

Flux Tube Entanglement Entropy Across Temperature, Geometry, and Gauge Group

Rocco Amorosso

with S. Syritsyn and R. Venugopalan



Papers Discussed

- R. Amorosso, S. Syritsyn, and R. Venugopalan, Entanglement entropy of a color flux tube in (2+1)D Yang-Mills theory, JHEP 12 177 (2024), arXiv:2410.00112 [hep-lat].
- R. Amorosso, S. Syritsyn, and R. Venugopalan, Entanglement Enabled Tomography of Flux Tubes in 2+1-D Yang-Mills Theory, to appear in JHEP (2026), arXiv:2601.17199 [hep-lat].
- R. Amorosso, S. Syritsyn, and R. Venugopalan, Upcoming paper on flux tube entanglement at finite temperature, (2026), arXiv:2610.???? [hep-lat].

Roadmap

- Introduction and past results from Flux Tube Entanglement Entropy (FTE²)
- Entanglement radius: an intrinsic “thickness” of the vibrating flux tube that can be measured through entanglement
- Preliminary finite temperature results

What is Flux Tube Entanglement Entropy?

- FTE^2 is the entanglement solely attributable to the color flux tube between a $q\bar{q}$ pair
- FTE^2 answers the question “How much of a region’s entanglement is due to the flux tube, and not the vacuum?”
- The flux tube’s contribution to the EE is isolated through vacuum subtraction
- Can be realized on the lattice using the replica trick and Polyakov loop correlators

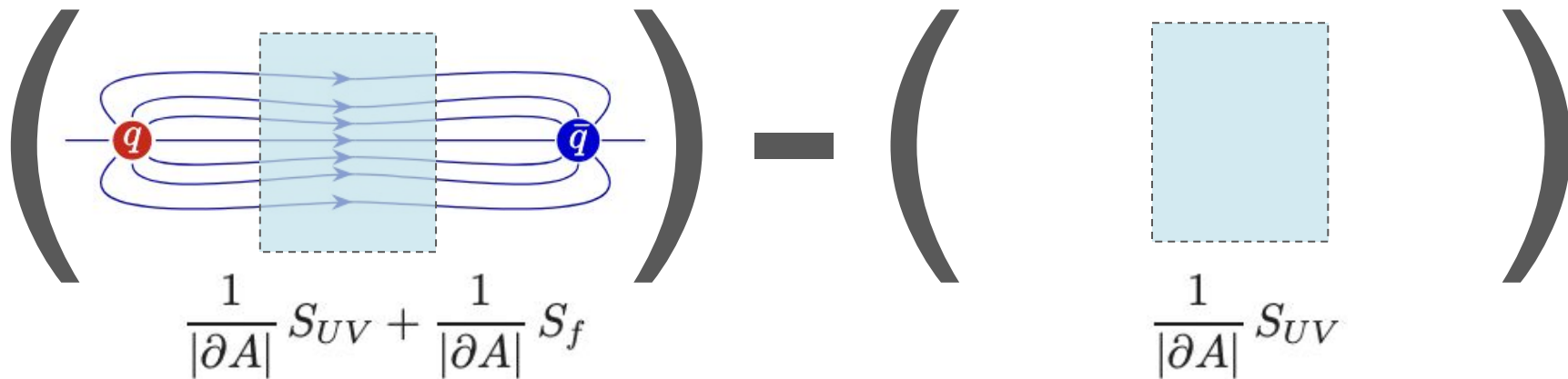
What is Flux Tube Entanglement Entropy?

$$\frac{1}{|\partial A|} S = \frac{1}{|\partial A|} S_{UV} + \frac{1}{|\partial A|} S_f$$

(P. V. Buividovich, M. I. Polikarpov arXiv:0802.4247,
S. Ryu et al arXiv:hep-th/0603001, T. Nishioka et al
arXiv:hep-th/0611035, I. Klebanov et al
arXiv:0709.2140)

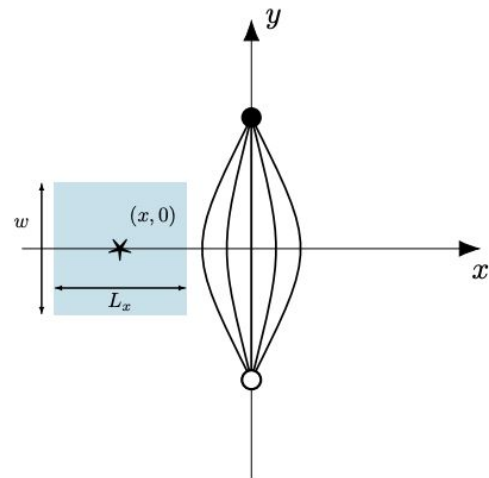
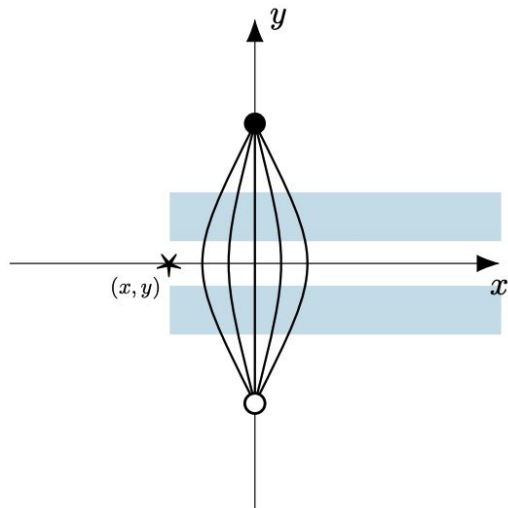
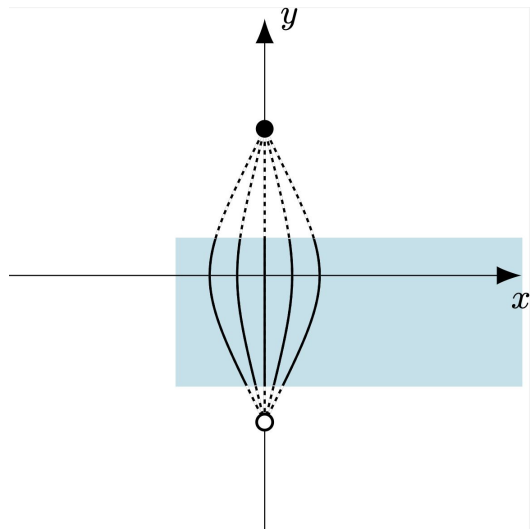
$$\tilde{S}_{|Q\bar{Q}}^{(q)} \equiv S_{|Q\bar{Q}}^{(q)} - S^{(q)}$$

Flux Tube Entanglement Entropy (FTE²)



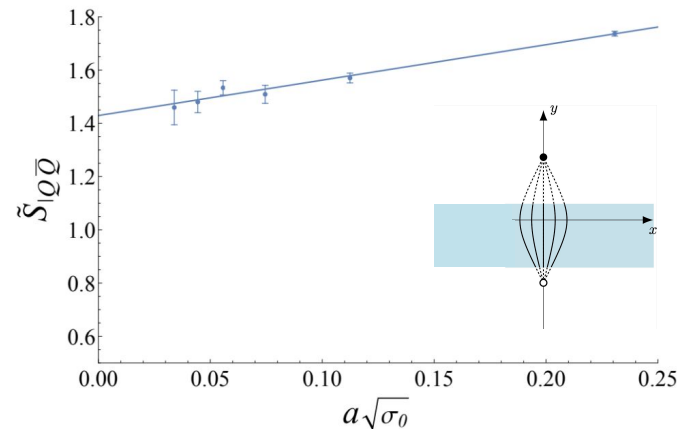
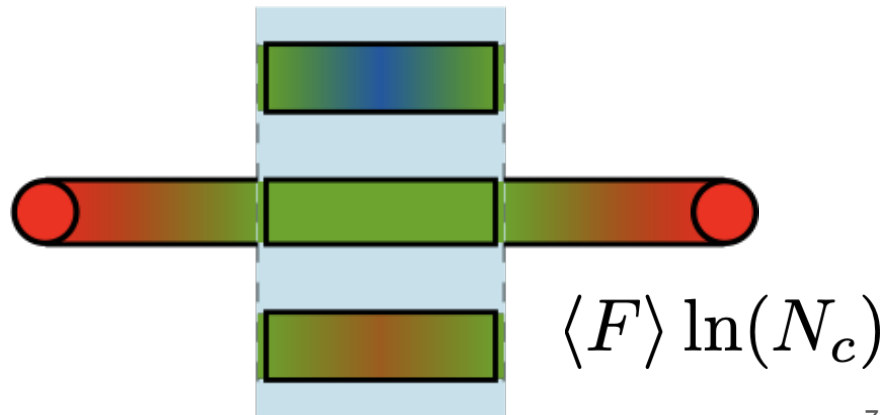
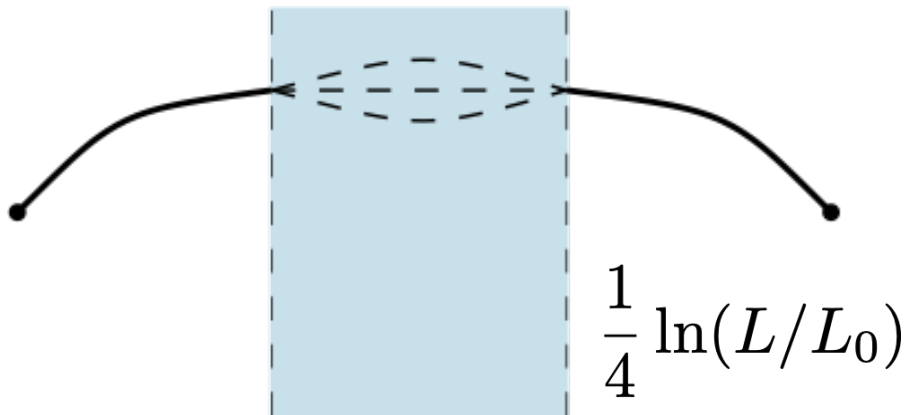
FTE² possibilities

- Can study different gauge groups, entangling regions V , temperatures, etc



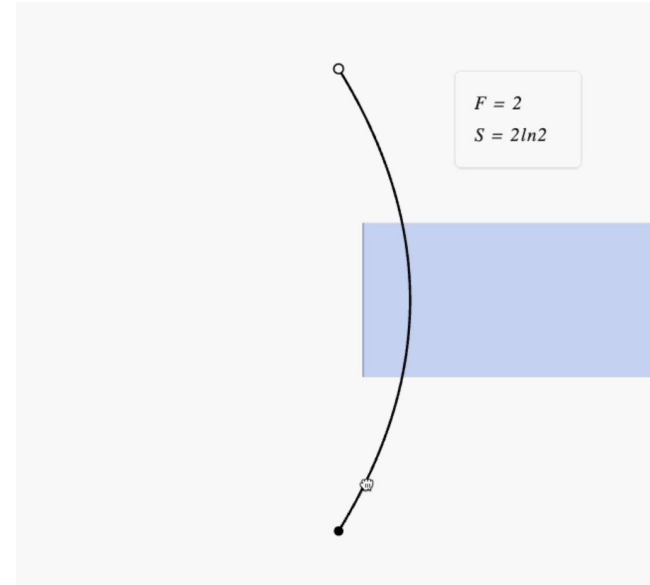
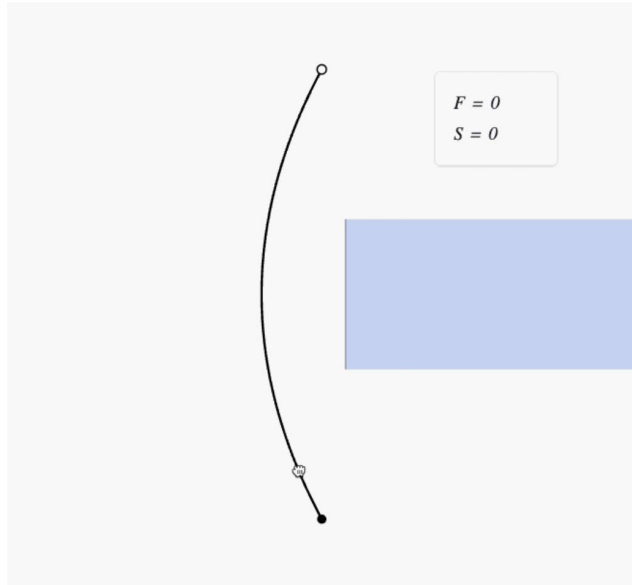
FTE² properties

- Finite, nonzero continuum limit
- Splits into two components following the string model
- Internal: Due to color degrees of freedom inside flux tube
- Vibrational: Due to mechanical vibrations of flux tube



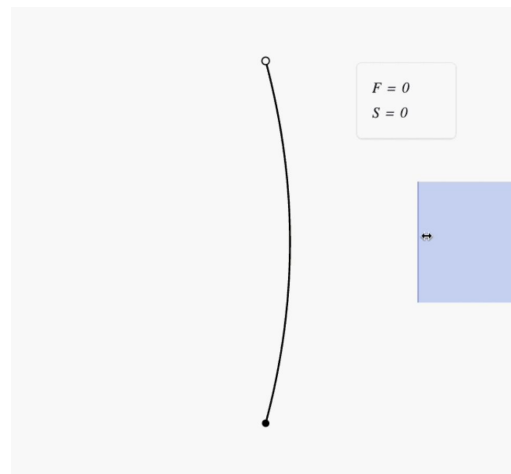
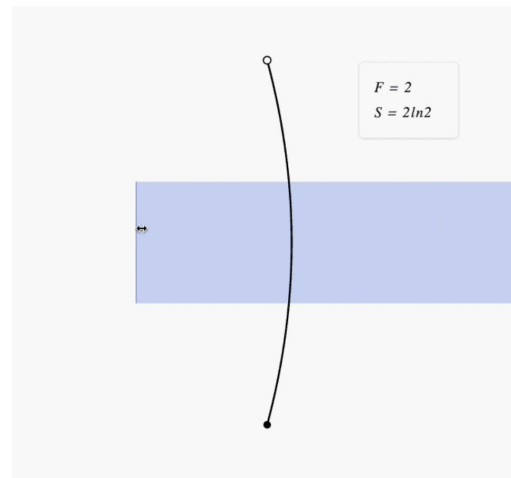
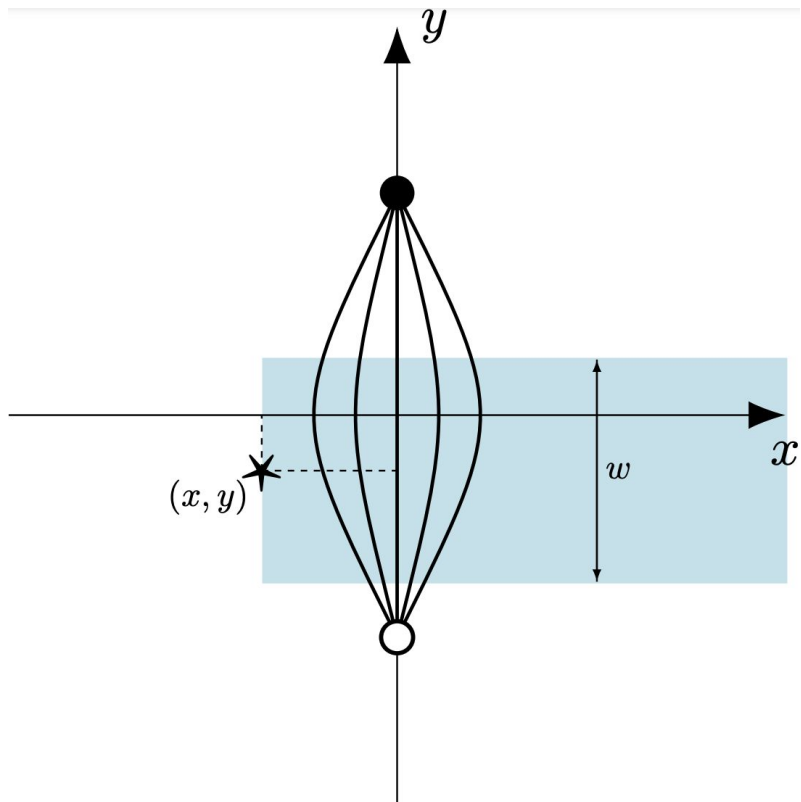
Internal Entropy

- Entanglement for every boundary crossing
- $\langle F \rangle \ln(N_c)$, where $\langle F \rangle$ is expected number of boundary crossings

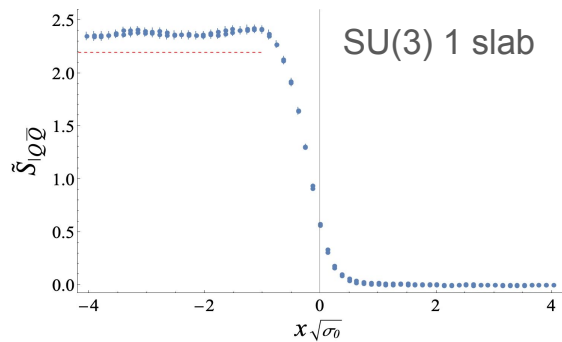
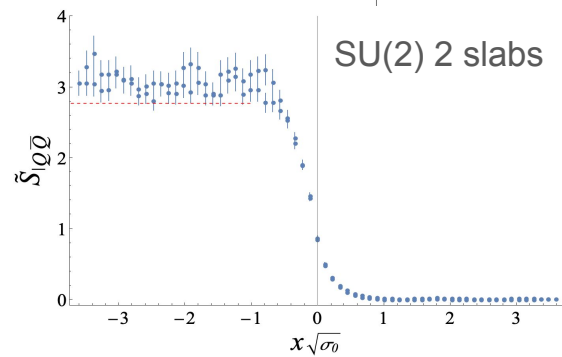
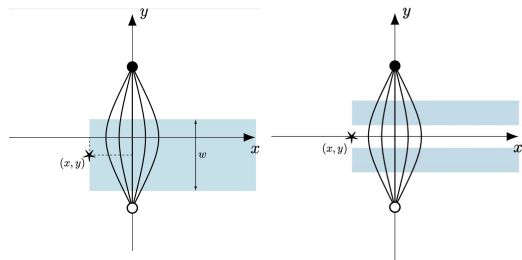
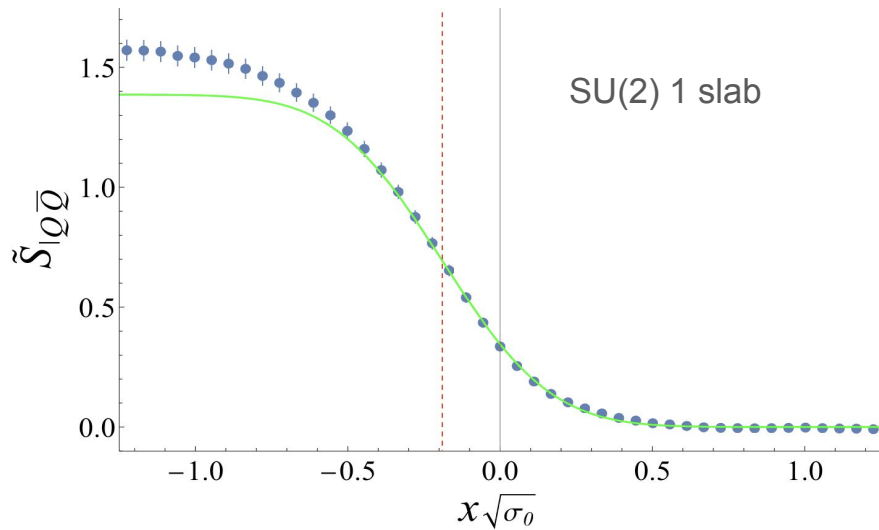


The Entanglement Radius

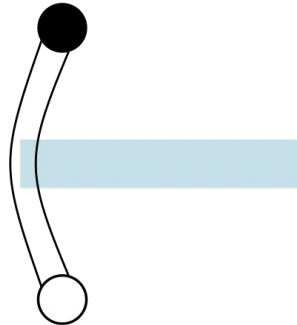
Lattice Setup: A = Half-slab



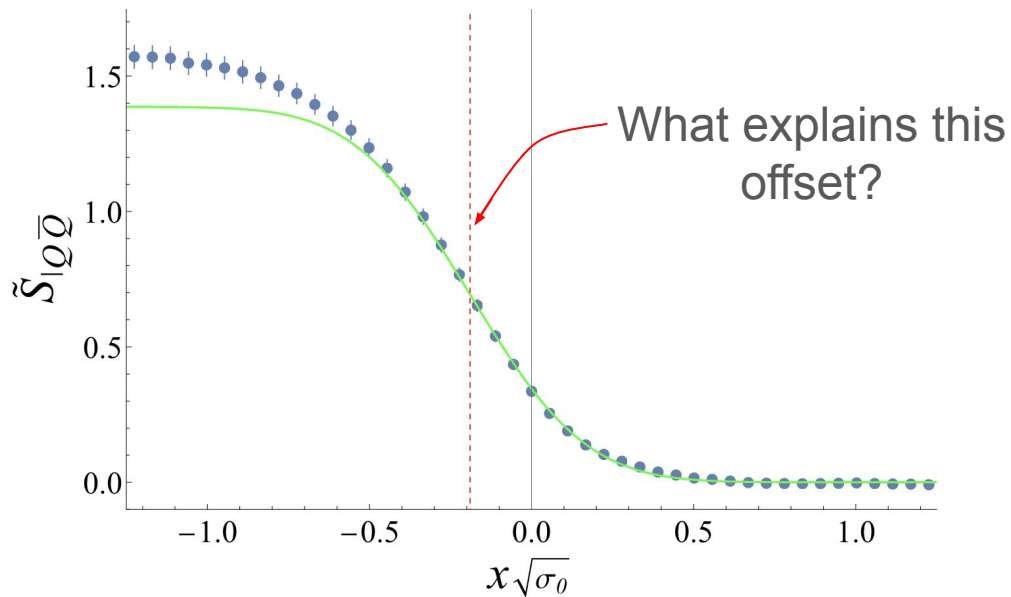
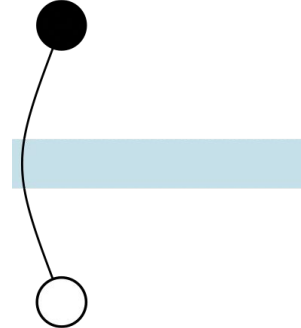
Internal Entropy



Entanglement Radius

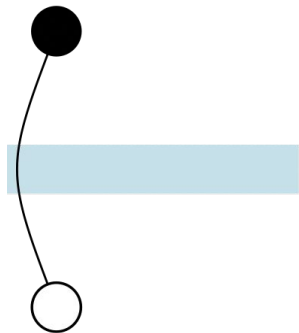


vs.



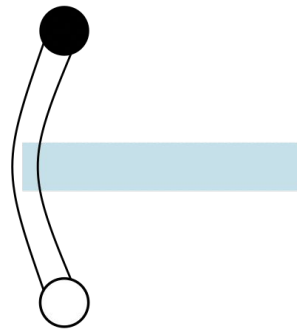
- Flux tube only seems to contribute to EE when center deeper inside region A than $\xi_0\sqrt{\sigma} \approx 0.2$
- ξ_0 has stable, non-zero continuum limit
- What is the meaning of this offset?

Entanglement Radius



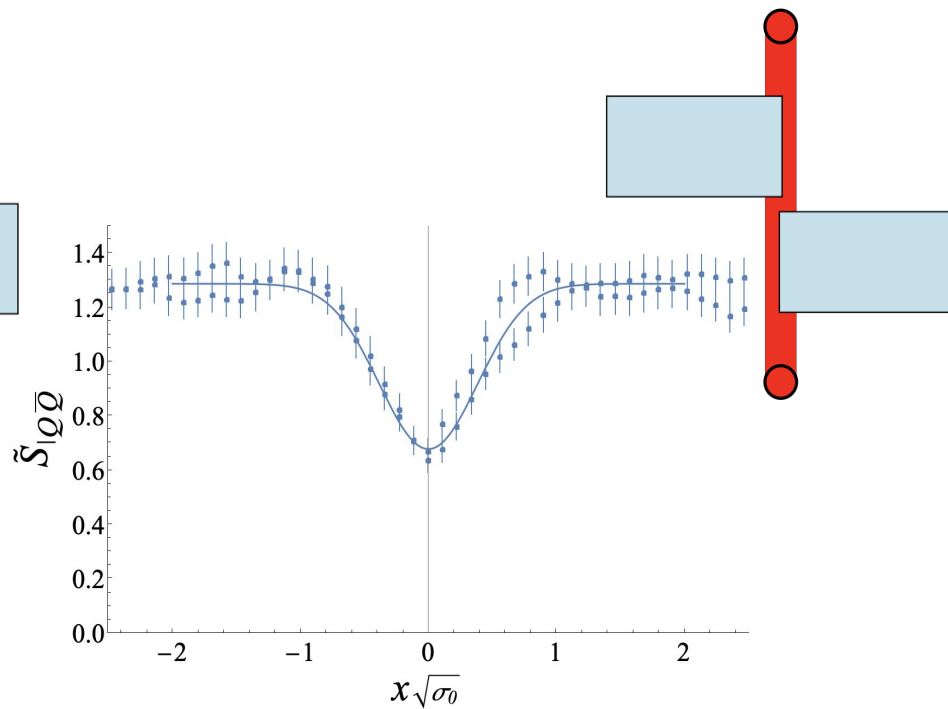
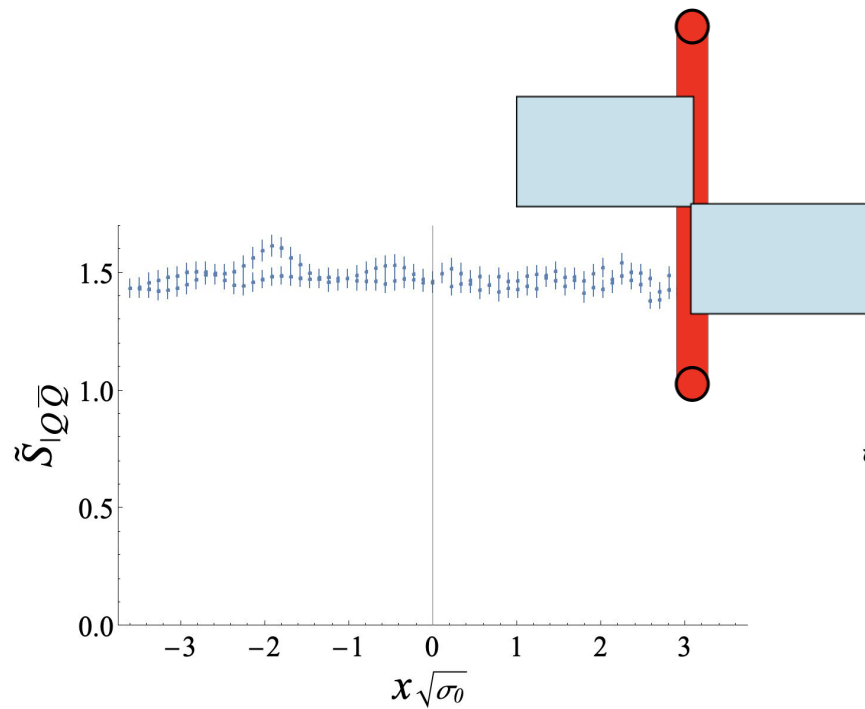
- Full boundary crossing
 - $F = 2$
- $F \ln(N_c)$ entropy contribution

vs.

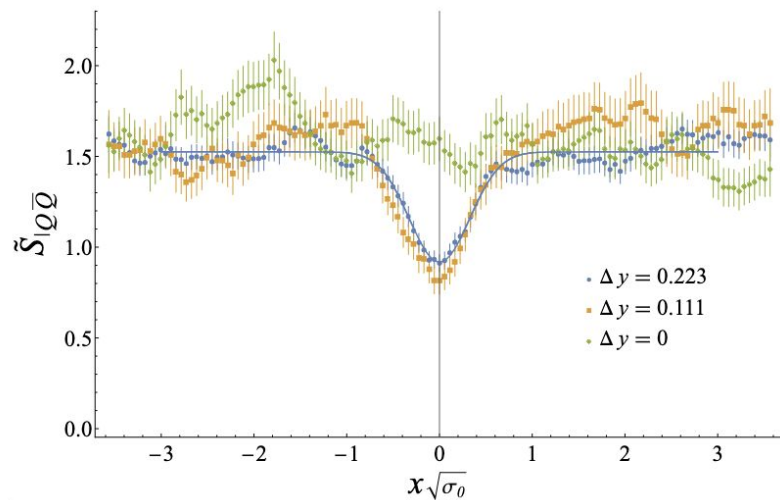
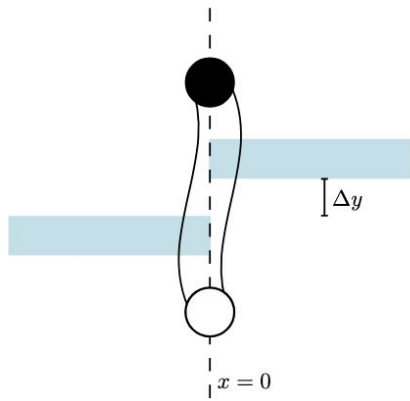
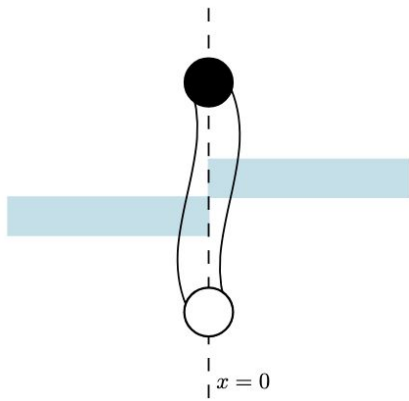


- Partial boundary crossing
 - $F = ?$
- If $F = 0$, could explain offset

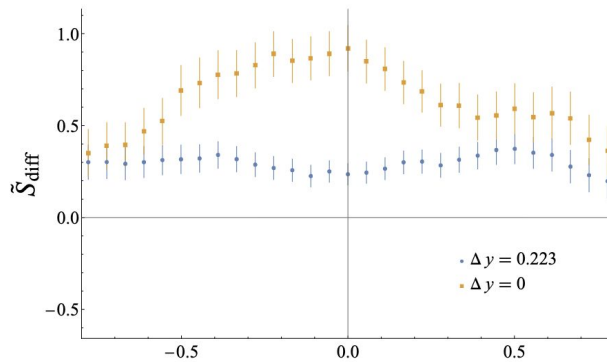
Entanglement Radius



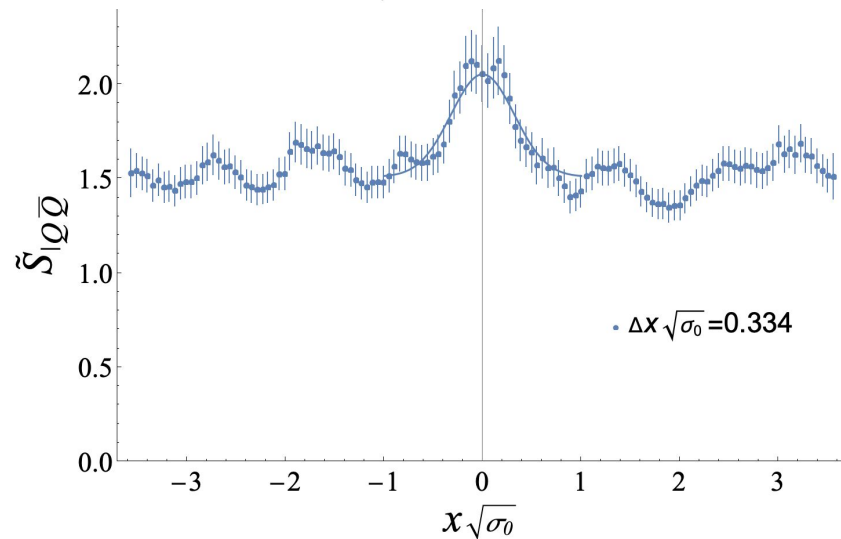
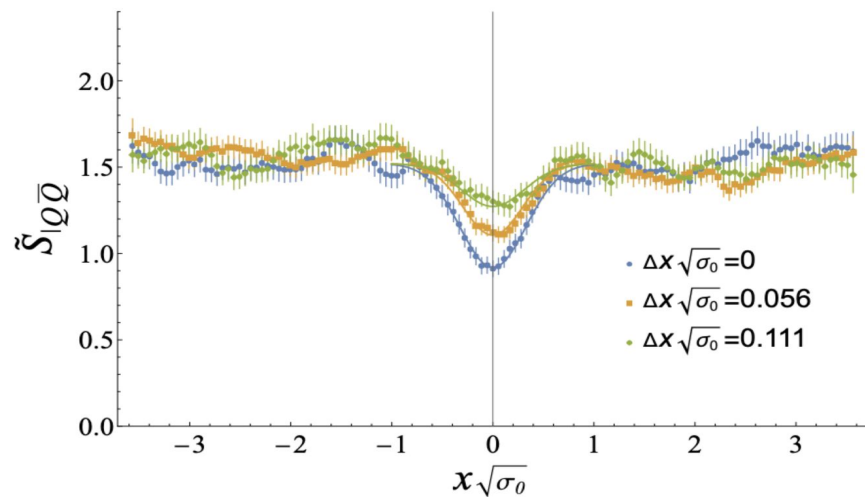
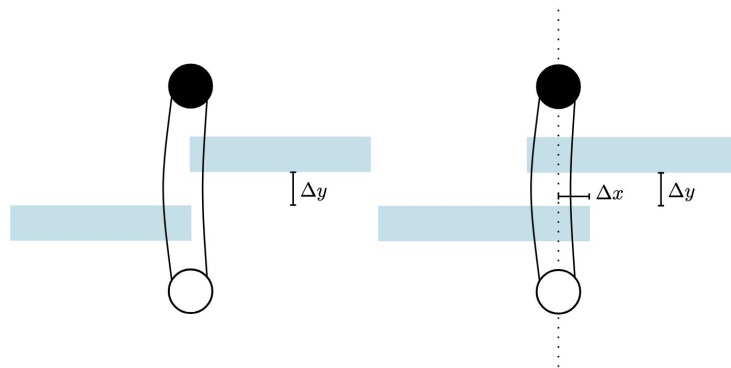
Entanglement Radius



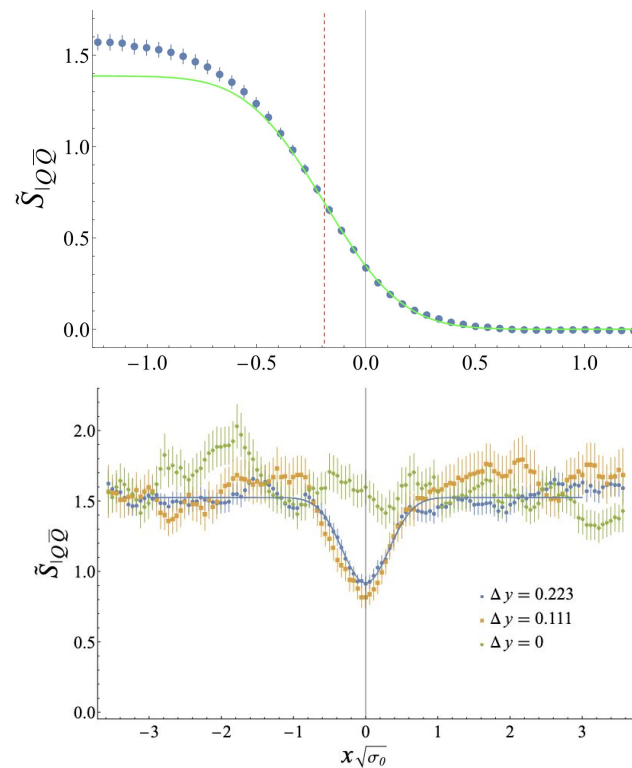
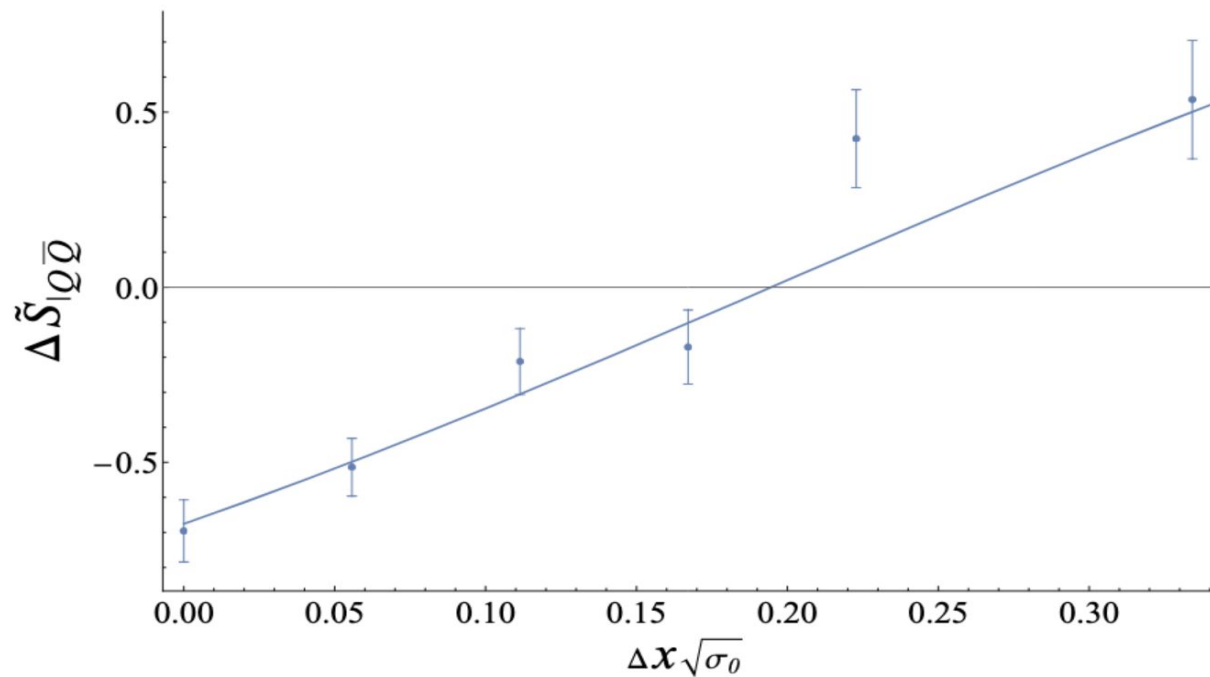
$$\tilde{S}_{\text{diff}}(x) \equiv \tilde{S}_{A \cup B}(x) - (\tilde{S}_A(x) + \tilde{S}_B(x))$$



Entanglement Radius

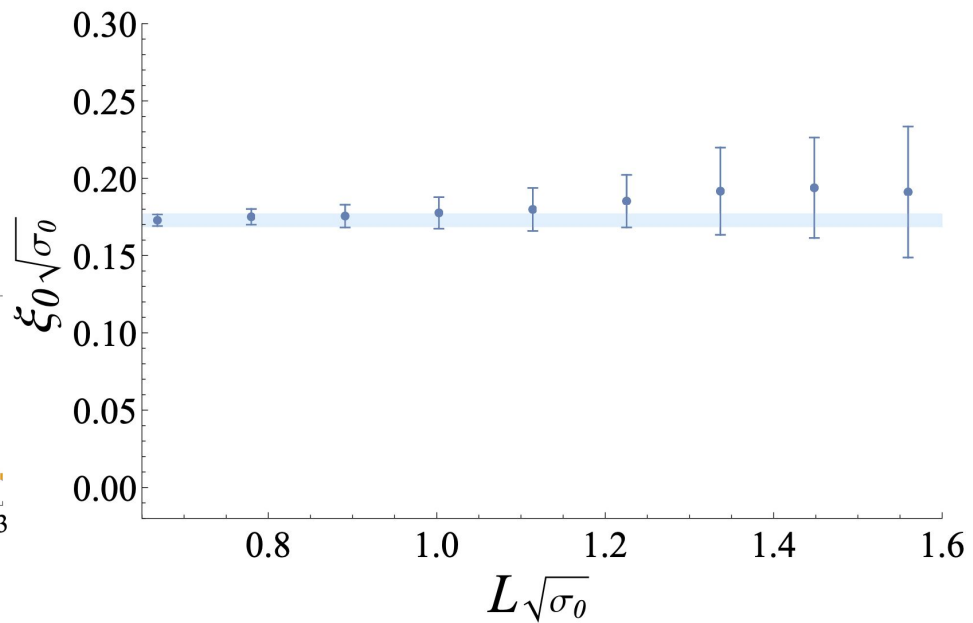
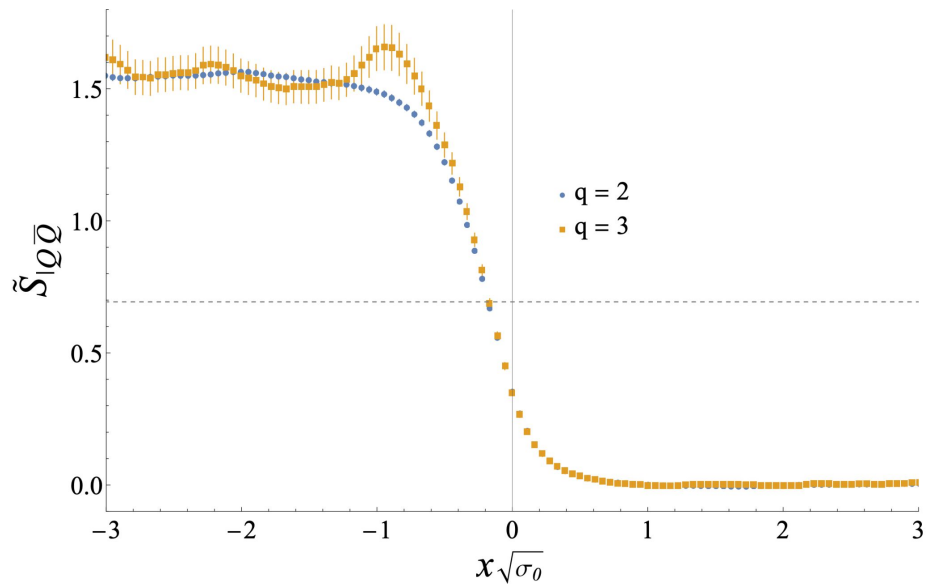


Entanglement Radius

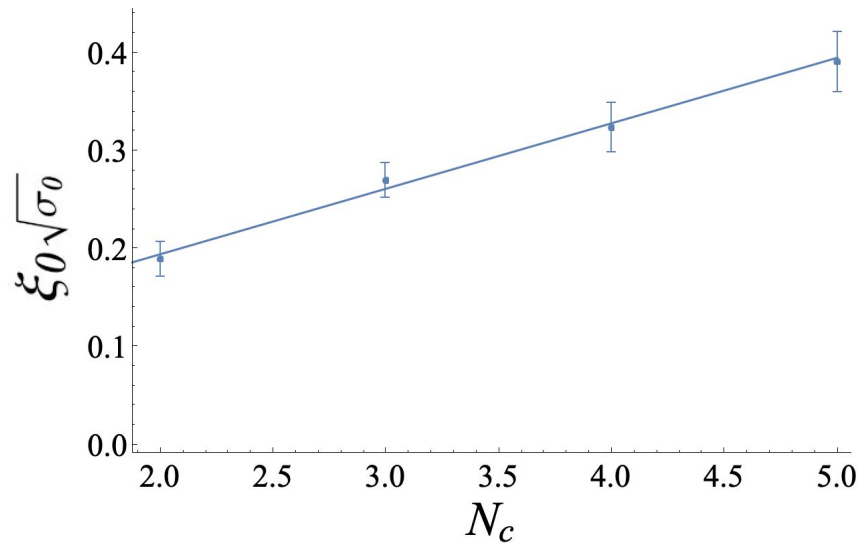
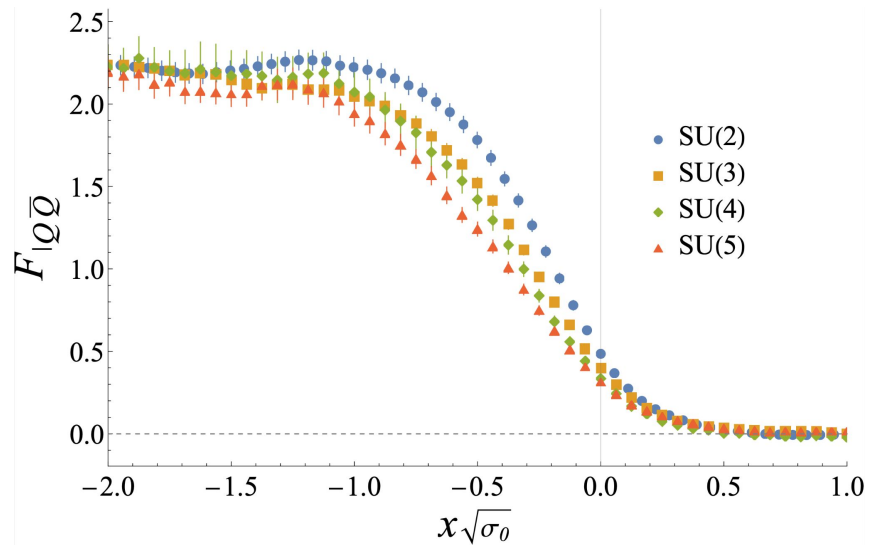


Takeaway: Entanglement Radius is consistent for multiple definitions.

Entanglement Radius

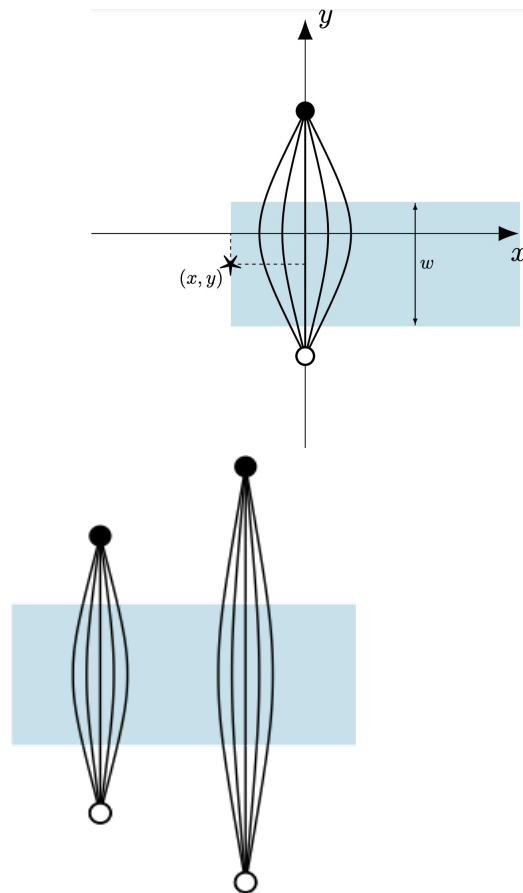
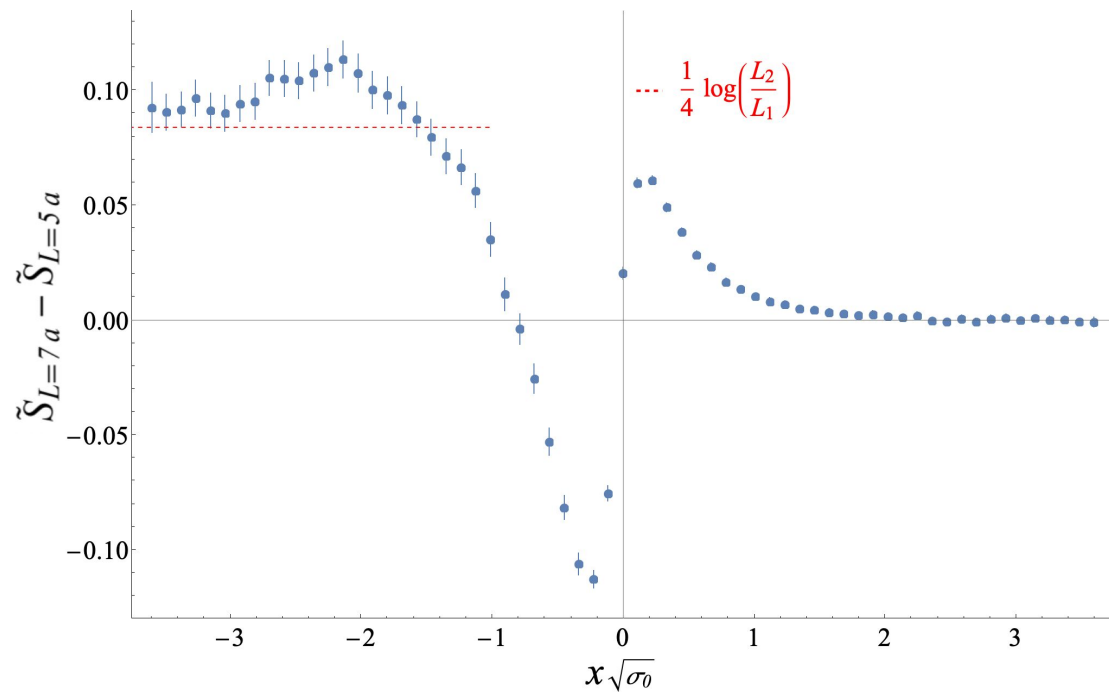


Entanglement Radius

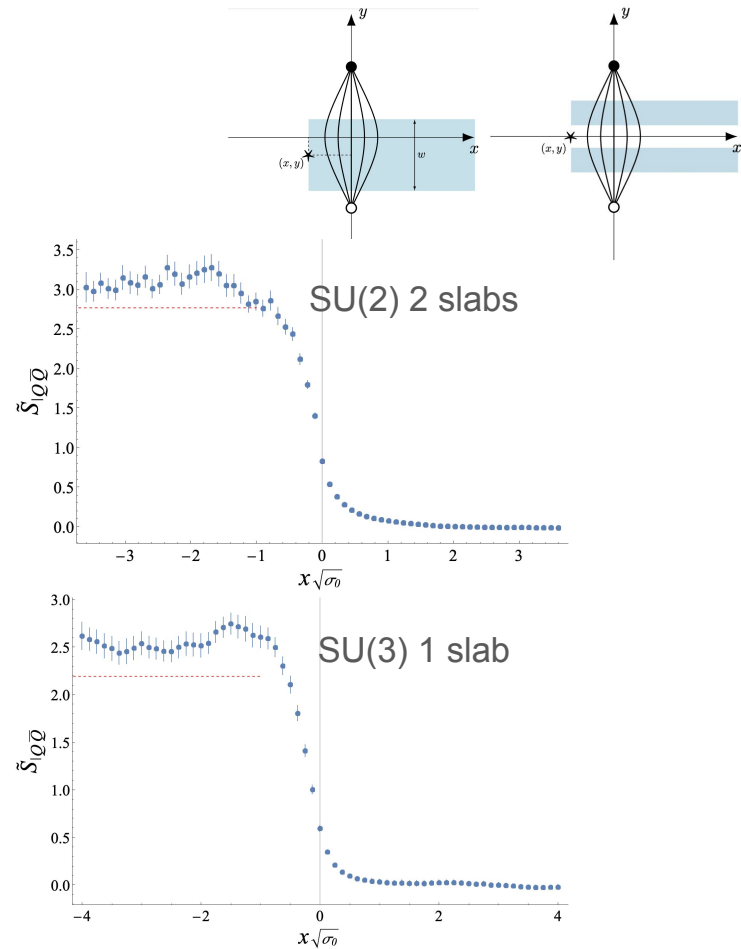
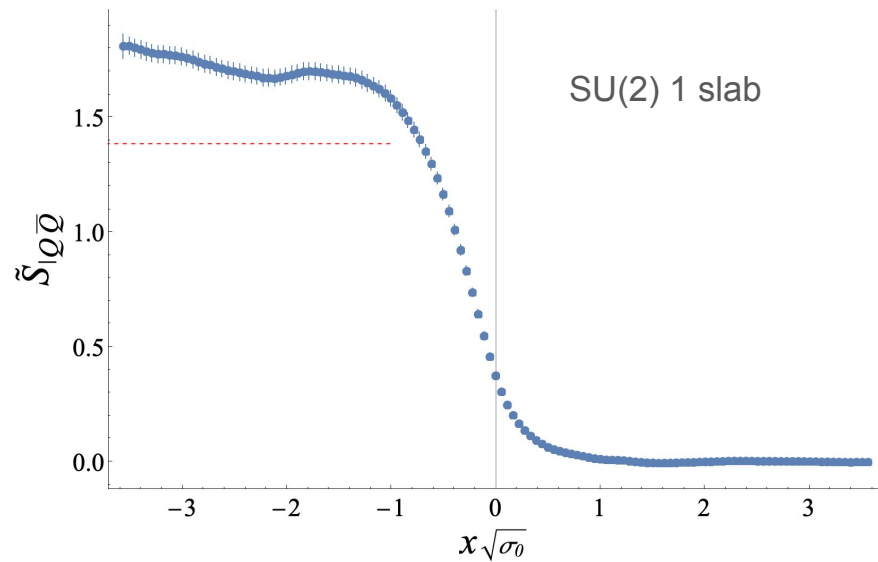


Finite Temperature

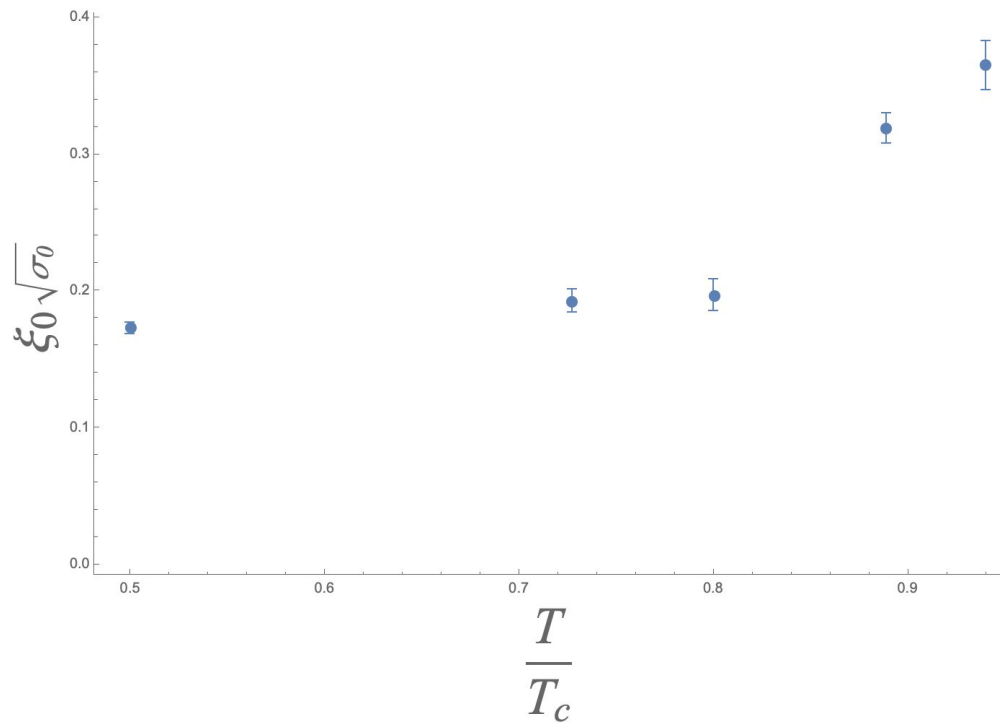
Finite Temperature ($0.8 T_c$)



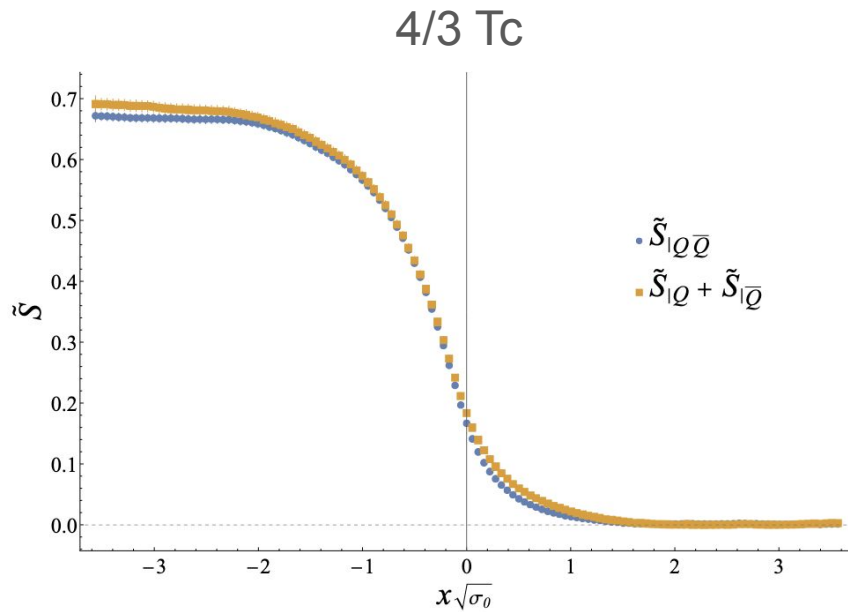
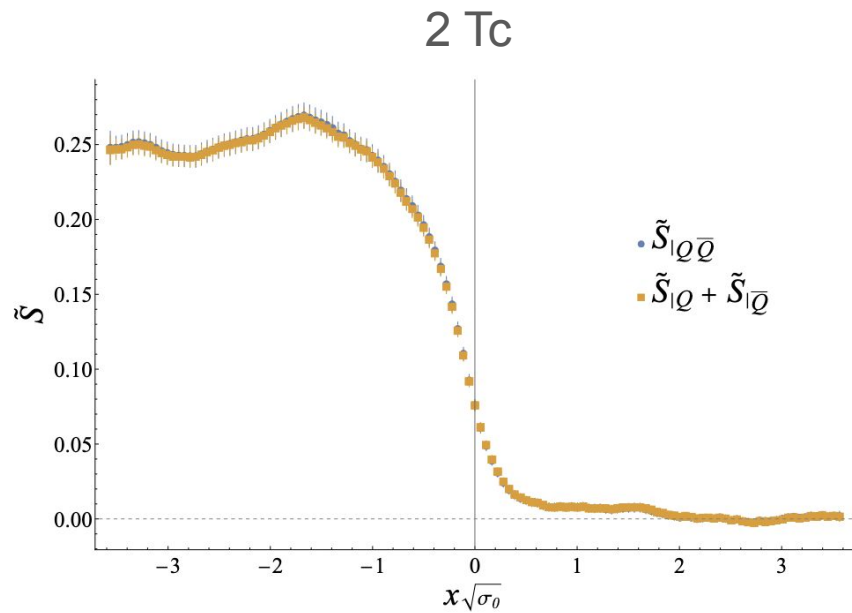
Finite Temperature ($0.8 T_c$)



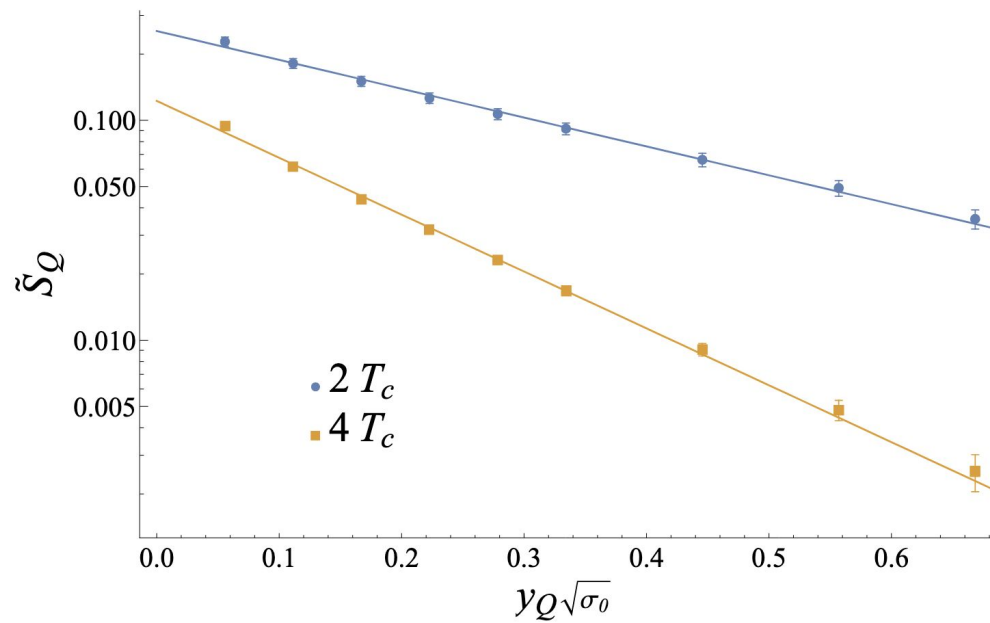
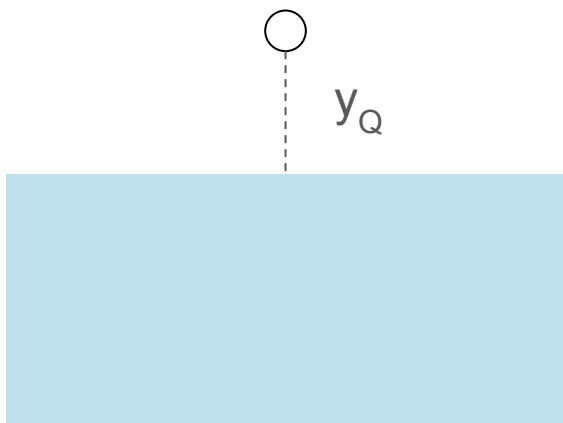
Entanglement radius at finite temperature



Deconfinement



Deconfinement

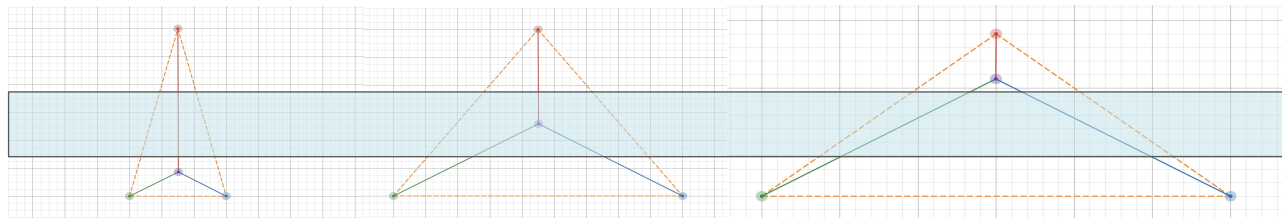
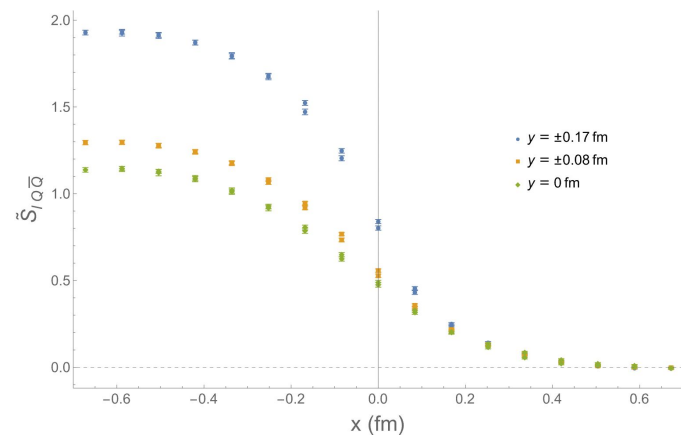
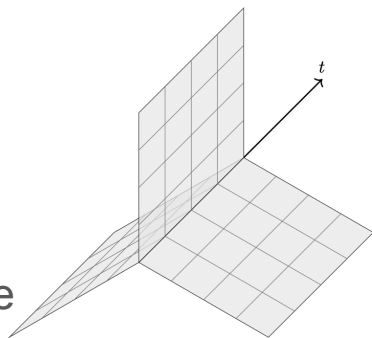


Conclusions

- Vacuum subtraction creates a UV-finite entanglement measure
- Flux tube largely behaves like a thin string with modifications
- Internal entropy follows $\langle F \rangle \ln N_c$ when F counts *full boundary crossings*
- Entanglement radius has continuum limit, is consistent across multiple ways of measuring it
- The flux tube still behaves as an effective string with entanglement radius at finite temperature
- Deconfinement produces much smaller flux tube entanglement

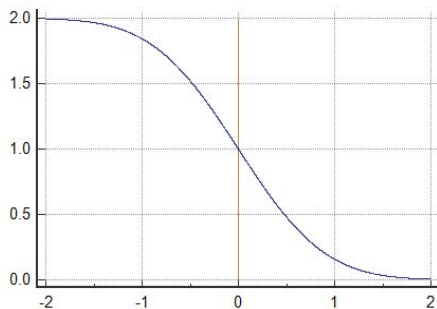
Future Directions

- FTE² with adjoint quarks
- FTE² with static baryons
- Making sense of topological nature of ξ
- Further extensions of FTE² (3+1, dynamical quarks)

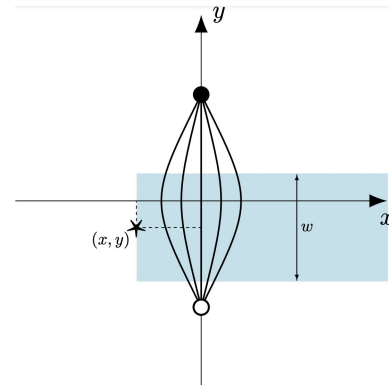
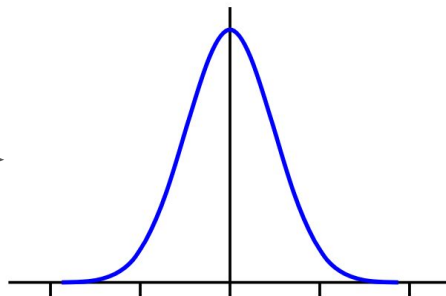


Backup

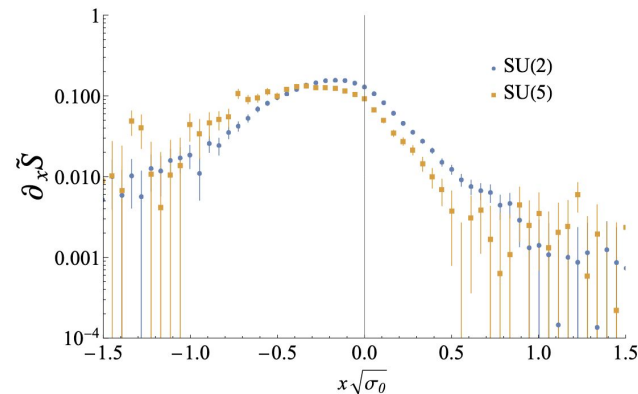
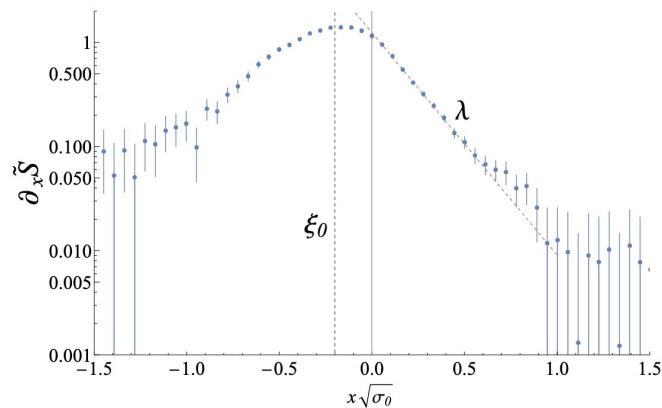
The Decay Length



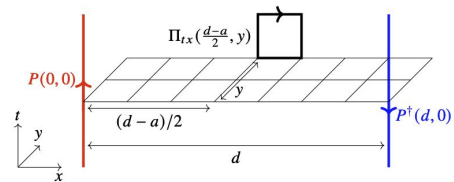
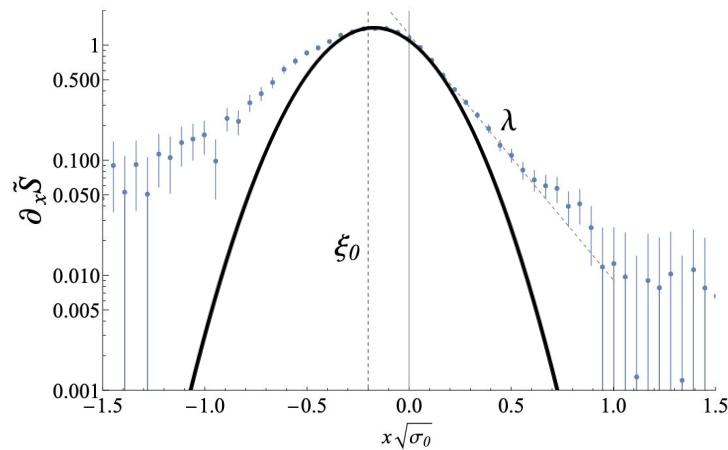
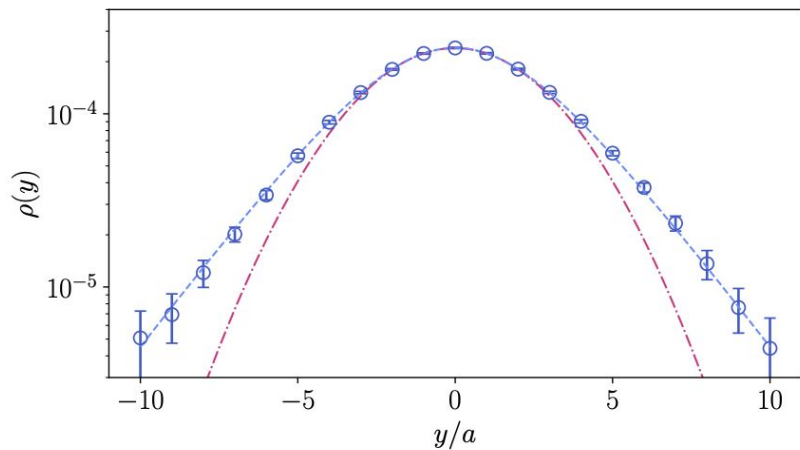
∂_x



New scale: λ



Decay Length vs Intrinsic Width



(L.Verzichelli et al arxiv:2501.01740)

Group	$\xi_0\sqrt{\sigma}$	$\lambda\sqrt{\sigma}$	$\sqrt{\sigma}/m_{0++}$
SU(2)	0.189(6)	0.223(15)	0.212(2)
SU(3)	0.273(12)	0.199(45)	0.231(2)
SU(4)	0.320(8)	0.218(31)	0.236(3)
SU(5)	0.394(9)	0.191(29)	0.239(3)

UV-finite Entanglement Entropy

$$\frac{1}{|\partial A|} S = \frac{1}{|\partial A|} S_{UV} + \frac{1}{|\partial A|} S_f$$

(P. V. Buividovich, M. I. Polikarpov arXiv:0802.4247)
(T. Nishioka, T Takayanagi arXiv:hep-th/0611035)

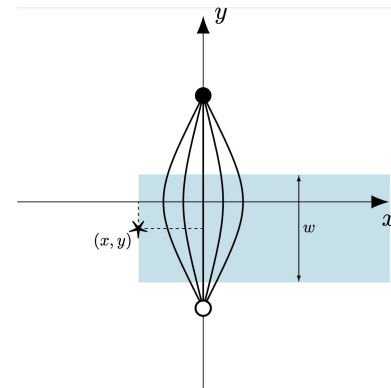
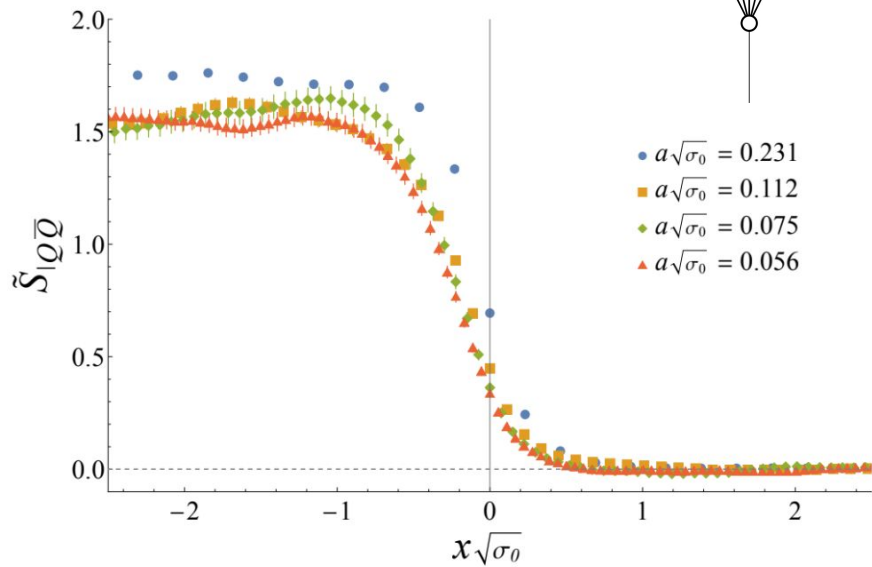
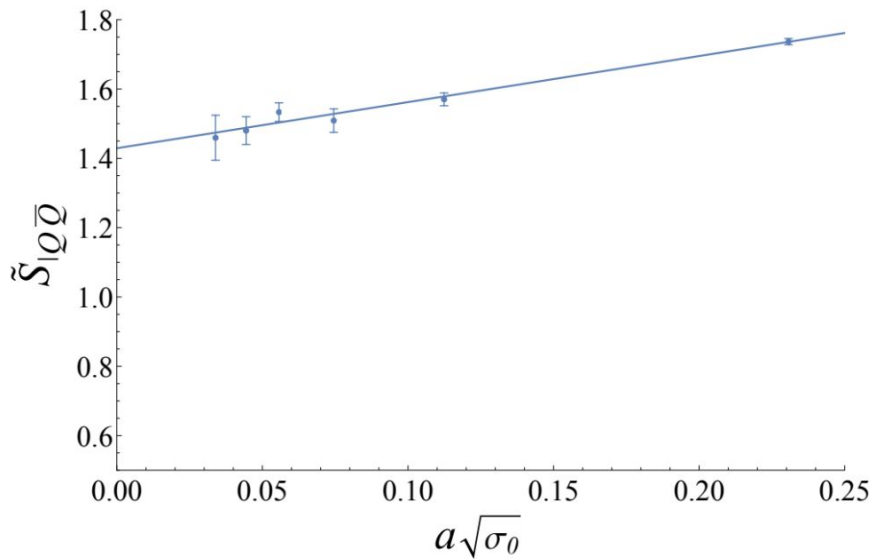
$$\tilde{S}_{|Q\bar{Q}}^{E(q)} \equiv S_{|Q\bar{Q}}^{E(q)} - S^{E(q)}$$

(Flux Tube Entanglement Entropy)

$$S_X^{E(q)} = -\frac{1}{q-1} \log \text{Tr}_A \left[(\text{Tr}_{\bar{A}} \rho_X)^q \right] = -\frac{1}{q-1} \log \frac{Z_X^{(q)}}{(Z_X)^q}$$

$$-\frac{1}{q-1} \log \frac{\langle \prod_{r=1}^q \text{Tr} P_0^{(r)} \text{Tr} P_{\vec{x}}^{(r)\dagger} \rangle}{\langle \text{Tr} P_0 \text{Tr} P_{\vec{x}}^\dagger \rangle^q} = -\frac{1}{q-1} \log \frac{Z_{|Q\bar{Q}}^{(q)} / Z^{(q)}}{(Z_{|Q\bar{Q}} / Z)^q} = -\frac{1}{q-1} \log \left(\frac{Z_{|Q\bar{Q}}^{(q)}}{(Z_{|Q\bar{Q}})^q} \cdot \frac{Z^q}{Z^{(q)}} \right) = S_{|Q\bar{Q}}^{E(q)} - S^{E(q)}$$

Finite Entanglement Entropy (Scaling Study)

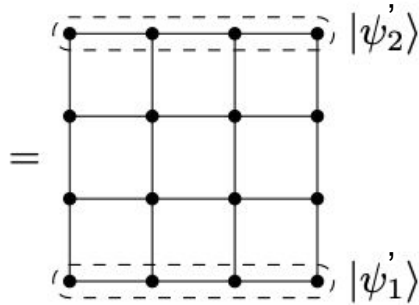


Takeaway: FTE^2 has finite non-zero continuum limit for $x \lesssim 0$

Reduced Density Matrix in Lattice Field Theory

$$\hat{\rho}_{\mathcal{A}} = \text{tr}_{\bar{\mathcal{A}}} (\hat{\rho})$$

$$\langle \psi'_1 | \hat{\rho} | \psi'_2 \rangle = \frac{1}{Z} \langle \psi'_1 | e^{-\beta \hat{H}} | \psi'_2 \rangle$$



$$= \sum_{|\psi\rangle} \langle \psi_1 | \hat{\rho}_{\mathcal{A}} | \psi_2 \rangle$$

A 4x4 grid of points representing a lattice. The top and bottom rows are enclosed in dashed boxes and labeled $|\psi_2\rangle$ and $|\psi_1\rangle$ respectively. The grid is divided into two regions: $\bar{\mathcal{A}}$ (the left half) and $\mathcal{A}$ (the right half). A summation symbol $\sum_{|\psi\rangle}$ is placed to the left of the grid.

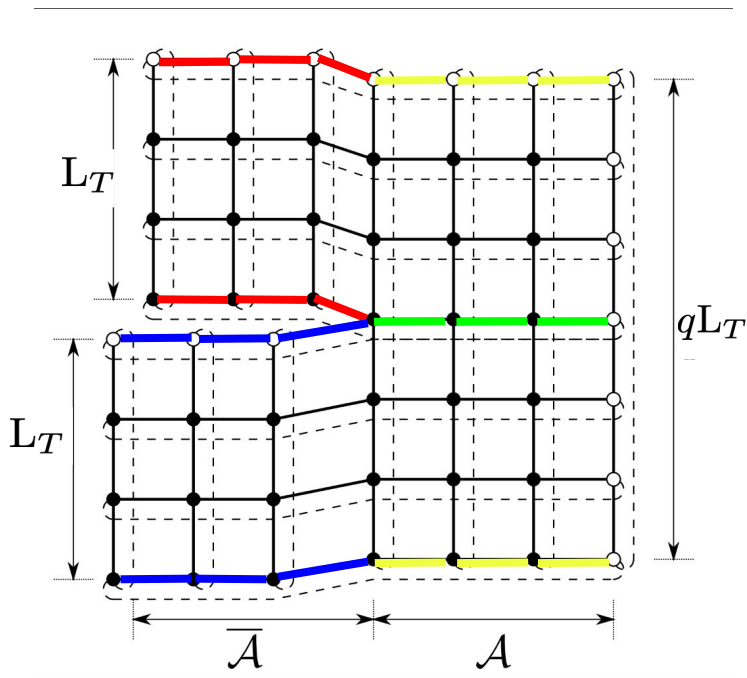
(A. Rabenstein et al arXiv:1812.04279)

Reduced Density Matrix squared

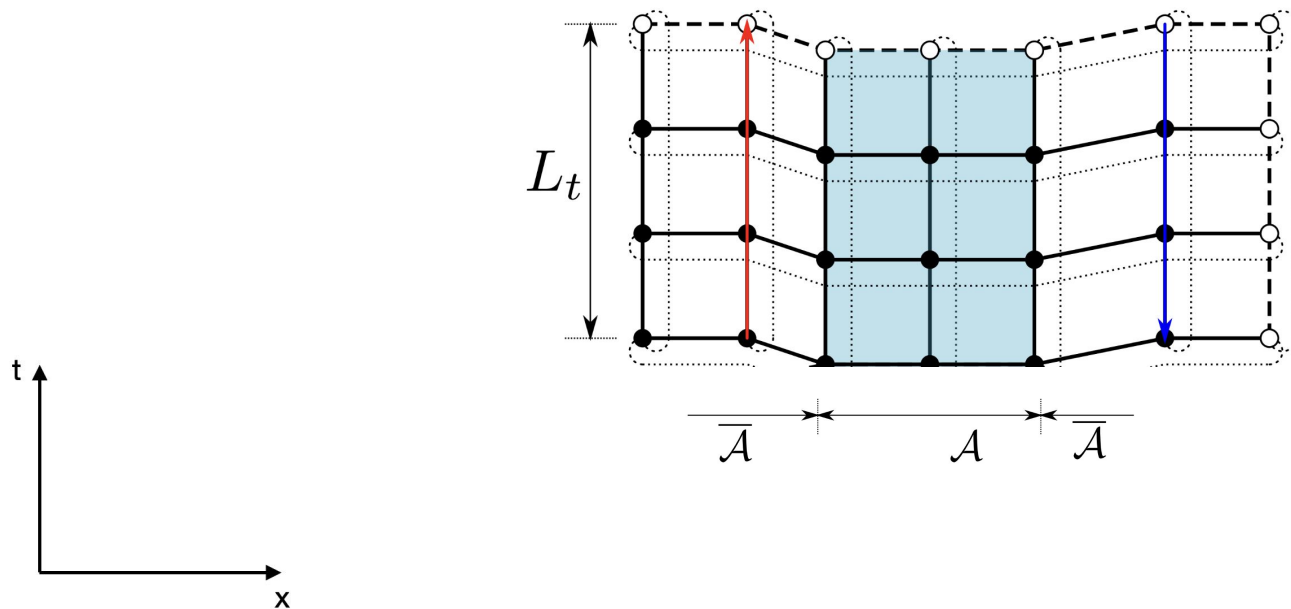
$$\langle \psi_1 | \hat{\rho}_{\mathcal{A}}^2 | \psi_2 \rangle = \langle \psi_1 | \hat{\rho}_{\mathcal{A}} | \psi_k \rangle \langle \psi_k | \hat{\rho}_{\mathcal{A}} | \psi_2 \rangle = \sum_{|\psi_k\rangle, |\psi_{\bar{\mathcal{A}}_1}\rangle, |\psi_{\bar{\mathcal{A}}_2}\rangle} \langle \psi_{\bar{\mathcal{A}}_2} | \langle \psi_{\bar{\mathcal{A}}_1} | \langle \psi_k | \hat{\rho}_{\mathcal{A}} | \psi_2 \rangle \hat{\rho}_{\mathcal{A}} | \psi_1 \rangle$$

Reduced Density Matrix squared

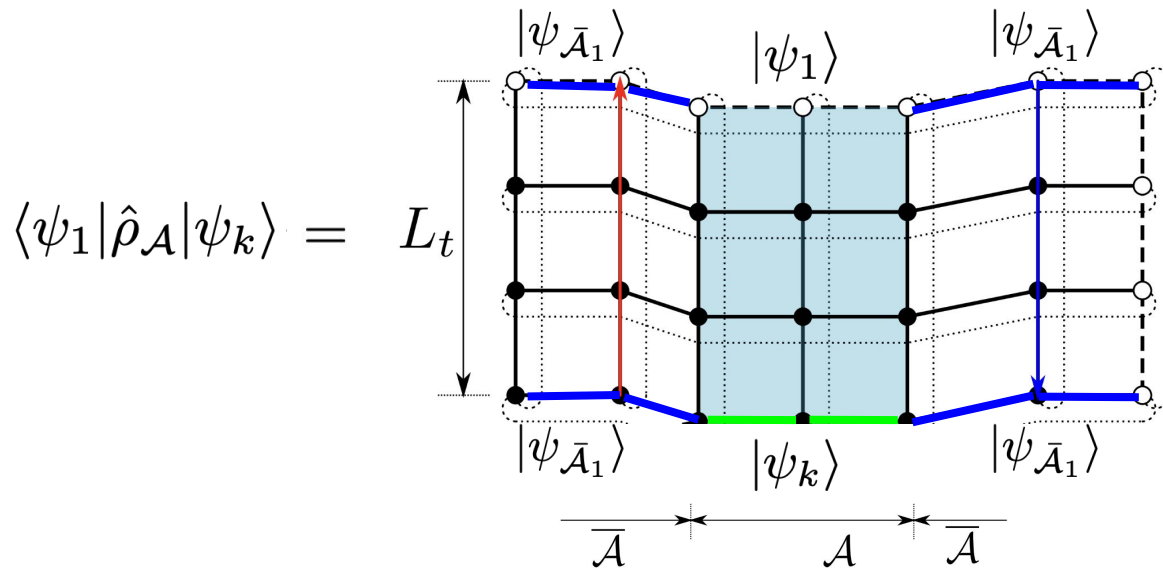
$$\begin{aligned}\text{Tr} \hat{\rho}_{\mathcal{A}}^2 &= \sum_{\psi_{\mathcal{A}}} \langle \psi_{\mathcal{A}} | \hat{\rho}_{\mathcal{A}}^2 | \psi_{\mathcal{A}} \rangle = \\ &= \frac{Z_2}{(Z_1)^2}\end{aligned}$$



Polyakov Lines

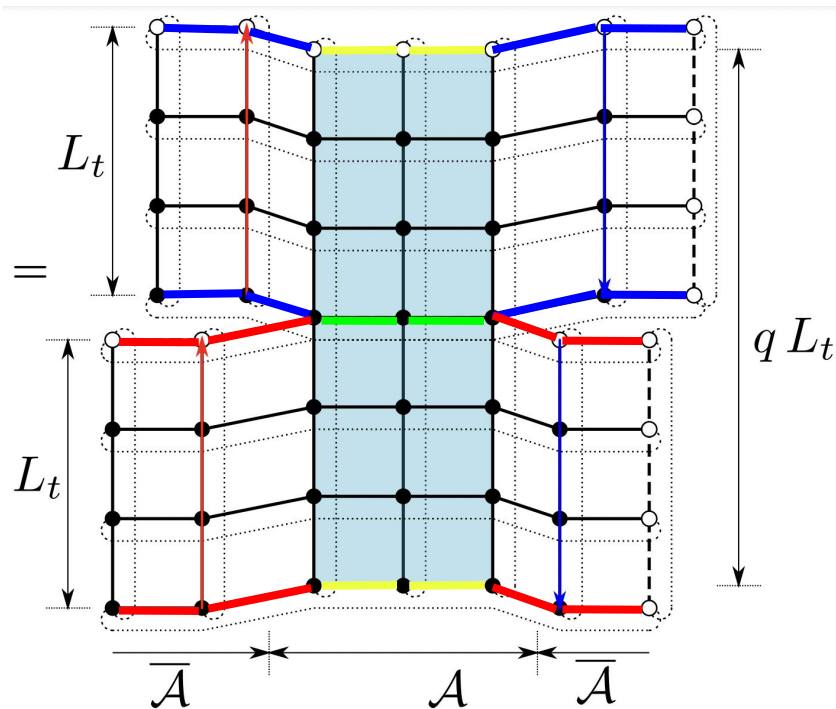
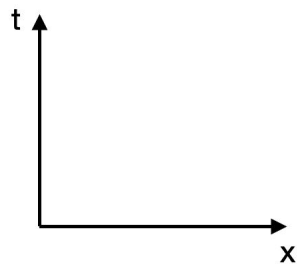


Polyakov Lines: Reduced Density Matrix



Polyakov Lines: Reduced Density Matrix squared

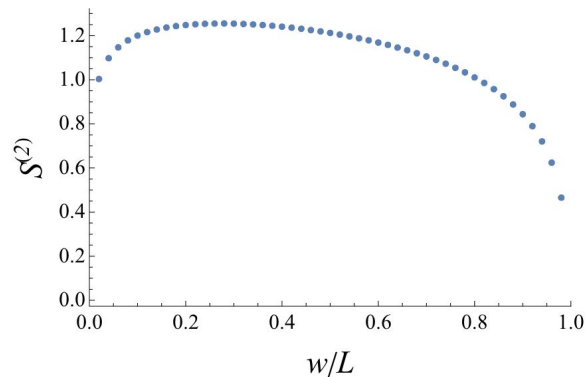
$$\text{Tr} \hat{\rho}_{\mathcal{A}}^2 = \sum_{\psi_{\mathcal{A}}} \langle \psi_{\mathcal{A}} | \hat{\rho}_{\mathcal{A}}^2 | \psi_{\mathcal{A}} \rangle =$$



String Model Expectations (Vibrational)

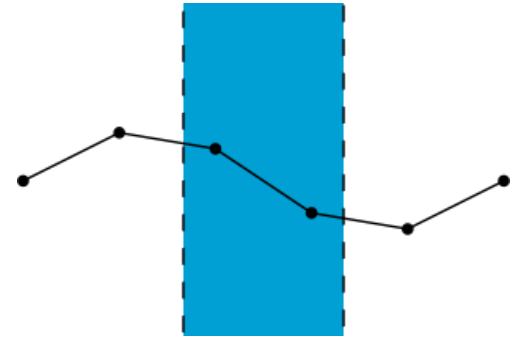
- Cut string into small pieces of size ϵ
- Numerically calculate entropy taking L/ϵ to infinity
- Splits into cutoff scaling and string proportion terms

$$S^{(2)} = \frac{1}{4} \log(L/\epsilon) + \text{finite contributions}$$

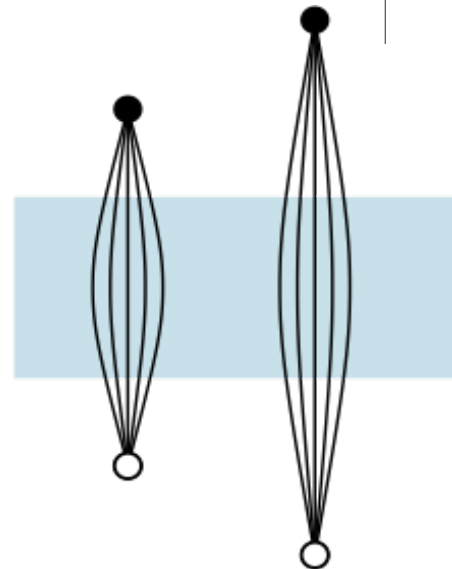
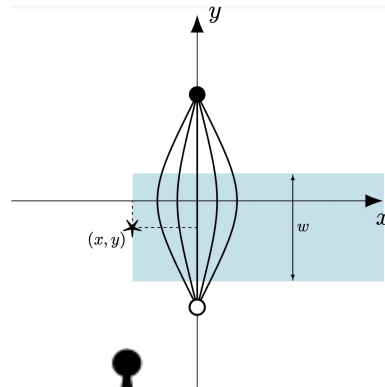
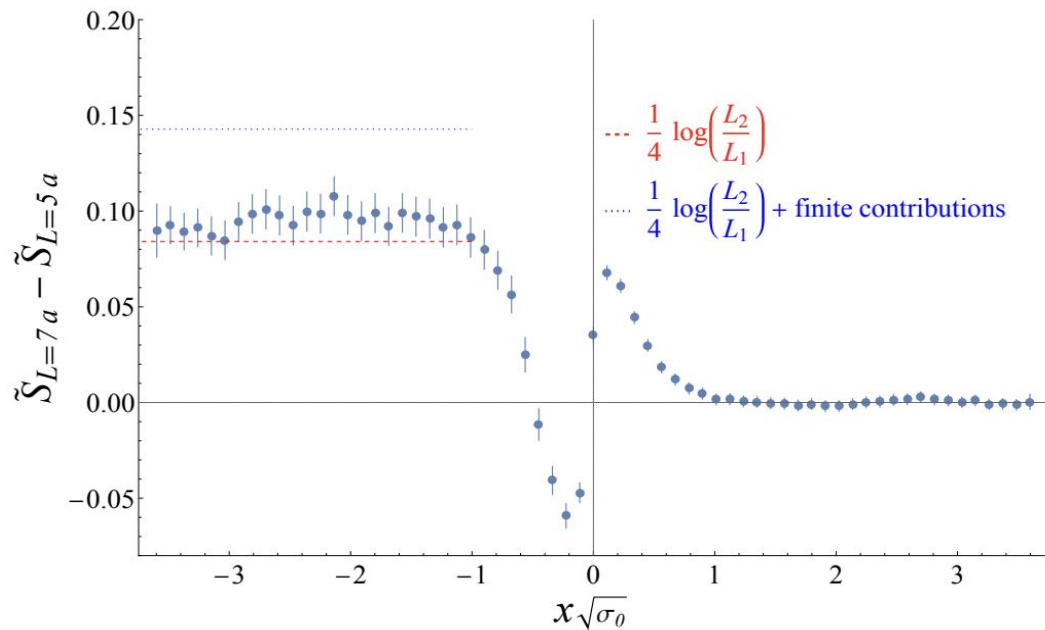


M. Luscher et al , Nucl. Phys. B 180, 1 (1981)

$$H = M^2 L + \frac{\pi}{2M^2 L} \int_0^\pi ds (p^2 + M^4 x'(s)^2)$$

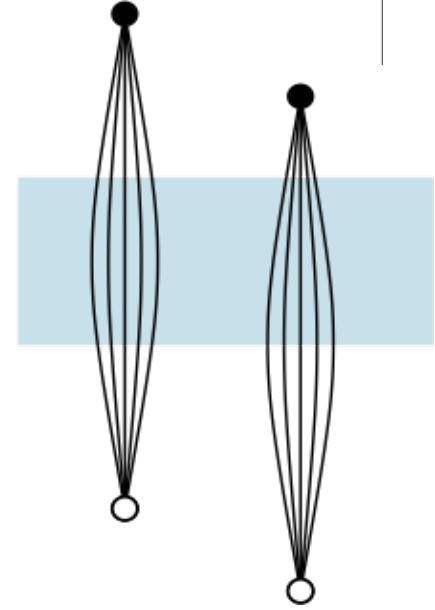
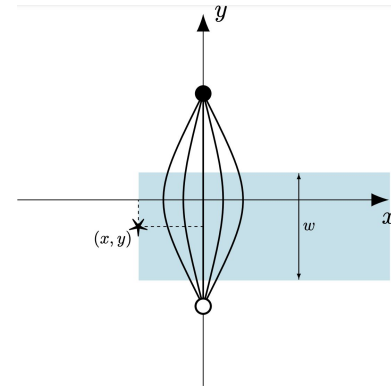
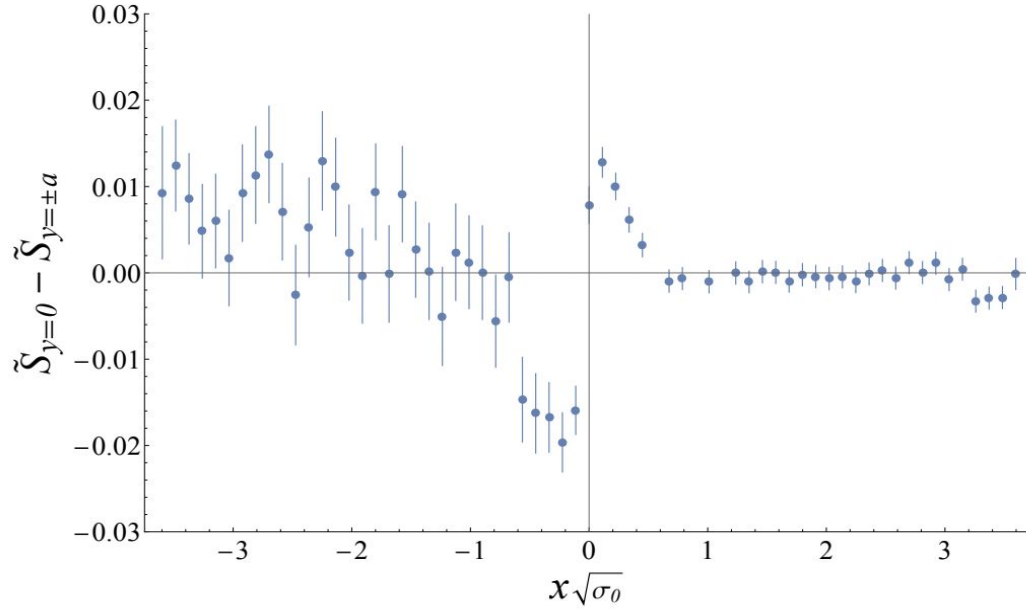


Entropy String Length Dependence



Takeaway: Present length scaling term, other finite contributions small/absent.

Entropy Location along the String Dependence



Takeaway: Finite contributions small/absent from vibrational entropy.