Searching for Dark Matter with Magnetic Resonance and Atomic Clocks

Yevgeny Stadnik

Humboldt Fellow

Johannes Gutenberg University, Mainz, Germany

Collaborators (Theory):

Victor Flambaum group (UNSW) Peter Wolf group (SYRTE)

Collaborators (Experiment):

Dmitry Budker group (Mainz) nEDM collaboration at PSI and Sussex

Quantum Sensors for Fundamental Physics, Oxford, October 2018

Overwhelming astrophysical evidence for existence of **dark matter** (~5 times more dark matter than ordinary matter).



Traditional "scattering-off-nuclei" searches for heavy WIMP dark matter particles ($m_{\chi} \sim \text{GeV}$) have not yet produced a strong positive result.



Traditional "scattering-off-nuclei" searches for heavy WIMP dark matter particles ($m_{\chi} \sim \text{GeV}$) have not yet produced a strong positive result.



 $\mathcal{M} \propto \left(e' \right)^2$

Traditional "scattering-off-nuclei" searches for heavy WIMP dark matter particles ($m_{\chi} \sim \text{GeV}$) have not yet produced a strong positive result.



 $\mathcal{M} \propto \left(e'
ight)^2$

 $\sum N => \frac{d\sigma}{d\Omega} \propto |\mathcal{M}|^2 \propto (e')^4$

Traditional "scattering-off-nuclei" searches for heavy WIMP dark matter particles ($m_{\chi} \sim \text{GeV}$) have not yet produced a strong positive result.



<u>Challenge</u>: Observable is <u>fourth power</u> in a small interaction constant (e' << 1)!

Traditional "scattering-off-nuclei" searches for heavy WIMP dark matter particles ($m_{\chi} \sim \text{GeV}$) have not yet produced a strong positive result.



Question: Can we instead look for effects of dark matter that are **first power** in the interaction constant?

• Low-mass spin-0 particles form a coherently oscillating classical field $\varphi(t) = \varphi_0 \cos(m_{\varphi}c^2t/\hbar)$, with energy density $<\rho_{\varphi}> \approx m_{\varphi}^2 \varphi_0^2/2$ ($\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3$)



- Low-mass spin-0 particles form a coherently oscillating classical field $\varphi(t) = \varphi_0 \cos(m_{\varphi}c^2t/\hbar)$, with energy density $<\rho_{\varphi}> \approx m_{\varphi}^2 \varphi_0^2/2 \ (\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3)$
- Coherently oscillating field, since *cold* ($E_{\varphi} \approx m_{\varphi}c^2$)

- Low-mass spin-0 particles form a coherently oscillating classical field $\varphi(t) = \varphi_0 \cos(m_{\varphi}c^2t/\hbar)$, with energy density $<\rho_{\varphi}> \approx m_{\varphi}^2 \varphi_0^2/2 \ (\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3)$
- Coherently oscillating field, since *cold* ($E_{\varphi} \approx m_{\varphi}c^2$)
- Classical field for $m_{\varphi} << 1 \text{ eV}$, since $n_{\varphi}(\lambda_{\mathrm{dB},\varphi}/2\pi)^3 >> 1$

- Low-mass spin-0 particles form a coherently oscillating classical field $\varphi(t) = \varphi_0 \cos(m_{\varphi}c^2t/\hbar)$, with energy density $<\rho_{\varphi}> \approx m_{\varphi}^2 \varphi_0^2/2 \ (\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3)$
- Coherently oscillating field, since *cold* ($E_{\varphi} \approx m_{\varphi}c^2$)
- Classical field for $m_{\varphi} << 1 \text{ eV}$, since $n_{\varphi}(\lambda_{\mathrm{dB},\varphi}/2\pi)^3 >> 1$
- Coherent + classical DM field = "Cosmic laser field"

- Low-mass spin-0 particles form a coherently oscillating classical field $\varphi(t) = \varphi_0 \cos(m_{\varphi}c^2t/\hbar)$, with energy density $<\rho_{\varphi}> \approx m_{\varphi}^2 \varphi_0^2/2 \ (\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3)$
- Coherently oscillating field, since *cold* ($E_{\varphi} \approx m_{\varphi}c^2$)
- Classical field for $m_{\varphi} << 1 \text{ eV}$, since $n_{\varphi}(\lambda_{\mathrm{dB},\varphi}/2\pi)^3 >> 1$
- Coherent + classical DM field = "Cosmic laser field"
- $10^{-22} \text{ eV} \leq m_{\varphi} << 1 \text{ eV} <=> 10^{-8} \text{ Hz} \leq f << 10^{14} \text{ Hz}$ \uparrow $\lambda_{\text{dB},\varphi} \leq L_{\text{dwarf galaxy}} \sim 1 \text{ kpc}$ Classical field

- Low-mass spin-0 particles form a coherently oscillating classical field $\varphi(t) = \varphi_0 \cos(m_{\varphi}c^2t/\hbar)$, with energy density $<\rho_{\varphi}> \approx m_{\varphi}^2 \varphi_0^2/2 \ (\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3)$
- Coherently oscillating field, since *cold* ($E_{\varphi} \approx m_{\varphi}c^2$)
- Classical field for $m_{\varphi} << 1 \text{ eV}$, since $n_{\varphi}(\lambda_{\mathrm{dB},\varphi}/2\pi)^3 >> 1$
- Coherent + classical DM field = "Cosmic laser field"
- $10^{-22} \text{ eV} \leq m_{\varphi} << 1 \text{ eV} <=> 10^{-8} \text{ Hz} \leq f << 10^{14} \text{ Hz}$ \uparrow $\lambda_{\text{dB},\varphi} \leq L_{\text{dwarf galaxy}} \sim 1 \text{ kpc}$ Classical field
 - $m_{\varphi} \sim 10^{-22} \text{ eV} \iff T \sim 1 \text{ year}$



 \rightarrow Time-varying

fundamental constants

- Atomic clocks
- Optical cavities
- Fifth-force searches
- Astrophysics (e.g., BBN)

- → Time-varying spindependent effects
 - Co-magnetometers
 - Nuclear magnetic resonance
 - Torsion pendula



QCD axion resolves strong CP problem

Pseudoscalars (Axions): $\varphi \xrightarrow{P} - \varphi$

→ Time-varying spindependent effects

- Co-magnetometers
- Nuclear magnetic resonance
 - Torsion pendula

"Axion Wind" Spin-Precession Effect

[Flambaum, talk at Patras Workshop, 2013], [Graham, Rajendran, PRD 88, 035023 (2013)], [Stadnik, Flambaum, PRD 89, 043522 (2014)]

 $D_{-}(f)$

725

$$\mathcal{L}_{aff} = -\frac{C_f}{2f_a} \partial_i [a_0 \cos(\varepsilon_a t - p_a \cdot x)] \bar{f} \gamma^i \gamma^5 f$$

$$=> H_{\text{eff}}(t) \simeq \sigma_f \cdot B_{\text{eff}} \sin(m_a t)$$

$$Pseudo-\text{magnetic field}^*$$

$$B_{\text{eff}} \propto v$$

* Compare with usual magnetic field: $H = -\mu_f \cdot B$

Nucleons: [Graham, Rajendran, *PRD* 84, 055013 (2011)] Atoms and molecules: [Stadnik, Flambaum, *PRD* 89, 043522 (2014)]

Electric Dipole Moment (EDM) = parity (P) and time-

reversal-invariance (T) violating electric moment



Nucleons: [Graham, Rajendran, *PRD* **84**, 055013 (2011)] Atoms and molecules: [Stadnik, Flambaum, *PRD* **89**, 043522 (2014)]

$$\mathcal{L}_{aGG} = \frac{C_G a_0 \cos(m_a t)}{f_a} \frac{g^2}{32\pi^2} G^a_{\mu\nu} \tilde{G}^{a\mu\nu}$$

Nucleons: [Graham, Rajendran, *PRD* **84**, 055013 (2011)] Atoms and molecules: [Stadnik, Flambaum, *PRD* **89**, 043522 (2014)]

$$\mathcal{L}_{aGG} = \frac{C_G a_0 \cos(m_a t)}{f_a} \frac{g^2}{32\pi^2} G^a_{\mu\nu} \tilde{G}^{a\mu\nu}$$





Nucleons: [Graham, Rajendran, *PRD* **84**, 055013 (2011)] Atoms and molecules: [Stadnik, Flambaum, *PRD* **89**, 043522 (2014)]

$$\mathcal{L}_{aGG} = \frac{C_G a_0 \cos(m_a t)}{f_a} \frac{g^2}{32\pi^2} G^a_{\mu\nu} \tilde{G}^{a\mu\nu}$$

Nucleon EDMs

CP-violating intranuclear forces



Nucleons: [Graham, Rajendran, *PRD* **84**, 055013 (2011)] Atoms and molecules: [Stadnik, Flambaum, *PRD* **89**, 043522 (2014)]

$$\mathcal{L}_{aGG} = \frac{C_G a_0 \cos(m_a t)}{f_a} \frac{g^2}{32\pi^2} G^a_{\mu\nu} \tilde{G}^{a\mu\nu}$$

Nucleon EDMs

CP-violating intranuclear forces



In nuclei, <u>tree-level</u> *CP*-violating intranuclear forces dominate over <u>loop-induced</u> nucleon EDMs (loop factor = $1/(8\pi^2)$).

Proposals: [Flambaum, talk at *Patras Workshop*, 2013; Stadnik, Flambaum, *PRD* **89**, 043522 (2014); arXiv:1511.04098; Stadnik, PhD Thesis (2017)]

Use *spin-polarised sources*: Atomic magnetometers, ultracold neutrons, torsion pendula

Proposals: [Flambaum, talk at *Patras Workshop*, 2013; Stadnik, Flambaum, *PRD* **89**, 043522 (2014); arXiv:1511.04098; Stadnik, PhD Thesis (2017)]

Use *spin-polarised sources*: Atomic magnetometers, ultracold neutrons, torsion pendula

Experiment (n/Hg): [nEDM collaboration, PRX 7, 041034 (2017)]

$$\frac{\nu_n}{\nu_{\rm Hg}} = \left| \frac{\mu_n B}{\mu_{\rm Hg} B} \right| + R(t)$$

$$\uparrow \qquad \uparrow \qquad \uparrow$$
Energy
$$\uparrow \qquad \sigma \qquad B \qquad \qquad B-{\rm field} \quad {\rm Axion \ DM} \\ effect \qquad effect$$

Proposals: [Flambaum, talk at *Patras Workshop*, 2013; Stadnik, Flambaum, *PRD* **89**, 043522 (2014); arXiv:1511.04098; Stadnik, PhD Thesis (2017)]

Use *spin-polarised sources*: Atomic magnetometers, ultracold neutrons, torsion pendula

Experiment (n/Hg): [nEDM collaboration, PRX 7, 041034 (2017)]

$$\frac{\nu_n}{\nu_{\rm Hg}} = \left| \frac{\mu_n R}{\mu_{\rm Hg} R} \right| + R(t)$$

$$\uparrow \qquad \uparrow \qquad \uparrow$$
Energy
$$\uparrow \qquad \sigma \qquad B$$

$$= field \quad Axion \, DM \\ effect \qquad effect$$

Proposals: [Flambaum, talk at *Patras Workshop*, 2013; Stadnik, Flambaum, *PRD* **89**, 043522 (2014); arXiv:1511.04098; Stadnik, PhD Thesis (2017)]

Use *spin-polarised sources*: Atomic magnetometers, ultracold neutrons, torsion pendula

Experiment (n/Hg): [nEDM collaboration, PRX 7, 041034 (2017)]

$$\frac{\nu_n}{\nu_{\rm Hg}} = \left| \frac{\mu_n R}{\mu_{\rm Hg} R} \right| + R(t) \qquad \qquad \mathbf{E} \quad \boldsymbol{\sigma} \quad \mathbf{B}$$
$$R_{\rm EDM}(t) \propto \cos(m_a t) \qquad \qquad \mathbf{f} \quad \mathbf{f} \quad$$

Proposals: [Flambaum, talk at *Patras Workshop*, 2013; Stadnik, Flambaum, *PRD* **89**, 043522 (2014); arXiv:1511.04098; Stadnik, PhD Thesis (2017)]

Use *spin-polarised sources*: Atomic magnetometers, ultracold neutrons, torsion pendula

Experiment (n/Hg): [nEDM collaboration, PRX 7, 041034 (2017)]

$$\frac{\nu_n}{\nu_{\rm Hg}} = \left| \frac{\mu_n R}{\mu_{\rm Hg} R} \right| + R(t)$$

$$R_{\rm EDM}(t) \propto \cos(m_a t)$$

$$R_{\rm wind}(t) \propto \sum_{i=1,2,3} A_i \sin(\omega_i t)$$

$$B_{\rm eff}$$

$$= m_a, \ \omega_2 = m_a + \Omega_{\rm sidereal}, \ \omega_3 = |m_a - \Omega_{\rm sidereal}|$$

$$Earth's rotation$$

 ω_1

Proposals: [CASPEr collaboration, Quantum Sci. Technol. 3, 014008 (2018)]

Use nuclear magnetic resonance ("sidebands" technique)

Experiment (Formic acid): [CASPEr collaboration, In preparation]



$$H_J \sim J I_H \cdot I_C$$

Experiment (Formic acid): [CASPEr collaboration, In preparation]



 $H_{I} \sim J I_{H} \cdot I_{C}$

• J-coupling only: H



Frequency [Hz]

Experiment (Formic acid): [CASPEr collaboration, In preparation]



Signal amplitude [a.u]

Experiment (Formic acid): [CASPEr collaboration, In preparation]



Proposals: [Budker, Graham, Ledbetter, Rajendran, A. O. Sushkov, PRX 4, 021030 (2014)]

Use nuclear magnetic resonance

Proposals: [Budker, Graham, Ledbetter, Rajendran, A. O. Sushkov, PRX 4, 021030 (2014)]

Use nuclear magnetic resonance

Traditional NMR



Resonance: $2\mu B_{ext} = \omega$

Proposals: [Budker, Graham, Ledbetter, Rajendran, A. O. Sushkov, PRX 4, 021030 (2014)]

Use nuclear magnetic resonance



Resonance: $2\mu B_{ext} = \omega$

Resonance: $2\mu B_{ext} \approx m_a$ Measure transverse magnetisation

Experiments

Co-magnetometry: $10^{-23} \text{ eV} < m_a < 10^{-17} \text{ eV}$

• n/Hg (PSI): [nEDM collaboration, *PRX* 7, 041034 (2017)]

"Sidebands" NMR: $10^{-16} \text{ eV} < m_a < 10^{-13} \text{ eV}$

• Formic acid (Mainz): [CASPEr collaboration, In preparation]

"Normal" NMR: $10^{-14} \text{ eV} < m_a < 10^{-7} \text{ eV}$

- Liquid Xe (Mainz)
- Pb in ferroelectric medium (Boston)

Constraints on Interaction of Axion Dark Matter with Gluons

nEDM constraints: [nEDM collaboration, PRX 7, 041034 (2017)]



3 orders of magnitude improvement!

Constraints on Interaction of Axion Dark Matter with Nucleons

v_n/v_{Hg} constraints: [nEDM collaboration, PRX 7, 041034 (2017)]

Formic acid NMR constraints (preliminary): [CASPEr collaboration, In preparation]

2 orders of magnitude improvement (laboratory bounds)!





\rightarrow Time-varying

fundamental constants

- Atomic clocks
- Optical cavities
- Fifth-force searches
- Astrophysics (e.g., BBN)

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)], [Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)], [Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

$$\mathcal{L}_f = -\frac{\phi^2}{(\Lambda'_f)^2} m_f \bar{f} f$$

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)], [Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

$$\mathcal{L}_f = -\frac{\phi^2}{(\Lambda'_f)^2} m_f \bar{f} f$$
 c.f. $\mathcal{L}_f^{SM} = -m_f \bar{f} f$

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)], [Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

$$\mathcal{L}_f = -\frac{\phi^2}{(\Lambda'_f)^2} m_f \bar{f} f \quad \text{c.f.} \quad \mathcal{L}_f^{\text{SM}} = -m_f \bar{f} f \quad => \quad m_f \to m_f \left[1 + \frac{\phi^2}{(\Lambda'_f)^2} \right]$$

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)], [Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

$$\mathcal{L}_f = -\frac{\phi^2}{(\Lambda'_f)^2} m_f \bar{f} f \quad \text{c.f.} \quad \mathcal{L}_f^{\text{SM}} = -m_f \bar{f} f \quad => \quad m_f \to m_f \left[1 + \frac{\phi^2}{(\Lambda'_f)^2} \right]$$
$$=> \frac{\delta m_f}{m_f} = \frac{\phi_0^2}{(\Lambda'_f)^2} \cos^2(m_\phi t)$$

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)], [Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

$$\mathcal{L}_{f} = -\frac{\phi^{2}}{(\Lambda'_{f})^{2}} m_{f} \bar{f} f \quad \text{c.f.} \quad \mathcal{L}_{f}^{\text{SM}} = -m_{f} \bar{f} f \quad => \quad m_{f} \to m_{f} \left[1 + \frac{\phi^{2}}{(\Lambda'_{f})^{2}} \right]$$
$$=> \frac{\delta m_{f}}{m_{f}} = \frac{\phi_{0}^{2}}{(\Lambda'_{f})^{2}} \cos^{2}(m_{\phi} t) = \left[\frac{\phi_{0}^{2}}{2(\Lambda'_{f})^{2}} + \frac{\phi_{0}^{2}}{2(\Lambda'_{f})^{2}} \cos(2m_{\phi} t) \right]$$

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)], [Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

$$\mathcal{L}_{f} = -\frac{\phi^{2}}{(\Lambda_{f}')^{2}} m_{f} \bar{f} f \quad \text{c.f.} \quad \mathcal{L}_{f}^{\text{SM}} = -m_{f} \bar{f} f \quad => \quad m_{f} \to m_{f} \left[1 + \frac{\phi^{2}}{(\Lambda_{f}')^{2}} \right]$$
$$= > \frac{\delta m_{f}}{m_{f}} = \frac{\phi_{0}^{2}}{(\Lambda_{f}')^{2}} \cos^{2}(m_{\phi}t) = \left[\frac{\phi_{0}^{2}}{2(\Lambda_{f}')^{2}} + \frac{\phi_{0}^{2}}{2(\Lambda_{f}')^{2}} \cos(2m_{\phi}t) \right]$$
$$\rho_{\phi} = \frac{m_{\phi}^{2}\phi_{0}^{2}}{2} \implies \phi_{0}^{2} \propto \rho_{\phi}$$

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)], [Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

$$\mathcal{L}_{f} = -\frac{\phi^{2}}{(\Lambda'_{f})^{2}} m_{f} \bar{f} f \quad \text{c.f.} \quad \mathcal{L}_{f}^{\text{SM}} = -m_{f} \bar{f} f \quad => \quad m_{f} \rightarrow m_{f} \left[1 + \frac{\phi^{2}}{(\Lambda'_{f})^{2}} \right] \\ => \frac{\delta m_{f}}{m_{f}} = \frac{\phi_{0}^{2}}{(\Lambda'_{f})^{2}} \cos^{2}(m_{\phi}t) = \underbrace{\frac{\phi_{0}^{2}}{2(\Lambda'_{f})^{2}}}_{2(\Lambda'_{f})^{2}} + \underbrace{\frac{\phi_{0}^{2}}{2(\Lambda'_{f})^{2}}}_{2(\Lambda'_{f})^{2}} \cos(2m_{\phi}t) \\ \text{'Slow' drifts [Astrophysics (high ρ_{DM}): BBN, CMB]}$$

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)], [Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

$$\mathcal{L}_{f} = -\frac{\phi^{2}}{(\Lambda_{f}')^{2}} m_{f} \bar{f} f \quad \text{c.f.} \quad \mathcal{L}_{f}^{\text{SM}} = -m_{f} \bar{f} f \quad => \quad m_{f} \rightarrow m_{f} \left[1 + \frac{\phi^{2}}{(\Lambda_{f}')^{2}} \right]$$
$$= > \frac{\delta m_{f}}{m_{f}} = \frac{\phi_{0}^{2}}{(\Lambda_{f}')^{2}} \cos^{2}(m_{\phi}t) = \underbrace{\frac{\phi_{0}^{2}}{2(\Lambda_{f}')^{2}}}_{2(\Lambda_{f}')^{2}} + \underbrace{\frac{\phi_{0}^{2}}{2(\Lambda_{f}')^{2}}}_{2(\Lambda_{f}')^{2}} \cos(2m_{\phi}t)$$
$$\text{'Slow' drifts [Astrophysics} (high \rho_{\text{DM}}): \text{BBN, CMB]}$$
$$+ \text{Gradients [Fifth-forces]}$$

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)], [Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]



Atomic Spectroscopy Searches for Oscillating Variations in Fundamental Constants due to Dark Matter

[Arvanitaki, Huang, Van Tilburg, PRD 91, 015015 (2015)], [Stadnik, Flambaum, PRL 114, 161301 (2015)]

$$\frac{\delta\left(\omega_{1}/\omega_{2}\right)}{\omega_{1}/\omega_{2}} \propto \sum_{X=\alpha, m_{e}/m_{p}, \dots} \begin{pmatrix} K_{X,1} - K_{X,2} \end{pmatrix} \cos\left(\omega t\right)$$

$$\uparrow \qquad \uparrow \qquad f$$
Sensitivity coefficients

 $\omega = m_{\varphi}$ (linear coupling) or $\omega = 2m_{\varphi}$ (quadratic coupling)

Atomic Spectroscopy Searches for Oscillating Variations in Fundamental Constants due to Dark Matter

[Arvanitaki, Huang, Van Tilburg, PRD 91, 015015 (2015)], [Stadnik, Flambaum, PRL 114, 161301 (2015)]

$$\frac{\delta\left(\omega_{1}/\omega_{2}\right)}{\omega_{1}/\omega_{2}} \propto \sum_{X=\alpha, m_{e}/m_{p}, \dots} \begin{pmatrix} K_{X,1} - K_{X,2} \end{pmatrix} \cos\left(\omega t\right)$$

$$A = \alpha, m_{e}/m_{p}, \dots \qquad A = \alpha, m_{e}/m_{p}, \dots$$
Sensitivity coefficients

 $\omega = m_{\varphi}$ (linear coupling) or $\omega = 2m_{\varphi}$ (quadratic coupling)

Sensitivity coefficients K_X calculated extensively by Flambaum group and co-workers (1998 – present), see the reviews
 [Flambaum, Dzuba, *Can. J. Phys.* 87, 25 (2009); *Hyperfine Interac.* 236, 79 (2015)]

Atomic Spectroscopy Searches for Oscillating Variations in Fundamental Constants due to Dark Matter

[Arvanitaki, Huang, Van Tilburg, PRD 91, 015015 (2015)], [Stadnik, Flambaum, PRL 114, 161301 (2015)]

$$\frac{\delta\left(\omega_{1}/\omega_{2}\right)}{\omega_{1}/\omega_{2}} \propto \sum_{X=\alpha, m_{e}/m_{p}, \dots} \begin{pmatrix} K_{X,1} - K_{X,2} \end{pmatrix} \cos\left(\omega t\right)$$

$$\uparrow \qquad \uparrow \qquad \uparrow \qquad \uparrow \qquad \uparrow \qquad \uparrow \qquad f$$
Sensitivity coefficients

 $\omega = m_{\varphi}$ (linear coupling) or $\omega = 2m_{\varphi}$ (quadratic coupling)

- Sensitivity coefficients K_X calculated extensively by Flambaum group and co-workers (1998 present), see the reviews
 [Flambaum, Dzuba, Can. J. Phys. 87, 25 (2009); Hyperfine Interac. 236, 79 (2015)]
- Precision of optical clocks approaching ~10⁻¹⁸ fractional level



Gravitational-wave detector (LIGO/Virgo), *L* ~ 4 km



Small-scale cavity, $L \sim 0.2 \text{ m}$

• Compare $L \sim Na_{\rm B}$ with λ

- Compare $L \sim Na_{\rm B}$ with λ
- For a "usual" atomic optical transition and in the nonrelativistic limit:*

$$\Phi = \frac{\omega L}{c} \propto \left(\frac{e^2}{a_{\rm B}\hbar}\right) \left(\frac{Na_{\rm B}}{c}\right) = N\alpha \implies \frac{\delta\Phi}{\Phi} \approx \frac{\delta\alpha}{\alpha}$$

* For numerical calculations, including (small) relativistic effects, see [Pasteka, Hao, Borschevsky, Flambaum, Schwerdtfeger, arXiv:1809.02863].

- Compare $L \sim Na_{\rm B}$ with λ
- For a "usual" atomic optical transition and in the nonrelativistic limit:*

$$\Phi = \frac{\omega L}{c} \propto \left(\frac{e^2}{a_{\rm B}\hbar}\right) \left(\frac{Na_{\rm B}}{c}\right) = N\alpha \implies \frac{\delta\Phi}{\Phi} \approx \frac{\delta\alpha}{\alpha}$$

• Multiple reflections of light beam enhance the effect $(N_{\rm eff} \sim 10^5 \text{ in small-scale interferometers with highly reflective mirrors; c.f. <math>N_{\rm eff} \sim 100$ in LIGO/Virgo)

* For numerical calculations, including (small) relativistic effects, see [Pasteka, Hao, Borschevsky, Flambaum, Schwerdtfeger, arXiv:1809.02863].

Experiments

Clock/clock comparisons: $10^{-23} \text{ eV} < m_{\varphi} < 10^{-16} \text{ eV}$

- Dy/Cs (Mainz): [Van Tilburg *et al.*, *PRL* 115, 011802 (2015)], [Stadnik, Flambaum, *PRL* 115, 201301 (2015)]
 - Rb/Cs (SYRTE): [Hees *et al.*, *PRL* 117, 061301 (2016)], [Stadnik, Flambaum, *PRA* 94, 022111 (2016)]
- Rb/Cs (GPS network)*: [Roberts et al., Nature Commun. 8, 1195 (2017)]
- Al⁺/Yb, Yb/Sr, Al⁺/Hg⁺ (NIST): [Hume, Leibrandt, Wineland *et al.*, In prep.]
 - Yb⁺(*E*3)/Sr (PTB): [Huntemann, Peik *et al.*, In preparation]

Clock/cavity comparisons: $10^{-20} \text{ eV} < m_{\varphi} < 10^{-15} \text{ eV}$

- Sr/ULE cavity (Torun)*: [Wcislo et al., Nature Astronomy 1, 0009 (2016)]
 - Sr/Si cavity (JILA): [Robinson, Ye et al., In preparation]
- * Searches for domain wall dark matter.

Constraints on Linear Interaction of Scalar Dark Matter with the Photon

Clock/clock (DM) constraints: [Van Tilburg *et al.*, *PRL* **115**, 011802 (2015)], [Hees *et al.*, *PRL* **117**, 061301 (2016)]; Clock/clock (fifth force) constraints: [Leefer *et al.*, *PRL* **117**, 271601 (2016)]

4 orders of magnitude improvement!



Constraints on Quadratic Interaction of Scalar Dark Matter with the Photon

Clock/clock + BBN constraints: [Stadnik, Flambaum, *PRL* **115**, 201301 (2015); *PRA* **94**, 022111 (2016)]; MICROSCOPE + Eöt-Wash constraints: [Hees *et al.*, *PRD* **98**, 064051 (2018)]

15 orders of magnitude improvement!



Summary

New classes of dark matter effects that are

first power in the underlying interaction constant

=> Up to 15 orders of magnitude improvement

with experiments based on magnetic resonance and atomic clock spectroscopy

Back-up Slides

Linear vs Quadratic Couplings [Hees, Minazzoli, Savalle, Stadnik, Wolf, PRD 98, 064051 (2018)] Consider the effect of a massive body (e.g., Earth) on the scalar DM field:



Gradients + screening/amplification



Gradients + screening/amplification

Constraints on Linear Interaction of Scalar Dark Matter with the Higgs Boson

Rb/Cs constraints:

[Stadnik, Flambaum, PRA 94, 022111 (2016)]

2 – 3 orders of magnitude improvement!



BBN Constraints on 'Slow' Drifts in Fundamental Constants due to Dark Matter [Stadnik, Flambaum, PRL 115, 201301 (2015)]

- Largest effects of DM in early Universe (highest $\rho_{\rm DM}$)
- Big Bang nucleosynthesis ($t_{weak} \approx 1s t_{BBN} \approx 3 min$)
- Primordial ⁴He abundance sensitive to *n/p* ratio (almost all neutrons bound in ⁴He after BBN)

$$\frac{\Delta Y_p(^{4}\text{He})}{Y_p(^{4}\text{He})} \approx \frac{\Delta (n/p)_{\text{weak}}}{(n/p)_{\text{weak}}} - \Delta \left[\int_{t_{\text{weak}}}^{t_{\text{BBN}}} \Gamma_n(t) dt \right]$$

$$p + e^- \rightleftharpoons n + \nu_e$$

$$n + e^+ \rightleftharpoons p + \bar{\nu}_e$$

$$n \to p + e^- + \bar{\nu}_e$$

Back-Reaction Effects in BBN

[Sörensen, Sibiryakov, Yu, PRELIMINARY – In preparation]

