

CONVEXITY OF THE EQUATION OF STATE IN MERGER SIMULATIONS AND GRAVITATIONAL WAVES

G. RIVIECCIO, A. RIOS, P. CERDÀ-DURÀN, M. RUIZ, J. A. FONT

(UV)

(UB)

(UV)

(UV)

(UV)

✉ giuseppe.rivieccio@uv.es

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Focus of the talk

GR, Guerra, Ruiz & Font (2024)

Gravitational wave imprints of non-convex dynamics in binary neutron star mergers

Phys. Rev. D 109, 064032



GR, Rios, Cerdà-Durán, Ruiz & Font (2026)

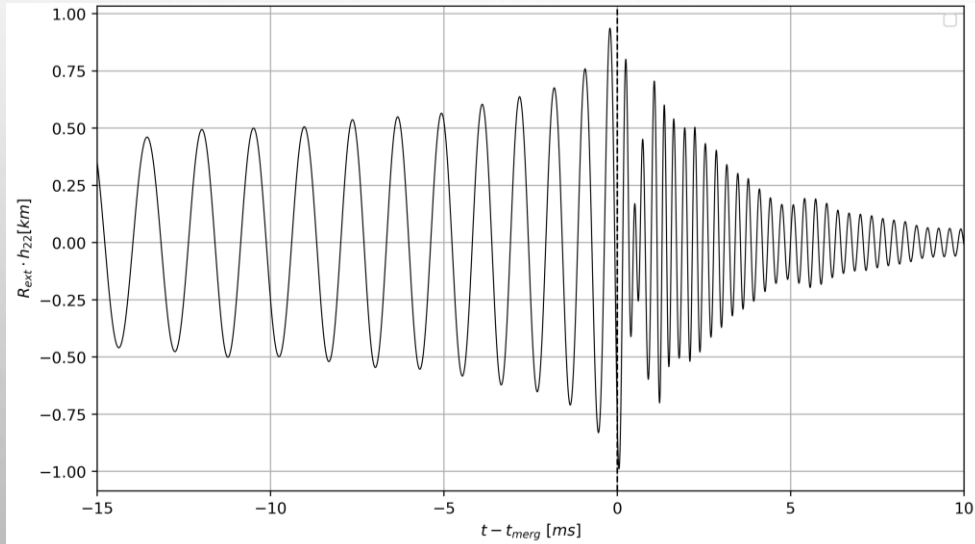
Convexity of parametrized phase transition and gravitational wave signature

Still loading

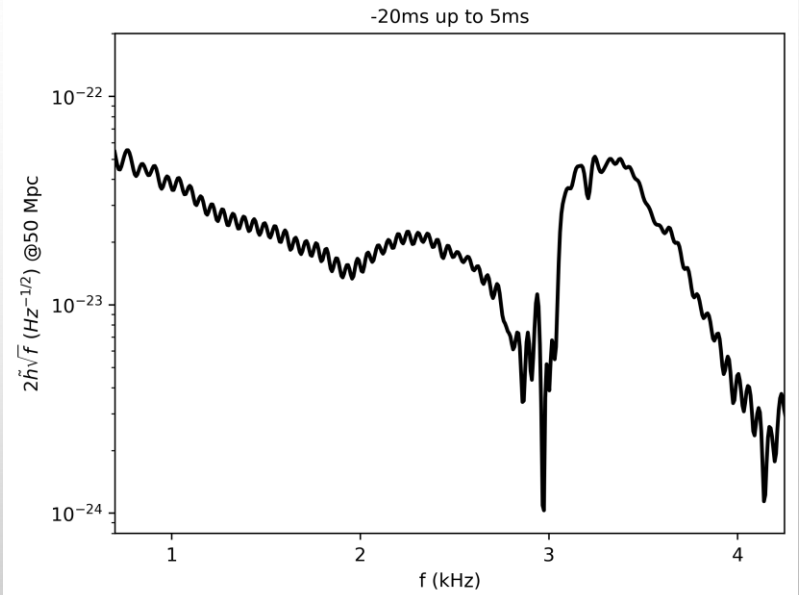


BINARY NEUTRON STARS (BNS) MERGER

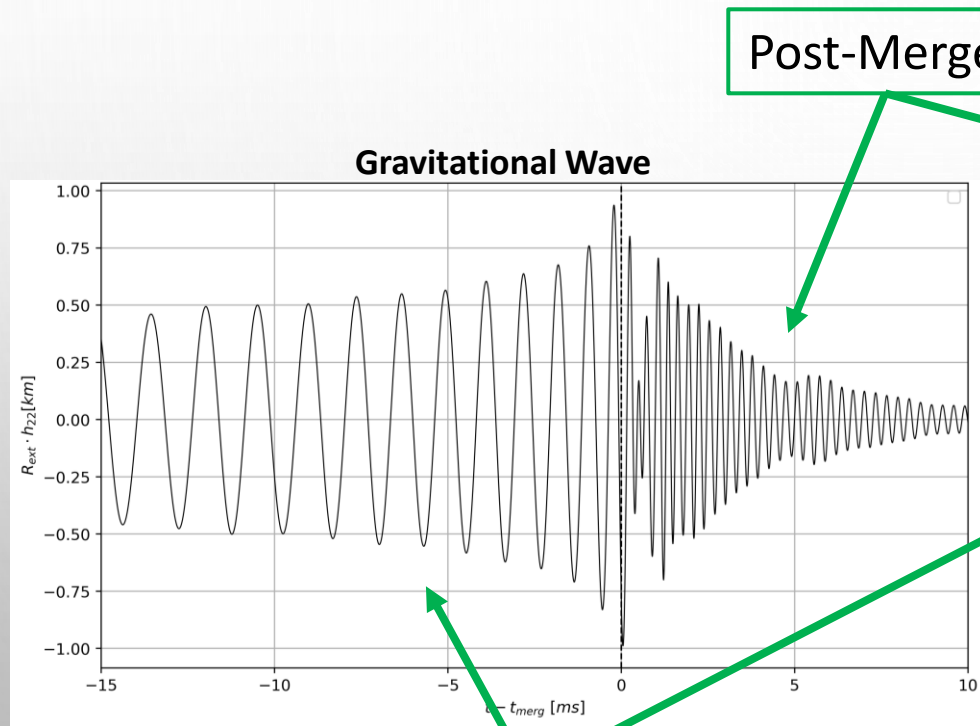
Gravitational Wave



Fourier Transform of Gravitational Wave



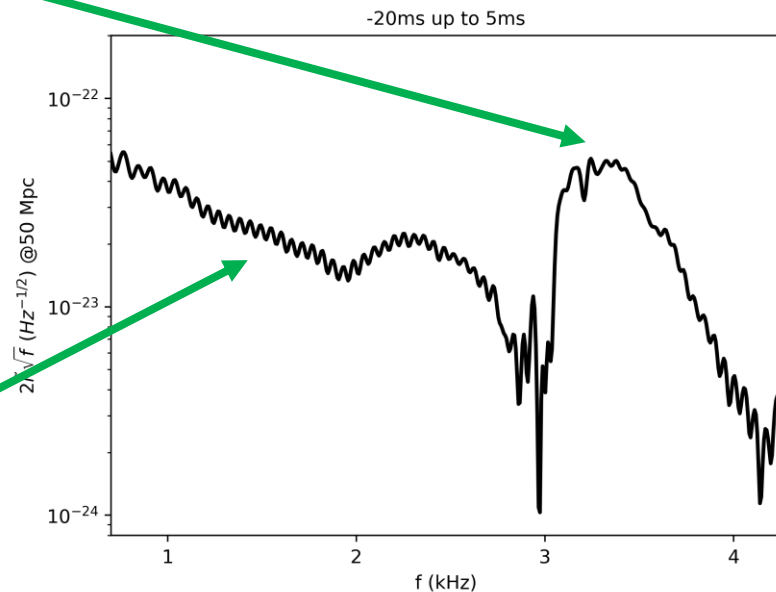
BINARY NEUTRON STARS (BNS) MERGER



Inspiral

Post-Merger

Fourier Transform of Gravitational Wave

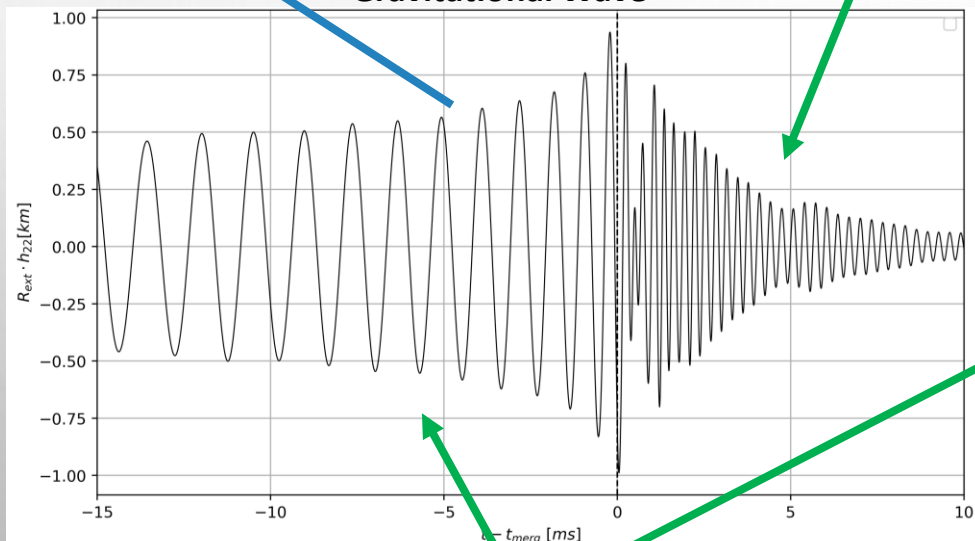


BINARY NEUTRON STARS (BNS) MERGER

Tidal deformability

Λ

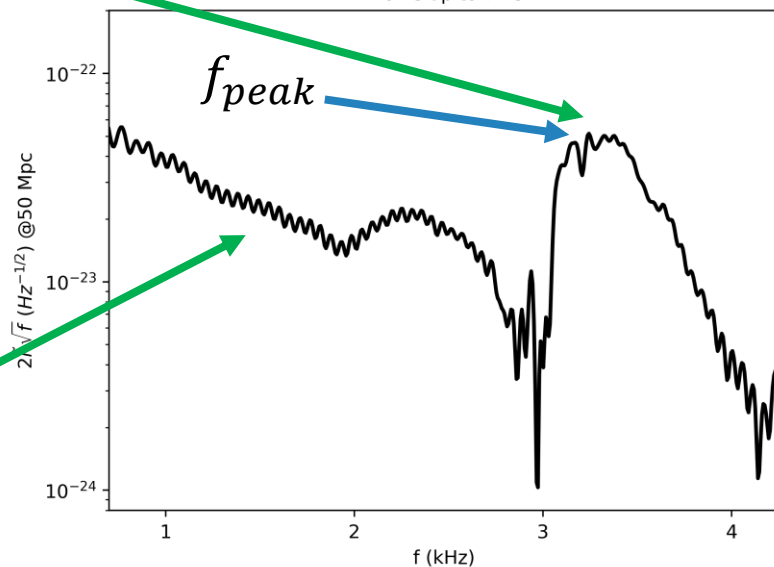
Gravitational Wave



Post-Merger

Fourier Transform of Gravitational Wave

-20ms up to 5ms



Inspiral

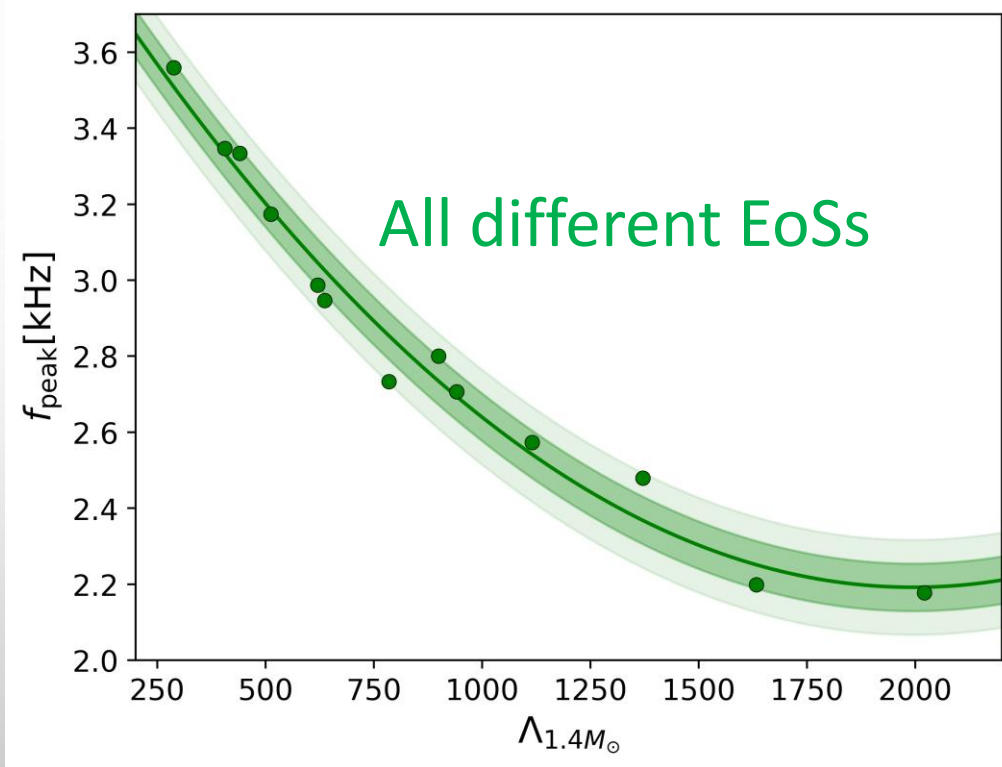
Quasi-Universal Relation

Λ := Tidal deformability
(Measured pre-merger)

f_{peak} := Characteristic Frequency
post-merger

We can find a relation that is
EoS-independent (almost)

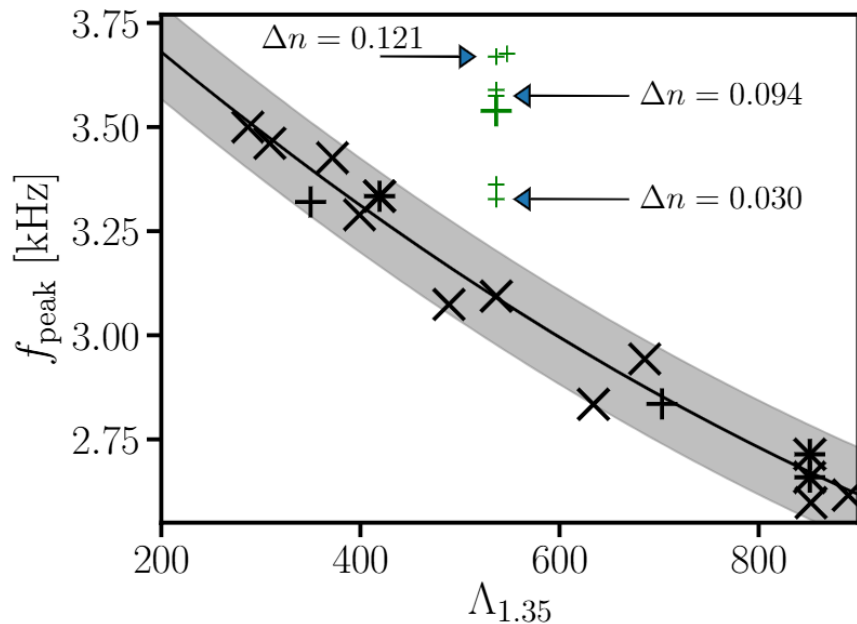
Useful to constrain the EoS



Quasi-Universal Relation including Phase Transition

First order Phase transition softens the EoS during the merger

Shift in f_{peak} is informative of a possible phase transition

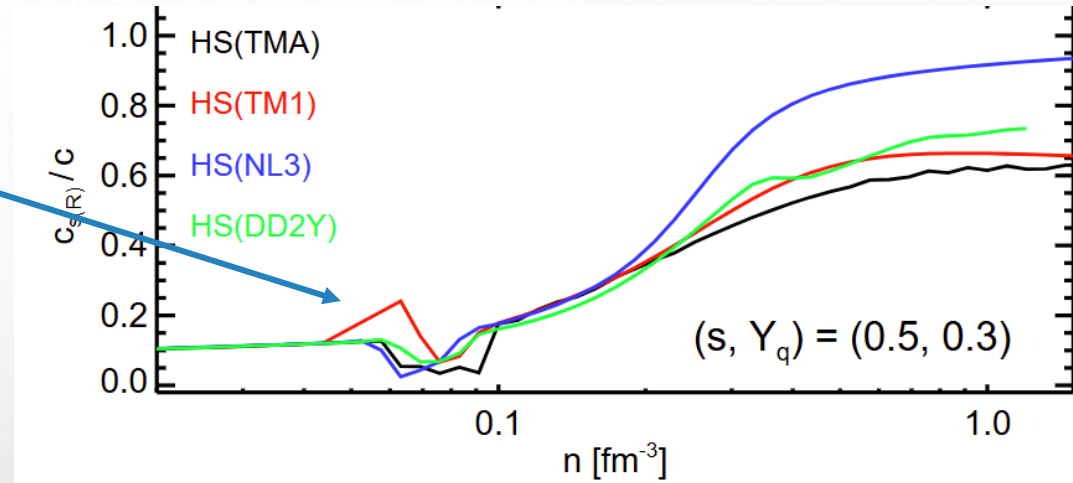


Bauswein et al. 2019
Blackmer et al. 2023,2025

Sound speed is non-monotonic

$$c_s^2 = \left. \frac{\partial P}{\partial \epsilon} \right|_s \text{ is non-monotonic}$$

Because of the drop
of the sound speed:



Aloy et al. (2019)

Phase Transition are in general
Nonconvex

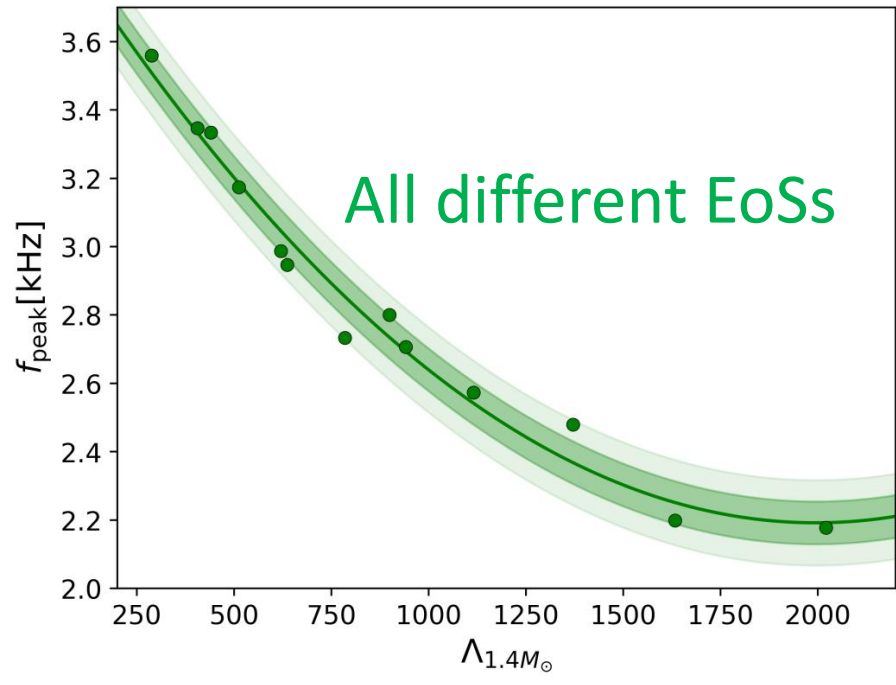
R Menikoff and B. Plohr Rev. Mod. Phys. 61, 75

Quasi-Universal Relation

Λ := Tidal deformability
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f_{peak} := Characteristic Frequency
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EoS-independent (almost)



G. Riveccio et al. (2024) Phys. Rev. D **109**, 064032

How does nonconvexity affect this relation?

Wave structure - Fundamental Derivative

$$G \equiv -\frac{1}{2} V \frac{\frac{\partial^2 P}{\partial V^2} \Big|_s}{\frac{\partial P}{\partial V} \Big|_s} \quad \text{Fundamental Derivative}$$

Depends only on EoS

Menikoff and Plohr (1986), Ibañez et al. (2013)

Wave structure - Fundamental Derivative

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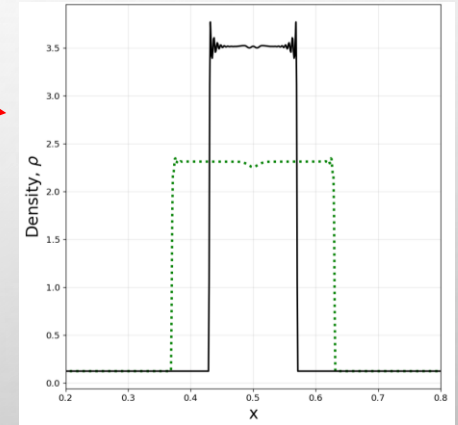
Depends only on EoS

Menikoff and Plohr (1986), Ibañez et al. (2013)

$G > 0$ convex

Standard wave propagation

Colliding Slabs

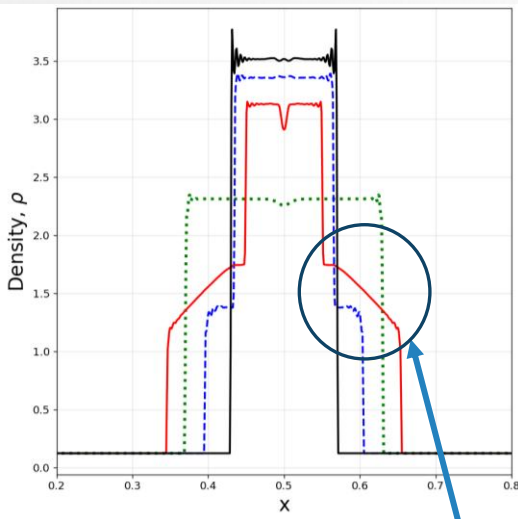


Wave structure - Fundamental Derivative

Depends only on EoS

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Menikoff and Plohr (1986), Ibañez et al. (2013)



Rarefaction waves
(Lower density but entropy increase)

$G > 0$ convex

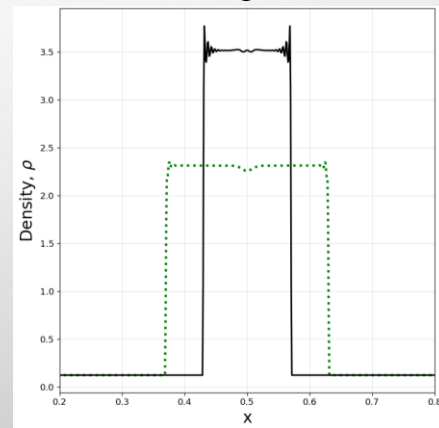
Standard wave propagation

$G < 0$ nonconvex!

Composite waves
Rarefaction wave attached to a shock front

Split waves
(e.g. Phase Transition)

Colliding Slabs



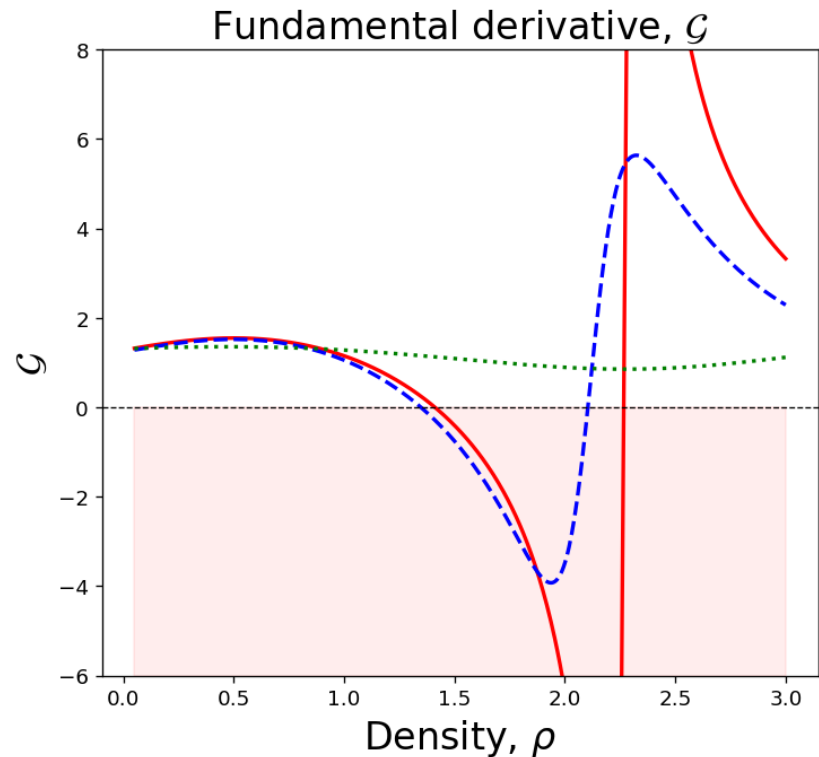
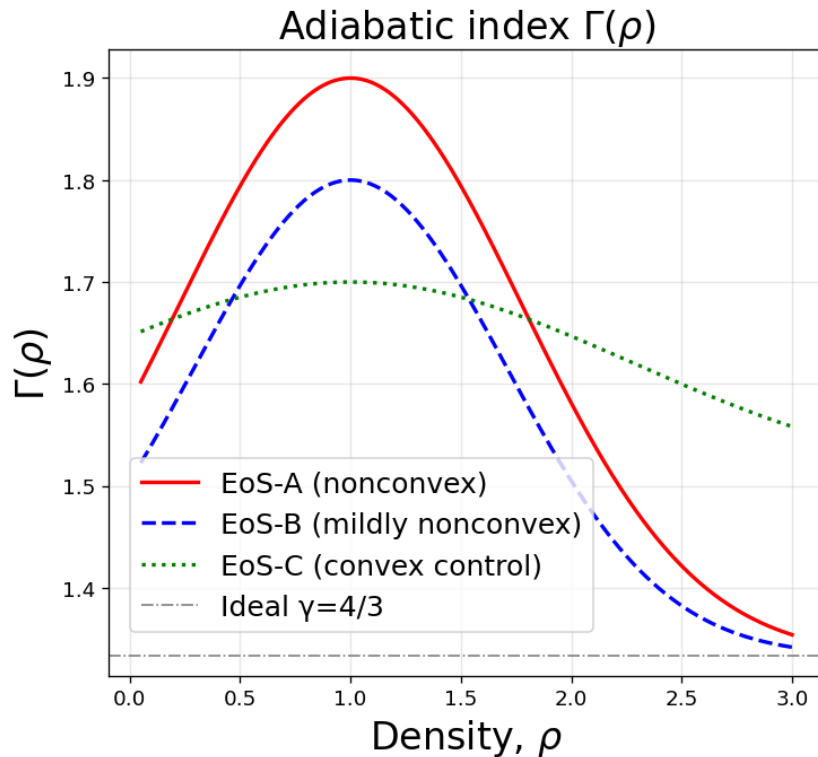
ASSUMING A Γ -LAW EOS

$$P = (\Gamma - 1) \rho \epsilon$$

Ideal gas : $\Gamma = \text{const}$

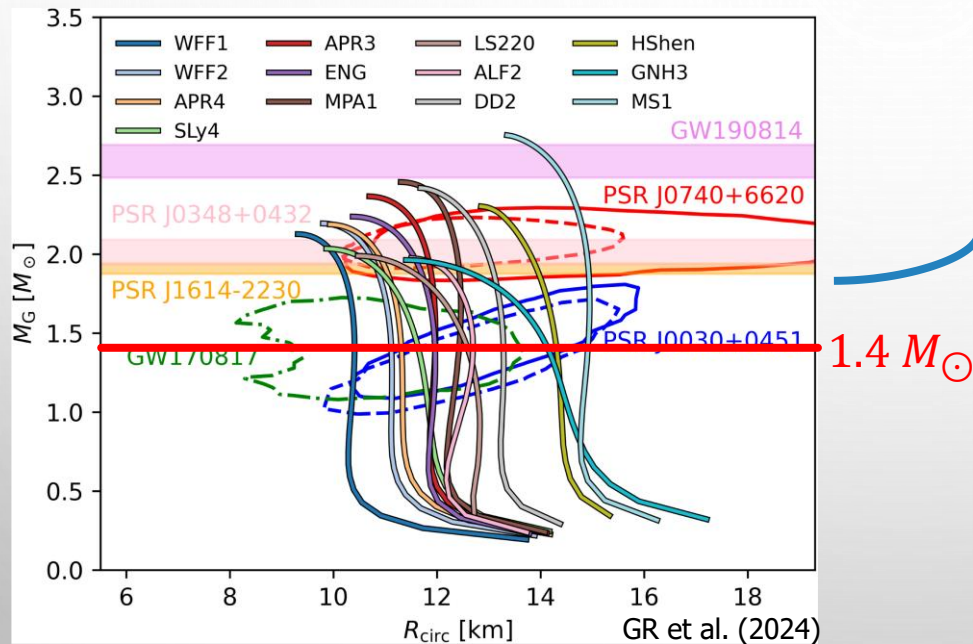
$$\text{GGL : } \Gamma = \Gamma_0 + (\Gamma_1 - \Gamma_0) e^{\frac{(\rho - \rho_1)^2}{2\sigma^2}}$$

(Gaussian Gamma Law)



WHAT ABOUT BNS SIMULATIONS?

Take initial Configurations at $1.4 M_{\odot}$



LORENE

<https://gitlab.in2p3.fr/lorene/Lorene>

Run GRHD simulations:

Einstein Toolkit + IllinoisGRMHD

www.einsteintoolkit.org

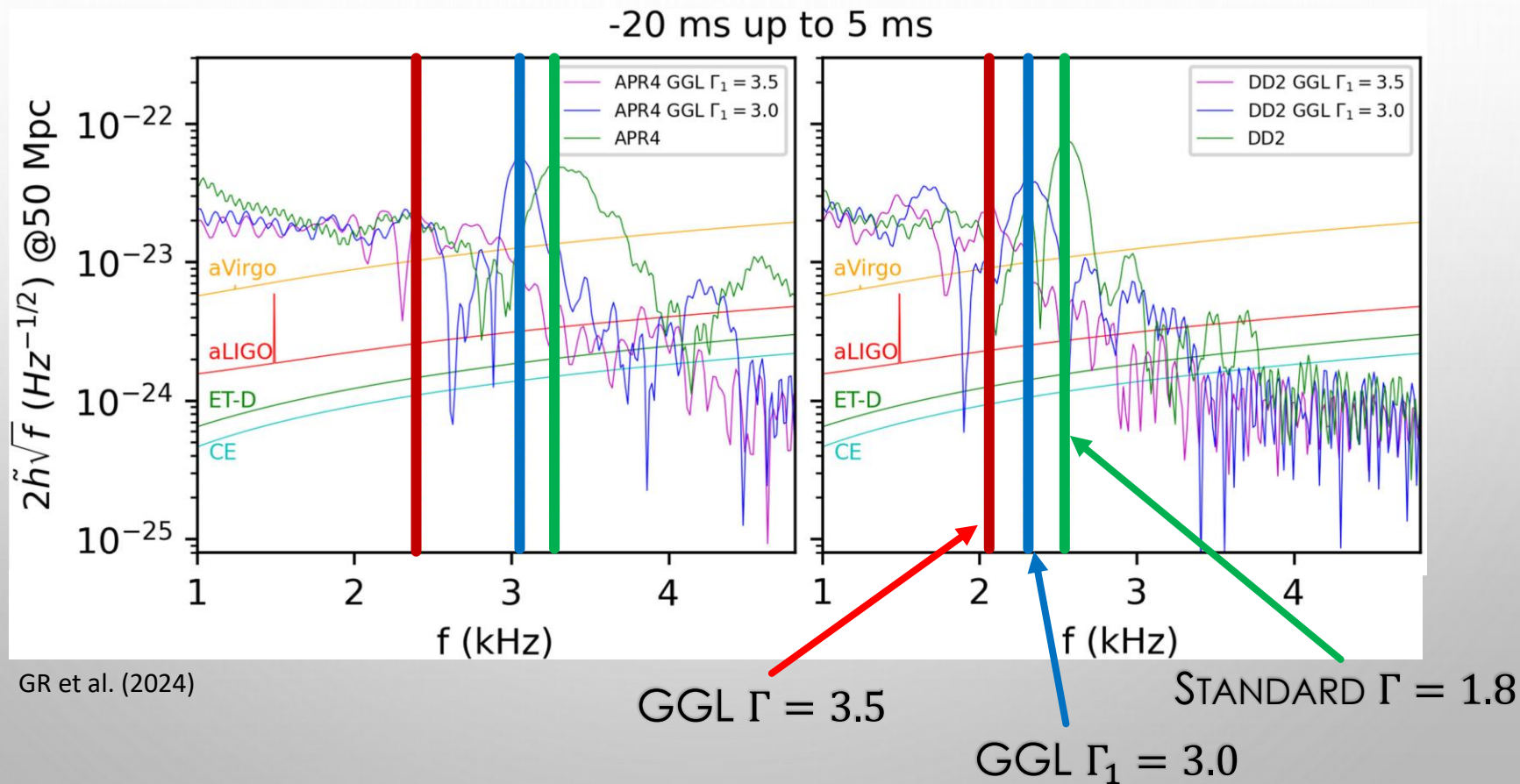


Piecewise Politrope + Γ -law

$$P = k\rho^n + (\Gamma_{th} - 1)\epsilon n$$

$$\Gamma_{th} = 1.8 \quad \text{or} \quad \Gamma_{th} = GGL$$

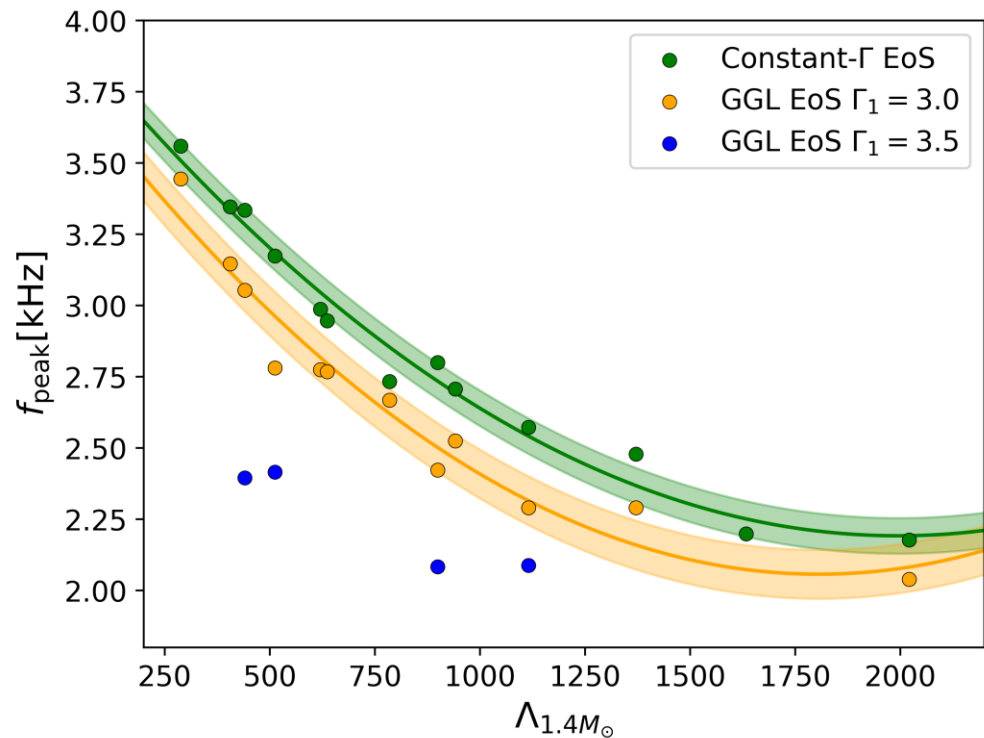
RESULTS: FREQUENCY PEAK SHIFT WITH STRONGER NONCONVEXITY



QUASI-UNIVERSAL RELATION

Shifted Quasi-Universal relation!

Stronger outliers for more nonconvexity



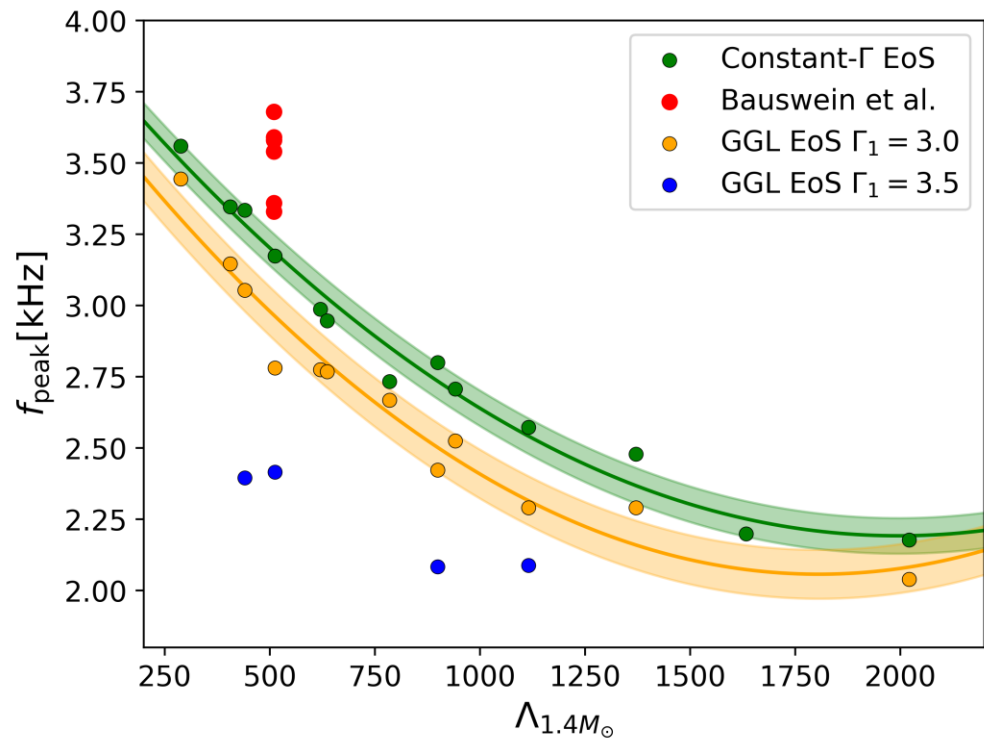
G. Riveccio et al. (2024) Phys. Rev. D **109**, 064032

QUASI-UNIVERSAL RELATION

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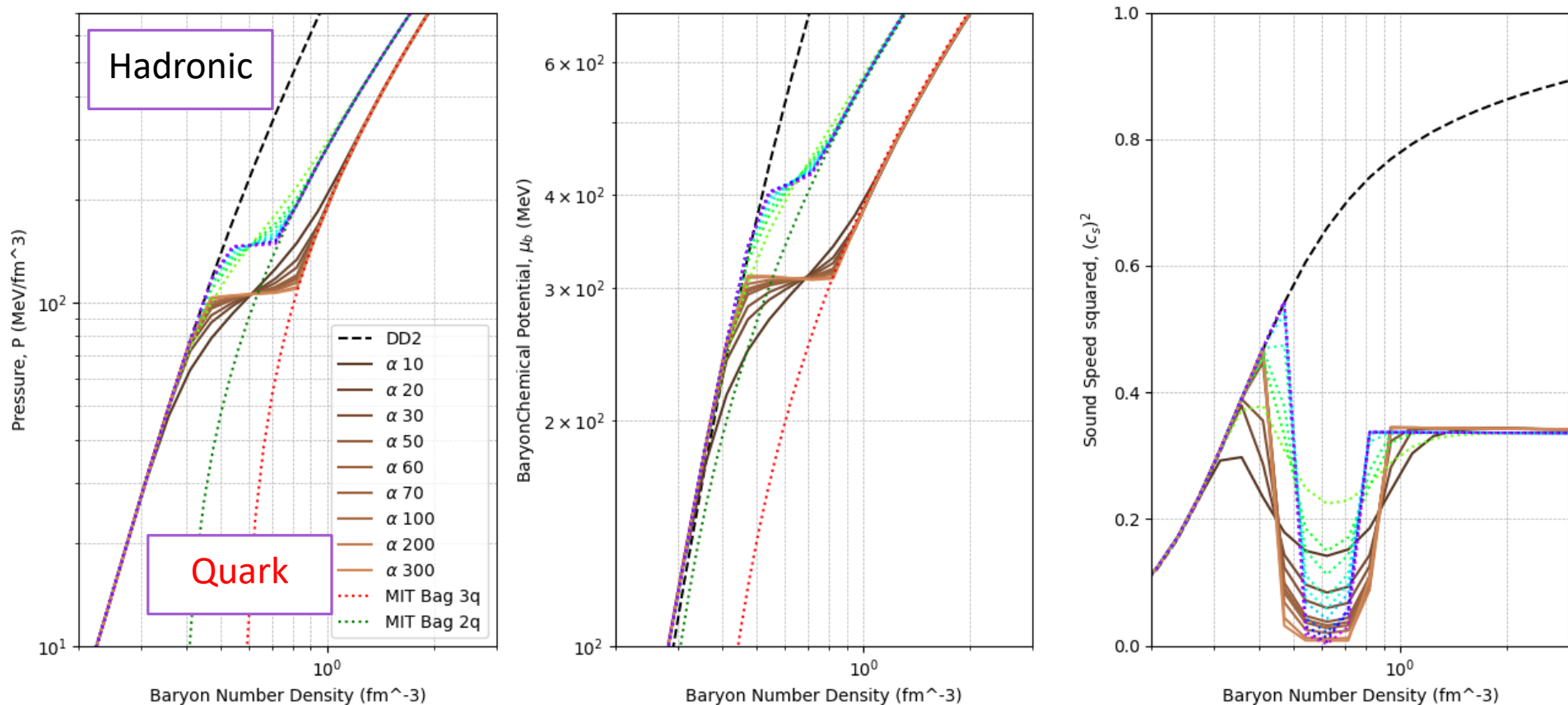
GGL EoS gives **lower frequencies** while phase transition has higher frequencies instead



G. Riveccio et al. (2024) Phys. Rev. D **109**, 064032

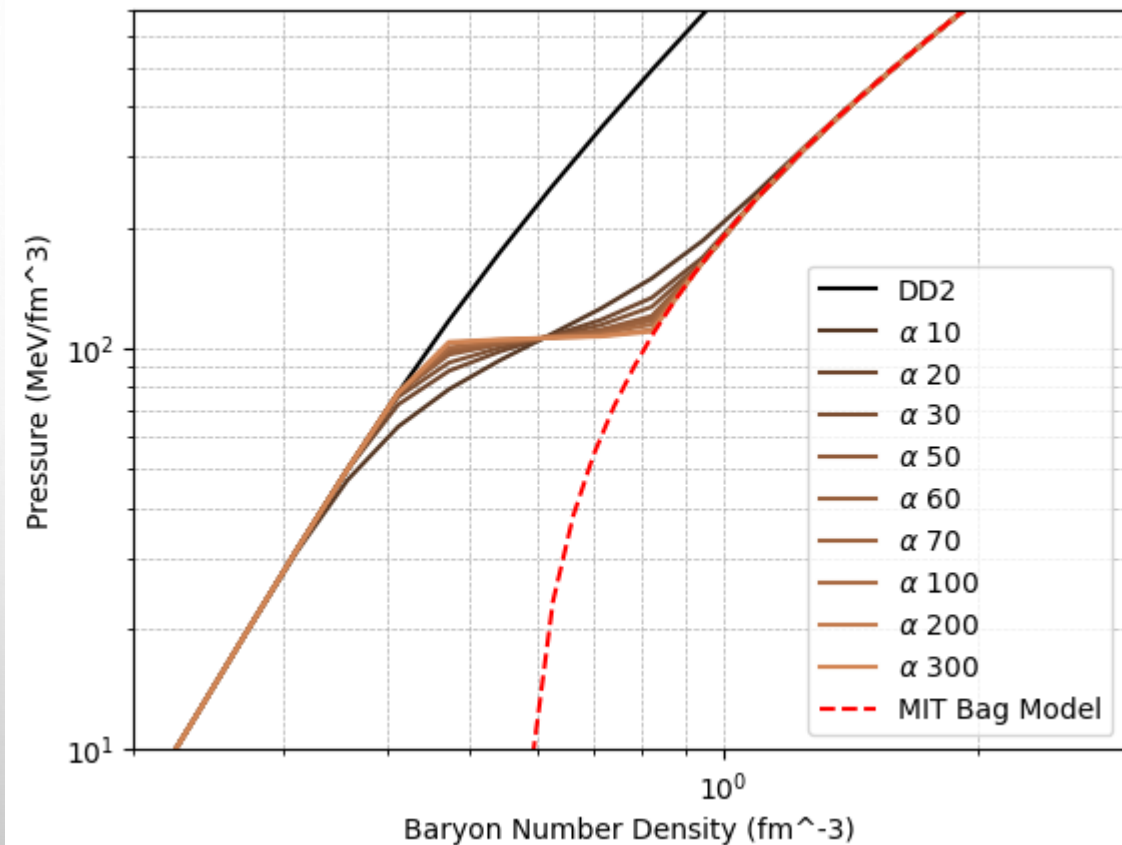
We need to have a phase transition with controlled convexity!

PARAMETRISED PHASE TRANSITION



PARAMETRISED PHASE TRANSITION

How do we build it?



PARAMETRISED PHASE TRANSITION

How do we build it?

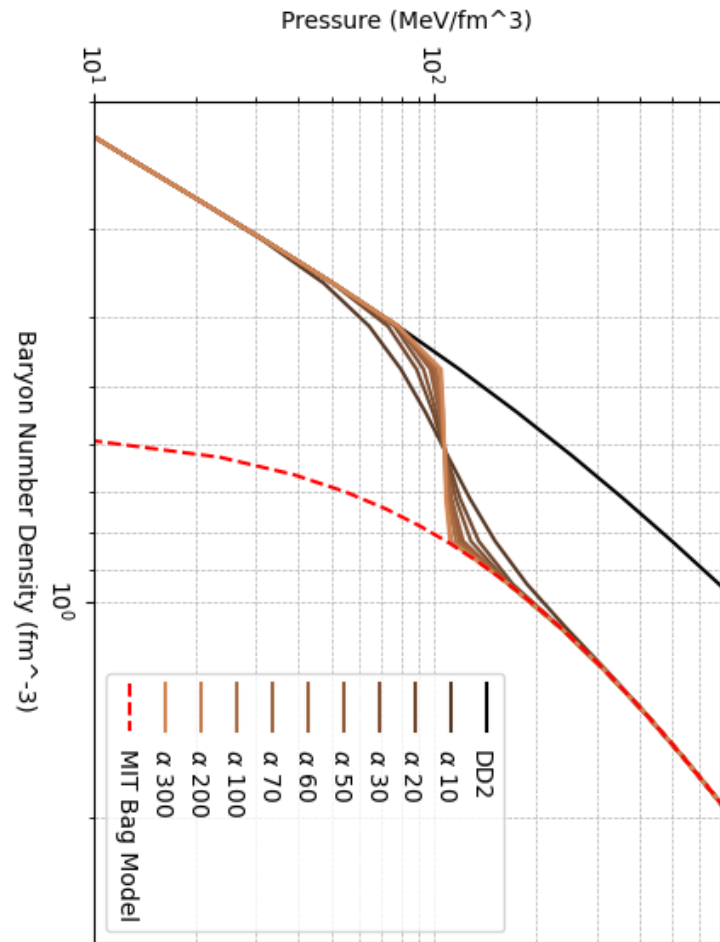
The inverted function is a step!
We apply a sigmoid.

$$n = n_{Nuc}\sigma + n_{Qrk}(1 - \sigma)$$

where

$$\sigma = \frac{1}{1 + e^{-\alpha(p-p_t)}}$$

α controls the sharpness
 p_t is the transition pressure



PARAMETRISED PHASE TRANSITION

How do we build it?

We apply the conditions

Mechanical equilibrium:

$$p_{Nuc}(n_{on}) = p_{Qrk}(n_{off})$$

Chemical equilibrium:

$$\mu_{g,Nuc}(n_{on}) = \mu_{g,Qrk}(n_{off})$$

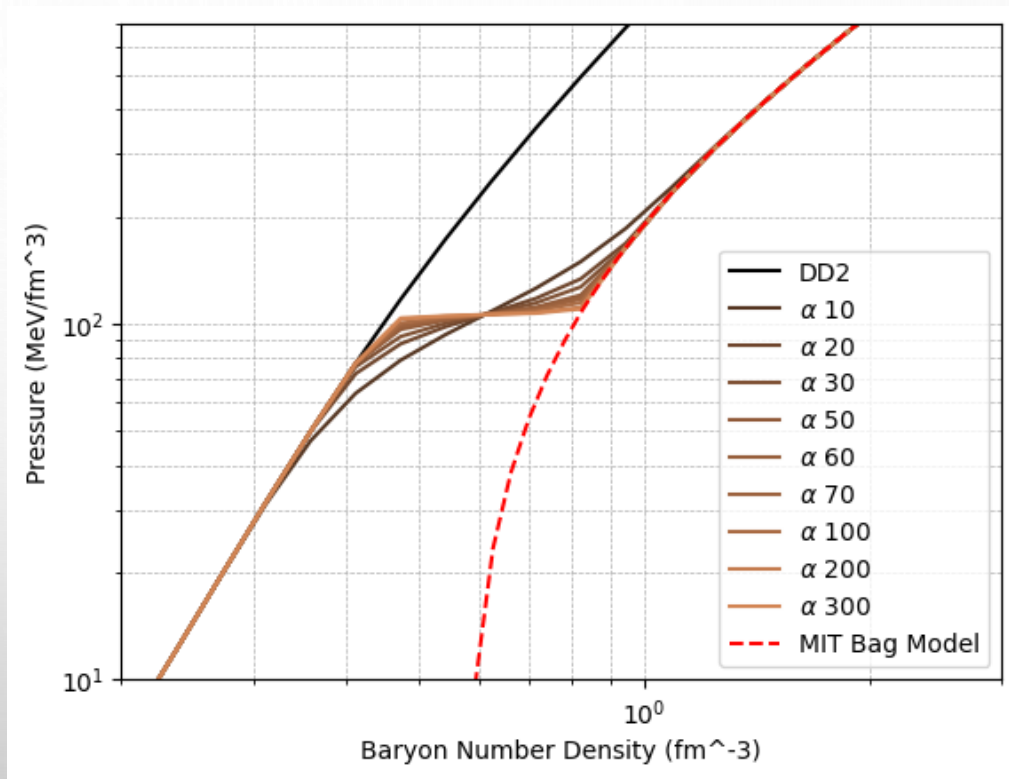
where

$$\mu_g = \mu_b + Y_q * \mu_q$$

or

$$\mu_g = \mu_b + Y_e * \mu_l$$

(if we do the transition with electrons)



PARAMETRISED PHASE TRANSITION

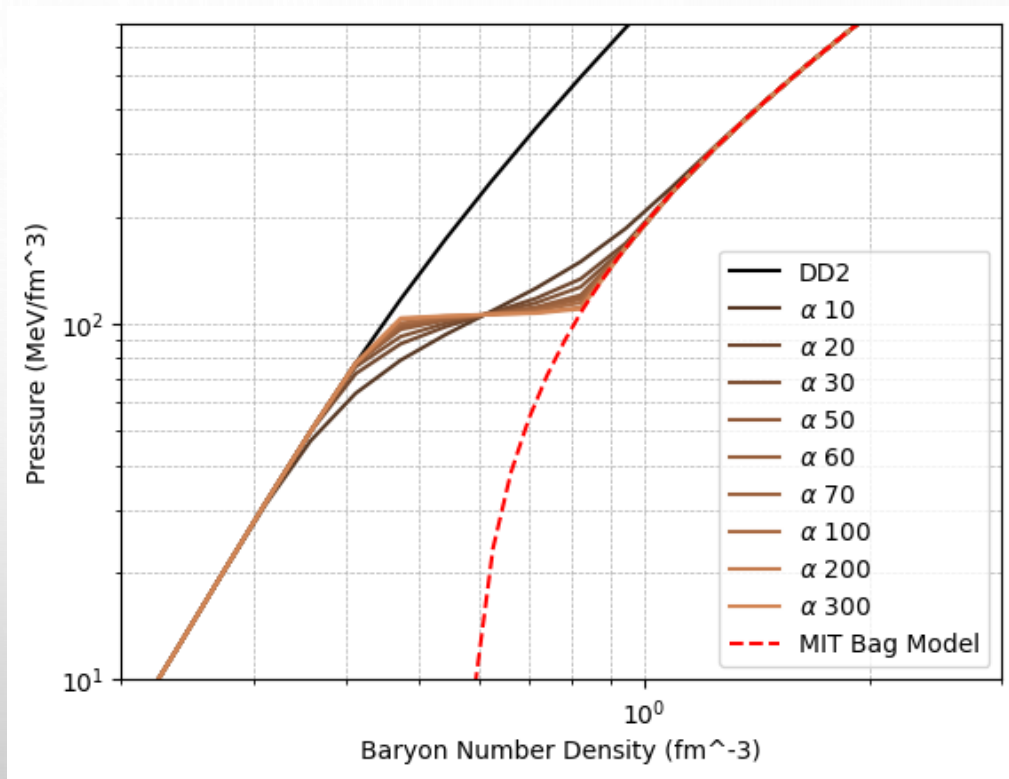
We have the pressure!

How do we get everything else?

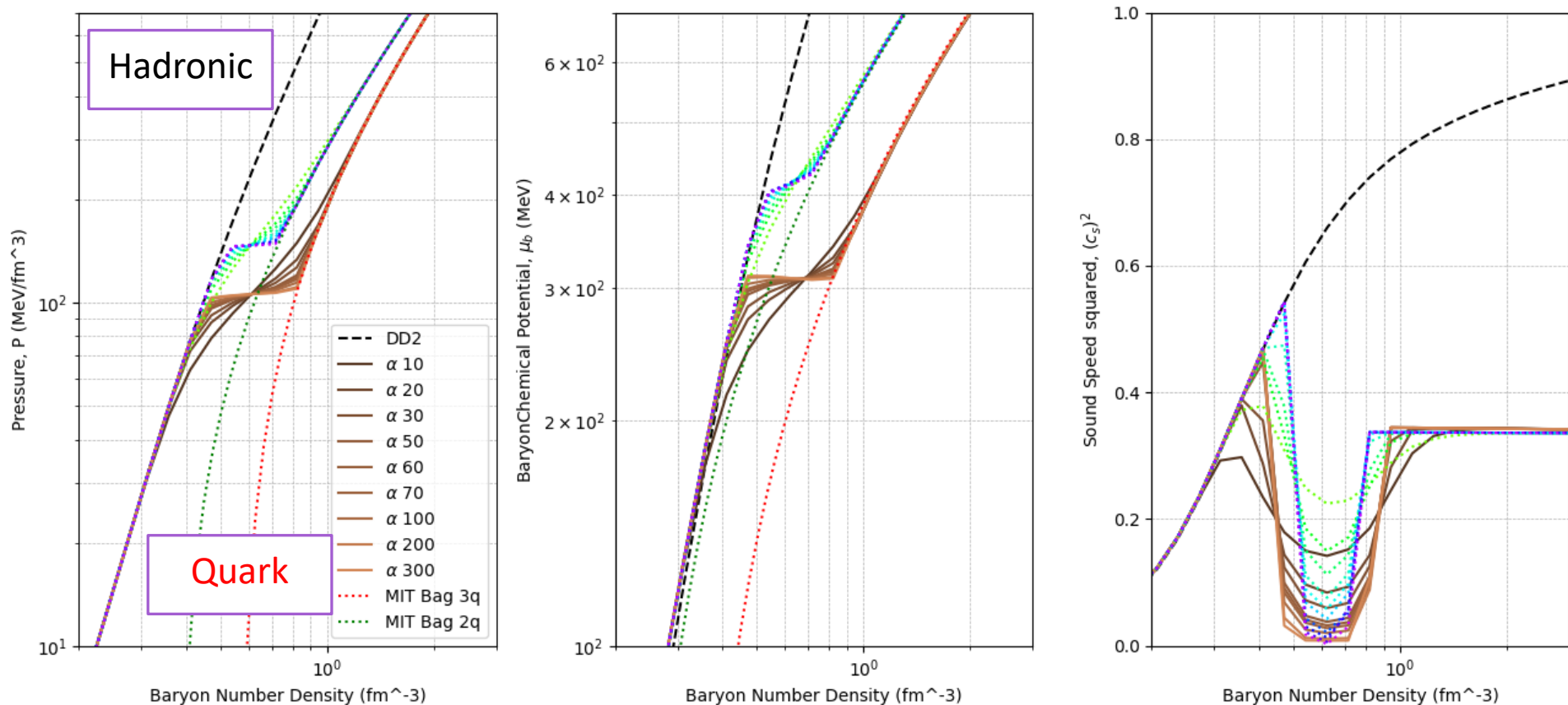
- Invert $P = n^2 \frac{\partial F}{\partial n}$
- From the free energy F we get everything else!

(Typel et al 2015)

To ensure thermodynamical consistency
Piecewise quintic Hermite polynomials



PARAMETRISED PHASE TRANSITION



Parametrised Phase Transition

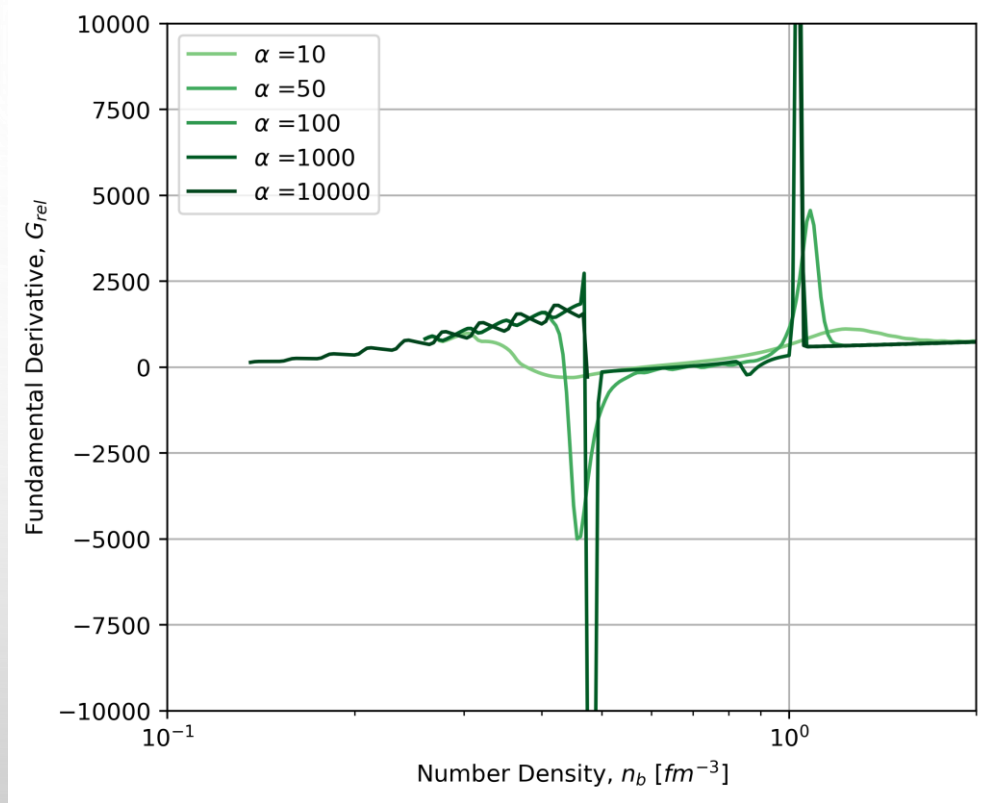
And the fundamental derivative?

We write

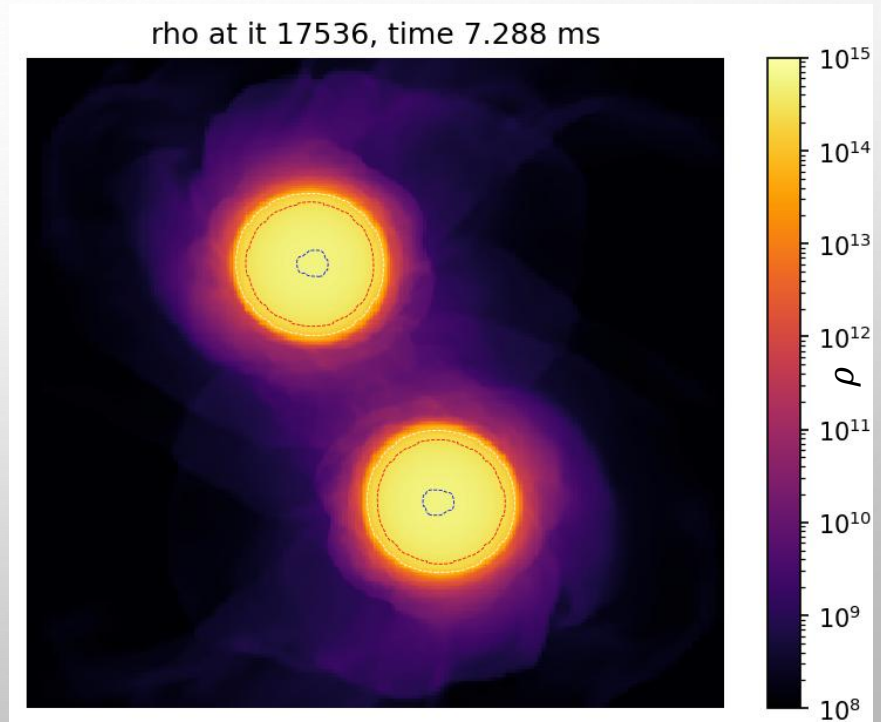
$$G = 1 + \frac{n_b}{2 c_s^2} \left. \frac{\partial c_{s,C}^2}{\partial n_b} \right|_s - \frac{3}{2} c_{s,R}^2$$

(C is Classical, R is Relativistic)

We will have a continuous crossover with still non-convex matter.

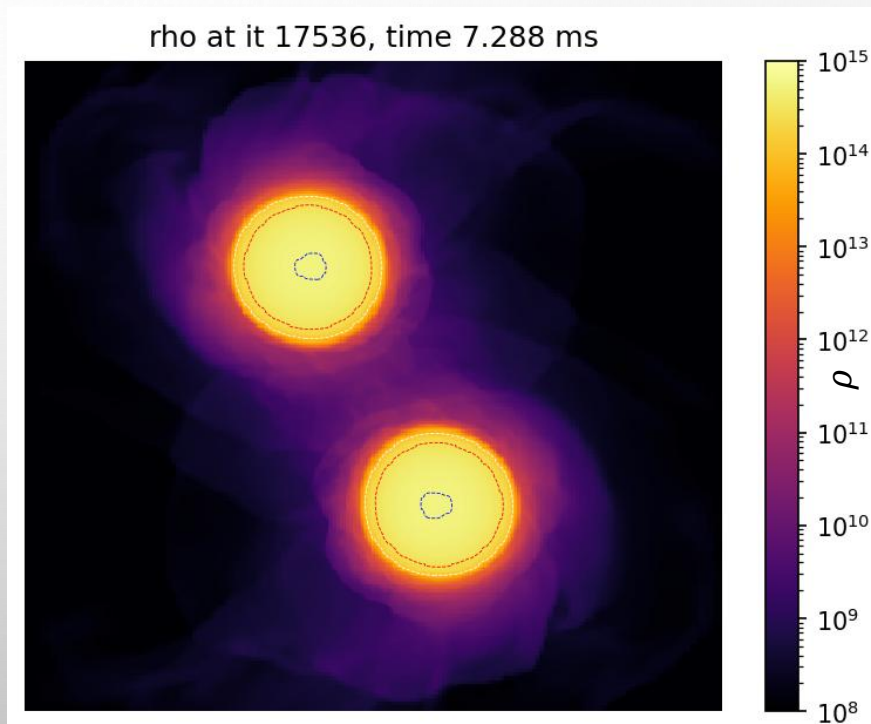
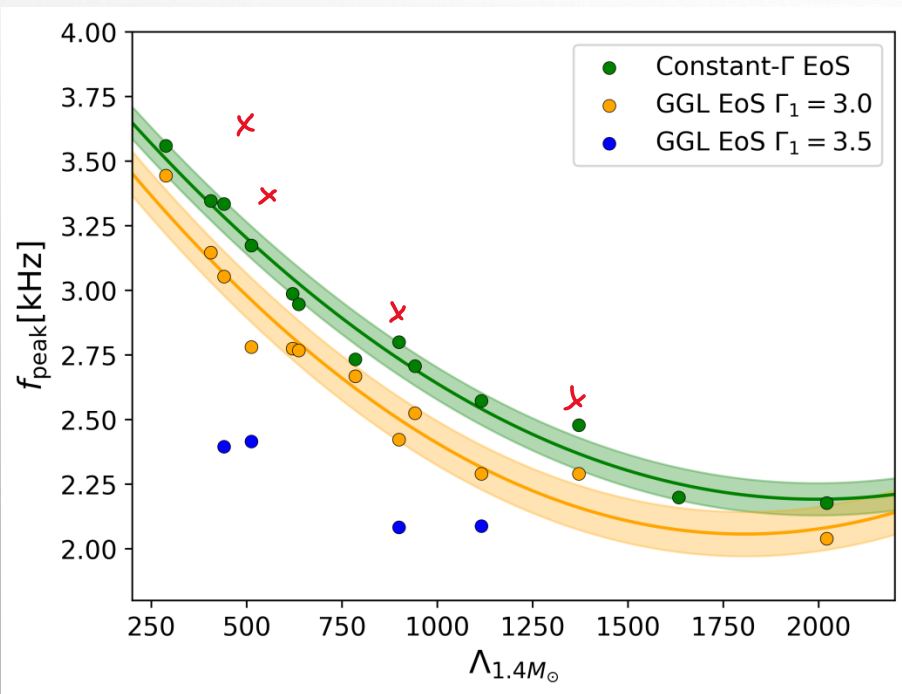


Simulations are running!



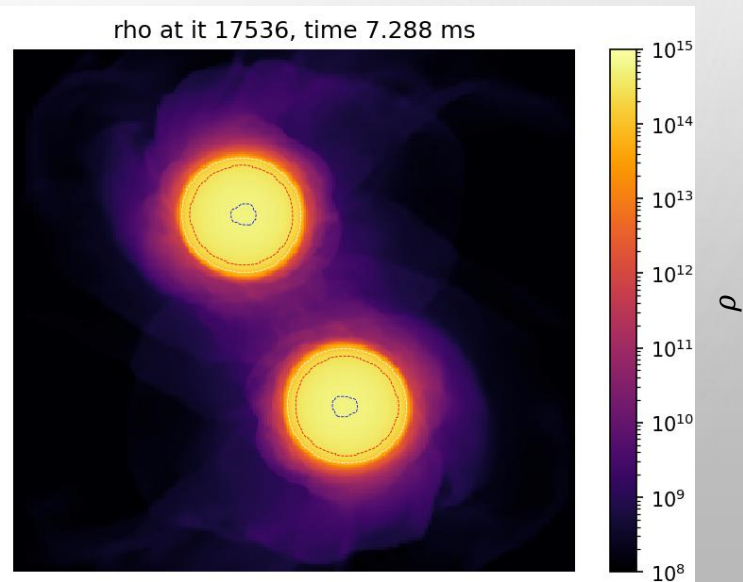
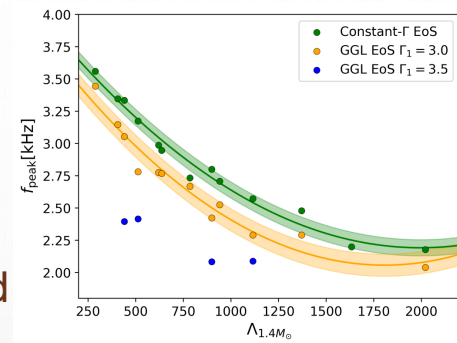
Simulations are running!

We expect to populate the high frequency side with both nonconvex crossover and first order phase transition.



Take away message!

- Convexity determines the wave structure and the evolution of the fluid
- Nonconvexity has an imprint in the gravitational wave!
- Phase transitions and crossovers are nonconvex
- Simulation are running to determine the definitive contribution of nonconvexity in BNS merger





M. Ruiz (UV)



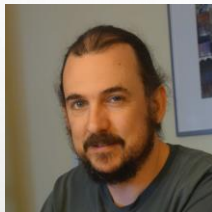
D. Guerra (UV)



A. J. Font (UV)



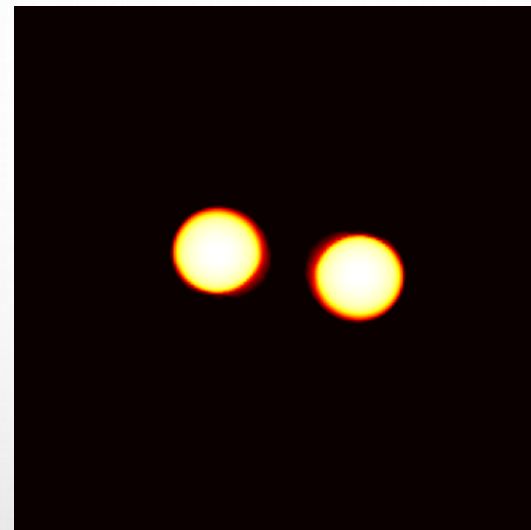
Arnau Rios (UB)



P. Cerdà-Durà (UV)

THANKS!

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