

Nucleon Electromagnetic Form Factors at the Physical Pion Mass from PNDME Collaboration

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Acknowledgements

PNDME and NME Collaborations

T. Bhattacharya, R. Gupta, E. Mereghetti (LANL)

V. Cirigliano (INT), B. Yoon, B. Joo (NVIDIA)

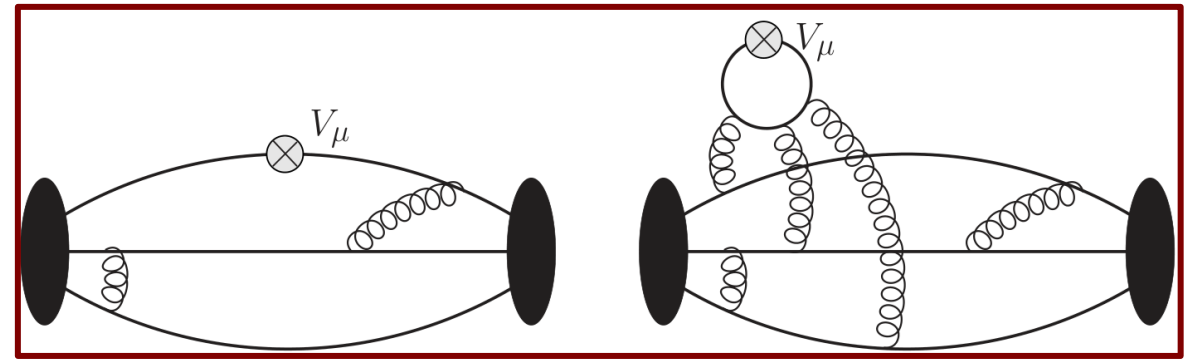
J. Yoo (Korea U.), H.-W. Lin (MSU), F. Winter (JLab)

- **Acknowledge the MILC collaboration for HISQ ensembles**
- The calculations used the CHROMA software suite.
- We thank DOE for computer time allocations at NERSC and OLCF.
- We thank the USQCD collaboration for computer time
- Institutional Computing at LANL for computer time

Electromagnetic Form Factors, G_E and G_M

- Dirac/Pauli Form factors, $F_{1,2}$ parametrize the matrix element of the electromagnetic current $V_\mu^{\text{em}} = \frac{2}{3}\bar{u}\gamma_\mu u - \frac{1}{3}\bar{d}\gamma_\mu d$ within nucleon ground states
$$\langle N(\mathbf{p}', s') | V_\mu^{\text{em}} | N(\mathbf{p}, s) \rangle = \bar{u}(\mathbf{p}', s') \left(F_1(q^2) \gamma_\mu + \frac{i\sigma_{\mu\nu} q^\nu}{2M_N} F_2(q^2) \right) u(\mathbf{p}, s)$$
- Sachs form factors, $G_E = F_1 - \frac{Q^2}{4M_N^2} F_2$, $G_M = F_1 + F_2$
- Isovector form factors $G_{E,M}^{u-d} = G_{E,M}^p - G_{E,M}^n$ assuming flavor SU(2) isospin symmetry,
$$\langle p | \bar{u}\gamma_\mu u - \bar{d}\gamma_\mu d | p \rangle = \langle p | V_\mu^{\text{em}} | p \rangle - \langle n | V_\mu^{\text{em}} | n \rangle$$

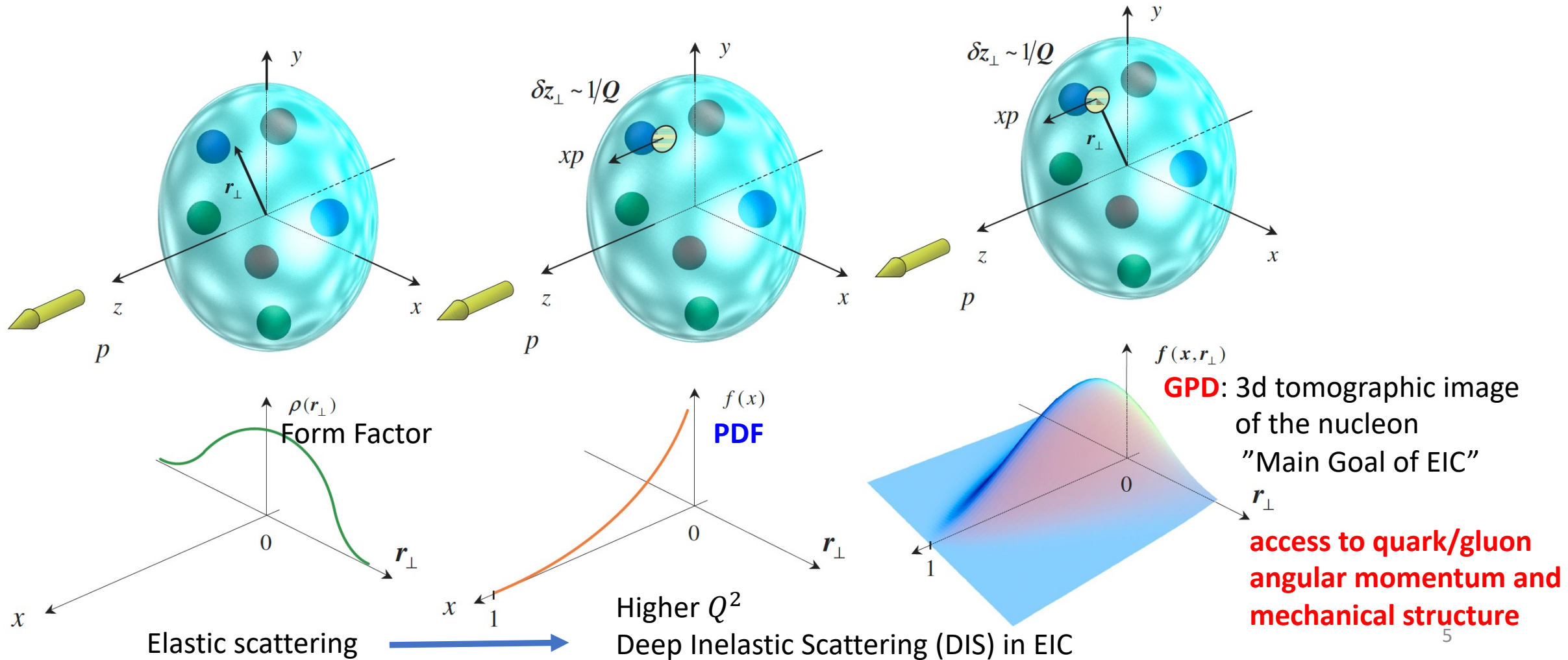
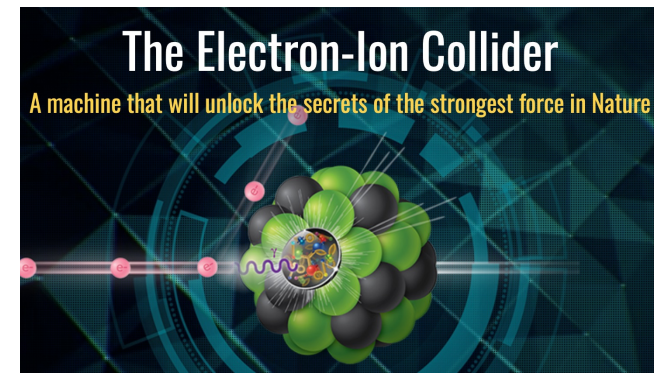
Isovector Charge Radii



- Proton electromagnetic form factor $G_{E,M}^p$ has contribution from both connected and disconnected diagrams
- We focus on the isovector form factor $G_{E,M}^{u-d} = G_{E,M}^{p-n}$ where the **disconnected contribution cancels** in the isospin symmetric limit.
- Charge radii squared

$$\langle r_E^2 \rangle^{u-d} = -6 \frac{d}{dQ^2} \left(\frac{G_{E,M}^{u-d}(Q^2)}{G_{E,M}^{u-d}(0)} \right) \Big|_{Q^2=0}$$

Beyond the nucleon form factor: Parton Distribution Function (PDF) Generalized Parton Distribution (GPD)



Three 2+1+1-flavor physical pion mass HISQ Ensembles

Ensemble ID	a [fm]	M_π [MeV]	$M_\pi L$	$L^3 T$	N_{conf}
a09m135	0.0873	135	3.95	$64^3 96$	1290 -> 6060 / 8000
a06m135	0.0564	135	3.95	$96^3 192$	675 -> 8269 / 10000
a04m135	0.0426	135	3.95	$144^3 288$	0 / 2000

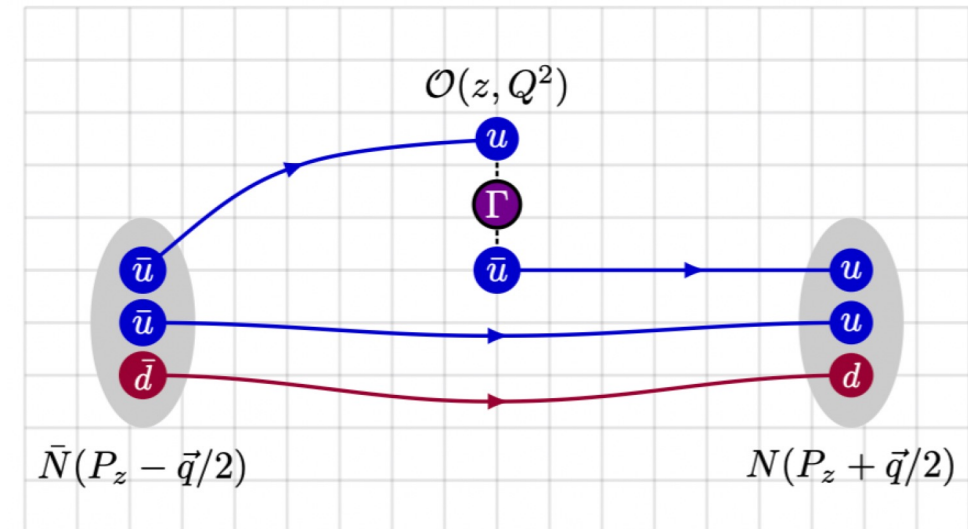
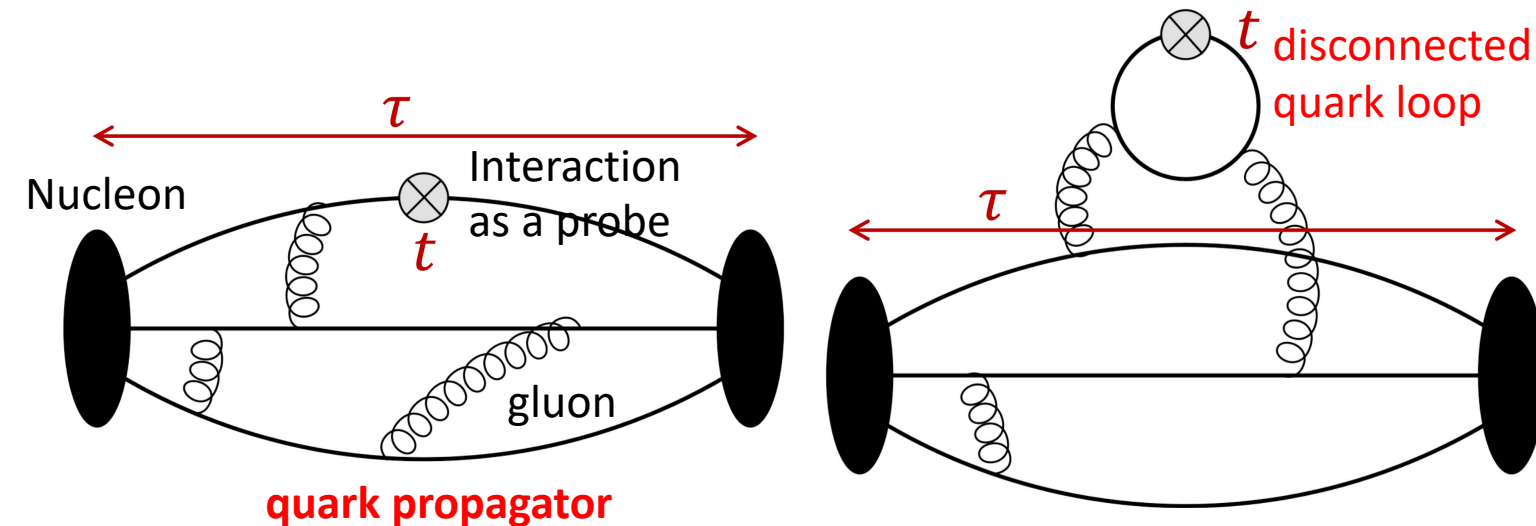
- Ensembles generated by MILC/CalLat/LANL Collaborations
- Improved gauge link smearing
- Precise quark mass parameter tuning
- Larger source-sink time separations to remove excited states
- Doubled range of momentum transfer $\vec{q} = 2\pi\vec{n}/L$ Largest spacelike 4-momentum transfer $Q^2 = -q^2$ increased from 0.45 GeV^2 to 0.82 GeV^2
- **Will improve all nucleon results in the continuum limit very significantly (2-3X)**

Nucleon structures from LQCD correlation functions:

$$C_{\Gamma}^q(t, \tau) = C_{\Gamma}^{q, \text{conn}}(t, \tau) + C_{\Gamma}^{q, \text{disc}}(t, \tau)$$

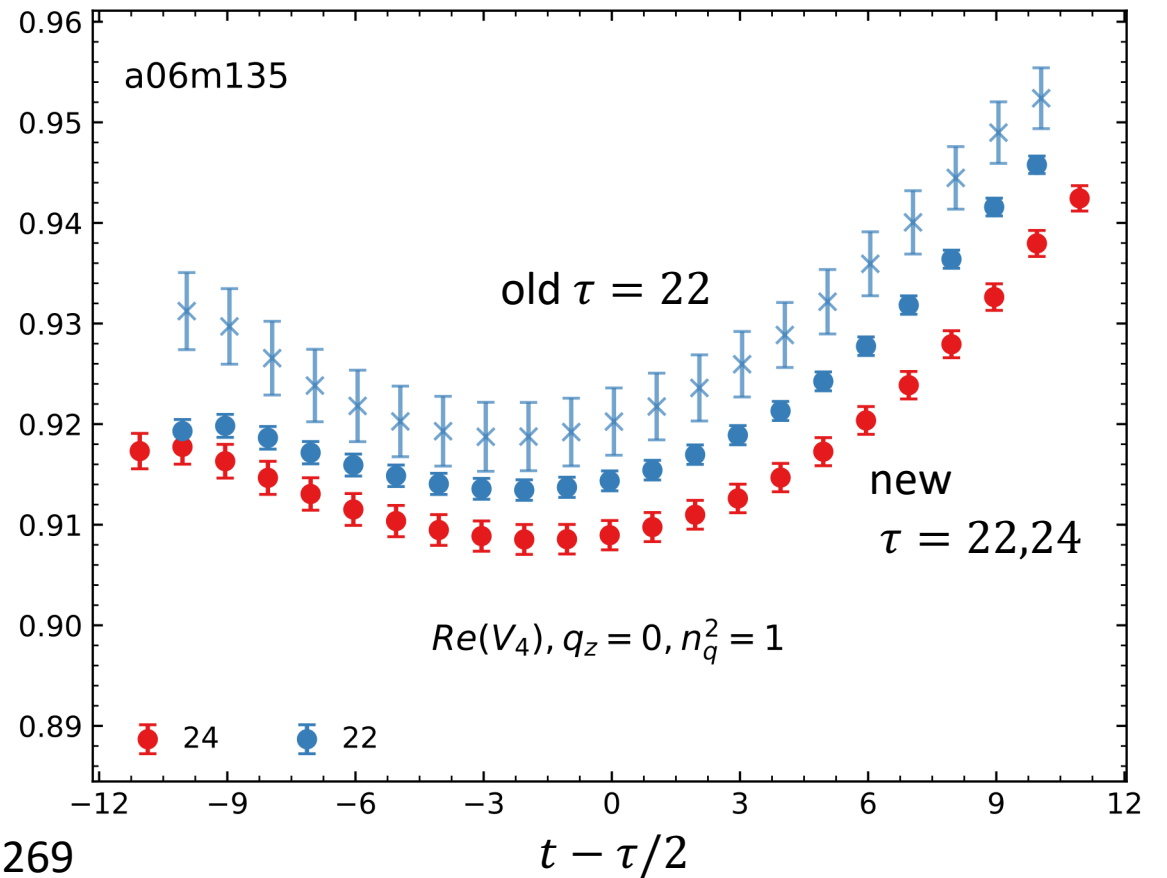
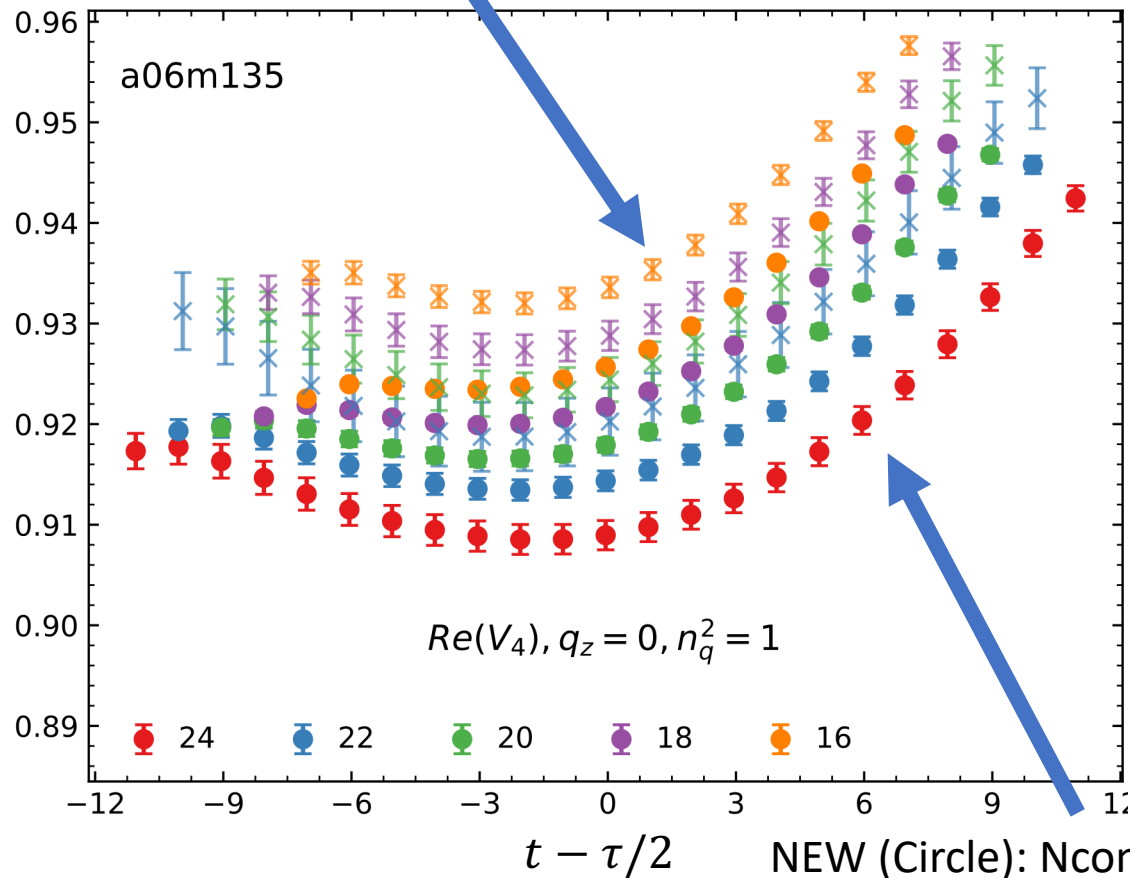
→ ***Charges, form factors, moments,
PDFs and GPDs***

$$\mathcal{R}_{\mathcal{O}}(t, \tau, \mathbf{p}, \mathbf{0}) = \frac{C_{\mathcal{O}}^{3\text{pt}}(t, \tau; \mathbf{p}, \mathbf{0})}{C^{2\text{pt}}(\tau, \mathbf{p})} \times \left[\frac{C^{2\text{pt}}(t, \mathbf{p}) C^{2\text{pt}}(\tau, \mathbf{p}) C^{2\text{pt}}(\tau - t, \mathbf{0})}{C^{2\text{pt}}(t, \mathbf{0}) C^{2\text{pt}}(\tau, \mathbf{0}) C^{2\text{pt}}(\tau - t, \mathbf{p})} \right]^{1/2}$$



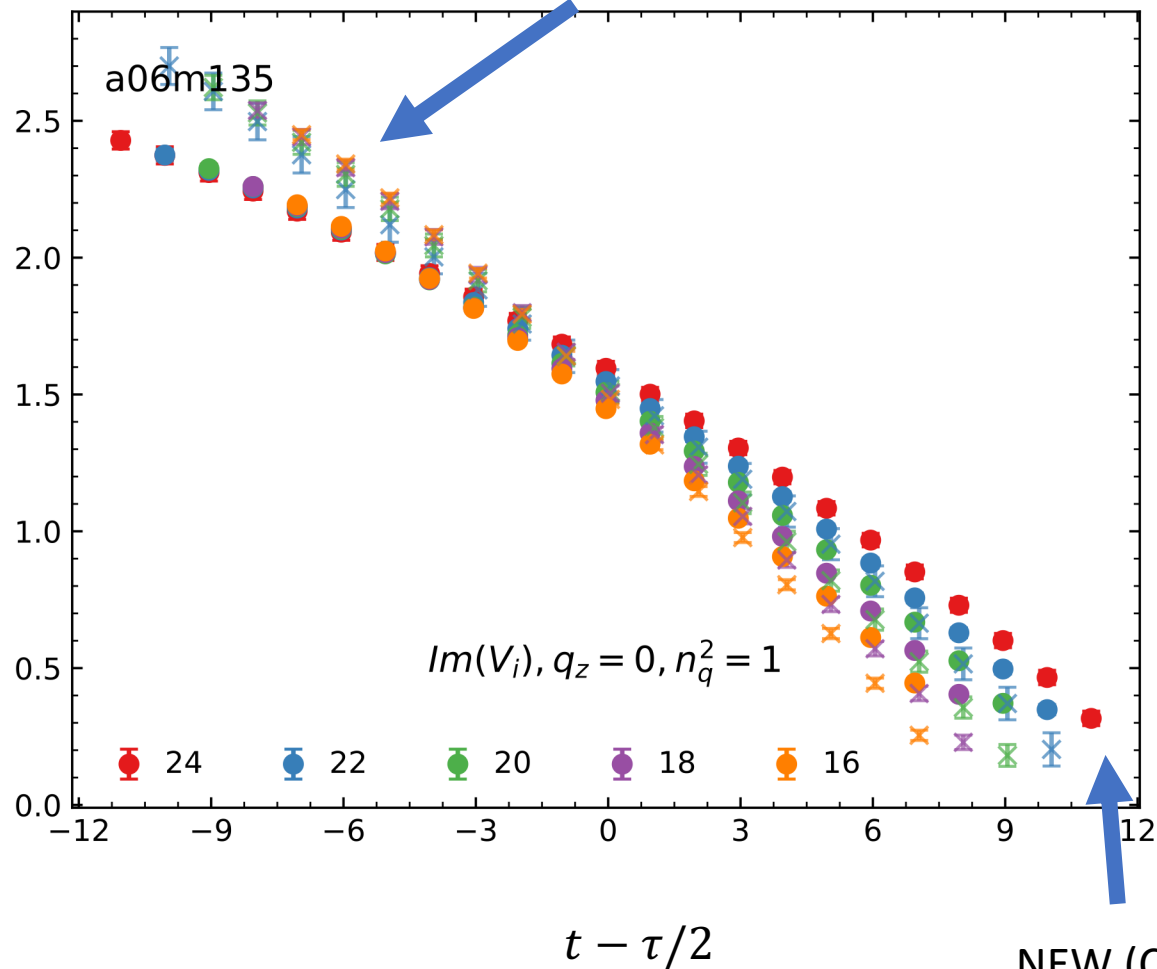
G_E : improvement in signal at $a \sim 0.06\text{fm}$, $|\vec{n}|^2 = 1$

OLD (Cross): Nconf=1074 (PNDME 20 used 675)

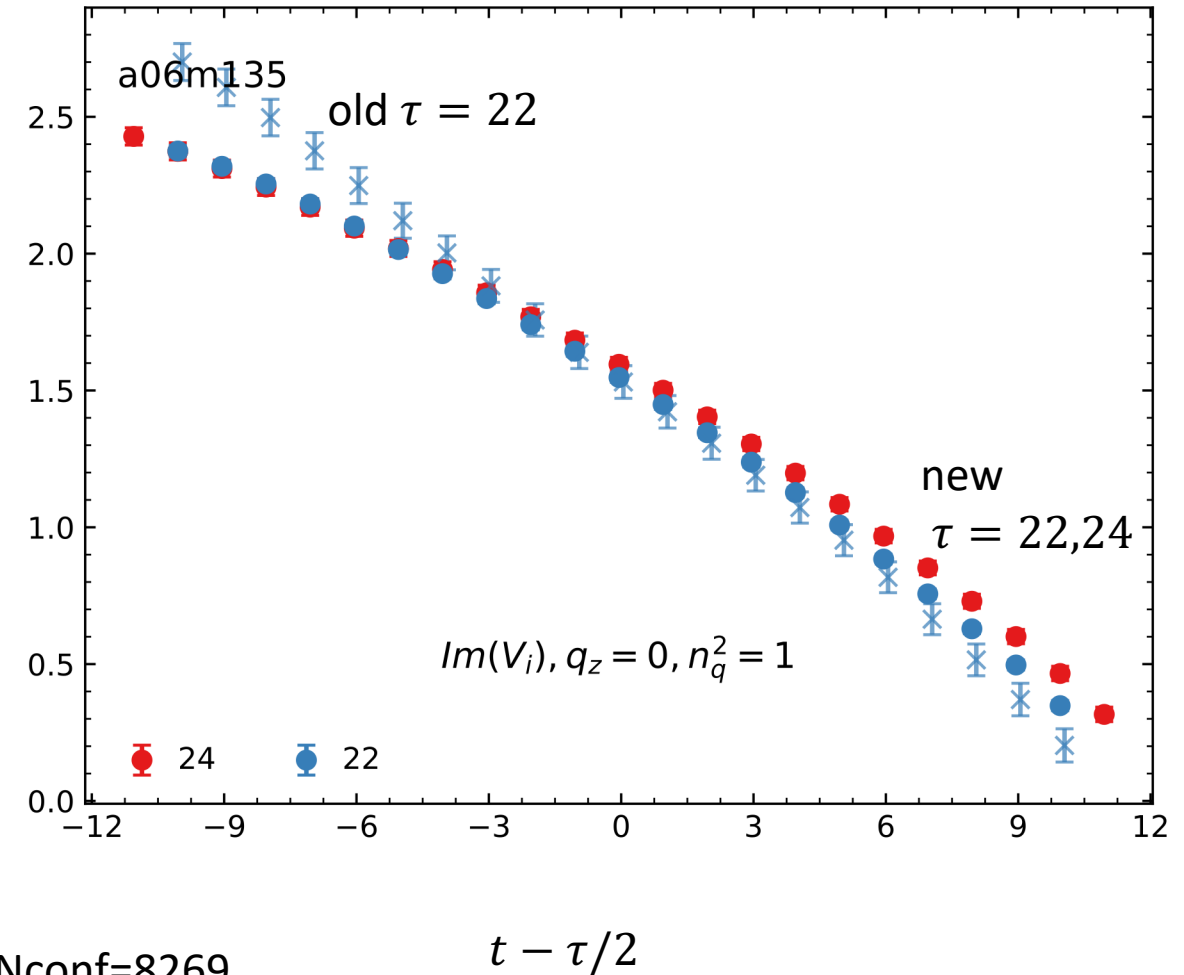


G_E : improvement in signal at $a \sim 0.06\text{fm}$, $|\vec{n}|^2 = 1$

OLD (Cross): Nconf=1074 (PNDME 20 used 675)

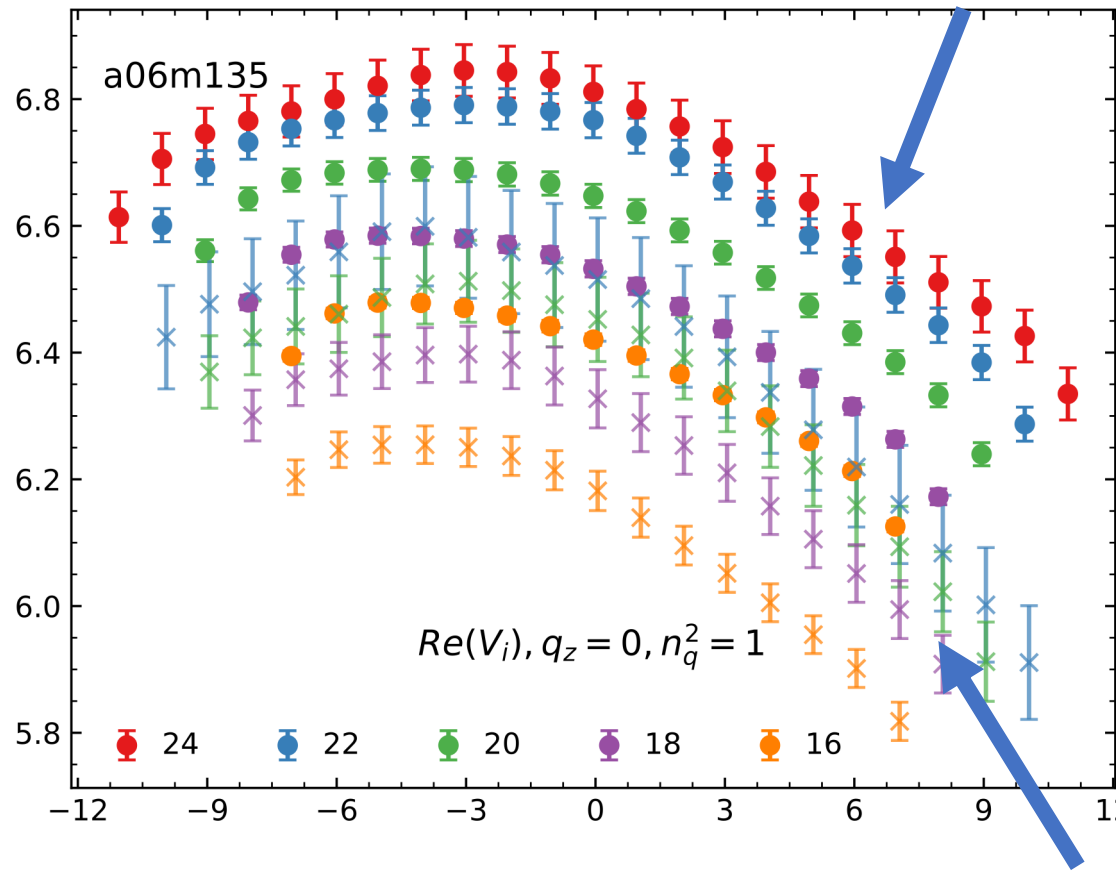


NEW (Circle): Nconf=8269



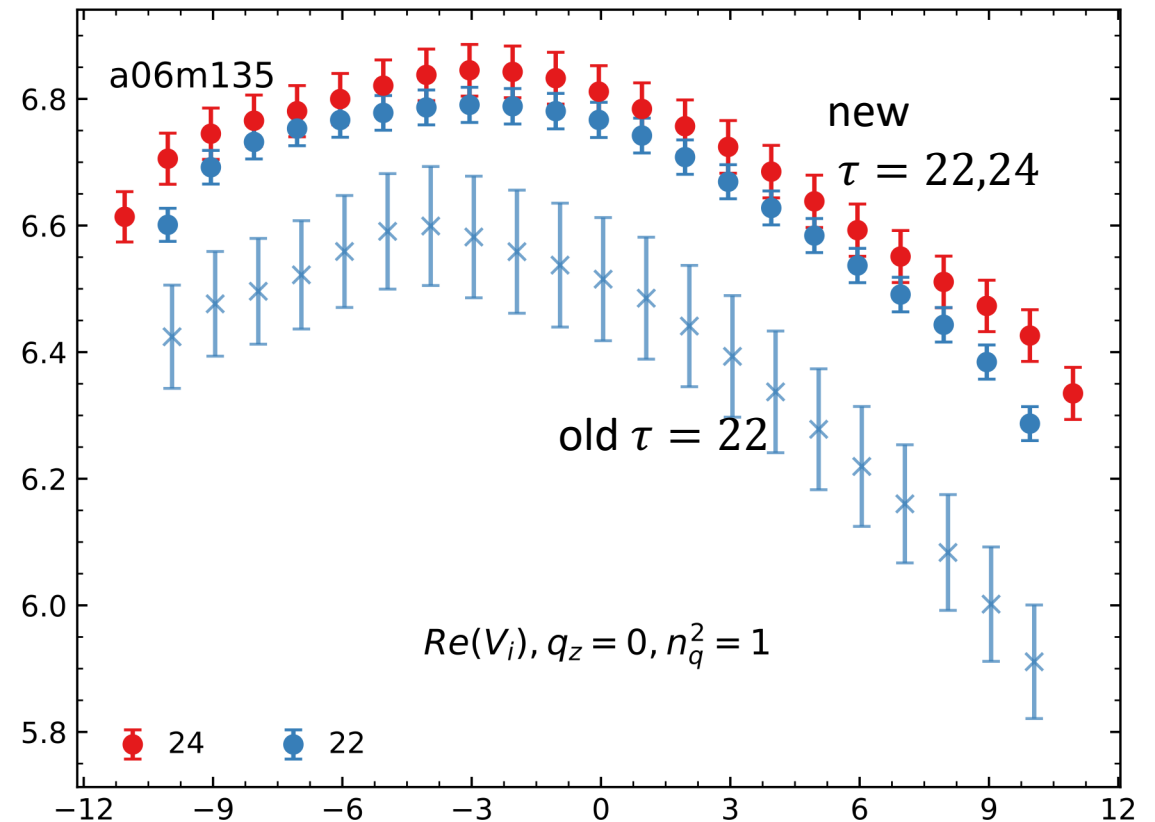
G_M : improvement in signal at $a \sim 0.06\text{fm}$, $|\vec{n}|^2 = 1$

NEW (Circle): Nconf=8269



$t - \tau/2$

OLD (Cross): Nconf=1074 (PNDME 20 used 675)



old $\tau = 22$

new
 $\tau = 22, 24$

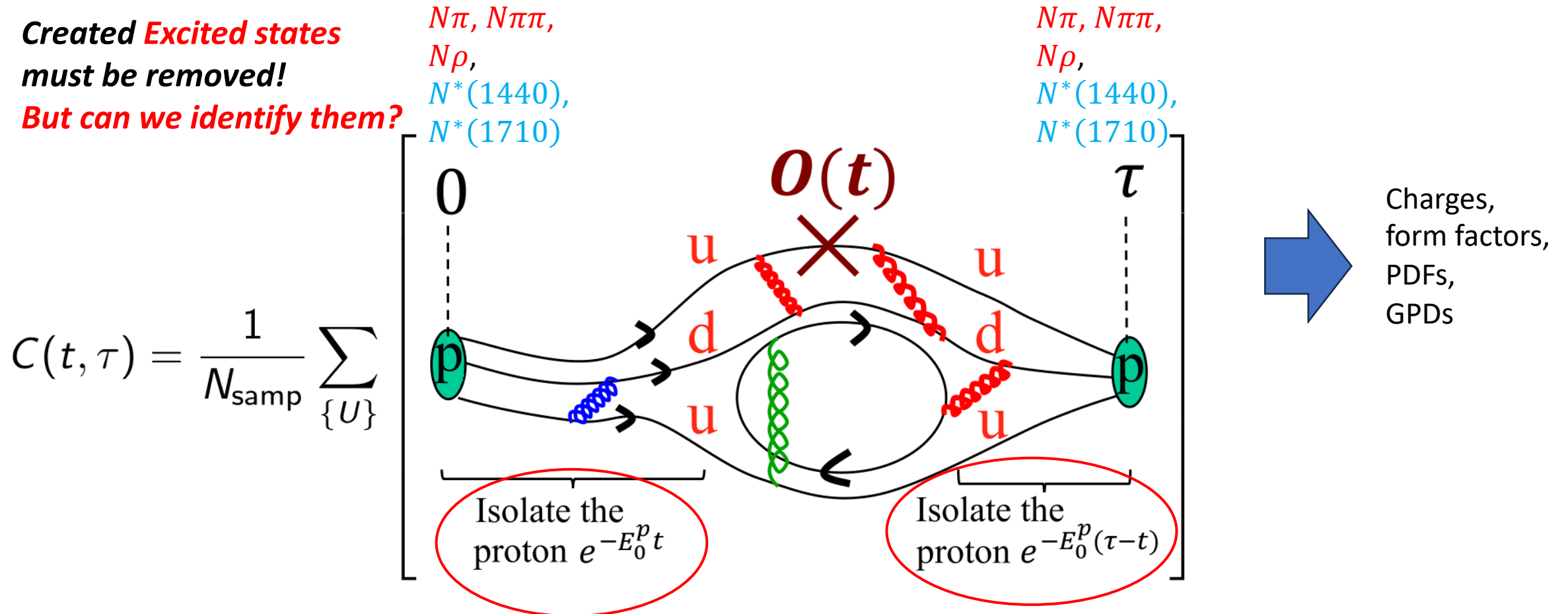
$Re(V_i), q_z = 0, n_q^2 = 1$

$t - \tau/2$

Remove Excited State (ES) effects from Nucleon Matrix Element

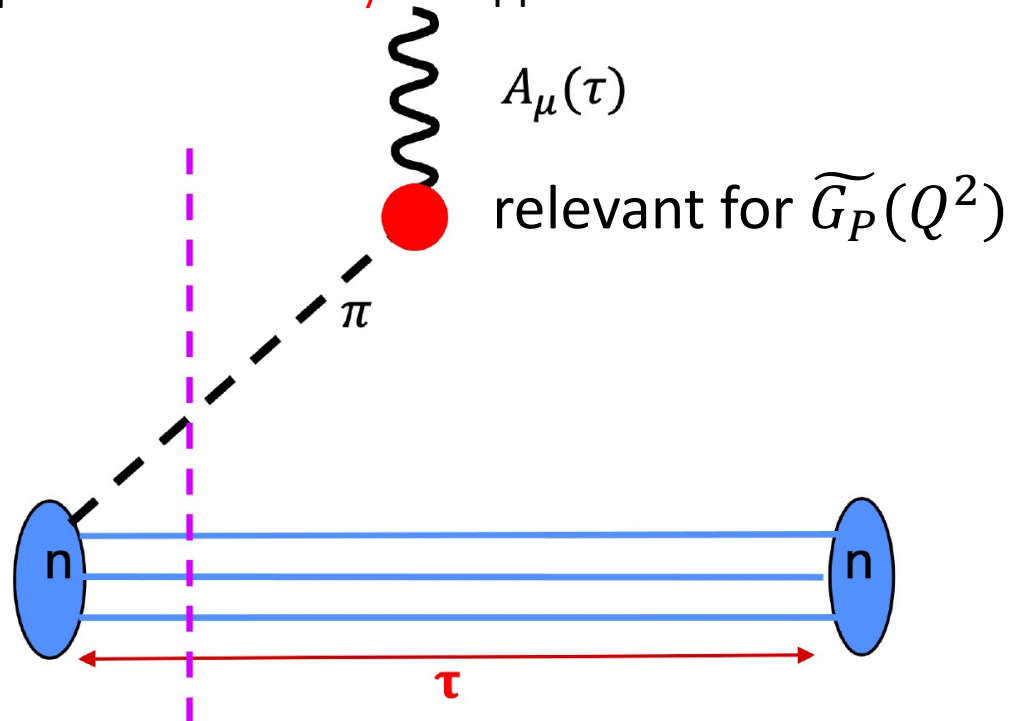
- Nucleon signal/noise decays $\propto e^{-(E-1.5M_\pi)\tau}$ with Euclidean time τ . Cannot go to large τ

Created **Excited states**
must be removed!
But can we identify them?



$N\pi$ -Excited state contribution can be enhanced

- To control ESC, we need to first determine the spectrum
- In nucleon 2pt functions, $1/L^3$ -suppressed coupling to $N\pi$ -state makes it challenging to extract its spectrum
- In certain 3pt functions (matrix elements), ChPT suggests $N\pi$ -state contributions are enhanced and compensate for the $1/L^3$ -suppression

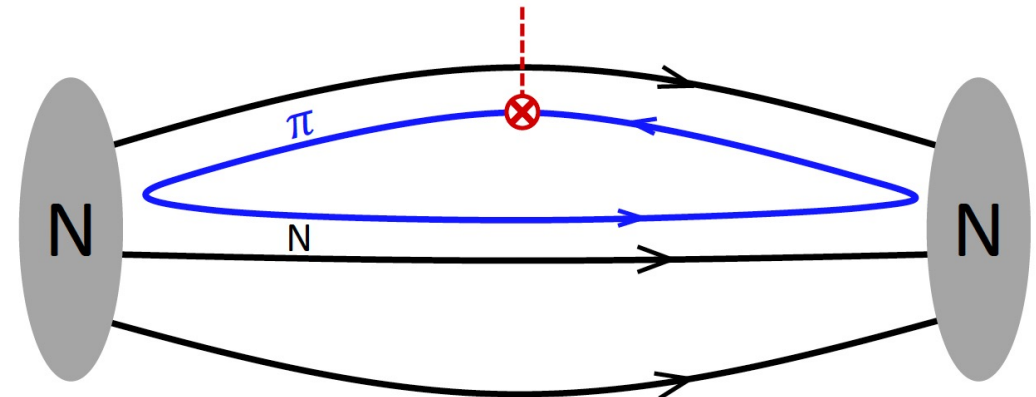


ChiPT: O.Bär, PRD99 054506 (2019), PRD 100 054507 (2019)

Lattice: Jang et al. PRL [PNDME 2020]

Park et al. PRD 105 054505 [NME 2022]

Scalar operator
relevant for $g_S^{u,d}$ and $\sigma_{\pi N} = m_{ud} g_S^{u+d}$

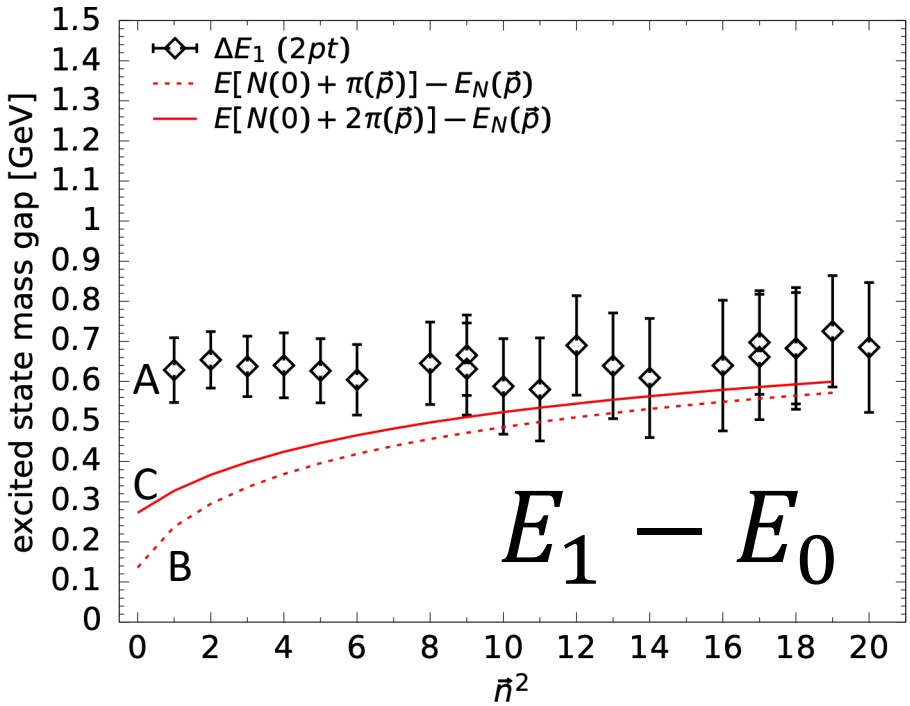
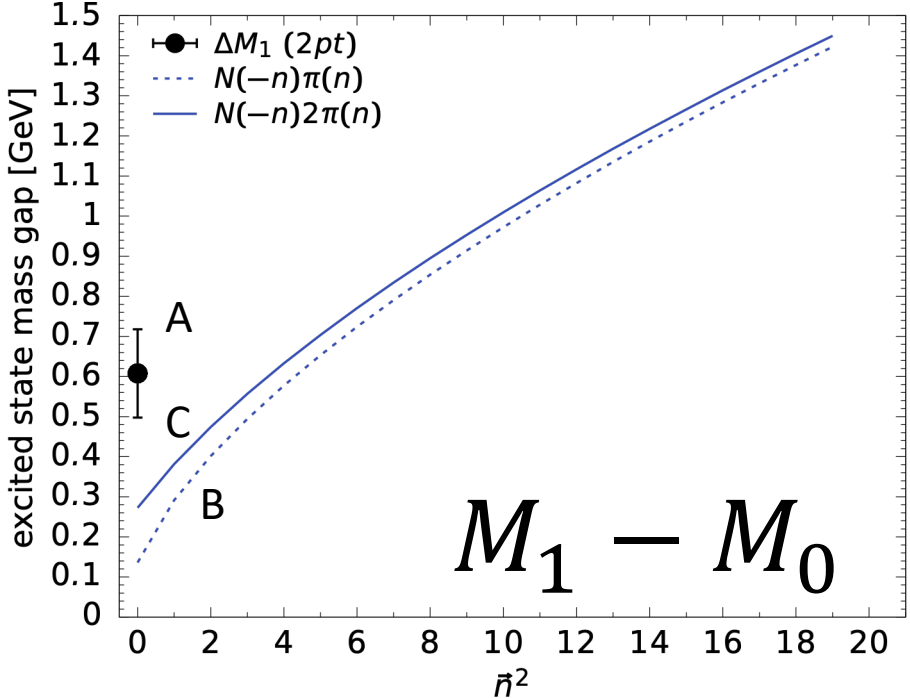


ChiPT and Lattice:
Gupta et al., [PNDME 2021]

3 Excited State Fit Strategies for M_1 and M_2

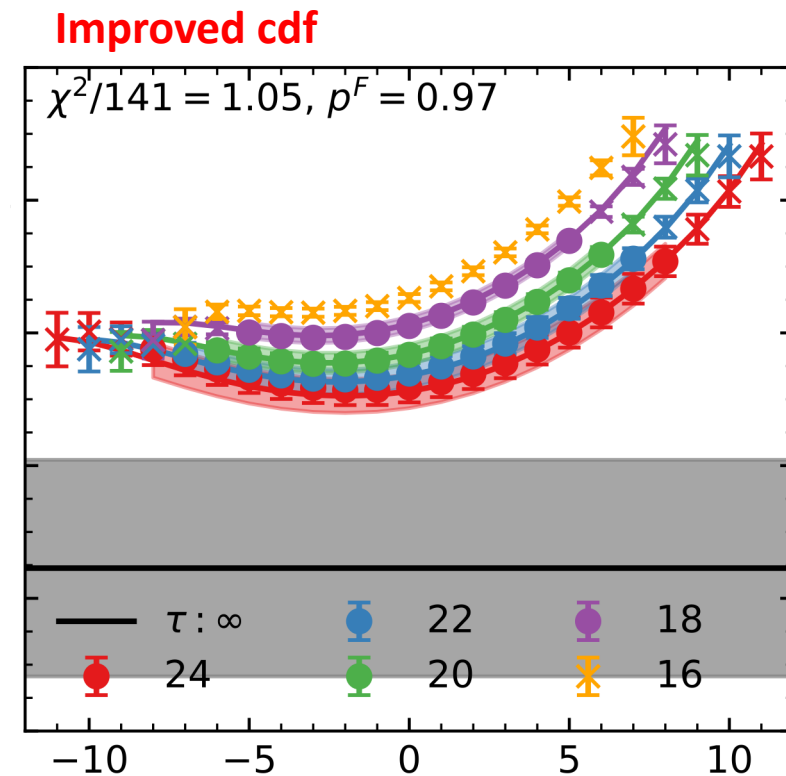
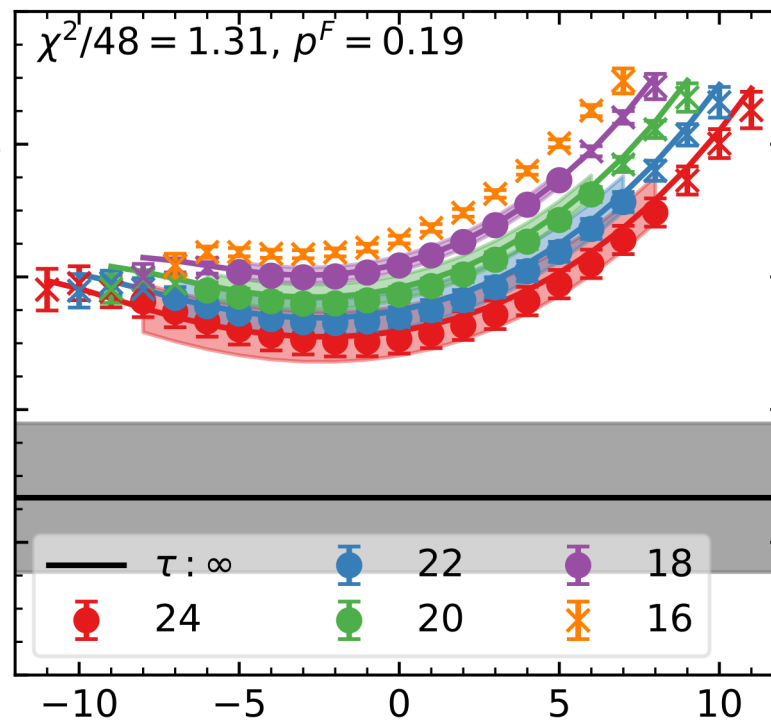
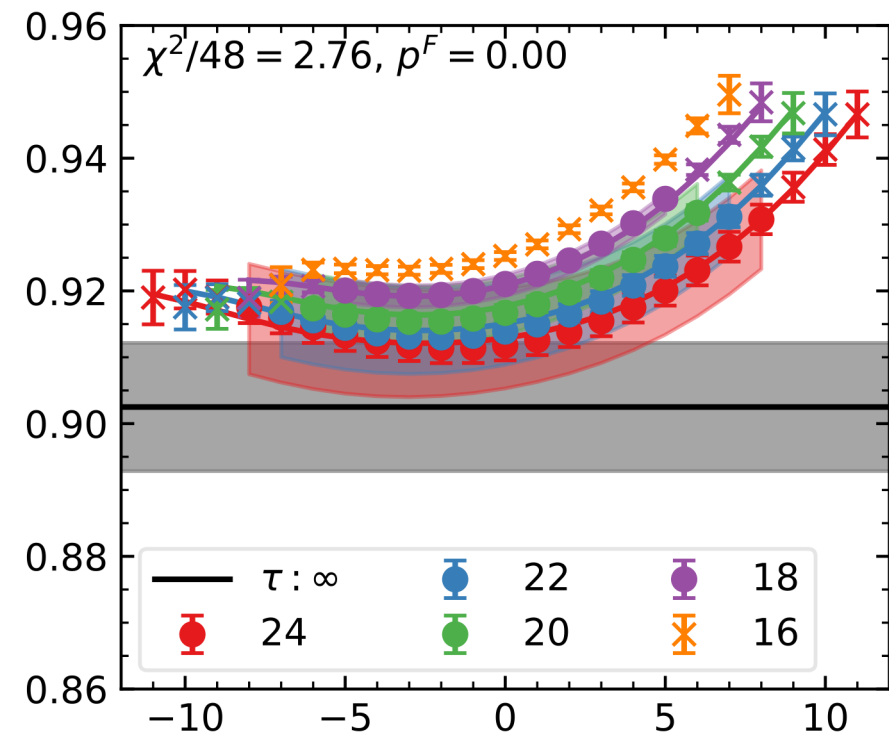
M_0 taken from 4-state fit

Strategy	$M_1 - M_0,$ $E_1 - E_0$	$M_2 - M_0,$ $E_2 - E_0$	Note
A	0.6 GeV	2.0 GeV	* Spectrum determined from 2pt fits with empirical Bayesian priors
B	Noninteracting $N - \pi$ $\sim 0.2\text{GeV}$	1.5 GeV	
C	Noninteracting $N - 2\pi$ $\sim 0.3\text{GeV}$	0.6 GeV	* Vector-meson dominance model, * Simultaneous fits to ReV_4, ImV_i, ReV_i



G_E from $Re(V_4)$

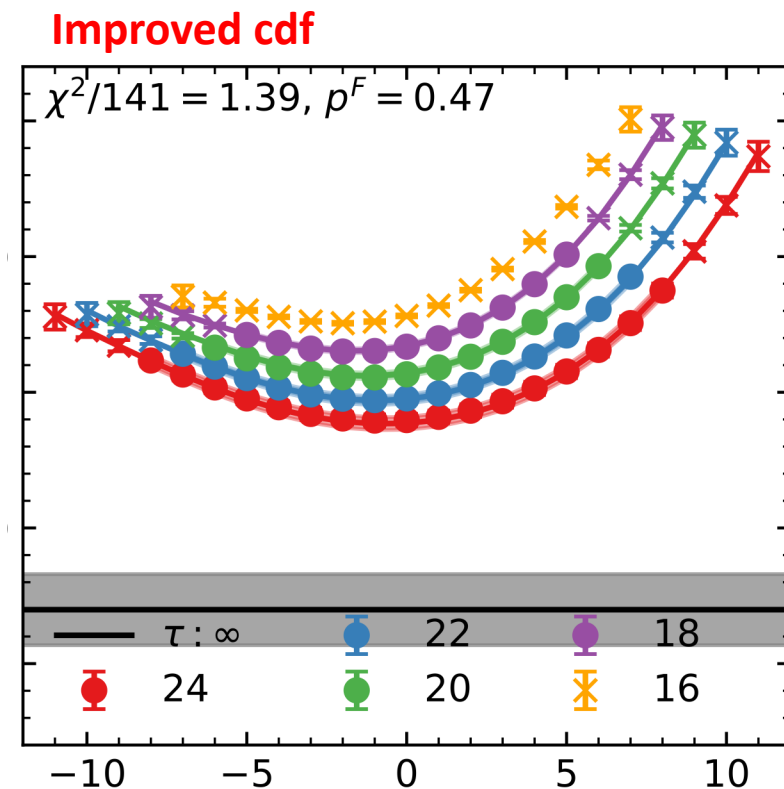
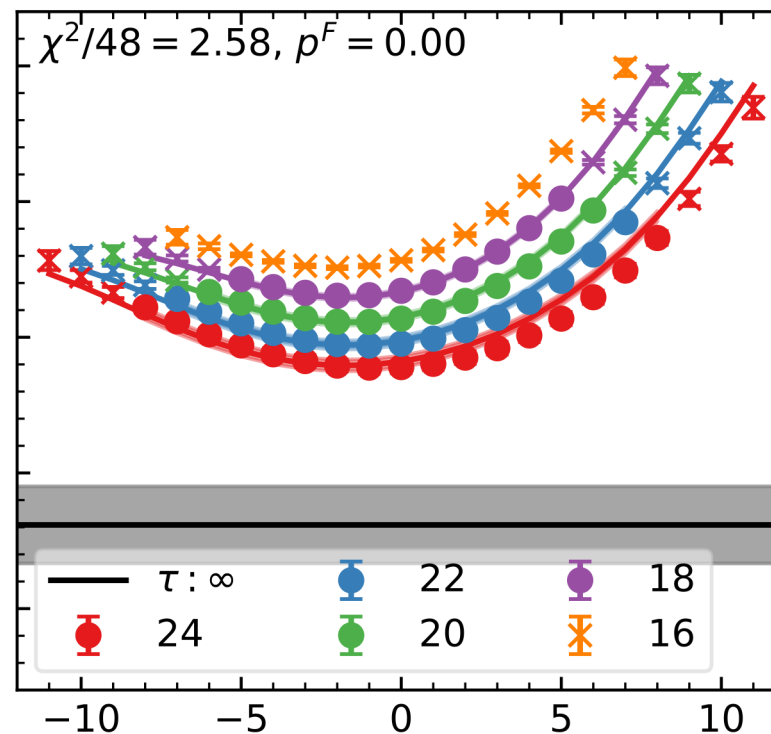
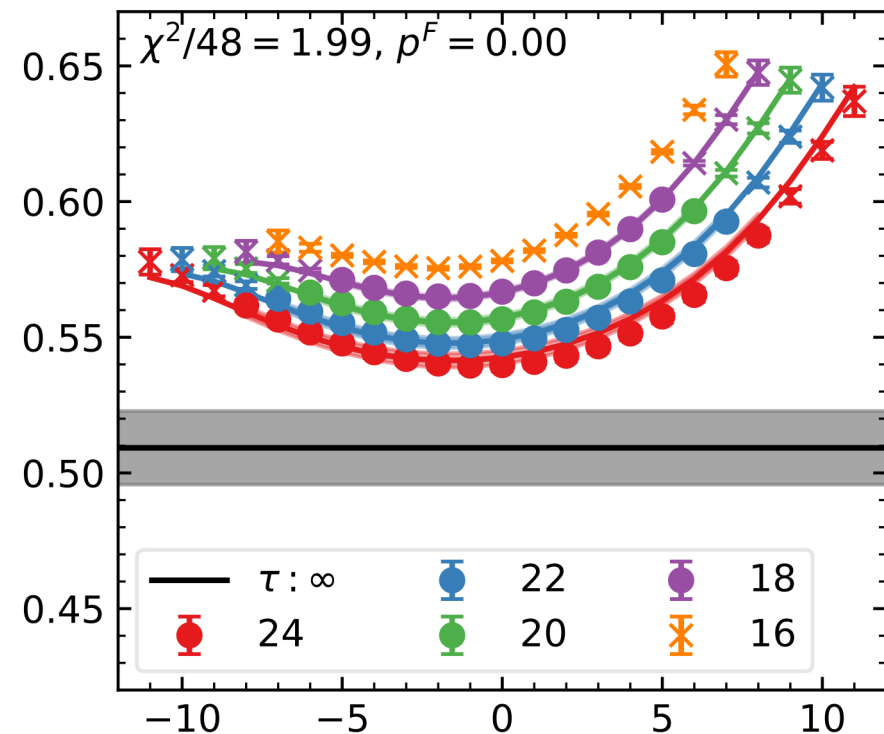
: χ^2 for the 3 strategies at $a \sim 0.06 \text{ fm}$, $|\vec{n}|^2 = 1$



Strategy	A	B	C
$M_1 - M_0, M_2 - M_0$	0.6 GeV, 2.0 GeV	0.2 GeV, 1.5 GeV	0.3 GeV, 0.6 GeV

G_E from $Re(V_4)$

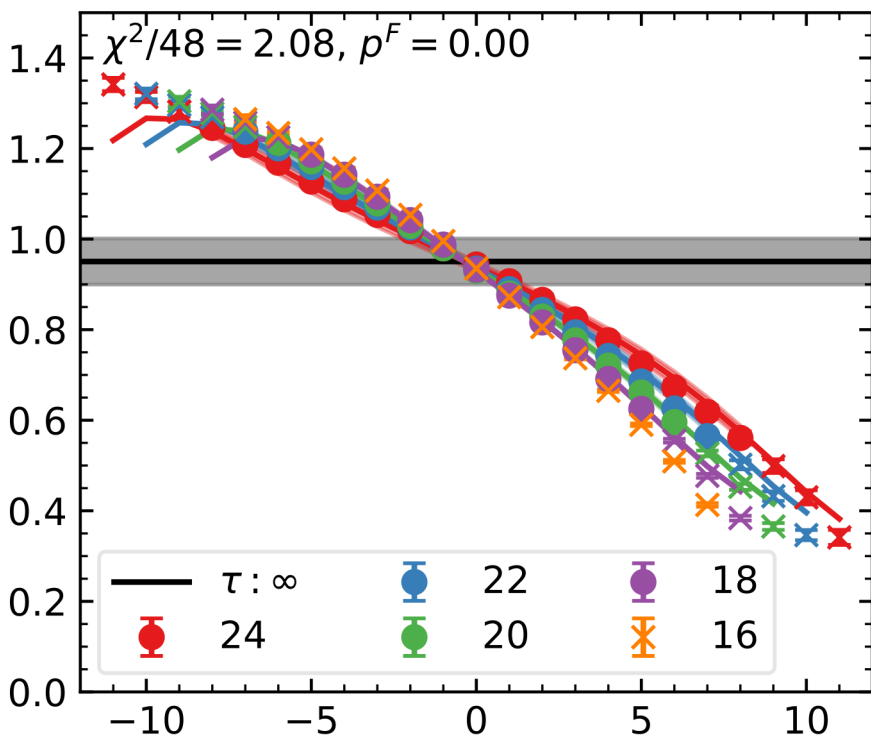
: χ^2 for the 3 strategies at $a \sim 0.06 \text{ fm}$, $|\vec{n}|^2 = 6$



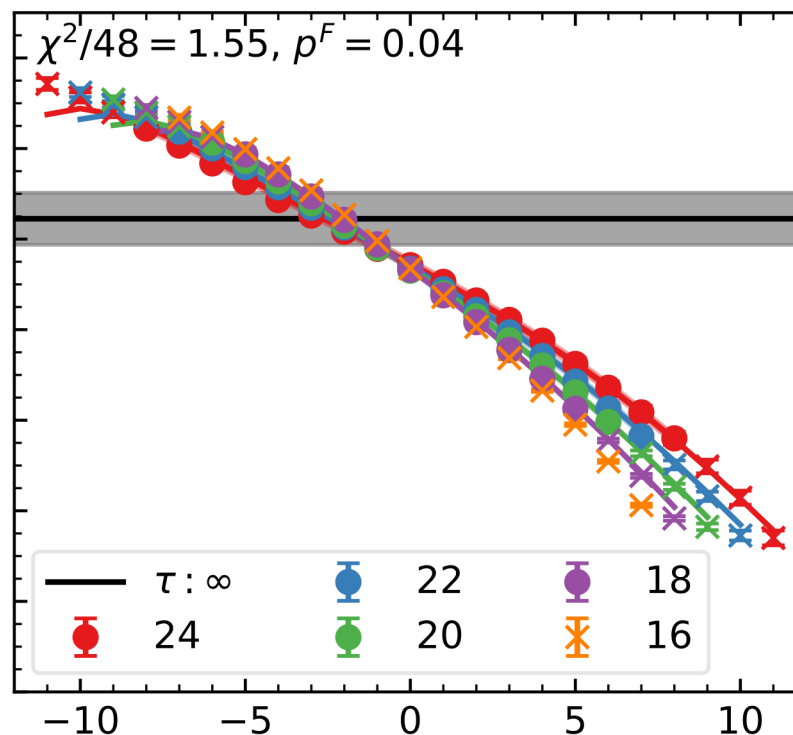
Strategy	A	B	C
$M_1 - M_0, M_2 - M_0$	0.6 GeV, 2.0 GeV	0.2 GeV, 1.5 GeV	0.3 GeV, 0.6 GeV

G_E from $Im(V_i)$

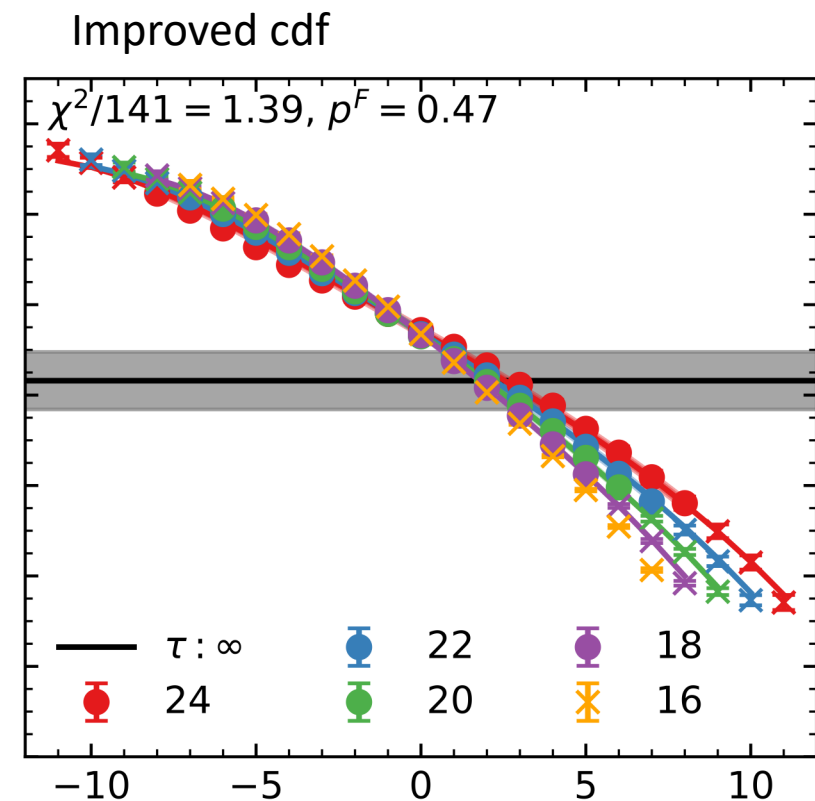
: χ^2 for the 3 strategies at $a \sim 0.06 \text{ fm}$, $|\vec{n}|^2 = 6$



GE(ReV4)=0.528(14)
GE(ImVi)=0.577(31) **inconsistent**



GE(ReV4)=0.500(15)
GE(ImVi)=0.629(32) **inconsistent**

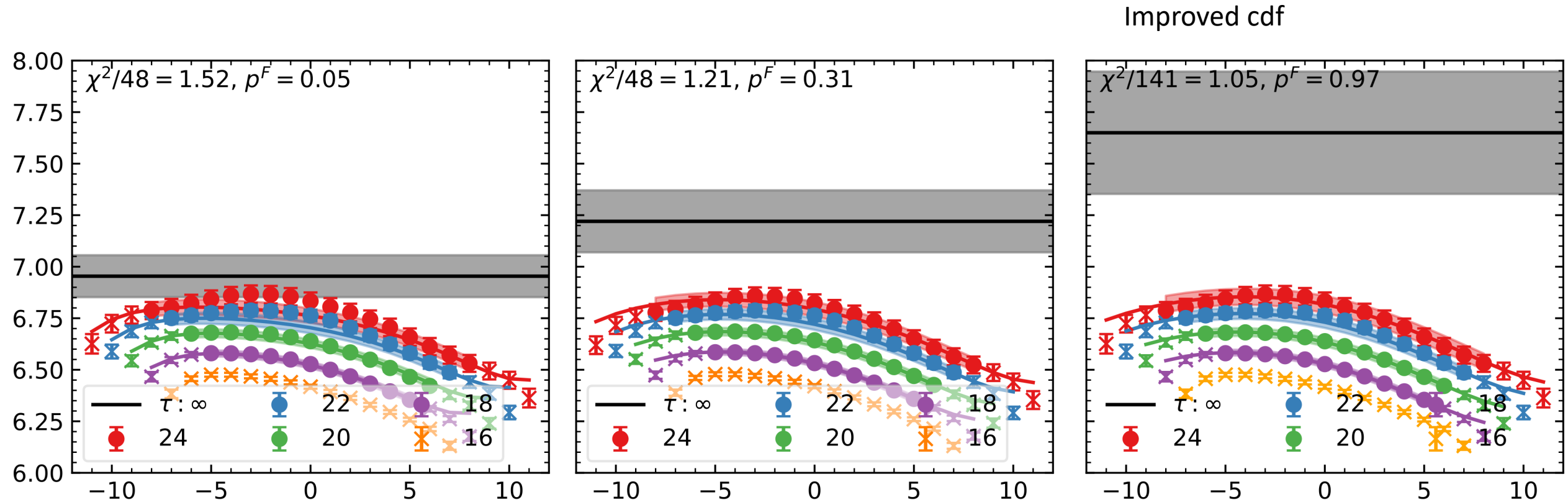


GE(ReV4)=0.487(14)
GE(ImVi)=0.505(39) **consistent**

Strategy	A	B	C
$M_1 - M_0, M_2 - M_0$	0.6 GeV, 2.0 GeV	0.2 GeV, 1.5 GeV	0.3 GeV, 0.6 GeV

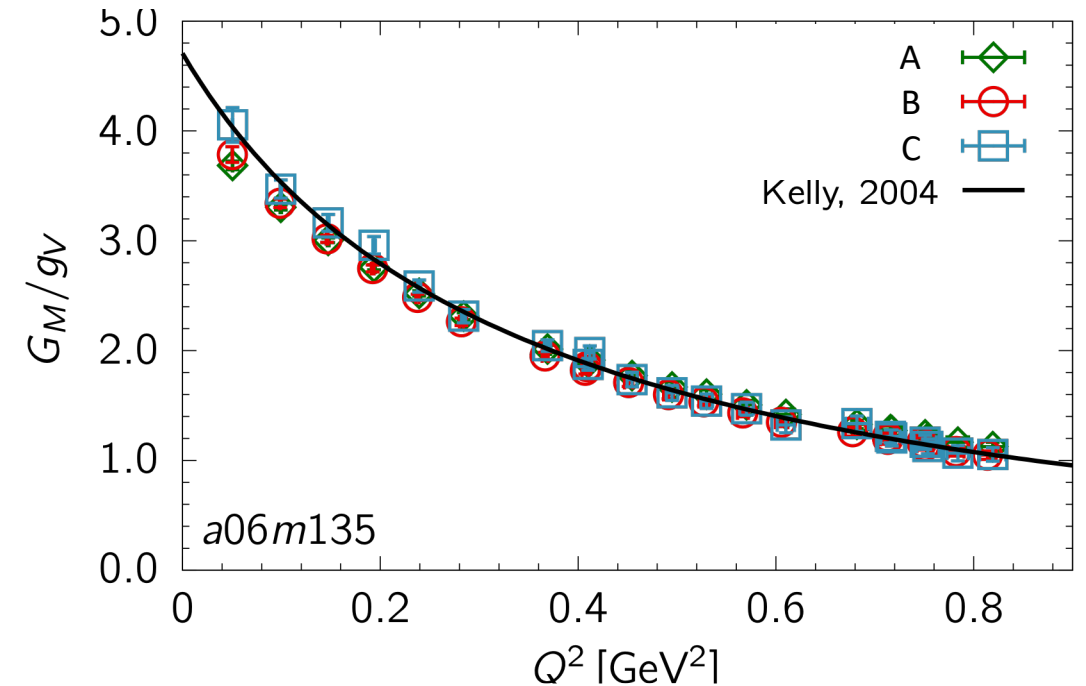
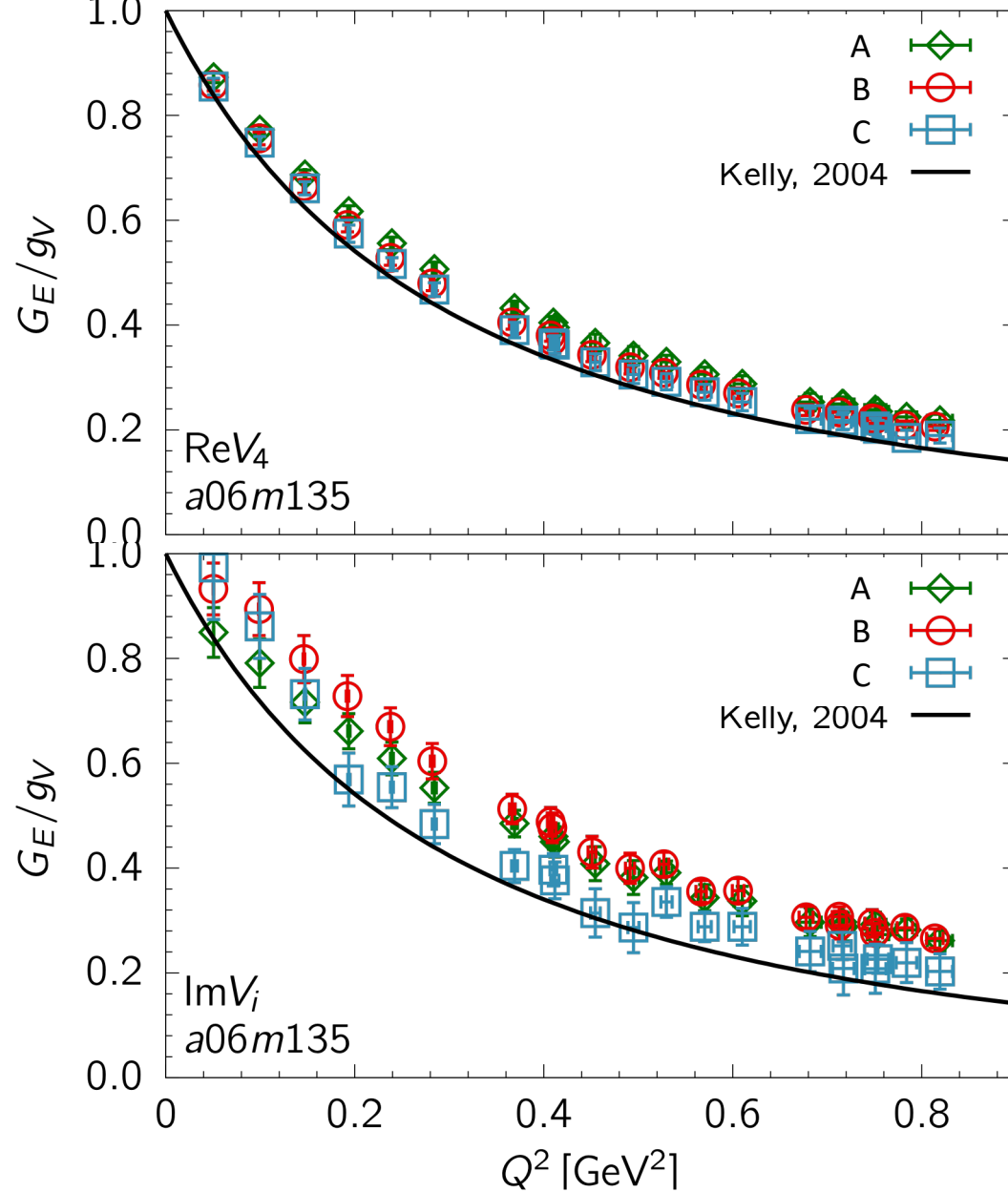
G_M from $Re(V_i)$

: χ^2 for the 3 strategies at $a \sim 0.06 \text{ fm}$, $|\vec{n}|^2 = 1$



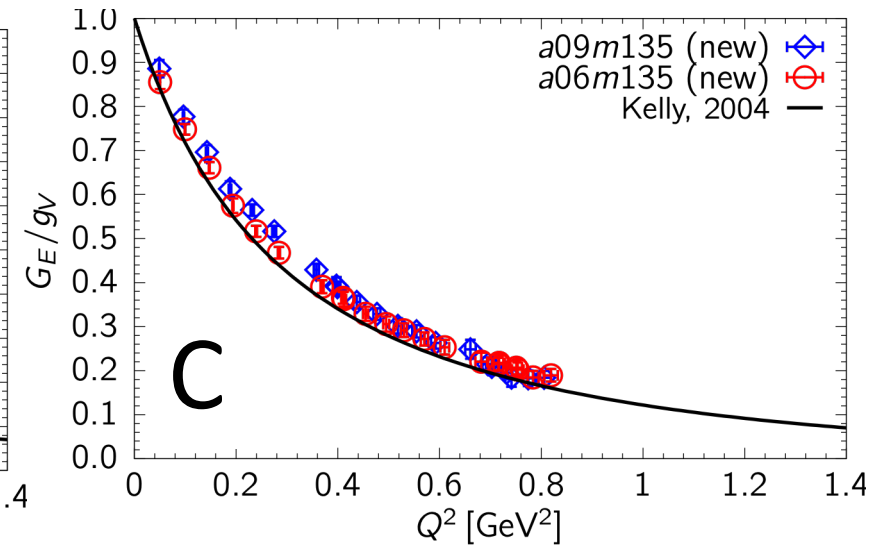
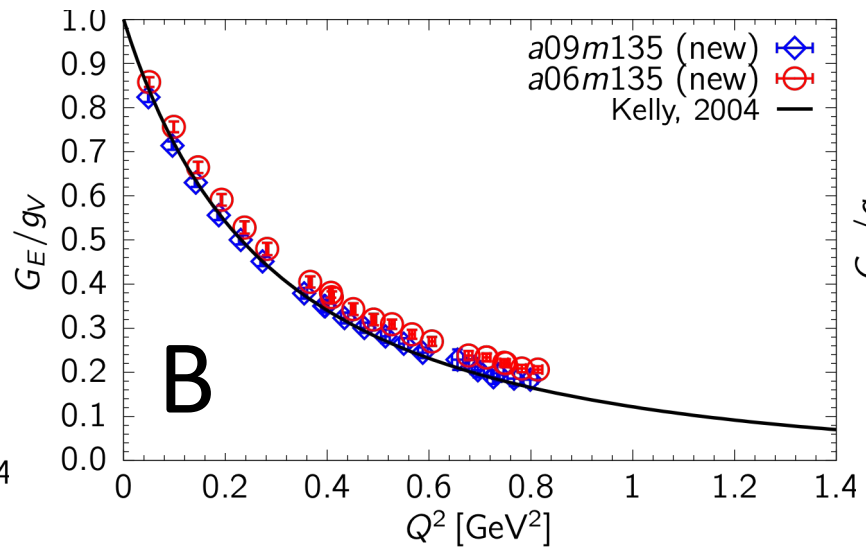
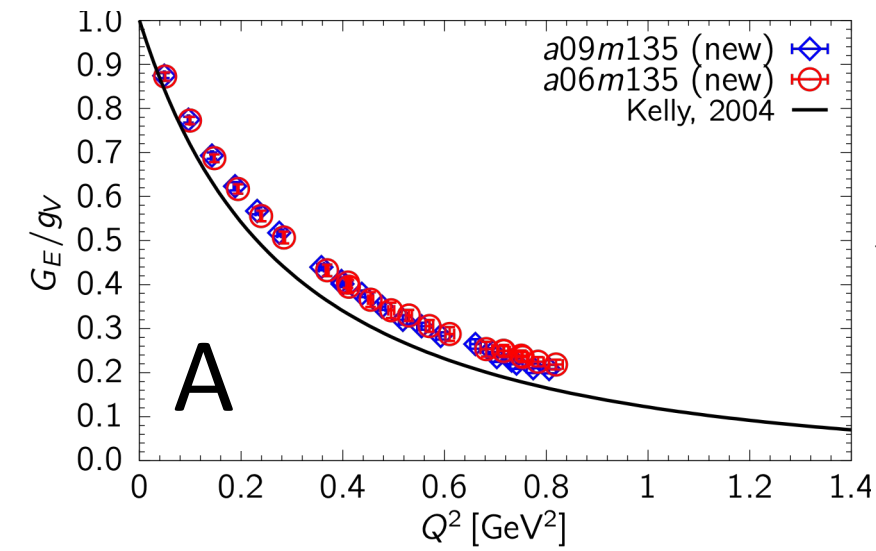
Strategy	A	B	C
$M_1 - M_0, M_2 - M_0$	0.6 GeV, 2.0 GeV	0.2 GeV, 1.5 GeV	0.3 GeV, 0.6 GeV

$G_E(Q^2), G_M(Q^2)$ at $a \sim 0.06\text{fm}$ comparison of ES strategies



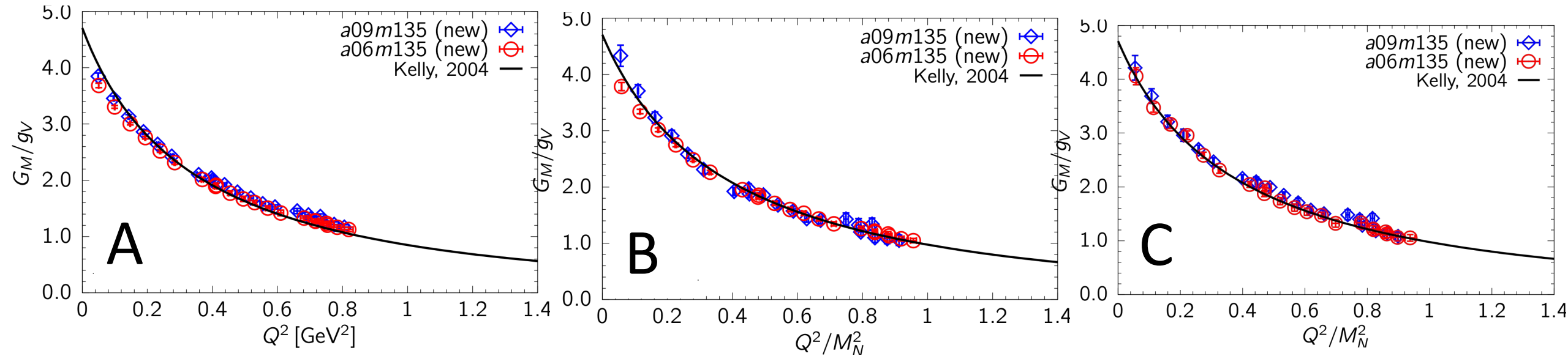
Strategy	A	B	C
$M_1 - M_0, M_2 - M_0$	0.6 GeV, 2.0 GeV	0.2 GeV, 1.5 GeV	0.3 GeV, 0.6 GeV

$G_E(Q^2)$ a-dependence



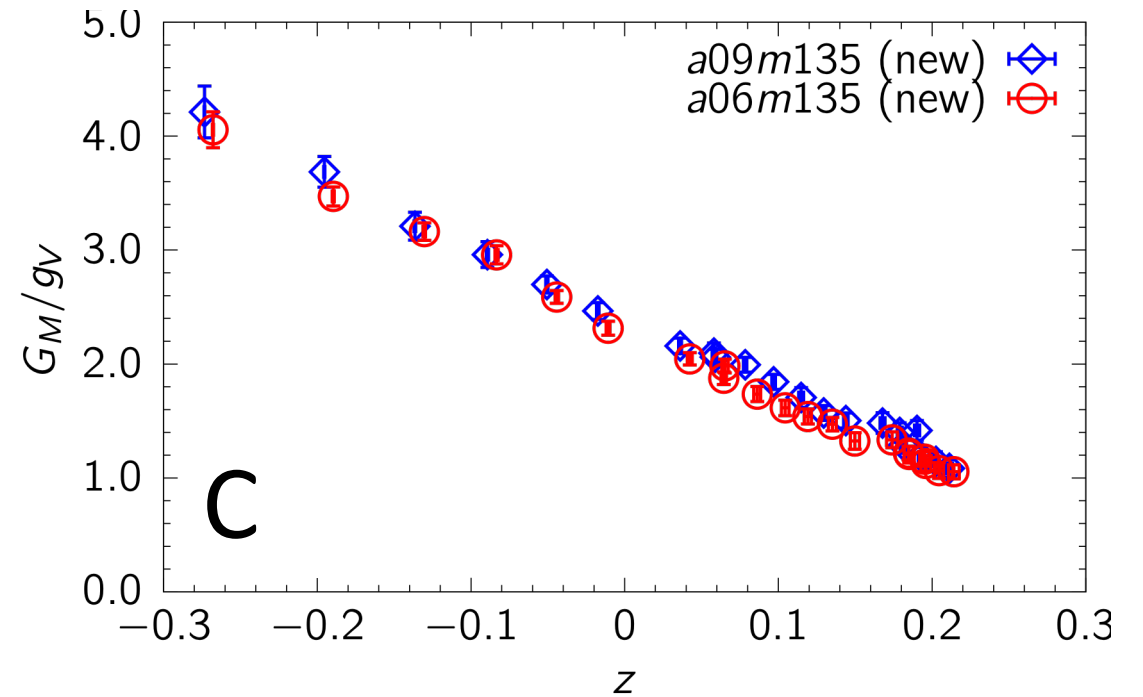
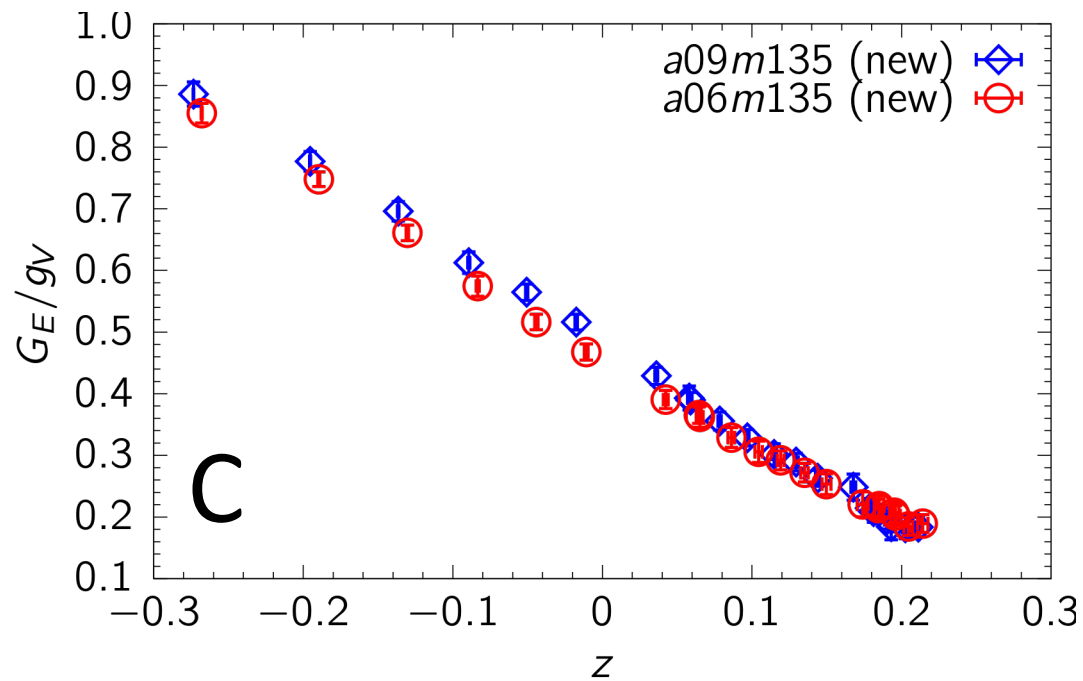
Strategy	A	B	C
$M_1 - M_0, M_2 - M_0$	0.6 GeV, 2.0 GeV	0.2 GeV, 1.5 GeV	0.3 GeV, 0.6 GeV

$G_M(Q^2)$ a-dependence

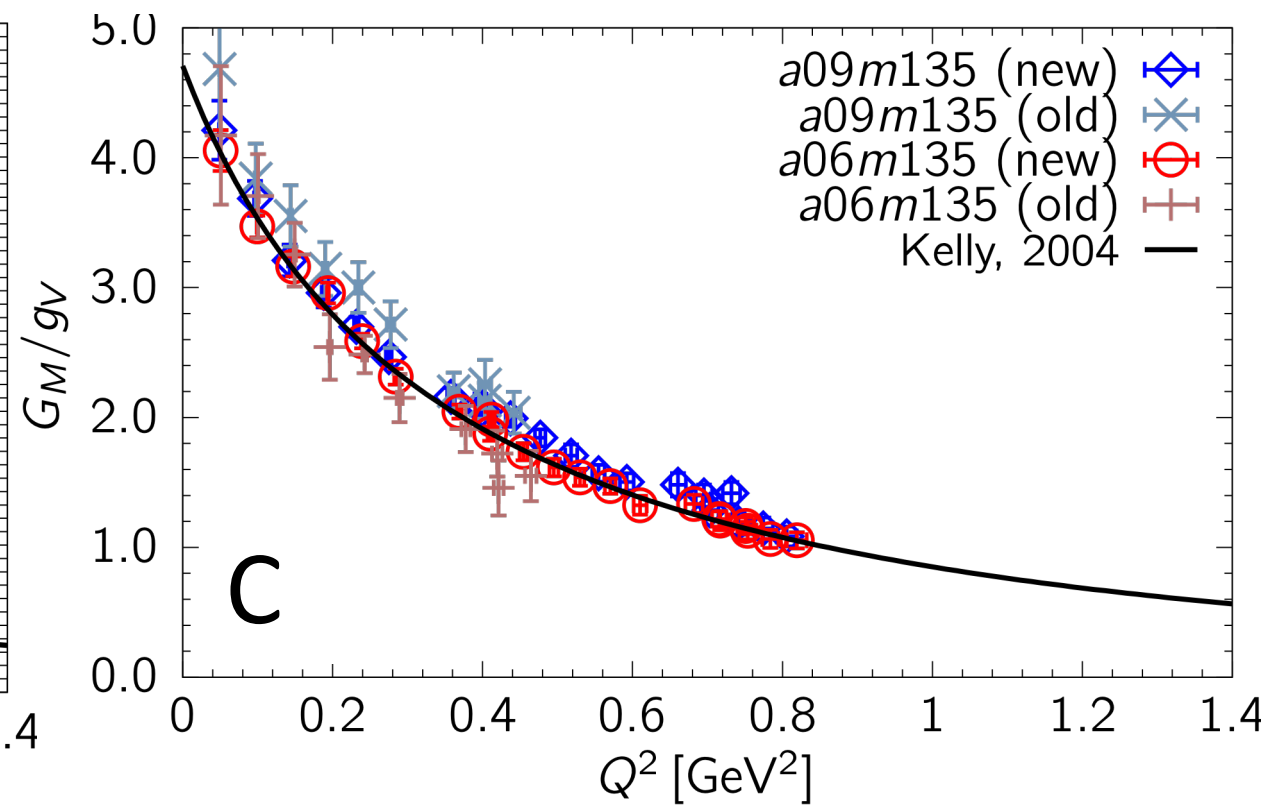
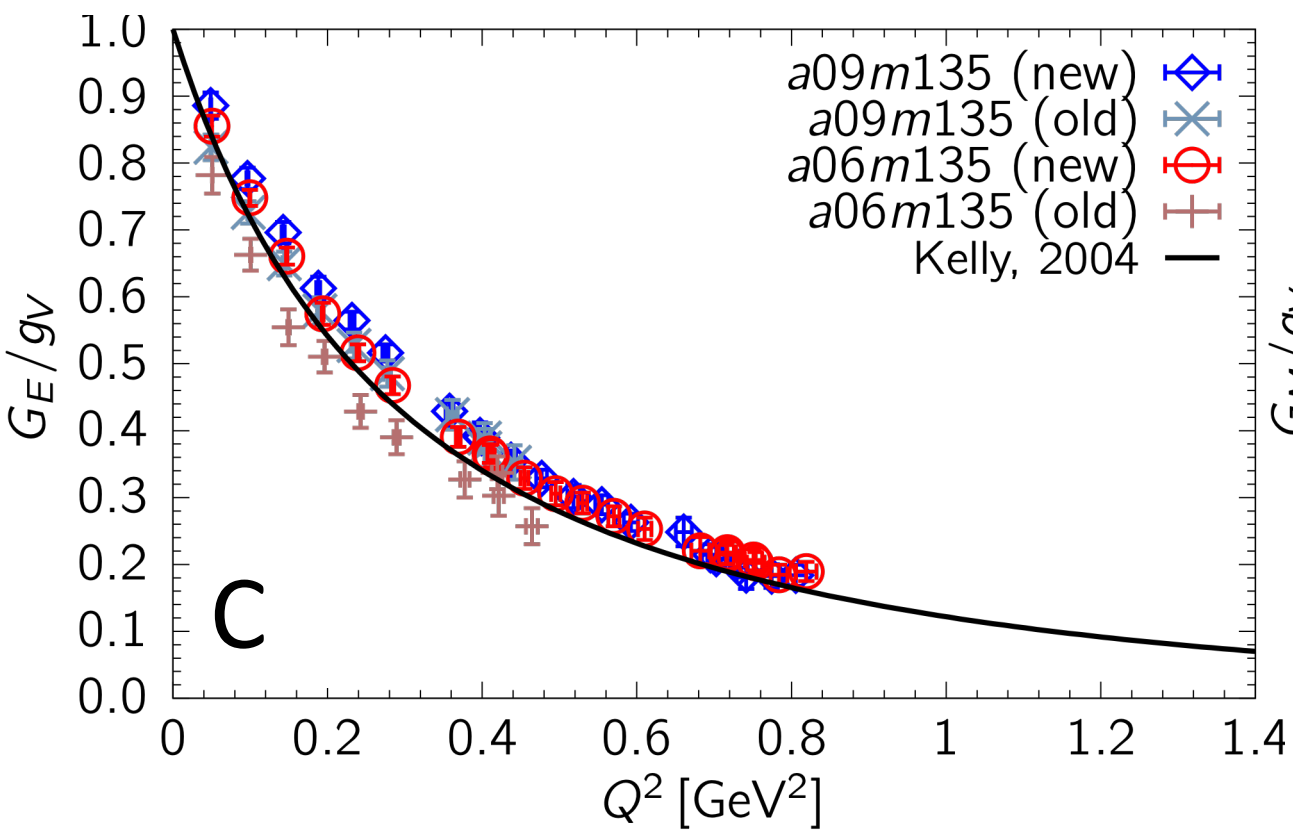


Strategy	A	B	C
$M_1 - M_0, M_2 - M_0$	0.6 GeV, 2.0 GeV	0.2 GeV, 1.5 GeV	0.3 GeV, 0.6 GeV

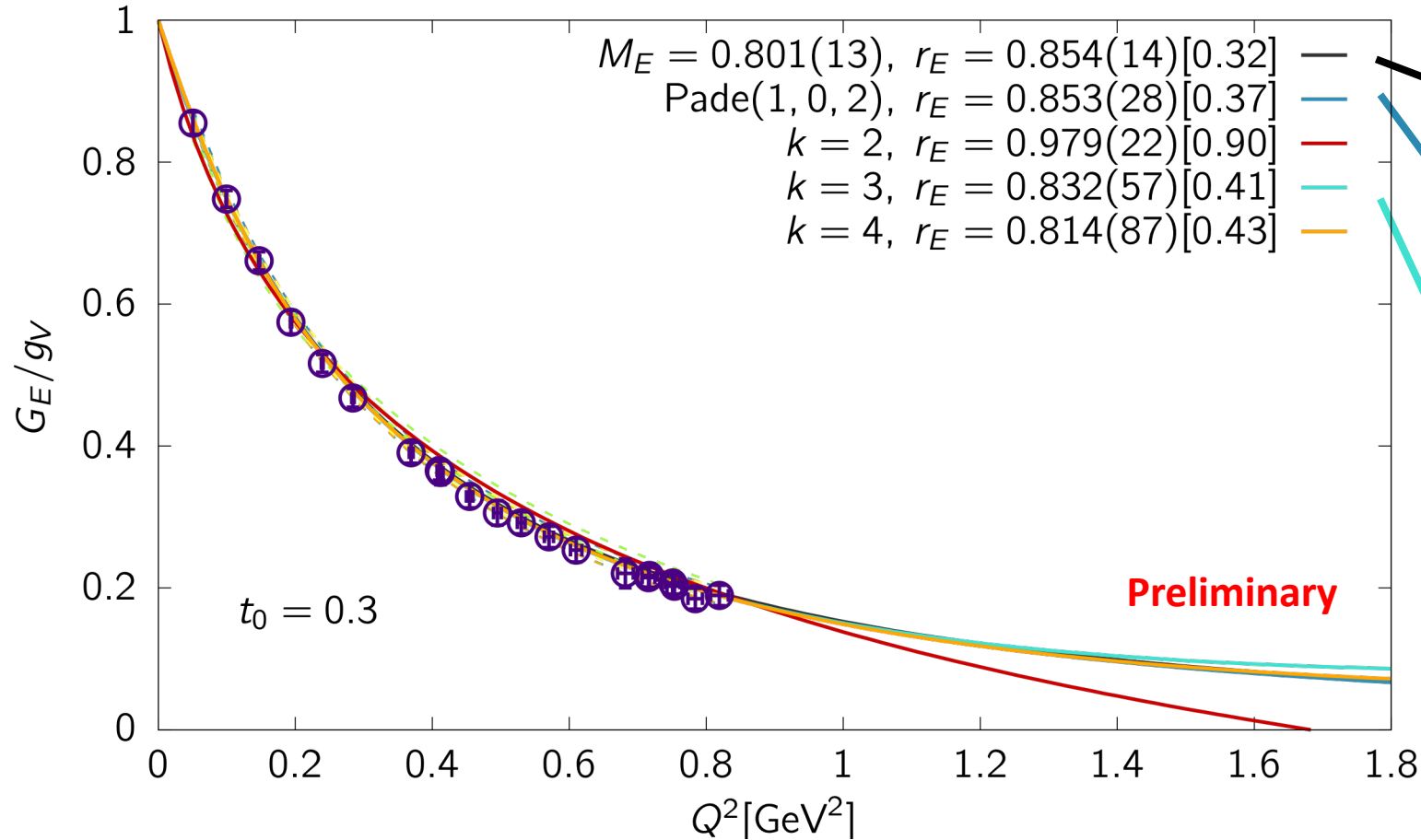
$G_E(Q^2)$, $G_M(Q^2)$ and z-expansion parameter



G_E^{u-d} , G_M^{u-d} : improvement



G_E^{u-d} : Examined Dipole, Pade and z -expansion fits

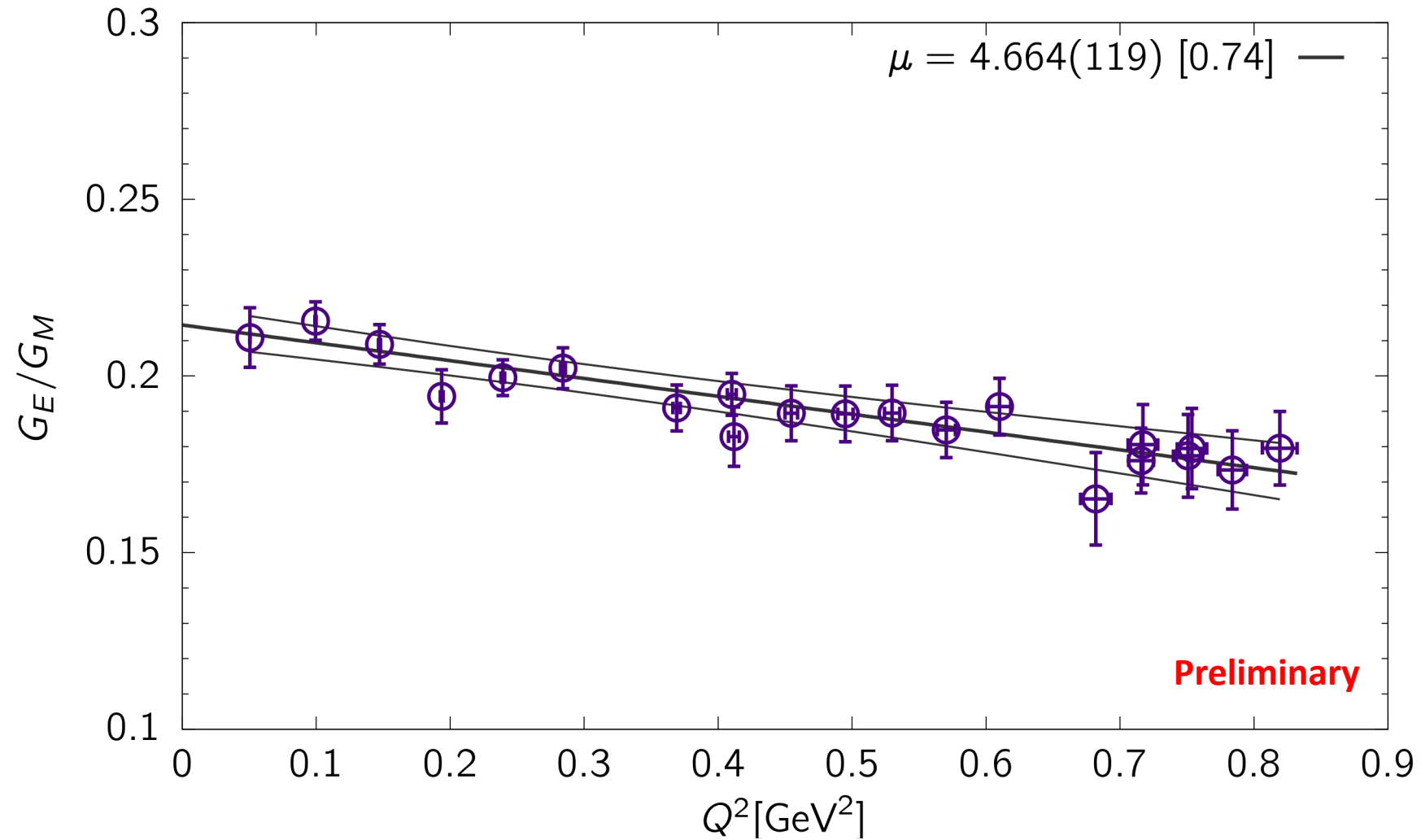


- Dipole: $\frac{G_E(0)/g_V}{(1+Q^2/M_E^2)}$

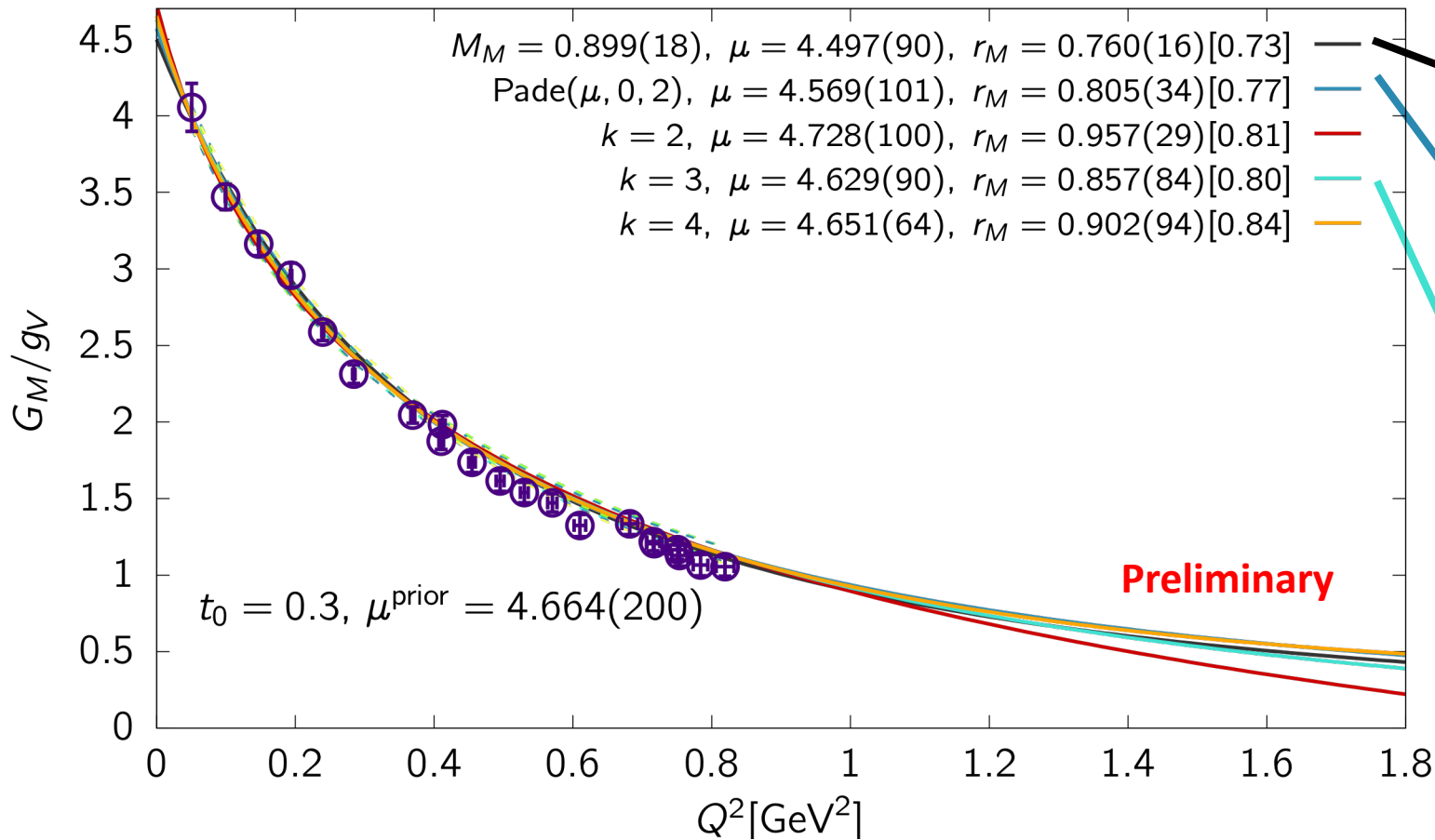
- Pade: $\frac{g}{1+b_1Q^2+b_2Q^4}$

- z^n -
expansion:
 $\sum_{k=0}^n a_k z(Q^2)^k$
Need z^3 to fit the data

G_E/G_M : linear fit



G_M^{u-d} : Examined Dipole, Pade and z -expansion fits

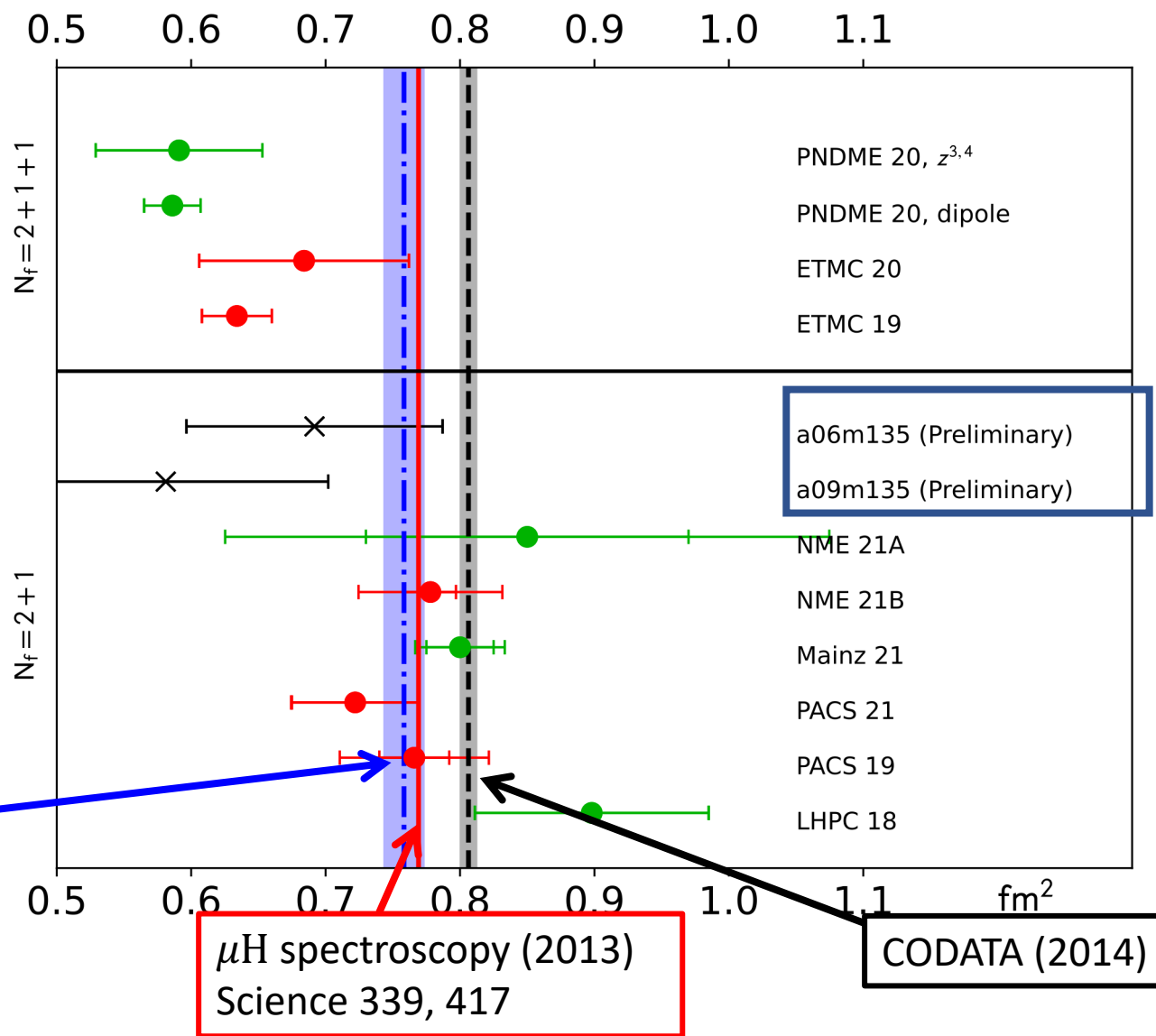


- Dipole: $\frac{G_E(0)/g_V}{(1+Q^2/M_E^2)}$

- Pade: $\frac{g}{1+b_1Q^2+b_2Q^4}$

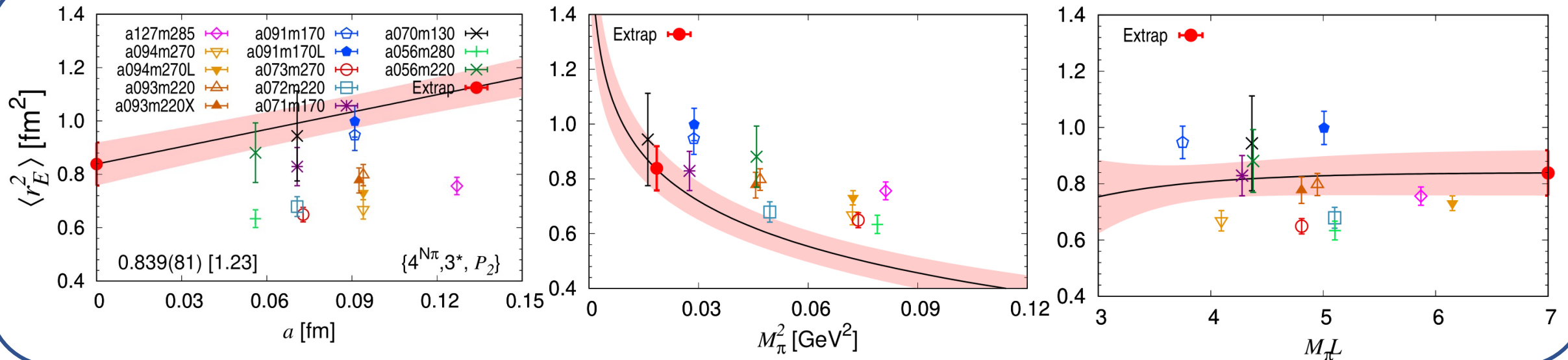
- z^n -
expansion:
 $\sum_{k=0}^n a_k z(Q^2)^k$
Need z^3 to fit the data

$$\langle r_E^2 \rangle^{u-d}$$



$\langle r_E^2 \rangle^{u-d}$: Chiral-Continuum-Finite Volume extrapolation

NME (2022), preliminary



Leading corrections included in the fit ansatz

$$c_1 + c_2 a + c_3 \log \frac{M_\pi^2}{\lambda^2} + c_4 \log \frac{M_\pi^2}{\lambda^2} e^{-M_\pi L}$$

Nucl.Phys.A635, 121 (1998)
 Nucl.Phys.A679, 698 (2001)
 Phys.Rev.D71, 034508 (2005)

Summary

- We are calculating isovector electromagnetic form factors $G_{E,M}^{u-d}(Q^2)$ using high statistics MILC HISQ ensembles at M_{Phys} and $a \approx 0.06$ and 0.09 as part of a comprehensive analysis of nucleon structure.
- Form factors are presented as a function of Q^2 over $0.04 < Q^2 < 0.8 \text{ GeV}^2$.
- Using the high-statistics data, we examined excited state (ES) fit strategies for different M_1 and M_2 based on χ^2 , and the result supports vector meson dominance i.e., the insertion of $2\pi(q)$ by the current.