

Parton level Monte Carlos, vector boson scattering and other fun work with Tao

Dieter Zeppenfeld
TaoFest, May 14-15, 2026, Pittsburgh

KIT Center Elementary Particle and Astroparticle Physics - KCETA



40 years ago to the day: 1986 Madison Workshop

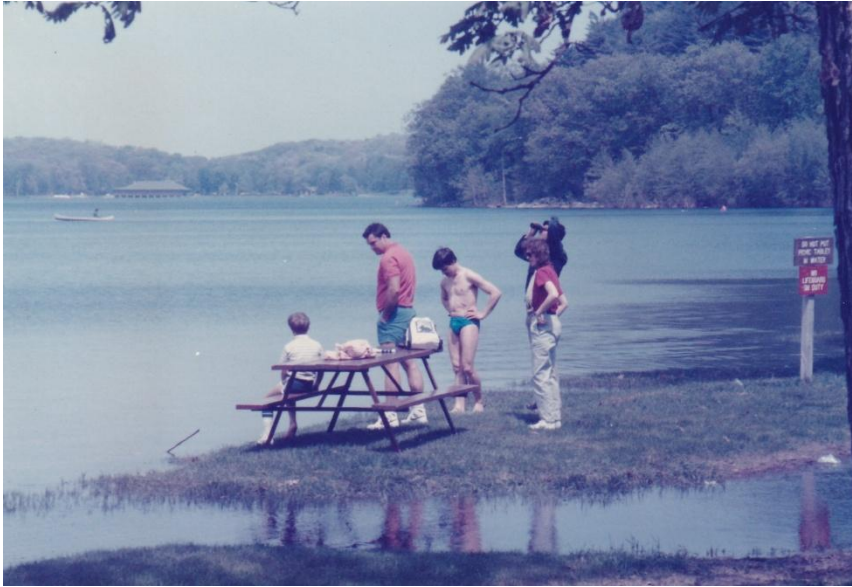
From Colliders to Supercolliders

First time I met Tao, who had helped organize the workshop

40 years ago to the day: 1986 Madison Workshop

From Colliders to Supercolliders

First time I met Tao, who had helped organize the workshop



Workshop excursion and picnic at Devil's Lake

What followed was a productive period of work with Tao and Vernon

Large- p_T Weak-Boson Production at the Fermilab Tevatron

V. Barger,⁽¹⁾ T. Han,⁽¹⁾ J. Ohnemus,⁽²⁾ and D. Zeppenfeld⁽¹⁾

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(Received 15 December 1988)

Perturbative QCD calculations of the processes $p\bar{p} \rightarrow W^\pm, Z + n$ jets ($n=1,2,3$) with $W \rightarrow e\nu$ and $Z \rightarrow e\bar{e}$ decays are presented for comparison with forthcoming Fermilab Tevatron data at $\sqrt{s} = 1.8$ TeV. The $W, Z + 2$ -jet results are compared with the corresponding $p\bar{p} \rightarrow WW, WZ$, and ZZ electroweak cross sections.

PACS numbers: 13.85.Qk, 12.38.Bx

- LO QCD simulations of W (or Z) production with additional jets
- Programming helicity amplitude tools (based on work with Kaoru)
- pre-HELAS, still the basic ingredient of VBFNLO amplitude calculations
- Important help from Bill Long to generate fast code

Including pheno work






Physics Letters B
Volume 220, Issue 3, 6 April 1989, Pages 464-470



Constraints on exotic particles from tevatron missing transverse momentum data


V. Barger ^a, K. Hagiwara ^{a, b}, T. Han ^a, D. Zeppenfeld ^a

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[https://doi.org/10.1016/0370-2693\(89\)90904-0](https://doi.org/10.1016/0370-2693(89)90904-0)

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Abstract

The search at the Tevatron for events with large missing transverse momentum ($p_{\text{stroke};\tau}$), which set mass bounds exceeding 73 GeV for both squarks and gluinos, implies significantly improved mass limits on other exotic particles. We study various sources for $p_{\text{stroke};\tau}$ events: production of leptoquarks or their fermionic partners in E_6 models, Majorana leptogluons in composite models of quarks and leptons, and cosmions in a model which explains the solar neutrino and dark matter problems. The CDF data can already exclude neutral leptogluons with masses below 110 GeV, while an integrated luminosity of order 2 pb^{-1} would set lower mass bounds of order 100 GeV for the other exotic particles.

...some of which
probably all 4 of us
have forgotten about
long time ago

More important: dijet + weak boson pairs → background to weak boson scattering

PHYSICAL REVIEW D

VOLUME 41, NUMBER 9

1 MAY 1990

Pair production of W^\pm , γ , and Z in association with jets

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J. Ohnemus

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(Received 22 November 1989)

Amplitudes for the production of electroweak-gauge-boson pairs (any combination of γ , Z , or W^\pm) in association with up to two jets are given. The gauge bosons may be either real or virtual with subsequent decays into charged or neutral leptons. The amplitudes are presented in a form directly amenable to numerical evaluation. Representative cross sections are given for Fermilab Tevatron, CERN Large Hadron Collider, and Superconducting Super Collider center-of-mass energies.

... and then we explored forward jet tagging

PHYSICAL REVIEW D

VOLUME 44, NUMBER 5

1 SEPTEMBER 1991

Comparative study of the benefits of forward jet tagging in heavy-Higgs-boson production at the Superconducting Super Collider

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(Received 10 April 1991)

The event rate for production of a Higgs boson of mass ~ 1 TeV with decay $H \rightarrow ZZ \rightarrow 4$ charged leptons is of order 25 events per year at standard Superconducting Super Collider luminosity and the QCD background is of comparable size. By tagging a *single* forward jet of energy $E_j > 1$ TeV and rapidity $2 < |\eta_j| < 5$ from the $qq \rightarrow qqZZ$ process, the QCD background can be essentially eliminated, with about 10 Higgs-boson signal events per year remaining, which amounts to 70% of the $qq \rightarrow qqZZ$ signal rate. The experimental separation of the vector-boson scattering subprocess is thereby possible.

Single-forward-jet tagging and central-jet vetoing to identify the leptonic WW decay mode of a heavy Higgs boson

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(Received 30 May 1991)

We study the extraction of the heavy-Higgs-boson signal $H \rightarrow W^+W^- \rightarrow \bar{\ell}\nu, \ell\bar{\nu}$ ($\ell = e$ or μ) from the standard-model background at hadron supercolliders. By tagging a single forward jet with energy $E_j > 3$ TeV and pseudorapidity $3 < |\eta_j| < 5$ and by vetoing central jets of transverse momenta $p_{Tj} > 60$ GeV in the pseudorapidity range $0 < |\eta_j| < 3$, the QCD WWj and $t\bar{t}j \rightarrow WWb\bar{b}j$ backgrounds are suppressed. For $m_H = 1$ TeV there are about 46 signal events from electroweak-vector-boson scattering (of which 36 events are of Higgs-boson origin) at the Superconducting Super Collider (SSC) for an integrated luminosity of 10 fb^{-1} and 10 other events from the WWj and $t\bar{t}j$ backgrounds for $m_t = 140$ GeV. The experimental separation of the vector-boson-scattering subprocess is thereby possible. At the CERN Large Hadron Collider, with an $E_j > 2$ TeV jet energy cut, all cross sections are about a factor of 10 below the SSC values.

PHYSICAL REVIEW D

VOLUME 46, NUMBER 5

1 SEPTEMBER 1992

Full tree-level calculation of the $qq \rightarrow qqWZ$ electroweak process at hadron supercolliders

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T. Han

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A. Stange and D. Zeppenfeld

Department of Physics, University of Wisconsin, Madison, Wisconsin 53706

(Received 23 December 1991)

We present a full standard model calculation of the $qq \rightarrow qqWZ$ electroweak process in pp collisions at the Superconducting Super Collider, including W and Z leptonic decay correlations. We also analyze the backgrounds to this signal from $Zt\bar{t} \rightarrow ZW + \text{jets}$ and $q\bar{q} \rightarrow WZ + \text{jets}$. Single forward jet tagging and vetoing of events with ≥ 2 central jets can suppress the backgrounds with little effect on the signal. With 10 fb^{-1} integrated luminosity we expect 50 signal events and 30 background events.

PACS number(s): 13.85.Qk, 12.15.Cc, 12.15.Ji, 14.80.Er

Finding the leptonic WW decay mode of a heavy Higgs boson at hadron supercolliders

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¹*Department of Physics, University of Wisconsin, Madison, Wisconsin 53706*

²*Department of Physics and Astronomy, Northwestern University, Evanston, Illinois 60208*

³*Fermi National Accelerator Laboratory, P. O. Box 500, Batavia, Illinois 60510*

(Received 13 May 1993)

We reanalyze the extraction of the heavy Higgs boson signal $H \rightarrow W^+W^- \rightarrow \bar{\ell}\nu, \ell\bar{\nu}$ ($\ell = e$ or μ) from the standard model background at hadron supercolliders, taking into account revised estimates of the top quark background. With new acceptance criteria the detection of the signal remains viable. Requiring a forward jet tag, a central jet veto, and a large relative transverse momentum of the two charged leptons yields $S/\sqrt{B} > 6$ for one year of running at the SSC or CERN LHC.

PACS number(s): 14.80.Gt, 13.85.Qk, 13.87.-a

At the time, weak boson fusion was thought to be useful for heavy Higgs search only. For light (SUSY) Higgs gluon fusion and Higgs-strahlung were the main anticipated search modes

Collaboration with Tao didn't quite stop when he moved to Davis

PHYSICAL REVIEW D

VOLUME 56, NUMBER 1

1 JULY 1997

$W\gamma\gamma$ production at the Fermilab Tevatron collider: Gauge invariance and radiation amplitude zero

U. Baur,¹ T. Han,² N. Kauer,³ R. Sobey,² and D. Zeppenfeld³

¹*Department of Physics, State University of New York, Buffalo, New York 14260*

²*Department of Physics, University of California, Davis, California 95616*

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(Received 19 February 1997)

The electroweak process $pp\bar{\ell}^{\pm}\nu\gamma\gamma$ is calculated at the tree level, including finite W width effects. In order to obtain a gauge-invariant amplitude, the imaginary parts of $WW\gamma$ triangle graphs and $WW\gamma\gamma$ box diagrams have to be included, in addition to resumming the imaginary contributions to the W vacuum polarization. We demonstrate the existence of a radiation amplitude zero in $pp\bar{\nu}\rightarrow W^{\pm}\gamma\gamma\rightarrow\ell^{\pm}\nu\gamma\gamma$, and discuss how it may be observed in correlations of the $\gamma\gamma$ and lepton rapidities at the Fermilab Tevatron. [S0556-2821(97)04813-3]

PACS number(s): 12.15.Lk, 13.85.Qk, 14.70.Fm

In our albums I did not find any photos of Tao from these first ten years....-

Barger/Olsson: Classical Mechanics at 25



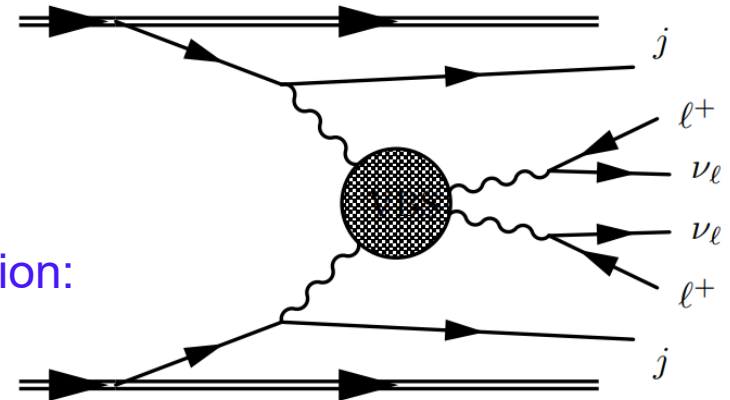
1998: celebrating 25 years
of a great classical mechanics
textbook



... building on work with Tao

Early work on forward jet tagging and central jet veto led to studies of vector boson fusion and weak boson scattering

- Basic process: $VV \rightarrow VV$
- accompanied by 2 quark jets
= tagging jets
- Identify weak bosons by decay leptons
(or hadronic V decay)
- Characteristic radiation pattern of gluon emission:
suppressed emission in central region
- Central jet veto suppresses background



Long standing problem: reliable modeling of $VV+3$ jets or $H+3$ jets for signal and background. Need NLO QCD calculations for these processes

- $H + 3$ jets (PhD thesis of Terrance Figy)
- $Z + 3$ jets is excellent lab for experimental verification of NLO expectations
- W^+W^++3 jets recently calculated by Barbara Jäger et al.

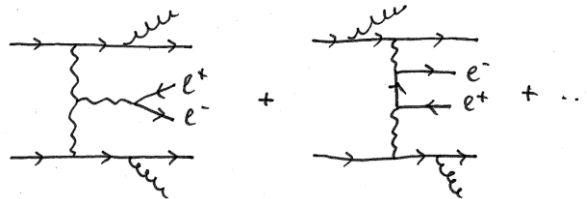
Speeding up NLO calculations: VBF Z+3jet as a case study

(work with T. Jezo and M. Löschner)

Many contributing subprocesses

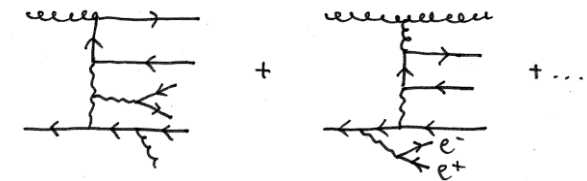
e.g. 1117 (118) Z+4 parton real emission processes with 5 (2) massless quark flavors

■ Quark initiated



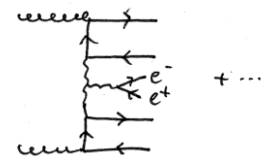
x4 for $q \rightarrow q\bar{q}$
6 distinct flavor combinations

■ qg initiated



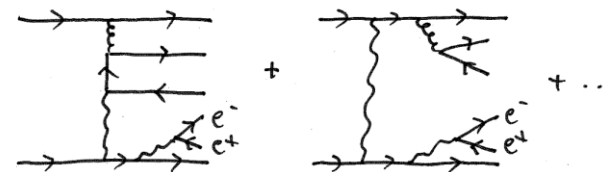
x4 for $q \rightarrow q\bar{q}$, $qg \rightarrow qg$
6 distinct flavor combinations

■ gg initiated



6 distinct flavor combinations

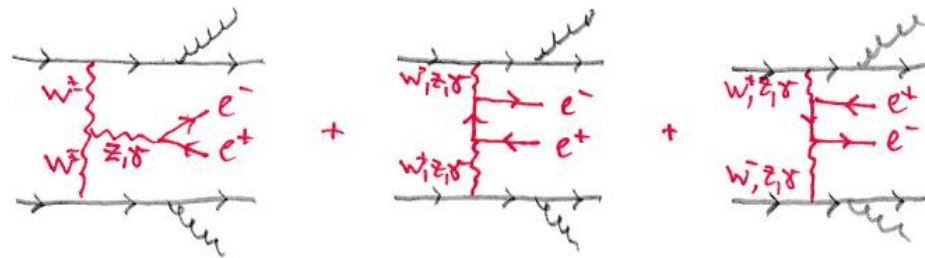
■ 6 quark processes



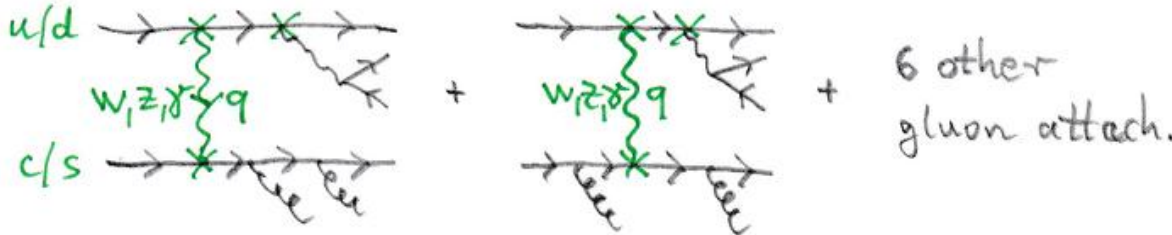
x4 for $q \rightarrow q\bar{q}$
16 distinct flavor combinations

- **Helas/MadGraph philosophy:**
 generate separate amplitude programs for each distinct flavor structure;
 do crossing by sign factors in separate numerical evaluations
 → no reuse of partial amplitudes
- **VBFNLO philosophy:** reuse partial amplitudes

- leptonic tensors
 for VBF part:
 $VV \rightarrow \text{leptons}$



- do spinor/Lorentz structure common to various flavor combinations only once.
 e.g.

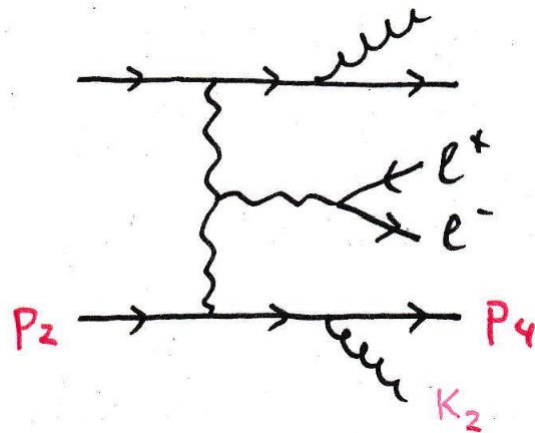


Multiply **scalar factor** containing flavor dependent couplings & boson propagators

$$M \sim g_z \cdot g_{gg} * \text{couplings (hel., flavor)} \frac{1}{q^2 - m^2}$$

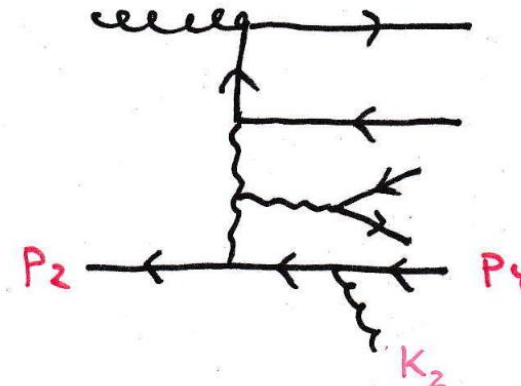
- Result: VBFNLO code for tree level $Z+4\text{quark}+2\text{gluon}$ amplitudes is 300 times faster than Helas routines generated with MadGraph-2
- Can we improve on this?
 - decay currents for $Z \rightarrow \text{leptons}$ and $\gamma \rightarrow \text{leptons}$ are proportional to each other for given fermion helicities
 - Charge conjugation relates quark- and antiquark-lines of opposite chirality: same except for sign factor

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 - decay currents for $Z \rightarrow \text{leptons}$ and $\gamma \rightarrow \text{leptons}$ are proportional to each other for given fermion helicities
 - Charge conjugation relates quark- and antiquark-lines of opposite chirality: same except for sign factor, for example



quark line

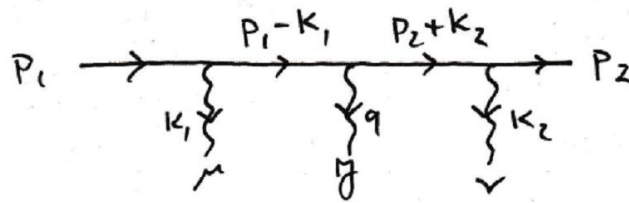
vs.



anti-quark line

Relating massless quark- to anti-quark-lines

Consider



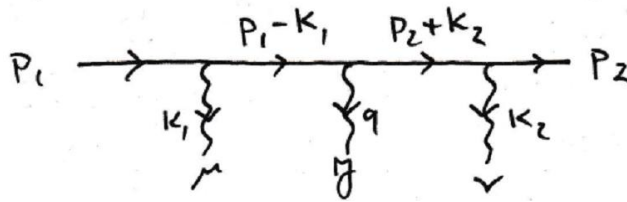
i.e.

$$\bar{u}(P_2) \gamma^\nu \frac{1}{\not{P}_2 + \not{k}_2} \not{P}_\sigma \frac{1}{\not{P}_1 - \not{k}_1} \gamma^\mu u(P_1)$$

of chirality σ : $P_\sigma = \frac{1}{2} (1 + \sigma \gamma_5)$

Relating massless quark- to anti-quark-lines

Consider



i.e.
$$\underbrace{\bar{u}(P_2)}_{-v^T(P_2)C^{-1}} \gamma^\nu \frac{1}{\cancel{P_2 + k_2}} \cancel{\not{P}_\sigma} \frac{1}{\cancel{P_1 - k_1}} \gamma^\mu \underbrace{u(P_1)}_{C\bar{v}^T(P_1)}$$

of chirality σ : $P_\sigma = \frac{1}{2}(1 + \sigma\gamma_5)$

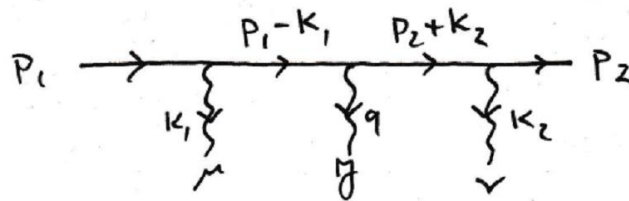
and use $C^{-1}\not{a}C = -\not{a}^T$

$$= -\bar{v}(P_2) (-\gamma^\mu) \frac{1}{-\cancel{P_1 + k_1}} P_\sigma (-\cancel{\not{P}}) \frac{1}{-\cancel{P_2 - k_2}} (-\gamma^\nu) v(P_2)$$

$$= \bar{v}(P_2) \gamma^\mu \frac{1}{\cancel{P_1 + k_1}} \cancel{\not{P}_\sigma} \frac{1}{\cancel{P_2 - k_2}} \gamma^\nu v(P_2)$$

Relating massless quark- to anti-quark-lines

Consider



i.e.

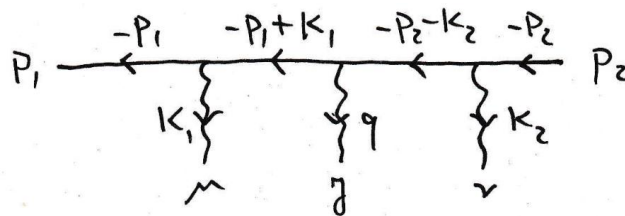
$$\bar{u}(P_2) \gamma^\nu \frac{1}{\not{P}_2 + \not{k}_2} \not{A} P_\sigma \frac{1}{\not{P}_1 - \not{k}_1} \gamma^\mu u(P_1)$$

$-v^T(P_2) C^{-1}$ $C \bar{v}^T(P_1)$

of chirality σ : $P_\sigma = \frac{1}{2} (1 + \sigma \gamma_5)$

and use $C^{-1} \not{a} C = -\not{a}^T$

$$= -\bar{v}(P_2) (-\gamma^\mu) \frac{1}{-\not{P}_1 + \not{k}_1} P_\sigma (-\not{A}) \frac{1}{-\not{P}_2 - \not{k}_2} (-\gamma^\nu) v(P_1)$$



$$= \bar{v}(P_1) \gamma^\mu \frac{1}{-\not{P}_1 + \not{k}_1} \not{A} P_\sigma \frac{1}{-\not{P}_2 - \not{k}_2} \gamma^\nu v(P_2)$$

fermion line (σ) = anti-fermion line ($-\sigma$) $\times (-1)^{\# \text{ of propagators}}$

Example: boxline routine in VBFNLO

- VBFNLO has dedicated routines for the sum of loop graphs correcting individual fermion lines, e.g. qqVg, qqgV, qqVV lines

$$\left[P_1 \rightarrow \begin{array}{c} \text{gluon loop} \\ \text{quark line} \end{array} \rightarrow P_2 + \begin{array}{c} \text{gluon loop} \\ \text{quark line} \end{array} + \begin{array}{c} \text{gluon loop} \\ \text{quark line} \end{array} + \begin{array}{c} \text{gluon loop} \\ \text{quark line} \end{array} \right]_{\sigma}$$

$$= (-1)^3 \left[P_1 \leftarrow \begin{array}{c} \text{gluon loop} \\ \text{quark line} \end{array} \leftarrow P_2 + \dots \right]_{-\sigma}$$

- Antiquark line with identical external momenta and gauge boson ordering does not need extra calculation
- Note sign difference of antiquark line with 3-gluon vertex (non-abelian box)
- Analogous results for quark lines with pentagons or hexagons
- Easy to implement when basic building blocks of amplitudes are currents or tensors representing entire quark lines

- Result: VBFNLO code for tree level $Z+4\text{quark}+2\text{gluon}$ amplitudes is 300 times faster than Helas routines generated with MadGraph-2
- Can we improve on this?
 - decay currents for $Z \rightarrow \text{leptons}$ and $\gamma \rightarrow \text{leptons}$ are proportional to each other for given fermion helicities
 - Charge conjugation relates quark- and antiquark-lines of opposite chirality: same except for sign factor
 - **This charge conjugation relation also holds for loop amplitudes**
 - **Pauli interference for identical fermions enters at sub-percent level \rightarrow calculate interference contributions separately with reduced statistics and neglect fermion interchange in main Monte Carlo run**
 - 1-loop corrections enter via interference with Born amplitude. After Catani-Seymour (CS) subtraction they are small compared to Born^2 and can be evaluated separately with substantially reduced statistics
 - **Born amplitudes for a number of CS subtractions from real emission are identical (78 distinct subtraction configurations for $Z+3\text{jet VBF}$ out of maximum total of 228): no need to recalculate the rest again**

Combining quark and anti-quark scattering

- Do spinor calculation of amplitude
 - without coupling factors and 1-particle reducible gauge boson propagators
 - for one direction of fermion number flow only
- This yields stripped amplitudes. Obtain anti-quark from quark diagrams (of opposite chirality) by sign factors for groups of diagrams with same number of propagators on individual quark lines
- Coupling×propagator factors depend on external quark flavors, helicities and momenta
- Choose momentum assignments to maximize relationships between different subamplitudes/Feynman diagrams
- Don't calculate the same thing twice

Summary

Implementation in VBFNLO code for Z+3jet at NLO QCD resulted in

- Loop contributions take only ~ 5 times longer than Born for same statistical error
- Additional factor 5 speed increase in cross section calculation of sum of $qq \rightarrow Zqqgg$ processes
- $>95\%$ of CPU time is spent in (CS subtracted) real emission contributions \Rightarrow speedup of entire NLO code by factor 5

Summary

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Could this be implemented in generators like MadGraph?

Message to Tao

I had a lot of fun working with you and discussing physics over the last 40 years

Messages to Tao

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discussing physics over the last 40 years

Retirement is great:

No committees

No writing of grant proposals

No department meetings

But the fun with physics continues

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Thank you!!

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