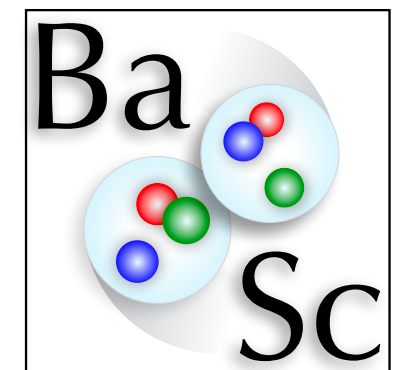


NN scattering amplitudes at light quark masses

Miguel Salg

For the Baryon Scattering (BaSc) Collaboration



Lattice 2026, University of Maryland, College Park, July 26 — August 1, 2026



Collaborators



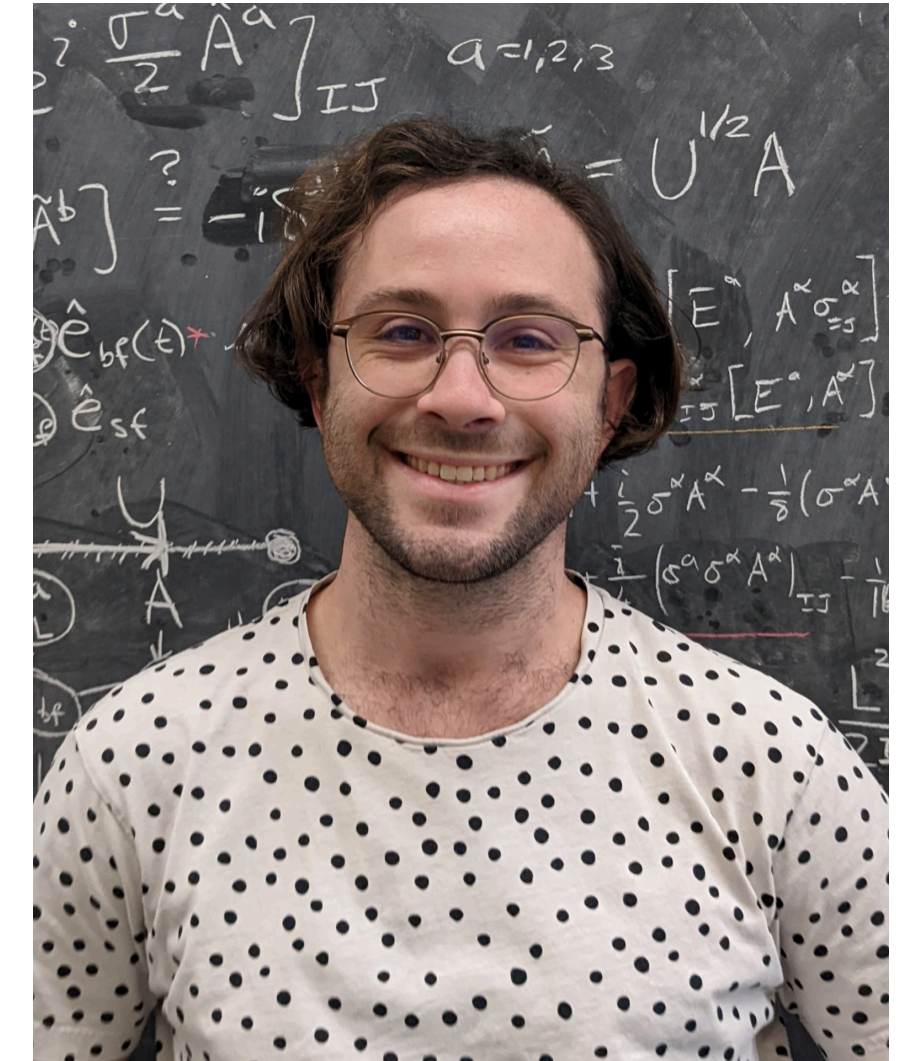
Jorge Baeza-Ballesteros



André Baião Raposo



Davide Laudicina



Malcolm Lazarow

John Bulava

Andrew Hanlon

Colin Morningstar

Fernando Romero-López

Hartmut Wittig

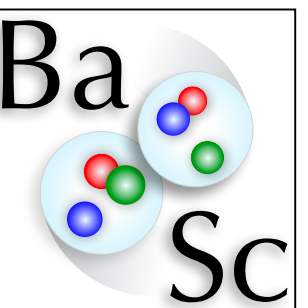
Jeremy Green

Daniel Mohler

Amy Nicholson

André Walker-Loud

+ BaSc collaboration



NN scattering from lattice QCD

Deuteron towards the physical pion mass

Bulava et al. (BaSc) 2026 [PRC 113, 024002]

- **Large pion masses** ($M_\pi \approx 710$ MeV): both $I = 0$ and $I = 1$ NN interaction **attractive**, but **not strong** enough to form bound states

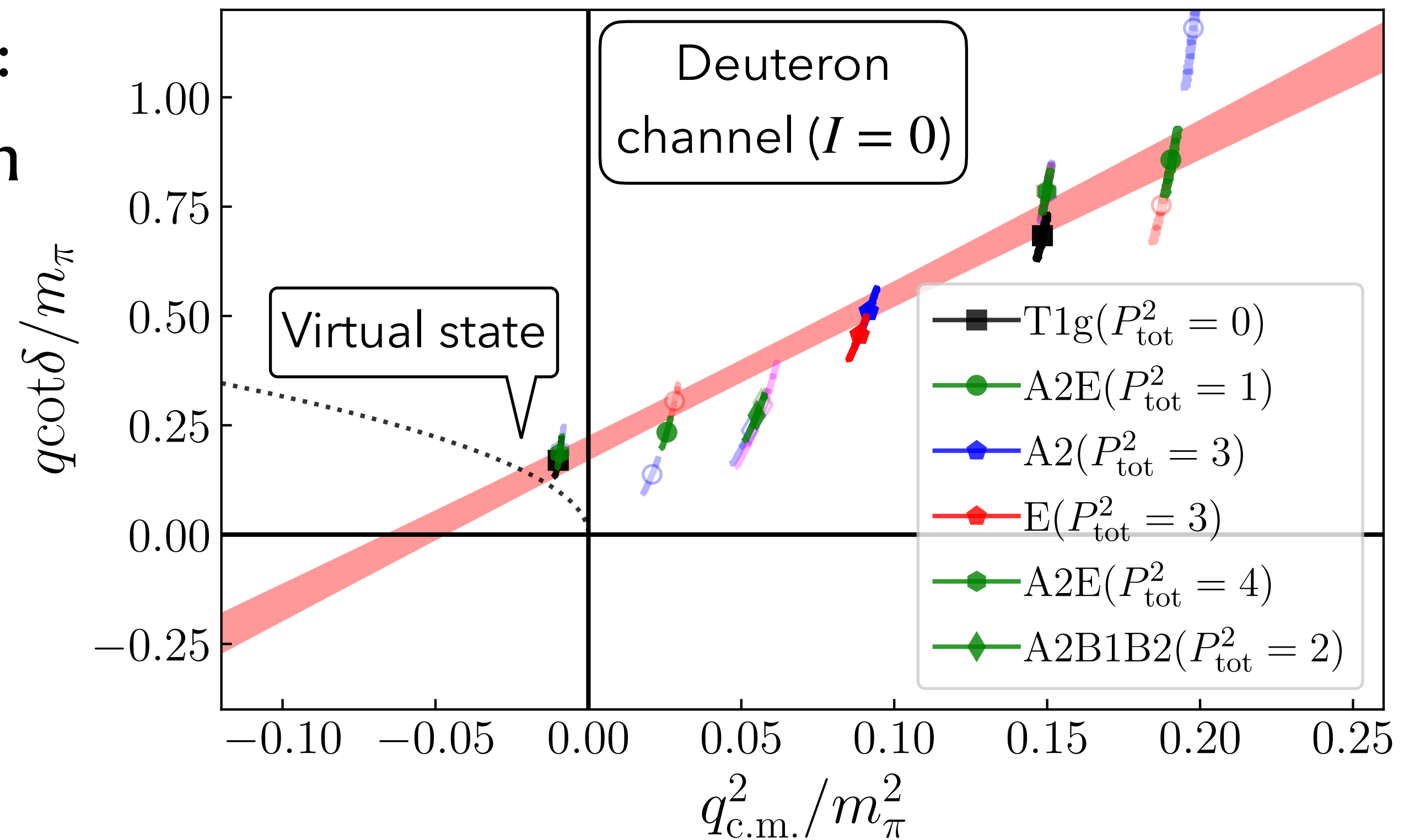
—> deuteron is **virtual state**

Inoue et al. (HAL QCD) 2012 [Nucl. Phys. A 881, 28]

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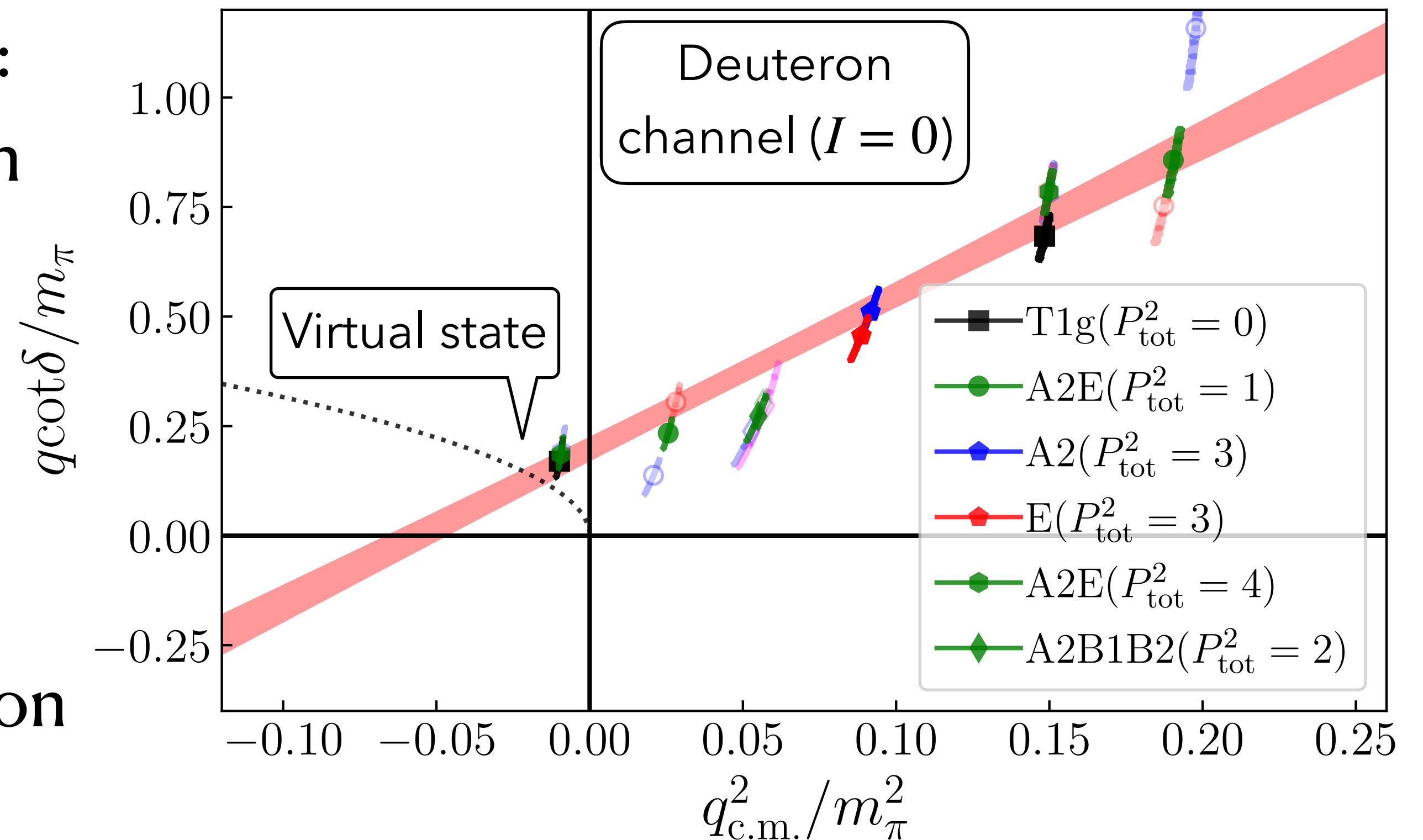
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- **Now**: preliminary results at $M_\pi = 280$ MeV (CLS D251) and $M_\pi = 200$ MeV (J250)



NN scattering from lattice QCD

Robert Perry, plenary Thu 09:30

Joseph Moscoso, Fri 15:00

William Detmold, Fri 16:10

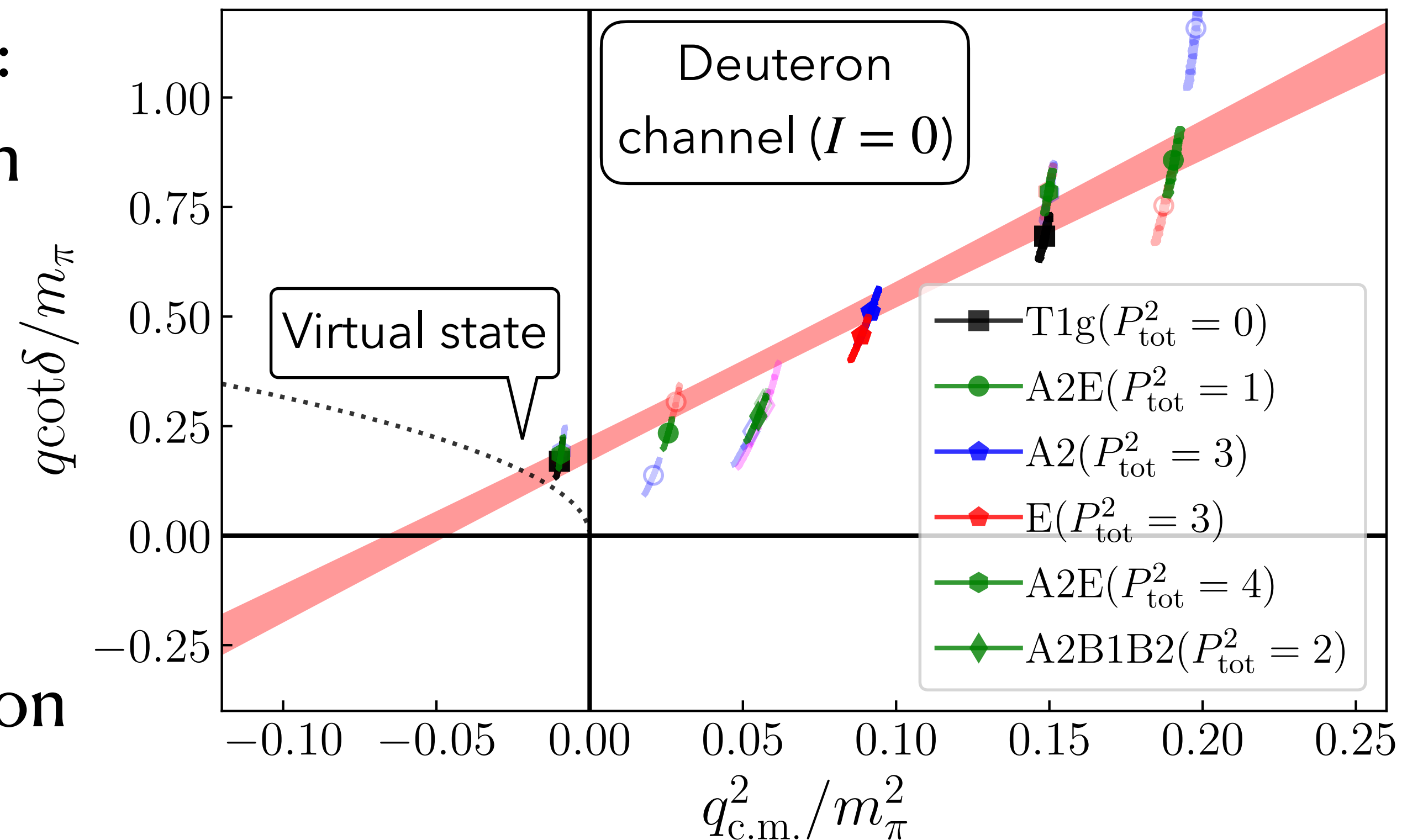
Michael Wagman, Fri 16:30

Jorge Baeza-Ballesteros, poster session A

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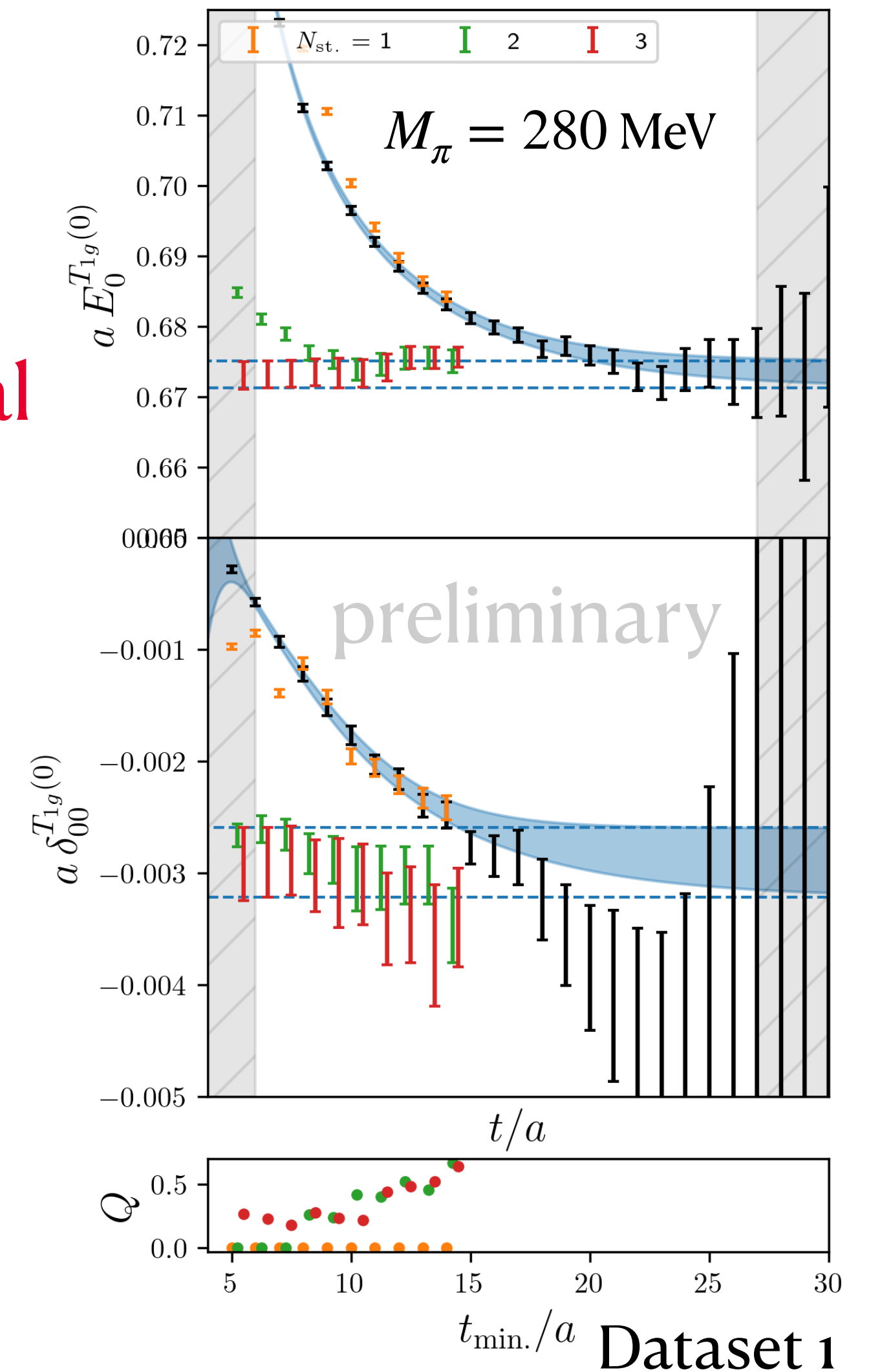
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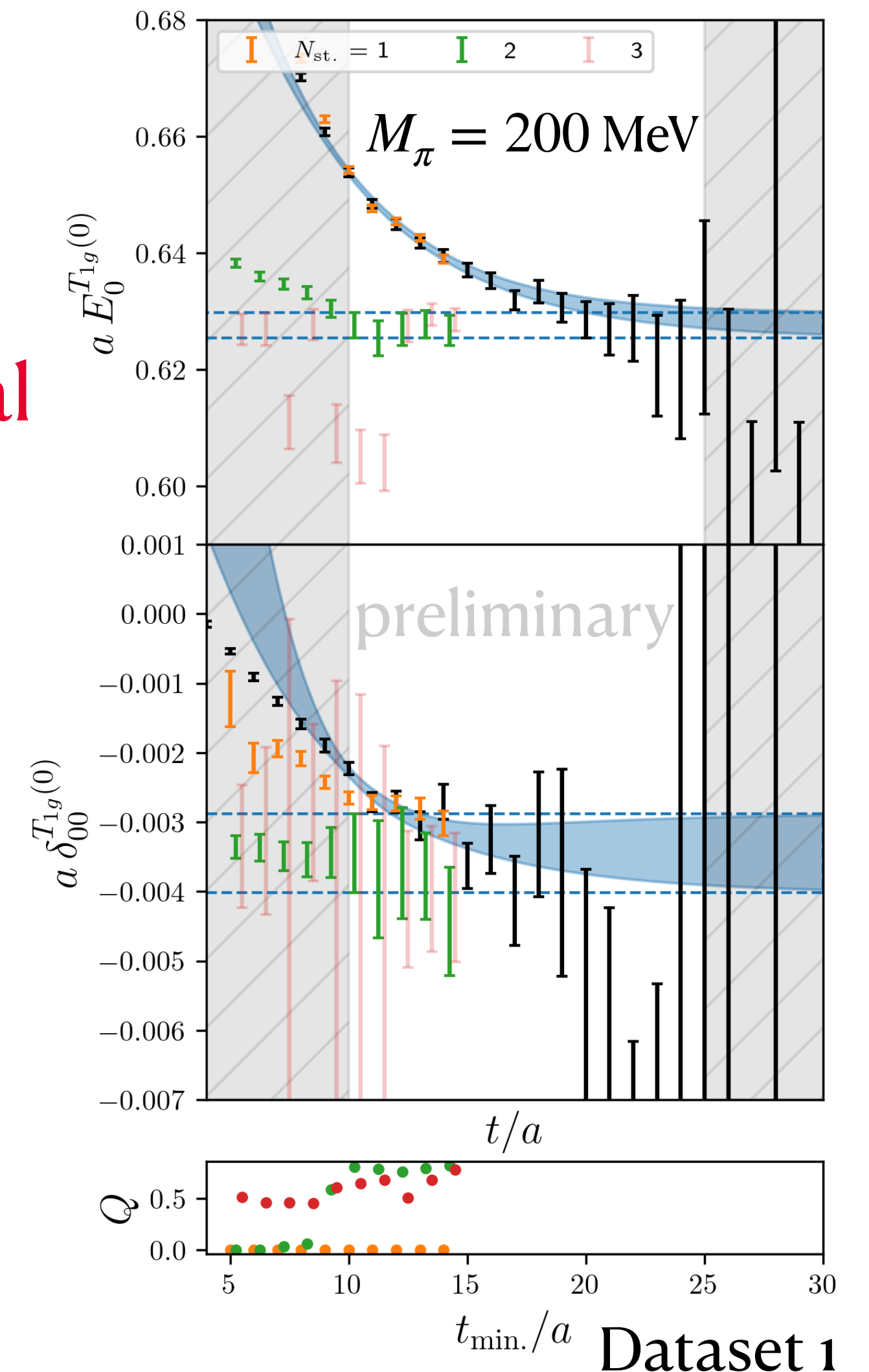
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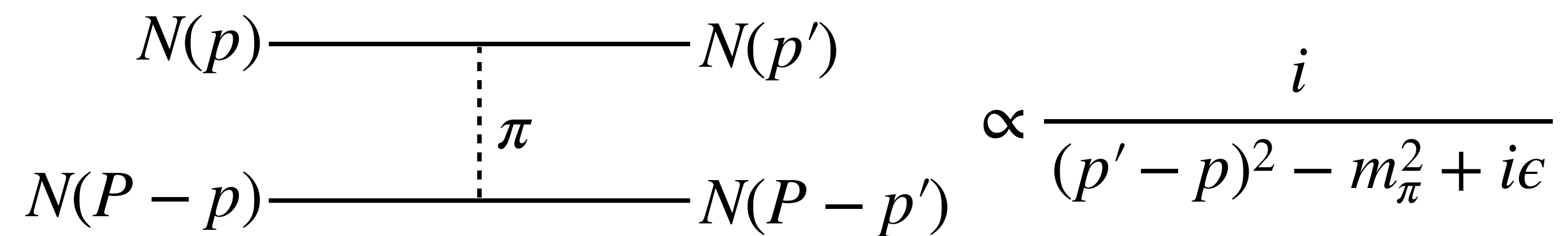
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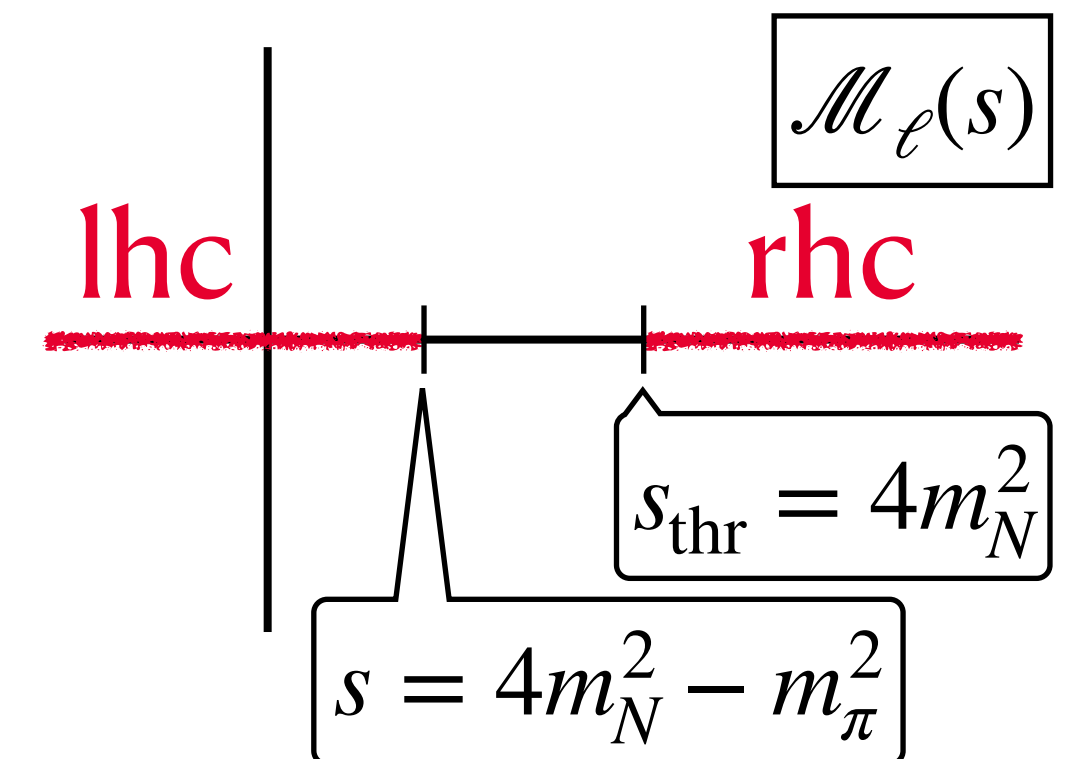


NN scattering: left-hand cuts

- t -channel exchange of light particles (here: pions) leads, after partial-wave projection, to **left-hand cut** (lhc) at $-\infty < E^2 < 4m_N^2 - m_\pi^2$ in scattering amplitude



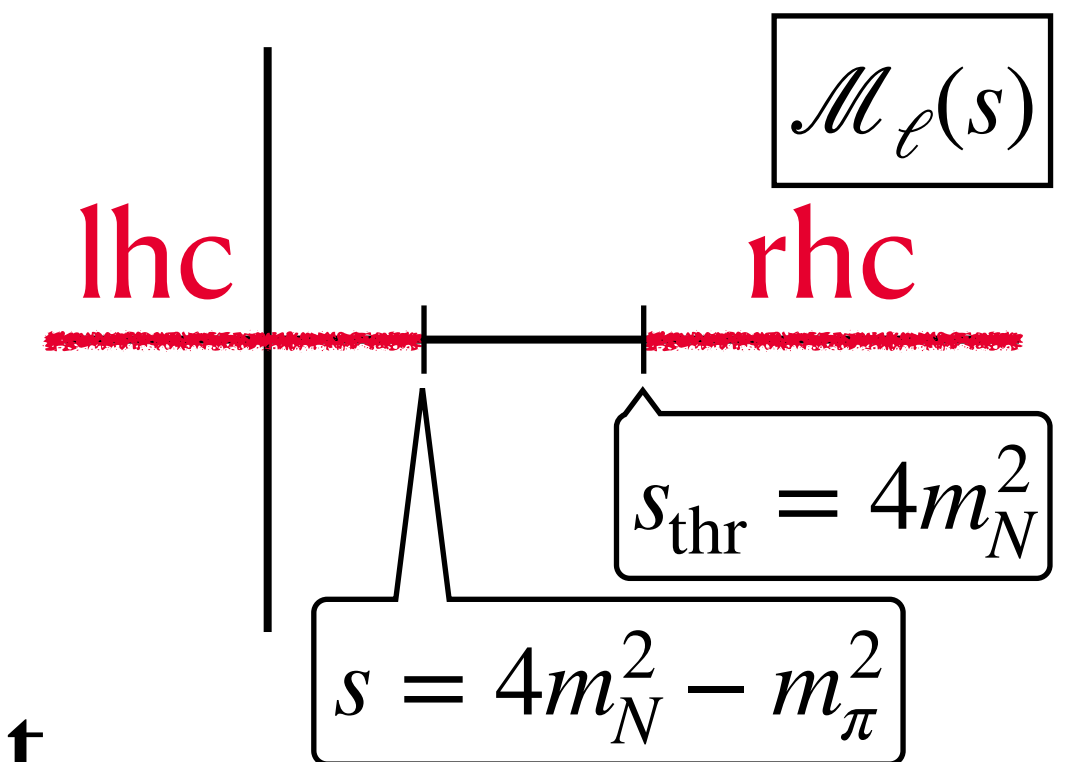
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- This induces **uncontrolled finite-volume effects** close to the cut
- Here: two **alternatives** that are valid in the presence of lhcs
 - **Relativistic** formalism based on all-orders perturbation theory
 - QC in **plane-wave basis** $\rightarrow$ based on Lippmann-Schwinger eq. with NR **EFT** potential

Baião Raposo and Hansen 2024 [JHEP 2024, 075]

Baião Raposo et al. 2025 [JHEP 2025, 186]

Meng and Epelbaum 2021 [JHEP 2021, 051]

Meng et al. 2024 [PRD 109, L071506]

Jorge Baeza-Ballesteros, poster session A

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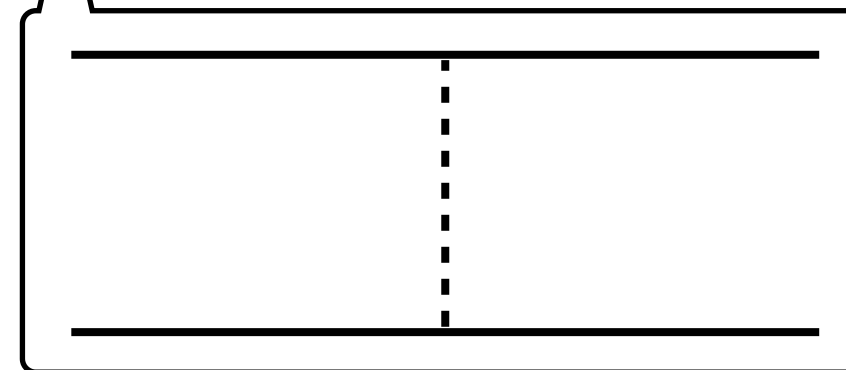
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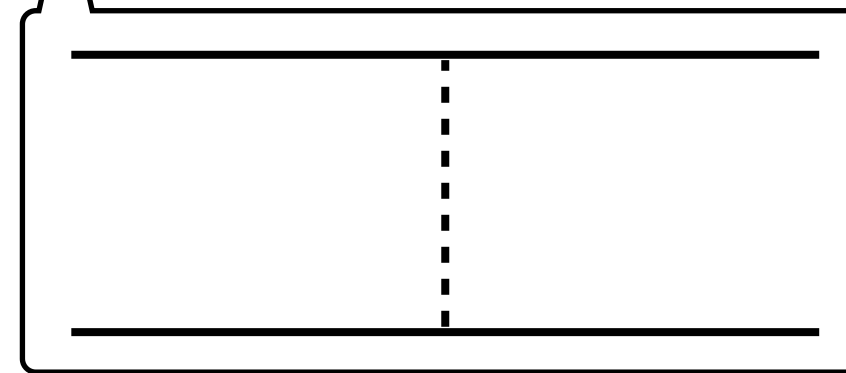
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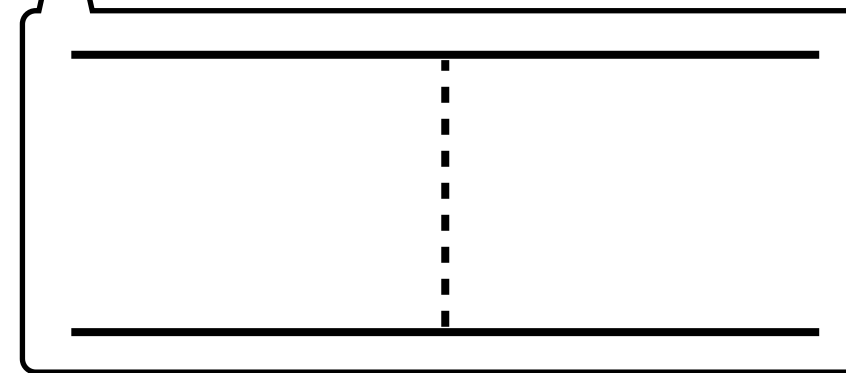
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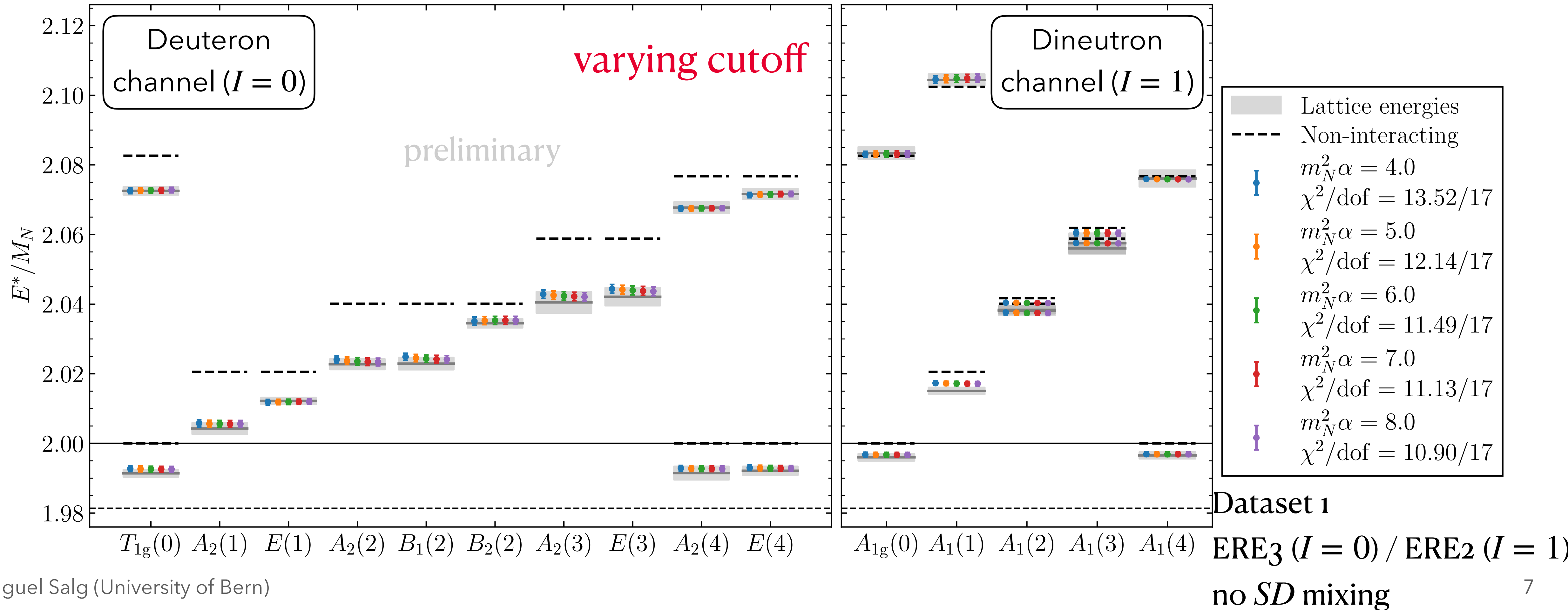


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- $\mathcal{C}_L$ contains inverse of matrix in plane-wave basis $\rightarrow$ numerically **more expensive**
- OPE** contributes to **all partial waves** since partial-wave projection is done only at the end

NN scattering on D251 ($M_\pi = 280$ MeV)

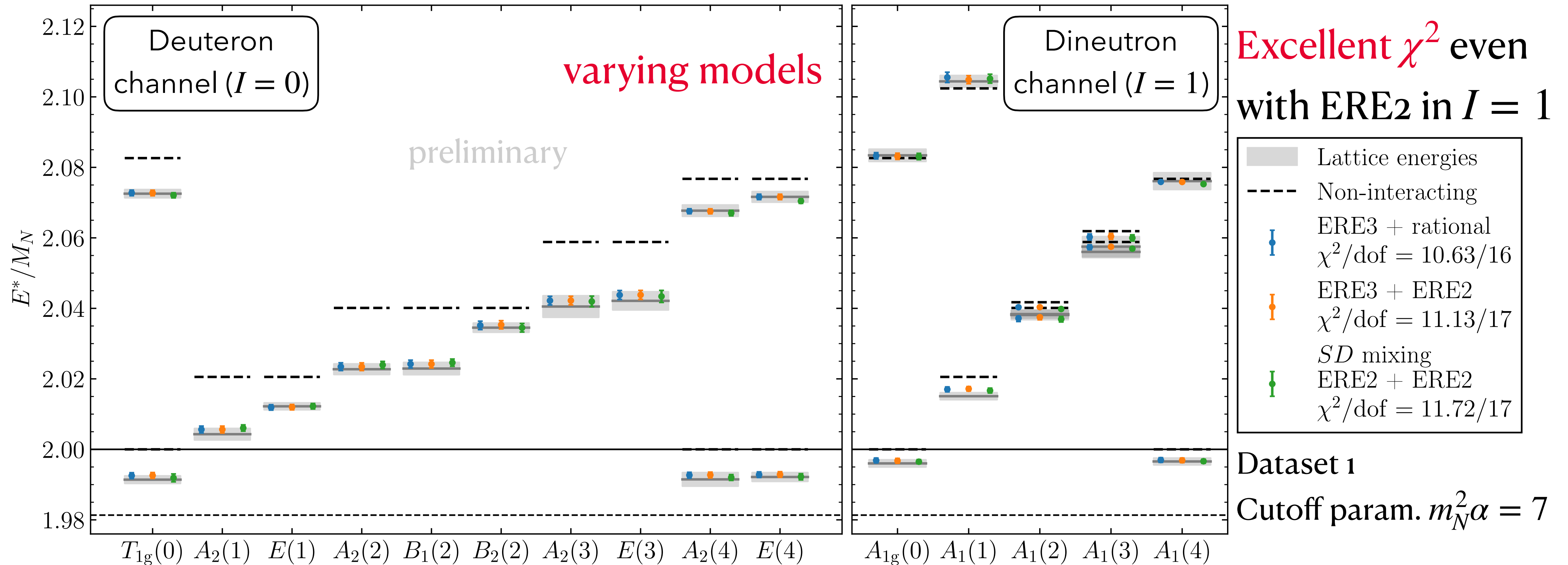
Finite-volume spectrum and fit stability (relativistic formalism)

Combined fit to both isospins with OPE coupling $g_{NN\pi}$ as common fit parameter



NN scattering on D251 ($M_\pi = 280 \text{ MeV}$)

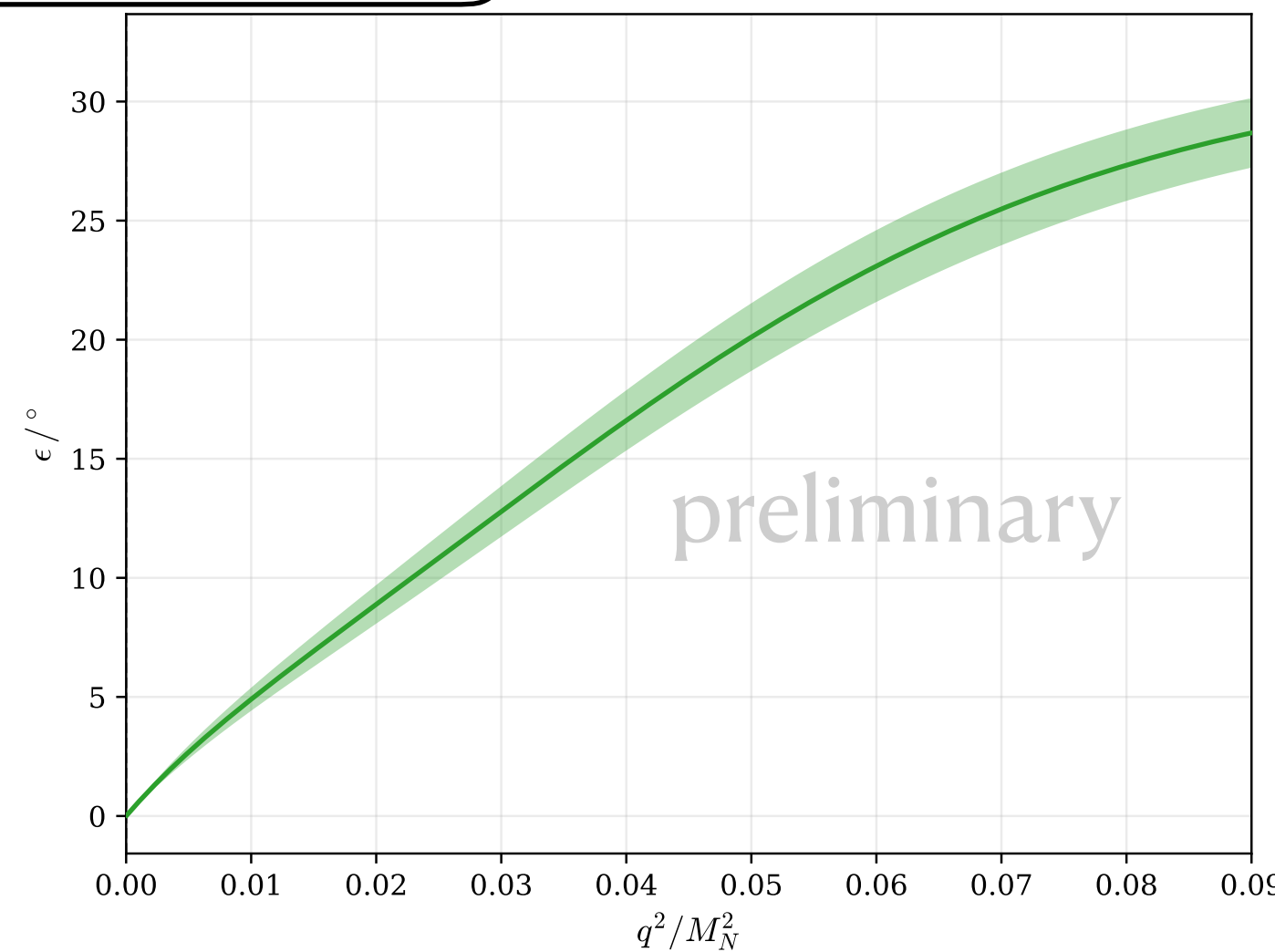
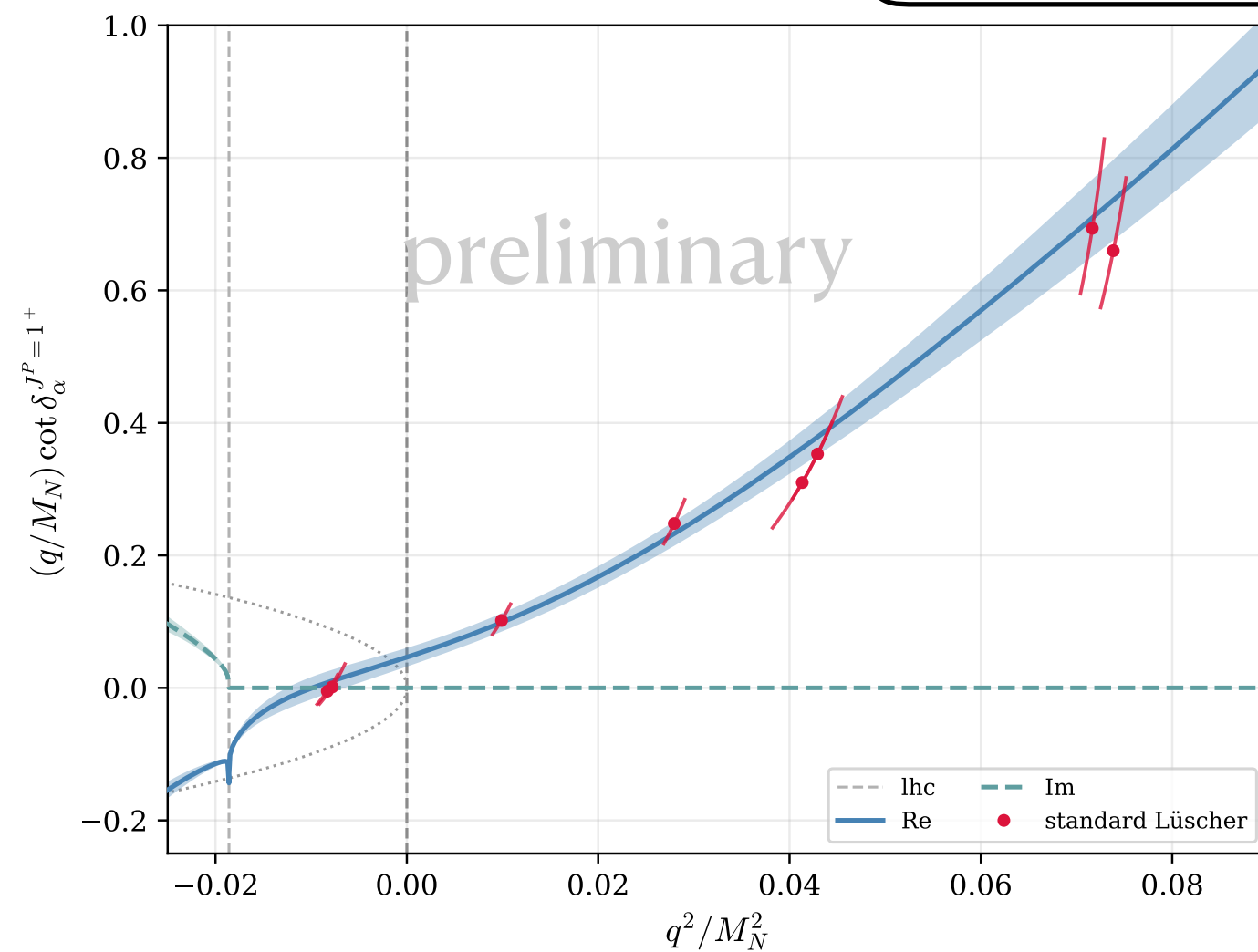
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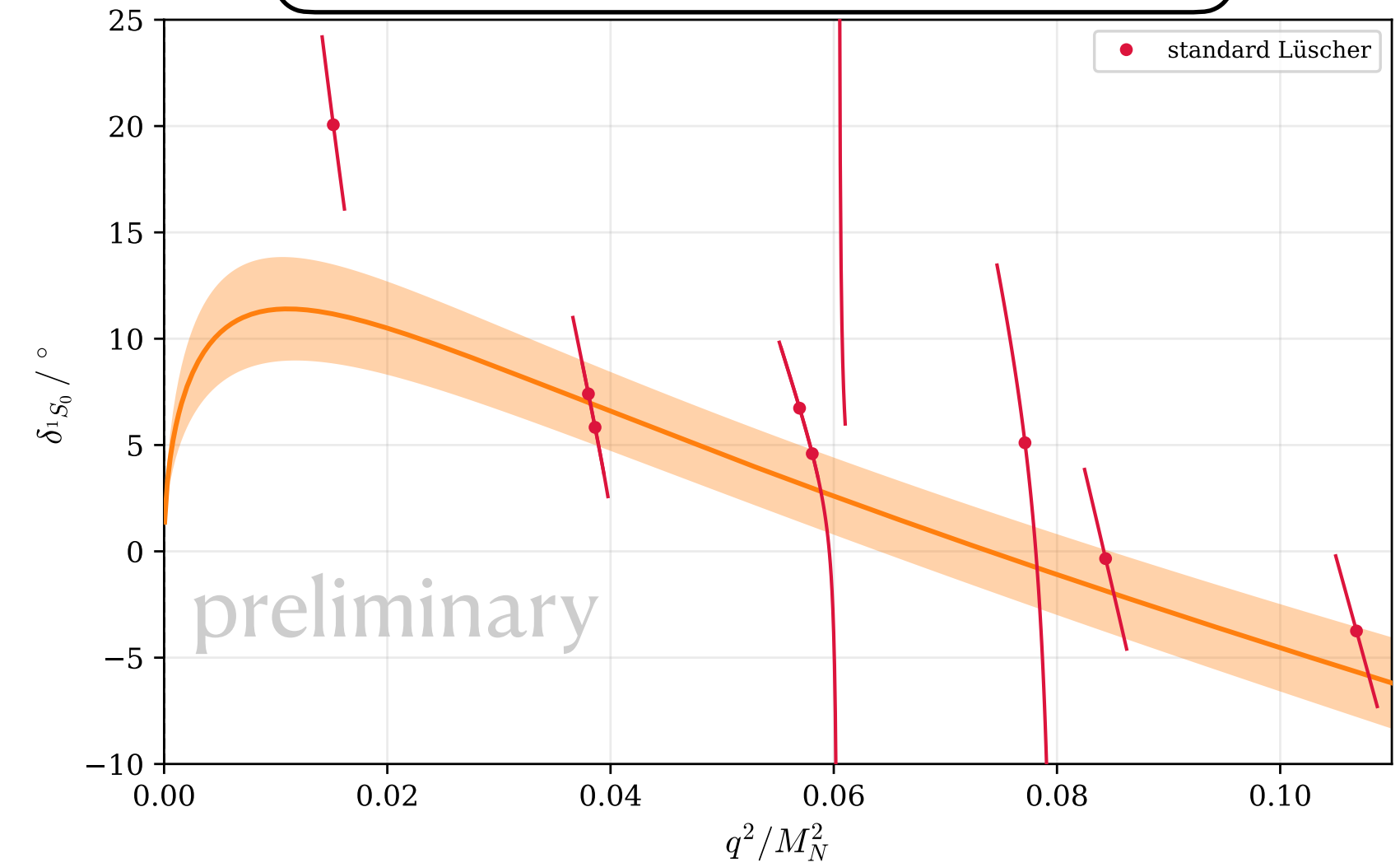
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Fit results for the phase shifts (relativistic formalism)

Deuteron channel ($I = 0$)



Dineutron channel ($I = 1$)



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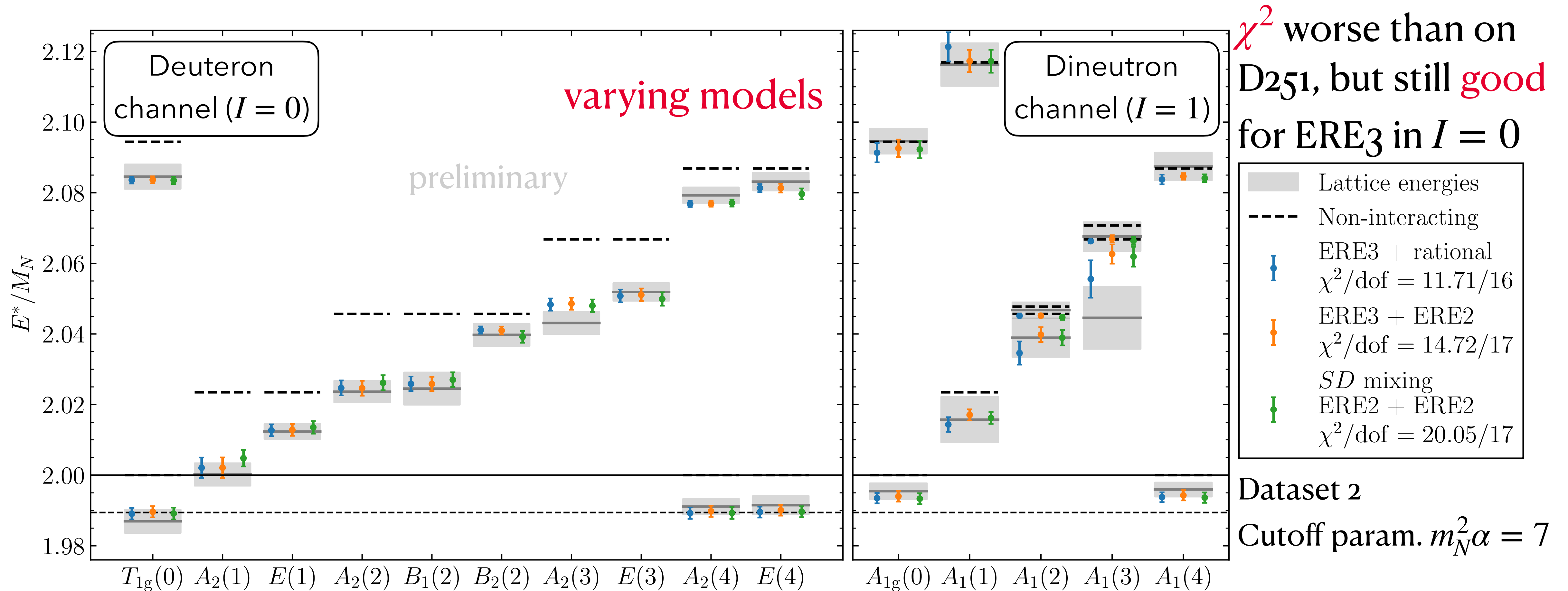
Dataset 1

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Cutoff param. $m_N^2 \alpha = 7$

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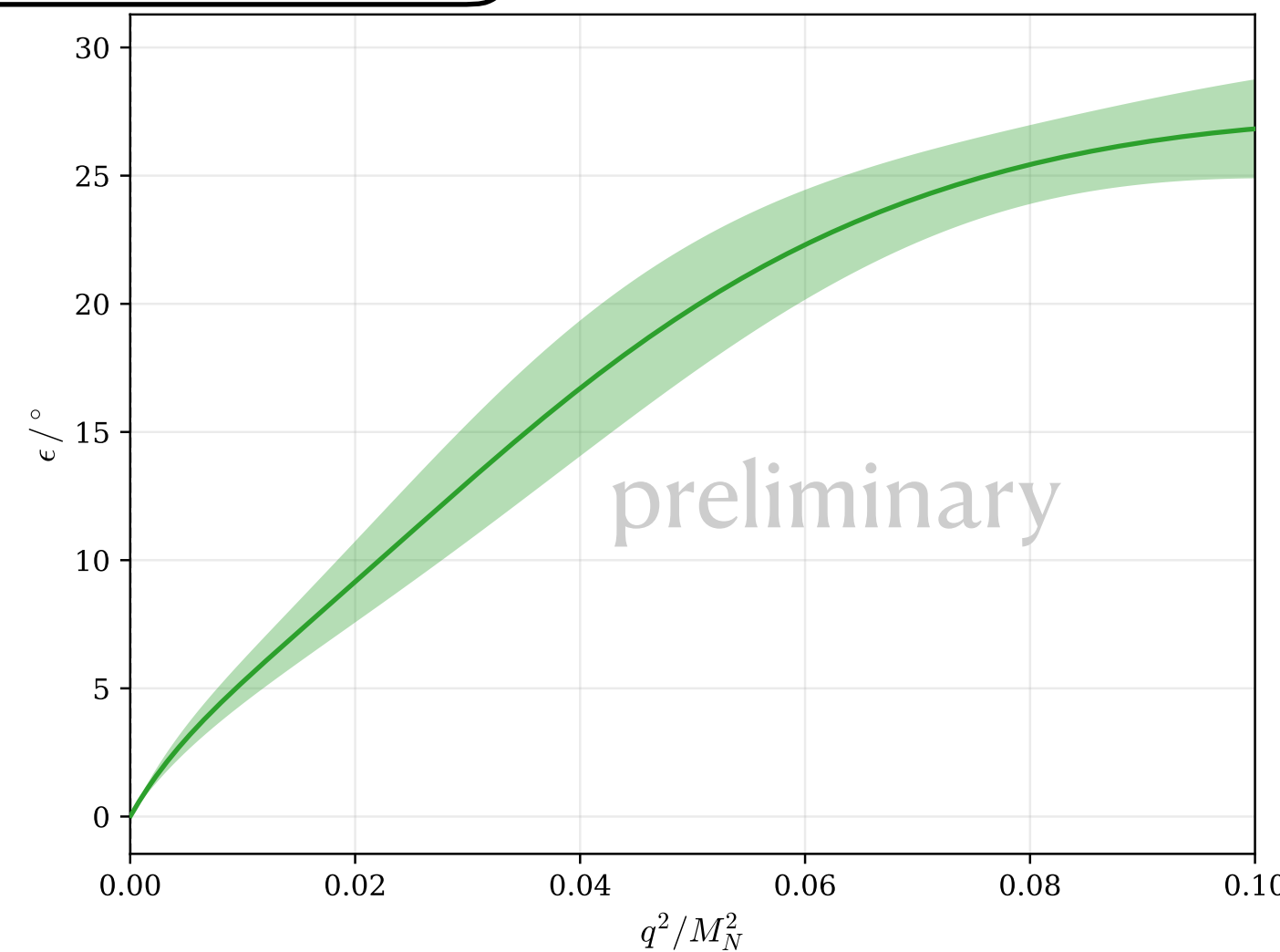
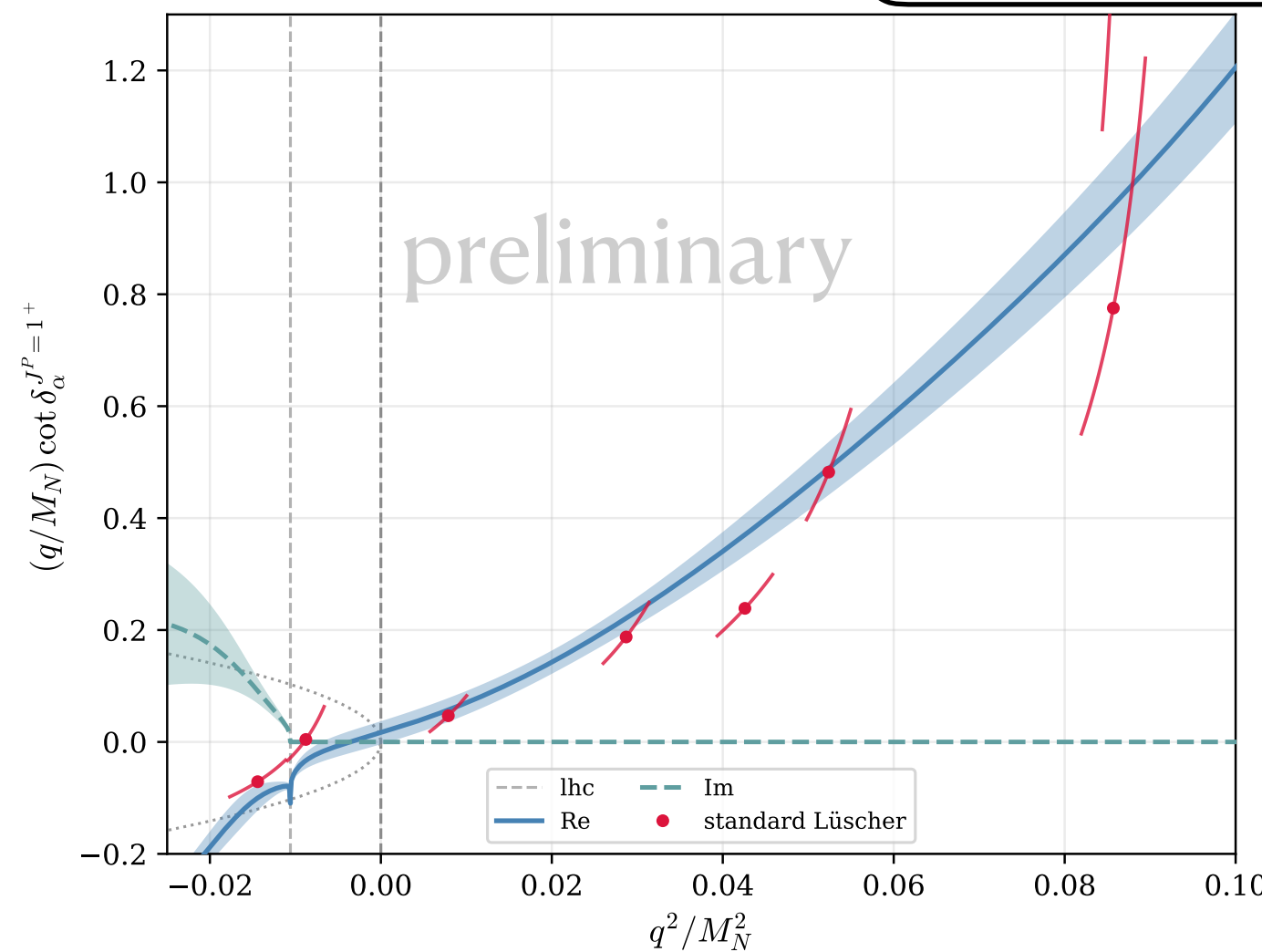
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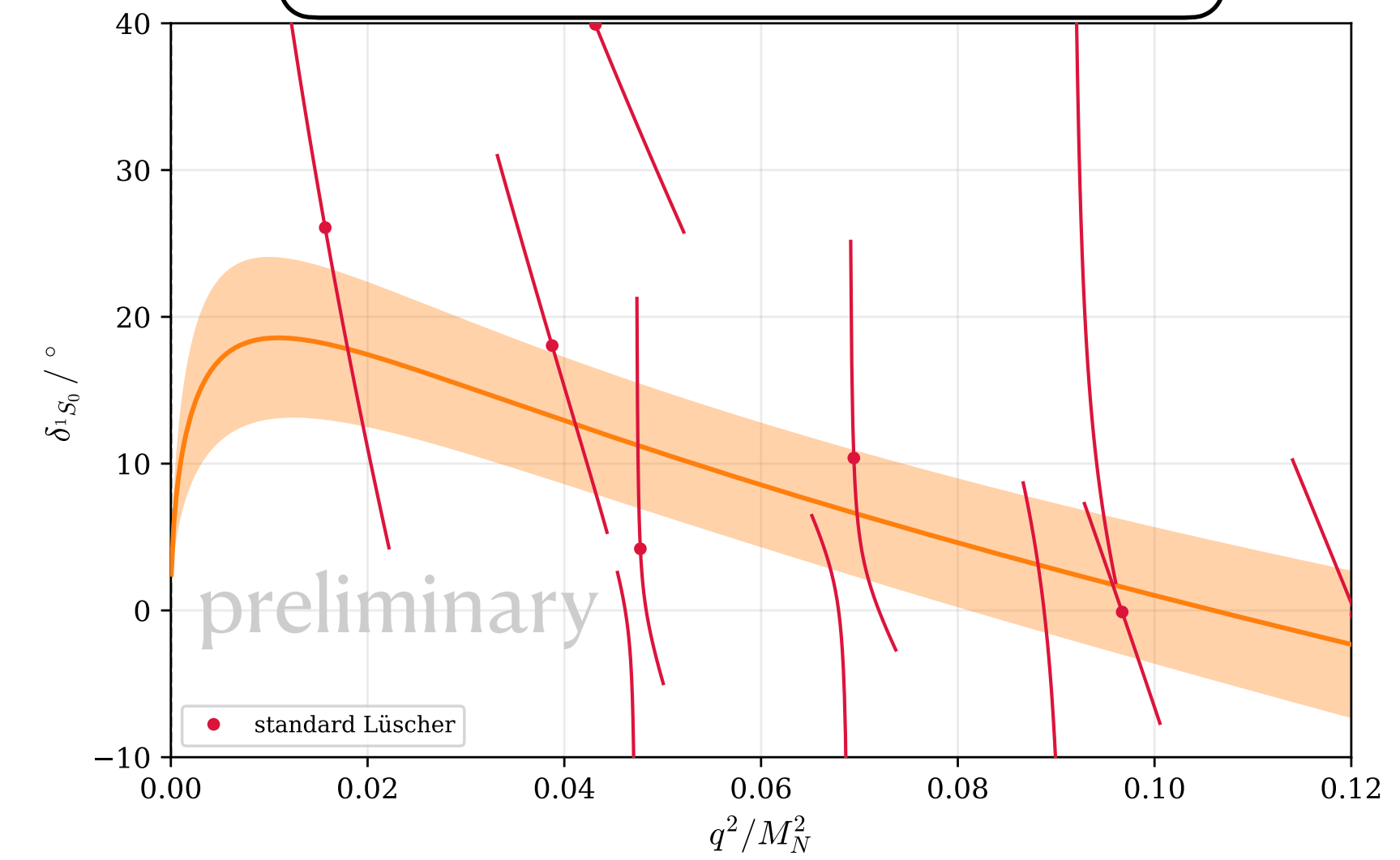
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NN scattering: left-hand cuts

Quantisation condition in the plane-wave basis Meng and Epelbaum 2021 [JHEP 2021, 051]
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Jorge Baeza-Ballesteros, poster session A

- Starting point: non-relativistic **Lippmann-Schwinger** equation (LSE),

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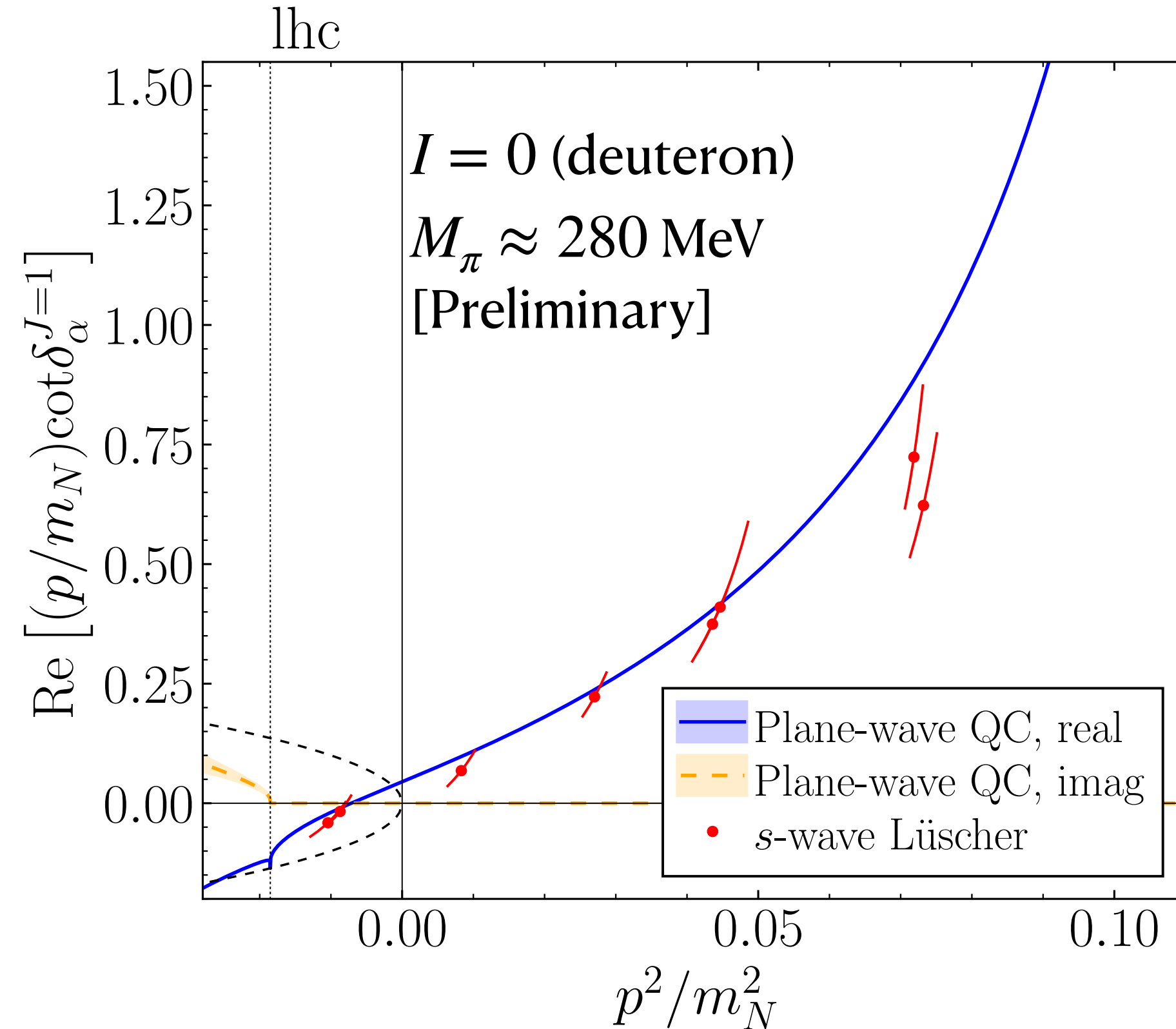
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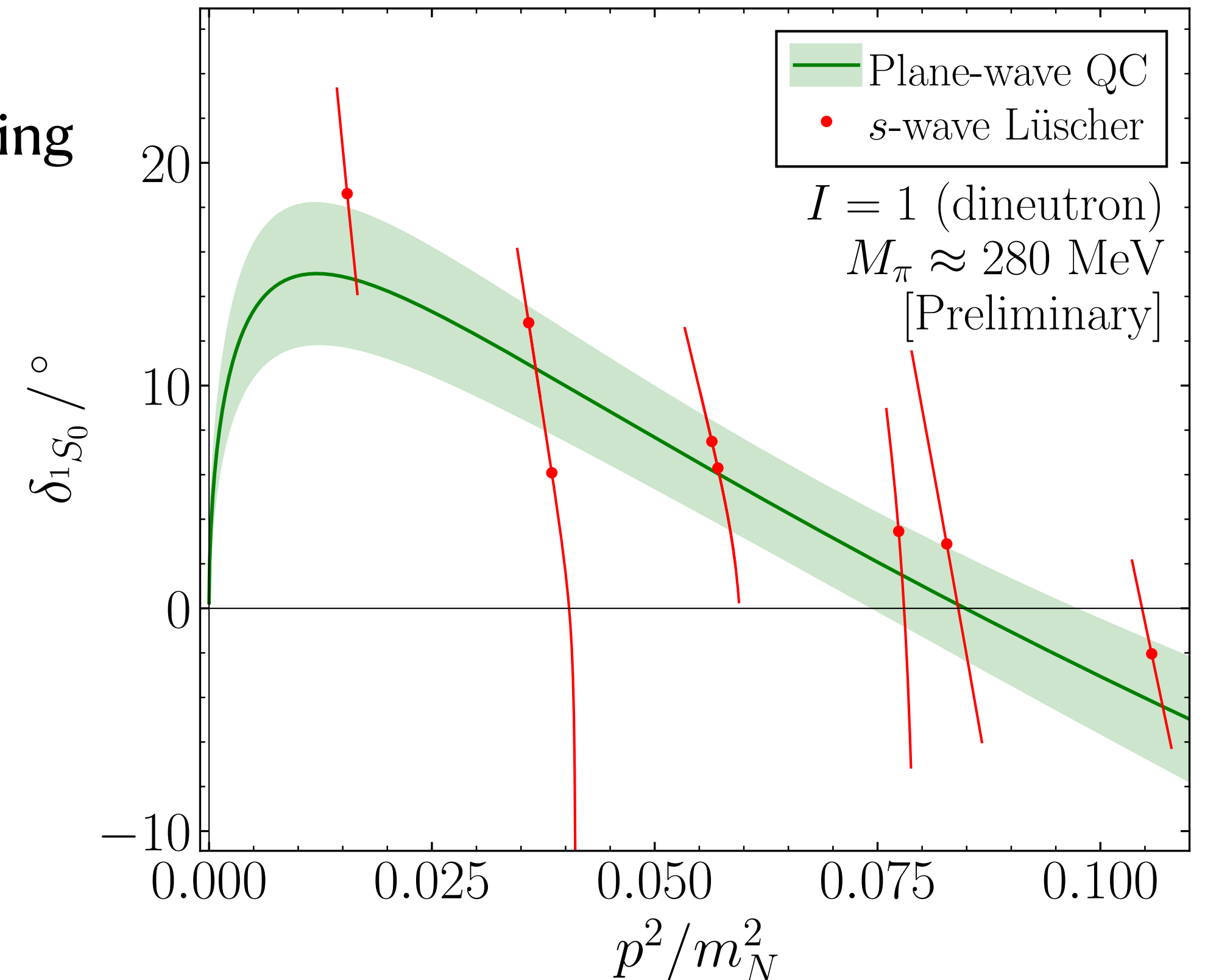
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NN scattering on D251 ($M_\pi = 280$ MeV)

Fit results for the phase shifts (plane-wave basis)



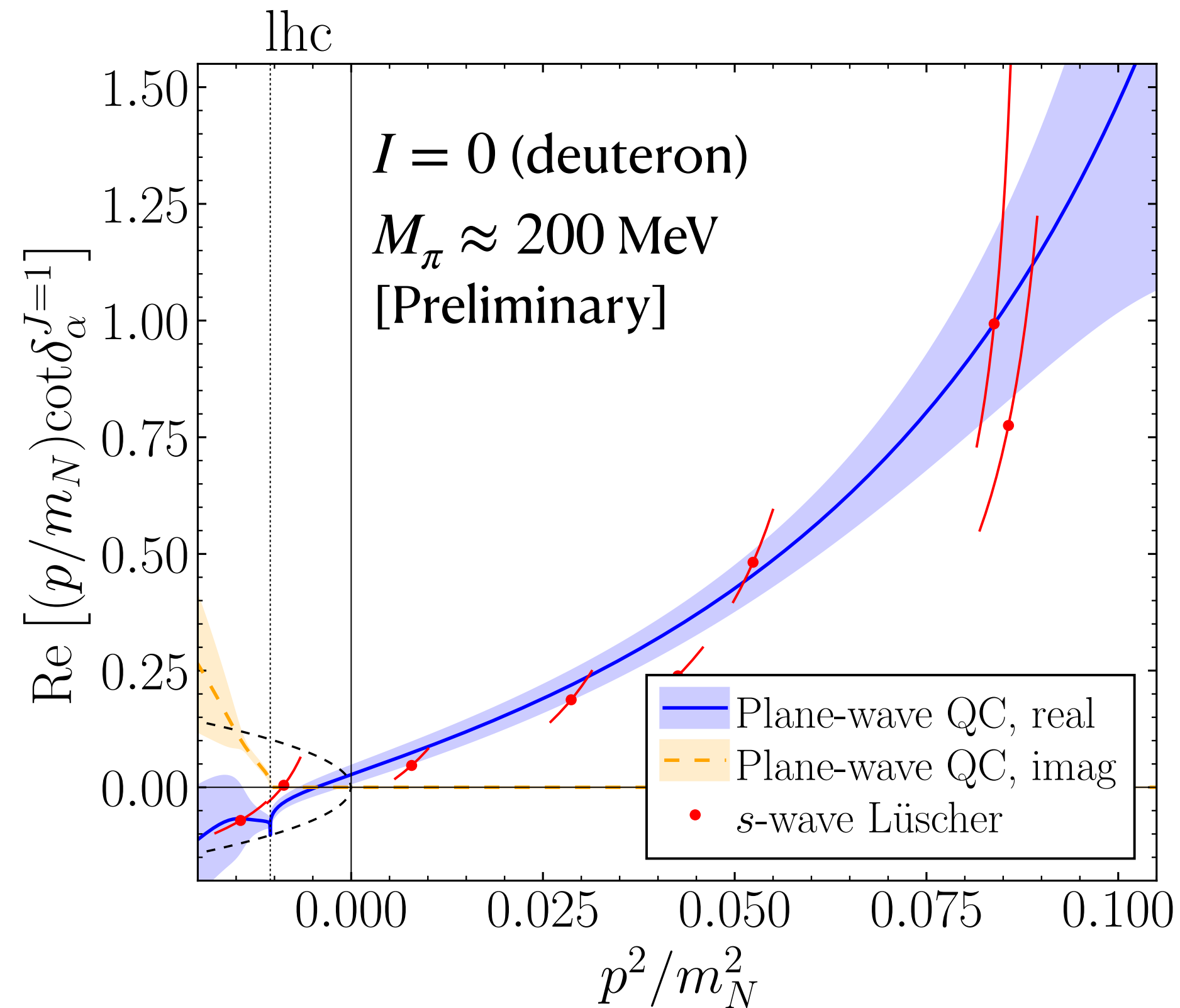
Dataset 2
NLO, no SD mixing
5 fit parameters
 $\Lambda = 700$ MeV



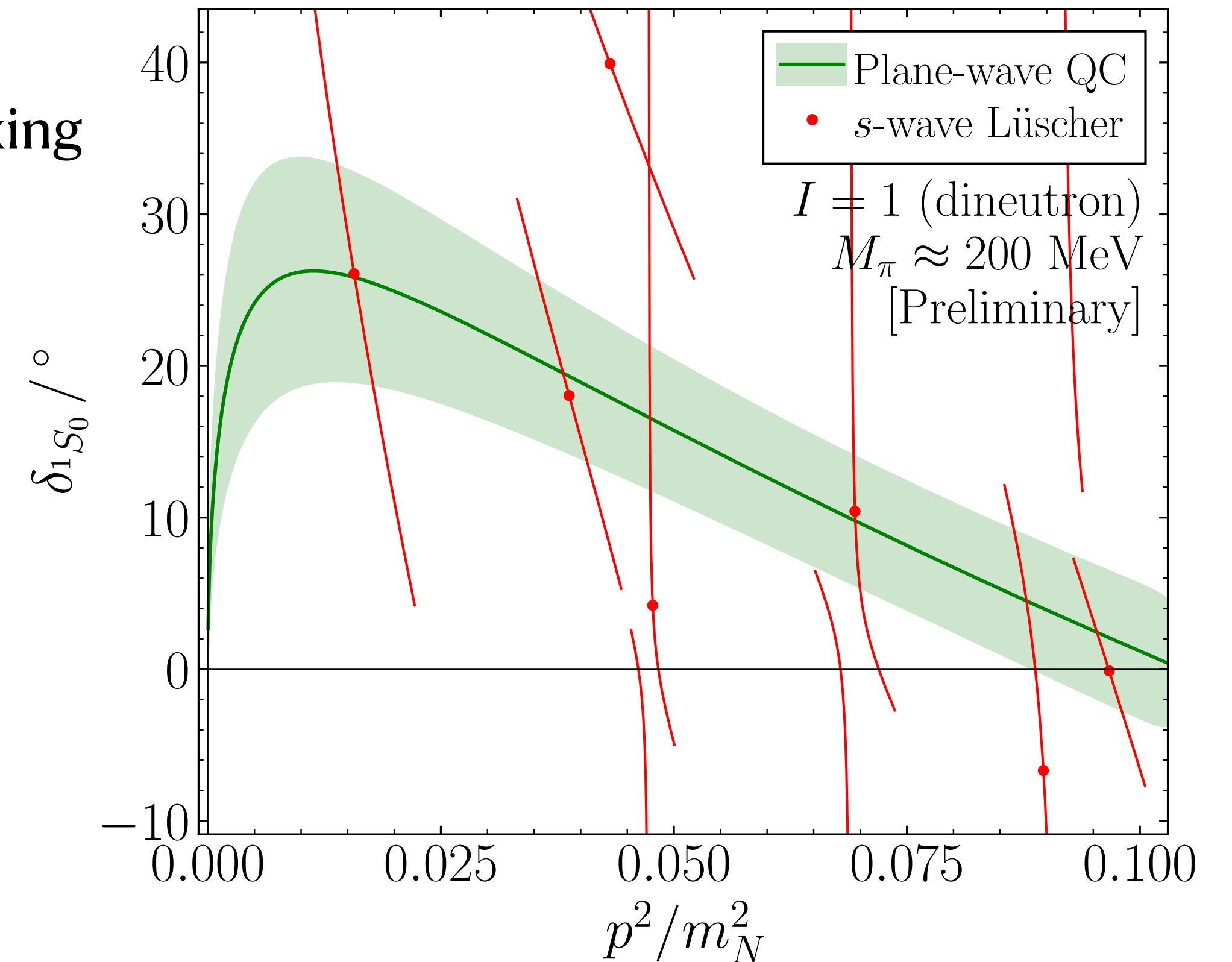
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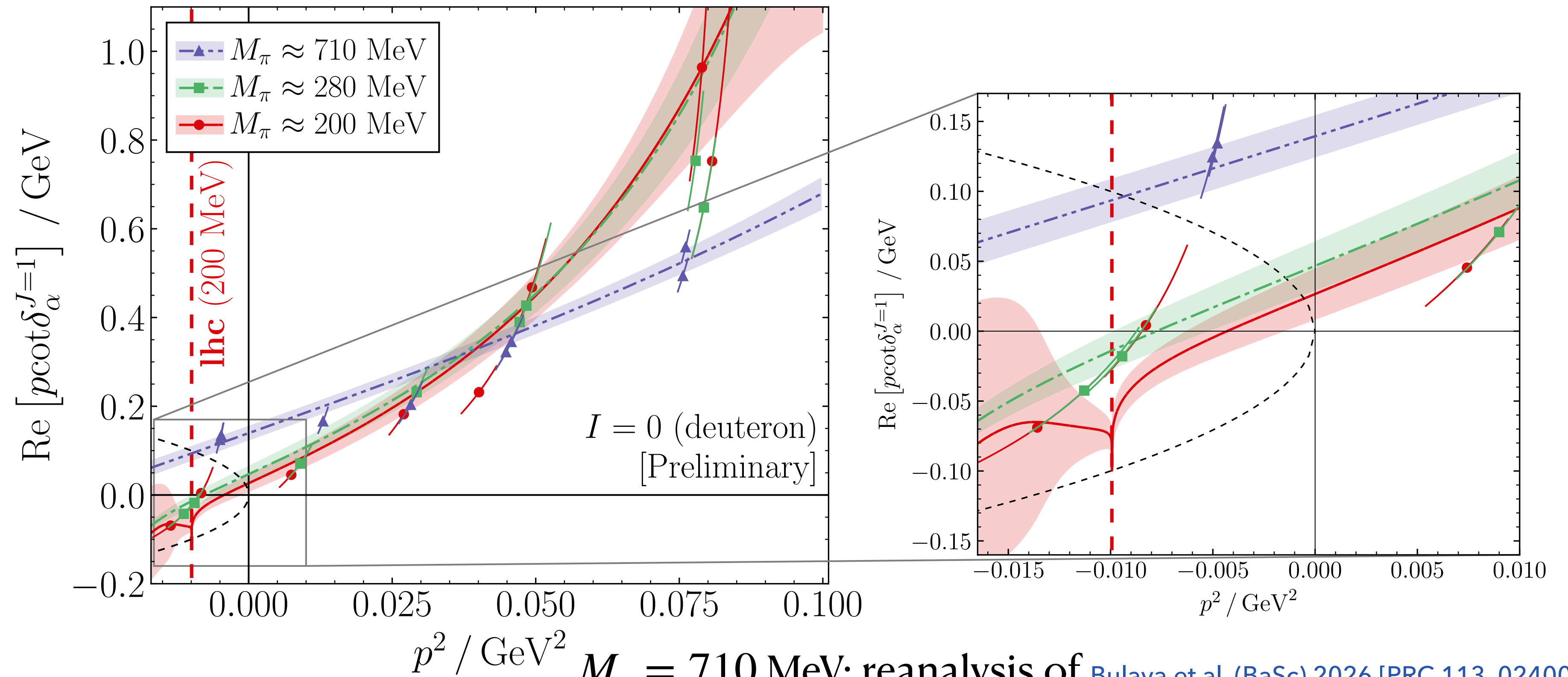


Lowest few energy levels in $I = 0$ (deuteron) severely affected by lhc

NN scattering: plane-wave basis

Comparison between the three ensembles for the deuteron

- Significant **lhc** effects around threshold for $M_\pi = 200$ MeV



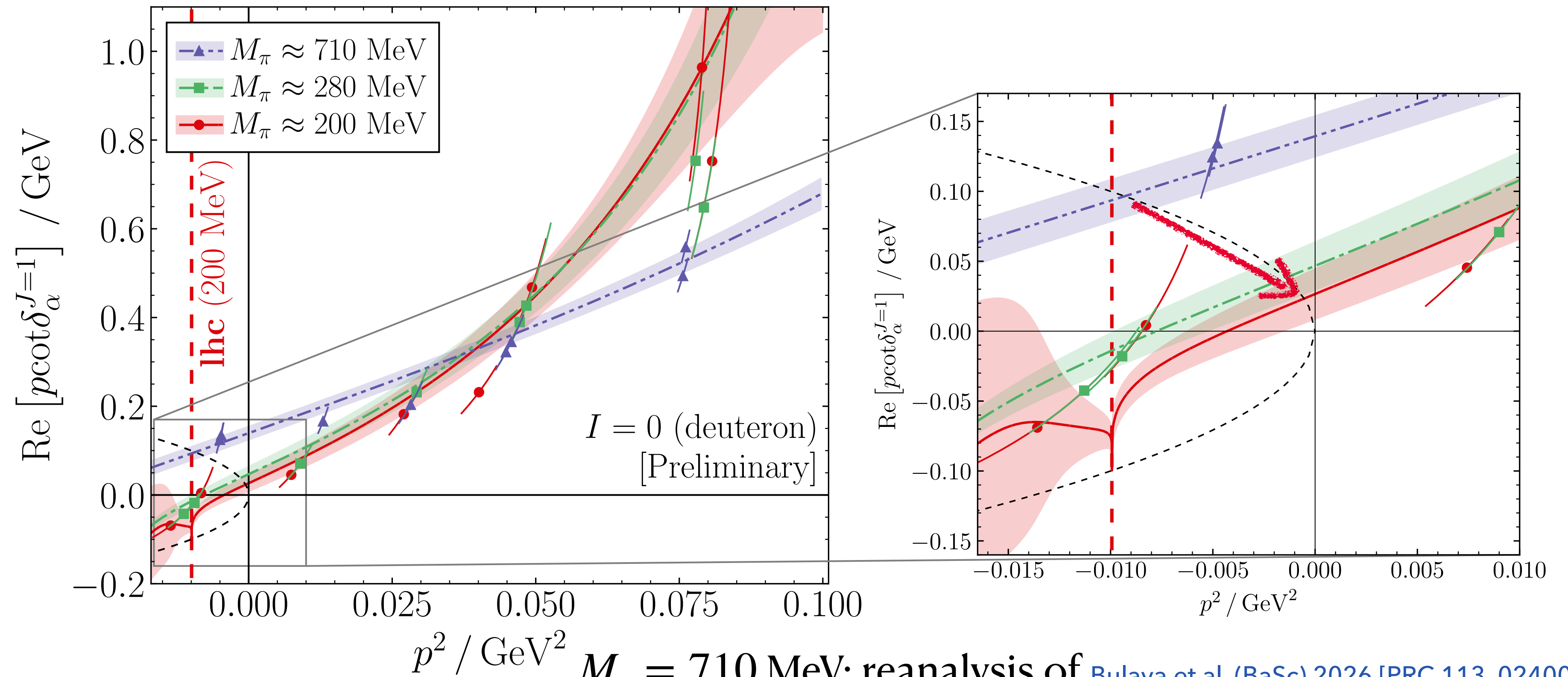
$M_\pi = 710$ MeV: reanalysis of [Bulava et al. \(BaSc\) 2026 \[PRC 113, 024002\]](#)

$M_\pi = 280$ MeV, 200 MeV: this work (dataset 2)

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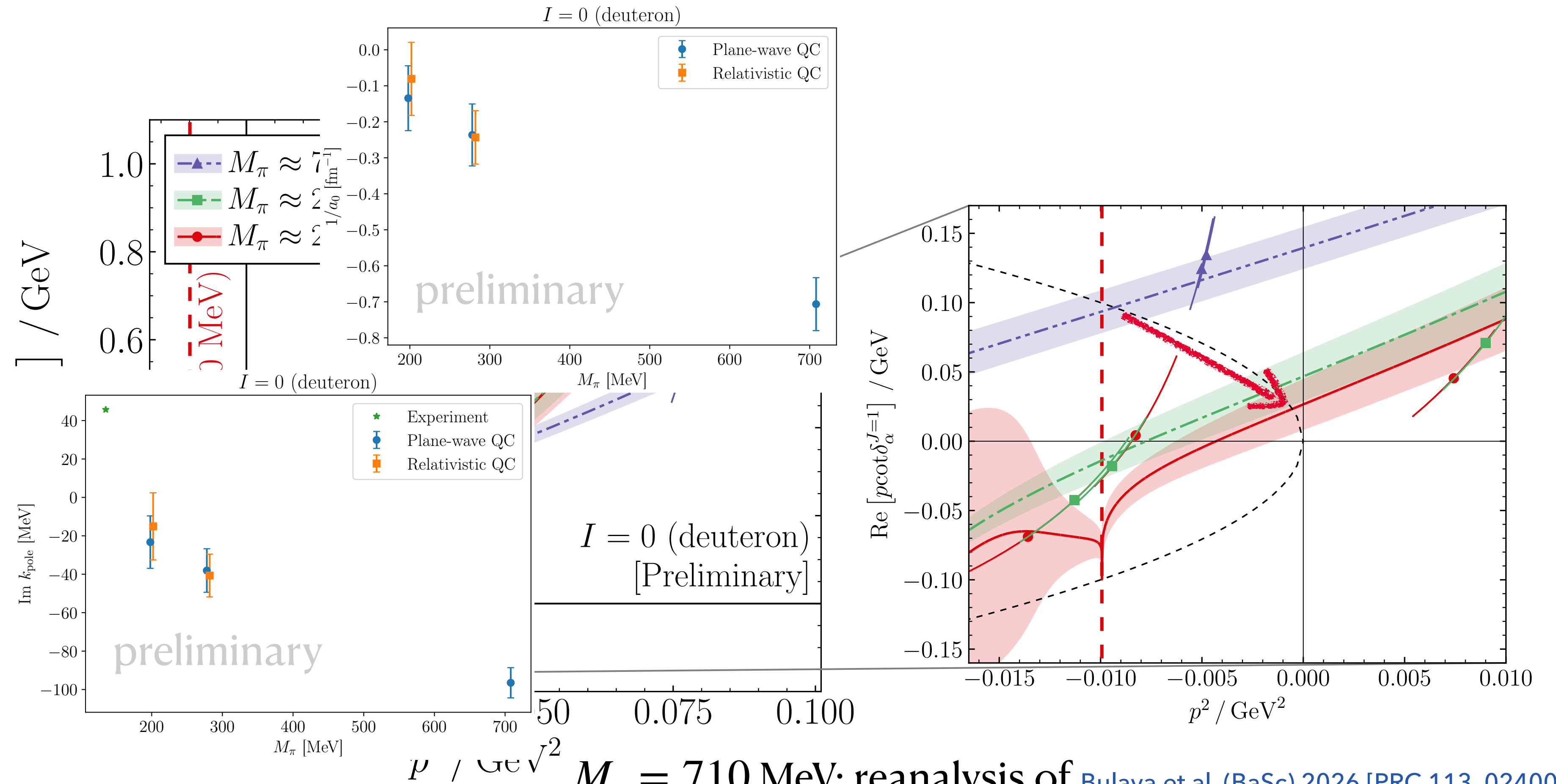
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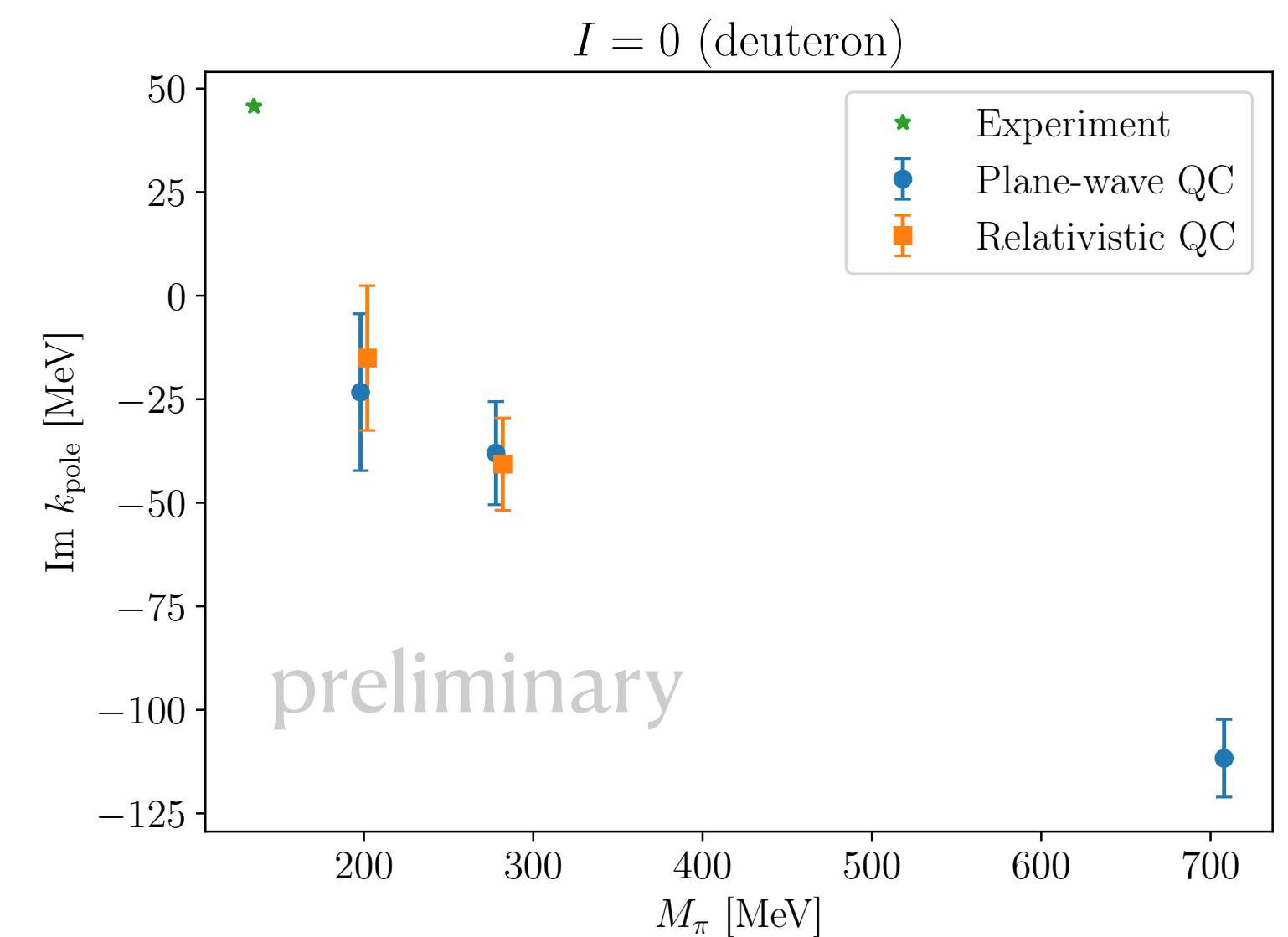


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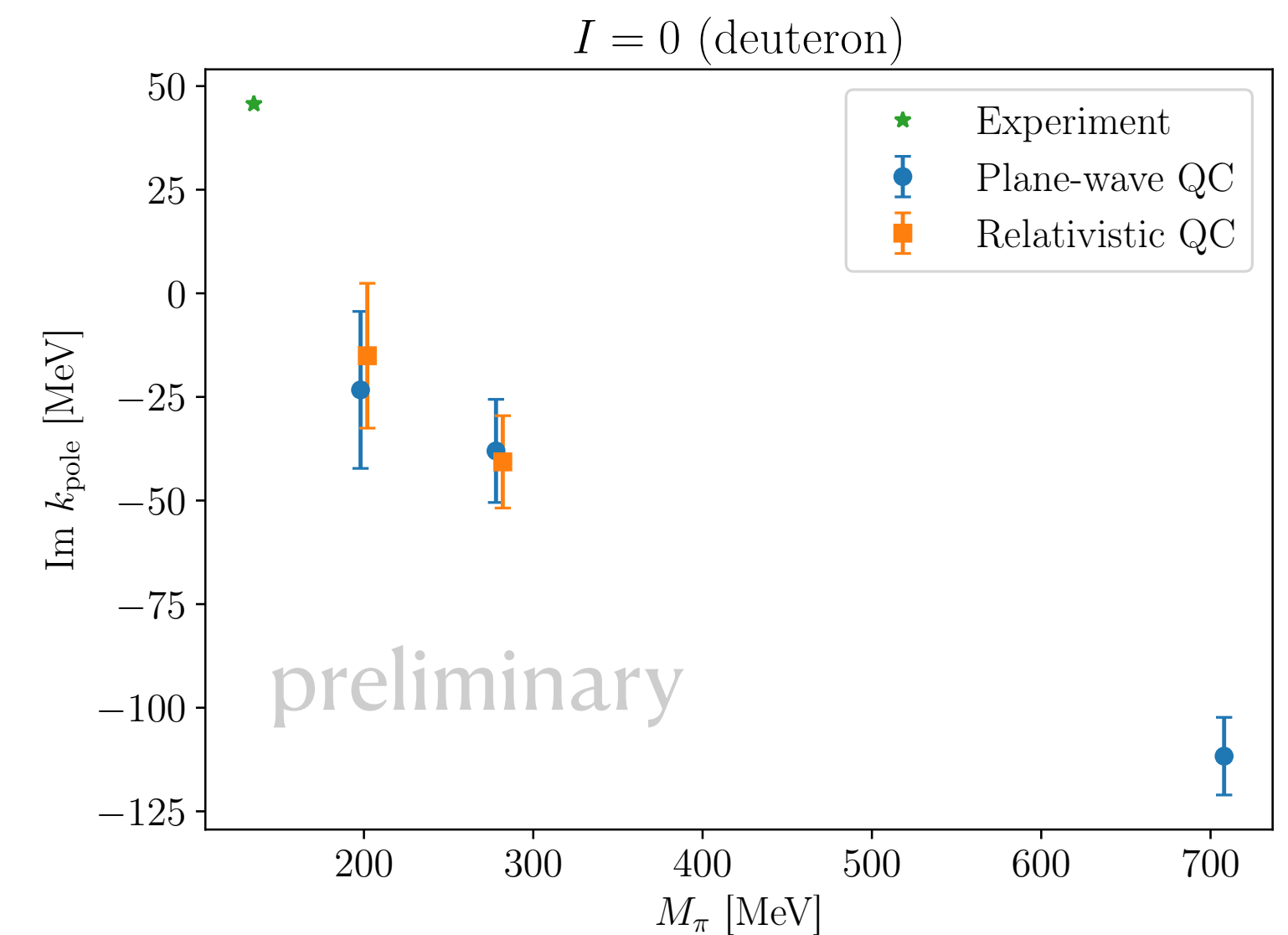
Summary and outlook

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 - Relativistic formalism
 - Plane-wave basis QC



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 - Plane-wave basis QC
- Future directions
 - **Even smaller pion masses** $\rightarrow$ find bound deuteron
 - **Continuum limit**

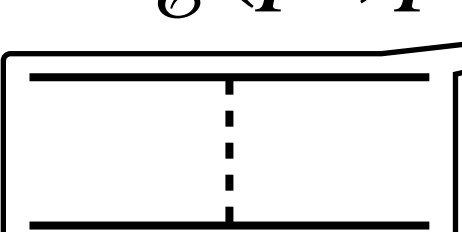


Backup

NN scattering: relativistic formalism

Obtaining the physical scattering amplitude Baião Raposo et al. 2025 [JHEP 2025, 186]

- Partial-wave projected infinite-volume **scattering amplitude** can be written as
 $\mathcal{M} = \mathcal{M}_{\mathcal{E}} + (1 + i\mathcal{L})(\mathcal{M}_0^{-1} + \mathcal{C})^{-1}(1 + i\mathcal{R})$, where $\mathcal{M}_0$ is amplitude corresponding to $\mathcal{K}_0$
- $\mathcal{M}_{\mathcal{E}}$ satisfies the **integral equation** Two-particle pole in IV

$$\mathcal{M}_{\mathcal{E}}(p', p) = i\mathcal{E}(p', p) + \int_0^\infty \frac{k^2 dk}{2\pi^2} i\mathcal{E}(p', k) i\Delta_{2,\infty}(k) i\mathcal{M}_{\mathcal{E}}(k, p) \rightarrow \text{solve as matrix equation}$$
-  $i\mathcal{L}(p) = \int_0^\infty \frac{k^2 dk}{2\pi^2} i\mathcal{M}_{\mathcal{E}}(p, k) i\Delta_{2,\infty}(k) \left(\frac{k}{q}\right)^\ell$, $i\mathcal{R}(p) = \int_0^\infty \frac{k^2 dk}{2\pi^2} \left(\frac{k}{q}\right)^{\ell'} i\Delta_{2,\infty}(k) i\mathcal{M}_{\mathcal{E}}(k, p)$,
 and $i\mathcal{C} = \int_0^\infty \frac{k'^2 dk'}{2\pi^2} \int_0^\infty \frac{k^2 dk}{2\pi^2} \left(\frac{k'}{q}\right)^{\ell'} i\Delta_{2,\infty}(k') i\mathcal{M}_{\mathcal{E}}(k', k) i\Delta_{2,\infty}(k) \left(\frac{k}{q}\right)^\ell$

NN scattering: relativistic formalism

Spectrum fit: Models for $I = 0$

- Without S , D -wave mixing: ERE3 for $\mathcal{K}_0$, $q \cot(\tilde{\delta}(E)) = -\frac{1}{A} + \frac{1}{2}Rq^2 + C^{(2)}q^4$
- With S , D -wave mixing: Blatt and Biedenharn 1952 [PR 86, 399], 1954 [PR 93, 1387] Blatt-Biedenharn parameterisation for modified K-matrix,

$$\mathcal{K}_0^{-1}(E) = Q\mathcal{R}^\dagger(\theta q^2) \begin{pmatrix} q \cot(\tilde{\delta}_\alpha(E)) & 0 \\ 0 & q \cot(\tilde{\delta}_\beta(E)) \end{pmatrix} \mathcal{R}(\theta q^2)Q$$
, where $Q_{\ell'\ell} = \delta_{\ell'\ell}q^\ell$, $\mathcal{R}$ is a rotation matrix, and we use $q \cot(\tilde{\delta}_\alpha(E)) = -\frac{1}{A_\alpha} + \frac{1}{2}R_\alpha q^2$ and $q^5 \cot(\tilde{\delta}_\beta(E)) \rightarrow \infty$
- A_α , R_α , $C^{(2)}$, and θ are parameters of $\mathcal{K}_0$ and thus inherently **scheme dependent**
- Additional fit parameter $g_{NN\pi}^2$ (shared with $I = 1$) is coefficient of $\overline{\text{---}} \text{---} \rightarrow$ **physical**

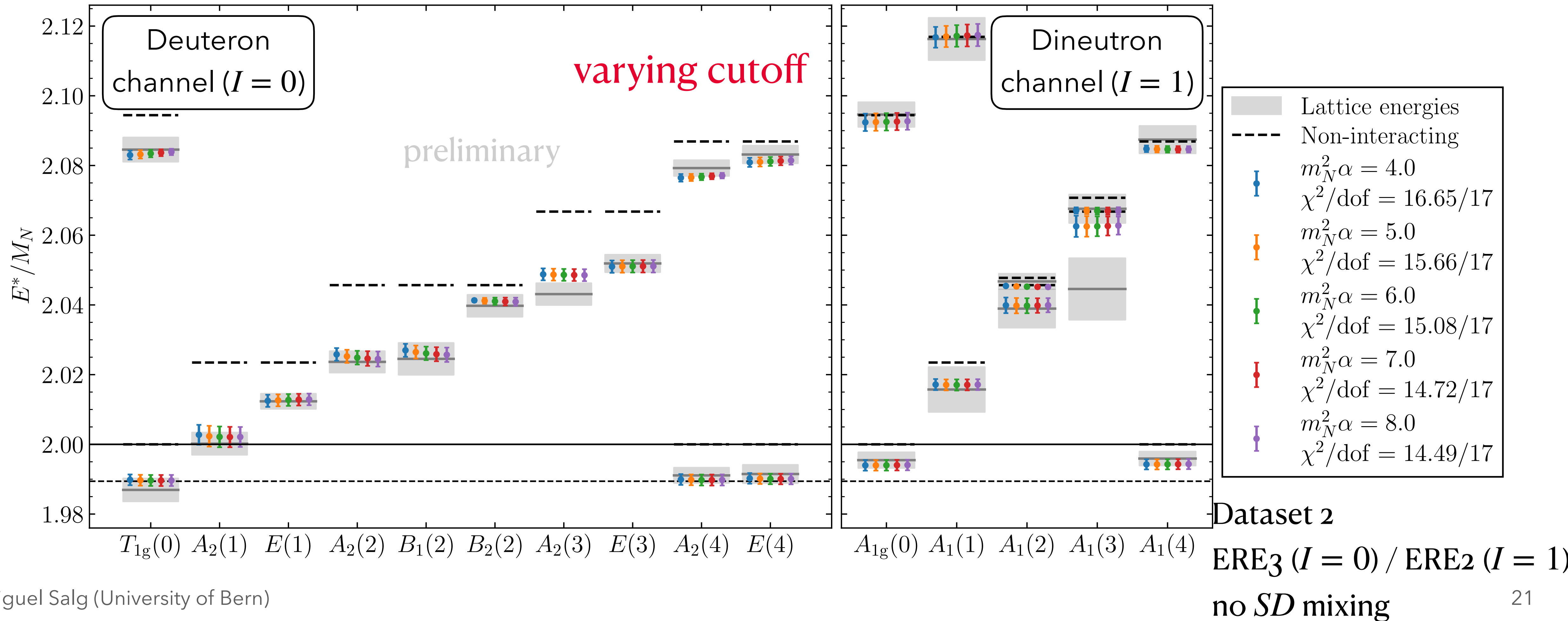
NN scattering: relativistic formalism

Spectrum fit: Models for $I = 1$

- No S, D -wave mixing, only A_1 -like irreps
- Two models for modified K-matrix
 - ERE2, $q \cot(\tilde{\delta}(E)) = -\frac{1}{A} + \frac{1}{2}Rq^2$
 - Rational, $\frac{1}{q} \tan(\tilde{\delta}(E)) = \frac{C_0 + C_1q^2}{1 + C_2q^2} \rightarrow$ allows for both zero and pole in $q \cot \tilde{\delta}$
- A, R , and C_i are parameters of $\mathcal{K}_0$ and thus inherently scheme dependent
- Additional fit parameter $g_{NN\pi}^2$ (shared with $I = 1$) is coefficient of $\overline{\pi\pi} \rightarrow$ physical

NN scattering on J250 ($M_\pi = 200$ MeV)

Finite-volume spectrum and fit stability (relativistic formalism)



NN scattering on D251 ($M_\pi = 280 \text{ MeV}$)

Cutoff dependence of fit parameters (relativistic formalism)

$\mathcal{C}_L$ and $F_{2,L}$
internally use
exponential **cutoff**,
 $\exp\{\alpha[(q^*)^2 - (\mathbf{k}^*)^2]\}$

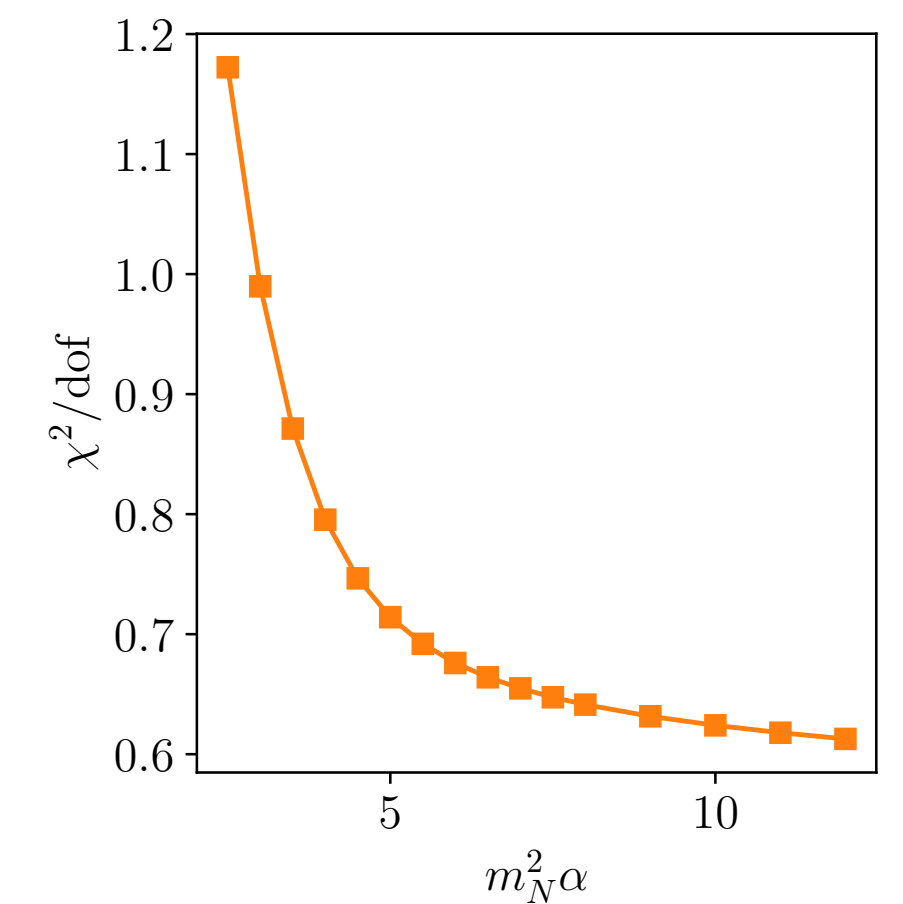
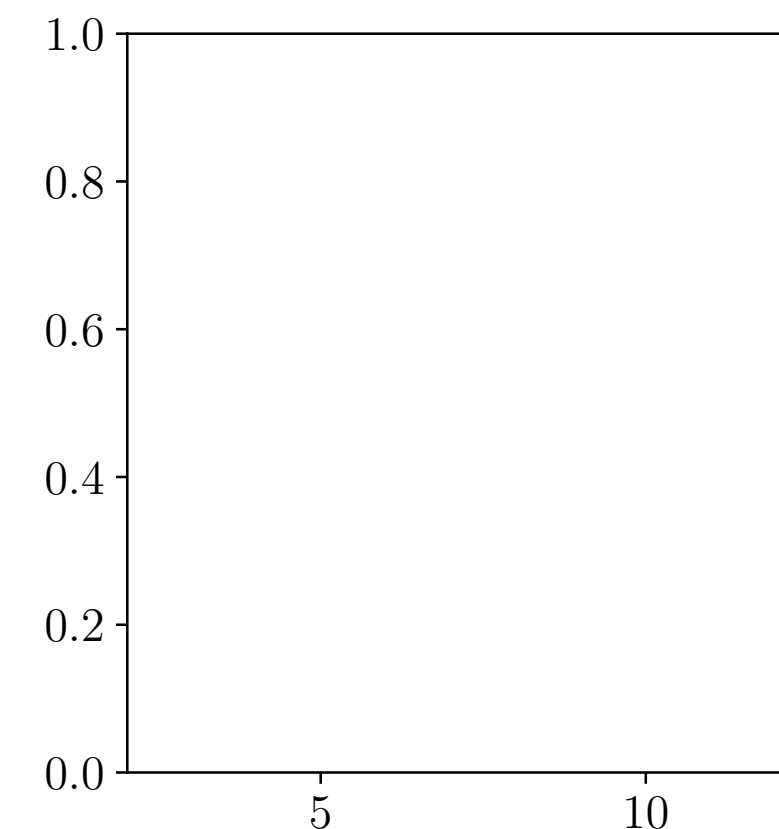
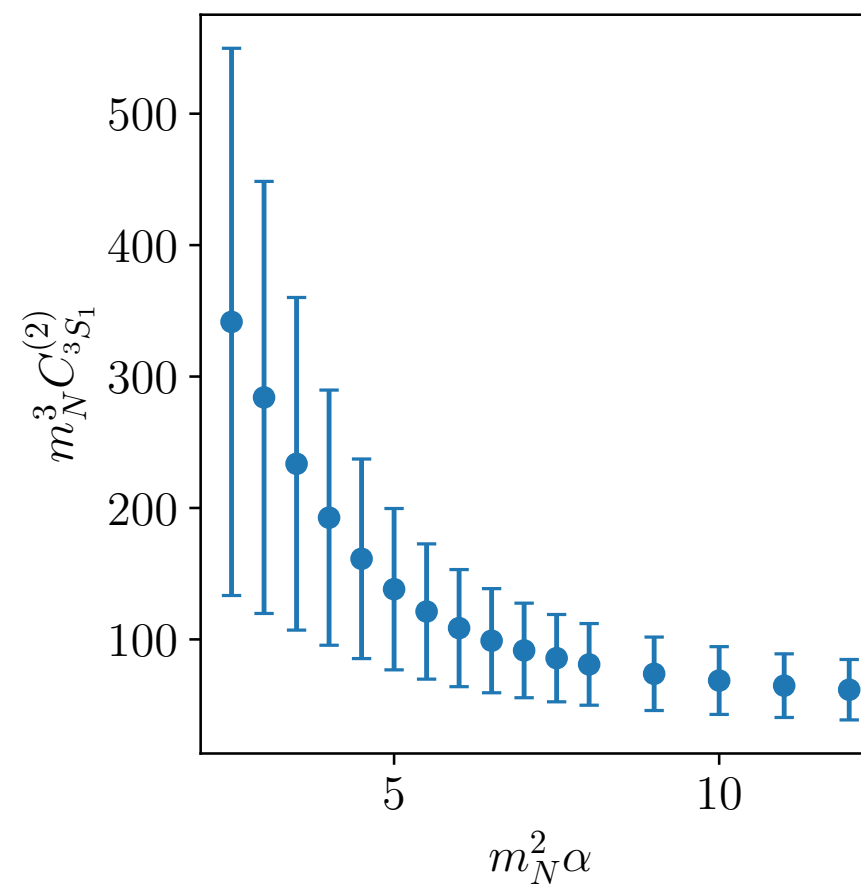
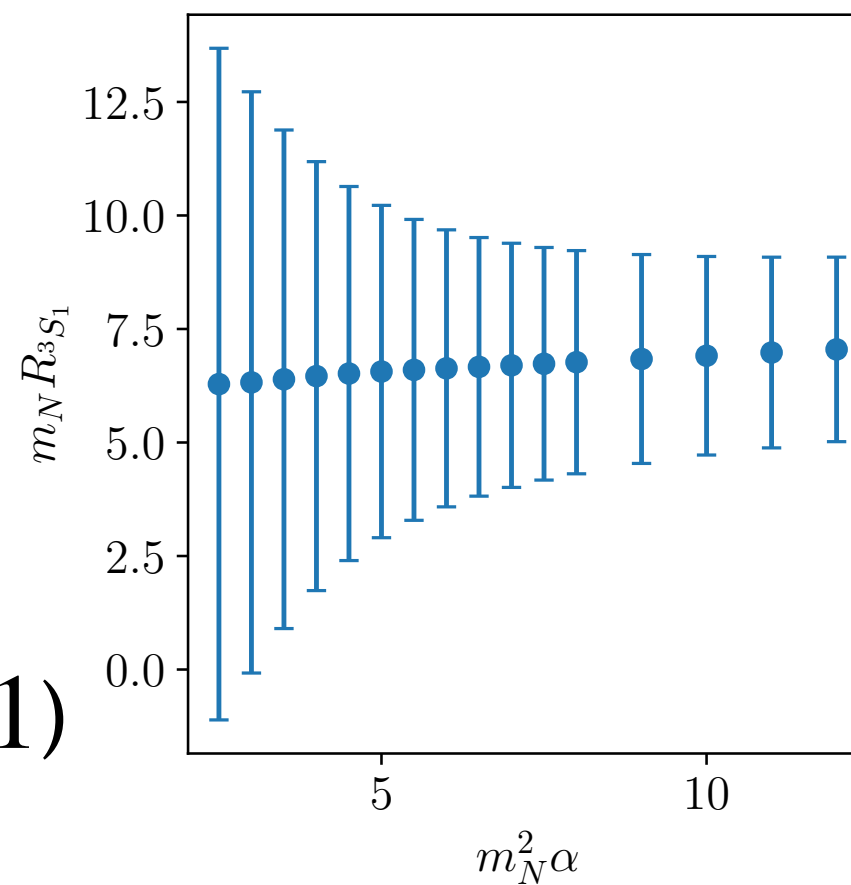
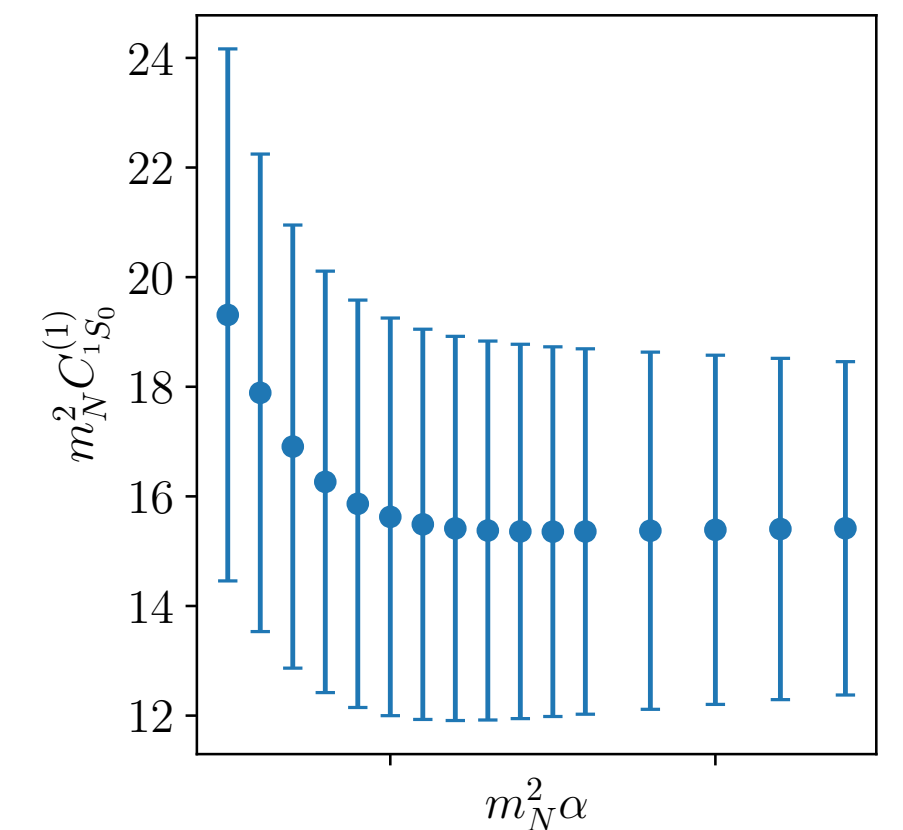
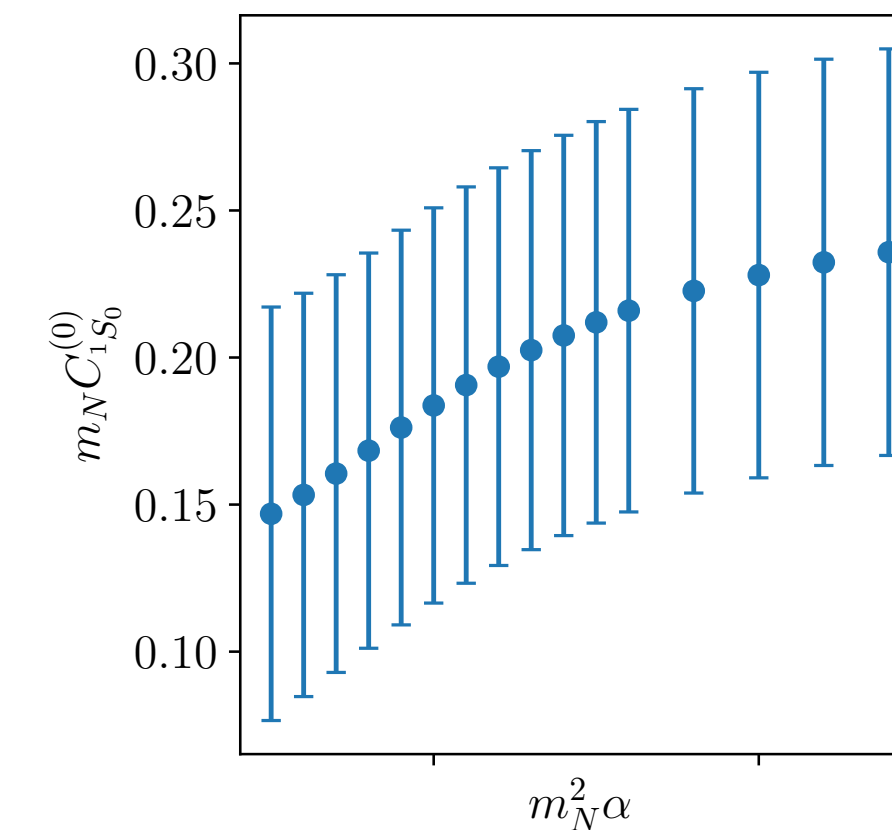
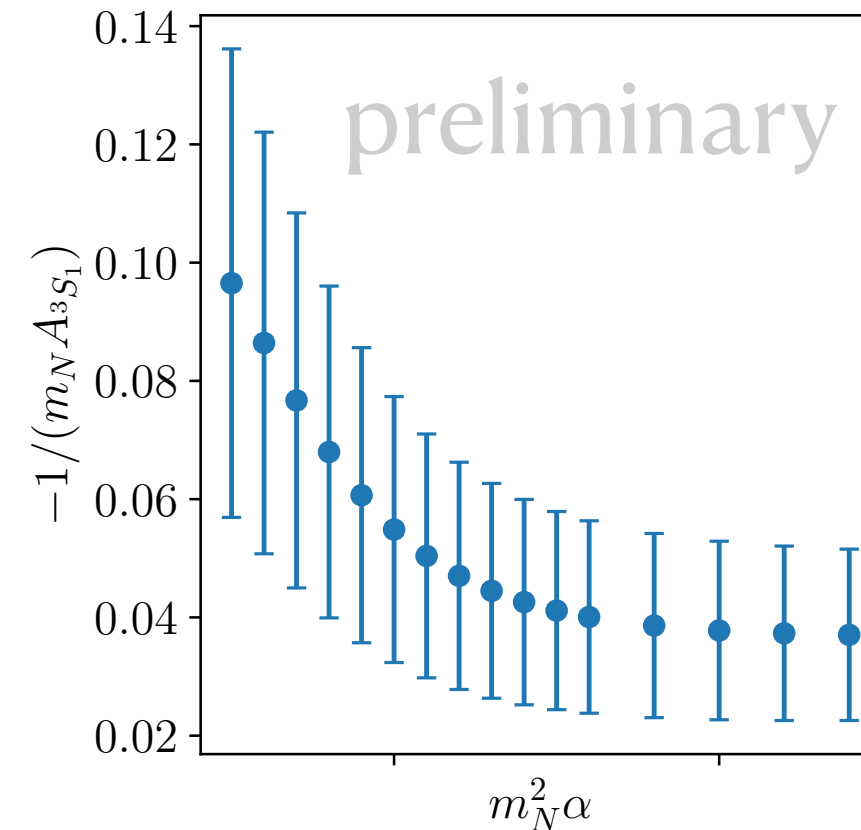
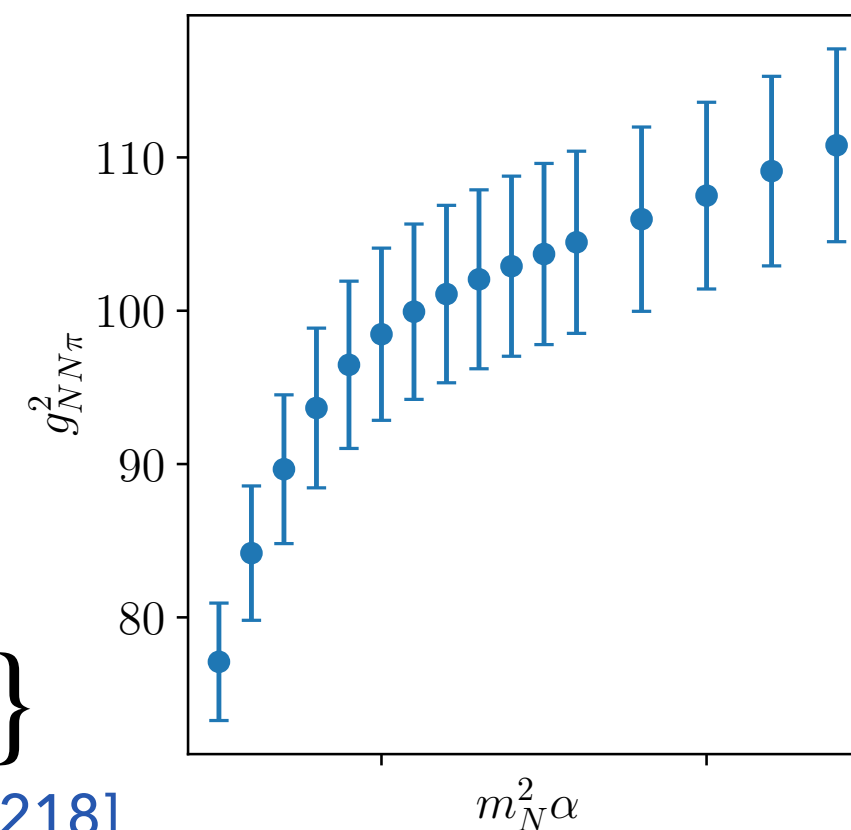
Kim, Sachrajda, and Sharpe 2005 [NPB 727, 218]

—> Larger α means
stricter cutoff

Dataset 1

ERE₃ ($I = 0$) / ERE₂ ($I = 1$)

no SD mixing



NN scattering on J250 ($M_\pi = 200 \text{ MeV}$)

Cutoff dependence of fit parameters (relativistic formalism)

$\mathcal{C}_L$ and $F_{2,L}$
internally use
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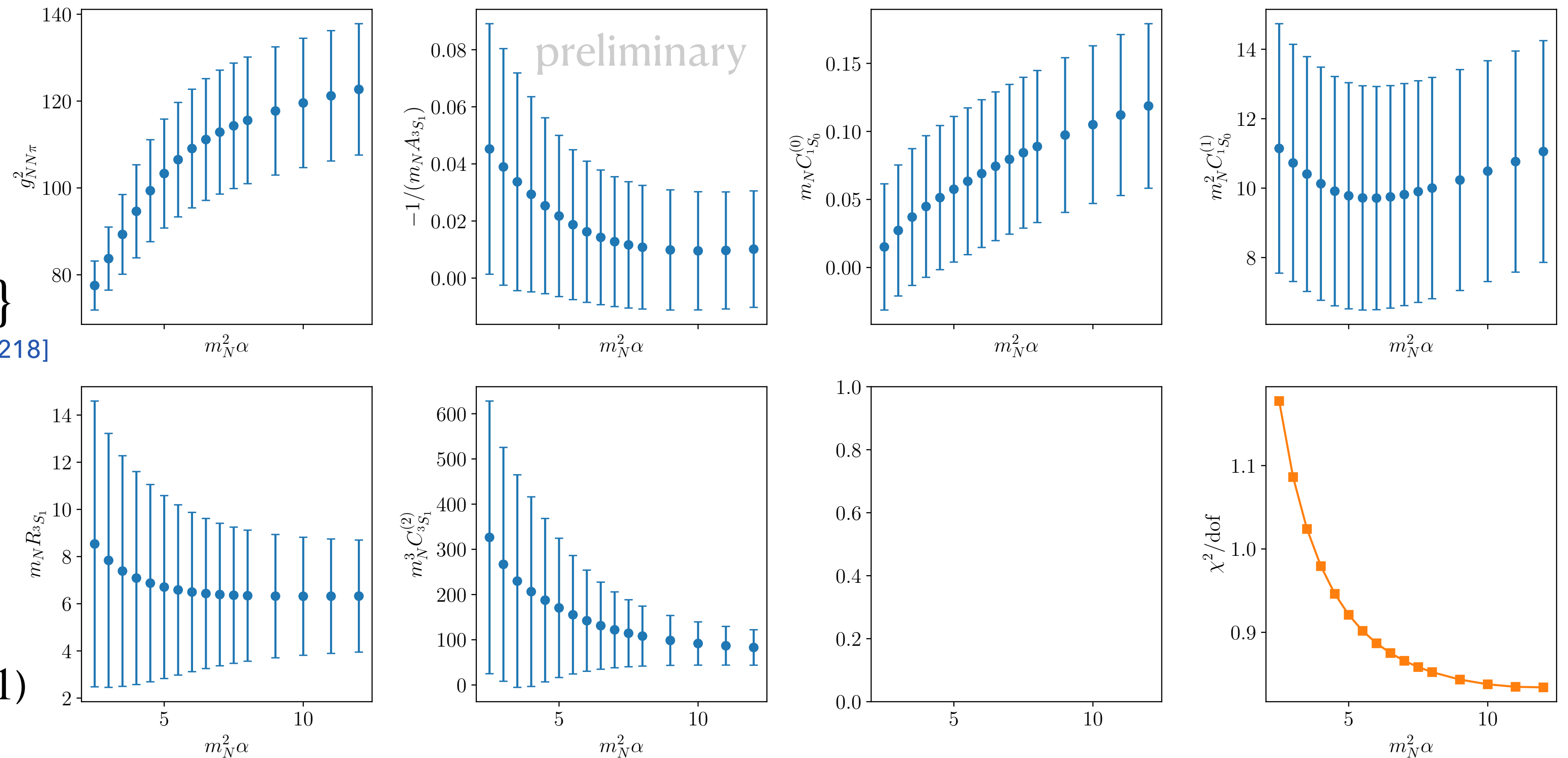
Kim, Sachrajda, and Sharpe 2005 [NPB 727, 218]

—> Larger α means
stricter cutoff

Dataset 2

ERE₃ ($I = 0$) / ERE₂ ($I = 1$)

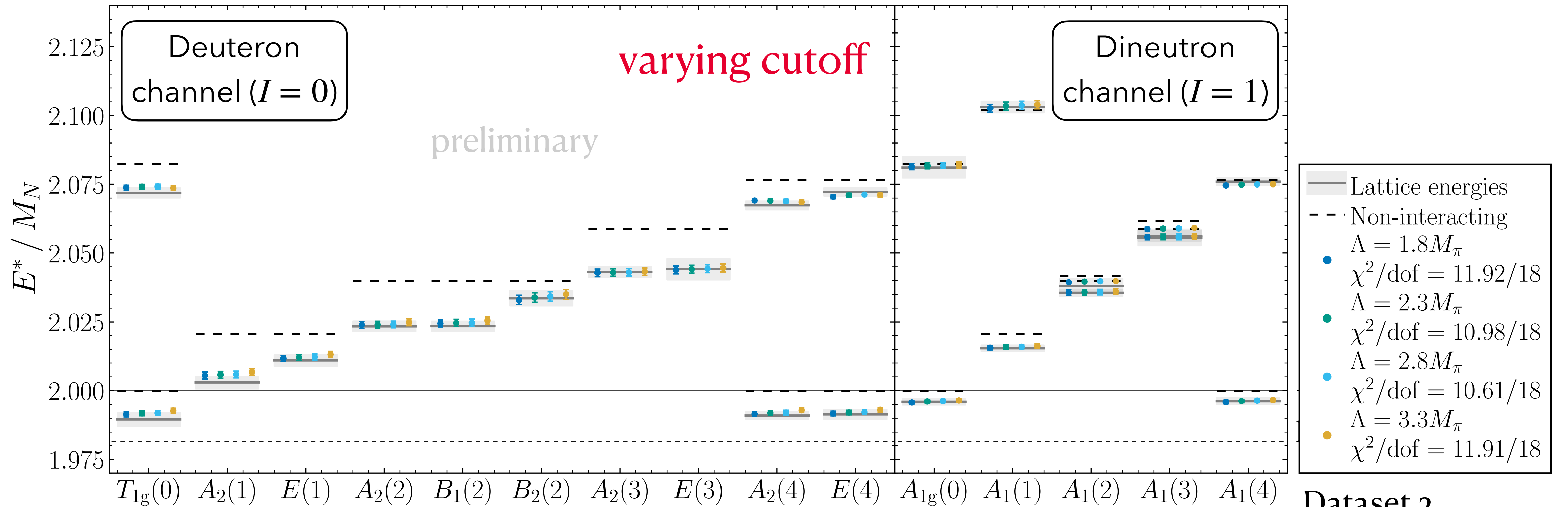
no SD mixing



NN scattering on D251 ($M_\pi = 280$ MeV)

Finite-volume spectrum and fit stability (plane-wave basis)

Combined fit to both isospins with OPE coupling $g_{NN\pi}$ as common fit parameter



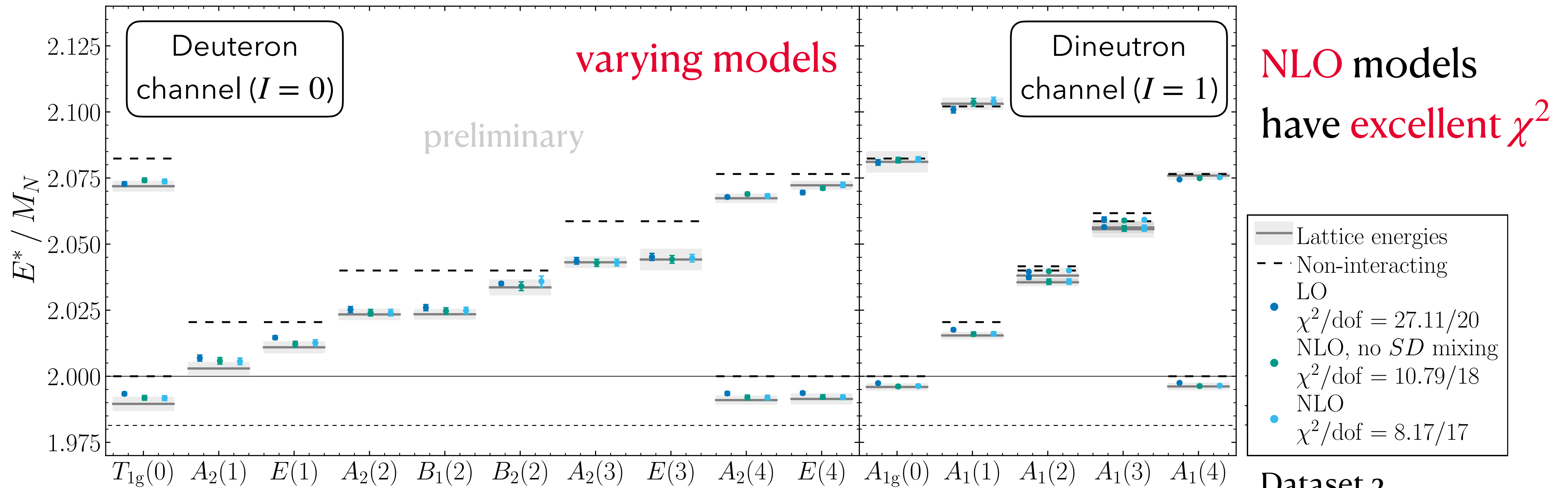
Dataset 2

NLO, no SD mixing 24

NN scattering on D251 ($M_\pi = 280$ MeV)

Finite-volume spectrum and fit stability (plane-wave basis)

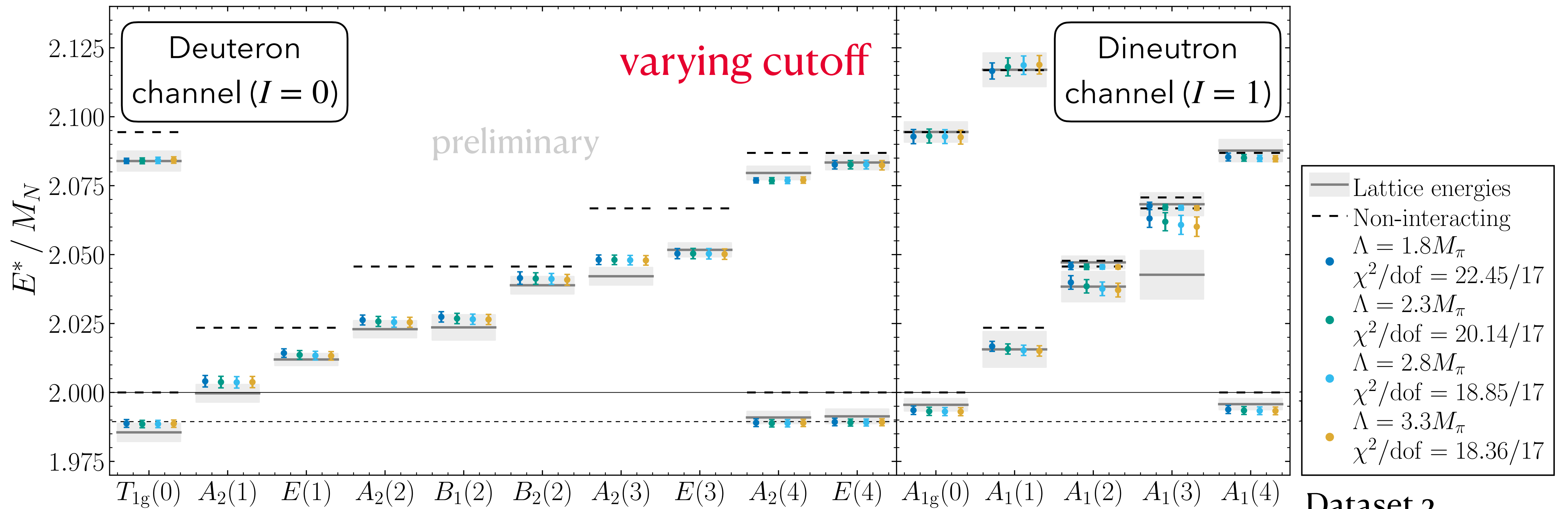
Combined fit to both isospins with OPE coupling $g_{NN\pi}$ as common fit parameter



NN scattering on J250 ($M_\pi = 200$ MeV)

Finite-volume spectrum and fit stability (plane-wave basis)

Combined fit to both isospins with OPE coupling $g_{NN\pi}$ as common fit parameter



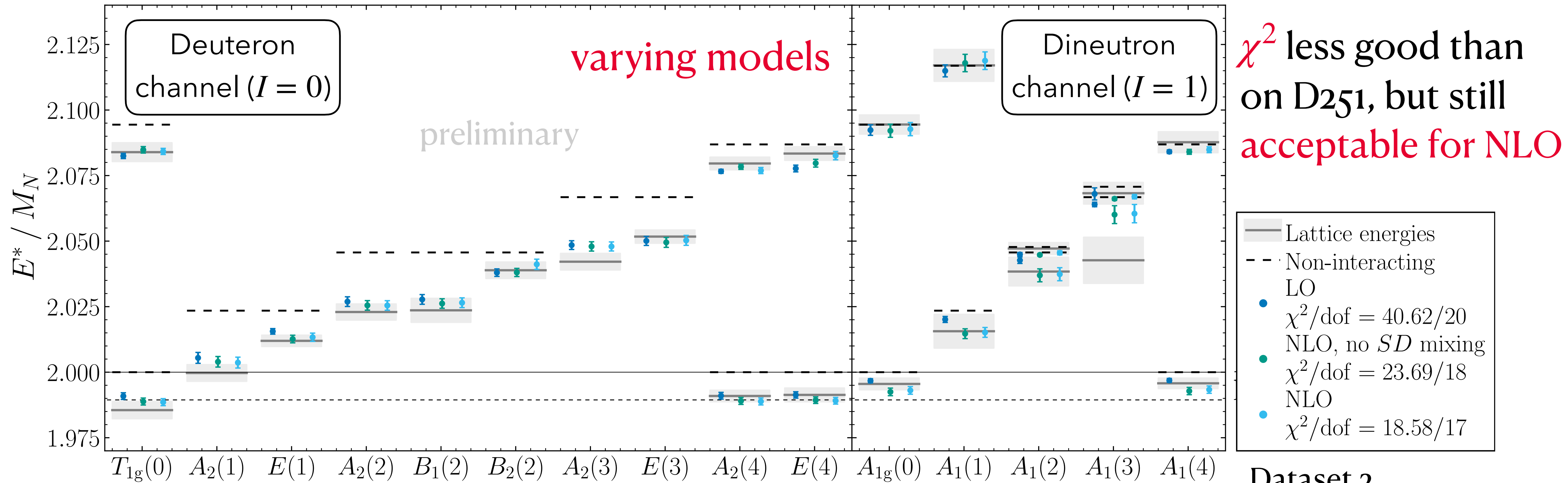
Dataset 2

NLO, with SD mixing₆

NN scattering on J250 ($M_\pi = 200$ MeV)

Finite-volume spectrum and fit stability (plane-wave basis)

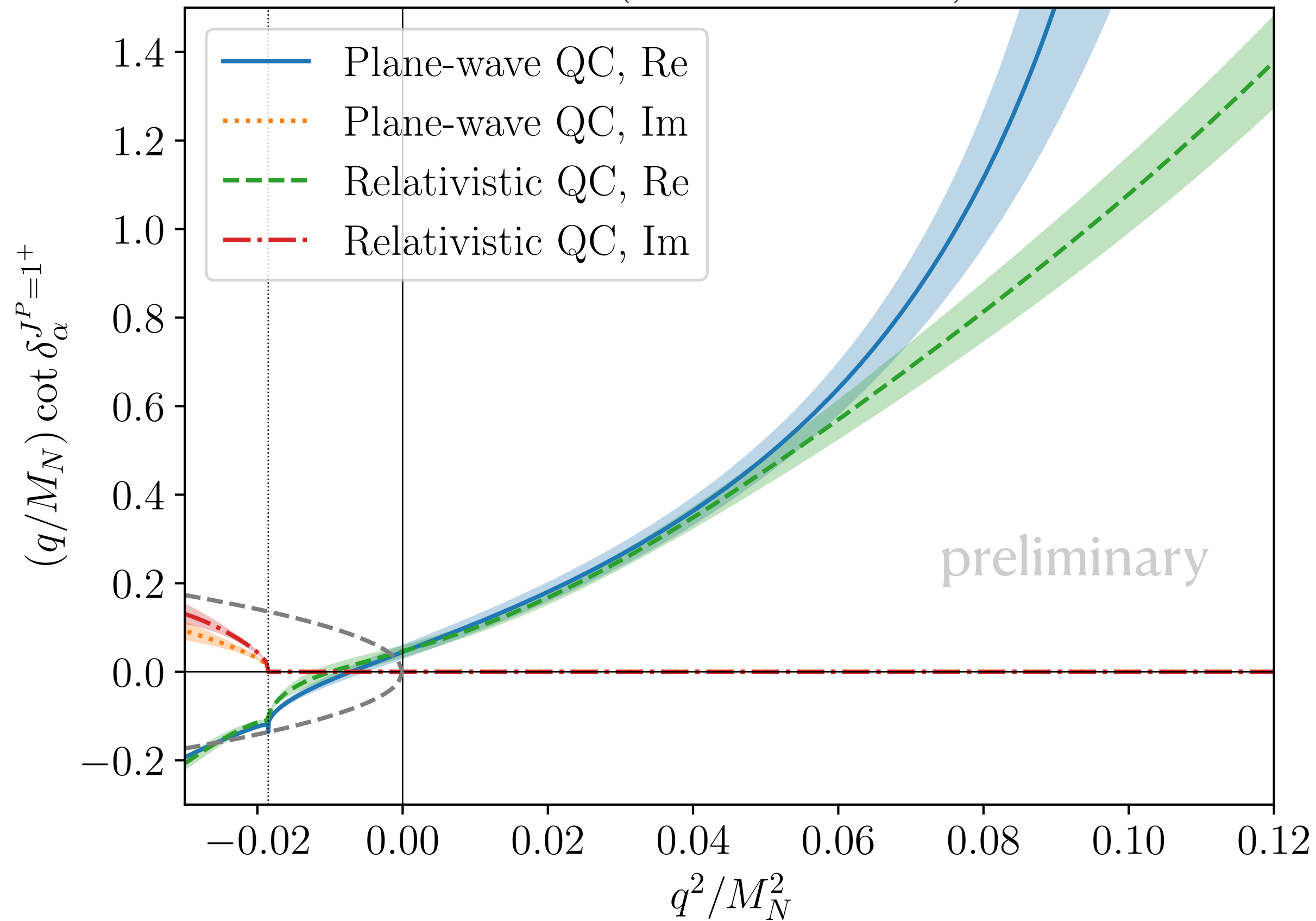
Combined fit to both isospins with OPE coupling $g_{NN\pi}$ as common fit parameter



NN scattering: $I = 0$ (deuteron)

Comparison of plane-wave and relativistic formalisms

D251 ($M_\pi = 280$ MeV)



J250 ($M_\pi = 200$ MeV)

