

Resolving vector-current spectral functions with singular value decomposition

Sakura Itatani (SOKENDAI)

collaboration with:

Shoji Hashimoto (KEK, SOKENDAI)

Ryutaro Tsuji (KEK)



Jul. 29 @LATTICE2026

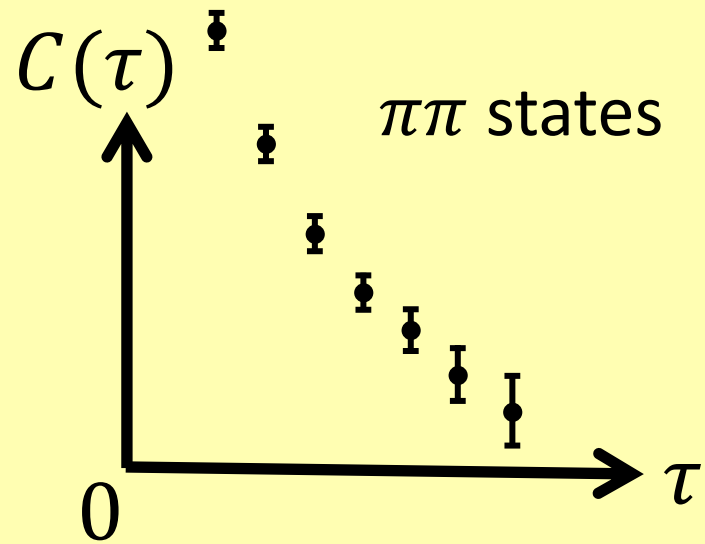
Analysis of vector-current correlator

1

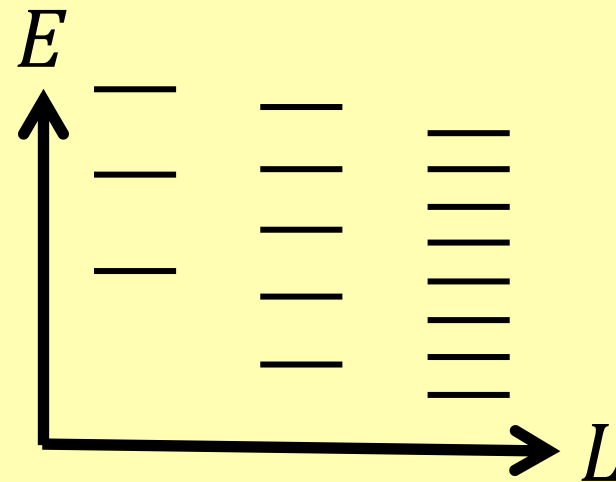
Standard method:

ρ = study of ρ -resonance i.e. $I=1 \pi\pi$ states

Lattice data



Energy spectrum



Lüscher's method

M. Lüscher [Nucl. Phys. B. 354 (1991)]

$$\det[F^{-1}(E, L) + \mathcal{M}(E)] = 0$$

- Phase shift
- Scattering length
- Decay width

In larger volumes, the energy level becomes denser.
It will be difficult to distinguish those states.

Alternative methods to study hadron spectroscopy?

Smearred spectral functions

Correlator: $C(\tau) = \int dE \rho(E) e^{-E\tau}$

1st-order moment: $\langle E \rangle(\tau) = \int dE \rho(E) E e^{-E\tau} / C(\tau)$

2nd-order moment: $\Delta E(\tau) = \sqrt{\langle (E - \langle E \rangle)^2 \rangle}$

Written as a smeared spectrum with different smearing **kernel**.

- For lattice determination of $\langle E \rangle$ and ΔE , we use SVD of $e^{-E\tau}$.
- Study of sensitivity to determine m_ρ and Γ_ρ .

$$e^{-E\tau} \xrightarrow{\text{SVD}} \sum_l U_l(\tau) \sigma_l V_l(E)$$

$$\sum_E V_{l'}(E) V_l(E) = \delta_{ll'}$$

$$\sum_\tau U_{l'}(\tau) U_l(\tau) = \delta_{ll'}$$

$$C(\tau) = \int dE \rho(E) e^{-E\tau}$$

$$= \sum_l U_l(\tau) \sigma_l \underbrace{\int dE \rho(E) V_l(E)}_{\equiv \rho_l}$$

R. Tsuji, S. Hashimoto [arXiv:2605.15674]

And talks by

R. Tsuji (Jul. 28, 15:20-)

S. Hashimoto (Jul.27, plenary)

$$\rho(E) = \sum_l \rho_l V_l(E),$$

$$\rho_l = \frac{1}{\sigma_l} \sum_\tau C^{\text{lat}}(\tau) U_l(\tau)$$

Construction of smeared spectrum with SVD 4

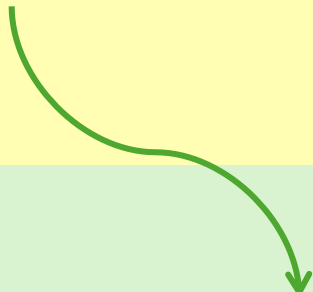
$$\langle E^n \rangle(\tau) = \frac{1}{C(\tau)} \int dE \ E^n \ \rho(E) \ e^{-E\tau}$$

$$\rho(E) = \sum_l \ \rho_l \ V_l(E)$$

$$= \frac{1}{C(\tau)} \sum_l \rho_l \int dE \ E^n \ V_l(E) \ e^{-E\tau}$$

SVD of $e^{-E\tau}$
+
Lattice data $C^{\text{lat}}(\tau)$

$\Rightarrow \rho_l = \frac{1}{\sigma_l} \sum_{\tau} C^{\text{lat}}(\tau) U_l(\tau) \Rightarrow \langle E^n \rangle(\tau)$



Breit-Wigner type

$$\delta^{\text{BW}}(s) = \frac{m_\rho}{m_\rho^2 - s} \Gamma_\rho(s)$$

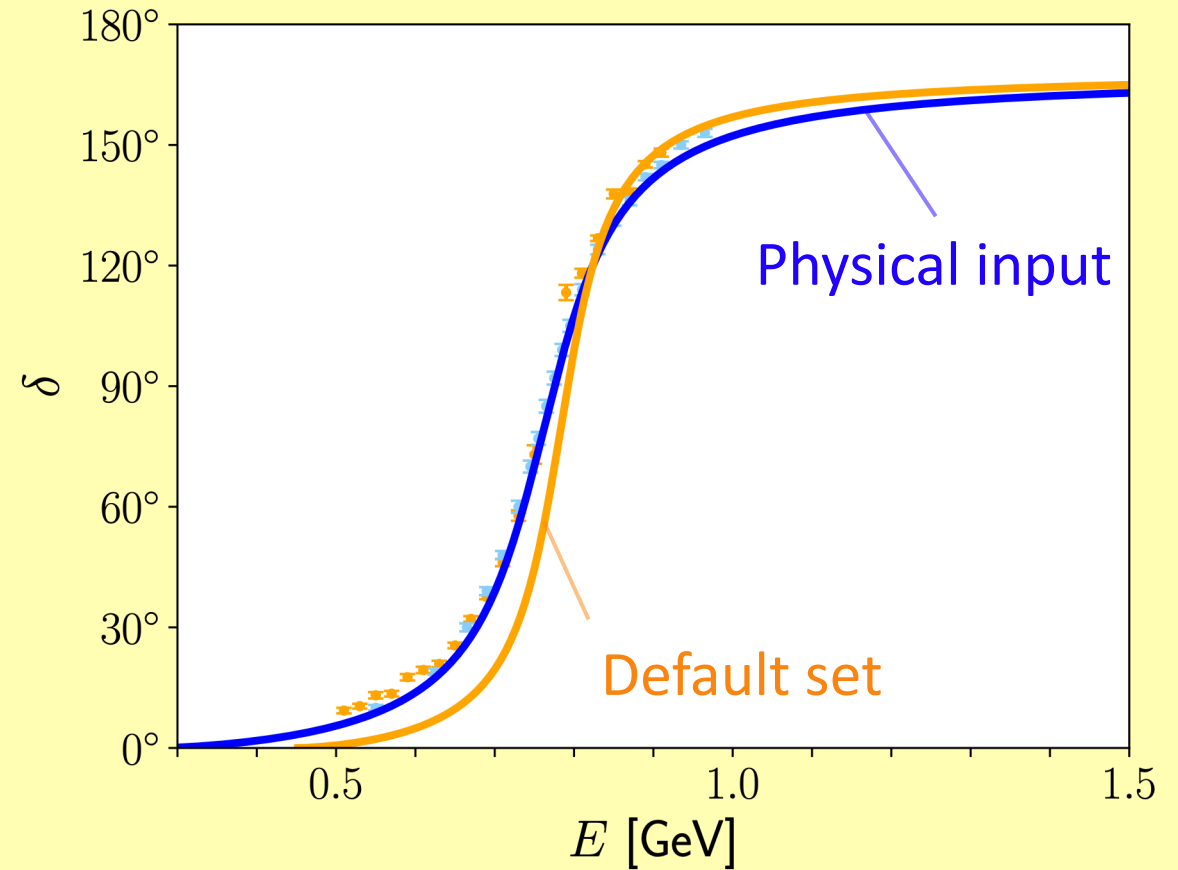
$$\Gamma_\rho(s) = \Gamma_\rho \left(\frac{s - 4m_\pi^2}{m_\rho^2 - 4m_\pi^2} \right)^{3/2} \frac{m_\rho}{\sqrt{s}}$$

To compare with our lattice:

JLQCD, domain-wall lattice,

$32^3 \times 64, 48^3 \times 96,$

$a \sim 0.08 \text{ fm}, m_\pi = 226 \text{ MeV}$



Model params.	m_π [MeV]	m_ρ [MeV]	Γ_ρ [MeV]	$\sqrt{\langle r_\pi^2 \rangle}$ [fm]
Physical	139.57	775.26	149.1	0.66
Our default set	226	792	98	0.635

Omnés representation

$\ln|F_\pi|(s)$

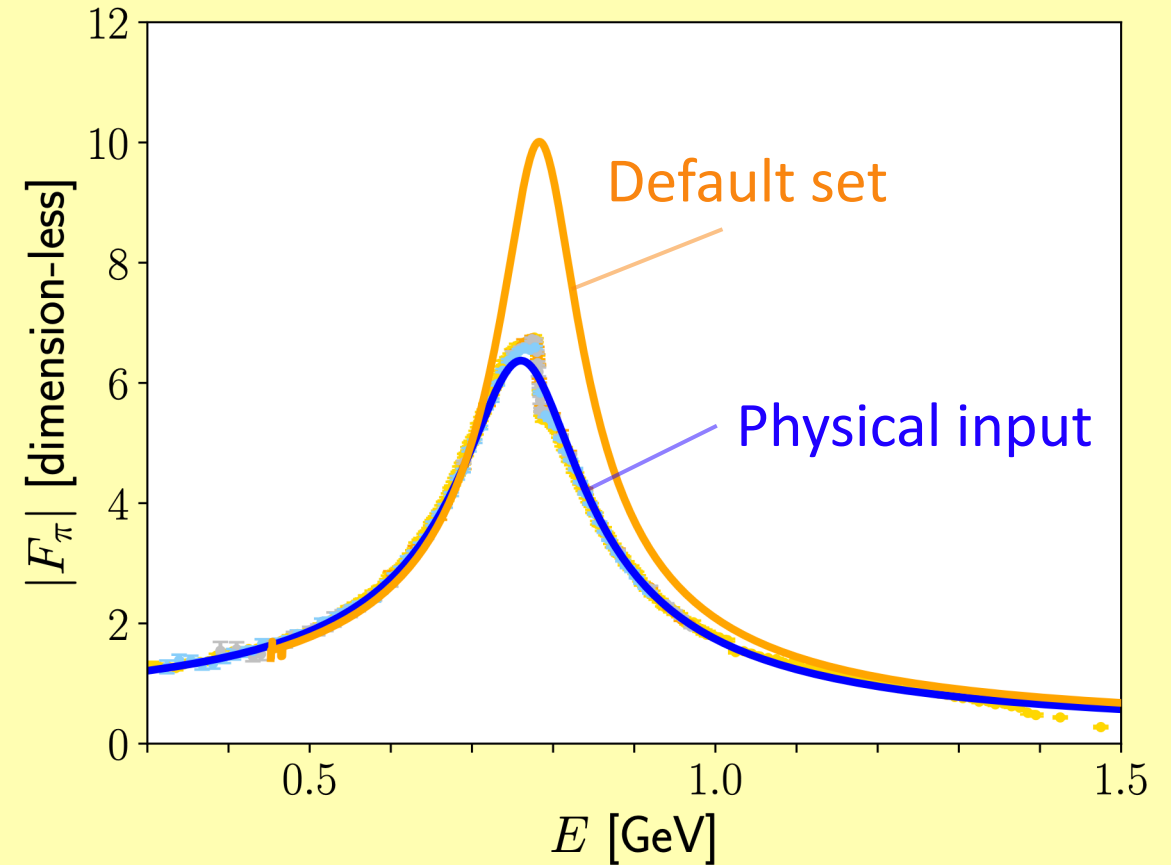
$$= \frac{1}{6} \langle r_\pi^2 \rangle s + \frac{s^2}{\pi} \int_{4m_\pi^2}^{\infty} ds' \frac{\delta^{\text{BW}}(s')}{s'(s' - s)}$$

To compare with our lattice:

JLQCD, domain-wall lattice,

$32^3 \times 64, 48^3 \times 96,$

$a \sim 0.08 \text{ fm}, m_\pi = 226 \text{ MeV}$



Model params.	m_π [MeV]	m_ρ [MeV]	Γ_ρ [MeV]	$\sqrt{\langle r_\pi^2 \rangle}$ [fm]
Physical	139.57	775.26	149.1	0.66
Our default set	226	792	98	0.635

Construction of the model correlator

7

Finite volume:

SI, H. Fukaya, S. Hashimoto [PRD **111**, 114507]

$$C(\tau, L) = \sum_n |\langle 0 | J_\mu | \pi\pi, n \rangle|^2 e^{-E_n \tau}$$

$$E_n = 2\sqrt{m_\pi^2 + k_n^2}$$

Lüscher's quantization condition

M. Lüscher [Nucl. Phys. B. **354** (1991)]

$$\det[F^{-1}(E, L) + \mathcal{M}(E)] = 0$$

Finite V

Infinite-V
amplitude

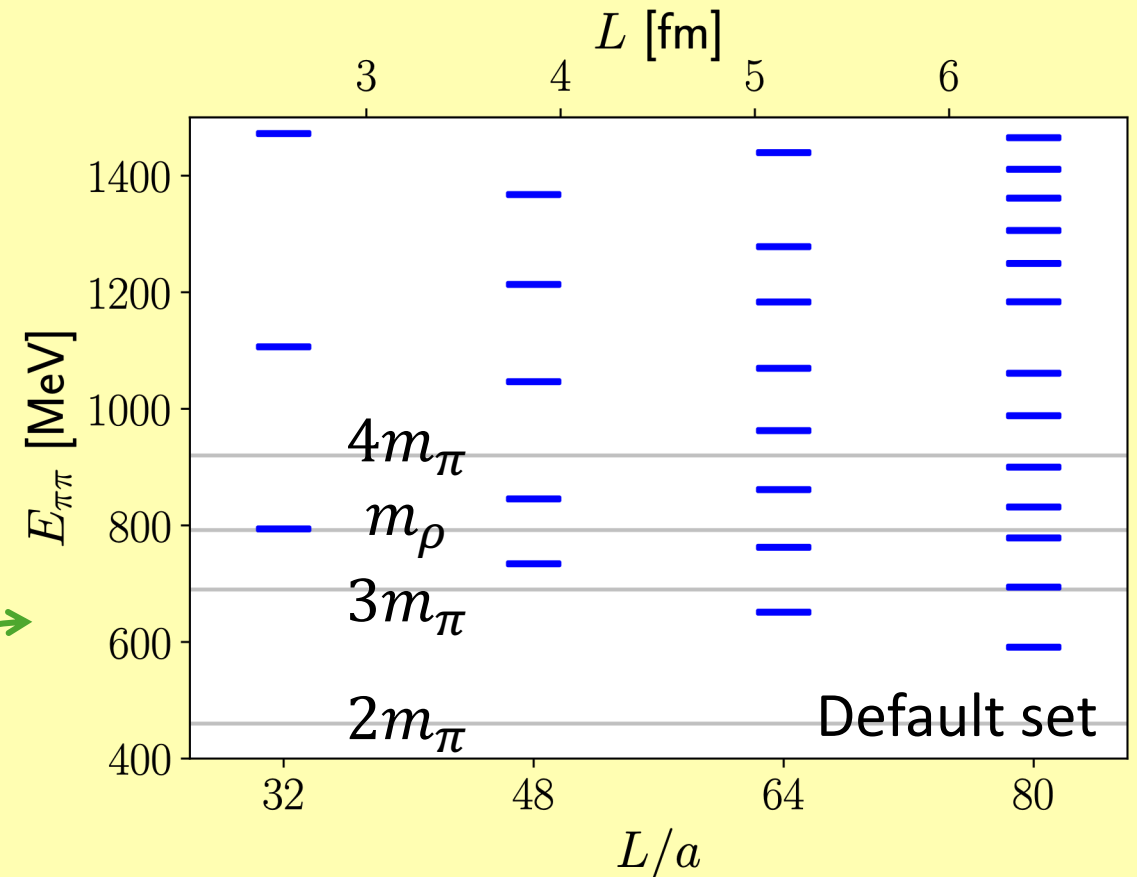
P-wave only:

$$\phi\left(\frac{kL}{2\pi}\right) + \delta^{\text{BW}}(k) = 0$$

Sym. of the system

Scattering
phase shift

solve



Finite volume:

SI, H. Fukaya, S. Hashimoto [PRD **111**, 114507]

$$C(\tau, L) = \sum_n |\langle 0 | J_\mu | \pi\pi, n \rangle|^2 e^{-E_n \tau}$$

$E_n = 2\sqrt{m_\pi^2 + k_n^2}$

Lellouch-Lüscher-Meyer formula

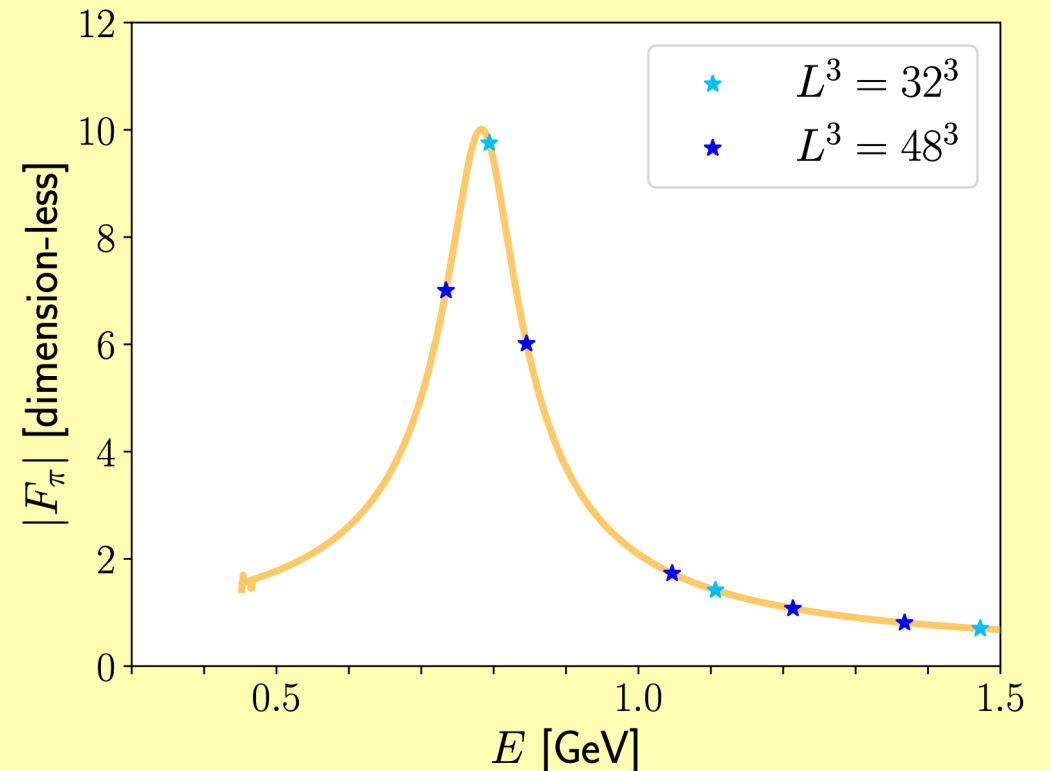
H. B. Meyer [Phys. Rev. L. **107** (2022)]

$$|F_\pi|^2 = (\text{factor}) \cdot |\langle 0 | J_\mu | \pi\pi, n \rangle|^2$$

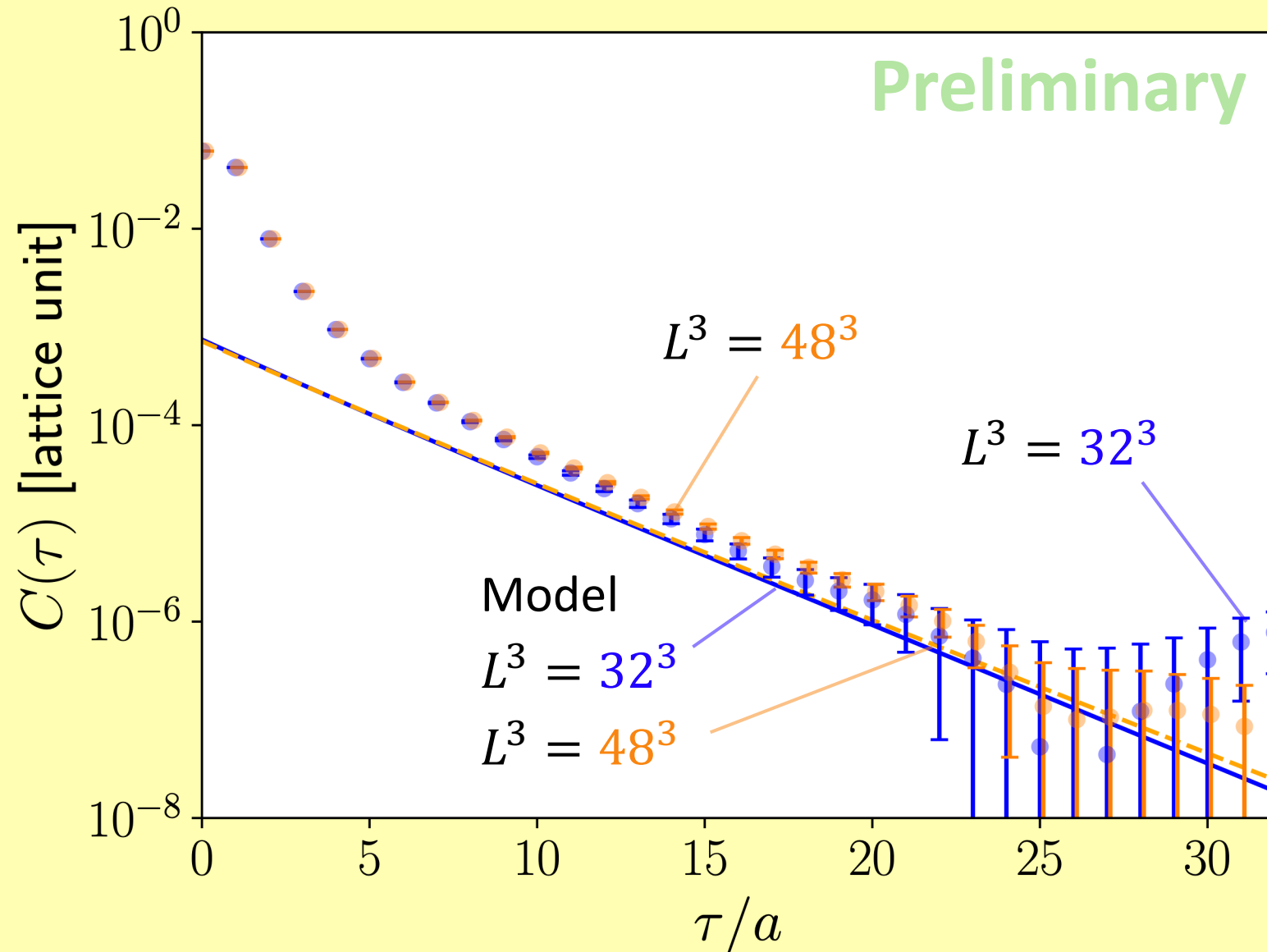
Infinite-V form factor Finite-V amplitude

- Cutoff at $E < 1.5$ GeV
- 4π or higher states ignored

... only large τ should be considered



Default model compared to the lattice data 9



$$C(\tau) = \int d\omega \rho(\omega) e^{-\omega\tau}$$

Lattice data:

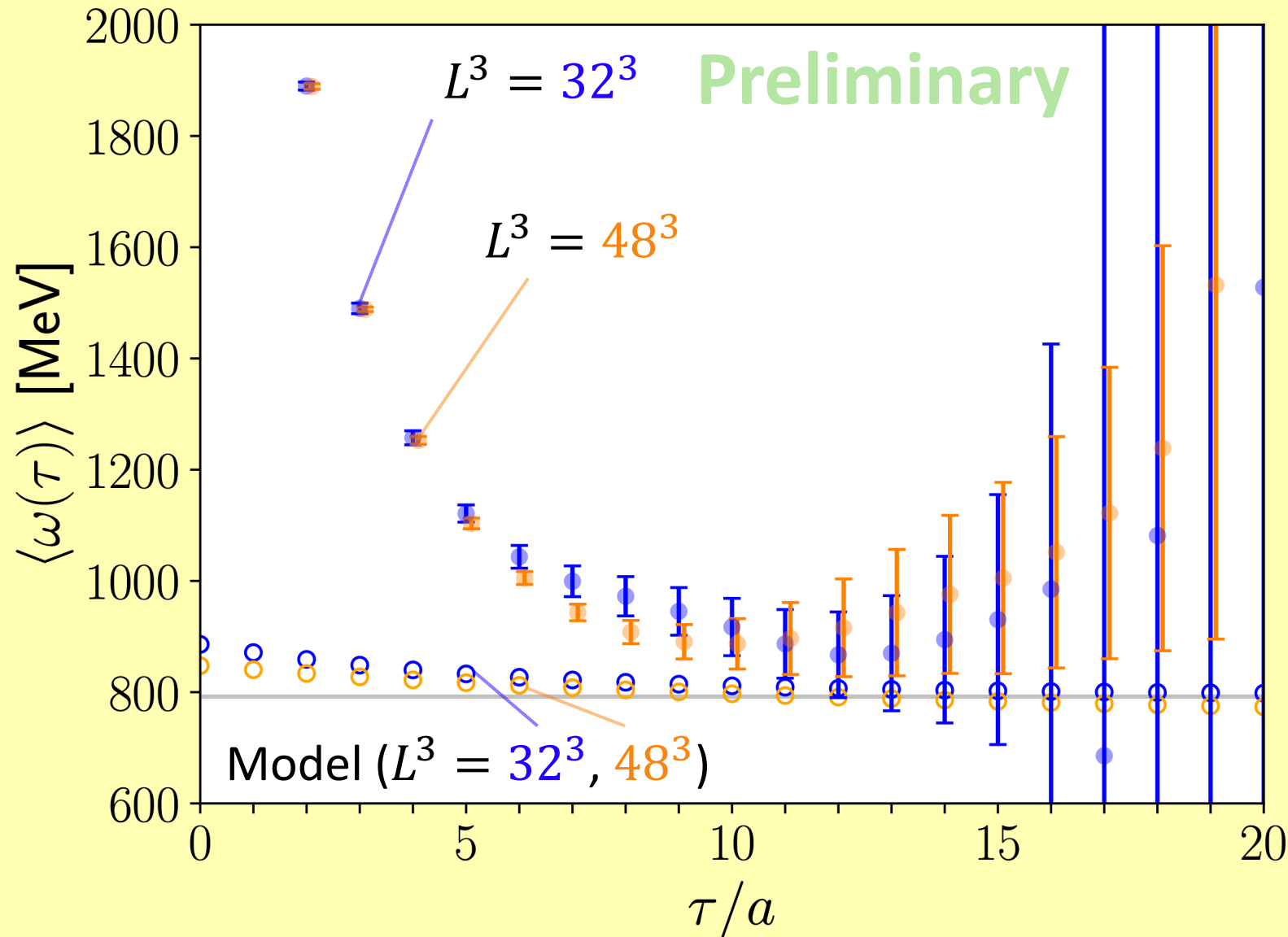
JLQCD, domain-wall lattice,
 $32^3 \times 64$, $48^3 \times 96$,

$m_\pi = 226\text{MeV}$,

$a \sim 0.08\text{ fm}$

conf = 100,

measurement
= 600 (32^3), 800 (48^3)



$$\langle \omega \rangle(\tau) = \int d\omega \rho(\omega) \omega e^{-\omega\tau}$$

Lattice data:

JLQCD, domain-wall lattice,
 $32^3 \times 64$, $48^3 \times 96$,

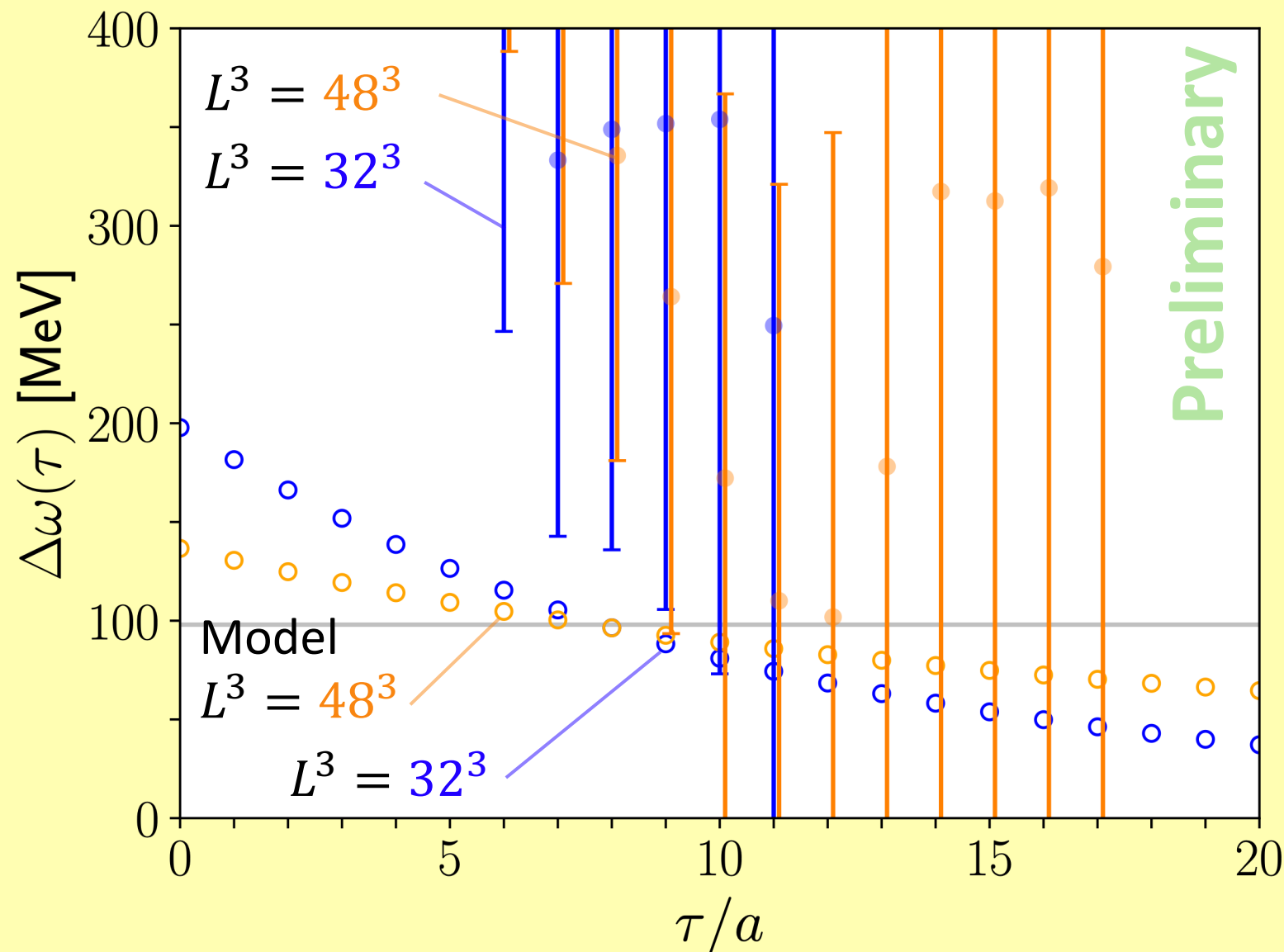
$m_\pi = 226 \text{ MeV}$,

$a \sim 0.08 \text{ fm}$

conf = 100,

measurement

= 600 (32^3), 800 (48^3)



$$\Delta\omega(\tau) = \sqrt{\langle\omega^2\rangle(\tau) - \langle\omega\rangle^2(\tau)}$$

Lattice data:

JLQCD, domain-wall lattice,

$32^3 \times 64$, $48^3 \times 96$,

$m_\pi = 226\text{MeV}$,

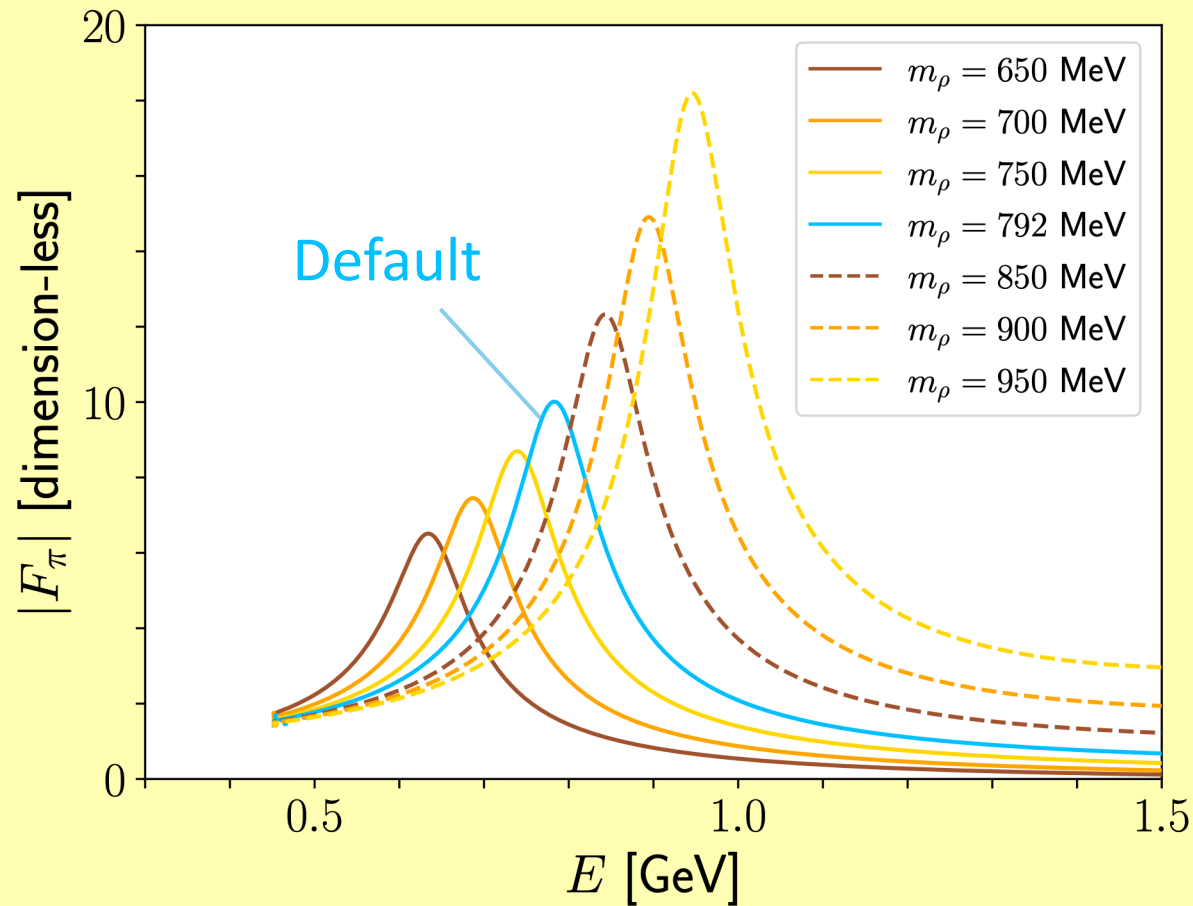
$a \sim 0.08$ fm

conf = 100,

measurement
= 600 (32^3), 800 (48^3)

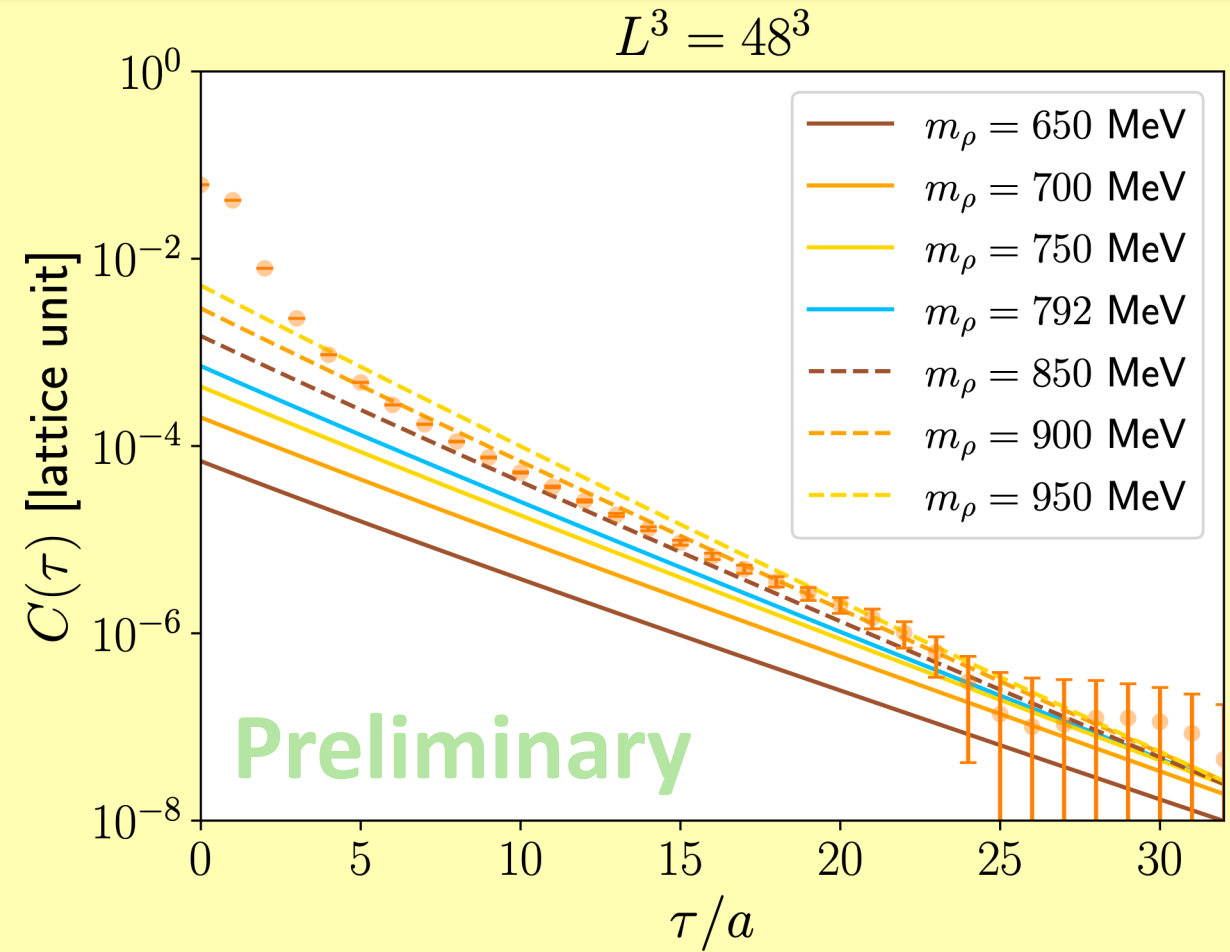
m_ρ -sensitivity of $\mathcal{C}(\tau)$

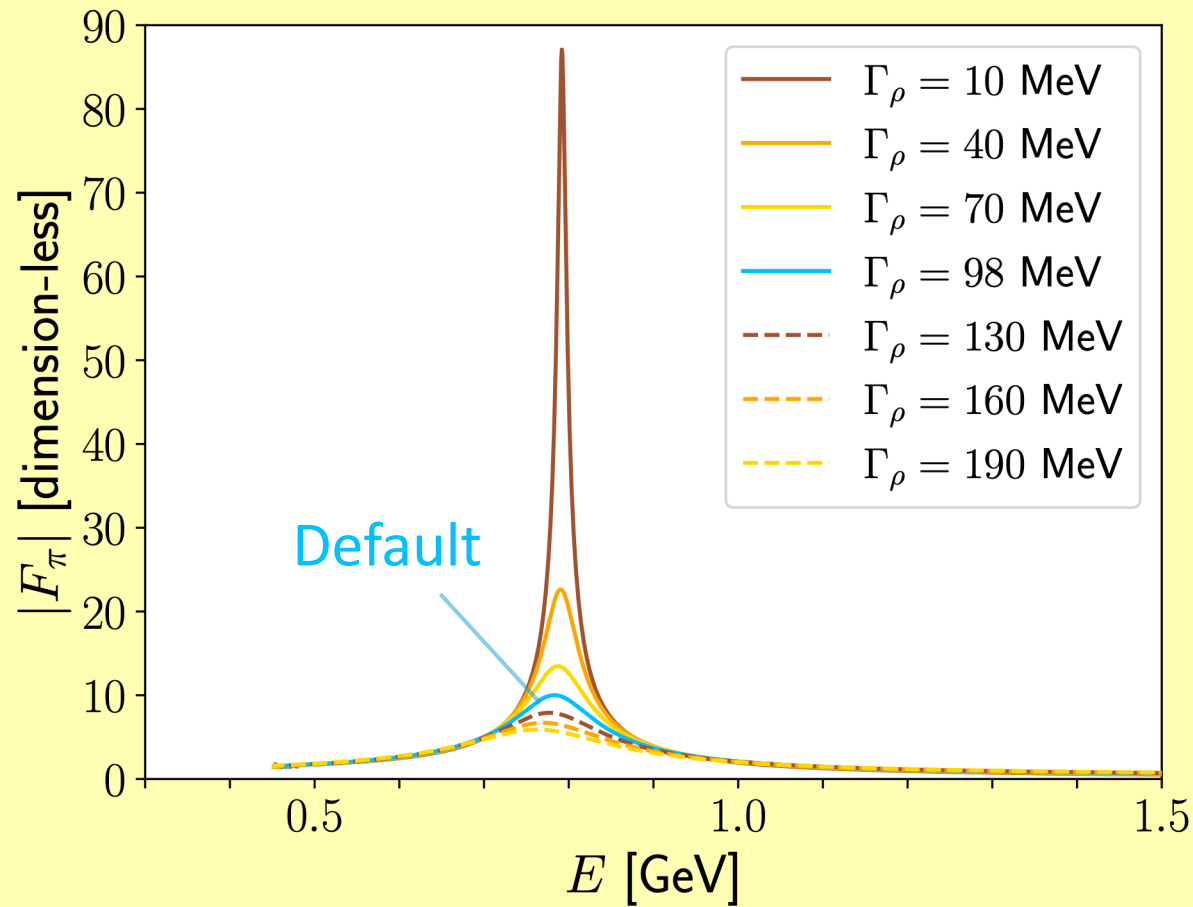
12



Default

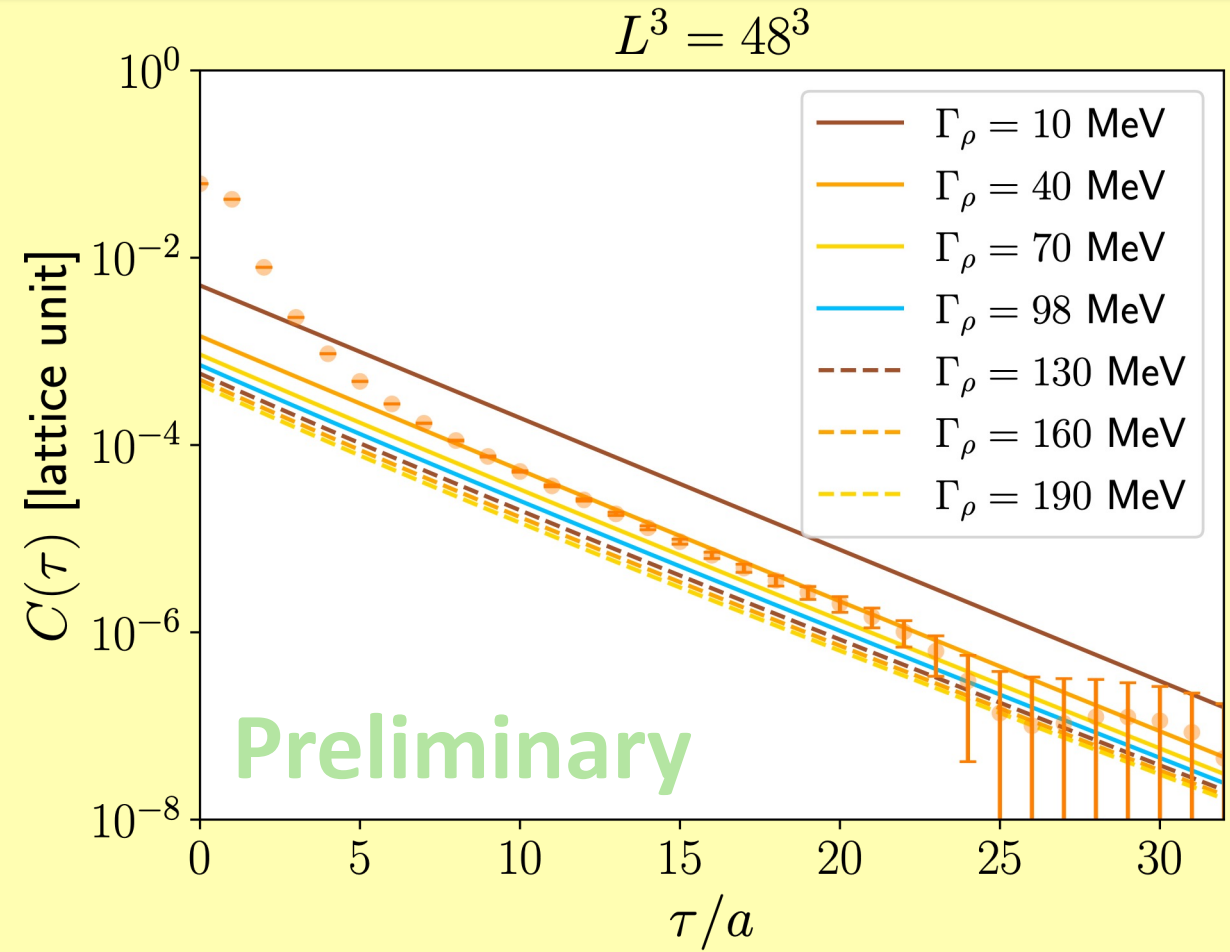
m_π [MeV]	m_ρ [MeV]	Γ_ρ [MeV]	$\sqrt{\langle r_\pi^2 \rangle}$ [fm]
226	792	98	0.635





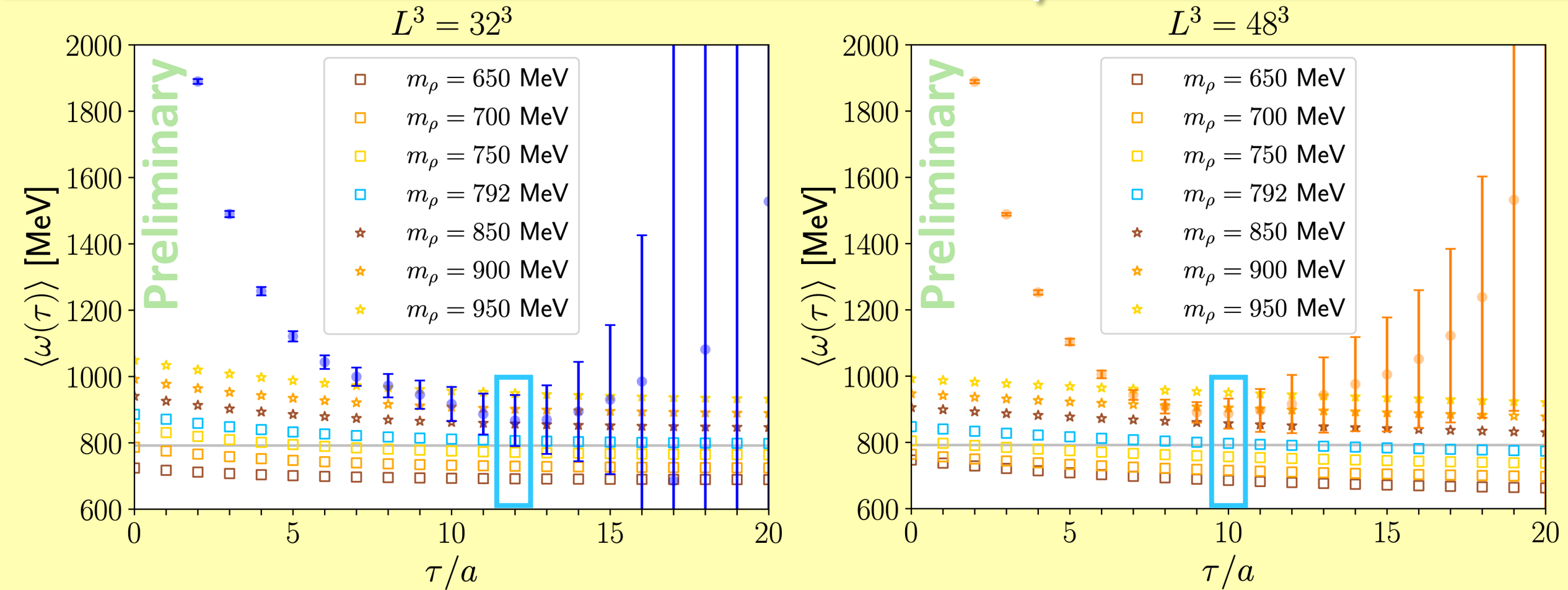
Default

m_π [MeV]	m_ρ [MeV]	Γ_ρ [MeV]	$\sqrt{\langle r_\pi^2 \rangle}$ [fm]
226	792	98	0.635



1st-order moment (different m_ρ s)

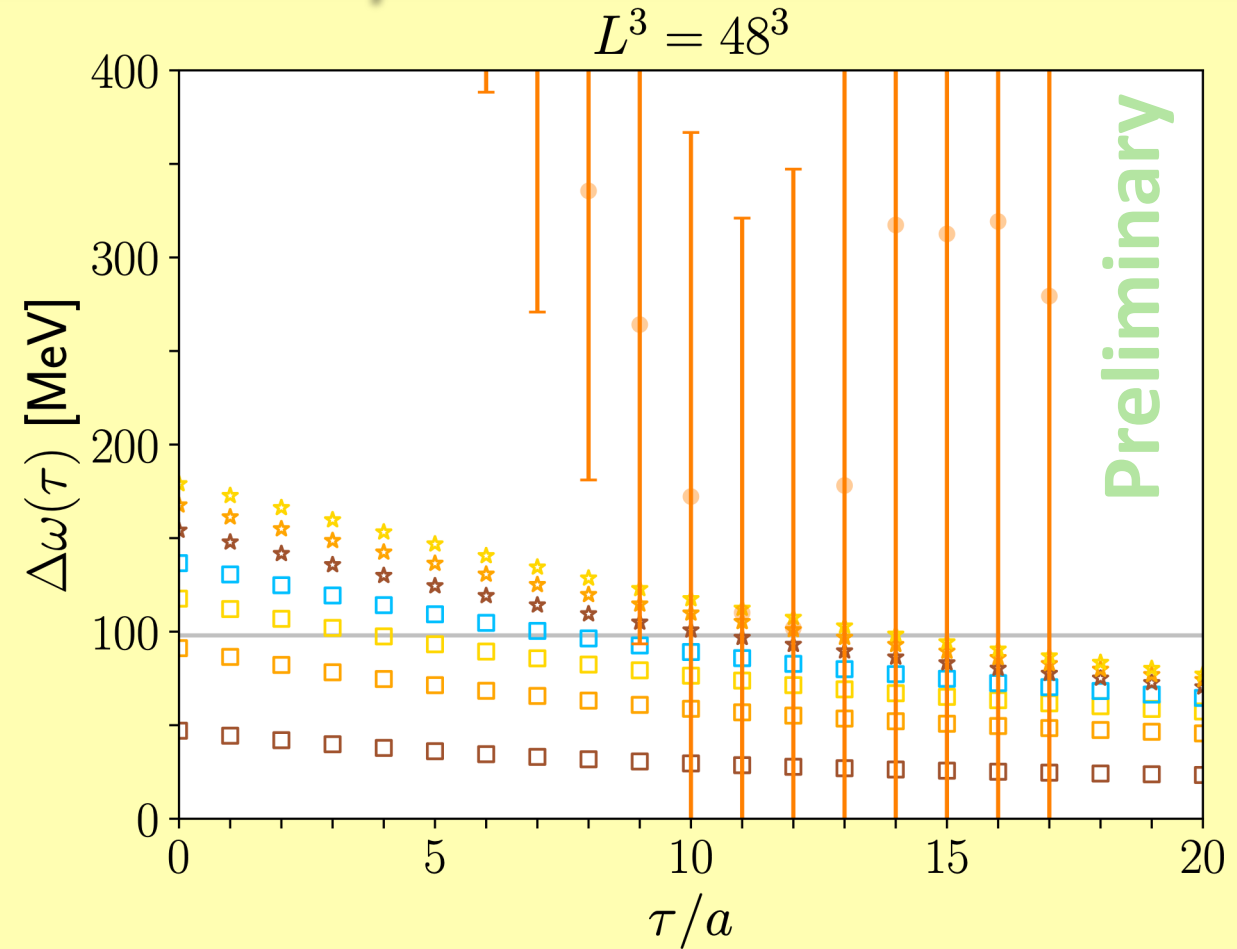
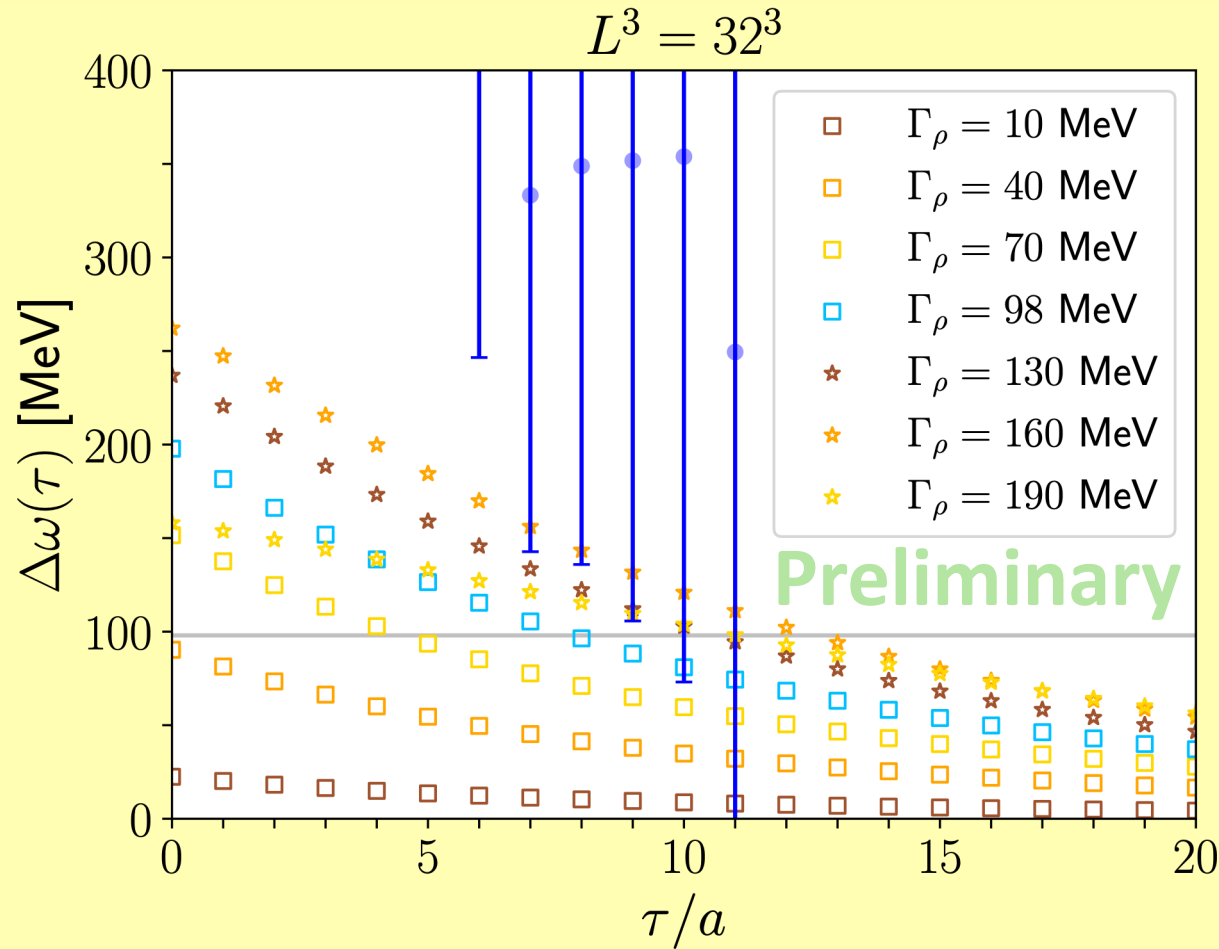
14



Estimation: $m_\rho \sim 900$ MeV

2nd-order moment (different Γ_ρ s)

15



Need more statistics

- Smeared spectrum ($\mathcal{C}(\tau)$, $\langle\omega\rangle(\tau)$, $\Delta\omega(\tau)$) constructed from a model for $\pi\pi$ states
- SVD of $e^{-E\tau}$ to obtain $\langle\omega\rangle(\tau)$ and $\Delta\omega(\tau)$ from the lattice correlator

As m_ρ and Γ_ρ increase...

- Sensitivity for m_ρ and Γ_ρ
 - More statistics is required to perform detailed study
 - m_ρ, Γ_ρ -fit will be performed to the lattice results

	m_ρ	Γ_ρ
$\mathcal{C}(\tau)$	↗	↘
$\langle\omega\rangle(\tau)$	↗	
$\Delta\omega(\tau)$		↗

Thank you for your attention.



Backup slides

$$C(\tau) = \int dE \rho(E) e^{-E\tau}$$

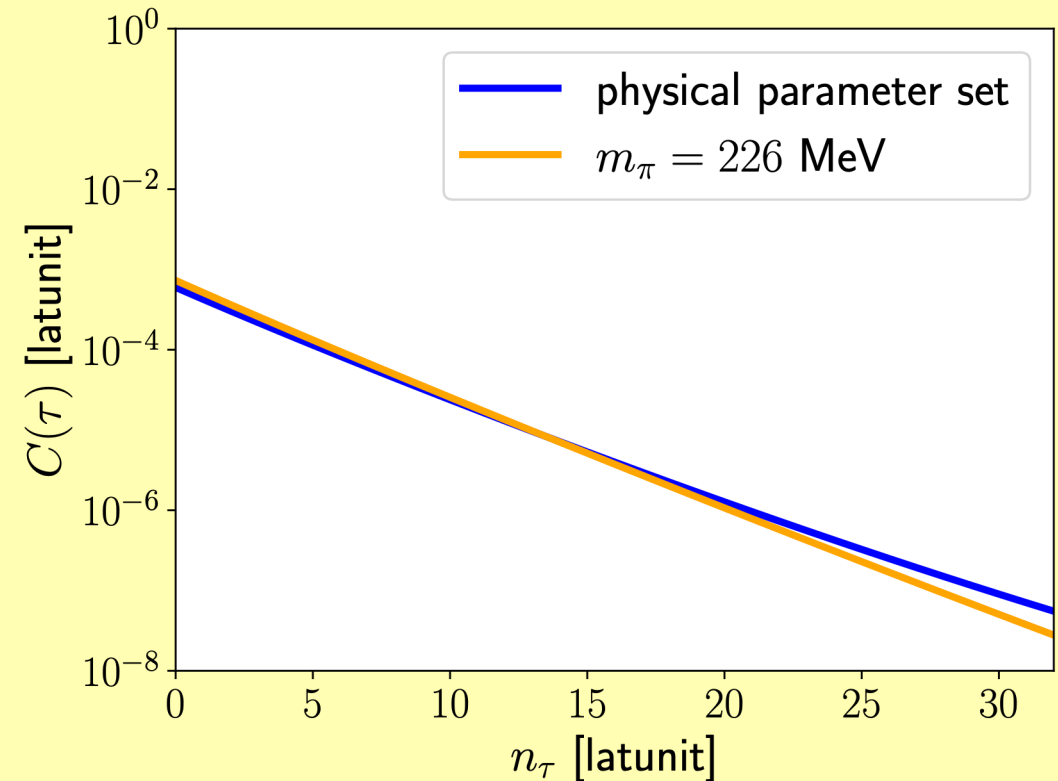
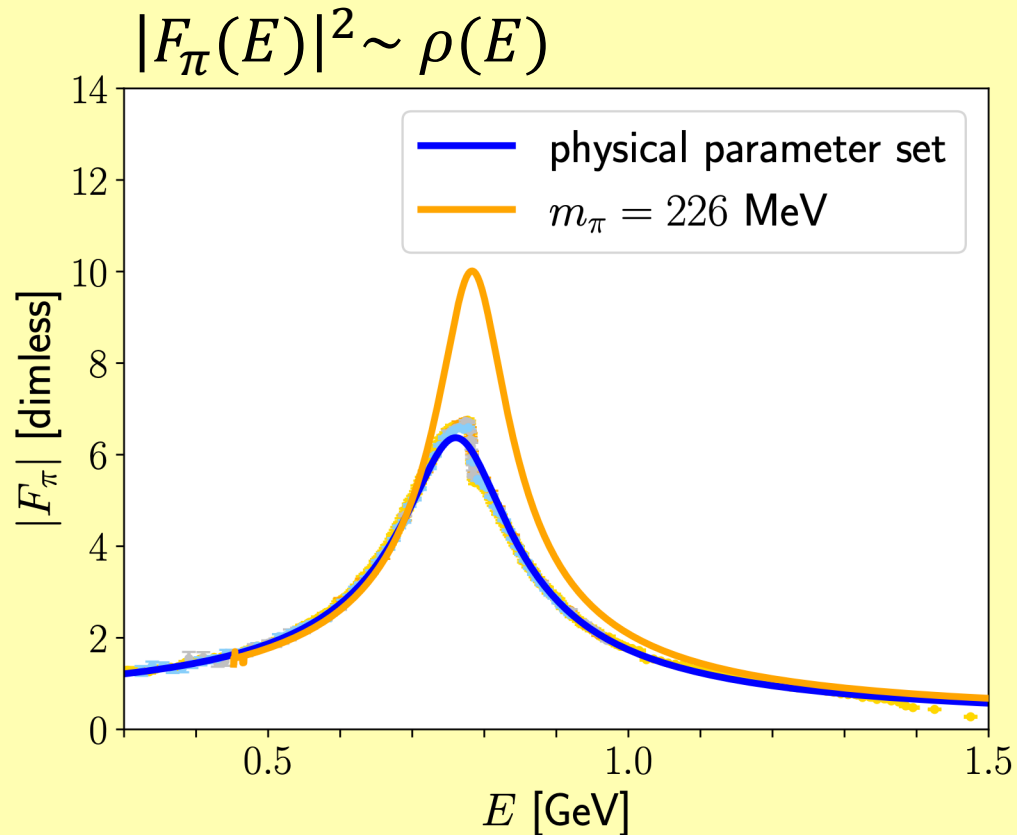
It is hard to obtain the spectral function $\rho(E)$ with a high resolution from the lattice data of $C(\tau)$.

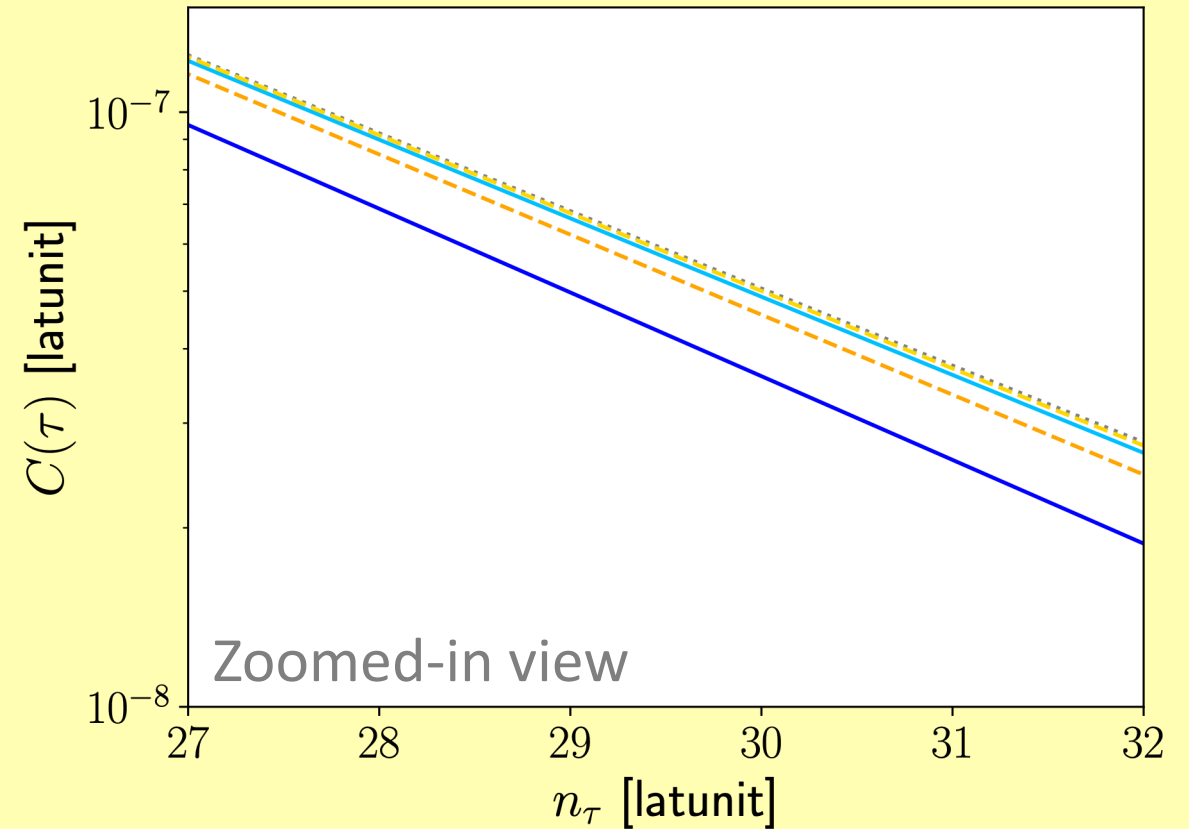
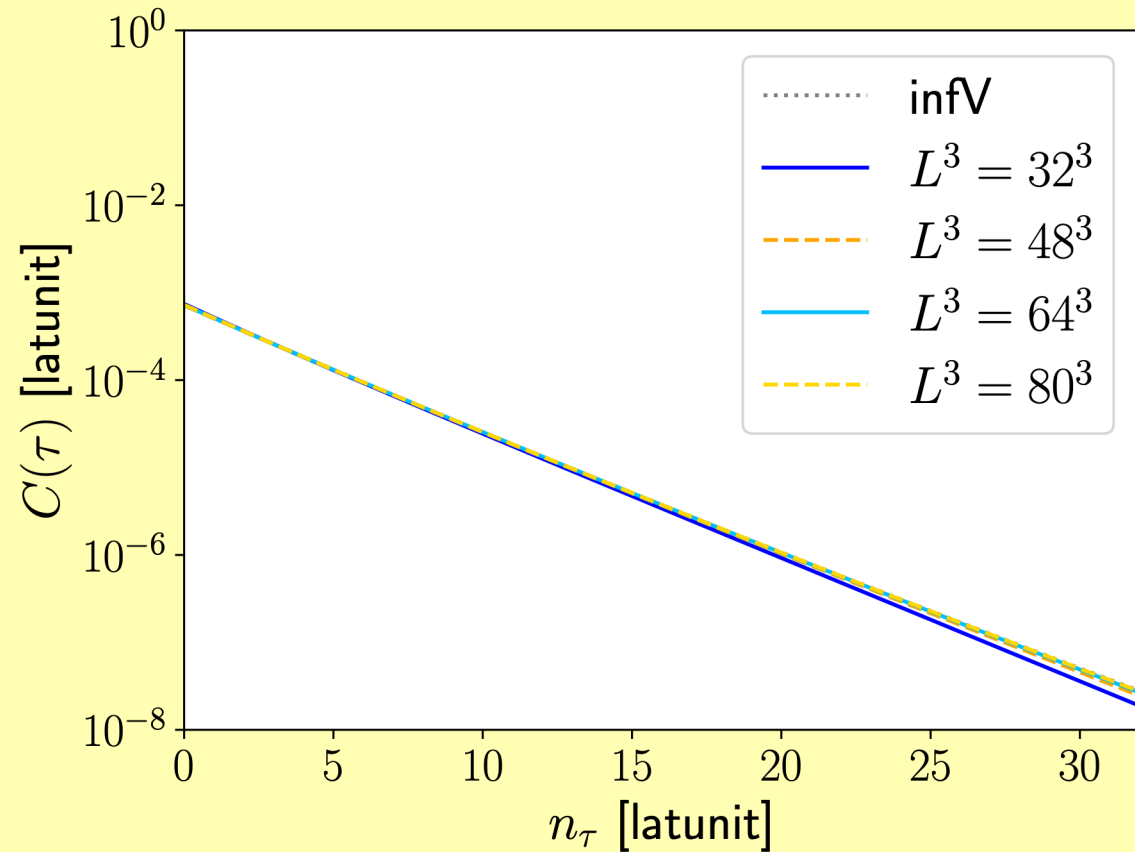
- Highly suppressed excited states by $e^{-E\tau}$
- A finite set of lattice data
- Finite precision

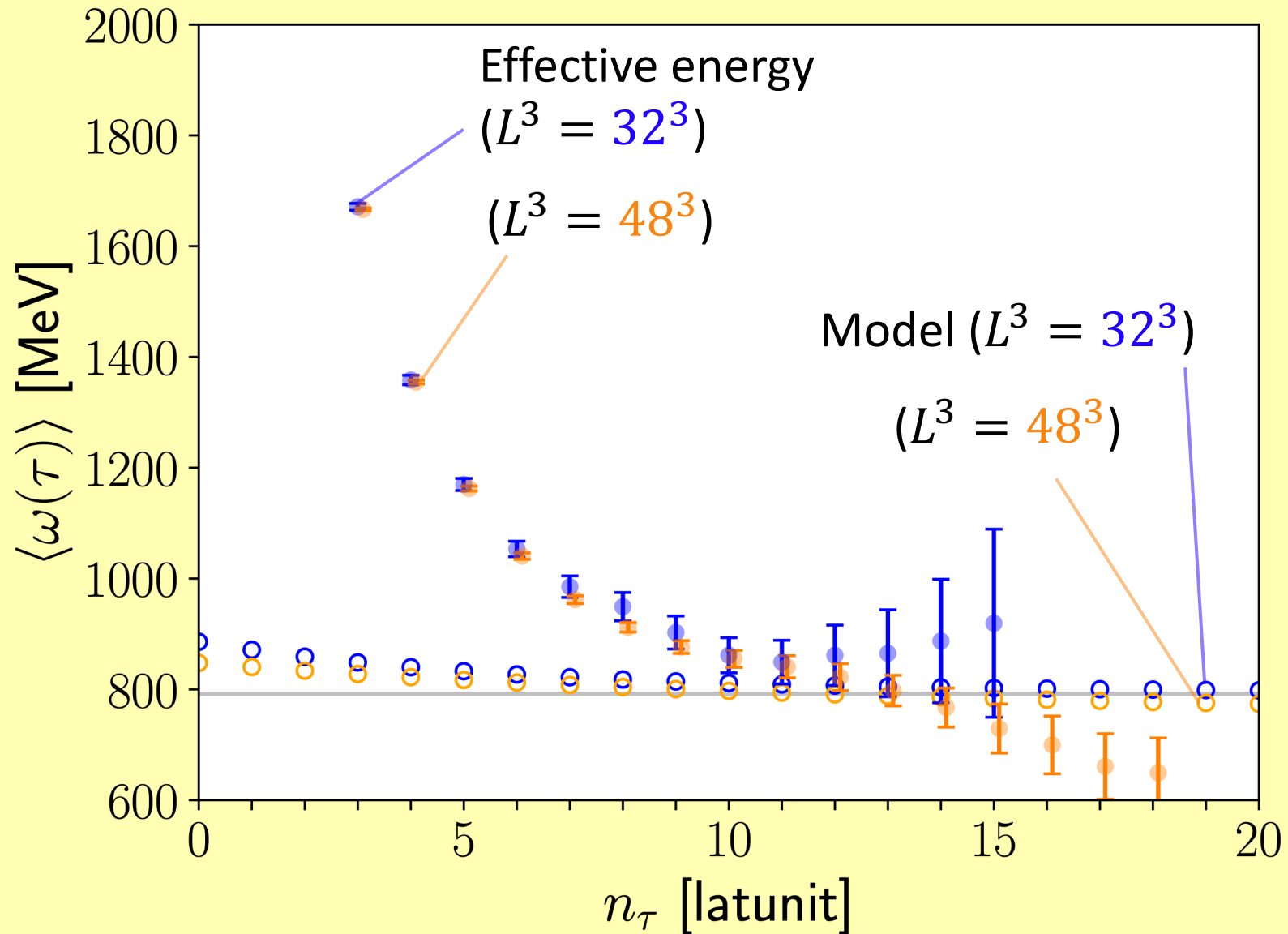
Consider a **smeared** spectral function.

Infinite volume:

$$C(\tau) = \int dE \, \rho^{\text{BW}}(E) \, e^{-E\tau}$$

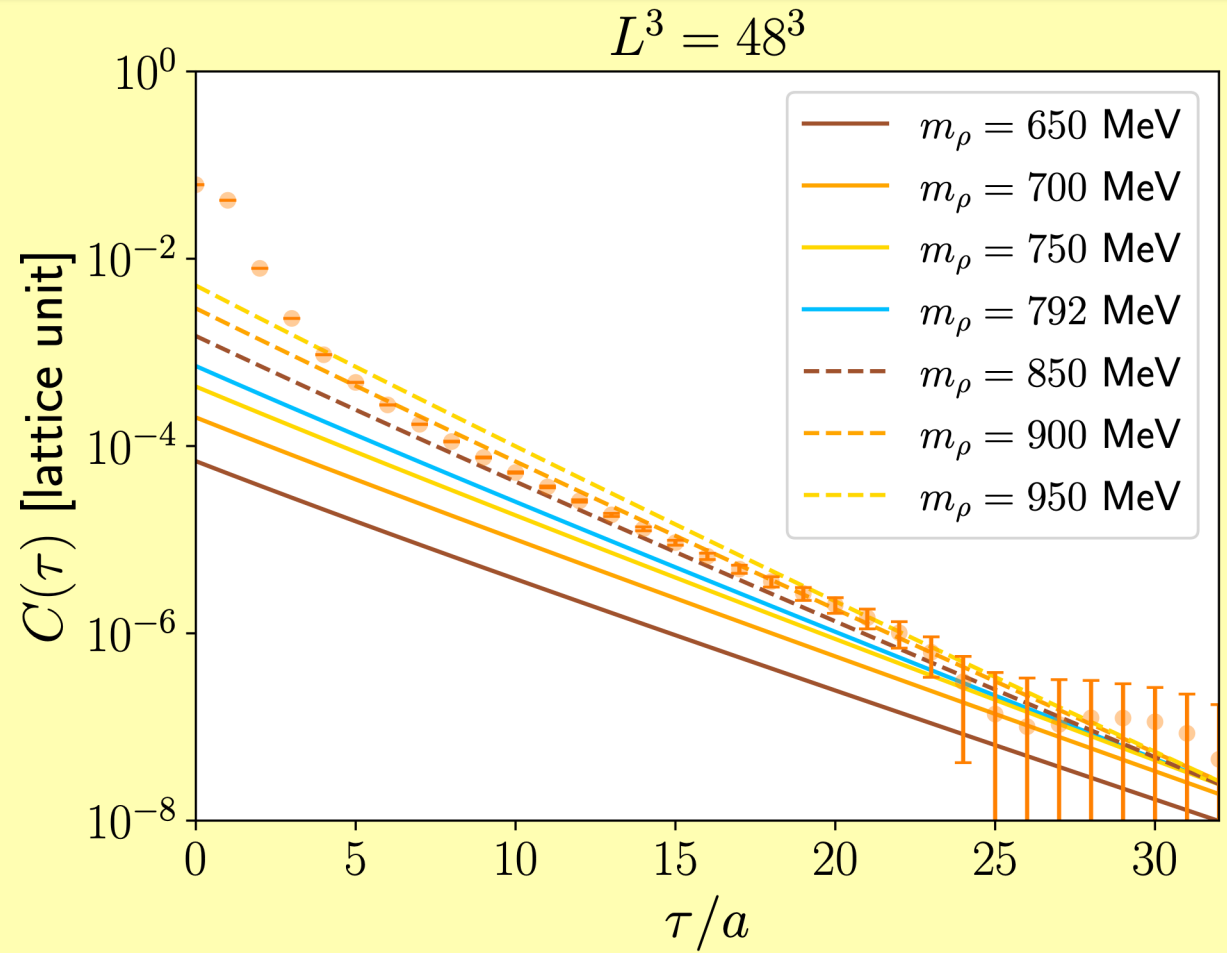
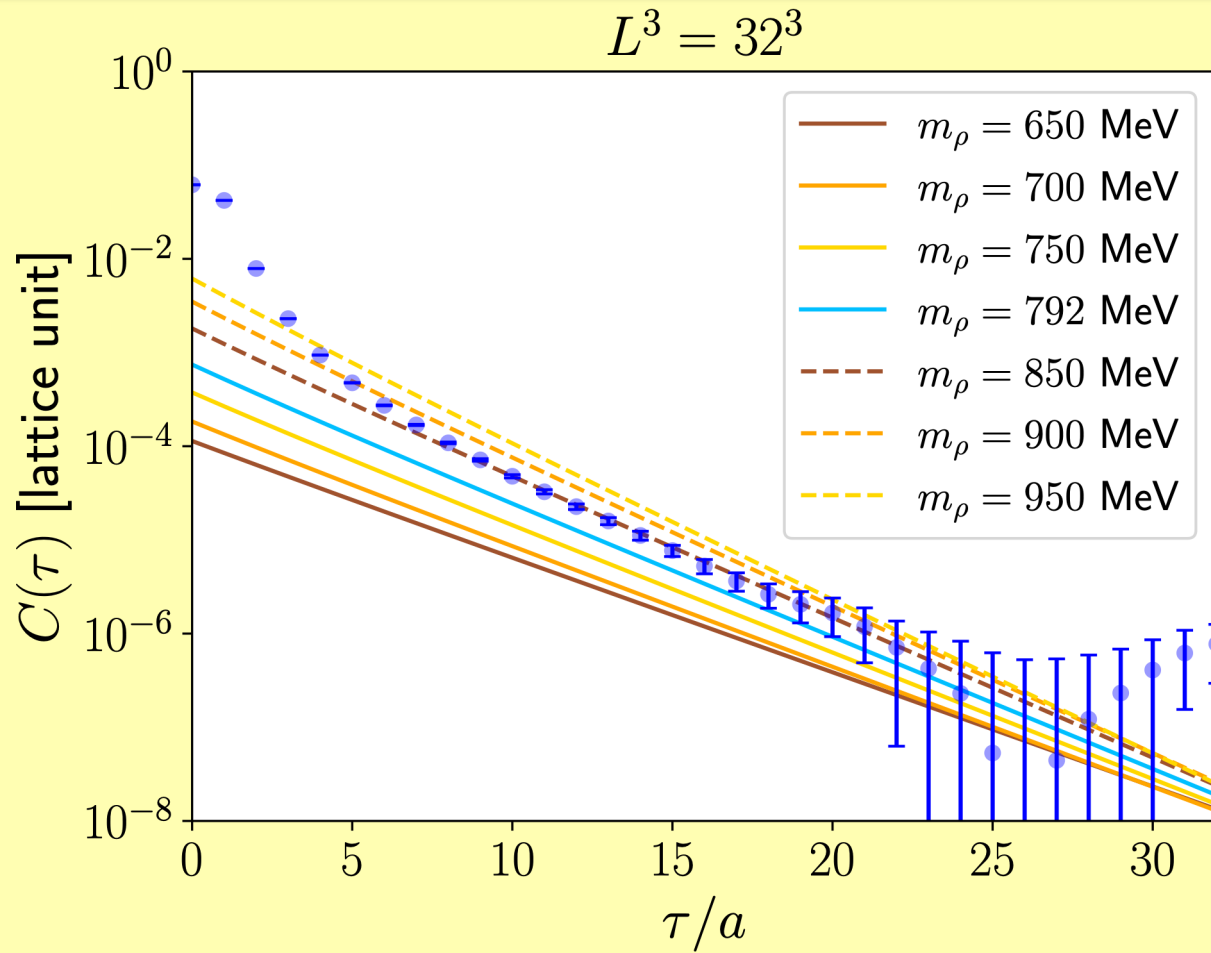






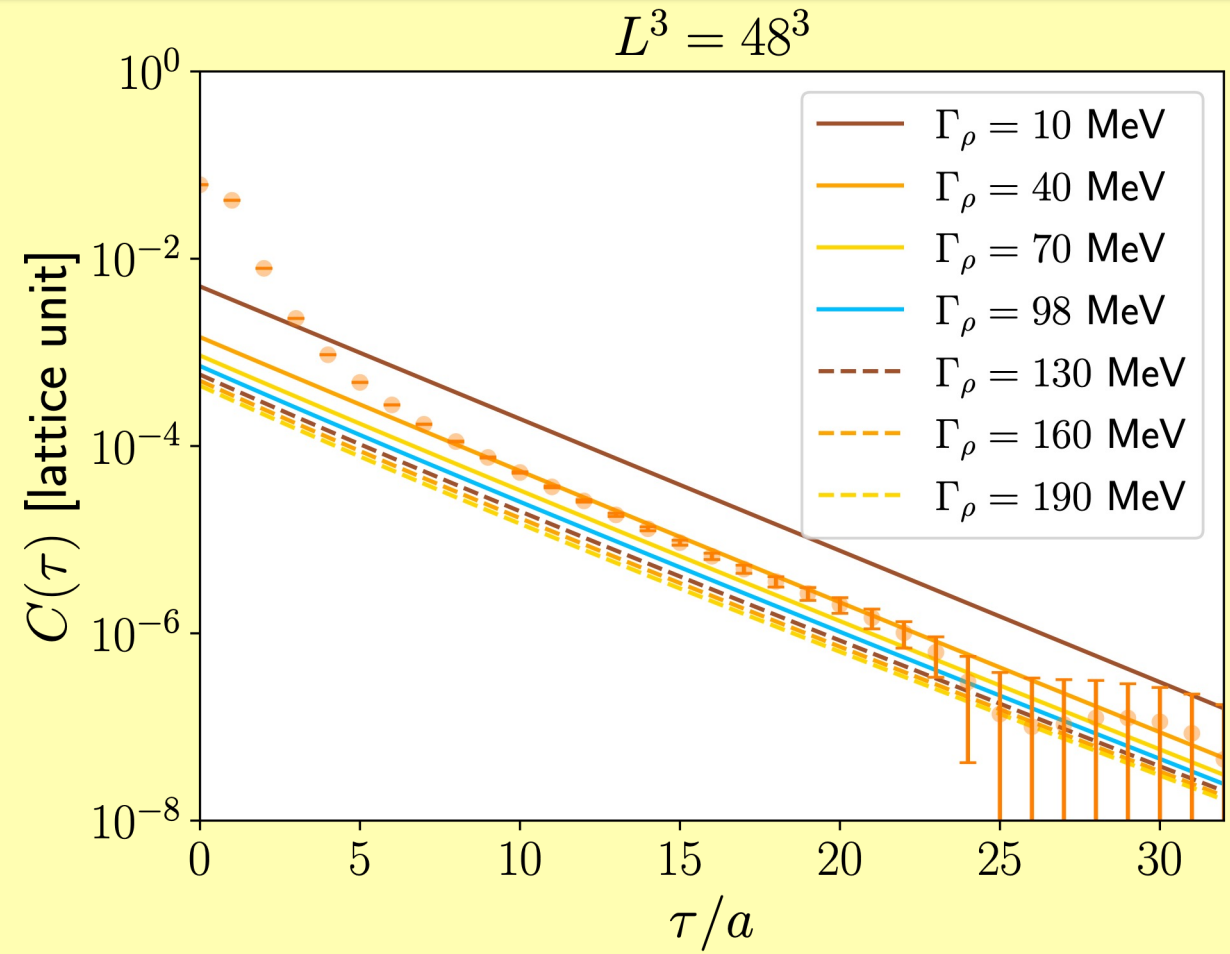
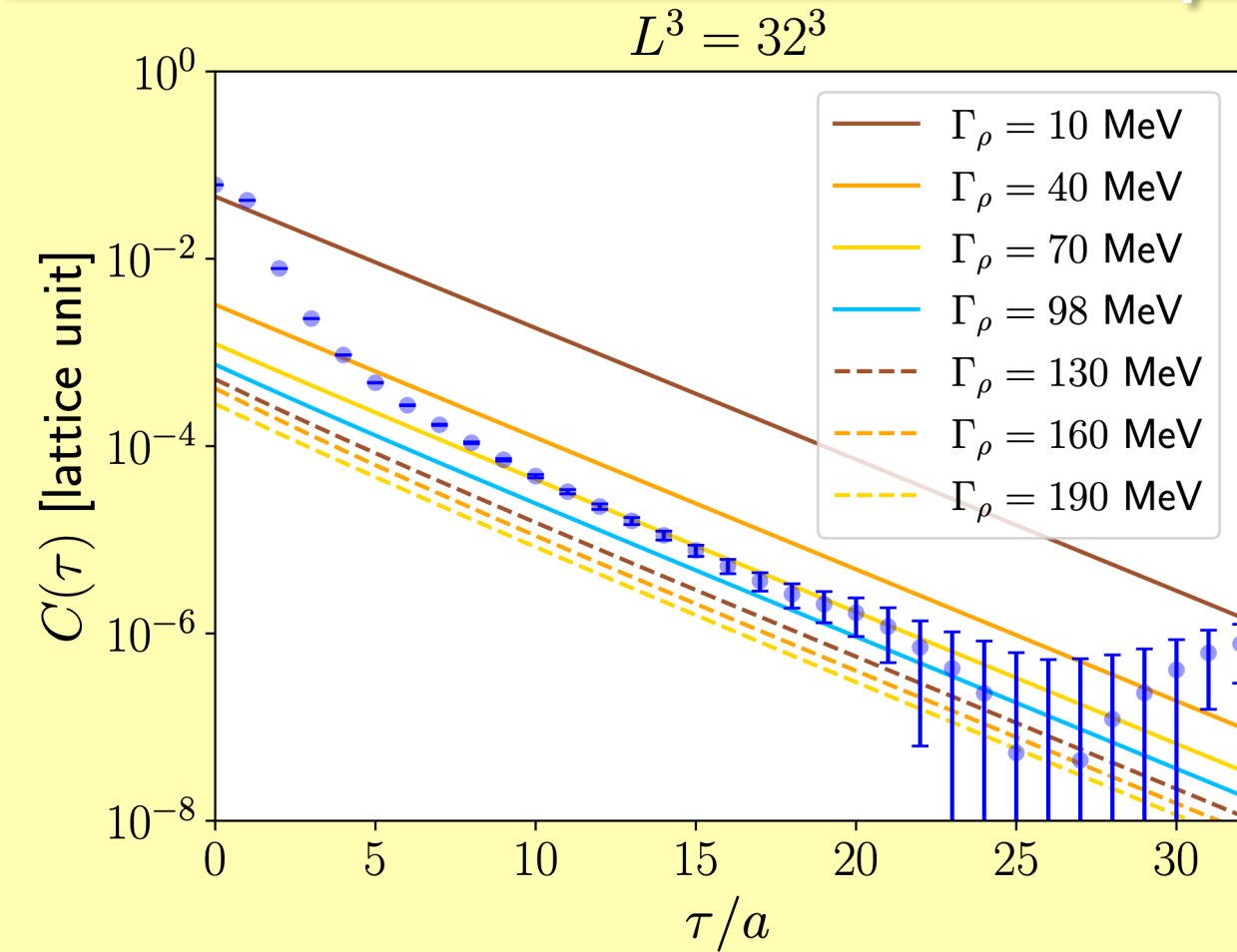
Correlator (different m_ρ s)

22



Correlator (different Γ_ρ s)

23



1st-order moment (different m_ρ s)

24

