

Neural Network Ground States of Lattice Gauge Theories

An introduction and overview

Thomas Spriggs, Lattice 2026, Maryland, USA

Quick disclaimer:

References are a guide for further reading

Incomplete

Neural Quantum States (NQS)

Carleo, G., & Troyer, M. (2017). Solving the quantum many-body problem with artificial neural networks. *Science*, 355(6325), 602-606.

Neural Quantum States (NQS)

- Simulate lattice gauge theories in **real, continuous** time

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Neural Quantum States (NQS)

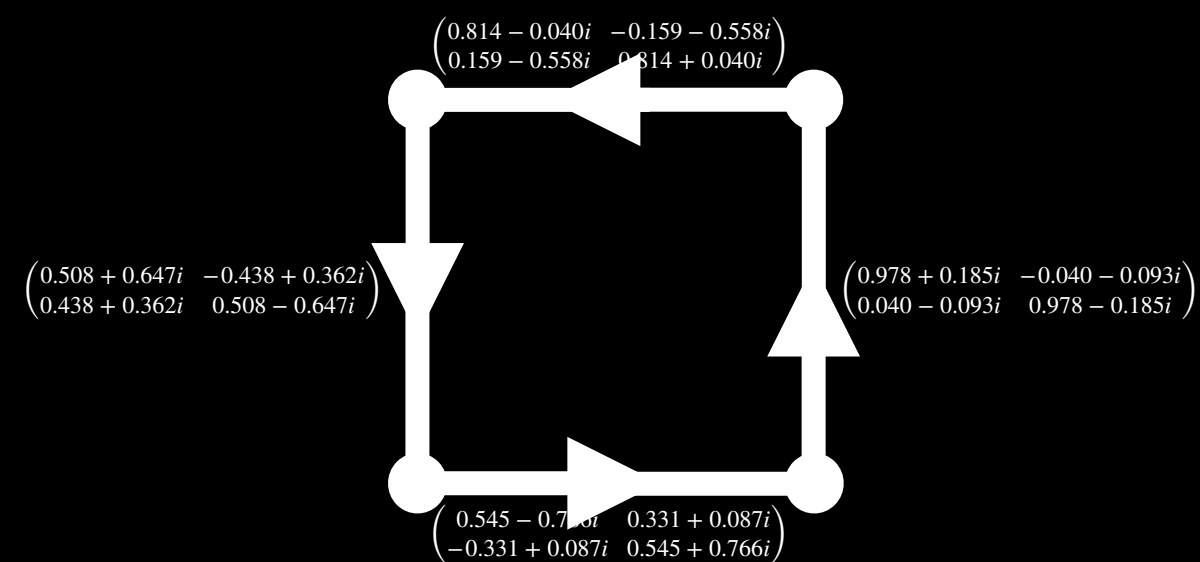
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- Already done for: $\mathbf{Z}_2$, $U(1)$, $SU(2)$ lattice gauge theories

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 - Hamiltonian framework
- Don't sample from the action
 - Learn the wavefunction
 - No sign problem*
- Already done for: $\mathbf{Z}_2$, $U(1)$, $SU(2)$ lattice gauge theories
 - Commonplace in condensed matter physics and quantum chemistry

Carleo, G., & Troyer, M. (2017). Solving the quantum many-body problem with artificial neural networks. *Science*, 355(6325), 602-606.

NQS workflow for LGTs



NN

$|\Psi_\theta\rangle$

NQS workflow for LGTs

NN
 $|\Psi_\theta\rangle$

$$|\langle U_0 | \Psi_\theta \rangle|^2$$

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NN
 $|\Psi_\theta\rangle$

$$|\langle U_0 | \Psi_\theta \rangle|^2 \quad |\langle U_1 | \Psi_\theta \rangle|^2$$

NQS workflow for LGTs

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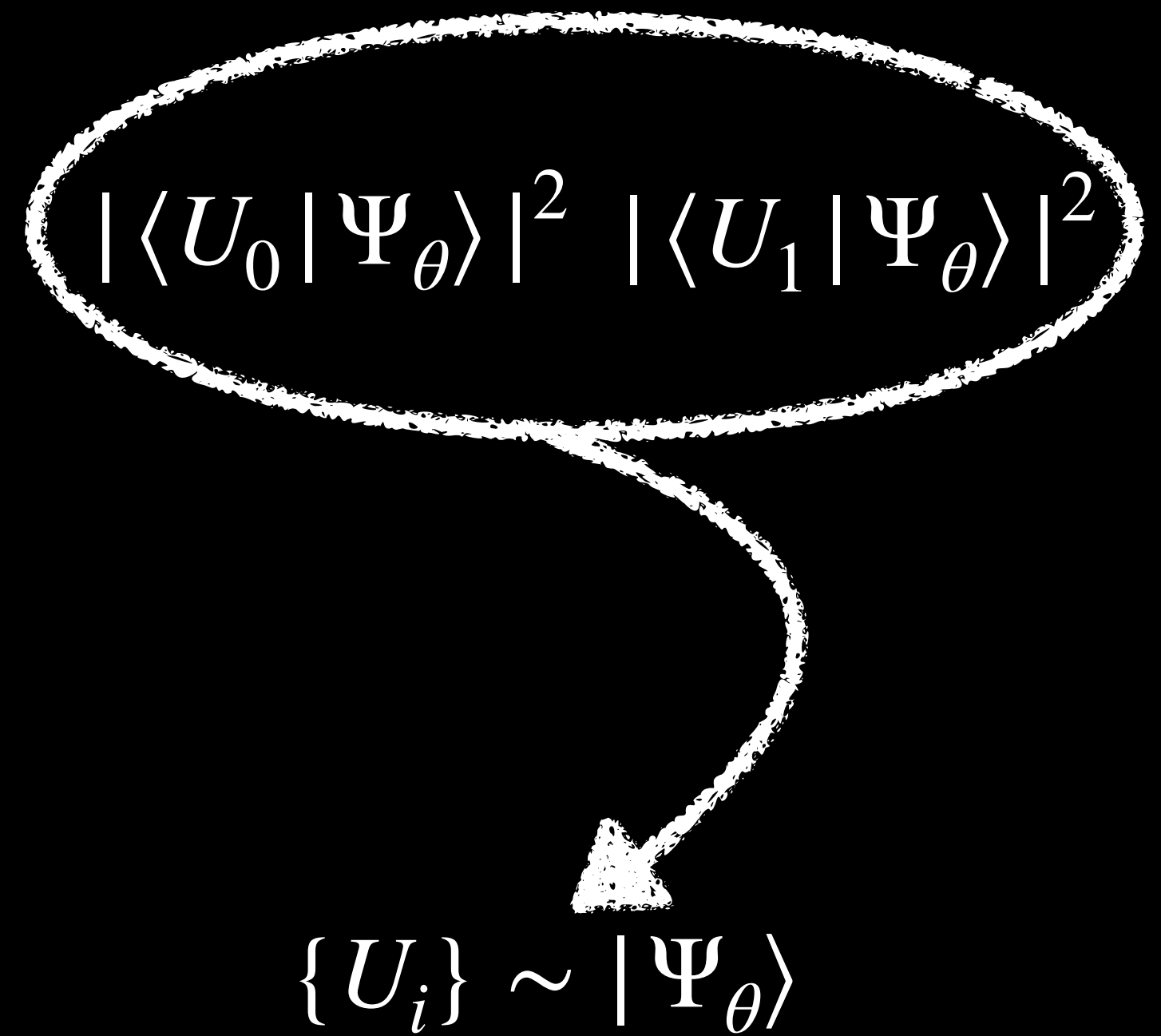
Accept/reject

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$$|\langle U_0 | \Psi_\theta \rangle|^2 \quad |\langle U_1 | \Psi_\theta \rangle|^2$$

$$\{U_i\} \sim |\Psi_\theta\rangle$$

$$E \approx \frac{1}{N} \sum_{U_i \sim |\langle \Psi_\theta | U \rangle|^2} \mathcal{H}_L(U_i)$$

NQS workflow for LGTs

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 $|\Psi_\theta\rangle$

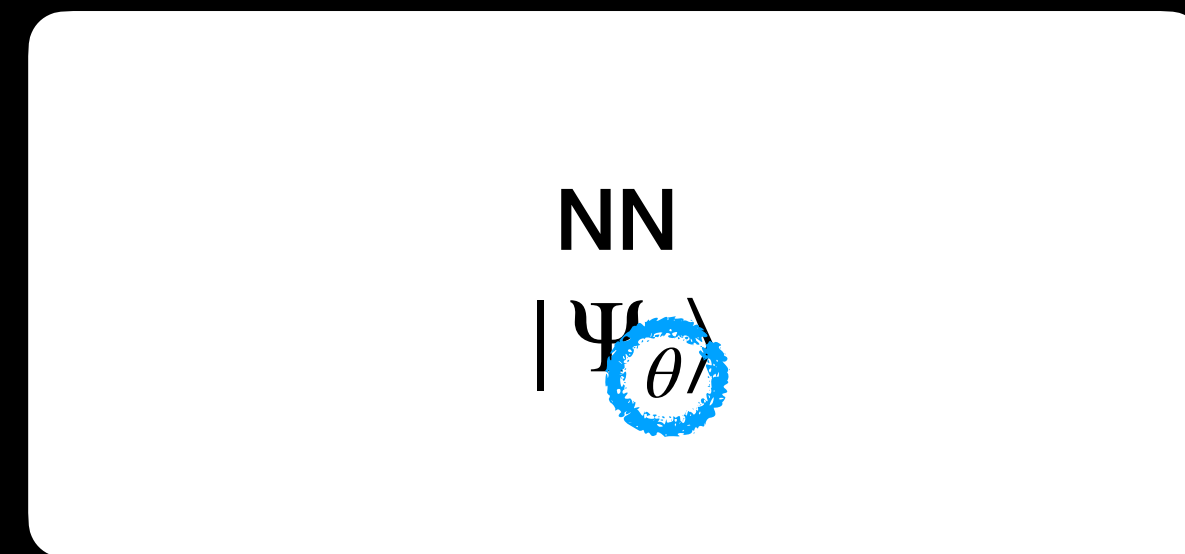
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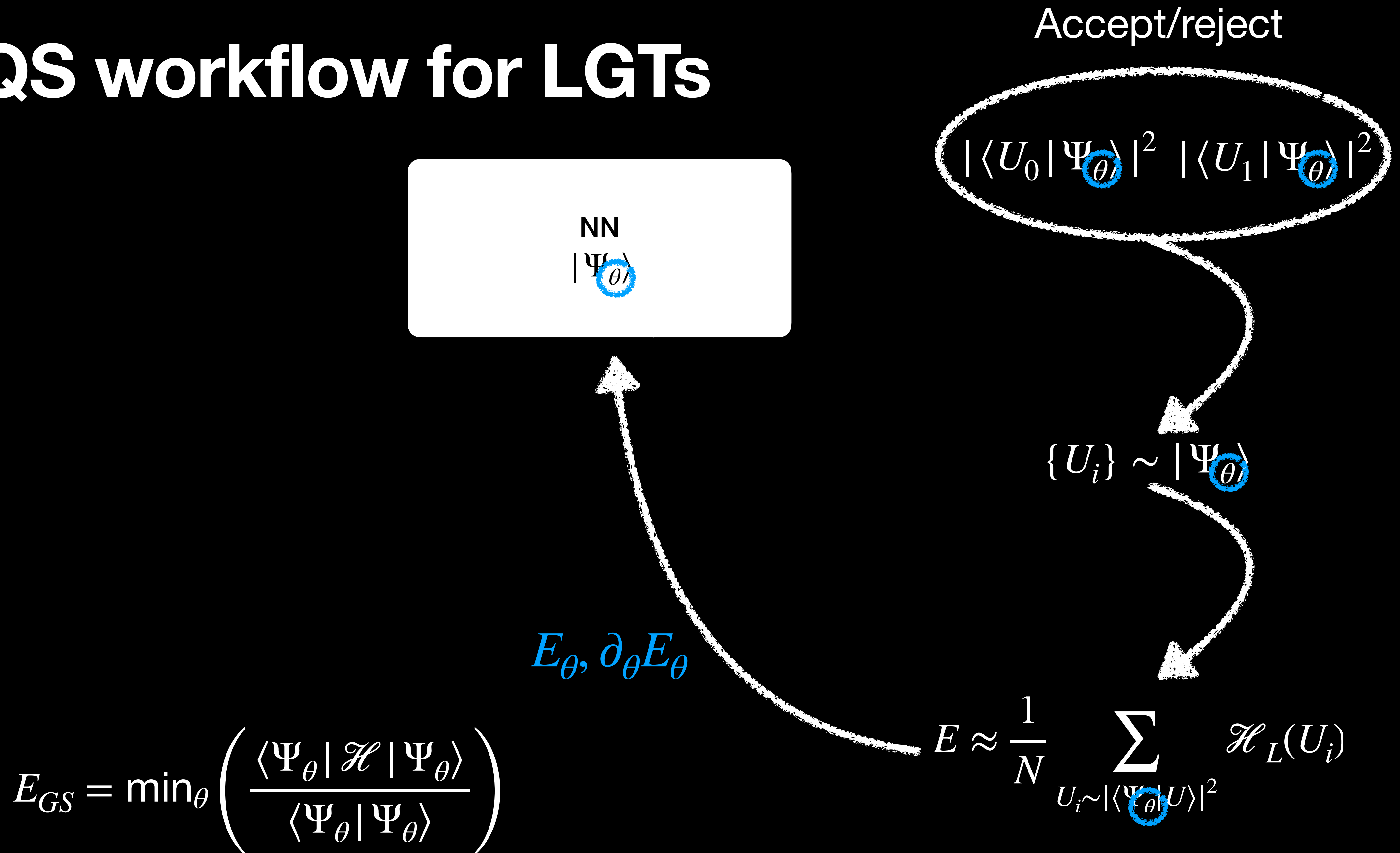
$$|\langle U_0 | \Psi_{\theta} \rangle|^2 \quad |\langle U_1 | \Psi_{\theta} \rangle|^2$$

$$\{U_i\} \sim |\Psi_{\theta}\rangle$$

$$E_{\theta}, \partial_{\theta} E_{\theta}$$

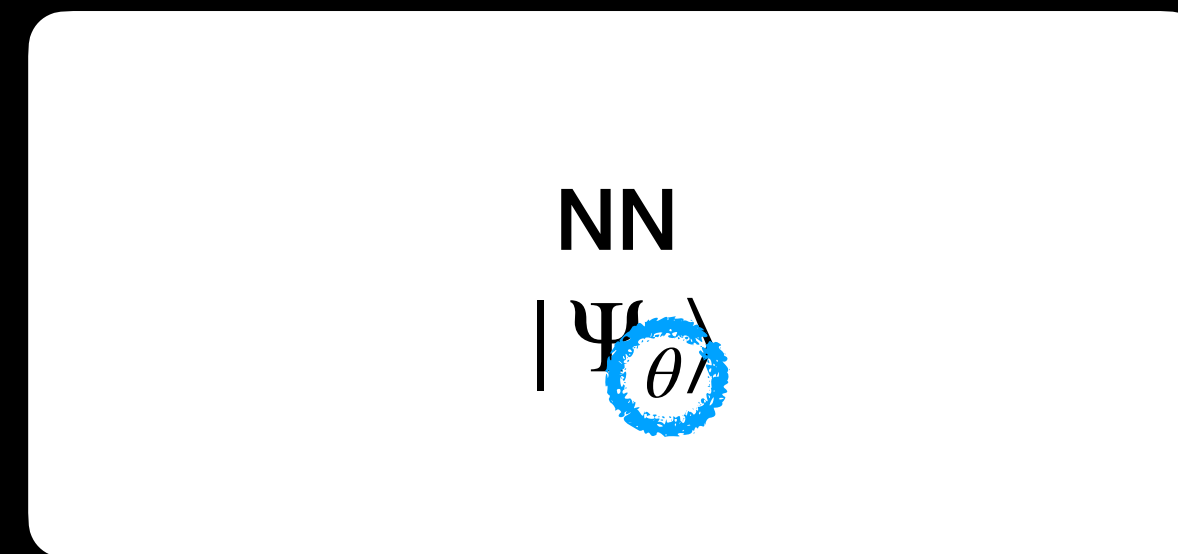
$$E \approx \frac{1}{N} \sum_{U_i \sim |\langle \Psi_{\theta} | U \rangle|^2} \mathcal{H}_L(U_i)$$

NQS workflow for LGTs



NQS workflow for LGTs

Accept/reject



$$|\langle U_0 | \Psi_\theta \rangle|^2 |\langle U_1 | \Psi_\theta \rangle|^2$$

$$\{U_i\} \sim |\Psi_\theta\rangle$$

Task: find function that returns ground state wavefunction amplitudes

$$E_\theta, \partial_\theta E_\theta$$

$$E_{GS} = \min_\theta \left(\frac{\langle \Psi_\theta | \mathcal{H} | \Psi_\theta \rangle}{\langle \Psi_\theta | \Psi_\theta \rangle} \right)$$

$$E \approx \frac{1}{N} \sum_{U_i \sim |\langle \Psi_\theta | U \rangle|^2} \mathcal{H}_L(U_i)$$

Lagrangians and Hamiltonians

Lagrangian vs Hamiltonian formalism

From an NQS perspective

Lagrangian vs Hamiltonian formalism

From an NQS perspective

- Continuum Dirac equation
 - Observables found through the path integral

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$$\int dx O(x) p(x) \approx \frac{1}{N} \sum_{x \sim p(x)} O(x)$$

Lagrangian vs Hamiltonian formalism

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$$\int dx O(x) p(x) \approx \frac{1}{N} \sum_{x \sim p(x)} O(x)$$

$$\int D[U] O[U] e^{S_E[U]} \approx \frac{1}{N} \sum_{U \sim e^{S_E[U]}} O[U]$$

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$$\int dx O(x) p(x) \approx \frac{1}{N} \sum_{x \sim p(x)} O(x)$$

Lagrangian vs Hamiltonian formalism

Trade off

Lose access to real time dynamics

Get a known expression for the probability distribution of gauge fields

Complex actions ruin this: e^{S_E} is no longer a probability

Lagrangian vs Hamiltonian formalism

Trade off

Keep access to real time dynamics

Have to learn the probability distribution (using a neural network)

Probability distribution is always real

Neural (network) wavefunctions

A quick aside into spin systems (easier to picture)

Some definitions: wavefunction

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- Let $|\Psi\rangle = \sum_i a_i |x_i\rangle$

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 - 2 qubits, spin basis: $|\Psi\rangle = a_0 |00\rangle + a_1 |01\rangle + a_2 |10\rangle + a_3 |11\rangle$

Some definitions: wavefunction

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 - “Finding ground state” == finding four* numbers

Some definitions: **variational** wavefunction

- Let $|\Psi\rangle = \sum_i a_i |x_i\rangle$
 - 2 qubits, spin basis: $|\Psi\rangle = a_0 |00\rangle + a_1 |01\rangle + a_2 |10\rangle + a_3 |11\rangle$
 - “Finding ground state” == finding **a function that gives** four* numbers

$$|\Psi_\theta\rangle = \sum_i f_\theta(x_i) |x_i\rangle$$

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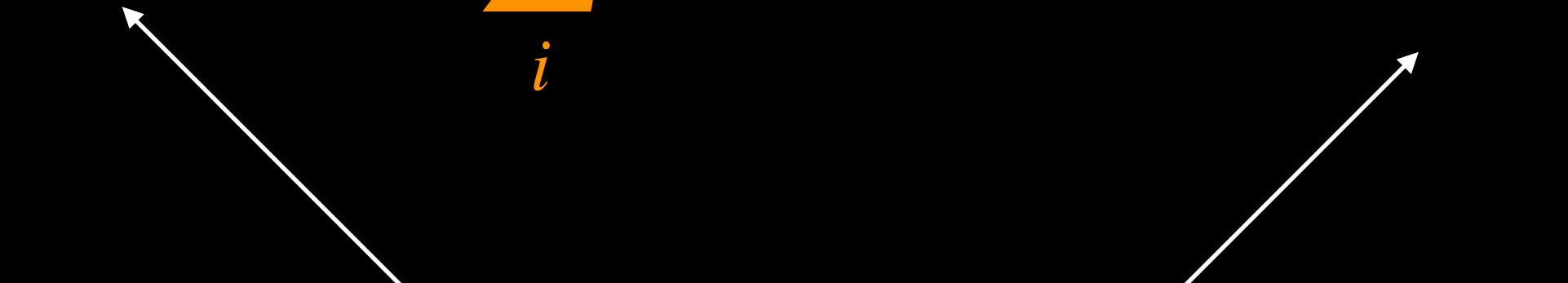
$$\langle x_0 | \Psi_\theta \rangle = \sum_i \langle x_0 | f_\theta(x_i) | x_i \rangle = f_\theta(x_0)$$

Some definitions: **variational** wavefunction

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“Configuration in, amplitude out”

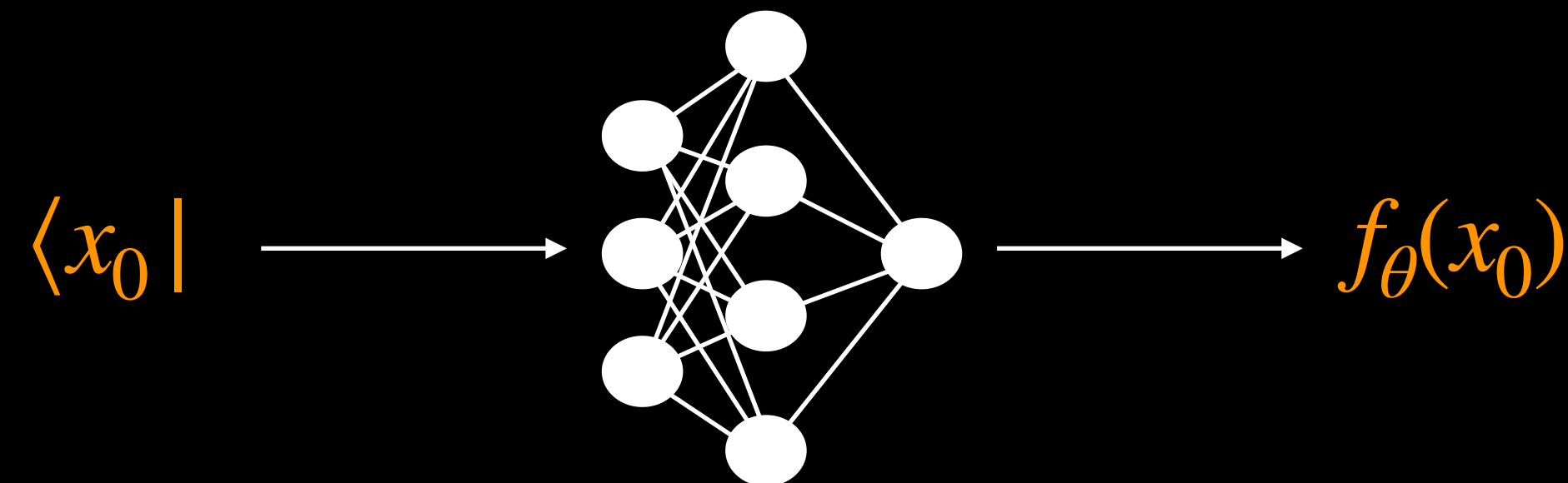


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“Configuration in, amplitude out”



Wavefunctions in gauge theories

Finding the probability measure

Take an un-normalised variational state, $|\Psi_\theta\rangle$, a Hilbert space, $\{|\vec{x}\rangle\}$, and a Hamiltonian, $\mathcal{H}$

$$E = \langle \mathcal{H} \rangle_\theta = \frac{\langle \Psi_\theta | \mathcal{H} | \Psi_\theta \rangle}{\langle \Psi_\theta | \Psi_\theta \rangle}$$

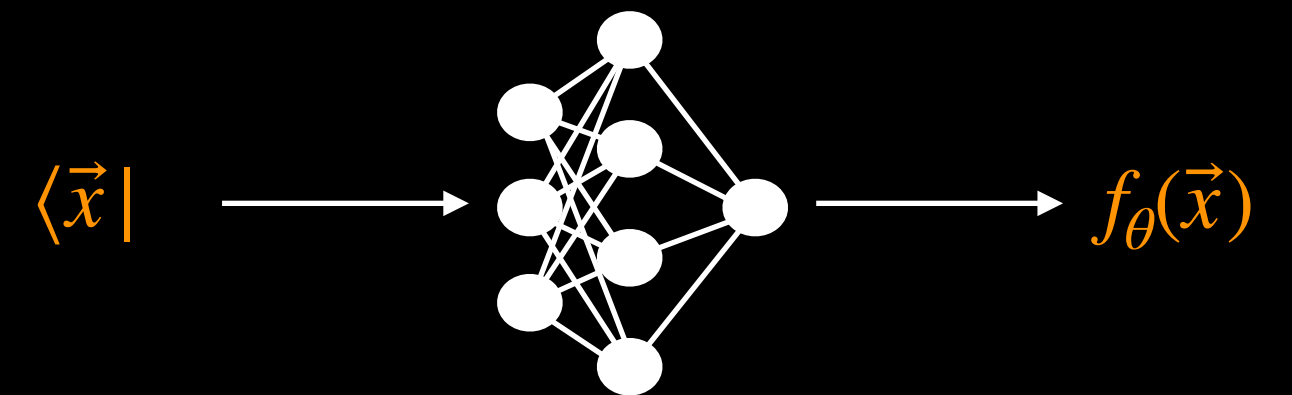
Introduce a complete set of states $\sum_{\vec{x}} |\vec{x}\rangle \langle \vec{x}|$

$$E = \frac{\sum_{\vec{x}} \langle \Psi_\theta | \vec{x} \rangle \langle \vec{x} | \mathcal{H} | \Psi_\theta \rangle}{\sum_{\vec{x}'} \langle \Psi_\theta | \vec{x}' \rangle \langle \vec{x}' | \Psi_\theta \rangle}$$

Multiply top and bottom by a factor of $\langle \vec{x} | \Psi_\theta \rangle$

$$E = \frac{\sum_{\vec{x}} |\langle \Psi_\theta | \vec{x} \rangle|^2 \langle \vec{x} | \mathcal{H} | \Psi_\theta \rangle}{\sum_{\vec{x}'} |\langle \Psi_\theta | \vec{x}' \rangle|^2 \langle \vec{x}' | \Psi_\theta \rangle} = \frac{\sum_{\vec{x}} p(\vec{x}) \mathcal{H}_L(\vec{x})}{\sum_{\vec{x}'} p(\vec{x}')} \approx \frac{1}{N} \sum_{\vec{x} \sim |\langle \Psi_\theta | \vec{x} \rangle|^2} \mathcal{H}_L(x)$$

Finding the probability measure



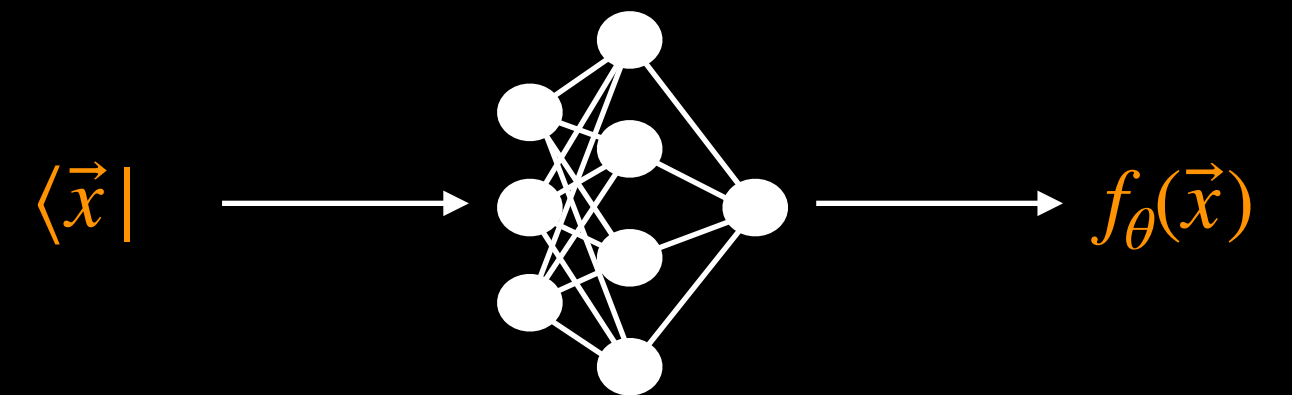
- We can use importance sampling
 - Draw samples according to

$$| \langle \Psi_{\theta} | \vec{x} \rangle |^2$$

- Measure **local** value of the observable on **each sample**

$$\mathcal{H}_L = \frac{\langle \vec{x} | \mathcal{H} | \Psi_{\theta} \rangle}{\langle \vec{x} | \Psi_{\theta} \rangle}$$

Finding the probability measure



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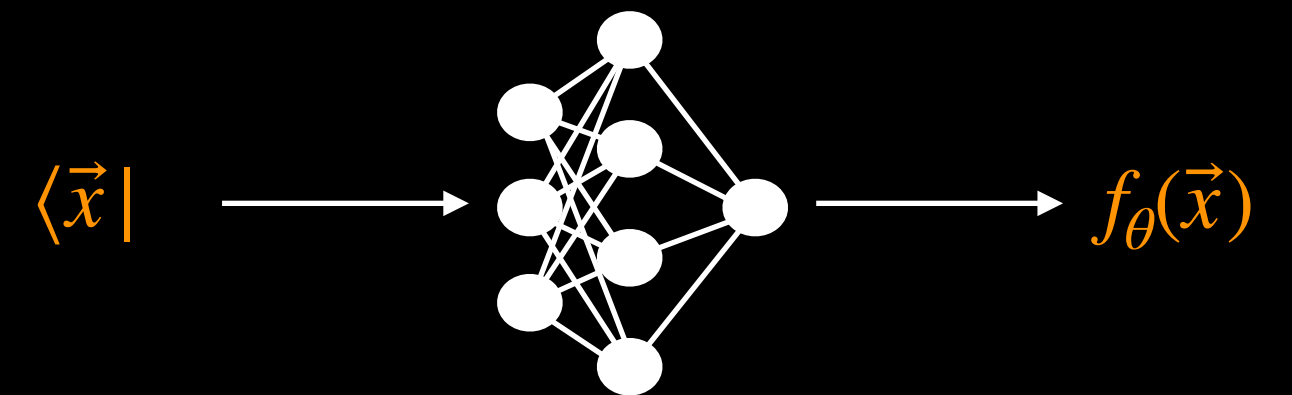
$$|\langle \Psi_{\theta} | \vec{x} \rangle|^2$$

Always real; no sign problem

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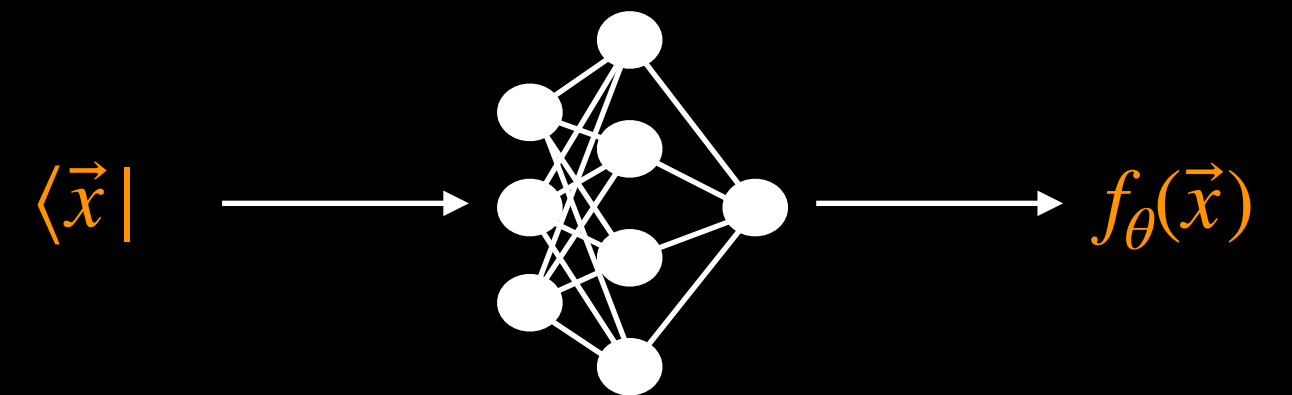
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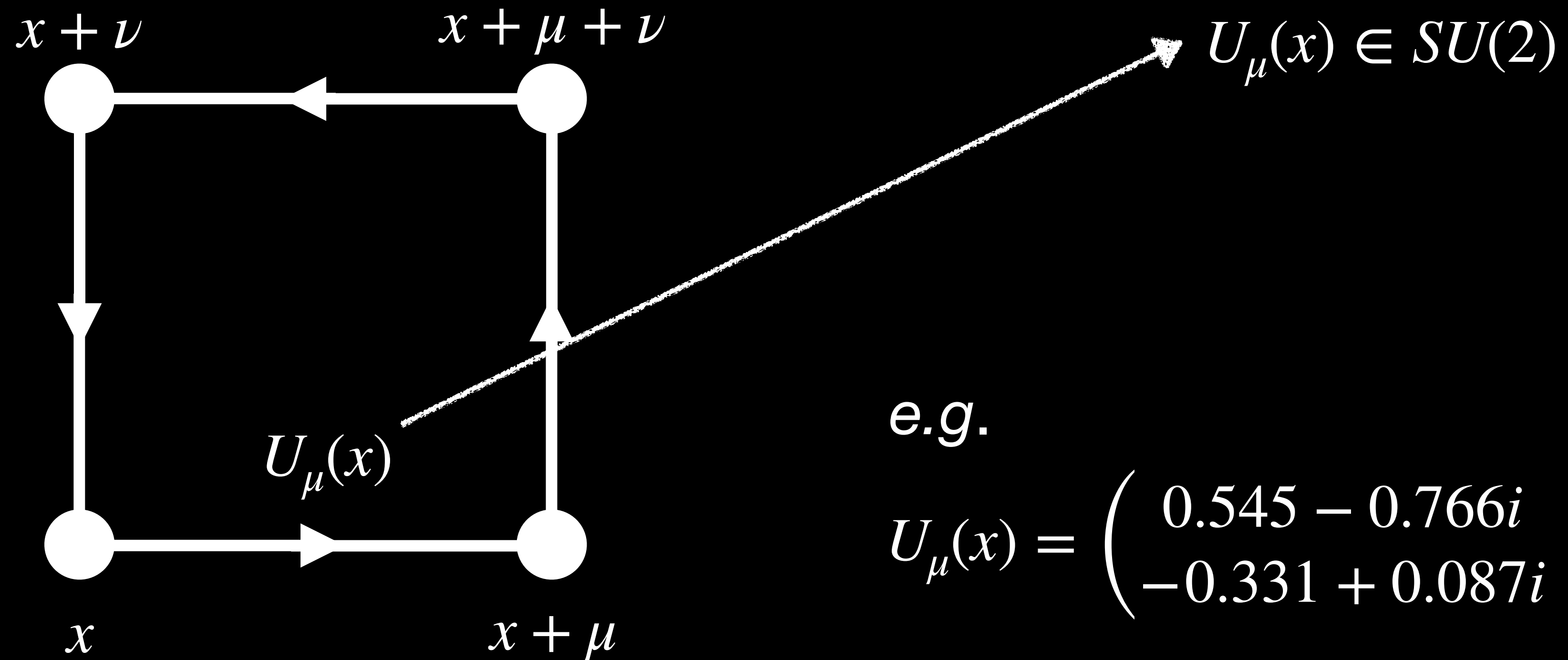
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$$\mathcal{H}_L = \frac{\langle \vec{x} | \mathcal{H} | \Psi_{\theta} \rangle}{\langle \vec{x} | \Psi_{\theta} \rangle}$$

$$\text{Energy} \approx \frac{1}{N} \sum_{\vec{x} \sim |\langle \Psi_{\theta} | \vec{x} \rangle|^2} \mathcal{H}_L(x)$$

Gauge-theory-aware neural networks

SU(2) example

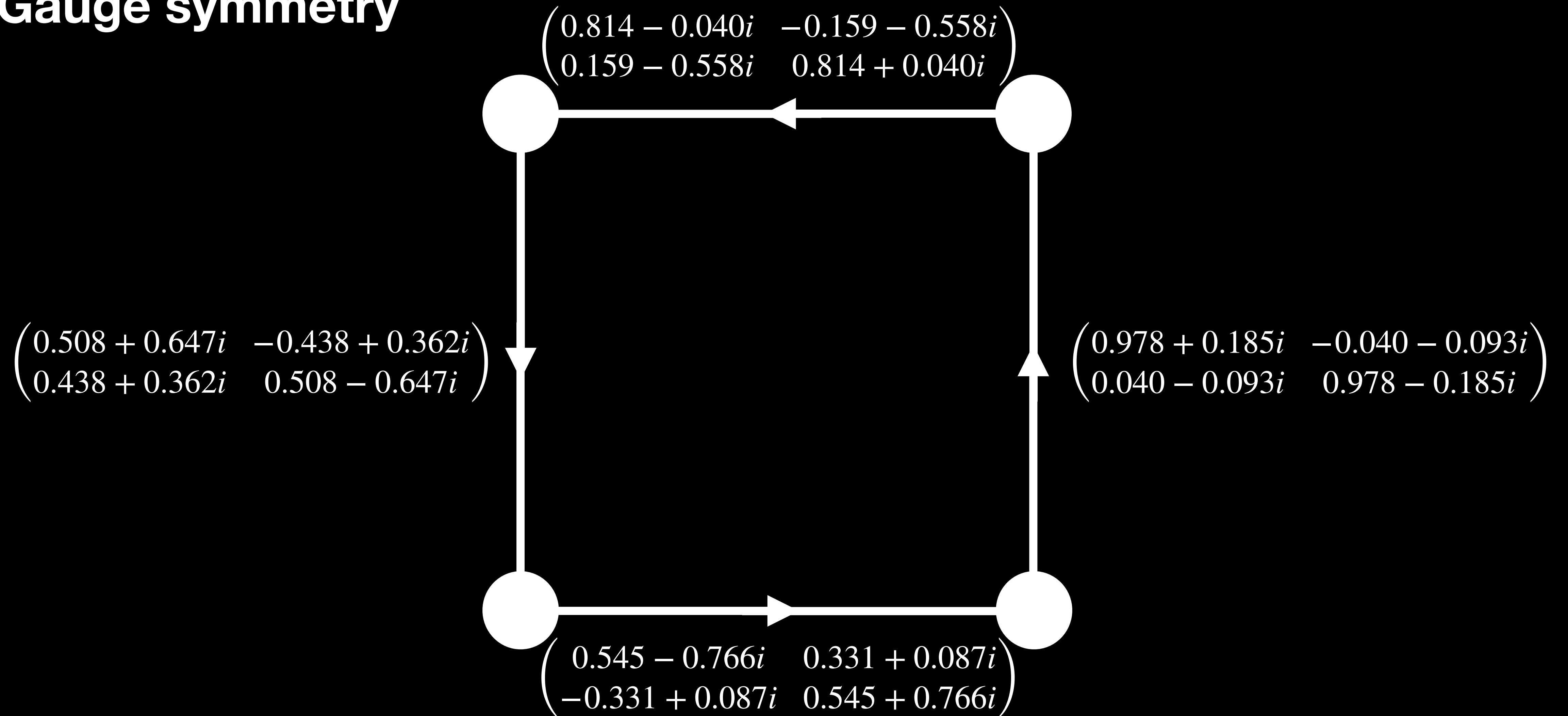


e.g.

$$U_\mu(x) = \begin{pmatrix} 0.545 - 0.766i & 0.331 + 0.087i \\ -0.331 + 0.087i & 0.545 + 0.766i \end{pmatrix}$$

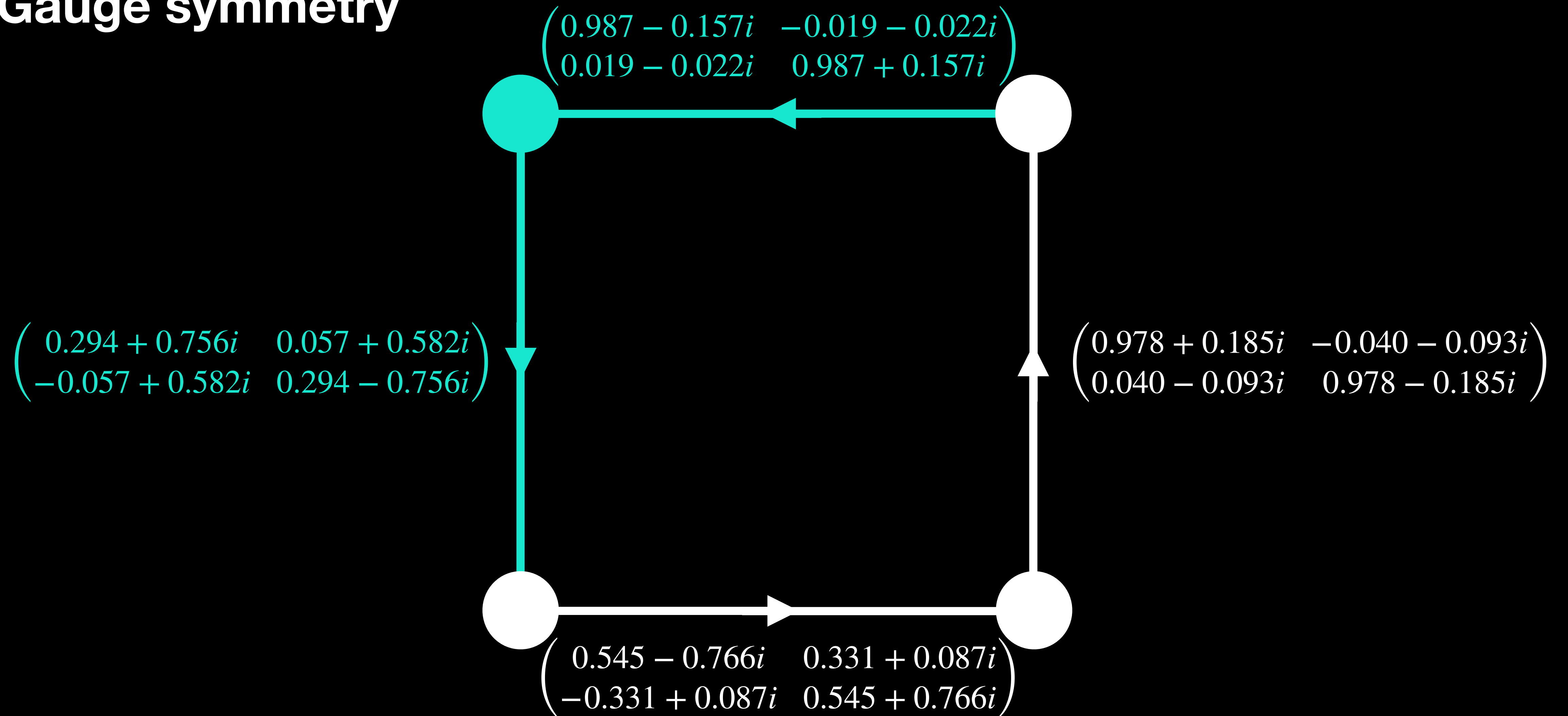
SU(2) example

Gauge symmetry



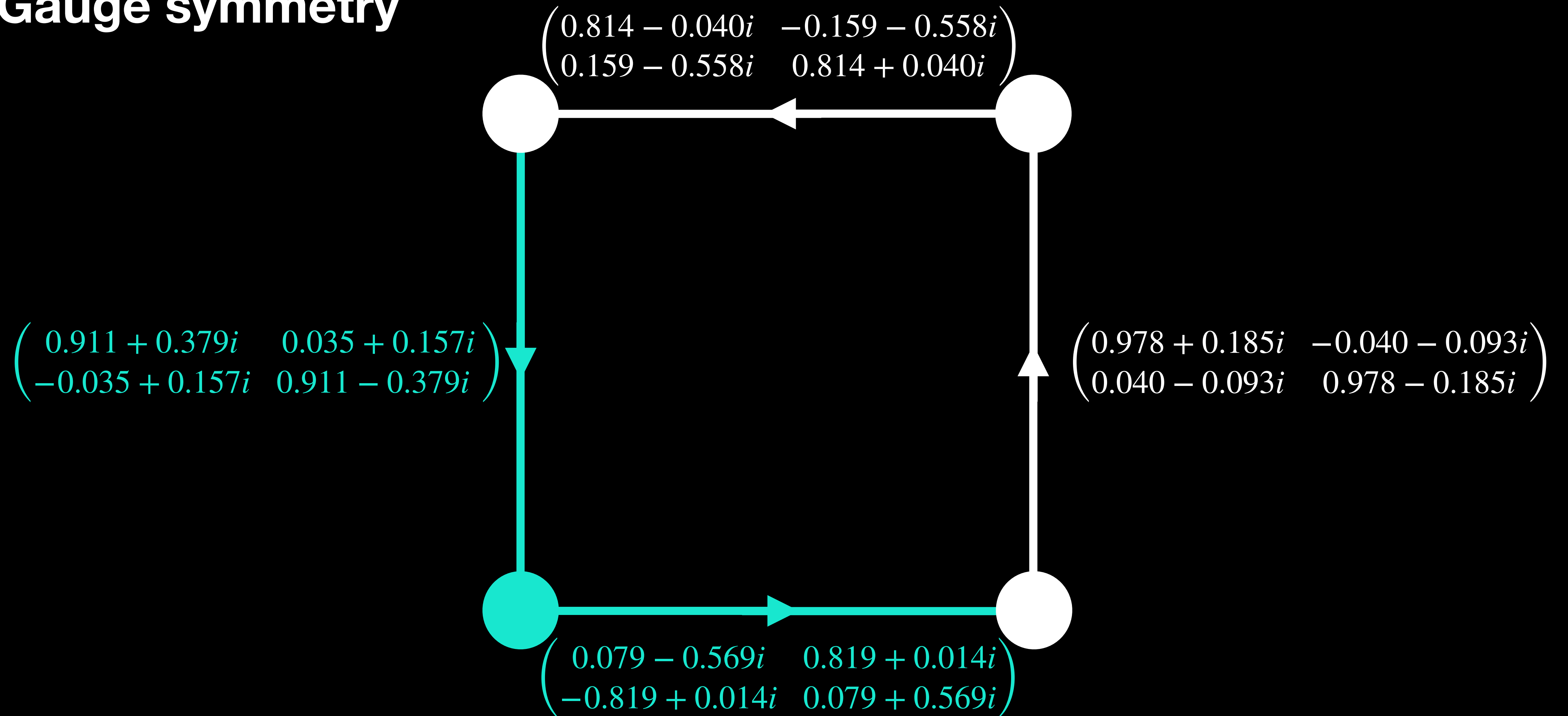
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Gauge symmetry



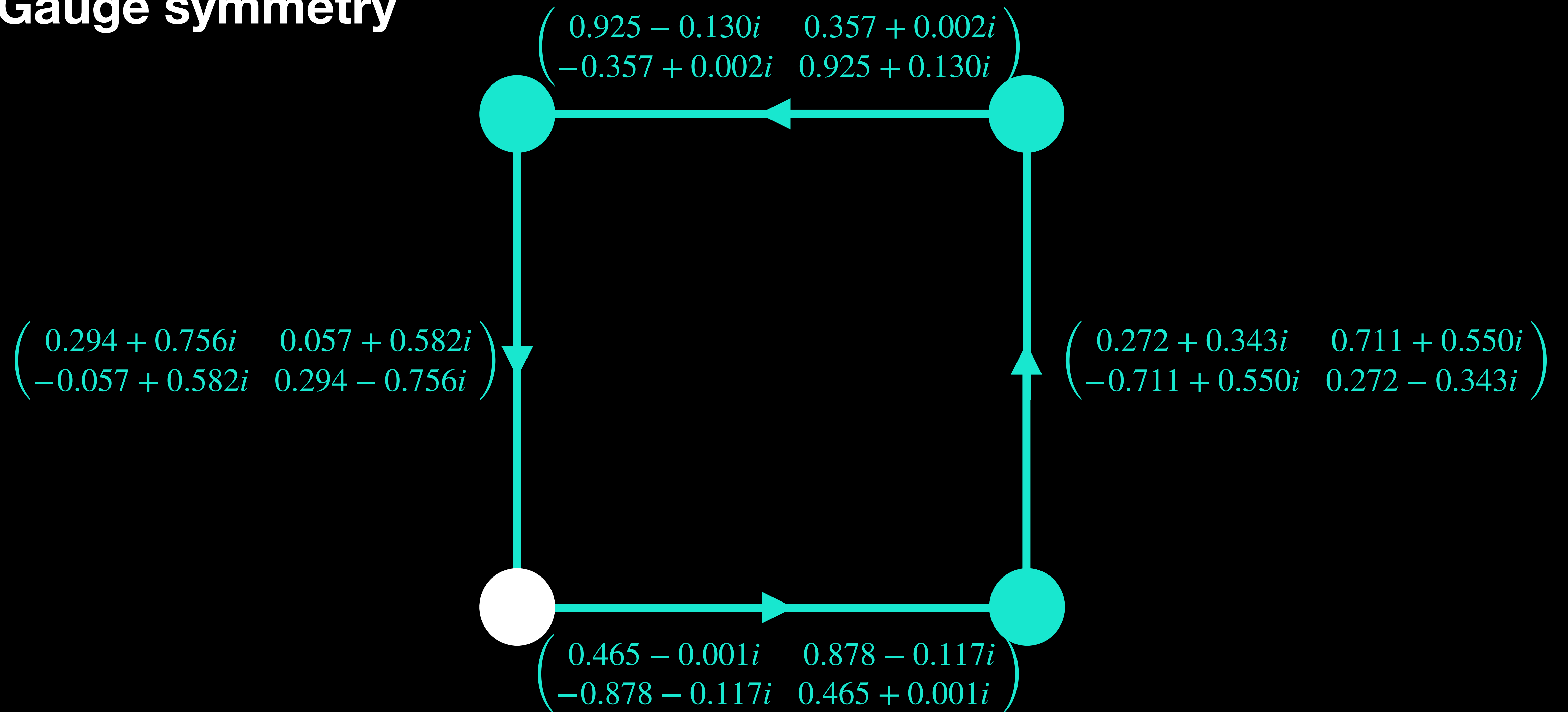
SU(2) example

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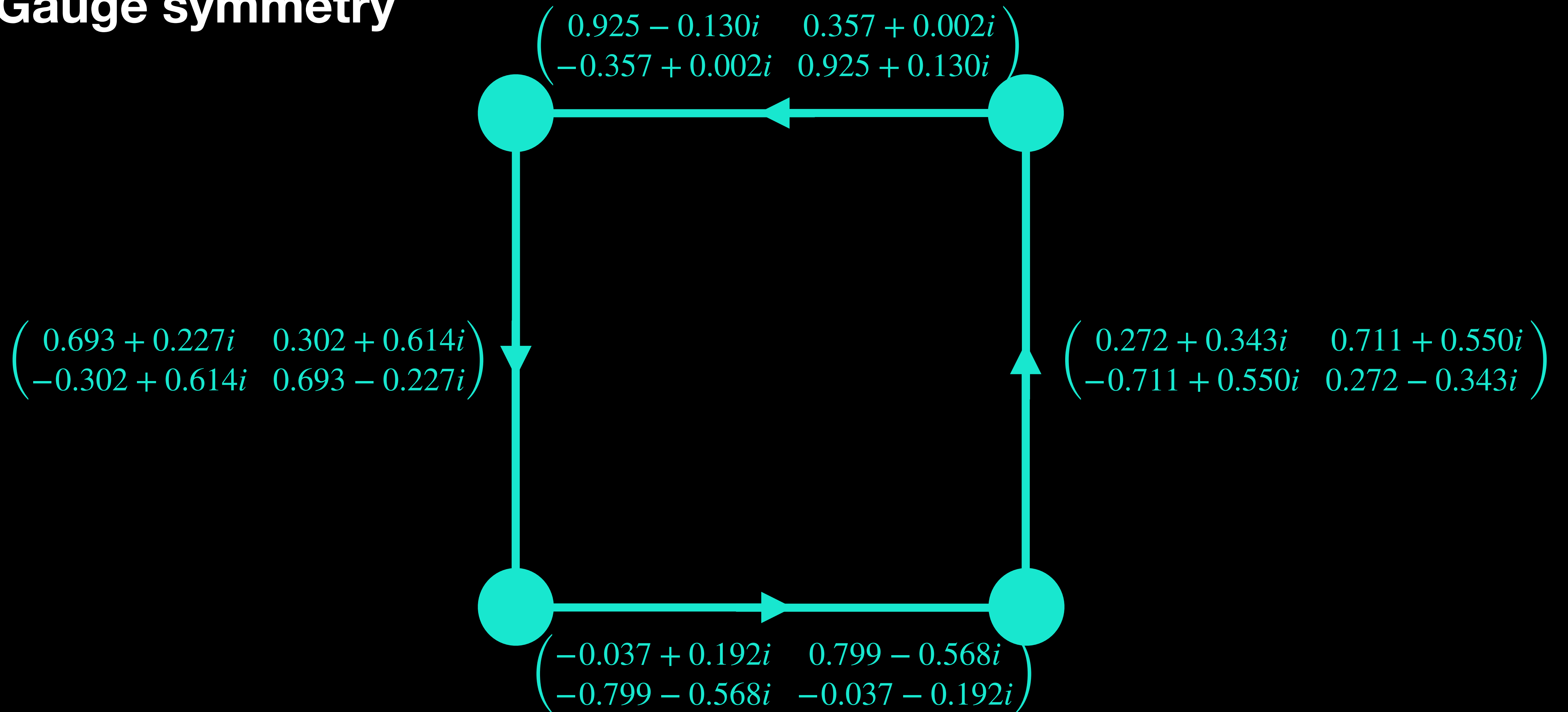
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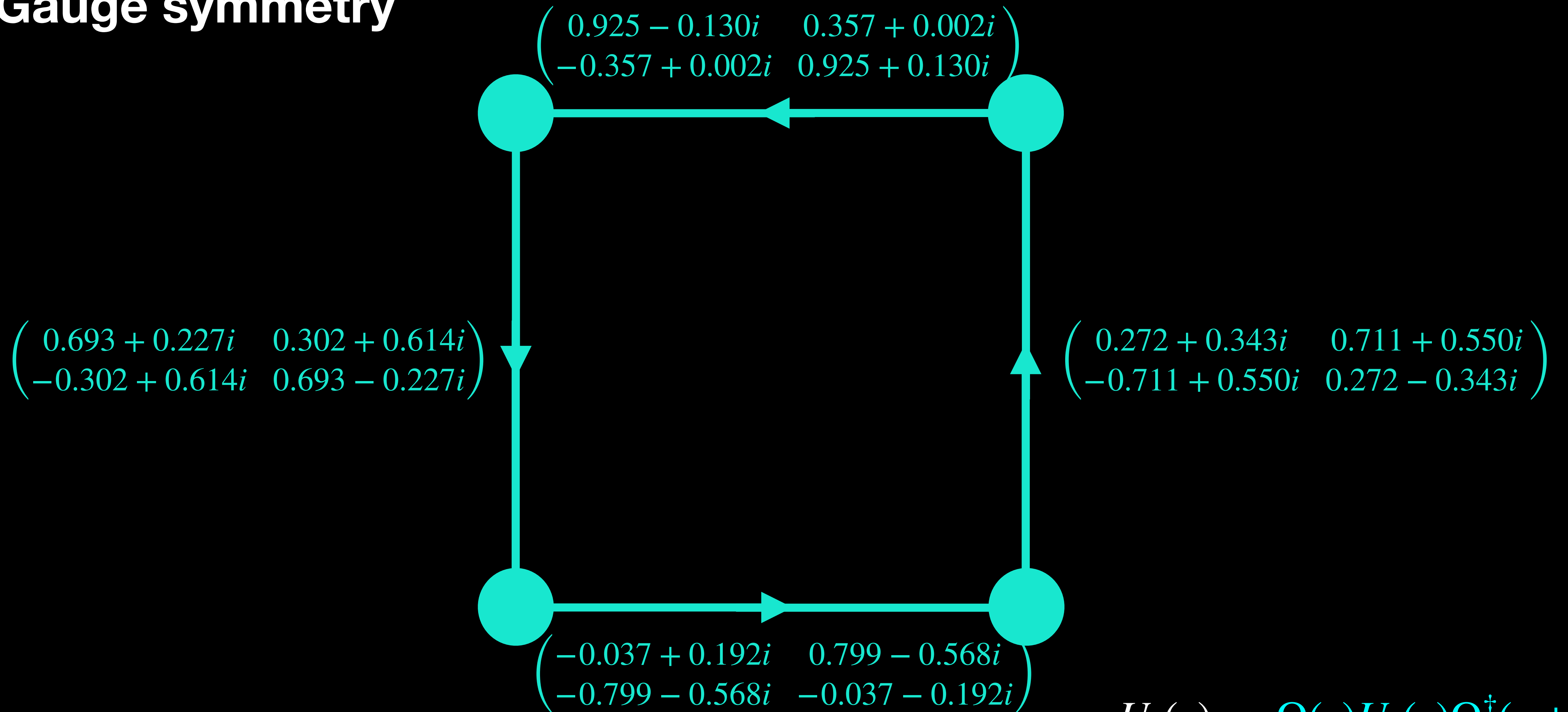
SU(2) example

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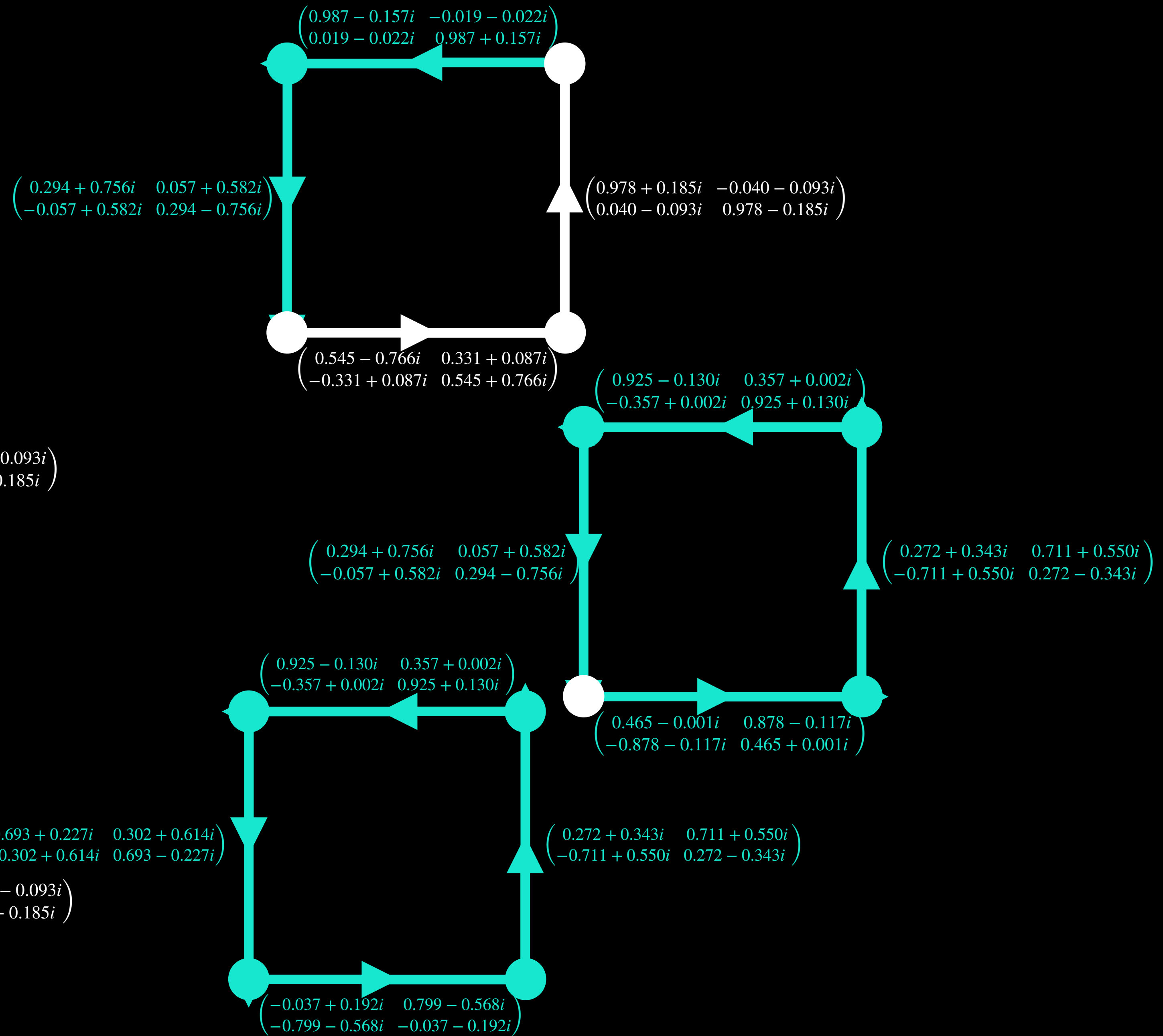
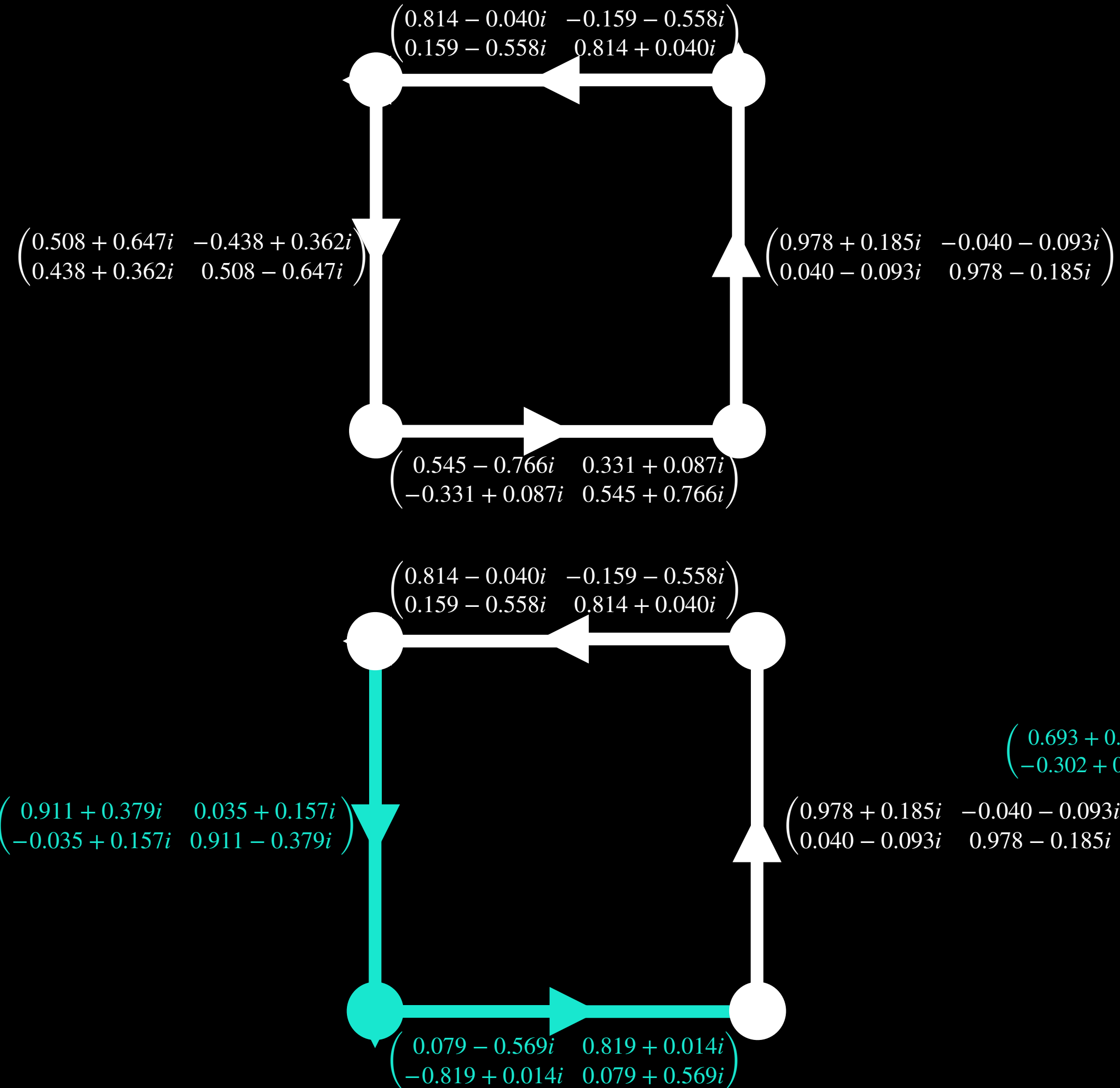
Gauge symmetry



$$U_\mu(x) \rightarrow \Omega(x) U_\mu(x) \Omega^\dagger(x + \mu)$$

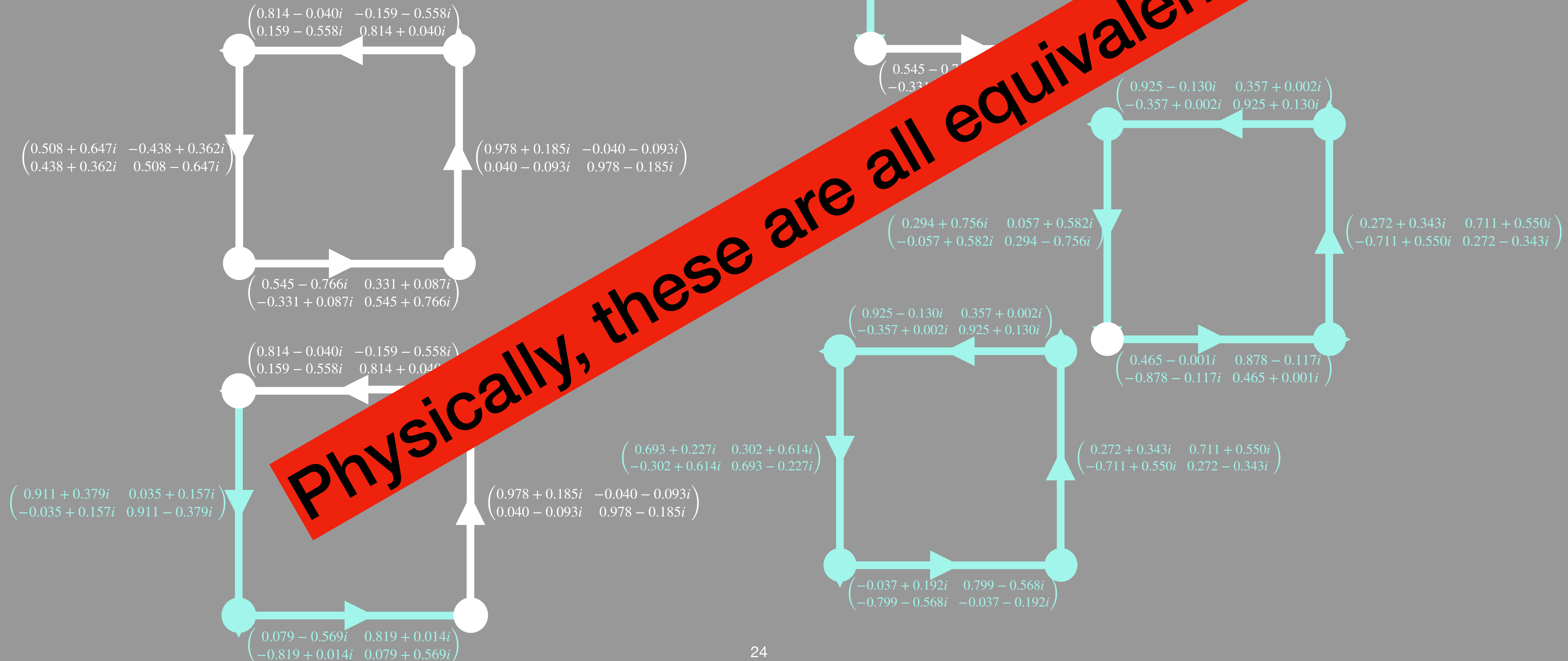
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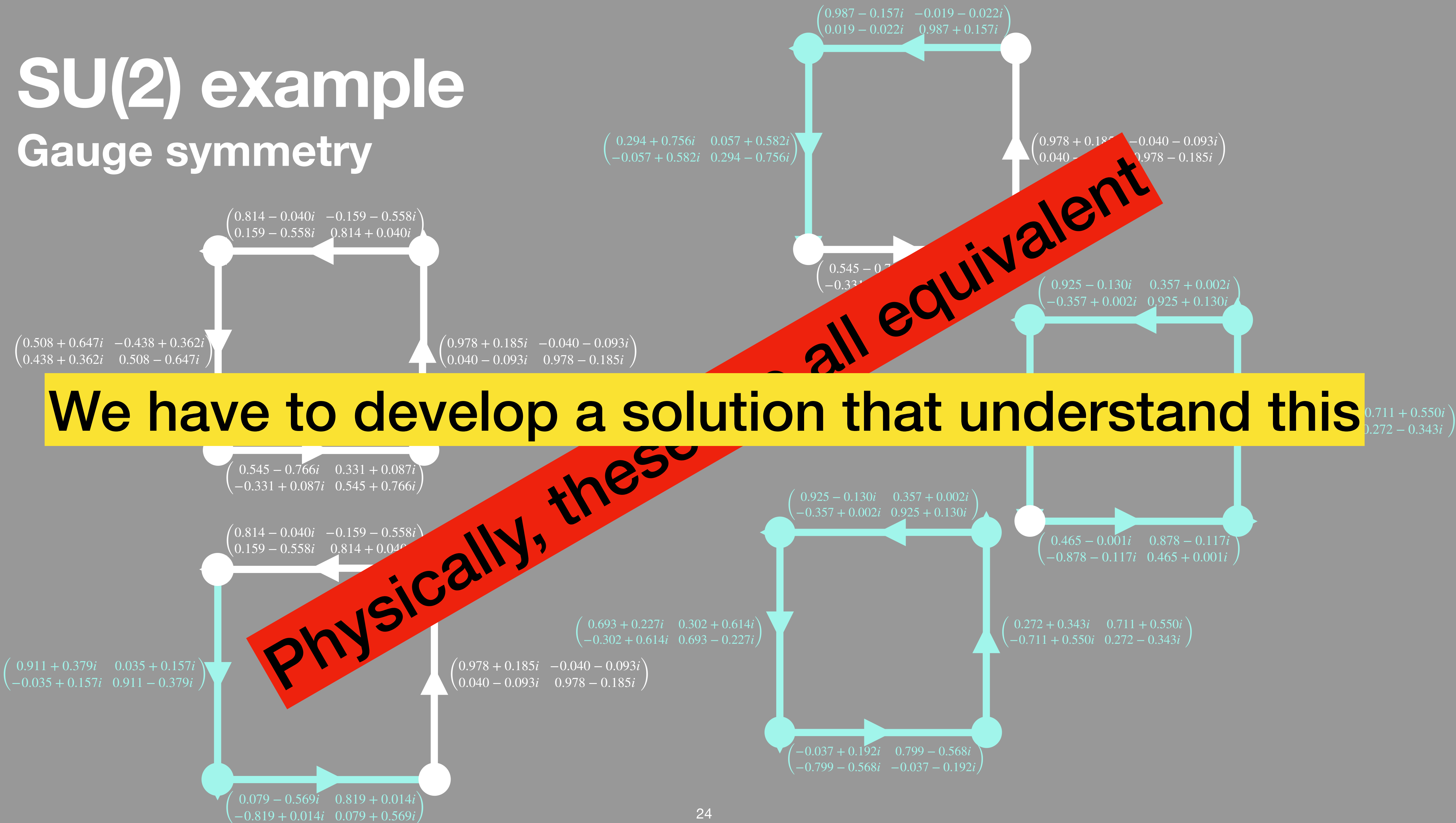
SU(2) example

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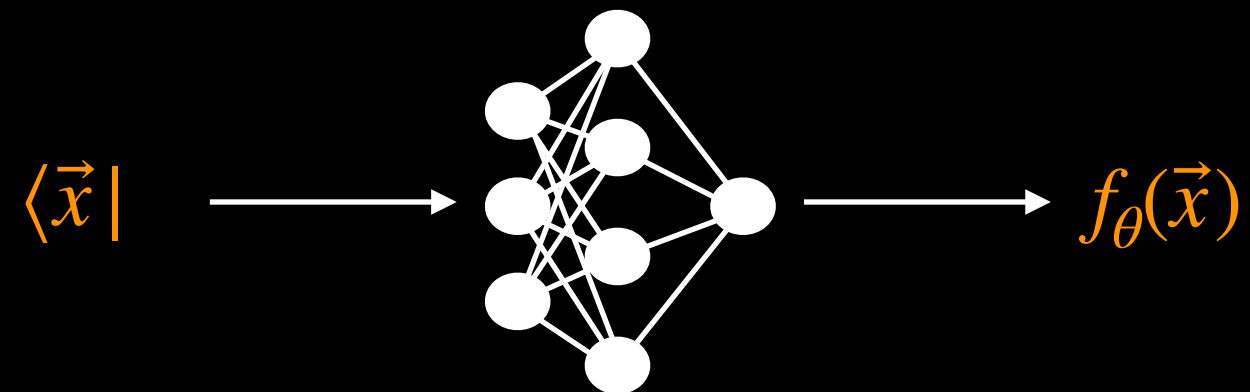
SU(2) example

Gauge symmetry



Some ansatze

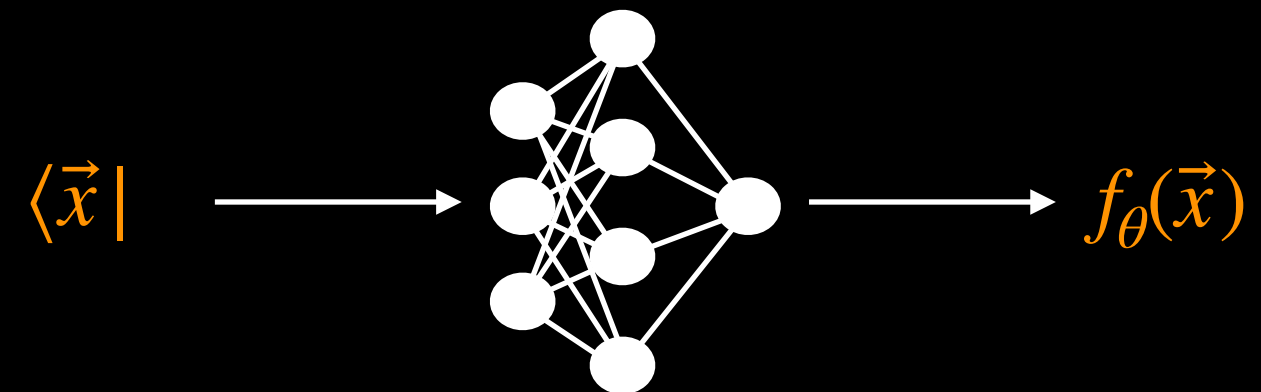
Physics-informed



Greensite, J. P. (1979). Calculation of the Yang-Mills vacuum wave functional. *Nuclear Physics B*, 158(2-3), 469-496.

Chin, S. A., Van Roosmalen, O. S., Umland, E. A., & Koonin, S. E. (1985). Exact ground-state properties of the SU (2) Hamiltonian lattice gauge theory. *Physical Review D*, 31(12), 3201.

Physics-informed



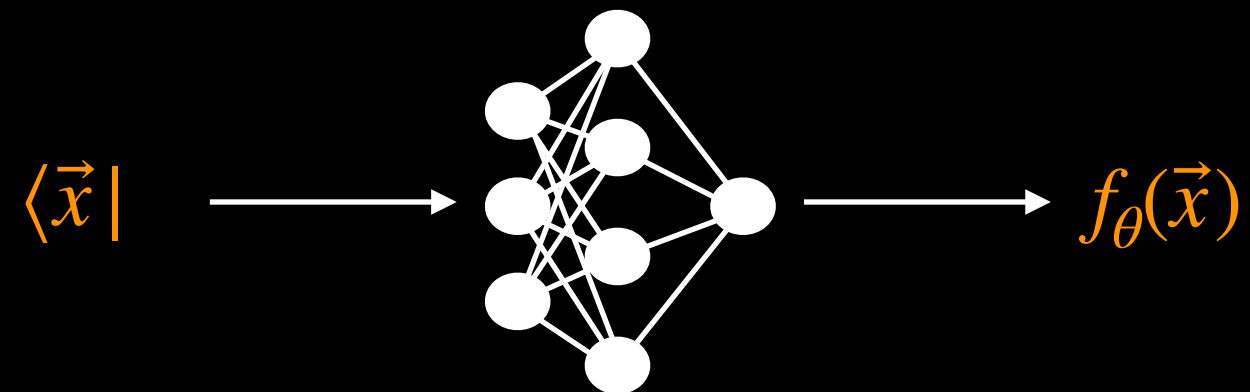
$$f(\mathbf{U}) = \prod_p e^{\alpha \frac{1}{2} \text{Tr}(P_p)}$$

$$\sum \text{Tr} \begin{array}{c} \begin{array}{ccc} \bullet & & \bullet \\ \downarrow U_{\nu}^{\dagger}(x) & & \uparrow U_{\nu}(x+\mu) \\ \bullet & & \bullet \\ \leftarrow U_{\mu}^{\dagger}(x+\nu) & & \rightarrow U_{\mu}(x) \end{array} \end{array}$$

Greensite, J. P. (1979). Calculation of the Yang-Mills vacuum wave functional. *Nuclear Physics B*, 158(2-3), 469-496.

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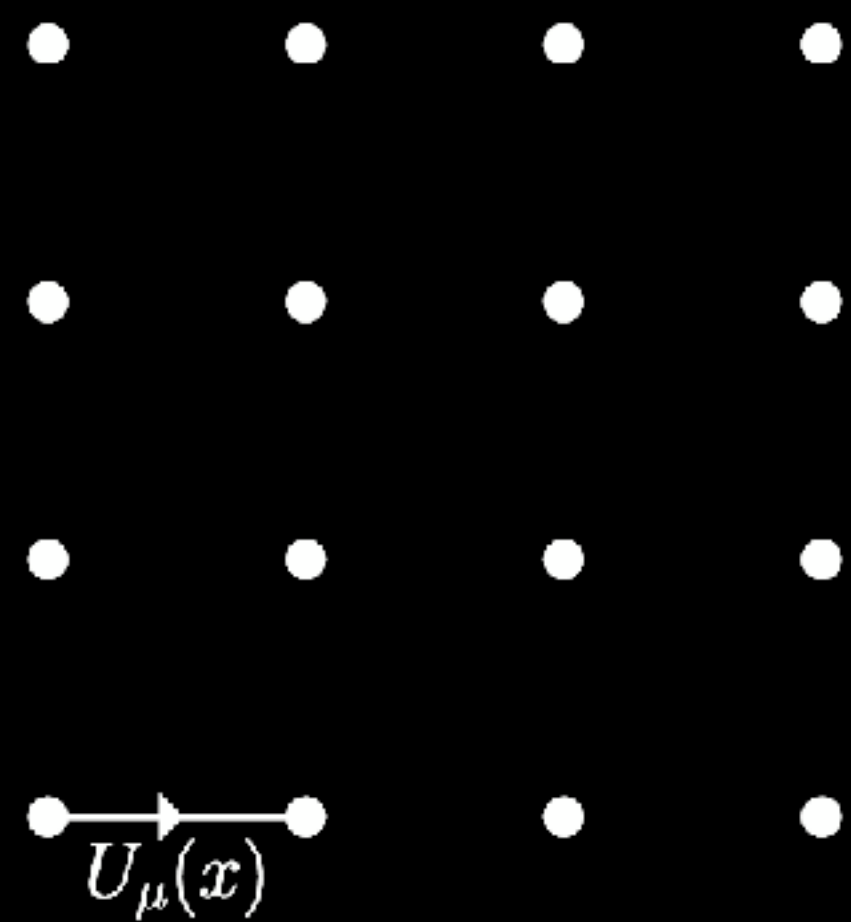
$$\sum \text{Tr} \begin{array}{c} \begin{array}{ccc} & U_{\mu}^{\dagger}(x + \nu) & \\ U_{\nu}^{\dagger}(x) \downarrow & \square & \uparrow U_{\nu}(x + \mu) \\ U_{\mu}(x) & & \end{array} \end{array}$$

Reliance on hand-crafted features

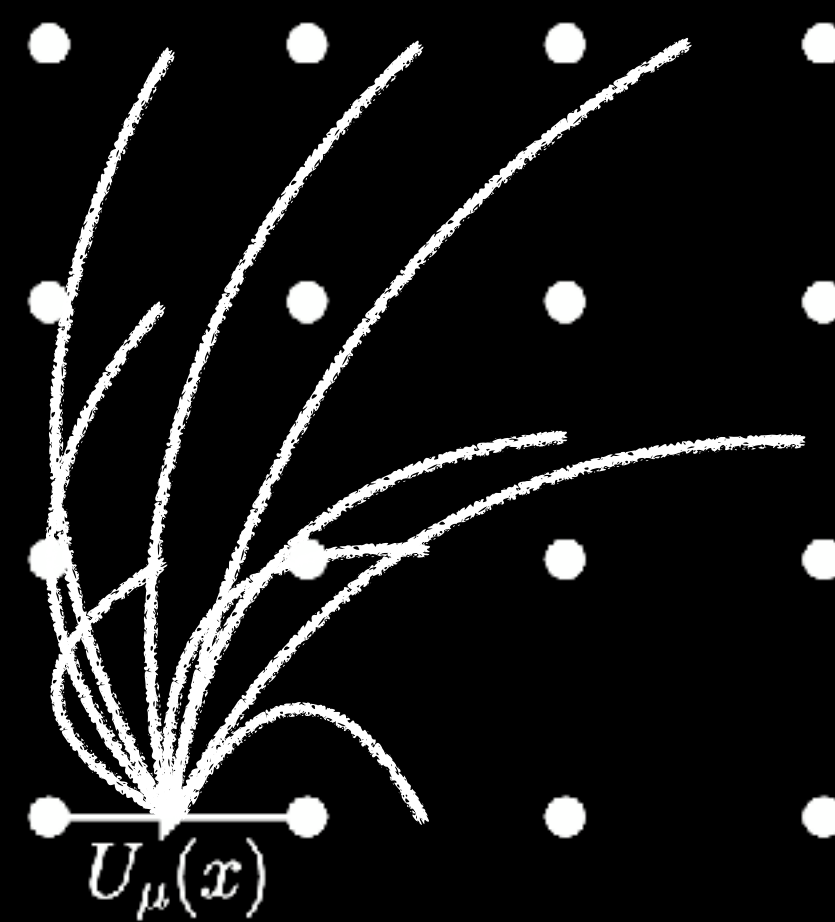
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Neural network + physics-informed

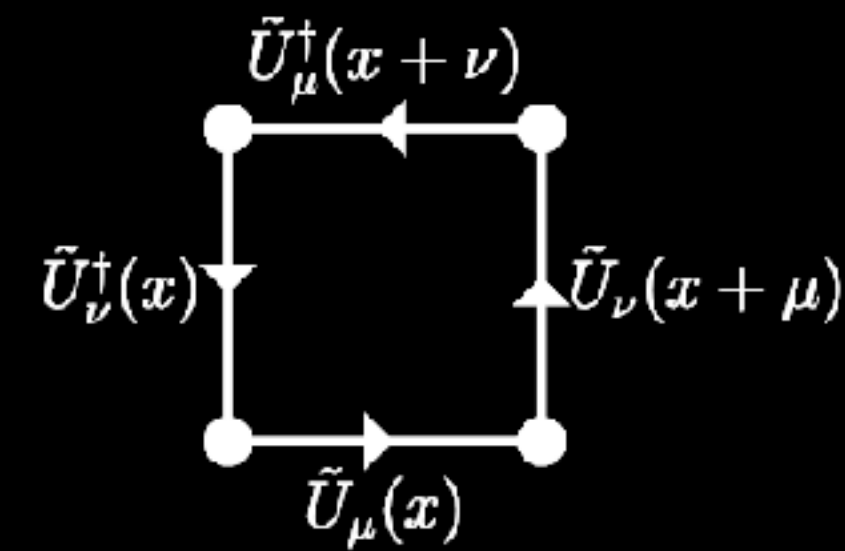


Neural
network
layers



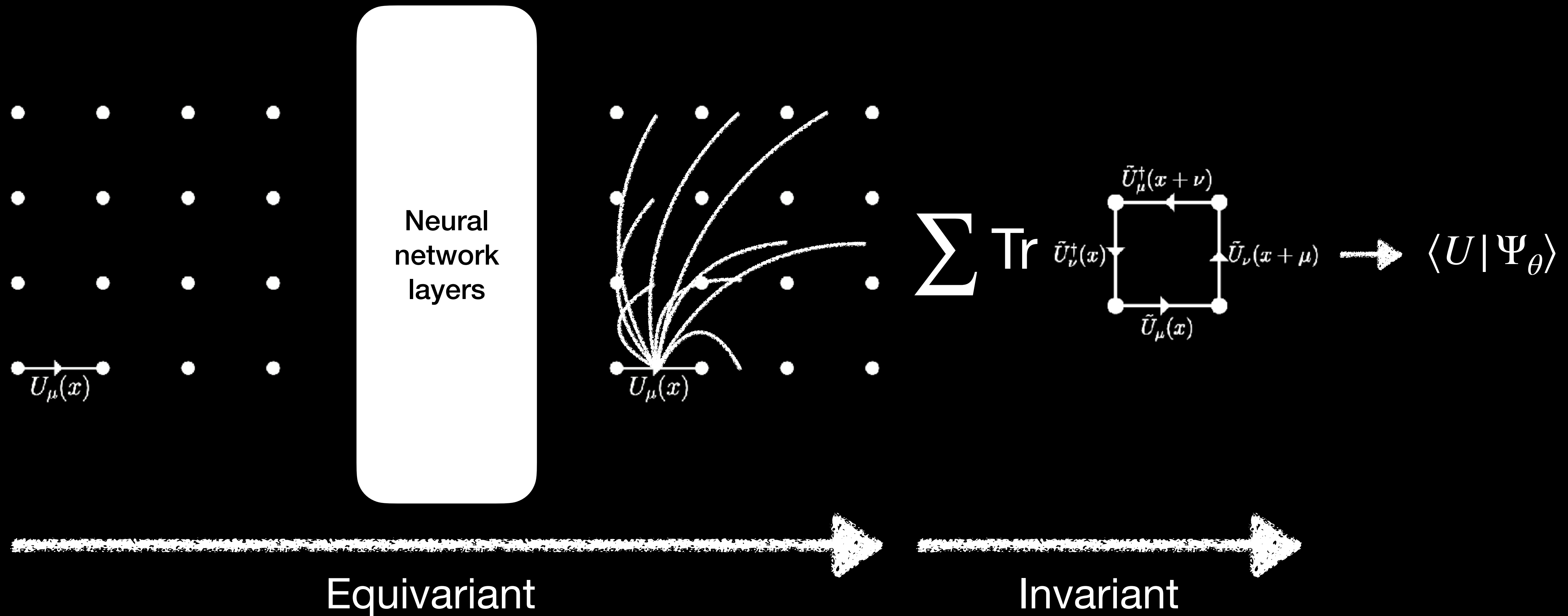
Σ

Tr

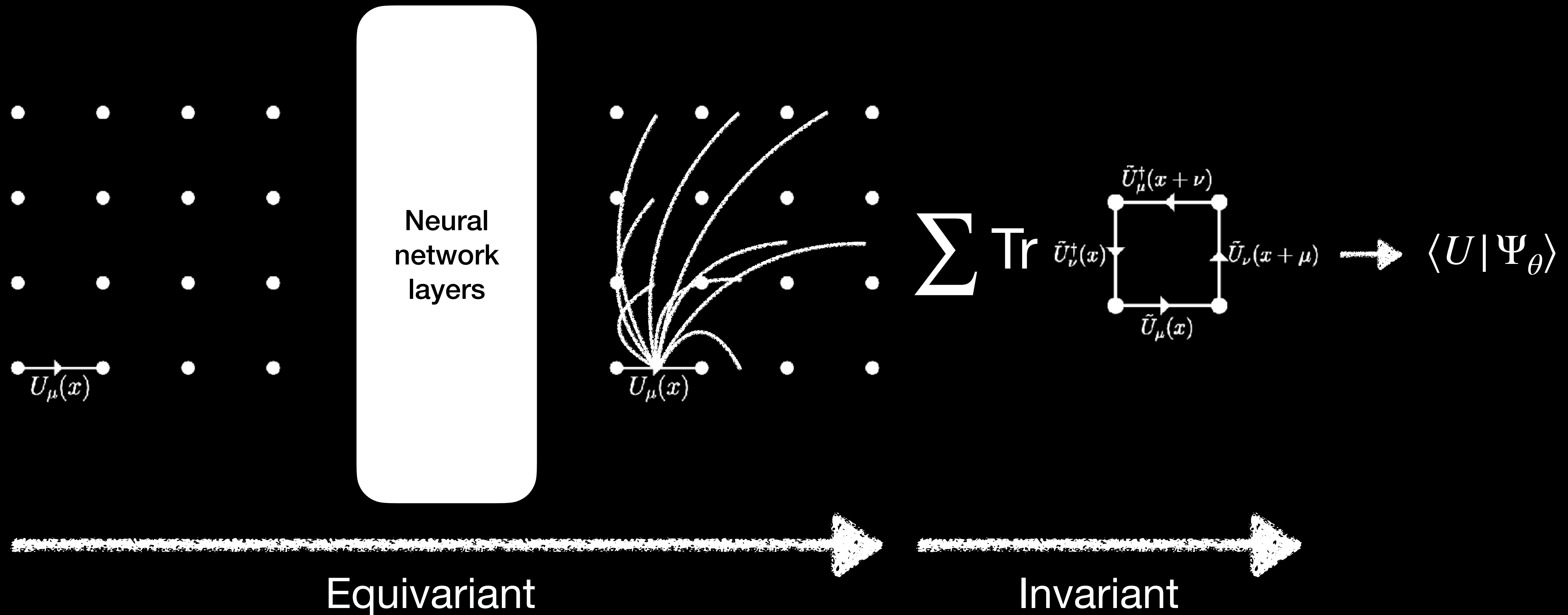


$\langle U | \Psi_\theta \rangle$

Neural network + physics-informed



Neural network + physics-informed

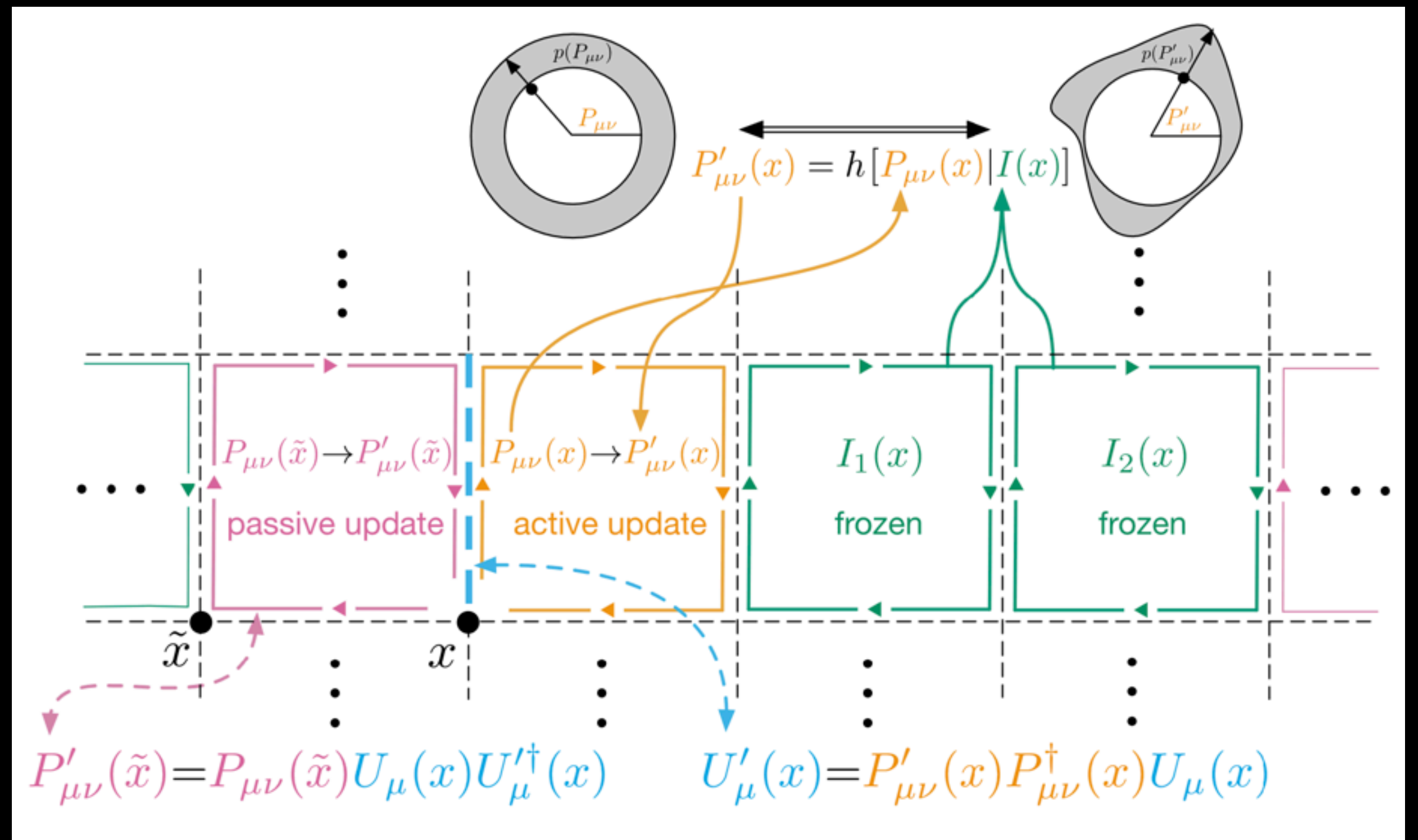


$$U_\mu(x) \rightarrow \Omega(x) U_\mu(x) \Omega^\dagger(x + \mu) \text{ before and after NN}$$

Choices for network

All gauge equivariant

- “Normalising flow networks”



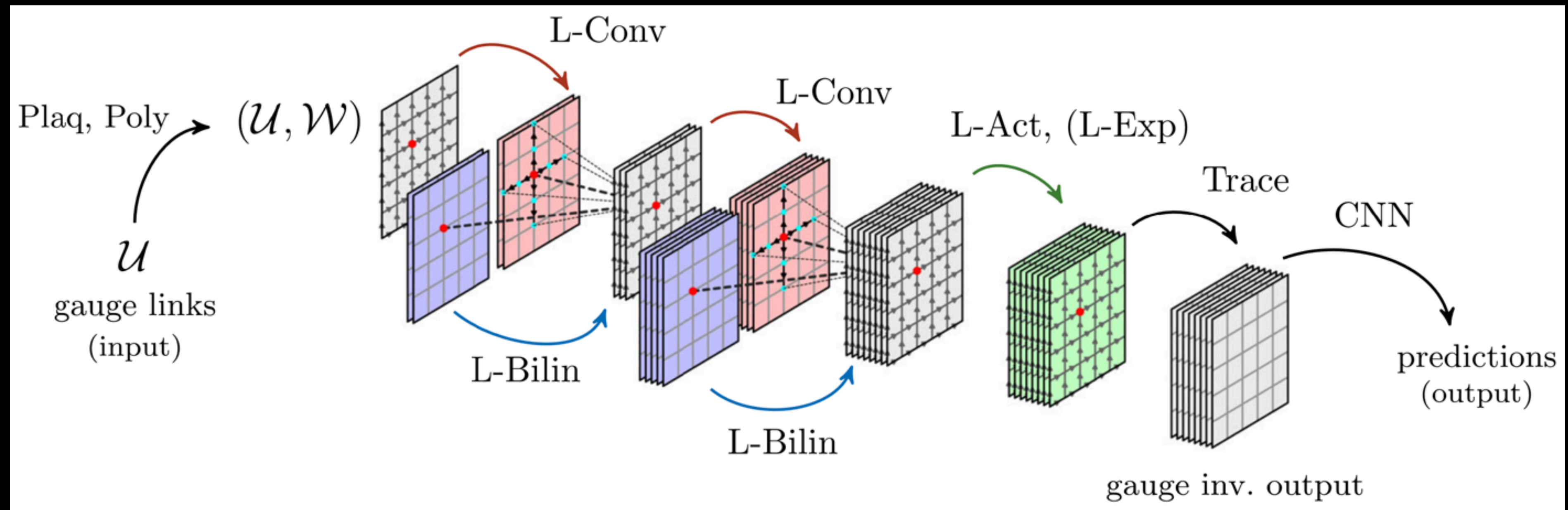
Boyda, D., Kanwar, G., Racanière, S., Rezende, D. J., Albergo, M. S., Cranmer, K., ... & Shanahan, P. E. (2021). Sampling using SU (N) gauge equivariant flows. *Physical Review D*, 103(7), 074504.

Kanwar, G., Albergo, M. S., Boyda, D., Cranmer, K., Hackett, D. C., Racaniere, S., ... & Shanahan, P. E. (2020). Equivariant flow-based sampling for lattice gauge theory. *Physical Review Letters*, 125(12), 121601.

Choices for network

All gauge equivariant

- Lattice Gauge Equivariant Convolutional Neural Networks (LCNN)



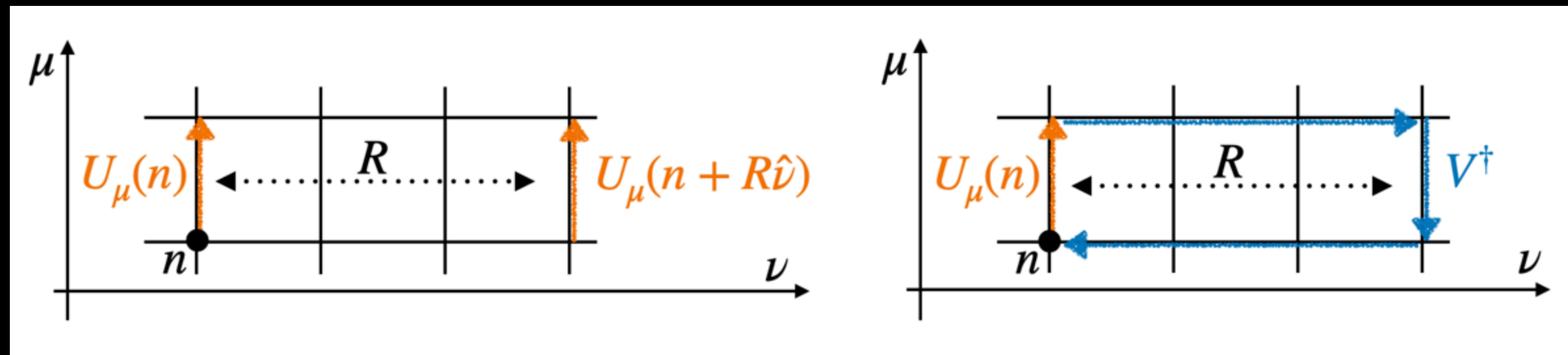
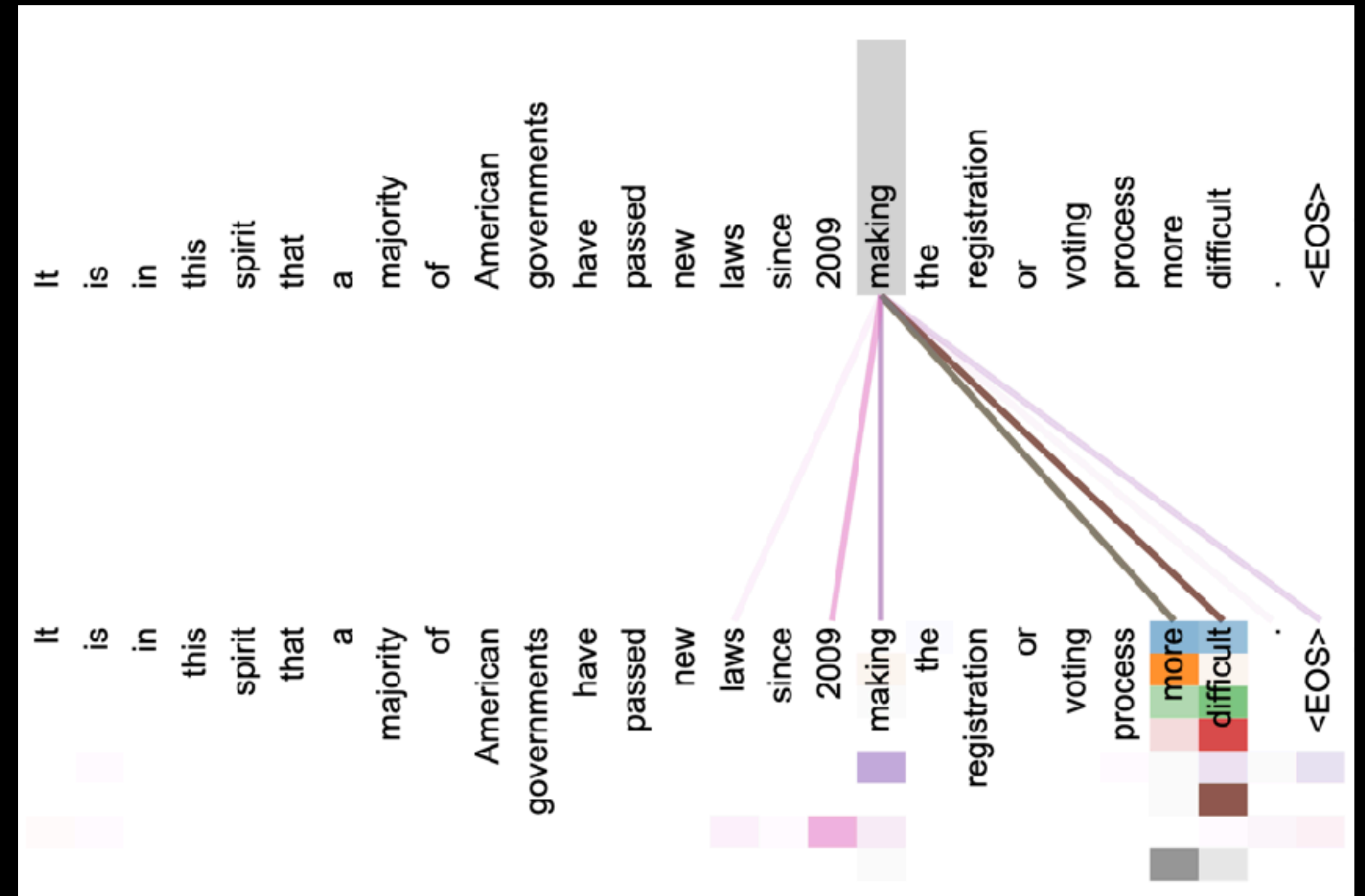
Favoni, M., Ipp, A., Müller, D. I., & Schuh, D. (2022). Lattice gauge equivariant convolutional neural networks. *Physical Review Letters*, 128(3), 032003.

Choices for network

All gauge equivariant

- Gauge equivariant transformer

Vaswani, A., ... & Polosukhin, I. (2017). Attention is all you need

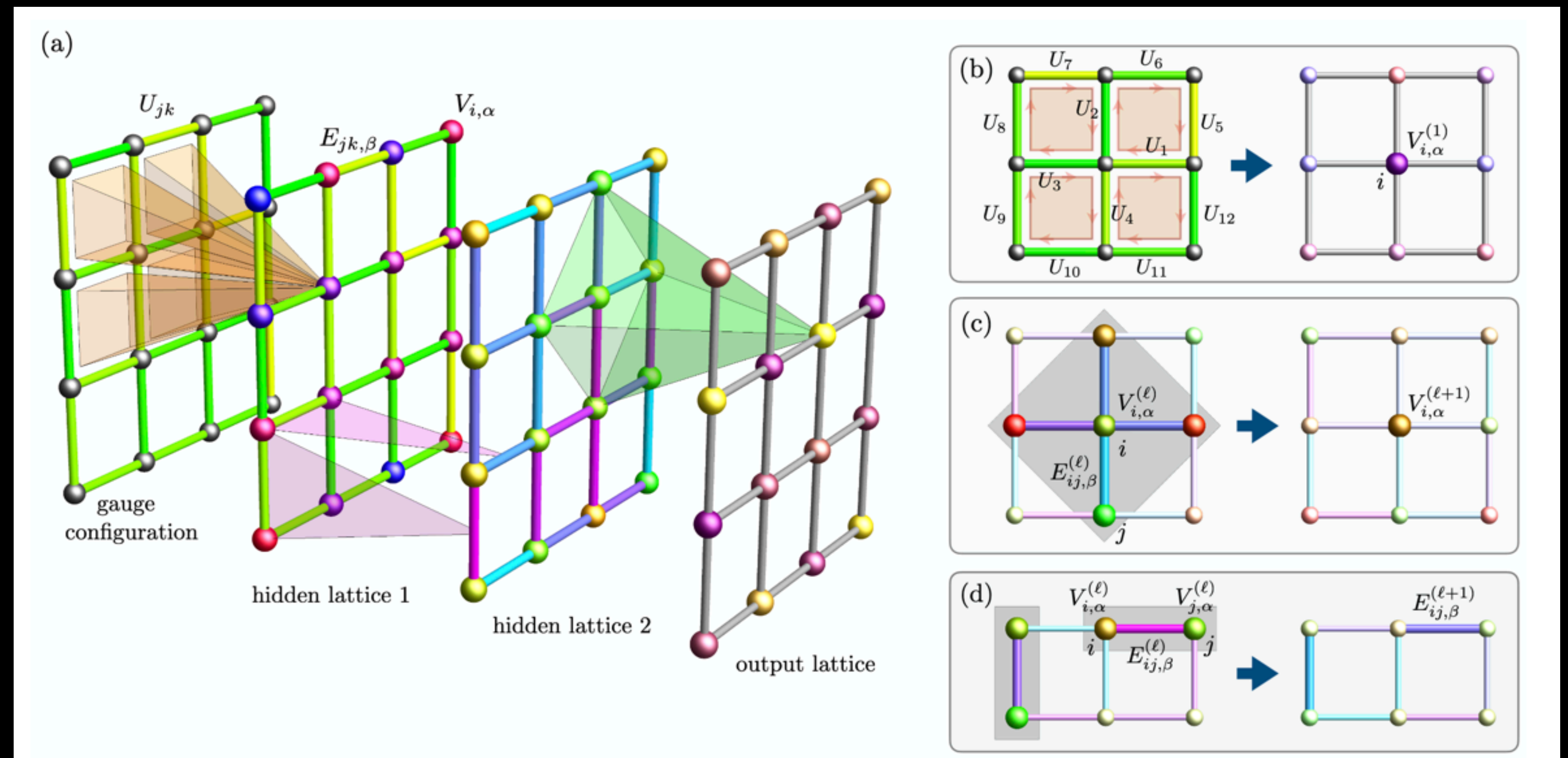


Nagai, Y., Ohno, H., & Tomiya, A. (2025). CASK: A Gauge Covariant Transformer for Lattice Gauge Theory. *arXiv preprint arXiv:2501.16955*.

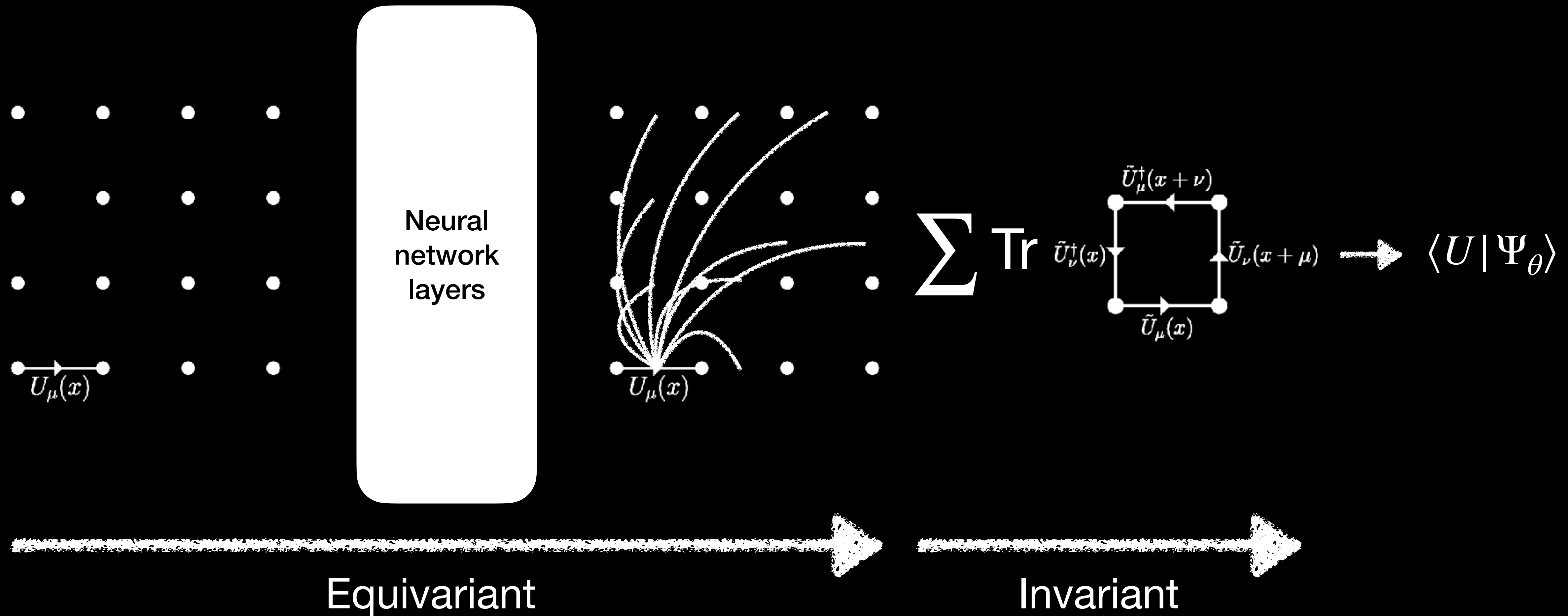
Choices for network

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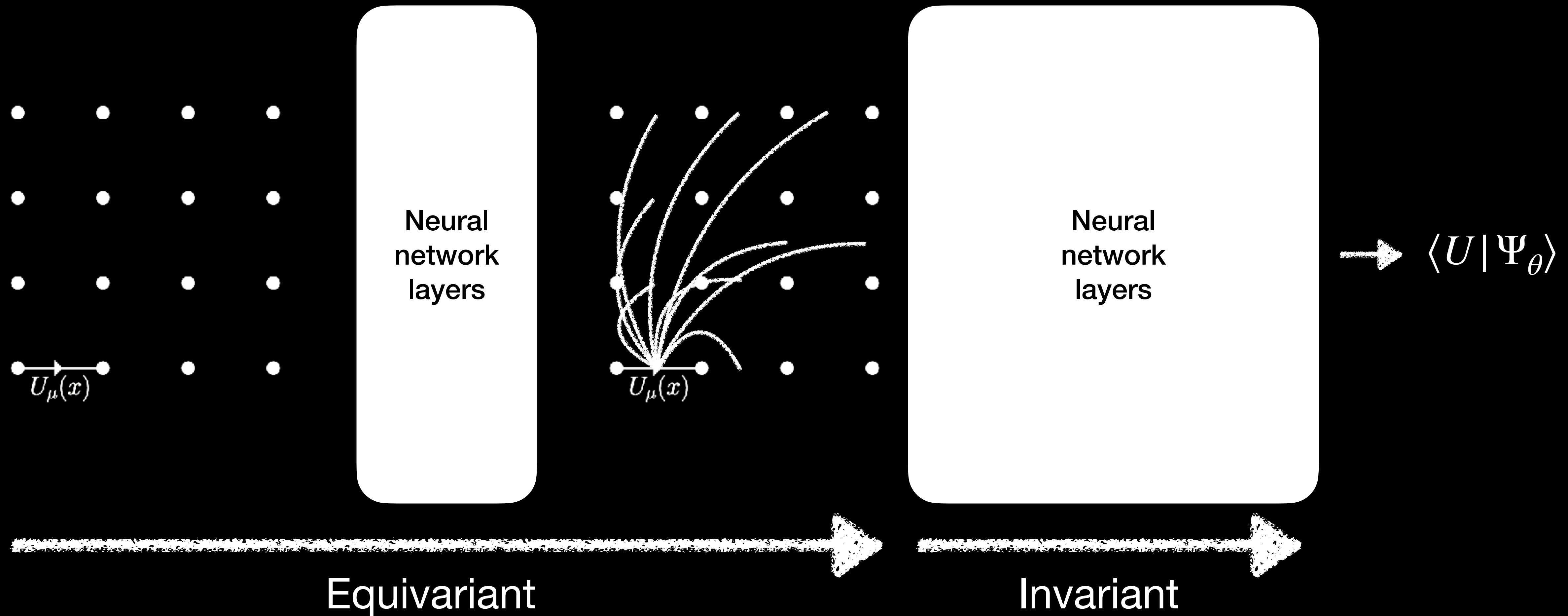
- Gauge equivariant graph neural network



Neural network + physics-informed



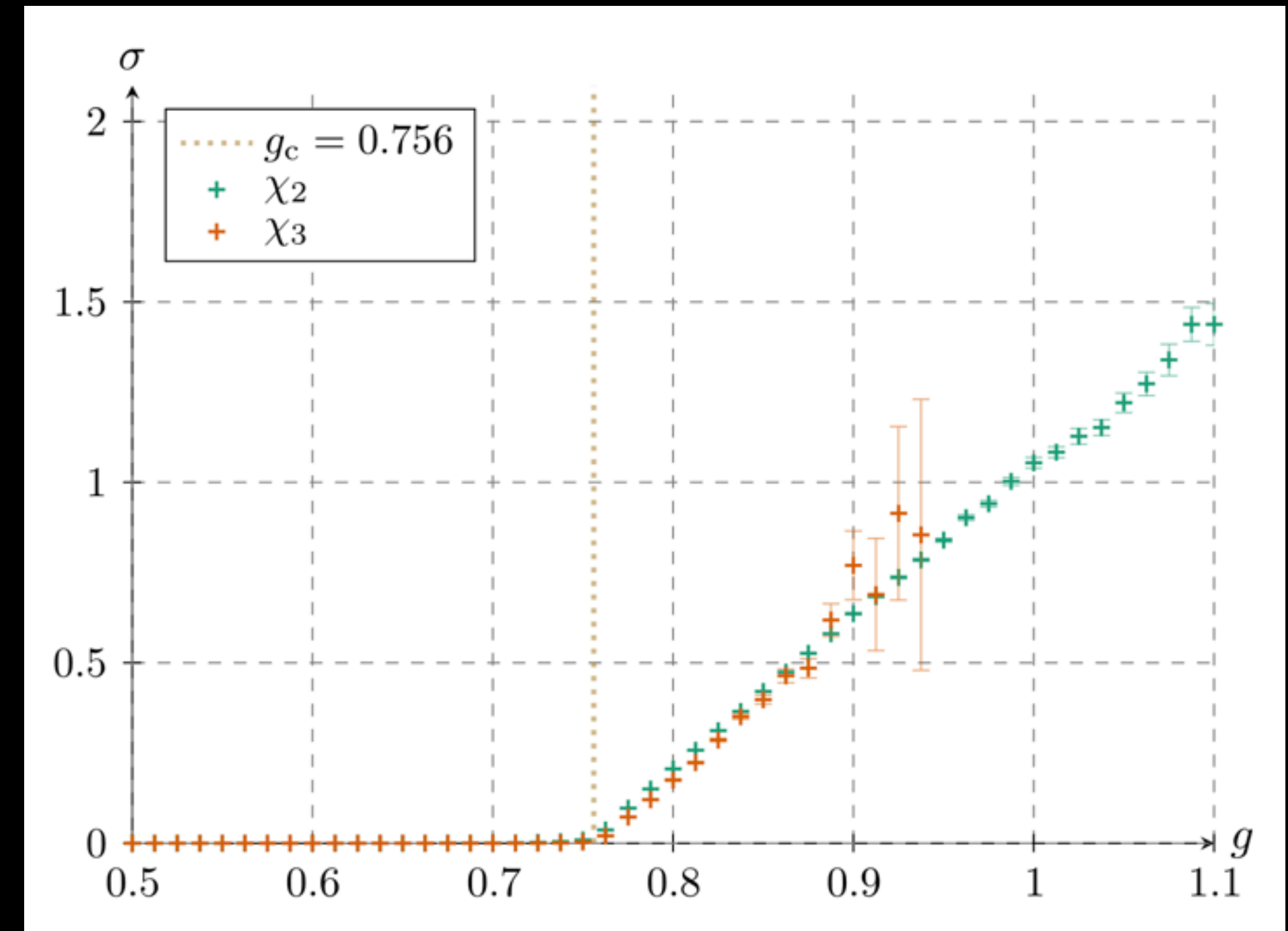
Neural network only



Existing works

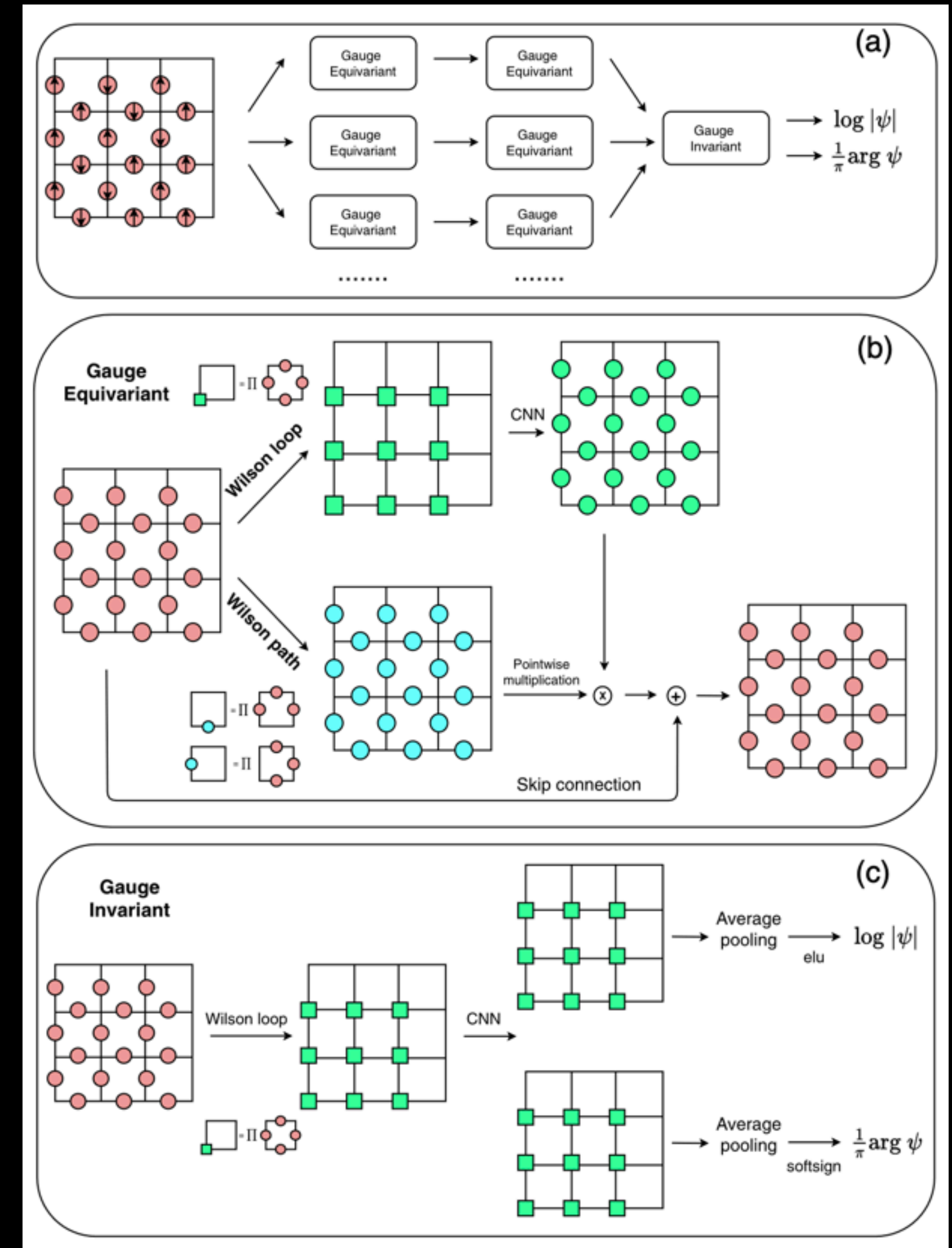
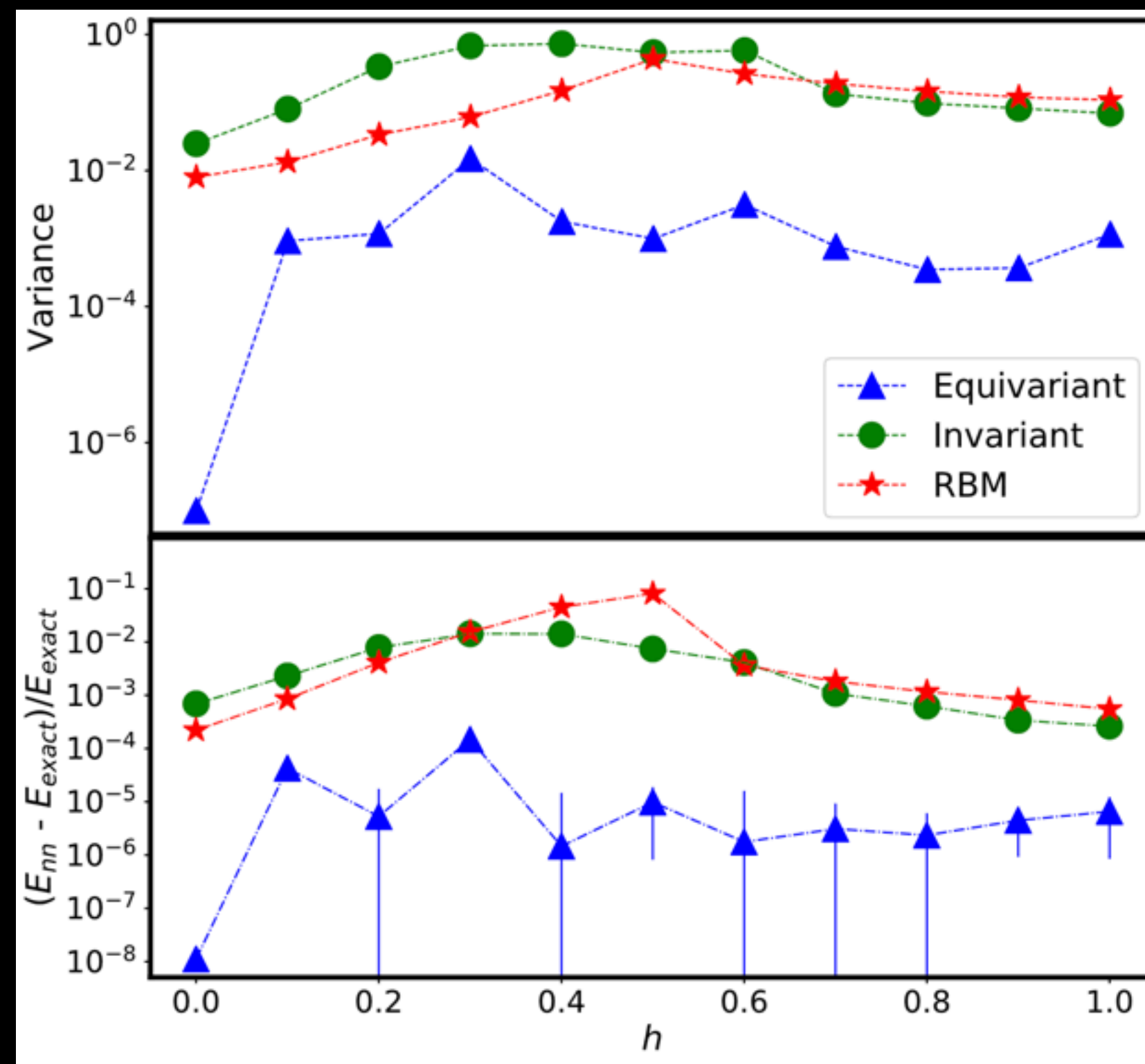
Z_2 pure gauge theory

- Using the LCNN architecture
- Detect phase transition
 - Perimeter $\rightarrow$ Area law scaling of Wilson loop



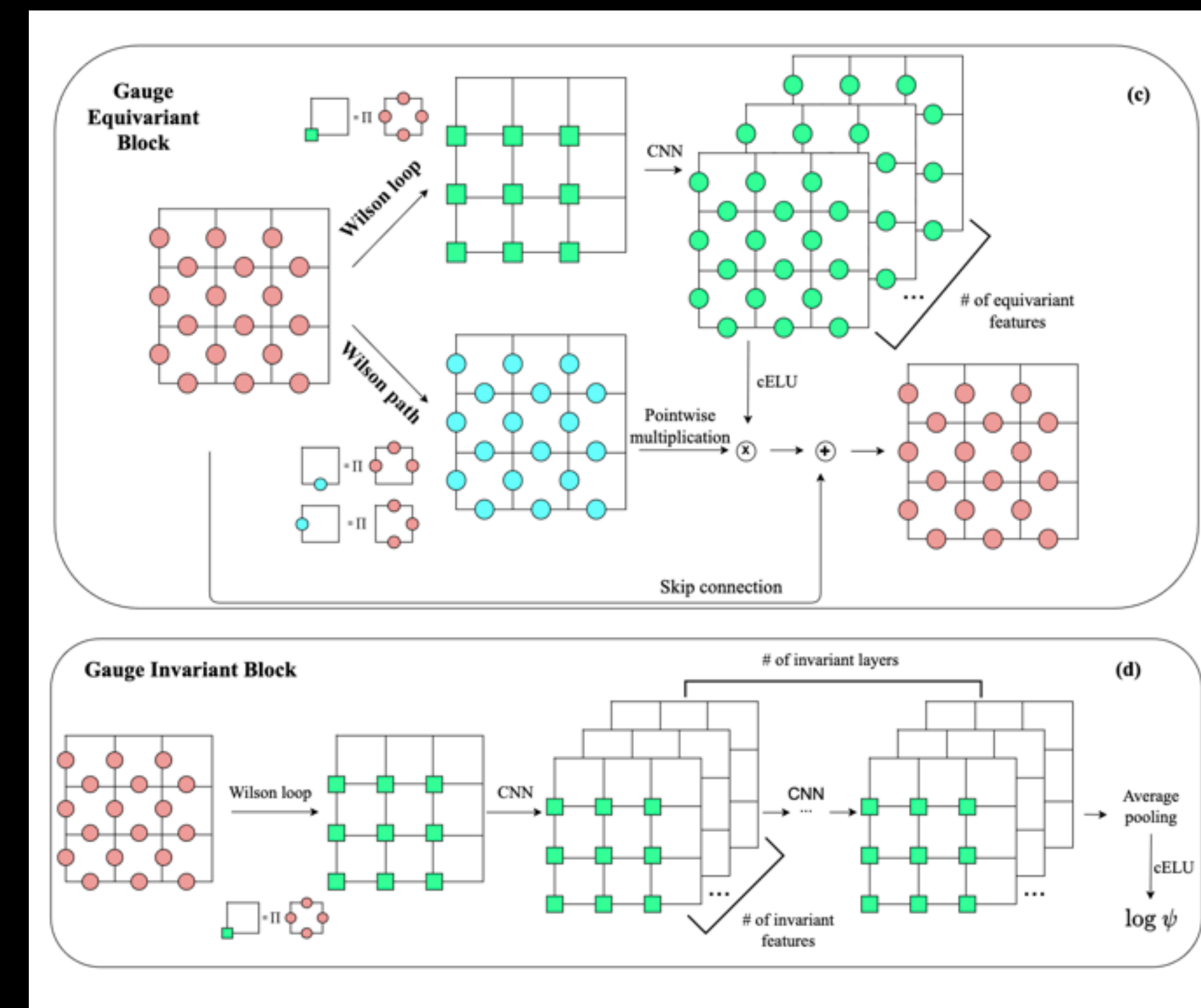
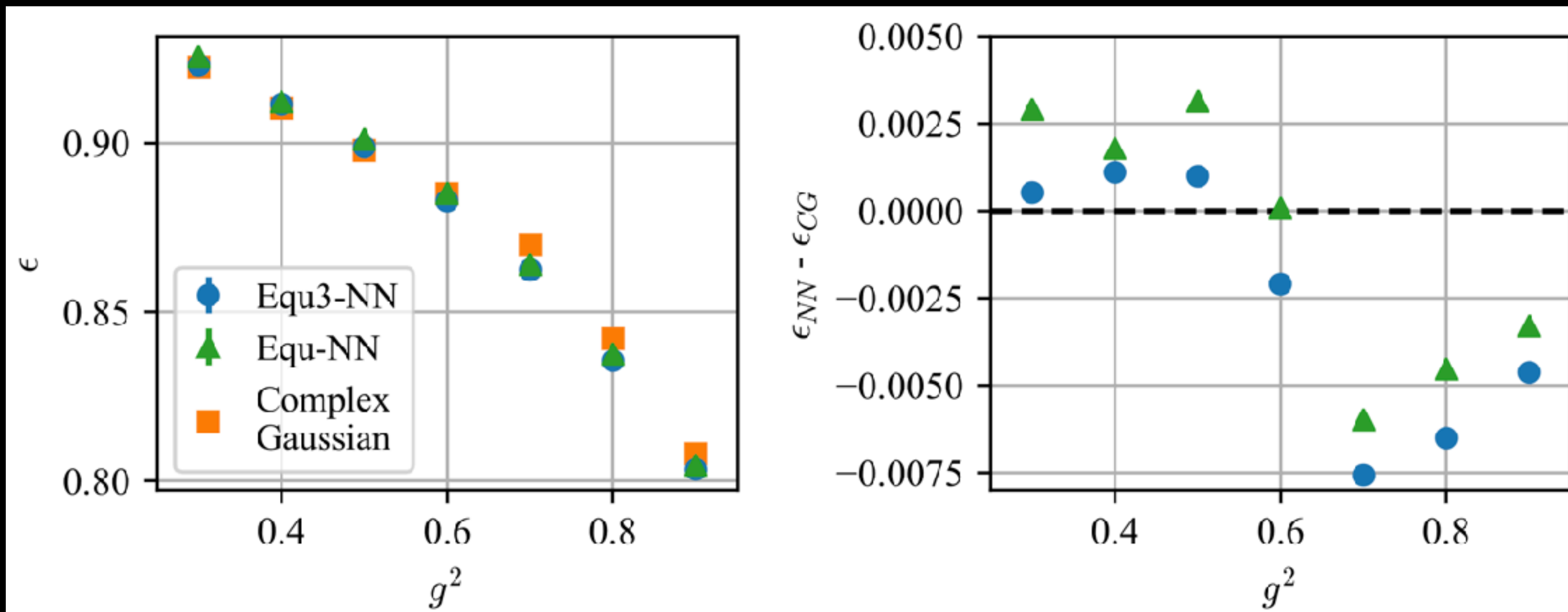
$\mathbb{Z}_2$ pure gauge theory

- Using a custom ansatz
- Energy comparable to exact solution



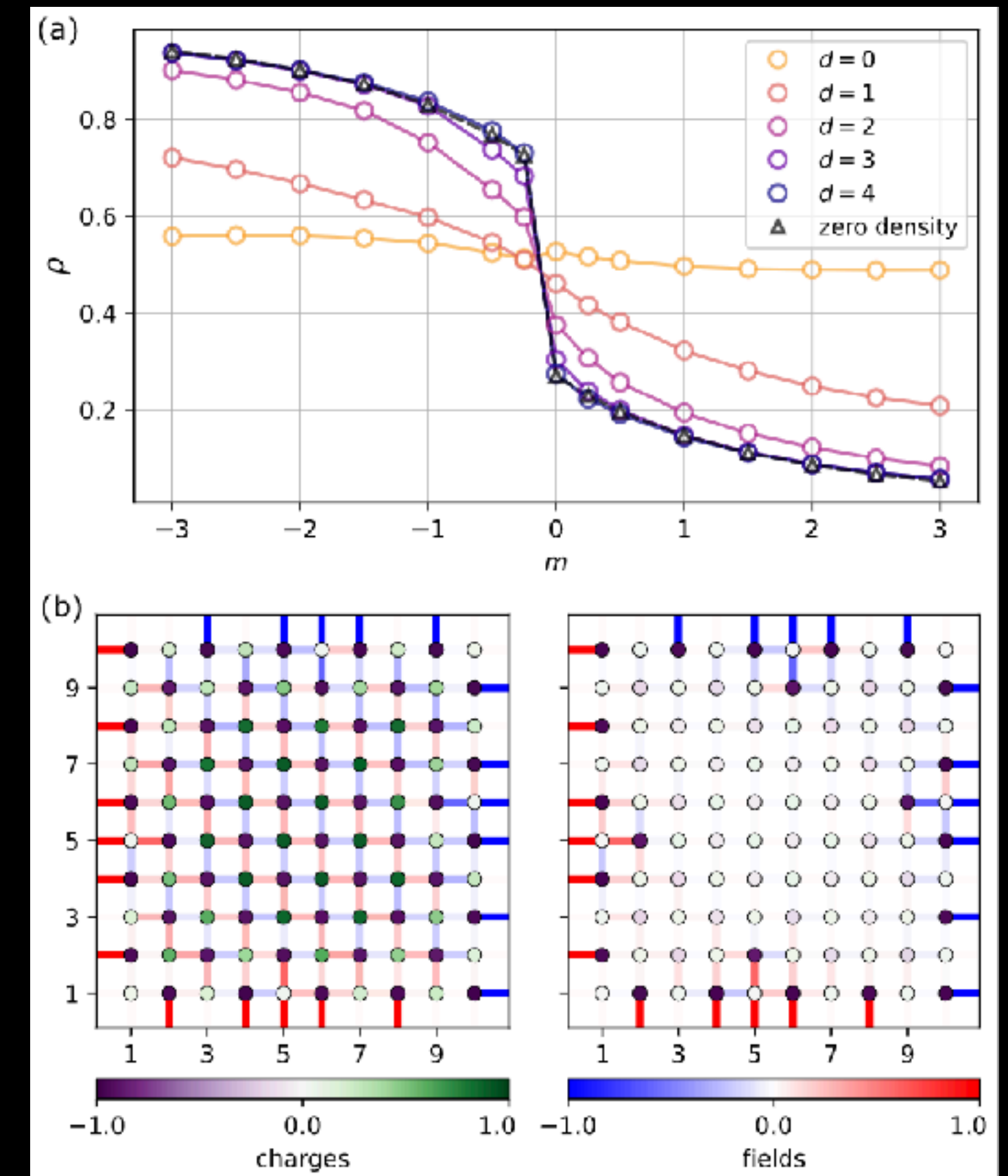
$U(1)$ pure gauge theory

Energy similar to physics informed ansatz



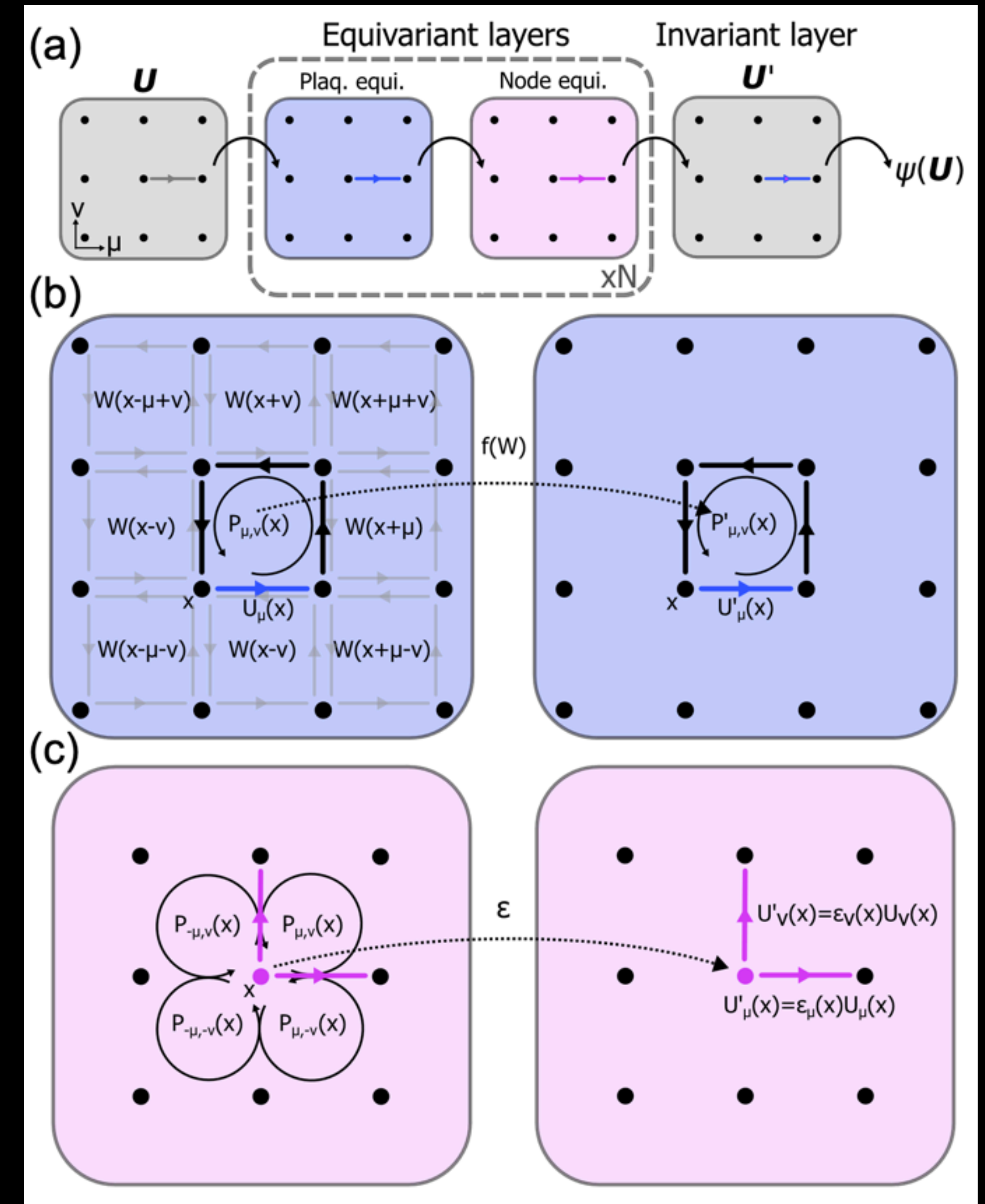
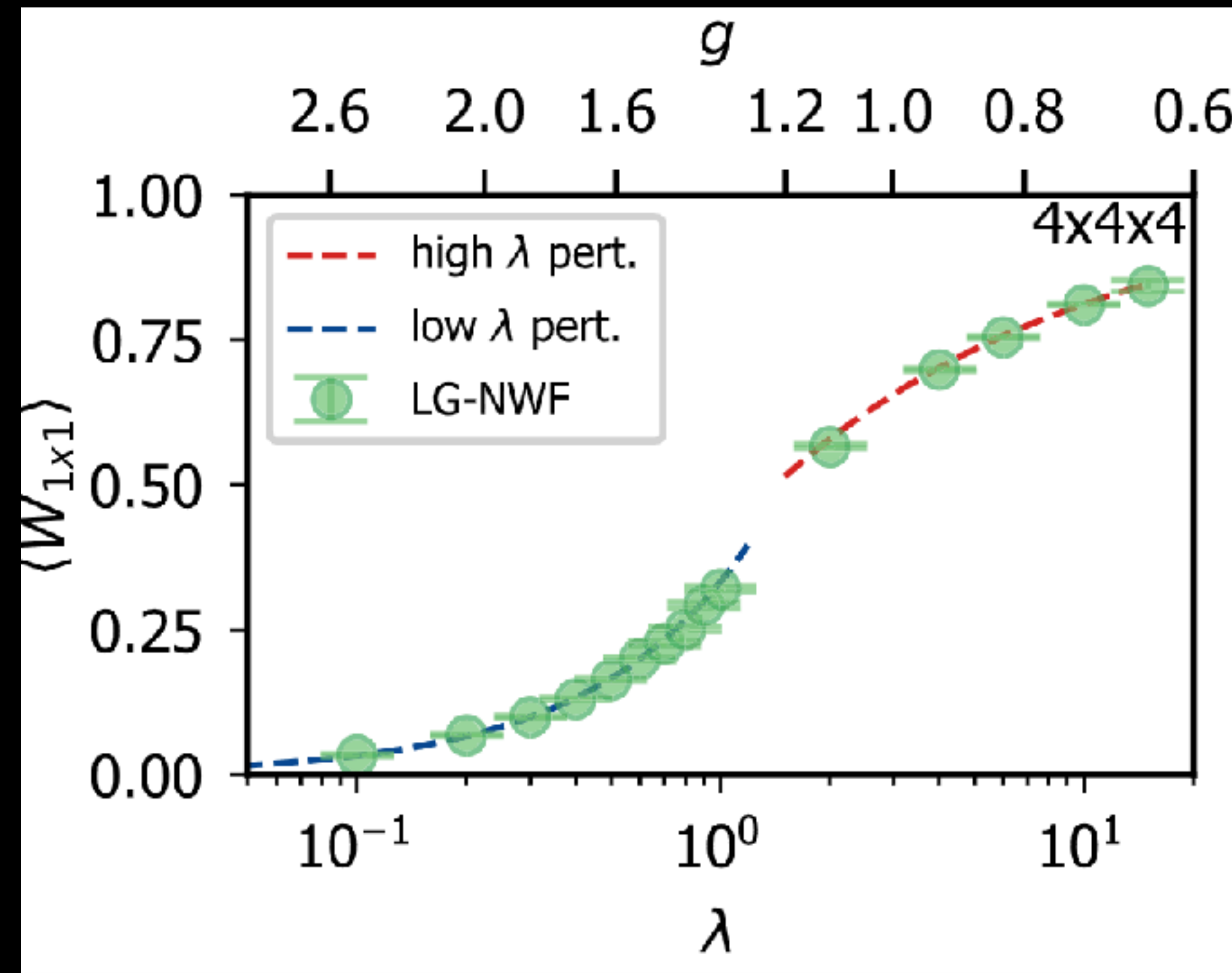
$U(1)$ gauge theory with matter

- **Autoregressive Ansatz** (no Markov chains)
- Detect/predict transitions in matter density



$SU(2)$ pure gauge theory

- Both LCNN and flow-like architectures
- 3D



Getting involved

NetKet:

<https://www.netket.org/>

- Great tutorials
- Very modular code
 - Mixed learning curve
 - Lots of code/feature reuse
- Regular updates & new features
- https://gitlab.com/QMAI/papers/su2_neural_wavefunctions

At Lattice 2026

Wednesday morning

Vacuum structure and confinement

Ground States of adjoint $SU(2)$ from Neural Newtork assisted Variational Montecarlo

Jul 29, 2026, 11:10 AM

20m

Thurgood Marshall (Adele H. Stamp Student Union)

Vacuum structure a...

Vacuum structure and ...

Speaker

David Ward (National Taiwan University)

Thursday afternoon

Algorithms and machine learning

Diagonalizing the Kogut-Susskind Hamiltonian with Physics-Informed Neural Networks

Jul 30, 2026, 4:10 PM

20m

Benjamin Banneker B (Adele H. Stamp Student Union)

Algorithms and artifi...

Algorithms and artifi...

Speaker

Simone Romiti

Neural Wavefunctions in Quantum Field Theory

Jul 30, 2026, 4:30 PM

20m

Benjamin Banneker B (Adele H. Stamp Student Union)

Contributed talk

Algorithms and artifi...

Algorithms and artifi...

Speaker

Suryansh Rajawat (University of Maryland)

3+1d $SU(3)$ Yang-Mills + theta-term with Neural Network Quantum States

Jul 30, 2026, 4:50 PM

20m

Benjamin Banneker B (Adele H. Stamp Student Union)

Contributed talk

Algorithms and artifi...

Algorithms and artifi...

Speaker

Joshua Lin (Argonne National Laboratory)



Eliska Greplova



Juan Carrasquilla



Jannes Nys

Thank you for listening

Backup slides

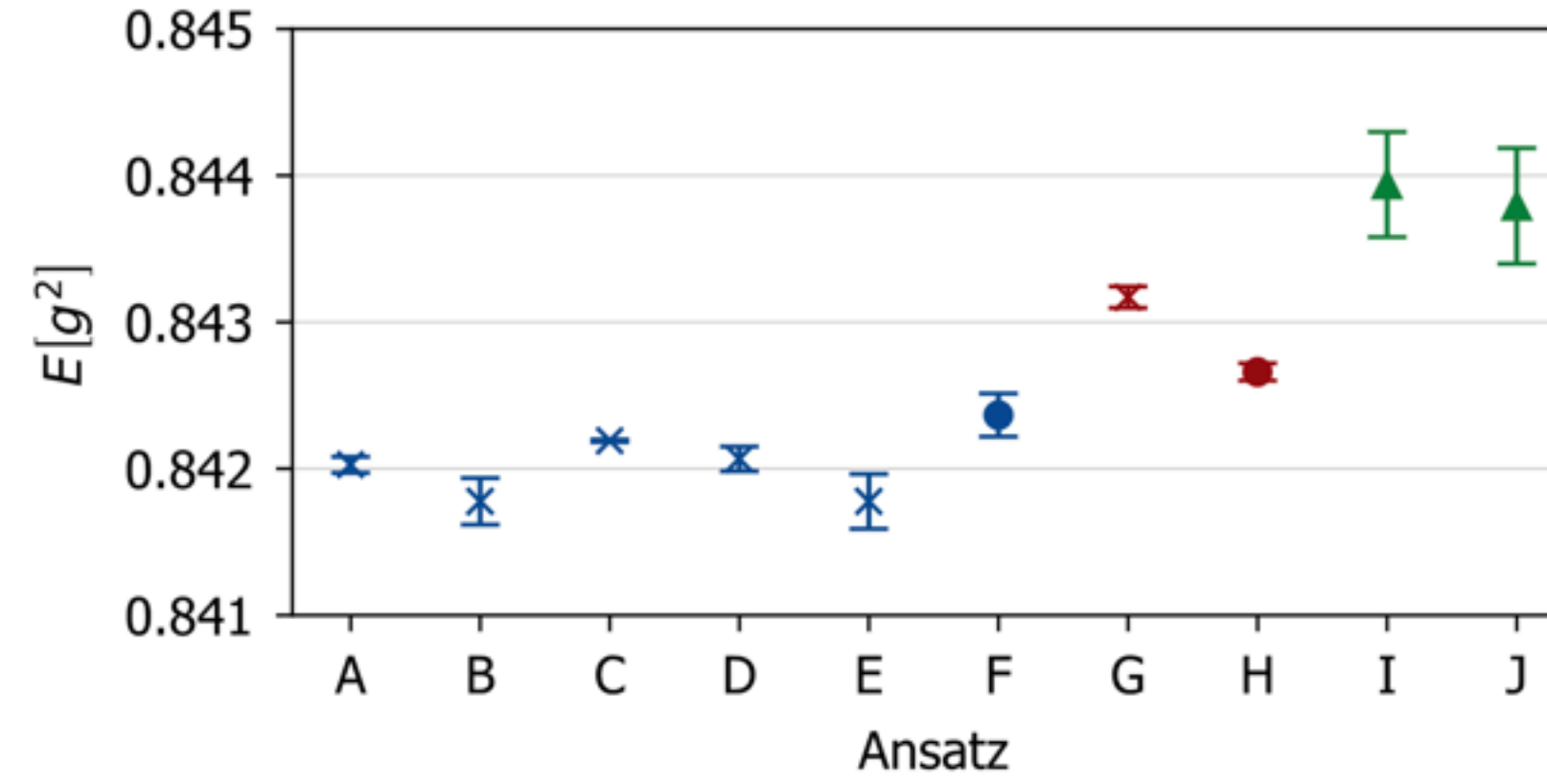


Figure S2. Variational ground state energy for a range of ansatze outlined in Table S1 for $\lambda = 1$. Errors are the Monte Carlo error of the energy over the final set of samples.

Figure S2 label	Plaquette layer			Node layer	Num. equivariant blocks	Final layer
	Architecture	Num. layers	Other			
A	CNN	1	K=3	Yes	1	Jastrow
B	CNN	1	K=1	Yes	1	Jastrow
C	CNN	2	K=3	Yes	1	Jastrow
D	CNN	1	K=3	Yes	2	Jastrow
E	CNN	1	K=3	Yes	1	FCN
F	ViT	1	NH=2, D=4	Yes	1	Jastrow
G	CNN	1	K=3	No	1	Jastrow
H	ViT	1	NH=2, D=4	No	1	Jastrow
I	None	N/A	N/A	No	0	Jastrow
J	None	N/A	N/A	No	0	FCN

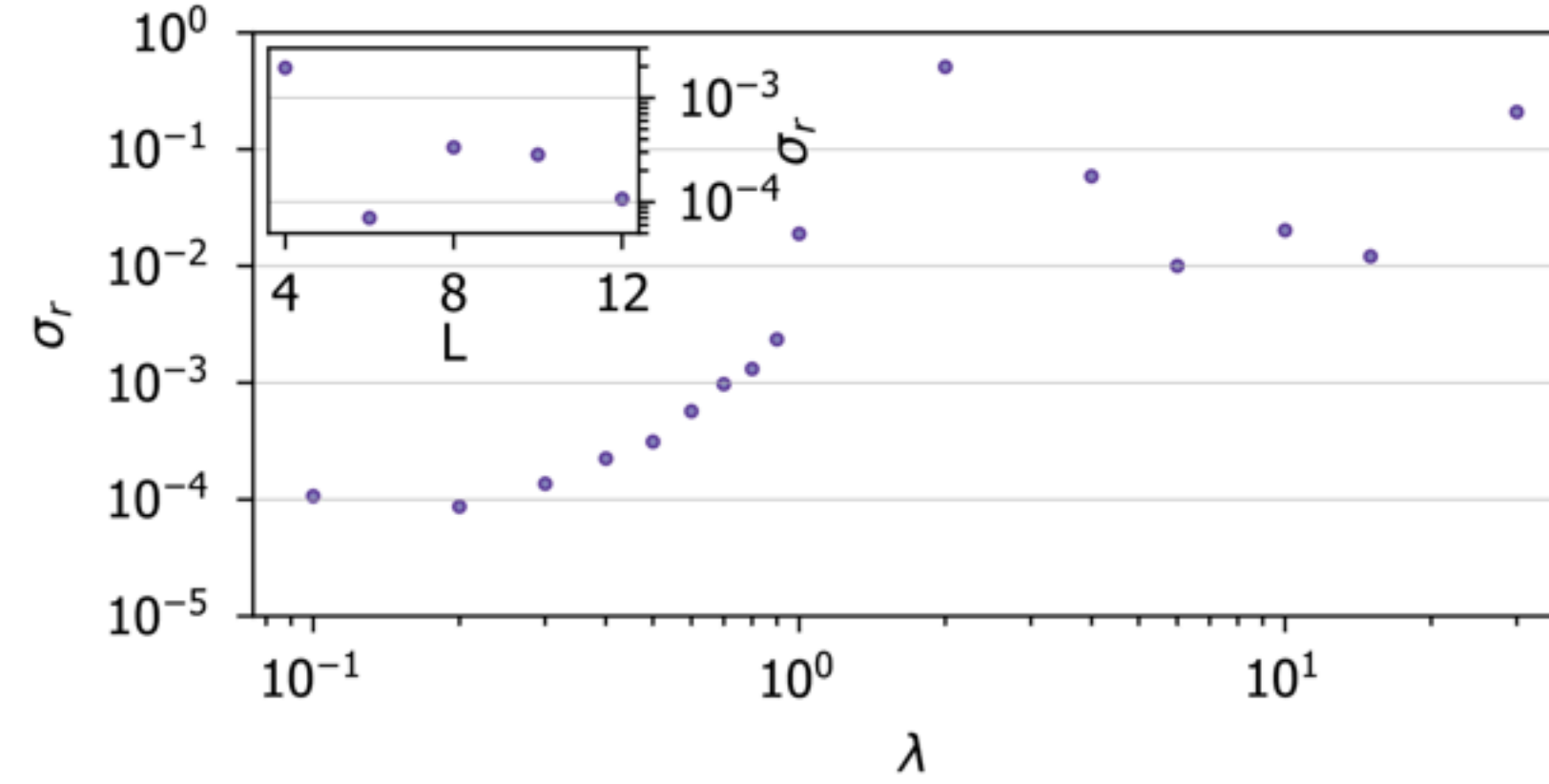


Figure S4. The rescaled variance of the energy, as in Eq. (D22), averaged over the final 20 iterations of VMC. The main panel shows results for 12×12 lattices at various values of λ ; the inset shows the rescaled variance but against system size ($L \times L$) with fixed $\lambda = 0.1$.

the version trained under the alternative Hamiltonian and then fine-tuned and they agree within errors.

During the VMC procedure we also compute the energy variance σ , defined as

$$\sigma = \frac{\langle \Psi | H^2 | \Psi \rangle}{\langle \Psi | \Psi \rangle} - \left(\frac{\langle \Psi | H | \Psi \rangle}{\langle \Psi | \Psi \rangle} \right)^2. \quad (\text{D21})$$

Different gauge groups

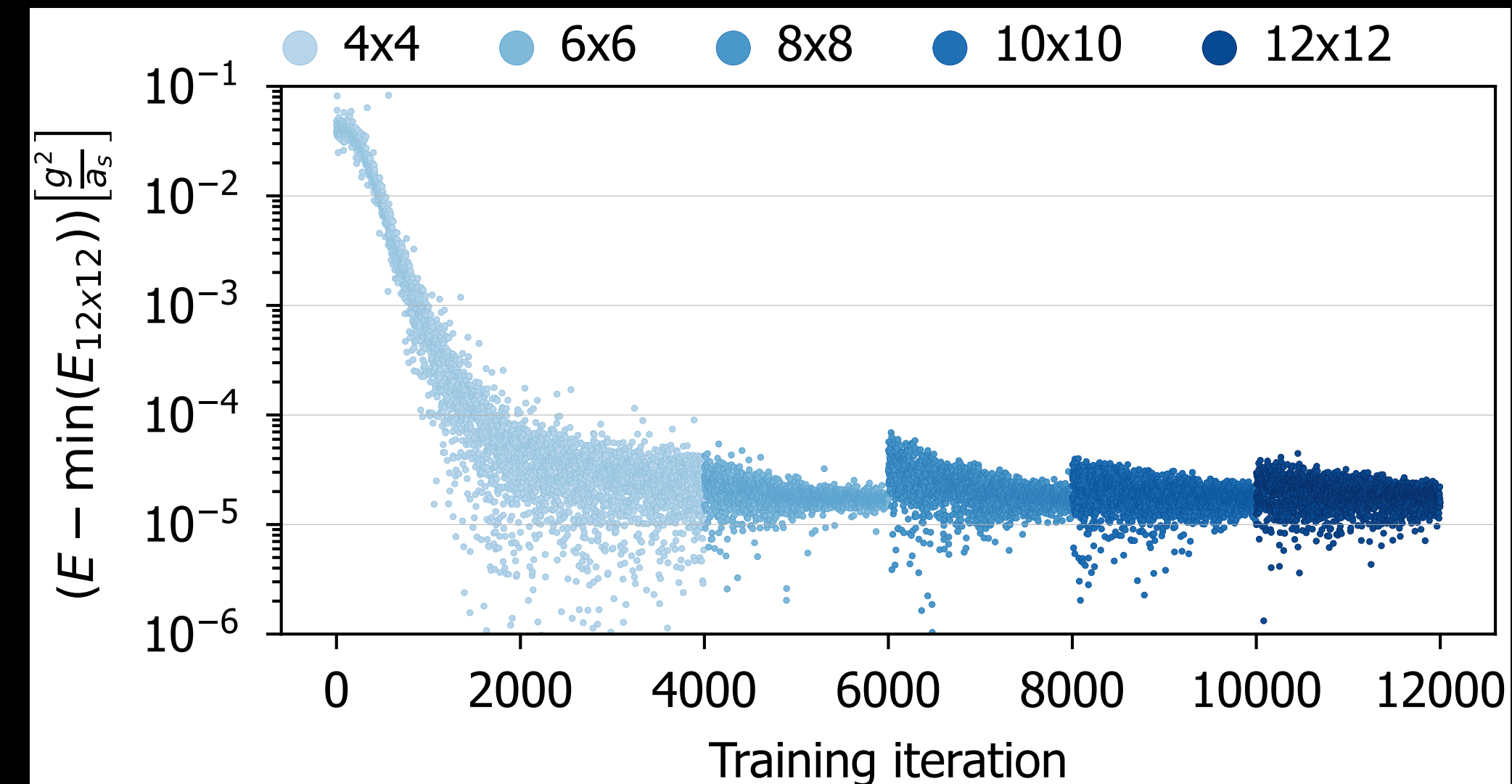
Do-able

- $U = U(\theta, \phi, \psi) \rightarrow 8$ angles
- $\nabla^2 \rightarrow$ More complicated
 - (possibly done but not explained in Chin, S. A., *et al.* Phys. Rev. D **37**, 3001)
- Note: canonical ordering of eigenvalues non-trivial in SU(3)
 - When do go $P_{\mu,\nu}(x) \rightarrow \lambda_P \rightarrow \lambda'_P \rightarrow P'_{\mu,\nu}(x)$ you have to order the eigenvalues before and after so they correspond to the same eigenvectors

Larger lattices

Trivial

- We did transfer learning in this work already
- Equivariant layers are local
 - Convolution layers, fixed kernel size
- Final layer is non-local
 - One new parameter for each new unique distance on the lattice
 - Can be initialised as ‘small’
- Can also be done in arbitrary dimension, 2 and 3 shown in paper



Fermions and time evolution

Hard, open study

- Both done in spin systems with neural wavefunctions
 - Even at the same time:
 - Nys, J., Pescia, G., Sinibaldi, A. and Carleo, G., 2024. Ab-initio variational wave functions for the time-dependent many-electron Schrödinger equation. *Nature communications*, 15(1), p.9404.
- Gauge invariance and representation for LGTs with fermions remain open questions
- Time evolution is more straightforward (only variational challenges, nothing conceptual)

Other applications

Done, but not for LGTs

- Physics distilled into neural network parameters, θ
 - Fewer parameters than basis states
- Infer phase transitions
 - Hernandez, V., **Spriggs**, T., Khaleefah, S., Greplova, E. ICLR Workshop on Neural Network Weights as a New Data Modality (2025). *quant-ph arXiv:2503.17140*.
- Excited states and phase transitions
 - Medvidović, M., Orfi, A., Carrasquilla, J., Sels, D. arXiv:2510.15030
- Compute non-stabiliserness
 - **Spriggs**, T., Ahmadi, A., Chen, B. and Greplova, E., 2025. *Machine Learning: Science and Technology*, 6(1), p.015042

Time evolution

- Time evolution of the quantum state $\rightarrow$ time evolution of parameters
- $|\Psi(t)\rangle = |\Phi_{\theta(t)}\rangle$
- Many approaches
 - TVMC, pTVMC, TreTVMC
- TVMC equates a small Trotter evolution of the true state to a Taylor expansion of the variational state:
 - $|\Psi(t + \delta t)\rangle = |\Phi_{\theta(t+\delta t)}\rangle \rightarrow \Psi(t)(1 - iH\delta t) = \Phi_{\theta(t)}(1 + \dot{\theta}t)$
 - Minimise Fubini-Study distance between these two expansions, solve for $\dot{\theta}$