

# Neural Network Ground States of Lattice Gauge Theories

An introduction and overview

Thomas Spriggs, Lattice 2026, Maryland, USA

[t.e.spriggs@tudelft.nl](mailto:t.e.spriggs@tudelft.nl)    <https://thomas-spriggs-codes.github.io/>

Quick disclaimer:

References are a guide for further reading

Incomplete

# Neural Quantum States (NQS)

Carleo, G., & Troyer, M. (2017). Solving the quantum many-body problem with artificial neural networks. *Science*, 355(6325), 602-606.



# Neural Quantum States (NQS)

- Simulate lattice gauge theories in **real, continuous** time

# Neural Quantum States (NQS)

- Simulate lattice gauge theories in **real, continuous** time
  - Hamiltonian framework

# Neural Quantum States (NQS)

- Simulate lattice gauge theories in **real, continuous** time
  - Hamiltonian framework
- Don't sample from the action



# Neural Quantum States (NQS)

- Simulate lattice gauge theories in **real, continuous** time
  - Hamiltonian framework
- Don't sample from the action
  - Learn the wavefunction

# Neural Quantum States (NQS)

- Simulate lattice gauge theories in **real, continuous** time
  - Hamiltonian framework
- Don't sample from the action
  - Learn the wavefunction
  - No sign problem\*



# Neural Quantum States (NQS)

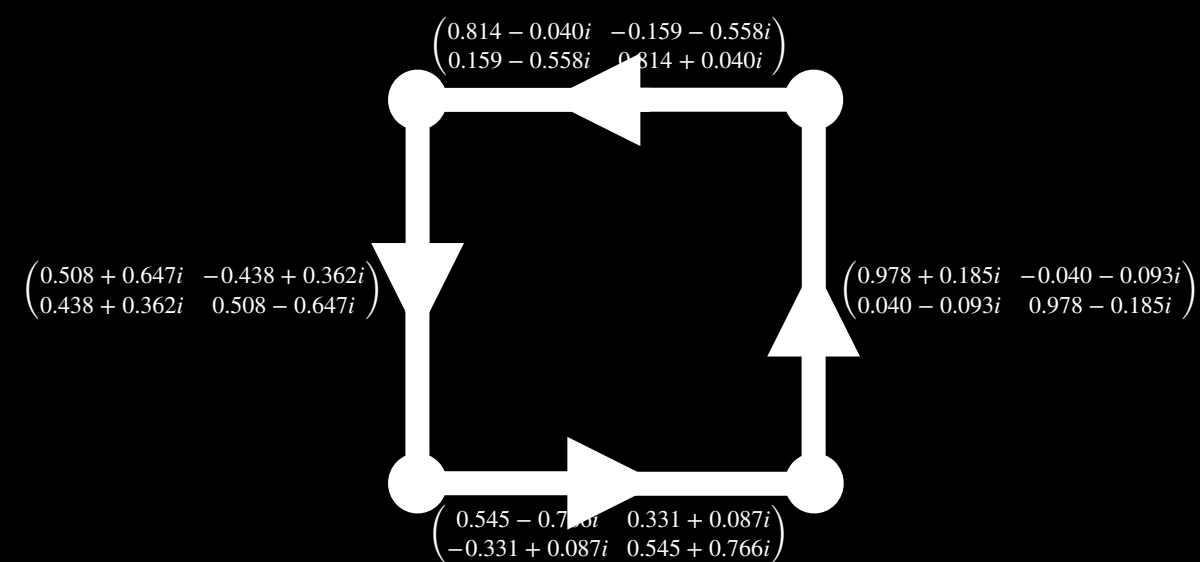
- Simulate lattice gauge theories in **real, continuous** time
  - Hamiltonian framework
- Don't sample from the action
  - Learn the wavefunction
  - No sign problem\*
- Already done for:  $\mathbf{Z}_2$ ,  $U(1)$ ,  $SU(2)$  lattice gauge theories

# Neural Quantum States (NQS)

- Simulate lattice gauge theories in **real, continuous** time
  - Hamiltonian framework
- Don't sample from the action
  - Learn the wavefunction
  - No sign problem\*
- Already done for:  $\mathbf{Z}_2$ ,  $U(1)$ ,  $SU(2)$  lattice gauge theories
  - Commonplace in condensed matter physics and quantum chemistry

Carleo, G., & Troyer, M. (2017). Solving the quantum many-body problem with artificial neural networks. *Science*, 355(6325), 602-606.

# NQS workflow for LGTs



$$\text{NN} \\ |\Psi_{\theta}\rangle$$



# NQS workflow for LGTs

NN  
 $|\Psi_\theta\rangle$

$$|\langle U_0 | \Psi_\theta \rangle|^2$$

# NQS workflow for LGTs

NN  
 $|\Psi_\theta\rangle$

$$|\langle U_0 | \Psi_\theta \rangle|^2 \quad |\langle U_1 | \Psi_\theta \rangle|^2$$

# NQS workflow for LGTs

NN  
 $|\Psi_\theta\rangle$

Accept/reject

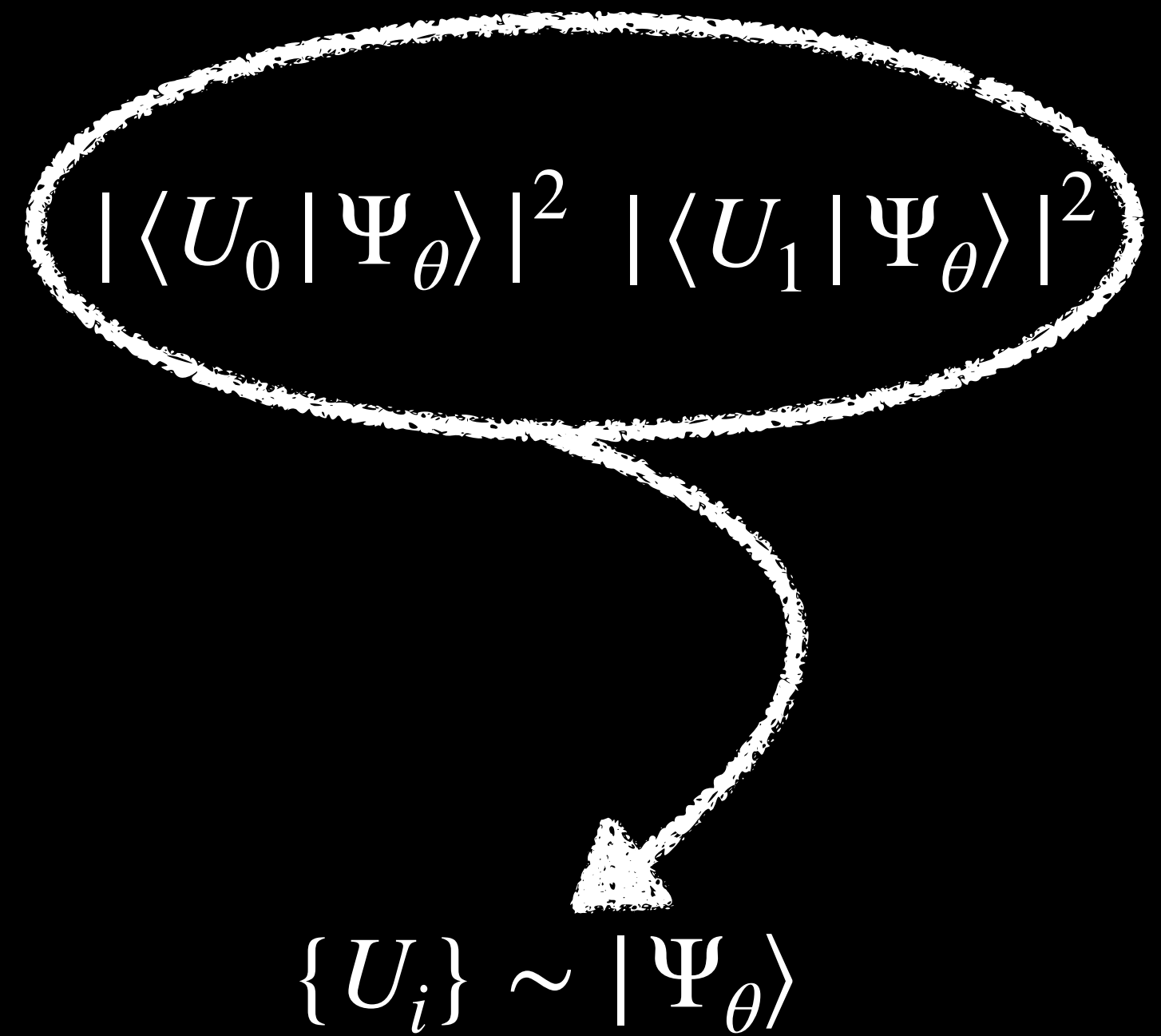
$$|\langle U_0 | \Psi_\theta \rangle|^2 \quad |\langle U_1 | \Psi_\theta \rangle|^2$$



# NQS workflow for LGTs

NN  
 $|\Psi_\theta\rangle$

Accept/reject



# NQS workflow for LGTs

NN  
 $|\Psi_\theta\rangle$

Accept/reject

$$|\langle U_0 | \Psi_\theta \rangle|^2 \quad |\langle U_1 | \Psi_\theta \rangle|^2$$

$$\{U_i\} \sim |\Psi_\theta\rangle$$

$$E \approx \frac{1}{N} \sum_{U_i \sim |\langle \Psi_\theta | U \rangle|^2} \mathcal{H}_L(U_i)$$



# NQS workflow for LGTs

NN  
 $|\Psi_\theta\rangle$

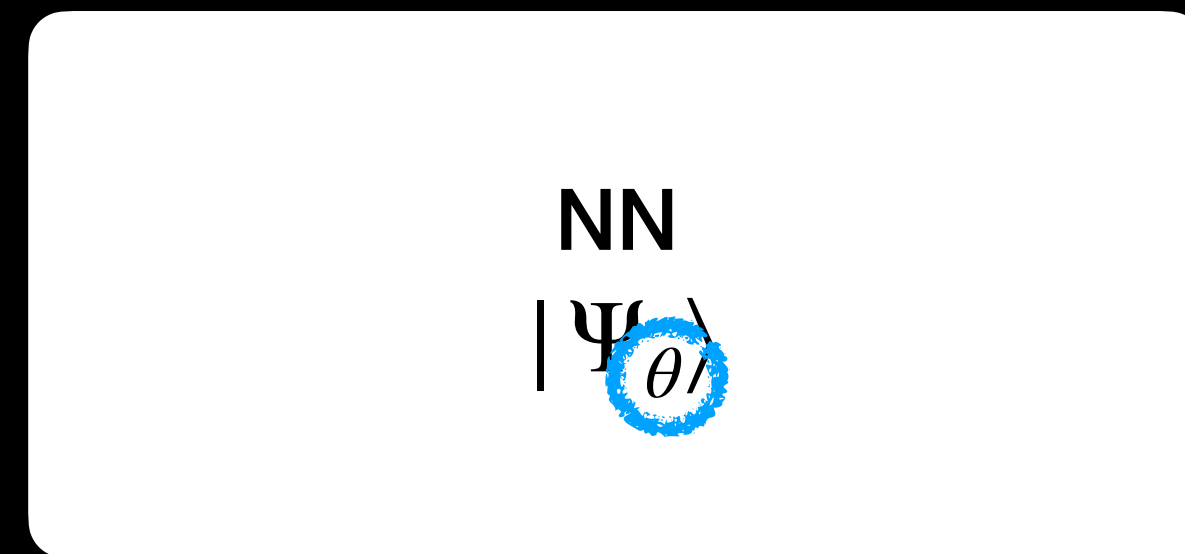
Accept/reject

$$|\langle U_0 | \Psi_\theta \rangle|^2 \quad |\langle U_1 | \Psi_\theta \rangle|^2$$

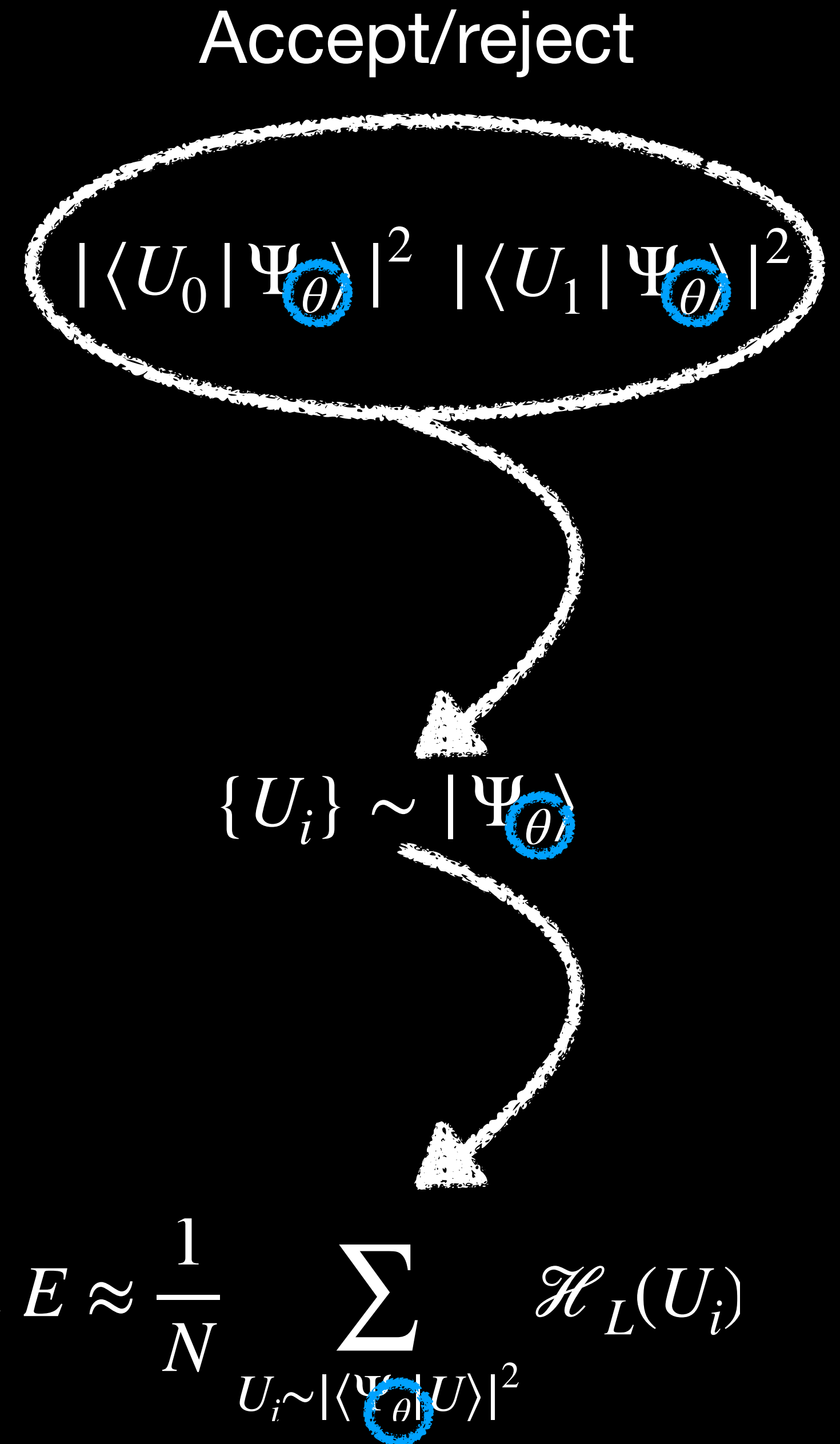
$$\{U_i\} \sim |\Psi_\theta\rangle$$

$$E \approx \frac{1}{N} \sum_{U_i \sim |\langle \Psi_\theta | U \rangle|^2} \mathcal{H}_L(U_i)$$

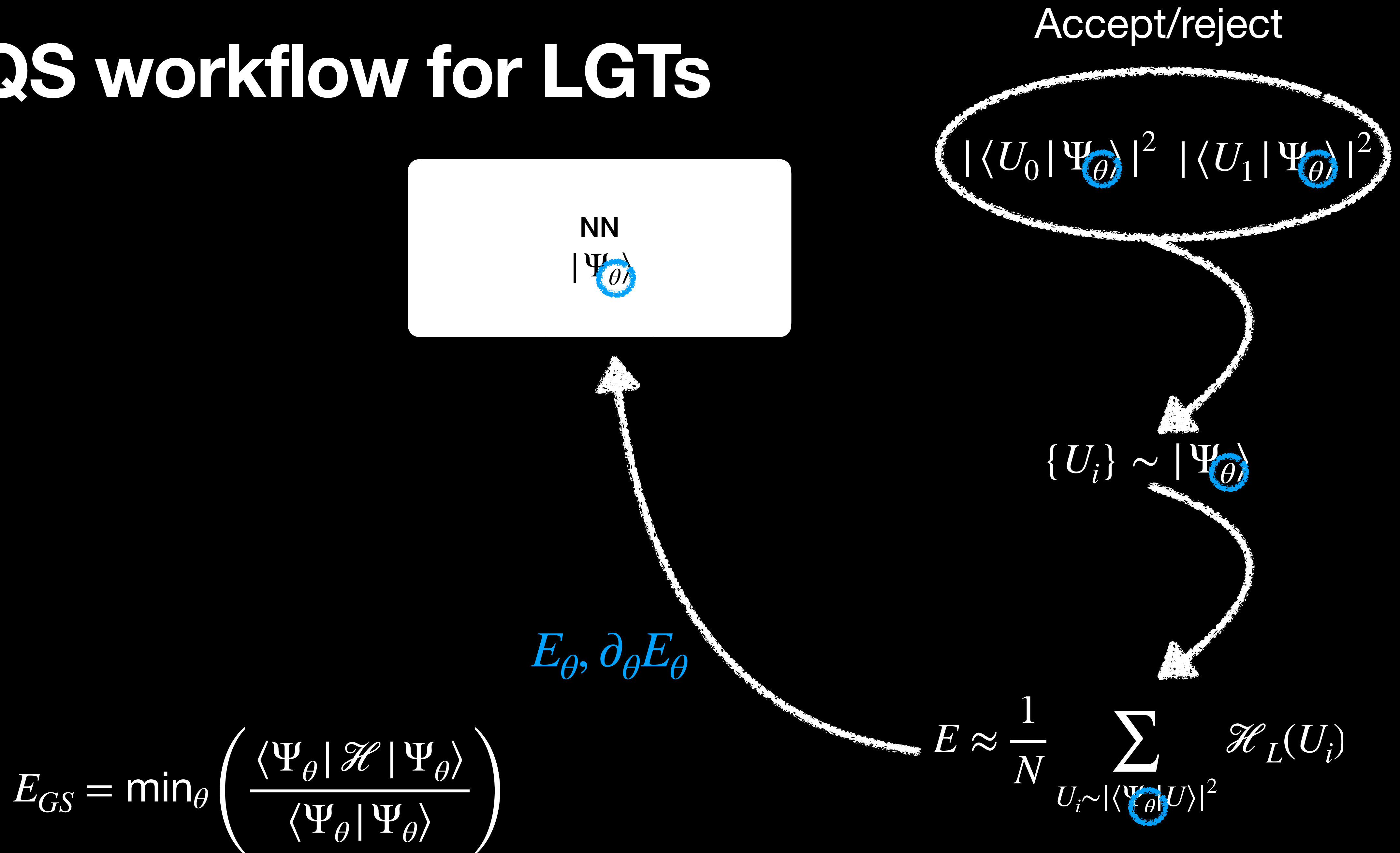
# NQS workflow for LGTs



$$E_\theta, \partial_\theta E_\theta$$



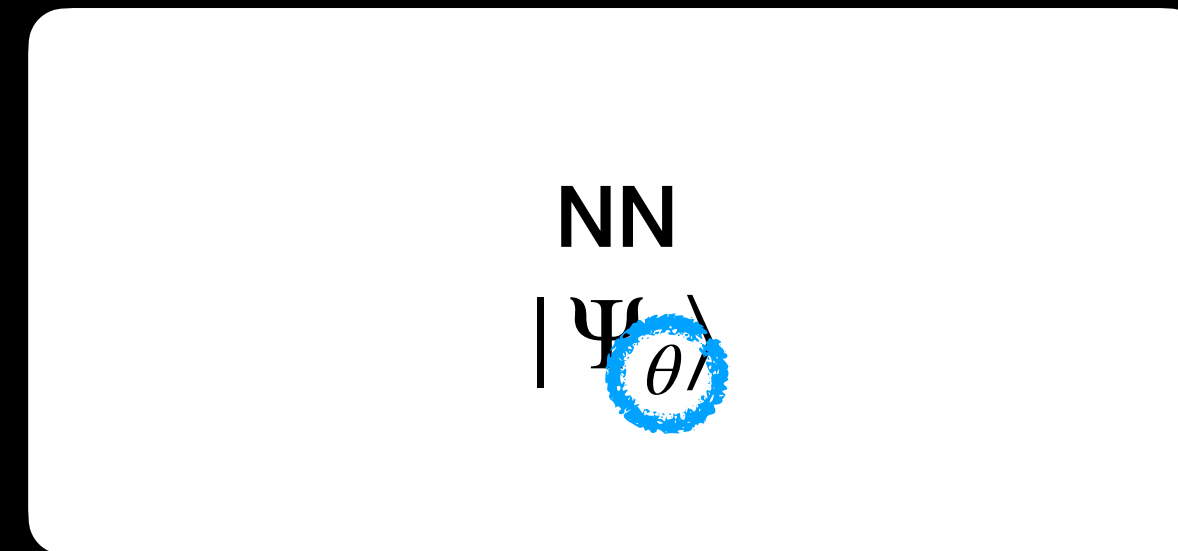
# NQS workflow for LGTs





# NQS workflow for LGTs

Accept/reject



$$|\langle U_0 | \Psi_\theta \rangle|^2 |\langle U_1 | \Psi_\theta \rangle|^2$$

$$\{U_i\} \sim |\Psi_\theta\rangle$$

Task: find function that returns ground state wavefunction amplitudes

$$E_\theta, \partial_\theta E_\theta$$

$$E_{GS} = \min_\theta \left( \frac{\langle \Psi_\theta | \mathcal{H} | \Psi_\theta \rangle}{\langle \Psi_\theta | \Psi_\theta \rangle} \right)$$

$$E \approx \frac{1}{N} \sum_{U_i \sim |\langle \Psi_\theta | U \rangle|^2} \mathcal{H}_L(U_i)$$

# Lagrangians and Hamiltonians

# Lagrangian vs Hamiltonian formalism

From an NQS perspective



# Lagrangian vs Hamiltonian formalism

From an NQS perspective

- Continuum Dirac equation
  - Observables found through the path integral

# Lagrangian vs Hamiltonian formalism

From an NQS perspective

- Continuum Dirac equation
  - Observables found through the path integral
- Discretise space: regularising theory



# Lagrangian vs Hamiltonian formalism

From an NQS perspective

- Continuum Dirac equation
  - Observables found through the path integral
- Discretise space: regularising theory
- Discretise time and Wick rotate ( $t \rightarrow i\tau$ )

# Lagrangian vs Hamiltonian formalism

From an NQS perspective

- Continuum Dirac equation
  - Observables found through the path integral
- Discretise space: regularising theory
- Discretise time and Wick rotate ( $t \rightarrow i\tau$ )

$$\int dx O(x) p(x) \approx \frac{1}{N} \sum_{x \sim p(x)} O(x)$$

# Lagrangian vs Hamiltonian formalism

## From an NQS perspective

- Continuum Dirac equation
  - Observables found through the path integral
- Discretise space: regularising theory
- Discretise time and Wick rotate ( $t \rightarrow i\tau$ )

$$\int dx O(x) p(x) \approx \frac{1}{N} \sum_{x \sim p(x)} O(x)$$

$$\int D[U] O[U] e^{S_E[U]} \approx \frac{1}{N} \sum_{U \sim e^{S_E[U]}} O[U]$$



# Lagrangian vs Hamiltonian formalism

From an NQS perspective

# Lagrangian vs Hamiltonian formalism

From an NQS perspective

- Continuum Dirac equation
  - Observables found through the path integral

# Lagrangian vs Hamiltonian formalism

From an NQS perspective

- Continuum Dirac equation
  - Observables found through the path integral
- Discretise space: regularising theory



# Lagrangian vs Hamiltonian formalism

From an NQS perspective

- Continuum Dirac equation
  - Observables found through the path integral
- Discretise space: regularising theory
- ~~Discretise time and Wick rotate ( $t \rightarrow i\tau$ )~~

# Lagrangian vs Hamiltonian formalism

From an NQS perspective

- Continuum Dirac equation
  - Observables found through the path integral
- Discretise space: regularising theory
- ~~Discretise time and Wick rotate ( $t \rightarrow i\tau$ )~~

$$\int dx O(x) p(x) \approx \frac{1}{N} \sum_{x \sim p(x)} O(x)$$



# Lagrangian vs Hamiltonian formalism

From an NQS perspective

- Continuum Dirac equation
  - Observables found through the path integral
- Discretise space: regularising theory
- ~~Discretise time and Wick rotate ( $t \rightarrow i\tau$ )~~

$$\int dx O(x) p(x) \approx \frac{1}{N} \sum_{x \sim p(x)} O(x)$$

# Lagrangian vs Hamiltonian formalism

Trade off

Lose access to real time dynamics

Get a known expression for the probability distribution of gauge fields

Complex actions ruin this:  $e^{S_E}$  is no longer a probability

# Lagrangian vs Hamiltonian formalism

Trade off

Keep access to real time dynamics

Have to learn the probability distribution (using a neural network)

Probability distribution is always real



# Neural (network) wavefunctions

A quick aside into spin systems (easier to picture)

# Some definitions: wavefunction

# Some definitions: wavefunction

- Let  $|\Psi\rangle = \sum_i a_i |x_i\rangle$



# Some definitions: wavefunction

- Let  $|\Psi\rangle = \sum_i a_i |x_i\rangle$ 
  - 2 qubits, spin basis:  $|\Psi\rangle = a_0 |00\rangle + a_1 |01\rangle + a_2 |10\rangle + a_3 |11\rangle$

# Some definitions: wavefunction

- Let  $|\Psi\rangle = \sum_i a_i |x_i\rangle$ 
  - 2 qubits, spin basis:  $|\Psi\rangle = a_0 |00\rangle + a_1 |01\rangle + a_2 |10\rangle + a_3 |11\rangle$
  - “Finding ground state” == finding four\* numbers



# Some definitions: **variational** wavefunction

- Let  $|\Psi\rangle = \sum_i a_i |x_i\rangle$ 
  - 2 qubits, spin basis:  $|\Psi\rangle = a_0 |00\rangle + a_1 |01\rangle + a_2 |10\rangle + a_3 |11\rangle$
  - “Finding ground state” == finding **a function that gives** four\* numbers

$$|\Psi_\theta\rangle = \sum_i f_\theta(x_i) |x_i\rangle$$

# Some definitions: **variational** wavefunction

$$|\Psi_{\theta}\rangle = \sum_i f_{\theta}(x_i) |x_i\rangle$$

# Some definitions: **variational** wavefunction

$$|\Psi_\theta\rangle = \sum_i f_\theta(x_i) |x_i\rangle$$

$$\langle x_0 | \Psi_\theta \rangle = \sum_i \langle x_0 | f_\theta(x_i) | x_i \rangle = f_\theta(x_0)$$

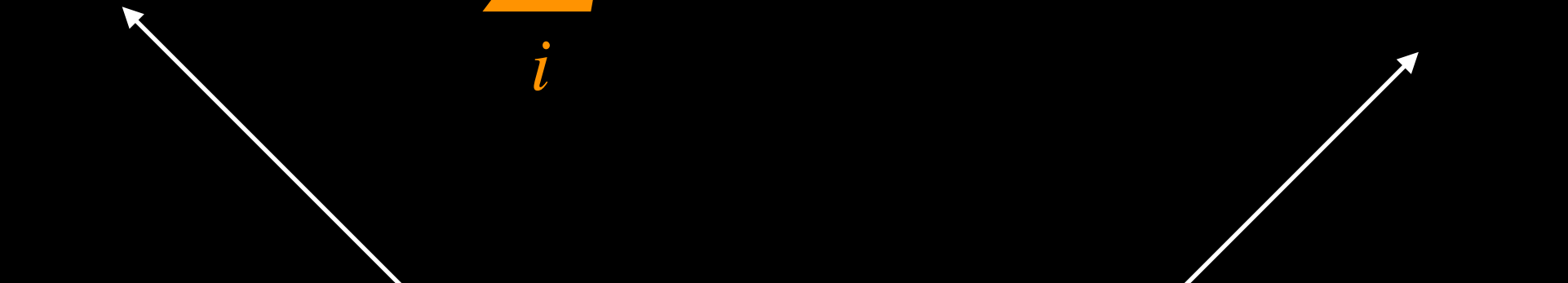


# Some definitions: **variational** wavefunction

$$|\Psi_\theta\rangle = \sum_i f_\theta(x_i) |x_i\rangle$$

$$\langle x_0 | \Psi_\theta \rangle = \sum_i \langle x_0 | f_\theta(x_i) | x_i \rangle = f_\theta(x_0)$$

“Configuration in, amplitude out”

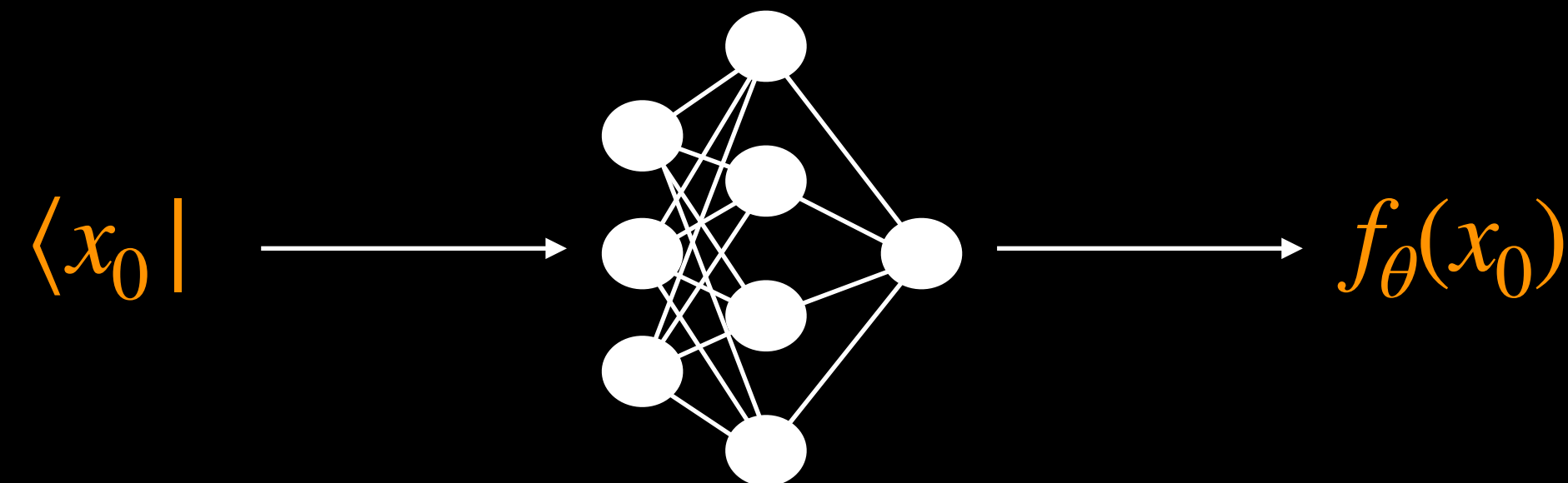
Two white arrows originate from the text. One arrow points from the word 'Configuration' to the bra index  $x_0$  in the equation. The other arrow points from the word 'amplitude' to the ket index  $x_i$  in the equation.

# Some definitions: **variational** wavefunction

$$|\Psi_\theta\rangle = \sum_i f_\theta(x_i) |x_i\rangle$$

$$\langle x_0 | \Psi_\theta \rangle = \sum_i \langle x_0 | f_\theta(x_i) | x_i \rangle = f_\theta(x_0)$$

“Configuration in, amplitude out”



# Wavefunctions in gauge theories



# Finding the probability measure

Take an un-normalised variational state,  $|\Psi_\theta\rangle$ , a Hilbert space,  $\{|\vec{x}\rangle\}$ , and a Hamiltonian,  $\mathcal{H}$

$$E = \langle \mathcal{H} \rangle_\theta = \frac{\langle \Psi_\theta | \mathcal{H} | \Psi_\theta \rangle}{\langle \Psi_\theta | \Psi_\theta \rangle}$$

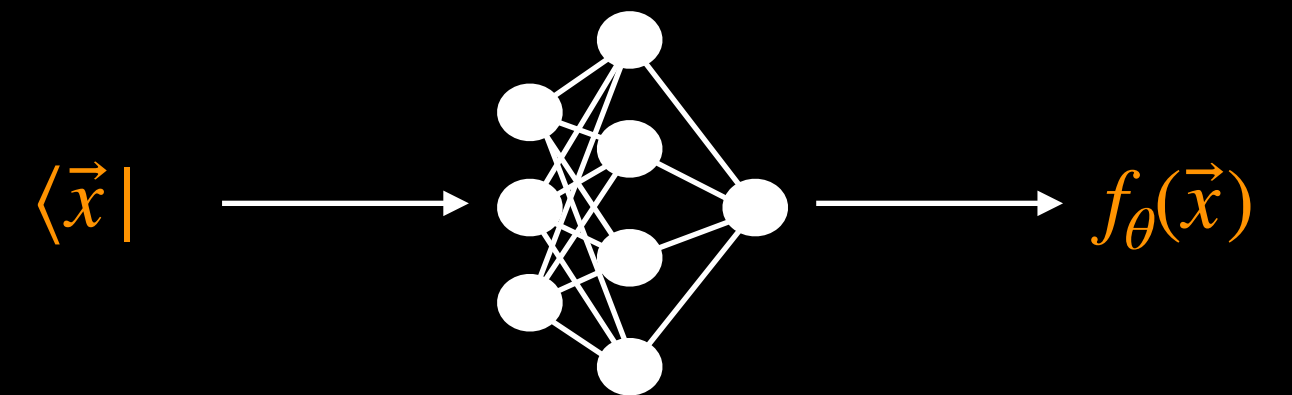
Introduce a complete set of states  $\sum_{\vec{x}} |\vec{x}\rangle \langle \vec{x}|$

$$E = \frac{\sum_{\vec{x}} \langle \Psi_\theta | \vec{x} \rangle \langle \vec{x} | \mathcal{H} | \Psi_\theta \rangle}{\sum_{\vec{x}'} \langle \Psi_\theta | \vec{x}' \rangle \langle \vec{x}' | \Psi_\theta \rangle}$$

Multiply top and bottom by a factor of  $\langle \vec{x} | \Psi_\theta \rangle$

$$E = \frac{\sum_{\vec{x}} |\langle \Psi_\theta | \vec{x} \rangle|^2 \langle \vec{x} | \mathcal{H} | \Psi_\theta \rangle}{\sum_{\vec{x}'} |\langle \Psi_\theta | \vec{x}' \rangle|^2 \langle \vec{x}' | \Psi_\theta \rangle} = \frac{\sum_{\vec{x}} p(\vec{x}) \mathcal{H}_L(\vec{x})}{\sum_{\vec{x}'} p(\vec{x}')} \approx \frac{1}{N} \sum_{\vec{x} \sim |\langle \Psi_\theta | \vec{x} \rangle|^2} \mathcal{H}_L(x)$$

# Finding the probability measure



- We can use importance sampling

- Draw samples according to

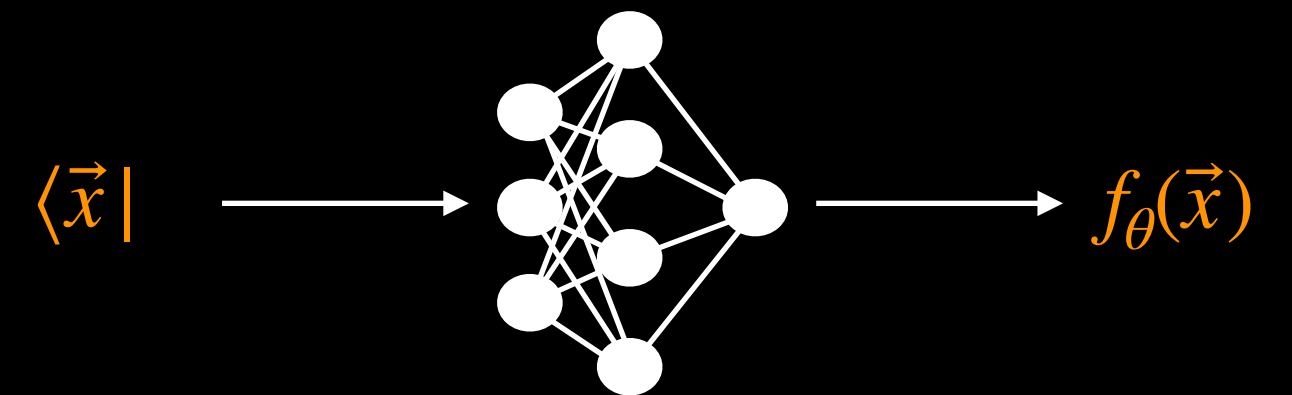
$$| \langle \Psi_{\theta} | \vec{x} \rangle |^2$$

- Measure **local** value of the observable on **each sample**

$$\mathcal{H}_L = \frac{\langle \vec{x} | \mathcal{H} | \Psi_{\theta} \rangle}{\langle \vec{x} | \Psi_{\theta} \rangle}$$



# Finding the probability measure



- We can use importance sampling

- Draw samples according to

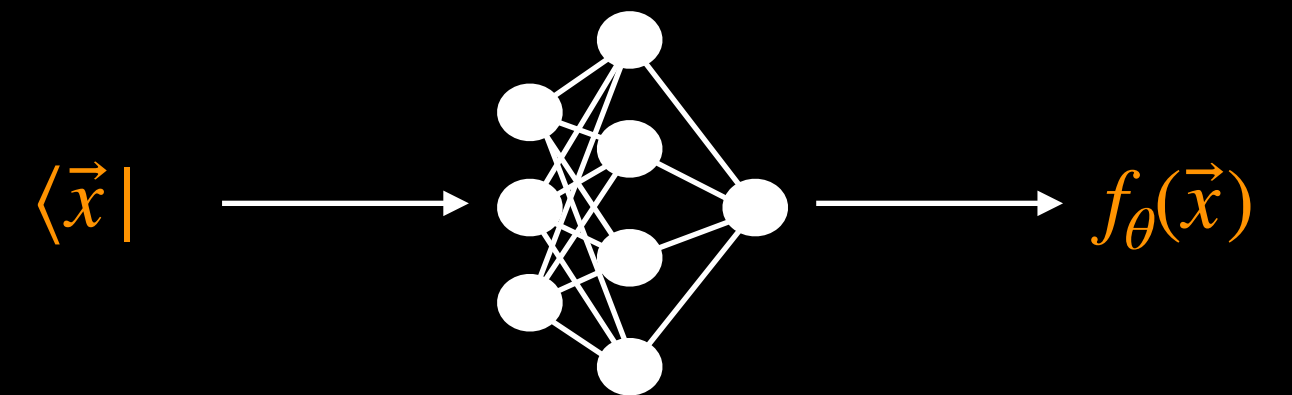
$$|\langle \Psi_{\theta} | \vec{x} \rangle|^2$$

**Always real; no sign problem**

- Measure **local** value of the observable on **each sample**

$$\mathcal{H}_L = \frac{\langle \vec{x} | \mathcal{H} | \Psi_{\theta} \rangle}{\langle \vec{x} | \Psi_{\theta} \rangle}$$

# Finding the probability measure



- We can use importance sampling
  - Draw samples according to

$$|\langle \Psi_{\theta} | \vec{x} \rangle|^2$$

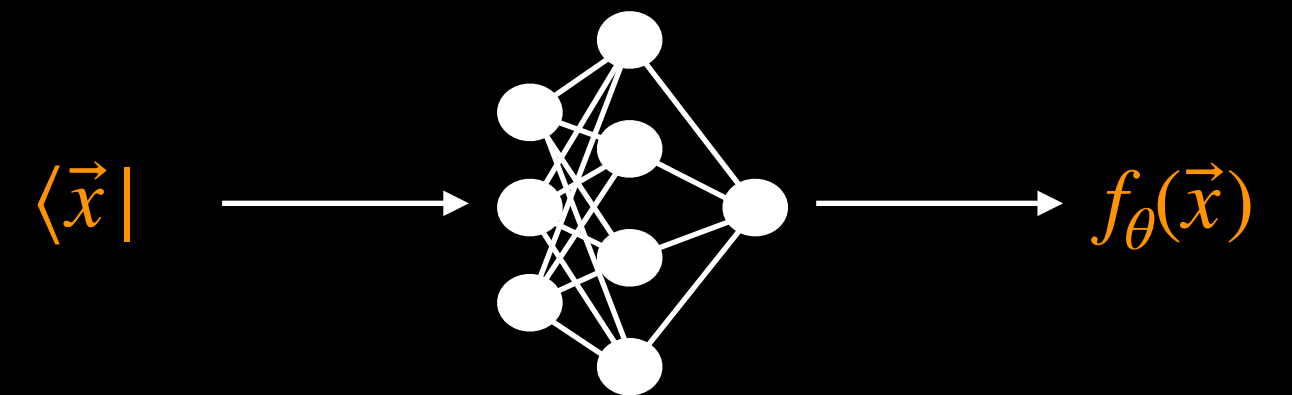
Always real; no sign problem

We do not know this distribution *a priori*

- Measure **local** value of the observable on **each sample**

$$\mathcal{H}_L = \frac{\langle \vec{x} | \mathcal{H} | \Psi_{\theta} \rangle}{\langle \vec{x} | \Psi_{\theta} \rangle}$$

# Finding the probability measure



- We can use importance sampling
  - Draw samples according to

$$|\langle \Psi_{\theta} | \vec{x} \rangle|^2$$

Always real; no sign problem

We do not know this distribution *a priori*

- Measure **local** value of the observable on **each sample**

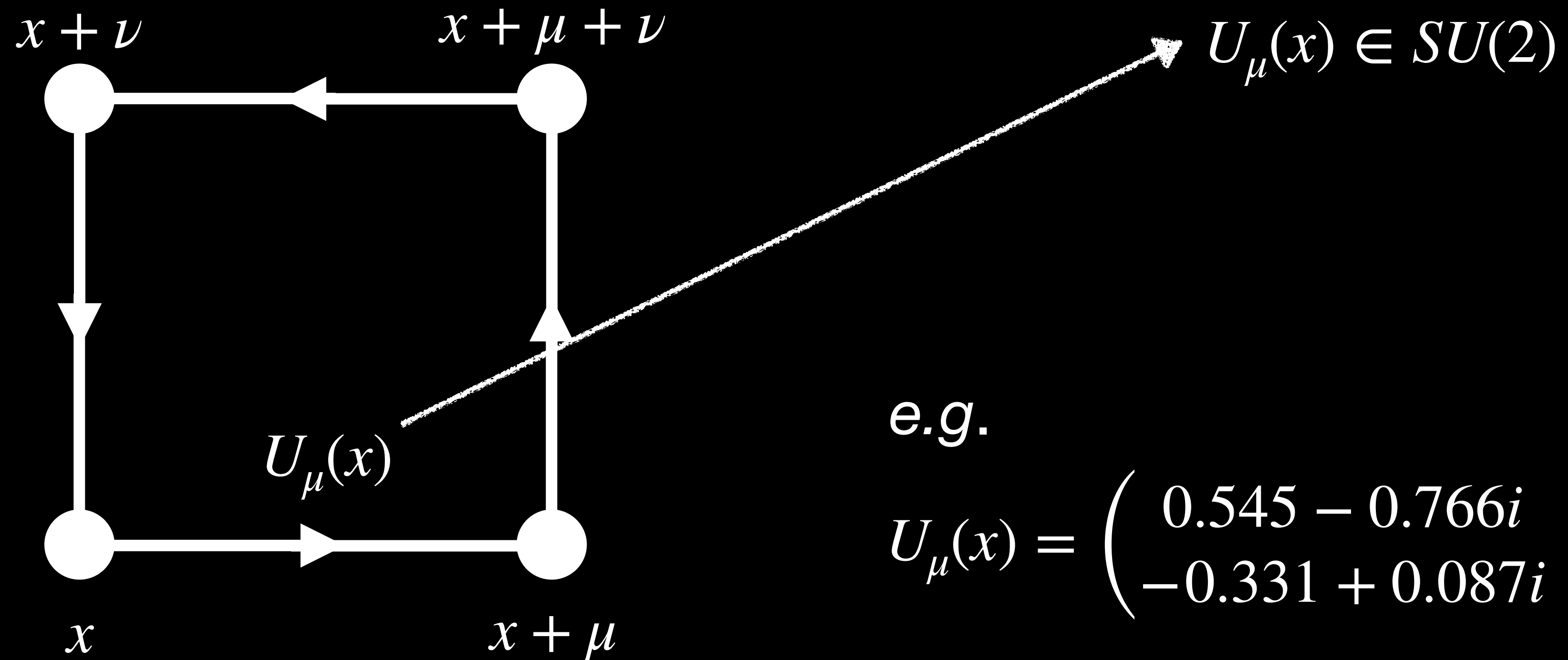
$$\mathcal{H}_L = \frac{\langle \vec{x} | \mathcal{H} | \Psi_{\theta} \rangle}{\langle \vec{x} | \Psi_{\theta} \rangle}$$

$$\text{Energy} \approx \frac{1}{N} \sum_{\vec{x} \sim |\langle \Psi_{\theta} | \vec{x} \rangle|^2} \mathcal{H}_L(x)$$



# Gauge-theory-aware neural networks

# SU(2) example

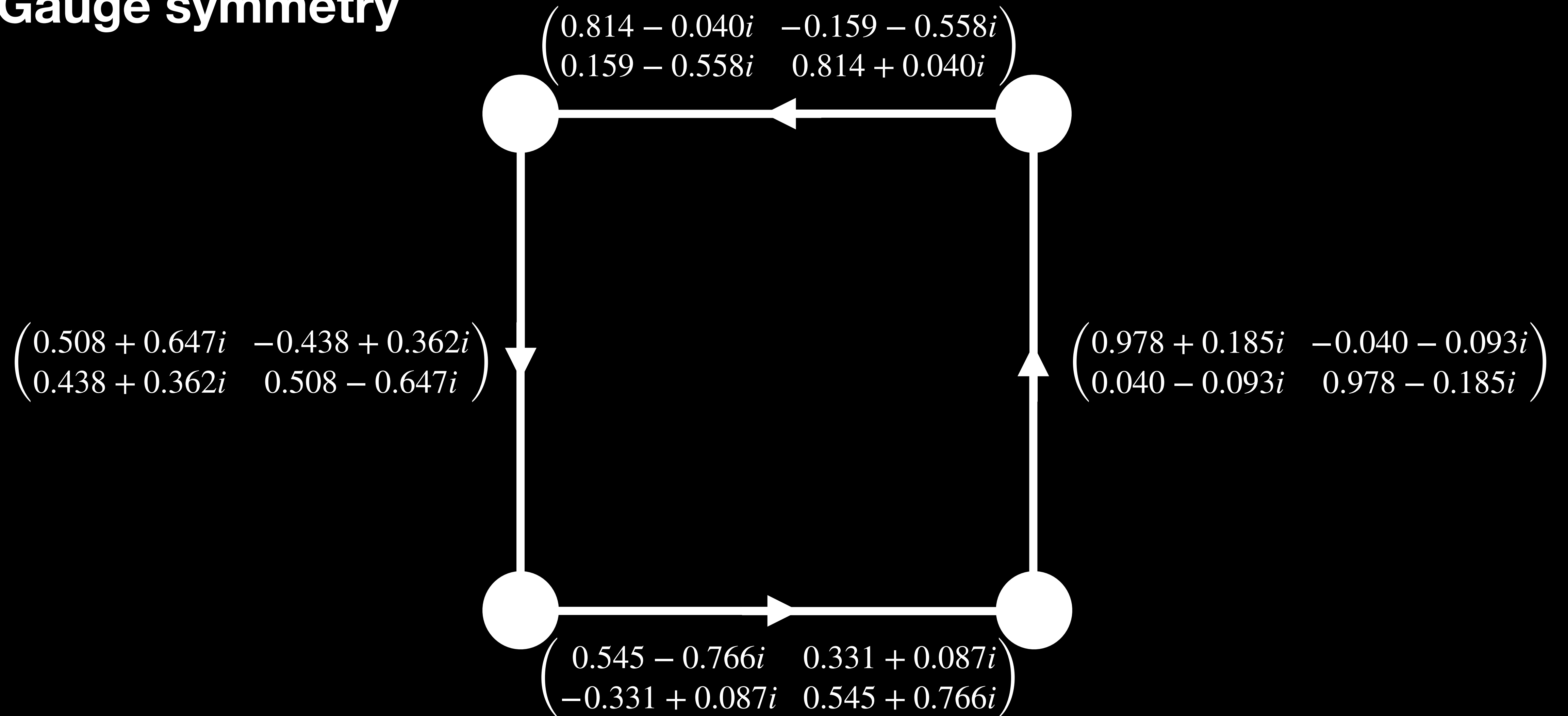


e.g.

$$U_\mu(x) = \begin{pmatrix} 0.545 - 0.766i & 0.331 + 0.087i \\ -0.331 + 0.087i & 0.545 + 0.766i \end{pmatrix}$$

# SU(2) example

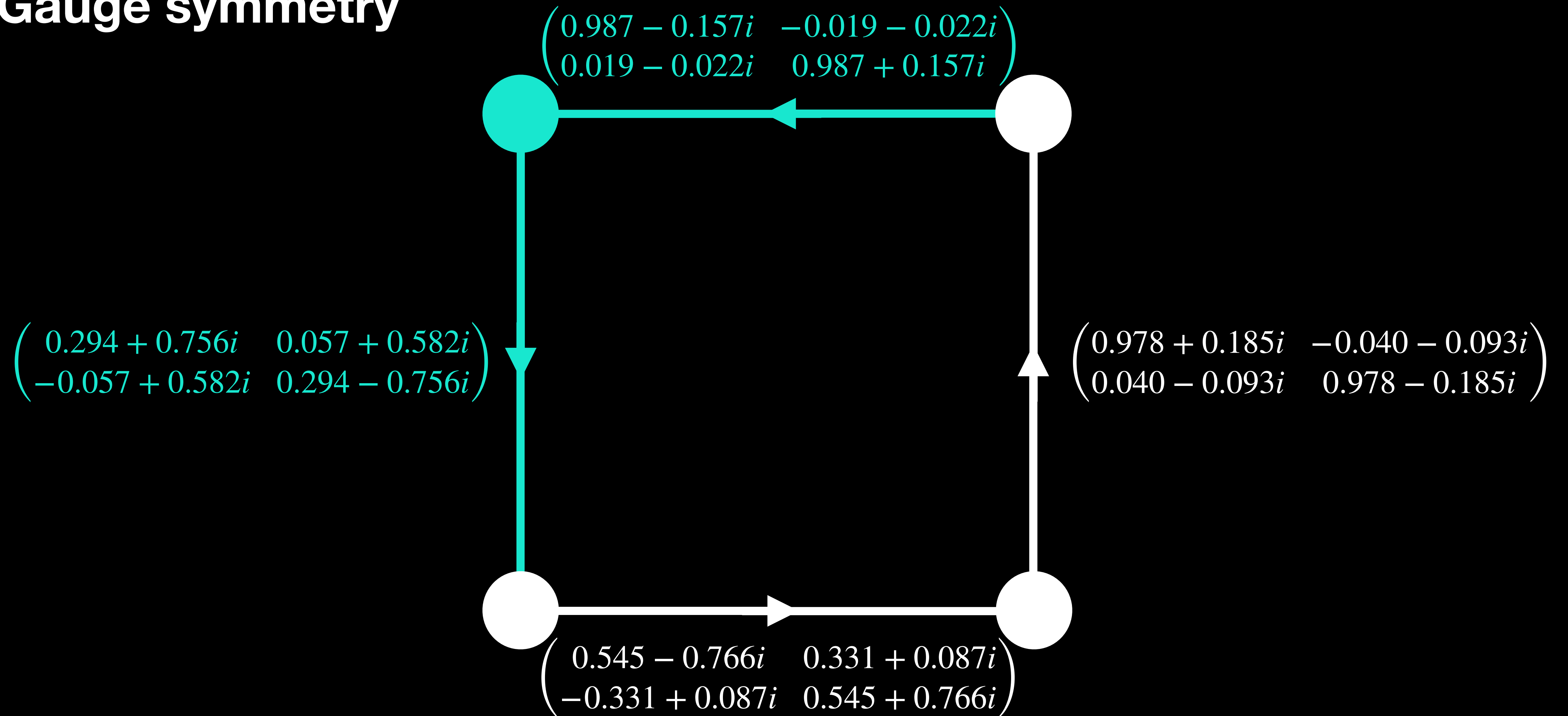
## Gauge symmetry





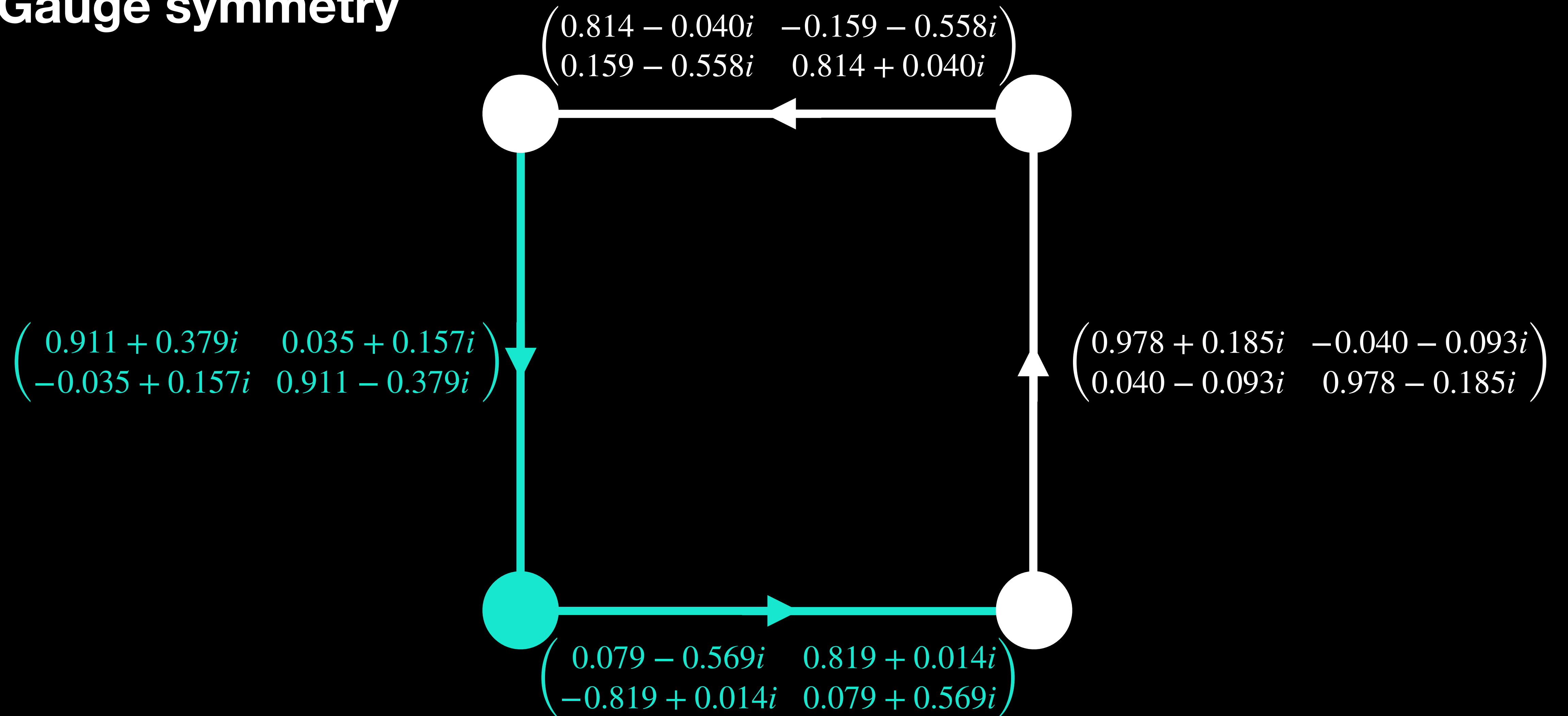
# SU(2) example

Gauge symmetry



# SU(2) example

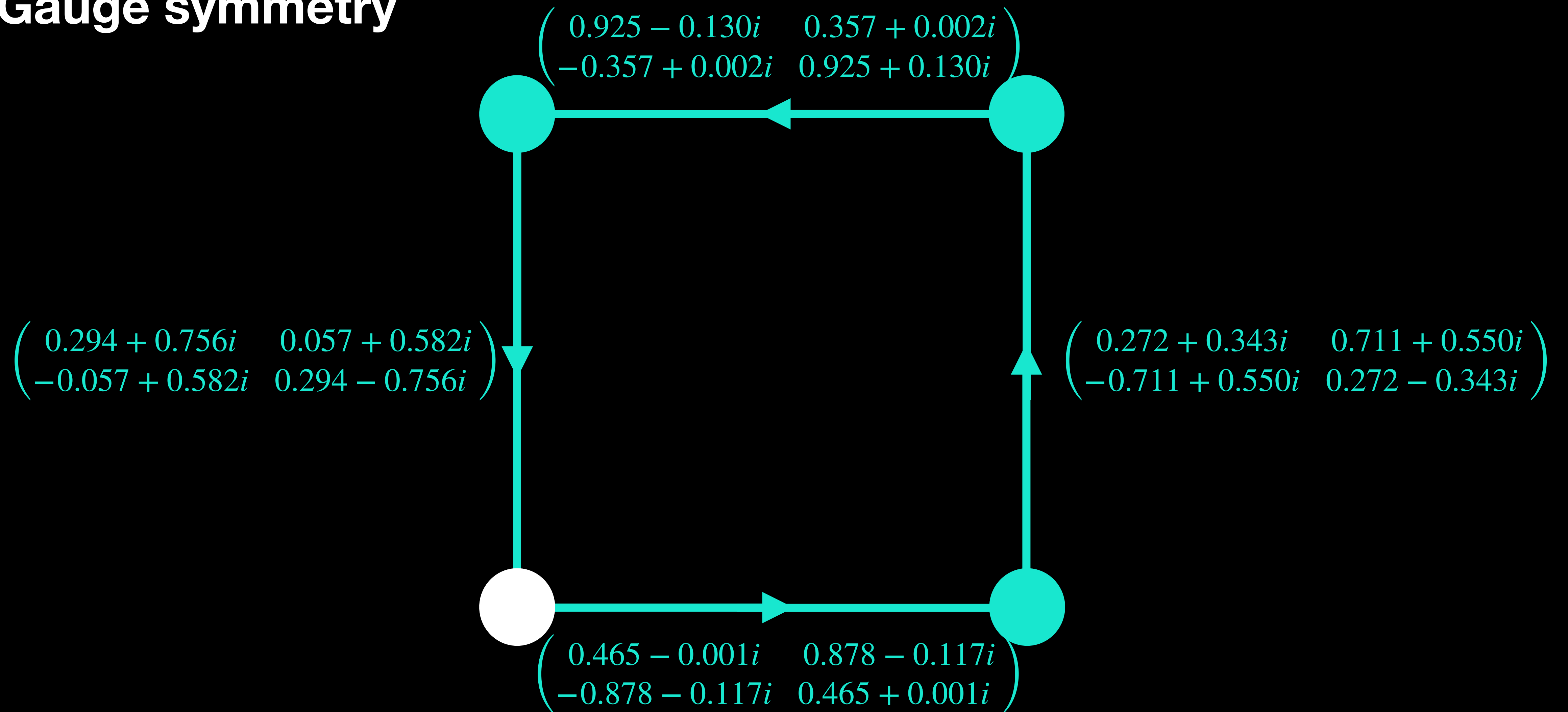
Gauge symmetry





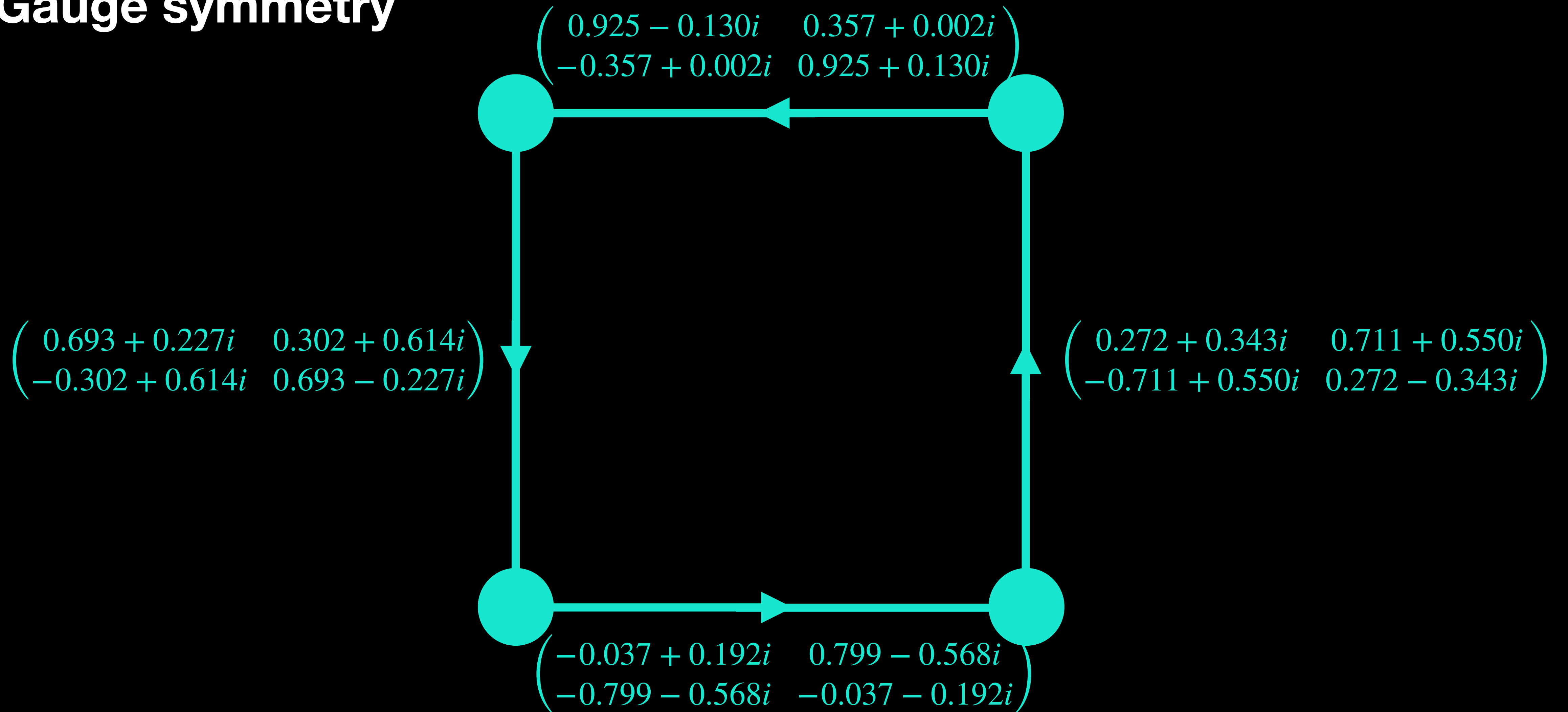
# SU(2) example

Gauge symmetry



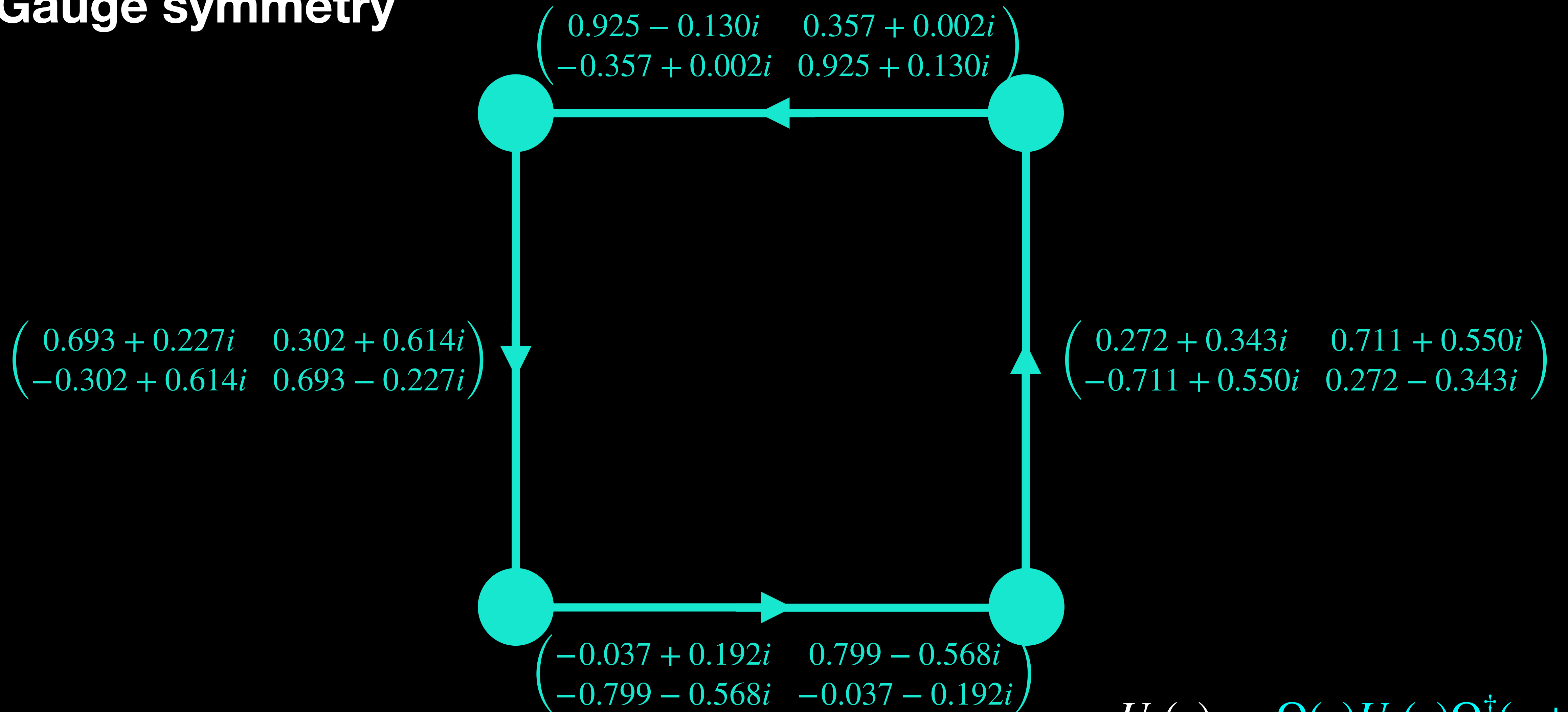
# SU(2) example

Gauge symmetry



# SU(2) example

## Gauge symmetry

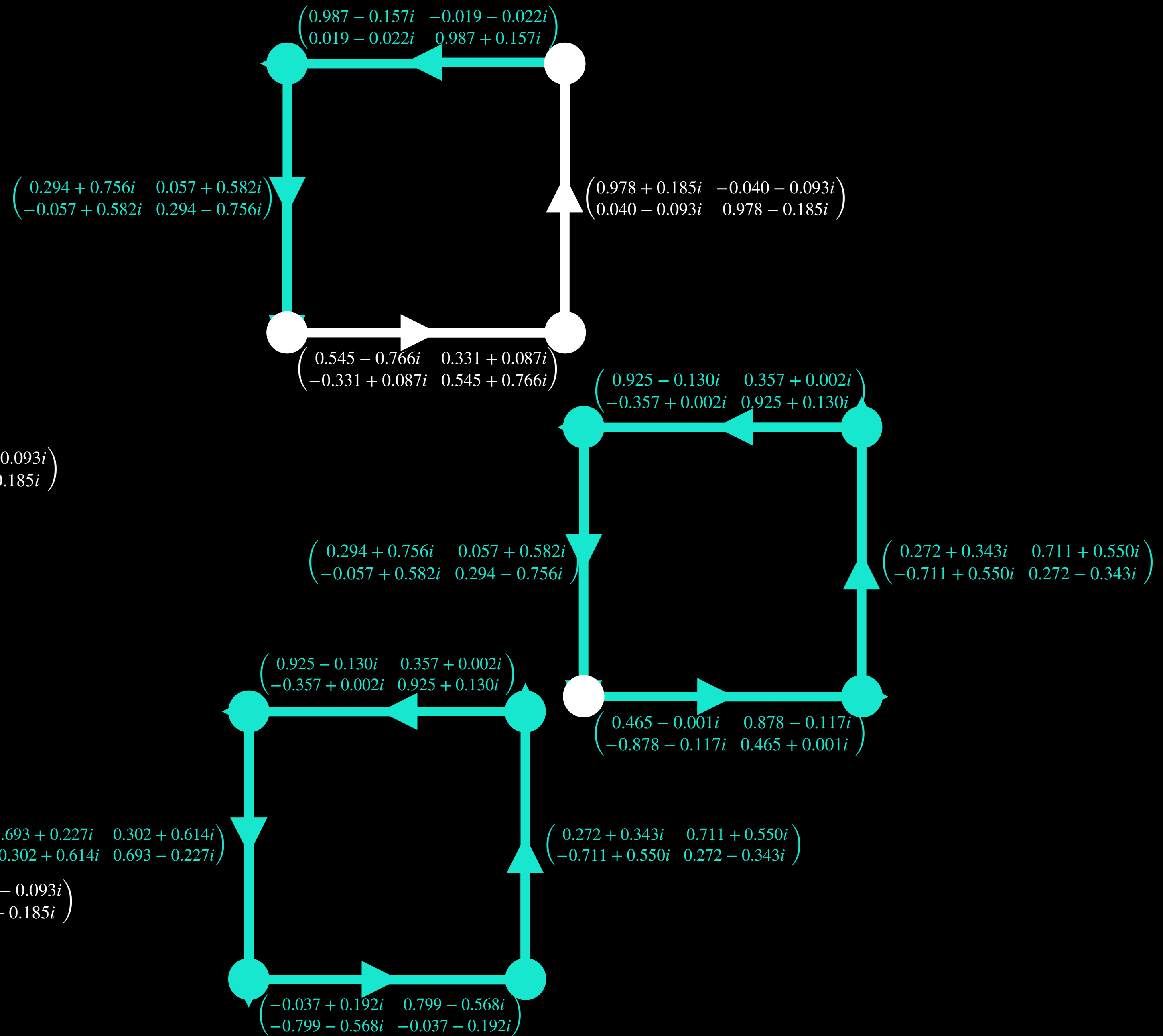
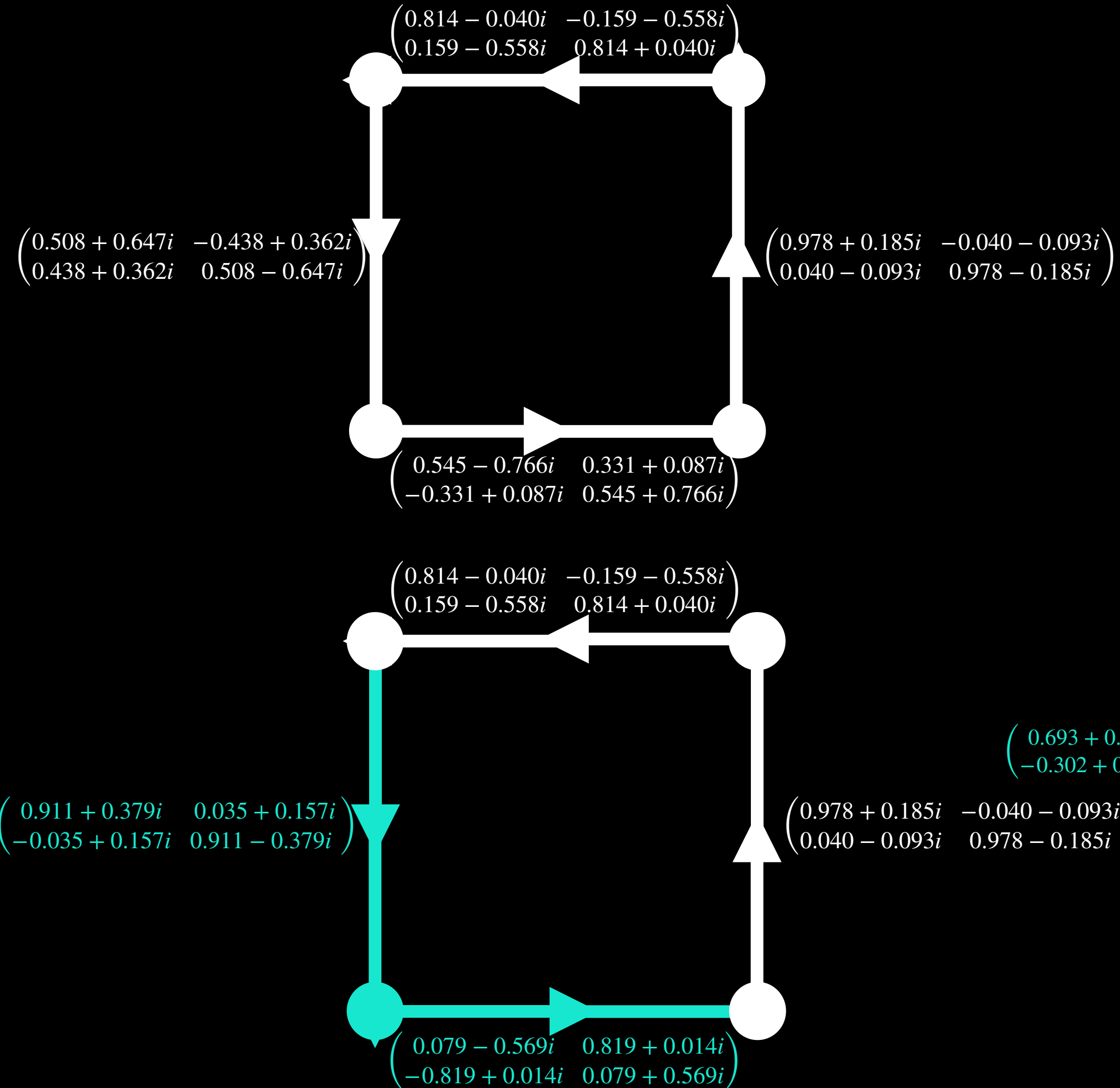


$$U_\mu(x) \rightarrow \Omega(x) U_\mu(x) \Omega^\dagger(x + \mu)$$



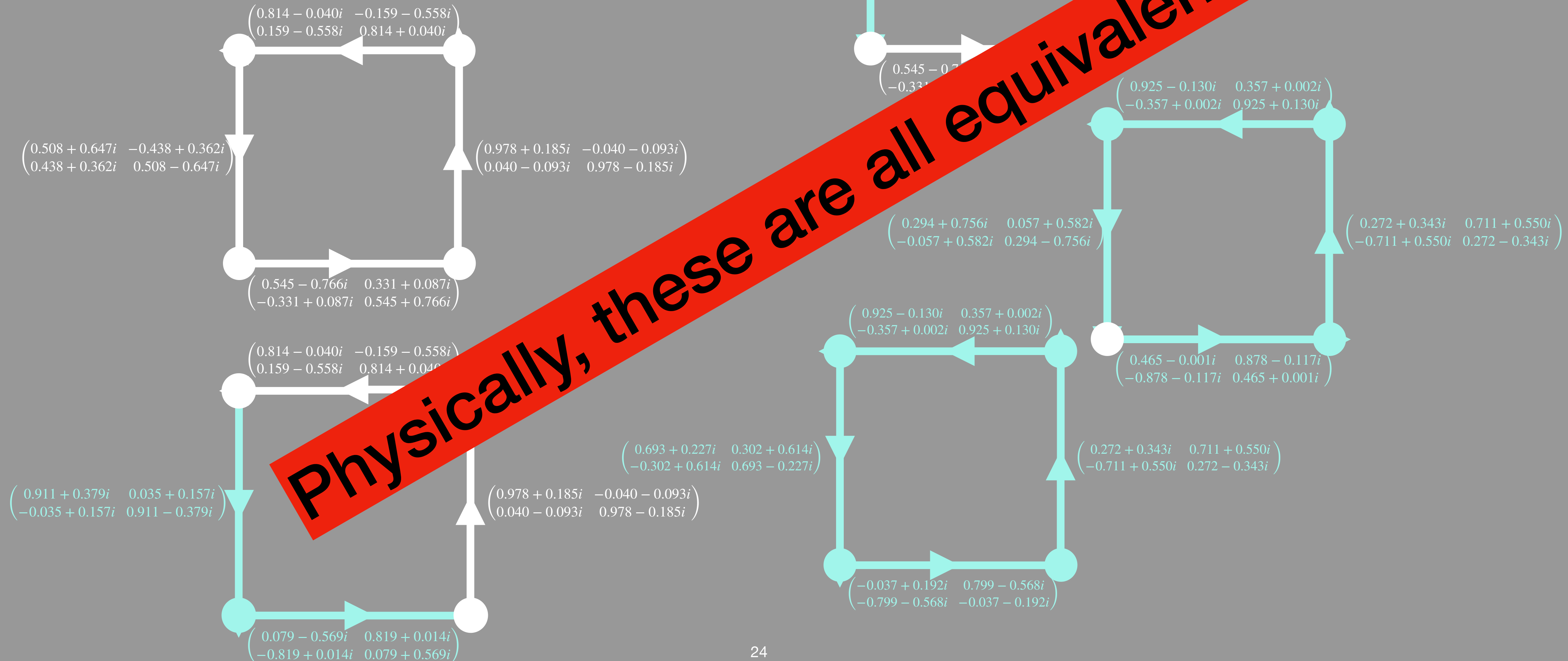
# SU(2) example

## Gauge symmetry



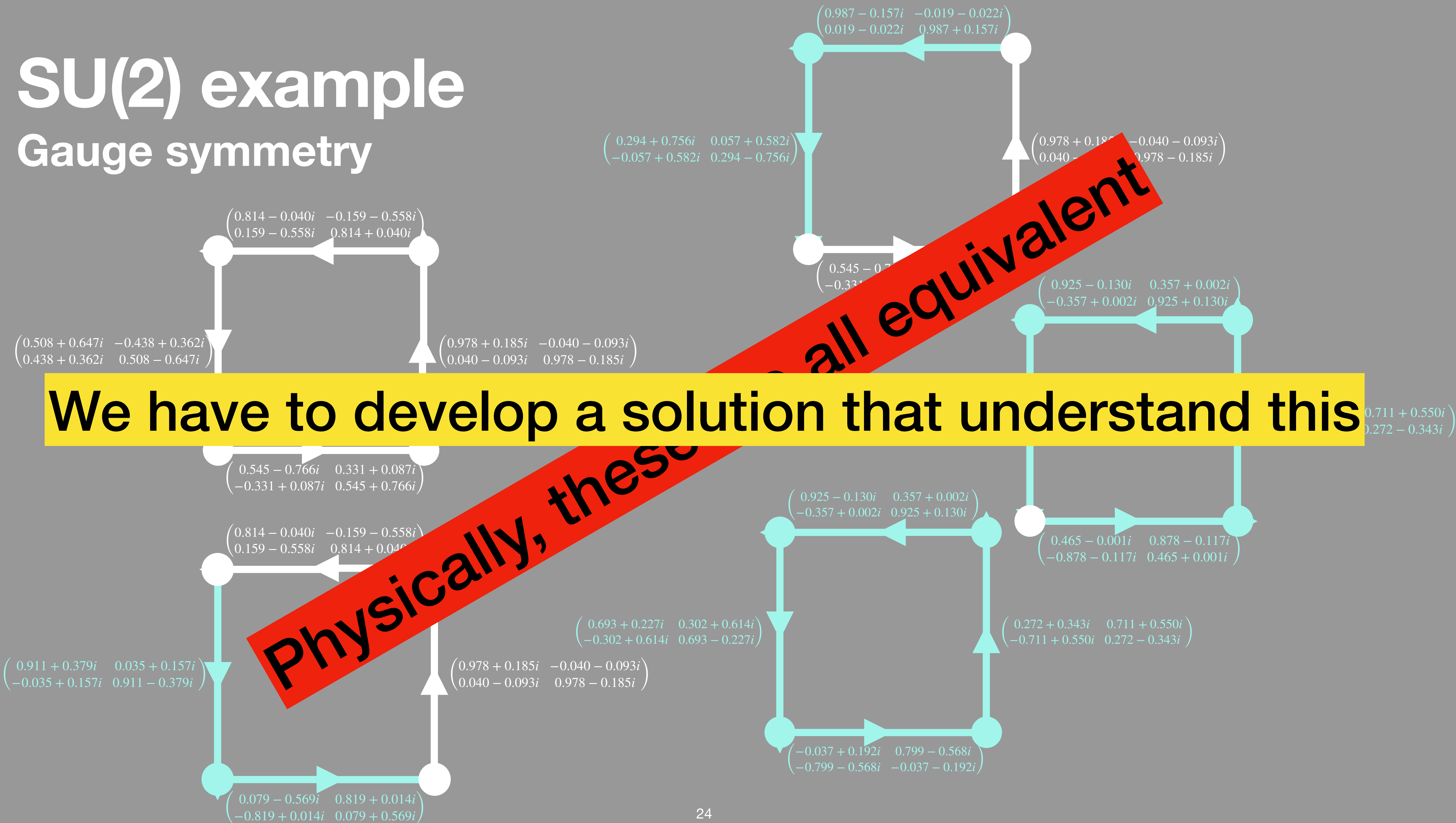
# SU(2) example

## Gauge symmetry



# SU(2) example

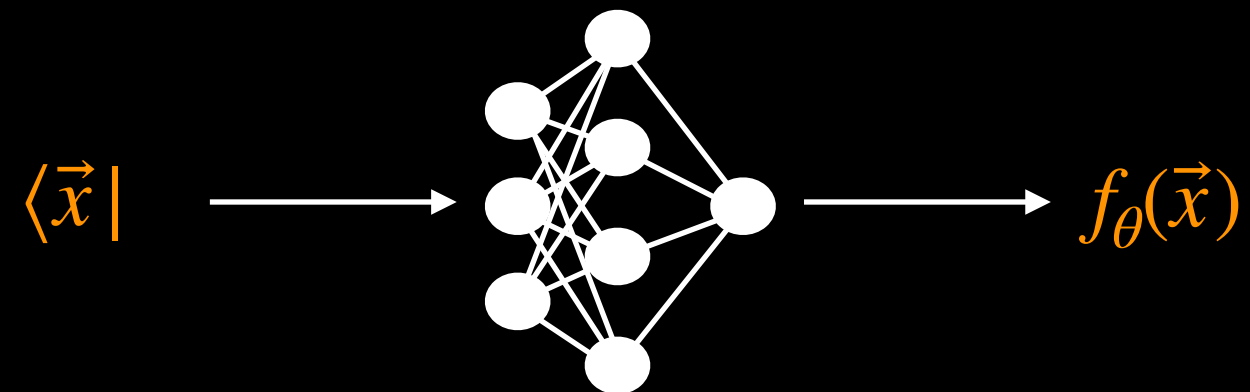
## Gauge symmetry





# Some ansatze

# Physics-informed

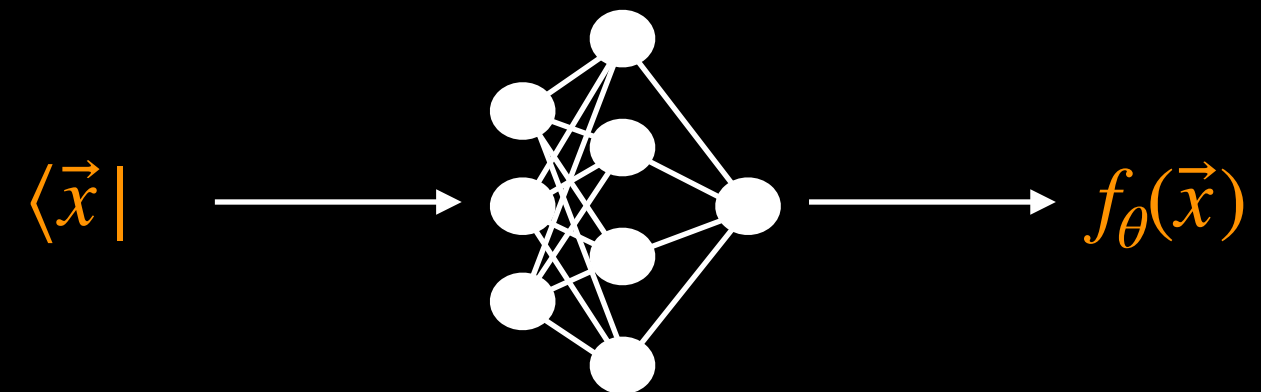


Greensite, J. P. (1979). Calculation of the Yang-Mills vacuum wave functional. *Nuclear Physics B*, 158(2-3), 469-496.

Chin, S. A., Van Roosmalen, O. S., Umland, E. A., & Koonin, S. E. (1985). Exact ground-state properties of the SU (2) Hamiltonian lattice gauge theory. *Physical Review D*, 31(12), 3201.



# Physics-informed



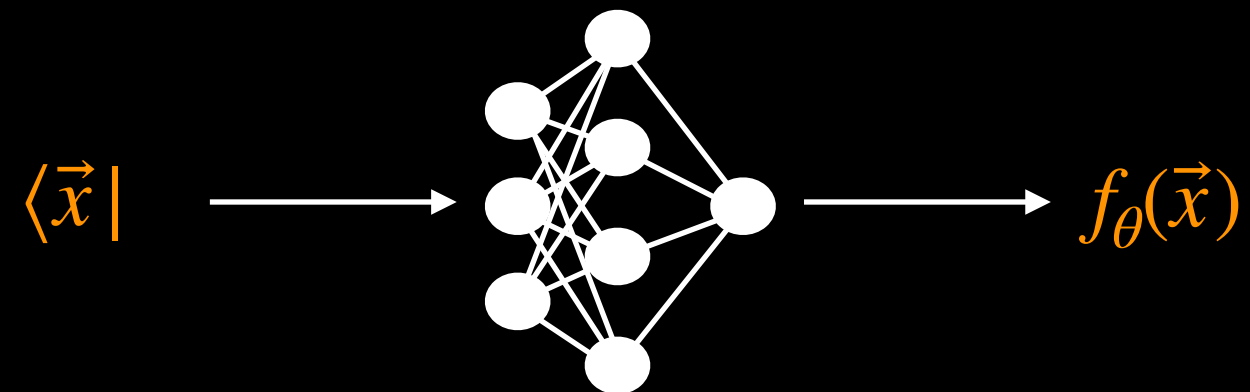
$$f(\mathbf{U}) = \prod_p e^{\alpha \frac{1}{2} \text{Tr}(P_p)}$$

$$\sum \text{Tr} \begin{array}{c} \begin{array}{ccc} & U_{\mu}^{\dagger}(x + \nu) & \\ U_{\nu}^{\dagger}(x) \downarrow & \square & \uparrow U_{\nu}(x + \mu) \\ U_{\mu}(x) & & \end{array} \end{array}$$

Greensite, J. P. (1979). Calculation of the Yang-Mills vacuum wave functional. *Nuclear Physics B*, 158(2-3), 469-496.

Chin, S. A., Van Roosmalen, O. S., Umland, E. A., & Koonin, S. E. (1985). Exact ground-state properties of the SU (2) Hamiltonian lattice gauge theory. *Physical Review D*, 31(12), 3201.

# Physics-informed



$$f(\mathbf{U}) = \prod_p e^{\alpha \frac{1}{2} \text{Tr}(P_p)}$$

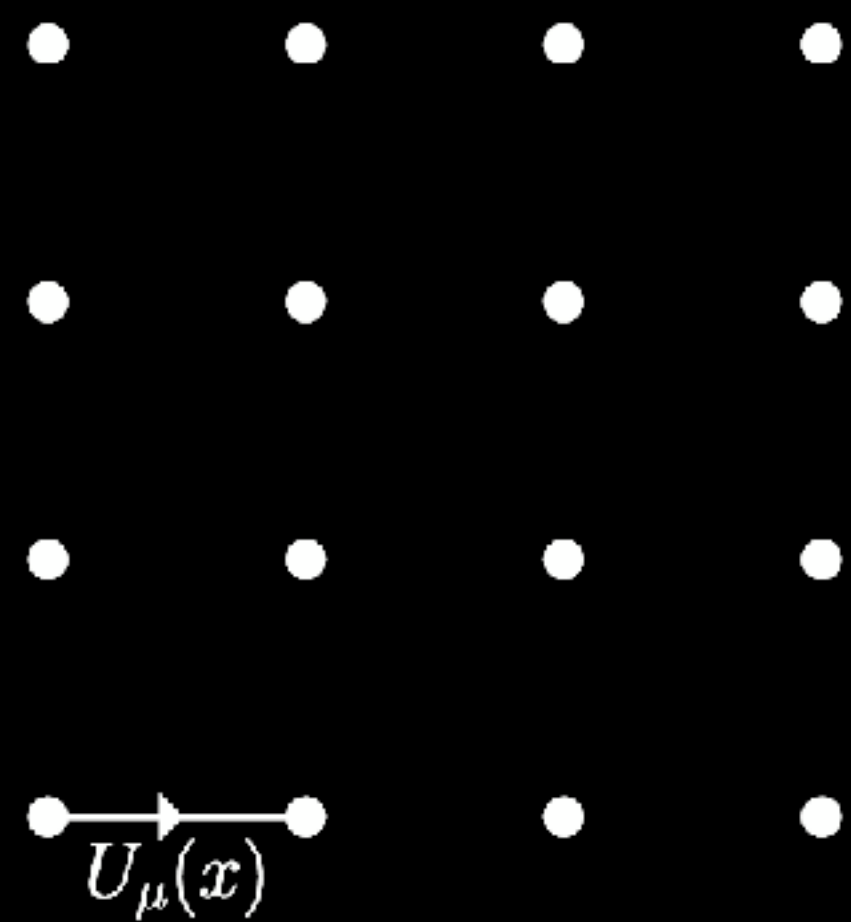
$$\sum \text{Tr} \begin{array}{c} \begin{array}{ccc} & U_{\mu}^{\dagger}(x + \nu) & \\ \downarrow U_{\nu}^{\dagger}(x) & \square & \uparrow U_{\nu}(x + \mu) \\ U_{\mu}(x) & & \end{array} \end{array}$$

Reliance on hand-crafted features

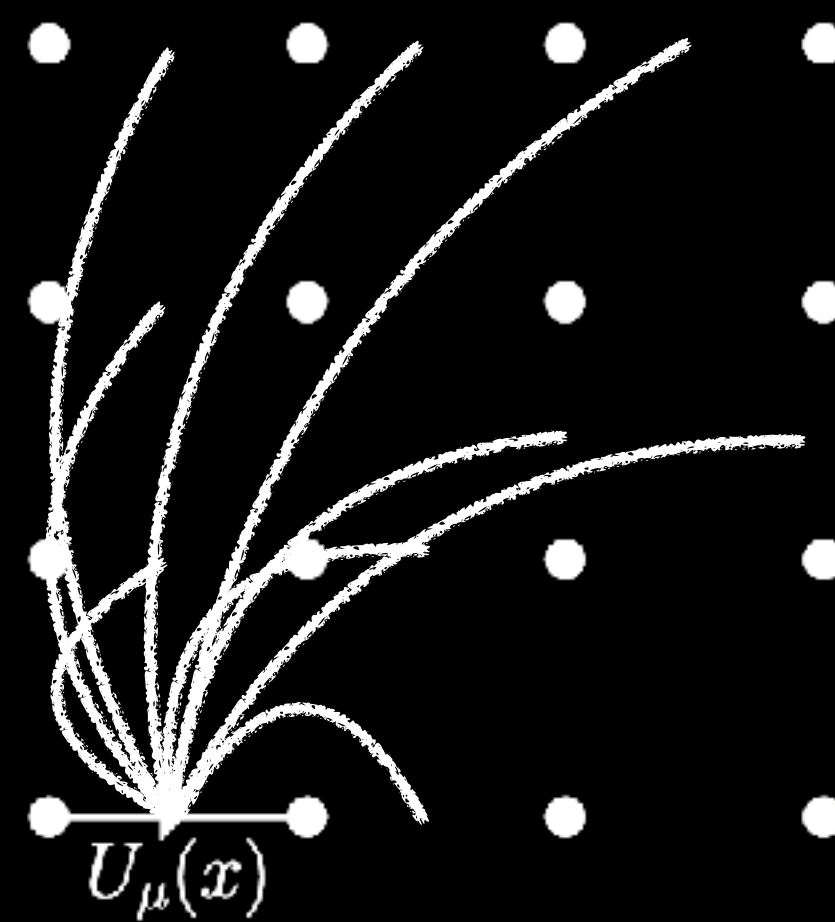
Greensite, J. P. (1979). Calculation of the Yang-Mills vacuum wave functional. *Nuclear Physics B*, 158(2-3), 469-496.

Chin, S. A., Van Roosmalen, O. S., Umland, E. A., & Koonin, S. E. (1985). Exact ground-state properties of the SU (2) Hamiltonian lattice gauge theory. *Physical Review D*, 31(12), 3201.

# Neural network + physics-informed

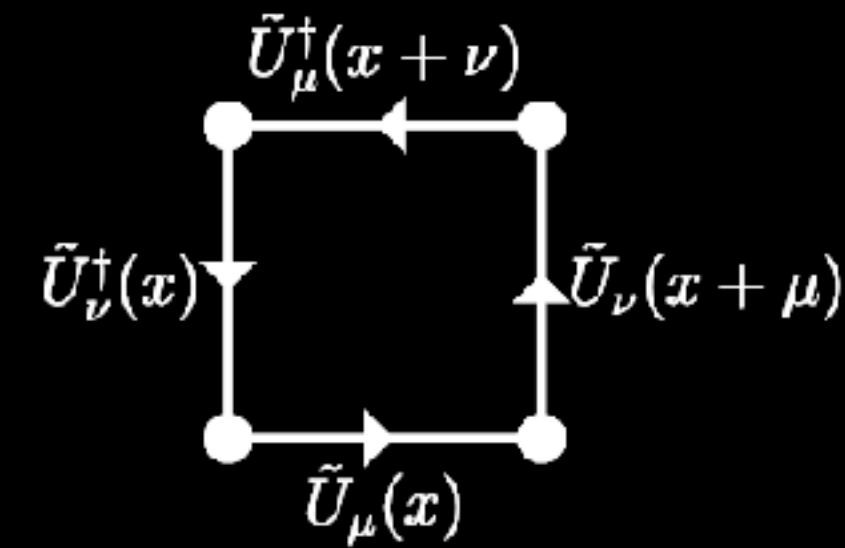


Neural  
network  
layers



$\Sigma$

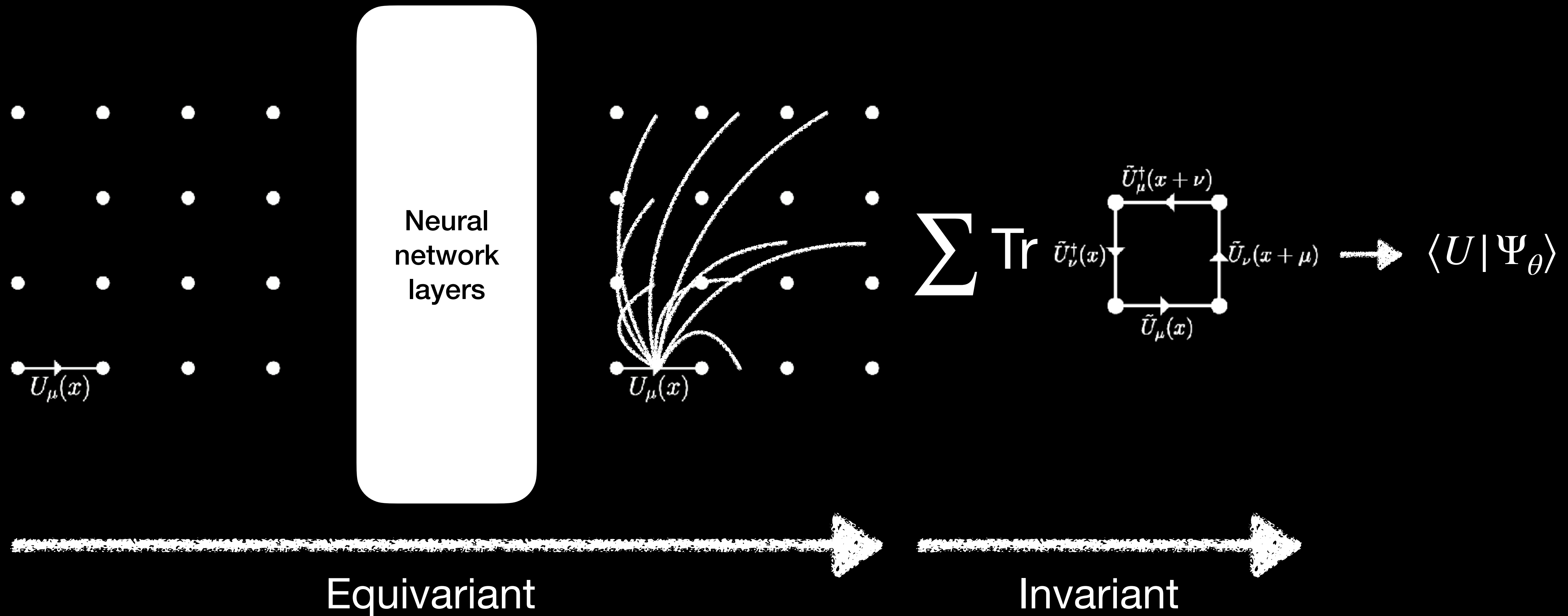
$\text{Tr}$



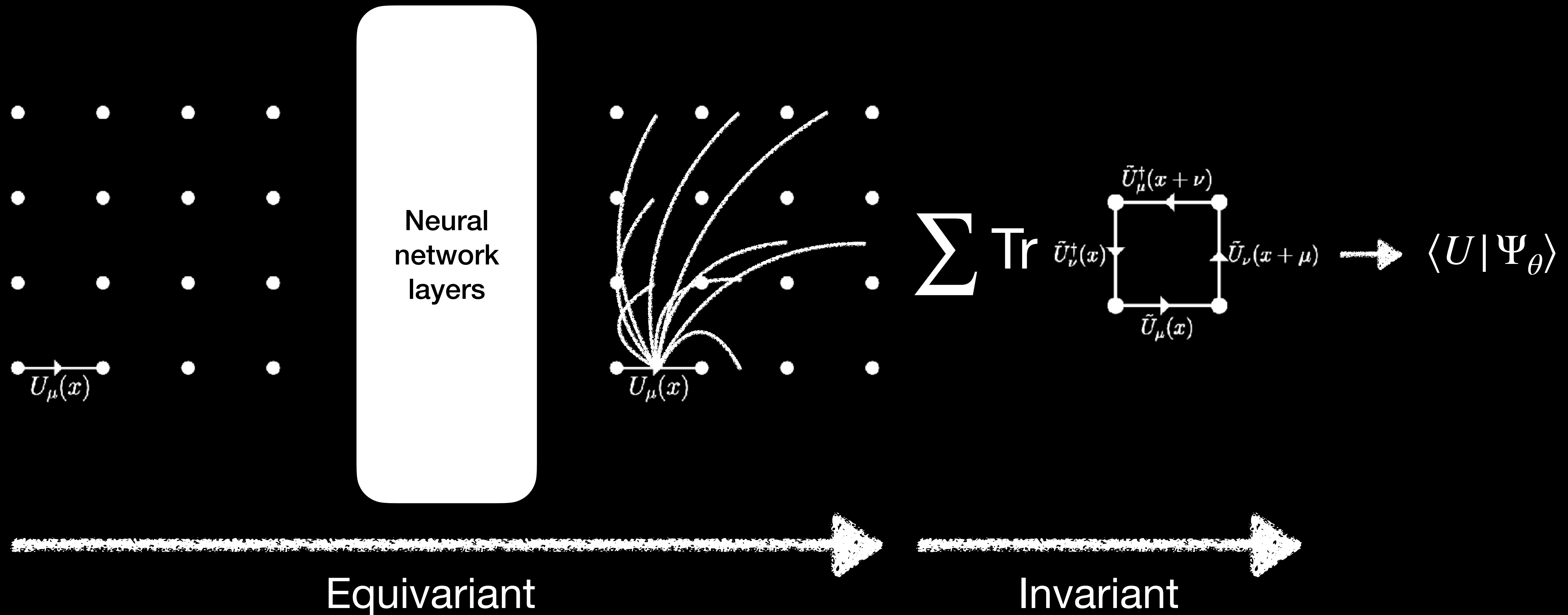
$\langle U | \Psi_\theta \rangle$



# Neural network + physics-informed



# Neural network + physics-informed

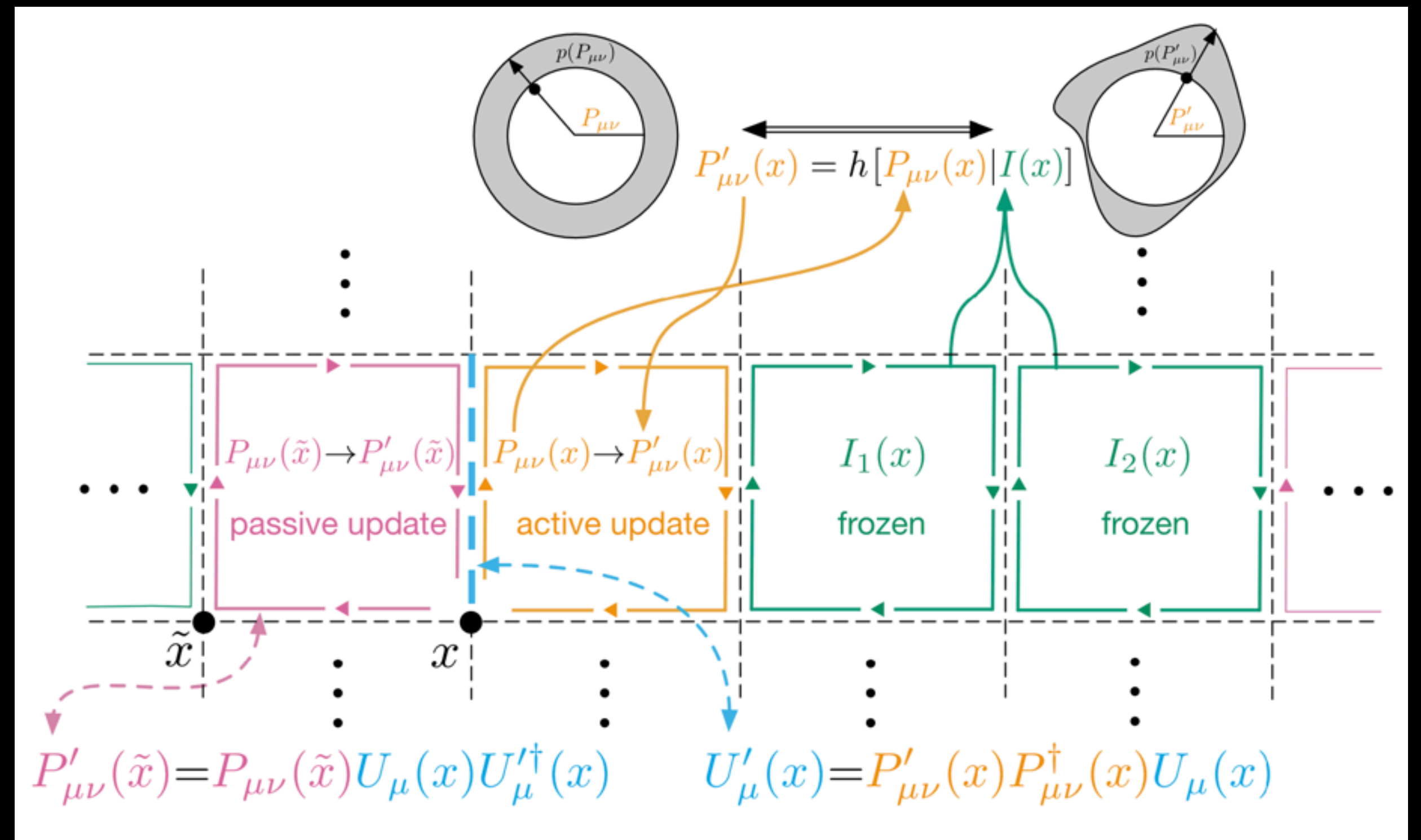


$$U_\mu(x) \rightarrow \Omega(x) U_\mu(x) \Omega^\dagger(x + \mu) \text{ before and after NN}$$

# Choices for network

## All gauge equivariant

- “Normalising flow networks”



Boyda, D., Kanwar, G., Racanière, S., Rezende, D. J., Albergo, M. S., Cranmer, K., ... & Shanahan, P. E. (2021). Sampling using SU (N) gauge equivariant flows. *Physical Review D*, 103(7), 074504.

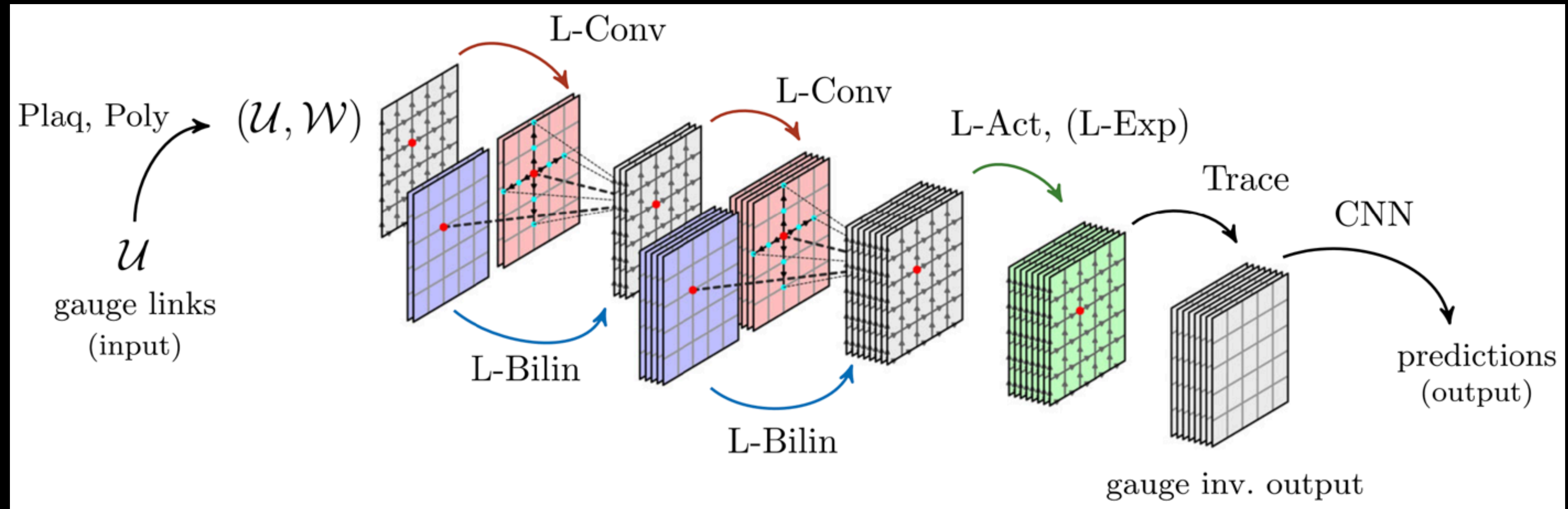
Kanwar, G., Albergo, M. S., Boyda, D., Cranmer, K., Hackett, D. C., Racaniere, S., ... & Shanahan, P. E. (2020). Equivariant flow-based sampling for lattice gauge theory. *Physical Review Letters*, 125(12), 121601.



# Choices for network

## All gauge equivariant

- Lattice Gauge Equivariant Convolutional Neural Networks (LCNN)



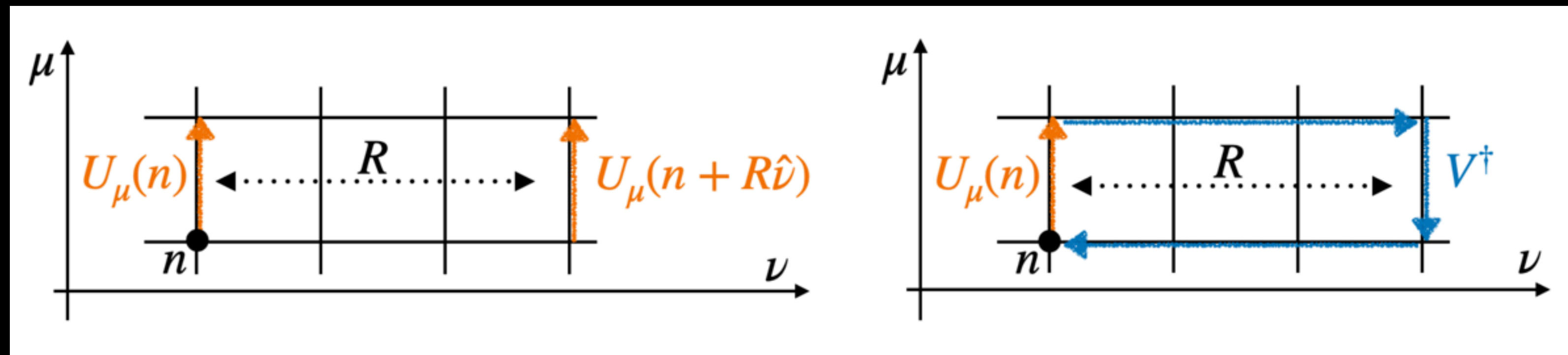
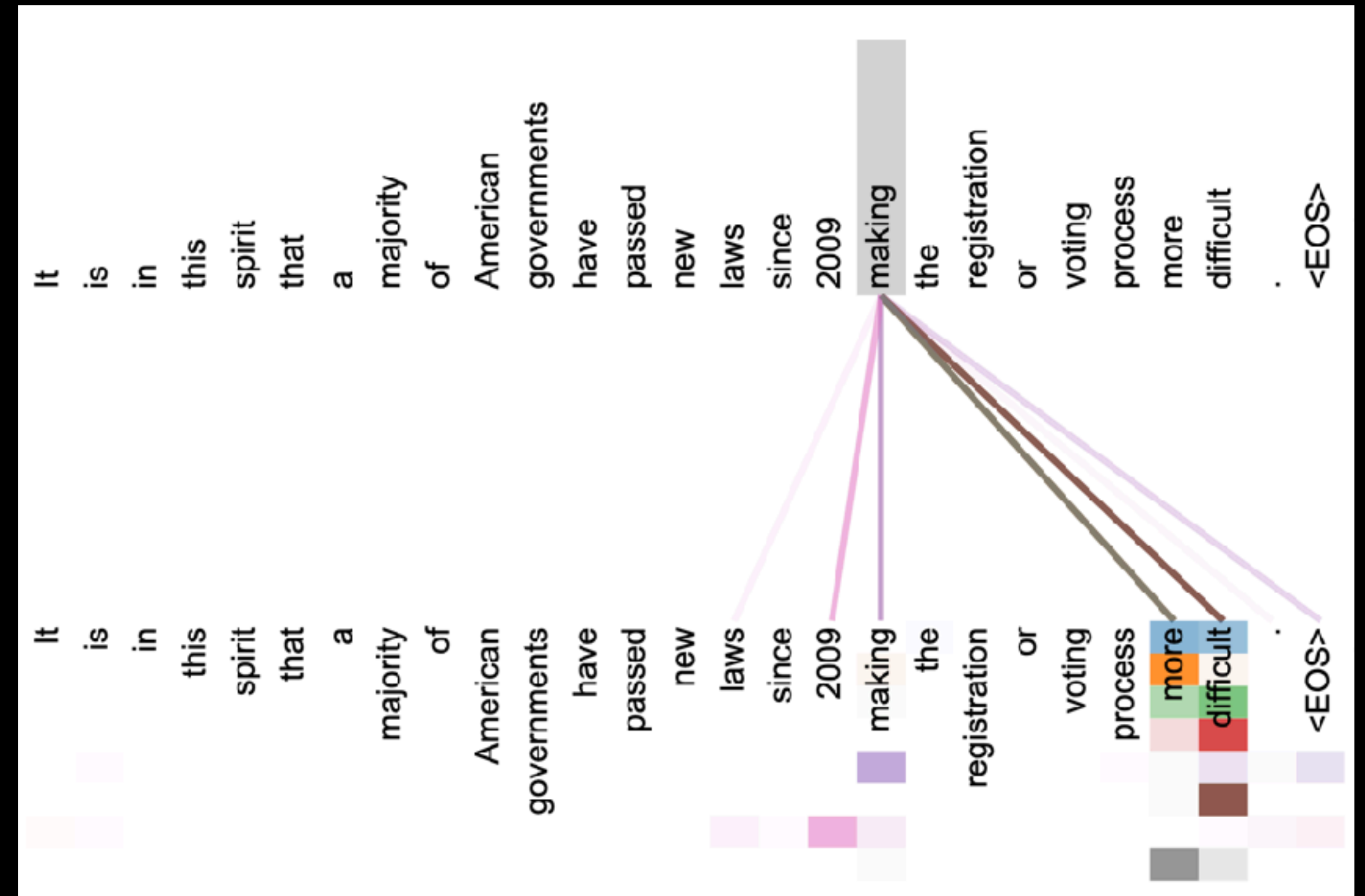
Favoni, M., Ipp, A., Müller, D. I., & Schuh, D. (2022). Lattice gauge equivariant convolutional neural networks. *Physical Review Letters*, 128(3), 032003.

# Choices for network

## All gauge equivariant

- Gauge equivariant transformer

Vaswani, A., ... & Polosukhin, I. (2017). Attention is all you need



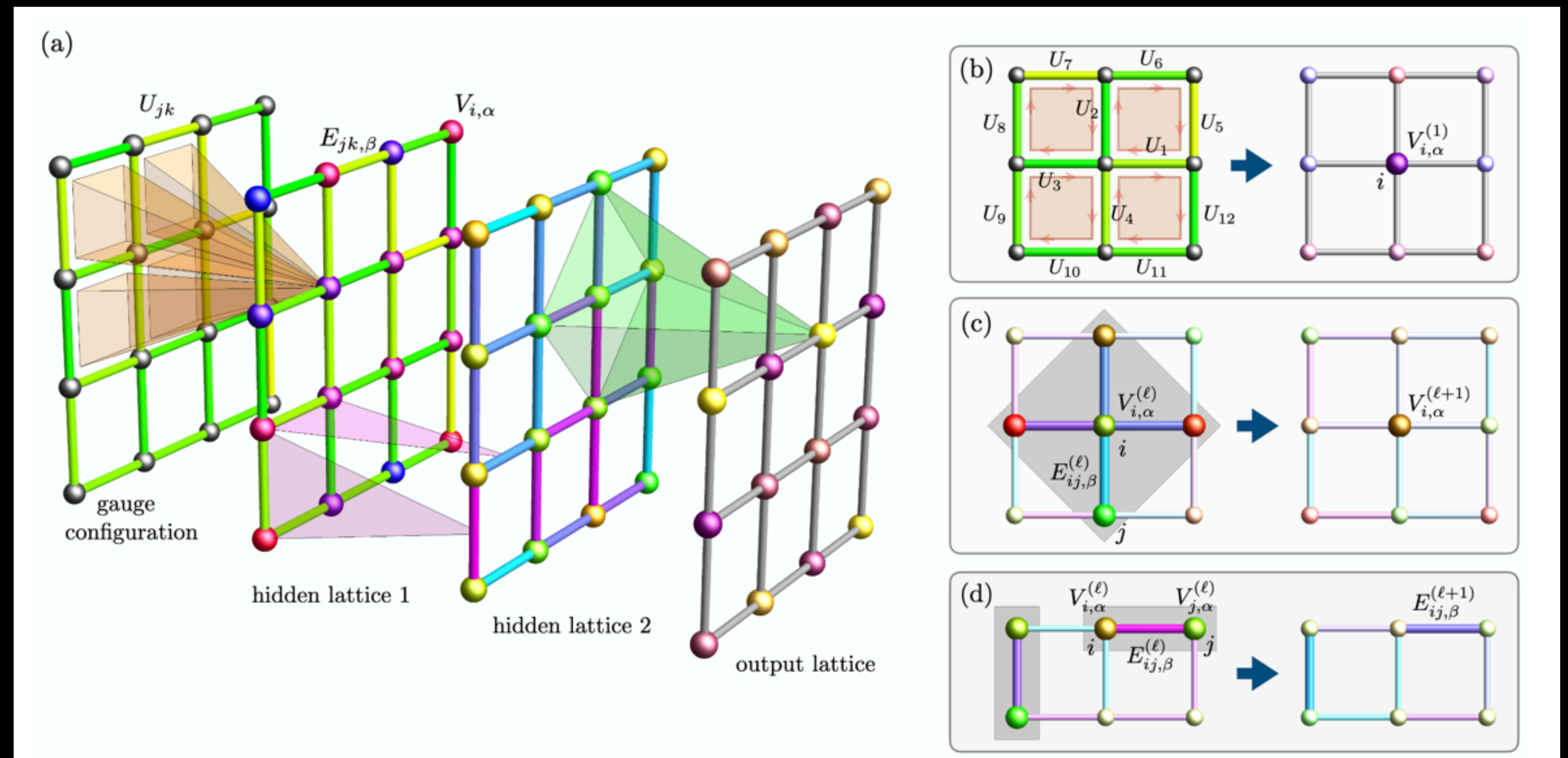
Nagai, Y., Ohno, H., & Tomiya, A. (2025). CASK: A Gauge Covariant Transformer for Lattice Gauge Theory. *arXiv preprint arXiv:2501.16955*.



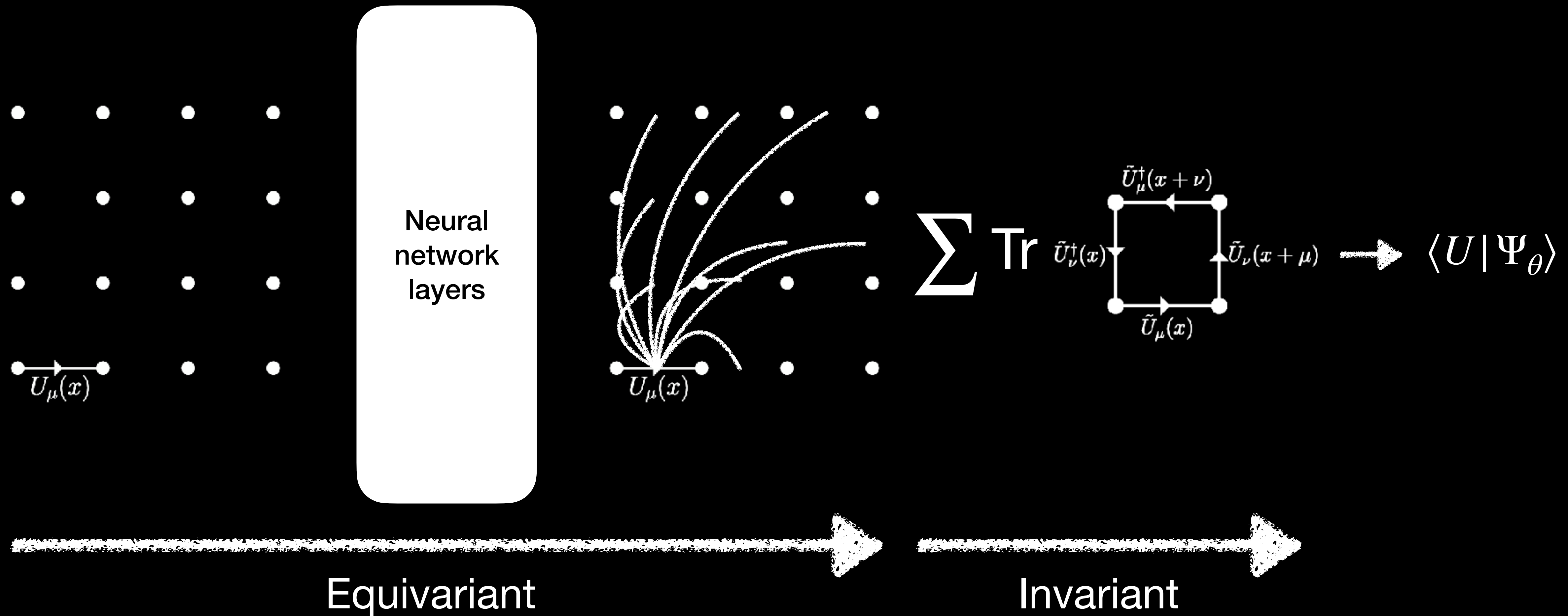
# Choices for network

## All gauge equivariant

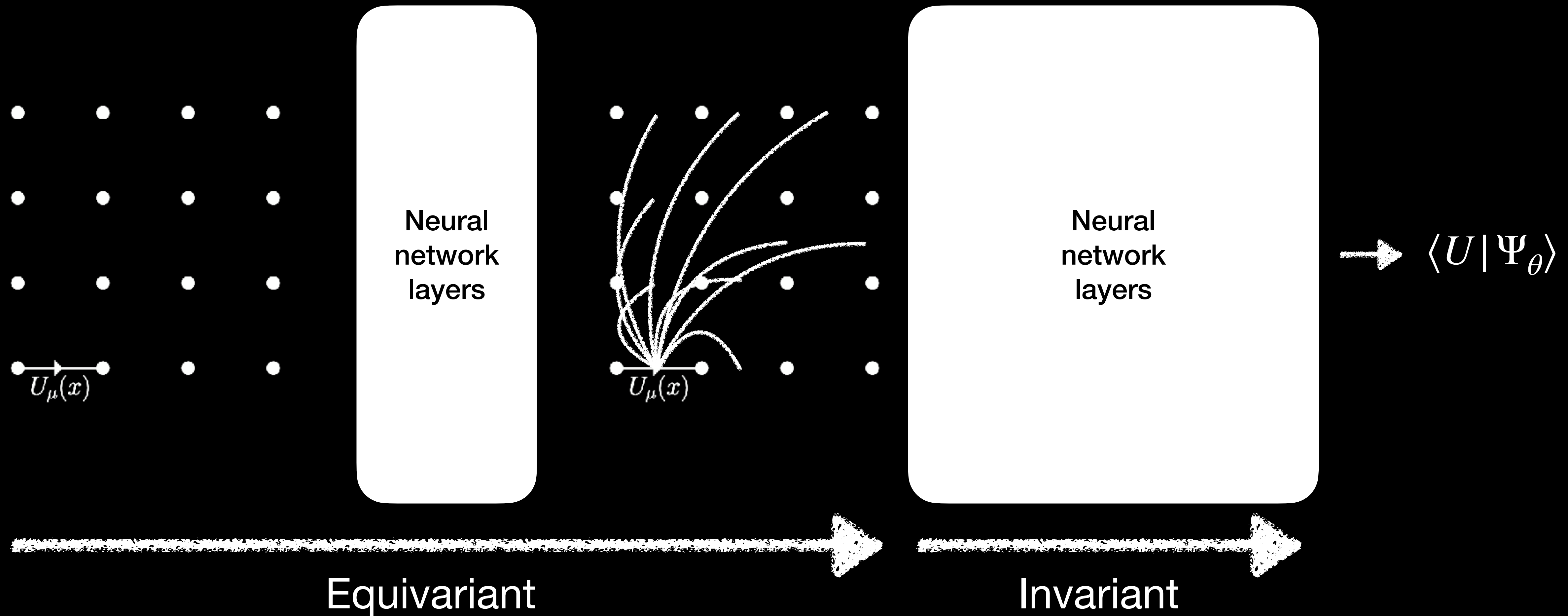
- Gauge equivariant graph neural network



# Neural network + physics-informed



# Neural network only

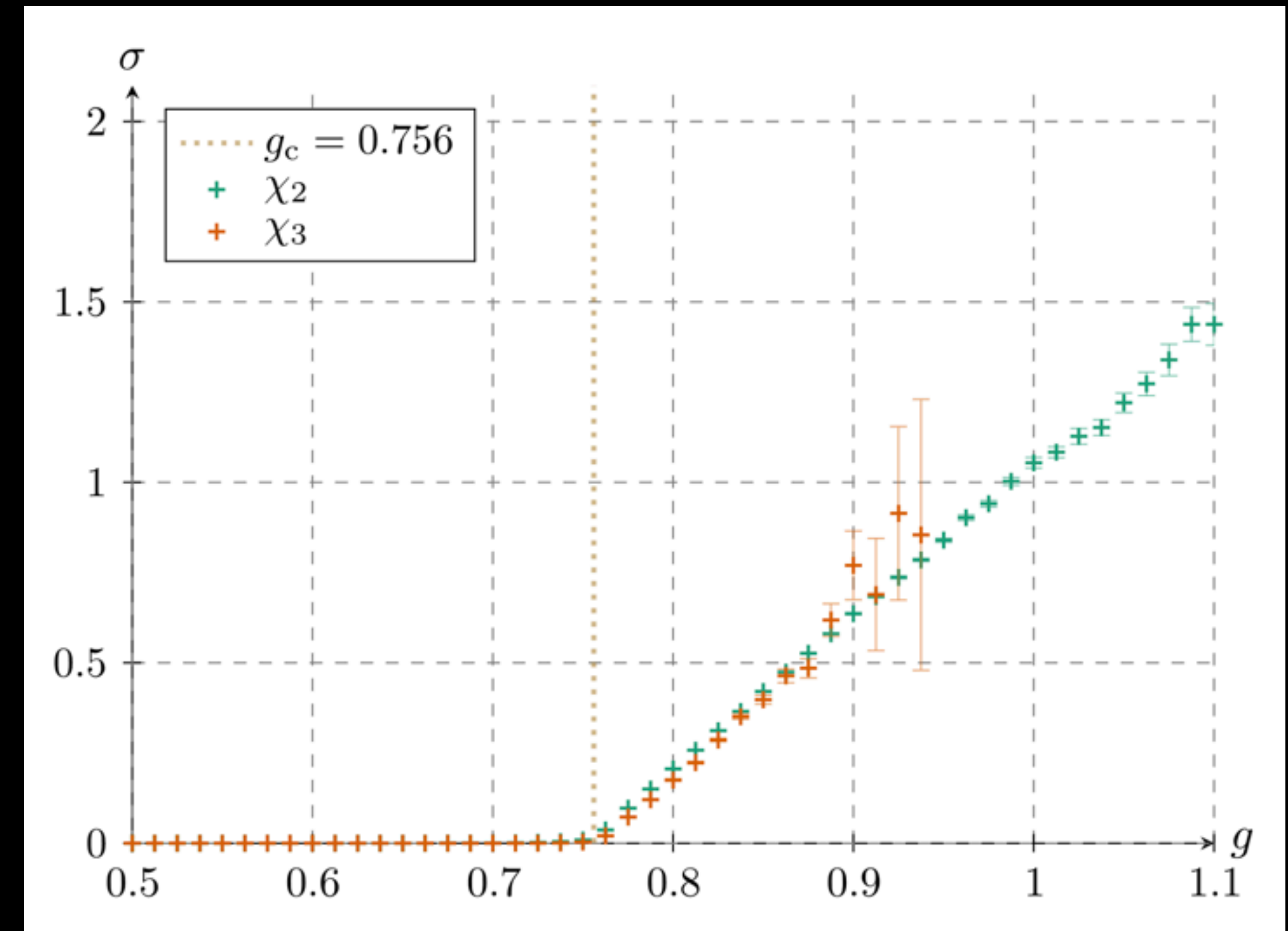




# Existing works

# $Z_2$ pure gauge theory

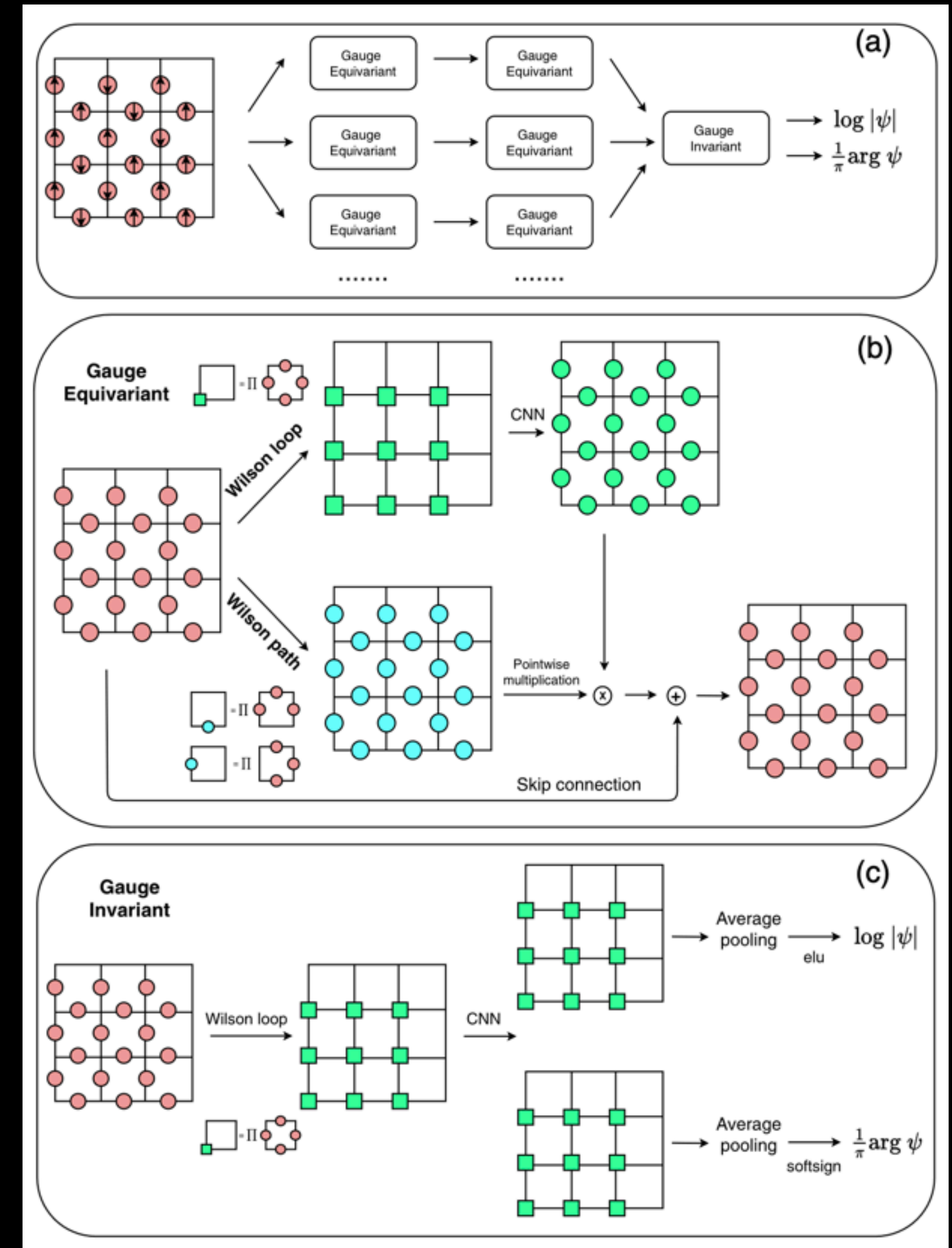
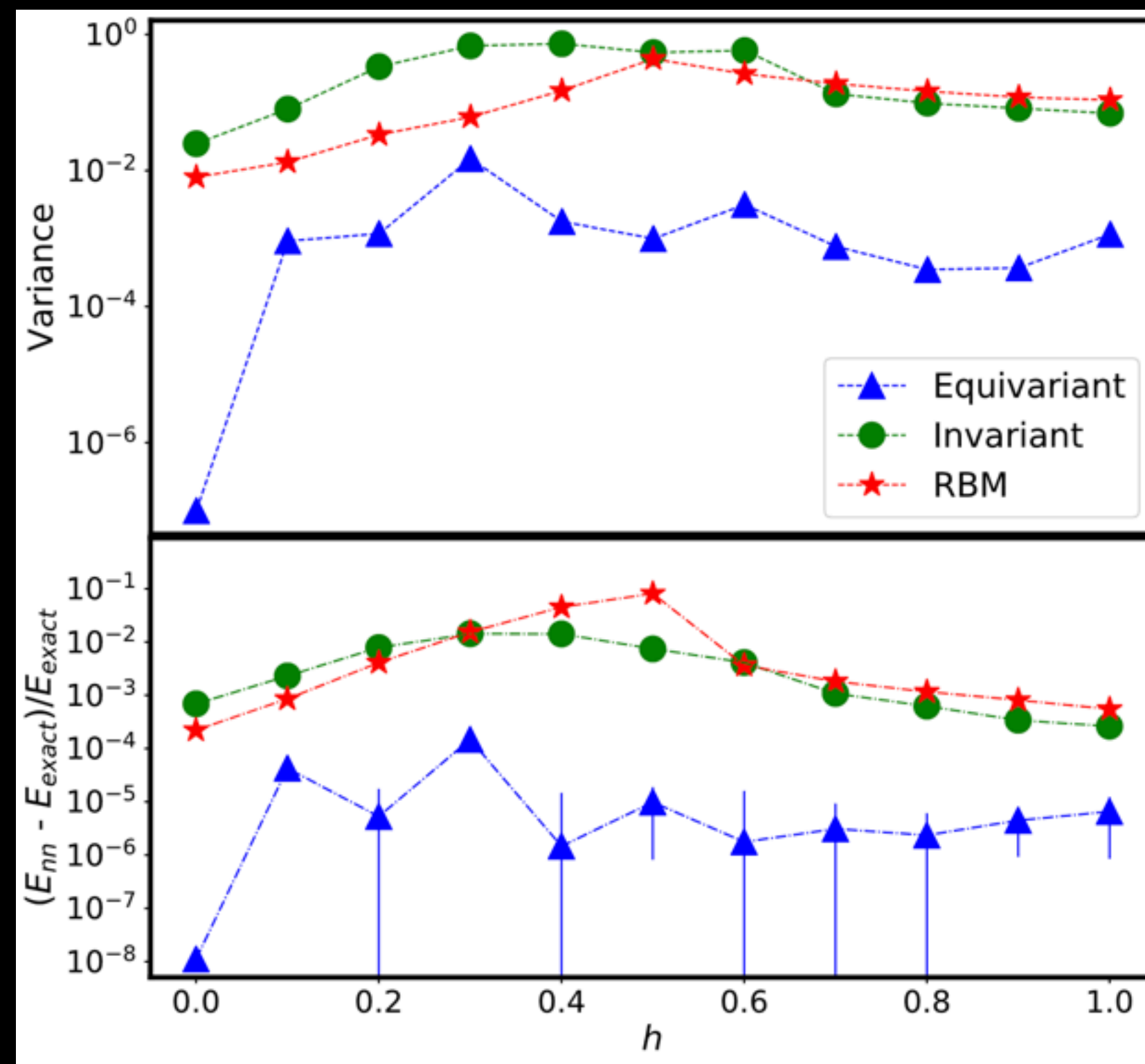
- Using the LCNN architecture
- Detect phase transition
  - Perimeter  $\rightarrow$  Area law scaling of Wilson loop





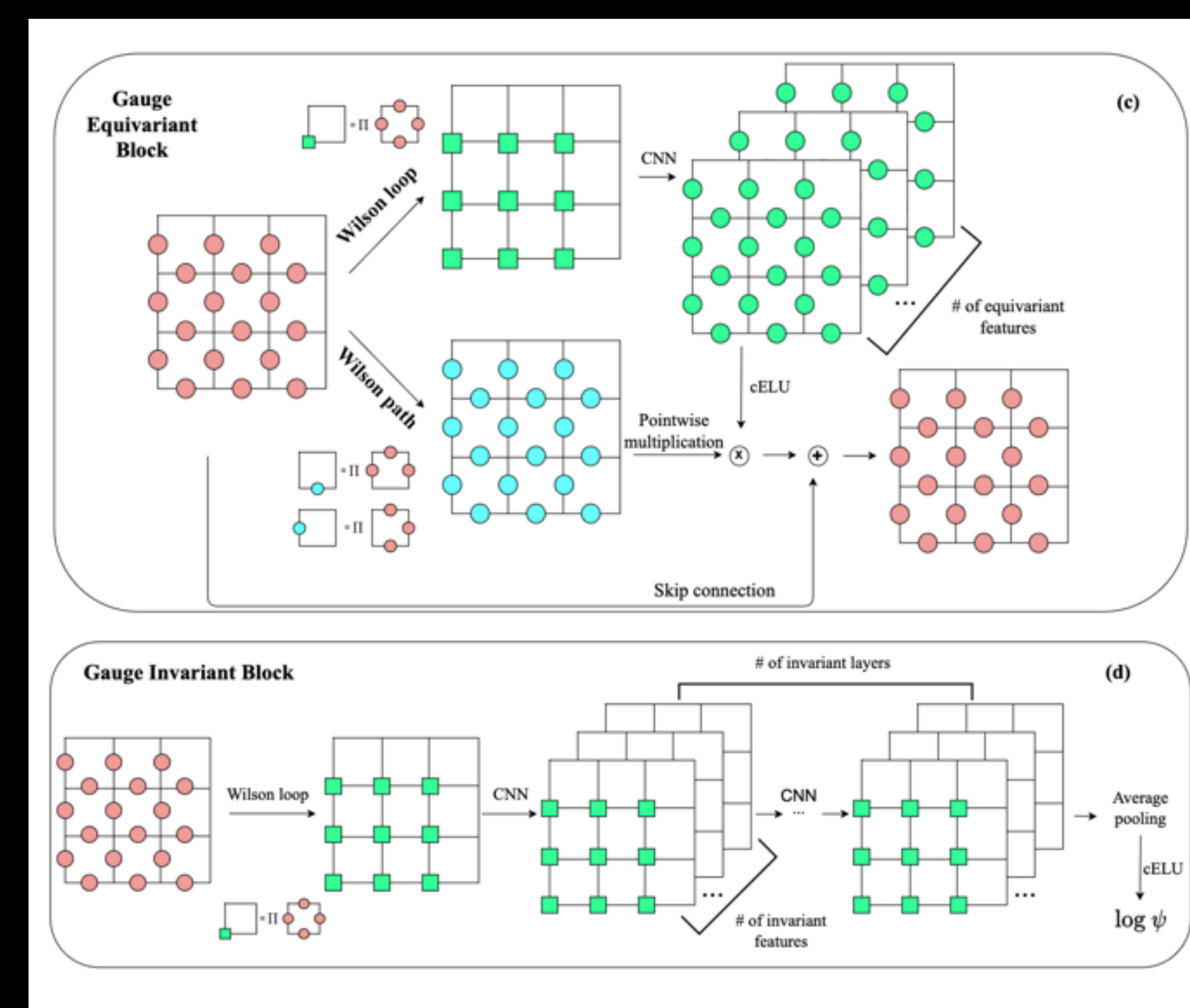
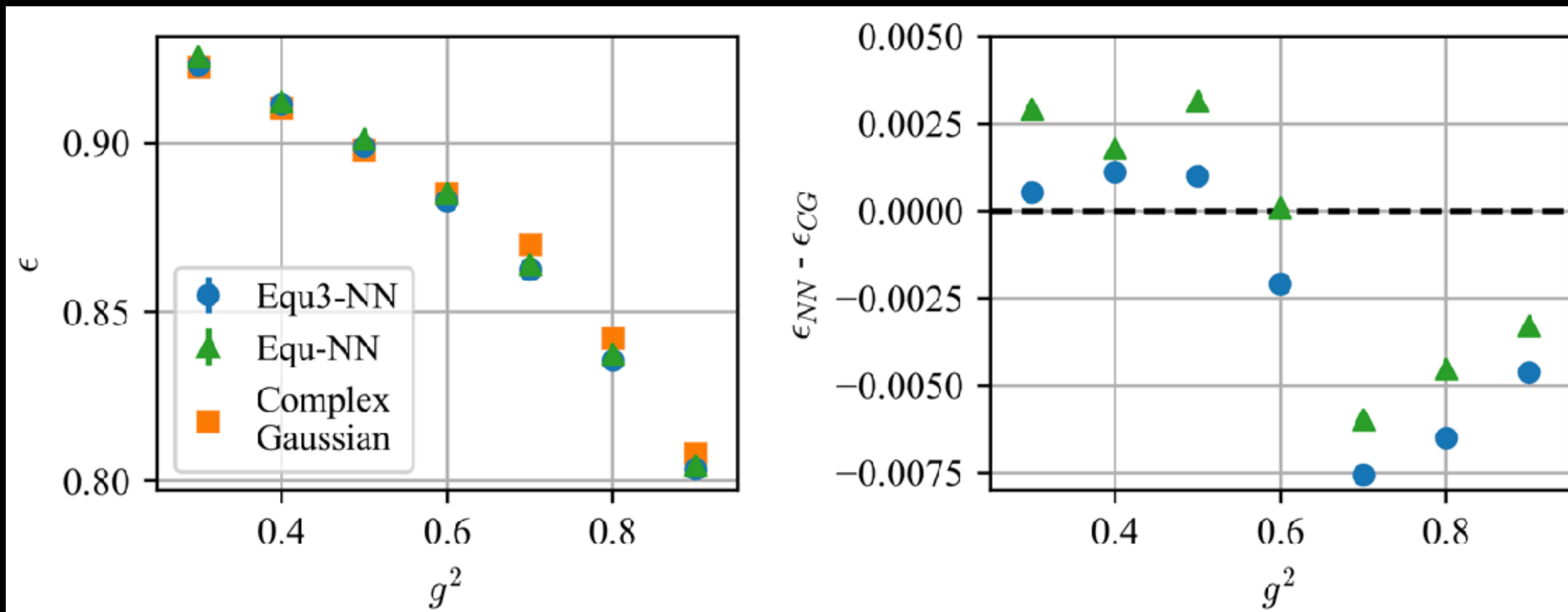
# $\mathbb{Z}_2$ pure gauge theory

- Using a custom ansatz
- Energy comparable to exact solution



# $U(1)$ pure gauge theory

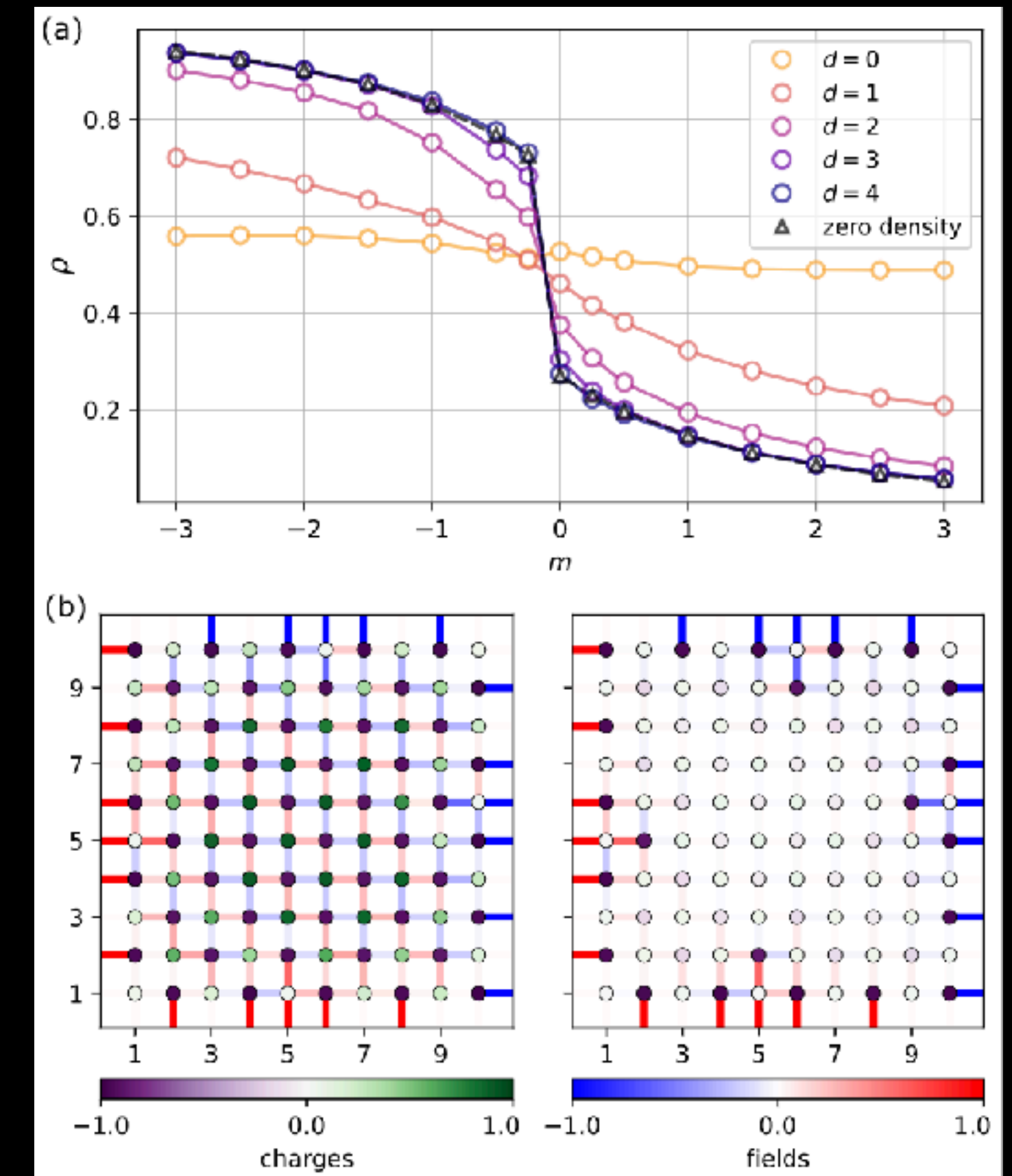
Energy similar to physics informed ansatz



Luo, D., Yuan, S., Stokes, J., & Clark, B. K. (2022). Gauge equivariant neural networks for 2+ 1d  $U(1)$  gauge theory simulations in hamiltonian formulation. *arXiv preprint arXiv:2211.03198*.

# $U(1)$ gauge theory with matter

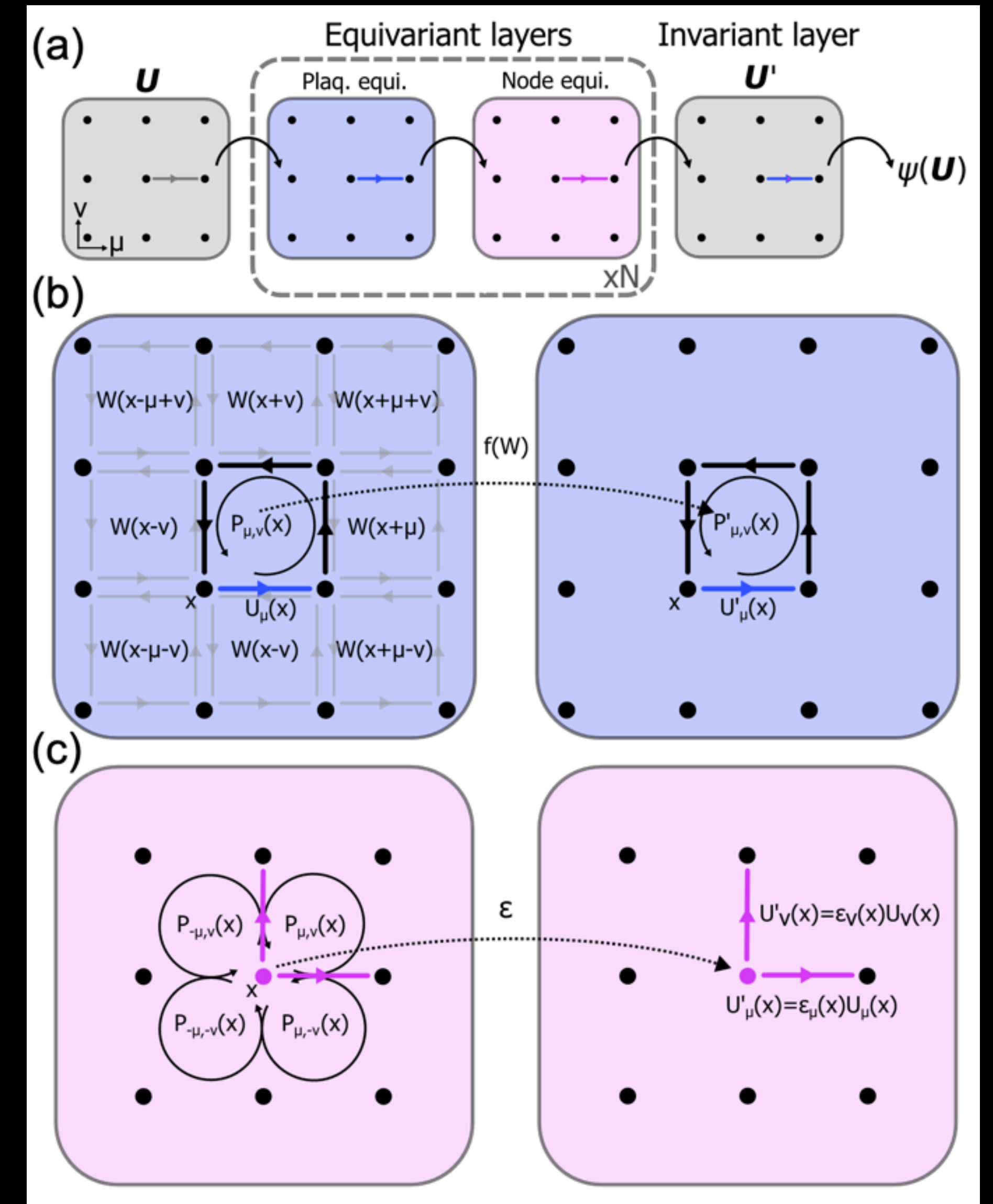
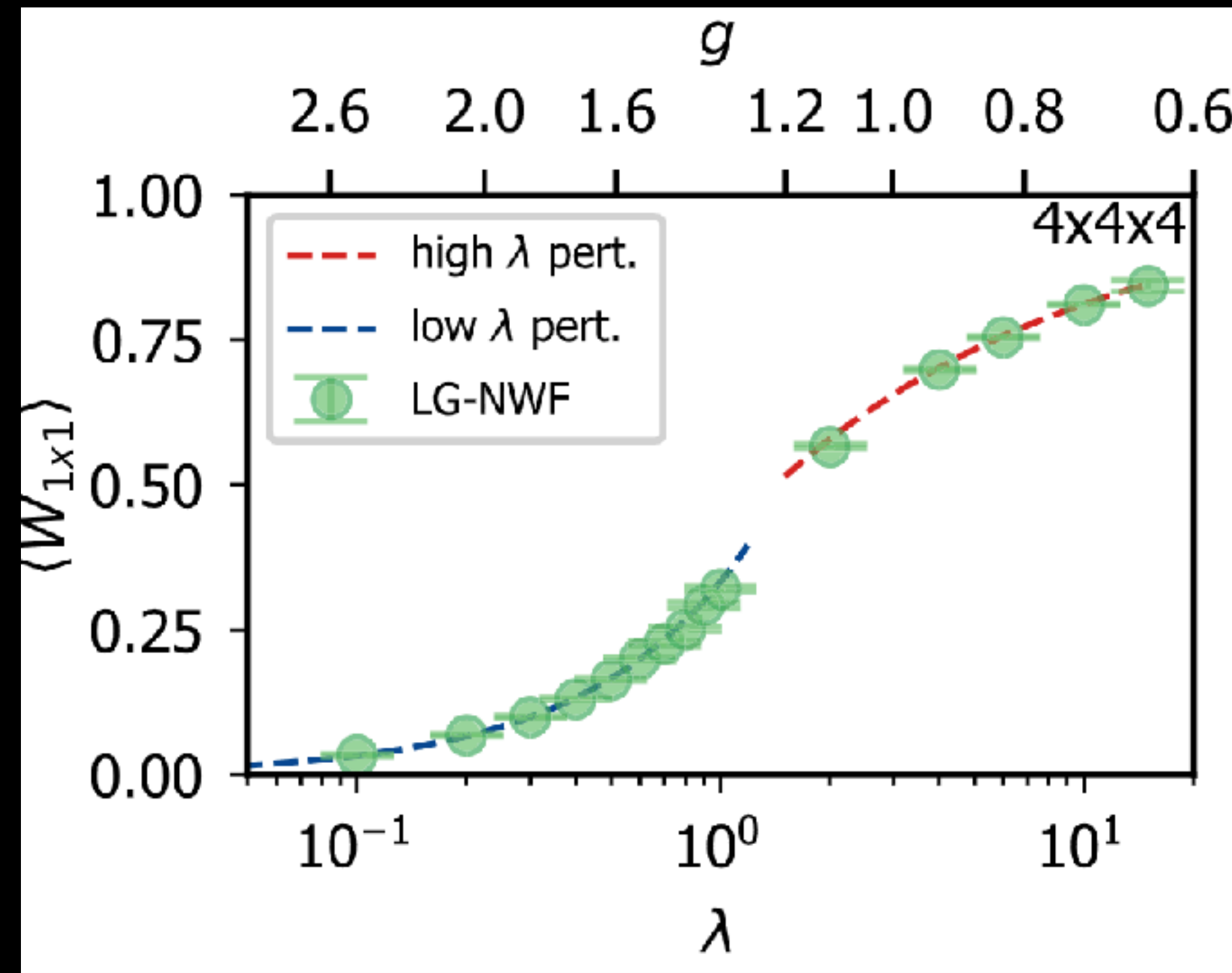
- **Autoregressive Ansatz** (no Markov chains)
- Detect/predict transitions in matter density





# $SU(2)$ pure gauge theory

- Both LCNN and flow-like architectures
- 3D



# Getting involved

# NetKet:

<https://www.netket.org/>

- Great tutorials
- Very modular code
  - Mixed learning curve
  - Lots of code/feature reuse
- Regular updates & new features
- [https://gitlab.com/QMAI/papers/su2\\_neural\\_wavefunctions](https://gitlab.com/QMAI/papers/su2_neural_wavefunctions)



**At Lattice 2026**

# Wednesday morning

## Vacuum structure and confinement

Ground States of adjoint  $SU(2)$  from Neural Newtork assisted Variational Montecarlo

Jul 29, 2026, 11:10 AM

20m

Thurgood Marshall (Adele H. Stamp Student Union)

Vacuum structure a...

Vacuum structure and ...

Speaker

David Ward (National Taiwan University)

# Thursday afternoon

## Algorithms and machine learning

Diagonalizing the Kogut-Susskind Hamiltonian with Physics-Informed Neural Networks

Jul 30, 2026, 4:10 PM

20m

Benjamin Baneker B (Adele H. Stamp Student Union)

Algorithms and artifi...

Algorithms and artifi...

Speaker

Simone Romiti

Neural Wavefunctions in Quantum Field Theory

Jul 30, 2026, 4:30 PM

20m

Benjamin Baneker B (Adele H. Stamp Student Union)

Contributed talk

Algorithms and artifi...

Algorithms and artifi...

Speaker

Suryansh Rajawat (University of Maryland)

3+1d  $SU(3)$  Yang-Mills + theta-term with Neural Network Quantum States

Jul 30, 2026, 4:50 PM

20m

Benjamin Baneker B (Adele H. Stamp Student Union)

Contributed talk

Algorithms and artifi...

Algorithms and artifi...

Speaker

Joshua Lin (Argonne National Laboratory)



Eliska Greplova



Juan Carrasquilla



Jannes Nys

# Thank you for listening



# Backup slides



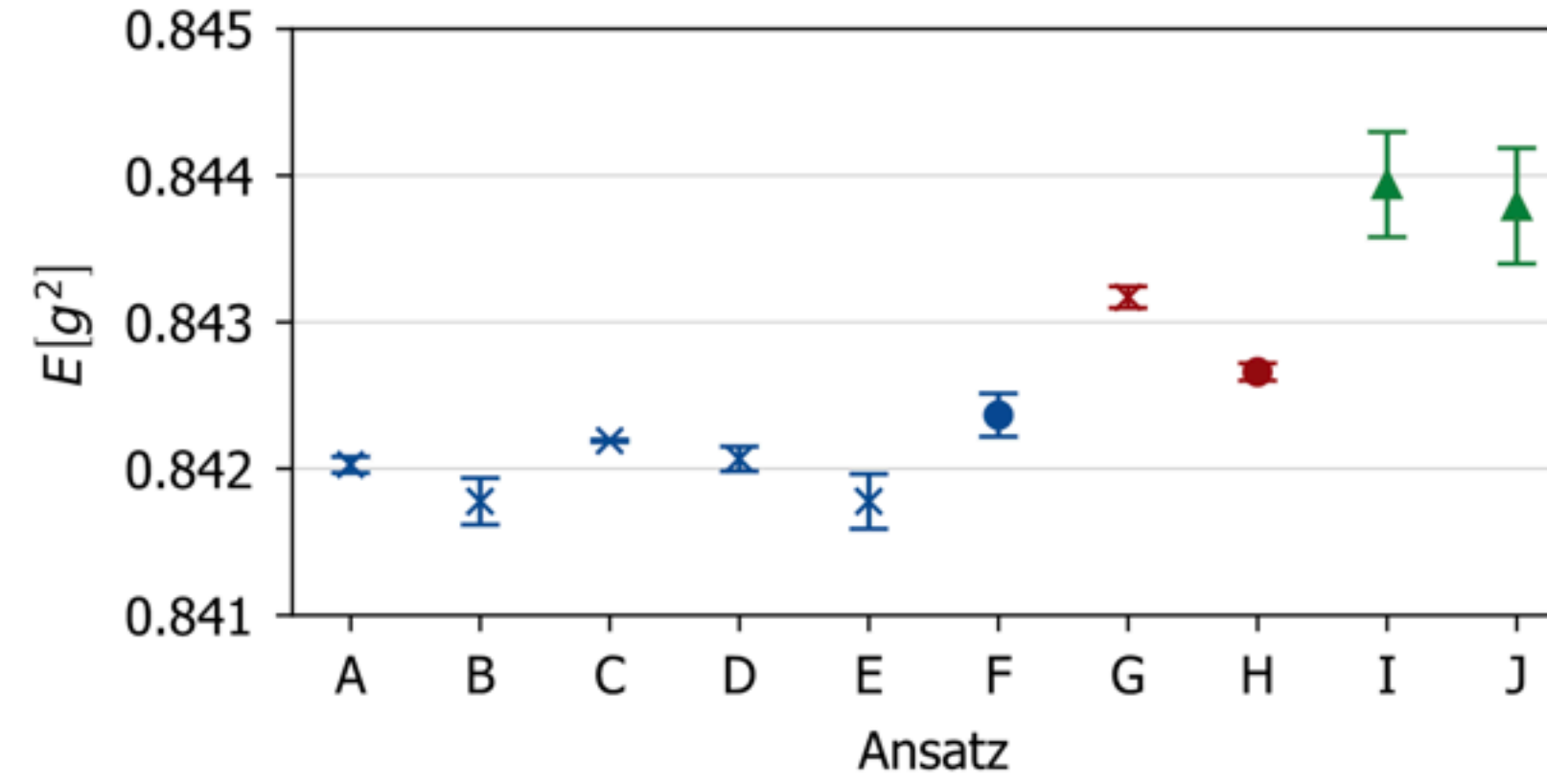


Figure S2. Variational ground state energy for a range of ansatze outlined in Table S1 for  $\lambda = 1$ . Errors are the Monte Carlo error of the energy over the final set of samples.

| Figure S2 label | Plaquette layer |             |           | Node layer | Num. equivariant blocks | Final layer |
|-----------------|-----------------|-------------|-----------|------------|-------------------------|-------------|
|                 | Architecture    | Num. layers | Other     |            |                         |             |
| A               | CNN             | 1           | K=3       | Yes        | 1                       | Jastrow     |
| B               | CNN             | 1           | K=1       | Yes        | 1                       | Jastrow     |
| C               | CNN             | 2           | K=3       | Yes        | 1                       | Jastrow     |
| D               | CNN             | 1           | K=3       | Yes        | 2                       | Jastrow     |
| E               | CNN             | 1           | K=3       | Yes        | 1                       | FCN         |
| F               | ViT             | 1           | NH=2, D=4 | Yes        | 1                       | Jastrow     |
| G               | CNN             | 1           | K=3       | No         | 1                       | Jastrow     |
| H               | ViT             | 1           | NH=2, D=4 | No         | 1                       | Jastrow     |
| I               | None            | N/A         | N/A       | No         | 0                       | Jastrow     |
| J               | None            | N/A         | N/A       | No         | 0                       | FCN         |



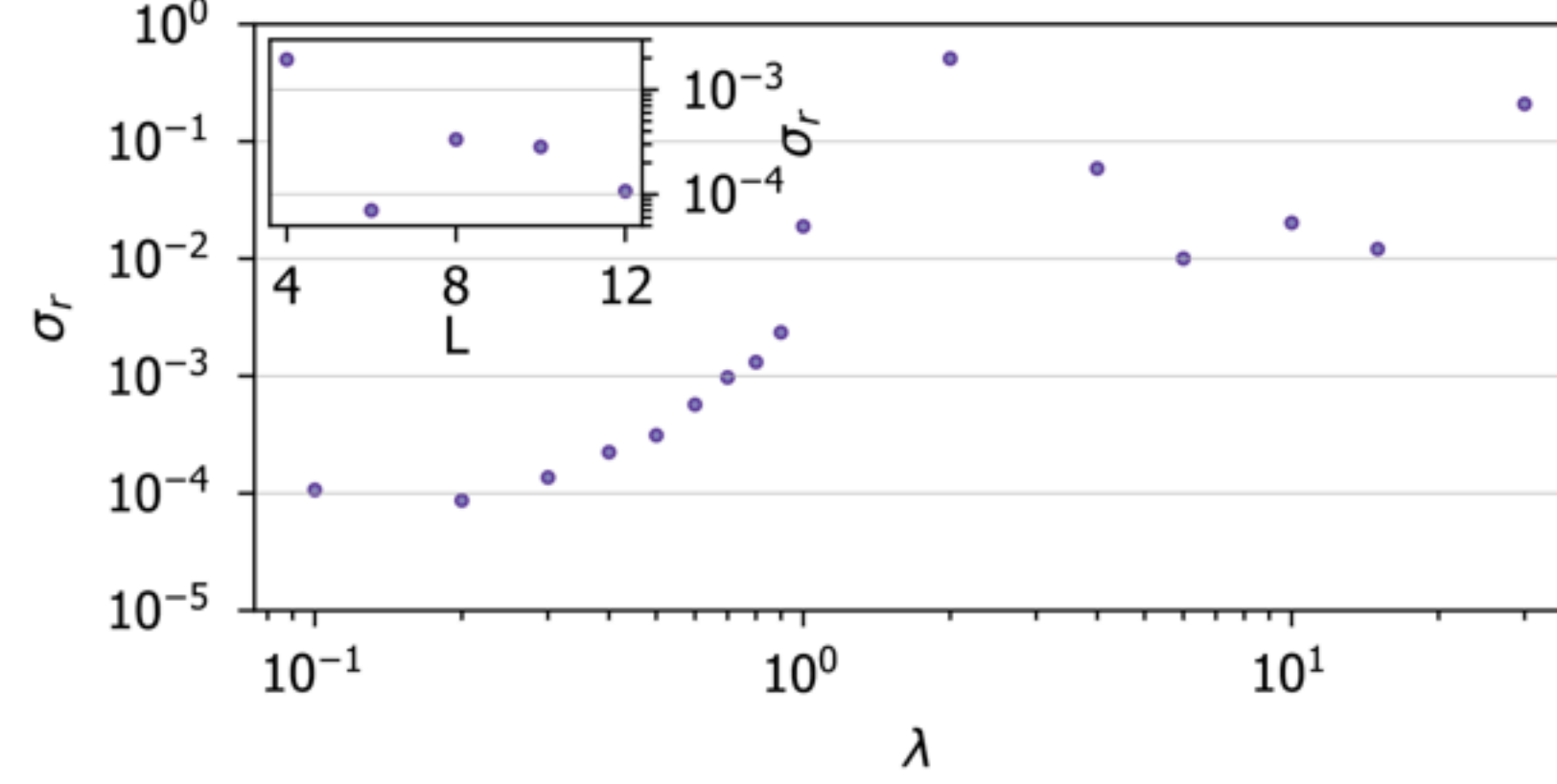


Figure S4. The rescaled variance of the energy, as in Eq. (D22), averaged over the final 20 iterations of VMC. The main panel shows results for  $12 \times 12$  lattices at various values of  $\lambda$ ; the inset shows the rescaled variance but against system size ( $L \times L$ ) with fixed  $\lambda = 0.1$ .

the version trained under the alternative Hamiltonian and then fine-tuned and they agree within errors.

During the VMC procedure we also compute the energy variance  $\sigma$ , defined as

$$\sigma = \frac{\langle \Psi | H^2 | \Psi \rangle}{\langle \Psi | \Psi \rangle} - \left( \frac{\langle \Psi | H | \Psi \rangle}{\langle \Psi | \Psi \rangle} \right)^2. \quad (\text{D21})$$

# Different gauge groups

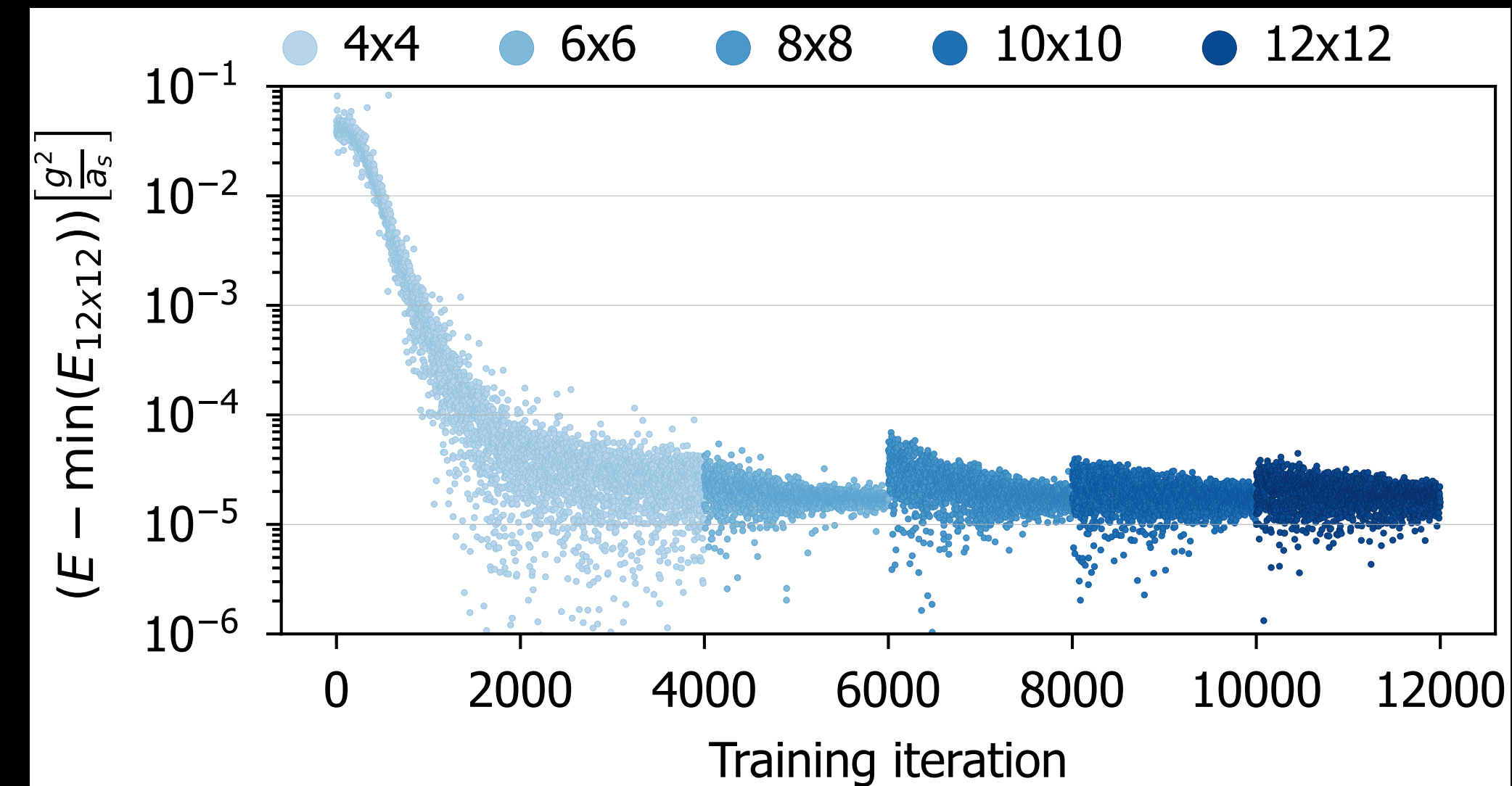
## Do-able

- $U = U(\theta, \phi, \psi) \rightarrow 8$  angles
- $\nabla^2 \rightarrow$  More complicated
  - (possibly done but not explained in Chin, S. A., *et al.* Phys. Rev. D **37**, 3001)
- Note: canonical ordering of eigenvalues non-trivial in SU(3)
  - When do go  $P_{\mu,\nu}(x) \rightarrow \lambda_P \rightarrow \lambda'_P \rightarrow P'_{\mu,\nu}(x)$  you have to order the eigenvalues before and after so they correspond to the same eigenvectors

# Larger lattices

## Trivial

- We did transfer learning in this work already
- Equivariant layers are local
  - Convolution layers, fixed kernel size
- Final layer is non-local
  - One new parameter for each new unique distance on the lattice
    - Can be initialised as ‘small’
- Can also be done in arbitrary dimension, 2 and 3 shown in paper





# Fermions and time evolution

## Hard, open study

- Both done in spin systems with neural wavefunctions
  - Even at the same time:
    - Nys, J., Pescia, G., Sinibaldi, A. and Carleo, G., 2024. Ab-initio variational wave functions for the time-dependent many-electron Schrödinger equation. *Nature communications*, 15(1), p.9404.
- Gauge invariance and representation for LGTs with fermions remain open questions
- Time evolution is more straightforward (only variational challenges, nothing conceptual)

# Other applications

## Done, but not for LGTs

- Physics distilled into neural network parameters,  $\theta$ 
  - Fewer parameters than basis states
- Infer phase transitions
  - Hernandez, V., **Spriggs**, T., Khaleefah, S., Greplova, E. ICLR Workshop on Neural Network Weights as a New Data Modality (2025). *quant-ph arXiv:2503.17140*.
- Excited states and phase transitions
  - Medvidović, M., Orfi, A., Carrasquilla, J., Sels, D. arXiv:2510.15030
- Compute non-stabiliserness
  - **Spriggs**, T., Ahmadi, A., Chen, B. and Greplova, E., 2025. *Machine Learning: Science and Technology*, 6(1), p.015042

# Time evolution

- Time evolution of the quantum state  $\rightarrow$  time evolution of parameters
- $|\Psi(t)\rangle = |\Phi_{\theta(t)}\rangle$
- Many approaches
  - TVMC, pTVMC, TreTVMC
- TVMC equates a small Trotter evolution of the true state to a Taylor expansion of the variational state:
  - $|\Psi(t + \delta t)\rangle = |\Phi_{\theta(t+\delta t)}\rangle \rightarrow \Psi(t)(1 - iH\delta t) = \Phi_{\theta(t)}(1 + \dot{\theta}t)$
  - Minimise Fubini-Study distance between these two expansions, solve for  $\dot{\theta}$