

Spectral density reconstruction from lattice QCD correlators: an approach via numerical anti-Laplace transform.

Analytical characterisation.

José Alejandro Vidal Maxiá

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In collaboration with F. Di Renzo, P. Dimopoulos, D. Gavriel & M. Aliberti

LATTICE 2026 contributed talk

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UNIVERSITÀ
DI PARMA



Istituto Nazionale di Fisica Nucleare



Talk overview

- Numerical approach to the inverse problem
- Singular value decomposition (SVD) key aspects
- Relation to other approaches
- Conclusions/outlook

Numerical approach to the inverse problem

Euclidean correlators as Laplace transform of spectral densities
(Kallen-Lehmann representation)

$$C(na) = \int_0^\infty \rho(\omega) e^{-\omega n} d\omega = \omega_0 \int_0^\infty \rho(\omega_0 \xi) e^{-\omega_0 \xi n} d\xi$$

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$$n \in \{n_0, \dots, N_t + n_0 - 1\}, n_0 \geq 1$$

$$k \in \{0, \dots, N_\omega - 1\}$$

$$C(na) = \int_0^\infty \rho(\omega) e^{-\omega n} d\omega = \omega_0 \int_0^\infty \rho(\omega_0 \xi) e^{-\omega_0 \xi n} d\xi \approx \omega_0 \sum_{k=0}^{N_\omega-1} w_k e^{-\xi_k(\omega_0 n-1)} \rho(\omega_0 \xi_k).$$

Gauss-Laguerre (GL) quadrature rule applied for
integrals of the form $\int_0^\infty f(x) e^{-x} dx$

(all quantities understood to be in lattice units)

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Gauss-Laguerre (GL) quadrature rule applied for
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We can thus write

$$C_n = \sum_k Z_{nk} \rho_k$$

w/

$$C_n \equiv C(na)$$

$$Z_{nk} \equiv \omega_0 w_k e^{-\xi_k(\omega_0 n-1)}$$

$$\rho_k \equiv \rho(\omega_0 \xi_k)$$

Numerical approach to the inverse problem

This inverse problem needs regularisation
(in our case Tikhonov prescription)

Minimise $(\mathbf{c} - \hat{\mathbf{Z}}\boldsymbol{\rho})^T \hat{\Sigma}^{-1}(\mathbf{c} - \hat{\mathbf{Z}}\boldsymbol{\rho}) + \alpha \|\boldsymbol{\rho}\|_2^2.$

$$\mathbf{c} \in \mathbb{R}^{N_t}, \boldsymbol{\rho} \in \mathbb{R}^{N_\omega}$$

$\hat{\Sigma}$ is the $N_t \times N_t$ correlation matrix
for the observable

for now we
set it to $\mathbb{I}_{N_t}$

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Tikhonov-regularised least-squares solution $\boldsymbol{\rho}^{(\alpha)} = (\hat{\mathbf{Z}}^T \hat{\Sigma}^{-1} \hat{\mathbf{Z}} + \alpha \mathbb{I}_{N_\omega})^{-1} \hat{\mathbf{Z}}^T \hat{\Sigma}^{-1} \mathbf{C}.$

for now we
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Regularisation induces smearing



regularised delta peaks for
discrete spectrum

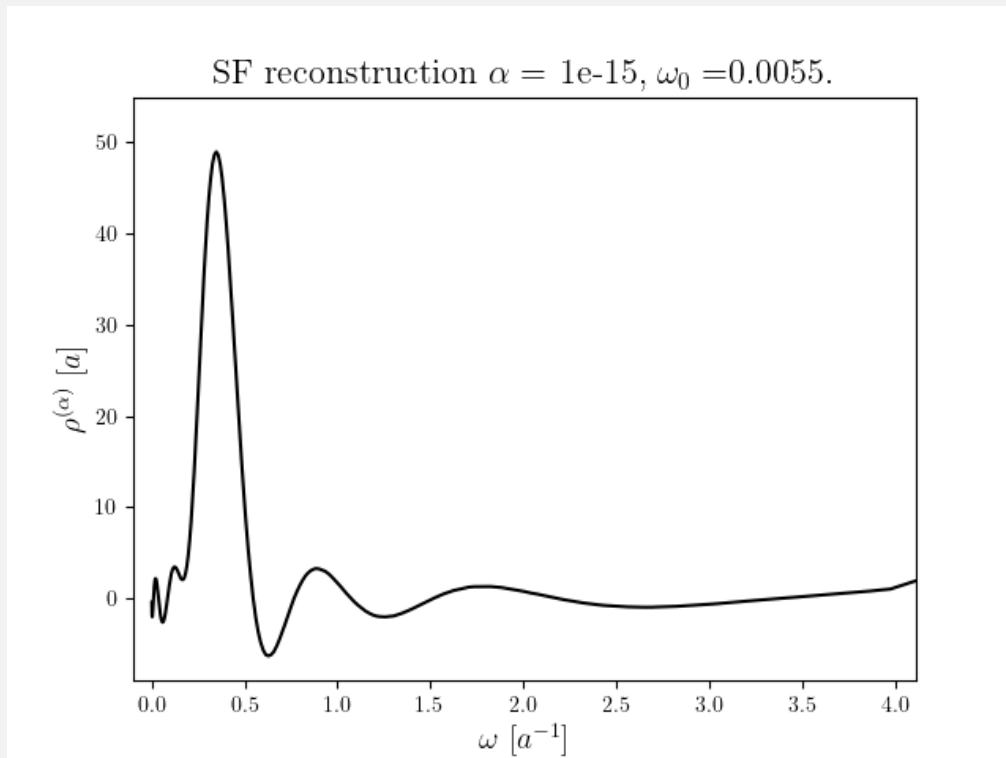
Numerical approach to the inverse problem

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Synthetic noiseless correlator $C_n = e^{-0.2n} + 2e^{-0.29n} + 3.5e^{-0.35n} + 3e^{-0.41n}$

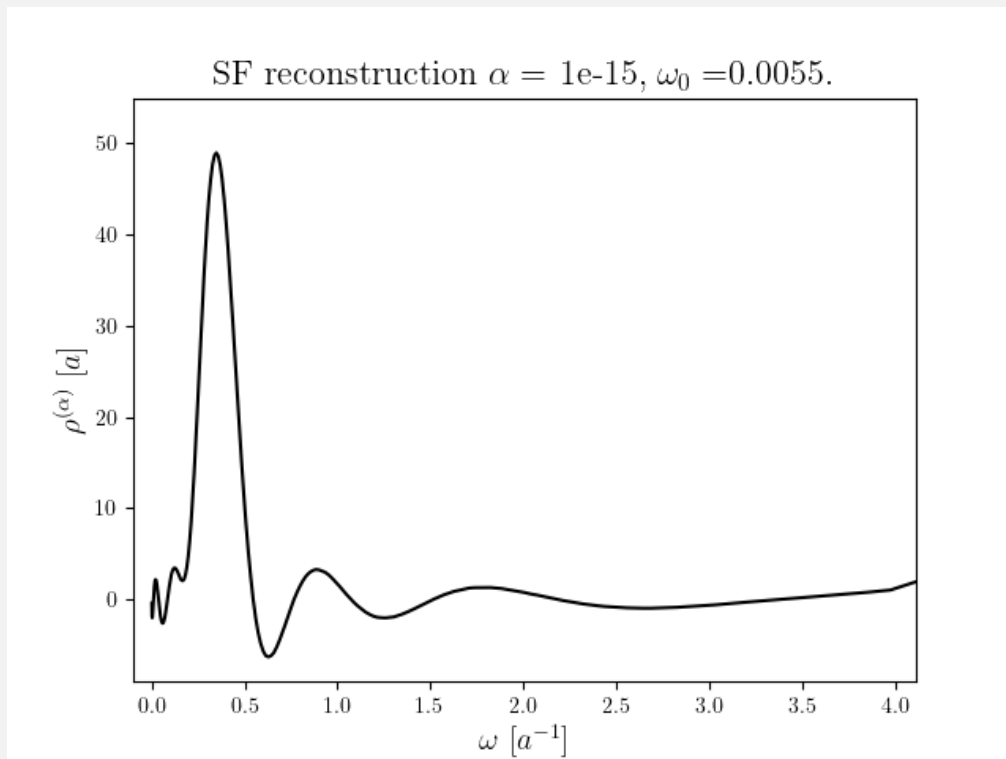
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$$\begin{aligned} N_\omega &= 195 \\ N_t &= 30 \\ n_0 &= 1 \end{aligned}$$

Many peaks clumped up

Structure **appears** 'hidden'

SVD key aspects

Singular Value Decomposition (SVD) is a central component of Laplace (**and others**) numerical inversion

[R.Tsuji, S.Hashimoto. arXiv:2605.15674v1](#)

[M. Bruno, L. Giusti, M. Saccardi, arXiv:2407.04141v1](#)

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Continuous time:

$$t \in [0, \infty), \quad \omega \in [0, \infty)$$

SVD of Laplace kernel

$$e^{-\omega t} = \int_{-\infty}^{+\infty} v_s(\omega) \lambda_s^* u_s^*(t) ds, \quad \text{w/ } |\lambda_s| = \sqrt{\frac{\pi}{\cosh(\pi s)}}$$

the expression is actually real, it can be expressed in a different choice of basis similar to Fourier transform

$$v_s(\omega) \text{ eigenfunctions of the Carlemann operator } \mathcal{H}(\omega, \omega') \equiv \int_0^\infty e^{-(\omega+\omega')t} dt = \frac{1}{\omega+\omega'}$$

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Simple relation among both (Mellin basis)

$$v_s(\omega) = u_s^*(\omega) = u_{-s}(\omega) = \frac{\omega^{-is}}{\sqrt{2\pi\omega}}$$

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equal (real) s. values but
spectrum is halved

$\bar{v}_s(\omega)$ eigenfunctions of the Carlemann operator $\bar{\mathcal{H}}(\omega, \omega') \equiv \sum_n e^{-n(\omega+\omega')} = \frac{e^{-(\omega+\omega')n_0}}{1-e^{-(\omega+\omega')}}$

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No longer simple relation between both bases

[R.Tsuji, S.Hashimoto. arXiv:2605.15674v1](#)

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$\bar{u}_s(n)$ and $\bar{v}_s(\omega)$ given by different
Hypergeometric functions ${}_qF_p$

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$$Z_{nk} = \omega_0 w_k e^{-\xi_k(\omega_0 n - 1)} = \sum_s U_s(n) \lambda_s V_s(\xi_k)$$

Recall our initial inverse problem

$$C_n = \sum_k Z_{nk} \rho_k$$

$V_s(\xi_k)$ eigenbasis of the ‘Carleman’ operator $\tilde{\mathcal{H}}(\xi_k, \xi_{k'}) \equiv (\hat{Z}^T \hat{Z})_{kk'} = \sum_s V_s(\xi_k) |\lambda_s|^2 V_s(\xi_{k'})$

Noise modifies singular values and basis vectors but the treatment is the same overall

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General solution in terms of SVD

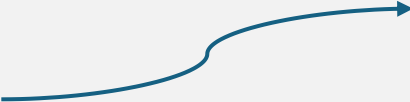
$$\rho_k = \sum_s V_s(\xi_k) \frac{\lambda_s}{\lambda_s^2 + \alpha} \sum_n U_s(n) C_n$$

$U_s(n)$ and $V_s(\xi_k)$ form two orthonormal bases of $\mathbb{R}^{N_t}$ and $\mathbb{R}^{N_\omega}$, resp.

SVD key aspects

Characterisation of lattice correlator signals

In a lattice-regulated Euclidean theory, we get sums of exponentials


$$C_n = \sum_q A_q e^{-nE_q} = \sum_q A_q e^{-n\omega_0 \bar{E}_q}$$

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Our anti-transformed signal will thus be of the form

$$\rho_k = \sum_q A_q \underbrace{\sum_s V_s(\xi_k) \frac{\lambda_s}{\lambda_s^2 + \alpha} \sum_n U_s(n) e^{-n\omega_0 \bar{E}_q}}_{\delta_k^{(\alpha)}(\bar{E}_q) \equiv \delta_{qk}^{(\alpha)}}$$

Sums of regularised delta peaks $\delta_k^{(\alpha)}(\bar{E}_q)$

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Take $\bar{E}_q = \xi_l$ from within the GL
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effective cutoff in eigenvalue sum

Under $\sum_k \omega_0 w_k e^{\xi_k} \rightarrow \int_0^\infty d\omega$

$$\delta_{kl}^{(\alpha)} = \frac{e^{-\xi_l}}{\omega_0 w_l} \sum_s V_s(\xi_k) \frac{\lambda_s^2}{\lambda_s^2 + \alpha} V_s(\xi_l) \xrightarrow{\alpha \rightarrow 0} \frac{e^{-\xi_l}}{\omega_0 w_l} \delta_{kl} \xrightarrow{N_\omega \rightarrow \infty} \delta(\xi_k - \xi_l)$$

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$\delta_{kl}^{(\alpha)}$ continuous in $\bar{E}_q$ (sum of exponentials) $\longrightarrow$ $\delta_{kq}^{(\alpha)}$ continuously extrapolates δ behaviour

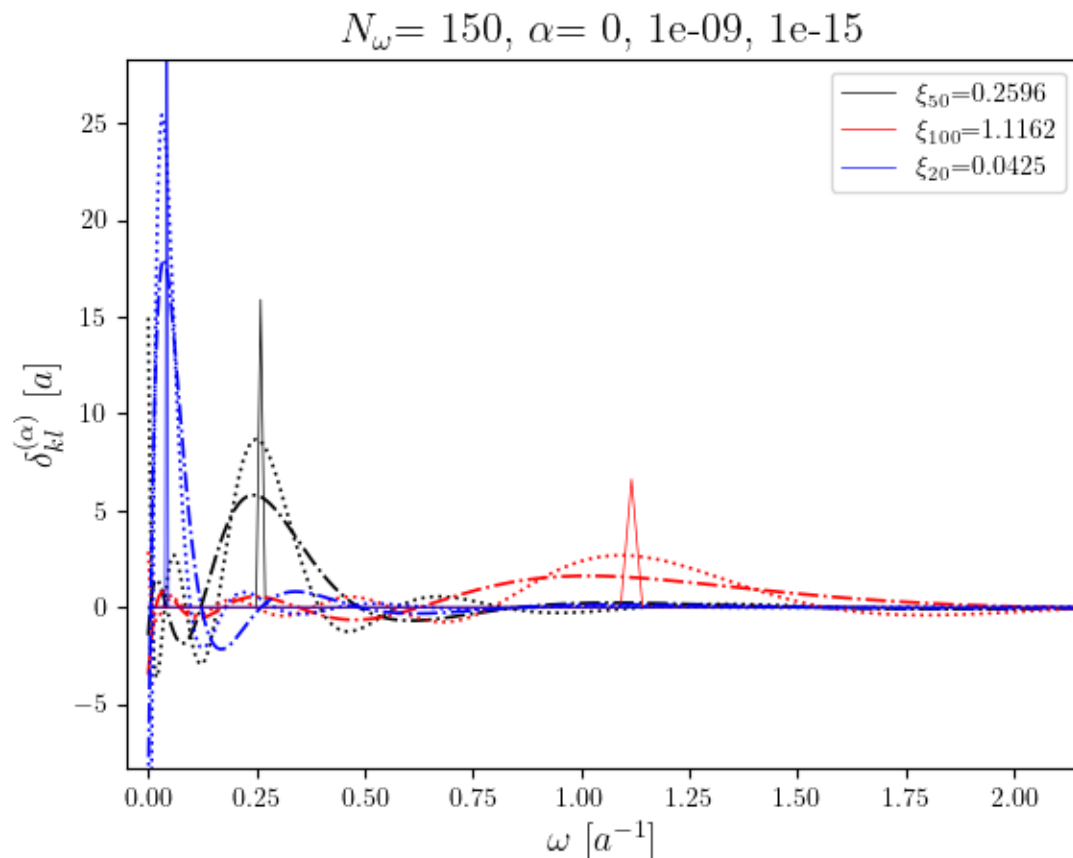
SVD key aspects

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Sums of regularised delta peaks.. **graphically**

$\delta^{(\alpha)}$ at quadrature points

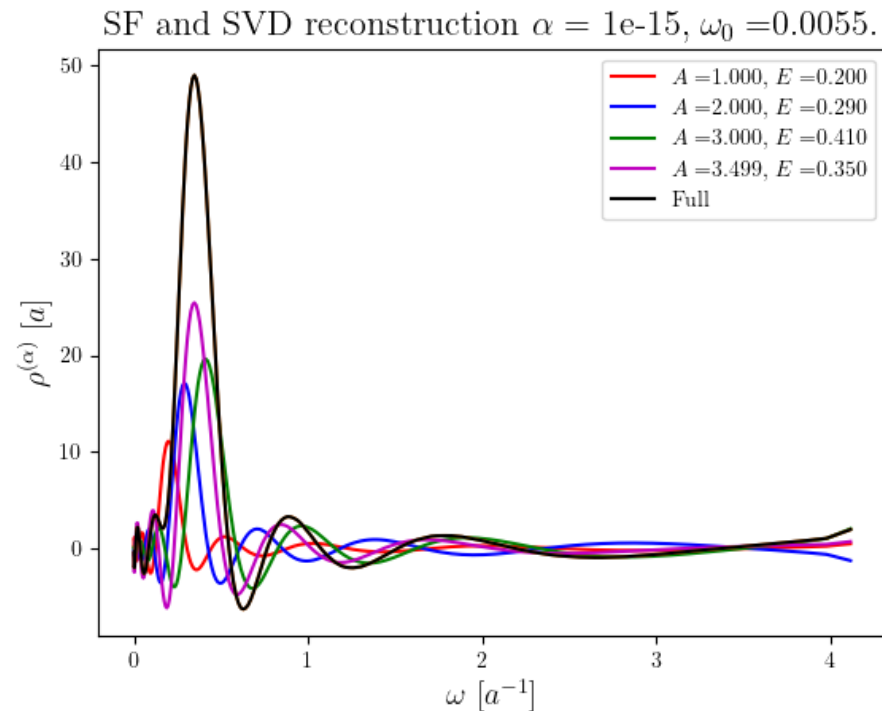
- **Continuous** lines are $\delta_{kl}^{(0)} = \frac{e^{-\xi_l}}{\omega_0 w_l} \delta_{kl}$ at different ξ_l
- GL quadrature points **not** evenly spaced
- Smaller α leads to more pronounced peaks



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Sums of regularised delta peaks.. **graphically**



$\delta^{(\alpha)}$ at arbitrary points

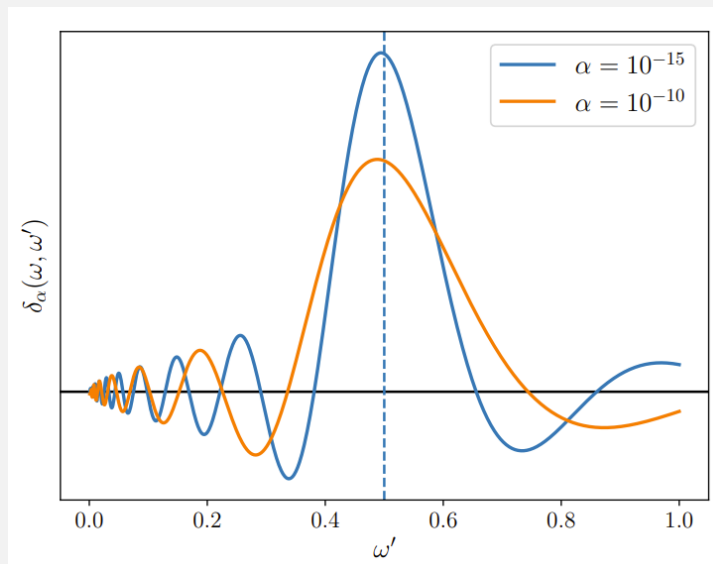
- Actual decomposition of previous reconstructed density
- Behaviour of δ peaks indeed like at quadrature points
- Subtle cancellations lead to overall shape

Relation to other approaches

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This kind of profile (**rather families**) already encountered in other occasions

Mellin transforms

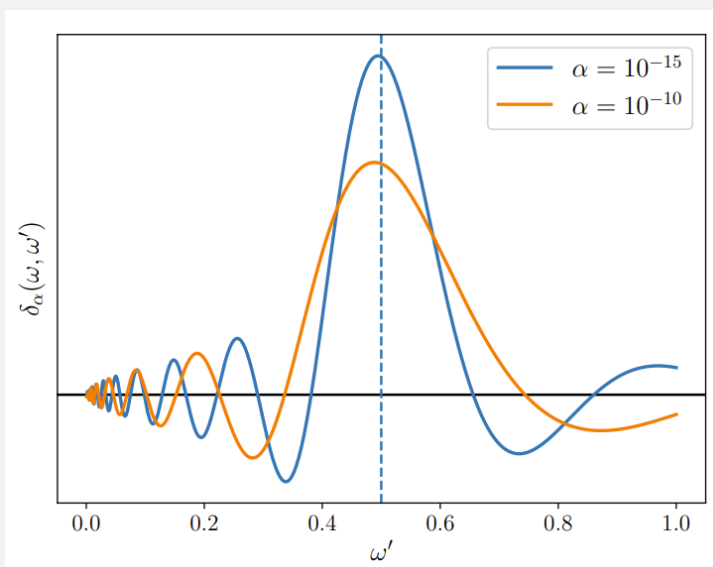


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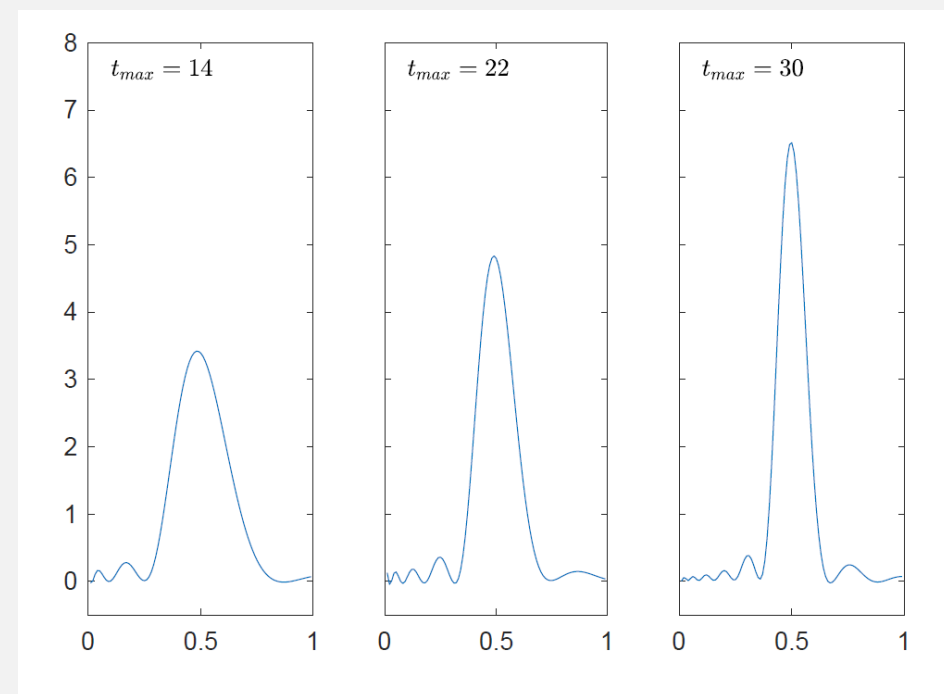
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Backus-Gilbert method



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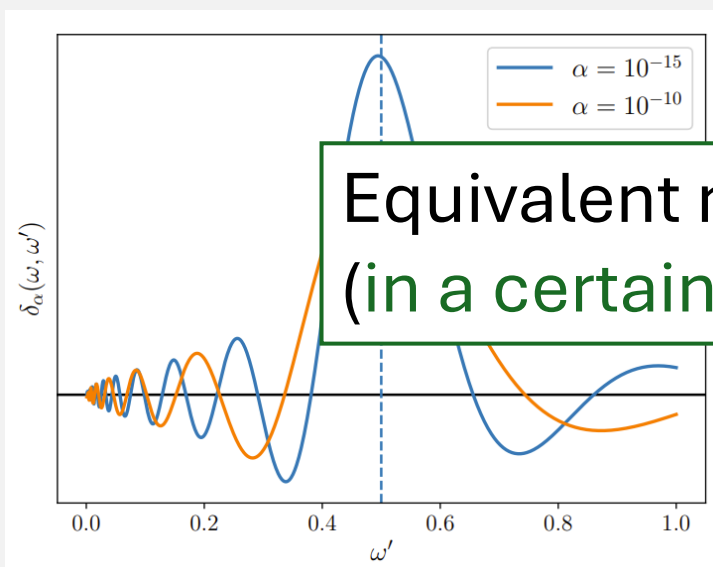
G. Backus and F. Gilbert, *The Resolving Power of Gross Earth Data*, Geophys. J. Int. **16** (1968), no. 2 169–205.

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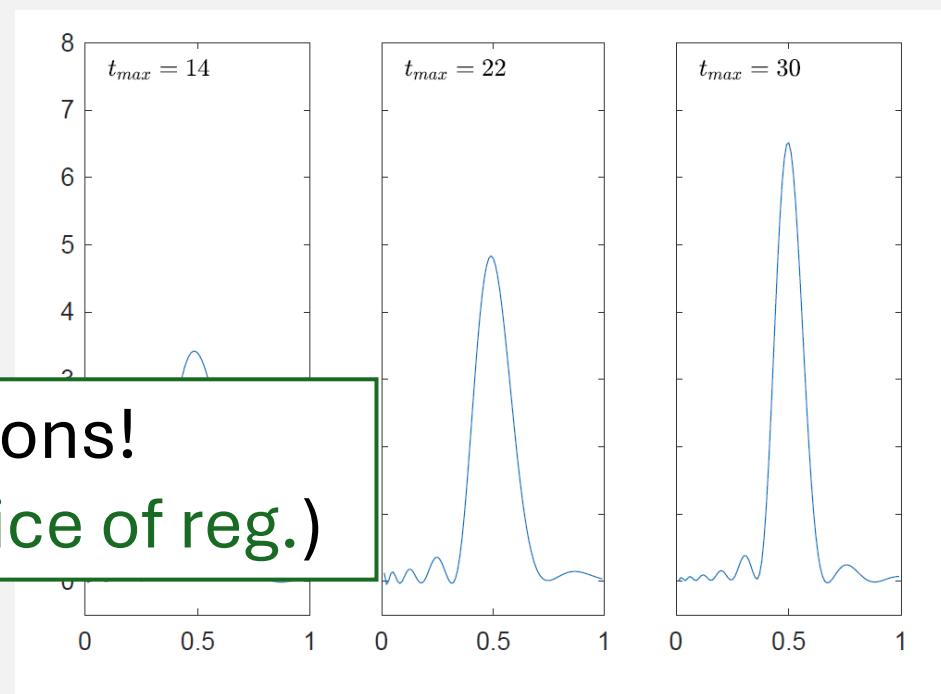
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Equivalent minimizations!
(in a certain limit/choice of reg.)

Backus-Gilbert method



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Conclusions

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- Regulated delta peaks can be easily characterized
- SVD powerful tool $\longrightarrow$ equivalent to our inverse problem
- Relation extends to other linear reconstruction approaches

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Outlook

- Errors need to be accounted for next
- Adequate minimization procedure underway for spectral reconstruction (energies, matrix elements)
- Interesting to consider generalizing to linear programming

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Goodbye

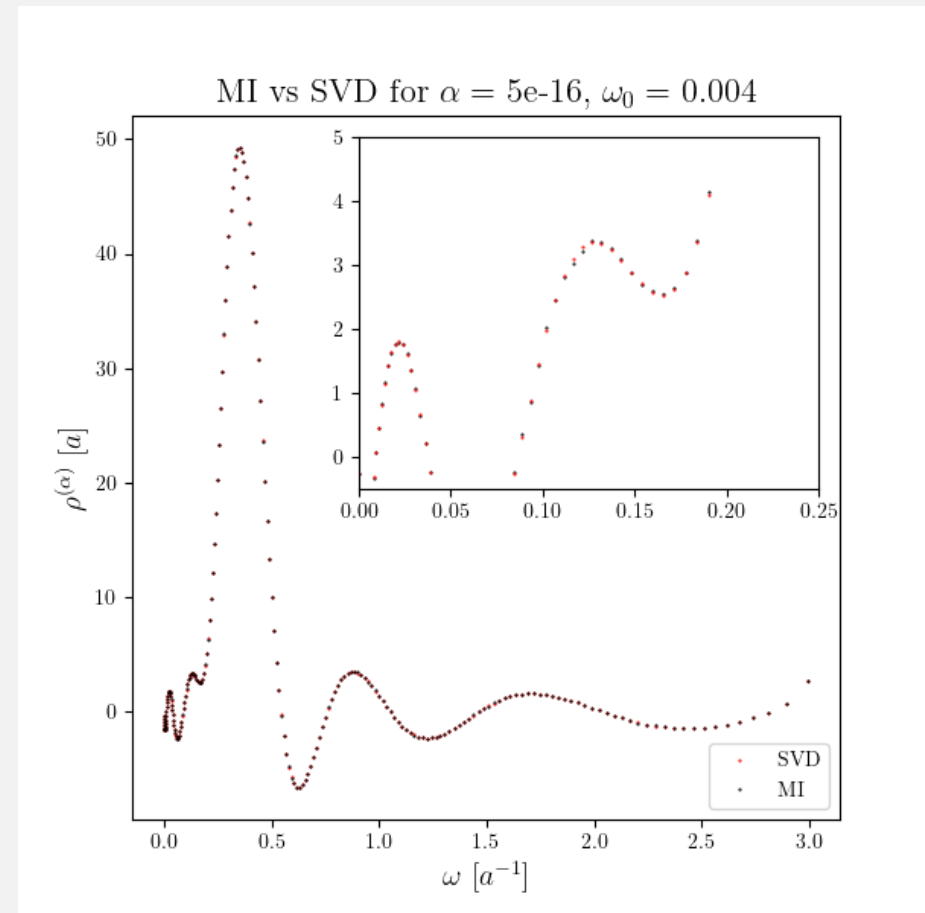
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**THANK YOU FOR THE
ATTENTION!**

Extra: equivalence MI vs SVD

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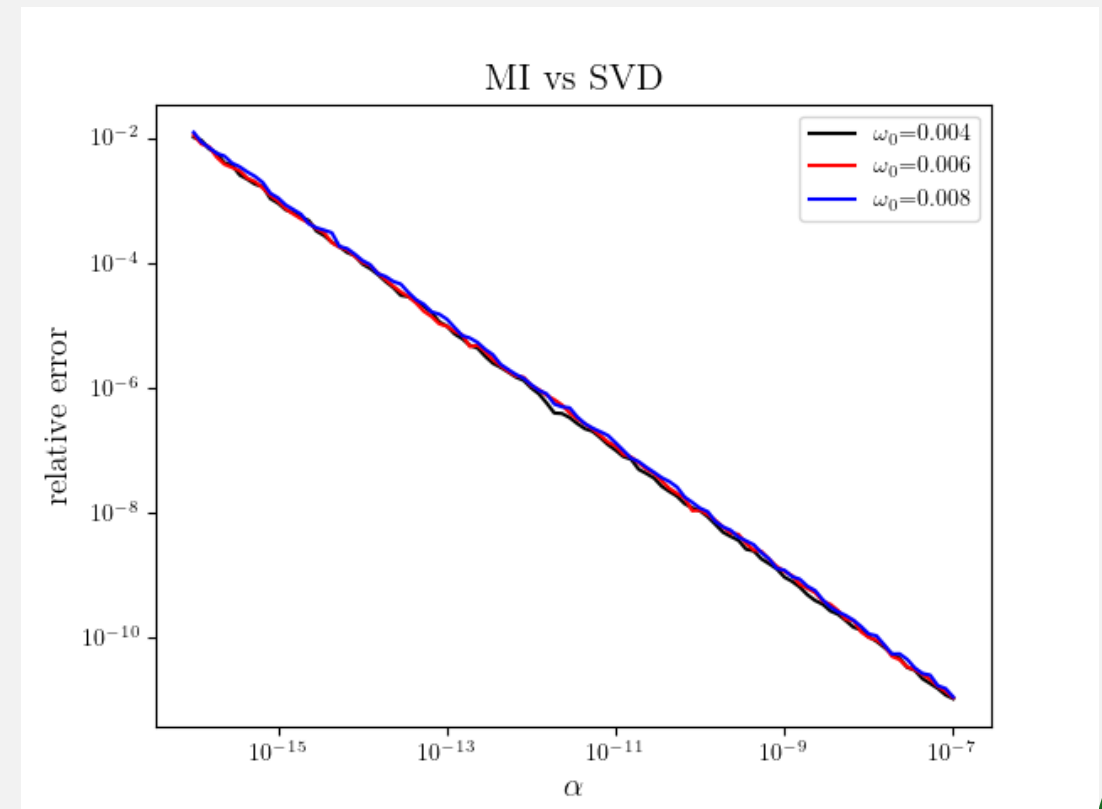
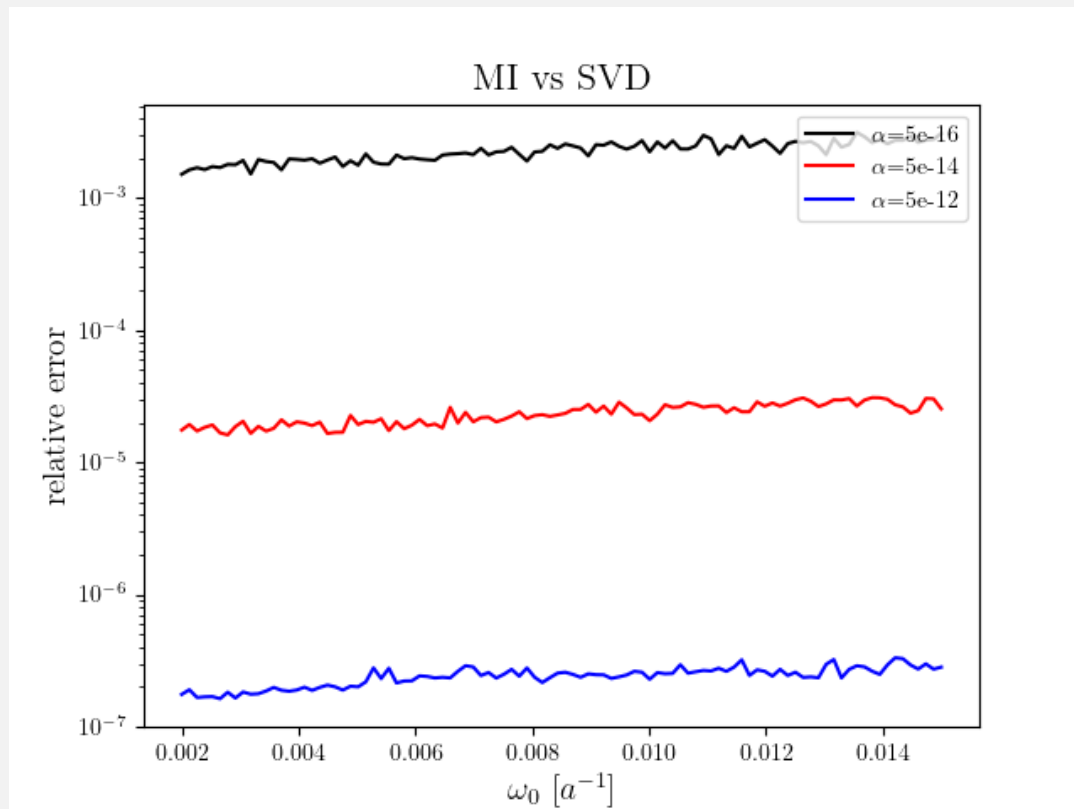
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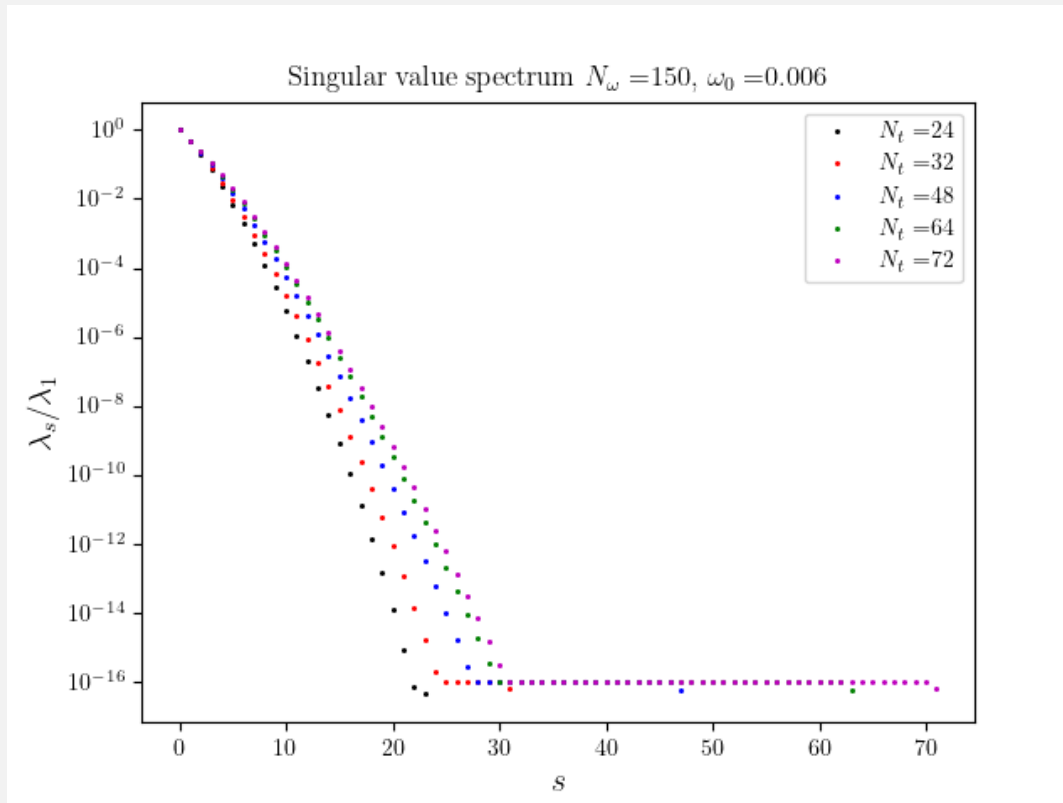
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Extra: SVD spectrum decay

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Maximum s available determines spectral
function resolution



It does depend on N_t ...

but not dramatically so