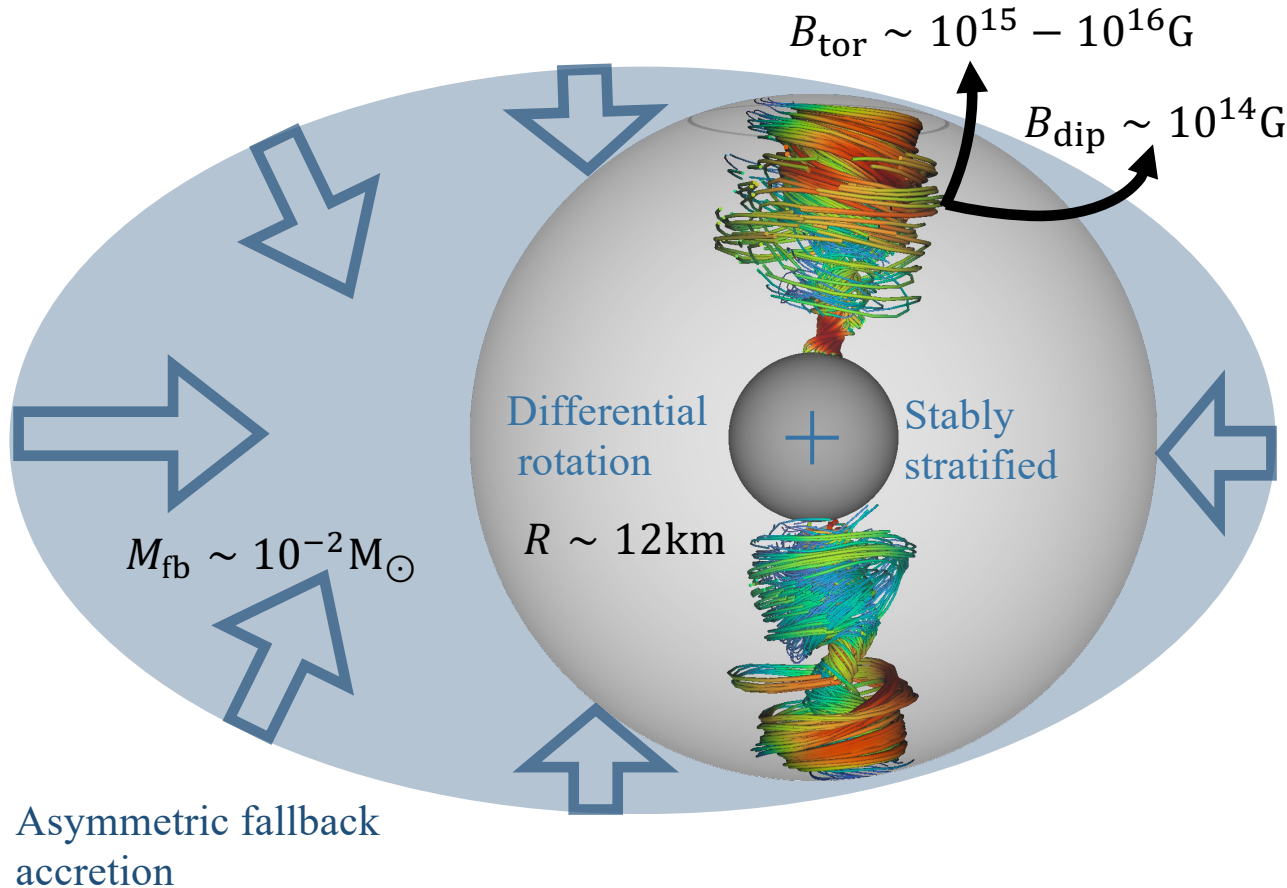
Barrère+ 22, 23, 25, 26

Chiral magnetic instability, Dehman & Pons 2024

# From proto-NS to NS simulations



# Tayler-Spruit dynamo: a new magnetar formation scenario

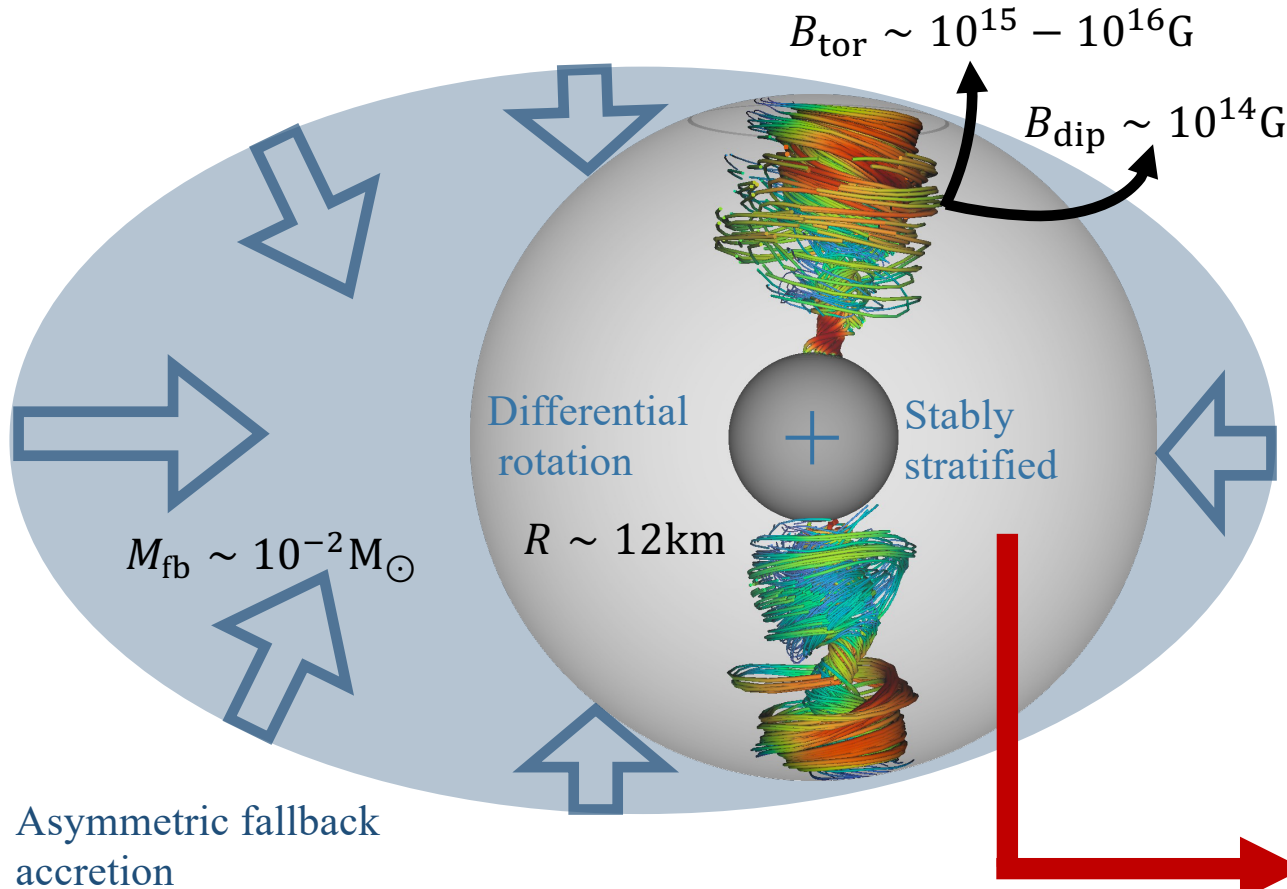


## Fallback in 3D CCSN simulations

- *Fallback* = matter that remains gravitationally bound during the explosion
- starts  $\sim 10\text{ s}$  after the core bounce
- $M_{\text{fallback}} \in [10^{-4}, 10^{-1}] M_{\odot}$
- lasts from minutes to hours
- Potential to spin up the proto-neutron star up to break up

(Chan+20, Janka+22)

# Tayler-Spruit dynamo: a new magnetar formation scenario



## Crucial for stellar physics

- **Angular momentum transport** in radiative zones
- **Only 2 analytical studies**
  - ➔ Spruit 02 & Fuller+19
  - ➔ different predictions for the saturation
- First numerical studies are **very new**
  - ➔ Petitdemange+23, Barrère+23

= “Tayler-Spruit” dynamo ?

# Modelling protoneutron star dynamo with the MagIC code



(Wicht+02, Gastine+12)

Pseudo-spectral numerical method in spherical geometry

Physical assumptions:

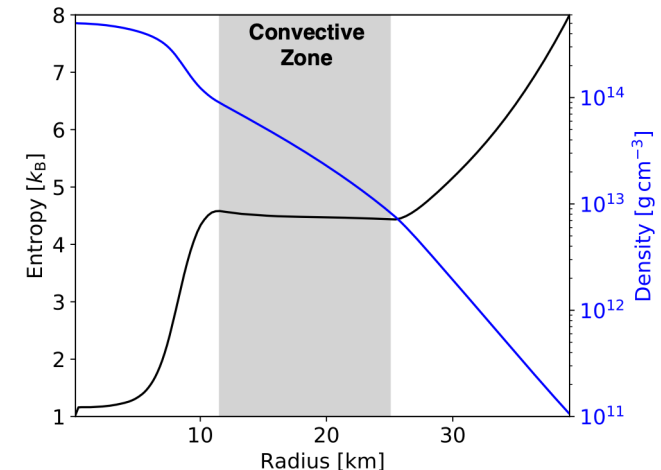
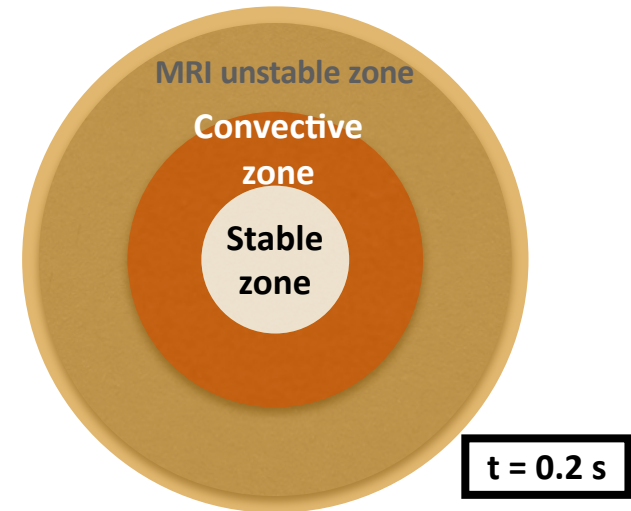
- Classical MHD equations, (pseudo-)Newtonian gravity
- Anelastic or Boussinesq approximation (sound proof)
- Ultra-high density equation of state

Ekman number :  $E = \frac{\nu}{\Omega d^2} = 10^{-5} - 10^{-3} > 10^{-4} - 10^{-8}$

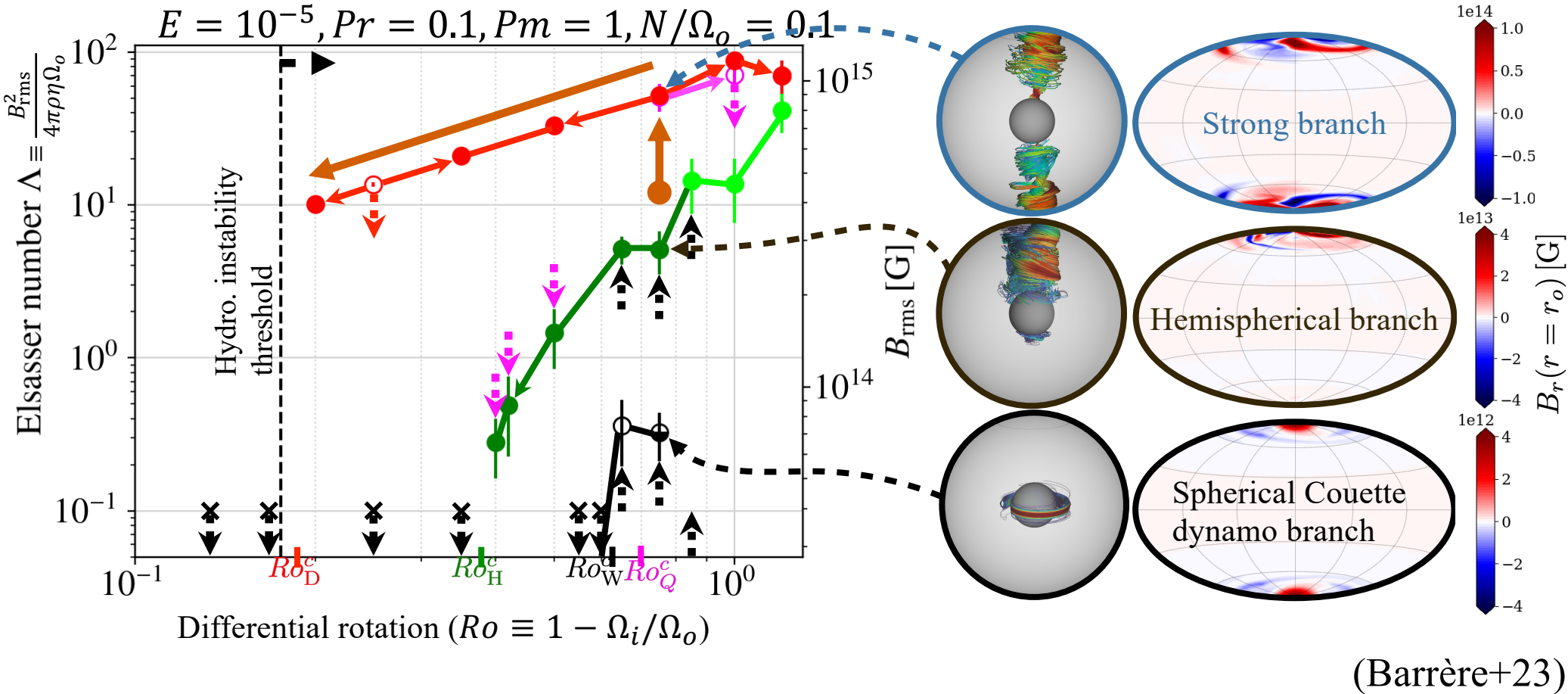
Rayleigh number :  $E = \frac{N^2 d^4}{\nu \kappa} = 10^6 - 10^7 < 10^8 - 10^{12}$

Magnetic Prandtl number :  $Pm = \nu/\eta \in [1, 16] \ll 10^{12}$

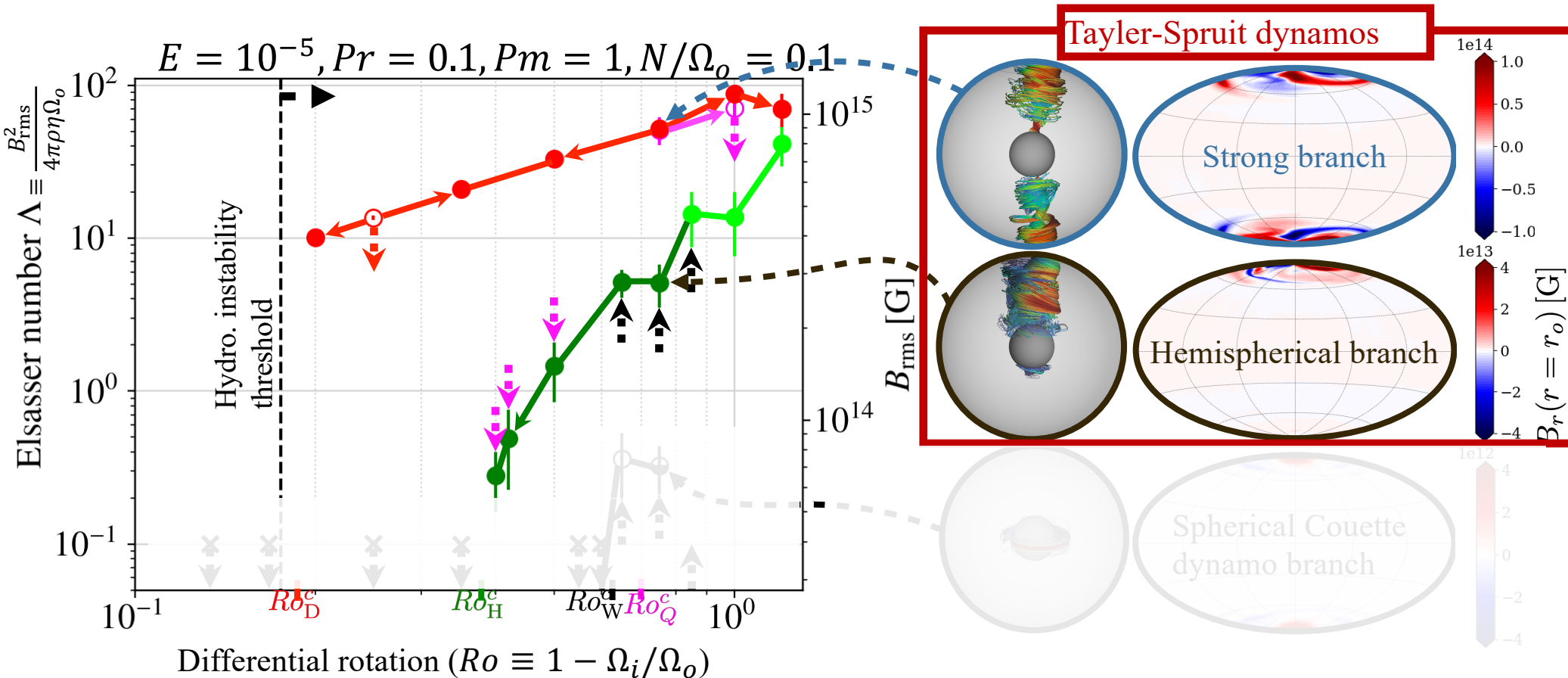
Protoneutron star structure



# Different dynamo branches

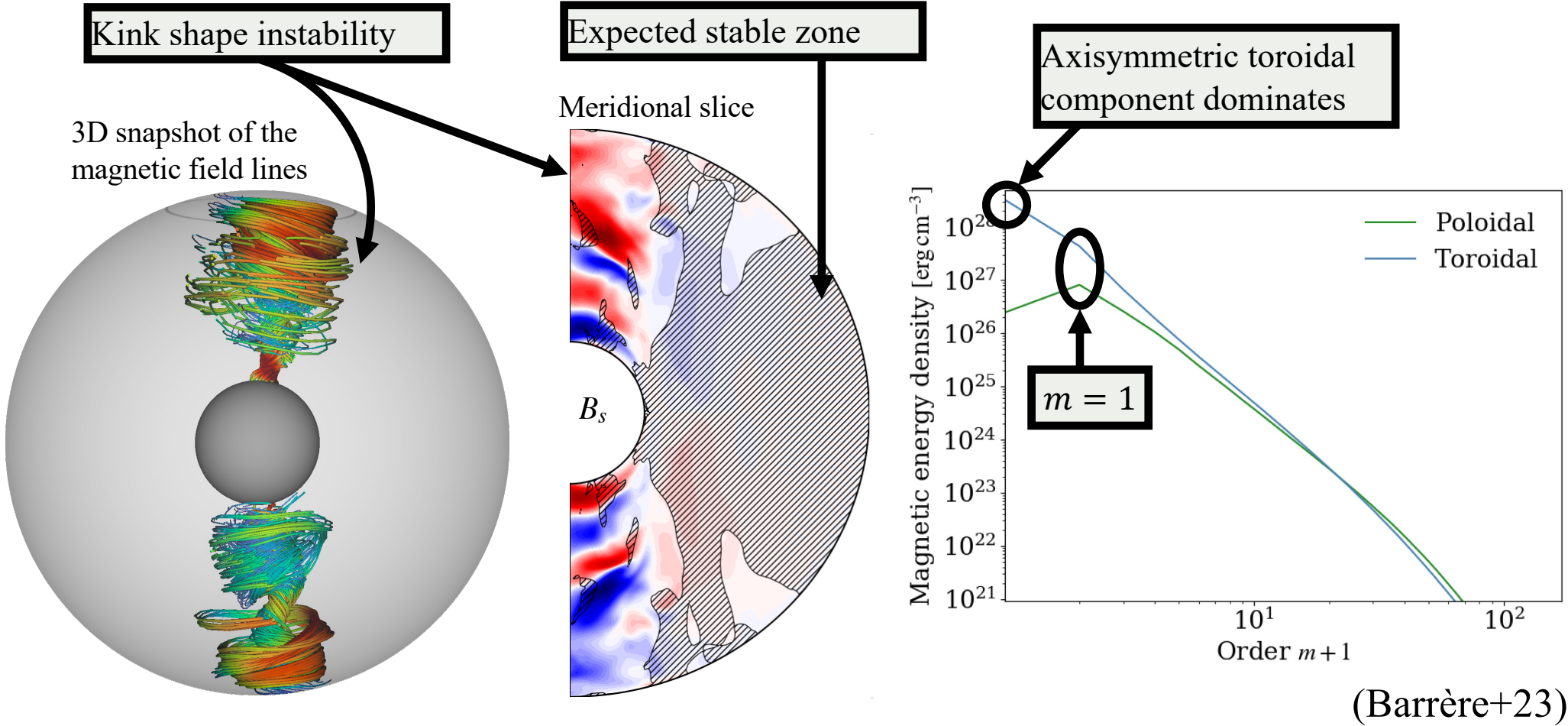


# Different dynamo branches



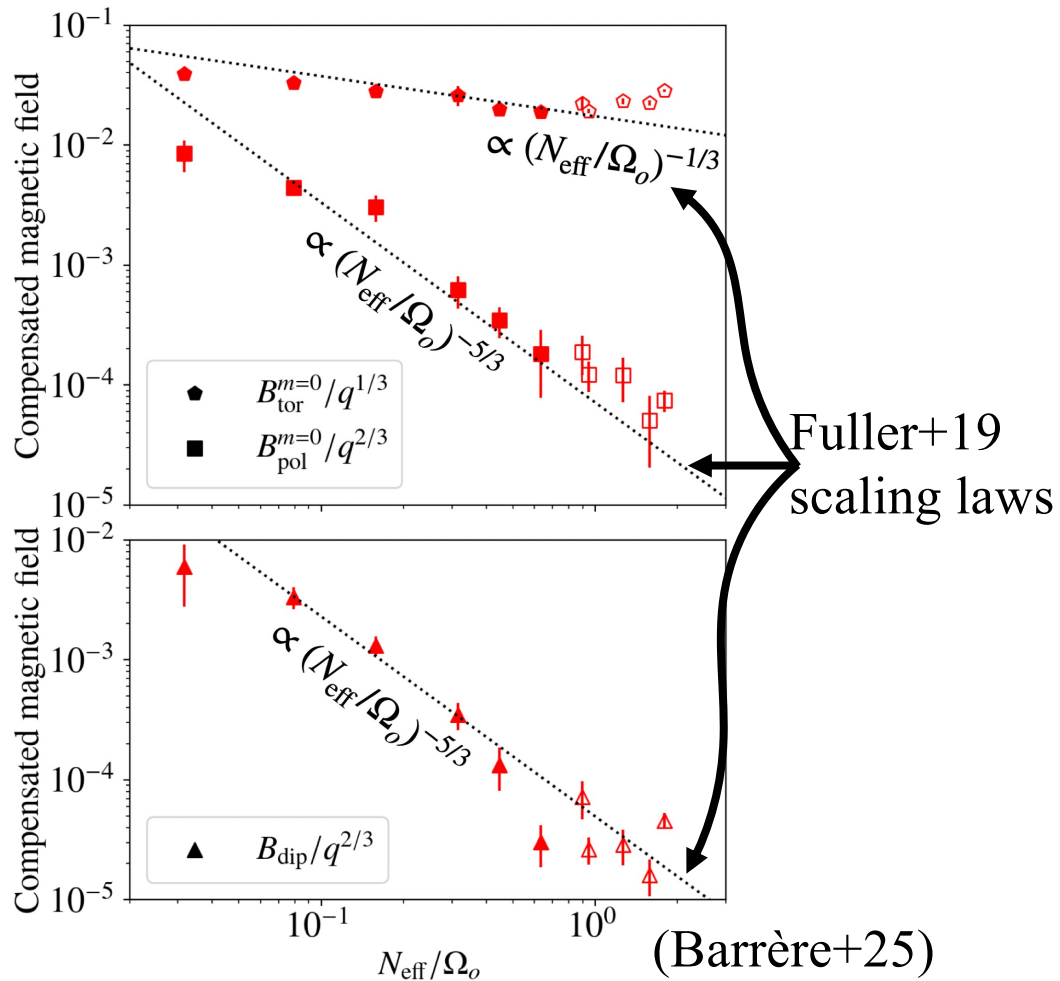
(Barrère+23)

# Magnetic field structure



(Barrère+23)

# Scaling laws



## Parametrisation

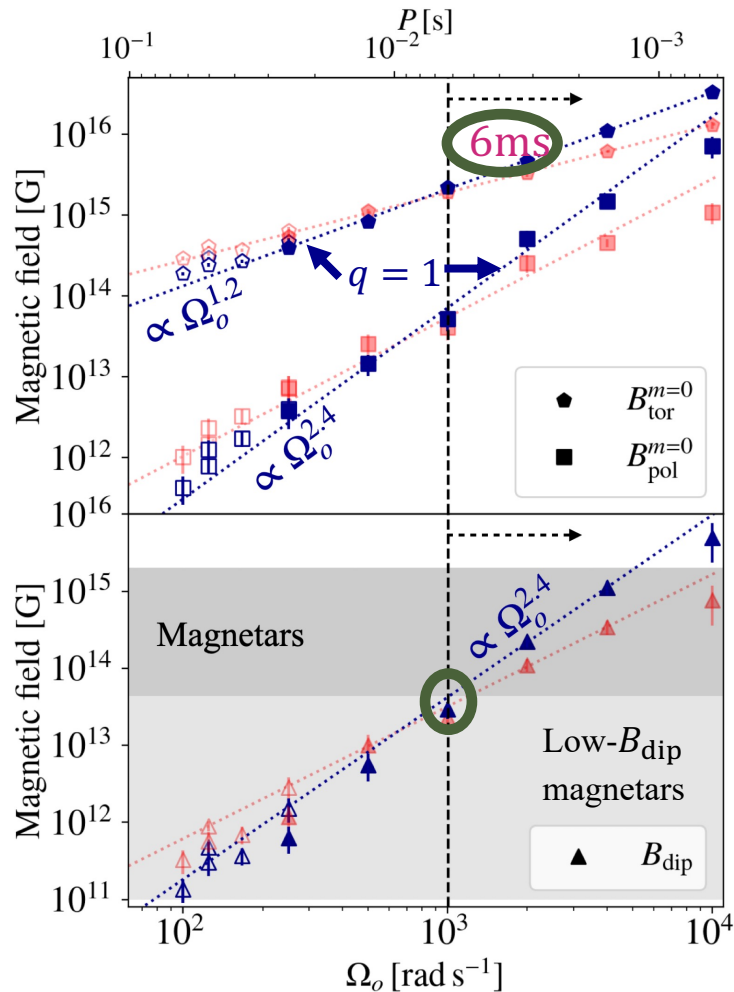
Dimensionless normalisation factor

$$B_{\text{tor}}^{m=0} = \alpha \sqrt{4\pi\rho r_o^2 \Omega_o} \left(\frac{q\Omega_o}{N_{\text{eff}}}\right)^{1/3}$$

$\alpha \sim 0.01$   
in our simulations

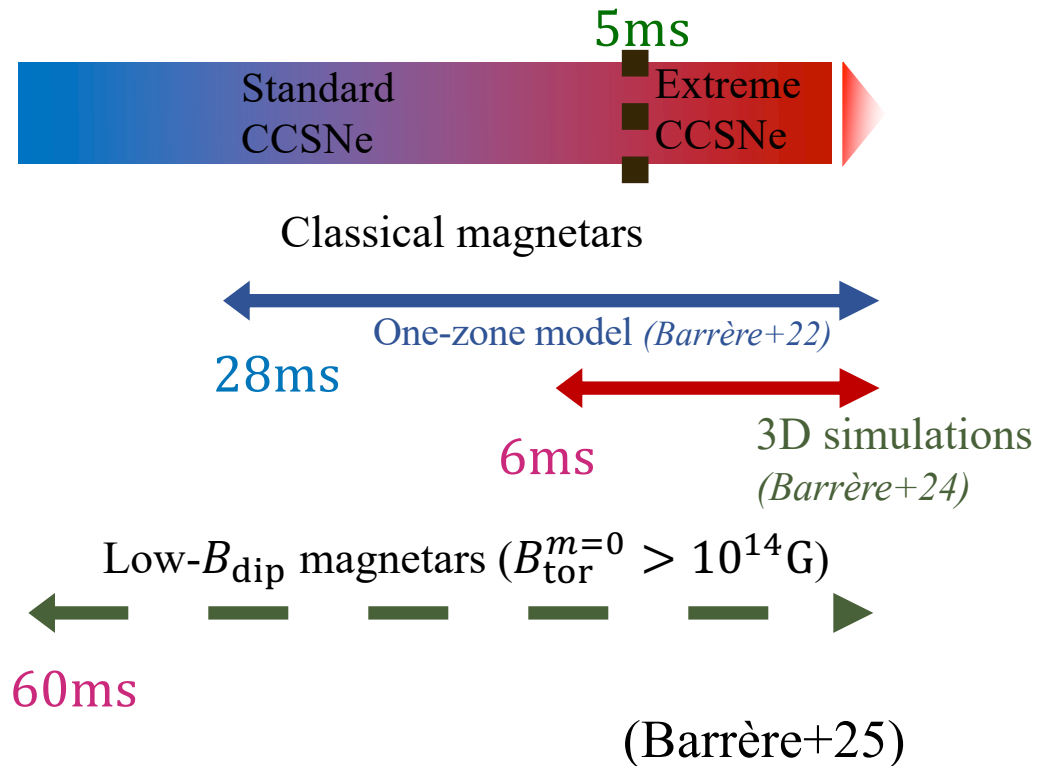
$\alpha \sim 0.25 - 1$   
Calibrated on asteroseismology  
in sub-/red giants  
(Fuller+19,+23,Eggenberger+19)

# Application to magnetar formation



Impose  $N = 1000\text{s}^{-1}$

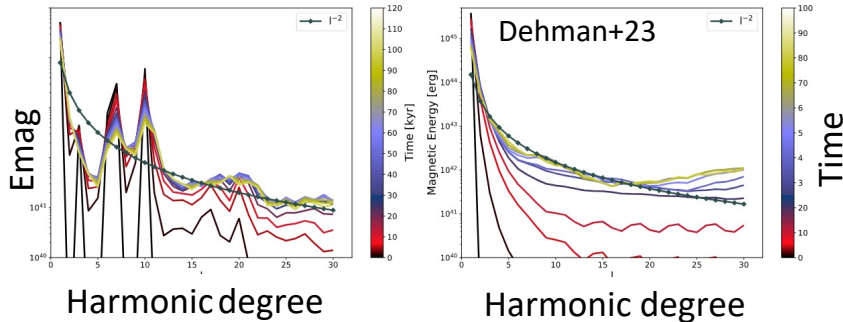
Rotation period  $P$



# Long term evolution: magneto-thermal evolution

## Neutron star crust

- ions are fixed ; currents generated by the flow of electrons
- Ohmic dissipation  $\tau_\eta = \frac{R^2}{\eta}$
- Hall drift  $\tau_H = \frac{4\pi n_0 e R^2}{c B_0}$
- $Ha = \tau_\eta / \tau_H \gg 1$
- Hall cascade partially keep B field structure



## Magnetic induction

Grandis+2020

$$\frac{\partial \vec{B}}{\partial t} = \underbrace{Se \nabla \left( \frac{1}{\mu} \right) \times \nabla T^2}_{\ll 1} + \underbrace{Ha \nabla \times \left[ \frac{1}{\mu^3} \vec{B} \times (\nabla \times \vec{B}) \right]}_{\text{Hall term}} - \underbrace{\nabla \times \left[ \frac{1}{\mu^2} \nabla \times \vec{B} \right]}_{\text{Diffusion term}}$$

## Thermal diffusion

$$\frac{1}{Ro} \frac{C_v}{T} \frac{\partial T^2}{\partial t} = \underbrace{\nabla \cdot (\mu^2 \hat{\chi}(B) \cdot \nabla T^2)}_{\text{Heat flux}} + \underbrace{\frac{Pe}{Se} \frac{|\nabla \times \vec{B}|^2}{\mu^2}}_{\text{Ohmic heating}} + \underbrace{Pe \mu (\nabla \times \vec{B}) \cdot \nabla \left( \frac{T^2}{\mu^2} \right)}_{\ll 1}$$

$$\text{Anisotropic heat transport } \hat{\chi}(B) = \frac{\delta_{ij} + Ha B_i B_j / \mu^2 - Ha \epsilon_{ijk} B_k / \mu}{1 + Ha^2 |B|^2 / \mu^2},$$

> reduced heat flow in the direction orthogonal to B field lines

# Implementing a PNS dynamo state as an initial condition

Taylor-Spruit dynamo

Neutron star crust

$$\frac{r_i}{r_o} = 0.25$$

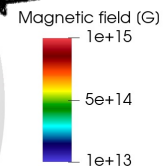
$$\frac{r_i}{r_o} = 0.9$$

## Method

- Assume a polynomial expansion for the poloidal and toroidal potentials  $b_{lm}^{\text{crust}}(r)$
- Determine the coefficients by matching the boundary conditions and imposing  $b_{lm}^{\text{crust}}(r^*) = b_{lm}^{\text{PNS}}(r^*)$ ,  $\forall l \leq 30$  at  $r^* = [0.93r_o, 0.96r_o, r_o]$

## Outer boundary

B : insulating b. c.  
T : blackbody emission



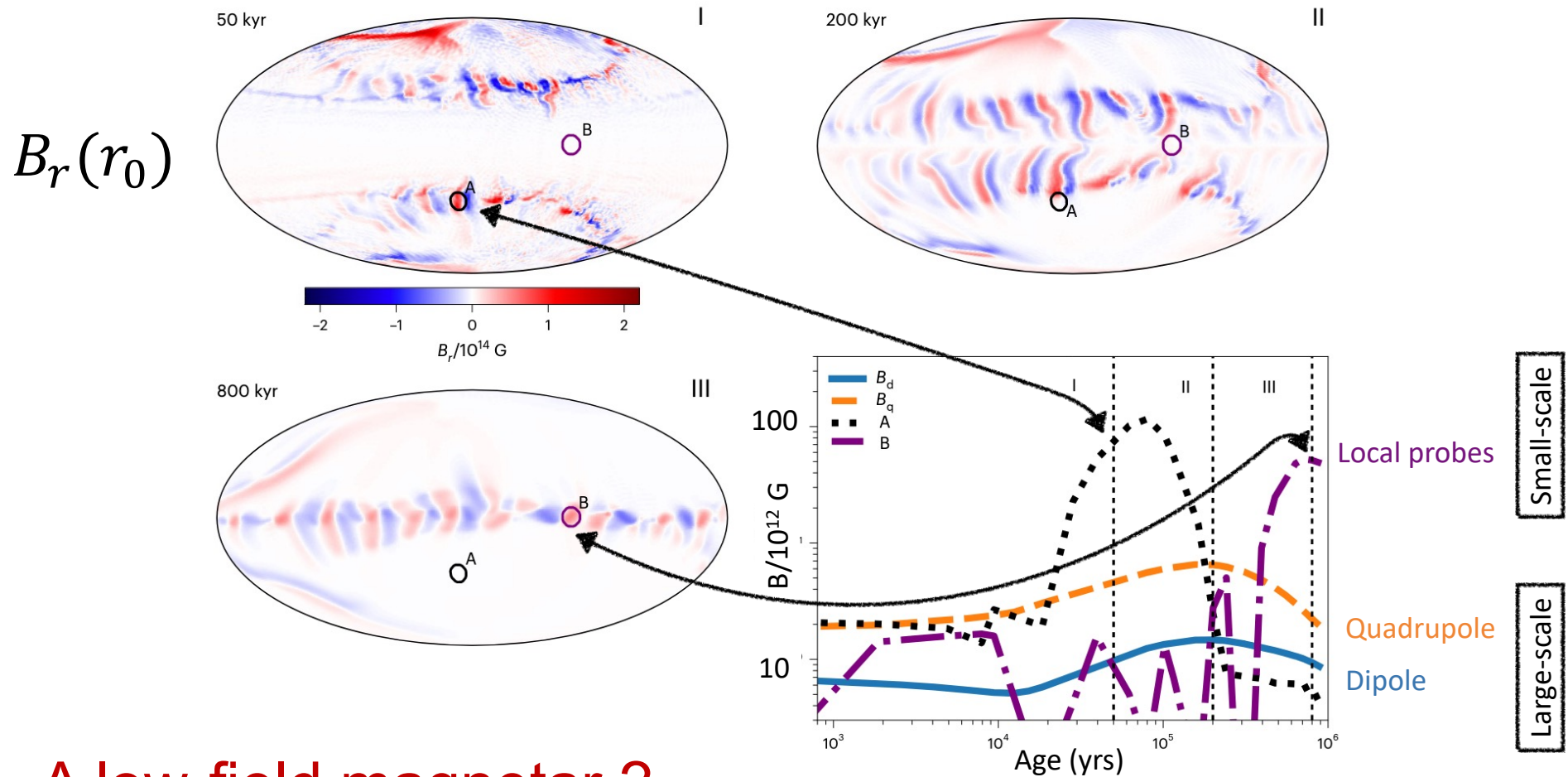
$d = 1\text{km}$

## Inner boundary

B : Meissner condition  
T : no flux  $\partial_r T = 0$

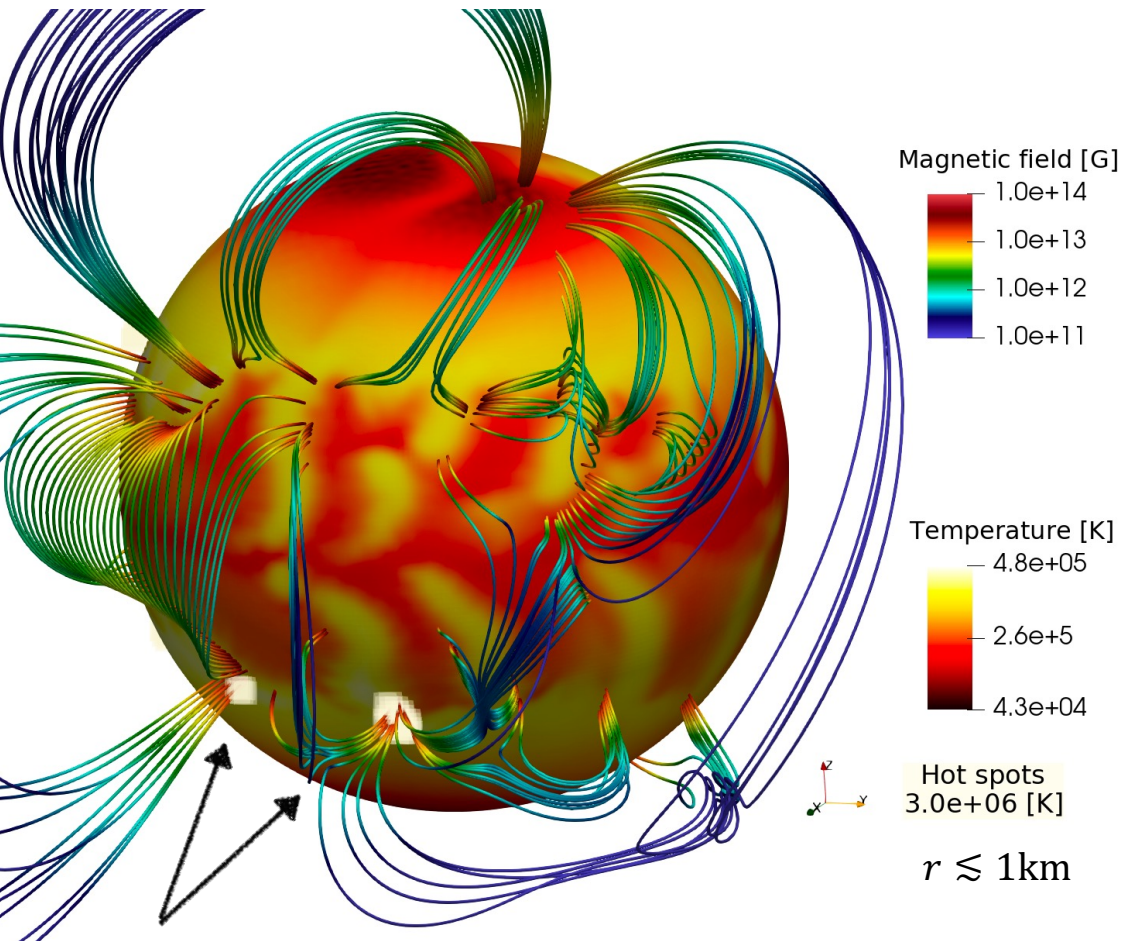
Simulation from  
Barrère+ 23

# Magnetic field long term evolution



A low-field magnetar ?

# External magnetic field at 200 kyr



B field creates large T variations

Footpoints with  $B_r > 10^{13} \text{ G}$  could produce hot spots heated by magnetospheric currents  
 $\Rightarrow$  X-ray light curve modulation

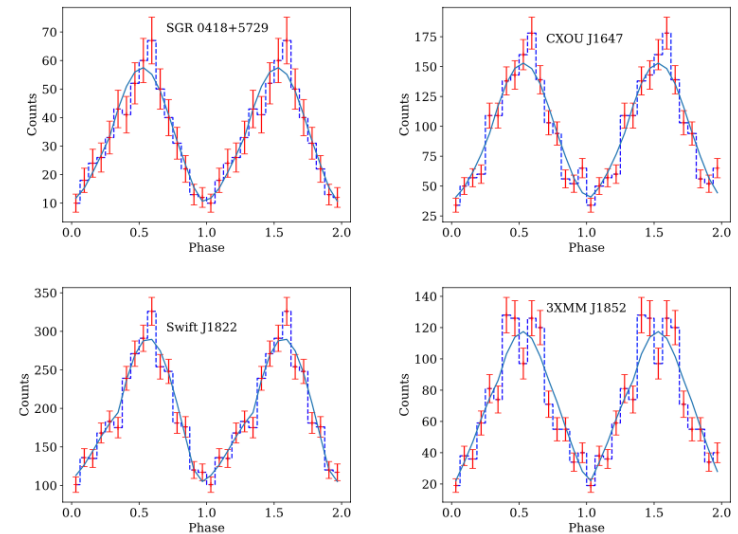
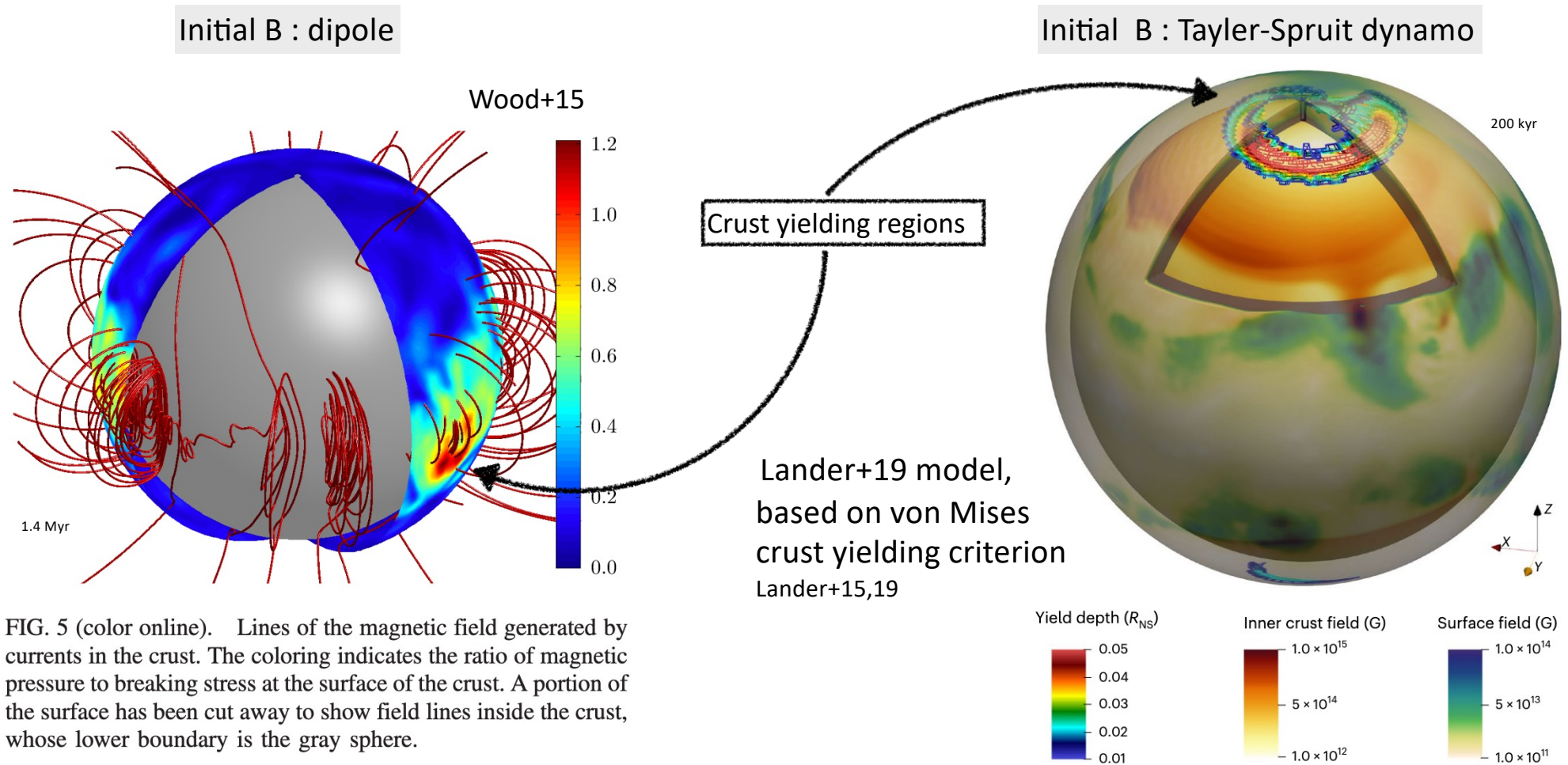


Figure 10: Observed soft X-ray lightcurves in the range of 0.3-2 keV for low-field magnetars in the quiescent state and the best fits (solid blue line). It is assumed that emission is produced by hot spots formed at the places with radial magnetic field exceeding  $7 \times 10^{13} \text{ G}$ .

# Crustal magnetic stresses

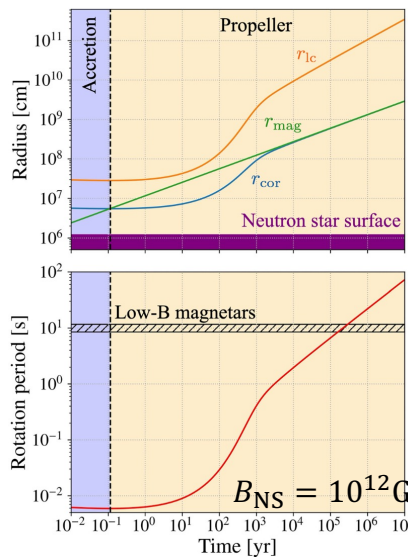


# Spin down

Taylor-Spruit dynamo invoke CCSN fallback  
 => rotational evolution controlled by interaction  
 with a fallback disc ( $M \sim 0.01M_{\odot}$ )

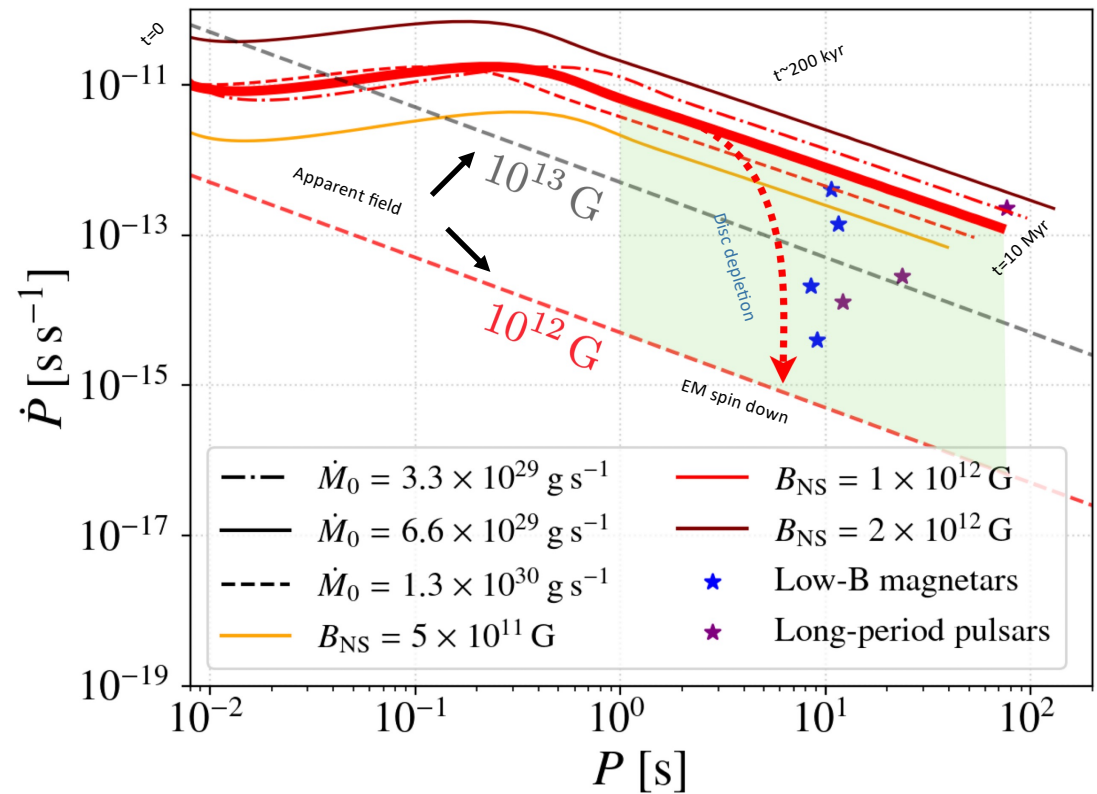
$$I_{\text{NS}} \dot{\Omega}_{\text{NS}} = N_{\text{acc}} + N_{\text{dip}} \quad (\text{Ronchi+22})$$

Accretion torque      Electromagnetic torque



Propeller regime  
 => NS spin down

$$P(t \sim 200\text{kyr}) \approx 10\text{s}$$



# Conclusions

From Tayler Spruit dynamo...

Classical magnetar possible for fast rotation ( $P < 6$  ms)  
=> tension with observational constraints

Strong azimuthal B field: could explain low-field magnetar

3D geometry : concentrated near pole

Scaling consistent with theoretical predictions (Fuller+19)  
but with a small normalisation factor

## Perspectives

- Realistic PNS structure
- Large magnetic Prandl number (Pm)

...to long term evolution

First proof of concept

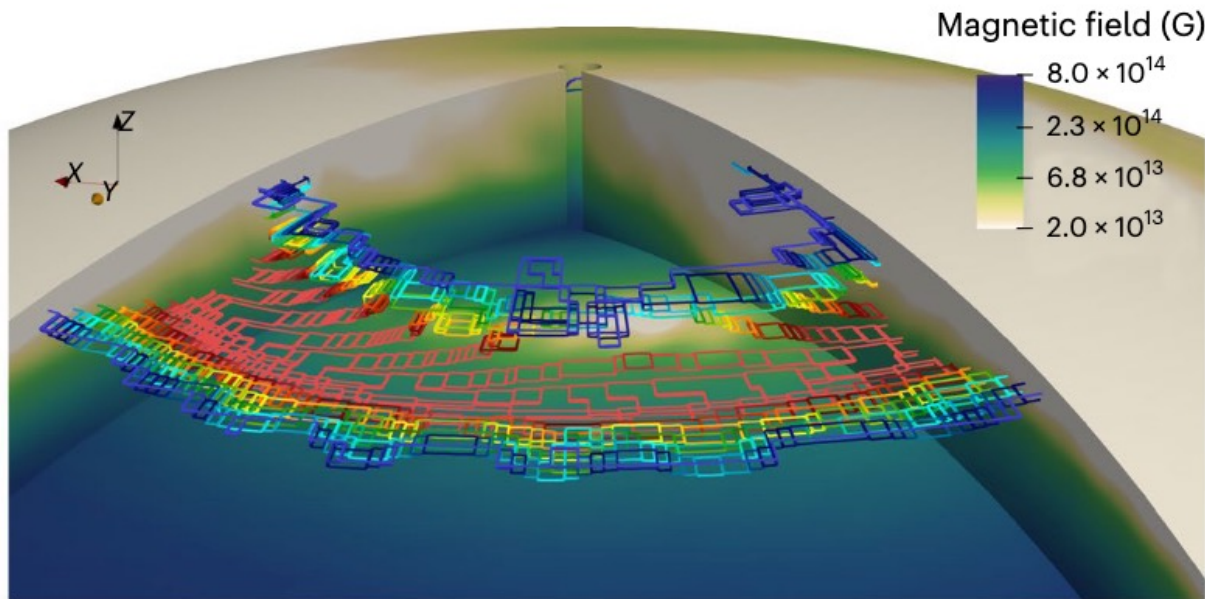
- Consistent with low-field magnetar:
- low dipole but locally strong B field
  - crust yield
  - slow rotation reproduced

## Perspectives

- Different dynamo initial conditions
- Better mapping from dynamo to crust : magnetic relaxation
- Below the crust : global models including neutron star core (Igoshev+23)  
Talk by Nicolas Moraga on Friday

Questions ?

# Powering magnetar X-ray activity



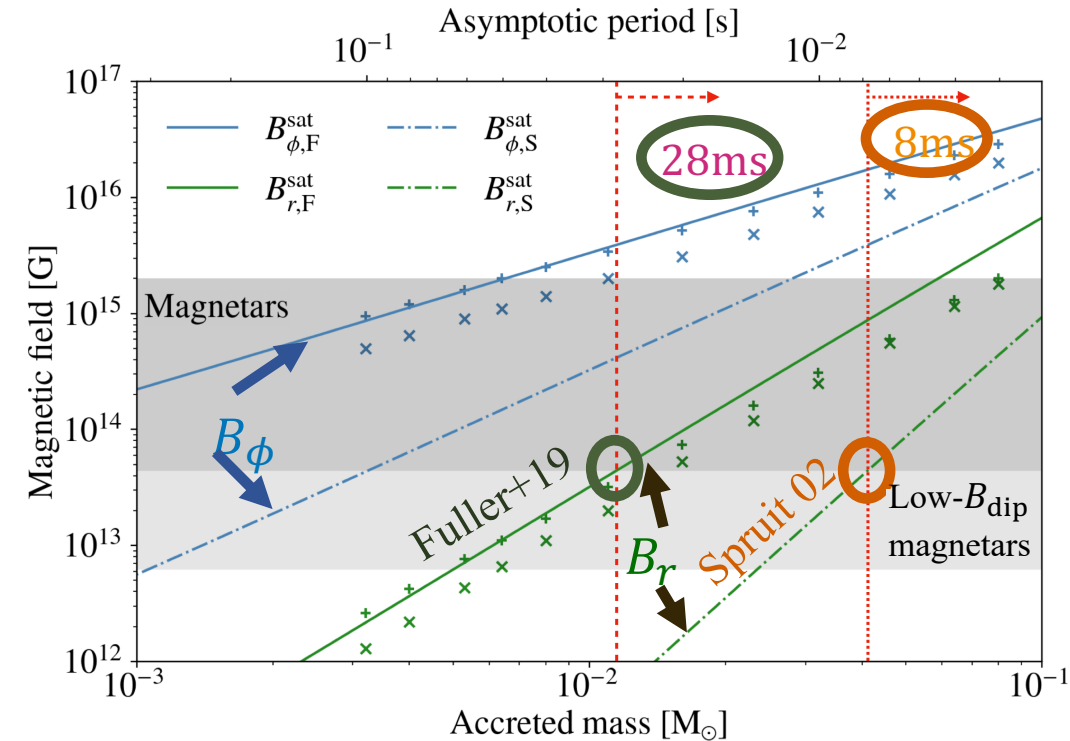
**Energy release in crustal failure (Thompson+95)**

$$E_{\text{out}} \lesssim 2 \times 10^{39} \text{erg} \left( \frac{l}{1 \text{km}} \right)^2 \left( \frac{|B|}{2 \times 10^{14} \text{G}} \right)^2$$

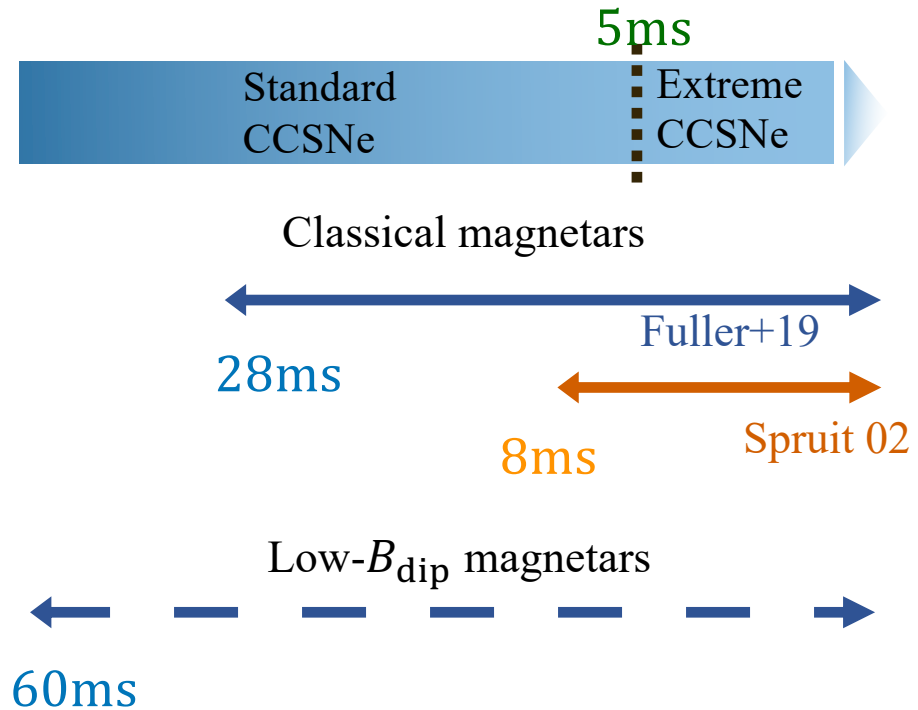
Low-field magnetar observations

$$E_{\text{burst}} \sim 10^{37} \text{erg} \quad (\text{Muno+07})$$

# Proof of concept: one-zone model



Rotation period  $P$



(Barrère+22)

# Impact of stratification on Tayler modes

