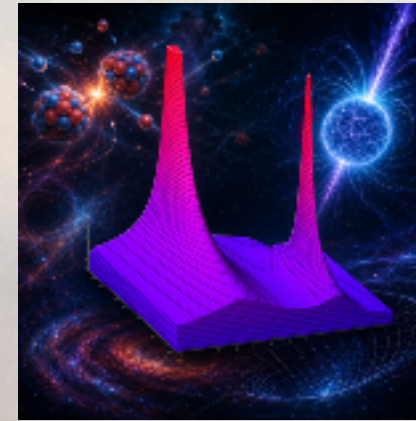


# Four Flavo(u)rs of MDWF

Lattice 2026  
26 July - 1 August

André Walker-Loud



CoSMoN Collaboration  
[[cosmon-collaboration.github.io](https://cosmon-collaboration.github.io)]



# Why a new action?

- ❑ I have personally been interested in generating DWF for  $> 20$  years [chiral symmetry]
  - ❑ “But they are expensive!” “And DWF don’t have multi-grid!” ...
- ❑ We have observed that calculations with MDWF in the valence sector require significantly less samples to achieve the same stochastic uncertainty [we believe this comes from a mix of  $\chi$ -symmetry and smearing]
- ❑ DWF appear to have milder discretization effects [fewer and coarser  $\{a\}$  to achieve same continuum?]
- ❑ DWF are significantly simpler to renormalize [less human time]
- ❑ DWF are more suited to heterogeneous architectures [more compute intense solves and greater percent wall clock time spent in highly optimized solve compared to the other pieces of measurement calcs]
- ❑ Generating physical pion mass ensembles is not an issue (aside from needing a machine larger than Perlmutter to fit within the queue)
- ❑ In aggregate, I suspect DWF are not more expensive than clover, especially when human time is included
- ❑ I was also interested in having a set of configurations to compute correlators involving baryons and scattering that were also free of the typical “politics” of use

# What is our action?

- Lüscher Weisz gauge action  $S_G = \frac{\beta}{2N_c} \sum (c_0 \square + c_1 \square) \quad c_0 = 1 - 8c_1 \quad c_1 = -\frac{1}{12}$
- Möbius Domain Wall Fermions  
 $b - c = 1$ : rescaled Shamir operator  $D^M(M_5) = \frac{(b + c)D^W(M_5)}{2 + (b - c)D^W(M_5)} = (b + c) \frac{D^W(M_5)}{2 + D^W(M_5)} = \alpha D^{\text{Shamir}}(M_5)$
- Dirac operator uses 8-steps of stout smeared gauge links with  $\rho = 0.125$   
[Morningstar & Peardon, PRD 69 2004)]  
 $V_\mu = U_\mu^{(8)}, \quad U_\mu^{(l+1)} = e^{i\rho Q_\mu^{(l)}} U_\mu^{(l)}, \quad D^M[V_\mu](M_5)$
- Most design decisions and all new software needed to generate our configurations are due to  
R. Jamie Hudspith — as well as all preliminary results shown here

# Why 8 steps of stout smearing?

- ❑ Seems like a lot of smearing! 🤖
- ❑ With MDWF/HISQ mixed-action calculations by CalLat with  $t_{gf} = 1 \times a^2$  [Action paper - 1701.07559]
  - $\pi^+ \leftarrow \pi^-$ : 1805.02634
  - $g_A$  @ 0.74% uncertainty: 1805.12130, 2104.05226, 2503.09891
  - $F_K/F_\pi$  @ 0.37% uncertainty: 2005.04795
  - $w_0$  scale @ 0.67% uncertainty: 2011.12166

We observed improved chiral properties of  $D^M[V_\mu]$

Passed several high-precision tests

No sign of over-smearing

- ❑ 8 steps of stout with  $\rho = 0.125$  is similar to gradient-flow of length 1 in lattice units [see perturbative equivalence in Ammer & Durr, PRD 110 (2024) [2406.03493]]
- ❑ Clearly - we need to perform several high-precision tests of the action - see later

# MDWF: $M_5, L_5, b(c), m_l, m_s, m_c$

- ❑ Parameters of the fermion action
- ❑  $M_5 = 1$  for all ensembles: as  $a \rightarrow 0$ , the  $(-1)^t$  modes that appear in temporal dependence functions decouple — [Syritsyn & Negele LATTICE2007 [0710.0425]]
- ❑ tune  $b$  to minimize  $m_{res}$  and set  $L_5$  such that  $\sqrt{t_0}m_{res} \lesssim 0.0009$ , or  $m_{res} \lesssim 1.5$  MeV
- ❑ For each  $\beta$ , tune  $m_s$  such that connected  $\eta_s$   $8t_0m_{\eta_s}^2 = 1.989(32)$  [Dowdall et al, 1110.6887]
- ❑ Set  $am_c = 11.8 \times am_s$  from FLAG determination of ratio
- ❑ Scan  $m_l$  to build up ensembles with different  $m_\pi$
- ❑ Most initial ensembles generated with  $m_\pi L \simeq 3.5$

| $\approx a/\text{fm}$ | $\beta$ | s.h | $b$  | $c$  | $L_5^*$ | $M_5$ | $L_{\text{traj}}^{**}$ | $(am_s, am_c)$   | $10^5 am_{res}$ |
|-----------------------|---------|-----|------|------|---------|-------|------------------------|------------------|-----------------|
| 0.12                  | 4.008   | a12 | 1.75 | 0.75 | 10      | 1.0   | 0.7                    | (0.0725,0.8555)  | 74(5)           |
| 0.10                  | 4.068   | a10 | 1.50 | 0.50 | 8       | 1.0   | 1.0                    | (0.0560,0.6608)  | 43(3)           |
| 0.08                  | 4.160   | a08 | 1.35 | 0.35 | 6       | 1.0   | 1.6                    | (0.0415,0.4897)* | 22(3)           |
| 0.06                  | 4.238   | a06 | 1.20 | 0.20 | 4       | 1.0   | 2.6                    | (0.0305,0.3599)  | 39(1)           |
| 0.05                  | 4.333   | a05 | 1.16 | 0.16 | 4       | 1.0   | 4.0                    | (0.0230,0.2714)  | 23(0)           |

\* Yes, we tried  $L_5 = 2$  😂

\*\* I will return to this

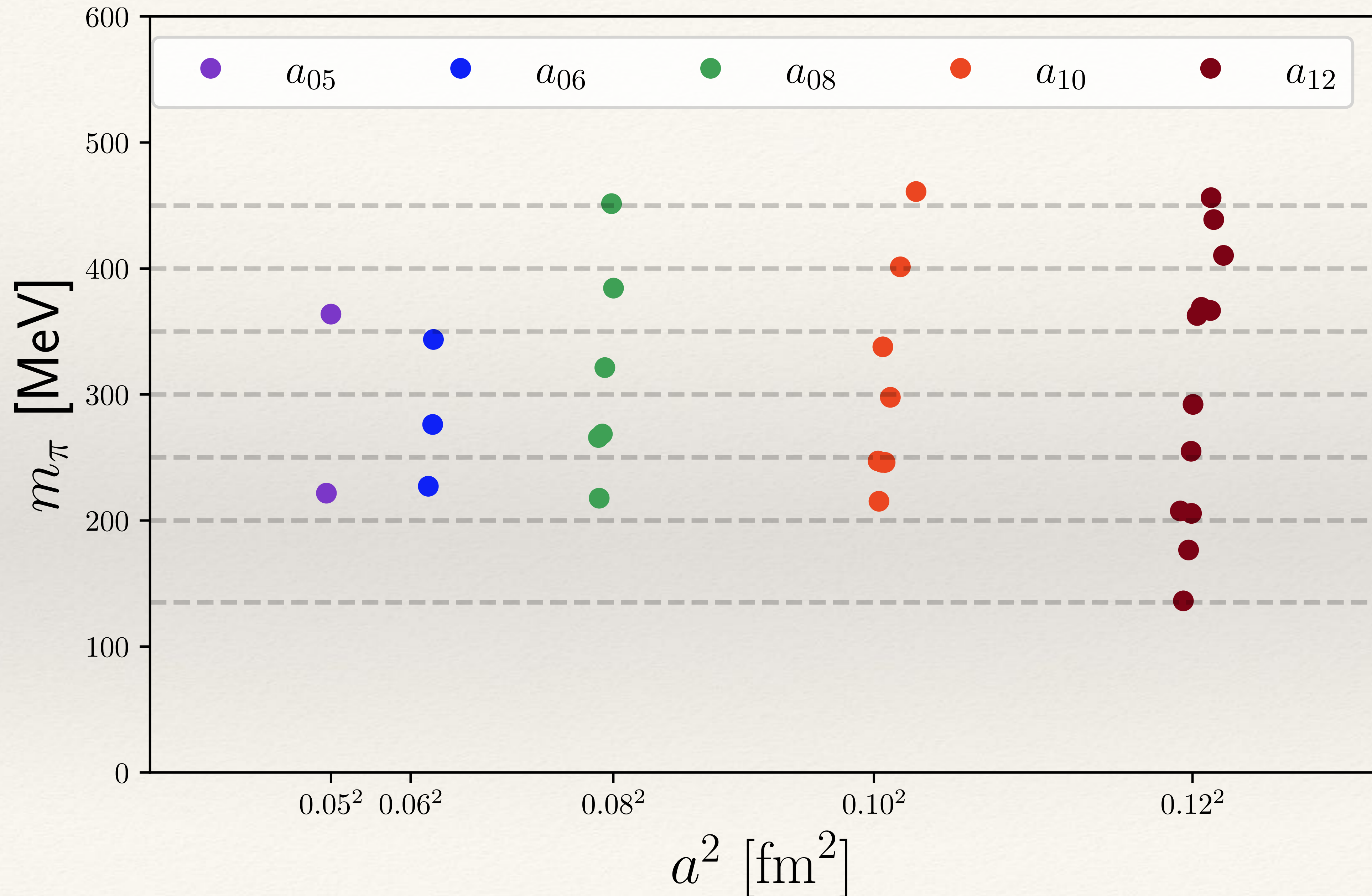
# Landscape of current ensembles

TABLE III. Volume and light quark mass in lattice units for the ensembles used in this work. Additionally, the gradient flow scales ( $t_0$ ,  $w_0$ ), topological charge ( $Q$ ) and their respective autocorrelation lengths are tabulated. The derived values of  $F_K/F_\pi$ ,  $\xi_\pi = m_\pi^2/(4\pi F_\pi)^2$ ,  $\xi_K = m_K^2/(4\pi F_\pi)^2$ ,  $m_\pi L$  and  $\epsilon_a^2 = a^2/(4w_0^2)$  are listed as well as the number of independent configurations used  $N_{\text{cfg}}$  and the number of time-translated sloppy solves  $N_{\text{src}}$ .

| ensemble                                   | $\frac{L}{a}^3 \times \frac{T}{a}$ | $am_l$ | $\sqrt{t_0}/a$ | $w_0/a$     | $Q$      | $\tau_{\sqrt{t_0}}, \tau_{w_0}, \tau_Q$ | $F_K/F_\pi$ | $\xi_\pi$   | $\xi_K$     | $m_\pi L$  | $\epsilon_a^2$ | $N_{\text{cfg}} \times N_{\text{src}}$ |
|--------------------------------------------|------------------------------------|--------|----------------|-------------|----------|-----------------------------------------|-------------|-------------|-------------|------------|----------------|----------------------------------------|
| $\beta = 4.008, a \approx 0.12 \text{ fm}$ |                                    |        |                |             |          |                                         |             |             |             |            |                |                                        |
| a12m456                                    | $16^3 \times 32$                   | 0.0300 | 1.1961(31)     | 1.4766(82)  | 0.3(2)   | (16,16,16)                              | 1.0800(20)  | 0.1262(19)  | 0.2106(28)  | 4.424(18)  | 0.1147(13)     | $173 \times 32$                        |
| a12m436                                    | $16^3 \times 32$                   | 0.0270 | 1.1946(20)     | 1.4766(82)  | 0.4(3)   | (16,24,24)                              | 1.0898(24)  | 0.1225(15)  | 0.2178(24)  | 4.255(16)  | 0.1141(15)     | $179 \times 32$                        |
| a12m410                                    | $16^3 \times 32$                   | 0.0235 | 1.1891(20)     | 1.4766(82)  | -0.2(3)  | (16,16,16)                              | 1.1012(25)  | 0.1091(17)  | 0.2150(27)  | 3.979(18)  | 0.1162(11)     | $118 \times 32$                        |
| a12m375                                    | $16^3 \times 32$                   | 0.0195 | 1.2016(40)     | 1.5086(132) | 0.2(3)   | (24,24,16)                              | 1.1109(33)  | 0.0918(17)  | 0.2070(30)  | 3.580(19)  | 0.1098(20)     | $106 \times 32$                        |
| a12m375M                                   | $20^3 \times 48$                   | 0.0195 | 1.1969(13)     | 1.5086(132) | 0.2(3)   | (24,24,16)                              | 1.1061(17)  | 0.09031(66) | 0.2048(12)  | 4.449(11)  | 0.11275(67)    | $204 \times 48$                        |
| a12m375L                                   | $24^3 \times 48$                   | 0.0195 | 1.2040(40)     | 1.5138(57)  | 0.4(5)   | (16,16,8)                               | 1.1007(20)  | 0.08770(76) | 0.2003(16)  | 5.276(11)  | 0.10909(85)    | $95 \times 32$                         |
| a12m363                                    | $16^3 \times 32$                   | 0.0190 | 1.1964(13)     | 1.4943(39)  | -0.1(2)  | (8,8,12)                                | 1.1198(32)  | 0.0959(17)  | 0.2221(30)  | 3.550(20)  | 0.11197(59)    | $235 \times 32$                        |
| a12m292                                    | $20^3 \times 40$                   | 0.0120 | 1.2064(7)      | 1.5258(27)  | 0.6(1)   | (8,16,8)                                | 1.1382(30)  | 0.06587(76) | 0.2149(19)  | 3.541(12)  | 0.10738(39)    | $353 \times 20$                        |
| a12m250                                    | $24^3 \times 48$                   | 0.0090 | 1.2075(17)     | 1.5319(46)  | -0.3(6)  | (16,16,16)                              | 1.1520(26)  | 0.05215(49) | 0.2137(16)  | 3.7086(93) | 0.10654(64)    | $116 \times 48$                        |
| a12m208                                    | $28^3 \times 64$                   | 0.0057 | 1.2138(6)      | 1.5528(28)  | -0.6(3)  | (12,24,8)                               | 1.1743(28)  | 0.03752(30) | 0.2204(14)  | 3.5223(73) | 0.10368(38)    | $176 \times 64$                        |
| a12m208M                                   | $32^3 \times 64$                   | 0.0057 | 1.2073(6)      | 1.5639(20)  | 0.6(5)   | (16,16,8)                               | 1.1606(21)  | 0.03548(23) | 0.2118(10)  | 3.9883(75) | 0.10584(28)    | $154 \times 64$                        |
| a12m176                                    | $32^3 \times 64$                   | 0.0040 | 1.2090(8)      | 1.5460(22)  | 0.3(4)   | (16,16,16)                              | 1.1792(20)  | 0.02720(20) | 0.2147(11)  | 3.4247(72) | 0.10460(31)    | $157 \times 64$                        |
| a12m135                                    | $48^3 \times 96$                   | 0.0020 | 1.2120(3)      | 1.5514(11)  | -0.1(10) | (12,24,12)                              | 1.1894(36)  | 0.01733(11) | 0.22251(99) | 3.9631(74) | 0.10386(21)    | $101 \times 48$                        |
| $\beta = 4.068, a \approx 0.10 \text{ fm}$ |                                    |        |                |             |          |                                         |             |             |             |            |                |                                        |
| a10m460                                    | $16^3 \times 32$                   | 0.0220 | 1.4076(19)     | 1.7175(48)  | -0.1(1)  | (12,12,24)                              | 1.0995(40)  | 0.1336(25)  | 0.2236(36)  | 3.745(22)  | 0.08475(47)    | $483 \times 32$                        |
| a10m401                                    | $20^3 \times 40$                   | 0.0185 | 1.4223(21)     | 1.7589(75)  | -0.1(3)  | (16,16,40)                              | 1.0963(25)  | 0.1010(15)  | 0.1975(24)  | 4.074(18)  | 0.08081(75)    | $128 \times 40$                        |
| a10m340                                    | $20^3 \times 40$                   | 0.0135 | 1.4392(30)     | 1.8227(106) | -0.5(3)  | (16,16,16)                              | 1.1259(51)  | 0.0822(19)  | 0.2090(37)  | 3.430(21)  | 0.07525(90)    | $150 \times 40$                        |
| a10m300                                    | $24^3 \times 48$                   | 0.0100 | 1.4319(17)     | 1.7915(77)  | -0.7(5)  | (20,40,40)                              | 1.1286(26)  | 0.06275(79) | 0.1938(18)  | 3.628(14)  | 0.07789(70)    | $174 \times 24$                        |
| a10m247                                    | $28^3 \times 64$                   | 0.0067 | 1.4439(11)     | 1.8276(27)  | -0.5(7)  | (12,12,24)                              | 1.1495(27)  | 0.04700(47) | 0.1979(13)  | 3.514(12)  | 0.07484(22)    | $115 \times 64$                        |
| a10m247M                                   | $32^3 \times 64$                   | 0.0067 | 1.4370(12)     | 1.8128(38)  | -0.5(3)  | (16,16,8)                               | 1.1464(31)  | 0.04659(46) | 0.1988(15)  | 3.998(13)  | 0.07607(32)    | $115 \times 32$                        |
| a10m214                                    | $32^3 \times 64$                   | 0.0050 | 1.4430(11)     | 1.8266(30)  | 0.5(4)   | (12,12,12)                              | 1.1646(39)  | 0.03770(48) | 0.2026(19)  | 3.497(12)  | 0.07493(25)    | $100 \times 32$                        |
| $\beta = 4.160, a \approx 0.08 \text{ fm}$ |                                    |        |                |             |          |                                         |             |             |             |            |                |                                        |
| a08m430                                    | $20^3 \times 40$                   | 0.0170 | 1.8136(51)     | 2.2291(146) | 0.4(2)   | (20,20,30)                              | 1.1242(58)  | 0.1378(41)  | 0.2396(62)  | 3.626(28)  | 0.05031(67)    | $233 \times 20$                        |
| a08m375                                    | $24^3 \times 48$                   | 0.0165 | 1.8097(46)     | 2.2108(118) | 0.0(5)   | (20,20,30)                              | 1.1195(44)  | 0.0948(22)  | 0.2049(38)  | 3.705(21)  | 0.05115(58)    | $112 \times 24$                        |
| a08m323                                    | $32^3 \times 64$                   | 0.0090 | 1.8268(14)     | 2.2637(35)  | -0.6(3)  | (15,15,30)                              | 1.1215(14)  | 0.06837(45) | 0.18874(98) | 4.1297(83) | 0.04879(15)    | $307 \times 64$                        |
| a08m271                                    | $32^3 \times 64$                   | 0.0060 | 1.8318(24)     | 2.2767(55)  | 0.1(5)   | (15,15,25)                              | 1.1580(37)  | 0.05323(80) | 0.2003(22)  | 3.453(13)  | 0.04823(23)    | $106 \times 32$                        |
| a08m271rs                                  | $32^3 \times 64^*$                 | 0.0060 | 1.8398(29)     | 2.3017(73)  | 0.6(7)   | (20,20,40)                              | 1.1538(32)  | 0.05315(65) | 0.1972(18)  | 3.417(12)  | 0.04719(31)    | $101 \times 64$                        |
| a08m220                                    | $40^3 \times 80$                   | 0.0039 | 1.8383(12)     | 2.2304(28)  | 0.1(8)   | (15,15,40)                              | 1.1703(51)  | 0.03777(38) | 0.2028(15)  | 3.4971(96) | 0.04708(12)    | $147 \times 20$                        |
| $\beta = 4.238, a \approx 0.06 \text{ fm}$ |                                    |        |                |             |          |                                         |             |             |             |            |                |                                        |
| a06m343                                    | $32^3 \times 64$                   | 0.0070 | 2.3142(49)     | 2.8498(164) | 0.4(3)   | (30,40,60)                              | 1.1379(34)  | 0.0815(14)  | 0.2023(27)  | 3.491(15)  | 0.03078(38)    | $172 \times 64$                        |
| a06m275                                    | $40^3 \times 80$                   | 0.0045 | 2.3173(20)     | 2.8617(76)  | 1.5(5)   | (10,20,30)                              | 1.1453(32)  | 0.05389(67) | 0.1879(18)  | 3.508(13)  | 0.03053(17)    | $149 \times 40$                        |
| $\beta = 4.333, a \approx 0.05 \text{ fm}$ |                                    |        |                |             |          |                                         |             |             |             |            |                |                                        |
| a05m364                                    | $40^3 \times 80$                   | 0.0060 | 2.8964(55)     | 3.5177(132) | 0.8(8)   | (30,30,250)                             | 1.1225(43)  | 0.0827(21)  | 0.1886(36)  | 3.675(26)  | 0.02020(15)    | $159 \times 20$                        |
| a05m218                                    | $64^3 \times 128$                  | 0.0021 | 2.9334(23)     | 3.6199(77)  | -1.3(12) | (15,20,40)                              | 1.1704(46)  | 0.03699(51) | 0.1914(20)  | 3.583(15)  | 0.019078(88)   | $102 \times 16$                        |

□ Ensembles from  
 $F_K/F_\pi$   
[2605.06560]

# Landscape of current ensembles



# HMC Integrator

- We use a 3-level integrator

$$\text{Det}(M) = \text{Sdet} \left( \frac{m_l}{h_1} \right) \boxed{\text{Sdet} \left( \frac{h_l}{h_2} \right)} \cdots \text{Sdet} \left( \frac{h_{n-1}}{h_n} \right) \text{Sdet} \left( \frac{h_n}{m_{\text{PV}}} \right) \text{Sdet}_{\text{EOFA}} \left( \frac{m_s}{m_c} \right) \text{Sdet} \left( \frac{m_c}{m_{\text{PV}}} \right) \Bigg| S_{\text{Gauge}}$$

Level 1

OMF2

$O(10)$  steps

Level 2

OMF2

1 step

Level 3

OMF4

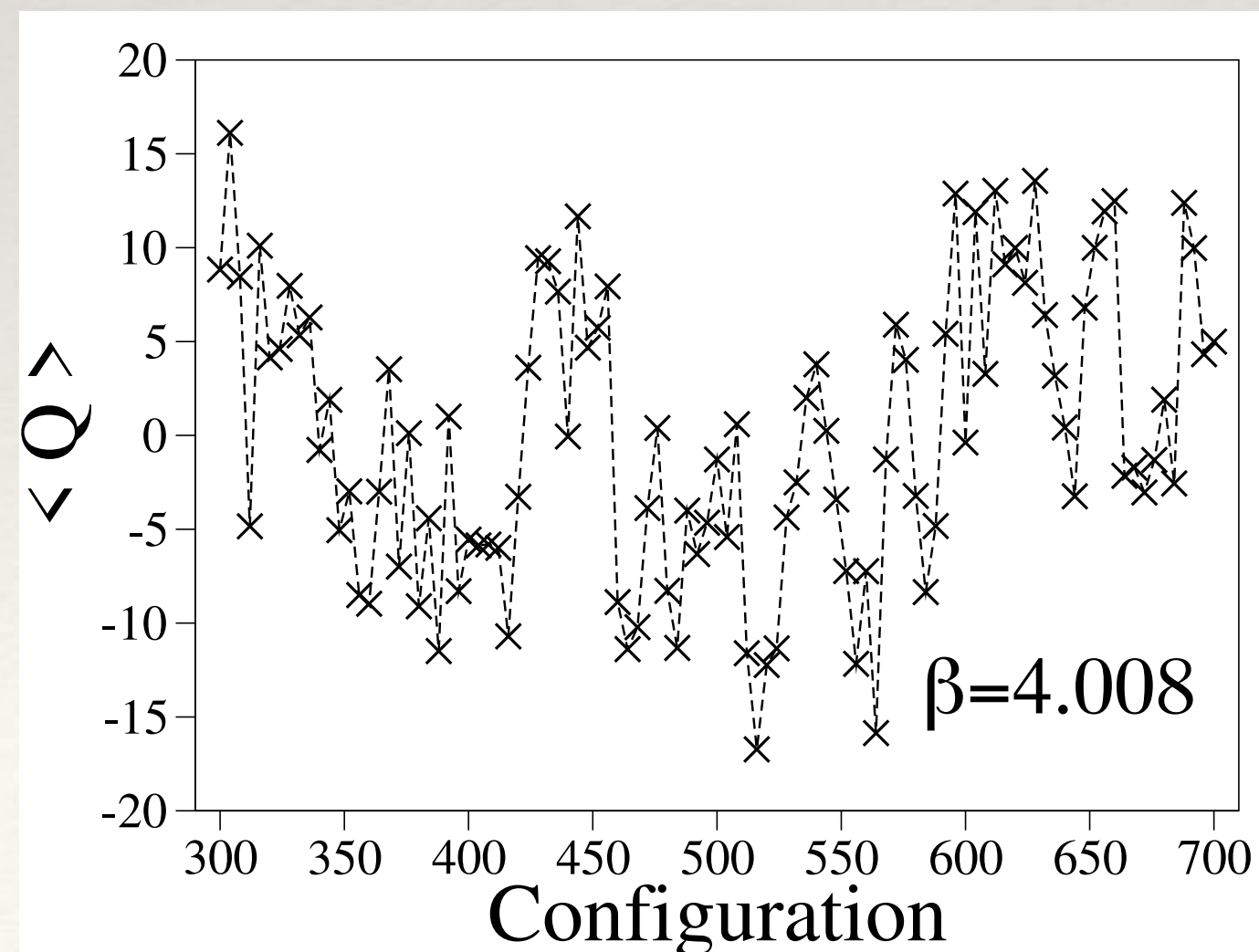
1 – 2 steps

- We tune the number and value of the Hasenbusch mass ratios to keep “Fdt” on each level  $\lesssim 0.4$   
Therefore, at least  $\text{Sdet} \left( \frac{m_l}{h_1} \right)$  is on Level 1, with additional Hasenbush ratios added to Level 1 as needed to keep Fdt small
- Acceptance is kept 70 %  $\lesssim$  acceptance rate  $\lesssim$  80 %
- For new physical pion mass ensembles with  $a \simeq \{0.10, 0.08\}$  fm and  $V = \{64^3 \times 128 \times 8, 80^3 \times 160 \times 6\}$ , the **Force-Gradient-Integrator** for the light quarks has become cost effective

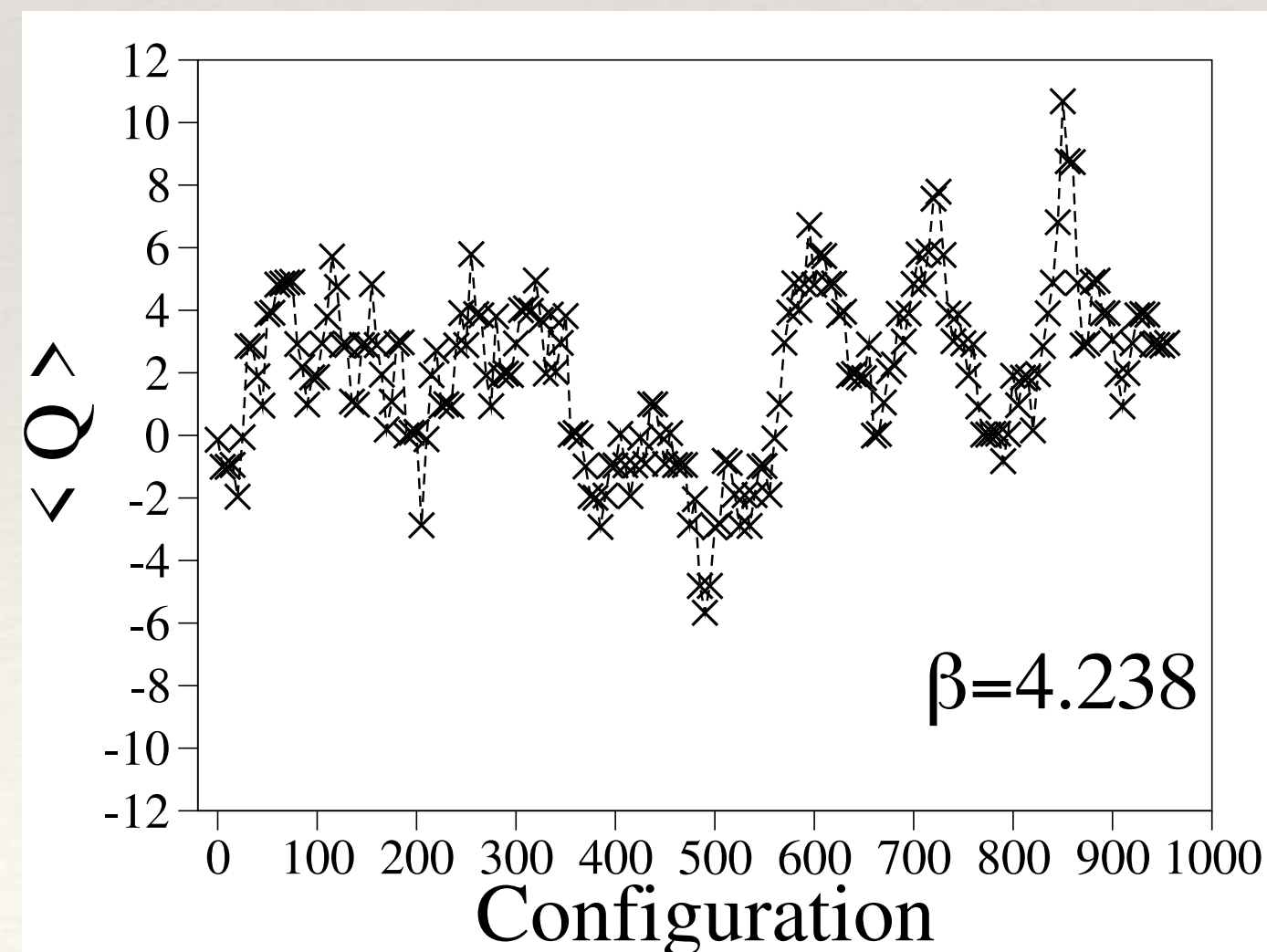
# Q

- ❑ Rather than keeping the MC integration time fixed as a function of  $\beta$ , the MD integration length is held fixed in physical units.
- ❑ We use  $t_0$  as a proxy for this physical scale and on the 2nd coarsest, we have  $L_{\text{traj}}^{\beta=4.068} = 1$
- ❑ This leads to setting  $L_{\text{traj}} = \{0.7, 1.0, 1.6, 2.6, 4.0\}$  on the  $a \simeq \{0.12, 0.10, 0.08, 0.06, 0.05\}$  fm ensembles
- ❑ We find that  $Q$  keeps fluctuating down to the finest ensemble  
We speculate (but need to do a careful comparison) that the longer integration time improves the fluctuations

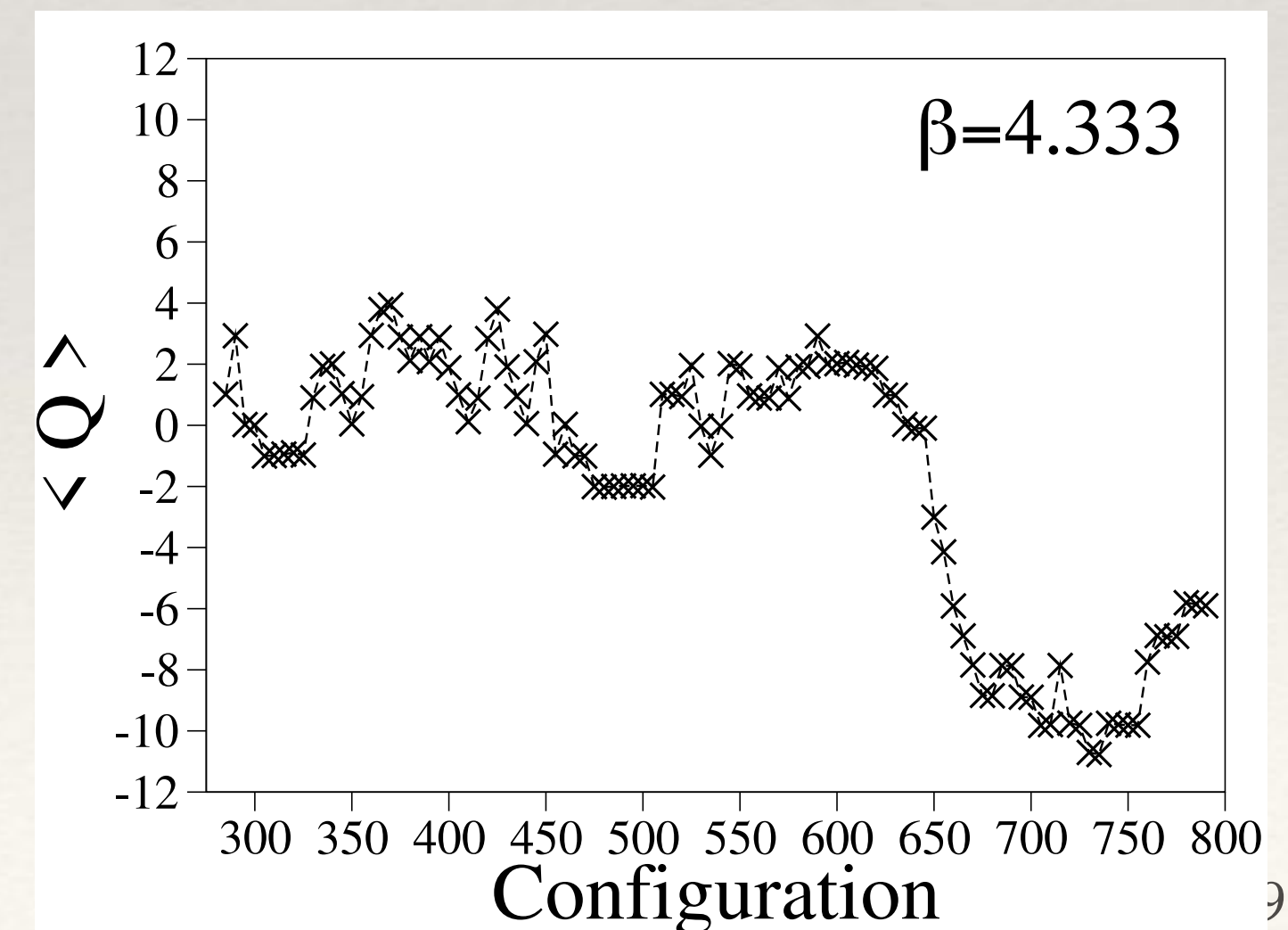
$a \simeq 0.12$  fm



$a \simeq 0.06$  fm



$a \simeq 0.05$  fm



# Topological Susceptibility: Preliminary

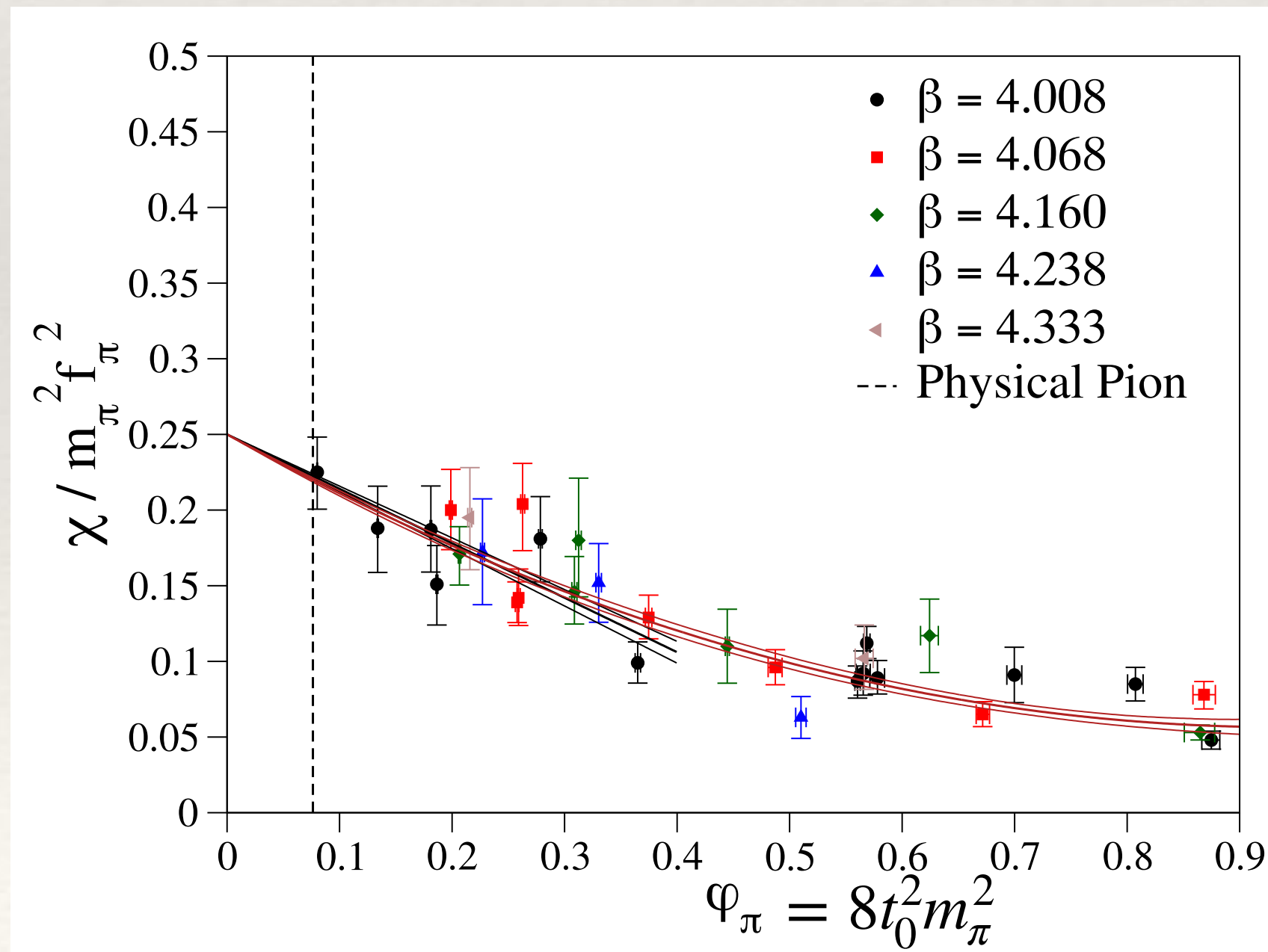
- Topological susceptibility is computed with the “slab” method

[Bietenholz, de Forcrand, Gerber 1509.06433]  $\chi(\tau) = C + \chi V \frac{\tau}{T} + \mathcal{O}(e^{-m_\eta L})$

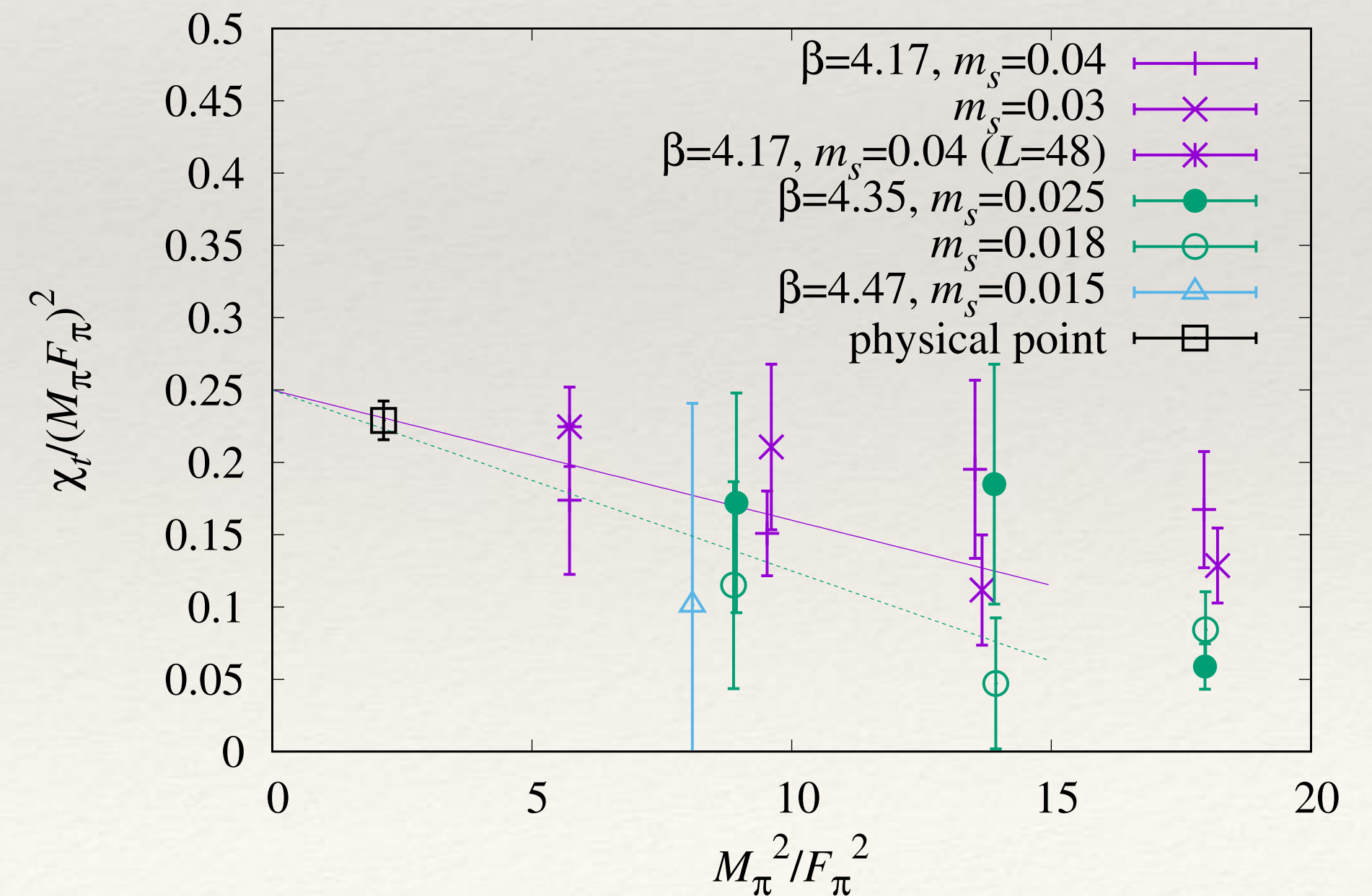
- 30 iterations of HYP smearing with  $\alpha_1 = 0.75$ ,  $\alpha_2 = 0.5$ ,  $\alpha_3 = 0.25$  are used

- $\chi = \frac{m_\pi^2 F_\pi^2}{4} \left( 1 + 2l' \frac{m_\pi^2}{F_\pi^2} + \mathcal{O}(m_\pi^4/F_\pi^4) \right)$

$$\chi = 0.223(4)_{\text{stat}}(4)_{\phi_\pi}$$

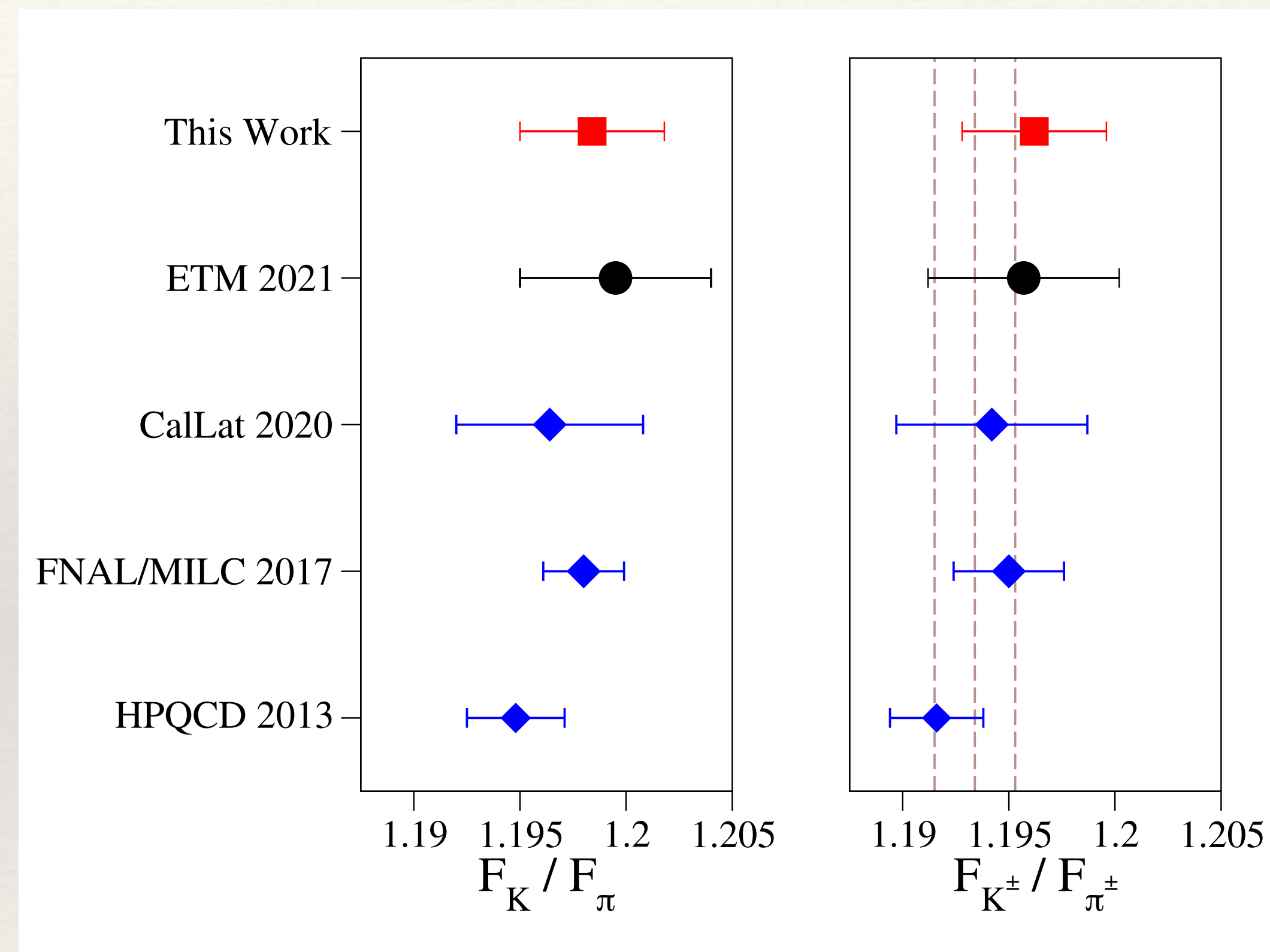
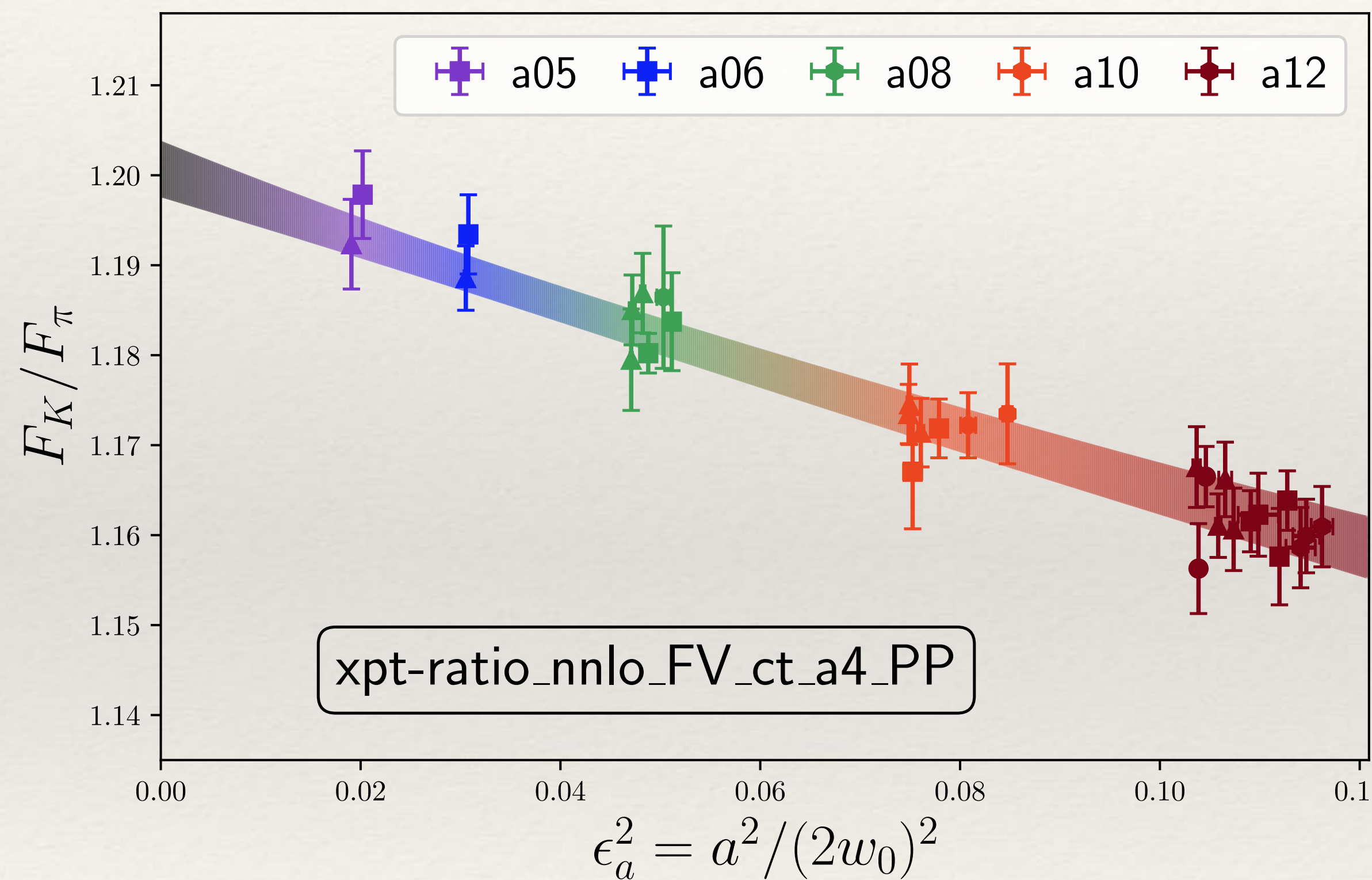


Aoki et al, JLQCD, PTEP 4 (2018) [1705.10906]



# $F_K/F_\pi$ as a precision test

□ We computed  $F_K/F_\pi$  as a high-precision test - [\[2605.06560\]](#) — see Zack Hall's talk after this one



# Charmed Quarks

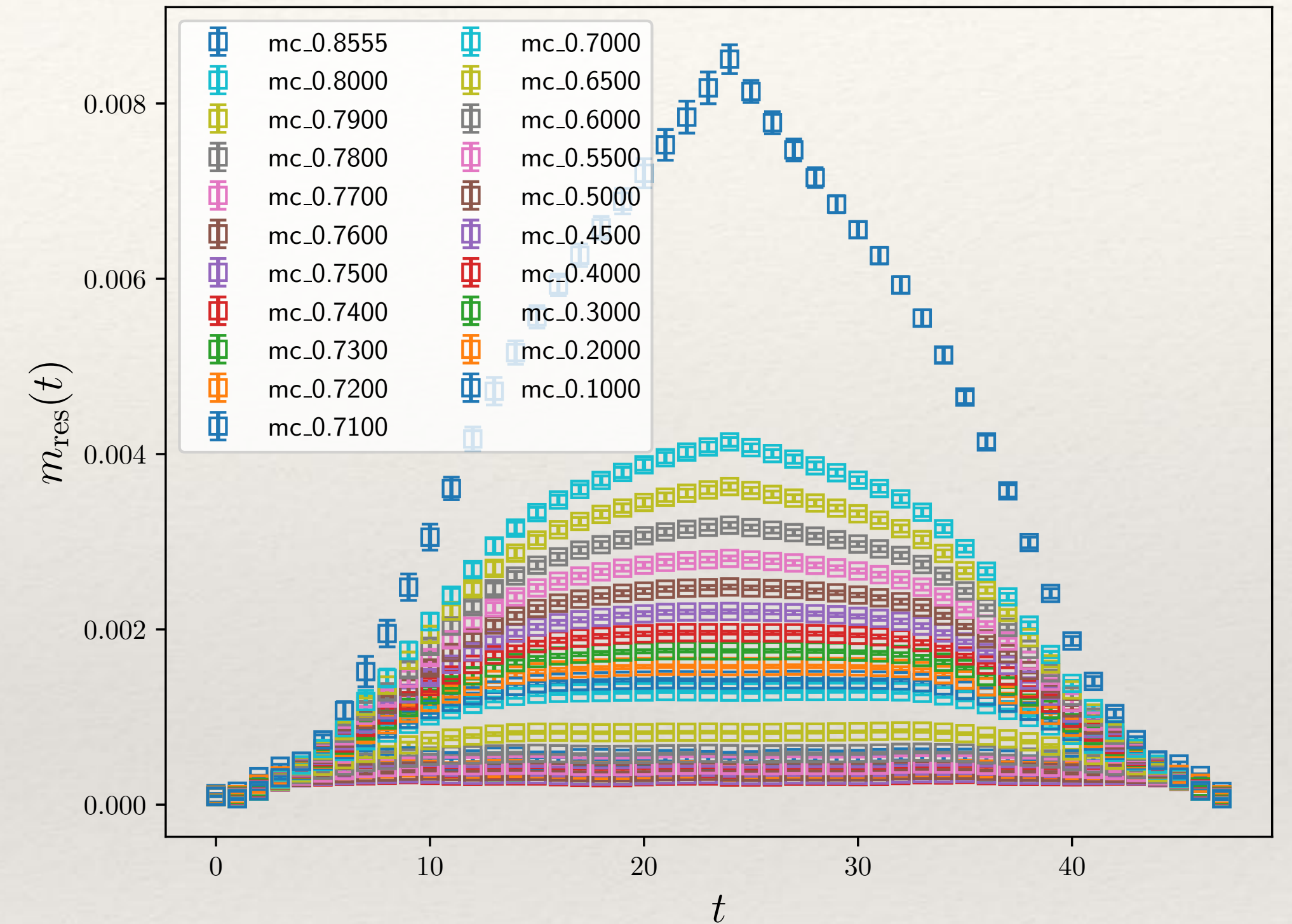
❑ On the coarsest ensemble, we observe some issues with the charmed fermion

❑  $m_{\text{res}}$  correlator does not plateau

❑  $\eta_{cc}$  correlators with  $A_4$  and  $P$  sinks  
do not exhibit same effective mass ground state

❑  $F_{D_s}/F_D$  not obviously wrong on coarsest lattice spacing

❑ We are investigating if we can extrapolate quantities such as  $m_{\text{res}}$  and the  $\eta_c$  mass using partially quenched valence charm mass from lighter up to sea value  $am_c = 0.8555$



# $\Omega_c, \Omega_{cc}$ : Preliminary

□ In contrast, the  $\Omega_c = (s^T C \gamma_i s) c$  and  $\Omega_{cc} = (c^T C \gamma_i c) s$  correlators show no issues

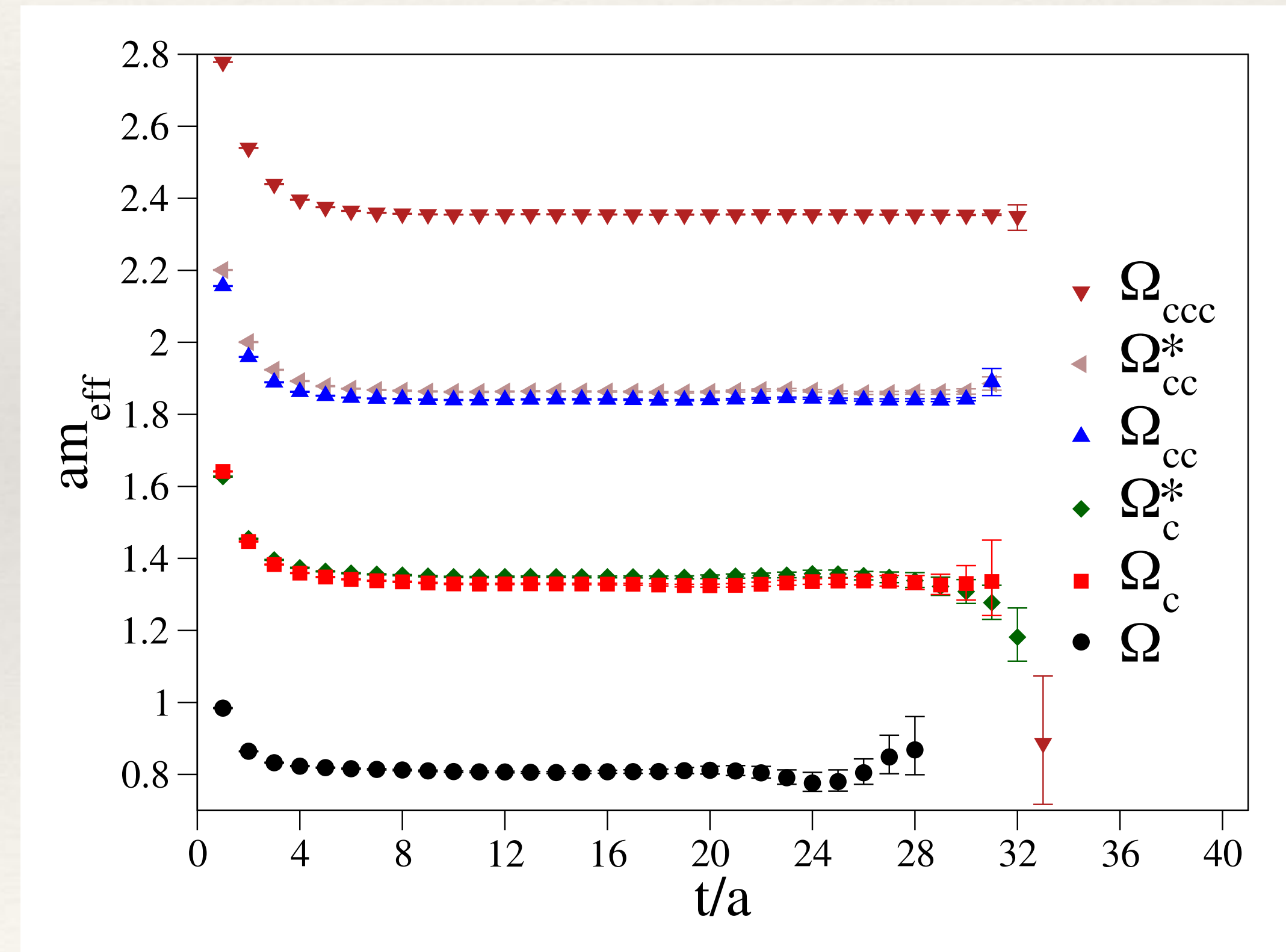
□ We explore the ratio of  $J = 3/2$  to  $J = 1/2$  correlators to fit the mass splitting

$$\frac{C^{3/2|+}(t)}{C^{1/2|+}(t)} \approx A e^{-\Delta t}$$

□ LHCb (Beauty 2026) have preliminary  $J = 1/2$   $\Omega_{cc}$  mass (3725.9(1.5) MeV) and  $J = 3/2$  unknown

□ The  $\Omega_c$  splitting is known from Belle 2009 and BaBar 2006:  $\Delta = 70.7(9)$  MeV

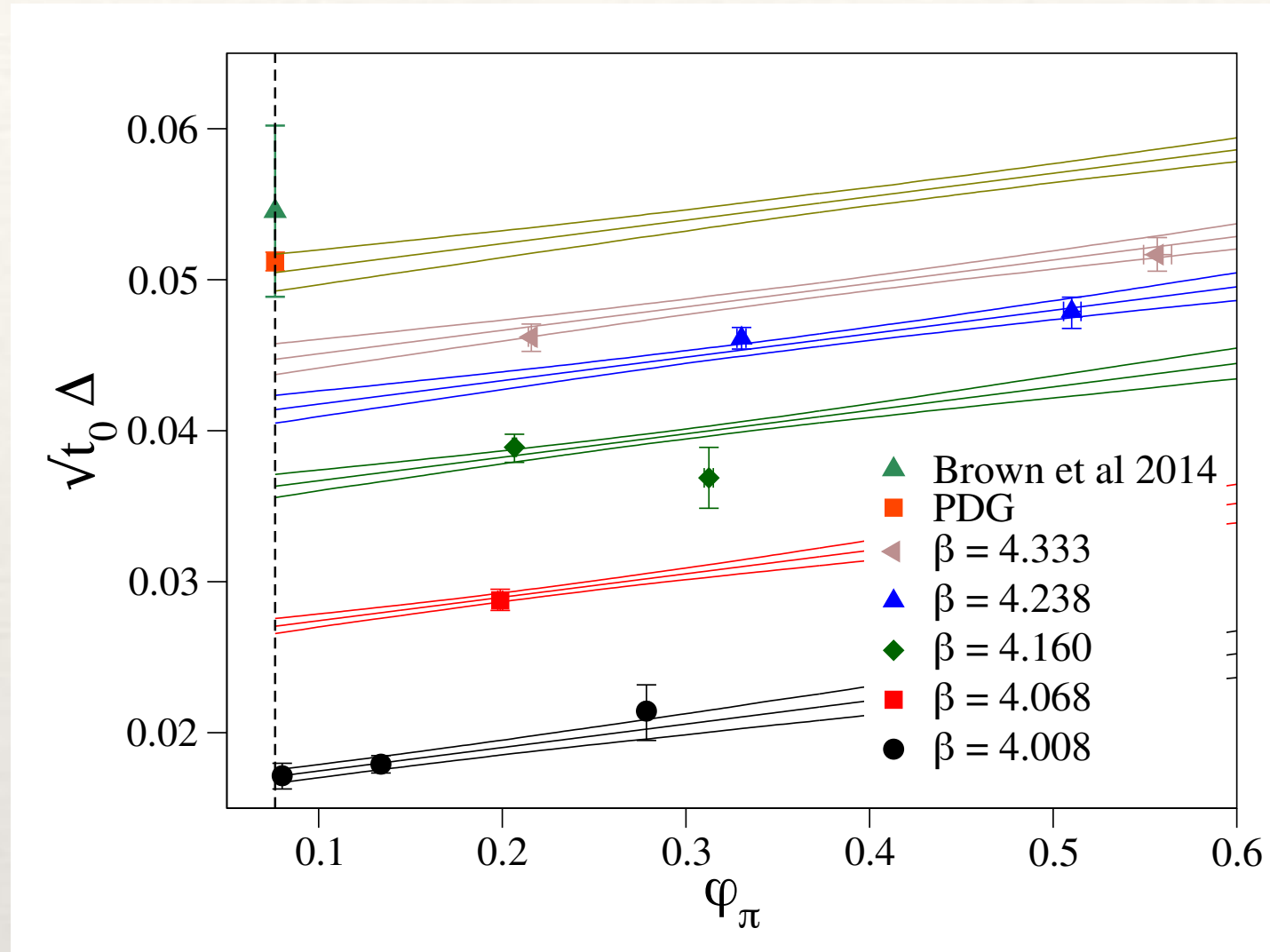
□ Right: sample effective mass plots from Coulomb gauge fixed wall sources with Gaussian profile for single to triply charmed  $\Omega$



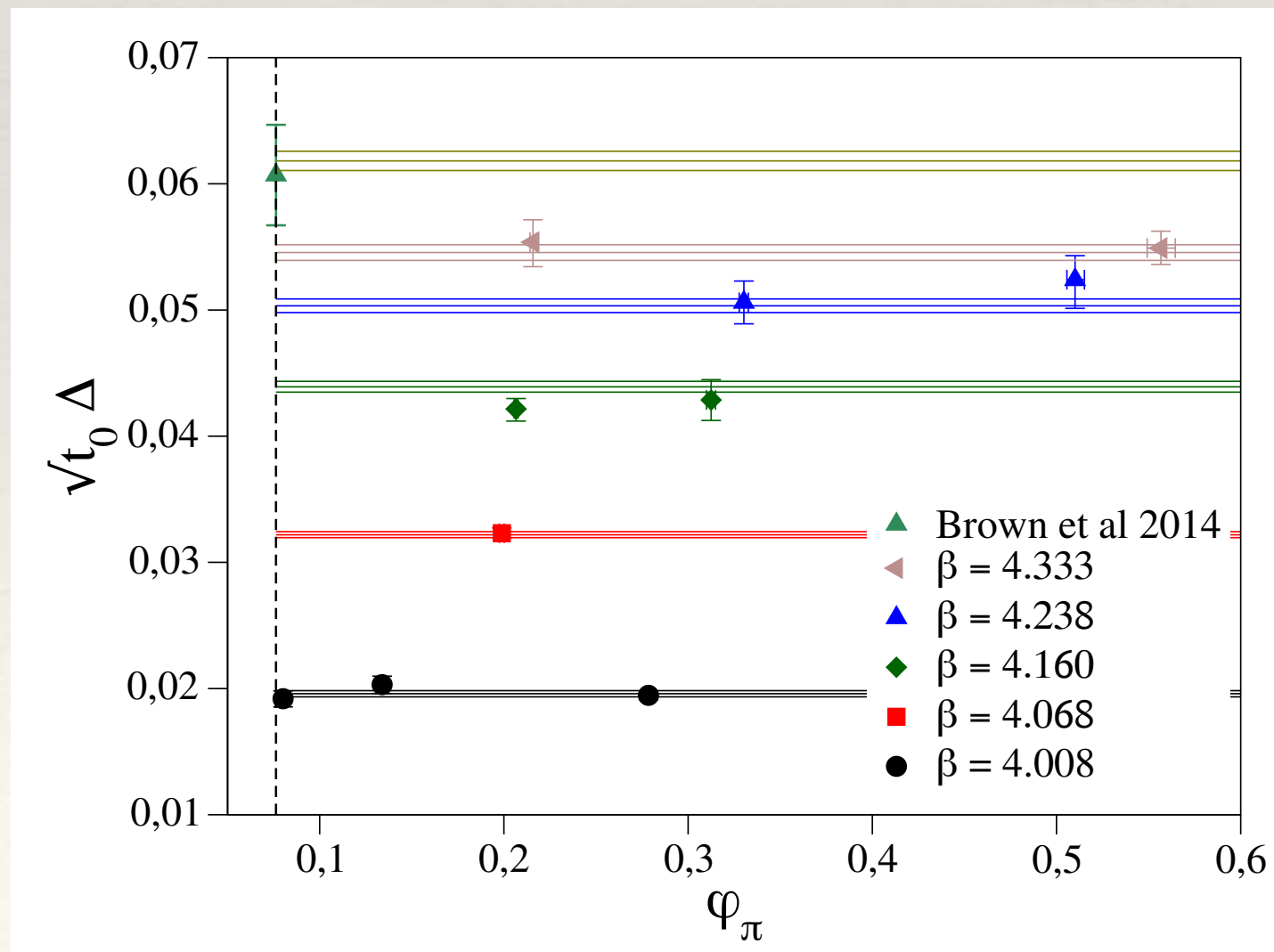
# $\Omega_c, \Omega_{cc}$ : Preliminary

□ Simple continuum + chiral fit to mass splittings

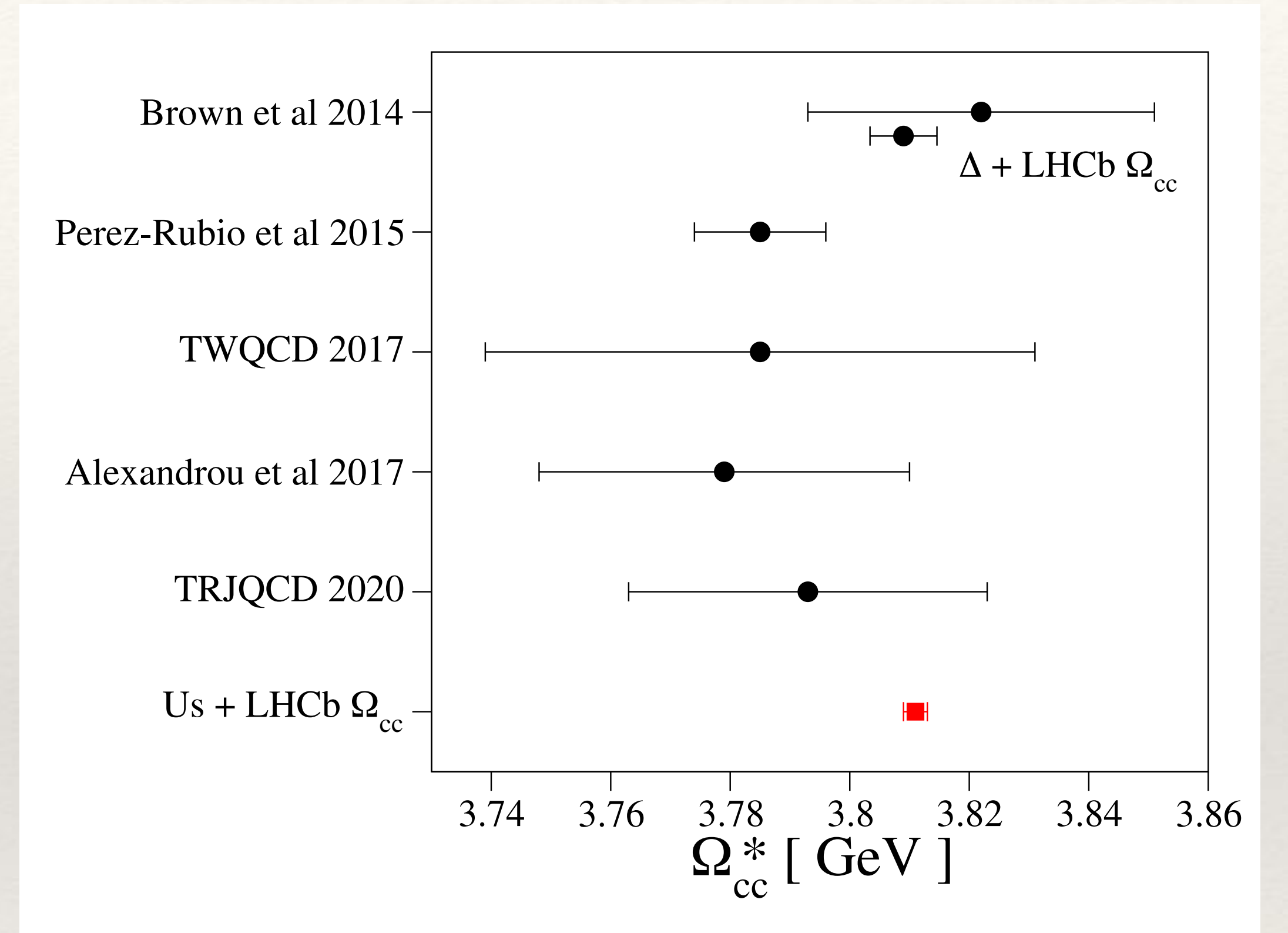
$$\sqrt{t_0} \Delta(\phi_\pi^2, a^2/t_0) = \sqrt{t_0^{\text{Cont}}} \Delta(\phi_\pi^{\text{Cont}}, 0) (1 + A(\phi_\pi - \phi_\pi^{\text{Cont}}) + Ba^2/t_0)$$



$\Delta_{\Omega_c}$



$\Delta_{\Omega_{cc}}$

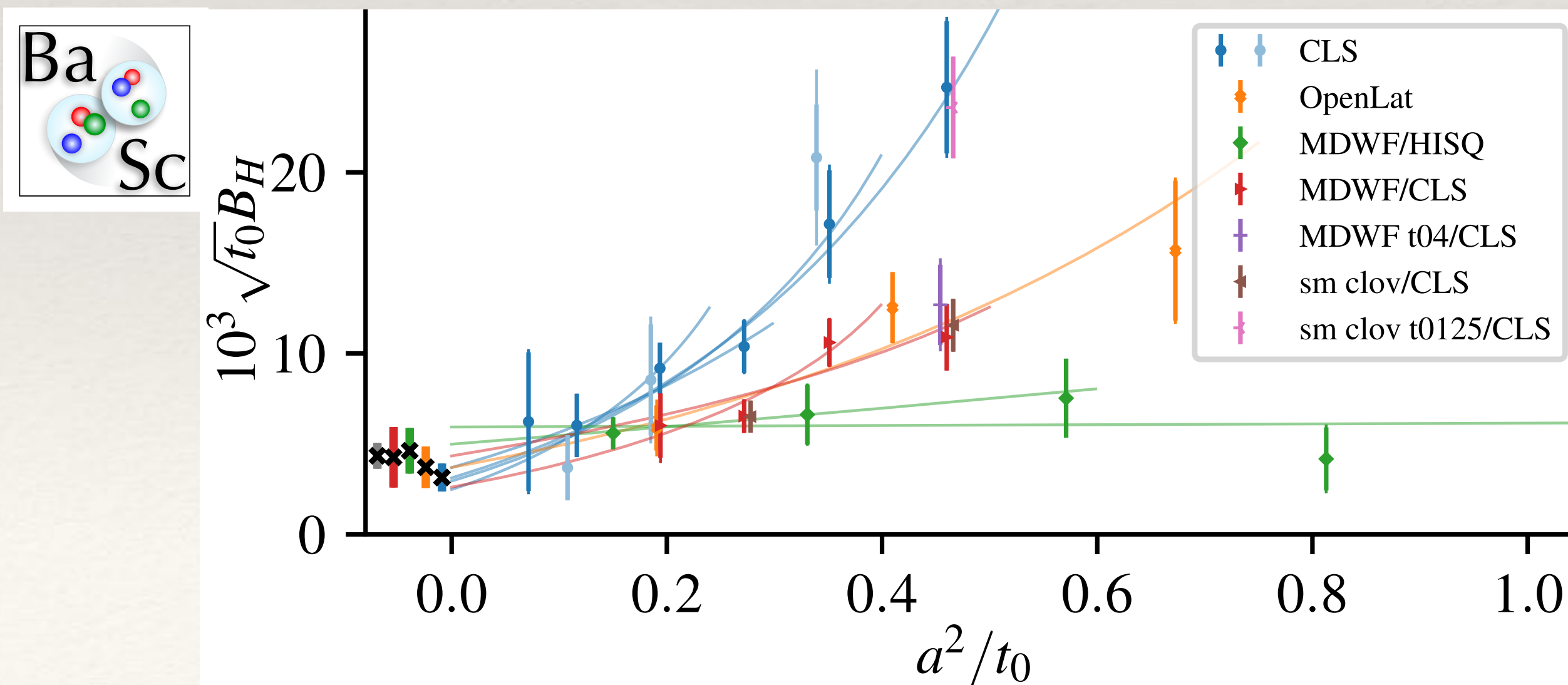


$$\Delta_{\Omega_{cc}} = 85.4(1.2) \text{ MeV}$$

$$M_{\Omega_{cc}^*} = 3811(2) \text{ MeV}$$

# Future studies of the action

- ❑ Spectroscopy study using  $\pi\pi$  systems is underway  
eg. “perambulators” computed by **T. Donovan/A. Hanlon**  
 $I = 0$   $\pi\pi$  correlators computed by **M. Salg**  
a12m375 ensemble
- ❑ Note:  $I = 0$  may require use of tetra-quark operators to get the correct spectrum — see  
**Hanlon et al, PRD 114 (2026) [2603.19192]**
- ❑ From our h-dibaryon continuum study with many actions we do not expect any surprises, but good to check



to be arXiv'd

figure: J. Green

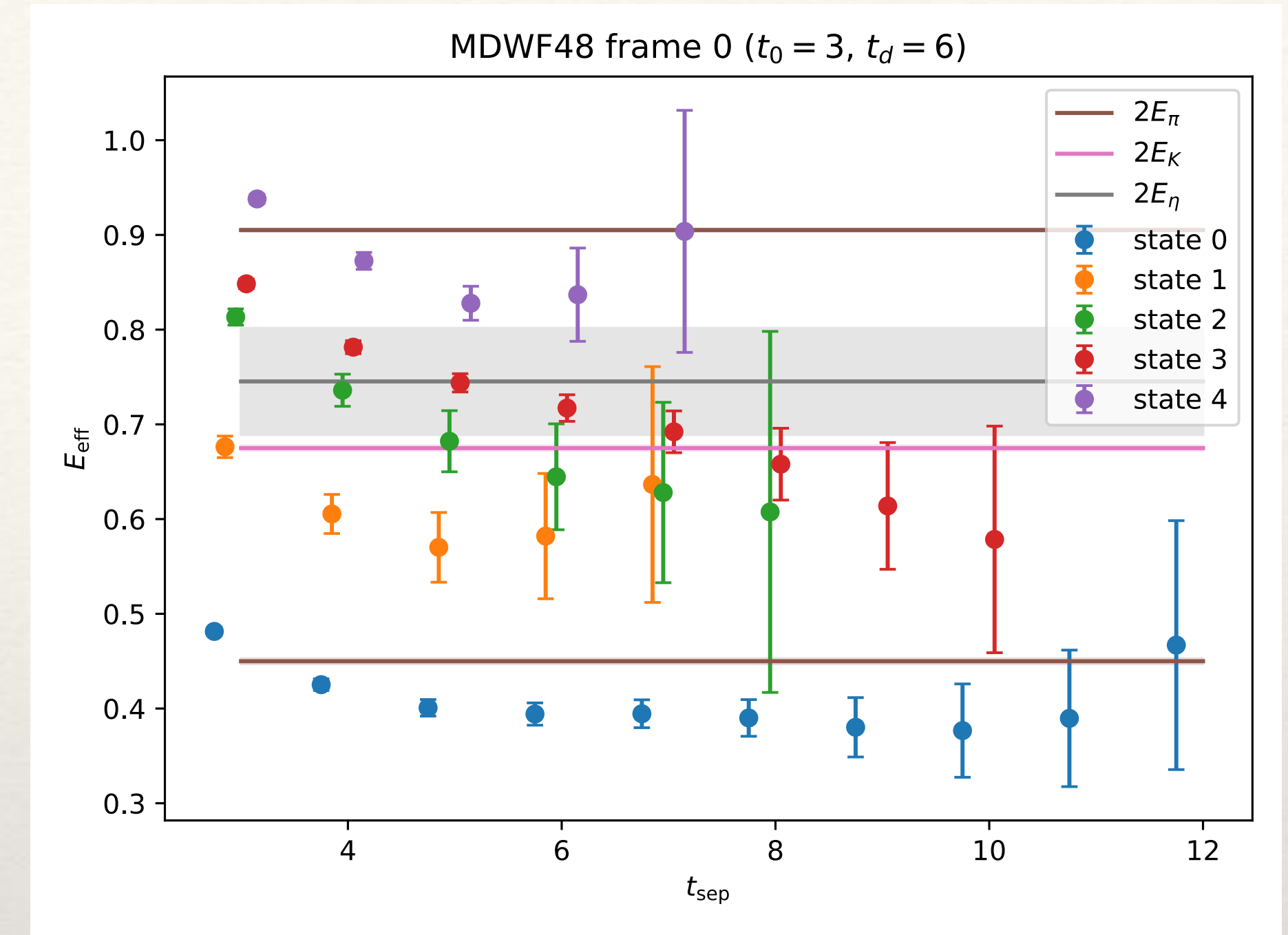


figure: M. Salg

# Future studies of the action

- ❑ An “action paper” is in preparation that will contain the details presented here as well as more
- ❑ High-precision calculation of  $F_{D_s}/F_D$  is underway
- ❑ We have started a pilot study of  $K \rightarrow \pi \ell \nu$  on a constant line of physics @  $m_\pi \approx 375$  MeV
- ❑ We have started generating two more physical pion mass ensembles and several larger volume ensembles at lighter pion mass
- ❑ We are generating SU(3)-flavor symmetric ensembles to carry out various studies
  - ❑  $t_0/t_1$
  - ❑ NPR
  - ❑ ...
- ❑ Assuming no “gotchas”, we hope to shift our full physics program to these ensembles

# Data availability

- ❑ All configurations from our first paper — 2605.06560 — are available via Globus from the Globus “[File Manager](#)”, search for the “[NERSC Gauge Connection](#)” (which is unfortunately down at the moment from the NERSC upgrade)

$F_K/F_\pi$  as a precision test of a new four flavor Domain Wall fermion action

Renwick J. Hudspith,<sup>1</sup> Nicolas Garron,<sup>2,3</sup> Zack Hall,<sup>4</sup> Andrew Hanlon,<sup>5</sup> Henry Monge-Camacho,<sup>6</sup> Colin Morningstar,<sup>1</sup> Amy Nicholson,<sup>7</sup> Dimitra A. Pefkou,<sup>8,4</sup> Thomas R. Richardson,<sup>8,4</sup> Fernando Romero-López,<sup>9</sup> Miguel Salg,<sup>9</sup> Wyatt A. Smith,<sup>10,4,8</sup> Pavlos Vranas,<sup>11,4</sup> André Walker-Loud,<sup>4,8</sup> and Bigeng Wang<sup>12,4</sup>

- ❑ We intend to keep dumping configs in this public domain as we use them in publications — feel encouraged to check them out!
- ❑ Configurations generated with a fork of [Grid](#) [Boyle, Yamaguchi, Cossu, Portelli, 1512.03487] <https://github.com/RJHudspith/Grid> that supports efficient force computations from stout smearing as well as different integrators on different levels
- ❑ Perambulators computed with [quda\\_laph](#) — [https://github.com/cosmon-collaboration/quda\\_laph](https://github.com/cosmon-collaboration/quda_laph)
- ❑ Pion/Kaon/Omega etc. 2pt functions computed with Wit [R.J. Hudspith, J. Green, S. Paul]

# Conclusions & Outlook

- ❑ We have generated a promising new 2+1+1 MDWF action
- ❑ I didn't emphasize, but, the 30 ensembles from our first paper were generated with a modest allocation on NERSC in  $\sim 1.2$  years — signaling a relatively inexpensive set of MDWF ensembles
- ❑ Our first high-precision scrutiny is available on the arXiv:  $F_K/F_\pi$  with more to come, eg.  $F_{D_s}/F_D$
- ❑ An action paper is being prepared
- ❑ The configurations are in the public domain
- ❑ If you have any questions, let us know!

*Thank You!*