

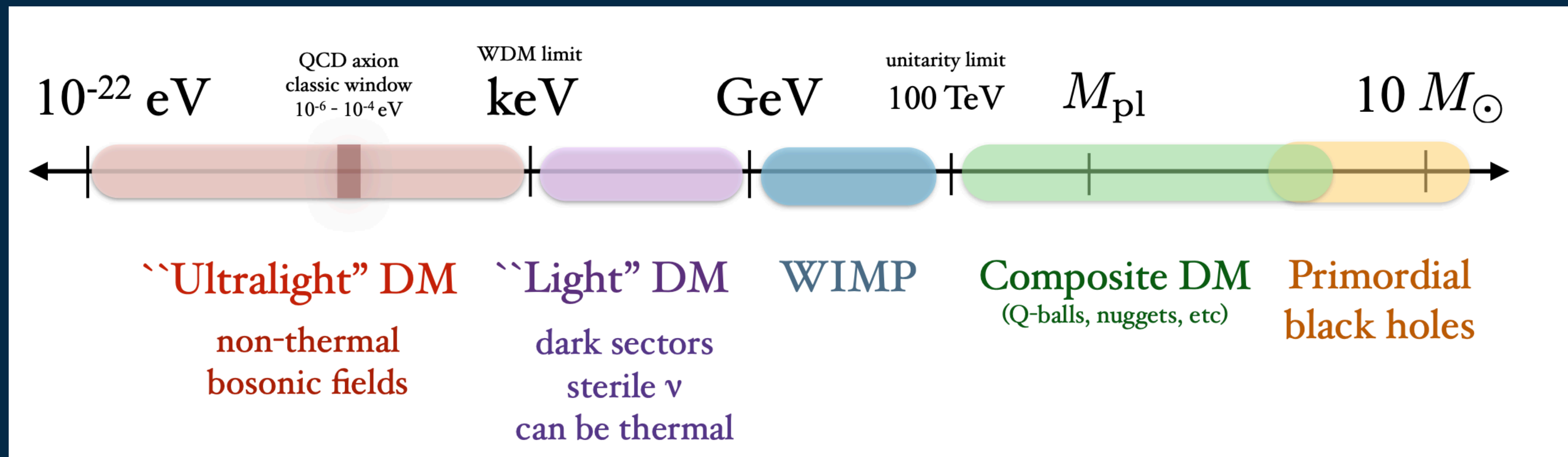
Hadronic Dark Matter Annihilation During Big Bang Nucleosynthesis

Based on work with Adam Ritz, [PRD 113, 035004/ arXiv 2510.11791]

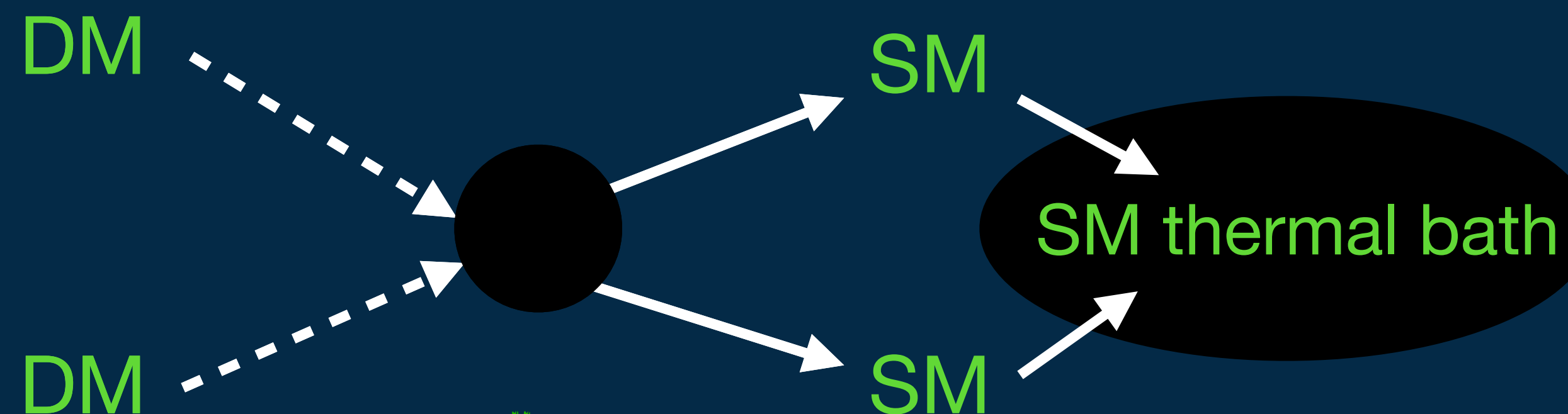
Afif Omar
July, 2026



Light Dark World 2026, Carleton University



T. Lin 1904.07915

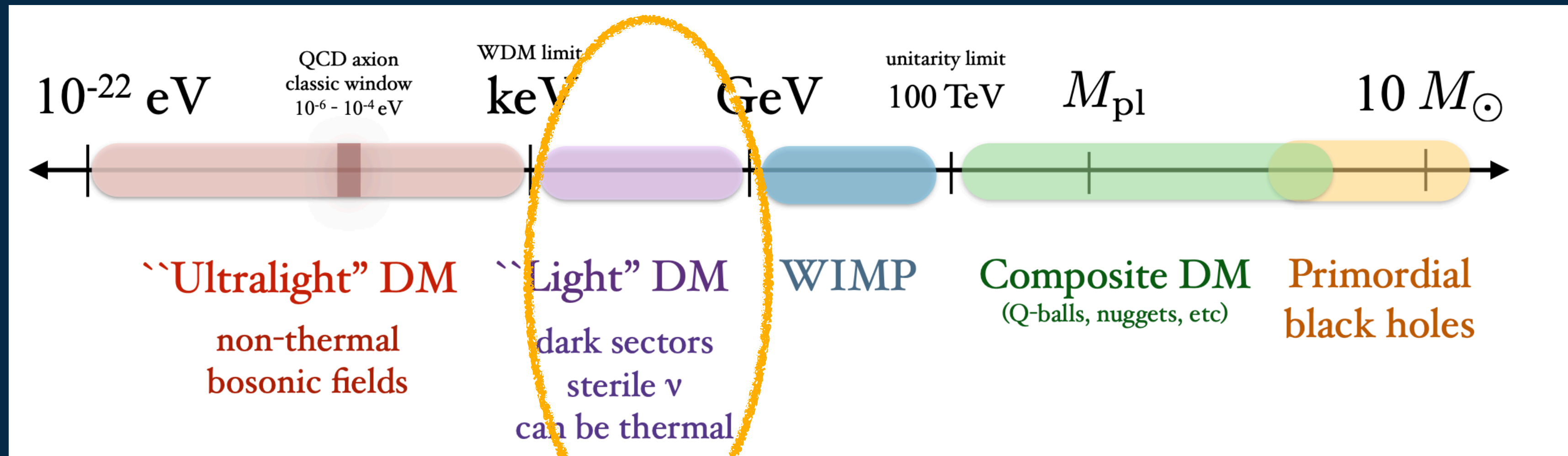


Velocity independent (*s*-wave)

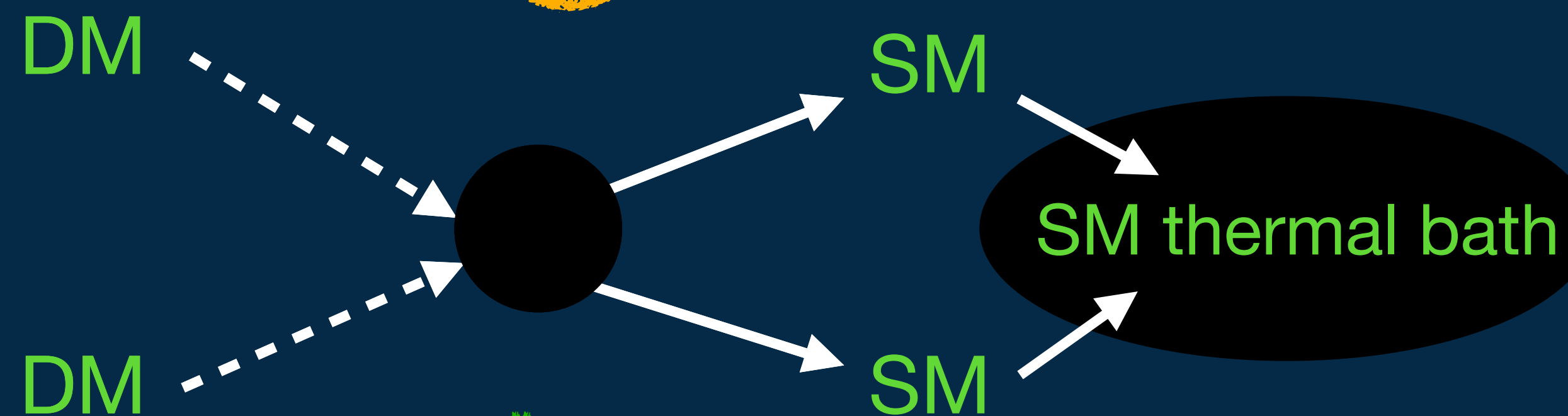
Ruled out by CMB

Velocity dependent (*p*-wave)

Evades CMB



T. Lin 1904.07915

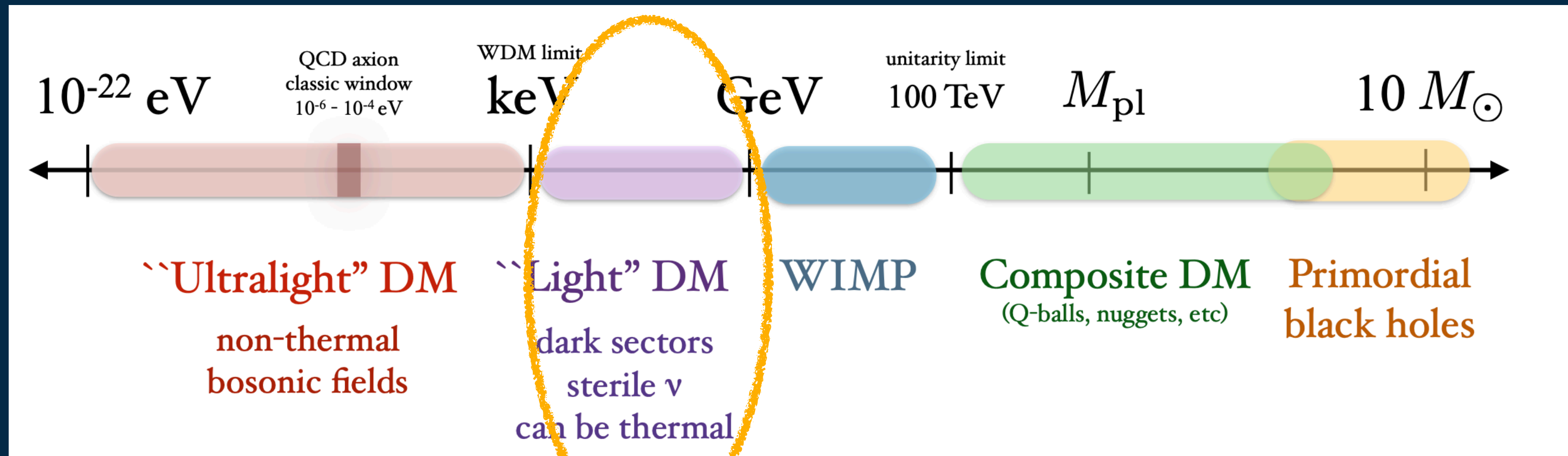


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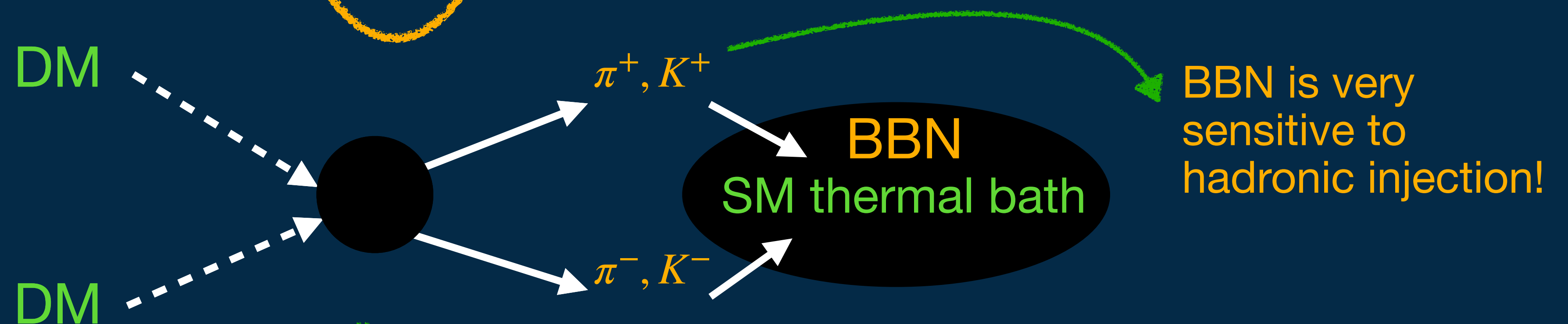
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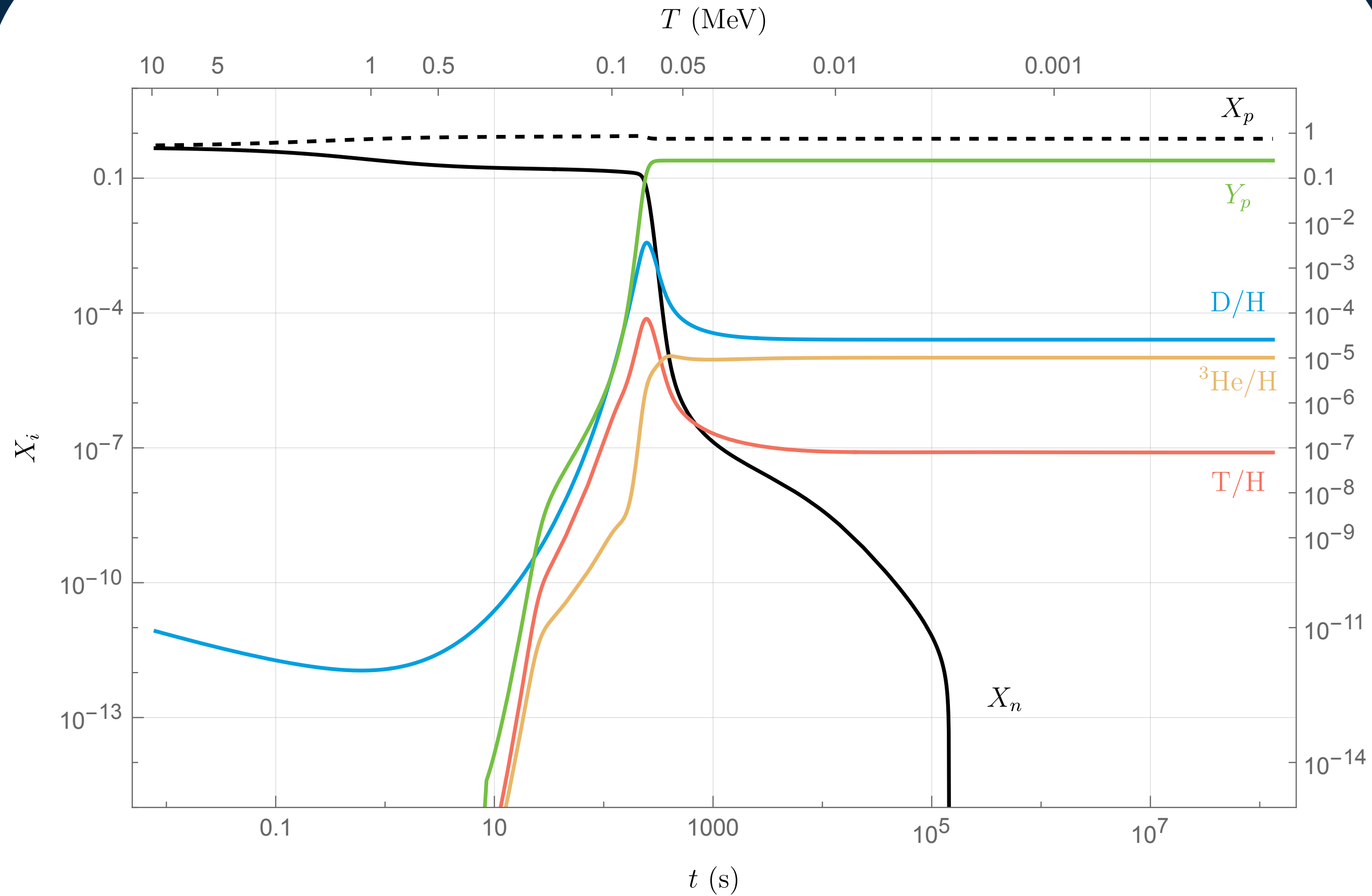


Velocity independent (*s*-wave)

Ruled out by CMB

Velocity dependent (*p*-wave)

Evades CMB



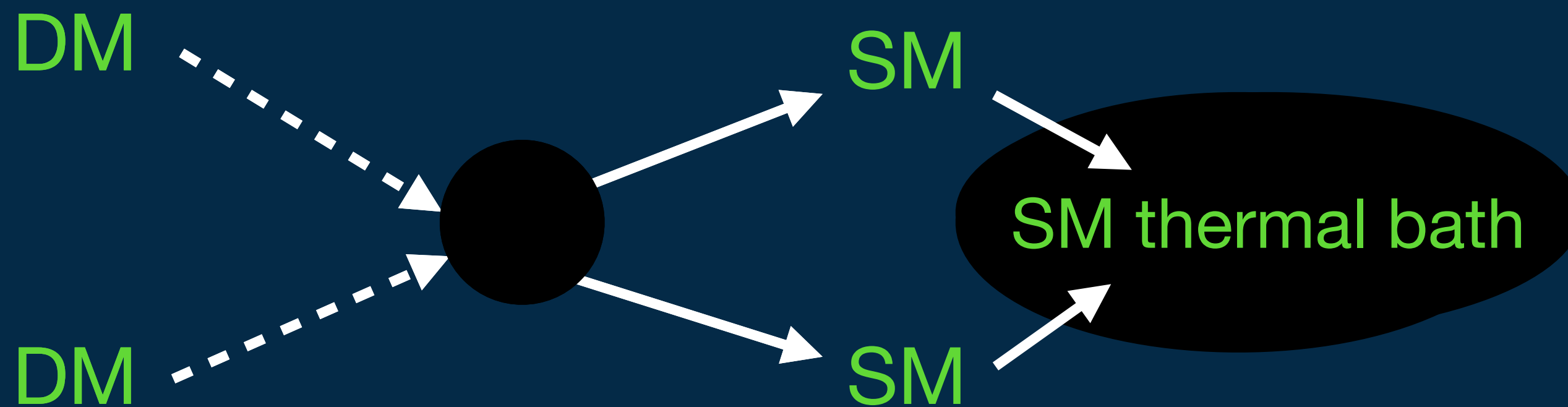
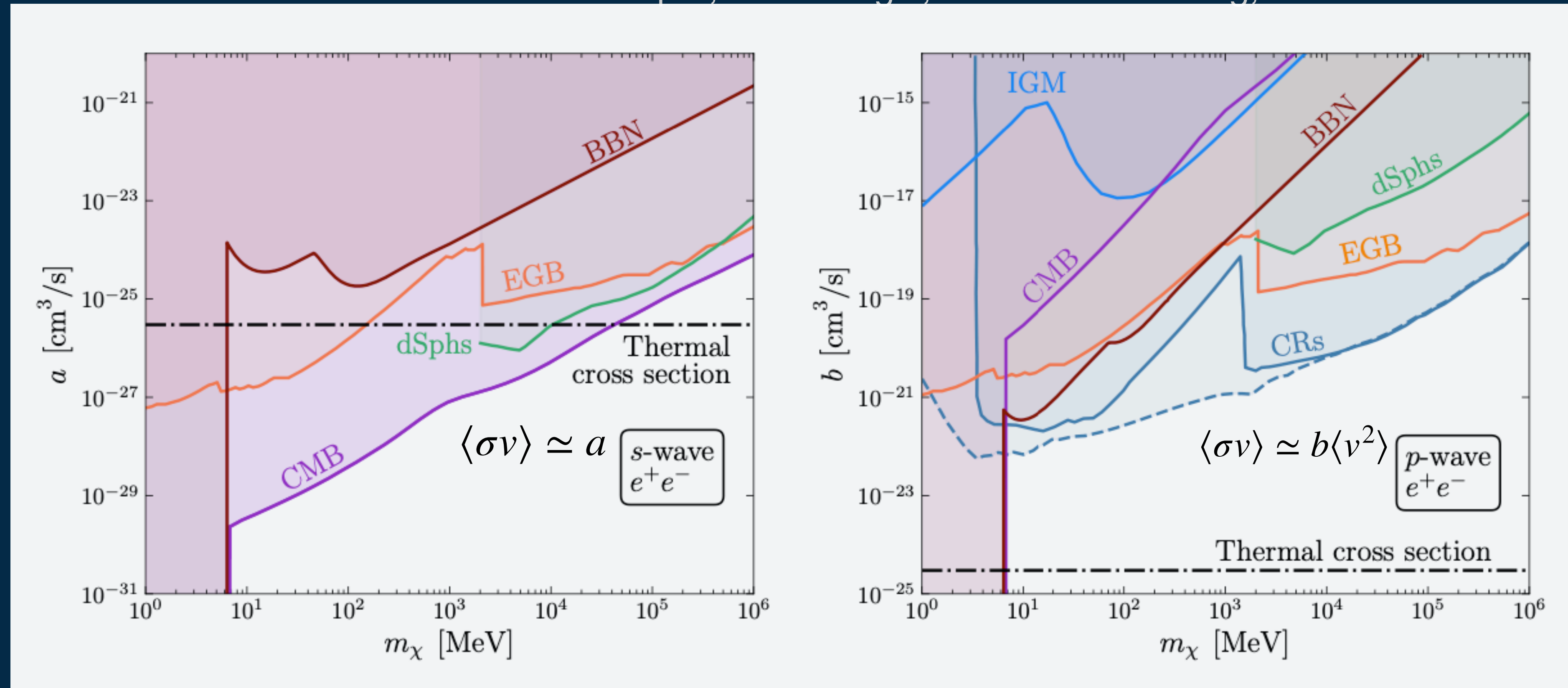
Inputs:

- Nuclear interaction rates
- Neutron lifetime and weak rates
- Hubble rate
- Ratio of baryons to photons
- ...

Output:

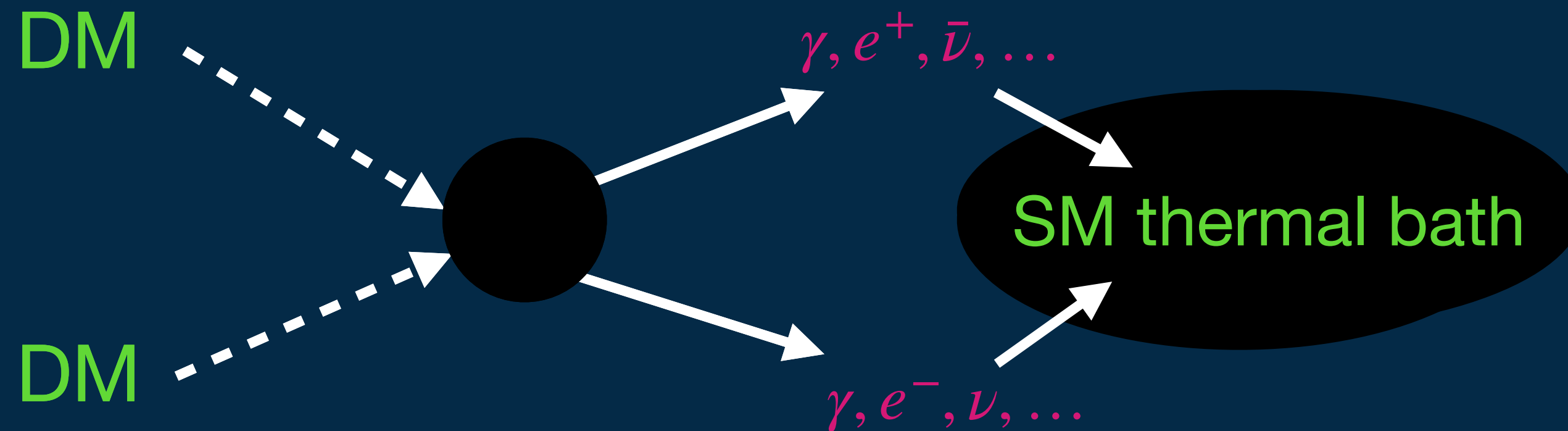
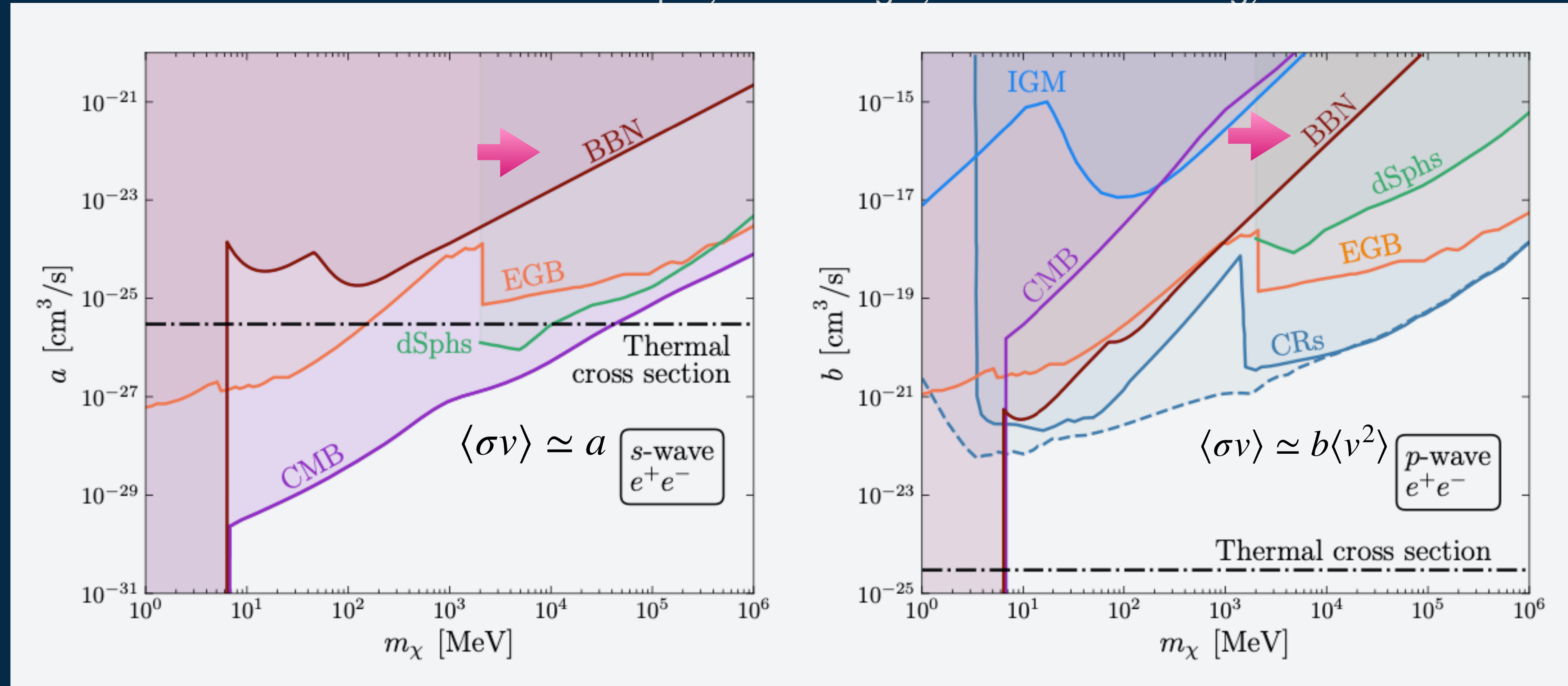
- Primordial nuclear abundances

Entropy injection
Photodisintegration
Hadronic injection?



Nuclear interaction rates
Neutron lifetime and weak rates
Hubble rate
Ratio of baryons to photons

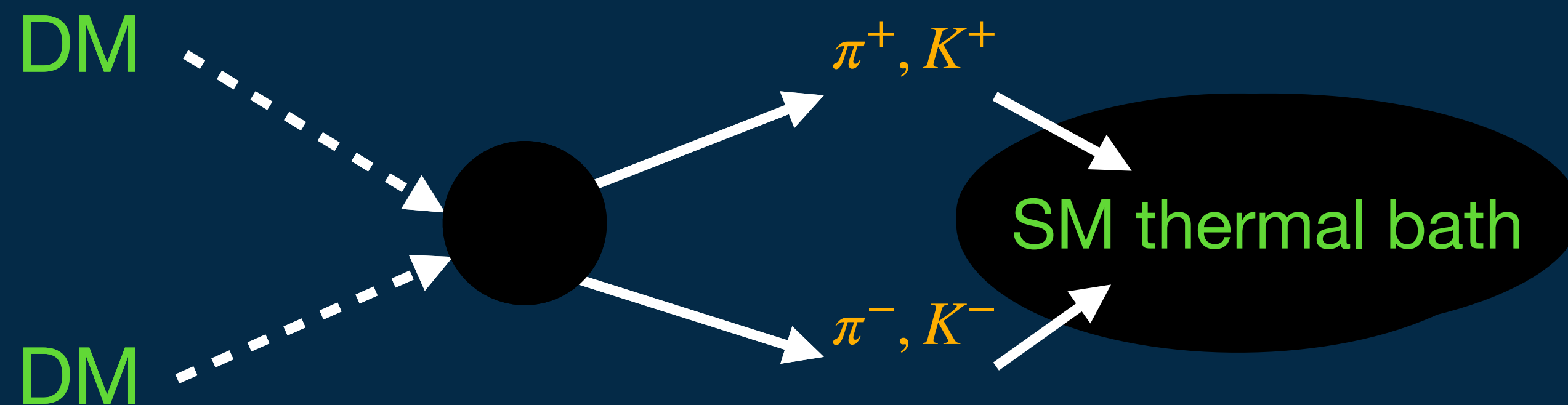
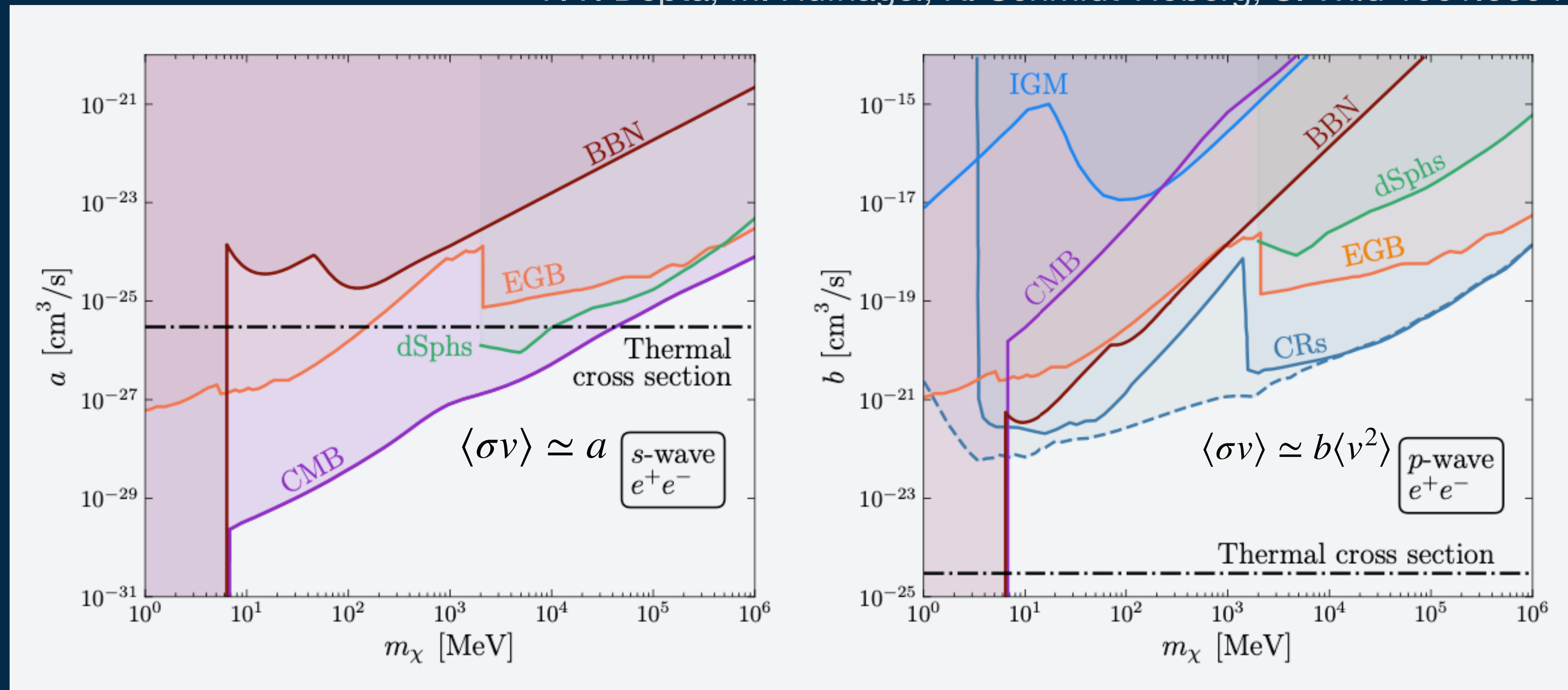
Entropy injection
 Photodisintegration
 Hadronic injection?



Nuclear interaction rates
 Neutron lifetime and weak rates
 Hubble rate
 Ratio of baryons to photons

Entropy injection
Photodisintegration
Hadronic injection?

With focus on p -wave to avoid CMB constraints



Nuclear interaction rates
Neutron lifetime and weak rates
Hubble rate
Ratio of baryons to photons

SBBN

$$\frac{dn_{n,p}}{dt} + 3Hn_{n,p} \propto \text{weak } n \leftrightarrow p \text{ rates}$$

$$\frac{dn_{\text{nucI}}}{dt} + 3Hn_{\text{nucI}} \propto \text{Standard nuclear rates}$$

*

$$\begin{aligned} \pi^- + p^+ &\rightarrow n + \gamma, & (\sigma v)_{pn\gamma}^{\pi^-} &\simeq 0.57 \text{ mb}, \\ \pi^- + p^+ &\rightarrow n + \pi^0, & (\sigma v)_{pn\pi^0}^{\pi^-} &\simeq 0.88 \text{ mb}, \\ \pi^+ + n &\rightarrow p + \pi^0, & (\sigma v)_{np\pi^0}^{\pi^+} &\simeq 1.7 \text{ mb}. \end{aligned}$$

M. Pospelov, J. Pradler 1006.4172

BBN + Pion injection

$$\frac{dn_{n,p}}{dt} + 3Hn_{n,p} \propto \text{weak } n \leftrightarrow p \text{ rates} + \pi^\pm \text{ interactions}^*$$

$$\frac{dn_{\text{nucI}}}{dt} + 3Hn_{\text{nucI}} \propto \text{Standard nuclear rates}^\dagger$$

$$\begin{aligned} \frac{dn_{\pi^\pm}}{dt} + 3Hn_{\pi^\pm} &\propto \text{DM injection} - \text{pion decay} \\ &\quad - \pi^\pm \text{ interactions}^* \end{aligned}$$

SBBN

$$\frac{dn_{n,p}}{dt} + 3Hn_{n,p} \propto \text{weak } n \leftrightarrow p \text{ rates}$$

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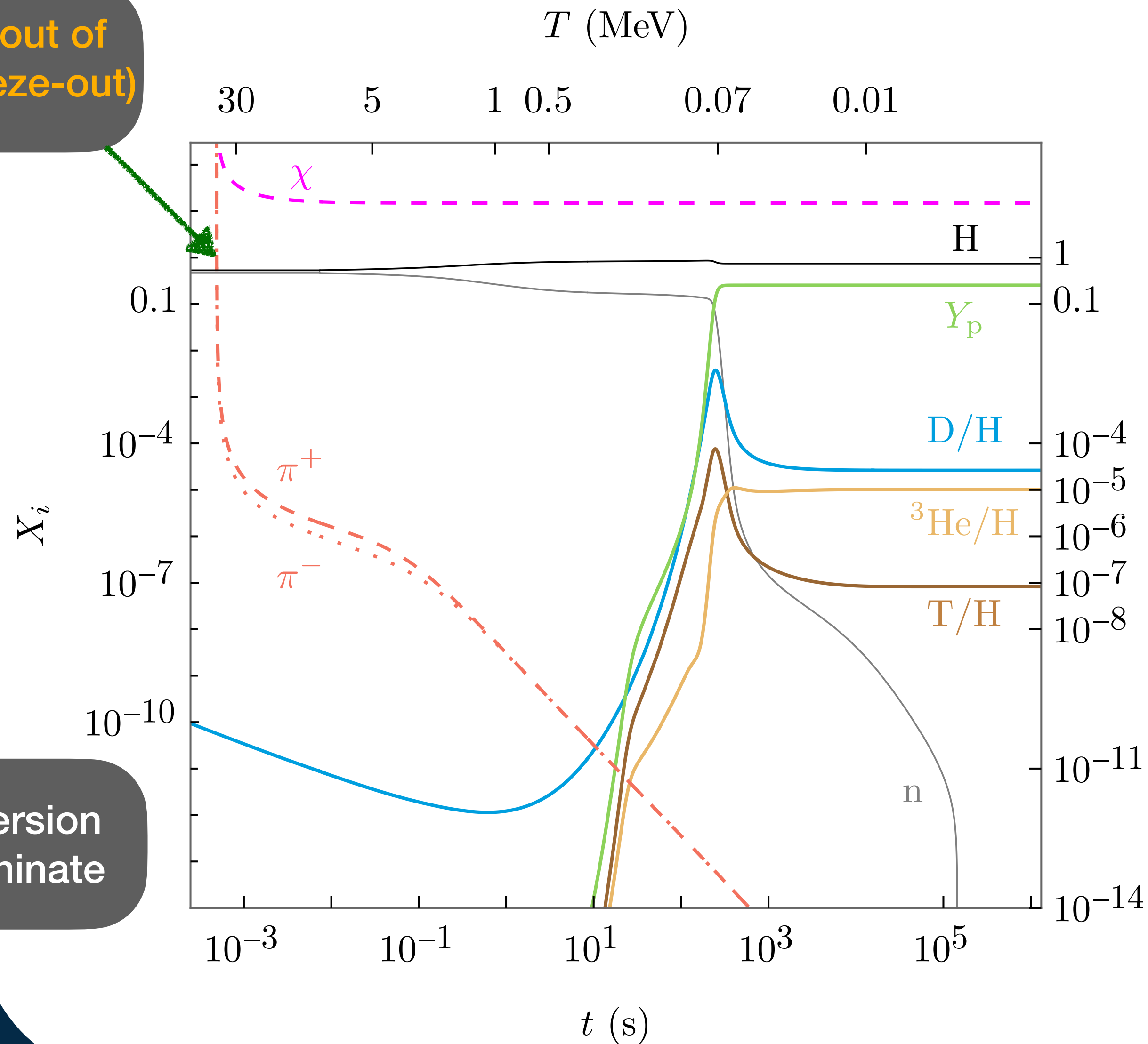
$$\sim \frac{1}{2}n_\chi^2 \langle \sigma v \rangle \quad \sim \frac{n_{\pi^\pm}}{\tau_\pi}$$

$$\frac{dn_{\pi^\pm}}{dt} + 3Hn_{\pi^\pm} \propto \text{DM injection} - \text{pion decay} - \pi^\pm \text{ interactions}^*$$

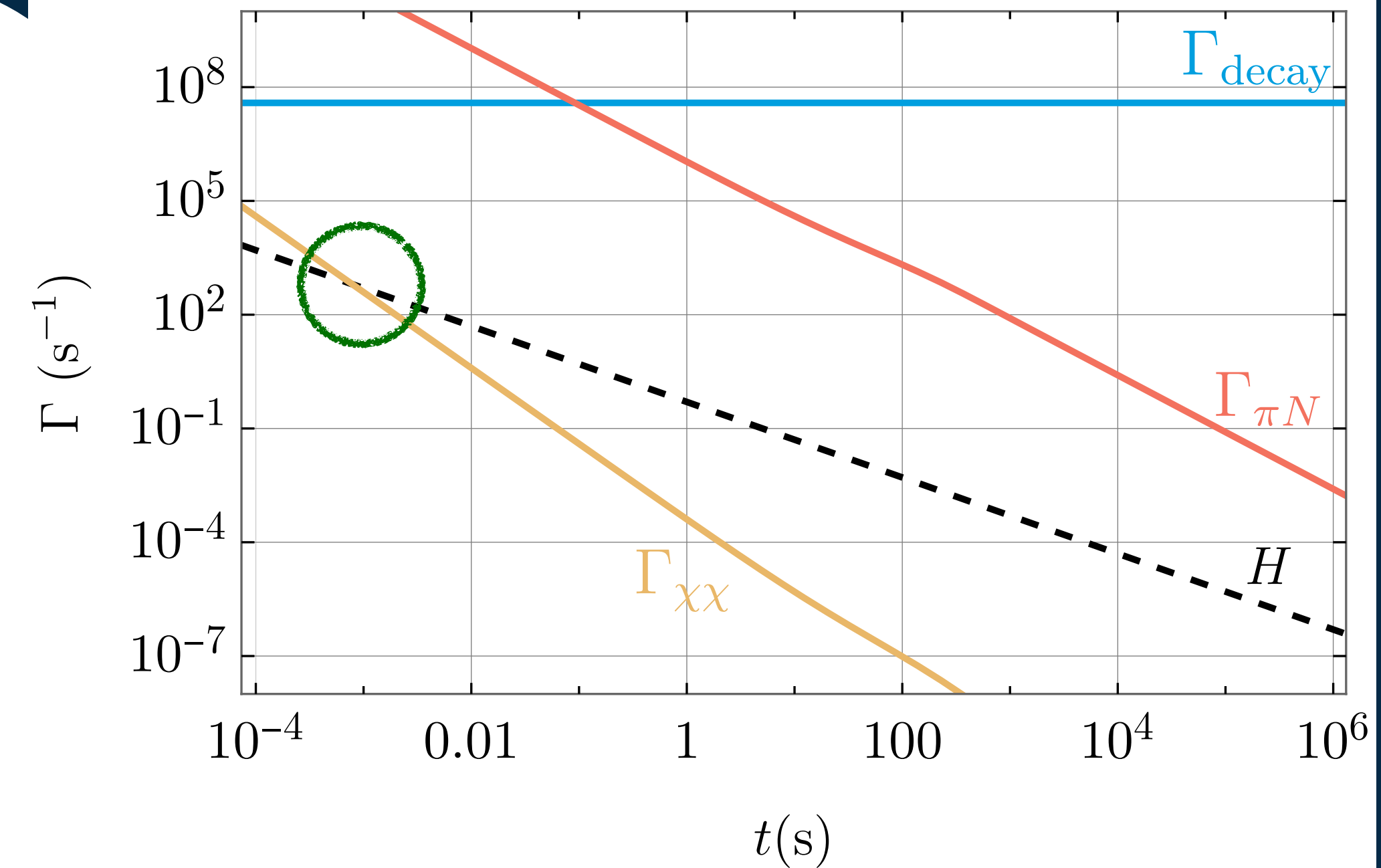
$$m_\chi = 400 \text{ MeV}$$

$$b = 10^{-23} \text{ cm}^3/\text{s}$$

Pions/DM fall out of equilibrium (freeze-out)



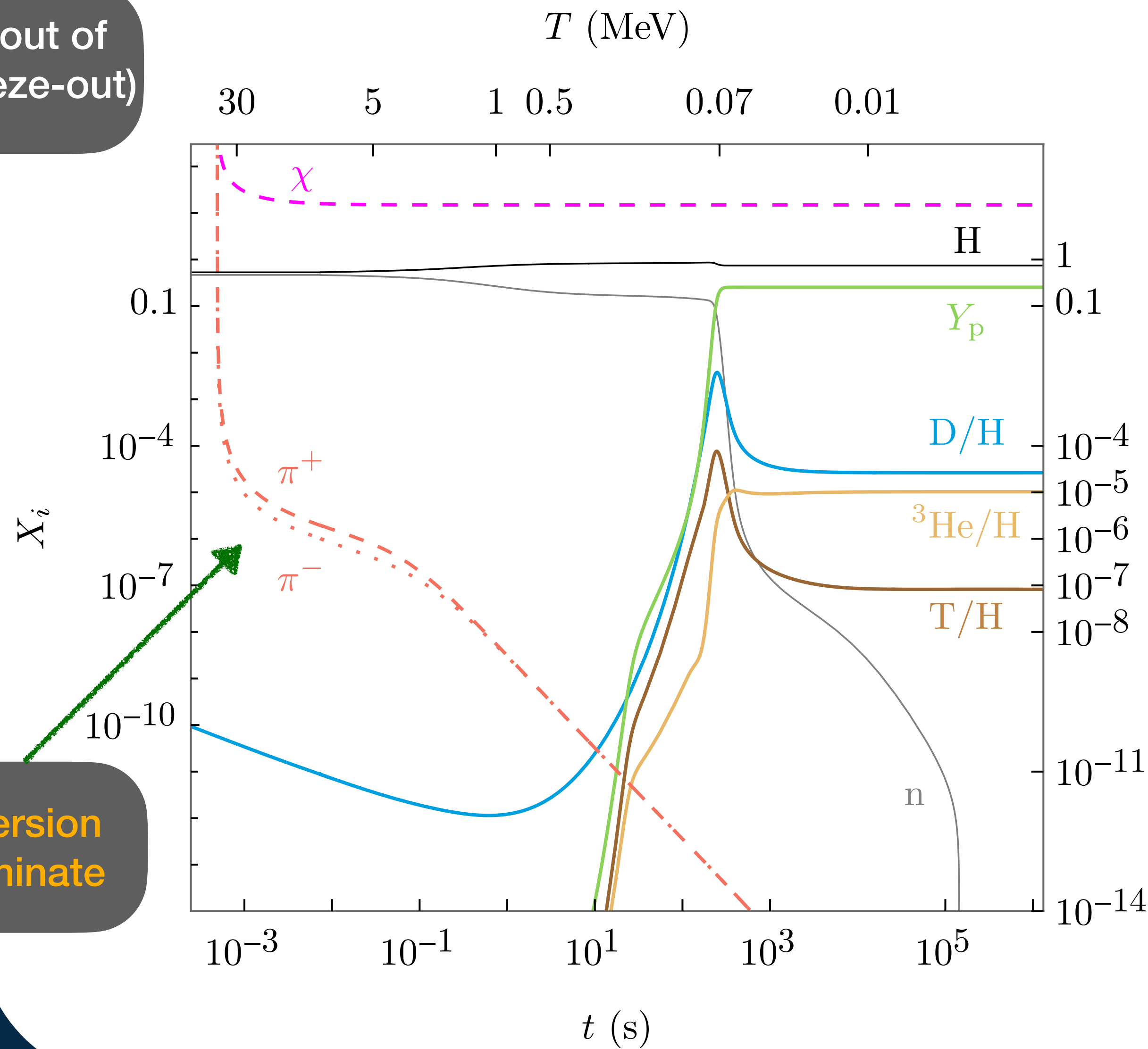
Charge conversion reactions dominate



AO, A. Ritz 2510.11791

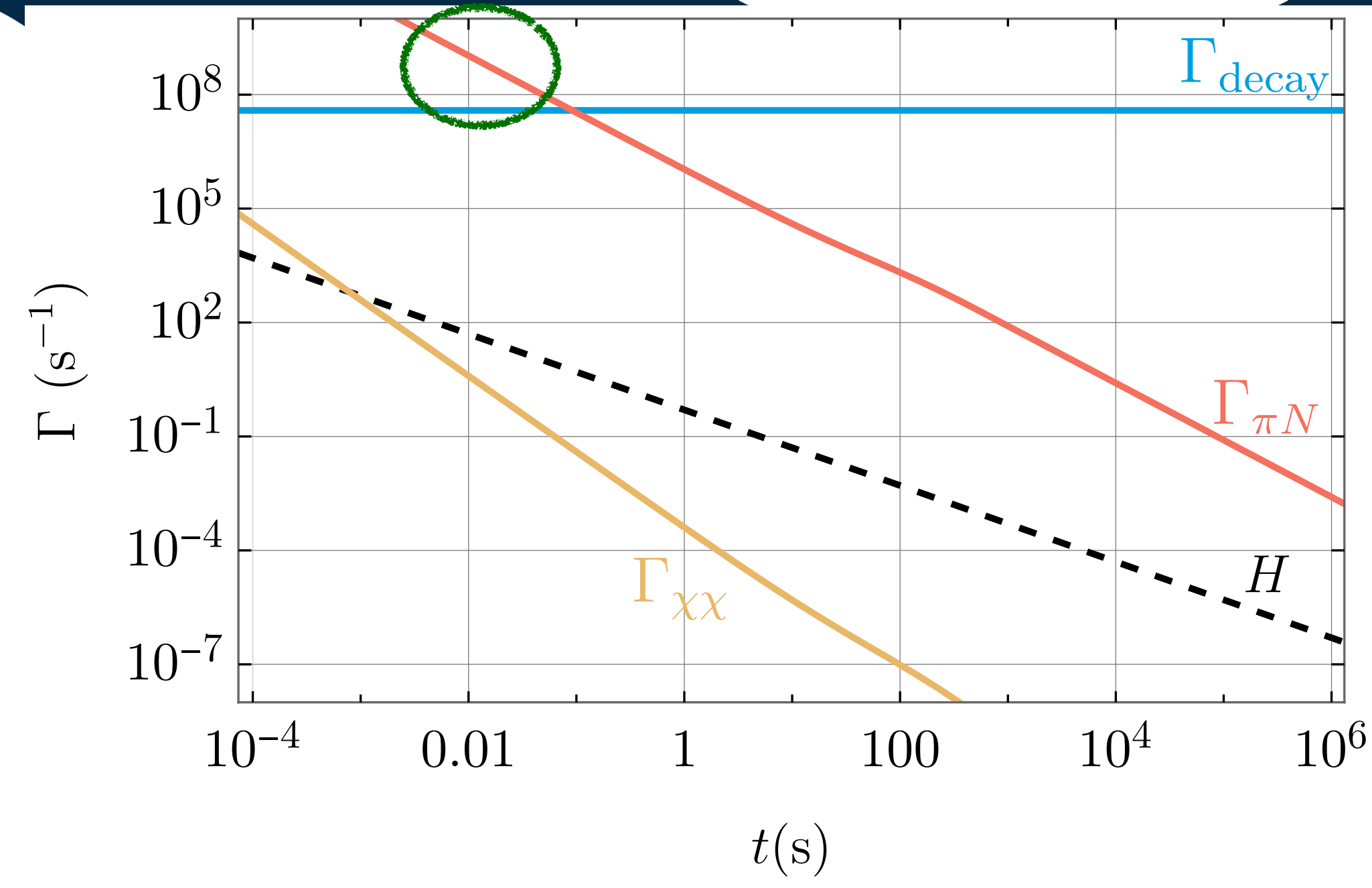
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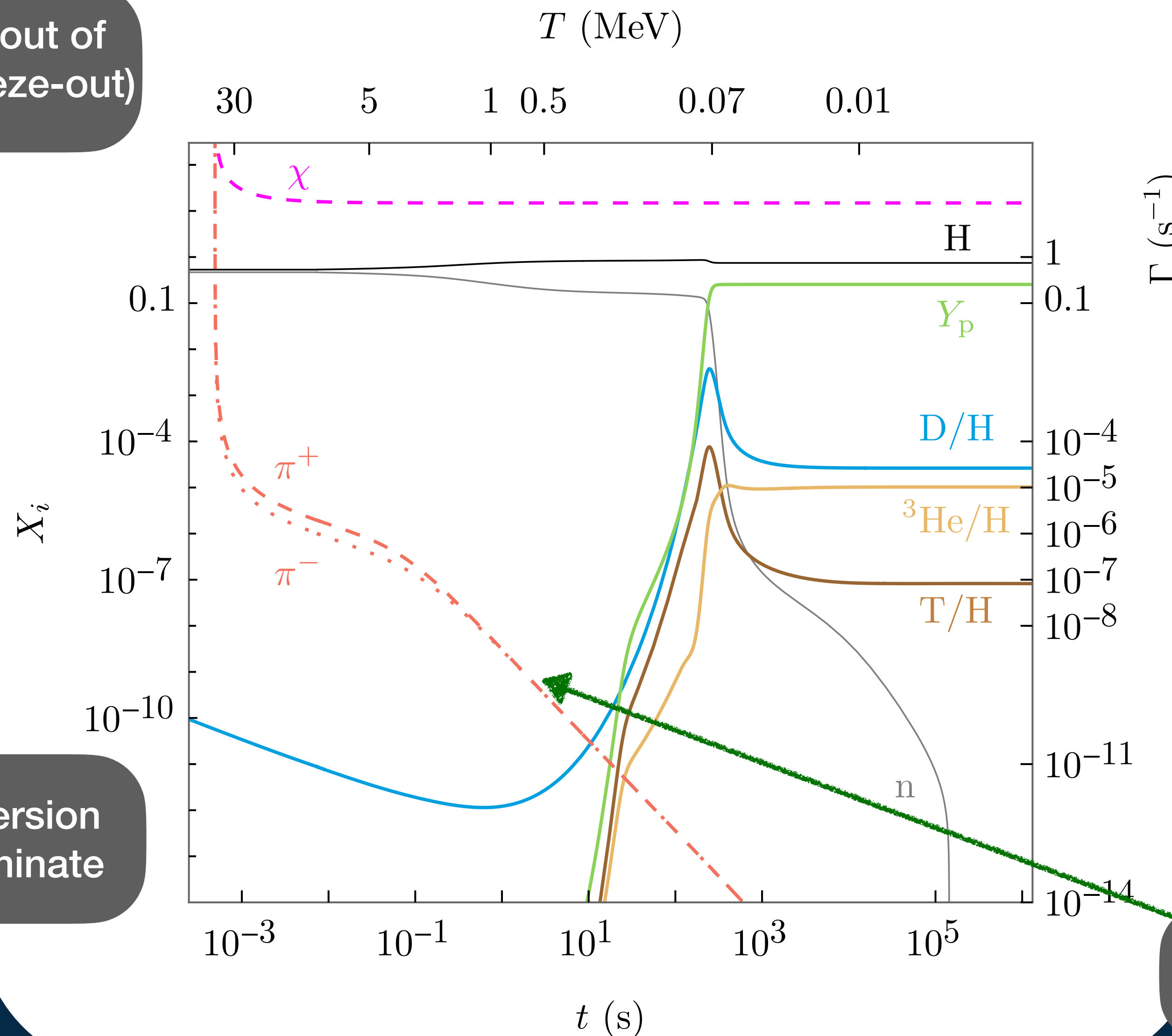


AO, A. Ritz 2510.11791

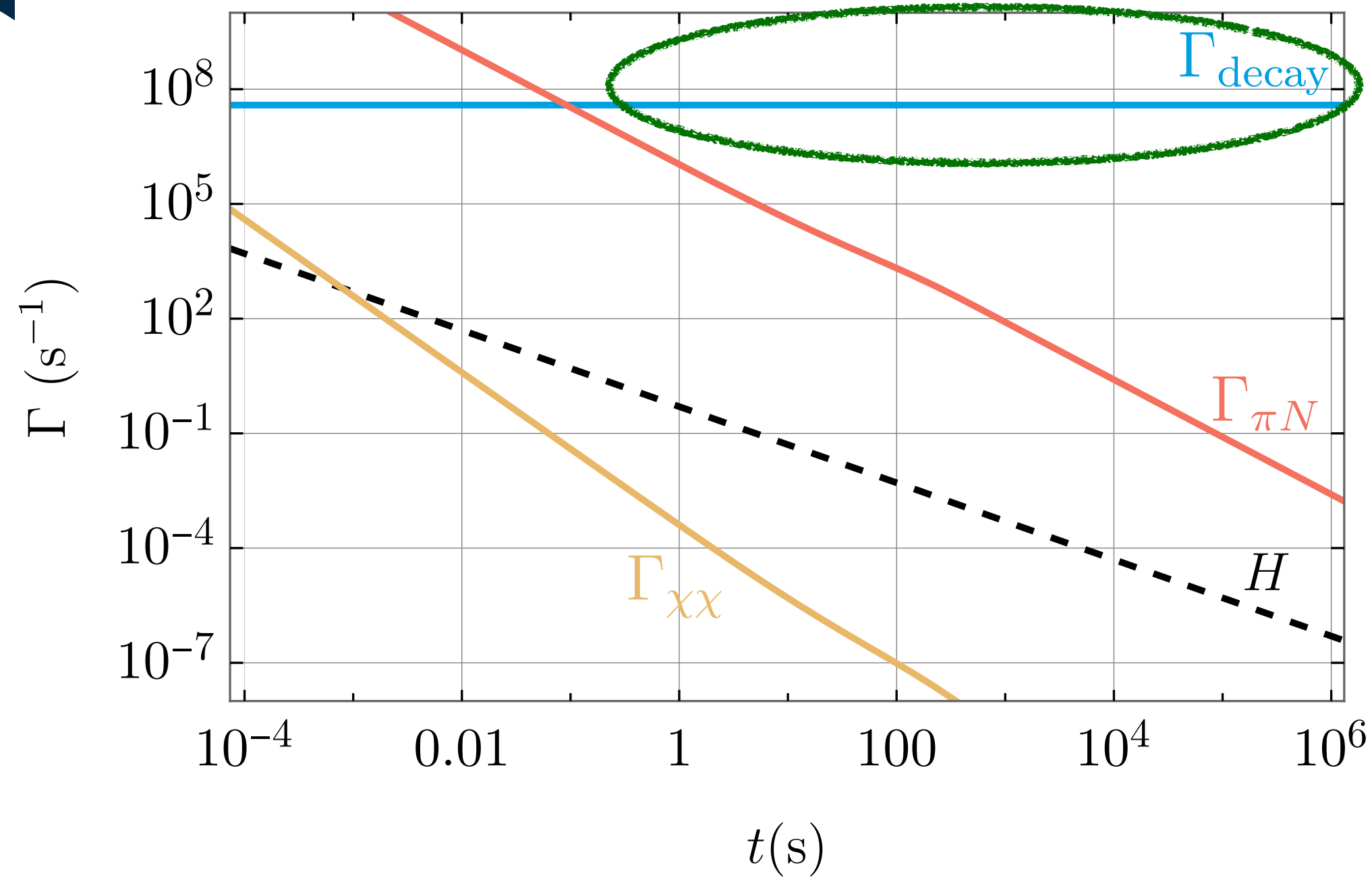
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AO, A. Ritz 2510.11791

Pion decay dominates

Compared with $\mathcal{O}(1)$ mb for $\pi^\pm$

Similar analysis as $\pi^\pm$

Higher mass threshold ($m_K \simeq 494$ MeV), but reaction rates are higher.

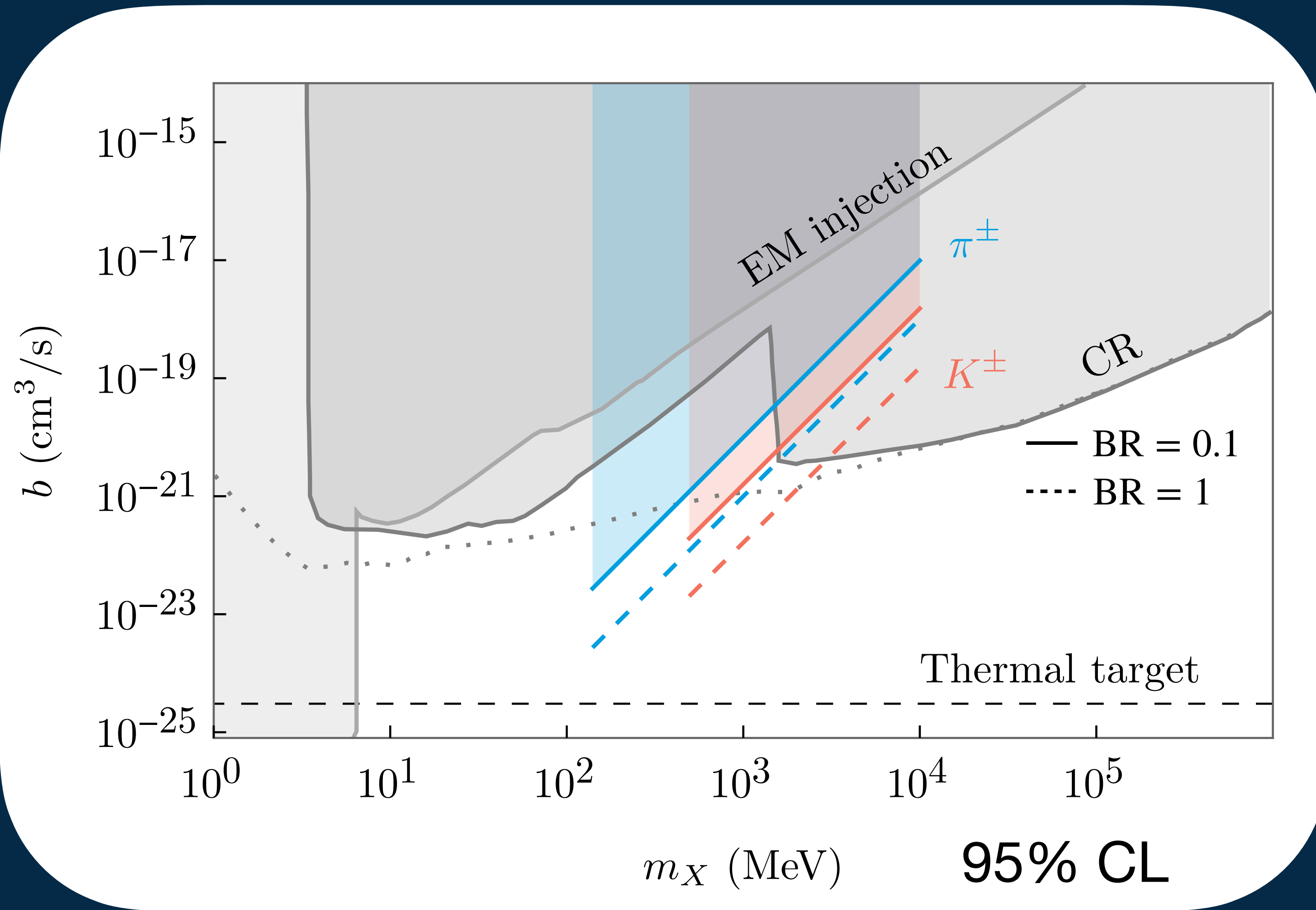
Must include $K^\pm$ decays to $\pi^\pm$.


$$\begin{aligned} K^- + p &\rightarrow n + X : & (\sigma v)_{pn}^{K^-} &\simeq 32 \text{ mb}, \\ K^- + n &\rightarrow p + X : & (\sigma v)_{np}^{K^-} &\simeq 13 \text{ mb}. \end{aligned}$$

M. Pospelov, J. Pradler 1006.4172

$$\begin{aligned} K^\pm &\rightarrow \pi^\pm \pi^0 (\pi^0) & \text{BR} &= 22.4\%, \\ K^\pm &\rightarrow \pi^\pm \pi^\pm \pi^\mp & \text{BR} &= 5.6\%. \end{aligned}$$

PDG (2024)

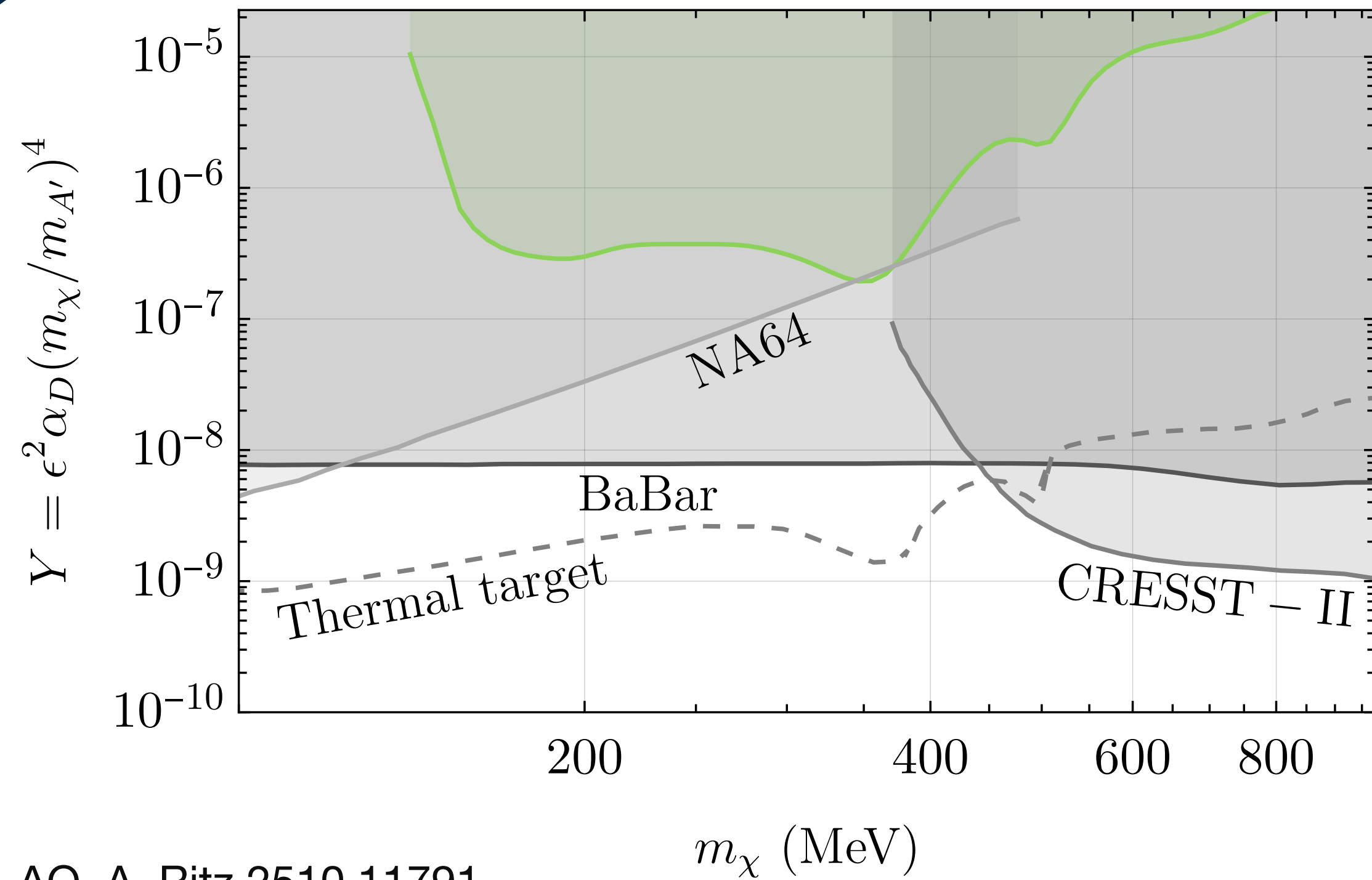
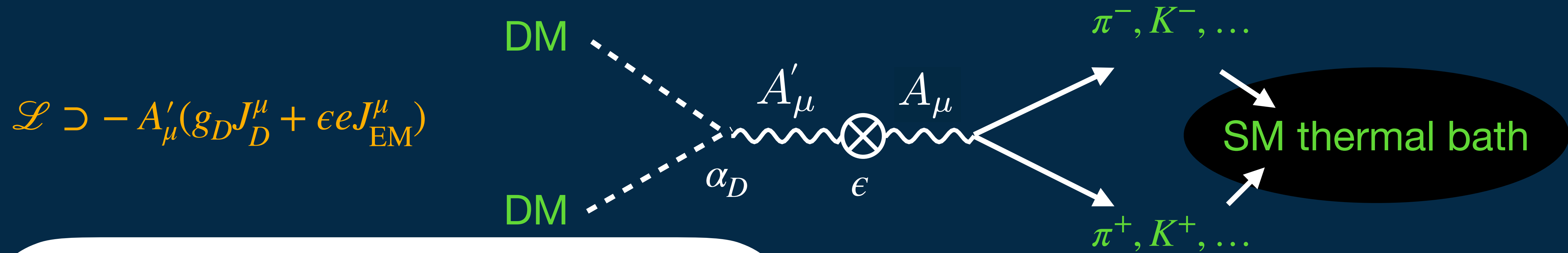


D and ^4He observations:

$$D/H \times 10^{-5} = 2.547 \pm 0.025$$

$$Y_p = 0.245 \pm 0.003$$

PDG (2024) and refs within



this model imposes a
BR less than 10%

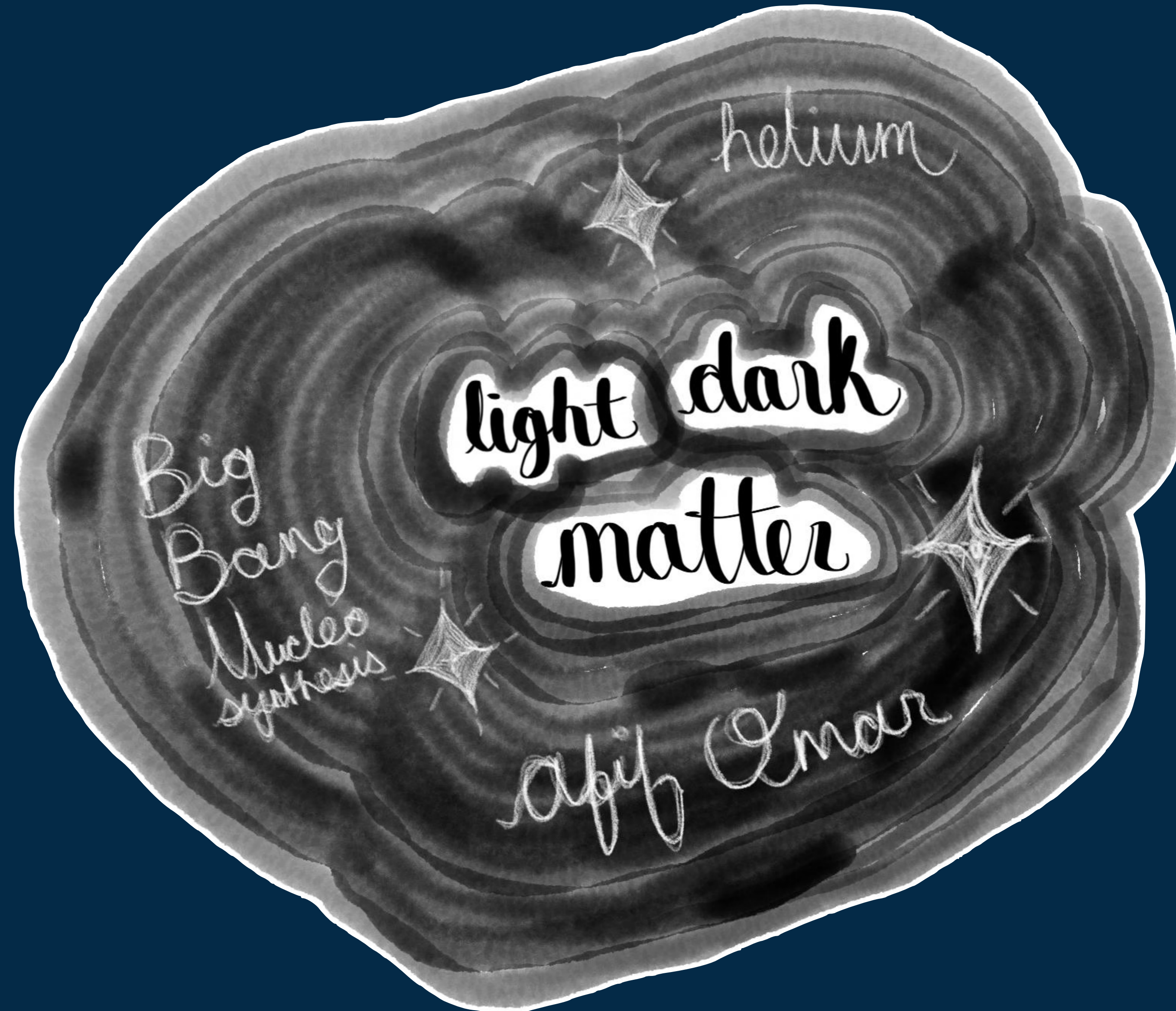
$$\alpha_D = 0.5$$

$$m_{A'}/m_\chi = 3$$

AO, A. Ritz 2510.11791

Takeaways

1. BBN is a viable probe of annihilating DM, especially with hadronic final states.
2. BBN sensitivity does not yet reach the thermal target. Improving D/H and Y_p precision could close that gap.
3. Pseudo-Dirac DM with a small mass splitting, Δm , evades CMB constraints. BBN could potentially impose important constraints on Δm .



Back up slides

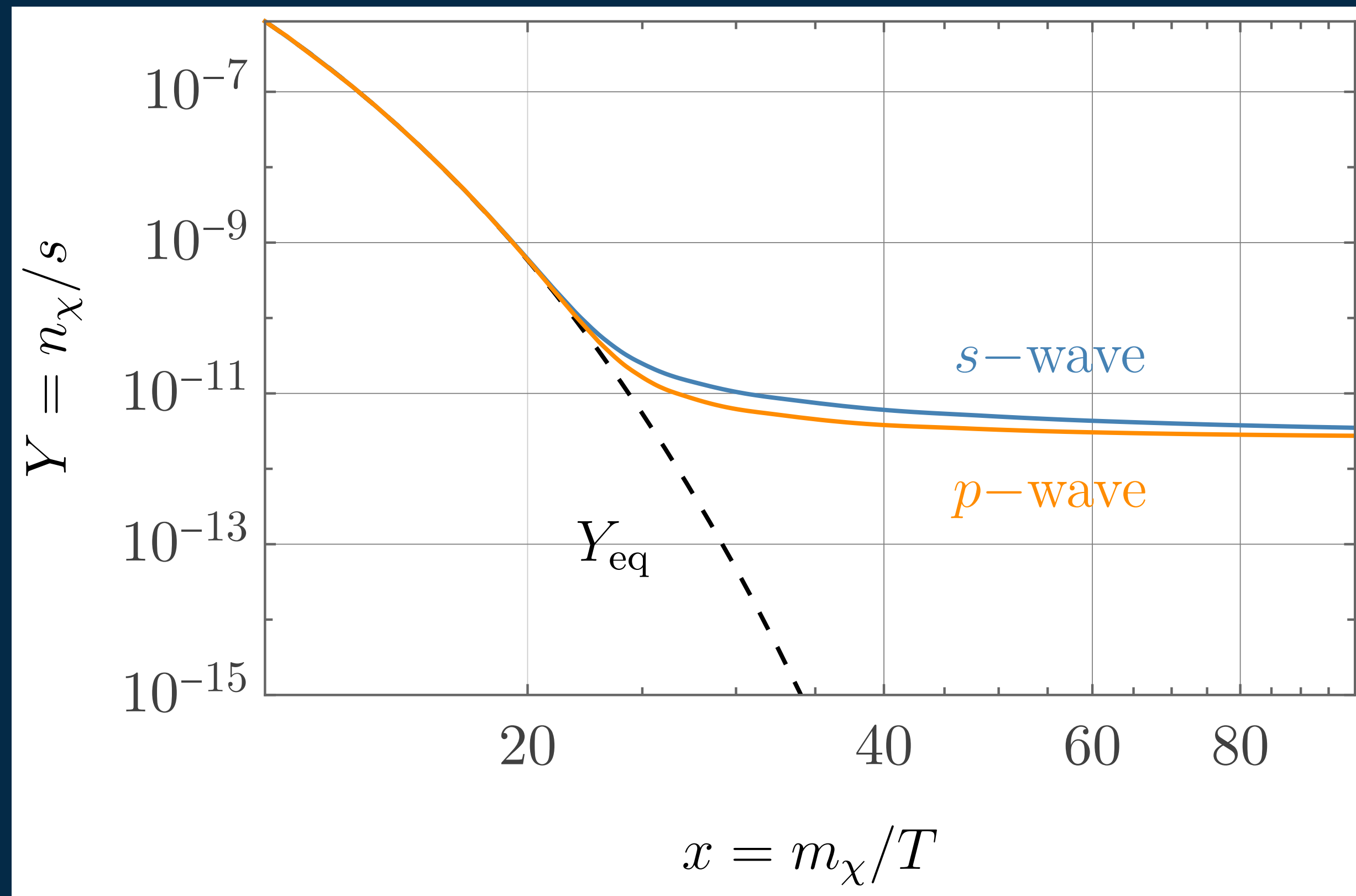
DM freeze-out

DM $\longleftrightarrow$ SM

Thermal equilibrium

Freeze-out DM $\longrightarrow$ SM

$$T \lesssim m_\chi$$



Velocity-independent

$$\langle \sigma v \rangle \simeq a$$

$$\Omega_{\text{DM}} h \simeq 0.12$$

Velocity-dependent

$$\langle \sigma v \rangle \simeq b \langle v^2 \rangle$$

Boltzmann suppression

$$\propto e^{-m_\chi/T}$$

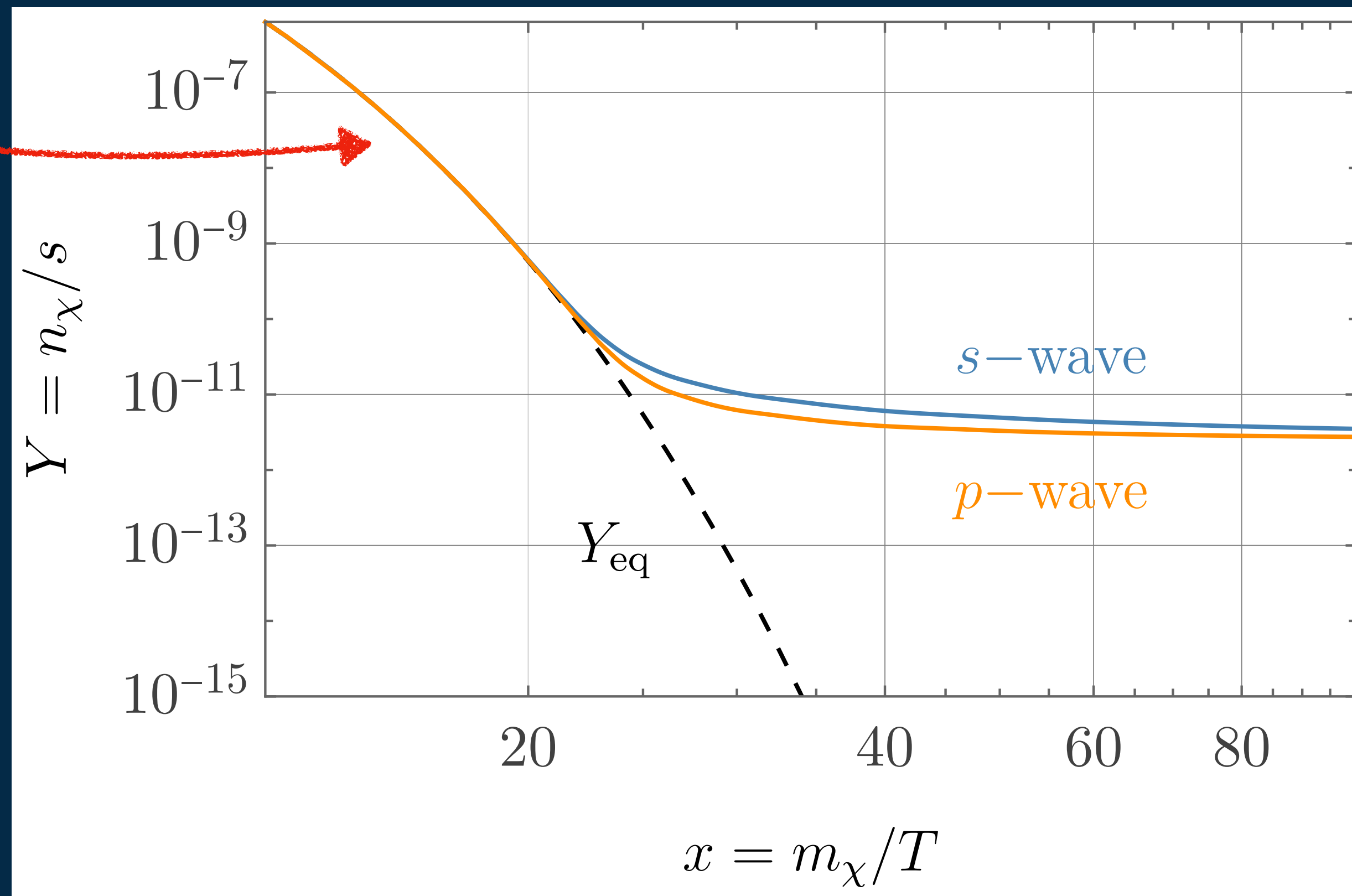
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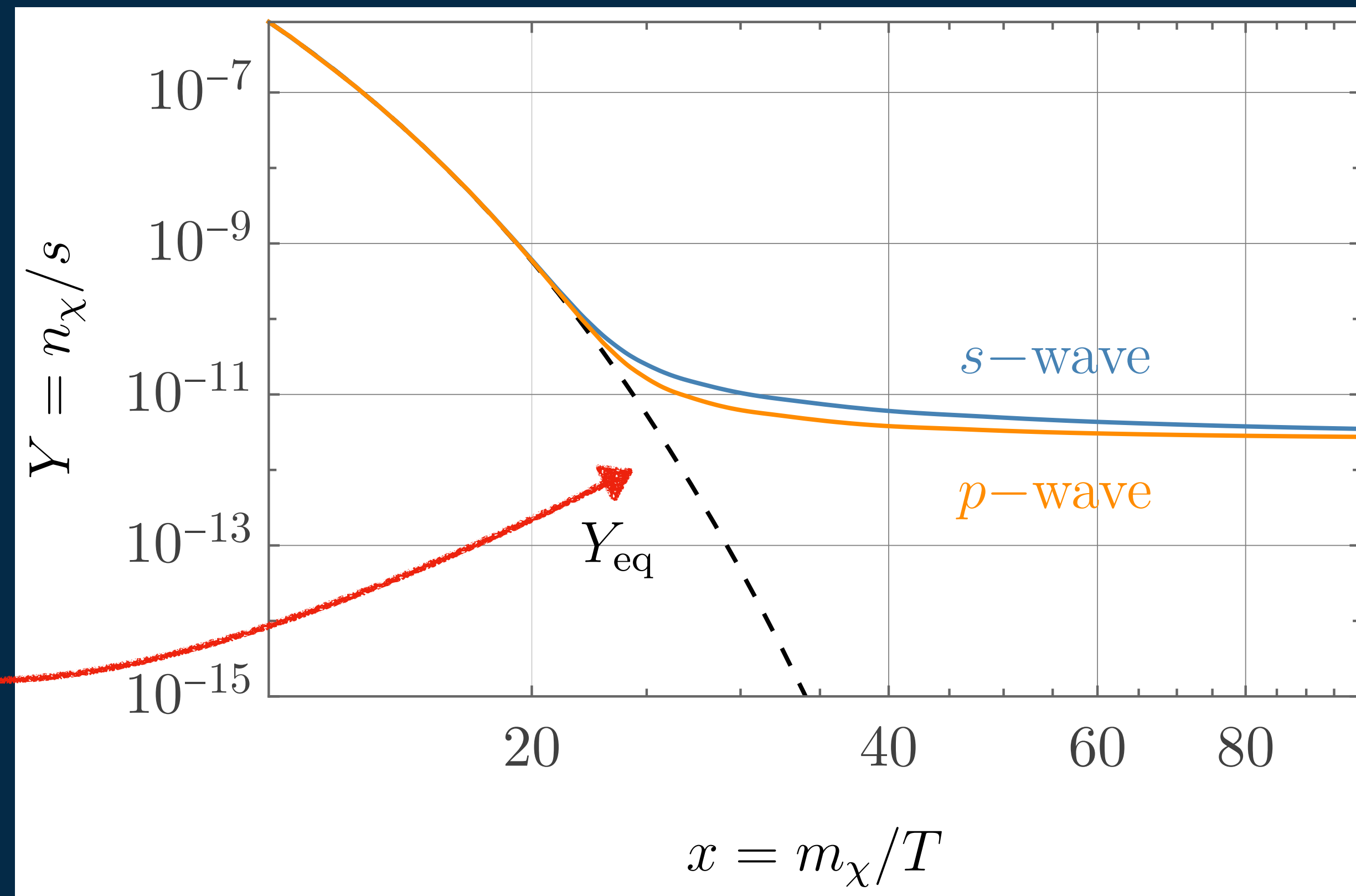
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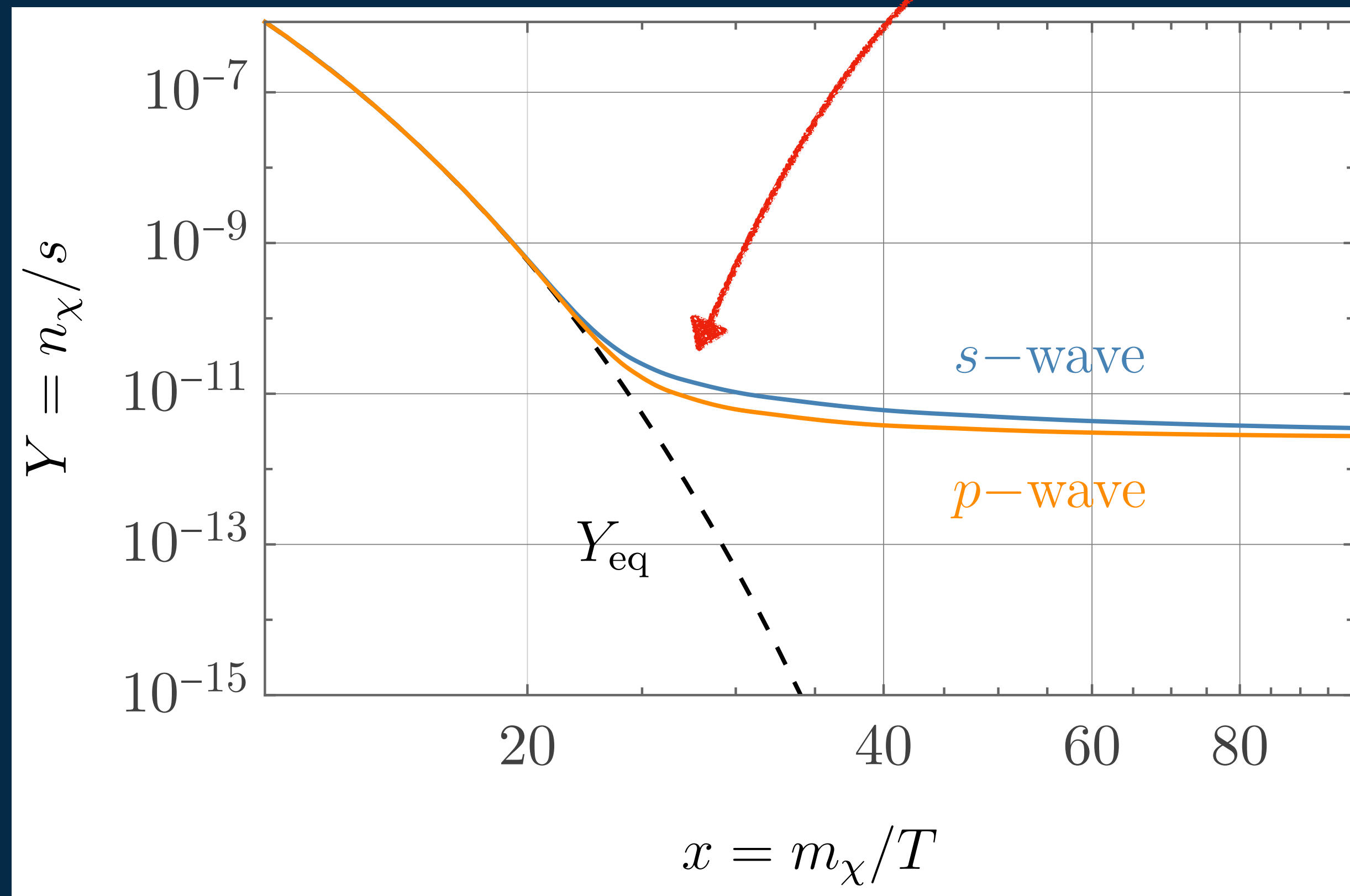
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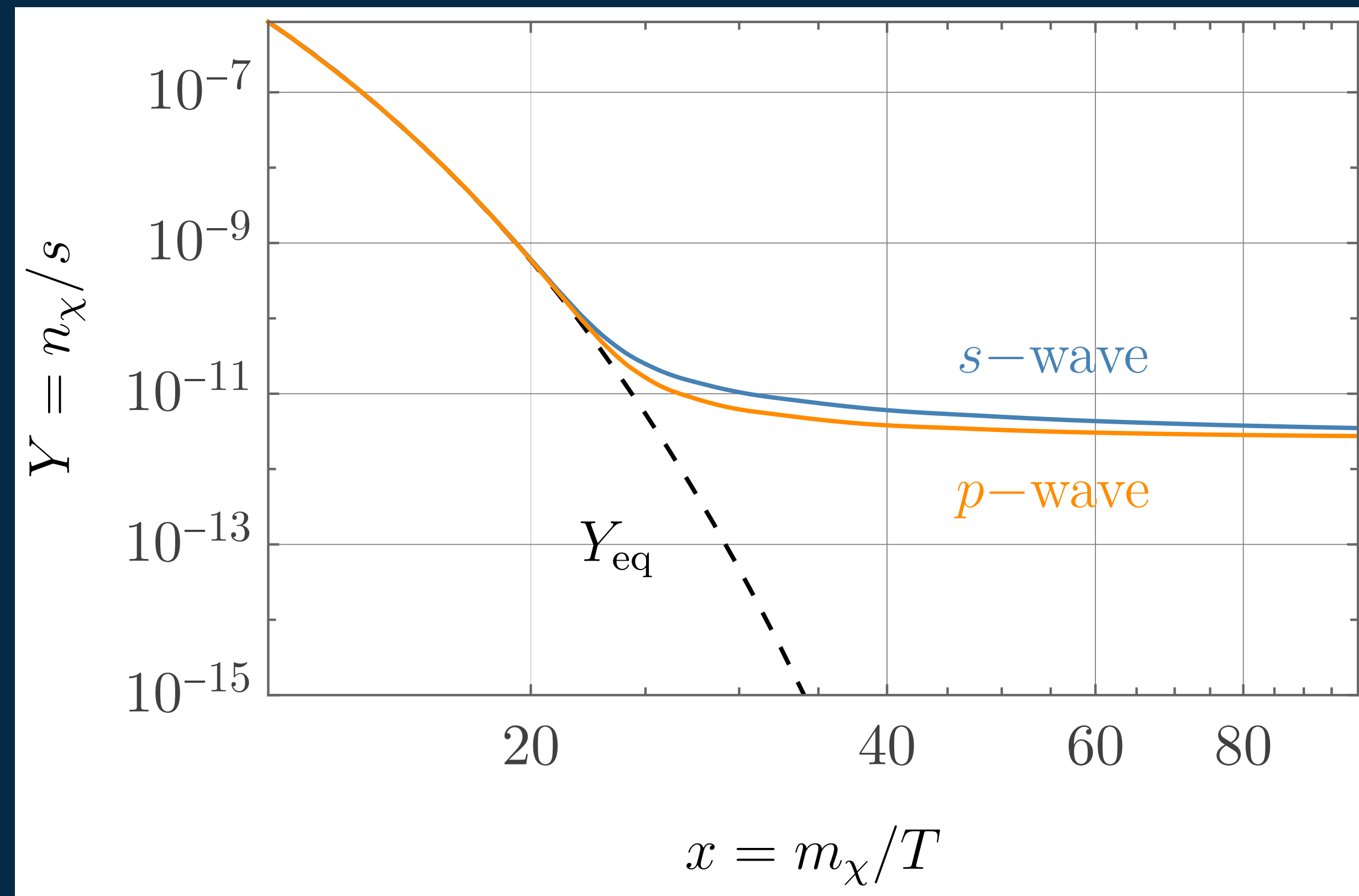
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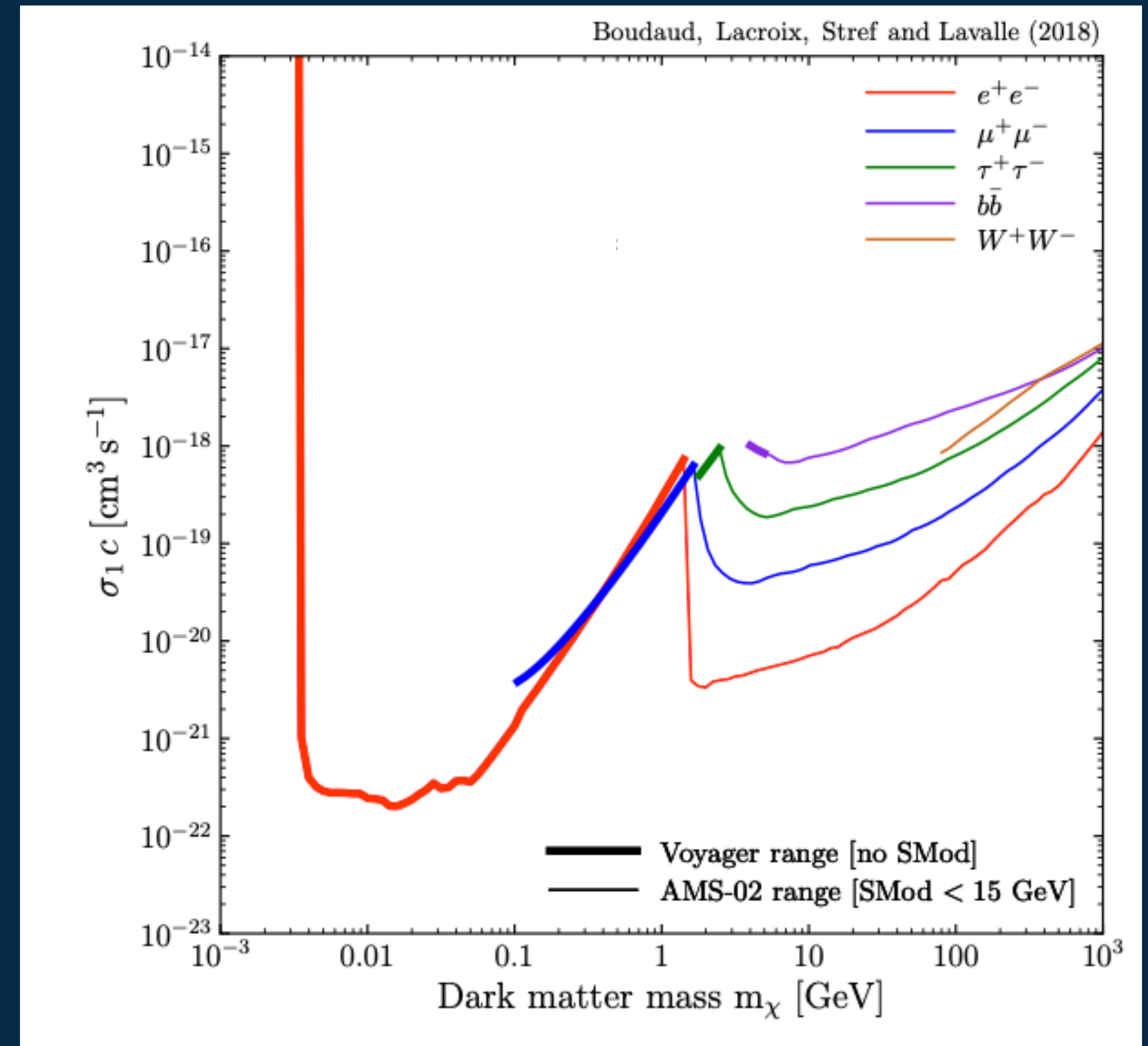
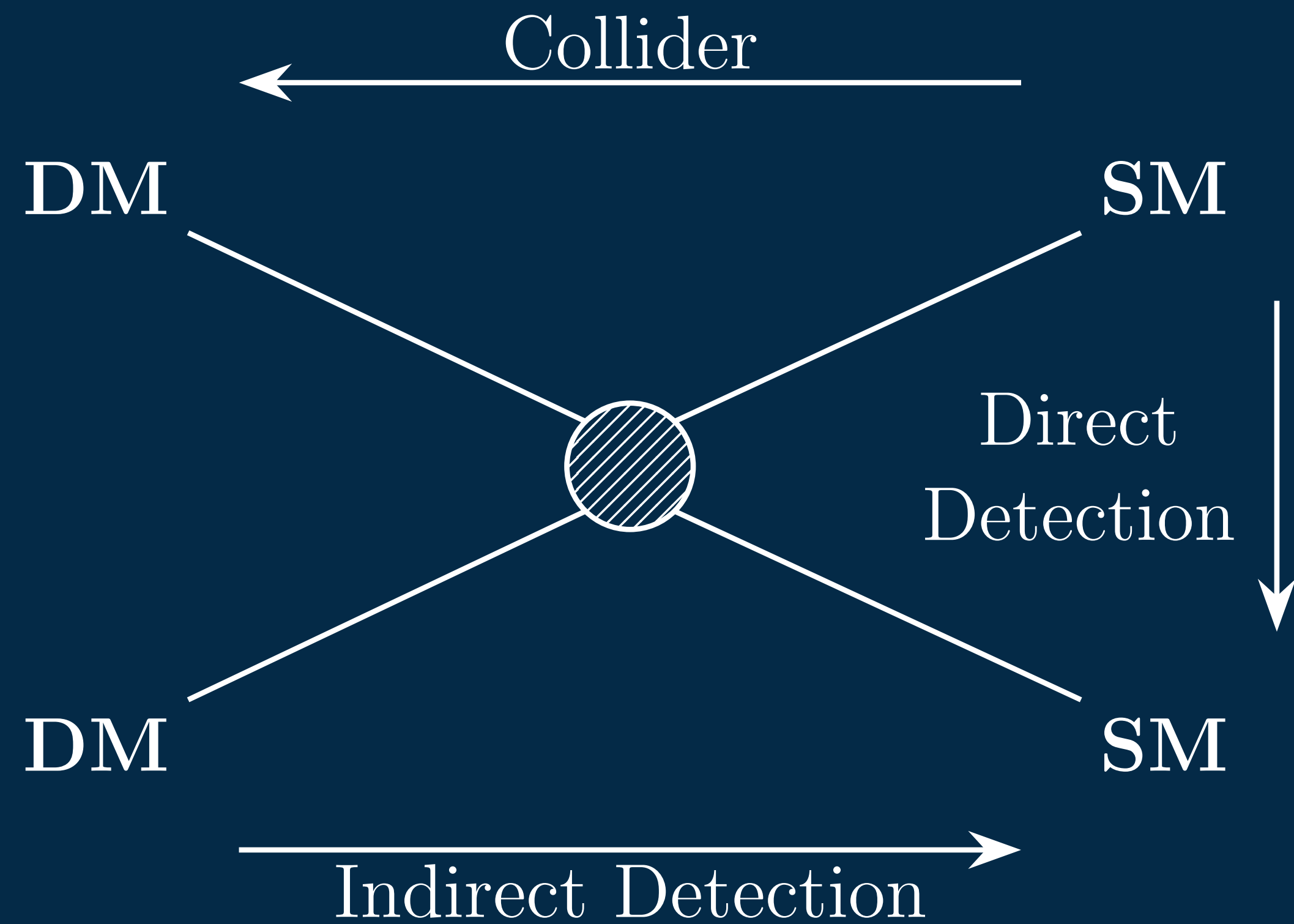
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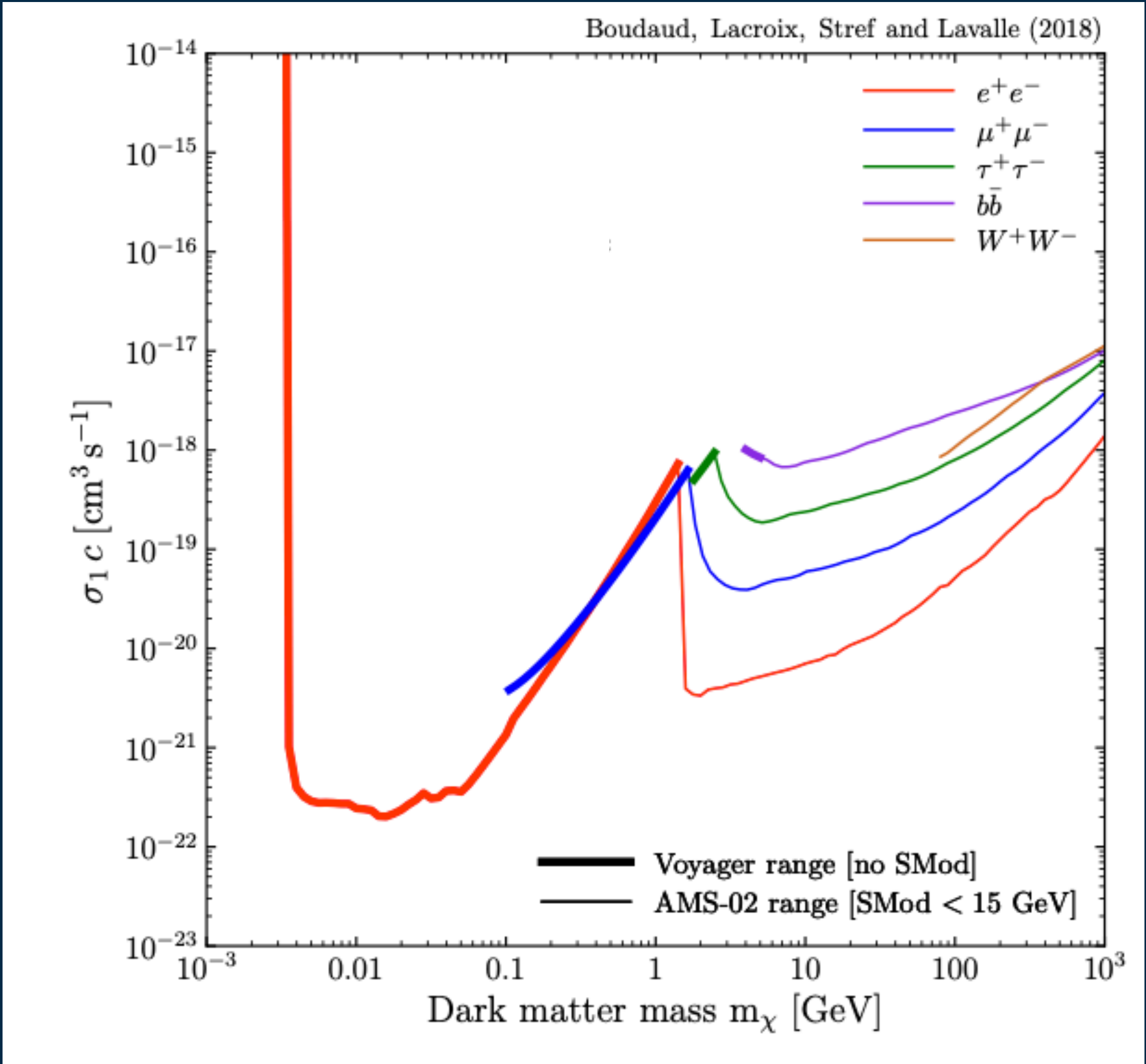
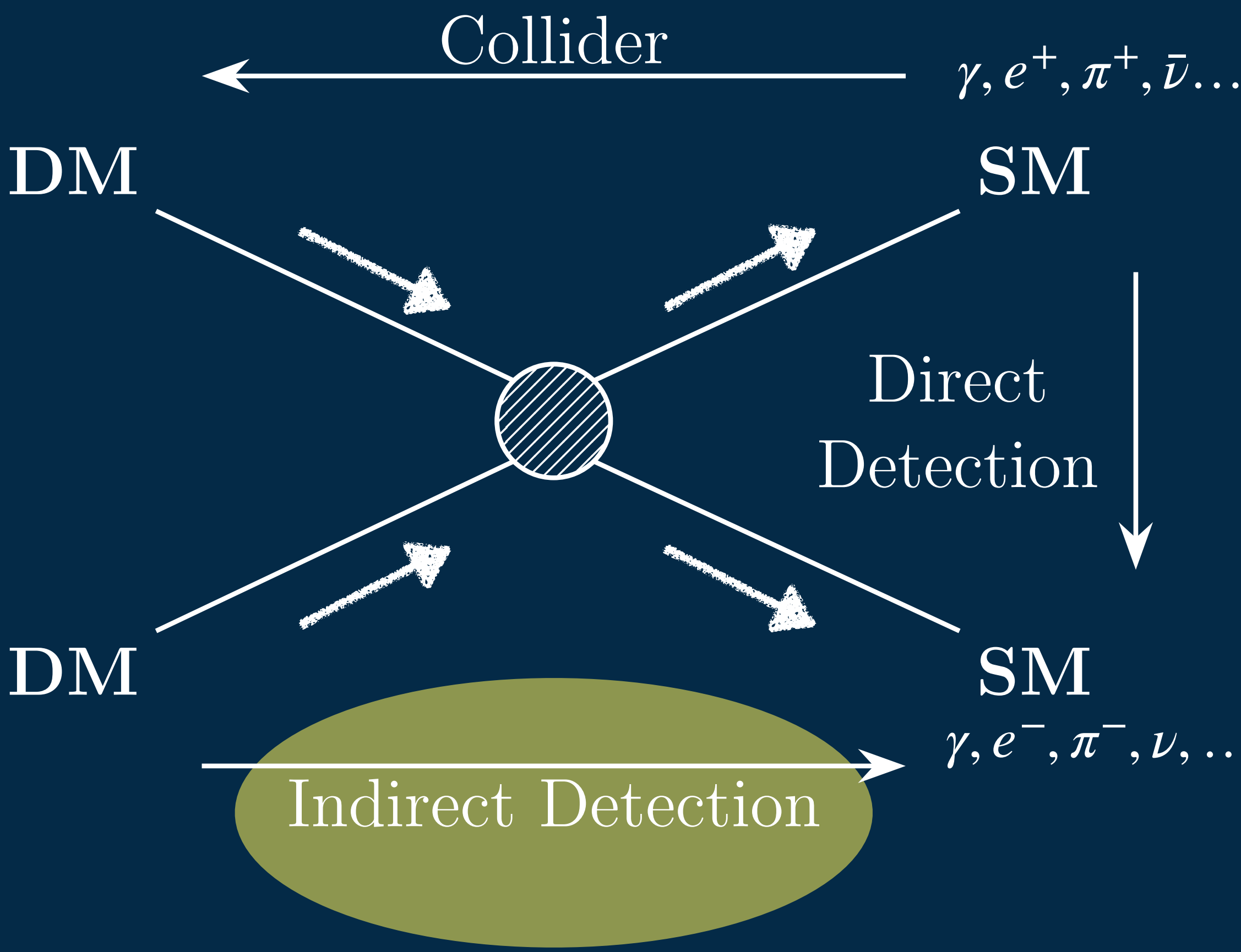
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DM detection



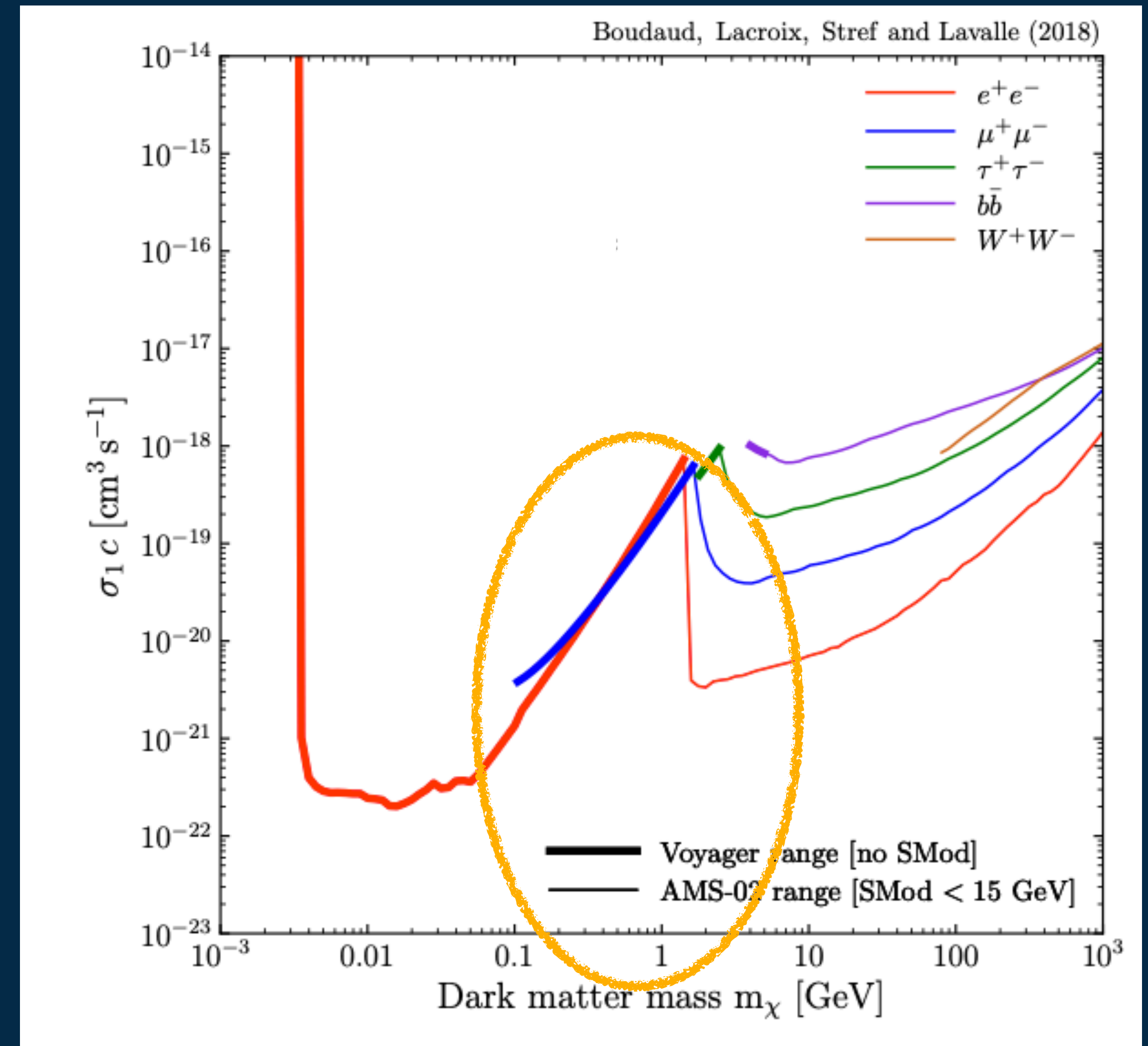
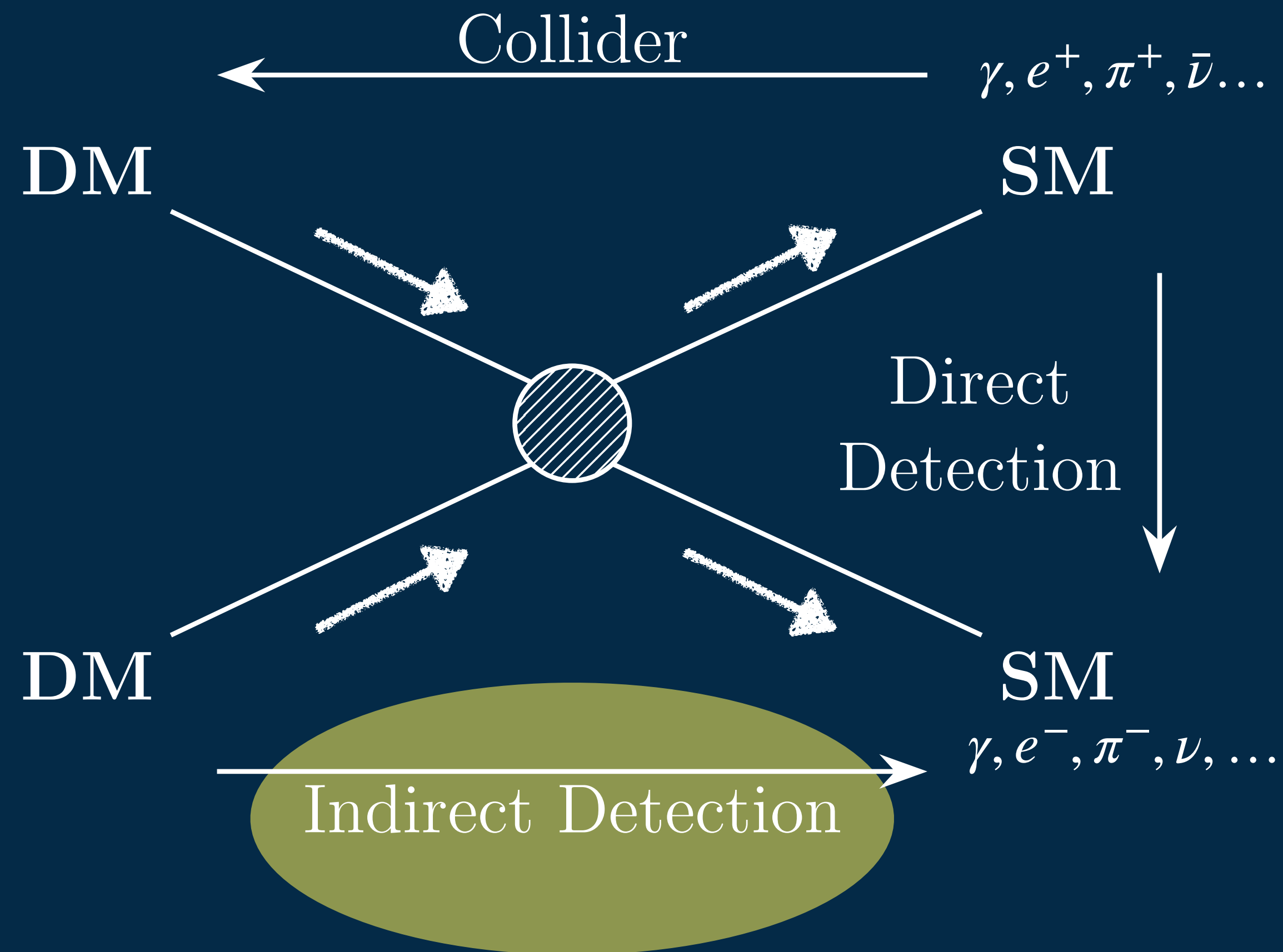
M. Boudaud, T. Lacroix, M. Stref, J. Lavalle 1810.01680

DM detection



M. Boudaud, T. Lacroix, M. Stref, J. Lavalle 1810.01680

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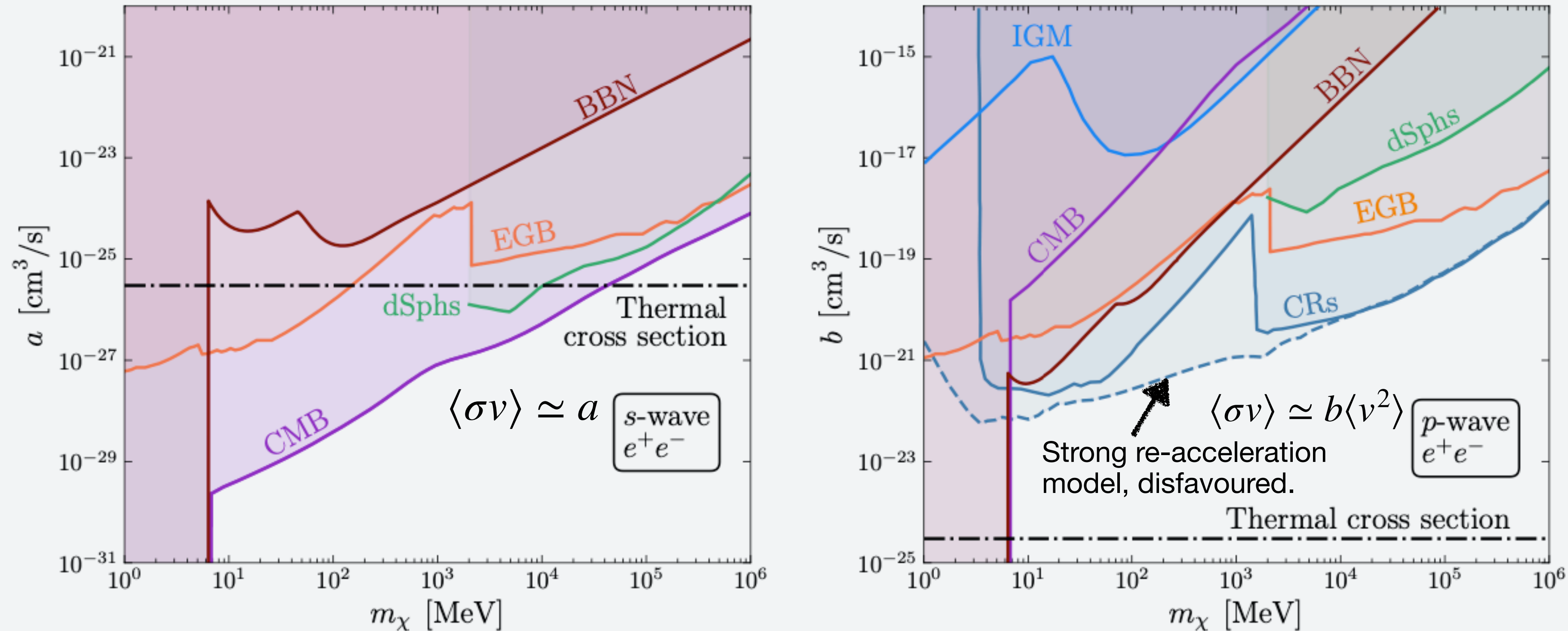


M. Boudaud, T. Lacroix, M. Stref, J. Lavalle 1810.01680

Could BBN be sensitive to this range?

How DM could influence BBN?

P. F. Depta, M. Hufnagel, K. Schmidt-Hoberg, S. Wild 1901.06944

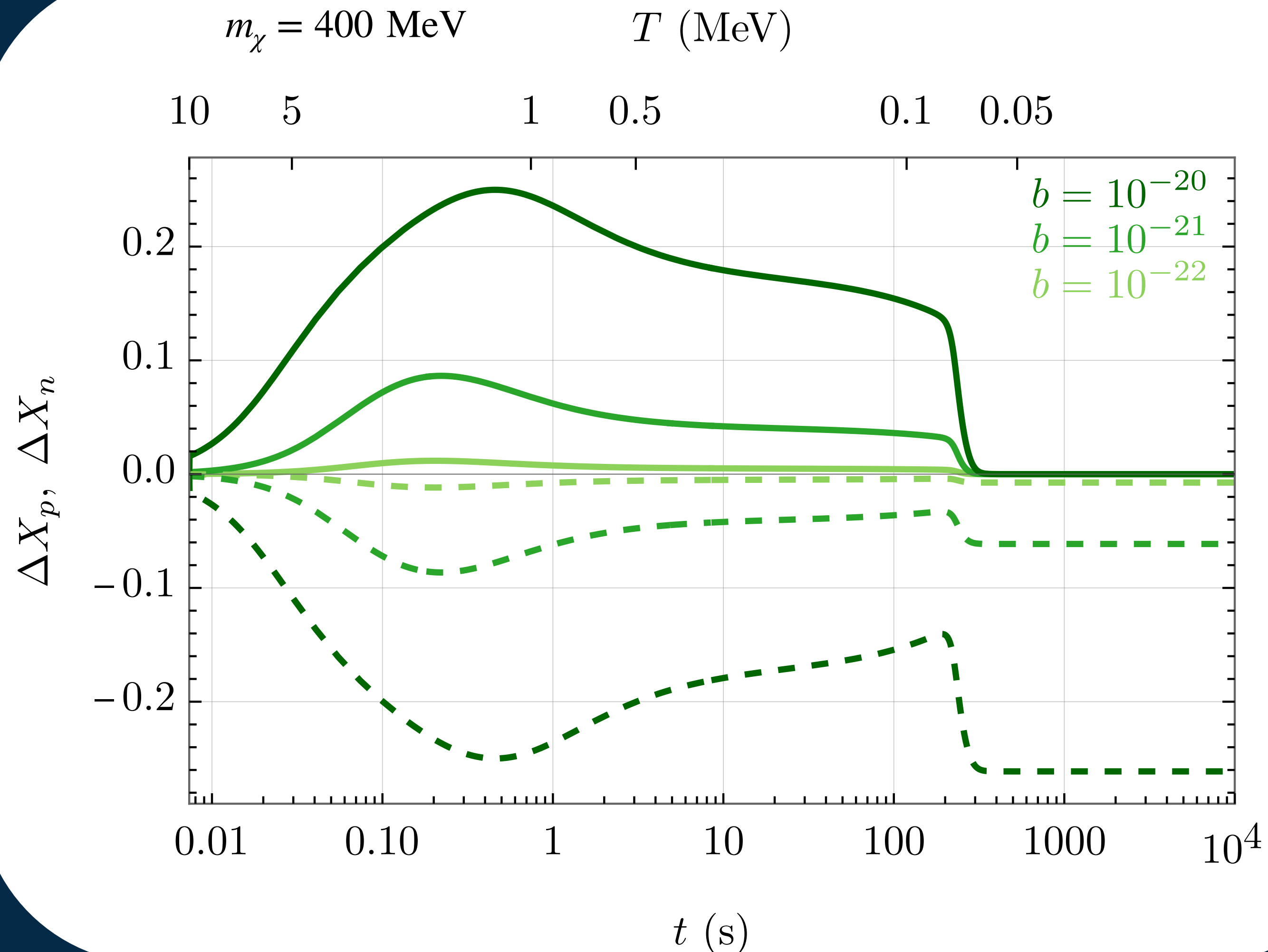


high-redshift intergalactic medium temperature
gamma-ray observations of MW-satellite dwarf galaxies
diffuse extragalactic background

Charged Pion Injection: some analysis considerations

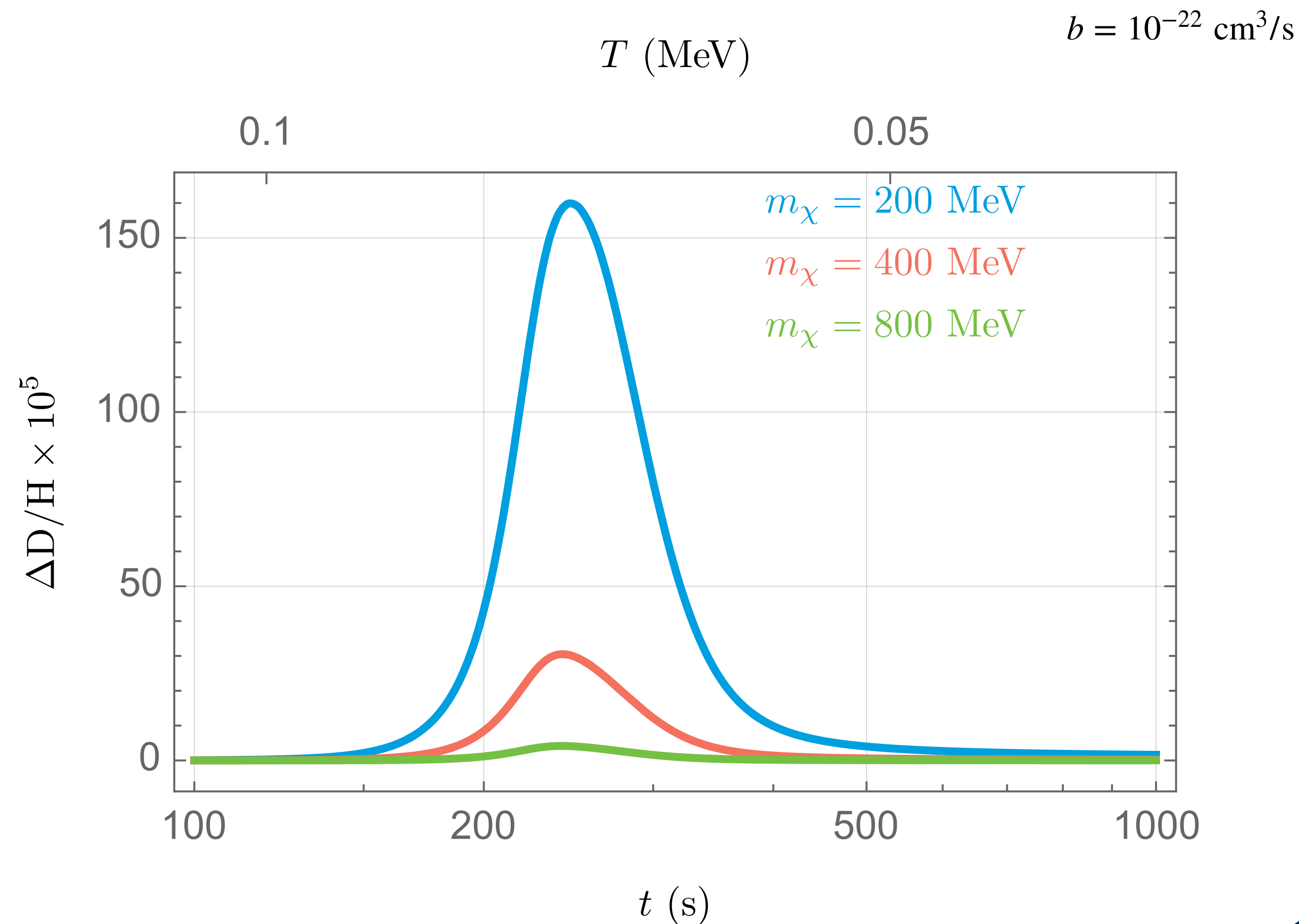
1. DM annihilates entirely into $\pi^\pm$
 2. $Y_\chi = Y_{\text{relic}}(m_\chi)$ remains constant.
 3. Neutrino decoupling is simplified by using $N_\nu \simeq 3.045$ as input.
 4. The thermal evolution is unaltered.
 5. The pions are very efficiently stopped
 6. The most important interactions are with protons and neutrons.
 7. $\pi^+\pi^-$ annihilations, back reactions to DM, and interactions with other mesons are all negligible.
- * $\pi^\pm \rightarrow \mu^\pm + \nu_\mu(\bar{\nu}_\mu)$ accounted for, but products are neglected.

Effects of $\pi^\pm$ injection on neutrons and protons



- Neutrons biased due to Coulomb enhanced $\pi^- + p \rightarrow n$.
- Excess neutrons are later synthesized into excess ^4He .

Generic constraints on velocity suppressed annihilation

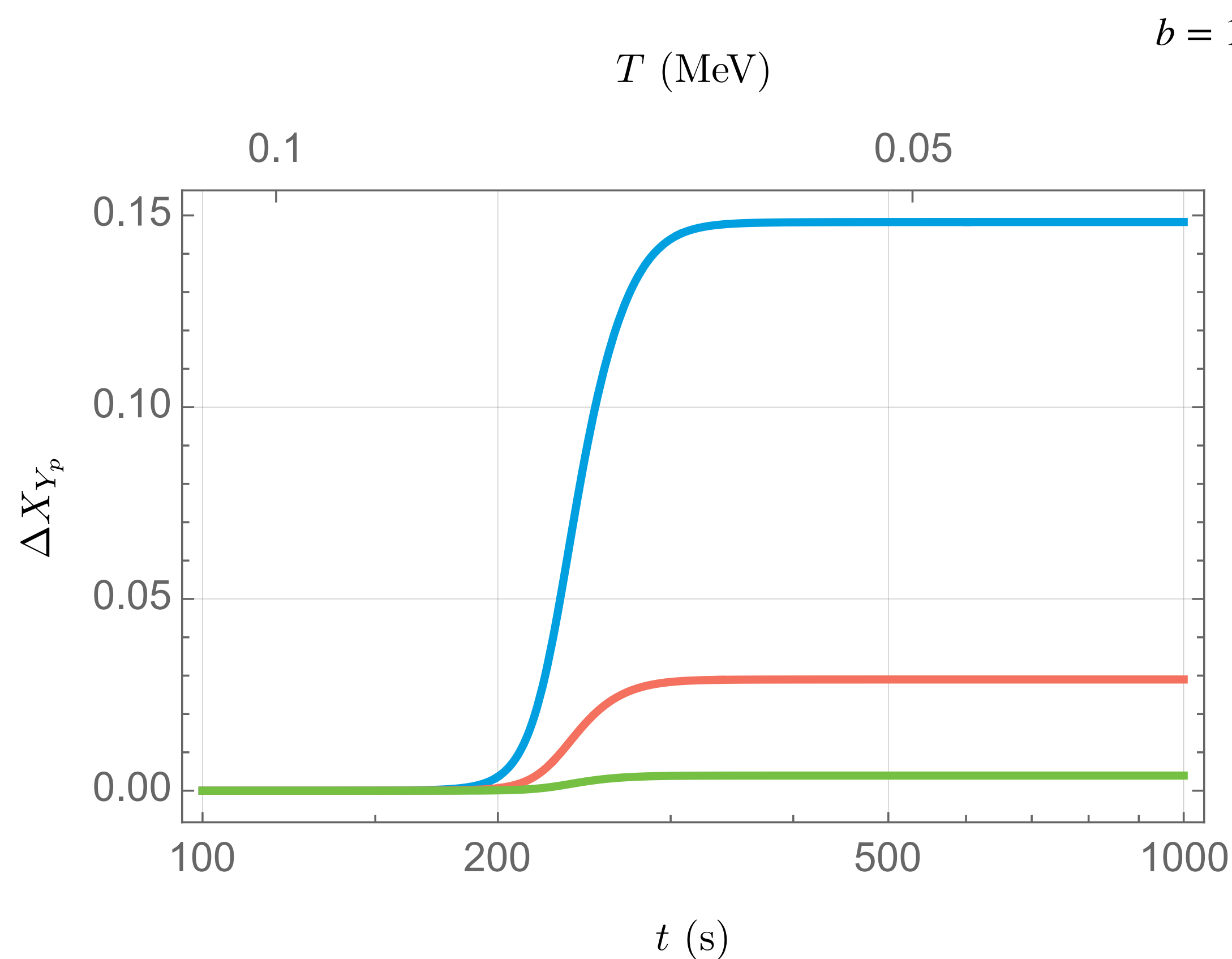


We can get constraints on DM interaction rate by demanding that D and ^4He predictions match observations*

$$D/H \times 10^{-5} = 2.547 \pm 0.025$$

$$Y_p = 0.245 \pm 0.003$$

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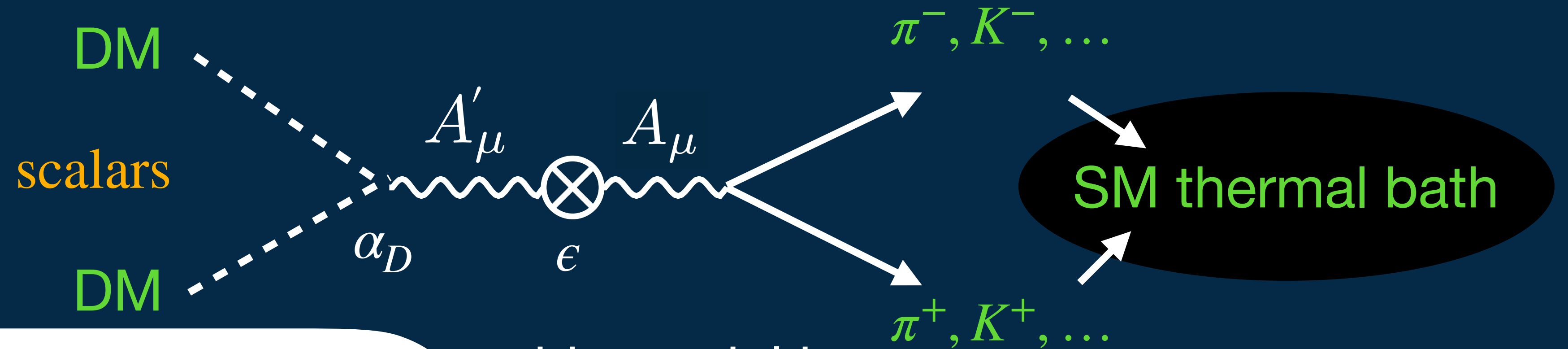
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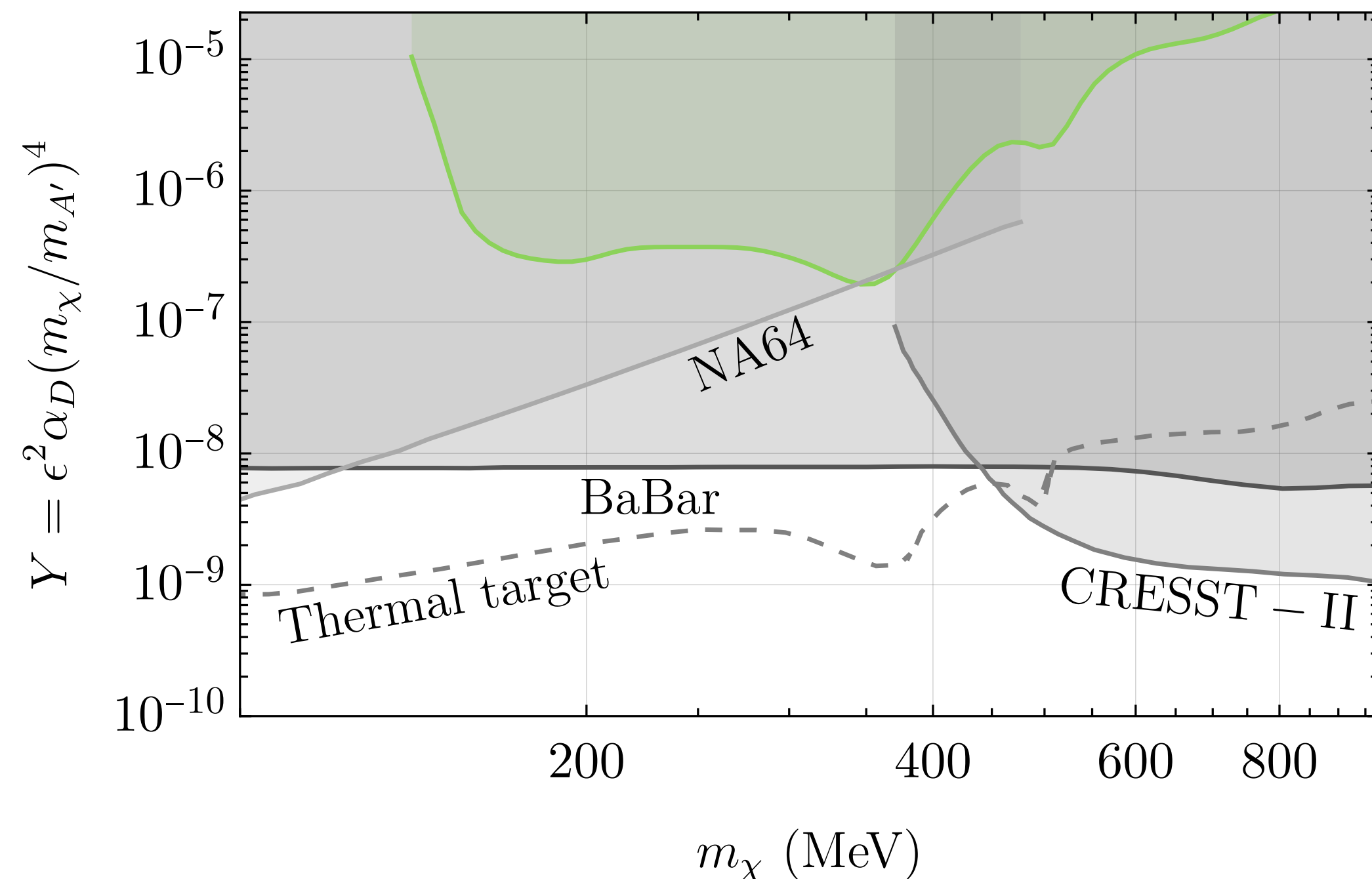
What about neutral mesons?

- Neutral pions have a very short lifetime ($\sim 10^{-18}$ s). The photons they decay to thermalize quickly and have no effect.
- K_S ($\tau \sim 10^{-10}$ s) decays to charged pions (BR $\sim 69\%$), resulting in a similar effect.
- K_L ($\tau \sim 10^{-7}$ s) has a per-particle effect on BBN that is comparable to charged kaons (via hyperon decays).
 - *Effects depend on the model's BR, but are limited to $\sim \mathcal{O}(1)$.
- In the dark photon model, including effects due to neutral mesons could potentially strengthen limits by $\sim 10\%$ in the high mass region (near the ϕ resonance, ~ 1 GeV), so our limits are somewhat conservative.

A DM benchmark model: additional comments



this model imposes a
BR less than 10%



- *The BRs are obtained from on shell DP decay.
- *Although the DP is off shell, we approximate the BR to correspond to the on shell branching for $m_{A'} = 2m_\chi$ (as opposed to $3m_\chi$)

$$\alpha_D = 0.5$$

$$m_{A'}/m_\chi = 3$$