



Impact of Octupole Deformation on the Nuclear Electromagnetic Response

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Octupole deformation and electromagnetic response

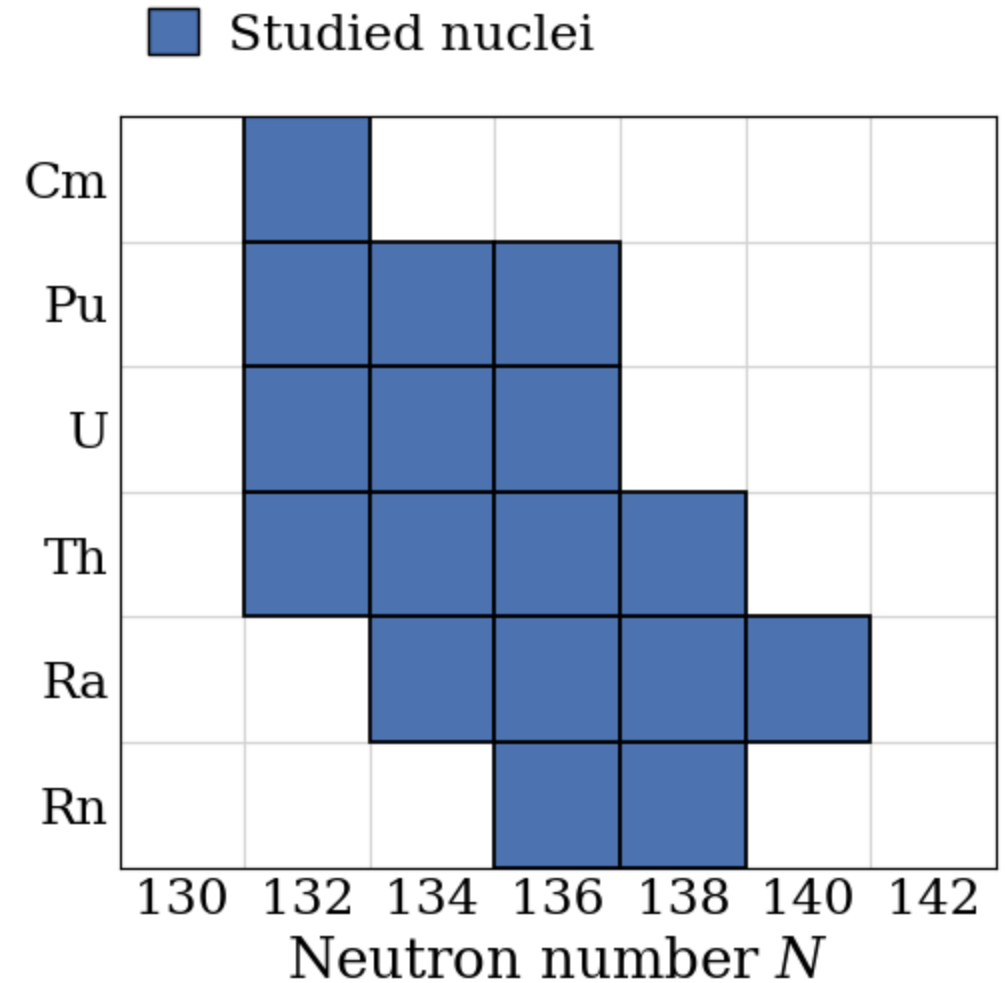


- Electromagnetic response
 - Provides a sensitive probe of nuclear structure
 - Collective excitations are sensitive to nuclear deformation
- Reflection-asymmetric (octupole-deformed) shapes are predicted in certain heavy nuclei
 - Strong experimental indications exist for such nuclei
- Octupole deformation breaks intrinsic parity symmetry
 - How does octupole deformation affect electromagnetic response across different multipole modes?

Scope of the study



- Focus on heavy, unstable nuclei with octupole-deformed ground states
 - Even-even isotopes with $Z = 86-96$ ($A = 222-230$)
- Electromagnetic response:
 - E1, M1, E2 (IS/IV), E3 (IS/IV)
 - Excitation energy range: 0–38 MeV

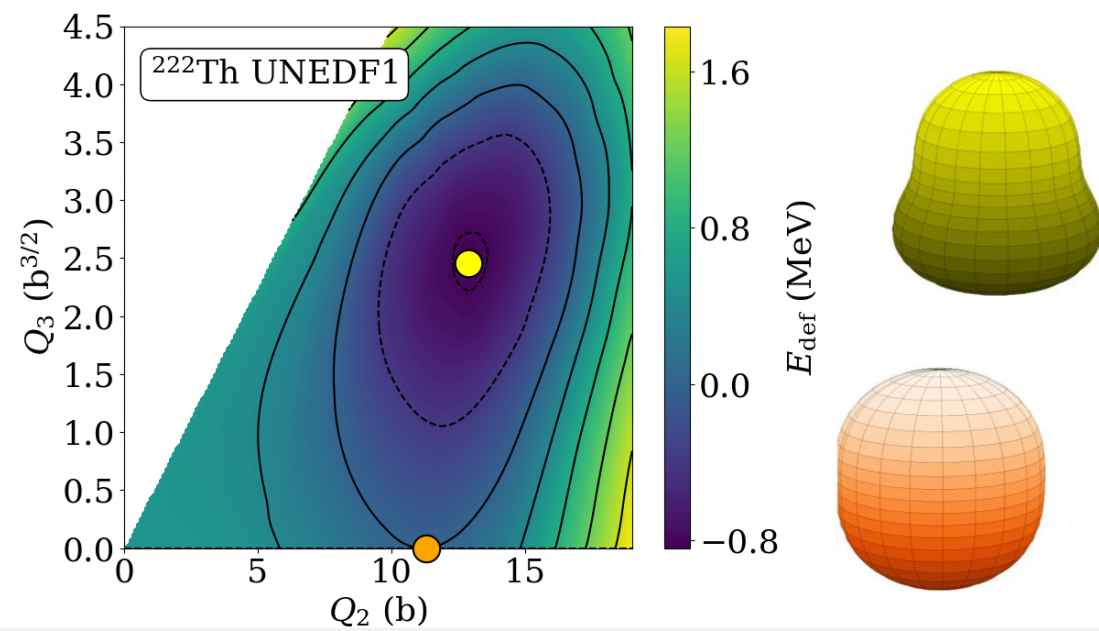


Based on:

M. Kanerva *et al.* Phys. Rev. C 113, 044320 (2026) – DOI: 10.1103/ddl7-y8xj (CC BY 4.0)

Ground states

- Framework: Skyrme–HFB
 - SkM*, SLy4, UNEDF1
- The potential energy surface is explored to identify two distinct ground state solutions:
 - One with conserved reflection symmetry (prolate, $Q_3 = 0$)
 - One with broken reflection symmetry (pear-shaped, $Q_3 \neq 0$)



Response to electromagnetic fields

- The obtained ground-state solutions are subjected to a polarizing external field $F(t, \omega)$
 - The linear response of HFB solution is determined by solving the FAM-QRPA equations:

$$\begin{aligned} (E_k + E_l - \omega)X_{lk}(\omega) + \delta H_{lk}^{20}(\omega) &= F_{lk}^{20}(\omega) \\ (E_k + E_l + \omega)Y_{lk}(\omega) + \delta H_{lk}^{02}(\omega) &= F_{lk}^{02}(\omega) \end{aligned}$$

Observables

- Transition strength

$$\begin{aligned} \frac{dB(\omega, \hat{F})}{d\omega} &= \sum_n |\langle \Phi_n | \hat{F}(\omega) | \Phi_0 \rangle|^2 \delta(\omega - E_n) \\ &= -\frac{1}{\pi} \text{Im Tr}[f(\omega)(UXV^T + V^*Y^T U^\dagger)] \end{aligned}$$

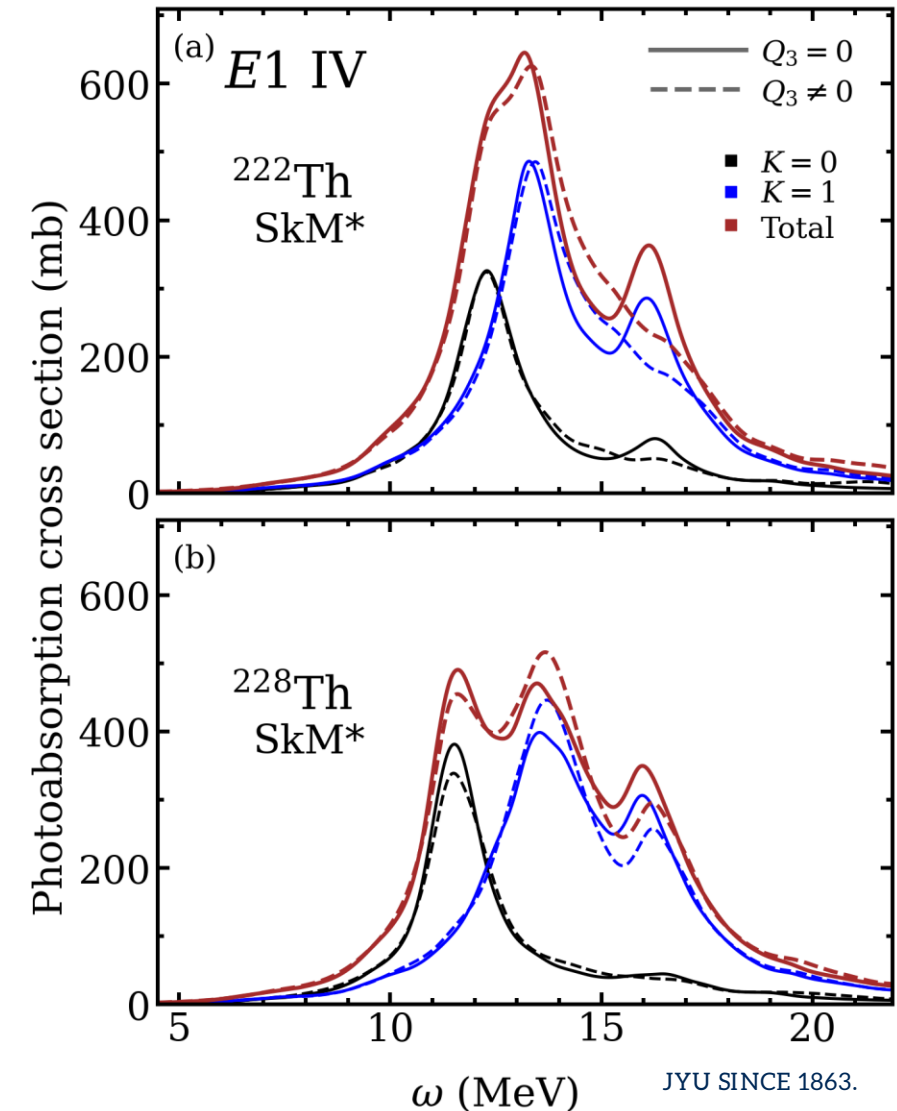
- Photo absorption cross-sections

$$\sigma_{\text{abs}}(\omega) = \frac{4\alpha\pi^2}{e^2} \omega \sum_{K=0,\pm 1} \frac{dB(\omega, \hat{F}_{E1}(K))}{d\omega}$$

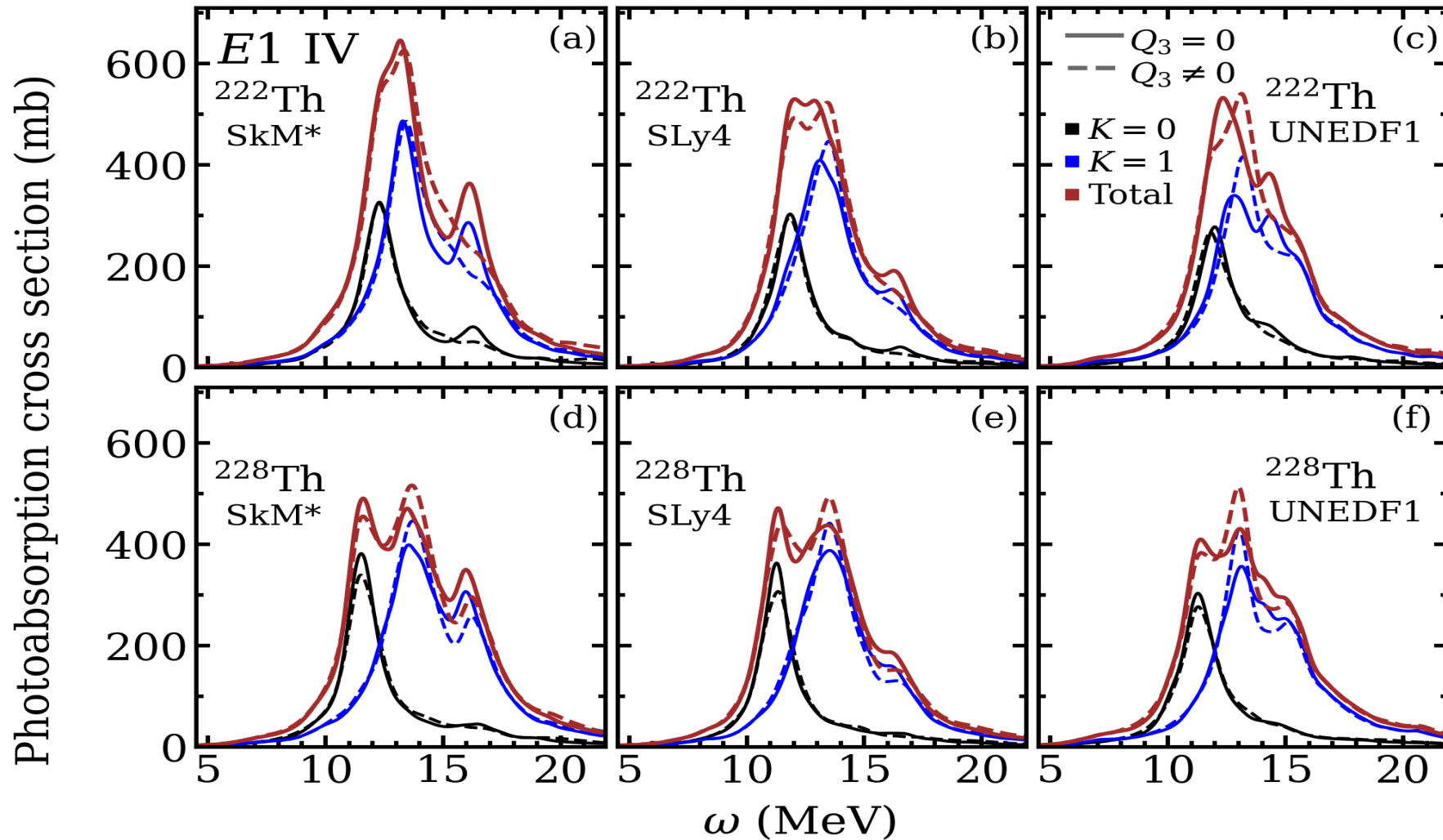
Electric dipole response



- Photoabsorption cross sections for two Th isotopes (^{222}Th and ^{228}Th) calculated with the SkM* functional
- Differences between the two deformed ground-state solutions are generally small
- A third peak around 16 MeV, prominent in SkM* results, is flattened in some octupole-deformed solutions
 - The suppression is not observed in all isotopes, suggesting that it is not a direct consequence of octupole deformation, but rather induced indirectly



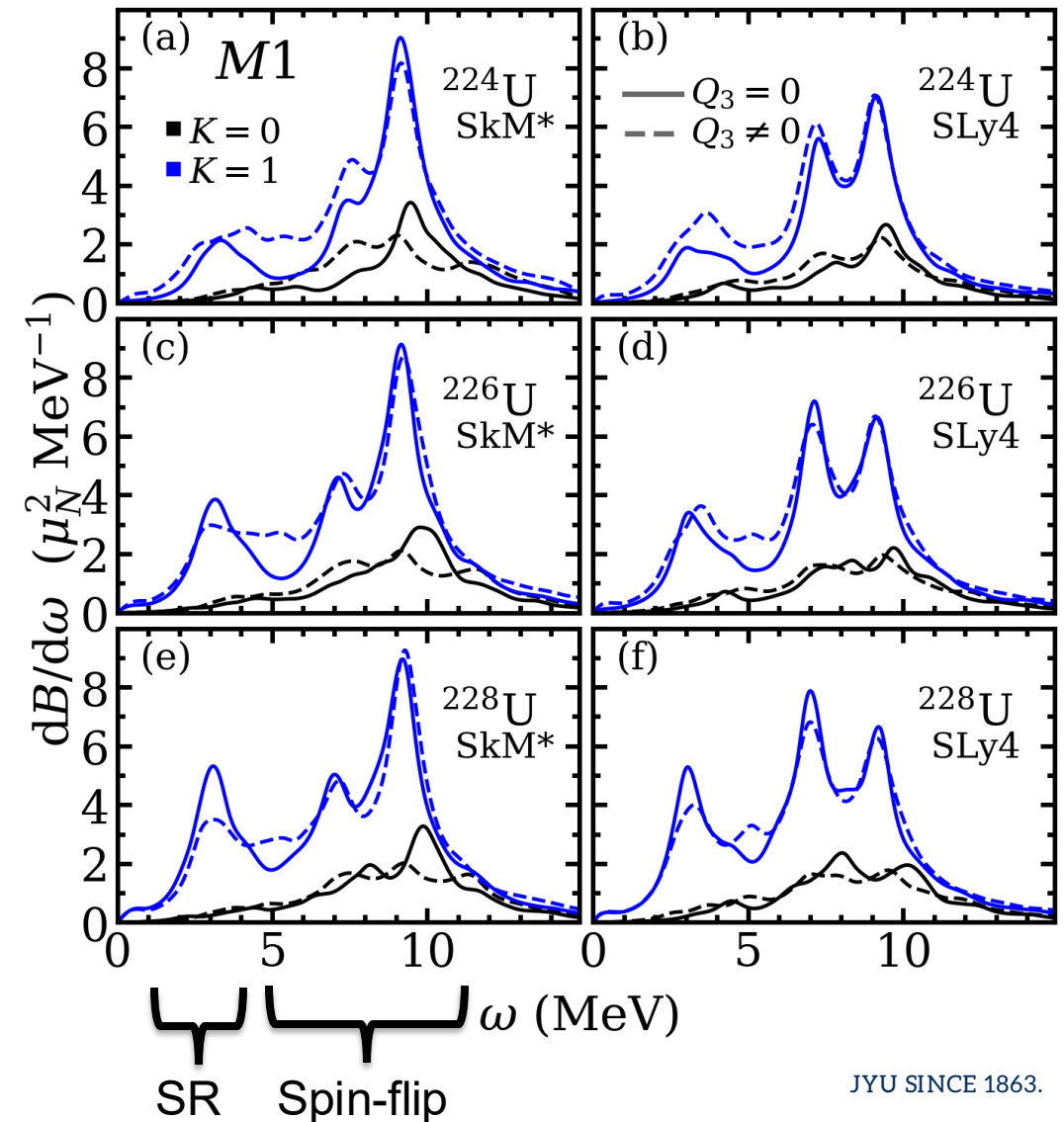
Electric dipole response – functional dependency



Magnetic dipole response



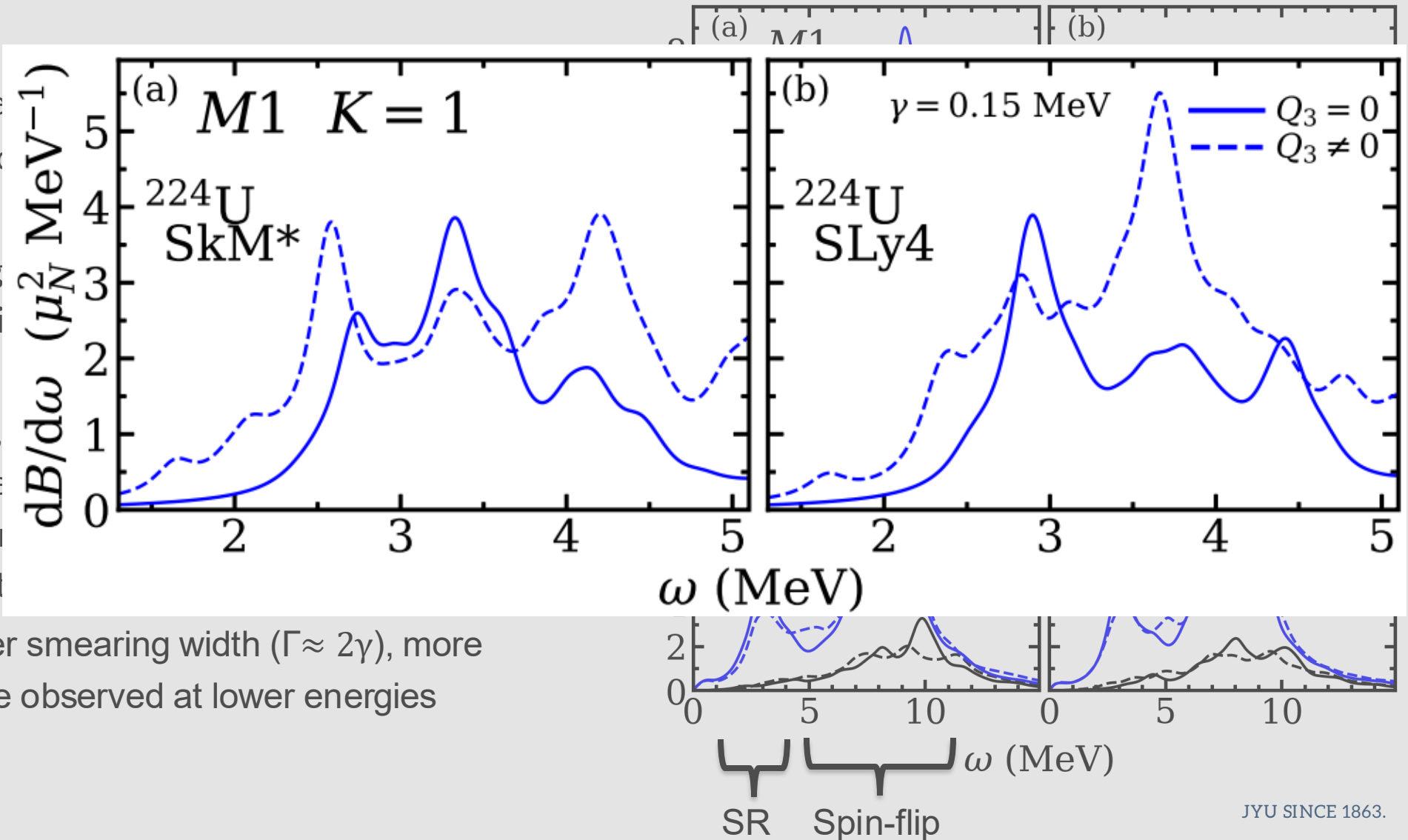
- $M1$ transition strength functions for three U isotopes (^{224}U , ^{226}U , ^{228}U) calculated with the SkM* and SLy4 functionals
- In the spin-flip region, the transition strength distributions are similar for different deformation minima
- At lower energies, especially around the expected scissors resonance (SR) (~ 3 MeV), noticeable differences appear between the two deformed ground-state solutions
 - With a smaller smearing width ($\Gamma \approx 2\gamma$), more details can be observed at lower energies



Magnetic dipole response



- M1 transition strength (^{224}U , ^{226}U , ^{228}U) calculated with Skyrme functionals
- In the spin-flip region, distributions are similar with minima
- At lower energies, scissors resonance differences appear between ground-state solutions

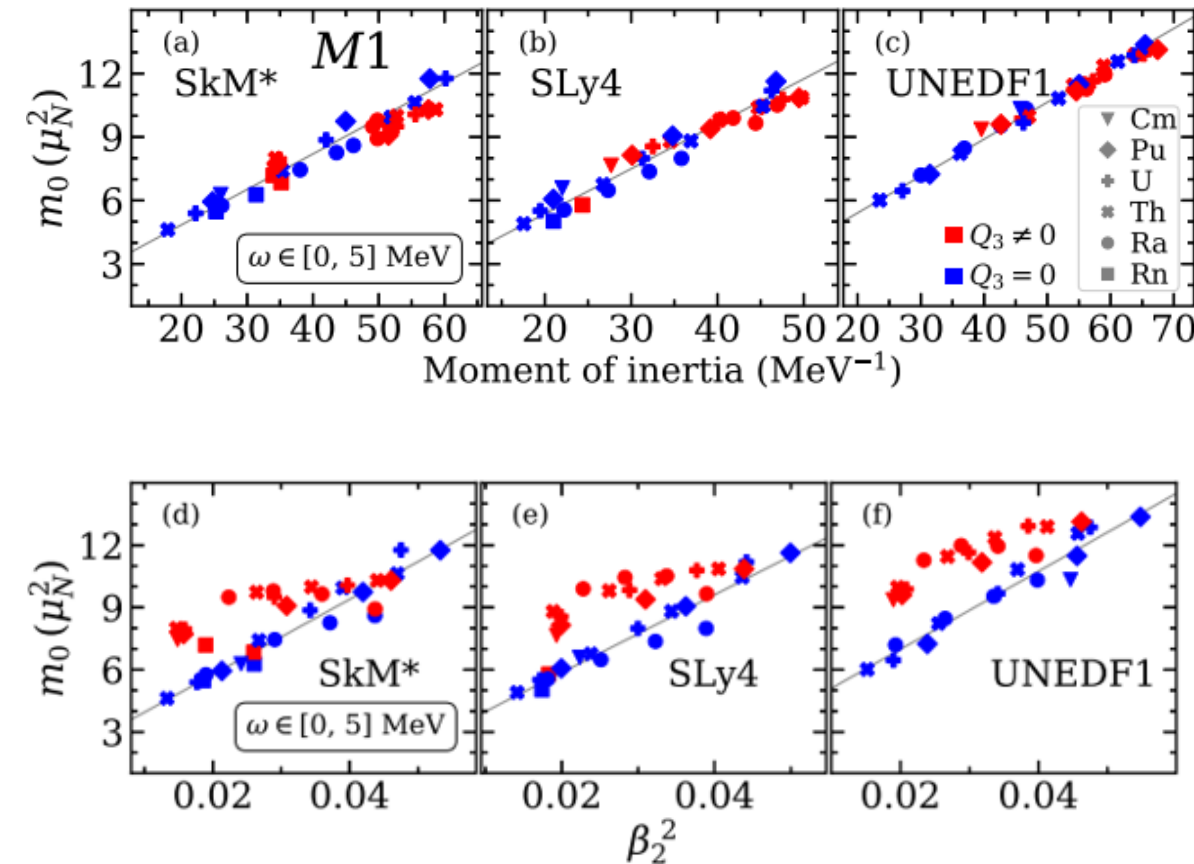


- With a smaller smearing width ($\Gamma \approx 2\gamma$), more details can be observed at lower energies

Magnetic dipole response – sum rules



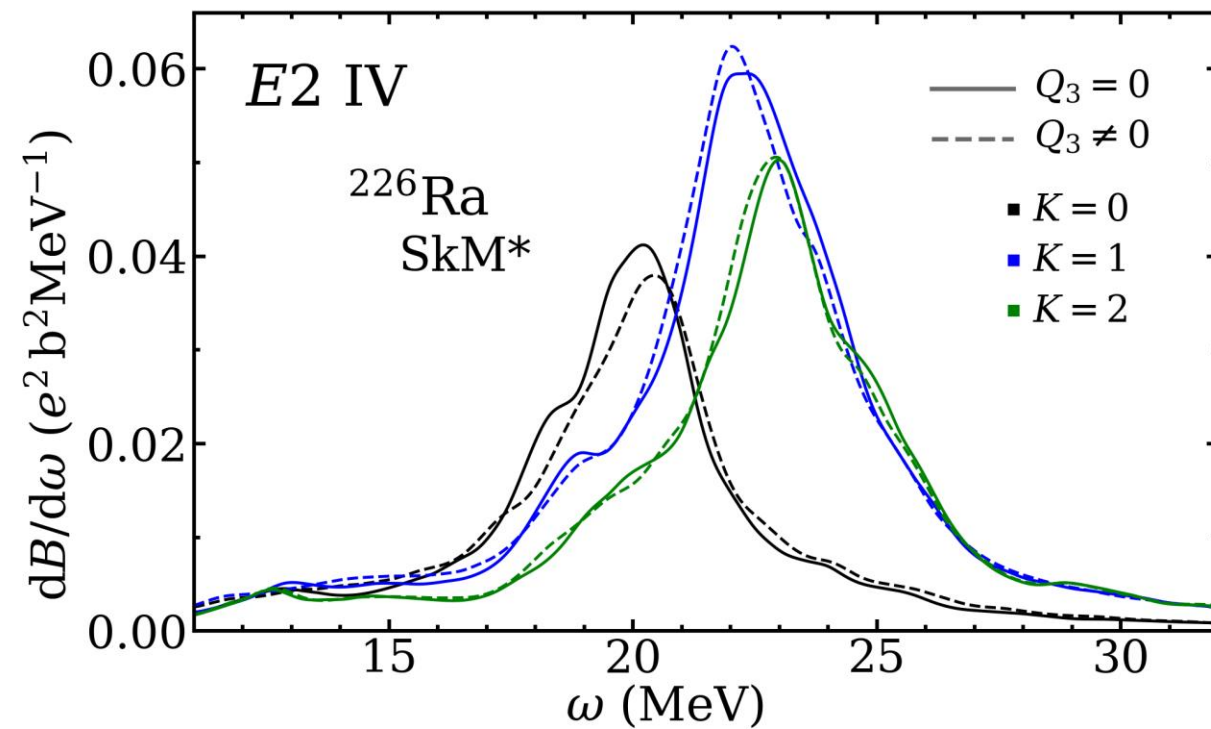
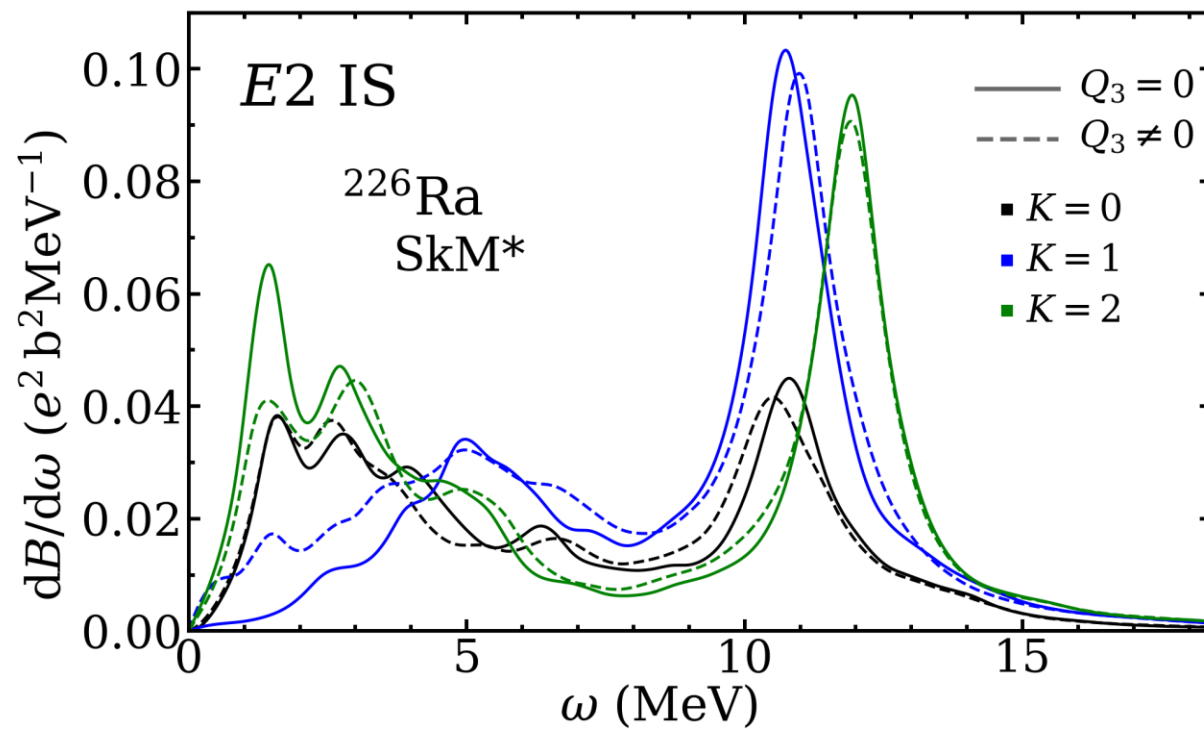
- Non-energy-weighted sum rule m_0 for low-energy transitions (0–5 MeV) in all studied nuclei with $A = 222–230$
- m_0 is expected to scale linearly with the moment of inertia
 - This trend is observed for both deformation minima (panels a–c)
 - Octupole-deformed solutions are concentrated at the upper end of the moment-of-inertia range
 - correspondingly larger m_0 values
- m_0 is also expected to scale with β_2^2
 - This correlation is weakened for octupole-deformed solutions (panels d–f)
 - could serve as a signature of octupole deformation in experiments



Electric quadrupole response



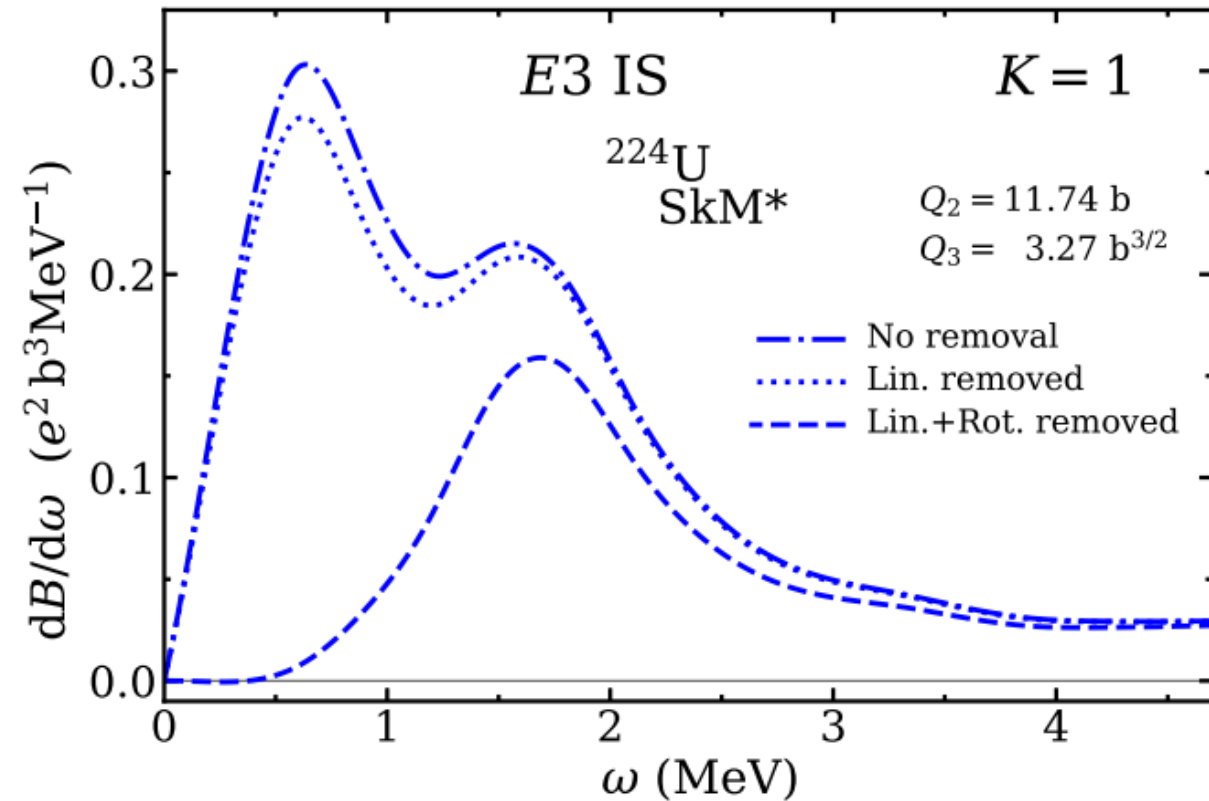
- Isoscalar (IS) and isovector (IV) $E2$ strength functions for ^{226}Ra calculated with the SkM* functional
 - The rotational spurious mode is removed in the IS $K=1$ channel
- Octupole deformation has little effect at resonance energies
- Differences appear at low energies (IS), which can be related to changes in the potential energy surface between the two deformation minima



Electric octupole response – rotational spurious mode



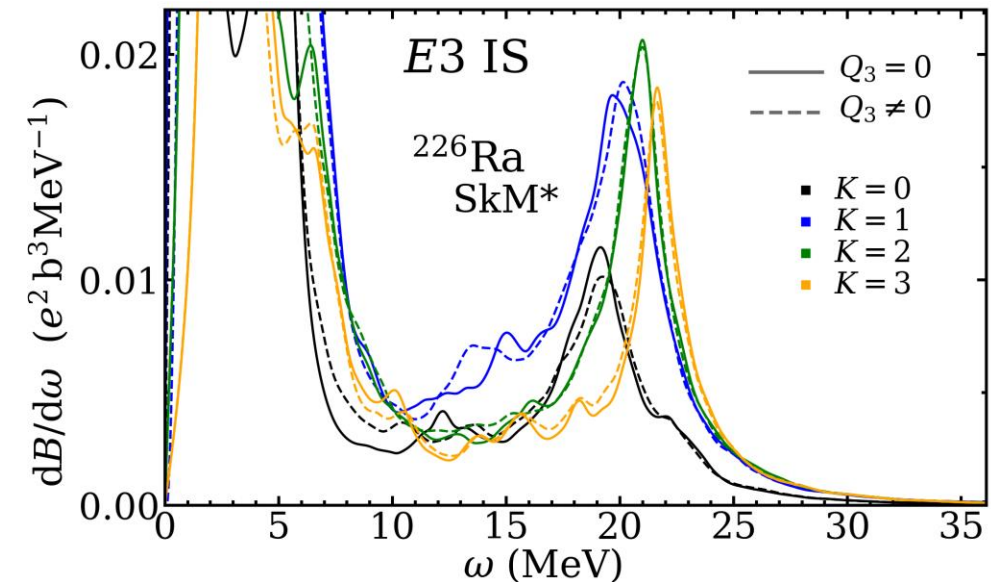
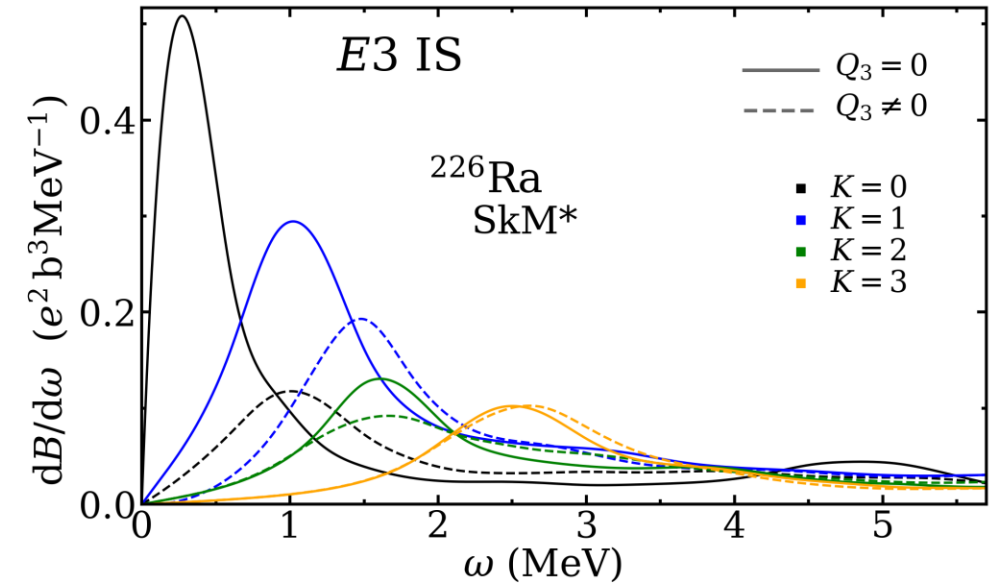
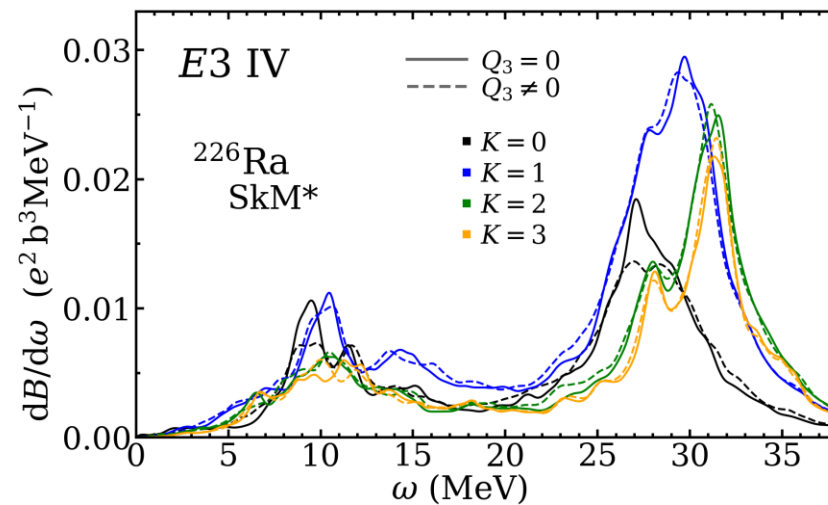
- Isoscalar $E3$ transition strength function calculated using **three different treatments of spurious contributions**:
 1. no removal
 2. removal of linear mode only
 3. removal of both linear and rotational modes
- The octupole-deformed solution exhibits a significant rotational spurious component due to broken parity symmetry
- This contribution must be removed in the IS $E3$ $K=1$ channel to obtain a physical strength distribution



Electric octupole response



- Isoscalar and isovector $E3$ strength functions for ^{226}Ra calculated with the SkM* functional
 - All relevant linear and rotational spurious contributions are removed in the $K=0$ and $K=1$ channels
- At resonance energies, octupole deformation has little effect on the strength distributions
- Differences appear at low energies in the IS channel
 - Again, these can be related to changes in the potential energy surfaces around the two deformation minima





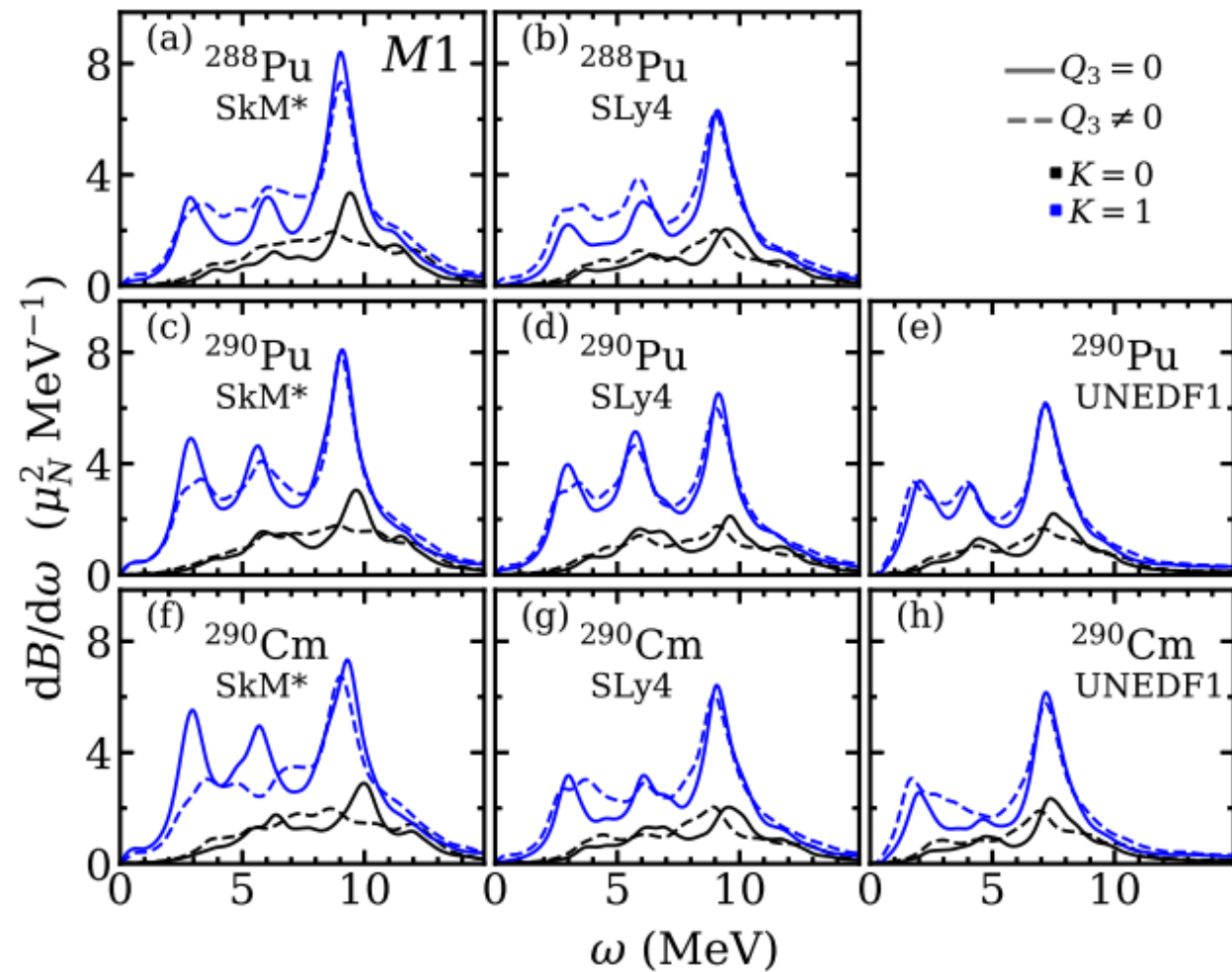
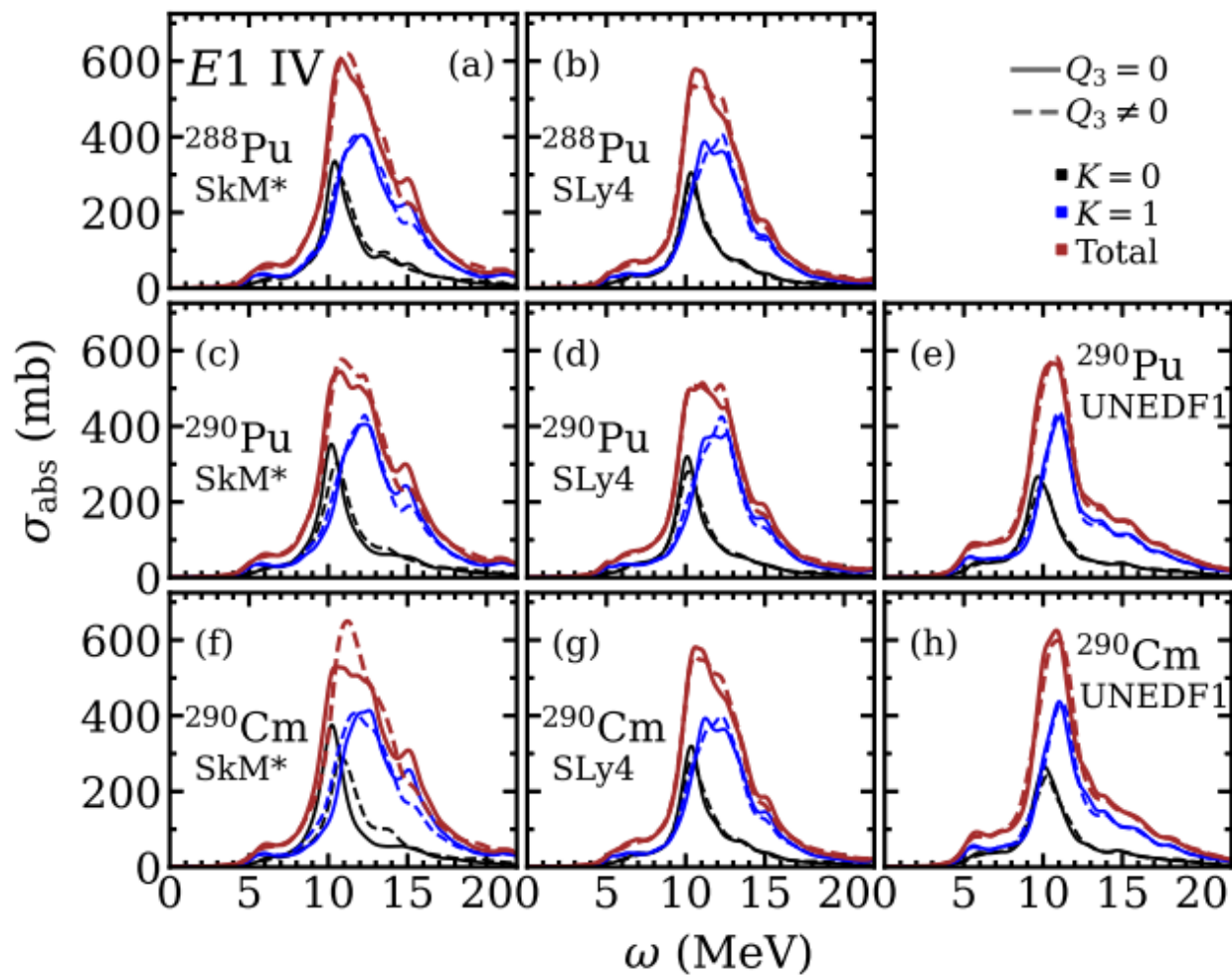
Conclusions

- **Octupole deformation** has a limited effect at resonance energies across all multipole modes
- **Low-energy M1 response** shows differences in strength distributions and total strength
- In parity-broken solutions, the rotational spurious mode must also be removed in the $K = 1^-$ channel (e.g., isoscalar E3)
- Observed trends are consistent across Skyrme functionals

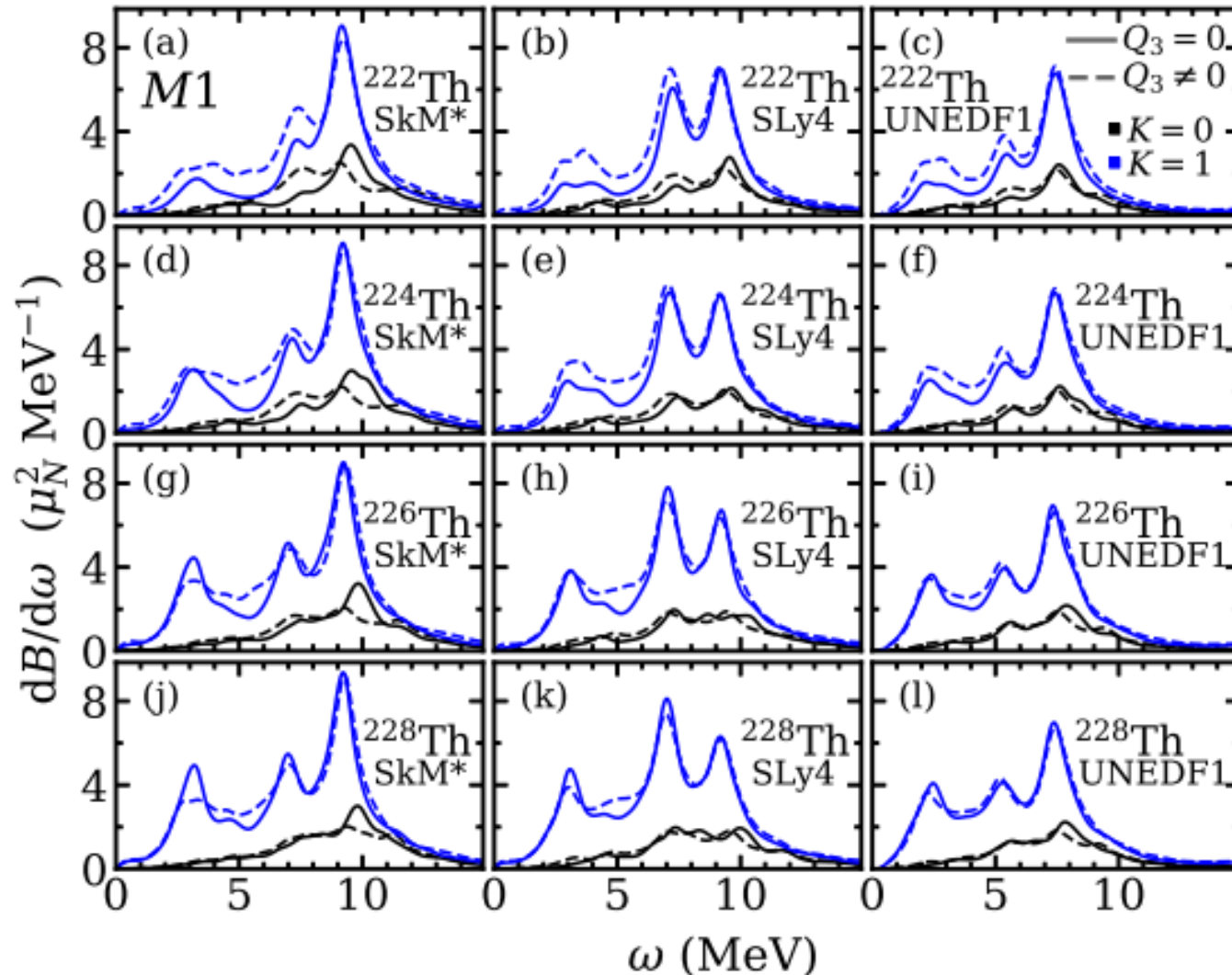
Outlook

- Open question: role of parity (and rotational) symmetry restoration
- Low-energy M1 excitation modes
 - Detailed role of octupole deformation in individual M1 modes
 - May provide an experimental signature of octupole deformation

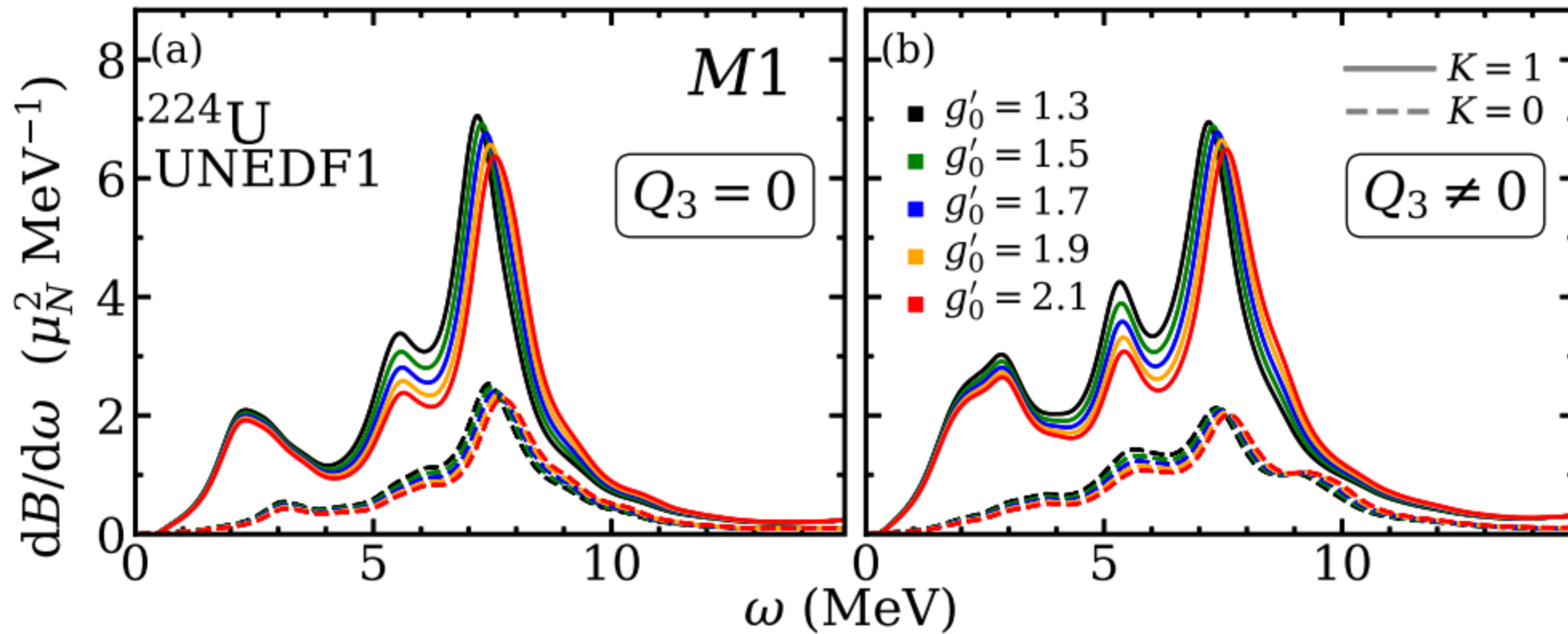
Electromagnetic response of $A = 288-290$ isotopes



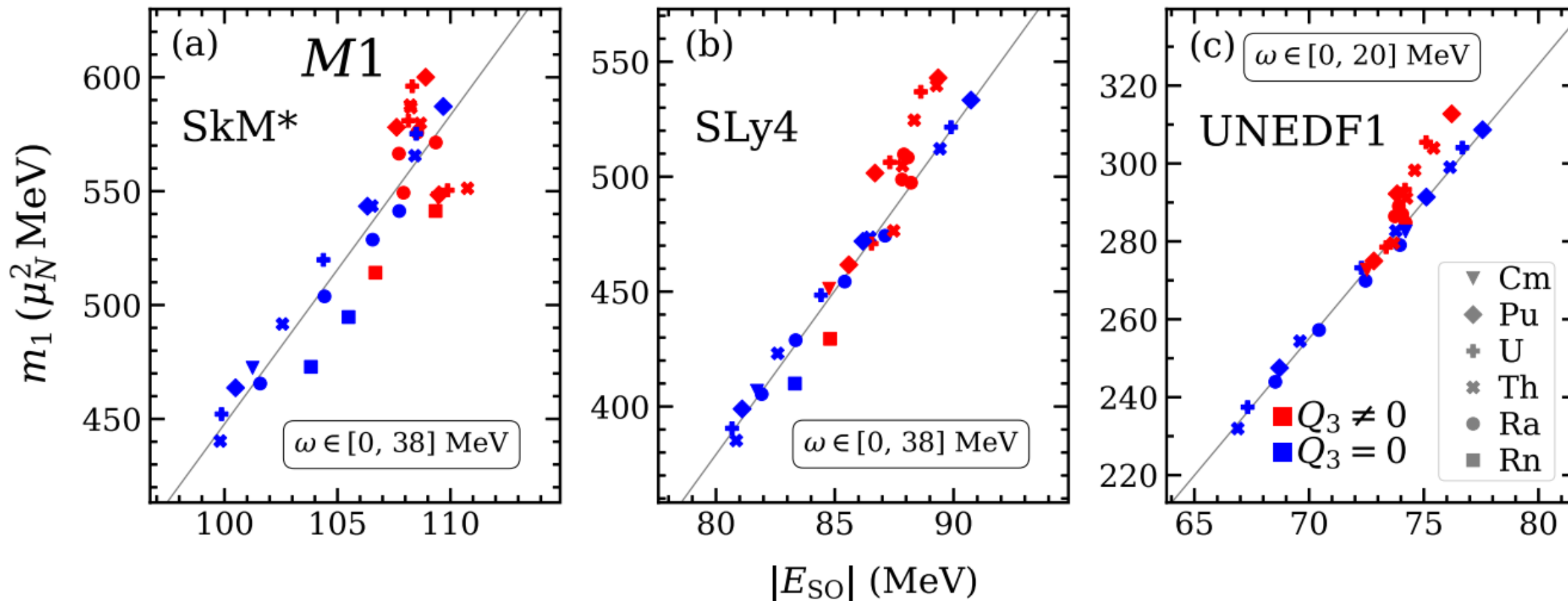
Magnetic dipole response in Th isotopes (SkM*, SLy4, UNEDF1)



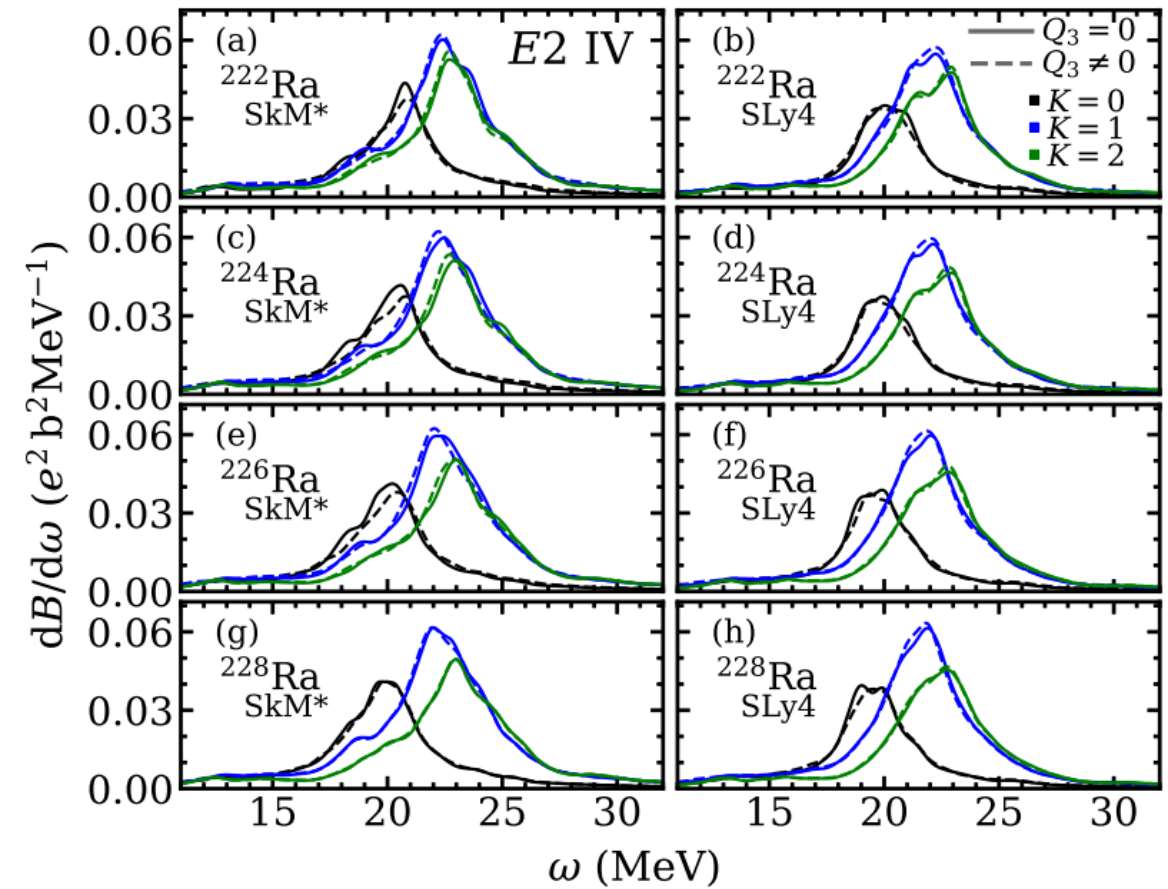
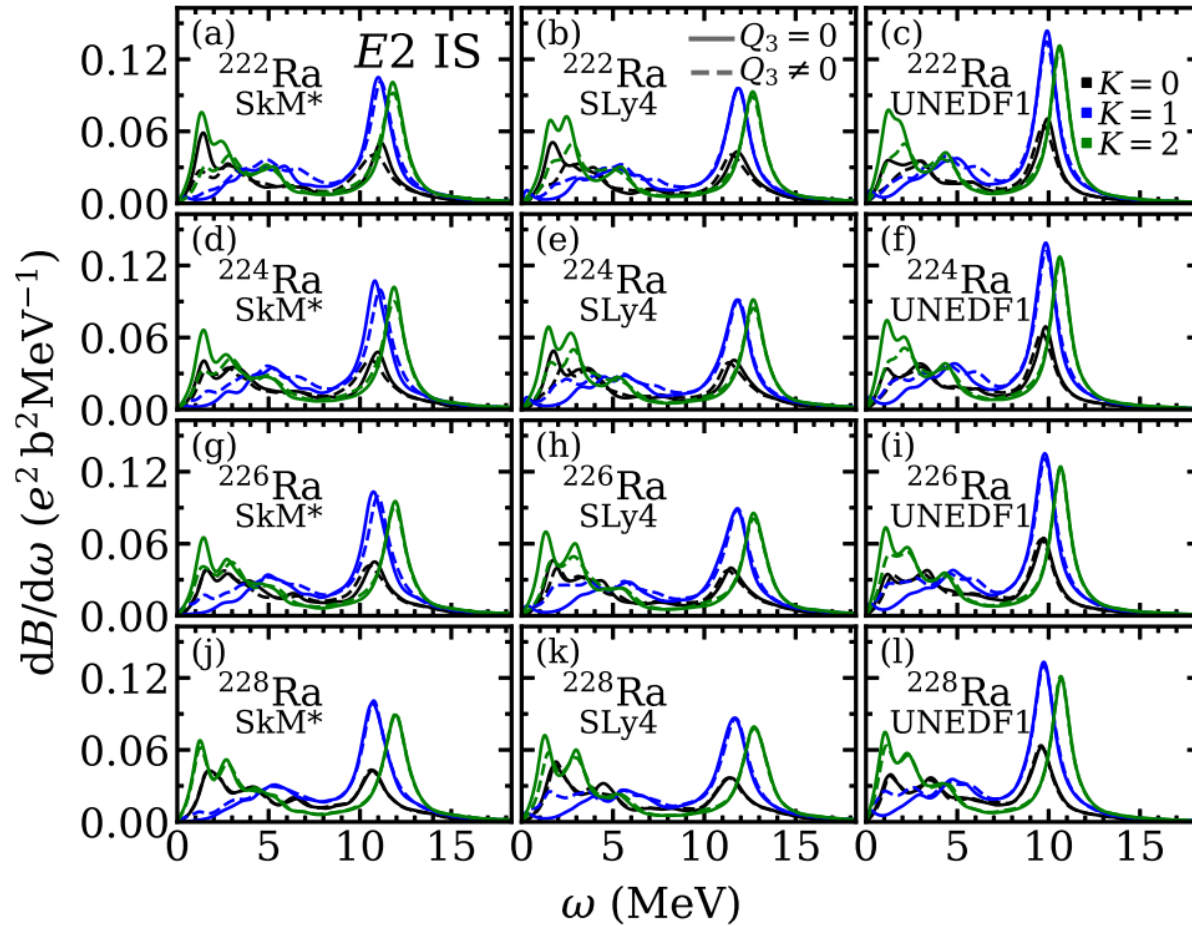
Magnetic dipole response – impact of Landau parameter



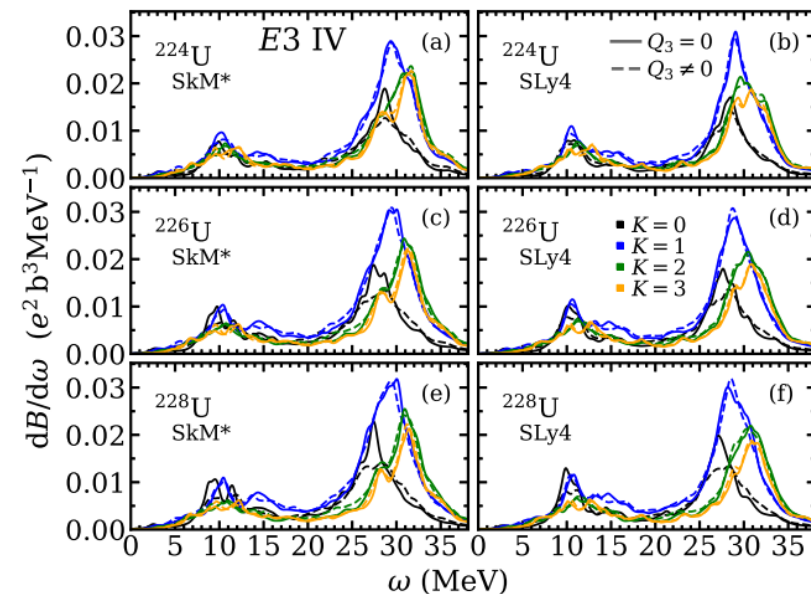
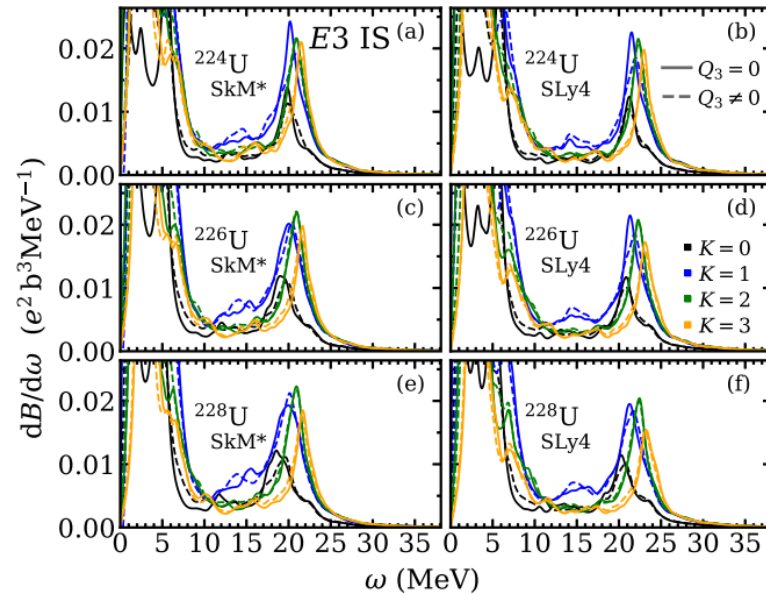
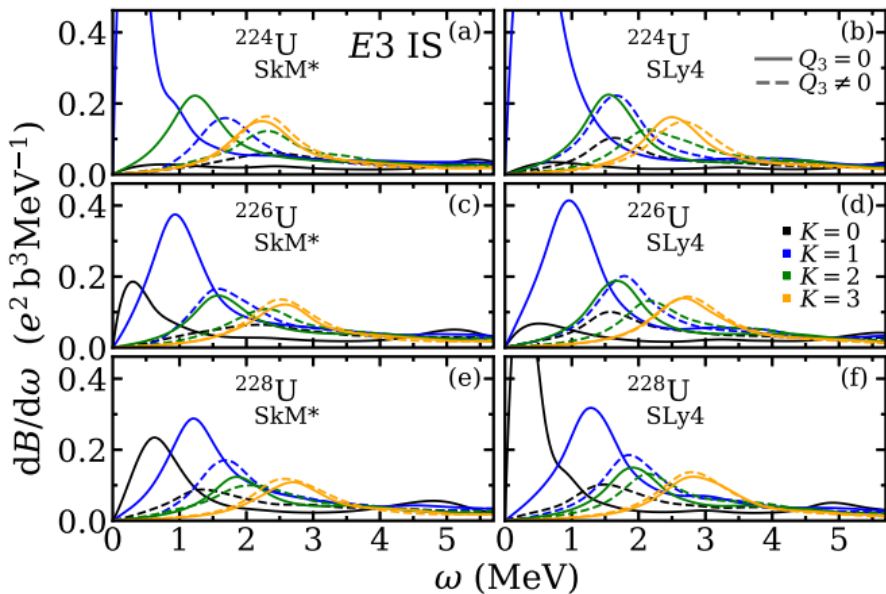
Magnetic dipole response – energy-weighted sum rules



Electric quadrupole response (IS and IV) in Ra isotopes



Electric octupole response (IS and IV) in Ra isotopes



Ground state properties



Nucleus	Force	Q_2' (b)	Q_2 (b)	Q_3 (b ^{3/2})	$E(Q_2', 0)$ (MeV)	$E_{\text{def}}(0, 0)$ (MeV)	$E_{\text{def}}(Q_2, Q_3)$ (MeV)	Nucleus	Force	Q_2' (b)	Q_2 (b)	Q_3 (b ^{3/2})	$E(Q_2', 0)$ (MeV)	$E_{\text{def}}(0, 0)$ (MeV)	$E_{\text{def}}(Q_2, Q_3)$ (MeV)
²²² Rn	SkM*	12.509	12.669	2.830	-1703.679	0.595	-0.438	²²⁴ U	SkM*	12.347	11.741	3.271	-1699.625	0.811	-1.648
	SLy4	12.112	12.427	1.127	-1699.971	1.197	-0.027		SLy4	12.173	13.171	2.596	-1703.603	0.735	-0.773
²²⁴ Rn	SkM*	15.075	15.082	1.792	-1714.463	1.690	-0.053	UNEDF1	12.831	13.504	2.752	-1709.794	0.786	-1.009	
²²² Ra	SkM*	12.615	13.780	3.311	-1701.435	0.807	-1.391	²²⁶ U	SkM*	17.556	16.236	3.692	-1714.514	2.969	-1.256
	SLy4	12.363	13.920	2.484	-1701.029	1.301	-0.700		SLy4	16.423	16.185	3.059	-1717.310	2.687	-0.853
²²⁴ Ra	UNEDF1	12.773	14.133	2.435	-1707.339	1.329	-0.611	UNEDF1	17.570	16.504	3.185	-1723.823	2.427	-0.989	
	SkM*	15.922	15.953	3.395	-1713.749	2.351	-0.919	²²⁸ U	SkM*	21.106	19.360	3.678	-1729.555	5.555	-0.366
	SLy4	14.809	15.788	2.707	-1711.932	2.599	-0.659		SLy4	20.363	18.854	3.345	-1730.966	4.905	-0.502
²²⁶ Ra	UNEDF1	15.245	15.972	2.556	-1718.647	2.417	-0.511	UNEDF1	21.163	19.103	3.415	-1737.786	4.377	-0.528	
	SkM*	18.343	18.121	2.988	-1725.840	3.946	-0.322	²²⁶ Pu	SkM*	13.774	11.869	3.402	-1702.691	1.161	-1.519
	SLy4	17.093	17.546	2.644	-1722.532	3.881	-0.423		SLy4	13.336	13.369	2.692	-1708.162	0.944	-0.659
²²⁸ Ra	UNEDF1	17.476	17.683	2.262	-1729.661	3.504	-0.266	UNEDF1	14.632	13.510	2.950	-1714.458	1.036	-0.960	
	SkM*	20.257	20.282	1.409	-1737.544	5.282	-0.026	²²⁸ Pu	SkM*	19.776	16.988	3.725	-1718.772	3.698	-0.834
	SLy4	19.107	19.162	2.044	-1732.786	5.082	-0.129		SLy4	18.371	17.025	3.076	-1723.076	3.198	-0.539
²²² Th	UNEDF1	19.399	19.378	1.617	-1740.286	4.484	-0.070	UNEDF1	20.685	17.291	3.302	-1729.784	3.118	-0.601	
	SkM*	10.521	11.125	3.106	-1695.131	0.218	-1.619	²³⁰ Pu	SkM*	22.713	21.176	3.371	-1734.851	6.607	-0.043
	SLy4	10.879	12.626	2.398	-1697.522	0.527	-0.717		SLy4	22.001	20.666	3.099	-1737.920	5.705	-0.165
²²⁴ Th	UNEDF1	11.292	12.888	2.453	-1703.599	0.563	-0.821	UNEDF1	23.101	21.265	3.228	-1744.877	5.359	-0.134	
	SkM*	15.239	15.242	3.559	-1708.708	1.944	-1.522	²⁸⁸ Pu	SkM*	20.266	15.931	5.402	-1992.829	0.190	-1.438
	SLy4	14.393	15.174	2.839	-1709.960	2.138	-0.943		SLy4	17.218	18.117	4.589	-1952.142	0.005	-0.849
²²⁶ Th	UNEDF1	14.934	15.387	2.872	-1716.354	1.907	-0.979	²⁹⁰ Pu	SkM*	28.325	22.099	5.798	-1997.272	1.956	-0.509
	SkM*	18.802	17.734	3.620	-1722.446	4.117	-0.783		SLy4	25.124	22.511	5.050	-1954.144	1.488	-0.569
	SLy4	17.635	17.376	3.115	-1722.243	3.690	-0.757	UNEDF1	23.298	20.739	4.605	-1980.658	0.249	-0.796	
²²⁸ Th	UNEDF1	18.340	17.566	3.084	-1728.973	3.428	-0.714	²²⁸ Cm	SkM*	14.887	11.699	3.513	-1704.331	1.456	-1.326
	SkM*	21.004	20.406	3.090	-1735.774	6.036	-0.163		SLy4	14.345	13.373	2.708	-1711.228	1.267	-0.457
	SLy4	20.245	19.578	3.113	-1734.235	5.437	-0.348	UNEDF1	20.438	13.273	3.138	-1717.665	1.424	-0.736	
²²⁸ Th	UNEDF1	20.783	19.797	2.962	-1741.264	4.739	-0.266	²⁹⁰ Cm	SkM*	32.332	13.996	5.551	-2015.804	0.860	-1.316
	SkM*	20.245	19.578	3.113	-1734.235	5.437	-0.348		SLy4	19.380	17.595	4.897	-1978.145	0.450	-0.910
	UNEDF1	20.783	19.797	2.962	-1741.264	4.739	-0.266	UNEDF1	16.680	15.428	4.605	-2002.682	0.224	-1.240	